

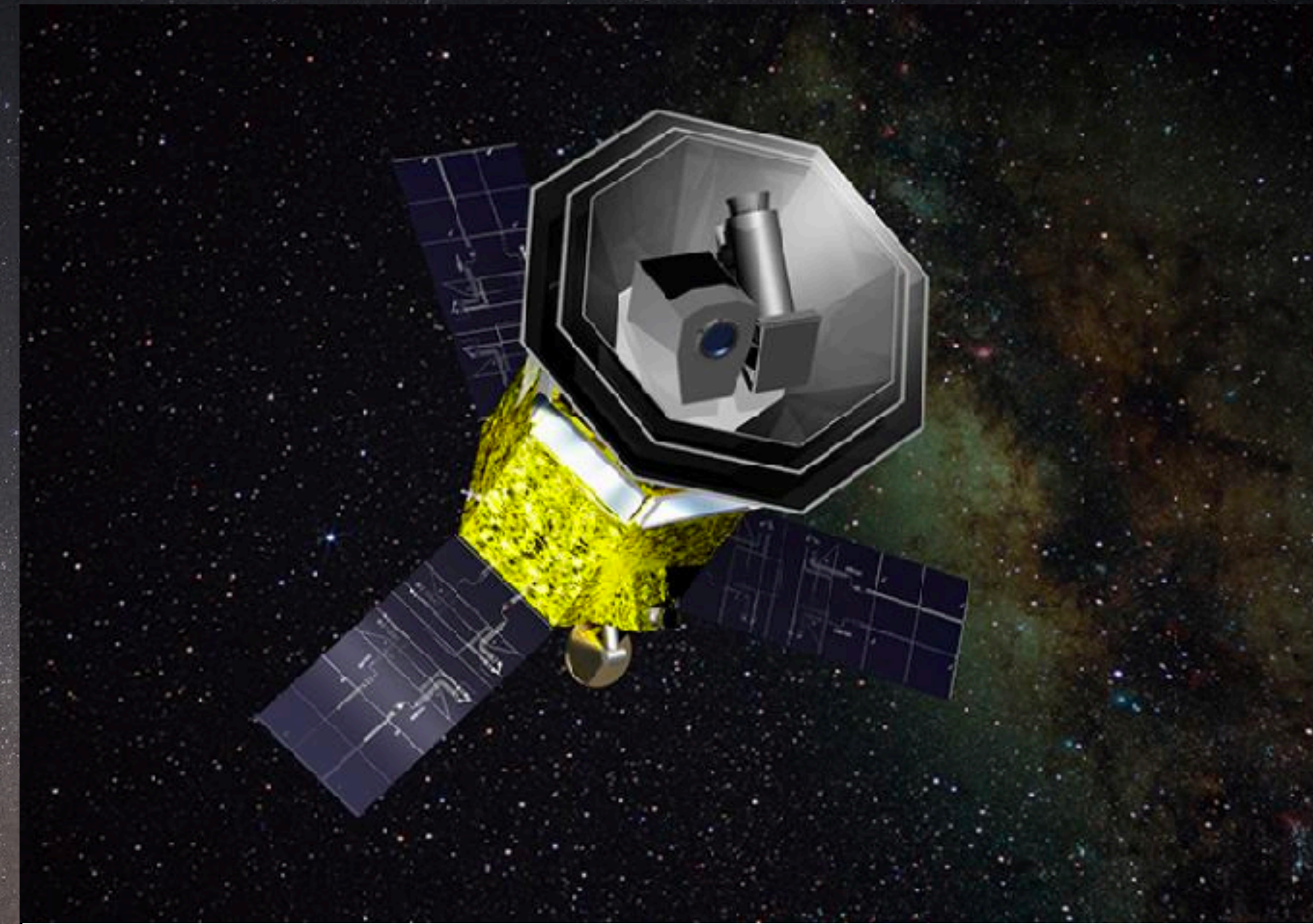
Moment expansion and the challenge of polarized dust SED complexity for B mode detections

Léo Vacher - Jonathan Aumont - Ludovic Montier - François Boulanger - Susanna Azzoni - Mathieu Remazeilles

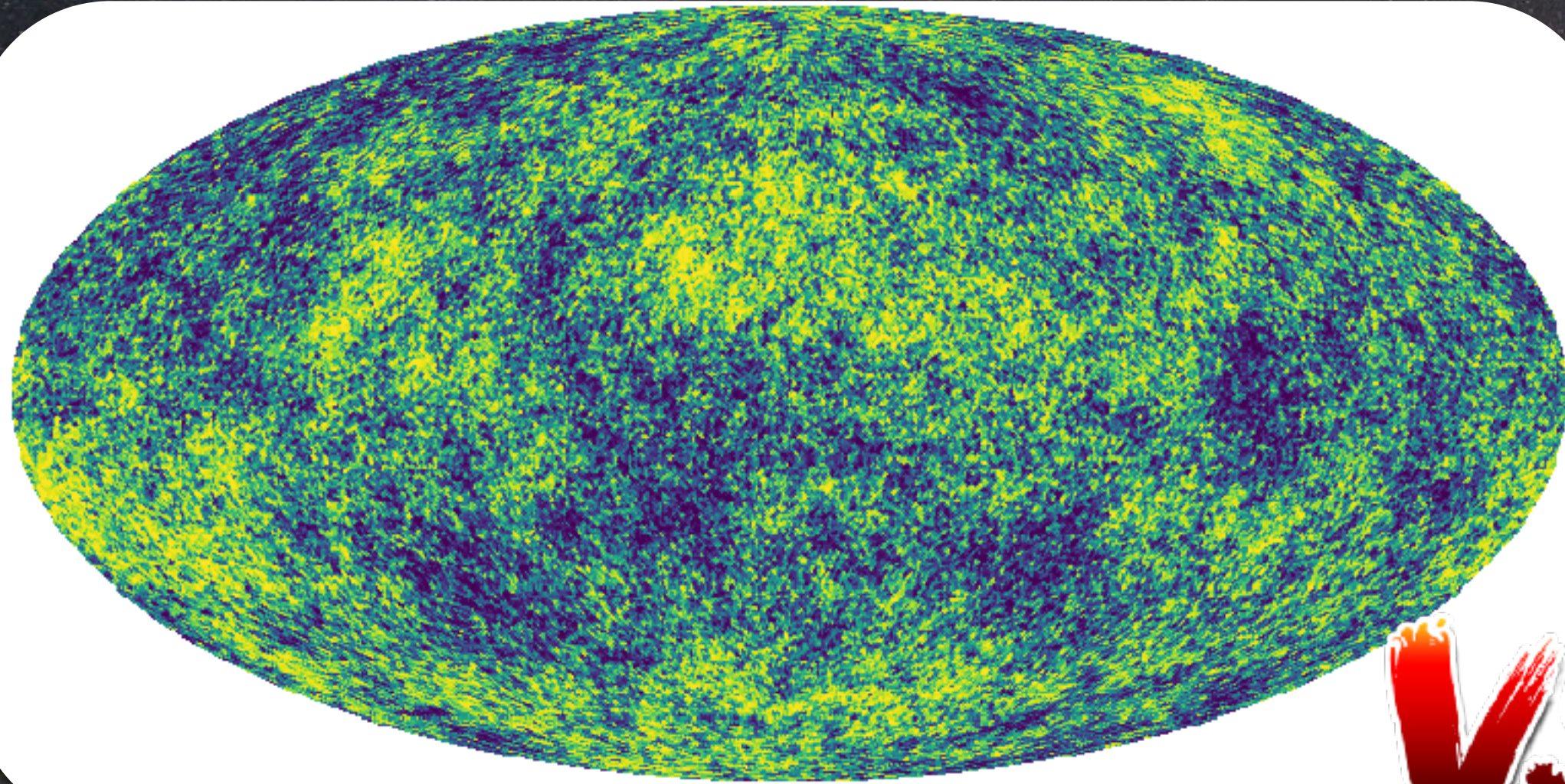


LiteBIRD and the *B*-modes quest

- **JAXA** project. Phase A **CNES**. **ESA, NASA, CSA** involved
- **Lite** (*Light*) satellite for the studies of **B**-mode polarization and *Inflation from cosmic background Radiation Detection*
- Build to reach $\delta r = 1 \times 10^{-3}$
- **3** telescopes LFT, MFT, HFT
- Expected in **2029** at **L2** for more than **3 years** of observation



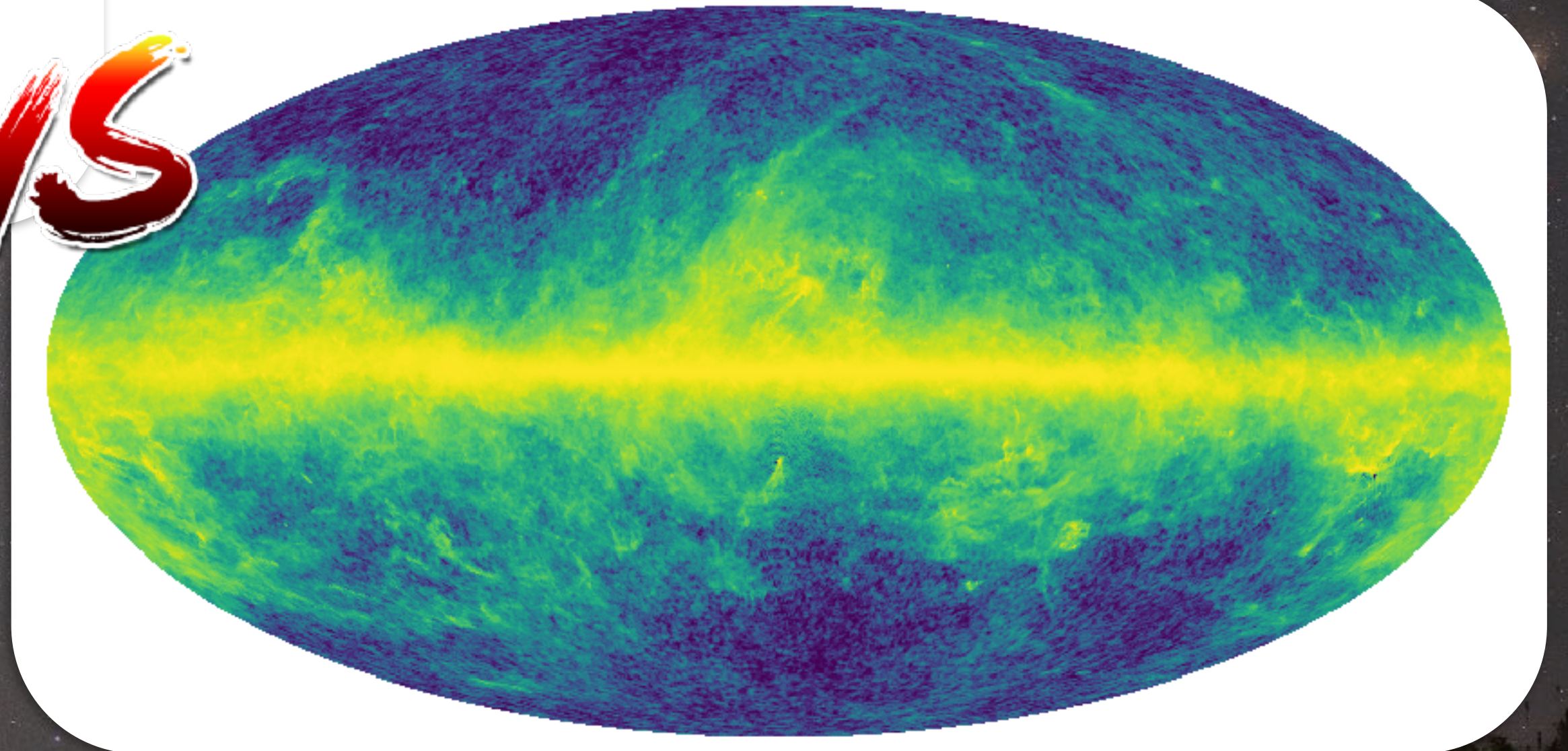
Foregrounds



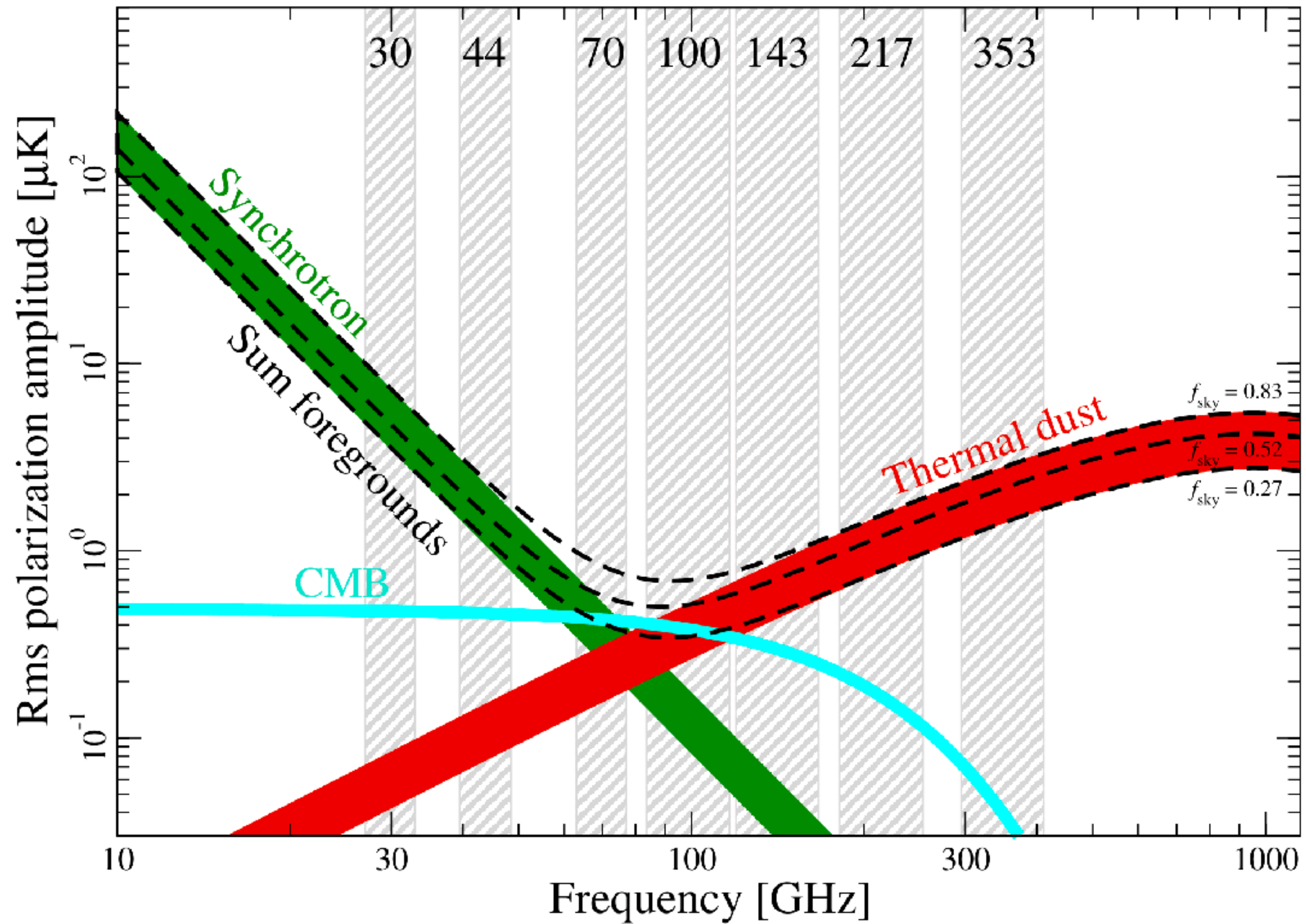
Astrophysical sources emitting mainly in CMB's wavelength interval :

- Dust thermal emission
- Synchrotron
- Free-Free/ Brehmstrahlung

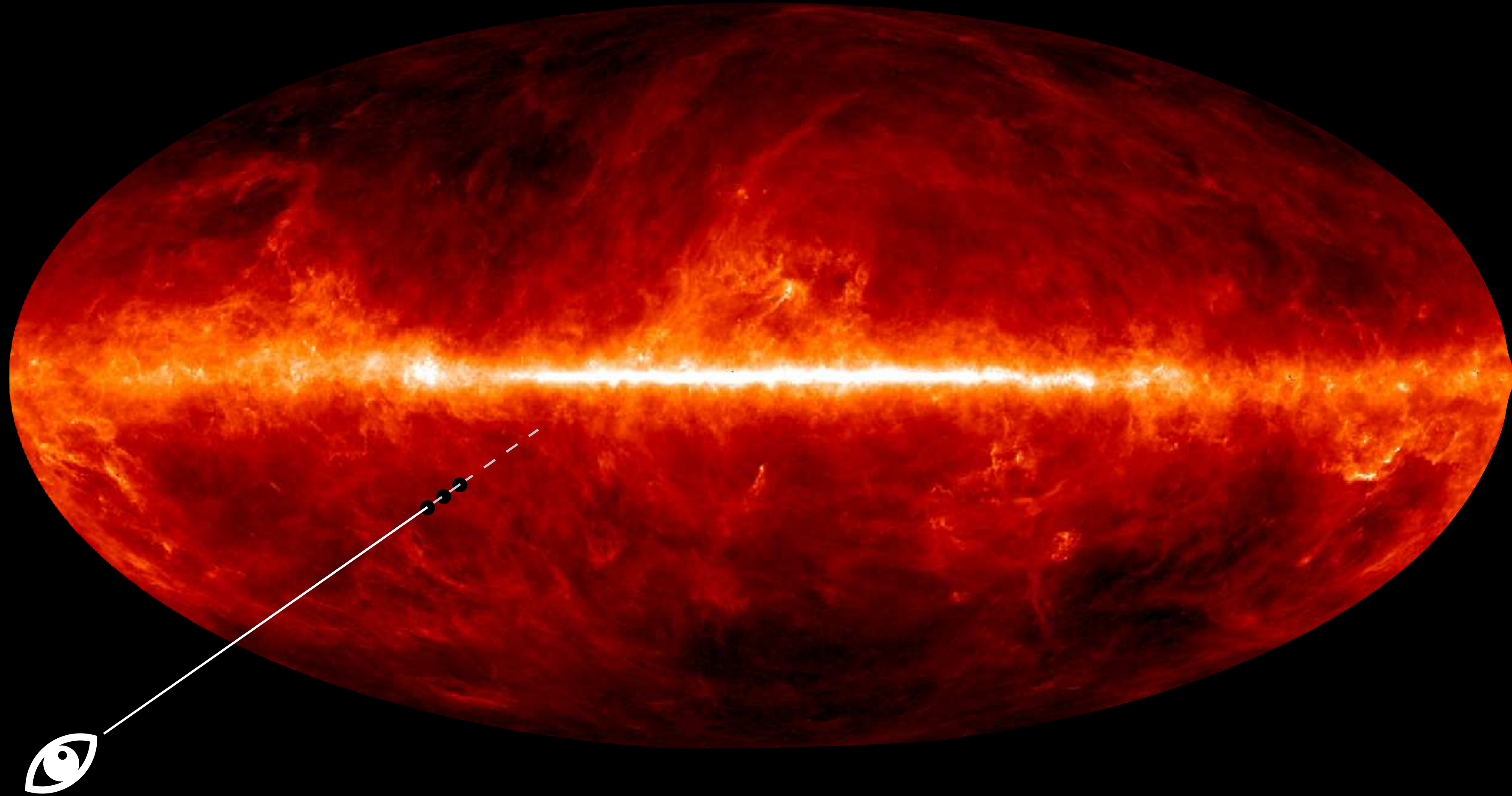
VS



I. Foregrounds : problematic



Dust — Averaging SEDs

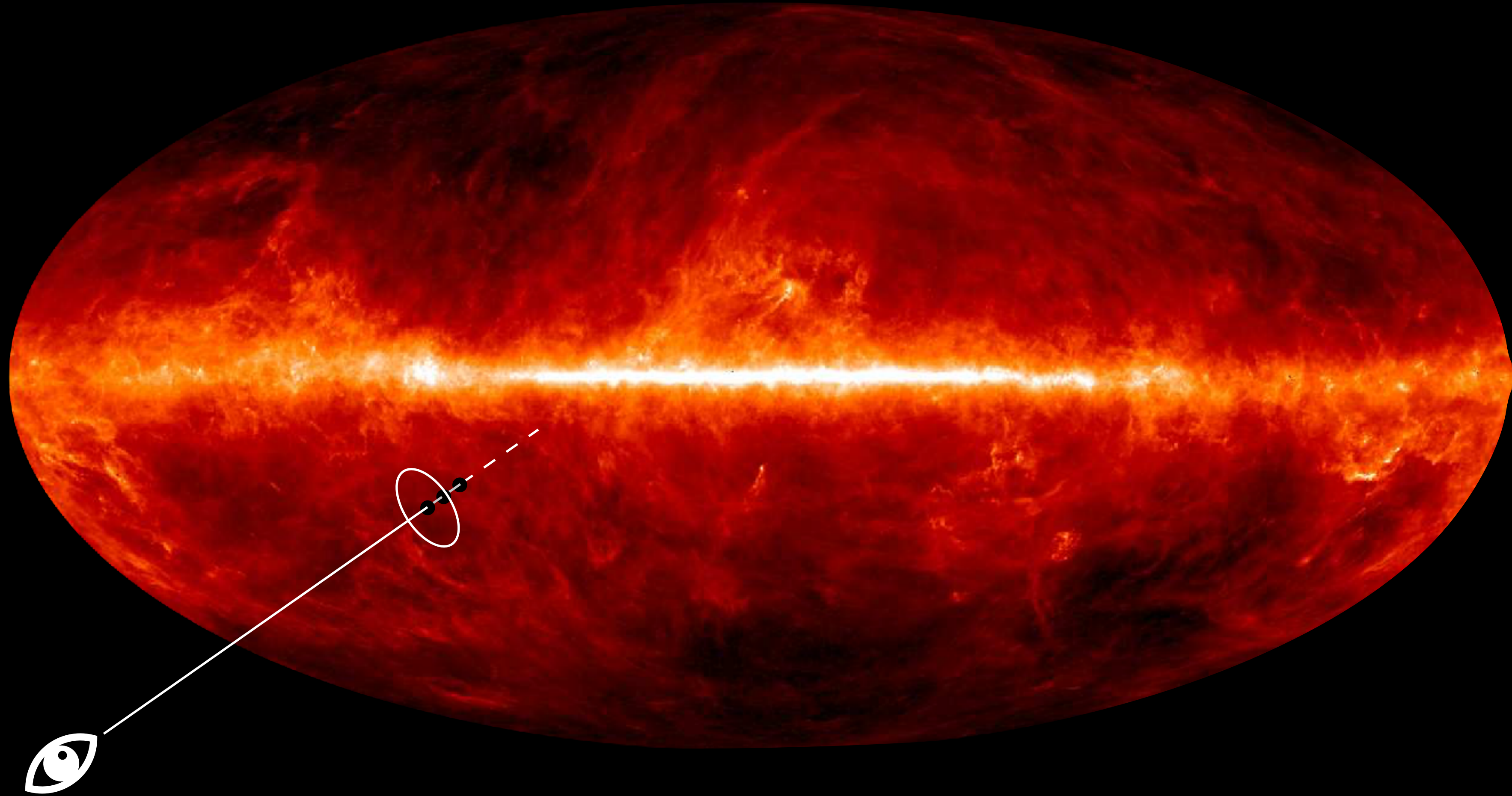


Modified black-body SED in every volume element

★ Line-of-sight average (*always there!*)

[Planck 2018 M]

Dust — Averaging SEDs

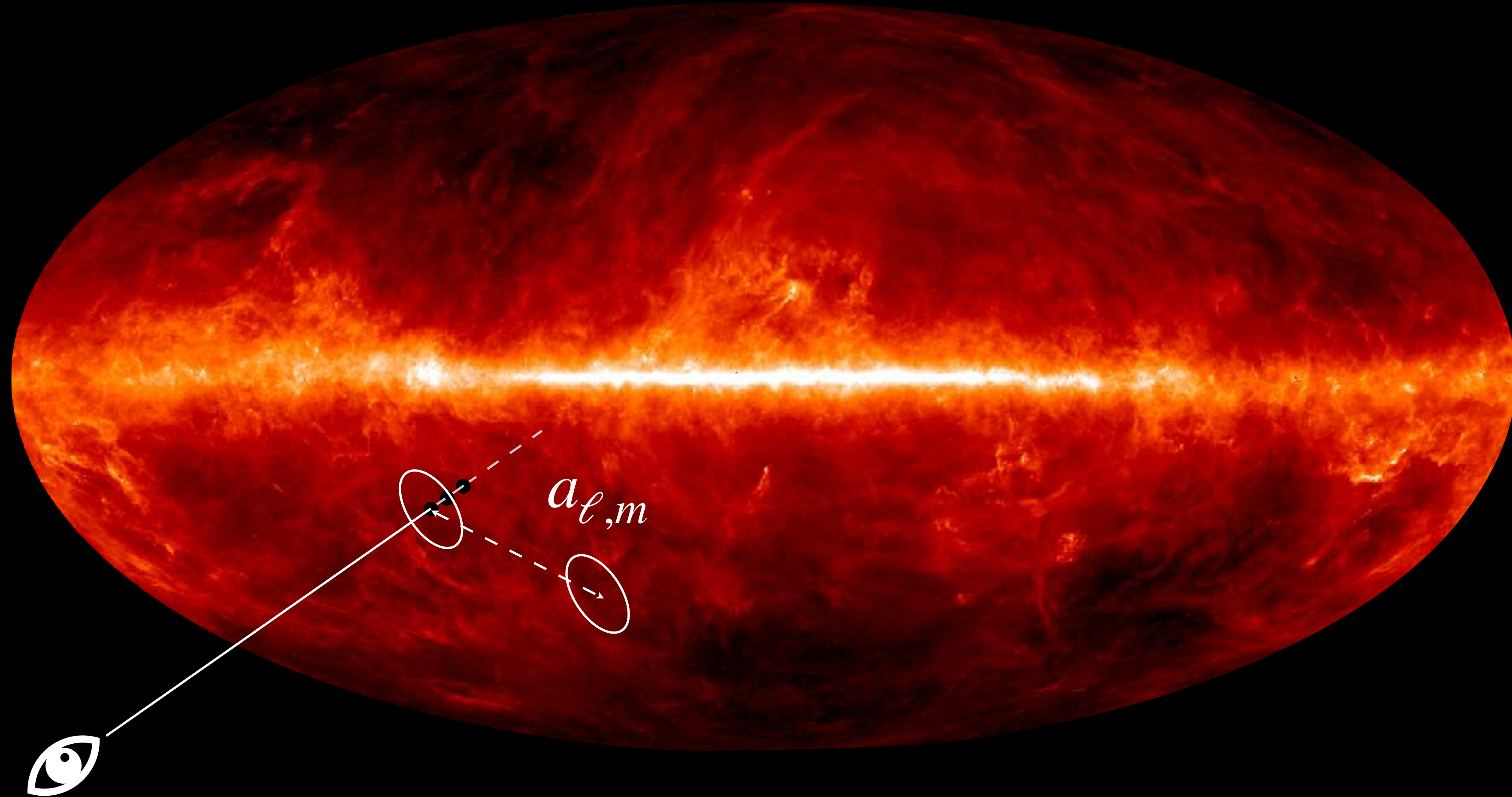


Modified black-body SED in every volume element

★ Line-of-sight average (*always there!*)

★ Experimental beam and frequency average

Dust — Averaging SEDs



[Planck 2018 M]

Modified black-body SED in every volume element

- ★ Line-of-sight average (*always there!*)
- ★ Experimental beam and frequency average
- ★ Map operations average (e.g., spherical harmonic expansion)

The modified black body

Canonical model : Empirical + Non-linear

Spectral index

Modified Black-Body

dust map

$$I_D(\nu, \vec{n}) = \left(\frac{\nu}{\nu_0}\right)^{\beta_0} \frac{B_\nu(T_0)}{B_{\nu_0}(T_0)} A(\vec{n}) = \frac{I_\nu(\beta_0, T_0)}{I_{\nu_0}(\beta_0, T_0)} A(\vec{n}).$$

Black body function

The moment expansion in pixel space

Taylor inspired expansion around the MBB in β (T fixed at T_0):

$$I_D(\nu, \vec{n}) = \frac{I_\nu(\beta_0, T_0)}{I_{\nu_0}(\beta_0, T_0)} \left[A(\vec{n}) + \omega_1(\vec{n}) \ln \left(\frac{\nu}{\nu_0} \right) + \frac{1}{2} \omega_2(\vec{n}) \ln^2 \left(\frac{\nu}{\nu_0} \right) + \frac{1}{6} \omega_3(\vec{n}) \ln^3 \left(\frac{\nu}{\nu_0} \right) + \dots \right].$$

MBB (order 0)

+ order 1

+ order 2

+ order 3 + ...

[Chluba et al., 2017]

The moment expansion in pixel space

« $\Omega(\nu, T)$ »

Taylor inspired expansion around the MBB in β and T :

$$I_D(\nu, \vec{n}) = \frac{I_\nu(\beta_0, T_0)}{I_{\nu_0}(\beta_0, T_0)} \left(1 + \omega_1^\beta \ln \left(\frac{\nu}{\nu_0} \right) + \frac{h\omega_1^T}{kT_0^2} \left(\frac{\nu e^{\frac{h\nu}{kT_0}}}{e^{\frac{h\nu}{kT_0}} - 1} - \frac{\nu_0 e^{\frac{h\nu_0}{kT_0}}}{e^{\frac{h\nu_0}{kT_0}} - 1} \right) + \dots \right)$$

MBB (order 0)

+ order 1 beta + order 1 T + ...

[Chluba et al., 2017]

The moment expansion in harmonic space

β only (Mangilli et al 2021)

β and T

$$\mathcal{D}_\ell(\nu_i \times \nu_j) = \frac{I_{\nu_i}(\beta_0(\ell), T_0(\ell)) I_{\nu_j}(\beta_0(\ell), T_0(\ell))}{I_{\nu_0}^2(\beta_0(\ell), T_0(\ell))} \times \left\{ \begin{array}{l} 0^{\text{th}} \text{ order } \left\{ \mathcal{D}_\ell^{A \times A} \right. \\ 1^{\text{st}} \text{ order } \left\{ \begin{array}{l} + \left[\ln\left(\frac{\nu_i}{\nu_0}\right) + \ln\left(\frac{\nu_j}{\nu_0}\right) \right] \mathcal{D}_\ell^{A \times \omega_1^\beta} \\ + \left[\ln\left(\frac{\nu_i}{\nu_0}\right) \ln\left(\frac{\nu_j}{\nu_0}\right) \right] \mathcal{D}_\ell^{\omega_1^\beta \times \omega_1^\beta} \end{array} \right. \\ 2^{\text{nd}} \text{ order } \left\{ \begin{array}{l} + \frac{1}{2} \left[\ln^2\left(\frac{\nu_i}{\nu_0}\right) + \ln^2\left(\frac{\nu_j}{\nu_0}\right) \right] \mathcal{D}_\ell^{A \times \omega_2^\beta} \\ + \frac{1}{2} \left[\ln\left(\frac{\nu_i}{\nu_0}\right) \ln^2\left(\frac{\nu_j}{\nu_0}\right) + \ln\left(\frac{\nu_j}{\nu_0}\right) \ln^2\left(\frac{\nu_i}{\nu_0}\right) \right] \mathcal{D}_\ell^{\omega_1^\beta \times \omega_2^\beta} \\ + \frac{1}{4} \left[\ln^2\left(\frac{\nu_i}{\nu_0}\right) \ln^2\left(\frac{\nu_j}{\nu_0}\right) \right] \mathcal{D}_\ell^{\omega_2^\beta \times \omega_2^\beta} \\ + \dots \end{array} \right. \end{array} \right\}.$$

$$\mathcal{D}_\ell(\nu_i \times \nu_j) = \frac{I_{\nu_i}(\beta_0(\ell), T_0(\ell)) I_{\nu_j}(\beta_0(\ell), T_0(\ell))}{I_{\nu_0}(\beta_0(\ell), T_0(\ell))^2} \cdot \left\{ \begin{array}{l} 0^{\text{th}} \text{ order } \left\{ \mathcal{D}_\ell^{A \times A} \right. \\ 1^{\text{st}} \text{ order } \left\{ \begin{array}{l} + \mathcal{D}_\ell^{A \times \omega_1^\beta} \left(\ln\left(\frac{\nu_i}{\nu_0}\right) + \ln\left(\frac{\nu_j}{\nu_0}\right) \right) \\ + \mathcal{D}_\ell^{\omega_1^\beta \times \omega_1^\beta} \left(\ln\left(\frac{\nu_i}{\nu_0}\right) \ln\left(\frac{\nu_j}{\nu_0}\right) \right) \\ + \mathcal{D}_\ell^{A \times \omega_1^T} (\Omega_i + \Omega_j - 2\Omega_0) \\ + \mathcal{D}_\ell^{\omega_1^\beta \times \omega_1^T} \left(\ln\left(\frac{\nu_j}{\nu_0}\right) (\Omega_i - \Omega_0) + \ln\left(\frac{\nu_i}{\nu_0}\right) (\Omega_j - \Omega_0) \right) \\ + \mathcal{D}_\ell^{\omega_1^T \times \omega_1^T} (\Omega_i - \Omega_0) (\Omega_j - \Omega_0) + \dots \end{array} \right. \end{array} \right\}$$

The moment expansion in harmonic space

β only (Mangilli et al)

β and T

$$\mathcal{D}_\ell = \frac{\ell(\ell+1)}{2\pi} \mathcal{C}_\ell$$

$$\mathcal{D}_\ell(\nu_i \times \nu_j) = \mathcal{D}_\ell(\text{Map}(\nu_i), \text{Map}(\nu_j))$$

$$\mathcal{D}_\ell(\nu_i \times \nu_j) = \frac{I_{\nu_i}(\beta_0(\ell), T_0(\ell)) I_{\nu_j}(\beta_0(\ell), T_0(\ell))}{I_{\nu_0}^2(\beta_0(\ell), T_0(\ell))}$$



$$\begin{aligned} &0^{\text{th}} \text{ order} \left\{ \mathcal{D}_\ell^{A \times A} \right. \\ &1^{\text{st}} \text{ order} \left\{ \begin{aligned} &+ \left[\ln\left(\frac{\nu_i}{\nu_0}\right) + \ln\left(\frac{\nu_j}{\nu_0}\right) \right] \mathcal{D}_\ell^{A \times \omega_1^\beta} \\ &+ \left[\ln\left(\frac{\nu_i}{\nu_0}\right) \ln\left(\frac{\nu_j}{\nu_0}\right) \right] \mathcal{D}_\ell^{\omega_1^\beta \times \omega_1^\beta} \end{aligned} \right. \\ &2^{\text{nd}} \text{ order} \left\{ \begin{aligned} &+ \frac{1}{2} \left[\ln^2\left(\frac{\nu_i}{\nu_0}\right) + \ln^2\left(\frac{\nu_j}{\nu_0}\right) \right] \mathcal{D}_\ell^{A \times \omega_2^\beta} \\ &+ \frac{1}{2} \left[\ln\left(\frac{\nu_i}{\nu_0}\right) \ln^2\left(\frac{\nu_j}{\nu_0}\right) + \ln\left(\frac{\nu_j}{\nu_0}\right) \ln^2\left(\frac{\nu_i}{\nu_0}\right) \right] \mathcal{D}_\ell^{\omega_1^\beta \times \omega_2^\beta} \\ &+ \frac{1}{4} \left[\ln^2\left(\frac{\nu_i}{\nu_0}\right) \ln^2\left(\frac{\nu_j}{\nu_0}\right) \right] \mathcal{D}_\ell^{\omega_2^\beta \times \omega_2^\beta} \\ &+ \dots \end{aligned} \right. \end{aligned}$$

$$\mathcal{D}_\ell(\nu_i \times \nu_j) = \frac{I_{\nu_i}(\beta_0(\ell), T_0(\ell)) I_{\nu_j}(\beta_0(\ell), T_0(\ell))}{I_{\nu_0}(\beta_0(\ell), T_0(\ell))^2} \cdot \left\{ \begin{aligned} &0^{\text{th}} \text{ order} \left\{ \mathcal{D}_\ell^{A \times A} \right. \\ &1^{\text{st}} \text{ order} \left\{ \begin{aligned} &+ \mathcal{D}_\ell^{A \times \omega_1^\beta} \left(\ln\left(\frac{\nu_i}{\nu_0}\right) + \ln\left(\frac{\nu_j}{\nu_0}\right) \right) \\ &+ \mathcal{D}_\ell^{\omega_1^\beta \times \omega_1^\beta} \left(\ln\left(\frac{\nu_i}{\nu_0}\right) \ln\left(\frac{\nu_j}{\nu_0}\right) \right) \\ &+ \mathcal{D}_\ell^{A \times \omega_1^T} (\Omega_i + \Omega_j - 2\Omega_0) \\ &+ \mathcal{D}_\ell^{\omega_1^\beta \times \omega_1^T} \left(\ln\left(\frac{\nu_j}{\nu_0}\right) (\Omega_i - \Omega_0) + \ln\left(\frac{\nu_i}{\nu_0}\right) (\Omega_j - \Omega_0) \right) \\ &+ \mathcal{D}_\ell^{\omega_1^T \times \omega_1^T} (\Omega_i - \Omega_0) (\Omega_j - \Omega_0) + \dots \end{aligned} \right. \end{aligned} \right.$$

The moment expansion in harmonic space

β only (Mangilli et al 2021)

β and T

$$\mathcal{D}_\ell(\nu_i \times \nu_j) = \frac{I_{\nu_i}(\beta_0) \dots I_{\nu_j}(\beta_0) \dots}{I_{\nu_0}(\beta_0) \dots I_{\nu_0}(\beta_0) \dots}$$

0th order { \mathcal{D}

1st order { +
+
+ }

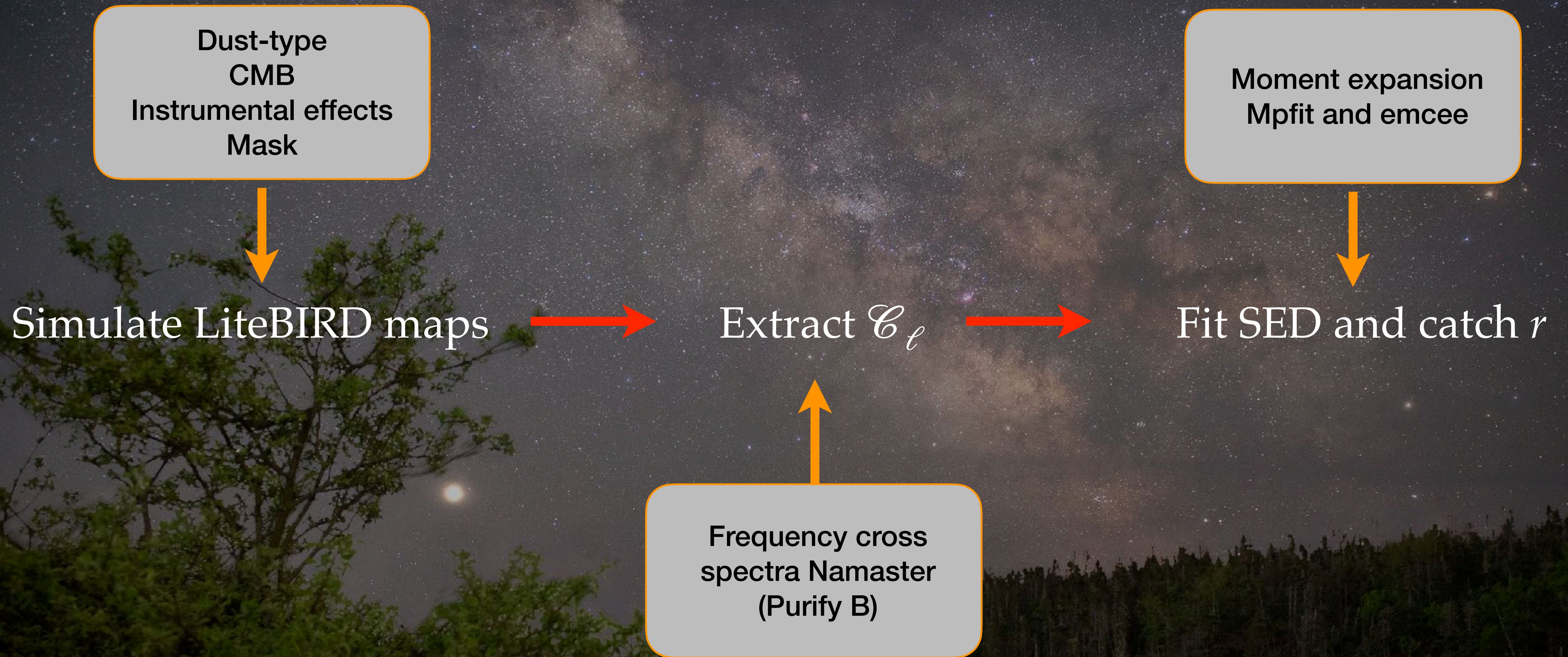
2nd order { +
+
+
+ ... }

- ★ New approach that allow to describe, at the angular power spectrum level, the dust SED distortions due to line-of-sight, beam and map operations averages
- ★ Naturally describes, in a model-independent way, the effect of spatial variations of the dust SED, additional dust components and in a more general way of dust frequency decorrelation!

$$\begin{aligned}
 & + \frac{1}{2} \left[\ln \left(\frac{\nu_i}{\nu_0} \right) \ln^2 \left(\frac{\nu_j}{\nu_0} \right) + \ln \left(\frac{\nu_j}{\nu_0} \right) \ln^2 \left(\frac{\nu_i}{\nu_0} \right) \right] \mathcal{D}_\ell^{\omega_1^\beta \times \omega_2^\beta} \\
 & + \frac{1}{4} \left[\ln^2 \left(\frac{\nu_i}{\nu_0} \right) \ln^2 \left(\frac{\nu_j}{\nu_0} \right) \right] \mathcal{D}_\ell^{\omega_2^\beta \times \omega_2^\beta} \\
 & + \dots
 \end{aligned}$$

$$\begin{aligned}
 \text{1st order} \left\{ \right. & + \mathcal{D}_\ell^{A \times \omega_1^T} (\Omega_i + \Omega_j - 2\Omega_0) \\
 & + \mathcal{D}_\ell^{\omega_1^\beta \times \omega_1^T} \left(\ln \left(\frac{\nu_j}{\nu_0} \right) (\Omega_i - \Omega_0) + \ln \left(\frac{\nu_i}{\nu_0} \right) (\Omega_j - \Omega_0) \right) \\
 & \left. + \mathcal{D}_\ell^{\omega_1^T \times \omega_1^T} (\Omega_i - \Omega_0) (\Omega_j - \Omega_0) + \dots \right\}
 \end{aligned}$$

Roadmap



LiteBIRD simulations with various dust components

dust models :

- ★ **d0** : $\beta = 1.54$, $T = 20$ K
- ★ **"d1T"** : $\beta(\vec{n})$, 22 K
- ★ **d1** : $\beta(\vec{n})$, $T(\vec{n})$

CMB :

- ★ With lensing **"c"** or without
- ★ $r_{\text{sim}} = 0$ and $r_{\text{sim}} = 0.01$



LiteBIRD simulations with various dust components

LiteBIRD:

- ★ *LiteBIRD noise* (IMO)
- ★ 9 highest frequencies (100-402 GHz)

Mask:

Planck **Mask** + cutoff + apodisation

Such that $f_{\text{sky}} = 0.7$

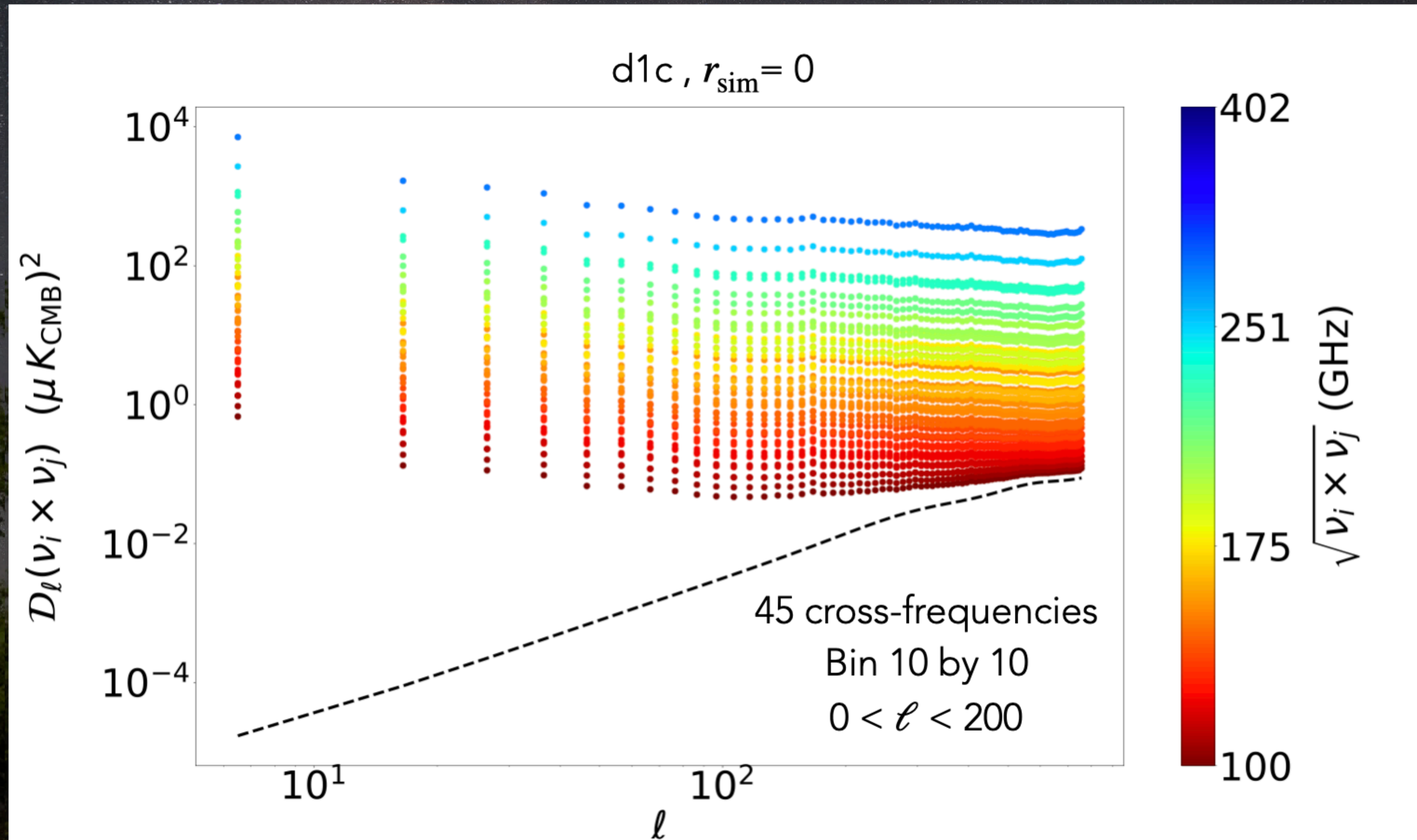
Map:

N_{side} = 256

500 simulations of each type



Estimation of the B-mode cross-frequency spectra (Namaster)



Best fit implementation

χ^2 minimisation using `mpfit` or `emcee`

$$\mathcal{D}_\ell^{\text{model}}(\nu_i \times \nu_j) = \mathcal{D}_\ell^{\text{dust}} \left(\beta(\ell), T_0(\ell), \mathcal{D}_\ell^{\text{ab}}(\nu_i \times \nu_j) \right) + \kappa \times \mathcal{D}_\ell^{\text{lensing}} + r \times \mathcal{D}_\ell^{\text{tensor}}$$

- ★ **Moment expansion** pushed at various order from 0 (MBB) to 3
- ★ $\beta(\ell)$ free at order zero and corrected through iterative process
- ★ $T_0(\ell)$ free at order zero then fixed at best fit value



Best fit implementation

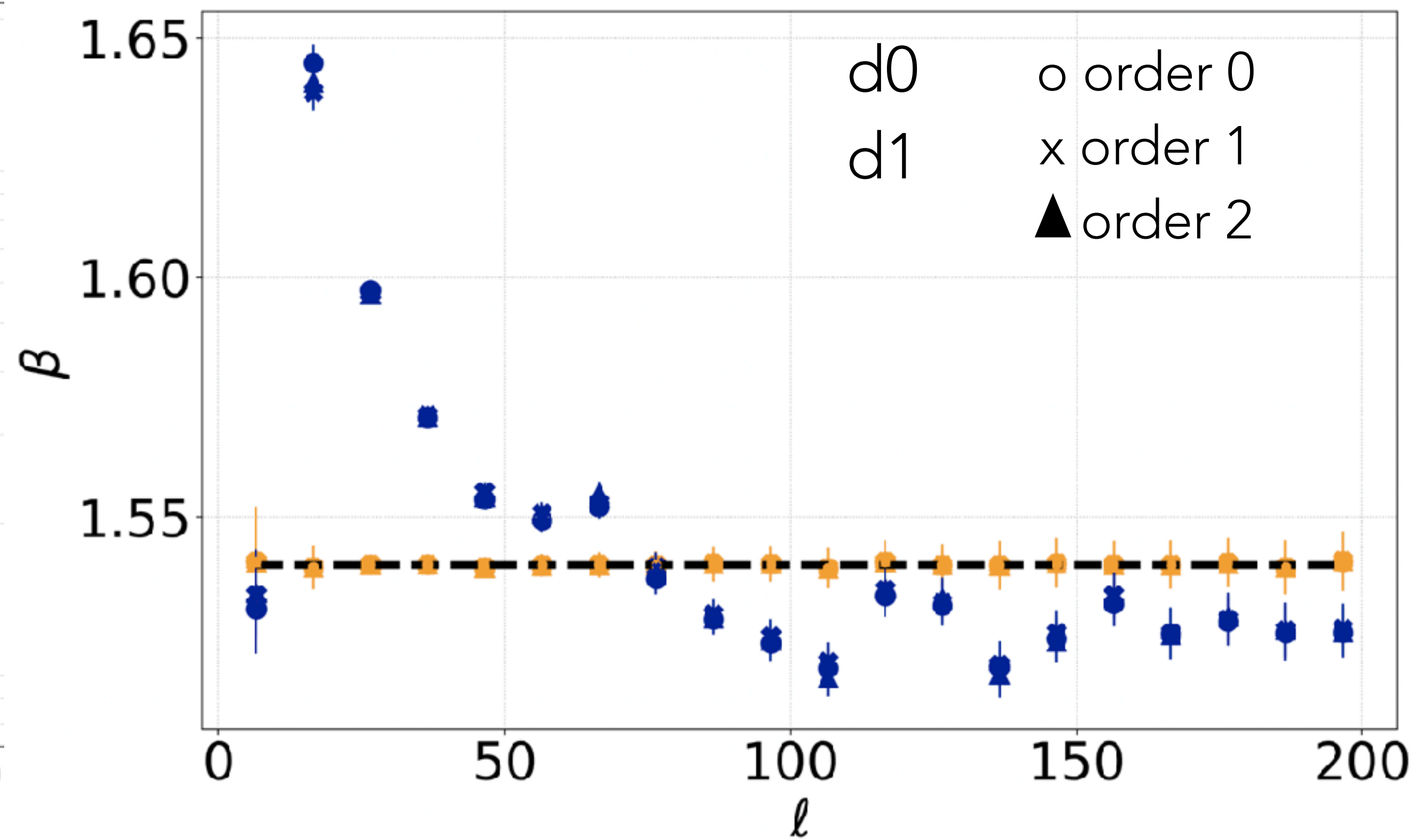
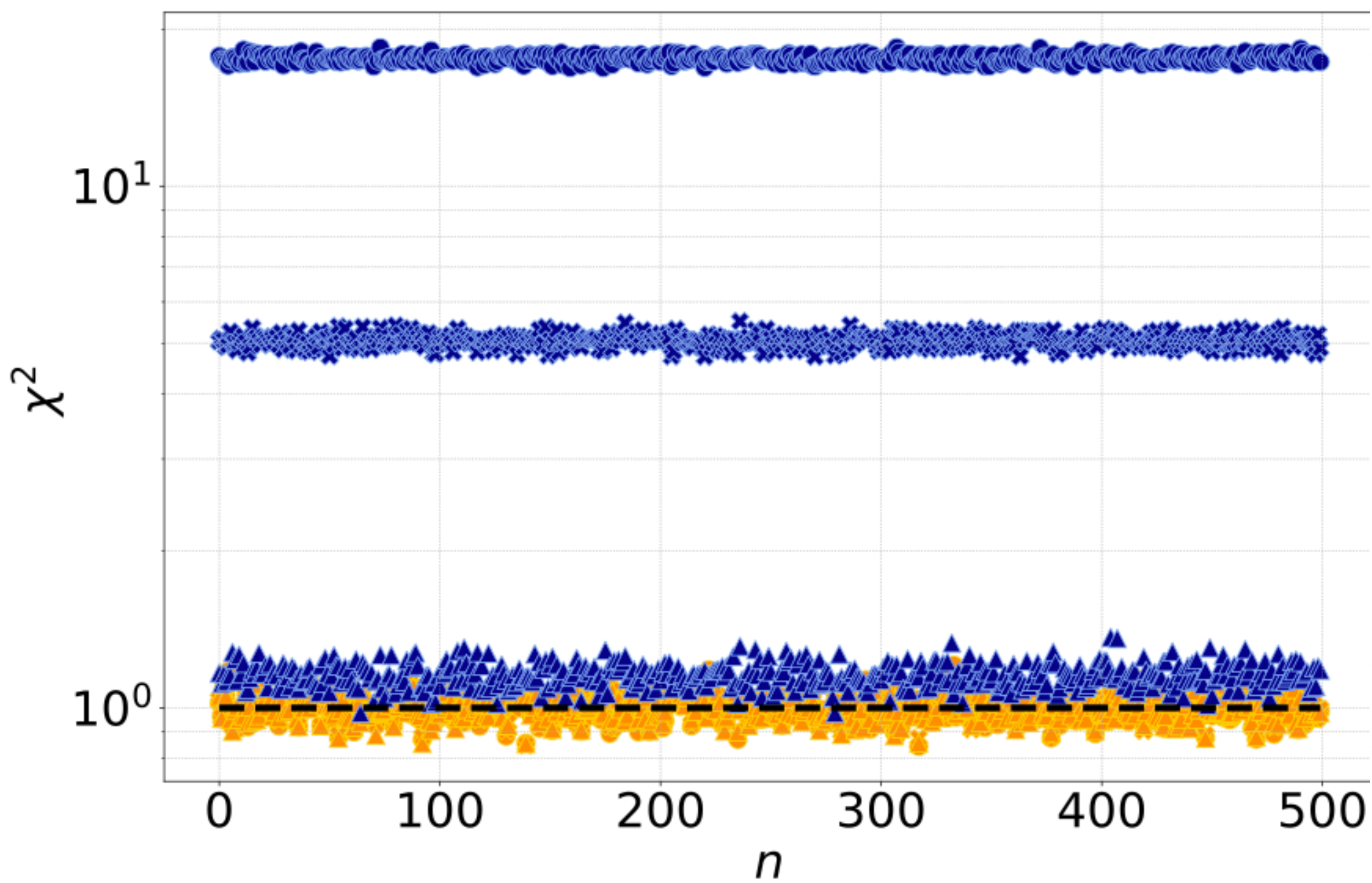
χ^2 minimisation using `mpfit` or `emcee`

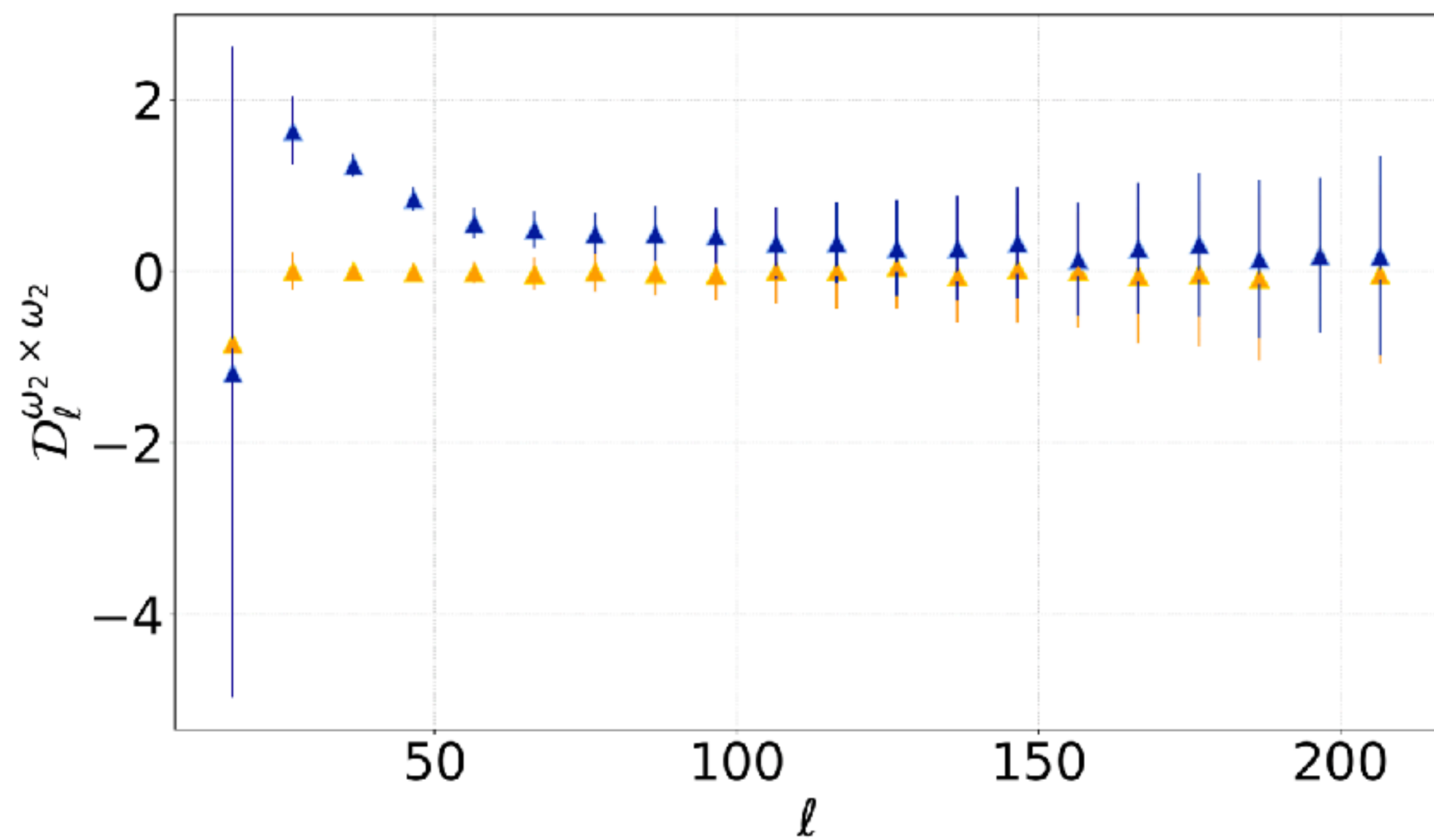
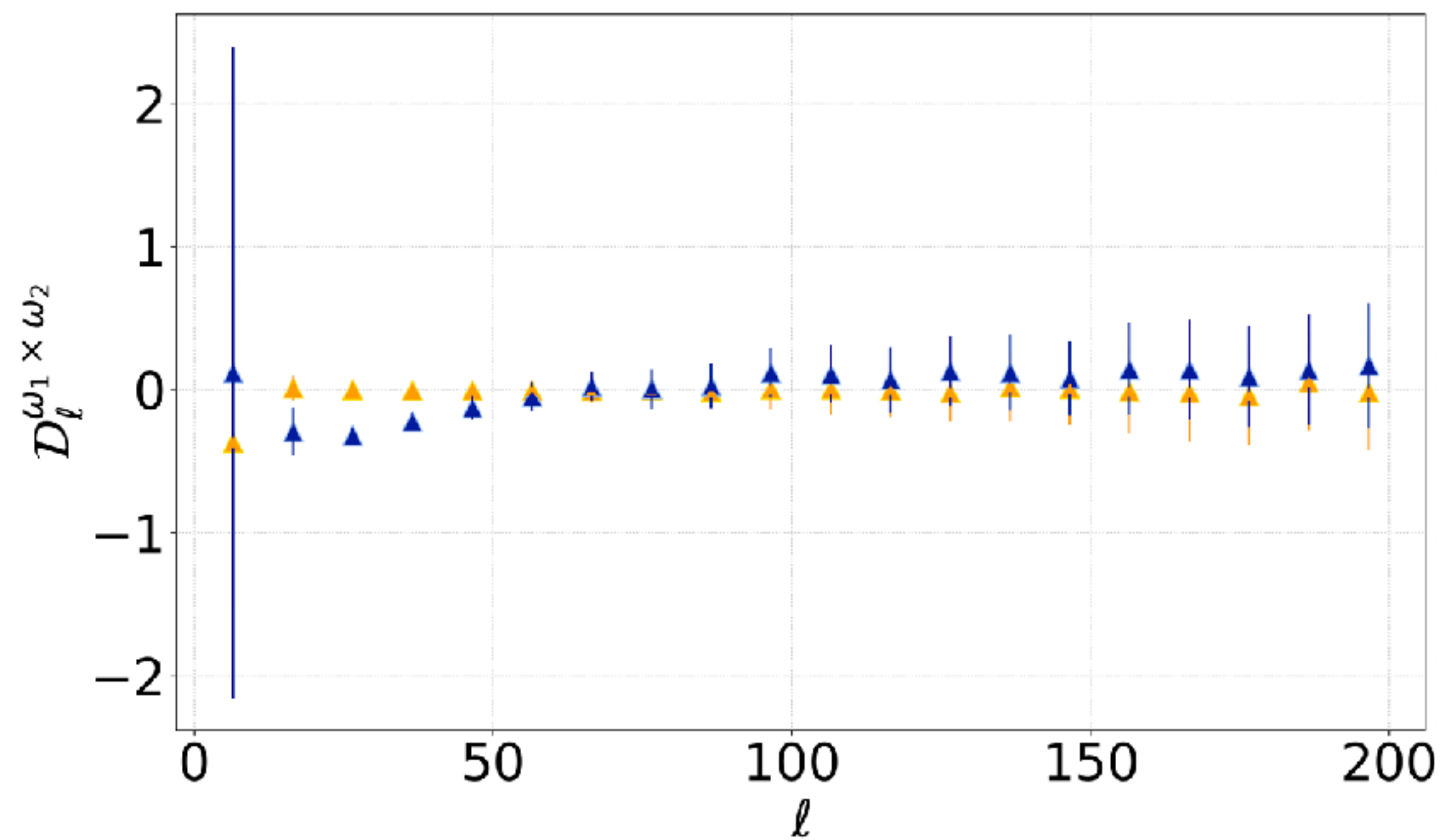
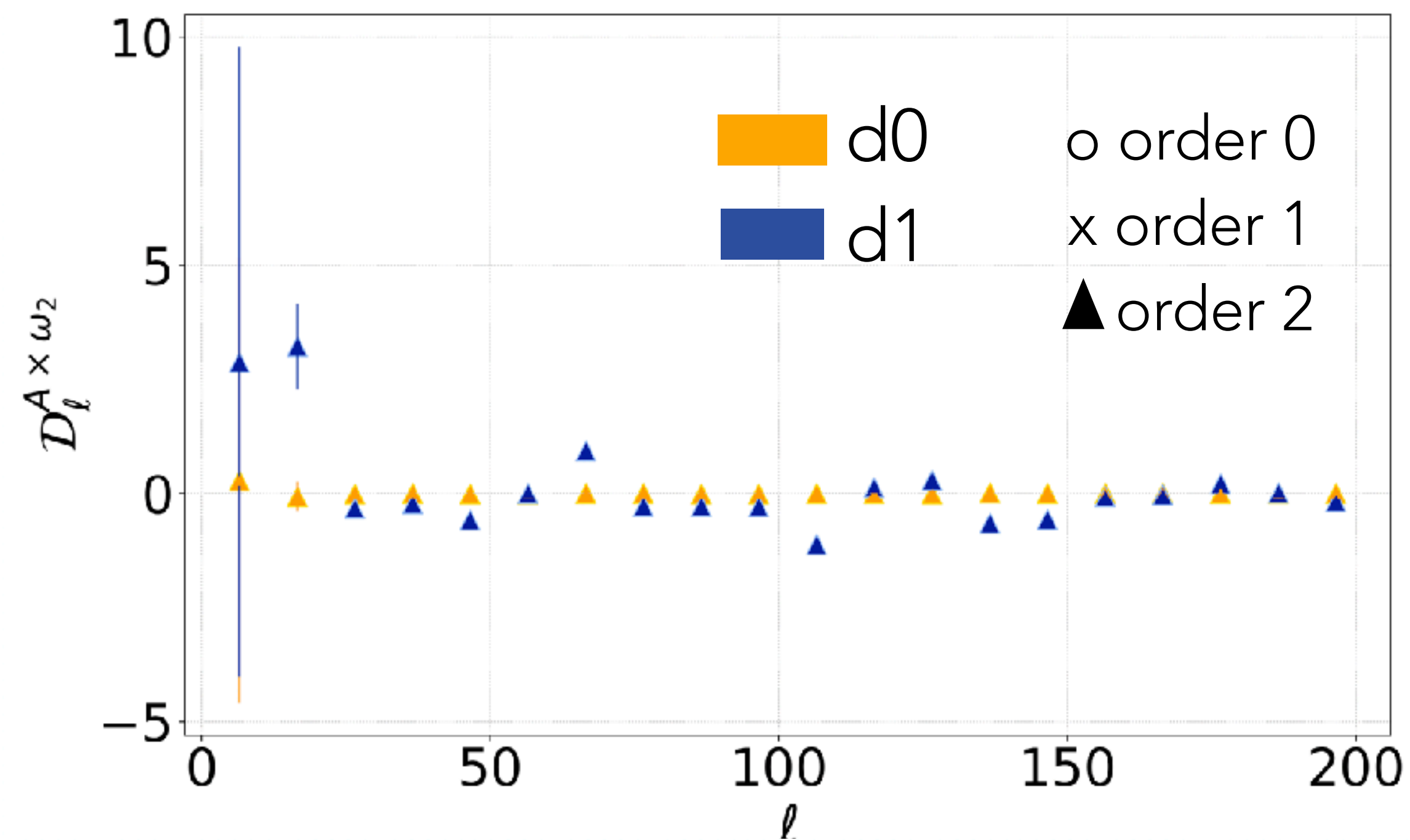
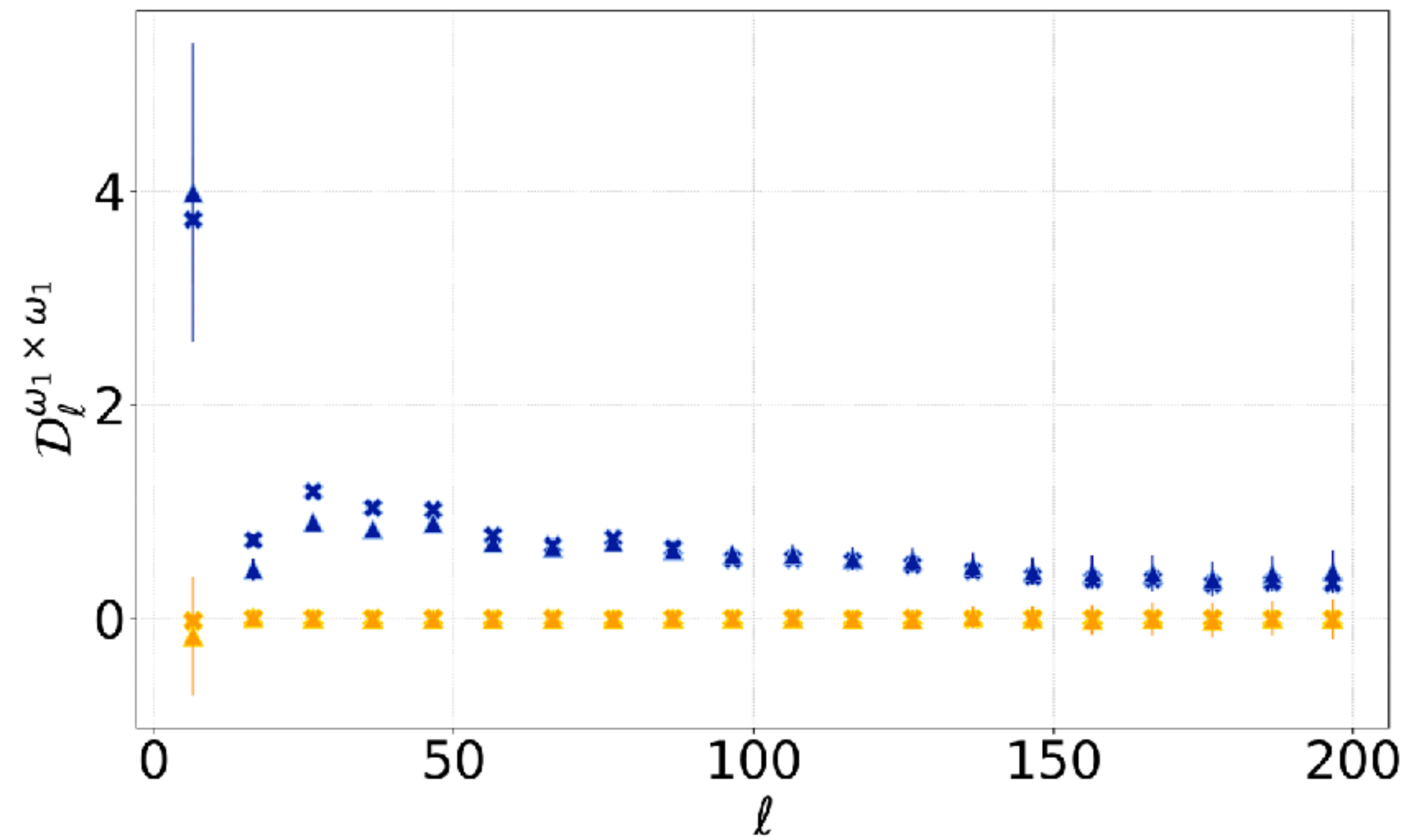
$$\mathcal{D}_\ell^{\text{model}}(\nu_i \times \nu_j) = \mathcal{D}_\ell^{\text{dust}} \left(\beta(\ell), T_0(\ell), \mathcal{D}_\ell^{\text{ab}}(\nu_i \times \nu_j) \right) + \kappa \times \mathcal{D}_\ell^{\text{lensing}} + r \times \mathcal{D}_\ell^{\text{tensor}}$$

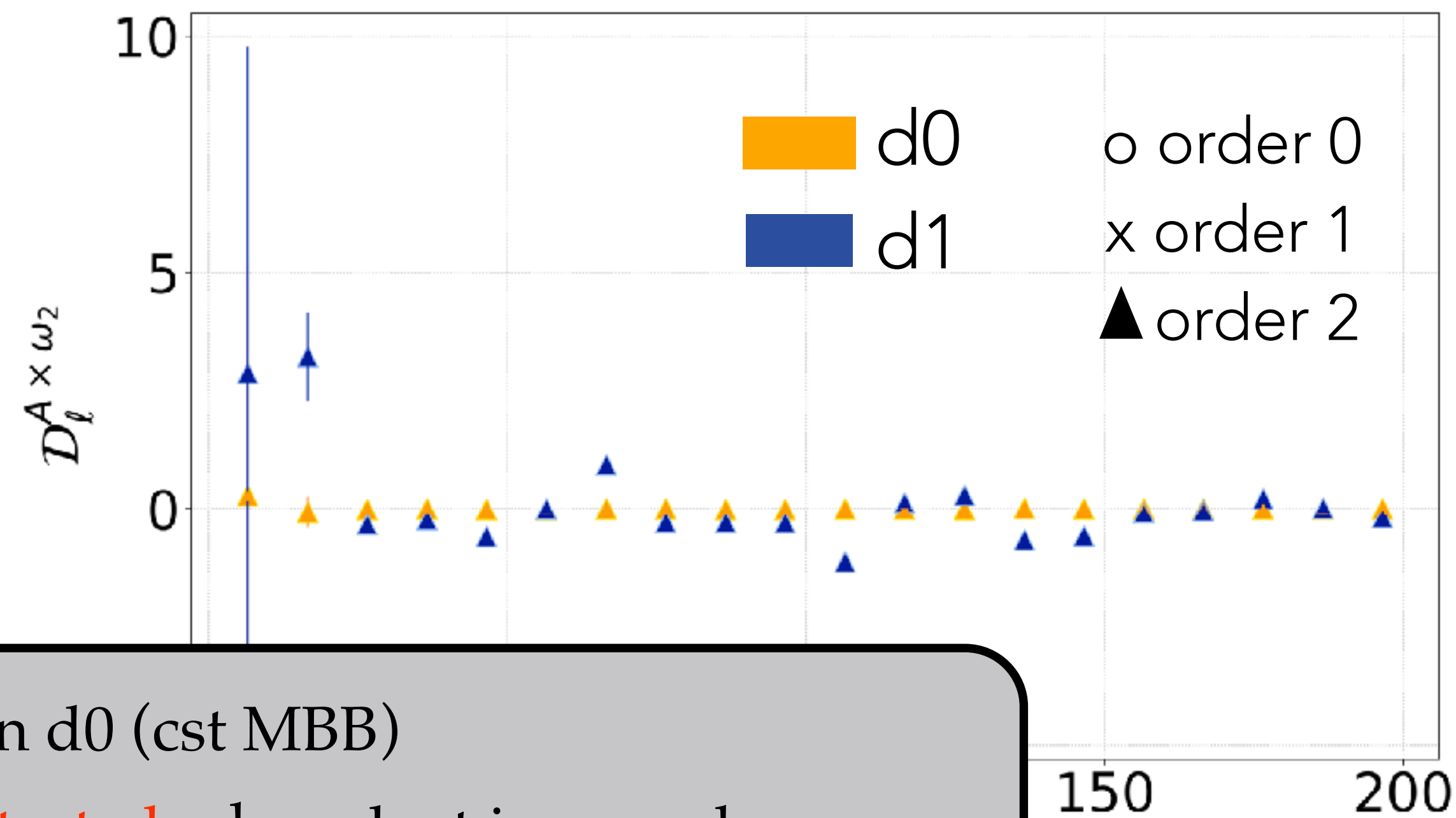
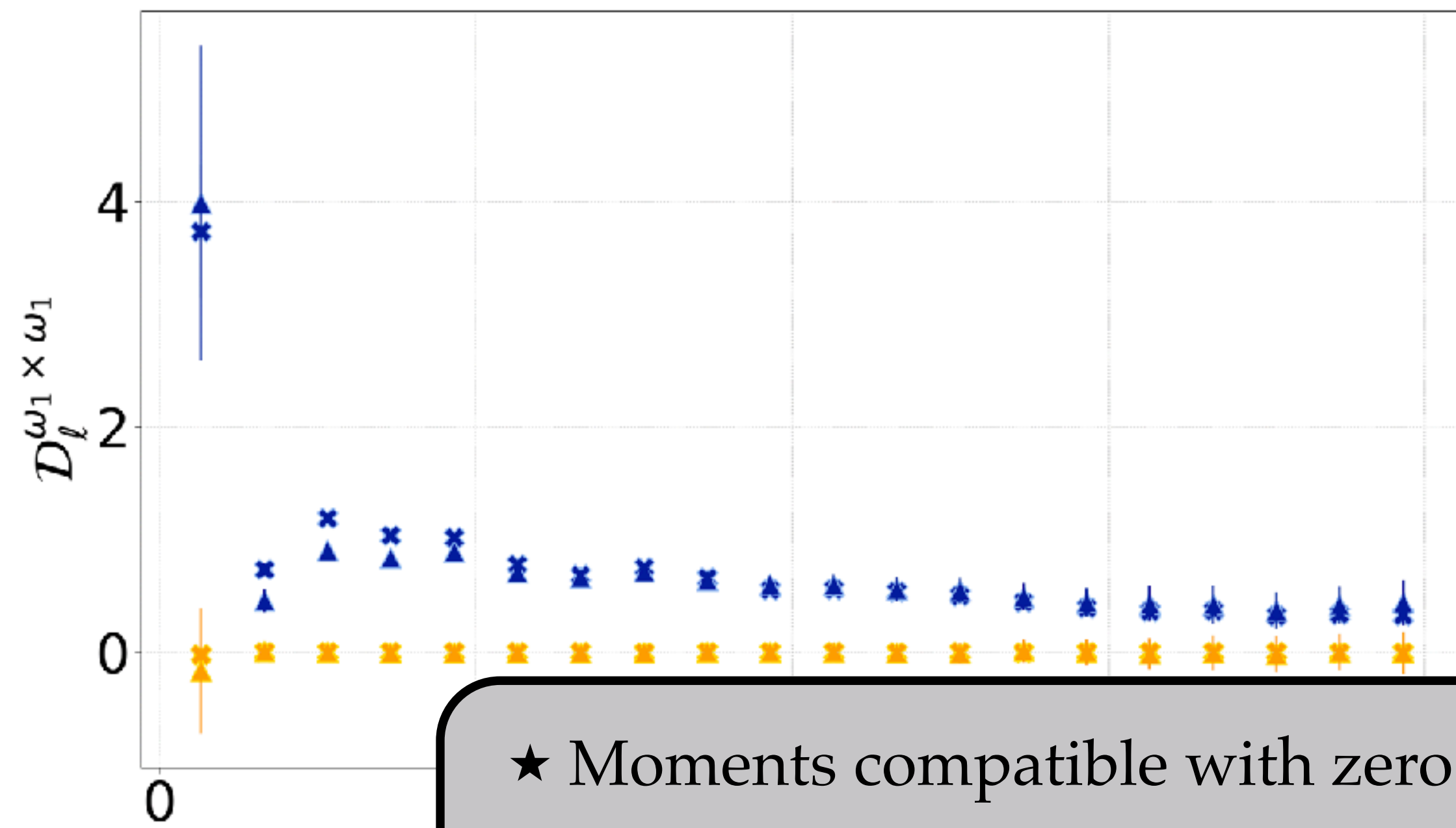
- ★ κ always fixed : 0 for d or 1 for dc
- ★ r fixed or free



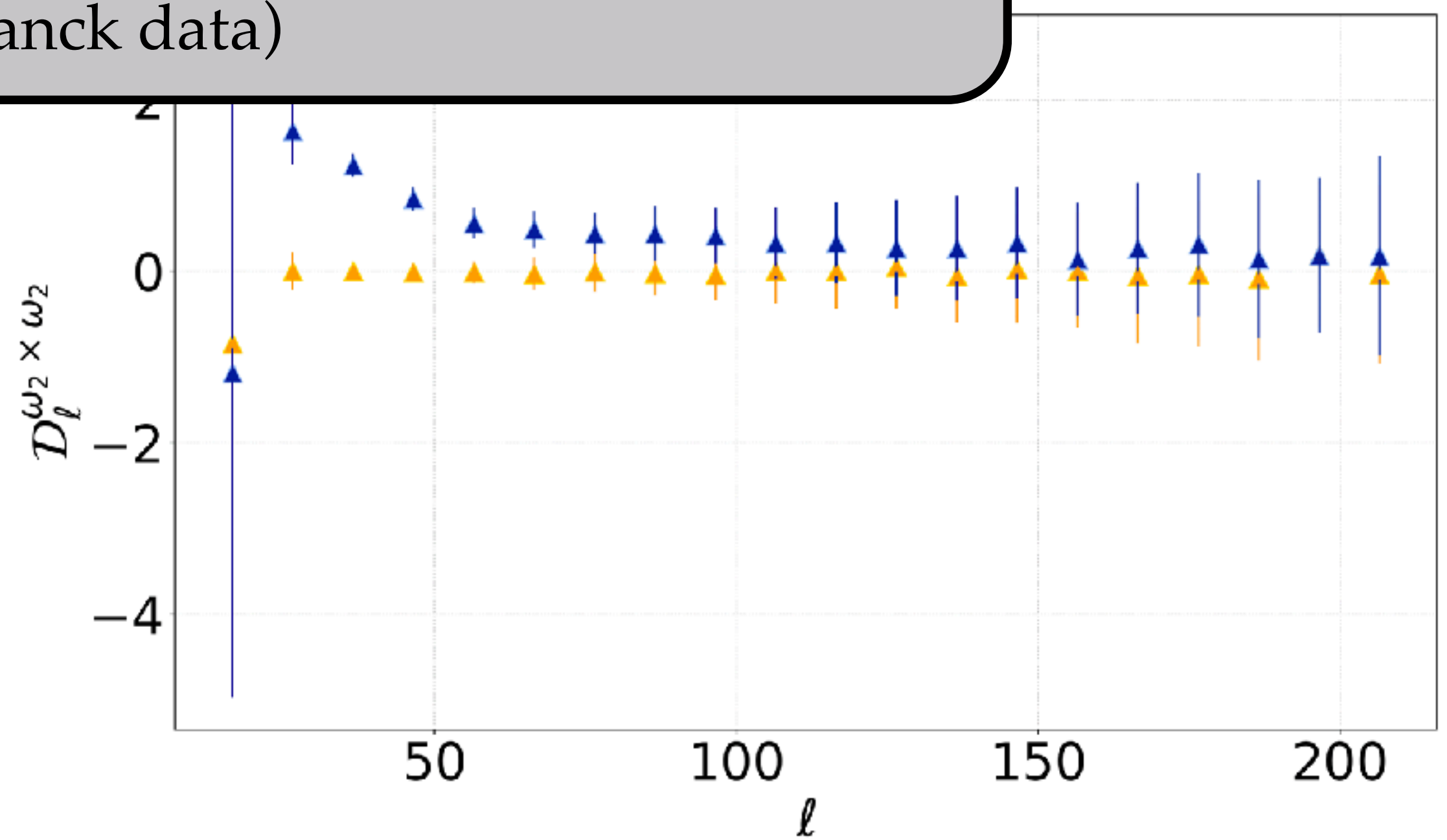
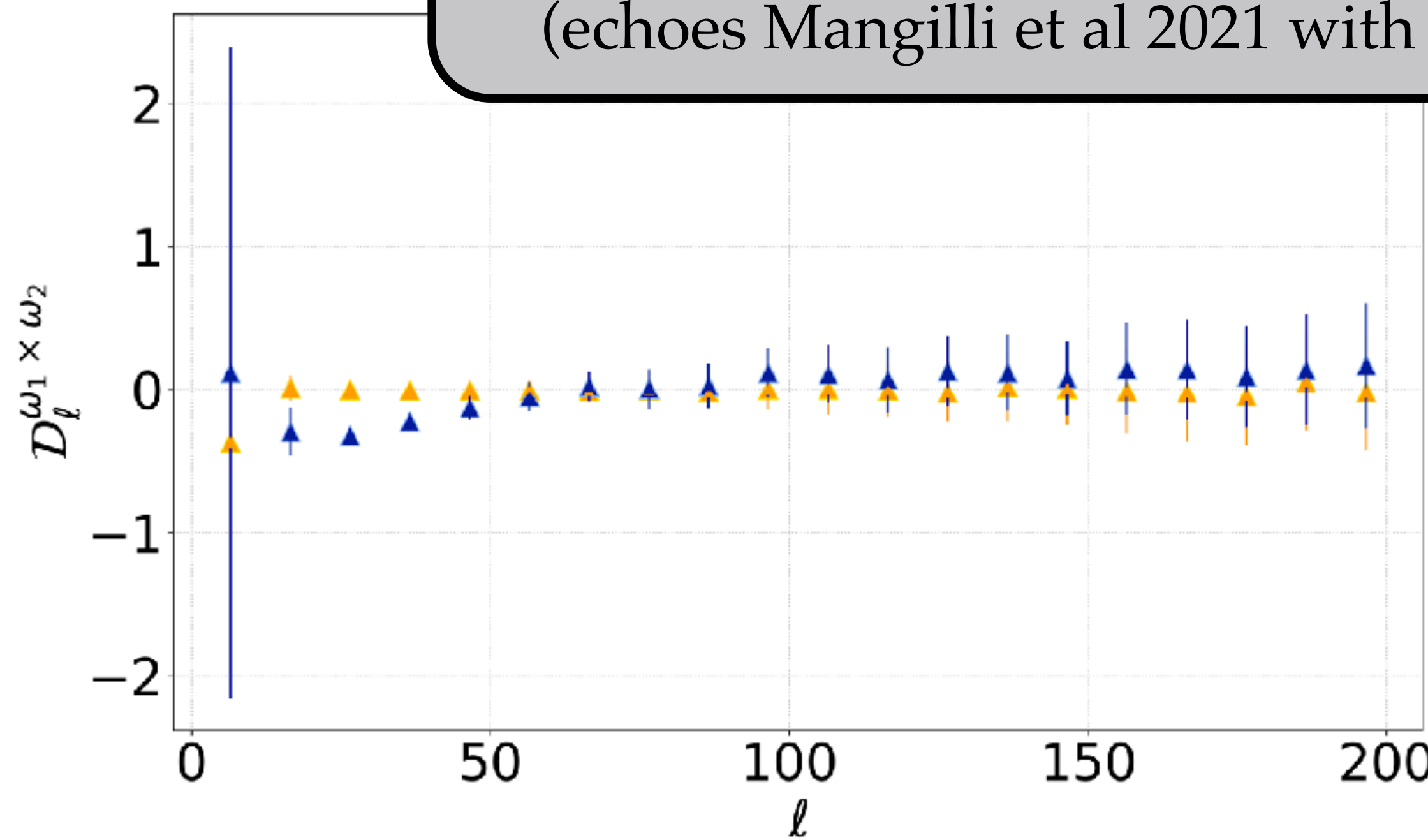
Results I : dust only ($\kappa=0, r=0$ fixed)





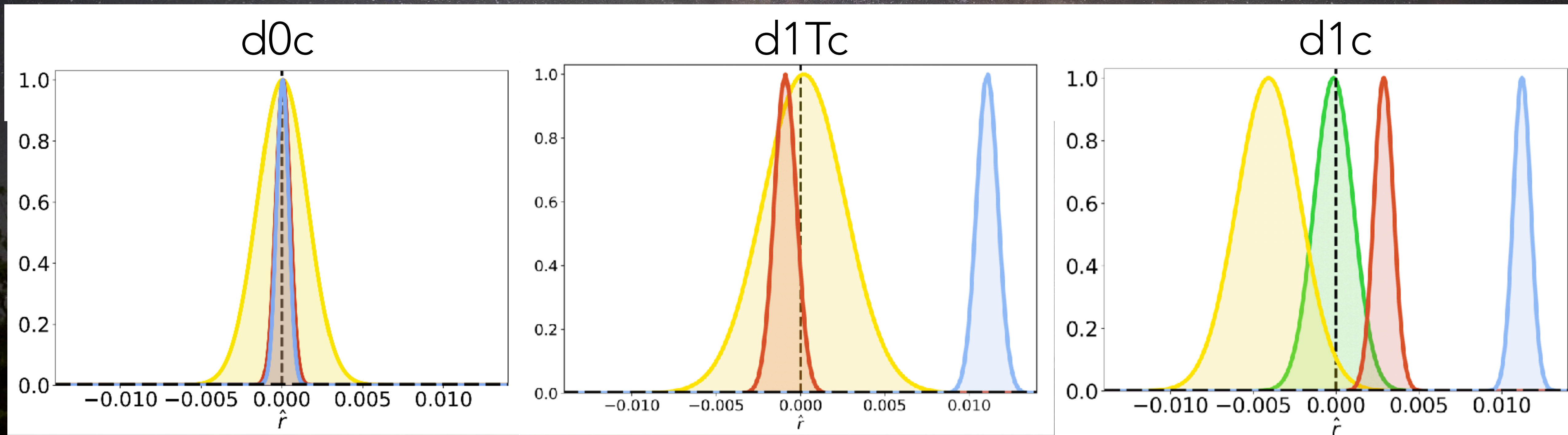


★ Moments compatible with zero in d0 (cst MBB)
 ★ Moments always **significantly detected** when dust in complex (echoes Mangilli et al 2021 with Planck data)



Results II : dust + cmb ($\kappa = 1$, r free, $r_{\text{sim}} = 0$)

■ Order 0 (MBB) ■ Order 1 beta ■ Order 1 beta ad T ■ Order 2 beta



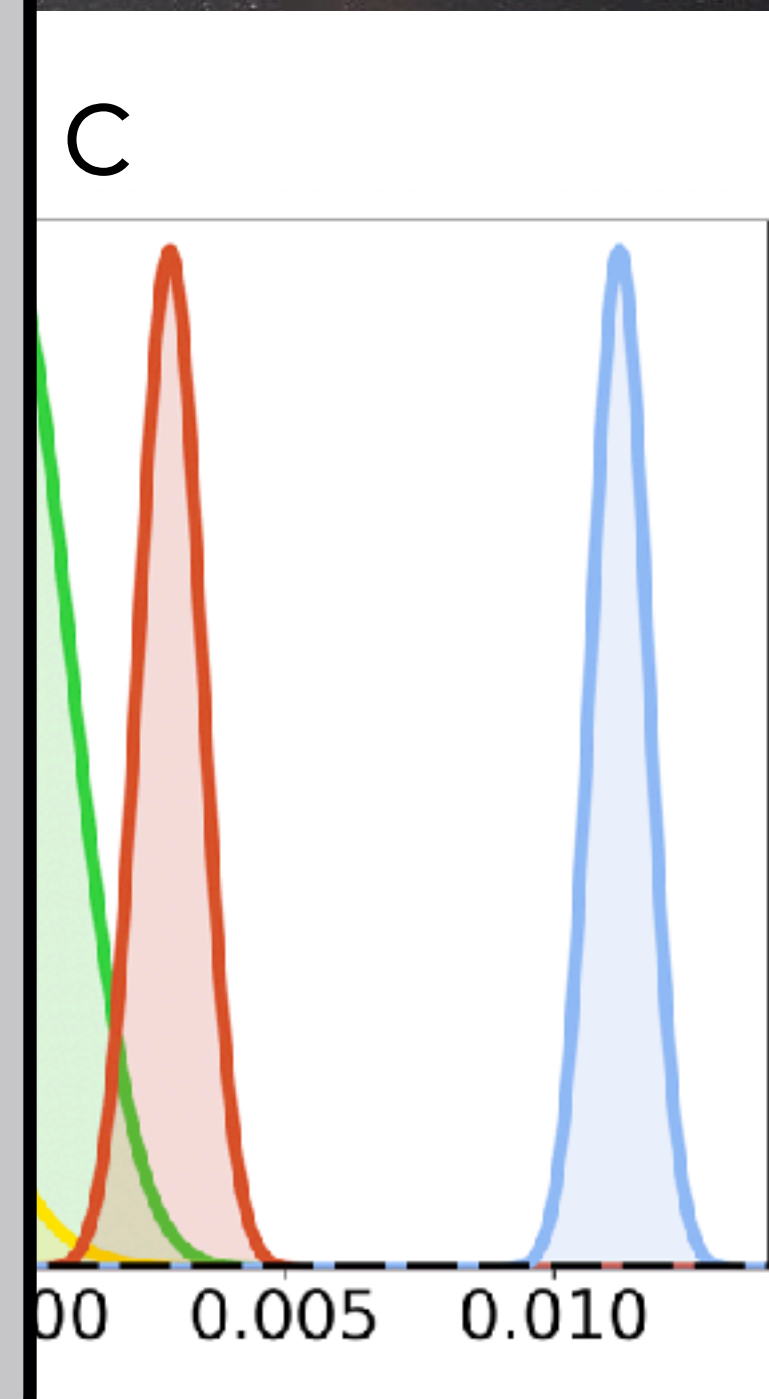
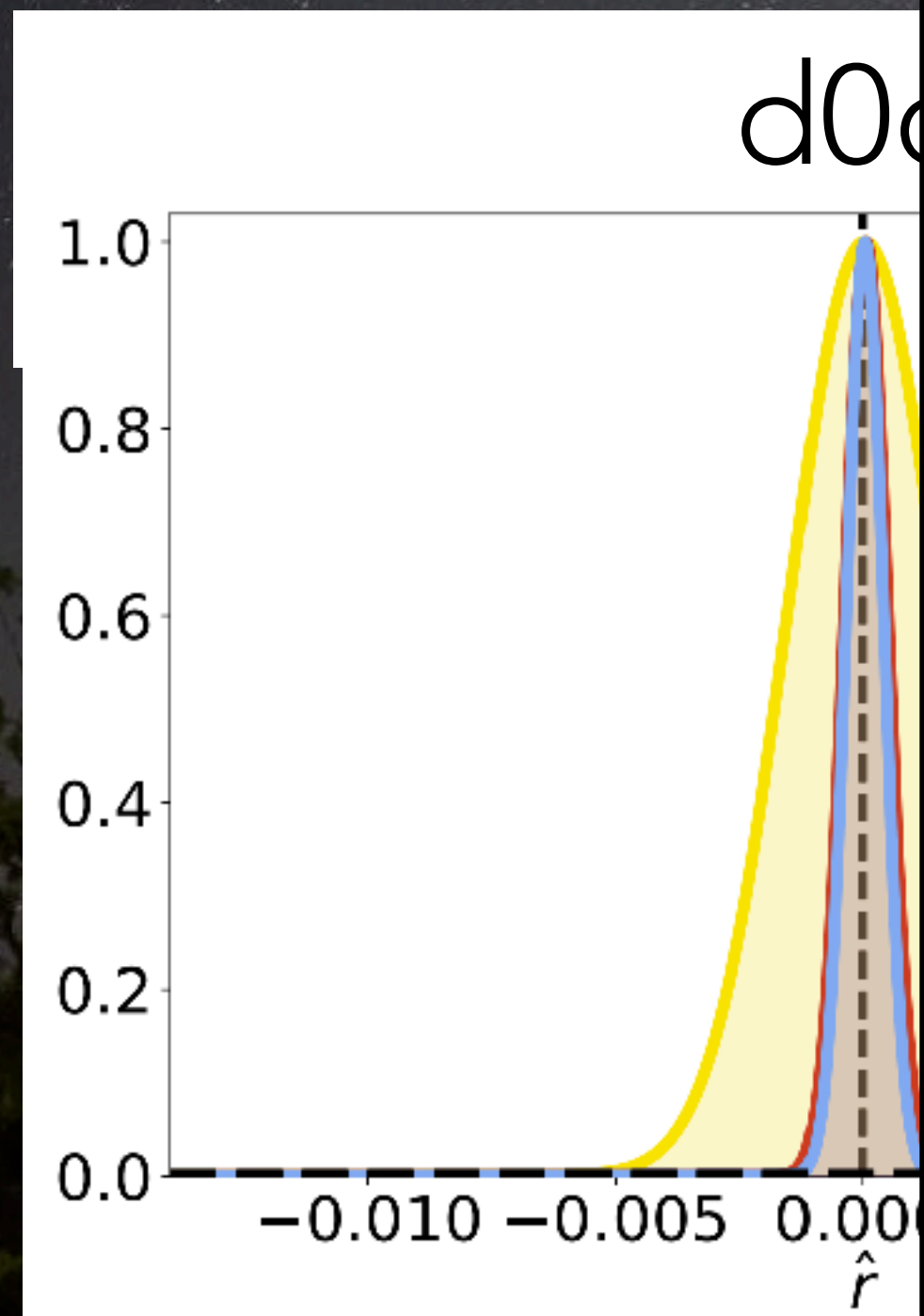
Results II : dust + cmb ($\kappa = 1, r$ free, $r_{\text{sim}} = 0$)

■ Order 0 (MBB)
 ■ Order 1 beta
 ■ Order 1 beta ad T
 ■ Order 2 beta

★ Results robust with $r_{\text{sim}} \neq 0 \rightarrow$ moment expansion can be trust to detect r

- ★ d0 MBB no bias + moment expansions no bias
- ★ d1T MBB biased + moment order 1 and 2 no bias
- ★ d1c : - Order 0 and 1 in beta with dispersion around noise (few 10^{-4})
 - Order 1 beta-T : no biais, dispersion around (1.1×10^{-3})
 - Order 2 double the dispersion + negative bias at order 2

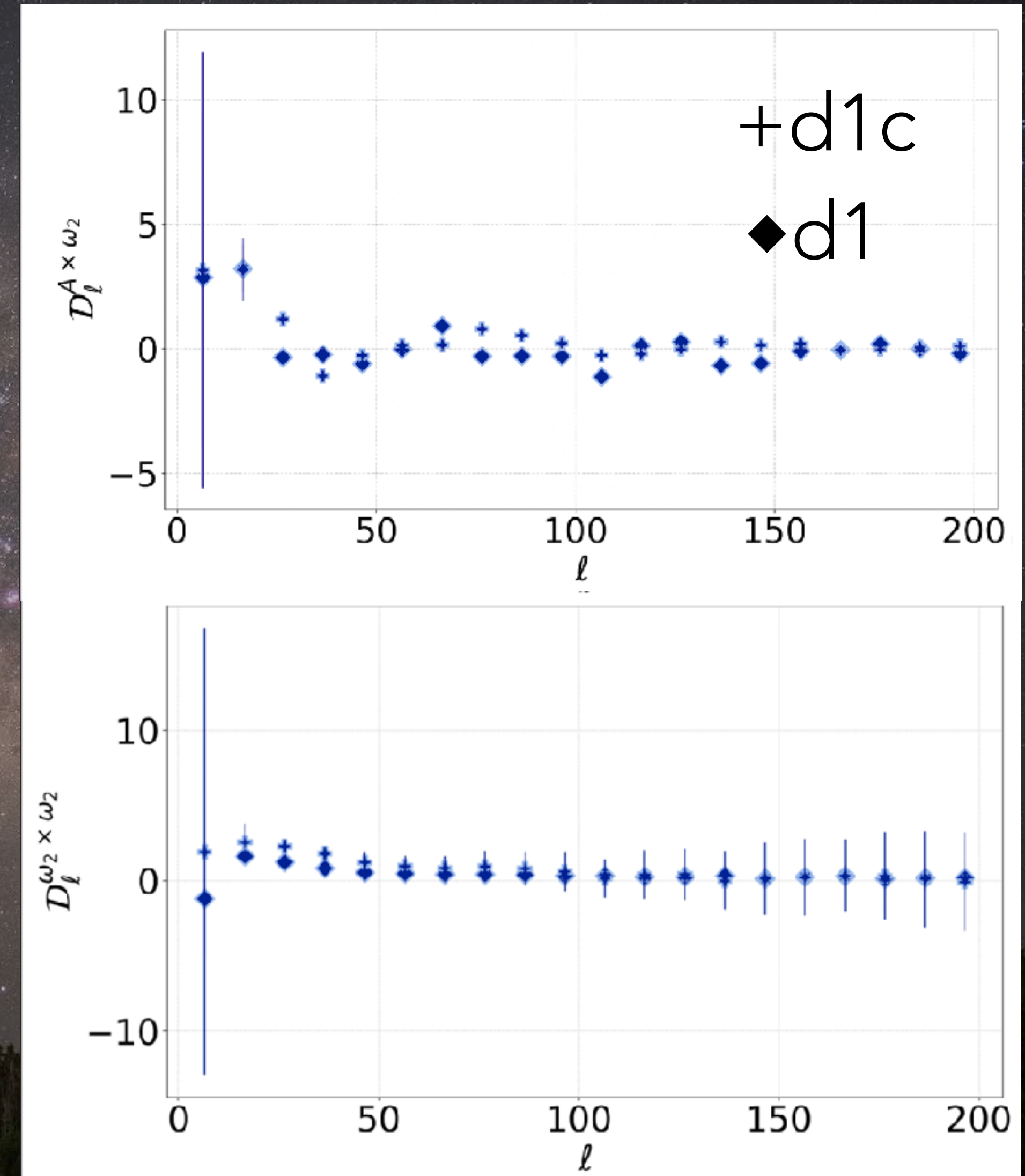
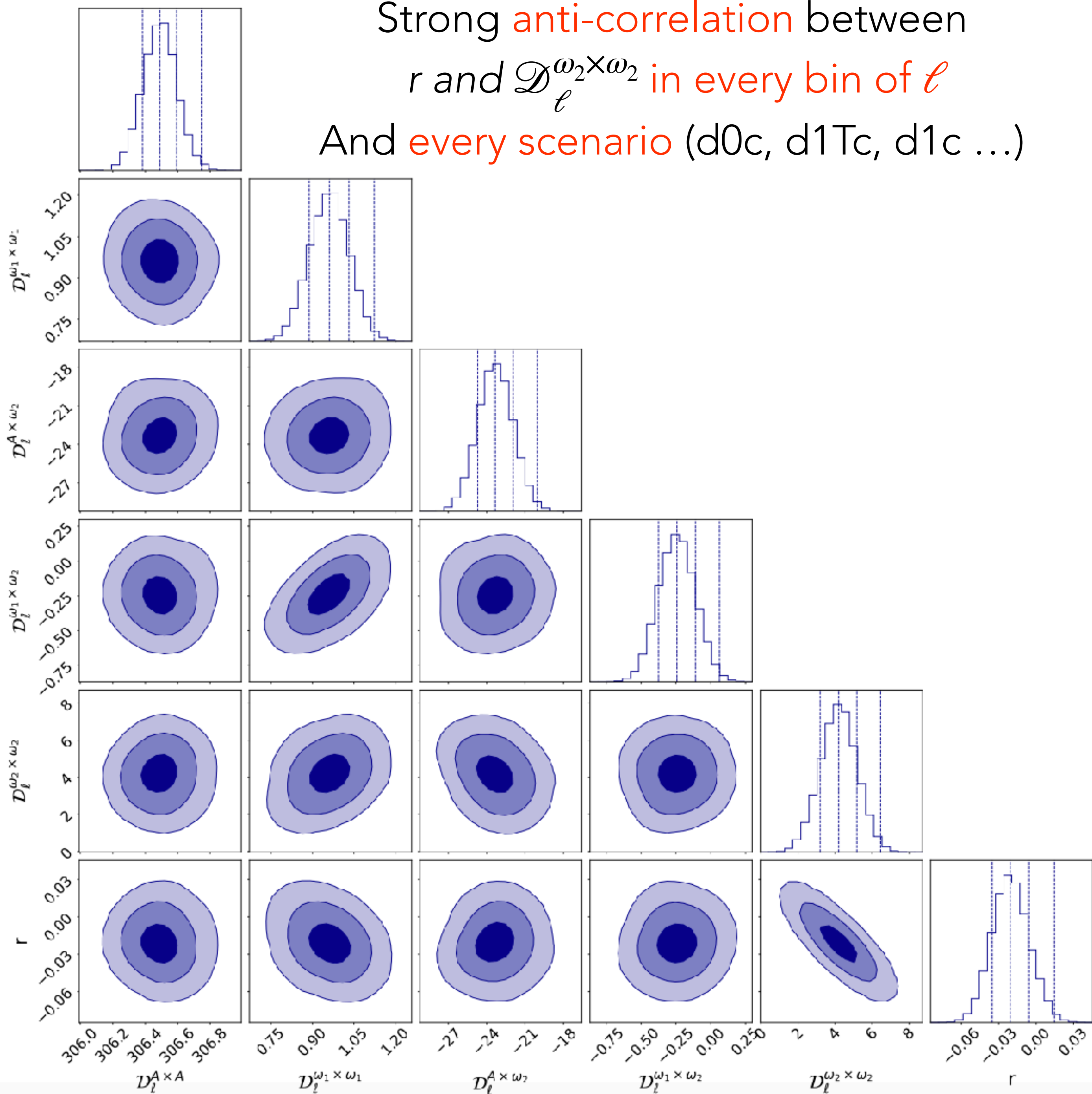
... Why ?



Strong anti-correlation between

r and $D_\ell^{\omega_2 \times \omega_2}$ in every bin of ℓ

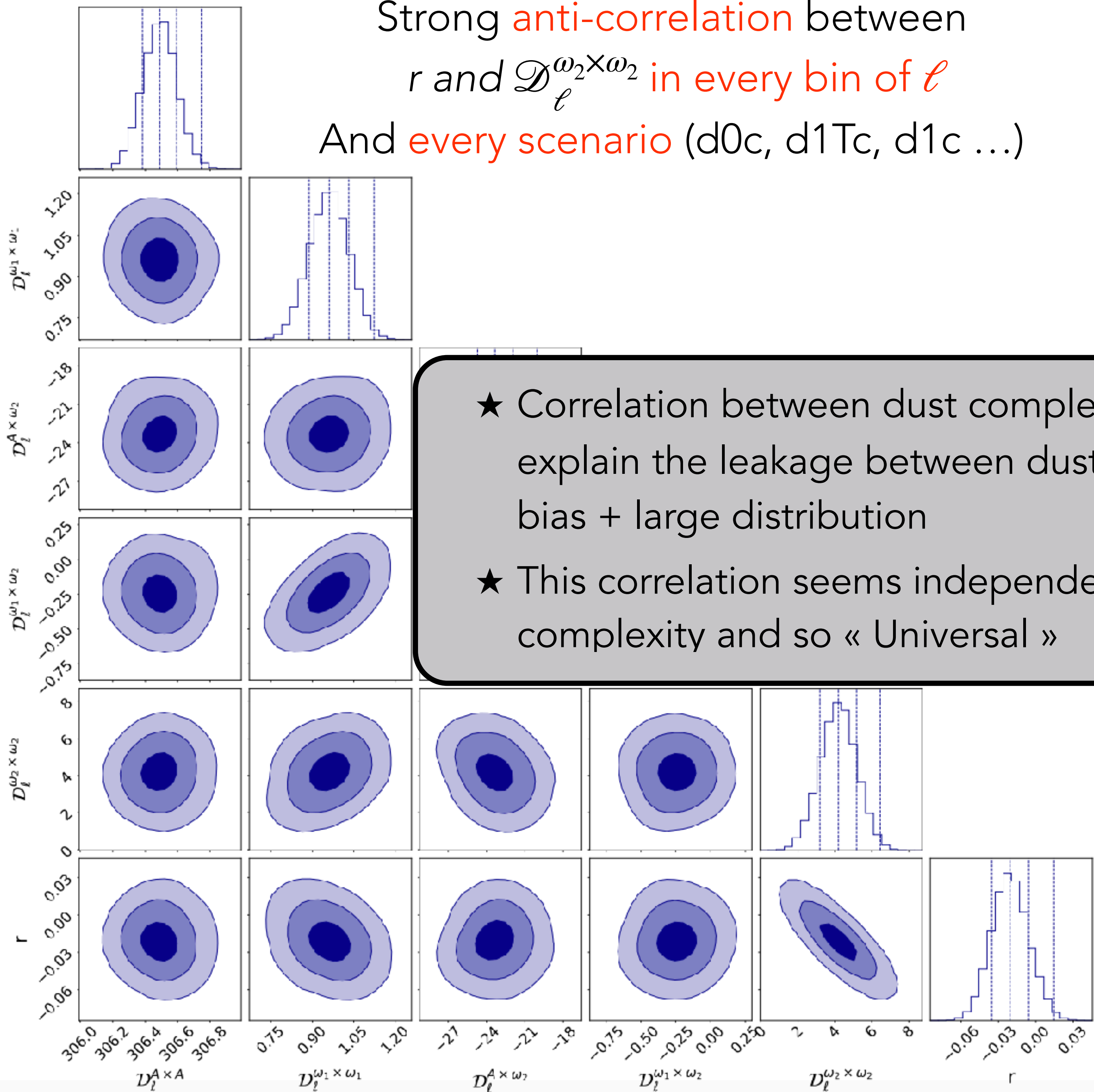
And every scenario (d0c, d1Tc, d1c ...)



Strong **anti-correlation** between

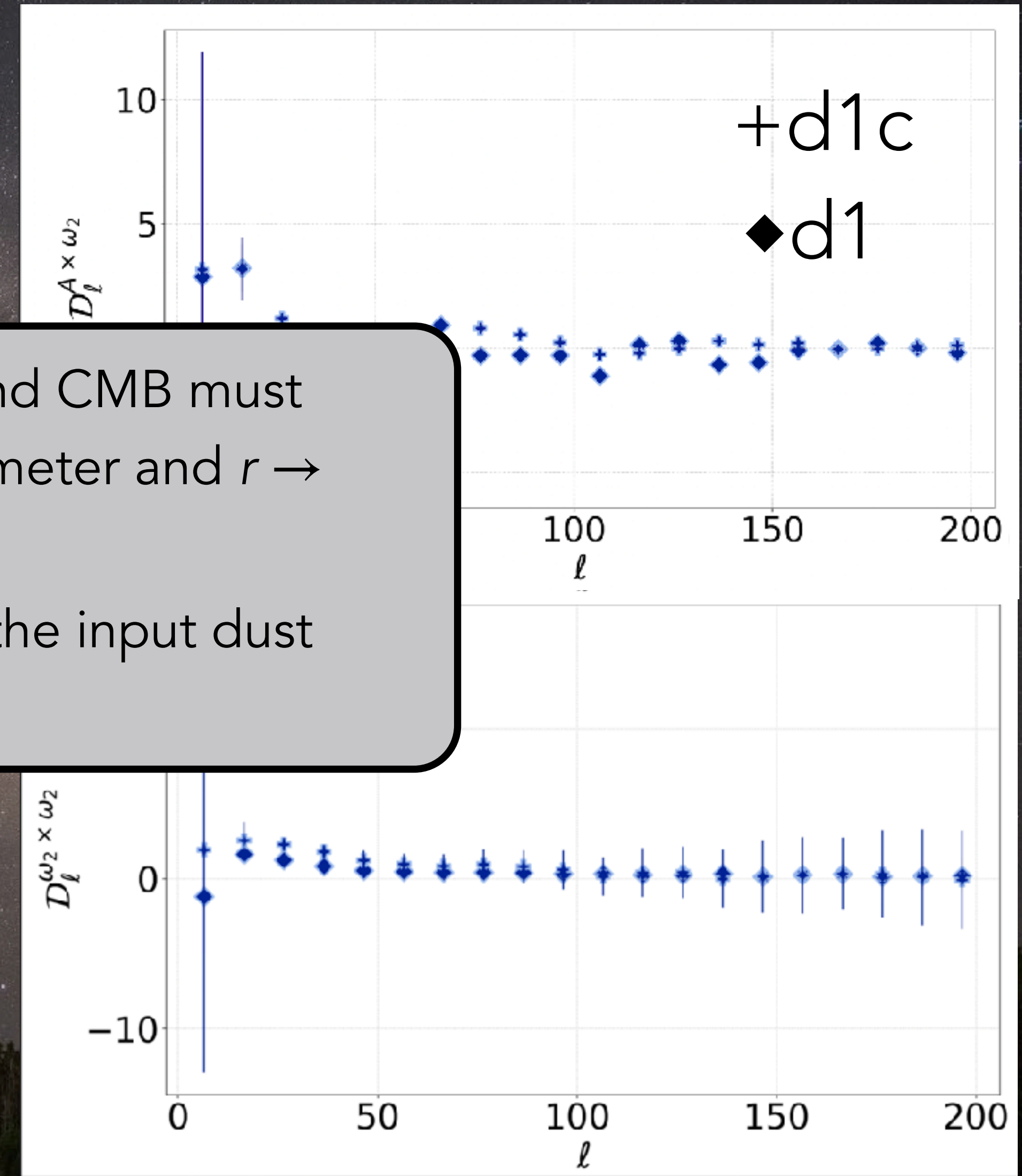
r and $D_\ell^{\omega_2 \times \omega_2}$ in every bin of ℓ

And **every scenario** (d0c, d1Tc, d1c ...)



★ Correlation between dust complexity and CMB must explain the leakage between dust parameter and $r \rightarrow$ bias + large distribution

★ This correlation seems independent of the input dust complexity and so « Universal »



What's next :

- ★ Update the **expansion around T**
- ★ Finalise our first paper ! (*Vacher et al* to come)
- ★ Introduce a formalism more **friendly with polarisation**
- ★ Consider also **synchrotron** and all frequency bands
- ★ Application to other instruments (SO (Azzoni et al 2021), QUBIC ongoing)
- ★ Take into account for **frequency variations** of SED parameters
 - + **non gaussian** dust emission
- ★ Couple moments with **other methods** ...
(already explored in pixel space (*Remazeilles et al 2017*))



Take away :

- ★ $\beta(\ell)$ and $T(\ell)$ can be fitted together with a MBB without biais/error-bar explosion
- ★ **Temperature** could be a critical parameter for LB
- ★ Using moment expansion at order 1, one can **reduce/absorb** the **bias on r**
- ★ d1 contains order 2 terms in power spectra level ...
- ★ **Degeneracies** between dust complexity and tensor to scalar ratio ! Prevent to take into account for the additional complexity properly

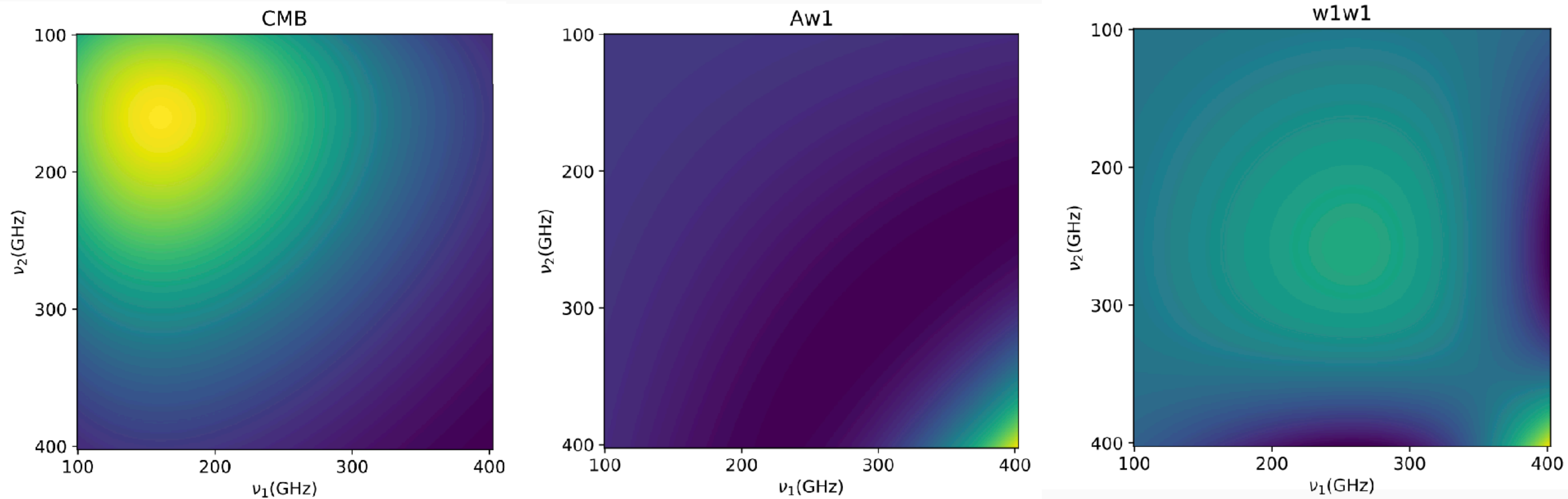
Could be problematic for component separation



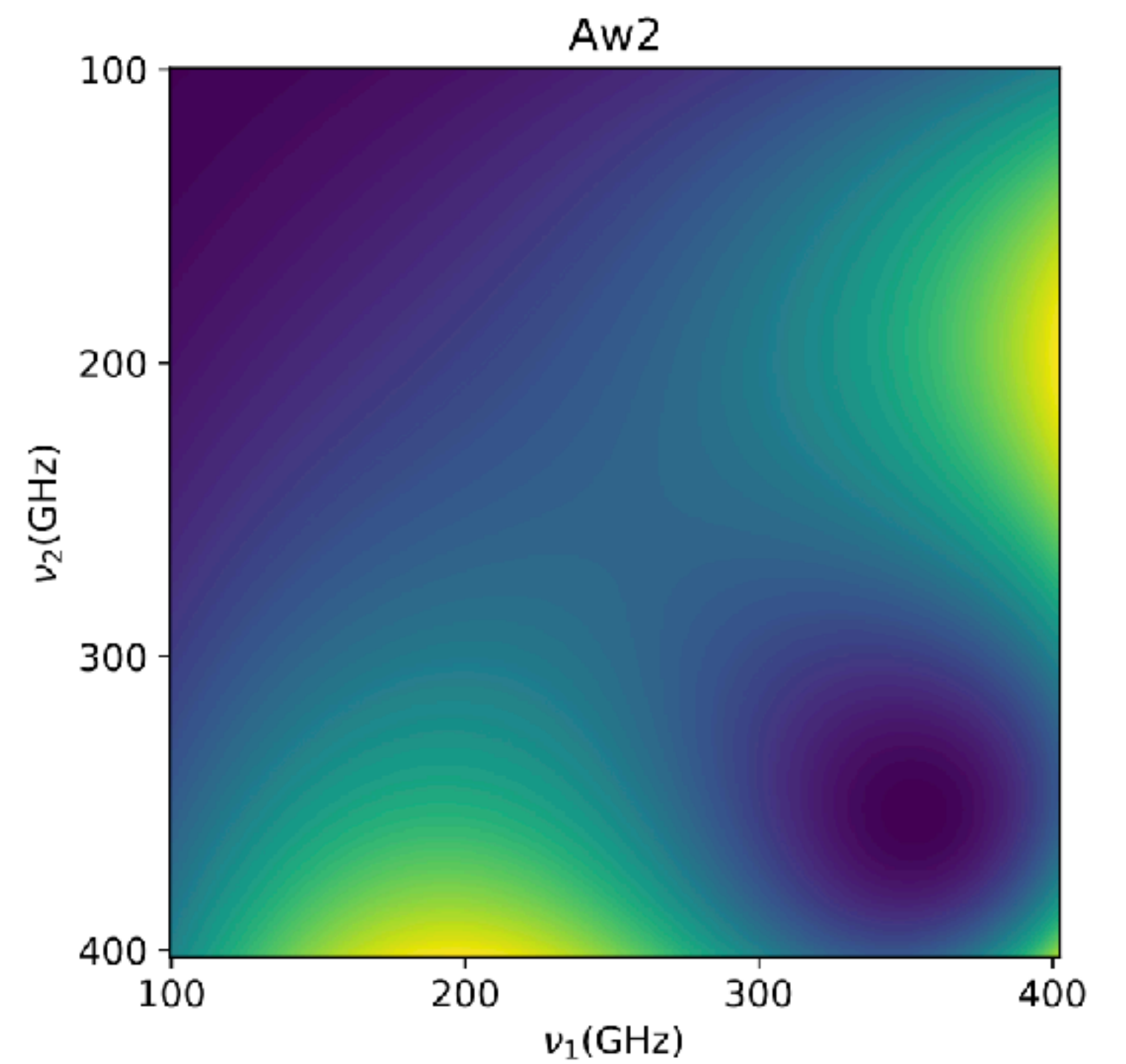
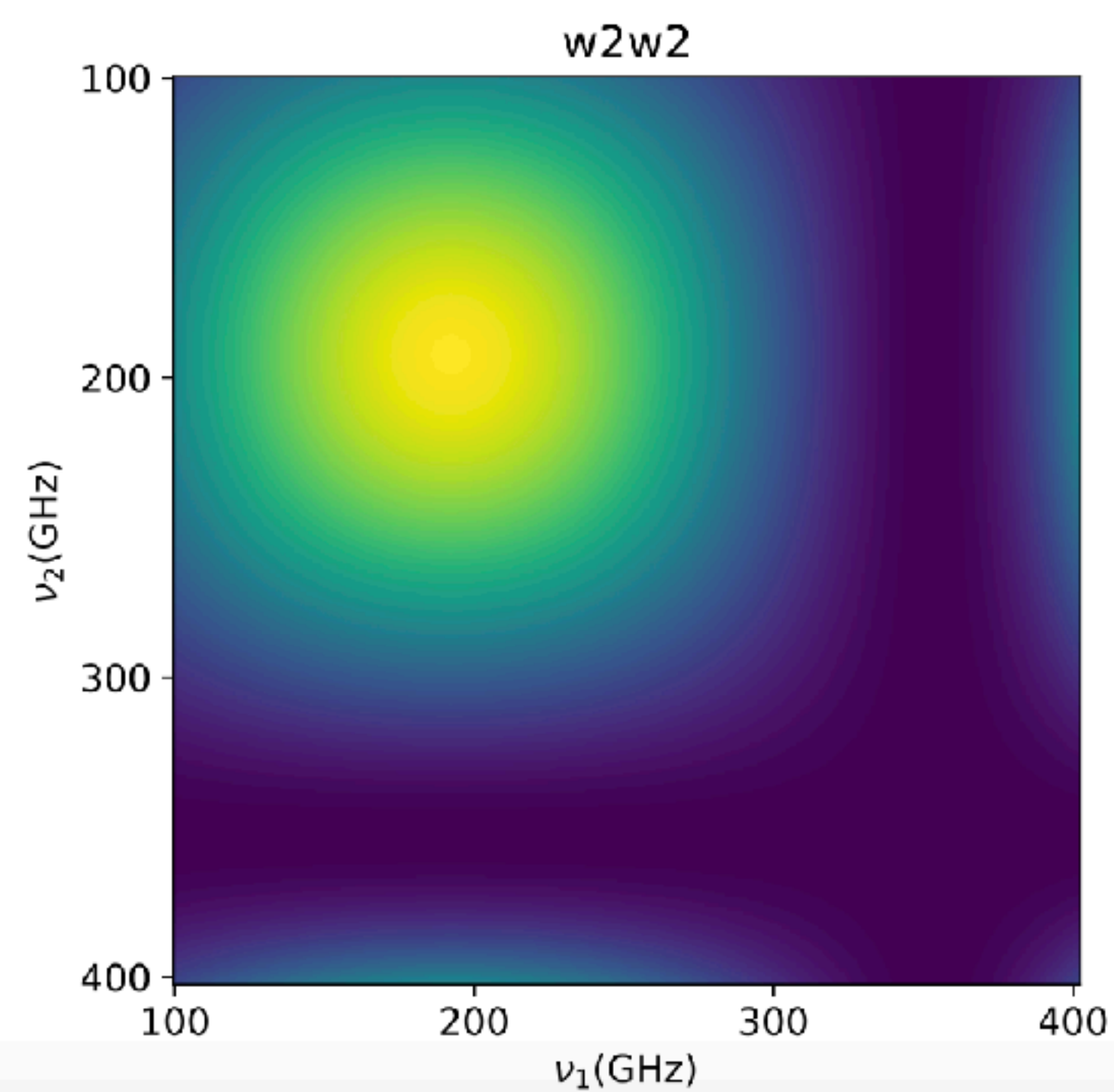
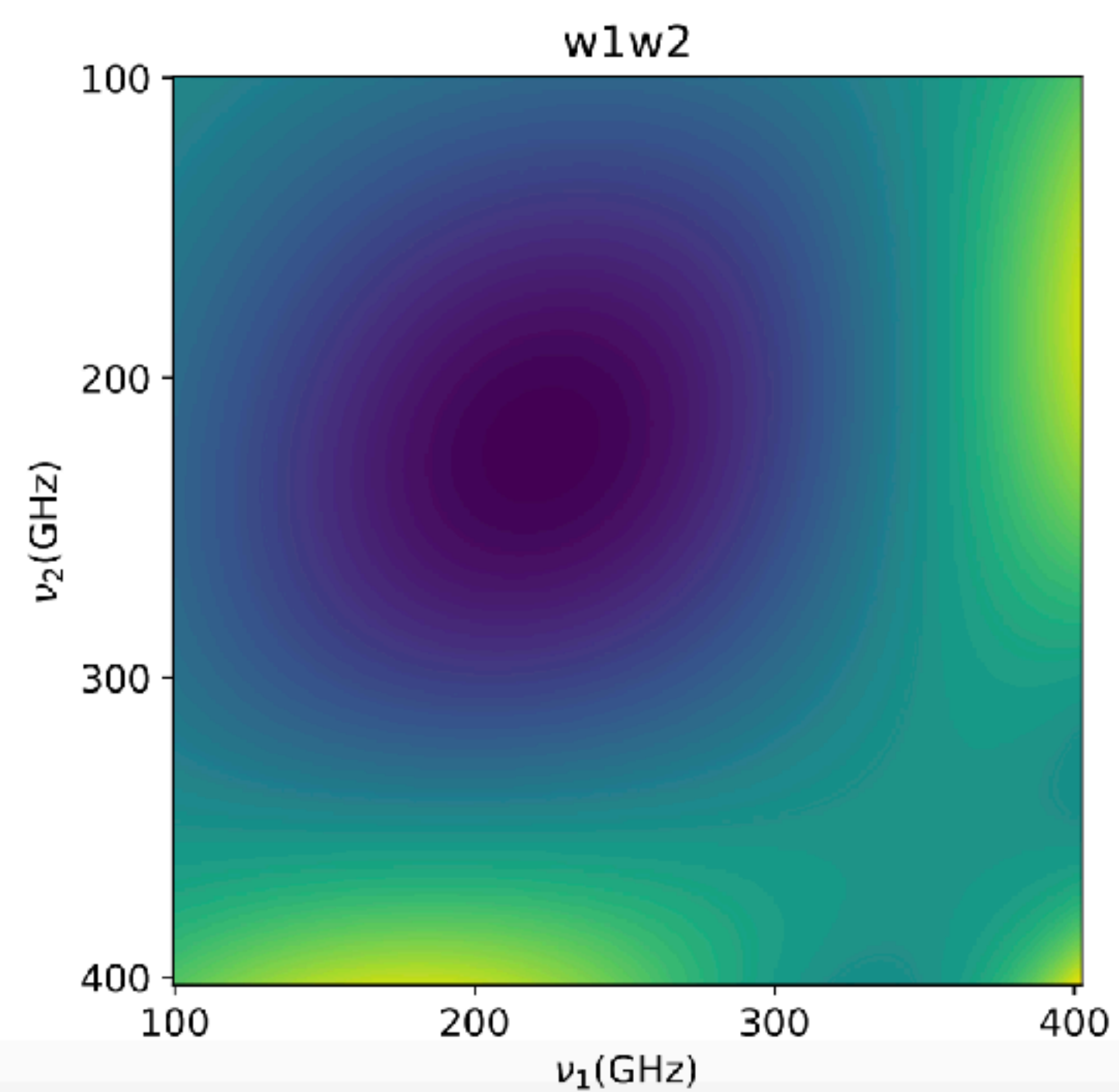
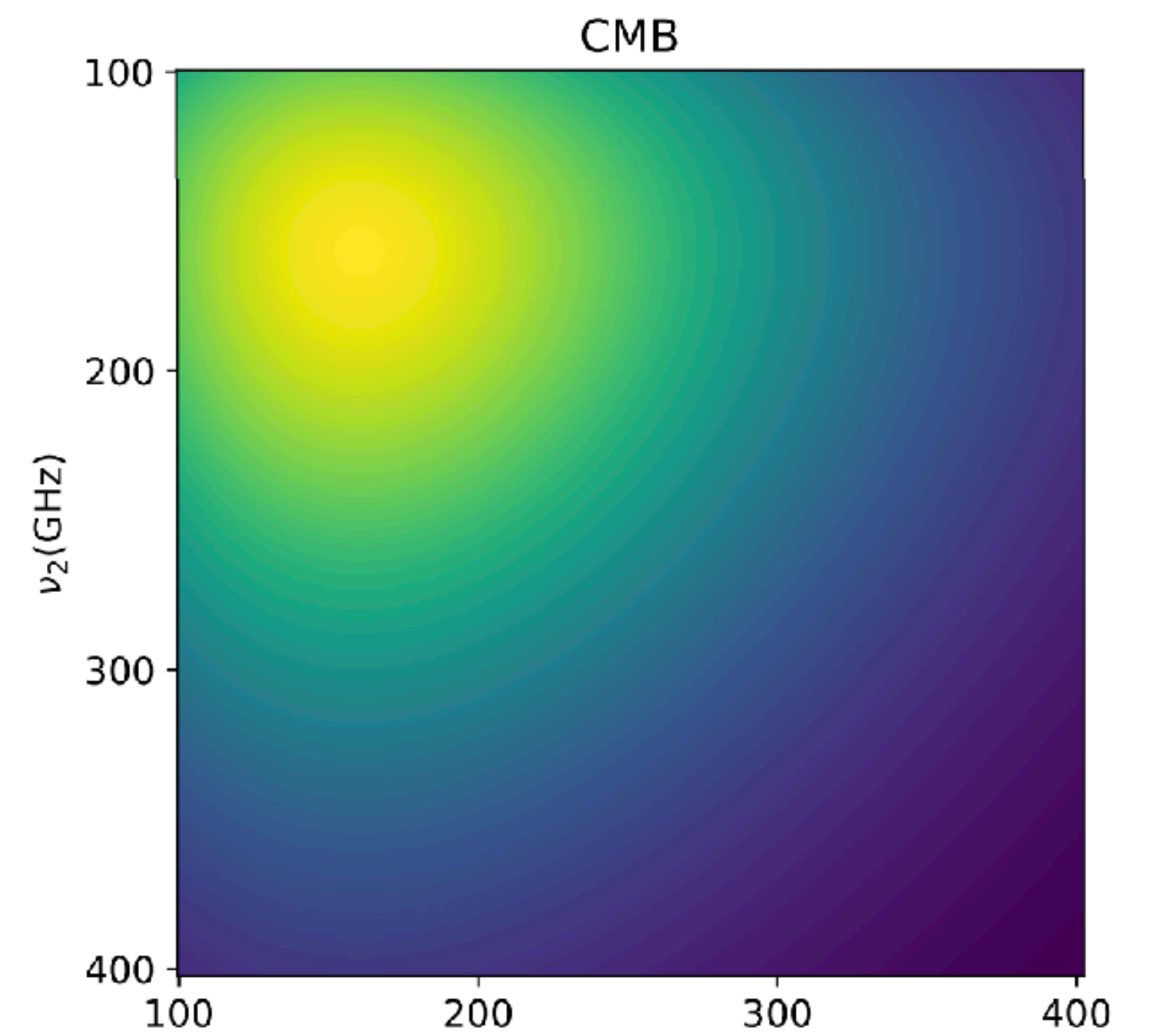
Thanks for listening !
Any questions ?



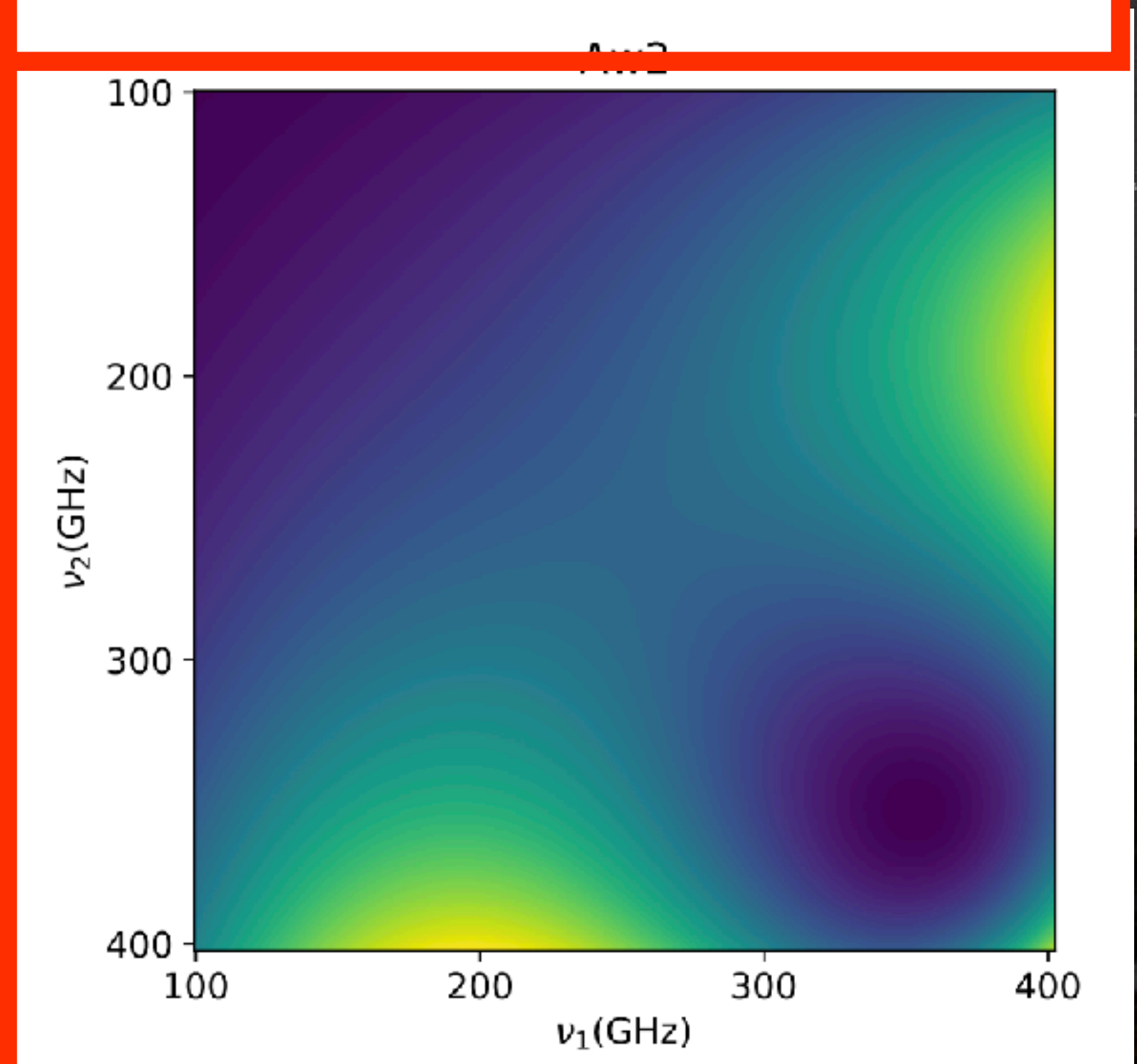
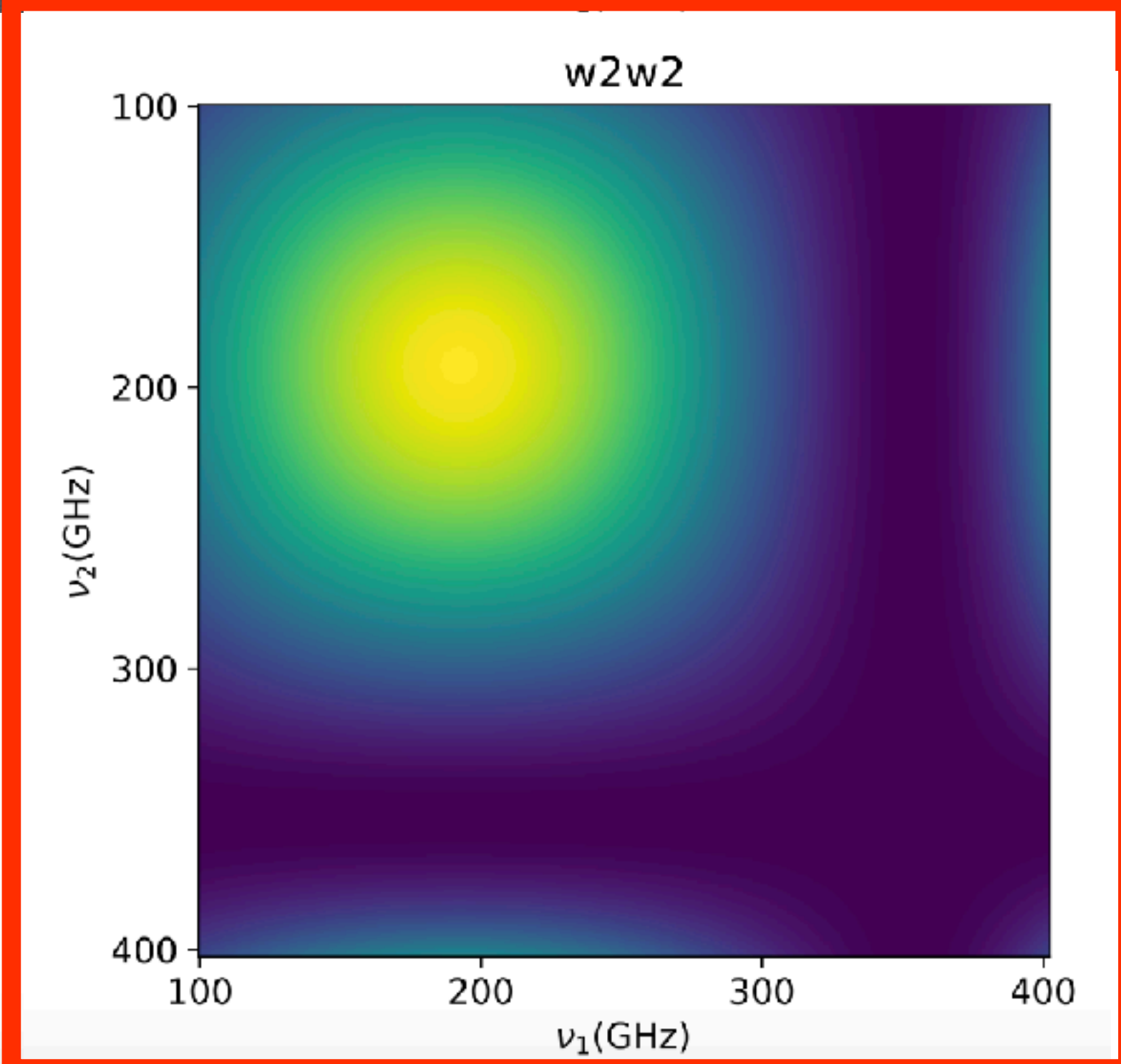
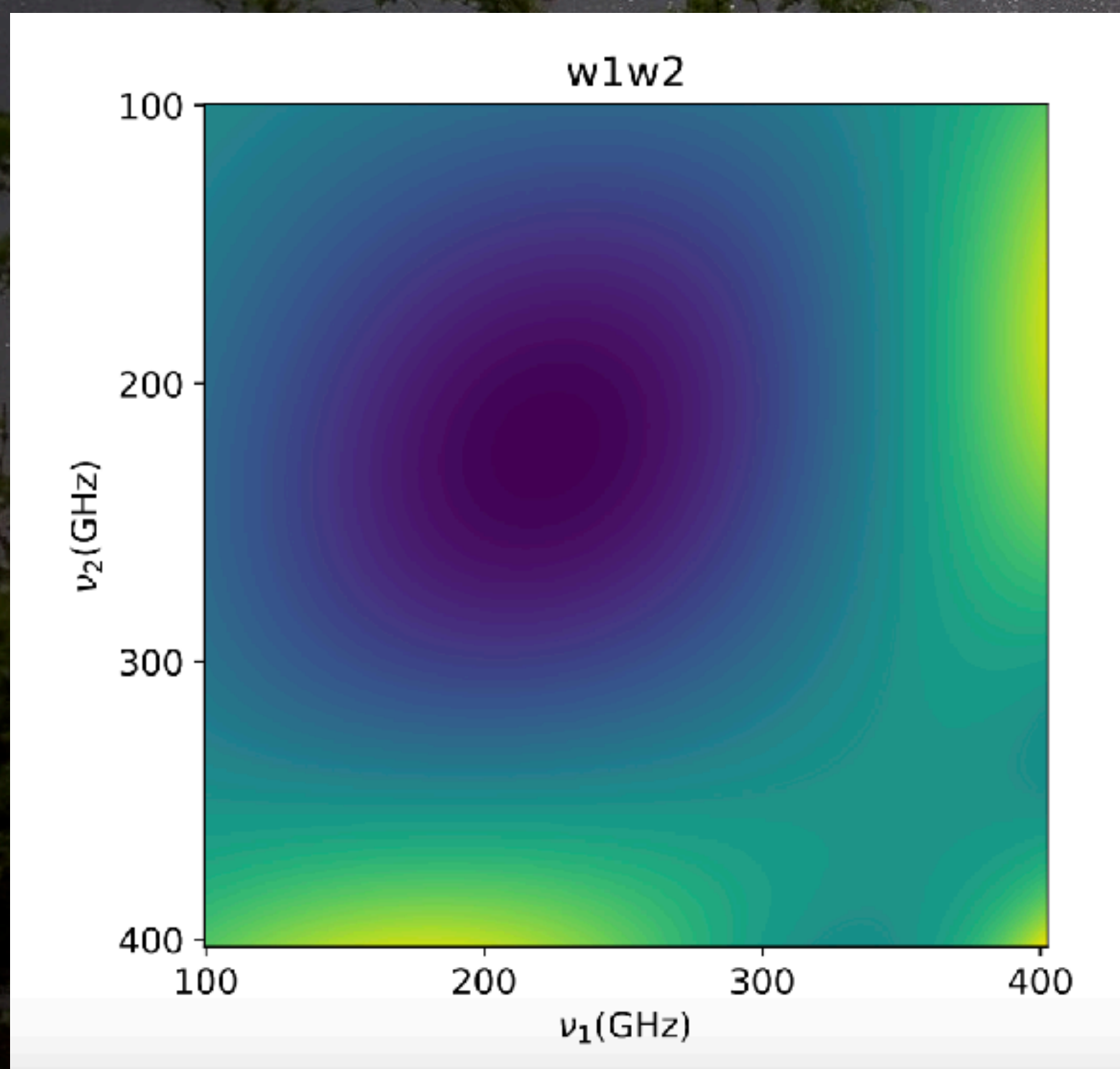
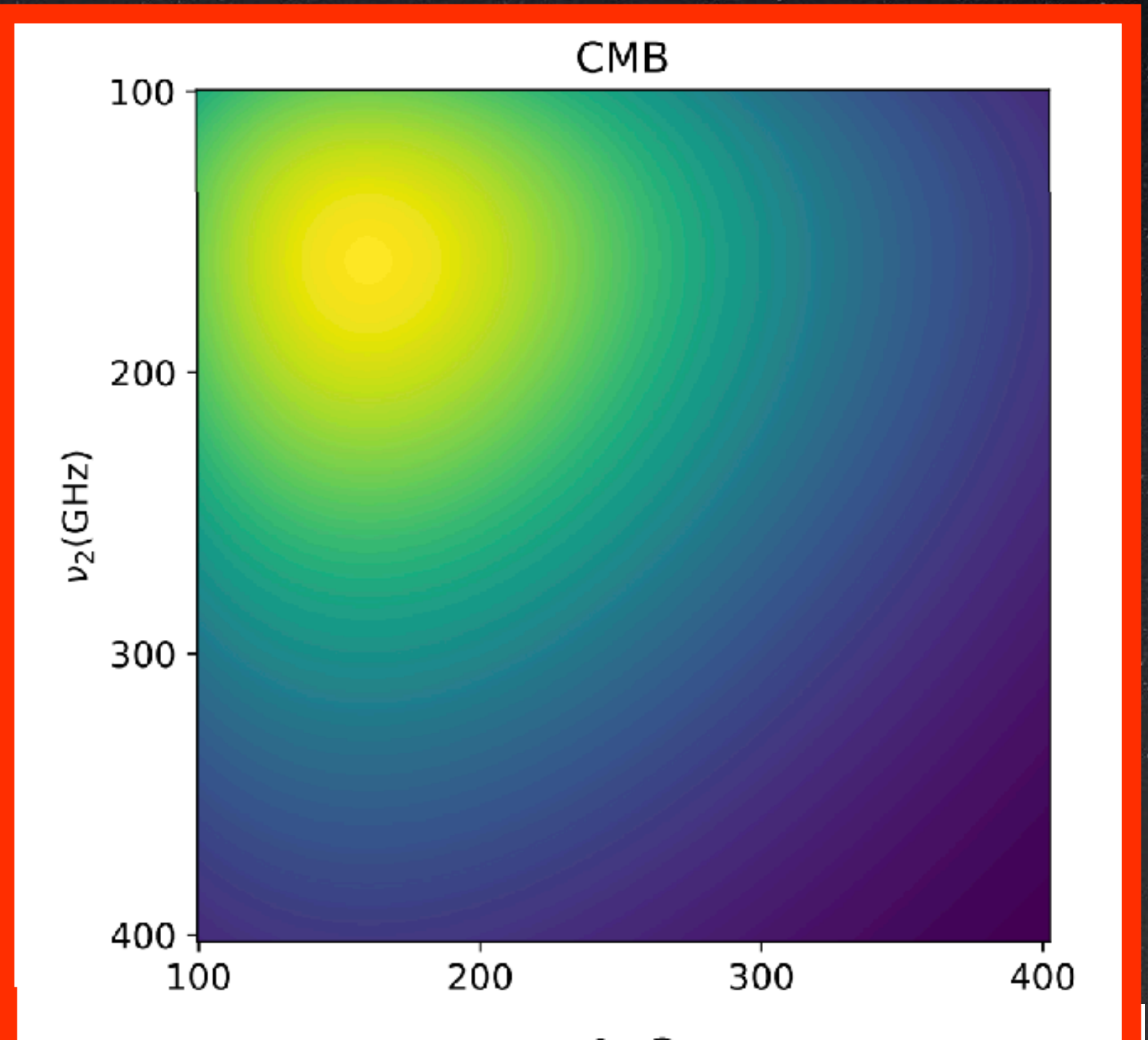
Correlations of order 1 coefficients with r



Correlations of order 2 coefficients with r



Correlations of order 2 coefficients with r



Correlations of order 2 terms with r

- ★ With LB freq. all order 2 terms are significantly correlated with the CMB
- ★ Correlation seems « universal » but exact value depends in a complex fashion on the instrumental parameters (sensitivity, frequency range ...)

