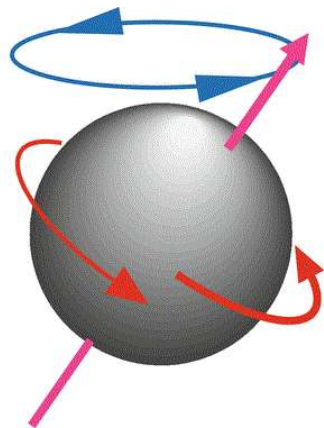


# Characterization of Single Crystalline FeV thin Films on GaAs Substrate by Ferromagnetic Resonance

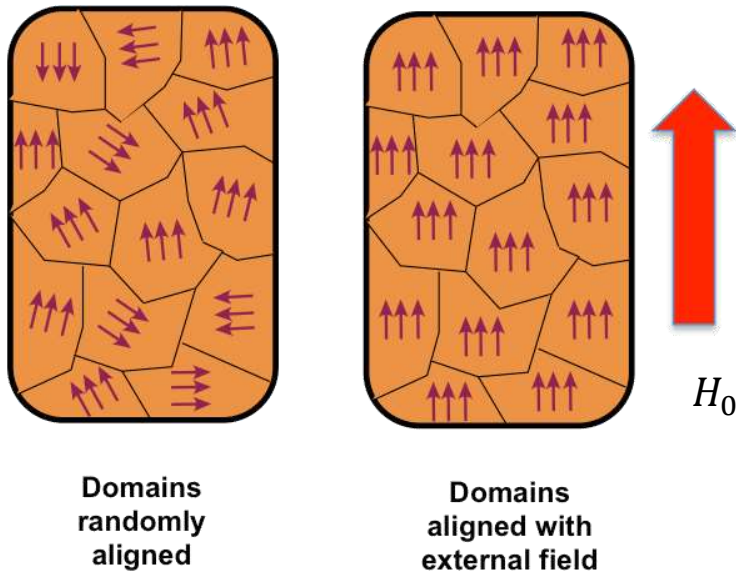


**Student Name - G.D.K. De Silva**  
**Supervisor - Matthieu Bailleul**

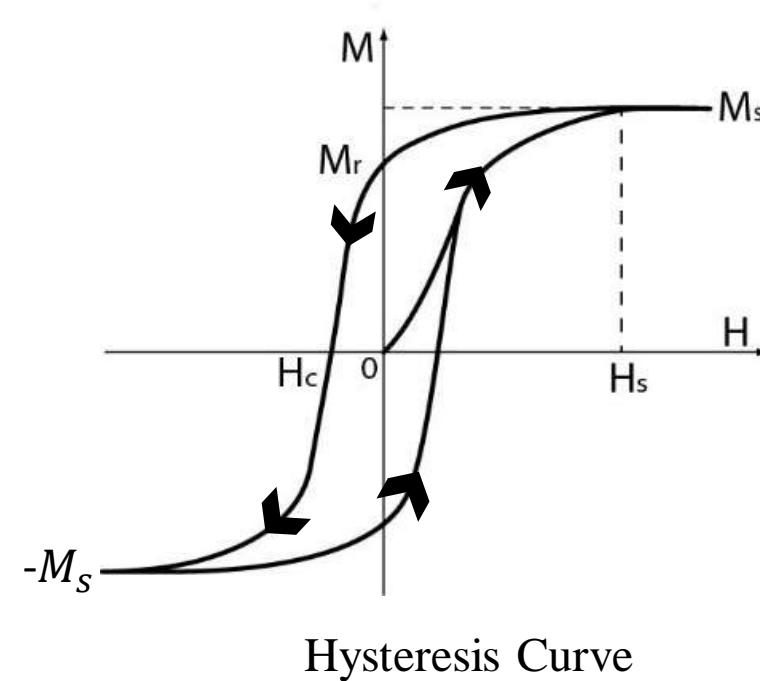
# Objectives

- ❖ Acquire experience of the Ferromagnetic Resonance (FMR) method and experimental setup.
- ❖ Understand the magnetic state of Fe and FeV thin films (20 nm) grown on GaAs substrate.

# Ferromagnetic Materials



- ❖ If no magnetic field is applied, Each Domain align in different directions
- ❖ If an external magnetic field is applied, domains align to the magnetic field



**Magnetization (M)** - The density of magnetic moments in a magnetic material

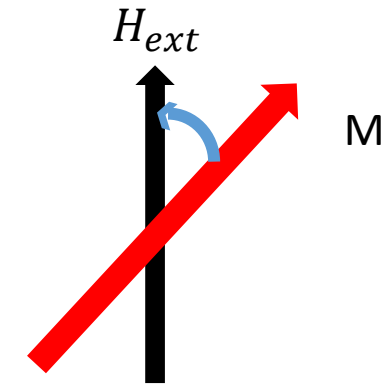
$$M = \frac{\sum m_i}{V}$$

**Saturation Magnetization ( $M_s$ )** - maximum magnetic moment per unit volume for a magnetic material

# Magnetic Energies

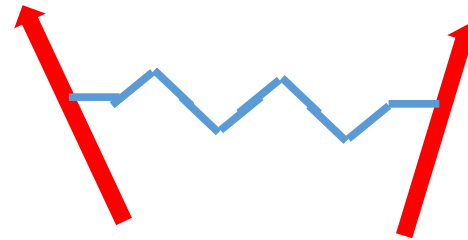
## ❖ Zeeman Energy

The interaction of the magnetization  $M$  with an **external magnetic field**  $H_{ext}$ .



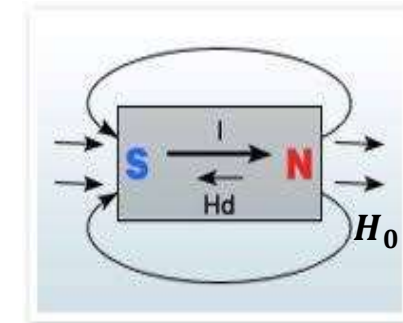
## ❖ Exchange Energy

Interaction Energy between **two Spins**



## ❖ Demagnetizing Field Energy

This is given by the **dipolar interaction** between magnetic moments in the material. This interaction creates a field that opposes the magnetization.



## ❖ Cubic Anisotropy Energy

Energy that depends on **orientation of the magnetization** with respect to the lattice symmetry direction of the material

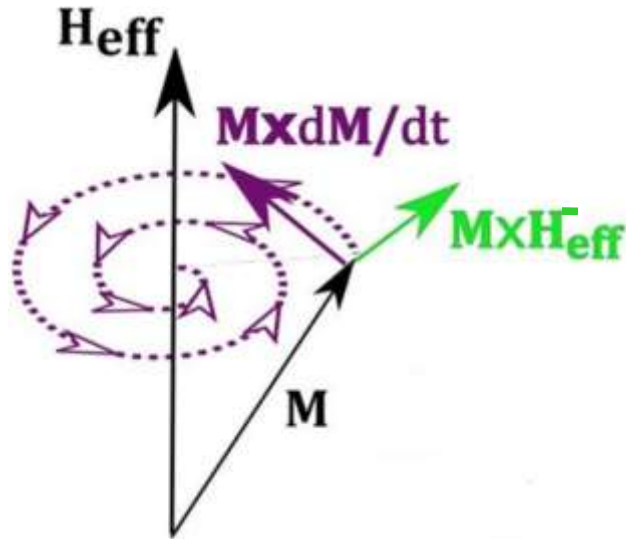


$$\mu_0 \vec{H}_{eff} = \vec{\nabla}_{\vec{M}} \epsilon_T$$

$$\vec{H}_{eff} = \vec{H}_{Ze} + \vec{H}_{Ex} + \vec{H}_{de} + \vec{H}_{ku}$$

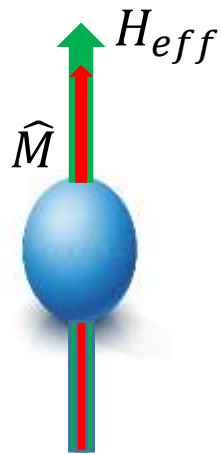
# Magnetization Dynamics

## Landau Lifshitz Gilbert equation

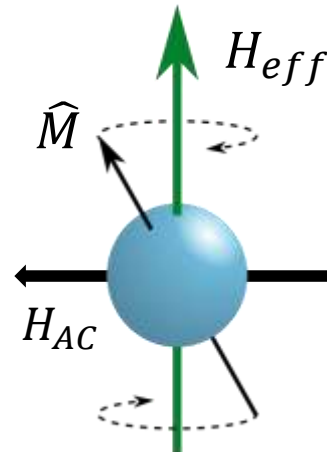


$$\frac{dM}{dt} = \underbrace{-\gamma M \times \mu_0 H_{eff}}_{\text{Steady State Precession}} + \underbrace{\alpha M \times \frac{dM}{dt}}_{\text{Dissipation Damping}}$$

$$f = \frac{\gamma \mu_0}{2\pi} \sqrt{(H_0 + H_x)(H_0 + H_y)}$$



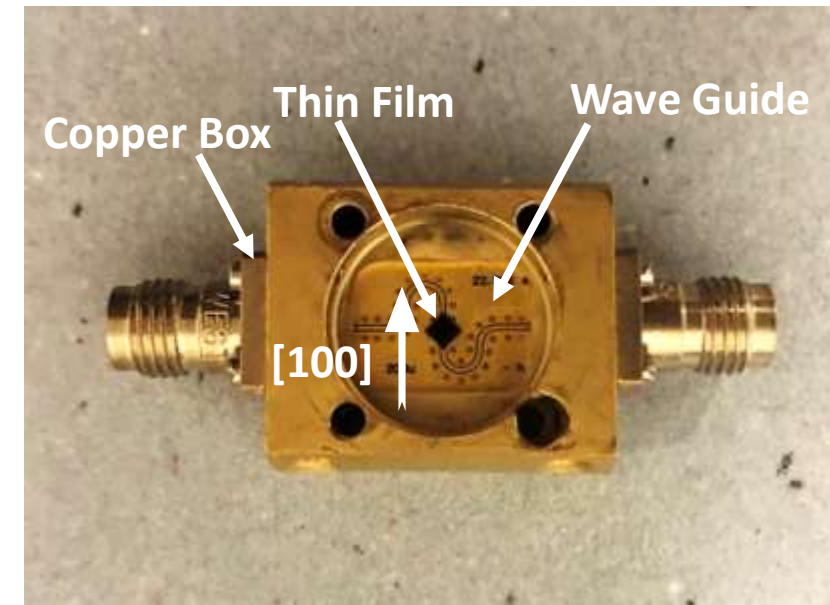
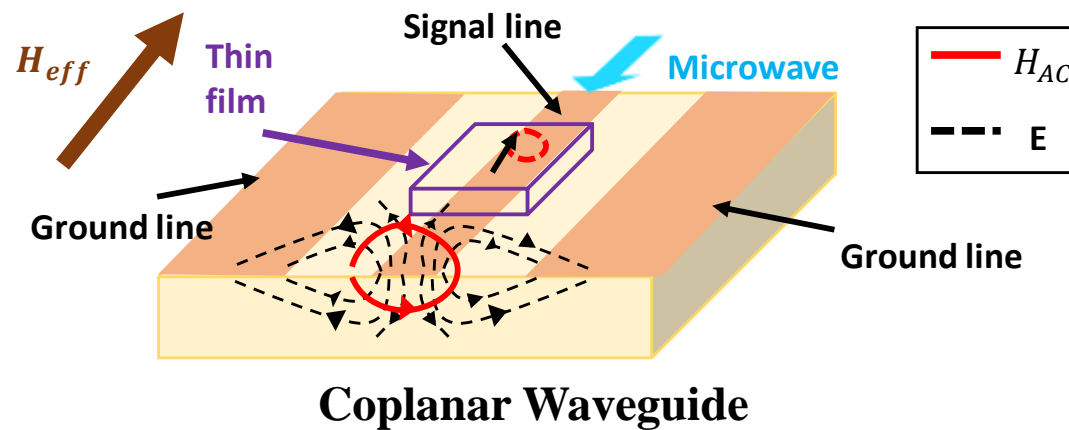
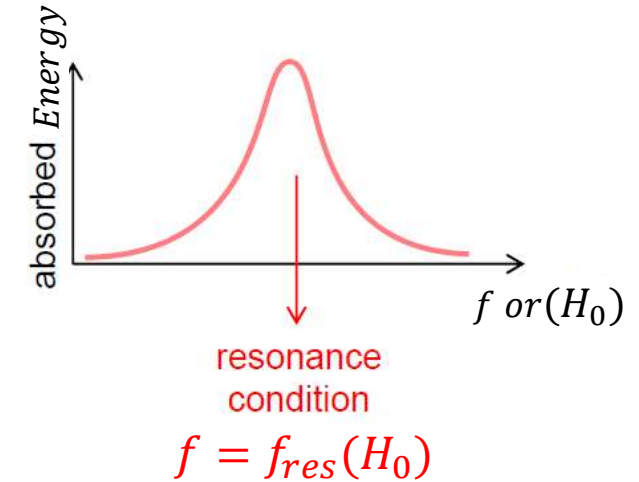
- ❖ At equilibrium, magnetization will align with Effective magnetic field.



- ❖ Applying a transversal alternating magnetic field  $H_{AC}$  will drive the precession of magnetization  $M$ .

# Ferromagnetic Resonance (FMR)

- ❖ Experimentally, the magnetization precession is driven by an external electromagnetic Wave (oscillating magnetic field)
- ❖ The Magnetic material absorbs energy from the microwave leading to the magnetization precession.
- ❖ The energy absorption will be maximum when:  
The frequency of the excitation wave = resonance frequency of the magnetization.

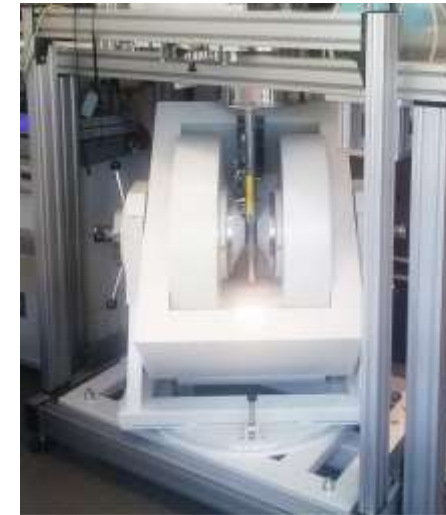
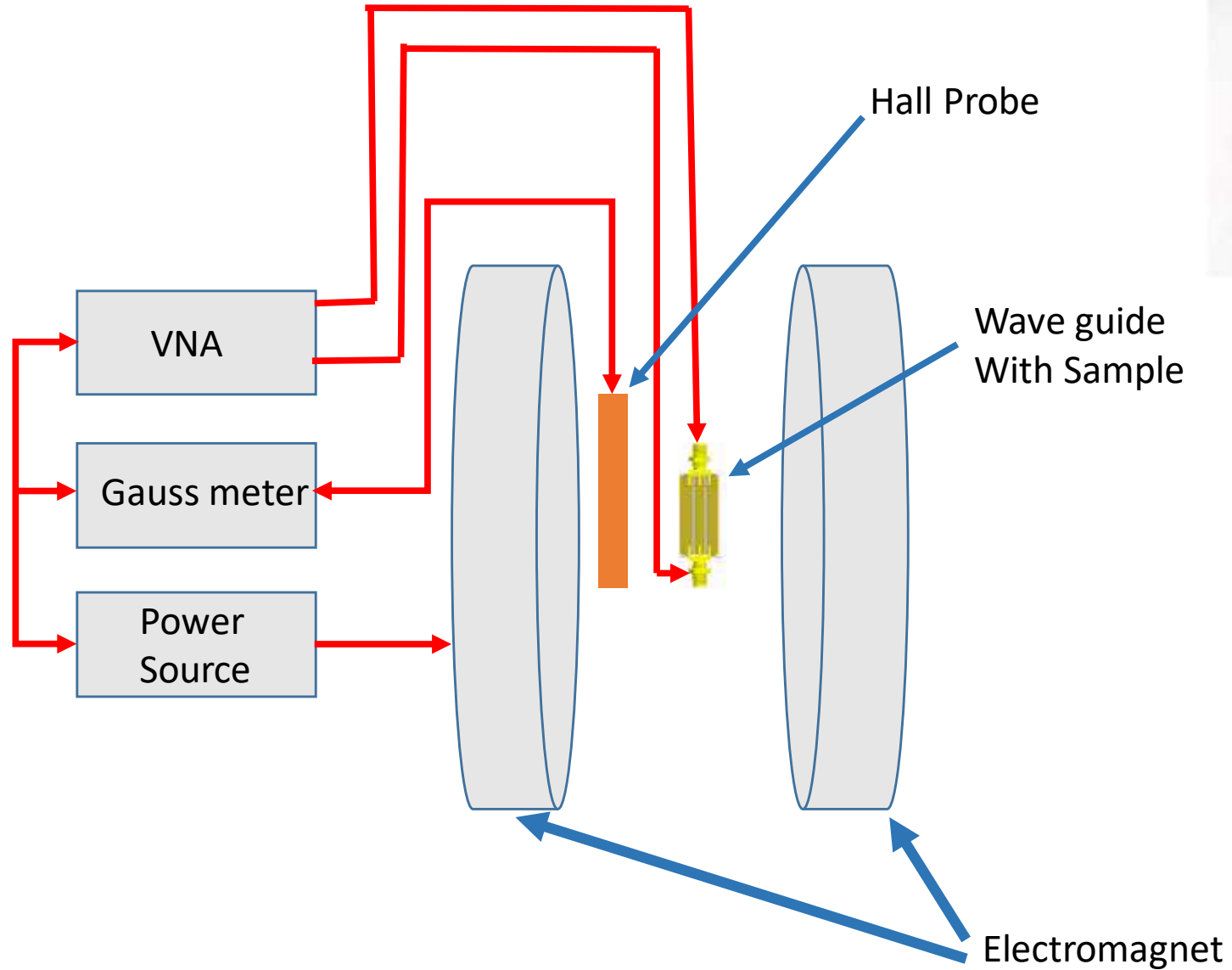


Coplanar Waveguide (inside the Copper Box) with a thin film on it.

# Experimental Setup



VNA (Vector Network Analyzer)



# $Fe_{1-x}V_x$ Thin Films (single Crystalline)

Grown by Molecular Beam Epitaxy (MBE)

by Matthias Kronseder

Al	3nm
$Fe_{1-x}V_x$	20nm
GaAs substrate	

Films under study

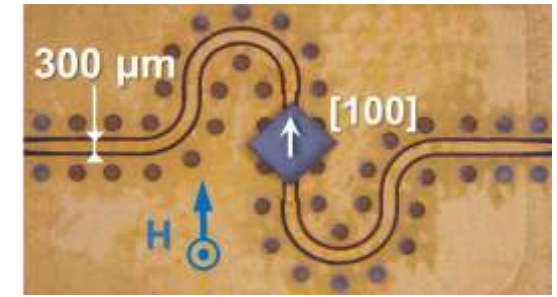
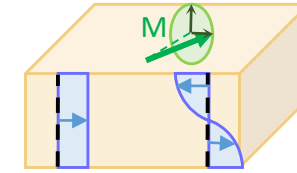
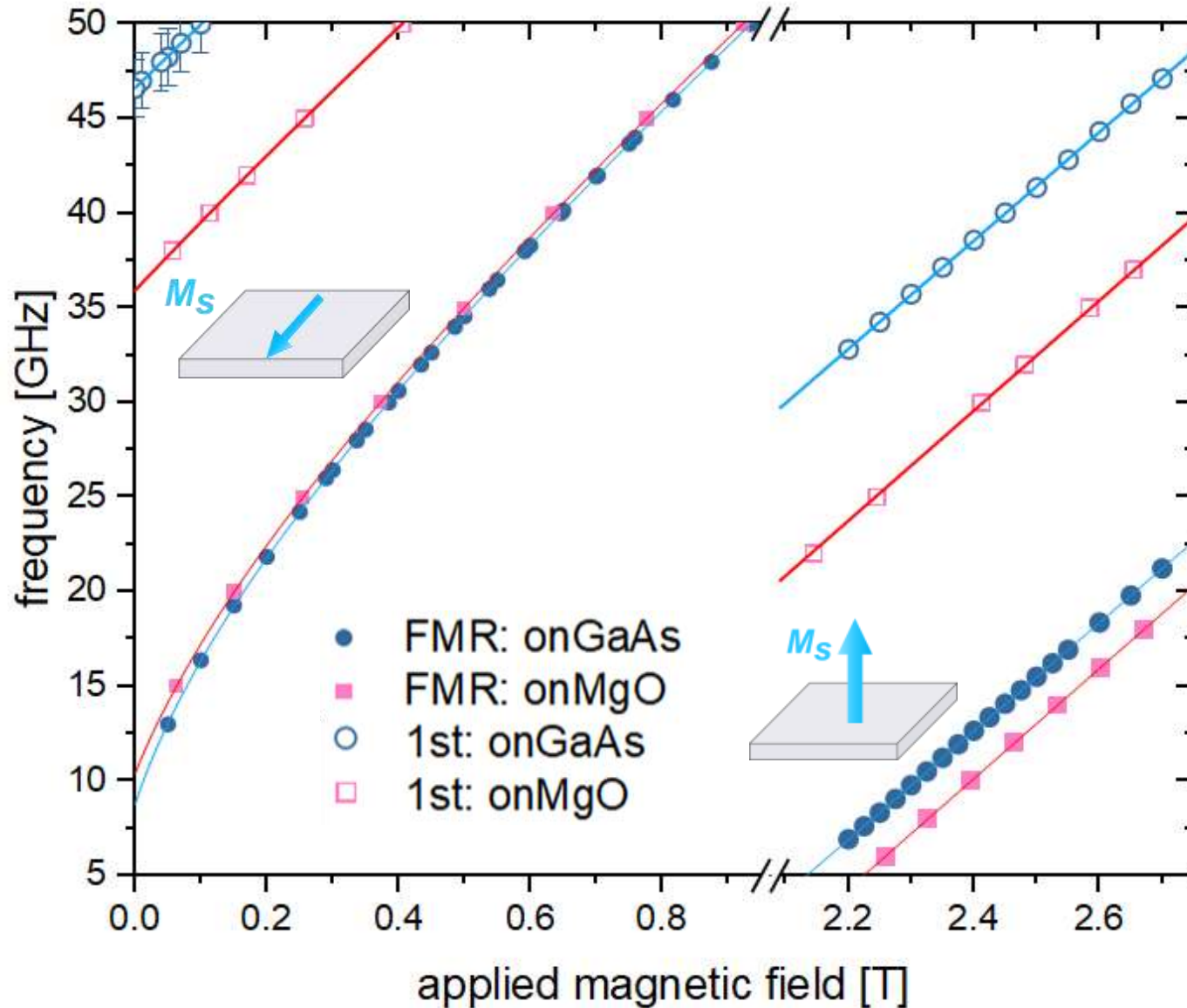
by David Halley (IPCMS)

Ti	4.5nm
MgO	8nm
$Fe_{1-x}V_x$	20nm
MgO	40nm
MgO substrate	

Reference Films



# Fe: FMR



$$f = \frac{\gamma\mu_0}{2\pi} \sqrt{(H_0 + H_x)(H_0 + H_y)}$$

(In Plane)

$$f = \frac{\gamma\mu_0}{2\pi} (H_0 - H_z)$$

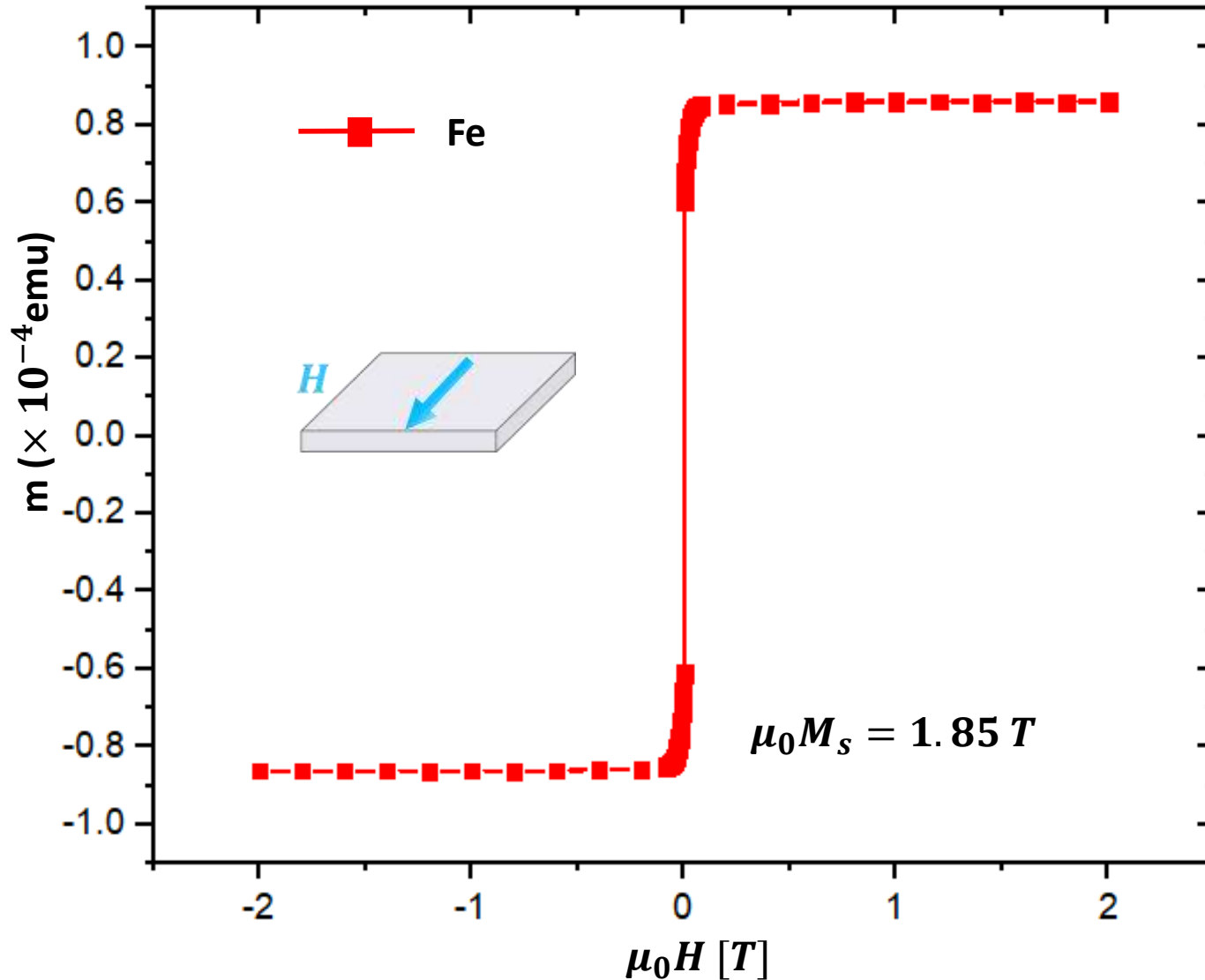
(Out Plane)

		on MgO (Reference)	on GaAs (Measured)	
$H_x$	→ $H_K$ [T]	0.06	0.04	↓
$H_y$	→ $M_{eff}$ [T]	2.07	2.07	-
$H_y, H_z$	→ $H_{ex}$ [T]	0.55	0.85*	↑

- Equal effective magnetization and increased exchange field\* suggest a thickness smaller than expected: [16 nm](#)

$$H_{ex} \propto \frac{1}{M_s t^2}$$

# Fe: SQUID (by Jerome Robert)

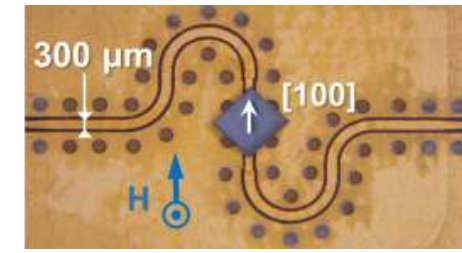
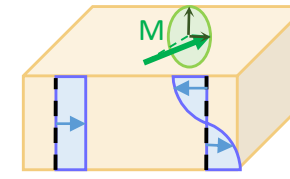
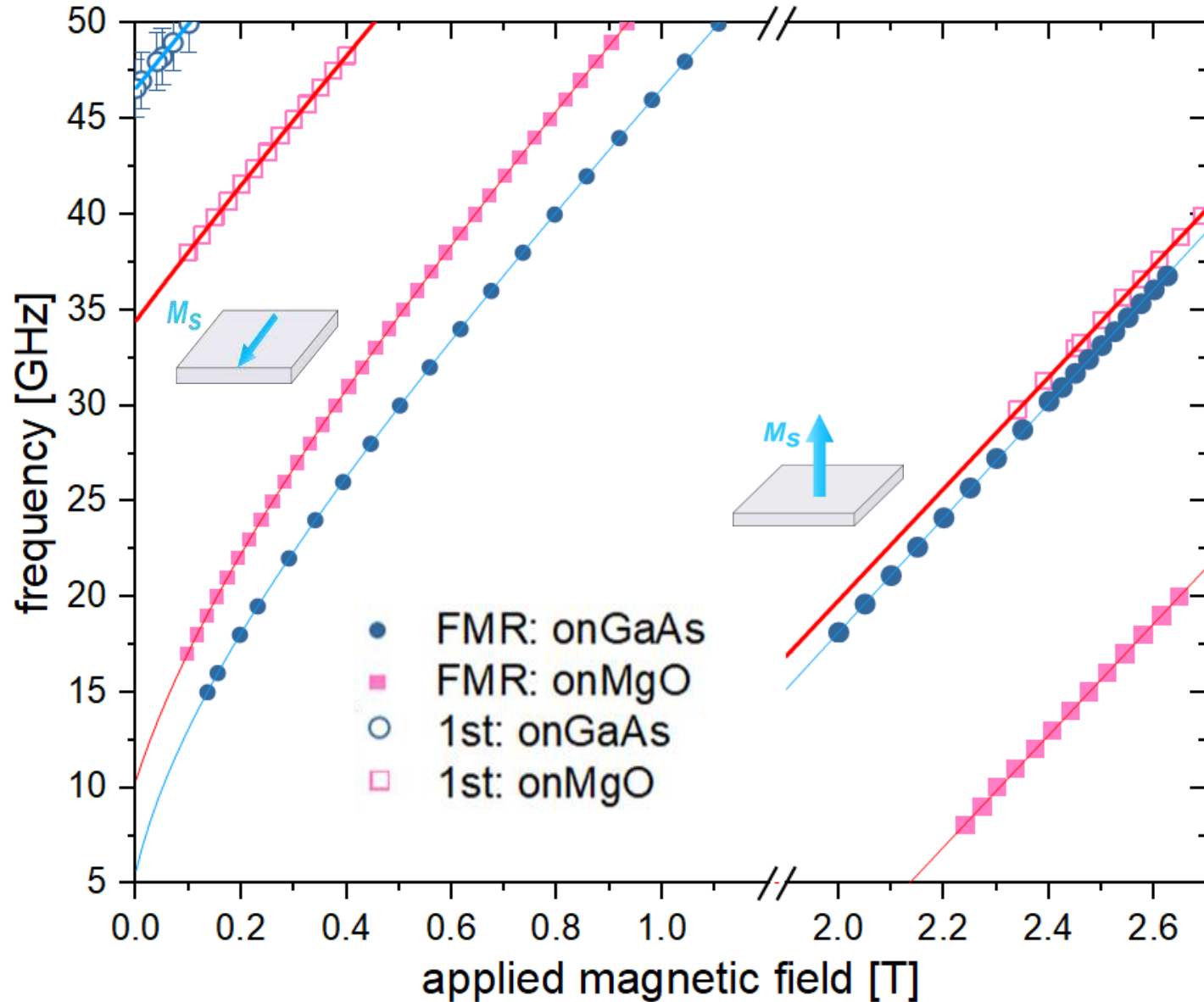


- Thickness = 20 nm is assumed for the calculation of  $M$ .
- **Expected  $M_s$ : 2.15 T**

$$M_s = \frac{\sum m_i}{V}$$

- Thickness smaller than expected?  
17 nm

# FeV: FMR



$$f = \frac{\gamma\mu_0}{2\pi} \sqrt{(H_0 + H_x)(H_0 + H_y)}$$

(In Plane)

$$f = \frac{\gamma\mu_0}{2\pi} (H_0 - H_z)$$

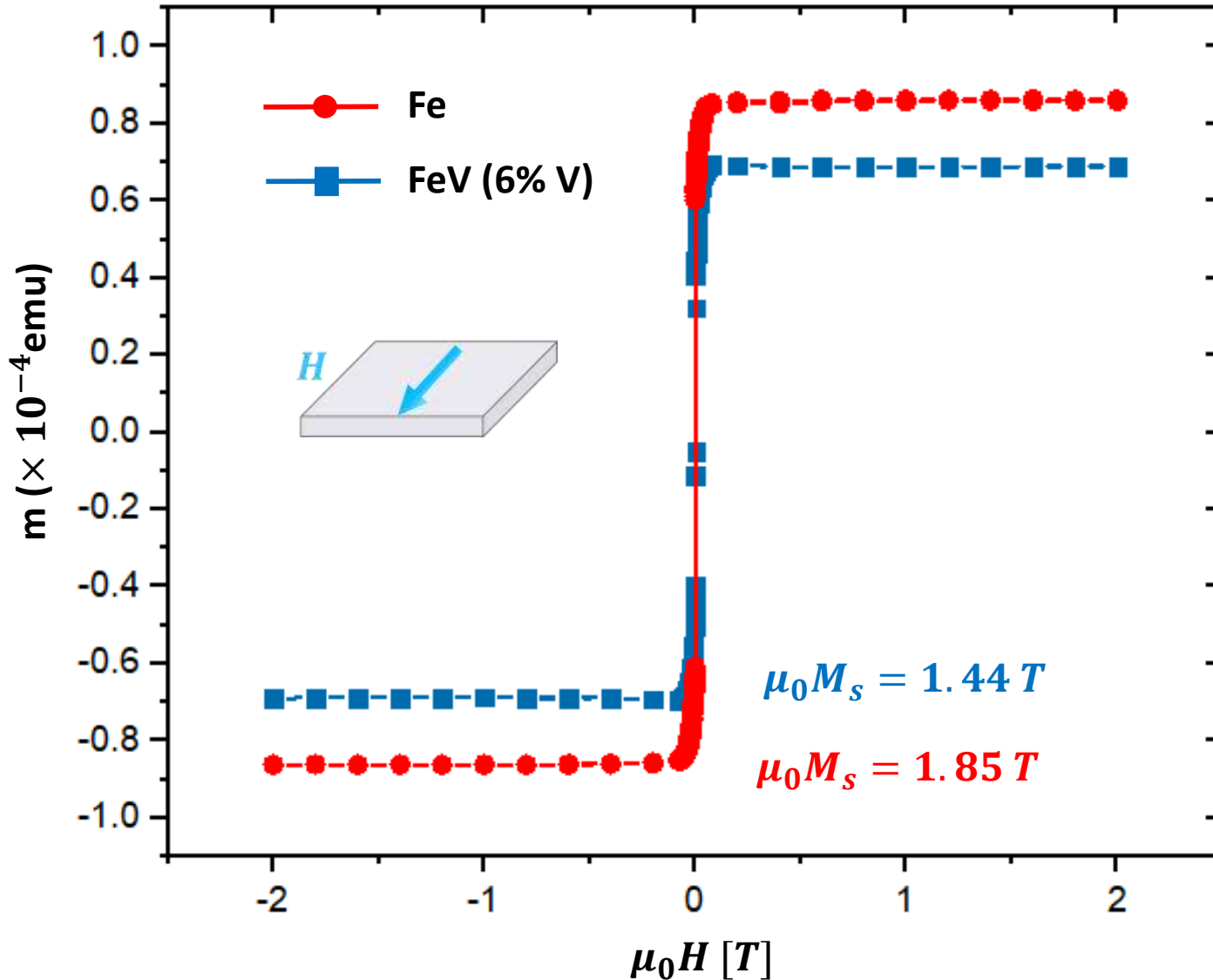
(Out Plane)

		on MgO (Reference)	on GaAs (Measured)	
$H_x$	$\rightarrow H_K$ [T]	0.06	0.03	$\downarrow$
$H_y$	$\rightarrow M_{eff}$ [T]	1.99	1.48	$\downarrow$
$H_y, H_z$	$\rightarrow H_{ex}$ [T]	0.53	0.85*	$\uparrow$

- Smaller effective magnetization and increased exchange field\* suggest:
  - Lower  $M_s$ . (Saturation Magnetization)
  - Or/And thickness smaller than expected: [16 nm](#)

$$H_{ex} \propto \frac{1}{M_s t^2}$$

# FeV (6% V): SQUID (by Jerome Robert)



- Thickness=20 nm is assumed for the calculation of M.
- **Expected  $M_s$ : 1.95 T**  
[Devolder, Appl. Phys. Lett. 103, 242410 (2013)]

$$M_s = \frac{\sum m_i}{V}$$

- 1) Thickness smaller than expected?  
15 nm?
- 2) V concentration higher than targeted?  
20-24%?

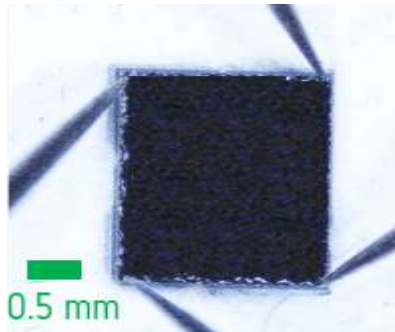
# Resistivity

Room Temperature Van der Pauw measurements,

For Fe,

sample	resistivity [ $\mu\Omega$ cm]		
	MgO (Reference)	GaAs (Measured)	
MC211111C_e4	11.4 – 13.3	16.4	↑

- Actual Thickness = 15 nm ?



For FeV (6%)

sample	resistivity [ $\mu\Omega$ cm]		
	MgO (Reference)	GaAs (Measured)	
MC211116A_a1	25	54.2	↑

- Actual Thickness = 9 nm ?

=> Increased concentration of impurities. (Higher than 6%)

# Conclusions

## For Fe;

1. All measurements agree with an actual thickness  $< 20$  nm.
2. Decrease of cubic anisotropy: An indication of impurities in the film?

## For FeV;

1. In this case, a smaller thickness may not be the only explanation.
2. Measurements also suggest an  $M_s$  smaller than expected:  
⇒ Larger concentration of V impurities.

- ❖ **So this characterization gave clues to that the samples were not in the required condition. Therefore, the characterization continued to understand the origin of the problems**
- ❖ **Now we know the films are actually 17-18 nm thick, and the V concentration is higher than expected (9.5 – 12 %). Also there is some oxygen inside.**
- ❖ **By considering all these things, growth condition need to be checked and Al capping layer should be also improved.**

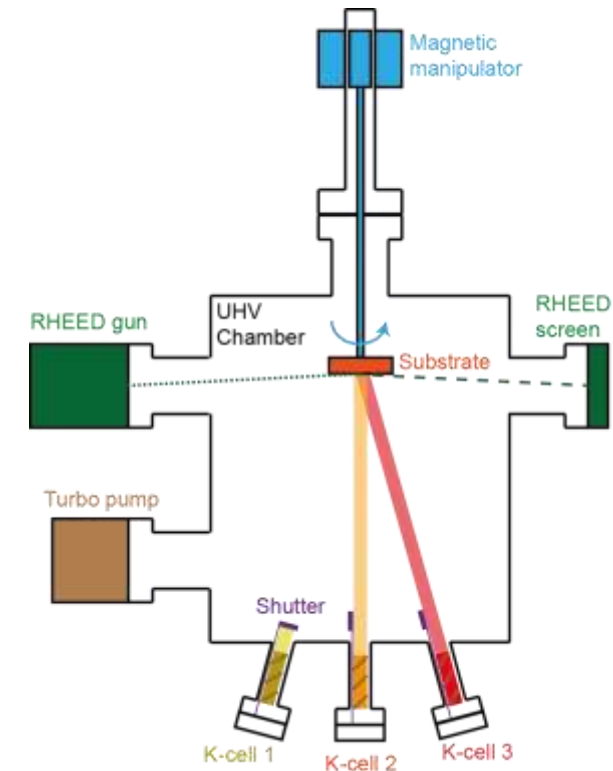
**THANK YOU**

# Backups

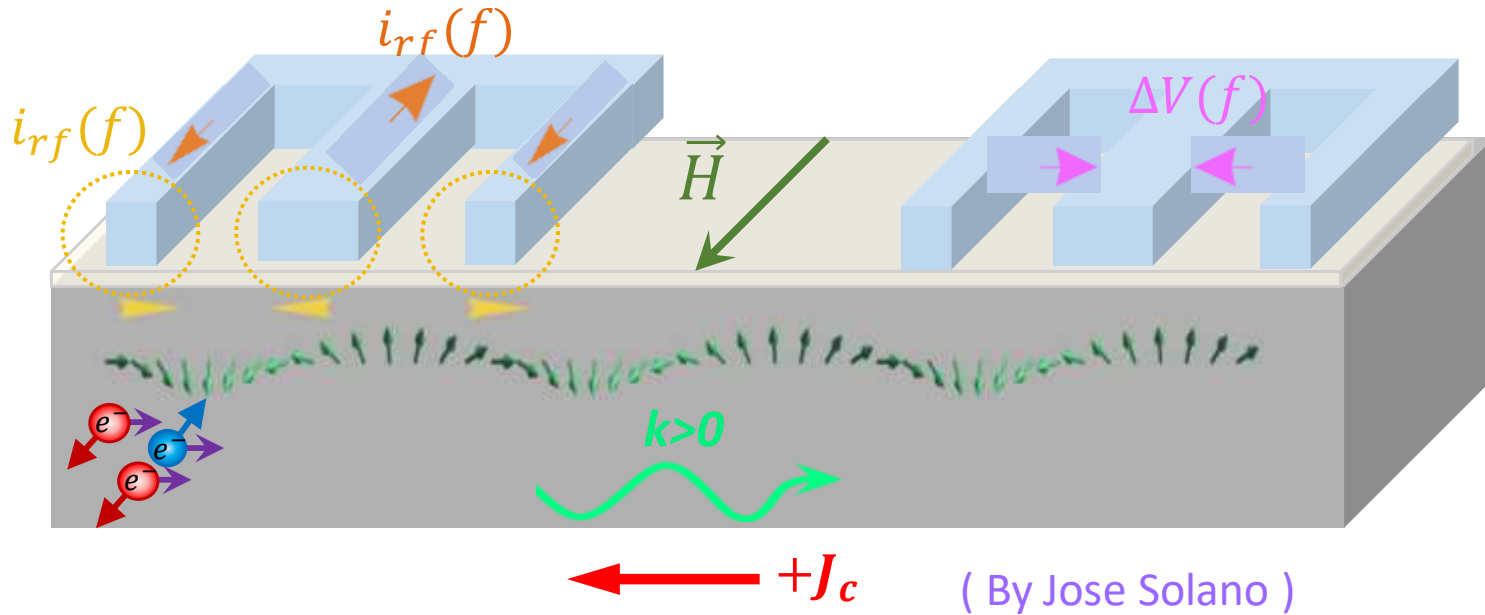


# MBE (Molecular Beam Epitaxy)

- ❖ MBE has evolved into one of the most widely used techniques for producing epitaxial layers of metals, insulators and superconductors as well.
- ❖ it consists essentially of atoms or clusters of atoms, which are produced by heating up a solid source.
- ❖ They then migrate in an UHV environment and impinge on a hot substrate surface, where they can diffuse and eventually incorporate into the growing film.



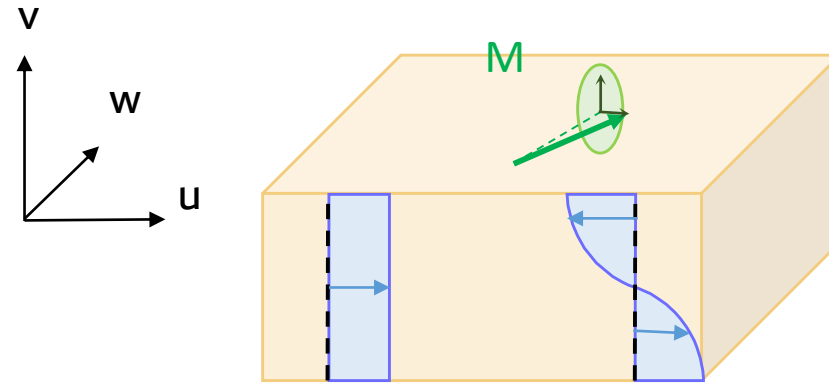
# Application – Spin Wave Doppler Shift



- ❖ Two antennas on top of the sample, exiting with respective magnetic fields.
- ❖ in magnetic materials, there is an unbalance between spin up and spin down electron densities. (two current models)
- ❖ That create an effective magnetic moment.
- ❖ That effective magnetic moment interact with the spin waves and produce a frequency shift.
- ❖ This shift depends on sign of the current or propagation direction of the spin wave.
- ❖ This effect doesn't change the magnitude of the magnetization. Only modifies its frequency.

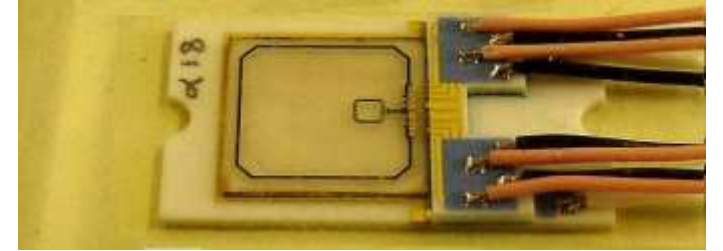
$$\frac{\partial \vec{M}}{\partial t} = -\gamma \mu_0 \vec{M} \times \vec{H} - u \frac{d\vec{M}}{dx}$$

# Spin Wave Modes

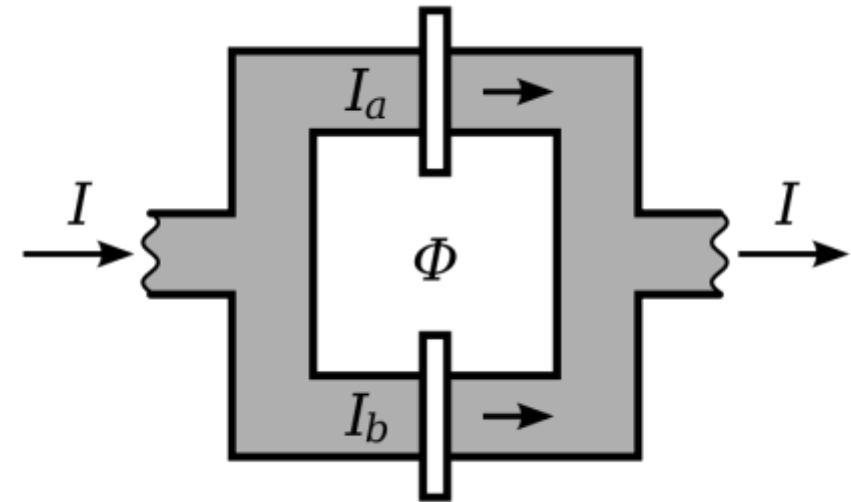


- ❖ They are thickness modes. Because of constrain thickness.
- ❖ Confinement of the magnetization oscillations leads to a discretization of the energy levels of the spin waves (spin wave modes become distinguishable)

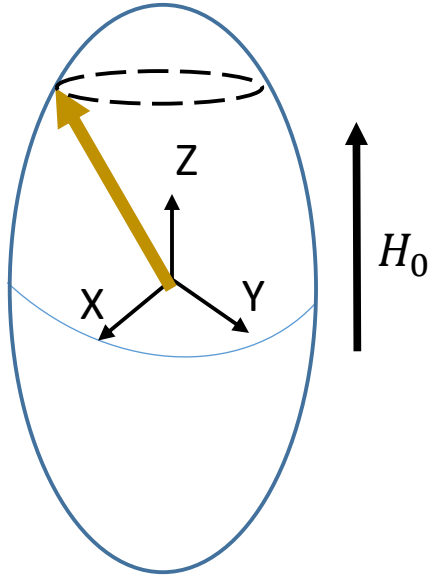
# SQUID (superconducting quantum interference device)



- ❖ The most sensitive magnetic flux detector is the superconducting quantum interference device SQUID.
- ❖ Contains two Josephson junctions (insulators) between two superconducting Wires.
- ❖ Classically, current not conducting through this.
- ❖ But in quantum mechanical limit, there is a probability for tunneling.
- ❖ It depends on temperature and amount of magnetic moments.



# Derivation of resonance frequency Equations



$$\frac{d\mathbf{M}}{dt} = -\gamma \mathbf{M} \times \mu_0 \mathbf{H}_{eff} \quad \longrightarrow \quad \text{Landau Lifshitz equation}$$

$$\vec{h} = \overline{\chi}_k^{-1} \vec{m} = \frac{1}{\gamma \mu_0 M_s} \begin{pmatrix} \omega_x & -i\omega \\ i\omega & \omega_y \end{pmatrix} \vec{m}$$

$$\omega_x = \gamma \mu_0 (H_0 + (N_x - N_z) M_s)$$

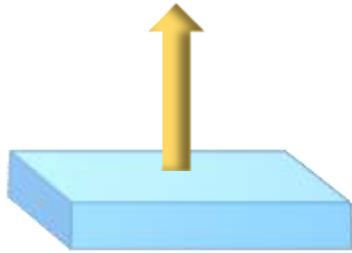
$$\omega_y = \gamma \mu_0 (H_0 + (N_y - N_z) M_s)$$

$$\omega = \sqrt{\omega_x \omega_y}$$

Kittel Formula

$$\omega_{res} = \gamma \mu_0 \left\{ [(H_0 + (N_x - N_z) M_s) \gamma \mu_0 (H_0 + (N_y - N_z) M_s)] \right\}^{1/2}$$

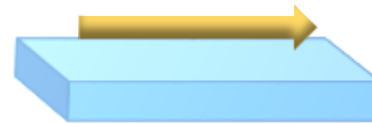
Out Plane



$$N_x = N_y = 0 \quad N_z = 1$$

$$\omega = \gamma \mu_0 (H_0 - M_s)$$

In Plane

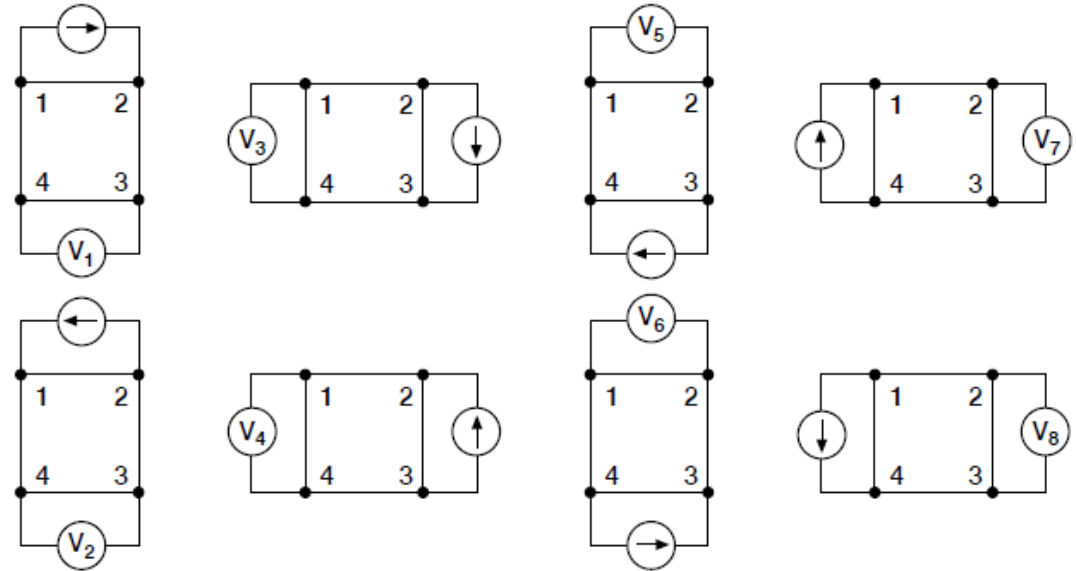
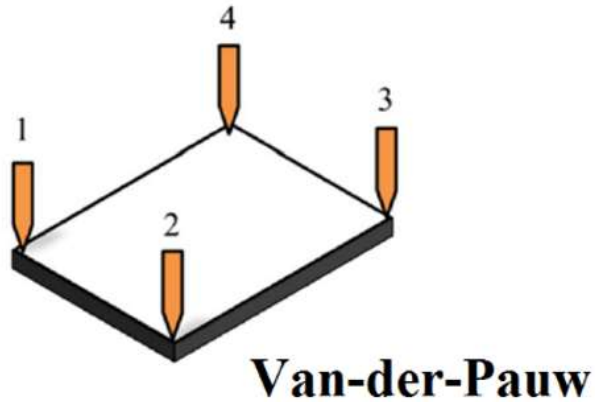


$$N_y = N_z = 0 \quad N_x = 1$$

$$\omega = \gamma \mu_0 \sqrt{H_0 (H_0 + M_s)}$$

# Van der Pauw Resistivity Measurement Method

- ❖ The van der Pauw method involves applying a current and measuring voltage using four small contacts on the perimeter of a flat, arbitrarily shaped sample of uniform thickness



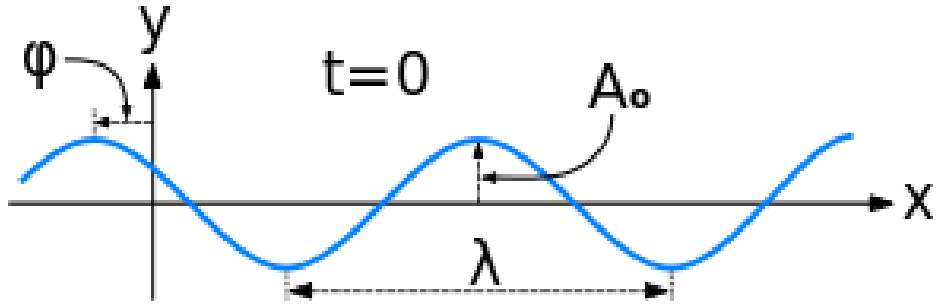
$$\rho_A = \frac{\pi}{\ln 2} f_A t_s \frac{(V_1 - V_2 + V_3 - V_4)}{4I}$$

$$\rho_B = \frac{\pi}{\ln 2} f_B t_s \frac{(V_5 - V_6 + V_7 - V_8)}{4I}$$

$$\rho_{AVG} = \frac{\rho_A + \rho_B}{2}$$

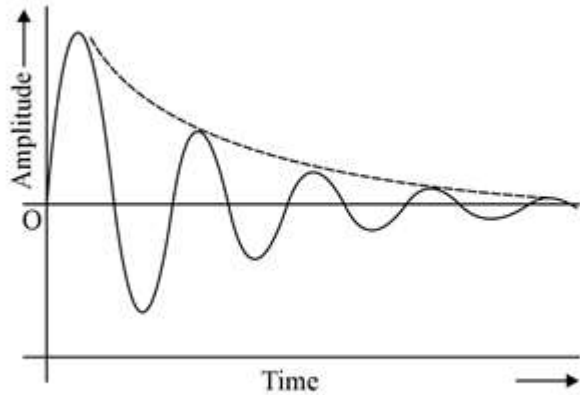
\*  $f_A, f_B$  = Geometrical Factors (Based on sample geometry)  
 ( $f_A = f_B = 1$  for perfect symmetry)

# Plane Waves



$$U(r, t) = Ae^{i(k.x - \omega t)}$$

$$U(r, t) = A\cos(k.x - \omega t) + iA\sin(k.x - \omega t)$$



$$U(r, t) = Ae^{-\lambda}e^{i(k.x - \omega t)}$$

## Landau Lifshitz Equation

- ❖ There is no damping term
- ❖ Signal width is zero

$$\frac{dM}{dt} = -\gamma M \times \mu_0 H_{eff}$$

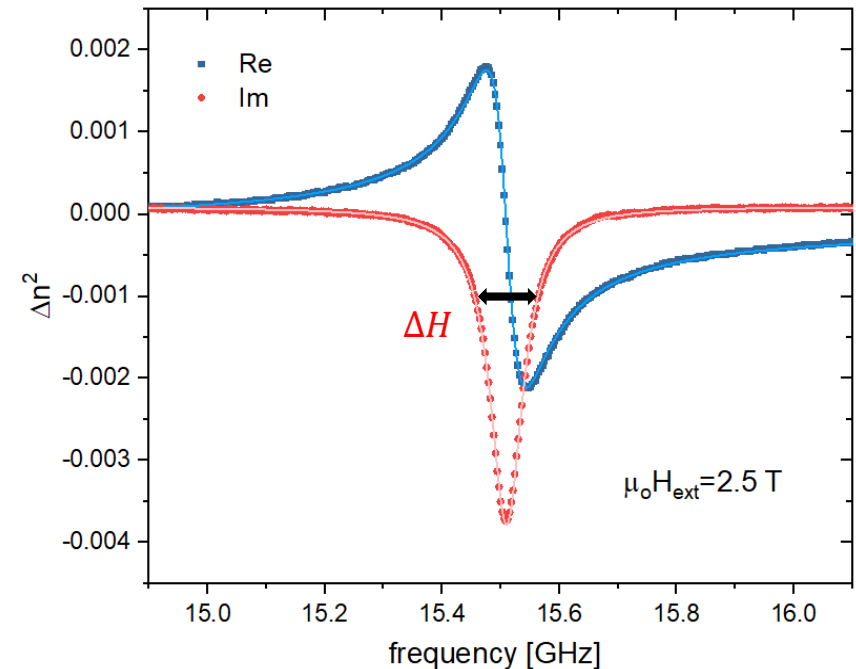
$$\omega = \gamma \mu_0 (H_0 - M_s)$$

$$\omega = \gamma \mu_0 \sqrt{H_0(H_0 + M_s)}$$

## Landau Lifshitz Gilbert equation

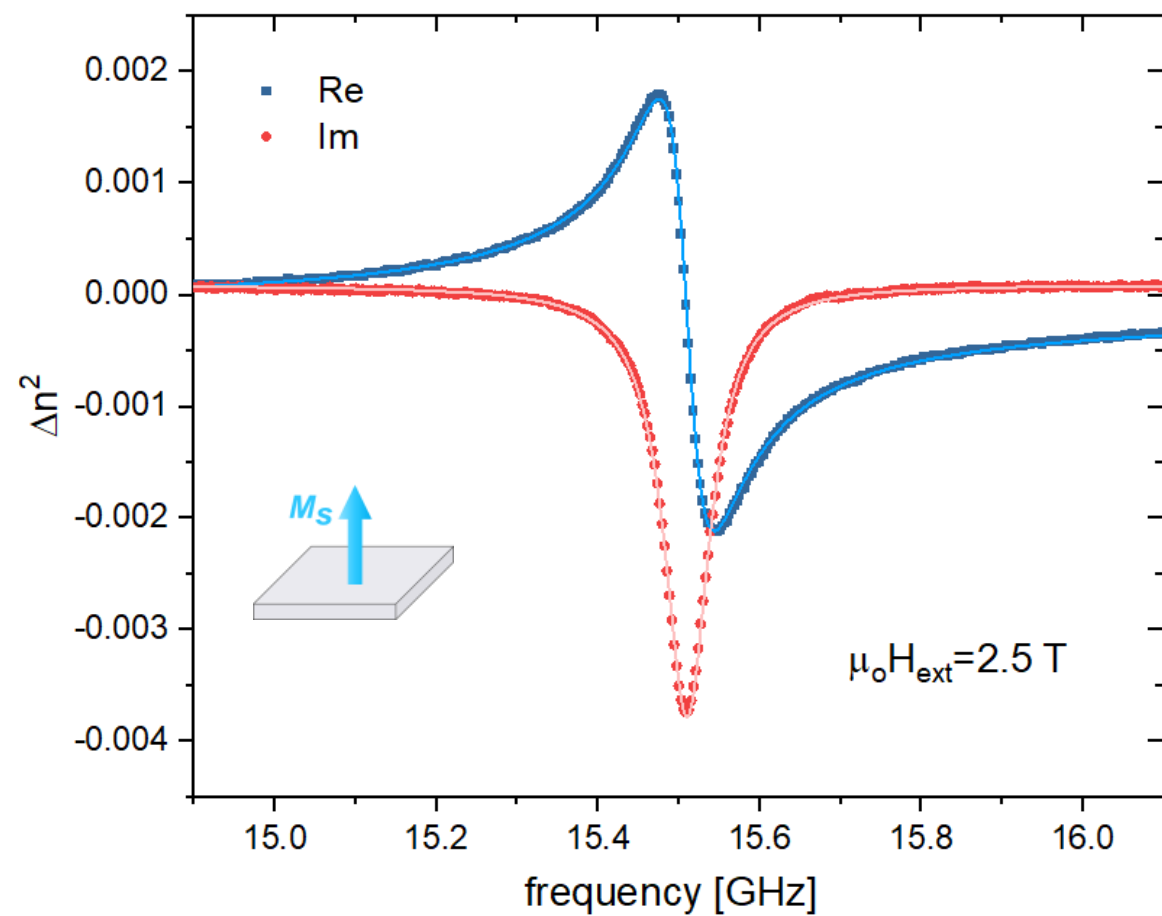
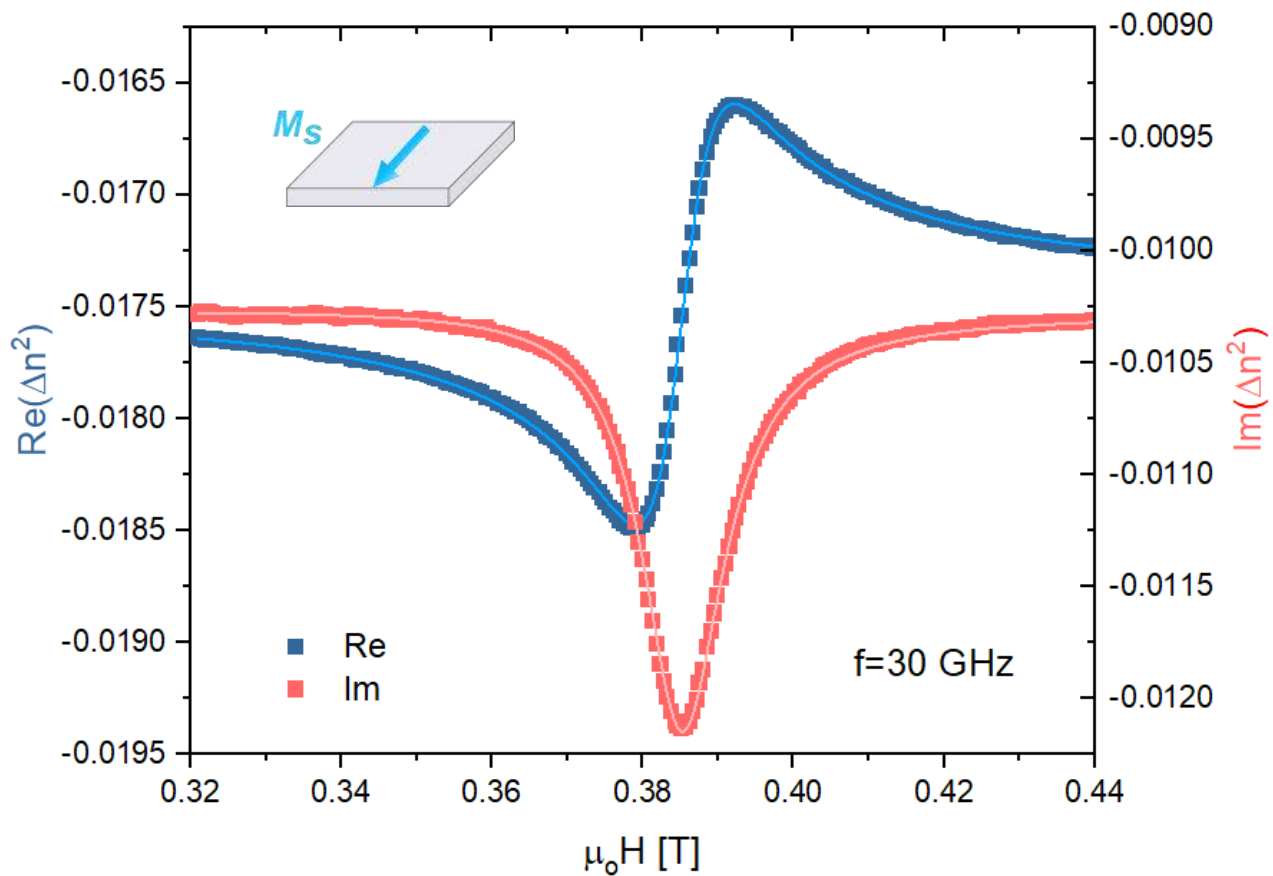
- ❖ The Damping term increase the width of the signal due to the dispersion of absorption.

$$\frac{dM}{dt} = -\gamma M \times \mu_0 H_{eff} + \alpha M \times \frac{dM}{dt}$$

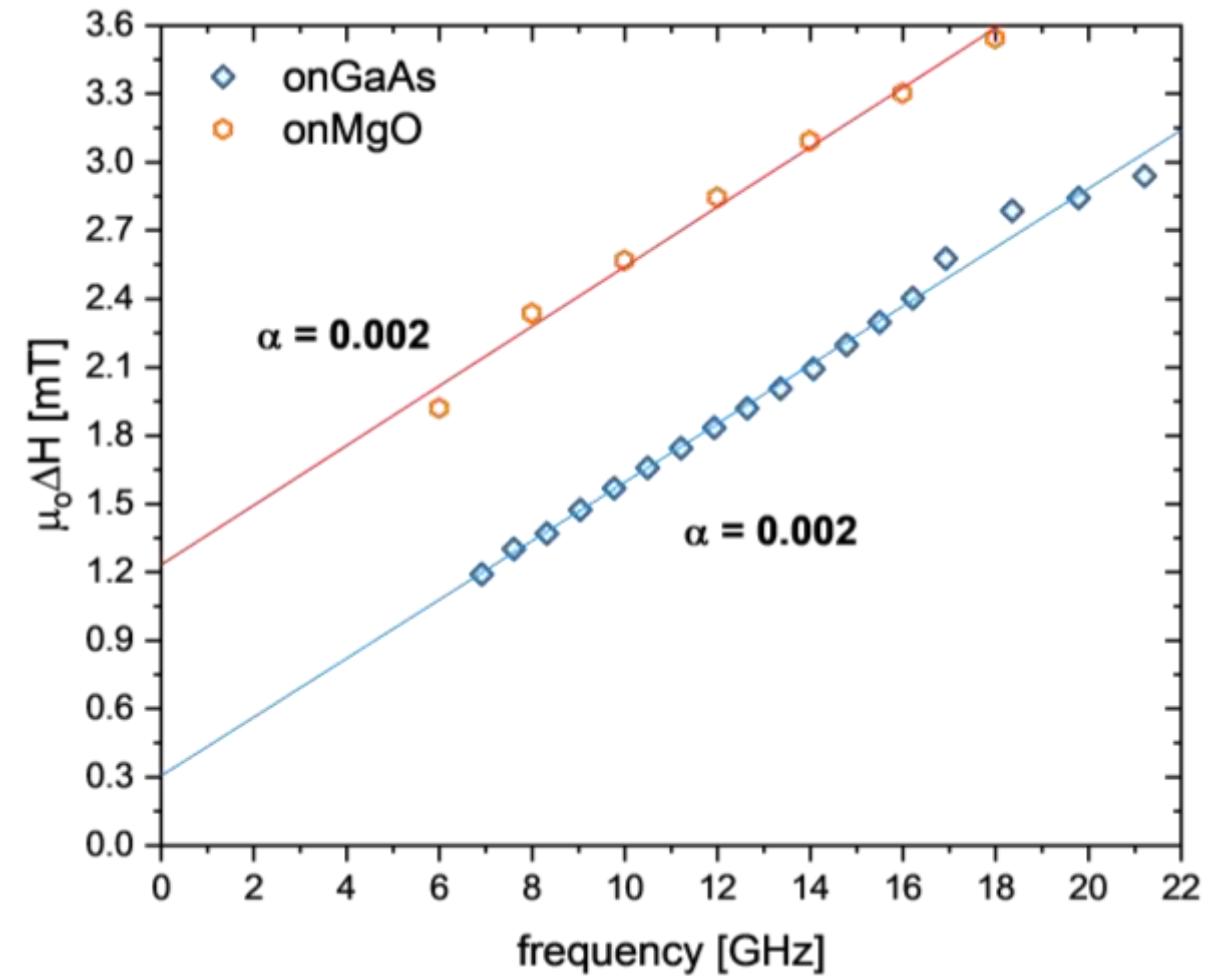
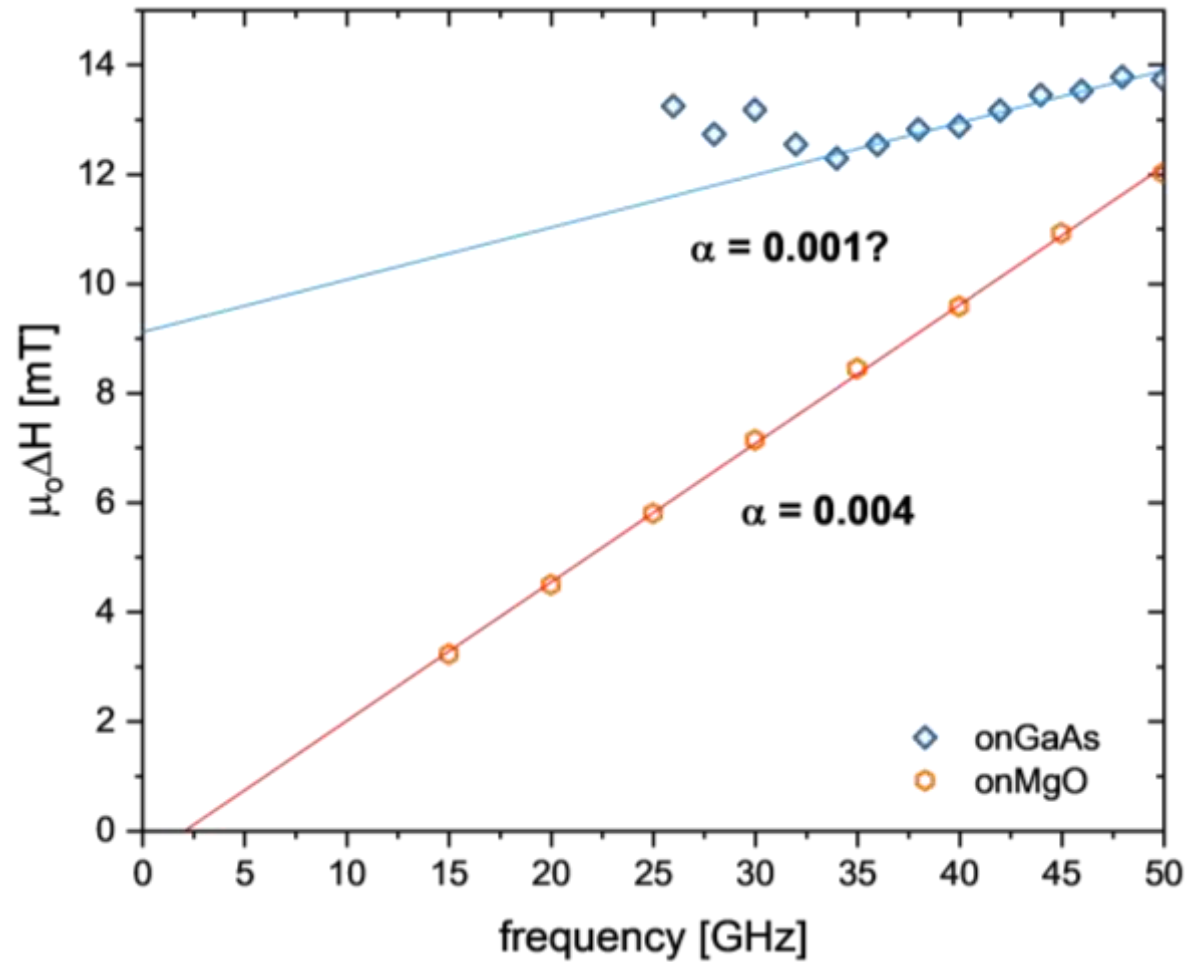




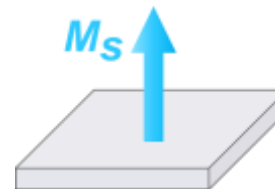
# Resonance Peaks



# Fe : Linewidth

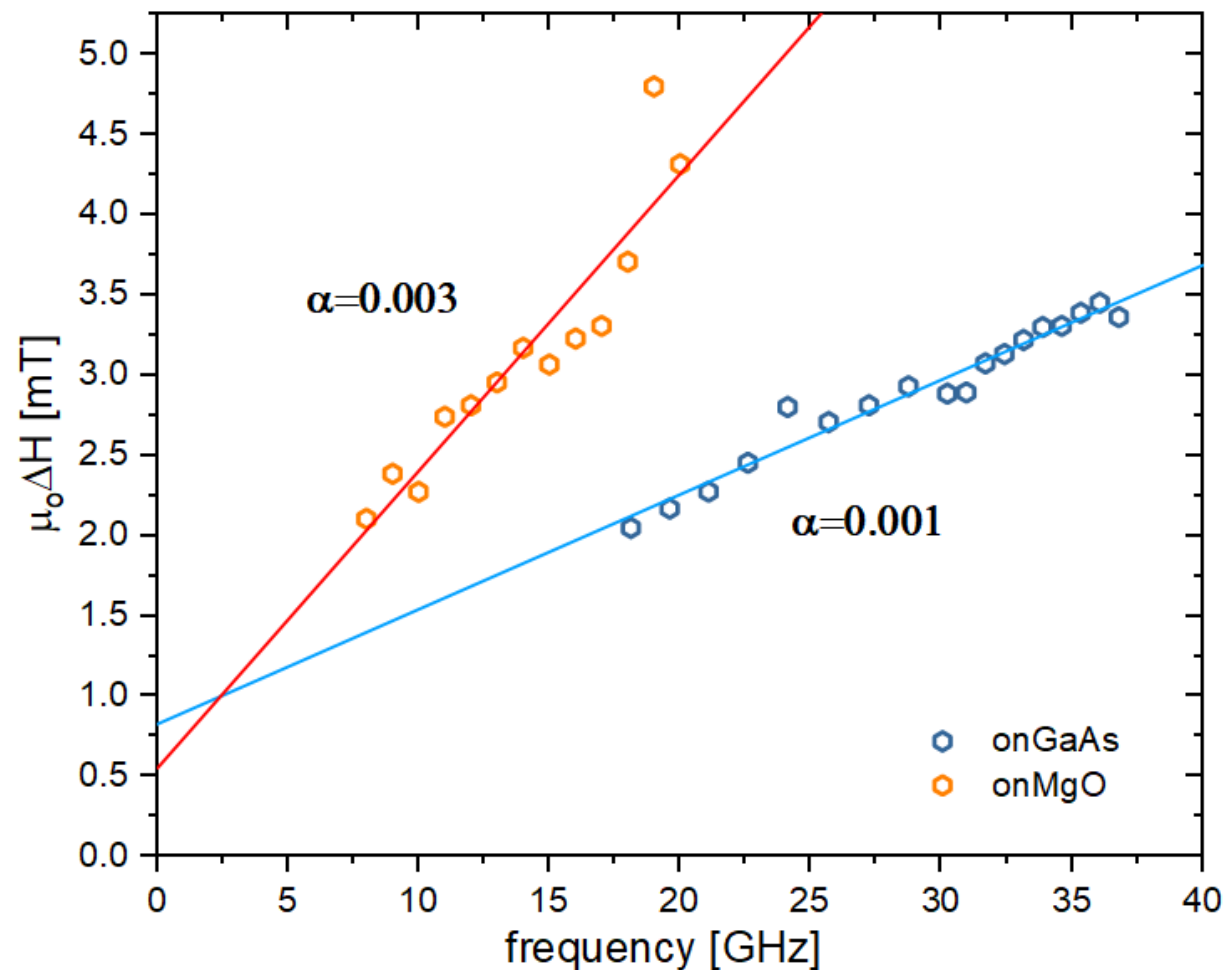
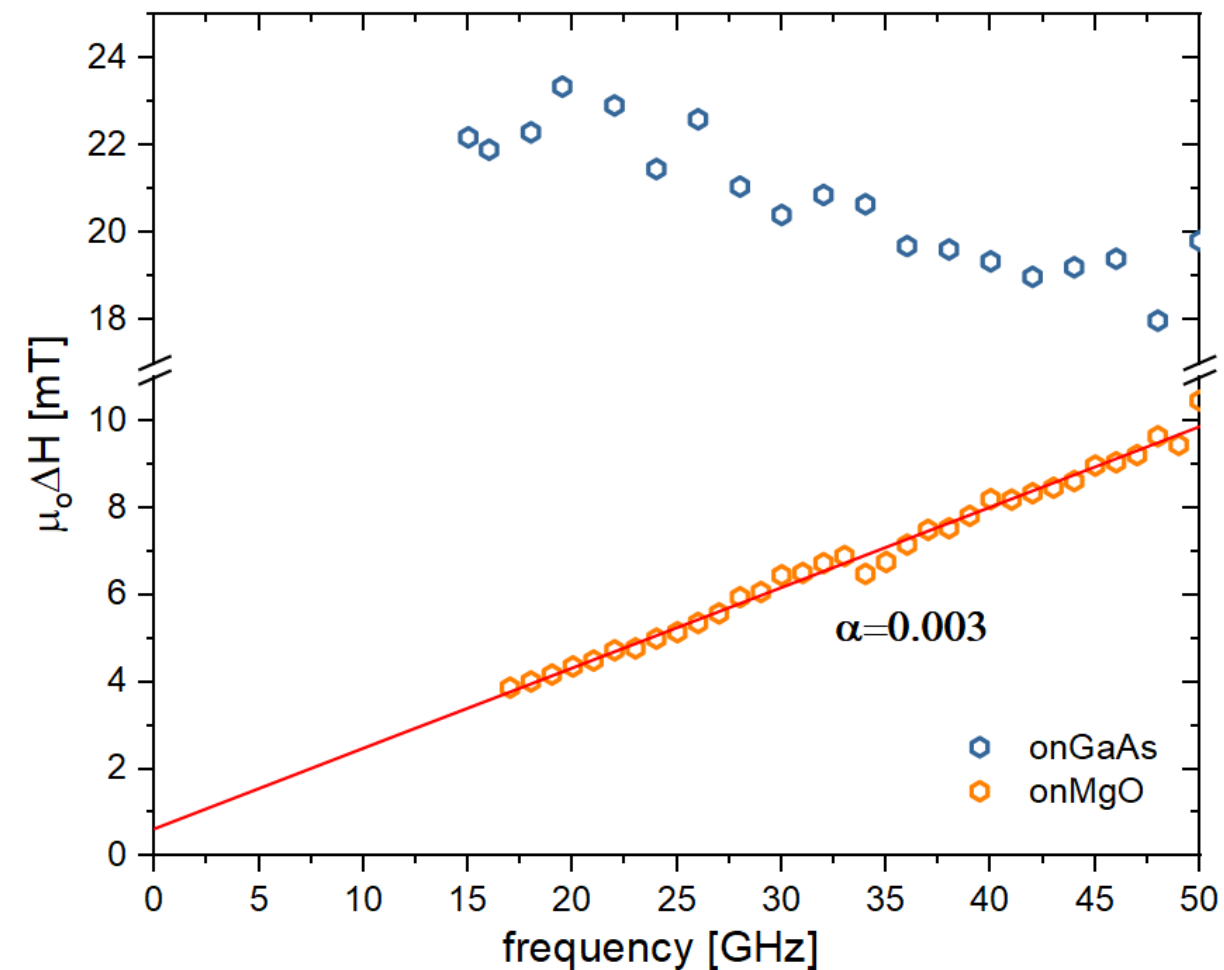


*MgO* -> *GaAs*  
 $\alpha$  0.004 -> 0.001?  
 $\Delta H_0$  -0.5 mT -> 9 mT



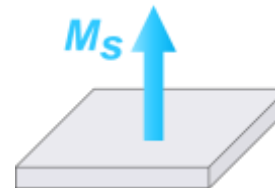
*MgO* -> *GaAs*  
 $\alpha$  0.002 -> 0.002  
 $\Delta H_0$  1.2 mT -> 0.3 mT

# FeV : Linewidth



*MgO* -> *GaAs*

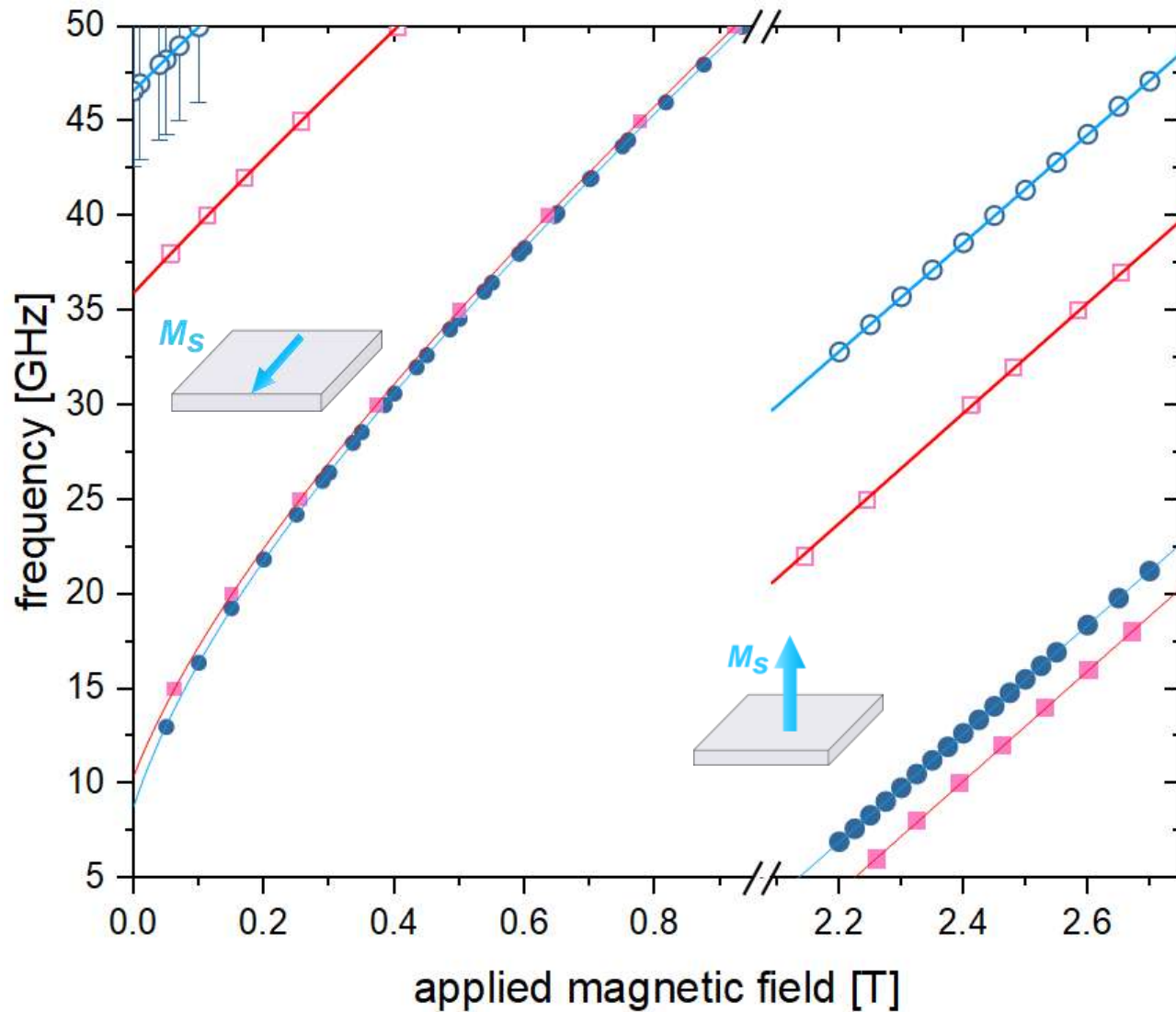
$\alpha$	0.003	->	-
$\Delta H_0$	0.6 mT	->	-



*MgO* -> *GaAs*

$\alpha$	0.003	->	0.001
$\Delta H_0$	0.5 mT	->	0.8 mT

# Fe: FMR



	on MgO	on GaAs	
$\gamma [GHz \cdot T^{-1}]$	29.1	28.6	↓
$H_K [T]$	0.06	0.04	↓
$M_{eff} [T]$	2.07	2.07	-
$H_{ex} [T]$	0.55	0.85*	↑
$H_{ex} + H_S [T]$	0.67	0.91	↑

*In Plane FMR: MgO -> GaAs*

$H_x = H_K$  0.06 T -> 0.04 T

$H_y = M_s + H_K - H_u - H_S$  2.08 T -> 2.18 T

*Out Plane FMR: MgO -> GaAs*

$H_z = M_s - H_K - H_u - H_S$  2.05 T -> 1.96 T

*In Plane :Second mode MgO -> GaAs*

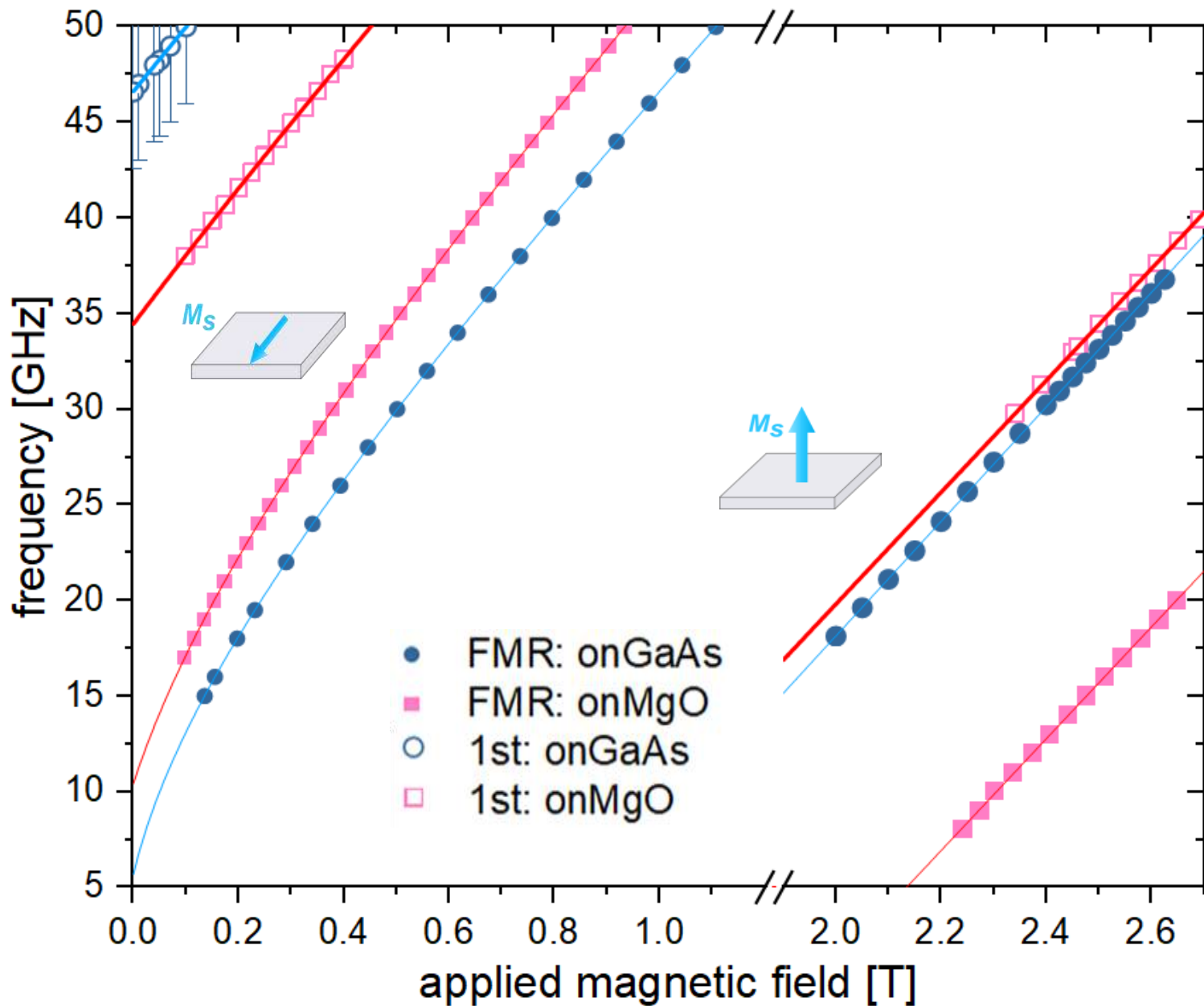
$H_K + H_{Ex}$  0.61 T -> 0.89 T

$M_s + H_K - H_u - 2H_S + H_{Ex}$  2.52 T -> 2.98 T

*Out Plane :Second mode MgO -> GaAs*

$M_s - H_K - H_u - 2H_S - H_{Ex}$  1.38 T -> 1.05 T

# FeV: FMR



	on MgO	on GaAs	
$\gamma$ [ $GHz \cdot T^{-1}$ ]	29.2	29.3	-
$H_K$ [T]	0.06	0.03	↓
$M_{eff}$ [T]	1.99	1.48	↓
$H_{ex}$ [T]	0.53	0.85*	↑

*In Plane FMR: MgO → GaAs*

$H_K$  0.06 T → 0.03 T  
 $M_S + H_K - H_u - H_S$  2.02 T → 1.57 T

*Out Plane FMR: MgO → GaAs*

$M_S - H_K - H_u - H_S$  1.96 T → 1.39 T

*In Plane :Second mode MgO → GaAs*

$H_K + H_{Ex}$  0.59 T → 0.89 T  
 $M_S + H_K - H_u - 2H_S + H_{Ex}$  2.39 T → 2.95 T

*Out Plane :Second mode MgO → GaAs*

$M_S - H_K - H_u - 2H_S - H_{Ex}$  1.32 T → -