



**CLUSTER OF EXCELLENCE**  
QUANTUM UNIVERSE

# Shift symmetries of flavorful axions

Quentin Bonnefoy  
(DESY Hamburg)

*LIO conference 2022 (IP2I Lyon)*  
*21/06/2022*

arXiv:2206.04182 [hep-ph]  
w/ C. Grojean and J. Kley

Why axions/ALPs ?

# Why axions/ALPs ?

- **unavoidable** : one of few light and interacting BSM modes, predicted by G/H, higher-dimensional models, string EFTs, ...

# Why axions/ALPs ?

- **unavoidable** : one of few light and interacting BSM modes, predicted by G/H, higher-dimensional models, string EFTs, ...
- **useful** : strong CP problem, dark matter, anomalies ...

# Why axions/ALPs ?

- **unavoidable** : one of few light and interacting BSM modes, predicted by G/H, higher-dimensional models, string EFTs, ...
- **useful** : strong CP problem, dark matter, anomalies ...
- **testable** : light and (weakly) interacting, many indirect effects (astro/cosmo/flavor/colliders/...), impressive experimental program, many sources, ...

# Why axions/ALPs ?

- **unavoidable** : one of few light and interacting BSM modes, predicted by G/H, higher-dimensional models, string EFTs, ...
- **useful** : strong CP problem, dark matter, anomalies ...
- **testable** : light and (weakly) interacting, many indirect effects (astro/cosmo/flavor/colliders/...), impressive experimental program, many sources, ...

**Theory effort** : studies of the UV landscape (light/heavy/decoupled/...) hand in hand with studies of **generic parametrizations** (EFTs)

# SM-axion EFT

# SM-axion EFT

**Goldstone boson :**  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $\partial_\mu a \bar{Q}_L \gamma^\mu Q_L$  , ...



# SM-axion EFT

**Goldstone boson** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $\partial_\mu a \bar{Q}_L \gamma^\mu Q_L$  , ...

QCD axion mass, ALP potential, no global symmetries... : **no such thing as a massless GB**. Shift (Peccei-Quinn) symmetries are **approximate**

# SM-axion EFT

**Goldstone boson :**  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $\partial_\mu a \bar{Q}_L \gamma^\mu Q_L$  , ...

QCD axion mass, ALP potential, no global symmetries... : **no such thing as a massless GB.** Shift (Peccei-Quinn) symmetries are **approximate**

**Light singlet :**  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $a \bar{Q}_L u_R \tilde{H}$  ,  $a |H|^2$  , ...

# SM-axion EFT

**Goldstone boson :**  $\mathcal{L} \supset a F^a \tilde{F}^a , \partial_\mu a \bar{Q}_L \gamma^\mu Q_L , \dots$

QCD axion mass, ALP potential, no global symmetries... : **no such thing as a massless GB.** Shift (Peccei-Quinn) symmetries are **approximate**

**Light singlet :**  $\mathcal{L} \supset a F^a \tilde{F}^a , a \bar{Q}_L u_R \tilde{H} , a |H|^2 , \dots$

**Measure of the breaking ?**

# SM-axion EFT

**Goldstone boson** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $\partial_\mu a \bar{Q}_L \gamma^\mu Q_L$  , ...

QCD axion mass, ALP potential, no global symmetries... : **no such thing as a massless GB**. Shift (Peccei-Quinn) symmetries are **approximate**

**Light singlet** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $a \bar{Q}_L u_R \tilde{H}$  ,  $a |H|^2$  , ...

**Measure of the breaking ?**

- perturbative :  $\lambda a |H|^2$  ,  $c_g a G^2$  , ...

# SM-axion EFT

**Goldstone boson** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $\partial_\mu a \bar{Q}_L \gamma^\mu Q_L$  , ...

QCD axion mass, ALP potential, no global symmetries... : **no such thing as a massless GB**. Shift (Peccei-Quinn) symmetries are **approximate**

**Light singlet** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $a \bar{Q}_L u_R \tilde{H}$  ,  $a |H|^2$  , ...

**Measure of the breaking ?**

- perturbative :  $\lambda a |H|^2$  ,  $c_g a G^2$  , ...
- non-perturbative :  $\tilde{c}_g a G \tilde{G}$  , ...

# SM-axion EFT

**Goldstone boson** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $\partial_\mu a \bar{Q}_L \gamma^\mu Q_L$  , ...

QCD axion mass, ALP potential, no global symmetries... : **no such thing as a massless GB**. Shift (Peccei-Quinn) symmetries are **approximate**

**Light singlet** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $a \bar{Q}_L u_R \tilde{H}$  ,  $a |H|^2$  , ...

**Measure of the breaking ?**

- perturbative :  $\lambda a |H|^2$  ,  $c_g a G^2$  , ...
- non-perturbative :  $\tilde{c}_g a G \tilde{G}$  , ...
- **what about fermions ?**

# SM-axion EFT

**Goldstone boson** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  ~~$\partial_\mu a \bar{Q}_L \gamma^\mu Q_L$~~  , ...

QCD axion mass, ALP potential, no global symmetries... : **no such thing as a massless GB**. Shift (Peccei-Quinn) symmetries are **approximate**

**Light singlet** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $a \bar{Q}_L u_R \tilde{H}$  ,  $a |H|^2$  , ...

**Measure of the breaking ?**

- perturbative :  $\lambda a |H|^2$  ,  $c_g a G^2$  , ...
- non-perturbative :  $\tilde{c}_g a G \tilde{G}$  , ...
- **what about fermions ?**

# SM-axion EFT

**Goldstone boson** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  ~~$\partial_\mu a \bar{Q}_L \gamma^\mu Q_L$~~  , ...

QCD axion mass, ALP potential, no global symmetries... : **no such thing as a massless GB**. Shift (Peccei-Quinn) symmetries are **approximate**

**Light singlet** :  $\mathcal{L} \supset a F^a \tilde{F}^a$  ,  $a \bar{Q}_L u_R \tilde{H}$  ,  $a |H|^2$  , ...

**Measure of the breaking** ?

- perturbative :  $\lambda a |H|^2$  ,  $c_g a G^2$  , ...

- non-perturbative :  $\tilde{c}_g a G \tilde{G}$  , ...

- **what about fermions** ?

? Flavor ?  
? Field redefinitions ?



# Fermion-induced shift symmetry breaking

$$\frac{a}{f} \left( \bar{Q} \tilde{Y}_u \tilde{H} u + \bar{Q} \tilde{Y}_d H d + \bar{L} \tilde{Y}_e H e + \text{h.c.} \right)$$

# Fermion-induced shift symmetry breaking

$$\frac{a}{f} \left( \bar{Q} \tilde{Y}_u \tilde{H} u + \bar{Q} \tilde{Y}_d H d + \bar{L} \tilde{Y}_e H e + \text{h.c.} \right) \xrightarrow[\text{PQ}]{\text{if}} \frac{\partial_\mu a}{f} \sum_{\psi \in \text{SM}} \bar{\psi} c_\psi \gamma^\mu \psi$$



# Fermion-induced shift symmetry breaking

$$\frac{a}{f} \left( \bar{Q} \tilde{Y}_u \tilde{H} u + \bar{Q} \tilde{Y}_d H d + \bar{L} \tilde{Y}_e H e + \text{h.c.} \right) \xrightarrow[\text{PQ}]{\text{if}} \frac{\partial_\mu a}{f} \sum_{\psi \in \text{SM}} \bar{\psi} c_\psi \gamma^\mu \psi$$

**approximate ?**

# Fermion-induced shift symmetry breaking

$$\frac{a}{f} \left( \bar{Q} \tilde{Y}_u \tilde{H} u + \bar{Q} \tilde{Y}_d H d + \bar{L} \tilde{Y}_e H e + \text{h.c.} \right) \xrightarrow[\text{PQ}]{\text{if}} \frac{\partial_\mu a}{f} \sum_{\psi \in \text{SM}} \bar{\psi} c_\psi \gamma^\mu \psi$$

**approximate ?**

Need **order parameters** !

# Fermion-induced shift symmetry breaking

$$\frac{a}{f} \left( \bar{Q} \tilde{Y}_u \tilde{H} u + \bar{Q} \tilde{Y}_d H d + \bar{L} \tilde{Y}_e H e + \text{h.c.} \right) \xrightarrow[\text{PQ}]{\text{if}} \frac{\partial_\mu a}{f} \sum_{\psi \in \text{SM}} \bar{\psi} c_\psi \gamma^\mu \psi$$

**approximate ?**

Need **order parameters** !

$$I_u^{(1)} = \text{ReTr} \left( \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(2)} = \text{ReTr} \left( X_u \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(3)} = \text{ReTr} \left( X_u^2 \tilde{Y}_u Y_u^\dagger \right),$$

$$I_d^{(1)} = \text{ReTr} \left( \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(2)} = \text{ReTr} \left( X_d \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(3)} = \text{ReTr} \left( X_d^2 \tilde{Y}_d Y_d^\dagger \right),$$

$$I_{ud}^{(1)} = \text{ReTr} \left( X_d \tilde{Y}_u Y_u^\dagger + X_u \tilde{Y}_d Y_d^\dagger \right),$$

**PQ exact  
iff**

$$I_{ud,u}^{(2)} = \text{ReTr} \left( X_u^2 \tilde{Y}_d Y_d^\dagger + \{X_u, X_d\} \tilde{Y}_u Y_u^\dagger \right), \quad = \mathbf{0}$$

$$I_{ud,d}^{(2)} = \text{ReTr} \left( X_d^2 \tilde{Y}_u Y_u^\dagger + \{X_u, X_d\} \tilde{Y}_d Y_d^\dagger \right), \quad (+e)$$

$$I_{ud}^{(3)} = \text{ReTr} \left( X_d X_u X_d \tilde{Y}_u Y_u^\dagger + X_u X_d X_u \tilde{Y}_d Y_d^\dagger \right)$$

$$I_{ud}^{(4)} = \text{ImTr} \left( [X_u, X_d]^2 \left( [X_d, \tilde{Y}_u Y_u^\dagger] - [X_u, \tilde{Y}_d Y_d^\dagger] \right) \right)$$

$$(X \equiv Y Y^\dagger)$$

# Fermion-induced shift symmetry breaking

$$\frac{a}{f} \left( \bar{Q} \tilde{Y}_u \tilde{H} u + \bar{Q} \tilde{Y}_d H d + \bar{L} \tilde{Y}_e H e + \text{h.c.} \right) \xrightarrow[\text{PQ}]{\text{if}} \frac{\partial_\mu a}{f} \sum_{\psi \in \text{SM}} \bar{\psi} c_\psi \gamma^\mu \psi$$

**approximate ?**

Need **order parameters** !

$$I_u^{(1)} = \text{ReTr} \left( \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(2)} = \text{ReTr} \left( X_u \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(3)} = \text{ReTr} \left( X_u^2 \tilde{Y}_u Y_u^\dagger \right),$$

$$I_d^{(1)} = \text{ReTr} \left( \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(2)} = \text{ReTr} \left( X_d \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(3)} = \text{ReTr} \left( X_d^2 \tilde{Y}_d Y_d^\dagger \right),$$

$$\tilde{Y}_u = i(Y_u c_u - c_Q Y_u) \quad \text{ReTr} \left( X_d \tilde{Y}_u Y_u^\dagger + X_u \tilde{Y}_d Y_d^\dagger \right),$$

$$[c_Q, Y_u Y_u^\dagger] = i(\tilde{Y}_u Y_u^\dagger + \text{h.c.}) \quad \text{Tr} \left( X_u^2 \tilde{Y}_d Y_d^\dagger + \{X_u, X_d\} \tilde{Y}_u Y_u^\dagger \right), \quad = \mathbf{0}$$

$$\quad \text{Tr} \left( X_d^2 \tilde{Y}_u Y_u^\dagger + \{X_u, X_d\} \tilde{Y}_d Y_d^\dagger \right), \quad (+e)$$

$$I_{ud}^{(3)} = \text{ReTr} \left( X_d X_u X_d \tilde{Y}_u Y_u^\dagger + X_u X_d X_u \tilde{Y}_d Y_d^\dagger \right)$$

$$I_{ud}^{(4)} = \text{ImTr} \left( [X_u, X_d]^2 \left( [X_d, \tilde{Y}_u Y_u^\dagger] - [X_u, \tilde{Y}_d Y_d^\dagger] \right) \right)$$

$$(X \equiv Y Y^\dagger)$$

# Fermion-induced shift symmetry breaking

$$\frac{a}{f} \bar{Q}_1 \tilde{y}_{uc} \tilde{H} u_2 \quad \mathbf{PQ} ?$$

$$I_u^{(1)} = \text{ReTr} \left( \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(2)} = \text{ReTr} \left( X_u \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(3)} = \text{ReTr} \left( X_u^2 \tilde{Y}_u Y_u^\dagger \right),$$

$$I_d^{(1)} = \text{ReTr} \left( \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(2)} = \text{ReTr} \left( X_d \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(3)} = \text{ReTr} \left( X_d^2 \tilde{Y}_d Y_d^\dagger \right),$$

$$I_{ud}^{(1)} = \text{ReTr} \left( X_d \tilde{Y}_u Y_u^\dagger + X_u \tilde{Y}_d Y_d^\dagger \right),$$

**PQ exact  
iff**

$$I_{ud,u}^{(2)} = \text{ReTr} \left( X_u^2 \tilde{Y}_d Y_d^\dagger + \{X_u, X_d\} \tilde{Y}_u Y_u^\dagger \right), \quad = \mathbf{0}$$

$$I_{ud,d}^{(2)} = \text{ReTr} \left( X_d^2 \tilde{Y}_u Y_u^\dagger + \{X_u, X_d\} \tilde{Y}_d Y_d^\dagger \right), \quad (+e)$$

$$I_{ud}^{(3)} = \text{ReTr} \left( X_d X_u X_d \tilde{Y}_u Y_u^\dagger + X_u X_d X_u \tilde{Y}_d Y_d^\dagger \right)$$

$$I_{ud}^{(4)} = \text{ImTr} \left( [X_u, X_d]^2 \left( [X_d, \tilde{Y}_u Y_u^\dagger] - [X_u, \tilde{Y}_d Y_d^\dagger] \right) \right)$$

$$(X \equiv Y Y^\dagger)$$



# Fermion-induced shift symmetry breaking

$\frac{a}{f} \bar{Q}_1 \tilde{y}_{uc} \tilde{H} u_2$  **PQ ?** Which flavor basis ? Fully specified ?

$$I_u^{(1)} = \text{ReTr} \left( \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(2)} = \text{ReTr} \left( X_u \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(3)} = \text{ReTr} \left( X_u^2 \tilde{Y}_u Y_u^\dagger \right),$$

$$I_d^{(1)} = \text{ReTr} \left( \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(2)} = \text{ReTr} \left( X_d \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(3)} = \text{ReTr} \left( X_d^2 \tilde{Y}_d Y_d^\dagger \right),$$

$$I_{ud}^{(1)} = \text{ReTr} \left( X_d \tilde{Y}_u Y_u^\dagger + X_u \tilde{Y}_d Y_d^\dagger \right),$$

**PQ exact  
iff**

$$I_{ud,u}^{(2)} = \text{ReTr} \left( X_u^2 \tilde{Y}_d Y_d^\dagger + \{X_u, X_d\} \tilde{Y}_u Y_u^\dagger \right), \quad = \mathbf{0}$$

$$I_{ud,d}^{(2)} = \text{ReTr} \left( X_d^2 \tilde{Y}_u Y_u^\dagger + \{X_u, X_d\} \tilde{Y}_d Y_d^\dagger \right), \quad (+e)$$

$$I_{ud}^{(3)} = \text{ReTr} \left( X_d X_u X_d \tilde{Y}_u Y_u^\dagger + X_u X_d X_u \tilde{Y}_d Y_d^\dagger \right)$$

$$I_{ud}^{(4)} = \text{ImTr} \left( [X_u, X_d]^2 \left( [X_d, \tilde{Y}_u Y_u^\dagger] - [X_u, \tilde{Y}_d Y_d^\dagger] \right) \right)$$

$$(X \equiv Y Y^\dagger)$$

# Fermion-induced shift symmetry breaking

$\frac{a}{f} \bar{Q}_1 \tilde{Y}_{uc} \tilde{H} u_2$  **PQ ?** Which flavor basis ? Fully specified ?

## Flavor-invariance

$$\begin{aligned}
 I_u^{(1)} &= \text{ReTr}(\tilde{Y}_u Y_u^\dagger), & I_u^{(2)} &= \text{ReTr}(X_u \tilde{Y}_u Y_u^\dagger), & I_u^{(3)} &= \text{ReTr}(X_u^2 \tilde{Y}_u Y_u^\dagger), \\
 I_d^{(1)} &= \text{ReTr}(\tilde{Y}_d Y_d^\dagger), & I_d^{(2)} &= \text{ReTr}(X_d \tilde{Y}_d Y_d^\dagger), & I_d^{(3)} &= \text{ReTr}(X_d^2 \tilde{Y}_d Y_d^\dagger), \\
 I_{ud}^{(1)} &= \text{ReTr}(X_d \tilde{Y}_u Y_u^\dagger + X_u \tilde{Y}_d Y_d^\dagger),
 \end{aligned}$$

**PQ exact**

$$I_{ud}^{(2)} = \text{ReTr}(X_u^2 \tilde{Y}_u Y_u^\dagger + \{X_u, X_d\} \tilde{Y}_u Y_u^\dagger) = \mathbf{0}$$

	$SU(3)_Q$	$SU(3)_u$	$SU(3)_d$	$SU(3)_L$	$SU(3)_e$
$Q_L$	<b>3</b>	<b>1</b>	<b>1</b>	<b>1</b>	<b>1</b>
$Y_u$	<b>3</b>	<b><math>\bar{3}</math></b>	<b>1</b>	<b>1</b>	<b>1</b>
$Y_d$	<b>3</b>	<b>1</b>	<b><math>\bar{3}</math></b>	<b>1</b>	<b>1</b>
$Y_e$	<b>1</b>	<b>1</b>	<b>1</b>	<b>3</b>	<b><math>\bar{3}</math></b>

$$\begin{aligned}
 & \{X_u, X_d\} \tilde{Y}_d Y_d^\dagger, & & (+e) \\
 & X_d X_u \tilde{Y}_d Y_d^\dagger) \\
 & \left[ X_u, \tilde{Y}_d Y_d^\dagger \right] - \left[ X_u, \tilde{Y}_d Y_d^\dagger \right]
 \end{aligned}$$

# Fermion-induced shift symmetry breaking

$$\frac{a}{f} \bar{Q}_1 \tilde{y}_{uc} \tilde{H} u_2$$

**PQ ?**

Which flavor basis ? Fully specified ?  
What are  $y_{uc}, \tilde{y}_{ut}, \tilde{y}_{dc}, \dots$  ?

**Flavor-invariance**

$$I_u^{(1)} = \text{ReTr} \left( \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(2)} = \text{ReTr} \left( X_u \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(3)} = \text{ReTr} \left( X_u^2 \tilde{Y}_u Y_u^\dagger \right),$$

$$I_d^{(1)} = \text{ReTr} \left( \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(2)} = \text{ReTr} \left( X_d \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(3)} = \text{ReTr} \left( X_d^2 \tilde{Y}_d Y_d^\dagger \right),$$

$$I_{ud}^{(1)} = \text{ReTr} \left( X_d \tilde{Y}_u Y_u^\dagger + X_u \tilde{Y}_d Y_d^\dagger \right),$$

$$I_{ud,u}^{(2)} = \text{ReTr} \left( X_u^2 \tilde{Y}_d Y_d^\dagger + \{X_u, X_d\} \tilde{Y}_u Y_u^\dagger \right), \quad = \mathbf{0}$$

$$I_{ud,d}^{(2)} = \text{ReTr} \left( X_d^2 \tilde{Y}_u Y_u^\dagger + \{X_u, X_d\} \tilde{Y}_d Y_d^\dagger \right), \quad (+e)$$

$$I_{ud}^{(3)} = \text{ReTr} \left( X_d X_u X_d \tilde{Y}_u Y_u^\dagger + X_u X_d X_u \tilde{Y}_d Y_d^\dagger \right)$$

$$I_{ud}^{(4)} = \text{ImTr} \left( [X_u, X_d]^2 \left( [X_d, \tilde{Y}_u Y_u^\dagger] - [X_u, \tilde{Y}_d Y_d^\dagger] \right) \right)$$

$$(X \equiv Y Y^\dagger)$$

# Fermion-induced shift symmetry breaking

$$\frac{a}{f} \bar{Q}_1 \tilde{y}_{uc} \tilde{H} u_2$$

**PQ ?**

Which flavor basis ? Fully specified ?  
What are  $y_{uc}, \tilde{y}_{ut}, \tilde{y}_{dc}, \dots$  ?

**Flavor-invariance**  
**Collective effect**

$$I_u^{(1)} = \text{ReTr}(\tilde{Y}_u Y_u^\dagger), \quad I_u^{(2)} = \text{ReTr}(X_u \tilde{Y}_u Y_u^\dagger), \quad I_u^{(3)} = \text{ReTr}(X_u^2 \tilde{Y}_u Y_u^\dagger),$$

$$I_d^{(1)} = \text{ReTr}(\tilde{Y}_d Y_d^\dagger), \quad I_d^{(2)} = \text{ReTr}(X_d \tilde{Y}_d Y_d^\dagger), \quad I_d^{(3)} = \text{ReTr}(X_d^2 \tilde{Y}_d Y_d^\dagger),$$

$$I_{ud}^{(1)} = \text{ReTr}(\tilde{Y}_u Y_u^\dagger + X_u \tilde{Y}_d Y_d^\dagger),$$

$$I_{ud,u}^{(2)} = \text{ReTr}(X_u^2 \tilde{Y}_d Y_d^\dagger + \{X_u, X_d\} \tilde{Y}_u Y_u^\dagger), \quad = \mathbf{0}$$

$$I_{ud,d}^{(2)} = \text{ReTr}(X_d^2 \tilde{Y}_u Y_u^\dagger + \{X_u, X_d\} \tilde{Y}_d Y_d^\dagger), \quad (+e)$$

$$I_{ud}^{(3)} = \text{ReTr}(X_d X_u X_d \tilde{Y}_u Y_u^\dagger + X_u X_d X_u \tilde{Y}_d Y_d^\dagger)$$

$$I_{ud}^{(4)} = \text{ImTr}([X_u, X_d]^2 ([X_d, \tilde{Y}_u Y_u^\dagger] - [X_u, \tilde{Y}_d Y_d^\dagger]))$$

$$(X \equiv Y Y^\dagger)$$

**PQ exact**  
**iff**

# Fermion-induced shift

Discussion similar for **CPV**  
**in the SM** and the  
**Jarlskog** invariant

[Jarlskog '85]

$$\frac{a}{f} \bar{Q}_1 \tilde{y}_{uc} \tilde{H} u_2$$

**PQ ?**

Which flavo  
 What are  $y_i$

**Flavor-invariance**  
**Collective effect**

$$I_u^{(1)} = \text{ReTr}(\tilde{Y}_u Y_u^\dagger), \quad I_u^{(2)} = \text{ReTr}(X_u \tilde{Y}_u Y_u^\dagger), \quad I_u^{(3)} = \text{ReTr}(X_u^2 \tilde{Y}_u Y_u^\dagger),$$

$$I_d^{(1)} = \text{ReTr}(\tilde{Y}_d Y_d^\dagger), \quad I_d^{(2)} = \text{ReTr}(X_d \tilde{Y}_d Y_d^\dagger), \quad I_d^{(3)} = \text{ReTr}(X_d^2 \tilde{Y}_d Y_d^\dagger),$$

$$I_{ud}^{(1)} = \text{ReTr}(X_d \tilde{Y}_u Y_u^\dagger + X_u \tilde{Y}_d Y_d^\dagger),$$

$$I_{ud,u}^{(2)} = \text{ReTr}(X_u^2 \tilde{Y}_d Y_d^\dagger + \{X_u, X_d\} \tilde{Y}_u Y_u^\dagger), \quad = \mathbf{0}$$

$$I_{ud,d}^{(2)} = \text{ReTr}(X_d^2 \tilde{Y}_u Y_u^\dagger + \{X_u, X_d\} \tilde{Y}_d Y_d^\dagger), \quad (+e)$$

$$I_{ud}^{(3)} = \text{ReTr}(X_d X_u X_d \tilde{Y}_u Y_u^\dagger + X_u X_d X_u \tilde{Y}_d Y_d^\dagger)$$

$$I_{ud}^{(4)} = \text{ImTr}([X_u, X_d]^2 ([X_d, \tilde{Y}_u Y_u^\dagger] - [X_u, \tilde{Y}_d Y_d^\dagger]))$$

$$(X \equiv Y Y^\dagger)$$

**PQ exact**  
**iff**

Matching to ~~PQ~~ UV models

**Invariants easily track features of ~~PQ~~**

# Matching to ~~PQ~~ UV models

## Invariants easily track features of ~~PQ~~

Ex : axiflavor/flaxion model

[Ema/Hamaguchi/Moroi/Nakayama '16  
Calibbi/Goertz/Redigolo/Ziegler/Zupan '16]

$$-\mathcal{L} = \alpha_{ij}^d \left( \frac{\phi}{M} \right)^{q_{Q_i} - q_{d_j}} \bar{Q}_i H d_j + \alpha_{ij}^u \left( \frac{\phi}{M} \right)^{q_{Q_i} - q_{u_j}} \bar{Q}_i \tilde{H} u_j + \alpha_{ij}^e \left( \frac{\phi}{M} \right)^{q_{L_i} - q_{e_j}} \bar{L}_i H e_j + \text{h.c.}$$

# Matching to ~~PQ~~ UV models

## Invariants easily track features of ~~PQ~~

Ex : axiflavor/flaxion model

[Ema/Hamaguchi/Moroi/Nakayama '16  
Calibbi/Goertz/Redigolo/Ziegler/Zupan '16]

$$-\mathcal{L} = \alpha_{ij}^d \left( \frac{\phi}{M} \right)^{q_{Q_i} - q_{d_j}} \bar{Q}_i H d_j + \alpha_{ij}^u \left( \frac{\phi}{M} \right)^{q_{Q_i} - q_{u_j}} \bar{Q}_i \tilde{H} u_j + \alpha_{ij}^e \left( \frac{\phi}{M} \right)^{q_{L_i} - q_{e_j}} \bar{L}_i H e_j + \text{h.c.}$$

$$\implies I_u^{(1)} = I_u^{(2)} = \dots = 0$$



# Matching to ~~PQ~~ UV models

## Invariants easily track features of ~~PQ~~

Ex : axiflavoron/flaxion model

[Ema/Hamaguchi/Moroi/Nakayama '16  
Calibbi/Goertz/Redigolo/Ziegler/Zupan '16]

$$-\mathcal{L} = \alpha_{ij}^d \left(\frac{\phi}{M}\right)^{q_{Q_i} - q_{d_j}} \bar{Q}_i H d_j + \alpha_{ij}^u \left(\frac{\phi}{M}\right)^{q_{Q_i} - q_{u_j}} \bar{Q}_i \tilde{H} u_j + \alpha_{ij}^e \left(\frac{\phi}{M}\right)^{q_{L_i} - q_{e_j}} \bar{L}_i H e_j + \text{h.c.}$$

$$\implies I_u^{(1)} = I_u^{(2)} = \dots = 0$$

$$\bullet -\mathcal{L}_{PQ} = \delta_{i1} \delta_{j1} \alpha' \left(\frac{\phi}{M}\right)^{q'_{Q_1} - q'_{u_1}} \bar{Q}_1 \tilde{H} u_1 + \text{h.c.} .$$

$$\implies I_u^{(1)}, I_u^{(2)}, \dots \propto q_{Q_1} - q_{u_1} - (q'_{Q_1} - q'_{u_1})$$

# Matching to ~~PQ~~ UV models

## Invariants easily track features of ~~PQ~~

Ex : axiflavoron/flaxion model

[Ema/Hamaguchi/Moroi/Nakayama '16  
Calibbi/Goertz/Redigolo/Ziegler/Zupan '16]

$$-\mathcal{L} = \alpha_{ij}^d \left(\frac{\phi}{M}\right)^{q_{Q_i} - q_{d_j}} \bar{Q}_i H d_j + \alpha_{ij}^u \left(\frac{\phi}{M}\right)^{q_{Q_i} - q_{u_j}} \bar{Q}_i \tilde{H} u_j + \alpha_{ij}^e \left(\frac{\phi}{M}\right)^{q_{L_i} - q_{e_j}} \bar{L}_i H e_j + \text{h.c.}$$

$$\implies I_u^{(1)} = I_u^{(2)} = \dots = 0$$

- $-\mathcal{L}_{PQ} = \delta_{i1} \delta_{j1} \alpha' \left(\frac{\phi}{M}\right)^{q'_{Q_1} - q'_{u_1}} \bar{Q}_1 \tilde{H} u_1 + \text{h.c.}$

$$\implies I_u^{(1)}, I_u^{(2)}, \dots \propto q_{Q_1} - q_{u_1} - (q'_{Q_1} - q'_{u_1})$$

- $\alpha_{ij}^u \left(\frac{\phi}{M}\right)^{q_{Q_i} - q_{u_j}} \bar{Q}_i \tilde{H} u_j \rightarrow \alpha_{ij}^u \left(\frac{\phi}{M}\right)^{q'_{Q_i} - q_{u_j}} \bar{Q}_i \tilde{H} u_j$

$$\implies I_u^{(1)} = \dots I_d^{(3)} = 0, \quad I_{ud}^{(1)} \propto q_{Q_1} - q_{Q'_1}$$

# Matching to PQ UV models

## Invariants easily track features of PQ

Ex : axiflavor/flaxion model

[Ema/Hamaguchi/Moroi/Nakayama '16  
Calibbi/Goertz/Redigolo/Ziegler/Zupan '16]

$$-\mathcal{L} = \alpha_{ij}^d \left(\frac{\phi}{M}\right)^{q_{Q_i} - q_{d_j}} \bar{Q}_i H d_j + \alpha_{ij}^u \left(\frac{\phi}{M}\right)^{q_{Q_i} - q_{u_j}} \bar{Q}_i \tilde{H} u_j + \alpha_{ij}^e \left(\frac{\phi}{M}\right)^{q_{L_i} - q_{e_j}} \bar{L}_i H e_j + \text{h.c.}$$

$$\implies I_u^{(1)} = I_u^{(2)} = \dots = 0$$

- $-\mathcal{L}_{PQ} = \delta_{i1} \delta_{j1} \alpha' \left(\frac{\phi}{M}\right)^{q'_{Q_1} - q'_{u_1}} \bar{Q}_1 \tilde{H} u_1 + \text{h.c.}$

$$\implies I_u^{(1)}, I_u^{(2)}, \dots \propto q_{Q_1} - q_{u_1} - (q'_{Q_1} - q'_{u_1})$$

- $\alpha_{ij}^u \left(\frac{\phi}{M}\right)^{q_{Q_i} - q_{u_j}} \bar{Q}_i \tilde{H} u_j \rightarrow \alpha_{ij}^u \left(\frac{\phi}{M}\right)^{q'_{Q_i} - q_{u_j}} \bar{Q}_i \tilde{H} u_j$

$$\implies I_u^{(1)} = \dots I_d^{(3)} = 0, \quad I_{ud}^{(1)} \propto q_{Q_1} - q_{Q'_1}$$

**Collective  
effect**

# RGEs of $\overline{\text{PQ}}$ invariants

$$\dot{I}_u^{(1)} = 2\gamma_u I_u^{(1)} + 6I_u^{(2)} - 3I_{ud}^{(1)} - 2\text{Tr}(X_u) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_u^{(2)} = 4\gamma_u I_u^{(2)} + 9I_u^{(3)} - 3I_{ud,u}^{(2)} - 2\text{Tr}(X_u^2) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_u^{(3)} = 6\gamma_u I_u^{(3)} + 12I_u^{(4)} - 3I'_u - 2\text{Tr}(X_u^3) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(1)} = 2\gamma_d I_d^{(1)} + 6I_d^{(2)} - 3I_{ud}^{(1)} + 2\text{Tr}(X_d) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(2)} = 4\gamma_d I_d^{(2)} + 9I_d^{(3)} - 3I_{ud,d}^{(2)} + 2\text{Tr}(X_d^2) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(3)} = 6\gamma_d I_d^{(3)} + 12I_d^{(4)} - 3I'_d + 2\text{Tr}(X_d^3) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud}^{(1)} = 2(\gamma_u + \gamma_d) I_{ud}^{(1)},$$

$$\dot{I}_{ud,u}^{(2)} = (4\gamma_u + 2\gamma_d) I_{ud,u}^{(2)} + 3I'_u - 6I_{ud}^{(3)} - 2\text{Tr}(X_u X_d X_u) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud,d}^{(2)} = (4\gamma_d + 2\gamma_u) I_{ud,d}^{(2)} + 3I'_d - 6I_{ud}^{(3)} + 2\text{Tr}(X_d X_u X_d) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud}^{(3)} = 4(\gamma_u + \gamma_d) I_{ud}^{(3)},$$

$$\dot{I}_{ud}^{(4)} = 6 \left( \gamma_u + \gamma_d + \frac{1}{2} \text{Tr}(X_u + X_d) \right) I_{ud}^{(4)} - \text{ImTr}([X_u, X_d]^3) (I_u^{(1)} + I_d^{(1)}).$$

# RGEs of $\overline{\text{PQ}}$ invariants

$$\dot{I}_u^{(1)} = 2\gamma_u I_u^{(1)} + 6I_u^{(2)} - 3I_{ud}^{(1)} - 2\text{Tr}(X_u) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_u^{(2)} = 4\gamma_u I_u^{(2)} + 9I_u^{(3)} - 3I_{ud,u}^{(2)} - 2\text{Tr}(X_u^2) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_u^{(3)} = 6\gamma_u I_u^{(3)} + 12I_u^{(4)} - 3I'_u - 2\text{Tr}(X_u^3) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(1)} = 2\gamma_d I_d^{(1)} + 6I_d^{(2)} - 3I_{ud}^{(1)} + 2\text{Tr}(X_d) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(2)} = 4\gamma_d I_d^{(2)} + 9I_d^{(3)} - 3I_{ud,d}^{(2)} + 2\text{Tr}(X_d^2) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(3)} = 6\gamma_d I_d^{(3)} + 12I_d^{(4)} - 3I'_d + 2\text{Tr}(X_d^3) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud}^{(1)} = 2(\gamma_u + \gamma_d) I_{ud}^{(1)},$$

$$\dot{I}_{ud,u}^{(2)} = (4\gamma_u + 2\gamma_d) I_{ud,u}^{(2)} + 3I'_u - 6I_{ud}^{(3)} - 2\text{Tr}(X_u X_d X_u) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud,d}^{(2)} = (4\gamma_d + 2\gamma_u) I_{ud,d}^{(2)} + 3I'_d - 6I_{ud}^{(3)} + 2\text{Tr}(X_d X_u X_d) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud}^{(3)} = 4(\gamma_u + \gamma_d) I_{ud}^{(3)},$$

$$\dot{I}_{ud}^{(4)} = 6 \left( \gamma_u + \gamma_d + \frac{1}{2} \text{Tr}(X_u + X_d) \right) I_{ud}^{(4)} - \text{ImTr}([X_u, X_d]^3) (I_u^{(1)} + I_d^{(1)}).$$

**Order  
parameters  
run into  
themselves**

# RGEs of $\overline{\text{PQ}}$ invariants

$$\dot{I}_u^{(1)} = 2\gamma_u I_u^{(1)} + 6I_u^{(2)} - 3I_{ud}^{(1)} - 2\text{Tr}(X_u) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_u^{(2)} = 4\gamma_u I_u^{(2)} + 9I_u^{(3)} - 3I_{ud,u}^{(2)} - 2\text{Tr}(X_u^2) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_u^{(3)} = 6\gamma_u I_u^{(3)} + 12I_u^{(4)} - 3I'_u - 2\text{Tr}(X_u^3) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(1)} = 2\gamma_d I_d^{(1)} + 6I_d^{(2)} - 3I_{ud}^{(1)} + 2\text{Tr}(X_d) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(2)} = 4\gamma_d I_d^{(2)} + 9I_d^{(3)} - 3I_{ud,d}^{(2)} + 2\text{Tr}(X_d^2) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(3)} = 6\gamma_d I_d^{(3)} + 12I_d^{(4)} - 3I'_d + 2\text{Tr}(X_d^3) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud}^{(1)} = 2(\gamma_u + \gamma_d) I_{ud}^{(1)},$$

$$\dot{I}_{ud,u}^{(2)} = (4\gamma_u + 2\gamma_d) I_{ud,u}^{(2)} + 3I'_u - 6I_{ud}^{(3)} - 2\text{Tr}(X_u X_d X_u) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud,d}^{(2)} = (4\gamma_d + 2\gamma_u) I_{ud,d}^{(2)} + 3I'_d - 6I_{ud}^{(3)} + 2\text{Tr}(X_d X_u X_d) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud}^{(3)} = 4(\gamma_u + \gamma_d) I_{ud}^{(3)},$$

$$\dot{I}_{ud}^{(4)} = 6 \left( \gamma_u + \gamma_d + \frac{1}{2} \text{Tr}(X_u + X_d) \right) I_{ud}^{(4)} - \text{ImTr}([X_u, X_d]^3) (I_u^{(1)} + I_d^{(1)}).$$

**Order  
parameters  
run into  
themselves**

$$I_u^{(4)} = \text{Tr}(X_u) I_u^{(3)}$$

$$- \frac{1}{2} \left( (\text{Tr} X_u)^2 - \text{Tr} X_u^2 \right) I_u^{(2)}$$

$$+ \frac{1}{6} \left( (\text{Tr} X_u)^3 - 3 \text{Tr} X_u^2 \text{Tr} X_u + 2 \text{Tr} X_u^3 \right) I_u^{(1)}$$

# RGEs of $\overline{\text{PQ}}$ invariants

$$\dot{I}_u^{(1)} = 2\gamma_u I_u^{(1)} + 6I_u^{(2)} - 3I_{ud}^{(1)} - 2\text{Tr}(X_u) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_u^{(2)} = 4\gamma_u I_u^{(2)} + 9I_u^{(3)} - 3I_{ud,u}^{(2)} - 2\text{Tr}(X_u^2) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_u^{(3)} = 6\gamma_u I_u^{(3)} + 12I_u^{(4)} - 3I_u' - 2\text{Tr}(X_u^3) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(1)} = 2\gamma_d I_d^{(1)} + 6I_d^{(2)} - 3I_{ud}^{(1)} + 2\text{Tr}(X_d) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(2)} = 4\gamma_d I_d^{(2)} + 9I_d^{(3)} - 3I_{ud,d}^{(2)} + 2\text{Tr}(X_d^2) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_d^{(3)} = 6\gamma_d I_d^{(3)} + 12I_d^{(4)} - 3I_d' + 2\text{Tr}(X_d^3) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud}^{(1)} = 2(\gamma_u + \gamma_d) I_{ud}^{(1)},$$

$$\dot{I}_{ud,u}^{(2)} = (4\gamma_u + 2\gamma_d) I_{ud,u}^{(2)} + 3I_u' - 6I_{ud}^{(3)} - 2\text{Tr}(X_u X_d X_u) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud,d}^{(2)} = (4\gamma_d + 2\gamma_u) I_{ud,d}^{(2)} + 3I_d' - 6I_{ud}^{(3)} + 2\text{Tr}(X_d X_u X_d) \left( I_e^{(1)} + 3(I_d^{(1)} - I_u^{(1)}) \right),$$

$$\dot{I}_{ud}^{(3)} = 4(\gamma_u + \gamma_d) I_{ud}^{(3)},$$

$$\dot{I}_{ud}^{(4)} = 6 \left( \gamma_u + \gamma_d + \frac{1}{2} \text{Tr}(X_u + X_d) \right) I_{ud}^{(4)} - \text{ImTr}([X_u, X_d]^3) (I_u^{(1)} + I_d^{(1)}).$$

**Order  
parameters  
run into  
themselves**

**Similar**

$$I_u^{(4)} = \text{Tr}(X_u) I_u^{(3)}$$

$$- \frac{1}{2} \left( (\text{Tr} X_u)^2 - \text{Tr} X_u^2 \right) I_u^{(2)}$$

$$+ \frac{1}{6} \left( (\text{Tr} X_u)^3 - 3 \text{Tr} X_u^2 \text{Tr} X_u + 2 \text{Tr} X_u^3 \right) I_u^{(1)}$$

# Applications

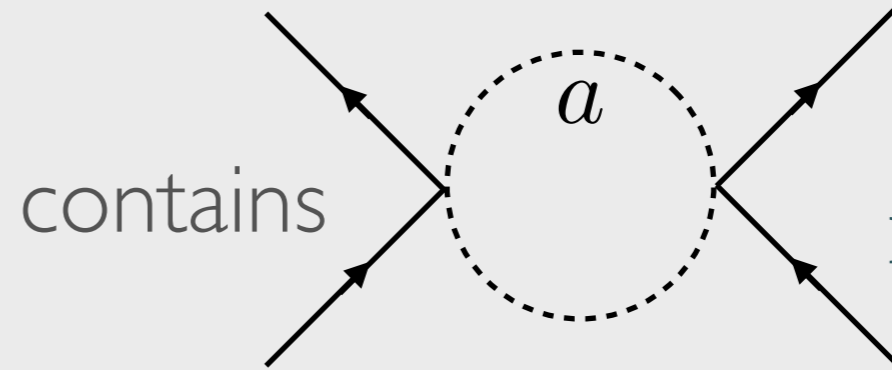
- **Algebraic formulae**



# Applications

- **Algebraic formulae**

Ex, SMEFT+axion RGEs :  $\frac{dC_{\bar{L}e\bar{d}Q}}{d\mu}$



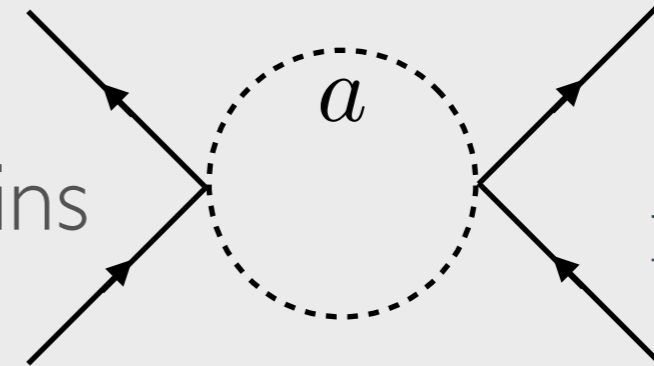
**[Galda,  
Neubert,  
Renner '21]**

# Applications

- **Algebraic formulae**

Ex, SMEFT+axion RGEs :  $\frac{dC_{\bar{L}e\bar{d}Q}}{d\mu}$

contains



[Galda,  
Neubert,  
Renner '21]

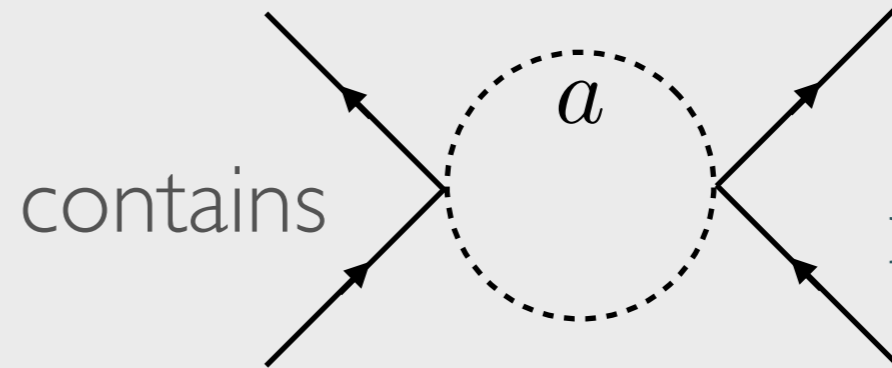
but  $\text{Re} \left( \frac{dC_{\bar{L}e\bar{d}Q}}{d\mu} Y_e^\dagger Y_d \right) \propto \cancel{PQ}$

i.e. runs as in the SMEFT for a GB

# Applications

- **Algebraic formulae**

Ex, SMEFT+axion RGEs :  $\frac{dC_{\bar{L}e\bar{d}Q}}{d\mu}$



[Galda,  
Neubert,  
Renner '21]

but  $\text{Re} \left( \frac{dC_{\bar{L}e\bar{d}Q}}{d\mu} Y_e^\dagger Y_d \right) \propto \cancel{PQ}$

i.e. runs as in the SMEFT for a GB

- **Low-energy constraints**

$$I_u^{(1)} = \text{ReTr} \left( \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(2)} = \text{ReTr} \left( X_u \tilde{Y}_u Y_u^\dagger \right), \quad I_u^{(3)} = \text{ReTr} \left( X_u^2 \tilde{Y}_u Y_u^\dagger \right),$$

$$I_d^{(1)} = \text{ReTr} \left( \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(2)} = \text{ReTr} \left( X_d \tilde{Y}_d Y_d^\dagger \right), \quad I_d^{(3)} = \text{ReTr} \left( X_d^2 \tilde{Y}_d Y_d^\dagger \right),$$

~~$$I_{ud}^{(1)} = \text{ReTr} \left( X_d \tilde{Y}_u Y_u^\dagger + X_u \tilde{Y}_d Y_d^\dagger \right),$$~~

~~$$I_{ud,u}^{(2)} = \text{ReTr} \left( X_u^2 \tilde{Y}_d Y_d^\dagger + \{X_u, X_d\} \tilde{Y}_u Y_u^\dagger \right),$$~~

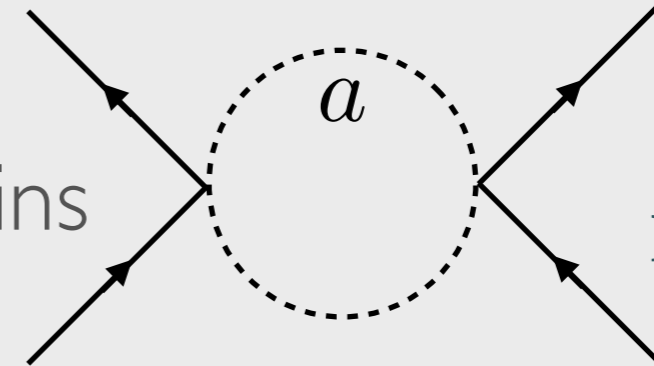
~~$$I_{ud,d}^{(2)} = \text{ReTr} \left( X_d^2 \tilde{Y}_u Y_u^\dagger + \{X_u, X_d\} \tilde{Y}_d Y_d^\dagger \right),$$~~

# Applications

- **Algebraic formulae**

Ex, SMEFT+axion RGEs :  $\frac{dC_{\bar{L}e\bar{d}Q}}{d\mu}$

contains



[Galda,  
Neubert,  
Renner '21]

but  $\text{Re} \left( \frac{dC_{\bar{L}e\bar{d}Q}}{d\mu} Y_e^\dagger Y_d \right) \propto \cancel{PQ}$

i.e. runs as in the SMEFT for a GB

- **Low-energy constraints**

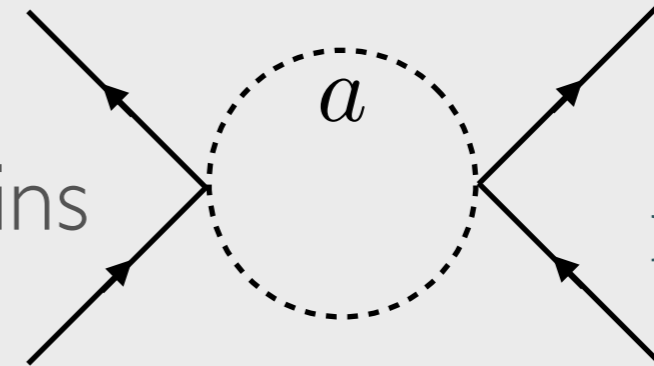
but  $\frac{d}{d\mu} \text{Re} \left( C_{\bar{u}\gamma^\mu d\bar{d}\gamma_\mu u} \left[ X_d \tilde{Y}_u Y_u^\dagger + (u \leftrightarrow d) \right] \right) = 0$  at one-loop  
&  $\mathcal{O} \left( \frac{1}{v^2 f} \right)$

# Applications

- **Algebraic formulae**

Ex, SMEFT+axion RGEs :  $\frac{dC_{\bar{L}e\bar{d}Q}}{d\mu}$

contains



[Galda,  
Neubert,  
Renner '21]

but  $\text{Re} \left( \frac{dC_{\bar{L}e\bar{d}Q}}{d\mu} Y_e^\dagger Y_d \right) \propto \cancel{PQ}$

i.e. runs as in the SMEFT for a GB

- **Low-energy constraints**

but  $\frac{d}{d\mu} \text{Re} \left( C_{\bar{u}\gamma^\mu d\bar{d}\gamma_\mu u} \left[ X_d \tilde{Y}_u Y_u^\dagger + (u \leftrightarrow d) \right] \right) = 0$  at one-loop  
&  $\mathcal{O} \left( \frac{1}{v^2 f} \right)$

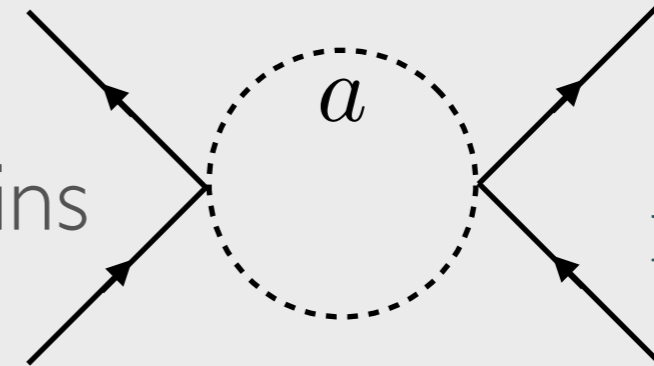
Use the constraints at low E !

# Applications

- **Algebraic formulae**

Ex, SMEFT+axion RGEs :  $\frac{dC_{\bar{L}e\bar{d}Q}}{d\mu}$

contains



[Galda,  
Neubert,  
Renner '21]

but  $\text{Re} \left( \frac{dC_{\bar{L}e\bar{d}Q}}{d\mu} Y_e^\dagger Y_d \right) \propto \cancel{PQ}$

i.e. runs as in the SMEFT for a GB

- **Low-energy constraints**

but  $\frac{d}{d\mu} \text{Re} \left( C_{\bar{u}\gamma^\mu d\bar{d}\gamma_\mu u} \left[ X_d \tilde{Y}_u Y_u^\dagger + (u \leftrightarrow d) \right] \right) = 0$  at one-loop  
&  $\mathcal{O} \left( \frac{1}{v^2 f} \right)$

Use the constraints at low E !

Ex, axion-induced EDMs :  $d_{\text{Hg}} \approx 4 \times 10^{-4} d_n$  for a GB

[Di Luzio, Gröber, Paradisi '20]

# Outlook

In the fermionic sector,  $PQ$  is a **collective effect**

# Outlook

In the fermionic sector,  $PQ$  is a **collective effect**

We captured it with **flavor-invariant order parameters**



# Outlook

In the fermionic sector,  $PQ$  is a **collective effect**

We captured it with **flavor-invariant order parameters**

This is confirmed by **matching to UV models** and obtaining **closed RGEs**

# Outlook

In the fermionic sector,  $\mathbb{PQ}$  is a **collective effect**

We captured it with **flavor-invariant order parameters**

This is confirmed by **matching to UV models** and obtaining **closed RGEs**

The invariants are algebraic : **sum-rules from shift-invariance**

# Outlook

In the fermionic sector,  $PQ$  is a **collective effect**

We captured it with **flavor-invariant order parameters**

This is confirmed by **matching to UV models** and obtaining **closed RGEs**

The invariants are algebraic : **sum-rules from shift-invariance**

## **Perspectives :**

- CPV
- flavor assumptions VS  $PQ$
- EFTs with non-linearly realized EW symmetry
- ...

THANK YOU

# Non-perturbative PQ

$$\frac{a}{f} \left( \bar{Q} \tilde{Y}_u \tilde{H} u + \bar{Q} \tilde{Y}_d H d + \bar{L} \tilde{Y}_e H e + \text{h.c.} \right) - \frac{C_g g_3^2}{16\pi^2} \frac{a}{f} G \tilde{G}$$

For a non-anomalous GB :

$$I_g = C_g - i \text{Tr} \left( Y_u^{-1} \tilde{Y}_u + (u \leftrightarrow d) \right) = 0$$

**Collective effect**

Ex : axiflavor/flaxion model,  $I_g = \text{anomaly polynomial}$

RG running :  $\frac{dI_g}{d\mu} = 0$  at one-loop