

# Low Q-value decays for absolute mass scale determination of the neutrino

T. Eronen, Z. Ge, J. Kostensalo, A. de Roubin, J. Suhonen



# The determination of the absolute mass scale of the neutrino

The determination of the **absolute mass scale of the neutrino** is one of the **burning question in physics**

Neutrino mixing experiments have revealed that neutrinos have a non-zero mass

↳ BUT: not able to gauge the absolute mass values!  
(only the difference of squared neutrino masses)

## How to pinpoint the (anti)neutrino mass?

- Precisely
- In a model independent way



## Kinematical studies

- $\beta^-$  decay
- Electron capture

# Direct (anti)neutrino mass measurements

## Single Electron Capture

- Decay energy shared between:
  - Emitted neutrino
  - Atomic excitation of the daughter atom
  - (recoiling nucleus)
- Experimental approach:
  - Low temperature calorimetry measurements (ECHO, HOLMES, NuMECS)  
*L. Gastaldo et al., Eur. Phys. J. Special Topics* **226**, 1623 (2017)  
*M. Faverzani et al., J. Low Temp. Phys.* **184**, 922 (2016)  
*A. Faessler et al., Phys. Rev. C* **91**, 045505 (2015)

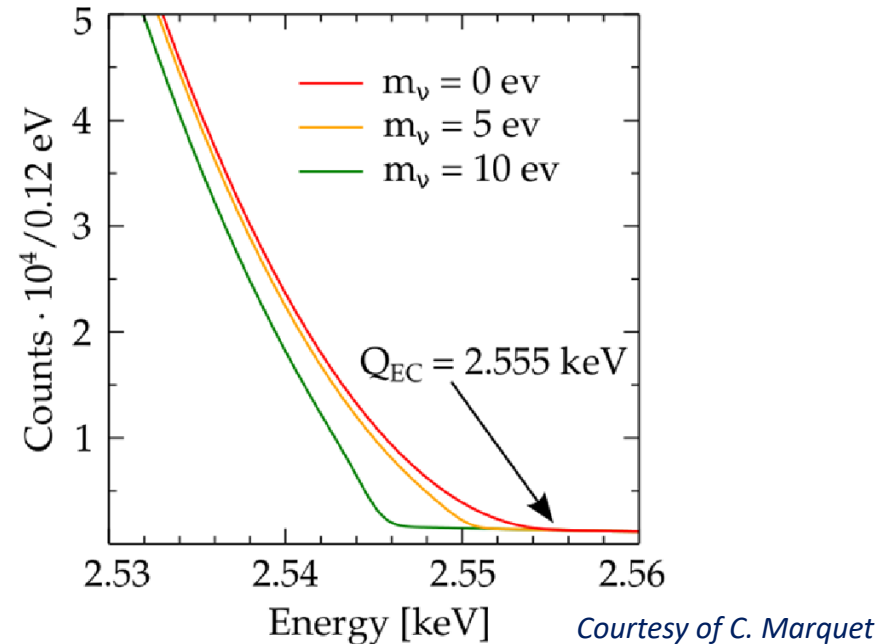
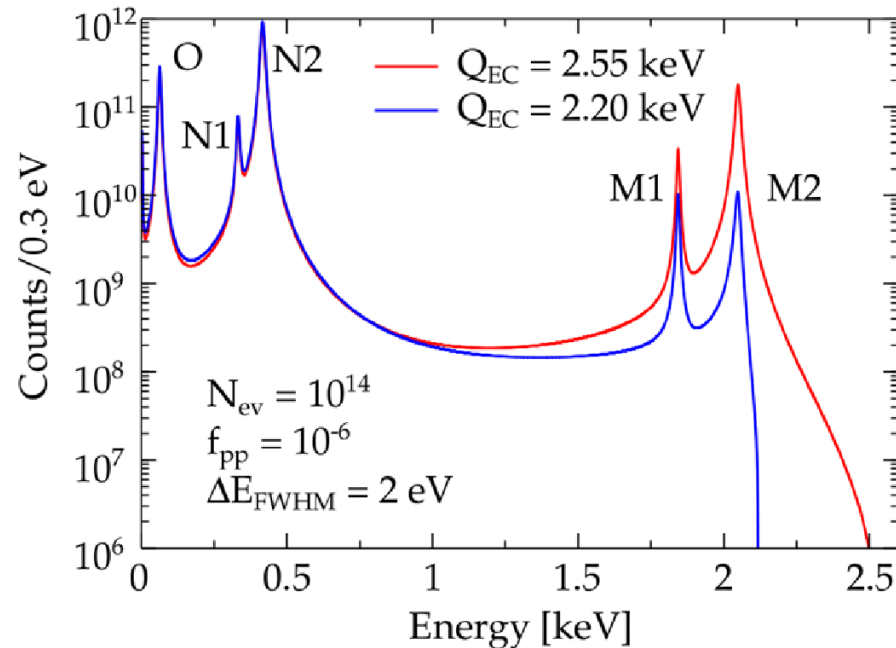
## Single $\beta^-$ decay

- Energy is distributed between the decay products
- Experimental approaches:
  - Electron spectroscopy (KATRIN)  
*C. Weinheimer, Prog. Part. Nucl. Phys.*, **48**, 141 (2002)  
*KATRIN collaboration, Nat. Phys.*, **18**, 160 (2022)
  - Cyclotron radiation emission spectroscopy (Super8)  
*B. Monreal et al. Mass. Phys. Rev., D* **80**, 051301 (2009)

**Interested in the endpoint of the energy spectrum!**

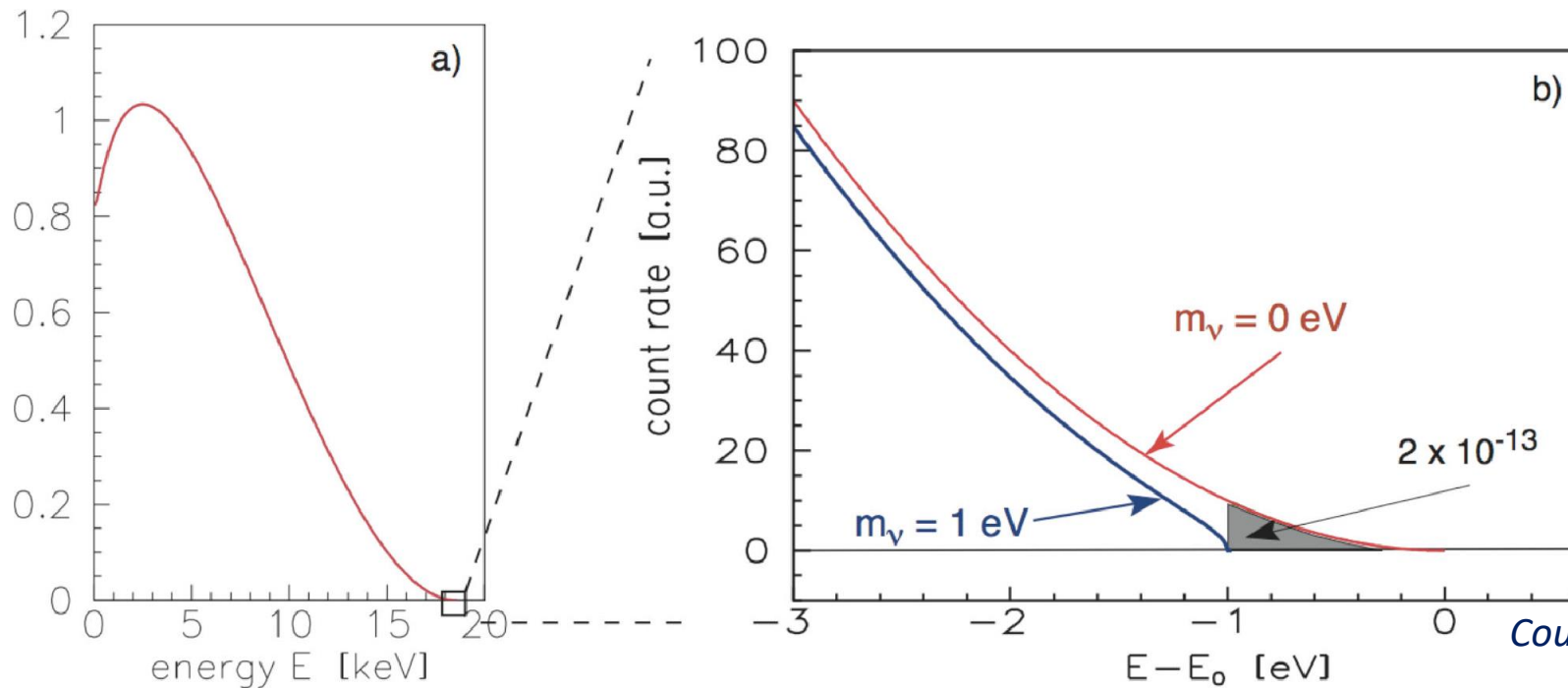
Need decay with a **small Q value** to maximize the sensitivity to the neutrino mass!

# Single electron capture



- Measurement of emitted particles and their energies following the EC
- The mismatch of energy available in decay and emitted energy reveals the neutrino mass
- Upper limit of the effective electron-neutrino mass  $m_{\nu_e} = 150 \text{ eV}/c^2$  (95% confidence level)

# Single $\beta^-$ decay



Courtesy of C. Marquet

- Measurement of emitted electron energy
- The maximum energy is linked to the neutrino mass
- Upper limit of the effective electron-neutrino mass  $m_{\nu_e} = 0.8 \text{ eV}/c^2$  (90% confidence level)

*KATRIN collaboration, Nat. Phys., 18, 160 (2022)*

# Ultra low Q-value $\beta$ decay for neutrino physics

**Interested in the endpoint of the energy spectrum!**

Need decay with a **small Q value** to maximize the sensitivity to the neutrino mass!



Smallest GS-to-GS Q-values around 1 keV

# Ultra low Q-value $\beta$ decay for neutrino physics

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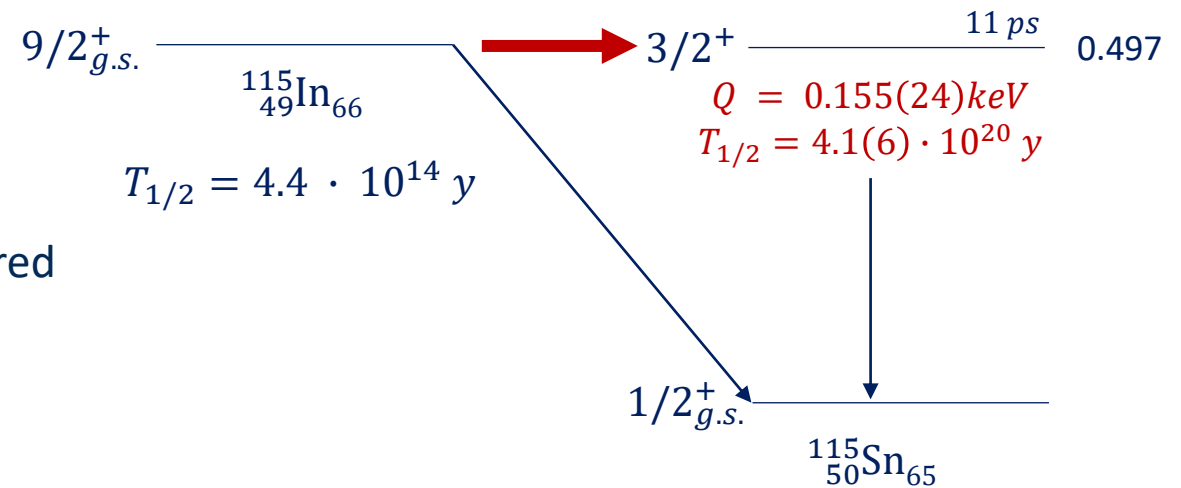
Smallest GS-to-GS Q-values around 1 keV

- Possibility to use decays with even smaller Q-values

- $^{115}\text{In}(9/2^+) \xrightarrow{\beta^-} ^{115}\text{Sn}(3/2^+)$
- $Q = 0.155(24) \text{ keV}$  → smallest  $\beta$  Q-value ever measured
- **A « detector » for the neutrino mass ?**

- Other possible low Q-values

- Many potential candidates
- Need to identify them with Penning traps



B.J. Mount, M. Redshav, E.G. Myers, *Phys. Rev. Lett.* 103 122502 (2009)  
J. S. E. Wieslander, *Phys. Rev. Lett.* 103, 122501 (2009)  
C.M. Cattadori, *Nucl. Phys. A*, 748, 333-347 (2005)

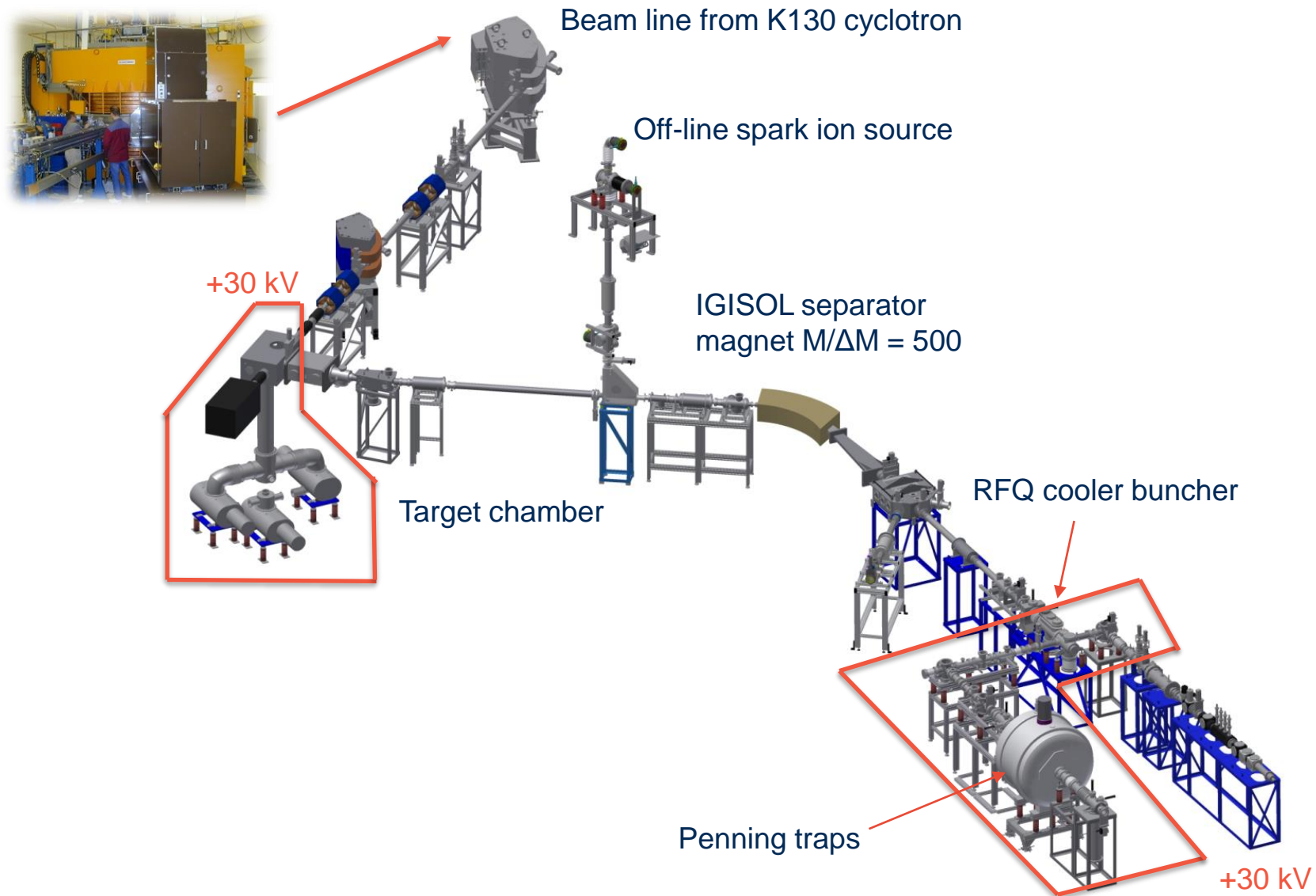


# Ultra low Q-value $\beta$ decay for neutrino physics

<b>Mother( <math>J^\pi</math> )</b> <i>Ground state</i>	<b>Daughter ( <math>J^\pi</math> )</b> <i>Excited state</i>	<b>E* (keV)</b>	<b>Q (keV)</b>	<b>Type of decay</b>
$^{135}\text{Cs}(7/2^+)$	$^{135}\text{Ba}(11/2^-)$	268.218(20)	0.5(12)	$\beta^-$
$^{159}\text{Dy}(3/2^-)$	$^{159}\text{Tb}(5/2^-)$	363.5449(14)	1.7(12)	EC
	$^{159}\text{Tb}(11/2^+)$	362.050(40)	3.2(12)	EC
$^{111}\text{In}(9/2^+)$	$^{111}\text{Cd}(7/2^+)$	853.94(7)	6.3(34)	EC
	$^{111}\text{Cd}(3/2^+)$	855.6(10)	4.6(36)	EC
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	$^{111}\text{Cd}(3/2^+)$	866.60(6)	-6.4(34)	EC
$^{131}\text{I}(7/2^+)$	$^{131}\text{Xe}(9/2^+)$	971.22(13)	-0.42(61)	$\beta^-$

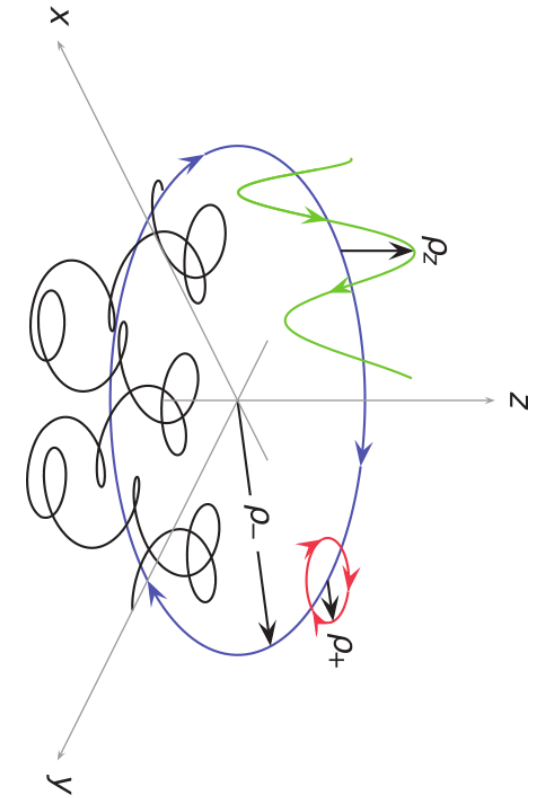
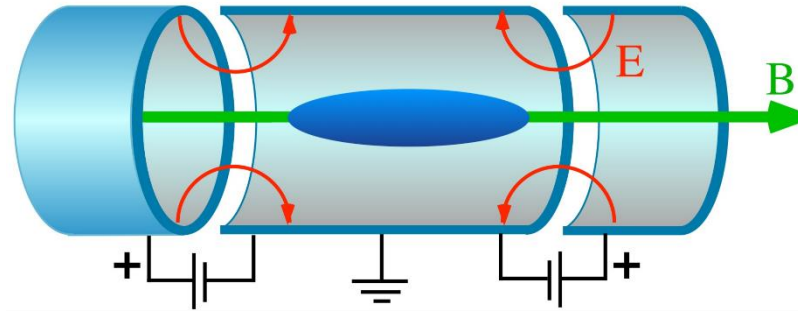


# The IGISOL-4 facility



# Penning trap

- \ Radial confinement
  - strong homogeneous magnetic field
- \ Axial confinement
  - electric field



## 3 ion motions

- Axial
- Magnetron
- Reduced cyclotron



## 3 ion frequencies

- $\nu_z$
- $\nu_-$
- $\nu_+$

- Invariance theorem

$$\nu_c^2 = \nu_-^2 + \nu_+^2 + \nu_z^2$$

- Cyclotron frequency

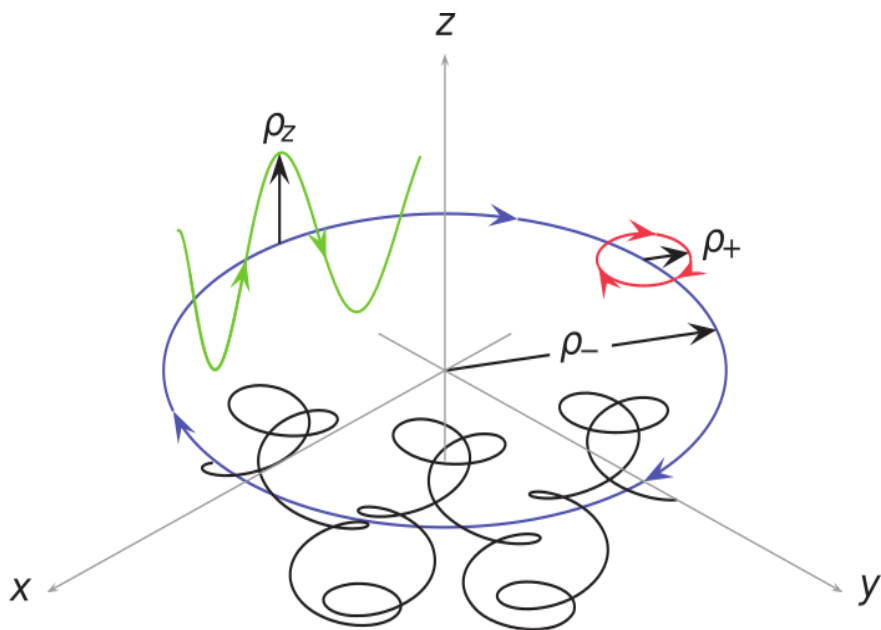
$$\nu_c = \nu_- + \nu_+$$

## Cyclotron frequency

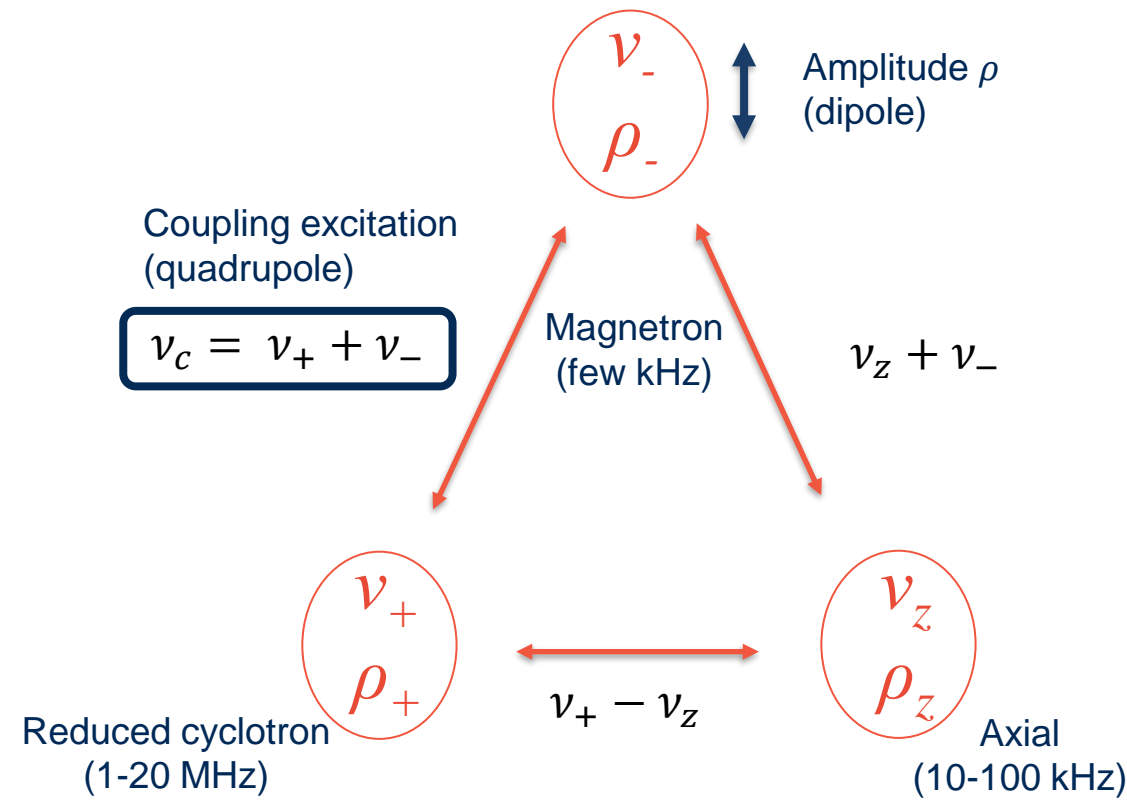
$$\nu_c = \frac{1}{2\pi} \frac{q}{m} B$$

q : electric charge  
 B : magnetic field  
 m : mass

# Penning trap



Control amplitudes  
  
 Convert motions



## 3 ion motions

- Axial
- Magnetron
- Reduced cyclotron



## 3 ion frequencies

- $\nu_z$
- $\nu_-$
- $\nu_+$

## Invariance theorem

$$\nu_c^2 = \nu_-^2 + \nu_+^2 + \nu_z^2$$

## Cyclotron frequency

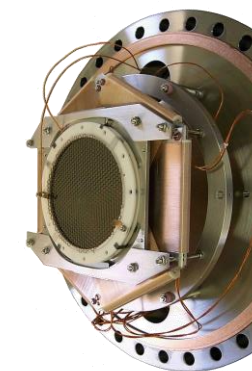
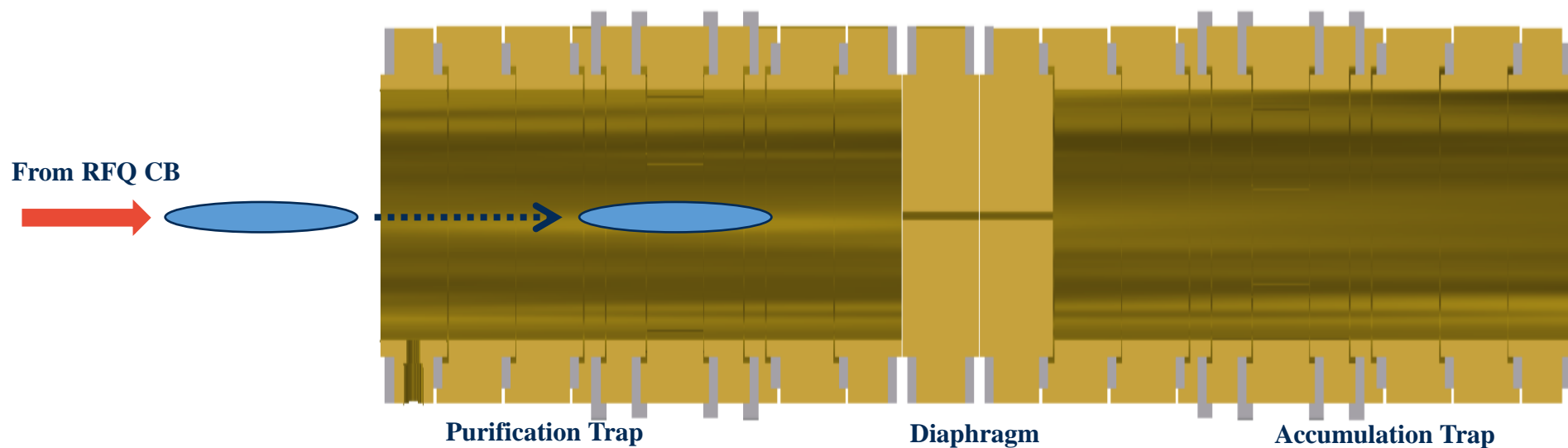
$$\nu_c = \nu_- + \nu_+$$

## Cyclotron frequency

$$\nu_c = \frac{1}{2\pi} \frac{q}{m} B$$

q : electric charge  
 B : magnetic field  
 m : mass

# Buffer gas cooling technique



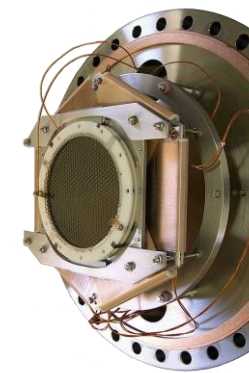
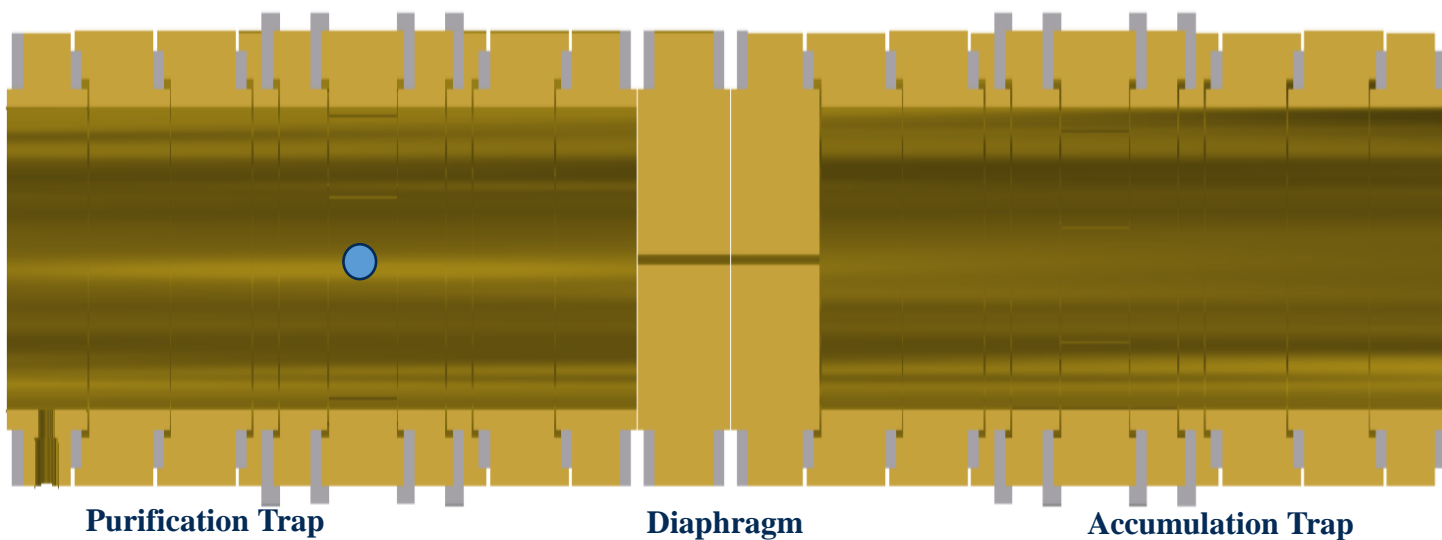
MicroChannel Plate detector  
(position sensitive)

Injection and capture of the ions

➤ 100  $\mu$ s

# Buffer gas cooling technique

From RFQ CB



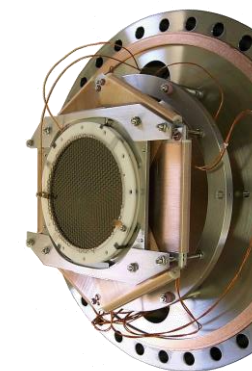
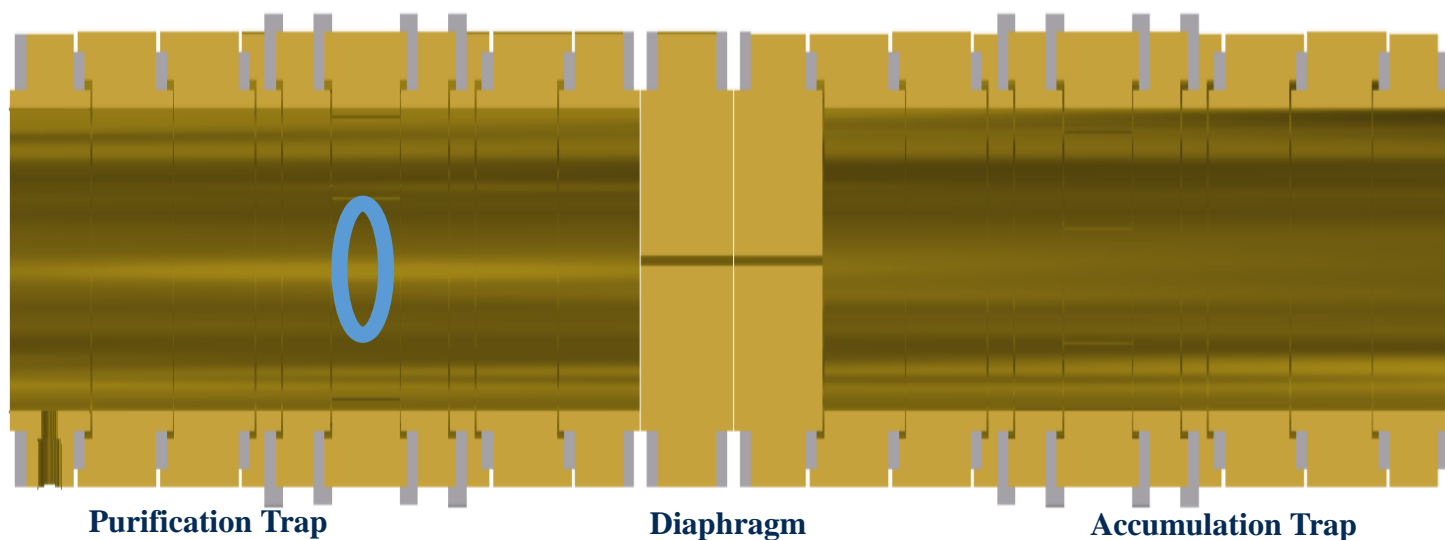
MicroChannel Plate detector  
(position sensitive)

Cooling of radial and axial motions

➤ 30 - 100 ms

# Buffer gas cooling technique

From RFQ CB



MicroChannel Plate detector  
(position sensitive)

## Buffer gas cooling technique

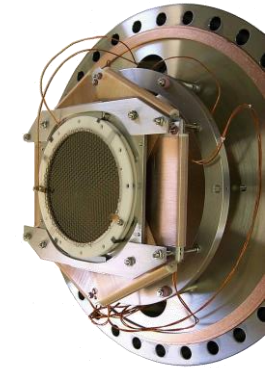
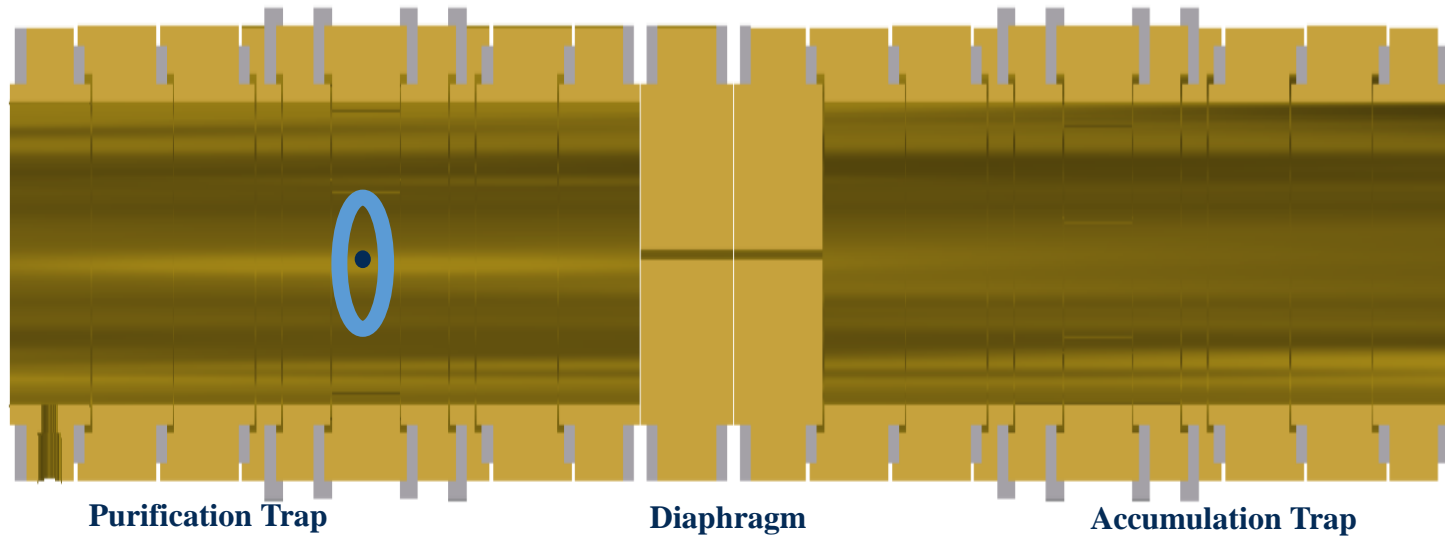
1. Magnetron excitation:  $\nu_{-}$   
➤ 10 ms



This excitation is mass **independent**

# Buffer gas cooling technique

From RFQ CB



MicroChannel Plate detector  
(position sensitive)

## Buffer gas cooling technique

1. Magnetron excitation:  $\nu_-$ 
  - 10 ms
2. Quadrupole excitation:  $\nu_c = \nu_+ + \nu_-$ 
  - Convert the radial magnetron motion ( $\nu_-$ ) into modified cyclotron motion ( $\nu_+$ )
  - 40 – 400 ms
  - Centering the « good guys »

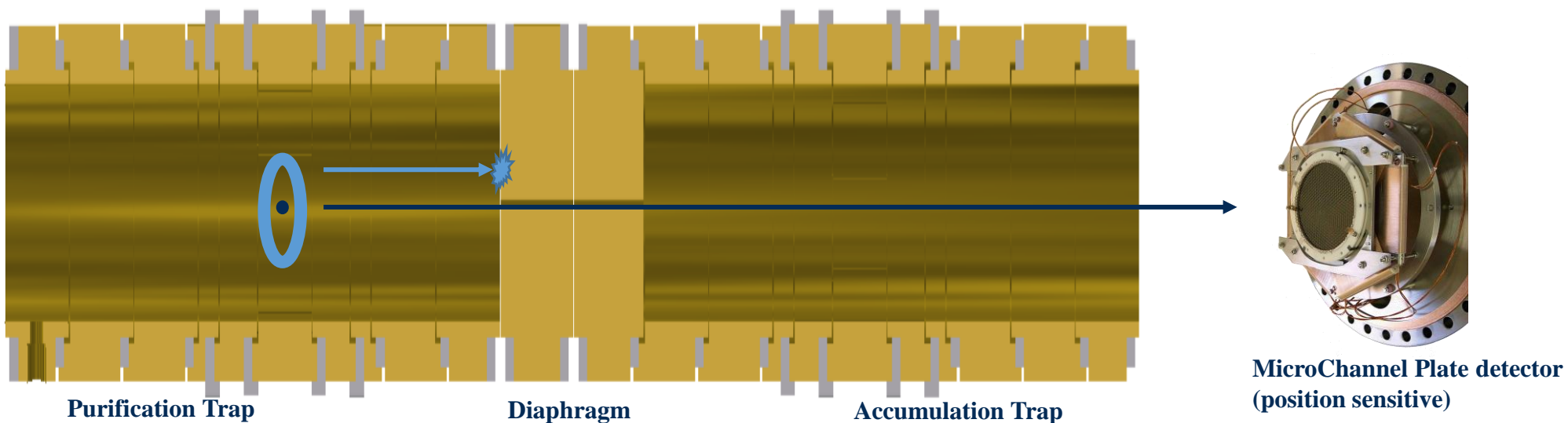


- \ This excitation is mass **dependent !!**
- \ Magnetron motion is amplified
- \ modified cyclotron motion is damped



# Buffer gas cooling technique

From RFQ CB



## Mass resolving power of Penning trap

$$R = \frac{M}{\Delta M} = \frac{\nu_c}{\Delta \nu_c} \propto \frac{1}{\Delta \nu_c} \left( \frac{q}{m} B \right) \rightarrow \text{Scaling}$$

Constant for each method 0.1 ... 100 Hz

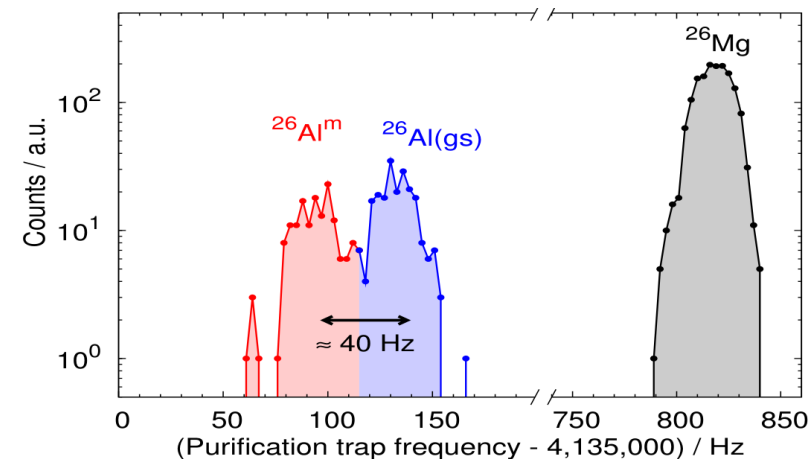
## Rule of thumb:

$B = 7 \text{ T}$ ,  $q = 1e$ ,  $m = 100u$   
 $\rightarrow \nu_c \sim 1 \text{ MHz}$ ,  $100 \text{ keV/Hz}$

## Buffer gas cooling technique

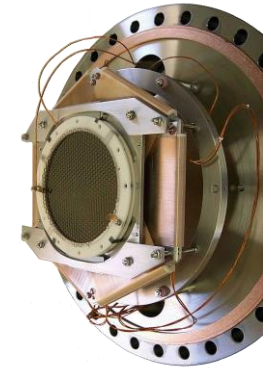
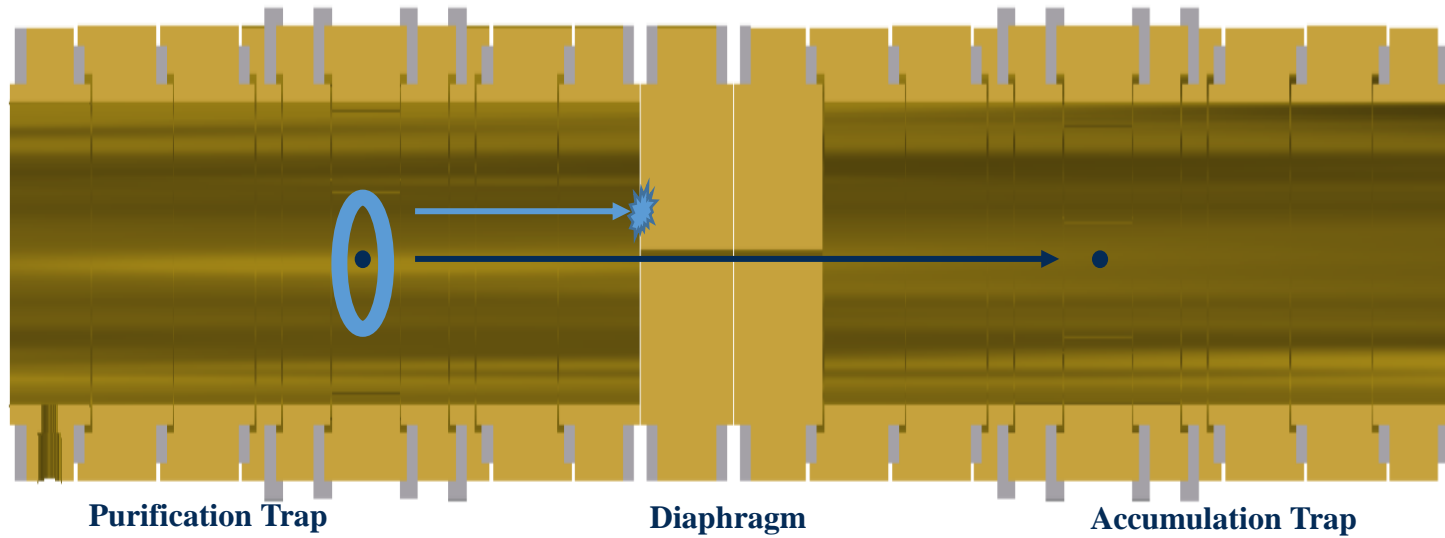
- Gas filled purification trap
- Separation
  - 10 ... 100 Hz
  - $R = 10^4 \dots 10^5$
- Also cooling

G. Savard et al., Phys. Lett. A **158** (1991) 247–252



# Time-of-flight ion cyclotron resonance technique (ToF-ICR)

From RFQ CB



MicroChannel Plate detector  
(position sensitive)

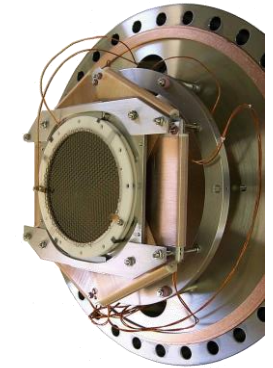
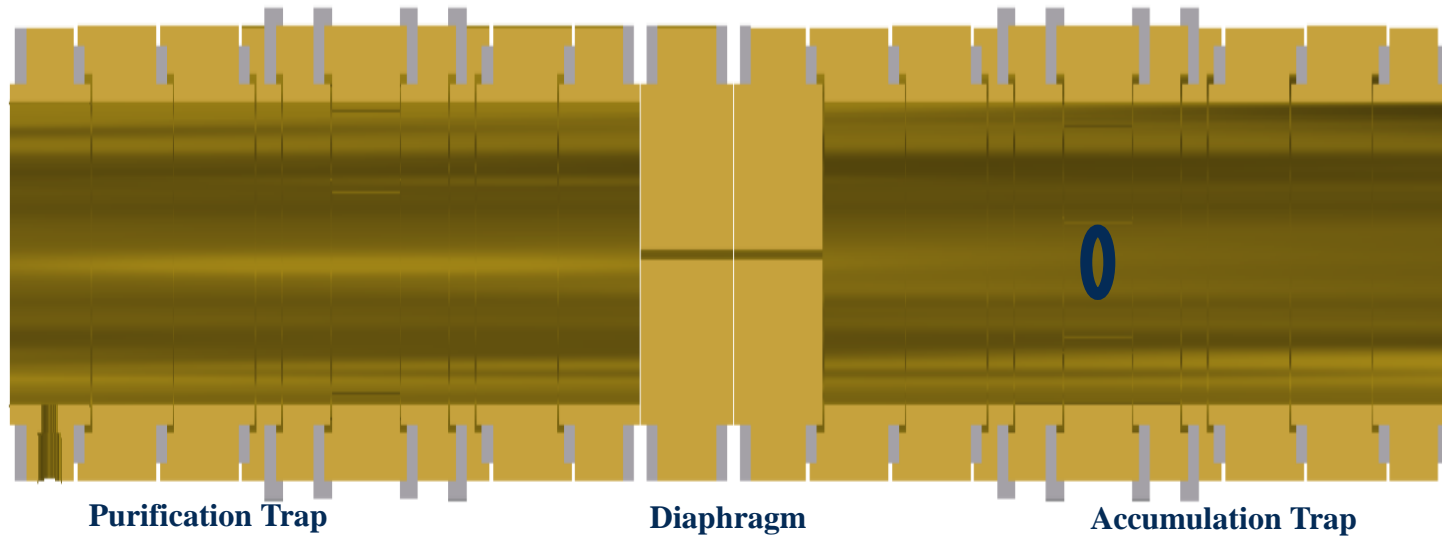
**ToF-ICR technique** (Time-of-Flight Ion Cyclotron Resonance)



Transfer of the clean ion sample to the measurement trap

# Time-of-flight ion cyclotron resonance technique (ToF-ICR)

From RFQ CB



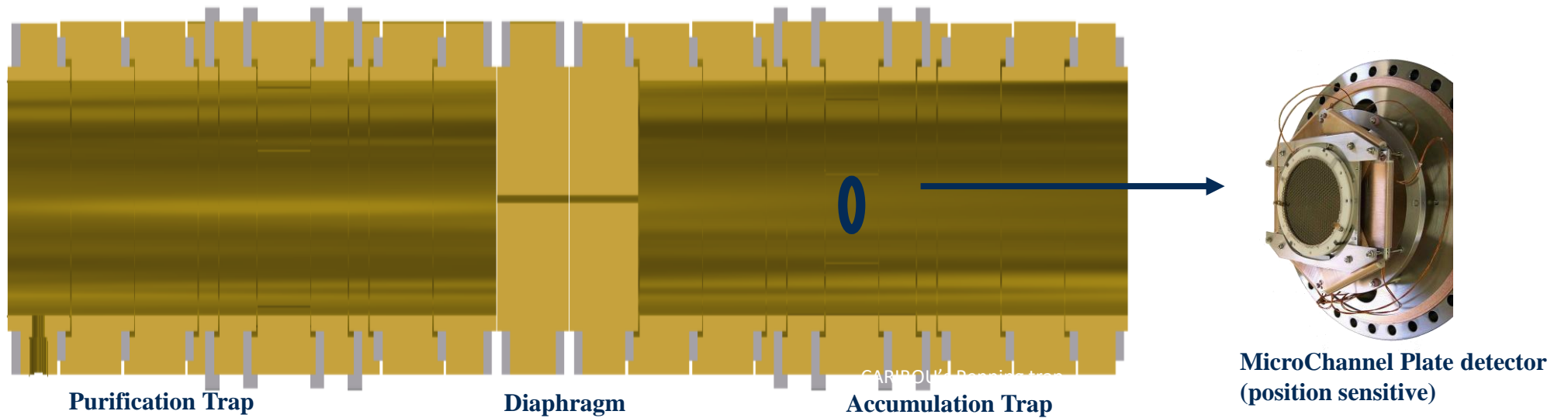
MicroChannel Plate detector  
(position sensitive)

## ToF-ICR technique

1. Magnetron excitation:  $\nu_z$ 
  - 10 ms
  - Give a slow radial motion to the ions

# Time-of-flight ion cyclotron resonance technique (ToF-ICR)

From RFQ CB



## ToF-ICR technique

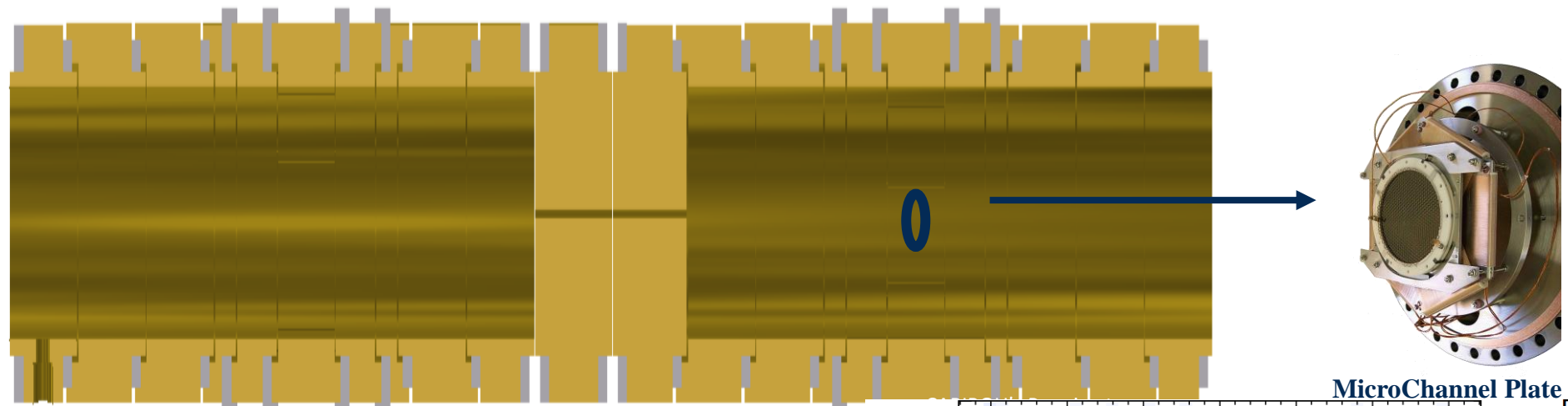
1. Magnetron excitation:  $\mathbf{v}_-$ 
  - 10 ms
  - Give a slow radial motion to the ions
2. Quadrupole excitation:  $\mathbf{v}_c = \mathbf{v}_+ + \mathbf{v}_-$ 
  - Convert the radial magnetron motion ( $\mathbf{v}_-$ ) into modified cyclotron motion ( $\mathbf{v}_+$ )
  - 40 – 400 ms



- \ The modified cyclotron motion is a fast motion
- \ The ions are not recentered (no gas)
- \ Extraction towards the MCP

# Time-of-flight ion cyclotron resonance technique (ToF-ICR)

From RFQ CB

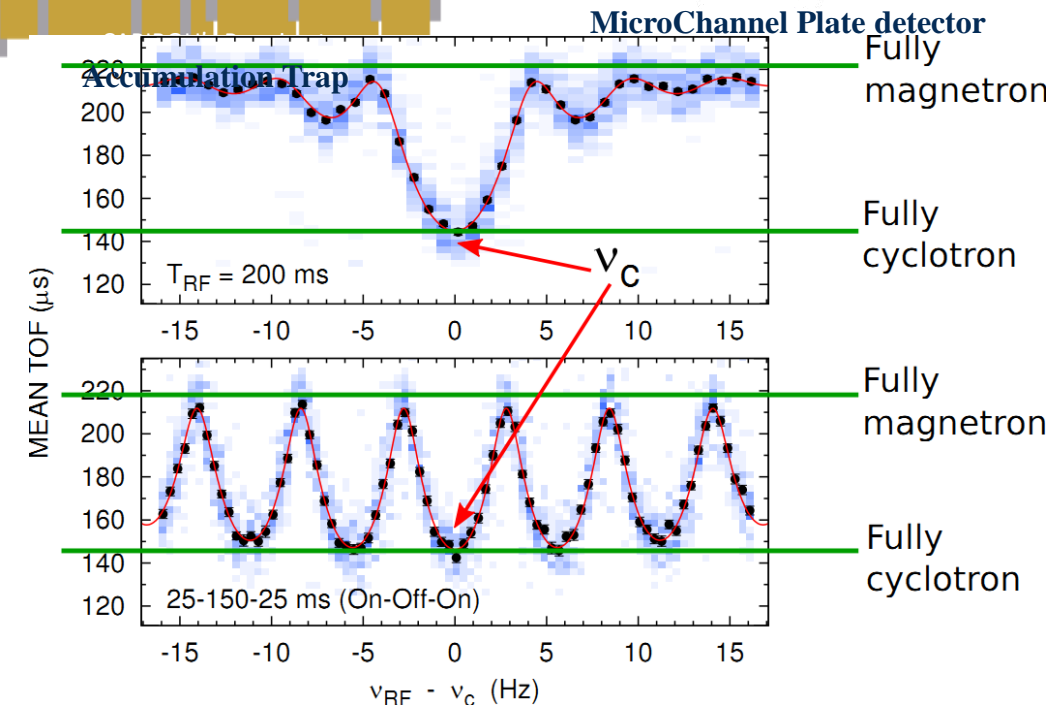


## ToF-ICR technique

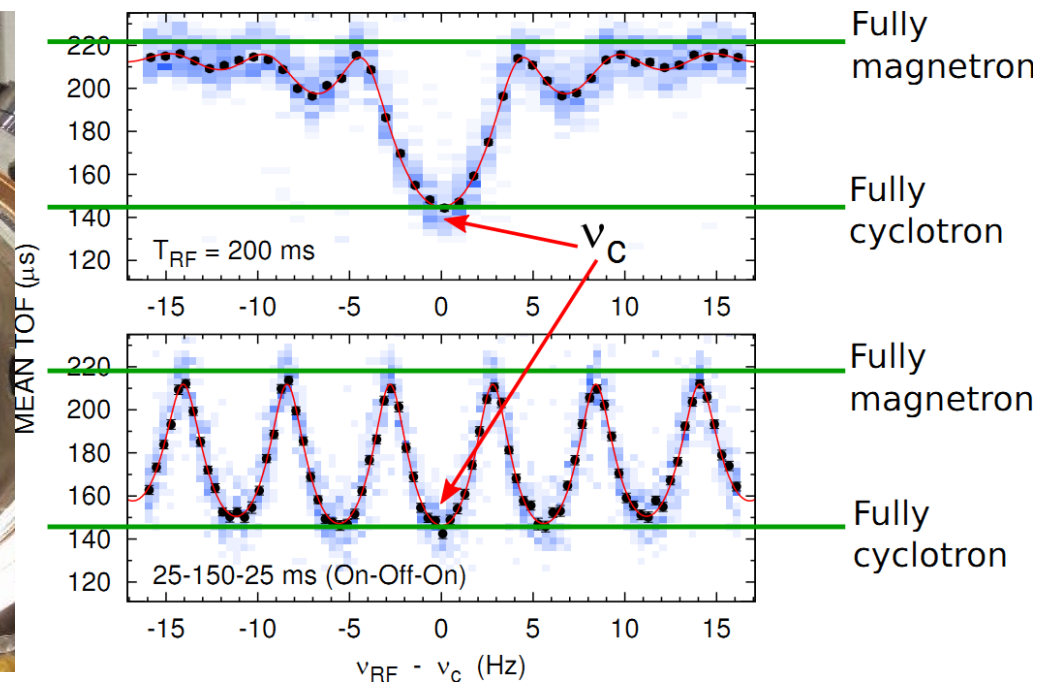
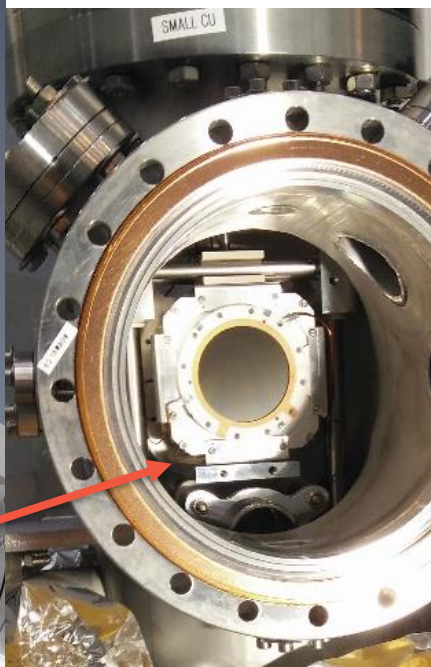
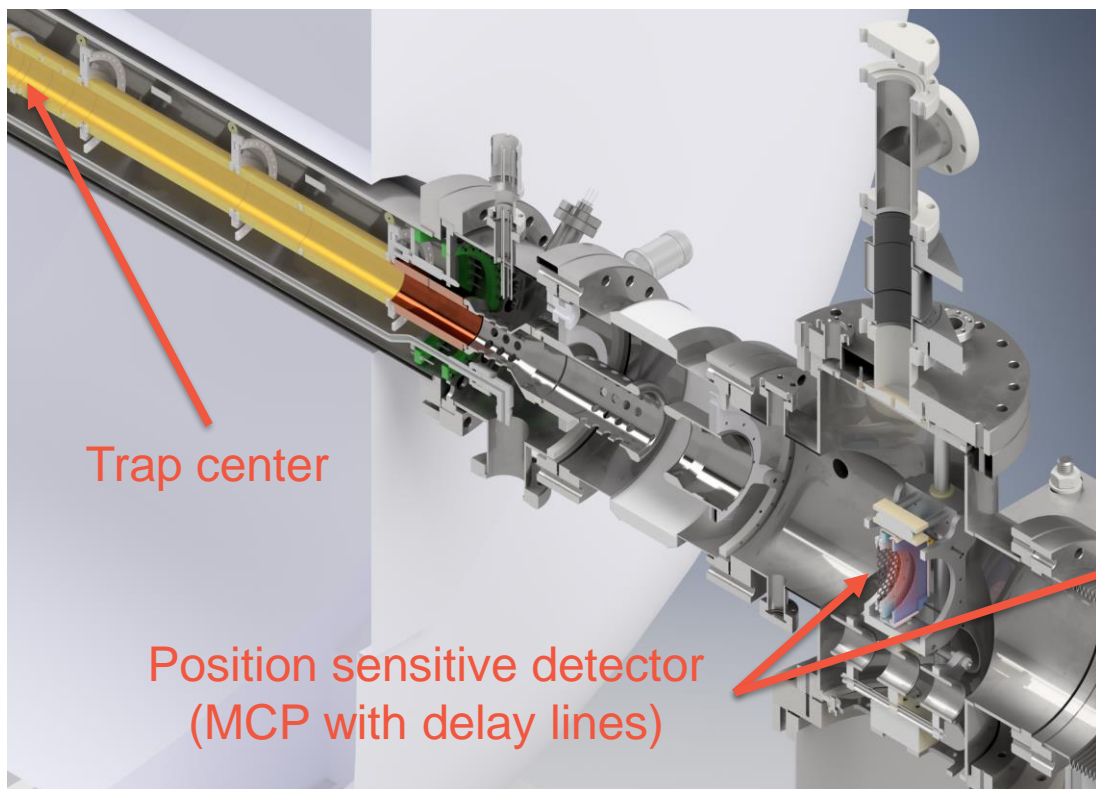
- Excitation around  $\nu_c$ 
  - / Large magnetic moment ( $\nu_+$  amplitude) for resonant ions
  - /  $\mu \sim \omega_+^2 \rho_+^2 - \omega_-^2 \rho_-^2$  ( $E_p = -\vec{\mu} \cdot \vec{B}$ )
- Ejection
  - / Radial kinetic energy change to axial energy
- ToF is reduced

## Ramsey ToF-ICR technique

- Improves precision, reduces measurement time  $\sim x3$

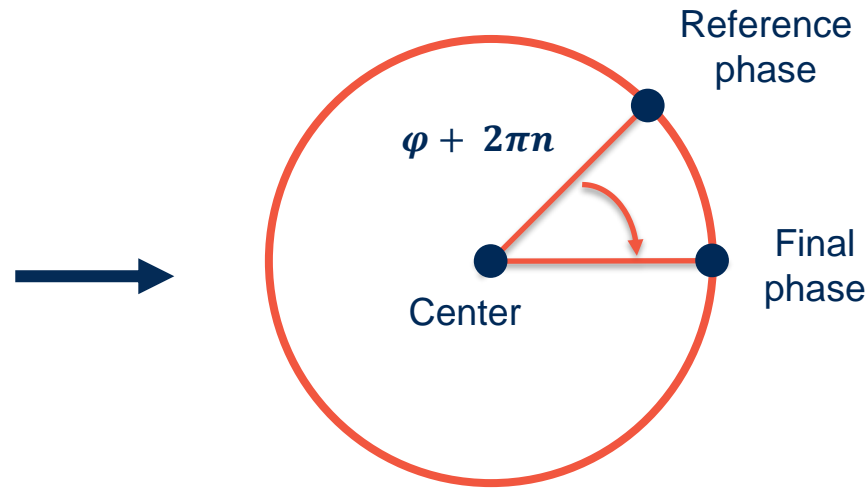
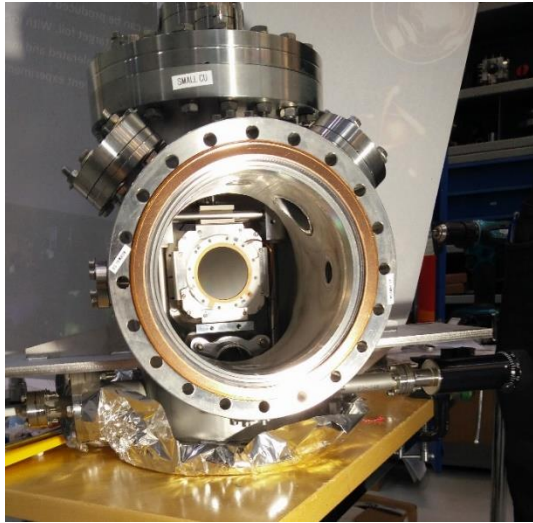


# Position sensitive detector

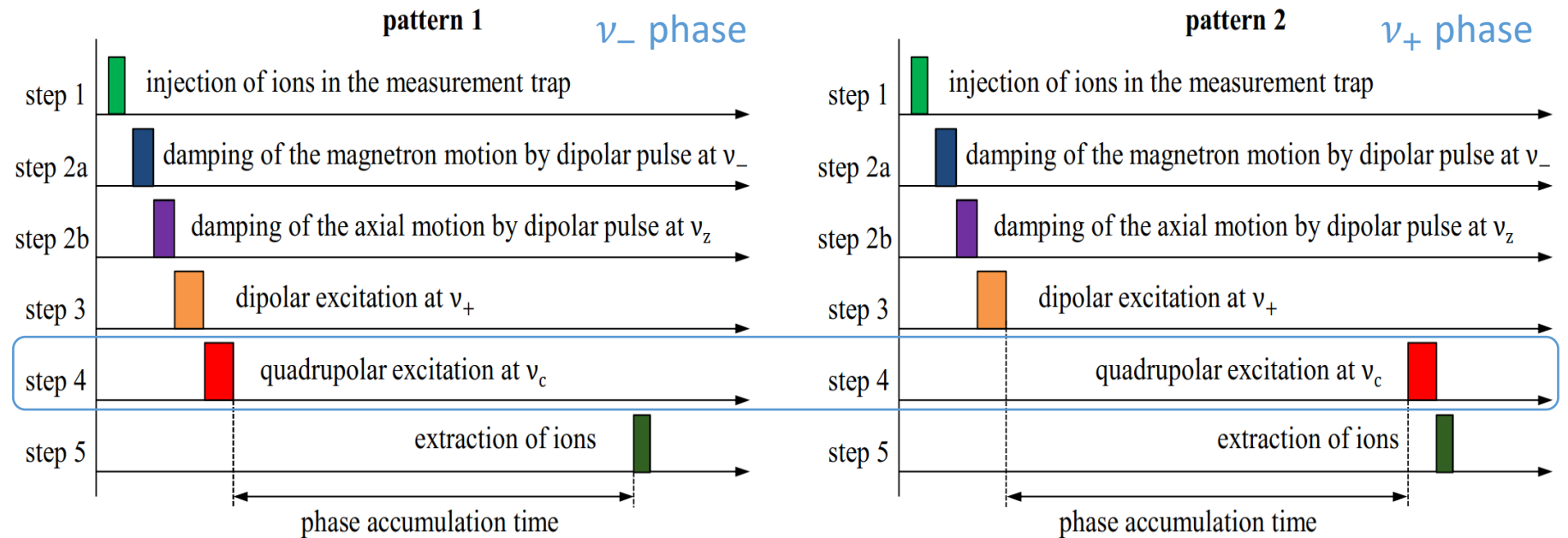




# Phase-imaging ion cyclotron resonance technique (PI-ICR)

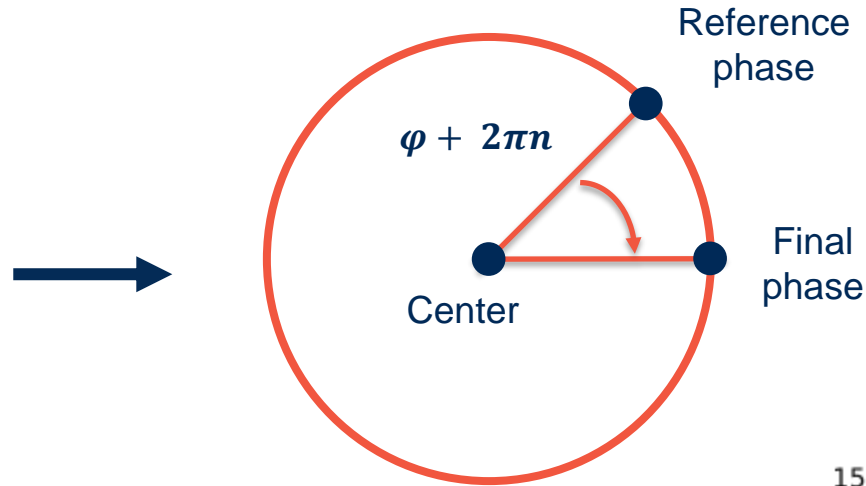
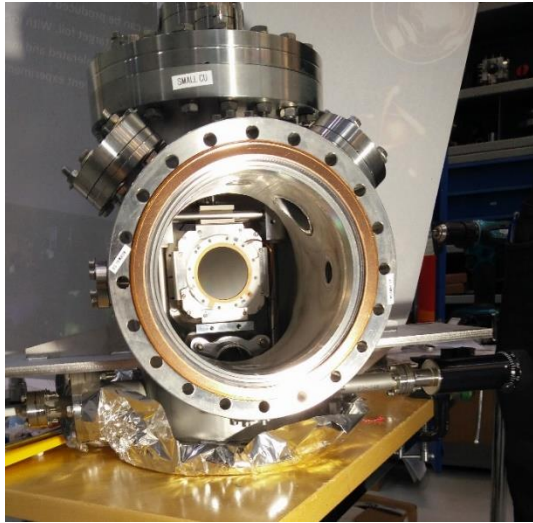


$$\nu = \frac{\varphi + 2\pi n}{2\pi t}$$





# Phase-imaging ion cyclotron resonance technique (PI-ICR)



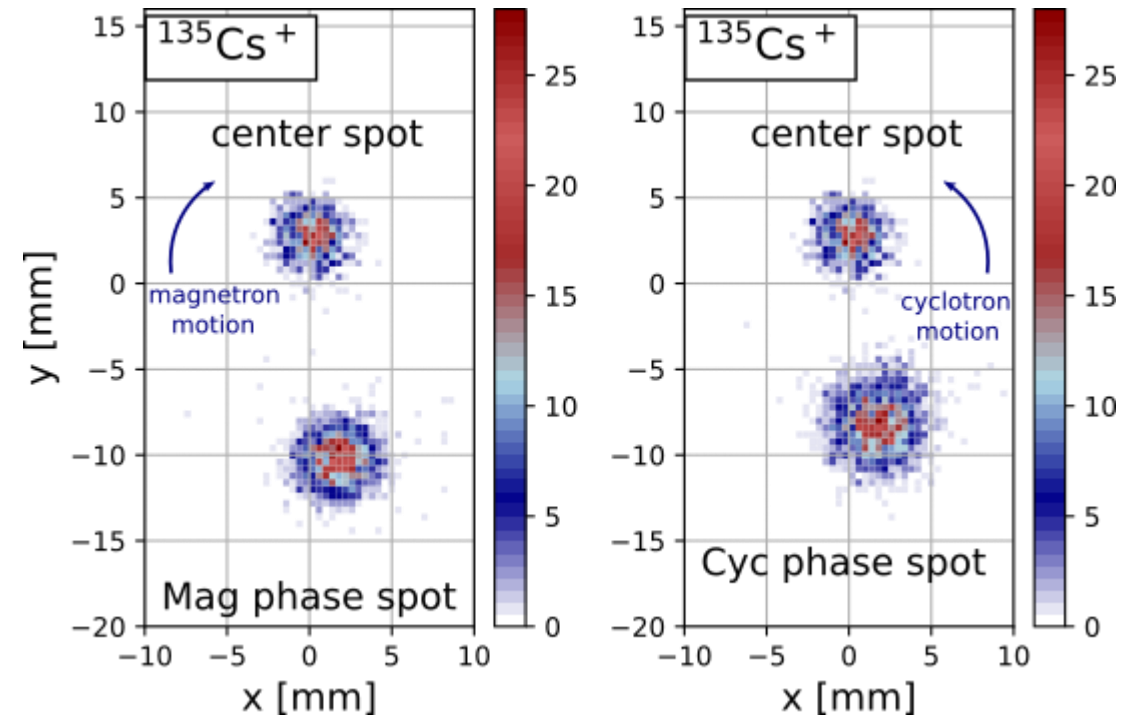
$$v = \frac{\varphi + 2\pi n}{2\pi t}$$

## Advantages of PI-ICR over TOF-ICR:

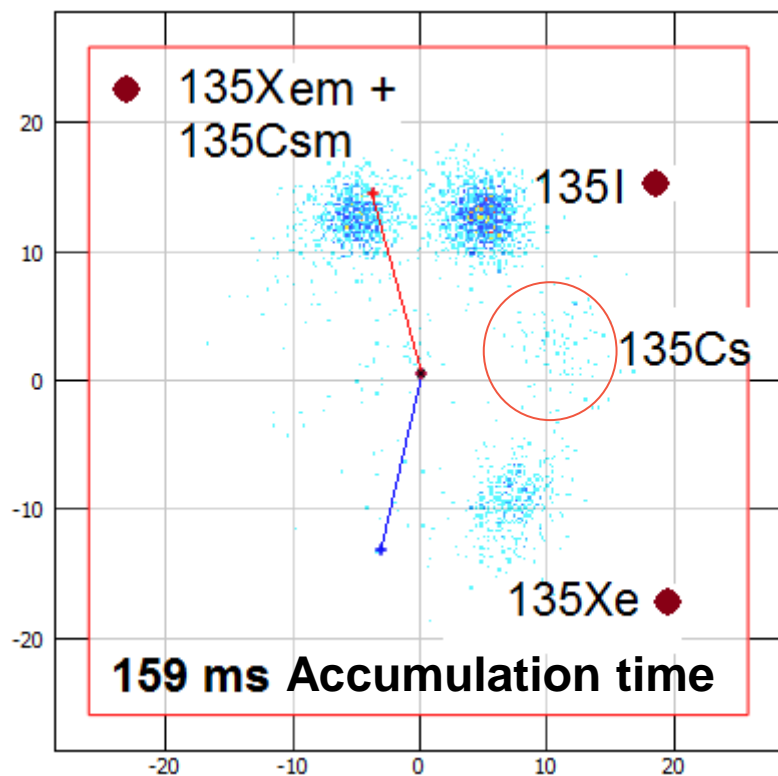
- No scanning of frequency, just data collection to one spot
- x40 faster x5 increase in precision

## Requirements:

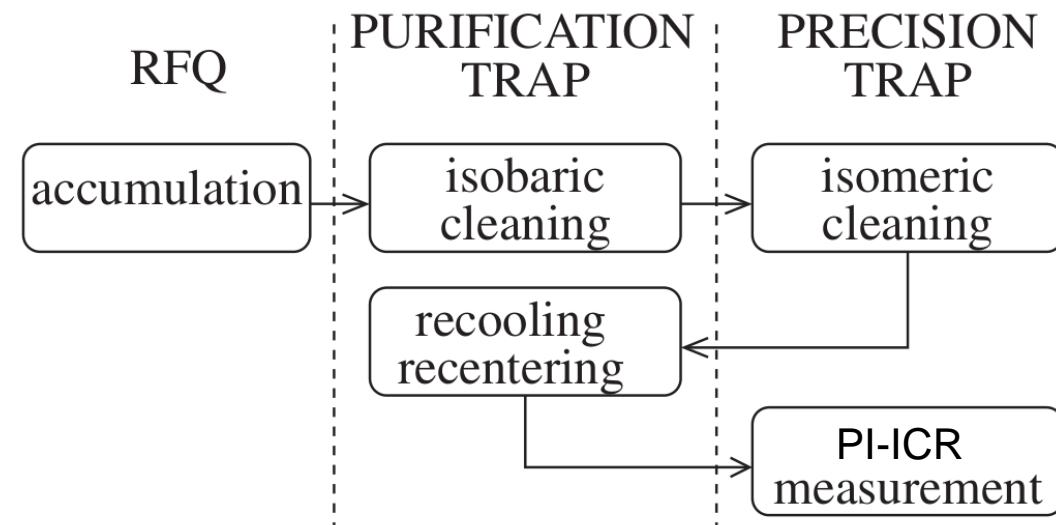
- Need to prepare the ion more carefully
- Damping of ion motions; longer setting up time
- Sensitive to voltage fluctuations



# Isomeric cleaning



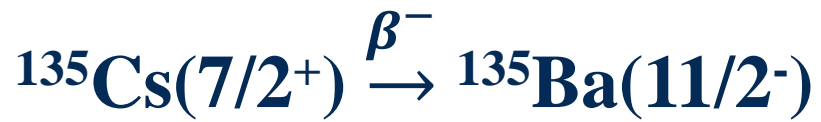
Ramsey cleaning



Same cleaning processed was needed for different measurements presented here

# Ultra low Q-value $\beta$ decay for neutrino physics

<b>Mother( <math>J^\pi</math> )</b> <i>Ground state</i>	<b>Daughter ( <math>J^\pi</math> )</b> <i>Excited state</i>	<b>E* (keV)</b>	<b>Q (keV)</b>	<b>Type of decay</b>
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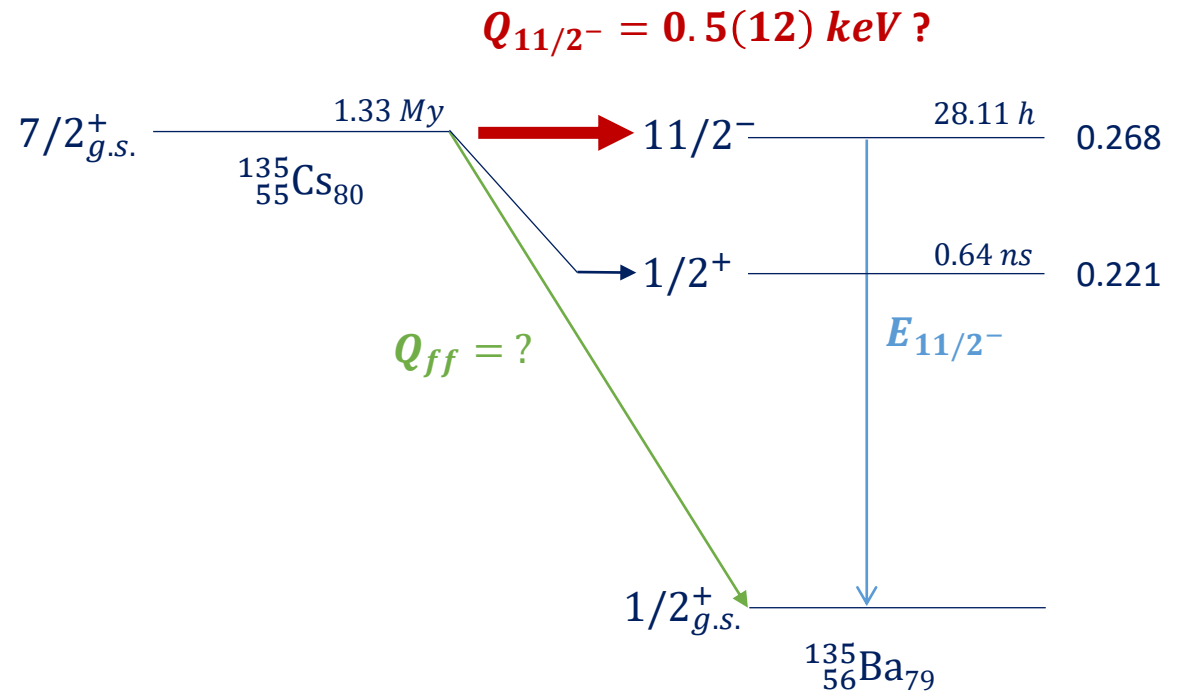
**But :** The transition is poorly known!

- Not yet proven that it is energetically possible :  $Q_{11/2^-} < 0$  ?
- Decay never observed directly

To determine  $Q_{11/2^-}$  we need:

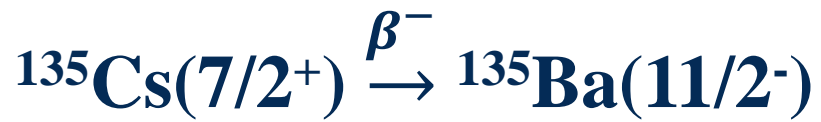
- The excitation energy  $E_{11/2^-} = 268.218(20) \text{ keV}$
- The ground-state-to-ground-state Q-value  $Q_{ff}$
- $Q_{ff} = 268.9(1.1) \text{ keV}$  [1]

$$Q_{ff} = 268.9(1.1) \text{ keV} \rightarrow Q_{11/2^-} = 0.5(1.2) \text{ keV}$$



**New measurement of  $Q_{ff}$  with Penning traps!!**

[1] M. Wang et al., Chin. Phys. C 41, 030003 (2017)



$$Q_{gs-gs} = 268.66(30) \text{ keV}$$

$$\rightarrow Q_{11/2^-} = 0.44(31) \text{ keV}$$

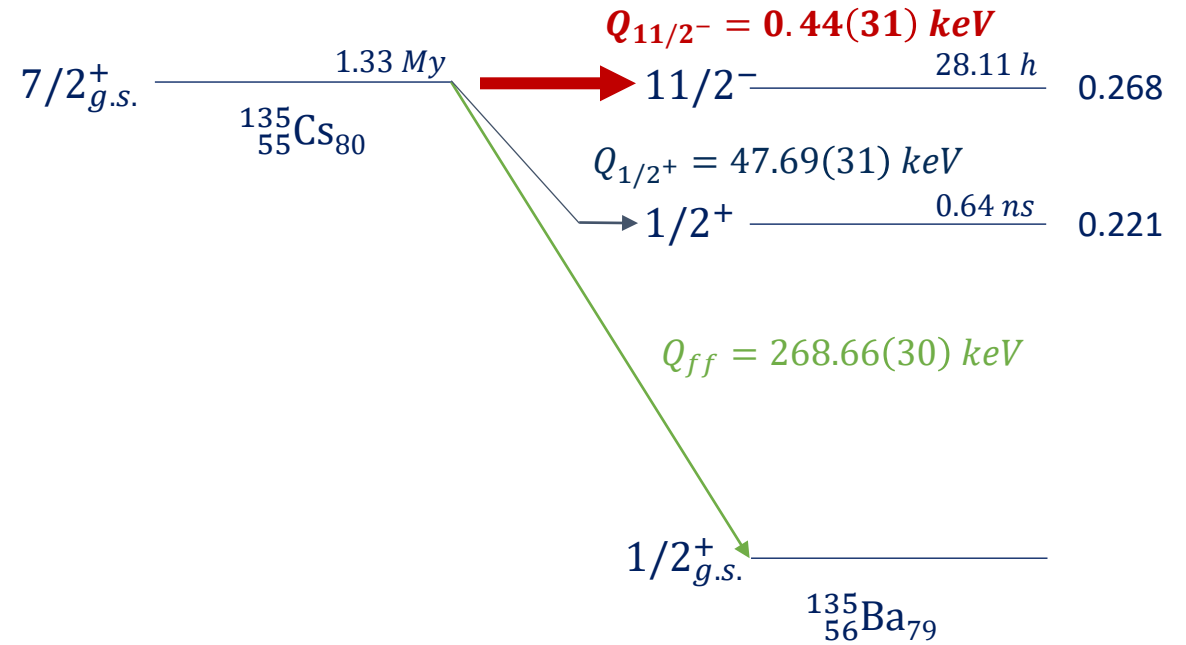
$$\rightarrow Q_{1/2^+} = 47.69(31) \text{ keV}$$

Half-lives calculations computed with NuShell@MSU

- $T_{1/2} (^{135}\text{Ba}(11/2^-)) = 8.2(32) \cdot 10^{11} \text{ y}$
- $T_{1/2} (^{135}\text{Ba}(1/2^+)) = 6.5(17) \cdot 10^{13} \text{ y}$



Branching ratio  $(0.04 - 16) \cdot 10^{-6}$



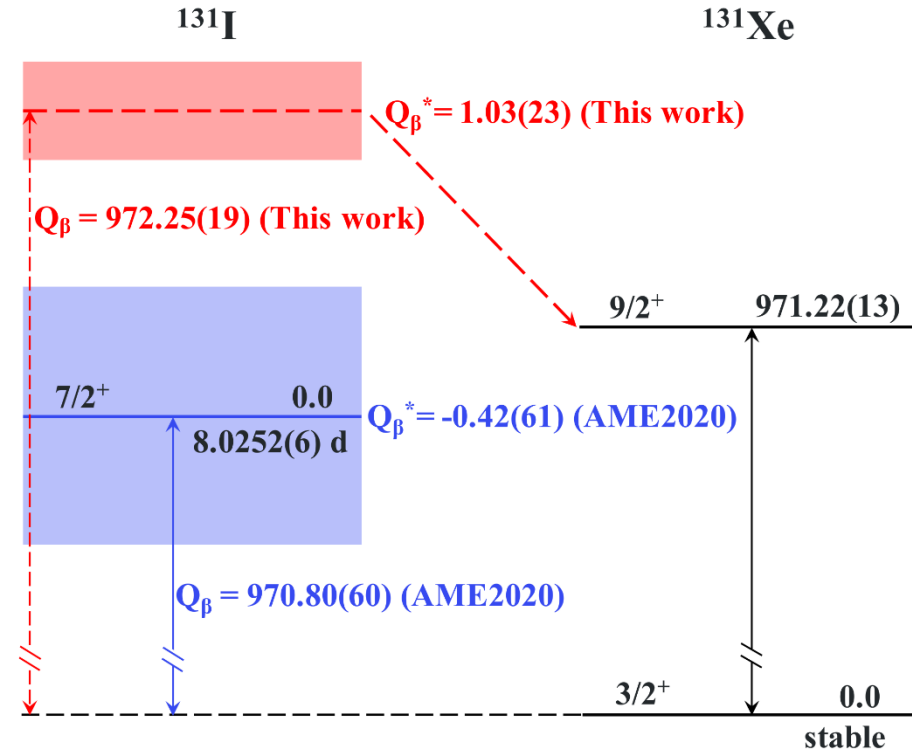
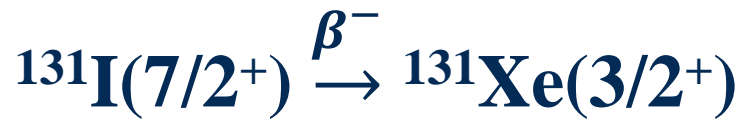
A. de Roubin et al., Phys. Rev. Lett. A **124**, 222503 (2020)

#### Conclusion:

- $^{135}\text{Cs}(7/2^+) \rightarrow ^{135}\text{Ba}(11/2^-) > 0$
- First direct Q value measurement
- Direct measurement of the antineutrino mass?

#### Expected challenges:

- Half-life of the  $11/2^-$  state too long
- Cannot be used in coincidence



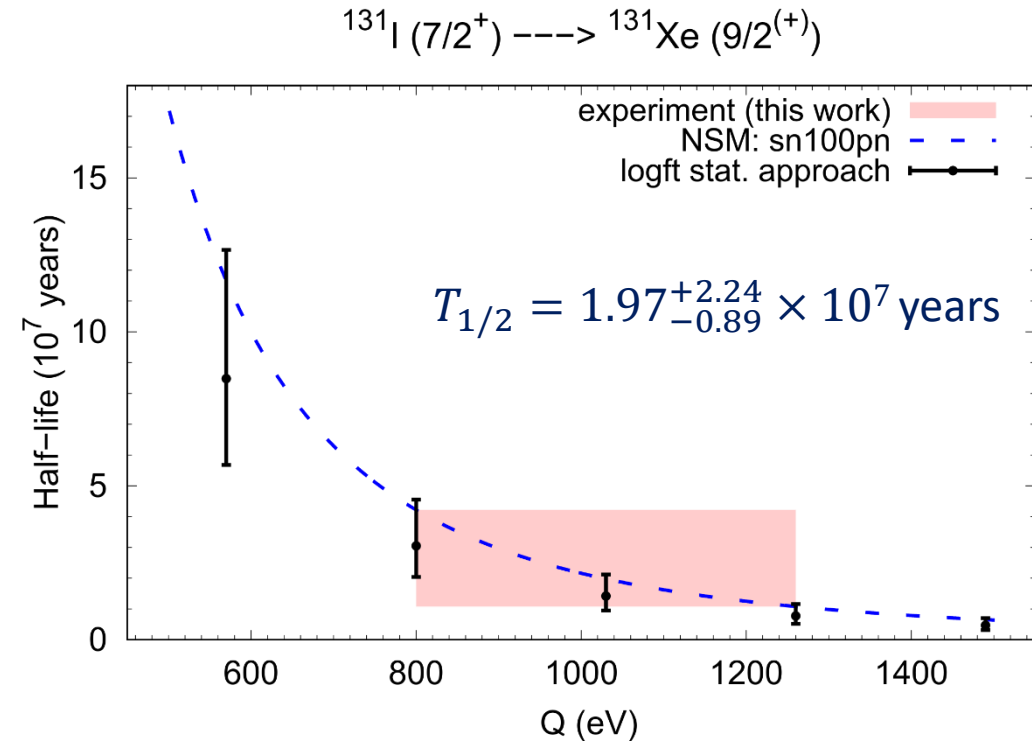
$$Q_{gs-gs} = 972.25(19) \text{ keV}$$

$$\rightarrow Q_{9/2^+} = 1.03(23) \text{ keV}$$

T. Eronen et al. accepted for publication in PLB

Half-lives calculations for this transition

- Nuclear shell model calculations with NuShellX  
*B. Brown, W. Rae, Nucl. Data Sheets 120, 115 (2014)*
- Partial half-life is extremely sensitive to the exact value of the decay energy

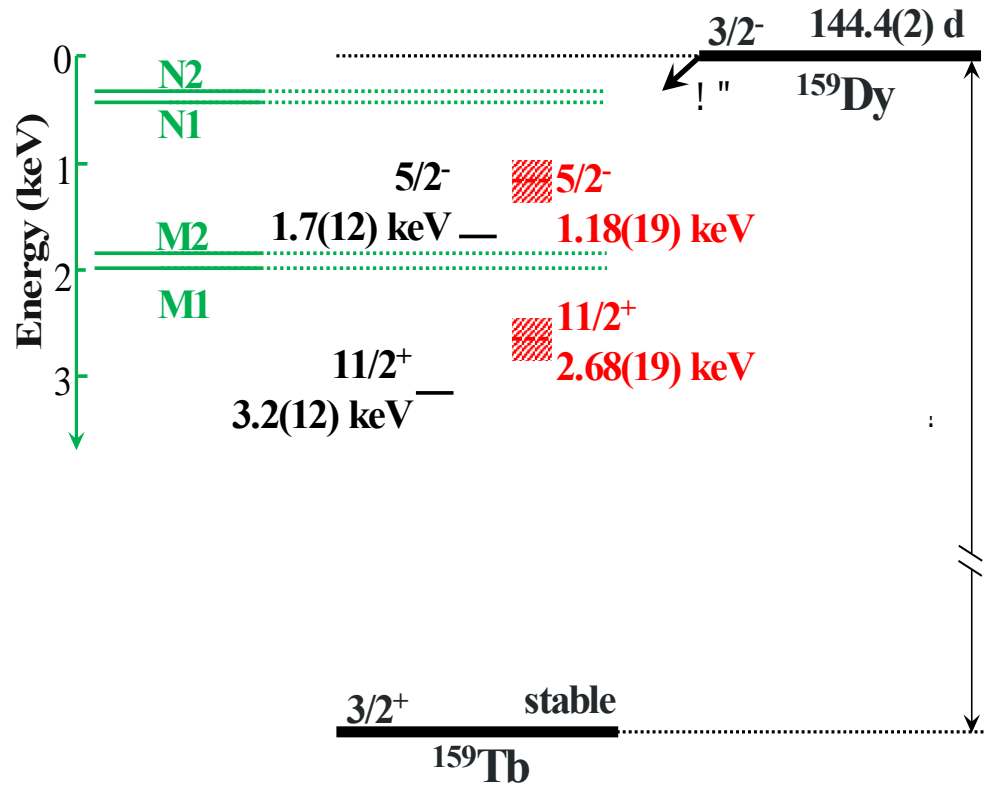


**Conclusion:**

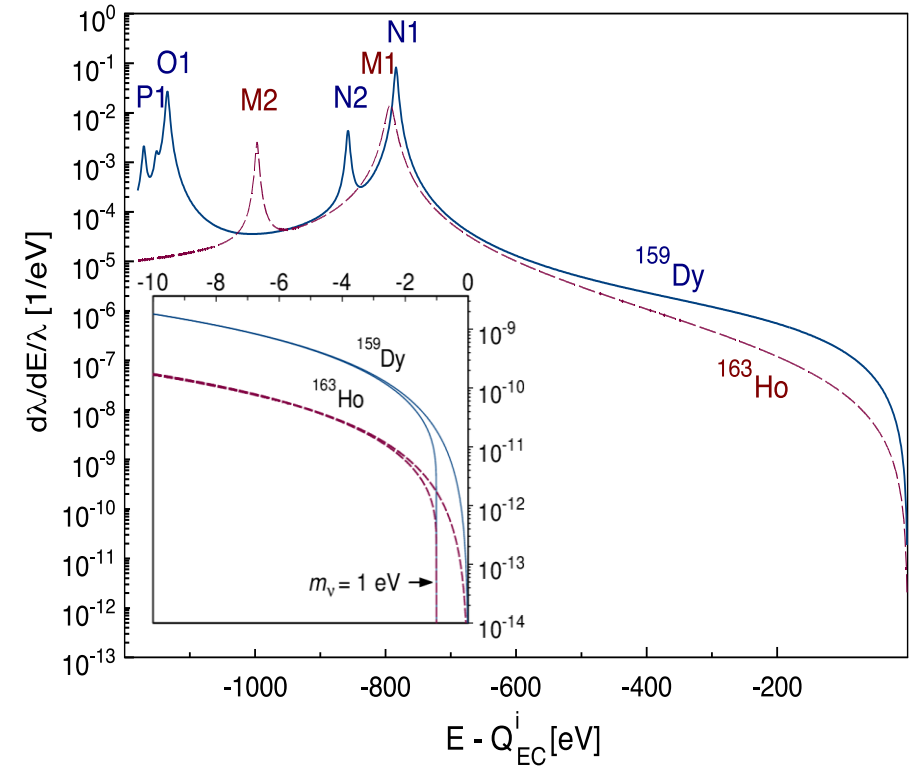
- The decay  $7/2^+ \rightarrow 9/2^+$  is of allowed type
- The transition has a simple universal shape

**Expected challenge:**

- The branching ratio  $\sim 10^{-10}$



EC spectrum of  $^{159}\text{Dy}(3/2^- \rightarrow 5/2^-)$  compared to  $^{163}\text{Ho}$  (Dirac-Hartree-Fock atomic many-body calculations)



Courtesy of Z. Ge

Z. Ge, T. Eronen, K. S. Tyrin et al., Phys. Rev. Lett. 127, 272301 (2021)

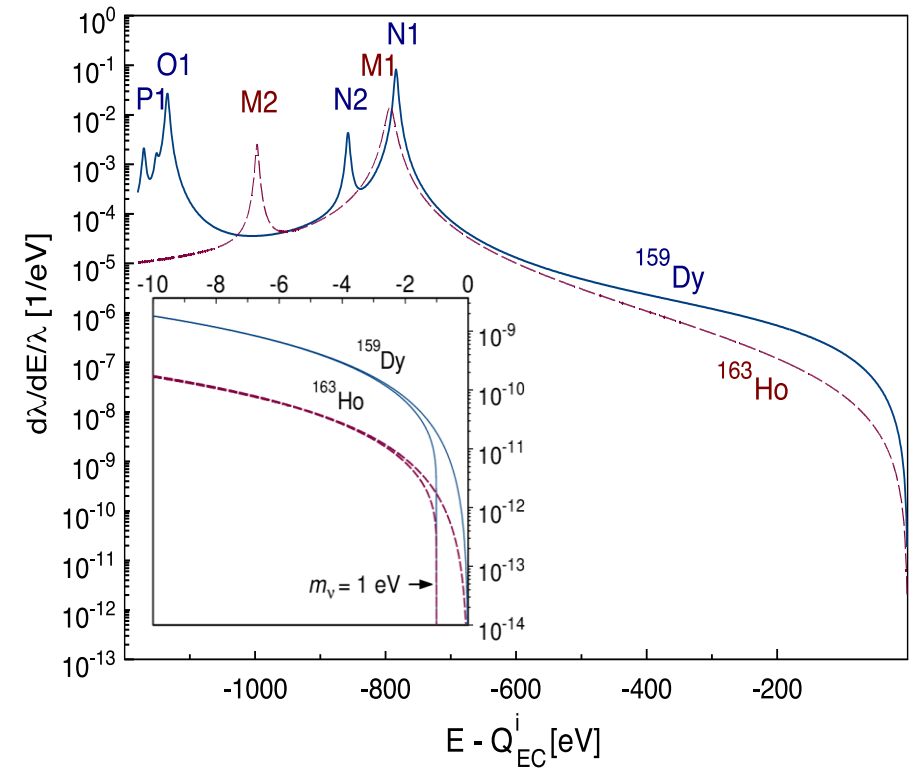




$Q_{EC}^* = 1.18(19)$  keV

- Lowest EC Q value
  - Lower than the GS-to-GS  $Q_{EC}$  of  $^{163}\text{Ho}$ , utilized in presently running or planned direct neutrino mass experiments
  - allowed transition
  - known branching ratio  $1.9(5) \times 10^{-6}$
  - Smallest EC Q value
- ⇒ The most promising gs-to-excited state transition for future calorimetric experiment

EC spectrum of  $^{159}\text{Dy}(3/2^- \rightarrow 5/2^-)$  compared to  $^{163}\text{Ho}$  (Dirac-Hartree-Fock atomic many-body calculations)



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Z. Ge, T. Eronen, K. S. Tyrin et al., *Phys. Rev. Lett.* 127, 272301 (2021)



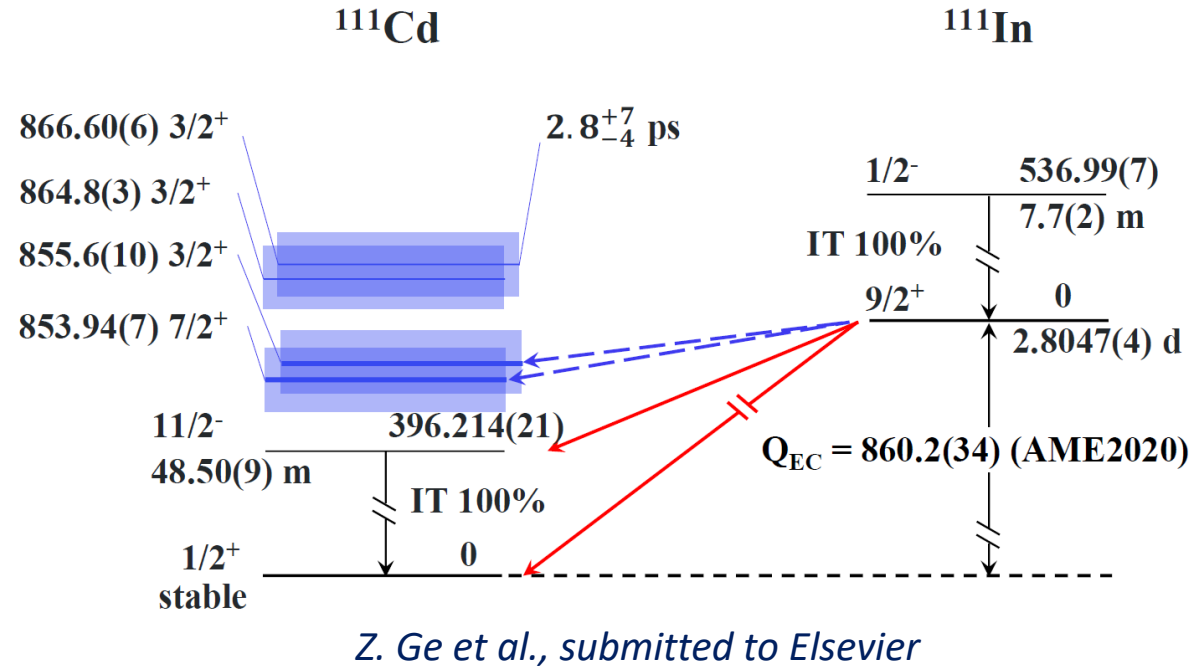
$$Q_{gs-gs} = 857.63(17) \text{ keV}$$

Mother ( $J^\pi$ ) Ground state	Daughter ( $J^\pi$ ) Excited state	$E^*$ (keV)	$Q$ (keV)	Decay type
$^{111}\text{In}(9/2^+)$	$^{111}\text{Cd}(7/2^+)$	853.94(7)	3.69(19)	allowed
	$^{111}\text{Cd}(3/2^+)$	855.6(10)	2.0(10)	2 <sup>nd</sup> FU
	$^{111}\text{Cd}(3/2^+)$	864.8(3)	-7.17(35)	
	$^{111}\text{Cd}(3/2^+)$	866.60(6)	-8.97(18)	

Estimated half-life for  $9/2^+ \rightarrow 3/2^+$  transition  $t_{1/2} = 5.6 \times 10^{17}$

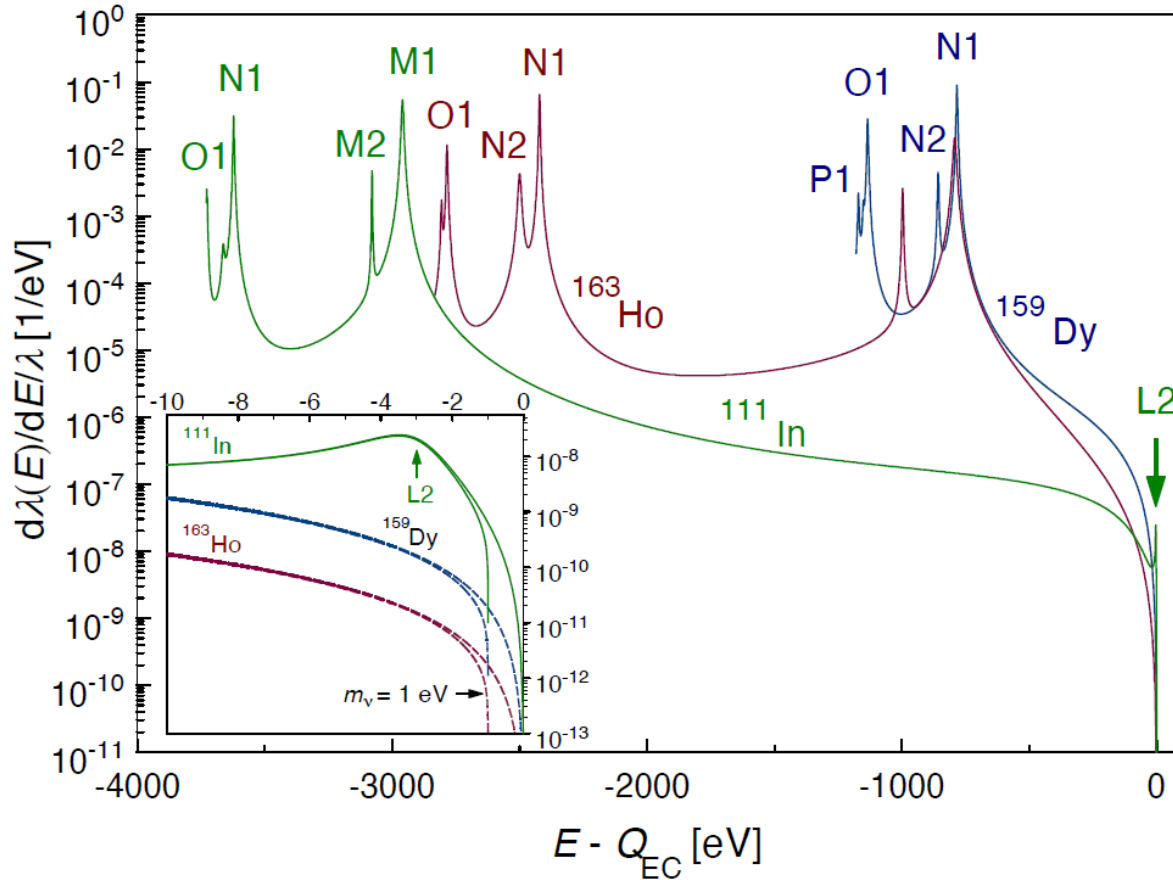
➤ Transition excluded!

Z. Ge et al., arXiv:2201.12573v1





EC spectrum of  $^{111}\text{In}(9/2^+ \rightarrow 5/2^-)$  compared to  $^{163}\text{Ho}$  and  $^{159}\text{Dy}$



- EC spectrum shape investigated with atomic Dirac-HF many-body calculations
- Presence of a L2 resonance near the end point (3 eV)
- Increases significantly the counting rate in the ROI

$Q = 2.0(10) \text{ keV}$

**Need more precise excitation energy measurement !**

Z. Ge et al., arXiv:2201.12573v1

# Summary of the measured Q-values of potential candidates

- List of measured promising low Q-value decay candidates for neutrino mass determination

Parent	T1/2	Daughter	E* (keV)	decay type	Q* (keV)	Decay	Q <sub>0</sub> (keV)	dQ <sub>0</sub> (keV)	
146Pm(3-)	5.53(5) y	146Nd(2+)	1470.63(6)	1st FNU	1.3(4.2)	EC	1472.000	4.000	
149Gd(7/2-)	9.28(10) dy	149Eu(5/2+)	1312(4)	1st FNU	2(6.4)	EC	1314.100	4.000	
155Tb(3/2+)	5.32(6) dy	155Gd{3/2+}	815.731(3)	Allowed{?}	4.2(10.1)	EC	820.000	10.000	
159Dy(3/2-)	144.4(2) dy	159Tb(5/2-)	363.5449(14)	Allowed	1.7(1.2)	EC	365.200	1.200	
Z. Ge, T. Eronen et al., Phys. Rev. Lett. 127, 272301		159Tb(11/2+)	362.050(40)	3rd FU	3.2(1.2)	EC	365.200	1.200	
		161Ho(5/2-)	18.479(4) hr	161Dy{7/2+}	858.502(7)	1st FNU	1.0(2.2)	EC	858.500
		161Dy{3/2-}	858.7919(18)	Allowed	-0.3(2.2)	EC	858.500	2.200	
Z. Ge, T. Eronen et al., PHYSICAL REVIEW C 103, 065502 (2021)	72As(2-)	26.0(1)h	72Ge{1}	4358.7(3)	Allowed{?}	-2.8(4.0)	EC	4356.000	4.000
			72Ge(3-)	3325.01(3)	Allowed	8.9(4.0)	β+	4356.000	4.000
			72Ge(2+)	3327(3)	1st FNU	6.9(5.0)	β+	4356.000	4.000
			72Ge{1+}	3338.0(3)	1st FNU{?}	-4.1(4.0)	β+	4356.000	4.000
			72Ge{2-}	3341.76(4)	Allowed{?}	-7.9(4.0)	β+	4356.000	4.000
159Gd(3/2-)	26.24(9) h	159Tb{1/2+}	971	1st FNU{?}	0.0(1.8)	β-	970.900	0.800	
77As(3/2-)	38.79(5) h	77Se(5/2+)	680.1035(17)	1st FNU	3.1(1.7)	β-	683.200	1.700	
76As(2-)	26.24(9) h	76Se{2-}	2968.4(7)	Allowed{?}	-7.8(1.1)	β-	2960.600	0.900	
153Tb(5/2+)	2.34(1)dy	153Gd(5/2-)	548.7645(18)	1st FNU	-1.2(4.0)	β+	1569.000	4.000	
		153Gd{5/2}	551.092(19)	Allowed{?}	-3.5(4.0)	β+	1569.000	4.000	
Submitted	3dy	111In(9/2+)	864.8(3)	2nd FU	-6.6(3.0)	EC	860.2	3.4	
		111Cd(3/2+)	864.8(3)	2nd FU	-4.6(3.0)	EC	860.2	3.4	
		111Cd(3/2+)	855.6(1.0)	2nd FU	4.6(3.2)	EC	860.2	3.4	
		111Cd(7/2+)	853.94(7)	Allowed	6.3(3.0)	EC	860.2	3.4	
131I(7/2+)	8dy	131Xe{9/2+}	971.22(13)	Allowed{?}	-0.42(0.61)	β-	970.80	0.60	
		131Xe(7/2+)	973.11(14)	Allowed	-2.31(0.62)	β-	970.80	0.60	
155Eu(5/2+)	5yr	155Gd(9/2-)	251.7056(10)	1st FU	0.1(1.8)	β-	252.00	2.40	

Q<sub>0</sub> from: M. Wang et al., Chinese Physics C 45, 030003 (2021)

E\* from: National nuclear data center, Available at <https://www.nndc.bnl.gov>

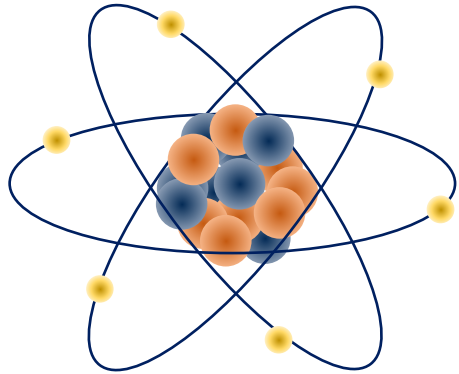
Courtesy of Z. Ge



Thanks a lot for your attention



# Binding energy



$$M_{Nuclei} = N \times M_{Neutron} + Z \times M_{Proton} - B(N, Z)/c^2$$

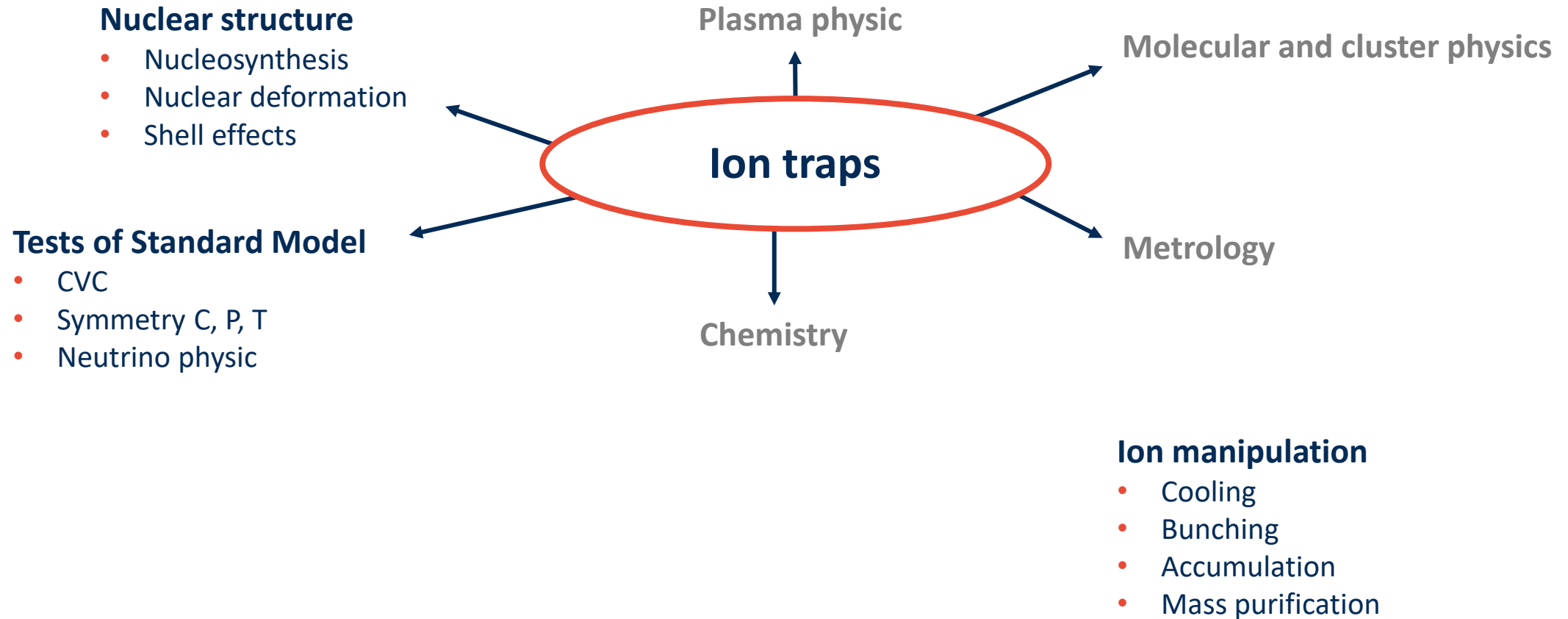
↳  $E = mc^2 \rightarrow$  The binding energy corresponds to the missing mass

## Measuring $B(N, Z)$

- $\delta M/M = 10^{-7} - 10^{-9}$  (For  $A = 100 u \rightarrow 100 GeV$ )
- Nuclear deformation  $\rightarrow \delta M/M = 10^{-7} \rightarrow 10 keV$
- Neutrino physic  $\rightarrow \delta M/M = 10^{-9} \rightarrow 100 eV$



# Ion traps





# The ECHo experiment

Electron capture of  $^{163}\text{Ho}$  with  $Q_{EC} = 2.8 \text{ keV}$

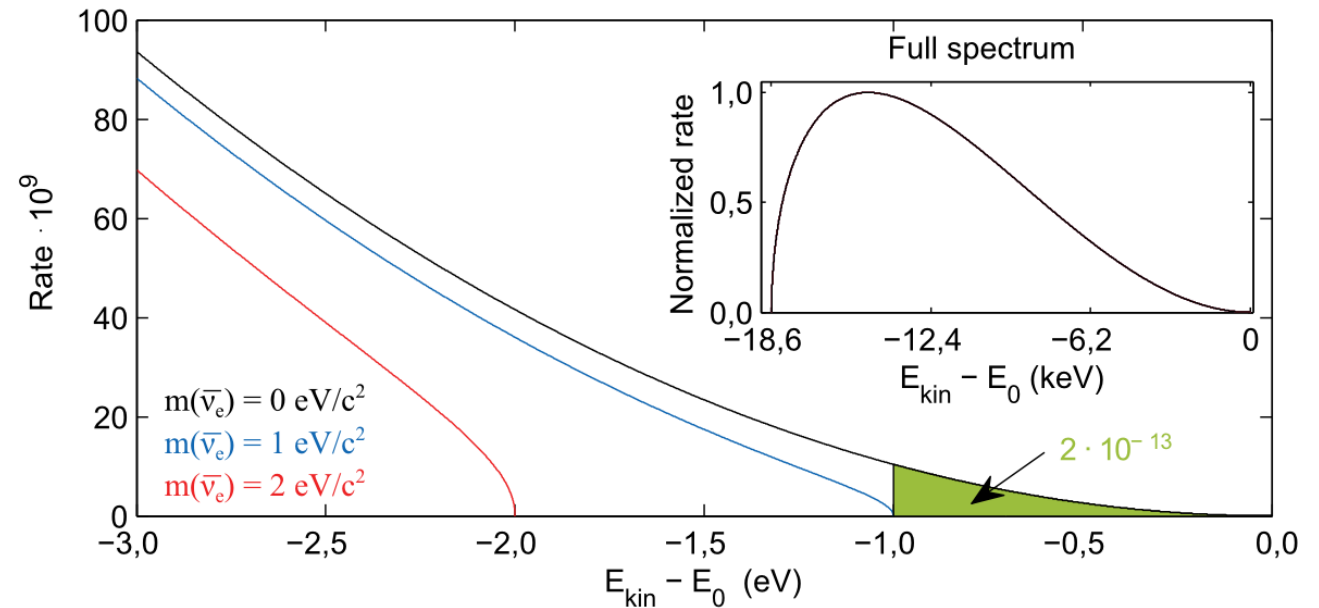
$^{163}\text{Ho}$  source is enclosed in large array of magnetic calorimeters

- Deposition of energy
- Temperature increased
- Change the magnetization of a paramagnetic temperature sensor
- Change the magnetic flux of a voltage convertor
- The output voltage signal is proportional to the deposited energy

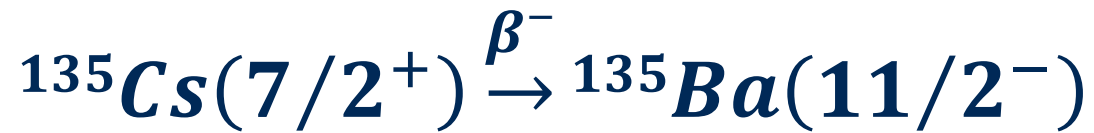
# The KATRIN experiment

- $\beta$  decay via the mass of the electron antineutrino
  - A slight distortion of the  $\beta$  decay spectrum relates directly to the electron-neutrino mass
- KATRIN [1] experiment
  - ${}^3\text{H}(1/2^+) \xrightarrow{\beta^-} {}^3\text{He}(1/2^+)$
  - Allowed decay
  - $Q = 18.5920(3) \text{ keV}$
- Low  $Q \rightarrow$  higher sensitivity

## Tritium $\beta$ decay electron spectrum



[1] C. Weinheimer, *Prog. Part. Nucl. Phys.* **48**, 141 (2002)



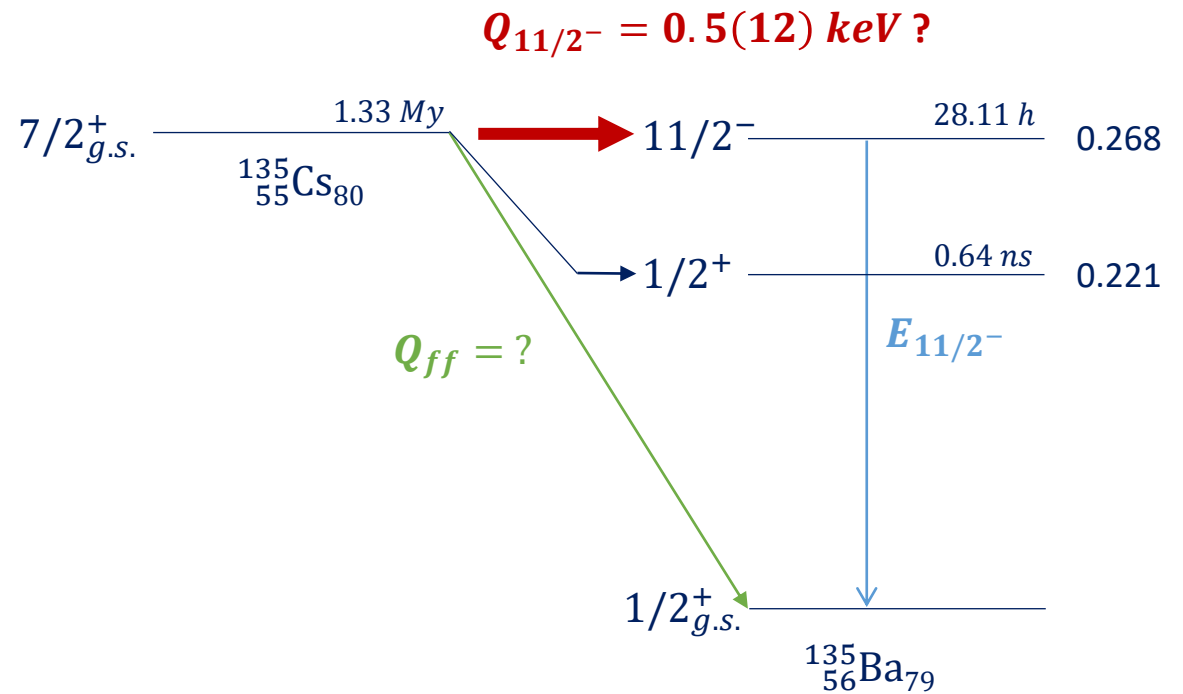
**But :** The transition is poorly known!

- Not yet proven that it is energetically possible :  $Q_{11/2^-} < 0$  ?
- Decay never observed directly

To determine  $Q_{11/2^-}$  we need:

- The excitation energy  $E_{11/2^-} = 268.218(20) \text{ keV}$
- The ground-state-to-ground-state Q-value  $Q_{ff}$
- $Q_{ff} = 268.9(1.1) \text{ keV}$  [1]

$$Q_{ff} = 268.9(1.1) \text{ keV} \rightarrow Q_{11/2^-} = 0.5(1.2) \text{ keV}$$



**New measurement of  $Q_{ff}$  with Penning traps!!**

[1] M. Wang et al., Chin. Phys. C 41, 030003 (2017)

# PI-ICR measurements

- \ Pure sample of  $^{135}\text{Cs}$  sent in the measurement trap
- \ Every few rounds, switched between  $^{135}\text{Cs}$  from IGISOL and  $^{135}\text{Ba}$  from the off-line spark source
- \ Measurement of the GS-to-GS Q-value performed with
  - TOF-ICR Ramsey technique
  - PI-ICR technique

