

# Nuisances in the fit: the nuclear model uncertainties in the new era of Spectral Function

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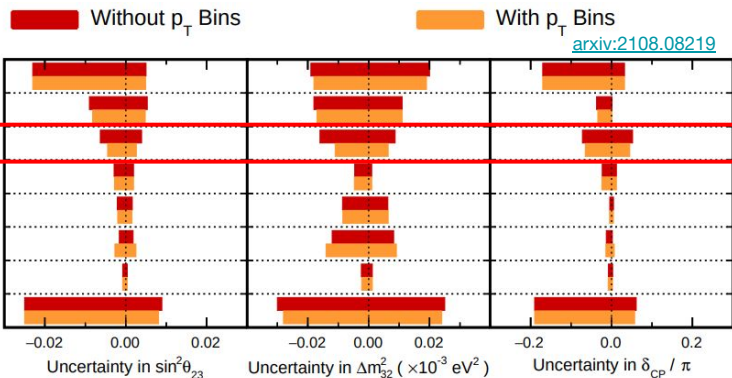
IRN Neutrino meeting  
LAPP, Annecy  
June 30<sup>th</sup>, 2022

1. Introduction
2. Neutrino-nucleus interaction models
3. Proposed systematic uncertainties
4. Validations of the systematic uncertainties with fits to cross section measurements
5. Projected constraints with T2K ND Upgrade
6. Summary and prospects

# 1. Introduction

# Neutrino oscillations: the precision era

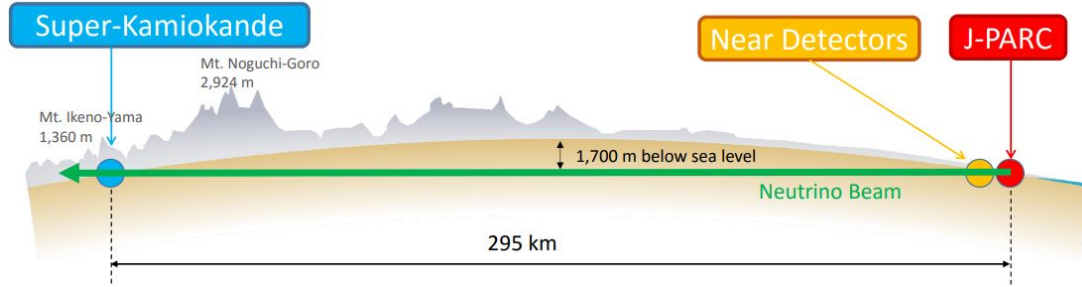
- First hints of CP violation
- Currently: NOvA, T2K (w/ ND upgrade)
- Future: DUNE, Hyper-Kamiokande
- Still limited by statistics, but not for long!



Phys.Rev.D 103, 112008 (2021)

Error source	1-Ring $\mu$		1-Ring $e$			
	FHC	RHC	FHC	RHC	FHC/ 1 d.e.	FHC/RHC
SK Detector	2.4	2.0	2.8	3.8	13.2	1.5
SK FSI+SI+PN	2.2	2.0	3.0	2.3	11.4	1.6
Flux + Xsec (ND constrained)	3.3	2.9	3.2	3.1	4.1	2.7
Nucleon Removal Energy	2.4	1.7	7.1	3.7	3.0	3.6
$\sigma(\nu_e)/\sigma(\bar{\nu}_e)$	0.0	0.0	2.6	1.5	2.6	3.0
NC1 $\gamma$	0.0	0.0	1.1	2.6	0.3	1.5
NC Other	0.3	0.3	0.2	0.3	1.0	0.2
$\sin^2 \theta_{23} + \Delta m^2_{21}$	0.0	0.0	0.5	0.3	0.5	2.0
$\sin^2 \theta_{13}$ PDG2018	0.0	0.0	2.6	2.4	2.6	1.1
All Systematics	5.1	4.5	8.8	7.1	18.4	6.0

⇒ Neutrino interaction uncertainties must be reduced!

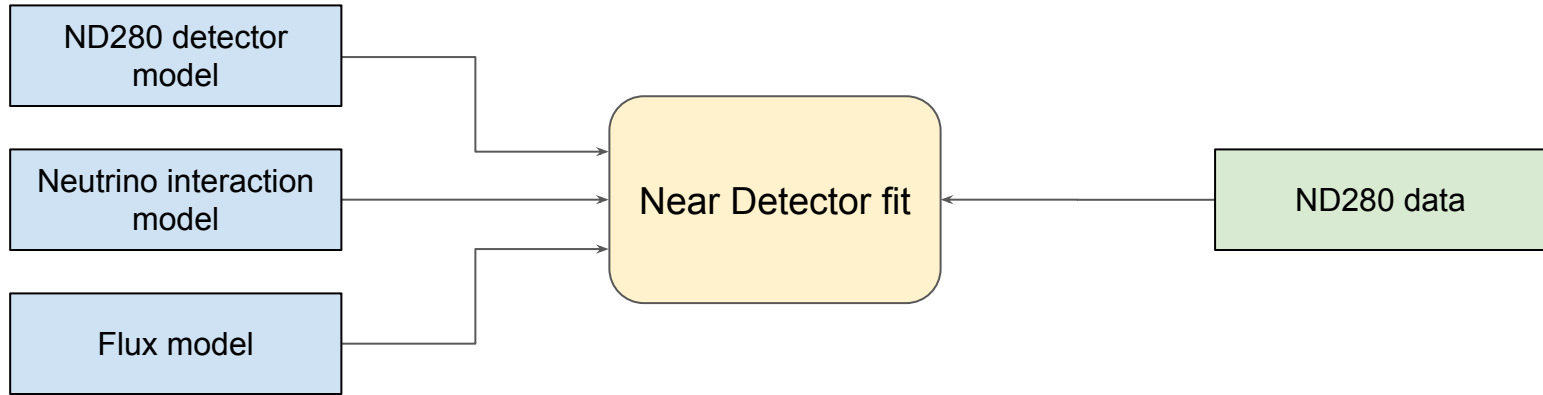


$$\begin{aligned}
 N_{\nu_\alpha}^{ND}(E_\nu) &= \Phi_{\nu_\alpha}^{ND}(E_\nu) \times \epsilon^{ND}(E_\nu) \times \sigma_{\nu_\alpha}^{ND}(E_\nu) \\
 N_{\nu_\beta}^{FD}(E_\nu) &= \Phi_{\nu_\beta}^{FD}(E_\nu) \times \epsilon^{FD}(E_\nu) \times \sigma_{\nu_\beta}^{FD}(E_\nu) \times P_{\nu_\alpha \rightarrow \nu_\beta}(E_\nu)
 \end{aligned}$$

Flux model
Detector model
Neutrino interaction model

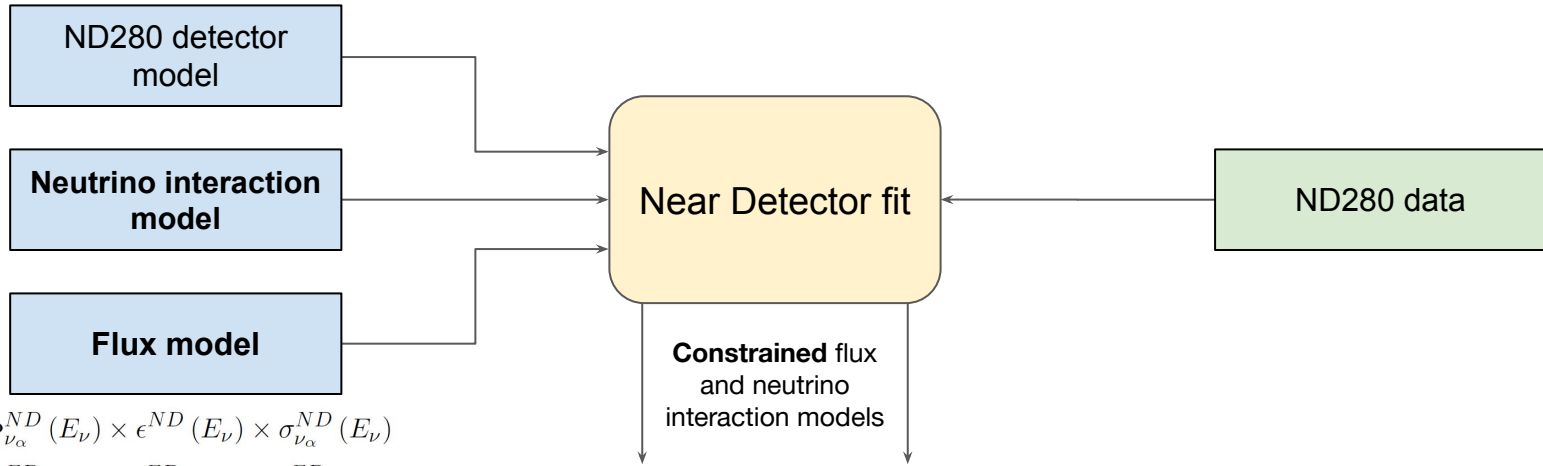
● **Far/Near ratio does not fully cancel systematic uncertainties:**

- Flux model different at ND vs. FD due to geometry and oscillation
- Different detectors, i.e. different acceptance, efficiencies, targets...
- Mainly muon neutrinos at ND interacting with CH → use model to infer interactions with electron neutrino interactions and with H<sub>2</sub>O



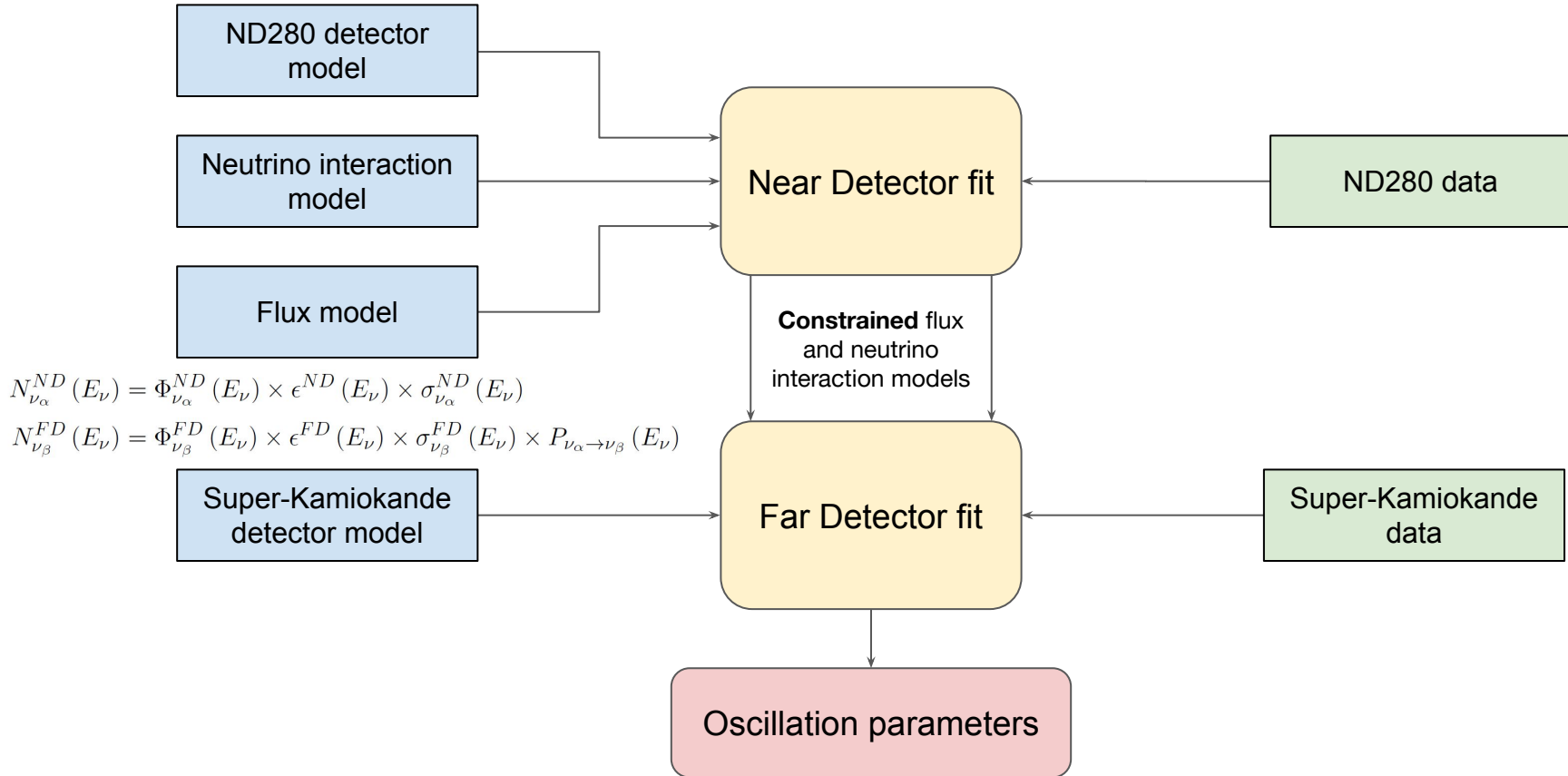
$$N_{\nu_\alpha}^{ND}(E_\nu) = \Phi_{\nu_\alpha}^{ND}(E_\nu) \times \epsilon^{ND}(E_\nu) \times \sigma_{\nu_\alpha}^{ND}(E_\nu)$$

$$N_{\nu_\beta}^{FD}(E_\nu) = \Phi_{\nu_\beta}^{FD}(E_\nu) \times \epsilon^{FD}(E_\nu) \times \sigma_{\nu_\beta}^{FD}(E_\nu) \times P_{\nu_\alpha \rightarrow \nu_\beta}(E_\nu)$$



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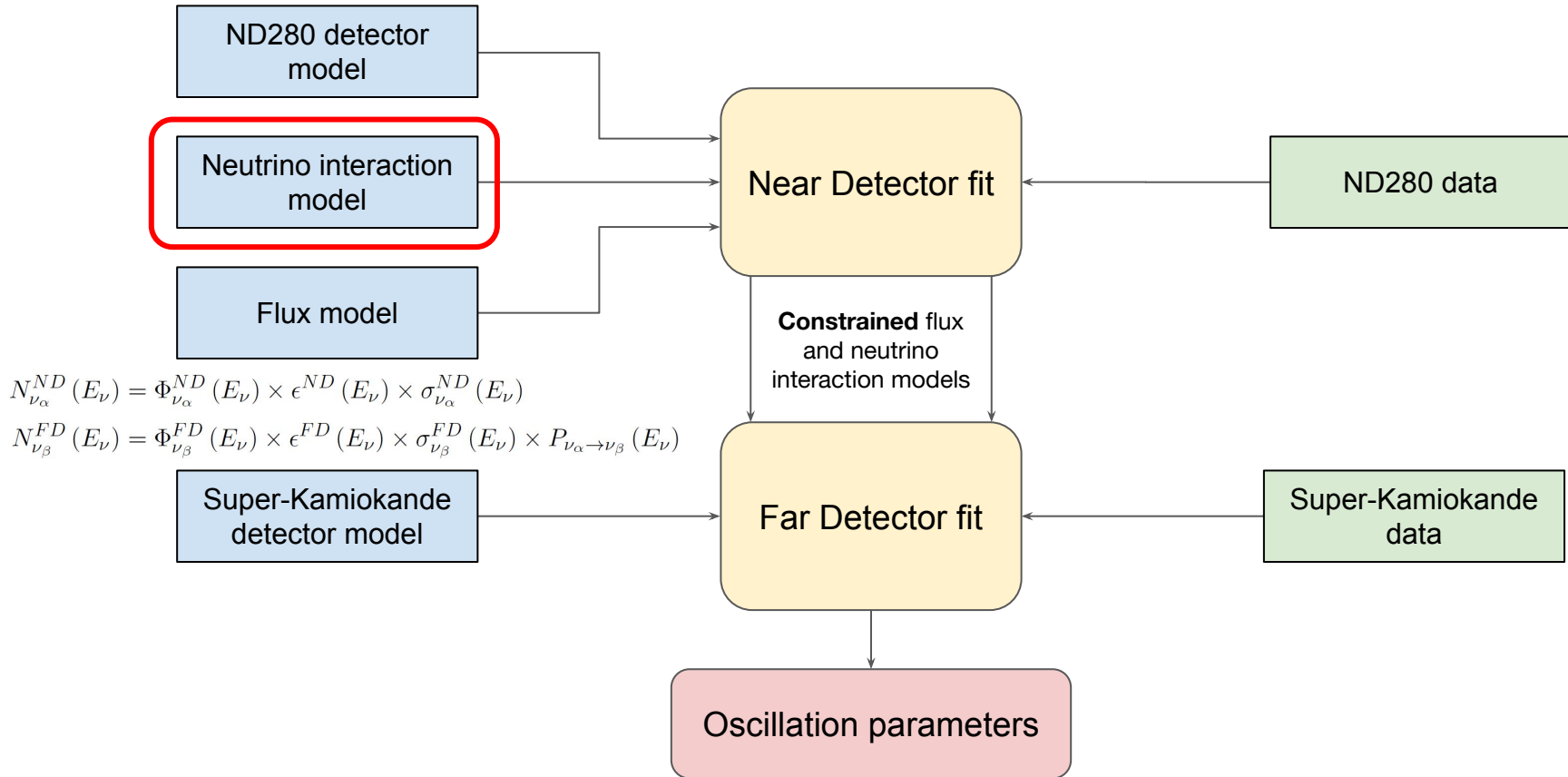
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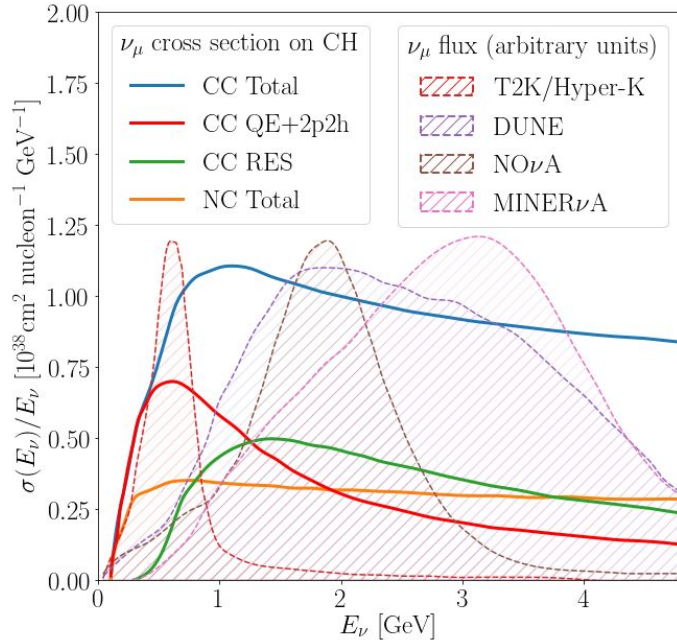




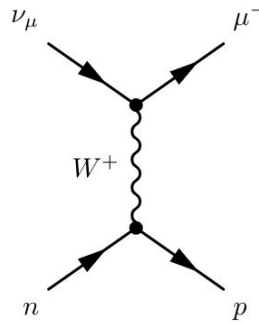


## 2. Neutrino-nucleus interaction models

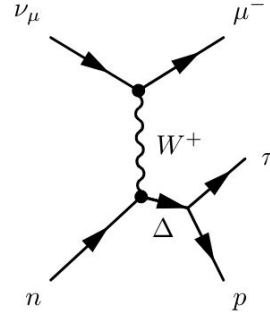




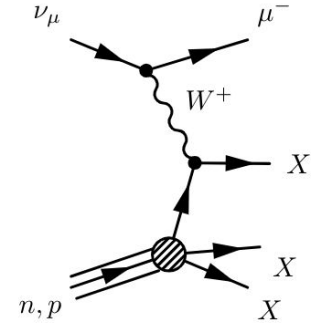
Quasi-elastic  
(CCQE)



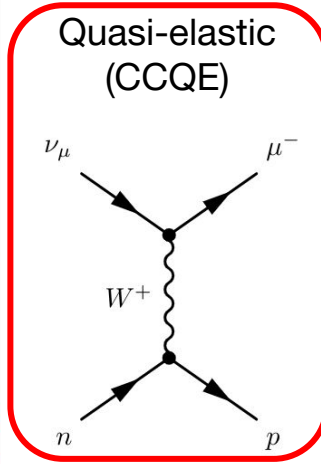
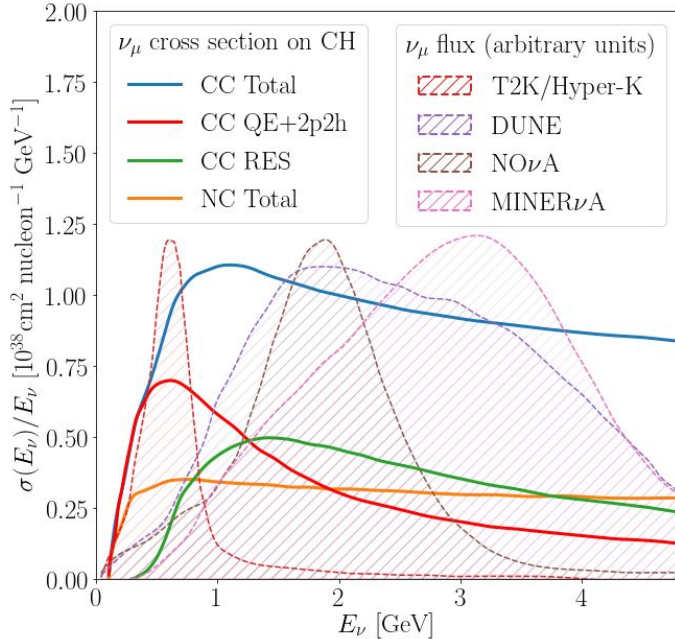
Resonant  
(CCRES)



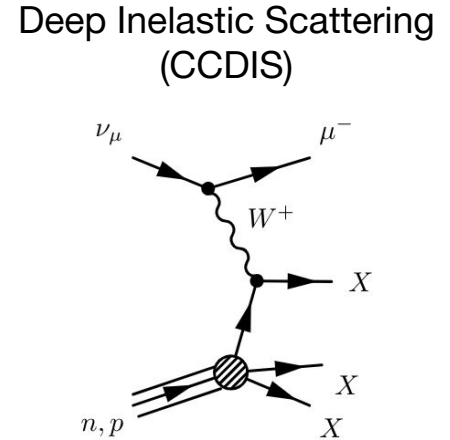
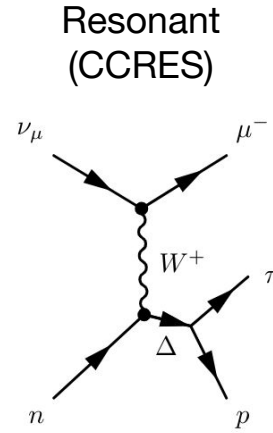
Deep Inelastic Scattering  
(CCDIS)



- In order to estimate neutrino energy, a good understanding of neutrino-nucleus interactions is necessary
- CCQE is the dominant interaction in T2K/Hyper-Kamiokande, and is a significant mode in NO $\nu$ A, MINER $\nu$ A and DUNE

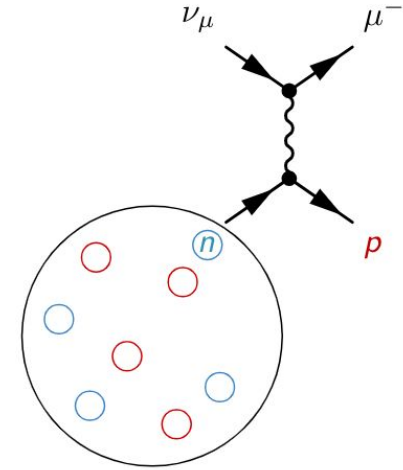


This talk



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- Neutrinos can interact with nucleons bound within nuclei (Carbon, Oxygen, Argon...)
- Initial state nucleons are non-static: Fermi motion
- How to model this?



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## Fermi gas

### Relativistic Fermi Gas (RFG)

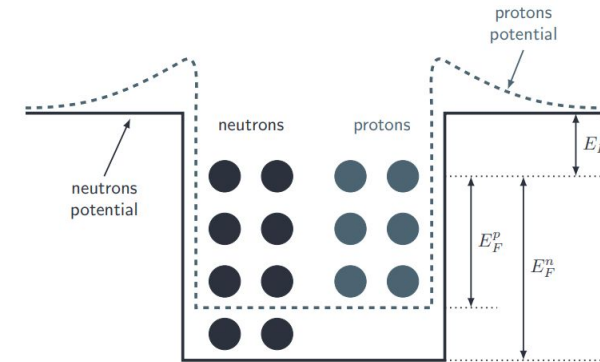
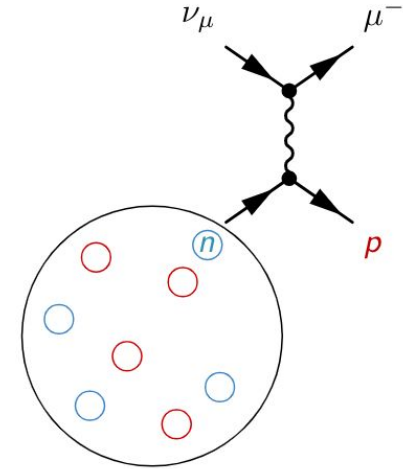
Nucleons move freely in a constant binding energy within the nuclear volume

$$p_F = \left(3\pi^2 \rho \frac{Z}{A}\right)^{1/3}$$

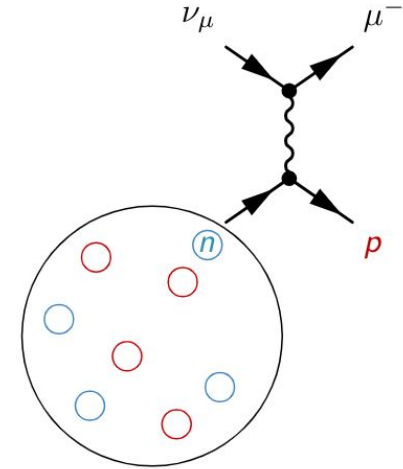
### Local Fermi Gas (LFG)

The nucleus is described with the local density approximation

$$p_F(r) = \left(3\pi^2 \rho(r) \frac{Z}{A}\right)^{1/3}$$



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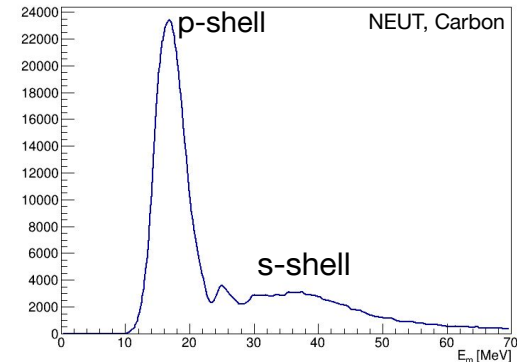


## Spectral Function (SF)

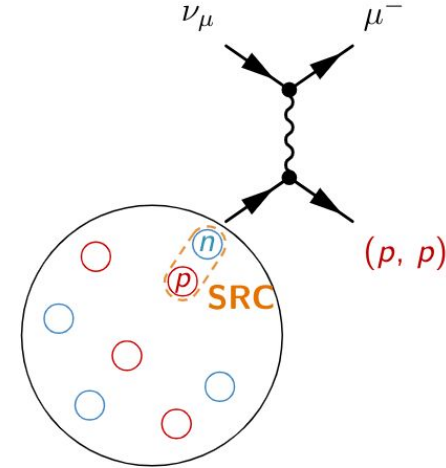
The probability of removing of a nucleon with momentum  $p_m$  and leaving residual nucleus with excitation energy  $E_m$

$$P(p_m, E_m) = P_{MF}(p_m, E_m) + P_{corr}(p_m, E_m)$$

Independent nucleons, moving in a mean-field potential within the shell-model picture → built from (e,e'p) data (~80%)  
→ **One outgoing nucleon is produced**



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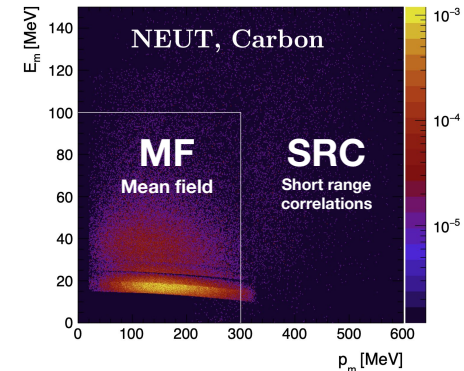
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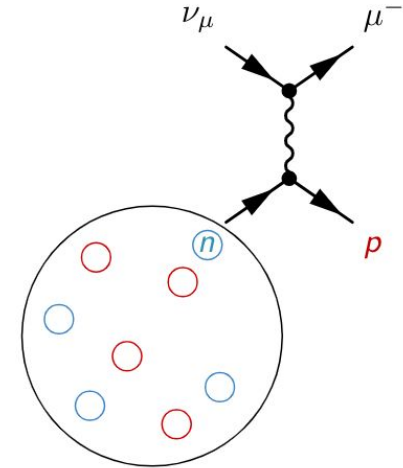
pairs of strongly-correlated nucleons (~20%)

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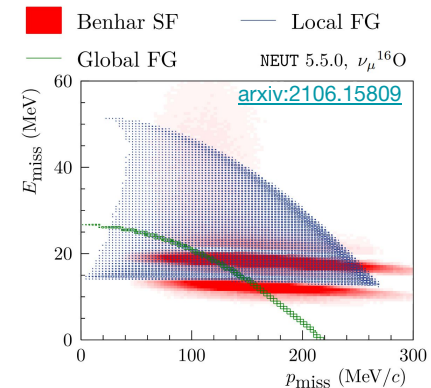
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This talk

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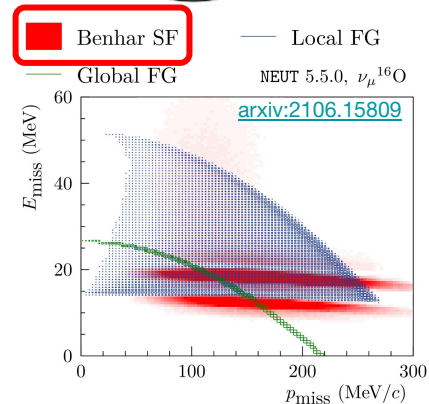
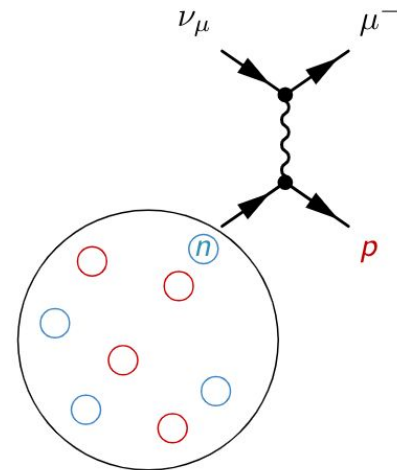
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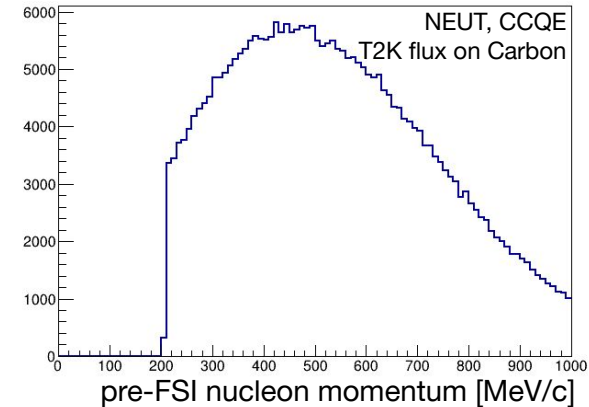
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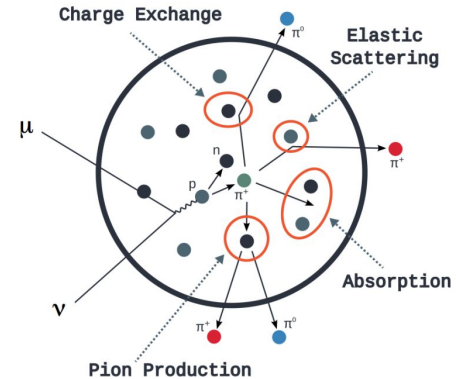
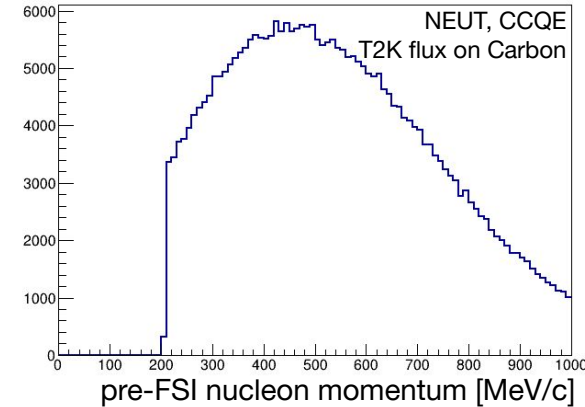
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- Pauli blocking (PB):
  - By Pauli principle, an interaction cannot occur if it leads to the creation of a nucleon in a state that is already occupied
  - Simple model: reject events with outgoing nucleon momentum below Fermi level

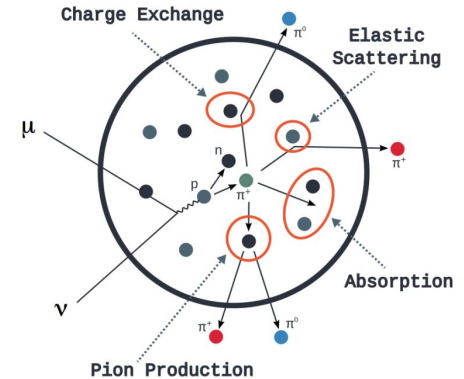
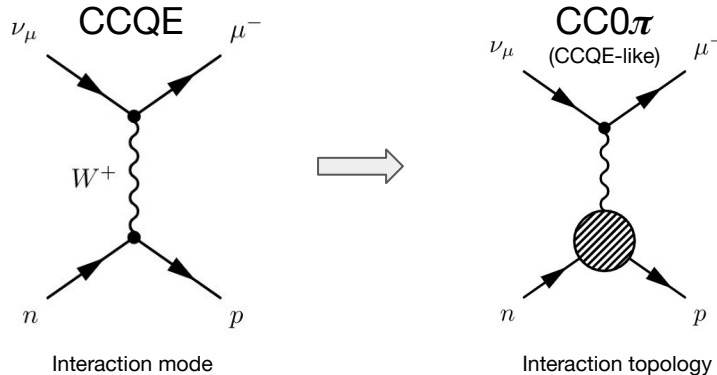
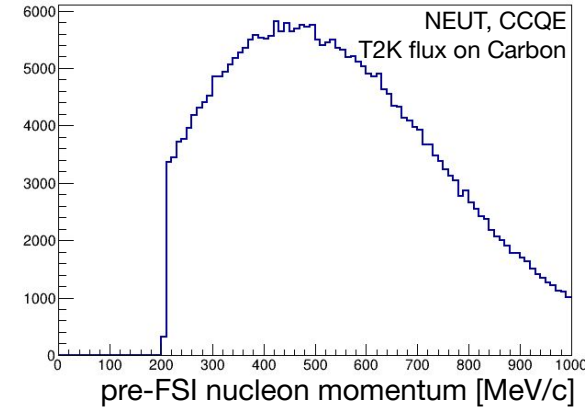


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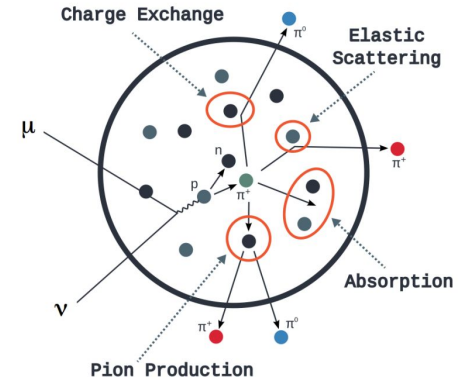
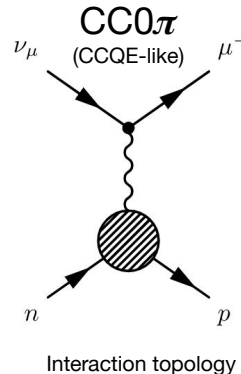
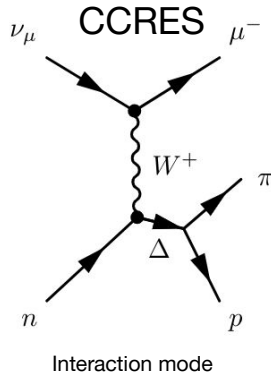
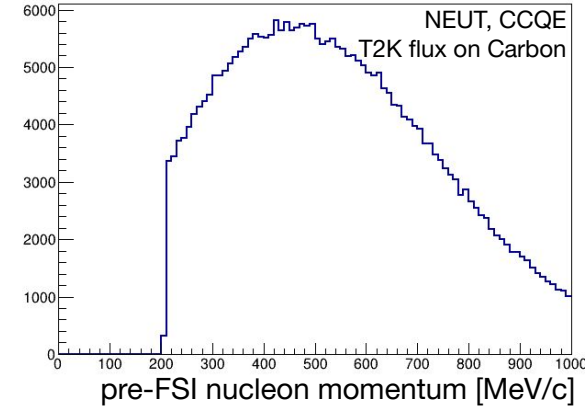
See Anna's presentation

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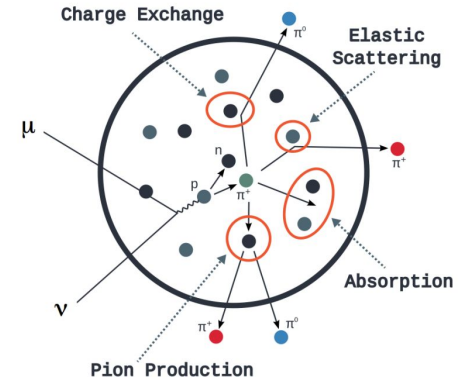
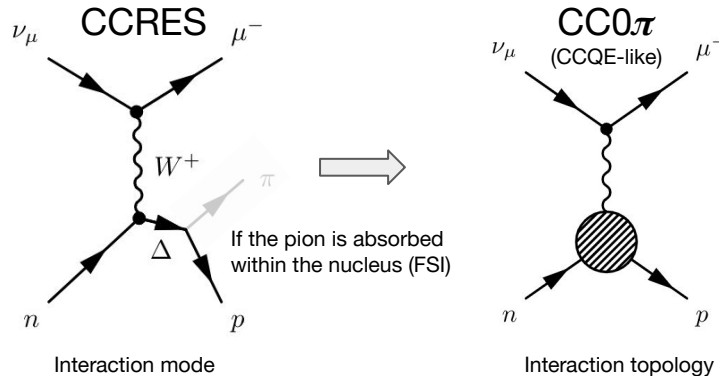
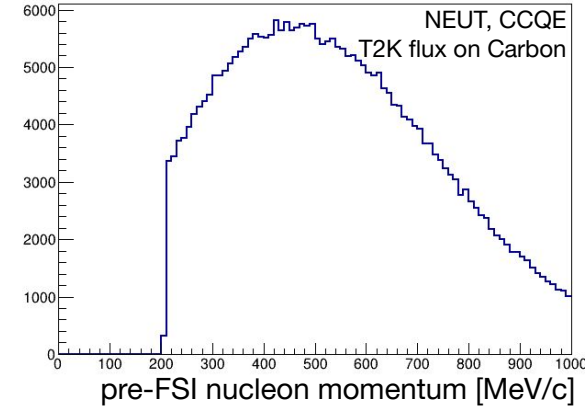
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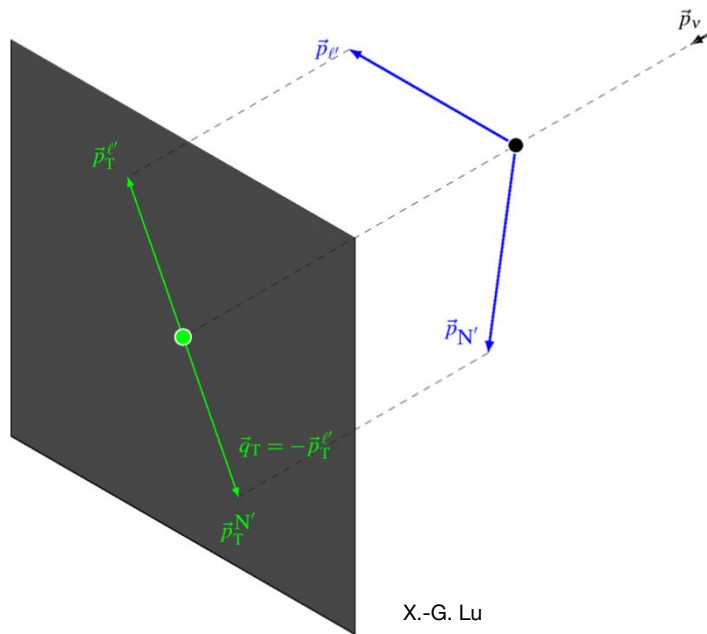
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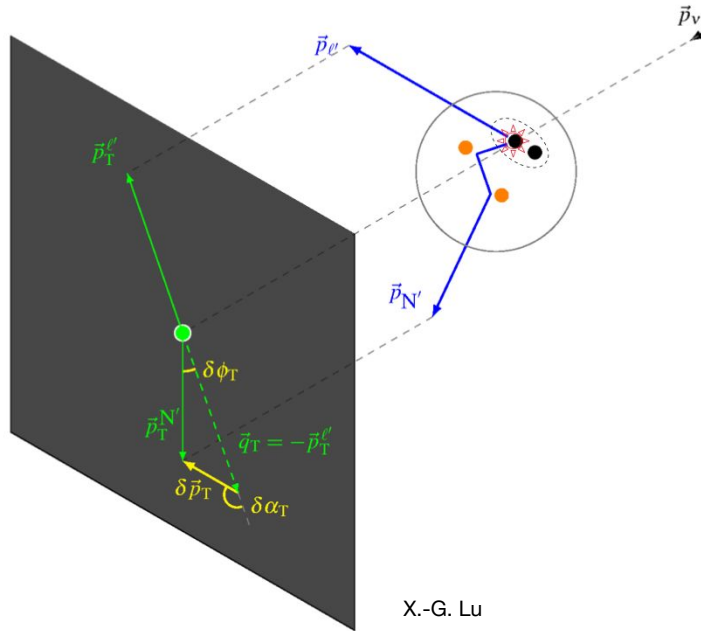
See Anna's presentation

Static nucleon target



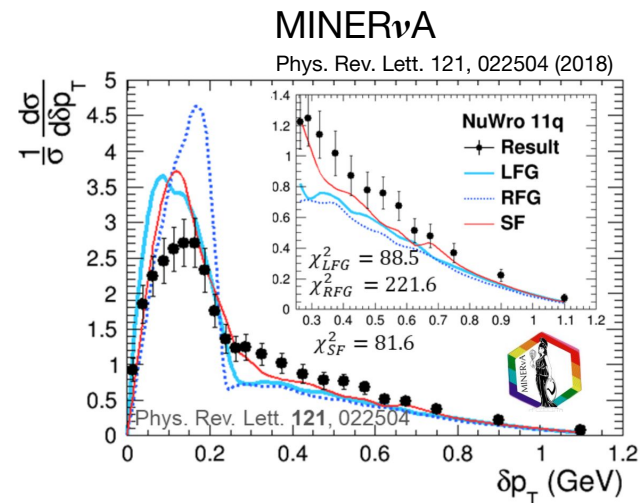
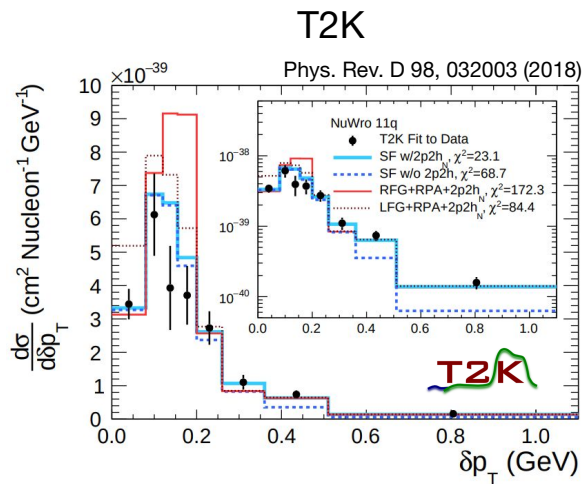
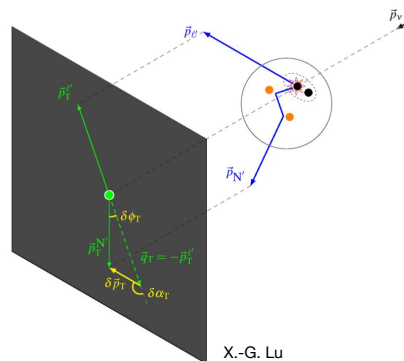


Nucleon bound within nuclear target

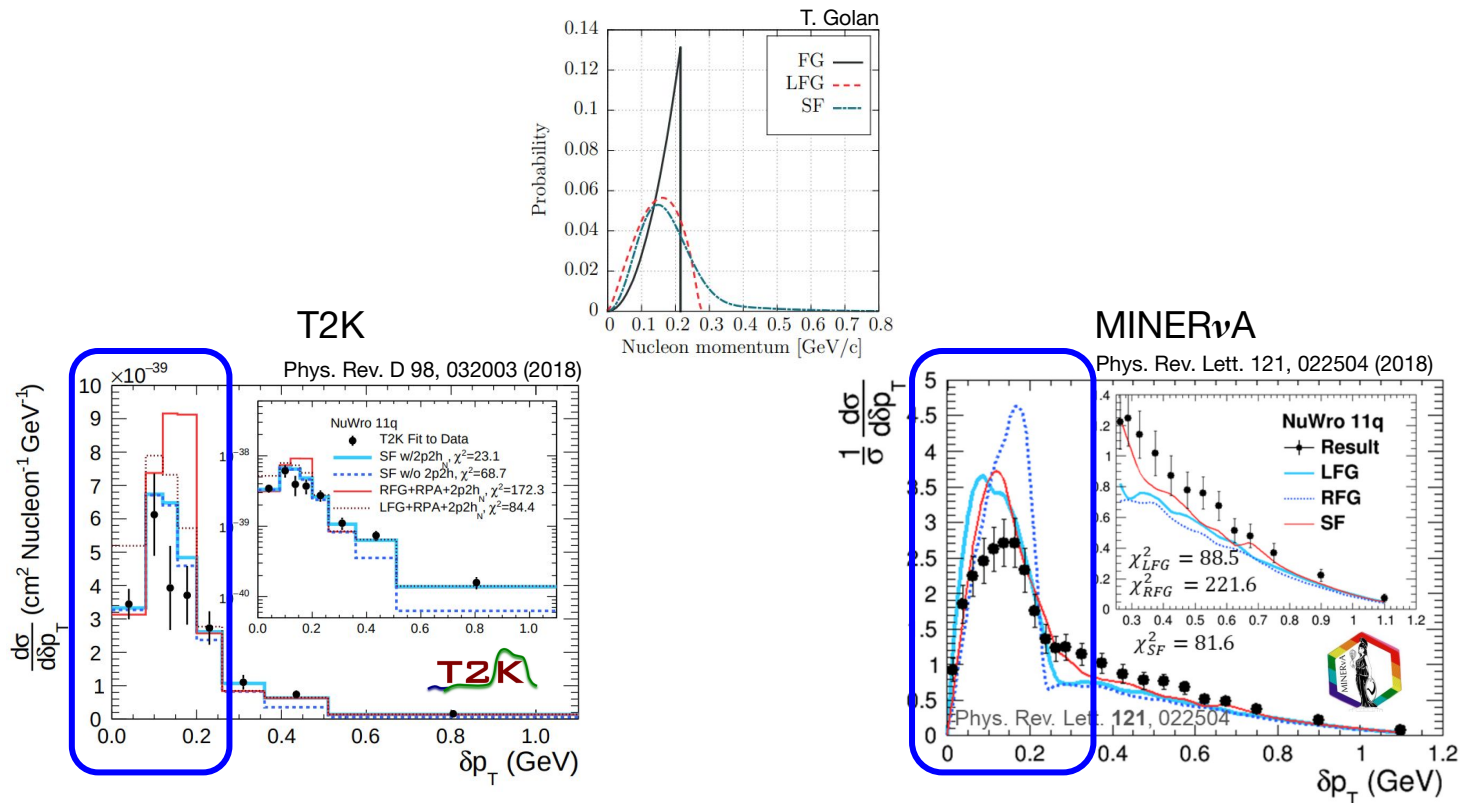


X.-G. Lu

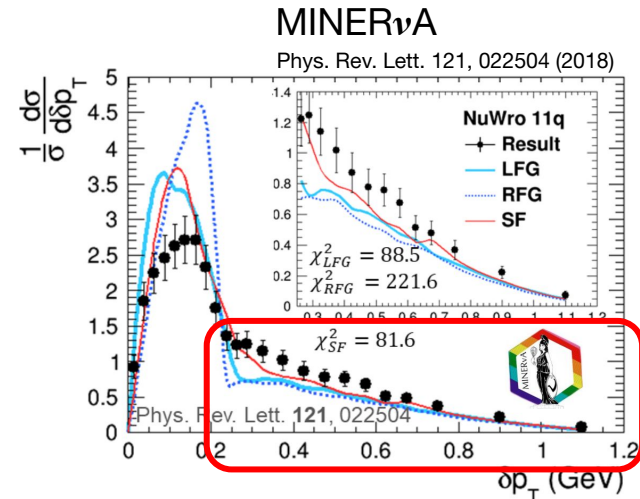
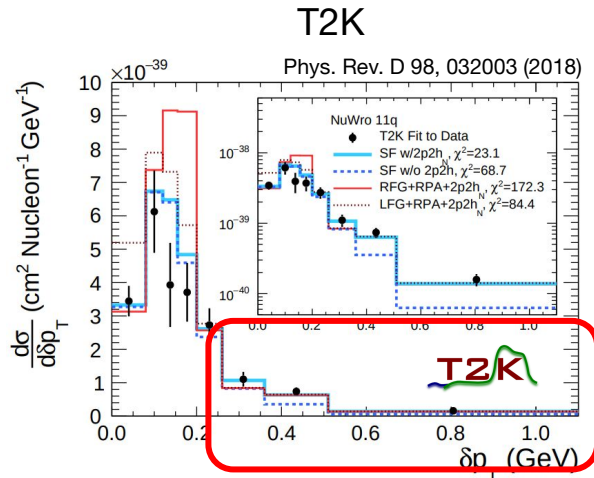
- Need the reconstruction of both muons and nucleons
- Probe nuclear effects (Fermi motion, FSI, ...)



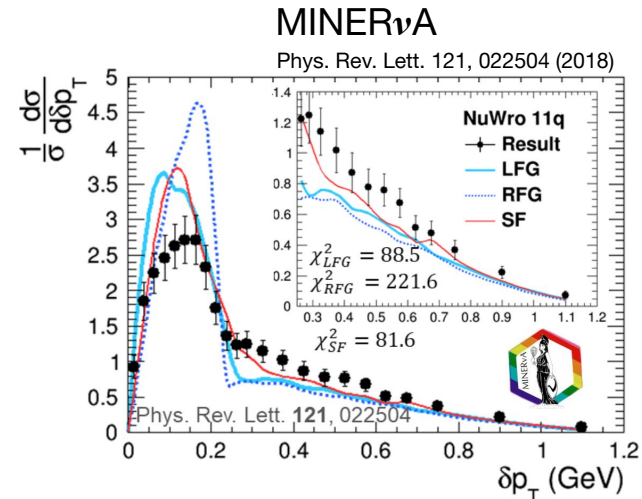
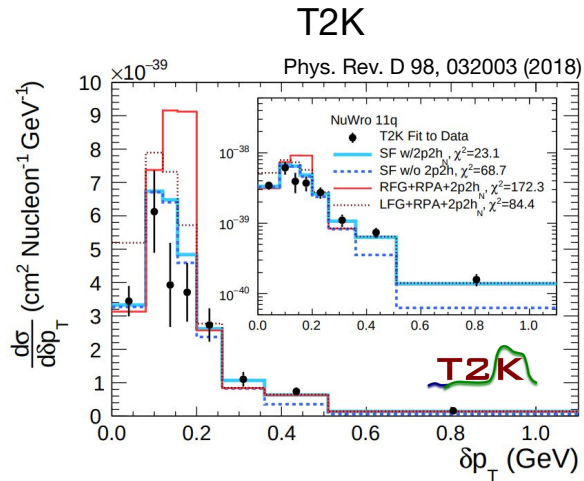
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- The tail of the distribution is sensitive to FSI, SRC, 2p2h



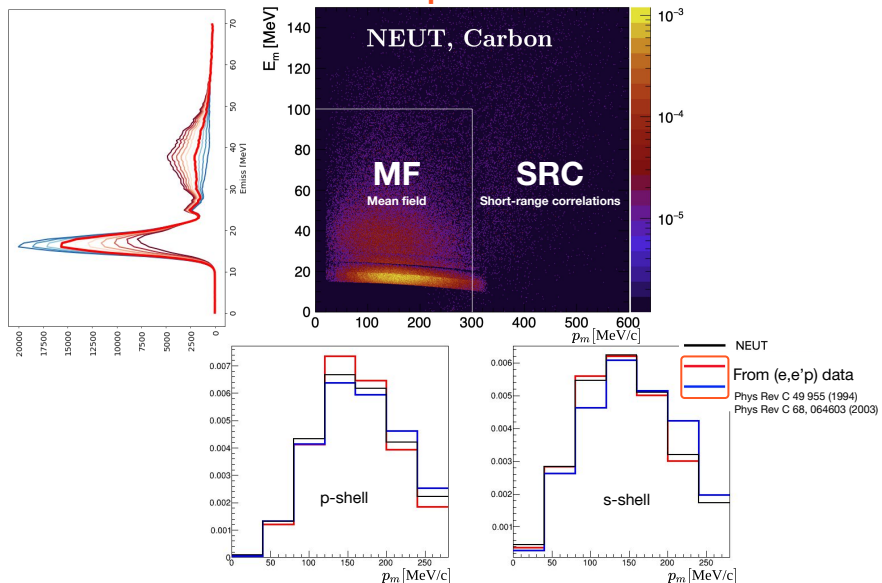
- The bulk of the distribution is sensitive to the initial state nucleon momentum
- The tail of the distribution is sensitive to FSI, SRC, 2p2h
- None of the models describe well the data...





**3. Proposed SF systematic uncertainties**

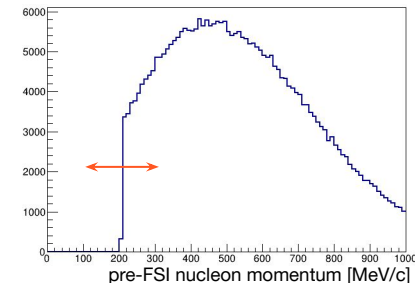
## Initial state parameters



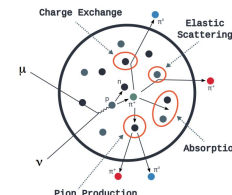
- Change the shape of the missing energy distribution by altering the shell occupancies
- Change the missing momentum distribution shape
- Change the contribution of SRC

## Pauli Blocking parameter

- Change Pauli Blocking threshold on the pre-FSI nucleon momentum



## FSI parameters



- Outgoing nucleon: Change the amount of events undergoing FSI w.r.t. no FSI (see Anna's presentation for ways to improve)
- Outgoing lepton: Change the lepton kinematics with an optical nuclear potential (Phys. Rev. D 91, 033005 (2015))



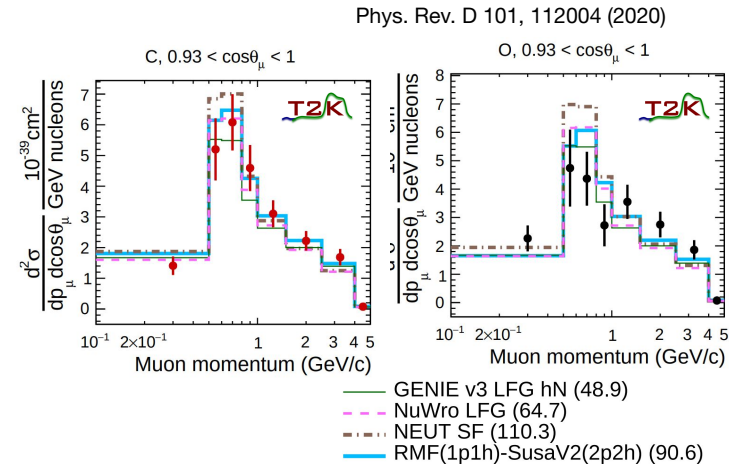
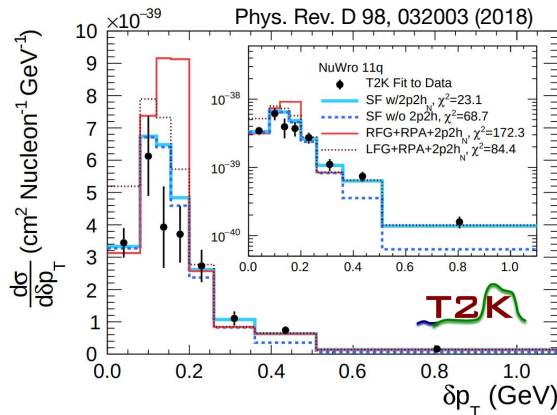
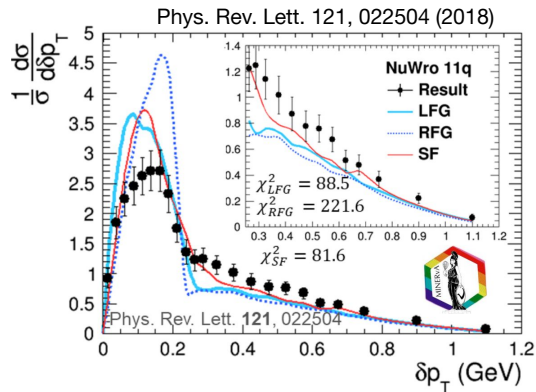
This parameterisation was implemented in [NUANCE](#) and applied on NEUT 5.4.0 neutrino event generator ([arxiv:2106.15809](#))



## 4. Validations with fits to available cross section measurements



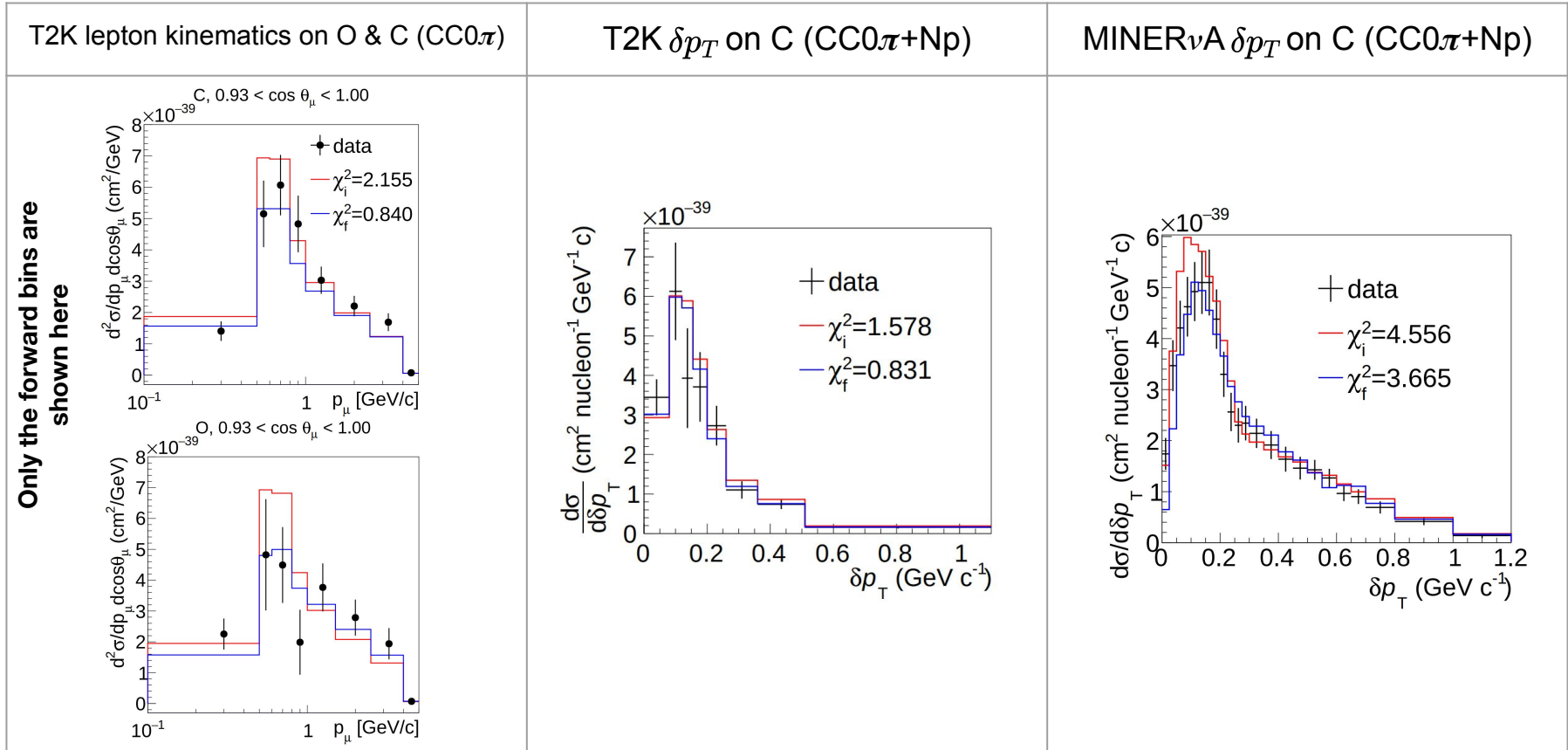
- Cross section measurements allow to test our neutrino interaction models
- Current models struggle to describe cross section data
- In this talk we explore the uncertainties on the models to see if they can accommodate these discrepancies, and also use the ND data to constrain them



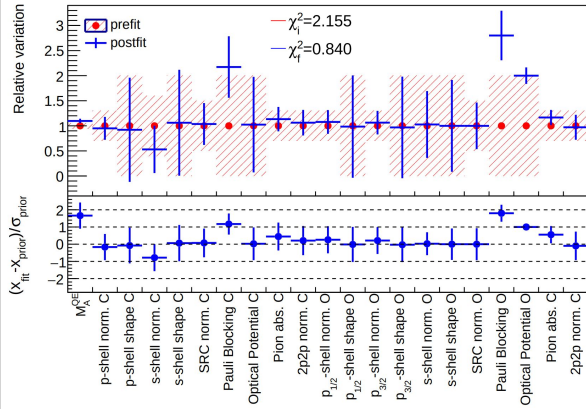
- Ingredients:
  - Model: NEUT with SF model (for Oxygen and Carbon)
  - Parameters: SF model parameters (+ normalisation parameters for other interactions)
  - Data: cross section measurements from T2K and MINERvA
- Chi-square(\*):  $\chi_{\text{data}}^2 = \sum_{1 \leq i, j \leq n} (B_i - B_i^{MC}) (M^{-1})_{ij} (B_j - B_j^{MC})$
- How is this parameterisation able to improve agreement with the data?
- [NUISANCE](#), which is a framework that aims to provide a coherent framework for comparing different neutrino event generators to external data, is used for this study



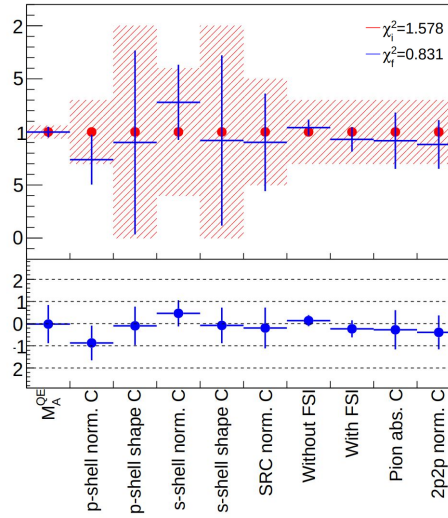
(\*) Peelle's Pertinent Puzzle was avoided using a different decomposition of the data histogram and covariance matrix, see back-up



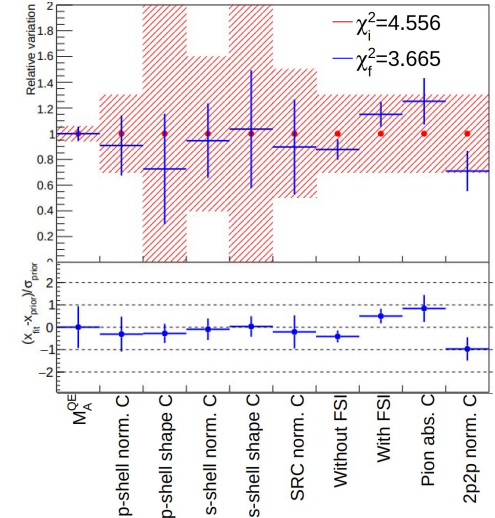
T2K lepton kinematics on O & C (CC0 $\pi$ )



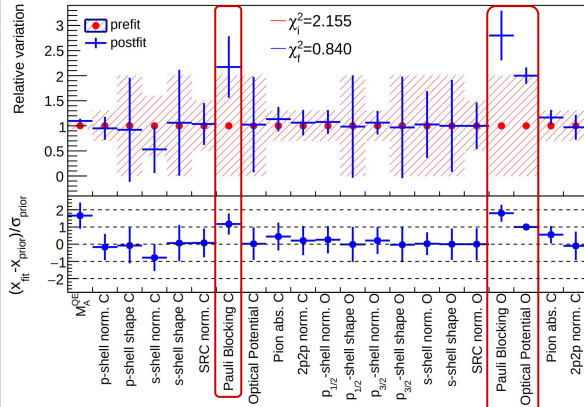
T2K  $\delta p_T$  on C (CC0 $\pi$ +Np)



MINER $\nu$ A  $\delta p_T$  on C (CC0 $\pi$ +Np)

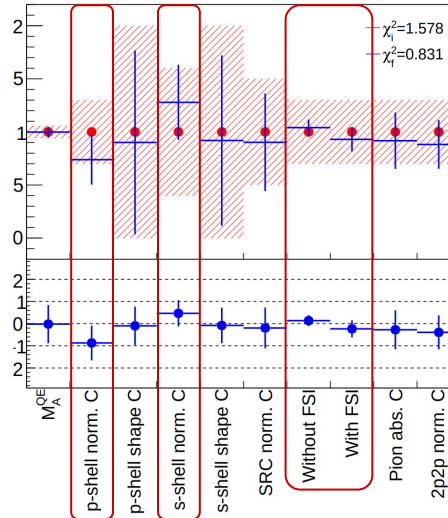


T2K lepton kinematics on O & C (CC0 $\pi$ )



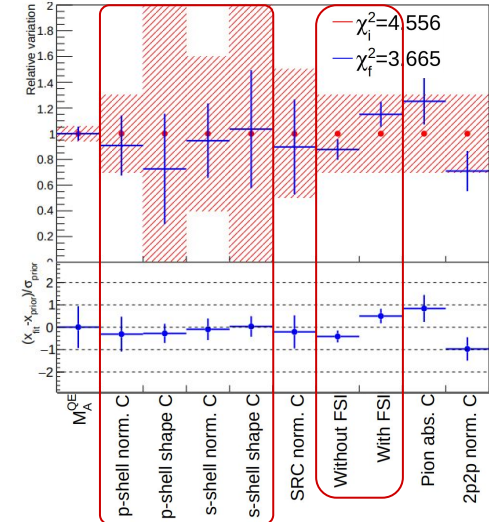
- The agreement with the data is driven by the Pauli Blocking and Optical Potential parameters
- Some sensitivity to the shell parameters can be noticed
- No sensitivity to the missing momentum shape parameters

T2K  $\delta p_T$  on C (CC0 $\pi$ +Np)



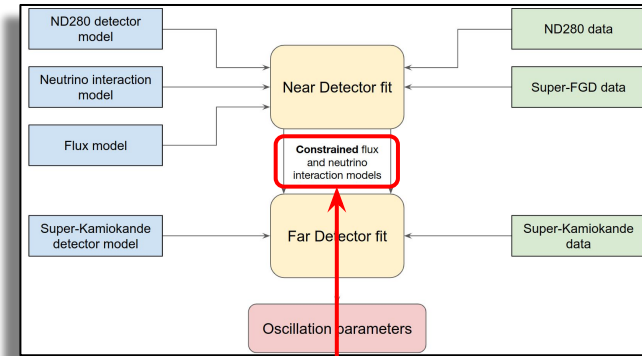
- The transverse momentum imbalance fits show more sensitivity to the initial state parameters as well as the FSI parameters
- The high postfit chi-square in MINER $\nu$ A data may suggest that the FSI model is insufficient

MINER $\nu$ A  $\delta p_T$  on C (CC0 $\pi$ +Np)



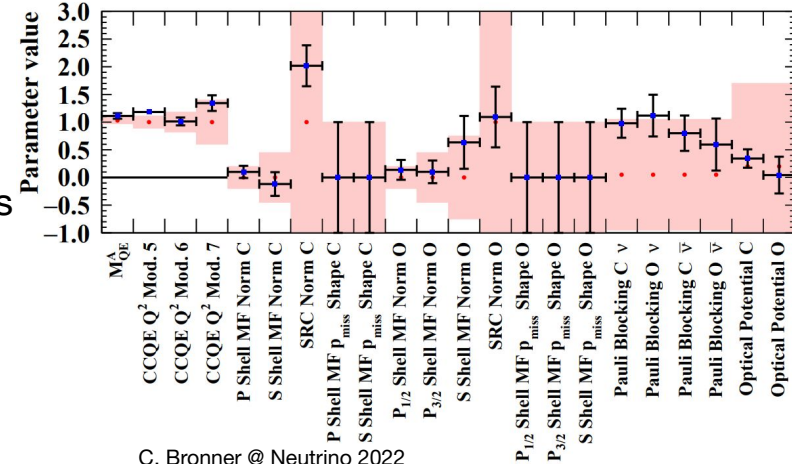
- Fits to **published cross section measurements** gave sensible results and validated the newly-introduced parameters

- Fits to **published cross section measurements** gave sensible results and validated the newly-introduced parameters
- This parameterisation was used in the latest **Oscillation Analysis (OA)** presented at Neutrino 2022 (which uses lepton kinematics as fitting variables)
- What we see:
  - All Pauli blocking parameters are pulled to higher values (suppression of low energy transfer region)
  - Some sensitivity to shell normalisation parameters
  - No sensitivity to the missing momentum shape uncertainties



CCQE Parameters

T2K Run1-10, 2022 Preliminary



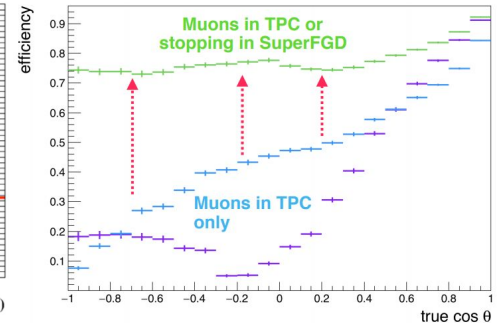
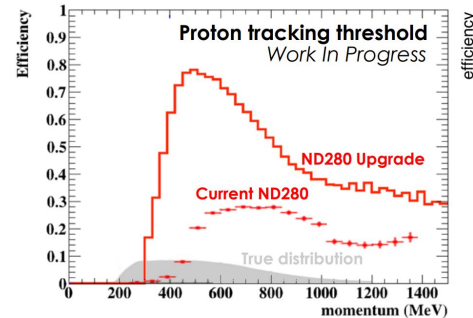
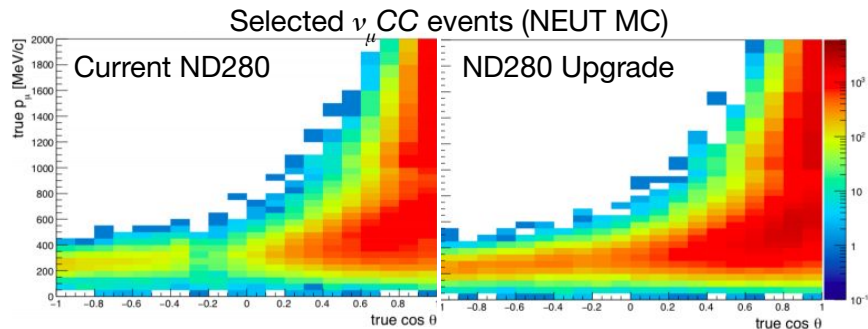
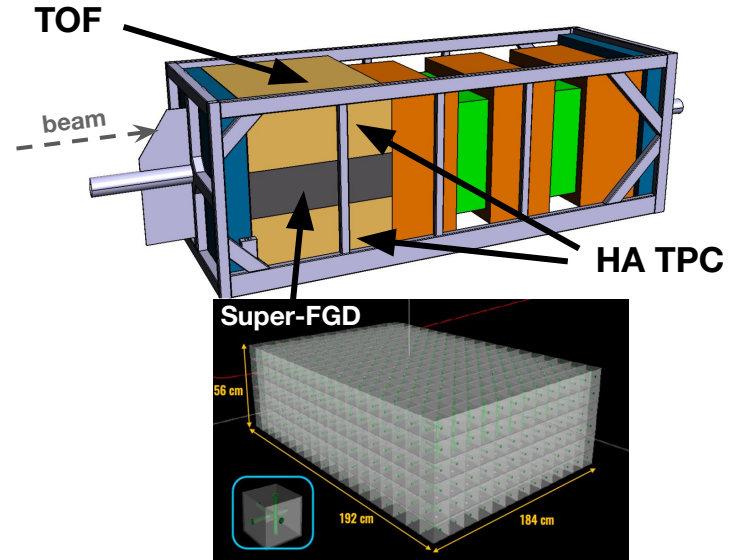
C. Bronner @ Neutrino 2022

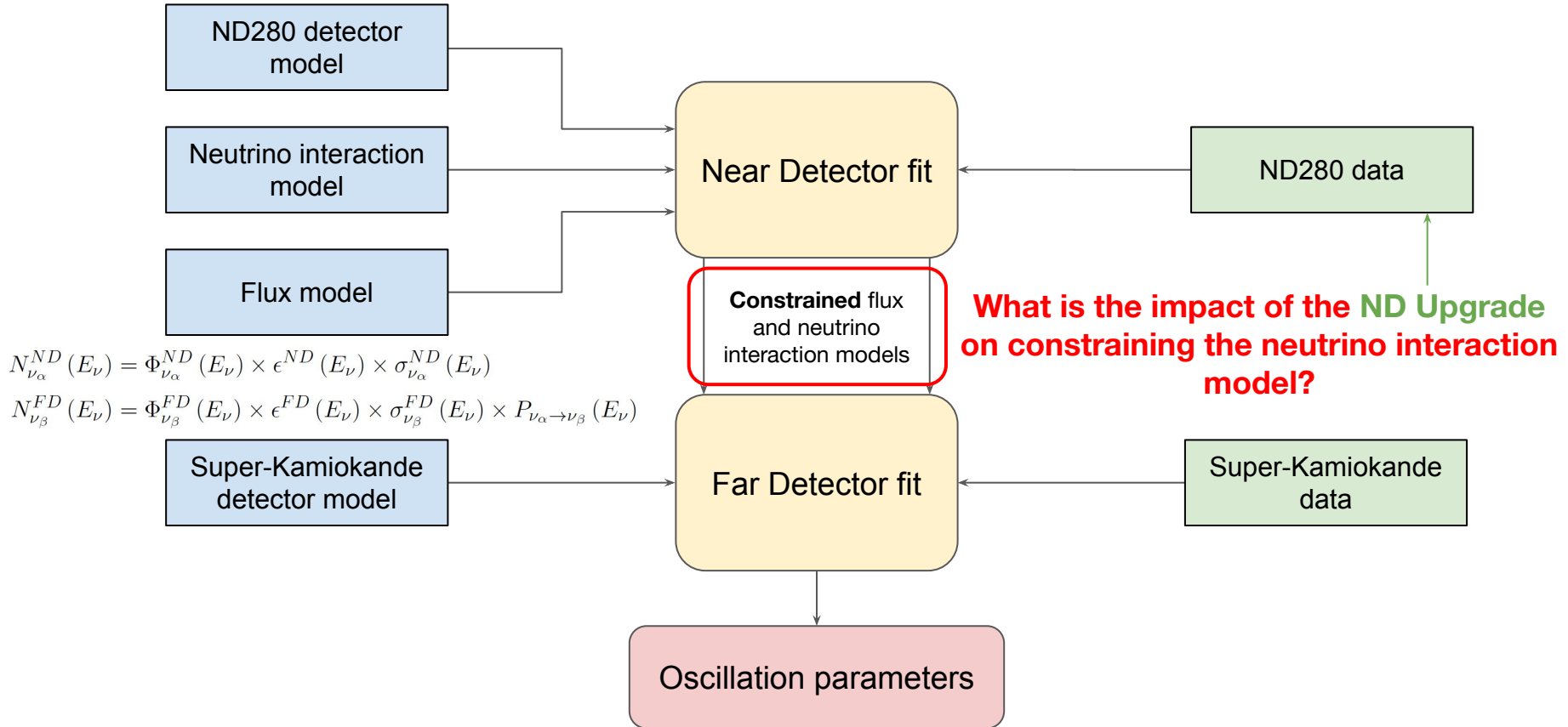


## 5. Projected constraints with T2K ND Upgrade



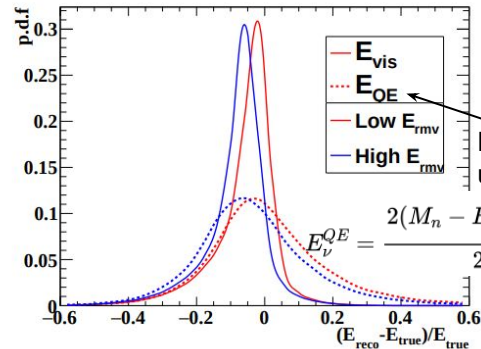
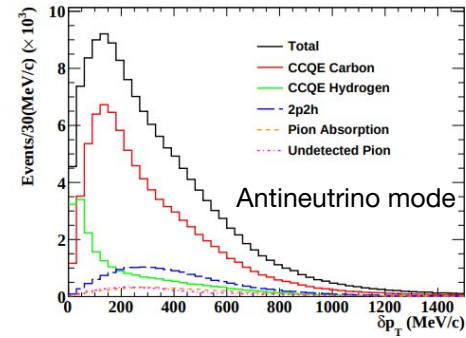
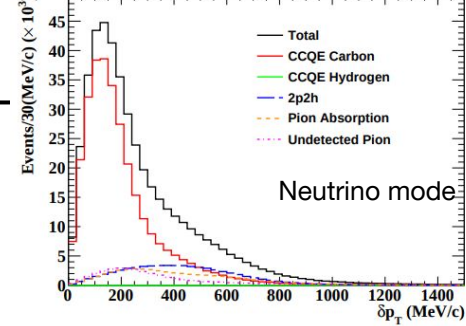
- Improved reconstruction at high and backward angles  $\rightarrow$  better constraints on the neutrino interaction model
- Increased target mass (x2 current ND280)  $\rightarrow$  more statistics
- Better reconstruction of outgoing nucleons  $\rightarrow$  access to new observables





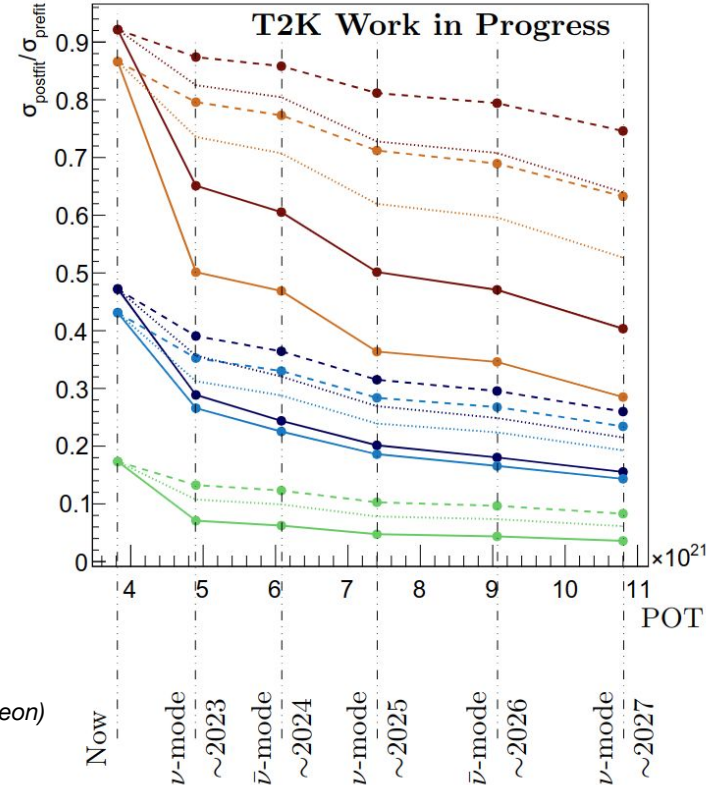
# Future Oscillation Analysis?

- Currently, T2K uses only lepton kinematics for the Oscillation Analysis (OA)
- With the ND280 Upgrade, we expect to obtain more precise measurements of the nucleons coming out from neutrino interactions → what will the impact be on the OA?
- With the nucleon information, we can introduce samples with new observables:
  - Transverse momentum imbalance
  - Visible energy:
    - $E_{\text{vis}} = E_{\mu} + T_p$  for neutrino interactions
    - $E_{\text{vis}} = E_{\mu} + T_n$  for antineutrino interactions
- We use T2K projections of POT assuming a scenario where nu and anti-nu beam modes are alternated on a yearly basis



# Expected improvement: Carbon SF parameters

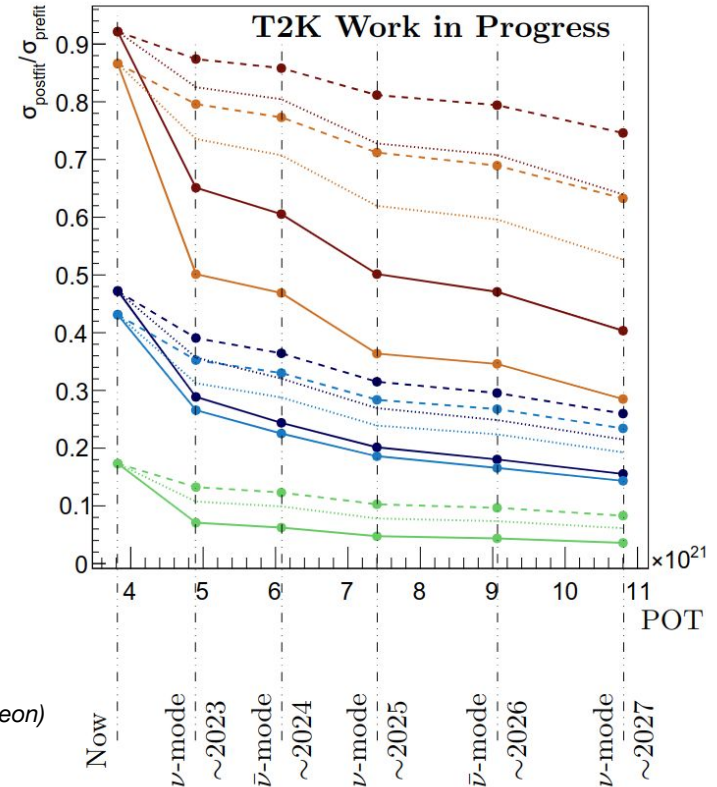
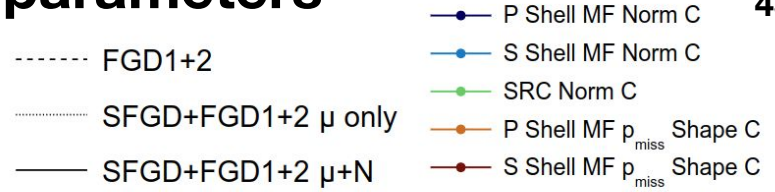
- FGD1+2
- SFGD+FGD1+2  $\mu$  only
- SFGD+FGD1+2  $\mu+N$
- P Shell MF Norm C
- S Shell MF Norm C
- SRC Norm C
- P Shell MF  $p_{miss}$  Shape C
- S Shell MF  $p_{miss}$  Shape C



- FGD1+2 : Current ND fit (no ND280 Upgrade)
- SFGD+FGD1+2  $\mu$  only : ND280 Upgrade using lepton kinematics only
- SFGD+FGD1+2  $\mu+N$  : ND280 Upgrade using (Evis,  $\delta p_T$ ) (when reconstructing a nucleon)

# Expected improvement: Carbon SF parameters

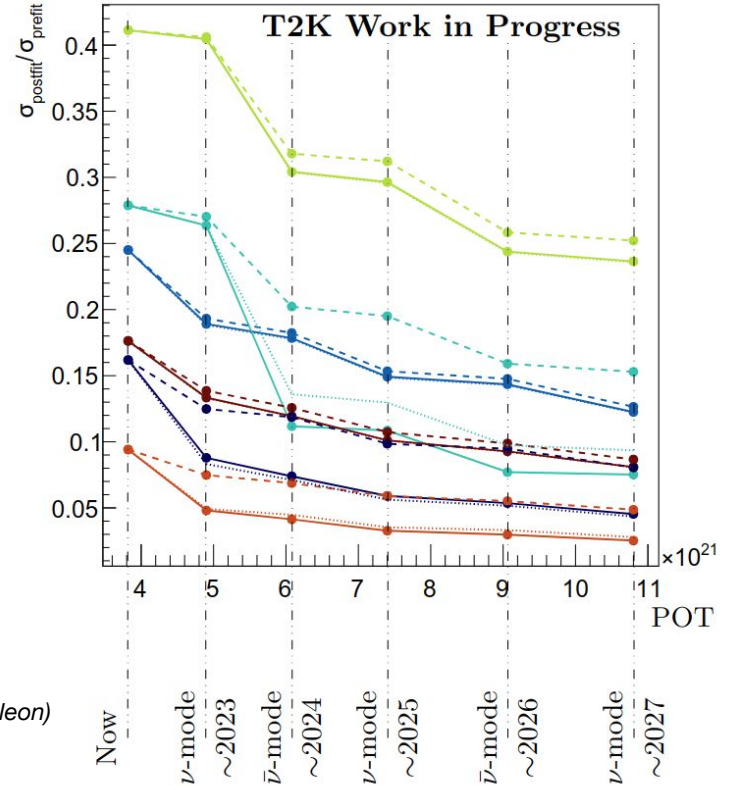
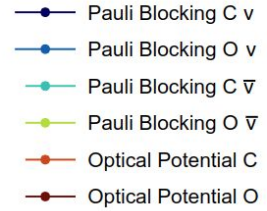
- Significant improvement with respect to the current ND configuration
- The use of nucleon information with (Evis,  $\delta p_T$ ) allows larger constraints especially on the  $p_{miss}$  shape parameters (these parameters are fixed in the current OA due to lack of sensitivity)



- - - - - FGD1+2 : Current ND fit (no ND280 Upgrade)  
 . . . . . SFGD+FGD1+2  $\mu$  only : ND280 Upgrade using lepton kinematics only  
 ——— SFGD+FGD1+2  $\mu$ +N : ND280 Upgrade using (Evis,  $\delta p_T$ ) (when reconstructing a nucleon)

# Expected improvement: PB and OP

- The increased statistics and the use of lepton kinematics when the nucleon is not reconstructed allow improved constraints on the Pauli Blocking and Optical Potential parameters (low energy transfer region)
- The upgrade has small impact on the oxygen parameters (SFGD is made of plastic scintillator)



----- FGD1+2 : Current ND fit (no ND280 Upgrade)

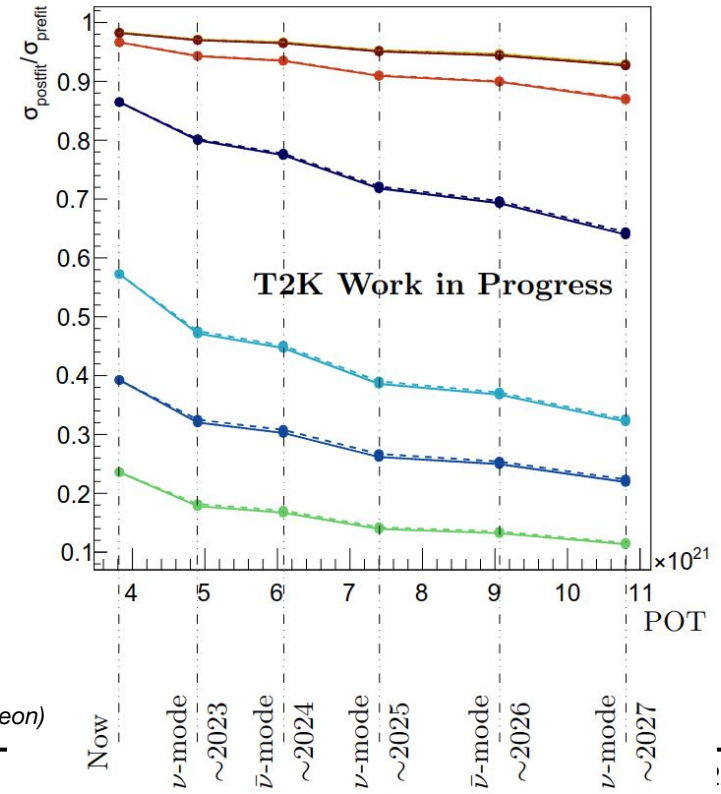
..... SFGD+FGD1+2 μ only : ND280 Upgrade using lepton kinematics only

— SFGD+FGD1+2 μ+N : ND280 Upgrade using (Evis, δp<sub>T</sub>) (when reconstructing a nucleon)

# Expected improvement: Oxygen SF parameters

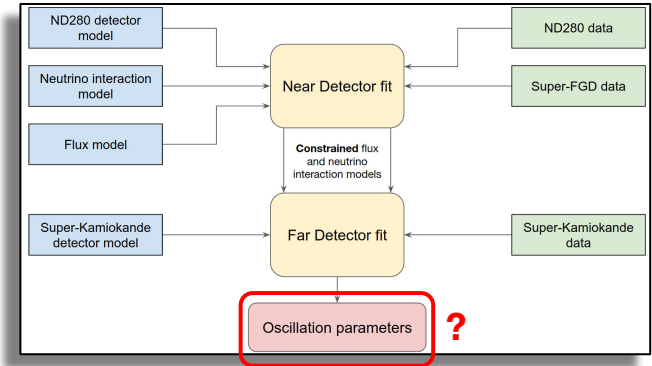
- With the presented parameterisation, the upgrade has little impact on the constraints of the Oxygen parameters:
  - SF model is built independently from electron scattering data for O and C
- This is still under study

- P<sub>1/2</sub> Shell MF Norm O
- P<sub>3/2</sub> Shell MF Norm O
- S Shell MF Norm O
- SRC Norm O
- P<sub>1/2</sub> Shell MF p<sub>miss</sub> Shape O
- P<sub>3/2</sub> Shell MF p<sub>miss</sub> Shape O
- S Shell MF p<sub>miss</sub> Shape O

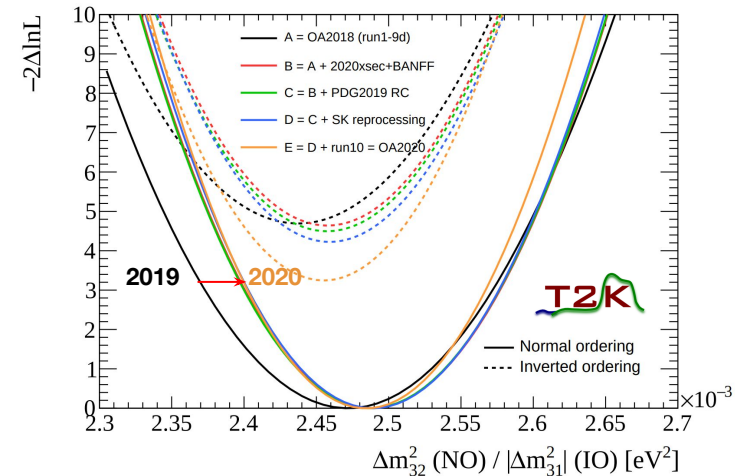


- FGD1+2 : Current ND fit (no ND280 Upgrade)
- ..... SFGD+FGD1+2 μ only : ND280 Upgrade using lepton kinematics only
- SFGD+FGD1+2 μ+N : ND280 Upgrade using (Evis, δp<sub>T</sub>) (when reconstructing a nucleon)

- We plan on evaluating the impact of the improved sensitivity with the ND Upgrade on the oscillation parameters using the FD fitter



- The 2020 results, which improved the cross section model, showed better constraints on  $\Delta m^2$





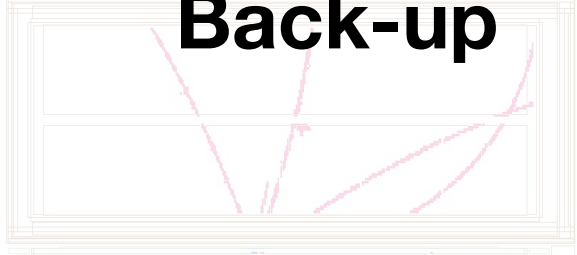
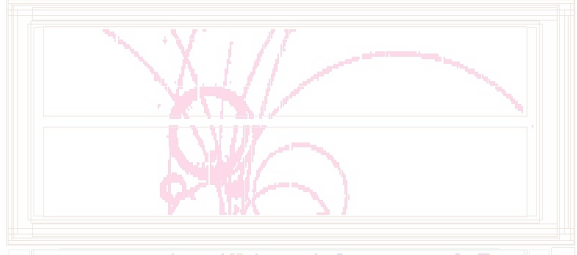
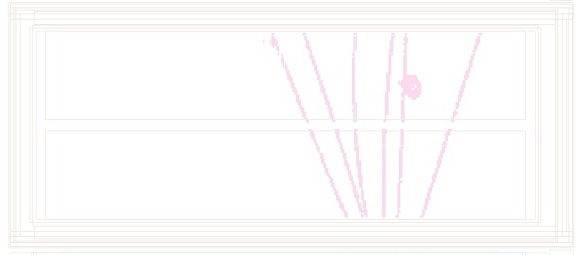
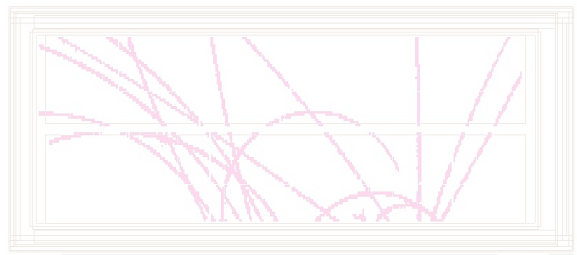


## 6. Summary and outlook

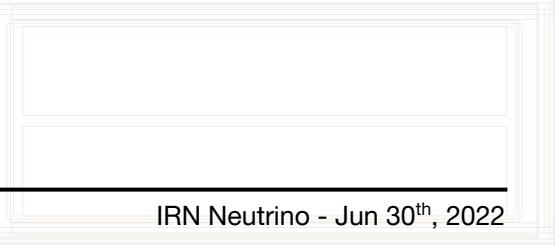
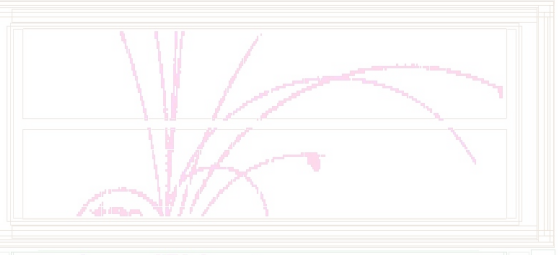
- A new set of physically-motivated uncertainties on the predictions of the CCQE Spectral Function model were introduced
- Fitting to external data shows great improvement for T2K cross section measurements, whereas MINERvA data is sensitive to this parameterisation but shows little improvement in the data/MC agreement
- The impact of adding SFGD samples on the ND fit is studied with:
  - The choice of the binning variables (lepton kinematics vs. (Evis,  $\delta pT$ )) has a significant impact on the constraints of some uncertainties (e.g. SF model parameters, 2p2h, nucleon FSI)
  - The impact on Oxygen parameters is limited since the current model for OA2022 does not account for O vs. C correlations

Next steps (non-exhaustive):

- What is the impact of the added SFGD samples on the oscillation parameters?
- What would the impact be if we add correlations between O and C?
- This parameterisation is a first step to account for the uncertainty on the nucleon production, it still needs further improvement especially for FSI (*see Anna's talk next*)



**Back-up**



- Attempts to fit neutrino cross-section data with a parameterisation of the interaction model gives us seemingly unphysically low normalisations PPP
- The reason this happens is subtle, but is related to the strong correlations in published covariance matrices and the corresponding “type” of Gaussian errors approximation
- In our standard approach we assume the absolute uncertainty on the cross section is independent of its normalisation
  - i.e. our uncertainties state that a 10 fb / GeV uncertainty on some bin remains at 10 fb / GeV even if we had underestimated our flux by 10% (and so the cross section is lower than measured)
  - This implies the relative uncertainty is larger if fitting to models that predict lower normalisations. This is what give us PPP
- We could alternatively suggest that it should be the relative uncertainty that is independent of its normalisation (D’Agostini does)
  - i.e. our uncertainties would state that a 10% uncertainty on some bin remains at 10% even if we had underestimated our flux by 10%
  - This implies the absolute uncertainty is larger if fitting to models that predict lower normalisations

- We can construct a covariance matrix that keeps the relative uncertainties constant when the normalization changes: a “Norm-Shape” covariance where one row contains the normalization of the data and the rest contains the shape

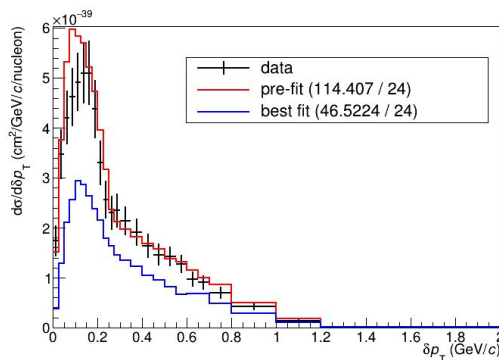
$$H_1 : \{B_1, \dots, B_n\} \rightarrow H_2 : \{C_1, \dots, C_n\} \quad C_i = \begin{cases} \frac{B_i}{\sum_k B_k} & , i < n \\ \sum_k B_k & , i = n \end{cases}$$

- We can obtain  $\text{Cov}[\{C_i\}]$  (norm-shape covariance) directly from the data covariance given by experiments  $\text{Cov}[\{B_i\}]$
- Perform the fit in this new basis using  $N = \text{Cov}[\{C_i\}]$

$$\chi_{\text{NS}}^2 = \sum_{1 \leq i, j \leq n} (C_i - C_i^{\text{MC}}) (N^{-1})_{ij} (C_j - C_j^{\text{MC}})$$

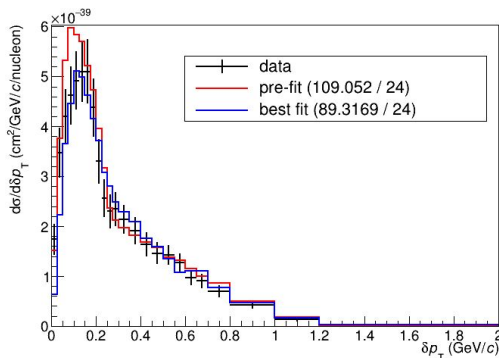
- Example of MINERvA dpt fit with the same parameterisation of slide 41:

Fit minimizing the usual chi2



	Normal chi2	NS chi2
<b>prefit</b>	<b>114.407</b>	109.052
<b>postfit</b>	<b>46.5224</b>	108.201

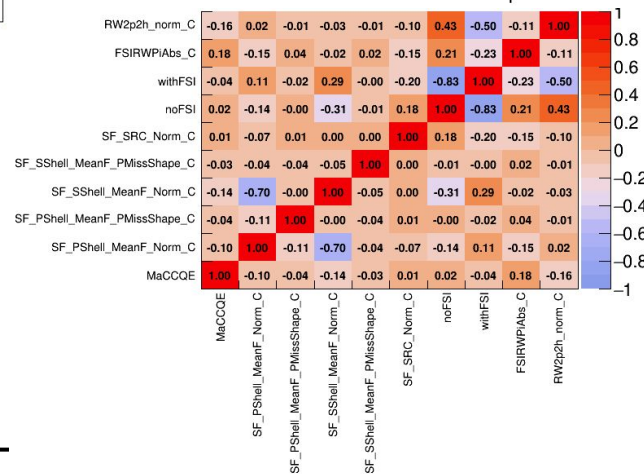
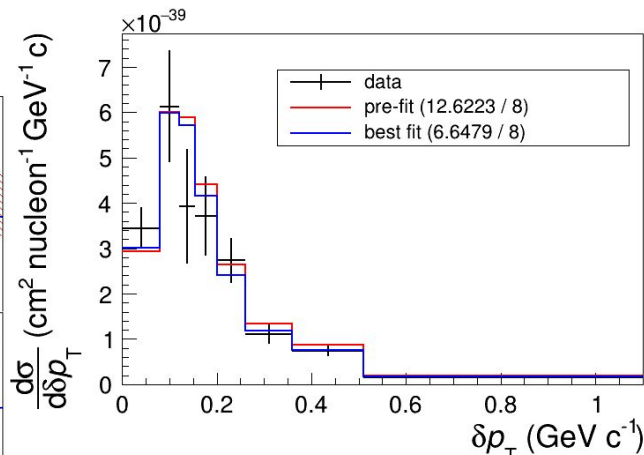
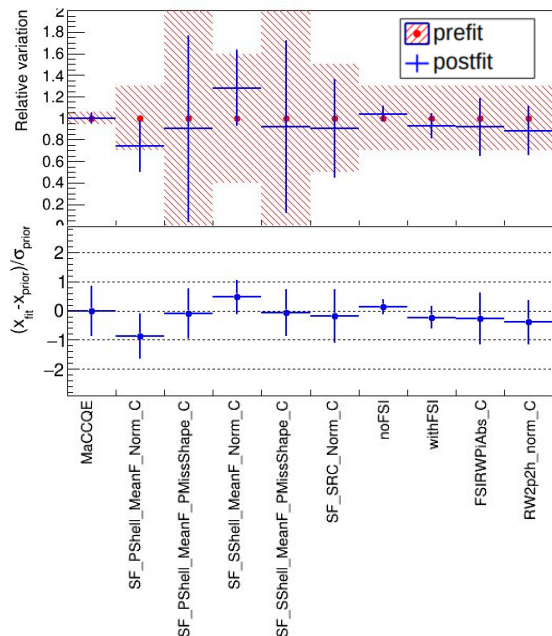
Fit minimizing the new chi2



	Normal chi2	NS chi2
<b>prefit</b>	114.407	<b>109.052</b>
<b>postfit</b>	85.2858	<b>89.3169</b>

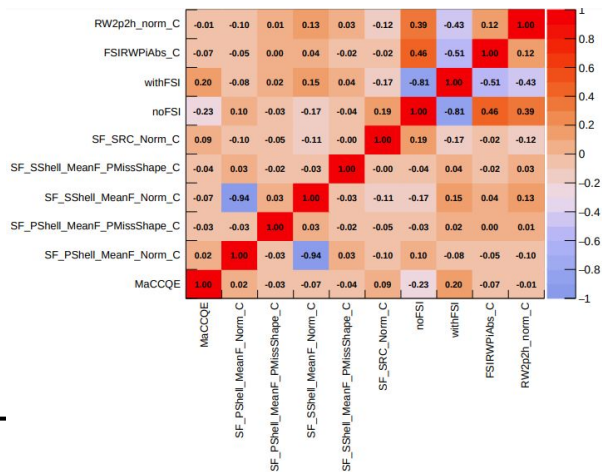
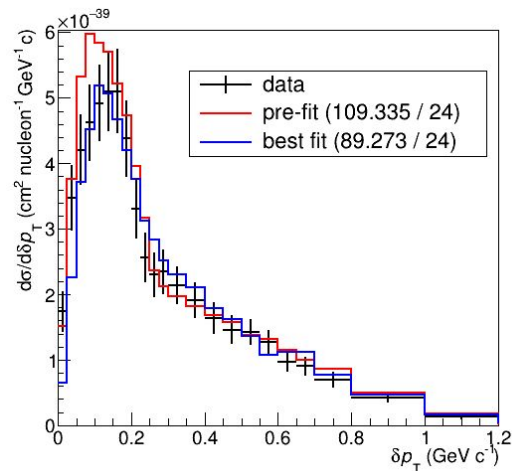
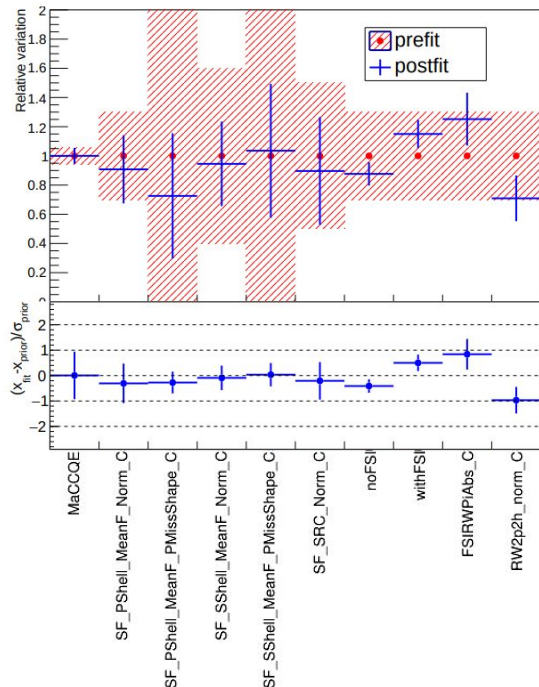
Parameters used in this fit:

- $M_A^{QE}$
- Shell occupancies
- $p_m$  shape
- SRC norm
- No FSI / with FSI (correlated)
- Pion absorption normalization
- 2p2h normalization



Parameters used in this fit:

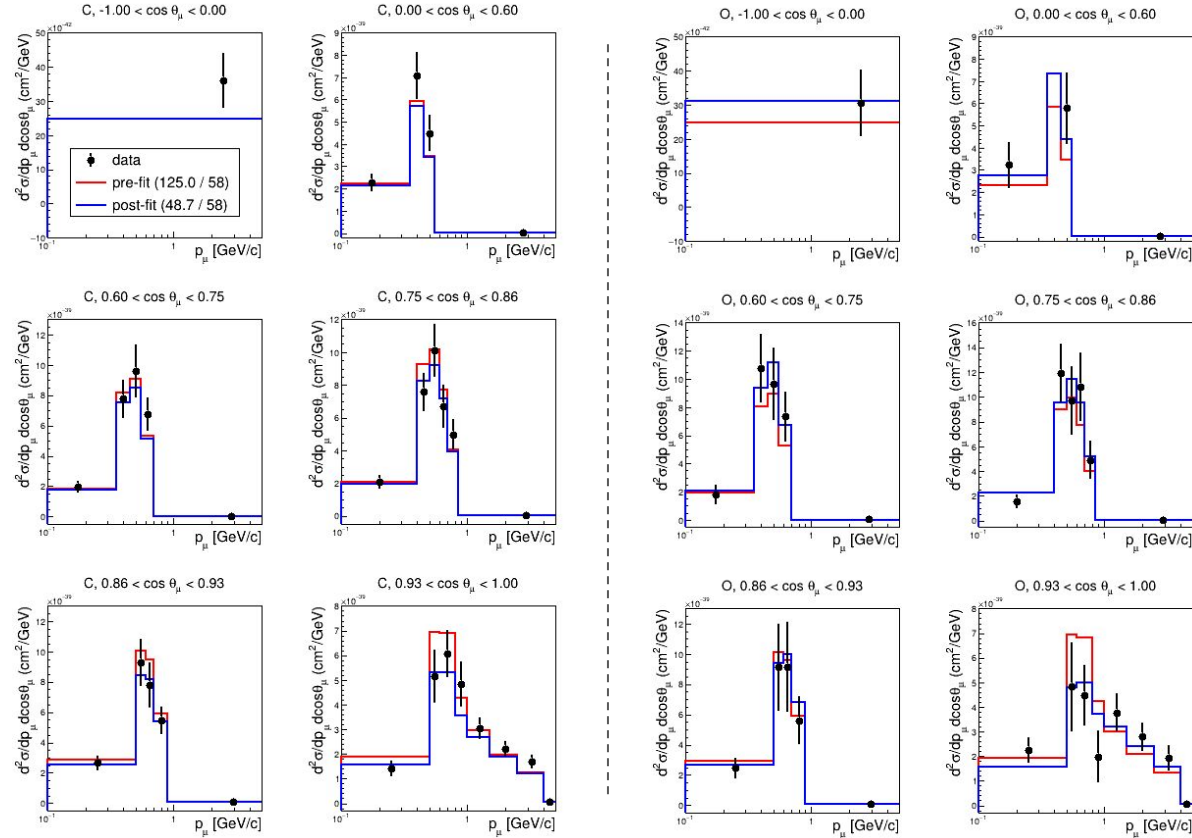
- $M_A^{QE}$
- Shell occupancies
- $p_m$  shape
- SRC norm
- No FSI / with FSI (correlated)
- Pion absorption normalization
- 2p2h normalization





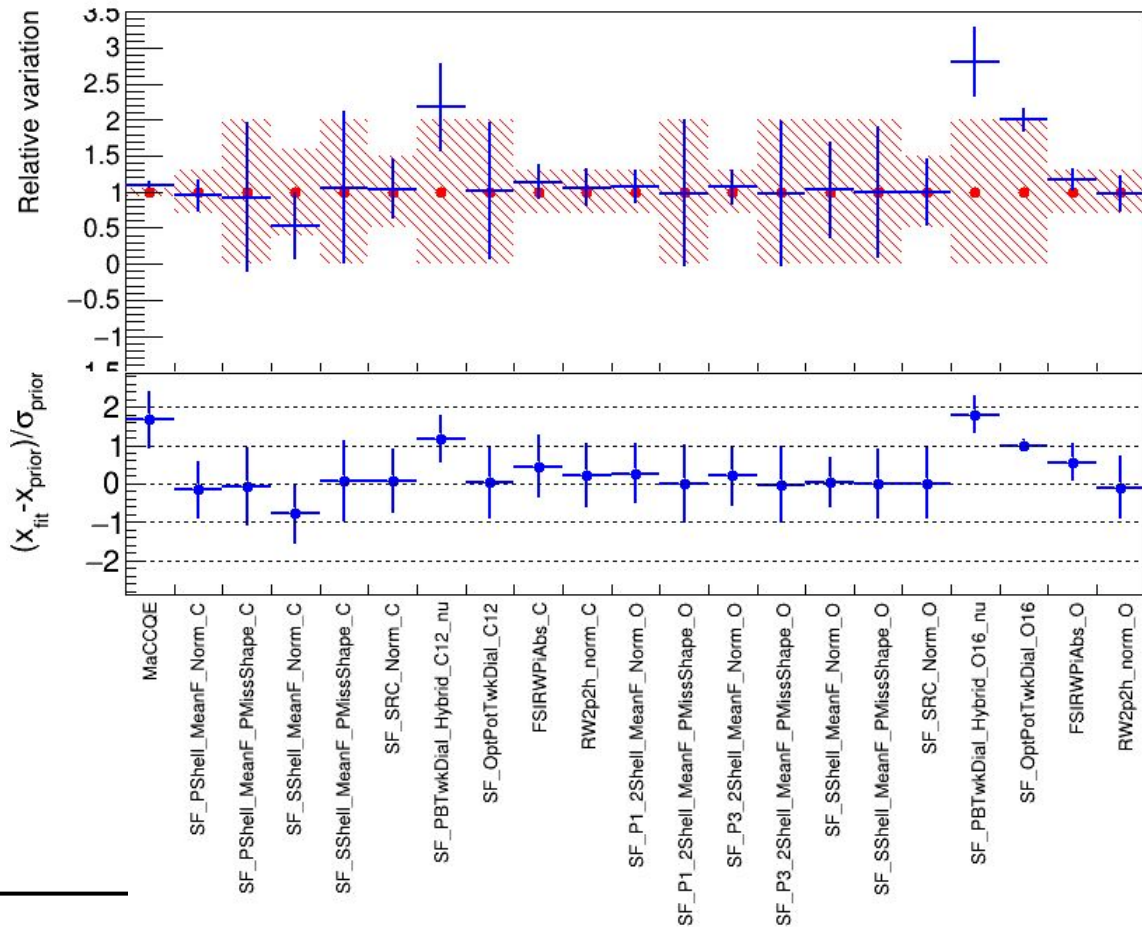
Parameters used in this fit:

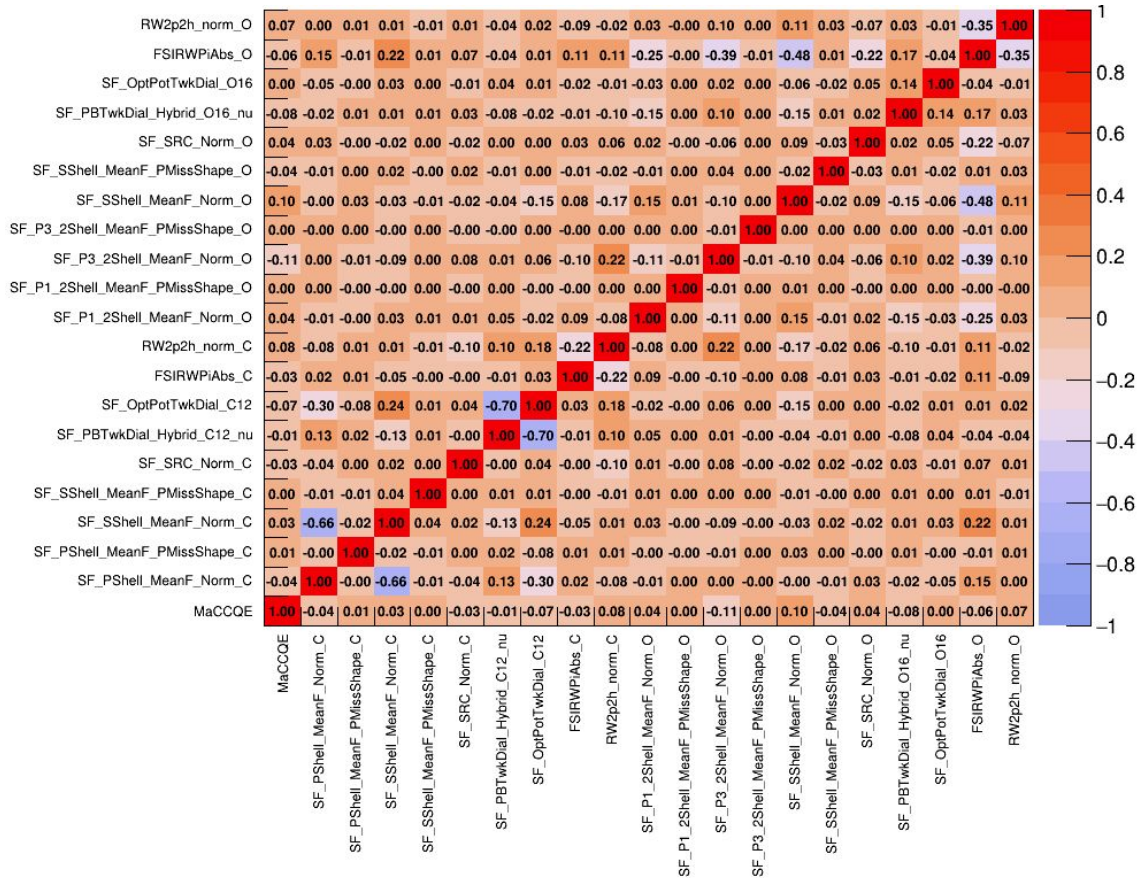
- $M_A^{QE}$
- Shell occupancies
- $p_m$  shape
- SRC norm
- Pauli blocking
- Optical potential correction
- Pion absorption normalization
- 2p2h normalization



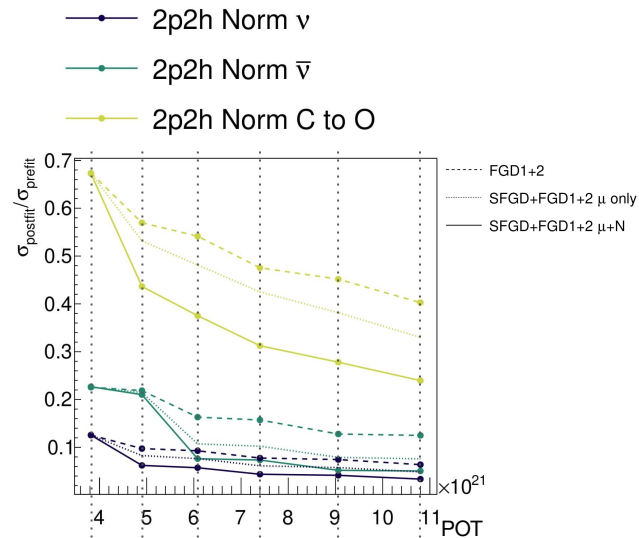
Parameters used in this fit:

- $M_A^{QE}$
- Shell occupancies
- $p_m$  shape
- SRC norm
- Pauli blocking
- Optical potential correction
- Pion absorption normalization
- 2p2h normalization



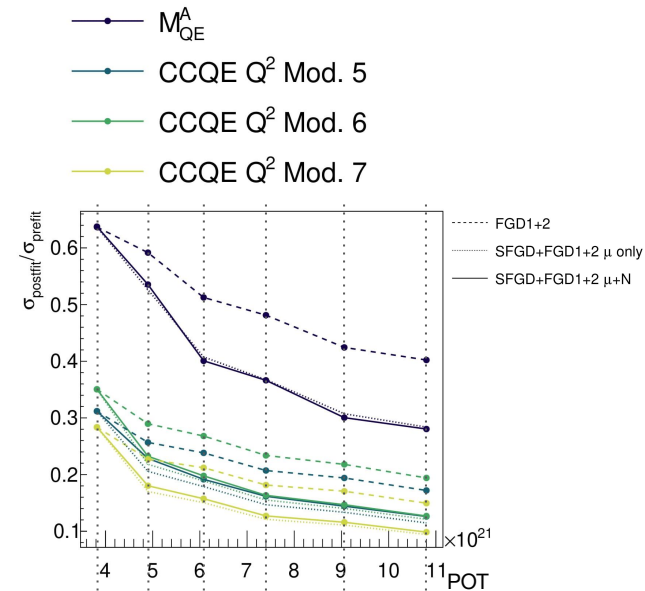


- 2p2h norm parameters also benefit from the Upgrade and the added nucleon information



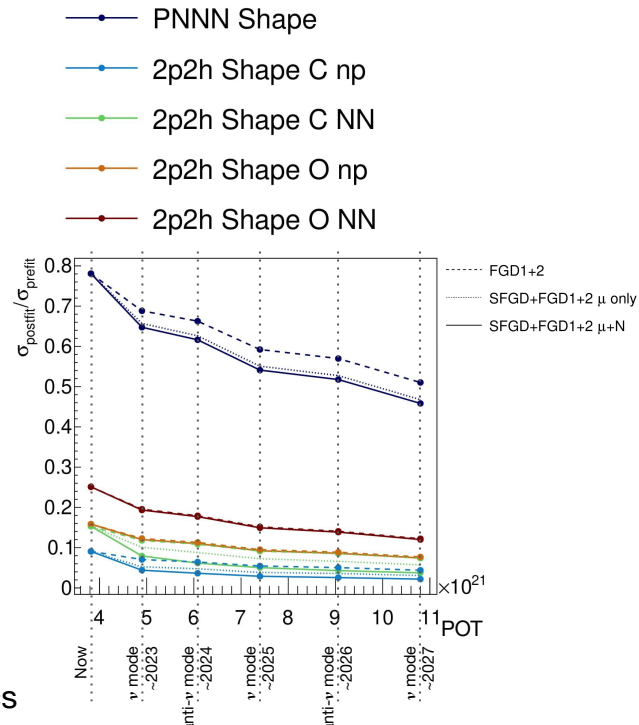
- FGD1+2 : Current ND fit, no additional samples
- SFGD+FGD1+2  $\mu$  only : Add to current ND fit SFGD samples binned in lepton kinematics
- SFGD+FGD1+2  $\mu$ +N : Add to current ND fit SFGD samples binned in (Evis,  $\delta p_T$ )

- MAQE has tight prefit constraints, the improvement comes mainly from the additional statistics
- SFGD+FGD1+2  $\mu$  only shows slightly better constraints for some of the high- $Q^2$  parameters



- FGD1+2 : Current ND fit, no additional samples
- SFGD+FGD1+2  $\mu$  only : Add to current ND fit SFGD samples binned in lepton kinematics
- SFGD+FGD1+2  $\mu$ +N : Add to current ND fit SFGD samples binned in (Evis,  $\delta p_T$ )

- The Upgrade has no significant impact on the Oxygen parameters
- The SFGD adds little more constraint on the 2p2h shape parameters → these parameters are designed to impact lepton kinematics only



----- FGD1+2 : Current ND fit, no additional samples

----- SFGD+FGD1+2  $\mu$  only : Add to current ND fit SFGD samples binned in lepton kinematics

----- SFGD+FGD1+2  $\mu$ +N : Add to current ND fit SFGD samples binned in (Evis,  $\delta p_T$ )

# Expected improvement: Flux parameters

- ND280 FHC  $\nu_\mu$  [0, 0.4 GeV]
- ND280 FHC  $\nu_\mu$  [0.4, 0.5 GeV]
- ND280 FHC  $\nu_\mu$  [0.5, 0.6 GeV]
- ND280 FHC  $\nu_\mu$  [0.6, 0.7 GeV]
- ND280 FHC  $\nu_\mu$  [0.7, 1 GeV]
- ND280 FHC  $\nu_\mu$  [1, 1.5 GeV]
- ND280 FHC  $\nu_\mu$  [1.5, 2.5 GeV]
- ND280 FHC  $\nu_\mu$  [2.5, 3.5 GeV]
- ND280 FHC  $\nu_\mu$  [3.5, 5 GeV]
- ND280 FHC  $\nu_\mu$  [5, 7 GeV]
- ND280 FHC  $\nu_\mu$  [7, 30 GeV]

- ND280 FHC  $\nu_e$  [0, 0.5 GeV]
- ND280 FHC  $\nu_e$  [0.5, 0.7 GeV]
- ND280 FHC  $\nu_e$  [0.7, 0.8 GeV]
- ND280 FHC  $\nu_e$  [0.8, 1.5 GeV]
- ND280 FHC  $\nu_e$  [1.5, 2.5 GeV]
- ND280 FHC  $\nu_e$  [2.5, 4 GeV]
- ND280 FHC  $\nu_e$  [4, 30 GeV]
- ND280 FHC  $\bar{\nu}_e$  [0, 2.5 GeV]
- ND280 FHC  $\bar{\nu}_e$  [2.5, 30 GeV]

- ND280 RHC  $\bar{\nu}_\mu$  [0, 0.4 GeV]
- ND280 RHC  $\bar{\nu}_\mu$  [0.4, 0.5 GeV]
- ND280 RHC  $\bar{\nu}_\mu$  [0.5, 0.6 GeV]
- ND280 RHC  $\bar{\nu}_\mu$  [0.6, 0.7 GeV]
- ND280 RHC  $\bar{\nu}_\mu$  [0.7, 1 GeV]
- ND280 RHC  $\bar{\nu}_\mu$  [1, 1.5 GeV]
- ND280 RHC  $\bar{\nu}_\mu$  [1.5, 2.5 GeV]
- ND280 RHC  $\bar{\nu}_\mu$  [2.5, 3.5 GeV]
- ND280 RHC  $\bar{\nu}_\mu$  [3.5, 5 GeV]
- ND280 RHC  $\bar{\nu}_\mu$  [5, 7 GeV]
- ND280 RHC  $\bar{\nu}_\mu$  [7, 30 GeV]

- FGD1+2
- SFGD+FGD1+2  $\mu$  only
- SFGD+FGD1+2  $\mu$ +N

- ND280 RHC  $\nu_\mu$  [0, 0.7 GeV]
- ND280 RHC  $\nu_\mu$  [0.7, 1 GeV]
- ND280 RHC  $\nu_\mu$  [1, 1.5 GeV]
- ND280 RHC  $\nu_\mu$  [1.5, 2.5 GeV]
- ND280 RHC  $\nu_\mu$  [2.5, 30 GeV]

