

# Big-bang nucleosynthesis

Cyril Pitrou



---

# The Origin of Chemical Elements

R. A. ALPHER\*

*Applied Physics Laboratory, The Johns Hopkins University,  
Silver Spring, Maryland*

AND

H. BETHE

*Cornell University, Ithaca, New York*

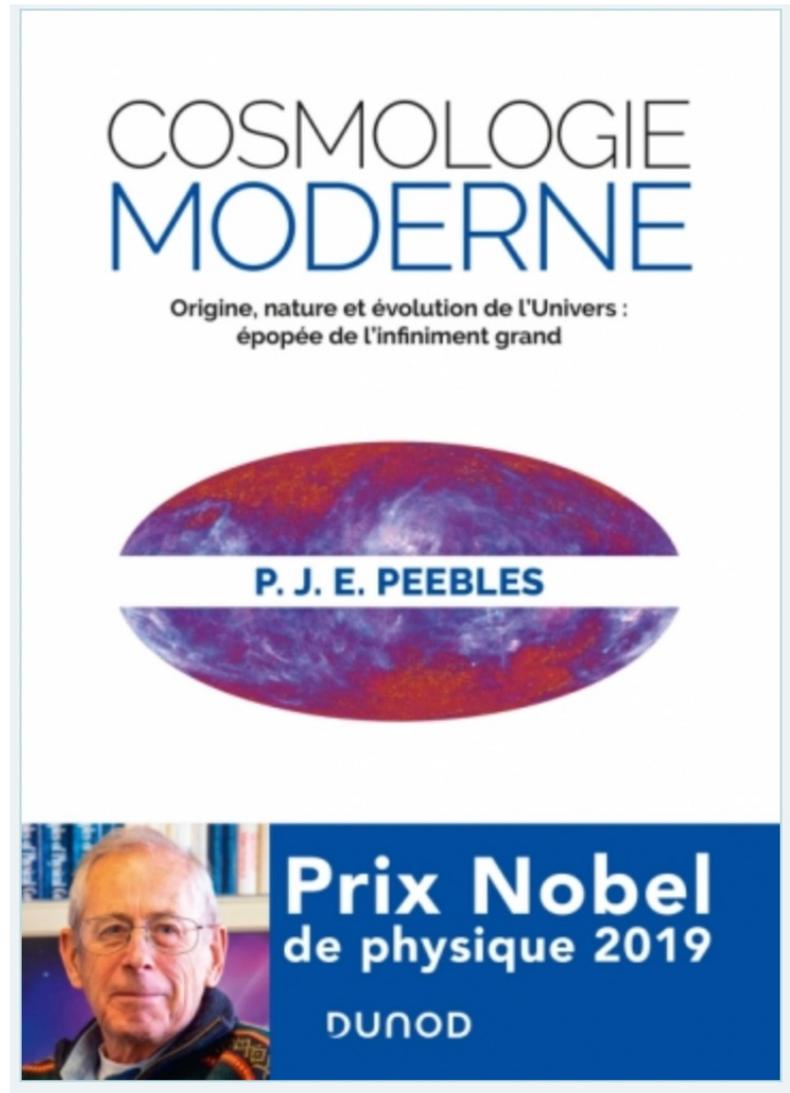
AND

G. GAMOW

*The George Washington University, Washington, D. C.*

February 18, 1948

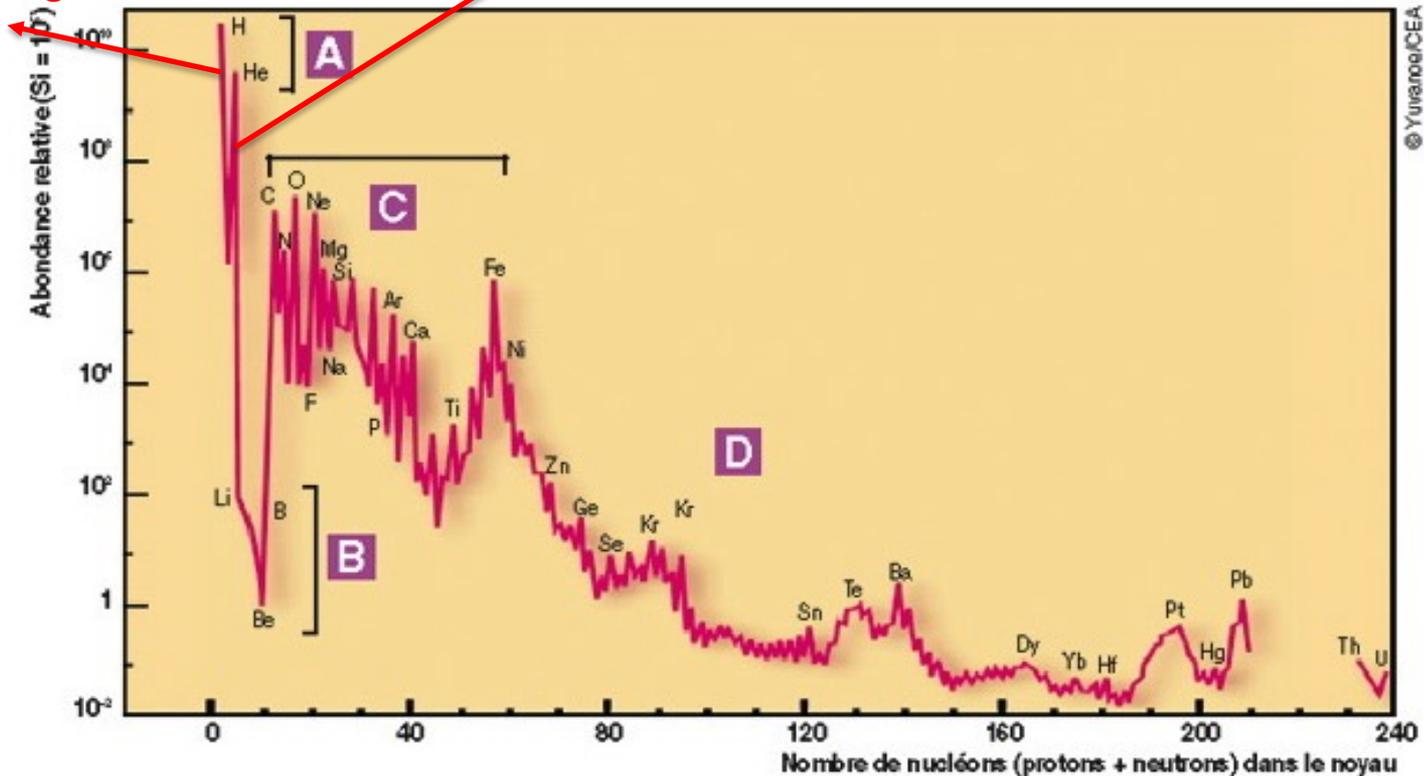
To know more about the History of BBN :



# ❑ Origin of elements

1. Very **light elements** : Primordial Universe
2. Some light elements : spallation from cosmic rays
3. Heavy elements : stars and star explosions

75% Hydrogen → 24% Helium



## Outline

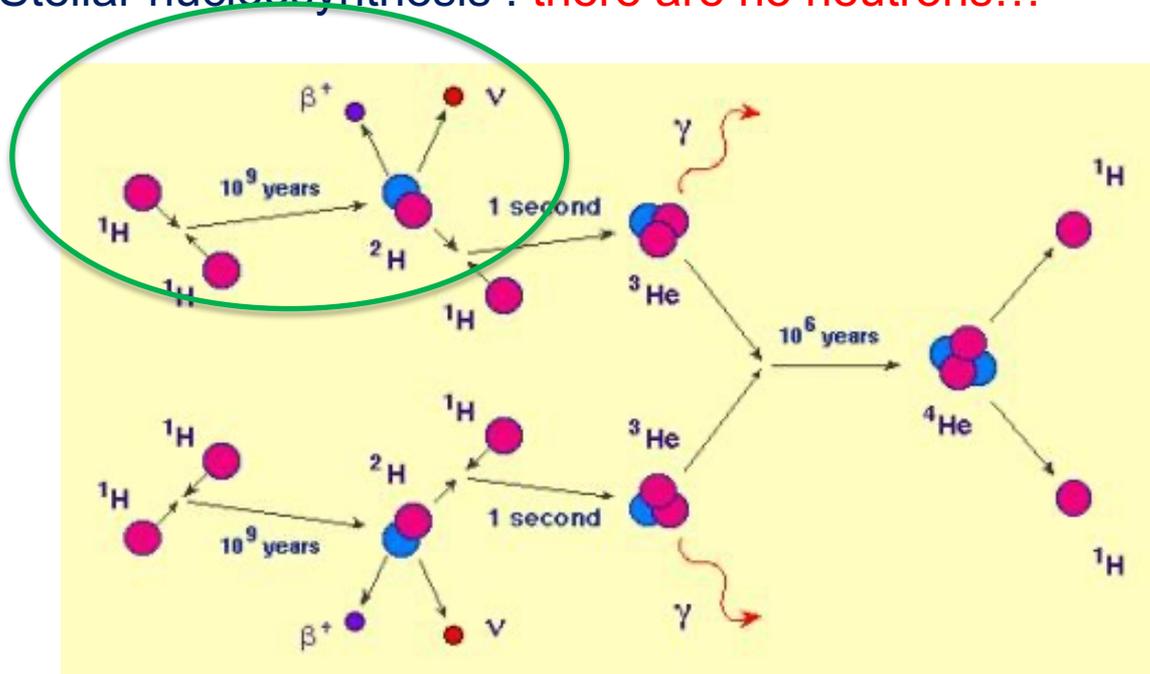
1. BBN in a nutshell
2. The hunt for precision
3. Cosmological parameters

# Outline

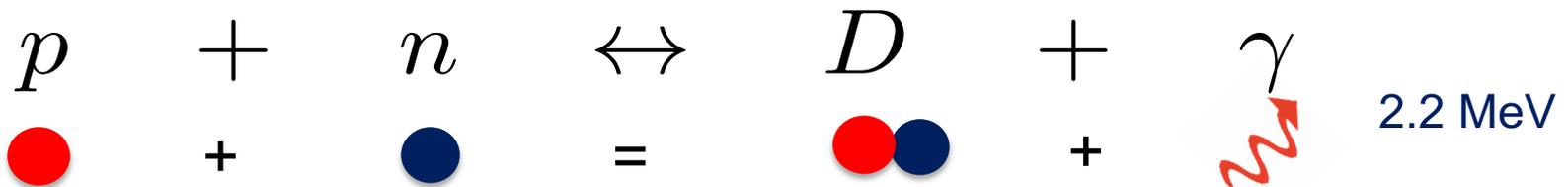
1. BBN in a nutshell
2. The Hunt for precision
3. Cosmological parameters

# □ Why primordial synthesis ?

## 1. Stellar nucleosynthesis : there are no neutrons...

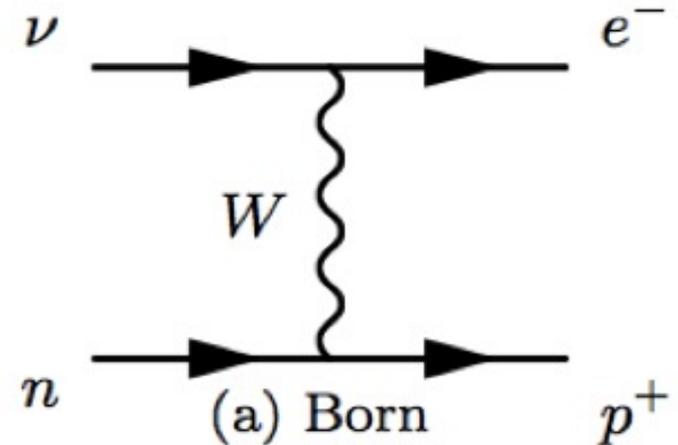
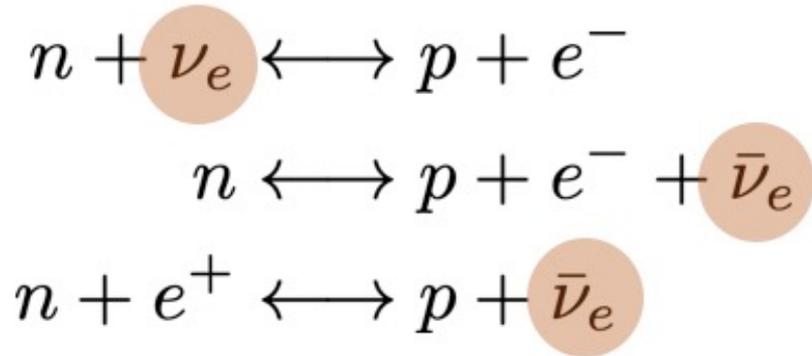


## 2. Primordial nucleosynthesis : there was plenty of neutrons !



When does it happen ? How many neutrons were available ?

## □ Neutron/proton conversions



- If enough interactions, then statistical equilibrium

$$n = e^{-\frac{E}{k_B T}}$$

Baryons are non-relativistic:  $E \simeq m$

$$m_p = 938.2 \text{ MeV}$$

$$m_n = m_p + 1.3 \text{ MeV}$$

Protons  $n_p = e^{-m_p/T}$

Neutrons  $n_n = e^{-m_n/T} = n_p e^{-(m_n - m_p)/T}$

□ Friedmann equation in radiation era

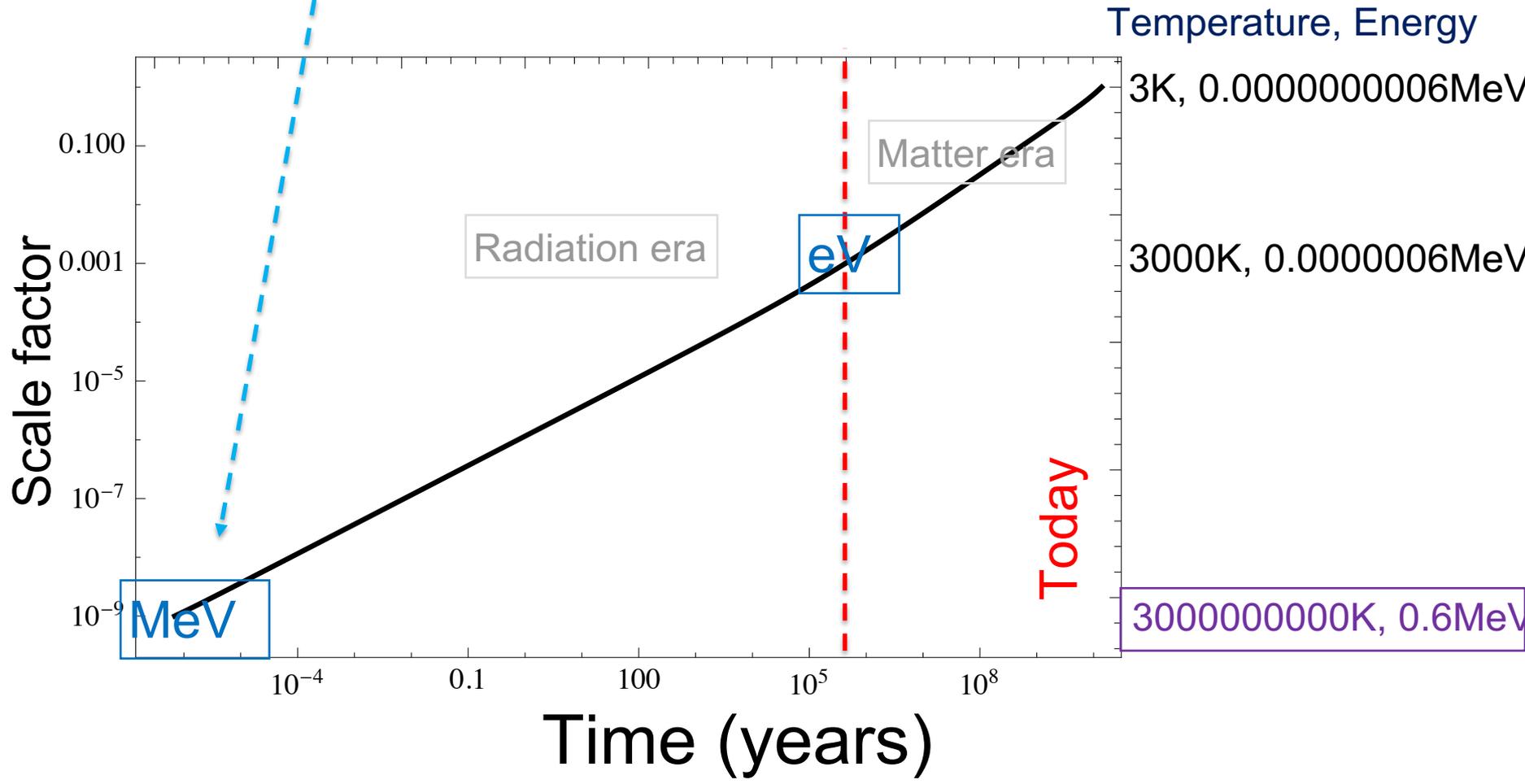
$$H^2 = \left( \frac{\dot{a}}{a} \right)^2 = \frac{8\pi G}{3} \rho$$
$$= \frac{8\pi G}{3} \left( \cancel{\frac{\rho_{\text{matter}}^{\text{today}}}{a^3}} + \frac{\rho_{\text{rad}}^{\text{today}}}{a^4} \right)$$

Radiation dominated universe  $a \propto t^{1/2}$

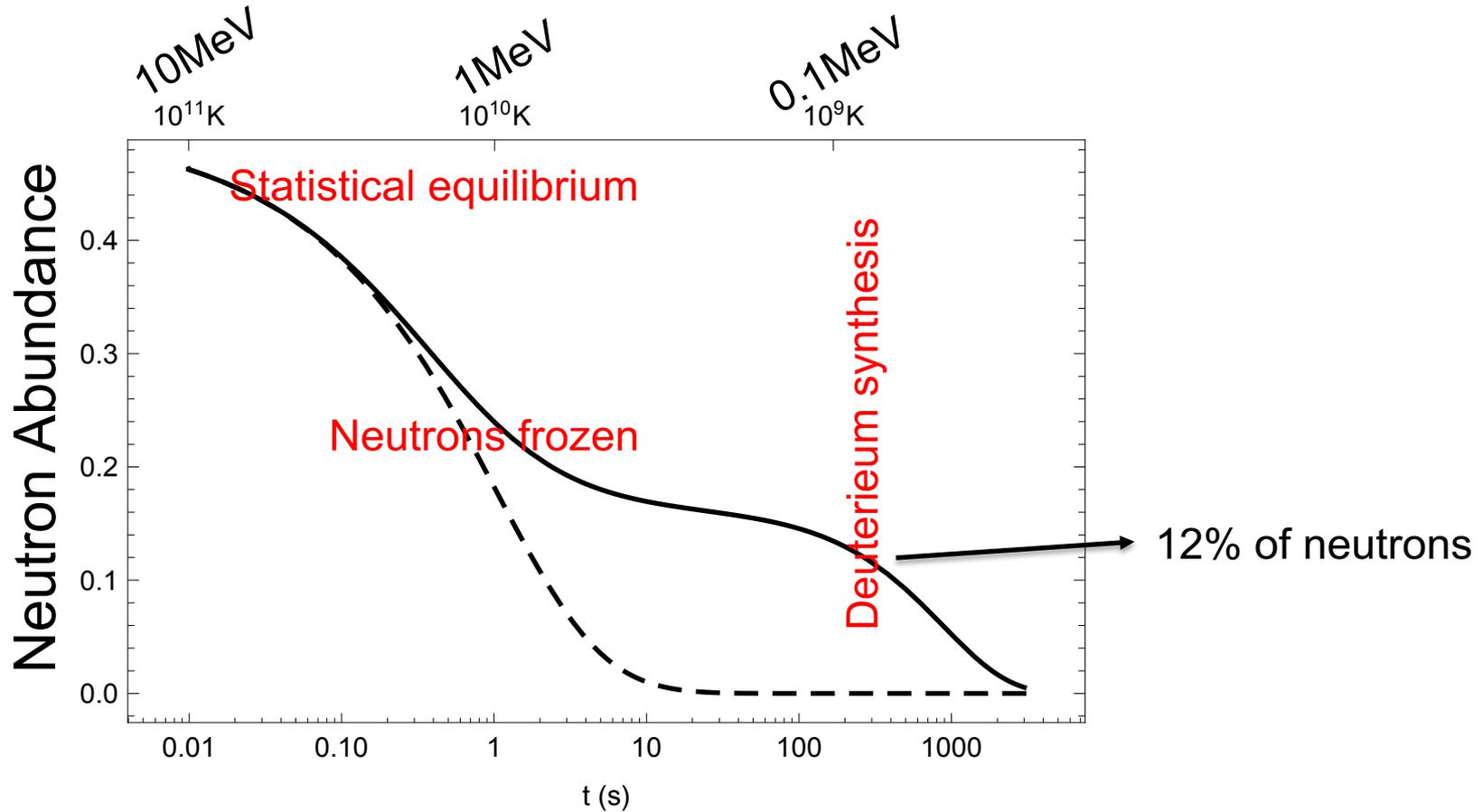
$$\rho_{\text{rad}} \propto T^4 \propto a^{-4}$$

$$T^2 \propto 1/t$$

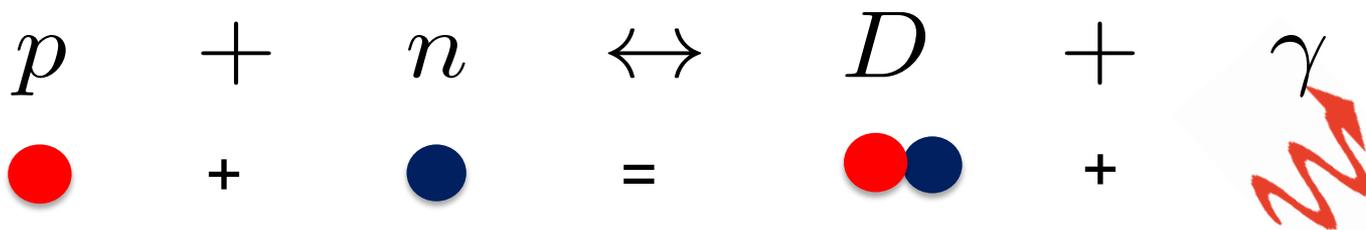
□ The MeV scale



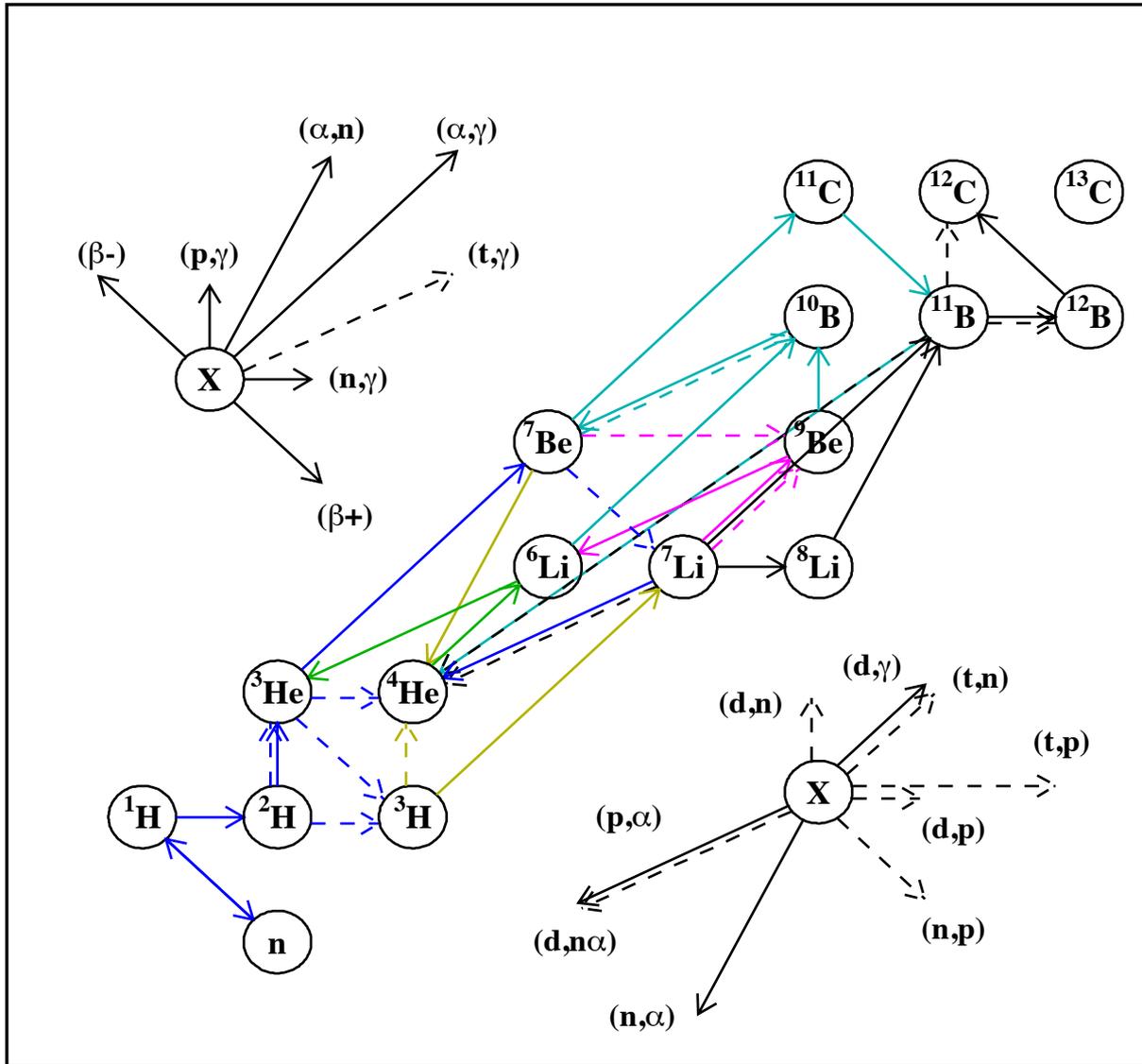
# Evolution of neutrons

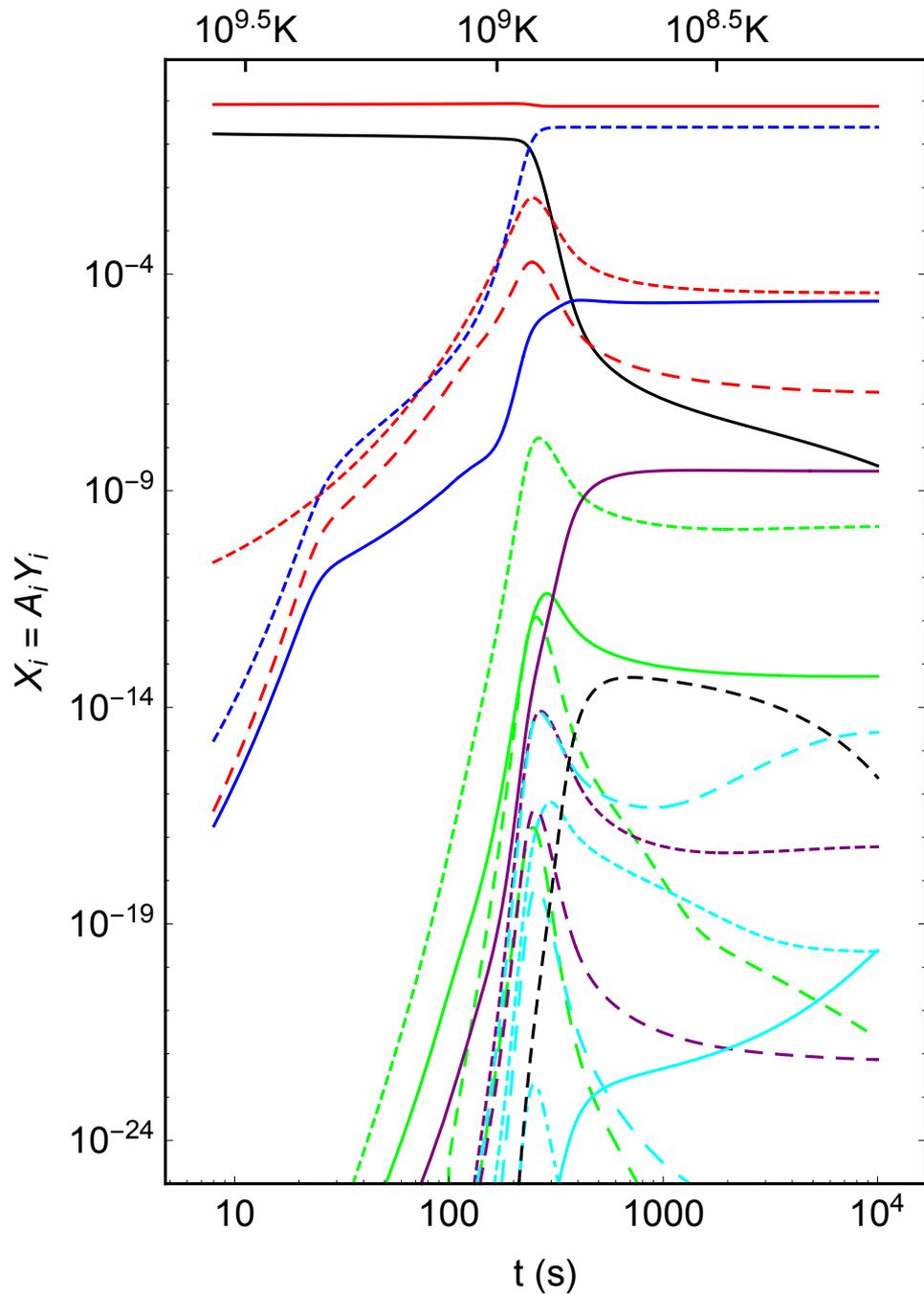


## Deuterium



# Nuclear network





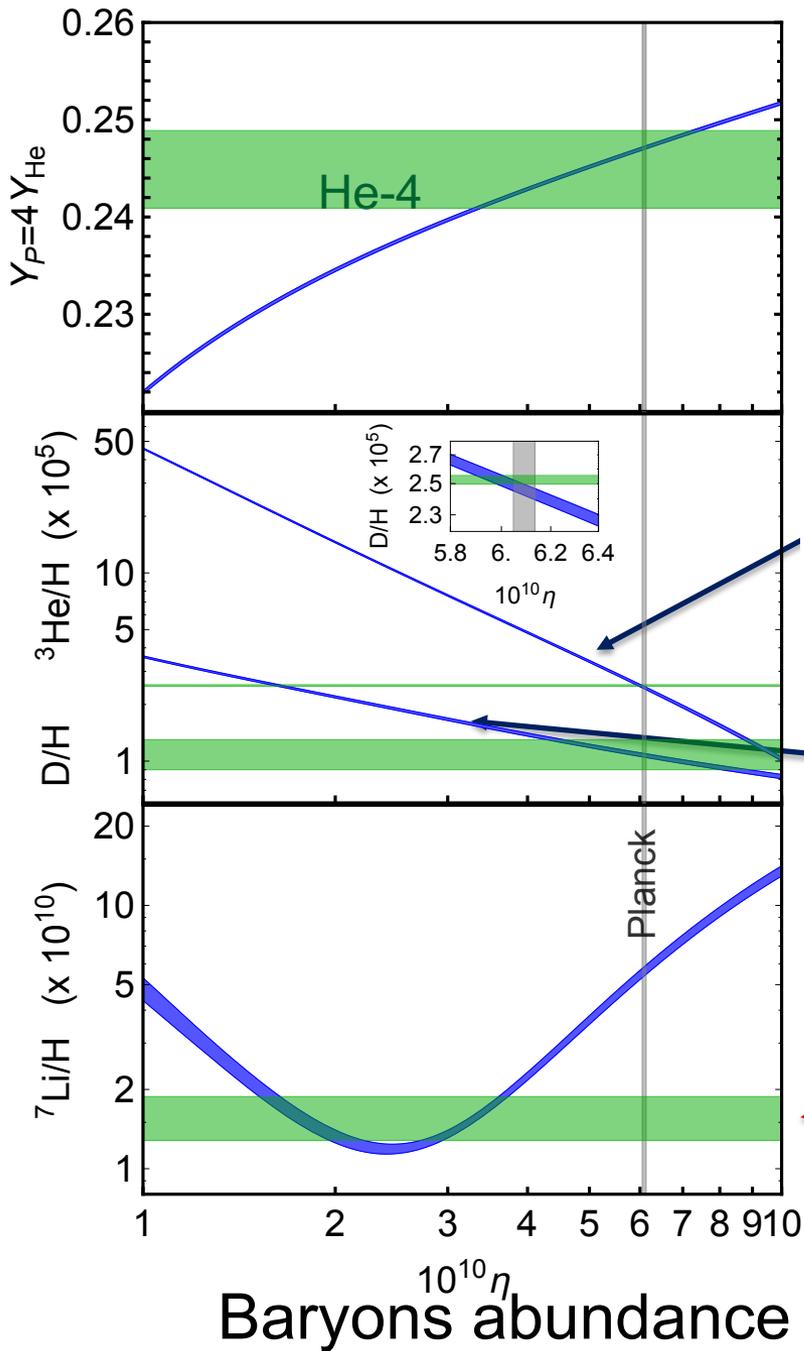
- n
- $^1\text{H}$
- -  $^2\text{H}$
- -  $^3\text{H}$
- $^3\text{He}$
- -  $^4\text{He}$
- $^6\text{Li}$
- -  $^7\text{Li}$
- -  $^8\text{Li}$
- -  $^9\text{Li}$
- $^7\text{Be}$
- -  $^9\text{Be}$
- -  $^{10}\text{Be}$
- $^8\text{B}$
- -  $^{10}\text{B}$
- -  $^{11}\text{B}$
- -  $^{12}\text{B}$
- -  $^{13}\text{B}$
- -  $^{11}\text{C}$

Public BBN codes:

**PRIMAT**

**AlterBBN**

**ParthENoPE**



*Aver et al. 2020*

$$Y_p = 0.2453 \pm 0.0034$$

**1.4 %**

Deuterium

$$D/H = (2.527 \pm 0.030) \times 10^{-5}$$

**1.2 %**

*Cooke et al. 2016*

He-3

Li-7

**Lithium problem**

# Outline

1. BBN in a nutshell
2. The Hunt for precision
3. Cosmological parameters

# Hunt for precision

- Plasma physics
- Neutrino physics
- Weak rates
- Nuclear rates

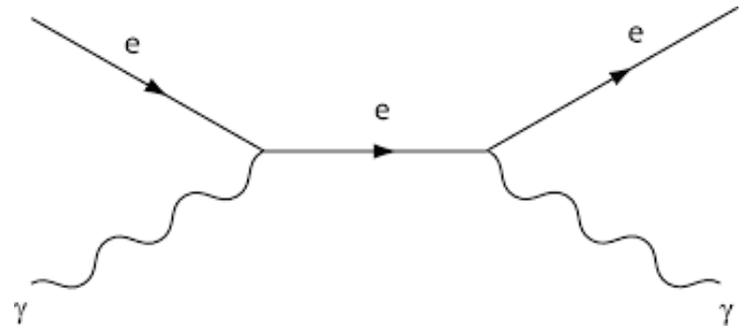
□ QED reactions :

- Momentum exchange :

$$\gamma + e^{\pm} \leftrightarrow \gamma + e^{\pm}$$

- Annihilations

$$\gamma + \gamma \leftrightarrow e^{\pm} + e^{\mp}$$



Collisions

$$\begin{cases} \dot{\rho}_{e^{\pm}} + 3H(\rho_{e^{\pm}} + P_{e^{\pm}}) = \dot{q}_{e^{\pm}} \\ \dot{\rho}_{\gamma} + 4H(\rho_{\gamma}) = \dot{q}_{\gamma} \end{cases}$$

□ Temperature evolution, *method 1*

$$\dot{q}_\gamma + \dot{q}_{e^\pm} = 0$$

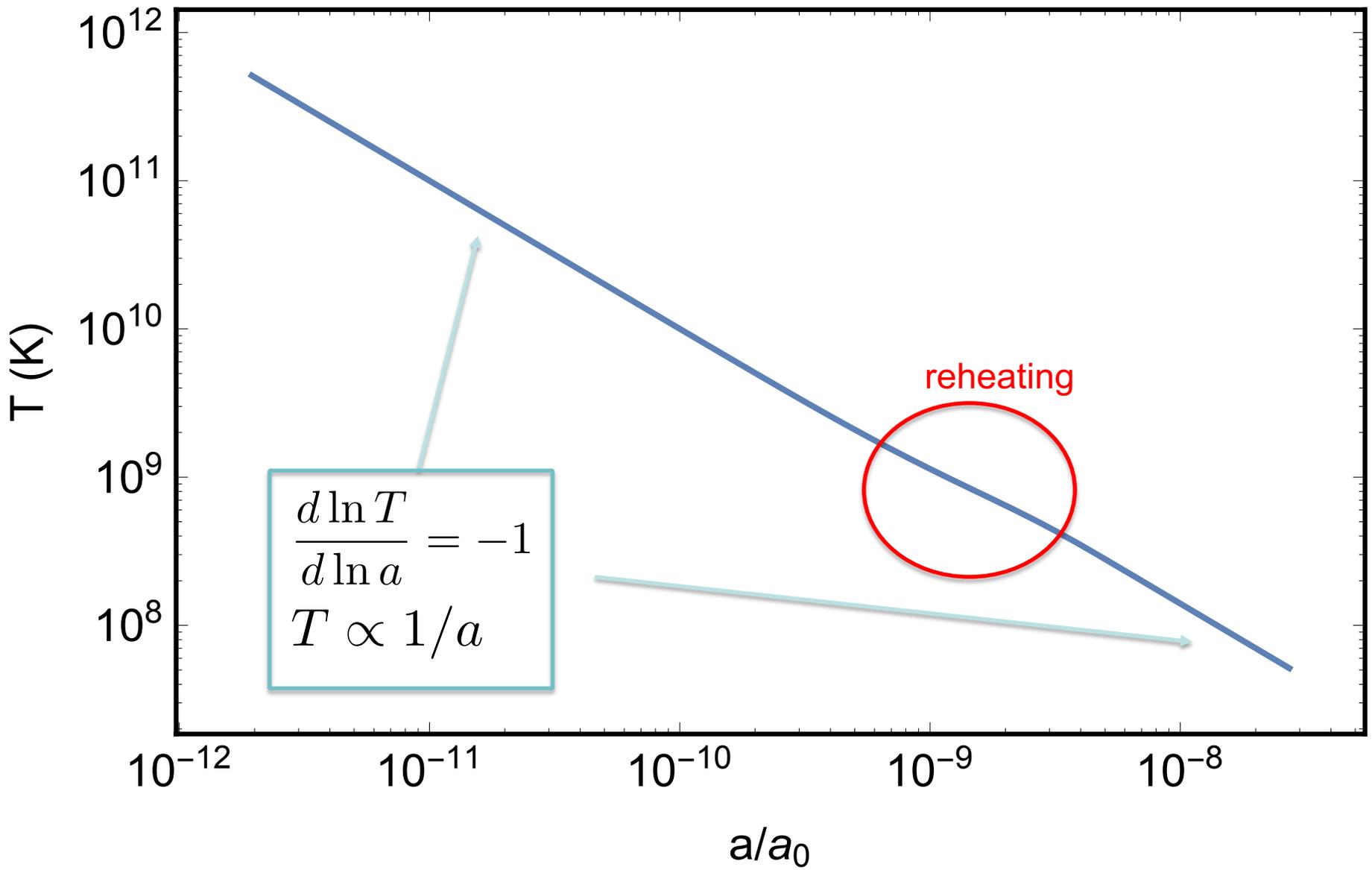
$$\dot{\rho}_{e^\pm} + \dot{\rho}_\gamma + 3H(\rho_{e^\pm} + P_{e^\pm}) + 4H\rho_\gamma = 0$$

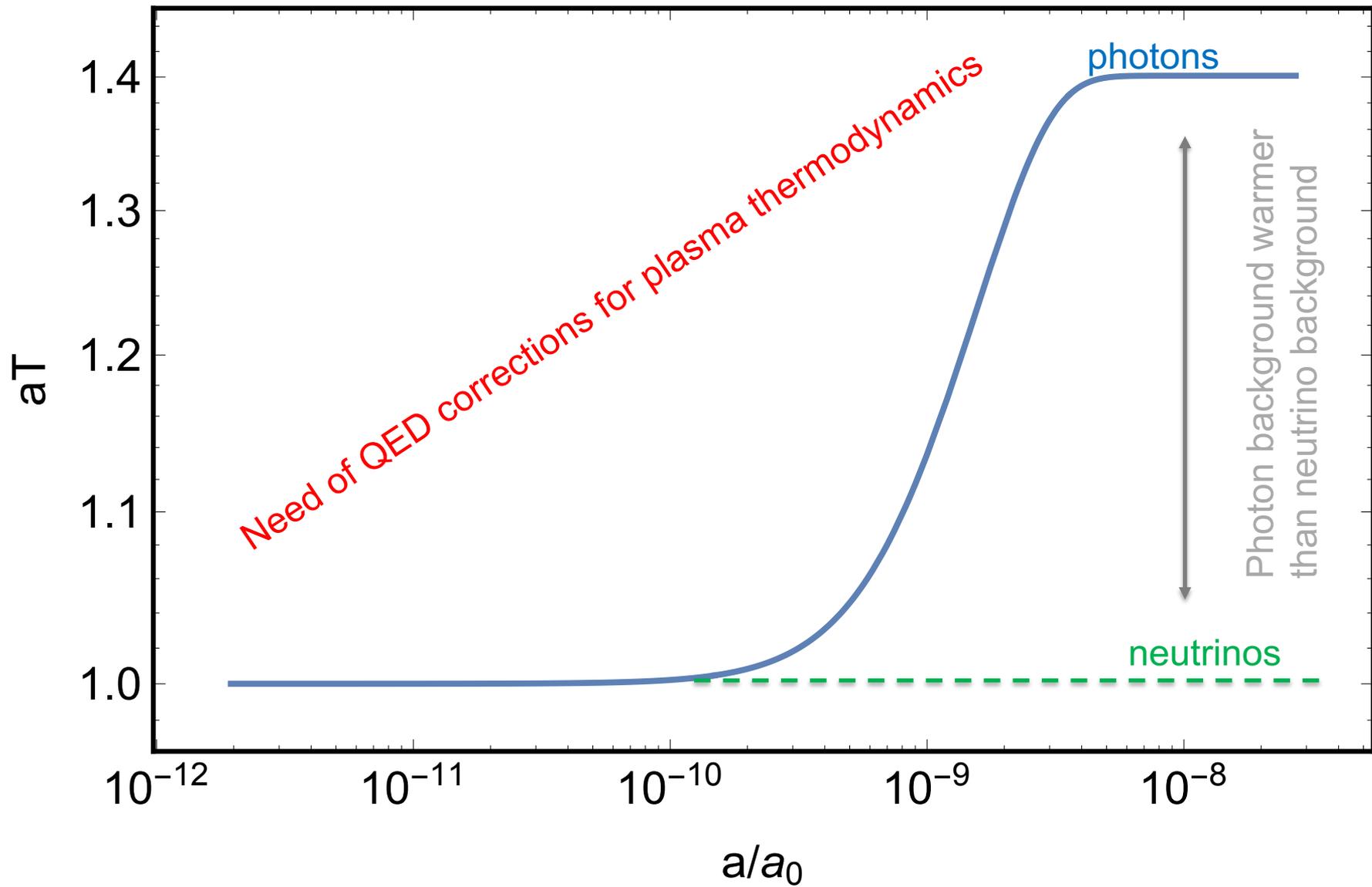
1. Using that thermodynamic quantities are functions of temperature only

$$\dot{\rho} = \dot{T} d\rho(T) / dT$$

2. Trading time for scale factor with :  $H = \frac{d \ln a}{dt}$

$$\frac{d \ln T}{d \ln a} = - \frac{3[\rho_{e^\pm}(T) + P_{e^\pm}(T)] + 4\rho_\gamma(T)}{T[\rho'_{e^\pm}(T) + \rho'_\gamma(T)]}$$





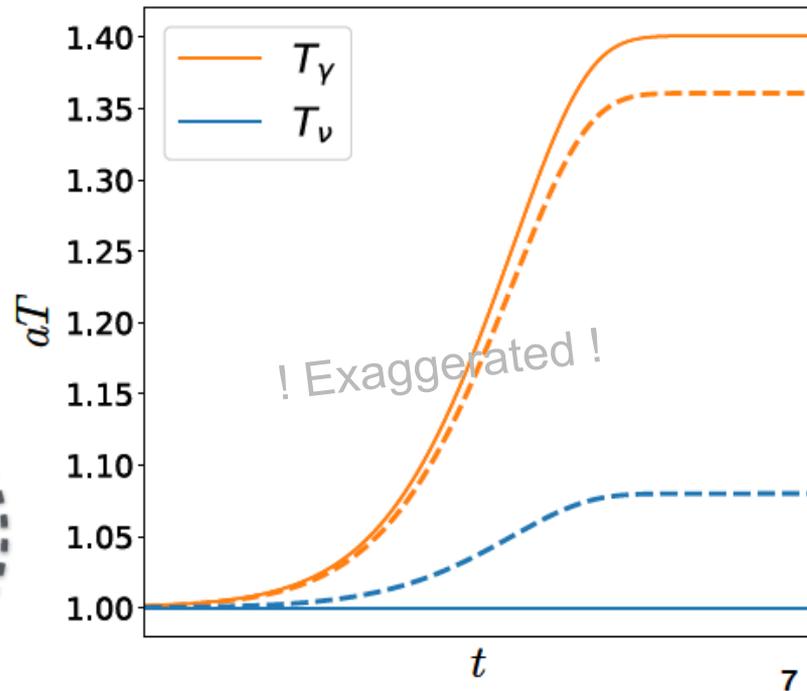
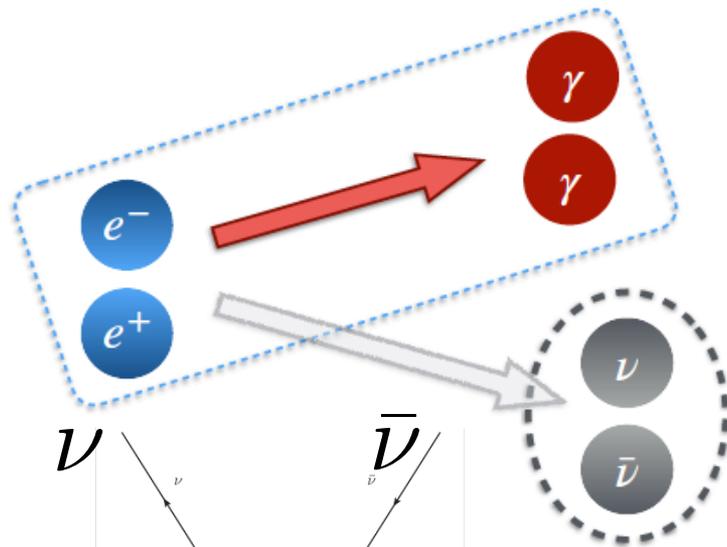
# Hunt for precision

- Plasma physics
- Neutrino physics
- Weak rates
- Nuclear rates

# Beyond the instantaneous decoupling approximation

- Overlap between decoupling and  $e^\pm$  annihilations

$\implies$  smaller  $T_\gamma$  and increased  $T_\nu$

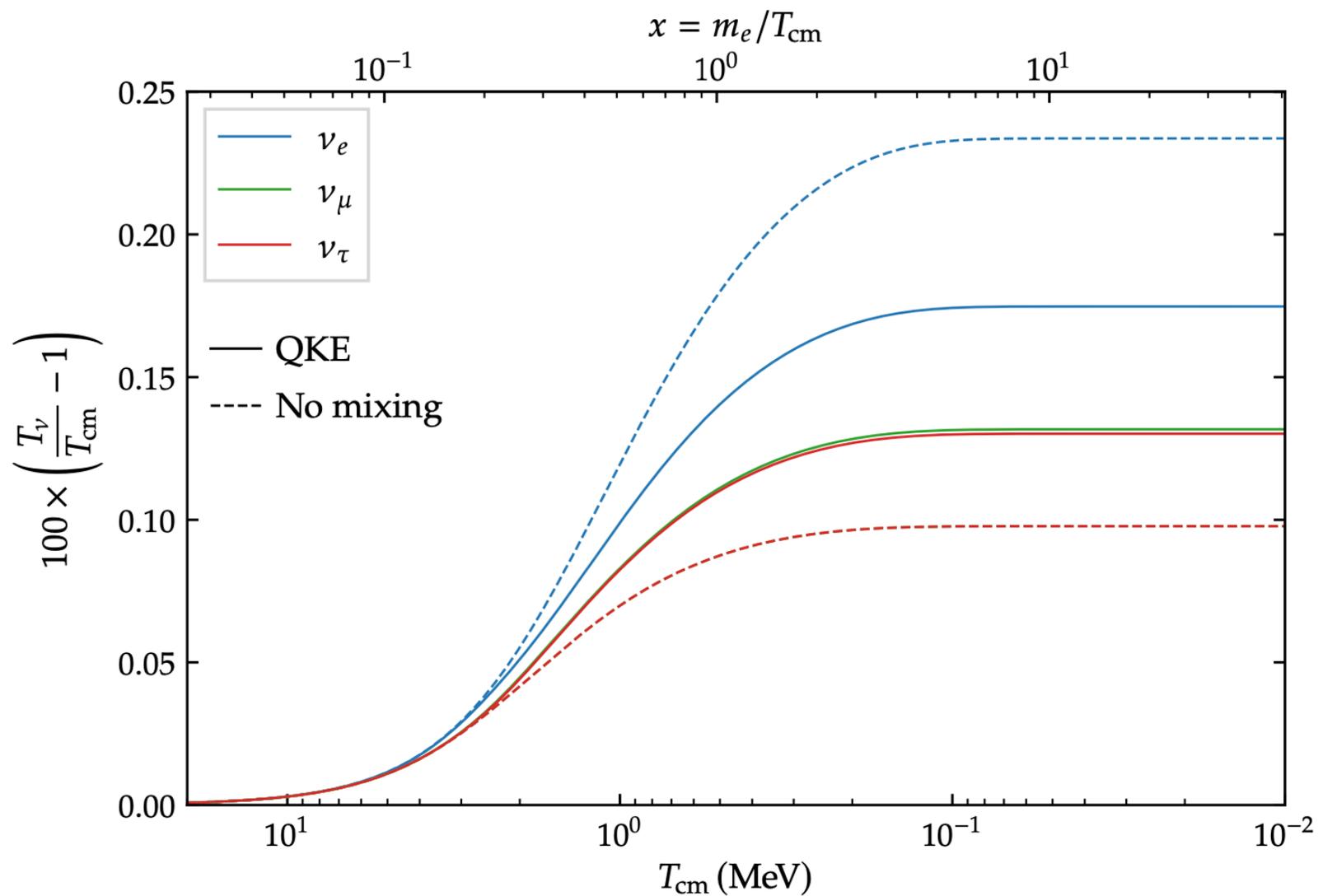


Froustey et al. 2020

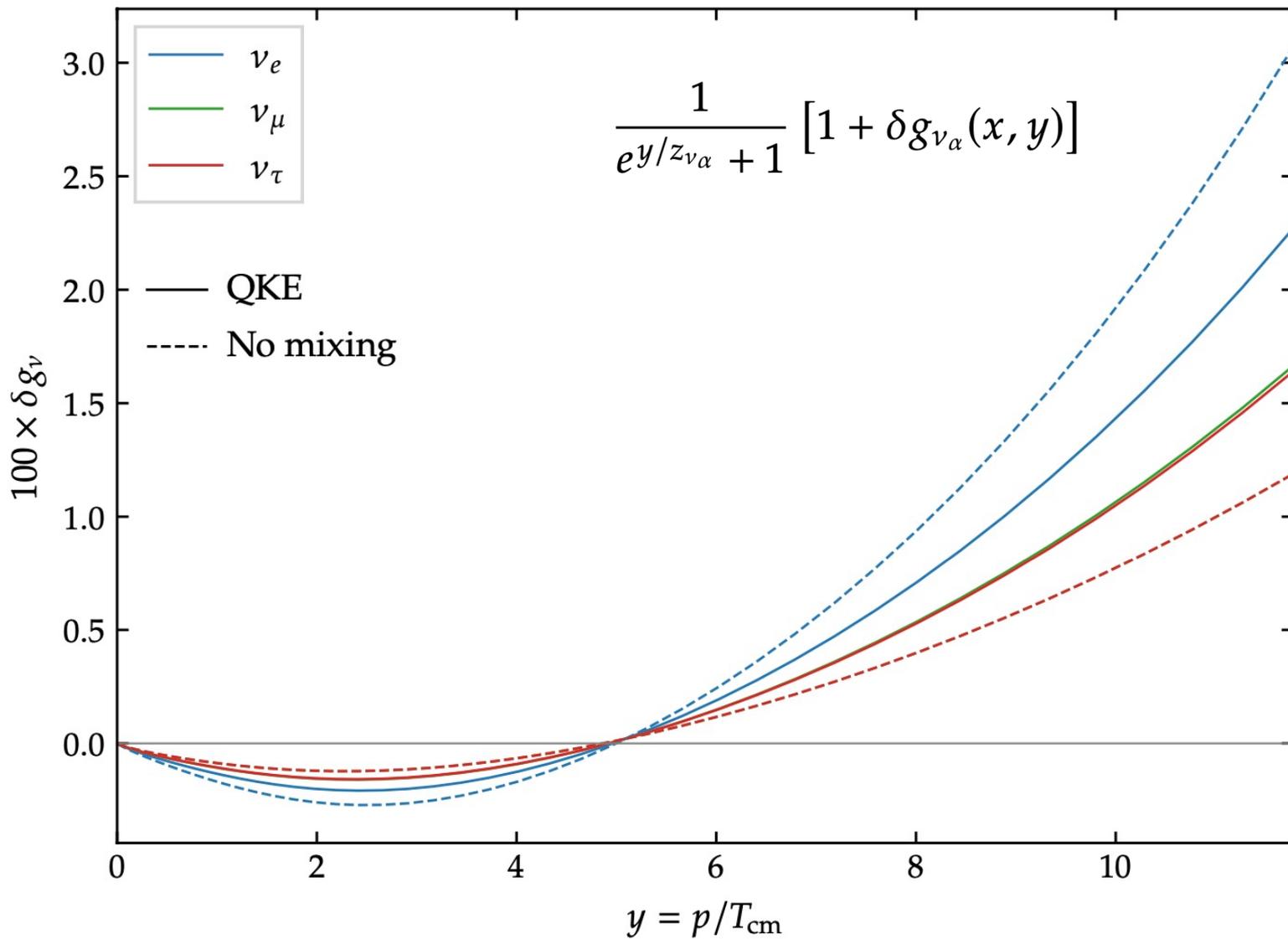
$$\rho_R = \rho_\nu + \rho_\gamma = \rho_\gamma \left( 1 + \frac{7}{8} N_{\text{eff}} \left( \frac{4}{11} \right)^{4/3} \right)$$

$$N_{\text{eff}} = 3.0440$$

# Neutrino mixing (aka oscillations) effects

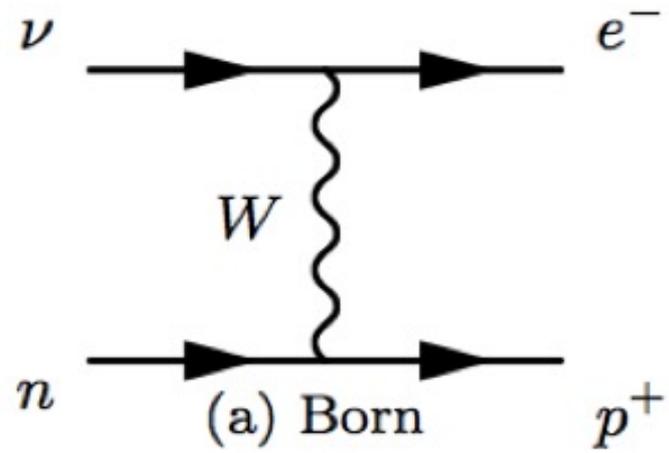


# Neutrino Fermi-Dirac spectrum distortions



# Hunt for precision

- Plasma physics
- Neutrino physics
- **Weak rates**
- Nuclear rates



# □ General expression of weak rates

$$\dot{n}_n + 3Hn_n = -n_n\Gamma_{n \rightarrow p} + n_p\Gamma_{p \rightarrow n}$$

$$\dot{n}_p + 3Hn_p = -n_p\Gamma_{p \rightarrow n} + n_n\Gamma_{n \rightarrow p}$$

Weak interaction Matrix element



$$n_n\Gamma = \int \Pi_i [d^3\mathbf{p}_i] (2\pi)^4 \delta^4(\underline{p}_n - \underline{p}_p + \alpha_\nu \underline{p}_\nu + \alpha_e \underline{p}_e) |M|^2 f_n(E_n) [1 - f_p(E_p)] f_\nu(\alpha_\nu E_\nu) f_e(\alpha_e E_e)$$

momentum conservation

$$\begin{cases} \alpha_i = 1 & \text{if initial particle} \\ \alpha_i = -1 & \text{if final particle} \end{cases}$$

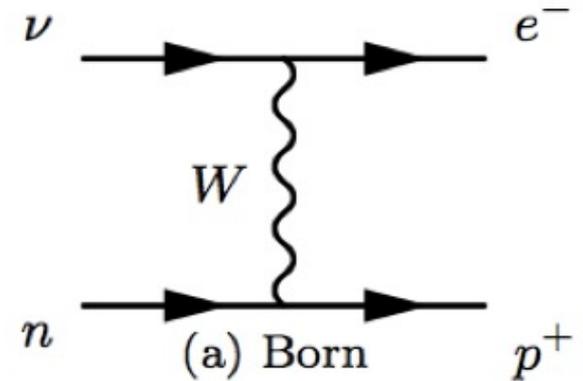
$$[d^3\mathbf{p}] \equiv \frac{d^3\mathbf{p}}{2E(2\pi)^3} \stackrel{\text{isotropy}}{=} \frac{4\pi p^2 dp}{2E(2\pi^3)}$$

$$g(-E) = 1 - g(E)$$

Fermi-Dirac Property

## □ Interaction Hamiltonian

$$\mathcal{H}_I = \frac{G_F}{\sqrt{2}} J_{e\nu}^\mu J_{pn, \mu}$$



$$J_{e\nu}^\mu = \bar{\nu} \gamma^\mu (1 - \gamma^5) e$$

$$J_{pn}^\mu = V_{ud} \bar{p} \left( \gamma^\mu (1 - g_A \gamma^5) + i \frac{f_{wm}}{m_N} 2 \Sigma^{\mu\nu} q_\nu \right) n$$

CKM angle

Axial current coupling

Weak-Magnetism

## □ BORN approximation

Simple integral on electron momentum :

$$\begin{aligned}\bar{\Gamma}_{n \rightarrow p} &= \bar{\Gamma}_{n \rightarrow p+e} + \bar{\Gamma}_{n+e \rightarrow p} \\ &= K \int_0^\infty p^2 dp [\chi_+(E) + \chi_+(-E)],\end{aligned}$$

Fermi-Dirac distributions

$$\chi_{\pm}(E) \equiv (E_{\nu}^{\mp})^2 g_{\nu}(E_{\nu}^{\mp}) g(-E),$$

$$E_{\nu}^{\mp} \equiv E \mp \Delta,$$

$$K \equiv \frac{4G_W^2(1 + 3g_A^2)}{(2\pi)^3}.$$

→ Axial coupling constant

$$G_W = G_F V_{ud} \longrightarrow \text{CKM angle}$$

## Neutron lifetime as a proxy

$$K \equiv \frac{4G_W^2(1 + 3g_A^2)}{(2\pi)^3}$$

$$K = 1/(\tau_n \lambda_0 m_e^5)$$

$$\lambda_0 \simeq 1.75474$$

*Best values are :*

$$\tau_n \simeq 879.4 \pm 0.6 \text{ s}$$

$$V_{ud} = 0.97420(20)$$

$$g_A = 1.2723(23)$$

## Neutron lifetime as a proxy

$$K \equiv \frac{4G_W^2(1 + 3g_A^2)}{(2\pi)^3}$$

$$K = 1/(\tau_n \lambda_0 m_e^5)$$

$$\lambda_0 \simeq 1.75474$$

*Best values are :*

$$\tau_n \simeq \del{879.4 \pm 0.6 \text{ s}}$$

$$V_{ud} = 0.97420(20)$$

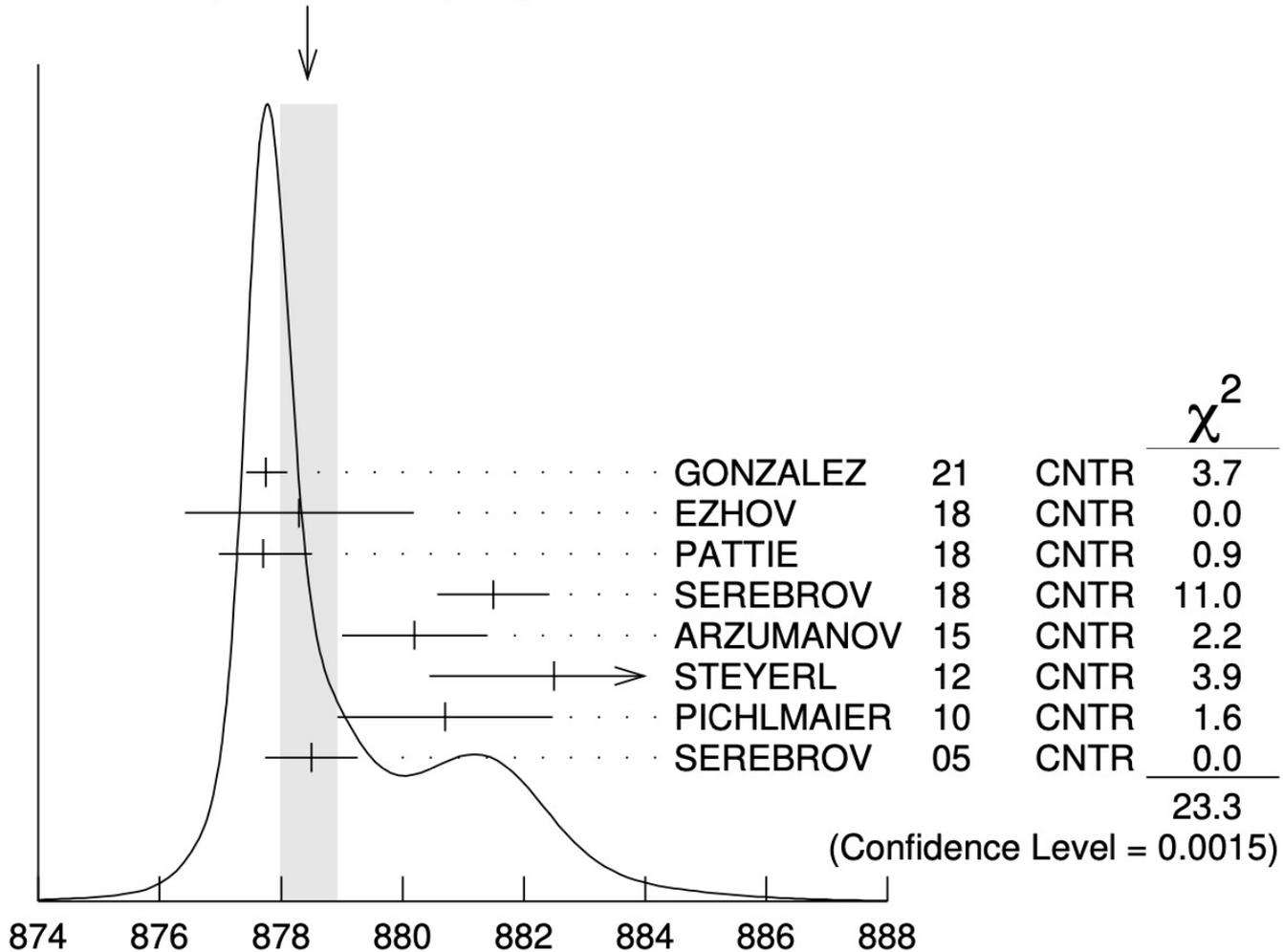
$$g_A = 1.2723(23)$$

PDG 2022

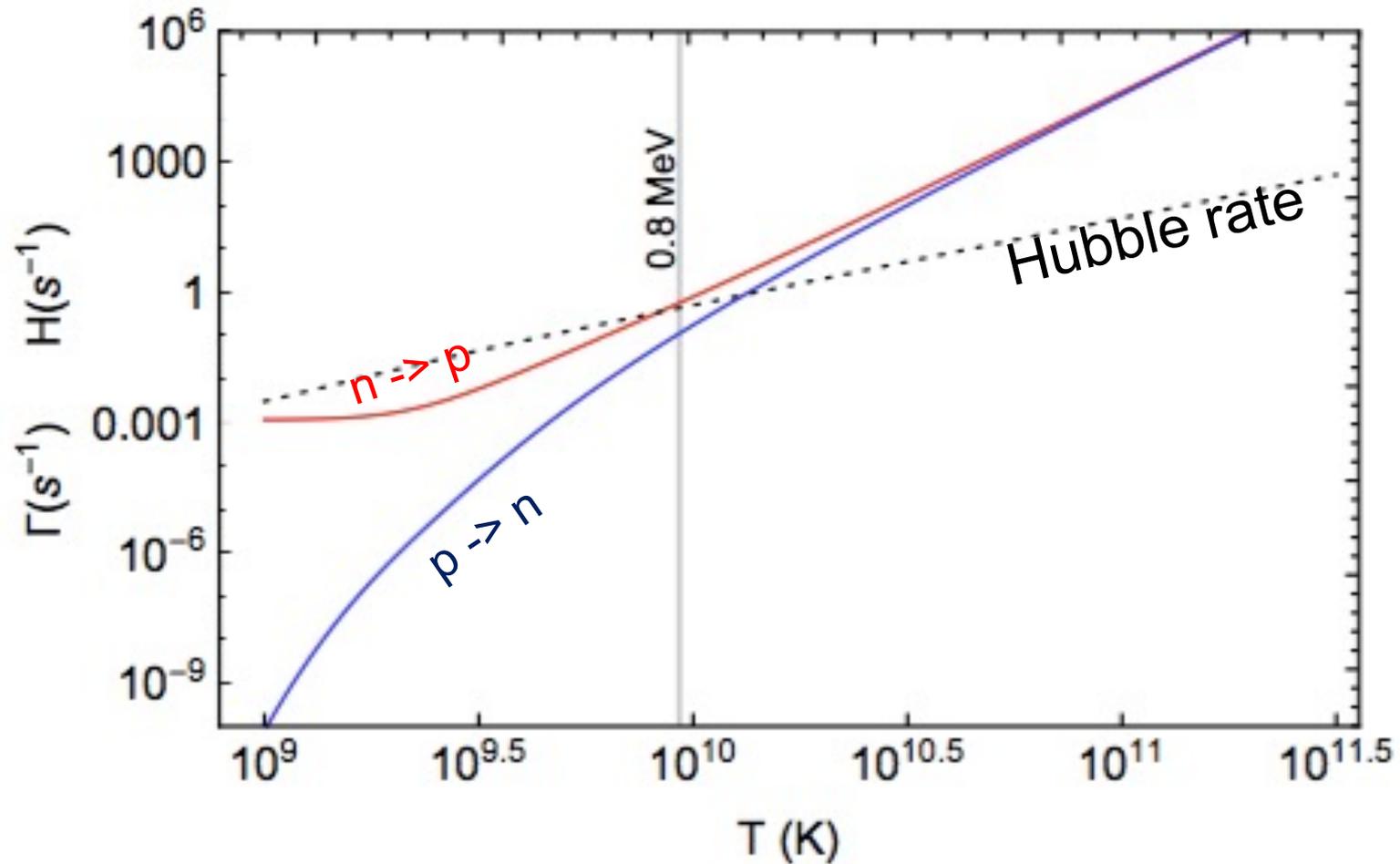
$$\tau_n = 878.4 \pm 0.5 \text{ s}$$

# Particle Data Group on neutron lifetime

WEIGHTED AVERAGE  
 $878.4 \pm 0.5$  (Error scaled by 1.8)

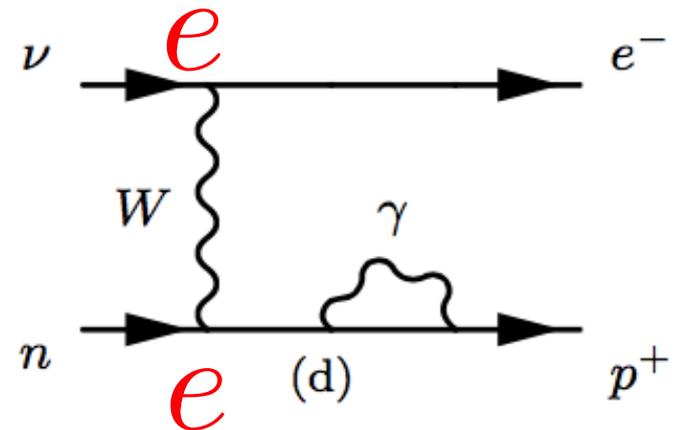
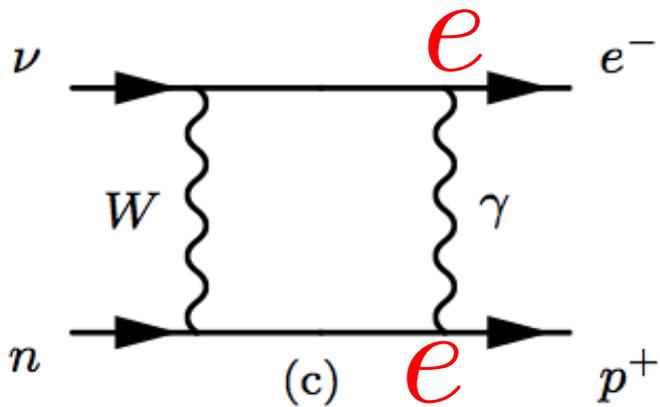
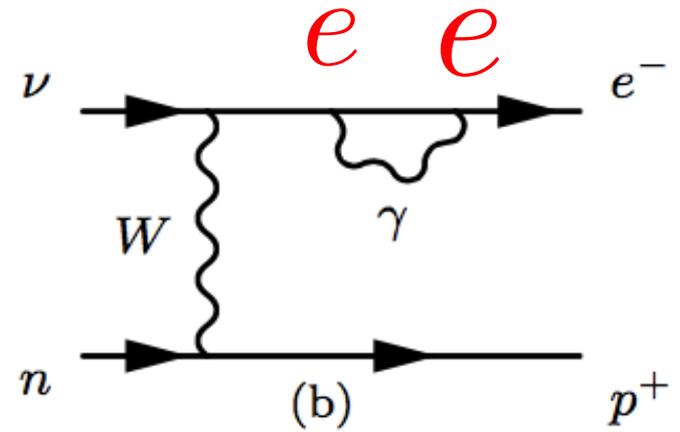
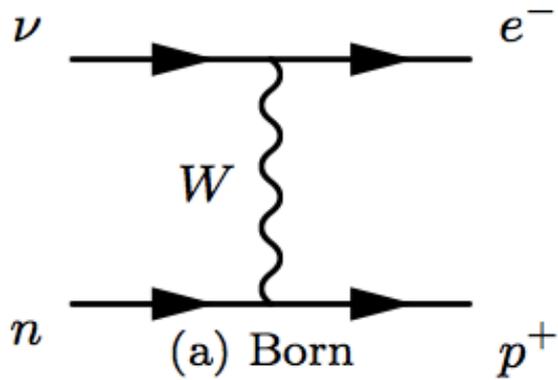


□ BORN approximation rates vs Hubble rate

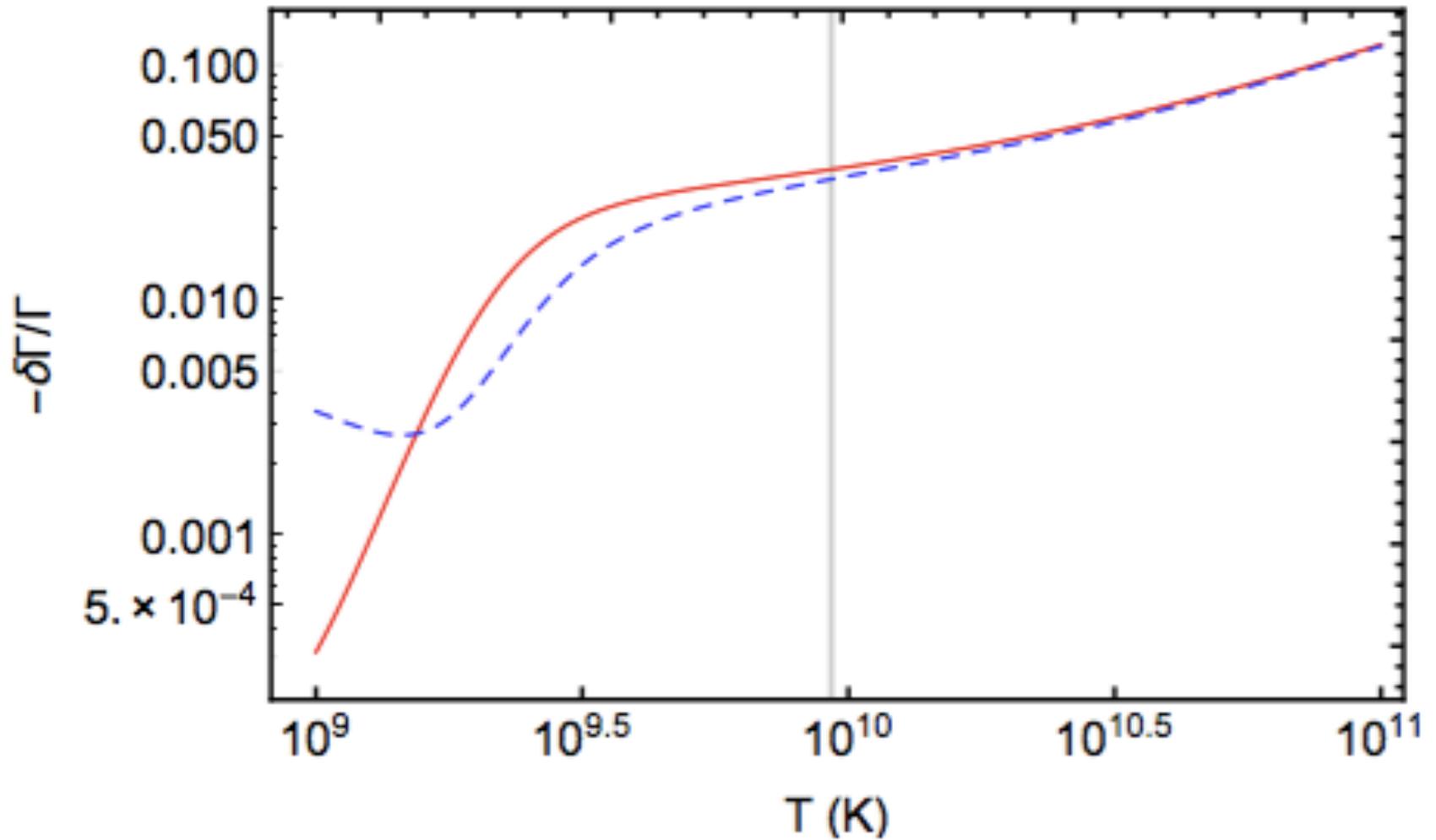


# □ Radiative corrections

$$\frac{e^2}{4\pi} = \alpha_{\text{FS}} \simeq \frac{1}{137}$$



□ Total corrections



# Hunt for precision

- Plasma physics
- Neutrino physics
- Weak rates
- Nuclear rates

## □ Evolution of abundances

$$\frac{dn_i}{dt} + 3Hn_i = \mathcal{J}_i \longrightarrow \text{Source from nuclear reactions}$$

$$\frac{dn_b}{dt} + 3Hn_b = 0, \quad \text{Baryons are only diluted}$$

$$Y_i \equiv n_i/n_b \longrightarrow \text{Removes dilution}$$

- Two-body reactions of the type  $i + j \leftrightarrow k + l$

$$\mathcal{J}_i \supset n_k n_l \gamma_{kl \rightarrow ij} - n_i n_j \gamma_{ij \rightarrow kl} \quad \gamma_{ij \rightarrow kl} \equiv \langle \sigma v \rangle_{ij \rightarrow kl}.$$

Average of cross-section over  
Maxwell-Boltzmann distribution

- General form

$$\dot{Y}_{i_1} = \sum_{i_2 \dots i_p, j_1 \dots j_q} N_{i_1} \left( \Gamma_{j_1 \dots j_q \rightarrow i_1 \dots i_p} \frac{Y_{j_1}^{N_{j_1}} \dots Y_{j_q}^{N_{j_q}}}{N_{j_1}! \dots N_{j_q}!} - \Gamma_{i_1 \dots i_p \rightarrow j_1 \dots j_q} \frac{Y_{i_1}^{N_{i_1}} \dots Y_{i_p}^{N_{i_p}}}{N_{i_1}! \dots N_{i_p}!} \right)$$

At nuclear statistical equilibrium (NSE), the densities of nuclei satisfy

$$n_i^{\text{NSE}} = \frac{g_i m_i^{3/2}}{2^{A_i}} \left( \frac{n_p}{m_p^{3/2}} \right)^{Z_i} \left( \frac{n_n}{m_n^{3/2}} \right)^{A_i - Z_i} \left( \frac{2\pi}{T} \right)^{\frac{3(A_i - 1)}{2}} e^{B_i/T}$$

Number protons

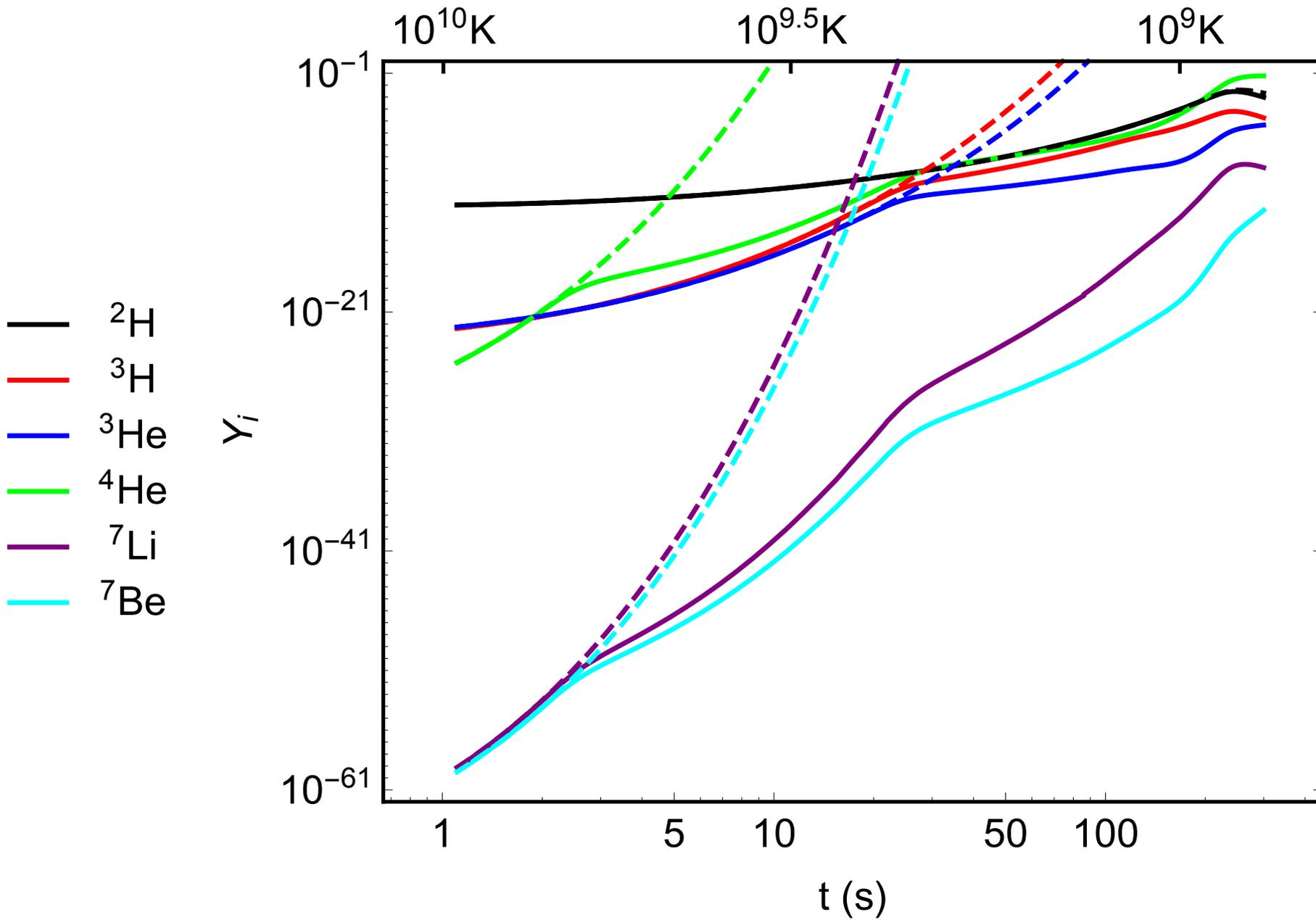
Number neutrons

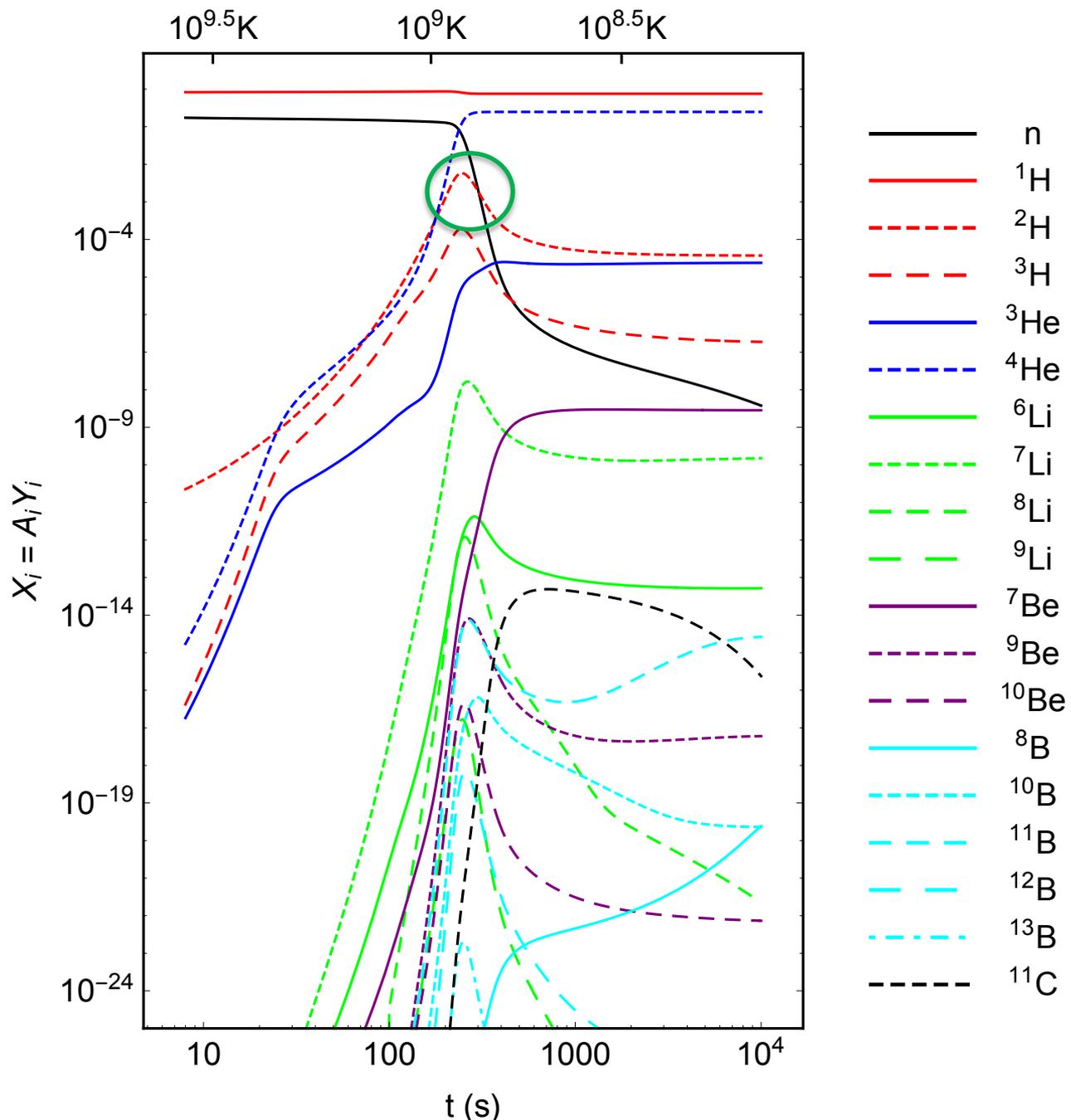
Binding energy

$$B_i \equiv Z_i m_p + (A_i - Z_i) m_n - m_i$$

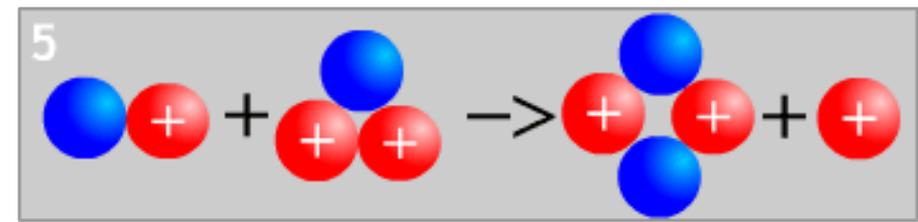
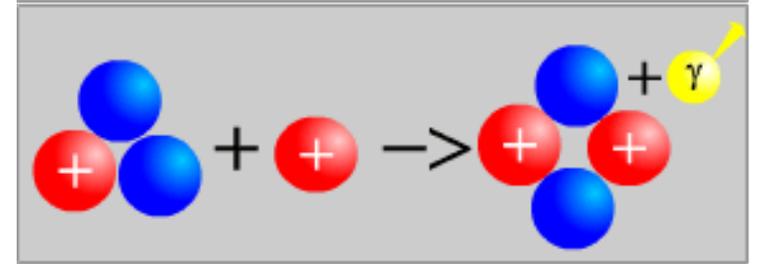
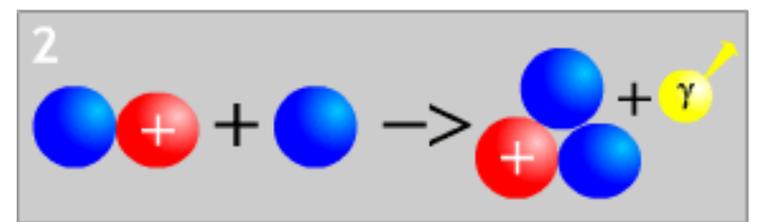
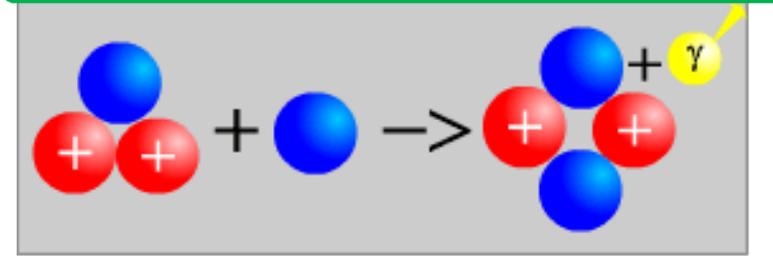
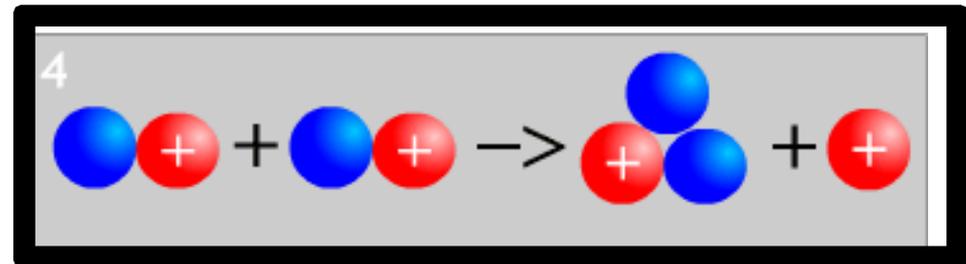
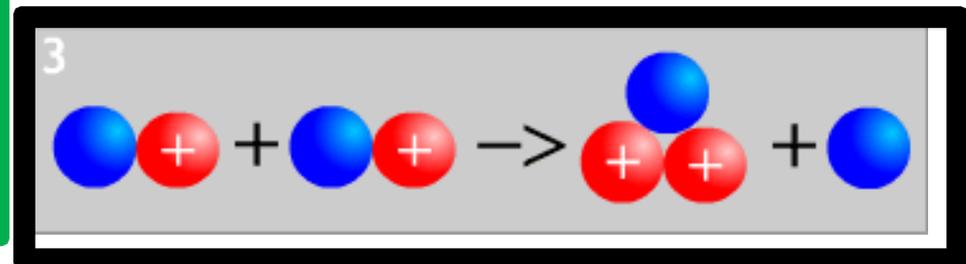
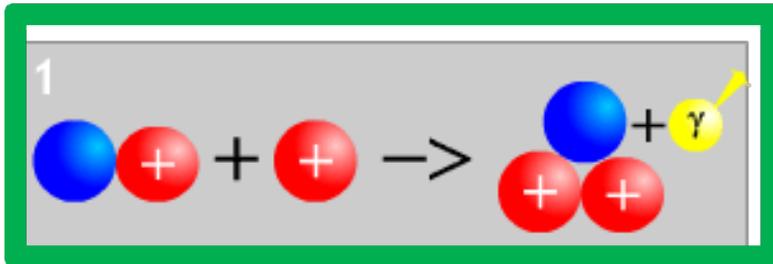
Demo tools :

- Non-relativistic number density  $n_i = g_i \left( \frac{m_i T}{2\pi} \right)^{3/2} e^{(\mu_i - m_i)/T}$
- Chemical equilibrium  $\mu_i = Z_i \mu_p + (A_i - Z_i) \mu_n$





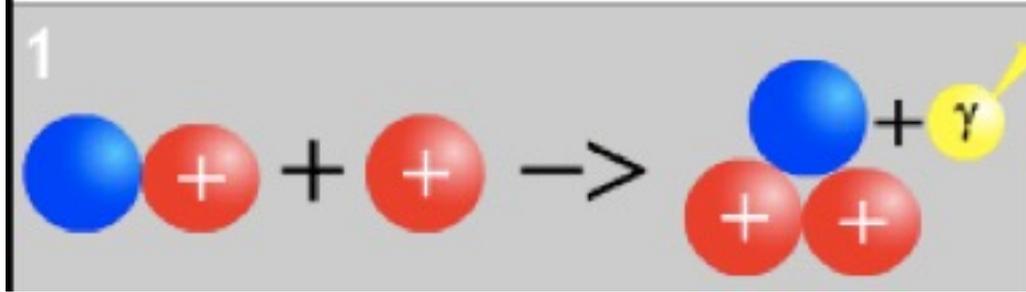
# Main reactions for D



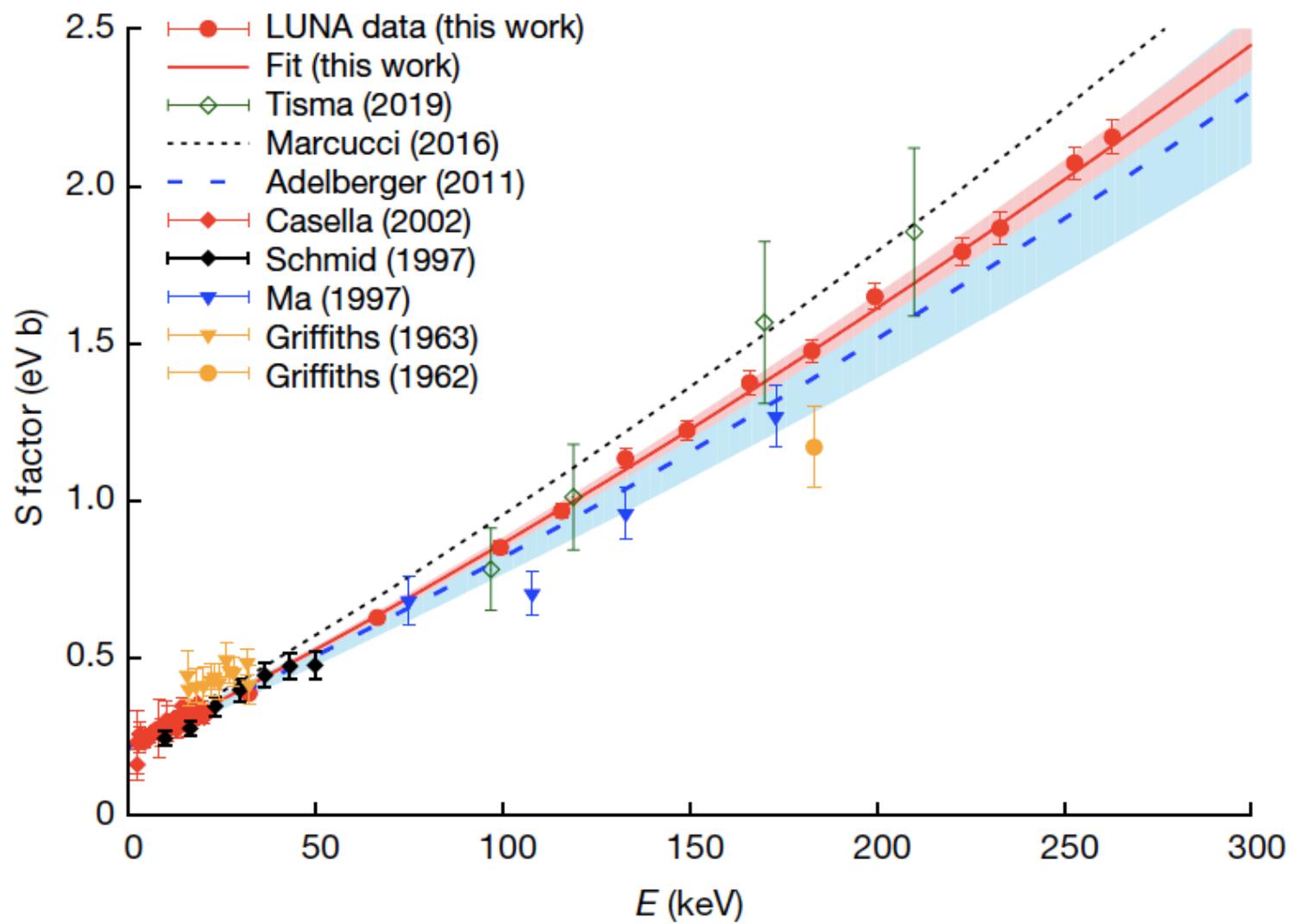
## Final deuterium sensitivity

$$\frac{\Delta(D/H)}{D/H} = -0.32 \frac{\Delta\langle\sigma v\rangle_{D(p,\gamma)^3\text{He}}}{\langle\sigma v\rangle_{D(p,\gamma)^3\text{He}}}$$

$$\frac{\Delta(D/H)}{D/H} = -0.54 \frac{\Delta\langle\sigma v\rangle_{D(d,n)^3\text{He}}}{\langle\sigma v\rangle_{D(d,n)^3\text{He}}} - 0.46 \frac{\Delta\langle\sigma v\rangle_{D(d,p)^3\text{H}}}{\langle\sigma v\rangle_{D(d,p)^3\text{H}}}$$

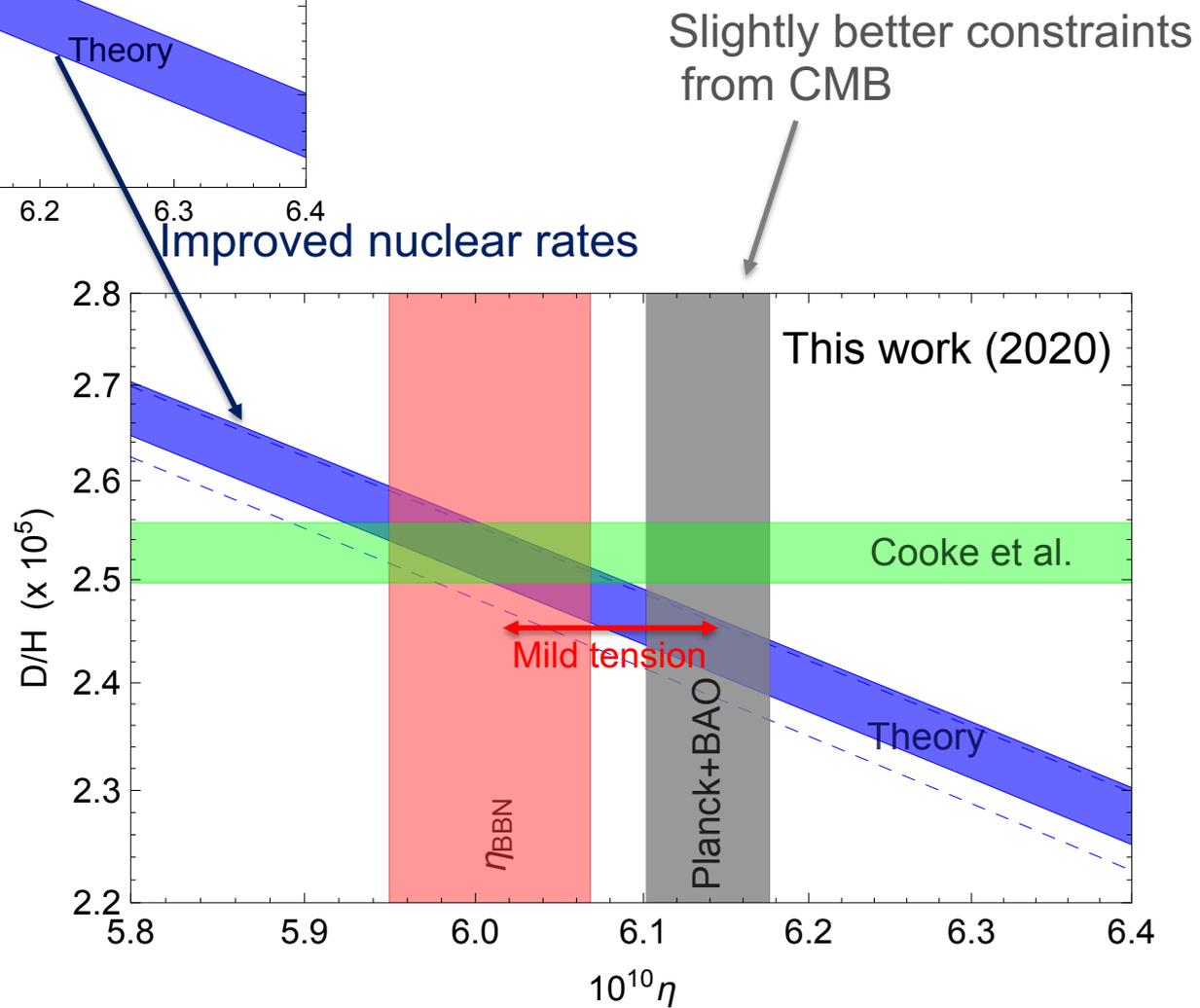
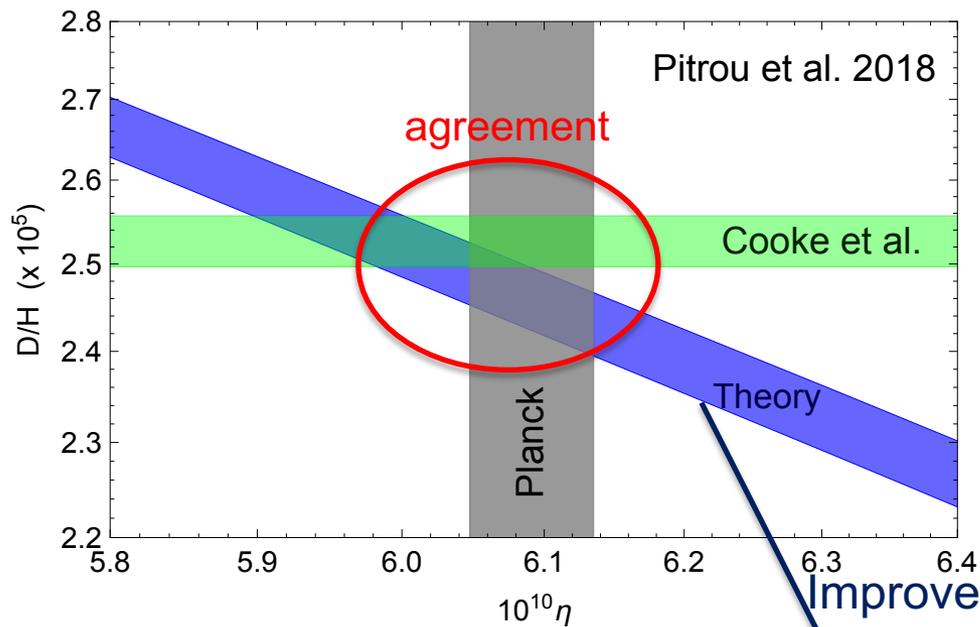


2020



# Outline

1. BBN in a nutshell
2. The Hunt for precision
3. **Cosmological parameters**



# A new tension in the cosmological model from primordial deuterium?

Cyril Pitrou,<sup>1\*</sup> Alain Coc,<sup>2</sup> Jean-Philippe Uzan,<sup>1</sup> Elisabeth Vangioni<sup>1</sup>

<sup>1</sup>*Institut d'Astrophysique de Paris, CNRS UMR 7095, 98 bis Bd Arago, 75014 Paris, France*

*Sorbonne Université, Institut Lagrange de Paris, 98 bis Bd Arago, 75014 Paris, France*

<sup>2</sup>*IJCLab, CNRS IN2P3, Université Paris-Saclay, Bâtiment 104, F-91405 Orsay Campus France*

---

## Primordial Deuterium after LUNA: concordances and error budget

**Ofelia Pisanti, Gianpiero Mangano, Gennaro Miele, and Pierpaolo Mazzella**

Dipartimento di Fisica E. Pancini, Università di Napoli Federico II, and INFN, Sezione di Napoli, Via Cintia, I-80126 Napoli, Italy

after the LUNA results, the value of Deuterium is quite precisely fixed, and points to a value of the baryon density in excellent agreement with the Planck result,

# The Impact of New $d(p, \gamma)^3\text{He}$ Rates on Big Bang Nucleosynthesis

Tsung-Han Yeh

*Department of Physics, University of Illinois, Urbana, IL 61801*

Keith A. Olive

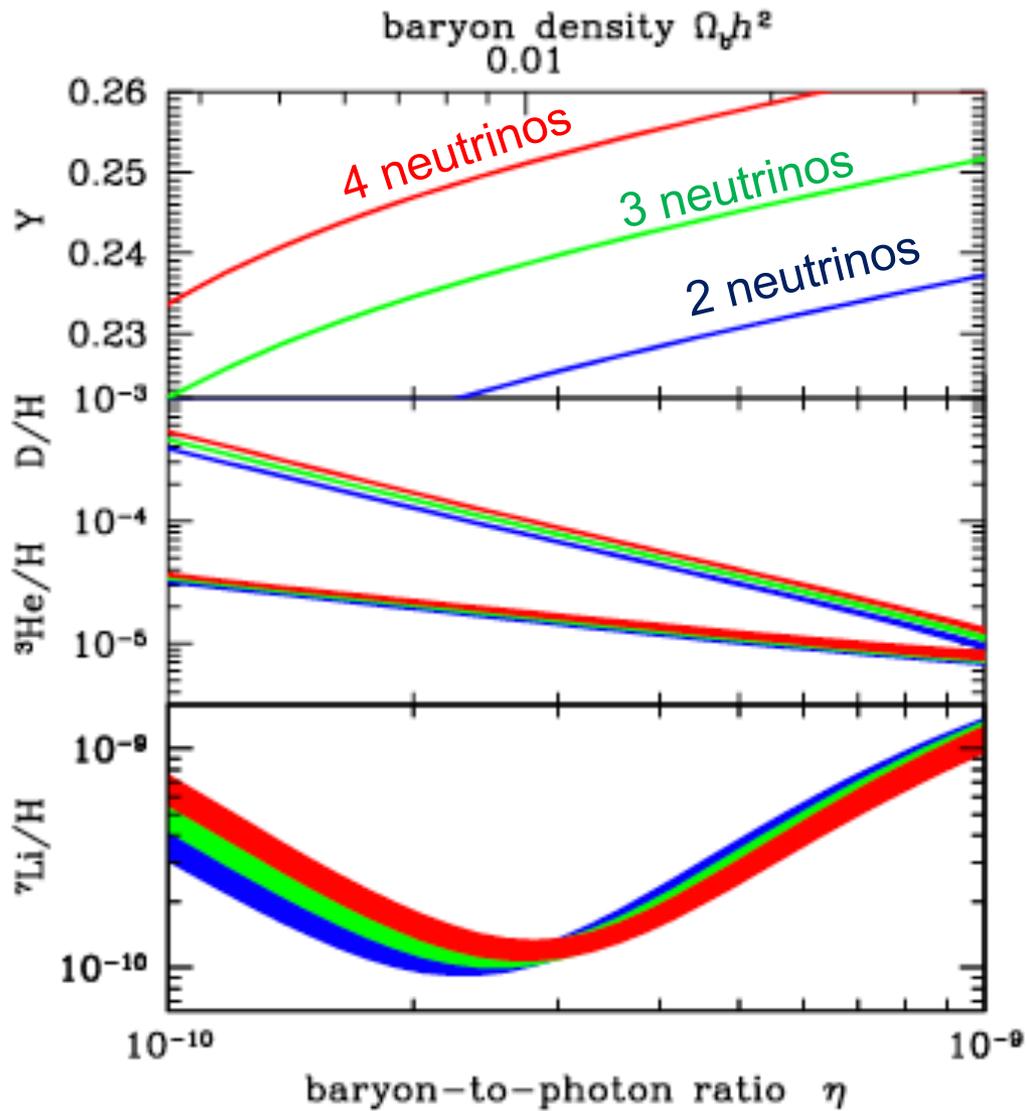
*William I. Fine Theoretical Physics Institute, School of Physics and Astronomy,  
University of Minnesota, Minneapolis, MN 55455, USA*

Brian D. Fields

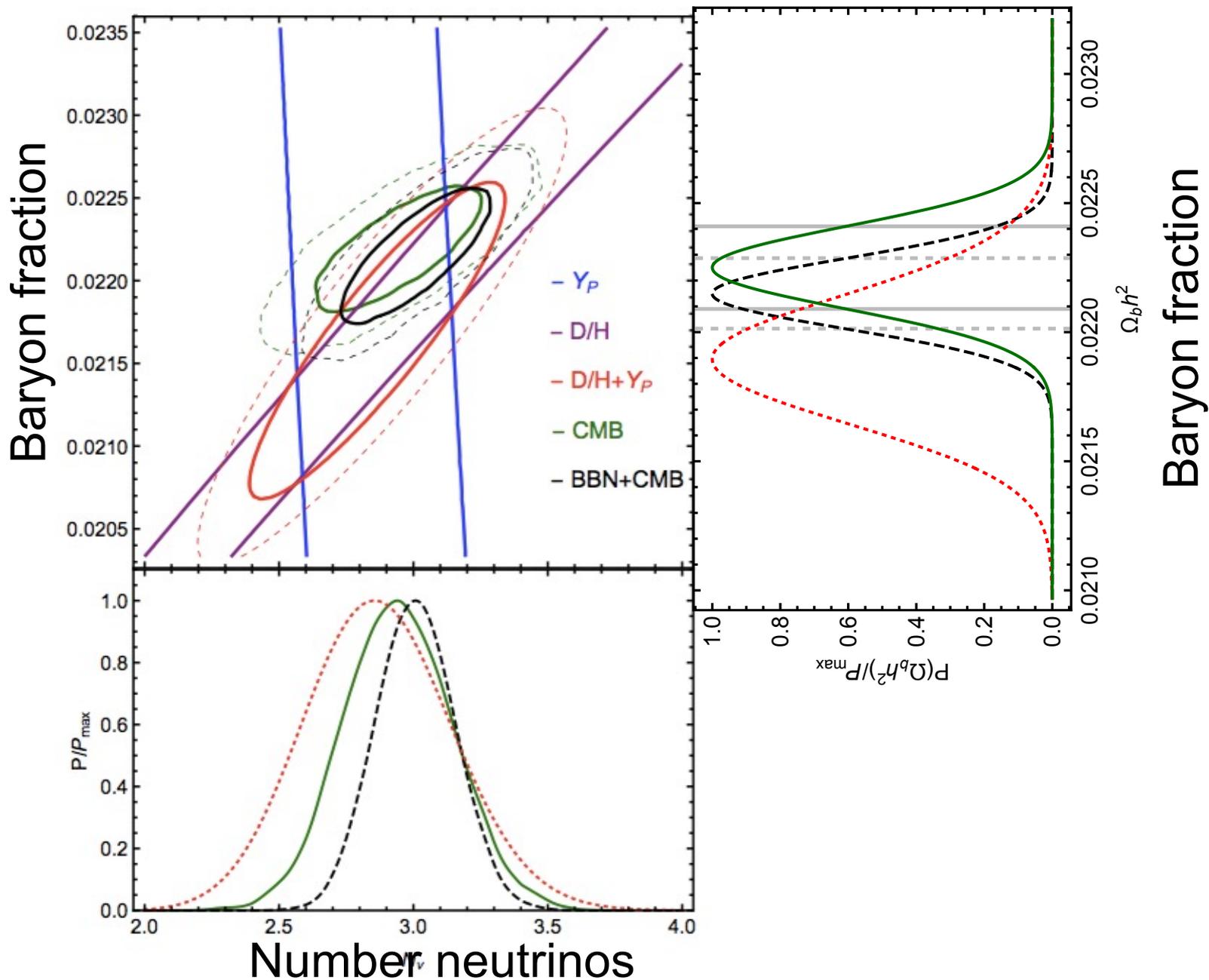
*Departments of Astronomy and of Physics,  
University of Illinois, Urbana, IL 61801*

Finally, we note that the observed deuterium abundance continues to be more precise than the BBN+CMB prediction, whose error budget is now dominated by  $d(d, n)^3\text{He}$  and  $d(d, p)^3\text{H}$ .

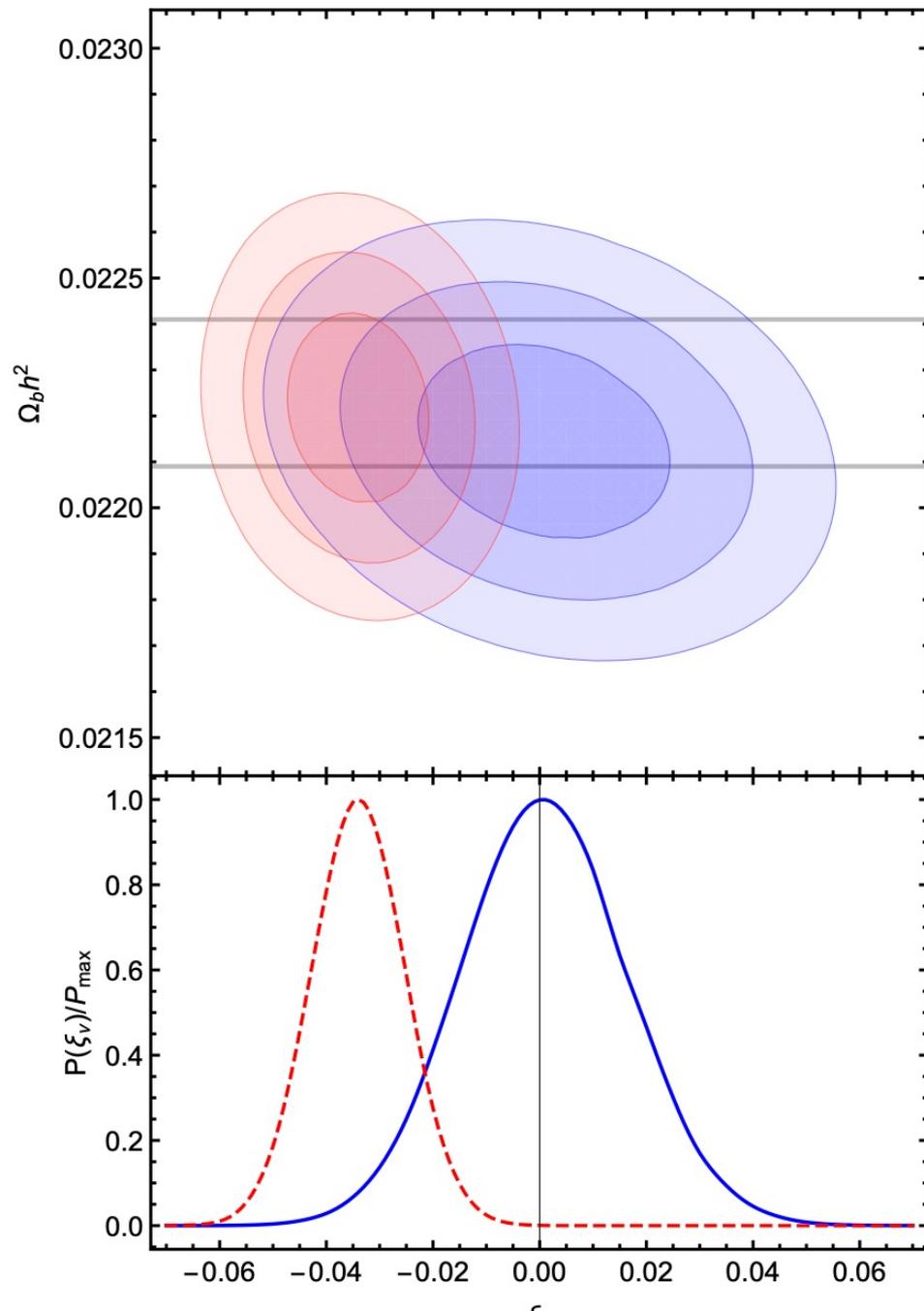
□ Dependence on number of neutrino species



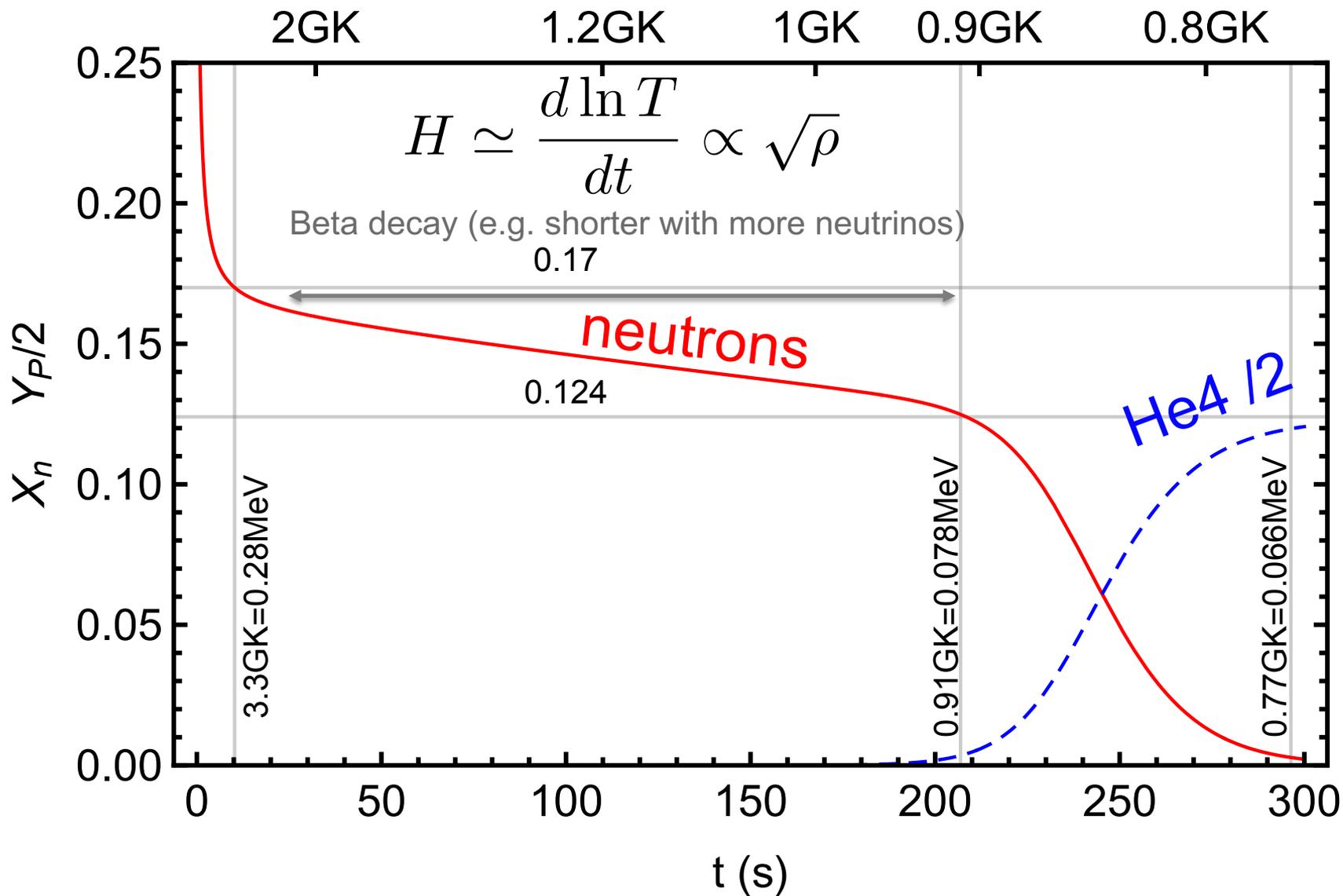
Cyburt et al. 2015

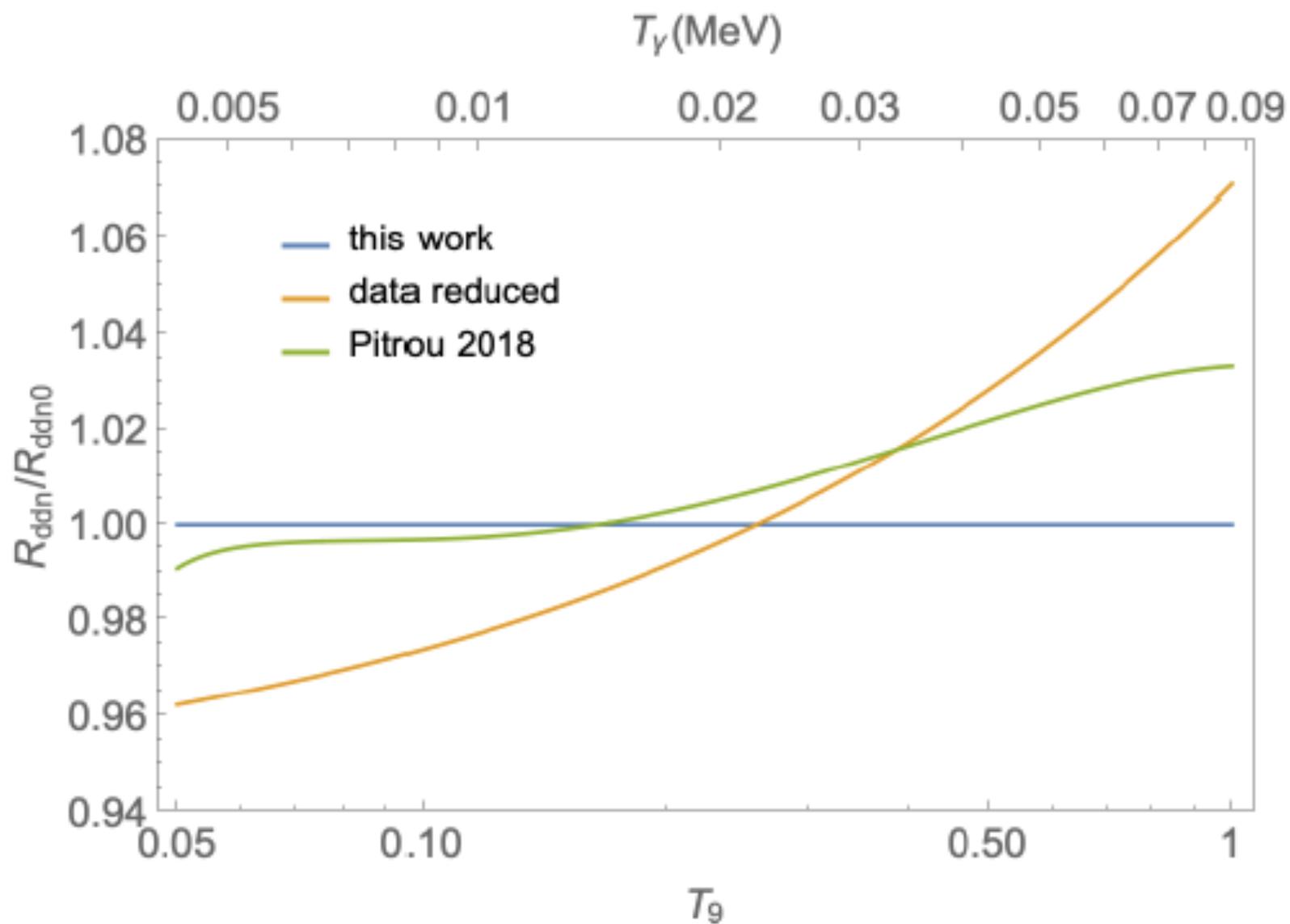


Thank you



# Effect on beta decay





$T_\gamma$ (MeV)

0.005

0.01

0.02

0.03

0.05

0.07

0.09

1.08

1.06

1.04

1.02

1.00

0.98

0.96

0.94

$R_{\text{ddp}}/R_{\text{ddp}0}$

- this work
- data reduced
- Pitrou 2018

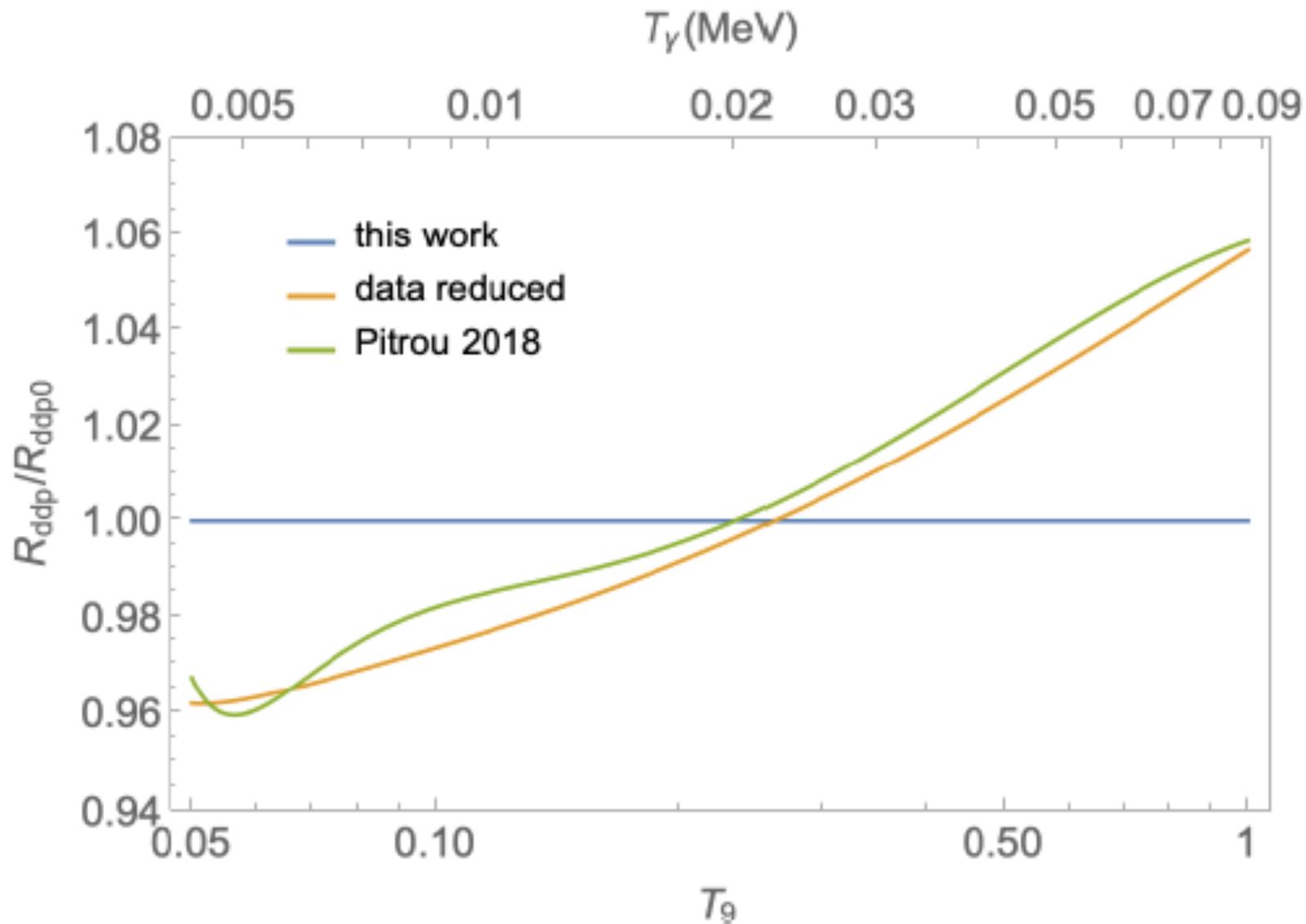
0.05

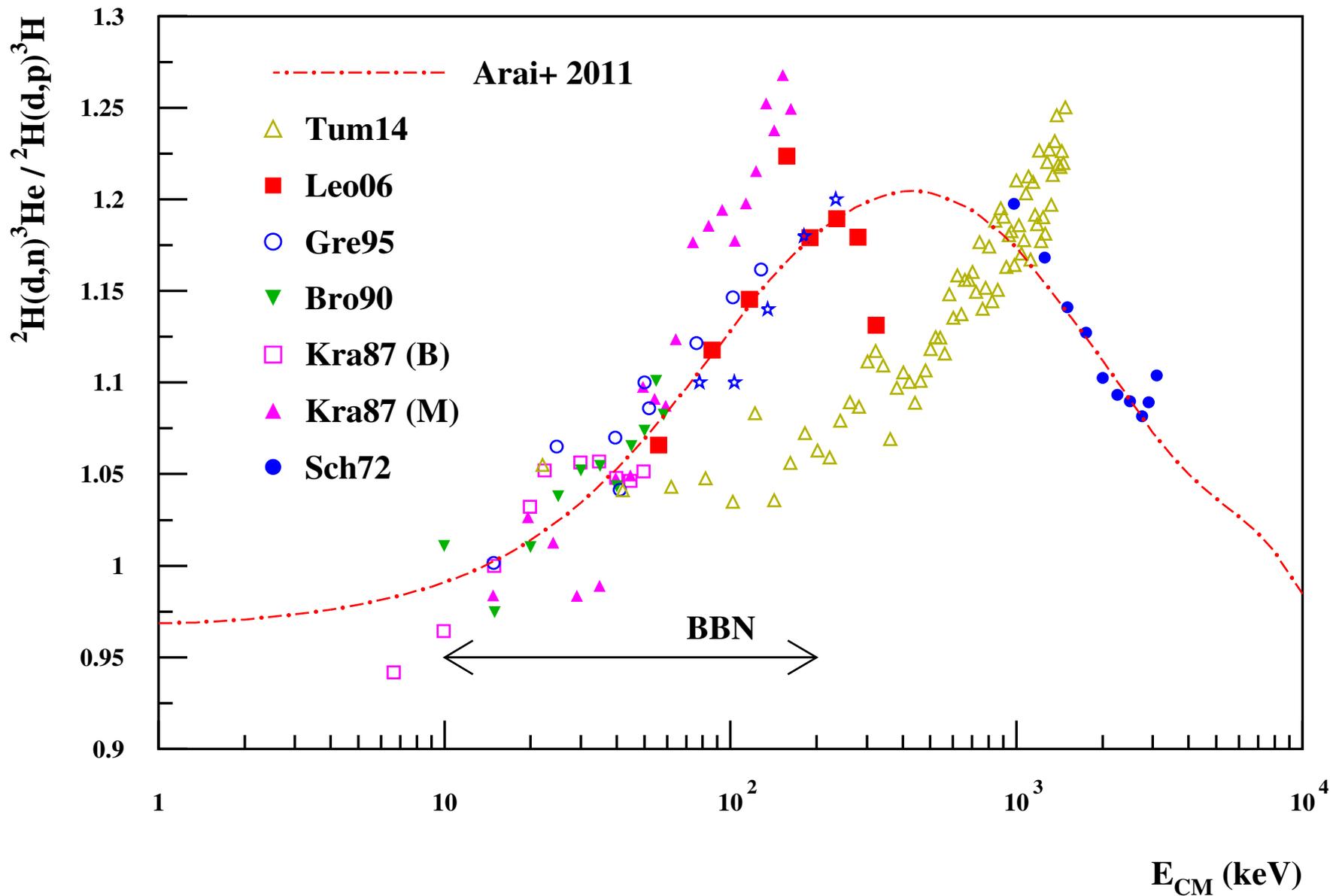
0.10

0.50

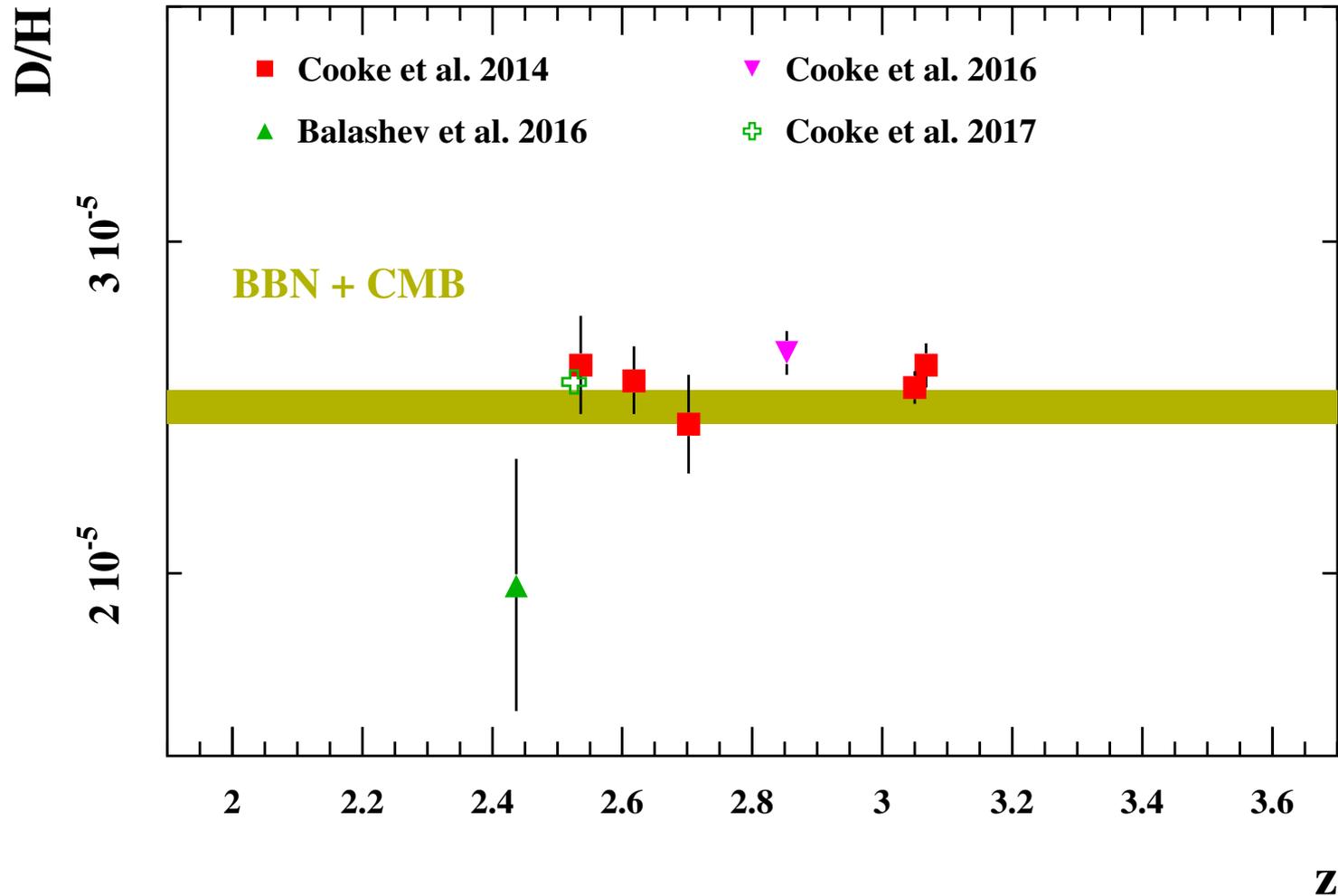
1

$T_9$





# Quasar absorption lines



# ☐ Helium production

