Algorithms for trajectography

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## main topics

- track finding
- track fitting
- progressive approach to Kalman Filter
- trajectory in a magnetic field
- vertex finding/fitting
- alignment/calibration




## What do we want ?

## measurements : gain of information

What was here?
vertices
3-momenta


> material : degradation of information

How to build the best estimator of the physical quantities?

## the ingredients <br> what is supposed to be known

- nature and precision of the measurements
- nature and magnitude of the "noises" in the matter (secondary interactions, multiple scattering, continuous energy loss) - equation of propagation (magnetic field)

Remarks: the nature of the particle (e, $\mu, \pi$, etc) may be unknown; the points above may depend on the mass hypothesis

## to be done

- grouping the local "hits" into track candidates (pattern recognition)
- fitting the parameters at origin (just after production)
if needed: iteration to solve the ambiguities
- inter/extrapolating to other detectors (RICH, muon chambers,...)
- if possible: information for particle identification ( $\mathrm{dE} / \mathrm{dx}, \ldots$ )
- finding primary/secondary vertices: topology and final fit



## pattern recognition vs final track fit

- aim of patt. rec.: find association of hits. The precision needed is the power of separation between hits, not the error on their position.
- the final track fit should give the best estimator, using a precise estimation of the positions of hits and the error on them, and the full covariance matrices of the track parameters.
- in practice, these tasks may interfere, and the whole procedure may be a more or less intricate combination of finding and fitting steps

Note: in many cases, the limiting factor is not the hit measurement error, but the noise (mainly multiple scattering). Do not be more royalist than the king !

## patt. rec. 1: extending tracks from seeds

 general principle: build seeds from a few shells, extrapolate to next shells as long as compatible hits are found tune criteria to:- accept a new point - confirm the track

- very flexible strategy (choice of shells for seeding, shell ordering,...)
- each new hit may be used to update the track parameters $\rightarrow$ better extrapolation
- may consists in successive passes, iterations, etc
- may need much tuning to optimize the trade-off between efficiency/ghost rate/speed



## patt. rec. 3: sample of routes

simulate trajectories of tracks of physical interest define the pattern of hits for each one collect enough patterns to cover the wanted phase space (e.g. $p_{t}>\min$ ) run time: flag the « filled» routes (flexible strategy to define the criteria of « filling »)


- OK for parallel computing with many small CPUs
- do not need any parameterization of trajectories
- large memory needed
- may produce multiple counting, ambiguities, ghosts


## pattern recognition in brief

- no universal solution: the procedure has to be adapted to the layout of the experiment
- in most cases, it consists of parallelizable sub-algorithms and more global cleaning steps (rejection of poor candidates, resolution of ambiguities)
- the best method is often a combination of different algorithms in successive steps
- the pattern recognition may internally use some track fitting procedures for a more precise discrimination and extrapolation. In general, the fit may be simplified
- machine learning may help to optimize the strategy




## gaussians in nD space

$$
\begin{aligned}
& \mathrm{G}(\mathbf{x})=\mathrm{K} \exp \left(-\Sigma \mathrm{W}_{\mathrm{ij}}\left(\mathrm{x}_{\mathrm{i}}-\mu_{\mathrm{i}}\right)\left(\mathrm{x}_{\mathrm{j}}-\mu_{\mathrm{j}}\right) / 2\right) \quad \mathrm{K}^{2}=\operatorname{det}(\mathrm{W}) /(2 \pi)^{\mathrm{n}} \\
& \text { covariance matrix } \mathrm{C}=\mathrm{W}^{-1}
\end{aligned}
$$

combining gaussians:
product: $\left(\boldsymbol{\mu}_{1}, \mathrm{~W}_{1}\right) \cdot\left(\boldsymbol{\mu}_{2}, \mathrm{~W}_{2}\right) \rightarrow\left(\mathrm{W}_{1}+\mathrm{W}_{2}\right)^{-1} \cdot\left(\mathrm{~W}_{1} \boldsymbol{\mu}_{1}+\mathrm{W}_{2} \boldsymbol{\mu}_{2}\right), \mathrm{W}_{1}+\mathrm{W}_{2}$ (« barycenter», addition of weight matrices)
convolution: $\left(\boldsymbol{\mu}_{1}, \mathrm{~W}_{1}\right) *\left(\boldsymbol{\mu}_{2}, \mathrm{~W}_{2}\right) \rightarrow \boldsymbol{\mu}_{1}+\boldsymbol{\mu}_{2},\left(\mathrm{~W}_{1}^{-1}+\mathrm{W}_{2}^{-1}\right)^{-1}$ (addition of biases, addition of covariance matrices)

$1 \sigma$ contours
quantitatively:
information = $1 /$ area ( $1 /$ volume in nD )


## a 1-parameter problem where is/was the flea?


a flea moves by jumps on $x$ axis; initial position: $x_{0}$ at each time step (independently):

- measurement (precision $\sigma$ )
- jump (standard deviation $\tau$ )
what is the "best" estimator of the position $x_{0}$ ? $x_{n}$ ?
intuitively :
- if $\sigma$ « $\tau$ : the instant one; the other ones are spoiled by the jumps
- if $\tau \sqrt{ } n<\sigma / V_{n}$ (that is $n \tau<\sigma$ ): the average of $n$ measurements
- intermediate case: not obvious; truncated mean? truncated weighted mean?
- the best linear estimator should be a weighted combination of the measurements

How to evaluate the weights ?

## The heavy optimal solution

One wants to estimate $\mathrm{x}_{0}$, accounting for the correlations between successive measurements:

```
x}\mp@subsup{}{0}{\mathrm{ mes }}=\mp@subsup{\textrm{x}}{0}{}+\mp@subsup{\varepsilon}{0}{
x}\mp@subsup{}{1}{mes}=\mp@subsup{x}{0}{}+\mp@subsup{\eta}{1}{}+\mp@subsup{\varepsilon}{1}{
x}\mp@subsup{}{2}{mes}=\mp@subsup{x}{0}{}+\mp@subsup{\eta}{1}{}+\mp@subsup{\eta}{2}{}+\mp@subsup{\varepsilon}{2}{
```

$\varepsilon_{\mathrm{k}}$ : meas. error at time $\mathrm{k} ; \eta_{\mathrm{k}}$ : jump at time k
covariance matrix $C$ of the deviations $\Delta x_{k}=x_{k}^{m e s}-x_{0}$ :
$\sigma^{2} 0 \quad 0 \quad 0 \ldots$
$0 \quad \sigma^{2}+\tau^{2} \quad \tau^{2} \quad \tau^{2} \ldots$
$0 \quad \tau^{2} \quad \sigma^{2}+2 \tau^{2} \quad 2 \tau^{2} \ldots$
$0 \quad \tau^{2} \quad 2 \tau^{2} \quad \sigma^{2}+3 \tau^{2} \quad 3 \tau^{2} \ldots$
$\chi^{2}=\Sigma\left(\mathrm{C}^{-1}\right)_{\mathrm{ij}} \Delta \mathrm{x}_{\mathrm{i}} \Delta \mathrm{x}_{\mathrm{j}} \rightarrow \mathrm{x}_{0}{ }^{\text {fit }}=\Sigma_{\mathrm{j}}\left(\mathrm{C}^{-1}\right)_{\mathrm{ij}} \mathrm{x}_{\mathrm{i}}^{\text {mes }}$
with $n$ measurements: matrix ( $n \times n$ ) to be inverted


## more (almost for free)

- final position $x_{n}$ :
forward filter (same procedure, going from 0 to $n$ )
- intermediate position $x_{k}$ (interpolation) : starting from both ends towards point $k$, combine independant backward and forward estimators $X_{n \rightarrow k}$ and $X_{0 \rightarrow k}$. $\mathrm{x}_{\mathrm{k}}{ }^{\text {mes }}$ may be omitted or included in one of them
(équivalent to the "smoother" in the kalmanian jargon)
- compatibility criterion : the variance of $x_{k}{ }^{\text {interp }}\left(w / o x_{k}{ }^{\text {mes }}\right)-x_{k}{ }^{\text {mes }}$ is $V$ (interp) $+\sigma^{2}$
- abnormal jump detected by comparing $X_{n \rightarrow k}-X_{0 \rightarrow k}$ to the predicted variance
in brief : with the forward filter and the backward filter (keeping the intermediate results) one can obtain all that

But: if one point is modified (e.g. one measurement added or removed), all following steps have to be redone).For example: if working on-the fly (incorporating measurements in real time), the backward filter would be heavy ... but probably useless




## linear approximation

## In real world : no exact linear model

possible solution:

- choose convenient parameters $\mathbf{p}$ (e.g. cartesian ou cylindrical coord.)
- define lines/surfaces (planes, cylinders,...) for measurements and material (the noise in a thin slice of material may be described by a matrix $\mathrm{C}_{\mathrm{b}}$ with a correlation between position and direction)
- define a reference trajectory $\mathbf{T}_{\text {ref }}$ close to the true one (from patt. rec. or preliminar fit)
- propagate the deviations $\delta \mathbf{p}$ of $\mathbf{p}$ from $\mathbf{T}_{\text {ref }}$ in the linear approximation:
$\mathrm{D}_{\mathrm{S} \rightarrow \mathrm{S}^{\prime}}=\partial\left(\delta \mathbf{p}^{\prime}\right) / \partial(\delta \mathbf{p})=\partial \mathbf{p}^{\prime} / \partial \mathbf{p}$ (jacobian matrix)
- apply the KF formalism; if needed, modify $\mathbf{T}_{\text {ref }}$ and iterate if the $\boldsymbol{\delta} \mathbf{p}$ are too large (it is also possible to change $\boldsymbol{T}_{\text {ref }}$ at some steps)


## a (false) technical problem: how to begin?

at start: insufficient information to define $\mathrm{p}_{0}$, and get inversible $\mathrm{C}_{0}, \mathrm{~W}_{0}$ example : the first measurement is x or a linear combination linéaire of x and $\mathrm{v} \rightarrow \mathrm{W}$ has a 0 eigenvalue (the p.d.f. is a stripe; $\mathrm{p}_{0}$ is degenerate along this stripe)
practically, the elementary matrix operations (convolution, propagation, product) are always possible :

- convolution : $\left(\mathrm{W}^{-1}+\mathrm{C}\right)^{-1}=(1+\mathrm{WC})^{-1} . \mathrm{W}$
$1+\mathrm{WC}$ is inversible in the useful cases
- propagation : $\mathrm{W}^{\prime}=\left(\mathrm{D}^{-1}\right)^{\mathrm{t}} . \mathrm{W} \cdot\left(\mathrm{D}^{-1}\right)$
- product : if $\mathrm{W}_{1}$ and/or $\mathrm{W}_{2}$ is singular, the system $\left(\mathrm{W}_{1}+\mathrm{W}_{2}\right) \mathrm{p}=\mathrm{W}_{1} \mathrm{p}_{1}+\mathrm{W}_{2} \mathrm{p}_{2}$ has a solution which does not depend on the choice of $p_{1}$ and $p_{2}$ on the axis of the stripes extreme case : parallel stripes : $p$ is undefined, and the result in again a stripe

```
one can use the weight matrices in all steps
```

usual method with the standard KF (using covariance matrices); start with large values in C.
but: possible problems of precision

## general case: 3D trajectory in Bfield (5 parameters)

which parameters?
it depends on the geometry of the tracking system
Examples:

- fixed target or endcap in a collider: surfaces: planes perpendicular to the beam (fixed $z$ )
- position: x,y
- direction: $\theta$ (or $\eta$ ) and $\phi$, or direction cosines $c_{x}, c_{y}$, or slopes $t_{x}=d x / d z, t_{y}=d y / d z$
- signed curvature ( $\mathrm{q} / \mathrm{R}$ or $\mathrm{q} / \mathrm{p}_{\mathrm{t}}$ ou $\mathrm{q} / \mathrm{p}$ )
- barrel in a collider, with $\mathbf{B}$ along z :
surfaces: cylinders (e.g. beam pipe + concentric shells) :
- position (angle $\Phi, z$ )
- direction (angles $\theta, \phi$ )
- curvature ( $q / R$ or $q / p_{t}$ ou $q / p$ )
procedure: same as before, with 5 -vectors for the state, $5 \times 5$ matrices for W,C,D


## "simple" measurement/noise

measurement of one coordinate, e.g. x:
$\mathrm{p}_{\text {meas }}=\left(\mathrm{x}_{\text {meas }}, 0,0,0,0\right) \quad \mathrm{W}_{\text {meas }}=\operatorname{diag}\left(1 / \sigma^{2}, 0,0,0,0\right)$
measurement of two coordinates $\mathrm{x}, \mathrm{y}$ :
$\mathrm{p}_{\text {meas }}=\left(\mathrm{x}_{\text {meas }}, \mathrm{y}_{\text {meas }}, 0,0,0\right) \quad \mathrm{W}_{\text {meas }}=\operatorname{diag}\left(1 / \sigma_{\mathrm{x}}{ }^{2}, 1 / \sigma_{\mathrm{y}}{ }^{2}, 0,0,0\right)$
scattering in a surface:
$\mathrm{C}_{\mathrm{ms}}=(2 \mathrm{x} 2)$ submatrix on $\mathrm{t}_{\mathrm{x}}, \mathrm{t}_{\mathrm{y}}$ (includes correlation)
scattering in a layer:
$\mathrm{C}_{\mathrm{ms}}=(4 \mathrm{x} 4)$ submatrix on $\mathrm{x}, \mathrm{y}, \mathrm{t}_{\mathrm{x}}, \mathrm{t}_{\mathrm{y}}$ (includes correlations)

## "oblique" measurements



- a combination is measured, e.g. $u=a x+b y$ (stereo) $\mathrm{w}_{\mathrm{u}}=1 / \sigma_{\mathrm{u}}^{2}$ ("weight" of the u measurement) contribution to $\chi^{2}$ : "stripe" in the ( $\mathrm{x}, \mathrm{y}$ ) plane $\mathrm{w}_{\mathrm{u}}\left(\mathrm{u}^{\text {mes }}-\mathrm{ax}-\mathrm{by}\right)^{2}=\left(\mathrm{x}-\mathrm{x}^{\text {mes }}, \mathrm{y}-\mathrm{y}^{\text {mes }}\right)^{\mathrm{t}} \mathrm{W}\left(\mathrm{x}-\mathrm{x}^{\text {mes }}, \mathrm{y}-\mathrm{y}^{\text {mes }}\right)$ $x^{\text {mes }}, y^{\text {mes }}$ : any point such that $a x^{\text {mes }}+b y^{\text {mes }}=u^{\text {mes }}$
$\mathrm{W}=(\mathrm{a}, \mathrm{b}) \cdot \mathrm{W}_{\mathrm{u}} \cdot(\mathrm{a}, \mathrm{b})^{\mathrm{t}}=1 / \sigma_{\mathrm{u}}^{2}\left(\mathrm{a}^{2} \mathrm{ab}, \mathrm{ab} \mathrm{b}^{2}\right)$ (matrix of rank 1)
- measurement in a detector which is oblique w.r.t the reference surface

trajectory of slope $\mathrm{a}=\mathrm{dx} / \mathrm{dz}$ measuring $\zeta$ (with error $\sigma$ ) in D amounts to measure $y=\zeta(\lambda+\mu a)$ with errror $|\lambda+\mu \mathrm{a}| . \sigma$
$\lambda, \mu:$ constants depending on geometry note: a is known at this stage (at least roughly)
general formulation for several measurements in the same detector:
contribution to $\chi^{2}=\left(\mathbf{p}-\mathbf{p}^{\text {mes }}\right)^{\mathrm{t}} \mathrm{W}_{\mathrm{p}}\left(\mathbf{p}-\mathbf{p}^{\text {mes }}\right)$ with $\mathrm{W}_{\mathrm{p}}=\mathrm{M}^{\mathrm{t}} \mathrm{W}_{\mathrm{m}} \mathrm{M}$
$=:$
$\mathrm{W}_{\mathbf{m}}$ : weight matrix of the measurements $\mathbf{m}$; M: dependence $\mathbf{d m} / \mathbf{d p}$


## exogenous measurements

some informations from non-trajectographic detectors may be injected at some stages on the filter:
examples:

- E measured in a calorimeter may be injected in the initial state of the backward filter as an estimator of $q / p$ (if the matching and the sign $q$ are inambiguous...)
- $\Delta \mathrm{E}$ mesured as a $\gamma$ energy in a calorimeter may be injected at an intermediate point or the trajectory (more delicate, but may be very useful for electrons...)


## not everything is gaussian in real world...

two kinds of "non-gaussianity"

- "short range" : e.g. measurement with uniform distribution in an interval smoothed by convolution (gaussian limit for large numbers)
- "with long tails": the gaussian limit may fail
practically, for charged particles :
- non-linearity in the propagation $\rightarrow$ distortion of the p.d.f.
- multiple sattering : low probability of a diffusion at large angle (à la Rutherford)
- energy loss:
. $\Delta \mathrm{E}$ through ionisation is almost déterministic, with small fluctuations . more violent occurrences : $\delta$-rays, and above all bremstrahlung (major problem for electrons)

If the gaussian approximation fails, what to do ?

## God's algorithm

5 -vector $\mathbf{p}$ to describe the state of the particle on a surface
chaining elementary operations on the p.d.f. $\mathrm{F}(\mathbf{p})$ :

- measurement (local) : multiplication by freas $(\mathrm{m}(\mathbf{p}))$
- noise (local) : convolution with $\mathrm{f}^{\text {floise }}(\mathbf{p})$
- propagation : changement of variables $\mathrm{F}(\mathbf{p}) \rightarrow \mathrm{Fr}^{\mathrm{pr}}\left(\mathbf{p}^{\mathrm{pr}}(\mathbf{p})\right)$ :
obvious difficulty: computing power needed for functions in a 5D space!

But : "On trouve avec le Ciel des accommodements" (Tartuffe)

## the gaussian sum

principle: approximation of $\mathrm{F}(\mathbf{p})$, $\mathrm{f}^{\text {meas }}$ et $\mathrm{f}^{\text {noise }}$ by a sum of gaussian functions
$F(\mathbf{p})=\Sigma \alpha_{\mathrm{i}} \mathrm{G}_{\mathrm{i}}(\mathbf{p})$ with $\mathrm{G}_{\mathrm{i}}(\mathbf{p})=\mathrm{C}_{\mathrm{i}} \exp \left(-\left(\mathbf{p}-\mathbf{p}_{\mathrm{i}}\right)^{\mathrm{t}} \mathrm{W}_{\mathrm{i}}\left(\mathbf{p}-\mathbf{p}_{\mathrm{i}}\right) / 2\right)$

- works well in many cases for $\mathrm{f}^{\text {meas }}$ et $\mathrm{f}^{\text {floise }}$ (function of 1 variable)
- $F$ is defined and positive everywhere if all $\alpha_{I}>0$, and it vanishes at infinity
- the operations (product, convolution, linear propagation) are easy and give again a sum of gaussians
product : $\left(\mathbf{p}_{1}, \mathrm{~W}_{1}\right) \times\left(\mathbf{p}_{2}, \mathrm{~W}_{2}\right)=\left(\left(\mathrm{W}_{1}+\mathrm{W}_{2}\right)^{-1}\left(\mathrm{~W}_{1} \mathbf{p}_{1}+\mathrm{W}_{2} \mathbf{p}_{2}\right), \mathrm{W}_{1}+\mathrm{W}_{2}\right)$
convolution: $\left(\mathbf{p}_{1}, \mathrm{~W}_{1}\right) *\left(\mathbf{p}_{2}, \mathrm{~W}_{2}\right)=\left(\mathbf{p}_{1}+\mathbf{p}_{2},\left(\mathrm{~W}_{1}^{-1}+\mathrm{W}_{2}^{-1}\right)^{-1}\right)$

But : the number of components increases multiplicatively possible remedies:

- suppress components of low amplitude
- merge nearby components into one
$\rightarrow$ to be optimized for each case, depending on the final impact on physics results
in practice: used mainly for electron trajectories


## propagation: the Runge-Kutta integration method

generic problem: $\quad y^{\prime}=f(t, y), \quad y\left(t_{0}\right)=y_{0} \quad$ solved by steps $h$ in $t$

1 step

$$
\begin{aligned}
& k_{1}=f\left(t_{n}, y_{n}\right) \\
& k_{2}=f\left(t_{n}+\frac{h}{2}, y_{n}+\frac{h}{2} k_{1}\right) \\
& k_{3}=f\left(t_{n}+\frac{h}{2}, y_{n}+\frac{h}{2} k_{2}\right) \\
& k_{4}=f\left(t_{n}+h, y_{n}+h k_{3}\right)
\end{aligned}
$$

steps along $z$ axis in a magnetic field RK applied to the state vector $\left(\mathrm{x}, \mathrm{y}, \mathrm{t}_{\mathrm{x}}, \mathrm{t}_{\mathrm{y}}\right)$

$$
\begin{aligned}
\frac{\mathrm{dx}}{\mathrm{dz}} & =t_{x} \\
\frac{\mathrm{dy}}{\mathrm{dz}} & =t_{y} \\
\frac{\mathrm{dt}_{\mathrm{x}}}{\mathrm{dz}} & =c \frac{q}{p} \sqrt{1+t_{x}^{2}+t_{y}^{2}}\left(t_{x} t_{y} B_{x}-\left(1+t_{x}^{2}\right) B_{y}+t_{y} B_{z}\right) \\
\frac{\mathrm{dt}_{\mathrm{y}}}{\mathrm{dz}} & =c \frac{q}{p} \sqrt{1+t_{x}^{2}+t_{y}^{2}}\left(\left(1+t_{y}^{2}\right) B_{x}-t_{x} t_{y} B_{y}-t_{x} B_{z}\right)
\end{aligned}
$$

## parameterized propagation

idea: instead of using RK extrapolation for every track, precompute formulae to get a faster execution principle:

- chose a few reference surfaces that will contain « nodes» of the Kalman Filter.
- to go from the initial surface $\Sigma_{\mathrm{i}}$ to the final one $\Sigma_{\mathrm{f}}$, express the state vector $\mathbf{S}_{\mathrm{f}}$ on $\Sigma_{\mathrm{f}}$ through analytical of tabulated functions of the components of the state vector $\mathbf{S}_{\mathrm{i}}$ on $\Sigma_{\mathrm{i}}$
guiding criteria
- at infinite momentum, the trajectory is a straight line
- so, we can try an expansion in powers of $q / p$ of $\Delta \mathbf{S}_{f}$, the difference between $\mathbf{S}_{f}$ and the straight line extrapolation
- the precision should be small compared to the other sources of error (mainly multiple scattering)
- the phase space may be reduced for trajectories close to the origin (particles for physics analysis)
first example in the « endcap » description ( $x, y, t_{x}, t_{y}, q / p$ at fixed $z$ ): propagate from $z_{i}=0$ to $z_{f}$
$-t_{x}$ and $t_{y}$ are bounded by the acceptance ;
- $x_{i}$ and $y_{i}$ are small, so terms at first order in $x_{i} y_{i}$ are sufficient


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## explicit formulae

$$
\Delta \mathbf{S}_{\mathbf{f}}=\sum_{k} \mathbf{A}_{k}\left(t_{x i}, t_{y i}\right)(q / p)^{k}+\sum_{k}\left(x_{i} \mathbf{B}_{k}\left(t_{x i}, t_{y i}\right)+y_{i} \mathbf{C}_{k}\left(t_{x i}, t_{y i}\right)\right)(q / p)^{k}
$$

this gives 4 expansions (for $x_{f} . x_{f}, t_{x f}, t_{x f}$ ), assuming $p$ to be constant, e.g. for $x_{f}$ :

$$
x_{f}=x_{i}+z_{f} t_{x i}+\sum_{k} A_{k}^{x}\left(t_{x i}, t_{y i}\right)(q / p)^{k}+\sum_{k}\left(x_{i} B_{k}^{x}\left(t_{x i}, t_{y i}\right)+y_{i} C_{k}^{x}\left(t_{x i}, t_{y i}\right)\right)(q / p)^{k}
$$

the coefficients $A, B, C$ may be tabulated or expressed as analytic functions of $t_{x i}, t_{y i}$

## byproducts

- jacobian matrix D: straightforward derivatives w.r.t. $x_{i}, y_{i}, q / p$, easy for $t_{x i}, t_{y i}$
- reverse propagation with the Newton-Raphson method:
starting from $S_{f}$, we want to find $S_{i}$ such that $S_{i} \rightarrow S_{f}$
if $\mathrm{S}_{\mathrm{i}}{ }^{0}$ is a good approximation, and $\mathrm{S}_{\mathrm{i}}{ }^{0} \rightarrow \mathrm{~S}_{\mathrm{f}}{ }^{0}$, then $\mathrm{S}_{\mathrm{f}} \approx \mathrm{S}_{\mathrm{f}}{ }^{0}+\mathrm{D} .\left(\mathrm{S}_{\mathrm{i}}-\mathrm{S}_{\mathrm{i}}{ }^{0}\right)$ so $\mathrm{S}_{\mathrm{i}} \approx \mathrm{S}_{\mathrm{i}}^{0}-\mathrm{D}^{-1}$. $\left(\mathrm{S}_{\mathrm{f}}-\mathrm{S}_{\mathrm{f}}^{0}\right)$
that is: we just need a direct propagation + a linear transform if needed: iterate (the convergence is very fast)
- propagation from $\mathrm{z}_{\mathrm{i}}$ to $\mathrm{z}_{\mathrm{f}}$ with $\mathrm{z}_{\mathrm{i}} \neq 0: \mathrm{z}_{\mathrm{i}} \rightarrow 0$ then $0 \rightarrow \mathrm{z}_{\mathrm{f}}$
 jacobian matrix $\mathrm{D}_{\mathrm{if}}=\mathrm{D}_{0 \mathrm{f}}{ }^{-1} \cdot \mathrm{D}_{\mathrm{i} 0}$
possible implementation: choose a few «main surfaces» for the full formulae and complement by short range extrapolation (1 step of RK or simpler local parameterization)

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## the full vertex fit

aim: use the convergence of trajectories to improve their reonstruction (add a virtual measurement and increase the lever arm)

first trial: fit the position as before, and introduce this point as an additional measurement to all tracks.
not optimal: this position is correlated to the other measurements on the track
second trial: iterative procedure: adjust alternatively the vertex position and the $\mathbf{p}_{\mathrm{i}}$ (3-momenta of the particles at the vertex) to fit the extrapolations to $\mathbf{q}_{\mathrm{i}}$ possible but the convergence may be slow (zig-zag path)

## the vertex fit as a hierarchical fit

"all in one" method: from a sample of $n$ trajectoires ( $\mathbf{q}_{\mathrm{i}}, \mathrm{W}_{\mathrm{i}}$ ) at initial point ( 5 n parameters) fit simultaneously $3 \mathrm{n}+3$ parameters with the constraint of convergence:

- the position $\mathbf{V}(\mathrm{X}, \mathrm{Y}, \mathrm{Z})$ of a common origin
- the 3-momenta $\mathbf{p}_{\mathrm{i}}$ of the particles at this point (or equivalently $\mathrm{q} / \mathrm{p}_{\mathrm{i}}, \theta_{\mathrm{i}}, \phi_{\mathrm{I}}$ )
tool : propagation function $\mathbf{q}=\mathbf{F}(\mathbf{V}, \mathbf{p})$ from vertex to initial point (simple if the initial point is close to the vertex, e.g. the perigee)

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formulation with a global }\mp@subsup{\chi}{}{2}\mathrm{ :
find }\mathbf{V}\mathrm{ and the }\mp@subsup{\textrm{p}}{\textrm{i}}{}\mathrm{ which minimize
\chi
```

a priori : problem in a space of dimension $3 n+3$
actually : hierarchical problem: 3 global param. +3 particular param. for each track
$\min \left(\chi^{2}\right)=\min \mid \mathbf{V}\left[\Sigma \min \mid p_{i}\left(\mathbf{q}_{\mathrm{i}}{ }^{\text {mes }}-\mathbf{F}\left(\mathbf{V}, \mathbf{p}_{\mathrm{i}}\right)\right)^{\mathrm{t}} \mathrm{W}_{\mathrm{i}}\left(\mathbf{q}_{\mathrm{i}}{ }^{\text {mes }}-\mathbf{F}\left(\mathbf{V}, \mathbf{p}_{\mathrm{i}}\right)\right)\right]$
the "internal" et "external" minimizations have dimension 3
Note: the "nesting" remains valid without the gaussian approximation
that is: you can use e.g. Minuit with a fon which itself calls n times Minuit (it
works actually !)

## other example of "hierarchical" fit (1)

sample of signals of the same shape, but with different amplitudes and dates: $\mathrm{S}(\mathrm{t})=\mathrm{A}_{\mathrm{i}} \mathrm{f}\left(\mathrm{t}-\mathrm{a}_{\mathrm{i}}\right)$; each one is measured at n times $\mathrm{t}_{\mathrm{k}} \rightarrow \mathrm{S}_{\mathrm{ik}}$ mes $=\mathrm{A}_{\mathrm{ik}} \mathrm{f}\left(\mathrm{t}_{\mathrm{k}}-\mathrm{a}_{\mathrm{i}}\right)+\varepsilon^{\text {mes }}$ the shape is defined by global parameters $p_{1}, p_{2}, \ldots$ to be fitted
e.g. here $f(t)=0$ for $t<0, \exp \left(-p_{1} t\right)-\exp \left(-p_{2} t\right)$ for $t>0$



how to extract $\mathrm{p}_{1}$ and $\mathrm{p}_{2}$ from these measured signals ?

## other example of "hierarchical" fit (2)

a set of events from the Surface Detector of AUGER (atmospheric showers)
signal in a tank at distance $r_{i}$ from shower axis: $S_{i}=A_{i} f\left(r_{i}\right)$

- global parameters $p, q$ for the shape, for example: $f(r)=1 / r^{p}\left(r+r_{1}\right)^{q}$
- individual parameters for each event: position ( $\mathrm{x}_{\mathrm{i}}, \mathrm{y}_{\mathrm{i}}$ ) of the core, amplitude $\mathrm{A}_{\mathrm{i}}$



## linearization

if $\mathbf{V} \approx \mathbf{V}_{0}$ (vertex) and $\mathbf{p}_{\mathrm{i}} \approx \mathbf{p}_{\mathrm{i} 0}$ (for every track):
$\mathbf{q}_{\mathrm{i}}=\mathbf{F}_{\mathrm{i}}\left(\mathbf{V}, \mathbf{p}_{\mathrm{i}}\right) \approx \mathbf{q}_{\mathrm{i} 0}+\mathrm{D}_{\mathrm{i}} \cdot\left(\mathbf{V}-\mathbf{V}_{0}\right)+\mathrm{E}_{\mathrm{i}} \cdot\left(\mathbf{p}_{\mathrm{i}}-\mathbf{p}_{\mathrm{i} 0}\right) \quad$ (short range propagation) $E_{i}$ et $D_{i}:(5 \times 3)$ matrices, simple to compute if $\mathbf{q}_{i}$ is at the perigee
setting $\Delta \mathbf{q}_{\mathrm{i}}=\mathbf{q}_{\mathrm{i}}^{\text {meas }}-\mathbf{q}_{\mathrm{i} 0}$, on can fit $\delta \mathbf{V}=\mathbf{V}-\mathbf{V}_{0}$ and the $\delta \mathbf{p}_{\mathrm{i}}=\mathbf{p}_{\mathrm{i}}-\mathbf{p}_{\mathrm{i} 0}$ to minimize $\chi^{2}=\Sigma\left(\Delta \mathbf{q}_{\mathrm{i}}-\mathrm{D}_{\mathrm{i}} \delta \mathbf{V}-\mathrm{E}_{\mathrm{i}} \delta \mathbf{p}_{\mathrm{i}}\right)^{\mathrm{t}} \mathrm{W}_{\mathrm{i}}\left(\Delta \mathbf{q}_{\mathrm{i}}-\mathrm{D}_{\mathrm{i}} \delta \mathbf{V}-\mathrm{E}_{\mathrm{i}} \delta \mathbf{p}_{\mathrm{i}}\right)$

- one block of 3 equations on the full set of parameters:
$\mathrm{A} \delta \mathbf{V}+\Sigma \mathrm{B}_{\mathrm{i}} \delta \mathbf{p}_{\mathrm{i}}=\mathbf{T}$ (1) with $\mathrm{A}=\Sigma \mathrm{D}_{\mathrm{i}}{ }^{\mathrm{t}} \mathrm{W}_{\mathrm{i}} \mathrm{D}_{\mathrm{i}}, \mathrm{B}_{\mathrm{i}}=\mathrm{D}_{\mathrm{i}}{ }^{\mathrm{t}} \mathrm{W}_{\mathrm{i}} \mathrm{E}_{\mathrm{i}}, \mathbf{T}=\Sigma \mathrm{D}_{\mathrm{i}}{ }^{\mathrm{t}} \mathrm{W}_{\mathrm{i}} \Delta \mathbf{q}_{\mathrm{i}}$
- $n$ blocks de 3 equations on $\mathbf{V}$ and one $\mathbf{p}_{i}$ :
$B_{i}{ }^{t} \delta \mathbf{V}+C_{i} \delta \mathbf{p}_{i}=\mathbf{U}_{i}$ (2) with $C_{i}=E_{i}^{t} W_{i} E_{i}, \mathbf{U}=\Sigma E_{i}^{t} W_{i} \Delta \mathbf{q}_{i}$

(sparse system by blocks $3 \times 3$ )


## resolution of the linear system

from equations (2) one can express the $\delta \mathbf{p}_{\mathrm{i}}$ as functions of $\delta \mathbf{V}$

$$
\begin{equation*}
\delta \mathbf{p}_{\mathrm{i}}=\mathrm{C}_{\mathrm{i}}^{-1}\left(\mathrm{U}_{\mathrm{i}}-\mathrm{B}_{\mathrm{i}}^{\mathrm{t}} \delta \mathbf{V}\right) \tag{3}
\end{equation*}
$$

injecting these expressions in (1) one obtains an equation in $\delta \mathbf{V}$ only

$$
\begin{equation*}
\left(\mathrm{A}-\Sigma \mathrm{B}_{\mathrm{i}} \mathrm{C}_{\mathrm{i}}^{-1} \mathrm{~B}_{\mathrm{i}}^{\mathrm{t}}\right) \delta \mathbf{V}=\mathrm{T}-\Sigma \mathrm{B}_{\mathrm{i}} \mathrm{C}_{\mathrm{i}}^{-1} \mathrm{U}_{\mathrm{i}} \tag{4}
\end{equation*}
$$

(4) gives $\delta \mathbf{V}$ then each of the equations (3) gives $\delta \mathbf{p}_{\mathrm{i}}$
as a bonus, we obtain also the full $(3 n+3) \times(3 n+3)$ covariance matrix $\ldots$

$$
\begin{aligned}
& \operatorname{cov}(\mathbf{V}, \mathbf{V})=\left(\mathrm{A}-\Sigma \mathrm{B}_{\mathrm{i}} \mathrm{C}_{\mathrm{i}}^{-1} \mathrm{~B}_{\mathrm{i}}^{\mathrm{t}}\right)^{-1} \\
& \operatorname{cov}\left(\mathbf{V}, \mathbf{p}_{\mathrm{i}}\right)=-\operatorname{cov}(\mathbf{V}, \mathbf{V}) \mathrm{B}_{\mathrm{i}} \mathrm{C}_{\mathrm{i}}^{-1} \\
& \operatorname{cov}\left(\mathbf{p}_{\mathrm{i}}, \mathbf{p}_{\mathrm{j}}\right)=\delta_{\mathrm{ij}} \mathrm{C}_{\mathrm{i}}^{-1}+\mathrm{C}_{\mathrm{i}}^{-1} \mathrm{~B}_{\mathrm{i}}^{\mathrm{t}} \operatorname{cov}(\mathbf{V}, \mathbf{V}) \mathrm{B}_{\mathrm{j}} \mathrm{C}_{\mathrm{j}}^{-1}
\end{aligned}
$$

note that this procedure introduces correlations between the 3-momenta of all particles in the vertex, to be used in principle in the physics analysis ...

## flexibility

## (adding or removing one particle)

to add a track (fitted as $\mathbf{q}_{\mathrm{n}+1}, \mathrm{~W}_{\mathrm{n}+1}$ ):

- add a triplet of parameters $\delta \mathbf{p}_{\mathrm{n}+1}$
- add in (1) $D_{n+1}^{t} W_{n+1} D_{n+1}$ to $A$, and one term $B_{n+1}=D_{n+1}^{t} W_{n+1} E_{n+1}$
- add in (2) one block of equations $\mathrm{B}_{\mathrm{n}+1}{ }^{\mathrm{t}} \delta \mathbf{V}+\mathrm{C}_{\mathrm{n}+1} \delta \mathbf{p}_{\mathrm{n}+1}=\mathbf{U}_{\mathrm{n}+1}$
taking as starting values the result of the fit with $n$ particles $\left(\mathbf{V}_{0}, \mathbf{p}_{\mathrm{i} 0}\right.$ for $\left.\mathrm{i}=1 \ldots \mathrm{n}\right)$ :
$\left(A+A_{n+1}\right) \delta V+\sum B_{i} \delta p_{i}=T_{n+1}$
$B_{i}^{t} \delta V+C_{i} \delta p_{i}=0$ for $i=1 \ldots n$
$B_{n+1}{ }^{\mathrm{t}} \boldsymbol{\delta} \mathbf{V}+\mathrm{C}_{\mathrm{n}+1} \delta \mathbf{p}_{\mathrm{n}+1}=\mathbf{U}_{\mathrm{n}+1}$
resolution:
$\left(A-\Sigma B_{i} C_{i}^{-1} B_{i}^{t}+A_{n+1}-B_{n+1} C_{n+1}^{-1} B_{n+1}^{t}\right) \quad \delta V=T_{n+1}-B_{n+1} C_{n+1}^{-1} U_{n+1}$ only the terms in red are computed : fast procedure $\rightarrow$ many combinations may be tried
removing a track $=$ adding it with a weight $-\mathrm{W}_{\mathrm{i}}$
Remark : the beam may be considered as a track to be added in a primary vertex (in general: very precise measurement of $x, y$, but $z$ is undefined)


## vertex fit with constraint(s)

examples:

- prompt or distant decay (neutral $\rightarrow+-$ ) with mass hypothesis
$-\gamma \rightarrow e^{+} e^{-}$with parallel tracks at the decay point;
in both cases: $\mathbf{p}$ points towards the main vertex (or just the beam line)
- more generally: combination of kinematical and geometrical constraints

Lagrange multipliers: universal tool
$\min \mid \mathbf{p}(\mathrm{F}(\mathbf{p}))$ with the constraint $\mathrm{C}(\mathbf{p})=0 \Leftrightarrow \min \mid \mathbf{p}, \lambda(\mathrm{F}(\mathbf{p})+\lambda \mathrm{C}(\mathbf{p}))$
easy to solve in the following approximation around the minimum (or maximum) :

- the $\chi^{2}$ or the log-likelihood is a quadratic function of the variations of parameters
- the constraint is linear
linear system $\rightarrow \mathbf{p}$ as a function of $\lambda$, then elimination of $\lambda$ with a linear equation
generalisation to several constraints:
$\min \mid \mathbf{p}, \lambda_{1}, \lambda_{2}, \ldots\left(\mathrm{~F}(\mathbf{p})+\lambda_{1} \mathrm{C}_{1}(\mathbf{p})+\lambda_{2} \mathrm{C}_{2}(\mathbf{p})+\ldots\right)$



## summary for track and vertex fit

- one can build a track fitting procedure by linking elementary operations on the local parameters trajectory (adding one measurement, adding one noise, propagation)
- when putting these operations in order, each step uses independent inputs
- in the linear approximation (almost always valid in useful cases), the steps are simple manipulations of 5 -vectors and ( $5 \times 5$ ) matrices
- in the gaussian approximation one can define quality tests in terms of Prob(chi2), either for the global fit, or for a given point (detection of outliers)
- exogenous measurements may be injected at some steps (e.g. detectable energy losses)
- if needed, some non gaussian effects may be taken into account (esp. for electrons)
- the track fit may be coupled to the pattern recognition to refine prediction to a layer (a large variety de strategies are possible)
- the vertex fit may be achieved in a fast procedure (CPU time proportional to the number of tracks) with flexibility (adding or removing a track is easy)
- geometrical and physical constraints may be added to improve the final reconstruction: invariant masses, combination of connected vertices in a decay tree

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GDR InF - tracking

## procedures of alignment

detector $=$ assembly of elements supposed to be rigid geometrical degrees of freedom for each element: translation, rotation; expansion, contraction?
first order: position of frames (« hardware » sensors) second order: fine corrections (position of sensitive elements) using signals from tracks (beam, cosmics, collision data)

## calibration

determine shape of signals, biases, measurement errors

- simulation
- external inputs (beam, cosmics, point sources, pulses on elements of the electronic chain)
- internal
- with or without B field


## 1. beam + external hodoscope

module to be internally aligned

the module should be moveable: impossible for big detectors

## 2. cosmic rays

mainly vertical (esp. in underground places) ; random impact position

may connect different modules of a big detector no hodoscope! weak modes may exist

## 3. internal track sample

large statistics, real time data, but useful tracks come mainly from origin

momentum dependent: again weak modes

## just for fun

when using a sample of tracks to make an alignment, you have to adjust:

- a few global parameters (the geometrical ones you want to obtain)
- individual parameters for each track (position, direction + curvature if magnetic field)
this is again a hierachical fit!



## examples of correlation between alignment and field map

in a field mainly along y axis:

- positive/negative particles get in average
negative/positive $\mathrm{x}, \mathrm{t}_{\mathrm{x}}$ : the $\mathrm{x}>0$ and $\mathrm{x}<0$ have opposite average charge populations
a positive $\delta B_{z}$ pushes both signs upwards partially compensated by pushing a z-plane downwards (negative $\Delta \mathrm{y}$ )

$\ldots=$
- an increase of $\left|\mathrm{B}_{\mathrm{y}}\right|$ increases the divergence
partially compensated by pushing a z-plane
- towards negative z
- if separate alignment of $x$-sides: pushing them inopposite directions
$-$an alignment by tracks may give different results depending on the range of momenta



## backup

## a problem of precision (LHCb)

trying to implement the Kalman Filter included in PrPixelTracking (Velo) in single precision on a GPU:

- discrepancies between the GPU and the CPU results, and between them and the weight/information algorithm, when applied to the same data
- more precisely: the discrepancies (on fitted position/slope, covariance matrix, chi2) decrease with the number of points in the track
- agreement between all versions in double precision, and between single and double with the weight formalism
- the discrepancies increase with the initial value given to $\operatorname{cov}(T x, T x)$ and $\operatorname{cov}(T y, T y)$ at the beginning of the loop over points


## origin and solution of the problem

```
// compute the prediction
```

const float $d z=$ zhit $-z$;
const float predx $=\mathrm{x}+\mathrm{dz} * \mathrm{tx}$;
const float dz_t_covTxTx $=\mathrm{dz} * \operatorname{covTxTx}$;
const float predcovXTx $=\operatorname{cov} X T x+d z \_t \_c o v T x T x ;$
const float $d x \_t \_c o v X T x=d z * \operatorname{covXTx}$;
const float predcovXX $=\operatorname{cov} X X+2 * d x \_t \_\operatorname{cov} X T x+d z * d z \_t \_c o v T x T x ;$
const float predcovTxTx $=$ covTxTx;
// compute the gain matrix
const float $\mathrm{R}=1.0 /(1.0 /$ whit + predcovXX);
const float $\mathrm{KX}=\operatorname{predcovXX} * \mathrm{R}$;
const float $\mathrm{KTx}=$ predcovXTx $* \mathrm{R}$;
// update the state vector
const float $r=x h i t-p r e d x ;$
$x=p r e d x+K x * r$;
$t x=t x+K T x * r$;
// update the covariance matrix. we can write it in many ways ...
$\operatorname{covXX} / *=$ predcovXX $-K x *$ predcovXX */ $=(1-K x) *$ predcovXX;
covXTx /*= predcovXTx - predcovXX * predcovXTx / R */ = (1-Kx) * predcovXT
covTxTx $=$ predcovTxTx $-K T x *$ predcovXTx;
// return the chi2
return $\mathrm{r} * \mathrm{r} * \mathrm{R}$;
at first point $\mathrm{C}_{\mathrm{xx}}=\sigma^{2}, \mathrm{C}_{\mathrm{TxTx}}=\mathbf{B i g}$ (in this code: $\mathbf{B i g}=1$ ) the loop (pred, upd, noise) begins at the second point with a nearly singular predicted covariance :
$\mathrm{C}^{\prime}{ }_{\mathrm{xx}}=\sigma^{2}+\mathbf{B i g}^{2} \Delta \mathrm{z}^{2}$
$\mathrm{C}^{\prime}{ }_{\mathrm{xTx}}=\boldsymbol{\operatorname { B i g }} \Delta \mathrm{z}, \mathrm{C}^{\prime}{ }_{\mathrm{TxTx}}=\mathbf{B i g}$
the «gain» business mixes Big and real quantities $\rightarrow$ rounding errors !
here: making $\operatorname{Big} \rightarrow \infty$ in the results after updating at point 2 :
$\mathrm{x}=\mathrm{x}_{2} \quad \mathrm{Tx}=\left(\mathrm{x}_{2}-\mathrm{x}_{1}\right) / \Delta \mathrm{z}$
$\operatorname{cov}=\left(\sigma^{2}, \sigma^{2} / \Delta \mathrm{z}, 2 \sigma^{2} / \Delta \mathrm{z}^{2}\right)$
$\chi^{2}=0$
the KF machinery was useless in the first step !
conclusion: do not rely blindly on black boxes put your eyes inside, and put your hands if needed

