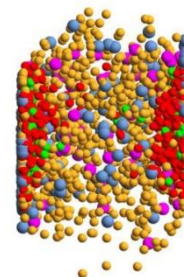


Dilepton production in transport models

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(GSI, Darmstadt & Uni. Frankfurt)



Workshop 'Prospects on various aspects of the dilepton
probe in hadronic physics'
November 25, 2021
IJClab, France



Electromagnetic probes

□ Dileptons and real photons :

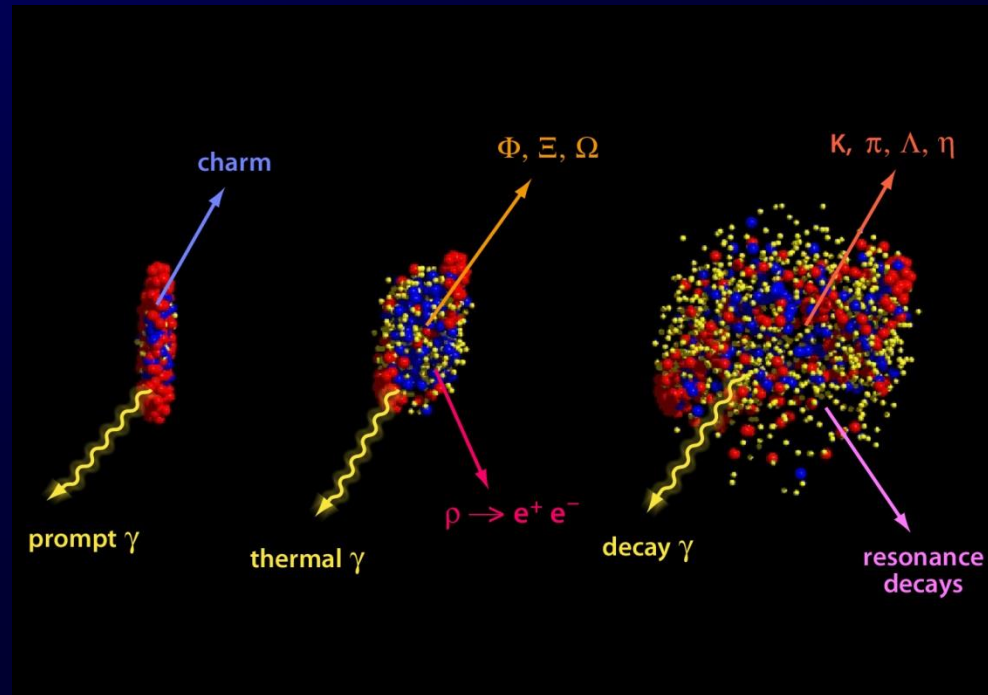
- not effected by final-state partonic and hadronic interactions
- promising signal of QGP – ‚thermal‘ photons and dileptons
- emitted from many production sources at different stages of the HIC

→ Requires **theoretical models** which describe the **dynamics** of heavy-ion collisions during the whole time evolution!



Microscopic transport approaches

Theoretical laboratory - PHSD





Parton-Hadron-String-Dynamics (PHSD)



PHSD is a non-equilibrium transport approach with

- explicit **phase transition** from hadronic to partonic degrees of freedom
- **IQCD EoS** for the partonic phase (‘crossover’ at finite μ_q)
- explicit **parton-parton interactions** - between quarks and gluons
- dynamical **hadronization** and hadronic interactions

□ **QGP phase** is described by the **Dynamical QuasiParticle Model (DQPM)** matched to reproduce lattice QCD

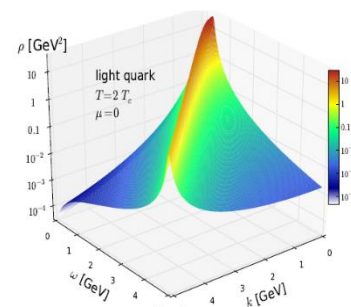
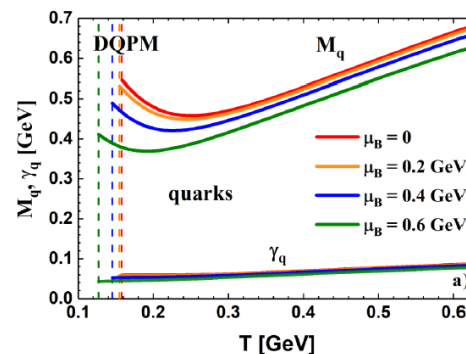
▪ **strongly interacting quasi-particles** (non-perturbative QCD): massive quarks and gluons (g, q, q_{bar}) with sizeable collisional widths in self-generated **mean-field potential**

▪ **Spectral functions:**

$$\rho_i(\omega, T, \mu_B) = \frac{4\omega\Gamma_i(T, \mu_B)}{(\omega^2 - \vec{p}^2 - M_i^2(T, \mu_B))^2 + 4\omega^2\Gamma_i^2(T, \mu_B)}$$

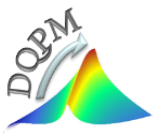
($i = q, \bar{q}, g$)

A. Peshier, W. Cassing, PRL 94 (2005) 172301;
W. Cassing, NPA 791 (2007) 365; NPA 793 (2007)



P. Moreau et al., PRC100 (2019) 014911;
O. Soloveva et al., PRC101 (2020) 045203

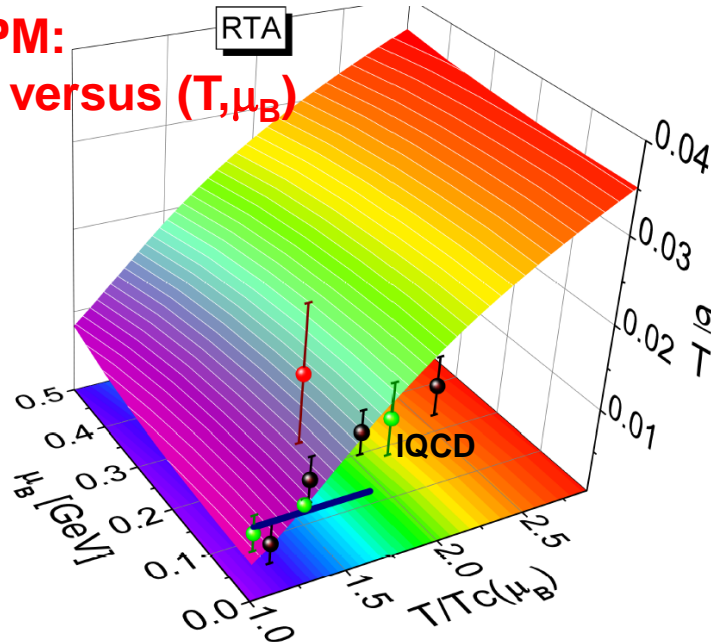
□ **Transport theory:** generalized off-shell transport equations based on the 1st order gradient expansion of Kadanoff-Baym equations (**applicable for strongly interacting systems!**)



Transport coefficients: electric conductivity σ_0/T

$\sigma_0 \rightarrow$ Probe of **electric properties of the QGP**

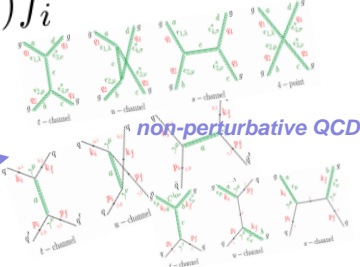
DQPM:
 σ_0/T versus (T, μ_B)



➤ Relaxation Time Approximation

$$\sigma_0^{\text{RTA}}(T, \mu_B) = \frac{e^2}{3T} \sum_{i=q, \bar{q}} q_i^2 \int \frac{d^3 p}{(2\pi)^3} \frac{\mathbf{p}^2}{E_i^2} \times \tau_i(\mathbf{p}, T, \mu_B) d_i (1 - f_i) f_i$$

relaxation time $\tau_i(\mathbf{p}, T, \mu_B) = \frac{1}{\Gamma_i(\mathbf{p}, T, \mu_B)}$
parton interaction rate

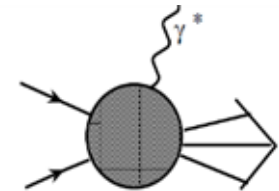


- the **QCD matter** even at $T \sim T_c$ is a much better electric conductor than **Cu or Ag** (at room temperature) by a factor of 500 !

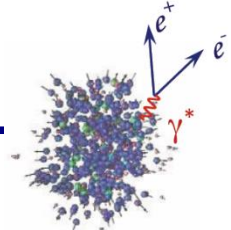
Exp. observables – photon and dilepton spectra

☐ **Photon emission:** rates for $q_0 \rightarrow 0$ are related to **electric conductivity σ_0**

$$q_0 \left. \frac{dR}{d^4 x d^3 q} \right|_{q_0 \rightarrow 0} = \frac{T}{4\pi^3} \sigma_0$$



Dilepton sources



from the QGP via partonic (q, qbar, g) interactions:

PHSD: non-perturbative QGP → DQPM (Dynamical Quasiparticle Model)



from hadronic sources:

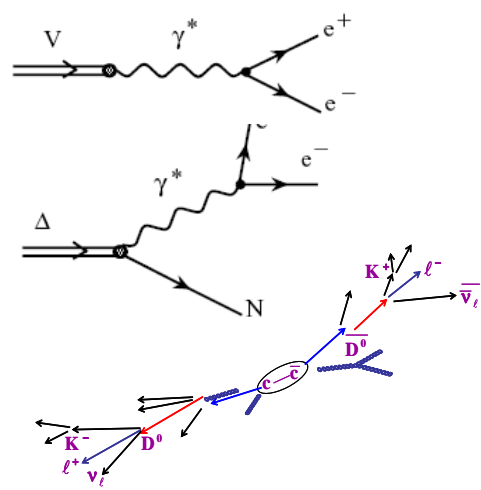
- direct decay of vector mesons ($\rho, \omega, \phi, J/\Psi, \Psi'$)

- Dalitz decay of mesons and baryons ($\pi^0, \eta, \Delta, \dots$)

- correlated D+Dbar pairs

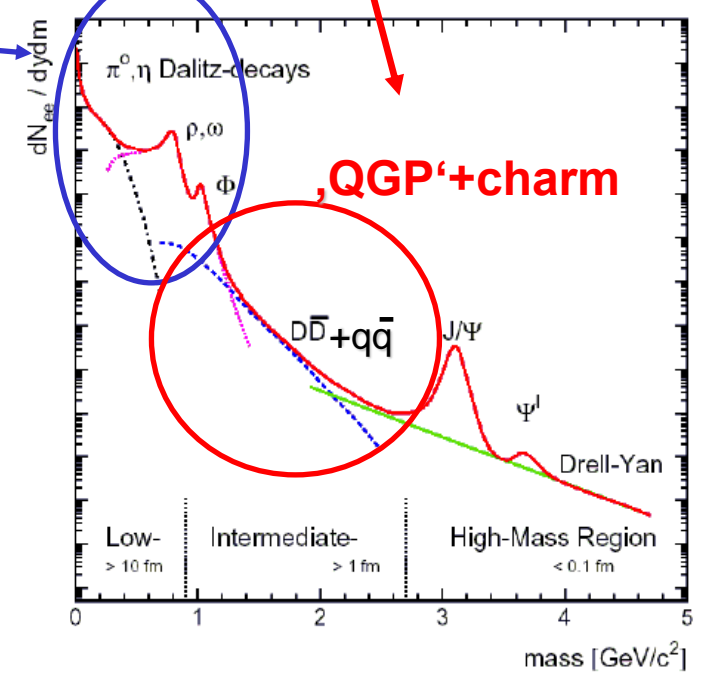
- radiation from multi-meson reactions

($\pi+\pi, \pi+\rho, \pi+\omega, \rho+\rho, \pi+a_1$) - , 4π '

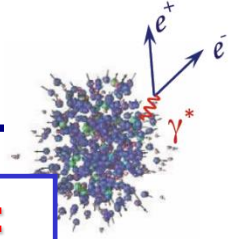


Hadronic decays

Plot from A. Drees

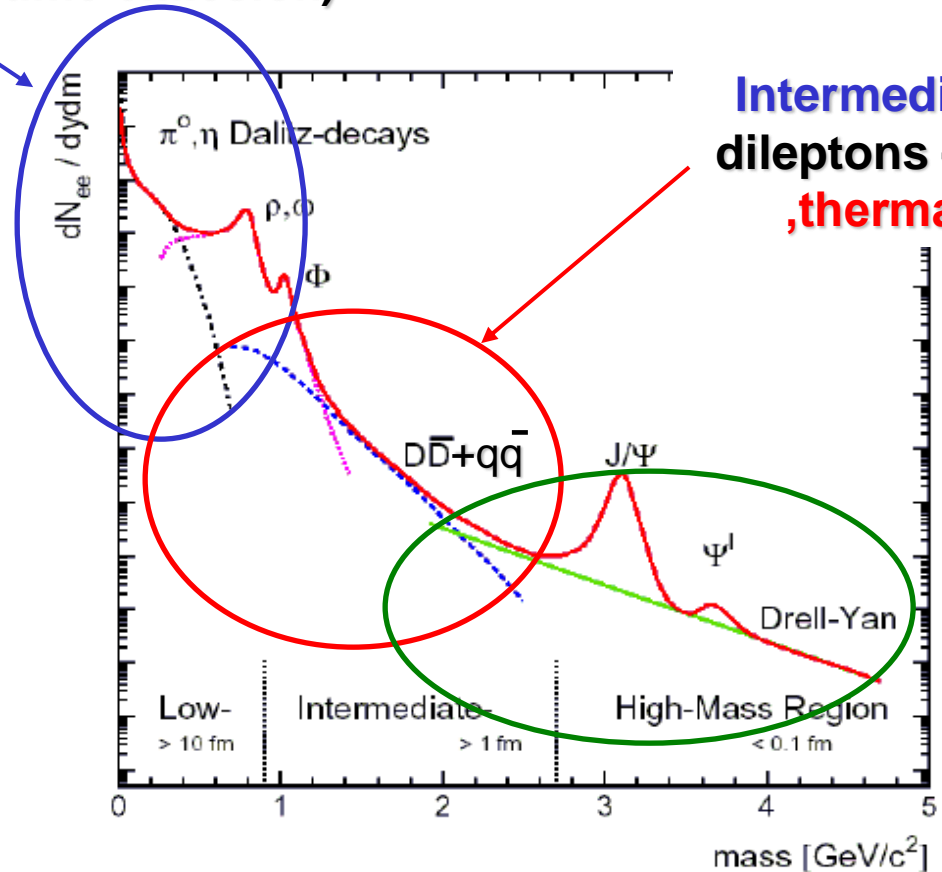


Physics with dileptons



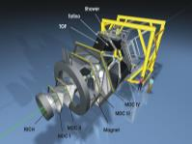
Low mass dileptons
-probe of **hadronic in-medium effects**
(late time emission)

Advantage of dileptons vs. photons:
additional „degree of freedom“ (M)
allows to disentangle various sources



Intermediate mass dileptons – probe of **„thermal QGP“**

High-mass dileptons
– probe of **pQGP and hard probes**
(early time emission)



Dileptons at SIS energies - HADES

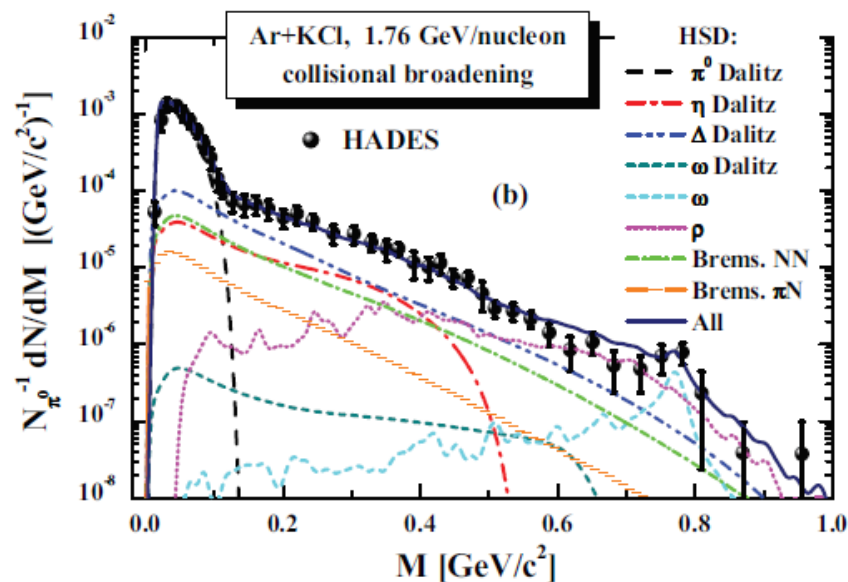
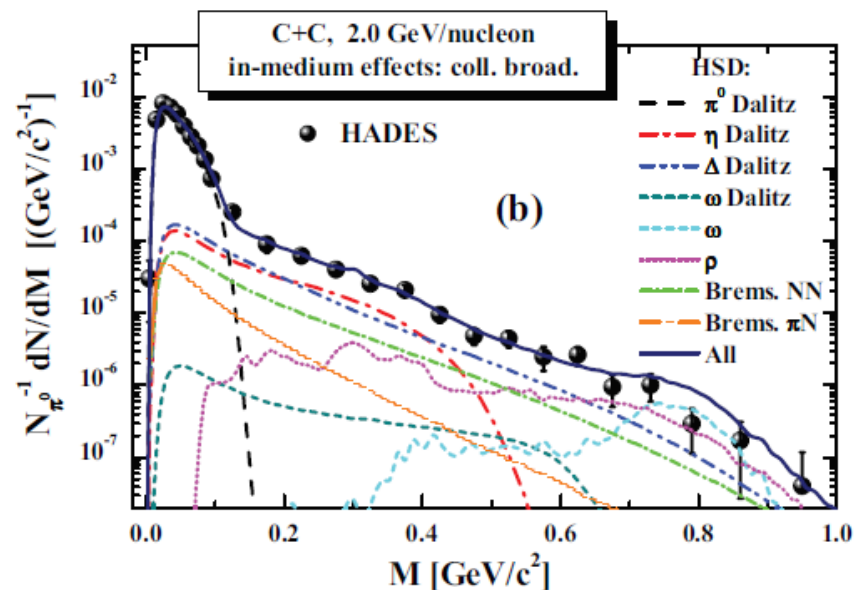
□ **HADES:** dilepton yield dN/dM scaled with the number of pions N_{π^0}

□ **Dominant hadronic sources for $M > m_{\pi^0}$:**

- η, Δ Dalitz decays
- NN bremsstrahlung
- direct ρ decay

➤ ρ meson = strongly interacting resonance
strong collisional broadening of the ρ width

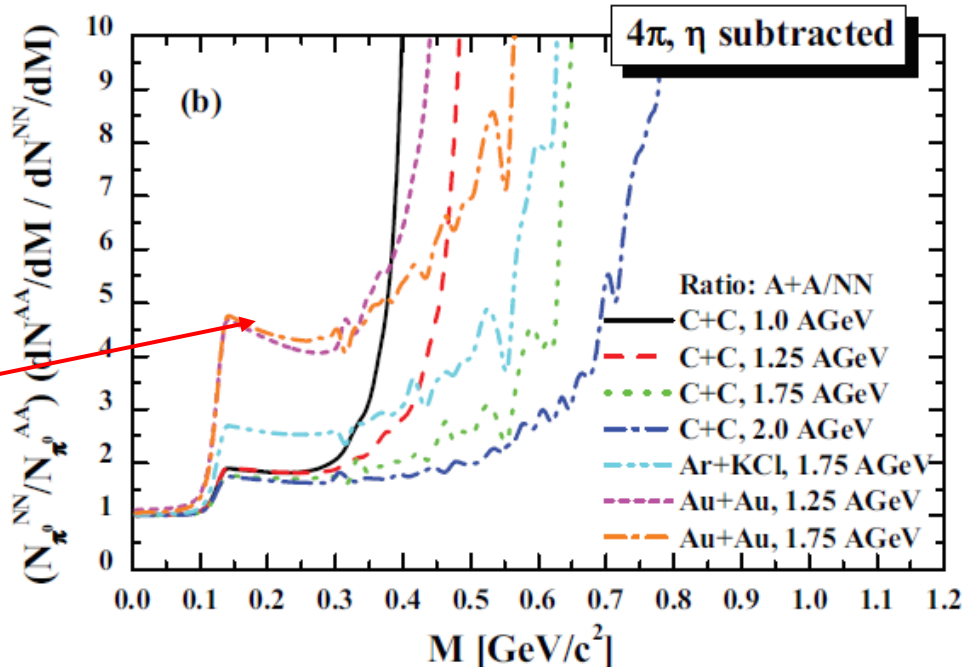
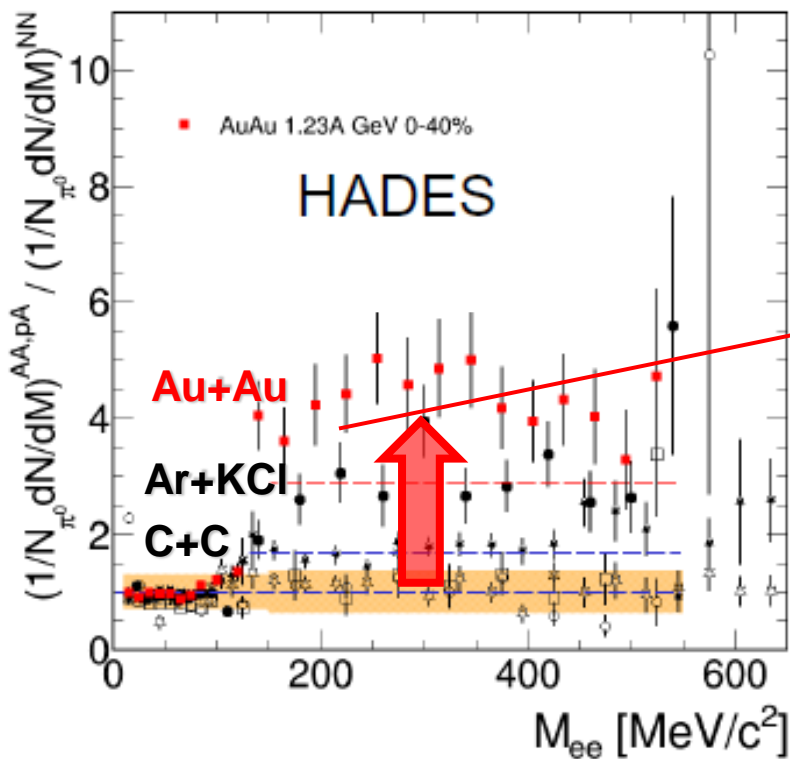
- In-medium effects are more pronounced for heavy systems such as Ar+KCl than C+C
- The peak at $M \sim 0.78$ GeV relates to ω/ρ mesons decaying in vacuum



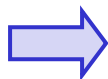
HADES, Nature Phys.15 (2019) 1040

▪ HSD predictions (2013)

HADES : Au+Au, 1.23 A GeV



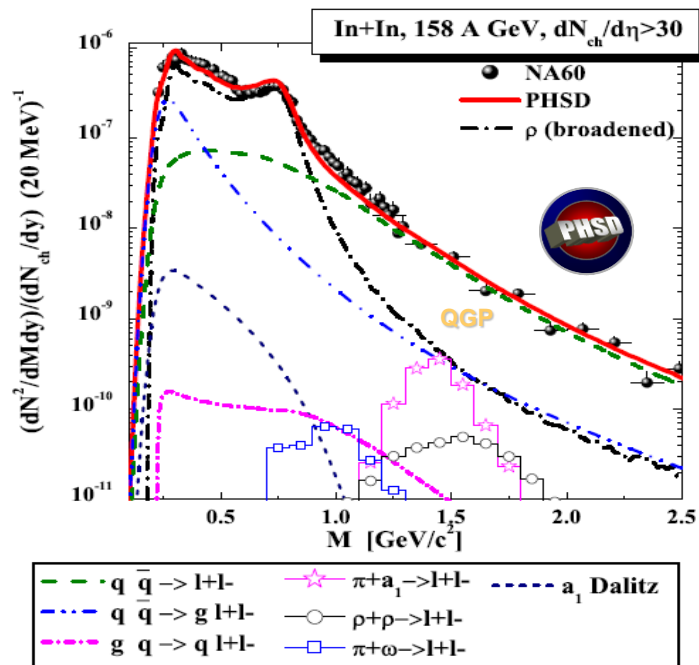
- Strong in-medium enhancement of dilepton yield in Au+Au vs. NN
- Increases with the system size!



- 1) the **multiple Δ regeneration** – dilepton emission from intermediate Δ 's which are part of the reaction cycles $\Delta \rightarrow \pi N$; $\pi N \rightarrow \Delta$ and $NN \rightarrow N\Delta$; $N\Delta \rightarrow NN$
- 2) the **pN bremsstrahlung** which scales with N_{bin} and not with N_{part} , i.e. pions

Lessons from SPS: NA60

□ Dilepton invariant mass spectra:



NA60: Eur. Phys. J. C 59 (2009) 607

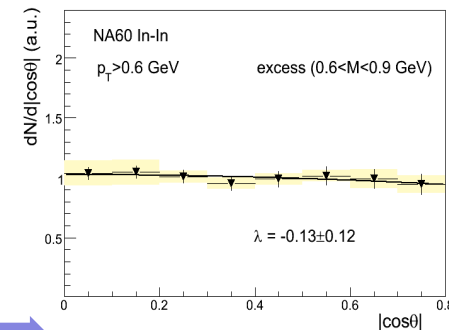
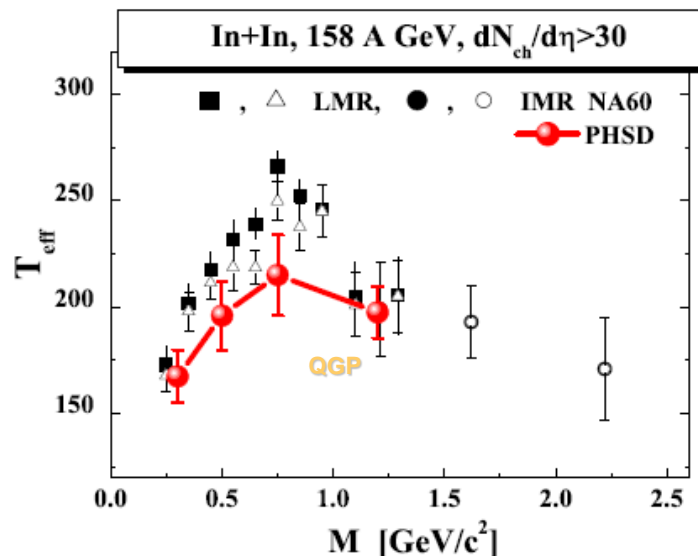
PHSD:
Linnyk et al, PRC 84 (2011) 054917

Message from SPS: (based on NA60 and CERES data)

- 1) Low mass spectra - evidence for the **in-medium broadening of ρ -mesons**
- 2) Intermediate mass spectra above 1 GeV - dominated by **partonic radiation**
- 3) The rise and fall of T_{eff} – evidence for the thermal **QGP radiation**
- 4) **Isotropic angular distribution** – indication for a **thermal origin of dimuons**

□ Inverse slope parameter T_{eff} :

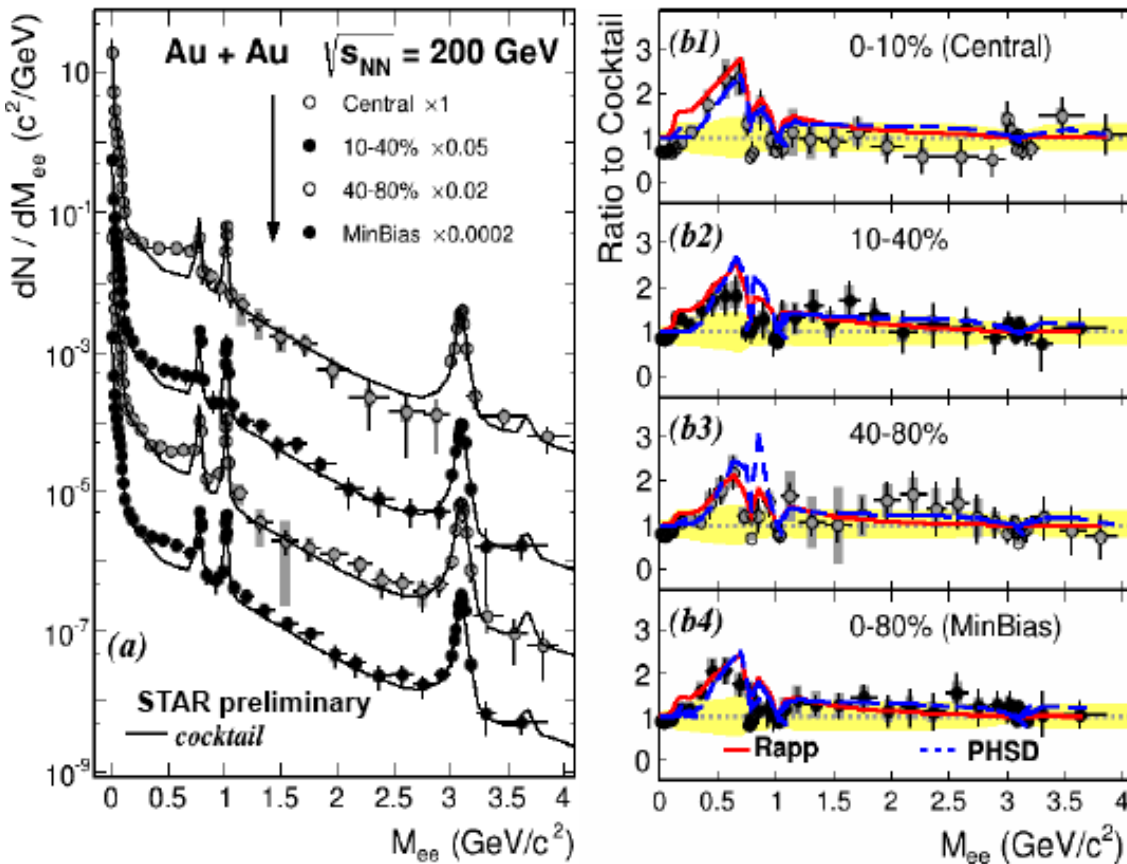
spectrum from QGP is softer than from hadronic phase since the QGP emission occurs dominantly before the collective radial flow has developed



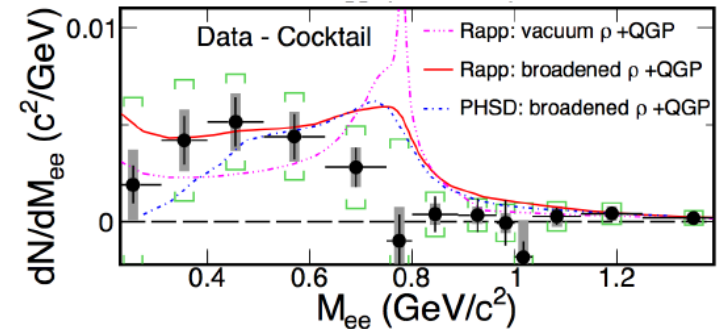
PRL 102 (2009) 222301

Dileptons at RHIC: STAR data vs model predictions

Centrality dependence of dilepton yield



Excess in low mass region, min. bias



Models:

- Fireball model – R. Rapp
- PHSD

Low masses:

collisional broadening of ρ

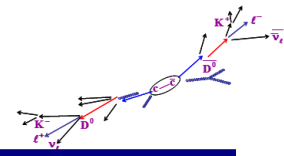
Intermediate masses:

QGP dominant

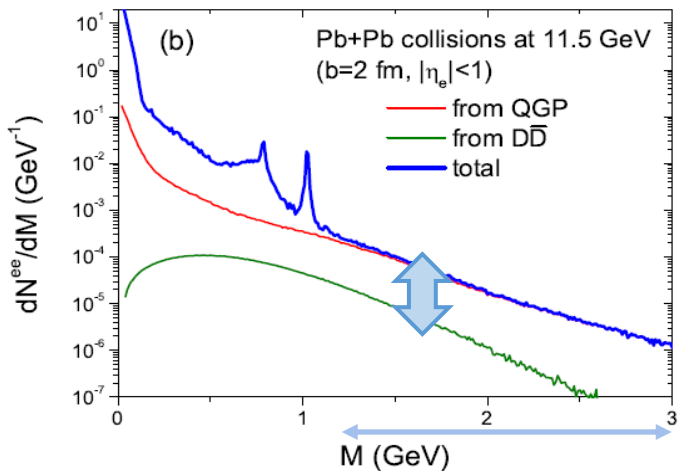
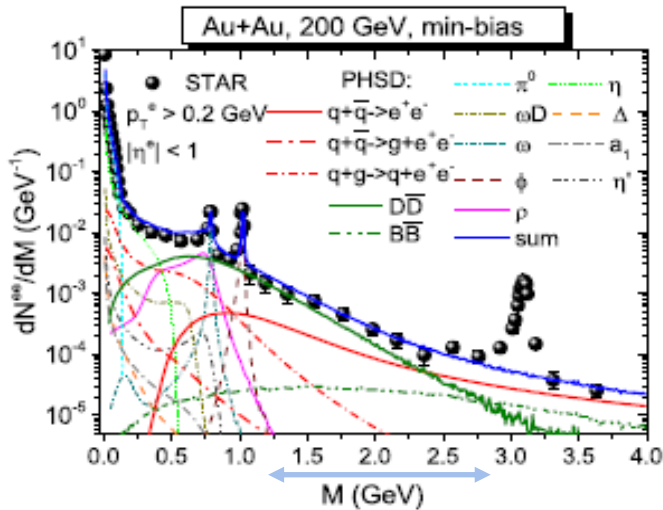
→ STAR data are described by models within a collisional broadening scenario for the vector meson spectral function + QGP

→ In-medium effects are stronger for the central A+A collisions

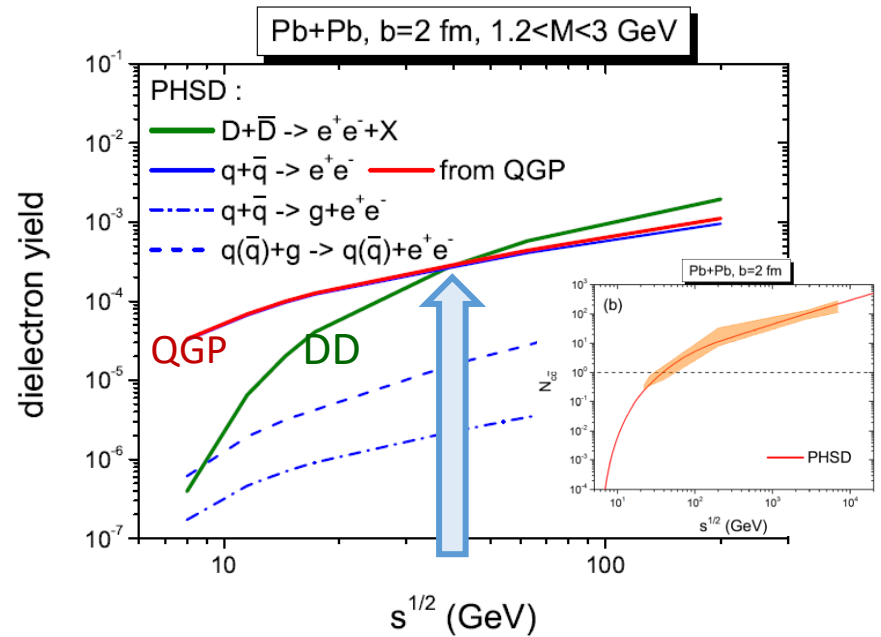
Dileptons: QGP vs charm



- STAR RHIC BES data and the ALICE data are described by PHSD within a **collisional broadening** scenario for the **vector meson** spectral functions + **QGP** + **correlated charm**



Excitation function of dilepton yield integrated for $1.2 < M < 3 \text{ GeV}$



- Dileptons from **QGP** overshadow charm dileptons **with decreasing beam energy!**

➔ Good perspectives for FAIR/NICA and RHIC BES!

Polarization: Dilepton anisotropy coefficients

E.B., M. Schafer, W. Cassing, U. Mosel, O.V. Teryaev, V. D. Toneev, Phys. Lett. B 348 (1995) 283; B 348 (1995) 325; B 362 (1995) 17, B376 (1996) 12; Z. Phys. C75 (1997)197

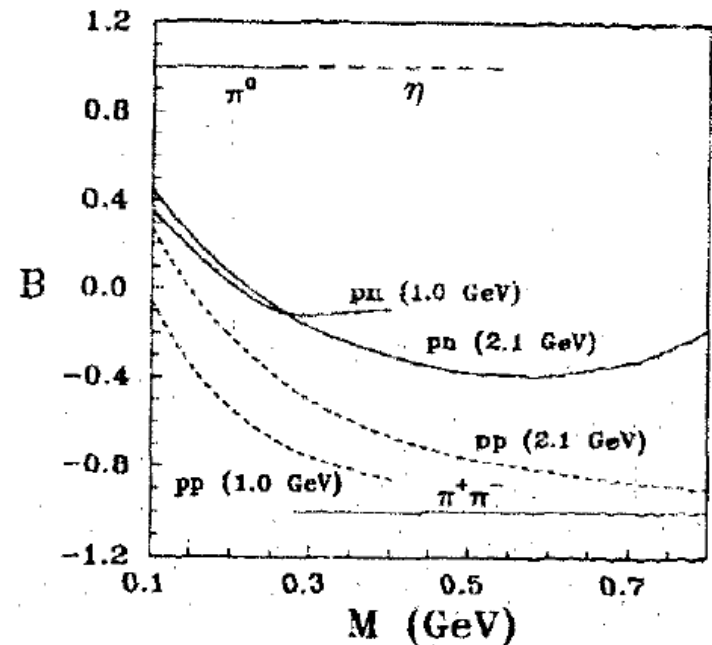
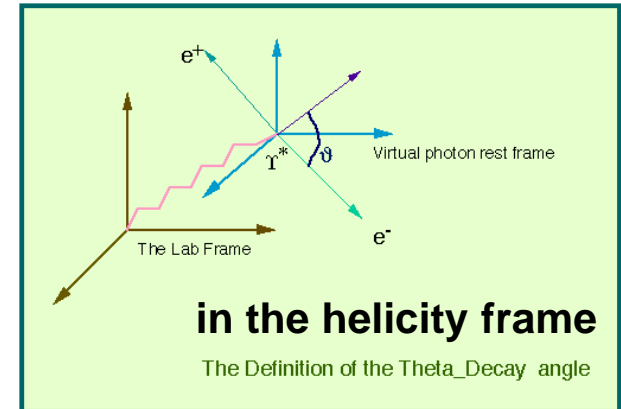
$$d\sigma/d(\cos\theta) \sim 1 + B \cos^2\theta$$

$$B = \frac{3\rho_{11} - 1}{1 - \rho_{11}}$$

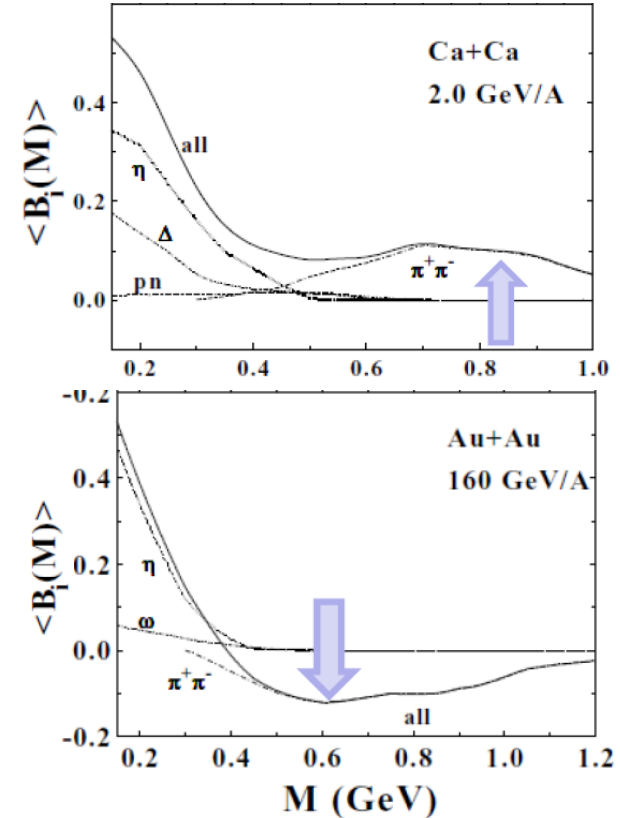
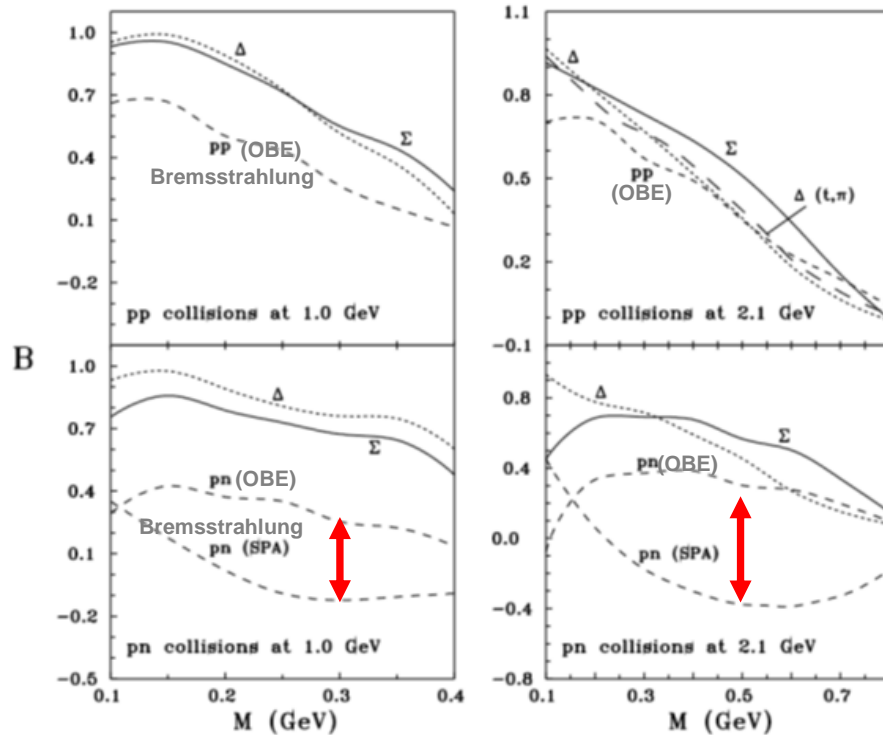
$$\rho_{00} + 2\rho_{11} = 1$$

Anisotropy coefficients for elementary channels:

- pseudoscalar mesons (e.g. π^0 and η):
B = +1
- vector mesons (e.g, ρ , ω and ϕ) from $NN \rightarrow VX$:
if no preferred spin orientation of VM
B = 0
- $\pi\pi$ annihilation: $\pi^+\pi^- \rightarrow \rho \rightarrow e^+e^-$:
 ρ wave ($L=1 \perp$ to $\pi\pi$ scattering plane)
B = -1
- Δ and N^* decays: **B $\neq 0$**
- NN and πN bremsstrahlung: **B $\neq 0$**



Dilepton anisotropy coefficients in N+N, A+A



N+N collisions:

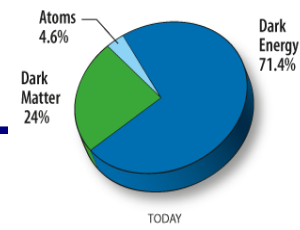
- B from Δ -decay and Bremsstrahlung (pp, pn) sensitive to the model details:
 - B from SPA $<$ B from OBE model
- Strong isospin dependence of B

A+A collisions:

- B from $\pi^+\pi^- \rightarrow \rho \rightarrow e^+e^-$ changes sign with increasing energy!
- Information on ρ polarization (depends on ρ production mechanism)

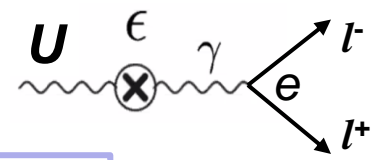
→ Opportunities for the HIC dilepton programs to study production mechanisms by polarizations!

Light dark photons searches with heavy ions



The '**vector**' portal : existence of a **U(1)-U(1)'** gauge symmetry group mixing

$$\mathcal{L}_{A'} = -\frac{1}{4}F'^{\mu\nu}F'_{\mu\nu} + \frac{1}{2} \frac{\epsilon}{\cos\theta_W} B^{\mu\nu}F'_{\mu\nu} - \frac{1}{2}m_{A'}^2 A'^{\mu}A'_{\mu}$$



Notation for 'dark photon': A' or U- boson

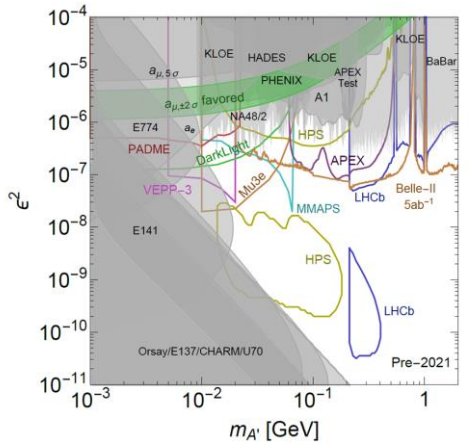
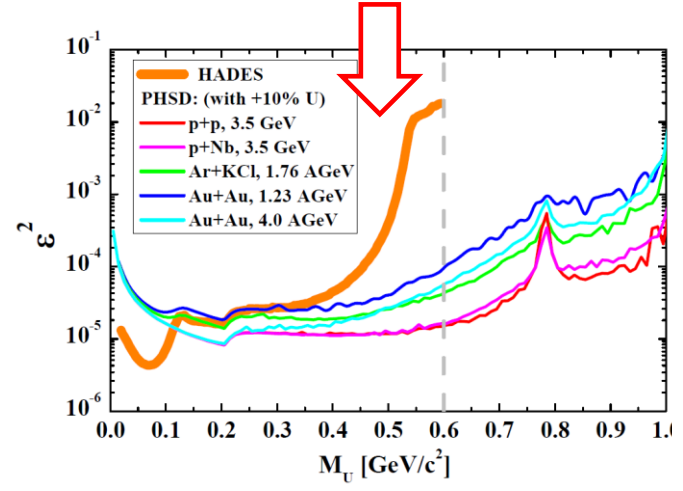
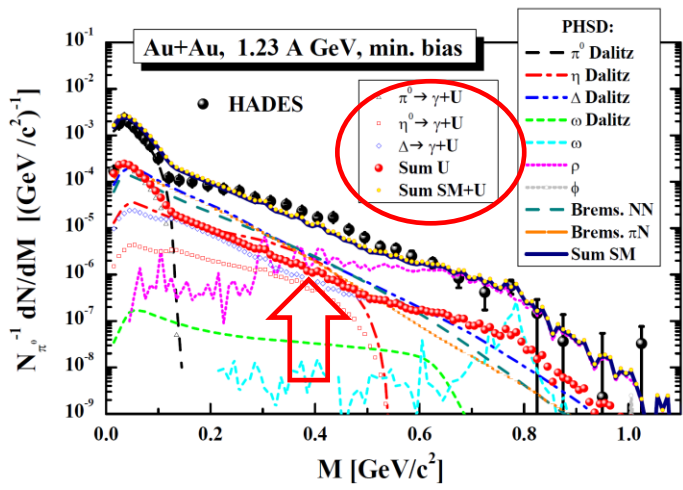
Unknown: kinetic mixing parameter ϵ and mass M_U

B. Holdom, PL B 166, 196 (1986)
B. Batell et al., PRD 80, 095024 (2009)



$\pi^0 \rightarrow \gamma + U,$
 $\eta \rightarrow \gamma + U, \quad U \rightarrow e^+e^-$
 $\Delta \rightarrow N + U$

The **upper limit for the kinetic mixing parameter $\epsilon^2(M_U)$** of light dark photons extracted from the **PHSD** dilepton spectra - with 10% allowed surplus of the total SM yield by an additional **DM** yield at given M:





Summary

I. SIS energies:

Strong in-medium enhancement of dilepton yield in Au+Au vs. NN (which increases with the system size) – due to the multiple Δ regeneration, pN bremsstrahlung and in-medium ρ -contribution

II. High energies:

□ Low dilepton masses:

Dilepton spectra show sizeable changes due to the in-medium effects – **modification of the properties of vector mesons** (as collisional broadening)

□ Intermediate dilepton masses $M > 1.2$ GeV :

- Dominant sources : **QGP** ($q\bar{q}$) and **correlated charm D/Dbar**
- Fraction of QGP **grows** with increasing energy; however, the relative contribution of QGP dileptons to dileptons from charm pairs increases with decreasing energy

→ **In-medium effects** can be observed at all energies from SIS to LHC

→ **QGP contribution overshines charm with decreasing energy**

Good perspectives for FAIR / NICA / RHIC BES !

→ Study of **polarization phenomena** with dileptons

→ Possibility to search for **dark photons** with dilepton heavy-ion experiments