

# On the Origin of Galactic Cosmic Rays

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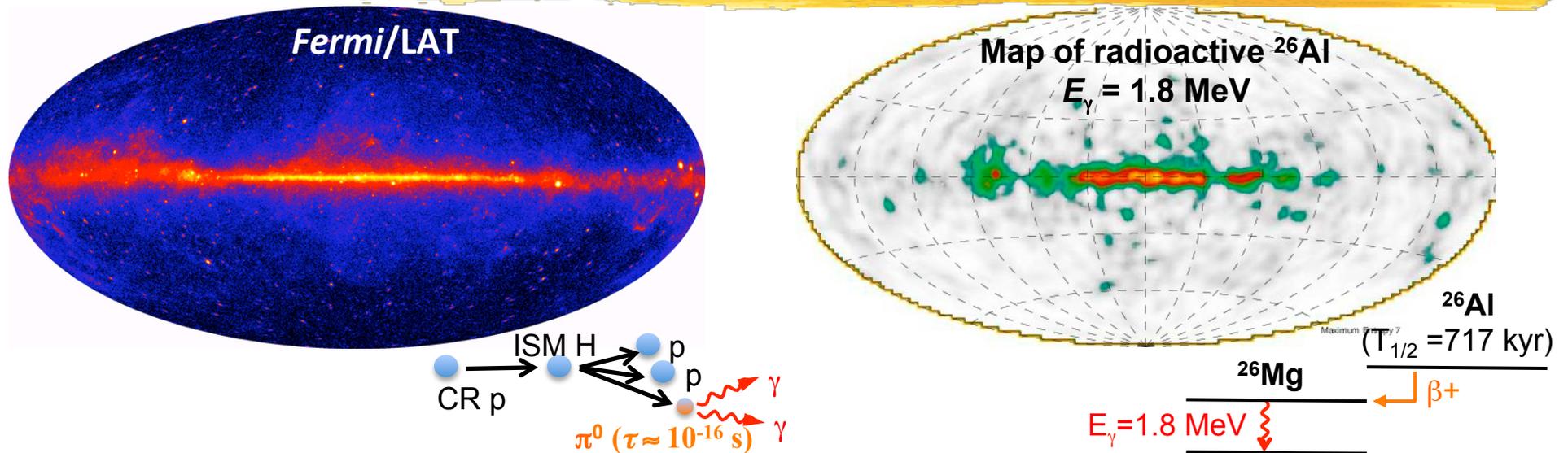
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# Galactic Cosmic Rays & supernova energetics

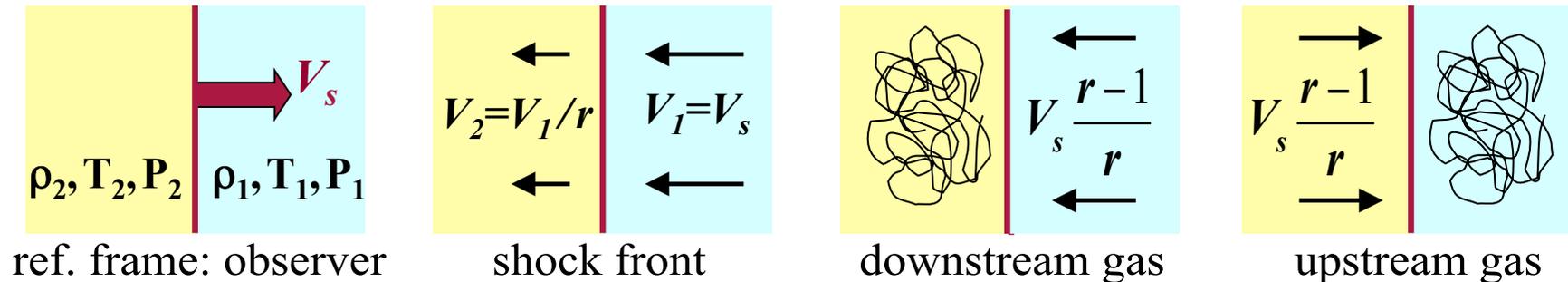
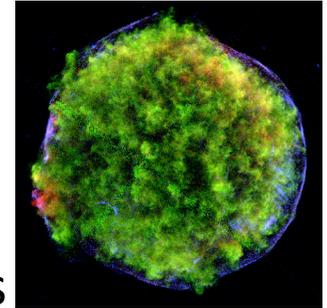
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- GCRs are thought to be produced in **supernova remnants** (Baade & Zwicky 1934)
- Consistent with the cosmic-ray power and supernova energetics:
  - ✓ Kinetic power of CRs injected in the Galaxy:  $L_{\text{CR}} = L_\gamma / R_\gamma \sim 10^{41} \text{ erg/s}$ , where  $R_\gamma \sim 0.004$  is the  $\gamma$ -ray radiation yield (= efficiency) for  $p + p \rightarrow \pi^0 + X$  and  $L_\gamma$  from  $\pi^0$  decay  $\sim 5 \times 10^{38} \text{ erg/s}$  (Fermi/LAT; see Strong et al. 2010)
  - ✓ Kinetic power supplied by supernovae:  $L_{\text{SN}} = E_{\text{SN}} \times f_{\text{SN}} \sim 6 \times 10^{41} \text{ erg/s}$ , where  $E_{\text{SN}} \sim 1.5 \times 10^{51} \text{ erg}$  is the mean energy of a SN and  $f_{\text{SN}} \sim 1.3 \text{ century}^{-1}$  is the SN rate in the Milky Way (from the observed mass of  $^{26}\text{Al}$ ; Diehl+ 2021)

# Diffusive shock acceleration in SN shocks

- High-velocity ejecta in supernova explosion:  $\sim 10\,000$  km/s
- **Strong shock**, with initial sonic Mach number  $M_S = V_s / c_s > 100$  with the sound speed  $c_s \approx 100 (T/10^6 \text{ K})^{0.5}$  km/s
- First-order Fermi (1949) acceleration process in SN shock waves (Krymskii 1977; Bell 1978; Axford et al. 1978; Blandford & Ostriker 1978)
- Particle diffusion on **magnetized turbulence** on both sides of the SN shock



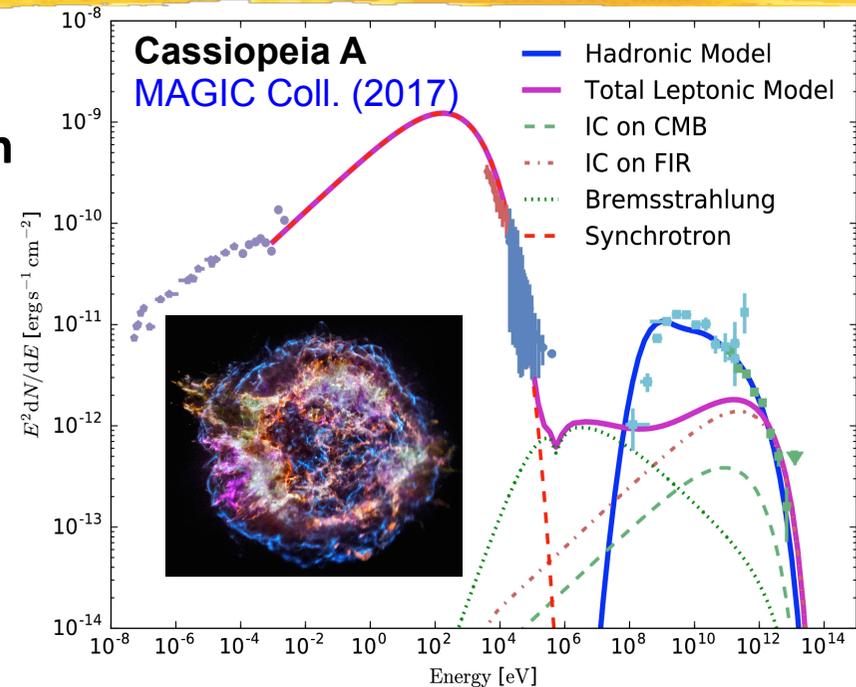
- Fractional momentum gain after each cycle up-down-up:  $\frac{\Delta p}{p} = \frac{4}{3} \frac{r-1}{r} \frac{V_s}{v}$
- Particle momentum spectrum:  $\frac{dN}{dp(p)} \propto p^{-q}$  with  $q = (r+2)/(r-1)$
- For a test-particle strong shock  $r = 4 \Rightarrow q = 2$ ; but observations suggest  $q = 2.1 \dots 2.4$  (?) (e.g. Gabici et al. 2019)

# Maximum CR energy & $\gamma$ -ray observations 4

- Maximum CR energy from the **rate of energy gain by diffusive shock acceleration** and the **finite age of SNR shock** (Lagage & Cesarsky 1983):

$$E_{\max} < 30 Z B_{\mu\text{G}} \text{ TeV}$$

- $\Rightarrow E_{\max}$  can reach the “knee” (3 PeV) if the B-field in the acceleration region is **amplified to  $B \sim 100 \mu\text{G}$**

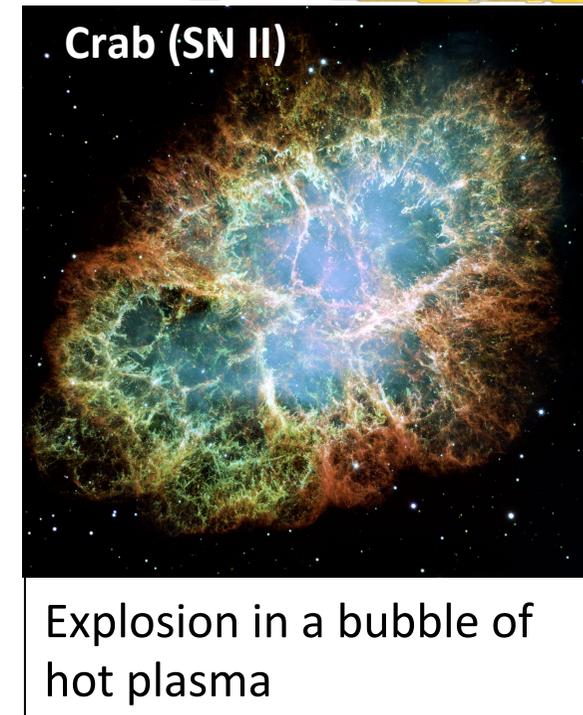
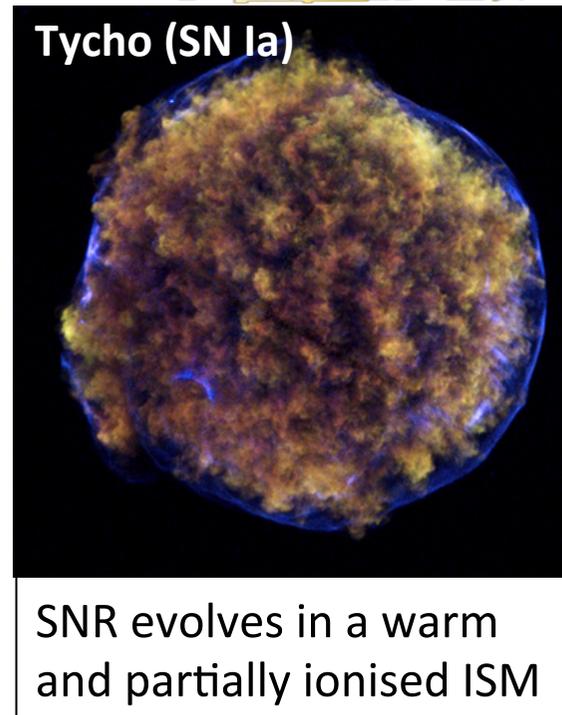


- Gamma-ray observations of young SN remnants** confirm that cosmic rays are produced in SN shocks and constrain the proton **maximum energy** (J. Devin’s talk)
- Predicted maximum energy of protons (Schure & Bell 2013) vs. observations:

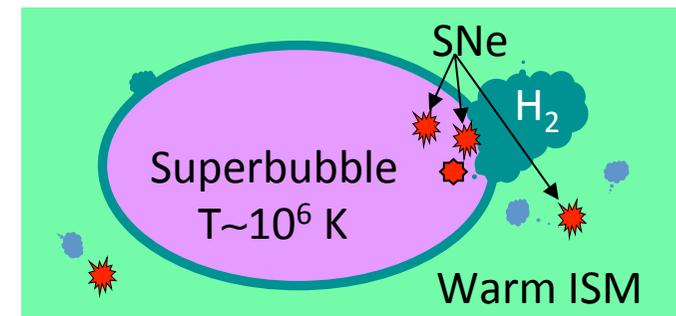
SNR Type	Age	$R_s$	$u_s$	$E_{\max}(\gamma\tau = 5)$	$B_{\text{sat}}$	$B_{\text{obs}}$	$E_{\max}(\text{obs})$
RSG (Cas A)	330 yr	2.2 pc	4900 km s <sup>-1</sup>	283 TeV	243 $\mu\text{G}$	210–230 $\mu\text{G}$	~12 TeV (MAGIC 2017)
Tycho	440 yr	3.2 pc	3900 km s <sup>-1</sup>	108 TeV	128 $\mu\text{G}$	200–230 $\mu\text{G}$	~10 TeV (VERITAS 2017)
SN1006	1000 yr	7.6 pc	4100 km s <sup>-1</sup>	<60 TeV	< 35 $\mu\text{G}$	80–150 $\mu\text{G}$	~50 TeV (Condon et al. 2016)



# SN distribution in the ISM phases



- Massive stars are born in **OB association** and their wind activities generate **superbubbles** (SBs) of hot plasma, where most core-collapse SNe explode ( $\sim 80\%$ ; [Lingenfelter & Higdon 2007](#))
- With 25% of Galactic SNe of **Type Ia** occurring in the **warm ISM**: **60% of SNe in hot SBs**, **40% in warm ISM** (28% in WNM, 12% in WIM)

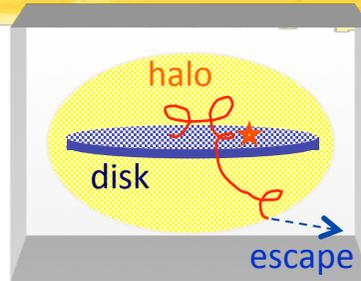


# Galactic cosmic-ray composition

Refs: Meyer, Drury & Ellison (1997); Ellison, Drury & Meyer (1997)

1. **Overabundance of elements with  $Z > 2$**  relative to H and He (as compared with the solar system composition)  
⇒ **Not necessarily**, because CR protons and  $\alpha$ -particles have different source spectra than the other elements (e.g. Tatischeff & Gabici 2018)
2. **Overabundance of refractory elements** over volatiles due to the more **efficient acceleration** of material locked in **dust grains**  
⇒ **OK, but which dust grains? From which ISM phase(s) are they accelerated?**
3. **Overabundance of the heavier volatile elements** compared to the lighter ones due to a **dependence of the acceleration efficiency on ion rigidity**  
⇒ **Expected from nonlinear DSA (Ellison+ 1981) and PIC simulations (Caprioli+ 2017), but ionisation states in shock precursors? Depends on the ISM phases**
4. **Overabundance of  $^{22}\text{Ne}$**  due to the acceleration of **Wolf-Rayet wind** material enriched in He-burning products  
⇒ **OK, but how exactly Wolf-Rayet wind material is incorporated in GCRs?**

# Protons, $\alpha$ -particles and O source spectra



- Fit to **Voyager 1** and **AMS-02** data using a **1D advection-diffusion model** with homogeneous diffusion for the GCR propagation (Evoli et al. 2019)
- **Broken power law source spectra** from a fit of propagated spectra to the data

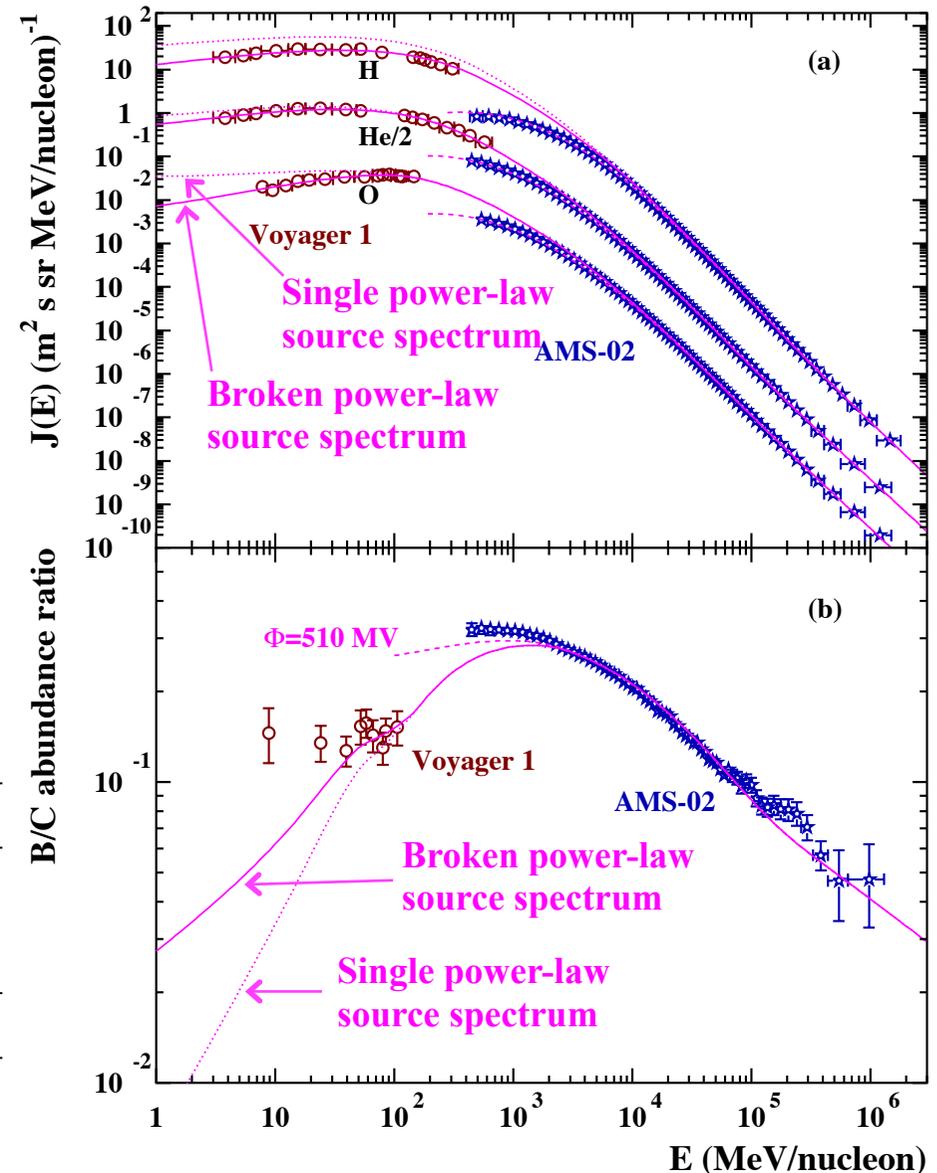
**Table 2.** CR source spectrum parameters (Eq. 2).

Parameter	H	He	O
$E_{\text{break}}$	$10 \pm 2$ GeV/n	$200^{+160}_{-120}$ MeV/n	$160^{+40}_{-30}$ MeV/n
$\gamma_{\text{l.e.}}$	$4.10 \pm 0.03$	$3.98^{+0.08}_{-0.20}$	$3.32^{+0.18}_{-0.24}$
$\gamma_{\text{h.e.}}^a$	4.31	4.21	4.26
$\chi^2_{\text{min}}^b$	16.0 for 13 d.o.f. <sup>c</sup>	7.3 for 14 d.o.f.	5.9 for 12 d.o.f.

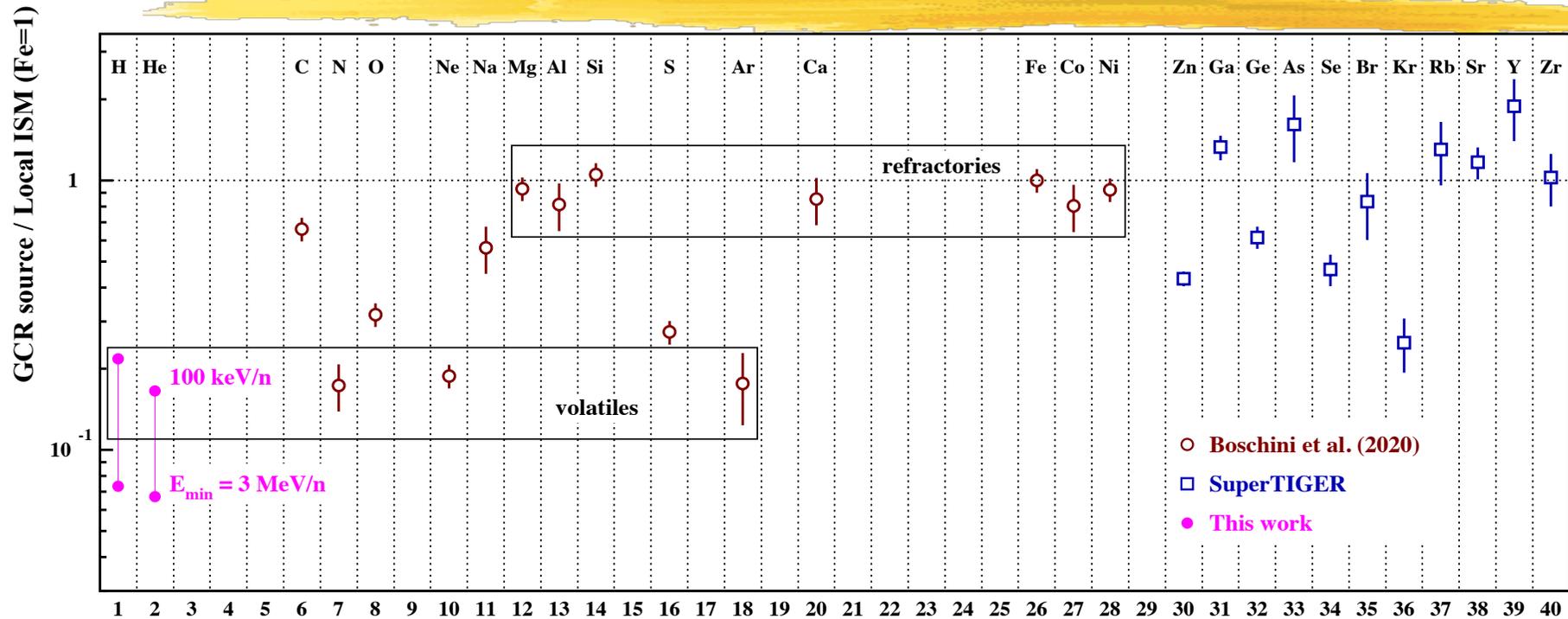
<sup>a</sup> Parameter fixed from Evoli et al. (2019).

<sup>b</sup> Minimum  $\chi^2$  from a fit of the propagated spectrum to Voyager 1 data.

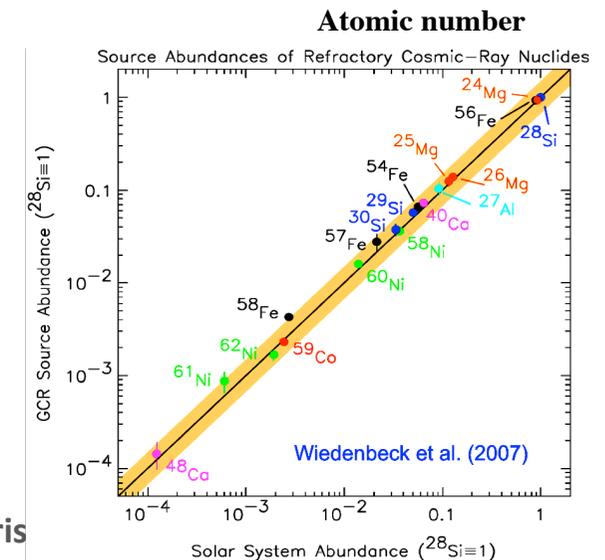
<sup>c</sup> d.o.f.: degrees of freedom.



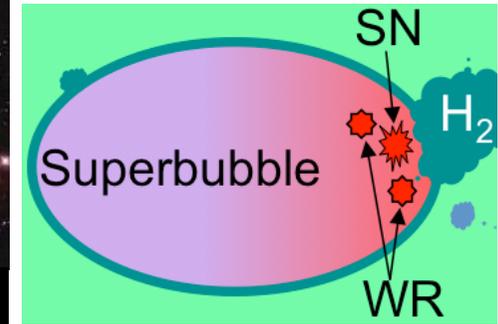
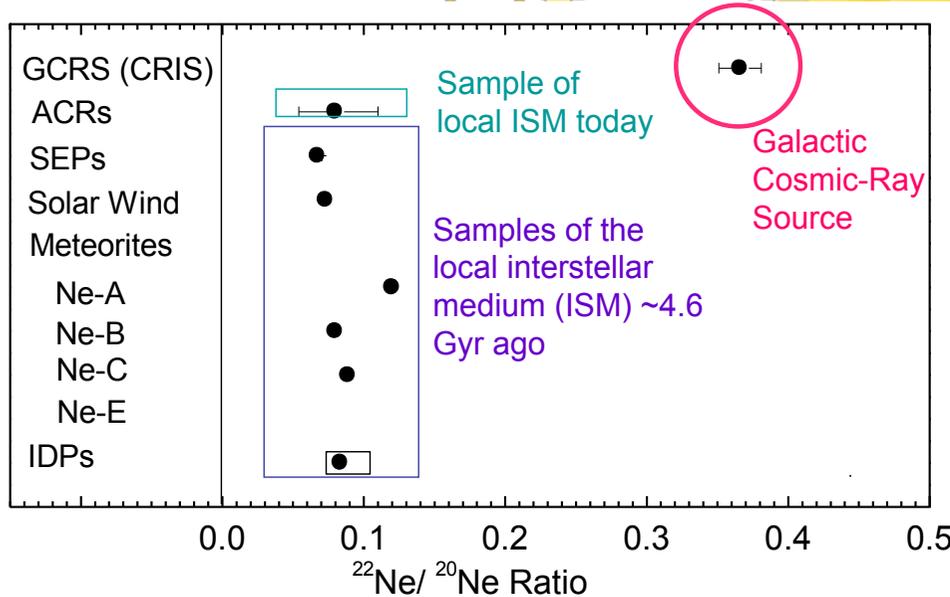
# GCR abundance data



- Integration of source spectra  $\Rightarrow$  p &  $\alpha$  abundances similar to those of the other volatiles N, Ne and Ar, provided that the **minimum CR source energy is of the order of a few hundred keV/n**
- Highly refractory elements Mg, Al, Si, Ca, Fe, Co, and Ni in solar system proportions  $\Rightarrow$  acceleration of various **dust grains of the ISM mix**

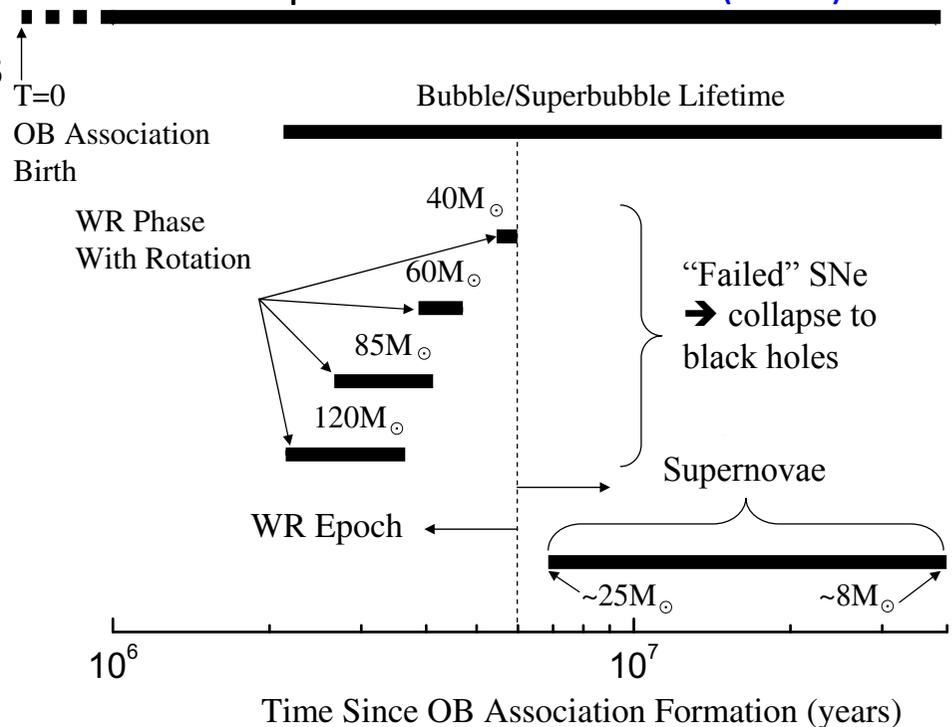


# $^{22}\text{Ne}$ abundance in GCRs



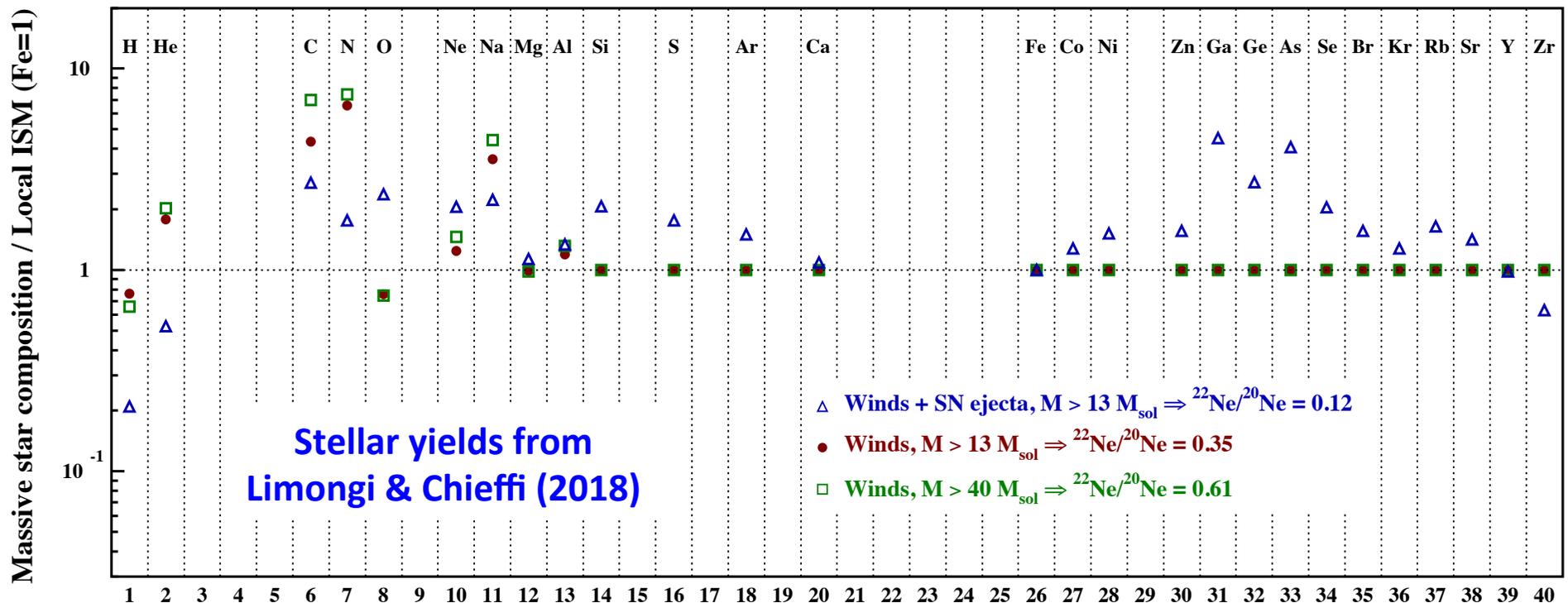
Adapted from Binns et al. (2008)

- GCR  $^{22}\text{Ne}/^{20}\text{Ne}$  ratio  $\approx 0.35$ , i.e. **~ 5 times the solar ratio** (Garcia-Munoz et al. 1970; Binns et al. 2005)
- Contribution to GCRs of **Wolf-Rayet wind material** ( $^{14}\text{N}(\alpha,\gamma)^{18}\text{F}(\beta^+)^{18}\text{O}(\alpha,\gamma)^{22}\text{Ne}$  during He burning)? (Cassé & Paul 1982)
- GCR origin in **superbubbles** enriched in  $^{22}\text{Ne}$  from winds of massive stars?



# GCR $^{22}\text{Ne}$ from enriched superbubble gas<sup>10</sup>

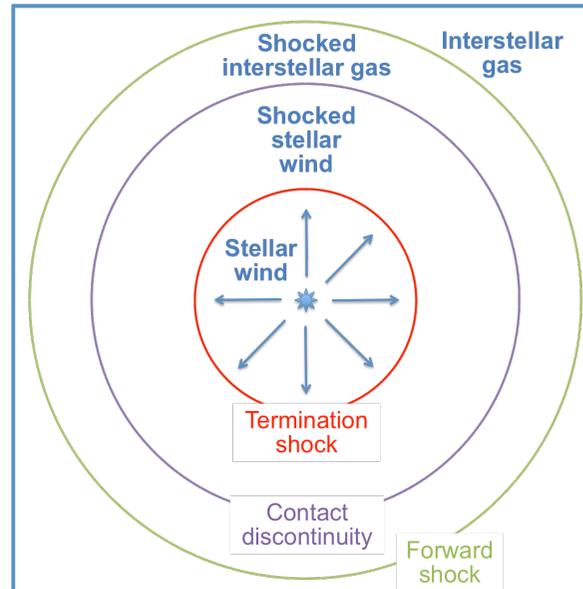
- Mix of massive star winds and SN ejecta in SB cores? **No,  $^{22}\text{Ne}/^{20}\text{Ne}=0.12$**   
Ratio close to solar, because massive star winds and SN ejecta are the principal sources of both  $^{20}\text{Ne}$  and  $^{22}\text{Ne}$  in the Universe (Prantzos 2012)
- **Only massive star winds in SB cores? No,  $^{22}\text{Ne}/^{20}\text{Ne}=0.35$**
- Winds from **very massive stars  $\geq 40 M_{\text{sol}}$ ?  $^{22}\text{Ne}/^{20}\text{Ne}=0.61$ , maybe...**



# GCR $^{22}\text{Ne}$ from wind termination shocks

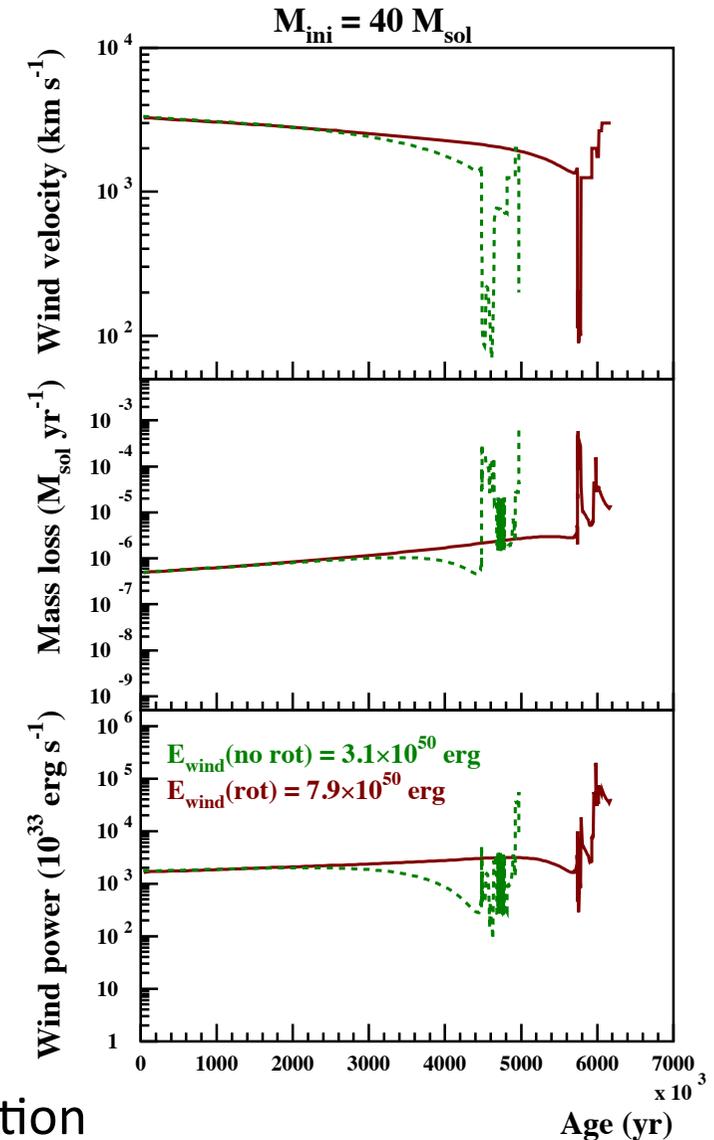
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- [Gupta et al. \(2020\)](#): WTSs can contribute **more than 25% of the CR production** in massive star clusters
- ⇒  $^{22}\text{Ne}$ -rich CR component (see also [Kalyashova et al. 2019](#))



- Time-dependent yields, mass loss rates & stellar types from the **Geneva Observatory database** (e.g. [Ekström et al. 2012](#))
- Instantaneous acceleration efficiency in WTS assumed to be proportional to the **wind mechanical power**

⇒  $^{22}\text{Ne}/^{20}\text{Ne}=1.56$  in the accelerated wind composition



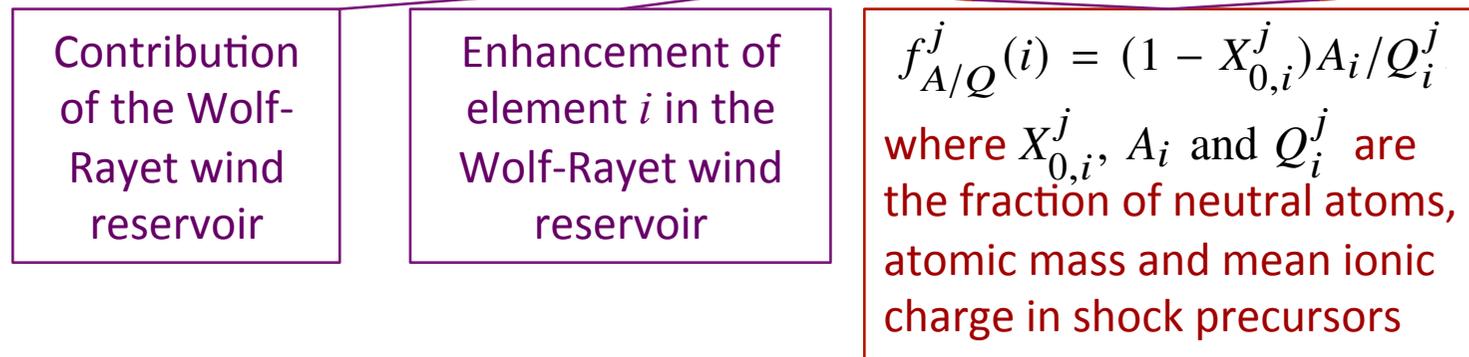
# GCR composition model

- Measured GCR source abundances:  $C_{\text{mes}}(i) = C_{\text{gas}}(i) + C_{\text{dust}}(i)$

- Dust contribution:  $C_{\text{dust}}(i) = \text{SC}(i) f_d(i) \epsilon_{\text{dust}}$



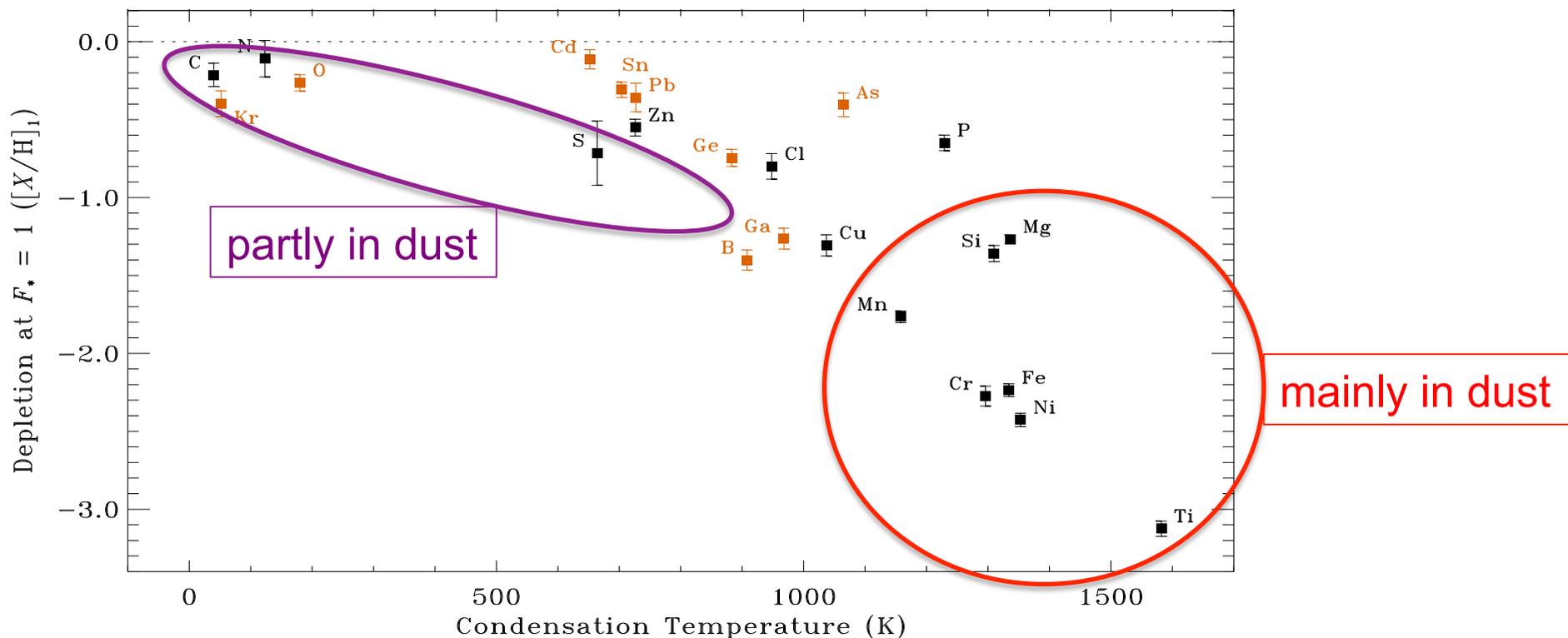
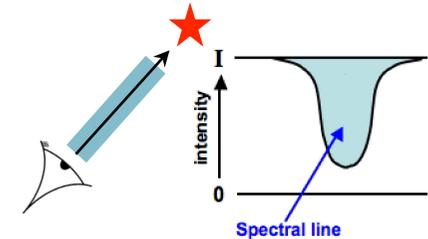
- Gas contribution:  $C_{\text{gas}}(i) = \text{SC}(i) (1 - f_d(i)) \epsilon_{\text{gas}} [x_w f_w(i) f_{A/Q}^w(i) + (1 - x_w) f_{A/Q}^{\text{SC}}(i)]$



- If the gas reservoir includes several phases of the ISM:  $f_{A/Q}^{\text{SC}}(i) = \sum_k a_k f_{A/Q}^{\text{SC},k}(i)$
- Fitting theoretical abundances to data to derive  $x_w$ ,  $\epsilon = \epsilon_{\text{dust}} / \epsilon_{\text{gas}}$ , as well as constraints on the **GCR source reservoirs** (e.g. their temperature)

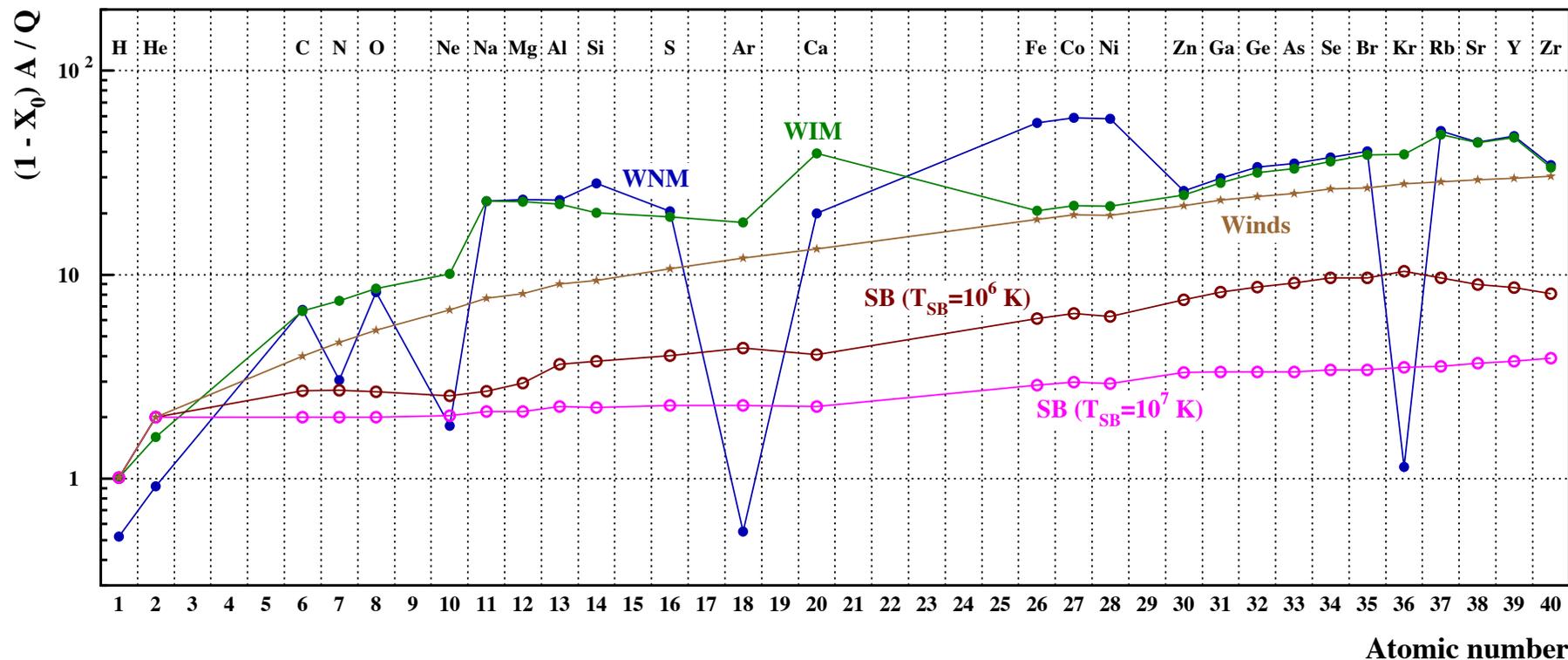
# Interstellar dust composition

- **Average fraction in dust for each element,  $f_d(i)$** , from
  - Gas-phase element **depletions** (Jenkins 2009, 2019; Ritchey et al. 2018)
  - The interstellar dust modeling framework **THEMIS** (Jones et al. 2017)
  - General properties of **primitive interplanetary dust**



# Ionisation states in shock precursors

- **Warm ISM:** Ionisation states of the WIM and the WNM from absorption/emission line measurements (e.g. Sembach et al. 2000; Madsen et al. 2006) + **photoionisation precursors** mainly produced by He I and He II photons from the thin ionisation zone behind the shock (Ghavamian et al. 2000; Medina et al. 2014)
- **Superbubbles:** collisional ionisation in a hot plasma (negligible photoionisation)
- **Stellar winds:** photoionisation by the EUV radiation of hot stars + EUV and X-rays from shocks in the winds => heavy elements mostly triply ionised (e.g. Hillier 2020)



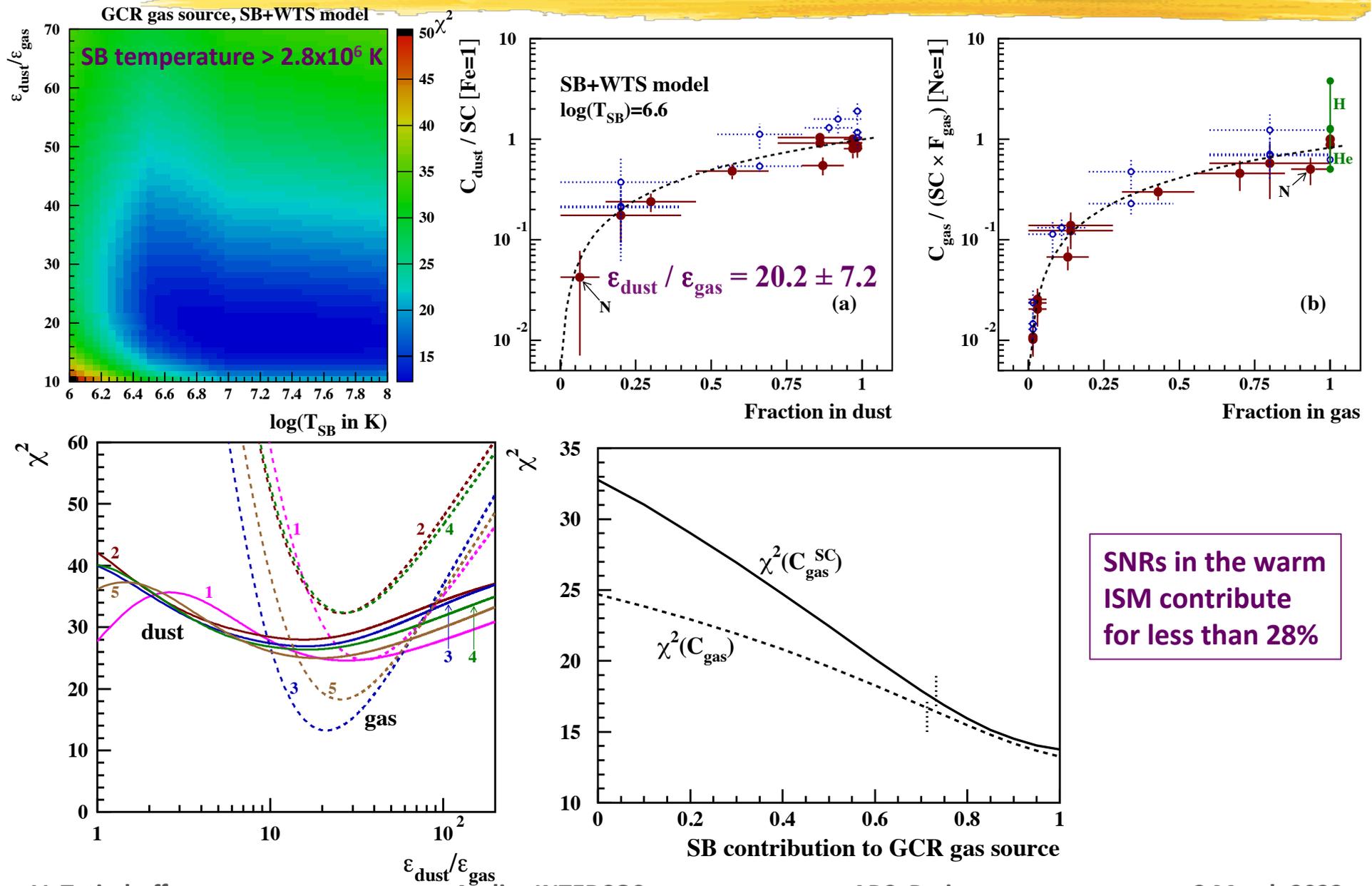
# Results of the GCR composition model

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	Model 1	Model 2	Model 3	Model 4	Model 5
GCR gas source of SC compo.	70% WNM, 30% WIM	SB	SB	60% SB, 28% WNM, 12% WIM	60% SB, 28% WNM, 12% WIM
$^{22}\text{Ne}$ -rich GCR gas source	Accelerated winds	Winds in SB	Accelerated winds	Winds in SB	Accelerated winds
SB temperature $\log(T_{\text{SB}})^a$	–	$6.50 \pm 0.25$	$> 6.45$	$6.5^{+0.3}_{-0.2}$	$> 6.35$
Relative eff. $\epsilon = \epsilon_{\text{dust}}/\epsilon_{\text{gas}}^b$	$33.8 \pm 13.4$	$26.0 \pm 13.2$	$17.9 \pm 9.7$	$27.0 \pm 13.2$	$22.8 \pm 10.6$
W.-R. wind contribution $x_w^c$	10.3%	48.9%	(5.1 – 6.1)%	$(55.6^{+1.3}_{-0.3})\%$	(7.3 – 7.9)%
$\chi_{\text{min}}^2$ (GCR dust source) <sup>d</sup>	24.6	26.9	25.9	26.0	24.8
$\chi_{\text{min}}^2$ (GCR gas source) <sup>e</sup>	24.7	31.1	12.2	31.4	16.7
SB temperature $\log(T_{\text{SB}})$	–	6.6 (fixed)	6.6 (fixed)	6.6 (fixed)	6.6 (fixed)
Relative eff. $\epsilon = \epsilon_{\text{dust}}/\epsilon_{\text{gas}}^b$	$33.8 \pm 13.4$	$23.2 \pm 9.4$	$20.2 \pm 7.2$	$24.6 \pm 10.2$	$24.4 \pm 9.2$
W.-R. wind contribution $x_w^c$	10.3%	48.9%	5.9%	56.0%	7.7%
$\chi_{\text{min}}^2$ (GCR dust source) <sup>d</sup>	24.6	28.0	26.9	26.4	25.0
$\chi_{\text{min}}^2$ (GCR gas source) <sup>e</sup>	24.7	32.3	13.2	32.4	18.3

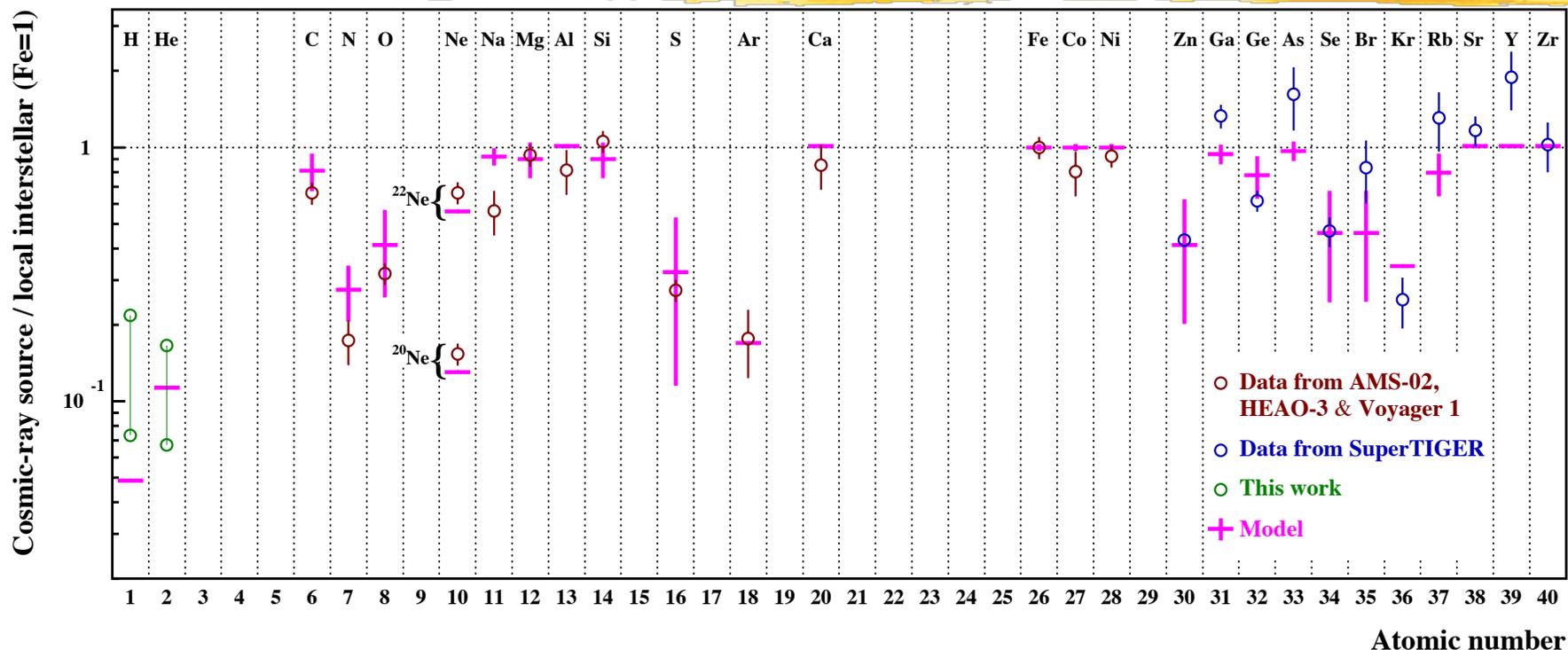
- Five models depending on the relative weights of the **ISM phases** in the GCR volatile production, and the **origin of GCR  $^{22}\text{Ne}$**
- Best-fit model: GCR volatiles accelerated in **superbubbles** +  $^{22}\text{Ne}$ -rich component from acceleration in **wind termination shocks** ( $x_w \approx 6\%$ )

# Results of the GCR composition model

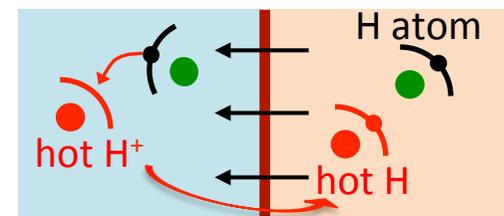


SNRs in the warm ISM contribute for less than 28%

# Origin of GCR volatiles in superbubbles

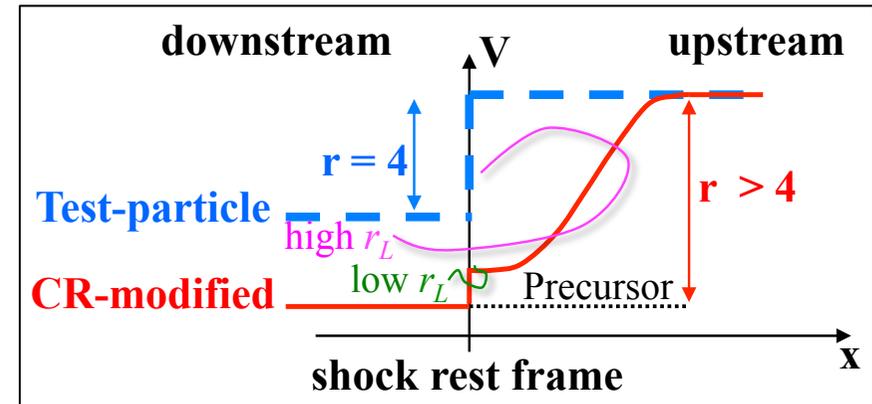
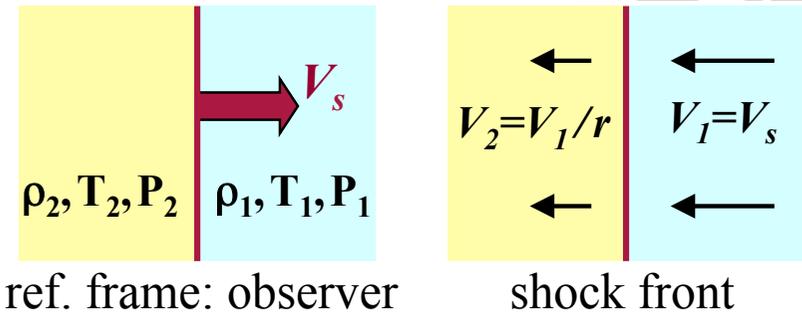


- SNRs in the warm ISM contribute to the GCR volatile composition for less than 28%, whereas ~ 40% of SNe occur in this phase and not in superbubbles (?)
- Effects of neutral atoms on the acceleration process: ion-neutral damping, **neutral return flux (Morlino et al. 2013)?**
  - ⇒ Heating of the upstream plasma
  - ⇒ Reduction of the shock Mach number
  - ⇒ **Reduction of the particle acceleration efficiency**



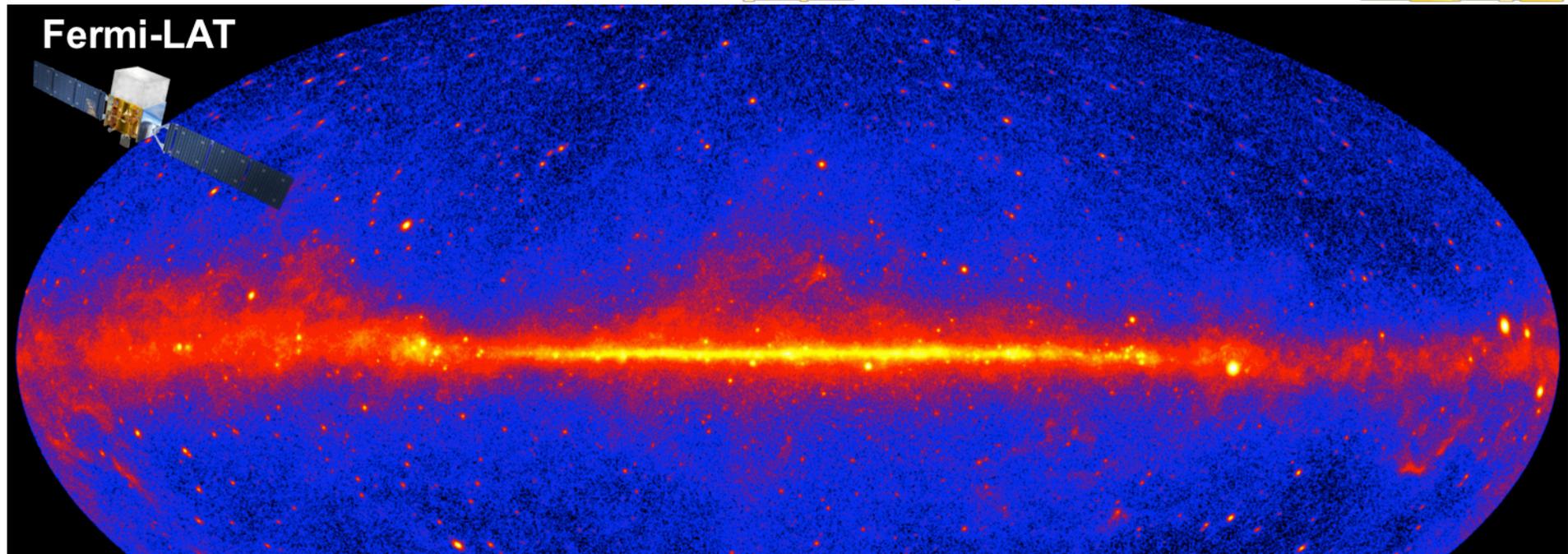
# Acceleration of dust grains

- **Preferential acceleration of dust grains** in SN shocks, because interstellar grains can have **very large**  $A/Q \sim 10^4 - 10^8$  and particles with a high rigidity ( $R \propto A/Q$ ) feel a larger  $\Delta V$  of the background plasma (Ellison et al. 1997, 1998)
  - i. Grain acceleration
  - ii. Grain sputtering with ambient atoms
  - iii. Injection of the sputtered ions in the acceleration process with the **supra-thermal velocity** of the parent grain



- **Can GCR refractories also be produced in superbubbles?**
  - ✓ **Dust is continuously replenished** in the Orion-Eradinus SB through thermal evaporation of embedded molecular clouds (Ochsendorf et al. 2015)
  - ✓ Timescale between two successive SNe in SBs,  $t \lesssim 1$  Myr, **shorter than the destruction timescale** of silicate grains against thermal sputtering (Tielens et al. 1994) - *not true for nano carbonaceous grains and PAH molecules* (to be studied)

# GCR acceleration efficiency



- Efficiency of GCR production from the  $\gamma$ -ray luminosity of the Milky Way and the proton source spectrum:  $N_{GCR}(p) \approx (0.2 - 1.5) \times 10^{45}$  protons  $s^{-1}$
- Estimating the mass of gas swept up by interstellar shocks, we get:
  - Acceleration efficiency of superbubble gas by SN shocks:  $\eta_{SB} \approx (0.4 - 2.3) \times 10^{-5}$
  - Acceleration efficiency of wind material by WTSs:  $\eta_{wind} \approx 0.8 \eta_{SB}$
  - Acceleration efficiency of GCR refractories from dust grains:  $\eta_{dust} \gtrsim 10^{-4}$

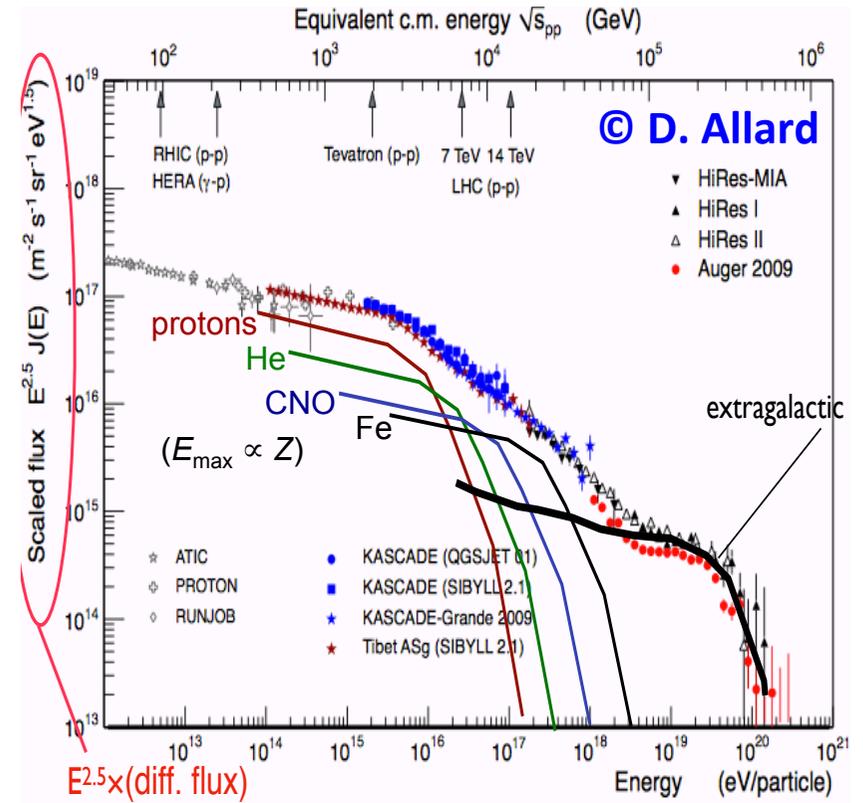
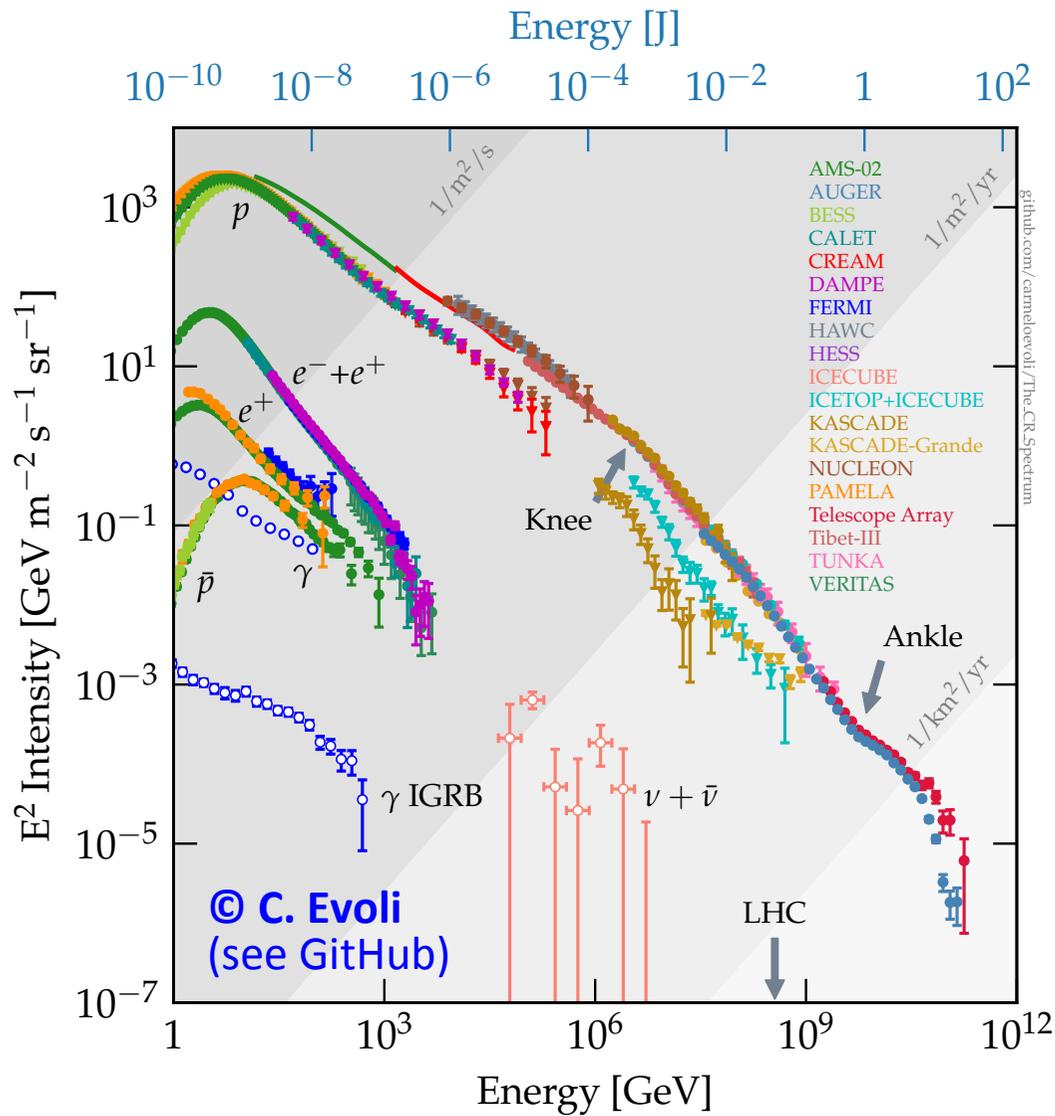
# Conclusions

- Measured source abundances of all primary and mostly primary CRs **from H to Zr are well explained**, including the overabundance of  $^{22}\text{Ne}$
- **No overabundance of elements with  $Z > 2$**  relative to H and He, if the **minimum CR source energy is of the order a few hundred keV nucleon<sup>-1</sup>**
  - ⇒ **Escape of low-energy CR from their sources? Source spectrum differences between p,  $\alpha$ -particles and heavy nuclei?**
- CR volatiles are mostly accelerated in **Galactic superbubbles**, from SN shocks sweeping up a plasma of  $T_{\text{SB}} > 2.8 \text{ MK}$ 
  - ⇒ **CR production in superbubbles up to and above  $3 \times 10^{15} \text{ eV}$  (Vieu et al. 2022)?**
- The overabundance of  $^{22}\text{Ne}$  is most likely due to a small ( $x_w \approx 6\%$ ) contribution of particle acceleration in **wind termination shocks** of massive stars
  - ⇒ **Diffusive shock acceleration in wind termination shocks (Morlino et al. 2021)**
- The GCR refractories, which originate from the **acceleration and sputtering of dust grains** in SNR shocks, may also be produced **in superbubbles**
  - ⇒ **Update of the grain acceleration model of Ellison et al. (1997) based on current knowledge of dust in the ISM (Cristofari et al., in preparation)**



# Extra slides

# Cosmic Ray Spectrum



Cosmic rays are of Galactic origin up to the “Knee” at  $\sim 3 \times 10^{15}$  eV (3 PeV) and probably even up to  $\sim 10^{18}$  eV

# GCR composition model

- Relative contributions of the ISM dust and gas reservoirs in the measured GCR abundances:

$$C_{\text{dust}}(i) = \frac{C_{\text{mes}}(i)}{\left(1 + \frac{1-f_d(i)}{f_d(i)} \cdot \frac{x_w f_w(i) f_{A/Q}^w(i) + (1-x_w) f_{A/Q}^{\text{SC}}(i)}{\epsilon}\right)},$$

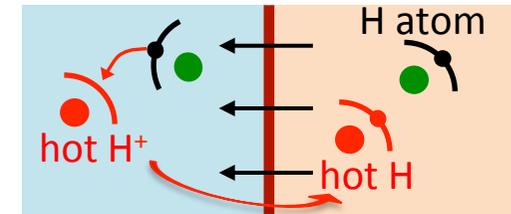
$$C_{\text{gas}}(i) = \frac{C_{\text{mes}}(i)}{\left(1 + \frac{f_d(i)}{1-f_d(i)} \cdot \frac{\epsilon}{x_w f_w(i) f_{A/Q}^w(i) + (1-x_w) f_{A/Q}^{\text{SC}}(i)}\right)},$$

- GCR abundances arising from the gas reservoir of standard cosmic composition:

$$C_{\text{gas}}^{\text{SC}}(i) = \frac{C_{\text{mes}}(i)}{\left(1 + \frac{f_d(i)}{1-f_d(i)} \cdot \frac{\epsilon}{(1-x_w) f_{A/Q}^{\text{SC}}(i)} + \frac{x_w f_w(i) f_{A/Q}^w(i)}{(1-x_w) f_{A/Q}^{\text{SC}}(i)}\right)}$$

# Effects of the neutral return flux

- Production of hot neutral hydrogen that may deposit kinetic energy in the shock upstream medium
  - ⇒ Heating of the upstream plasma
  - ⇒ Reduction of the shock Mach number
  - ⇒ **Reduction of the particle acceleration efficiency**



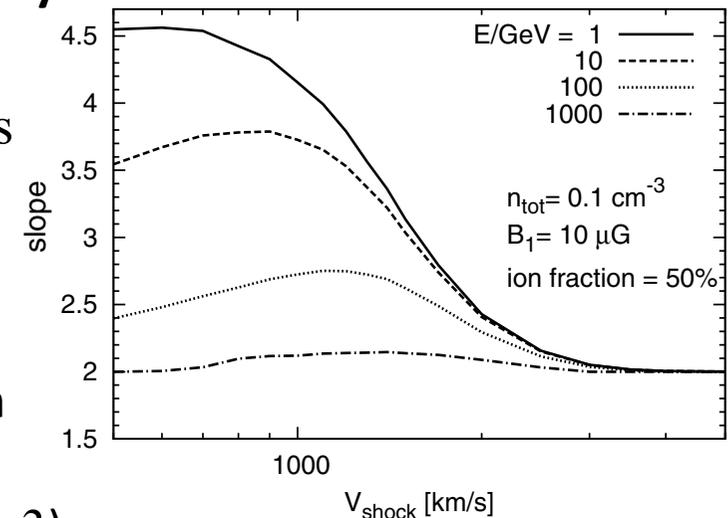
- For a neutral fraction of  $\sim 50\%$  (WNM), significant effects of the NRF for shock velocities  $\lesssim 3000$  km/s (Blasi et al. 2012; Morlino et al. 2013)

- Assuming CR production only for  $V_s > V_{s,\min}$  with
  - $V_{s,\min}$ (WNM) = 3000 km/s
  - $V_{s,\min}$ (WIM) = 300 km/s (ineffective acceleration in radiative shocks; Raymond et al. 2020)
  - $V_{s,\min}$ (SB) = 600 km/s ( $\Rightarrow$  Mach number  $M_{S,\min} = 3$ )

- Swept-up masses in the adiabatic stage:  $M_{s.-u.}^k \approx 68 M_\odot \left( \frac{E_{SN}}{10^{51} \text{ erg}} \right) \left( \frac{V_{s,\min}^k}{1000 \text{ km s}^{-1}} \right)^{-2}$

⇒ Relative contribution of the ISM phases to the GCR volatile production:

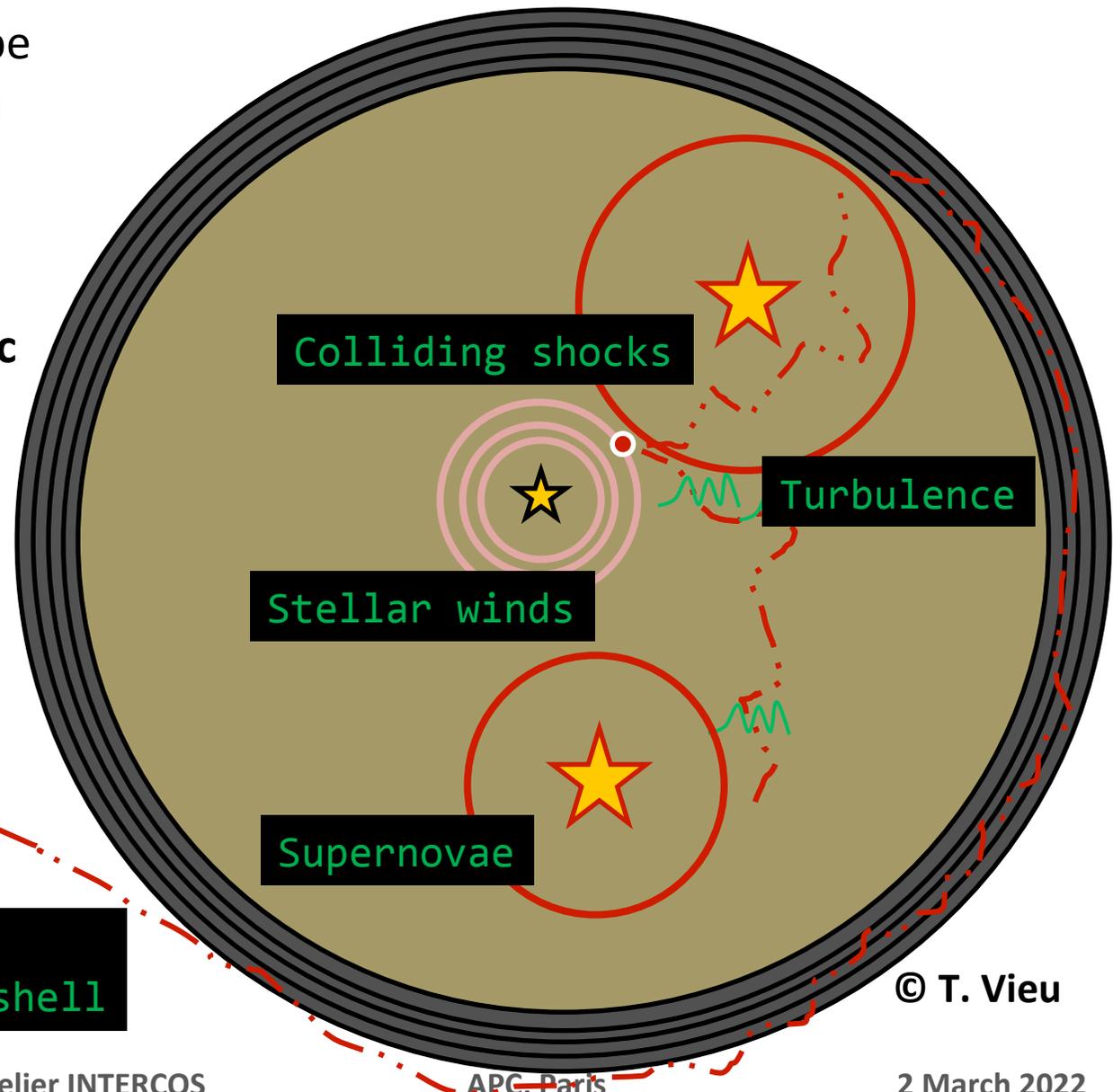
$$f_{SB}=55\%, f_{WIM}=44\%, f_{WNM}=1\% \Rightarrow \text{no good fit to the data } (\chi^2_{\min}=24.0)$$



# Cosmic-ray production in superbubbles

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- **Collective effects** should be taken into account in such energetics environments: **shock collisions**, particle **reacceleration** in multiple shocks and **hydromagnetic turbulence**... (Thibault Vieu's PhD thesis 2021)
- Production of cosmic rays **up to the knee of the spectrum and beyond?**



Propagation through the shell

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