

Fixing the dynamics in scalar-Gauss-Bonnet gravity

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Based on the works made in collaboration with Miguel Bezares, Enrico Barausse and Luis Lehner
[arXiv 2206.00014](https://arxiv.org/abs/2206.00014), accepted for publication in PRD

Why gravity **beyond** GR in a nutshell

usually based on two main arguments:

- no description of gravitational effects in **quantum environments**
- still confused by **dark energy** and **dark matter**

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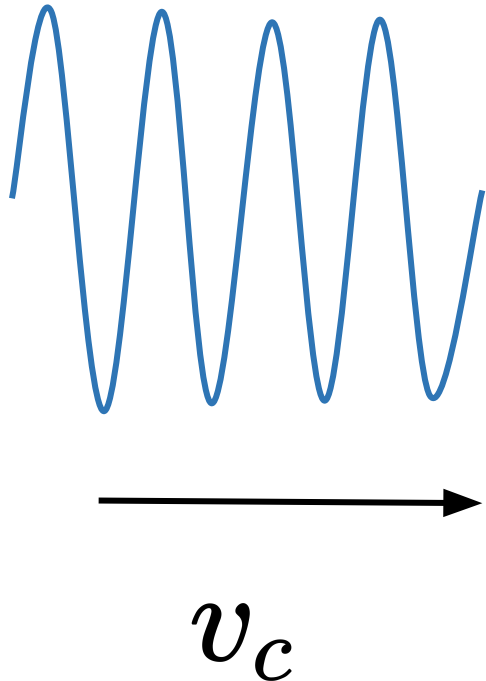
a different perspective: **testing GR with black holes**

Well-posedness beyond GR

- necessary to make prediction with numerical evolutions
- complex mathematical problem (not entering into details)
- focus on characteristic speeds
- derivative interactions can make them diverge or become imaginary dynamically

Well-posedness **beyond** GR

Characteristic speeds



- Take a wave
- Let it propagate in the theory
- **Characteristic speeds** set by feature of the theory: most important is the **highest-derivative operators**
- Depends on the dynamical evolution of the system under consideration

Example: Scalar Gauss-Bonnet gravity

$$\mathcal{G} = R_{abcd}R^{abcd} - 4R_{ab}R^{ab} + R^2$$

The action

$$S_{GB} = \frac{1}{2} \int \sqrt{-g} \left[R - \frac{1}{2} \partial_a \phi \partial^a \phi + f(\phi) \mathcal{G} \right]$$

The equations of motion

$$R_{ab} - \frac{1}{2} g_{ab} R = T_{ab}^{(\phi)} + T_{ab}^{(\mathcal{G})}$$

$$\square \phi = -f'(\phi) \mathcal{G}$$

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Scalar Gauss-Bonnet gravity

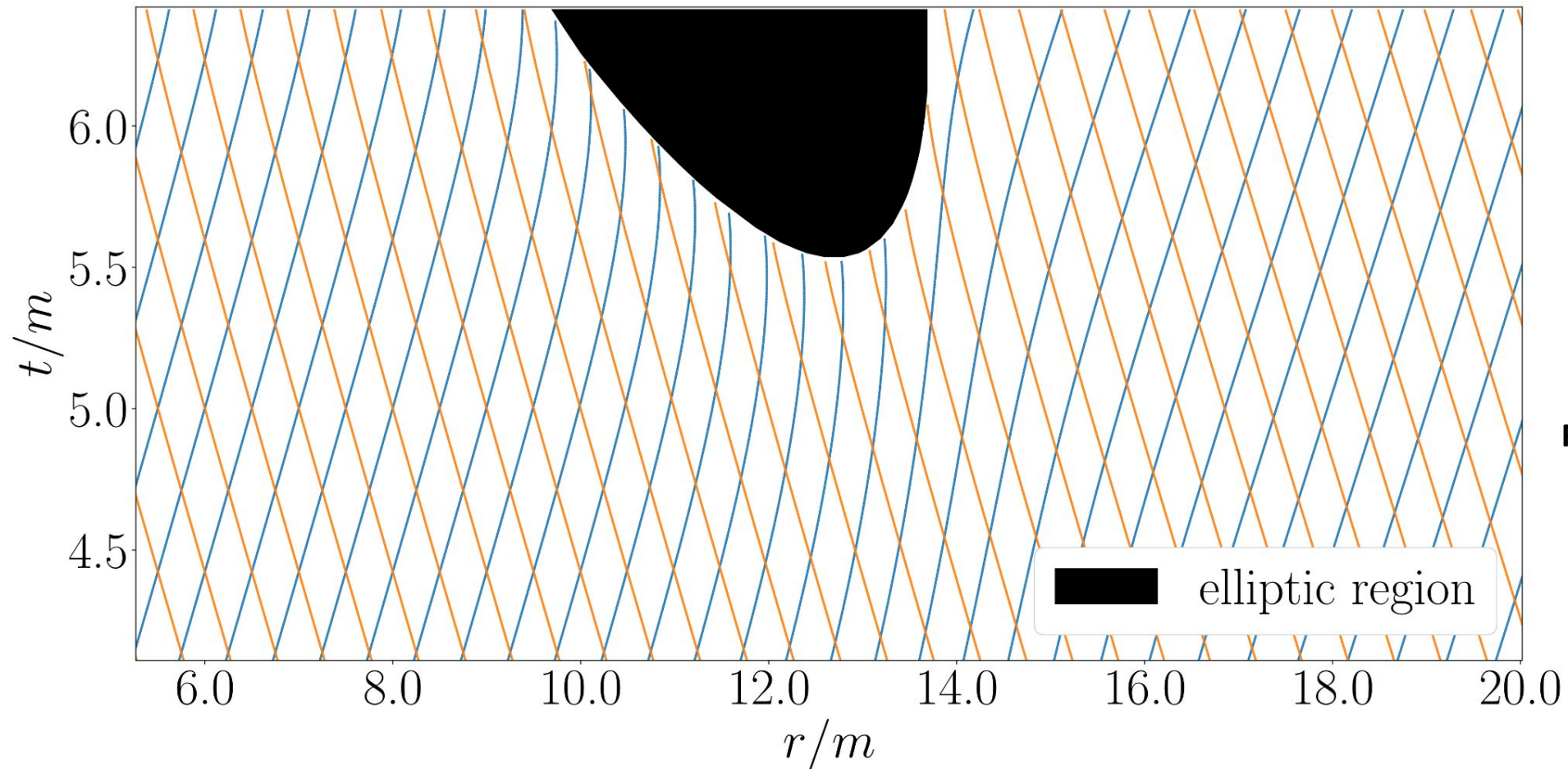
$$\mathcal{G} \propto \partial^2 g$$

$$T_{ab}^{\mathcal{G}} \propto \partial^2 g, \partial^2 \phi$$

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Ripley, Pretorius 2019

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Fixing Scalar Gauss-Bonnet gravity

Proposed in Cayuso, Ortiz, Lehner 2017

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$$R_{ab} - \frac{1}{2}g_{ab}R = T_{ab}^{(\phi)} + \cancel{T_{ab}^{(\mathcal{G})}} + \Gamma_{ab}$$

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The equations of motion

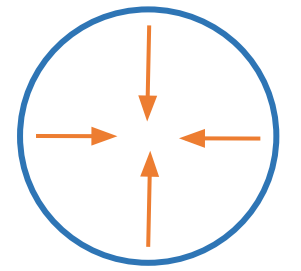
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$$\xi\square\Gamma_{ab} = \Gamma_{ab} - T_{ab}^{(\mathcal{G})}$$

$$\xi\square\Sigma = \Sigma + f'(\phi)\mathcal{G}$$

Spherical collapse in fixed sGB

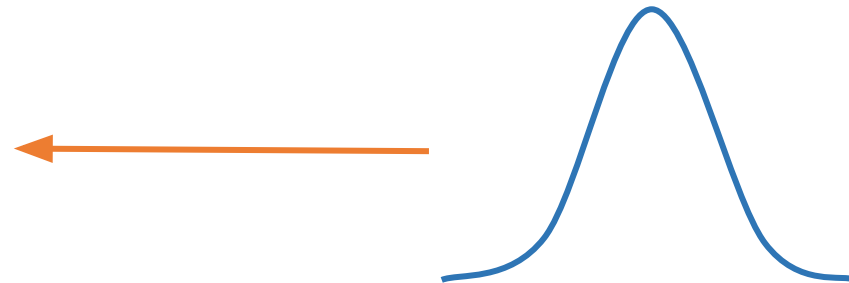


Assume spherical symmetry for metric and scalar

$$ds^2 = -e^{2A(t,r)} dt^2 + e^{2B(t,r)} dr^2 + r^2 d\Omega^2$$

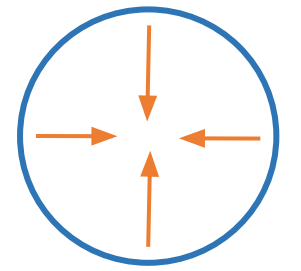
$$\phi = \phi(t, r)$$

Let evolve a gaussian scalar profile under gravity to form an apparent horizon



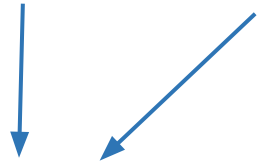
Spherical collapse in fixed sGB

Characteristic speeds



Full theory

$$v_{\pm} = \mathcal{C} \pm \sqrt{\mathcal{D}}$$



Complicated functions with derivatives of metric and scalar

Fixed theory

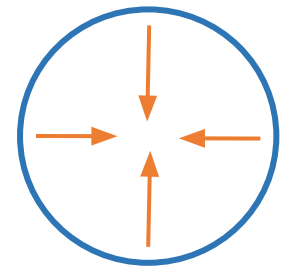
$$v_{\pm} = \pm e^{A-B}$$



Always real and finite

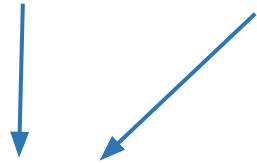
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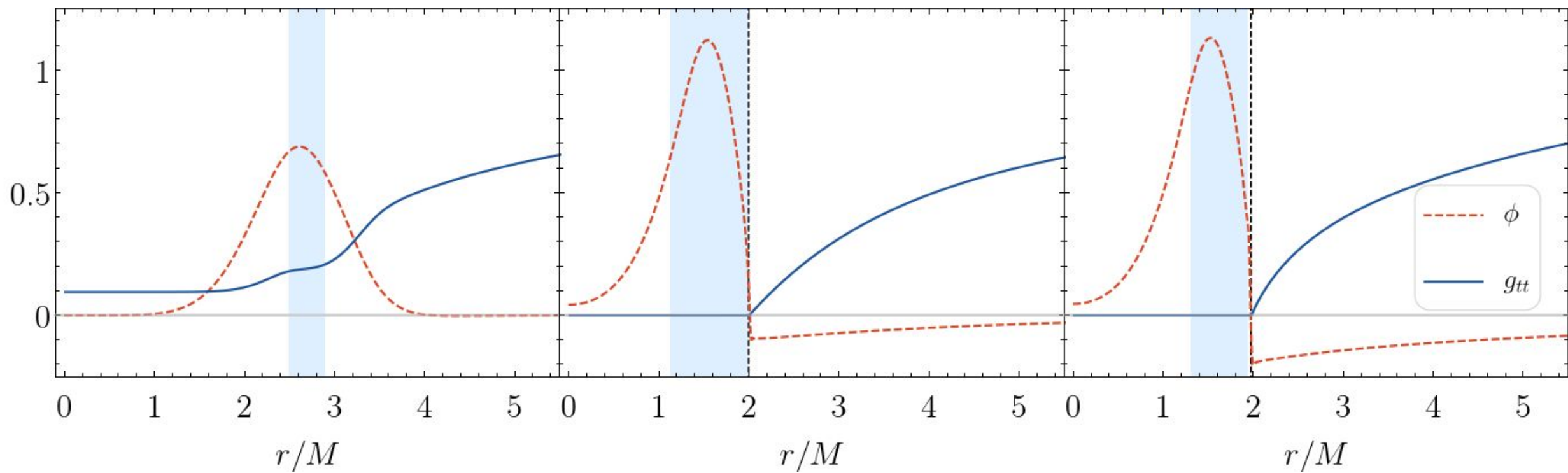
Complicated functions with derivatives of metric and scalar

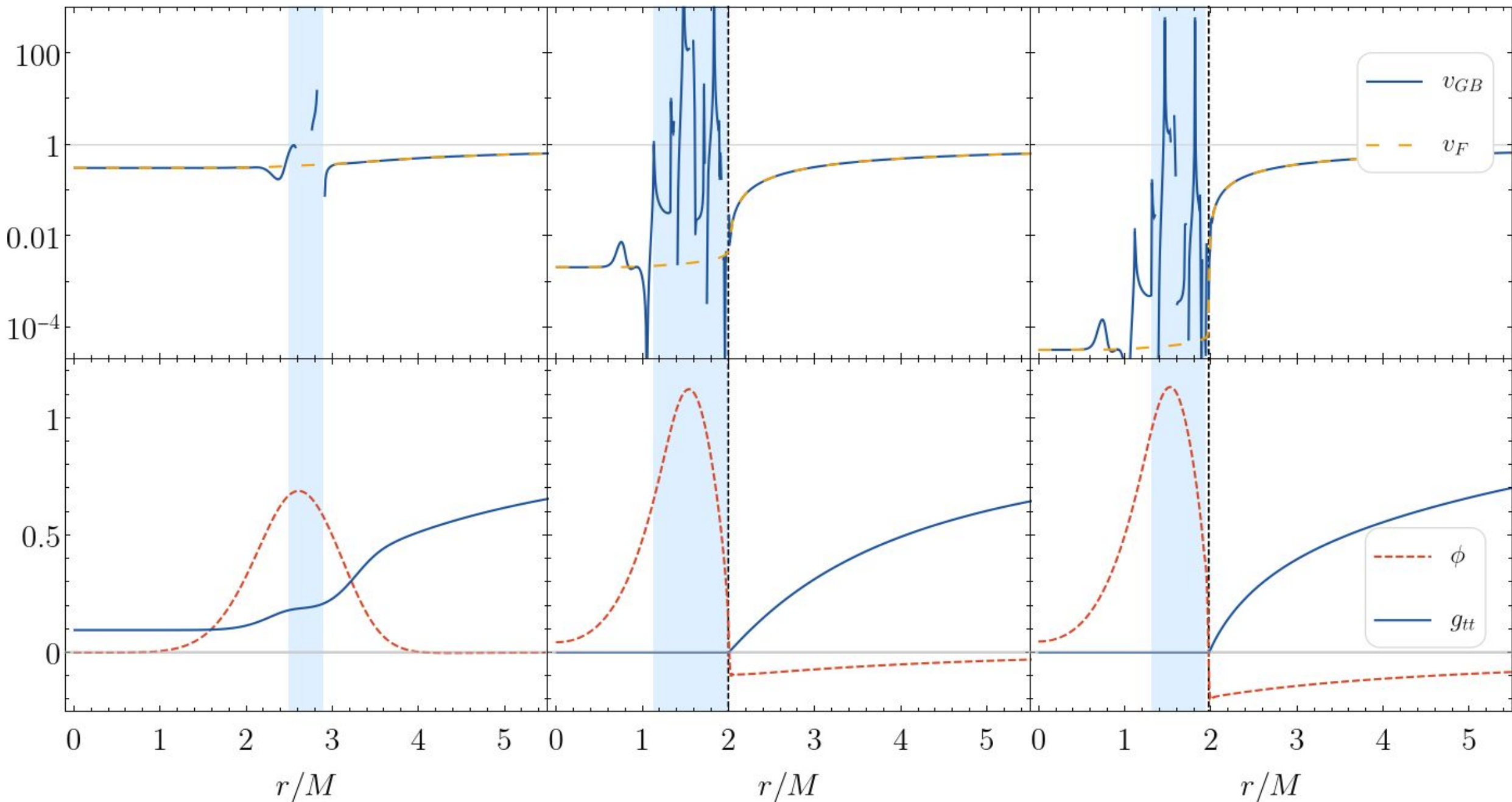
Fixed theory

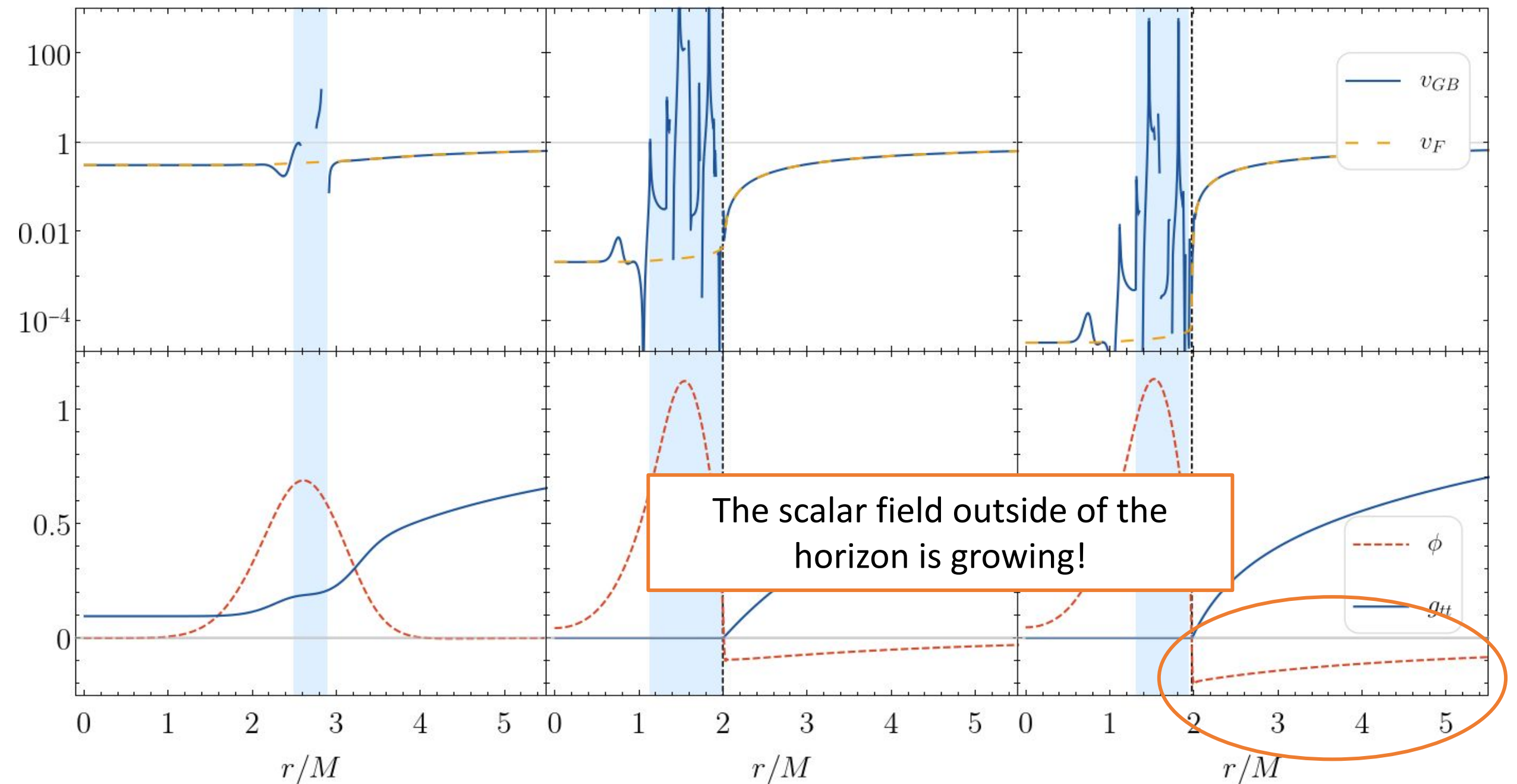
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Always real and finite

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$$S_{GB} = \frac{1}{2} \int \sqrt{-g} \left[R - \frac{1}{2} \partial_a \phi \partial^a \phi + f(\phi) \mathcal{G} \right]$$

$$f(\phi) = \frac{1}{8} \left[\eta \phi^2 + \frac{1}{2} \zeta \phi^4 \right]$$

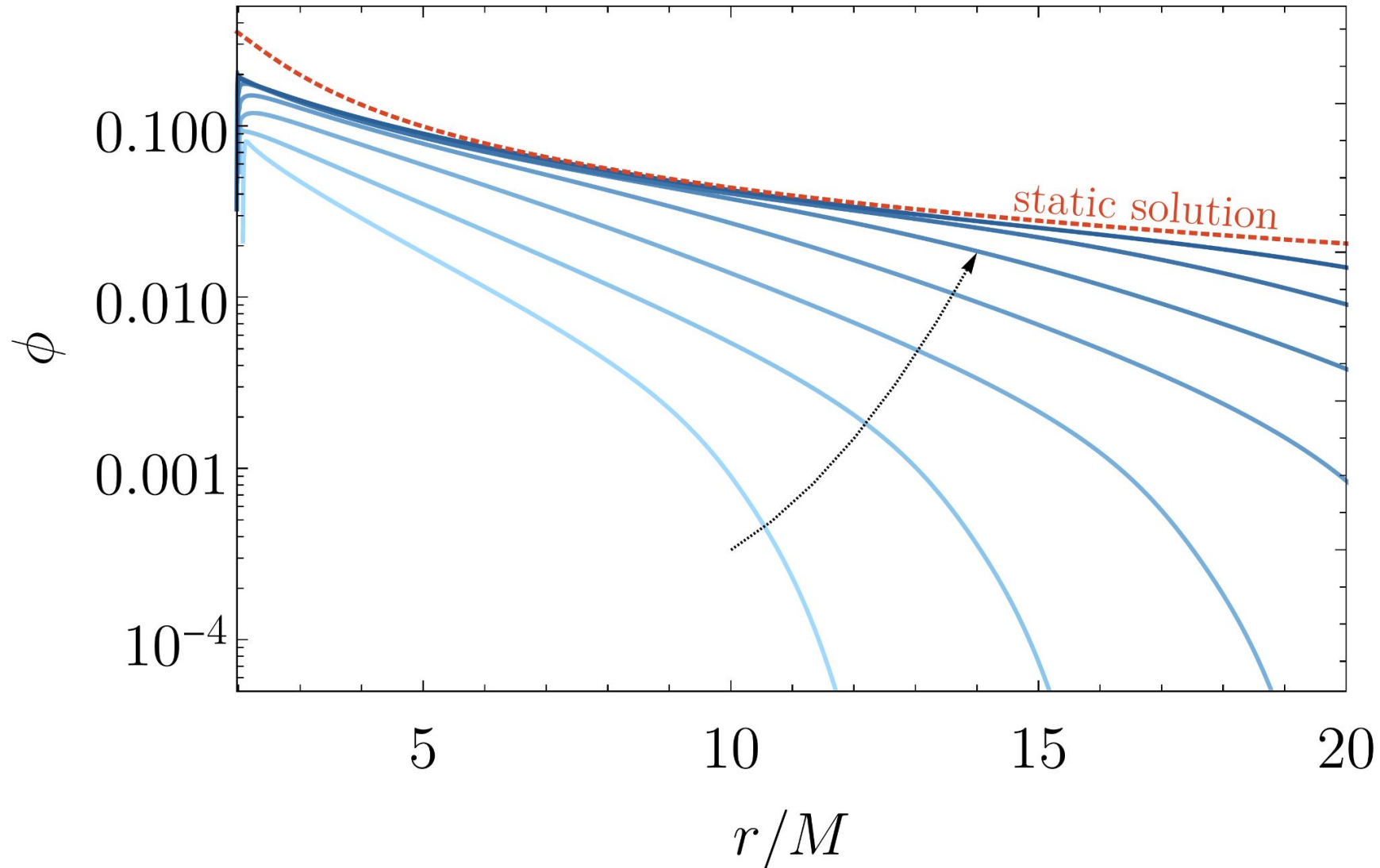
BH scalarization

[Silva, Macedo, Sotiriou, Gualtieri, Sakstein, Berti 2019]

$$\eta = 6$$

$$\zeta = -6\eta$$

$$\xi = \eta$$



Conclusions

- Beyond GR theories can be *ill-posed*
- Possible solution → *fixing-the-equations* [Cayuso, Ortiz, Lehner 2017]
- Spherical collapse faithful in the beginning and in the endstate

Open issues

- Application of the method to 3+1 problems (BH-BH or NS-NS binaries) for the generation of GW templates
- Can we trust the *fixing* in the dynamical region?
- Are all the *fixing* the same?

Thanks for the attention =D