

# Geometrical Aspects of Supergravity

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2021

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# From Supersymmetry to Supergravity

- Haag et al. :  $\mathfrak{g} = \mathfrak{iso}(1, 3) \oplus \mathfrak{g}_c \oplus \{Q_\alpha, \bar{Q}^{\dot{\alpha}}\}$ .
- Supersymmetry transformations generated by **spin- $\frac{1}{2}$**  operators  $Q_\alpha$  and  $\bar{Q}^{\dot{\alpha}}$ .
- **Anticommuting** parameters :  $\epsilon$  and  $\bar{\epsilon}$ .

## Composition of two Susy transformations :

$$[\epsilon^\alpha Q_\alpha, \bar{\epsilon}_{\dot{\alpha}} \bar{Q}^{\dot{\alpha}}] = 2i\epsilon^\alpha \sigma^\mu_{\alpha\dot{\alpha}} \bar{\epsilon}^{\dot{\alpha}} P_\mu \Rightarrow \text{space time translation.}$$

- Promoting **global** Susy to **local** :  $\epsilon \rightarrow \epsilon(x)$  :
  - ★ Composition of two Susy transformations generates a **Local translation**.
  - ★ Local Supersymmetry  $\equiv$  Supergravity

# The Equivalence Principle

## Equivalence Principle

The **local** equivalence of inertial and gravitational force

- In the instantaneous local referential frame  $\mathcal{R}_0(p)$  where gravity is eliminated : we recover **local Lorentz symmetry** !
- Span a **Flat tangent superspace** at each point of the curved superspace

# Dynamical Gauge Superfields

- **Superconnection** :  $\Omega_{\tilde{M}\tilde{M}}^N \rightarrow$ 
  - ★  $\mathcal{L}_{\text{Sugra}}$  Invariance under local Lorentz transformations
  - ★ Covariant derivatives :  $\mathcal{D}_{\tilde{M}}X^M = \partial_{\tilde{M}}X^M + \Omega_{\tilde{M}N}^M X^N$
- **Vielbein** :  $E_{\tilde{M}}^M \rightarrow$ 
  - ★  $\mathcal{L}_{\text{Sugra}}$  Invariance under superdiffeomorphism
  - ★  $E_{\tilde{M}}^M$  Curved superspace  $\rightarrow$  Flat tangent space
  - ★  $E_M^{\tilde{M}}$  Flat tangent space  $\rightarrow$  Curved superspace

## Example

$$\mathcal{D}_{\tilde{M}}X^M \rightarrow \mathcal{D}_N X^M = E_N^{\tilde{M}} \mathcal{D}_{\tilde{M}}X^M$$

# The Covariant Derivatives, the generators of Supergravity

## Generators of the Supergravity superalgebra

$$[\mathcal{D}_M, \mathcal{D}_N]_{|M||N|} = T_{MN}{}^Q \mathcal{D}_Q + \frac{1}{2} R_{MN\alpha\beta} J^{\beta\alpha}$$

- Torsion and Curvature tensor  $\propto E_{\tilde{M}}{}^M$  and  $\Omega_{\tilde{M}M}{}^N$
- Closes under Superdiffeomorphism + Lorentz group transformation

# The Bianchi Identities

- Bianchi Identities :

$$\begin{aligned} 0 = & (-)^{|M_1||M_3|} [\mathcal{D}_{M_1}, [\mathcal{D}_{M_2}, \mathcal{D}_{M_3}]_{|M_2||M_3|}]_{|M_1|(|M_2|+|M_3|)} \\ & + (-)^{|M_2||M_1|} [\mathcal{D}_{M_2}, [\mathcal{D}_{M_3}, \mathcal{D}_{M_1}]_{|M_3||M_1|}]_{|M_2|(|M_3|+|M_1|)} \\ & + (-)^{|M_3||M_2|} [\mathcal{D}_{M_3}, [\mathcal{D}_{M_1}, \mathcal{D}_{M_2}]_{|M_1||M_2|}]_{|M_3|(|M_1|+|M_2|)} \end{aligned}$$

- $M_i = \alpha, \dot{\alpha}, \mu, i=1,2,3 \rightarrow 40$  Bianchi identities

# The Supergravity Constrains

## Violation of the Equivalence Principle

- $\exists \mathcal{R}$  such that  $\Omega_{\tilde{M}\tilde{M}}^N = 0 \Rightarrow \emptyset$  Curvature
- $[\mathcal{D}_M, \mathcal{D}_N]_{|M||N|} = T_{MN}^Q \mathcal{D}_Q \neq \text{Susy Algebra}$
- $T_{MN}^Q$  violates the Equivalence Principle !

## Supergravity Constrains

$$\begin{aligned}T_{\alpha\dot{\alpha}}^{\mu} &= T_{\dot{\alpha}\alpha}^{\mu} = -2i\sigma^{\mu}_{\alpha\dot{\alpha}}, \\T_{\underline{\alpha}\underline{\beta}}^{\underline{\gamma}} &= 0, & T_{\alpha\beta}^{\mu} &= T_{\dot{\alpha}\dot{\beta}}^{\mu} = 0, \\T_{\underline{\alpha}\mu}^{\nu} &= T_{\mu\underline{\alpha}}^{\nu} = 0, & T_{\mu\nu}^{\rho} &= 0.\end{aligned}$$



# Bianchi Identities + Supergravity constrains $\rightarrow$ Reduction of the Bianchi identities :

- (1)  $I_{\dot{\alpha}\dot{\beta}\mu\nu} : R_{\dot{\alpha}\dot{\beta}\mu\nu} = 2i\sigma_{\nu\gamma\dot{\alpha}} T_{\dot{\beta}\mu}{}^{\gamma} + 2i\sigma_{\nu\gamma\dot{\beta}} T_{\dot{\alpha}\mu}{}^{\gamma},$
- (2)  $I_{\alpha\beta\dot{\gamma}\dot{\delta}} : R_{\alpha\beta\dot{\gamma}\dot{\delta}} = 2i\sigma^{\mu}{}_{\beta\dot{\gamma}} T_{\alpha\mu\dot{\delta}} + 2i\sigma^{\mu}{}_{\alpha\dot{\gamma}} T_{\beta\mu\dot{\delta}},$
- (3)  $I_{\dot{\alpha}\dot{\beta}\dot{\gamma}\dot{\delta}} : R_{\dot{\alpha}\dot{\beta}\dot{\gamma}\dot{\delta}} + R_{\dot{\beta}\dot{\gamma}\dot{\alpha}\dot{\delta}} + R_{\dot{\gamma}\dot{\alpha}\dot{\beta}\dot{\delta}} = 0,$
- (4)  $I_{\alpha\dot{\beta}\gamma\delta} : R_{\alpha\dot{\beta}\gamma\delta} + R_{\dot{\beta}\gamma\alpha\delta} = 2i\sigma^{\mu}{}_{\alpha\dot{\beta}} T_{\gamma\mu\delta} + 2i\sigma^{\mu}{}_{\gamma\dot{\beta}} T_{\alpha\mu\delta},$
- (5)  $I_{\alpha\dot{\beta}\mu\nu} : R_{\alpha\dot{\beta}\mu\nu} = 2i\sigma_{\nu\gamma\dot{\beta}} T_{\alpha\mu}{}^{\gamma} - 2i\sigma_{\nu\alpha\dot{\gamma}} T_{\dot{\beta}\mu}{}^{\dot{\gamma}}$
- (6)  $I_{\dot{\alpha}\dot{\beta}\mu\gamma} : \bar{\mathcal{D}}_{\dot{\alpha}} T_{\dot{\beta}\mu\gamma} + \bar{\mathcal{D}}_{\dot{\beta}} T_{\dot{\alpha}\mu\gamma} = 0,$
- (7)  $I_{\alpha\mu\nu\rho} : R_{\alpha\mu\nu\rho} + R_{\nu\alpha\rho\mu} = 2i\sigma_{\rho\alpha\dot{\beta}} T_{\mu\nu}{}^{\dot{\beta}},$
- (8)  $I_{\alpha\mu\beta\gamma} : R_{\alpha\mu\beta\gamma} + R_{\beta\mu\alpha\gamma} = \mathcal{D}_{\alpha} T_{\beta\mu\gamma} + \mathcal{D}_{\beta} T_{\alpha\mu\gamma},$
- (9)  $I_{\mu\alpha\dot{\beta}\dot{\gamma}} : R_{\mu\alpha\dot{\beta}\dot{\gamma}} = -\mathcal{D}_{\alpha} T_{\dot{\beta}\mu\dot{\gamma}} - \bar{\mathcal{D}}_{\dot{\beta}} T_{\alpha\mu\dot{\gamma}} - 2i\sigma^{\nu}{}_{\alpha\dot{\beta}} T_{\nu\mu\dot{\gamma}},$
- (10)  $I_{\mu\nu\dot{\alpha}\dot{\beta}} : R_{\mu\nu\dot{\alpha}\dot{\beta}} = -\mathcal{D}_{\mu} T_{\nu\dot{\alpha}\dot{\beta}} - \mathcal{D}_{\nu} T_{\dot{\alpha}\mu\dot{\beta}} - \bar{\mathcal{D}}_{\dot{\alpha}} T_{\mu\nu\dot{\beta}} + T_{\nu\dot{\alpha}}{}^{\underline{\gamma}} T_{\underline{\gamma}\mu\dot{\beta}} + T_{\dot{\alpha}\mu}{}^{\underline{\gamma}} T_{\underline{\gamma}\nu\dot{\beta}},$
- (11)  $I_{\mu\nu\rho\sigma} : R_{\mu\nu\rho\sigma} + R_{\nu\rho\mu\sigma} + R_{\rho\mu\nu\sigma} = 0,$
- (12)  $I_{\mu\nu\dot{\alpha}\dot{\beta}} : -\mathcal{D}_{\mu} T_{\nu\dot{\alpha}\dot{\beta}} - \mathcal{D}_{\nu} T_{\dot{\alpha}\mu\dot{\beta}} - \bar{\mathcal{D}}_{\dot{\alpha}} T_{\mu\nu\dot{\beta}} + T_{\nu\dot{\alpha}}{}^{\underline{\gamma}} T_{\underline{\gamma}\mu\dot{\beta}} + T_{\dot{\alpha}\mu}{}^{\underline{\gamma}} T_{\underline{\gamma}\nu\dot{\beta}} = 0,$
- (13)  $I_{\mu\nu\rho\alpha} : -\mathcal{D}_{\mu} T_{\nu\rho\alpha} - \mathcal{D}_{\nu} T_{\rho\mu\alpha} - \mathcal{D}_{\rho} T_{\mu\nu\alpha} + T_{\mu\nu}{}^{\underline{\gamma}} T_{\underline{\gamma}\rho\alpha} + T_{\nu\rho}{}^{\underline{\gamma}} T_{\underline{\gamma}\mu\alpha} + T_{\rho\mu}{}^{\underline{\gamma}} T_{\underline{\gamma}\nu\alpha} = 0,$

# Resolution of the 13 Bianchi identities :

- 1) Vector - Bispinor Isomorphism :  $v_\mu \rightarrow v_{\alpha\dot{\alpha}}$
  - 2) Decomposition of the curvature and torsion tensor in  $SL(2, \mathbb{C})$  irreducible representations
- The equations splits in different symmetry classes !
- ★ Result :  $R_{MNPQ}$  and  $T_{MNP} \propto 3$  superfields only :  $\mathcal{R}$ ,  $G_{\alpha\dot{\alpha}}$ ,  $W_{(\alpha\beta\gamma)}$ .

# Explicit Supergravity Algebra

$$\{\mathcal{D}_\alpha, \mathcal{D}_\beta\} = -\frac{1}{2}R_{\alpha\beta\gamma\delta}J^{\delta\gamma},$$

$$\{\bar{\mathcal{D}}_{\dot{\alpha}}, \bar{\mathcal{D}}_{\dot{\beta}}\} = -\frac{1}{2}R_{\dot{\alpha}\dot{\beta}\dot{\gamma}\dot{\delta}}J^{\dot{\delta}\dot{\gamma}},$$

$$\{\mathcal{D}_\alpha, \bar{\mathcal{D}}_{\dot{\alpha}}\} = T_{\alpha\dot{\alpha}}{}^\mu \mathcal{D}_\mu - \frac{1}{2}R_{\alpha\dot{\alpha}\beta\gamma}J^{\gamma\beta} - \frac{1}{2}R_{\alpha\dot{\alpha}\dot{\beta}\dot{\gamma}}J^{\dot{\gamma}\dot{\beta}},$$

$$[\mathcal{D}_\alpha, \mathcal{D}_\mu] = T_{\alpha\mu}{}^\beta \mathcal{D}_\beta - T_{\alpha\mu}{}^{\dot{\beta}} \bar{\mathcal{D}}_{\dot{\beta}} - \frac{1}{2}R_{\alpha\mu\beta\gamma}J^{\gamma\beta} - \frac{1}{2}R_{\alpha\mu\dot{\beta}\dot{\gamma}}J^{\dot{\gamma}\dot{\beta}},$$

$$[\bar{\mathcal{D}}_{\dot{\alpha}}, \mathcal{D}_\mu] = T_{\dot{\alpha}\mu}{}^\beta \mathcal{D}_\beta + T_{\dot{\alpha}\mu}{}^{\dot{\beta}} \bar{\mathcal{D}}_{\dot{\beta}} - \frac{1}{2}R_{\dot{\alpha}\mu\beta\gamma}J^{\gamma\beta} - \frac{1}{2}R_{\dot{\alpha}\mu\dot{\beta}\dot{\gamma}}J^{\dot{\gamma}\dot{\beta}},$$

$$[\mathcal{D}_\mu, \mathcal{D}_\nu] = T_{\mu\nu}{}^\beta \mathcal{D}_\beta - T_{\mu\nu}{}^{\dot{\beta}} \bar{\mathcal{D}}_{\dot{\beta}} - \frac{1}{2}R_{\mu\nu\beta\gamma}J^{\gamma\beta} - \frac{1}{2}R_{\mu\nu\dot{\beta}\dot{\gamma}}J^{\dot{\gamma}\dot{\beta}}.$$

# The Physical degrees of freedom of Supergravity

- $E_{\tilde{M}}^M$  has  $8 \times 8 \times 16 = 1024$  components

## General Change of Coordinates

$$z^{\tilde{M}} \rightarrow z^{\tilde{M}} + \xi^{\tilde{M}}(z),$$

- $z^{\tilde{M}}$  : Coordinates of the curved superspace
- Transformations parameters  $\xi^{\tilde{M}}(z)$  : superfields depending on the coordinates  $z^{\tilde{M}}$ .
- ★ Physical degrees of freedom :
  - Graviton :  $e_{\tilde{\mu}}^{\mu}$
  - Gravitino :  $(\psi_{\mu}^{\tilde{\alpha}}, \bar{\psi}_{\tilde{\mu}\tilde{\alpha}})$
  - $\mathcal{R}$ ,  $G_{\alpha\tilde{\alpha}}$  give us the two auxiliary field  $b_{\mu}$  and  $M$  to close the supermultiplet off shell.

## Annexe A : How do we get the Supergravity Constrains ?

- ★ In the referential frame where  $\Omega_{\tilde{M}\tilde{M}}^N = 0 \Rightarrow \emptyset$  Curvature because

$$R_{MNS}^Q = (-)^{|\tilde{M}|(|\tilde{N}|+|N|)} E_M^{\tilde{M}} E_N^{\tilde{N}} [\partial_{\tilde{M}} \Omega_{\tilde{N}S}^Q - \Omega_{\tilde{M}S}^R \Omega_{\tilde{N}R}^Q - (-)^{|\tilde{M}||\tilde{N}|} (\partial_{\tilde{N}} \Omega_{\tilde{M}S}^Q - \Omega_{\tilde{N}S}^R \Omega_{\tilde{M}R}^Q)].$$

- ★ The covariant derivatives reduce to  $\mathcal{D}_{\tilde{M}} = \partial_{\tilde{M}}$ .  
Introducing the Lorentz indices, we obtain the Covariant derivatives of Susy :

$$E_{\mu}^{\tilde{M}} \partial_{\tilde{M}} = D_{\mu} = \partial_{\mu}$$

$$E_{\alpha}^{\tilde{M}} \partial_{\tilde{M}} = D_{\alpha} = \partial_{\alpha} + i\sigma_{\alpha\dot{\alpha}}^{\mu} \bar{\theta}^{\dot{\alpha}} \partial_{\mu}$$

$$E_{\dot{\alpha}}^{\tilde{M}} \partial_{\tilde{M}} = \bar{D}_{\dot{\alpha}} = \bar{\partial}_{\dot{\alpha}} - i\theta^{\alpha} \sigma_{\alpha\dot{\alpha}}^{\mu} \partial_{\mu},$$

# Annexe A : How do we get the Supergravity Constraints ?

★ The superalgebra reduces to

$$\{D_\alpha, \bar{D}_{\dot{\alpha}}\} = T_{\alpha\dot{\alpha}}{}^\beta D_\beta + T_{\alpha\dot{\alpha}\dot{\beta}} \bar{D}^{\dot{\beta}} + T_{\alpha\dot{\alpha}}{}^\mu \partial_\mu$$

$$\{D_\alpha, D_\beta\} = T_{\alpha\beta}{}^\gamma D_\gamma + T_{\alpha\beta\dot{\beta}} \bar{D}^{\dot{\beta}} + T_{\alpha\beta}{}^\mu \partial_\mu$$

$$\{\bar{D}_{\dot{\alpha}}, \bar{D}_{\dot{\beta}}\} = T_{\dot{\alpha}\dot{\beta}}{}^\beta D_\beta + T_{\dot{\alpha}\dot{\beta}\dot{\gamma}} \bar{D}^{\dot{\gamma}} + T_{\dot{\alpha}\dot{\beta}}{}^\mu \partial_\mu$$

★ Algebra of the covariant derivatives in Supersymmetry :

$$\{D_\alpha, \bar{D}_{\dot{\alpha}}\} = -2i\sigma^\mu_{\alpha\dot{\alpha}} \partial_\mu$$

$$\{D_\alpha, D_\beta\} = \{\bar{D}_{\dot{\alpha}}, \bar{D}_{\dot{\beta}}\} = 0.$$

⇒ Constrain the Torsion tensor components !

$$T_{\alpha\beta}{}^\gamma = 0, \quad T_{\alpha\beta}{}^\mu = T_{\dot{\alpha}\dot{\beta}}{}^\mu = 0,$$

$$T_{\alpha\dot{\alpha}}{}^\mu = T_{\dot{\alpha}\alpha}{}^\mu = -2i\sigma^\mu_{\alpha\dot{\alpha}}.$$