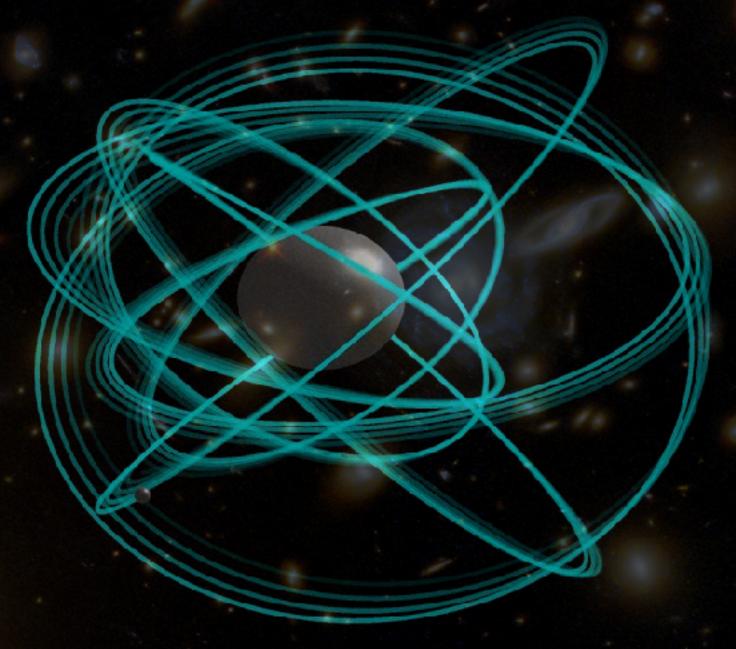


Cosmology with LISA standard sirens and their host galaxies

Orsay - 20/06/2022



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in collaboration with:

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UNDERSTANDING THE COSMIC EXPANSION HISTORY

• According to the standard cosmological model, we can describe the cosmic expansion history in terms of some **cosmological parameters**:

$$\Omega = \{H_0, \Omega_m, \Omega_{\Lambda}, \dots\}$$

Observing astrophysical objects in a flat FLRW metric:

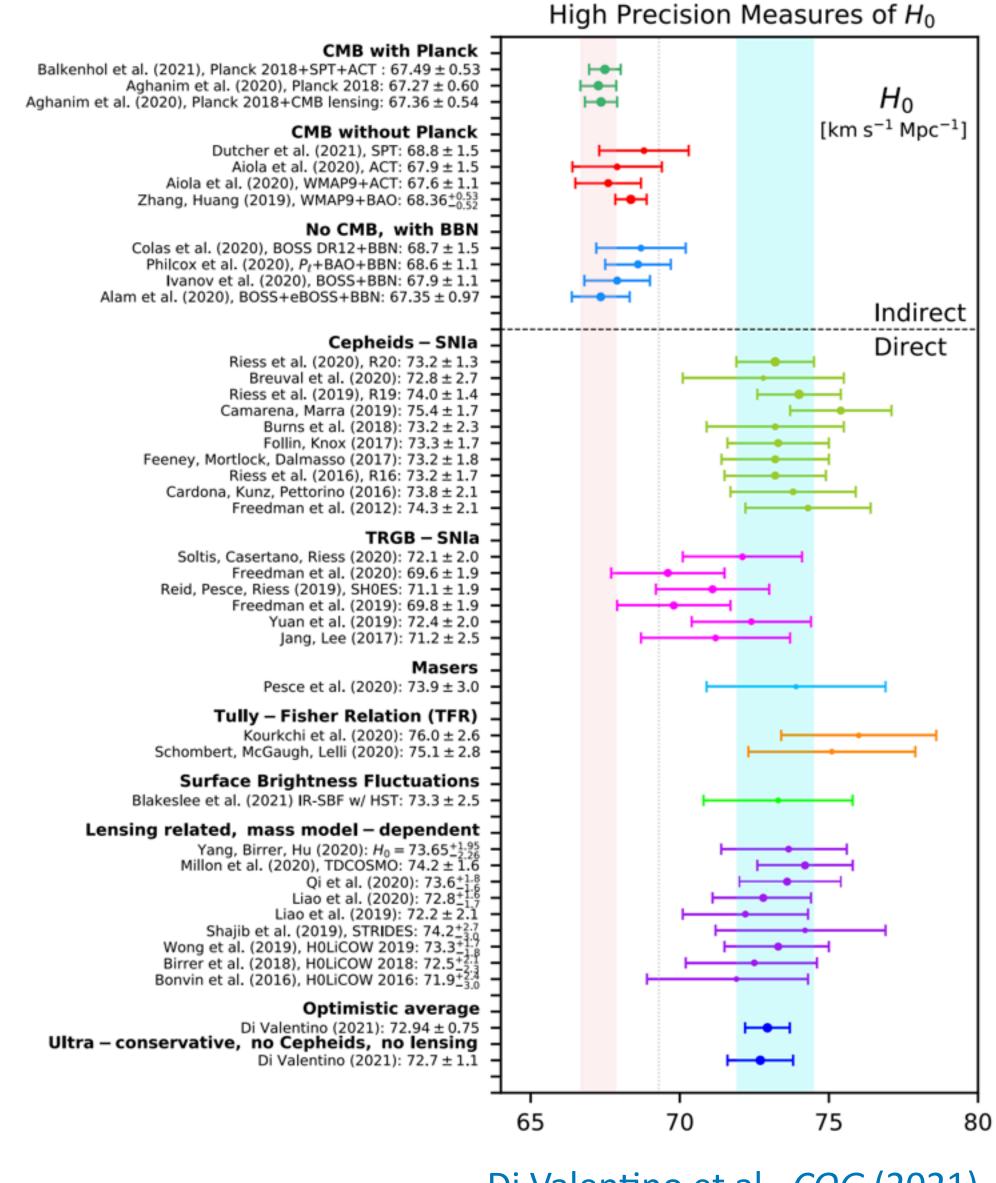
$$d_{L}(\Omega, z) = \frac{c(1+z)}{H_{0}} \int_{0}^{z} \frac{dz'}{\sqrt{\Omega_{m}(1+z')^{3} + \Omega_{\Lambda}(1+z')^{3(1+w_{0}+w_{a})}e^{-3\frac{w_{a}z'}{1+z'}}}}$$

MEASURING THE HUBBLE CONSTANT

 Different Hubble constant measurements do not agree with each other

• 4σ to 6σ disagreement between 'early time' vs 'late time' estimates

 Many proposals to resolve the Hubble puzzle, the matter is still under debate



Di Valentino et al., CQG (2021)

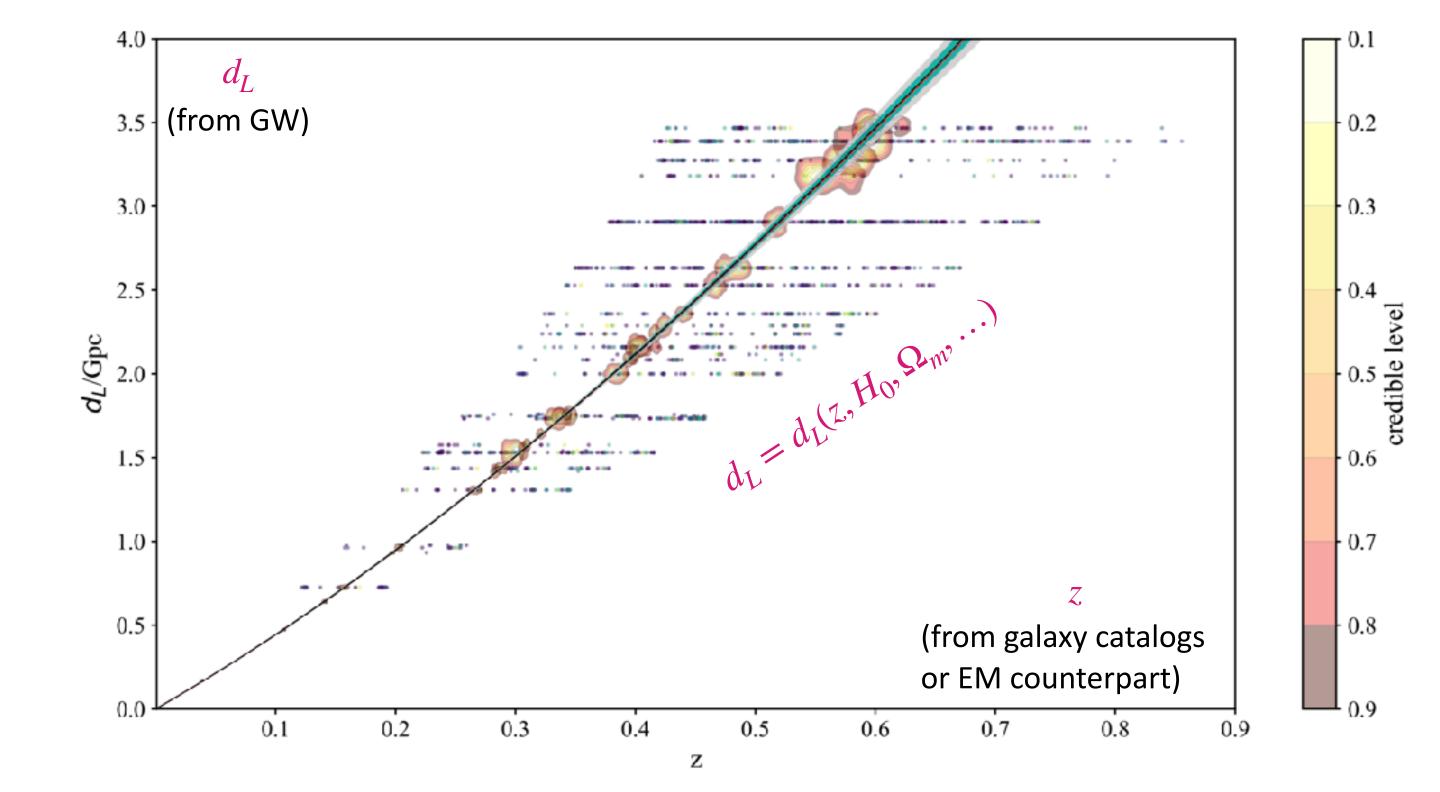
WHY GW COSMOLOGY?

Schutz, Nature (1986)

GWs are "self-calibrated": $h \sim d_L^{-1}$

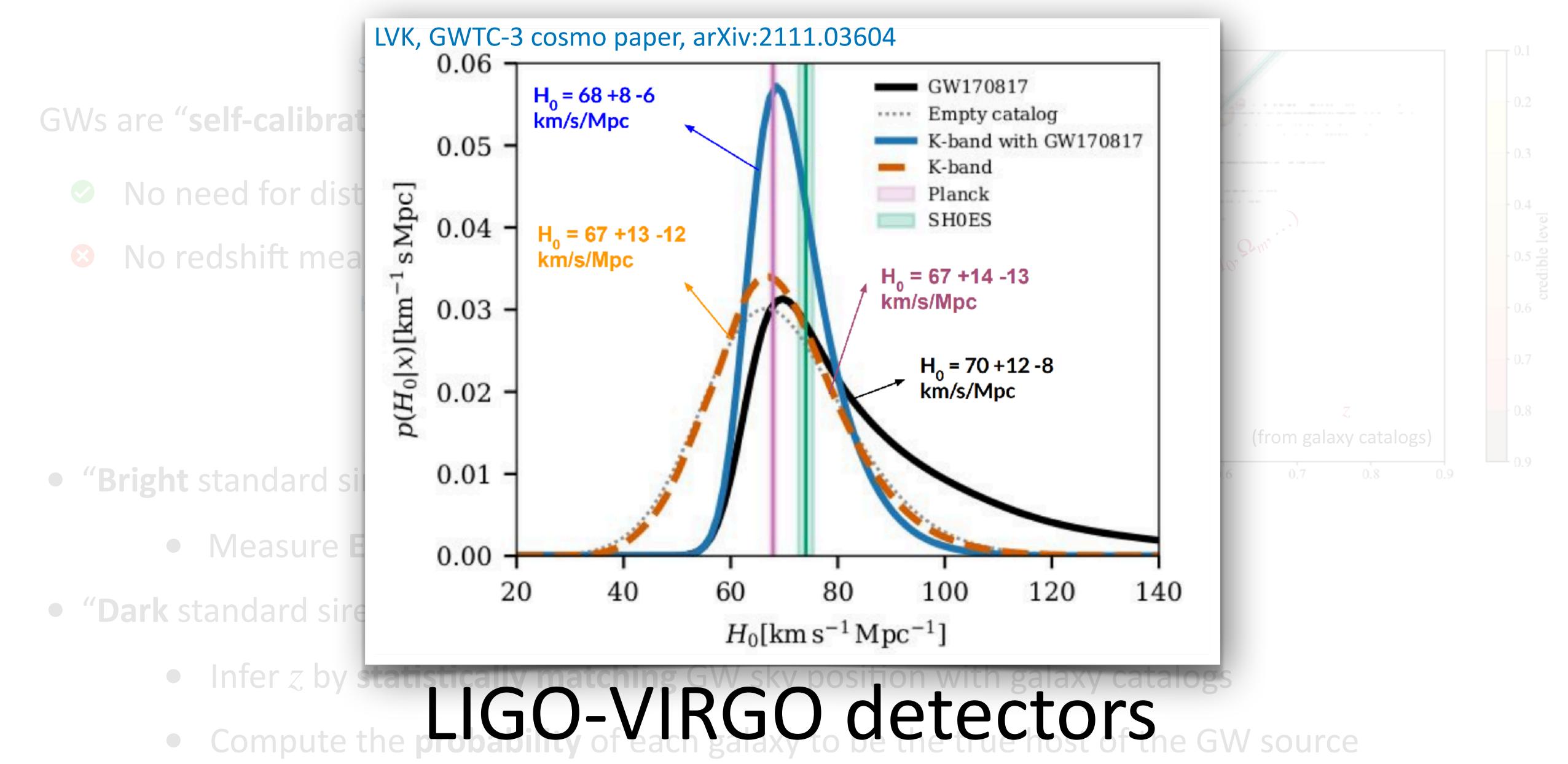
- No need for distance scale ladder
- No redshift measurement from GWs

Krolak, Schutz, GRG (1987)



- "Bright standard sirens":
 - Measure **EM counterpart** of the galaxy host and obtain z
- "Dark standard sirens":
 - ullet Infer z by **statistically matching** GW sky position with galaxy catalogs
 - Compute the probability of each galaxy to be the true host of the GW source

WHY GW COSMOLOGY?

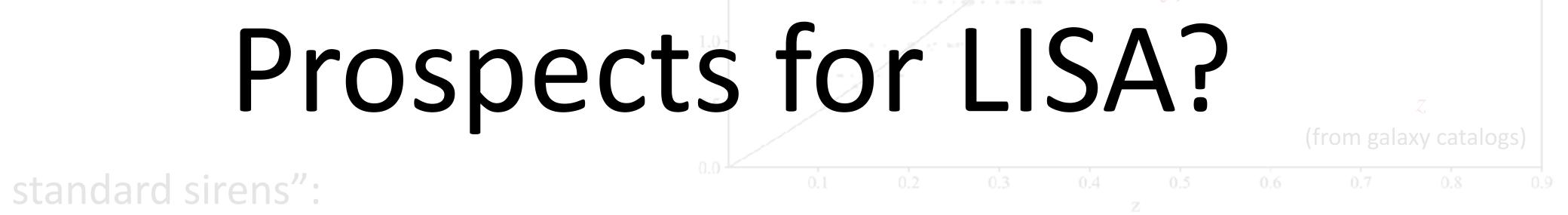


WHY GW COSMOLOGY?

(from LISA)

GWs are "self-calibrated": $h \sim d_{\rm T}^{-1}$

- No need for distance scale ladder
- No redshift measurement from GWs

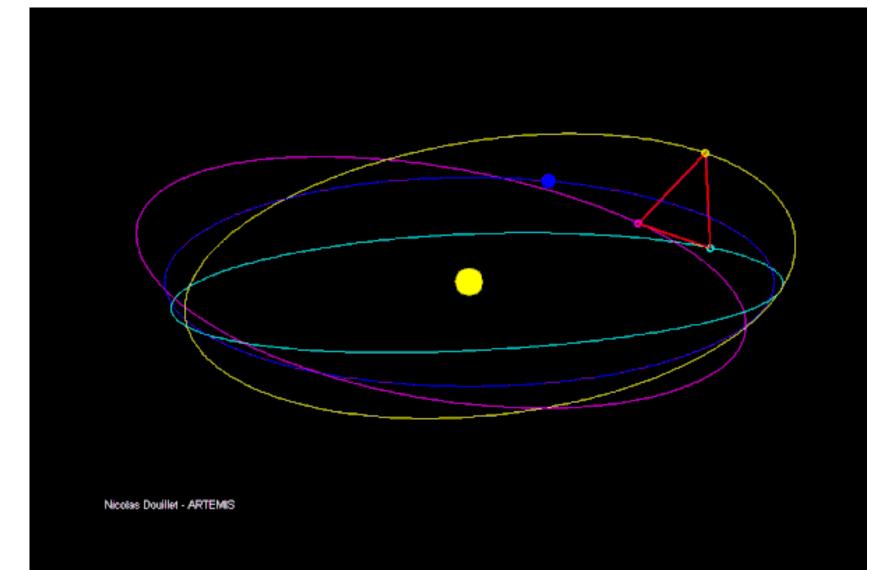


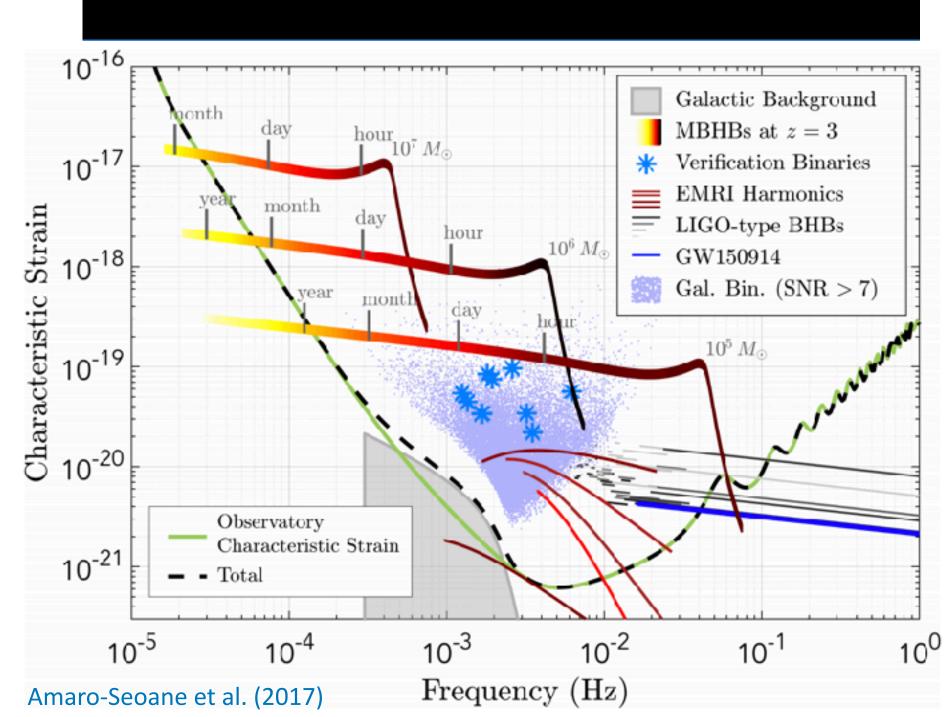
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LASER INTERFEROMETER SPACE ANTENNA

LISA will be the first space-based
 GW detector (expected launch in 2034)

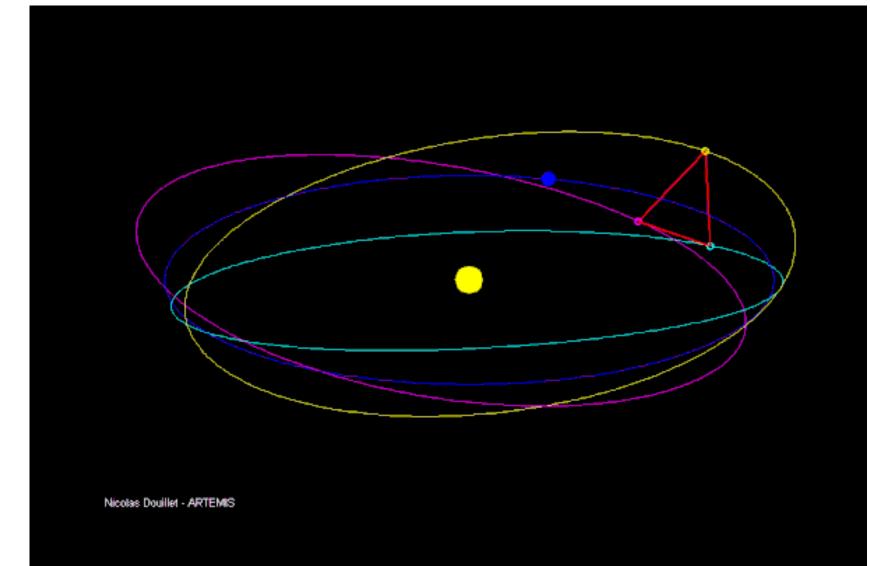
- LISA will observe GWs in a yet unexplored frequency range $(10^{-4} 10^{-1} \text{ Hz})$
- LISA will detect compact binary coalescences up to very high redshift

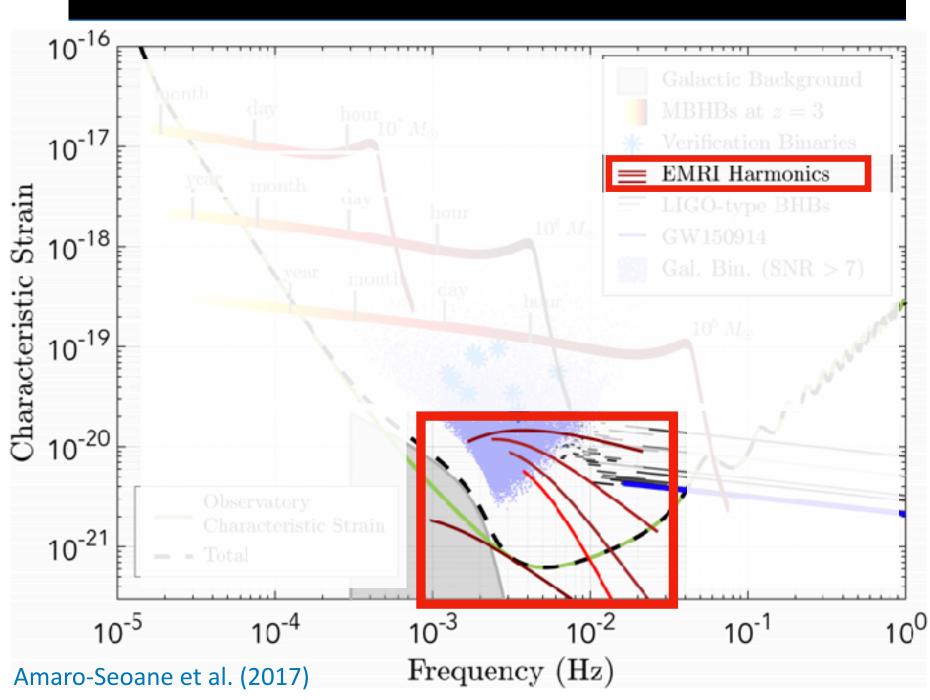




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EXTREME MASS-RATIO INSPIRALS

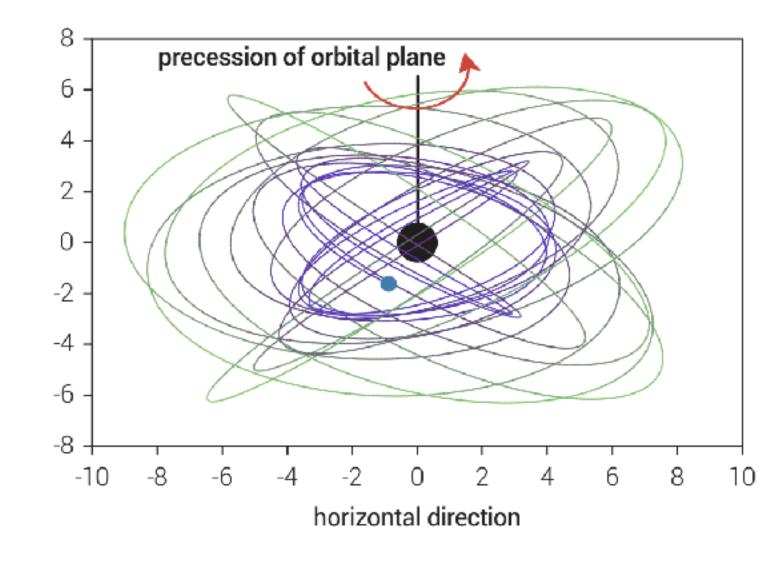
Binary systems with mass-ratio $q \sim 10^{-6} - 10^{-3}$

- Massive BH $(10^4 M_{\odot} 10^7 M_{\odot})$
- Compact object $(10 M_{\odot})$

Slow inspiral, $10^4 - 10^5$ orbital cycles in the final year before plunge

Extremely accurate measurements of the system parameters

No EM counterpart



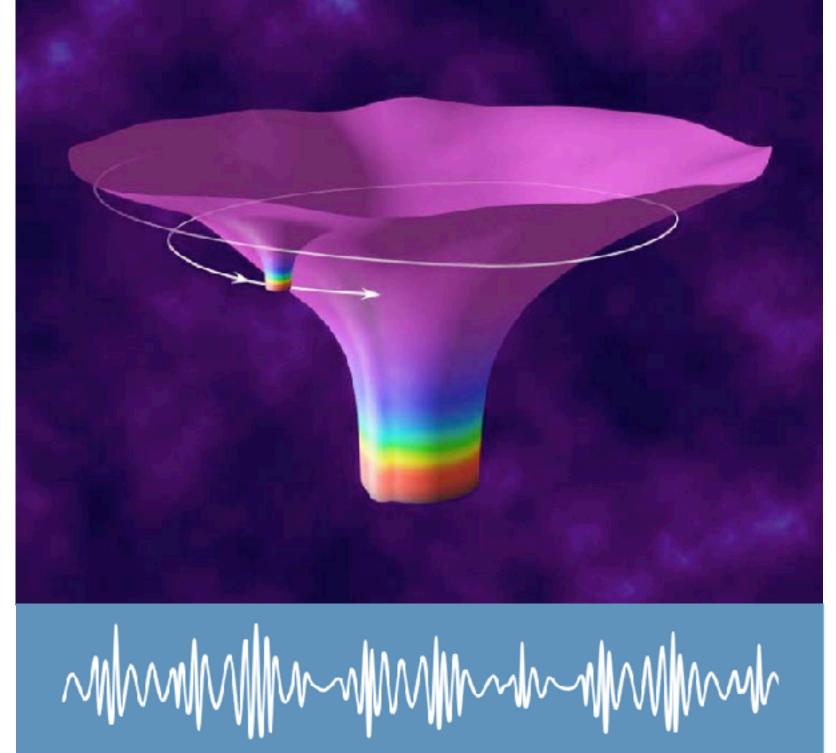


Figure 4: An artist's impression of the spacetime of an extreme-massratio inspiral and a representative waveform of the expected gravitational waves. A smaller black hole orbits around a supermassive black hole. *Credit: NASA*.

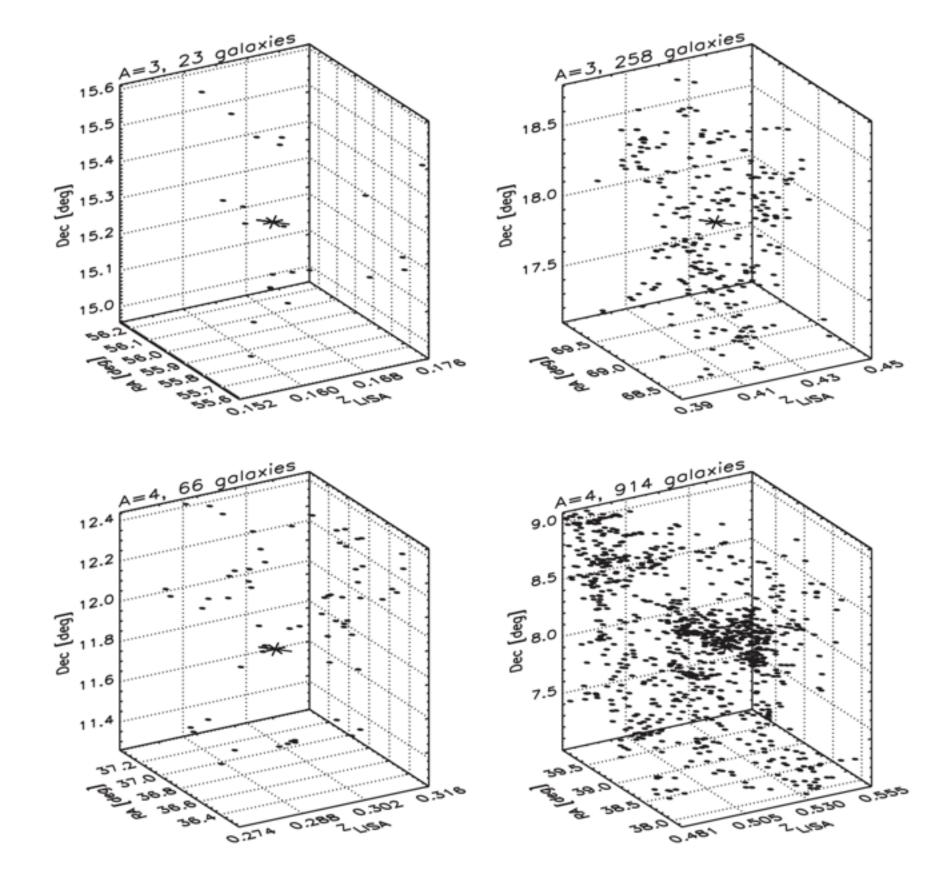
eLISA White Paper, arXiv:1305.5720

PREVIOUS STUDIES

Macleod, Hogan, PRD (2008):

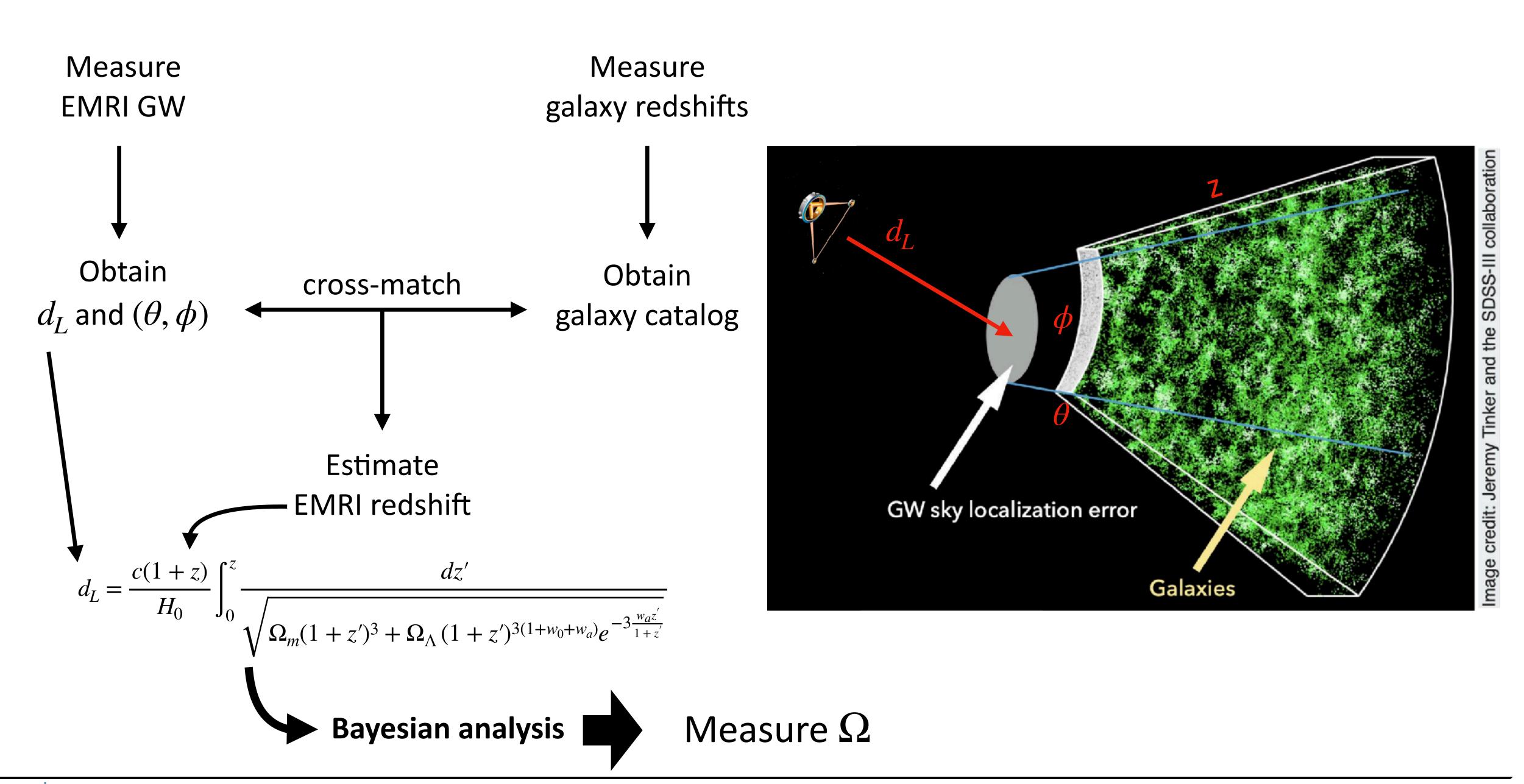
 H_0 at 1% with 20 EMRIs at z < 0.5

- assume only linear cosmic expansion
- assume old 5 Gm LISA configuration
- no PE on the GW signals
- no Bayesian inference framework



Macleod, Hogan, PRD (2008)

EMRIS AS DARK STANDARD SIRENS



HOW MANY EMRIS WILL WE OBSERVE?

EMRI rates span 2-3 orders of magnitudes, reflecting variations in:

- MBH population: semi-analytic models, realistic/pessimistic
- Stellar clusters distributions around MBHs
- EMRI's orbit parameters

Model	Mass function	MBH spin	Cusp erosion	$M - \sigma$ relation	$N_{ m p}$	$\mathrm{CO}_{\mathrm{mass}\;[M_{\odot}]}$	Total	EMRI rate [yr ⁻¹] Detected (AKK)	Detected (AKS)
M1	Barausse12	a98	yes	Gultekin09	10	10	1600	294	189
M2	Barausse12	a98	yes	KormendyHo13	10	10	1400	220	146
M3	${\bf Barausse 12}$	a98	yes	GrahamScott13	10	10	2770	809	440
M4	${\bf Barausse 12}$	a98	yes	Gultekin09	10	30	520 (620)	260	221
M5	Gair10	a98	no	Gultekin09	10	10	140	47	15
M6	${\bf Barausse 12}$	a98	no	Gultekin09	10	10	2080	479	261
M7	${\bf Barausse 12}$	a98	yes	Gultekin09	0	10	15800	2712	1765
M8	${\bf Barausse 12}$	a98	yes	Gultekin09	100	10	180	35	24
M9	${\bf Barausse 12}$	aflat	yes	Gultekin09	10	10	1530	217	177
M10	${\bf Barausse 12}$	a0	yes	Gultekin09	10	10	1520	188	188
M11	Gair10	a0	no	Gultekin09	100	10	13	1	1
M12	Barausse12	a98	no	Gultekin09	0	10	20000	4219	2279

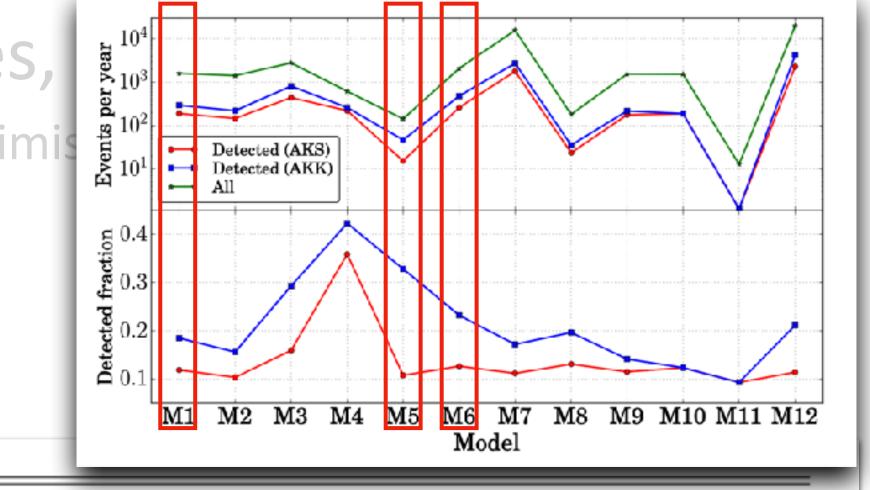
Babak et al., PRD (2017)



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		M11	Gair10	a0	no	Gultekin09	100	10	13	1	1
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Babak et al., PRD (2017)

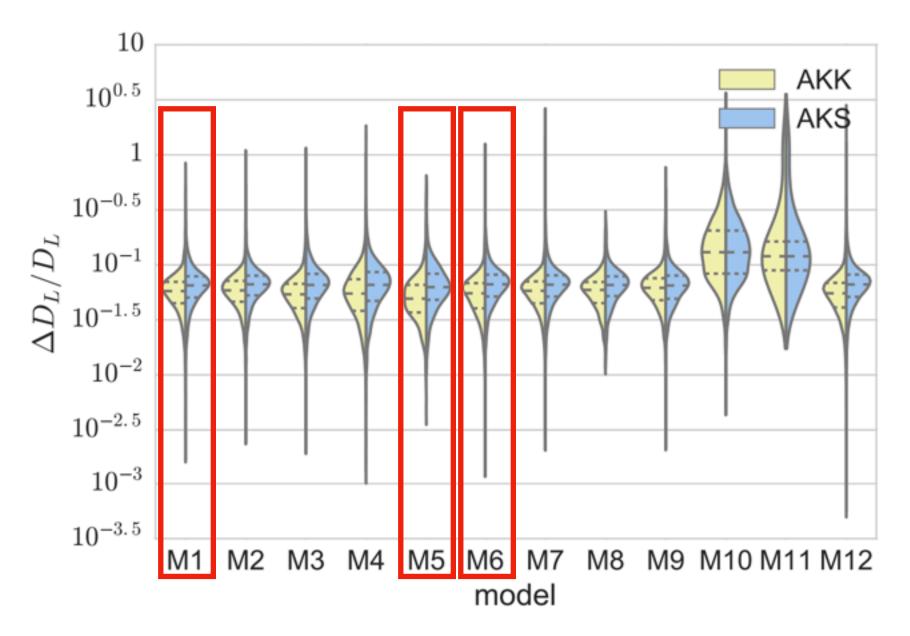
HOW WELL CAN WE LOCALIZE EMRIs?

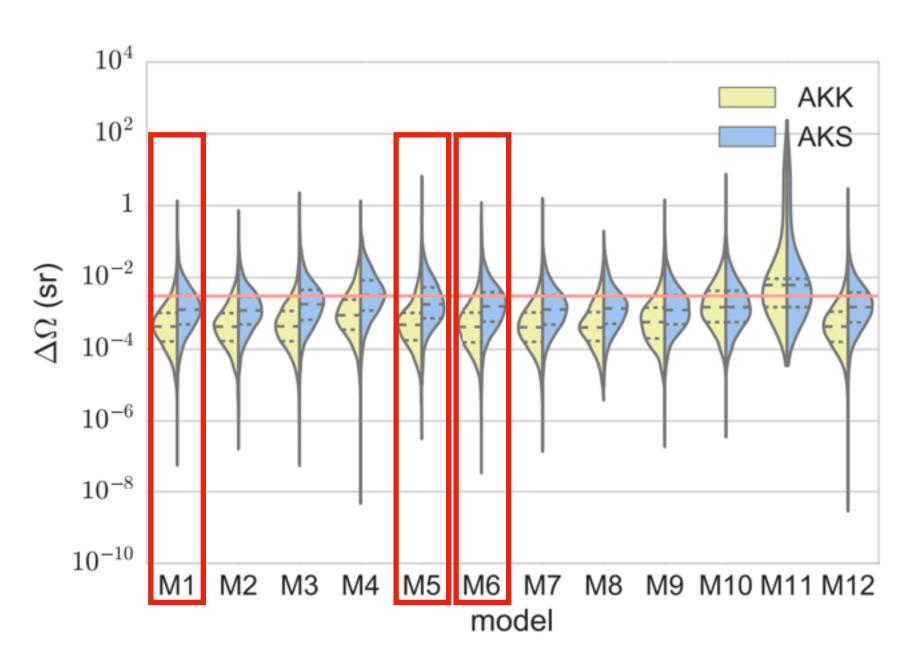
EMRI PE: catalogs of Babak et al. provide best estimates and uncertainties for:

$$d_L \pm \sigma_{d_L} \qquad \phi \pm \sigma_{\phi} \qquad \theta \pm \sigma_{\theta}$$

$$\Delta d_L/d_L \sim 10^{-1}$$

 $\Delta \Omega/\Omega \sim 10 \, \mathrm{deg}^2$





Babak et al., PRD (2017)

ERROR-BOXES

Flux-limited, full-sky galaxy simulations of Henriques et al., *MNRAS* (2012)

based on the Millennium Run

Springel et al., *Nature* (2005)

For a given cosmology:

$$\hat{d}_L \pm \Delta \hat{d}_L \longrightarrow z \pm \Delta z$$

Assuming cosmological priors:

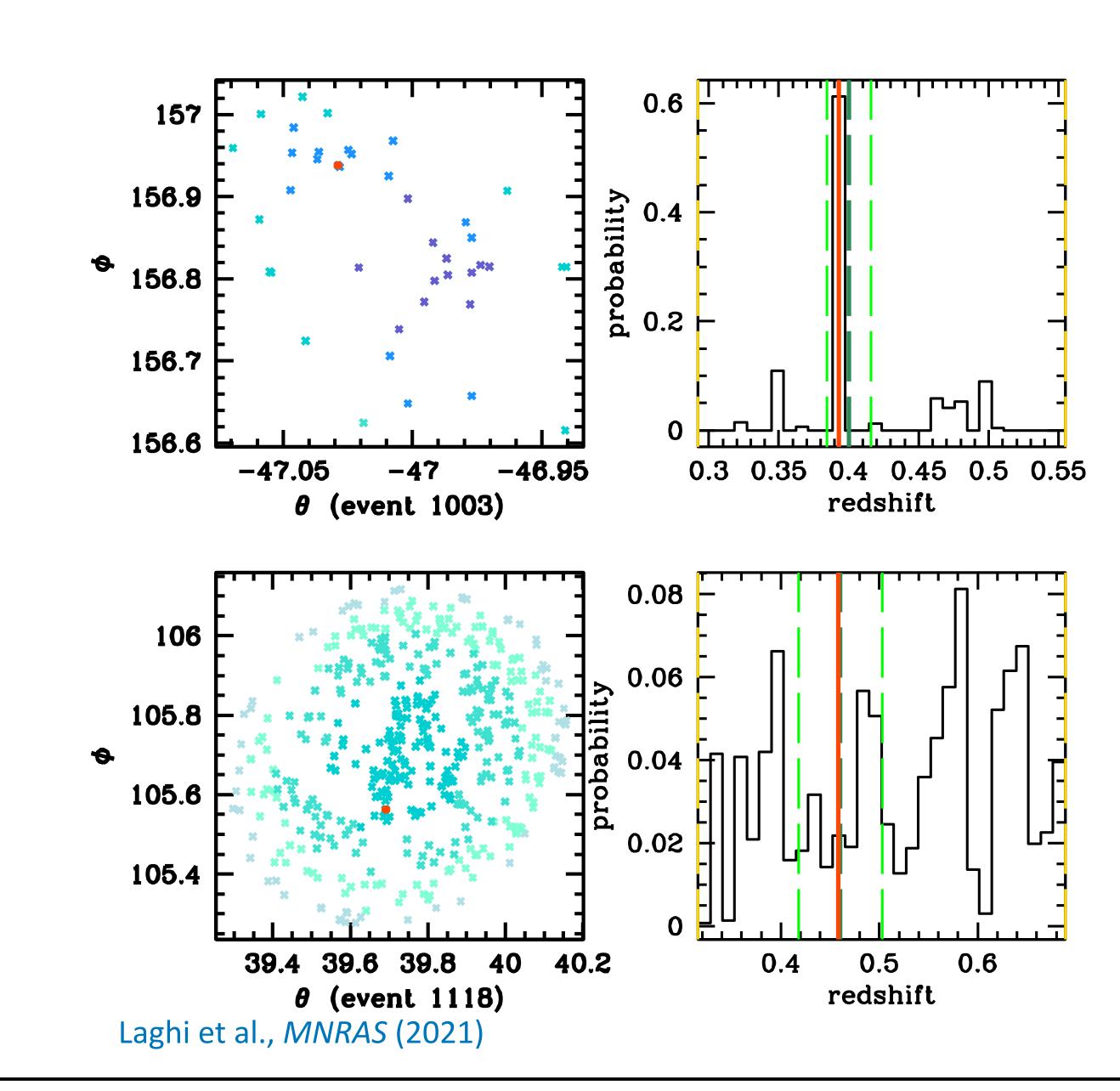
$$[z^-, z^+]$$

Accounting for galaxy peculiar velocities:

$$[z^{-} - \Delta z_{v_p}^{-}, z^{+} + \Delta z_{v_p}^{+}]$$

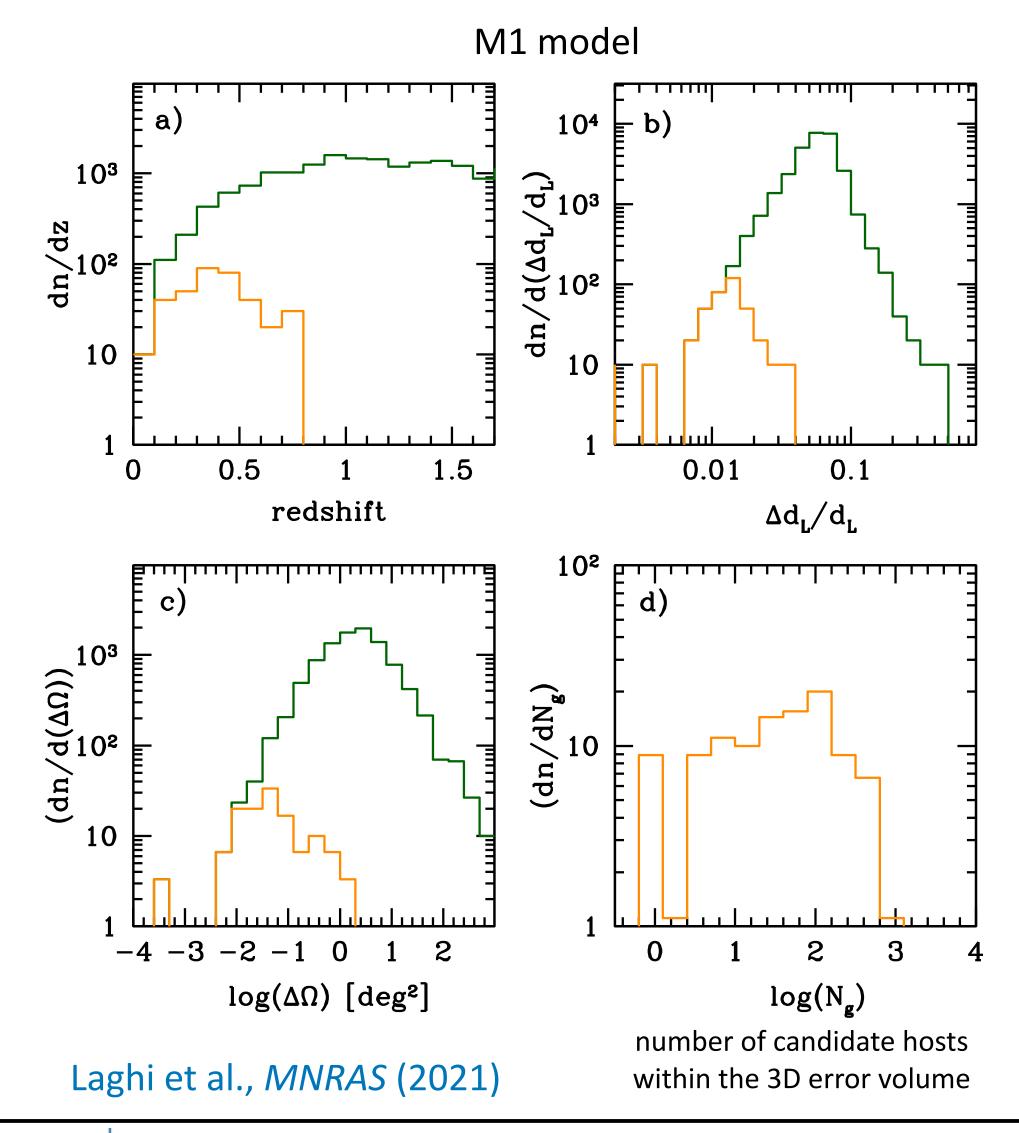
• EMRI error box:

$$\Delta\Omega_{sky} \times [z^- - \Delta z_{v_p}^-, z^+ + \Delta z_{v_p}^+]$$

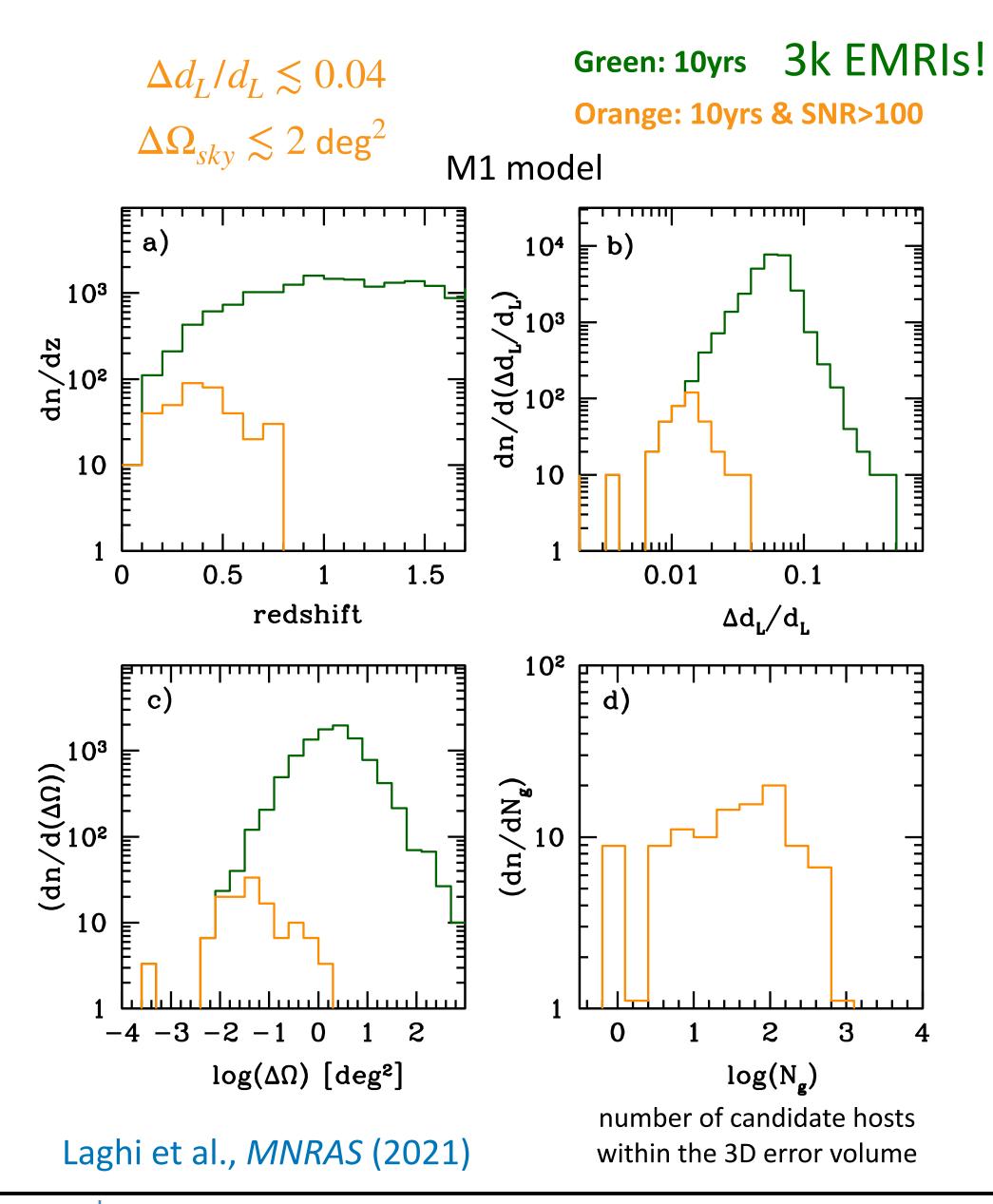


SELECTING EVENTS

Green: 10yrs 3k EMRIs!



SELECTING EVENTS

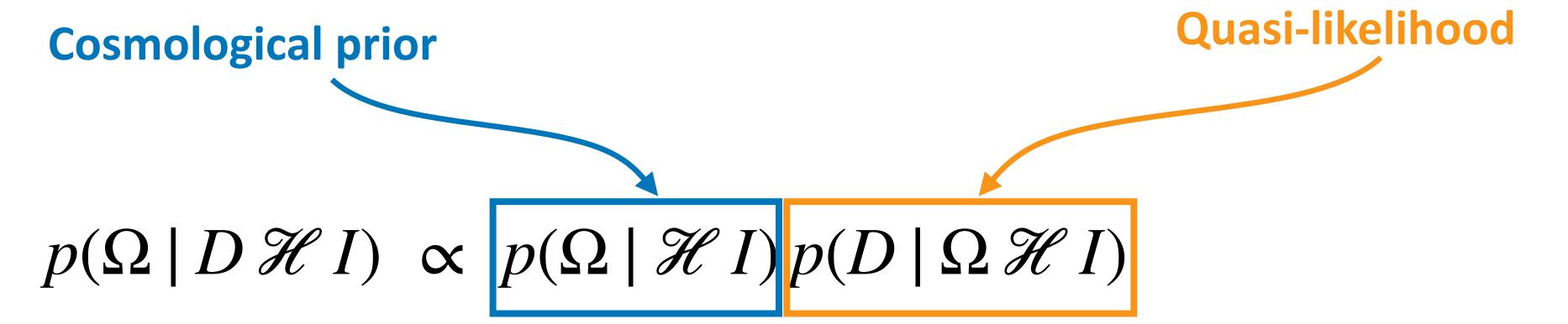


Require **SNR>100**:

- Well-localised, most-informative events
- Few hosts per error-box

(SNR>100)	Events
M5 (pessimistic)	O(5)
M1 (fiducial)	O(30)
M6 (optimistic)	O(70)

BAYESIAN INFERENCE

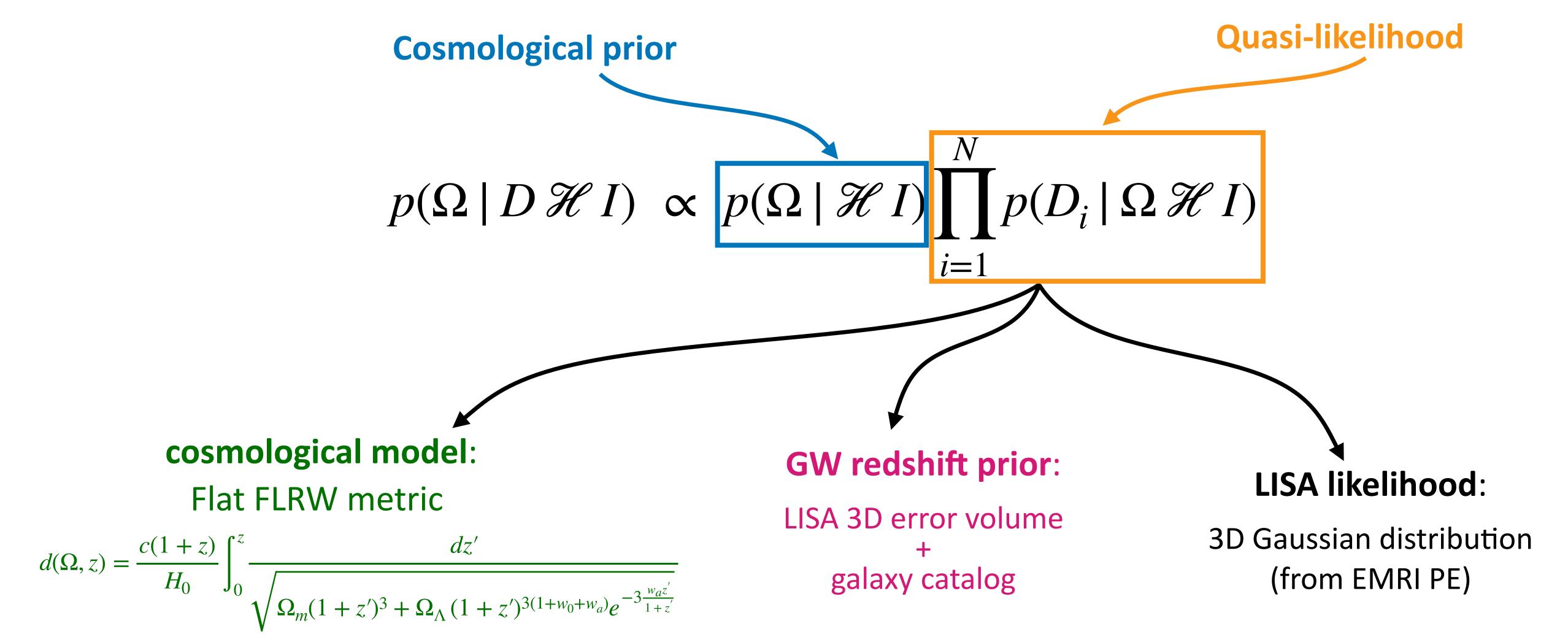


D = GW data

 $\mathcal{H} = cosmological model$

I =any information available

BAYESIAN INFERENCE



$$D = GW data$$

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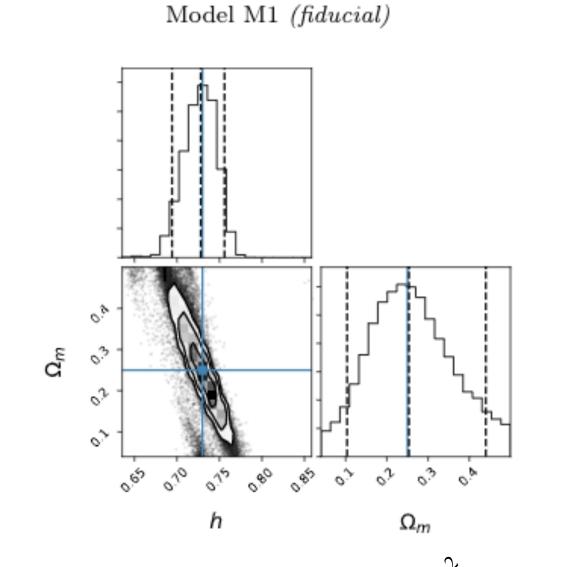


RESULTS: ACDM

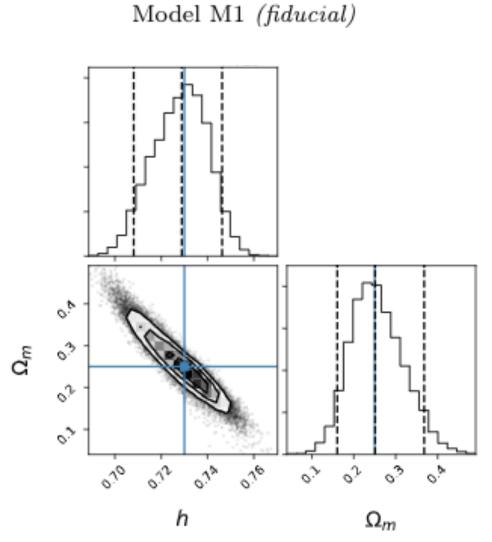
EMRIs will be excellent probes of H_0

h accuracy (90% CI) 1-6%

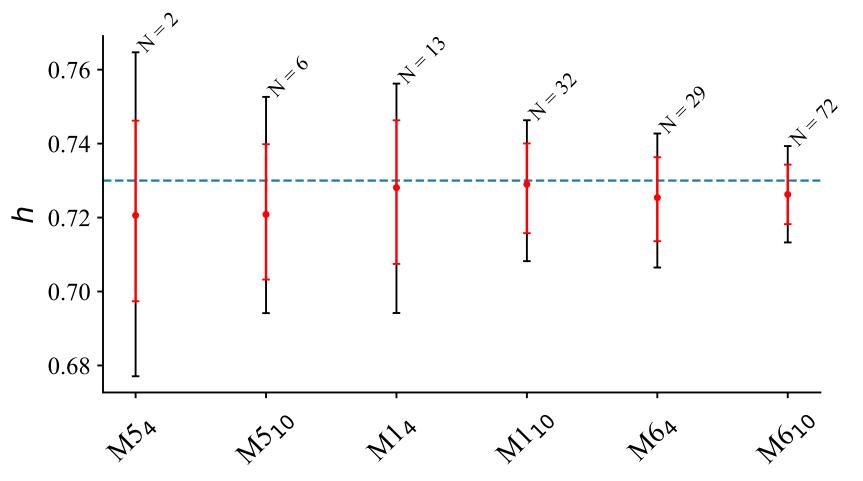
 Ω_m accuracy (90% CI) 25% at most

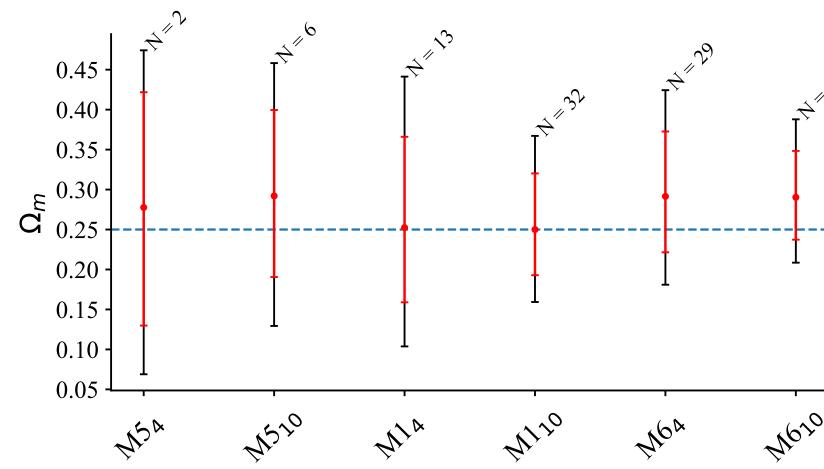


 Λ CDM, 4yr



 Λ CDM, 10yr





Analysis done with cosmoLISA

Del Pozzo, Laghi [https://github.com/wdpozzo/cosmolisa]

 $h = H_0/100 \, \mathrm{km}^{-1} \mathrm{s} \, \mathrm{Mpc}$

RESULTS: DE

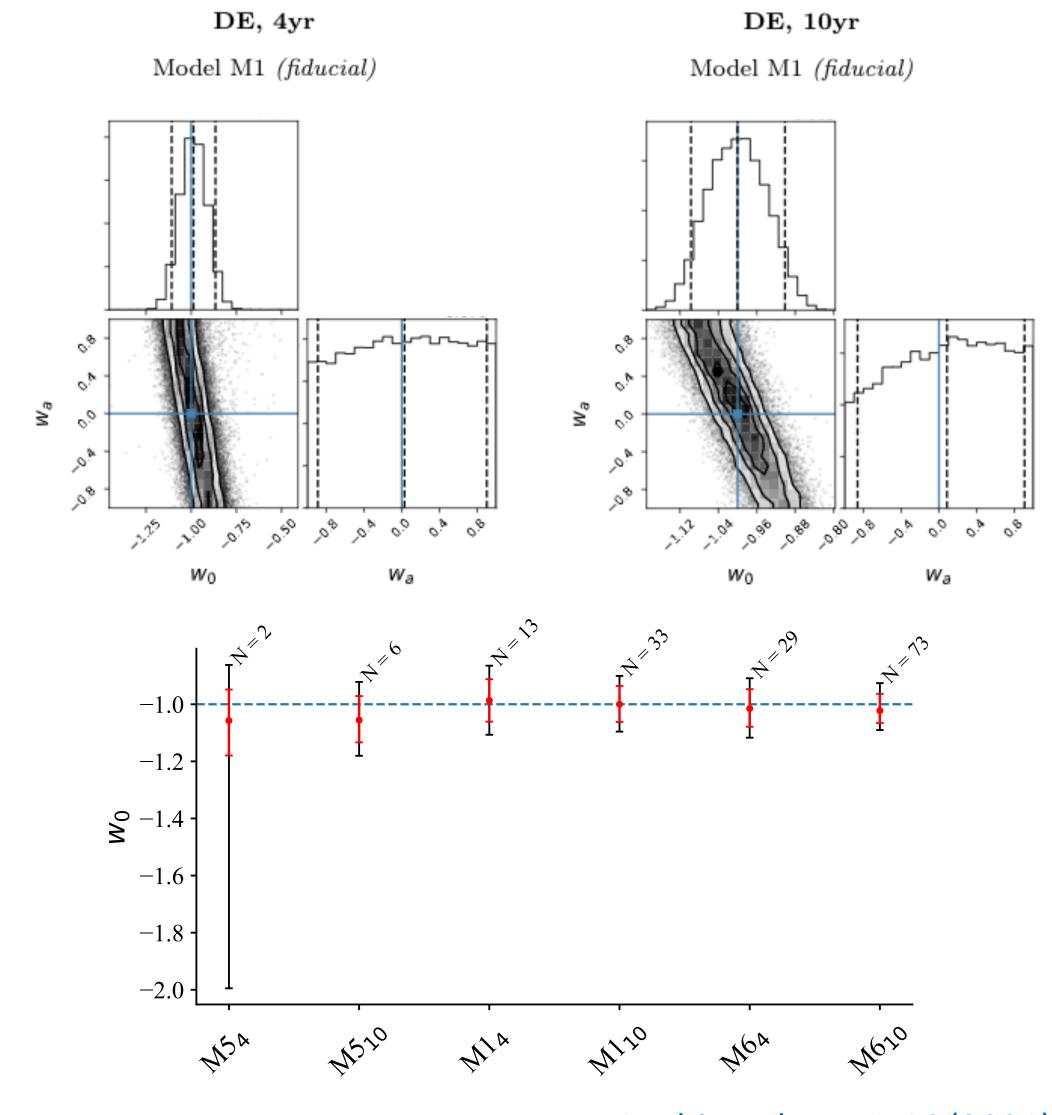
EMRIs can constrain w_0

$$w(z) = w_0 + w_a z/(1+z)$$

 w_0 accuracy (90% CI) 10% at most

Analysis done with cosmoLISA

Del Pozzo, Laghi [https://github.com/wdpozzo/cosmolisa]



Laghi et al., MNRAS (2021)

ANOTHER LISA CBC SOURCE: MASSIVE BLACK HOLE BINARIES

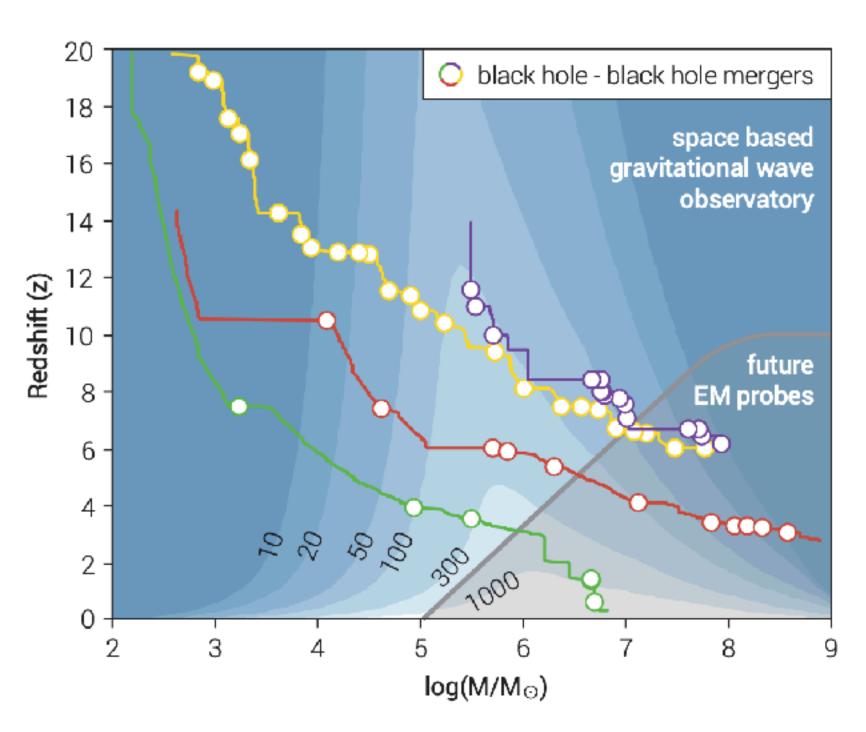


Figure 2: Constant-contour levels of the sky and polarisation angle-averaged SNR for eLISA, for equal mass non-spinning binaries as a function of their total rest frame mass, M, and cosmological redshift, z. The tracks represent the mass-redshift evolution of selected supermassive black holes: two possible evolutionary paths for a black hole powering a $z \sim 6$ QSO (starting from a massive seed, blue curve, or from a Pop III seed from a collapsed metal-free star, yellow curve); a typical $10^9 M_{\odot}$ black hole in a giant elliptical galaxy (red curve); and a Milky Way-like black hole (green curve). Circles mark black hole-black hole mergers occurring along the way. These were obtained using state of the art semi-analytical merger tree models [65]. The grey transparent area in the bottom right corner roughly identifies the parameter space for which massive black holes might power phenomena that will likely be observable by future electromagnetic probes.

eLISA White Paper, arXiv:1305.5720

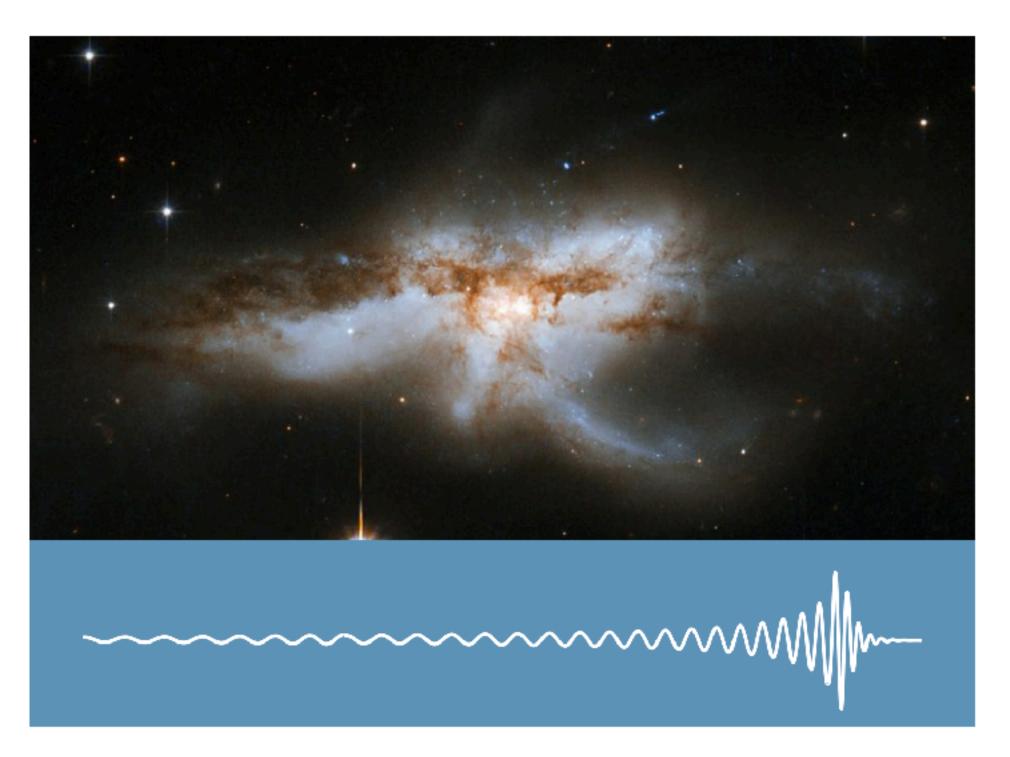
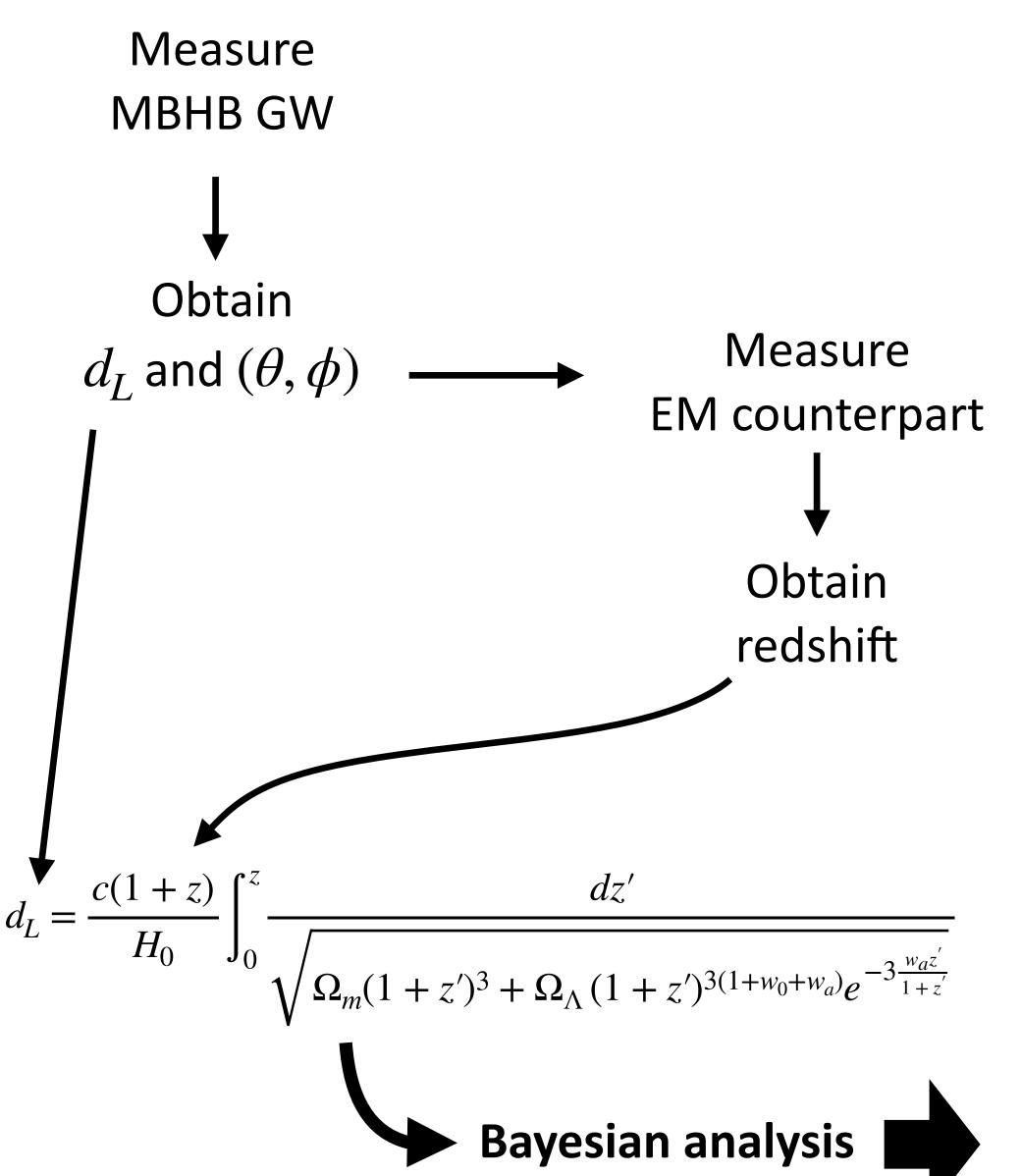
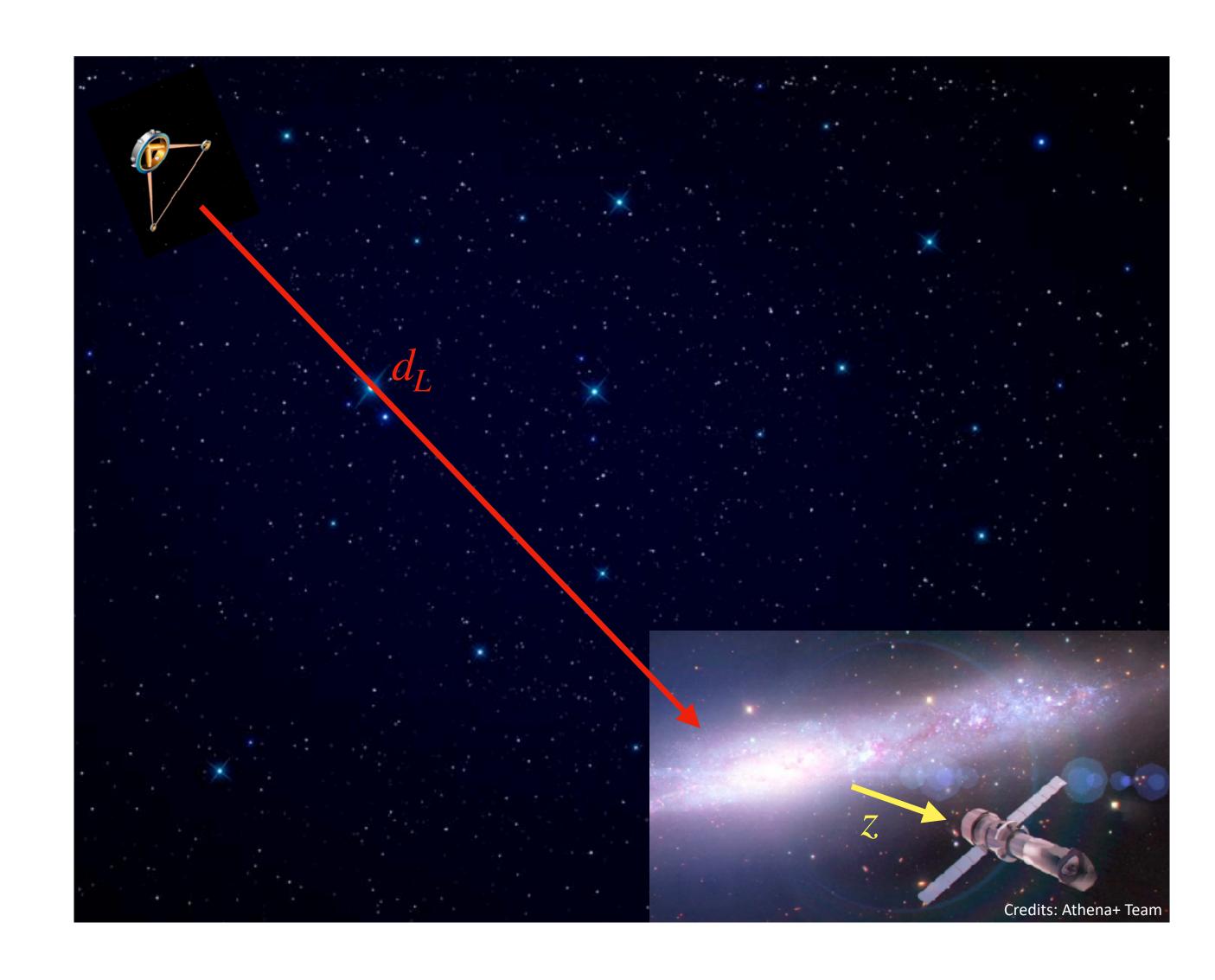


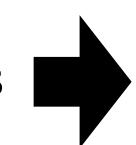
Figure 1: Merging galaxy NGC 6240 and a representative waveform of the expected gravitational waves from the coalesence of two supermassive black holes. Observations with NASA's Chandra X-ray observatory have disclosed two giant black holes inside NGC 6240. They will drift toward one another and eventually merge into a larger black hole. *Credit:* NASA, ESA, the Hubble Heritage (STScI/AURA)-ESA/Hubble Collaboration, and A. Evans (University of Virginia, Charlottesville/NRAO/Stony Brook University).

MBHBs AS BRIGHT STANDARD SIRENS



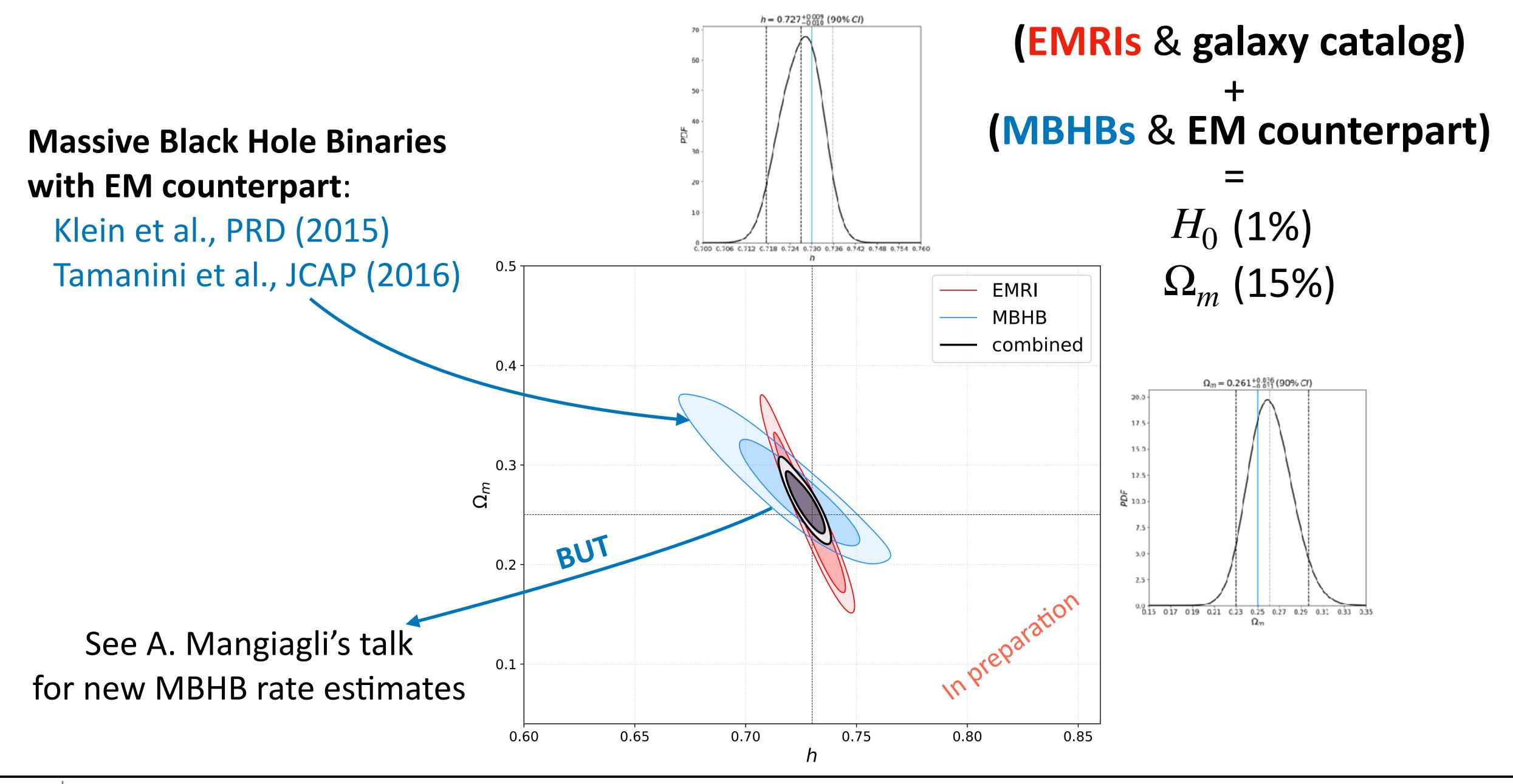






Measure Ω

EMRIS + MBHBs FOR 10 YEARS OF OBSERVATION





LISA COSMOLOGICAL ENCHILADA

