

Local maxima of white noise spectrograms and interlacing of zeros in higher Landau Levels

Luis Daniel Abreu

NuHAG, Faculty of Mathematics
University of Vienna

dpp-fermions-2022/Lyon
9/6/2022

FWF

Der Wissenschaftsfonds.



**universität
wien**

Landau Levels Operators

Landau Levels Operators

1-Landau Levels raising operator:

$$\nabla_z^+ = \partial_z - \bar{z}$$

Landau Levels Operators

1-Landau Levels raising operator:

$$\nabla_z^+ = \partial_z - \bar{z}$$

$$\nabla_z^+ z^n \bar{z}^m = n z^{n-1} \bar{z}^m - z^n \bar{z}^{m+1}$$

Landau Levels Operators

1-Landau Levels raising operator:

$$\nabla_z^+ = \partial_z - \bar{z}$$

$$\nabla_z^+ z^n \bar{z}^m = n z^{n-1} \bar{z}^m - z^n \bar{z}^{m+1}$$

2-Landau Levels lowering operator:

$$\nabla_z^- = \partial_{\bar{z}}$$

Landau Levels Operators

1-Landau Levels raising operator:

$$\nabla_z^+ = \partial_z - \bar{z}$$

$$\nabla_z^+ z^n \bar{z}^m = n z^{n-1} \bar{z}^m - z^n \bar{z}^{m+1}$$

2-Landau Levels lowering operator:

$$\nabla_z^- = \partial_{\bar{z}}$$

$$\nabla_z^- z^n \bar{z}^m = m z^n \bar{z}^{m-1}$$

Landau Levels Operators

1-Landau Levels raising operator:

$$\nabla_z^+ = \partial_z - \bar{z}$$

$$\nabla_z^+ z^n \bar{z}^m = n z^{n-1} \bar{z}^m - z^n \bar{z}^{m+1}$$

2-Landau Levels lowering operator:

$$\nabla_z^- = \partial_{\bar{z}}$$

$$\nabla_z^- z^n \bar{z}^m = m z^n \bar{z}^{m-1}$$

3-Landau Levels: $\sigma(L_z) = \{n : n \in \mathbb{N}_0\}$

$$L_z := -\partial_z \partial_{\bar{z}} + \bar{z} \partial_{\bar{z}} = -\nabla_z^+ \nabla_z^-$$

- Eigenspaces with reproducing kernel:

$$K_n(z, w) = L_n(|z - w|^2) e^{z\bar{w}}$$

Critical points of white noise spectrograms

- Gaussian Entire Function:

$$f_0(z) = \sum_{k=0}^{\infty} \zeta_k \frac{z^k}{\sqrt{k!}}$$

- Spectrogram of white noise (Bardenet, Flamant, Chainais)

$$\text{Spec}_{h_0} \mathcal{N}(z) = \left| V_{h_0} \mathcal{N}\left(\frac{\bar{z}}{\pi}\right) \right|^2 = \left| e^{-|z|^2/2} f_0(z) \right|^2$$

Critical points of white noise spectrograms

- Gaussian Entire Function:

$$f_0(z) = \sum_{k=0}^{\infty} \zeta_k \frac{z^k}{\sqrt{k!}}$$

- Spectrogram of white noise (Bardenet, Flamant, Chainais)

$$\text{Spec}_{h_0} \mathcal{N}(z) = \left| V_{h_0} \mathcal{N}\left(\frac{\bar{z}}{\pi}\right) \right|^2 = \left| e^{-|z|^2/2} f_0(z) \right|^2$$

- Critical points log-spectrogram

$$0 = \nabla \log \left| f_0(z) e^{-\frac{|z|^2}{2}} \right| = \frac{\partial_z f_0(z)}{f_0(z)} + \frac{\partial_{\bar{z}} f_0(z)}{f_0(z)} - \bar{z} = \frac{\partial_z f_0(z)}{f_0(z)} - \bar{z}$$

Critical points of white noise spectrograms

- Gaussian Entire Function:

$$f_0(z) = \sum_{k=0}^{\infty} \zeta_k \frac{z^k}{\sqrt{k!}}$$

- Spectrogram of white noise (Bardenet, Flamant, Chainais)

$$\text{Spec}_{h_0} \mathcal{N}(z) = \left| V_{h_0} \mathcal{N}\left(\frac{\bar{z}}{\pi}\right) \right|^2 = \left| e^{-|z|^2/2} f_0(z) \right|^2$$

- Critical points log-spectrogram

$$0 = \nabla \log \left| f_0(z) e^{-\frac{|z|^2}{2}} \right| = \frac{\partial_z f_0(z)}{f_0(z)} + \frac{\partial_{\bar{z}} f_0(z)}{f_0(z)} - \bar{z} = \frac{\partial_z f_0(z)}{f_0(z)} - \bar{z}$$

- Critical points equation (LL raising operator)

$$\partial_z f_0(z) - \bar{z} f_0(z) = 0 \iff \nabla_z^+ f_0(z) = 0$$

- 1-point intensities: expected number of points in $\Omega \subset \mathbb{C}$:

$$\mathbb{E}(\mathcal{X}(\Omega)) = \int_{\Omega} \rho^{(1)}(z) \frac{dz}{\pi}.$$

- 1-point intensities: expected number of points in $\Omega \subset \mathbb{C}$:

$$\mathbb{E}(\mathcal{X}(\Omega)) = \int_{\Omega} \rho^{(1)}(z) \frac{dz}{\pi}.$$

- $\rho_{point_process, n}^{(1)}(z)$

- 1-point intensities: expected number of points in $\Omega \subset \mathbb{C}$:

$$\mathbb{E}(\mathcal{X}(\Omega)) = \int_{\Omega} \rho^{(1)}(z) \frac{dz}{\pi}.$$

- $\rho_{point_process, n}^{(1)}(z)$
- n is the Landau Level
- $n = 0 \Rightarrow$ Entire Functions

Notation for 1-point intensities

- 1-point intensities: expected number of points in $\Omega \subset \mathbb{C}$:

$$\mathbb{E}(\mathcal{X}(\Omega)) = \int_{\Omega} \rho^{(1)}(z) \frac{dz}{\pi}.$$

- $\rho_{point_process,n}^{(1)}(z)$
- n is the Landau Level
- $n = 0 \Rightarrow$ Entire Functions
- Example: zeros of Gaussian Entire Function

$$\rho_{zeros,0}^{(1)}(z) = 1$$

Theorem (A., 2020)

$$\rho_{critical,0}^{(1)}(z) = 5/3, \rho_{loc\ max,0}^{(1)}(z) = 1/3 \text{ and } \rho_{saddle,0}^{(1)}(z) = 4/3.$$

Theorem (A., 2020)

$\rho_{critical,0}^{(1)}(z) = 5/3$, $\rho_{loc\ max,0}^{(1)}(z) = 1/3$ and $\rho_{saddle,0}^{(1)}(z) = 4/3$.

Proof idea At a critical point z_c ,

$$|f_0(z_c)|^2 \det D^{Hess}(\log |f_0(z)e^{-\frac{|z|^2}{2}}|)(z_c) = Jac(\nabla_z^+ f_0)(z_c)$$

- $Jac(\nabla_z^+ f_0)(z_c) > 0 \Rightarrow z_c$ maximum (minima are zeros)
- $Jac(\nabla_z^+ f_0)(z_c) < 0 \Rightarrow z_c$ saddle point
- Kac-Rice formulas:

$$\rho_{loc\ max,0}^{(1)}(z) = \mathbb{E}(|Jac(\nabla_z^+ f_0) : Jac(\nabla_z^+ f_0) > 0| \delta_0(\nabla_z^+ f_0)) = \frac{1}{3}$$

$$\rho_{saddle,0}^{(1)}(z) = \mathbb{E}(|Jac(\nabla_z^+ f_0) : Jac(\nabla_z^+ f_0) < 0| \delta_0(\nabla_z^+ f_0)) = \frac{4}{3}$$

Flandrin's honeycomb

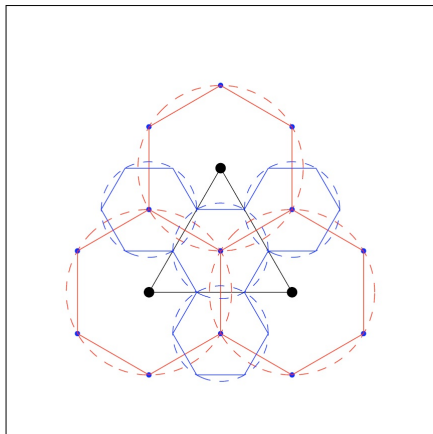


Figure: Flandrin's honeycomb: the area of the hexagons centered at local maxima (big dots) is three times the area of the hexagons centered at zeros (small dots)

Gaussian function in first Landau Level

- Gaussian Entire Function:

$$f_0(z) = \sum_{k=0}^{\infty} \frac{\zeta_k}{\sqrt{k!}} z^k$$

- Gaussian function in first Landau Level:

$$f_1(z, \bar{z}) = \nabla_z^+ f_0(z) = \sum_{k=0}^{\infty} \zeta_k \frac{kz^{k-1} - z^k \bar{z}}{\sqrt{k!}}$$

- Zeros:

$$\rho_{zeros,1}^{(1)}(z) = \rho_{critical,0}^{(1)}(z) = \frac{5}{3}$$

- Relation with spectrogram of white noise:

$$Spec_{h_1} \mathcal{N}(z) = \left| V_{h_1} \mathcal{N}\left(\frac{\bar{z}}{\pi}\right) \right|^2 = \left| e^{-|z|^2/2} f_1(z, \bar{z}) \right|^2$$

Gaussian function in first Landau Level

- Gaussian Entire Function:

$$f_0(z) = \sum_{k=0}^{\infty} \frac{\zeta_k}{\sqrt{k!}} z^k$$

- Gaussian function in first Landau Level:

$$f_1(z, \bar{z}) = \nabla_z^+ f_0(z) = \sum_{k=0}^{\infty} \zeta_k \frac{kz^{k-1} - z^k \bar{z}}{\sqrt{k!}}$$

- Zeros:

$$\rho_{zeros,1}^{(1)}(z) = \rho_{critical,0}^{(1)}(z) = \frac{5}{3}$$

- Relation with spectrogram of white noise:

$$Spec_{h_1} \mathcal{N}(z) = \left| V_{h_1} \mathcal{N}\left(\frac{\bar{z}}{\pi}\right) \right|^2 = \left| e^{-|z|^2/2} f_1(z, \bar{z}) \right|^2$$

- Interlacing property of local maxima

Maxima of $Spec_{h_0} \mathcal{N}(z)$ separate maxima of $Spec_{h_1} \mathcal{N}(z)$

Gaussian functions in higher Landau Levels

- Gaussian Entire Function:

$$f_0(z) = \sum_{k=0}^{\infty} \zeta_k \frac{z^k}{\sqrt{k!}}$$

- Gaussian (Polyanalytic) Function in higher Landau levels:

$$f_n(z, \bar{z}) := \frac{1}{\sqrt{n!}} (\nabla_z^+)^n f_0(z) = \sum_{k=0}^{\infty} \zeta_k \frac{(\nabla_z^+)^n z^k}{\sqrt{n!k!}}$$

- Orthogonal Gaussian functions:

$$\int_{\mathbb{C}} f_n(z, \bar{z}) \overline{f_{n'}(z, \bar{z})} e^{-|z|^2} dz = \delta_{nn'}.$$

- Correlation kernels

$$K_n(z, w) = L_n(|z - w|^2) e^{z\bar{w}}$$

Complex Hermite Polynomials

- Gaussian (Polyanalytic) Function in higher Landau levels:

$$f_n(z, \bar{z}) := \frac{1}{\sqrt{n!}} (\nabla_z^+)^n f_0(z) = \sum_{k=0}^{\infty} \zeta_k H_{k,n}(z, \bar{z})$$

Complex Hermite Polynomials

- Gaussian (Polyanalytic) Function in higher Landau levels:

$$f_n(z, \bar{z}) := \frac{1}{\sqrt{n!}} (\nabla_z^+)^n f_0(z) = \sum_{k=0}^{\infty} \zeta_k H_{k,n}(z, \bar{z})$$

- $H_{k,n}(z, \bar{z})$ are the complex Hermite polynomials:

$$H_{k,n}(z, \bar{z}) = \begin{cases} \sqrt{\frac{n!}{k!}} z^{k-n} L_n^{k-n}(|z|^2), & k > n \geq 0, \\ (-1)^{n-k} \sqrt{\frac{k!}{n!}} \bar{z}^{n-k} L_k^{n-k}(|z|^2), & 0 \leq k \leq n, \end{cases}$$

- Doubly-indexed orthogonality relation

$$\int_{\mathbb{C}} H_{k,n}(z, \bar{z}) \overline{H_{k',n'}(z, \bar{z})} e^{-|z|^2} dz = \delta_{kk'} \delta_{nn'}$$

- Orthogonality in concentric discs

$$\int_{D(0,R)} H_{k,n}(z, \bar{z}) \overline{H_{k',n'}(z, \bar{z})} e^{-|z|^2} dz = C_{n,k,k'}(R) \delta_{kk'}$$

- Recall:

$$\nabla_z^+ f_0(z) = f_1(z, \bar{z}) \implies \rho_{zeros,1}^{(1)}(z) = \rho_{critical,0}^{(1)}(z) = \frac{5}{3}$$

Zeros in higher Landau Levels

- Recall:

$$\nabla_z^+ f_0(z) = f_1(z, \bar{z}) \implies \rho_{zeros,1}^{(1)}(z) = \rho_{critical,0}^{(1)}(z) = \frac{5}{3}$$

- Relation with STFT of white noise

$$V_{h_n} \mathcal{W}\left(\frac{\bar{z}}{\pi}\right) = e^{ix\xi} e^{-|z|^2/2} f_n(z, \bar{z})$$

Zeros in higher Landau Levels

- Recall:

$$\nabla_z^+ f_0(z) = f_1(z, \bar{z}) \implies \rho_{zeros,1}^{(1)}(z) = \rho_{critical,0}^{(1)}(z) = \frac{5}{3}$$

- Relation with STFT of white noise

$$V_{h_n} \mathcal{W}\left(\frac{\bar{z}}{\pi}\right) = e^{ix\xi} e^{-|z|^2/2} f_n(z, \bar{z})$$

- Gaussian Weyl-Heisenberg function

Theorem (Haimi, Koliander, Romero)

$$\rho_{zeros,n}^{(1)}(z) = n + 1/2 + \frac{1}{4n+2}$$

Critical points in higher Landau Levels

Theorem (A., Shirai, 2022)

$$\rho_{critical,n}^{(1)}(z) = \frac{1}{2} \left(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3} \right)$$

Critical points in higher Landau Levels

Theorem (A., Shirai, 2022)

$$\rho_{critical,n}^{(1)}(z) = \frac{1}{2}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3})$$

$$\rho_{saddle,n}^{(1)}(z) = \frac{1}{4}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) + \frac{1}{2}$$

$$\rho_{extremal,n}^{(1)}(z) = \frac{1}{4}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) - \frac{1}{2}$$

Critical points in higher Landau Levels

Theorem (A., Shirai, 2022)

$$\begin{aligned}\rho_{critical,n}^{(1)}(z) &= \frac{1}{2}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) \\ \rho_{saddle,n}^{(1)}(z) &= \frac{1}{4}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) + \frac{1}{2} \\ \rho_{extremal,n}^{(1)}(z) &= \frac{1}{4}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) - \frac{1}{2}\end{aligned}$$

Proof Sketch Critical points equation

$$0 = \nabla \log \left| f_n(z, \bar{z}) e^{-\frac{|z|^2}{2}} \right| = \frac{\partial_z f_n(z, \bar{z})}{f_n(z, \bar{z})} + \frac{\partial_{\bar{z}} f_n(z, \bar{z})}{f_n(z, \bar{z})} - \bar{z}$$

Critical points in higher Landau Levels

Theorem (A., Shirai, 2022)

$$\begin{aligned}\rho_{critical,n}^{(1)}(z) &= \frac{1}{2}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) \\ \rho_{saddle,n}^{(1)}(z) &= \frac{1}{4}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) + \frac{1}{2} \\ \rho_{extremal,n}^{(1)}(z) &= \frac{1}{4}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) - \frac{1}{2}\end{aligned}$$

Proof Sketch Critical points equation

$$0 = \nabla \log \left| f_n(z, \bar{z}) e^{-\frac{|z|^2}{2}} \right| = \frac{\partial_z f_n(z, \bar{z})}{f_n(z, \bar{z})} + \frac{\partial_{\bar{z}} f_n(z, \bar{z})}{f_n(z, \bar{z})} - \bar{z}$$

$$\nabla_z^+ f_n(z, \bar{z}) + \nabla_{\bar{z}}^- f_n(z, \bar{z}) = 0$$

Critical points in higher Landau Levels

Theorem (A., Shirai, 2022)

$$\begin{aligned}\rho_{critical,n}^{(1)}(z) &= \frac{1}{2}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) \\ \rho_{saddle,n}^{(1)}(z) &= \frac{1}{4}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) + \frac{1}{2} \\ \rho_{extremal,n}^{(1)}(z) &= \frac{1}{4}(2n^2 + 4n + 3 + \frac{1}{2n^2 + 4n + 3}) - \frac{1}{2}\end{aligned}$$

Proof Sketch Critical points equation

$$0 = \nabla \log \left| f_n(z, \bar{z}) e^{-\frac{|z|^2}{2}} \right| = \frac{\partial_z f_n(z, \bar{z})}{f_n(z, \bar{z})} + \frac{\partial_{\bar{z}} f_n(z, \bar{z})}{f_n(z, \bar{z})} - \bar{z}$$

$$\nabla_z^+ f_n(z, \bar{z}) + \nabla_z^- f_n(z, \bar{z}) = 0$$

- STFT of white noise (Gaussian Weyl-Heisenberg function)

$$\nabla_z^+ f_n(z, \bar{z}) + \nabla_z^- f_n(z, \bar{z}) = e^{ix\xi} e^{-|z|^2/2} V_{\sqrt{nh_{n-1} - \sqrt{n+1}h_{n+1}}} \mathcal{N}\left(\frac{\bar{z}}{\sqrt{\pi}}\right)$$

- Average of Hermite spectrograms of single realization of \mathcal{S}

$$\frac{1}{N} \sum_{n=0}^{N-1} \text{Spec}_{h_n} \mathcal{S}(z) = \frac{1}{N} \sum_{n=0}^{N-1} \left| V_{h_n} \mathcal{S}\left(\frac{\bar{z}}{\pi}\right) \right|^2$$

- White noise case (conjecture):

$$\lim_{n \rightarrow \infty} \frac{1}{N} \sum_{n=0}^{N-1} \left| V_{h_n} \mathcal{N}\left(\frac{\bar{z}}{\pi}\right) \right|^2 = 1$$

- Average of Hermite spectrograms of single realization of \mathcal{S}

$$\frac{1}{N} \sum_{n=0}^{N-1} \text{Spec}_{h_n} \mathcal{S}(z) = \frac{1}{N} \sum_{n=0}^{N-1} \left| V_{h_n} \mathcal{S}\left(\frac{\bar{z}}{\pi}\right) \right|^2$$

- White noise case (conjecture):

$$\lim_{n \rightarrow \infty} \frac{1}{N} \sum_{n=0}^{N-1} \left| V_{h_n} \mathcal{N}\left(\frac{\bar{z}}{\pi}\right) \right|^2 = 1$$

- Reminiscent of Pythagoras theorem

$$\sin^2 x + \cos^2 x = 1$$

Pythagoras theorem

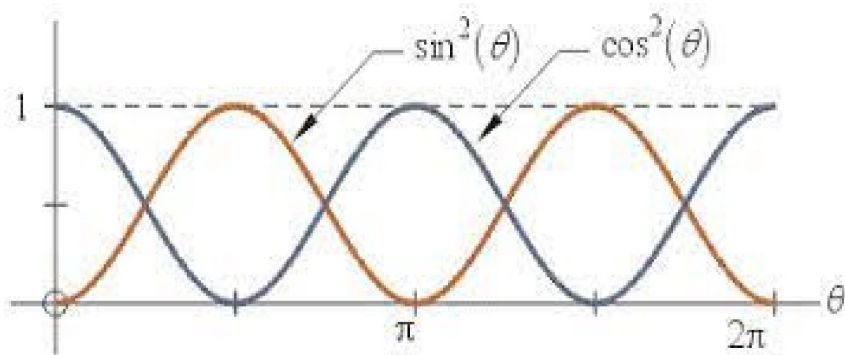


Figure: Energy allocation of $\sin^2 x$ in the line complements energy allocation of $\cos^2 x$; Maxima and zeros of $\sin x$ and $\cos x$ are interlaced.

Pythagoras theorem

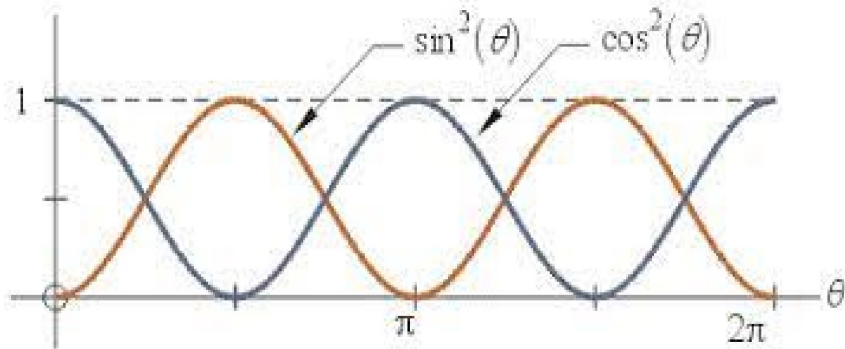


Figure: Energy allocation of $\sin^2 x$ in the line complements energy allocation of $\cos^2 x$; Maxima and zeros of $\sin x$ and $\cos x$ are interlaced.

- Do maxima and zeros of $\text{Spec}_{h_n} \mathcal{N}(z)$ and $\text{Spec}_{h_{n+k}} \mathcal{N}(z)$ tend to interlace?

Interlacing of Hermite spectrogram energy

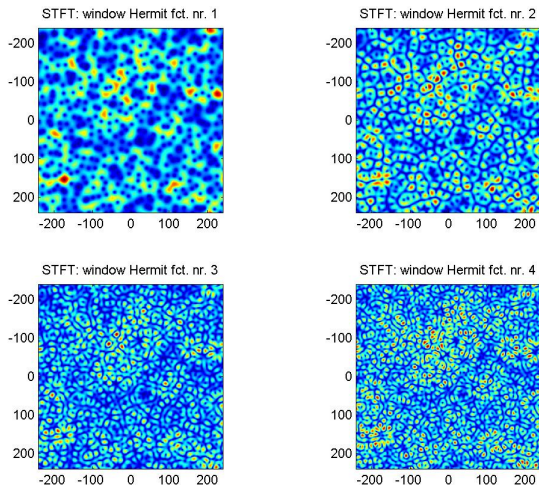


Figure: Orange: high energy (loud); Dark blue: low energy (silent)

Hermite spectrogram maxima (Flandrin)

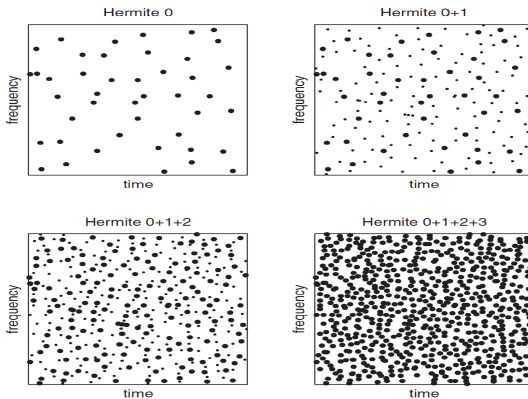


Figure: Maxima become more numerous and occur in different regions

- Maxima seem to interlace (seems hopeless to prove)

Hermite spectrogram maxima (Flandrin)

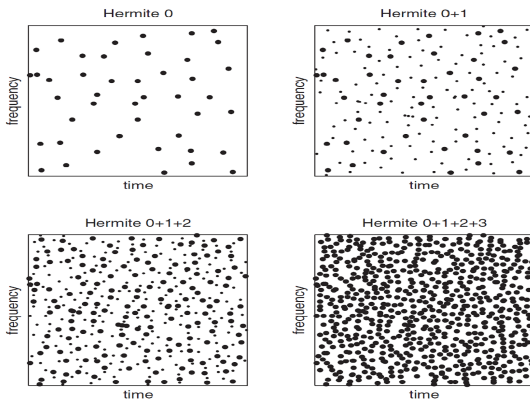


Figure: Maxima become more numerous and occur in different regions

- Maxima seem to interlace (seems hopeless to prove)
- What about zeros? Do zeros interlace? Repulsion?

Correlations between zeros of different Landau levels

- How strong is the orthogonality?

$$\int_{\mathbb{C}} V_{h_n} \mathcal{W}(z) \overline{V_{h_{n+k}} \mathcal{W}(z)} dz = \delta_{n,n+k}$$

- Energy of $Spec_{h_n} \mathcal{N}(z)$ and $Spec_{h_{n+k}} \mathcal{N}(z)$ placed in complementary regions? (repulsion?)

Correlations between zeros of different Landau levels

- How strong is the orthogonality?

$$\int_{\mathbb{C}} V_{h_n} \mathcal{W}(z) \overline{V_{h_{n+k}} \mathcal{W}(z)} dz = \delta_{n,n+k}$$

- Energy of $\text{Spec}_{h_n} \mathcal{N}(z)$ and $\text{Spec}_{h_{n+k}} \mathcal{N}(z)$ placed in complementary regions? (repulsion?)
- $\rho_{n,n+k}^{(2)}(z, w) =$ correlation zeros $f_n(z, \bar{z})$, zeros of $f_{n+k}(z, \bar{z})$.
- *Pair correlation function* $g_{n,n+1}(z, w)$:

$$g_{n,n+1}(z, w) = \frac{\rho_{n,n+k}^{(2)}(z, w)}{\rho_{\text{zeros},n}^{(1)}(z) \rho_{\text{zeros},n+k}^{(1)}(z)} = \begin{cases} < 1 \rightarrow & \text{Repulsion} \\ = 1 \rightarrow & \text{Neutrality} \\ > 1 \rightarrow & \text{Attraction} \end{cases}$$

Theorem (Short range correlations A., Shirai 2022)

As $w \rightarrow z$,

$$g_{n,n+1}(z, w) \sim \frac{n(n+2)}{(n+1/2 + \frac{1}{4n+2})(n+3/2 + \frac{1}{4n+4})} < 1$$

$$g_{n,n+k}(z, w) \sim 1, \quad k > 1$$

Theorem (Short range correlations A., Shirai 2022)

As $w \rightarrow z$,

$$g_{n,n+1}(z, w) \sim \frac{n(n+2)}{(n+1/2 + \frac{1}{4n+2})(n+3/2 + \frac{1}{4n+4})} < 1$$

$$g_{n,n+k}(z, w) \sim 1, \quad k > 1$$

- Repulsion between zeros of $V_{h_n} \mathcal{N}$ and $V_{h_{n+1}} \mathcal{N}$.
- Zeros of $V_{h_n} \mathcal{N}$ and $V_{h_{n+1}} \mathcal{N}$ tend to interlace.
- For $k > 1$, zeros of $V_{h_n} \mathcal{N}$ ‘don’t feel’ zeros of $V_{h_{n+k}} \mathcal{N}$.

Theorem (Short range correlations A., Shirai 2022)

As $w \rightarrow z$,

$$g_{n,n+1}(z, w) \sim \frac{n(n+2)}{(n+1/2 + \frac{1}{4n+2})(n+3/2 + \frac{1}{4n+4})} < 1$$

$$g_{n,n+k}(z, w) \sim 1, \quad k > 1$$

- Repulsion between zeros of $V_{h_n} \mathcal{N}$ and $V_{h_{n+1}} \mathcal{N}$.
- Zeros of $V_{h_n} \mathcal{N}$ and $V_{h_{n+1}} \mathcal{N}$ tend to interlace.
- For $k > 1$, zeros of $V_{h_n} \mathcal{N}$ ‘don’t feel’ zeros of $V_{h_{n+k}} \mathcal{N}$.
- Landau Levels too far apart?

Theorem (Short range correlations A., Shirai 2022)

As $w \rightarrow z$,

$$g_{n,n+1}(z, w) \sim \frac{n(n+2)}{(n+1/2 + \frac{1}{4n+2})(n+3/2 + \frac{1}{4n+4})} < 1$$

$$g_{n,n+k}(z, w) \sim 1, \quad k > 1$$

- Repulsion between zeros of $V_{h_n} \mathcal{N}$ and $V_{h_{n+1}} \mathcal{N}$.
- Zeros of $V_{h_n} \mathcal{N}$ and $V_{h_{n+1}} \mathcal{N}$ tend to interlace.
- For $k > 1$, zeros of $V_{h_n} \mathcal{N}$ ‘don’t feel’ zeros of $V_{h_{n+k}} \mathcal{N}$.
- Landau Levels too far apart?
- Inheritance of interlacing properties of h_n and h_{n+1} ?

Ideas in the proof

- For the proof, we use the formulas

$$\nabla^+ f_n = \sqrt{n+1} f_{n+1}, \quad \nabla^- f_n = -\sqrt{n} f_{n-1}.$$

- And rewrite the 6-dimensional Gaussian vector

$$\begin{aligned} & (f_n(z), (\nabla^+)^k f_n(w), \nabla^+ f_n(z), \nabla^- f_n(z), (\nabla^+)^{k+1} f_n(w), \nabla^- (\nabla^+)^k f_n(w)) \\ &= (f_n(z), \sqrt{(n+1)_k} f_{n+k}(w), \sqrt{n+1} f_{n+1}(z), -\sqrt{n} f_{n-1}(z), \\ & \quad \sqrt{(n+1)_{k+1}} f_{n+k+1}(w), -\sqrt{n+k} \sqrt{(n+1)_k} f_{n+k-1}(w)) \end{aligned}$$

Ideas in the proof

- For the proof, we use the formulas

$$\nabla^+ f_n = \sqrt{n+1} f_{n+1}, \quad \nabla^- f_n = -\sqrt{n} f_{n-1}.$$

- And rewrite the 6-dimensional Gaussian vector

$$\begin{aligned} & (f_n(z), (\nabla^+)^k f_n(w), \nabla^+ f_n(z), \nabla^- f_n(z), (\nabla^+)^{k+1} f_n(w), \nabla^- (\nabla^+)^k f_n(w)) \\ &= (f_n(z), \sqrt{(n+1)_k} f_{n+k}(w), \sqrt{n+1} f_{n+1}(z), -\sqrt{n} f_{n-1}(z), \\ & \quad \sqrt{(n+1)_{k+1}} f_{n+k+1}(w), -\sqrt{n+k} \sqrt{(n+1)_k} f_{n+k-1}(w)) \end{aligned}$$

- The resulting Kac-Rice formulas are intractable for general (z, w)
- As $w \rightarrow z$ several entries of the correlation matrix become 0





- No surprise in the long range correlations:




Theorem (Long range correlations A., Shirai 2022)

As $r \rightarrow \infty$,

$$g_{n,n+1}(r, 0) \rightarrow 1.$$

- zeros of $V_{h_n} \mathcal{N}$ ‘don’t feel’ zeros of $V_{h_{n+1}} \mathcal{N}$ when they are asymptotically far from each other.

-  P Flandrin, *Explorations in Time-frequency Analysis*, Cambridge University Press, 2018.
-  L. D. Abreu, *Local maxima of white noise spectrograms and Gaussian Entire Functions*, preprint RG (2020).
-  L. D. Abreu, T. Shirai, *Gaussian Random Functions in higher Landau levels and white noise spectrograms*, in preparation.
-  A. Haimi, G. Koliander, J. L. Romero, *Zeros of Gaussian Weyl-Heisenberg functions and hyperuniformity of charge*, arXiv preprint arXiv:2012.12298, (2020), Journal of Statistical Physics 187, 22 (2022).

-  R. Feng, *Correlations between zeros and critical points of random analytic functions*. Trans. Amer. Math. Soc., 371 (8), 5247-5265, (2019).
-  M. R. Douglas, B. Shiffman, S. Zelditch, *Critical points and supersymmetric vacua I*, Comm. Math. Phys. 252 (1-3), 325-358, (2004).
-  M. Sodin, B. Tsirelson, *Random complex zeroes, I. Asymptotic normality*. Israel J. of Math., 144 (1), 125-149, (2004).

A. Haimi, H. Hedenmalm, *The polyanalytic Ginibre ensembles*, J. Stat. Phys., 153(1):10–47, (2013).

T. Shirai, *Ginibre-type point processes and their asymptotic behavior*, J. Math. Soc. Japan 67 , 2, 763-787, (2015)

A. Haimi, H. Hedenmalm, *Asymptotic expansions of polyanalytic Bergman kernels*. J. Funct. Anal., 267, 4667-4731 (2014).

L. D. Abreu, J. M. Pereira, J. L. Romero, and S. Torquato. *The Weyl-Heisenberg ensemble: hyperuniformity and higher Landau levels*. J. Stat. Mech. Theor. Exp., 043103, (2017).

A. Haimi, A. Wennman, *A Central limit theorem for fluctuations in Polyanalytic Ginibre ensembles*. Inter. Math. Res. Notices, (5), 1350-1372, (2019).

L. D. Abreu, K. Gröchenig, J. L. Romero, *Harmonic analysis in phase space and finite Weyl-Heisenberg ensembles*. J. Stat. Phys., vol. 174, 5, 1104–1136, (2019).

M. Fenzl, G. Lambert, *Precise deviations for disk counting statistics of invariant determinantal processes*, Inter. Math. Res. Notices, (10) 7420-7494, (2022).

Thanks!

Thanks!