# LIGO-Virgo-KAGRA O3 cosmology results <u>Konstantin Leyde</u>

11.10.2022 GdR France Ondes Gravitationelles





2111.03604

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### From data to the $H_0$ posterior



# Cosmology with gravitational waves

#### IcaroGW

- Deduce redshift from joint fit of the source frame mass and cosmological parameters
- Marginalize over mass population

#### •GWcosmo

- Assumption: GW sources in galaxies
- Statistical redshift association from galaxy catalogs
- Fixed mass distribution







### Gravitational wave parameters



- GW frequency is shifted to lower values by the expansion
- Redshift information is degenerate with other variables

$$m^{(d)} = (1+z)m^{(s)}$$



# Source frame population

Assumption of mass model  $\rightarrow$  statistical measurement of redshift

$$m^{(d)} = (1+z)m^{(s)} \rightarrow z = \frac{m^{(d)}}{m^{(s)}}$$

- Joint fit of cosmological parameters and mass population models (Taylor et al. 2012, Taylor and Gair 2012, Farr et al. 2019, You et al. 2020
- Strong correlation between  $H_0$  and the characteristic mass scales

$$m^{(s)} = \frac{m^{(d)}}{1 + d_L H_0/c} \qquad z \approx \frac{d_L H_0}{c}$$



## The source mass population model



#### $p_{\rm pop}(\theta \mid \Lambda)$





### GWCOSMO (Gray et al. 1908.06050)

- Assumption: GW sources are found in galaxies
  - Possible GW hosts from galaxy catalogs (Schutz 1986)
  - Give importance to galaxy according to luminosity
- Input from GW side:
  - Sky position
  - Luminosity distance
  - Detector frame masses
- Input from galaxy catalog:
  - Sky position
  - Luminosity
  - Redshift
- Challenge: Galaxy catalogs incompleteness
  - Calculate selection effects:
    - Host galaxy is observed (in catalog) or not

utz 1986) osity



10	100	1000
Galaxy nu	mber (per pixel)	

Gray et al. 2111.04629





# Galaxy catalog used for GWcosmo

- All-sky catalog: Glade (Dálya et al. 2018)
- Partial coverage, but deeper in redshift: DES (Drlica-Wagner et al. 2018)





### Results IcaroGW





## Results GWcosmo

- Fix the mass distribution to the values obtained by Icarogw
- Result is strongly dependent on population assumption





Glade+  $H_0 = 67^{+13}_{-12} \,\mathrm{km}\,\mathrm{s}^{-1}\,\mathrm{Mpc}^{-1}$ K-band  $H_0 = 67^{+8}_{-6} \,\mathrm{km}\,\mathrm{s}^{-1}\,\mathrm{Mpc}^{-1}$ 

Glade+, K-band With GW170817



## Conclusions

- IcaroGW (no EM information) method allows to simultaneously constrain cosmological and population parameters
- GWcosmo uses galaxy catalog information to constrain  $H_0 \rightarrow$  Better constraint on  $H_0$
- Main result from O3

$$H_0 = 67^{+8}_{-6} \,\mathrm{km}\,\mathrm{s}^{-1}\,\mathrm{Mpc}^{-1}$$
 wi

- Strong degeneracies between the rate evolution  $\gamma$ , the overall rate of events  $R_0$ , the Hubble constant  $H_0$ 
  - $\rightarrow$  Marginalize over population assumptions

ith Glade+, K-band and GW170817



### Relaxing an assumption:

#### Modified propagation equation for gravitational waves

$$h''_{A} + 2\mathcal{H}(1 - \delta(\eta))h'_{A} + k^{2}h_{A} = 0$$

- $\delta$  the friction term
- DHOST, 1510.06930, 1703.03797, 1707.03625)
- Results in a modified gravitational wave distance (gravitational wave and electromagnetic distance do not coincide)
- **Testable** with gravitational wave observations

. k the wave vector, A the GW polarisation,  $\eta$  the conformal time,  $\mathcal{H}=\overset{a'}{-}$  and

• Appears in some modified gravity theories (e.g. beyond Horndeski 1404.6495,

12



## **Binary black hole dark siren** population analysis with modified gravity Konstantin Leyde, Simone Mastrogiovanni, Danièle Steer, Eric Chassande-Mottin, Christos Karathanasis





11.10.2022 Session Cosmologie du GdR Ondes Gravitationnelles





### Assumption on the modifications of GR

$$h_A'' + 2\mathcal{H}(1 - \delta(\eta))h_A' + k^2h_A = 0$$

Phenomenological model 1906.01593

$$\Xi_{0} \text{ characterises}$$
early time behaviour
$$d_{L}^{\text{GW}} = d_{L}^{\text{EM}} \left( \Xi_{0} + \frac{1 - \Xi_{0}}{(1 + z)^{n}} \right)$$

$$n \text{ characterises}$$

Assumptions: No modifications of the waveform during the  $\bullet$ inspiral phase and cosmological background is unchanged

GR:  $\Xi_0 = 1$ 

haracterises the ition from early to late times





### Results with O3 data

- GR:  $\Xi_0 = 1$
- For all modified gravity models: compatible with their GR values at 90% confidence level (for Multi Peak)



Multi peak mass model, varying SNR cut

	Broken Power Law	Multi Peak	$Power \ Law + Peak$	Truncated
D	$6^{+2}_{-2}$	$5^{+3}_{-1}$	$5^{+3}_{-1}$	$4.5\substack{+3.1 \\ -0.8}$
$\Xi_0$	$1.6^{+1.3}_{-0.8}$	$1.4^{+1.1}_{-0.7}$	$1.3^{+1.2}_{-0.7}$	$0.6^{+1.4}_{-0.2}$
$c_M$	$1.0\substack{+2.3 \\ -2.6}$	$0.5\substack{+2.5 \\ -2.4}$	$0.1\substack{+2.7 \\ -2.1}$	$-2^{+3}_{-1}$

#### 60 BBH events, SNR > 10, IFAR > 4 yr

42 BBH events, SNR > 11, IFAR > 4 yr

	Broken Power Law	Multi Peak	$Power \ Law + Peak$	Truncated
D	$4.7\substack{+2.9 \\ -0.9}$	$4.6\substack{+2.6 \\ -0.8}$	$4.7\substack{+2.7 \\ -0.9}$	$5^{+3}_{-1}$
$\Xi_0$	$2^{+3}_{-1}$	$2^{+4}_{-1}$	$2^{+3}_{-1}$	$0.7\substack{+3.0 \\ -0.4}$
$c_M$	$0.5\substack{+4.1 \\ -4.2}$	$1^{+4}_{-5}$	$1^{+4}_{-4}$	$-3^{+5}_{-2}$

35 BBH events, SNR > 12, IFAR > 4 yr

	Broken Power Law	Multi Peak	$Power \ Law + Peak$	Truncated
D	$5^{+3}_{-1}$	$4.6\substack{+2.9 \\ -0.9}$	$4.8^{+2.9}_{-1.0}$	$5^{+3}_{-1}$
$\Xi_0$	$1.2^{+1.4}_{-0.7}$	$1.4^{+1.8}_{-0.8}$	$1.4^{+1.8}_{-0.8}$	$0.8\substack{+2.0\-0.5}$
$c_M$	$-0.1^{+2.8}_{-3.0}$	$0.3^{+3.2}_{-3.3}$	$0.4^{+3.2}_{-3.0}$	$-2^{+5}_{-3}$

#### See also: Mancarella et al. 2112.05728

### Degeneracies



- GR:  $\Xi_0 = 1$
- Gravity deviation parameter  $\Xi_0$ strongly degenerate with the redshift distribution parameter  $\gamma$



# Conclusions

- parameters
  - Implication of O3 : bright sirens are rare
- O3 data favours GR over all modified gravity models investigated
- Study impact of mass models on the measurement of  $\Xi_0$  and on  $H_0$
- Strong degeneracies between the rate evolution  $\gamma$ , the overall rate of events  $R_0$ , the Hubble constant  $H_0$  and friction amplitude  $\Xi_0$ 
  - Assumptions on astrophysics can bias this measurement
  - $\rightarrow$  Marginalize over population assumptions
- Constrain  $\Xi_0$  to 50 % with O4 and 20 % with O4+O5

Method allows to simultaneously constrain modified gravity, cosmological and population



# Thank you!

### **Questions?**

# Forecast (with O4 + O5)

### The effect of a prior on the cosmological values

 $p(\Xi_0)$ 

#### Wide: Agnostic priors for the cosmological parameters

#### Planck: Priors from the Planck estimate for the cosmological parameters

	Agnostic	From Planck
$H_0$	$\mathcal{U}(30,130)$	$\mathcal{U}(66.07, 68.47)$
$\Omega_M$	$\mathcal{U}(0.05, 0.4)$	$\mathcal{U}(0.3082, 0.3250)$



# Back up slides

#### Statistical framework of IcaroGW (Mastrogiovanni et al. 2103.14663)

et al. 2007.05579)

$$p(\Lambda|\{x\}) \propto p(\Lambda) \prod_{j=1}^{N_{obs}} \frac{\int p(x_j|\theta_j) p_{pop}(\theta_j|\Lambda) d\theta_j}{\int p_{det}(\theta_j) p_{pop}(\theta_j|\Lambda) d\theta_j}$$

- Metaparameters  $\Lambda$ : population parameters, cosmological parameters, ...
- GW data  $\{x\}$
- Source parameters  $\theta = \{m_{1,2}^{(d)}, d_L^{GW}, \dots\}$
- GW likelihood  $p(x_i | \Lambda, \theta)$ , obtained from posterior samples
- Population assumption  $p_{pop}(\theta | \Lambda)$
- Detection probability  $p_{det}(\theta)$

Bayesian analysis with selection effects (Mandel et al. 1809.02063, Thrane and Talbot 1809.02293, Vitale

![](_page_21_Picture_13.jpeg)

![](_page_21_Picture_14.jpeg)

### Statistical framework of IcaroGW (Mastrogiovanni et al. 2103.14663) Bayesian analysis with selection effects $\prod_{j=1}^{N_{obs}} \frac{\int p(x_j | \theta_j) p_{pop}(\theta_j | \Lambda) d\theta_j}{\int p_{det}(\theta_j) p_{pop}(\theta_j | \Lambda) d\theta_j}$

$$p(\Lambda|\{x\}) \propto p(\Lambda) \prod_{i=1}^{N_{\text{obs}}}$$

- considered

Only events passing threshold (on signal to noise ratio or false alarm rate) are

• Numerical evaluation of  $p_{det}(\theta)$ : produce a set of events and label them either "detected" or "undetected" (passing SNR threshold and IFAR threshold)

![](_page_22_Picture_8.jpeg)

#### Statistical framework of GWcosmo (Gray et al. 1908.06050)

Bayesian analysis with selection effects

![](_page_23_Figure_2.jpeg)

- Assumes a fixed mass and redshift distribution
- $\hat{d}$ : "event is detected"
- $N_{\rm obs}$  number of observed events

	Observed events				
$N_{obs}$ $\prod_{i=1}^{N}$	$\sum_{g \in [G,\bar{G}]}$	<i>р</i> (х	$G_{\rm GW,i} \mid H_0, \Lambda_m, Q$	$g)p(g   H_0, \Lambda_m, \Lambda_z, \hat{d})$	
xy ?			One event posterior		
utio	ition		Hidden here selection effe	are the ects	

![](_page_23_Picture_9.jpeg)

![](_page_24_Figure_0.jpeg)

- Assumptions for selection effects lacksquare
  - Apparent magnitude threshold of the galaxy catalog
  - **Redshift distribution of galaxies**
  - Luminosity distribution of galaxies (e.g. Schechter function) •
- Pixelated approach: Treat selection effects as non-uniform in the sky

(Gray et al. 1908.06050)

Probability of being in catalog

Probability of **not** being in catalog

$$(\Lambda_{m,z}, G)p(G | H_0, \Lambda_{m,z}, \hat{d}) + p(x_{GW,i} | H_0, \Lambda_{m,z}, \bar{G})p(\bar{G} | H_0, \Lambda_{m,z}, \bar{G})p(\bar{G} | H_0, \Lambda_{m,z}, \bar{G})$$

![](_page_24_Figure_13.jpeg)

![](_page_24_Picture_16.jpeg)

## Results with O3 data

Bayes factor:

 $\frac{p(\text{data} | \text{model}_1)}{p(\text{data} | \text{model}_2)}$ 

- GR:  $\Xi_0 = 1$
- Compare Bayes factors
   → Multi Peak + General Relativity is
   preferred
  - Consistent results for all 3 SNR cuts

	$60~\mathrm{BBH}$ events, $\mathrm{SNR} > 10,\mathrm{IFAR} > 4\mathrm{yr}$				
		Broken Power Law	Multi Peak	$Power \ Law + Peak$	Trunca
	$\mathbf{GR}$	-2.4	0.0	-1.2	-6.3
_	D	-2.0	-0.2	-1.7	-6.4
	$\Xi_0$	-3.2	-0.9	-2.1	-6.8
	$c_M$	-3.0	-1.0	-2.1	-6.5
42 BBH events, SNR $>11,\mathrm{IFAR}>4\mathrm{yr}$					
		Broken Power Law	Multi Peak	$Power \ Law + Peak$	Trunca

0.0

-0.0

-0.6

-0.9

-0.8

-0.9

-1.4

-1.7

35 BBH events, SNR > 12, IFAR > 4 yr

	Broken Power Law	Multi Peak	$Power \ Law + Peak$	Truncat
$\mathbf{GR}$	-1.2	0.0	-1.1	-2.6
D	-1.1	-0.4	-1.2	-2.8
$\Xi_0$	-2.1	-1.0	-1.9	-3.3
$c_M$	-1.9	-1.2	-1.9	-3.1

GR

D

 $\Xi_0$ 

 $c_M$ 

-1.5

-1.5

-1.9

-1.9

![](_page_25_Figure_10.jpeg)

![](_page_25_Figure_11.jpeg)

![](_page_25_Figure_12.jpeg)

### Two limits...

#### "GWcosmo"

- Close-by signals, well-localised (very few compatible galaxies)
  - $H_0$  posterior is independent of the mass distribution assumed

![](_page_26_Figure_4.jpeg)

#### <u>"IcaroGW"</u>

![](_page_26_Figure_6.jpeg)

![](_page_26_Picture_8.jpeg)

![](_page_26_Figure_9.jpeg)

### The $\Xi_0$ paramatrization

Model	$\Xi_0-1$	n	Refs.
HS $f(R)$ gravity	$rac{1}{2}f_{R0}$	$rac{3( ilde{n}+1)\Omega_m}{4-3\Omega_m}$	[68]
Designer $f(R)$ gravity	$-0.24\Omega_m^{0.76}B_0$	$3.1\Omega_m^{0.24}$	[69]
Jordan–Brans–Dicke	$rac{1}{2}\delta\phi_0$	$rac{3( ilde{n}+1)\Omega_m}{4-3\Omega_m}$	[70]
Galileon cosmology	$rac{eta \phi_0}{2 M_{ m Pl}}$	$rac{\dot{\phi}_0}{H_0\phi}$	[71]
$lpha_M = lpha_{M0} a^{ ilde n}$	$rac{lpha_{M0}}{2 ilde{n}}$	$ ilde{n}$	[67]
$lpha_M = lpha_{M0} rac{\Omega_\Lambda(a)}{\Omega_\Lambda}$	$-rac{lpha_{M0}}{6\Omega_\Lambda}\ln\Omega_m$	$-rac{3\Omega_{\Lambda}}{\ln\Omega_{m}}$	[67, 72]
$\Omega = 1 + \Omega_+ a^{\tilde{n}}$	$rac{1}{2}\Omega_+$	${ ilde n}$	[6]
Minimal self-acceleration	$\lambda \left( \ln a_{acc} + \frac{C}{2} \chi_{acc} \right)$	$rac{C/H_0-2}{\ln a_{acc}^2-C\chi_{acc}}$	[66]

Table 1.

#### Table 1 of 1906.01593

Mapping of the parametrisation in Eq. (2.31) to a number of frequently studied, representative modified gravity models embedded in the Horndeski action (3.1) with luminal speed of gravitational waves. For simplicity, we have employed the approximations  $\alpha_{M0} \ll 1$  (and  $n \sim 1$ ).

### DHOST Lagrangian

$$S[g_{\mu\nu},\phi] = \int d^{4}x \sqrt{-g} \left[ F_{0}(\phi,X) + F_{1}(\phi,X)\Box\phi + F_{2}(\phi,X)R + \sum_{I=1}^{5} A_{I}(\phi,X)L_{I}^{(2)} \right]$$

$$L_{1}^{(2)} \equiv \phi^{\mu\nu}\phi_{\mu\nu}, \quad L_{2}^{(2)} \equiv (\phi_{\nu}{}^{\nu})^{2}, \quad L_{3}^{(2)} \equiv \phi_{\nu}{}^{\nu}\phi^{\rho}\phi_{\rho\sigma}\phi^{\sigma}, \quad L_{4}^{(2)} \equiv \phi^{\mu}\phi_{\mu\nu}\phi^{\nu\rho}\phi_{\rho}, \quad L_{5}^{(2)} \equiv (\phi^{\rho}\phi_{\rho\sigma}\phi^{\sigma})^{2}, \quad \text{where } \phi_{\mu\nu} = \nabla_{\nu}\nabla_{\mu}\phi, \text{ and } \phi_{\mu} = \nabla_{\mu}\phi.$$

$$F_{1} = 0, \quad F_{2} = \frac{M_{0}^{2}}{2} + c_{4}X^{2}, \quad A_{3} = -8c_{4} - \beta.$$

$$g_{1,015}^{(2)} = (c_{2} \rightarrow 3., c_{3} \rightarrow 5., c_{4} \rightarrow 1., \beta \rightarrow -5.28333) = 0$$

$$(c_{2} \rightarrow 3., c_{3} \rightarrow 5., c_{4} \rightarrow 1., \beta \rightarrow -5.28333) = 0$$

#### An ex

$$S[g_{\mu\nu}, \phi] = \int d^{4}x \sqrt{-g} \left[ F_{0}(\phi, X) + F_{1}(\phi, X) \Box \phi + F_{2}(\phi, X)R + \sum_{I=1}^{5} A_{I}(\phi, X)L_{I}^{(2)} \right]$$

$$L_{1}^{(2)} \equiv \phi^{\mu\nu}\phi_{\mu\nu}, \quad L_{2}^{(2)} \equiv (\phi_{\nu}{}^{\nu})^{2}, \quad L_{3}^{(2)} \equiv \phi_{\nu}{}^{\nu}\phi^{\rho}\phi_{\rho\sigma}\phi^{\sigma}, \\ L_{4}^{(2)} \equiv \phi^{\mu}\phi_{\mu\nu}\phi^{\nu\rho}\phi_{\rho}, \quad L_{5}^{(2)} \equiv (\phi^{\rho}\phi_{\rho\sigma}\phi^{\sigma})^{2},$$
where  $\phi_{\mu\nu} = \nabla_{\nu}\nabla_{\mu}\phi$ , and  $\phi_{\mu} = \nabla_{\mu}\phi$ .  

$$F_{0} = c_{2}X, \quad F_{1} = 0, \quad F_{2} = \frac{M_{0}^{2}}{2} + c_{4}X^{2}, \quad A_{3} = -8c_{4} - \beta.$$

$$\frac{d_{0}^{GW}}{d_{L}^{EM}} = \frac{1.015}{1.010}$$

$$\frac{d_{0}^{GW}}{d_{L}^{EM}} = \frac{1.025}{1.010}$$

$$\frac{d_{0}^{GW}}{d_{L}^{EM}} = \frac{1.025}{1.010}$$

![](_page_28_Figure_5.jpeg)

![](_page_28_Figure_6.jpeg)

#### Spectral Sirens: Cosmology from the Full Mass Distribution of Compact Binaries

Jose María Ezquiaga<sup>1</sup>, \* and Daniel E. Holz<sup>1,2</sup>

<sup>1</sup>Kavli Institute for Cosmological Physics and Enrico Fermi Institute, The University of Chicago, Chicago, IL 60637, USA <sup>2</sup>Department of Physics, Department of Astronomy & Astrophysics, The University of Chicago, Chicago, IL 60637, USA (Dated: August 15, 2022)

We explore the use of the mass spectrum of neutron stars and black holes in gravitational-wave compact binary sources as a cosmological probe. These standard siren sources provide direct measurements of luminosity distance. In addition, features in the mass distribution, such as mass gaps or peaks, will redshift, and thus provide independent constraints on their redshift distribution. We argue that the entire mass spectrum should be utilized to provide cosmological constraints. For example, we find that the mass spectrum of LIGO–Virgo–KAGRA events introduces at least five independent mass "features": the upper and lower edges of the pair instability supernova (PISN) gap, the upper and lower edges of the neutron star-black hole gap, and the minimum neutron star mass. We find that although the PISN gap dominates the cosmological inference with current detectors (2G), as shown in previous work, it is the lower mass gap that will provide the most powerful constraints in the era of Cosmic Explorer and Einstein Telescope (3G). By using the full mass distribution, we demonstrate that degeneracies between mass evolution and cosmological evolution can be broken, unless an astrophysical conspiracy shifts all features of the full mass distribution simultaneously following the (non-trivial) Hubble diagram evolution. We find that this self-calibrating "spectral siren" method has the potential to provide precision constraints of both cosmology and the evolution of the mass distribution, with 2G achieving better than 10% precision on H(z) at

 $z \lesssim 1$  within a year, and 3G reaching  $\lesssim 1\%$  at  $z \gtrsim 2$  within one month.

#### Future prospects

![](_page_29_Figure_8.jpeg)

sub-percent precision is possible with 1 month of observations.

![](_page_29_Picture_10.jpeg)

#### Future prospects: evolving mass distribution

![](_page_30_Figure_1.jpeg)

$$m_{1z}$$

 $m_{
m edge,ev}(z) = m_{
m edge} + \gamma \cdot z$ 

![](_page_30_Figure_5.jpeg)

![](_page_30_Picture_7.jpeg)