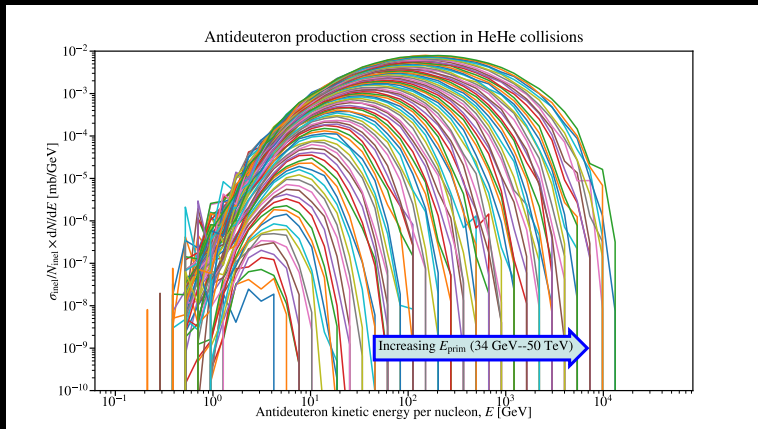


Antinuclei in the Milky Way

Michael Kachelrieß (NTNU, Trondheim)



with Jonas Tjemsland and Sergey Ostapchenko

Eur.Phys.J.A 56 (2020) 1, JCAP 08 (2020) 048, Eur.Phys.J.A 57 (2021) 5, 167, PoS TOOLS2020 (2021) 006,...

Outline of the talk

1 Introduction

- ▶ Motivation: **why antinuclei?**
 - ★ Probe of quark-gluon plasma
 - ★ **Signature of dark matter**
- ▶ Physical basis of **coalescence approach**

2 Coalescence models and antinuclei production

- ▶ Coalescence in momentum space
- ▶ Coalescence in phase space

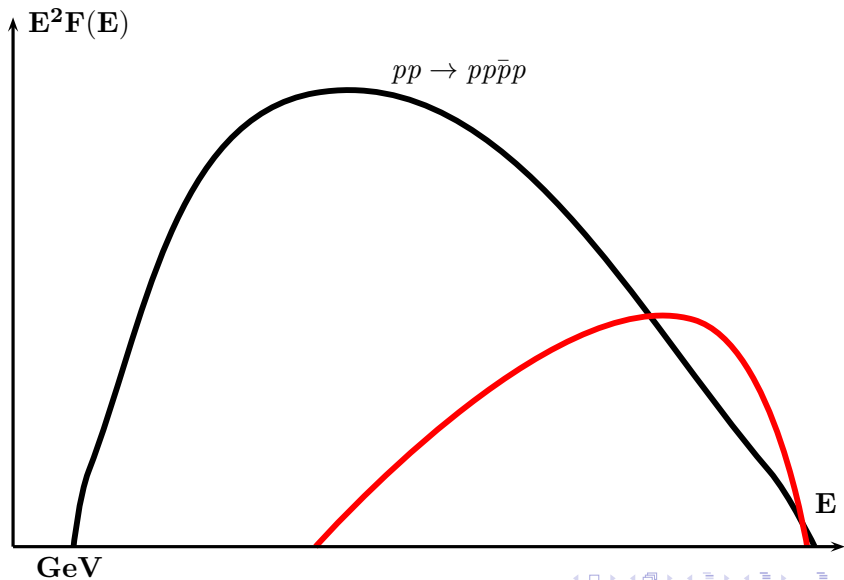
3 Antinuclei fluxes and detection prospects

- ▶ Boosting **anti-helium** fluxes?

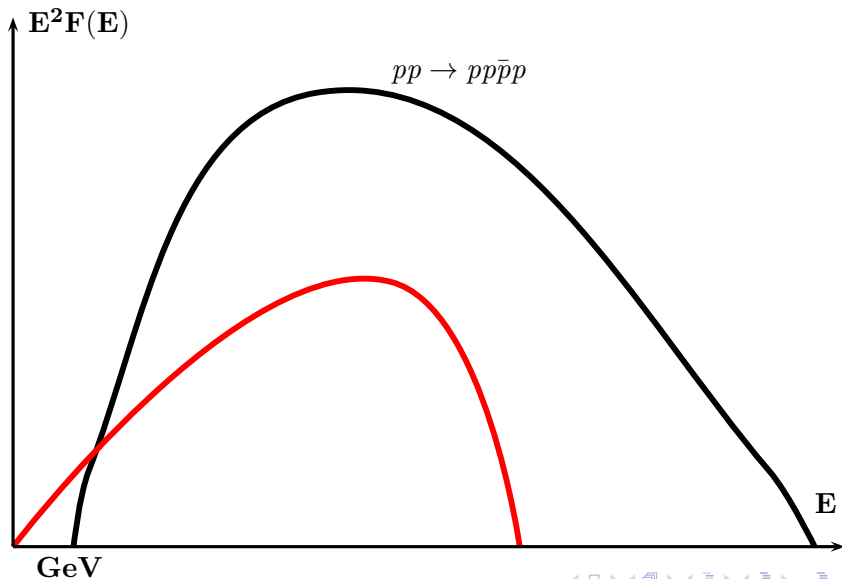
4 Conclusions

Indirect dark matter searches:

- intense search for DM since $\gtrsim 30$ years:
 - ▶ γ -ray line, γ -rays from dwarf galaxies,...
 - ▶ neutrinos from the Sun
 - ▶ excess of antimatter: positrons, antiprotons,...
- no clear evidence found yet...
- signal hidden below astrophysical backgrounds?

Antiprotons from DM: close to M_X 

Antiprotons from DM: below threshold for CR secondaries



Antideuterons as signature for DM

- **antiprotons**: two interesting energy regions
 - ▶ high energy: increasing \bar{p}/p ratio
 - ▶ low energy: **high threshold** for $pp \rightarrow \bar{p}ppp$ reduces CR background

Antideuterons as signature for DM

- antiprotons: two interesting energy regions
 - ▶ high energy: increasing \bar{p}/p ratio
 - ▶ low energy: high threshold for $pp \rightarrow \bar{p}ppp$ reduces CR background
- Donati, Fornengo, Salati '99: **threshold** for $pp \rightarrow pppn\bar{p}\bar{n}$ even higher

Antideuterons as signature for DM

- antiprotons: two interesting energy regions
 - ▶ high energy: increasing \bar{p}/p ratio
 - ▶ low energy: high threshold for $pp \rightarrow \bar{p}ppp$ reduces CR background
 - Donati, Fornengo, Salati '99: threshold for $pp \rightarrow pppn\bar{p}\bar{n}$ even higher
- ⇒ antideuterons **promising signature for DM**
- up-to **10 events/yr** in AMS-02

Antideuterons as signature for DM

- antiprotons: two interesting energy regions
 - ▶ high energy: increasing \bar{p}/p ratio
 - ▶ low energy: high threshold for $pp \rightarrow \bar{p}ppp$ reduces CR background
- Donati, Fornengo, Salati '99: threshold for $pp \rightarrow pppn\bar{p}\bar{n}$ even higher

⇒ antideuterons promising signature for DM

- up-to 10 events/yr in AMS-02
- **today**'s expectation?

Why to look into antimatter now?



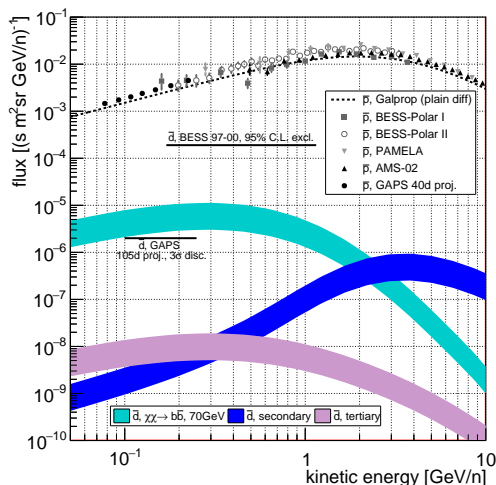
Why to look into antinuclei now?

- AMS-02: six ${}^3\overline{\text{He}}$ and two ${}^4\overline{\text{He}}$ candidates (2019)



Why to look into antinuclei now?

- AMS-02: six ${}^3\overline{\text{He}}$ and two ${}^4\overline{\text{He}}$ candidates (2019)
- GAPS: flights scheduled next antarctic summer



Formation of light nuclei

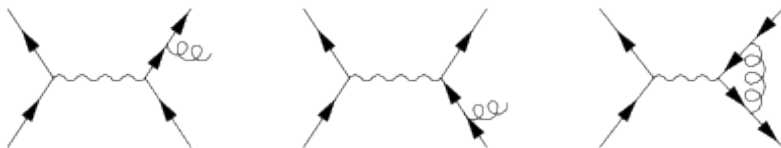
interested in **various types of reactions:**

- $DM+DM \rightarrow X\bar{d}$
- $e^+e^- \rightarrow X\bar{d}$
- $pp \rightarrow X\bar{d}$
- $Ap \rightarrow X\bar{d}$
- $AA \rightarrow X\bar{d}$

different physics, different communities \Rightarrow different approaches

Simplest case: $e^+e^- \rightarrow \text{hadrons}$

- **hard interaction:** LO or NLO matrix element



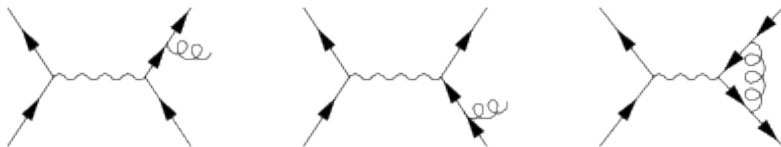
- perturbative parton cascade

▶ ordered in virtualities and angles $s \simeq Q_1^2 > Q_2^2 > \dots > Q_{\min}^2 \gg \Lambda_{\text{QCD}}^2$

- hadronisation volume in cms: $\sigma_{\parallel} \sim 1/(\gamma m_p)$, $\sigma_{\perp} \sim 1/\Lambda_{\text{QCD}}$

Simplest case: $e^+e^- \rightarrow \text{hadrons}$

- hard interaction: LO or NLO matrix element



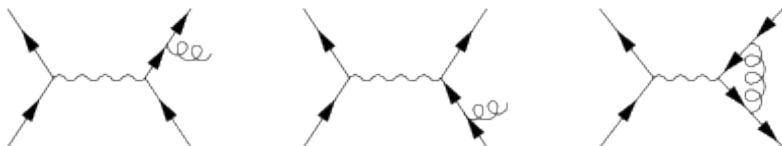
- perturbative parton cascade

▶ ordered in virtualities and angles $s \simeq Q_1^2 > Q_2^2 > \dots > Q_{\min}^2 \gg \Lambda_{\text{QCD}}^2$

- hadronisation volume in cms: $\sigma_{\parallel} \sim 1/(\gamma m_p)$, $\sigma_{\perp} \sim 1/\Lambda_{\text{QCD}}$

Simplest case: $e^+e^- \rightarrow \text{hadrons}$

- hard interaction: LO or NLO matrix element



- perturbative parton cascade

▶ ordered in virtualities and angles $s \simeq Q_1^2 > Q_2^2 > \dots > Q_{\min}^2 \gg \Lambda_{\text{QCD}}^2$

- **hadronisation** volume in cms: $\sigma_{\parallel} \sim 1/(\gamma m_p)$, $\sigma_{\perp} \sim 1/\Lambda_{\text{QCD}}$

General picture:

- separation of scales:

- ▶ $B_d \simeq 2 \text{ MeV} \ll \Lambda_{\text{QCD}}, T_{\text{QCD}}$

⇒ coalescence happens after hadronisation

- semiclassical picture: $\bar{p}(\mathbf{x}, \mathbf{p})$ and $\bar{n}(\mathbf{x}', \mathbf{p}')$ form an antideuteron, if “close” in phase-space

- approximations:

- ▶ DM: coalescence in momentum space: $V \ll 4\pi R_d^3/3$

- ⇒ $f_{\bar{p}}(\mathbf{x}, \mathbf{p}) \simeq \delta(\mathbf{x} - \mathbf{x}_0) f_{\bar{p}}(\mathbf{p})$

- ▶ Heavy-ion: coalescence in coordinate space: $V \gg 4\pi R_d^3/3$

- ▶ combination?

General picture:

- separation of scales:
 - ▶ $B_d \simeq 2 \text{ MeV} \ll \Lambda_{\text{QCD}}, T_{\text{QCD}}$

⇒ **coalescence happens after hadronisation**

- semiclassical picture: $\bar{p}(\mathbf{x}, \mathbf{p})$ and $\bar{n}(\mathbf{x}', \mathbf{p}')$ form an antideuteron, if “close” in phase-space
- approximations:
 - ▶ DM: coalescence in momentum space: $V \ll 4\pi R_d^3/3$
 $\Rightarrow f_{\bar{p}}(\mathbf{x}, \mathbf{p}) \simeq \delta(\mathbf{x} - \mathbf{x}_0) f_{\bar{p}}(\mathbf{p})$
 - ▶ Heavy-ion: coalescence in coordinate space: $V \gg 4\pi R_d^3/3$
 - ▶ combination?

General picture:

- separation of scales:

- ▶ $B_d \simeq 2 \text{ MeV} \ll \Lambda_{\text{QCD}}, T_{\text{QCD}}$

⇒ coalescence happens after hadronisation

- **semiclassical picture:** $\bar{p}(\mathbf{x}, \mathbf{p})$ and $\bar{n}(\mathbf{x}', \mathbf{p}')$ form an antideuteron, if “close” in phase-space

- approximations:

- ▶ DM: coalescence in momentum space: $V \ll 4\pi R_d^3/3$

- ⇒ $f_{\bar{p}}(\mathbf{x}, \mathbf{p}) \simeq \delta(\mathbf{x} - \mathbf{x}_0) f_{\bar{p}}(\mathbf{p})$

- ▶ Heavy-ion: coalescence in coordinate space: $V \gg 4\pi R_d^3/3$

- ▶ combination?

General picture:

- separation of scales:

- ▶ $B_d \simeq 2 \text{ MeV} \ll \Lambda_{\text{QCD}}, T_{\text{QCD}}$

⇒ coalescence happens after hadronisation

- semiclassical picture: $\bar{p}(\mathbf{x}, \mathbf{p})$ and $\bar{n}(\mathbf{x}', \mathbf{p}')$ form an antideuteron, if “close” in phase-space

- approximations:

- ▶ DM: **coalescence in momentum space**: $V \ll 4\pi R_d^3/3$

- ⇒ $f_{\bar{p}}(\mathbf{x}, \mathbf{p}) \simeq \delta(\mathbf{x} - \mathbf{x}_0) f_{\bar{p}}(\mathbf{p})$

- ▶ Heavy-ion: coalescence in coordinate space: $V \gg 4\pi R_d^3/3$

- ▶ combination?

General picture:

- separation of scales:

- ▶ $B_d \simeq 2 \text{ MeV} \ll \Lambda_{\text{QCD}}, T_{\text{QCD}}$

⇒ coalescence happens after hadronisation

- semiclassical picture: $\bar{p}(\mathbf{x}, \mathbf{p})$ and $\bar{n}(\mathbf{x}', \mathbf{p}')$ form an antideuteron, if “close” in phase-space

- approximations:

- ▶ DM: coalescence in momentum space: $V \ll 4\pi R_d^3/3$

- ⇒ $f_{\bar{p}}(\mathbf{x}, \mathbf{p}) \simeq \delta(\mathbf{x} - \mathbf{x}_0) f_{\bar{p}}(\mathbf{p})$

- ▶ Heavy-ion: **coalescence in coordinate space**: $V \gg 4\pi R_d^3/3$

- ▶ combination?

General picture:

- separation of scales:

- ▶ $B_d \simeq 2 \text{ MeV} \ll \Lambda_{\text{QCD}}, T_{\text{QCD}}$

⇒ coalescence happens after hadronisation

- semiclassical picture: $\bar{p}(\mathbf{x}, \mathbf{p})$ and $\bar{n}(\mathbf{x}', \mathbf{p}')$ form an antideuteron, if “close” in phase-space

- approximations:

- ▶ DM: coalescence in momentum space: $V \ll 4\pi R_d^3/3$

- ⇒ $f_{\bar{p}}(\mathbf{x}, \mathbf{p}) \simeq \delta(\mathbf{x} - \mathbf{x}_0) f_{\bar{p}}(\mathbf{p})$

- ▶ Heavy-ion: coalescence in coordinate space: $V \gg 4\pi R_d^3/3$

- ▶ combination?

Coalescence model in momentum space

- all **nucleons** with **cms** momentum difference $\Delta p < p_0$ form a **nucleus**

Coalescence model in momentum space

- all nucleons with cms momentum difference $\Delta p < p_0$ form a nucleus
- for **pp** and **pA collisions**: “fireball” **isotropic distribution** d^3N/dp^3

Coalescence model in momentum space

- all nucleons with cms momentum difference $\Delta p < p_0$ form a nucleus
- for pp and pA collisions: “fireball” isotropic distribution d^3N/dp^3
- antideuterons \sim antiprotons²

$$\frac{dN_{\bar{d}}}{dT} = \frac{p_0^3}{6} \frac{m_d}{m_N^2} \frac{1}{\sqrt{T_{\bar{d}}^2 + 2m_d T_{\bar{d}}}} \left(\left. \frac{dN_{\bar{N}}}{dT} \right|_{T_{\bar{d}}=T_{\bar{N}}/2} \right)^2$$

Coalescence model in momentum space

- all nucleons with cms momentum difference $\Delta p < p_0$ form a nucleus
- for pp and pA collisions: “fireball” isotropic distribution d^3N/dp^3
- antideuterons \sim antiprotons²

$$\frac{dN_{\bar{d}}}{dT} = \frac{p_0^3 m_d}{6 m_N^2} \frac{1}{\sqrt{T_d^2 + 2m_d T_d}} \left(\frac{dN_{\bar{N}}}{dT} \Big|_{T_{\bar{d}}=T_{\bar{N}}/2} \right)^2$$

- more general: **antinuclei A** $\sim B_A$ **antiprotons^A** with

$$B_A = A \left(\frac{4\pi p_0^3}{3 m_N} \right)^{A-1}$$

\Rightarrow **strong hierarchy** $\bar{p} \gg \bar{d} \gg \overline{{}^3\text{He}} \gg \overline{{}^4\text{He}}$

Coalescence model in momentum space

- all nucleons with cms momentum difference $\Delta p < p_0$ form a nucleus
- for pp and pA collisions: “fireball” isotropic distribution d^3N/dp^3
- antideuterons \sim antiprotons²

$$\frac{dN_{\bar{d}}}{dT} = \frac{p_0^3}{6} \frac{m_d}{m_N^2} \frac{1}{\sqrt{T_d^2 + 2m_d T_d}} \left(\left. \frac{dN_{\bar{N}}}{dT} \right|_{T_d = T_{\bar{N}}/2} \right)^2$$

- consider instead **DM annihilation** $XX \rightarrow W^+ W^-$:
- **decay products** of W are **boosted** in cone with $\vartheta \sim m_W/m_X$

Coalescence model in momentum space

- all nucleons with cms momentum difference $\Delta p < p_0$ form a nucleus
- for pp and pA collisions: “fireball” isotropic distribution d^3N/dp^3
- antideuterons \sim antiprotons²

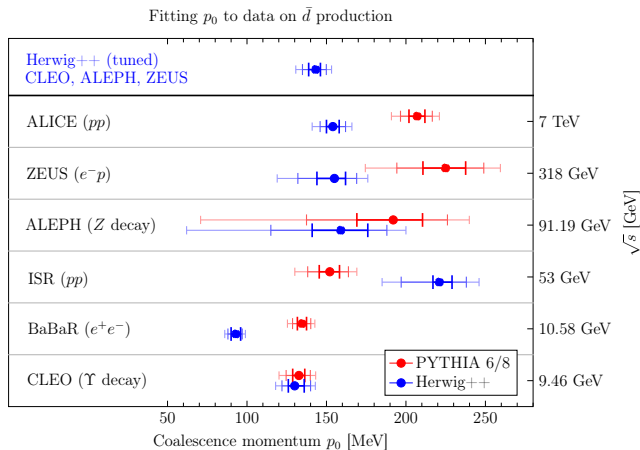
$$\frac{dN_{\bar{d}}}{dT} = \frac{p_0^3}{6} \frac{m_d}{m_N^2} \frac{1}{\sqrt{T_d^2 + 2m_d T_d}} \left(\left. \frac{dN_{\bar{N}}}{dT} \right|_{T_d = T_{\bar{N}}/2} \right)^2$$

- consider instead DM annihilation $XX \rightarrow W^+ W^-$:
- decay products of W are boosted in cone with $\vartheta \sim m_W/m_X$

\Rightarrow requires **momentum correlation on event-by-event basis**

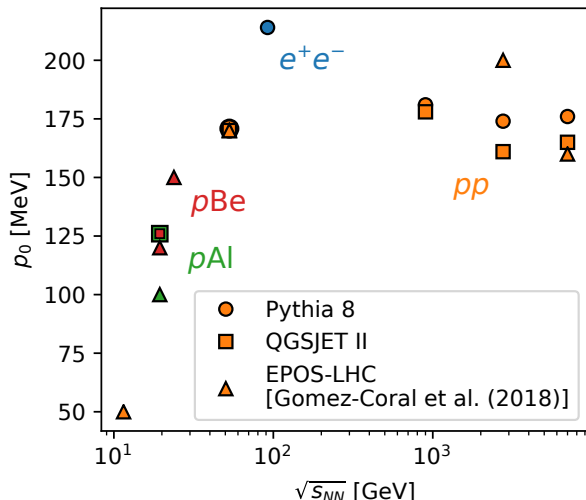
Problems of this approach:

- discrepancies in p_0 between reactions & MC simulations



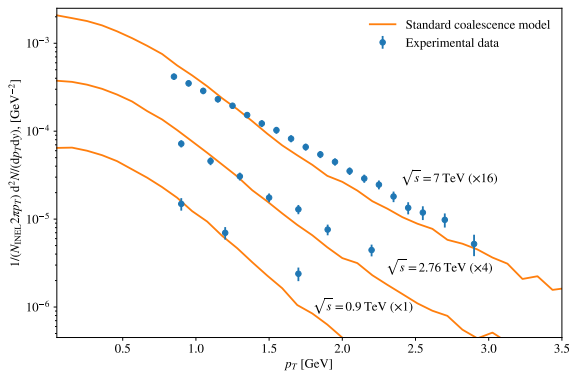
Problems of this approach:

- discrepancies in p_0 between reactions & MC simulations
- energy dependence of p_0 ?



Problems of this approach:

- discrepancies in p_0 between reactions & MC simulations
- energy dependence of p_0 ?
- bad fit of ALICE p_{\perp} spectra



Solution: use Wigner functions with momentum correlation

- two-body **Wigner function** $W(x, p)$ contains **full quantum mechanical information** of a system
- probability distributions follow as

$$\int dx W(x, p) = \phi^*(p) \phi(p), \quad \int \frac{dp}{2\pi} W(x, p) = \psi^*(x) \psi(x)$$

- use momentum distribution $G_{np}(\mathbf{p}_n, \mathbf{p}_p)$ from Monte Carlo
- add Gaussian guess for spatial distribution

Solution: use Wigner functions with momentum correlation

- two-body Wigner function $W(x, p)$ contains full quantum mechanical information of a system
- probability distributions follow as

$$\int dx W(x, p) = \phi^*(p) \phi(p), \quad \int \frac{dp}{2\pi} W(x, p) = \psi^*(x) \psi(x)$$

- use momentum distribution $G_{np}(\mathbf{p}_n, \mathbf{p}_p)$ from Monte Carlo
- add Gaussian guess for spatial distribution

Solution: use Wigner functions with momentum correlation

- two-body Wigner function $W(x, p)$ contains full quantum mechanical information of a system
- probability distributions follow as

$$\int dx W(x, p) = \phi^*(p) \phi(p), \quad \int \frac{dp}{2\pi} W(x, p) = \psi^*(x) \psi(x)$$

- use **momentum distribution** $G_{np}(\mathbf{p}_n, \mathbf{p}_p)$ from Monte Carlo
- add **Gaussian** guess for **spatial distribution**

Solution: use Wigner functions with momentum correlation

- two-body Wigner function $W(x, p)$ contains full quantum mechanical information of a system
- probability distributions follow as

$$\int dx W(x, p) = \phi^*(p) \phi(p), \quad \int \frac{dp}{2\pi} W(x, p) = \psi^*(x) \psi(x)$$

- use momentum distribution $G_{np}(\mathbf{p}_n, \mathbf{p}_p)$ from Monte Carlo
- add Gaussian guess for spatial distribution
- use **connection to density matrix**

$$\langle \psi(\mathbf{x})^\dagger \psi(\mathbf{x}') \rangle = \int \frac{dp}{2\pi} W\left(\mathbf{p}, \frac{\mathbf{x} + \mathbf{x}'}{2}\right) \exp[i\mathbf{p} \cdot (\mathbf{x} - \mathbf{x}')]$$

Evaluation using Monte Carlo correlations

- standard QM using **density matrices**

$$\frac{d^3 N_d}{dP_d^3} = \text{tr}\{\rho_d \rho_{\text{nucl}}\}$$

Evaluation using Monte Carlo correlations

- standard QM using density matrices

$$\frac{d^3 N_d}{dP_d^3} = \text{tr}\{\rho_d \rho_{\text{nucl}}\}$$

with

- ▶ deuteron density matrix $\rho_d = |\phi_d\rangle \langle \phi_d|$
- ▶ two-nucleon density matrix $\rho_{\text{nucl}} = |\psi_p \psi_n\rangle \langle \psi_n \psi_p|$

Evaluation using Monte Carlo correlations

- standard QM using density matrices

$$\frac{d^3 N_d}{dP_d^3} = \text{tr}\{\rho_d \rho_{\text{nucl}}\}$$

- few simple steps later:

$$\frac{d^3 N_d}{dP_d^3} = \frac{3\zeta}{(2\pi)^6} \int d^3 q e^{-q^2 d^2} G_{np}(+\mathbf{q}, -\mathbf{q}),$$

with

$$\zeta \equiv \left(\frac{d^2}{d^2 + 4\sigma^2} \right)^{3/2} \leq 1$$

Evaluation using Monte Carlo correlations

- standard QM using density matrices

$$\frac{d^3 N_d}{dP_d^3} = \text{tr}\{\rho_d \rho_{\text{nucl}}\}$$

- few simple steps later:

$$\frac{d^3 N_d}{dP_d^3} = \frac{3\zeta}{(2\pi)^6} \int d^3 q e^{-q^2 d^2} G_{np}(+\mathbf{q}, -\mathbf{q}),$$

with

$$\zeta \equiv \left(\frac{d^2}{d^2 + 4\sigma^2} \right)^{3/2} \leq 1$$

- “usual MC momentum approach” would be recovered for
 - ▶ $\sigma \ll d \Rightarrow \zeta \rightarrow 1$
 - ▶ $e^{-q^2 d^2} \rightarrow \vartheta(q - q_{\text{max}})$

Evaluation using Monte Carlo correlations

- standard QM using density matrices

$$\frac{d^3 N_d}{dP_d^3} = \text{tr}\{\rho_d \rho_{\text{nucl}}\}$$

- few simple steps later:

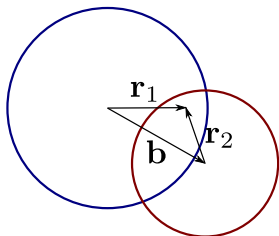
$$\frac{d^3 N_d}{dP_d^3} = \frac{3\zeta}{(2\pi)^6} \int d^3 q e^{-q^2 d^2} G_{np}(+\mathbf{q}, -\mathbf{q}),$$

with

$$\zeta \equiv \left(\frac{d^2}{d^2 + 4\sigma^2} \right)^{3/2} \leq 1$$

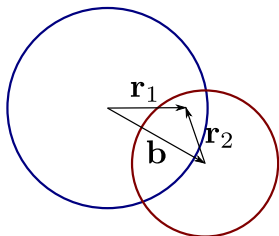
- “usual MC momentum approach” would be recovered for
 - ▶ $\sigma \ll d \Rightarrow \zeta \rightarrow 1$
 - ▶ $e^{-q^2 d^2} \rightarrow \vartheta(q - q_{\text{max}})$
- **fraction $\bar{d}/(\bar{p} + \bar{n})$ is bounded**

Generalising to Ap and AA collisions



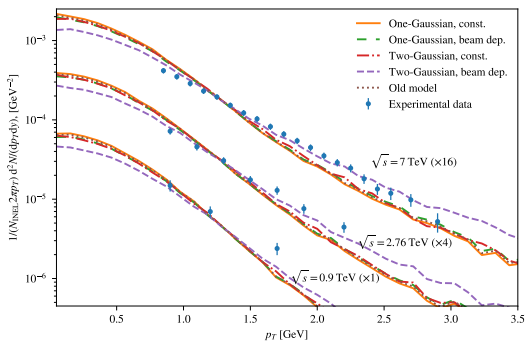
- ▶ **parton cloud** distributed within R_p or R_A
- ▶ **multiple parton interactions**
- ▶ **cluster** can form from **different parton interactions**

Generalising to Ap and AA collisions



- ▶ parton cloud distributed within R_p or R_A
 - ▶ multiple parton interactions
 - ▶ cluster can form from different parton interactions
- using Gaussian profiles:
 - ▶ pp : $\sigma^{pp} = \sqrt{2}\sigma^{e^+e^-}$

Comparison with ALICE and LEP data



Best fit values for spatial extension σ : (using PYTHIA)

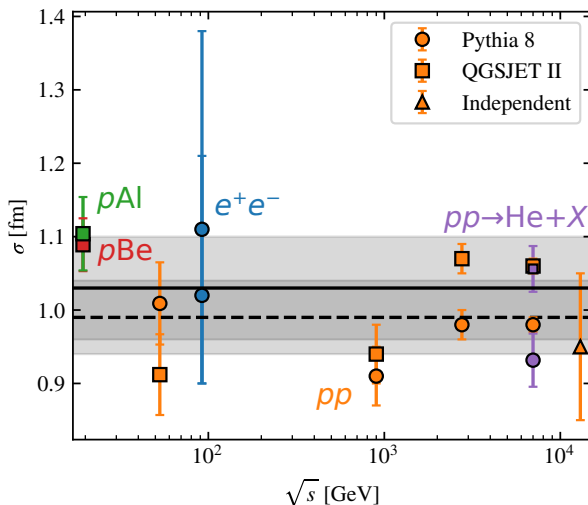
- $\sigma^{pp} = (7.6 \pm 0.1)/\text{GeV}$
- $\sigma^{e^+e^-} = (5.3^{+1.0}_{-0.6})/\text{GeV}$

Comparison with experimental data on pp and Ap:

- assume $R_A \simeq a_0 A^{1/3}$ with $\sigma^{pp} \simeq a_0$ as fit parameter

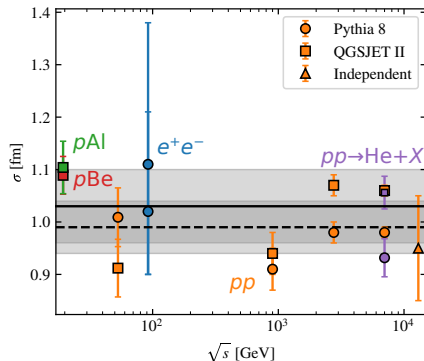
Comparison with experimental data on pp and A_p :

- assume $R_A \simeq a_0 A^{1/3}$ with $\sigma^{pp} \simeq a_0$ as fit parameter

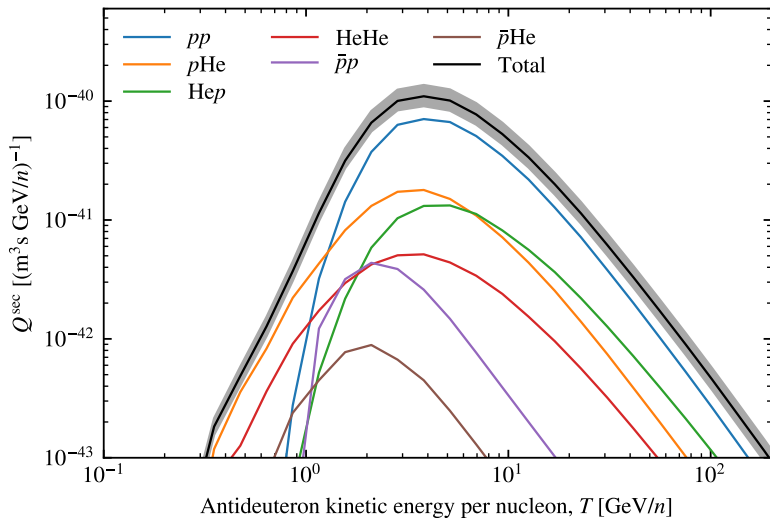


Comparison with experimental data on pp and Ap:

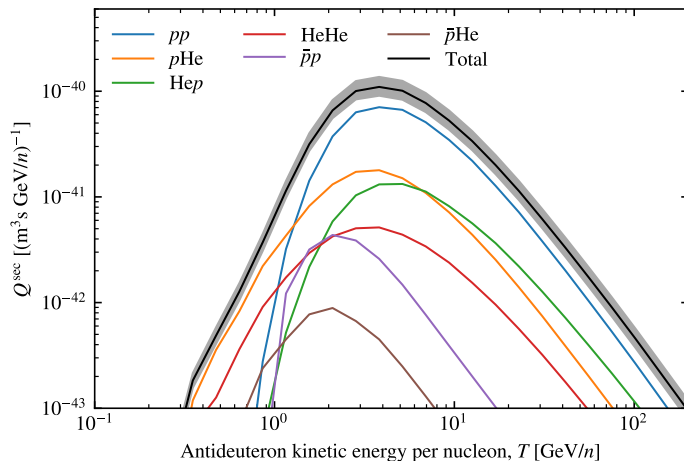
- assume $R_A \simeq a_0 A^{1/3}$ with $\sigma^{pp} \simeq a_0$ as fit parameter



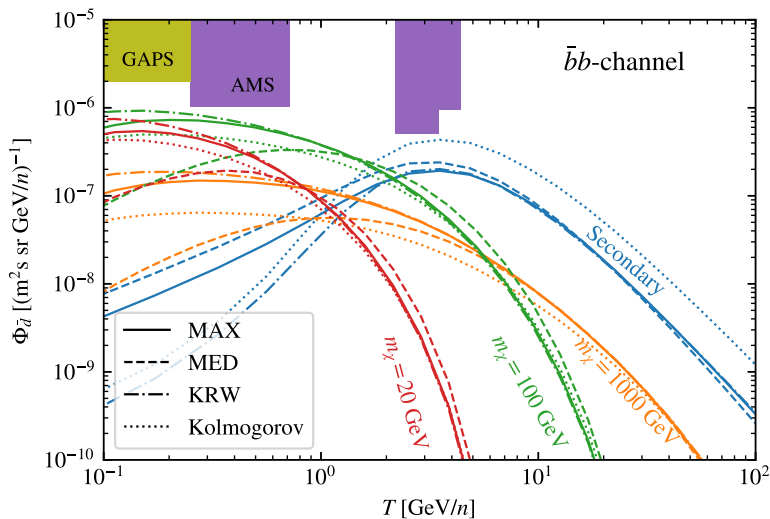
- good agreement with expectation $\sigma^{pp} \sim 1$ fm
- independent of energy and reaction type

Source term Q^{sec} for secondary production of \bar{d} :

Source term Q^{sec} for secondary production of \bar{d} :



- lower threshold in pA reactions \Rightarrow dominate at low T

Antideuteron flux: secondaries plus $\bar{b}b$ channel

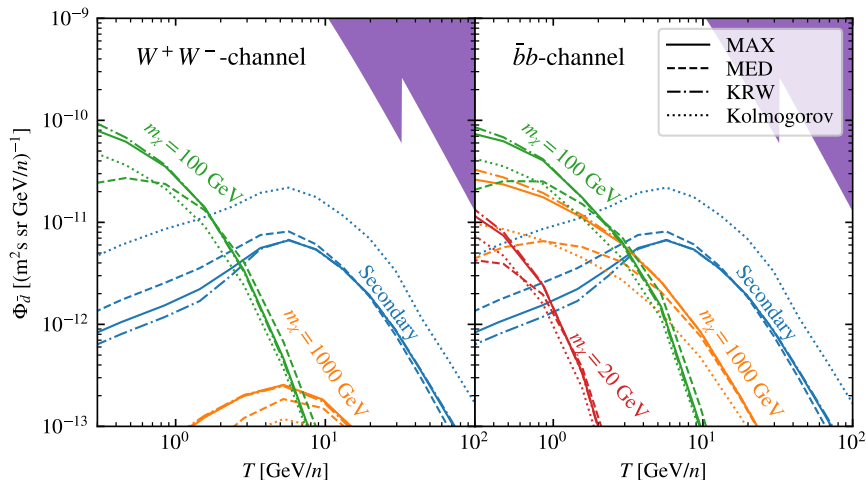
Antideuteron flux: secondaries plus

- DM signal at low-energies background-free

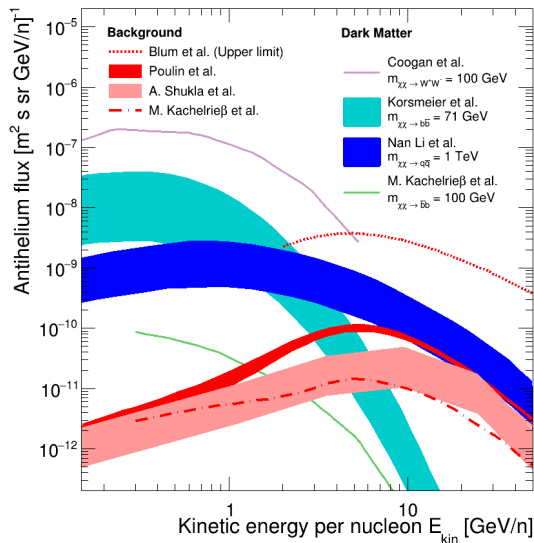
Antideuteron flux: secondaries plus

- DM signal at low-energies background-free
- \bar{d} flux limited by \bar{p}

Antihelium-3 flux:



Antihelium-3: Comparison to other results



Boosting (and shifting) the He flux?

- change cosmology: **inhomogenous barygenesis**
 - ⇒ **anti-stars** in Milky Way
 - ▶ **acceleration mechanism: anti-SNe, anti-SNR?**

[Dolgov, Silk '93, Poulin et al. '19]

Boosting (and shifting) the He flux?

- change cosmology: inhomogenous barygenesis
 - ⇒ anti-stars in Milky Way
 - ▶ acceleration mechanism: anti-SNe, anti-SNR?
- change **particle physics**:

[Dolgov, Silk '93, Poulin et al. '19]

$$n_d(\mathbf{p}) \propto n_n^2(\mathbf{p})$$

$$n_{3\text{He}}(\mathbf{p}) \propto n_n^3(\mathbf{p})$$

$$n_{4\text{He}}(\mathbf{p}) \propto n_n^4(\mathbf{p})$$

Boosting (and shifting) the He flux?

- change cosmology: inhomogenous baryogenesis
 - ⇒ anti-stars in Milky Way
 - ▶ acceleration mechanism: anti-SNe, anti-SNR?
- change particle physics:

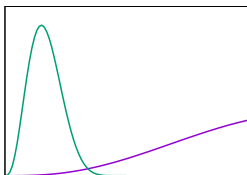
[Dolgov, Silk '93, Poulin et al. '19]

$$n_d(\mathbf{p}) \propto n_n^2(\mathbf{p})$$

$$n_{3\text{He}}(\mathbf{p}) \propto n_n^3(\mathbf{p})$$

$$n_{4\text{He}}(\mathbf{p}) \propto n_n^4(\mathbf{p})$$

⇒ need to **compress** $n_n(\mathbf{p})$:



Boosting the He flux – particle physics

- $m_{\text{DM}} = (1 + \varepsilon)m_{3\text{He}}$
- involve $\bar{\Lambda}_b$ decays
- strongly coupled DM sector

Can $\bar{\Lambda}_b$ decays boost $\overline{{}^3\text{He}}$ from DM?

[Winkler, Linden '21]

- **Majorana** DM: $\sigma_{\text{ann}} \propto m_f^2 \Rightarrow$ couples mainly to **b quarks** for $m_X < m_Z$

Can $\bar{\Lambda}_b$ decays boost $\overline{{}^3\text{He}}$ from DM?

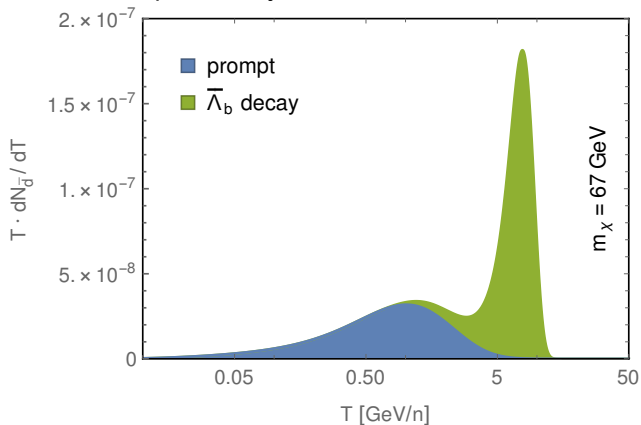
[Winkler, Linden '21]

- Majorana DM: $\sigma_{\text{ann}} \propto m_f^2 \Rightarrow$ couples mainly to b quarks for $m_X < m_Z$
- mass of $\bar{\Lambda}_b$ is close to $5m_N \Rightarrow$ relative momentum small \Rightarrow large coalescence probability for ${}^3\text{He}$

Can $\bar{\Lambda}_b$ decays boost $\overline{{}^3\text{He}}$ from DM?

[Winkler, Linden '21]

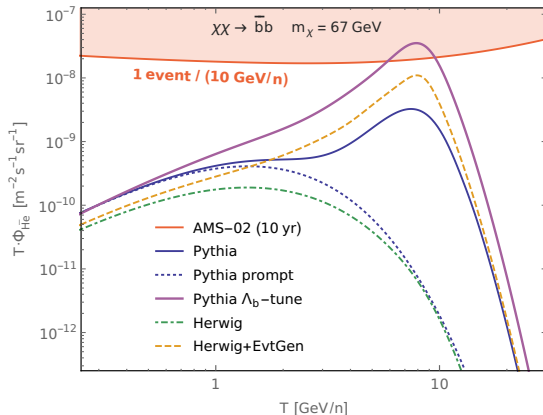
- Majorana DM: $\sigma_{\text{ann}} \propto m_f^2 \Rightarrow$ couples mainly to b quarks for $m_X < m_Z$
- mass of $\bar{\Lambda}_b$ is close to $5m_N \Rightarrow$ relative momentum small \Rightarrow large coalescence probability for $\overline{{}^3\text{He}}$



Can $\bar{\Lambda}_b$ decays boost ${}^3\text{He}$ from DM?

[Winkler, Linden '21]

- Majorana DM: $\sigma_{\text{ann}} \propto m_f^2 \Rightarrow$ couples mainly to b quarks for $m_X < m_Z$
- mass of $\bar{\Lambda}_b$ is close to $5m_N \Rightarrow$ relative momentum small \Rightarrow large coalescence probability for ${}^3\text{He}$



Can $\bar{\Lambda}_b$ decays boost $\overline{{}^3\text{He}}$ from DM?

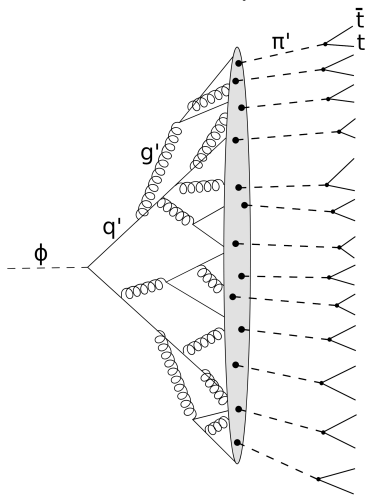
[Winkler, Linden '21]

- Majorana DM: $\sigma_{\text{ann}} \propto m_f^2 \Rightarrow$ couples mainly to b quarks for $m_X < m_Z$
- mass of $\bar{\Lambda}_b$ is close to $5m_N \Rightarrow$ relative momentum small \Rightarrow large coalescence probability for ${}^3\text{He}$
- **no:**
 - ▶ Λ_b tune of Pythia is excluded
 - ▶ Pythia overestimates $\text{BR}(\Lambda_b \rightarrow \bar{u}du(ud_0))$
- + can be tested by LHCb

Strongly coupled DM sector

[Winkler, de la Torre, Linden '22]

- dark QCD cascade produces ~ 1000 top-quarks with $\gamma \simeq \text{few}$



Strongly coupled DM sector

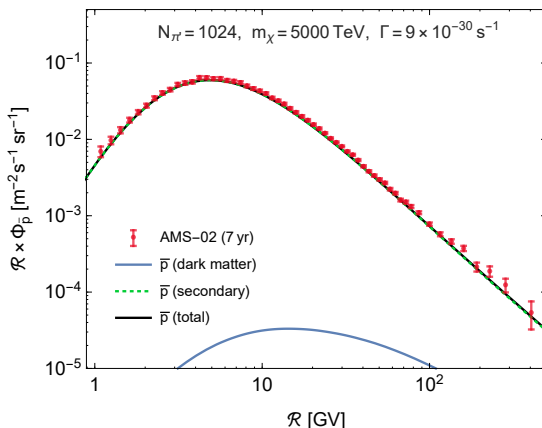
[Winkler, de la Torre, Linden '22]

- dark QCD cascade produces ~ 1000 top-quarks with $\gamma \simeq \text{few}$
- independent QCD cascades, **all nucleons can coalesce**

Strongly coupled DM sector

[Winkler, de la Torre, Linden '22]

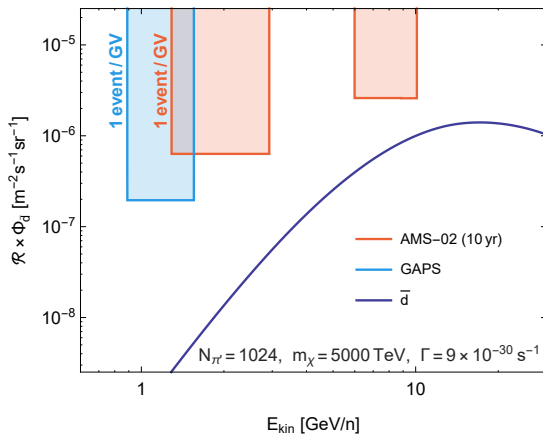
- dark QCD cascade produces ~ 1000 top-quarks with $\gamma \simeq \text{few}$
- independent QCD cascades, all nucleons can coalesce



Strongly coupled DM sector

[Winkler, de la Torre, Linden '22]

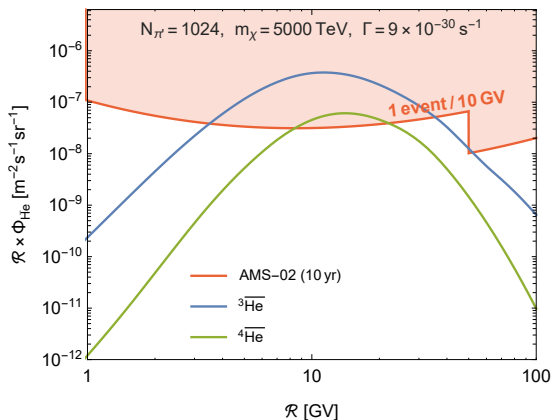
- dark QCD cascade produces ~ 1000 top-quarks with $\gamma \simeq \text{few}$
- independent QCD cascades, all nucleons can coalesce



Strongly coupled DM sector

[Winkler, de la Torre, Linden '22]

- dark QCD cascade produces ~ 1000 top-quarks with $\gamma \simeq \text{few}$
- independent QCD cascades, all nucleons can coalesce



Strongly coupled DM sector

[Winkler, de la Torre, Linden '22]

- dark QCD cascade produces ~ 1000 top-quarks with $\gamma \simeq \text{few}$
- independent QCD cascades, all nucleons can coalesce
- problem: **1000 top within $1/\Lambda_{\text{QCD}'} \simeq 1/\text{TeV} \Rightarrow \text{QGP}$**

Conclusions

- 1 Formation of light antinuclei is interesting in itself:
 - ▶ inclusion of two-particle momentum correlations necessary
 - ▶ reaction-dependent size of source is important
 - ▶ how to deal with spatial correlations?
 - ▶ when are collective effects important?
- 2 Coalescence in phasespace – WiFunC model:
 - ▶ consistent description of various reactions
- 3 Antinuclei are a useful tool searching for new physics
 - ▶ antideuterons as signal for DM
 - ▶ strong hierarchy of fluxes as function of A
 - ▶ antihelium-3 and especially antihelium-4 challenging
- 4 Upgrade of AMS-02 in 2023, extension of ISS