



# Joint galaxy clustering & lensing cosmological analysis

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# Lensing : coherent distortion of light rays

Credits: Matthew Becker, KITP workshop 2013

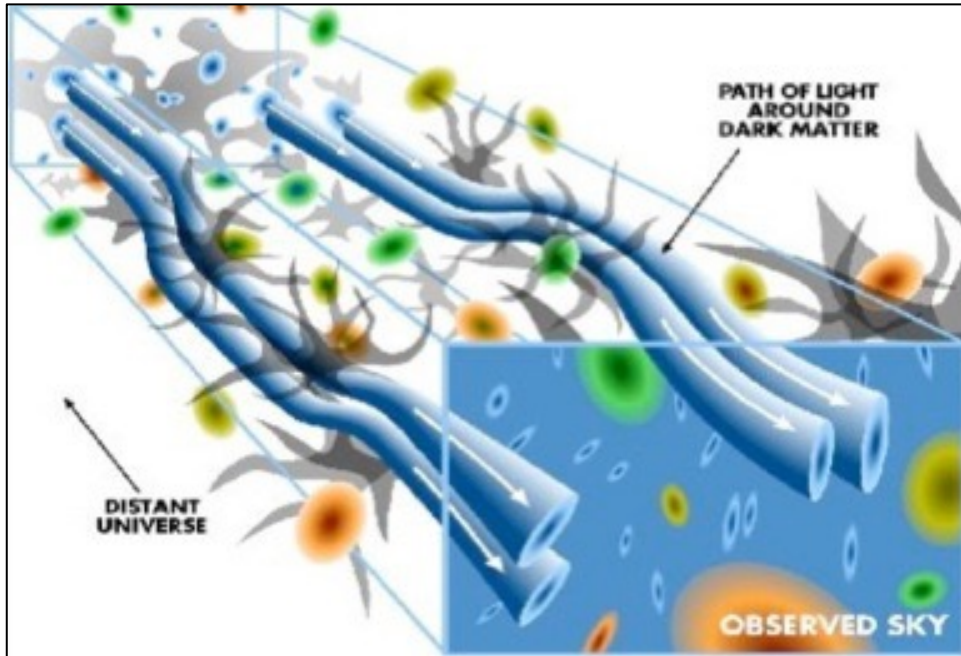
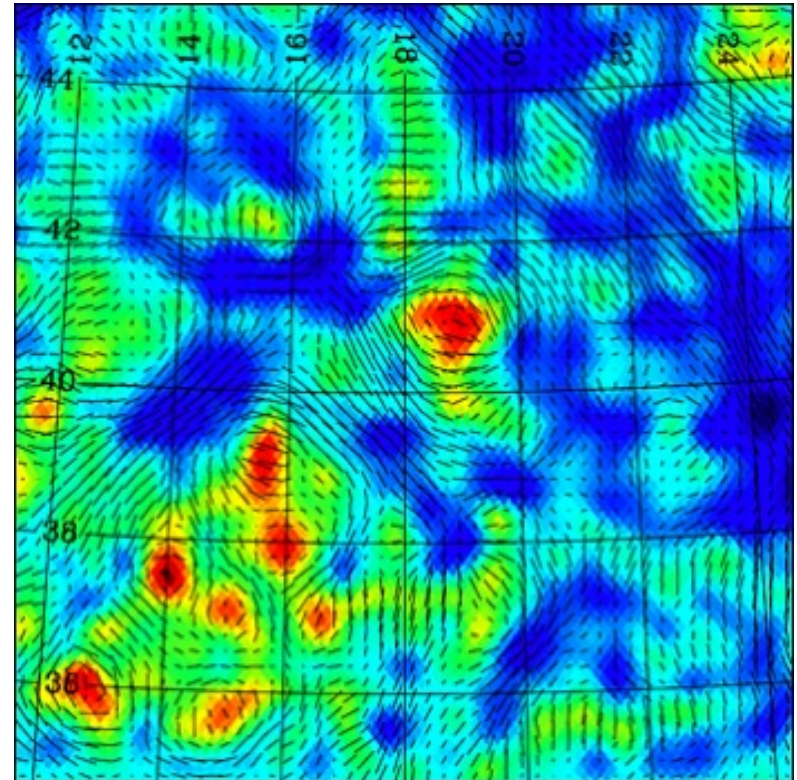


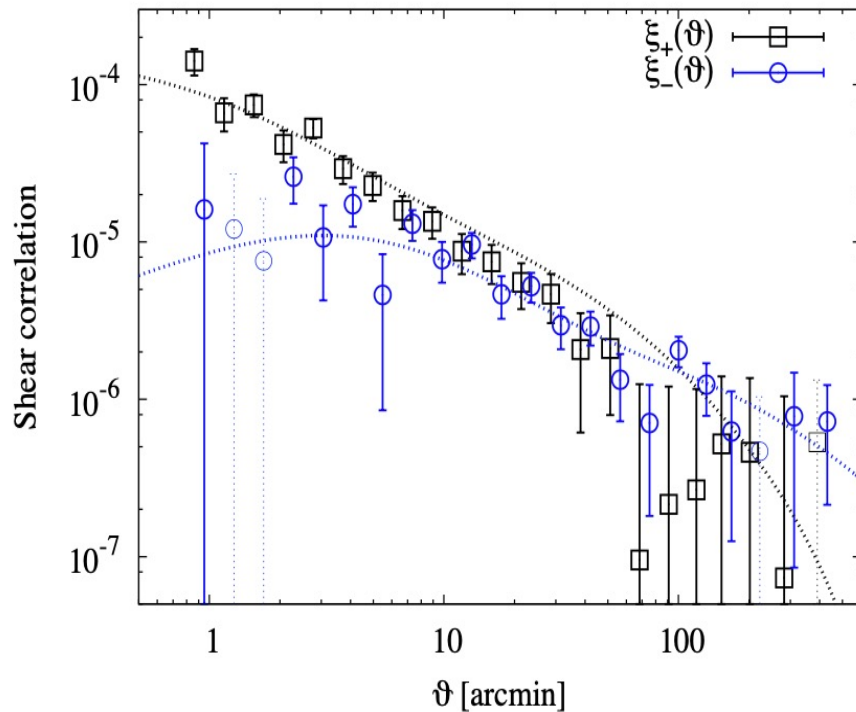
Illustration of the weak-lensing effect  
Tiny effect 1% => cosmic shear



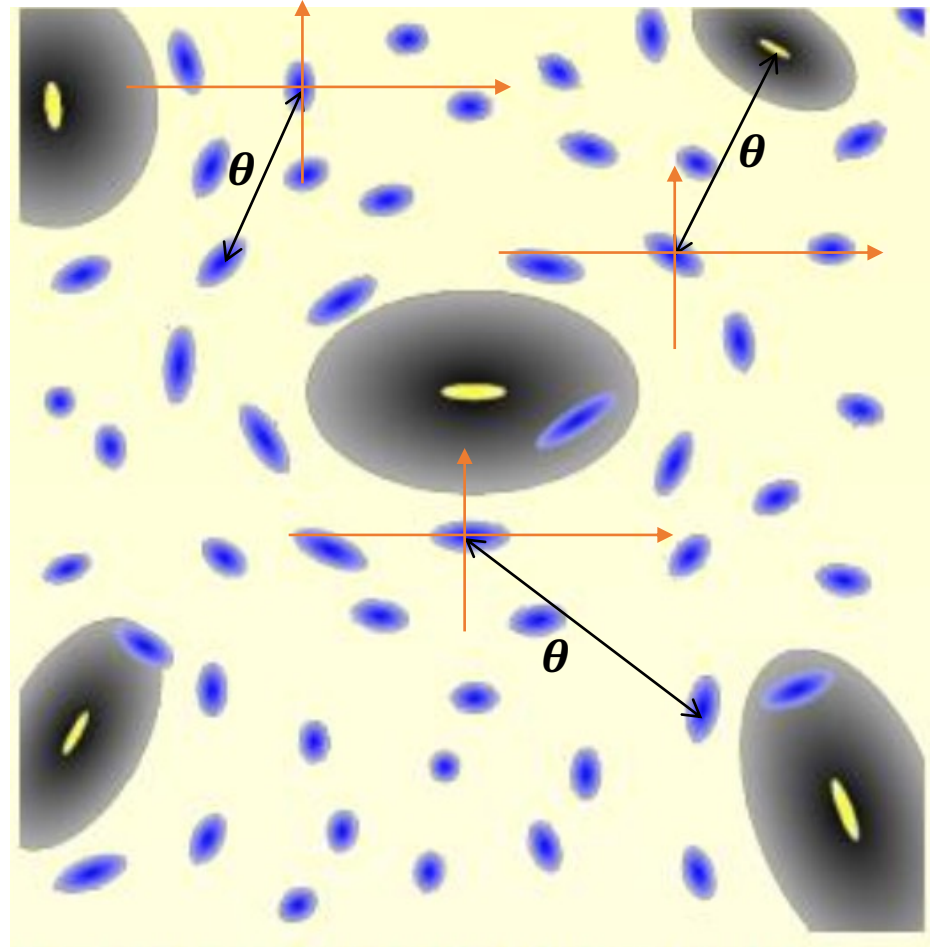
Observation of galaxy clusters (in red)  
and voids (in blue)

# Cosmic-shear estimator

We measure the average signal as a function of separation



Kilbinger et al. 2012 (CFHTLenS)



$$\text{Cosmic shear estimator } \xi_{\pm}(\theta) = \langle e_1 e_1 \rangle \pm \langle e_2 e_2 \rangle(\theta) \quad 3$$

# Cosmic shear equations

The effective projected convergence in physical units

$$\kappa = \Sigma / \Sigma_{crit} = \nabla\varphi / 2 \rightarrow \kappa_{eff}(\theta, D_s) = \frac{1}{c^2} \int_0^{D_s} dD_d \frac{D_d D_{ds}}{D_s} \nabla\Phi(D_d\theta, D_d)$$

In comoving units, and as a function of density contrast

$$\kappa_{eff}(\theta, w_s) = \frac{3H_0^2\Omega_m}{2c^2} \int_0^{w_s} dw_d n_d(w_d) \frac{w_d(w_s - w_d)}{w_s} \frac{\delta[w_d\theta, w_d]}{a(w_d)}$$

Power-spectrum

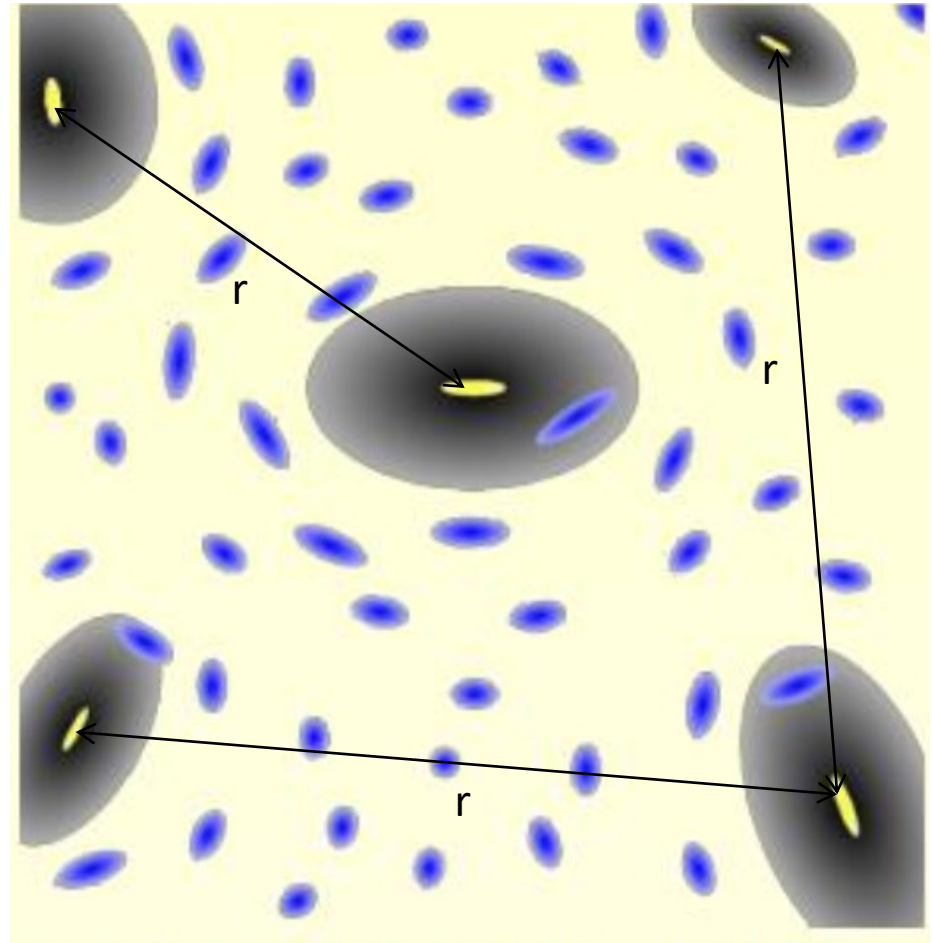
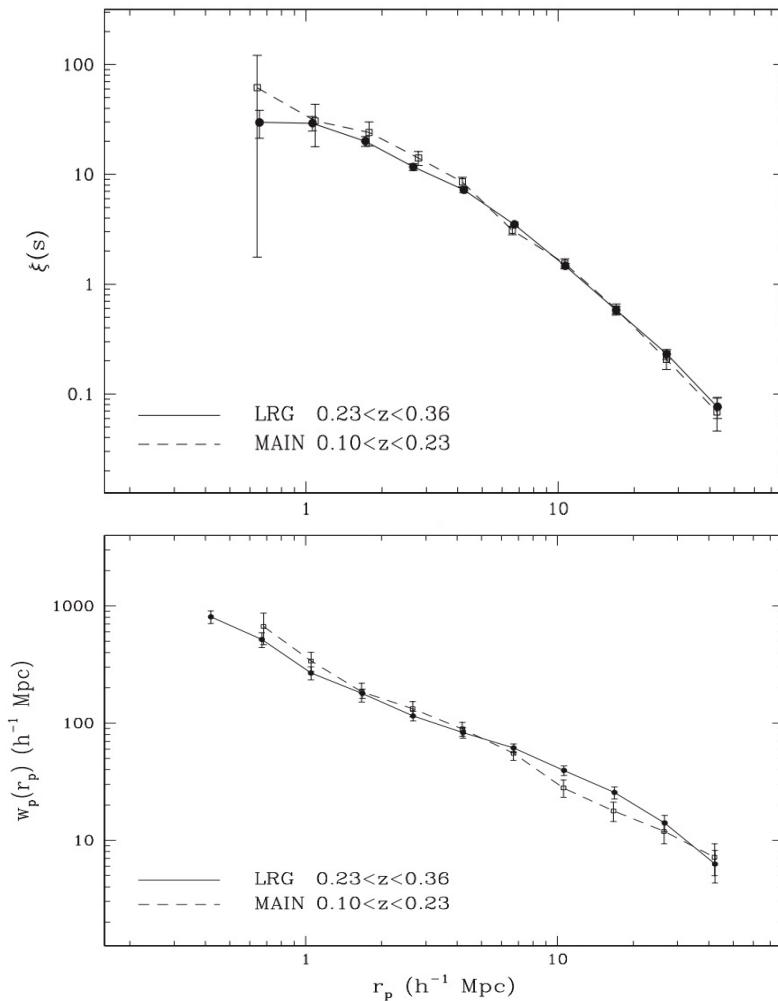
$$C_\kappa(\ell) = \int_0^{w_H} dw \frac{q_d(w)q_s(w)}{w^2} P_\delta\left(\frac{\ell}{w}, w\right)$$

$$\text{with } q_i(w) = \frac{3H_0^2\Omega_m}{2c^2} \frac{w}{a(w)} \int_w^{w_H} dw' n_i(w') \frac{w'-w}{w'} \text{ and } P_\delta(k) \propto \sigma_8^2$$

Correlation function

$$\xi_\pm(\theta) = \frac{1}{2\pi} \int d\ell C_\kappa(\ell) J_{0,4}(\ell\theta)$$

# Clustering estimators



# GGL in angular or comoving scales

Power-spectrum

$$C_{g\gamma}(\ell) = \int_0^{w_H} dw \frac{n_d(w) q_s(w)}{w^2} \text{ br } P_\delta\left(\frac{\ell}{w}, w\right)$$

Correlation functions

$$\gamma_t(\theta) = \frac{1}{2\pi} \int d\ell C_{g\gamma}(\ell) J_2(\ell\theta)$$

$$\Sigma(R) = \Omega_m \rho_c \int \xi_{gm}(\sqrt{R^2 + w^2}) dw$$

Differential surface mass density

$$\Delta\Sigma(R) = \gamma_t(w_d\theta) \Sigma_{\text{crit}}(w_d)$$

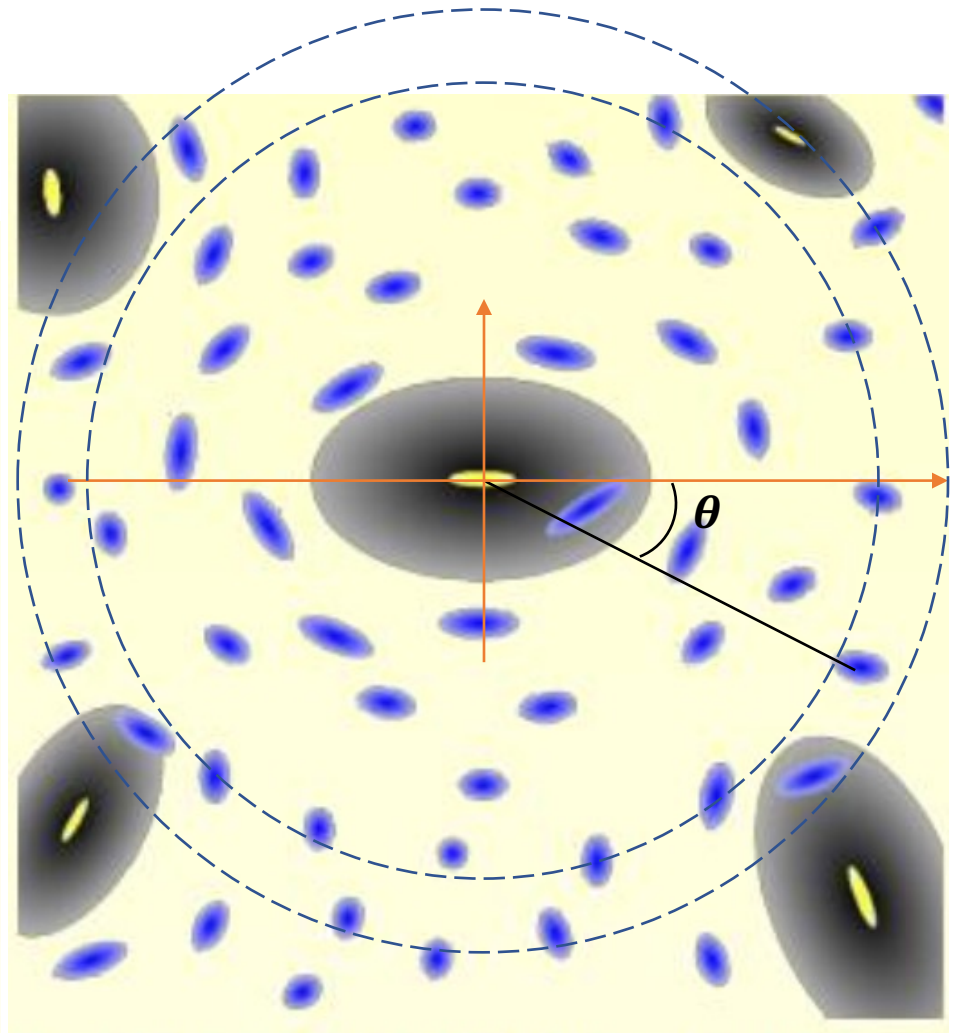
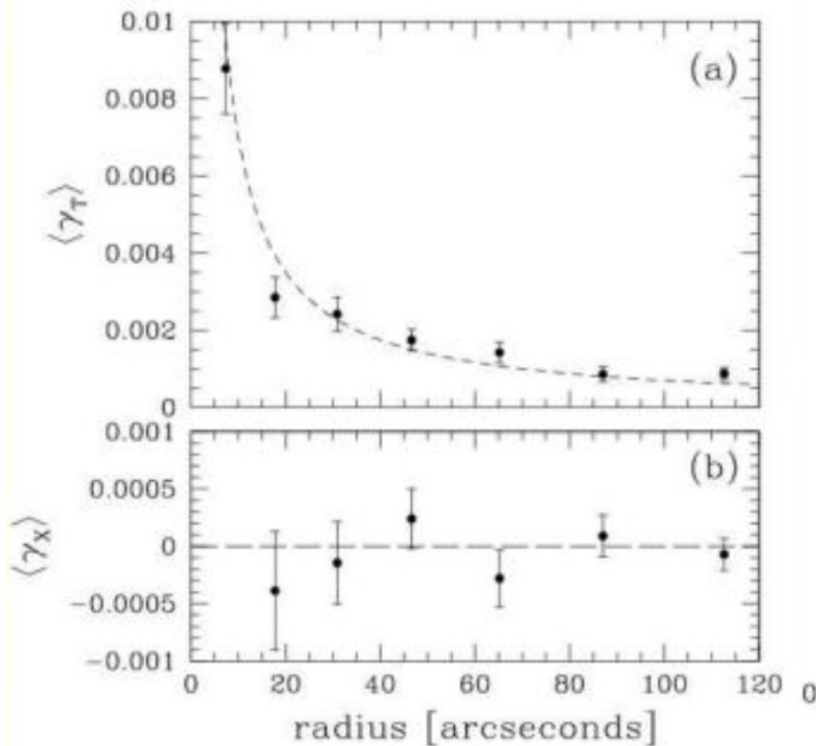
$$\Delta\Sigma(R) = \overline{\Sigma_{gm}}(R) - \Sigma_{gm}(R)$$

$$\text{with } \Sigma_{\text{crit}}^{-1}(w_d) = \int_{w_d}^{w_H} dw_s n_s(w_s) \frac{1}{\Sigma_{\text{crit}}(w_d, w_s)} \quad \text{and} \quad \Sigma_{\text{crit}}(w_d, w_s) = \frac{c^2}{4\pi G} \frac{w_s}{w_d w_{ds}}$$

=> This is what people now use in wide field analysis

# Galaxy-Galaxy lensing estimator

We measure the average signal around many lenses

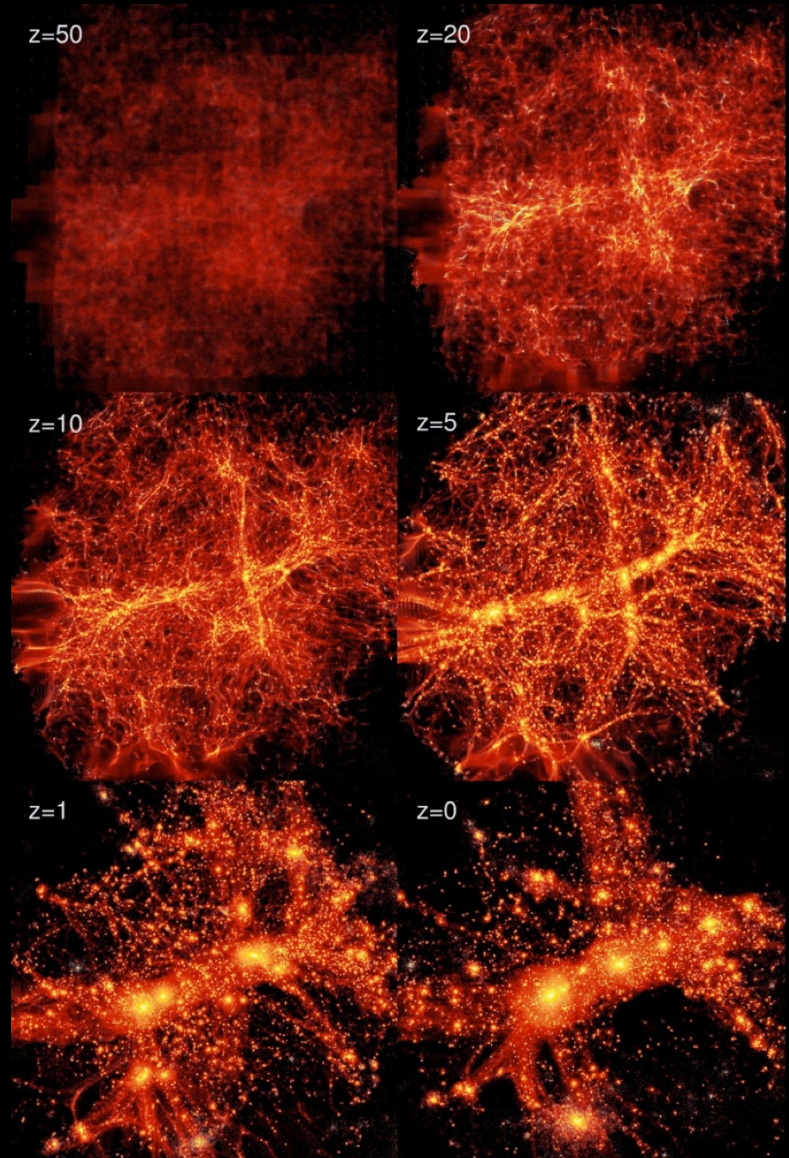
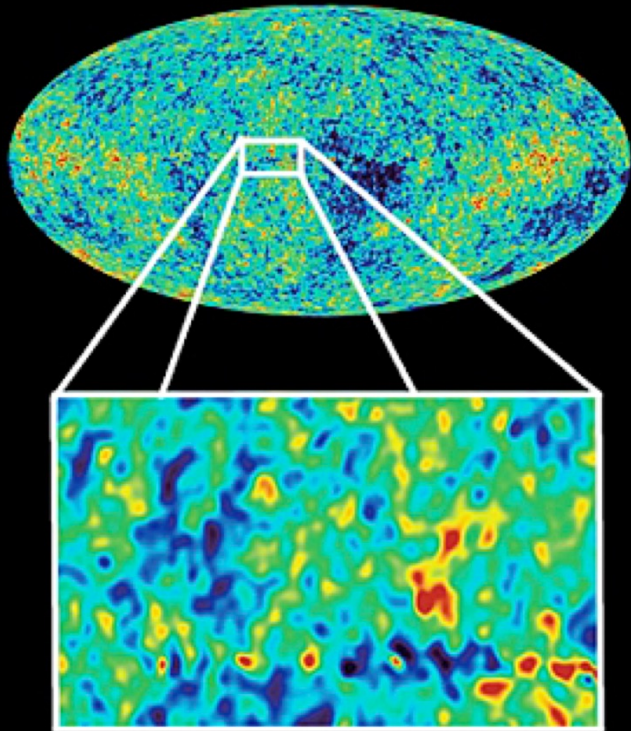


Hoekstra et al. 2004

$$\text{GGL estimator } \langle \gamma_t \rangle = -\langle e_1 \cos(2\theta) + e_2 \sin(2\theta) \rangle$$

# Growth of perturbations

Measure the fluctuations of density as a function of redshift ( $z$ ) and scale ( $k$ )





# Redshift space distortion

Galaxies possess coherent “peculiar velocities” on top of the overall cosmological expansion



# Redshift space distortion

- These velocities are driven by the matter distribution, according to **gravitational physics**

- For example in **linear perturbation theory**:

$$\theta = \vec{\nabla} \cdot \vec{v} = -f \delta_m$$

- in terms of the **growth rate**  $f = d(\ln G) / d(\ln a)$

$G(a)$  : **Growth factor** of the Universe

- The dependence of  $f$  on scale and time is a key discriminator between gravity models

# Lensing in GR

In the perturbed Friedman-Robertson-Walker metric

$$ds^2 = (1 + 2\Psi)dt^2 - a^2(1 + 2\Phi)dx^2$$

  
time

  
space

$a(t)$  : scale factor  $1 \rightarrow 0$  (Today  $\rightarrow$  Big Bang)

$\Phi$  and  $\Psi$  : Bardeen potentials. In GR  $\Phi = -\Psi$

Lensing is a projected effect  $\Rightarrow$  sensitive to  $\nabla^2(\Phi - \Psi)$   
along the line-of-sight

# Testing GR with RSD + Lensing

1. Smoking gun observational estimator (Zhang et al. 2007)

$$E_G = \frac{\text{Lensing}}{\text{RSD}} = \frac{\nabla^2(\Phi - \Psi)}{3H_0^2 a^{-1} f \delta_m} \rightarrow \frac{\Omega_m}{f} \text{ in GR}$$

2. Phenomenological model (Amendola et al. 2008)

$$\begin{aligned} 2\nabla^2\Psi &= 8\pi G a^2 (1 + \mu) \rho_m \delta_m \\ 2\nabla^2(\Phi - \Psi) &= 8\pi G a^2 (1 + \Sigma) \rho_m \delta_m \end{aligned}$$

# Historical review before Stage III

Combined clustering and lensing on wide field surveys :

Hoekstra et al. 2002 on 50deg<sup>2</sup> field (RCS+VIRMOS-DESCART surveys) to study galaxy bias and galaxy-lensing correlation factor  $r$ . They used  $M_{\text{ap}}(\boldsymbol{\theta})$ ,  $N_{\text{ap}}(\boldsymbol{\theta})$  and  $MN_{\text{ap}}(\boldsymbol{\theta})$

Sheldon et al. 2004 used SDSS (3800 deg<sup>2</sup>) and also found  $r \sim 1$ . They used  $\Delta\Sigma$  and  $w_p(r_p)$

Simon et al. 2007 used GaBoDS (15 deg<sup>2</sup>) to measure bias at redshift  $z \sim 0.6$ . They used  $M_{\text{ap}}(\boldsymbol{\theta})$ ,  $N_{\text{ap}}(\boldsymbol{\theta})$  and  $MN_{\text{ap}}(\boldsymbol{\theta})$ .

Reyes et al. 2010 used SDSS to measure  $E_G$ . They used  $\Delta\Sigma$  and  $w_p(r_p)$  and  $\beta$  from Tegmark et al. 2006

Jullo et al. 2012 used COSMOS (1 deg<sup>2</sup>) to measure bias up to  $z \sim 1$ . They used  $M_{\text{ap}}(\boldsymbol{\theta})$ ,  $N_{\text{ap}}(\boldsymbol{\theta})$  and  $MN_{\text{ap}}(\boldsymbol{\theta})$ .

Mandelbaum et al. 2012 used SDSS to constrain cosmological parameters. They used  $\Delta\Sigma$  and  $w_p(r_p)$ .

Leauthaud et al. 2011 used COSMOS (1 deg<sup>2</sup>) to study the SHMR up to  $z \sim 1$ . They used  $\Delta\Sigma$  and  $w_p(r_p)$ .

Coupon et al. 2014 used CFHTLens/VIPERS (23.1 deg<sup>2</sup>) to study the SHMR up to  $z \sim 0.8$ . They used  $\Delta\Sigma$  and  $w_p(r_p)$ .

More et al. 2014 used CFHTLens+BOSS (105 deg<sup>2</sup>) to estimate  $\Omega_m$  and  $\sigma_8$ . They used  $\Delta\Sigma$  and  $w_p(r_p)$ .

Leauthaud et al. 2016 used CFHTLens+Stripe82+BOSS (250 deg<sup>2</sup>) and found small value of  $S_8$ . They used  $\Delta\Sigma$  and  $w_p(r_p)$

Blake et al. 2016 used CFHTLens+RCS+WiggleZ+BOSS (466 deg<sup>2</sup>) to measure  $E_G$  at  $z \sim 0.5$ . They used  $\Delta\Sigma$  and RSD. Linear bias.

de la Torre et al. 2017 used CFHTLens+VIPERS (23.5 deg<sup>2</sup>) to measure  $E_G$  at  $z \sim 0.8$ . They used  $\Delta\Sigma$  and RSD. Non linear bias.

Amon et al. 2017 used KiDS+2dFLenS+GAMA+BOSS (350deg<sup>2</sup>) to measure  $E_G$  up to  $z < 0.9$ . They used  $\Delta\Sigma$  and RSD. Linear bias

Jullo et al. 2019 used CFHTLens+Stripe82+CMASS (250 deg<sup>2</sup>) to measure  $E_G$  at  $z \sim 0.5$ . They used  $\Delta\Sigma$  and RSD. Non linear bias

# Historical review before Stage III Summary

About 20 years of debate whether using  $\gamma_t(\boldsymbol{\theta})$  or  $\Delta\Sigma(R)$

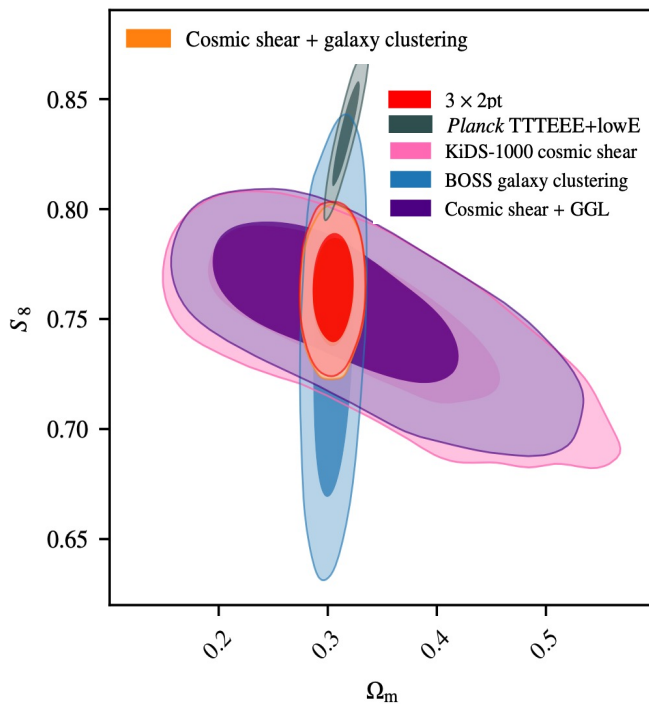
- It depends on the survey depth (Shirasaki et al. 2018)
- For magnification bias, it's cleaner to use  $\gamma_t(\boldsymbol{\theta})$  (personal opinion)

Long lasting use of lensing mass aperture  $M_{\text{ap}}(\boldsymbol{\theta})$  for cosmic-shear. Now  $\xi_{\pm}(\boldsymbol{\theta})$  used instead.

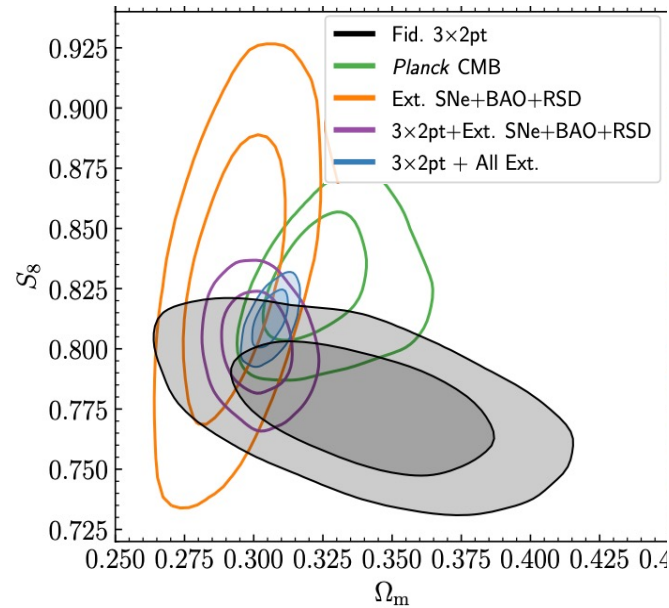
Joint RSD and lensing full-scale modeling with Blake et al. (2016) and de la Torre et al. (2017)

# STAGE III 3x2pt results

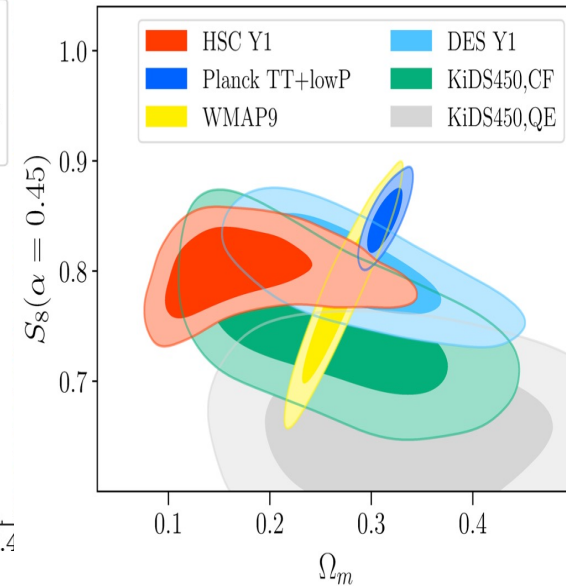
KiDS -- Heymans et al. 2021



Dark Energy Survey collaboration 2021

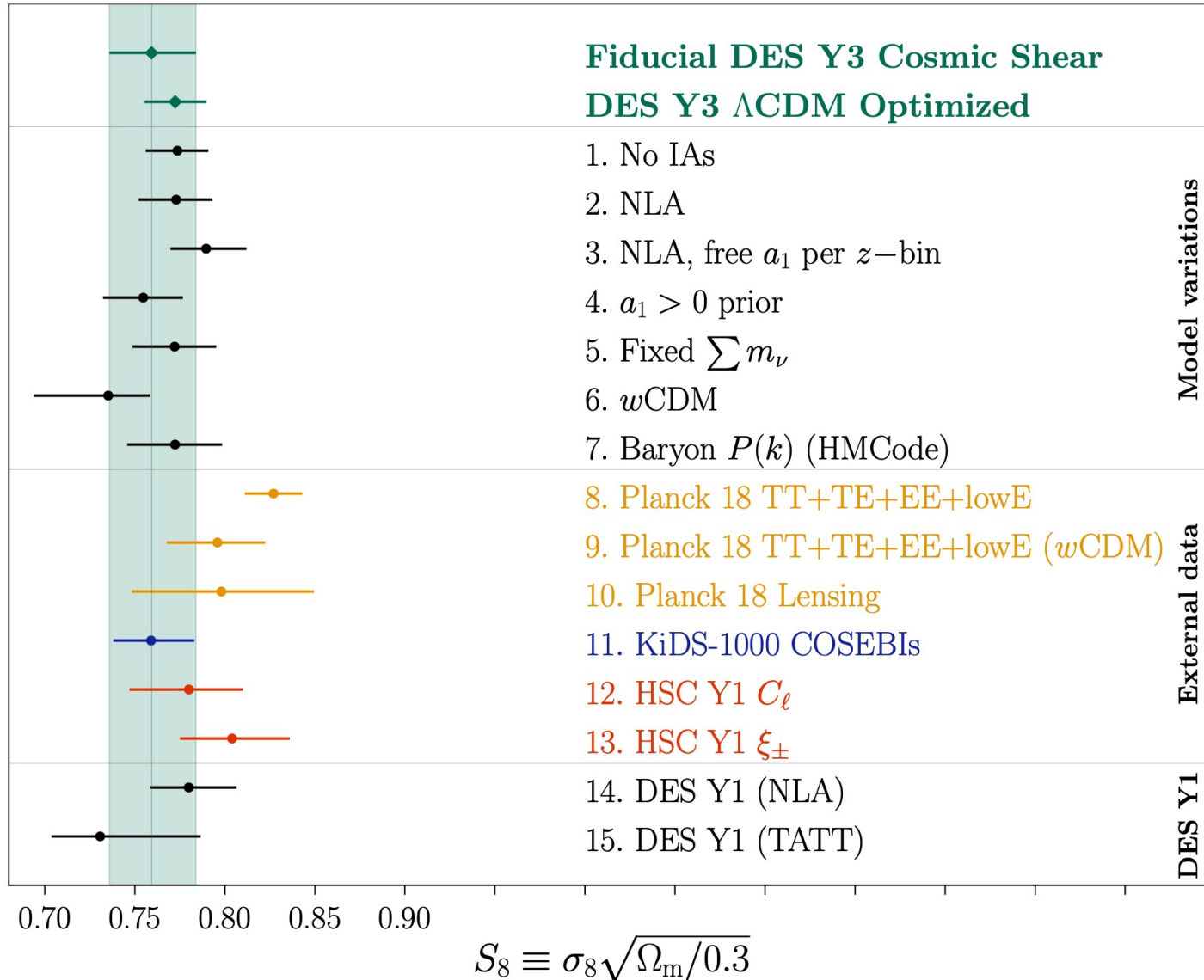


HSC -- Hikage et al. 2019



# STAGE III 3x2pt results

Secco et al. 2022

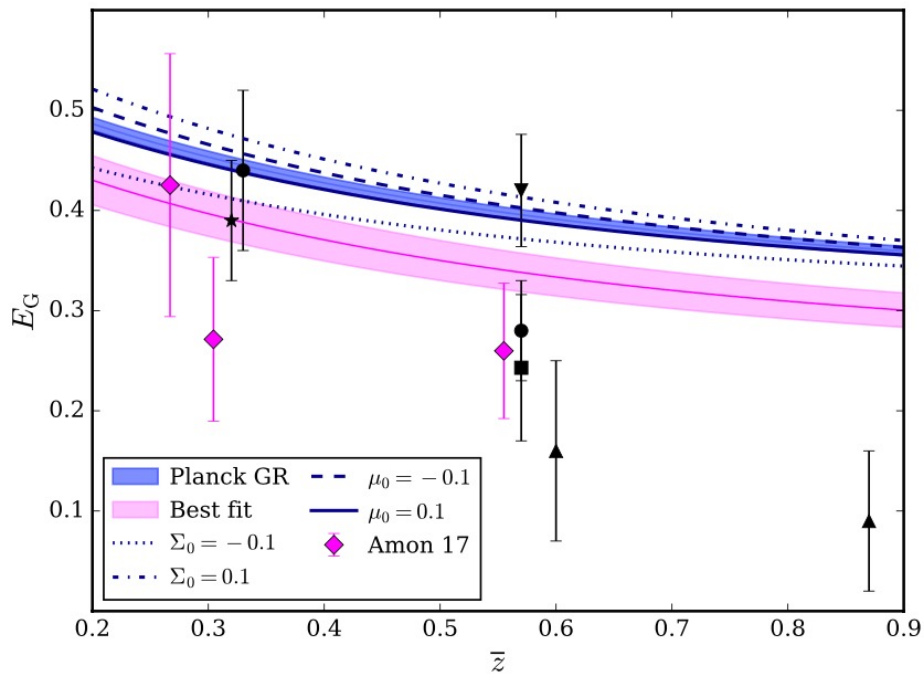




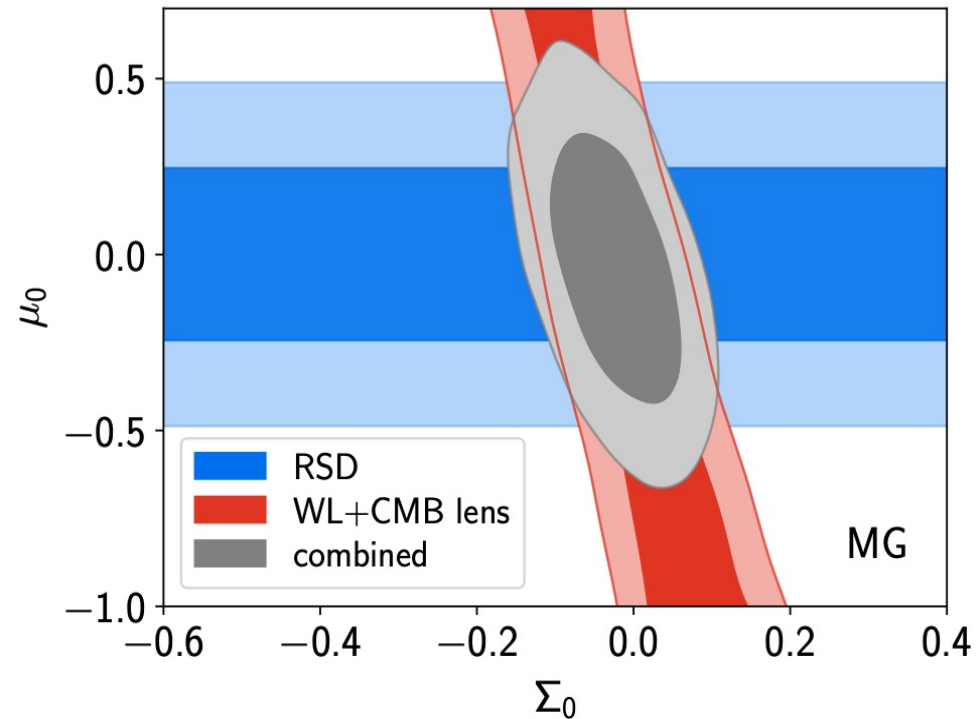
# Joint lensing & RSD

## Latest constraints on MG

KiDS collaboration, Amon et al. 2017



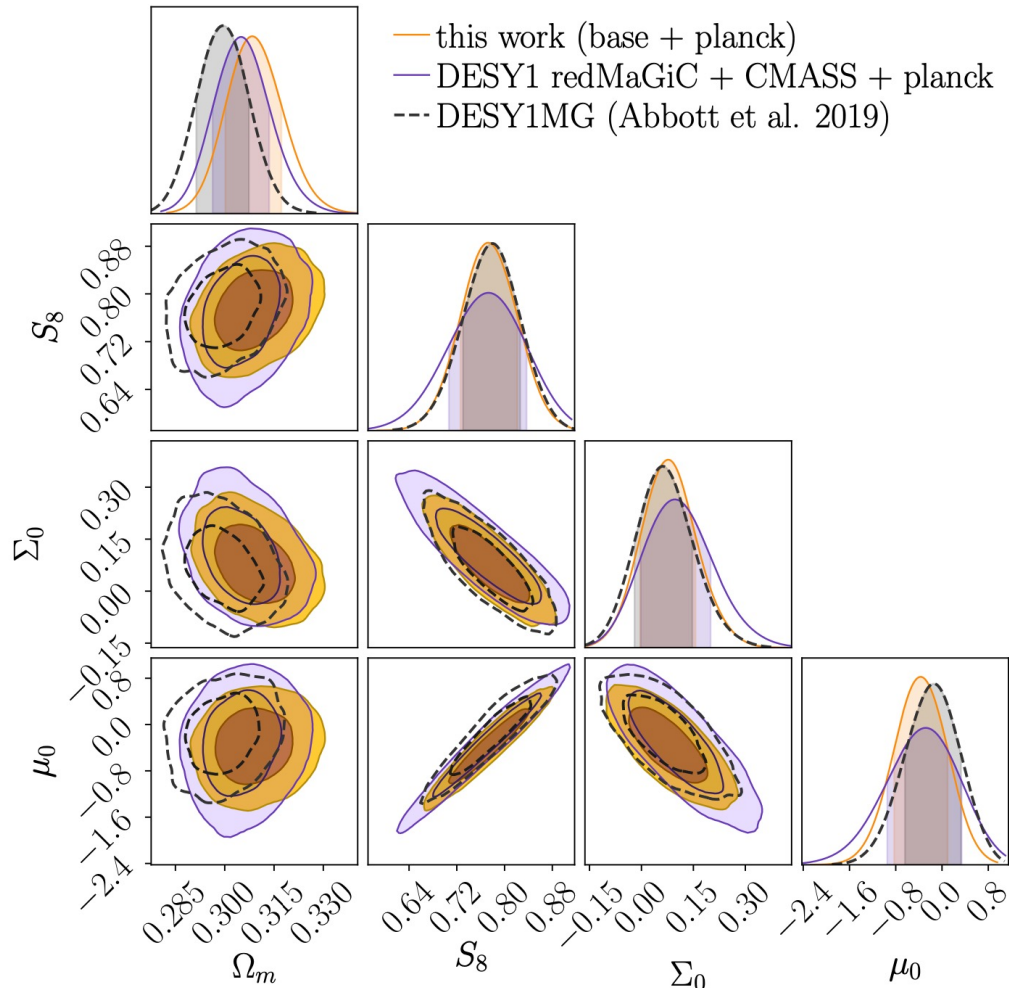
eBOSS collaboration 2020



# Joint lensing & RSD

## DES-Y1 MG results

Lee et al. 2021



=> In DES-Y3, since  $S_8$  is a bit larger then  $\Sigma_0$  and  $\mu_0$  should be more in agreement with GR

# Some systematic errors to deal with

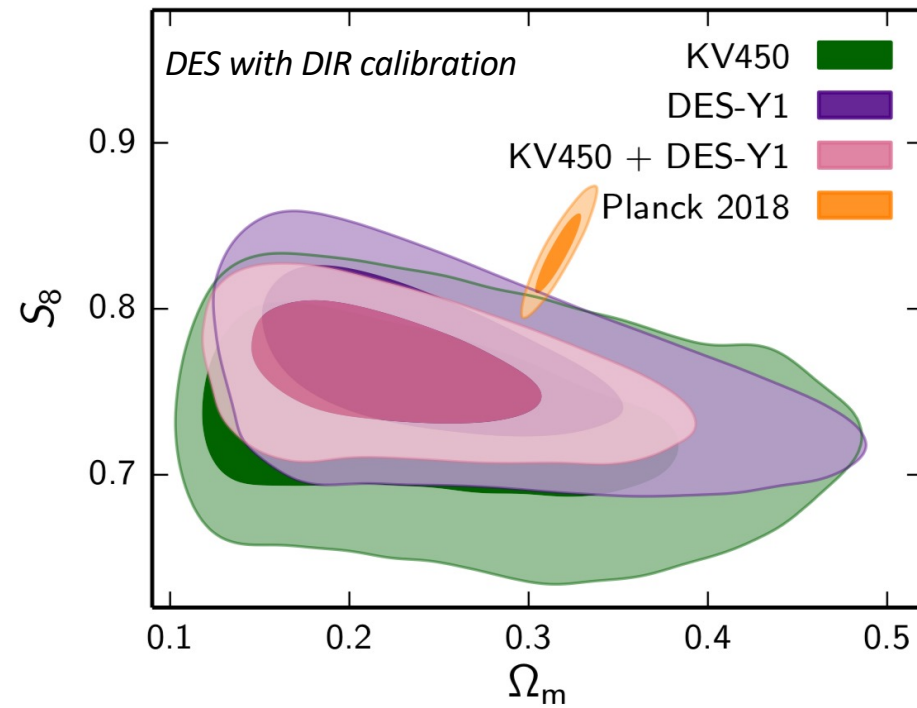
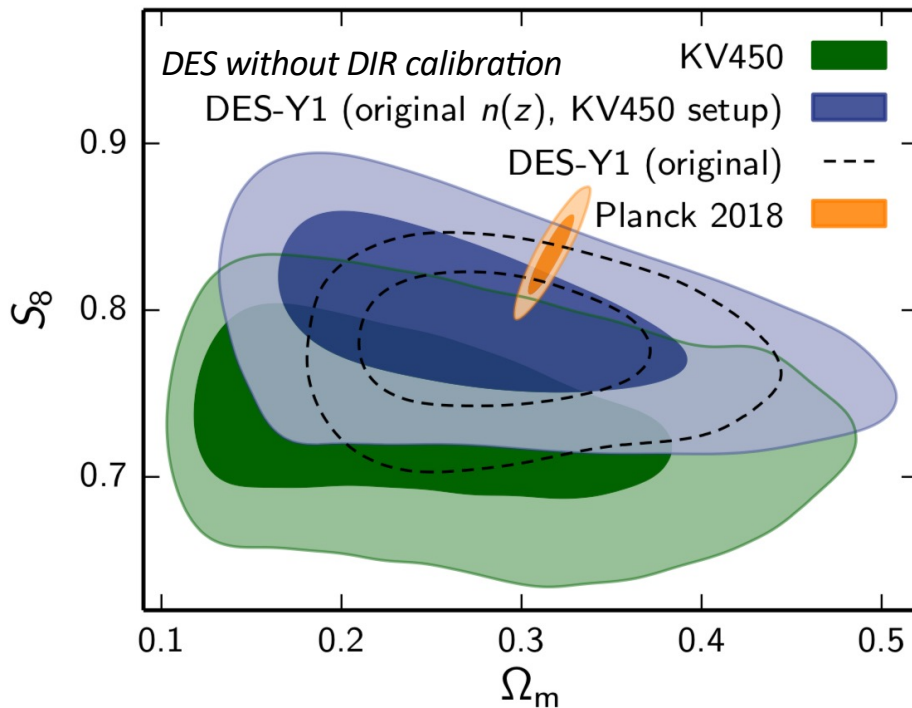
- Redshift estimations
- Shape noise
- Modeling issues:
  - Non-linear bias modeling
  - Magnification bias
  - Galaxy Halo Connection biases

In HSC Y3 (Li et al 2022) just identified a couple of issues that impede the final cosmological analysis:

- i) PSF model shape residual
- ii) Star-galaxy shape correlation additive systematics

# Revealing redshift bias

Joudaki et al. 2019



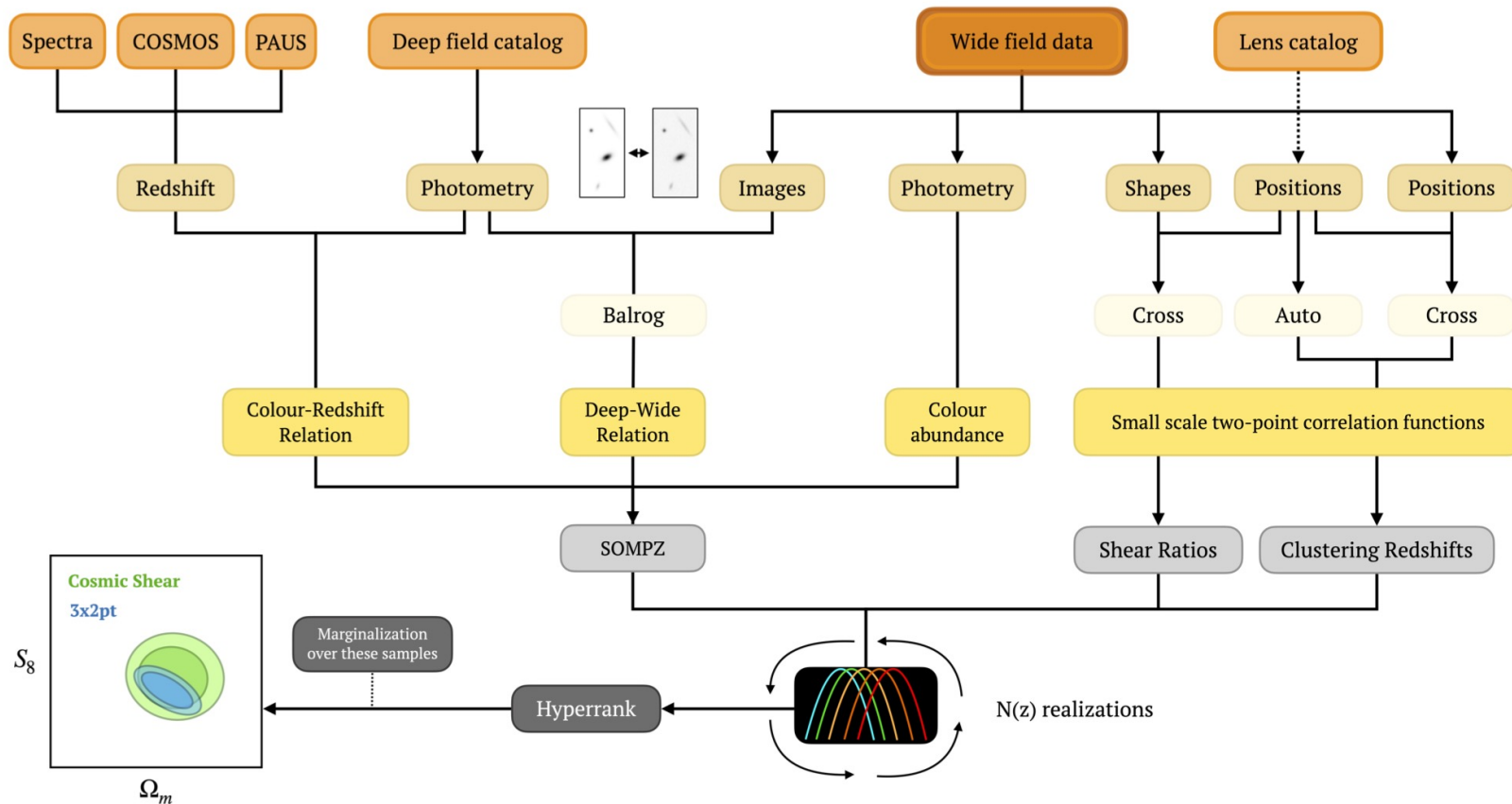
⇒ The calibration of the DES-Y1 data with spectroscopic redshift (DIR method) as used in the KIDS analysis leads to **agreement between the 2 surveys**

⇒ Enhanced discrepancy at  $2.5\sigma$  with Planck result => **new cosmology?**

# DES-Y3 analysis

## Photometric redshift calibration

Myles et al. 2021

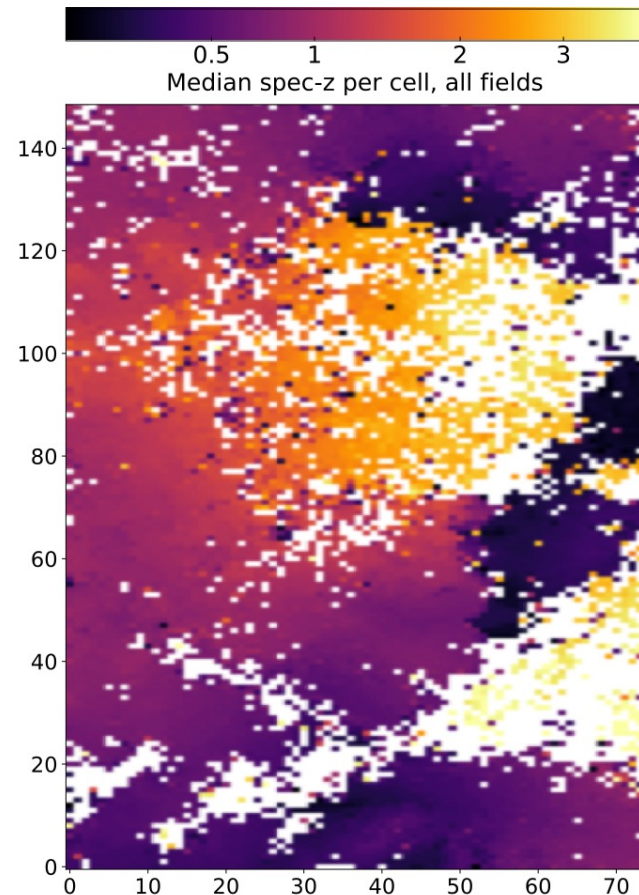


# Collective spectroscopic redshift effort

Master et al. 2019

Goal: calibrate Euclid & WFIRST

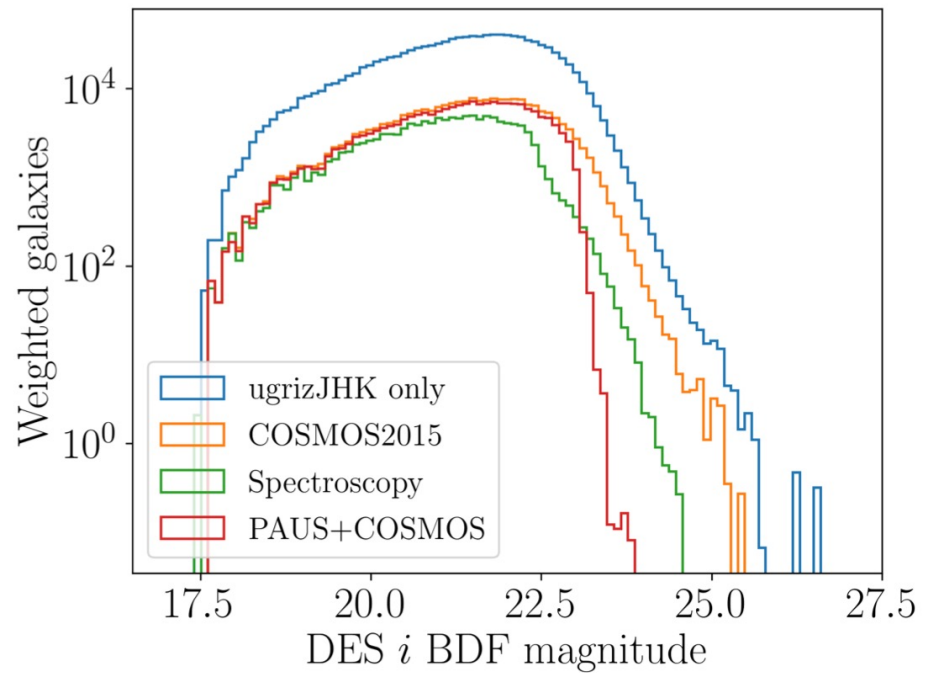
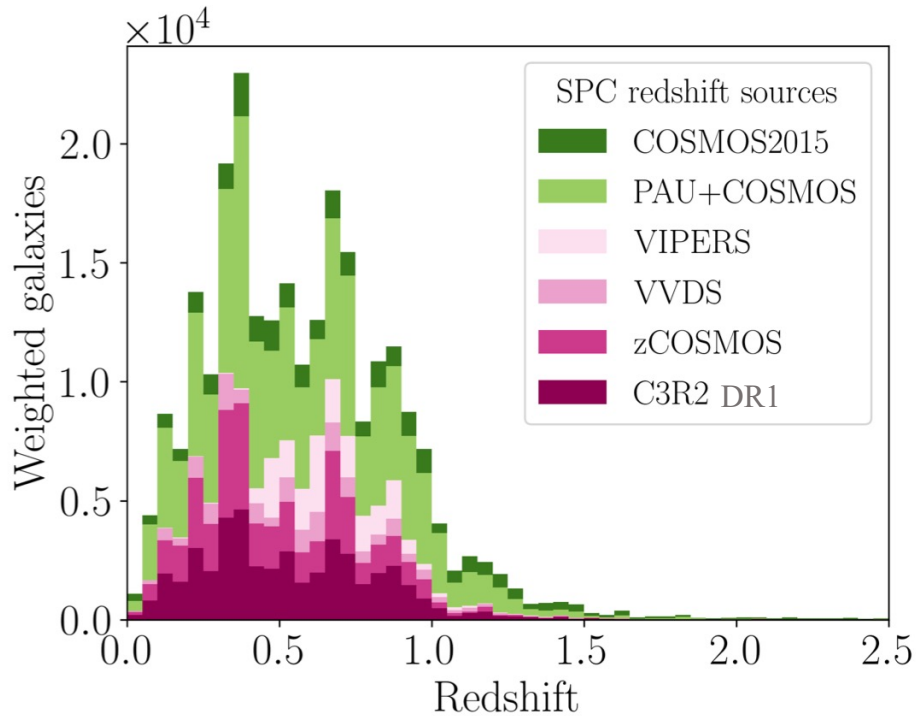
- Observed fields
  - VVDS-2h, COSMOS, EGS
- Keck observations
  - DEIMOS, LRIS, MOSFIRE
- Current status
  - DR1 1283 redshifts
  - DR2 4454 redshifts



# DES-Y3 analysis

## Spectroscopic calibration samples

Myles et al. 2021

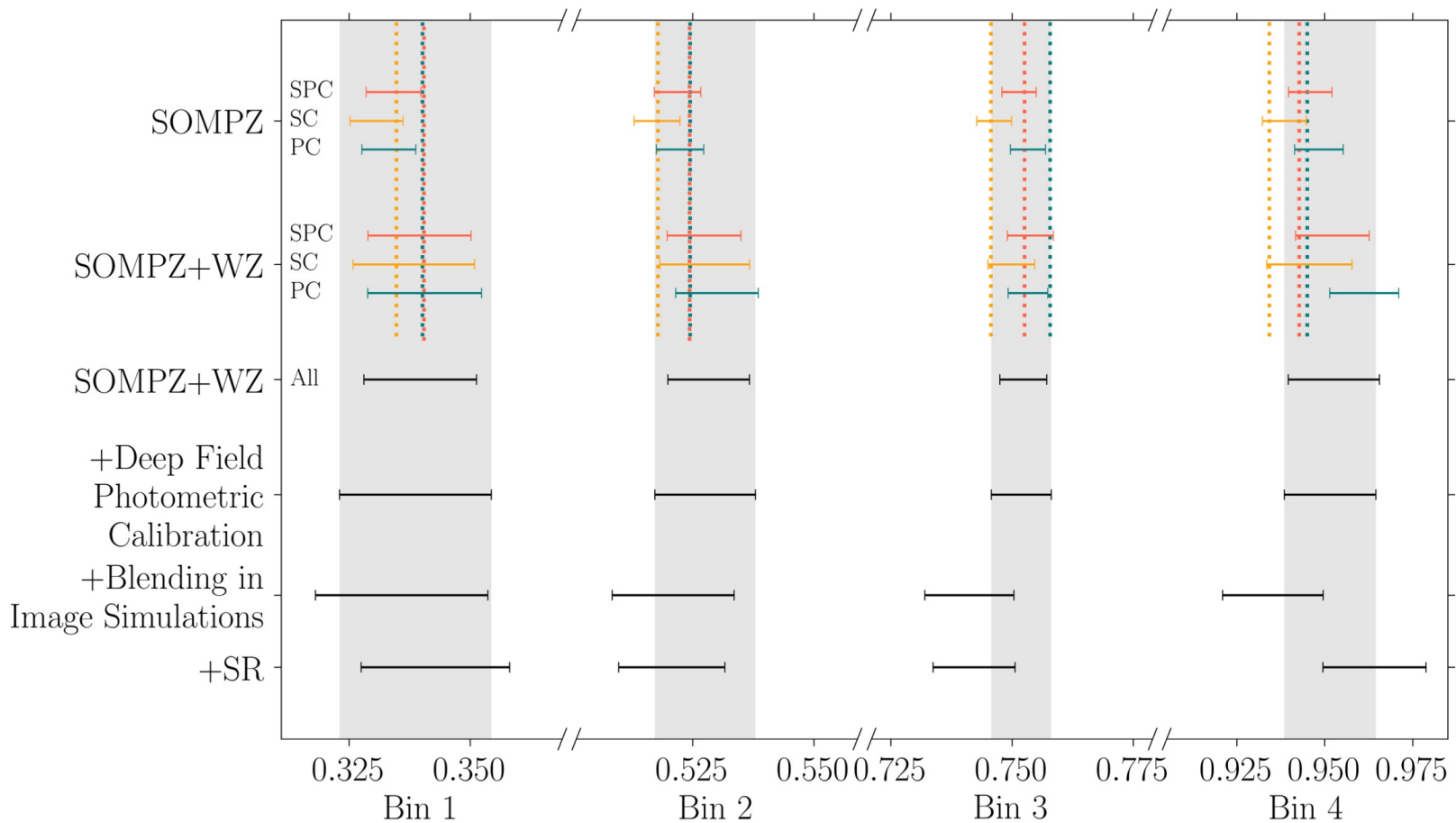


- ⇒ Several problems of completeness, hence the multiple techniques of calibrations
- ⇒ It seems not a good idea to use BOSS+eBOSS for redshift calibration

# DES-Y3 analysis

## Everytime adding more systematics

Myles et al. 2021

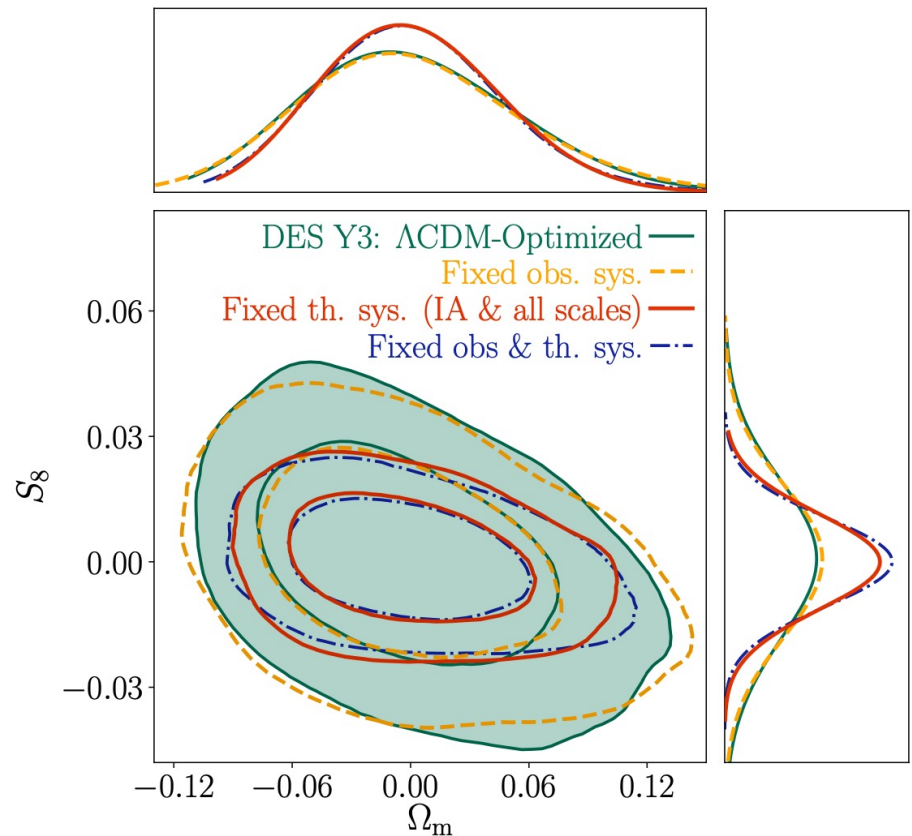
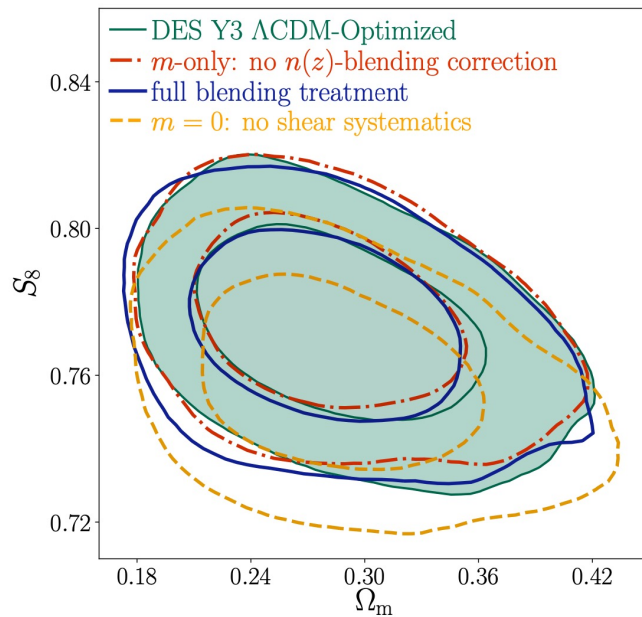




# DES-Y3 analysis

## Cosmological biases

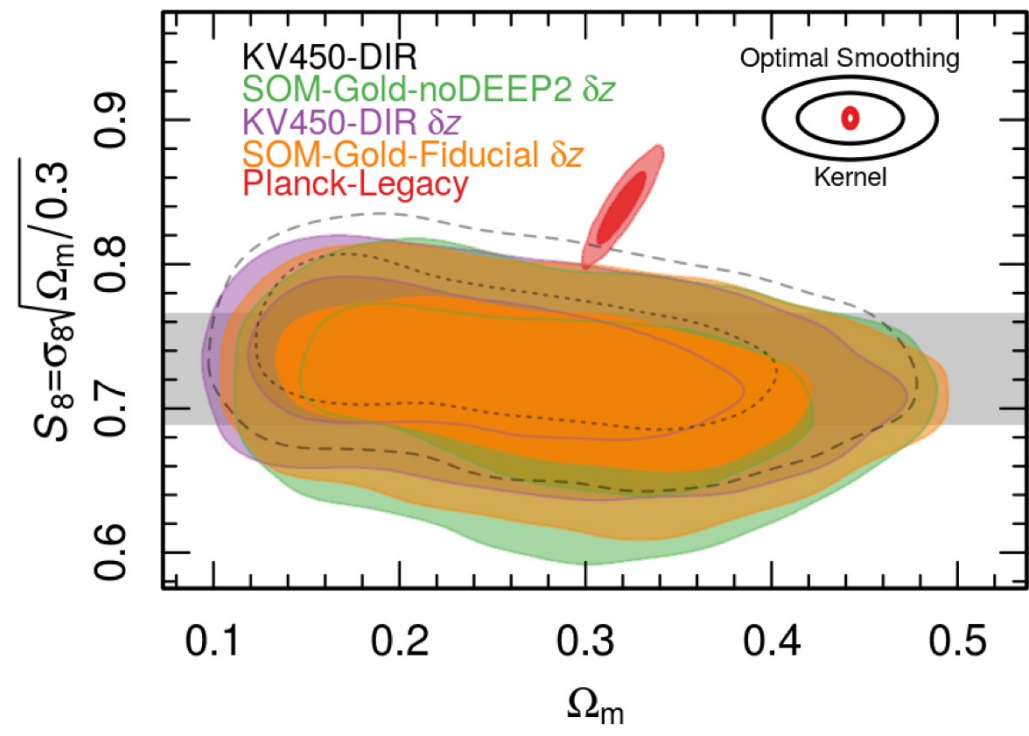
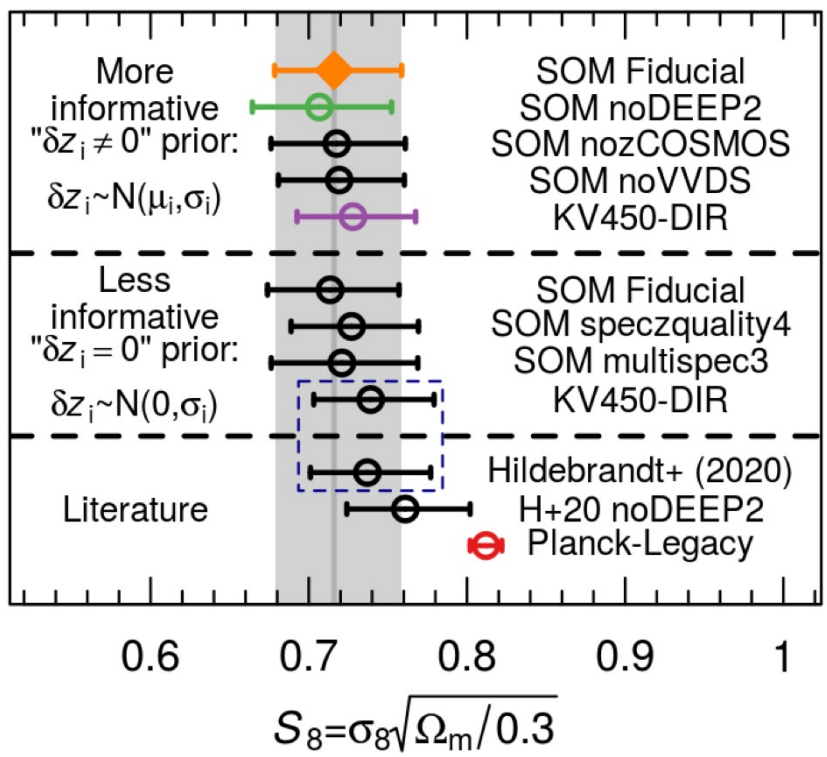
Amon et al. 2021



# KiDS analysis

## Redshift calibration and cosmological biases

Wright et al. 2021



Spectroscopic calibration datasets: CDFS, zCOSMOS, DEEP2, G15Deep, VVDS

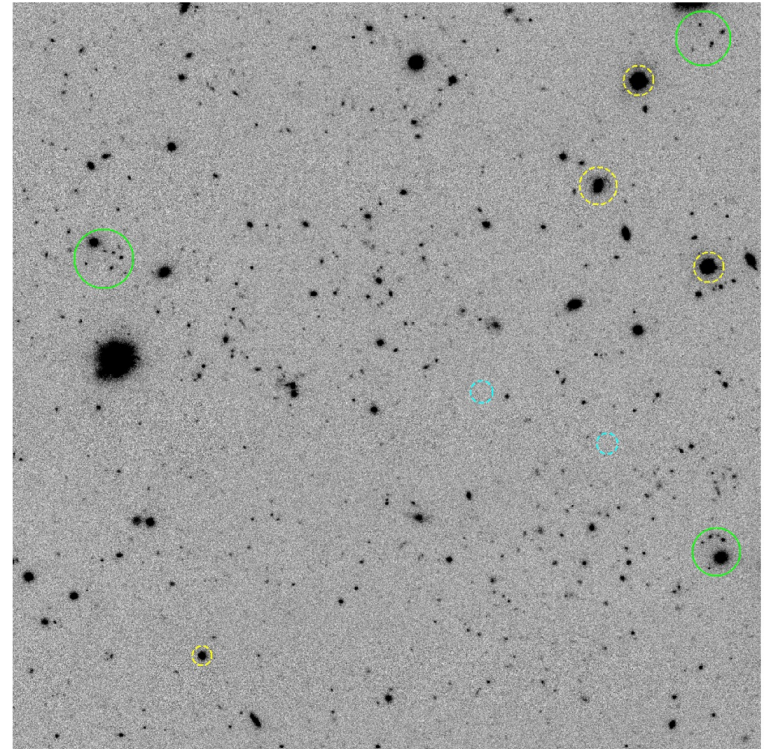
Photometric noise and spectroscopic selection effects contribute equally to the scatter

# KiDS analysis

## Image simulations

Kannawadi et al. 2019

- Realistic simulations of the VST r-band images (HST ACS input morphology)
- Observation depth variation
- Shear calibration for each tomographic bin
- Photometric redshifts calibration (nine-band photometry per galaxy)



# Modeling issues

## Pushing to small scales

In DES and KIDS analysis, they assume a linear galaxy-bias model, e.g.

$$\gamma_t^{ij}(\theta) = b^i(1 + m^j) \int \frac{dl l}{2\pi} J_2(l\theta) \int d\chi n_1^i(z(\chi)) \times \frac{q_s^j(\chi)}{H(z)\chi^2} P_{\text{NL}}\left(\frac{l+1/2}{\chi}, z(\chi)\right), \quad ($$

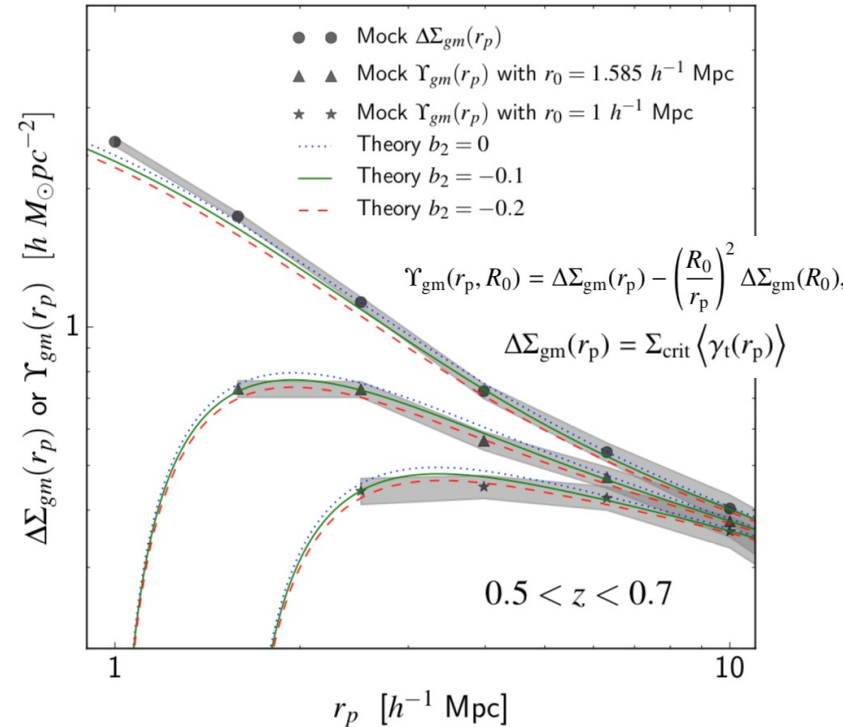
⇒ They put the complexity in  $P_{\text{NL}}$  with e.g. emulators

There are alternative models, but with more free parameters

$$\delta_g(\mathbf{x}) = b_1 \delta(\mathbf{x}) + \frac{1}{2} b_2 [\delta^2(\mathbf{x}) - \sigma^2] + \frac{1}{2} b_{s^2} [s^2(\mathbf{x}) - \langle s^2 \rangle] + O(s^3(\mathbf{x})), \quad \text{McDonald \& Roy 2009}$$

⇒ Being implemented for Euclid now as well

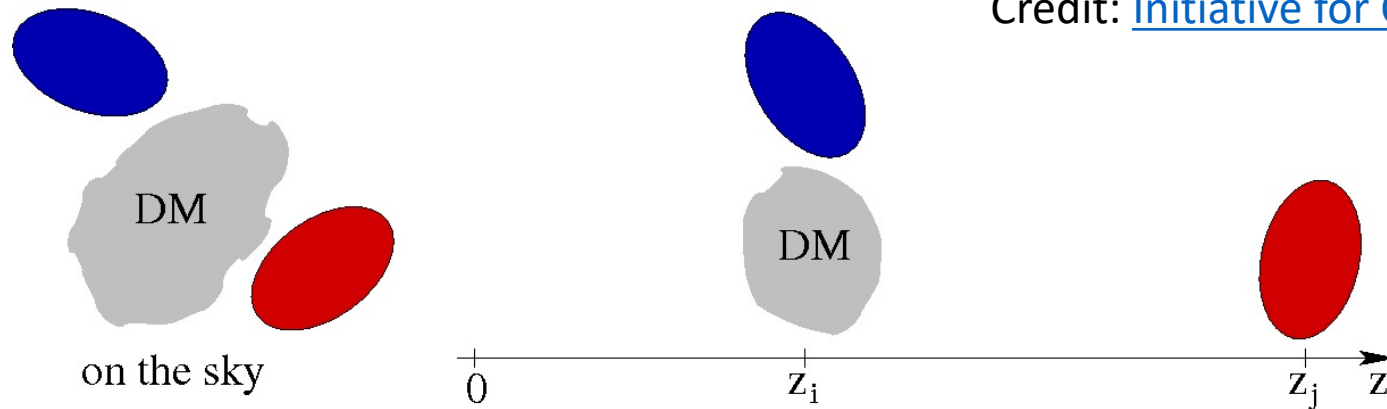
de la Torre et al. 2017



# Modeling issues

## Intrinsic alignment

- Intrinsic alignment tells us about galaxy formation (e.g. in filaments, [Hirata et al. 2004](#), [Chisari et al. 2016](#)). It is a contaminant in cosmic shear analysis. It is quite negligible in GGL analysis (Amon et al. 2022)



### Two types of contribution

- Fake correlation between galaxies infalling in the same halo : **II signal**  
=> More important (1–10%) when  $z_i \sim z_j$ .
- Fake correlation between infalling galaxies and background galaxies: **GI signal**  
=> More important ( $\sim 5\%$ ) when  $z_j \gg z_i$

# Modeling issues

## Boost factor

Mitigation solution proposed in Mandelbaum et al. 2013, Simet et al. 2016

- Account for intrinsic alignment (IA) and increase of sources density in high-density regions compared to a random distribution of lenses

$$B(R) = \frac{\sum_{ls} w_l w_s}{\sum_{rs} w_r w_s}$$

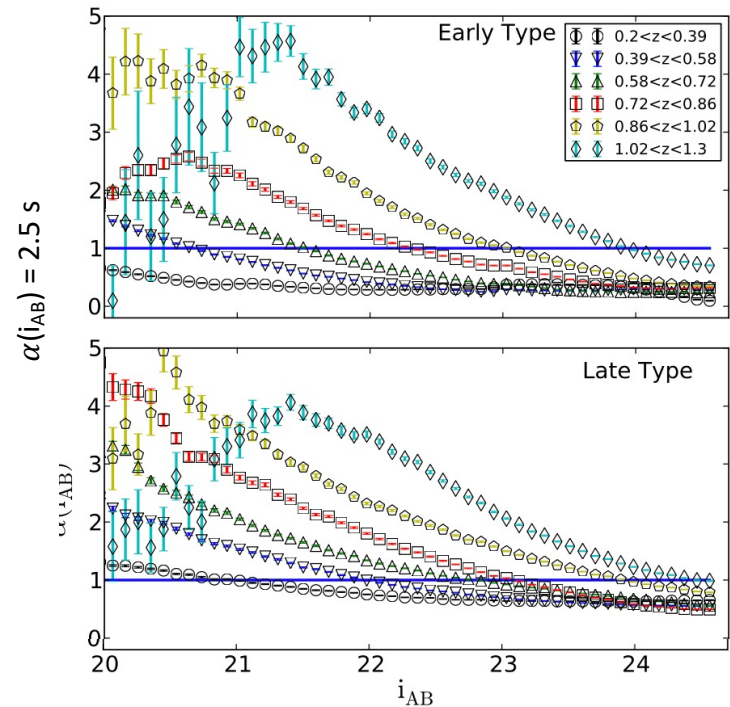
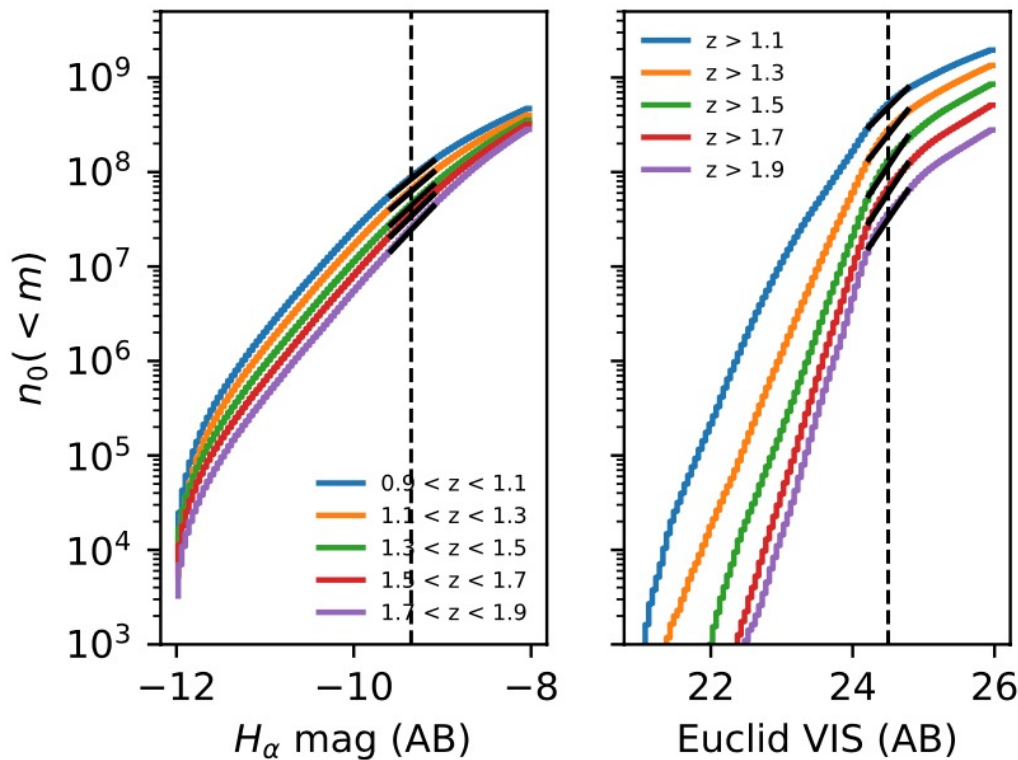
$$\Delta\Sigma(R) = B(R)\Delta\Sigma_l(R) - \Delta\Sigma_r(R)$$

- Not to confuse with magnification bias (all scales effect)

# Magnification effect description

Duncan et al. 2013

Euclid Flagship counts (Jullo et al. in prep)

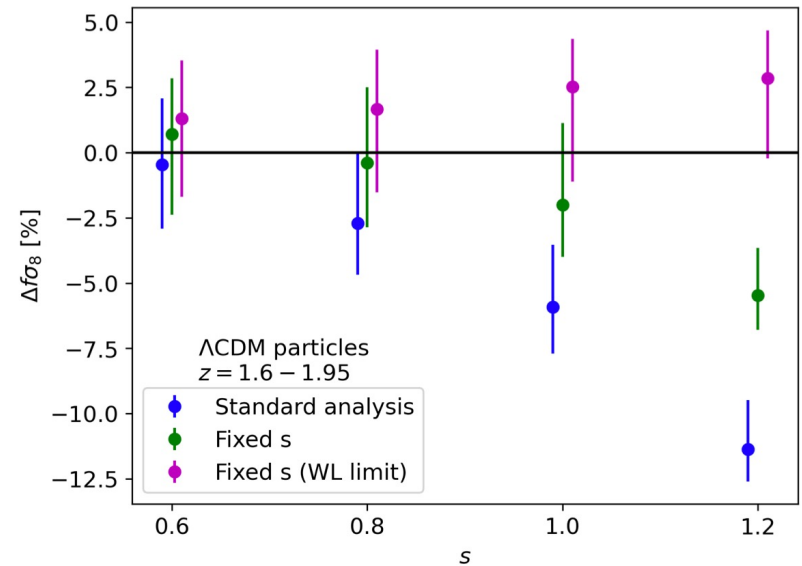
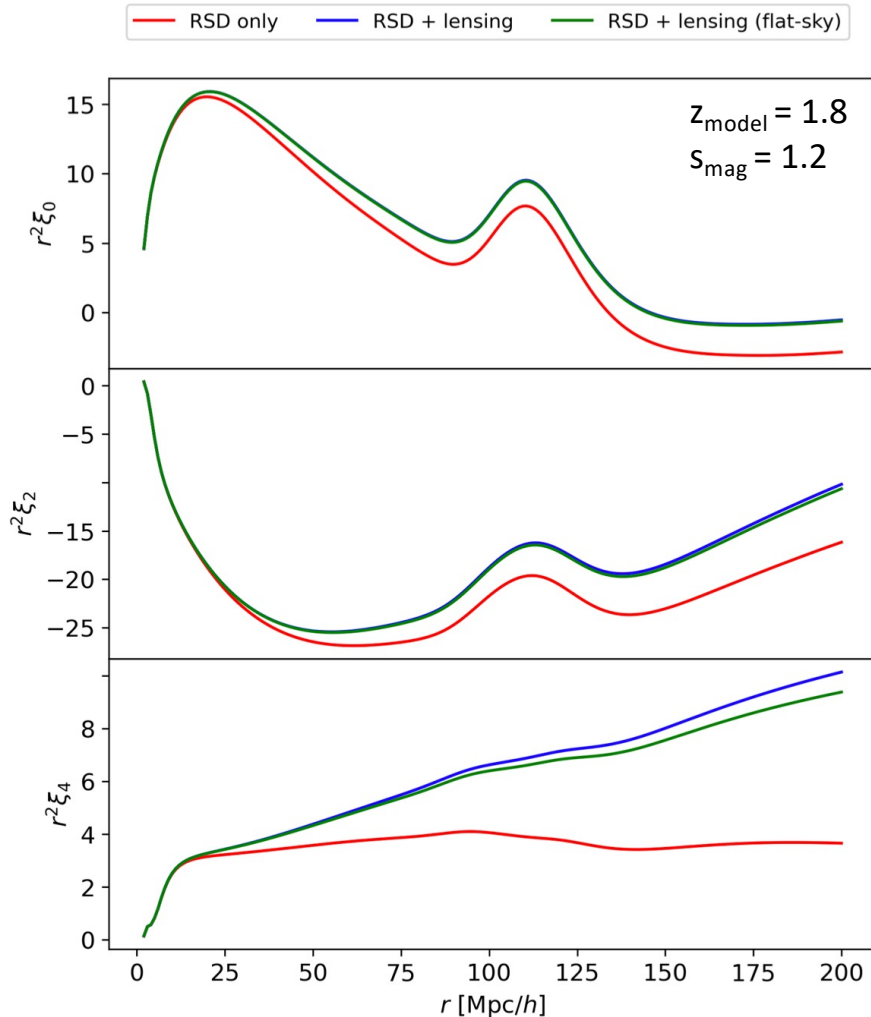


$$s = \left. \frac{d \log_{10} N(< \text{mag})}{d \text{mag}} \right|_{\text{mag}=\text{maglim}}$$

If  $s = 0.4 \Rightarrow$  no magnification bias, because lensed area compensated by number counts

# Magnification effect in Euclid RSD analysis

Breton et al. 2022



=> Important impact in Euclid RSD analysis  
 => For DESI, how much is it for QSO and Ly $\alpha$  forest?



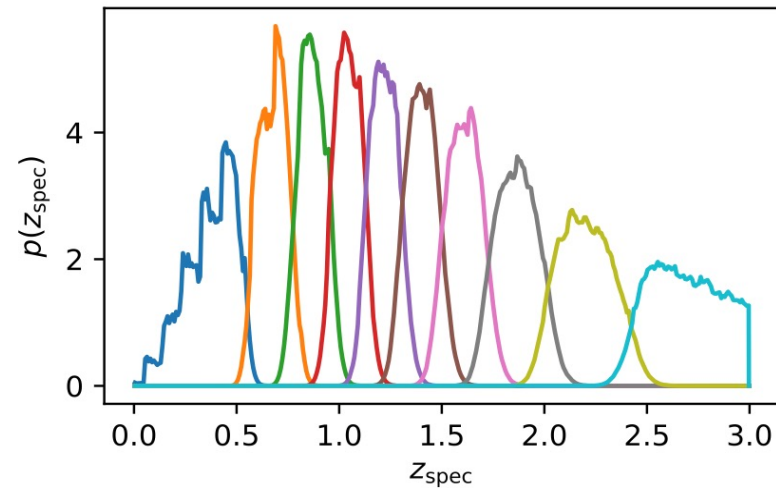
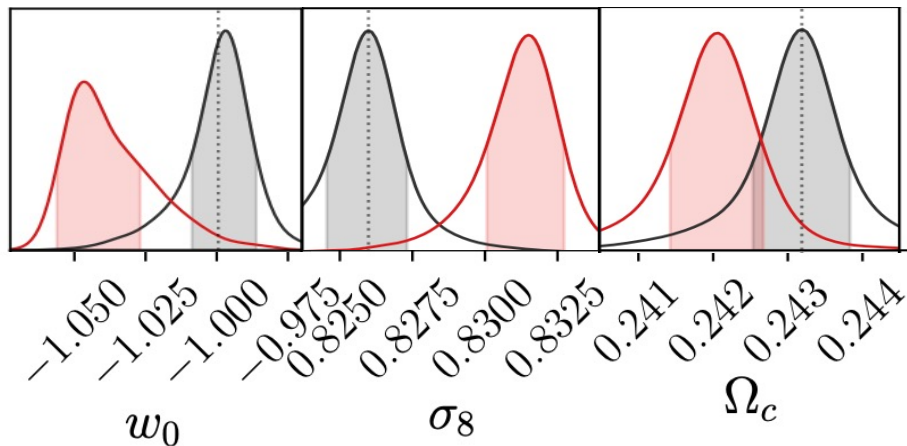
# Magnification effect in Euclid

## 3x2pt analysis

Duncan et al. 2022

### Analysis details

- Estimators used  $w(\boldsymbol{\theta})$ ,  $\gamma t(\boldsymbol{\theta})$  and  $\xi_{\pm}(\boldsymbol{\theta})$
- Euclid like density of sources at mag=24
- Magnification bias  $s = 0.52$  (i.e.  $\alpha=1.3$ , Deshpande et al. 2020)

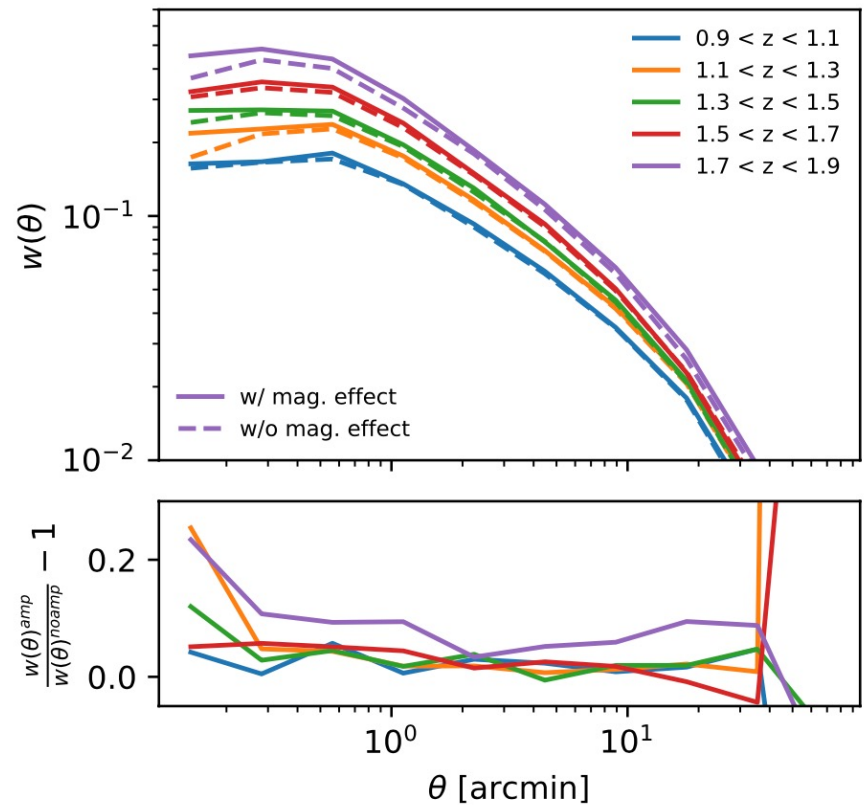
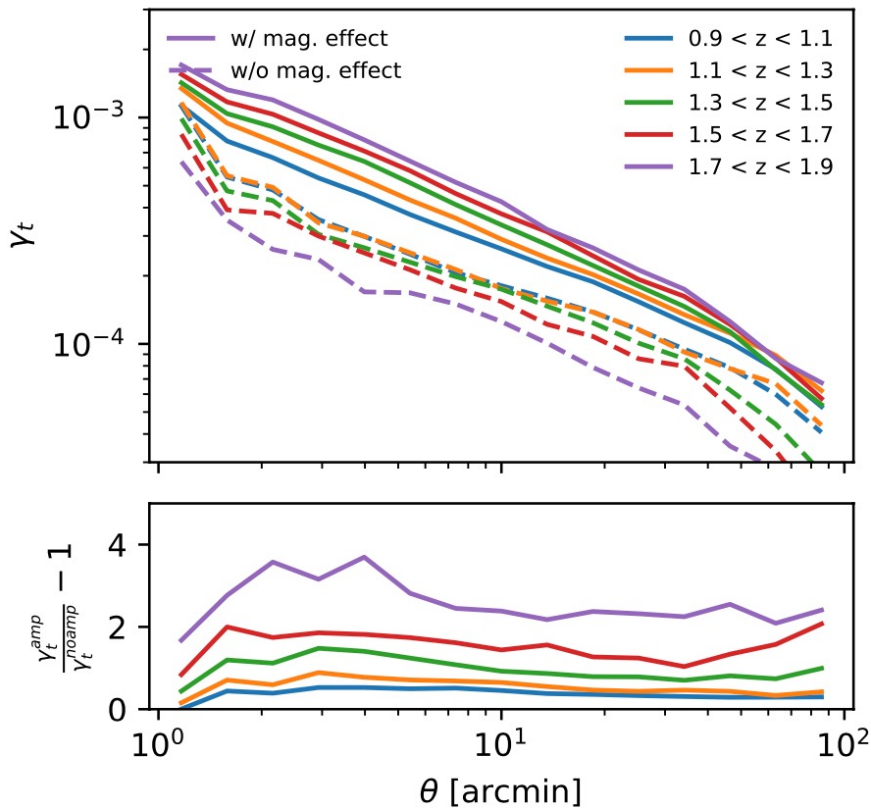


Based on SLICS simulations

=> Significant impact of lensing magnification bias on cosmological parameters in Euclid

# Magnification effect in Euclid GGL with the spectroscopic sample

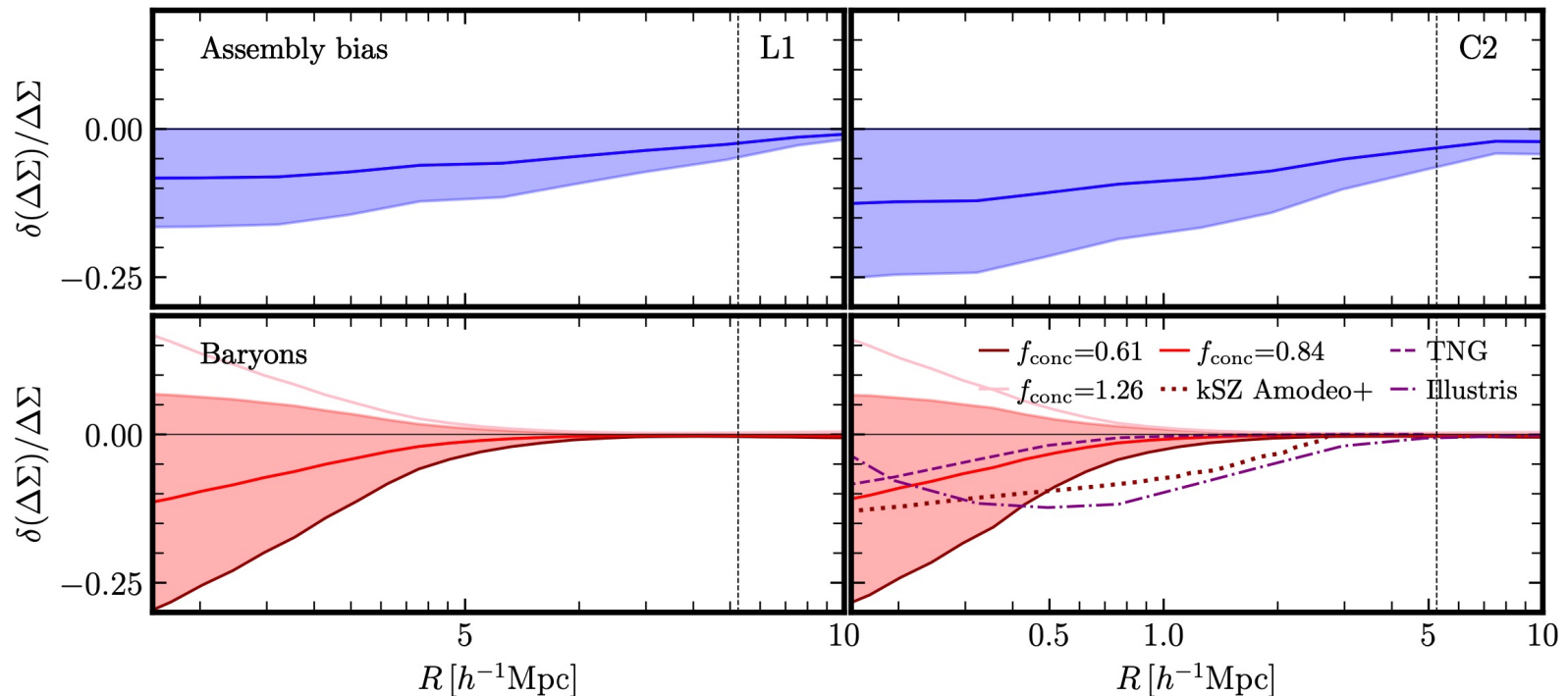
Jullo et al. in prep



# Galaxy Halo Connection

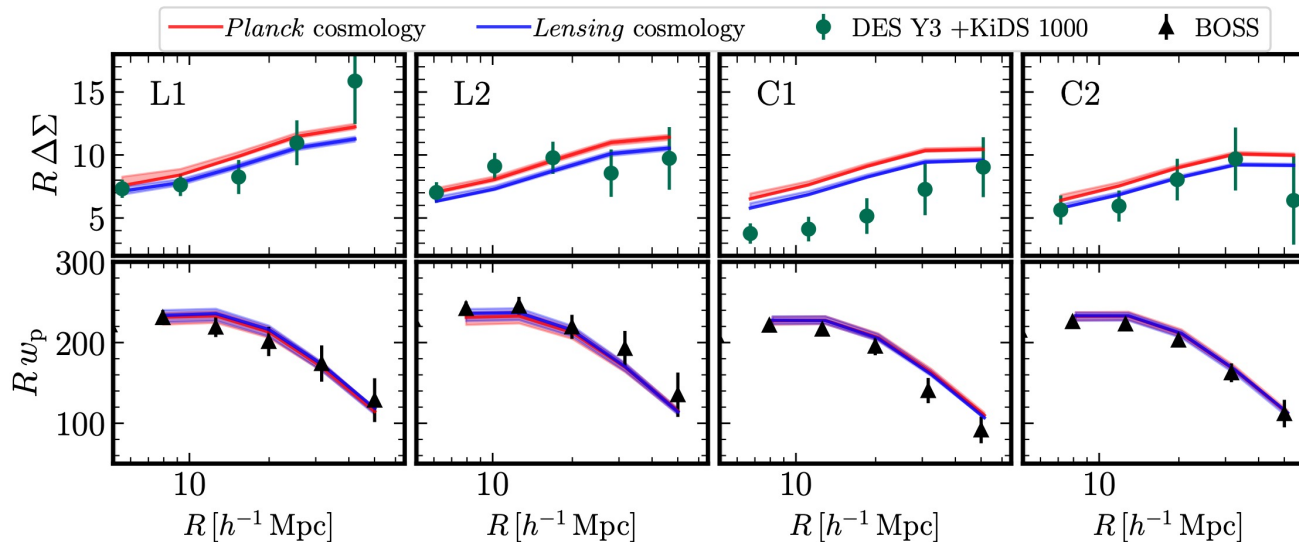
## Assembly bias

*Consistent lensing and clustering in a low- $S_8$  Universe with BOSS, DES Year 3, HSC Year 1 and KiDS-1000, Amon et al. 2022*



# Galaxy Halo Connection

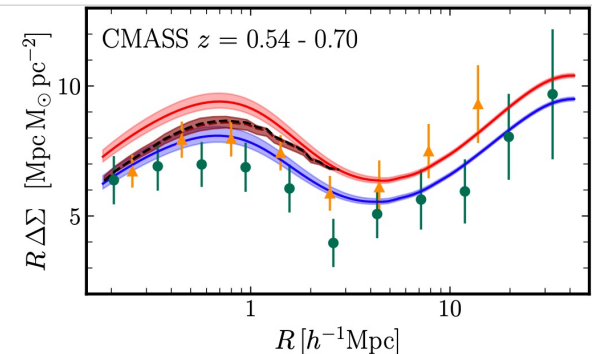
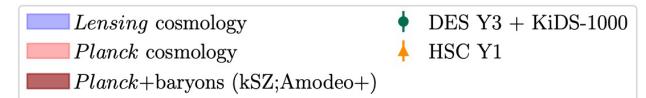
Amon et al. 2022



## Error budget:

- Assembly bias  $\sim 15\%$ , but negligible at  $R > 5.25 h \text{ Mpc}^{-1}$
- Baryons bias  $\sim 10\%$ , but negligible at  $R > 1 h \text{ Mpc}^{-1}$

=> Remaining problem with CMASS C2 sample at about  $5 \text{ Mpc}/h$



# Conclusion

Back in 2006: Dark Energy Task Force (Albrecht et al)

« *If* the systematic errors are at or below the level asserted by the proponents, it is likely to be the most powerful individual Stage-IV technique »

16 years later:

- Is the lensing low-S8 issue real? Current studies require large field coverage => lensing with Euclid

In the future (higher redshift & more precision), there is no other option than introducing lensing in clustering analysis (magnification bias)