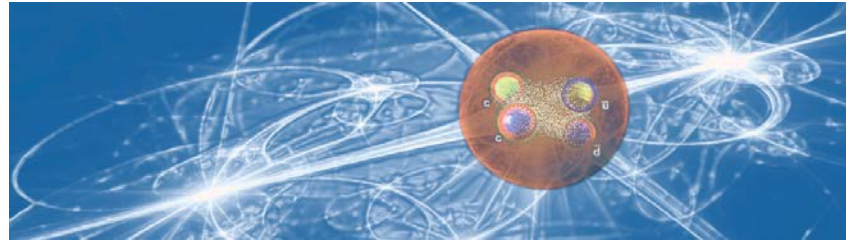


RECENT RESULTS FROM THE LATTICE ON STRONG-INTERACTION-STABLE HEAVY TETRAQUARKS



IP2I WORKSHOP ON DOUBLE CHARM TETRAQUARK AND OTHER EXOTICS

Lyon, November 22-23, 2021

KM, B. Colquhoun, A. Francis, R. Hudspith, R. Lewis

(C)FHLM: PRD102 (2020) 114506 (+earlier, in progress) [Francis, Hudspith, Lewis, Colquhoun +KM]

(J)MP: arXiv:2111.01147 [hep-lat] [Mathur, Padmanath]; PRD99 (2019) 034507 [+Junnarkar]

LMPW: arXiv:2108.10704 [hep-lat]; PRD100 (2019) 014503 [Leskovec, Meinel, Pflaumer, Wagner]

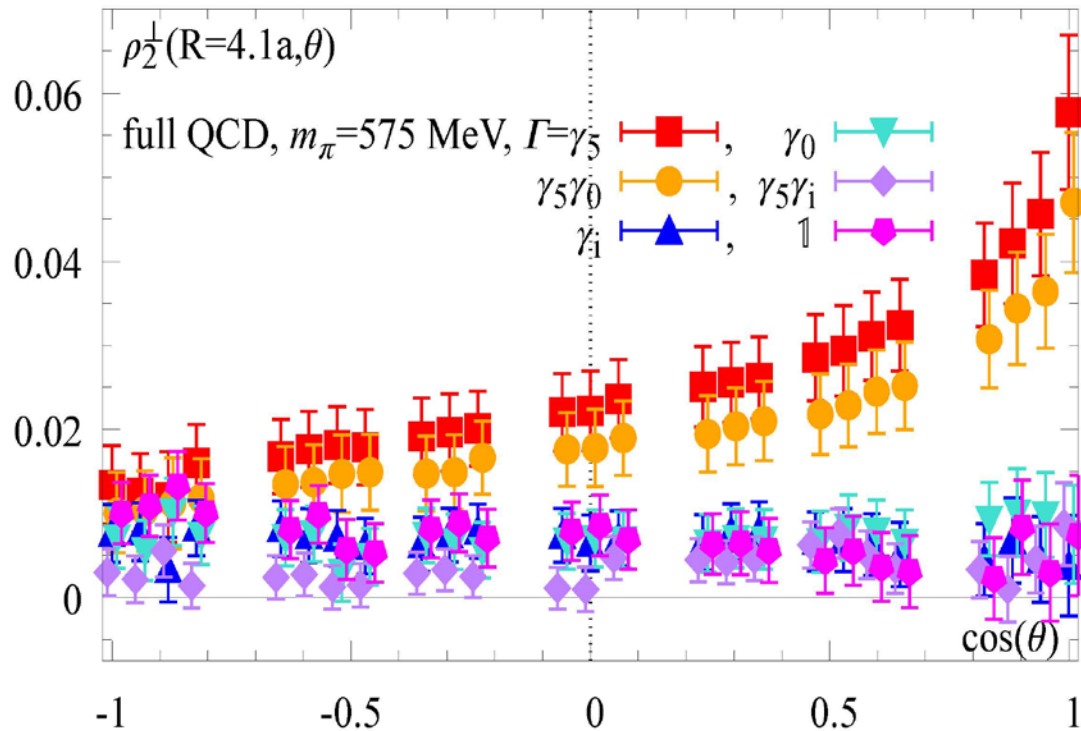
SPECIAL FEATURES OF DOUBLY HEAVY TETRAQUARK CHANNELS

- Attractive interactions for localized $\overline{Q}\overline{Q}'qq'$ configurations absent in well-separated two-heavy-meson states
 - Color Coulomb for $\overline{Q}\overline{Q}'$ in 3_c ($\Rightarrow qq'$ in $\overline{3}_c$)
 - Light quark spin-dependent for qq' in “good light diquark” ($J^P=0^+$, $\overline{3}_F, \overline{3}_c$) configuration (as per heavy baryon splittings)
- Attractive long-distance π -exchange in $V_H P_H, J^P=1^+$ channels (Manohar, Wise)
- Channels favorable to possible doubly heavy tetraquark binding
 - For $Q'=Q$: qq' in good light diquark; QQ with $J_{QQ}=1$ (no spatial excitation), 3_c : **overall $J^P=1^+$, $\overline{3}_F$**
 - For $Q \neq Q'$: qq' in good light diquark; QQ' with $J_{QQ'}=0$ or 1 (no spatial excitation), 3_c : **overall $J^P=0^+$ or 1^+ , $\overline{3}_F$**

A FEW ADDITIONAL DETAILS/OBSERVATIONS

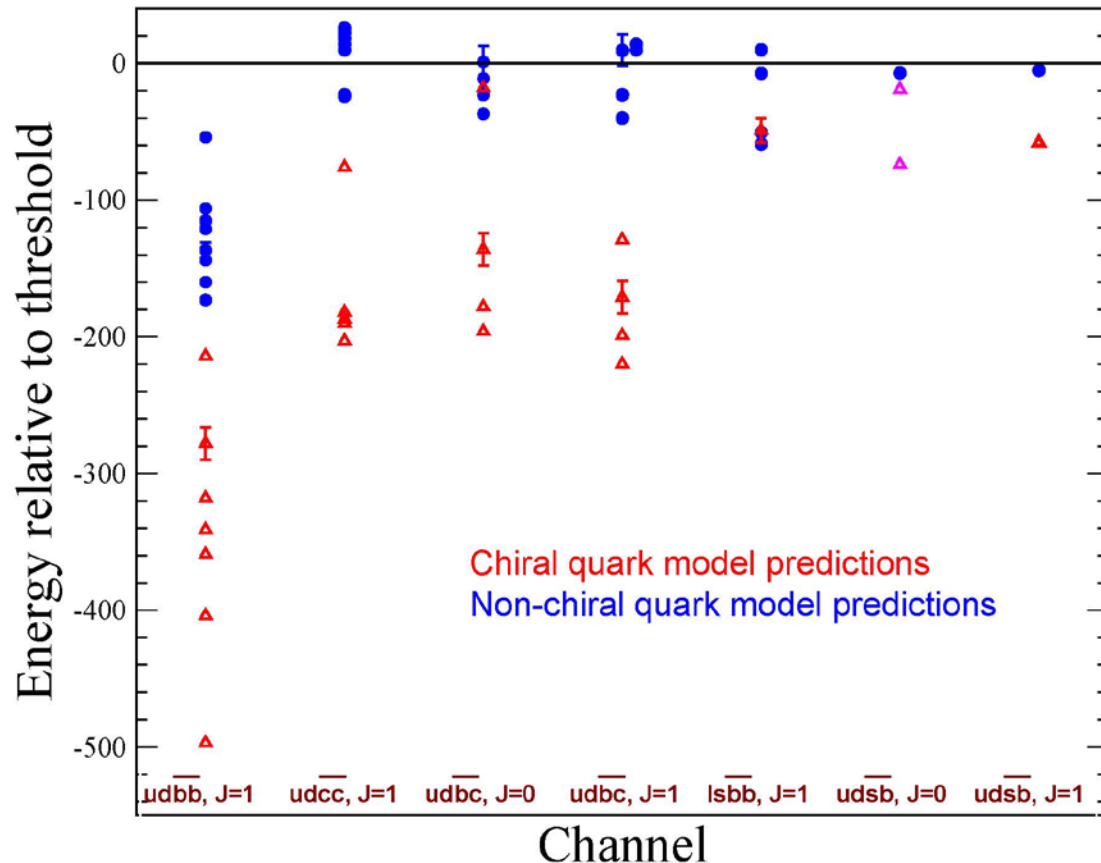
“Reality” of the good light diquark attraction: light-quark spatial correlations from the lattice

arXiv: 2106.09080



- **Colour Coulomb attraction proportional to $\overline{QQ'}$ reduced mass**
 - greatest for bb, then bc, then cc
 - bound $J^P=1^+$ doubly heavy QQ tetraquark in QCD for sufficiently heavy Q
- **Heavy baryon splittings $\Rightarrow$ increasing good light diquark attraction with decreasing m_q**
 - light m_π , short extrapolation to physical m_π important on lattice
 - potentially relevant for lattice if state weakly bound at physical m_π
- $\Sigma_b - \Lambda_b > \Sigma_c - \Lambda_c \Rightarrow$ residual $1/m_h$ -suppressed heavy-light interaction repulsive

EXTENSIVE PAST INVESTIGATIONS USING QCD-INSPIRED MODELS (since mid-80s; many experts at this workshop)



- Model parameters fixed from fits to ordinary meson, baryon spectrum
- $\Rightarrow$ model $qq\bar{6}_c, q\bar{q}8_c$ interactions (present in multiquark sector) entirely unconstrained
- differing model spin-flavor-color dependences $\Rightarrow$ different tetraquark results (e.g., figure)
 - $\triangleright$ chiral quark models (flavor dependent, color-independent spin-spin interactions)
 - $\triangleright$ c.f. non-chiral quark models (OGE-like color-dependent spin-spin interactions)
- **Experimental/lattice tetraquark channel results can rule out one (or both)**

DOUBLY HEAVY TETRAQUARKS ON THE LATTICE

- **Generic correlator (matrix) in a given channel**

$$C_{\mathcal{O}_1\mathcal{O}_2}(p,t) = \sum_x e^{ip \cdot x} \langle \mathcal{O}_1(x,t) \mathcal{O}_2(0,0)^\dagger \rangle ,$$

$$= \sum_n \langle 0 | \mathcal{O}_1 | n \rangle \langle n | \mathcal{O}_2 | 0 \rangle e^{-E_n(p)t}$$

- **$\{\mathcal{O}_k\}$: operators with channel quantum numbers**
e.g. for $Q=Q'=b$, $qq'=ud$, $J^P=1^+$: “diquark-antidiquark”,
“meson-meson” (“ $B B^*$ ”), ...

$$D(x) = (u_a^\alpha(x))^T (C \gamma_5)^{\alpha\beta} q_b^\beta(x)$$

$$\times \bar{b}_a^\kappa(x) (C \gamma_i)^{\kappa\rho} (\bar{b}_b^\rho(x))^T$$

$$M(x) = \bar{b}_a^\alpha(x) \gamma_5^{\alpha\beta} u_a^\beta(x) \bar{b}_b^\kappa(x) \gamma_i^{\kappa\rho} d_b^\rho(x)$$

$$- \bar{b}_a^\alpha(x) \gamma_5^{\alpha\beta} d_a^\beta(x) \bar{b}_b^\kappa(x) \gamma_i^{\kappa\rho} u_b^\rho(x)$$

- **Zero-momentum projection, $C(t) \sim \exp(-M_0 t)$ at large Euclidean time, t**

Generalities/practicalities

- Exponential signal decay $\leftrightarrow$ growth of noise at large t
- Effective mass: $\log[C(t)/C(t+a)] \rightarrow (aM_0)$ at large t
- **Goal: ground-state effective mass plateaus before signal overwhelmed by noise**
- Overlaps operator-dependent, hence want operator(s) with
 - good ground-state overlap
 - reduced excited-state overlap
- **Smearing of quark fields (e.g. at source) improves ground state overlap** (our results: Coulomb gauge-fixed wall sources, local (point) or “box” sinks $\leftrightarrow$ $C(t)$ from wall-local or wall-box propagators)

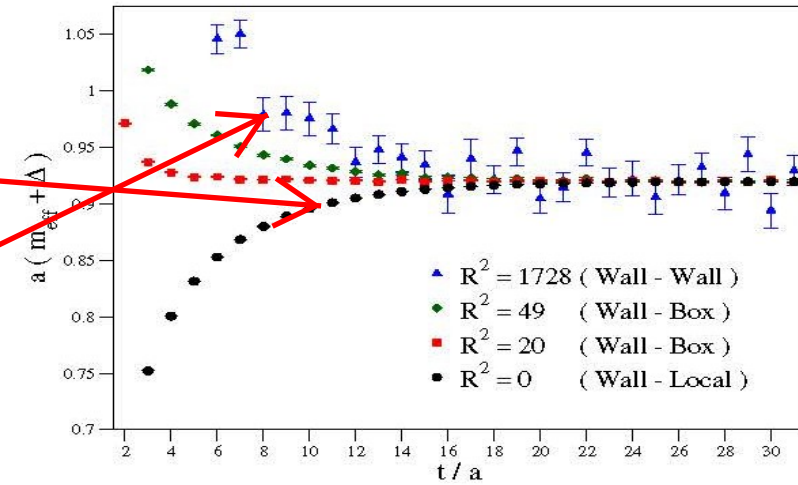
BOX SINK IMPROVEMENT

- **Wall-local propagators** (sink location (x,t)): $S_{WL}(x,t)$
 - Typically: approach ground state plateaus from below (negative excited-state contamination)
 - Late t plateaus in tetraquark channels

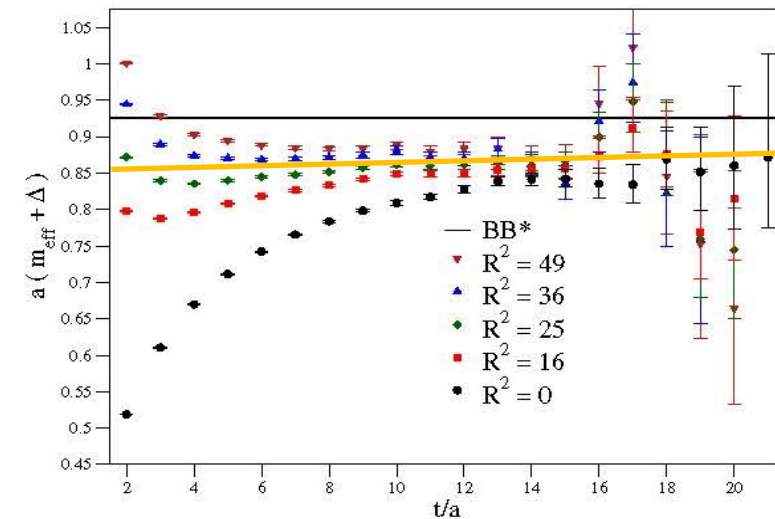
- **Wall-wall propagators**: $S_{WW}(t) = \sum_x S_{WL}(x, t)$
 - Advantage: ground state plateaus for diagonal correlator elements approached from above (positive excited state contamination)
 - Disadvantage: very noisy

- **Wall-box propagators**: intermediate between WL, WW: $S_{WB}^R(x,t) = \frac{1}{N} \sum_{\{r\} < R} S_{WL}(x + r, t)$. Expect
 - Intermediate behavior/reduced excited state contamination
 - Optimal improvement for $R \approx$ ground state size

- E.g., the B_c channel



- E.g., the $ud\bar{b}\bar{b}$ $I=0, J^P=1^+$ channel

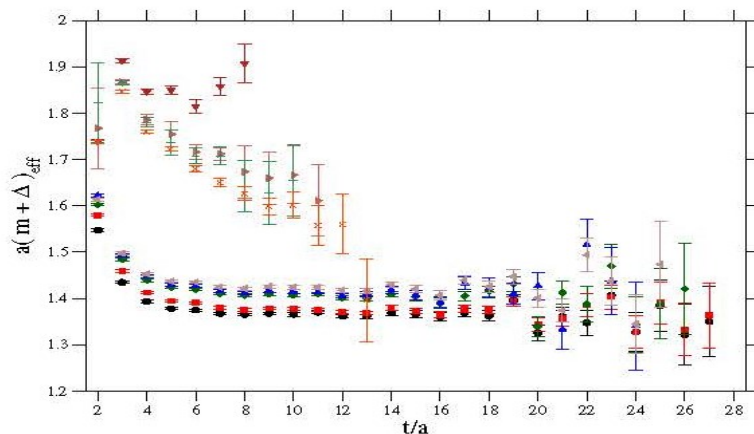


NEAR-PHYSICAL-POINT CHANNEL-SCAN FOR DEEPLY BOUND TETRAQUARK STATES

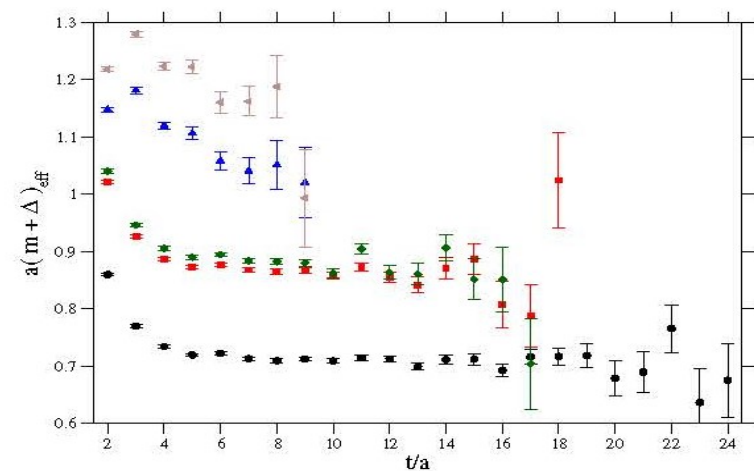
- CFHLM: PRD 102 (2020) 114506 [arXiv:2006.14294]
- Ensemble details
 - $n_f=2+1$, Wilson clover, Iwasaki gauge
 - NRQCD for b, relativistic heavy quark action for c
 - $48^3 \times 64$, $1/a=2.194(10)$ GeV ($a \sim 0.09$ fm); larger volume extension of PACS-CS $32^3 \times 64$ ensembles, same parameters
 - $m_\pi=192$ MeV, $m_\pi L = 4.2$
 - 122 configurations, 8 source locations per configuration
- **Conclusions: tetraquarks bound by more than ~ 10 -20 MeV ruled out in all channels except for $J^P=1^+$, $\bar{3}_F$ multiplet ($I=0$ $ud\bar{b}\bar{b}$ and $I=1/2$ $ls\bar{b}\bar{b}$ states)**
- more on the doubly bottom states later; first the other channels

THE $I=0$ $ud\bar{b}c$ $J^P=0^+$ and 1^+ CHANNELS

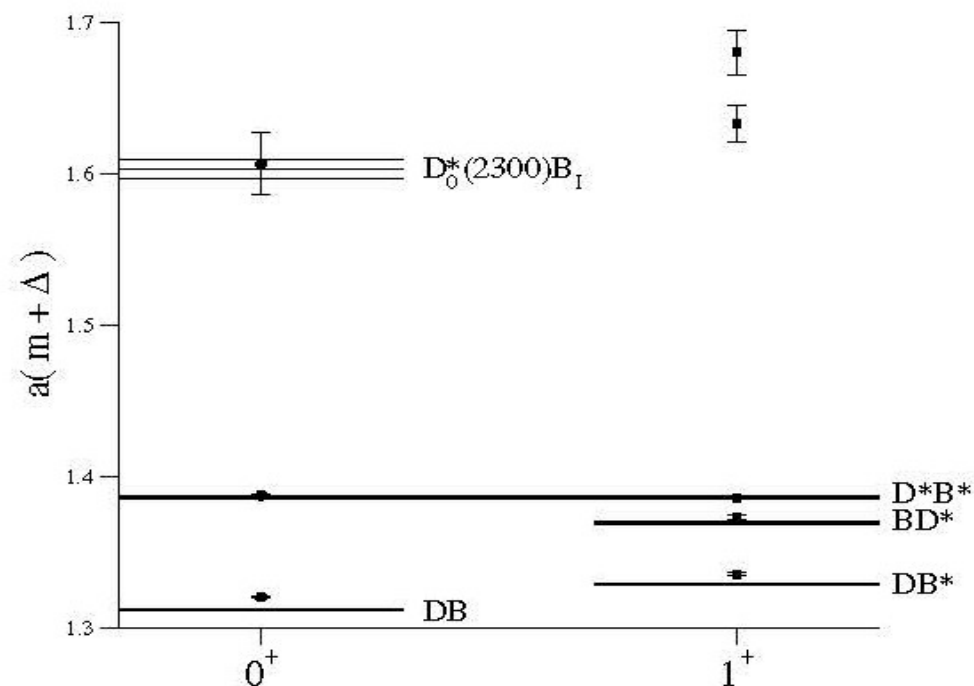
• $J^P=0^+$ effective mass



• $J^P=1^+$ effective mass



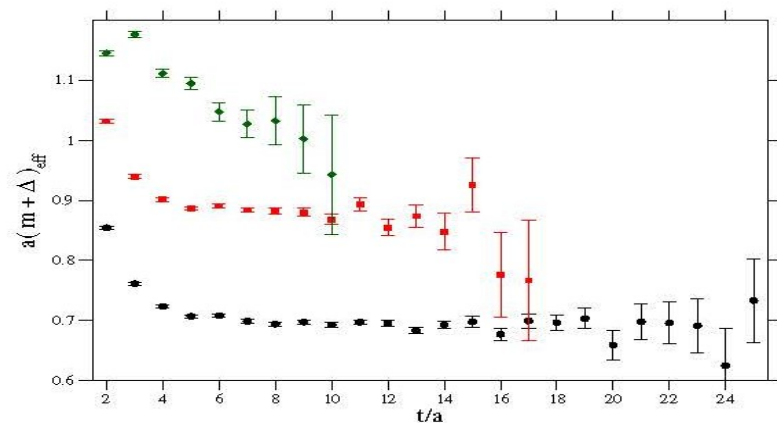
Spectra



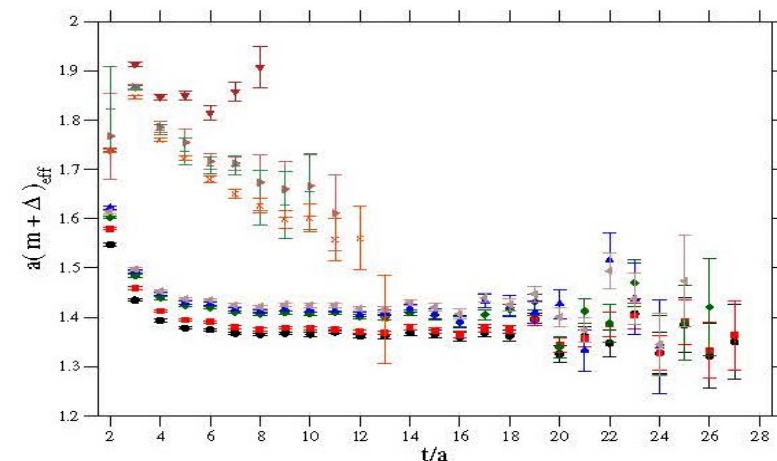
- $J^P=0^+, 1^+$ both unbound
- c.f. LMPW (preliminary): both unbound
- c.f. MP (preliminary) : $J^P=1^+$ $BE \simeq 20$ MeV

THE $I=1/2$ $Is\bar{b}c$ $J^P=0^+$ and 1^+ CHANNELS

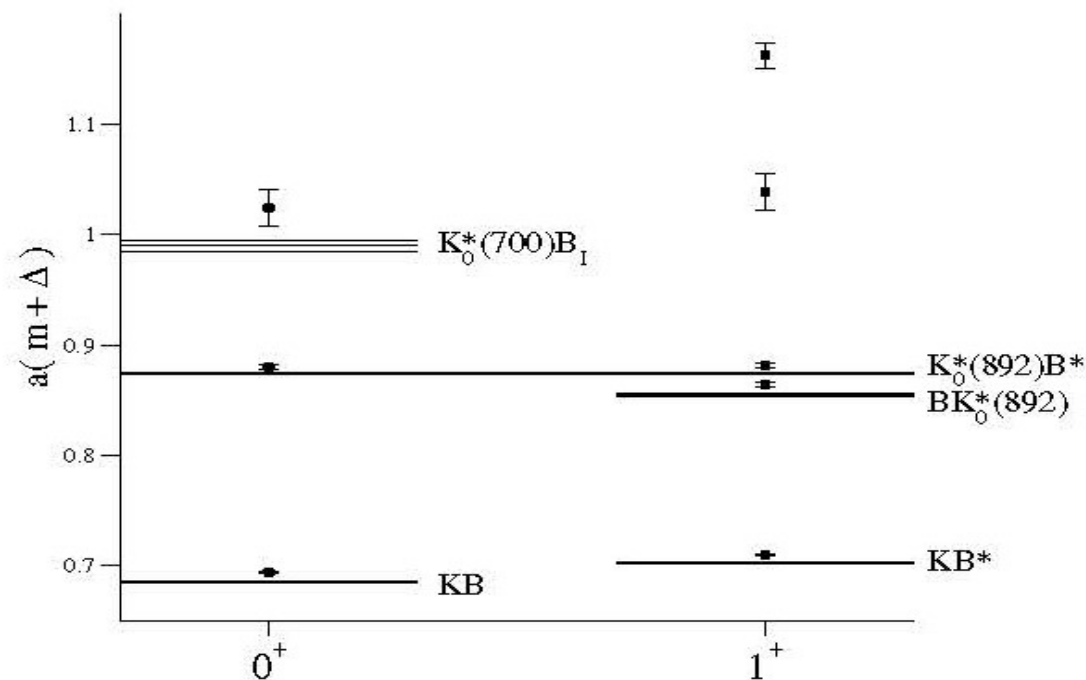
- $J^P=0^+$ effective mass



- $J^P=1^+$ effective mass



Spectra

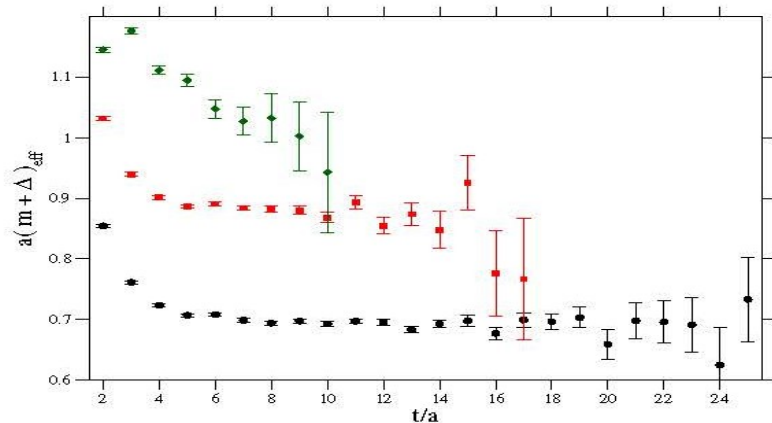


➤ $J^P=0^+, 1^+$ both unbound

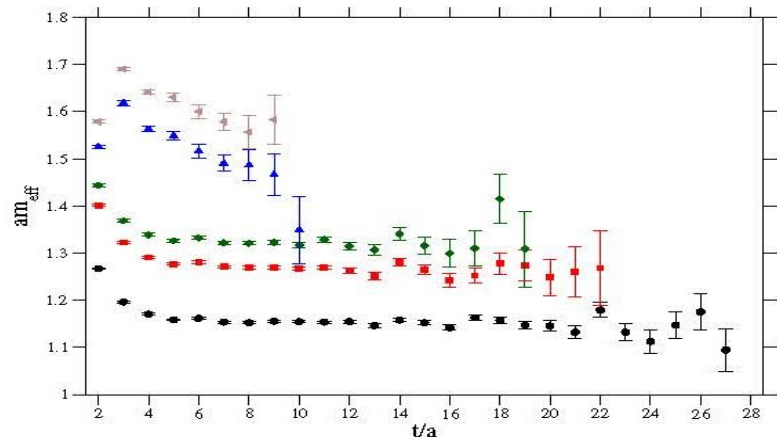
➤ c.f. MP (preliminary): $J^P=1^+$ BE $\simeq 10$ MeV

THE $I=0$ $uds\bar{b}$ $J^P=0^+$ and 1^+ CHANNELS

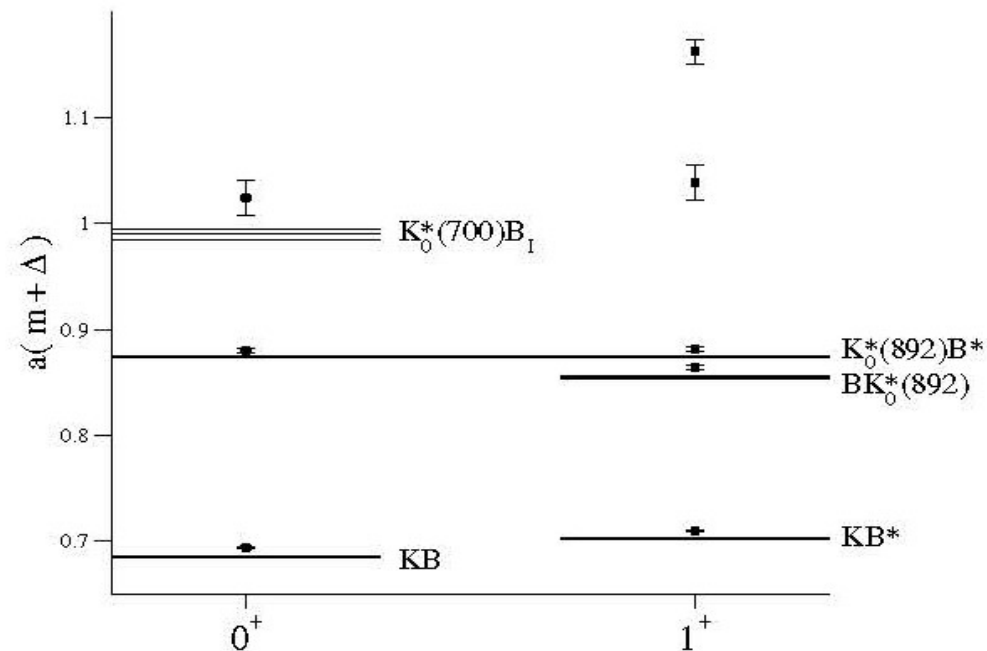
- $J^P=0^+$ effective mass



- $J^P=1^+$ effective mass



Spectra

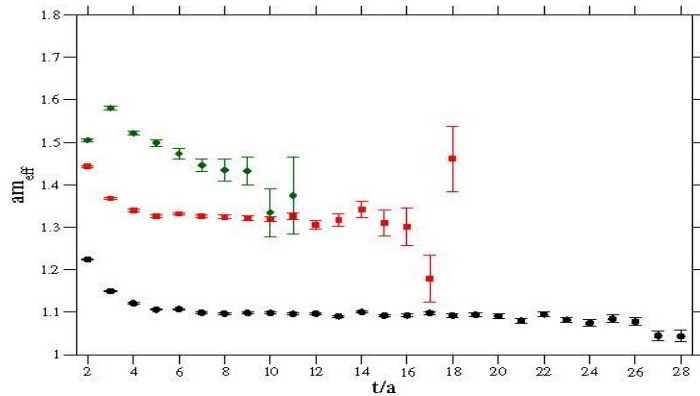


➤ $J^P=0^+, 1^+$ both unbound

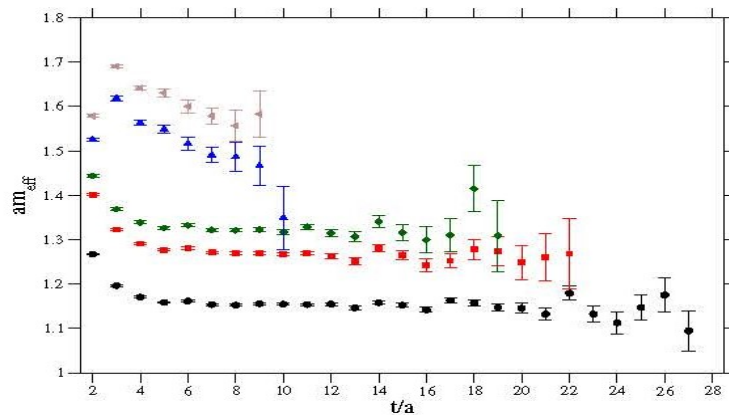
Original motivation: possible $SU(3)_F$ $uds\bar{b}$ partner of D_0 's putative $us\bar{d}\bar{b}$ $X(5568)$

THE $I=0$ $ud\bar{s}c$ $J^P=0^+$ and 1^+ CHANNELS

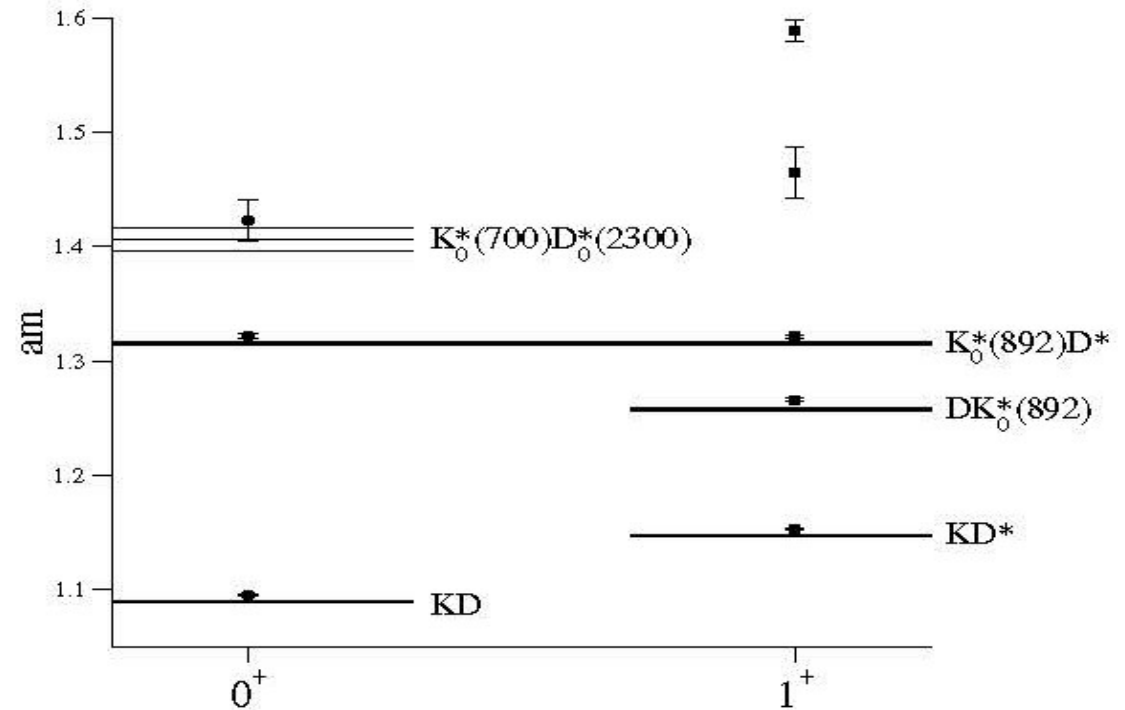
- $J^P=0^+$ effective mass



- $J^P=1^+$ effective mass



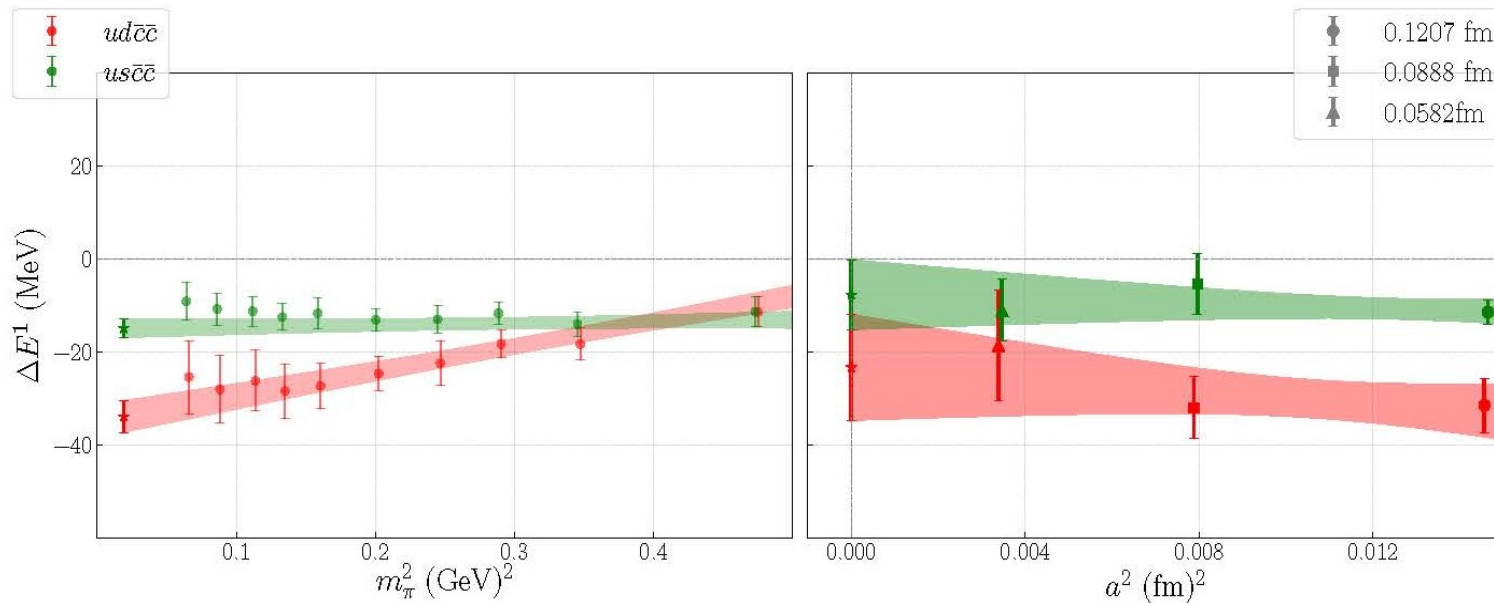
Spectra



➤ $J^P=0^+, 1^+$ both unbound

THE $J^P=1^+$, $I=0$ $ud\bar{c}\bar{c}$ CHANNEL (other lattice groups)

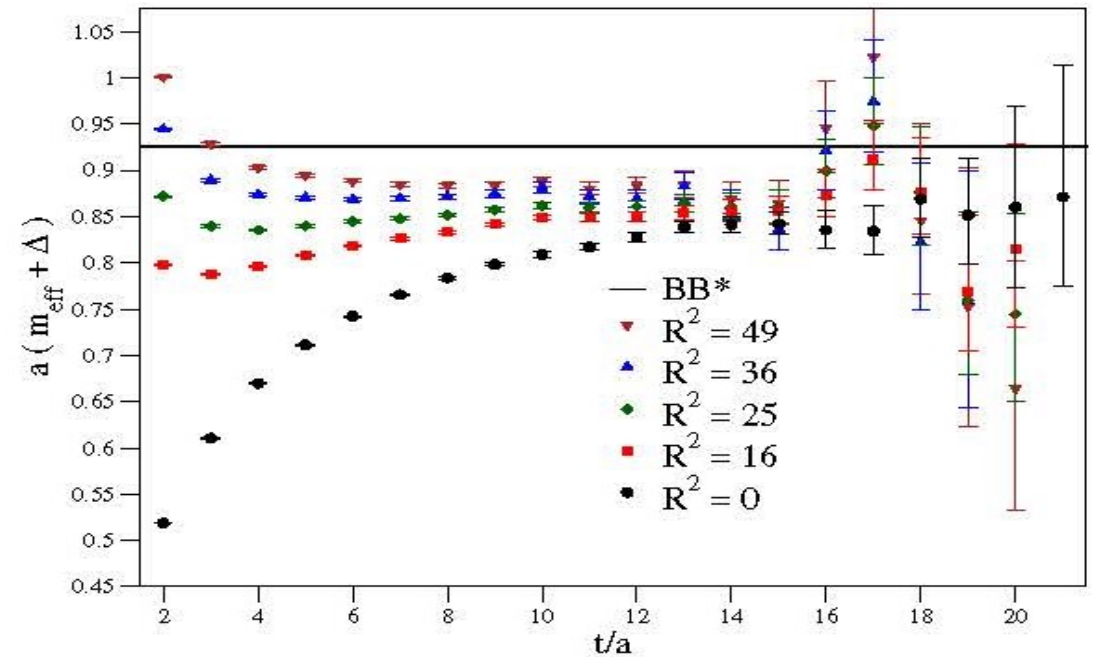
- **The LHCb experimental T_{cc} discovery (and motivation for this workshop!)**
 - very weakly bound (essentially at threshold)
 - consistent with favored $I=0$, $J^P=1^+$
- **HadSpec 2017 lattice [JHEP 11 (2017) 033; arXiv:1709.01417]:** no evidence for bound state or narrow resonance at $m_\pi=391$ MeV
- **JMP [PRD99 (2019) 034507]:** 23(11) MeV binding (also 8(8) MeV for $I=1/2$ $ls\bar{c}\bar{c}$) with caution re possible residual FV effects



THE DOUBLY BOTTOM $J^P=1^+, \bar{3}_F$ CHANNELS

- CFHLM: PRD 102 (2020) 114506+in progress: updates, improves wall-local study of PRL118 (2017) 142001)
 - PACS-CS $n_f=2+1$, Wilson clover, Iwasaki gauge, $a=0.09$ fm, $32^3 \times 48$ of earlier PRL, plus new larger-volume $48^3 \times 64$ ensembles with $m_\pi < 200$ MeV
 - Box-sink construction dramatically improves the ground state signal c.f. earlier PRL results
- Other groups
 - (J)MP: $n_f=2+1+1$ HISQ, wall-local, a -dependent minimum m_π : 153, 345, 545 MeV for $a \sim 0.12, 0.09, 0.06$ fm (near physical point for coarsest)
 - LMPW: RBC/UKQCD $n_f=2+1$ DWF, Iwasaki gauge ensembles, 2 with physical m_π ; smeared source, local sink;; 3 local operators, 2 non-local “meson-meson” for improved overlap with threshold 2-meson state

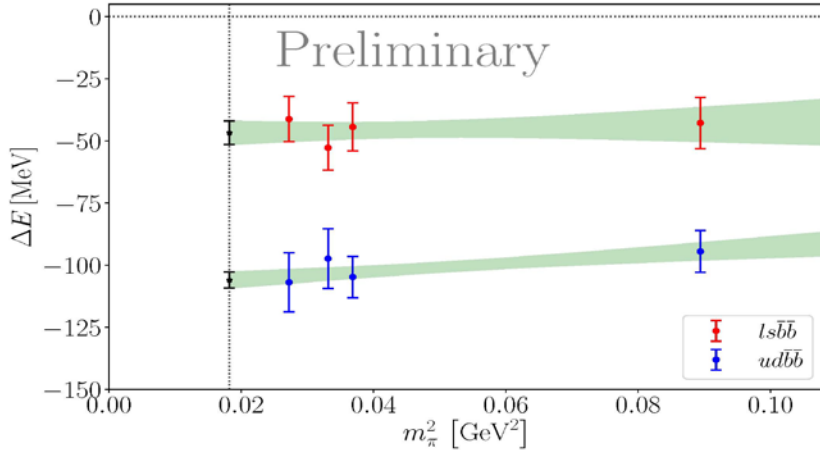
- Example of box-sink improvement for $ud\bar{b}\bar{b}$ $l=0, J^P=1^+$ channel, $m_\pi=192$ MeV



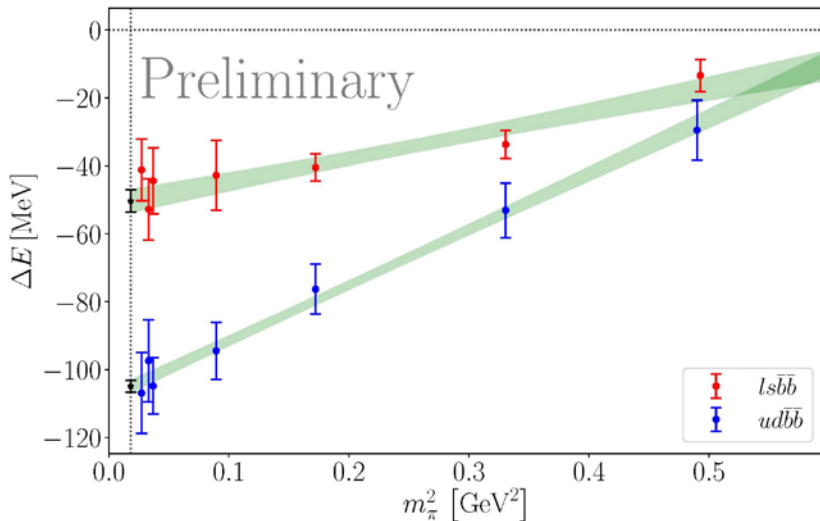
- Q: wall-local/late-plateau an issue for (J)MP?

CFHLM UPDATED $J^P=1^+$, $I=0$ $ud\bar{b}\bar{b}$, $I=1/2$ $ls\bar{b}\bar{b}$ BINDING ENERGIES

- Fits with $m_\pi < 300$ MeV



- Fits with $m_\pi < 702$ MeV



PRELIMINARY UPDATES

Errors: statistical, optimal box-sink size, R, choice only at present (i.e., preliminary)

- From $m_\pi < 300$ MeV fits

- $I=0$ $ud\bar{b}\bar{b}$: 109(5) MeV

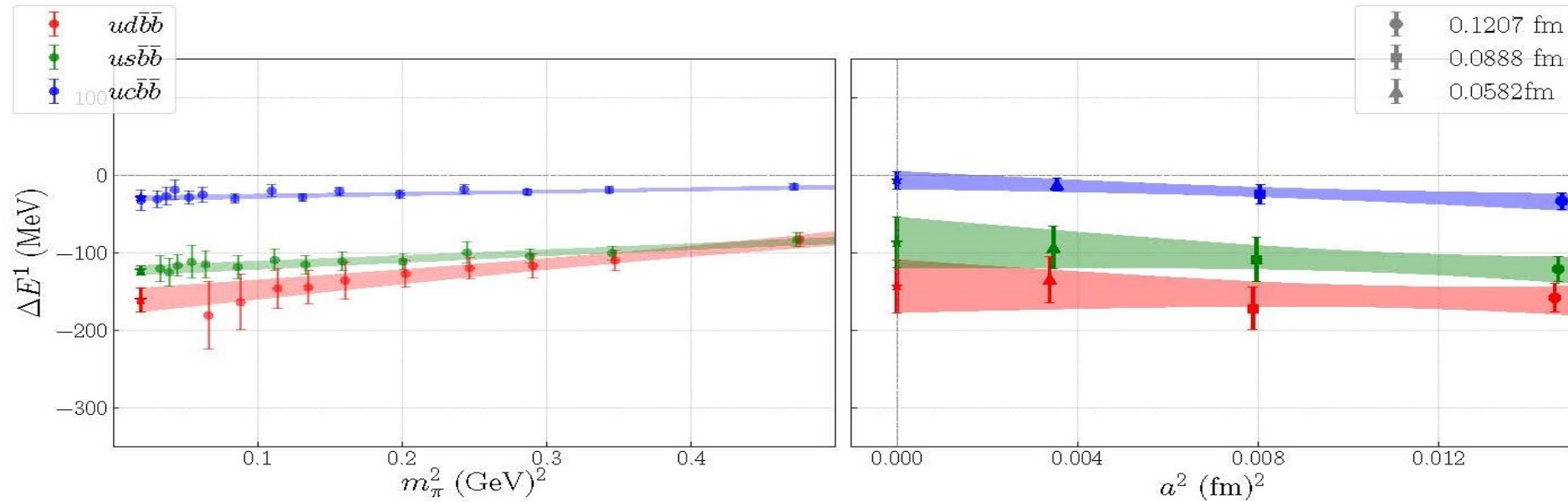
- $I=1/2$ $ls\bar{b}\bar{b}$: 48(7) MeV

- Extending fits to $m_\pi < 702$ MeV

- $I=0$ $ud\bar{b}\bar{b}$: 109 → 108 MeV

- $I=1/2$ $ls\bar{b}\bar{b}$: 48 → 52 MeV

JMP $J^P=1^+$, $I=0$ $ud\bar{b}\bar{b}$, $I=1/2$ $ls\bar{b}\bar{b}$ BINDING ENERGIES

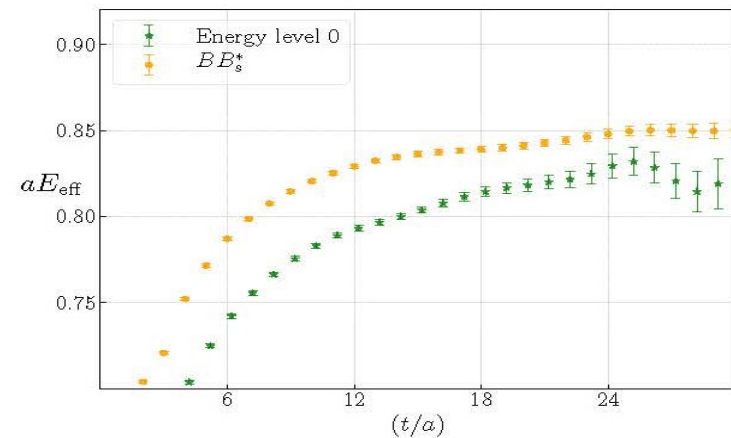


- Binding energies**

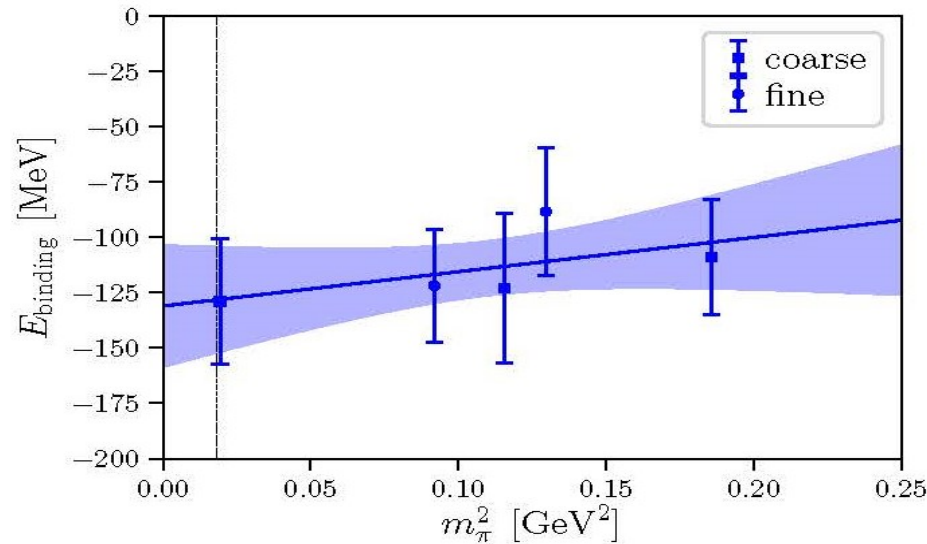
- $I=0$ $ud\bar{b}\bar{b}$: 143(34) MeV

- $I=1/2$ $ls\bar{b}\bar{b}$: 87(32) MeV

- Effective mass plateaus? E.g. $m_\pi=575$ MeV $ls\bar{b}\bar{b}$ example**



LMPW $J^P=1^+$, $I=0$ $ud\bar{b}\bar{b}$, $I=1/2$ $ls\bar{b}\bar{b}$ BINDING ENERGIES

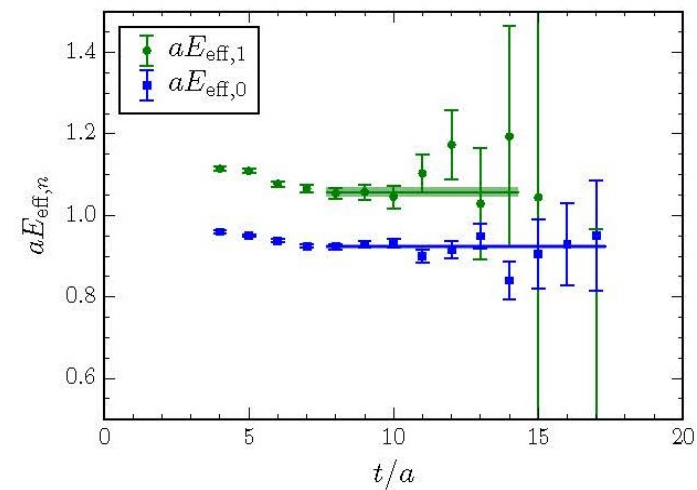


$I=0$ $ud\bar{b}\bar{b}$ physical
point extrapolation

- Binding energies

- $I=0$ $ud\bar{b}\bar{b}$: 128(26) MeV (PRD100 (2019) 014503)
- $I=1/2$ $ls\bar{b}\bar{b}$: ~ 80 MeV, ~ 25 MeV uncertainty (preliminary Lattice 2021 report, arXiv:2108.10704)

- $I=0$ $ud\bar{b}\bar{b}$ $m_\pi=340$ MeV effective mass e.g.



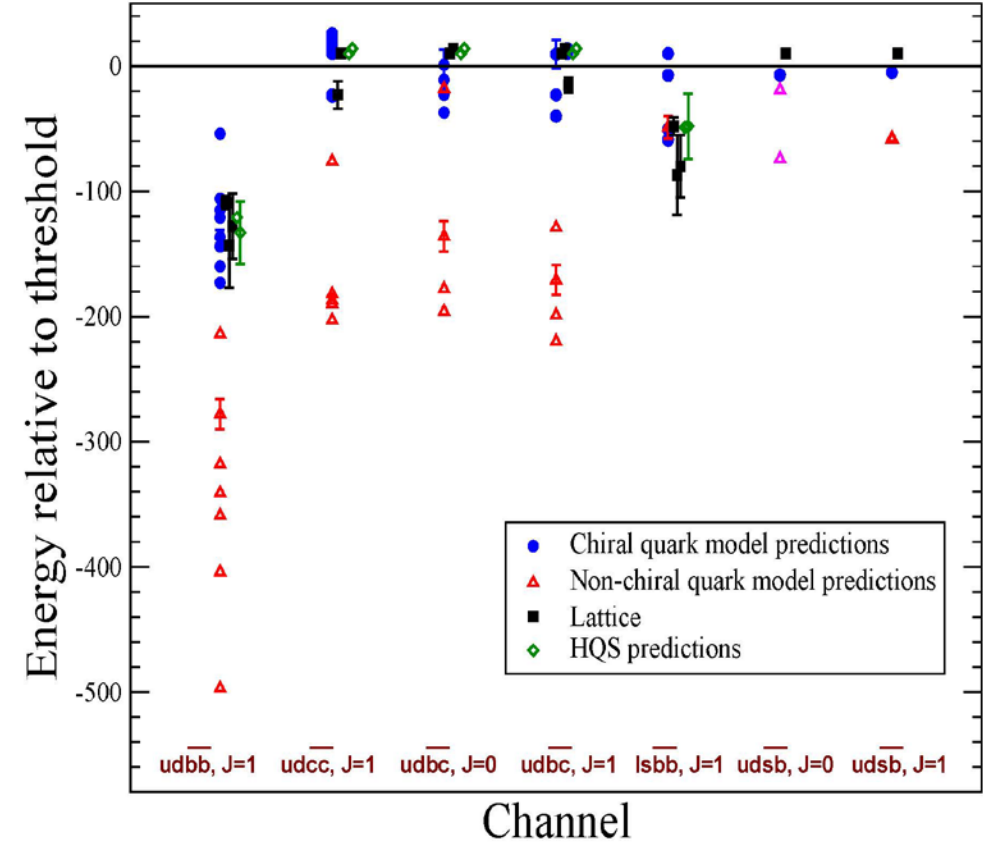
CONCLUSIONS/EXPECTATIONS/SPECULATIONS (1)

- Lattice tetraquark binding energy summary

- Clear evidence for strong-interaction-stable $J^P=1^+$, $I=0$ $ud\bar{b}\bar{b}$ state: bound by >100 MeV, hence with only weak decays (experimental detection challenge)
- Clear evidence for strong-interaction-stable $J^P=1^+$, $I=1/2$ $ls\bar{b}\bar{b}$
 - ❖ **possibility of < 45 MeV binding, $\sim 100\%$ BF EM decay to $BB_s\gamma$** (CFHLM lattice, Eichten-Quigg/Braaten-He-Mohapatra HQS results)
 - ❖ binding > 45 MeV however also possible (JMP, preliminary LMPW): weak decays only, multiple modes, small BFs
- No states with several x 10 MeV binding in other channels in CFHLM study

CONCLUSIONS/EXPECTATIONS/SPECULATIONS (2)

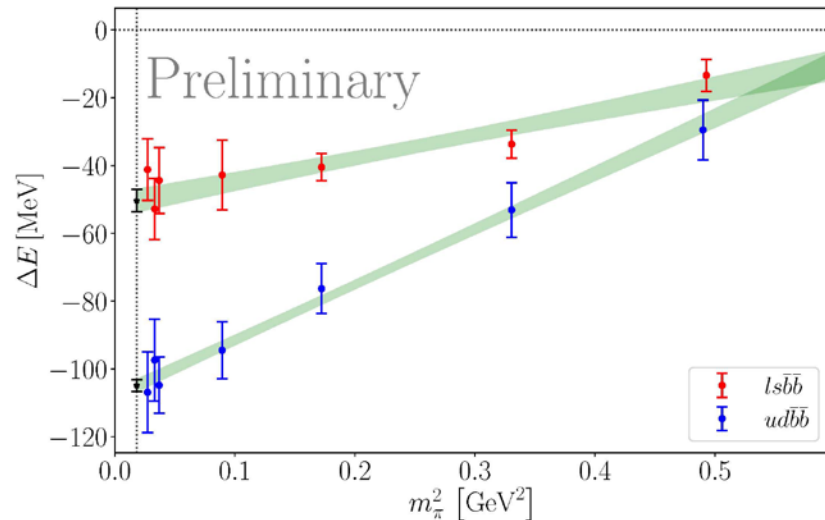
- **Lattice/LHCb T_{cc} : tests of model predictions**
 - Chiral quark models: predictions incompatible with lattice (and HQS) results/overbind T_{cc} c.f. LHCb
 - Non-chiral quark models (OGE-like spin-spin interactions): compatible with weakly bound T_{cc} , predictions in other channels in reasonable agreement with lattice (and HQS)
- **Conclusion: chiral quark model picture of effective low-energy interactions (in current forms) ruled out by multi-quark sector tests**



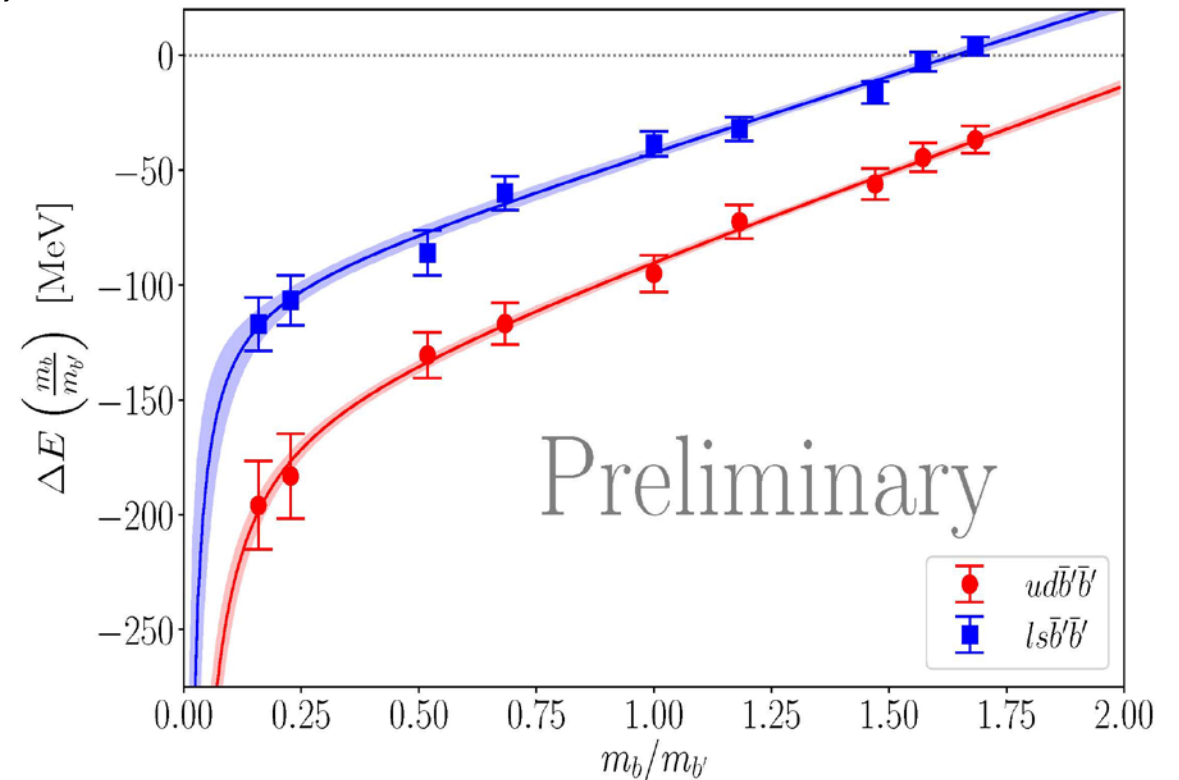
CONCLUSIONS/EXPECTATIONS/SPECULATIONS (3)

- Key physical mechanisms at play in doubly heavy tetraquark binding

- Heavy-heavy color Coulomb attraction: as expected + evidence from variable b' mass study (figure)
- Spin-dependent good-light-diquark attraction (increasing binding with decreasing light-quark mass)



- Variable b mass study: the heavy-heavy color Coulomb ($m_\pi=299 \text{ MeV}$)



CONCLUSIONS/EXPECTATIONS/SPECULATIONS (4)

- **A possible $I=0$, $J^P=1^+$ $ud\bar{b}\bar{c}$ state?**
 - very weak $I=0$, $J^P=1^+$ $ud\bar{c}\bar{c}$, T_{cc} binding suggests molecular-type component, importance of long-distance π -exchange force
 - same π -exchange force operative in analogous $I=0$, $J^P=1^+$ $ud\bar{b}\bar{c}$ channel
 - in addition, short-range attraction likely stronger for $ud\bar{b}\bar{c}$ than $ud\bar{c}\bar{c}$ (more attractive color Coulomb, similar good-light-diquark attraction)
 - CFHLM, LMPW, MP lattice studies all concur: no deep binding
 - suggests likelihood of $I=0$, $J^P=1^+$ T_{bc} analogue of T_{cc} bound by < 45 MeV, hence EM decay, likely to $BD\gamma$ (unless weakly bound $I=0$, $J^P=0^+$ tetraquark also exists)
 - **experimental search for correlated BD feasible?**

Thanks for your attention

in memory of Gabriel Karl and Simon Eidelman

BACKUP SLIDES (1)

EXPECTATIONS FROM HEAVY BARYON SPLITTINGS

- Observed splittings

- $\Sigma_b - \Lambda_b = 194 \text{ MeV}$

- $\Sigma_c - \Lambda_c = 167 \text{ MeV}$

- $\Xi_b' - \Xi_b = 142 \text{ MeV}$

- $\Xi_c' - \Xi_c = 109 \text{ MeV}$

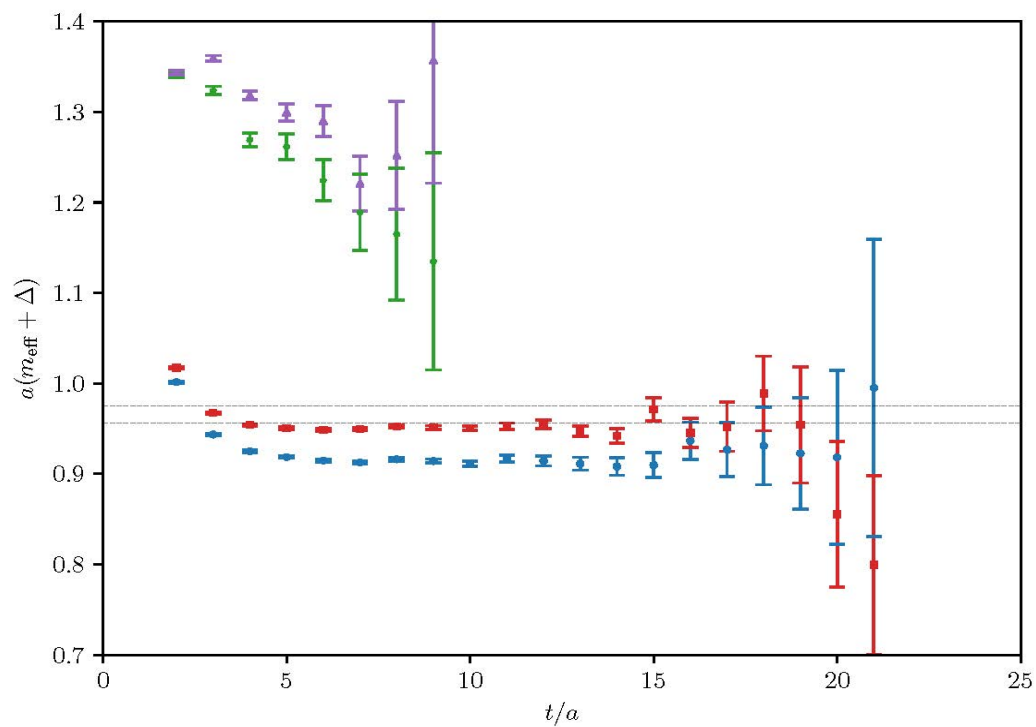
- good ud , ℓs diquark attraction relative to corresponding spin averages: ~ 145 , $105 \text{ MeV} \Rightarrow$ *increased* attraction with *decreased* m_q

- $h=c$ vs. $h=b$: residual ($\propto 1/m_h$) light-heavy interactions eat into good diquark attraction with decreasing m_h

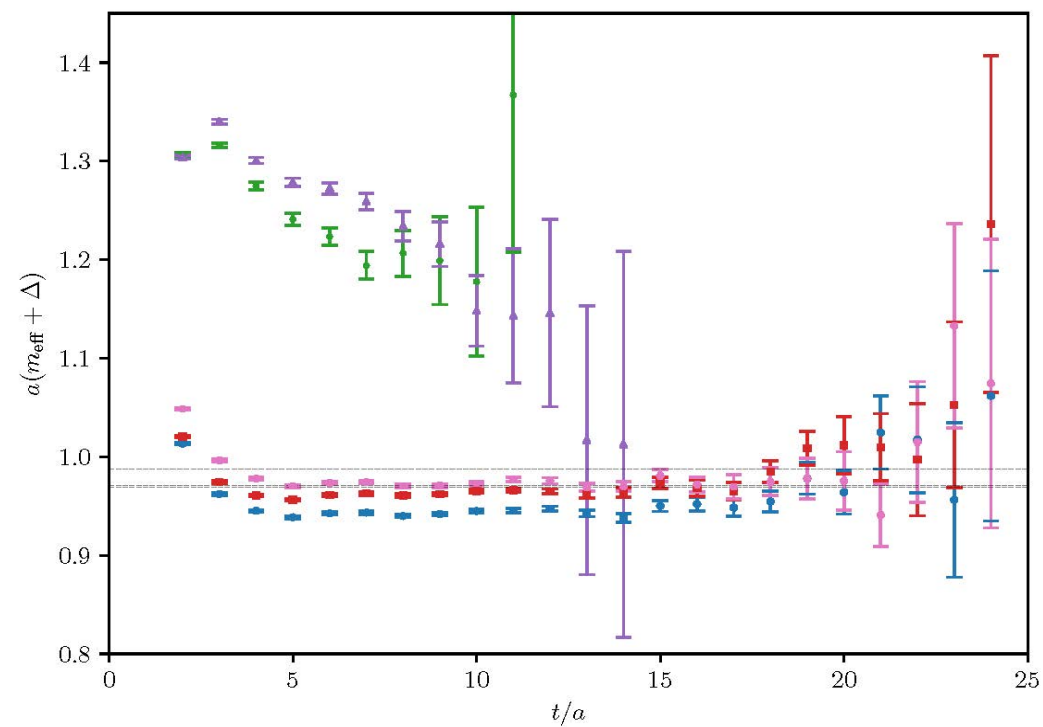
BACKUP SLIDES (2a)

- CFHLM effective mass plot examples
 - PACS-CS E5 ensemble ($m_\pi=299$ MeV)

$I=0, J^P=1^+ \text{ } u d \bar{b} \bar{b}$



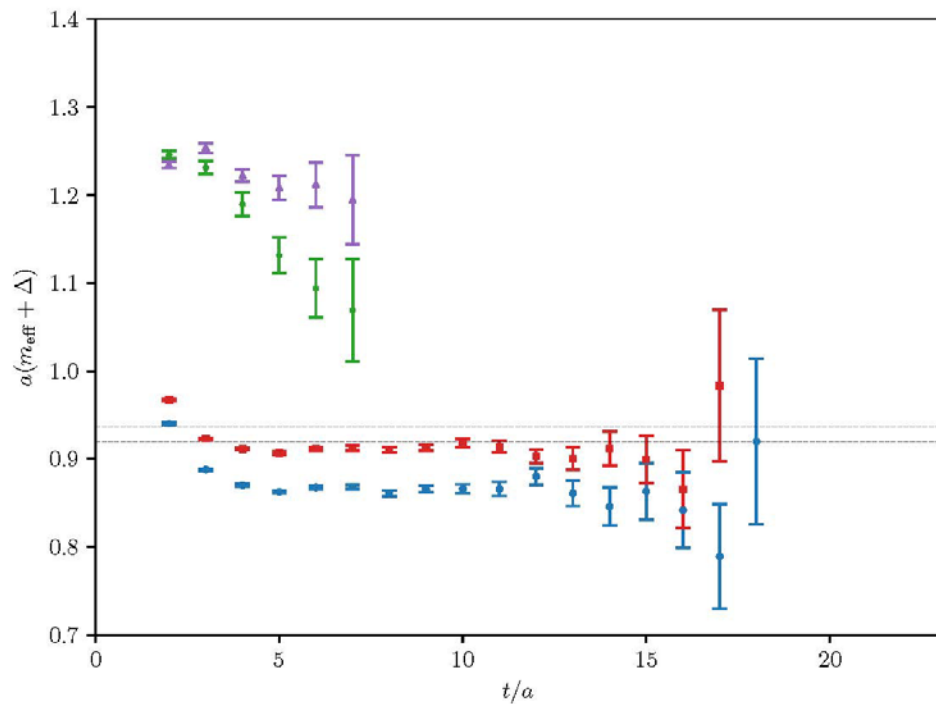
$I=0, J^P=1^+ \text{ } l s \bar{b} \bar{b}$



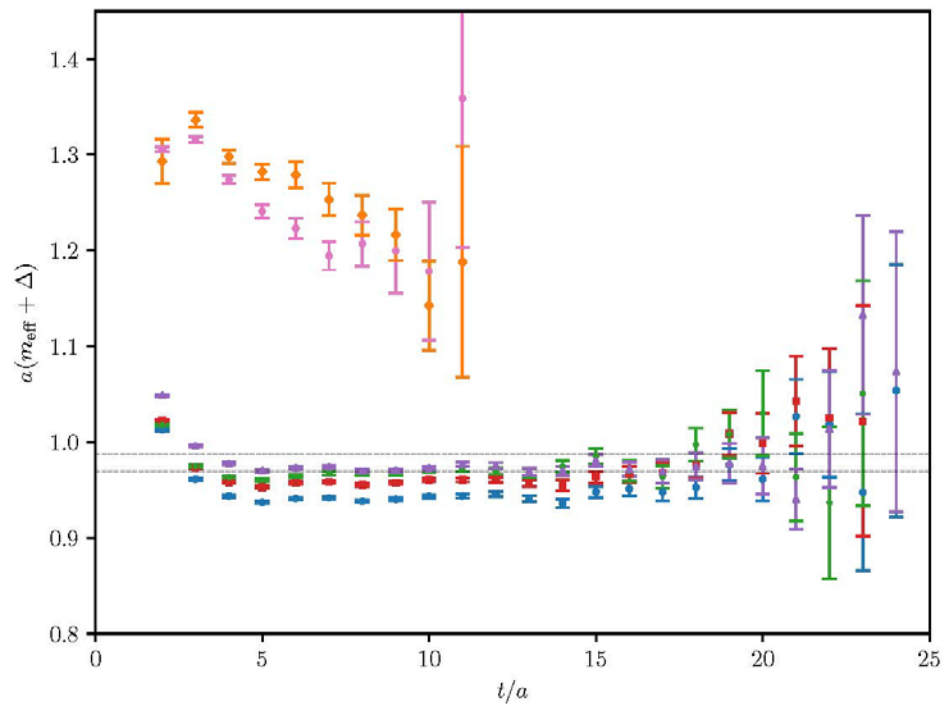
BACKUP SLIDES (2b)

- CFHLM effective mass plot examples
 - new $48^3 \times 64$ $\kappa_1=0.13779$ ensemble ($m_\pi=165$ MeV)

$I=0, J^P=1^+ \text{ } u d \bar{b} \bar{b}$



$I=0, J^P=1^+ \text{ } l s \bar{b} \bar{b}$



BACKUP SLIDES (3)

- LMPW 2019, $I=0$, $J^P=1^+$ $ud\overline{b}\overline{b}$: possible residual fit window dependence?
 - $m_\pi=340$ MeV multi-exponential fits,, variable windows, variable # of exponentials

