

# SIGNATURES OF THE YANG-MILLS DECONFINEMENT TRANSITION FROM THE GLUON TWO-POINT CORRELATOR

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**arXiv:2206.03841 [hep-ph]**

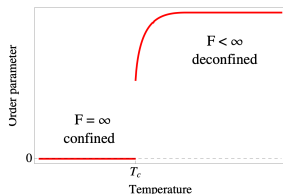
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# INTRODUCTION

- At some very high temperature  $T_c$ , hadrons become free quarks and gluons  $\rightarrow$  **quark-gluon plasma**.
- An order parameter for this transition is the **Polyakov loop**:

$$\ell \sim \langle P e^{i \int_0^\beta d\tau A_0(\tau, x)} \rangle \sim e^{-\beta F}$$



- Under center symmetry,  $\ell \rightarrow Z_N \ell$ , so deconfinement is signaled by a broken center symmetry (in pure Yang-Mills).
- Confirmed by lattice data: second order transition for  $SU(2)$ , first order for  $SU(3)$ .

# ENCODING OF THE TRANSITION

Polyakov loop:

$$\ell \sim \langle P e^{i \int_0^\beta d\tau A_0(\tau, x)} \rangle \sim e^{-\beta F}.$$

Because the Polyakov loop is related to  $A_0$ , it is expected that the transition is encoded in (the tower of)

$$\langle A_0 \rangle, \langle A_0 A_0 \rangle, \dots, \langle A_0^n \rangle.$$

For the appropriate choice of gauge, can the transition be reflected in the lowest order correlators?

# LANDAU GAUGE CORRELATOR

In principle:

- $\langle A \rangle$  is found by minimizing the effective action  $\Gamma[A]$ . It represents the state of the system.  $\langle A_0 \rangle \rightarrow$  order parameter.
- The two-point correlator derives from the effective action

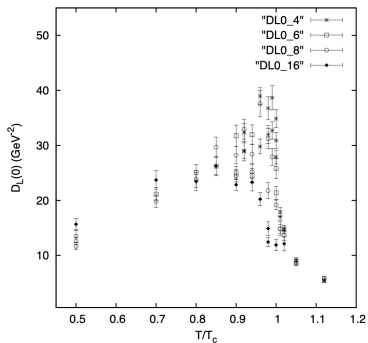
$$1 \left/ \frac{\partial^2 \Gamma}{\partial A^2} \right|_{A=\langle A \rangle} = \langle AA \rangle_c,$$

so for  $SU(2)$ ,  $\langle A_0 A_0 \rangle$  should diverge at  $T_c$ .

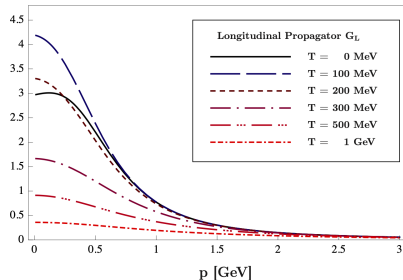
In practice:

- In the Landau gauge,  $\partial_\mu A_\mu = 0$ , then  $\langle A_0 \rangle = 0$ .  $\rightarrow$  no order parameter.
- No evidence of divergence of  $\langle A_0 A_0 \rangle$  was found on the (gauge-fixed) lattice and in the continuum.

# SU(2) LANDAU GAUGE CORRELATORS



Electric susceptibility (zero momentum longitudinal propagator)



Longitudinal gluon propagator

\*T. Mendes and A. Cucchieri, PoS LATTICE2014, 183 (2015).

\*L. Fister and J. M. Pawłowski, [arXiv:1112.5440 [hep-ph]](2012).

# BACKGROUND FIELD GAUGES

- In the **Landau gauge** the effective action is not center-symmetric  $\Gamma[A] \neq \Gamma[A^U]$ .
- Introducing a background field  $\bar{A}$ , the effective action is gauge-invariant  $\Gamma_{\bar{A}}[A] = \Gamma_{\bar{A}^U}[A^U]$ . Landau-deWitt gauge:

$$\bar{D}_\mu(A_\mu - \bar{A}_\mu) = 0, \quad \text{with } \bar{D}_\mu \equiv \partial_\mu - [\bar{A}_\mu, \cdot].$$

- In the **Background field effective action** one looks at  $\tilde{\Gamma}[\bar{A}] = \Gamma_{\bar{A}}[\bar{A}]$ .
  - ▶ The two-point function  $\langle AA \rangle_c$  is not directly accessible from  $\tilde{\Gamma}_{\bar{A}}$ .
  - ▶ Relies on the strict background independence of  $\tilde{\Gamma}[\bar{A}]$ , which is not easy to maintain in the presence of truncations.
- We propose the **Center-symmetric Landau gauge**, which fixes  $\bar{A} = \bar{A}_c$ . Then  $\Gamma_{\bar{A}}[A] = \Gamma_{\bar{A}^U c}[A^U c]$ .

# OUR SETUP

- We work in the Landau-deWitt gauge with a background field  $\bar{A}_\mu$ :

$$\bar{D}_\mu(A_\mu - \bar{A}_\mu) = 0, \quad \text{with } \bar{D}_\mu \equiv \partial_\mu - [\bar{A}_\mu, \cdot].$$

- We take  $A_\mu^a$  and  $\bar{A}_\mu^a$  in the temporal direction,  $\propto \delta_{\mu 0}$ , and along the diagonal color directions ( $\sigma^3$  for  $SU(2)$ ,  $(\lambda^3, \lambda^8)$  for  $SU(3)$ ), so that  $\Gamma[A, \bar{A}] \propto V(A, \bar{A})$ .
- The center-symmetric values for  $A_\mu^a$  and  $\bar{A}_\mu^a$  are found by Weyl chambers. For example in  $SU(2)$ ,  $A_{c,\mu} = \bar{A}_{c,\mu} = \delta_{\mu 0} \frac{T}{g} \pi \frac{\sigma^3}{2}$ .
- We fix  $\bar{A} = \bar{A}_c$ : **Center-symmetric Landau gauge**.  
Center-symmetric phase when  $\langle A \rangle = A_c$ ,  $\rightarrow$  **order parameter**.

## CURCI-FERRARI MODEL

We have computed  $\langle A \rangle$  and  $\langle A(0, \rho)A(0, -\rho) \rangle$  up to first loop order in the finite temperature Curci-Ferrari model:

$$S = S_{YM} + S_{gf} + \int_{x,\tau} \frac{m^2}{2} (A_\mu^a - \bar{A}_\mu^a)^2$$

Several motivations:

- Perturbative gauge-fixed Yang-Mills **breaks down** at low energies (Landau pole, Gribov copies...), there is no analytical model for this region.
- A gluon mass term seems to dominate the (unknown) gauge-fixed action in the IR; decoupling behaviour on the lattice. CF could be an **effective model**.
- The CF model is renormalizable, avoids the Landau pole and lifts the degeneracy between Gribov copies. Perturbative window into non-perturbative region.



# RESULTS - $T_c$ (MeV)

	Lattice	FRG-BG <sup>1</sup>	CF-BG, 1-lp <sup>2</sup>	CF-BG, 2-lp <sup>3</sup>	CF-CS, 1-lp
SU(2)	295	230	238	284	<b>265</b>
SU(3)	270	275	185	254	<b>267</b>


BG: Background effective action

CS: Centrosymmetric Landau gauge

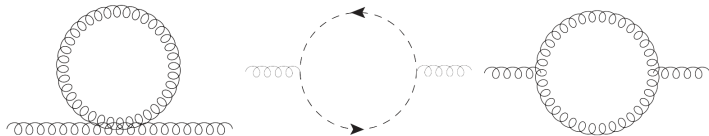
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<sup>1</sup>L. Fister and J. M. Pawłowski, Phys.Rev. D88 (2013) 045010

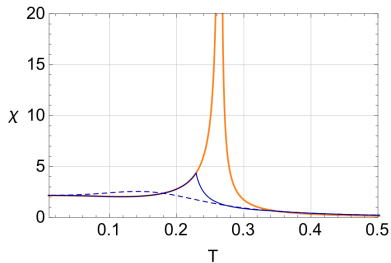
<sup>2</sup>U. Reinosa, J. Serreau, M. Tissier and N. Wschebor, Phys.Lett. B742 (2015) 61-68.

<sup>3</sup>U. Reinosa, J. Serreau, M. Tissier and N. Wschebor, Phys.Rev. D93 (2016) 105002. 

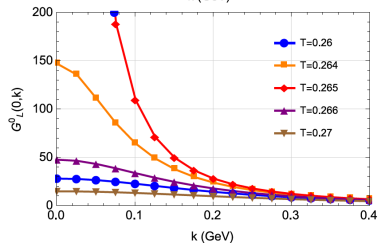
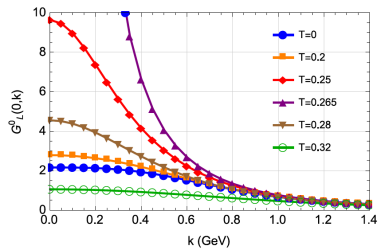
# FEYNMAN DIAGRAMS GLUON PROPAGATOR



# RESULTS: SU(2) GLUON PROPAGATOR



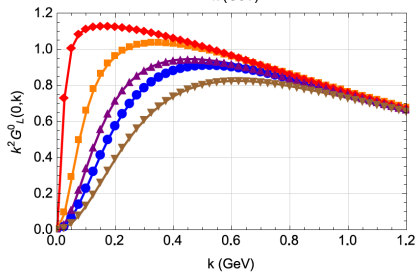
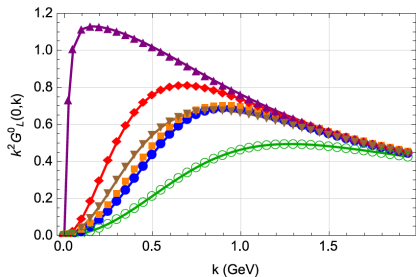
Landau gauge  
Background field effective action  
Centersymmetric Landau gauge



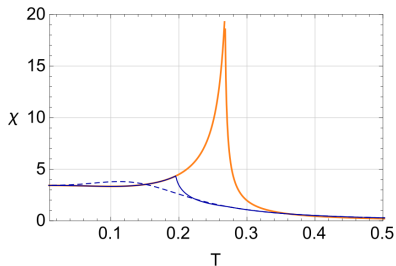
$m=0.68$  GeV,  $\mu=1$  GeV,  $g=7.5$

\*DvE, U. Reinosa, J. Serreau and M. Tissier, SciPost Phys. 12, 087 (2022).

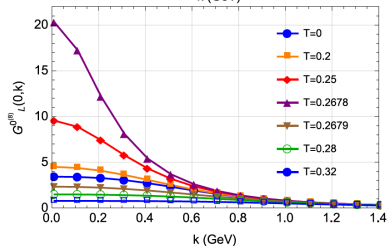
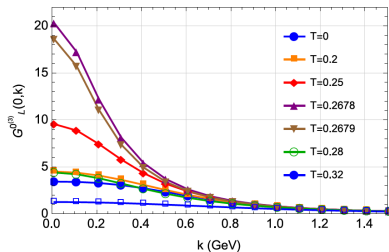
# SU(2) DRESSING FUNCTION



# RESULTS: SU(3) GLUON PROPAGATOR



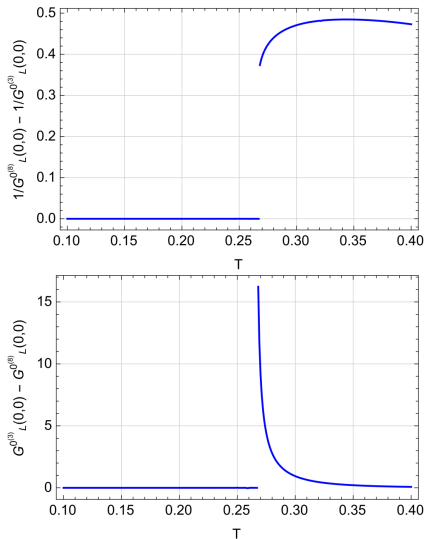
Landau gauge  
Background field effective action  
Centersymmetric Landau gauge



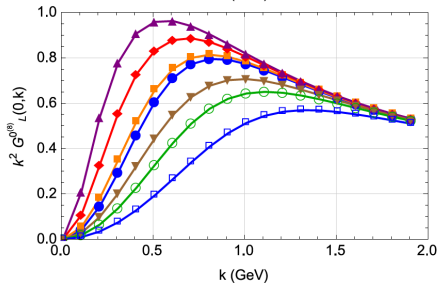
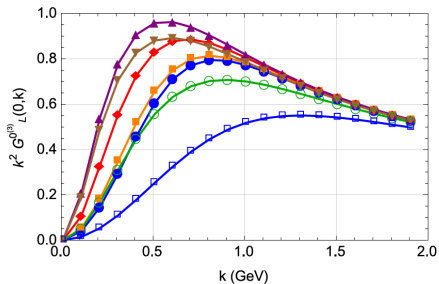
$m=0.54$  GeV,  $\mu=1$  GeV,  $g=4.9$

\*DvE, U. Reinosa, J. Serreau and M. Tissier, SciPost Phys. 12, 087 (2022).

# SU(3) PROPAGATOR DIFFERENCE



# SU(3) DRESSING FUNCTION

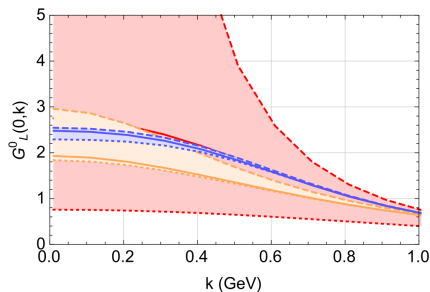


# CONCLUSION AND OUTLOOK

- We have performed, for the first time, calculations of the gluon one- and two-point correlator in the center-symmetric Landau gauge.
- We find a good agreement with lattice data for  $T_c$ .
- We find that for  $SU(2)$ , the deconfinement transition is signaled by a **divergence** of the longitudinal gluon propagator for  $k \rightarrow 0$ .
- For  $SU(3)$ , the difference between the propagators in the neutral color mode is an order parameter for the transition.
- This model can be tested on the lattice by changing the boundary conditions in the Landau gauge [with O. Oliveira and P. Silva].
- Ideas for future works: RG improvement, transversal propagator and two-loop calculation.



# DISCUSSION

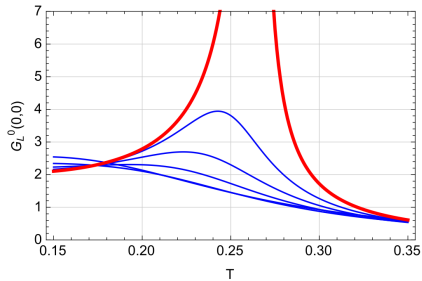


Landau gauge

Self-consistent Landau gauge

Center-symmetric Landau Gauge

# DISCUSSION



# DISCUSSION

