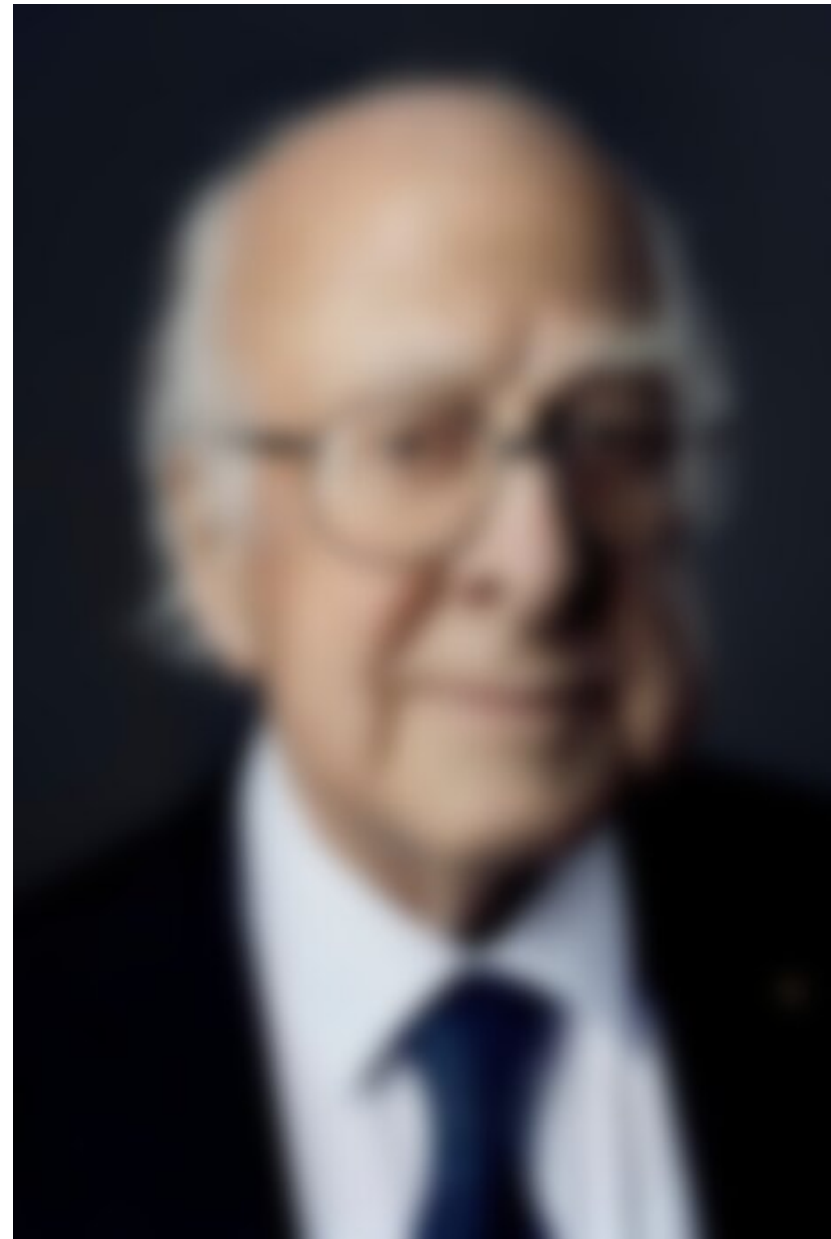


10 years of Higgs boson in ATLAS:

from discovery

to precision



Giovanni Marchiori (APC)

IRN Terascale
18 October 2022



10 years ago..

- 4th July 2012 @CERN: (experimental) birth of the Higgs boson

<https://indico.cern.ch/event/197461/>



*F. Gianotti
(ATLAS)*

*R. Heuer J. Incandela
(CERN) (CMS)*

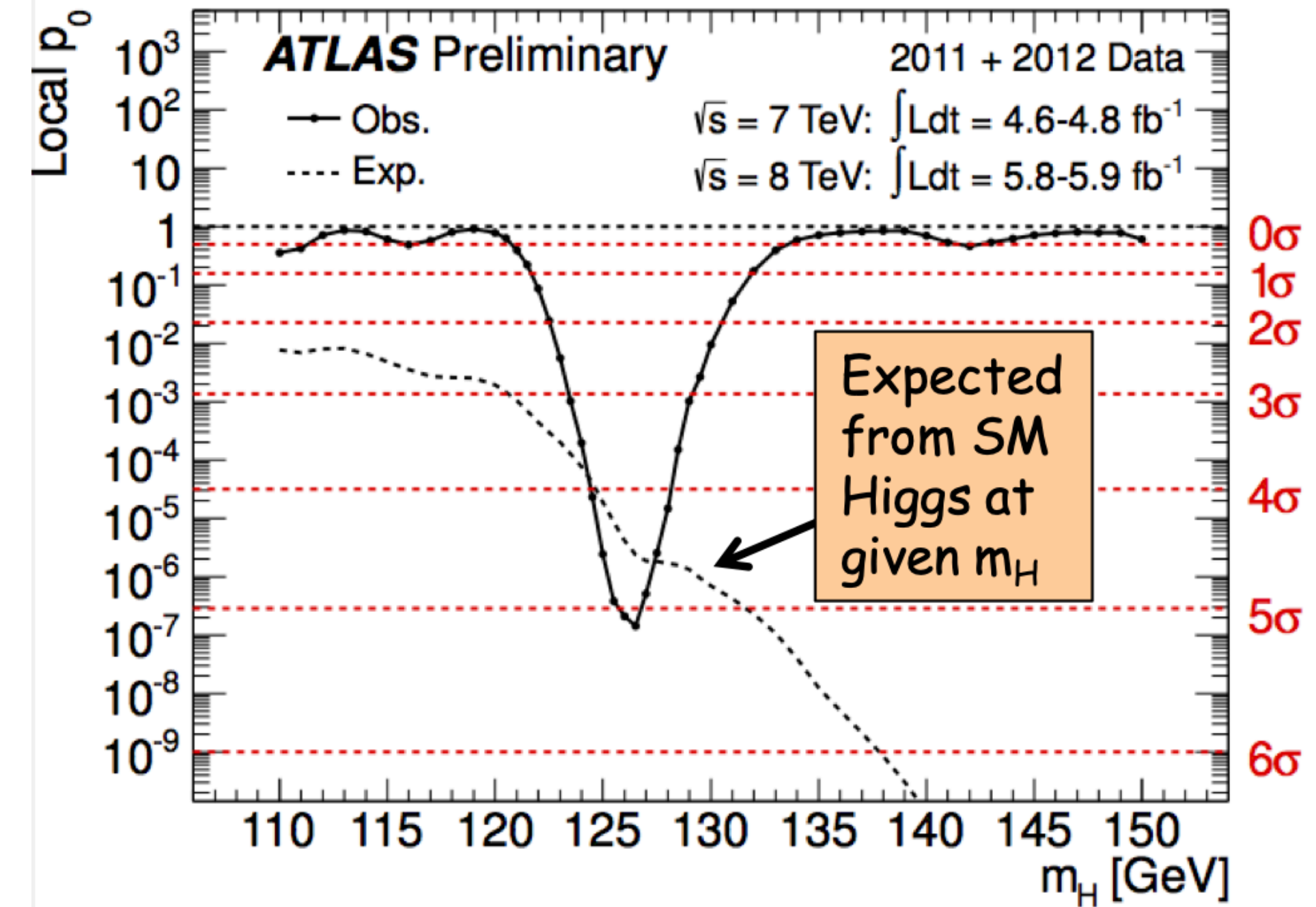
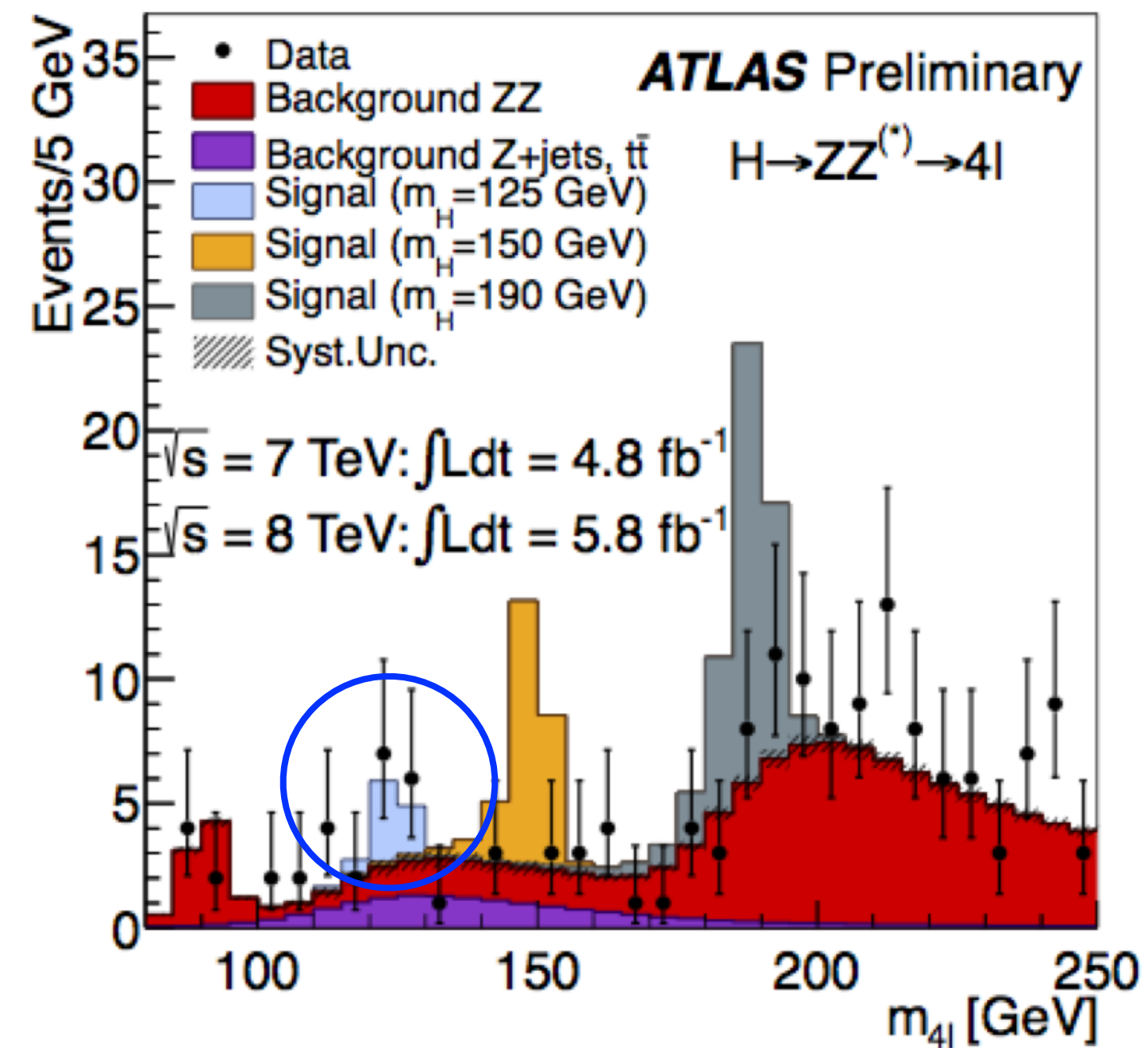
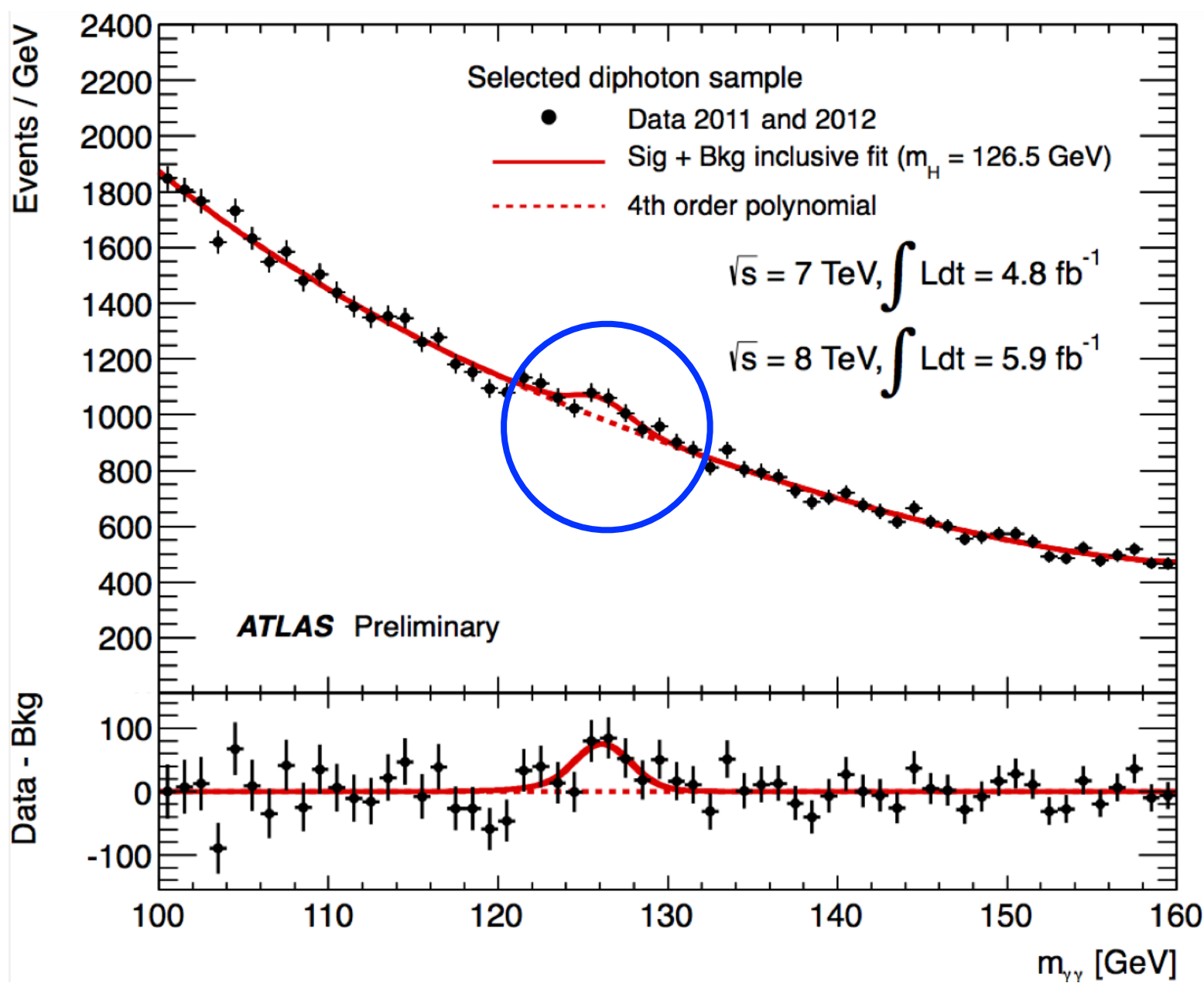


F. Englert

P. Higgs

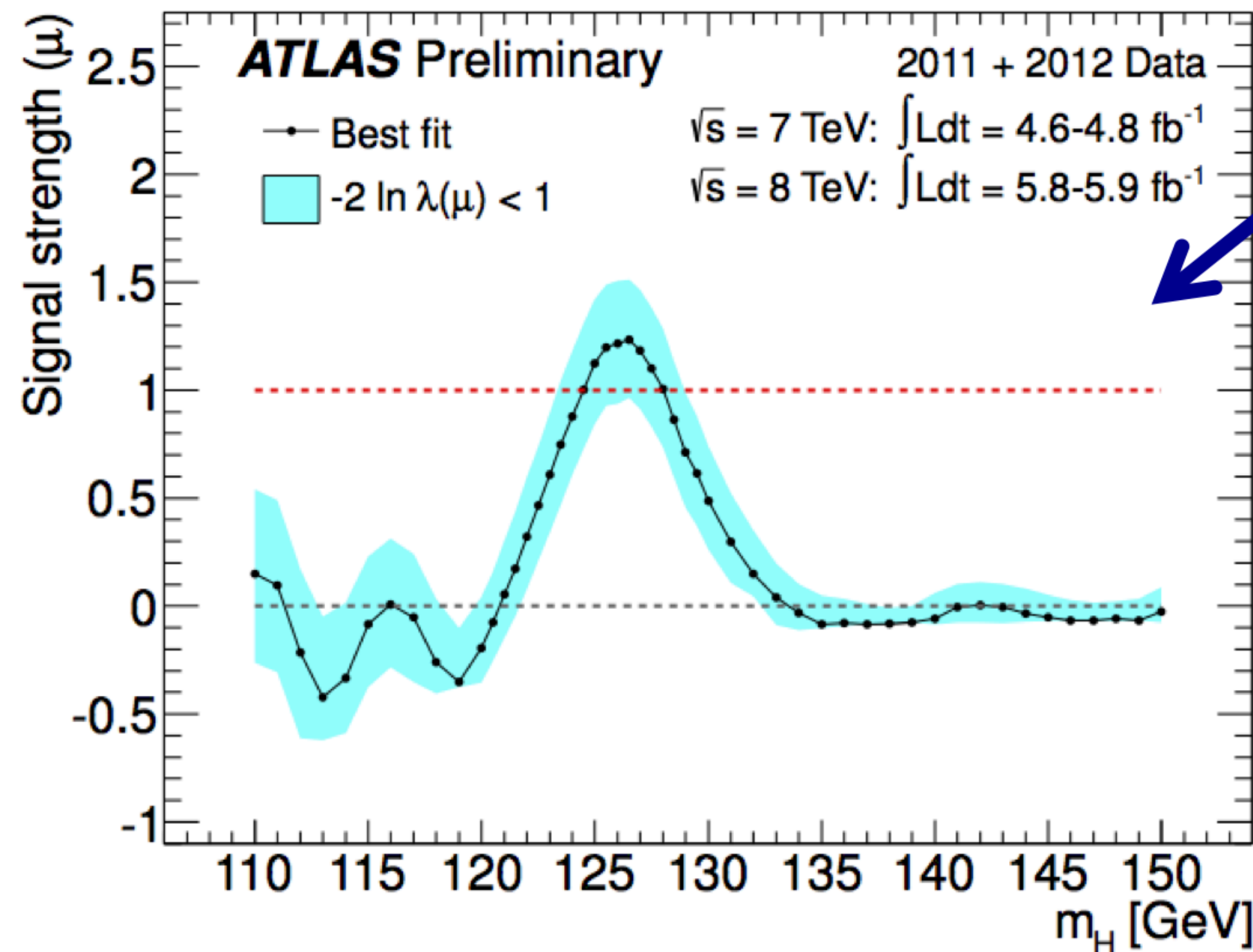
The Higgs boson discovery

- Analysis of up to 5/fb @ 7 TeV + 5/fb @ 8 TeV
- Main sensitivity / evidences from cleanest decay channels: $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^{(*)} \rightarrow 4l$ (full dataset)



The Higgs boson discovery

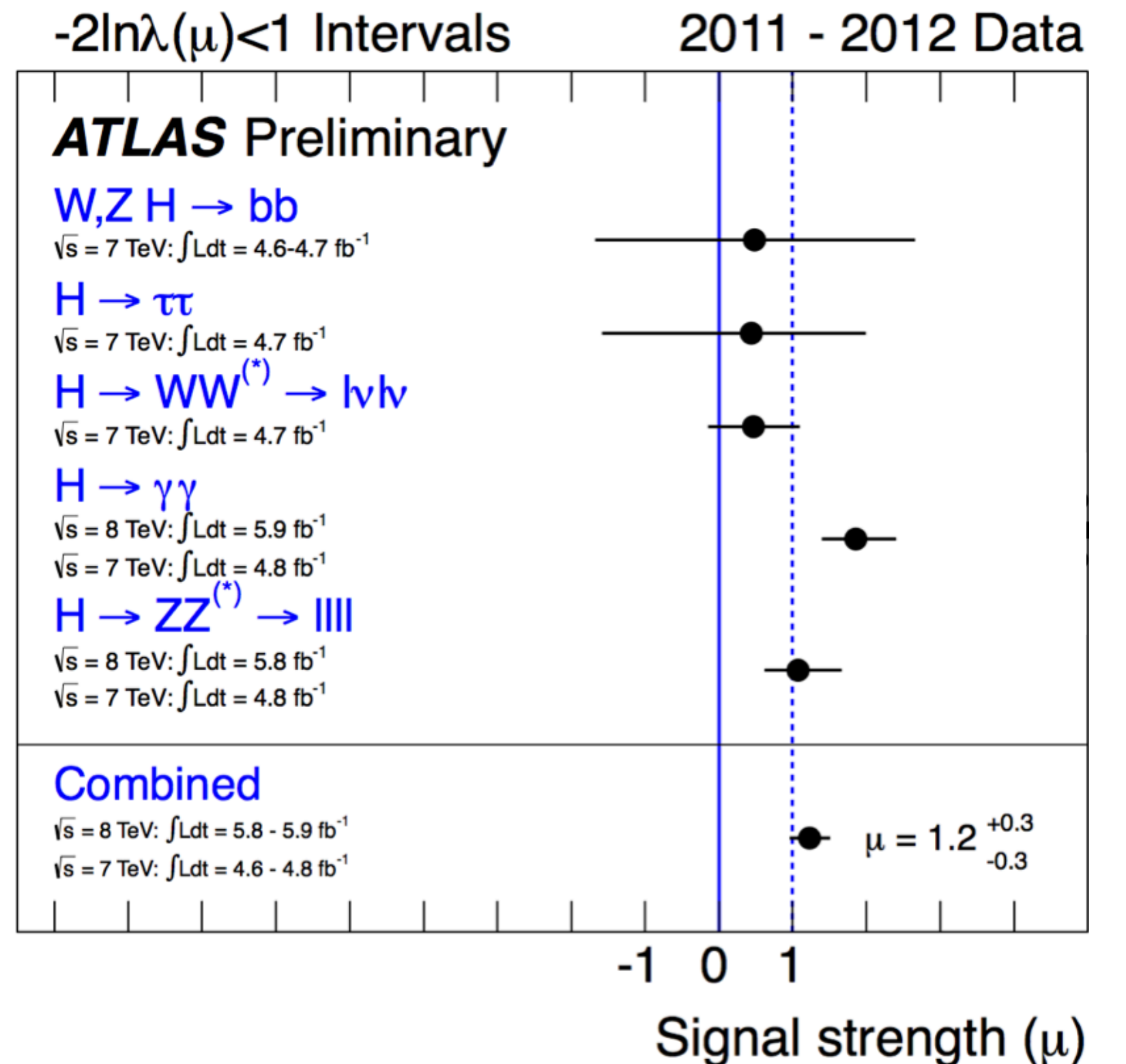
- Analysis of up to 5/fb @ 7 TeV + 5/fb @ 8 TeV
- Main sensitivity / evidences from cleanest decay channels: $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^{(*)} \rightarrow 4l$ (full dataset)
- Some additional sensitivity from additional decay channels. All “signal strengths” in agreement with Standard Model (SM) of 1, O(25%) relative uncertainty after combination



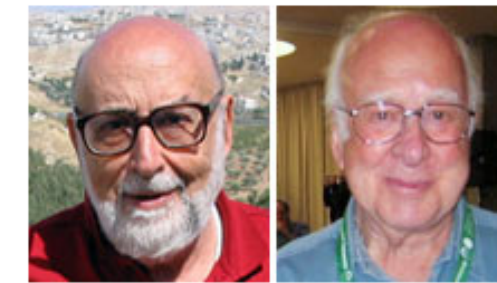
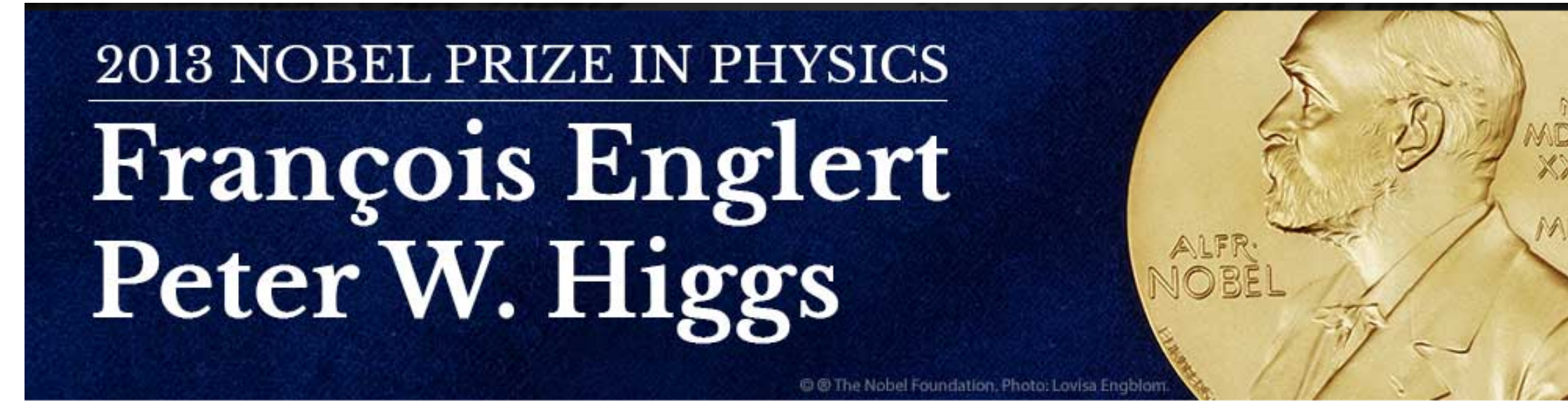
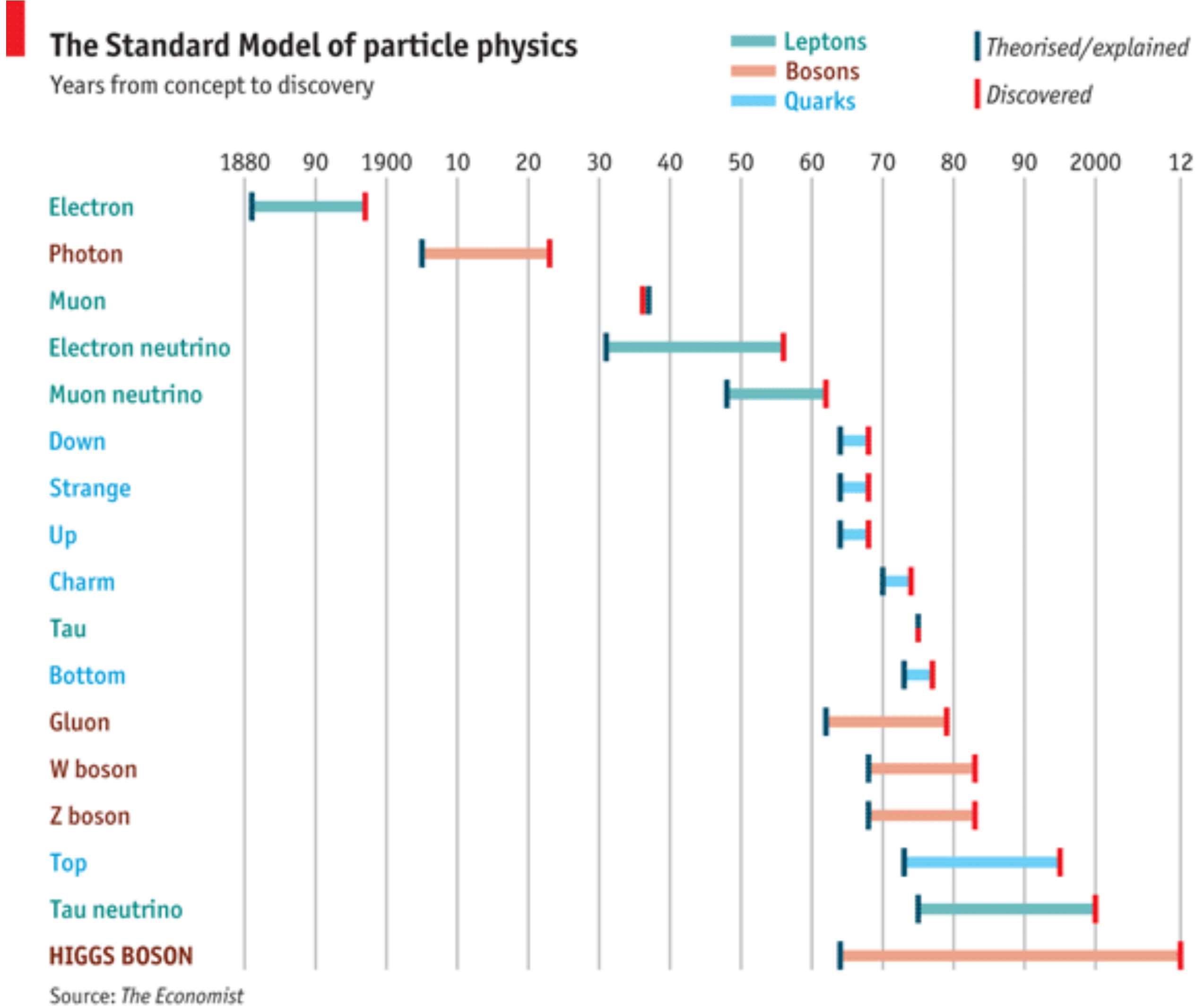
“signal strength”

Best-fit value at 126.5 GeV:
 $\mu = 1.2 \pm 0.3$

$$\mu = \frac{(\sigma \cdot BR)_{obs}}{(\sigma \cdot BR)_{SM}}$$



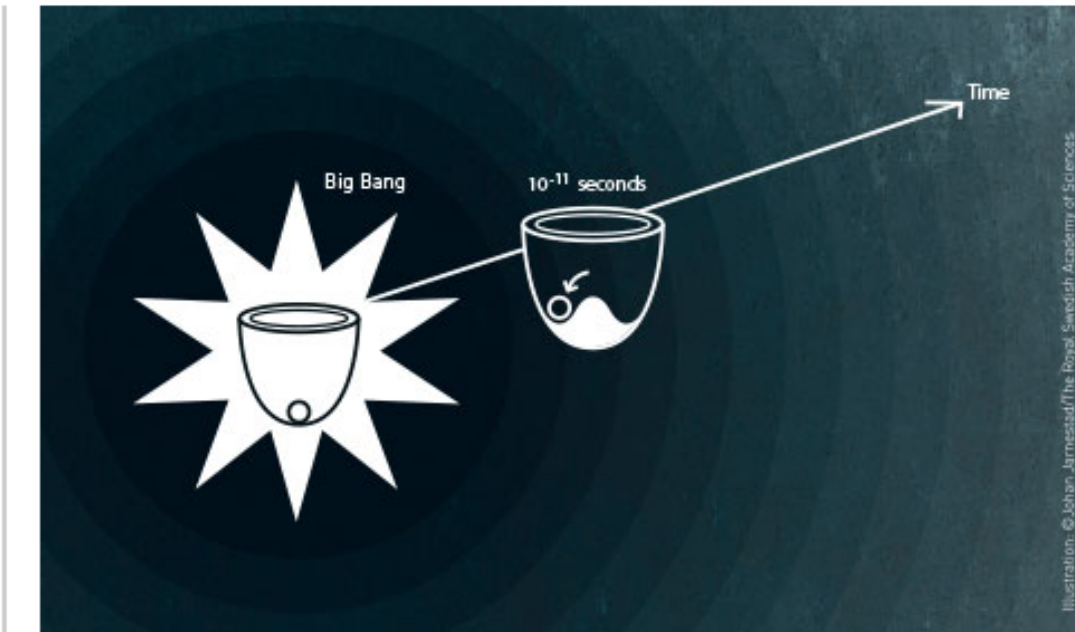
A long journey...



F. Englert and P. Higgs
Photo: Wikimedia Commons

2013 Nobel Prize in Physics

The Nobel Prize in Physics 2013 was awarded jointly to **François Englert** and **Peter W. Higgs** "for the theoretical discovery of a mechanism that contributes to our



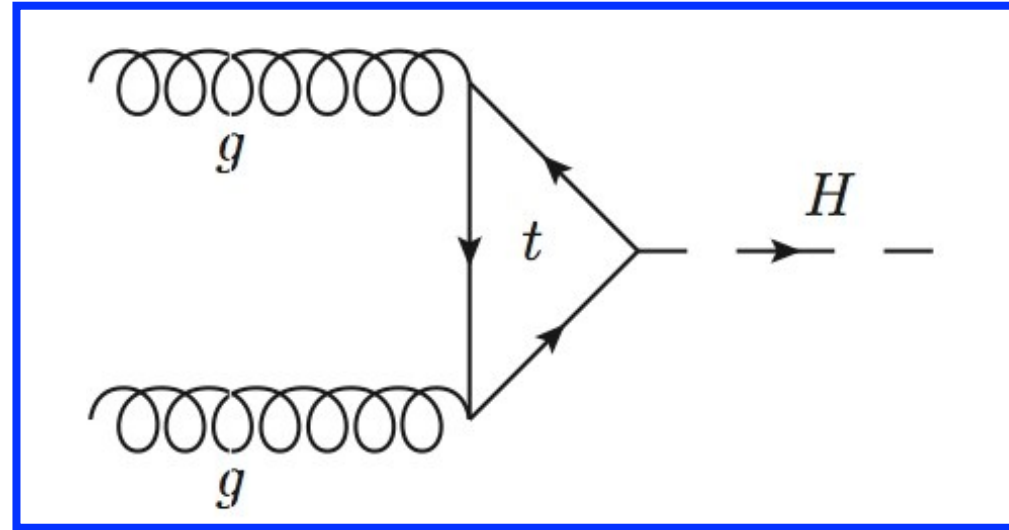
What Happened after the Big Bang?

Announcements of the 2013 Nobel Prizes

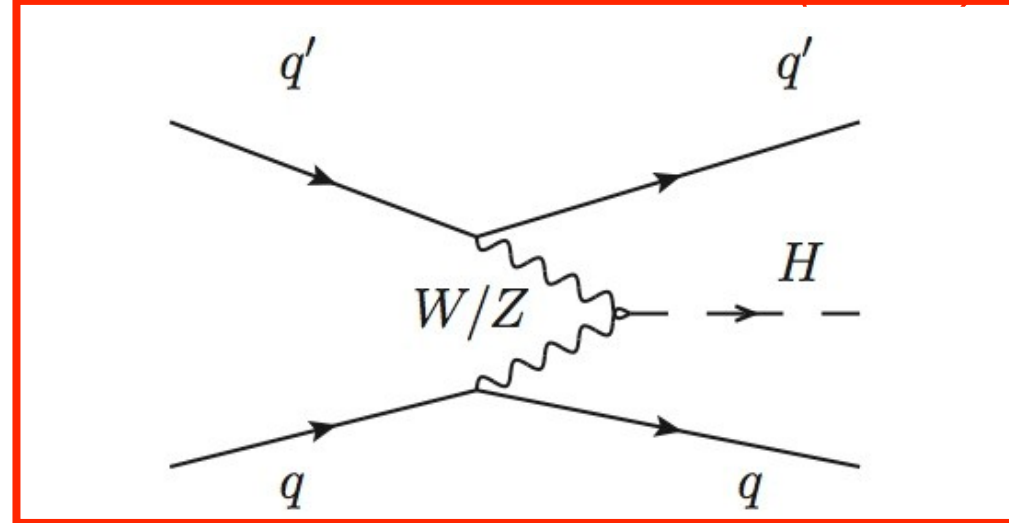
- Physiology or Medicine:
Announced Monday 7 October
- Physics:
Tuesday 8 October, 11:45 a.m. CET at the earliest
- Chemistry:
Wednesday 9 October, 11:45 a.m. CET at the earliest
- Literature:
Thursday 10 October 1.00 p.m. CET
- Peace:

Producing a Higgs boson at the LHC

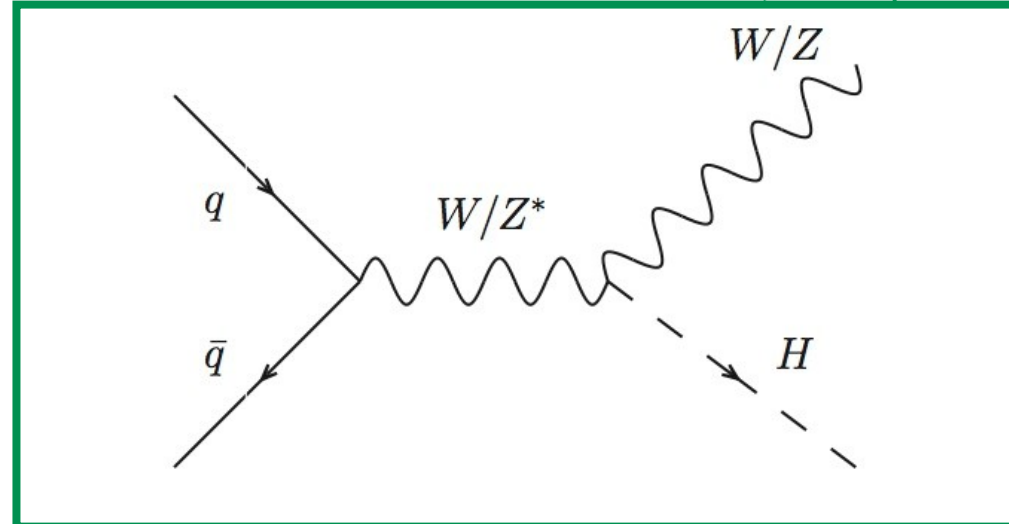
gluon-gluon fusion (ggF)



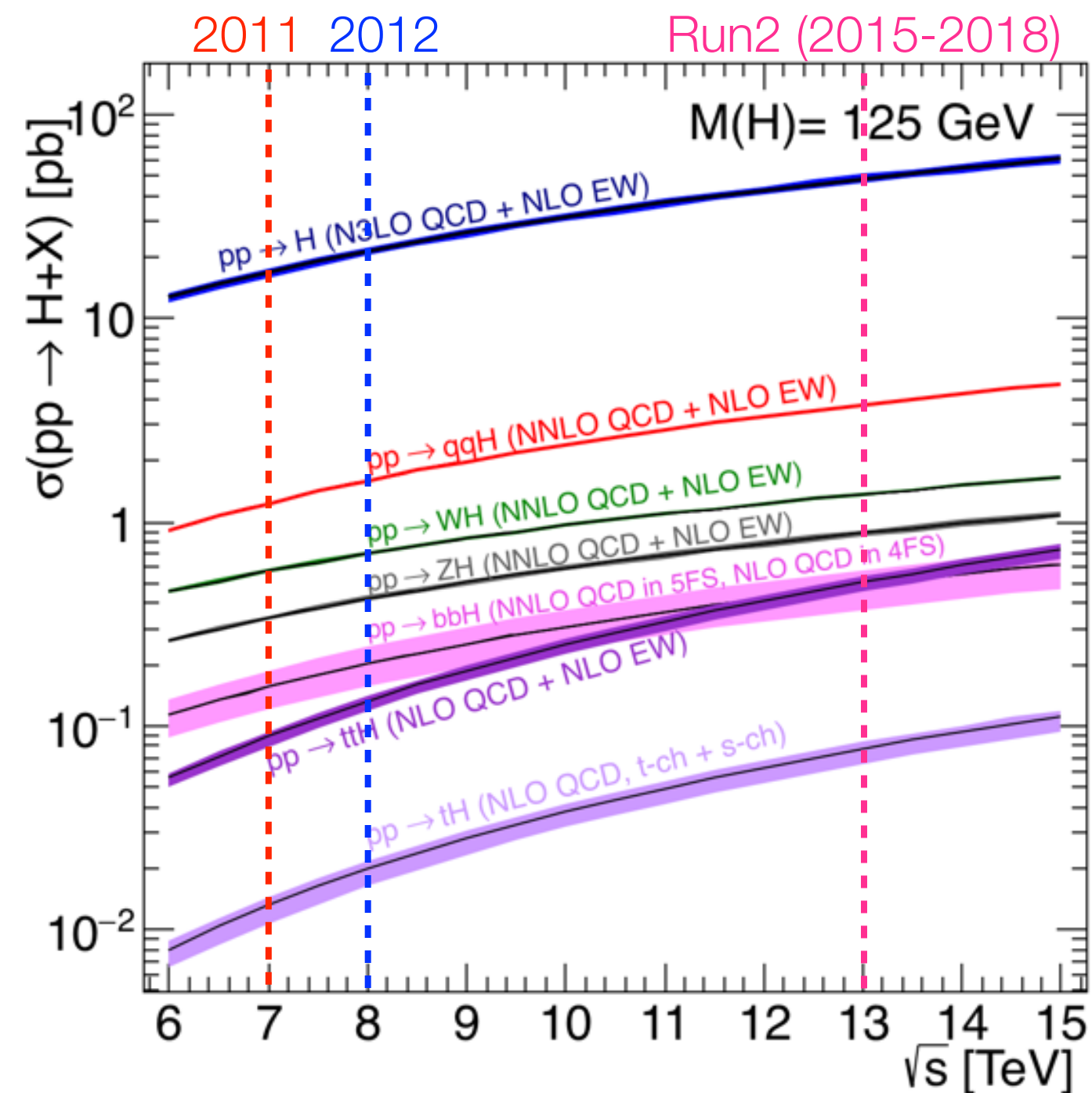
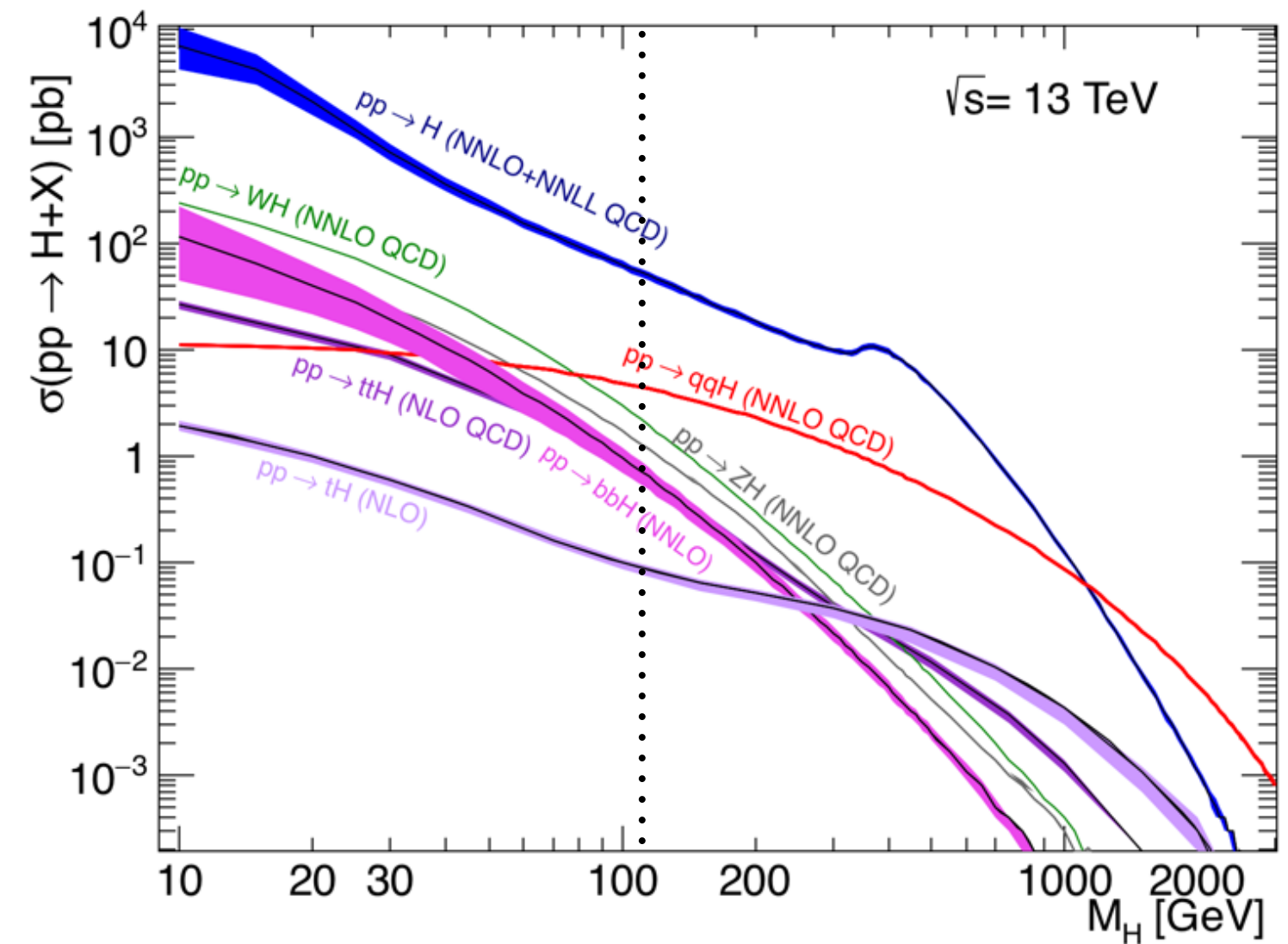
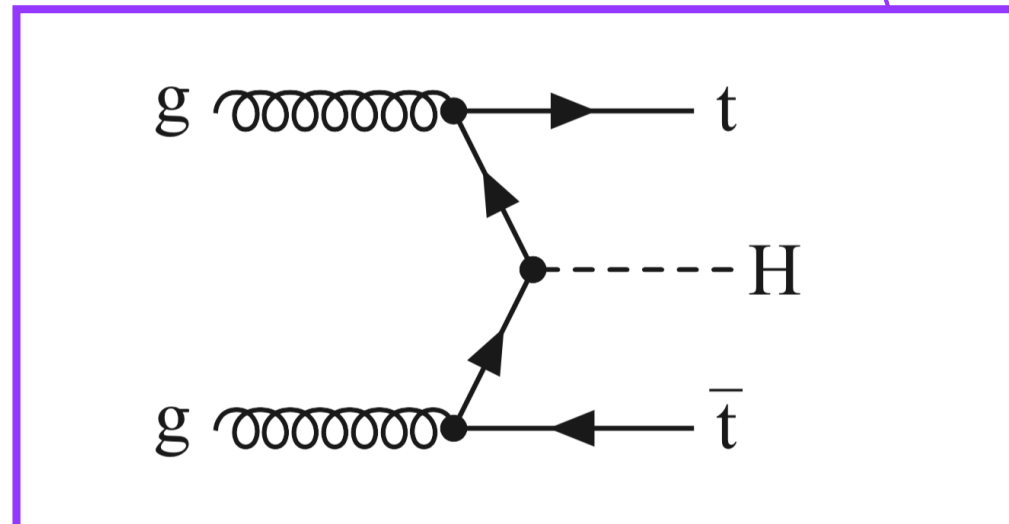
vector boson fusion (VBF)



associated with V=W, Z (VH)



associated with ttbar (ttH)

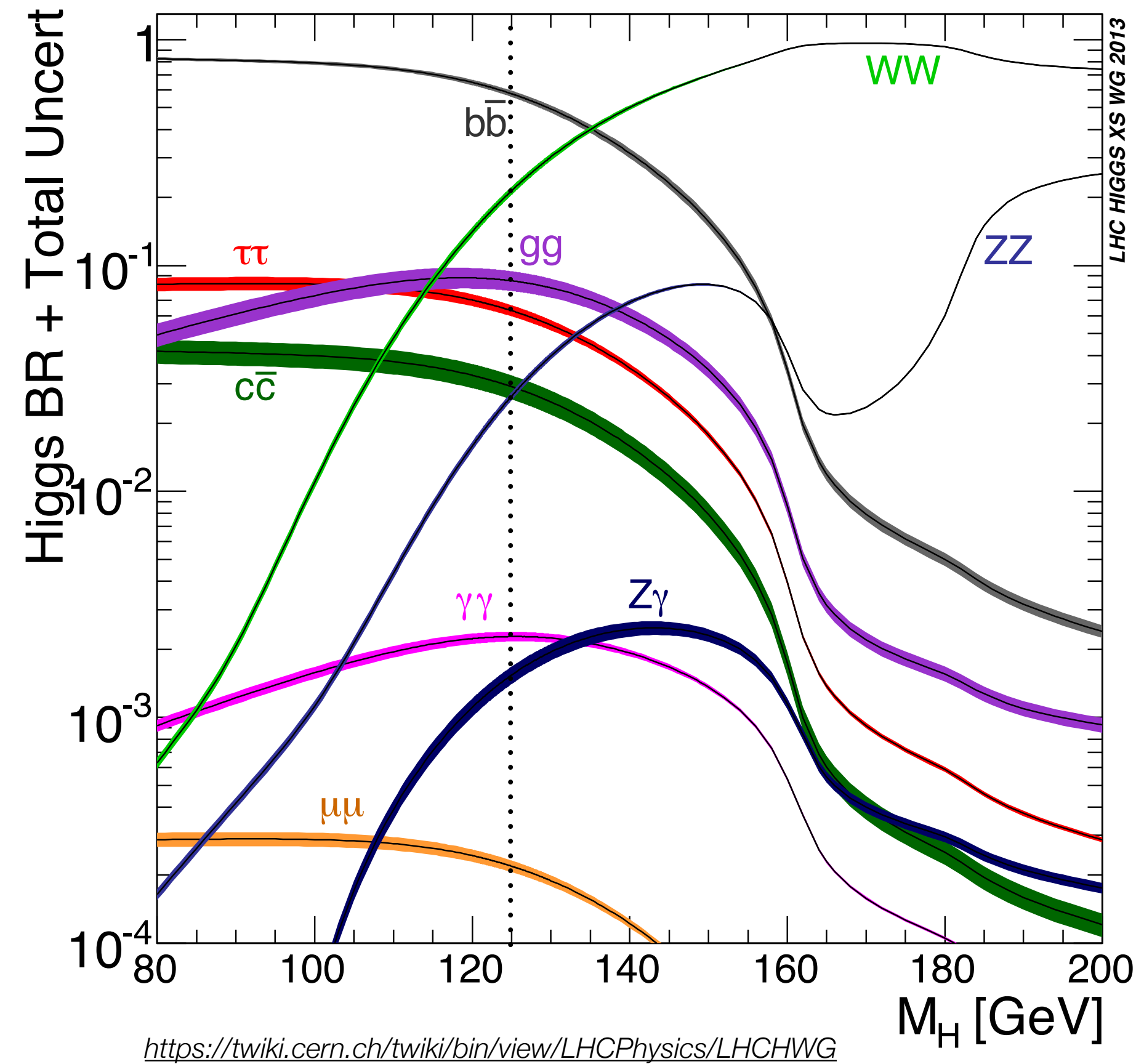


	σ $\sqrt{s}=13$ TeV $m=125$ GeV	N(Higgs) in 140/fb
ggF	49 pb	6.9M
VBF	3.8 pb	530k
VH	2.3 pb	320k
ttH	0.5 pb	70k
TOTAL	56 pb	7.8M

~600k Higgs bosons produced in Run1, 8M in Run2

⇒ LHC = 'Large Higgs Creator'!

Higgs boson decays

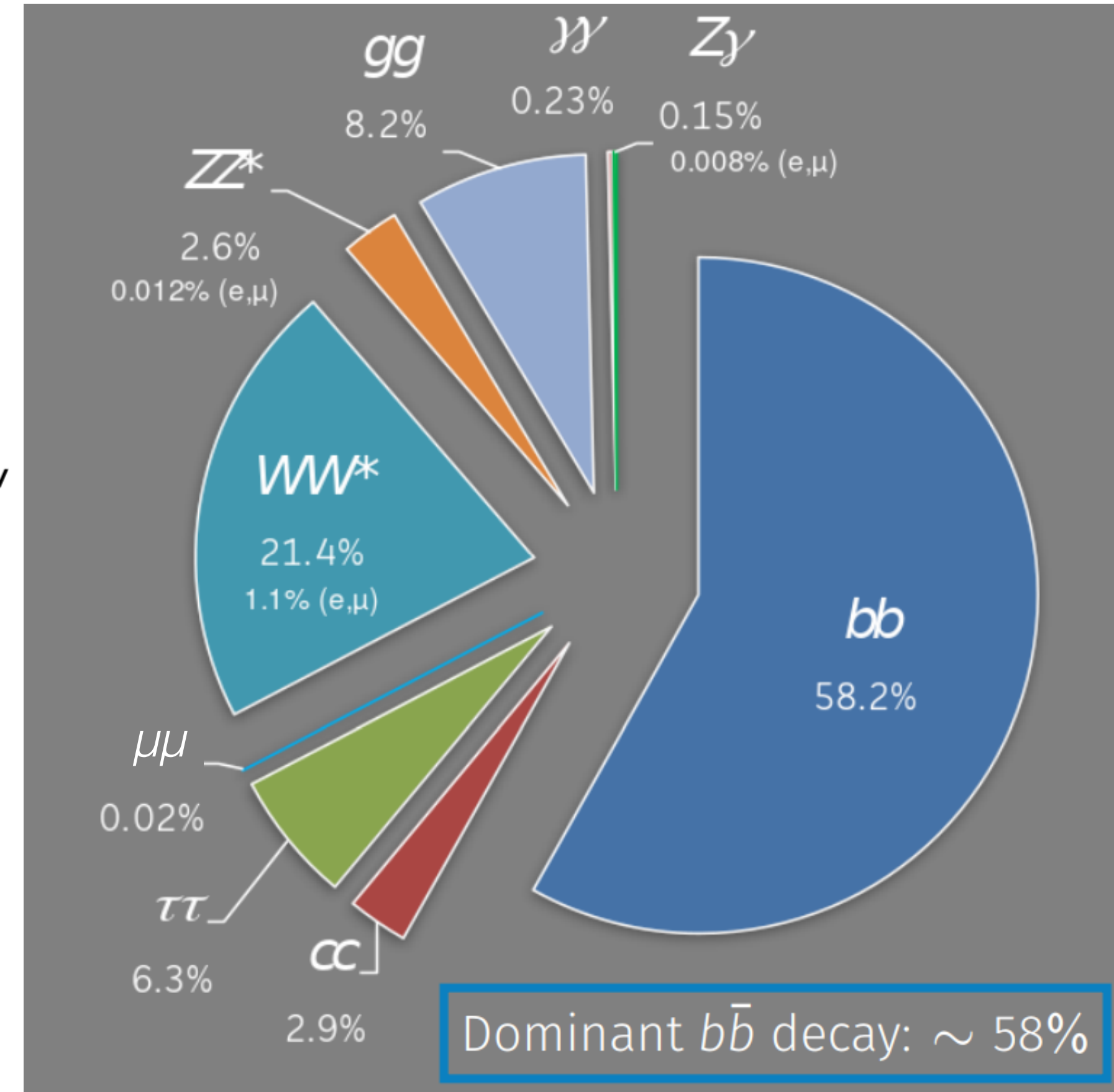


Δ
ZZ, $\gamma\gamma$: high mass resolution channels mass and precise differential measurements

WW: High BR, but low mass resolution

$\mu\mu$: very small BR, but access to coupling to 2nd generation fermions

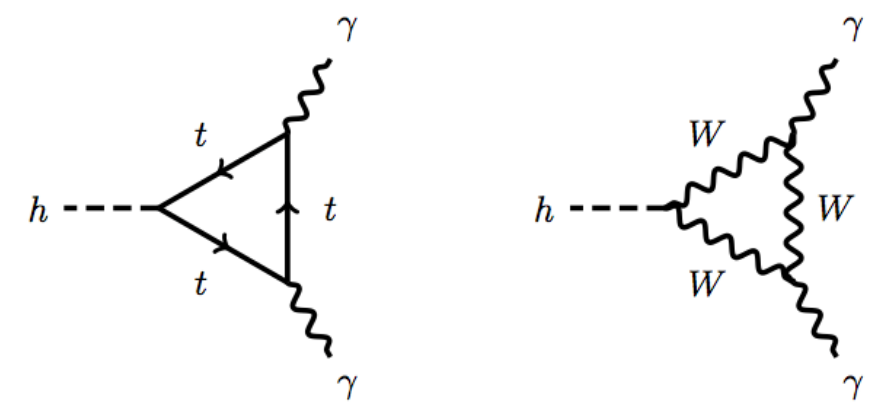
gg: high BR, but very poor q/g discrimination



bb, $\tau\tau$: high BR, but low S/B, important to directly probe Higgs boson coupling to fermions

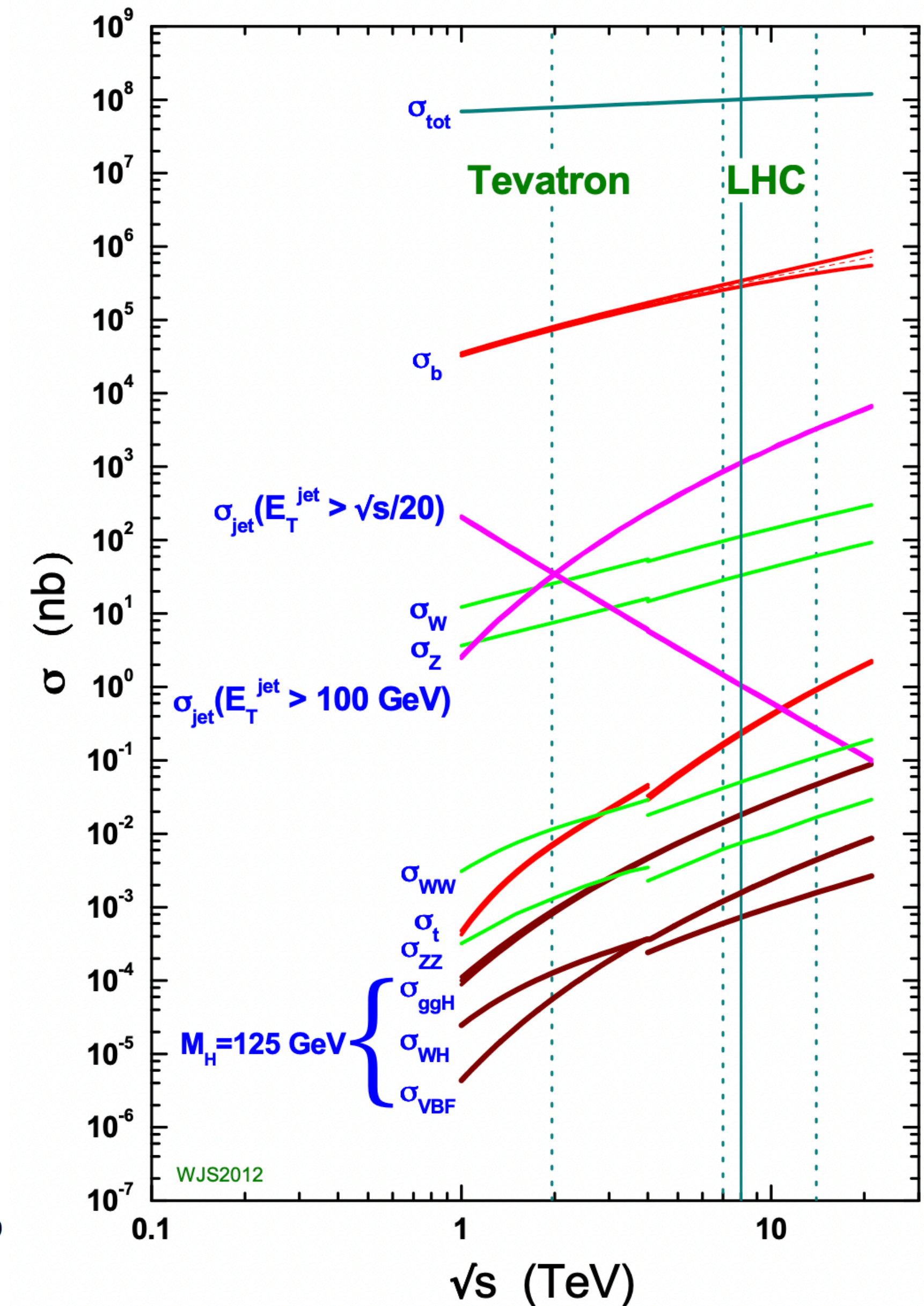
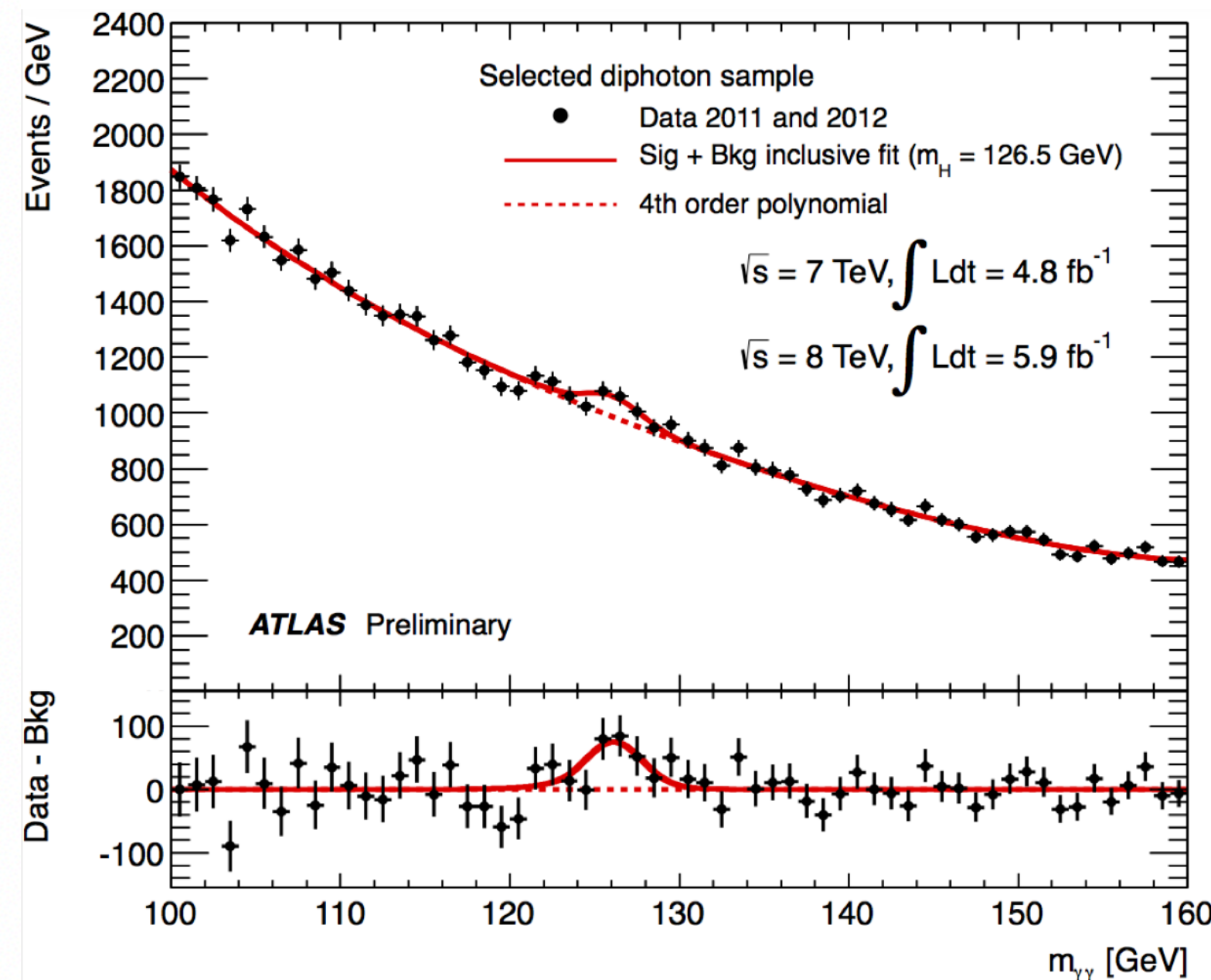
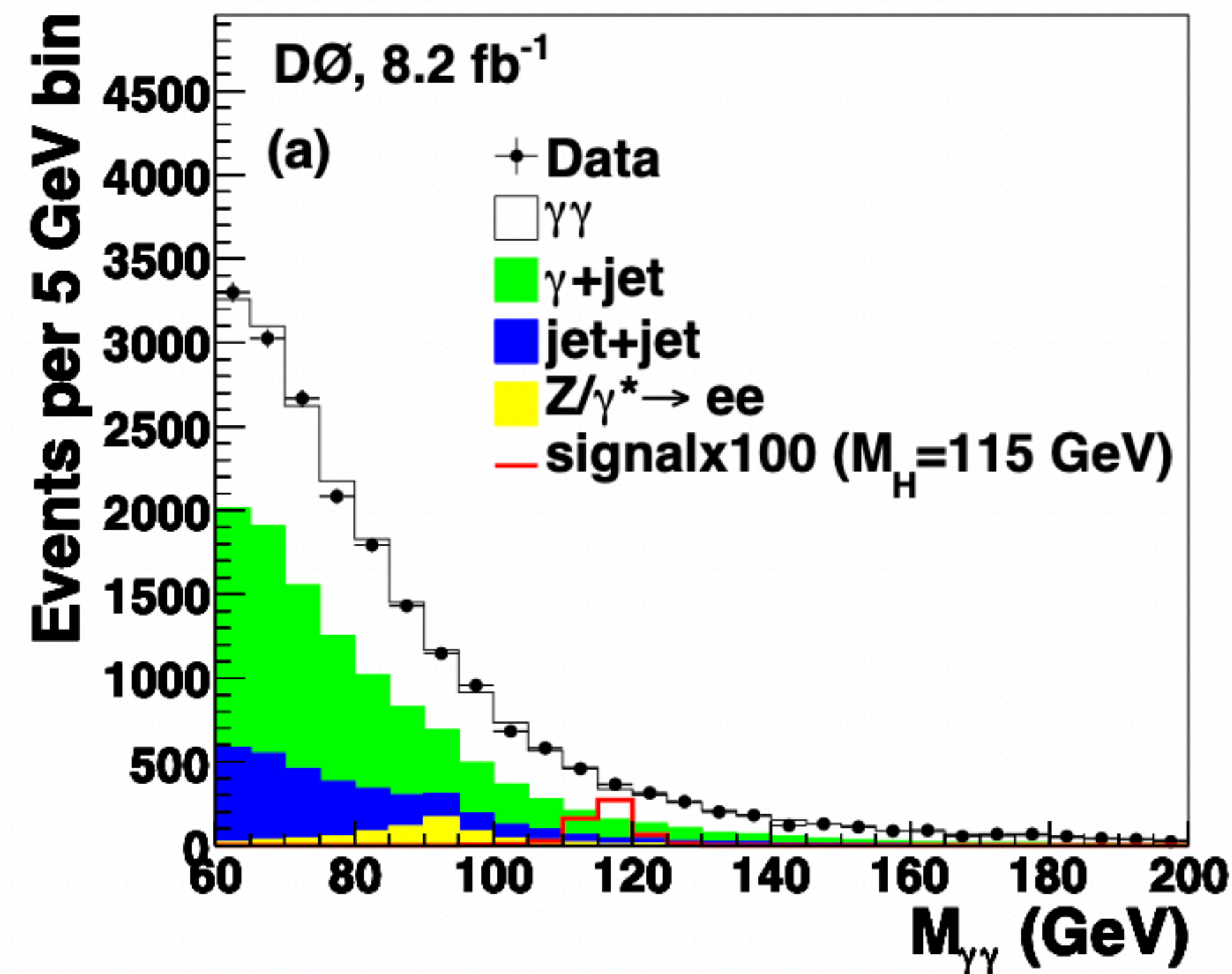
cc: decent BR, but poor q,g rejection => challenging for coupling to 2nd gen fermions

Loop-induced effective photon coupling:



Ingredients for the discovery - or why the LHC beat the Tevatron

- **Large production cross section:** LHC $\sim 30\times$ Tevatron
 - ➔ 30x larger yield at time of discovery (\sim same luminosity, 10/fb)
- **State-of-the-art detectors** with efficient and precise reconstruction and identification of leptons, photons & jets
 - ➔ 2x efficiency with S/B \sim an order of magnitude larger in $H \rightarrow \gamma\gamma$
- Use of **multivariate techniques** (BDTs) to “tag” rarer production modes with better S/B
 - ➔ 20-50% improvements in analysis sensitivity



Fast forward to 2022...

- 10 years after the Higgs boson discovery, we have been able to understand many more of its properties and compare them to the predictions of the Standard Model:

- Mass
- Width
- Spin & parity
- Main decay and production modes
- Rare decays
- Couplings to other particles
- Differential cross sections
- Self-coupling
- beyond-SM interactions (anomalous couplings to SM particles, CP violation, coupling to dark matter sector..)

PDG 2012

Higgs Bosons — H^0 and H^\pm , Searches for

The July 2012 news about Higgs searches is described in the addendum to the Higgs review in the data listings, but is not reflected here.

The limits for H_1^0 and A^0 refer to the m_h^{\max} benchmark scenario for the supersymmetric parameters.

H^0 Mass $m > 115.5$ and none 127–600 GeV, CL = 95%

H_1^0 in Supersymmetric Models ($m_{H_1^0} < m_{H_2^0}$)

Mass $m > 92.8$ GeV, CL = 95%

PDG 2022

H^0

$J = 0$

Mass $m = 125.25 \pm 0.17$ GeV ($S = 1.5$)
 Full width $\Gamma = 3.2^{+2.8}_{-2.2}$ MeV (assumes equal on-shell and off-shell effective couplings)

H^0 Signal Strengths in Different Channels

Combined Final States = 1.13 ± 0.06

$W W^* = 1.19 \pm 0.12$

$Z Z^* = 1.01 \pm 0.07$

$\gamma\gamma = 1.10 \pm 0.07$

$c\bar{c}$ Final State = 37 ± 20

$b\bar{b} = 0.98 \pm 0.12$

$\mu^+\mu^- = 1.19 \pm 0.34$

$\tau^+\tau^- = 1.15^{+0.16}_{-0.15}$

$Z\gamma < 3.6$, CL = 95%

$\gamma^*\gamma$ Final State = 1.5 ± 0.5

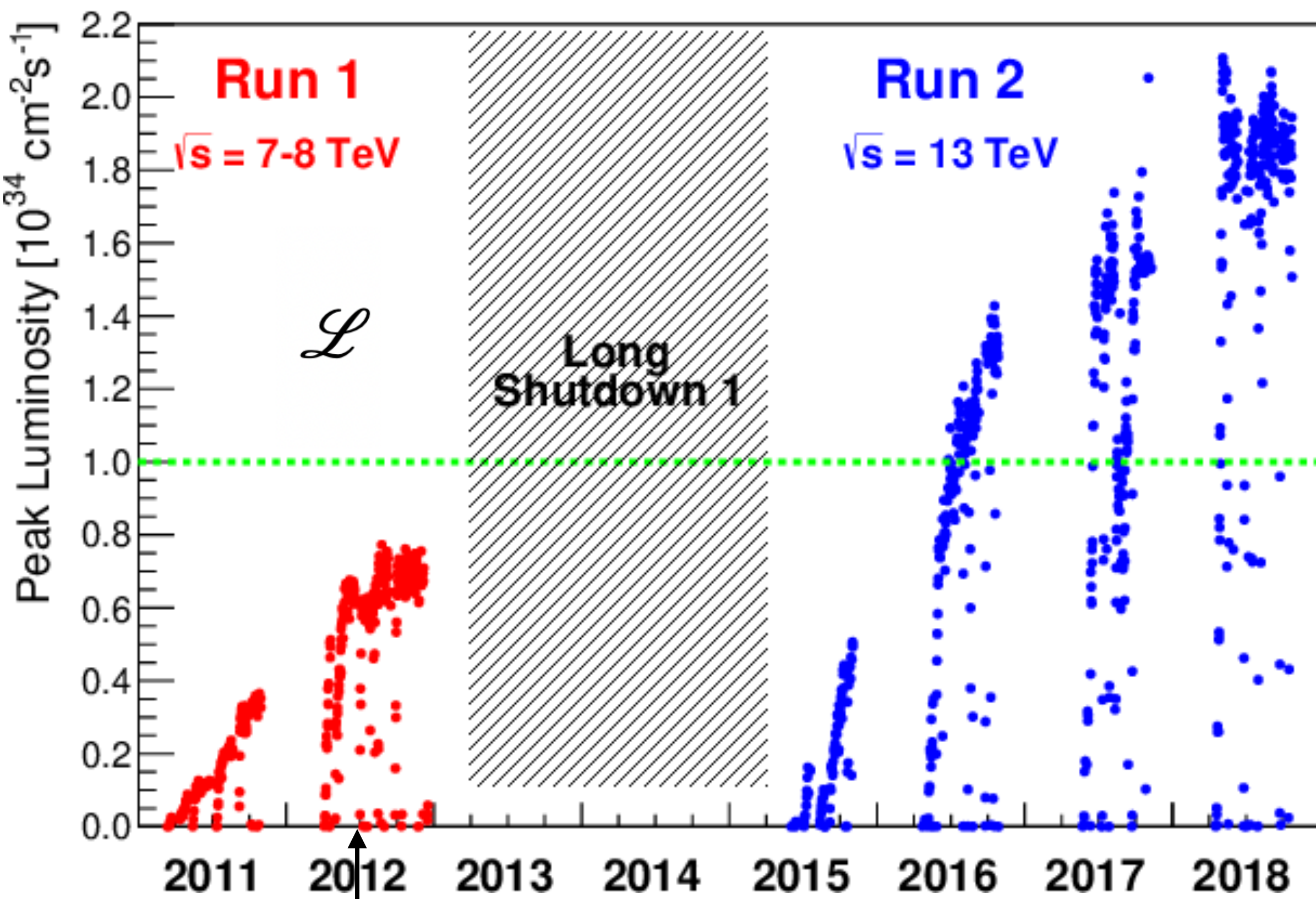
$t\bar{t}H^0$ Production = 1.10 ± 0.18

tH^0 production = 6 ± 4

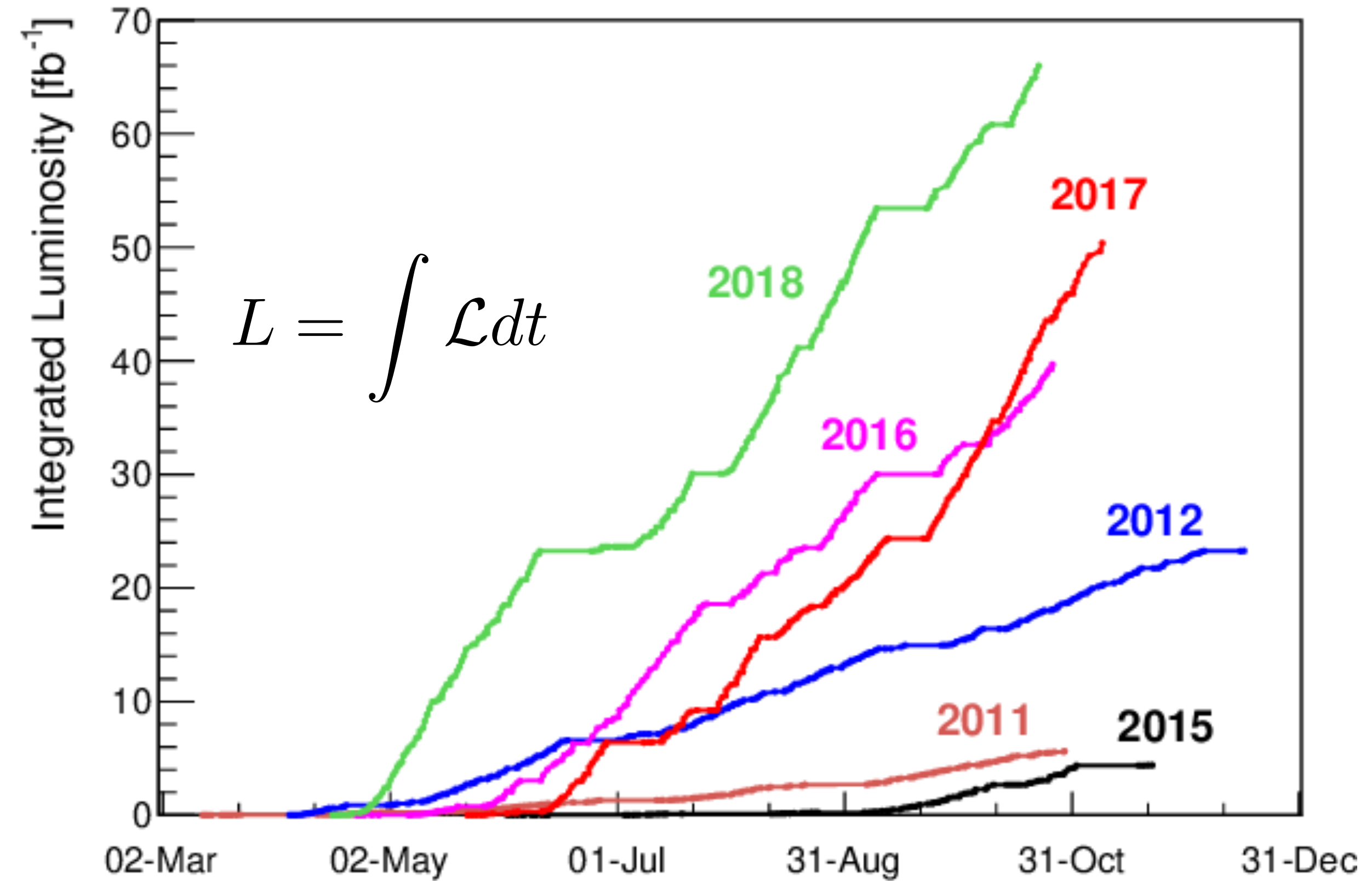
H^0 Production Cross Section in pp Collisions at $\sqrt{s} = 13$ TeV = 56 ± 4 pb

LHC luminosity evolution

$$N_H = L\sigma_H$$



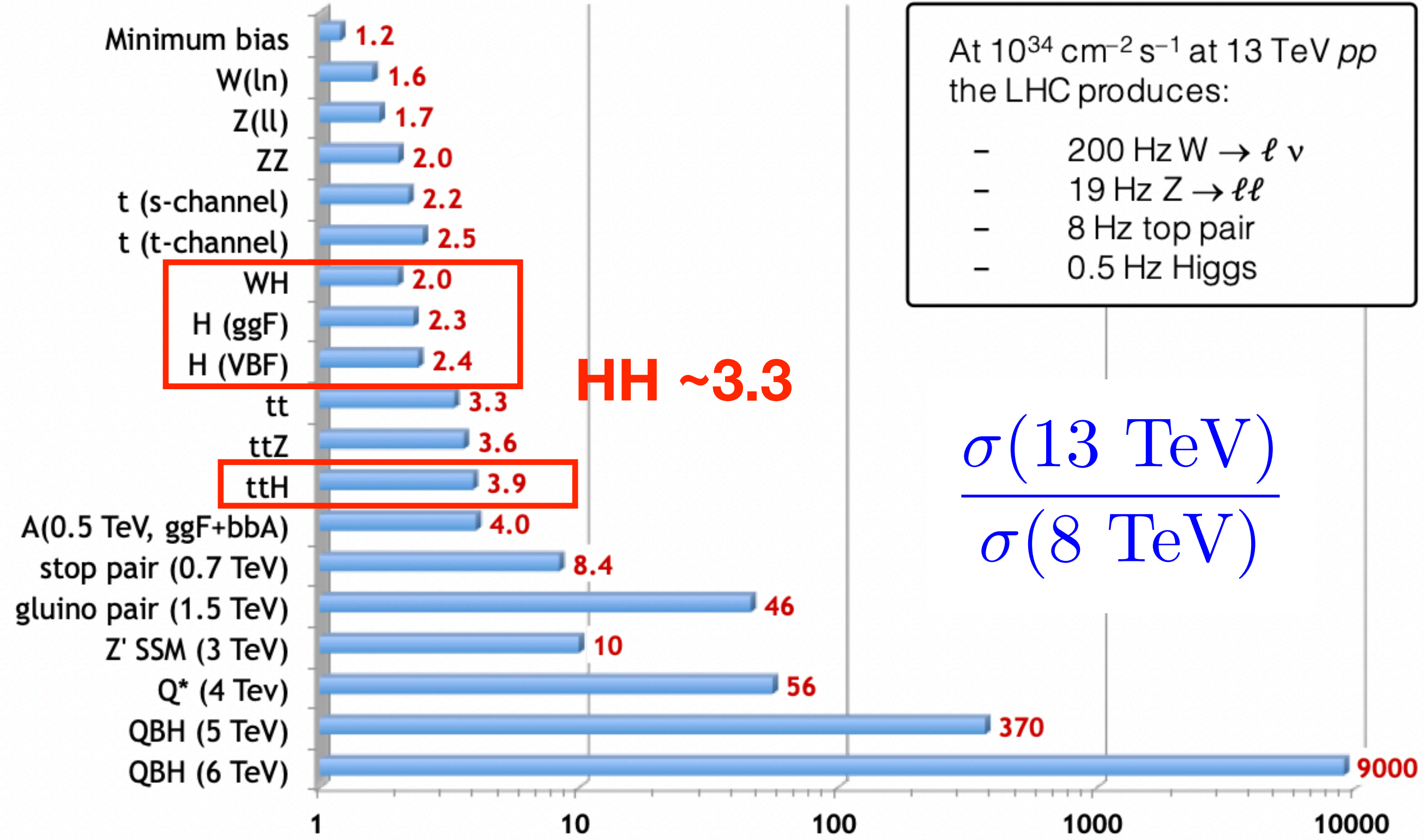
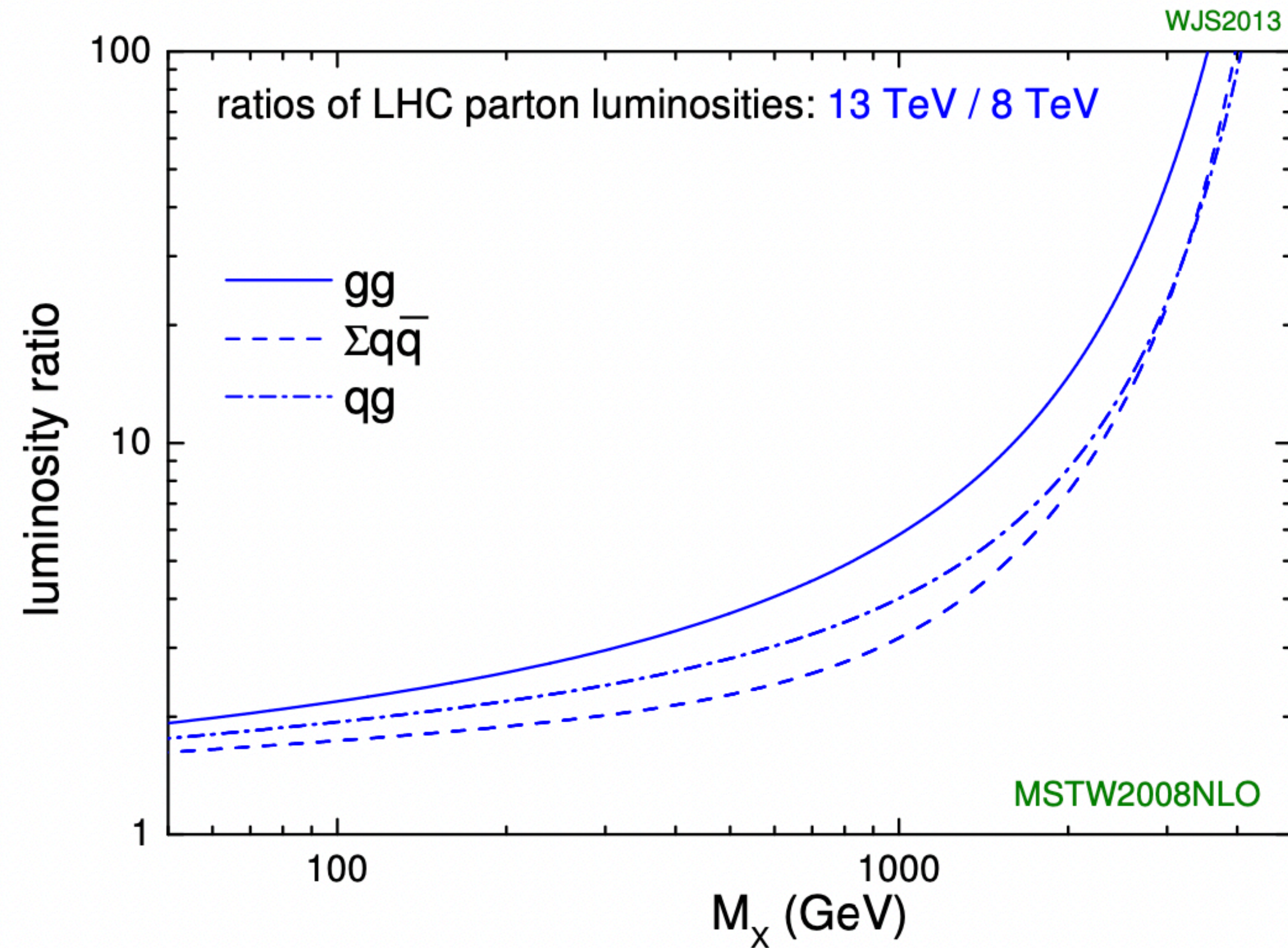
discovery



$L = 25/\text{fb}$ at 7-8 TeV in Run1
 $L = 140/\text{fb}$ at 13 TeV in Run2

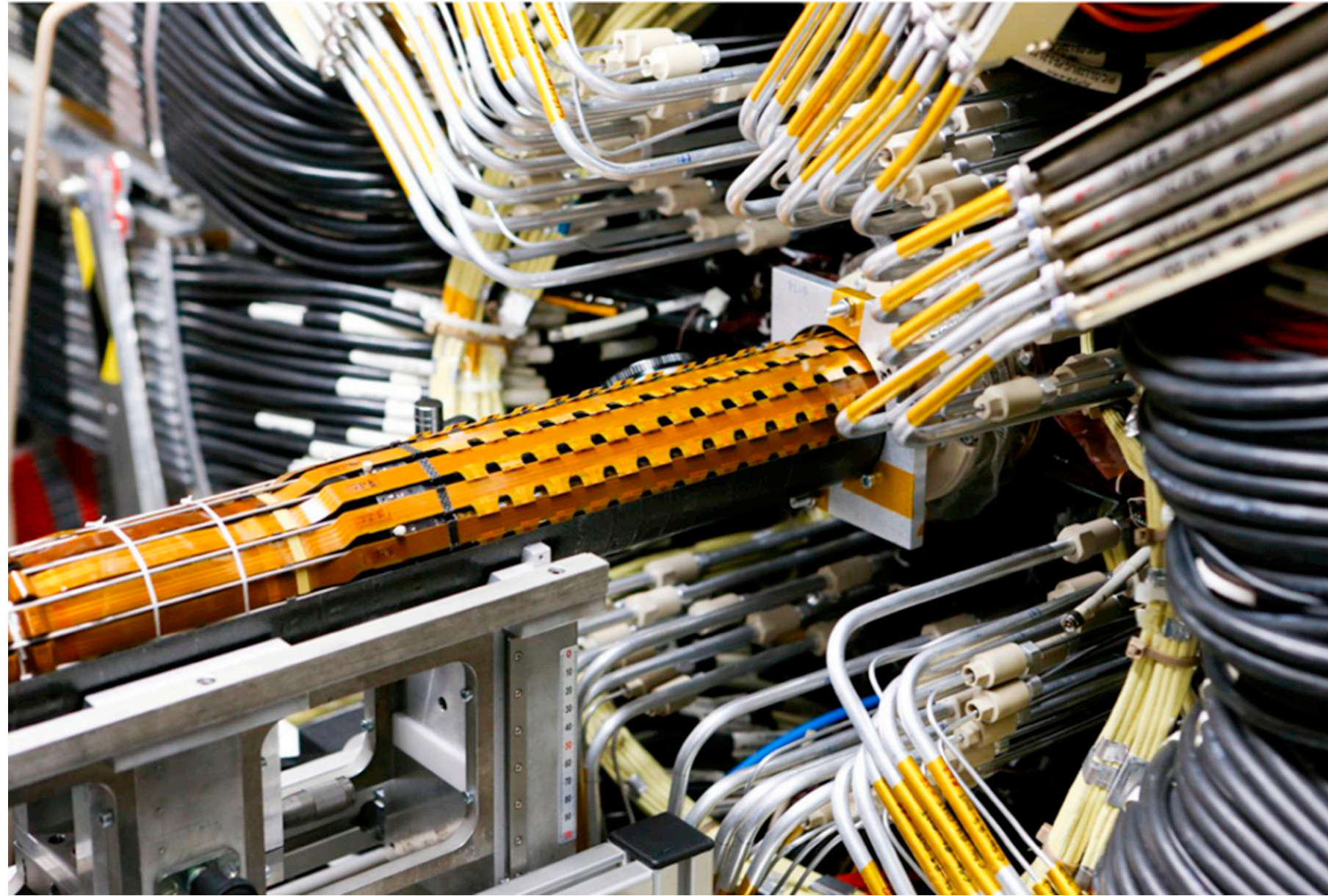
Higgs boson cross sections: from 8 to 13 TeV

$$N_H = L\sigma_H$$

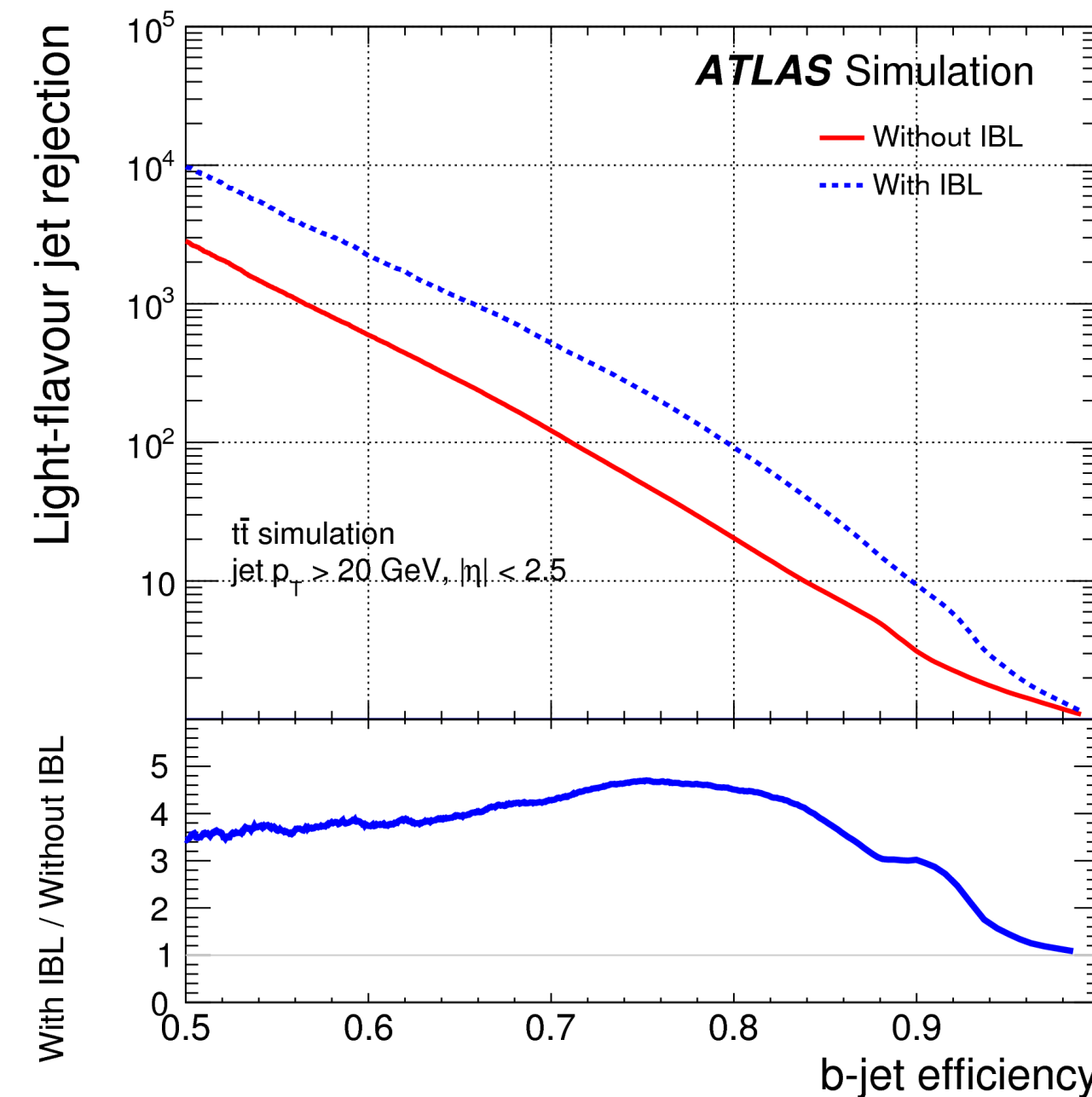
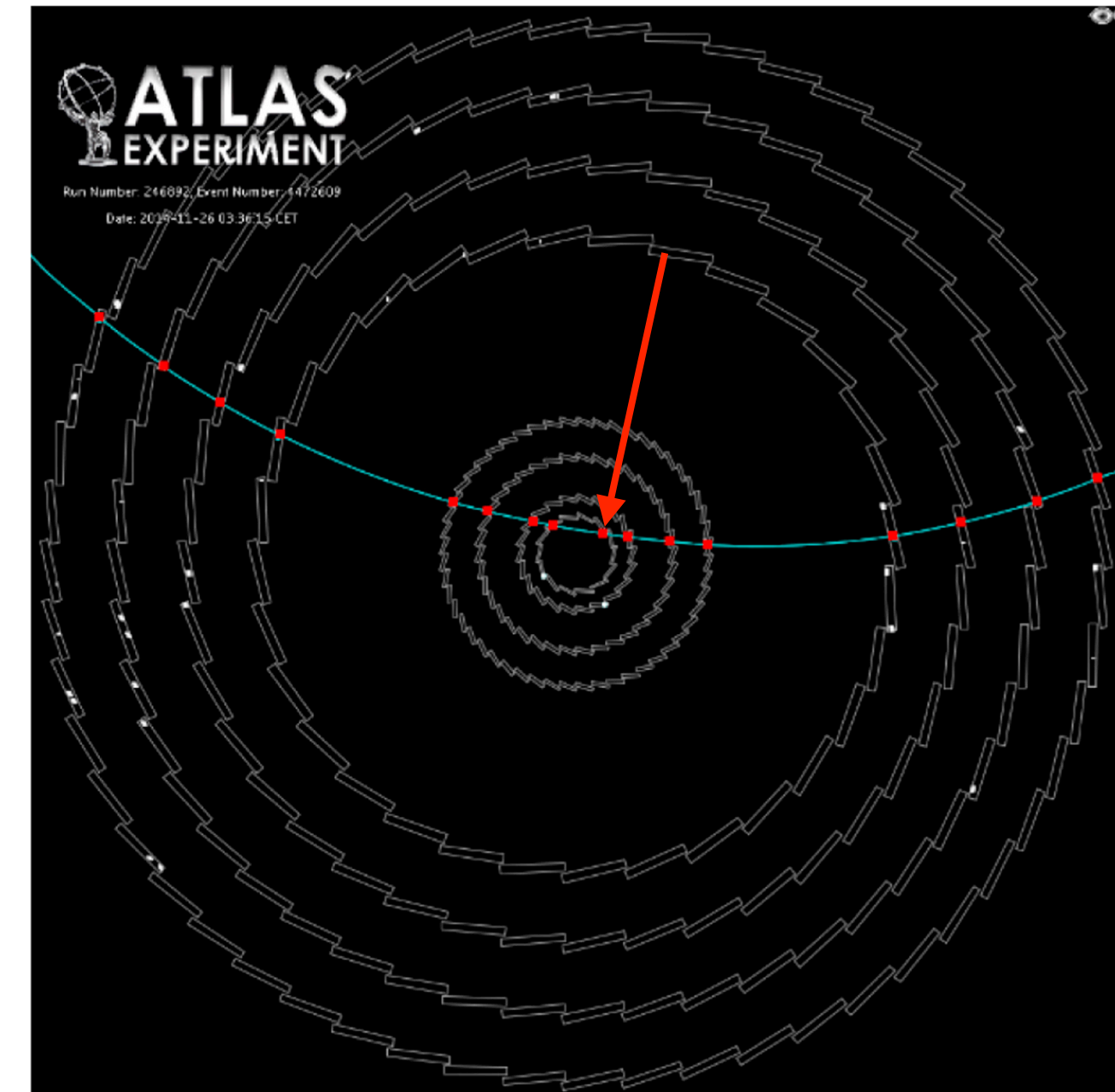
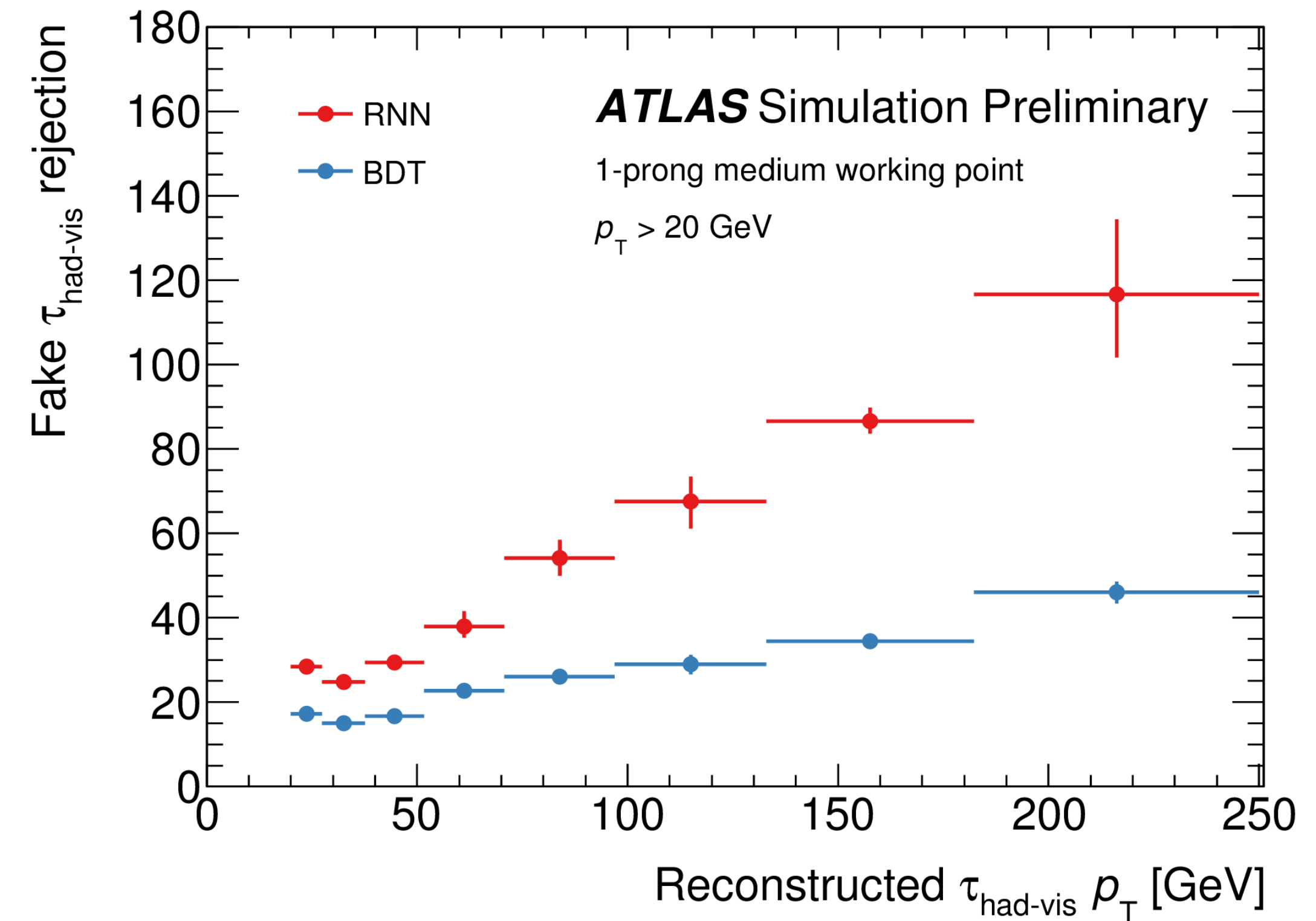


Improved detectors & analysis techniques

Insertable b-layer (IBL) and flavour tagging



(Recursive) neural networks for identification of hadronic tau decays



Improved theory predictions

higher-order calculations (ggF: NNLO -> N3LO)

NNLO+NNLL					
Cross Section (pb)	+QCD Scale %	-QCD Scale %	$\pm(\text{PDF}+\alpha_s)$ %	$\pm\text{PDF}$ %	$\pm\alpha_s$ %
4.414E+01	+7.6	-8.1	± 3.1	± 1.8	± 2.5

↓

N3LO						
Cross Section (pb)	+Theory %	-Theory %	TH Gaussian %	$\pm(\text{PDF}+\alpha_s)$ %	$\pm\text{PDF}$ %	$\pm\alpha_s$ %
4.858E+01	+4.6	-6.7	± 3.9	± 3.2	± 1.9	± 2.6

state-of-the-art MC generators (at least NLO; NNLO in ggF)

NNLOPS simulation of Higgs boson production

Keith Hamilton,^{a,b,1} Paolo Nason,^c Emanuele Re^d and Giulia Zanderighi^d

^aDepartment of Physics and Astronomy, University College London, London, WC1E 6BT, U.K.
^bTheory Division, CERN, CH-1211, Geneva 23, Switzerland
^cINFN, Sezione di Milano Bicocca, Piazza della Scienza 3, 20126 Milan, Italy
^dRudolf Peierls Centre for Theoretical Physics, University of Oxford, 1 Keble Road, U.K.

E-mail: keith.hamilton@ucl.ac.uk, paolo.nason@mib.infn.it, e.re@physics.ox.ac.uk, g.zanderighi@physics.ox.ac.uk

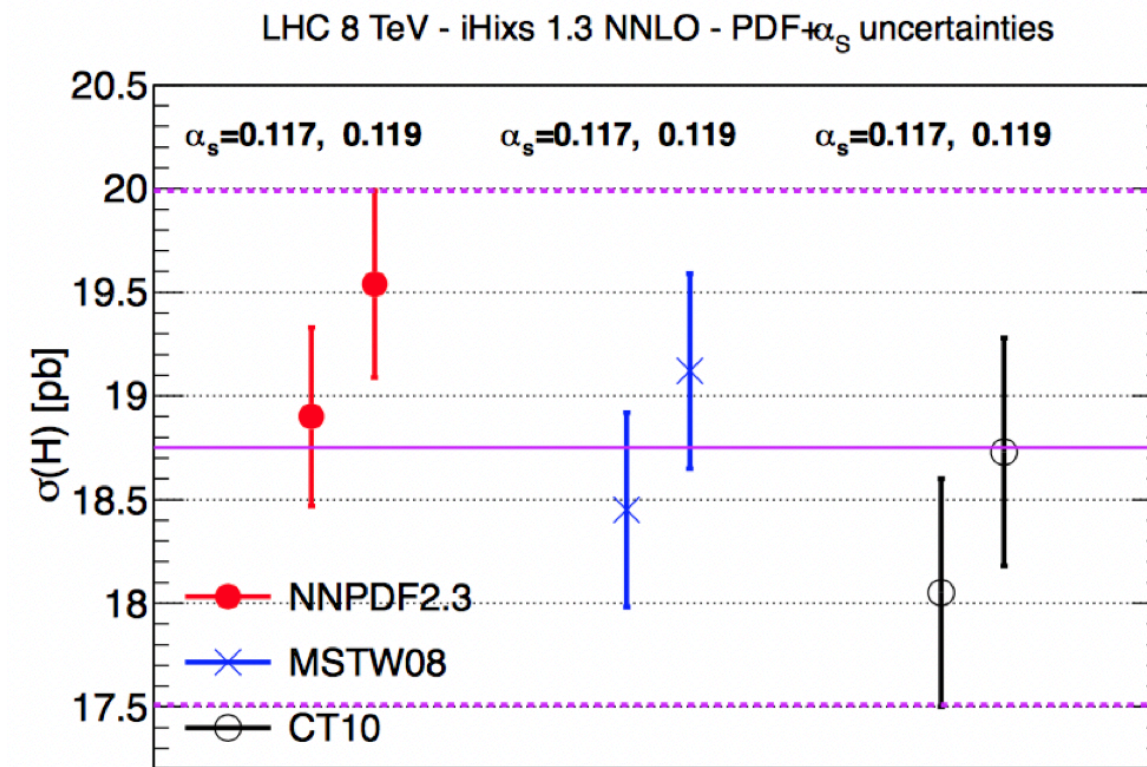
ABSTRACT: We detail a simulation of Higgs boson production via gluon fusion, accurate at next-to-next-to-leading order in the strong coupling, including matching to a parton

JHEP10(2013)222

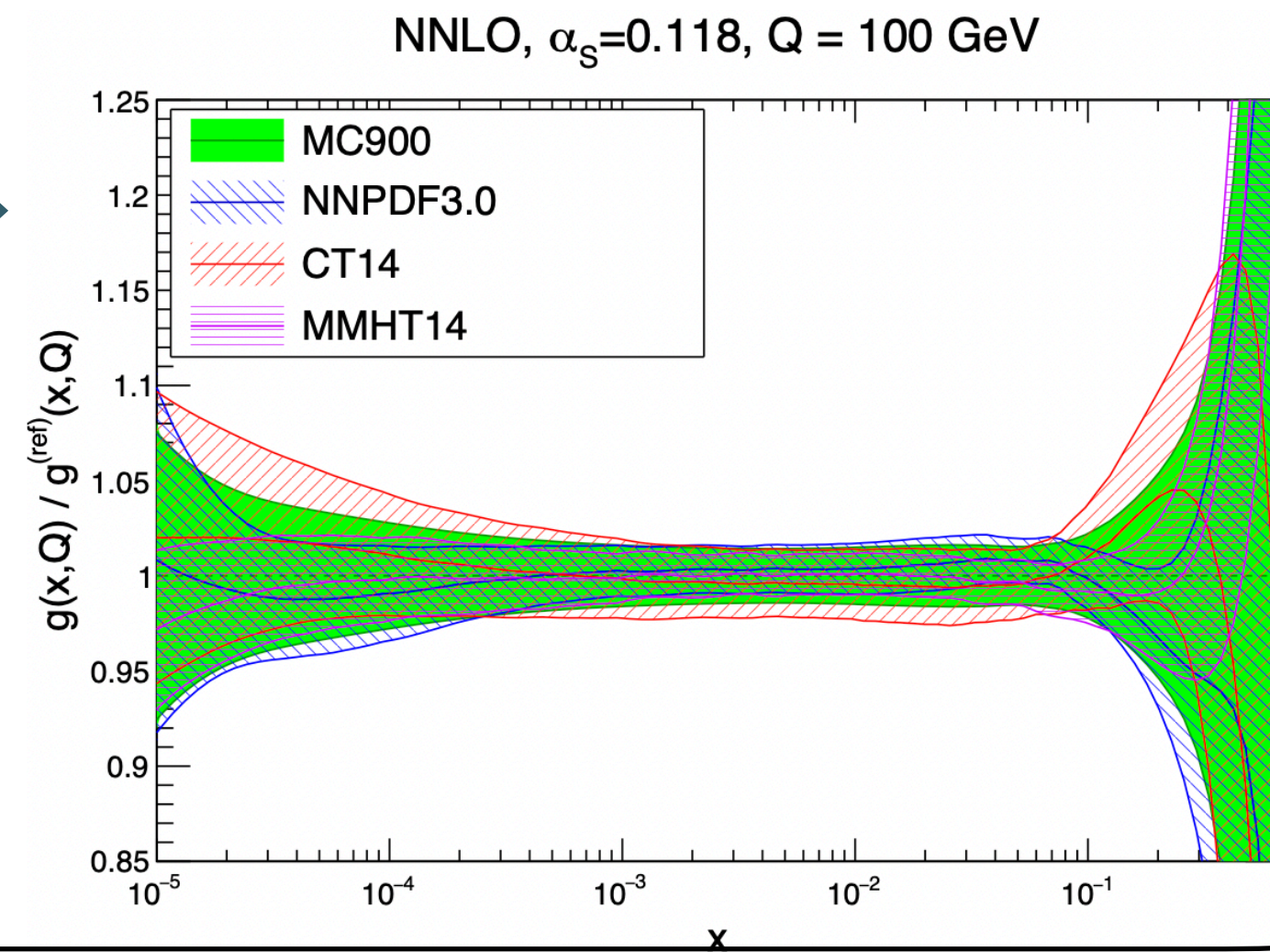
improved PDF sets (including data from LHC Run1)

THE OLD PDF4LHC RECOMMENDATION

HIGGS IN GLUON FUSION



Using the full envelope was conservative but statically not optimal



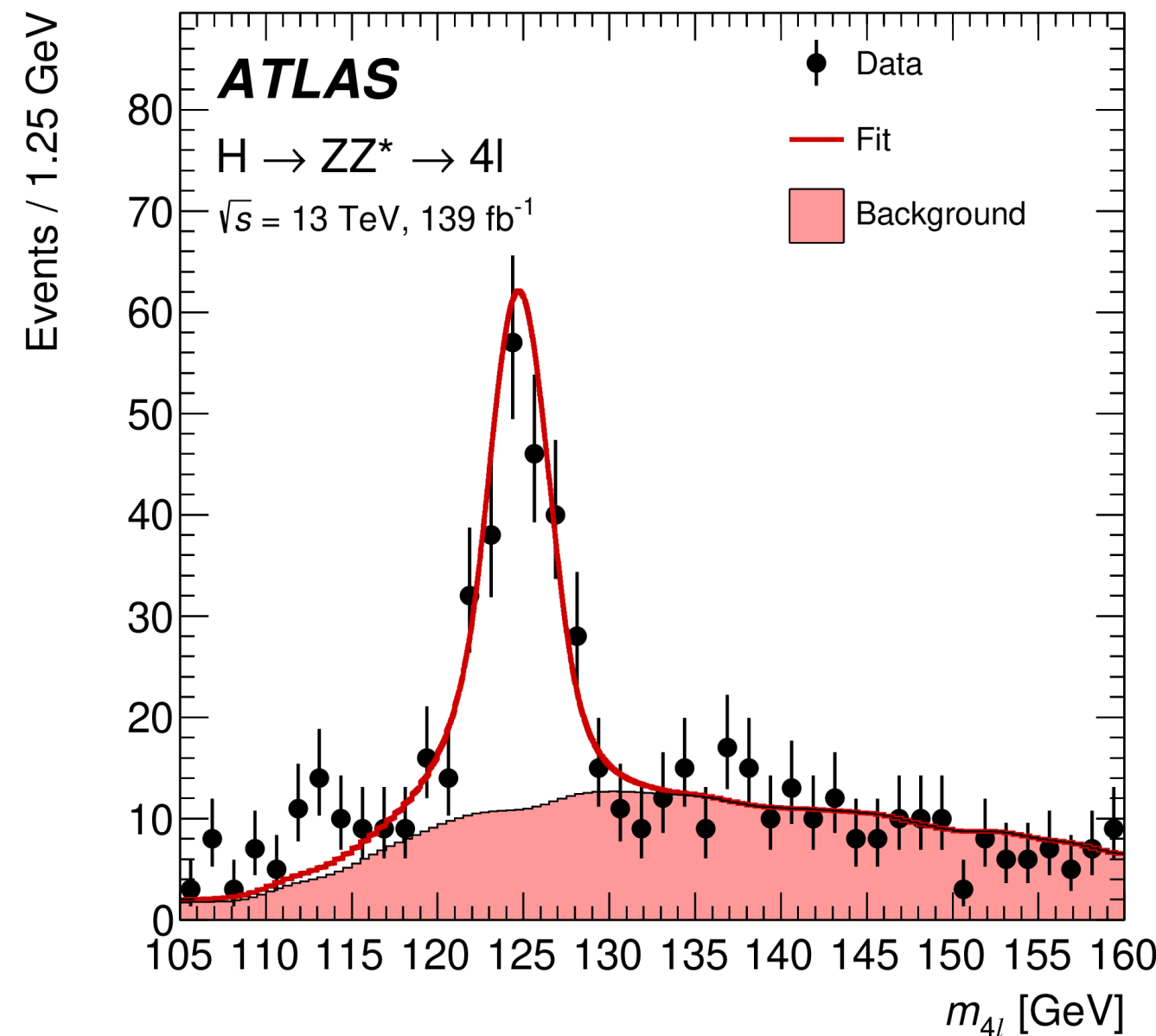
=> smaller modelling uncertainties in measurements, and reduced uncertainties in TH predictions compared to data

Higgs boson mass

- Measured in high-resolution channels ($\gamma\gamma$, $ZZ \rightarrow 4l$) from the position of the invariant mass peak
- Precision limited by statistical (ZZ) or experimental ($\gamma\gamma$) systematic uncertainties
- Requires precise lepton and photon energy calibration (use control samples such as $Z \rightarrow ee, \mu\mu$)

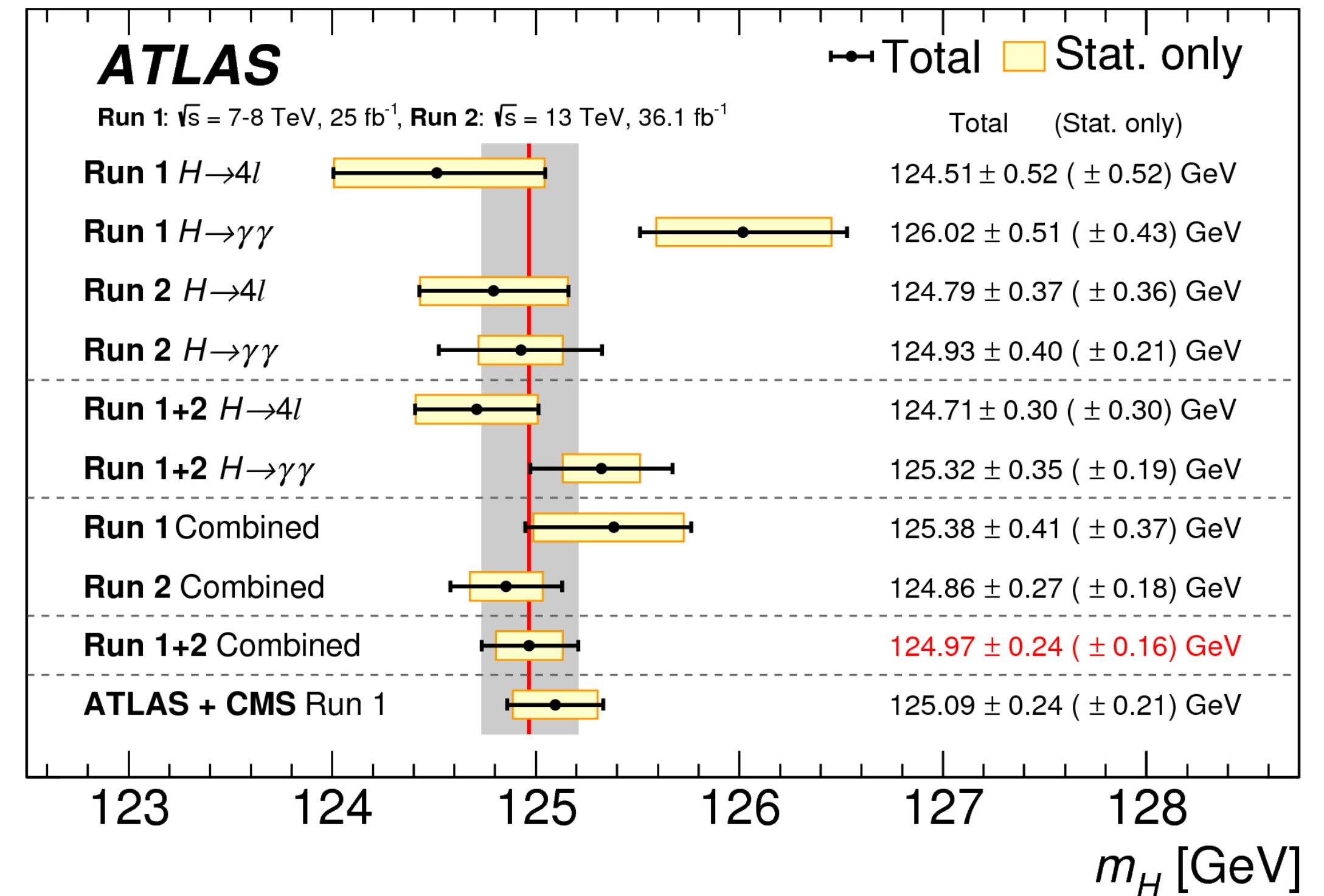
$H \rightarrow ZZ^* \rightarrow 4l$, 139/fb

[arXiv:2207.00320](https://arxiv.org/abs/2207.00320), submitted to PLB



$H \rightarrow ZZ^* \rightarrow 4l + H \rightarrow \gamma\gamma$, 36/fb

[Phys. Lett. B 784 \(2018\) 345](https://arxiv.org/abs/1708.01712)



$$m_H = 124.94 \pm 0.17(\text{stat.}) \pm 0.03(\text{syst.}) \text{ GeV}$$

- Run1 + full Run2: 0.14% precision

- Run1 + partial Run2: 0.19% precision

One of the most precisely measured electroweak parameters!

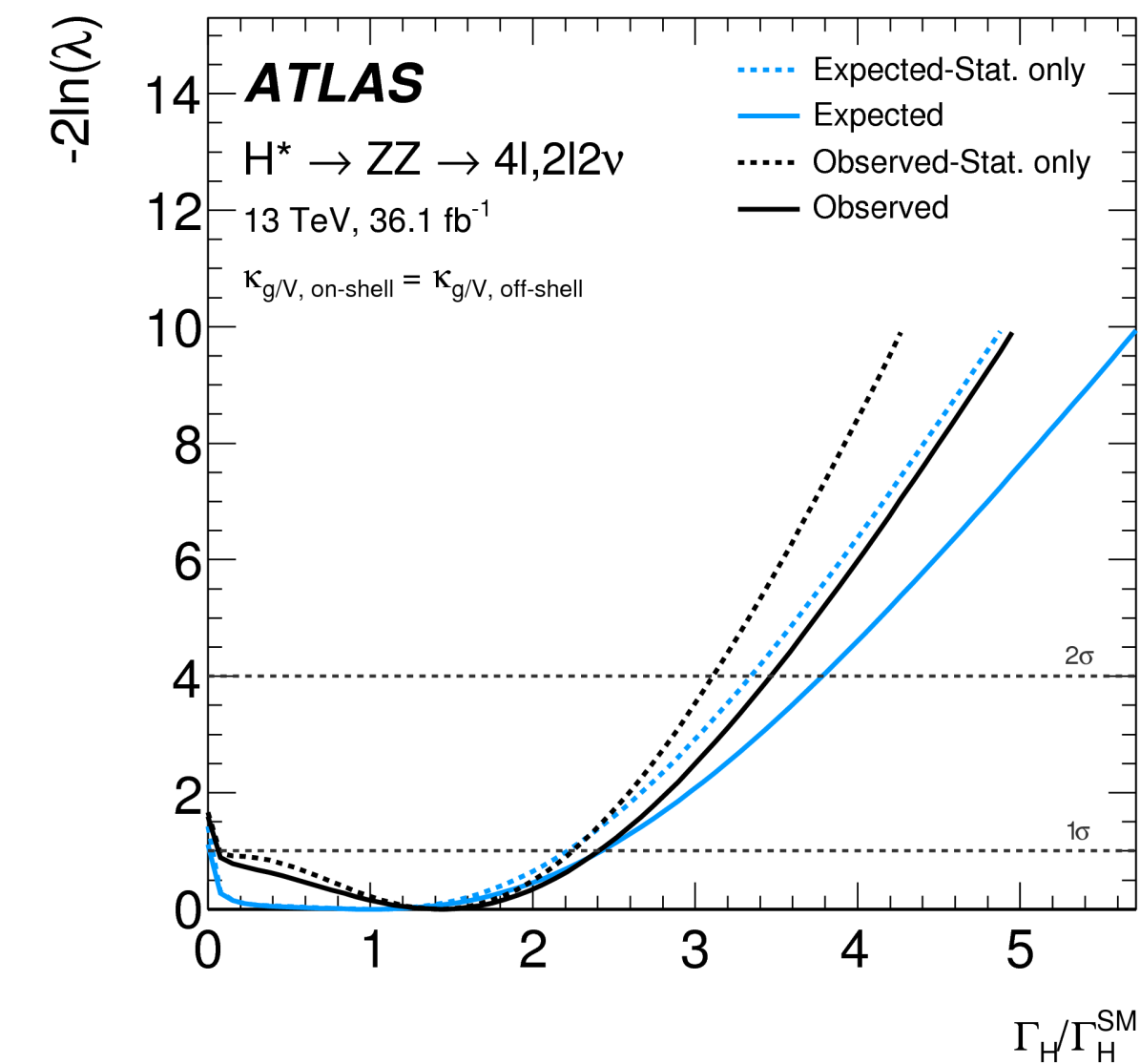
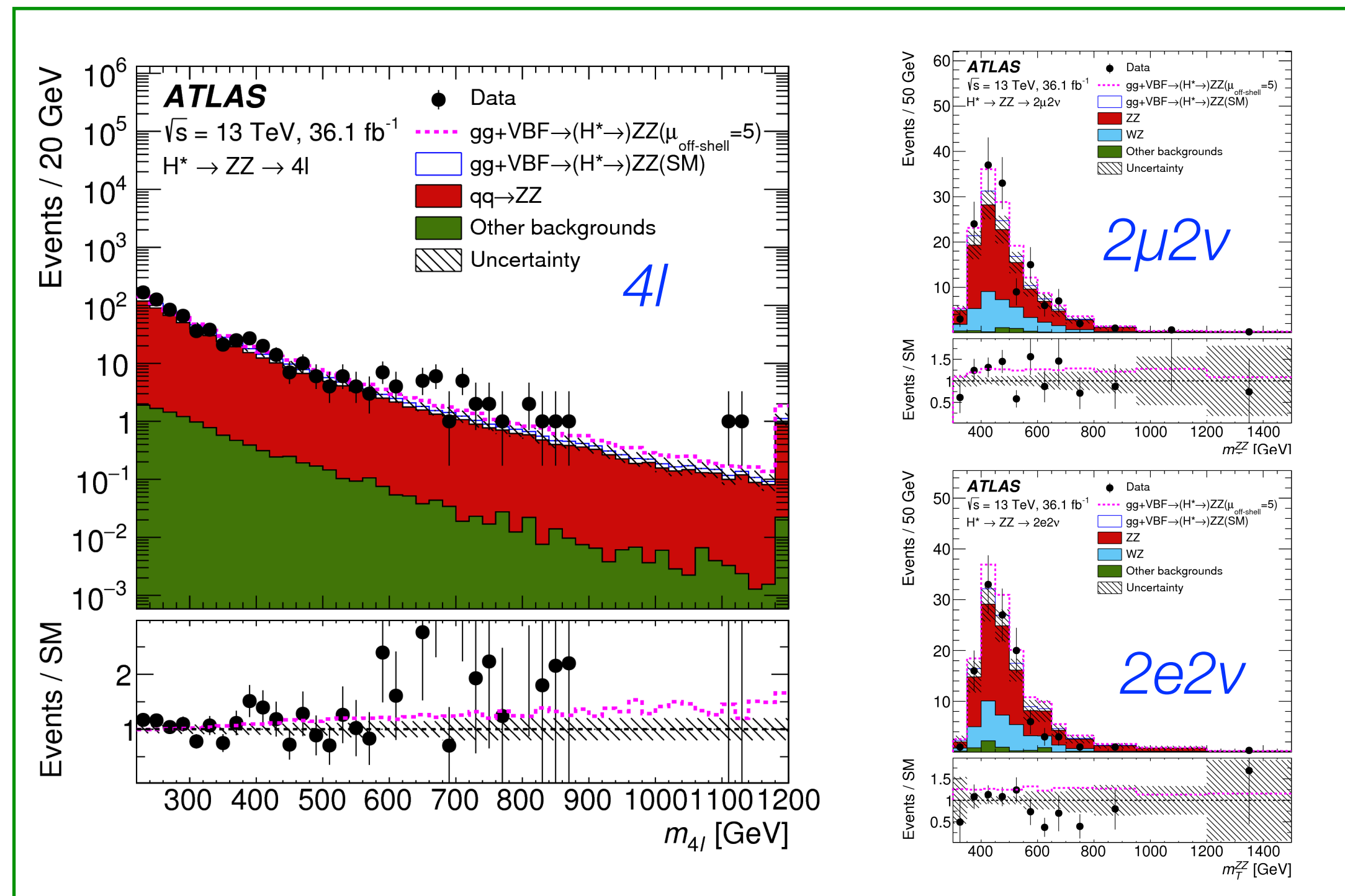
Higgs boson width

- SM Higgs boson width (**4.1 MeV**) \ll experimental resolution (**1–2 GeV**) \Rightarrow too small to be measured directly (direct limits ~ 1 GeV)
- Can be inferred from ratio of off-shell/on-shell $pp \rightarrow H^* \rightarrow ZZ$ (or WW) xsections

$$\mu_{\text{on-shell}} = \frac{\sigma_{\text{on-shell}}^{gg \rightarrow H \rightarrow ZZ^*}}{\sigma_{\text{on-shell,SM}}^{gg \rightarrow H \rightarrow ZZ^*}} = \frac{\kappa_{g,\text{on-shell}}^2 \cdot \kappa_{Z,\text{on-shell}}^2}{\Gamma_H / \Gamma_H^{\text{SM}}}, \quad \mu_{\text{off-shell}} = \frac{\sigma_{\text{off-shell}}^{gg \rightarrow H^* \rightarrow ZZ}}{\sigma_{\text{off-shell,SM}}^{gg \rightarrow H^* \rightarrow ZZ}} = \kappa_{g,\text{off-shell}}^2 \cdot \kappa_{Z,\text{off-shell}}^2,$$

with some assumptions:

- running of the couplings as in the SM: $\kappa_{\text{off-shell}} = \kappa_{\text{on-shell}}$
- no new signals in the search region, apart from a possibly enhanced off-shell Higgs contribution



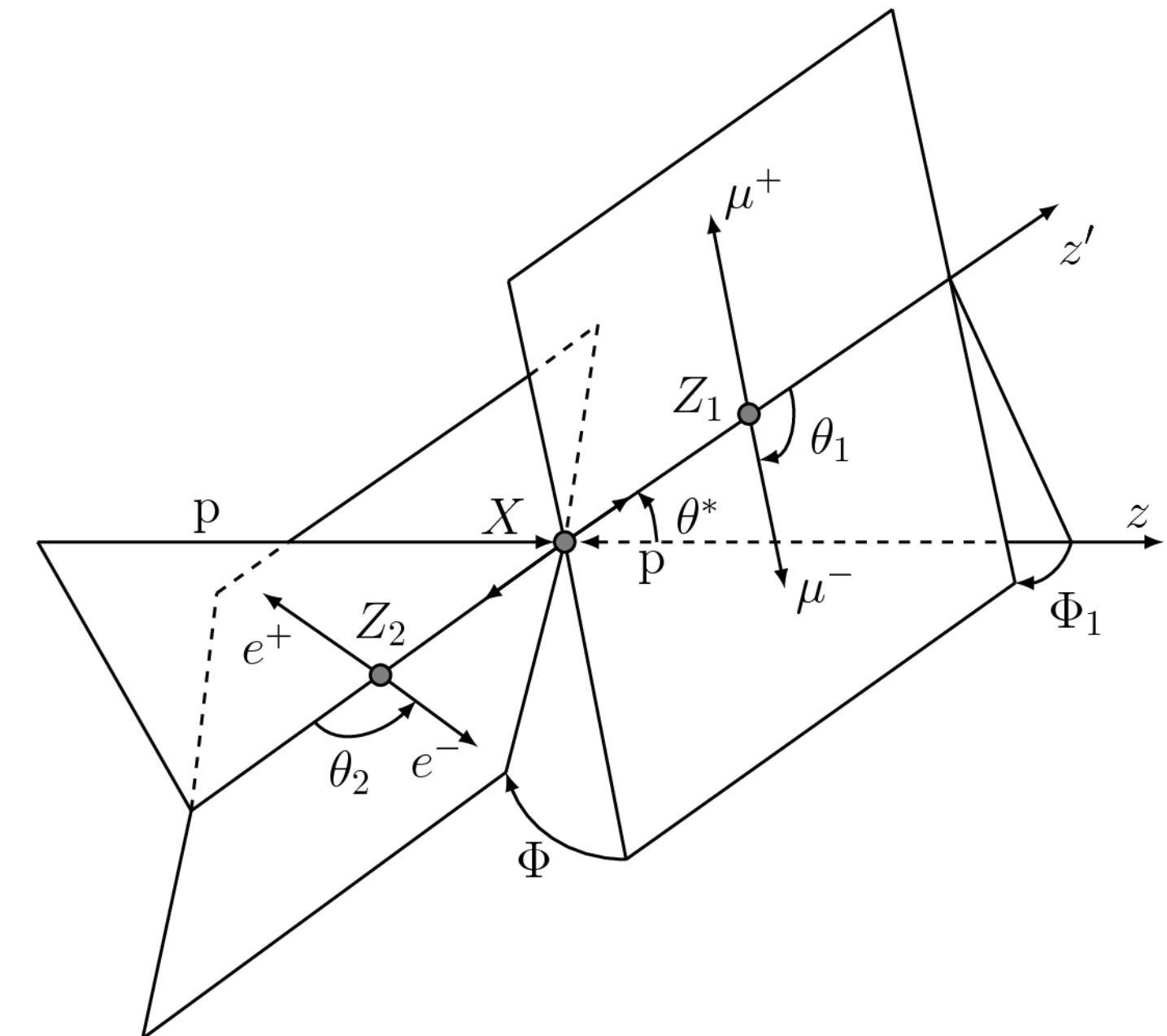
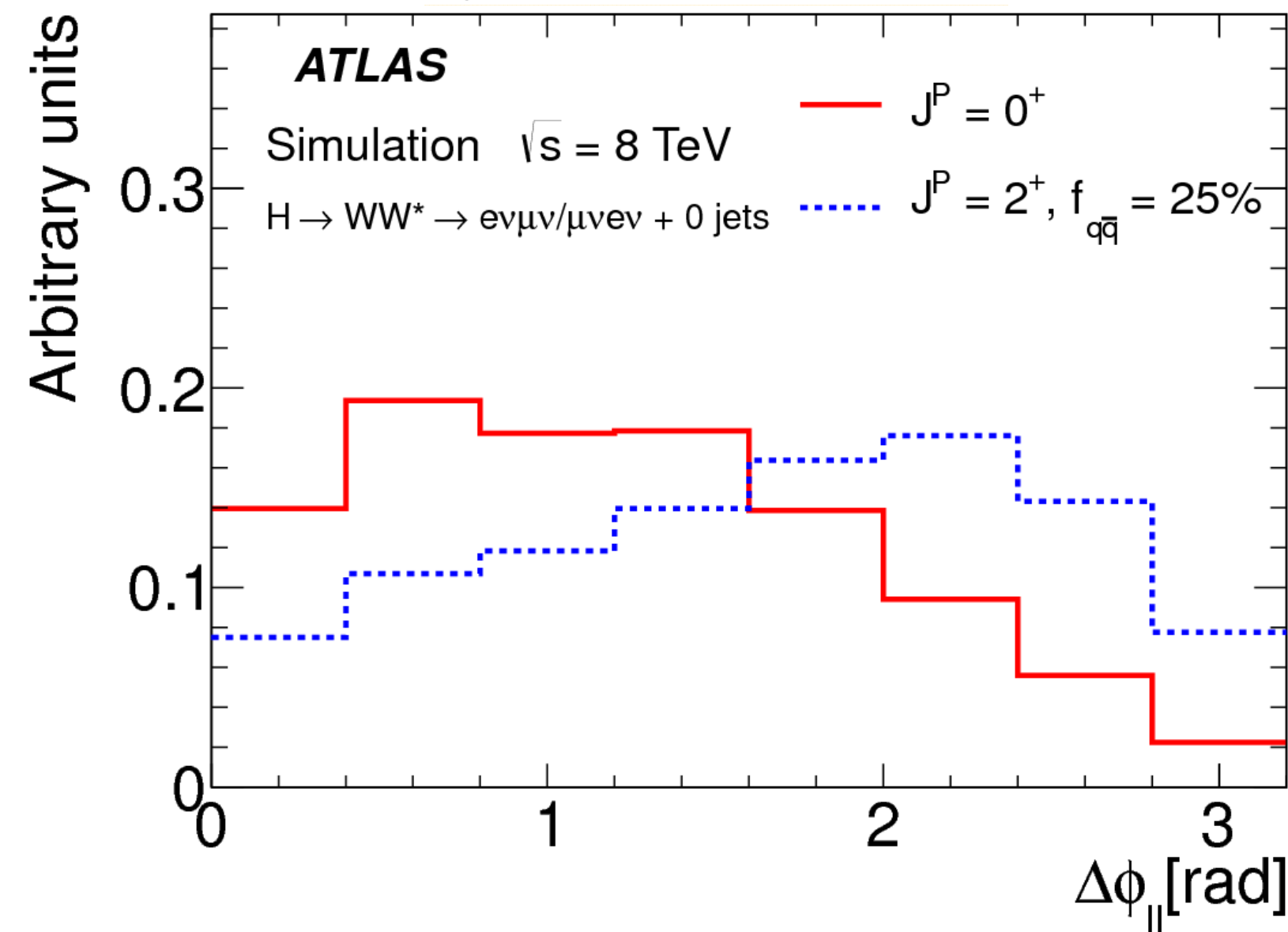
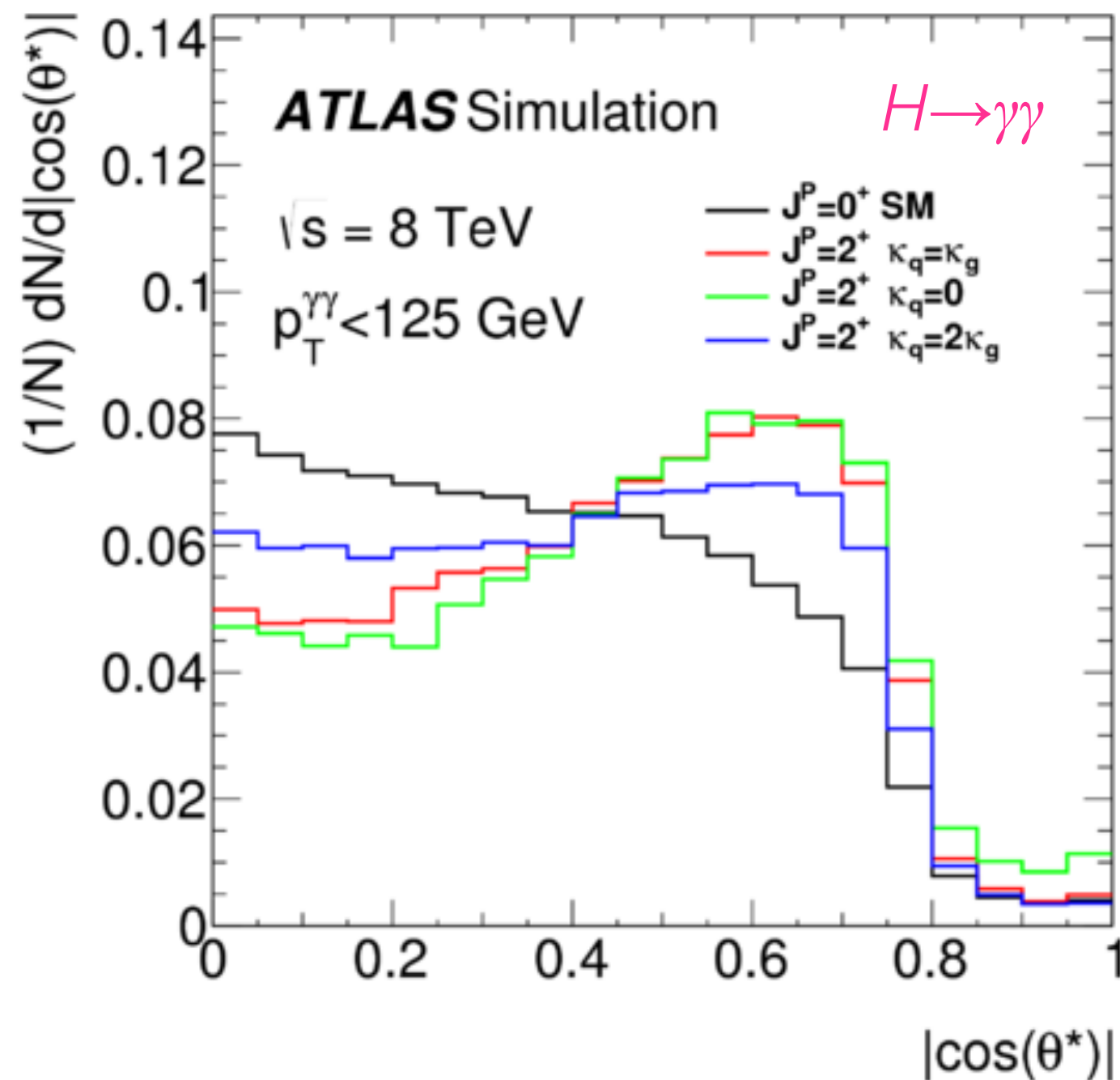
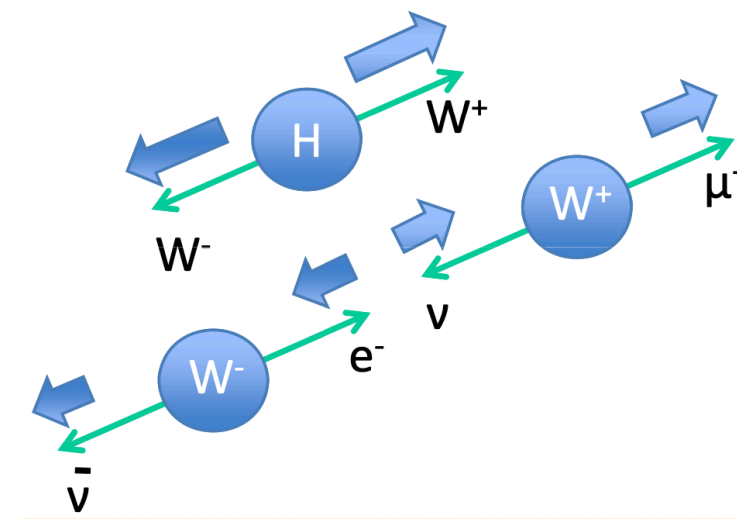
[Phys. Lett. B 786 \(2018\) 223](#)

95% CL upper limits		Observed	Expected
$\mu_{\text{off-shell}}$	$ZZ \rightarrow 4\ell$ analysis	4.5	4.3
	$ZZ \rightarrow 2\ell 2\nu$ analysis	5.3	4.4
	Combined	3.8	3.4
$\Gamma_H / \Gamma_H^{\text{SM}}$	Combined	3.5	3.7

25% precision at HL-LHC... but model-dependent!

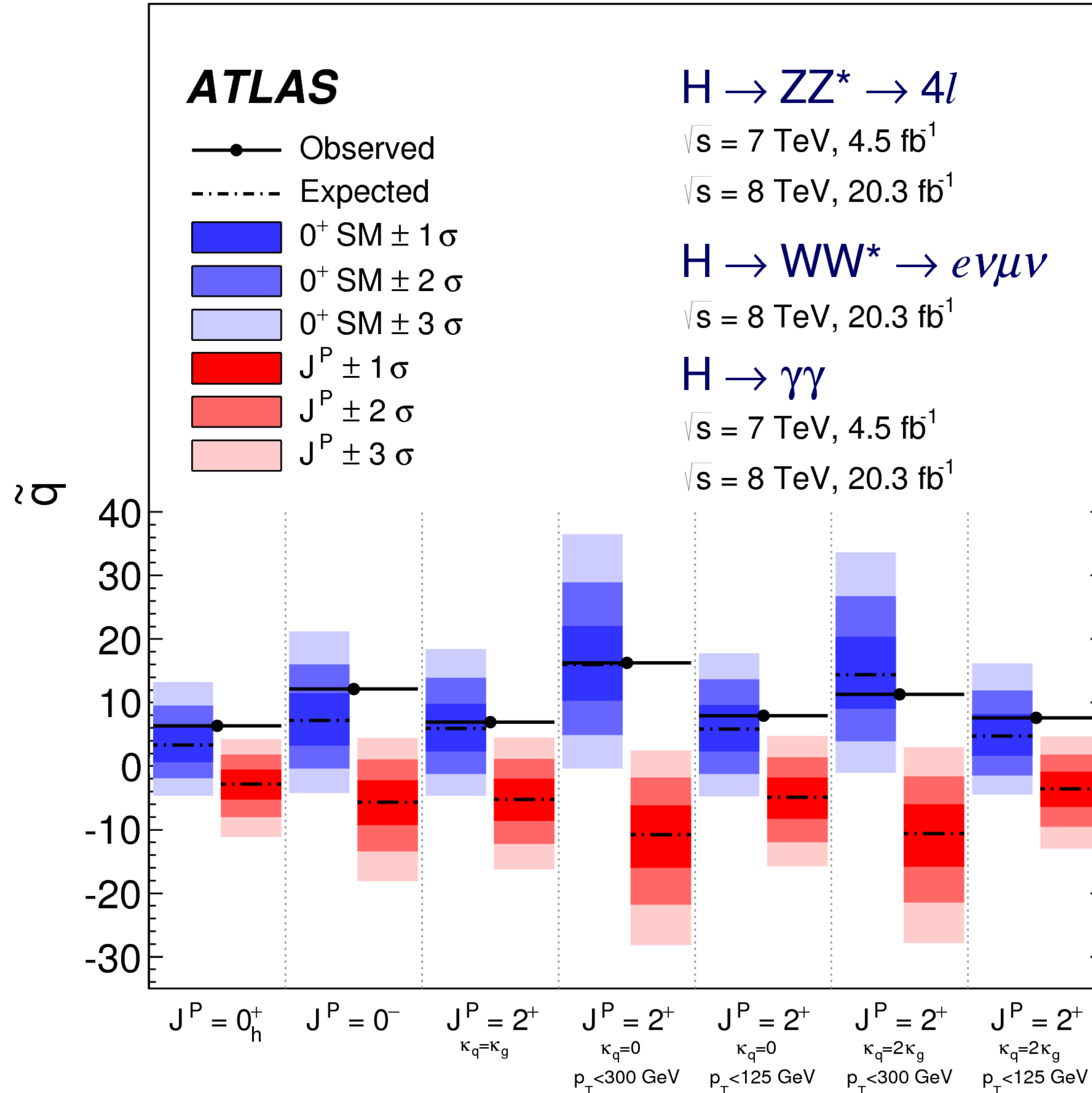
Higgs boson spin and parity

- **Spin 1 forbidden** by observation of $H \rightarrow \gamma\gamma$ decay (Landau-Yang's theorem)
- **$J^P = 0^-, 0^+$ non-SM** (different tensor structure of the HVV couplings), and various graviton-like **2^+** scenarios are **tested** one-by-one against the SM 0^+ hypothesis exploiting **angular distributions that are sensitive to J^P**
 - Polar angle θ^* of photon in $\gamma\gamma$ CM frame (flat for spin-0, quadratic in $\cos\theta^*$ for spin 2)
 - Azimuthal opening angle between two leptons in $H \rightarrow WW$ (small for spin-0, large for spin-2 due to W coupling to left-handed fermions)
 - Decay angles and dilepton invariant masses in $H \rightarrow ZZ \rightarrow 4l$



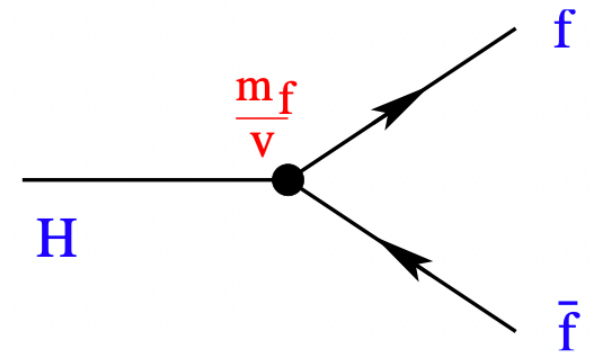
Higgs boson spin and parity

[Eur. Phys. J. C75 \(2015\) 476](#)

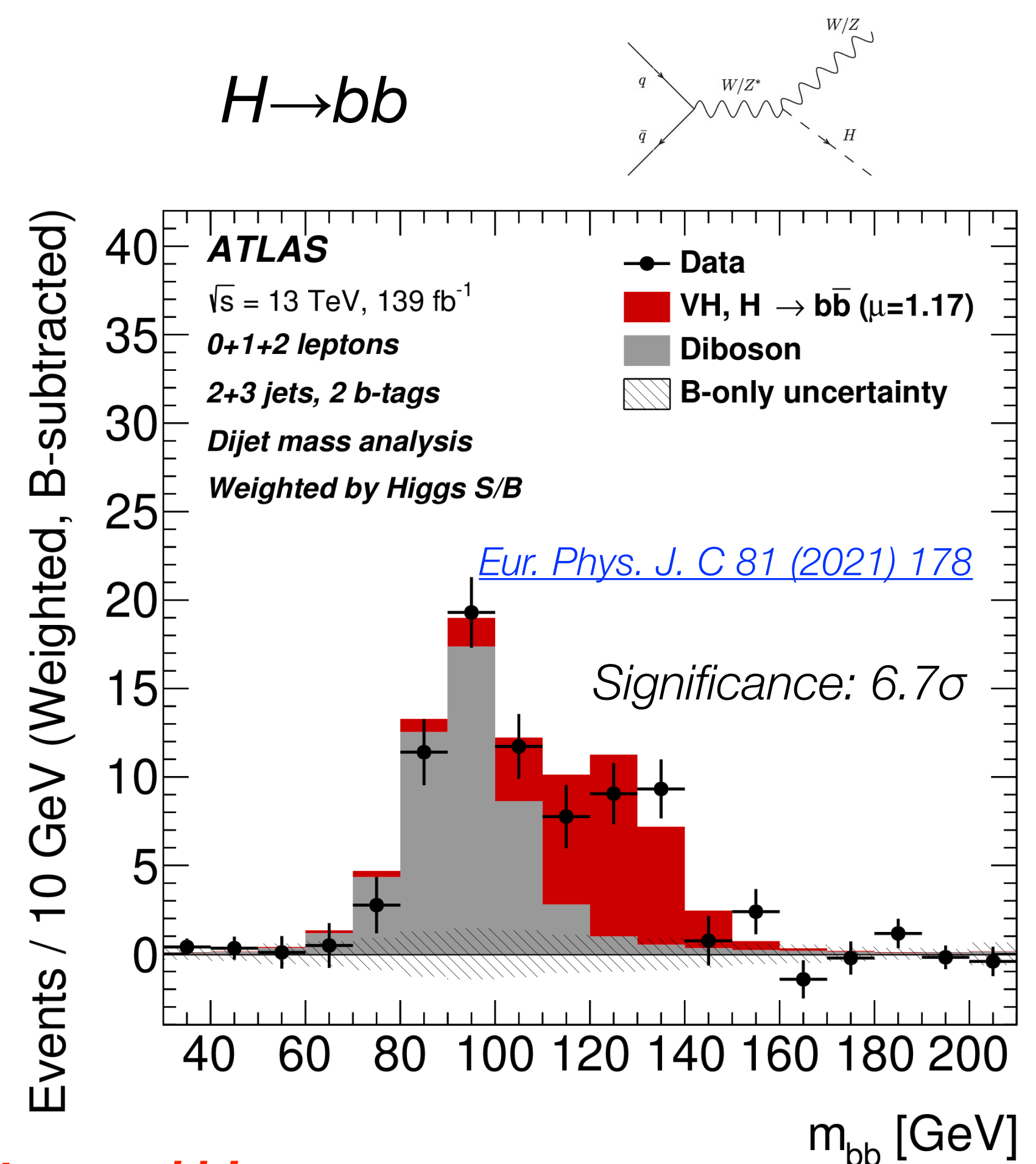
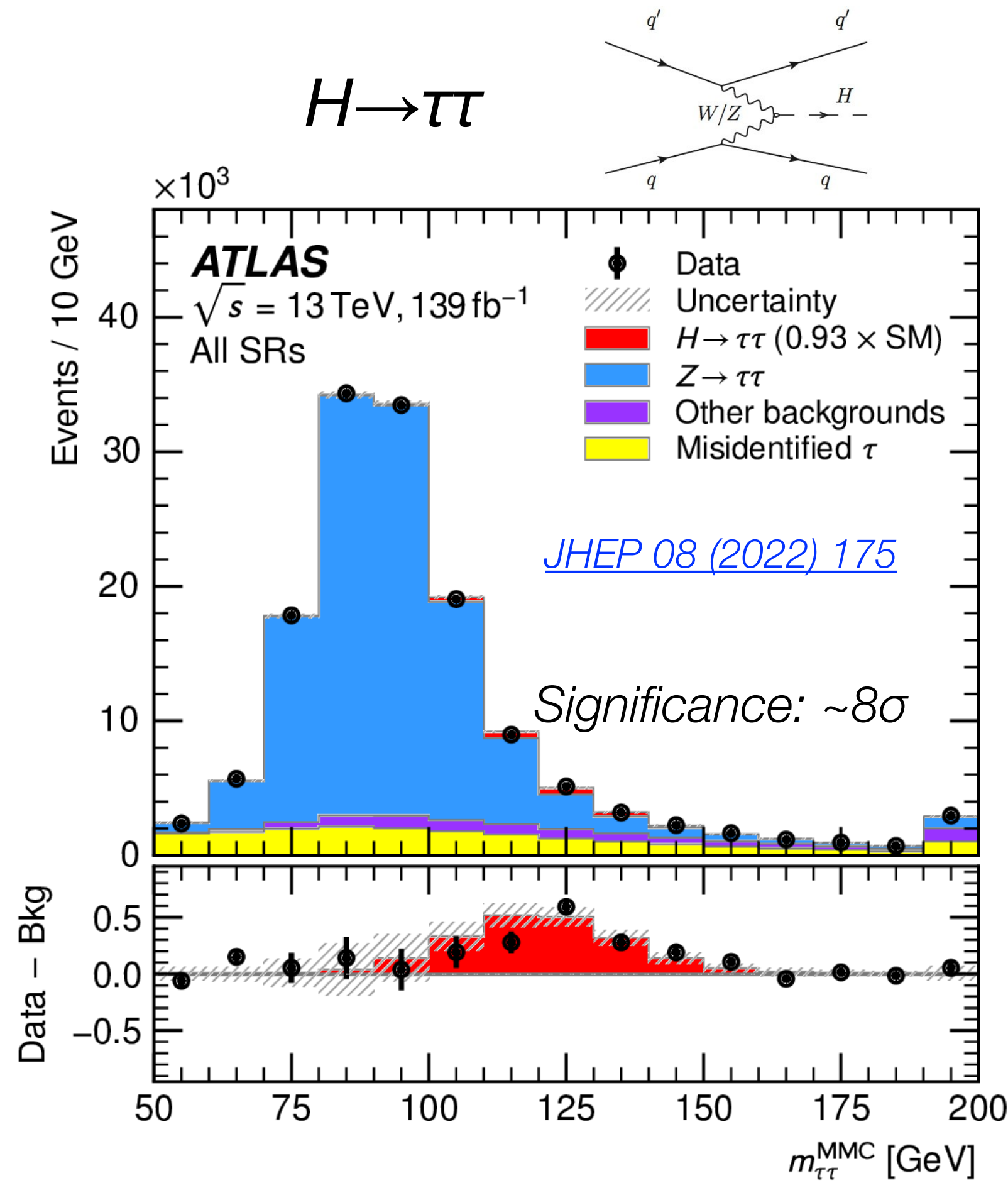


All alternative hypotheses disfavoured at $> 3\sigma$

Observation of Higgs boson decays to τ -leptons and to b-quarks



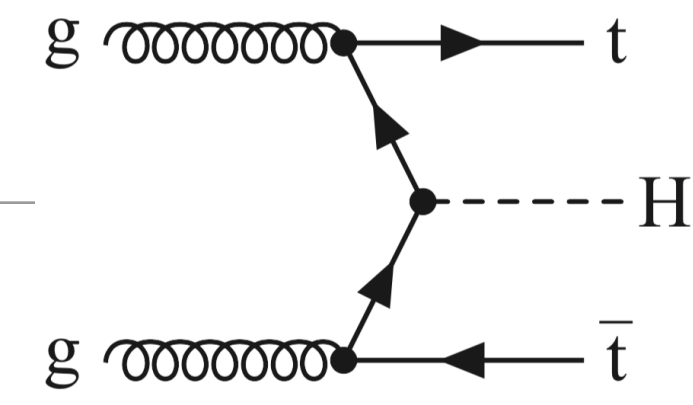
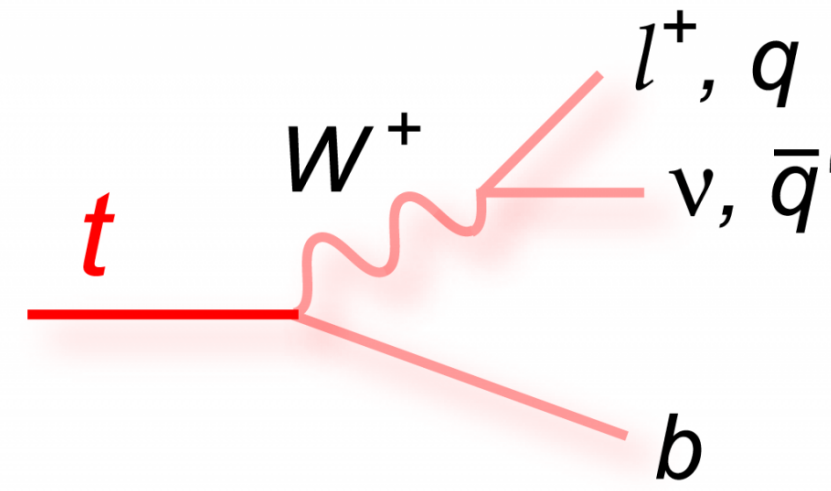
- First observation with partial Run2 datasets. More detailed/granular studies performed with full dataset
- Best sensitivity provided by production modes with lower x-section but much better bkg rejection than gluon fusion
 - VBF ($\sim 8\%$ of σ_H) for $H \rightarrow \tau\tau$, $V(\rightarrow \text{leptons})H$ ($\sim 0.9\%$ of σ_H) for $H \rightarrow bb$
- Large sensitivity boost from use of multivariate techniques for object reconstruction and S/B discrimination in Run2 analyses



Evidence of Yukawa couplings to τ and b!

Observation of Higgs boson production with ttbar pairs

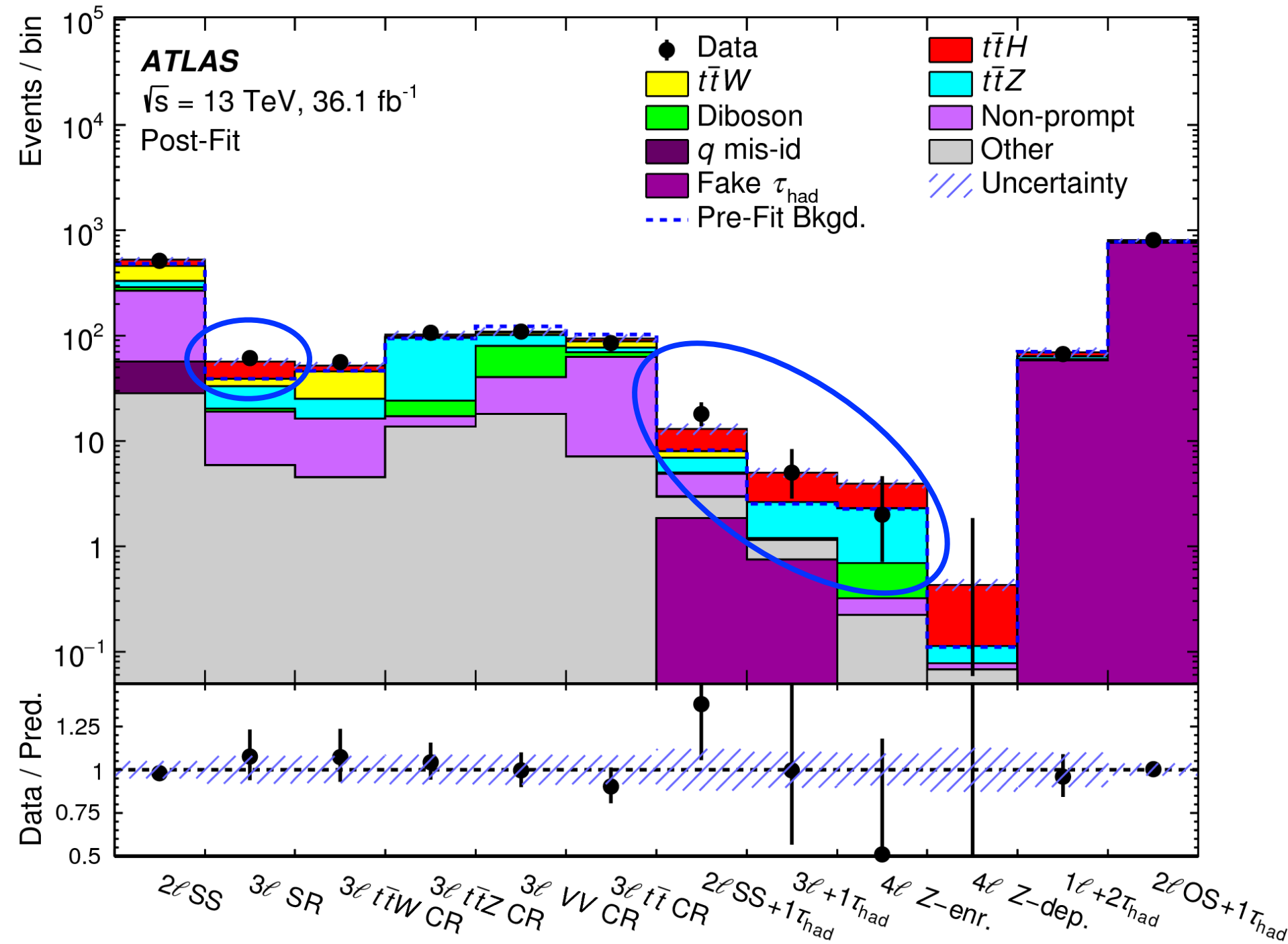
- ttbar pair identified by presence of b-jets, large jet multiplicity, possibly leptons and missing momentum
- Best sensitivity provided by decay modes with leptons (WW, $\tau\tau$) or photons



Analysis	Integrated luminosity [fb ⁻¹]	Expected significance	Observed significance
$H \rightarrow \gamma\gamma$	79.8	3.7 σ	4.1 σ
$H \rightarrow$ multilepton	36.1	2.8 σ	4.1 σ
$H \rightarrow b\bar{b}$	36.1	1.6 σ	1.4 σ
$H \rightarrow ZZ^* \rightarrow 4\ell$	79.8	1.2 σ	0 σ
Combined (13 TeV)	36.1–79.8	4.9 σ	5.8 σ
Combined (7, 8, 13 TeV)	4.5, 20.3, 36.1–79.8	5.1 σ	6.3 σ

Multileptons

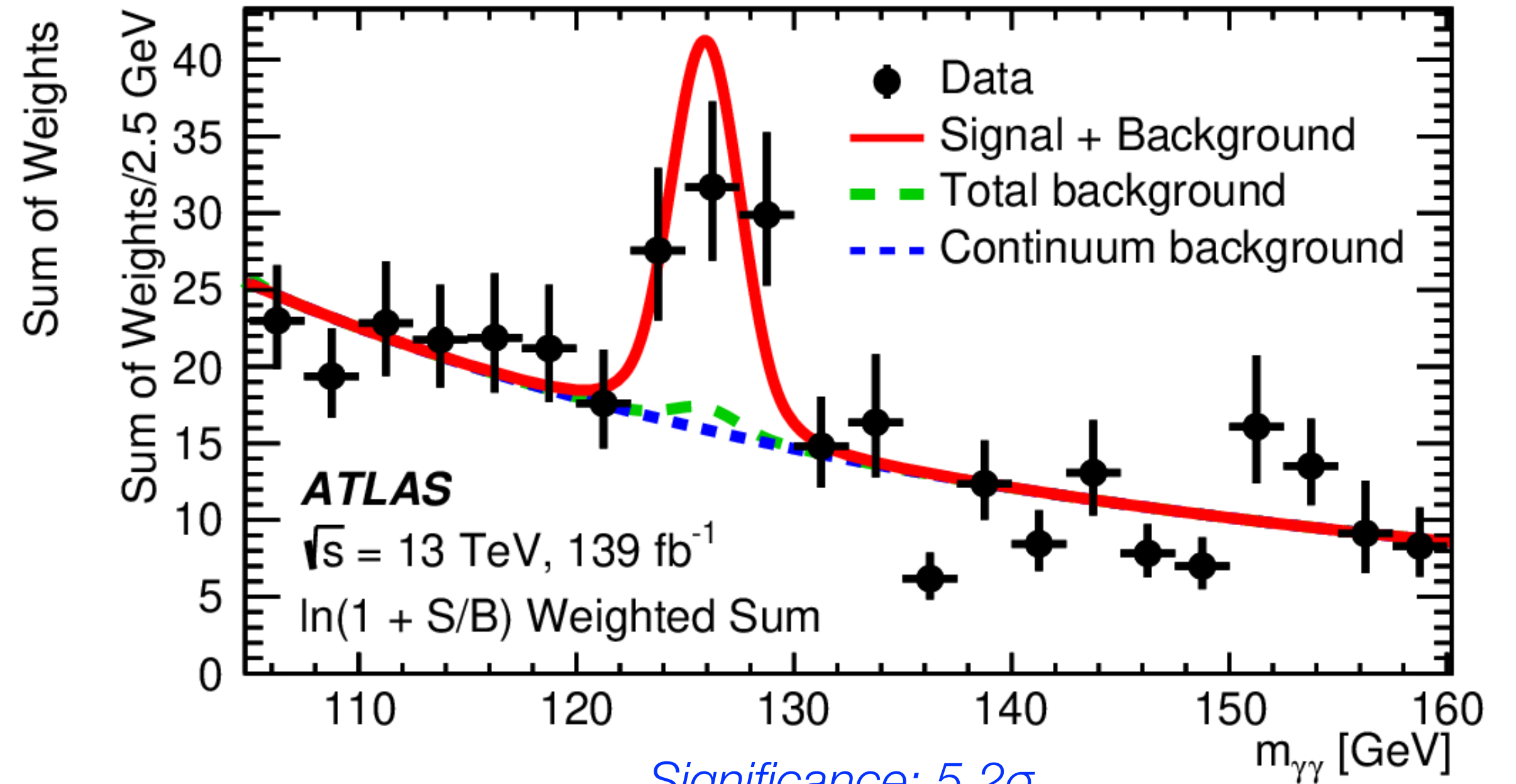
[Phys. Rev. D 97 \(2018\) 072003](#)



Significance: 4.1 σ

$H \rightarrow \gamma\gamma$

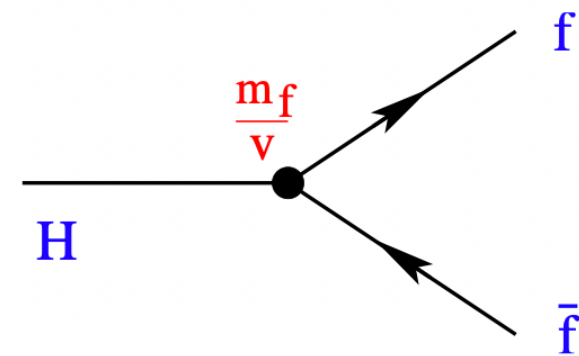
[Phys. Rev. Lett. 125 \(2020\) 061802](#)



Significance: 5.2 σ

Direct evidence of Yukawa couplings to the top quark!

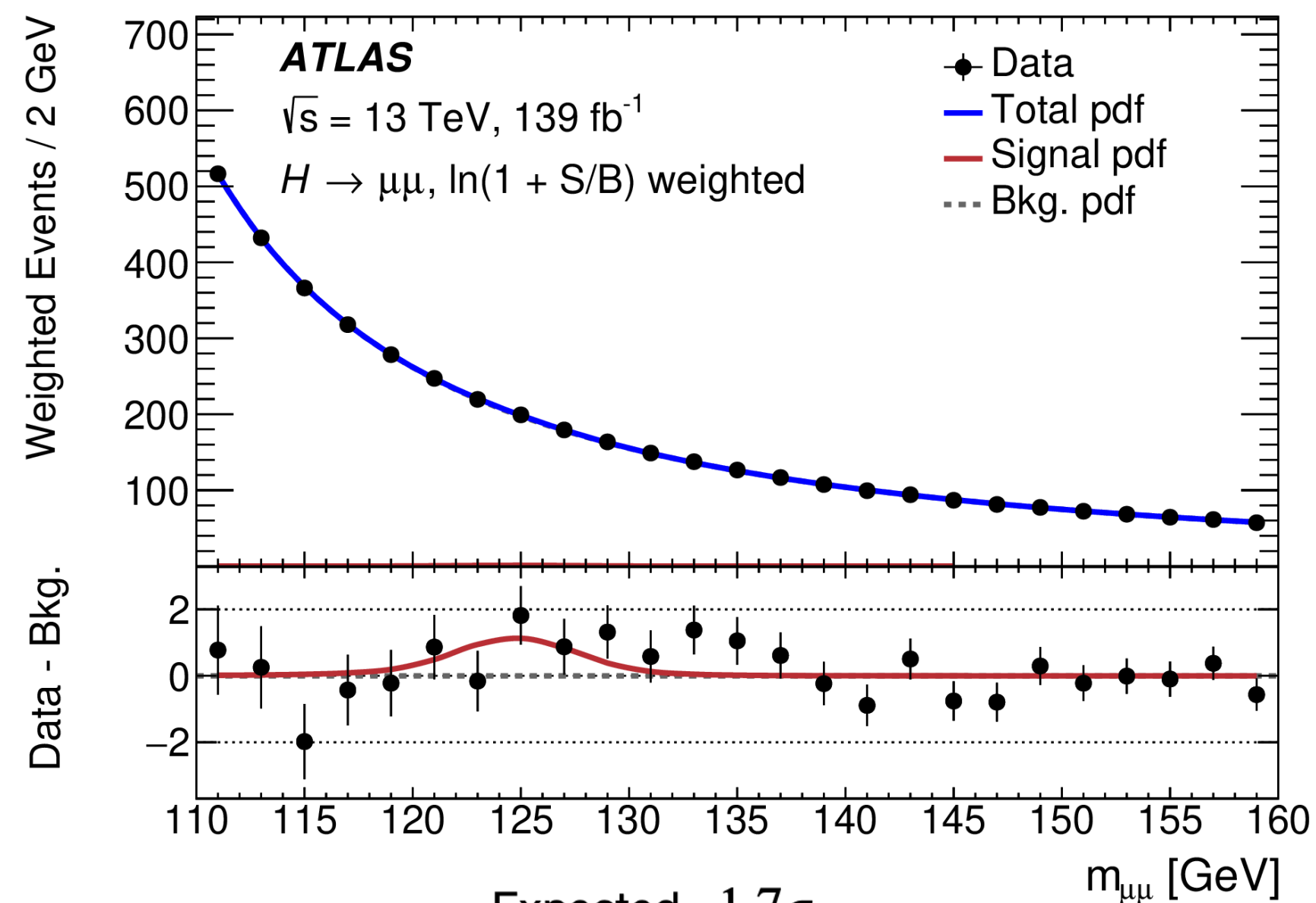
The challenging Yukawa couplings to 2nd-generation fermions



- $H \rightarrow \mu\mu$: very small BR (0.02%), important background from $Z^{(*)} \rightarrow \mu\mu$, but good resolution
- $H \rightarrow cc$: small BR (2%), poor resolution. Large bkg from QCD \Rightarrow search in $V(\rightarrow \text{leptons})H$. Poor c-tagging \Rightarrow bkg from $H \rightarrow bb$

$H \rightarrow \mu\mu$

[Phys. Lett. B 812 \(2021\) 135980](#)



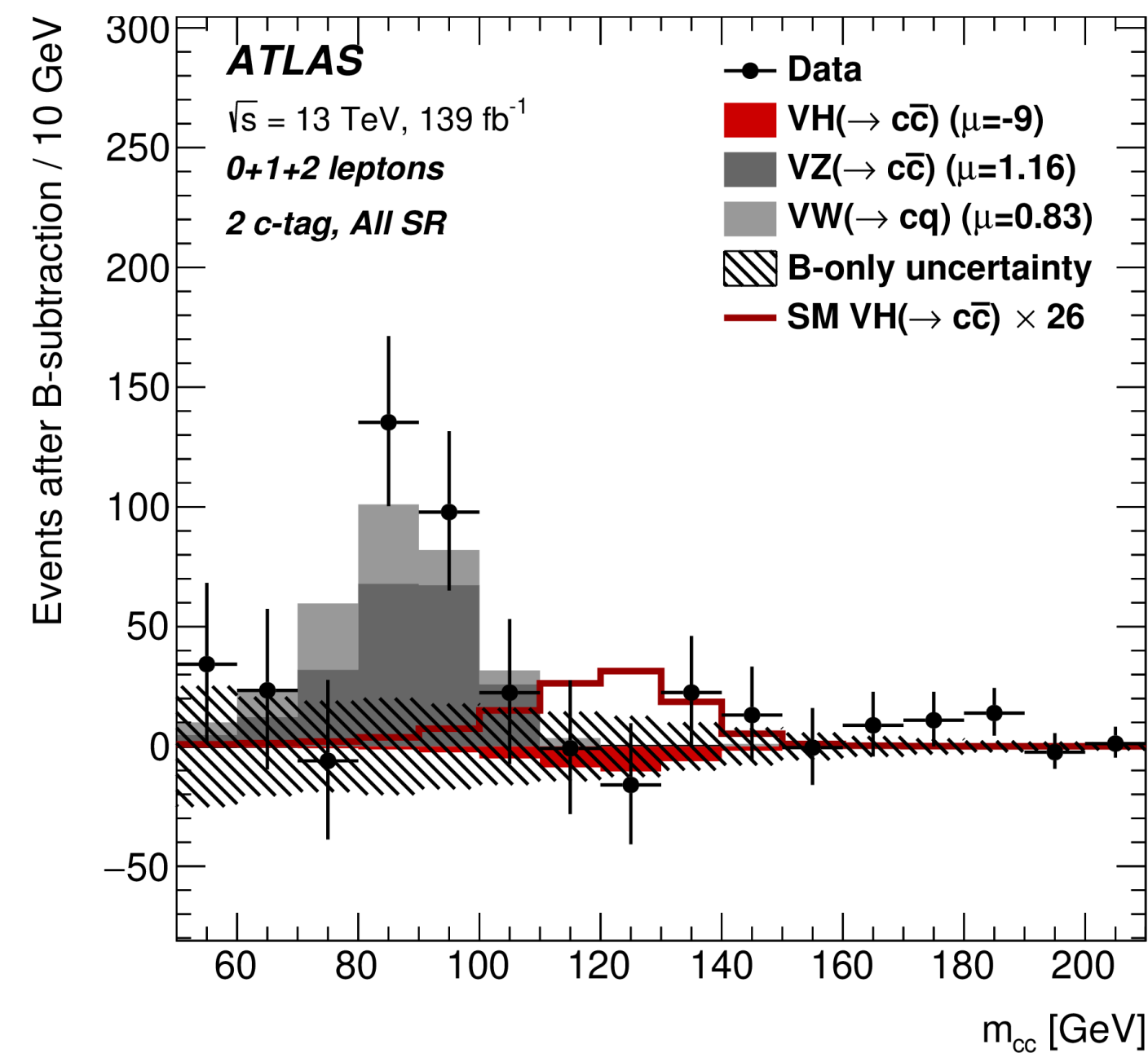
Expected 1.7σ

Observed 2.0σ

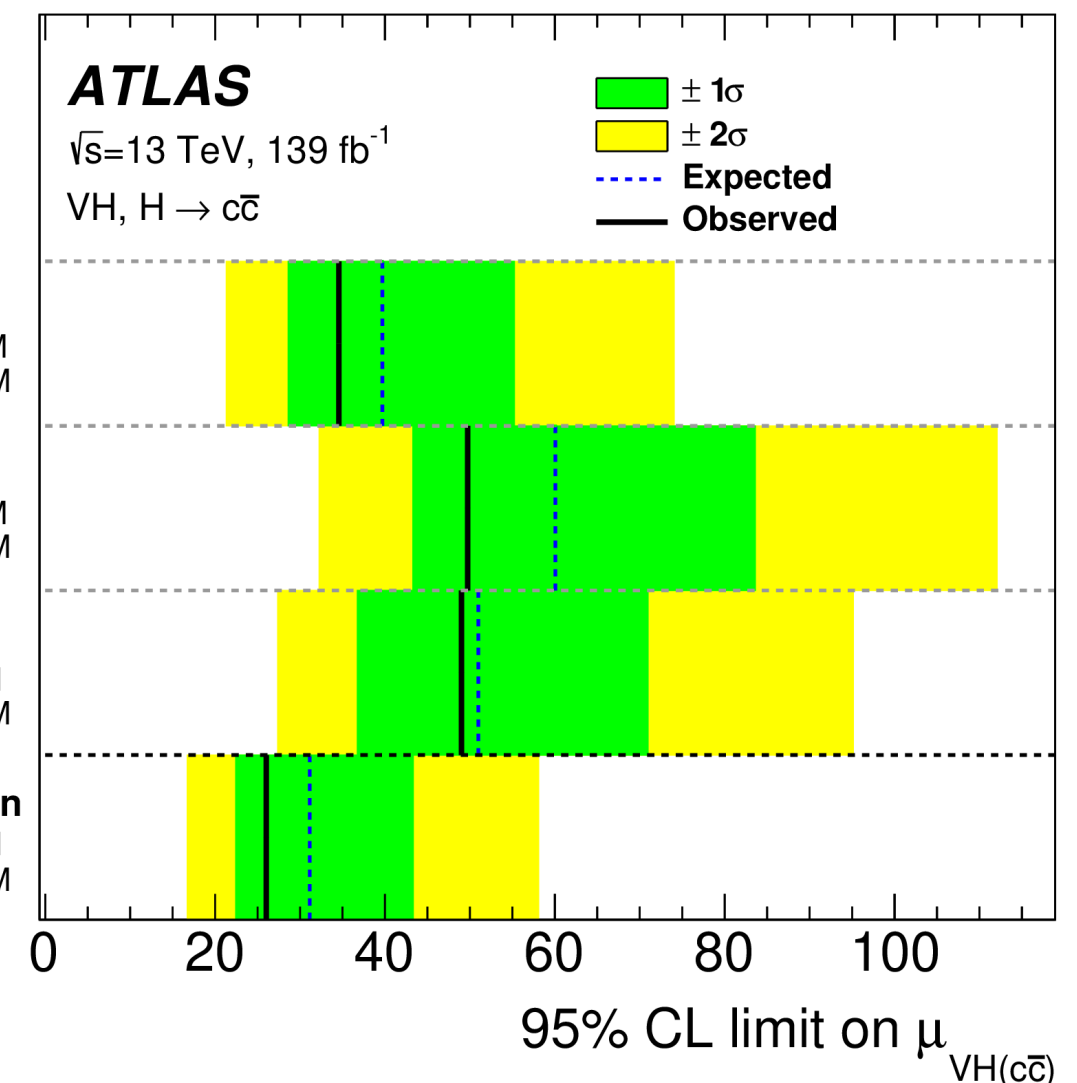
$$\mu = 1.2 \pm 0.6$$

$H \rightarrow cc$

[Eur. Phys. J. C 82 \(2022\) 717](#)



$$|\kappa_c| < 8.5 \text{ (12.4)}$$

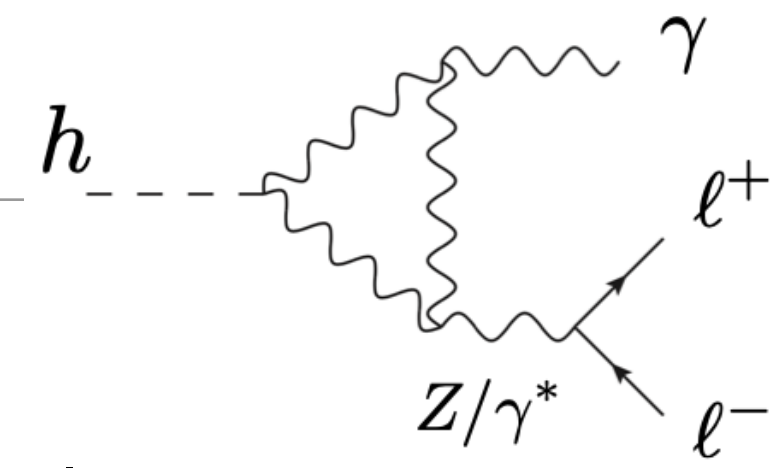


3σ evidence in CMS! Expect observation in Run3

Still a long way before the observation... maybe at HL-LHC?

(Decays to 1st generation fermions (ee) also searched for but no evidence found and UL set at $7 \cdot 10^4$ the SM prediction of $5 \cdot 10^{-9}$)

Rare (resonant and non resonant) decays to $l^+l^-\gamma$



- $H \rightarrow Z\gamma \rightarrow ll\gamma$ ($l=e,\mu$): $\sim 4.4\%$ of $BR(\gamma\gamma)$

- Tensor coupling, not measured yet: $|H^2|W_{\mu\nu}^a W^{\mu\nu a}$
- Large $Z\gamma$ background, low-momentum leptons and photons

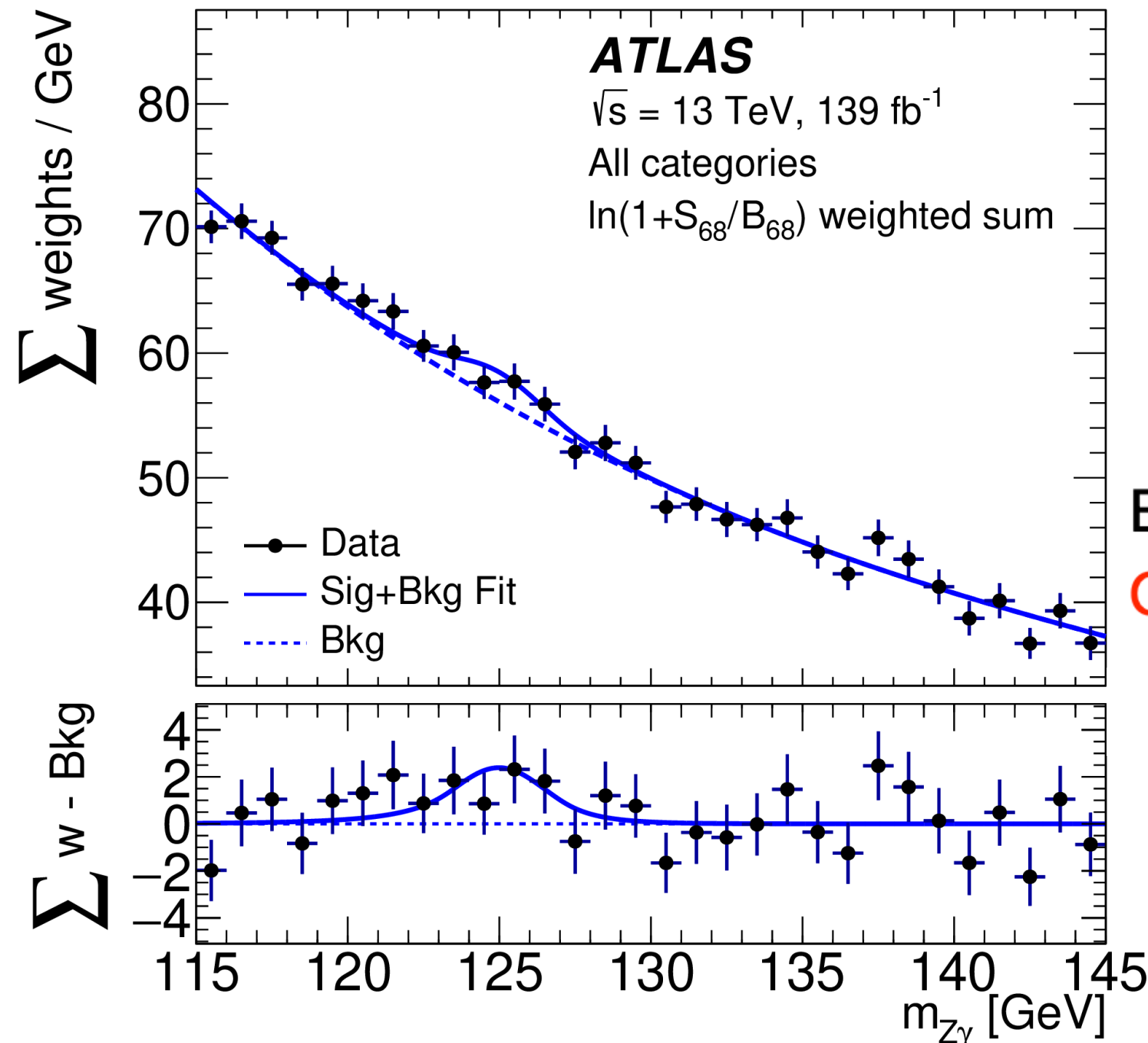
- $H \rightarrow ll\gamma$ ($l=e,\mu$), $m_{ll} < 30$ GeV: $\sim 5\%$ of $BR(\gamma\gamma)$

- Dedicated reconstruction of very close-by electrons (EM showers partially overlapping in the calorimeter)

Potential BSM physics that could explain flavour anomalies could also modify these rates

$H \rightarrow Z\gamma \rightarrow ll\gamma$

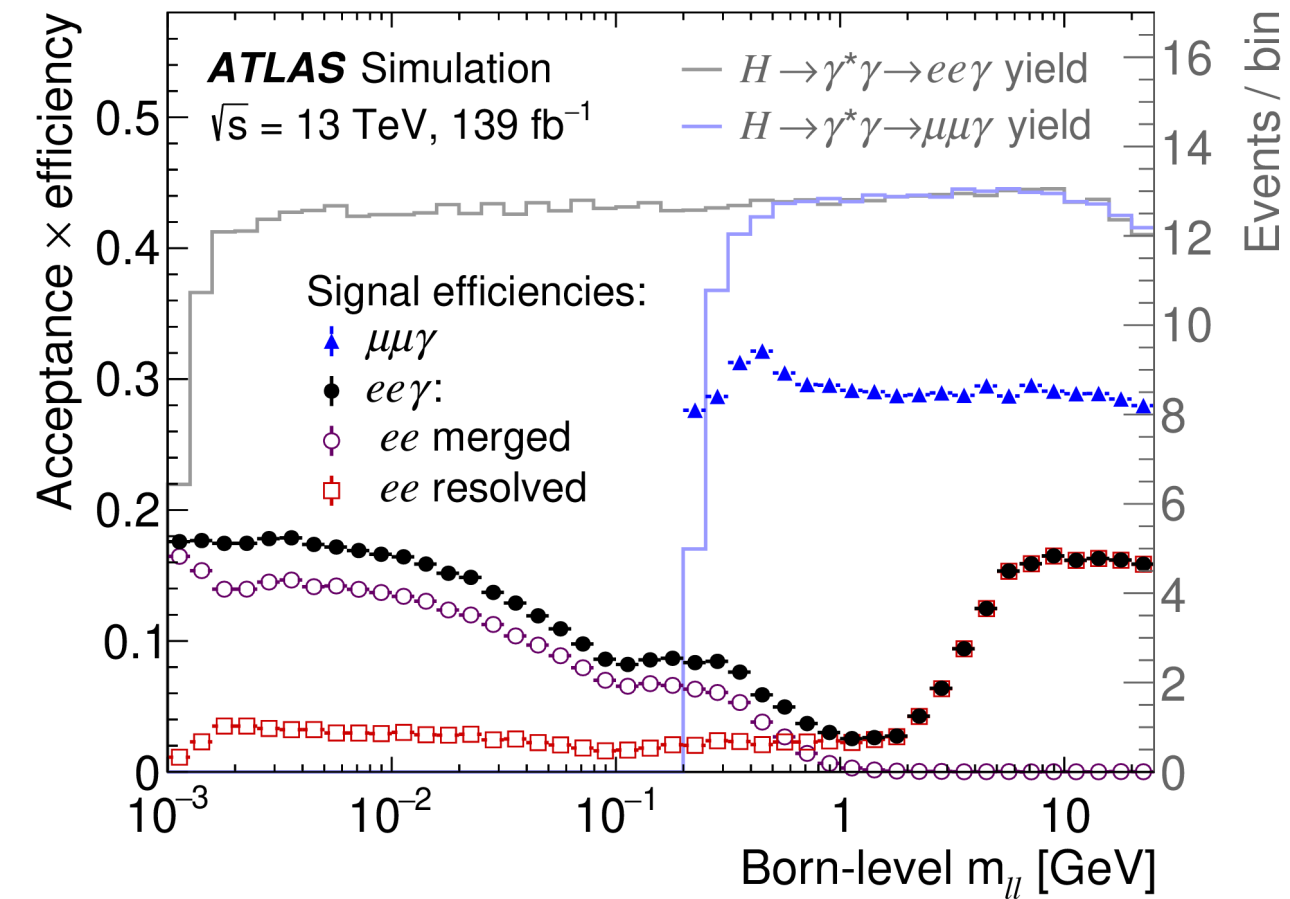
[Phys. Lett. B 809 \(2020\) 135754](#)



Similar excess in CMS.
 Potential evidence in Run3?

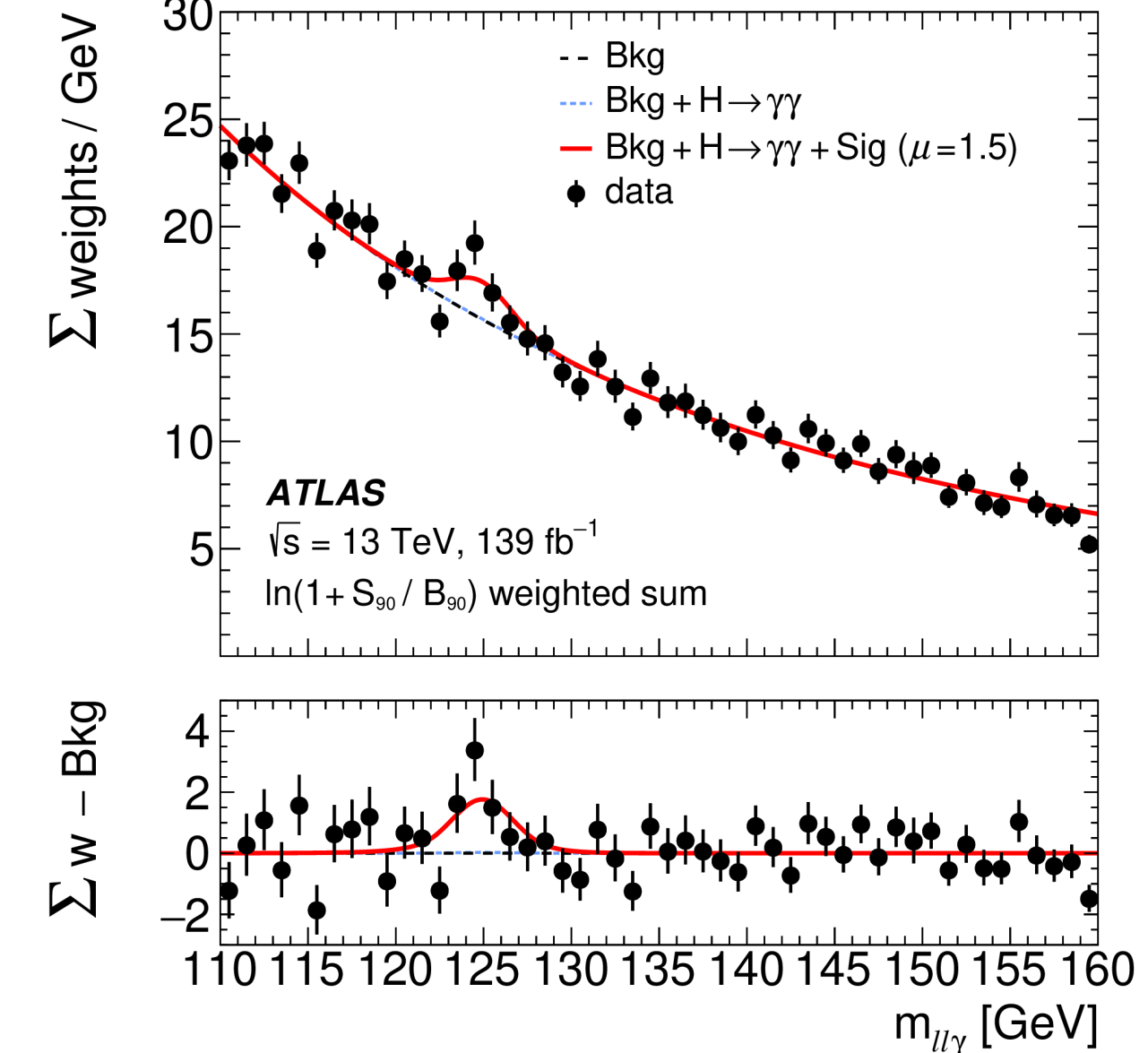
$H \rightarrow \gamma^*\gamma \rightarrow ll\gamma$

[Phys. Lett. B 819 \(2021\) 136412](#)



Expected 2.1σ
 Observed 3.2σ
 $\mu = 1.5 \pm 0.5$

First evidence! Keep watching with more data to look for SM deviations (compositeness, CP violation)



More precision in the bosonic channel, and putting all together

- Measurements of $\sigma \cdot \text{BR}$ in the full phase space, per production process and in simplified fiducial volumes (“bins”) of the phase space close have been performed in the 5 main Higgs boson decay channels with the full Run2 data: bb , $\tau\tau$, and the “discovery channels” $\gamma\gamma$, $ZZ \rightarrow 4l$ and $WW \rightarrow e\nu\mu\nu$ with the full Run2 data

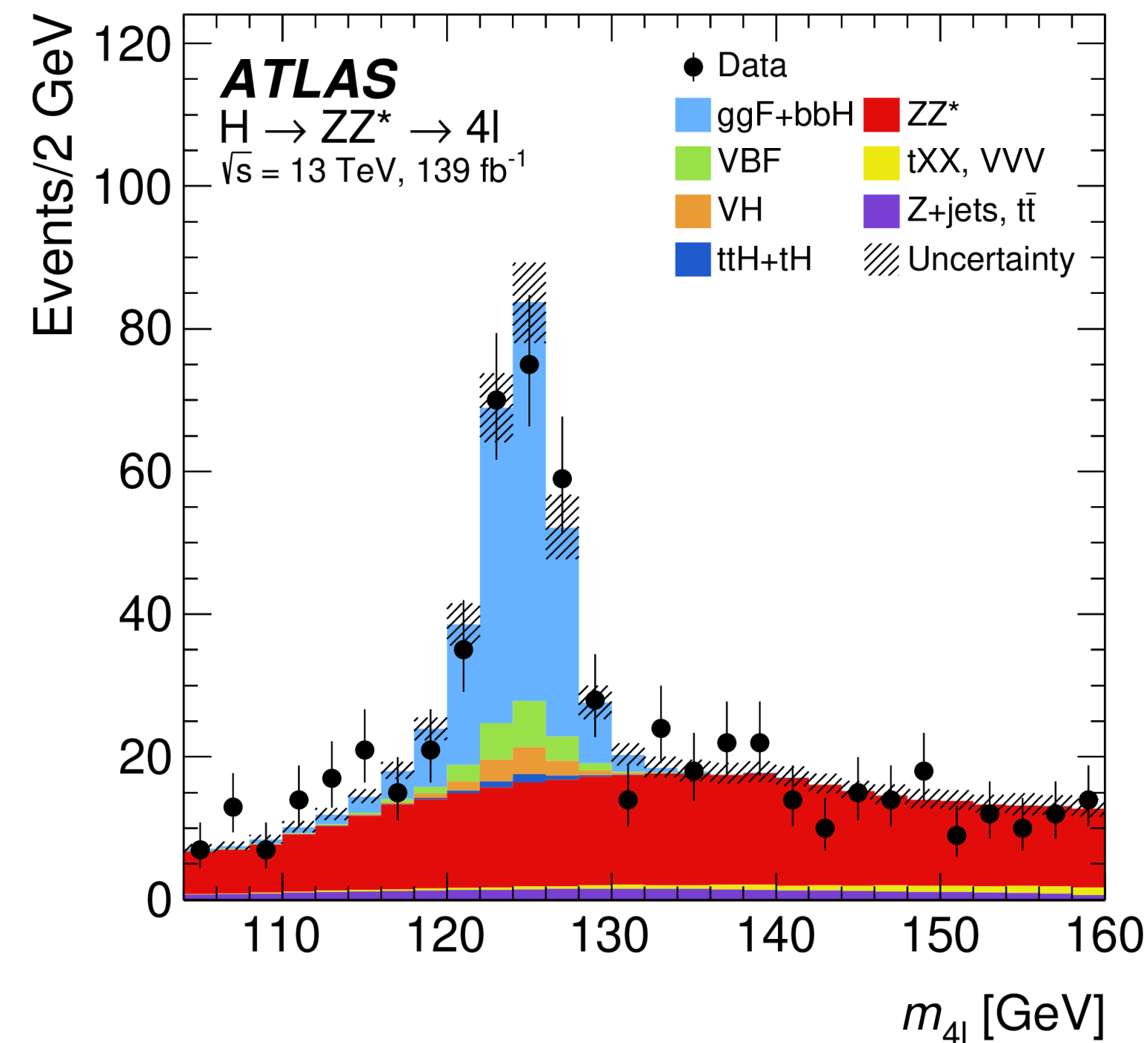
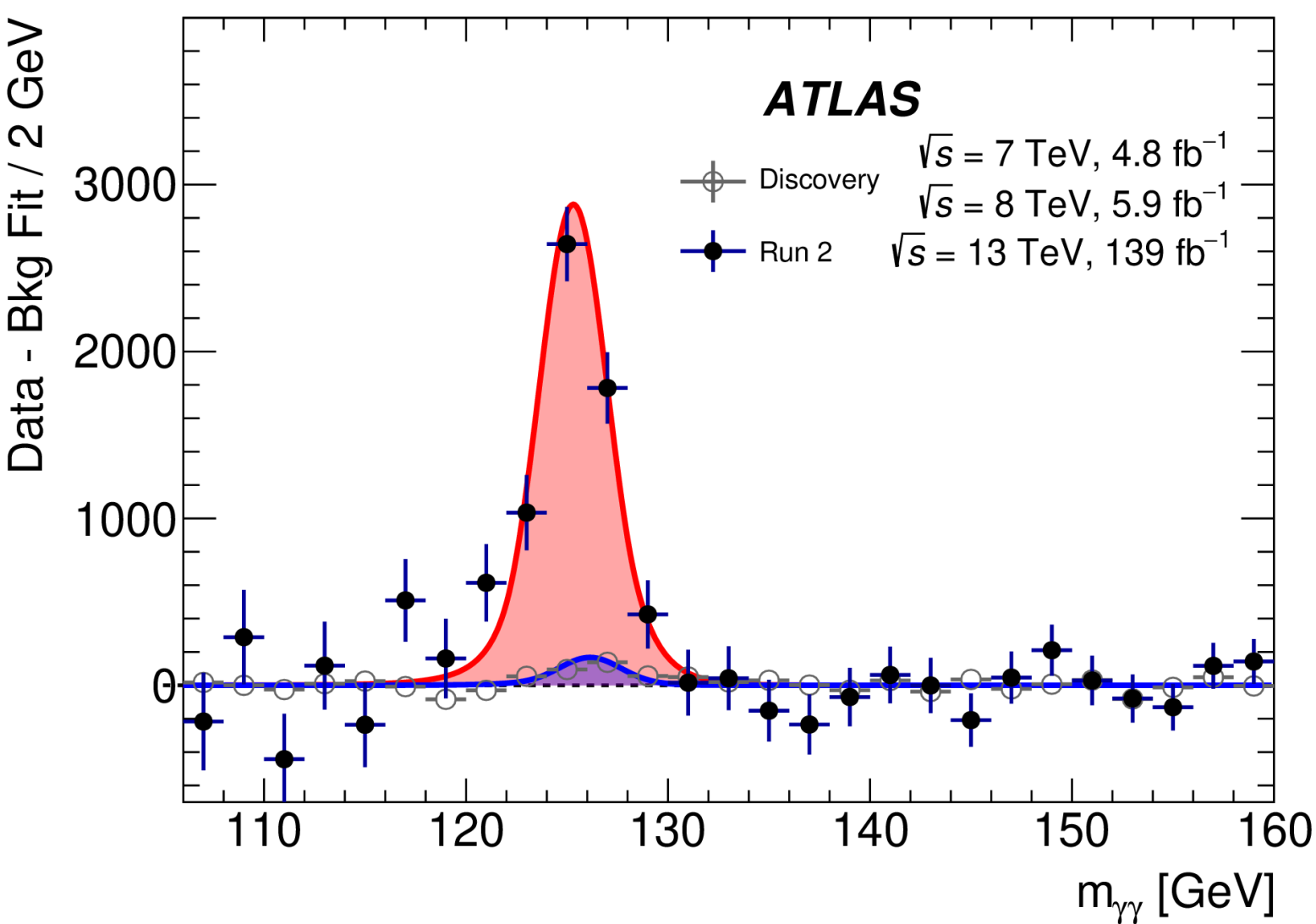
e.g. $H \rightarrow \gamma\gamma$ (Phys. Rev. D 98 (2018) 052005)

- Exploit different signatures of main production modes, define event categories enriched in one particular production mode or bin
- Simultaneous fit to event yields in various categories allows measurement of signal strengths for each production and decay mode or bin

$\gamma\gamma$: [arXiv:2207.00348](https://arxiv.org/abs/2207.00348)

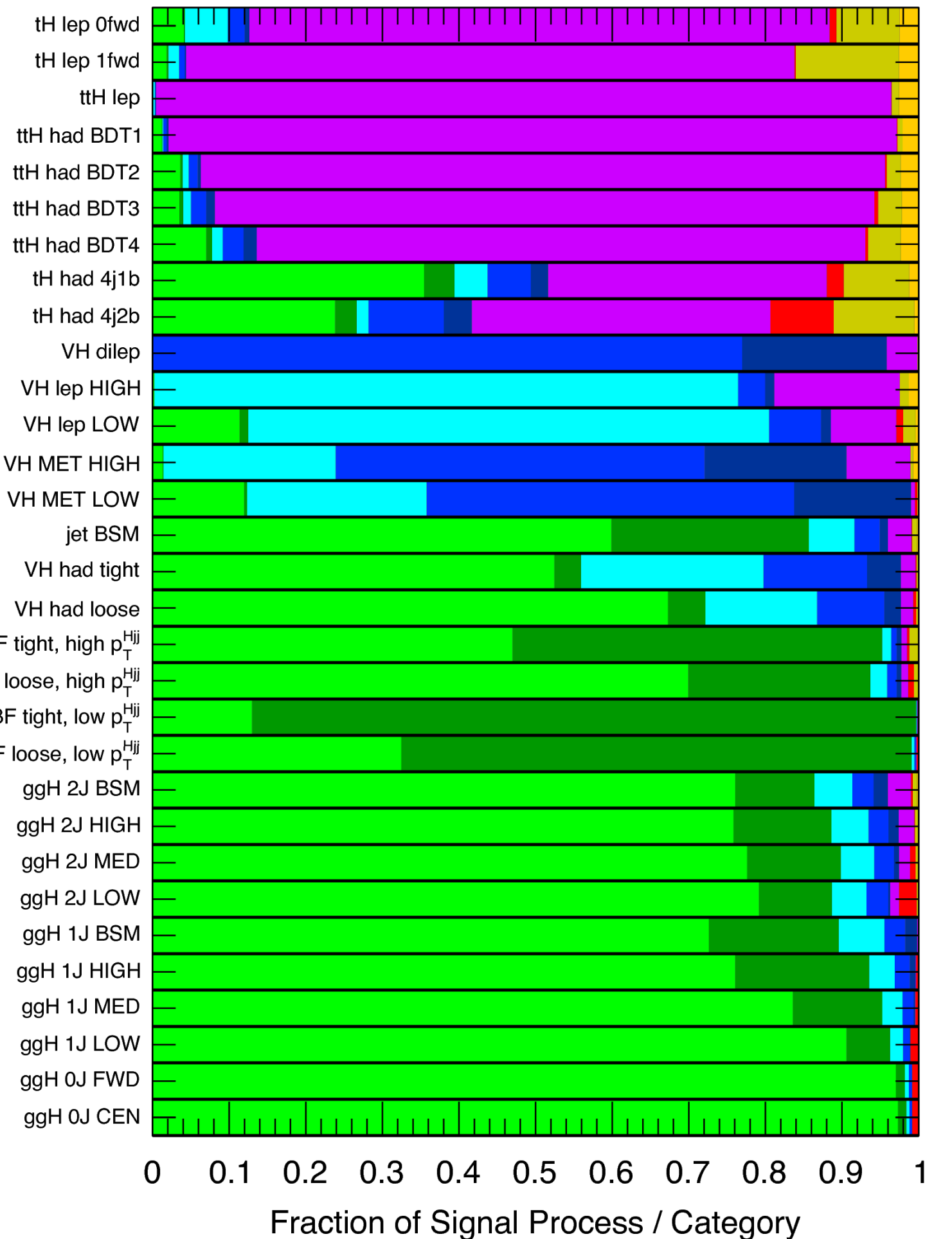
ZZ : [Eur. Phys. J. C 80 \(2020\) 957](https://arxiv.org/abs/2002.08864)

WW : [arxiv:2207.00338](https://arxiv.org/abs/2207.00338)



Legend for ATLAS Simulation: ggH (green), VBF (dark green), WH (cyan), ZH (blue), ggZH (dark blue), ttH (magenta), bbH (red), tHq (yellow-green), tHW (yellow)

ATLAS Simulation $H \rightarrow \gamma\gamma, m_H = 125.09 \text{ GeV}$



ttH-tag

WH/ZH-tag

VBF-tag

ggF-tag

More precision in the bosonic channel, and putting all together

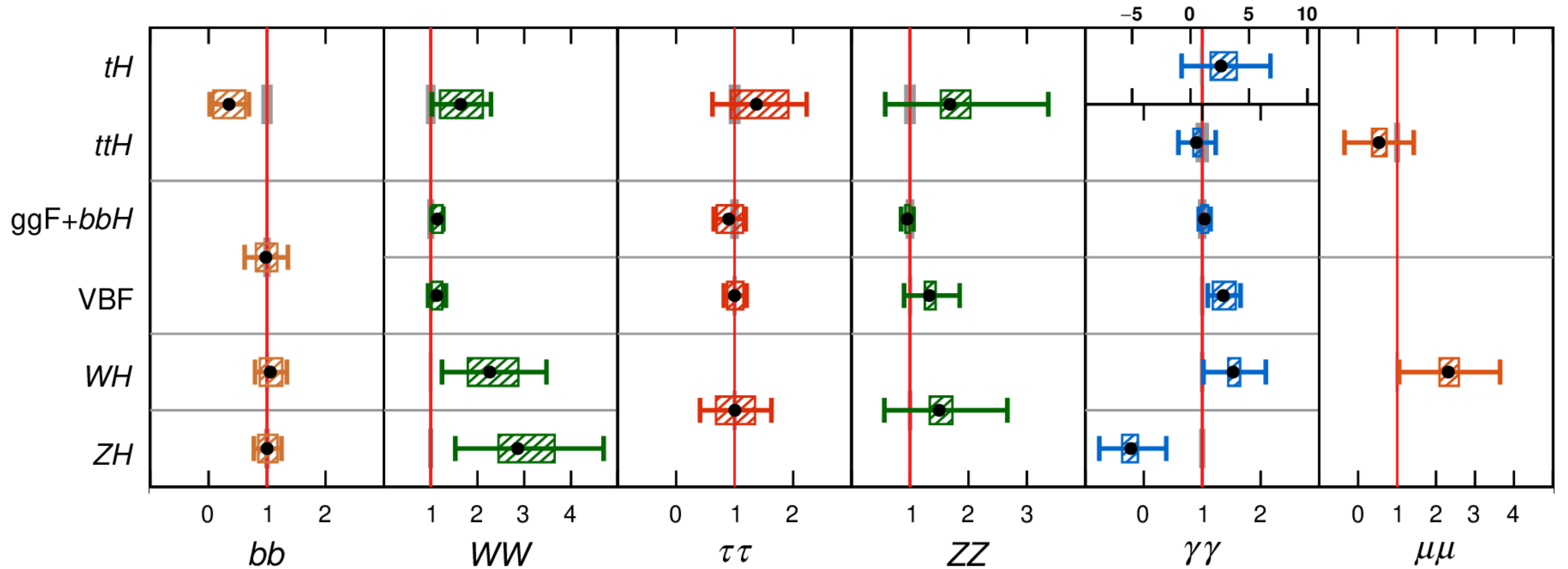
[Nature 607, 52 \(2022\)](#)

ATLAS Run 2

● Data (Total uncertainty)

▨ Syst. uncertainty

■ SM prediction



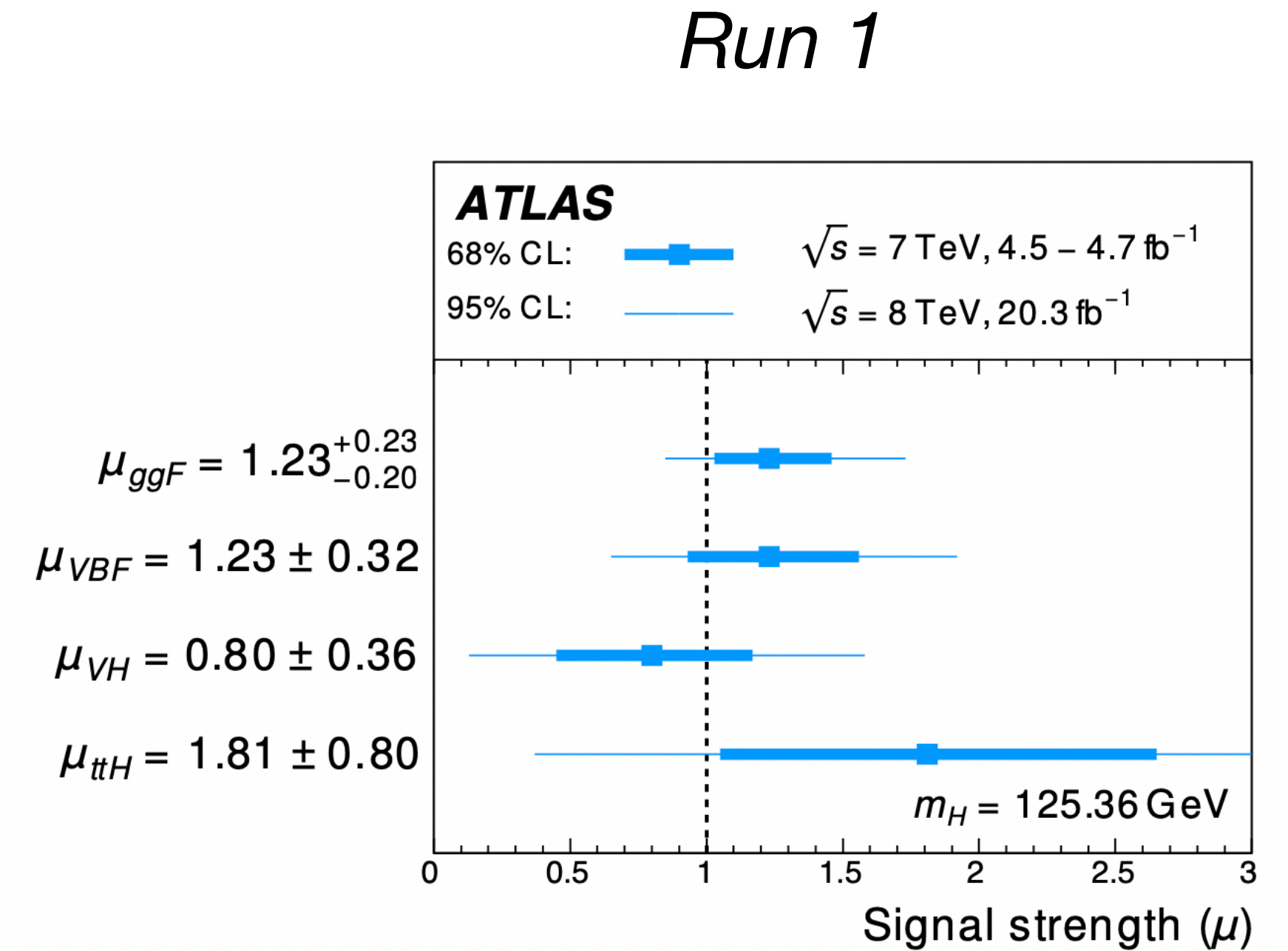
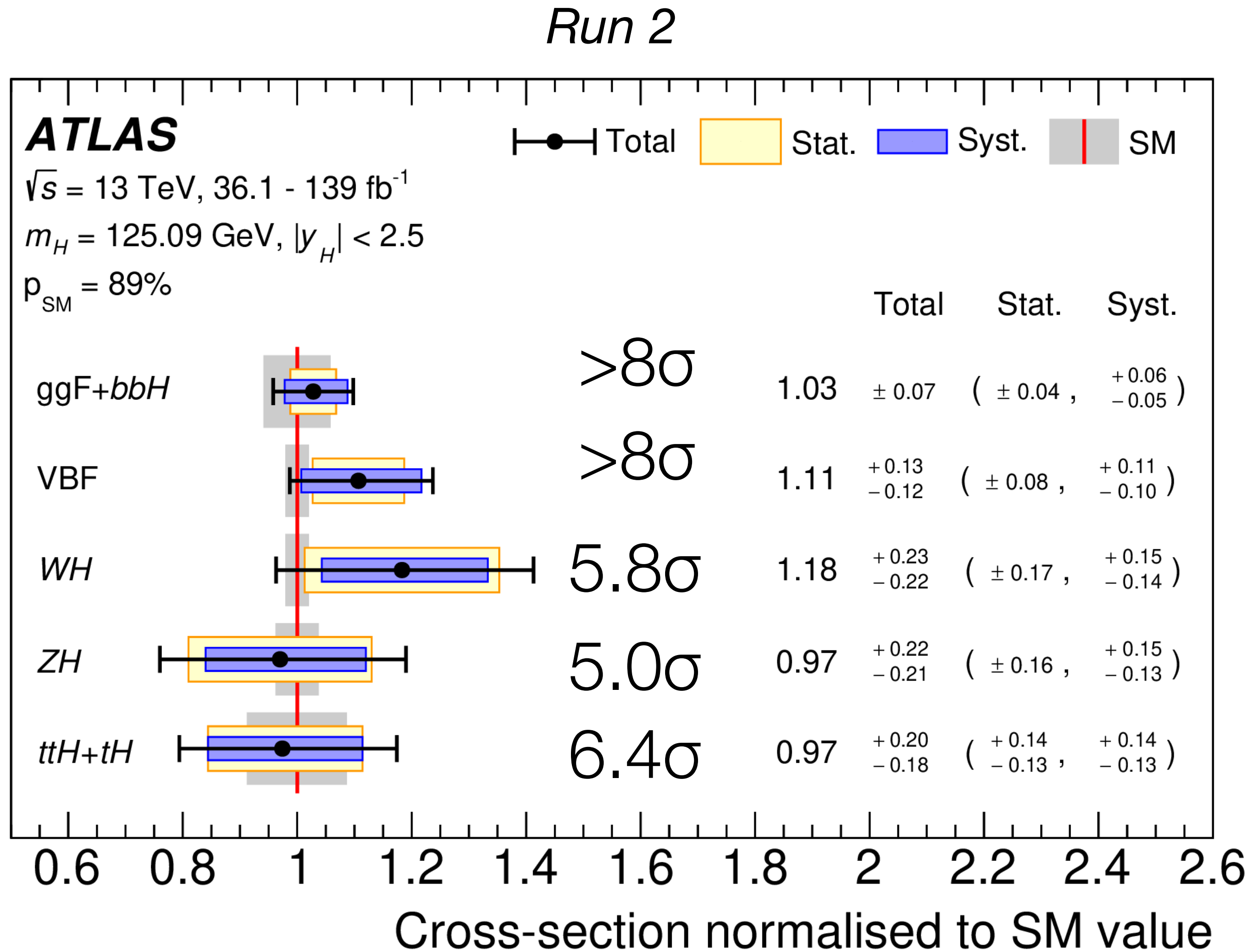
p-value = 72%

$\sigma \times B$ normalized to SM prediction

Measurements of production mode cross-sections

- Assume SM branching ratios and measure production mode cross-sections

[Nature 607, 52 \(2022\)](#)

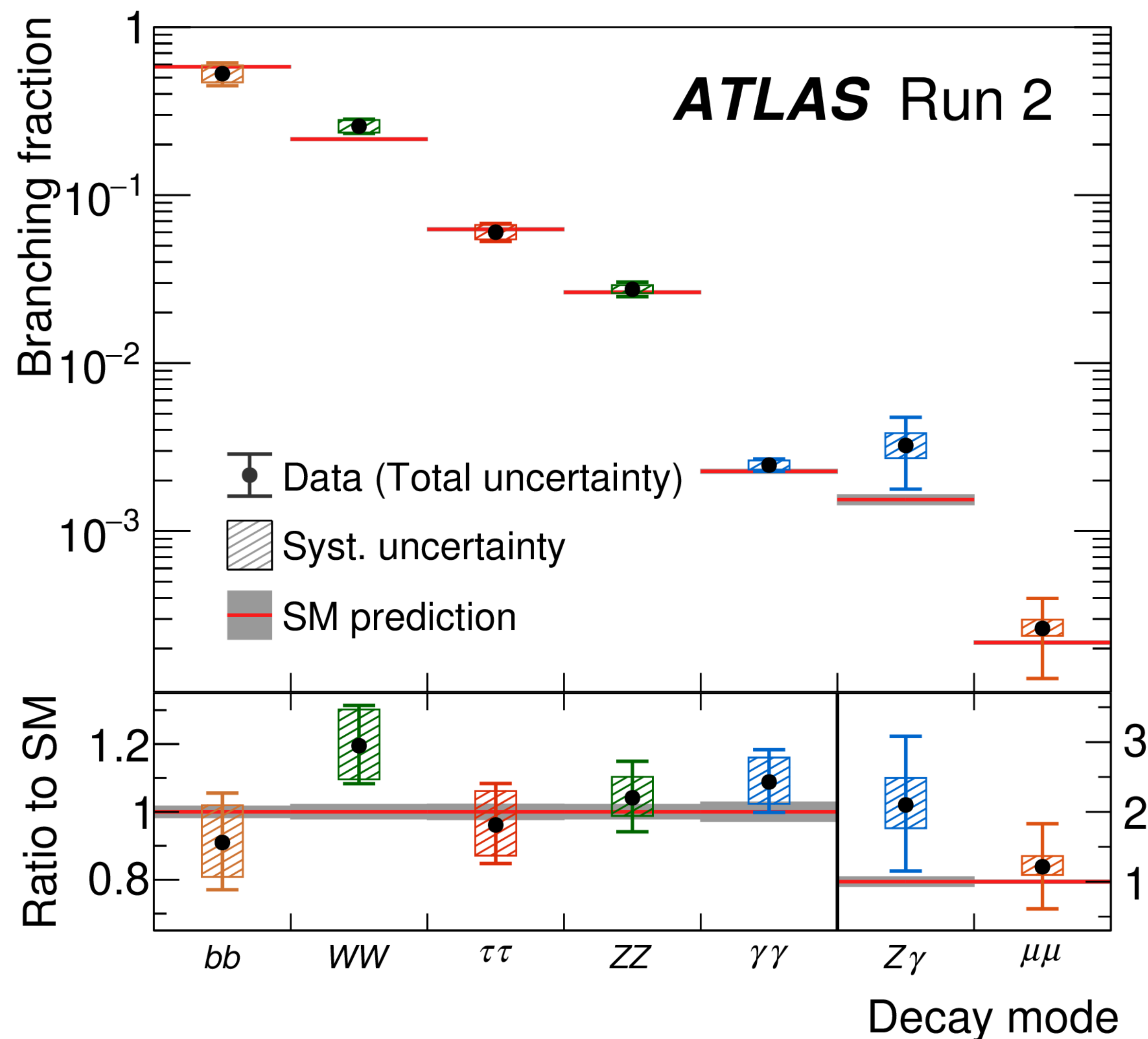


Process	VBF	ttH	WH	ZH	VH
Observed	4.3	2.5	2.1	0.9	2.6
Expected	3.8	1.5	2.0	2.1	3.1

Observation of all main production modes and 3-5x reduction in uncertainties compared to Run1 (precision 7-22%)

Measurement of decay branching ratios

- Assume SM production mode cross-sections and measure branching ratios



Run 1 [Nature 607, 52 \(2022\)](#)

ATLAS
 $m_H = 125.36$ GeV

	$\sigma(\text{stat.})$	$\sigma(\text{sys inc. theory})$	$\sigma(\text{theory})$
$H \rightarrow \gamma\gamma$ $\mu = 1.17^{+0.28}_{-0.26}$	+0.23 -0.23 +0.16 -0.11 +0.12 -0.08		
$H \rightarrow ZZ^*$ $\mu = 1.46^{+0.40}_{-0.34}$	+0.35 -0.31 +0.19 -0.13 +0.18 -0.11		
$H \rightarrow WW^*$ $\mu = 1.18^{+0.24}_{-0.21}$	+0.16 -0.16 +0.17 -0.14 +0.13 -0.09		
$H \rightarrow \tau\tau$ $\mu = 1.44^{+0.42}_{-0.37}$	+0.30 -0.29 +0.29 -0.23 +0.16 -0.10		
$H \rightarrow bb$ $\mu = 0.63^{+0.39}_{-0.37}$	+0.31 -0.30 +0.24 -0.23 +0.09 -0.07		
$H \rightarrow \mu\mu$ $\mu = -0.7^{+3.7}_{-3.7}$	+3.6 -3.6 +0.5 -0.7 +0.4 -0.4		
$H \rightarrow Z\gamma$ $\mu = 2.7^{+4.6}_{-4.5}$	+4.3 -4.2 +1.7 -1.3 +1.1 -0.3		

Uncertainties in main decay channels now at 10-12% level

Inclusive signal strength

- Assuming a global scale factor affecting equally all production modes and decay channels:

[Nature 607, 52 \(2022\)](#)

$$\mu_{if} = \left(\sigma_i / \sigma_i^{\text{SM}} \right) \times \left(B_f / B_f^{\text{SM}} \right) = \mu$$

- Run2 result:

$$\mu = 1.05 \pm 0.06 = 1.05 \pm 0.03 \text{ (stat.)} \pm 0.03 \text{ (exp.)} \pm 0.04 \text{ (sig. th.)} \pm 0.02 \text{ (bkg. th.)}.$$

- Compare to Run1:

$$\mu = 1.18^{+0.15}_{-0.14} = 1.18 \pm 0.10 \text{ (stat.)} \pm 0.07 \text{ (syst.)}^{+0.08}_{-0.07} \text{ (theo.)},$$

Improvement in relative precision: 14% (Run 1) → 6% (Run 2)

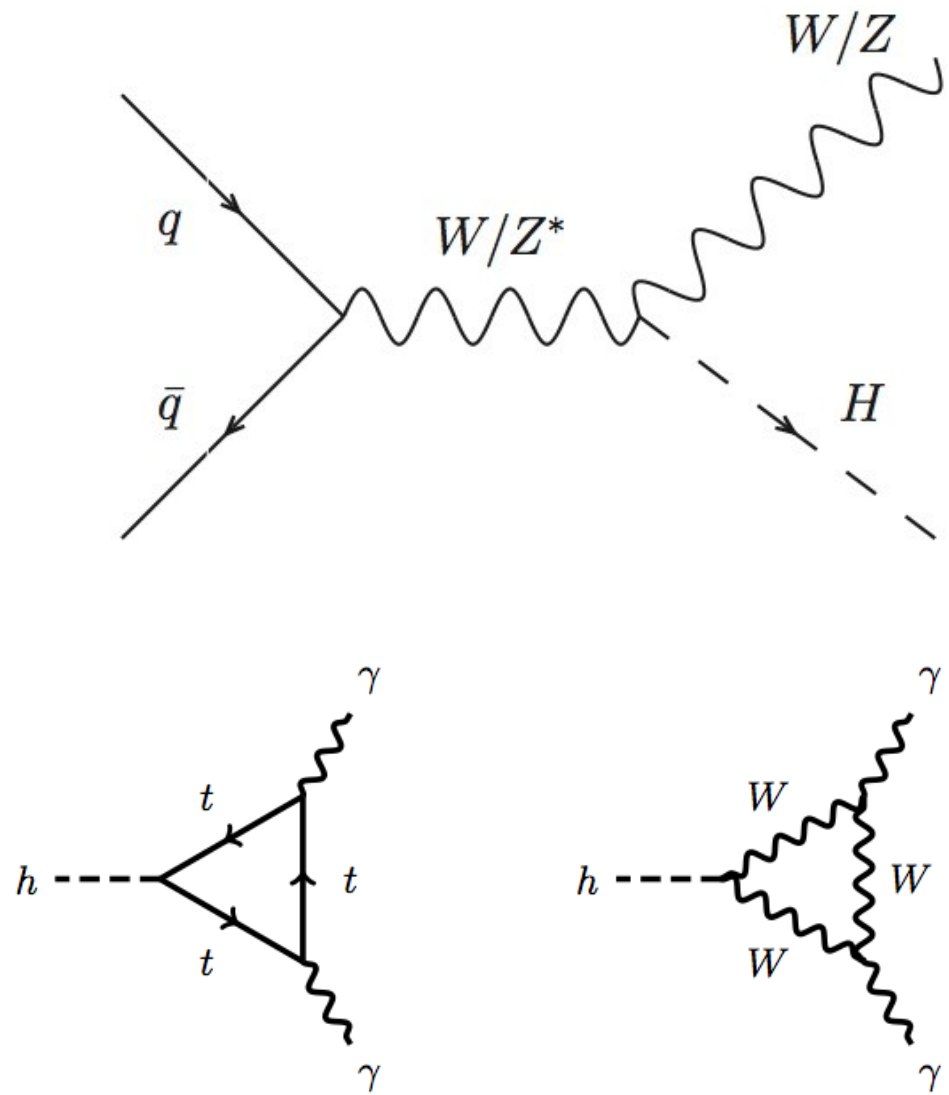
Theory uncertainty: 7% (Run 1) → 4% (Run 2)

Measurement of Higgs boson couplings to other particles

- Assuming that the signals in the different channels are due to a single, narrow, CP-even resonance, the signal strengths can be parametrised in terms of coupling scaling factors κ

$$(\sigma \cdot \text{BR})(i \rightarrow H \rightarrow f) = \frac{\sigma_i^{\text{SM}} \kappa_i^2 \cdot \Gamma_f^{\text{SM}} \kappa_f^2}{\Gamma_H^{\text{SM}} \kappa_H^2}$$

- A couple of examples:



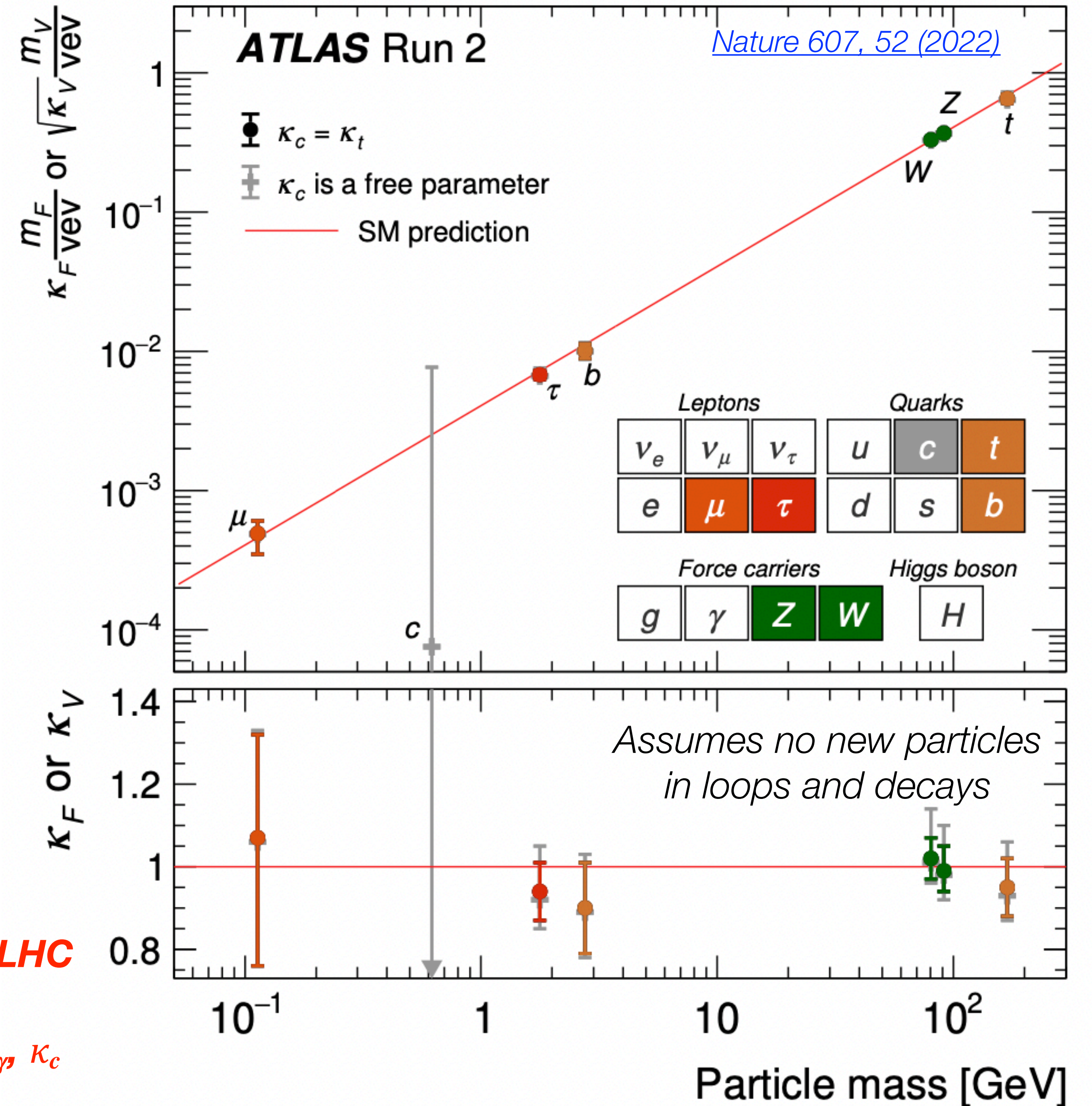
$$\sigma \propto \kappa_W^2, \kappa_Z^2$$

$$\Gamma_{\gamma\gamma} \propto 1.59 \kappa_W^2 + 0.07 \kappa_t^2 - 0.67 \kappa_W \kappa_t$$

All measured couplings to fermions and bosons agree with the SM

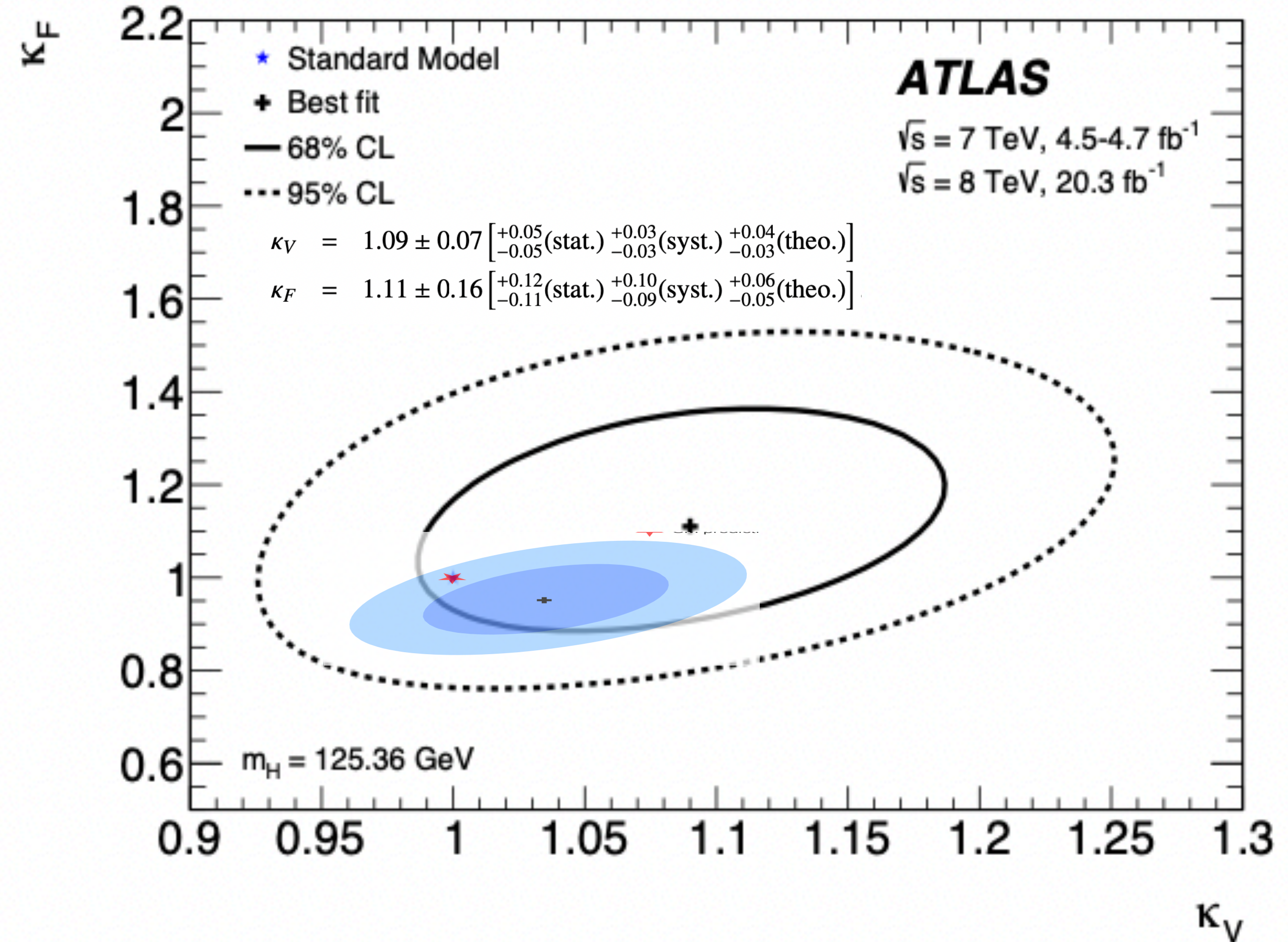
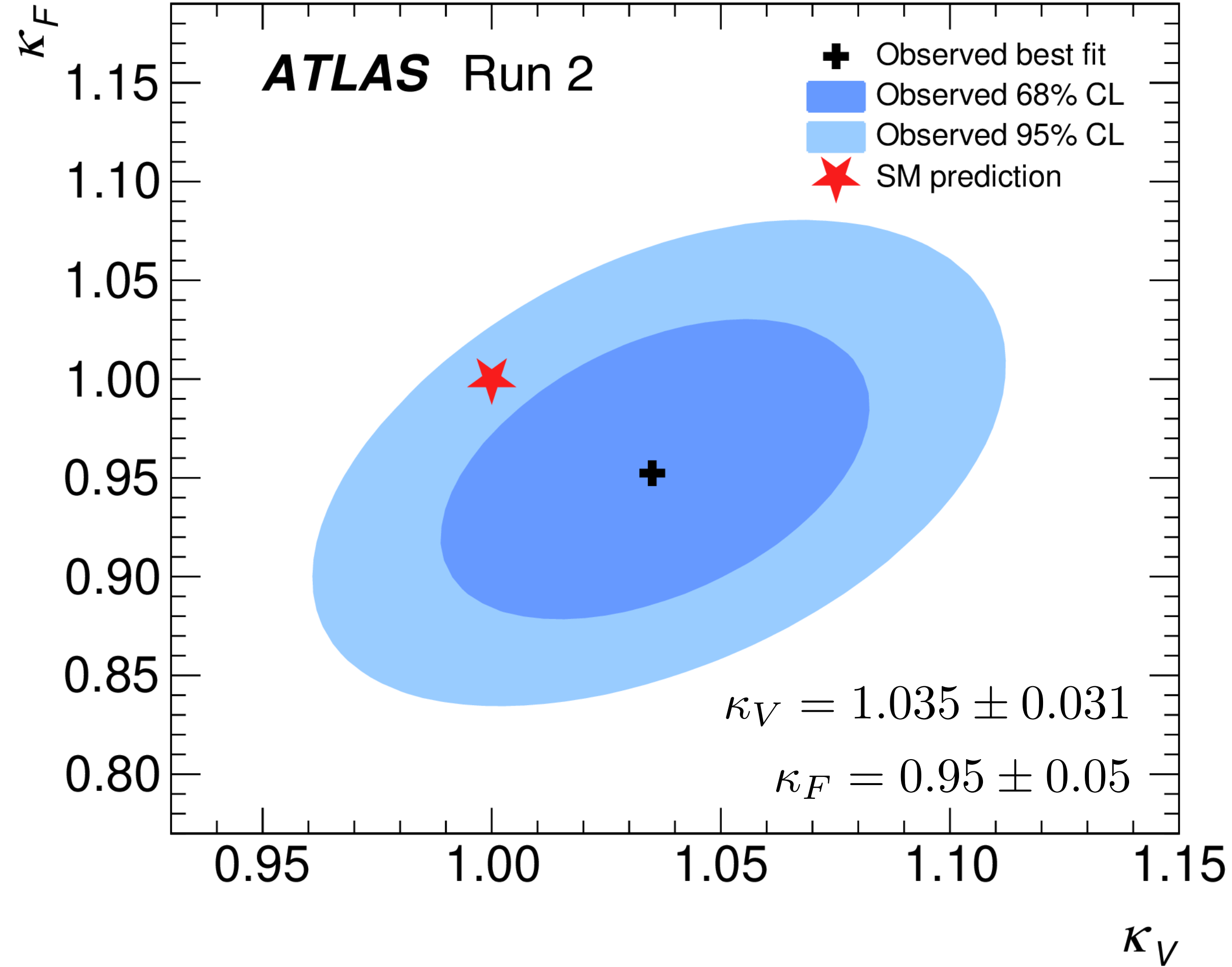
Current uncertainties (5-11%, 30% for μ) to be reduced by x3 at HL-LHC

In all scenarios, statistical ~ systematic uncertainty except for $\kappa_{\mu\mu}, \kappa_{Z\gamma}, \kappa_c$



Measurement of Higgs boson couplings to other particles

[Nature 607, 52 \(2022\)](#)



Measurement of “simplified template cross-sections” (STXS)

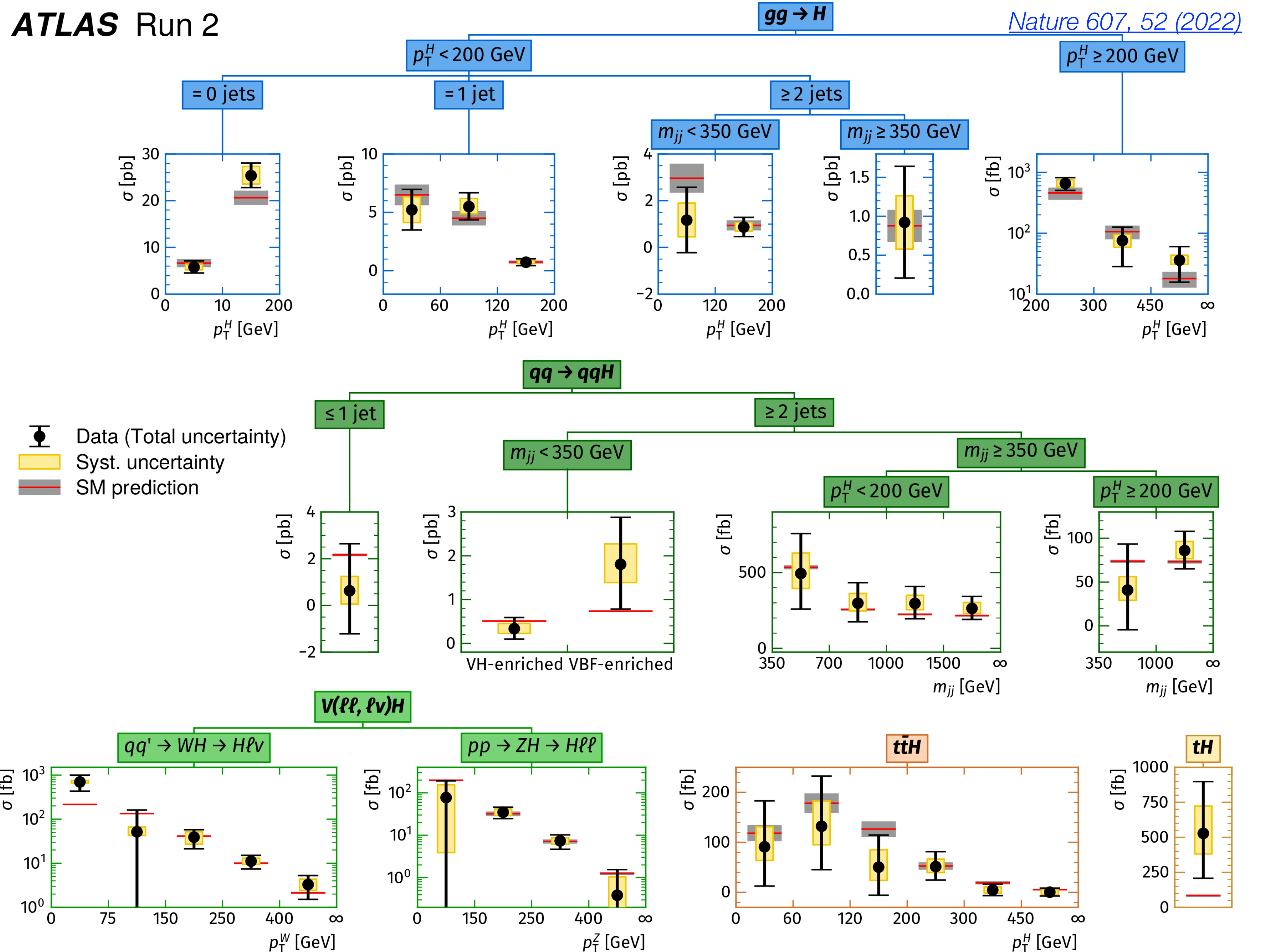
- Measurements of $\sigma \cdot \text{BR}$ in simplified fiducial volumes defined based on particle-level quantities such as p_{T}^H , $N(\text{jets})$, m_{jj} , ..
- Reduce extrapolation uncertainties to full phase space
- Provide less model-dependent measurements that can be used even in the future for setting constraints on SM (with more precise calculations) or other theories
- Allow searching more in detail / systematic way for deviations from SM due to anomalous couplings through e.g. EFT-based interpretation

Good agreement with SM (p-value 94%)

Most measurement statistically limited

Further granularity possible w/ more data

ATLAS Run 2



Nature 607, 52 (2022)

Effective field theory interpretation

- Effective Lagrangian built with SM fields at lowest order in expansion in $1/\Lambda$ (Λ = new physics energy scale):

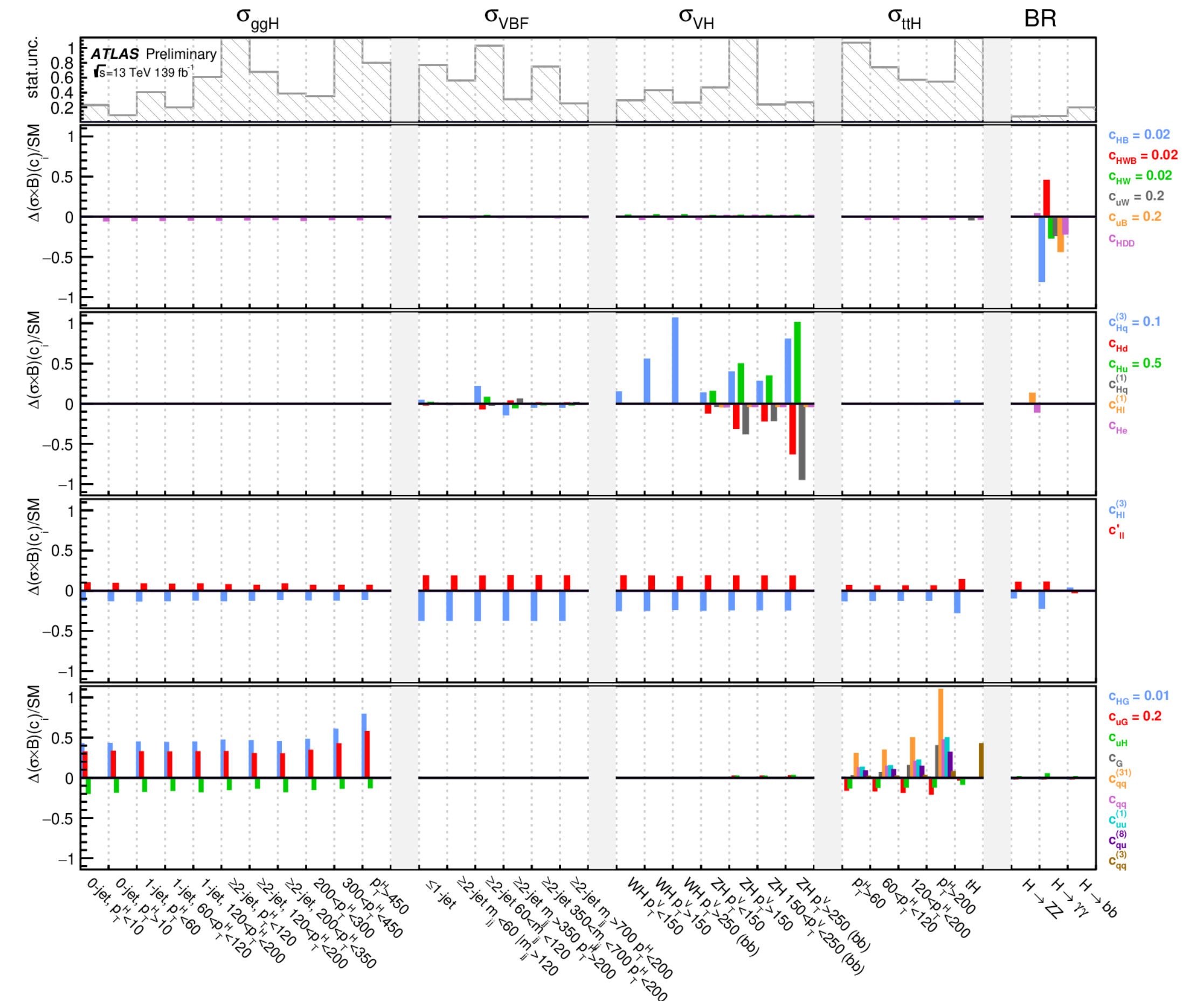
[ATLAS-CONF-2020-053](#)

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{C_i^{(d)}}{\Lambda^{(d-4)}} \mathcal{O}_i^{(d)} \quad \text{for } d > 4.$$

$d=6, \Lambda=1 \text{ TeV}$

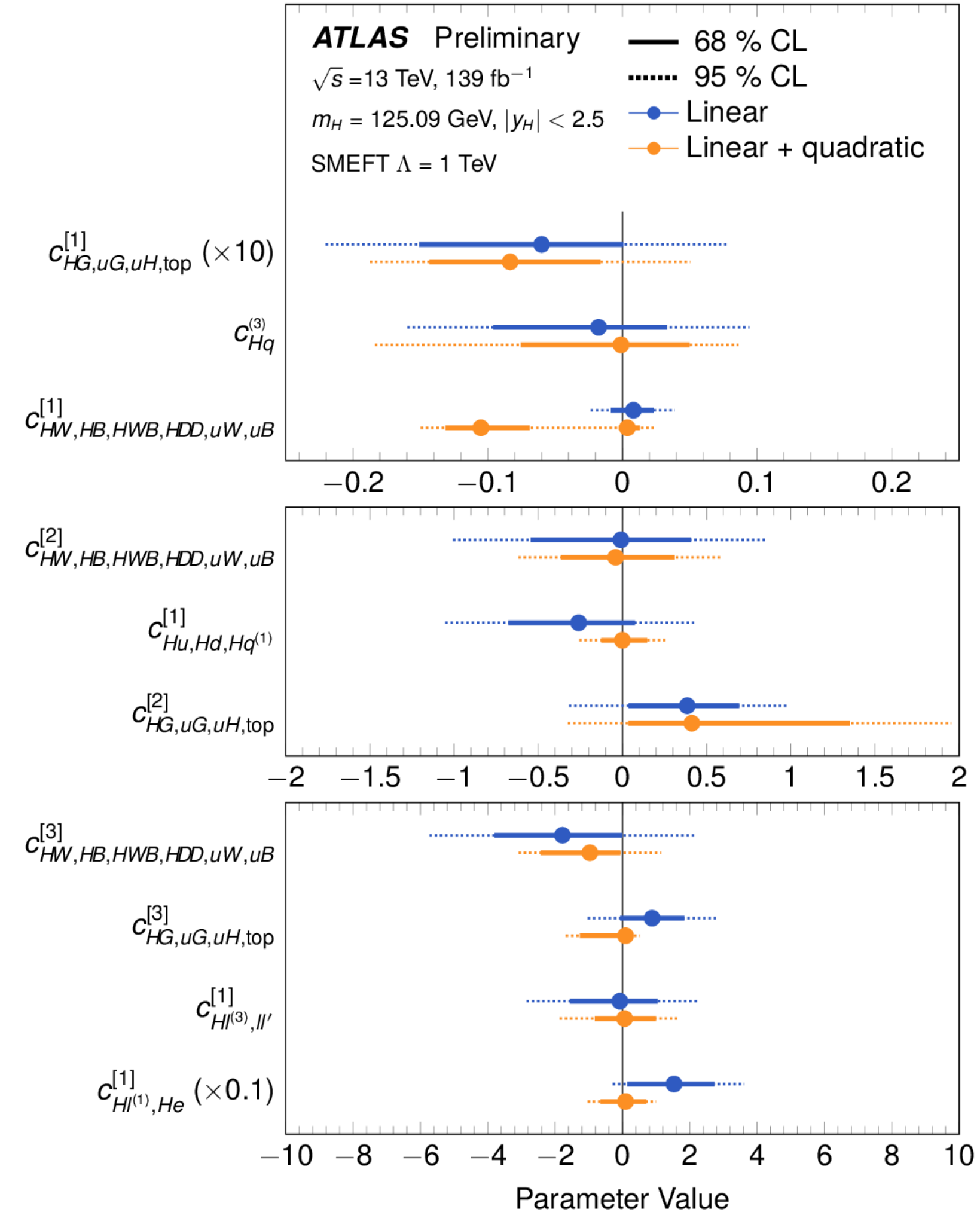
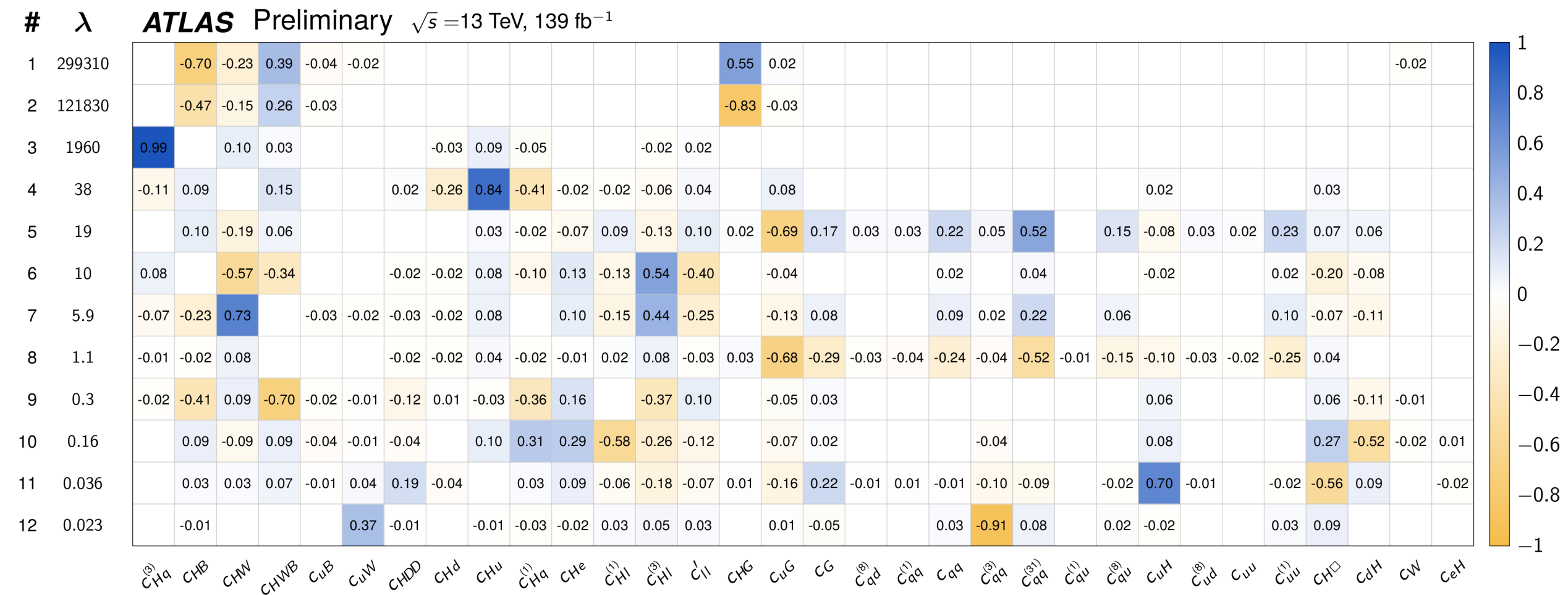
- Several operators induce anomalous couplings between the Higgs and the SM fields that modify the STXS predictions

Wilson coefficient	Operator	Wilson coefficient	Operator
$c_{H\Box}$	$(H^\dagger H)\Box(H^\dagger H)$	c_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{H} G_{\mu\nu}^A$
c_{HDD}	$(H^\dagger D^\mu H)^* (H^\dagger D_\mu H)$	c_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{H} W_{\mu\nu}^I$
c_{HG}	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	c_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{H} B_{\mu\nu}$
c_{HB}	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	c'_{ll}	$(\bar{l}_p \gamma_\mu l_t)(\bar{l}_r \gamma^\mu l_s)$
c_{HW}	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	$c_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_t)(\bar{q}_r \gamma^\mu q_s)$
c_{HWB}	$H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu}$	$c_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$
c_{eH}	$(H^\dagger H)(\bar{l}_p e_r H)$	c_{qq}	$(\bar{q}_p \gamma_\mu q_t)(\bar{q}_r \gamma^\mu q_s)$
c_{uH}	$(H^\dagger H)(\bar{q}_p u_r \tilde{H})$	$c_{qq}^{(31)}$	$(\bar{q}_p \gamma_\mu \tau^I q_t)(\bar{q}_r \gamma^\mu \tau^I q_s)$
c_{dH}	$(H^\dagger H)(\bar{q}_p d_r \tilde{H})$	c_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$
$c_{Hl}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{l}_p \gamma^\mu l_r)$	$c_{uu}^{(1)}$	$(\bar{u}_p \gamma_\mu u_t)(\bar{u}_r \gamma^\mu u_s)$
$c_{Hl}^{(3)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{l}_p \tau^I \gamma^\mu l_r)$	$c_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_t)(\bar{u}_r \gamma^\mu u_s)$
c_{He}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{e}_p \gamma^\mu e_r)$	$c_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$
$c_{Hq}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu q_r)$	$c_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
$c_{Hq}^{(3)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{q}_p \tau^I \gamma^\mu q_r)$	$c_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$
c_{Hu}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{u}_p \gamma^\mu u_r)$	c_W	$\epsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$
c_{Hd}	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{d}_p \gamma^\mu d_r)$	c_G	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$



Effective field theory interpretation (II)

- Not enough measurements to constrain all coefficients => do PCA analysis to find rotated basis of operators with maximum sensitivity



**Expect to see more of this in the future
 Combination with EW/top measurements for
 even more global interpretation**

Going even more fiducial: (fiducial) differential cross-sections

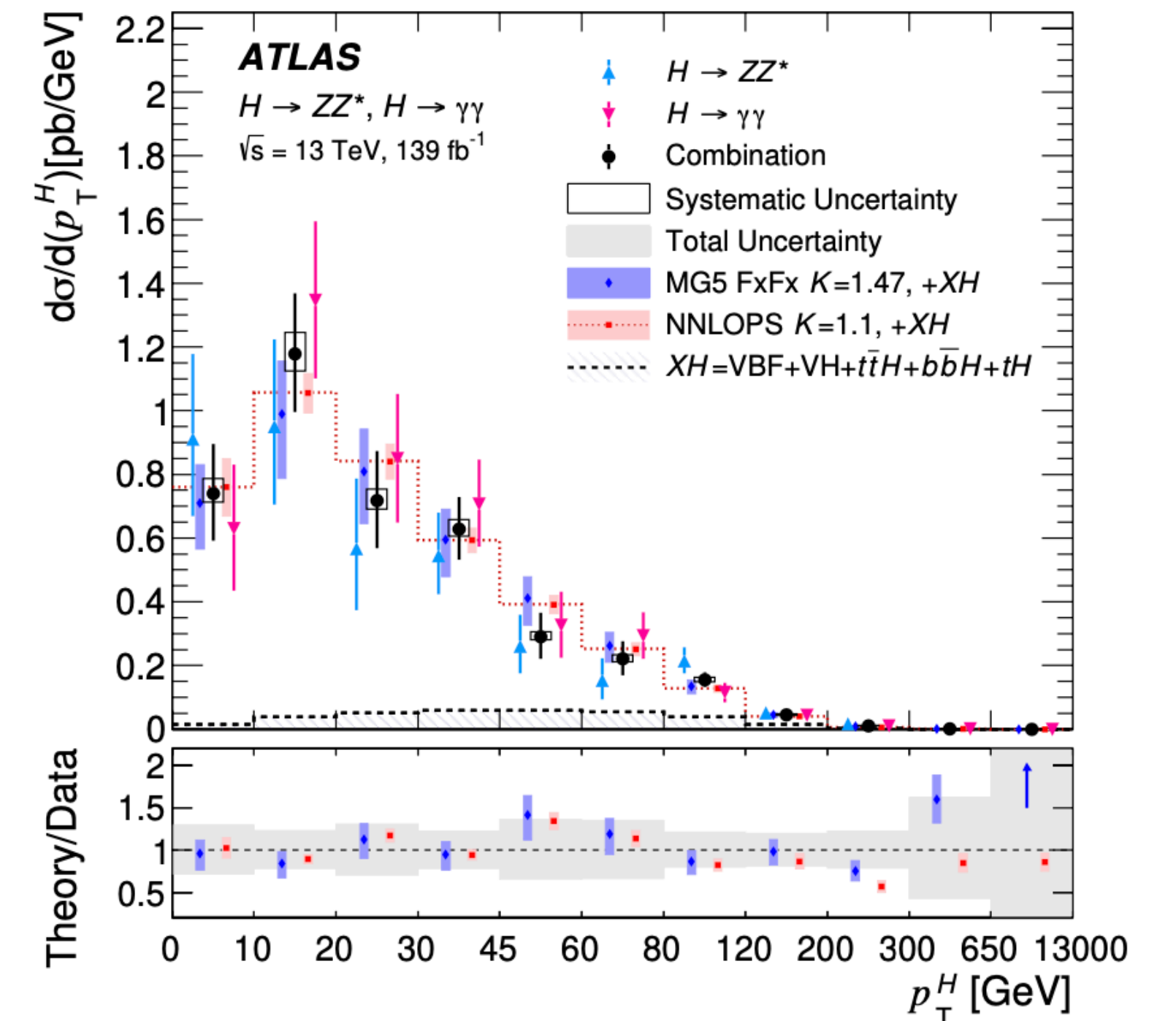
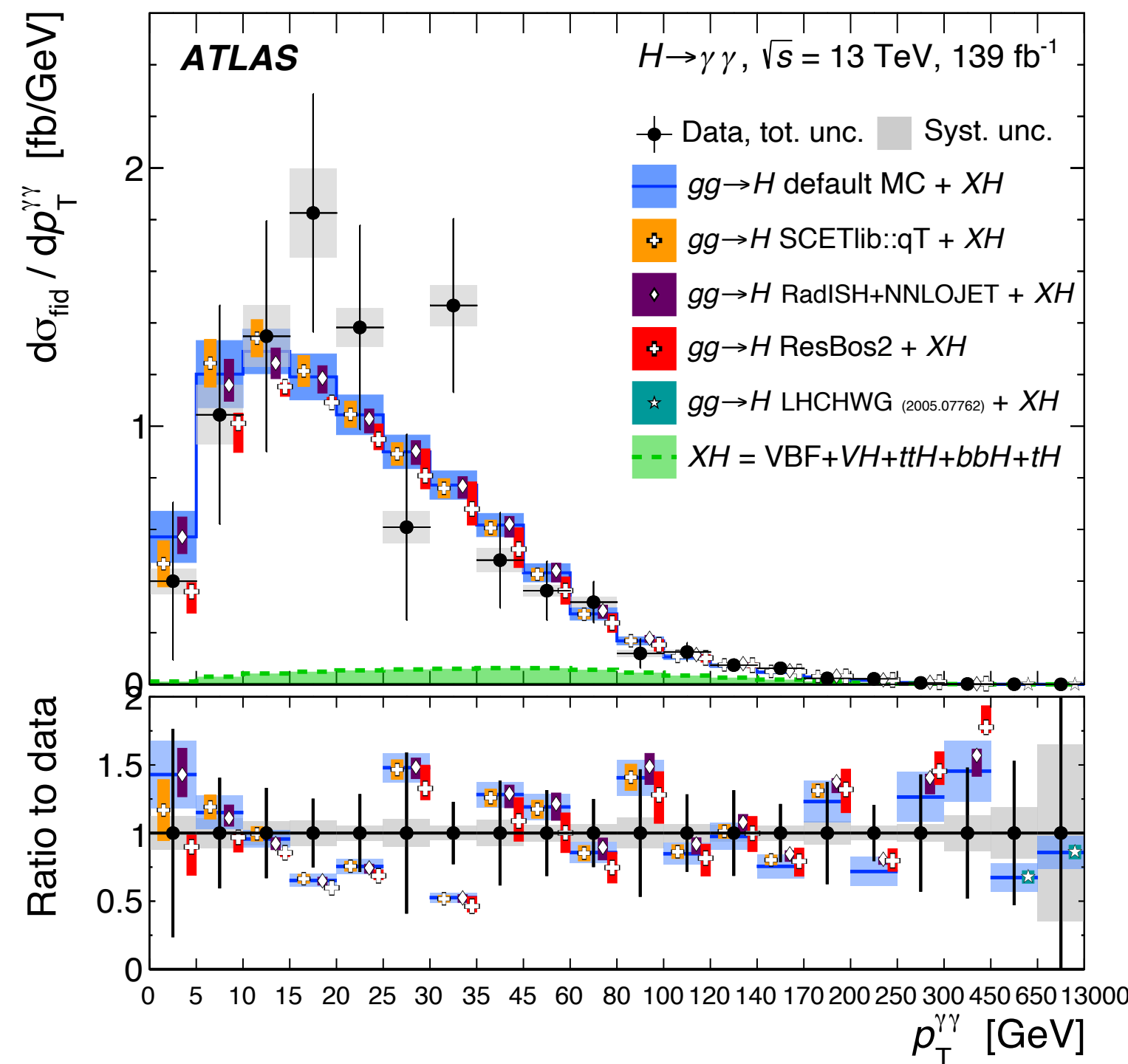
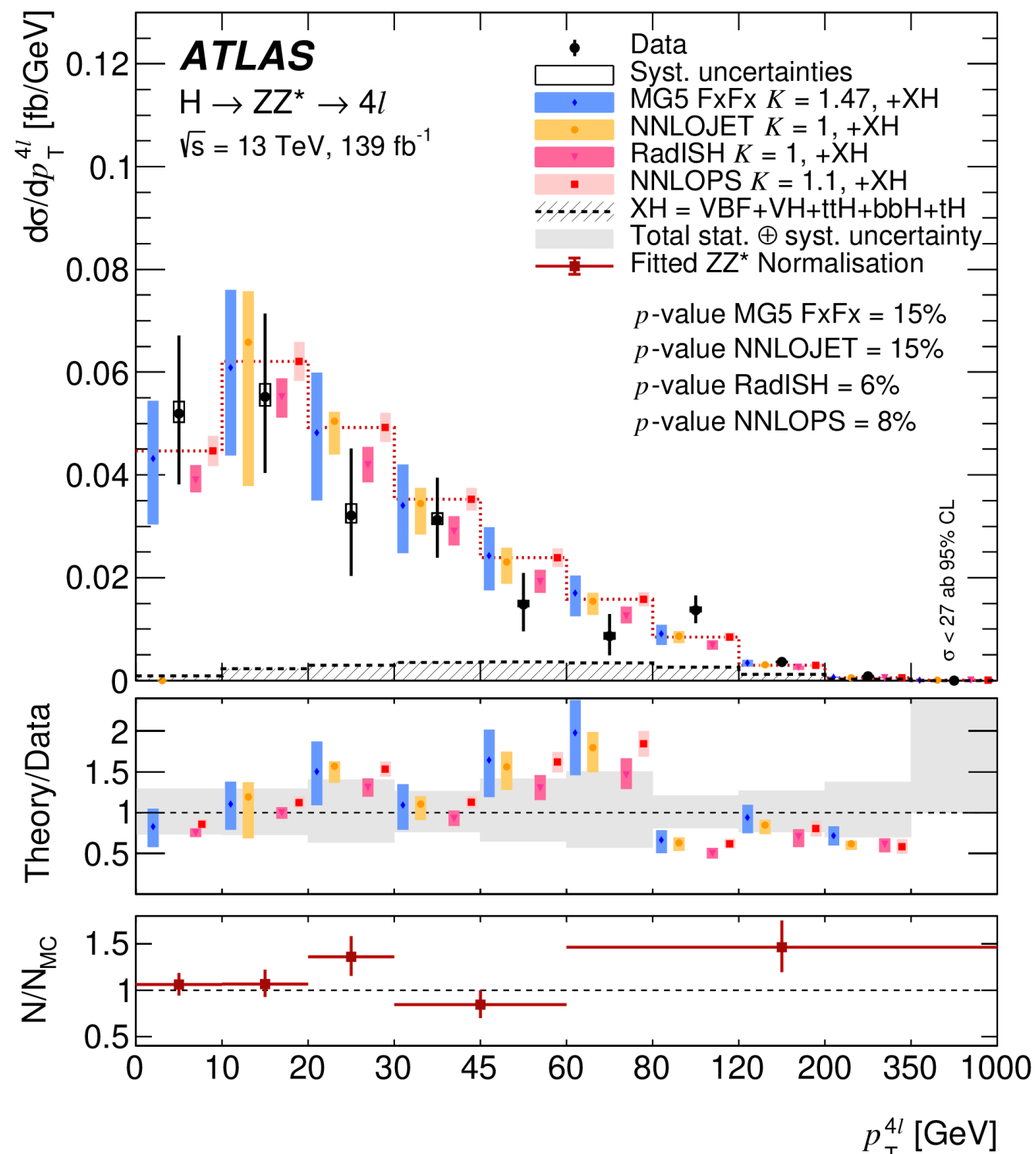
- More granular measurement as a function of several observables characterising the kinematics of Higgs boson production (Higgs p_T , N_{jets} , leading-jet p_T , invariant mass of leading and subleading jets (if present), ...)
- ➔ Measured observables can be sensitive to production mode xsection ratios / spin / CP / Higgs boson couplings ..
- In fiducial phase space, very close to experimental selection
- ➔ Minimise model dependence from extrapolation from selected to full phase space; efficiency similar for all production modes

$H \rightarrow 4l$ [Eur. Phys. J. C 80 \(2020\) 942](#)

$H \rightarrow \gamma\gamma$ [JHEP 08 \(2022\) 027](#)

Combination (in full phase space)

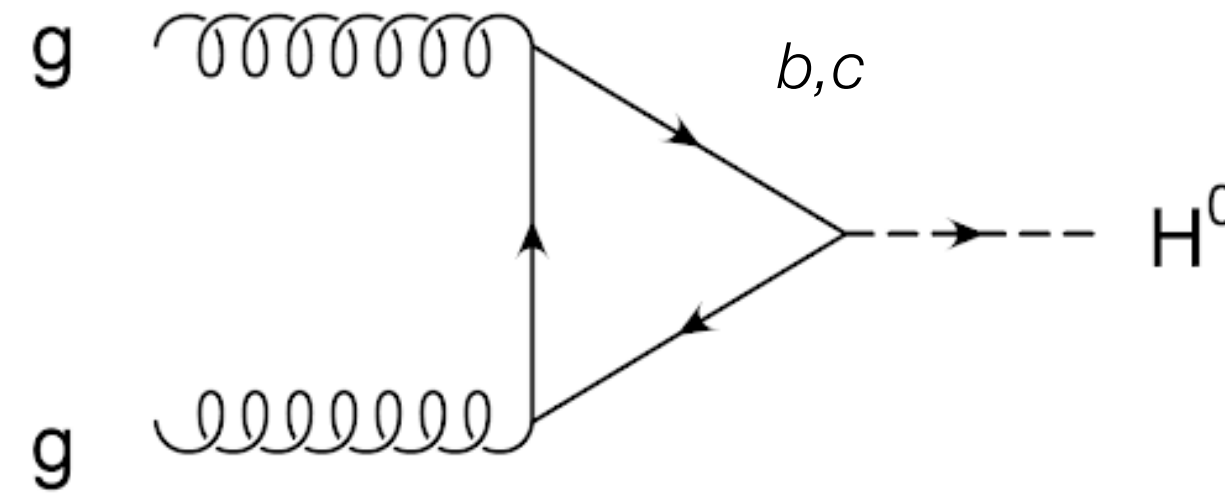
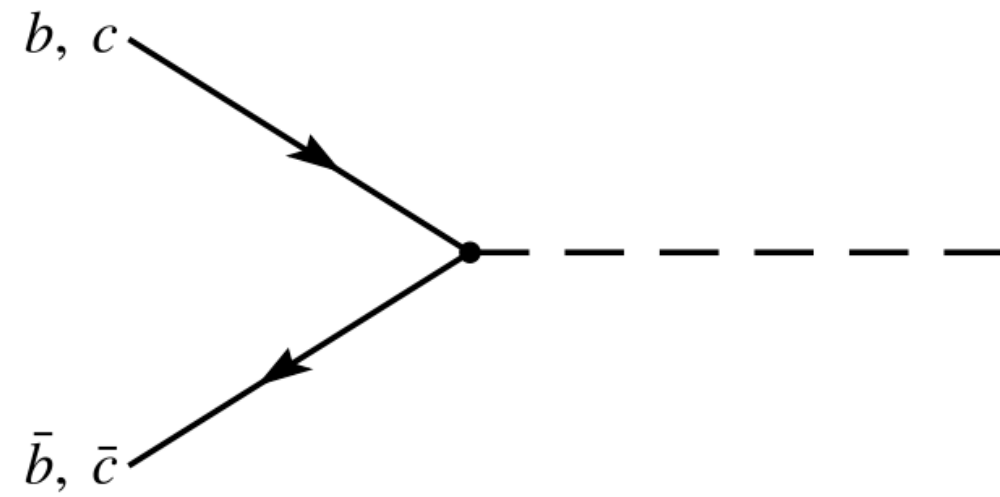
[arXiv:2207.08615](#)



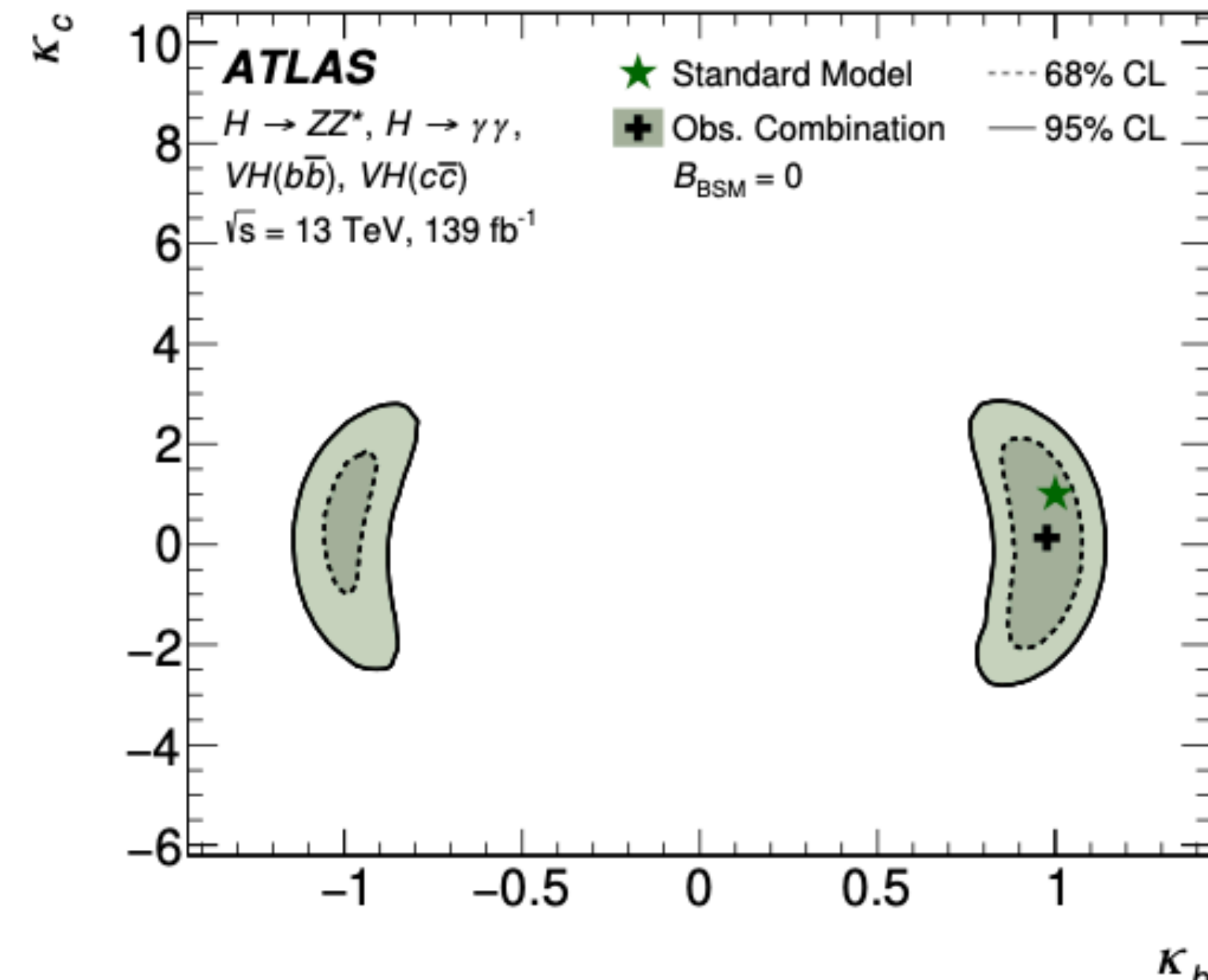
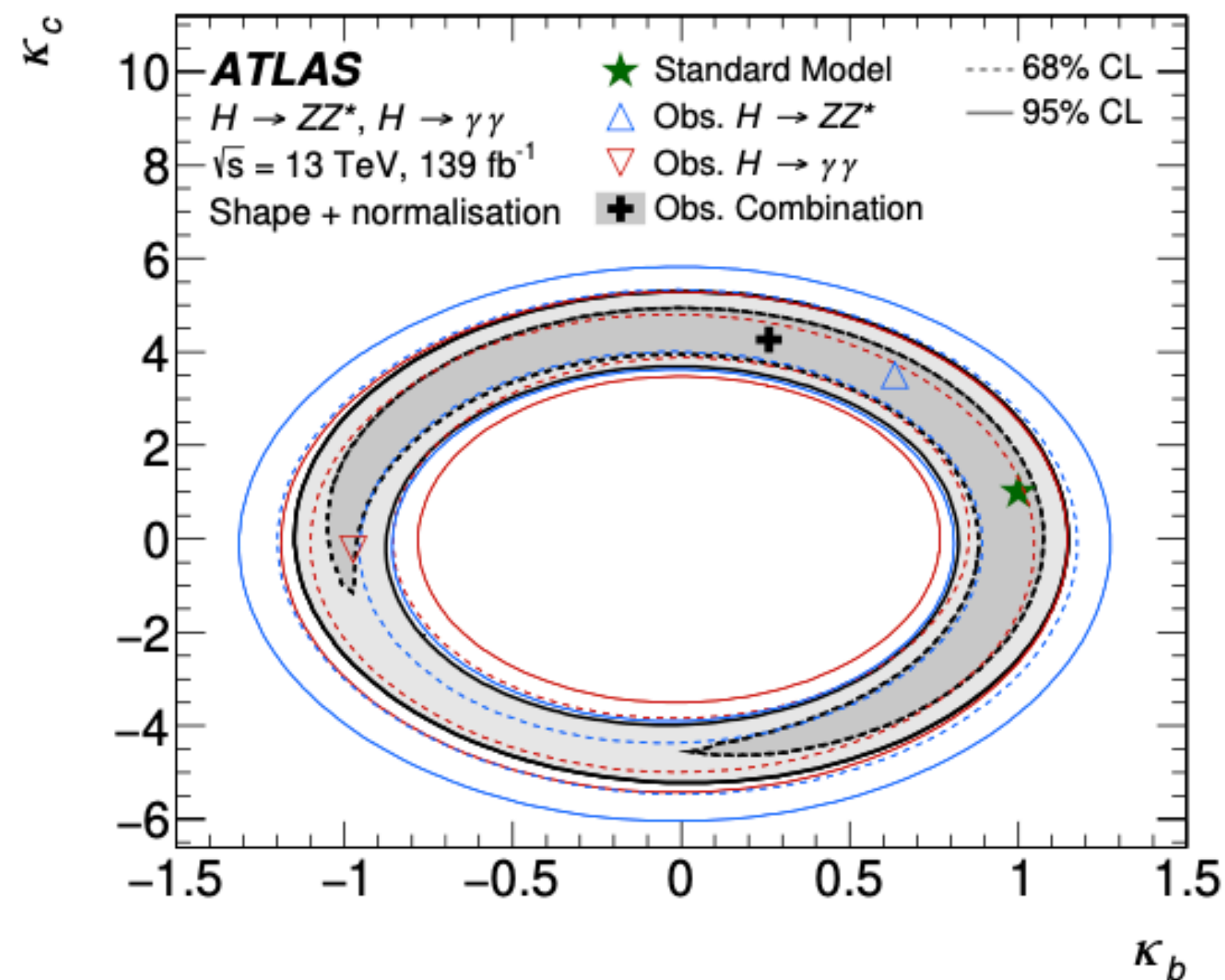
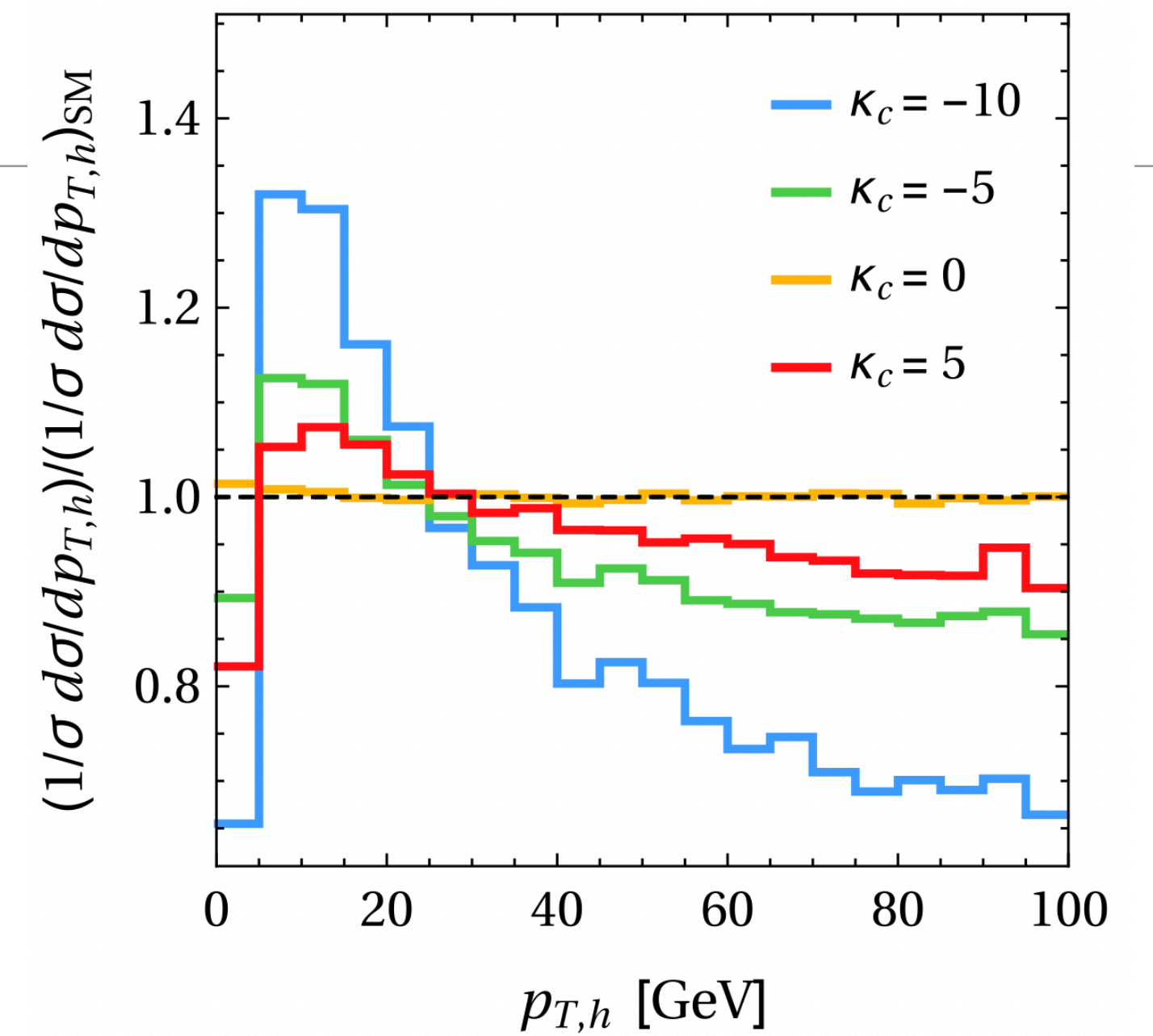
Good agreement with SM, measurements statistically limited

Differential cross sections: κ_b, κ_c interpretation

- Constrain Higgs b, c couplings indirectly from observed $p_T(H)$ spectra



[arXiv:2207.08615](https://arxiv.org/abs/2207.08615)



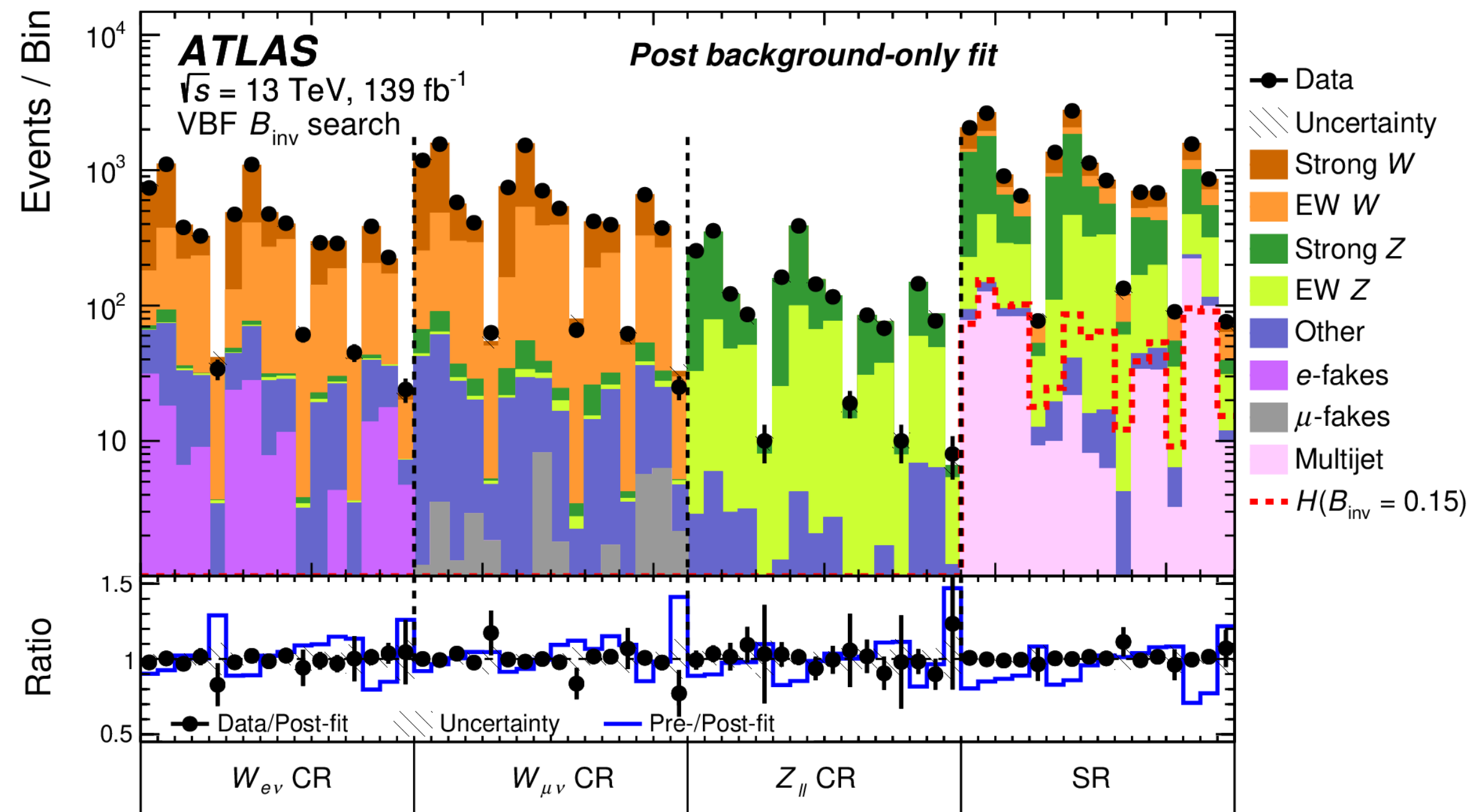
Constraints on κ_c similar to those from direct searches

Combination with direct constraints from $VH(qq)$ ($|\kappa_c| < 2.5$)

Exotic Higgs decays?

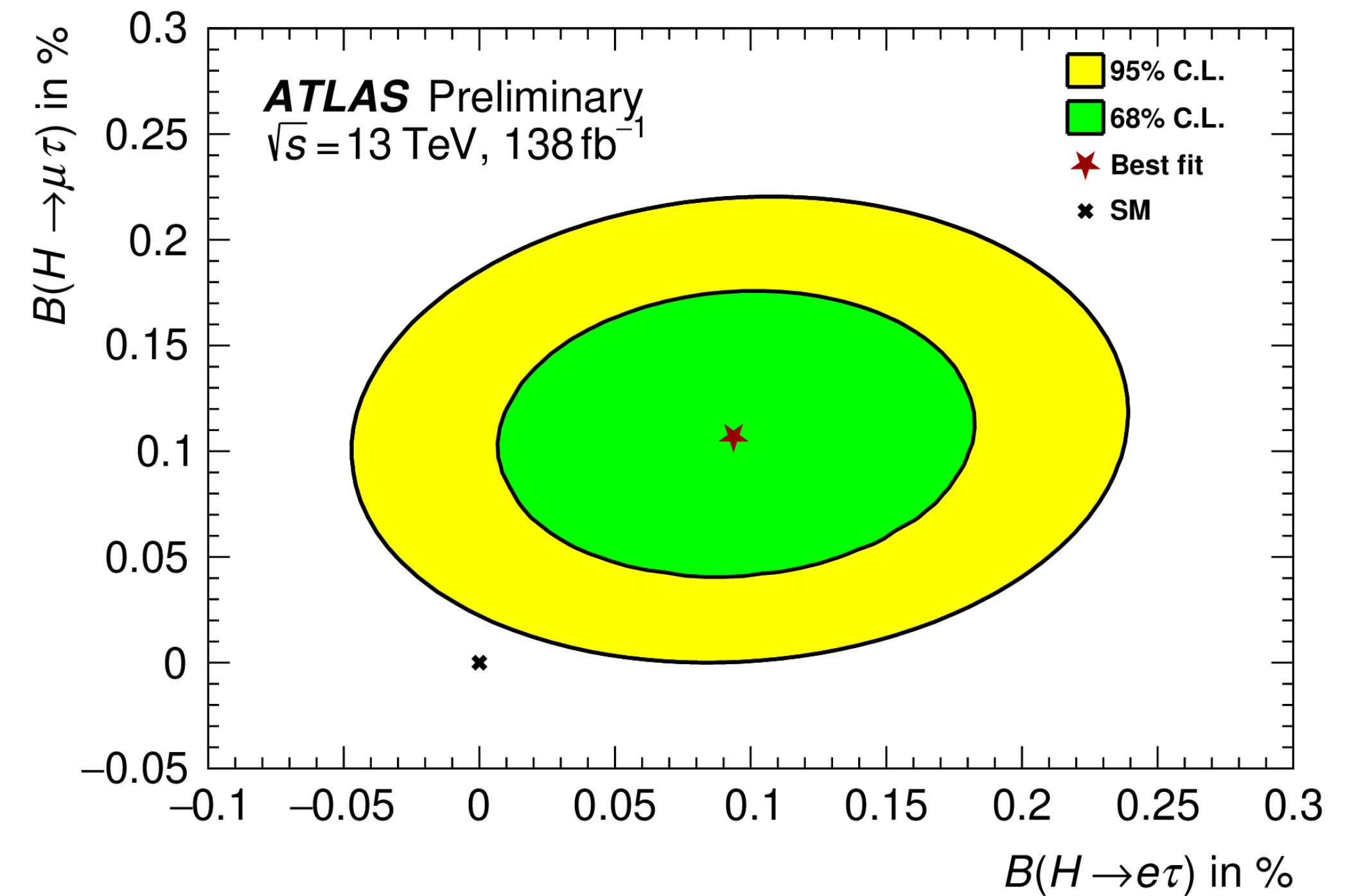
Invisible decays (e.g. to dark matter)

[JHEP 08 \(2022\) 104](#)



Lepton-flavour violating decays

[ATLAS-CONF-2022-060](#)

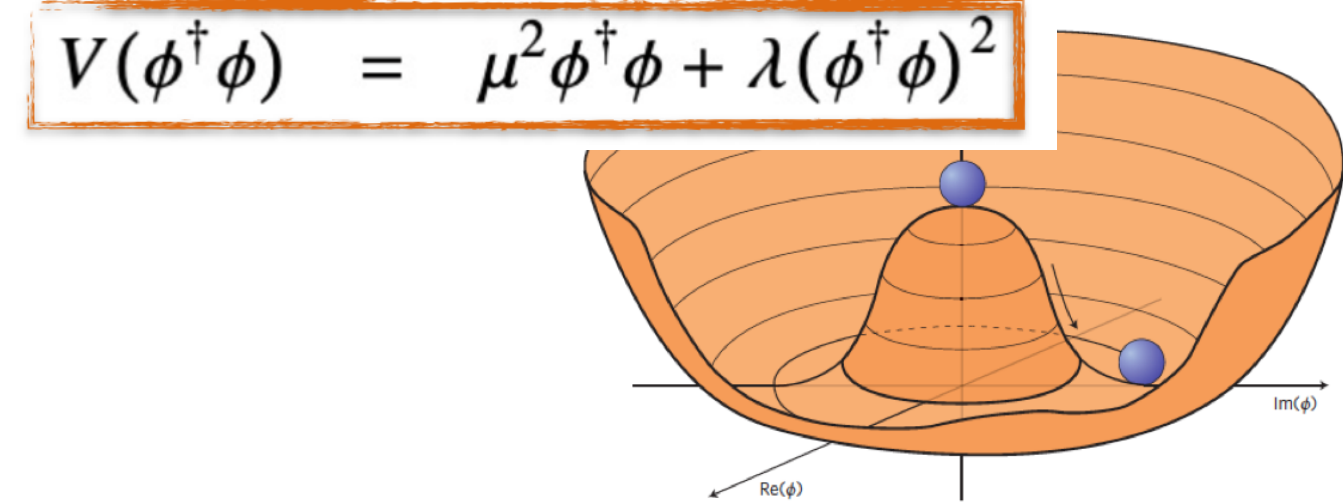
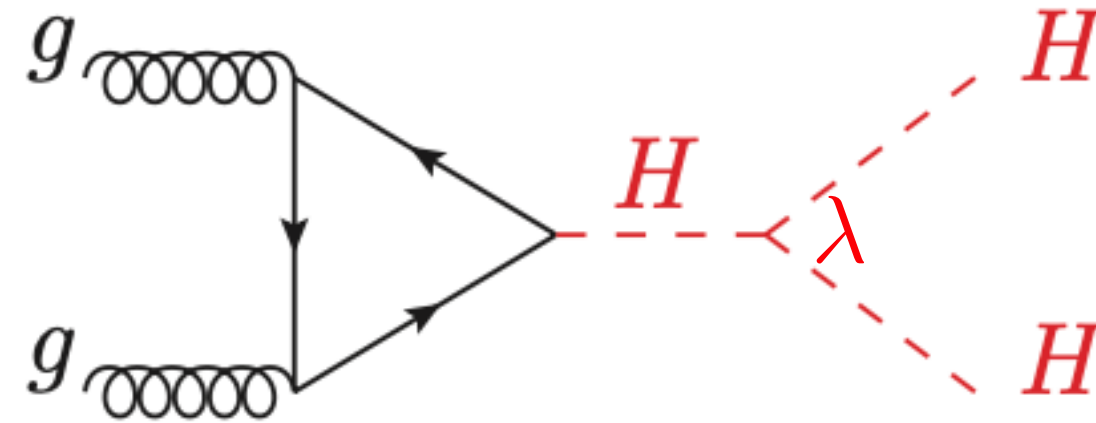


BR inv < 15% in VBF channel (Run2 only), ~11% including others

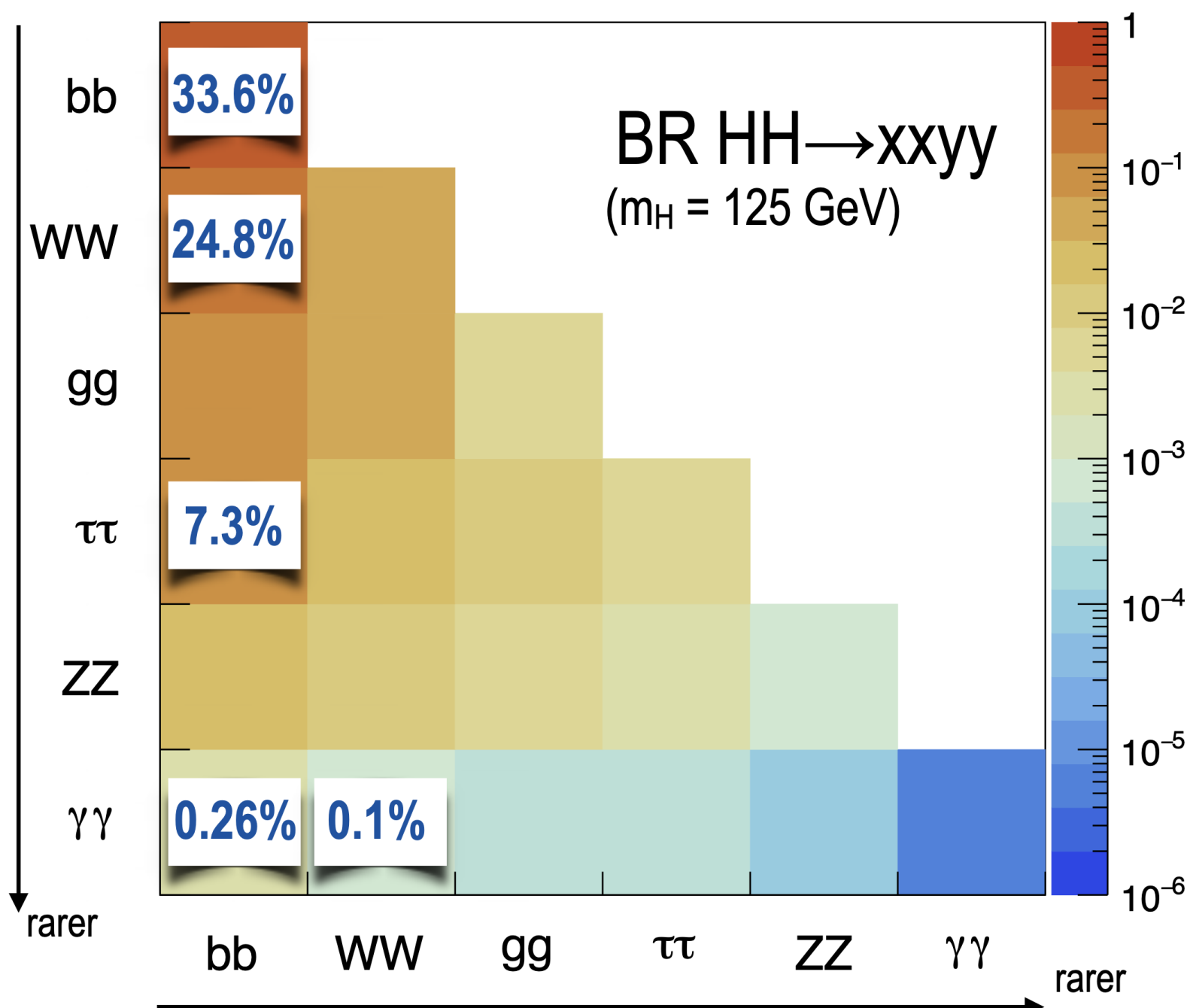
Upper limits at per mille level

Does the Higgs couple to itself?

- One of the novelties of Run2 was the **systematic search for di-Higgs production** in as many decay channels as possible
- Double-Higgs production = direct probe of Higgs self-coupling $\lambda \Rightarrow$ **crucial for determining shape of Higgs field potential**



- **Tiny cross section** $\sim 1/1000$ of Higgs production (33 fb at 13 TeV) \Rightarrow **extremely challenging!**
- **Multiple topologies investigated.** No significant signal seen yet \Rightarrow upper limits!



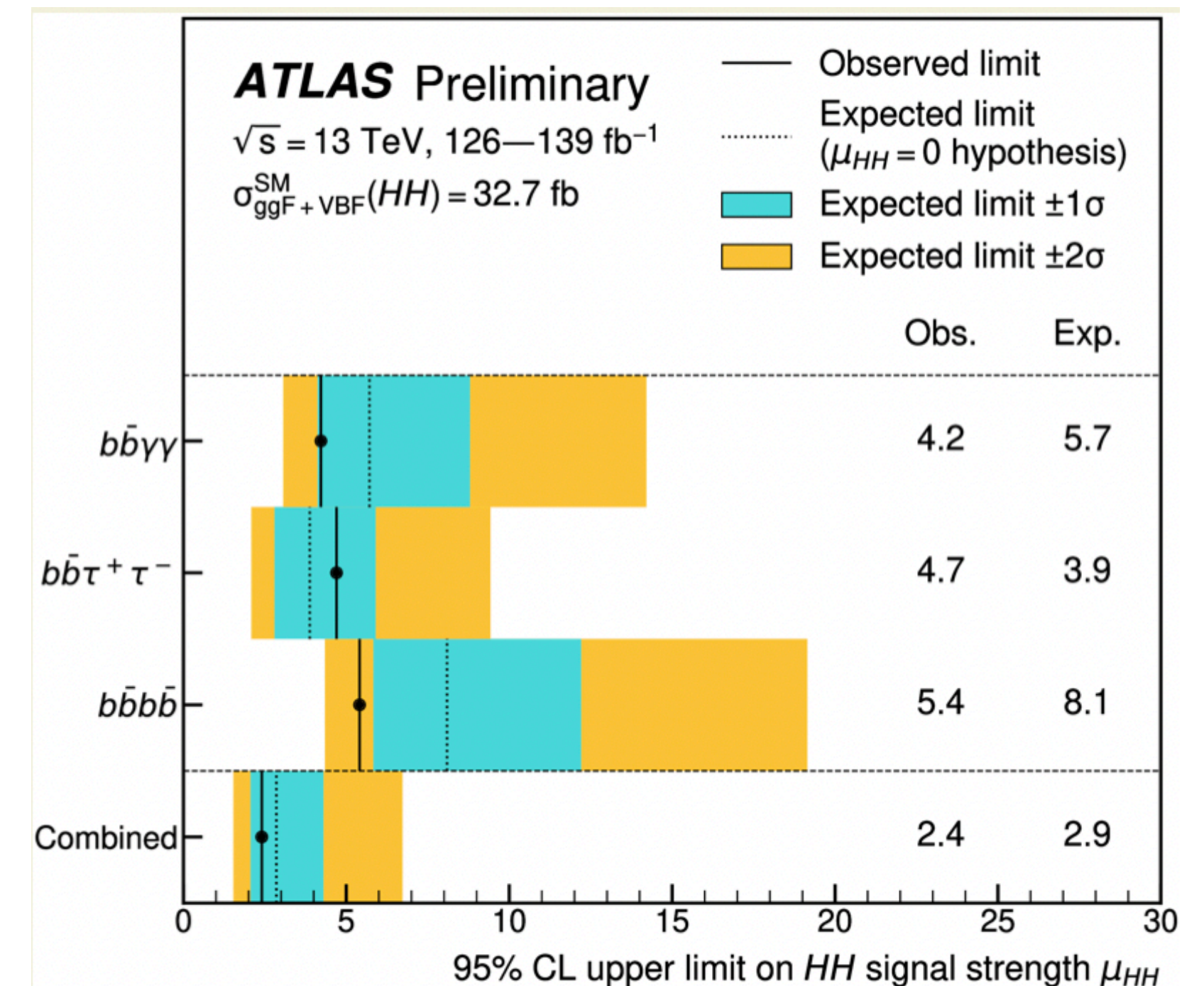
bbγγ/bbττ/4b

bbγγ: arXiv:2112.11876
bbbb: ATLAS-CONF-2022-035
bbττ: ATLAS-CONF-2021-030

combination

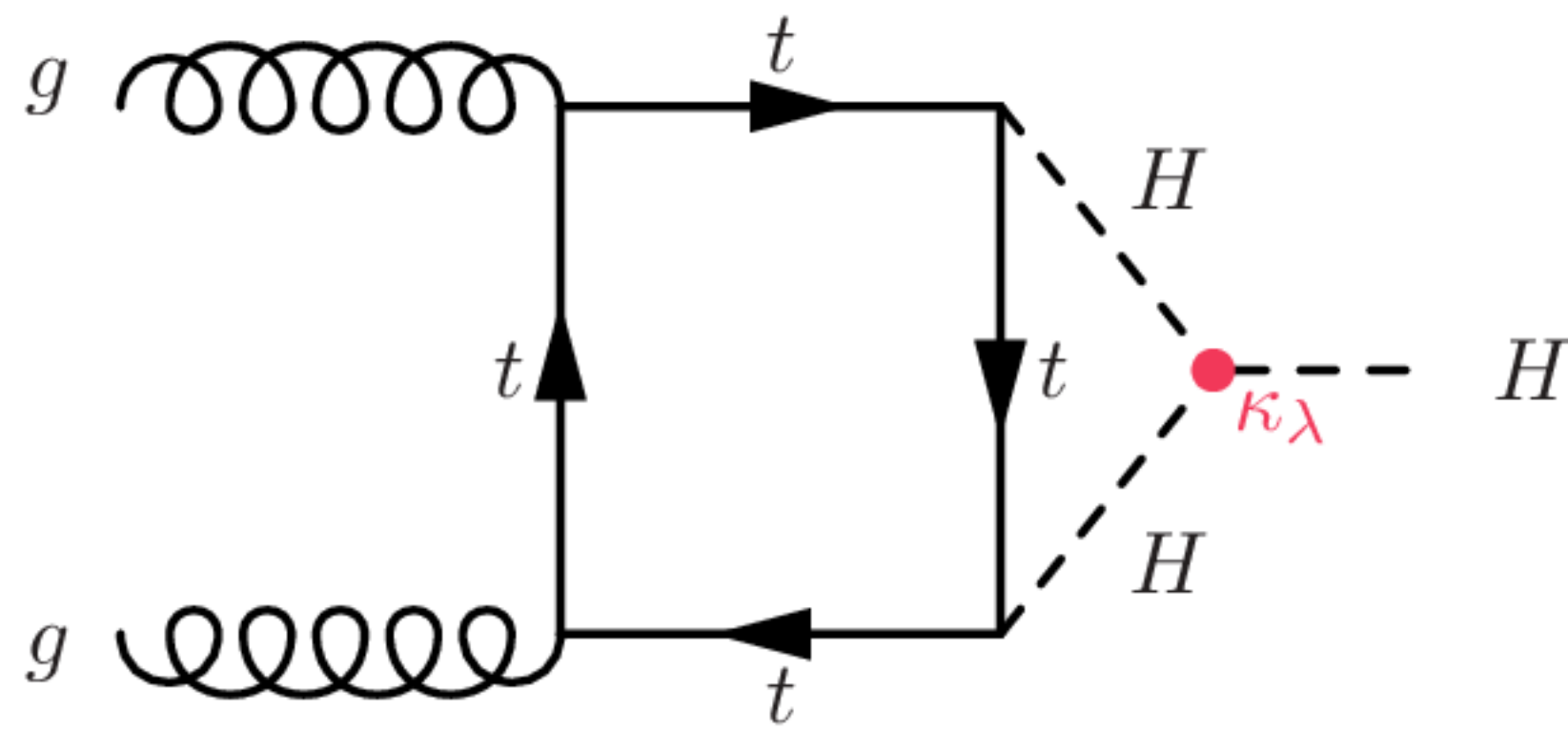
ATLAS-CONF-2022-050

$\mu < 2.4$ (2.9)



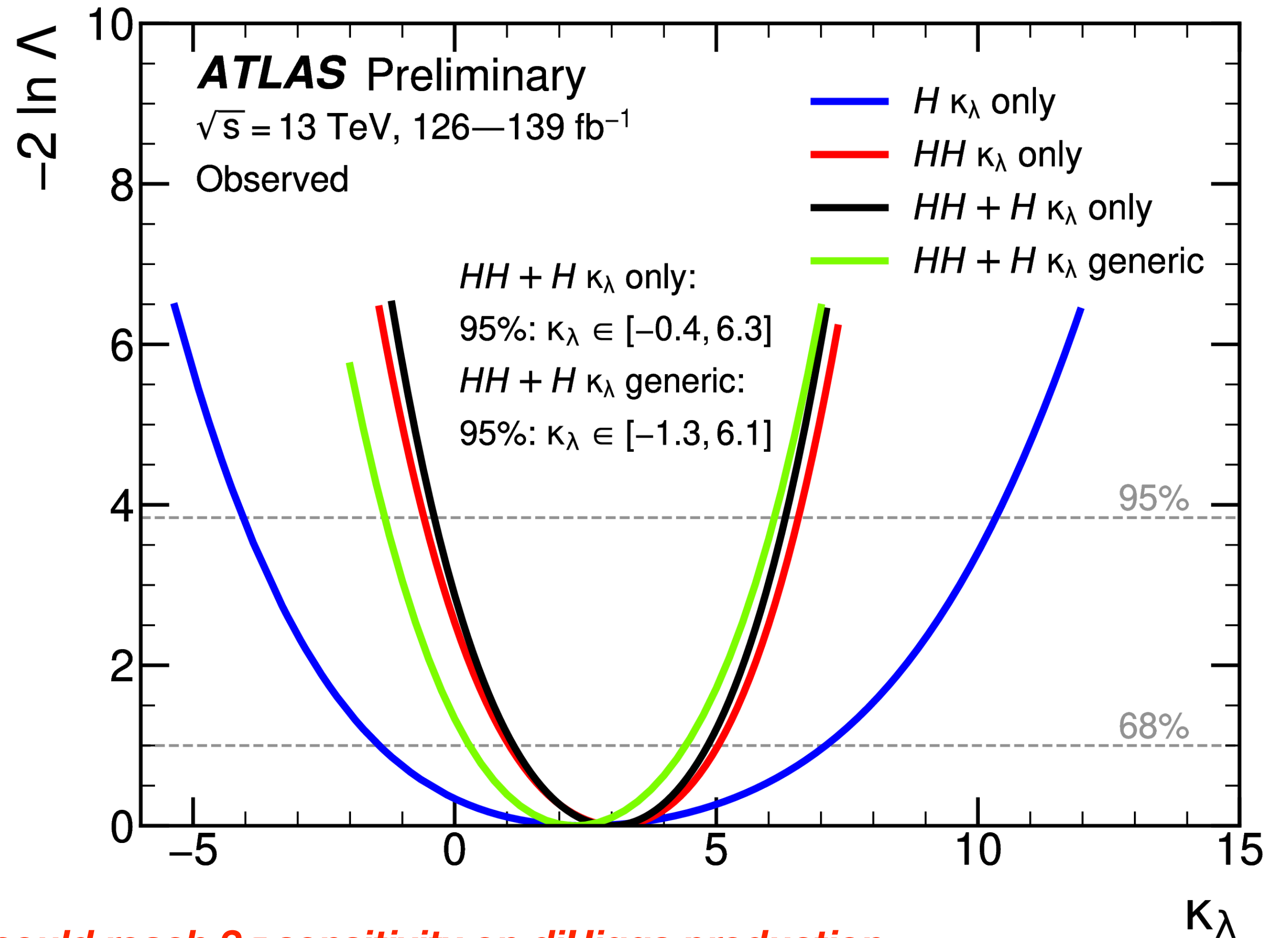
Higgs boson self-coupling

- Taking into account the λ dependence of the cross-section, limits on λ are determined
- Large correlation with other Higgs couplings (in particular to top) resolved by combining with single-Higgs measurements
 - Little increase in sensitivity on λ also provided by single-H xsection measurements from triple-Higgs vertex diagrams (NLO corrections)



Profile κ_λ only: $-0.4 < \kappa_\lambda < 6.3$ (95%CL)

Profile $\kappa_\lambda, \kappa_t, \kappa_V, \kappa_b, \kappa_\tau$: $-1.3 < \kappa_\lambda < 6.1$ (95% CL)



With the other channels + full Run 3 ATLAS+CMS could reach 2σ sensitivity on diHiggs production

Conclusion

- 10 years after its landmark discovery, the Higgs boson is now a grown-up kid
- Its characteristics resemble remarkably the SM predictions
 - Though accuracy of some measurement is still $O(10\%)$ or worse
- The Higgs physics programme of ATLAS and the (HL)-LHC is only in its infancy
 - Expect $\sim 3x$ the current data adding Run3 (2022-2025) and $20x$ by the end of the HL-LHC
 - On the watchlist (among others): rare and invisible decays, di-Higgs production (self-coupling), CP violation in Higgs interactions, ...



Thanks for listening and many thanks to the organisers!



Extra material

Inputs for the global coupling combination

Decay mode	Targeted production processes	\mathcal{L} [fb ⁻¹]	Ref.	Fits deployed in
$H \rightarrow \gamma\gamma$	ggF, VBF, $WH, ZH, t\bar{t}H, tH$	139	[31]	All
$H \rightarrow ZZ$	ggF, VBF, $WH + ZH, t\bar{t}H + tH$	139	[28]	All
	$t\bar{t}H + tH$ (multilepton)	36.1	[39]	All but fit of kinematics
$H \rightarrow WW$	ggF, VBF	139	[29]	All
	WH, ZH	36.1	[30]	All but fit of kinematics
	$t\bar{t}H + tH$ (multilepton)	36.1	[39]	All but fit of kinematics
$H \rightarrow Z\gamma$	inclusive	139	[32]	All but fit of kinematics
$H \rightarrow b\bar{b}$	WH, ZH	139	[33, 34]	All
	VBF	126	[35]	All
	$t\bar{t}H + tH$	139	[36]	All
	inclusive	139	[37]	Only for fit of kinematics
$H \rightarrow \tau\tau$	ggF, VBF, $WH + ZH, t\bar{t}H + tH$	139	[38]	All
	$t\bar{t}H + tH$ (multilepton)	36.1	[39]	All but fit of kinematics
$H \rightarrow \mu\mu$	ggF + $t\bar{t}H + tH, VBF + WH + ZH$	139	[40]	All but fit of kinematics
$H \rightarrow c\bar{c}$	$WH + ZH$	139	[41]	Only for free-floating κ_c
$H \rightarrow$ invisible	VBF	139	[42]	κ models with B_u & B_{inv} .
	ZH	139	[43]	κ models with B_u & B_{inv} .

Higgs boson selection efficiency

H(125 GeV) — approximate numbers

Channel	Produced		Selected	Mass resolution
$H \rightarrow \gamma\gamma$	18,200		6,440	1–2%
$H \rightarrow ZZ^*$	210,000	($\rightarrow 4\ell$)	210	1–2%
$H \rightarrow WW^*$	1,680,000	($\rightarrow 2\ell 2\nu$)	5,880	20%
$H \rightarrow \tau\tau$	490,000		2,380	15%
$H \rightarrow bb$	4,480,000		9,240	10%

A. Hoecker, CERN

- Out of 8M Higgs boson events only about 20k are selected and used to study the properties of the Higgs boson!

Higgs couplings fit in the presence of extra particles in loops or invisible decays

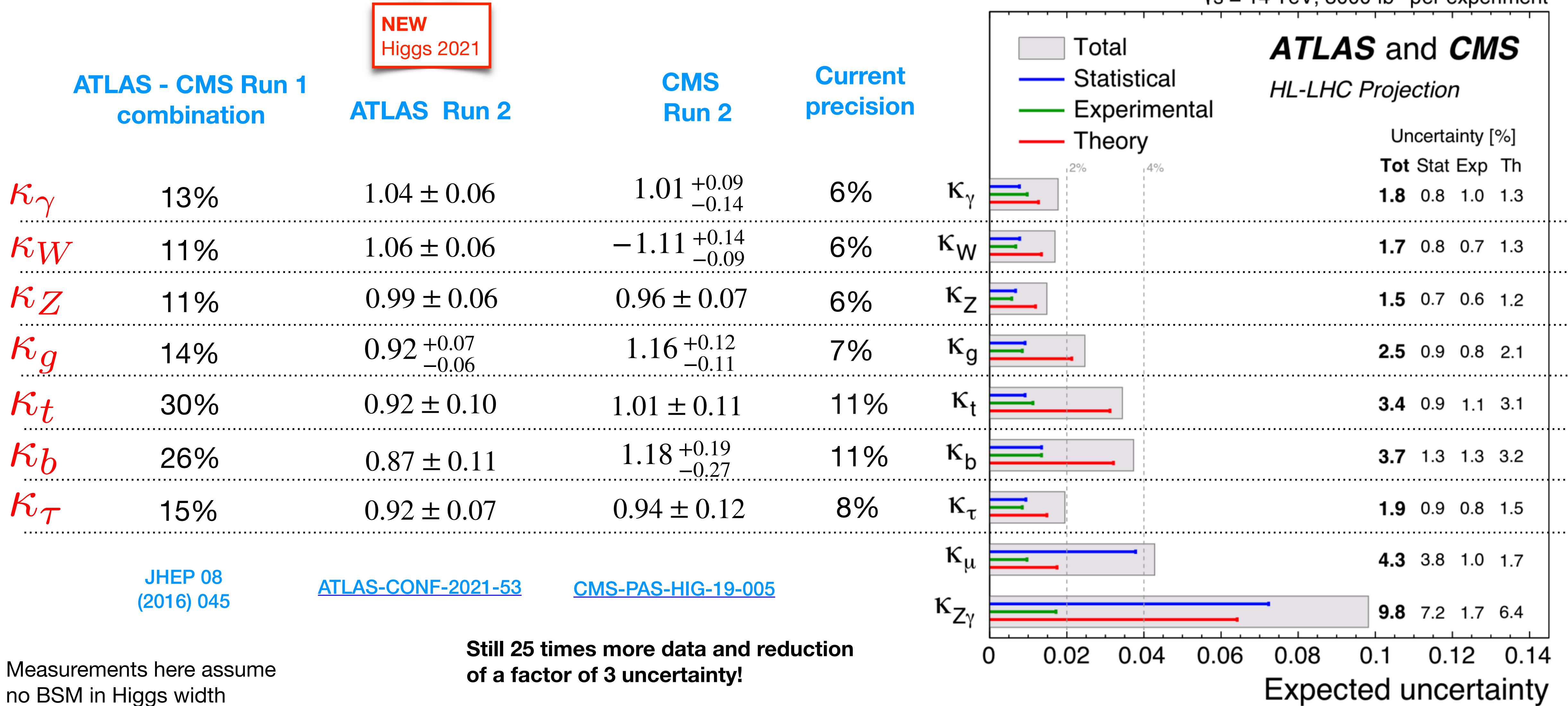
- Left: invisible/undetected BR fixed to zero
- Right: invisible/undetected BR floating
 - Use BR(inv) experimental upper limit
 - Assume $\kappa_W, \kappa_Z \leq 1$

	(a) $B_{inv.} = B_{u.} = 0$	(b) $B_{inv.}$ free, $B_{u.} \geq 0, \kappa_{W,Z} \leq 1$
κ_Z	$0.99^{+0.06}_{-0.06}$	$0.98^{+0.02}_{-0.05}$
κ_W	$1.05^{+0.06}_{-0.06}$	$1.00_{-0.02}$
κ_t	$0.94^{+0.11}_{-0.11}$	$0.94^{+0.11}_{-0.11}$
κ_b	$0.89^{+0.11}_{-0.11}$	$0.82^{+0.09}_{-0.08}$
κ_τ	$0.93^{+0.07}_{-0.07}$	$0.91^{+0.07}_{-0.06}$
κ_μ	$1.06^{+0.25}_{-0.30}$	$1.04^{+0.23}_{-0.30}$
κ_g	$0.95^{+0.07}_{-0.07}$	$0.94^{+0.07}_{-0.06}$
κ_γ	$1.01^{+0.06}_{-0.06}$	$0.98^{+0.05}_{-0.05}$
$\kappa_{Z\gamma}$	$1.38^{+0.31}_{-0.37}$	$1.35^{+0.29}_{-0.36}$
$B_{inv.}$	-	< 0.13
$B_{u.}$	-	< 0.12

HL-LHC Higgs couplings projections

M. Kado, Higgs 2021

$\sqrt{s} = 14 \text{ TeV}, 3000 \text{ fb}^{-1}$ per experiment



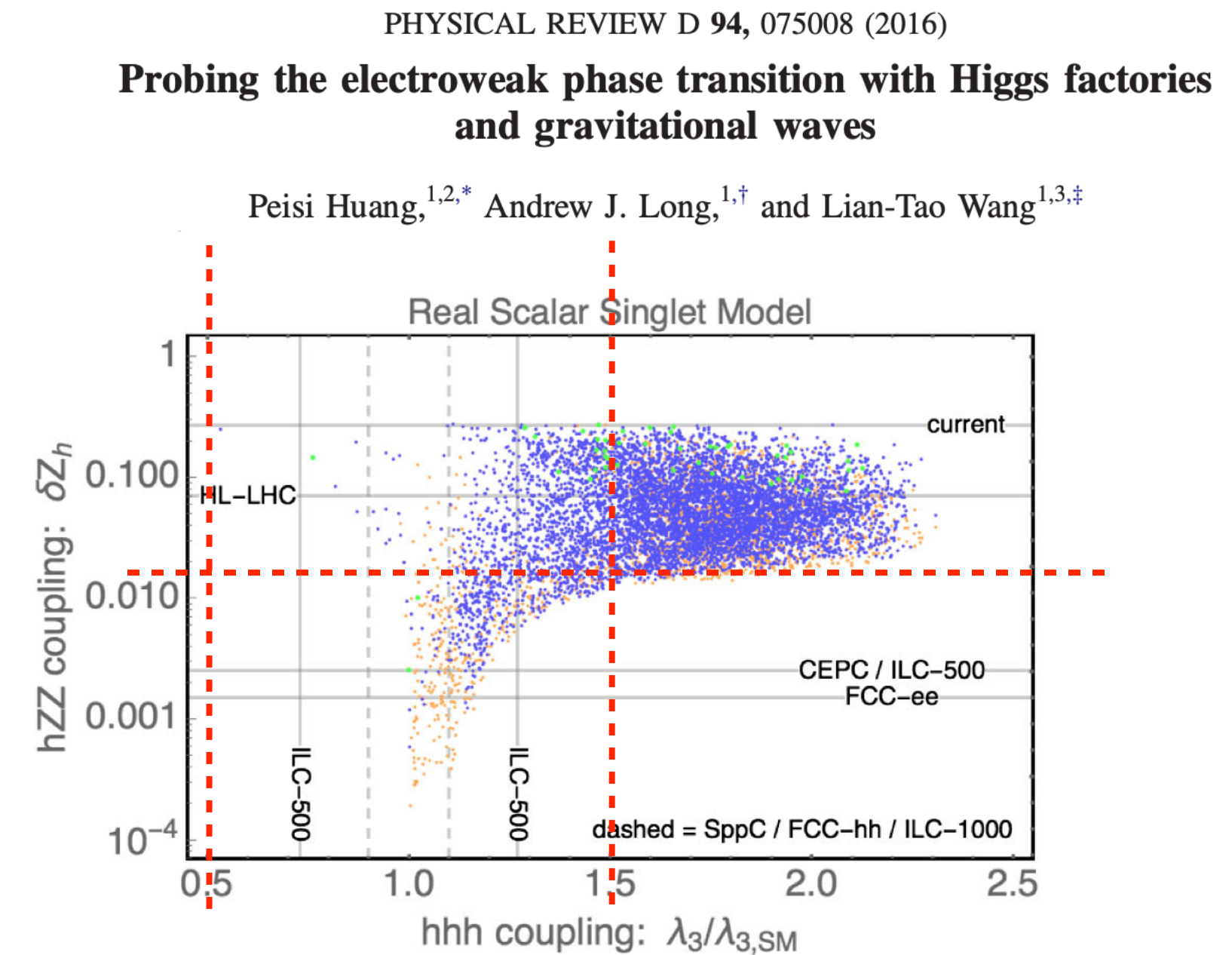
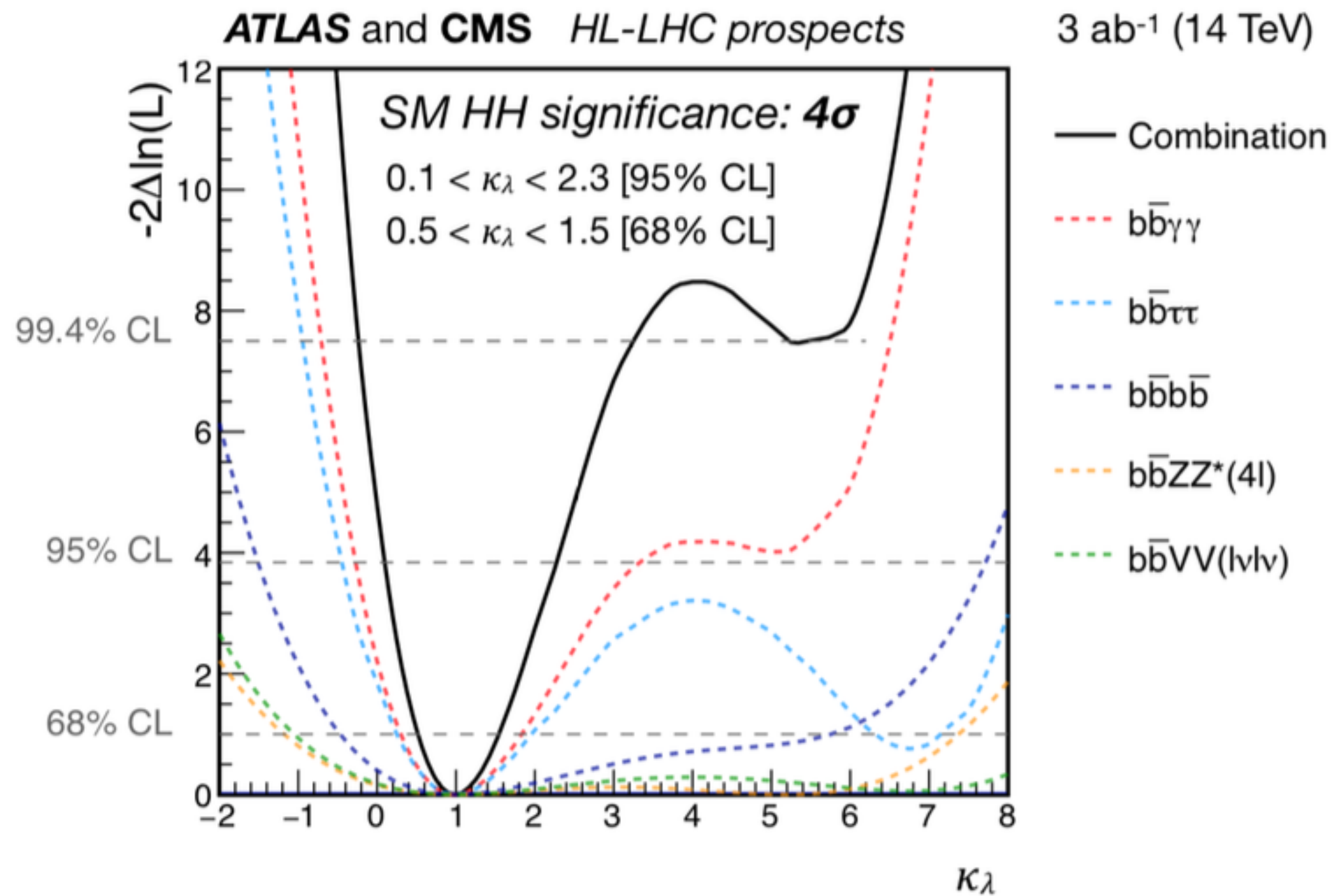
Measurements here assume no BSM in Higgs width

Still 25 times more data and reduction of a factor of 3 uncertainty!

Higgs boson self-coupling

- HL-LHC projection (~2038, ~20x more data than Run2): $0.5 < \kappa_\lambda < 1.5$
 - could constrain models which predict strong first-order electroweak phase transitions
 - complementary to information provided by gravitational waves detected by space-based interferometers

[Phys. Rev. D 94 \(2016\) 7, 075008](#)

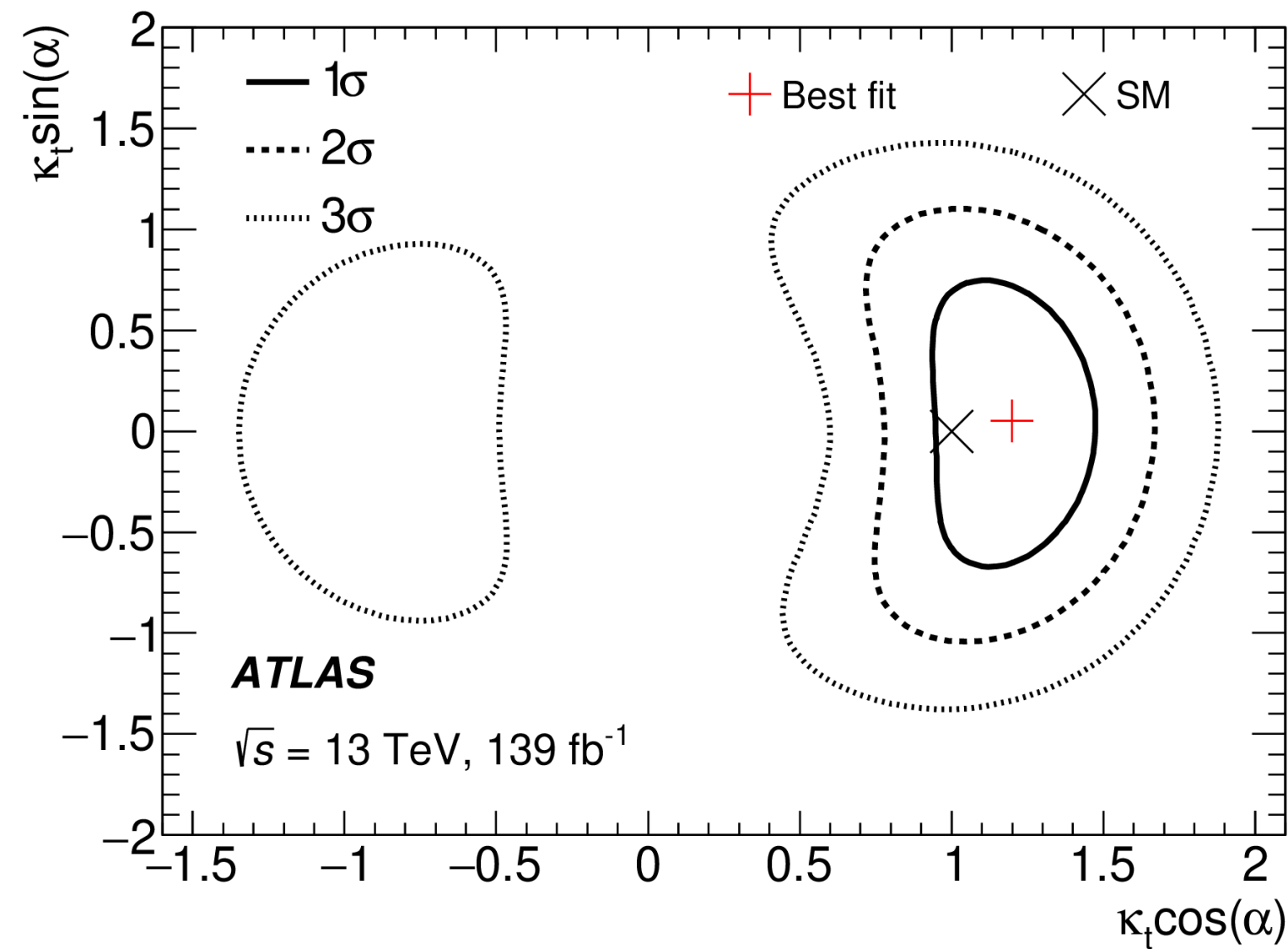


l. Parameter space scan for the singlet model of II A. An orange point indicates a first-order phase transition, a blue point indicates a strongly first-order phase transition (3.4), and a green point indicates a very strong first-order phase transition with potentially detectable gravitational wave signal at eLISA. The right panels shows the predicted gravitational wave spectrum today along with the projected sensitivity of eLISA [2].

Higgs boson CP

- CPV in t-H interactions with ttH($\gamma\gamma$): [Phys. Rev. Lett. 125, 061802 \(2020\)](#)

$$\mathcal{L}_{t\bar{t}H} = -\kappa'_t y_t \phi \bar{\psi}_t (\cos \alpha + i\gamma_5 \sin \alpha) \psi_t$$

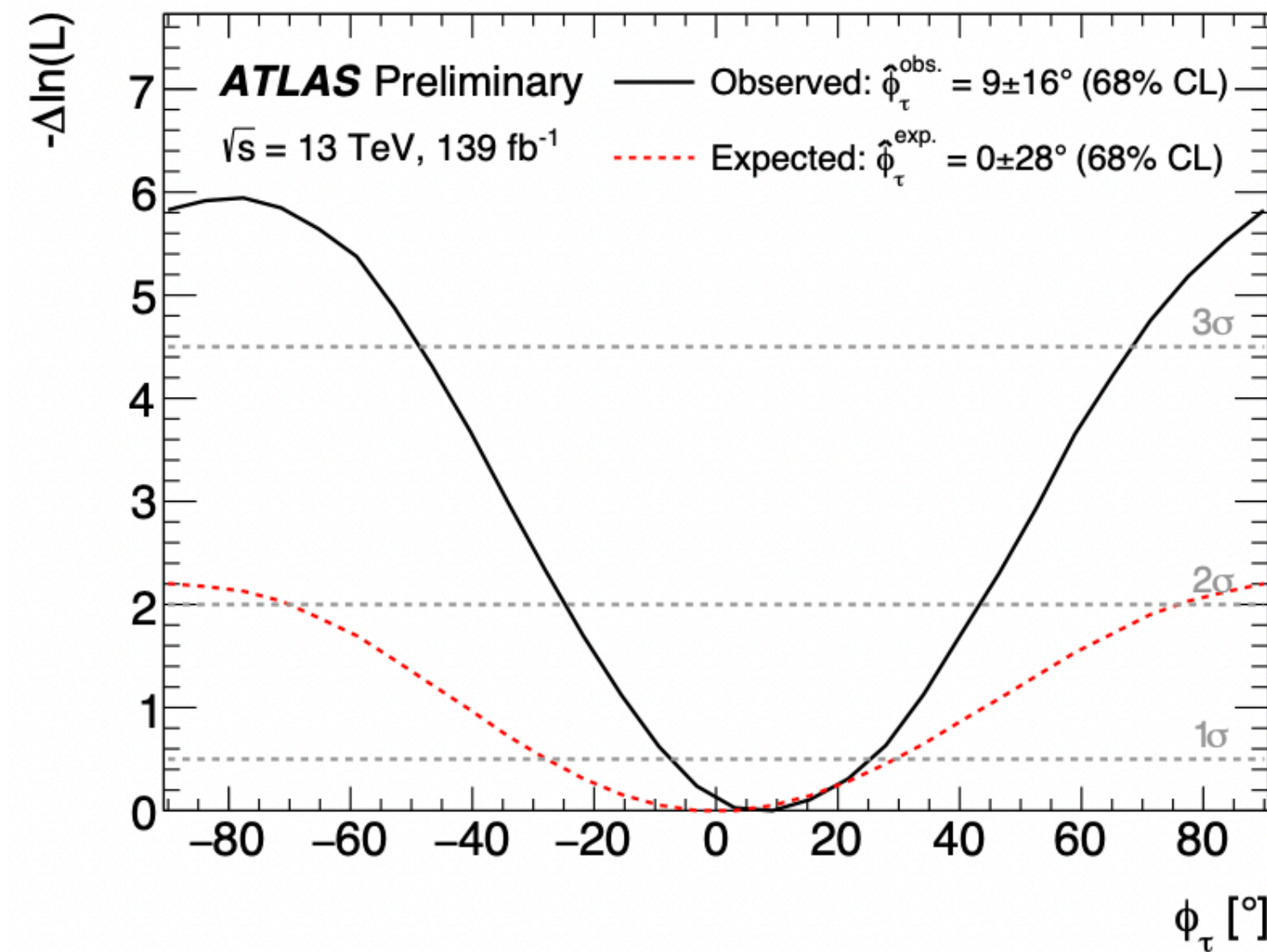


pure CP-odd coupling
excluded at 3.9σ

$|\alpha| > 43$ deg excluded @ 95% CL.

- CPV in τ -H interactions with $H \rightarrow \tau\tau$: [ATLAS-CONF-2022-032](#)

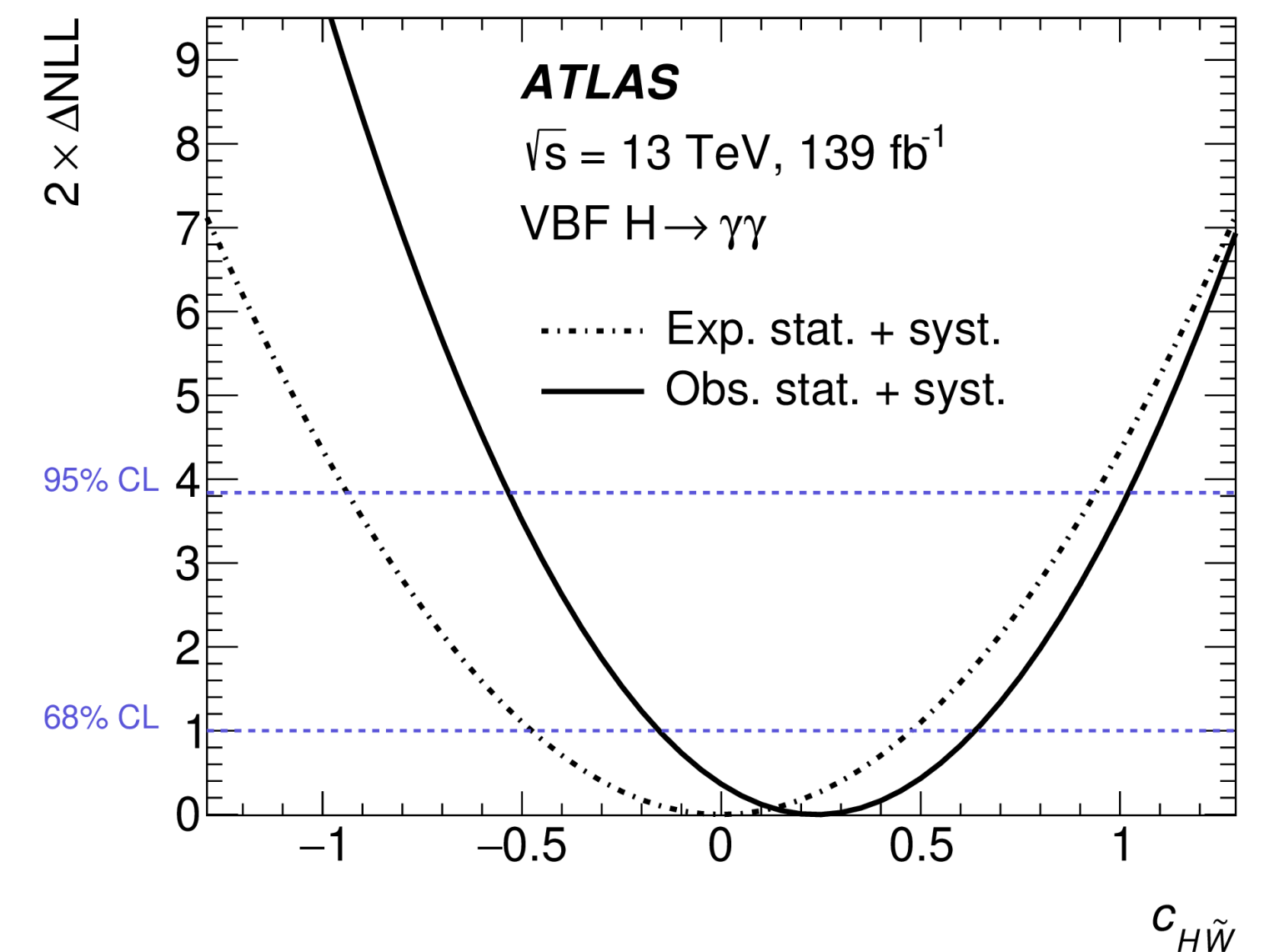
$$\mathcal{L}_{H\tau\tau} = -\frac{m_\tau}{\nu} \kappa_\tau (\cos \phi_\tau \bar{\tau}\tau + \sin \phi_\tau \bar{\tau}i\gamma_5\tau) H$$



$\phi_\tau = 9 \pm 16^\circ$ (28° exp)

pure CP-odd coupling
excluded at 3.4σ

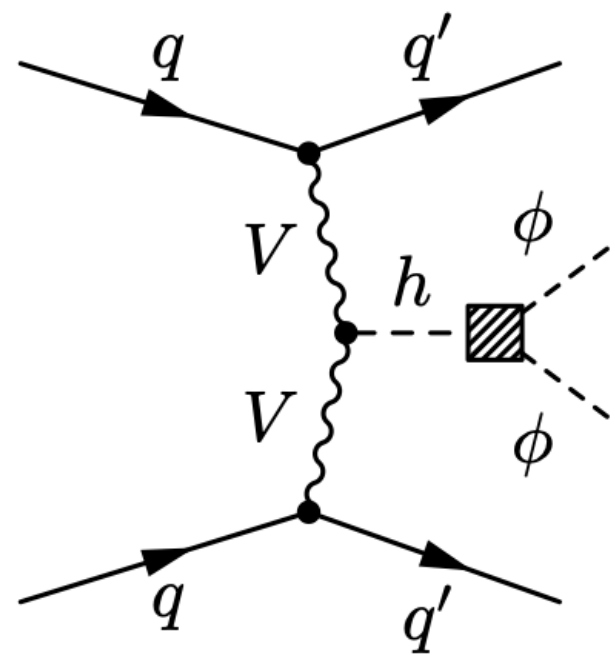
- CPV in HVV interactions with VBF: [arXiv:2208.02338](#)



	68% (exp.)	95% (exp.)	68% (obs.)	95% (obs.)
$c_{H\tilde{W}}$ (inter. only)	[-0.48, 0.48]	[-0.94, 0.94]	[-0.16, 0.64]	[-0.53, 1.02]
$c_{H\tilde{W}}$ (inter.+quad.)	[-0.48, 0.48]	[-0.95, 0.95]	[-0.15, 0.67]	[-0.55, 1.07]

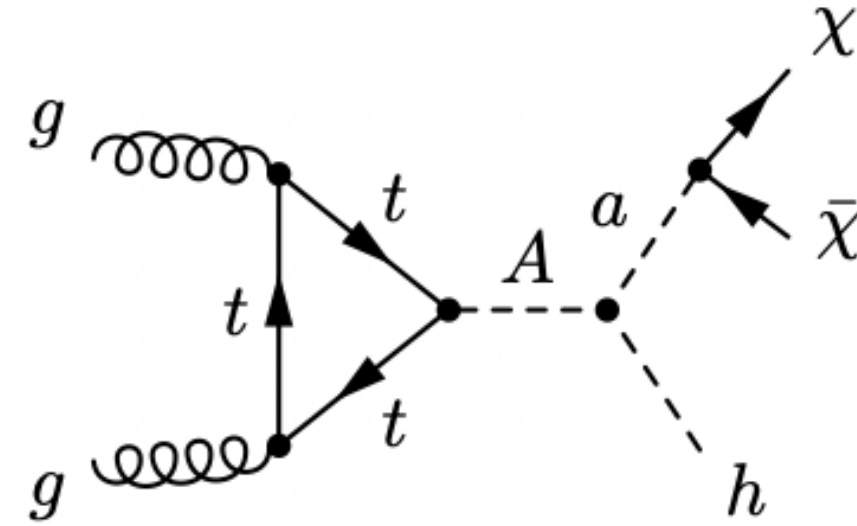
Higgs boson as a portal to dark matter?

- IF dark matter is composed of neutral, weakly-interacting, massive particle (WIMP), and the dark sector is coupled to the Higgs, Higgs physics at collider experiments could shed light on it (for a review see e.g. <https://arxiv.org/abs/2109.13597>)

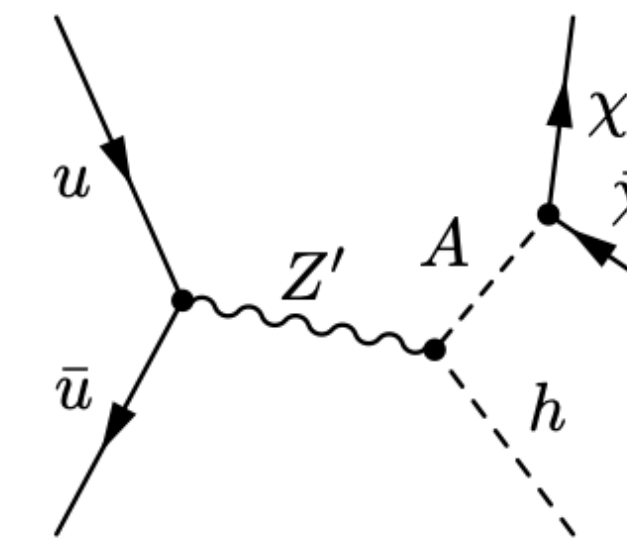


(a) Higgs portal ($m_{DM} < m_H/2$)

invisible Higgs decays
 $\Rightarrow BR(H \rightarrow \text{invisible})$



(b) Portals with extended Higgs sectors
 (e.g.: 2HDM+a)

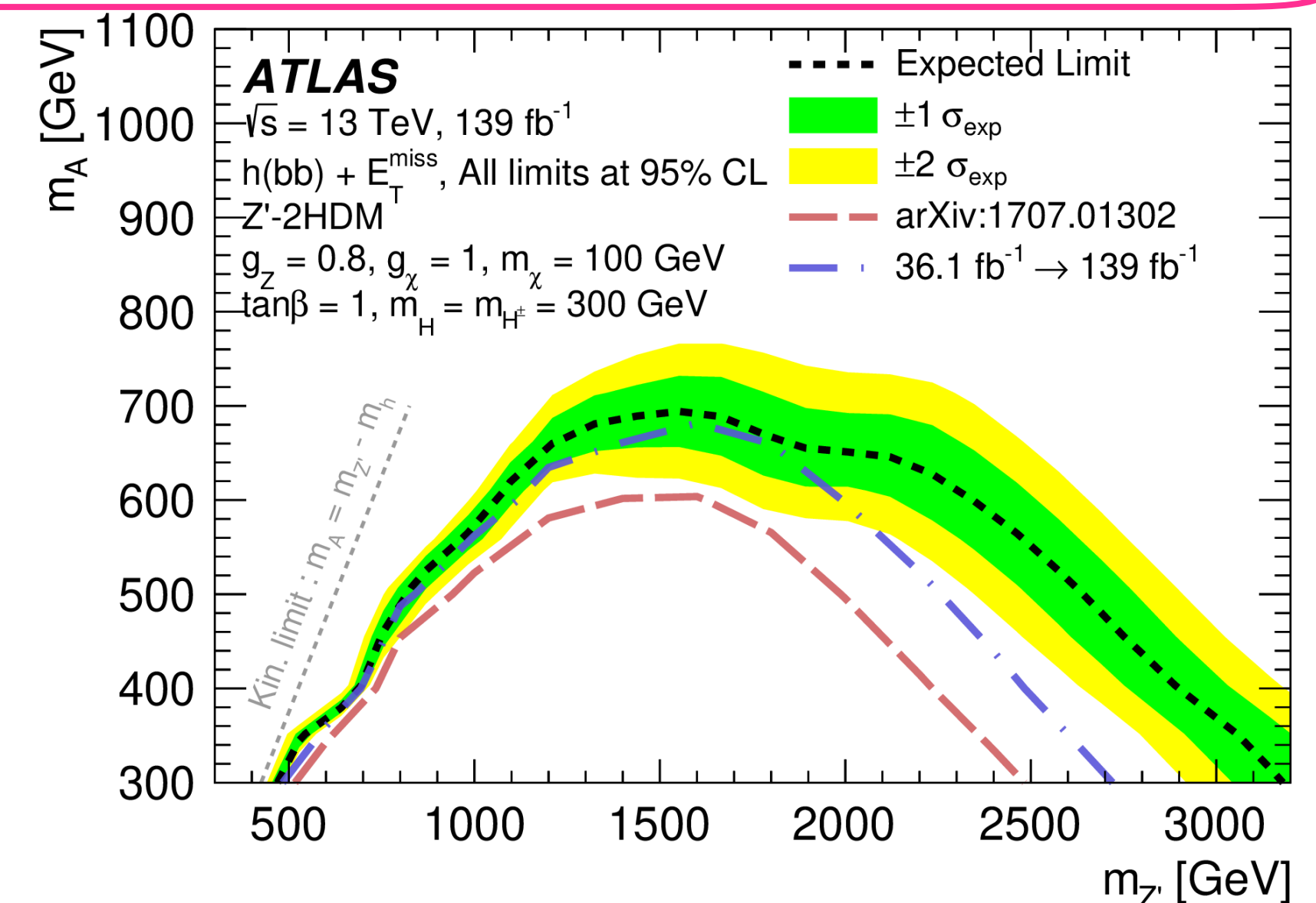
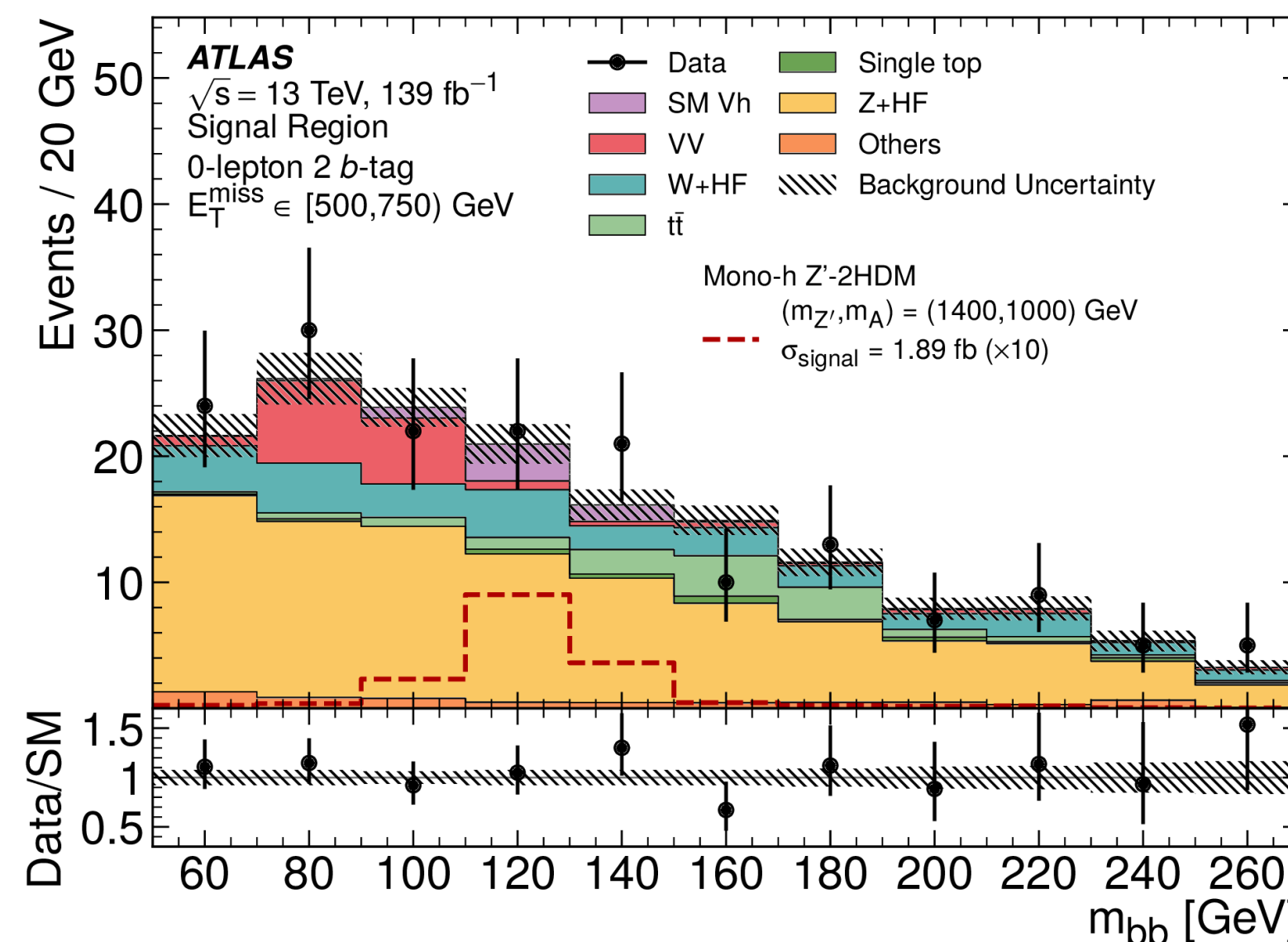


(c) Portals with extended Higgs and gauge sectors
 (e.g.: 2HDM+Z')

“mono-Higgs”: Higgs recoils against invisible particles (large missing momentum)
 (+ searches for heavier Higgs bosons and Z' resonances)

mono-H(bb)

[JHEP 11 \(2021\) 209](https://arxiv.org/abs/2109.13597)



Higgs invisible decays and WIMP-nucleon cross-section

Higgs portal model: the Higgs boson is assumed to be the only mediator in the WIMP–nucleon scattering
 The upper limit on the invisible BR is converted to an upper limit on the partial width for the decay:

$$\Gamma_H^{\text{inv}} = \frac{\text{BF}(H \rightarrow \text{invisible})}{1 - \text{BF}(H \rightarrow \text{invisible})} \times \Gamma_H,$$

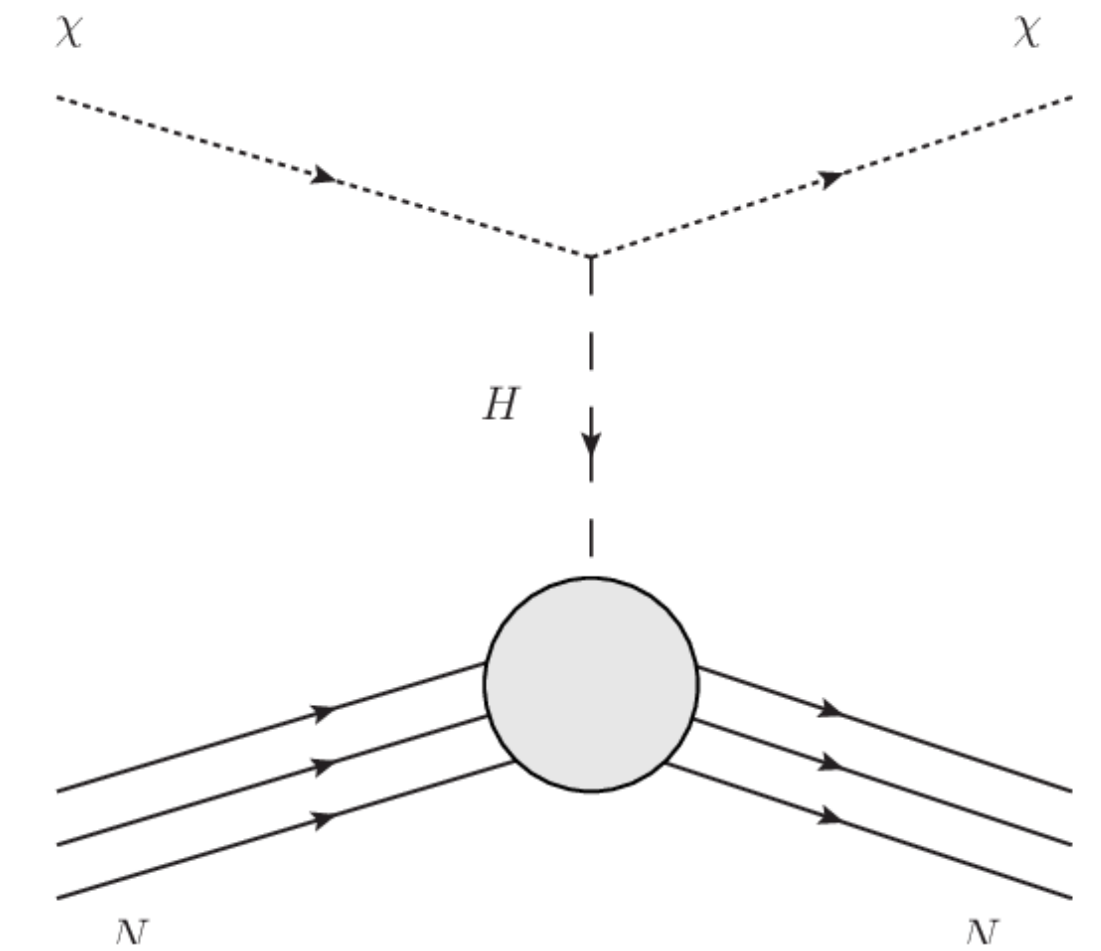
This in turn implies an upper limit on the Higgs-DM coupling λ :

$$\Gamma_{H \rightarrow SS}^{\text{inv}} = \frac{\lambda_{HSS}^2 v^2 \beta_S}{64\pi m_H},$$

$$\Gamma_{H \rightarrow VV}^{\text{inv}} = \frac{\lambda_{HVV}^2 v^2 m_H^3 \beta_V}{256\pi m_V^4} \left(1 - 4 \frac{m_V^2}{m_H^2} + 12 \frac{m_V^4}{m_H^4} \right),$$

$$\Gamma_{H \rightarrow ff}^{\text{inv}} = \frac{\lambda_{Hff}^2 v^2 m_H \beta_f^3}{32\pi \Lambda^2},$$

$$\beta_\chi = \sqrt{1 - 4m_\chi^2/m_H^2} \quad (\chi = S, V, f),$$



which can then be converted into a limit on the WIMP-proton cross-section:

$$\sigma_{SN}^{\text{SI}} = \frac{\lambda_{HSS}^2}{16\pi m_H^4} \frac{m_N^4 f_N^2}{(m_S + m_N)^2},$$

$$\sigma_{VN}^{\text{SI}} = \frac{\lambda_{HVV}^2}{16\pi m_H^4} \frac{m_N^4 f_N^2}{(m_V + m_N)^2},$$

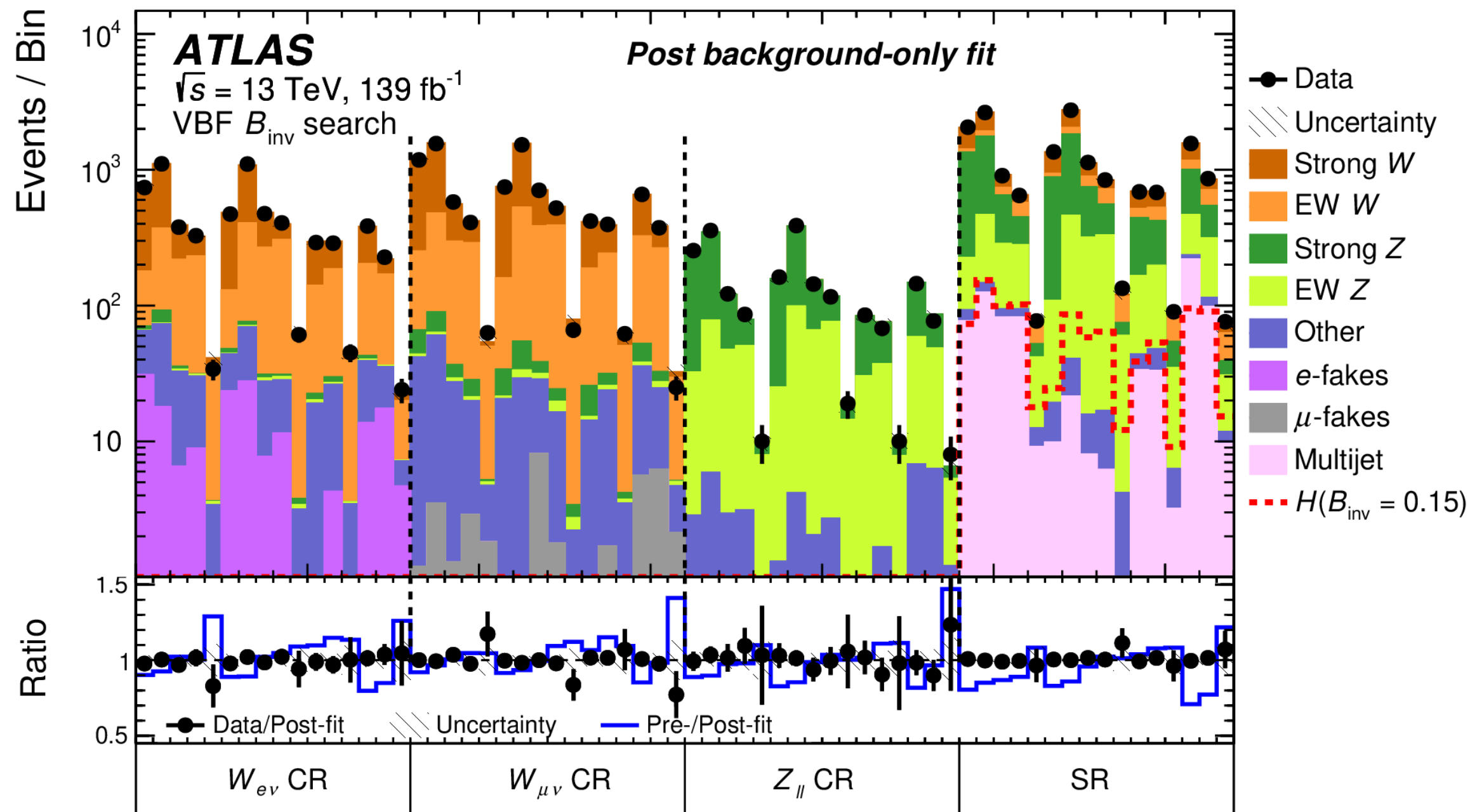
$$\sigma_{fN}^{\text{SI}} = \frac{\lambda_{Hff}^2}{4\pi \Lambda^2 m_H^4} \frac{m_N^4 m_f^2 f_N^2}{(m_f + m_N)^2},$$

Vacuum expectation value	$v/\sqrt{2}$	174 GeV
Higgs boson mass	m_H	125 GeV
Higgs boson width	Γ_H	4.07 MeV
Nucleon mass	m_N	939 MeV
Higgs–nucleon coupling form factor	f_N	$0.33^{+0.30}_{-0.07}$

Higgs boson invisible decays

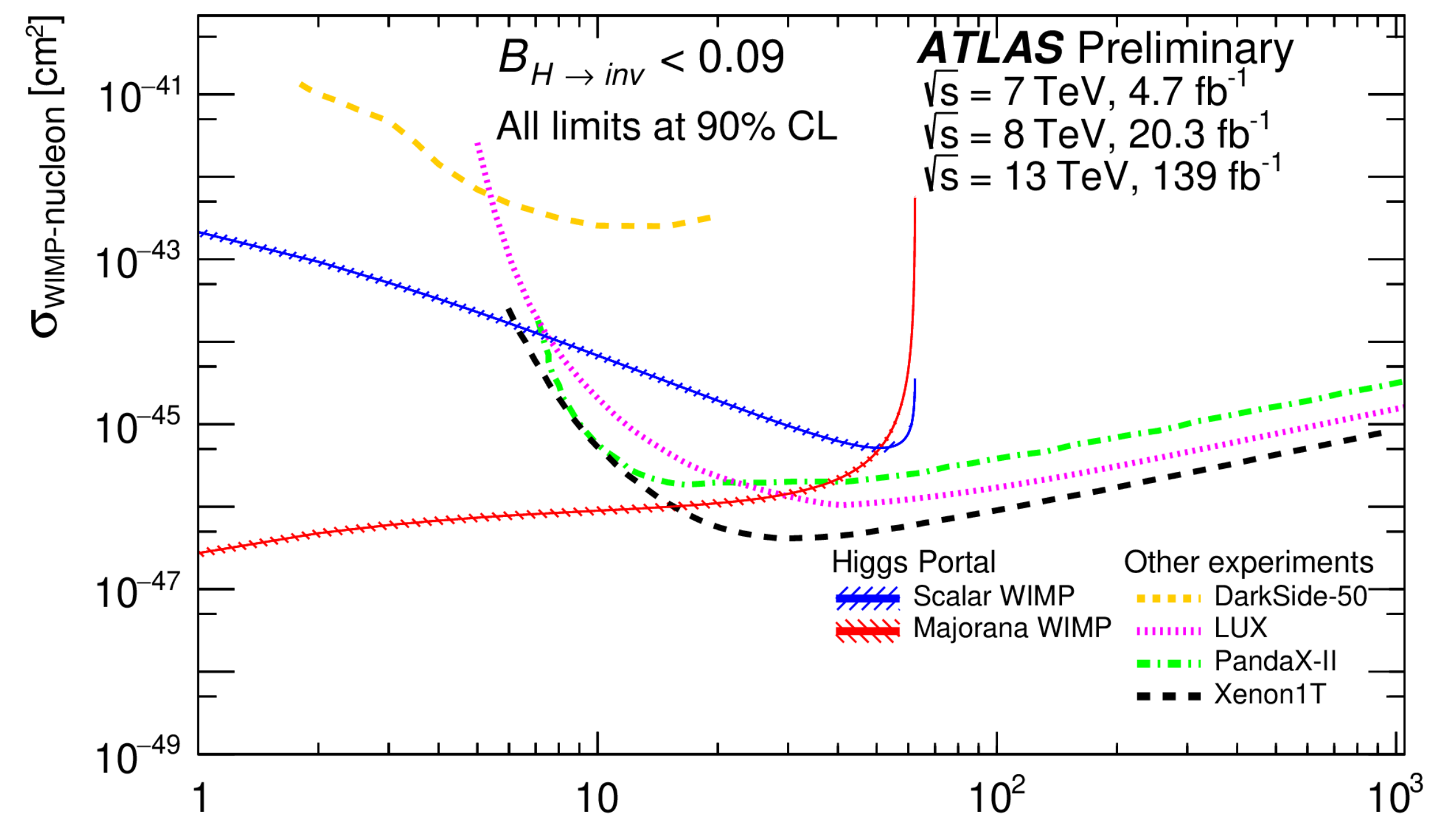
- $BR(H \rightarrow \text{invisible}) = 0.1\%$ in SM ($H \rightarrow ZZ^{(*)} \rightarrow 4\nu$), too small for detection (poor missing momentum resolution, backgrounds from $Z \rightarrow \nu\nu$)
- If Higgs decays to dark matter particles, $BR(H \rightarrow \text{invisible})$ could be enhanced to a detectable level
- Most sensitive signature: VBF-tagged jets + large missing momentum
- Further combination with other searches (such as $V(\text{lep})H$, $V(\text{had})H$, ..)

VBF $H \rightarrow \text{invisible}$ (Run2) [JHEP 08 \(2022\) 104](#)



$BR(H \rightarrow \text{invisible}) < 14.5\% @ 95\% CL$

Combination [ATLAS-CONF-2020-052](#)



$BR < 11\% @ 95\% CL$ ($< 9\%$ at 90%CL) $m_{\text{WIMP}} [\text{GeV}]$
Constraints on spin-independent WIMP-nucleon x-section complementary to direct searches