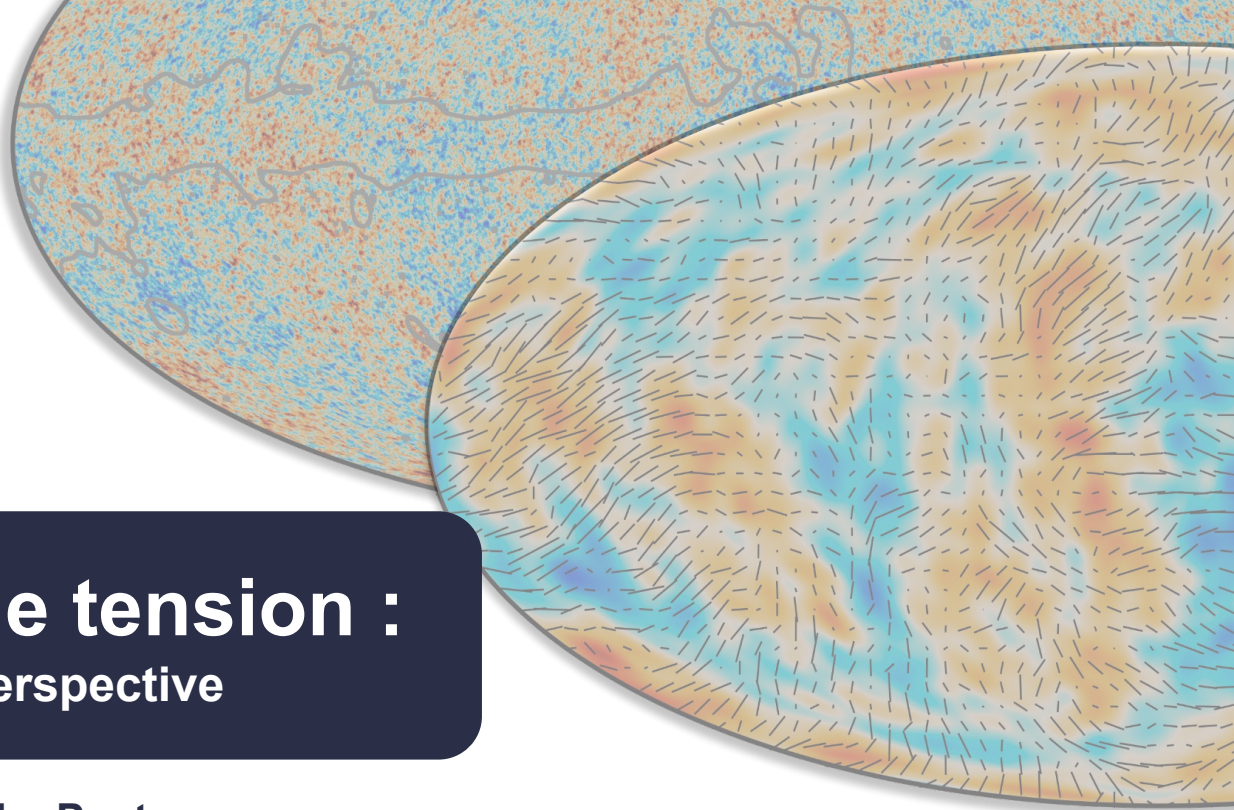


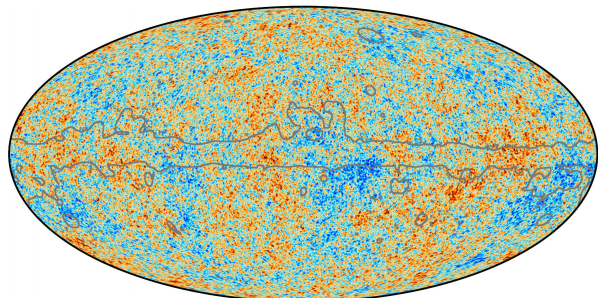
The Hubble tension : a CMB perspective

Adrien La Posta
IJClab



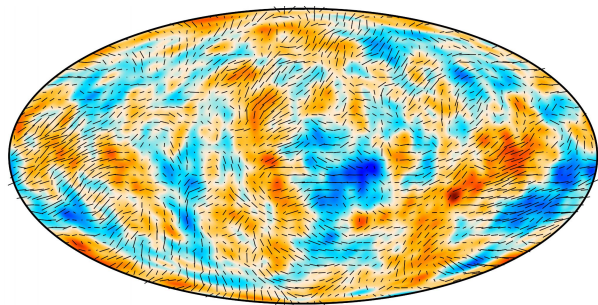
Measuring H_0 from the CMB

Temperature

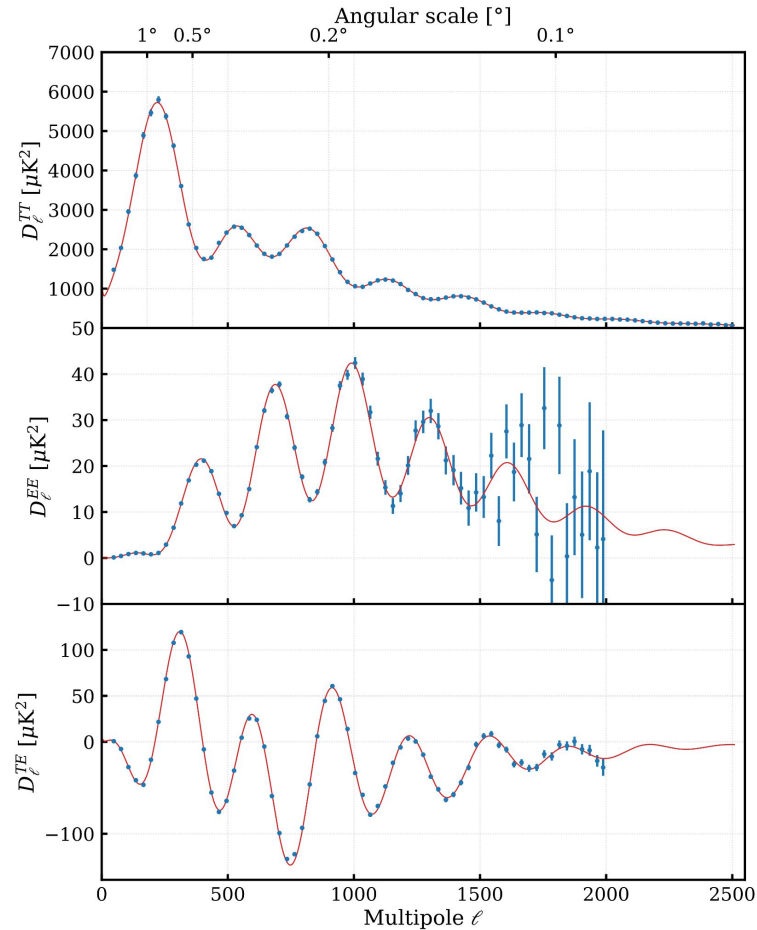


-300 300 μK

Polarization E-modes

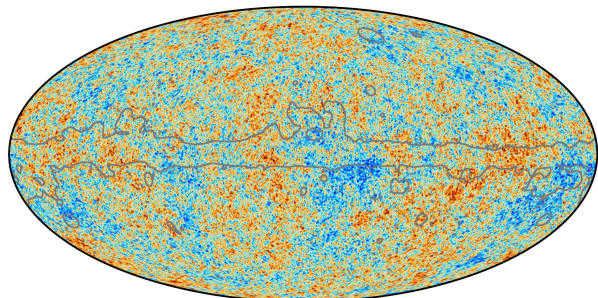


-160 160 μK



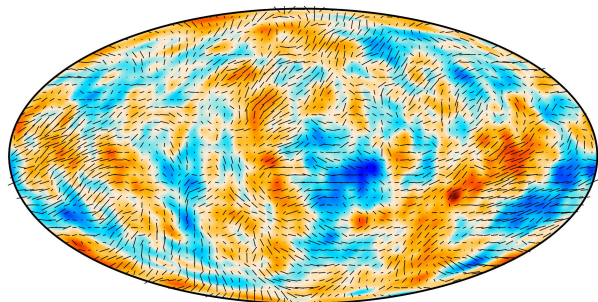
Measuring H_0 from the CMB

Temperature

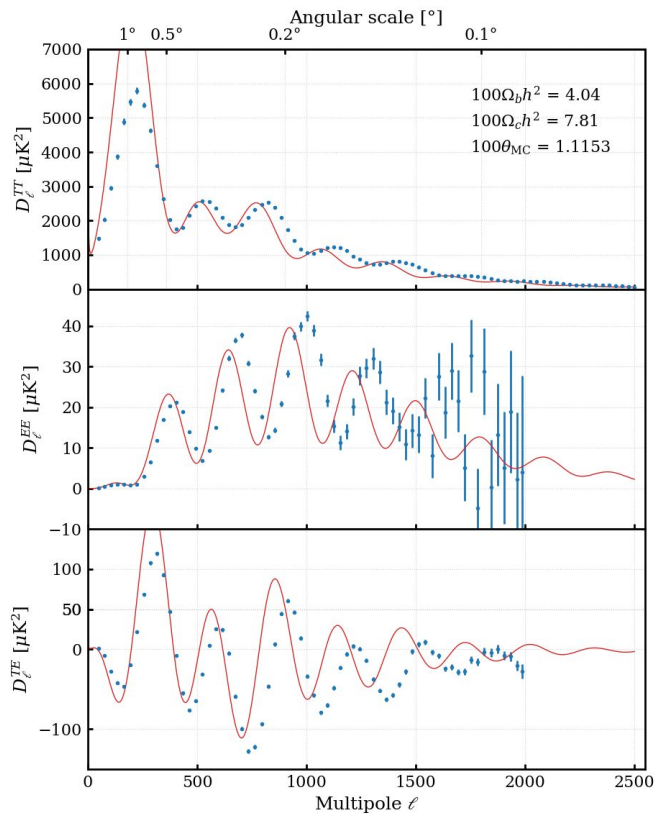


-300 300 μK

Polarization E-modes



-160 160 μK

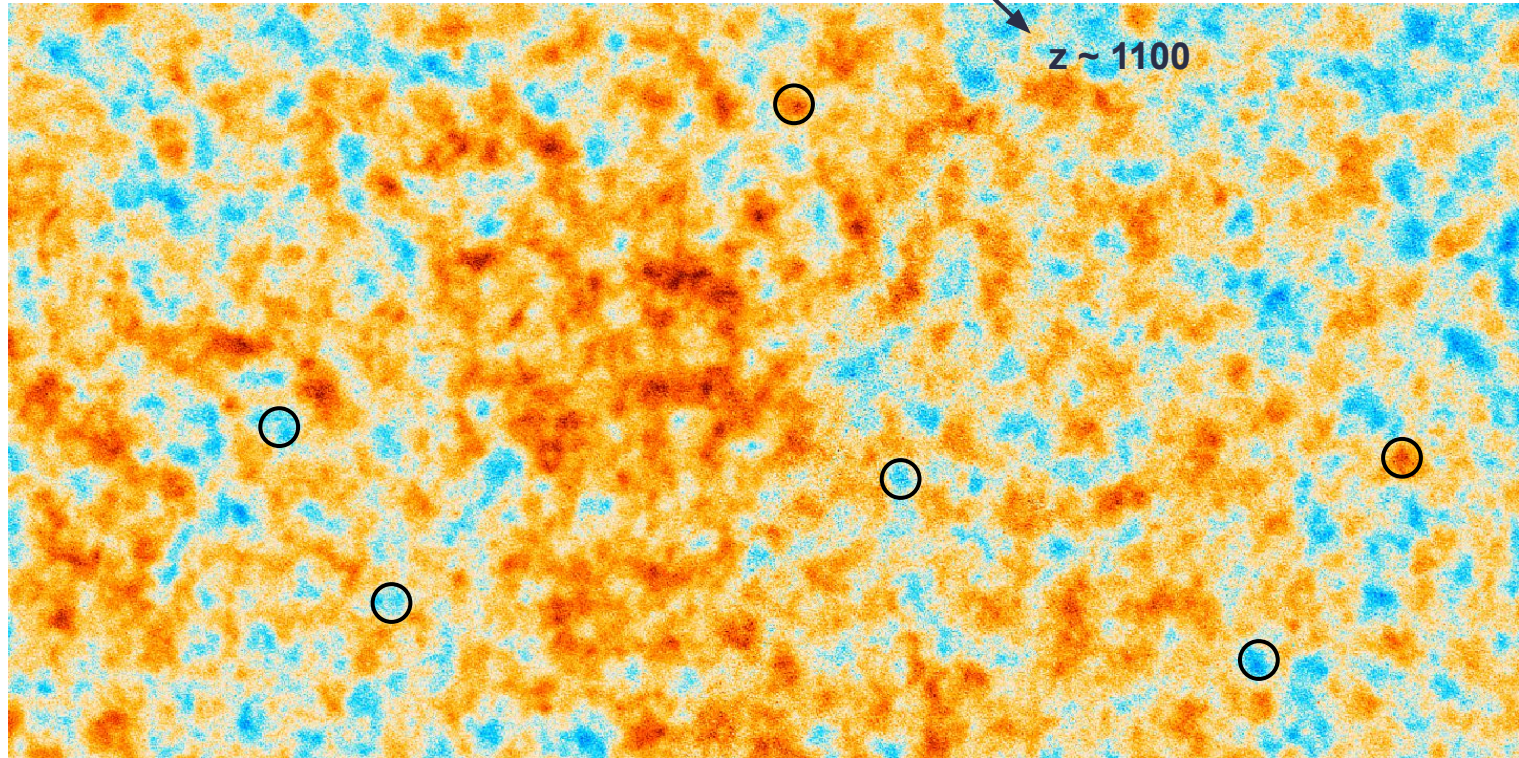


$\rightarrow \theta_* \rho_b^0 \rho_c^0$

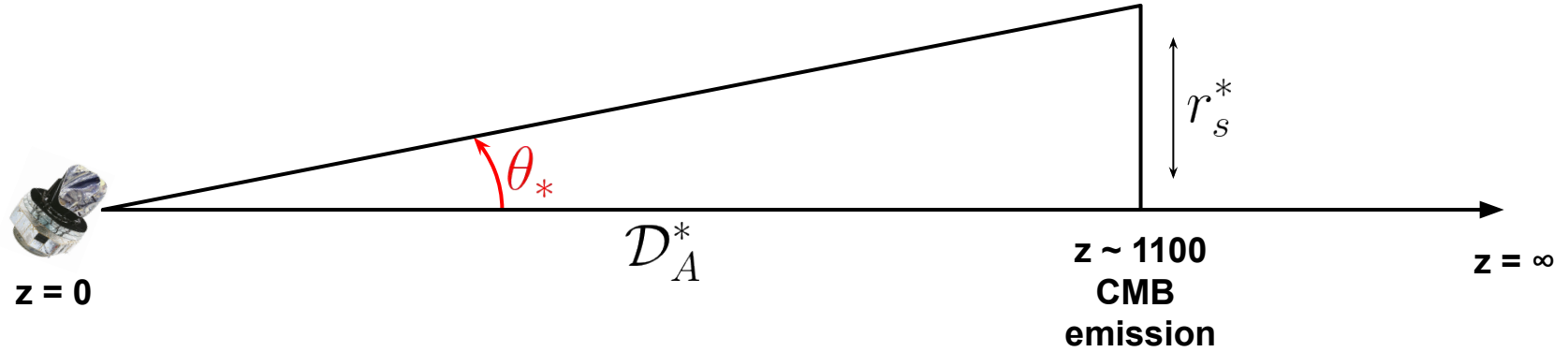
CMB standard ruler : size of the sound horizon at decoupling imprinted in the CMB radiation

Measuring H_0 from the CMB

CMB standard ruler : size of the sound horizon at decoupling imprinted in the CMB radiation

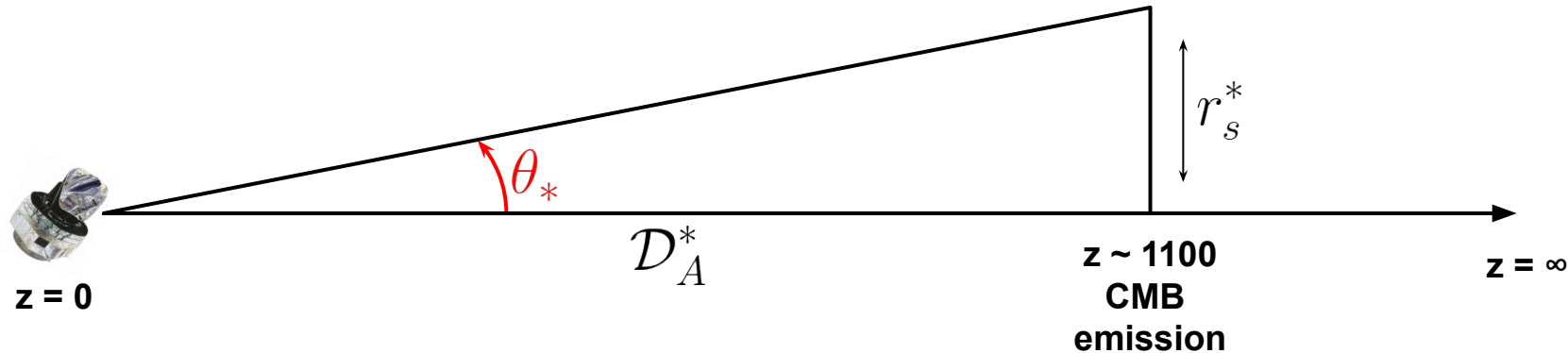


Measuring H_0 from the CMB



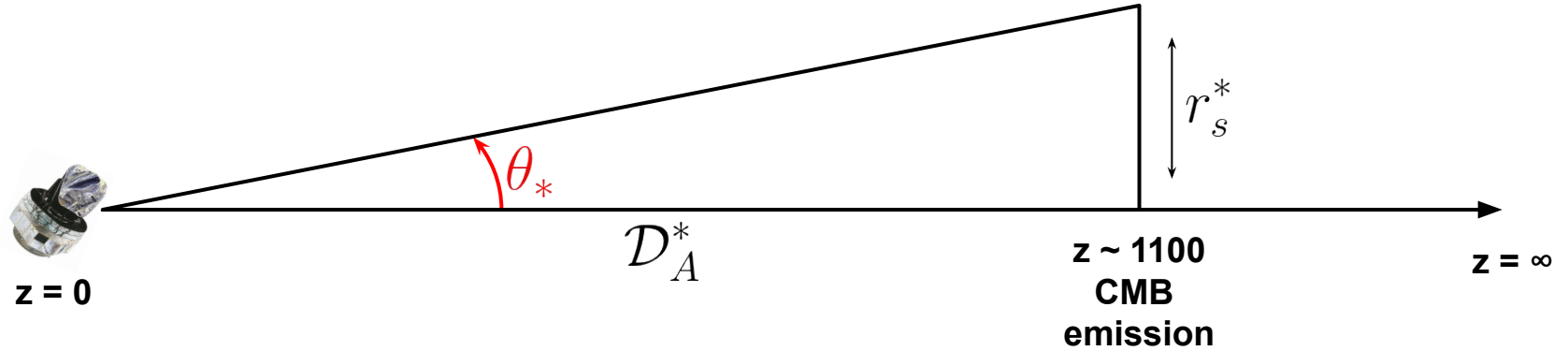
$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*}$$

Measuring H_0 from the CMB



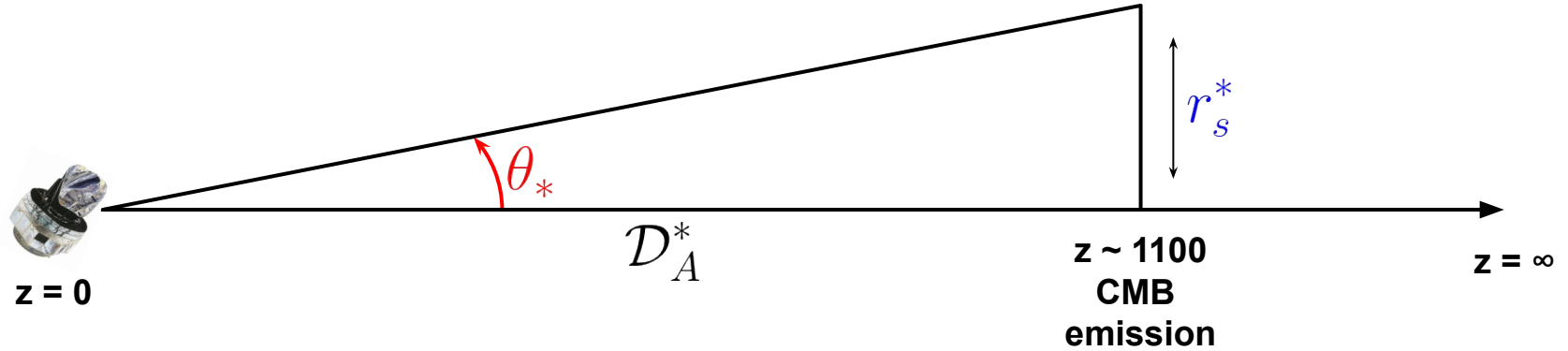
$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \rightarrow r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

Measuring H_0 from the CMB



$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*} \rightarrow r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z) \rightarrow c_s(z) = c \sqrt{\frac{1}{3 [1 + 3\rho_b^0 / 4\rho_\gamma^0 (1+z)^{-1}]}}$$
$$H_{\text{early}}^2(z) = \frac{8\pi G}{3} [\rho_r^0 (1+z)^4 + (\rho_b^0 + \rho_c^0) (1+z)^3]$$

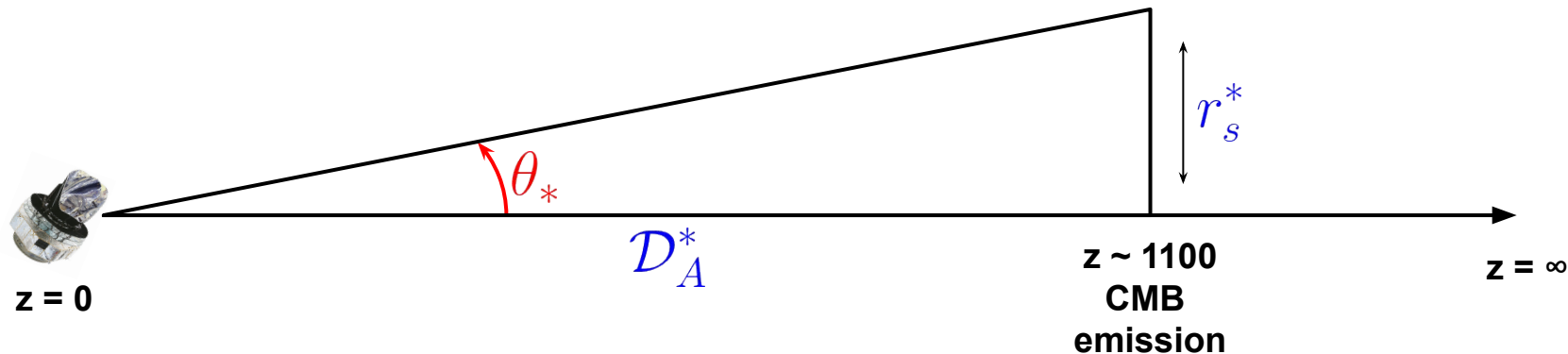
Measuring H_0 from the CMB



Now \mathcal{D}_A^* is known

$$\theta_* = \frac{r_s^*}{\mathcal{D}_A^*}$$

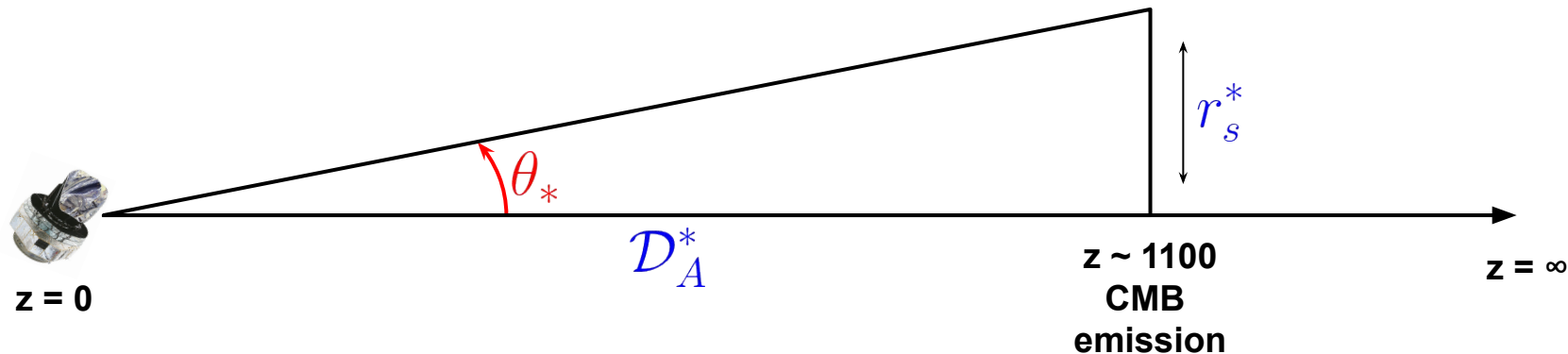
Measuring H_0 from the CMB



Now D_A^* is known

$$\theta_* = \frac{r_s^*}{D_A^*} \rightarrow D_A^* = c \int_0^{z^*} \frac{dz}{H(z)}$$
$$H_{\text{late}}^2(z) = \frac{8\pi G}{3} [(\rho_b^0 + \rho_c^0)(1+z)^3 + \rho_\Lambda]$$

Measuring H_0 from the CMB



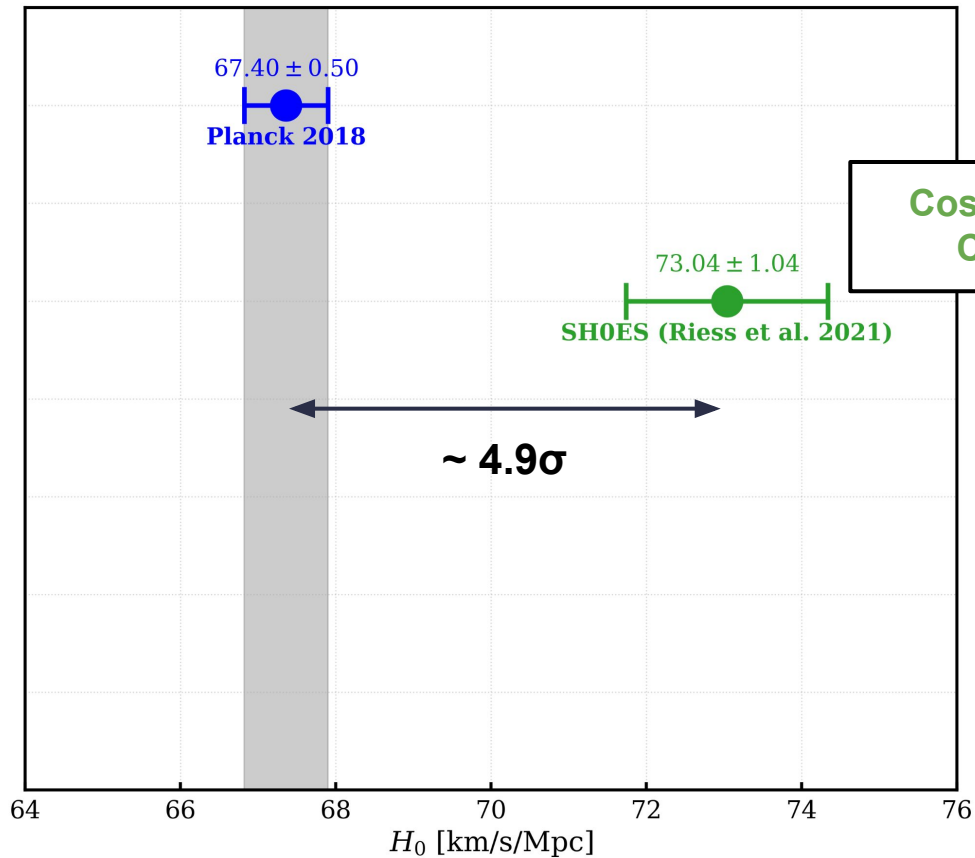
Now D_A^* is known

$$\theta_* = \frac{r_s^*}{D_A^*} \rightarrow D_A^* = c \int_0^{z^*} \frac{dz}{H(z)}$$

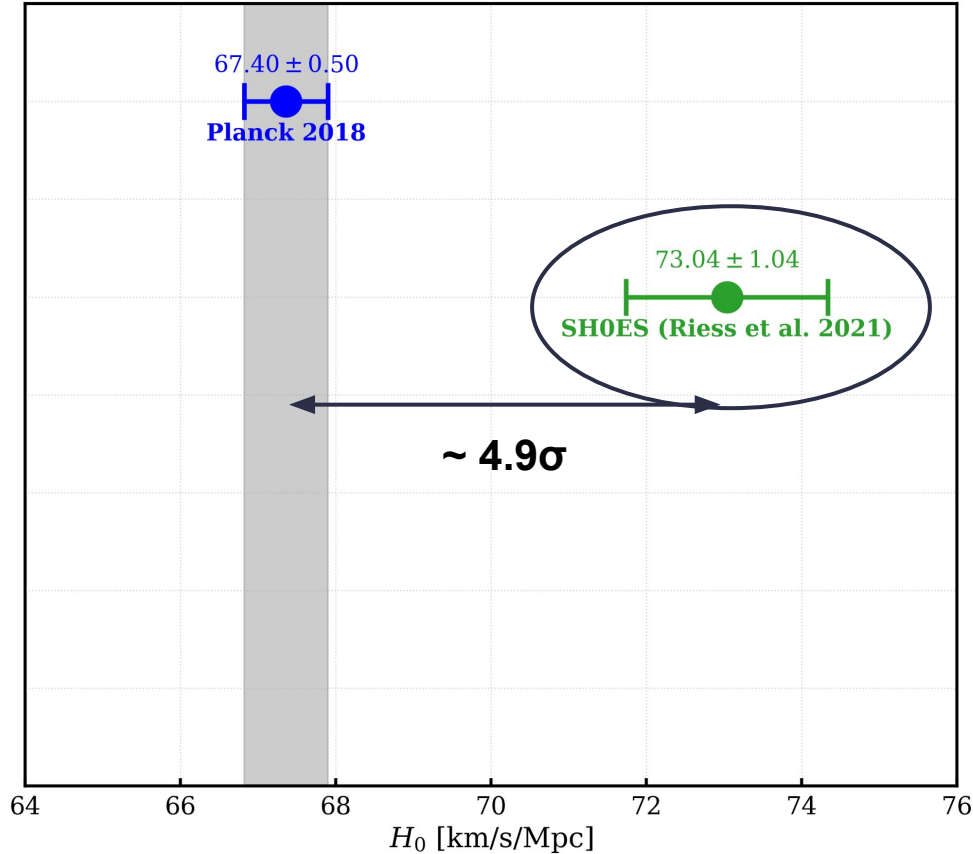
$$H_0^2 = \frac{8\pi G}{3} [\rho_b^0 + \rho_c^0 + \rho_\Lambda]$$

$$H_{\text{late}}^2(z) = \frac{8\pi G}{3} [(\rho_b^0 + \rho_c^0)(1+z)^3 + \rho_\Lambda]$$

The Hubble tension as of today

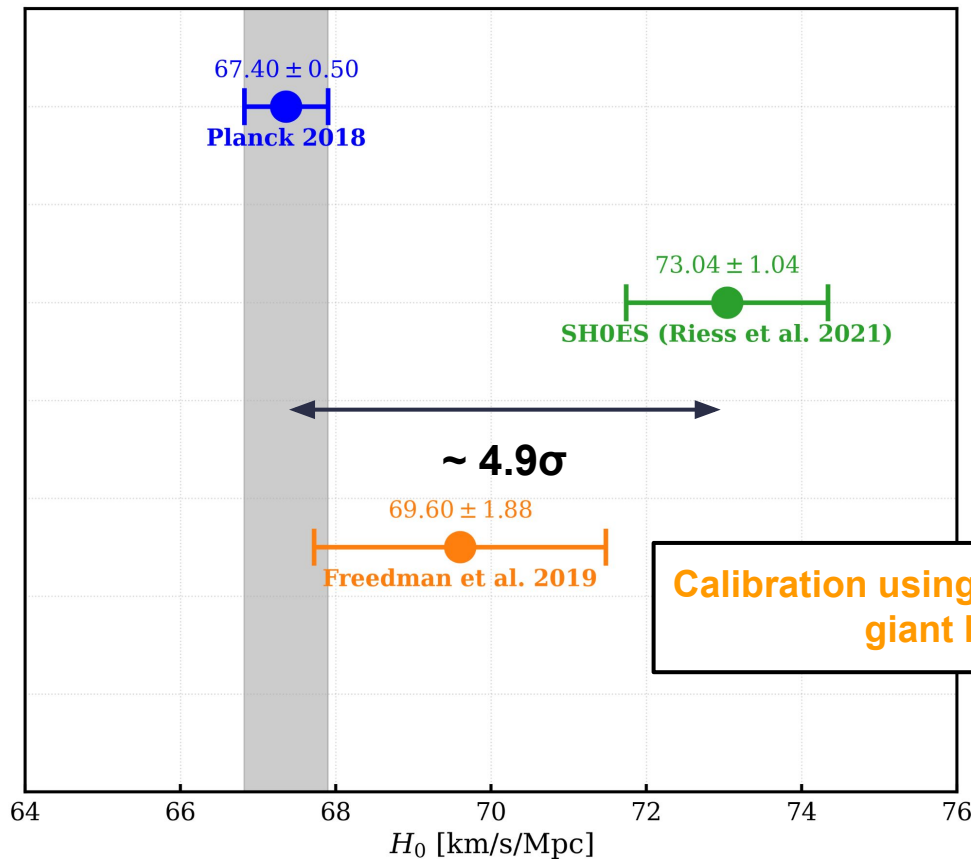


Cosmic distance ladder using
Cepheid calibrated SNIa



Option 1

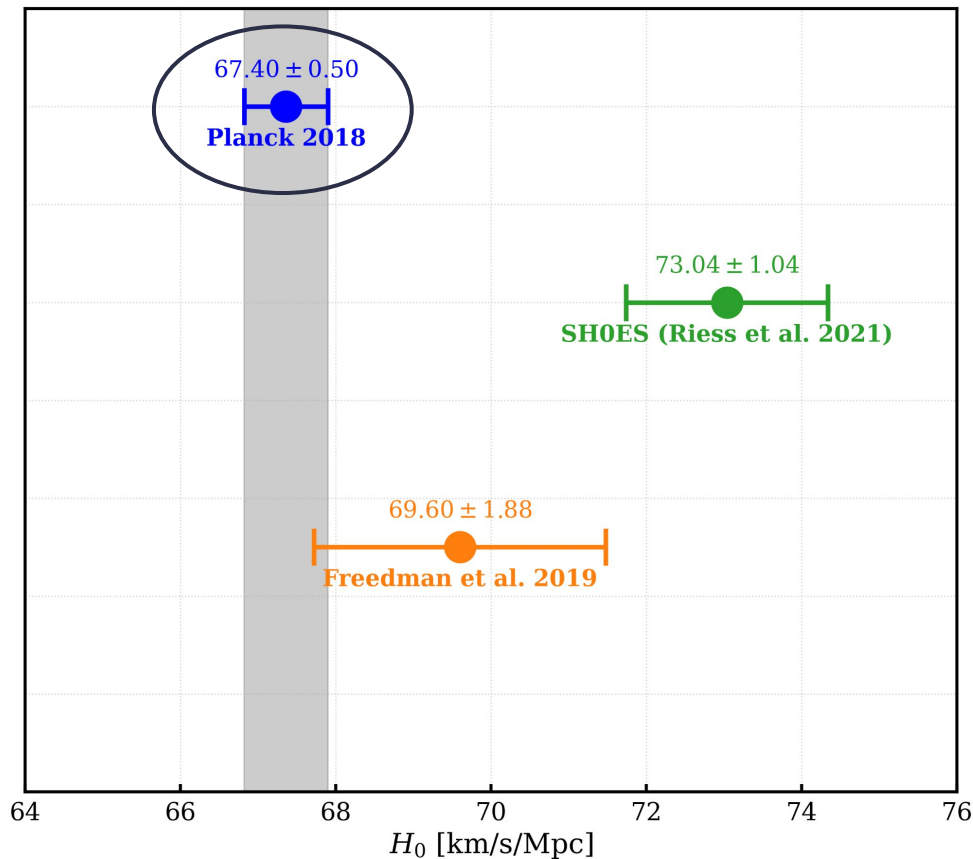
Astrophysical biases affecting the local measurement of H_0



Option 1

Astrophysical biases affecting the local measurement of H_0

Calibration using the tip of the red giant branch

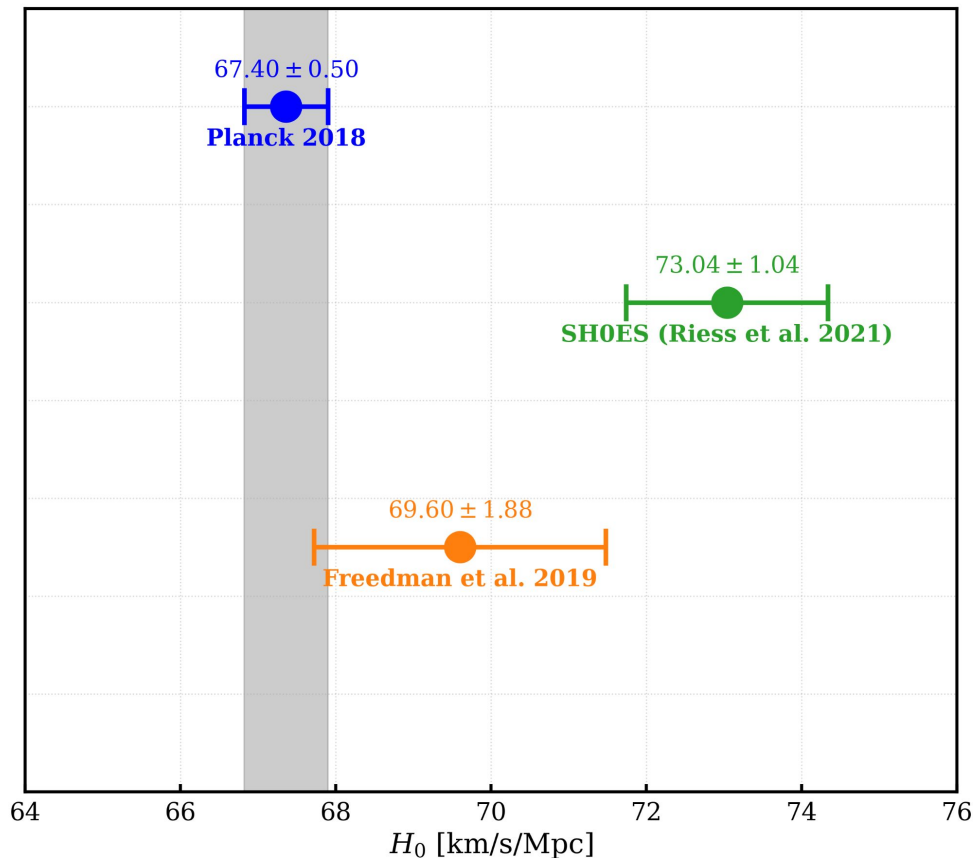


Option 1

Astrophysical biases affecting the local measurement of H_0

Option 2

Instrumental systematic effect biasing the value of H_0 inferred from the CMB



Option 1

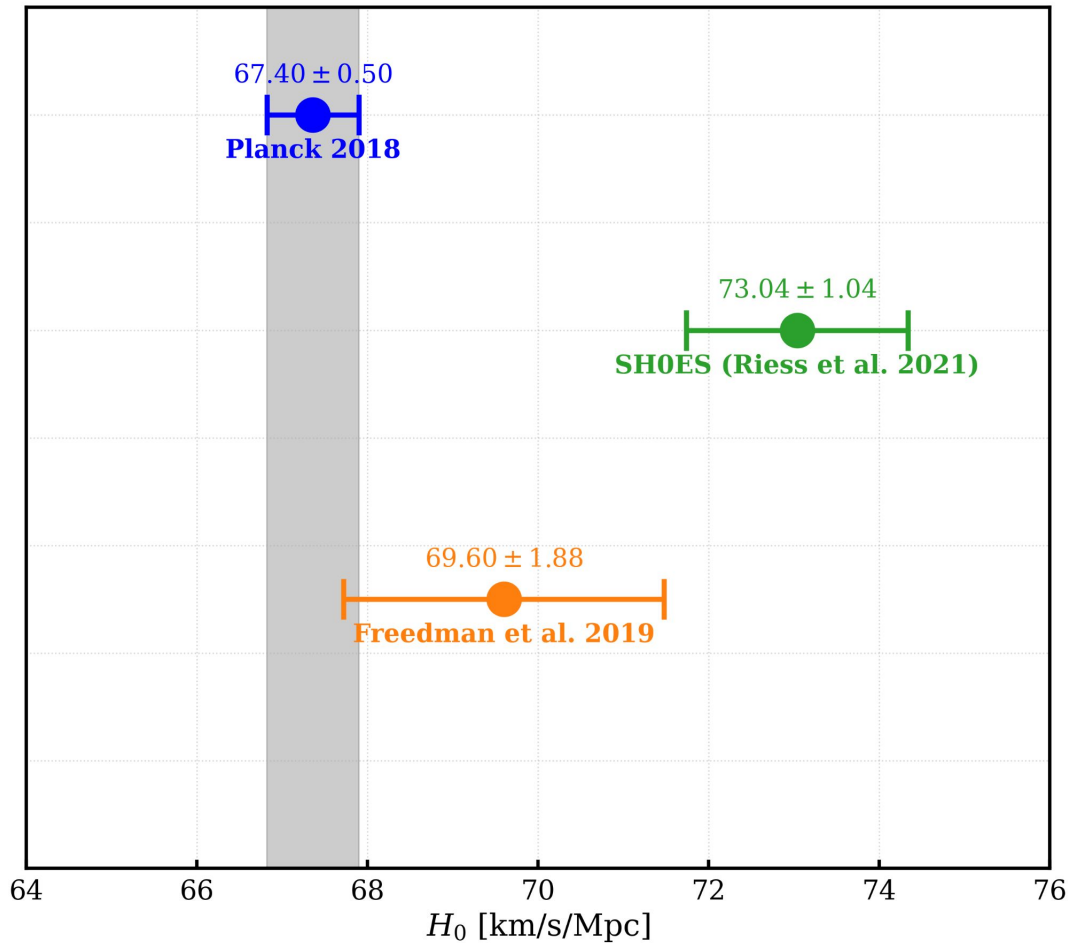
Astrophysical biases affecting the local measurement of H_0

Option 2

Instrumental systematic effect biasing the value of H_0 inferred from the CMB

Option 3

Physics beyond Λ CDM



Option 1

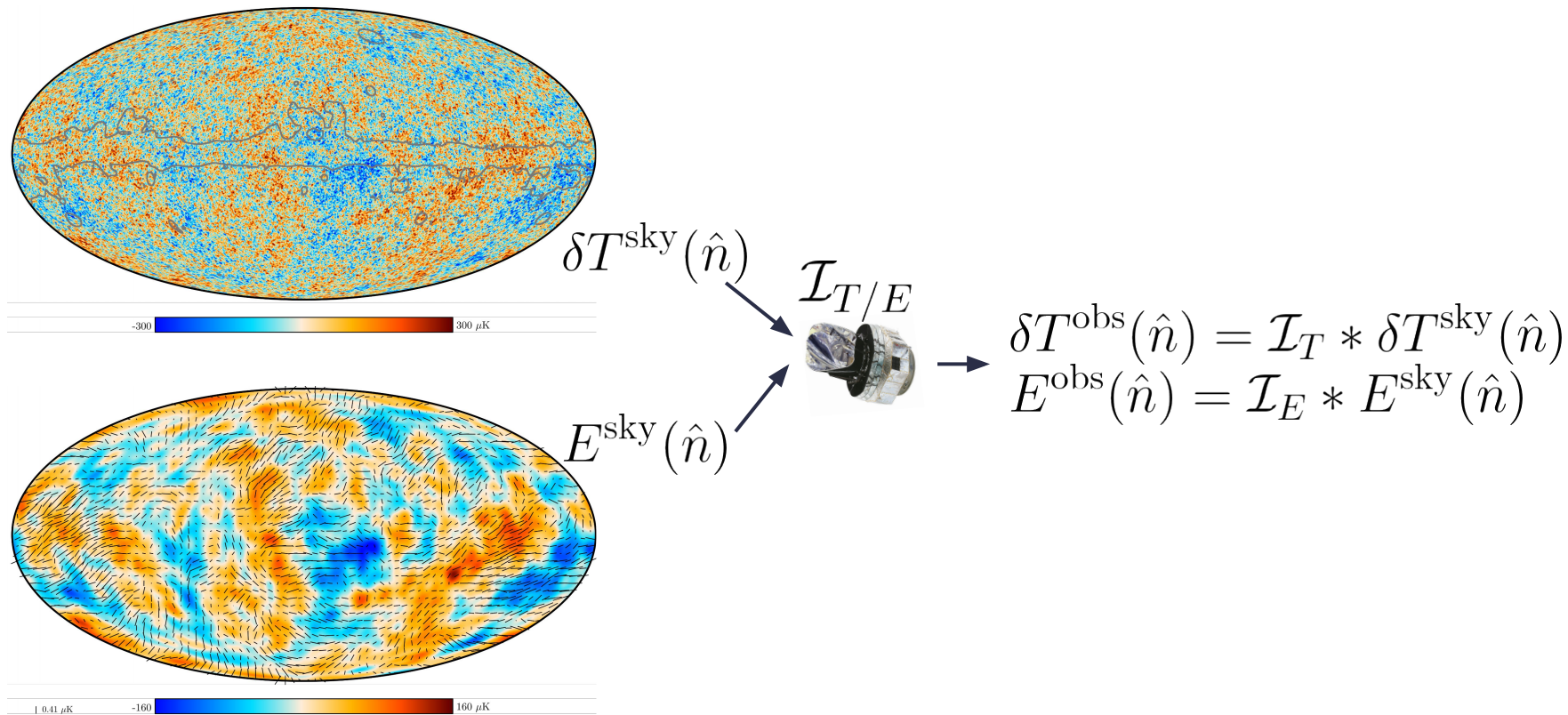
Astrophysical biases affecting
the local measurement of H_0

Option 2

Instrumental systematic effect
biasing the value of H_0
inferred from the CMB

Option 3

Physics beyond Λ CDM



$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

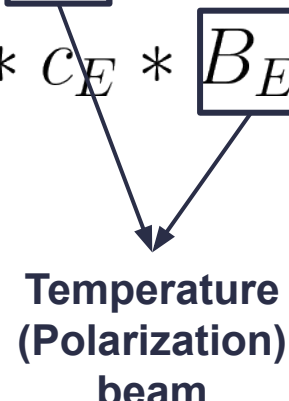
$$\mathcal{I}_T = \mathcal{F}_T * c * B_T$$

$$\mathcal{I}_E = \mathcal{F}_E * c * c_E * B_E$$

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)

$$\mathcal{I}_T = \mathcal{F}_T * c * B_T$$
$$\mathcal{I}_E = \mathcal{F}_E * c * c_E * B_E$$


Temperature
(Polarization)
beam

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration

$$\mathcal{I}_T = \mathcal{F}_T * \boxed{C} * B_T$$

$$\mathcal{I}_E = \mathcal{F}_E * \boxed{C} * c_E * B_E$$

Global
calibration

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration
- Polarization efficiency

$$\mathcal{I}_T = \mathcal{F}_T * c * B_T$$

$$\mathcal{I}_E = \mathcal{F}_E * c * \boxed{c_E} * B_E$$



**Polarization
efficiency**

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration
- Polarization efficiency
- Transfer functions (map-making)

$$\mathcal{I}_T = \boxed{\mathcal{F}_T} * c * B_T$$
$$\mathcal{I}_E = \boxed{\mathcal{F}_E} * c * c_E * B_E$$

Temperature
(polarization)
transfer function

$$\delta T^{\text{obs}}(\hat{n}) = \mathcal{I}_T * \delta T^{\text{sky}}(\hat{n})$$

$$E^{\text{obs}}(\hat{n}) = \mathcal{I}_E * E^{\text{sky}}(\hat{n})$$

- Finite angular resolution (beams)
- Calibration
- Polarization efficiency
- Transfer functions (map-making)

**These instrumental effects are
multiplicative in harmonic space**

$$C_{\ell}^{TT,\text{obs}} = (\mathcal{F}_{\ell}^T)^2 c^2 (B_{\ell}^T)^2 C_{\ell}^{TT}$$

$$C_{\ell}^{EE,\text{obs}} = (\mathcal{F}_{\ell}^E)^2 c^2 c_E^2 (B_{\ell}^E)^2 C_{\ell}^{EE}$$

$$C_{\ell}^{TE,\text{obs}} = \mathcal{F}_{\ell}^T \mathcal{F}_{\ell}^E c^2 c_E B_{\ell}^T B_{\ell}^E C_{\ell}^{EE}$$

$$\mathcal{R}_l^{TE} = \frac{\langle a_{lm}^T a_{lm}^{E*} \rangle}{\sqrt{\langle a_{lm}^T a_{lm}^{T*} \rangle \langle a_{lm}^E a_{lm}^{E*} \rangle}} = \frac{C_l^{TE}}{\sqrt{C_l^{TT} C_l^{EE}}}$$

$$\mathcal{R}_\ell^{TE} = \frac{\langle a_{\ell m}^T a_{\ell m}^{E*} \rangle}{\sqrt{\langle a_{\ell m}^T a_{\ell m}^{T*} \rangle \langle a_{\ell m}^E a_{\ell m}^{E*} \rangle}} = \frac{C_\ell^{TE}}{\sqrt{C_\ell^{TT} C_\ell^{EE}}}$$

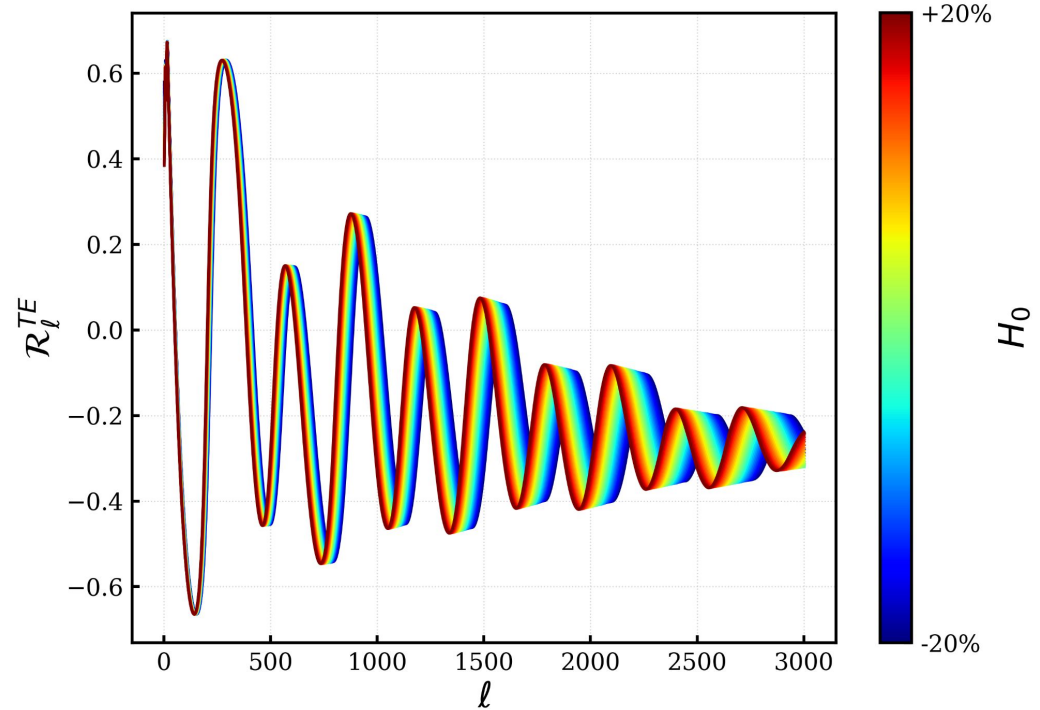
$$\mathcal{R}_\ell^{TE, \text{obs}} = \frac{\mathcal{F}_\ell^T \mathcal{F}_\ell^E c^2 c_E B_\ell^T B_\ell^E C_\ell^{TE}}{\sqrt{(\mathcal{F}_\ell^T)^2 c^2 (B_\ell^T)^2 C_\ell^{TT} \times (\mathcal{F}_\ell^E)^2 c^2 c_E^2 (B_\ell^E)^2 C_\ell^{EE}}}$$

$$\mathcal{R}_l^{TE} = \frac{\langle a_{lm}^T a_{lm}^{E*} \rangle}{\sqrt{\langle a_{lm}^T a_{lm}^{T*} \rangle \langle a_{lm}^E a_{lm}^{E*} \rangle}} = \frac{C_l^{TE}}{\sqrt{C_l^{TT} C_l^{EE}}}$$

$$\mathcal{R}_l^{TE, \text{obs}} = \frac{\cancel{\mathcal{F}_l^T \mathcal{F}_l^E} \cancel{c^2 c_E} \cancel{B_l^T B_l^E} C_l^{TE}}{\sqrt{(\cancel{\mathcal{F}_l^T})^2 \cancel{c^2} (\cancel{B_l^T})^2 C_l^{TT} \times (\cancel{\mathcal{F}_l^E})^2 \cancel{c^2} \cancel{c_E^2} (\cancel{B_l^E})^2 C_l^{EE}}} = \mathcal{R}_l^{TE}$$

- **The correlation coefficient is an observable insensitive to multiplicative biases**
 - **unbiased constraints on cosmological parameters**

- The correlation coefficient is an observable insensitive to multiplicative biases
→ unbiased constraints on cosmological parameters
- Particularly sensitive to H_0



Gaussian likelihood based on a Planck multifrequency likelihood (HiLLiPoP)

3 frequencies : **100**, **143** and **217** GHz

$$\ln \mathcal{L} \simeq -\frac{1}{2} (\Delta \mathcal{R}^{\text{vec}})^{\text{T}} \mathbf{\Xi}^{-1} (\Delta \mathcal{R}^{\text{vec}})$$

Gaussian likelihood based on a Planck multifrequency likelihood (HiLLiPoP)

3 frequencies : **100**, **143** and **217** GHz

$$\ln \mathcal{L} \simeq -\frac{1}{2} \boxed{\Delta \mathcal{R}^{\text{vec}}}^T \mathbf{\Xi}^{-1} (\Delta \mathcal{R}^{\text{vec}})$$

$$\Delta \mathcal{R}_\ell^{TE, \nu_1 \times \nu_2} = \hat{\mathcal{R}}_\ell^{TE, \nu_1 \times \nu_2} - \mathcal{R}_\ell^{TE, \nu_1 \times \nu_2, \text{model}}$$

Gaussian likelihood based on a Planck multifrequency likelihood (HiLLiPoP)

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$$\ln \mathcal{L} \simeq -\frac{1}{2} (\Delta \mathcal{R}^{\text{vec}})^T \mathbf{\Xi}^{-1} (\Delta \mathcal{R}^{\text{vec}})$$

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**Unbiased
estimator (data)**

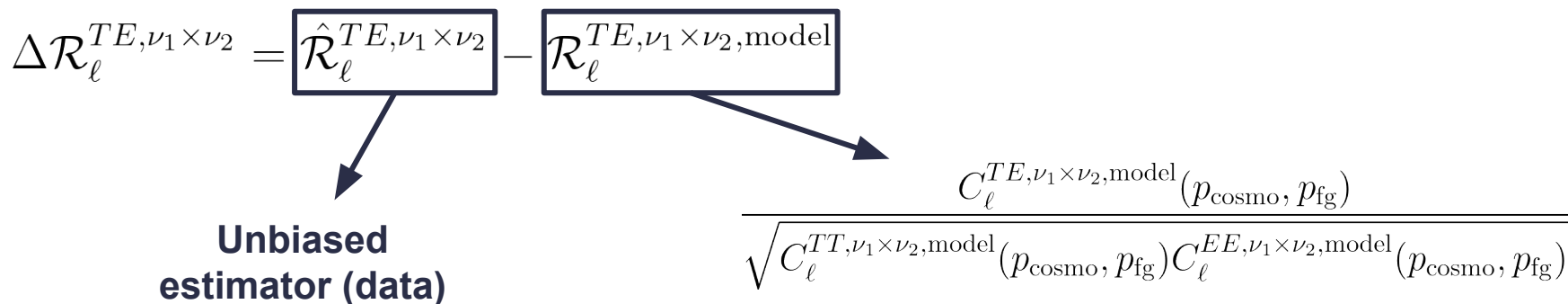
Gaussian likelihood based on a Planck multifrequency likelihood (HiLLiPoP)

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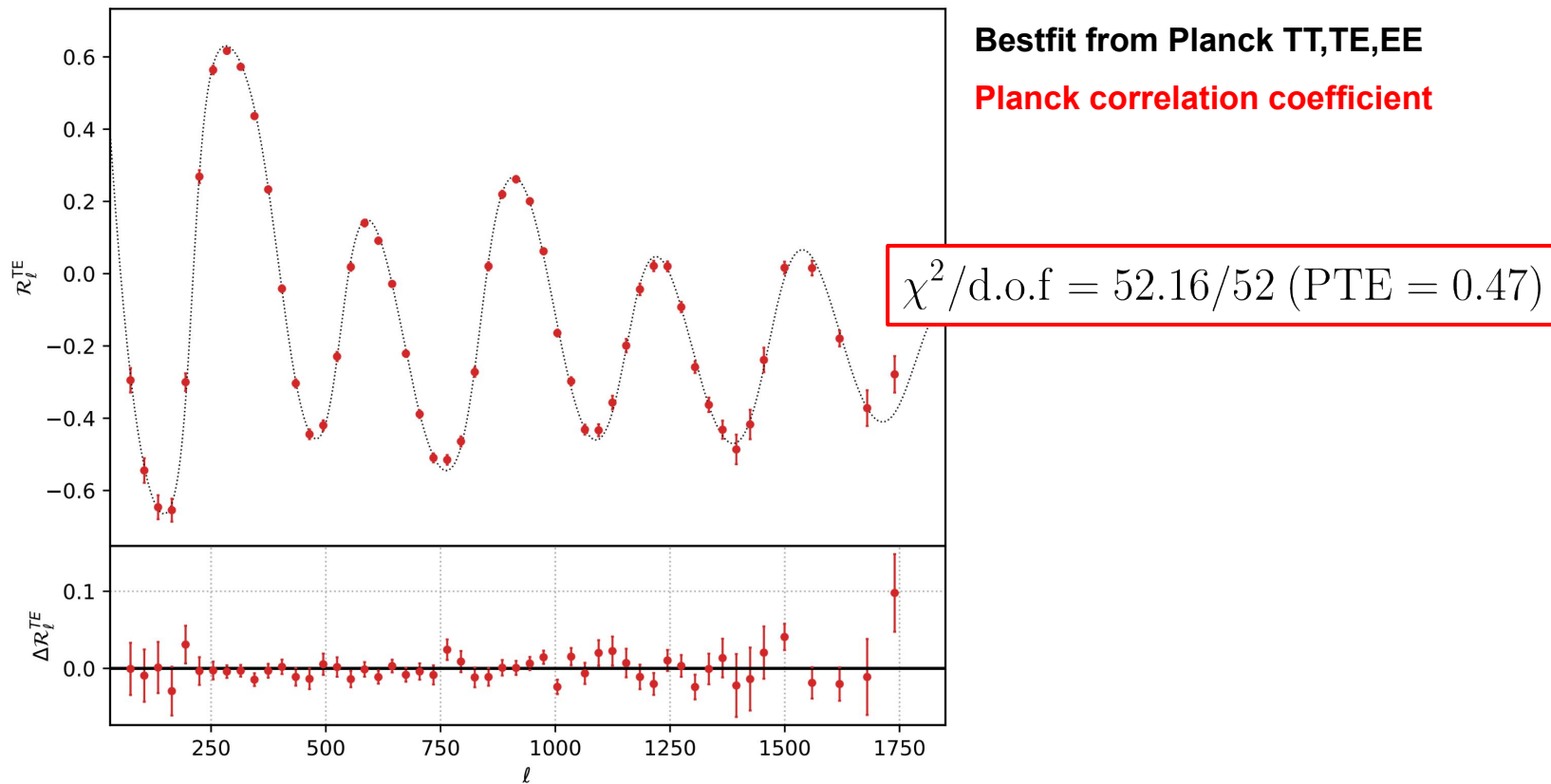
$$\ln \mathcal{L} \simeq -\frac{1}{2} (\Delta \mathcal{R}^{\text{vec}})^T \mathbf{\Xi}^{-1} (\Delta \mathcal{R}^{\text{vec}})$$

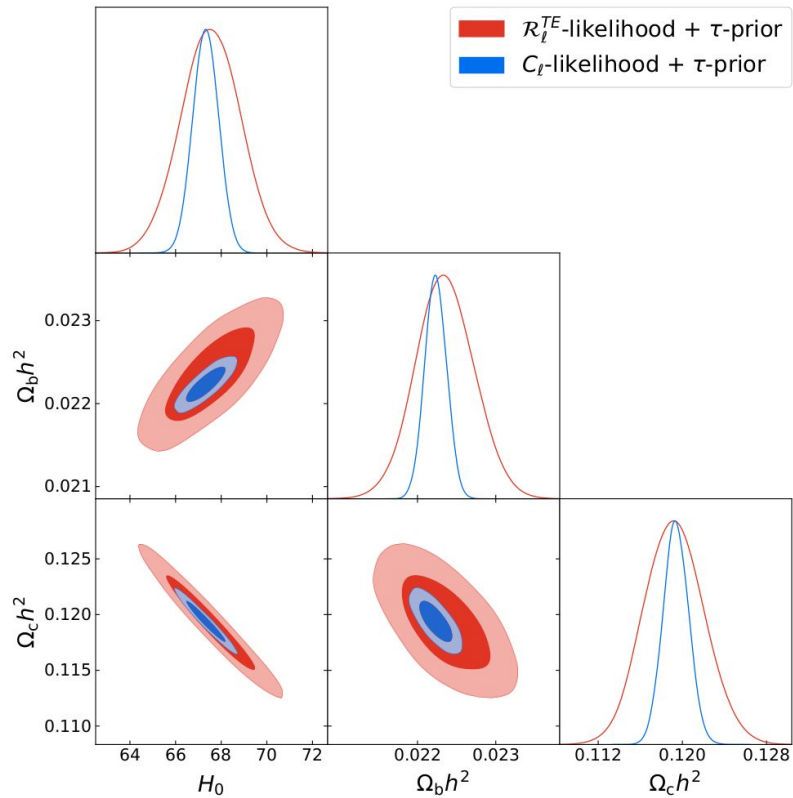
$$\Delta \mathcal{R}_\ell^{TE, \nu_1 \times \nu_2} = \boxed{\hat{\mathcal{R}}_\ell^{TE, \nu_1 \times \nu_2}} - \boxed{\mathcal{R}_\ell^{TE, \nu_1 \times \nu_2, \text{model}}}$$

Unbiased estimator (data)

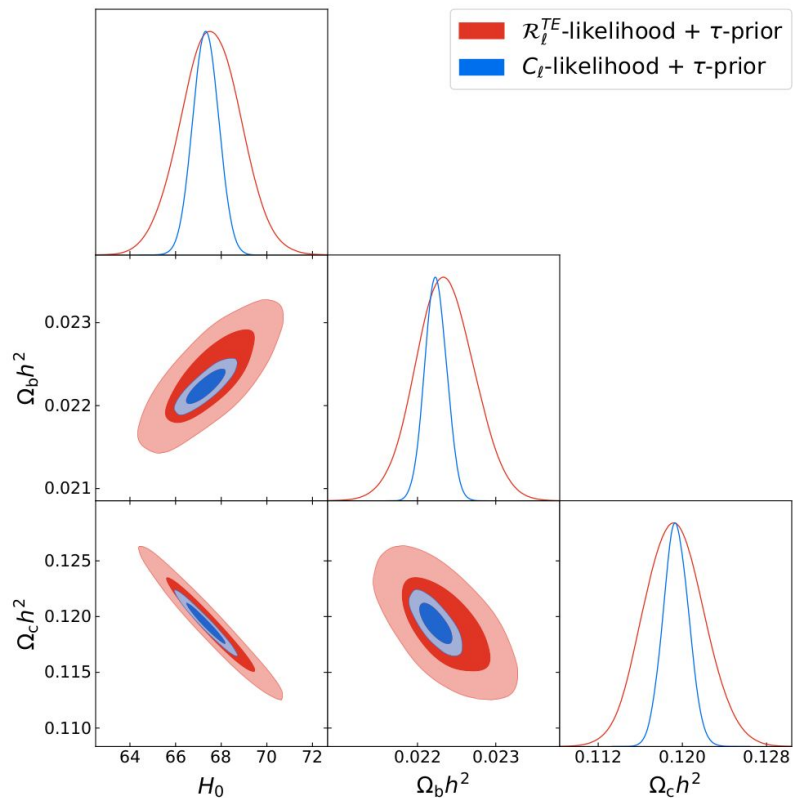
$$\frac{C_\ell^{TE, \nu_1 \times \nu_2, \text{model}}(p_{\text{cosmo}}, p_{\text{fg}})}{\sqrt{C_\ell^{TT, \nu_1 \times \nu_2, \text{model}}(p_{\text{cosmo}}, p_{\text{fg}}) C_\ell^{EE, \nu_1 \times \nu_2, \text{model}}(p_{\text{cosmo}}, p_{\text{fg}})}}$$


Planck correlation coefficient



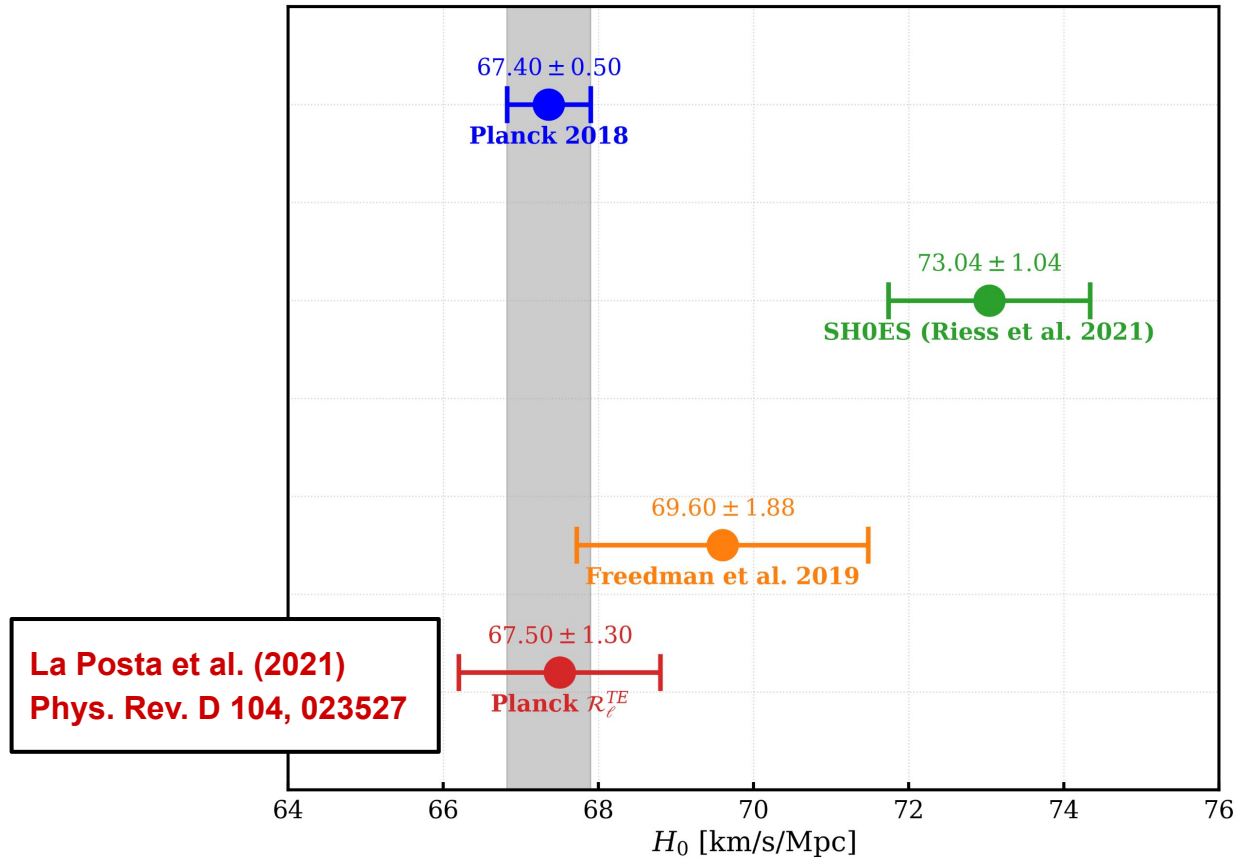


$$H_0 = 67.5 \pm 1.3 \text{ [km/s/Mpc]}$$

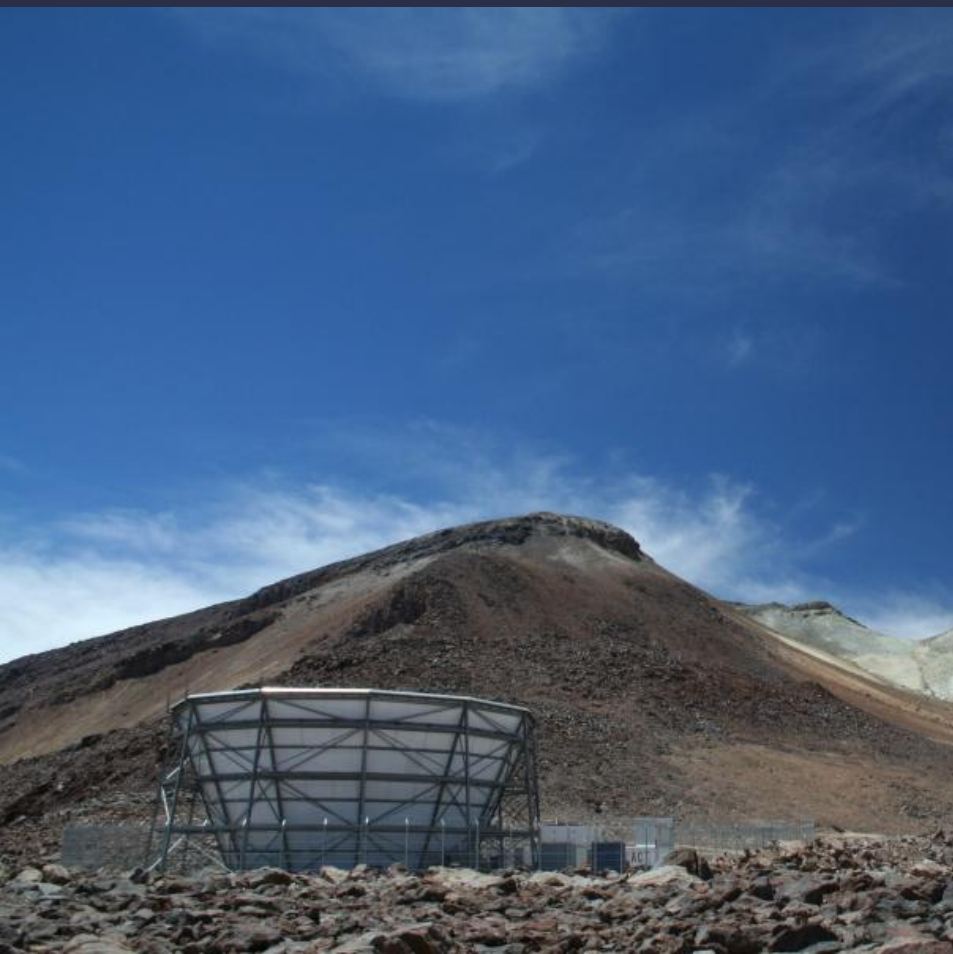


3.3 σ away from the latest
SH0ES measurement

$$H_0 = 67.5 \pm 1.3 \text{ [km/s/Mpc]}$$



Independent measurements of H_0 from the ground



Atacama Cosmology Telescope

6m telescope in the Atacama desert
(Chile ~5000m high)

ACT DR4 (**Choi+ 2020, Aiola+ 2020**)

data collected from 2013 to 2016

Cosmological analysis on ~5400 deg²

observed at 98 and 150 GHz

South Pole Telescope

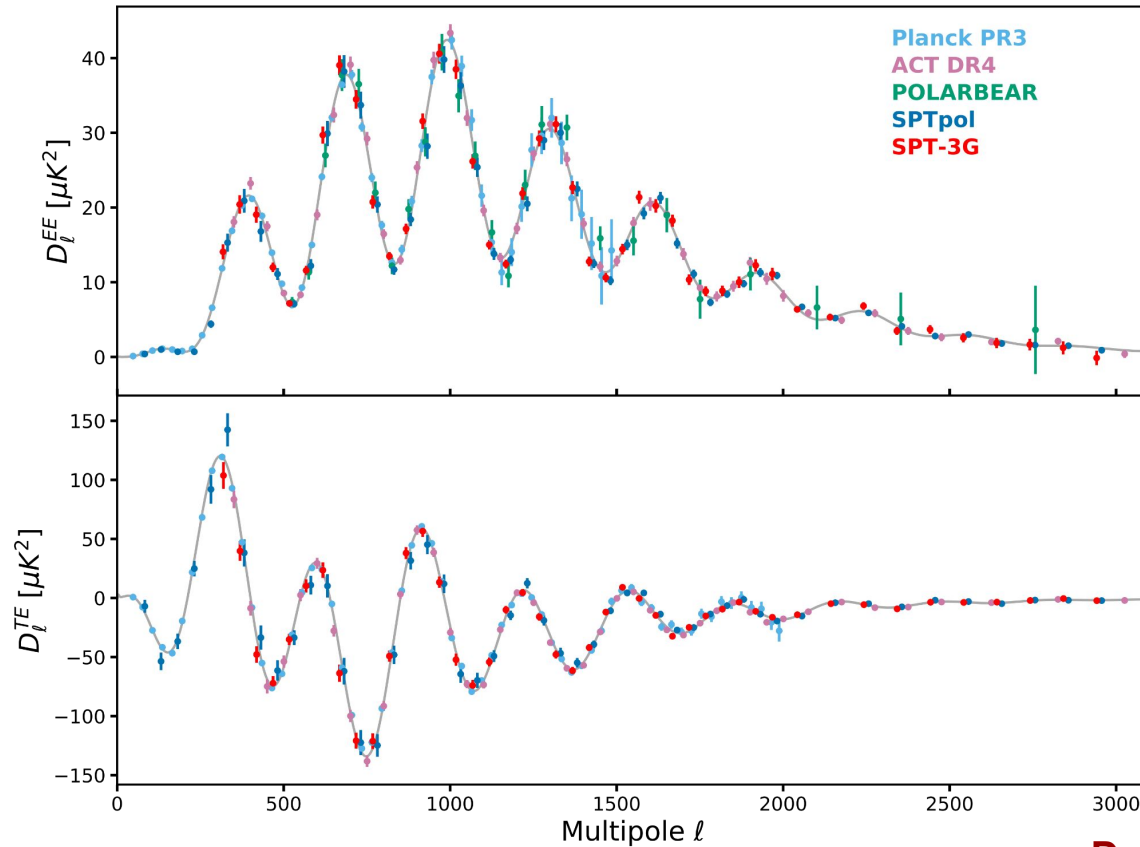
10m primary mirror
(South Pole ~2800m high)

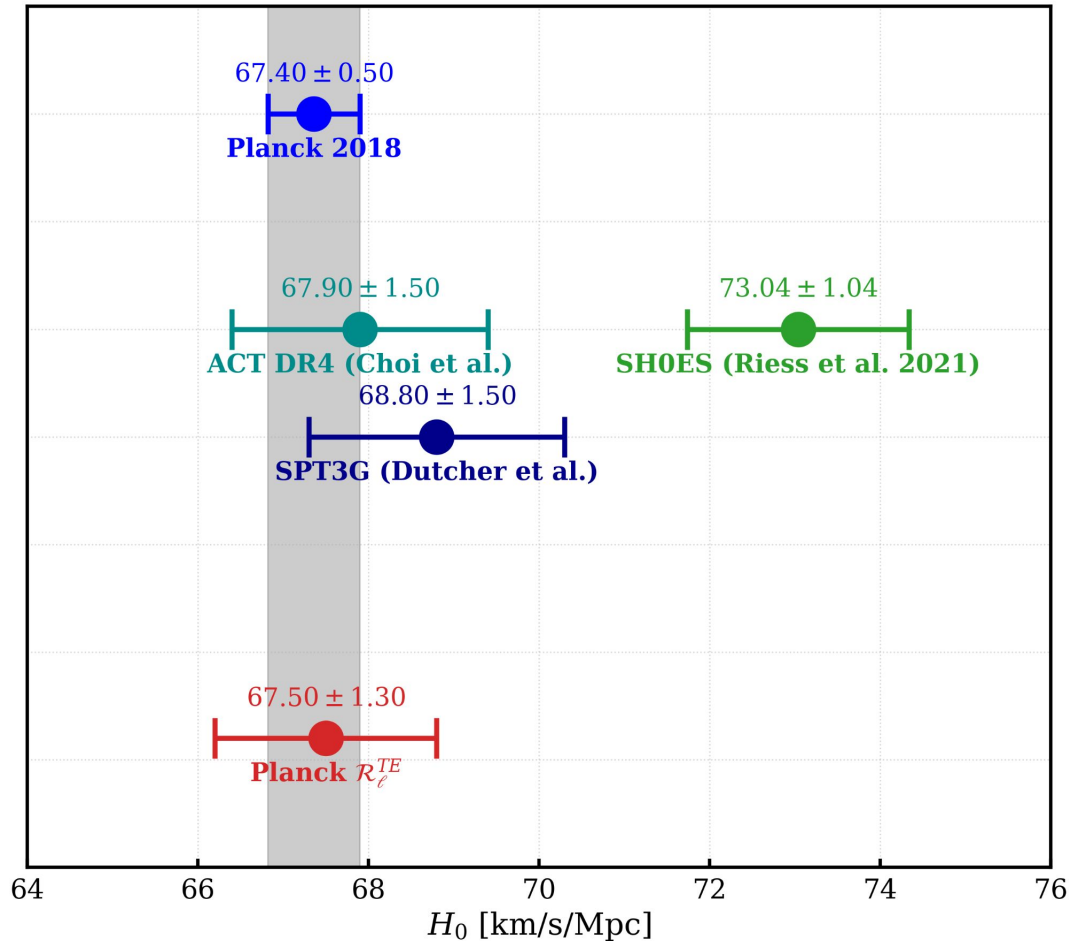
SPT-3G results (**Dutcher+ 2021**)

4 month period in 2018

Cosmological analysis on ~1500 deg²

observed at 95, 150 and 220 GHz





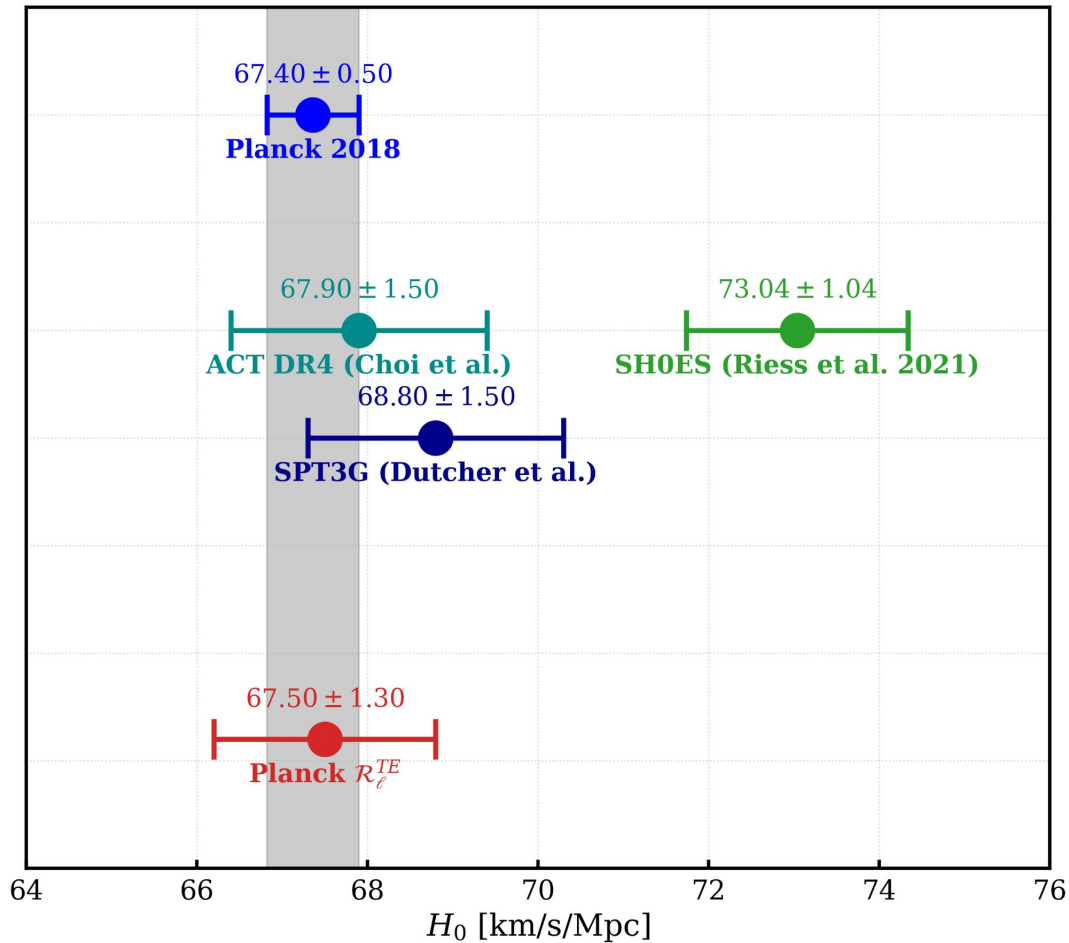
Option 2

Instrumental systematic effect
biasing the value of H_0
inferred from the CMB



Hard to shift the CMB inferred H_0
with a systematic effect :

- Independent measurements from Planck, ACT and SPT
- Constraint from the correlation coefficient, robust against multiplicative systematics



Option 1

Astrophysical biases affecting
the local measurement of H_0

Option 2

Instrumental systematic effect
biasing the value of H_0
inferred from the CMB

Option 3

Physics beyond Λ CDM

Motivation : obtain a higher value of H_0 from the CMB

Motivation : obtain a higher value of H_0 from the CMB \longrightarrow lower \mathcal{D}_A^*

Motivation : obtain a higher value of H_0 from the CMB \longrightarrow lower D_A^*

$$\theta_* = \frac{r_s^*}{D_A^*} \longrightarrow \text{Decrease } r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

θ_* is fixed by observations

$$\frac{3H_{\text{early}}^2(z)}{8\pi G} = \rho_r(z) + \rho_m(z)$$


Motivation : obtain a higher value of H_0 from the CMB \longrightarrow lower D_A^*

$$\theta_* = \frac{r_s^*}{D_A^*} \longrightarrow \text{Decrease } r_s^* = \int_{z^*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

θ_* is fixed by observations

$$\frac{3H_{\text{early}}^2(z)}{8\pi G} = \rho_r(z) + \rho_m(z) + \rho_{\text{EDE}}(z)$$

The EDE component is described as a scalar field ϕ (Poulin+ 2019, Smith+ 2019)

Background evolution : $\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$  axion-like potential

$$V(\phi) = m^2 f^2 \left[1 - \cos \left(\frac{\phi}{f} \right) \right]^3$$

The EDE component is described as a scalar field ϕ (Poulin+ 2019, Smith+ 2019)

Background evolution : $\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$

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m : mass parameter

The EDE component is described as a scalar field ϕ (Poulin+ 2019, Smith+ 2019)

Background evolution : $\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$

$$V(\phi) = m^2 f^2 \left[1 - \cos \left(\frac{\phi}{f} \right) \right]^3$$

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The EDE component is described as a scalar field ϕ (Poulin+ 2019, Smith+ 2019)

Background evolution : $\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$

$$V(\phi) = m^2 f^2 \left[1 - \cos\left(\frac{\phi}{f}\right) \right]^3$$

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f : decay constant

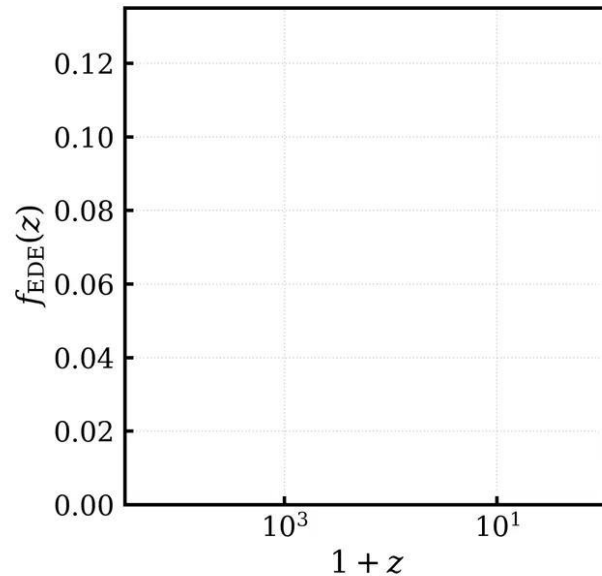
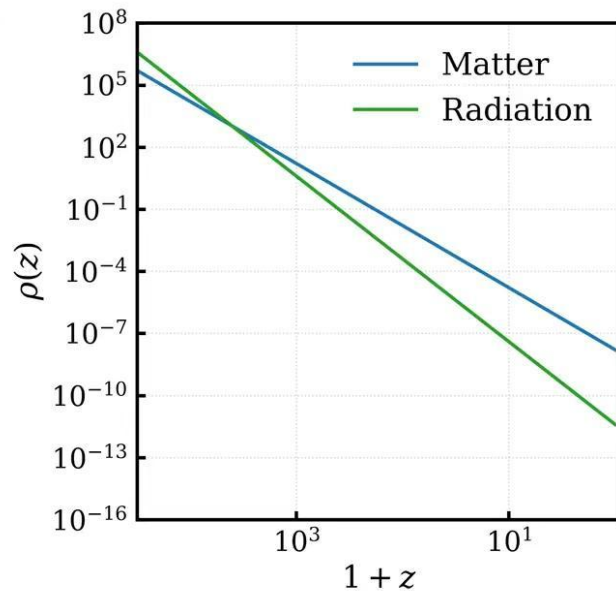
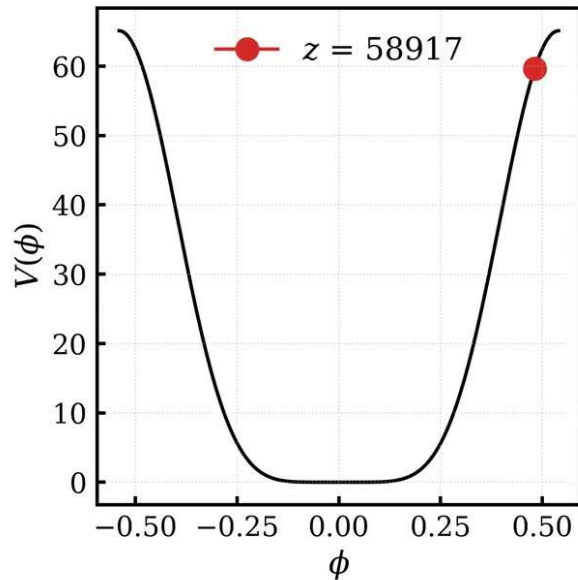
ϕ_i : initial field value

Early Dark Energy : frozen at early times

$$\ddot{\phi} + \boxed{3H\dot{\phi}} + V'(\phi) = 0$$

The field is initially frozen due to
Hubble friction ($H \gg m$)

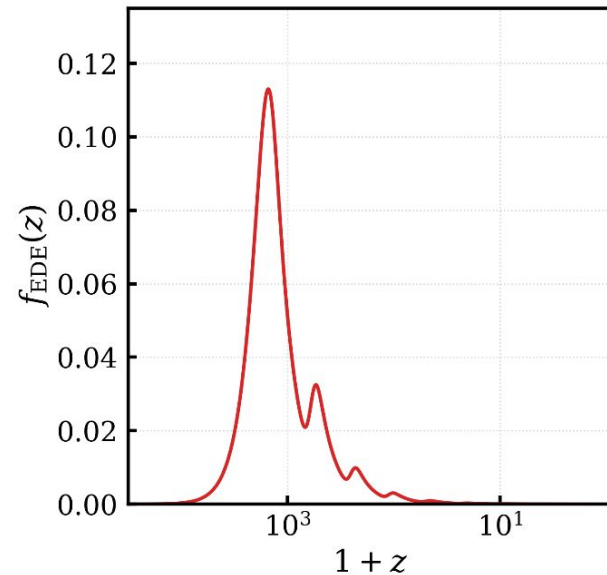
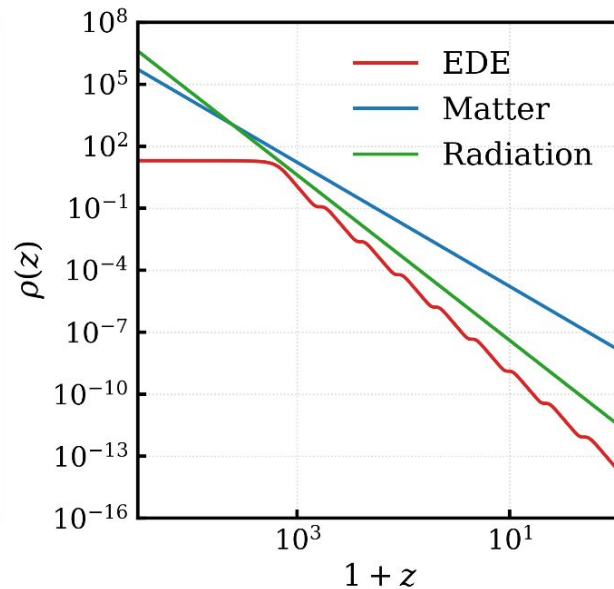
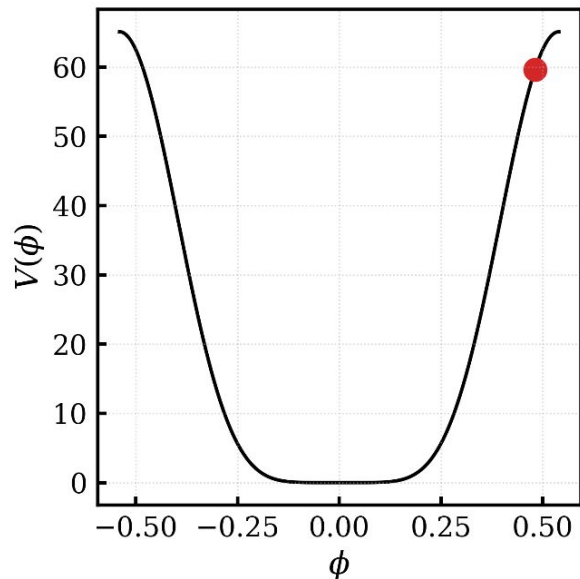
acts as dark energy ($w = -1$)



$$\ddot{\phi} + \boxed{3H\dot{\phi}} + V'(\phi) = 0$$

The field is initially frozen due to Hubble friction ($H \gg m$)

acts as dark energy ($w = -1$)

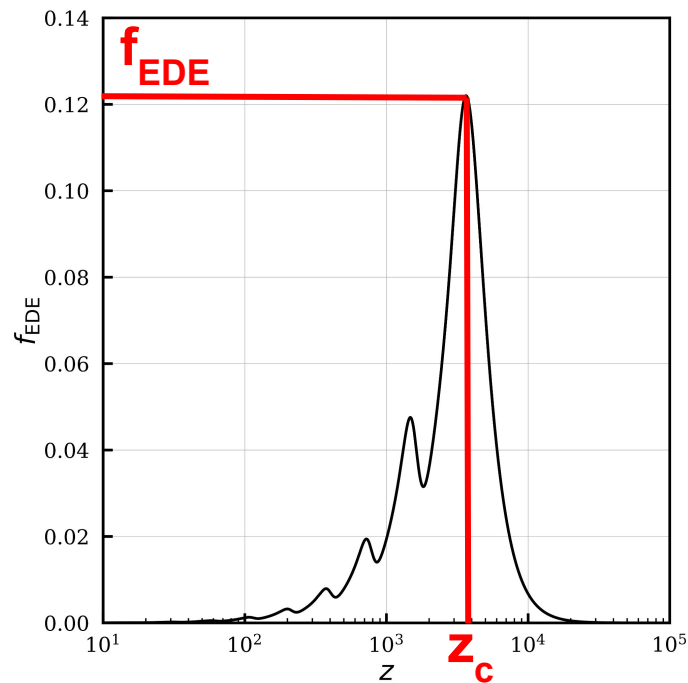


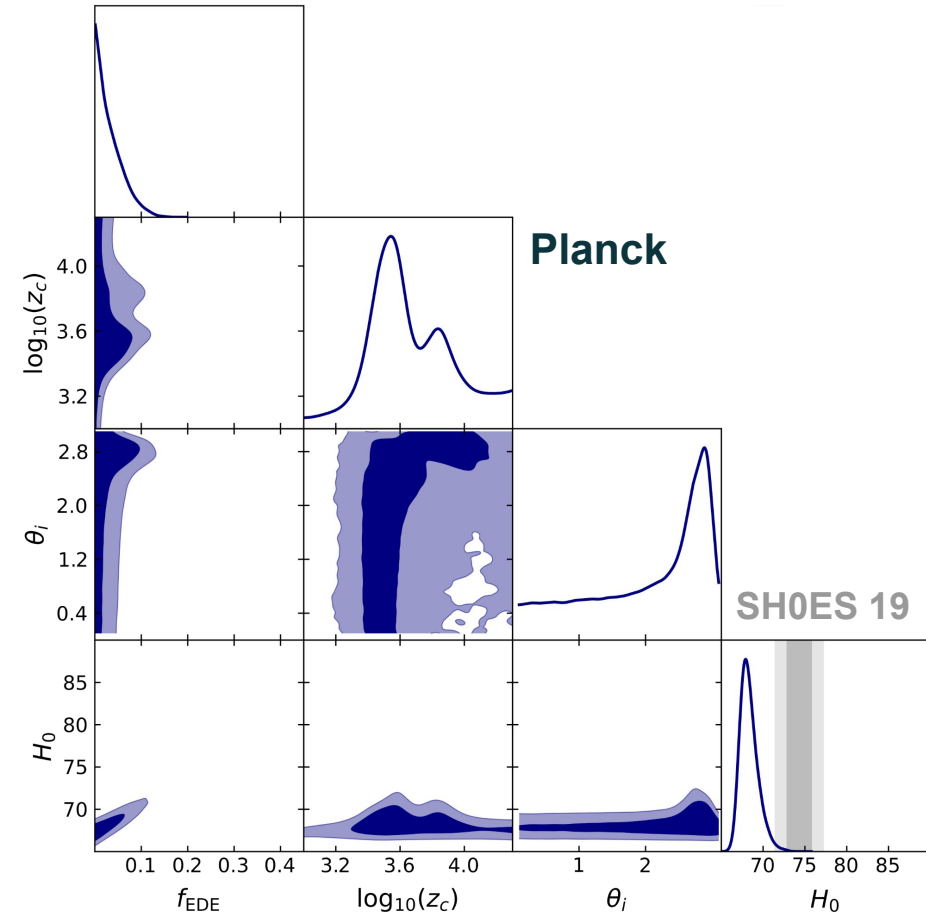
Early Dark Energy : phenomenological parametrization 24

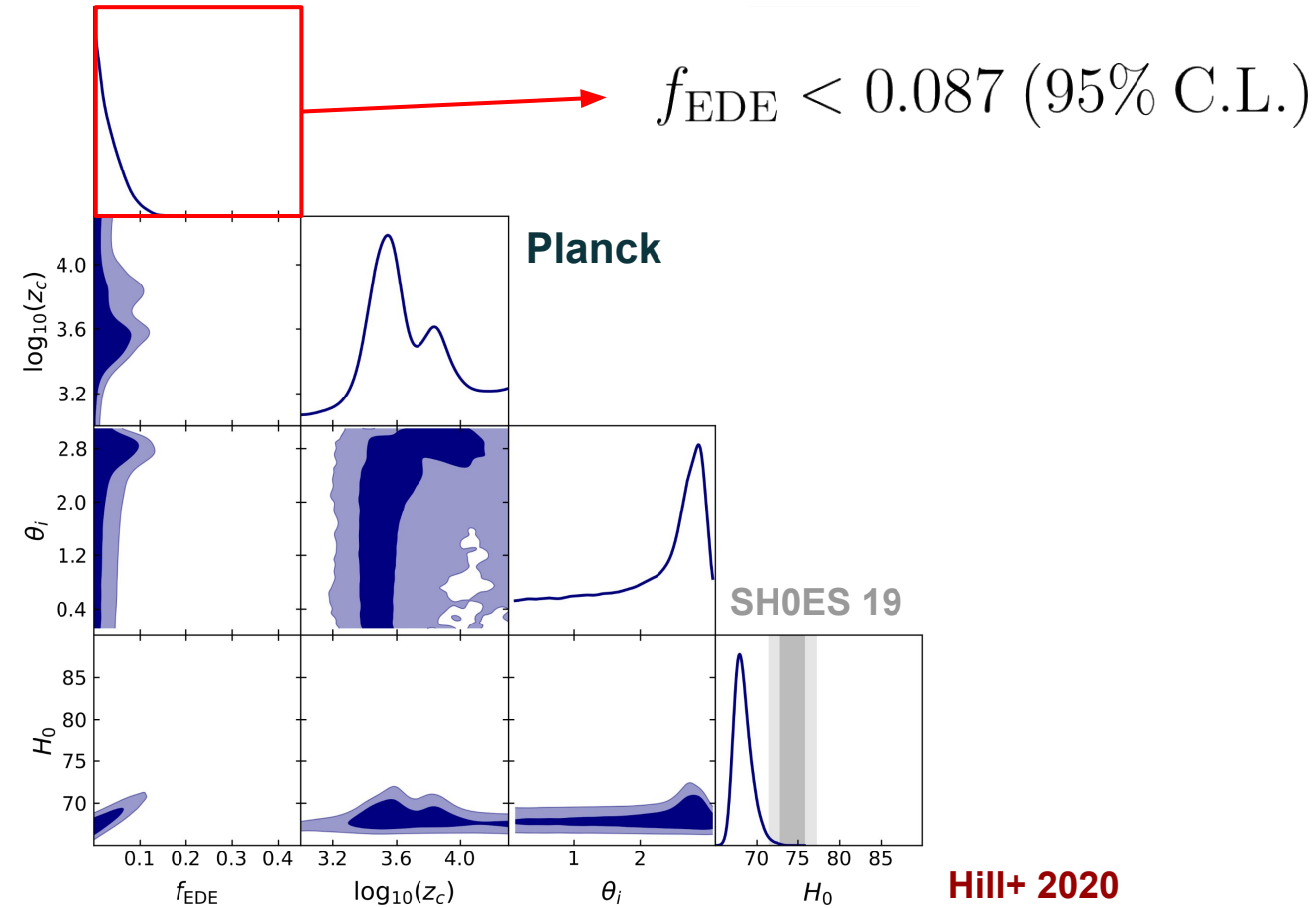
$$(m, f, \phi_i) \longrightarrow (f_{\text{EDE}}, z_c, \phi_i)$$

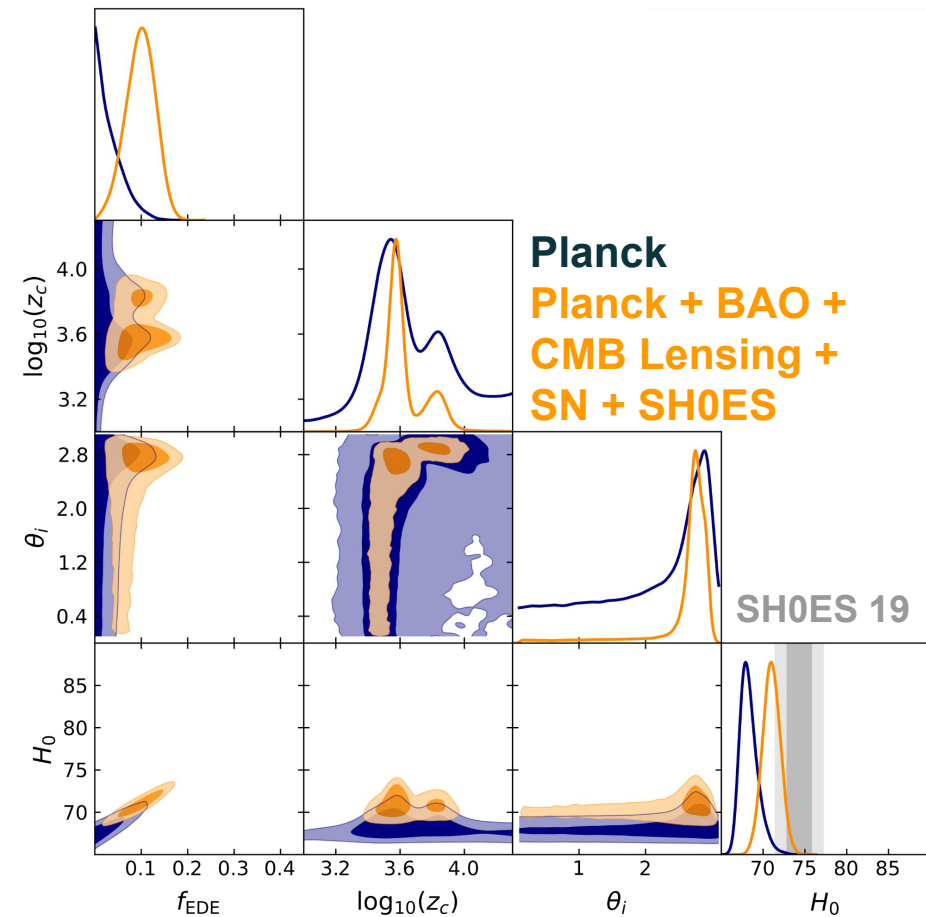
Field parameters

Phenomenological parametrization

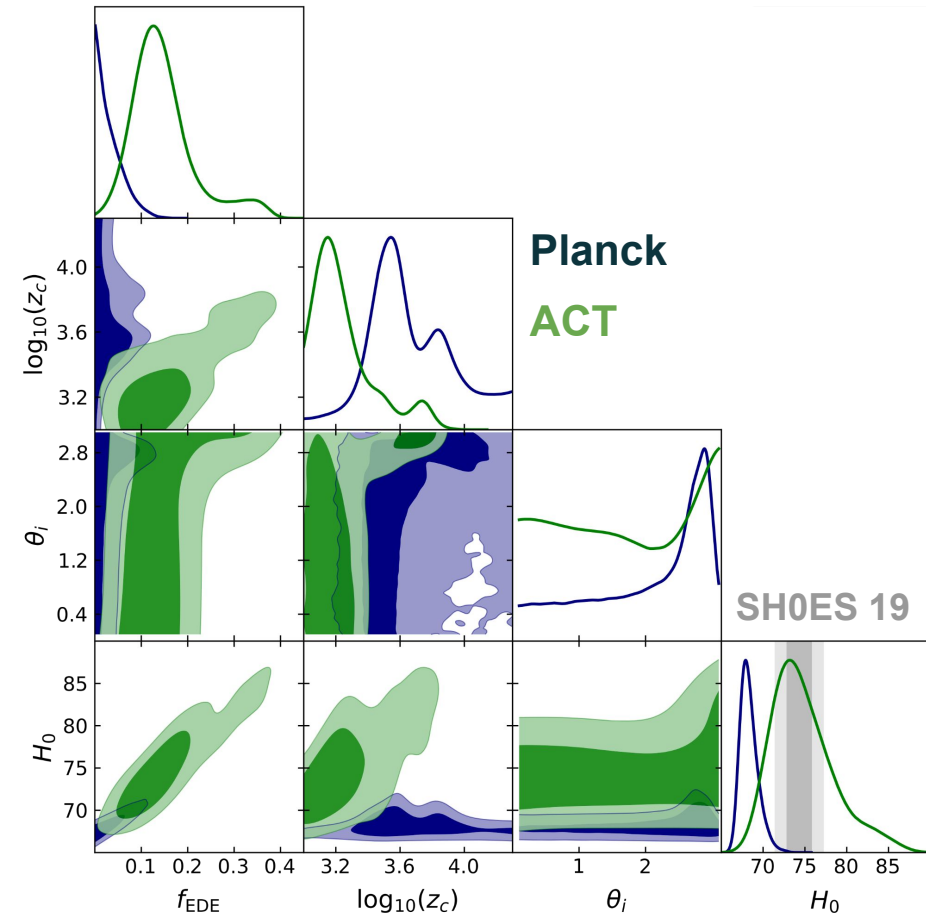


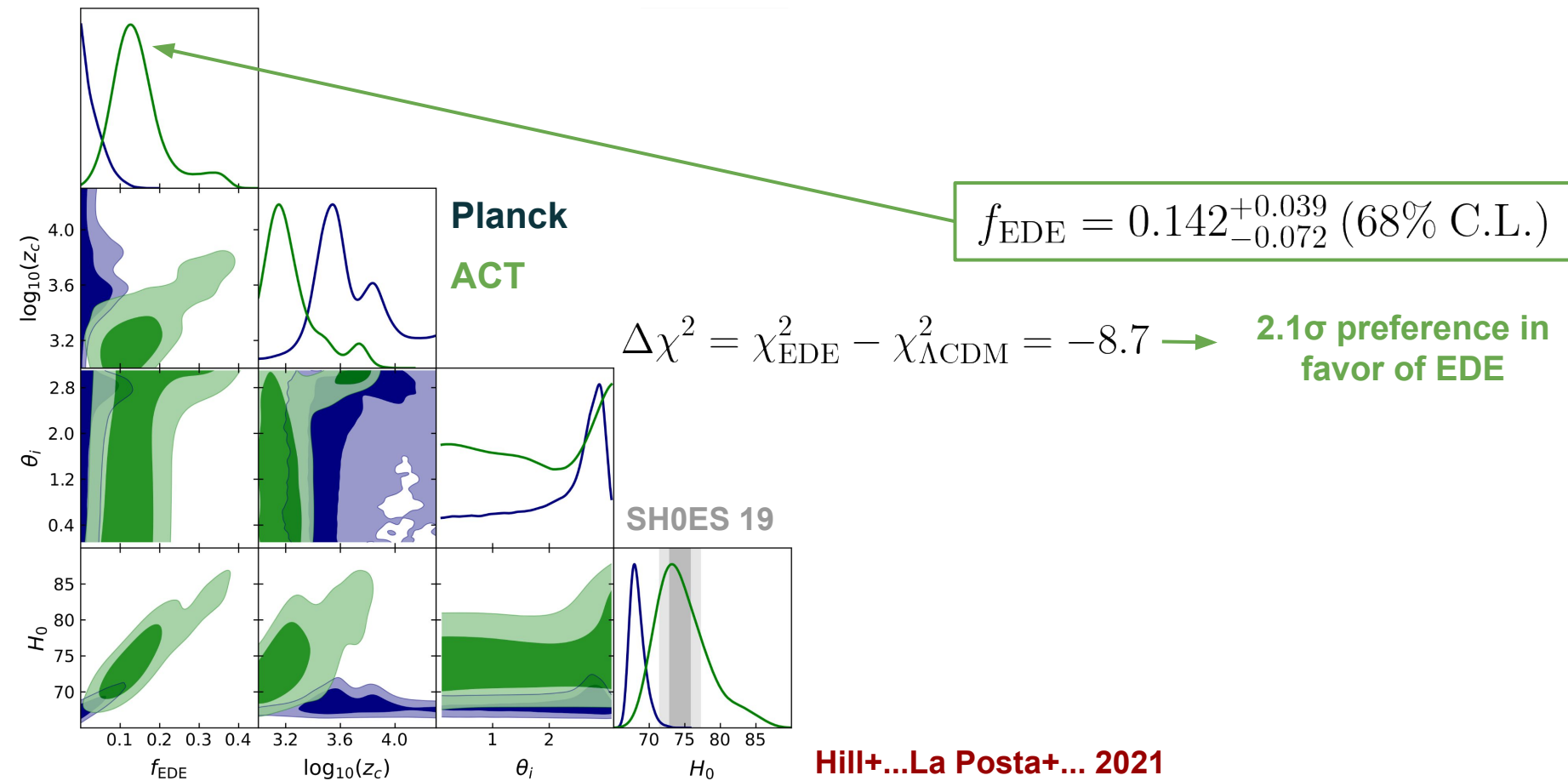


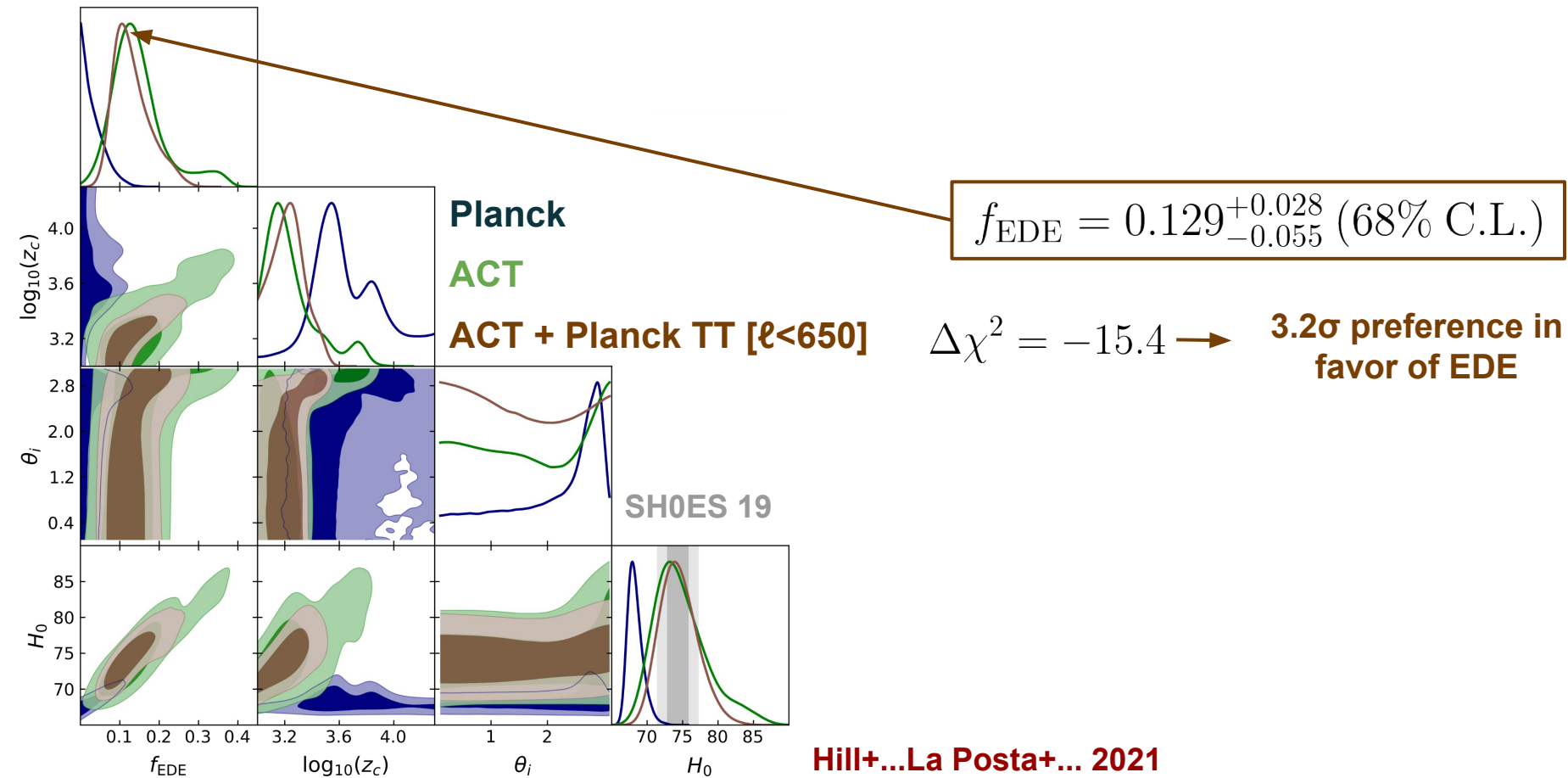




Poulin+ 2019, Smith+ 2019, Hill+ 2020

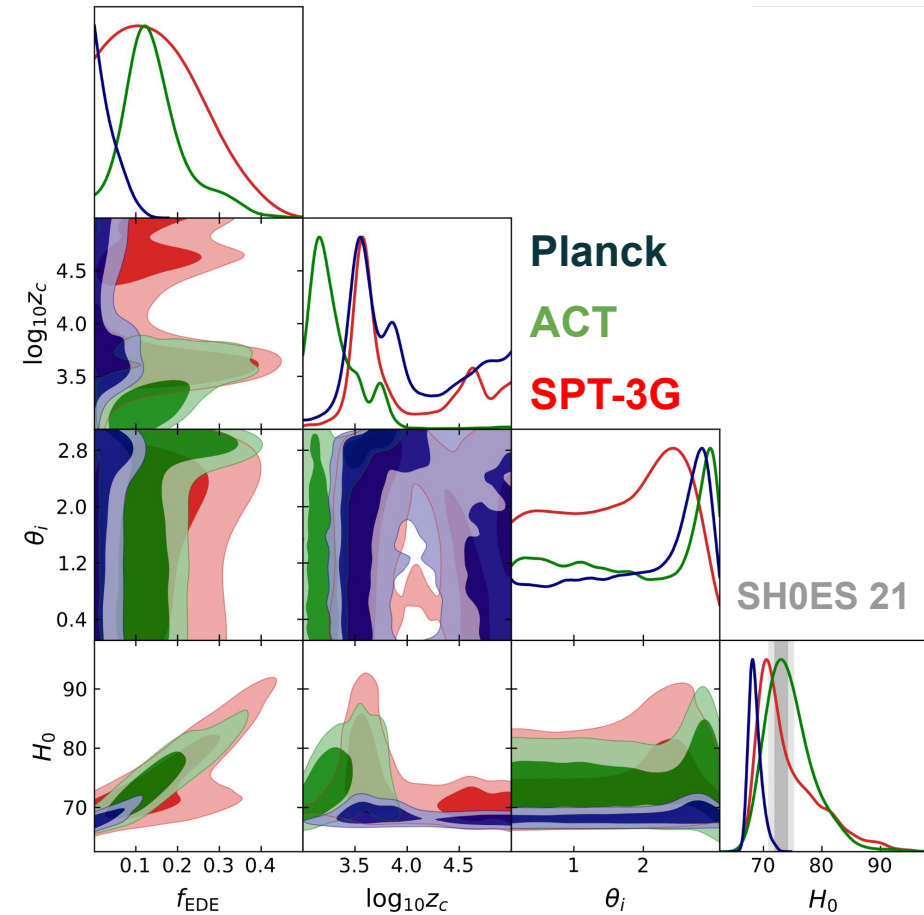


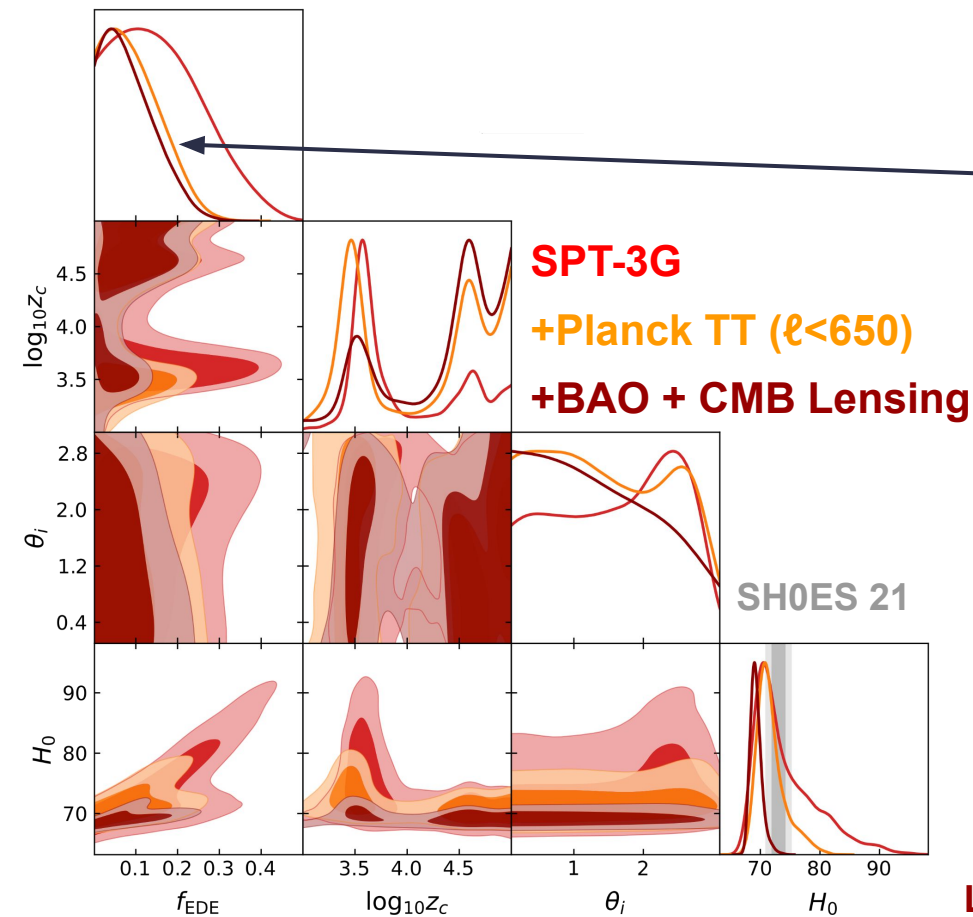




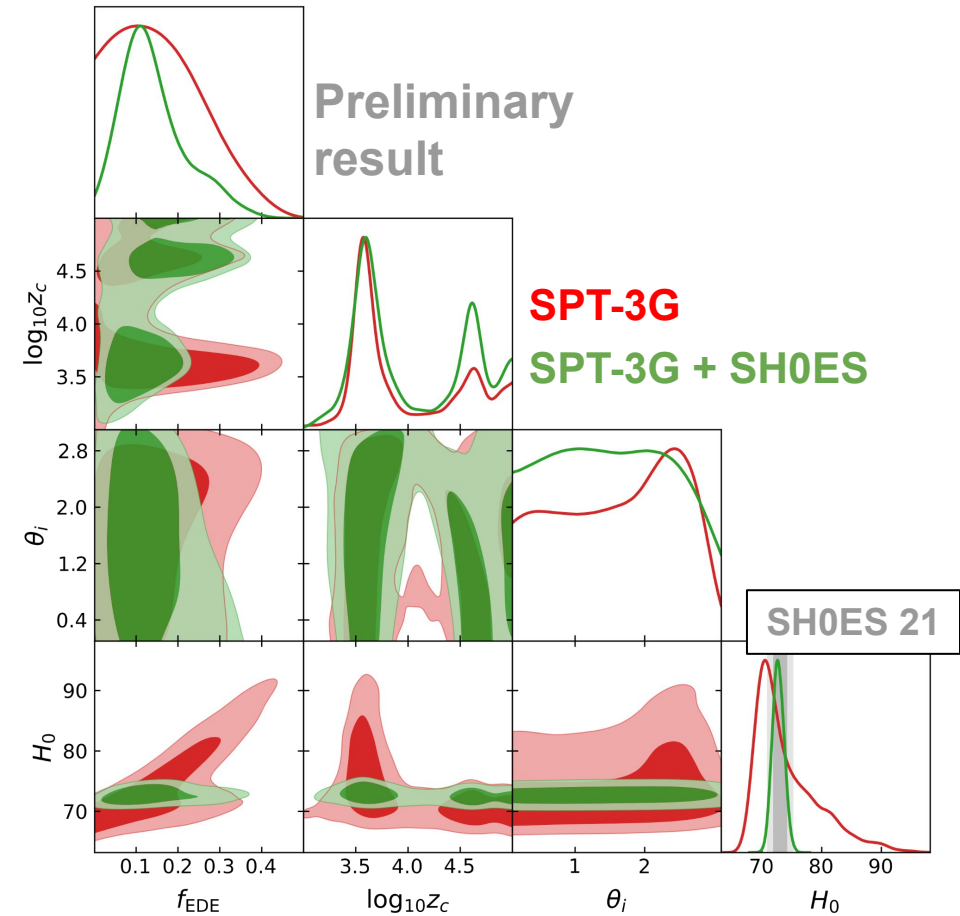
- Planck data alone don't favor high f_{EDE} values (Hill+ 2020)
- Planck data in combination with SH0ES show a preference for non-zero f_{EDE} (Poulin+ 2019, Smith+ 2019)
- ACT data alone favors EDE over Λ CDM (Hill+...La Posta+... 2021)

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- Motivates an analysis of EDE with public SPT-3G data





We tighten the constraint on f_{EDE} when we combine **SPT3G** and **Planck TT ($\ell < 650$)** or when we add **LSS probes**

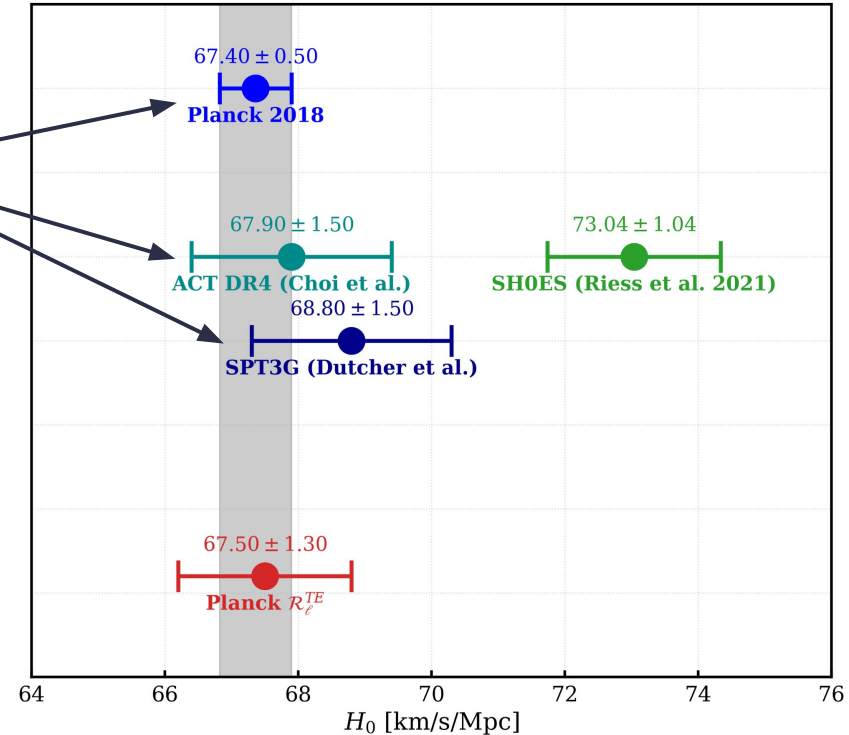


$$\Delta\chi_{\text{SPT-3G}}^2 = -6.3$$

improvement of the fit to
SPT-3G data (with
respect to ΛCDM)

Option 2 : Systematics in CMB data

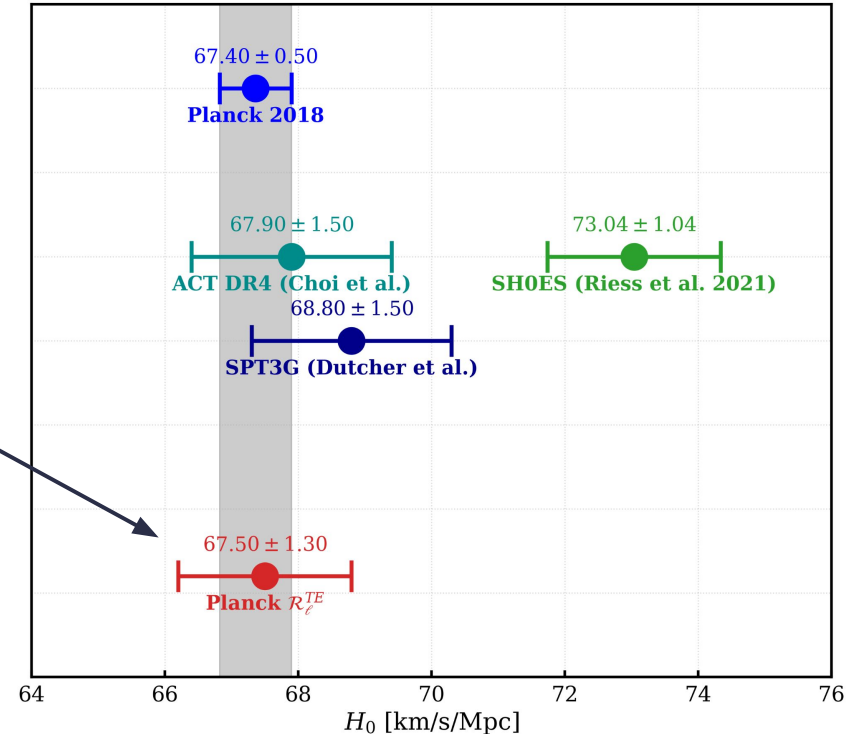
- 3 independent measurements of H_0



Option 2 : Systematics in CMB data

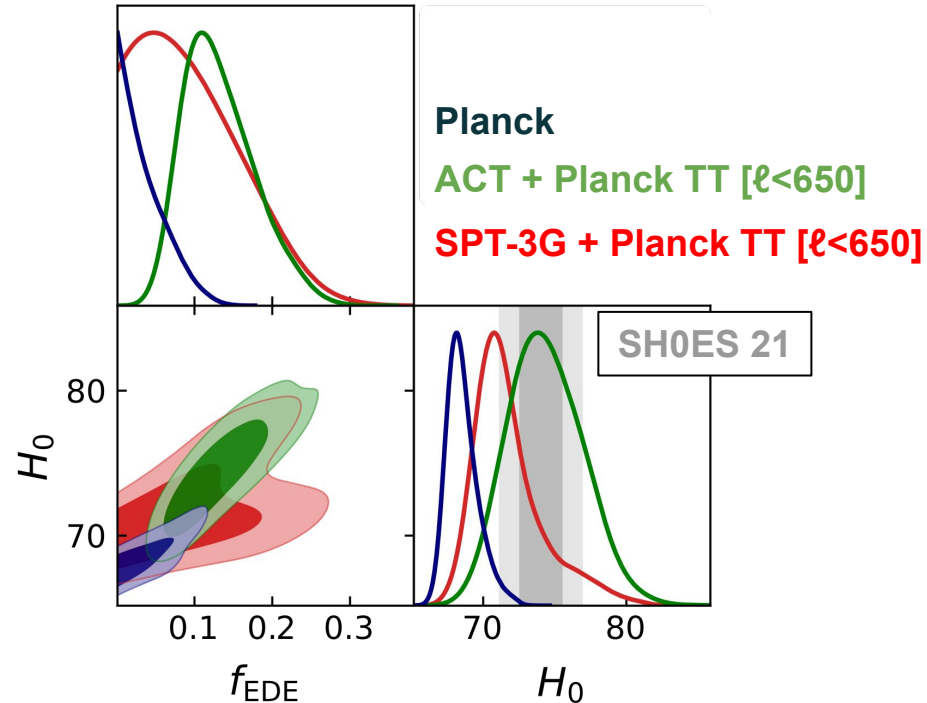
- 3 independent measurements of H_0
- Constraint from the **correlation coefficient** : insensitive to multiplicative systematic effects

It's hard to solve the Hubble tension with systematics in CMB data



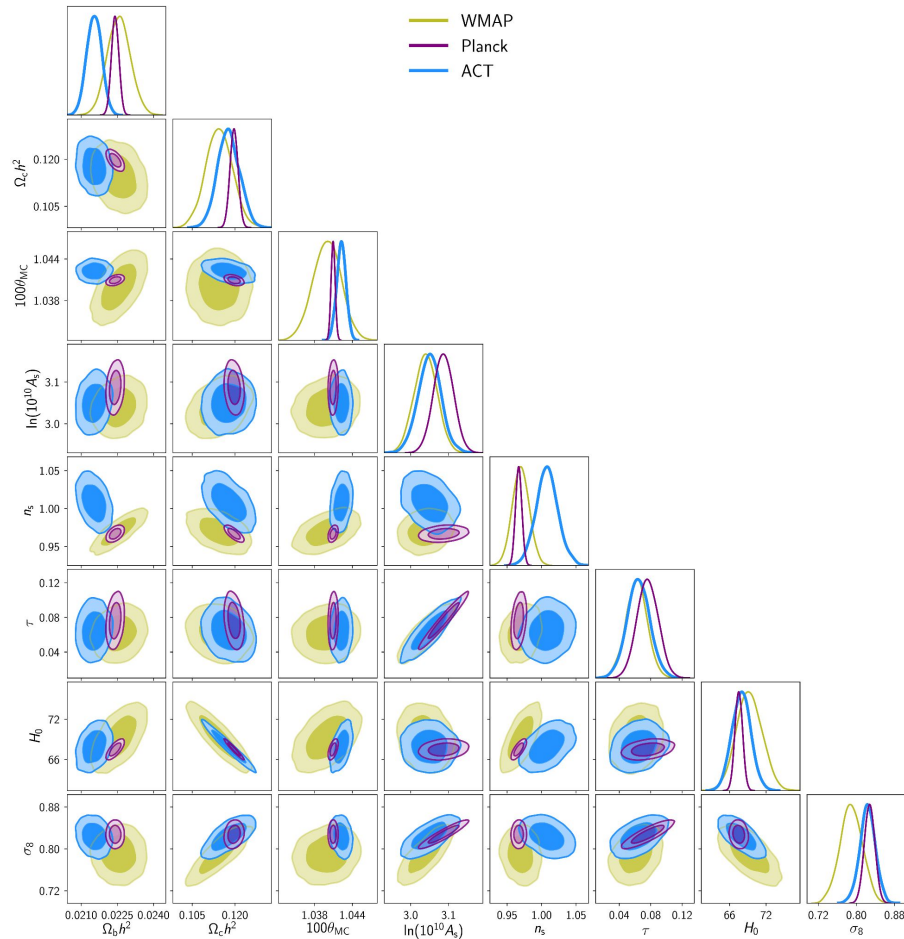
Option 3 : Beyond Λ CDM physics - Early Dark Energy

- Planck data alone do not favor high f_{EDE} values
- Planck + SH0ES show a preference for $f_{\text{EDE}} \sim 10\%$
- ACT DR4 data favors EDE over Λ CDM (with $f_{\text{EDE}} \sim 10\%$)
- SPT-3G is not as constraining as ACT and Planck : but sees some degree of EDE when combined with SH0ES

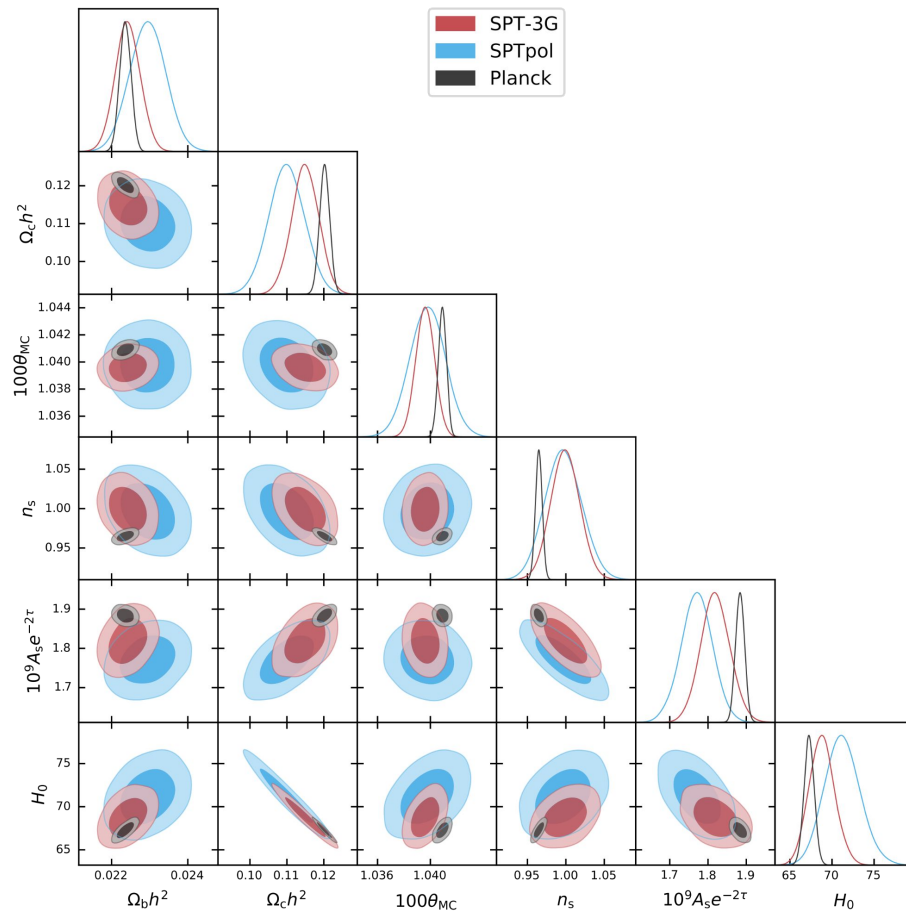


Extra-slides

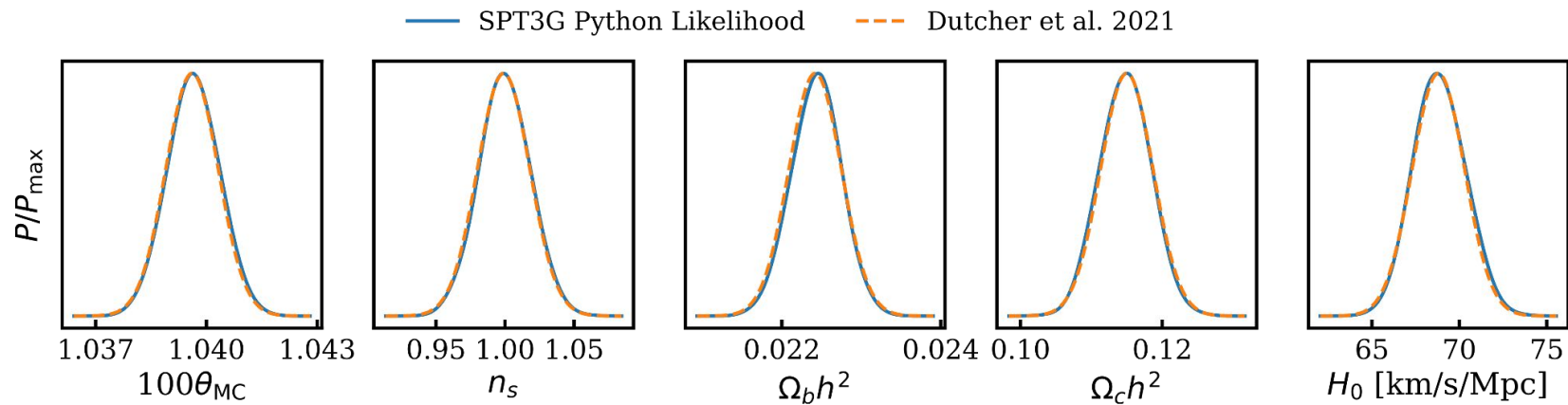
ACTPol cosmology



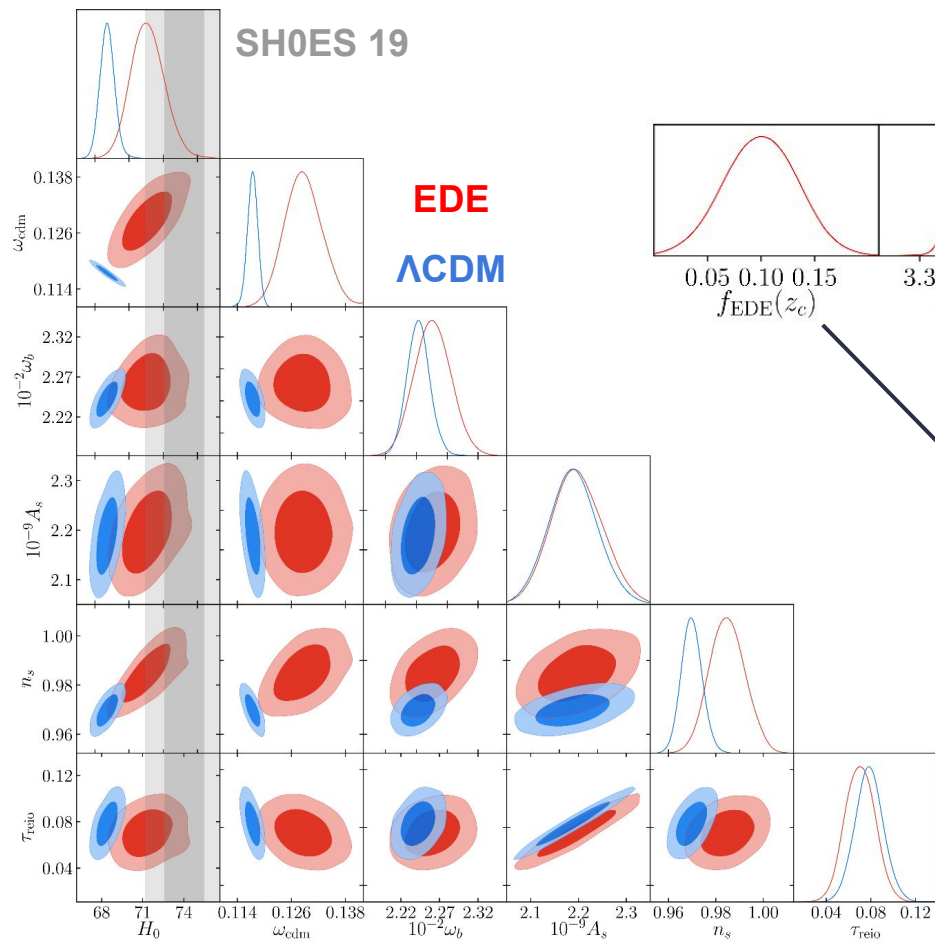
SPT-3G cosmology



SPT-3G python implementation



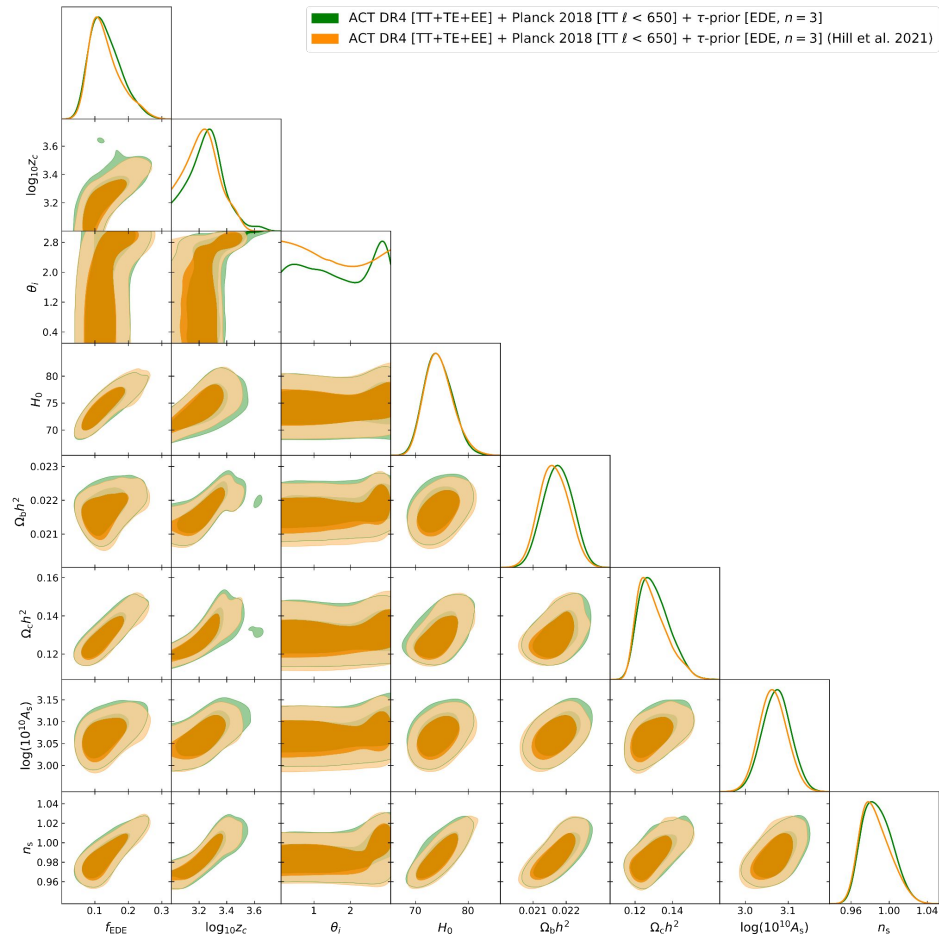
Results for a combination of Planck and SH0ES data

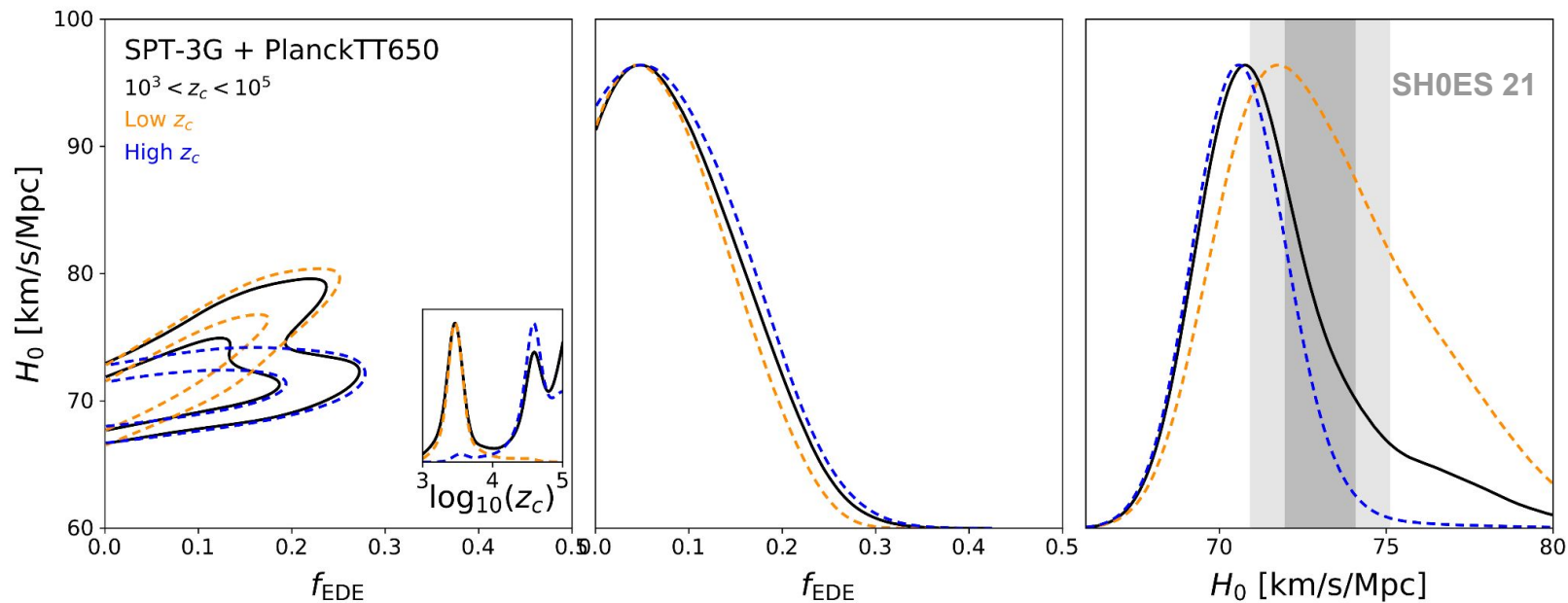


~10% of EDE is required to alleviate the tension

Poulin+2019, Smith+ 2019

CAMB/CLASS EDE models





Combining with other CMB datasets

