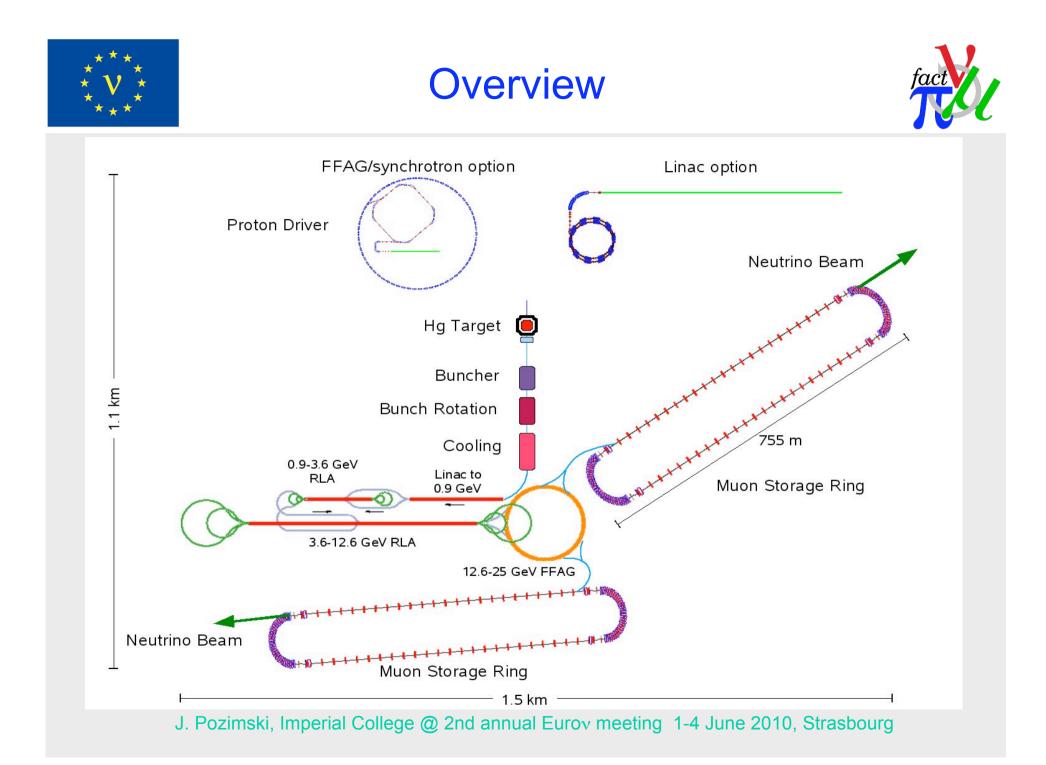




WP3 summary







R&D for the proton driver is decoupled from IDS as a hosting lab specific solution is assumed,.....but required beam parameters on target have been defined. Within Eurov the proton driver is part of the super beam work package. As the proton driver costing is strongly related to the hosting lab the following conveners have agreed to contribute:

CERN LINAC 4 / SPL : Roland Garoby

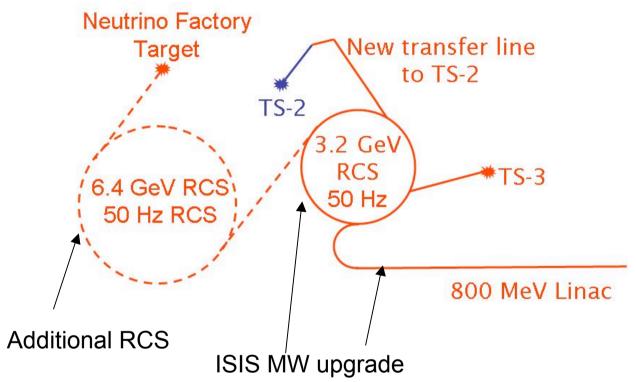
Fermilab Project X : Keith Gollwitzer

RAL - ISIS upgrade : John Thomason



NF- Proton driver and ISIS upgrade fact

Jaroslaw Pasternak / John Thomason



- Based on MW ISIS upgrade with 0.8 GeV linac and 3.2 GeV RCS.
- Assumes a sharing of the beam power at 3.2 GeV between the two facilities
- Requires additional RCS machine in order to meet the power and energy needs of the Neutrino Factory
- Both facilities can have the same ion source, RFQ, chopper, linac, H⁻ injection, accumulation and acceleration to 3.2 GeV J. Pozimski, Imperial College @ 2nd annual Eurov meeting 1-4 June 2010, Strasbourg





High power SPL could be the proton driver for a NF.

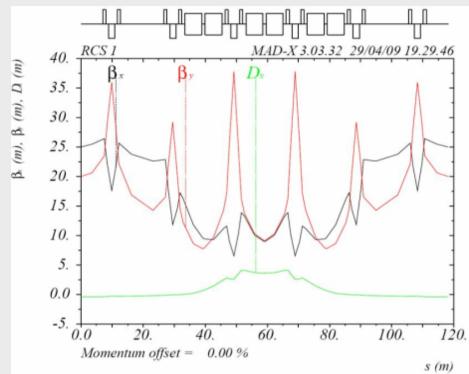
- Lattice design for accumulator and bunch compressor available.
- Detailed particle dynamics studies of accumulator and bunch compressor successfully performed.
- 3 bunch scenario is able to fulfil NF requirements.
- Costing for SPL finished, this has also to be performed for the rings.



Preliminary design of the second RCS



Jaroslaw Pasternak



- Lattice may allow for flexibility in gamma transition choice (even with beam).
- Ring is overdesigned in order to allow for 10.3 GeV.
- Optimised solution for 6.4 GeV is in preparation!

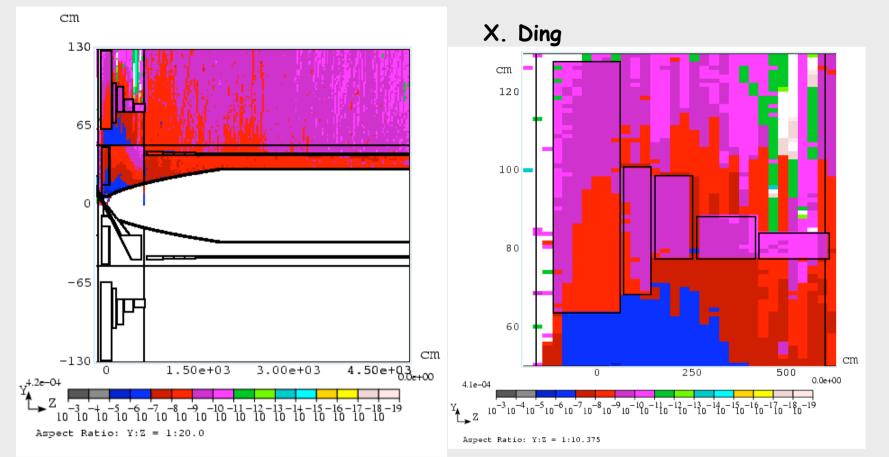
Number of superperiods	6
Circumference	708.788 m
Harmonic number	6
RF frequency	2.4717-2.5289 MHz
Betatron tunes (Q_H, Q_V)	(7.81, 7.78)
Gamma transition	7.9056
Beam power at 6.4 GeV	4 MW for 2 bunches
Bunch area	1.8 eVs
Δp/p at 3.2 GeV	5.3 10 ⁻³
Injection / extraction energy	3.2 / 6.4 [10.3] GeV
Repetition rate	50 Hz
Max B field in dipoles	1.2 T (at 10.3 GeV)
Length of long drift	12 m

Parameters of 6.4 (10.3) GeV RCS



Liquid mercury target -Energy deposition in the target area





Enhanced shield can decrease the power deposition in SC1 coil from 22.1kW to 4.8kW. By replacing the Res Sol by WC shield, the power deposition in SC1 coil can be decreased further to 1.3kW.







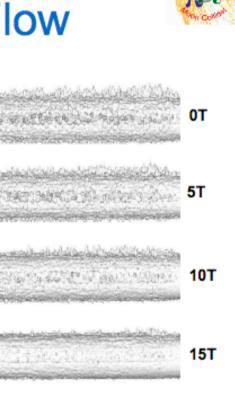
Previous Work -Jet Flow

- Free Jet + MHD + EDP
 - Numerical Study
 - R. Samulyak et al.

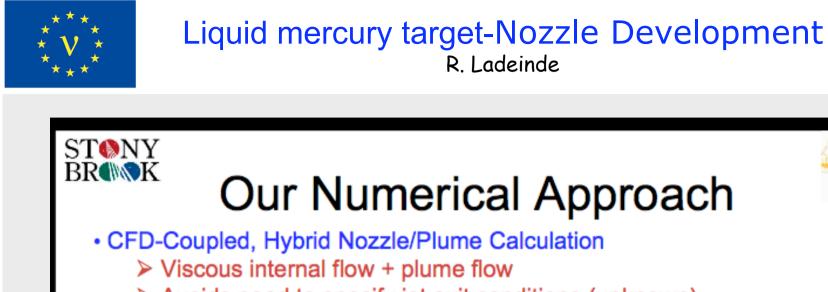
FronTier Code based on Front Tracking

- Surface instability
- MHD stabilizing effect
- Filament velocity in the simulations was about 25% smaller than the experimental value

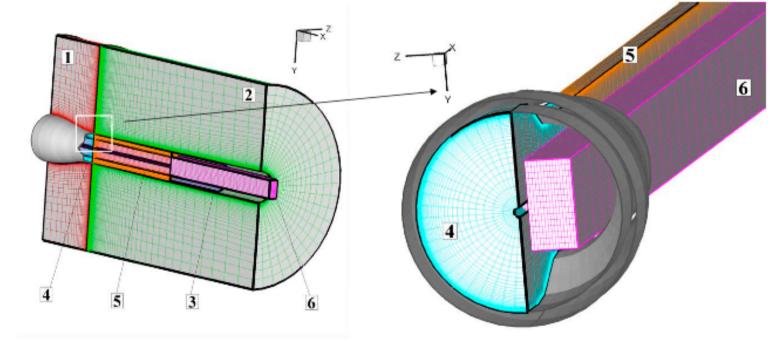
magnetic field	$5\mathrm{T}$	10T	15T
experiments	54 m/s	50 m/s	35 m/s
simulations	36 m/s	27 m/s	22 m/s



Surface filaments at 150 µs under longitudinal magnetic field

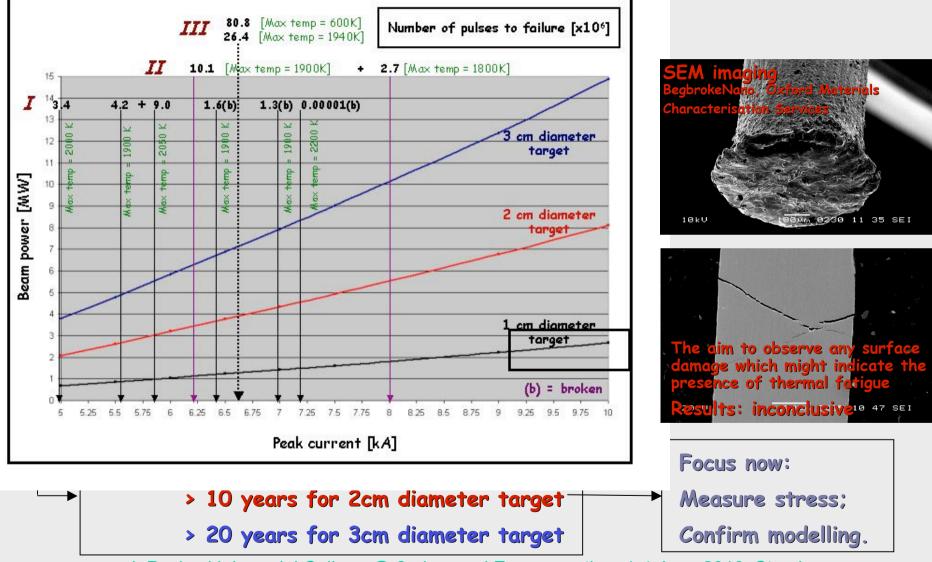


Avoids need to specify jet exit conditions (unknown)





Rob Edgecock



J. Pozimski, IrBerier de lower remperature! meeting 1-4 June 2010, Strasbourg





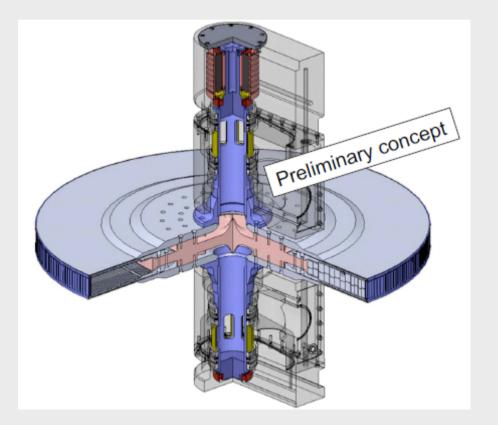


Solid target station

Rob Edgecock



- Current option: a wheel being investigated now
- Several already used, but most relevant: design study



Horizontal for compatibility with baseline target station

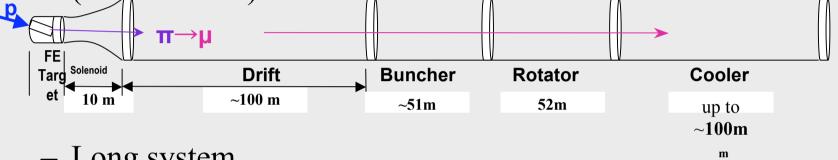


Compact Muon front end

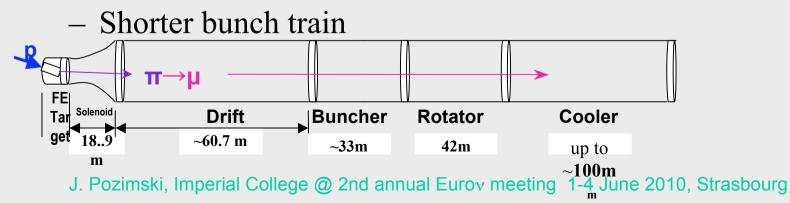




- ISS study based on $n_B = 18$ (280 MeV/c to 154 MeV/c)
 - Buncher 0 to 12MV/m; Rotator 12.5MV/m, B=1.75T (201.25 MHz)



- Long system,
- Try shorter version $n_B = 10$ (233 MeV/c to 154 MeV/c)
 - slightly lower fields (1.5T, 15MV/m)
 - Buncher 0 to 9 MV/m, Rotator 12MV/m



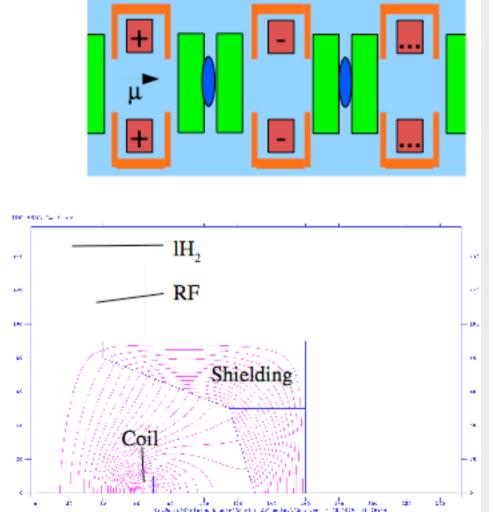


Shielded RF cooling lattice





- Increase cell length to remove RF from fringe fields
 - Add shielding using iron or bucking coils
- Look at cooling section
 - This is where the RF is most limited
 - This is where optics are most demanding
- How well can we cool in this shielded scenario?
- How well can we optimise the cooling lattice?
- Try to keep RF cavities in < 0.1 T fields
- Liquid Hydrogen absorbers



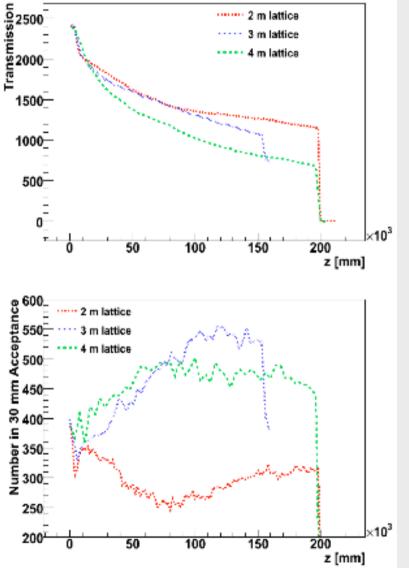


Shielded RF cooling lattice

C. Rogers

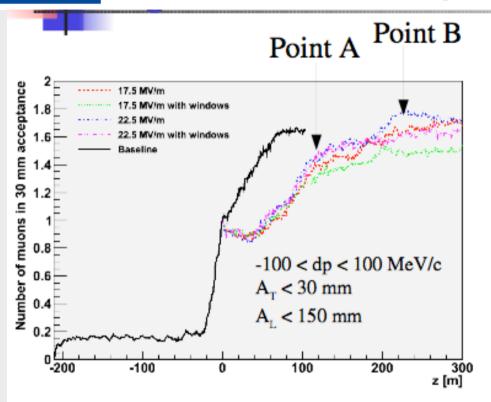


- Cell length optimisation
 - Simulated using long coil option
 - Race between RF packing fraction and β function
 - Higher RF packing => quicker cooling
 - Shorter lattice => lower β function (better equilibrium emittance)
- 3m lattice is optimal
 - Worry about initial beam loss
 - Nb low statistics
 - Get ~ 40 % with long coil (a bit more optimisation is possible)
- Case for beta tapering?



Higher momentum beam







- Fairly large transmission losses
 - >~ 50%
- Most of the remaining beam is inside the 30 mm acceptance
- Getting increase in rate of ~ 70 %
 - But with more hardware
 - Performance quite similar to baseline
- If I stop at point A I use roughly the same amount of hardware as the baseline (RF packing fraction ~ 1/2 that of the baseline)
 - And lose a few muons
- I can recover baseline performance if I go to Point B
 - But those last few muons are expensive!
 - J. Pozimski, Imperial College @ 2nd annual Eurov meeting 1-4 June 2010, Strasbourg



Front end work at CERN

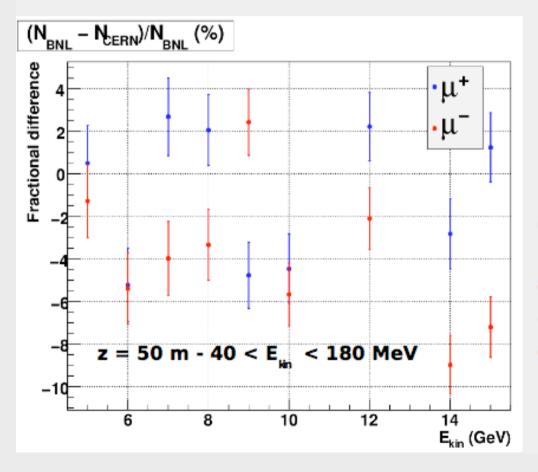
G. Prior



Particles production (coll. with BNL)

Tested the exact same input files on 2 different machines:

- SL4 2.6.9-89.0.11.EL.smp 32x little endian (BNL)
- SLC4 2.6.9-89.0.18.EL.cern 32x little endian (CERN)



Error =
$$\sqrt{(N.s_2 - s_1^2)/(N-1)}$$

 $s_2 = \sum w_i^2 \qquad s_1 = \sum w_i$

Difference in the yield HAS to come from different code versions: need to be confirmed by the MARS developers as we don't have access to the source files.

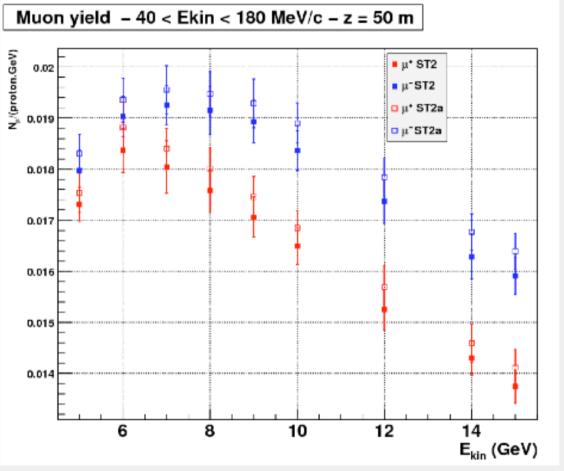


Front end work at CERN

G. Prior



ST2/ST2a comparison



Yield = MEAN Error = STD

ST2a slightly better than ST2 (within the errors bars).

Advantage of a higher field at end of taper but ST2 magnet configuratior in a more mature design.

Optimum at 7-8 GeV and yield difference between muon signs is likely to be model dependent.

June 3rd 2010



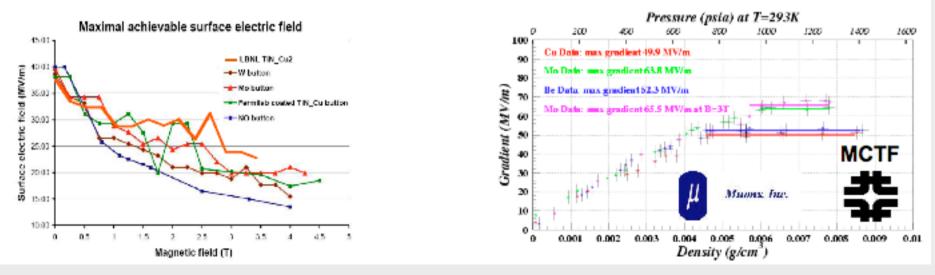
High pressure RF

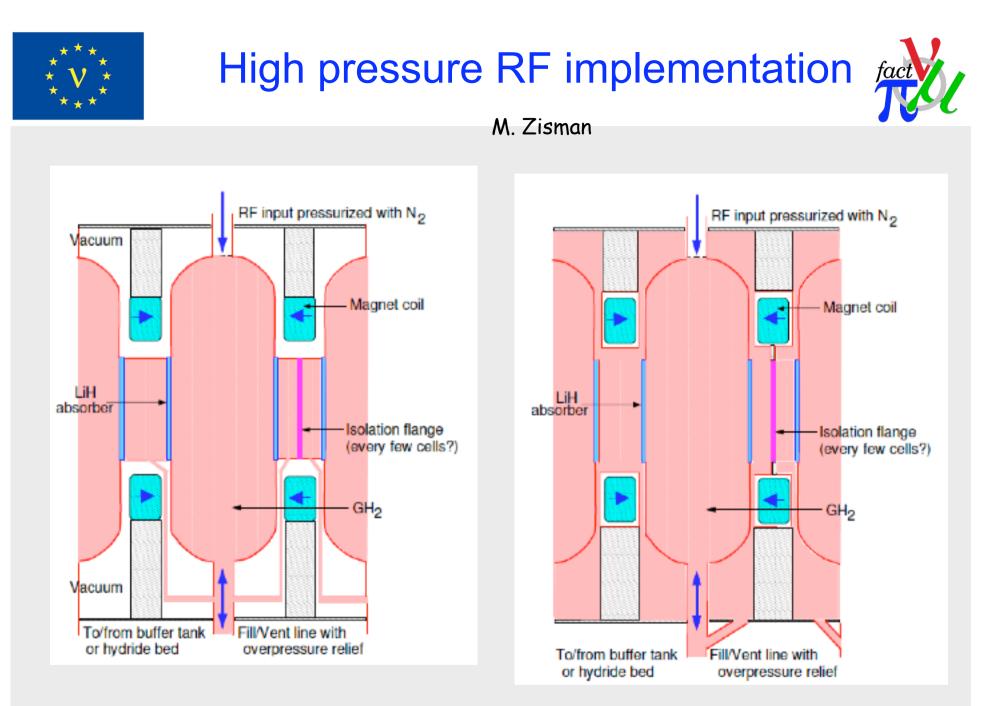
M. Zisman



- We have evidence that vacuum RF cavity gradient performance degrades in a strong magnetic field
 - alternative approach of HPRF does not
 though it has other potential issues
- It seems prudent to begin investigating the technical aspects of implementing HPRF in a linear cooling channel

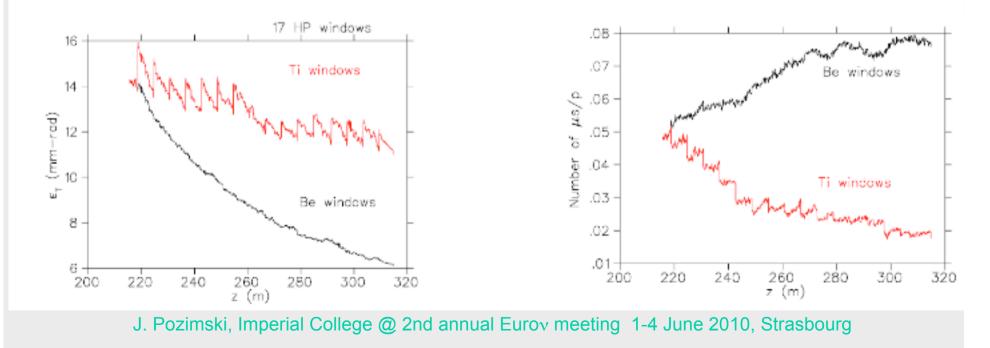
— minimizes changes in cooling channel layout and hardware

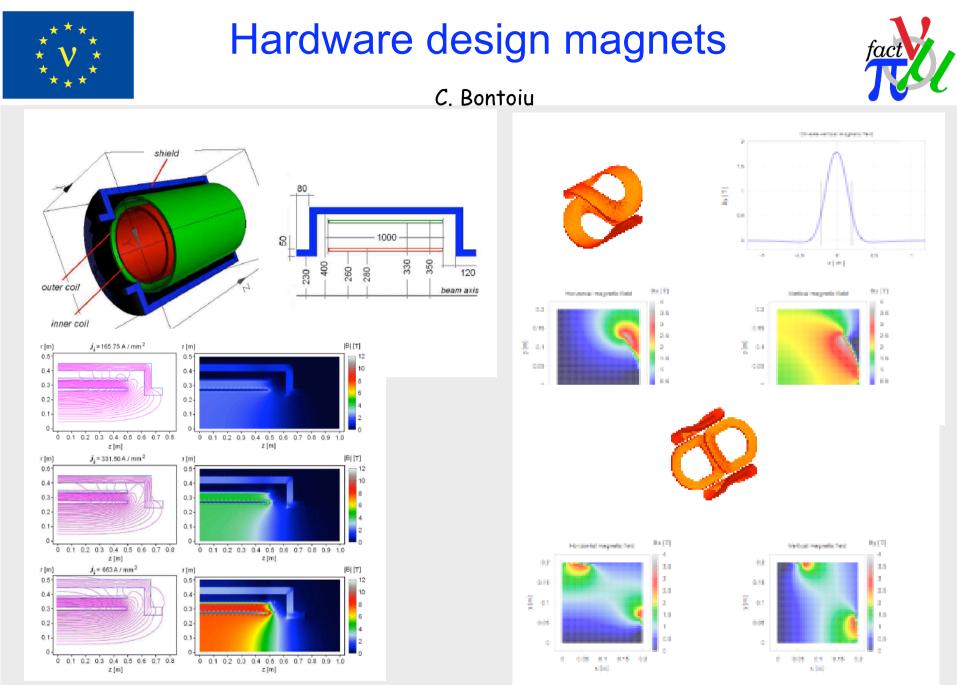






- To make sure we were not fooling ourselves, ran cases with both Be and Ti
 - the difference is obvious
- Will next look at Al and AlBeMet windows as time permits
 - Al is okay in terms of hydrogen embrittlement; not yet sure about other materials





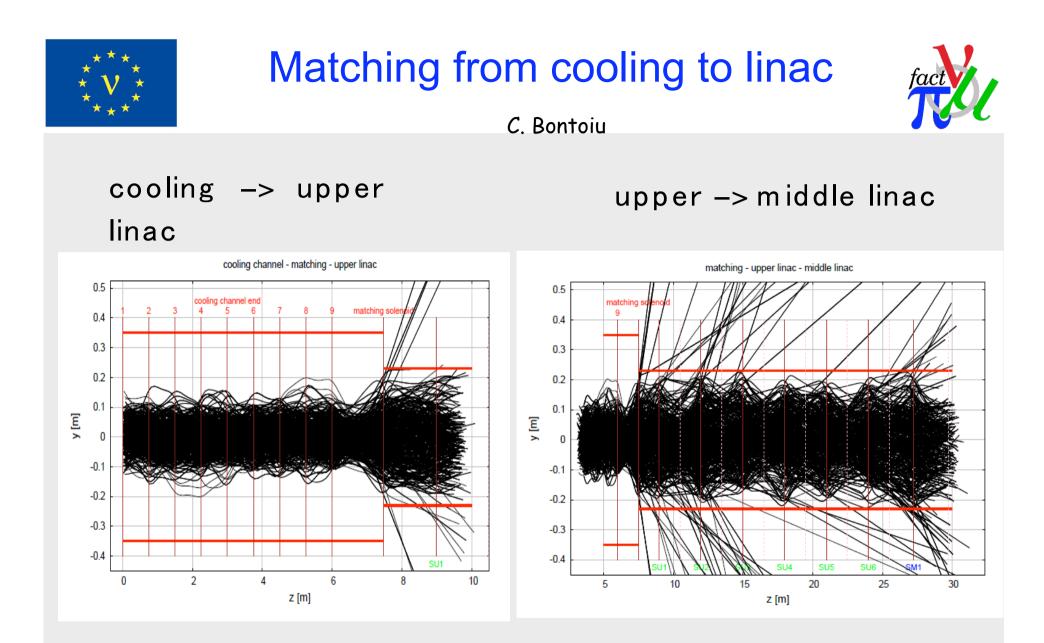
Hardware design RF cavities

C. Bontoiu / J. Pozimski



Normalized on-axis electric field

						1	$\begin{array}{c} \beta_{geom} = 0.9 \ (a) & & \\ \beta_{geom} = 0.9 \ (b) & & \\ \beta_{geom} = 1 & & \\ Study \ II & & \\ \end{array}$
Parameter	$\beta_{geom} = 1$	$\beta_{geom} = 0.9$ (a)	$\beta_{geom} = 0.9$ (b)	Study II	E [a.u.]	0.8	Study II
l_{cav} [m] r [m] f_0 [MHz] Q [10 ⁹]	$\begin{array}{c} 0.7448 \\ 0.6854 \\ 201.247 \\ 24.67 \end{array}$	0.67034 0.7042 201.251 19.6	0.67034 0.6804 201.255 18.8	$\begin{array}{c} 0.8282 \\ 0.6641 \\ 198.575 \\ 26.7 \end{array}$	ш	0.2	-1 -0.5 0 0.5 1 1.5
$ \begin{array}{c} {\rm T} \\ \hat{E} \ [{\rm MV/m}] \\ \bar{E} \ [{\rm MV/m}] \end{array} $	0.650 26.17 20.62	0.716 27.19 20.81	0.726 27.83 20.53	0.591 26.38 20.42			z [m]
$ E _{surface}^{max} [MV/m]$ $ H _{surface}^{max} [kA/m]$ $U [J]$ $\int_{-\infty}^{+\infty} E(0,z) \cos[\omega t(z)] dz$	21.70 48.06 712 8.6142	24.87 58.53 772 9.0081	29.45 61.92 797 9.1336	19.75 45.00 747 8.8466		1.8 1.6 1.4 1.2	$\beta_{geom} = 0.9 (a) \qquad \qquad$
$\int_{-\infty}^{+\infty} D(0,2) \cos[\omega t(2)] dz$ $\int_{-l_{cav}/2}^{+l_{cav}/2} E(0,z) \cos[\omega t(z)] dz$ $\int_{+l_{cav}/2}^{+\infty} E(0,z) \cos[\omega t(z)] dz$ correction [%]	10.0000 -0.69204 -13.841	10.0000 -0.49594 -9.9188	9.9999 -0.43320 -8.6639	10.0000 -0.75676 -15.135	۲۹	1 0.8 0.6 0.2	
						0.2	

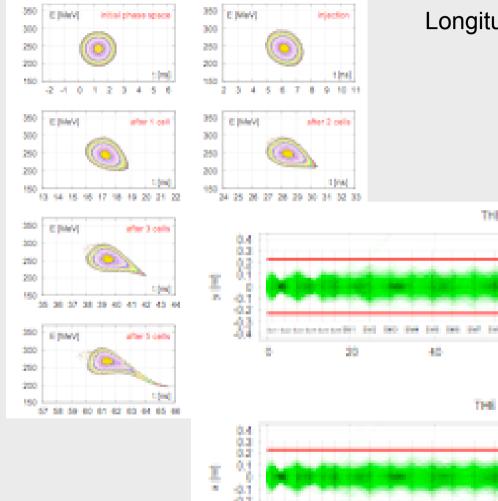




Particle tracking - Linac

C. Bontoiu

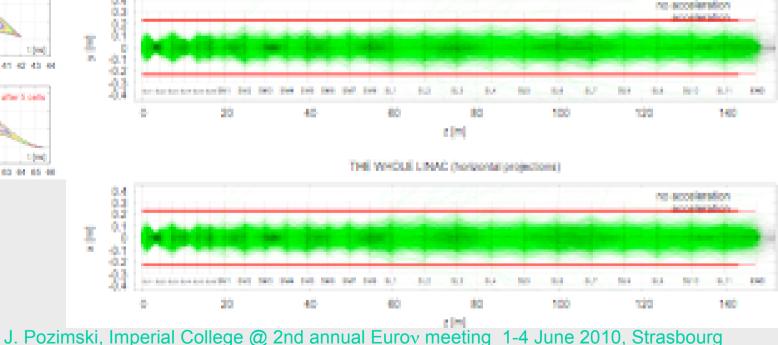


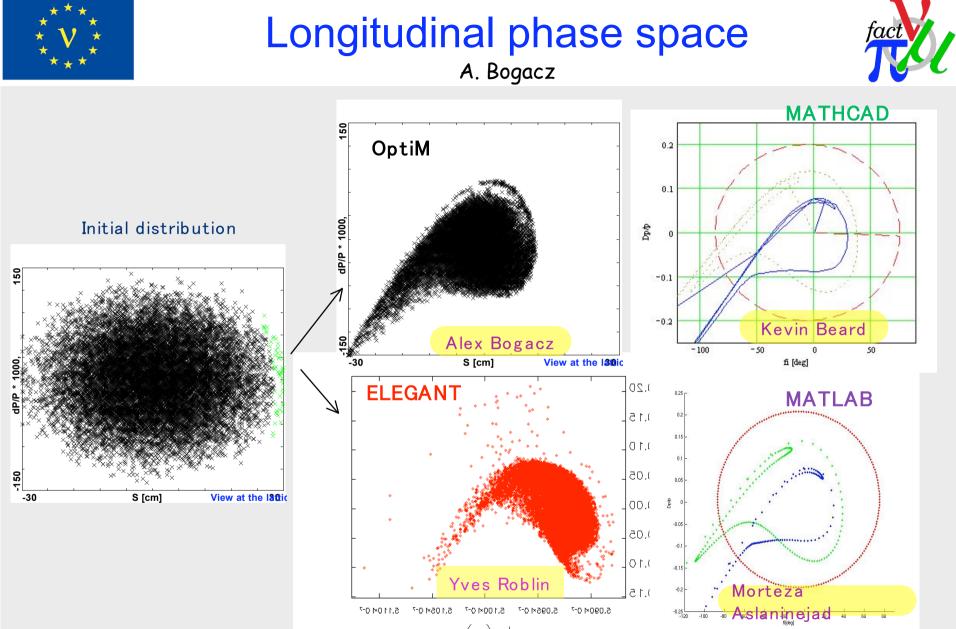


Longitudinal phase space

Transversal projections

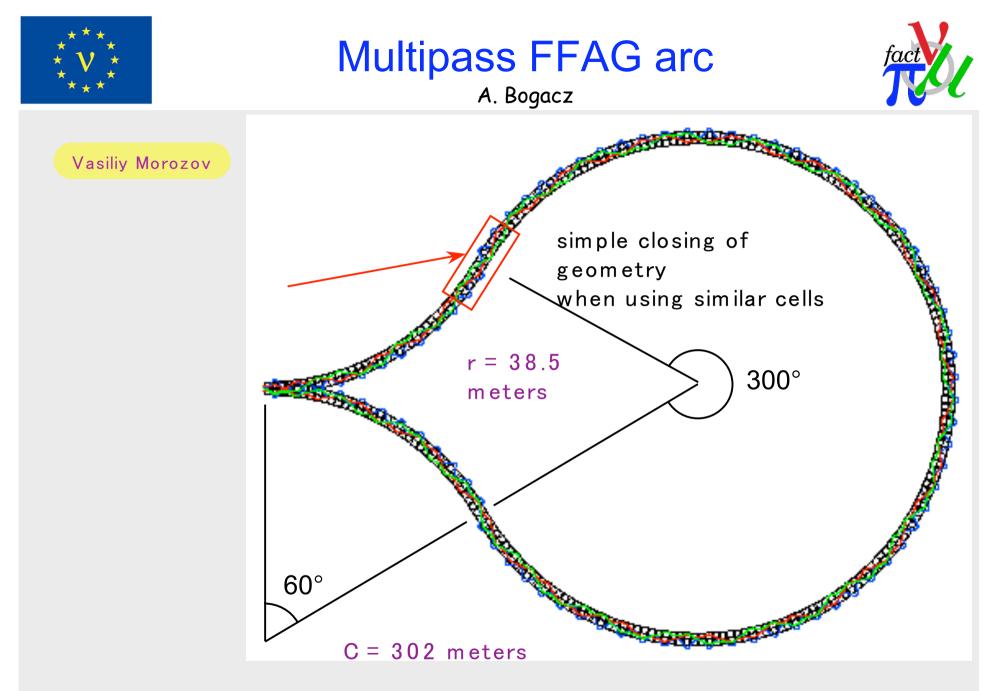
THE WHOLE LINAC (vertical projections)





(s) t

watch-point phase space---input: rlaonly.ele lattice: rla.lte



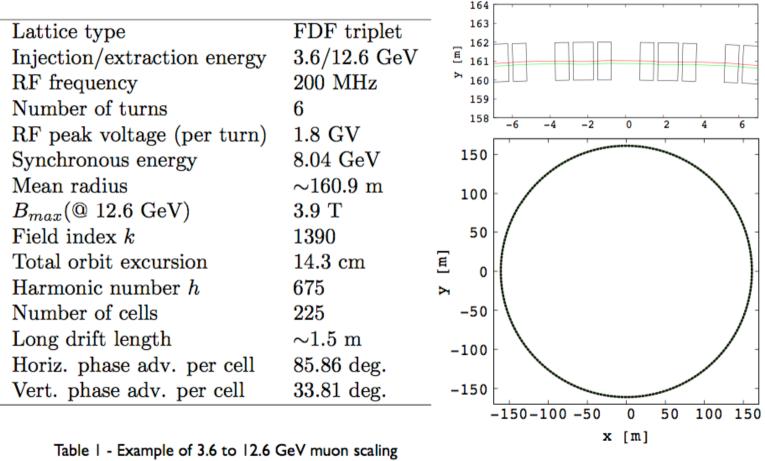


Scaling FFAG alternative for RLA 2

T. Planche

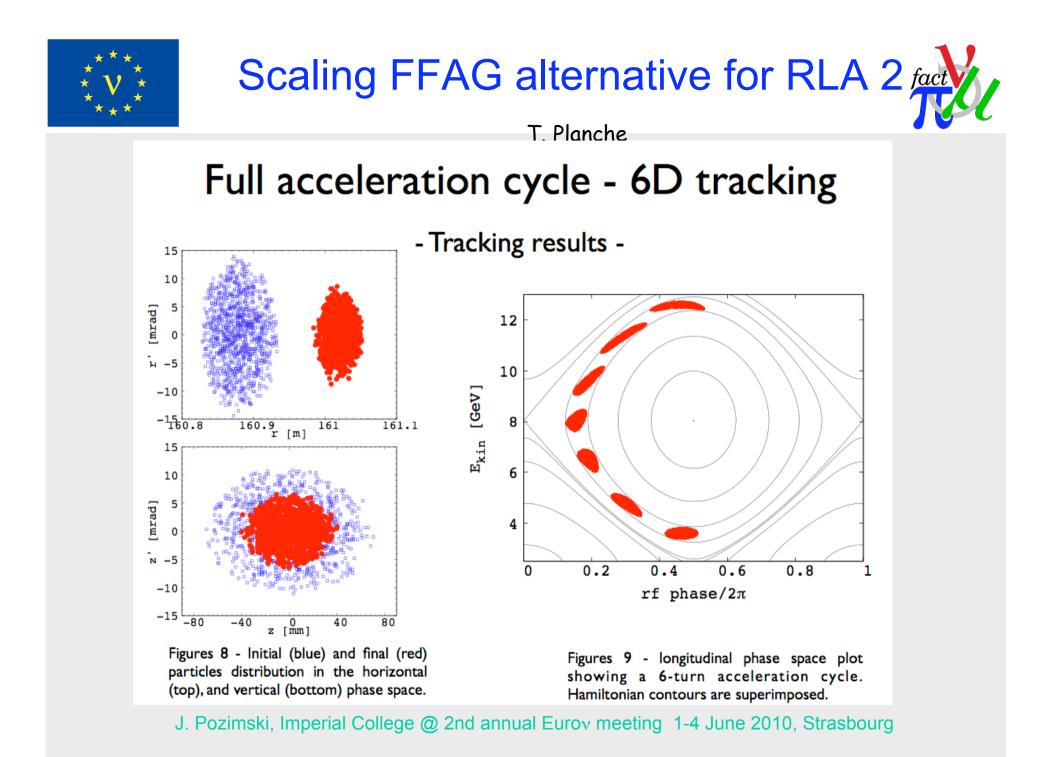


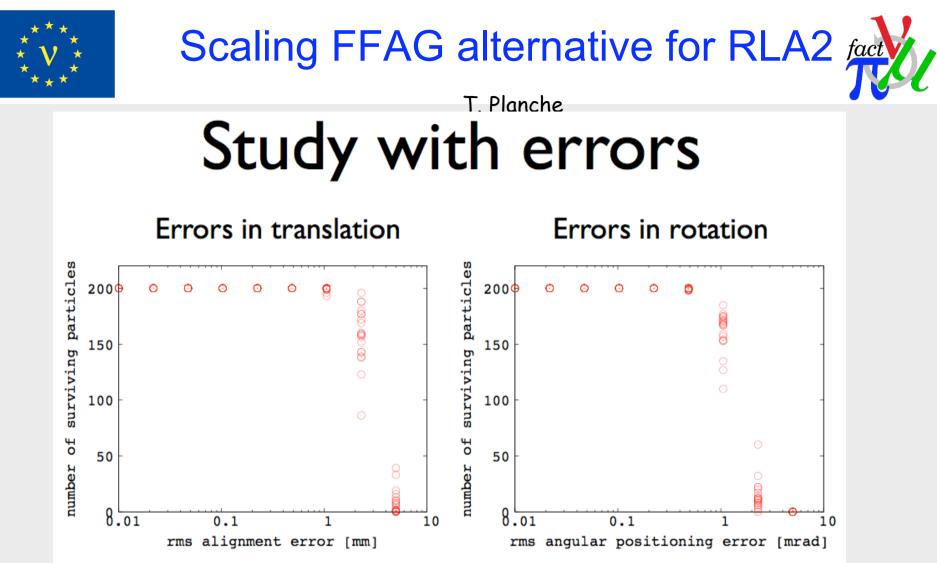
Parameters of a 3.6 to 12.6 GeV muon ring



FFAG ring parameters.

Figure 2 - Ring layout.





Figures 10 - Number of surviving particles depending on the rms error. 20 different lattices have been generated and tested for each value of rms error.

More tolerant to errors than the NS-FFAG!



FFAG Lattice & EMMA



S. Berg / R. Edgecock / F. Meot

Most promising due to longer drifts for RF and kickers : FDFCC

Minimum E 12.6 GeV Maximum E 25 GeV Cells 64 Circumference 546 m Cavity cells 88 RF Voltage 1119 MV Decay 5.6% Turns 11.8 D Radius 11.5 cm F Radius 15.3 cm D Field 6.5 T F Field 3.6 T D Length 1.96 m F Length 1.29 m Long drift 3 m Short drift 0.5 m







S. Machida

 Reduces time of flight dependence on transverse amplitude ○ Downsides Reduction in dynamic aperture Increase in magnet aper tures (cost) Modification of time of flight vs. energy Still studying optimal choice for this Some modest correction likely included



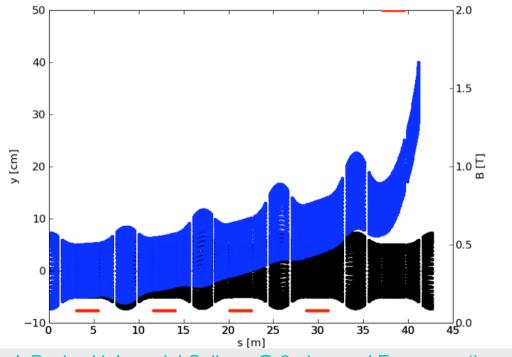
Extraction particle dynamics



J. Pasternak / D. Kelliher

FDFCC - Vertical extraction

- 4 kickers at 0.078 T in consecutive drifts
- Septum at 2T
- Several magnets require large aperture
- Extracted beam 17cm from magnet axis of F after septum





Injection & Extraction studies

J. Pasternak / D. Kelliher



Summary of parameters for injection/extraction (updated lattices)

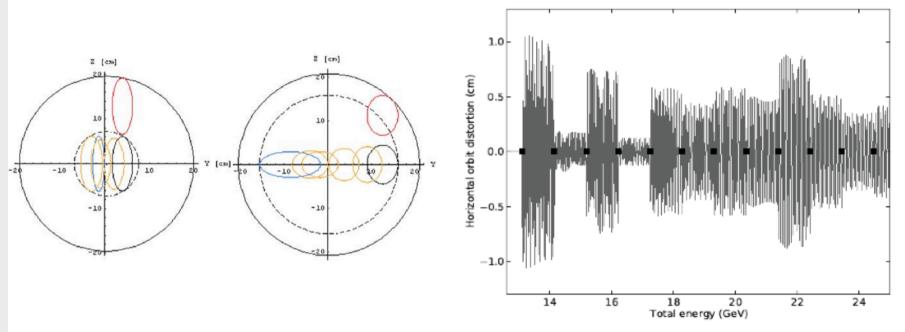
	Triplet 1		FOD	0	Triplet 2	
Scheme	FDFC Injection	FDFC Extraction	FCDC Injection	FCDC Extraction	FDFCC Injection	FDFCC Extraction
Plane	Horizontal	Vertical	Vertical	Vertical	Horizontal	Vertical
No. Kickers	6	6	6	6	3	4
Kicker top field [T]	0.103	0.103	0.113	0.101	0.089	0.078
Septum top field [T]	>3	>4	>3	>4	2	4
Kicker/Sept um length [m]	1.4	1.4	1.4	1.4	2.4	2.4
Cells needed	8	8	5	5	5	6



Extraction particle dynamics



J. Pasternak / / D. Kelliher Effects of special magnets



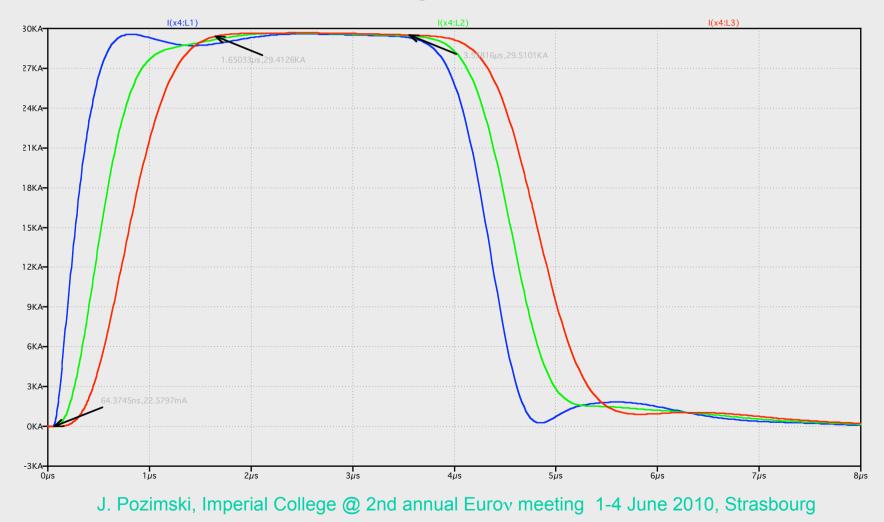
- Beams close to septum push the magnet aperture.
- Special magnets with higher aperture is needed in injection/extraction regions.
- Those magnets introduce the ring lattice symmetry breaking, which can cause accelerated orbit distortion.
- Current studies show that the effect is not dramatic, but more simulations are reqired.

Kicker electronics simulations



J. Pasternak

Current pulses in 3 kicker sections – "travelling wave" using PSPice

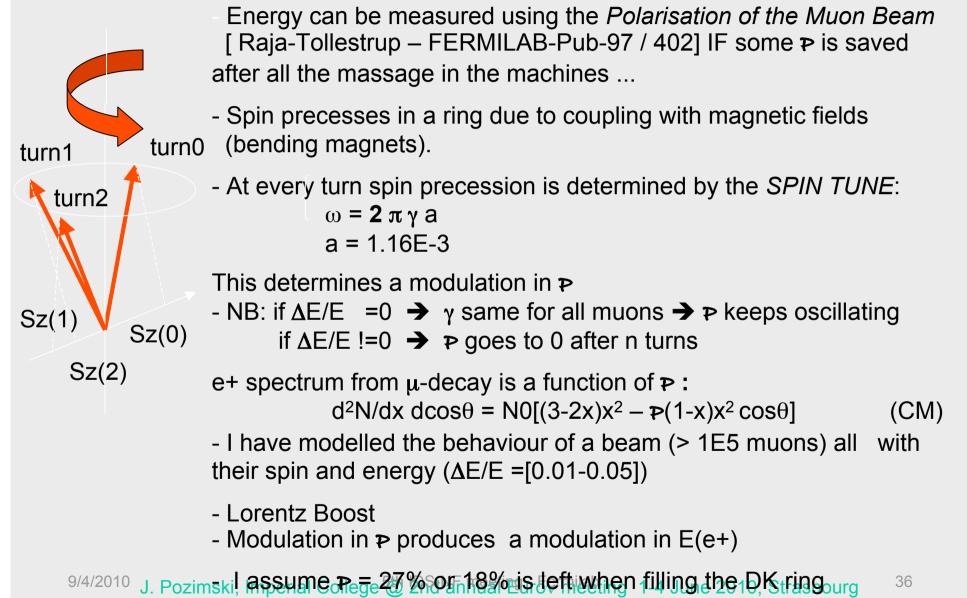




Ring instrumentation - Energy

M. Apollonio



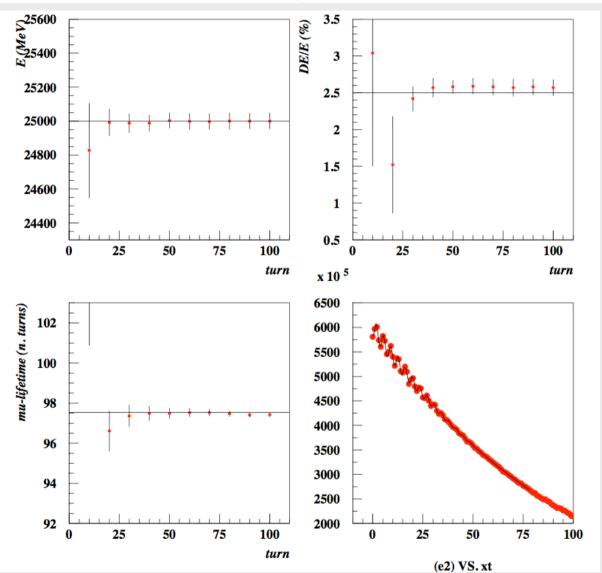




Simulation results



M. Apollonio





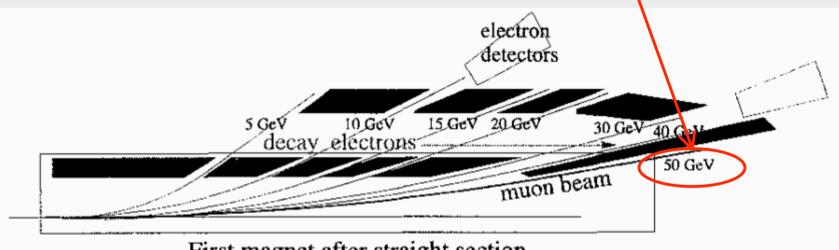
Implementation

fact

M. Apollonio

It has been suggested [Blondel – ECFA 99-197(1999)] to use the *first bending magnet* after the decay straight section to SELECT electron energy bins: what does that mean today with a realistic lattice (25 GeV)?

In fact electron is emitted ~parallel to μ (due to the high γ)



First magnet after straight section

The spectral power of the 1st magnet depends on its *FIELD* and *LENGTH*

A *G4Beamline simulation* can tell us where electrons impinge after decaying somewhere along the orbit





- Sub groups presented detailed plans for convergence.
- How do we organize the decision making process ?
- How do we handle alternatives and fall back options in the IDS ?
- IDR chapters should be prepared to be ready for submission to editors after RAL plenary (1st October)
- For IDS "mixed" costing according to readiness of the hardware design (global to detailed)