

RECENT PROGRESS IN HYDRODYNAMICS FOR HEAVY-ION COLLISIONS

BAYESIAN INFERENCE AND RECENT PHENOMENOLOGICAL CONSTRAINTS ON QGP PROPERTIES

Matthew Luzum

Based on:

JETSCAPE, Phys.Rev.C 103 (2021) 5, 054904; Phys.Rev.Lett. 126 (2021) 24, 242301
Nijs, et. al, Phys.Rev.Lett. 126 (2021) 20, 202301; Phys.Rev.C 103 (2021) 5, 054909 (Trajectum 1)
Nijs, et. al, arXiv:2110.13153 (Trajectum 2)

University of São Paulo

Strong and Electro-Weak Matter 2022
June 20

OUTLINE

- 1 BAYESIAN INFERENCE PRIMER
- 2 HYDRODYNAMIC MODEL AND PARAMETERIZATION
- 3 RESULTS AND DISCUSSION
- 4 SUMMARY / CONCLUSIONS

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BAYESIAN BASICS:

- **Parameter estimation**: Given a model, what parameter values are compatible with experimental data, and with what precision can we determine them?
- Can answer with **Bayesian inference** — ideal for detailed and systematic treatment of uncertainty
- Experimental data (D) and parameters (p) are each associated with probability distributions
- **Bayes' theorem** relates conditional probabilities. E.g., $\Pr(D|p)$ is the probability of D , given p .
- The probability that both D and p are true is

$$\Pr(p \& D) = \Pr(p) \times \Pr(D|p) = \Pr(D) \times \Pr(p|D)$$

prior \times likelihood = evidence \times posterior

- We typically want to know $\Pr(p|D) \propto \Pr(p) \Pr(D|p)$
- \implies need to choose a **prior** $\Pr(p)$ and compute the **likelihood** $\Pr(D|p)$ from comparison with data

$$\Pr(D|p) \propto e^{-\chi^2/2}$$

$$\text{with } \chi^2 = (D - \text{Model}(p))^T \Sigma^{-1} (D - \text{Model}(p))$$

and Σ = uncertainty covariance (exp. and theor.)

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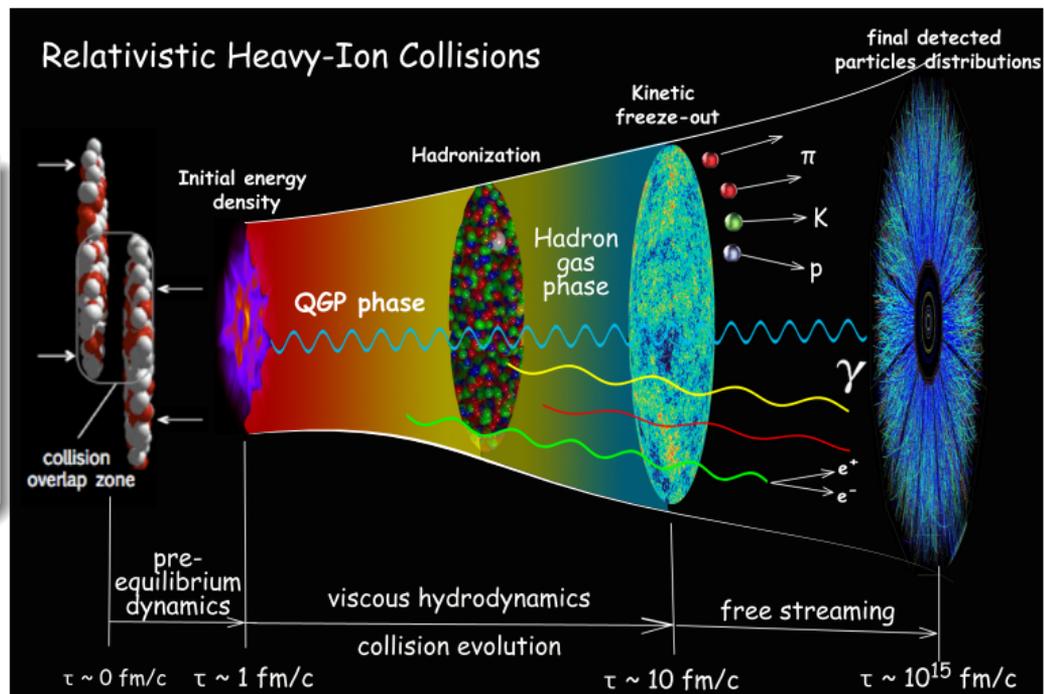
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TIME LINE OF HEAVY-ION COLLISION

Collision model

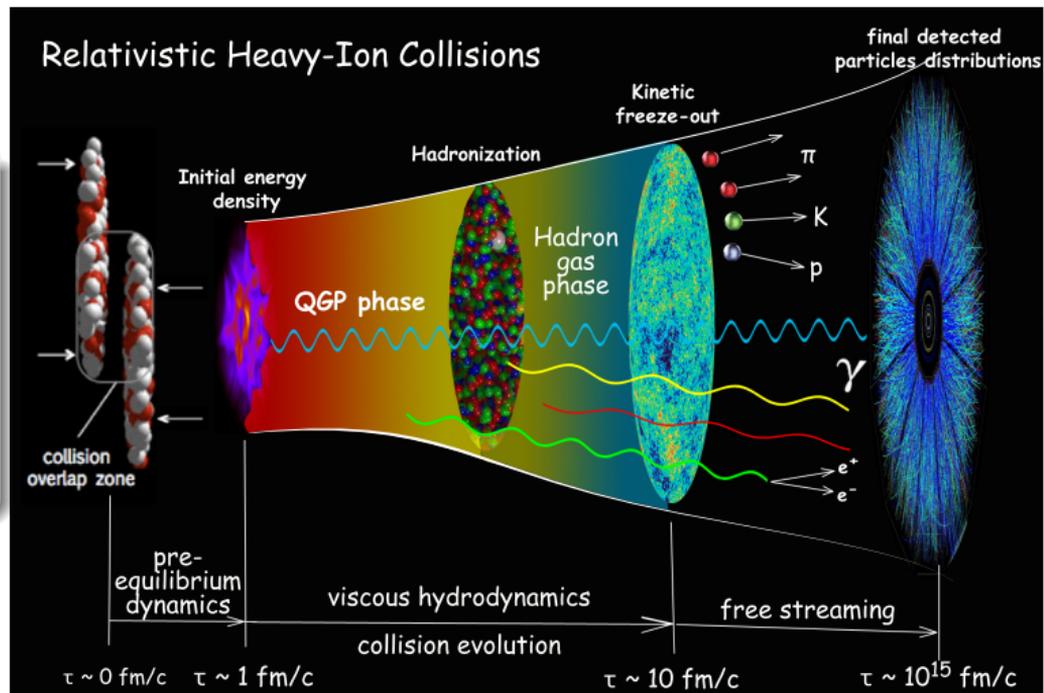
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- Initial scattering
- Hydrodynamization
- Relativistic Fluid
 - Quark-Gluon Plasma
 - Hadrons
- Hadronic scattering



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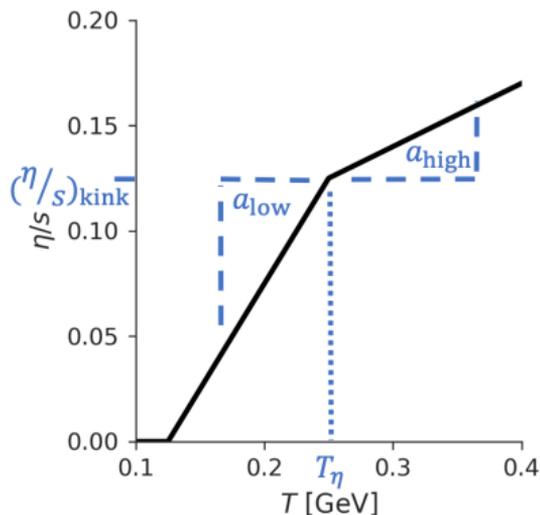
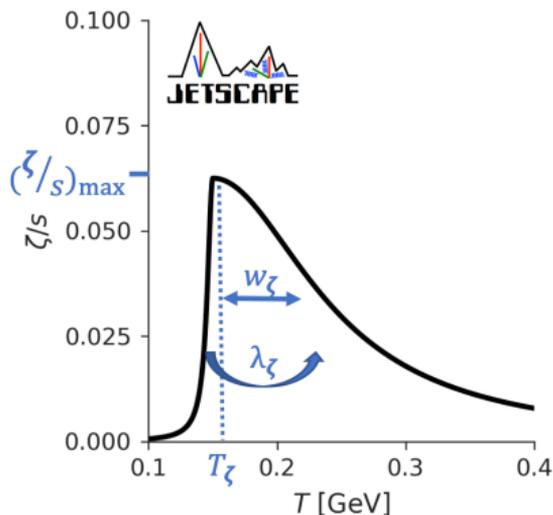
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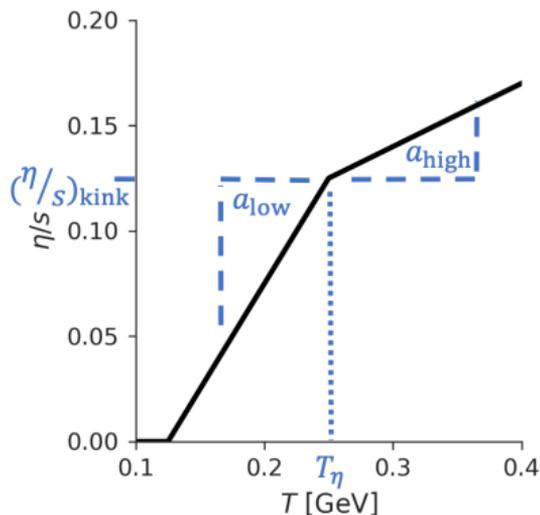
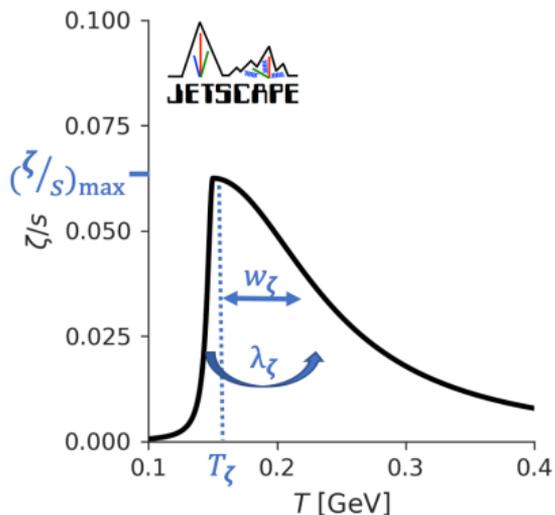
HYDRODYNAMICS

- Main workhorse: 2nd order relativistic viscous hydrodynamics
- Equation of state from Lattice $\epsilon(p)$
- Unknown quantities: transport coefficients
- Shear $\frac{\eta}{s}(T)$ and bulk viscosity $\frac{\zeta}{s}(T)$; second order transport coefficients $\tau_\pi, \tau_\Pi, \tau_{\pi\pi}$ (Trajectum)



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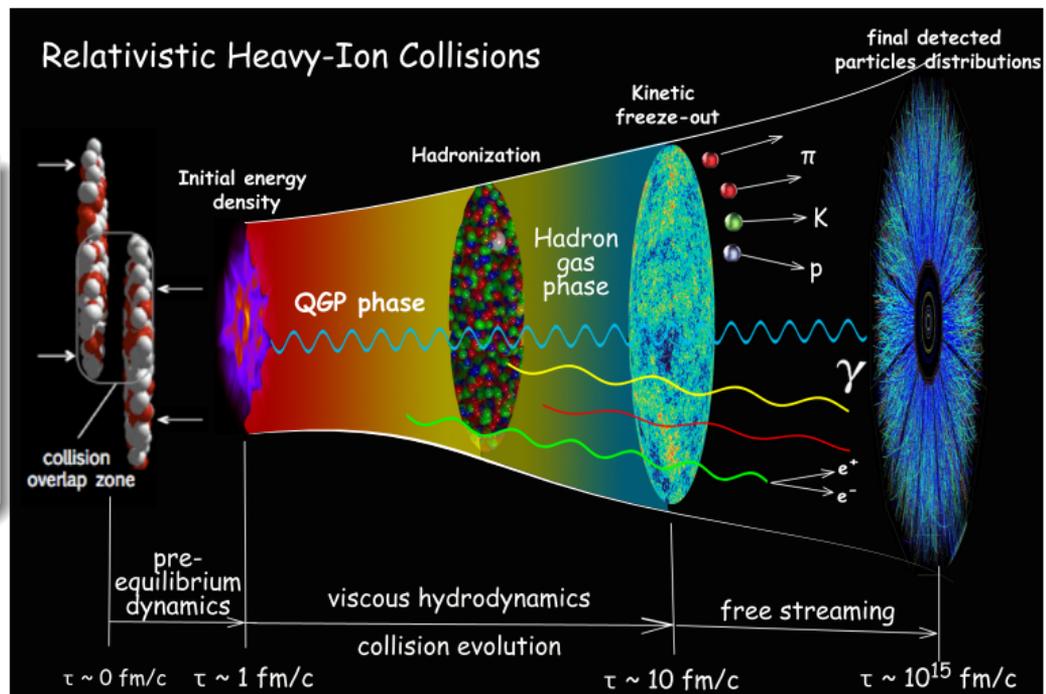
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INITIAL CONDITIONS FOR HYDRODYNAMICS

- Nucleus
 - Nucleon positions sampled from Woods-Saxon distribution
 - Reject nuclei with pairs of nucleons closer than d_{min}
- Trento
 - Boost invariant
 - Participants determined by impact-parameter dependent cross section with nucleon width parameter w
 - Energy density at time $\tau = 0^+$ proportional to generalized mean of nuclear thickness functions multiplied by a random fluctuation γ of variance σ_k^2 .

$$\tau\epsilon(\mathbf{x}) = NT_R(\mathbf{x}_\perp; \rho) = N \left(\frac{T_A^p(\mathbf{x}_\perp) + T_B^p(\mathbf{x}_\perp)}{2} \right)^{1/p}$$

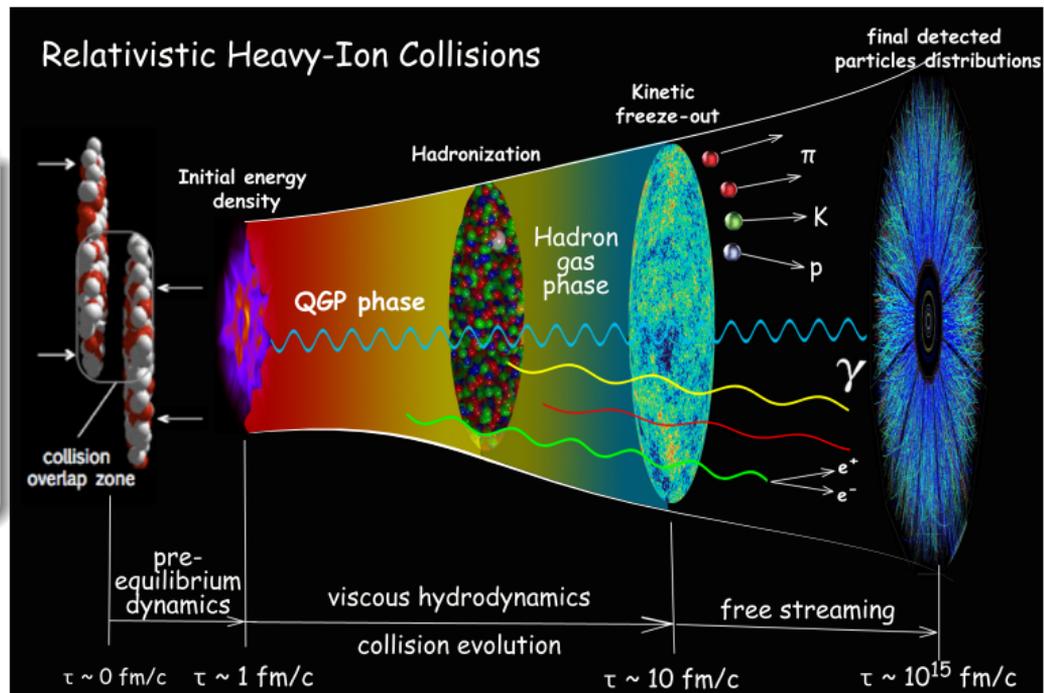
$$T_A(\mathbf{x}_\perp) = \sum_{i \in A} \gamma_i \rho(\mathbf{x}_\perp - \mathbf{x}_{i,\perp})$$

- Free steaming
 - Energy spreads out isotropically with transverse velocity $v = 1$ (JETSCAPE) or $v \leq 1$ (Trajectum) for time τ_{fs} , which can depend on energy via exponent α (JETSCAPE)
 - Full energy-momentum tensor at τ_{fs} used as initial condition for hydro

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HADRONIC AFTERBURNER

- Switch from fluid description to particles (hadrons) at temperature T_{sw}
- Equilibrium distribution function given by kinetic theory, but **viscous corrections** non-universal
- Estimate uncertainty via 3 models
 - Grad (JETSCAPE)
 - Chapman-Enskog (JETSCAPE)
 - Pratt-Torrieri-Bernhard (Trajectum & JETSCAPE)
- Collisions and decays via SMASH (JETSCAPE & Trajectum 1) or UrQMD (Trajectum 2)

SYSTEMS AND OBSERVABLES

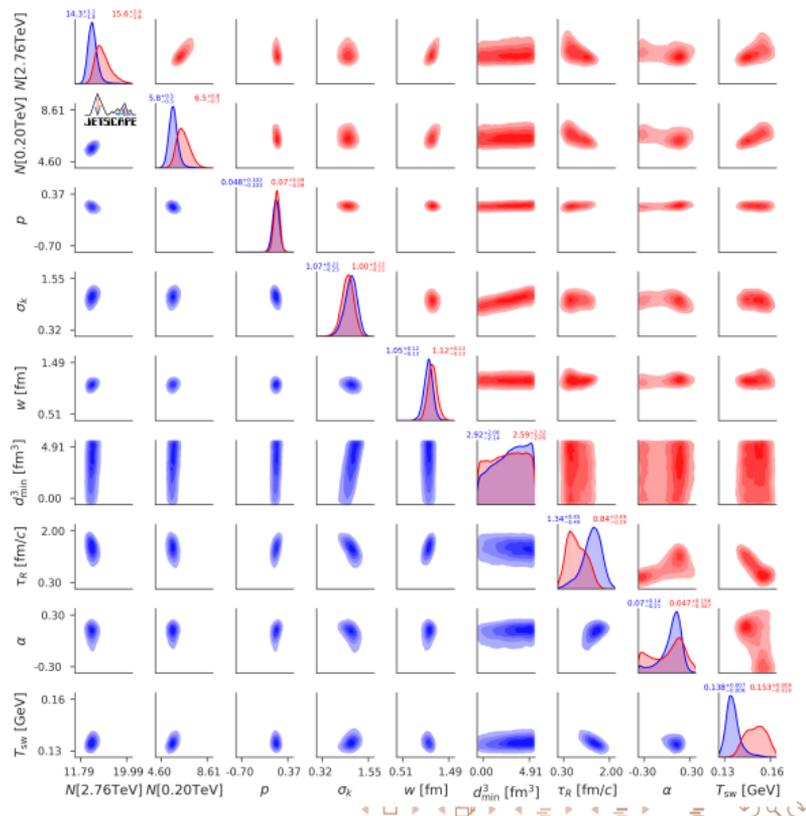
- Systems
 - Trajectum 1 (pPb & PbPb)
 - Trajectum 2 (PbPb)
 - JETSCAPE (PbPb & AuAu)
- Observables
 - Charged hadrons
 - Multiplicity $dN_{ch}/d\eta$
 - Transverse energy $dE_T/d\eta$
 - p_T fluctuations $\delta p_T / \langle p_T \rangle$
 - Integrated anisotropic flow (JETSCAPE) $v_2\{2\}$, $v_3\{2\}$, $v_4\{2\}$
 - Identified hadrons (pion, kaon, proton)
 - Yield dN/dy
 - $\langle p_T \rangle$
 - Differential anisotropic flow (Trajectum) $v_2\{2\}(p_T)$, $v_3\{2\}(p_T)$
 - p_T spectra (Trajectum)

OUTLINE

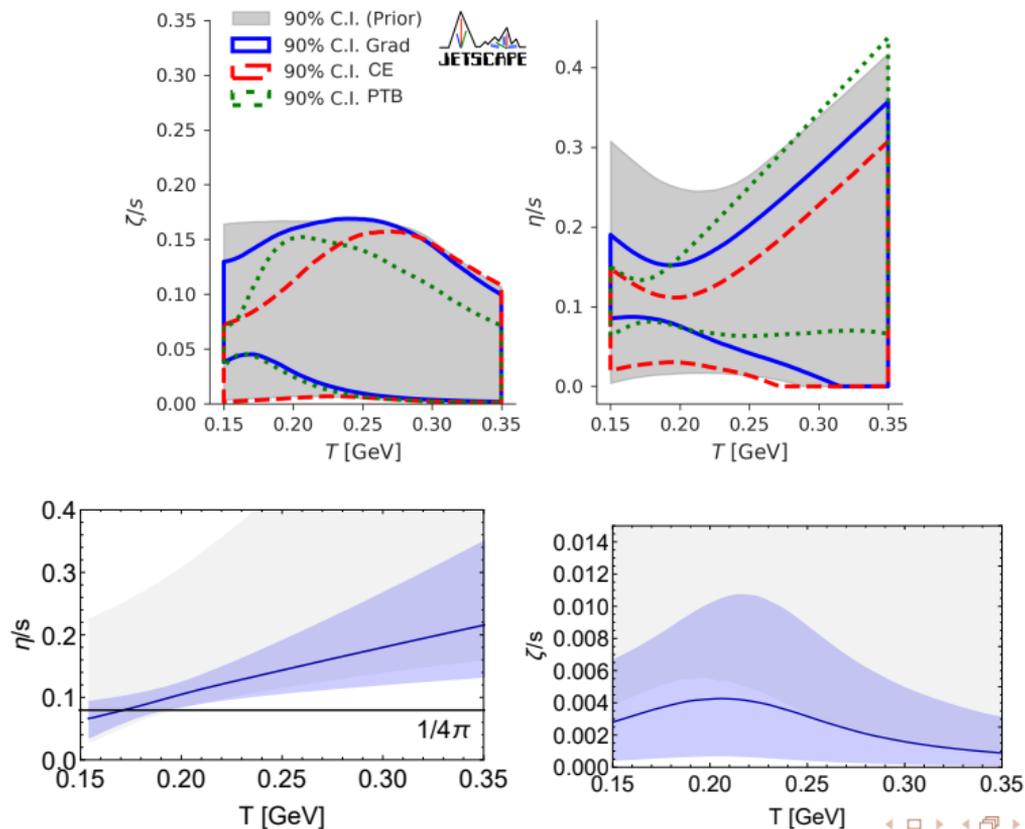
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POSTERIOR

- After a lot of work (emulator validation, closure tests, etc.) one of the main results is the multi-dimensional posterior $\Pr(p|D)$
- Visualize by marginalizing posterior over various parameters
- Point of maximum probability is the Maximum a Posteriori (MAP)



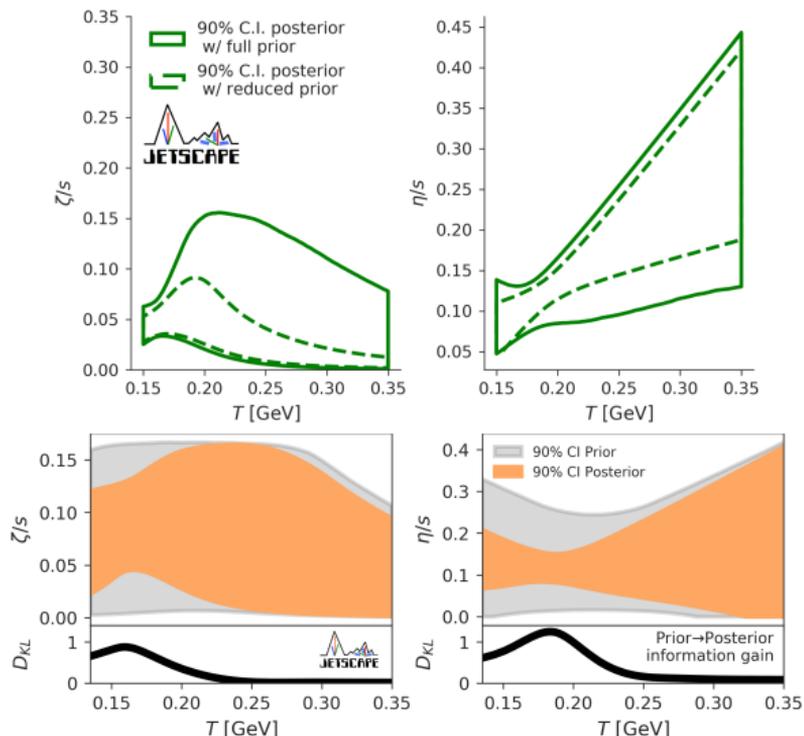
MARGINALIZED POSTERIOR: SHEAR AND BULK VISCOSITY



PRIOR $\Pr(p)$ AND INFORMATION GAIN D_{KL}

- $\Pr(p|D) \propto \Pr(p) \Pr(D|p)$
- The prior $\Pr(p)$ represents knowledge or belief about parameters **before** measurement
- There doesn't exist a neutral or uninformed choice
- The choice of prior can significantly affect the posterior
- Should compare prior and posterior. Can quantify the **information gain**

$$D_{KL} \equiv \sum_p \Pr(p) \log \left[\frac{\Pr(p)}{\Pr(p|D)} \right]$$

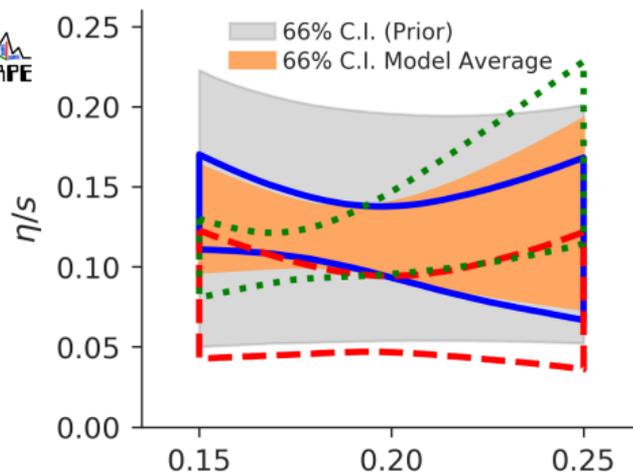
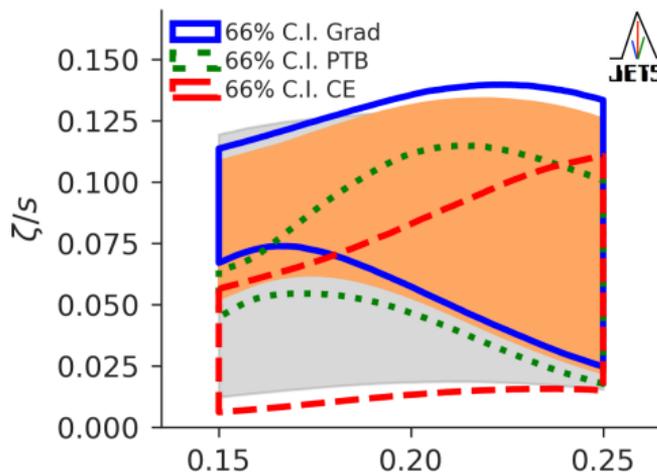


BAYESIAN MODEL AVERAGING

- Results are always interpreted in the context of a particular model — if something is missing from a model, this error does not appear in the results
- Can compare multiple models with **Bayesian evidence** $\Pr^{(i)}(D) = \int dp \Pr^{(i)}(D|p) \Pr(p)$

- E.g., models for the hadron distribution at hydro→kinetic theory transition
- Grad:PTB:CE \simeq 5000:2000:1 \implies CE disfavored by data
- Probability-weighted **Bayesian model average**:

$$\Pr_{\text{BMA}}(p, D) \propto \sum_i \Pr^{(i)}(D) \Pr^{(i)}(p|D)$$



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CONCLUSIONS

- A lot of tools and techniques have been developed for sophisticated and detailed extraction of information from data-to-model comparisons
- We now have better knowledge of QGP properties than at any time in the past, but improvements to the simulation models will make them more robust and accurate. For example:
 - Systematic exploration of viscous corrections to the freeze-out distribution function and other aspects of final hadronic evolution
 - Improvement of initial state / hydrodynamization model — is the Trento + free streaming model sufficiently flexible?
 - Inclusion of other data (more observables, more collision systems, rapidity-dependent dynamics)

EXTRA SLIDES

MODEL PARAMETERS

- A state-of-the-art simulation model was used with 18 parameters, including a parameterized initial condition, viscous hydrodynamic evolution (including temperature dependent shear viscosity and bulk viscosity), and hadron cascade afterburner
- Separate simulations were done with different models for particlization at the end of hydro (“Grad”, “PTB”, “CE”)

Parameter	Symbol	Prior	Parameter	Symbol	Prior
Norm. Pb-Pb 2.76 TeV	$N[2.76 \text{ TeV}]$	[10, 20]	temperature of (η/s) kink	T_η	[0.13, 0.3] GeV
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generalized mean	ρ	[-0.7, 0.7]	low temp. slope of (η/s)	a_{low}	[-2, 1] GeV ⁻¹
nucleon width	w	[0.5, 1.5] fm	high temp. slope of (η/s)	a_{high}	[-1, 2] GeV ⁻¹
min. dist. btw. nucleons	d_{min}^3	[0, 1.7 ³] fm ³	shear relaxation time factor	b_π	[2, 8]
multiplicity fluctuation	σ_k	[0.3, 2.0]	maximum of (ζ/s)	$(\zeta/s)_{\text{max}}$	[0.01, 0.25]
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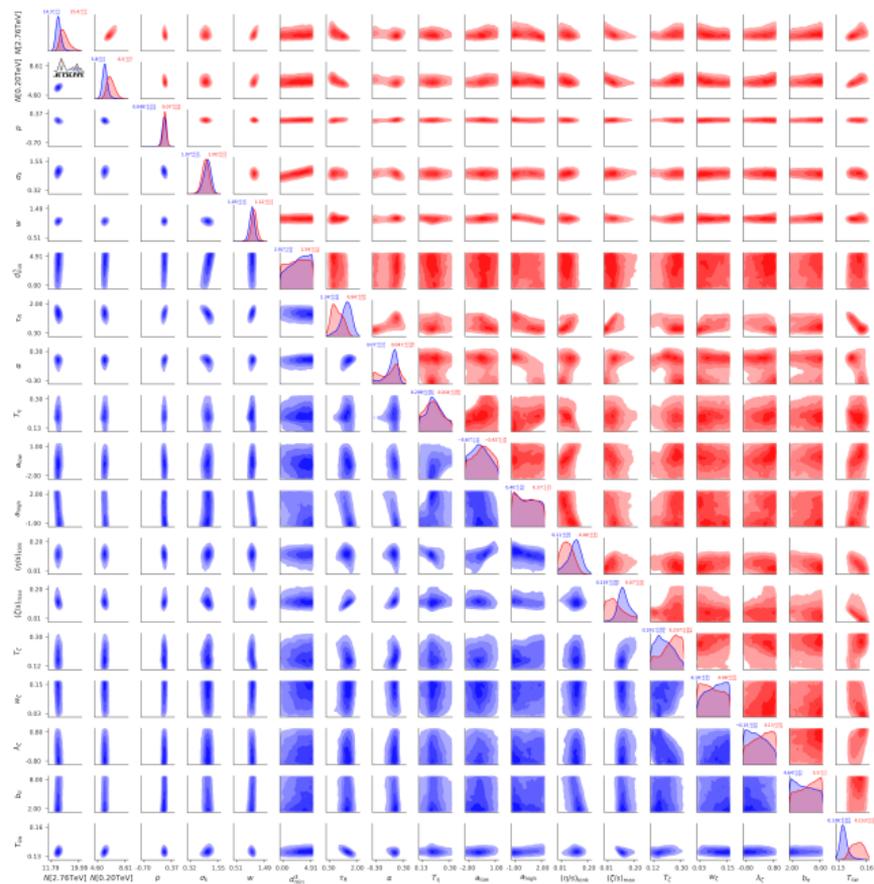
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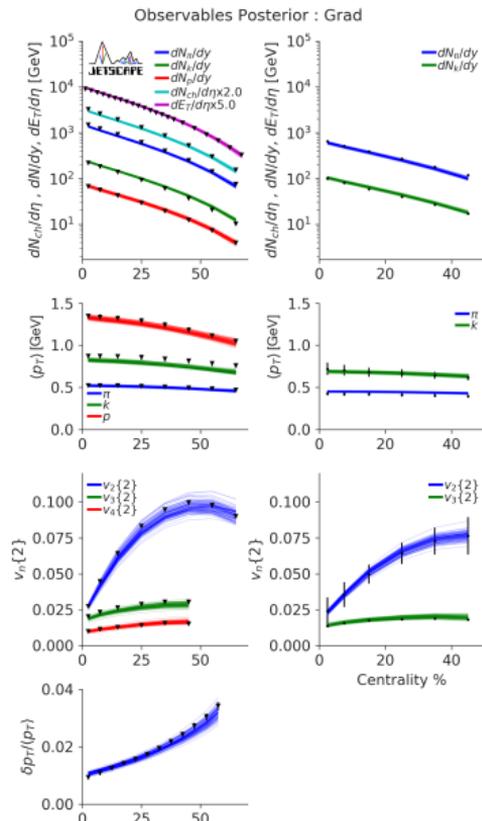
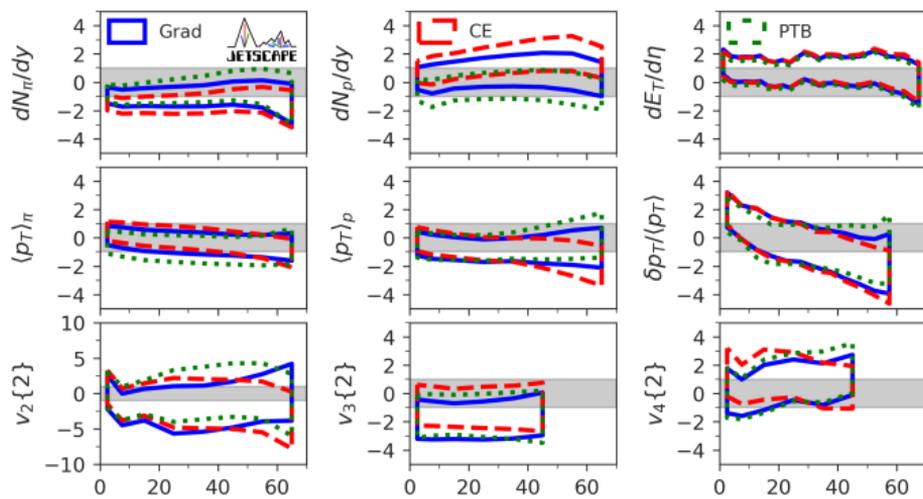
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EVALUATING MODEL SUCCESS

- Posterior does not tell the overall quality of model/fit (only *relative* quality at different parameter points)
- Must evaluate success of model separately
- E.g., direct observable comparison of posterior predictive distributions (right), or discrepancy relative to experimental uncertainty (below):



DIRECTED STUDY EXAMPLE: DEUTERONS (ARXIV:2203.08286)

- Bayesian methods can be used for smaller, directed studies
- Heavier particles such as deuterons have a larger sensitivity to bulk viscosity
- \Rightarrow Deuteron measurements can be used to better constrain ζ/s

