# **Theoretical Study of the QCD Phase Diagram at High Densities**

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# <span id="page-2-0"></span>**[Quantum Chromodynamics](#page-2-0)**





At low-energy QCD become hard  $\rightarrow$  non-perturbative.

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**Chiral symmetry breaking:** Quarks are much more massive in the vacuum.









#### **Chiral Condensate** σ



 $\langle 0 | \overline{q} q | 0 \rangle \neq 0$ 



 $\langle 0 | \overline{q} q | 0 \rangle \neq 0$ 

 $\langle 0 | q^T q | 0 \rangle \neq 0$ 



**Main idea**: remove gluons from the theory and replace them by mesons and diquarks.



Meson and diquark can develop **vacuum expectation values** → order parameter for the phases.

# <span id="page-16-0"></span>**[Phase Structure of the](#page-16-0) [Quark-Meson-Diquark model](#page-16-0)**

**Goal:** Compute the partition function *Z*.

**1st method:** Mean-field Approximation (MFA)

$$
Z = \int \mathcal{D}\sigma \mathcal{D}\pi \mathcal{D}\Delta \mathcal{D}\bar{q} \mathcal{D}q \ e^{-\beta(H-\mu N)}
$$

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$$
Z = \int \mathcal{D}\sigma \mathcal{D}\pi \mathcal{D}\Delta \mathcal{D}\bar{q} \mathcal{D}q e^{-\beta(H-\mu N)}
$$
  
MPA  
 $\sigma$ ,  $\Delta$  and  $\pi$  are found by minimising  
the thermodynamic potential

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$$
Z = \int \mathcal{D}\sigma \mathcal{D}\pi \mathcal{D}\Delta \mathcal{D}\bar{q} \mathcal{D}q e^{-\beta(H-\mu N)} = e^{-\beta V\Omega}
$$
  
MFA  

$$
\downarrow \qquad \qquad \downarrow \qquad \qquad \downarrow
$$
  
thermodynamic  
the thermodynamic potential  
potential

#### **2nd method:** Functional Renormalization Group (FRG)

• Take into account bosonic fluctuations in a **non-perturbative** way.

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- Take into account bosonic fluctuations in a **non-perturbative** way.
- Integrate fluctuations momentum-shell by momentum-shell.
- Diquark not included yet but give insight into the importance of fluctuation.



#### **Equation of State**



<span id="page-25-0"></span>**[Neutron Stars](#page-25-0)**

## **Origin of a Neutron Star**





## **Origin of a Neutron Star**







$$
\frac{dp}{dr} = -\frac{(\epsilon(r) + p(r)) [m(r) + 4\pi r^3 p(r)]}{r(r - 2Gm(r))}
$$

$$
\frac{dm}{dr} = 4\pi r^2 \epsilon(r)
$$

## **Tolman-Oppenheimer-Volkoff Equation**





#### **Mass-radius Relationship**



**Tidal deformability** λ: Ability of a star to deform under an external gravitational field.

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**How to measure it**:



## **Tidal Deformability**



#### **Summary:**

- Phase strucutre of the quark-meson-diquark model and results on neutron star observables.
- Maybe visible effect of diquarks with high coupling?
- Strong effect of mesonic fluctuations.

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- Phase strucutre of the quark-meson-diquark model and results on neutron star observables.
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## **Outlook:**

- Implement the FRG with diquarks.
- Build hybrid model of neutron star including neutrons.

# **Thank you!**

$$
\mathcal{L}_{QMD} = \overline{q} \left( -i \partial - \mu \gamma^0 \right) q + g_m \overline{q} (\sigma + i \gamma_5 \vec{\pi} \cdot \vec{\tau}) q
$$
  
+ 
$$
ig_{\Delta} \left( \Delta^* \overline{q} \gamma_5 \tau_2 \lambda_2 C \overline{q}^T + \Delta q^T C \gamma_5 \tau_2 \lambda_2 q \right)
$$
  
+ 
$$
\frac{1}{2} (\partial_{\mu} \sigma) (\partial^{\mu} \sigma) + \frac{1}{2} (\partial_{\mu} \vec{\pi}) (\partial^{\mu} \vec{\pi}) + (\partial_{\mu} \Delta) (\partial^{\mu} \Delta^*) + U(\sigma, \vec{\pi}, \Delta)
$$

$$
\mathcal{L}_{\text{QMD}} = \overline{\overline{q} \left( -i \partial\!\!\!/- \mu \gamma^0 \right)} q \; + \; g_m \overline{q} (\sigma + i \gamma_5 \vec{\pi} \cdot \vec{\tau}) q
$$

+ 
$$
ig_{\Delta}(\Delta^*\overline{q}\gamma_5\tau_2\lambda_2 C\overline{q}^T + \Delta q^T C\gamma_5\tau_2\lambda_2 q)
$$
  
+  $\frac{1}{2}(\partial_{\mu}\sigma)(\partial^{\mu}\sigma) + \frac{1}{2}(\partial_{\mu}\vec{\pi})(\partial^{\mu}\vec{\pi}) + (\partial_{\mu}\Delta)(\partial^{\mu}\Delta^*) + U(\sigma, \vec{\pi}, \Delta)$ 

$$
\mathcal{L}_{QMD} = \overline{q} \left( -i \partial \!\!\! / \partial - \mu \gamma^0 \right) q \; + \; \overline{g_m \overline{q} (\sigma + i \gamma_5 \overline{\pi} \cdot \overline{\tau}) q}
$$

Interaction quark-diquark

$$
+ig_{\varDelta}\left(\varDelta^{*}\overline{q}\gamma_{5}\tau_{2}\lambda_{2}C\overline{q}^{T}+\varDelta q^{T}C\gamma_{5}\tau_{2}\lambda_{2}q\right)
$$

$$
+\frac{1}{2}(\partial_{\mu}\sigma)(\partial^{\mu}\sigma)+\frac{1}{2}(\partial_{\mu}\vec{\pi})(\partial^{\mu}\vec{\pi})+(\partial_{\mu}\Delta)(\partial^{\mu}\Delta^{*})+U(\sigma,\vec{\pi},\Delta)
$$

$$
\mathcal{L}_{QMD} = \overline{\overline{q}(-i\partial - \mu \gamma^0)q} + \overline{g_m \overline{q}(\sigma + i\gamma_5 \vec{\pi} \cdot \vec{\tau})q}
$$

Interaction quark-diquark

$$
+ \left[i g_\varDelta \left( \varDelta^* \overline{q} \gamma_5 \tau_2 \lambda_2 C \overline{q}^T + \varDelta q^T C \gamma_5 \tau_2 \lambda_2 q \right)\right]
$$

$$
+\frac{1}{2}(\partial_{\mu}\sigma)\left(\partial^{\mu}\sigma\right)+\frac{1}{2}(\partial_{\mu}\vec{\pi})\left(\partial^{\mu}\vec{\pi}\right)+\frac{(\partial_{\mu}\Delta)\left(\partial^{\mu}\Delta^{*}\right)}{\text{diquark dynamics}}+U(\sigma,\vec{\pi},\Delta)
$$

meson dynamics

diquark dynamics

×

$$
\mathcal{L}_{QMD} = \overline{q} \left( -i \vartheta - \mu \gamma^0 \right) q + \left[ g_m \overline{q} (\sigma + i \gamma_5 \vec{\pi} \cdot \vec{\tau}) q \right]
$$
\nInteraction quark-diquark

\nInteraction quark-diquark

\ninteraction quark-diquark

\nthe reaction quark-diquark

\nthe second term

\nfunction

\n
$$
+ ig_{\Delta} \left( \Delta^* \overline{q} \gamma_5 \tau_2 \lambda_2 C \overline{q}^T + \Delta q^T C \gamma_5 \tau_2 \lambda_2 q \right)
$$
\nmeson and diquark self interaction

\n
$$
+ \frac{1}{2} (\partial_{\mu} \sigma) \left( \partial^{\mu} \sigma \right) + \frac{1}{2} (\partial_{\mu} \vec{\pi}) \left( \partial^{\mu} \vec{\pi} \right) + \left( \partial_{\mu} \Delta \right) \left( \partial^{\mu} \Delta^* \right) + U(\sigma, \vec{\pi}, \Delta)
$$
\nmeson dynamics