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Search for light vector boson using J/Ψ at BESIII and Belle II

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Introduction

- The Standard Model (SM) is a successful theory describing the physics at particle level and their interactions.

- There have been discussions of extending the SM by gauging the lepton number, e.g. $L_\mu - L_\tau$ or $L_e - L_\tau$, mainly intending to explain the muon anomalous magnetic moments $(g-2)_\mu$.

- This gives rise to a **leptophilic light vector boson X.**

- The X boson may couple to the SM quark sector via interactions with heavy vector-like fermions mixing with SM quark.

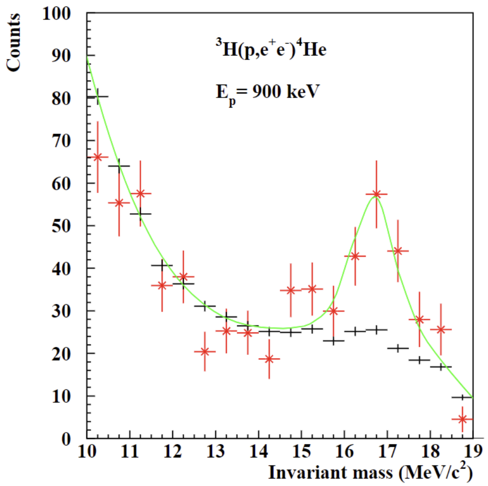


FIG. 3. *Invariant mass distribution derived for the 20.49 MeV transition in* ^{4}*He.*

- Recent result of $^8\mathrm{Be}^*$ anomaly from Atomki experiment prefers a 17 MeV vector boson that couples to the electrons, u- and d- quarks.

- High luminosity lepton colliders, such as **BESIII** and Belle II, provide less QCD background than hadron colliders, making them ideal environments to search for sub-GeV particles with feeble couplings to SM particles.

Introduction

- In this work, we focus on the light vector boson search in association with J/Ψ at BESIII and Belle II

- At BESIII, if the vector boson is lighter than about 110 MeV, it can be produced through $J/\Psi \to \eta_c + X$ followed by the $X \to e^+e^-$ or $X \to \nu\bar{\nu}$ decays.

- At Belle II, the process $e^+e^- \to \ell^+\ell^- J/\Psi \to \ell^+\ell^-\eta_c X \to \ell^+\ell^-\eta_c e^+e^-$, in which J/Ψ and η_c are inferred by the recoil masses of $\ell^+\ell^-$ and $\ell^+\ell^-e^+e^-$, respectively

J/Ψ decay channel

- The alternative channel at Belle II is $e^+e^- \to X + J/\Psi$, which is only relevant to the X boson-electron coupling.

- Due to higher center-of-mass (CM) energy and J/Ψ mass, the boosted X can travel several millimeters before it decays into e^+e^-, thereby resulting in effective suppression of background.

J/Ψ associated channel

Model – with constraints

- The vectorlike interactions of the X boson with the SM fermions, f, are introduced by the effective Lagrangian:

$$\mathcal{L} \supset -eX_\mu \sum_f \varepsilon_f \bar{f}\gamma^\mu f$$

- If the new boson X is responsible for the recent Atomki anomaly via the process ${}^8\mathrm{Be} + X \to {}^8\mathrm{Be} + e^+e^-$, its mass should be $m_X \simeq 17\ \mathrm{MeV}$ and couples to the first generation quarks with the coupling strengths:

$$|\varepsilon_u + \varepsilon_d| \simeq 3.7 \times 10^{-3}\,,$$

- For the couplings to first generation quarks, the strong constraint from NA48/2 for $\pi^0 \to X\gamma$ requires protophobic condition

$$|2\varepsilon_u + \varepsilon_d| < 8 \times 10^{-4}$$

- Taking both relations into account, we finally get the preferred value for up-type and down-type quark couplings:

$$\varepsilon_u \simeq \pm 3.7 \times 10^{-3},\ \varepsilon_d \simeq \mp 7.4 \times 10^{-3},$$

- The coupling to the leptons, especially to electron, are stringently constrained by the beam dump experiment SLAC E141 and the anomalous magnetic moment of the electron $(g-2)$

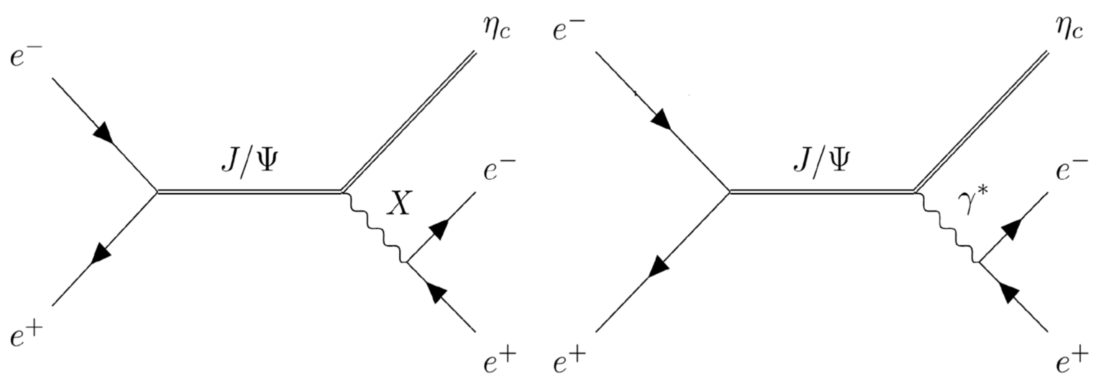
$$4.2 \times 10^{-4} \lesssim |\varepsilon_e| \lesssim 1.4 \times 10^{-3}$$

Model J/Ψ decay channel

- Here, we assume the X boson has universal coupling to each quark generation, so that if $2m_e \lesssim m_X \lesssim m_{J/\Psi} - m_{\eta_c} \simeq 113 \text{ MeV}$, the decay process $J/\Psi \to \eta_c + X \to \eta_c + e^+e^-$ is kinematic allowed and can be used to search for the 17 MeV or other light vector bosons.

- Lepton colliders such as **BESIII** and Belle II can copiously produce J/Ψ and therefore are sensitive to the J/Ψ rare decay channels.

- The dominating background comes from the off-shell photon contribution, $J/\Psi \to \eta_c + \gamma^* \to \eta_c + e^+e^-$



Model J/Ψ decay channel BESIII

- In order to exclusively reconstruct the $J/\Psi \to \eta_c e^+ e^-$ decays, we have to consider η_c decay modes that can be fully reconstructed with a reasonable background contamination.

TABLE I. The branching fractions of η_c decay modes with corresponding efficiencies.

Branching Ratio	Detection efficiency
$\mathrm{Br}(\eta_c \to K^+ K^- \pi^0) = (1.15 \pm 0.12)\%$	18.82%
$\mathrm{Br}(\eta_c \to K_S^0 K^\pm \pi^\mp) = (2.60 \pm 0.21)\%$	21.22%
$\mathrm{Br}(\eta_c \to 2(\pi^+ \pi^- \pi^0)) = (15.2 \pm 1.8)\%$	3.07%

- The overall efficiency ϵ of the above three η_c modes is obtained by adding the individual efficiencies weighted by their corresponding branching fractions: $\epsilon = 1.23\%$
- And taking 17 MeV X boson for $^8\mathrm{Be}^*$ anomaly, we can summarize the significance in Table II

$$S = N_{J/\Psi} \times \frac{\int_{(m_X-\sigma_m)^2}^{(m_X+\sigma_m)^2} dq^2 \frac{d\Gamma_{\eta_c X^*}}{dq^2}}{\Gamma_{J/\Psi}}$$

$$B = N_{J/\Psi} \times \frac{\int_{(m_X-\sigma_m)^2}^{(m_X+\sigma_m)^2} dq^2 \frac{d\Gamma_{\eta_c \gamma^*}}{dq^2}}{\Gamma_{J/\Psi}}$$

$$\Gamma_{J/\Psi} = 92.9\,\mathrm{keV}$$

TABLE II. For $N_{J/\Psi} = 10^{11}$, and chose favoured parameters $\varepsilon_c = \varepsilon_u = 3.7 \times 10^{-3}$, $\varepsilon_e = 10^{-3}$, $m_X = 17$ MeV for $^8\mathrm{Be}^*$ anomaly, the significances of signal to background from $J/\Psi \to \eta_c e^+ e^-$ with various energy resolutions of detector and 1.23% η_c reconstruction efficiency.

	σ_m=1 MeV	σ_m=2 MeV	σ_m=5 MeV	σ_m=10 MeV	σ_m=15 MeV
S	188	263	277	277	277
B	3686	7399	18989	42436	87640
$S/\sqrt{B}$	3.10	3.06	2.01	1.34	0.94

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Model J/Ψ decay channel BESIII

- For general light vector boson searches through $J/\Psi \to \eta_c e^+ e^-$, we show the variation of the expected significance over $(m_X, \varepsilon_c, \varepsilon_e)$
- The red boxes indicate the preferred regions for $^8\mathrm{Be}^*$ anomaly.

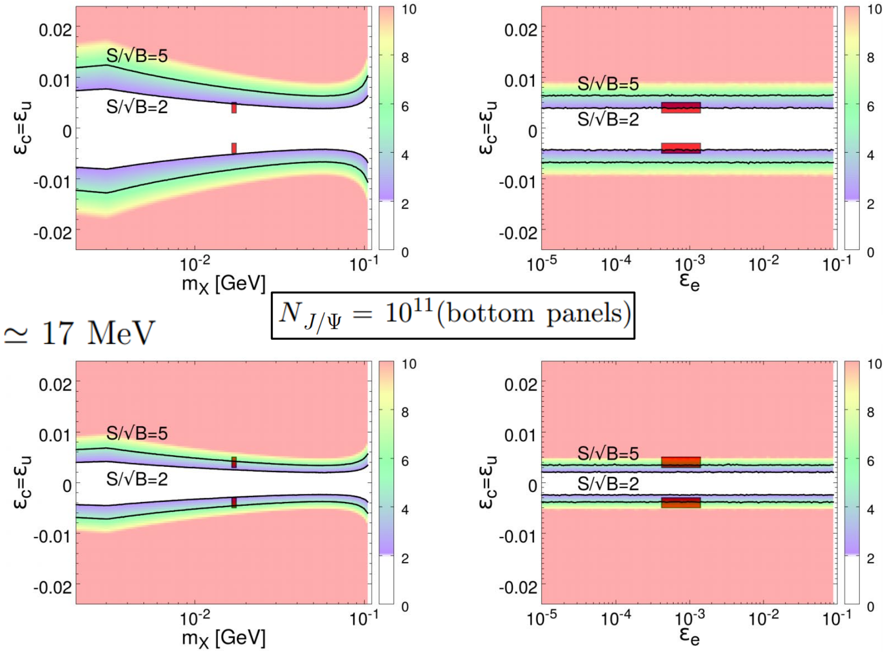
$|\varepsilon_c| \gtrsim 5 \times 10^{-3}$ at $m_X \simeq 17$ MeV

$N_{J/\Psi} = 10^{10}$(upper panels)

$|\varepsilon_c| \gtrsim 3 \times 10^{-3}$ at $m_X \simeq 60$ MeV

$|\varepsilon_c| \gtrsim 3 \times 10^{-3}$ at $m_X \simeq 17$ MeV

$N_{J/\Psi} = 10^{11}$(bottom panels)



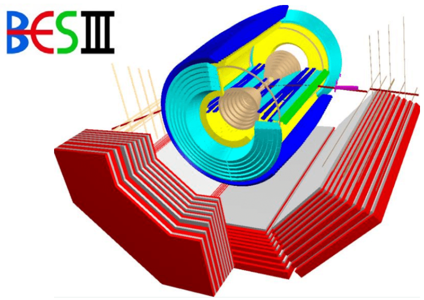
- The alternative way to reconstruct η_c from $J/\Psi \to \eta_c e^+ e^-$ at **BESIII** is using the recoil of e^+e^- .
- With an improvement of low-energy electron identification in the future, the **BESIII** with $N_{J/\Psi} = 10^{11}$ can reach the sensitivity of $|\varepsilon_c| \simeq 10^{-3}$.

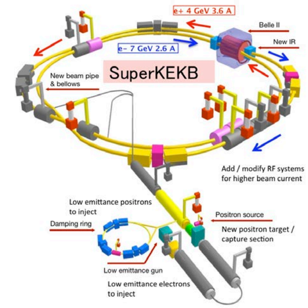
Model

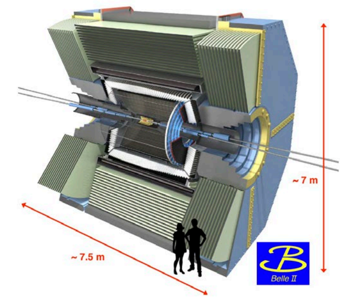
J/Ψ decay channel

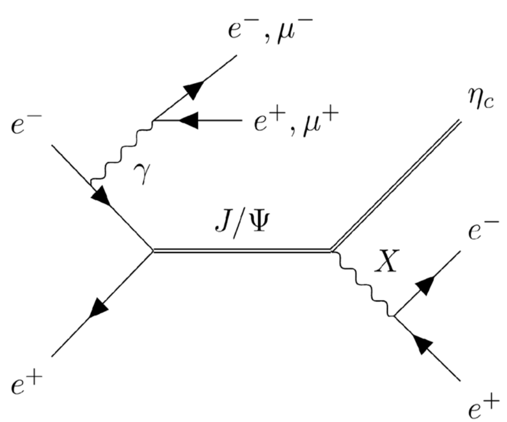
Belle II

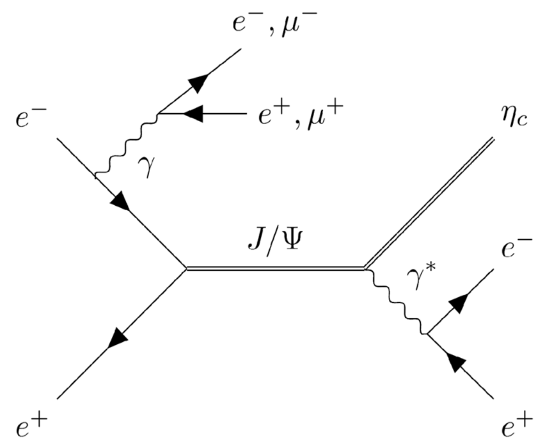
Cross section estimation











For vector meson J/Ψ, the partial width is given by the formula

$$\Gamma_{J/\Psi \to e^+e^-} = \frac{g^2_{J/\Psi ee}}{12\pi} m_{J/\Psi} \left(1 + \frac{2m_e^2}{m^2_{J/\Psi}}\right) \sqrt{1 - \frac{4m_e^2}{m^2_{J/\Psi}}}$$

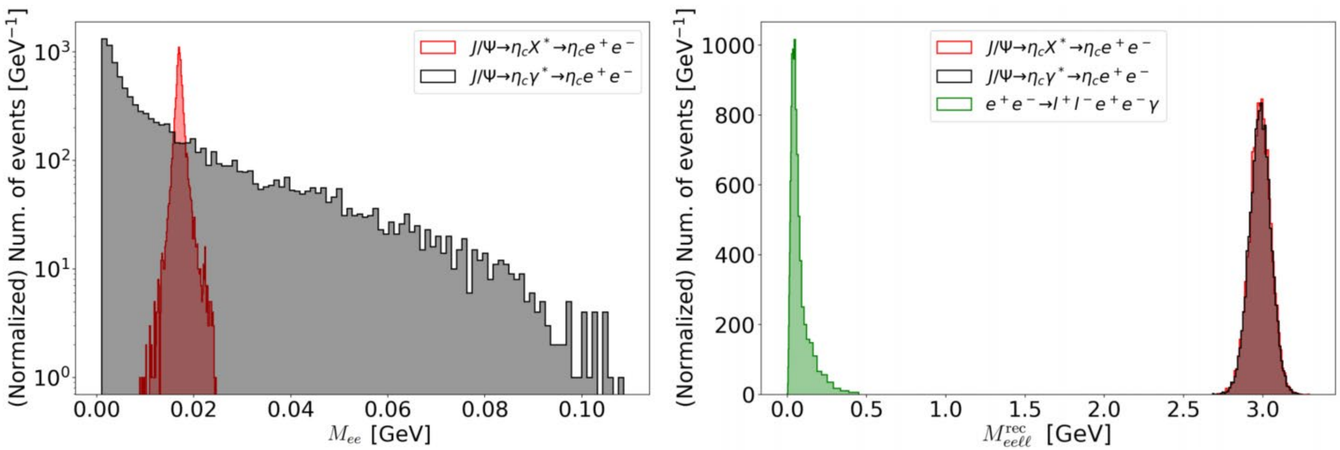
$$\Gamma_{J/\Psi \to e^+e^-} = 5.53 \ \text{keV} \Longrightarrow g_{J/\Psi ee} = 8.2048 \times 10^{-3}$$

Model J/Ψ decay channel Belle II

With the design integrated luminosity $L = 50\,\mathrm{ab}^{-1}$, we estimate $N_{J/\Psi} = 1.75 \times 10^7$ events for $e^+e^- \to \gamma^* + J/\Psi \to \ell^+\ell^- J/\Psi$ at Belle II.

$$S = L \times \sigma(e^+e^- \to \ell^+\ell^- J/\Psi) \times \mathrm{Br}(J/\Psi \to \eta_c X* \to \eta_c e^+e^-) \simeq 28.2 \left(\frac{\varepsilon_c}{10^{-2}}\right)^2$$

$$B = L \times \sigma(e^+e^- \to \ell^+\ell^- J/\Psi) \times \mathrm{Br}(J/\Psi \to \eta_c \gamma* \to \eta_c e^+e^-) \simeq 1772$$



MG5_aMC@NLO
+
FeynRules

- The e^+e^- invariant mass (left) and $e^+e^-\ell^+\ell^-$ recoil mass (right) distributions for the parton level Monte-Carlo simulation data with the smearing effect.

- We give the Gaussian smearing effect with the momentum resolution

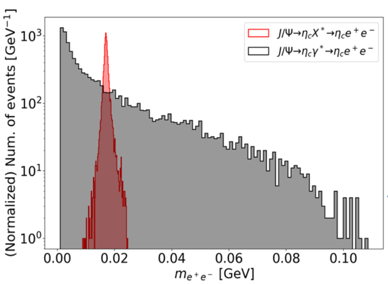
$$\sigma_{p_{\ell^\pm}}/p_{\ell^\pm} = 0.005$$

on the parton level data for our event analysis.

Model J/Ψ decay channel Belle II

[Baseline Cuts]

- To simulate the effects of the Belle II detector, we apply the following baseline cuts: $|\eta^*_{l^\pm}| \le 1.60$ in the CM frame, $|E_{\mu^\pm}| \ge 0.6$ GeV, and $|E_{e^\pm}| \ge 0.06$ GeV in the lab frame.



Cuts	B	S
Processes	$\eta_c\gamma^* \to \eta_c ee$	$\eta_c X \to \eta_c ee$
	100000	100000
Baseline Cuts	7170	6290
$\lvert M^{\rm rec}_{ee\ell\ell} - m_{\eta_c}\rvert \le 200$ MeV	7071	6219
$\lvert M_{ee} - m_X\rvert \le 2$ MeV	**377**	**5880**

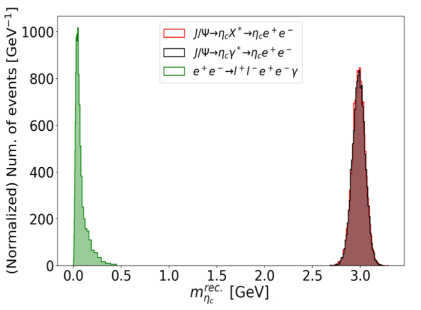


Table 3. Signal and background events of $e^+e^- \to \ell^+\ell^-e^+e^-\eta_c$ after cuts at Belle II.

- Sensitivity

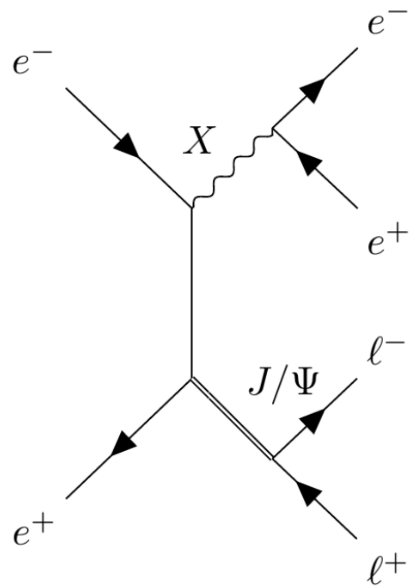
Luminosity	50 ab^{-1}	100 ab^{-1}	200 ab^{-1}
$\lvert\varepsilon_c\rvert$	$\gtrsim 1.76\times10^{-2}$	$\gtrsim 1.48\times10^{-2}$	$\gtrsim 1.24\times10^{-2}$

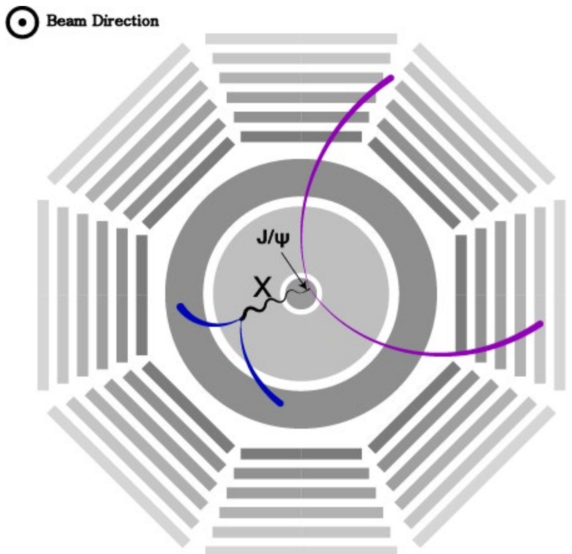
Table 4. Sensitivities on ε_c of 17 MeV X boson from $\ell^+\ell^-J/\Psi \to \ell^+\ell^-e^+e^-\eta_c$ search at Belle II with luminosities 50, 100, and 200 ab^{-1}. Here we require $S/\sqrt{B} = 2$.

Model J/Ψ associated channel Belle II

$$e^+e^- \to X + J/\Psi \to e^+e^- + J/\Psi$$

- The X boson can be boosted from the process $e^+e^- \to X + J/\Psi$ and travels several millimeters before decaying into e^+e^- in the Belle II detector.





- If the displaced vertex is between $2\ \mathrm{mm} \leq d_{xy} \leq 8\ \mathrm{mm}$*, which is inside the beam pipe, and outside the interaction region, it provides excellent vertex reconstruction and almost free from SM backgrounds.

* M. Duerr, T. Ferber, C. Hearty, F. Kahlhoefer, K. Schmidt-Hoberg and P. Tunney, JHEP 304 2002, 039 (2020)

Model J/Ψ associated channel Belle II

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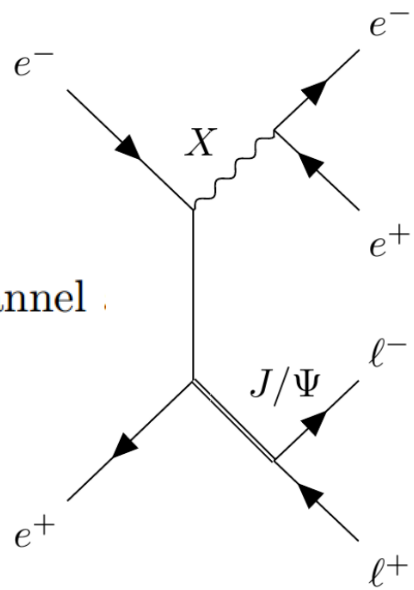
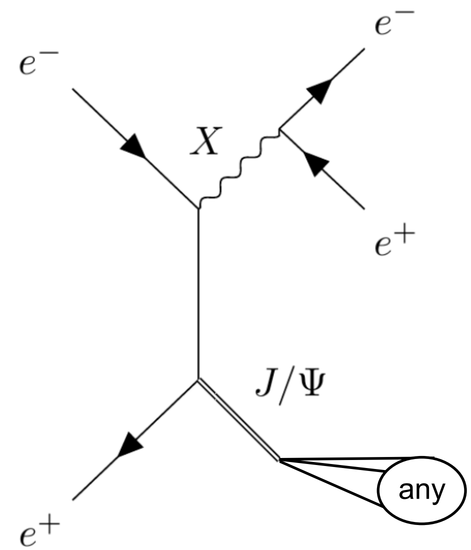
- With the baseline cuts and $2\ \mathrm{mm} \leq d_{xy} \leq 8\ \mathrm{mm}$, we estimate the signal sensitivity by considering two cases:
  (i) explicitly reconstructing $J/\Psi \to \ell^+\ell^-$ ($e^+e^-\ell^+\ell^-$ channel)
  (ii) using the recoil mass of $X \to e^+e^-$ to infer J/Ψ (e^+e^- channel)

$e^+e^-\ell^+\ell^-$ channel

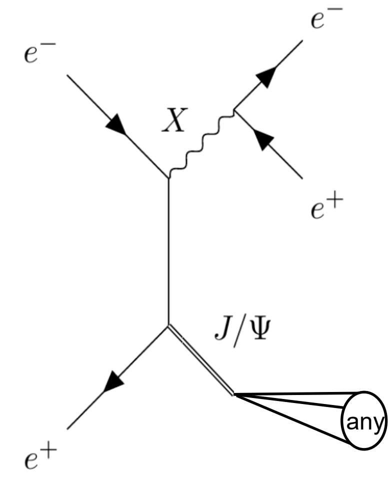




e^+e^- channel

$$\sigma(e^+e^- \to X + J/\Psi) = 2.77 \times 10^{-2} \times \left(\frac{\varepsilon_e}{10^{-3}}\right)^2\ \mathrm{fb}$$

Model J/Ψ associated channel Belle II



e^+e^- channel

- The displaced e^+e^- vertex searches can probe the 17 MeV X boson in the region

$$2.5 \times 10^{-4} \leq \varepsilon_e \leq 8.0 \times 10^{-4}$$

with significance larger than **2** by assuming near-zero background $B = 0.1$ and it covers the ε_e region preferred by Atomki.

$\varepsilon_e/10^{-4}$	8.0	7.0	5.0	4.0	3.0	2.0	1.0
Baseline Cuts(%)	17.6	17.6	17.6	17.6	17.6	17.6	17.6
$2\text{mm} < d_{xy} < 8\text{mm}(\%)$	1.6	5.3	12.3	12.9	7.4	2.3	0.5
N_S	14.6	35.7	42.7	28.7	9.23	1.28	0.07
Significance $(B = 0.1)$	$> 5\sigma$					2.2σ	0.4σ
Significance $(B = 1)$	$> 5\sigma$					1.6σ	0.9σ

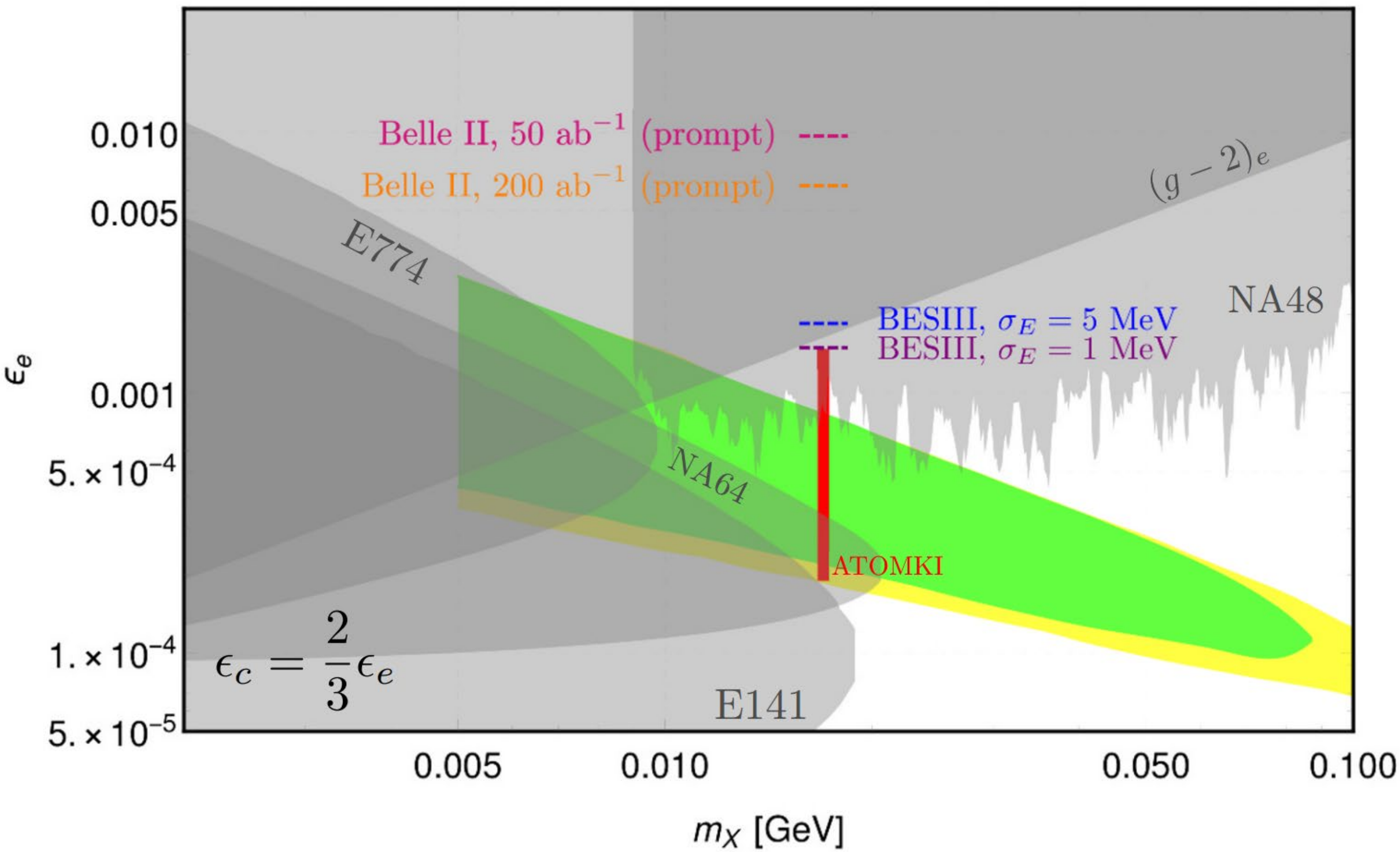
Table 6. The same as table 5, but using the e^+e^- channel.

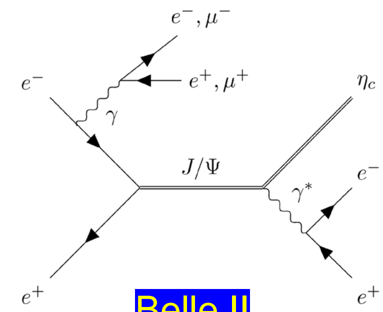
- While we expect less than **one** signal event with the currently available Belle data sample of $1\ \text{ab}^{-1}$, we can start exploring the Atomki preferred region within a few years once Belle II accumulates data sample of $10\ \text{ab}^{-1}$ or more.

Model J/Ψ decay channel J/Ψ associated channel

- The yellow (green) contour corresponds ≥ 2σ significance assuming SM background $B = 0.1$ $(B = 1)$ from e^+e^- channel at Belle II with 50 ab^{-1} luminosity, which probes rest of the favor parameter region of Atomki (red vertical band).

e^+e^- channel

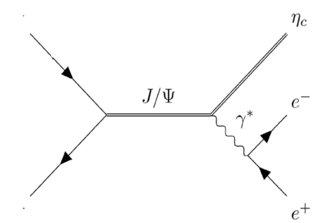


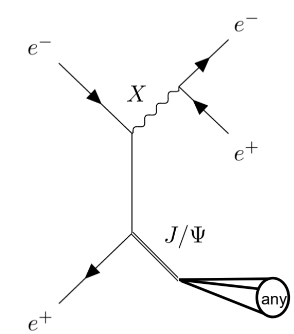


Belle II

J/Ψ decay channel

BESIII





Belle II

J/Ψ associated channel

- This study can probe the parameter region of $5\ \mathrm{MeV} \leq m_X \leq 100\ \mathrm{MeV}$ and $1.0\times10^{-4} \leq \varepsilon_e \leq 3\times10^{-3}$, which have not been constrained by any existing experiments.

Summary

- We propose several studies using J/Ψ at lepton colliders such as Belle II and BESIII, to search for light vector boson around the mass range suggested by the $^8\mathrm{Be}^*$ anomaly of the ATOMKI experiment.

- At BESIII, the $J/\Psi \to \eta_c X \to \eta_c e^+ e^-$ channel with the currently available sample of $N_{J/\Psi} = 10^{10}$ and effective η_c reconstruction efficiency of 1.23%, we can exclude the region $|\varepsilon_c| \gtrsim 5 \times 10^{-3}$ for $m_X = 17~\mathrm{MeV}$.

- On the other hand at Belle II with higher CM energy, we propose to study the process $e^+ e^- \to \ell^+ \ell^- J/\Psi$ followed by $J/\Psi \to \eta_c X \to \eta_c e^+ e^-$ and this channel can yield the sensitivity of $|\varepsilon_c| \gtrsim 1.8 \times 10^{-2}$ at $m_X = 17~\mathrm{MeV}$.

- Alternatively, we can study the process $e^+ e^- \to X + J/\Psi \to e^+ e^- \ell^+ \ell^-$ at Belle II and the X boson is boosted and produce displaced vertex of $X \to e^+ e^-$ which is longer than several millimeters.

- Selecting the $2~\mathrm{mm} \leq d_{xy} \leq 8~\mathrm{mm}$ window and requiring $> 2\sigma$ significance, it gives the sensitivity $2.0 \times 10^{-4} \leq |\varepsilon_e| \leq 8.0 \times 10^{-4}$ at $m_X = 17~\mathrm{MeV}$ for $50~\mathrm{ab}^{-1}$ luminosity and covers most of the favored signal region from the claimed $^8\mathrm{Be}^*$ anomaly.

- Extending the range of the X boson mass, this method can cover the unprecedented parameter space of $9~\mathrm{MeV} \leq m_X \leq 100~\mathrm{MeV}$ and $1.0 \times 10^{-4} \leq |\varepsilon_e| \leq 10^{-3}$.

Thank you for your attention

Appendix

Model J/Ψ decay channel

$\Gamma(J/\Psi \to \eta_c\gamma) = 1.5793$ keV, $m_{J/\Psi} = 3.0969$ GeV, and $m_{\eta_c} = 2.9839$ GeV

- Starting from the partial width of J/Ψ radiative decay

$$\Gamma(J/\Psi \to \eta_c\gamma) = \frac{1}{3}\frac{\alpha_{\rm EM}(m_{J/\Psi}^2 - m_{\eta_c}^2)^3}{8m_{J/\Psi}^3}|f_{VP}(0)|^2$$

- Differential widths for off-shell photon and X boson are

$$\frac{d\Gamma_{\eta_c\gamma^*}}{dq^2\Gamma_{J/\Psi\to\eta_c\gamma}} = |F_{VP}(q^2)|^2 \times F_{\rm QED}(q^2)$$

$$\frac{d\Gamma_{\eta_c X^*}}{dq^2\Gamma_{J/\Psi\to\eta_c\gamma}} = |F_{VP}(q^2)|^2 \times F_X(q^2)\,,$$

where $\Gamma_{\eta_c\gamma^*}$ and $\Gamma_{\eta_c X^*}$ are respectively partial decay widths of $J/\Psi \to \eta_c\gamma^* \to \eta_c e^+e^-$ and $J/\Psi \to \eta_c X^* \to \eta_c e^+e^-$

$$(2m_e)^2 \le q^2 = m_{e^+e^-}^2 \le (m_{J/\Psi} - m_{\eta_c})^2$$

$$F_{VP}(q^2) \equiv f_{VP}(q^2)/f_{VP}(0) = 1/(1 - \tfrac{q^2}{\Lambda^2})$$

from the pole approximation with pole mass $\Lambda = m_{\psi'} = 3.686097$ GeV for J/Ψ.

- The $F_{\rm QED}$ includes amplitude square and phase space is for off-shell photon, while F_X is adopted for off-shell X boson case as

$$F_X(q^2) = \frac{\alpha_{\rm EM}(\varepsilon_c \cdot \varepsilon_e)^2}{3\pi}\left(\frac{q^2}{[(q^2 - m_X^2)^2 + m_X^2\Gamma_X^2]}\right)$$
$$\times\left(1 - \frac{4m_e^2}{q^2}\right)^{1/2}\left(1 + \frac{2m_e^2}{q^2}\right)\left[\left(1 + \frac{q^2}{m_{J/\Psi}^2 - m_{\eta_c}^2}\right)^2 - \frac{4m_{J/\Psi}^2 q^2}{(m_{J/\Psi}^2 - m_{\eta_c}^2)^2}\right]^{3/2}$$

Model

J/Ψ decay channel

$\varepsilon_e \simeq 10^{-3}$

$\Gamma_X \ll m_X$

$$F_X(q^2) = \frac{\alpha_{\rm EM}(\varepsilon_c \cdot \varepsilon_e)^2}{3\pi}\left(\frac{q^2}{[(q^2-m_X^2)^2+m_X^2\Gamma_X^2]}\right) \times \left(1-\frac{4m_e^2}{q^2}\right)^{1/2}\left(1+\frac{2m_e^2}{q^2}\right)\left[\left(1+\frac{q^2}{m_{J/\Psi}^2-m_{\eta_c}^2}\right)^2-\frac{4m_{J/\Psi}^2q^2}{(m_{J/\Psi}^2-m_{\eta_c}^2)^2}\right]^{3/2}$$

$$\Gamma_{X\to e^+e^-} = \frac{\varepsilon_e^2\alpha_{\rm EM}m_X}{3}\left(1+\frac{2m_e^2}{m_X^2}\right)\sqrt{1-\frac{4m_e^2}{m_X^2}}$$

favoured coupling values $\varepsilon_c = \varepsilon_u = 3.7\times 10^{-3}$, $\varepsilon_e = 10^{-3}$ fixing $m_X = 17$ MeV

$$\mathrm{Br}(J/\Psi \to \eta_c X^* \to \eta_c e^+e^-) = 1.64\times 10^{-6}\left(\frac{\varepsilon_c}{10^{-2}}\right)^2$$

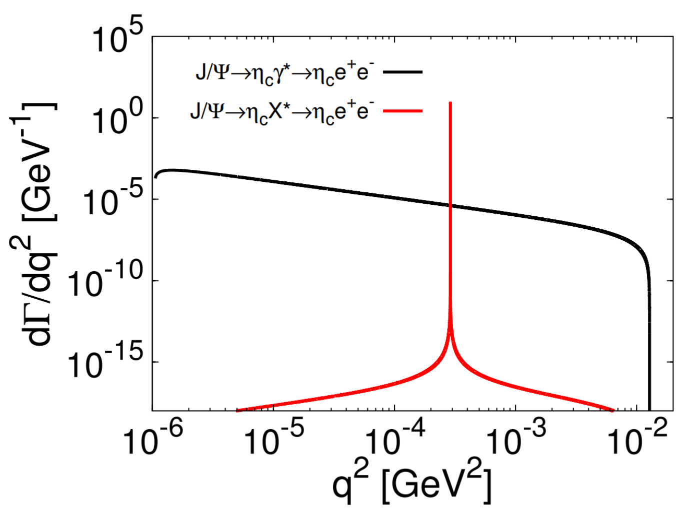
$$\mathrm{Br}(J/\Psi \to \eta_c \gamma^* \to \eta_c e^+e^-) = 1.03\times 10^{-4}$$

- The signal events and background can be statistically separated by the e^+e^- invariant mass distributions: signals show resonance peak at m_X, while this background has smooth distributions.

$$\frac{d\Gamma_{\eta_c\gamma^*}}{dq^2\Gamma_{J/\Psi\to\eta_c\gamma}} = |F_{VP}(q^2)|^2 \times F_{\rm QED}(q^2)$$

$$\frac{d\Gamma_{\eta_c X^*}}{dq^2\Gamma_{J/\Psi\to\eta_c\gamma}} = |F_{VP}(q^2)|^2 \times F_X(q^2)\,,$$

$$q^2 \subset [(m_X-\sigma_m)^2, (m_X+\sigma_m)^2]$$



$$S = N_{J/\Psi}\times\frac{\int_{(m_X-\sigma_m)^2}^{(m_X+\sigma_m)^2} dq^2\frac{d\Gamma_{\eta_c X^*}}{dq^2}}{\Gamma_{J/\Psi}}$$

$$B = N_{J/\Psi}\times\frac{\int_{(m_X-\sigma_m)^2}^{(m_X+\sigma_m)^2} dq^2\frac{d\Gamma_{\eta_c \gamma^*}}{dq^2}}{\Gamma_{J/\Psi}}$$

FIG. 1. The e^+e^- invariant mass distributions of signal $J/\Psi\to\eta_c X^*\to\eta_c e^+e^-$ and background $J/\Psi\to\eta_c\gamma^*\to\eta_c e^+e^-$, where $q^2 \equiv m_{e^+e^-}^2$. Input parameters are $m_X = 17$ MeV, $\varepsilon_c = 3.7\times10^{-3}$, and $\varepsilon_e = 10^{-3}$.

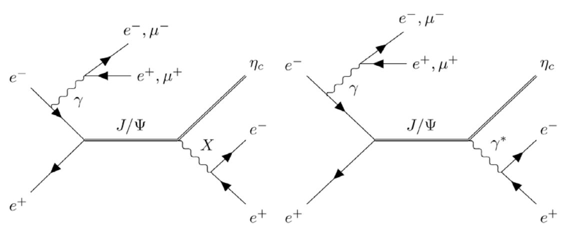
Model J/Ψ decay channel Belle II

Then the production cross sections for the background at Belle II yield

$$\sigma(e^+e^- \to \gamma^* + J/\Psi \to e^+e^- J/\Psi) = 286 \text{ fb}$$
$$\sigma(e^+e^- \to \gamma^* + J/\Psi \to \mu^+\mu^- J/\Psi) = 58.4 \text{ fb}$$

With the design integrated luminosity $L = 50\,\text{ab}^{-1}$, we estimate $N_{J/\Psi} = 1.75 \times 10^7$ events for $e^+e^- \to \gamma^* + J/\Psi \to \ell^+\ell^- J/\Psi$ at Belle II.

$$S = L \times \sigma(e^+e^- \to \ell^+\ell^- J/\Psi) \times \text{Br}(J/\Psi \to \eta_c X* \to \eta_c e^+e^-) \simeq 28.2 \left(\frac{\varepsilon_c}{10^{-2}}\right)^2$$
$$B = L \times \sigma(e^+e^- \to \ell^+\ell^- J/\Psi) \times \text{Br}(J/\Psi \to \eta_c \gamma* \to \eta_c e^+e^-) \simeq 1772$$



MG5_AMC@NLO
+
FEYNRULES

The amplitude of the electromagnetic Dalitz decay, $V \to Pe^+e^-$ can be written in a Lorentz-invariant form

$$T(V \to Pe^+e^-) = 4\pi\alpha_{\text{EM}} f_{VP} \epsilon^{\mu\nu\rho\sigma} p_\mu q_\nu \epsilon_\rho \frac{1}{q^2} \bar{u}_1 \gamma_\sigma \nu_2$$

$$\mathcal{L} \supset f_{VP}(-2\sqrt{\pi\alpha_{\text{EM}}}\partial_\mu P \partial_\nu V_\rho \epsilon^{\mu\nu\rho\sigma} A_\sigma - g_{Xc}\partial_\mu P \partial_\nu V_\rho \epsilon^{\mu\nu\rho\sigma} X_\sigma) - g_{eV}\bar{e}\gamma^\mu e V_\mu - g_{Xe}\bar{e}\gamma^\mu e X_\mu$$

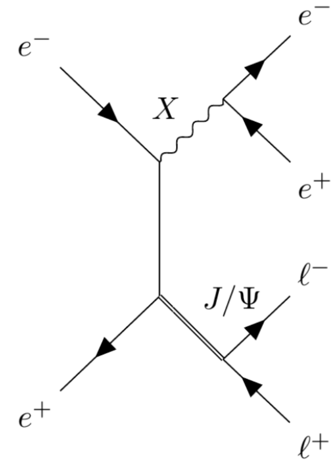
Model J/Ψ associated channel Belle II

- The signal efficiencies and expected significances for various assumed values of ε_e , according to the $50\ \mathrm{ab}^{-1}$ luminosity at Belle II

$J/\psi(1S)$ $I^G(J^{PC}) = 0^-(1^{--})$

Mass $m = 3096.900 \pm 0.006$ MeV
Full width $\Gamma = 92.9 \pm 2.8$ keV (S = 1.1)
$\Gamma_{ee} = 5.55 \pm 0.14 \pm 0.02$ keV

$J/\psi(1S)$ DECAY MODES		Fraction (Γ_i/Γ)	Scale factor/ Confidence level	p (MeV/c)
hadrons		(87.7 ± 0.5) %		–
virtual $\gamma \to$ hadrons		(13.50 ± 0.30) %		–
ggg		(64.1 ± 1.0) %		–
γgg		(8.8 ± 1.1) %		–
e^+e^-		(5.971 ± 0.032) %		1548
$e^+e^-\gamma$	[rraa]	$(8.8 \pm 1.4) \times 10^{-3}$		1548
$\mu^+\mu^-$		(5.961 ± 0.033) %		1545



$e^+e^-\ell^+\ell^-$ channel

TABLE V. The percentage pass baseline cuts mentioned in Section.IV B and the flight distance $2\ \mathrm{mm} \leq d_{xy} \leq 8$ mm cuts for several values of ε_e of 17 MeV X boson. N_S is the number of signal events from the $e^+e^-\ell^+\ell^-$ channel at Belle II with $50\ \mathrm{ab}^{-1}$ luminosity, and $\mathbb{S}_{B=1}$ ($\mathbb{S}_{B=0.1}$) is the expected significance assuming 1 (0.1) event of background in the analysis channel after all cuts.

$\varepsilon_e/10^{-4}$	8.0	7.0	6.0	5.0	4.5	4.0	3.0
Baseline Cuts(%)	13.8	13.8	13.8	13.8	13.8	13.8	13.8
$2\mathrm{mm} < d_{xy} < 8\mathrm{mm}$ (%)	1.5	4.7	7.4	10.1	11.0	10.1	5.2
N_S	1.60	3.85	4.42	4.18	3.69	2.69	0.78
$\mathbb{S}_{B=0.1}$	2.4σ	4.6σ	5.0σ	4.8σ	4.5σ	3.6σ	1.5σ
$\mathbb{S}_{B=1}$	1.6σ	2.9σ	3.2σ	3.1σ	2.8σ	2.3σ	1.2σ