A. Montanari (presenter), E. Moulin and D. Malyshev on behalf of the H.E.S.S. Collaboration

IRN Terascale @ Bonn – 29th March 2022

Dark Matter annihilation signals' search in the H.E.S.S. Inner Galaxy Survey

REF. A. Montanari et al. on behalf of the H.E.S.S. Collaboration, POS(ICRC2021)511





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Introduction: WIMPs & Indirect Dark Matter search in gamma rays





Evidence

Candidates... Mass - M⊙ [eV] Primordial black holes DM **10**³⁰ - M⊳ **10**²⁰ **WIMPs 10**10 ν_s 1 Ŵ **10**-10 Axion-like particles **10**-20



. . . .

e⁺.W⁺.q,.







Candidates ...







- Cold Dark Matter paradigm
- Focusing on Weakly Interactive Massive Particles (WIMPs)
- WIMPs created thermally in the Early Universe
 - Annihilation cross section expected for thermal WIMPs ($\langle \sigma v \rangle_{th} = 3x10^{-26} \text{ cm}^3 \text{ s}^{-1}$).

• WIMPs can self-annihilate and produce Standard Model particles in the final states









- Cold Dark Matter paradigm
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 WIMPs can self-annihilate and produce Standard Model particles in the final states eventually detectable by satellite (*Fermi-LAT*) and ground-based experiments (HAWC, H.E.S.S., MAGIC, VERITAS).







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- WIMPs can self-annihilate and produce \rightarrow gamma-rays eventually detectable by H.E.S.S.
- Assuming annihilation process almost at rest;
 → A smoking-gun signature for DM is a very distinct energy cut-off, close to the DM particle mass.
- Gamma-ray flux expected from DM annihilations:









- WIMPs can self-annihilate and produce \rightarrow gamma-rays eventually detectable by H.E.S.S.
- Assuming annihilation process almost at rest;
 → A smoking-gun signature for DM is a very distinct energy cut-off, close to the DM particle mass.
- Gamma-ray flux expected from DM annihilations:

$$\frac{d\phi_{\gamma}}{dE}(E_{\gamma},\Delta\Omega) = \frac{\langle \sigma v \rangle}{8\pi m_{\rm DM}^2} \sum_{f} Br_{f} \frac{dN_{f}}{dE_{\gamma}} J(\Delta\Omega)$$
Astrophysical
factor $J(\Delta\Omega) = \int \rho^2 (r(s,\theta)) ds d\Omega$:

- Model needed for the density profile;
- Dependence on dark matter halo modeling.









Galaxy satellites of the Milky Way

- Many of them within the 100 kpc from GC: lower signal than from the GC
- o Low astrophysical background

Cosmological simulation of a Milky Way-like galaxy Aquarius, Springel et al., Nature 2008







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Dark Matter subhalos in the Milky Way halo

0

0

Lower signal than the GC region No other wavelengths conventional counterpart No conventional astrophyisical background

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Galaxy satellites of the Milky Way

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Galactic Centre (GC)

 \bigcirc

- Proximity (~8kpc)
- Possibly brightest source of DM annihilation signals: DM profile: core? cusp?
- High astrophysical
 bck / source confusion

Dark Matter subhalos in the Milky Way halo

Lower signal than the GC region No other wavelengths conventional counterpart No conventional astrophyisical background

Inner Galactic halo

Large statistics

0

Galactic diffuse
 background

Cosmological simulation of a Milky Way-like galaxy Aquarius, Springel et al., Nature 2008













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- > Aim: to provide **unprecedented sensitivity** to DM signals in the GC region.





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- Dataset: 2014-2020 observations of the GC region with the full five-telescopes H.E.S.S. array.
- 2014-2020 exposure map with IGS pointing positions:
 - Exposure up to $b \approx 6^{\circ}$;
 - Total 546 hours of high-quality data;









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- 2014-2020 exposure map with IGS pointing positions:
 - Exposure up to $b \approx 6^{\circ}$;
 - Total 546 hours of high-quality data;
 - 25 regions of interest (ROI) defined to search for DM: 0.1°-width open rings;
 - Set of exclusion regions to avoid gamma-ray contamination in the ROIs.













- Definition of the ON region: 25 ROI.
- Reflected background method:
 - OFF region:
 - Symmetric to the ON region wrt the pointing position
 Same FoV and acceptance;
 - The excluded regions are cut symmetrically

 Same solid angle size;
 - Cut overlapping areas and areas where OFF is closer to GC than the ON:
 - The DM signal in the ON region is always higher ² than in the OFF region.
- Repeated for all the 25 ROI and over the ~1300 runs.









• 2D binned Poisson likelihood function exploits spatial and spectral DM features: bins in energy (i) and space (j):

 $\mathcal{L}_{i,j}(N_{S,ij}, N_{B,ij}, \beta_{ij} | N_{ON,ij}, N_{OFF,ij}, \alpha_j) = \frac{[\beta_{ij}(N_{S,ij} + N_{B,ij})]^{N_{ON,ij}}}{N_{ON,ij}!} e^{-\beta_{ij}(N_{S,ij} + N_{B,ij})} \frac{[\beta_{ij}(N'_{S,ij} + \alpha_j N_{B,ij})]^{N_{OFF,ij}}}{N_{OFF,ij}!} e^{-\beta_{ij}(N'_{S,ij} + \alpha_j N_{B,ij})} e^{-\beta_{ij}(N'_{S,ij} + \alpha_j N_{B,ij$

- Total likelihood function: $\mathcal{L} = \prod \mathcal{L}_{i,j}$
- $N_{ON,ij}$ and $N_{OFF,ij} \rightarrow$ number of measured events in spatial ON and OFF regions;
- $N_{S,ij} + N_{B,ij} \rightarrow$ expected total number of events in the spatial ON region;
- $N'_{S,ij} + \alpha_j N_{B,ij} \rightarrow$ expected total number of events in the spatial OFF region;
- $\alpha_j = \frac{\Delta \Omega_{ON}}{\Delta \Omega_{OFF}} \rightarrow$ ratio between angular size of ON and OFF regions.







 2D binned Poisson likelihood function exploits spatial and spectral DM features: bins in energy (i) and space (j):

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- Total likelihood function: $\mathcal{L} = \prod \mathcal{L}_{i,j}$
- The systematic uncertainties can be included via a nuisance parameter;

Refs: Silverwood, et al, JCAP03, 055 (2015); Lefranc, et al. Phys. Rev. D91, 122003 (2015); CTA Dark Matter Programme (2019)

- A value of 1% is used for the determination of the limits: σ_{β} =0.01
- The value of $\boldsymbol{\beta}$ is determined via conditional maximization
 - β is computed for each energy and spatial bins, i.e., $\beta_{i,j}$.







 2D binned Poisson likelihood function exploits spatial and spectral DM features: bins in energy (i) and space (j):

 $\mathcal{L}_{i,j}(N_{S,ij}, N_{B,ij}, \beta_{ij} | N_{ON,ij}, N_{OFF,ij}, \alpha_j) = \frac{[\beta_{ij}(N_{S,ij} + N_{B,ij})]^{N_{ON,ij}}}{N_{ON,ij}!} e^{-\beta_{ij}(N_{S,ij} + N_{B,ij})} \frac{[\beta_{ij}(N'_{S,ij} + \alpha_j N_{B,ij})]^{N_{OFF,ij}}}{N_{OFF,ij}!} e^{-\beta_{ij}(N'_{S,ij} + \alpha_j N_{B,ij})} e^{\frac{-(1 - \beta_{ij})^2}{\sigma_{\beta_{ij}}}} e^{-\beta_{ij}(N_{S,ij} + \alpha_j N_{B,ij})} \frac{[\beta_{ij}(N_{S,ij} + \alpha_j N_{B,ij})]^{N_{OFF,ij}}}{N_{OFF,ij}!} e^{-\beta_{ij}(N_{S,ij} + \alpha_j N_{B,ij})} e^{\frac{-(1 - \beta_{ij})^2}{\sigma_{\beta_{ij}}}} e^{-\beta_{ij}(N_{S,ij} + \alpha_j N_{B,ij})} \frac{[\beta_{ij}(N_{S,ij} + \alpha_j N_{B,ij})]^{N_{OFF,ij}}}{N_{OFF,ij}!} e^{-\beta_{ij}(N_{S,ij} + \alpha_j N_{B,ij})} e^{\frac{-(1 - \beta_{ij})^2}{\sigma_{\beta_{ij}}}} e^{\frac{-(1 - \beta_{ij})^2}{N_{OFF,ij}}} e^{\frac{-(1 - \beta_{ij})^2}{N_{OFF,ij}}} e^{\frac{-(1 - \beta_{ij})^2}{N_{OFF,ij}}} e^{\frac{-(1 - \beta_{ij})^2}{\sigma_{\beta_{ij}}}} e^{\frac{-(1 - \beta_{ij})^2}{N_{OFF,ij}}} e^{\frac{-(1 - \beta_{ij})^2}{N_{OFF,ij}}}} e^{\frac{-(1 - \beta_{ij})^2}{N_{OFF,ij}}} e^{\frac{-(1 - \beta_{ij})^2}{N_{OFF,ij}}}} e^{\frac{-(1 - \beta_{ij})^2}{N_{OFF,ij}}} e^{\frac{-(1 - \beta_{ij})^2}{N_{OFF,ij}}}} e^{\frac{-(1 - \beta_{ij$

- Total likelihood function: $\mathcal{L} = \prod \mathcal{L}_{i,j}$
- In absence of any significant excess in the FoV:
 - → 95% C.L. upper limits on the free parameter <ov> from a log-likelihood ratio test statistics (TS). Ref. Cowan, G., Cranmer, K., Gross, E. *et al. Eur. Phys. J. C* 71, 1554 (2011)
- Computation of expected limits and containment bands:
 - Independent Poisson realizations for the ON and OFF measurements;
 - \rightarrow mean and std deviation derived from the distribution of the obtained $\langle \sigma v \rangle$ values.











Computation of upper limits on $\langle \sigma v \rangle$

- No significant excess in the FoV: \rightarrow 95% C.L. upper limits on < σ v> from the TS;
- H.E.S.S. upper limits;
- Independent Poisson realizations for N_{ON} and N_{OFF} in the computation of the expected limits;
- Containment bands plotted at 1σ and 2σ level;
- Systematic uncertainty included in the limits via a nuisance parameter in the likelihood function.







- H.E.S.S. upper limits.
- Fermi-LAT dSph and GC, HAWC dSph and GC, MAGIC Segue 1, PLANCK CMB, H.E.S.S. GC (2016) and this work.
- \rightarrow Most constraining limits in the TeV-energy range.













- IGS campaign with pointing positions up to 3.2° is very fruitful:
 - Around 546 hours of high-quality data from 2014 to 2020.
- Computation of 95% C.L. expected and observed limits including systematic uncertainty.
- VHE observations of the GC region are unique for the study of the WIMP paradigm.
- With the unprecedented IGS dataset:
 → strongest constraints obtained in the TeV mass range.
- Limits are computed in other channels \rightarrow can probe the thermal relic scale.
- The IGS is one of the legacy of the H.E.S.S. collaboration and it paves the way for CTA.





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