

The H_0 Olympics: a fair ranking of proposed models

Guillermo Franco Abellán

Laboratoire Univers et Particules de Montpellier

Based on:

arXiv:2107.10291, accepted in Physics Reports

In collaboration with [Nils Schöneberg](#), [Andrea Pérez Sánchez](#), [Samuel J. Witte](#), [Vivian Poulin](#) and [Julien Lesgourgues](#)



Tensions in cosmology

With the era of precision cosmology, several discrepancies have emerged

- S_8 with weak-lensing data ($2-3\sigma$)
[KiDS-1000 2007.15632](#)
- H_0 with local measurements (5σ)
[Riess++ 2012.08534](#)

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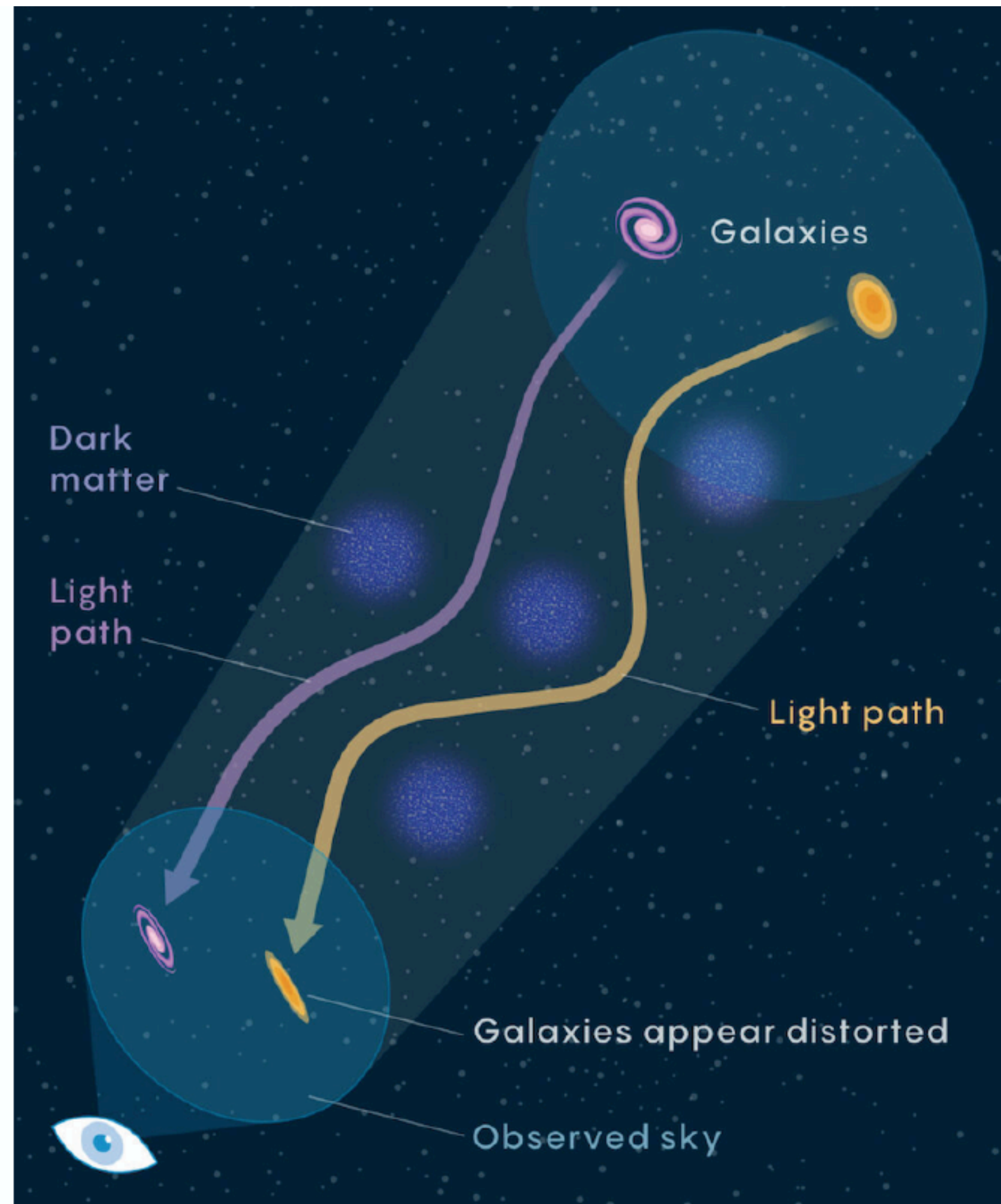
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- H_0 with local measurements (5σ)
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Unaccounted systematics?

Physics beyond Λ CDM?

The S_8 tension

Weak-lensing surveys are mainly sensible to $S_8 \equiv \sigma_8 \sqrt{\Omega_m/0.3}$



KiDS+BOSS+2dfLenS*:

$$S_8 = 0.766^{+0.020}_{-0.014}$$

Planck (*under* Λ CDM):

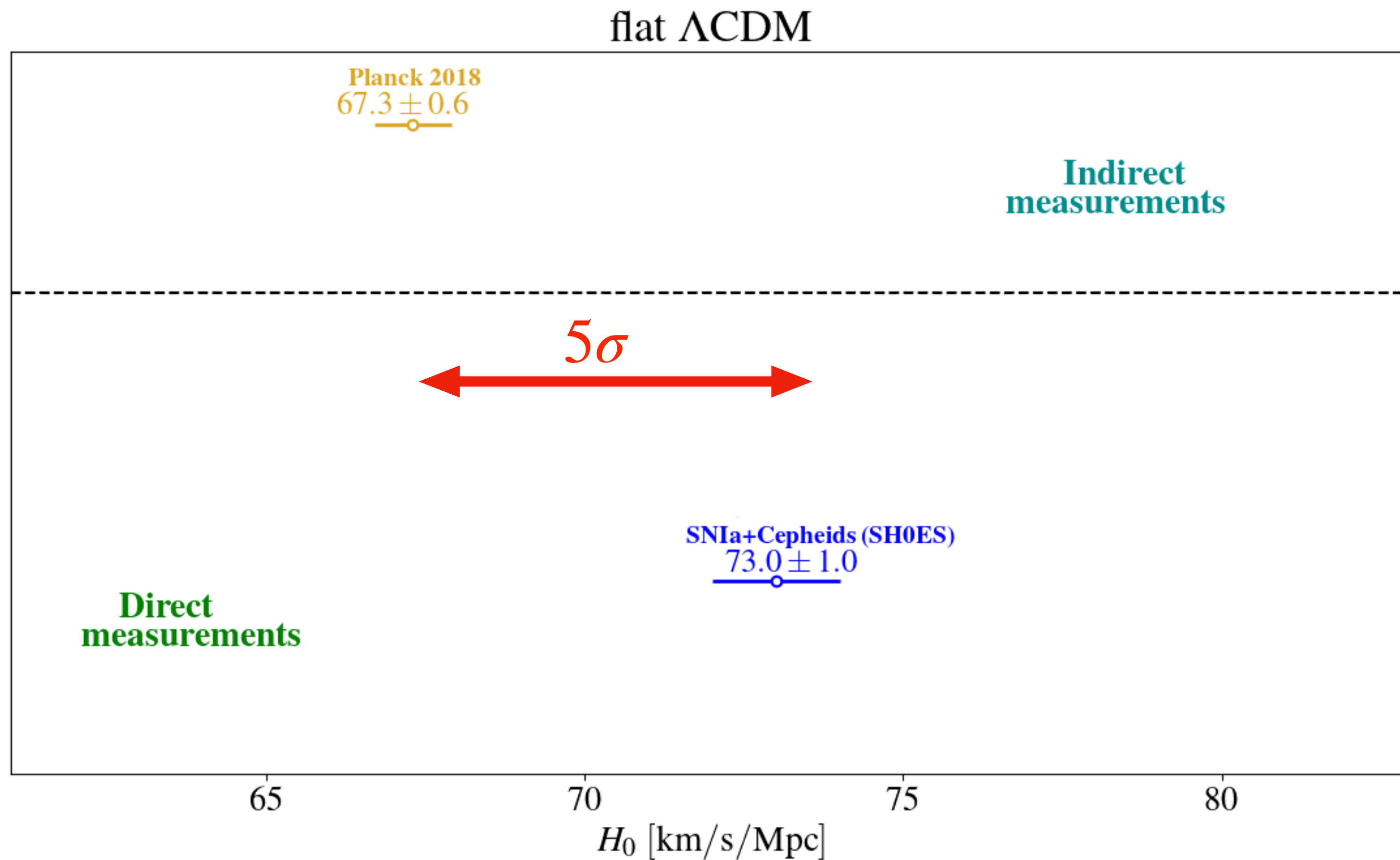
$$S_8 = 0.830 \pm 0.013$$

→ **$\sim 2 - 3\sigma$ tension**

*Other surveys such as DES, CFHTLenS or HSC yield similar results

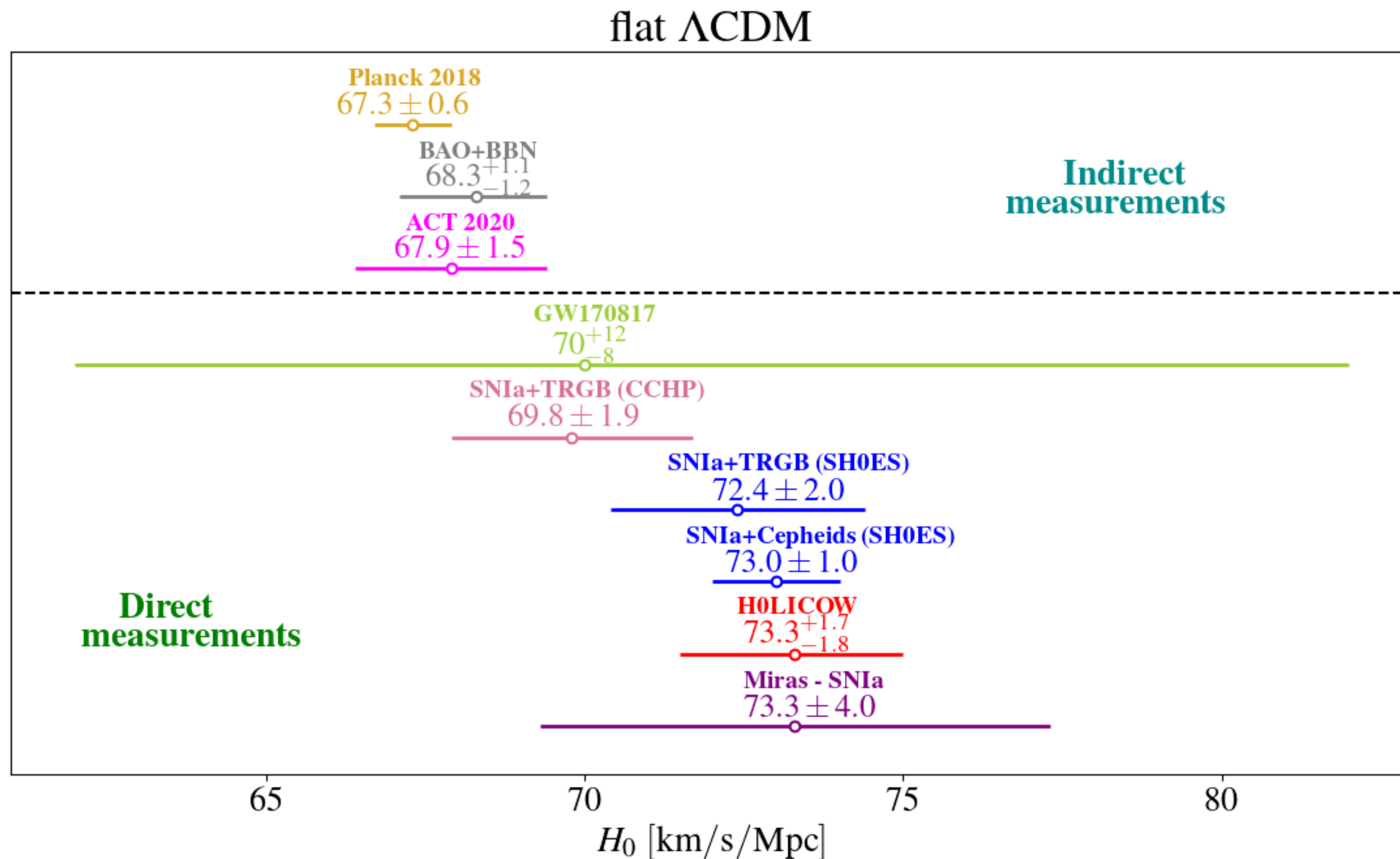
The H_0 tension

Planck (*under Λ CDM*) and SHoES measurements are now in **5σ tension** !

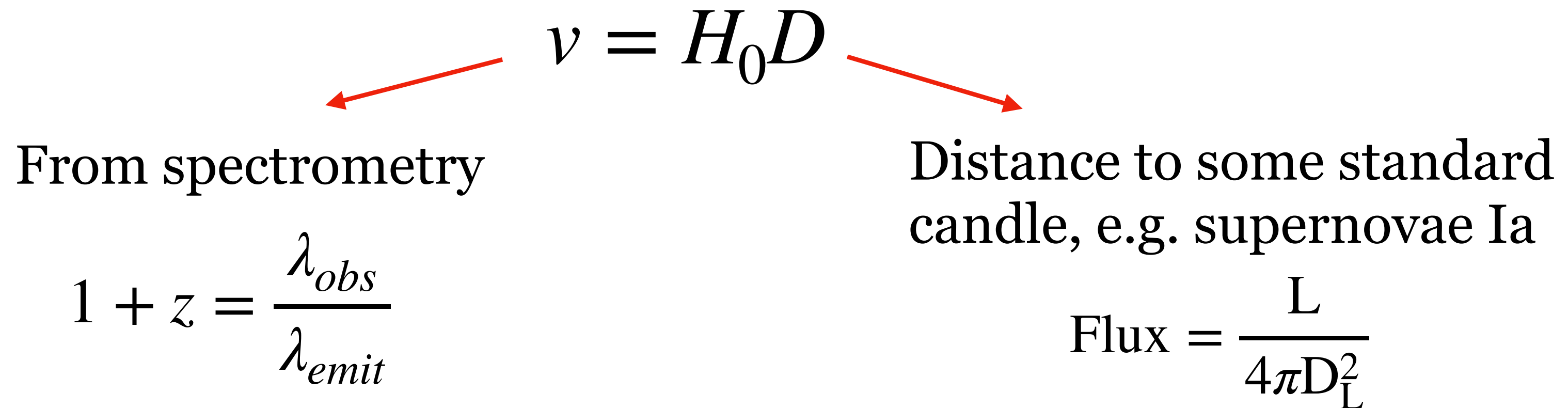


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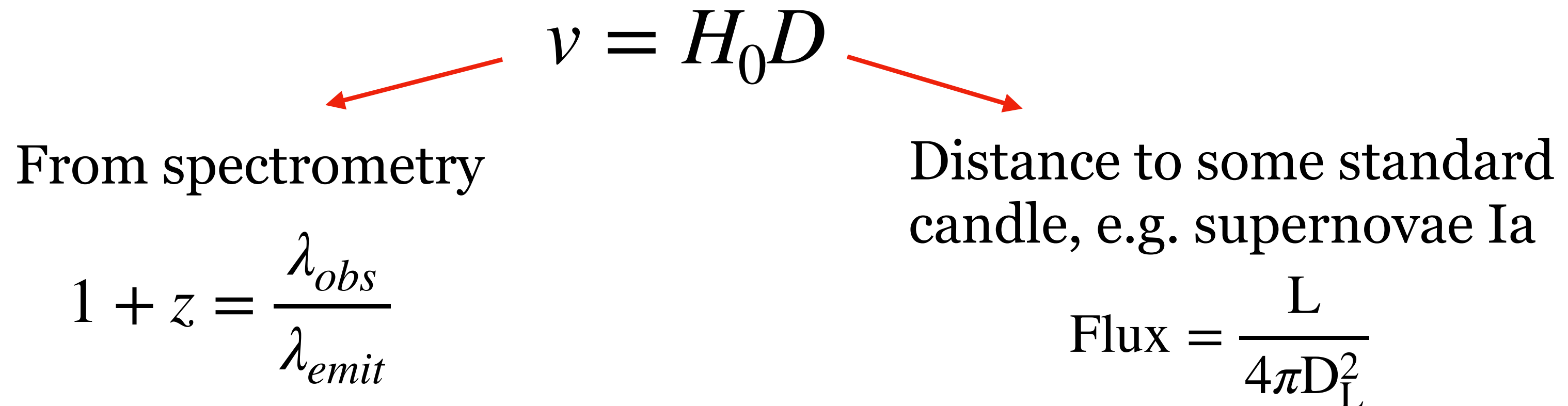
Planck (*under Λ CDM*) and SHoES measurements are now in **5 σ tension** !
High- and low-redshift probes are typically discrepant



How does SH0ES determine H_0 ?



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Focus on small z^* , for which distances are approx. **model-independent**

$$D_L = (1 + z) \int_0^z \frac{cdz'}{H(z')} \xrightarrow{z \ll 1} czH_0^{-1} \simeq vH_0^{-1}$$

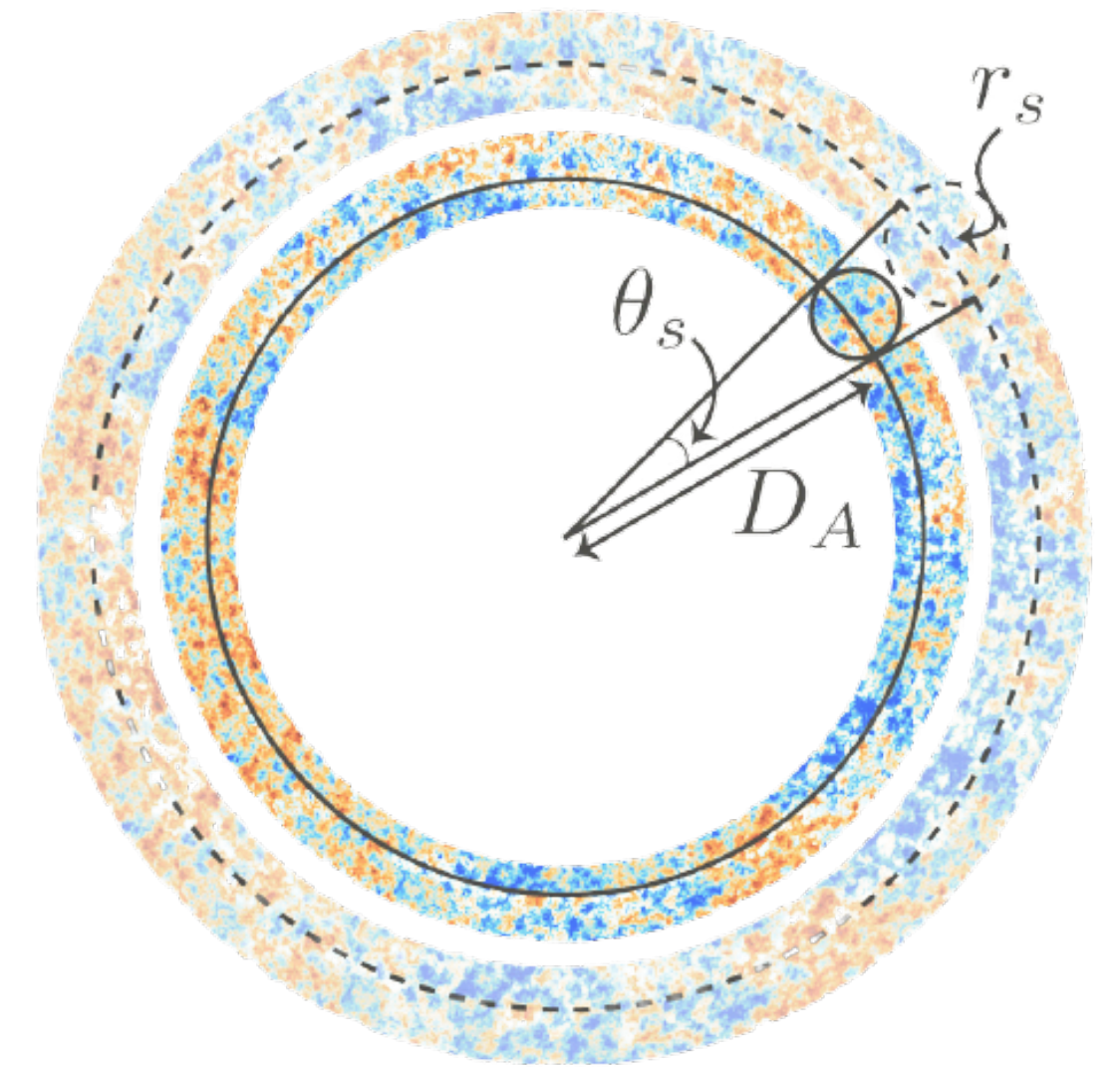
$$\text{where } H^2(z) = \frac{8\pi G}{3} \sum_i \rho_i(z)$$

*But not too small, to make sure peculiar velocities are negligible

How does Planck determine H_0 ?

Angular size of the sound horizon is measured at the 0.04 % precision

$$\theta_s = \frac{r_s(z_{\text{rec}})}{D_A(z_{\text{rec}})} = \frac{\int_0^{\tau_{\text{rec}}} c_s(\tau) d\tau}{\int_{\tau_{\text{rec}}}^{\tau_0} c d\tau}$$



T. Smith

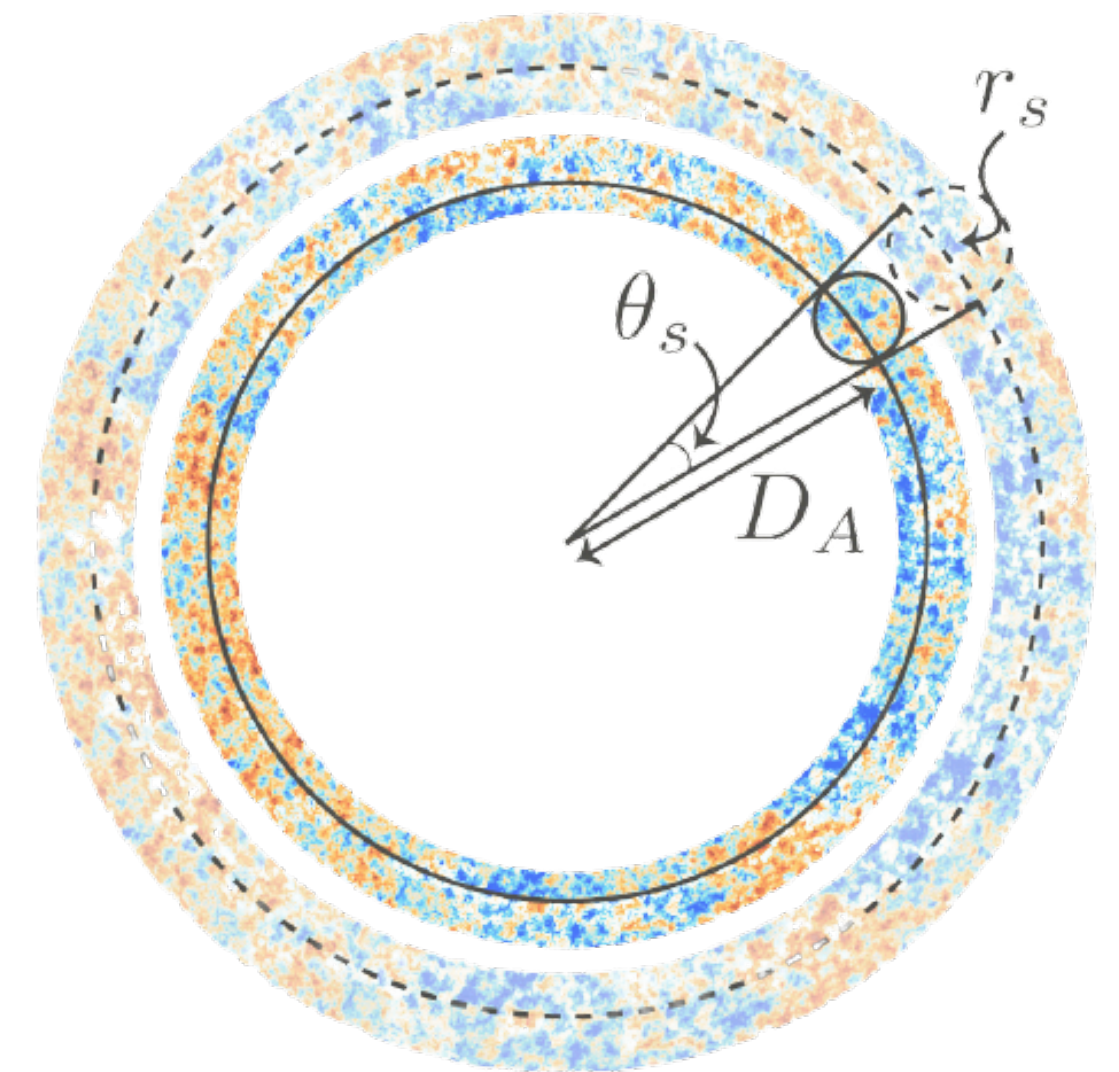
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with $D_A \propto 1/H_0 = 1/\sqrt{\rho_{\text{tot}}(0)}$

model prediction of r_s + measurement of $\theta_s \longrightarrow H_0$



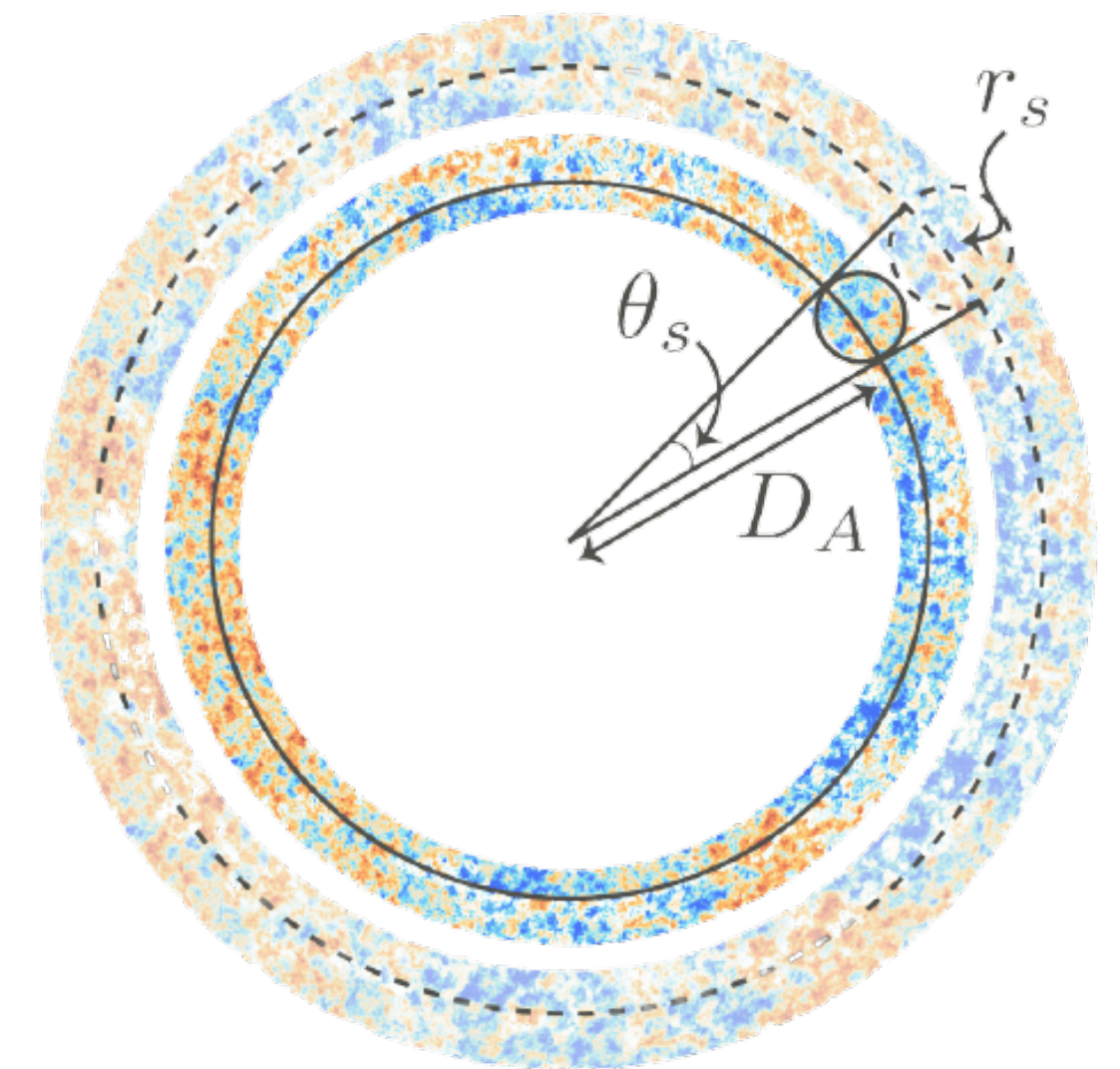
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Early-time solutions

Decrease $r_s(z_{\text{rec}})$ at fixed θ_s to decrease $D_A(z_{\text{rec}})$ and increase H_0

Ex : $\Delta N_{\text{eff}} > 0$

Late-time solutions

$r_s(z_{\text{rec}})$ and $D_A(z_{\text{rec}})$ are fixed, but $D_A(z < z_{\text{rec}})$ is changed to allow higher H_0

Ex : $w < -1$

Lost in the landscape of solutions

- Cosmological tensions have become a **very hot topic** (specially the H_0 tension)

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- [Di Valentino, Mena++ 2103.01183](#) \longrightarrow recent review of solutions, **more than 1000 refs!**

Early Dark Energy Can Resolve The Hubble Tension

Vivian Poulin¹, Tristan L. Smith², Tanvi Karwal¹, and Marc Kamionkowski¹

Relieving the Hubble tension with primordial magnetic fields

Karsten Jedamzik¹ and Levon Pogosian^{2,3}

The Neutrino Puzzle: Anomalies, Interactions, and Cosmological Tensions

Christina D. Kreisch,^{1,*} Francis-Yan Cyr-Racine,^{2,3,†} and Olivier Doré⁴

Rock 'n' Roll Solutions to the Hubble Tension

Prateek Agrawal¹, Francis-Yan Cyr-Racine^{1,2}, David Pinner^{1,3}, and Lisa Randall¹

The Hubble Tension as a Hint of Leptogenesis and Neutrino Mass Generation

Miguel Escudero^{1,*} and Samuel J. Witte^{2,†}

Can interacting dark energy solve the H_0 tension?

Eleonora Di Valentino,^{1,2,*} Alessandro Melchiorri,^{3,†} and Olga Mena^{4,‡}

Dark matter decaying in the late Universe can relieve the H_0 tension

Kyriakos Vattis, Savvas M. Koushiappas, and Abraham Loeb

A Simple Phenomenological Emergent Dark Energy Model can Resolve the Hubble Tension

XIAOLEI LI^{1,2} AND ARMAN SHAFIELOO^{1,3}

Early recombination as a solution to the H_0 tension

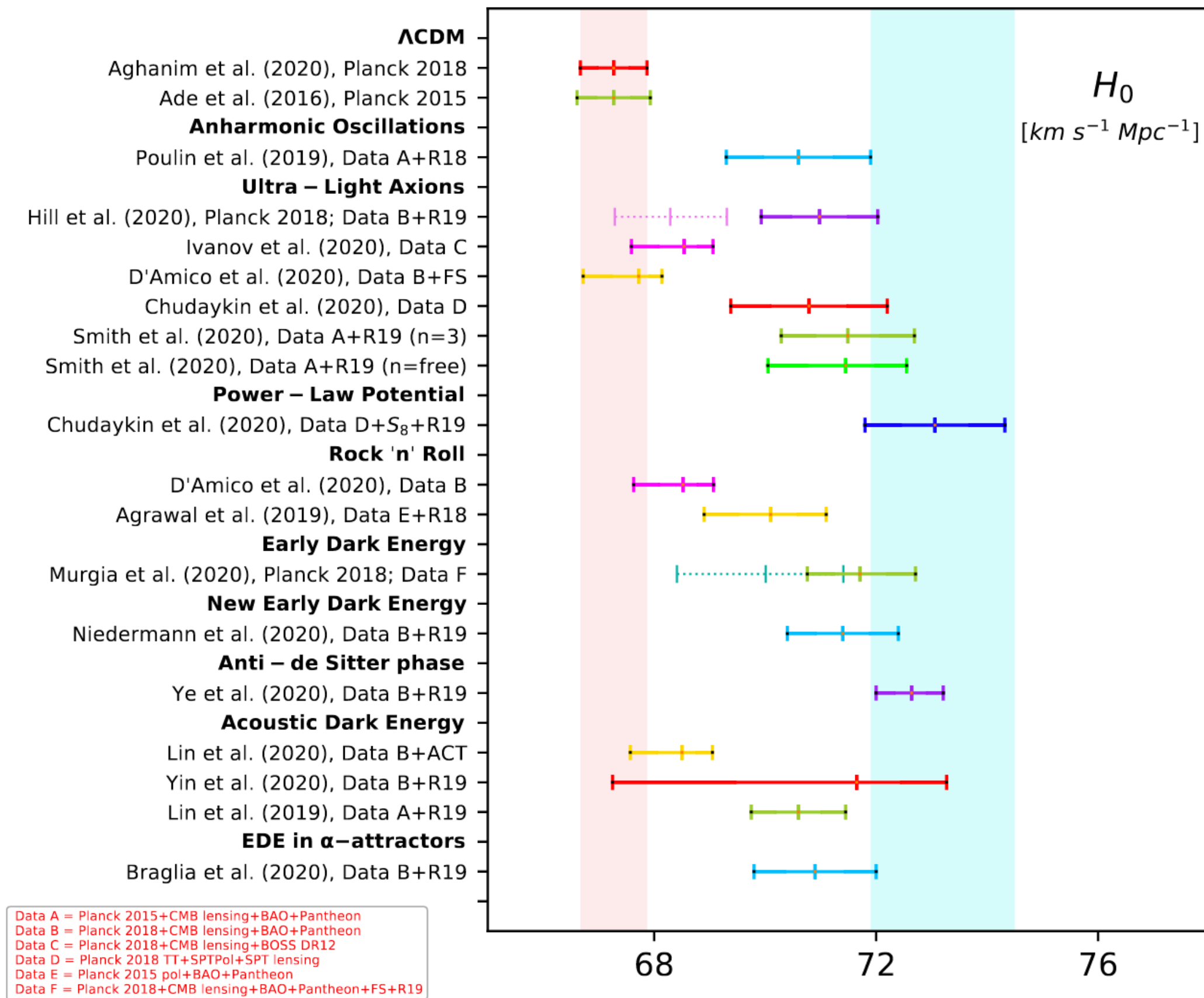
Toyokazu Sekiguchi^{1,*} and Tomo Takahashi^{2,†}

Early modified gravity in light of the H_0 tension and LSS data

Matteo Braglia,^{1,2,3,*} Mario Ballardini,^{1,2,3,†} Fabio Finelli,^{2,3,‡} and Kazuya Koyama^{4,§}

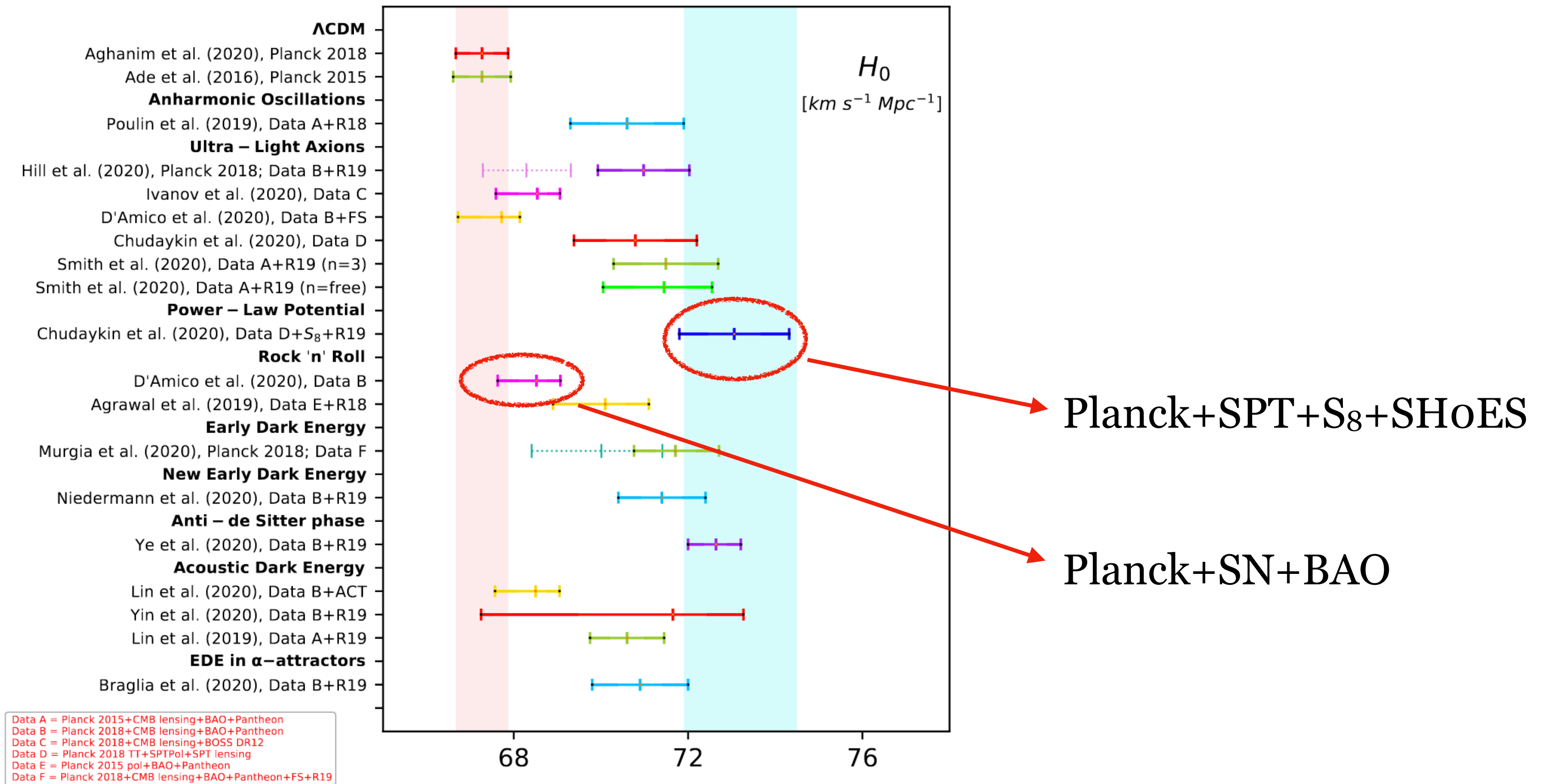
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It proves difficult to compare success of the different proposed solutions, since authors typically use **differing and incomplete combinations of data**



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Goal: Take a representative **sample of proposed solutions**, and quantify the **relative success** of each using certain metrics and a wide array of data

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Early universe

with Dark radiation

- Free-streaming DR (ΔN_{eff})
- Self-interacting DR (ΔN_{fluid})
- Mixed DR ($\Delta N_{\text{eff}} + \Delta N_{\text{fluid}}$)
- DM-DR interactions
- Self-interacting ν_s
- Majoron- ν_s interactions

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no Dark radiation

- Primordial B fields
- Varying m_e
- Varying $m_e + \Omega_k$
- Early Dark Energy (EDE)
- New Early Dark Energy (NEDE)
- Early Modified Gravity (EMG)

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Late universe

- CPL dark energy
- Phenomenological Emergent Dark Energy (PEDE)
- Modified PEDE
- Fraction DM \rightarrow DR
- DM \rightarrow DR + WDM

Model-independent treatment of the SH0ES data

The cosmic distance ladder method *doesn't directly measure H_0* .

It directly measures the intrinsic magnitude of SNIa M_b at redshifts $0.02 \leq z \leq 0.15$, and then infers H_0 by comparing with the *apparent SNIa magnitudes m*

$$m(z) = M_b + 25 - 5 \text{Log}_{10} H_0 + 5 \text{Log}_{10} (\hat{D}_L(z))$$

where

$$\hat{D}_L(z) \simeq z \left(1 + (1 - q_0) \frac{z}{2} - \frac{1}{6} (1 - q_0 - 3q_0^2 + j_0) z^2 \right)$$

Depends on the model!

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Quantifying model success

Criterion 1: Can we get high values of H_0 (or M_b) from a data combination D not including a SHoES prior?

Gaussian tension GT

$$\frac{\bar{x}_D - \bar{x}_{SHoES}}{\sqrt{\sigma_D^2 + \sigma_{SHoES}^2}} \text{ for } x = M_b$$

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Caveats:

- Only valid for gaussian posteriors ✘
- Doesn't quantify quality of the fit ✘

Quantifying model success

Criterion 2: Can we get a good fit to all the data in a given model?

Q_{DMAP} tension

$$\sqrt{\chi_{\min, D+\text{SH0ES}}^2 - \chi_{\min, D}^2}$$

[Raveri&Hu 1806.04649](#)

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[Raveri&Hu 1806.04649](#)

We demand $Q_{\text{DMAP}} < 3\sigma$

Caveats:

- Accounts for non-gaussianity of posteriors 
- Doesn't account for effects of over-fitting 

Quantifying model success

Criterion 3: Is a model M favoured over Λ CDM?

Akaike Information Criterion Δ AIC

$$\chi_{\min, M}^2 - \chi_{\min, \Lambda\text{CDM}}^2 + 2(N_M - N_{\Lambda\text{CDM}})$$

We demand $\Delta\text{AIC} < -6.91$ *

*Corresponds to weak preference according to Jeffrey's scale

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Caveats:

- Simple to use and prior-independent 

*Corresponds to weak preference according to Jeffrey's scale

Steps of the contest

1

Compare **all models** against

- Planck 2018 TTTEEE+lensing
- BAO (BOSS DR12+MGS+6dFGS)
- Pantheon SNIa catalog
- SHoES

Steps of the contest

②

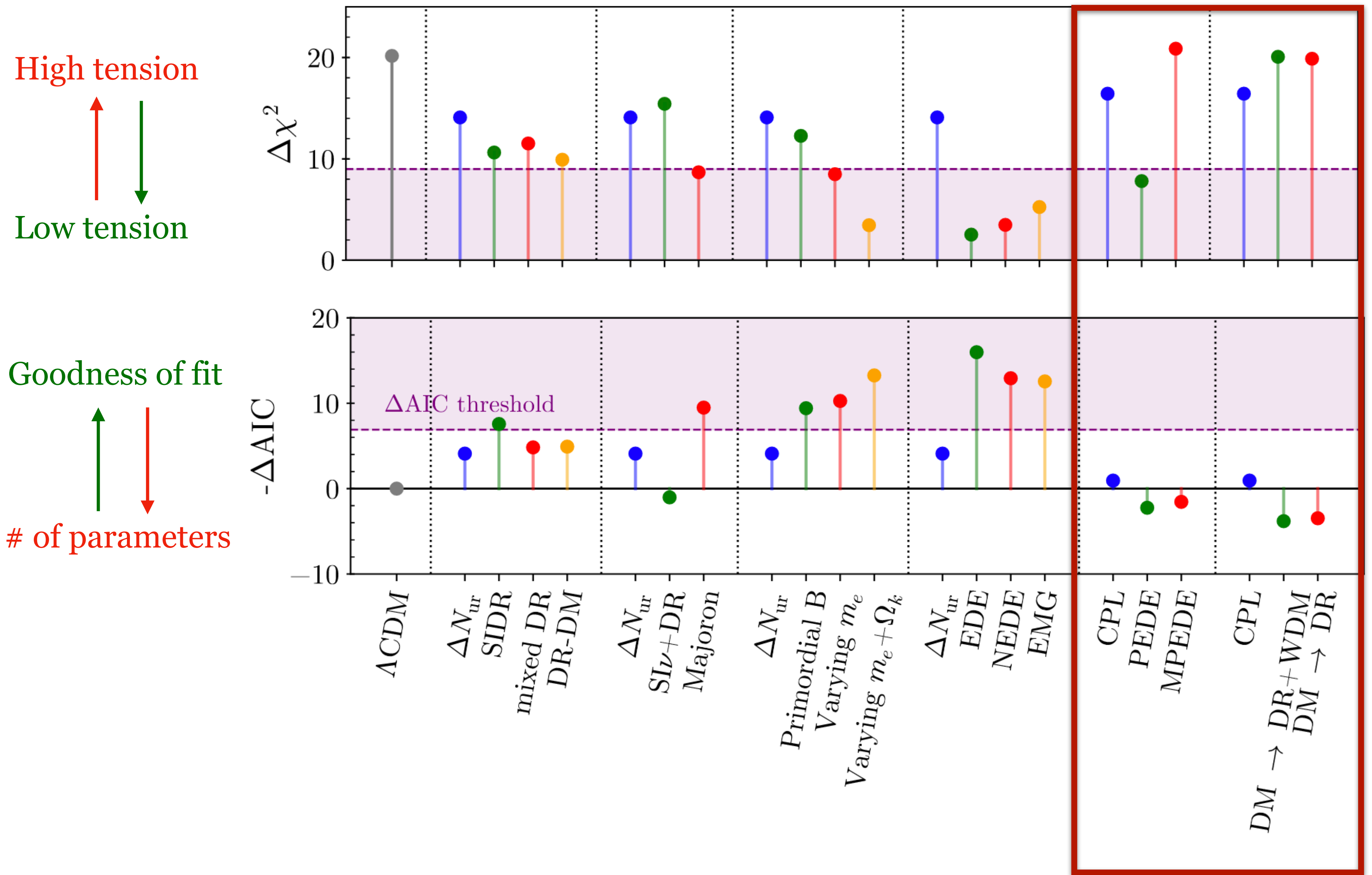
As long as $\Delta AIC < 0$, models go into **finalist** if criterium 2 or 3 are satisfied

Steps of the contest

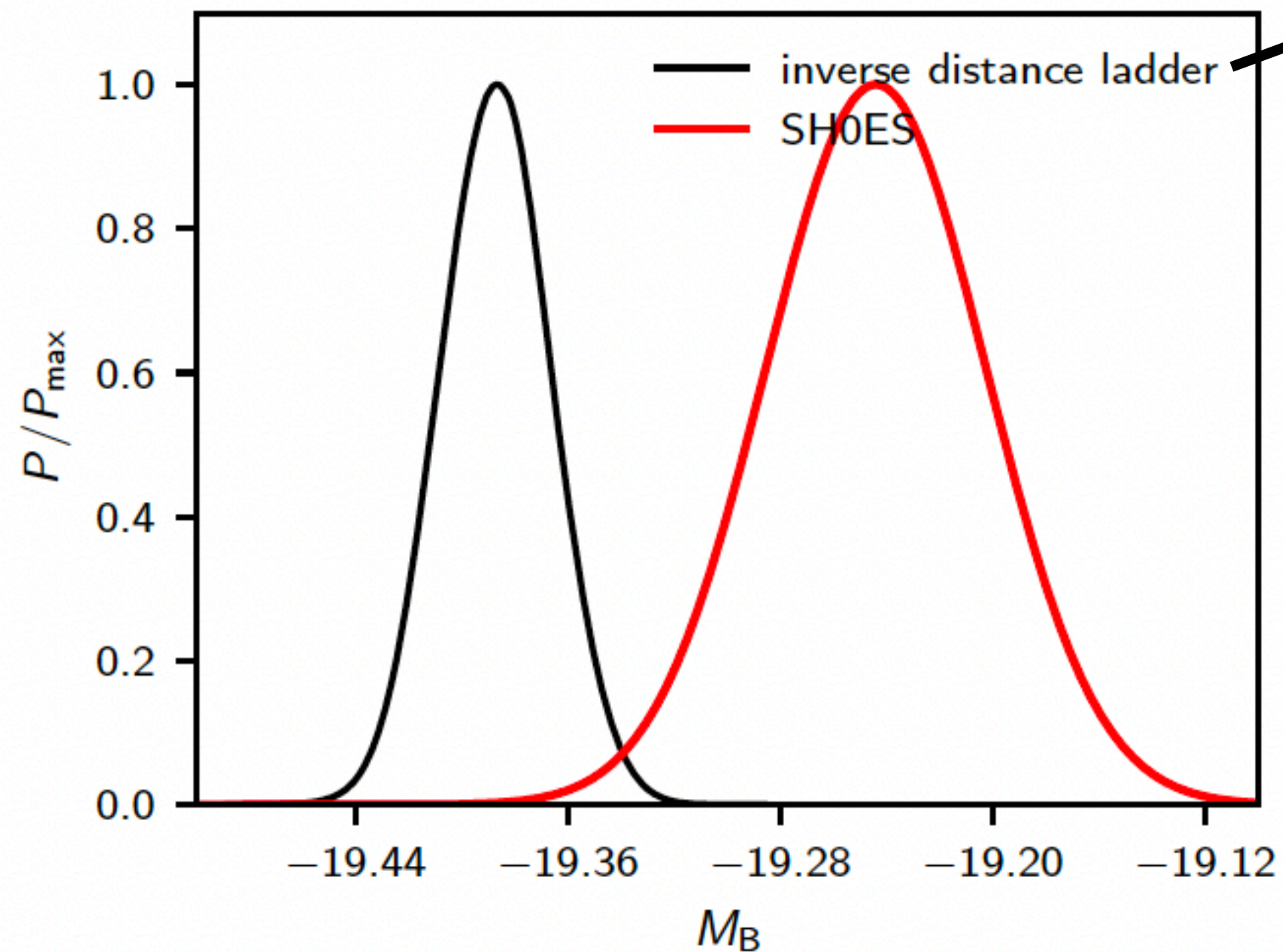
③

Finalists receive **bronze**, **silver** or **golden** medals if they satisfy **one**, **two** or **three** criteria, respectively

Results: late-time solutions



Late-time solutions are disfavoured by BAO+SN Ia



Efstathiou 2103.08723

Given r_s , obtain D_A using BAO data

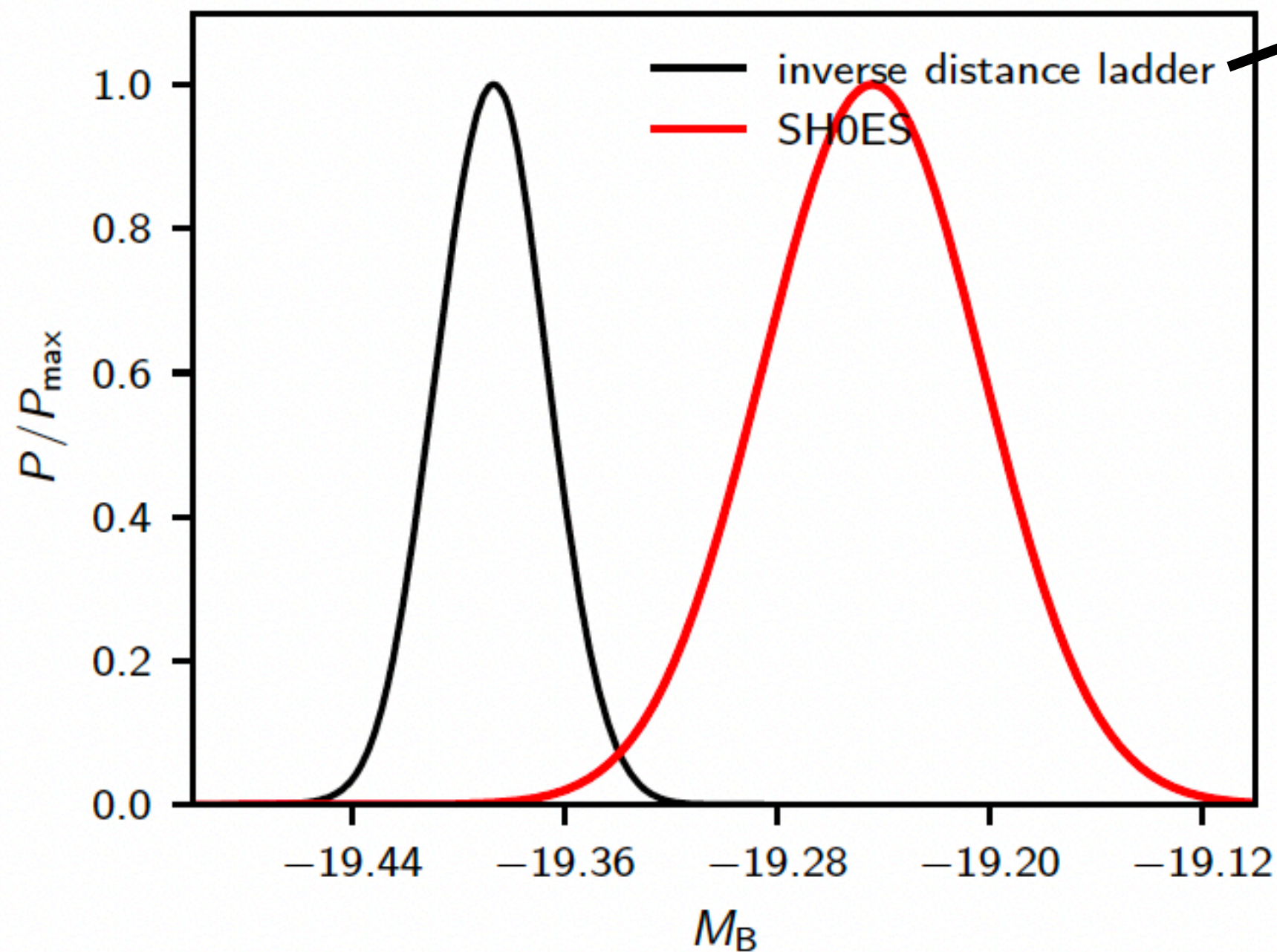
$$\theta_d(z)^\perp = \frac{r_s(z_{\text{drag}})}{D_A(z)}, \quad \theta_d(z)^\parallel = r_s(z_{\text{drag}})H(z)$$

$$D_L(z) = D_A(z)(1+z)^2$$

Obtain M_b from calibration const. of SNIa

$$m(z) = 5\text{Log}_{10}D_L(z) + \text{const}$$

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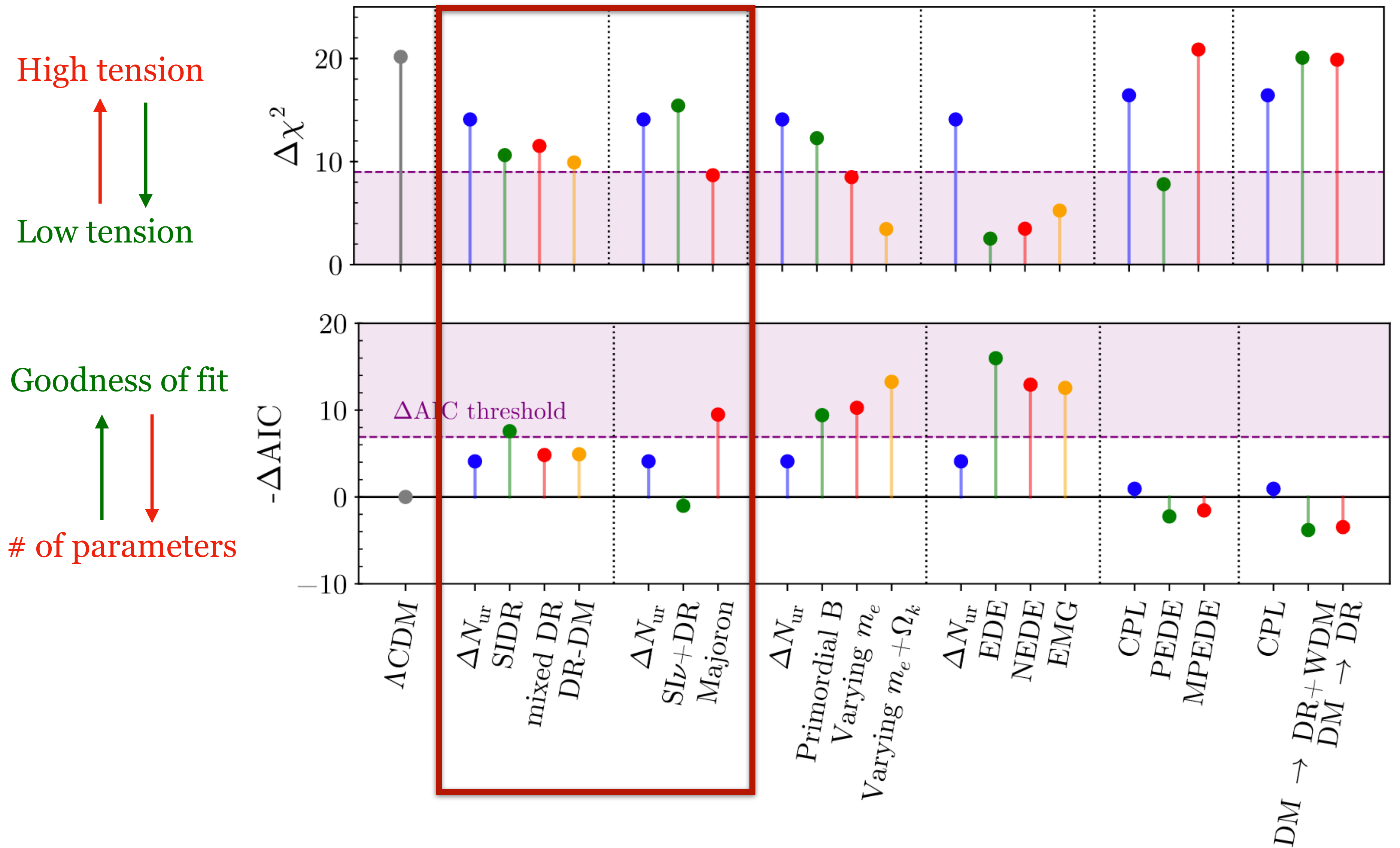
Efstathiou 2103.08723

For $r_s^{\Lambda\text{CDM}} = 147$ Mpc, **inverse distance ladder** disagrees with SH0ES

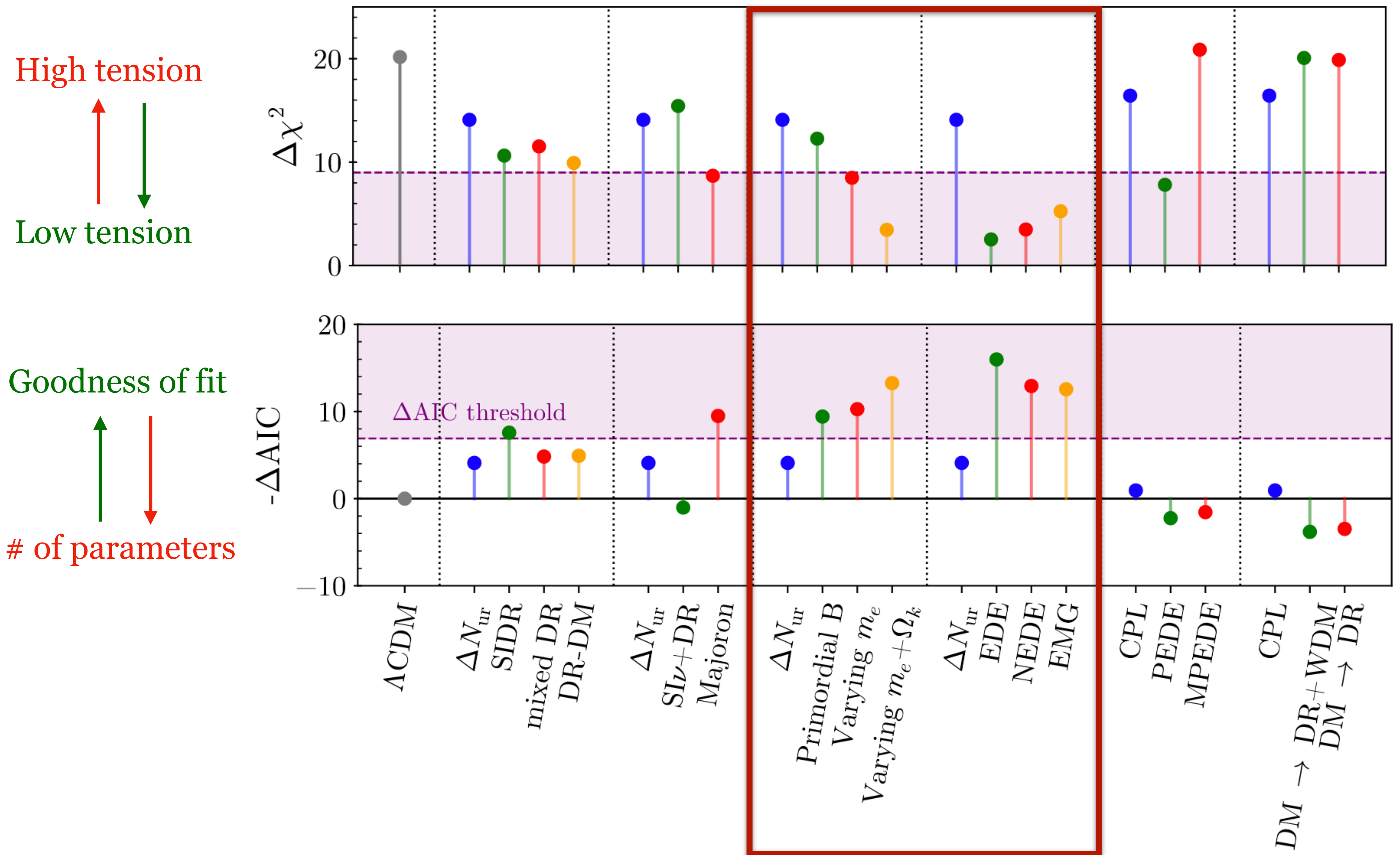
To make the two determinations agree, one is forced to **reduce** r_s

Ex: *Early Dark Energy or varying electron mass*

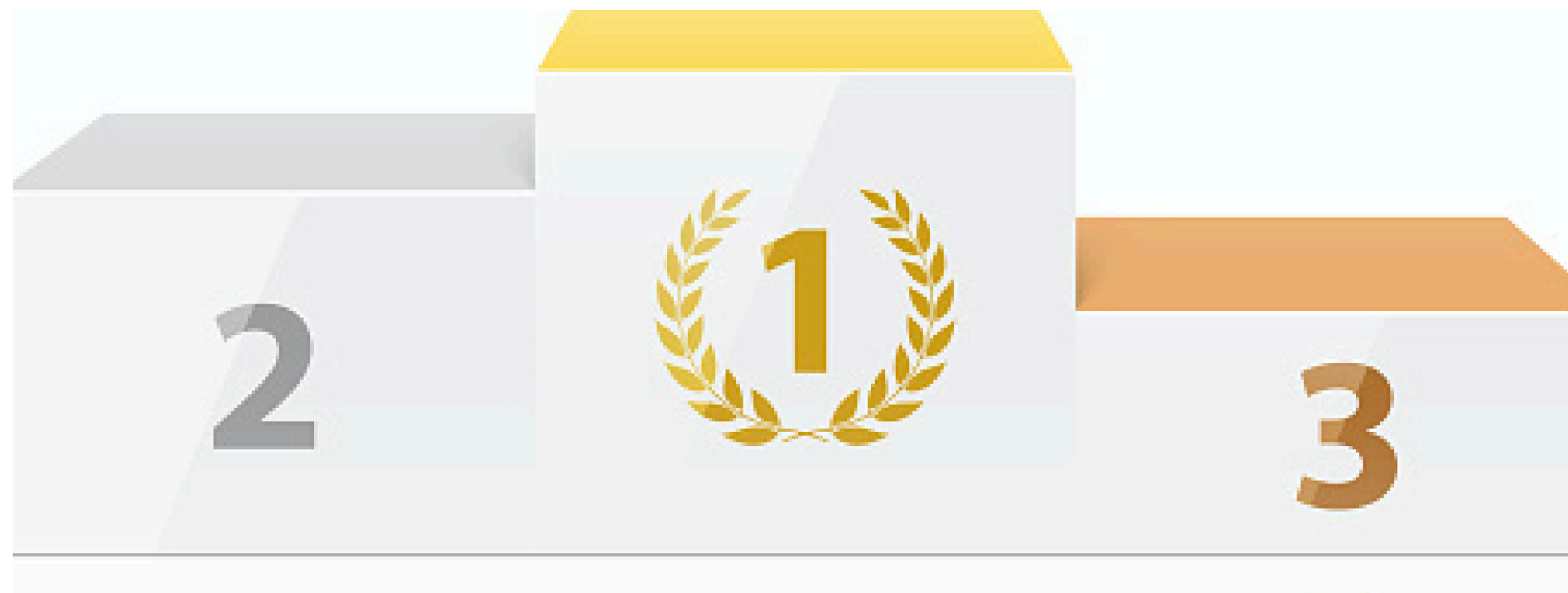
Results: early-time solutions **with** Dark Radiation



Results: early-time solutions **without** Dark Radiation



Results of the contest



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Results of the contest

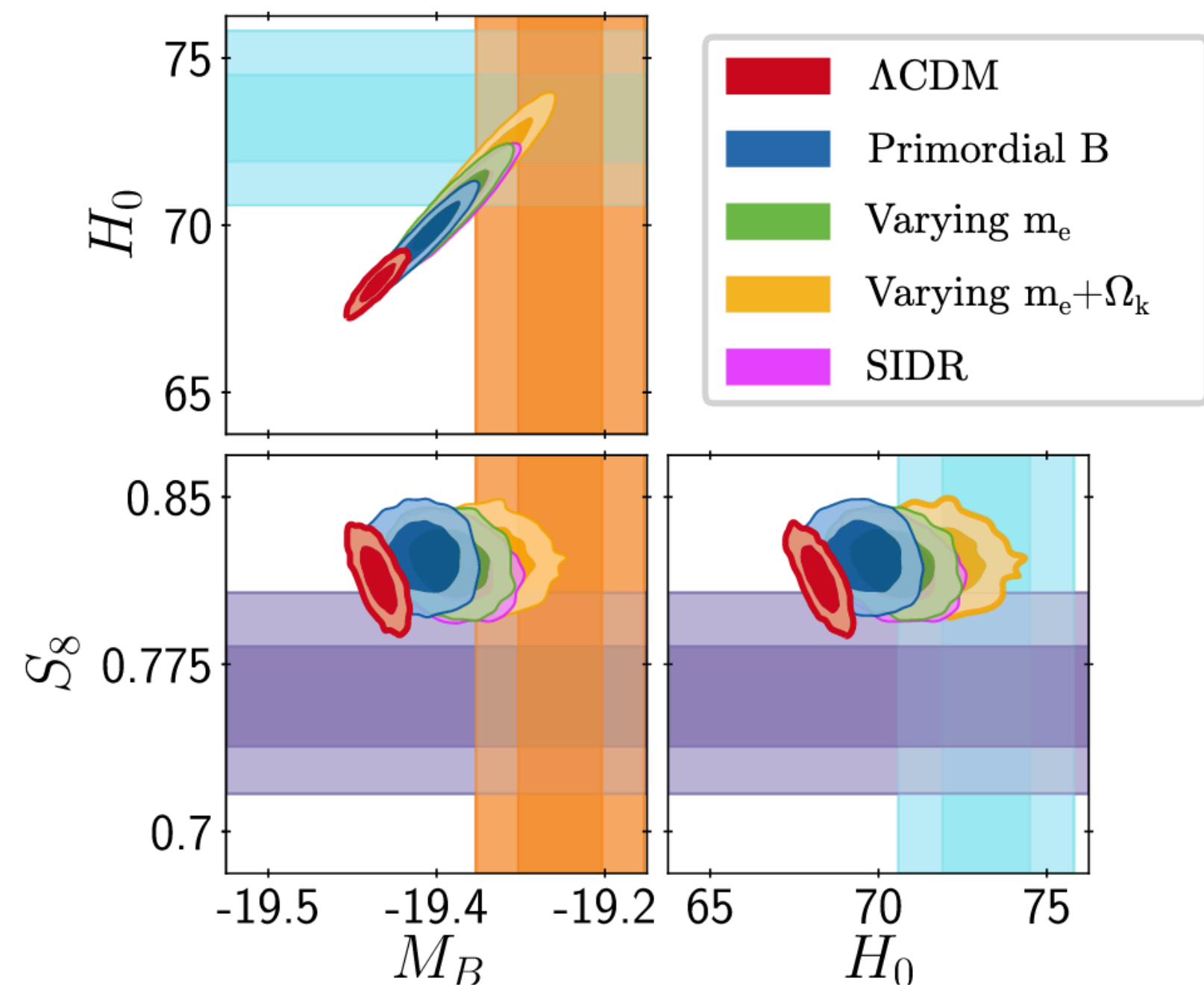
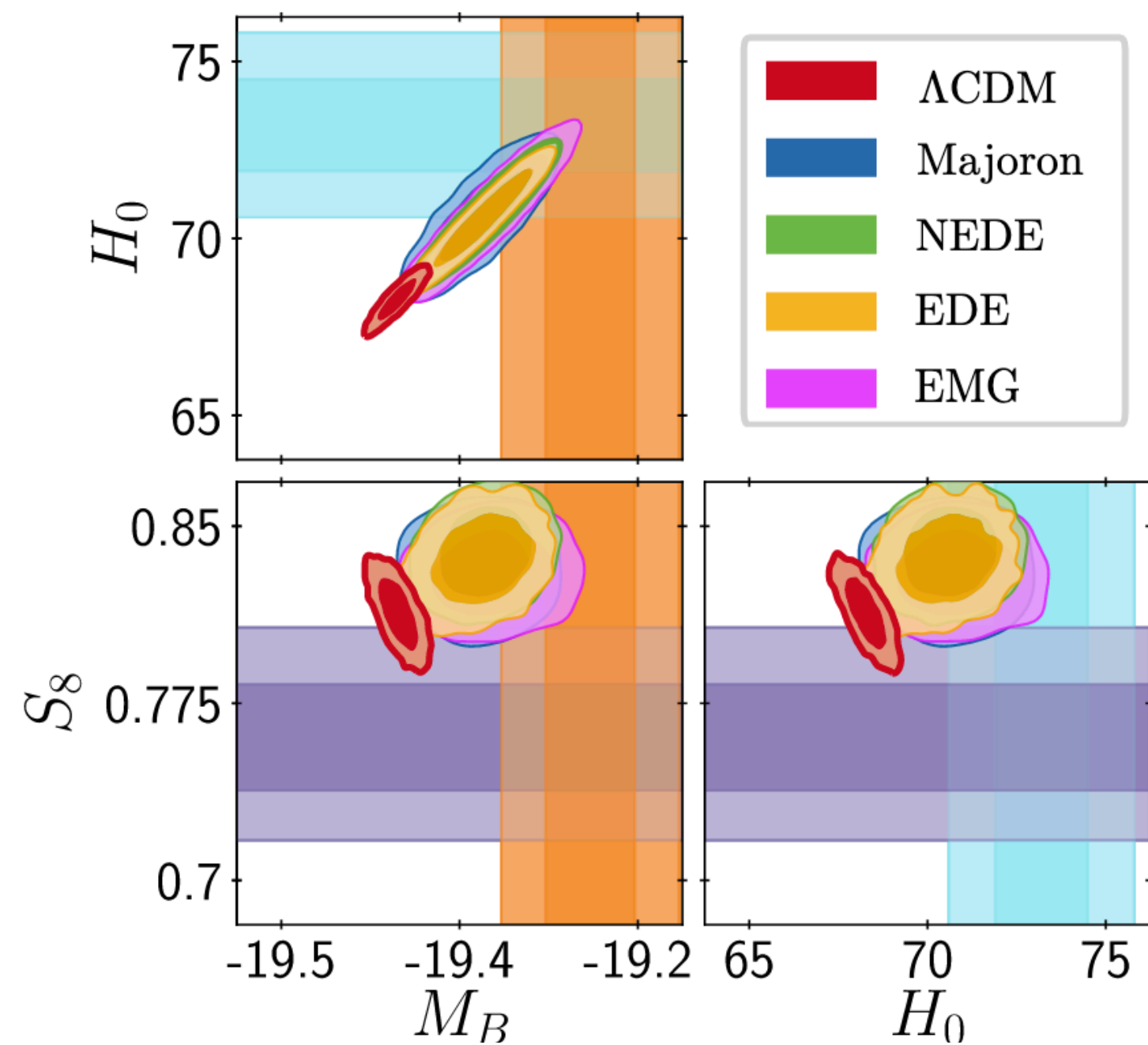


Results of the contest



Results of the contest

Unfortunately, the most successful models face strong **fine-tuning** problems, and are unable to explain the **S_8 tension**



Conclusions

- Λ CDM currently shows a 5σ H_0 tension and a $2-3\sigma$ S_8 tension, which could offer an interesting window to the yet unknown dark sector.

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- None of these successful models is able to relieve the S_8 tension. However, resolutions of these tensions might lie in different sectors ($H_0 \longleftrightarrow$ new background contribution, $S_8 \longleftrightarrow$ new perturbation properties).

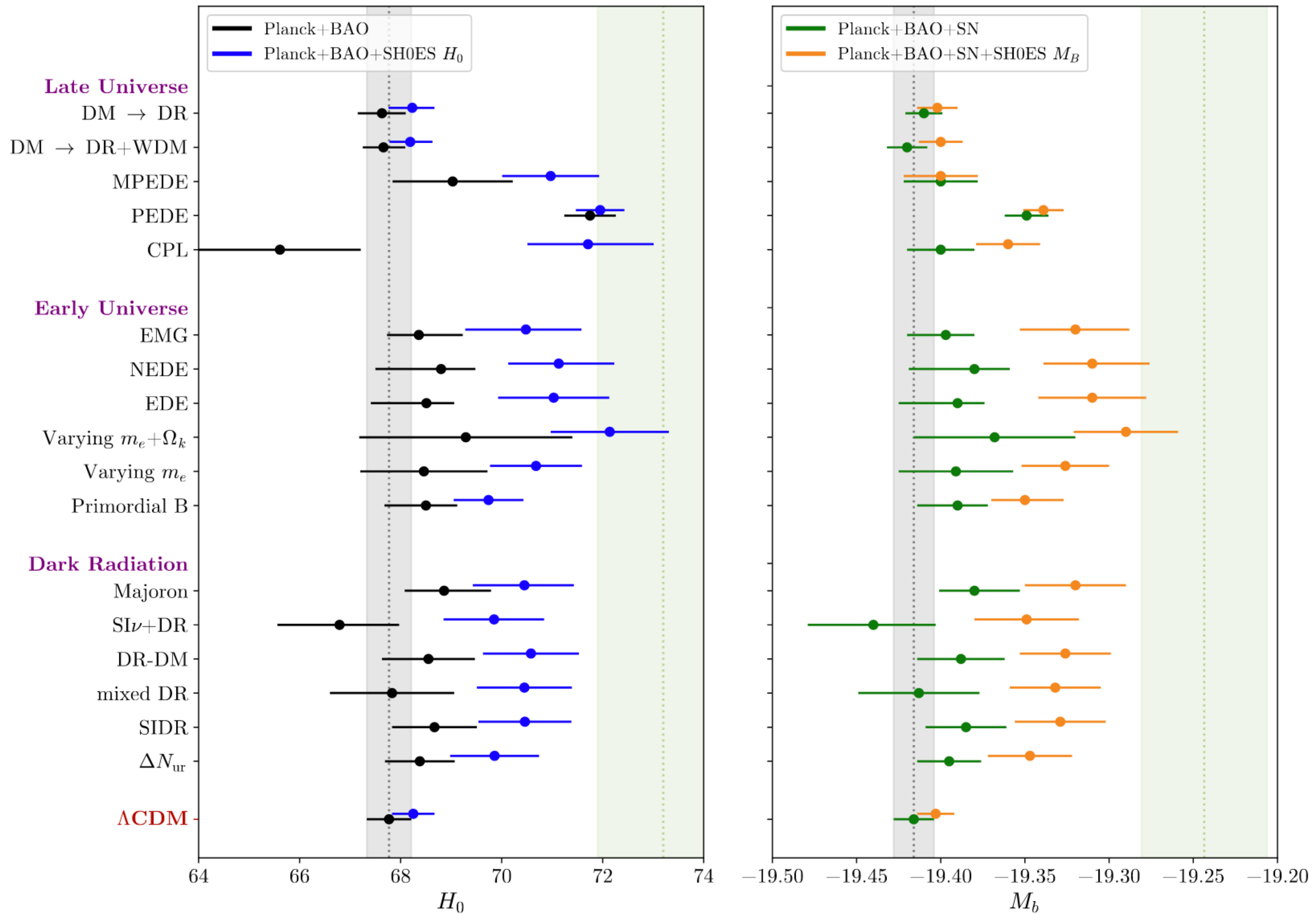
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We might be on the verge of the discovery of a rich dark sector!

BACK-UP SLIDES

Reconstructed values of H_0



H₀ Olympics: testing against other datasets

Role of Planck data: We replaced Planck by WMAP+ACT and BBN+BAO

→ No significant changes (*notable exceptions are EDE and NEDE*)

Adding extra datasets: We included data from Cosmic Chronometers, Redshift-Space-Distortions and BAO Ly- α .

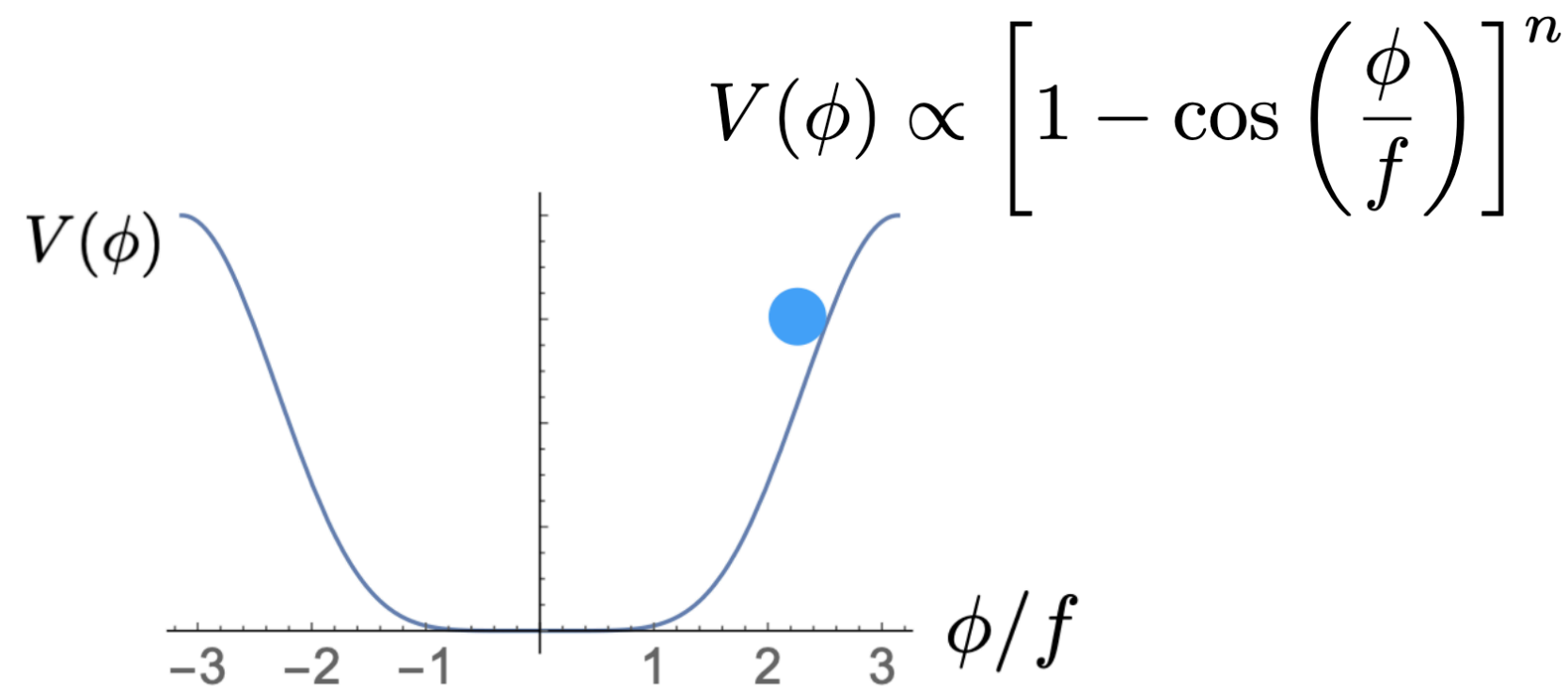
→ No huge impact, but decreases performance of finalist models

Early Dark Energy

Scalar field initially frozen, then dilutes away equal or faster than radiation

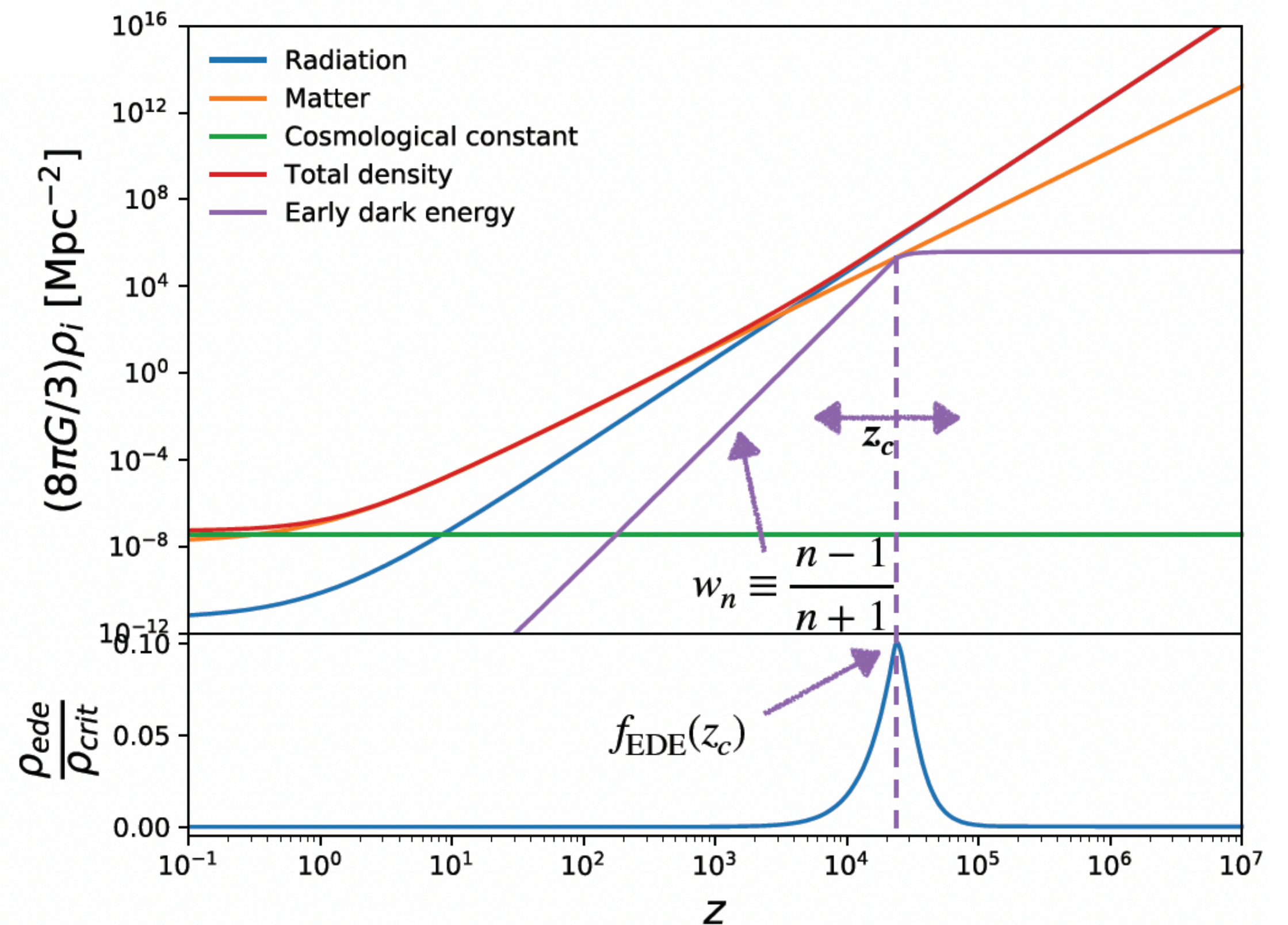
$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$$

+ perturbed linear eqs.



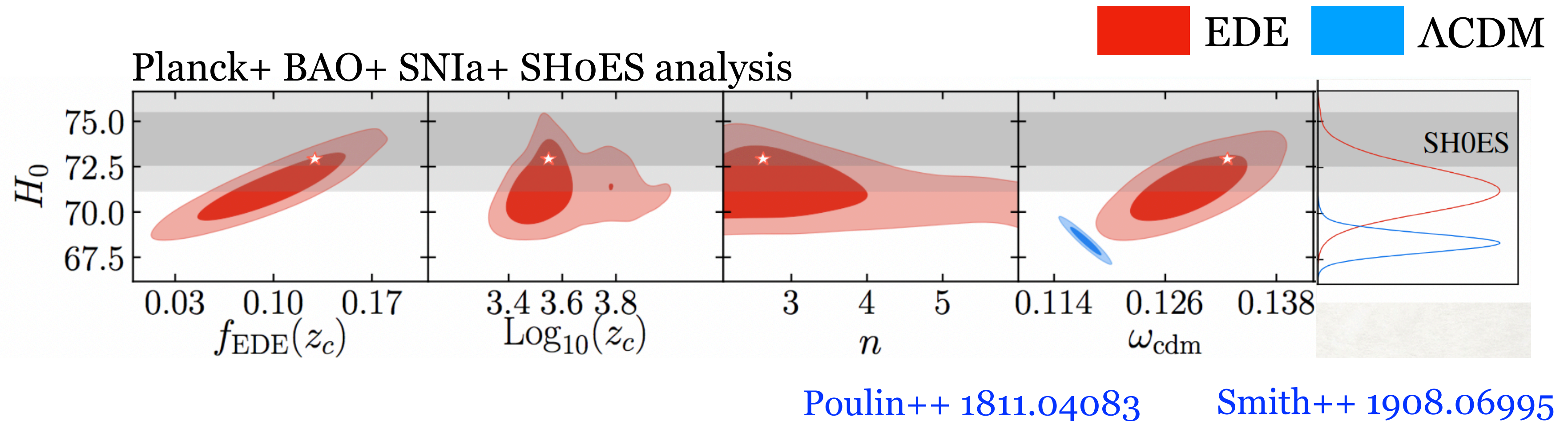
The model is fully specified by

$$\{f_{\text{EDE}}(z_c), z_c, n, \phi_i\}$$



Early Dark Energy

Early Dark Energy **can resolve the H_0 tension** if $f_{\text{EDE}}(z_c) \sim 10\%$ for $z_c \sim z_{\text{eq}}$



Some caveats

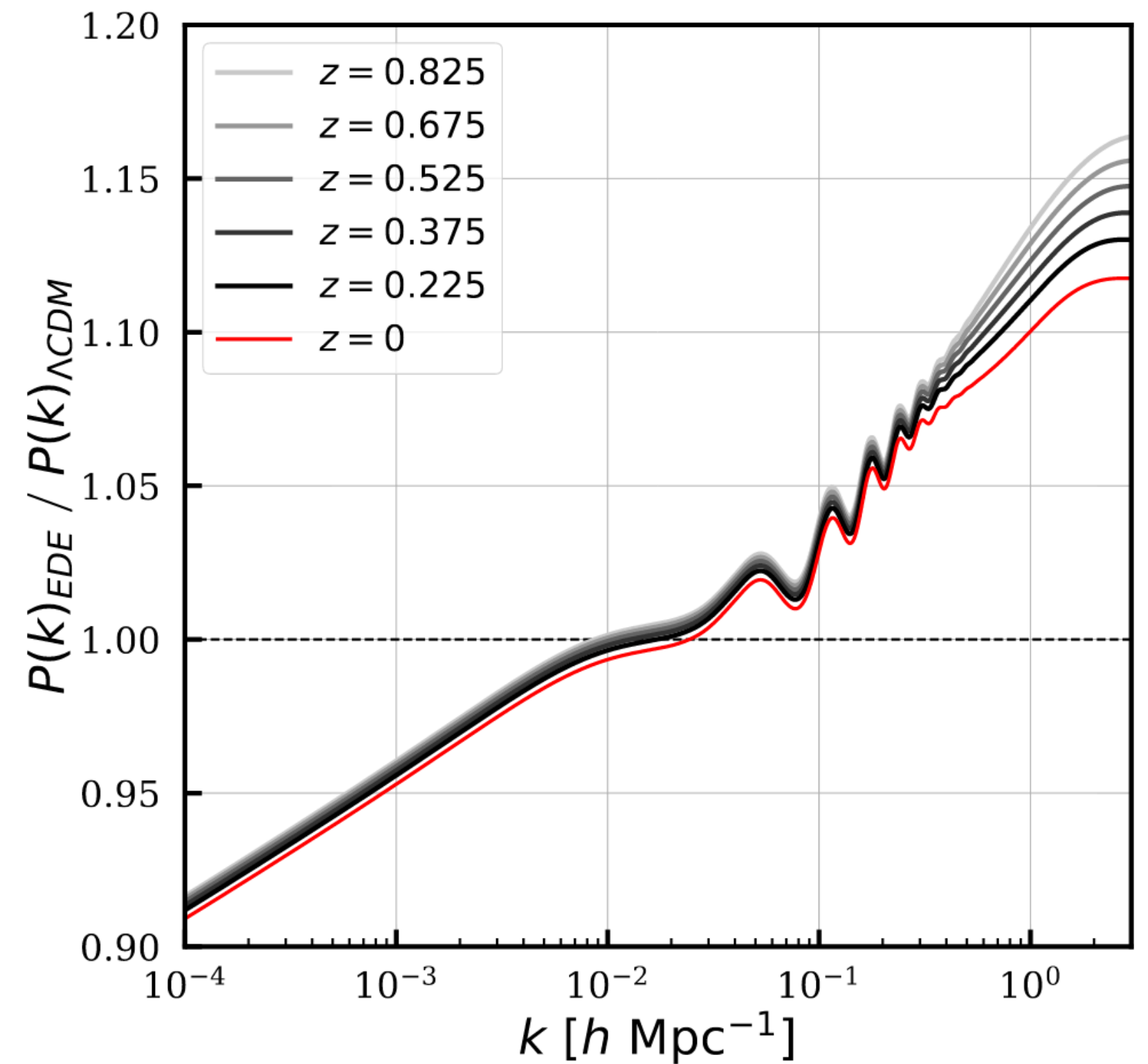
1. *Very fine tuned?*

→ Proposed connexions of EDE with neutrino sector and present DE
 Sakstein++ 1911.11760 Freese++ 2102.13655

2. Increased value of $\omega_{\text{cdm}} = \Omega_{\text{cdm}} h^2$, *exacerbates S_8 tension*
 Jedamzik++ 2010.04158.

Is EDE solution ruled out?

EDE solution **increases power at small k** (with a corresponding increase in S_8), rising mild tension with Large Scale Structure (LSS) data



Hill++ 2003.07355

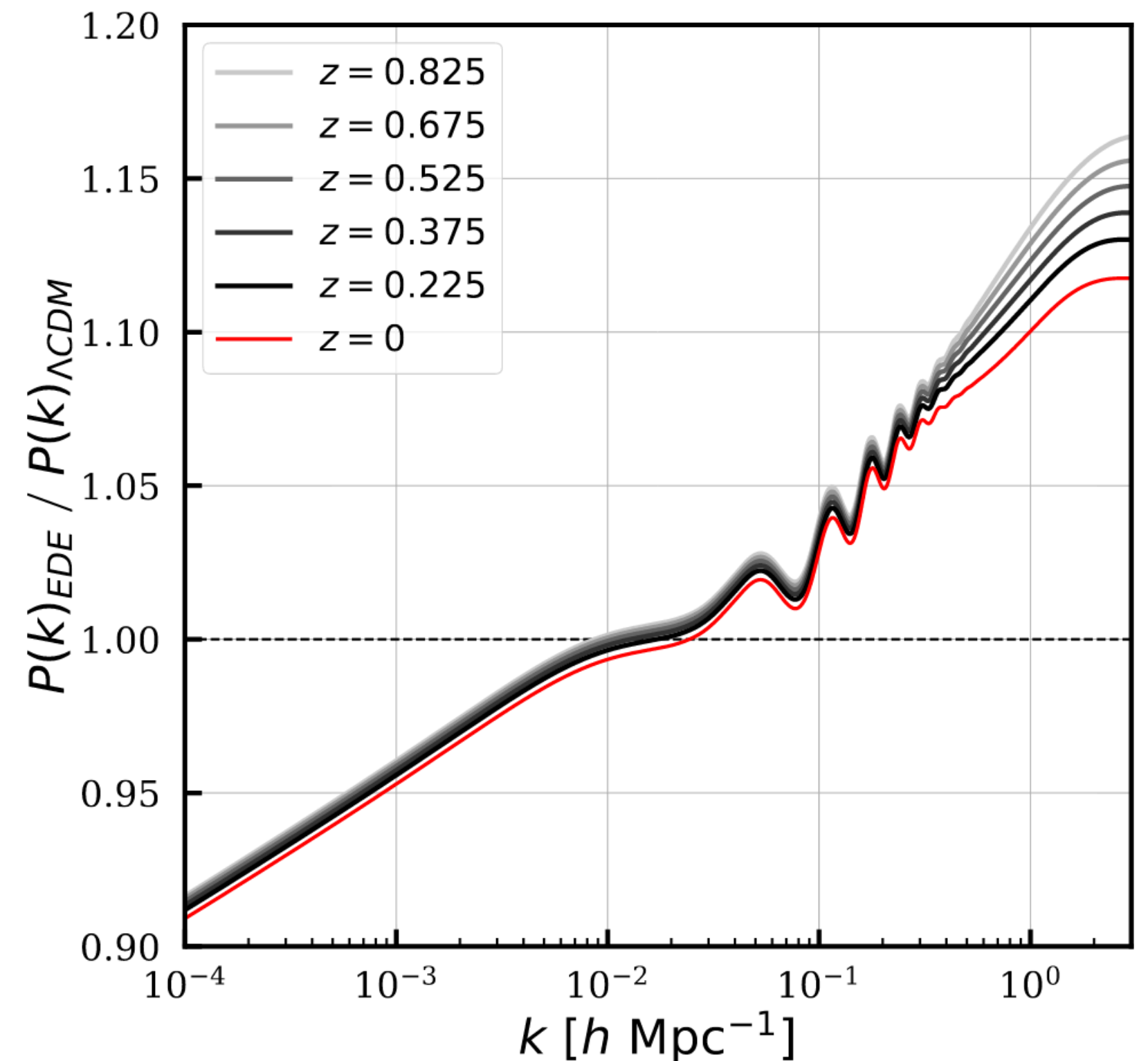
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When **LSS data** is added to analysis, EDE **detection is reduced** from 3σ to 2σ

In addition, EDE is **not detected from Planck data alone**

D'amico++ 2006.12420
Ivanov++ 2006.11235



Hill++ 2003.07355

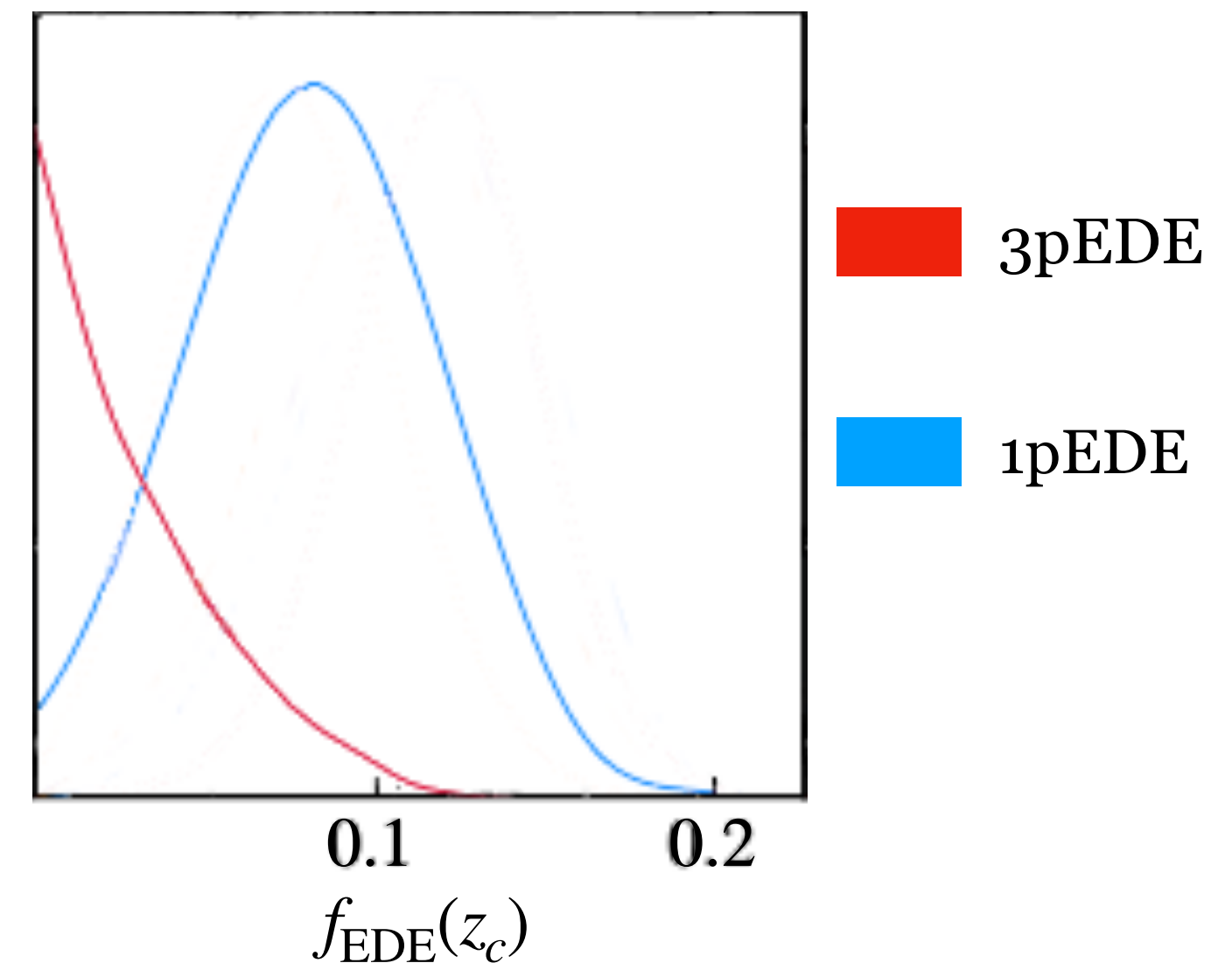
Answer: no, EDE solution is still robust

1. Why EDE is not detected from Planck alone?

χ^2 degeneracy in Planck between Λ CDM and EDE :

For $f_{\text{EDE}} \lesssim 4\%$, parameters z_c and ϕ_i become irrelevant, so posteriors are naturally weighted towards Λ CDM

Planck 2018



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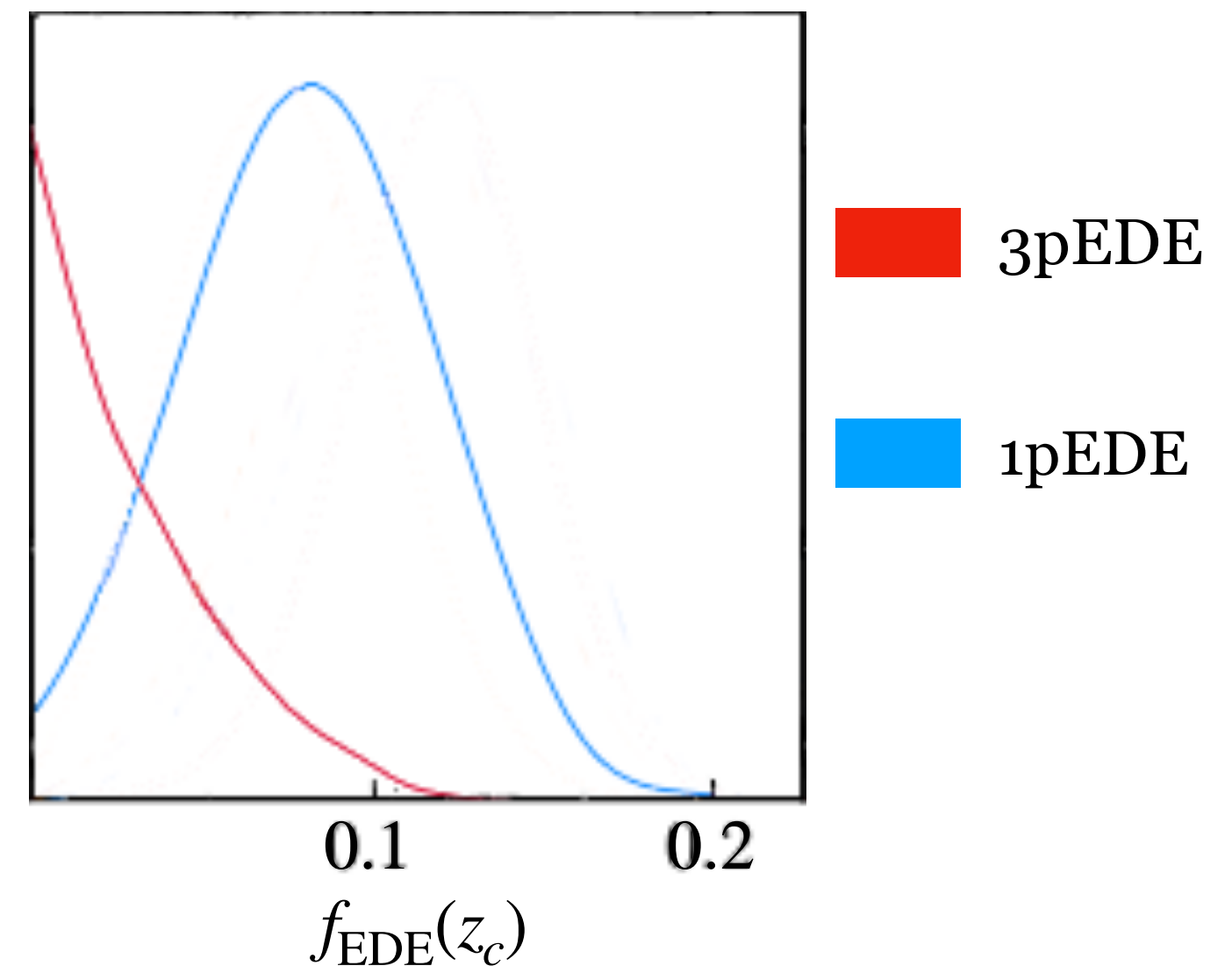
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For $f_{\text{EDE}} \lesssim 4\%$, parameters z_c and ϕ_i become irrelevant, so posteriors are naturally weighted towards Λ CDM

To avoid this Bayesian volume effect, consider a **1 parameter EDE model (1pEDE)**:

Fix z_c and ϕ_i and let f_{EDE} free to vary

Planck 2018



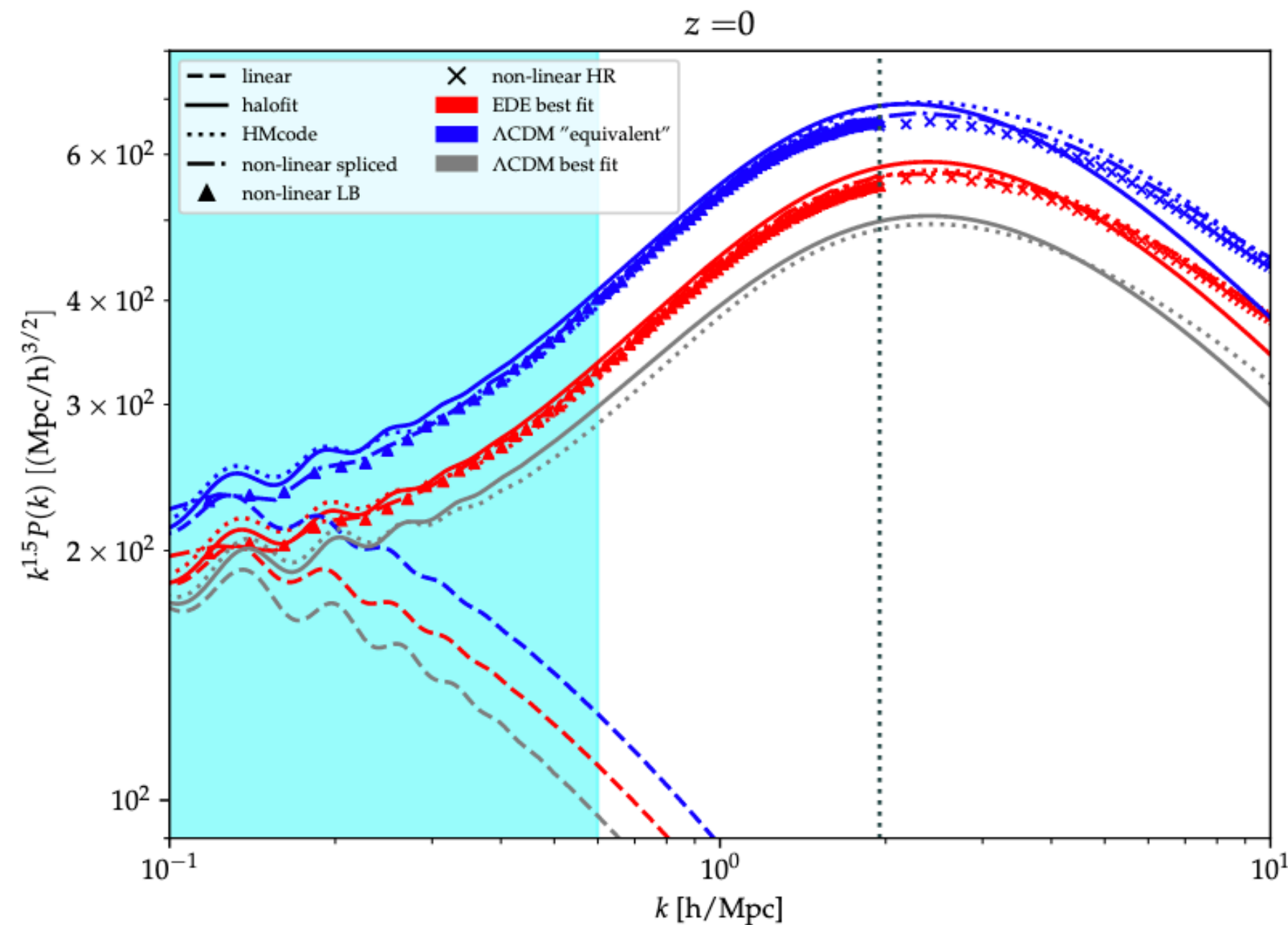
Within 1pEDE, we get a 2σ detection of EDE from *Planck data alone*

$$f_{\text{EDE}} = 0.08 \pm 0.04$$

$$H_0 = 70 \pm 1.5 \text{ km/s/Mpc}$$

Answer: no, EDE solution is still robust

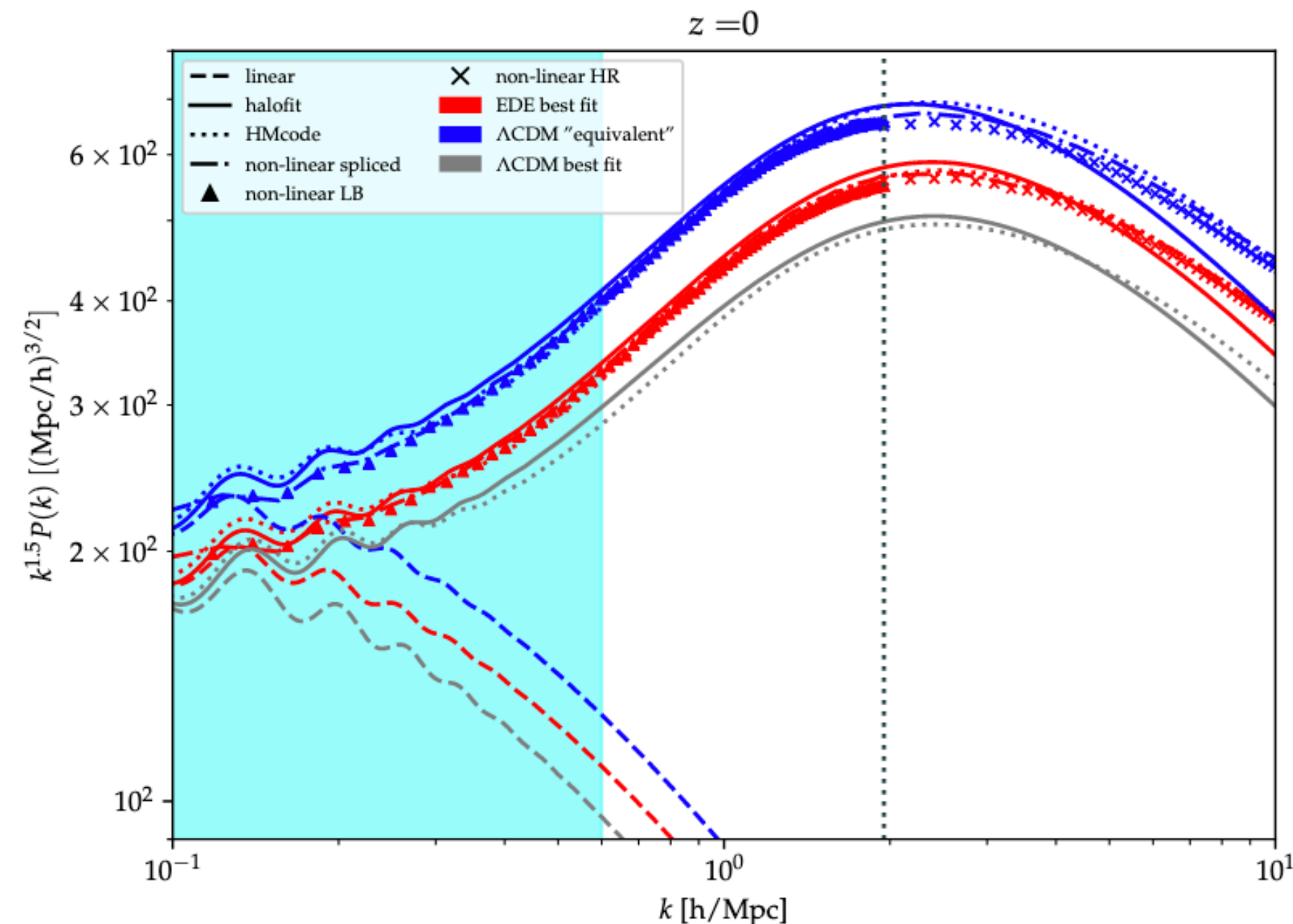
2. Is LSS data constraining enough to rule out EDE?



EDE non-linear $P(k)^*$ from halofit agrees well with results from N-body simulations

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EDE **non-linear** $P(k)^*$ from **halofit** agrees well with results from **N-body simulations**

1pEDE tested against **Planck+BAO+SNIa+SHoEs** and WL data from **KiDS/Viking+DES**: S_8 tension persists, but **fit is not significantly degraded wrt Λ CDM**, and solution to the H_0 tension survives

Murgia, GFA, Poulin 2107.10291

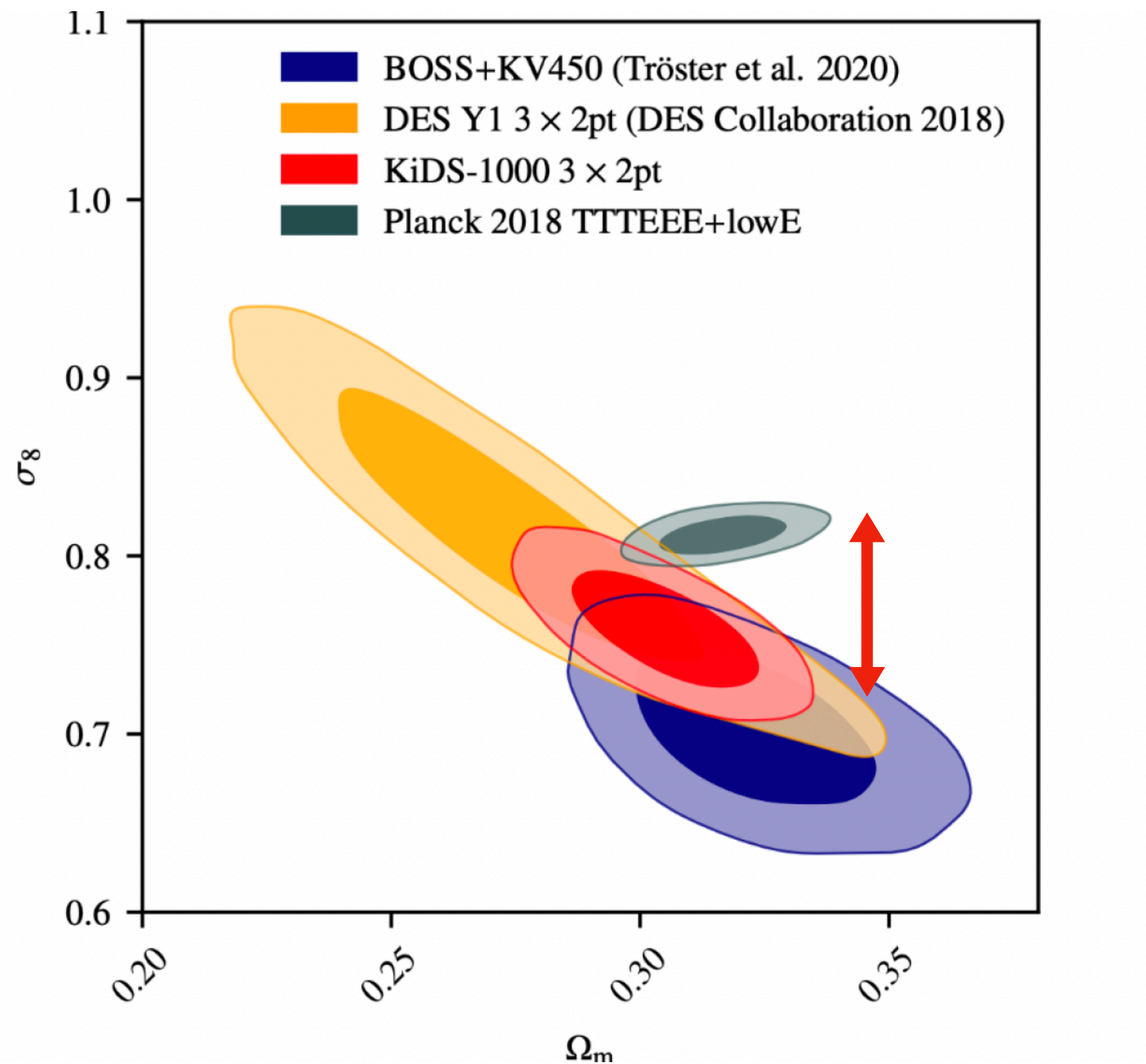
$$f_{\text{EDE}} = 0.09^{+0.03}_{-0.02} \quad H_0 = 71.3 \pm 0.9 \text{ km/s/Mpc}$$

*Intrinsic effect of EDE is a power suppression, but the shift of the Λ CDM params. leads to an enhancement **32**

What is needed to resolve the S_8 tension?

Di Valentino++ 2008.11285

$$S_8 \equiv \sigma_8 \sqrt{\Omega_m / 0.3}$$



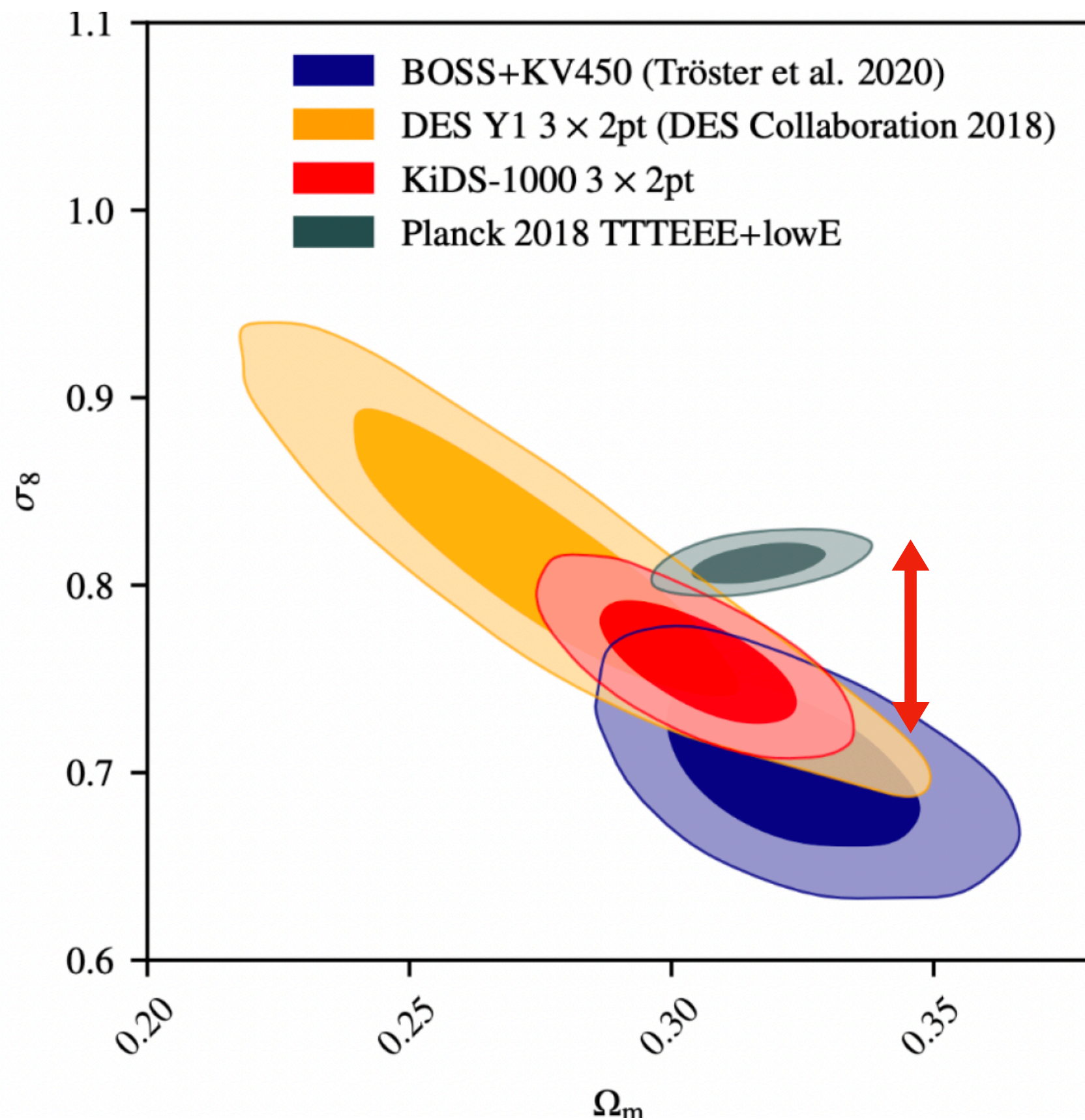
Ω_m should be left unchanged

$$\sigma_8 = \int P_m(k, z = 0) W_R^2(k) d \ln k$$

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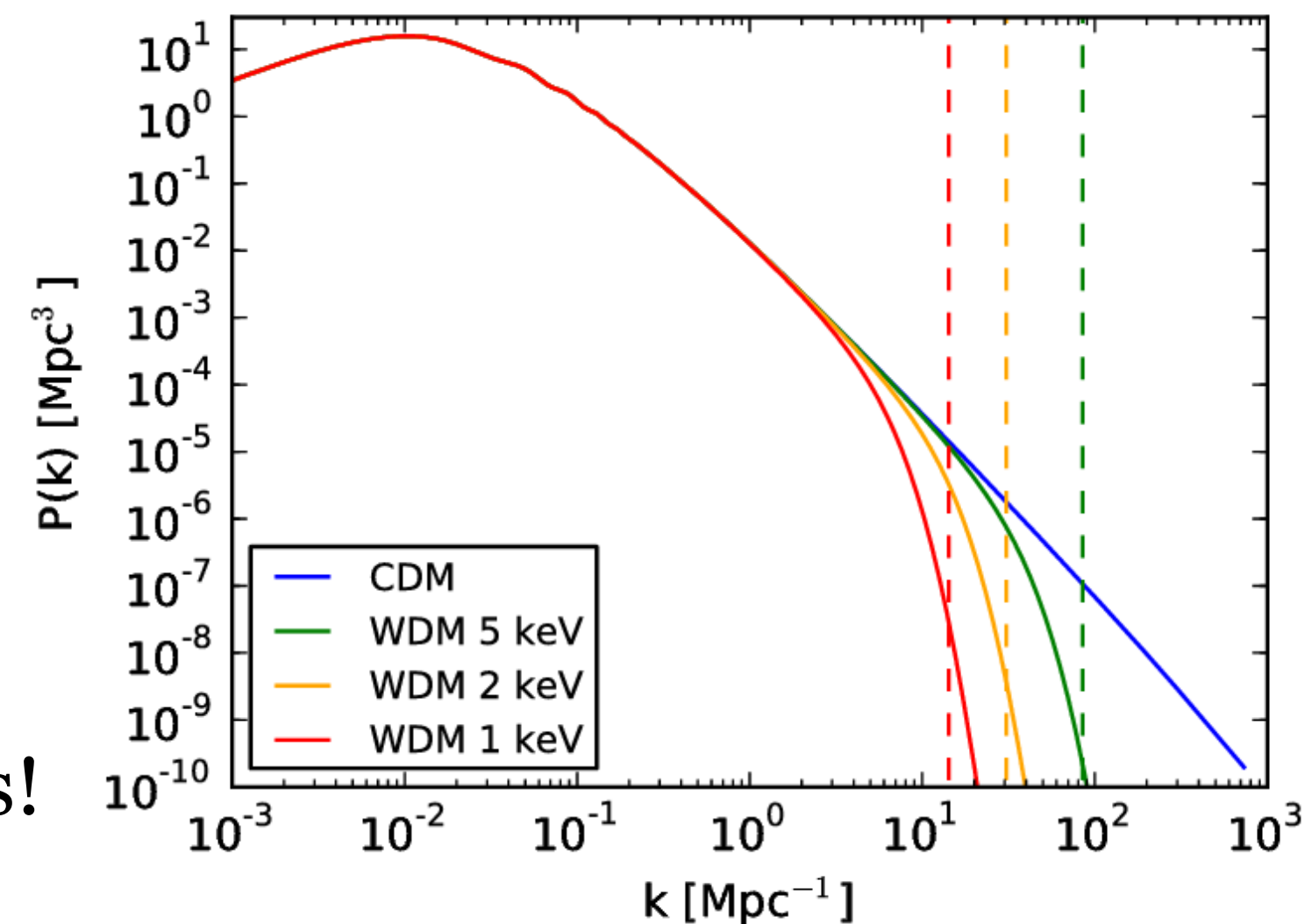


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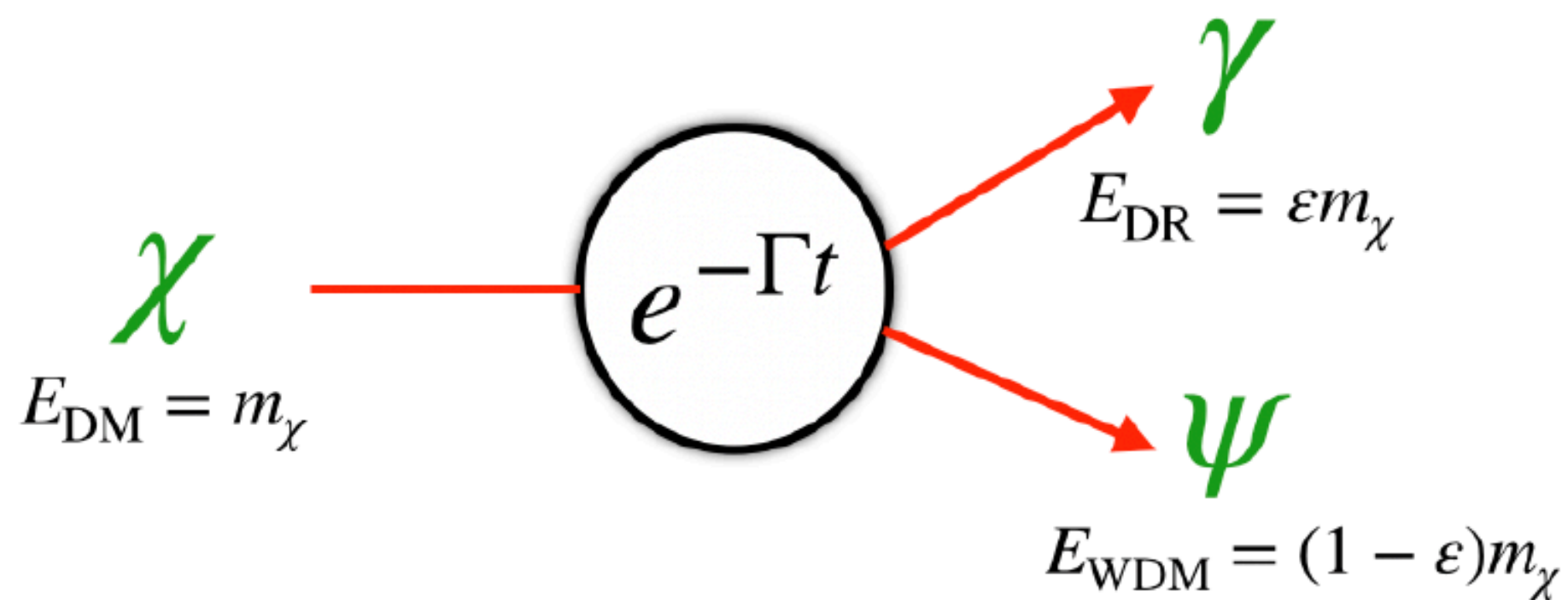
Need to **suppress power** at scales $k \sim 0.1 - 1 h/\text{Mpc}$



Ex: Warm Dark Matter
Very constrained by many probes!

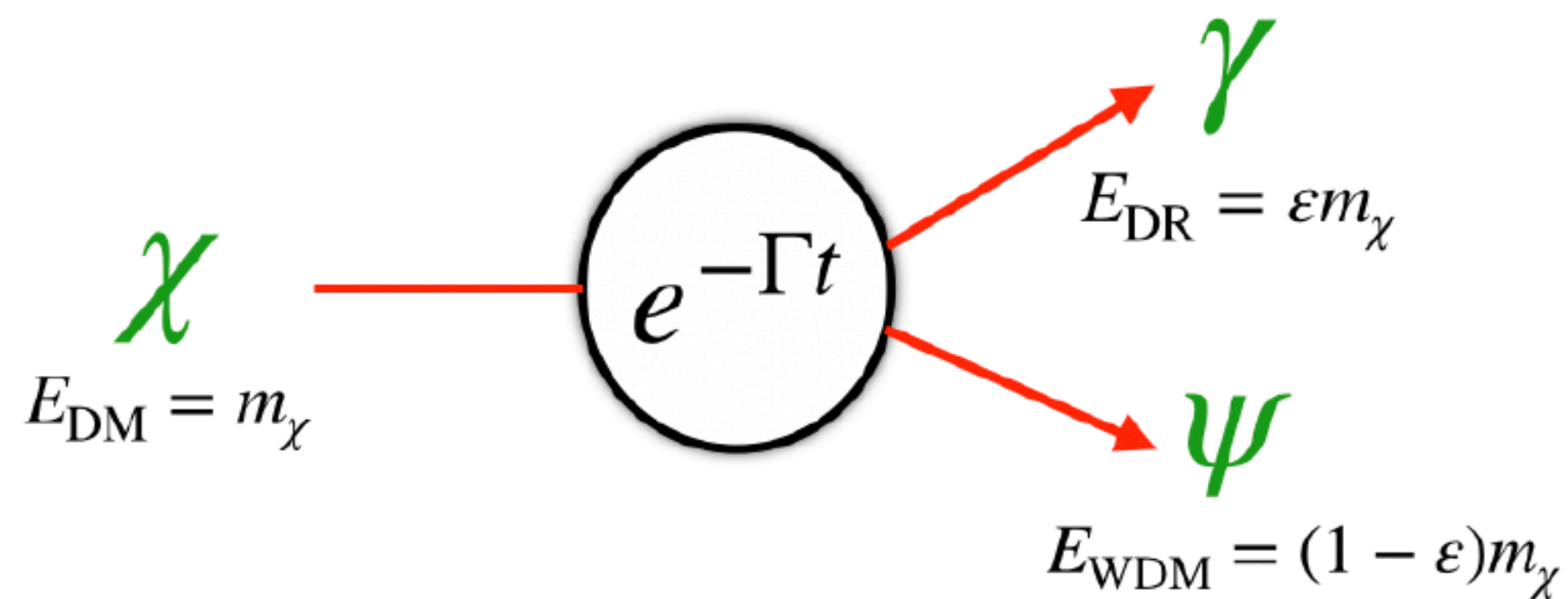
2-body Dark Matter decay

We explore DM decays to massless (**Dark Radiation**) and massive (**Warm Dark Matter**) particles, $\chi(\text{DM}) \rightarrow \gamma(\text{DR}) + \psi(\text{WDM})$



2-body Dark Matter decay

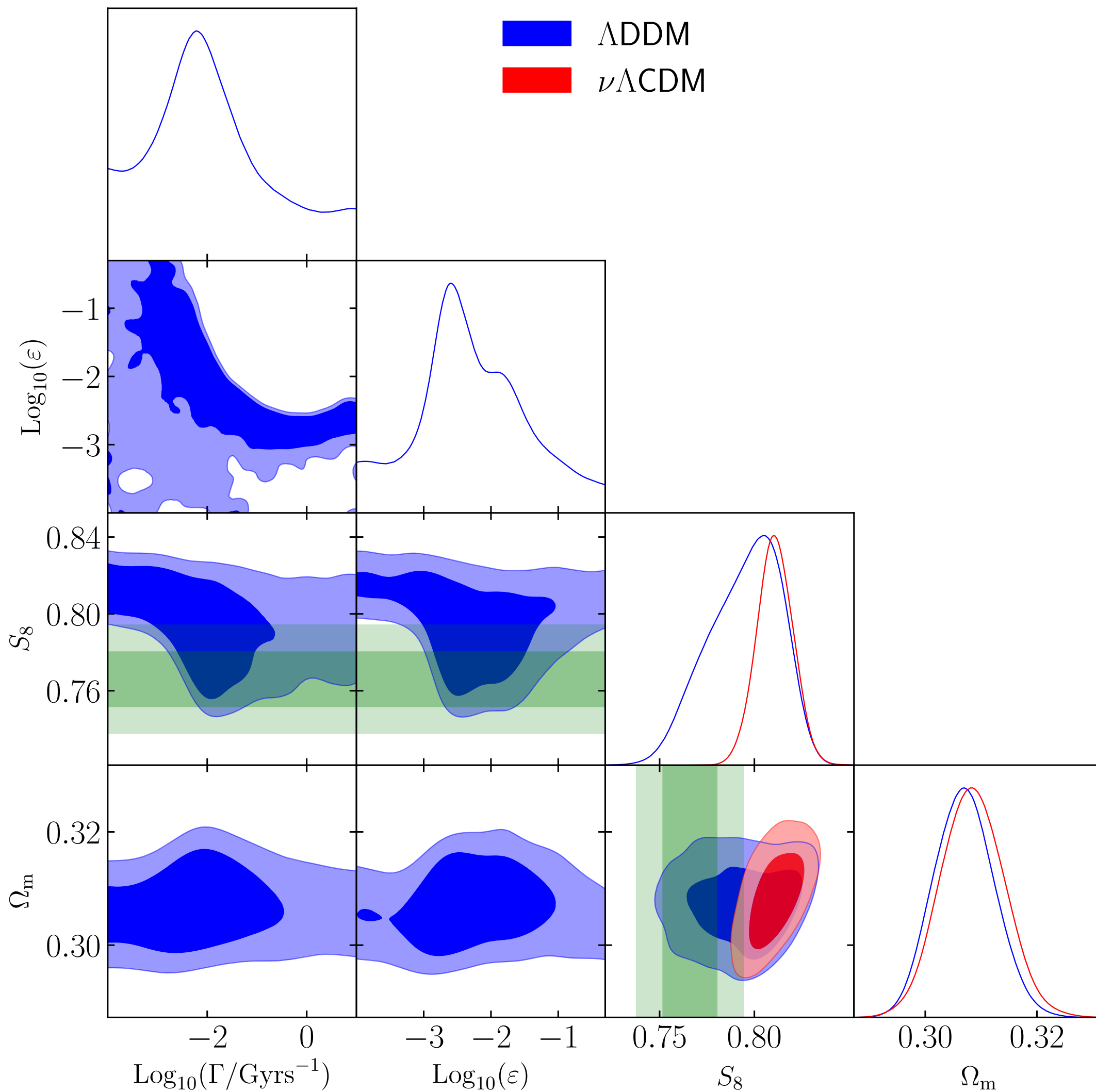
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The model is fully specified by:

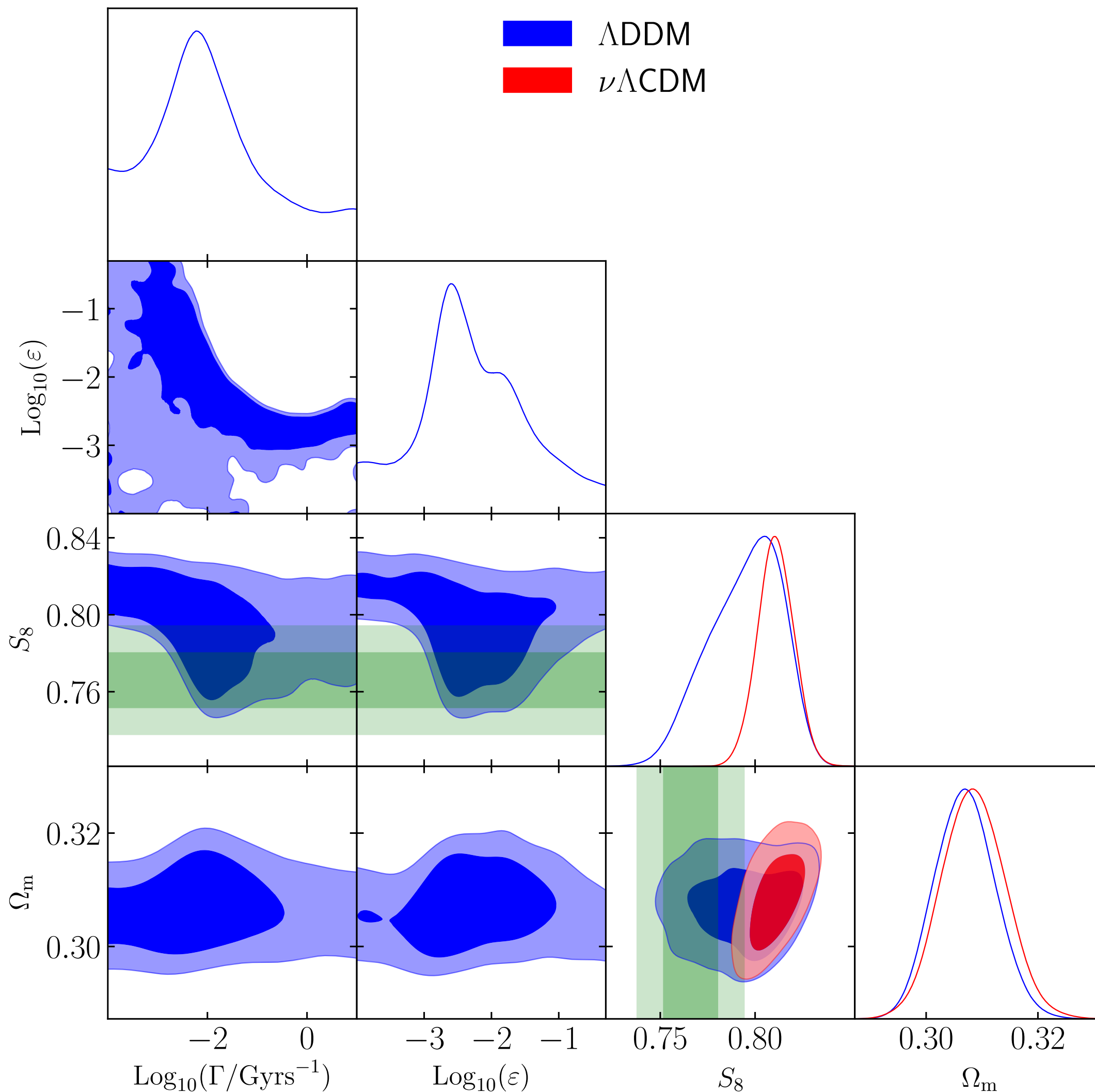
$$\{\Gamma, \varepsilon\} \quad \text{where} \quad \varepsilon = \frac{1}{2} \left(1 - \frac{m_\psi^2}{m_\chi^2} \right) \begin{cases} = \mathbf{0} & \text{for } \Lambda\text{CDM} \\ = \mathbf{1/2} & \text{for DM} \rightarrow \text{DR} \end{cases}$$

Explaining the S_8 tension



- MCMC analysis using Planck+BAO+SNIa+prior on S_8 from KIDS+BOSS+2dfLenS

Explaining the S_8 tension



- MCMC analysis using Planck+BAO+SNIa+prior on S_8 from KIDS+BOSS+2dfLenS
- Reconstructed S_8 values are in excellent agreement with WL data!

	$\nu\Lambda$ CDM	Λ DDM
χ^2_{CMB}	1015.9	1015.2
$\chi^2_{S_8}$	5.64	0.002

$$\longrightarrow \Delta\chi^2_{\text{min}} \simeq -5.5$$

$$\Gamma^{-1} \simeq 55 (\epsilon/0.007)^{1.4} \text{ Gyr}$$

Why does the 2-body DM decay work better than massive neutrinos?

The 2-body decay gives a better fit thanks to the **time-dependence of the power suppression** and the cut-off scale

