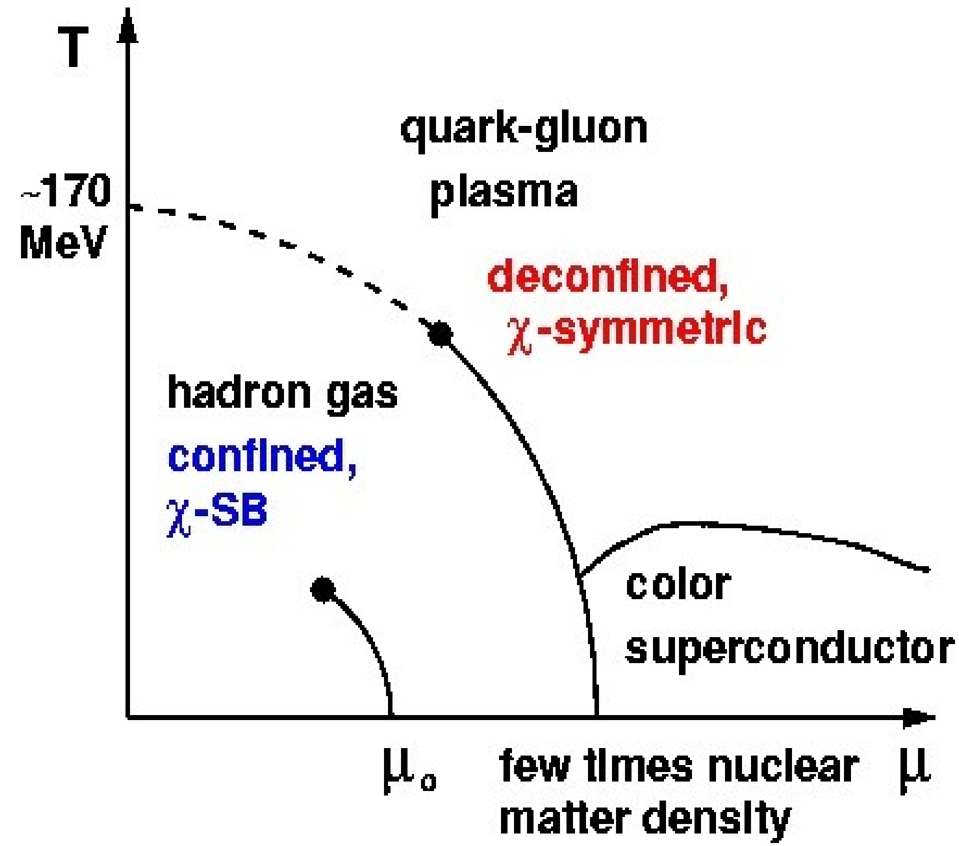
Exploring the Phase Diagram of Strongly Interacting Matter

from a wealth of experimental results
my personal view of
- highlights
- physics insights
- perspectives



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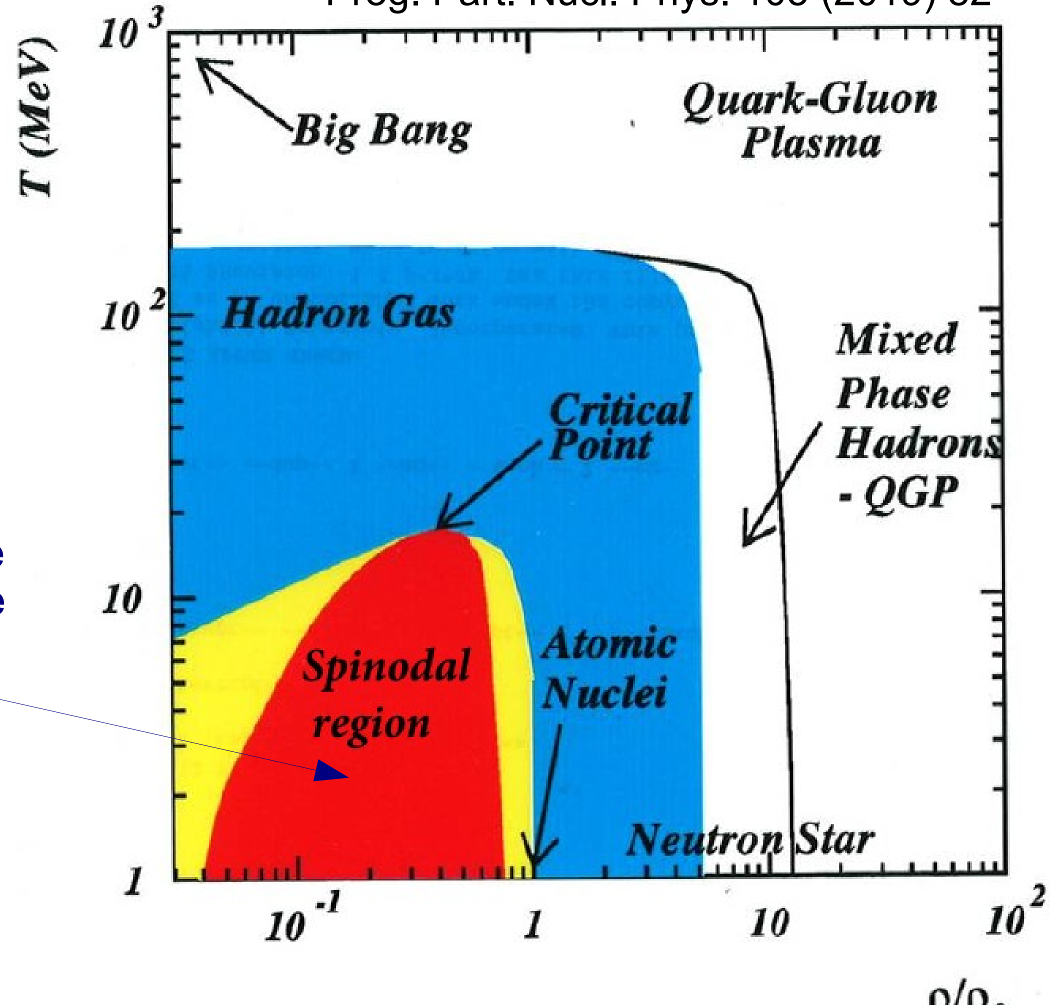
Johanna Stachel, Phys. Inst., Univ. Heidelberg
International Symposium From Particles to the Universe
Paris, December 10, 2021

the QCD phase diagram, as seen from a low energy nuclear physics viewpoint

the beginnings of research into the nuclear liquid gas phase transition:
G. Bertsch and P. J. Siemens, Nuclear Fragmentation, PLB 126 (1983) 9

"we argue that fragmentation will occur when the nuclear matter has expanded adiabatically to the onset of hydrodynamic instability $\partial P/\partial V_{|S} = 0$ "

B. Borderie, J. Frankland, Prog. Part. Nucl. Phys. 105 (2019) 82

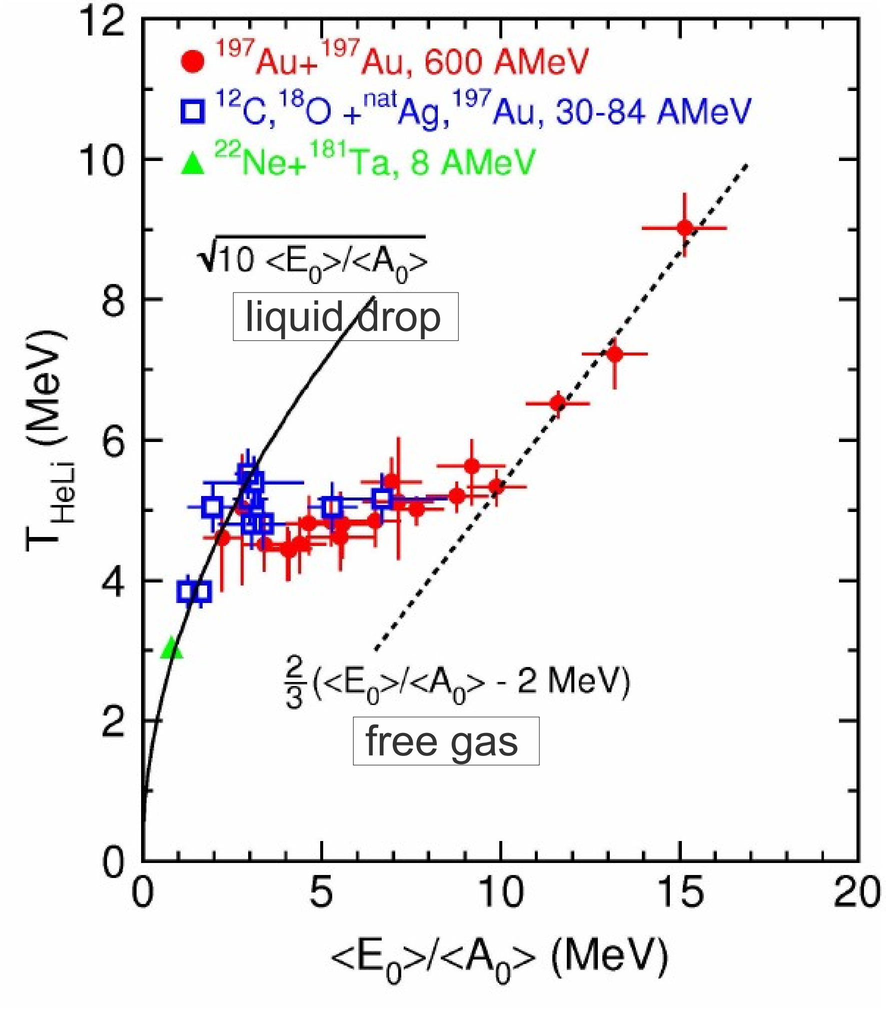


$T_{crit} \approx 16$ MeV $\rho_{crit} \approx 0.4\rho_0$

The caloric curve for the liquid gas phase transition

measurement with ALADiN

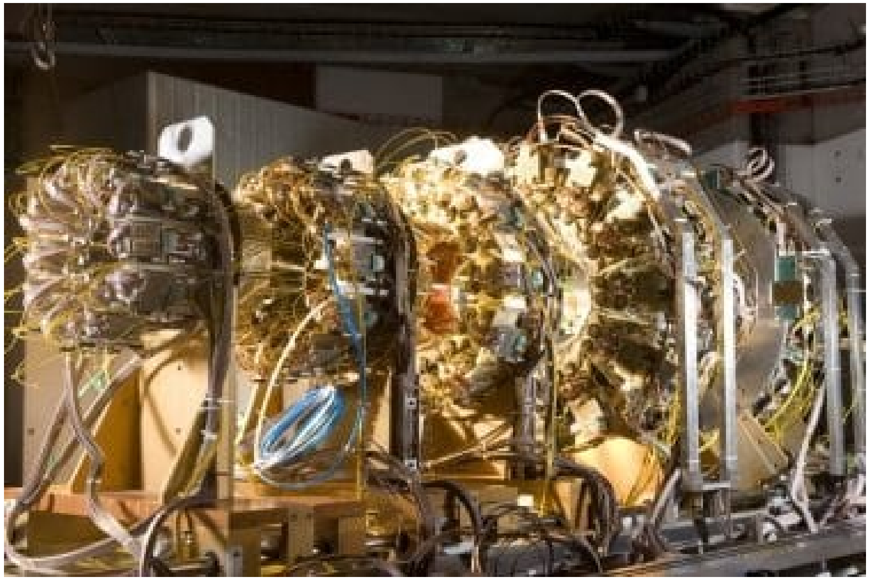
J. Pochodzalla, U. Lynen at al. Phys.Rev.Lett. 75 (1995) 1040



- temperatures from yield ratios He and Li
- excitation energies per nucleon of primary fragments from fragment and neutron distributions

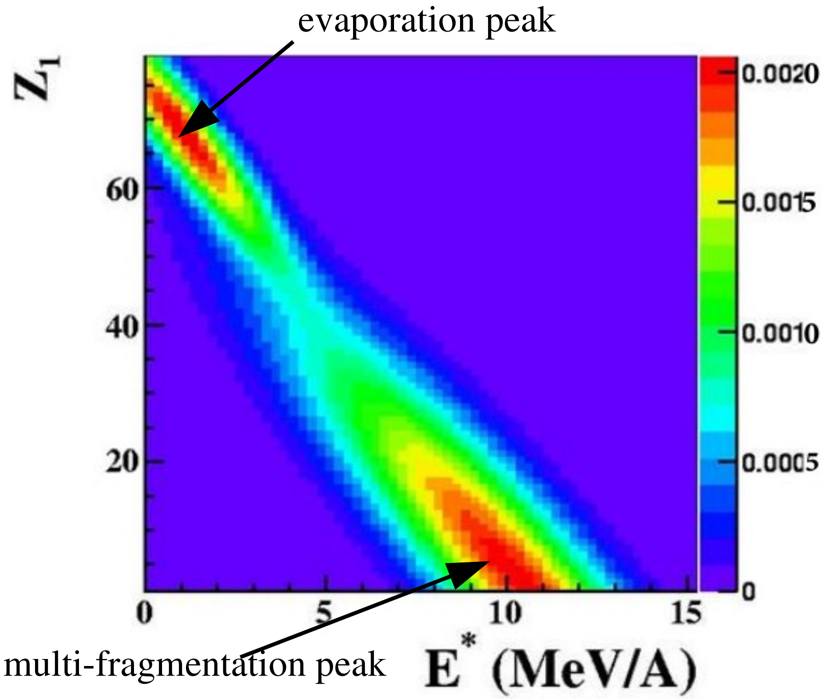
Bi-modality and latent heat of liquid-gas phase transition

a major accomplishement by the INDRA and ALADiN collaborations, data obtained in Au-Au collisions at 60-100 MeV/nucleon at GSI



bi-modality as signal for phase transition in intermediate energy nuclear collisions

proposed by: P. Chomaz, M. Colonna, J. Randrup, Phys. Rep. 389 (2004) 263



size of largest fragment serves as order parameter of 1st oder phase transition between gas and liquid phase

$$\Delta E = E_g - E_l = 8.1(\pm 0.4)_{stat}(+1.2 - 0.9)_{syst} \text{ MeV/A}$$

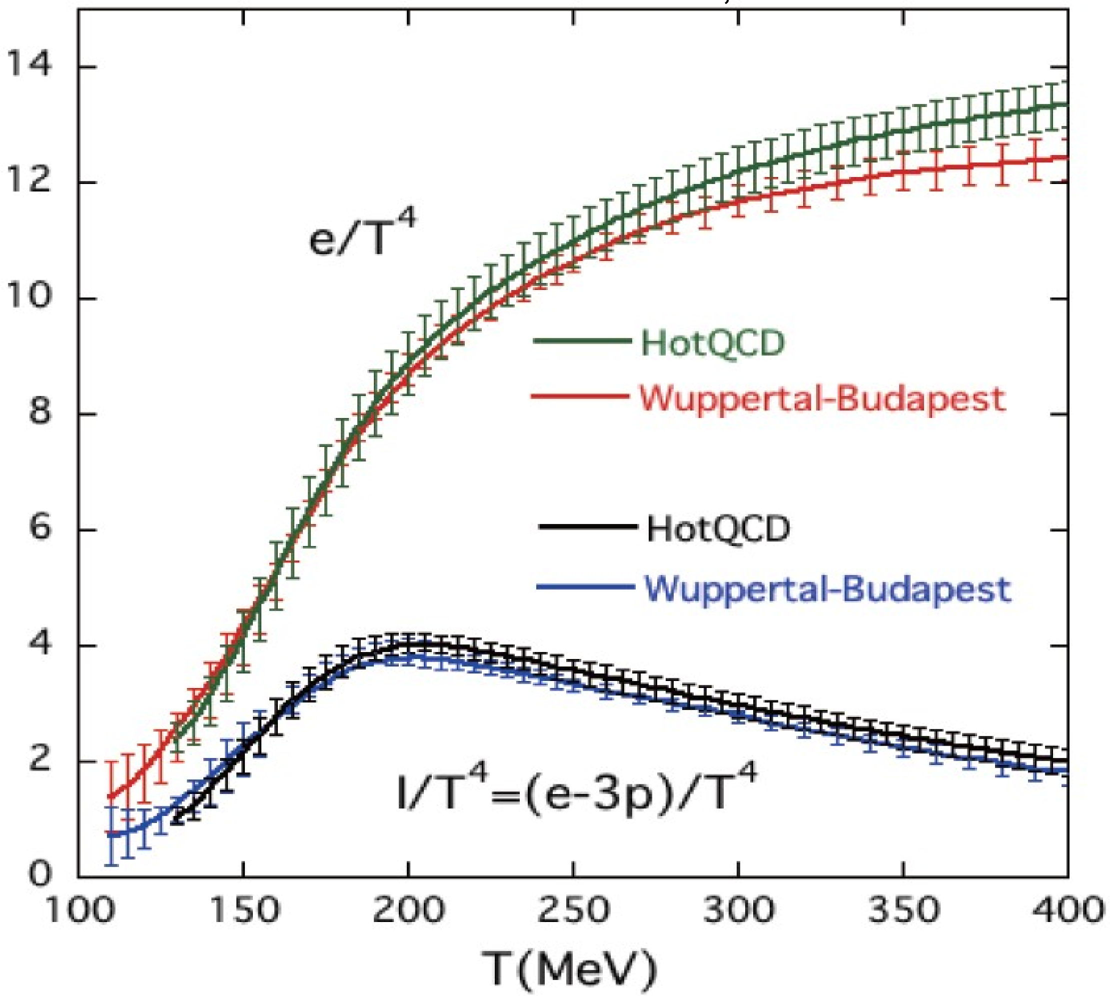
latent heat = energy difference between gas and liquid phase

E. Bonnet PRL 103 (2009) 072701

Equation of state of hot QCD matter in lattice QCD

computation of QCD EoS one of the major goals in lQCD community since 1980

A.Ukawa, arXiv:1501.04215



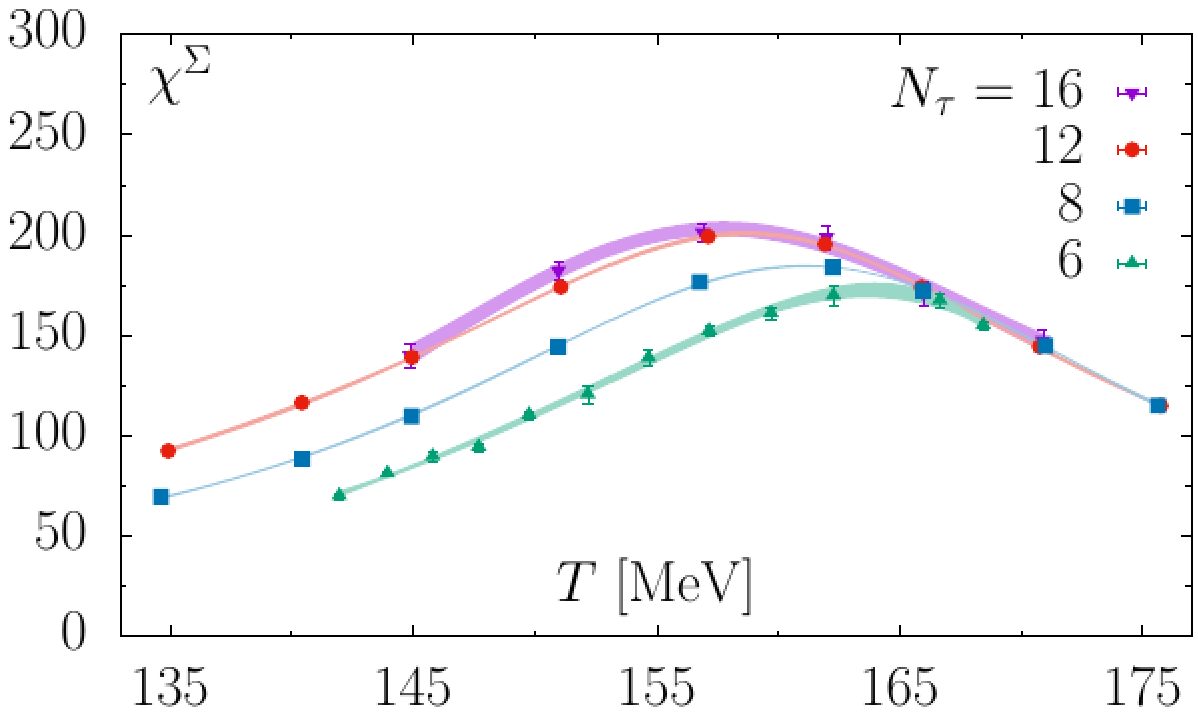
consolidated results on EoS from different groups, extrapolated to continuum and chiral limit

rapid rise of energy density (normalized to T^4 rise for relativistic gas)
- signals rapid increase in degrees of freedom due to transition from hadrons to quarks and gluons
- lQCD points to continuous cross over transition

Measure for chiral symmetry restoration in lQCD

order parameter: chiral condensate, its susceptibility peaks at T_c

S.Borsanyi et al. Wuppertal-Budapest Coll., PRL 125 (2020) 052001
A.Bazavov et al. HotQCD Coll., PLB795 (2019) 15



$$\langle \bar{\Psi}\Psi \rangle = \frac{T}{V}\frac{\partial \ln Z}{\partial m}$$

$$\chi_{\bar{\Psi}\Psi} = \frac{T}{V}\frac{\partial^2 \ln Z}{\partial m^2}$$

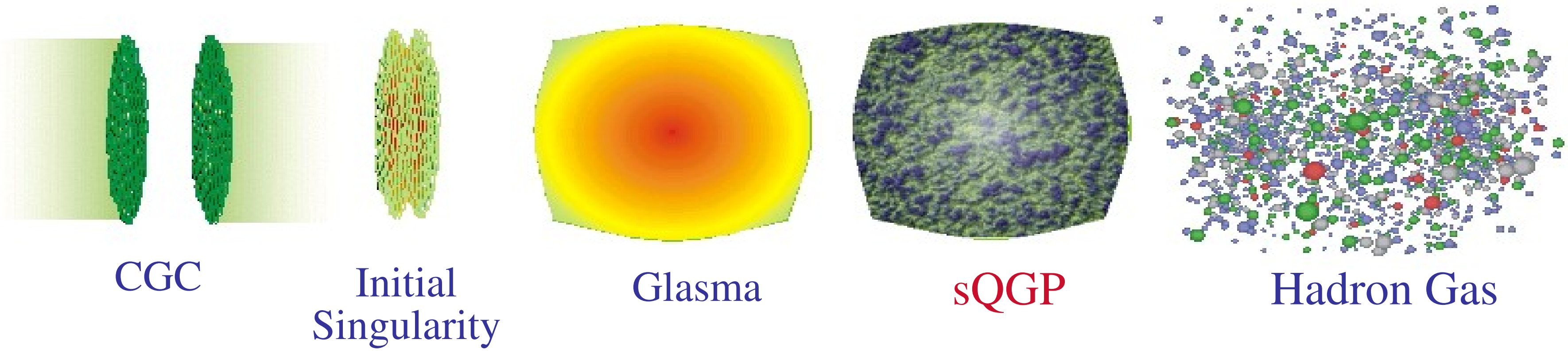
comparing different measures: pseudo-critical temperature for the chiral phase transition T_C = 156.5 +- 1.5 MeV

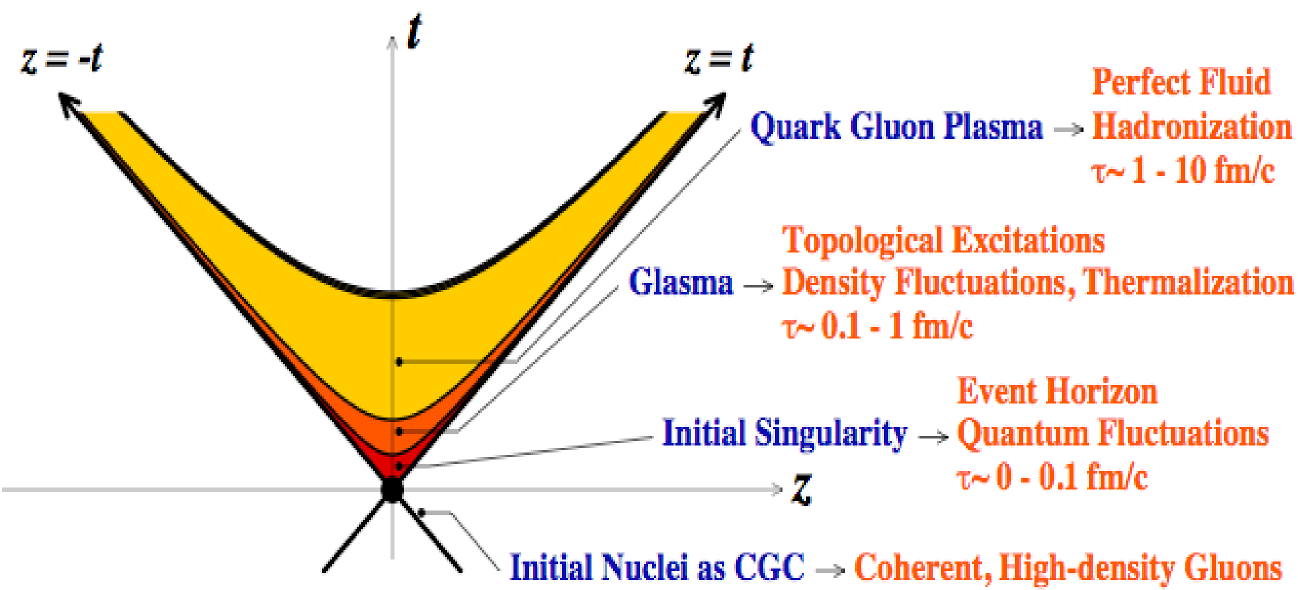
Experimental program

QGP and phase diagram studied in high energy collisions of nuclei

accelerator	years	$\sqrt{s_{NN}}$	large exp.
AGS	1986 - 2002	2.7 - 4.8 GeV	5
SPS	since 1986	6.2 - 19.3 GeV	7
RHIC	since 2000	7.0 – 200 GeV	4
LHC	since 2009	2.76 – 5.02 TeV	3 (4)

Space-time evolution of a relativistic nuclear collision at LHC energy



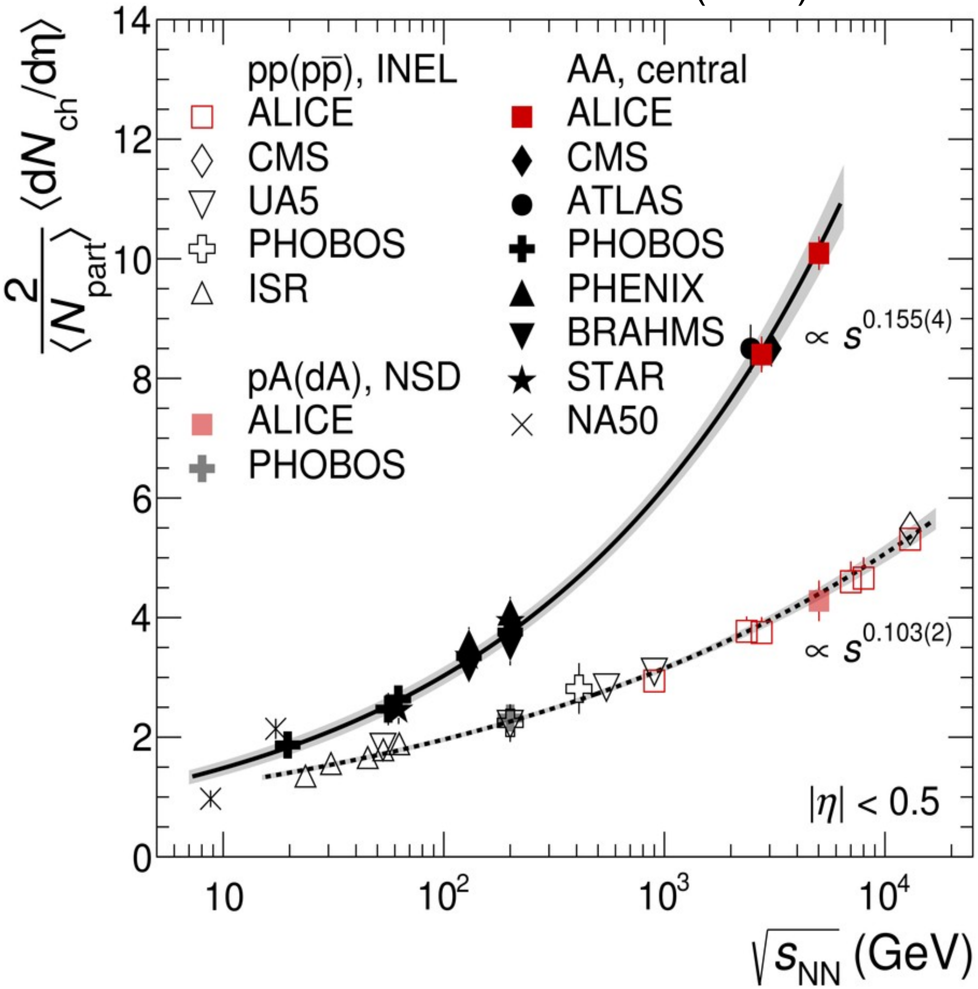


similar to early universe, fluctuations observed in the much later phase may allow to deduce early singularities

one possible view (courtesy L. McLerran)

Charged particle production

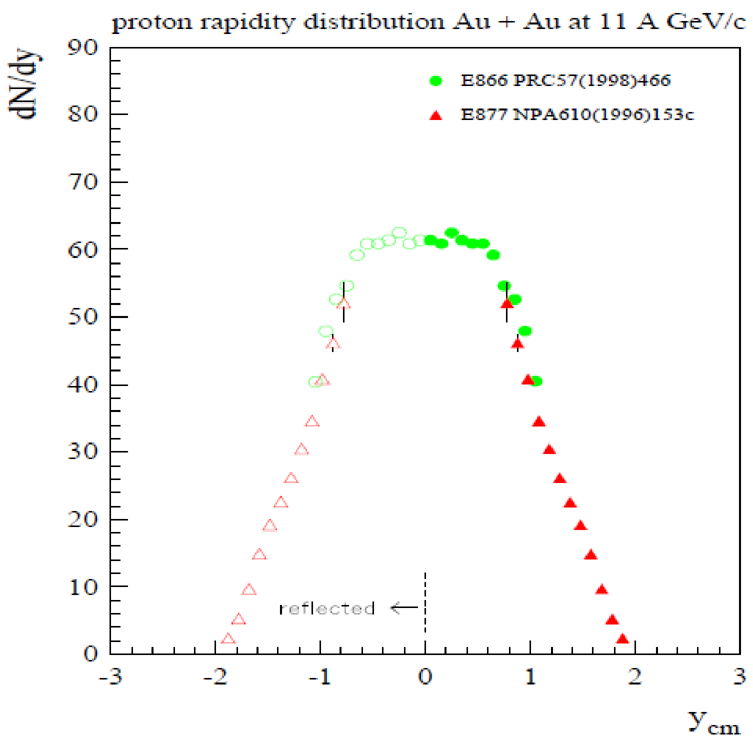
ALICE Coll. PRL116 (2016) 222302



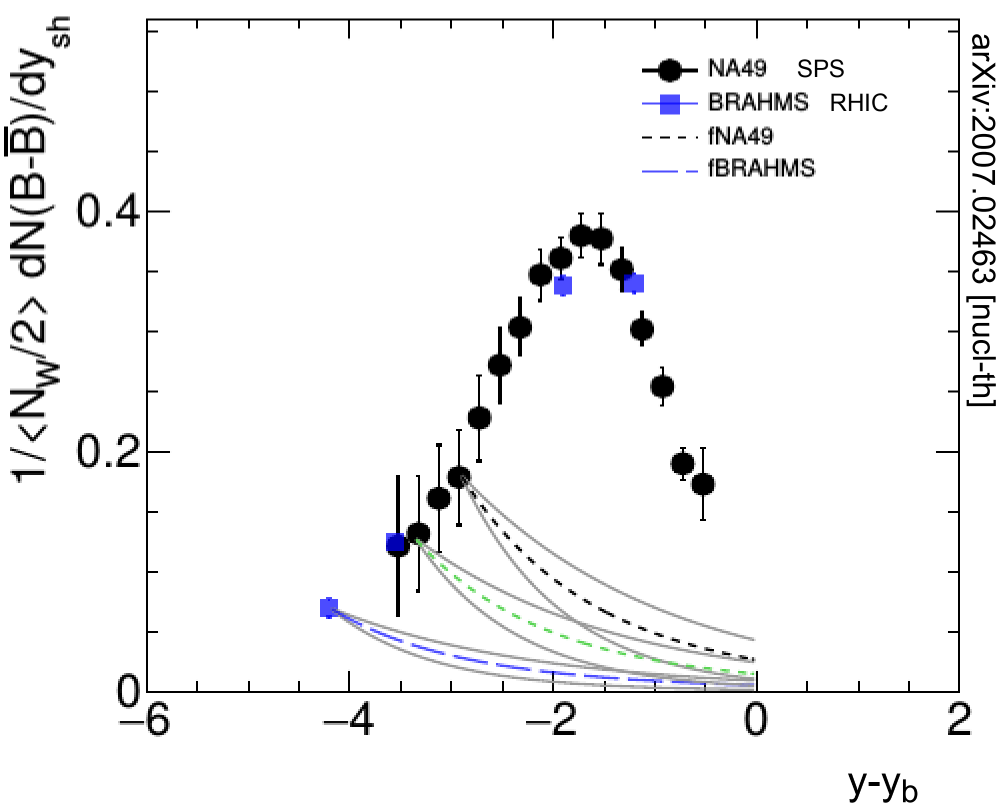
increase in nuclear collisions much faster with $\sqrt{s}$ than in pp

→ larger fractional energy loss in nuclear collision

Nuclear stopping power



AGS: nuclei stop each other completely $\Delta y = 1.7$
SPS: slight onset of transparency $\Delta y = 2.0$
RHIC: 'limiting fragmentation' $\Delta y = 2.0$
implying fraction $1\text{-}\exp(-\Delta y)$ = 86% E_{loss}
energy deposit in central fireball
in pp (Fermilab data): $\Delta y = 0.95 \triangleq 60\%\ E_{loss}$



Initial Energy Density

$$\epsilon_0 = \mathrm{d}E_t/\mathrm{d}y/A_t \times \mathrm{d}y/\mathrm{d}z = \langle m_t \rangle 1.5 \mathrm{d}N_{ch}/\mathrm{d}y/A_t \times \mathrm{d}y/\mathrm{d}z$$

Bjorken formula* using Jacobian $dy/dz=1/\tau_0$
typically evaluated at τ_0 = 1 fm/c

	$\sqrt{s_{NN}}$ (GeV)	dE_t/dy (GeV)	ε_0 (GeV/fm^3)	T (GeV)
AGS	4.8	200	1.4	0.17
SPS	17.2	450	3.0	0.21
RHIC	200	600	5.5	0.30
at $\tau_0 = 1/p_{sat} = 0.14$ fm/c			40	0.49
LHC	2760	1755	11.7	0.36
at $\tau_0 = 1/p_{sat} = 0.08$ fm/c			146	0.68

all above lQCD result for pseudo-critical energy density and temperature

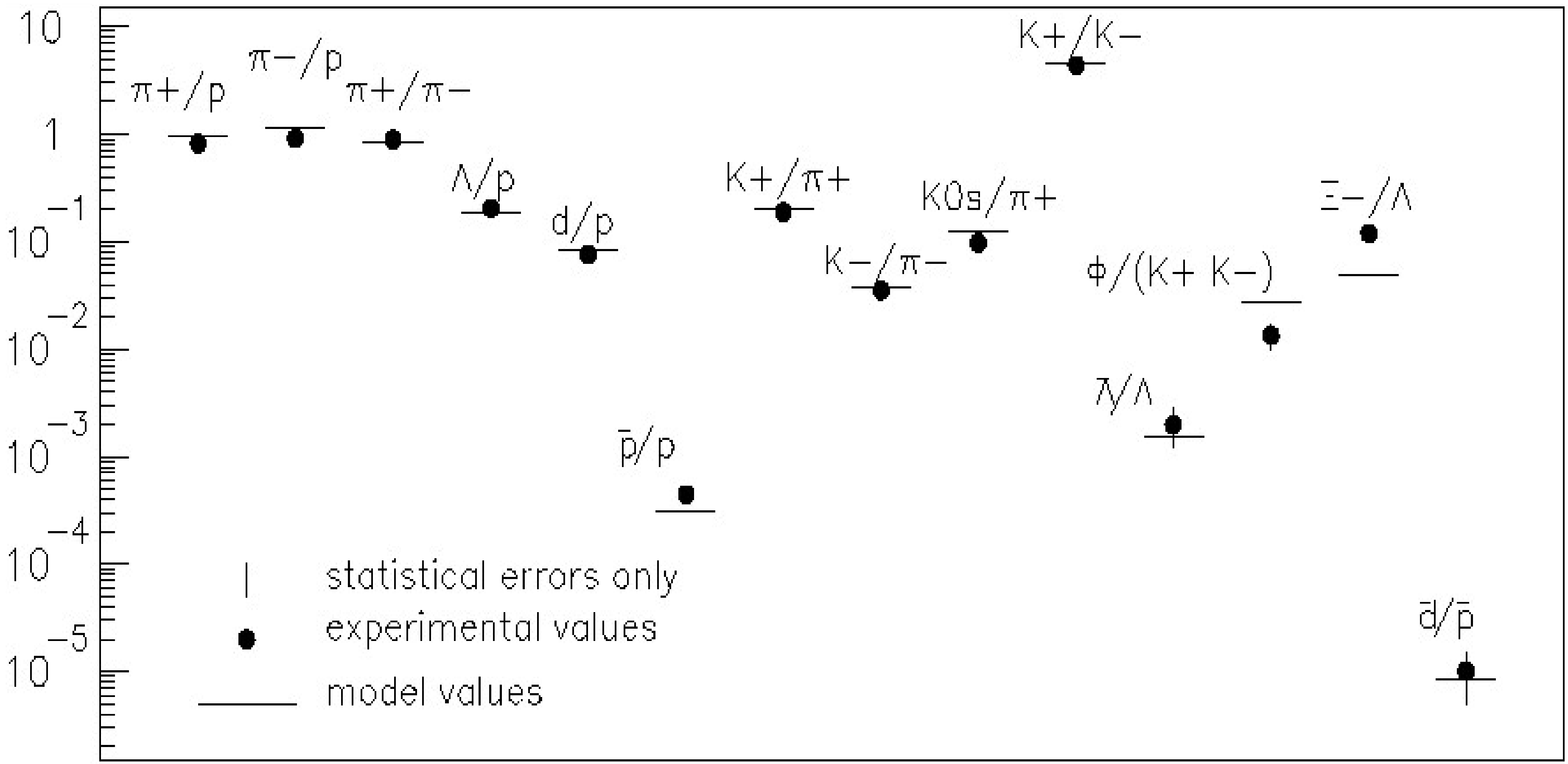
* these are lower bounds; if during expansion work is done (pdV) initial energy density higher (indications from hydrodynamics at LHC: factor 3)

Hadronization of the fireball

hadro-chemical freeze-out at phase boundary between QGP and hadronic matter

First measurement of a comprehensive set of hadrons at BNL AGS by 1993

14.6 A GeV/c central Si + Au collisions – combined data by E802, E810, E814



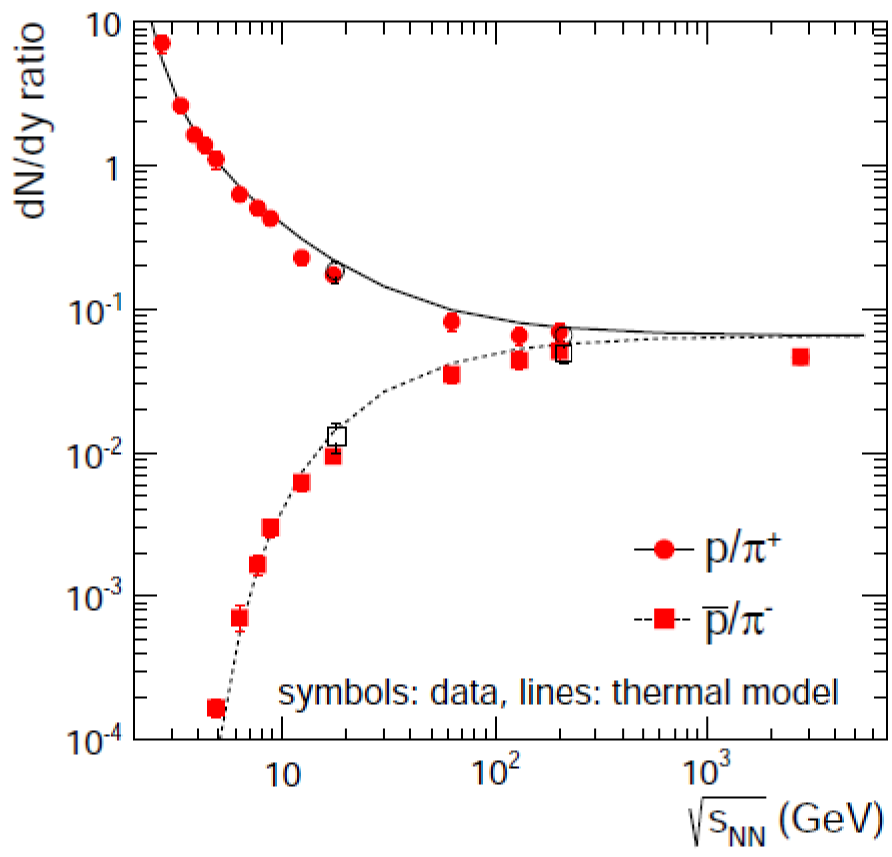
first successful application of statistical hadronization model (grand canonical ensemble) - 2 fit parameters dynamic range: 9 orders of magnitude! no deviation

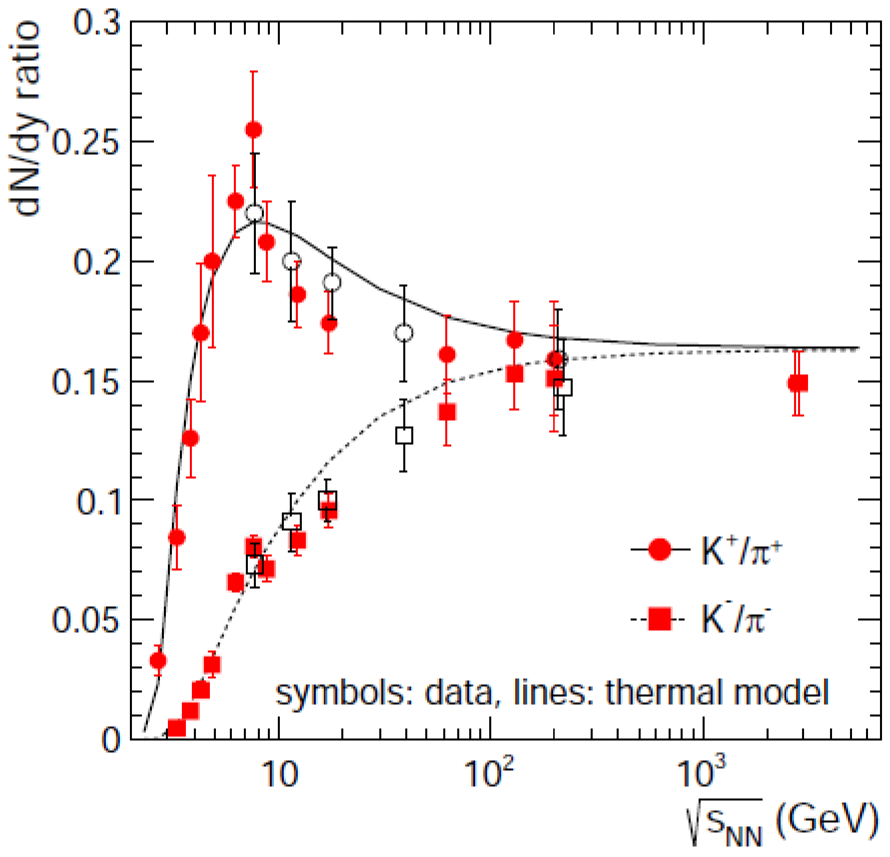
P. Braun-Munzinger, J. Stachel, J.P. Wessels, N. Xu, PLB344 (1995) 43

Beam energy dependence of hadron yields in AuAu and PbPb collisions from AGS to LHC

fits work equally well at higher beam energies
following the obtained T and μ_b evolution, features of proton/pion and kaon/pion ratios reproduced in detail

A. Andronic, P. Braun-Munzinger, J.Stachel, PLB 673 (2009) 142





Production of hadrons and (anti-)nuclei at LHC

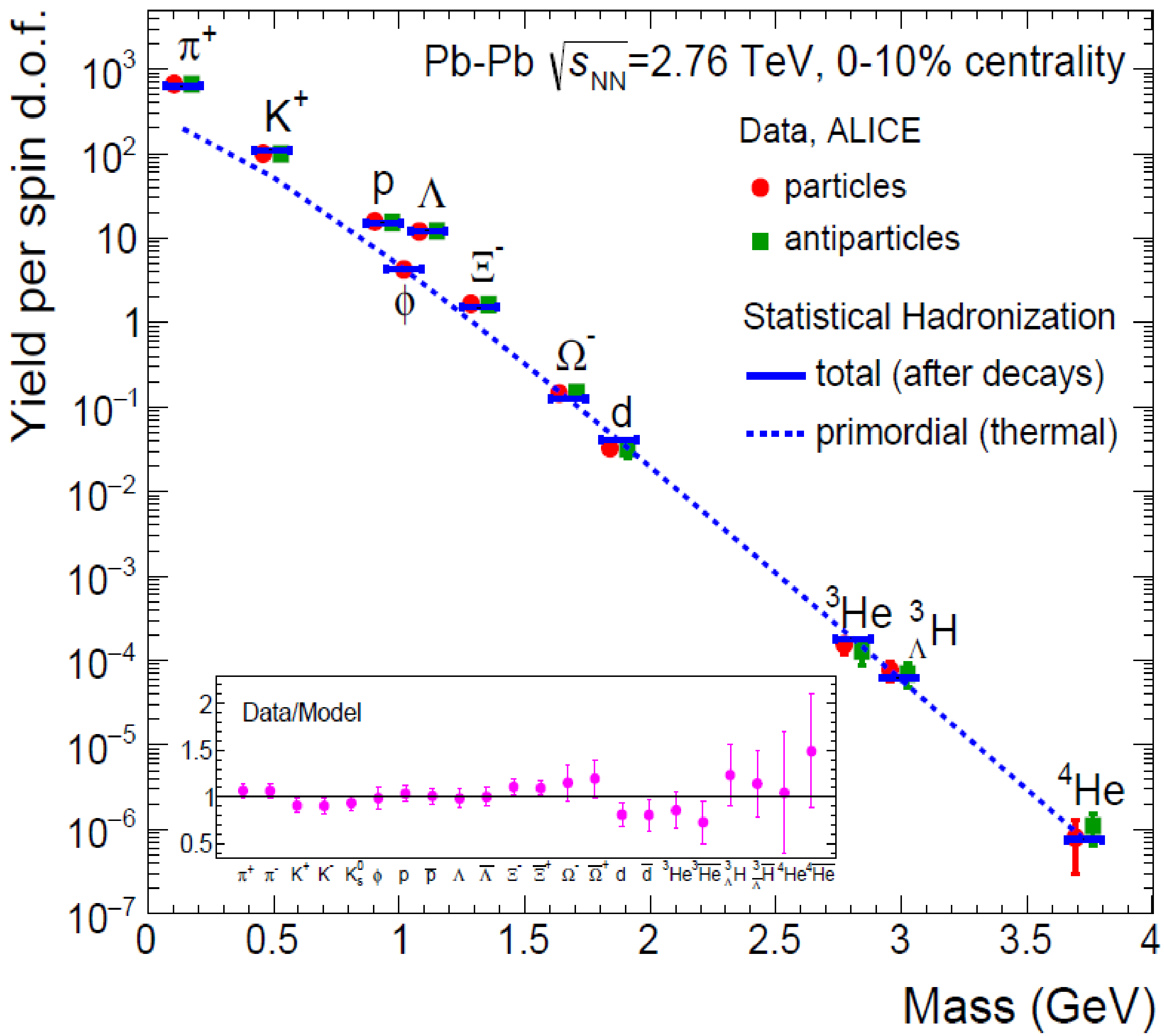
1 free parameter: temperature T

T = 156.5 ± 1.5 MeV

agreement over 9 orders of magnitude with QCD statistical operator prediction (- strong decays need to be added)

- matter and antimatter are formed in equal portions at LHC
- even large very fragile hypernuclei follow the same systematics

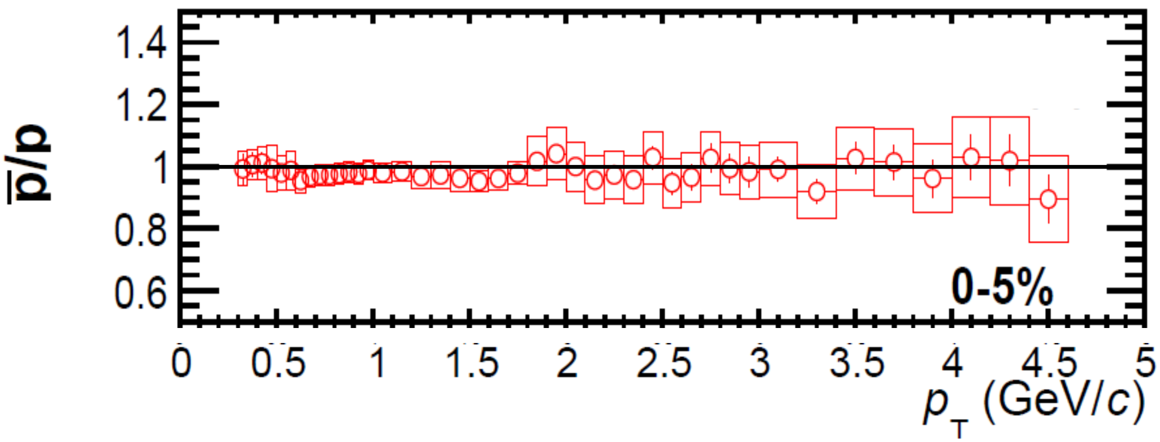
A. Andronic, P. Braun-Munzinger, K. Redlich, J. Stachel, Nature 561 (2018) 321



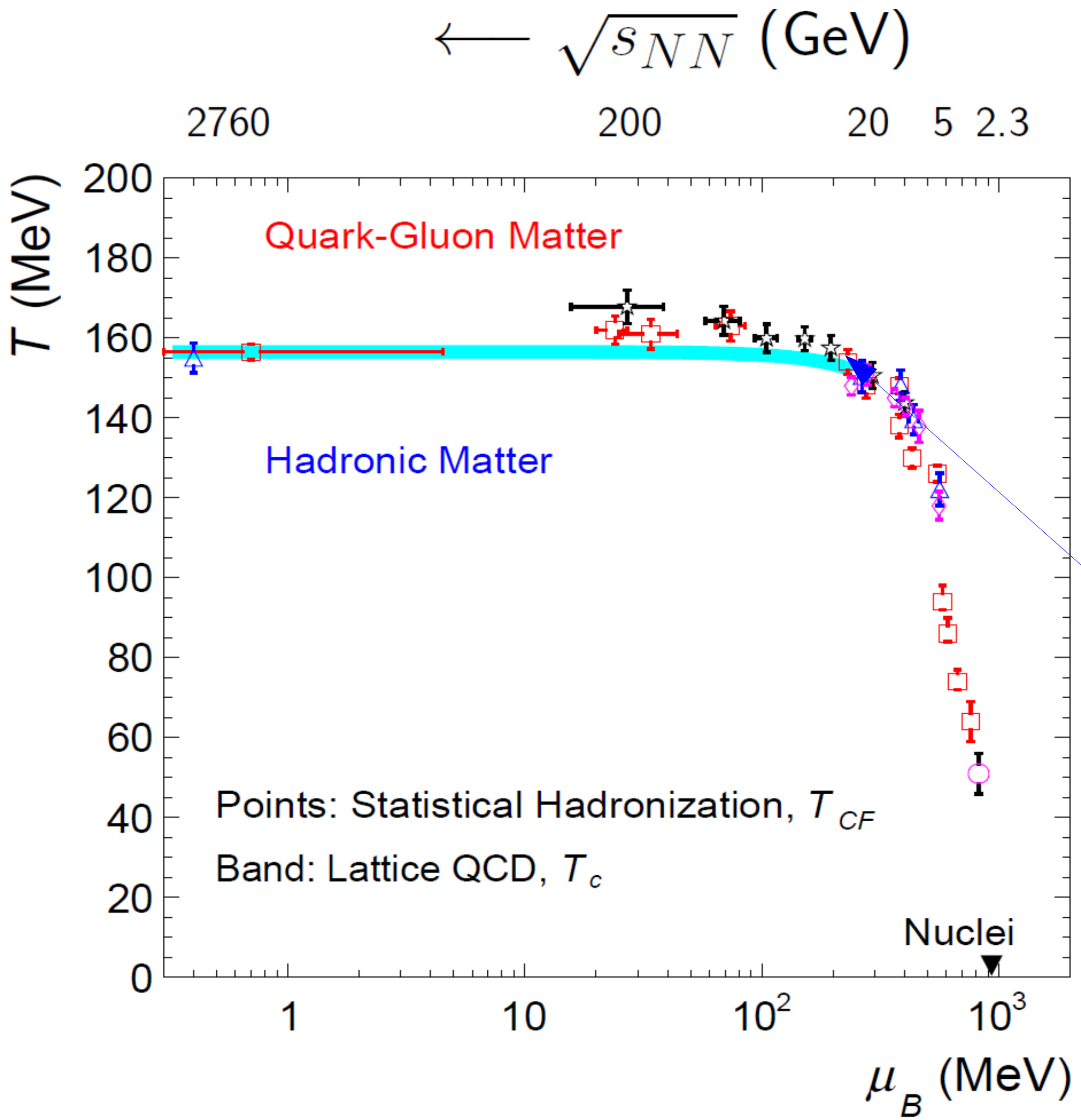
Biggest difference LHC compared to lower energies

- matter and anti-matter produced in equal proportions at LHC
- consistent with net-baryon free central region, (μ_b < 1 MeV) similar to early universe

ALICE 2.76 TeV PbPb



Energy dependence of temperature and baryochem pot.



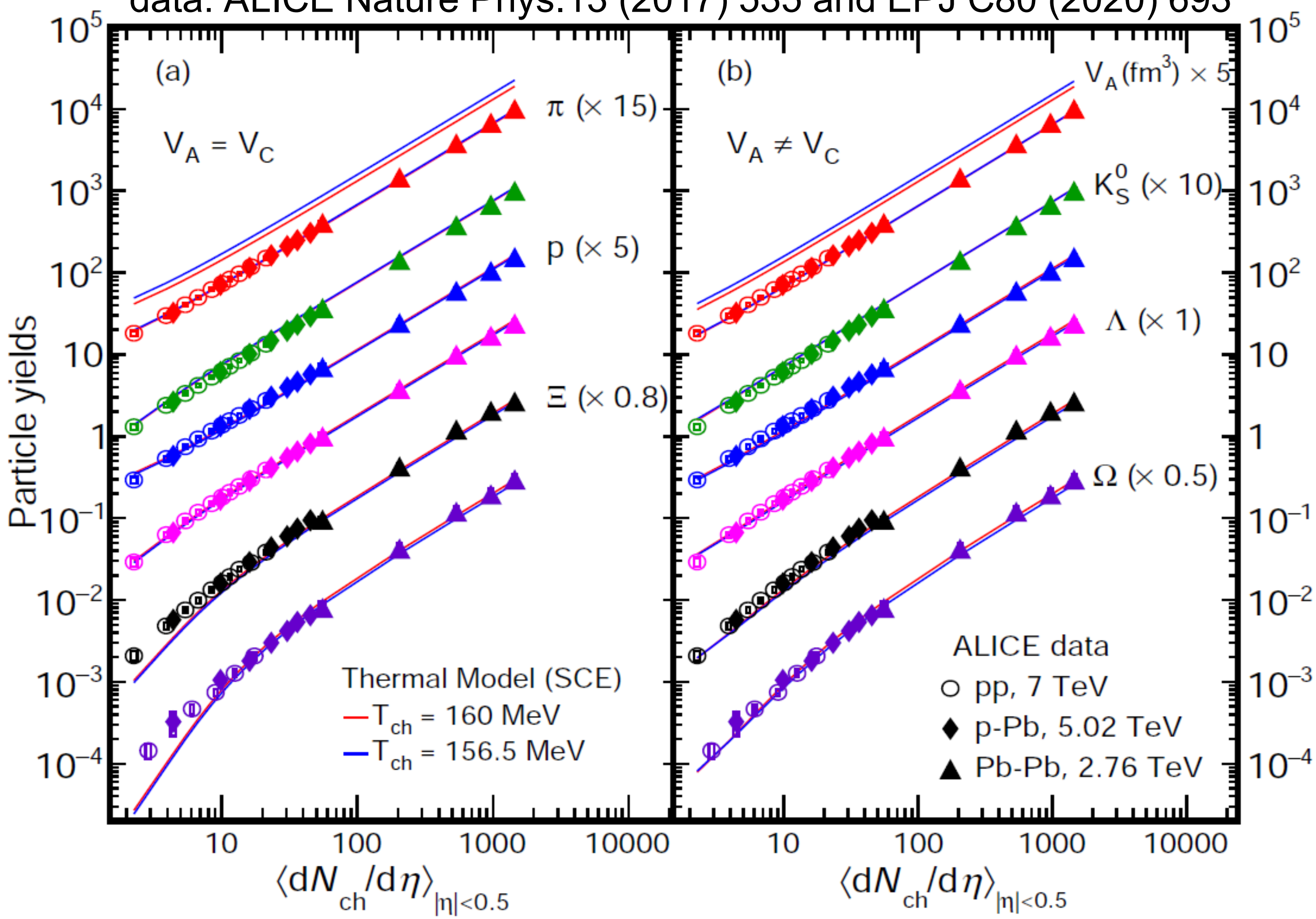
hadron yields for Pb-Pb central collisions from LHC down to RHIC, SPS, AGS and even SIS energies well described by a statistical ensemble

- there is a limiting temperature for a hadronic system
T_{lim} = 159 ± 3 MeV
reached for $\sqrt{s_{NN}} \geq 20$ GeV

- T at LHC in exact agreement with the pseudo-critical temperature
T_c = 156.5 ± 1.5 MeV from lQCD
A. Bazavov et al. PLB 795 (2019) 15

from pp to Pb-Pb collisions: smooth evolution with system size

data: ALICE Nature Phys.13 (2017) 535 and EPJ C80 (2020) 693



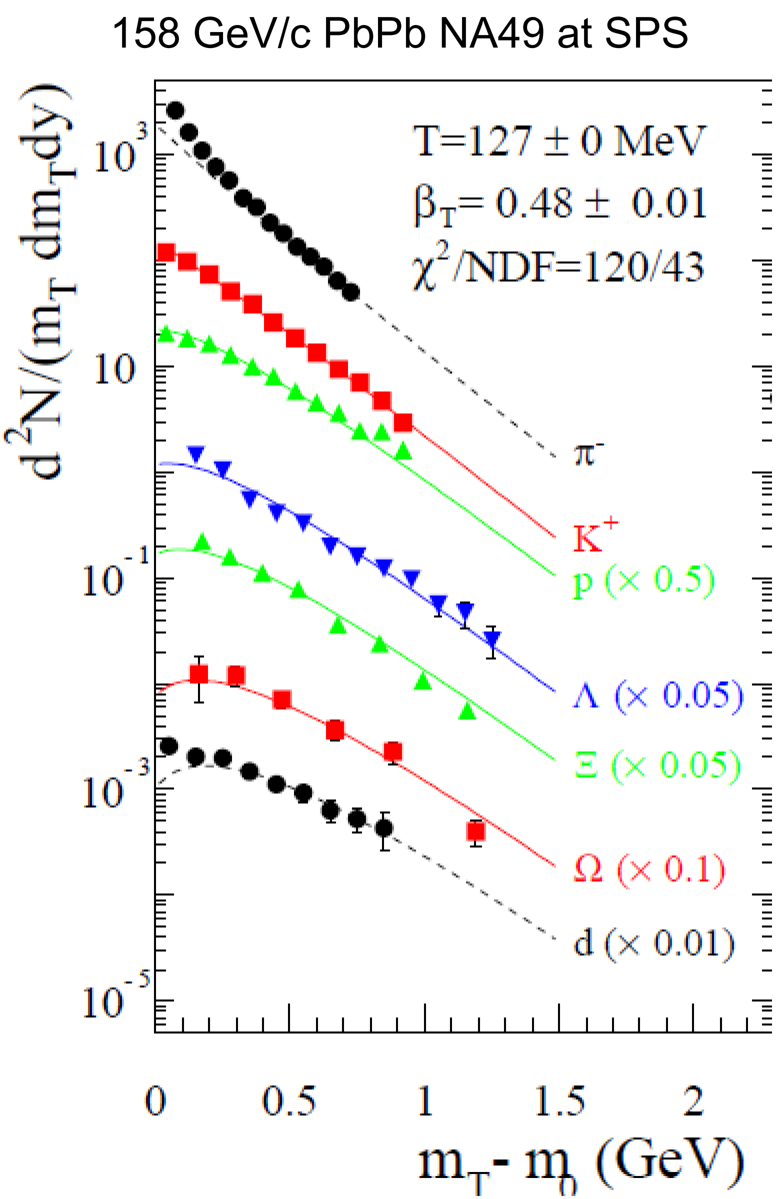
universal hadronization can be described with few parameters in addition to T and μ_B → transition from canonical to grand-canonical thermodynamics

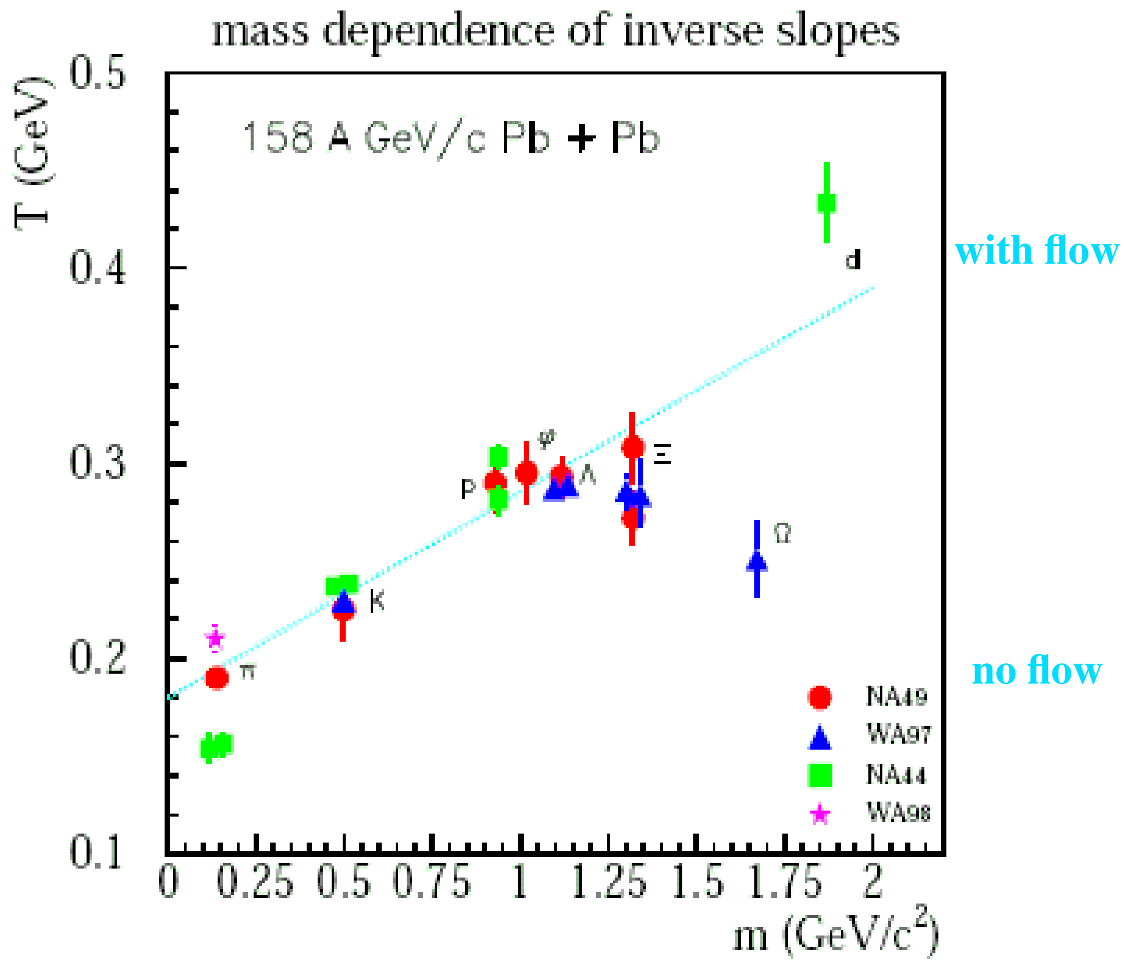
J. Cleymans, P.M. Lo, K. Redlich, N. Sharma, PRC 103 (2021) 014904

Hadron spectra and correlations

- reveal in addition to kinetic freeze-out temperature
  strong collective expansion
- survival of early fluctuations
- transport parameters of the QGP

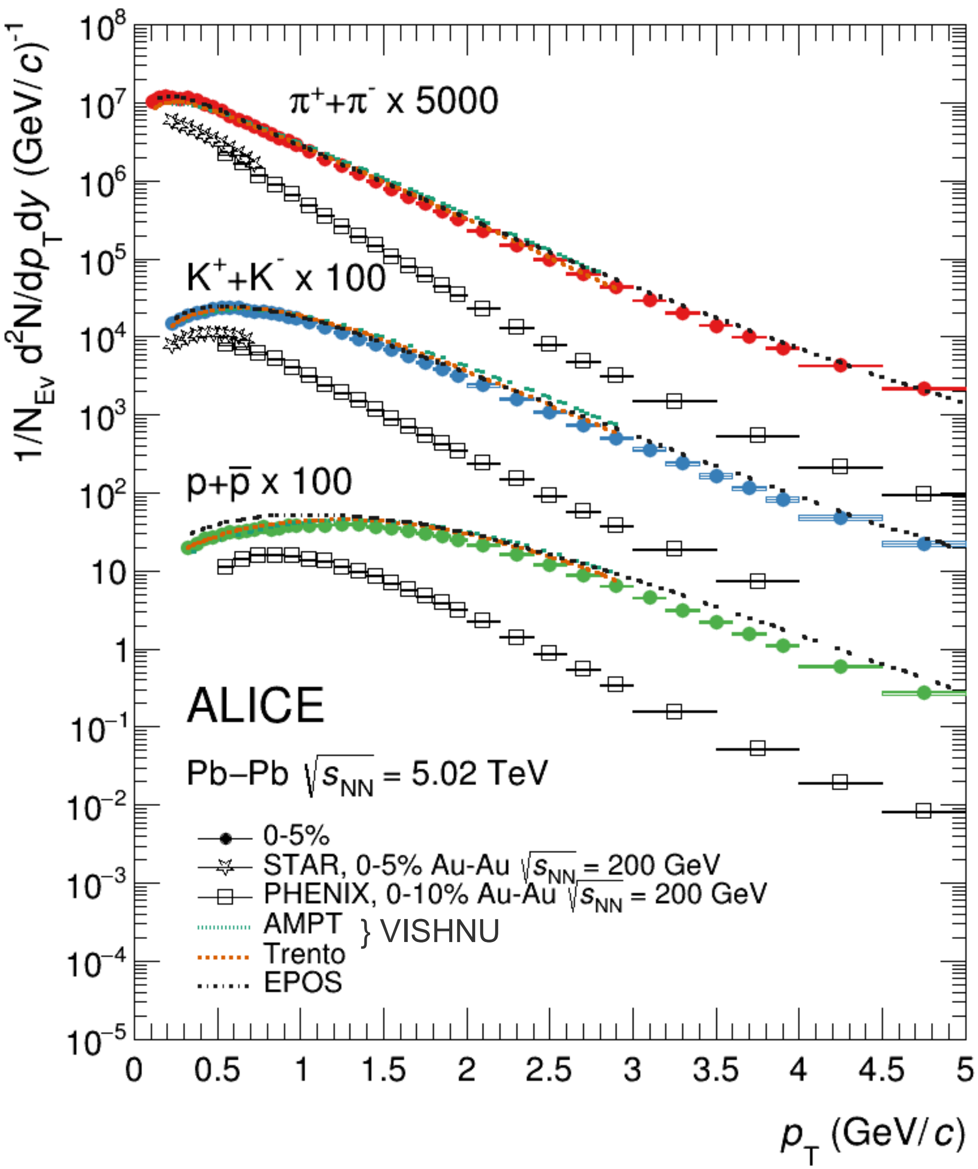
Spectra of identified hadrons at SPS





strong (linear) mass dependence of spectral slopes:
superposition of random thermal motion and
ordered collective expansion (flow) - $\beta_T > 0.5$

Spectra of identified hadrons at RHIC and LHC



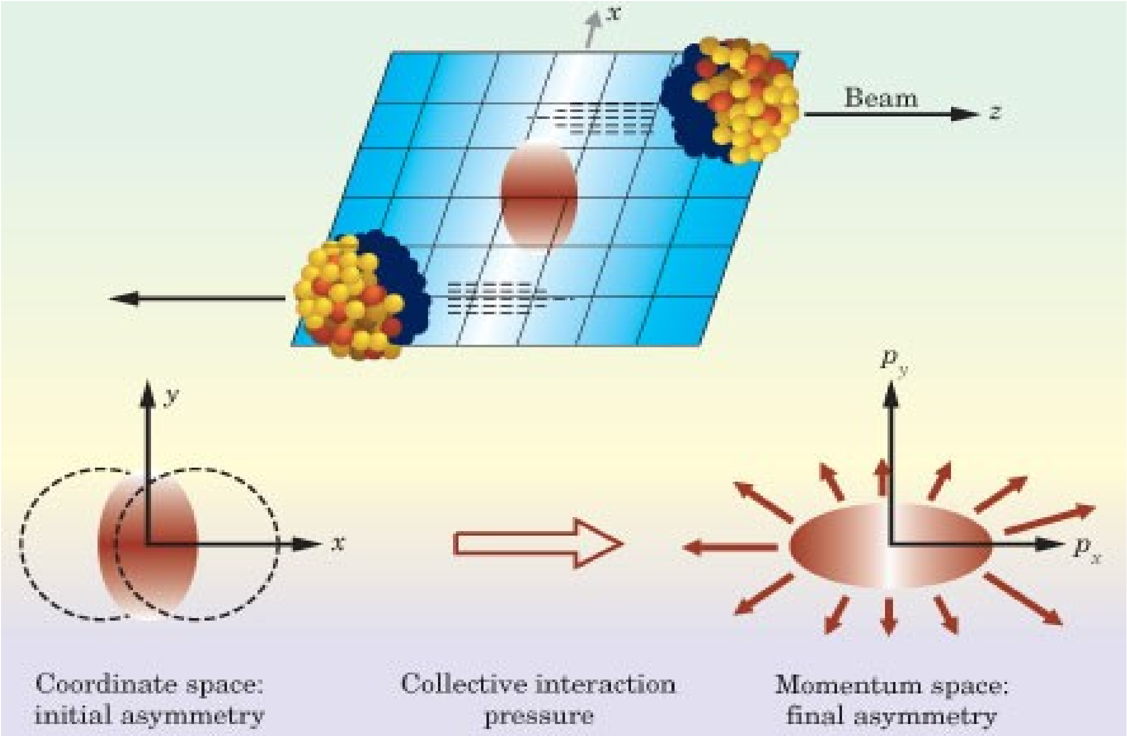
spectral shapes exhibit even stronger mass dependence
- characteristic for hydrodynamic expansion

indicate at LHC significantly larger expansion velocity than at RHIC

captured well by hydrodynamic expansion
(EPOS, K. Werner et al. Subatech)

expansion velocity at surface: ¾ c

Azimuthal anisotropy of transverse spectra



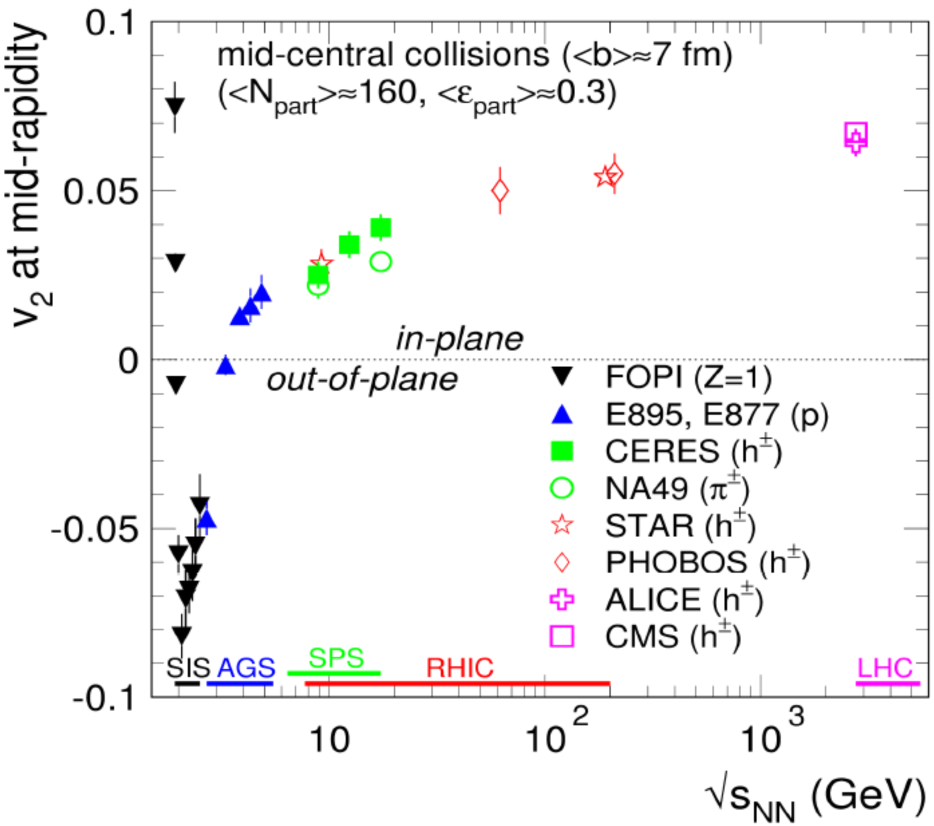
Fourier decomposition of momentum distributions rel. to reaction plane:

$$\frac{dN}{dp_t\, dy\, d\phi} = N_0 \cdot \left[1 + \sum_{i=1} 2\, v_i(y, p_t) \cos(i\phi) \right]$$

quadrupole component v_2
"elliptic flow"

the v_n are the equivalent of the power spectrum of cosmic microwave rad.

Elliptic flow as function of collision energy



- effect of expansion (positive v_2) seen from top AGS energy upwards
- at lower energy: shadowing by fragments
- first discovered as tiny 2% effect by E877 in 1993

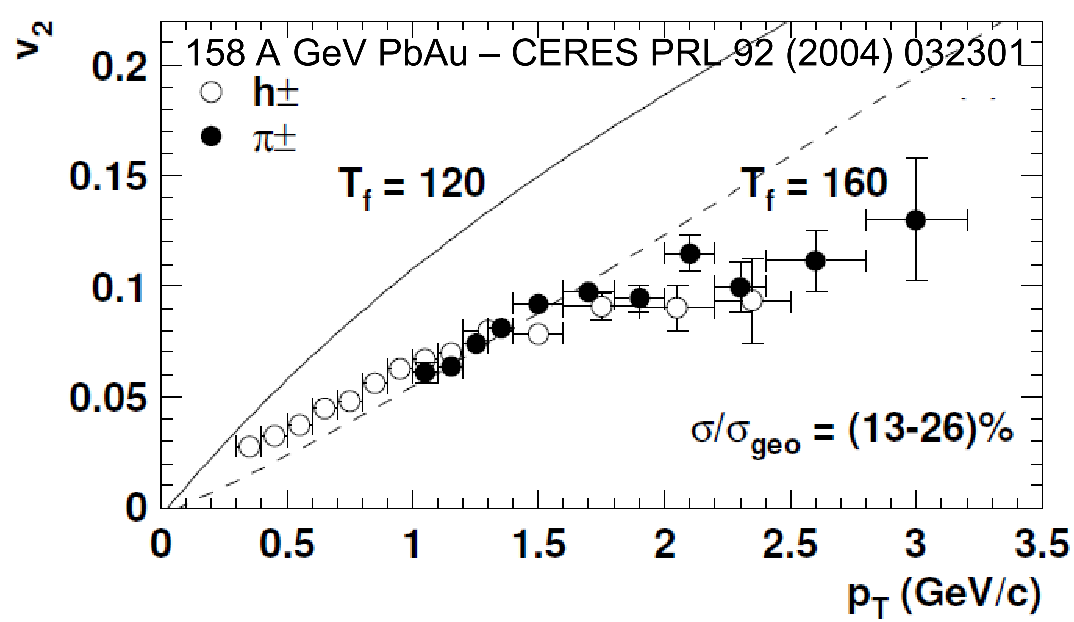
very much inspired by

Anisotropy as a signature of transverse collective flow

Jean-Yves Ollitrault (Saclay) (Mar 20, 1992)

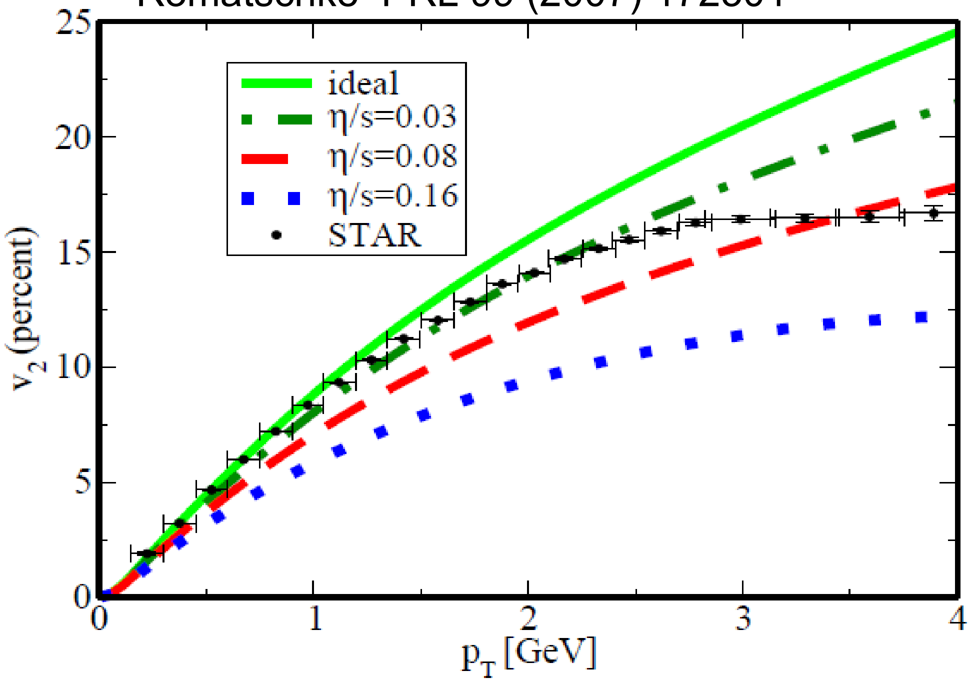
Published in: *Phys.Rev.D* 46 (1992) 229-245

at top SPS energy, modelling with ideal relativistic hydrodynamics close to exp. data



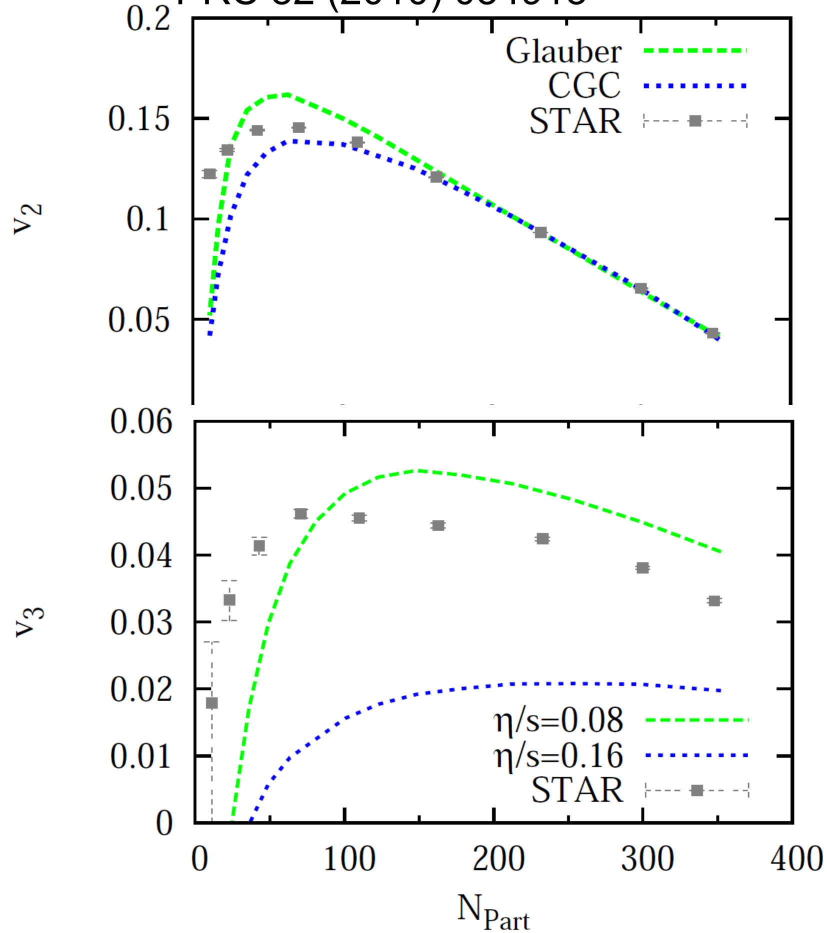
Discovery of RHIC: paradigm of QGP as near ideal liquid

Romatschke PRL 99 (2007) 172301



how perfect is the fluid observed at RHIC?
very small ratio of shear viscosity to entropy density η/s describes data

Alver, Gumbeaud, Luzum, Ollitrault, PRC 82 (2010) 034913



in hydro regime v_2 driven by
- initial condition and
- properties of the liquid as η/s

→ ambiguity between the two can be resolved by correlating observables

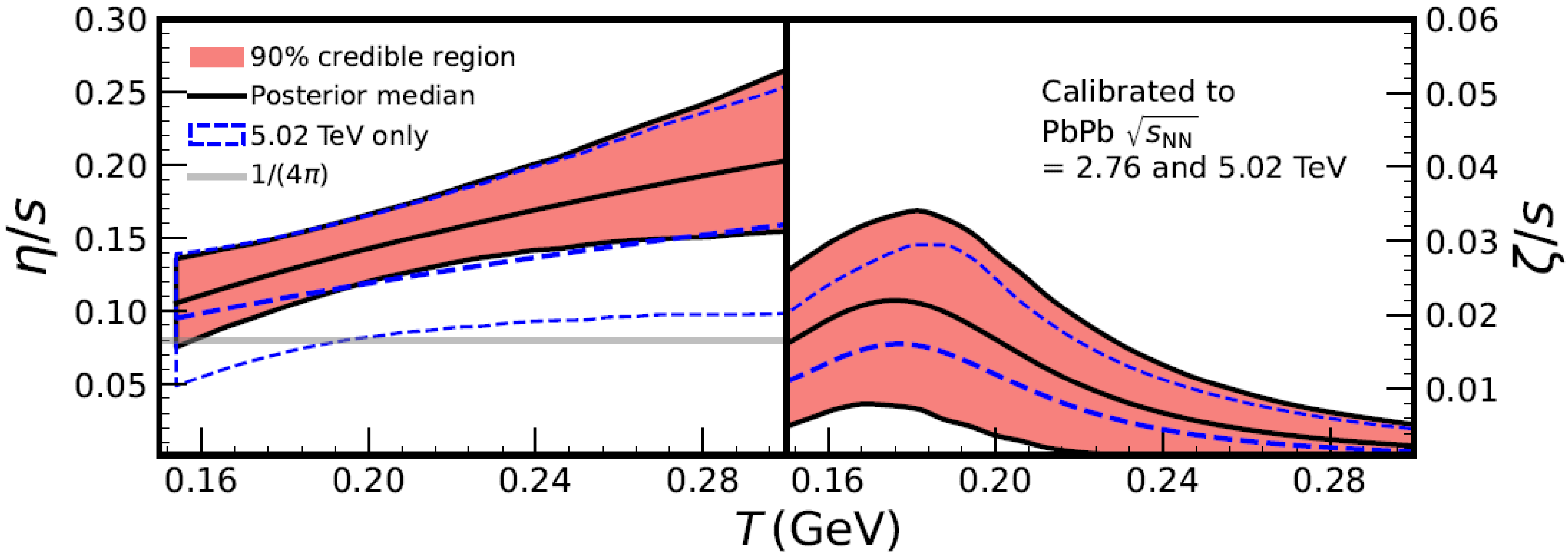
Constraining initial condition and QGP medium properties

much higher precision can be obtained from cumulants defined in terms of multiparticle azimuthal correlations

N. Borghini, P.M. Dinh, J.Y. Ollitrault, PRC 64 (2001) 054901 $\rightarrow$ visionary paper guiding the way

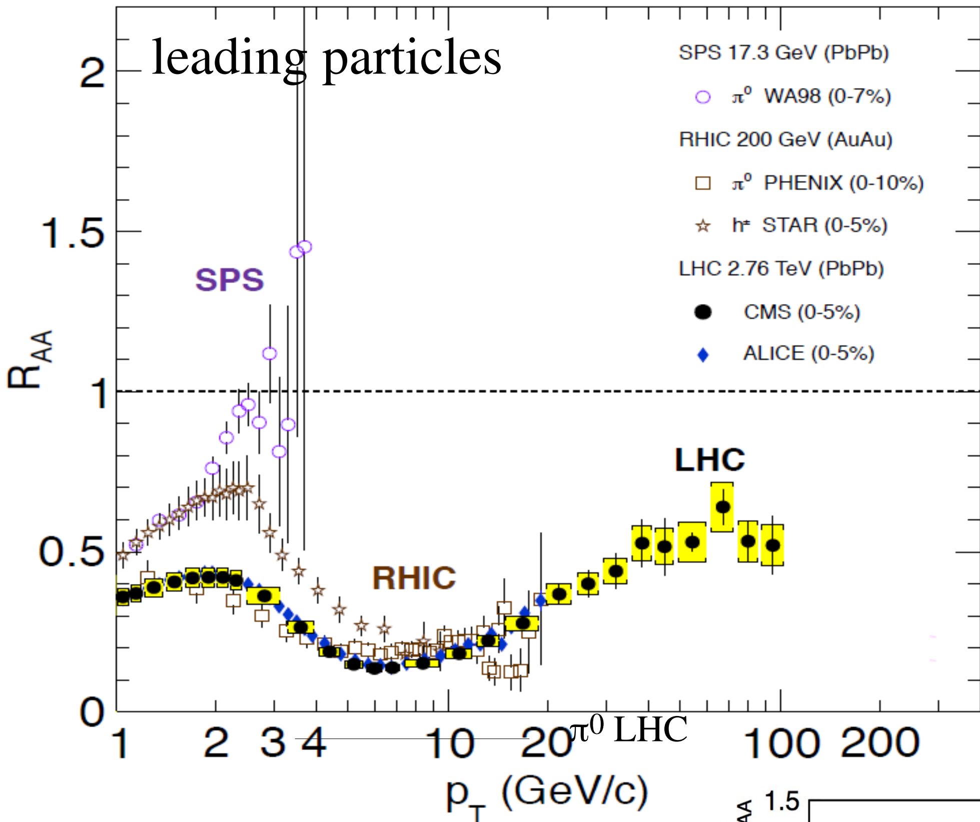
LHC: global Bayesian analysis of such new collective flow observables in PbPb from ALICE

J. Parkkila et al., arXiv: 2111.08145



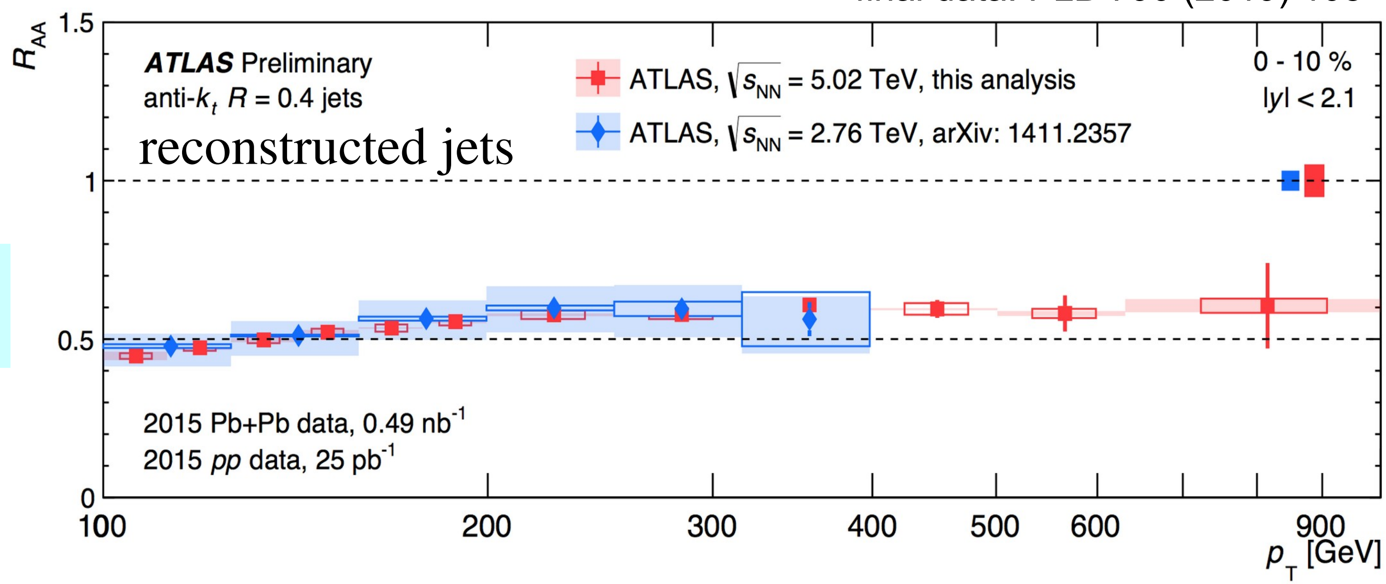
near T_C, shear viscosity/entropy density close to AdS/CFT lower bound $1/4\pi$ rising with temperature in QGP – bulk viscosity/entropy dens. peaks near T_C

Jet quenching – parton energy loss in QGP



- suppression of leading particles first observed at RHIC
- still stronger at LHC
- upturn beyond 7 GeV new at LHC
- levels off at 0.5

final data: PLB 790 (2019) 108



jet suppression at level 0.6 out to 1 TeV

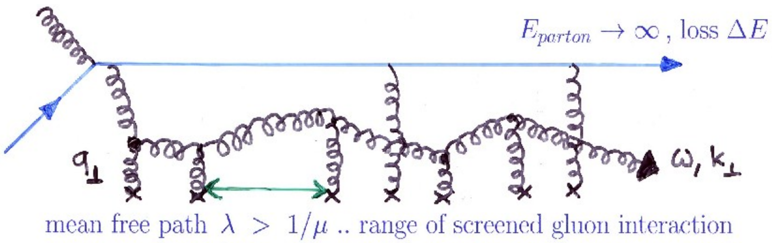
Extracting the jet quenching parameter

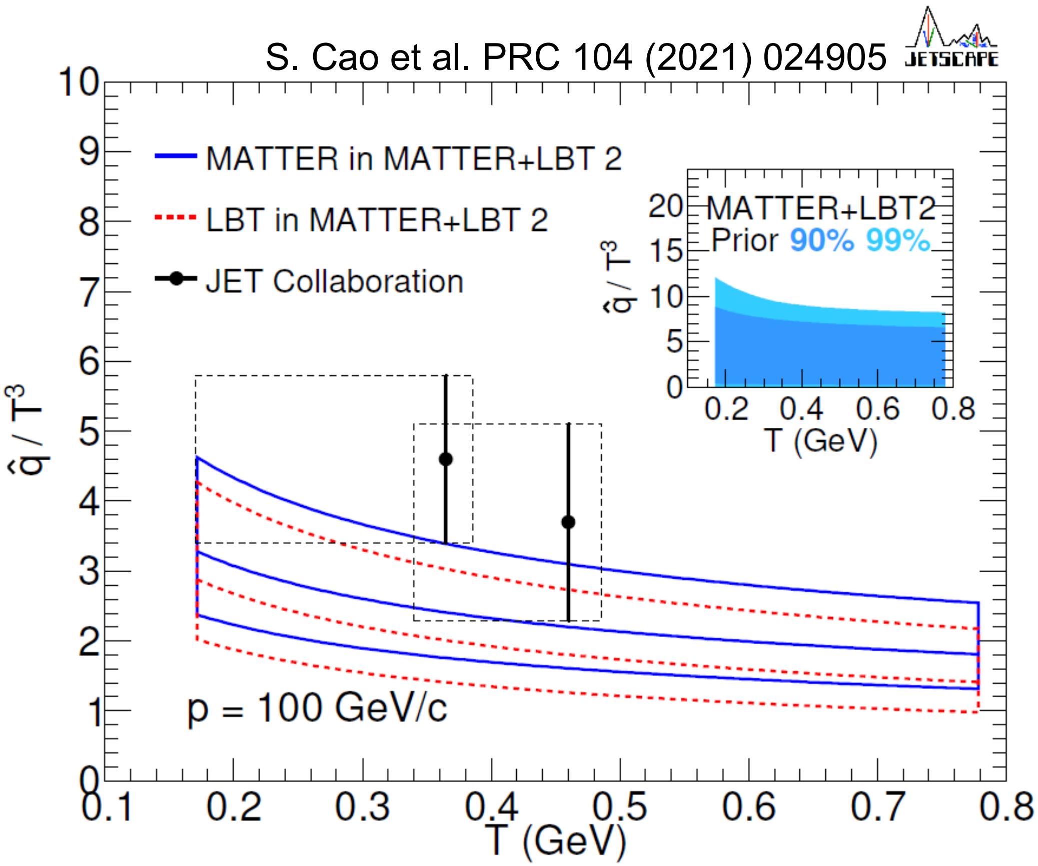
prediction: H. Baier, Y.L. Dokshitzer, A.H. Mueller, S. Peigne, D. Schiff, NPB 483 (1997) 291 and 484 (1997) 265

$$dE/dx \propto \rho\, \sigma \langle k_t^2 \rangle\, L$$

density of color charge carriers

transport coefficient $\hat{q} \propto \rho\, \sigma \langle k_t^2 \rangle$





determine transport coefficient from comparing a combined model of splitting of high virtuality partons (MATTER) and scattering between jet partons and a thermal QGP (LBT) to inclusive hadron R_{AA} data for RHIC and LHC (Bayesian parameter estimation)

obtain

$\hat{q}$ = 0.7±0.3 GeV^2/fm at T = 400 MeV

factor 20-40 larger than in cold nuclear matter (from DIS)!

Charmonia as a probe of deconfinement

the original idea: implant charmonia into the QGP and observe their modification (Debye screening of QCD), in terms of suppressed production in nucleus-nucleus collisions – sequential melting

T. Matsui and H. Satz PLB 178 (1986) 416

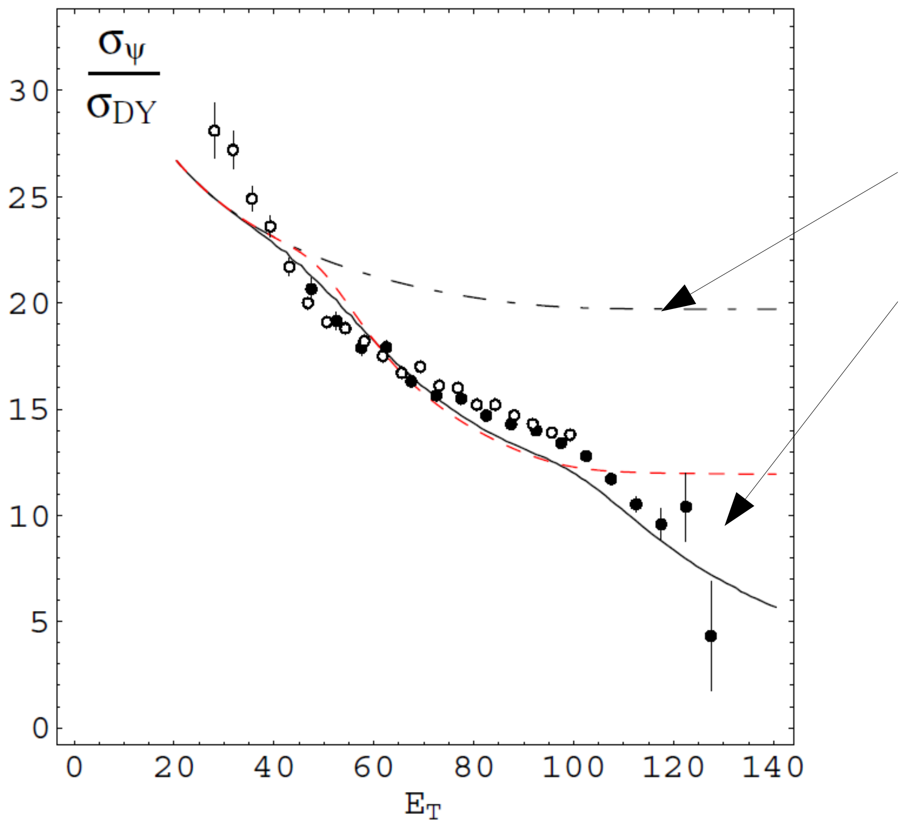
"If high energy heavy-ion collisions lead to the formation of a hot quark-gluon-plasma, then color screening prevents $c\bar{c}$ binding in the deconfined interior of the interaction region. ... It is concluded that J/ψ suppression in nuclear collisions should provide an unambiguous signature of quark-gluon-plasma formation."

First J/ψ suppression in nuclear collisions at SPS

key measurements by NA38 find suppression for O and S induced collisions
QM 1991 conf.: data on photon, hadron, and nucleus-nucleus coll. described by nuclear absorption

C. Gerschel and J. Hűfner, Z.Physik C56 (1992) 171

data: NA50 Phys. Lett. B447 (2000) 28



finally observations NA50:
- in pp, pA and light nuclei, suppression pattern consistent with absorption on (cold) nuclear matter 4.3±0.5 mb
- in central collisions of PbPb much stronger suppression

data described by dissolution of J/ψ at critical density $n_c = 3.7/fm^2$ — — —
& including energy density fluct. ——

J.P. Blaizot, P.M. Dinh, J.Y. Ollitrault PRL 85 (2000) 4012

Charmonium formation at hadronization: extension of statistical model to include charmed hadrons

new insight:
QGP screens all charmonia, but charm quarks remain in the fireball
charmonium production takes place at the phase boundary
➡ enhanced production at colliders – signal for deconfinement

P. Braun-Munzinger, J. Stachel, PLB 490 (2000) 196

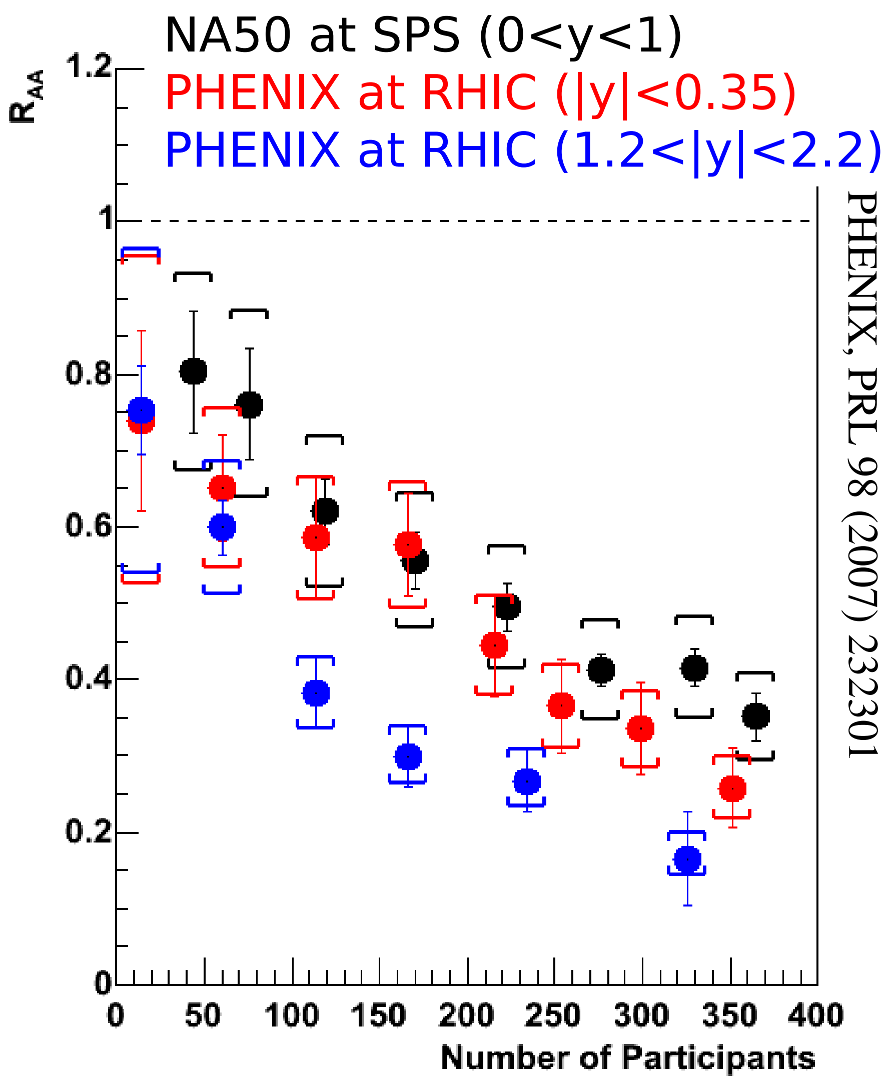
technically:

- assume: all charm quarks are produced in initial hard scattering
  number not changed in QGP
  $N_{c\bar{c}}^{direct}$ from data (total charm cross section) or from pQCD

- hadronization at T_c following grand canonical statistical model used for hadrons with light valence quarks (canonical corr. if needed)
  technically number of charm quarks fixed by a charm-balance equation containing fugacity g_c (no free parameter)

$$N_{c\bar{c}}^{direct} = \frac{1}{2} g_c V \left(\sum_i n_{D_i}^{therm} + n_{\Lambda_i}^{therm}\right) + g_c^2 V \left(\sum_i n_{\psi_i}^{therm}\right) + \ldots$$

J/ψ suppression at RHIC – 200 A GeV AuAu



PHENIX talk at QM2006:

suppression patterns are remarkably similar at SPS and RHIC!

cold matter suppression larger at SPS, hot matter suppression larger at RHIC, balance?

recombination cancels additional suppression at RHIC?

how did we get so “lucky”?

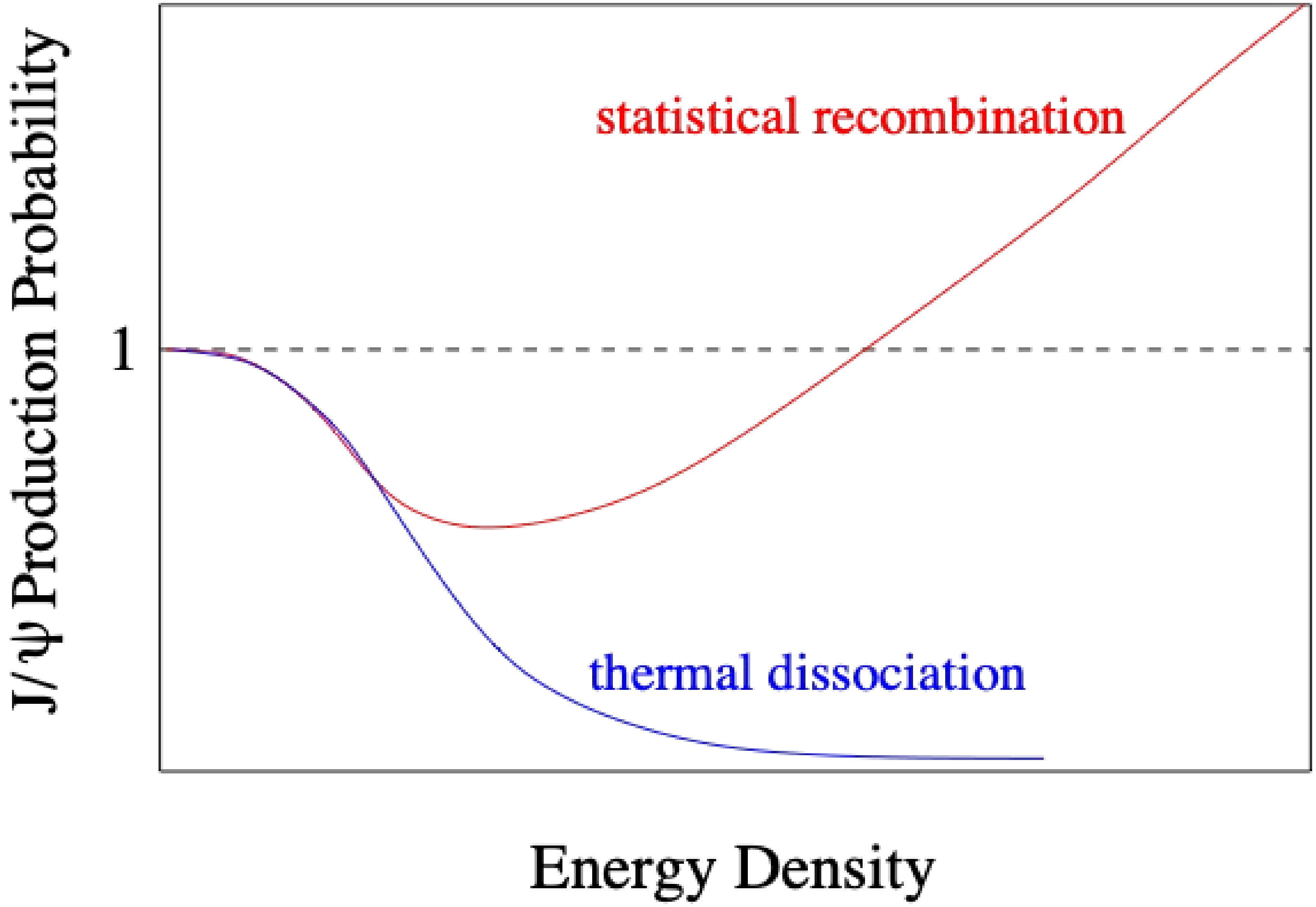
data could be indeed described by statistical hadronization using pQCD charm cross section

A.Andronic, P.Braun-Munzinger, K.Redlich, J.Stachel, PLB 652 (2007)259

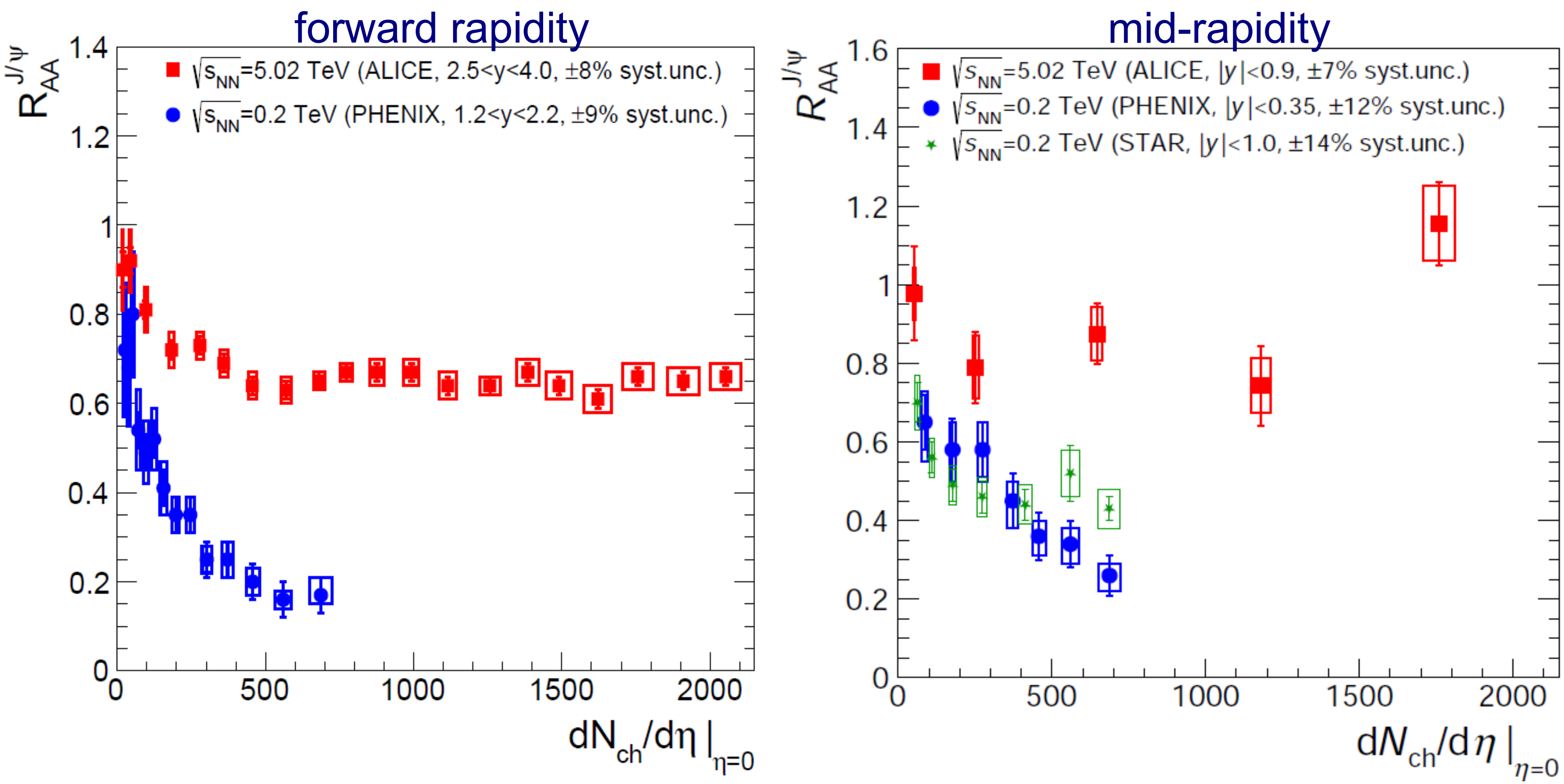
Expectations for LHC

2 possibilities:

Kluberg, Satz, arXiv:0901.3831

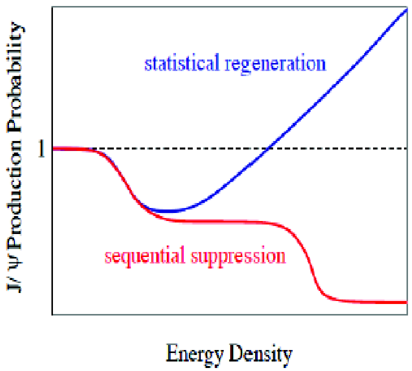


J/ψ production in PbPb collisions: LHC relative to RHIC



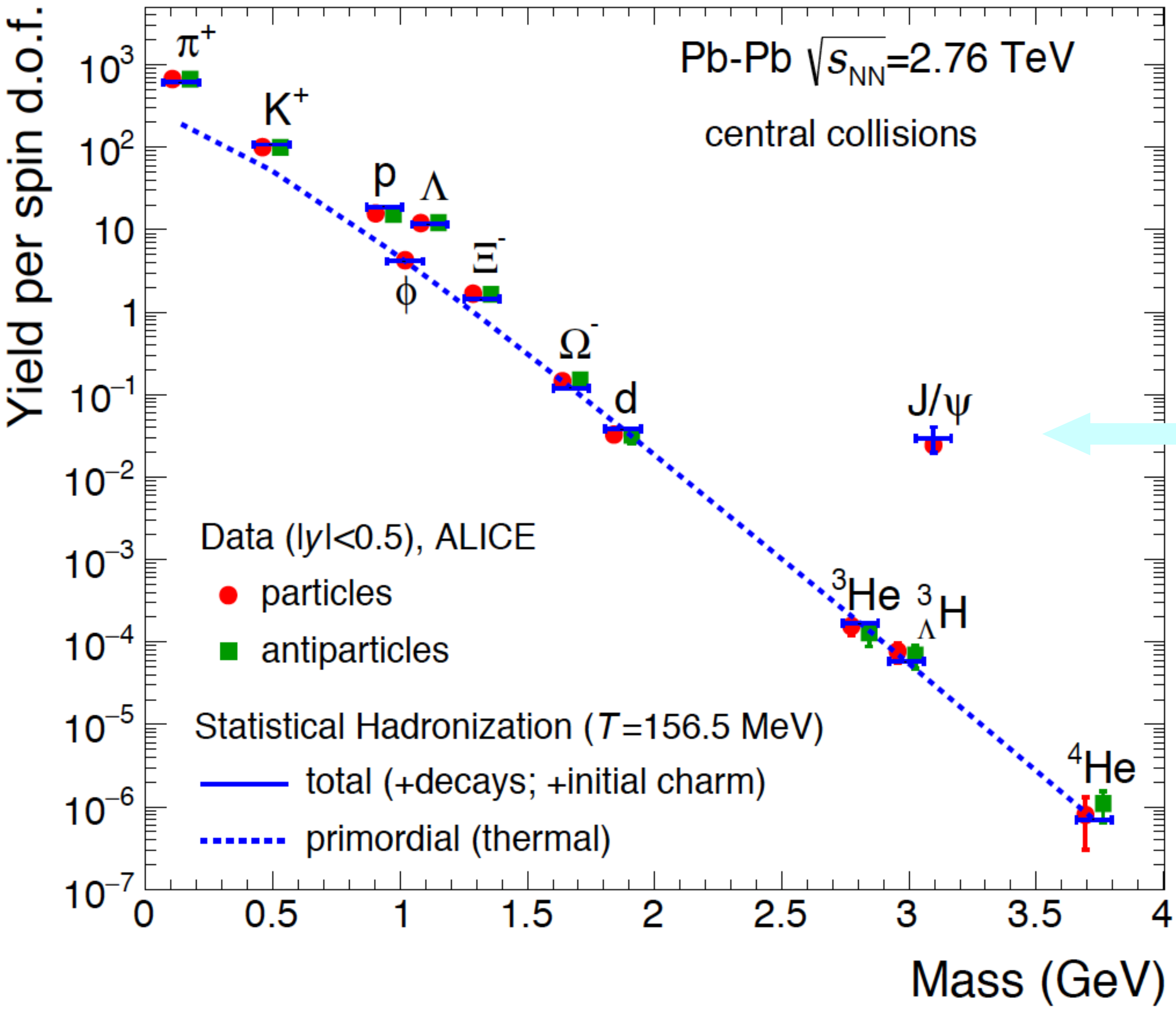
energy density →

melting scenario not observed
rather: enhancement with increasing energy density!
(from RHIC to LHC and from forward to mid-rapidity)



J/ψ overpopulation due to hard production of charm and statistical hadronization of deconfined quarks

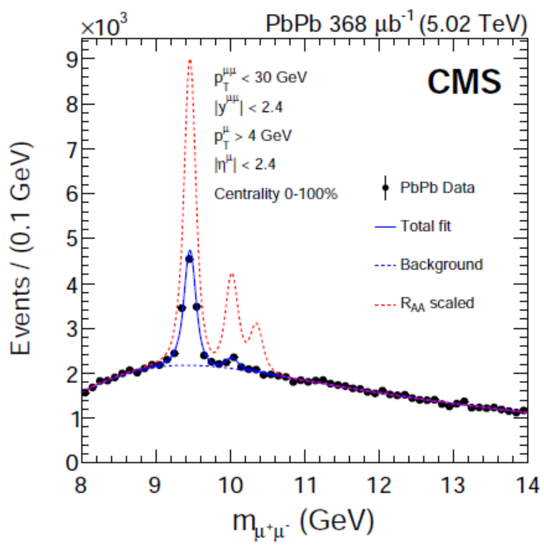
A.Andronic et al., PLB 797 (2019) 134836

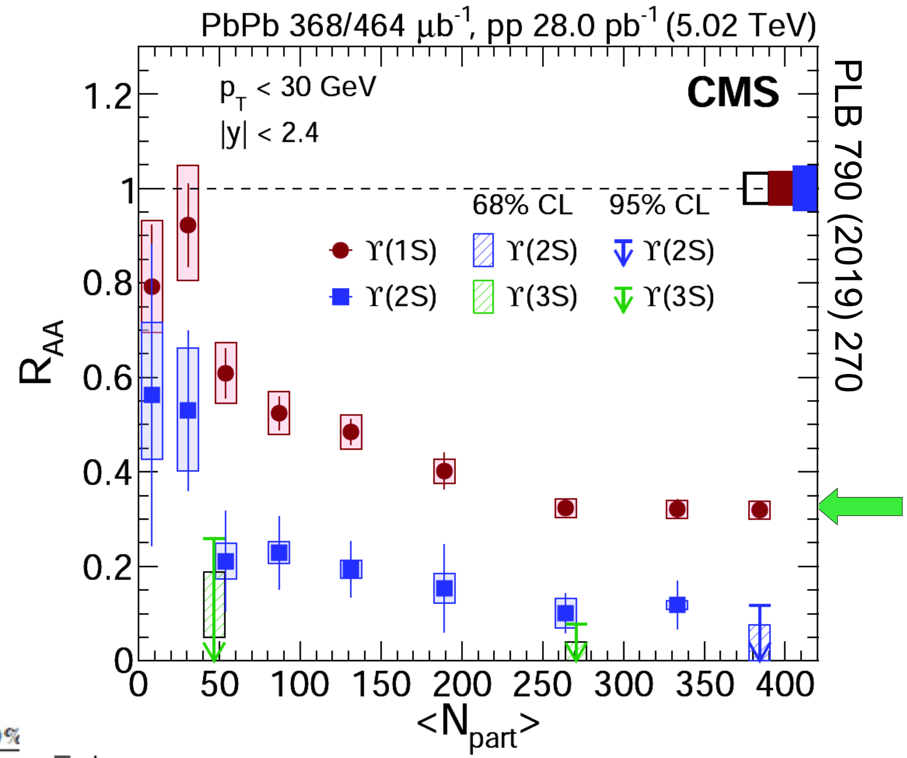


J/ψ enhanced compared to other M = 3 GeV hadrons by factor g_c^2 = 900 relative to purely thermal yield

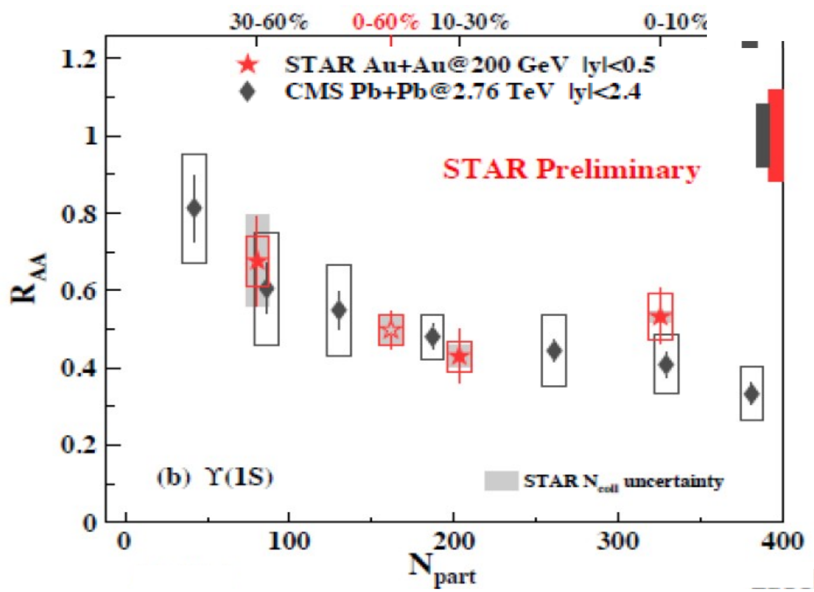
quantitative agreement with hadronization of deconfined thermalized charm quarks

Suppression of Upsilon states





not consistent with just excited state suppression (LHCb data: only 25 % feed-down in pp at LHC)



genuine ϒ suppression
- real and imaginary part of potential at finite temperature play a role
- similarity of RHIC and LHC suppression reminiscent of SPS and RHIC for J/ψ
- possibility of statistical hadronization?

Outlook – an incomplete, but maybe realistic wishlist for the coming decade

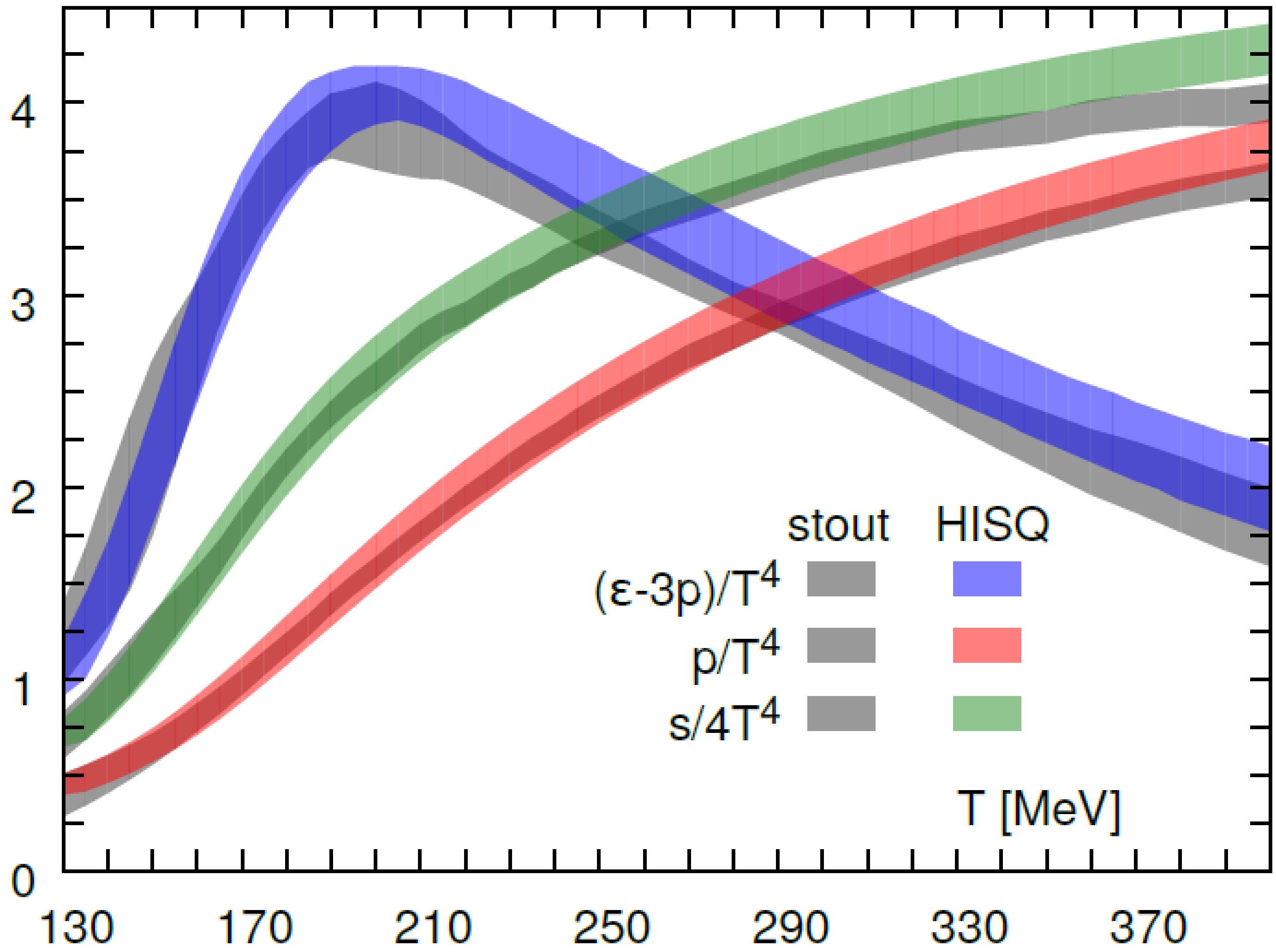
Important measurements to come:
- open charm cross section in PbPb down to pt=0, baseline for J/ψ
- higher charmonium states to determine T for deconfinement transition
- to what degree does beauty thermalize in QGP, baseline for Upsilon understanding
- direct photons (real and virtual) and their azimuthal asymmetries with larger significance, thermal evolution of QGP
- low mass dilepton pairs and rho spectral function, chiral symmetry restoration
- fluctuations of conserved charges as sign of critical behavior
  - at LHC due to proximity to O(4) critical region
  - at lower energies due to possible critical endpoint in phase diagram

Tasks for theory:
- the way to thermalization: from overpopulated gluon fields to hydrodynamics to hadronization
- determination of temperature dependent transport coefficients from exp. data
- first principles computation of transport coefficients

backup

Alternative for lattice QCD EoS

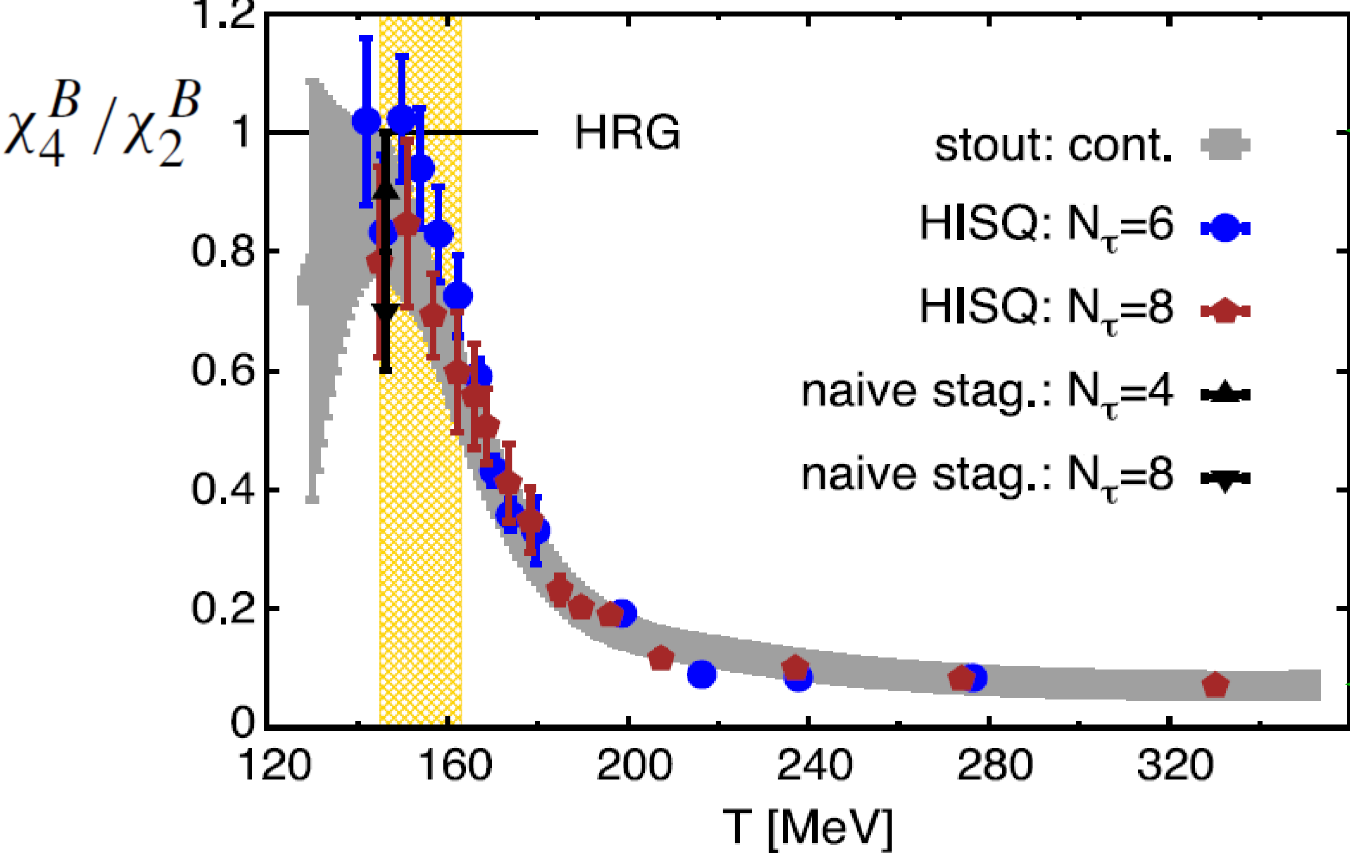
from Bazavov arXiv: 1407.6387



Measure of deconfinement in lQCD

$\chi_4^B/\chi_2^B \propto$ baryon number^2

H.T.Ding, NPA931 (2014) 52



confined: 1

measure suggested by Ejiri, Karsch, Redlich (2006)

deconfined: $6/9\pi^2$

rapid drop suggests: chiral cross over and deconfinement appear in the same narrow temperature range

Analysis of hadron yields: the statistical model – grand canonical

partition function: $$\ln Z_i = \frac{V g_i}{2\pi^2} \int_0^\infty \pm p^2 dp \ln(1 \pm \exp(-(E_i - \mu_i)/T))$$

particle densities: $$n_i = N/V = -\frac{T}{V} \frac{\partial \ln Z_i}{\partial \mu} = \frac{g_i}{2\pi^2} \int_0^\infty \frac{p^2 \, \mathrm{d}p}{\exp((E_i - \mu_i)/T) \pm 1}$$

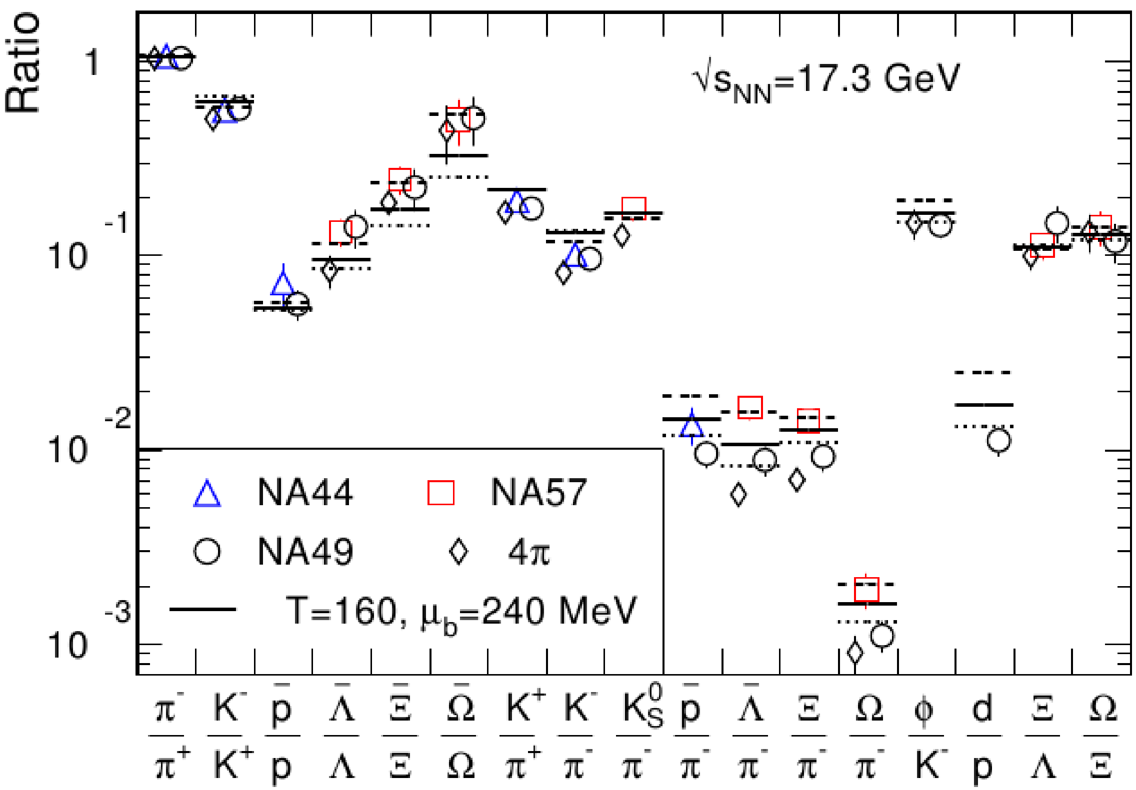
for every conserved quantum number there is a chemical potential:

$$\mu_i = \mu_B B_i + \mu_S S_i + \mu_{I_3} I_i^3$$

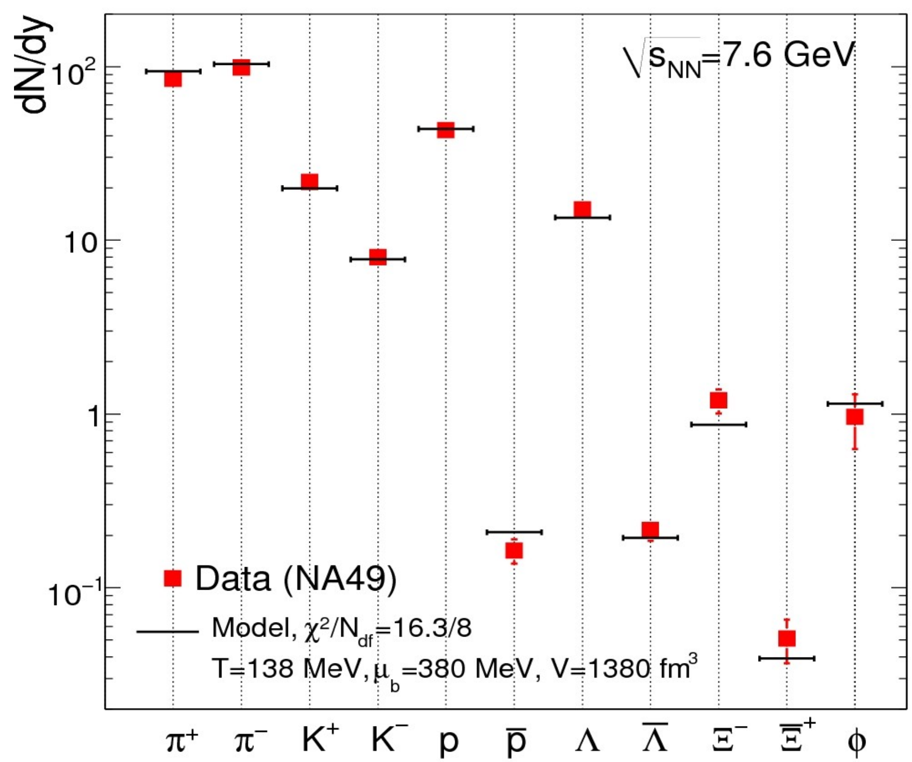
but can use conservation laws to constrain V, μ_S, μ_{I_3}

→ Fit at a given √s provides values for T and μ_b

SPS Pb + Pb data and thermal model



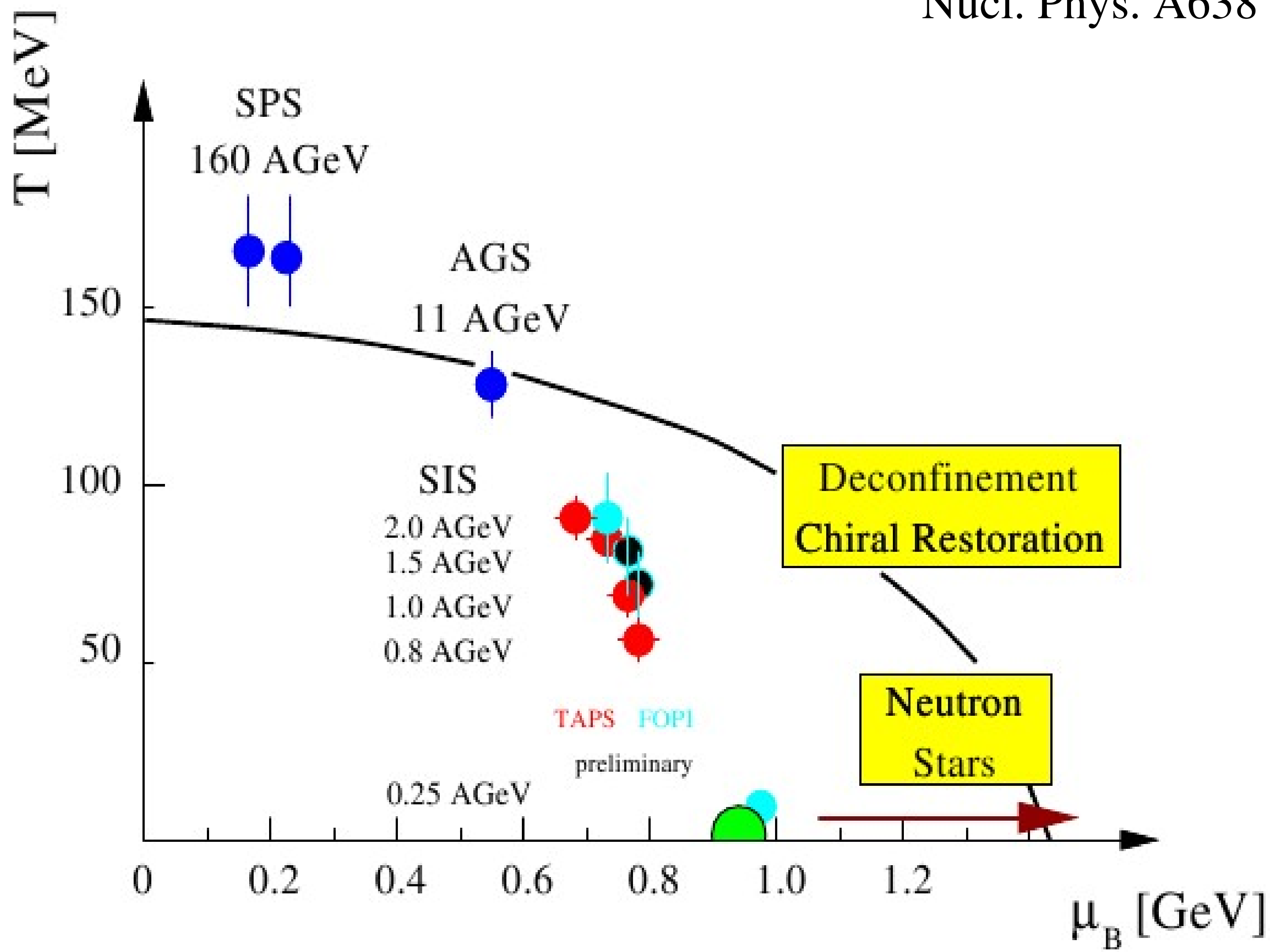
full energy – largest amount of data but also some problems in data revealed



figures from A. Andronic, P. Braun-Munzinger, J.S. Nucl. Phys. A772 (2006) 167

First phase diagram with experimental points

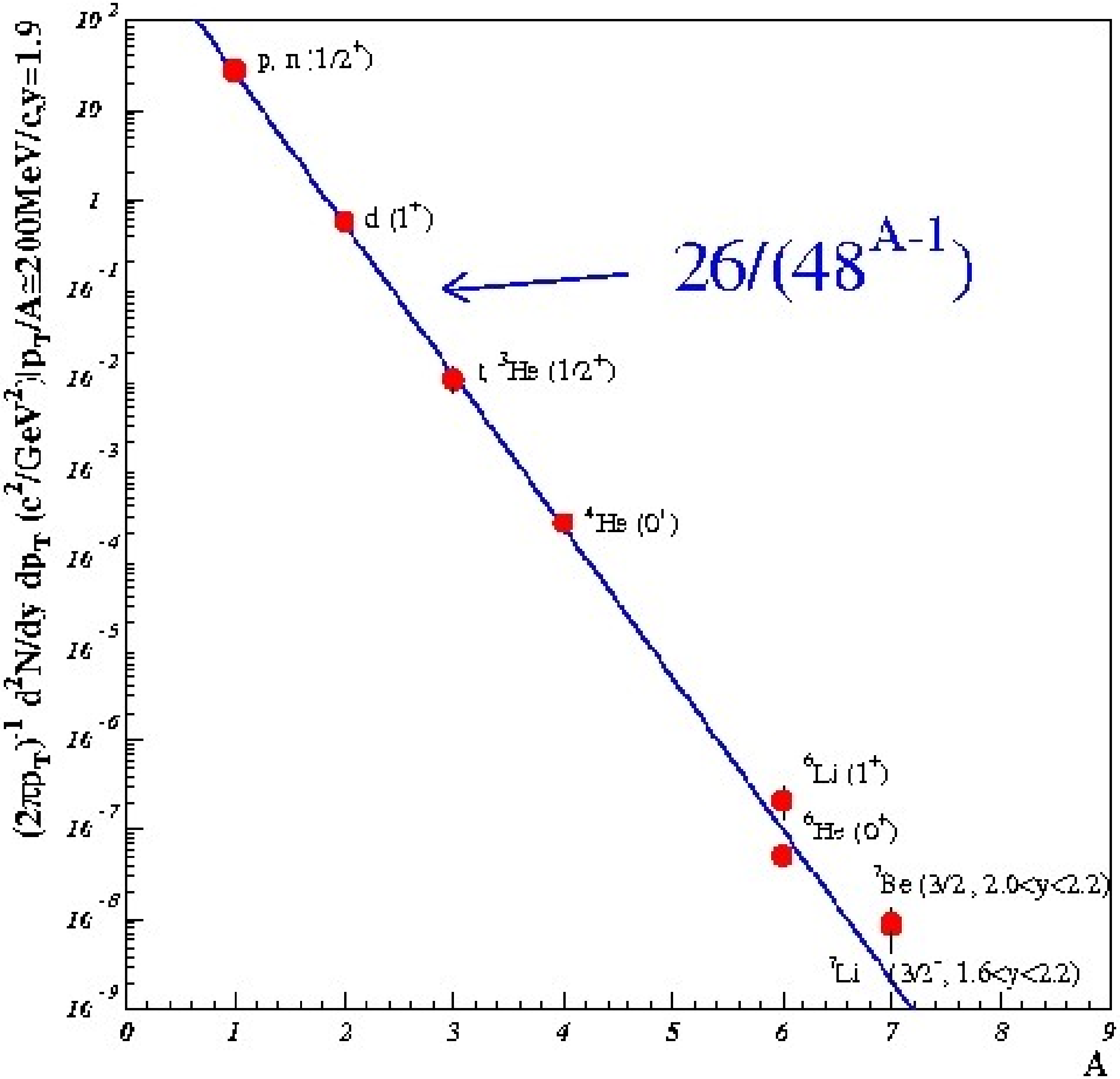
P.Braun-Munzinger and J. Stachel, nucl-th/9803015,
Nucl. Phys. A638 (1998) 3



Production of light nuclei and antinuclei at the AGS

E864 Coll., Phys. Rev. C61 (2000) 064908

<u>data cover 10 oom!</u>
addition of every nucleon
-> penalty factor $R_p = 48$
but data are at very low pt
use m-dependent slopes following
systematics up to deuteron
-> $R_p = 26$



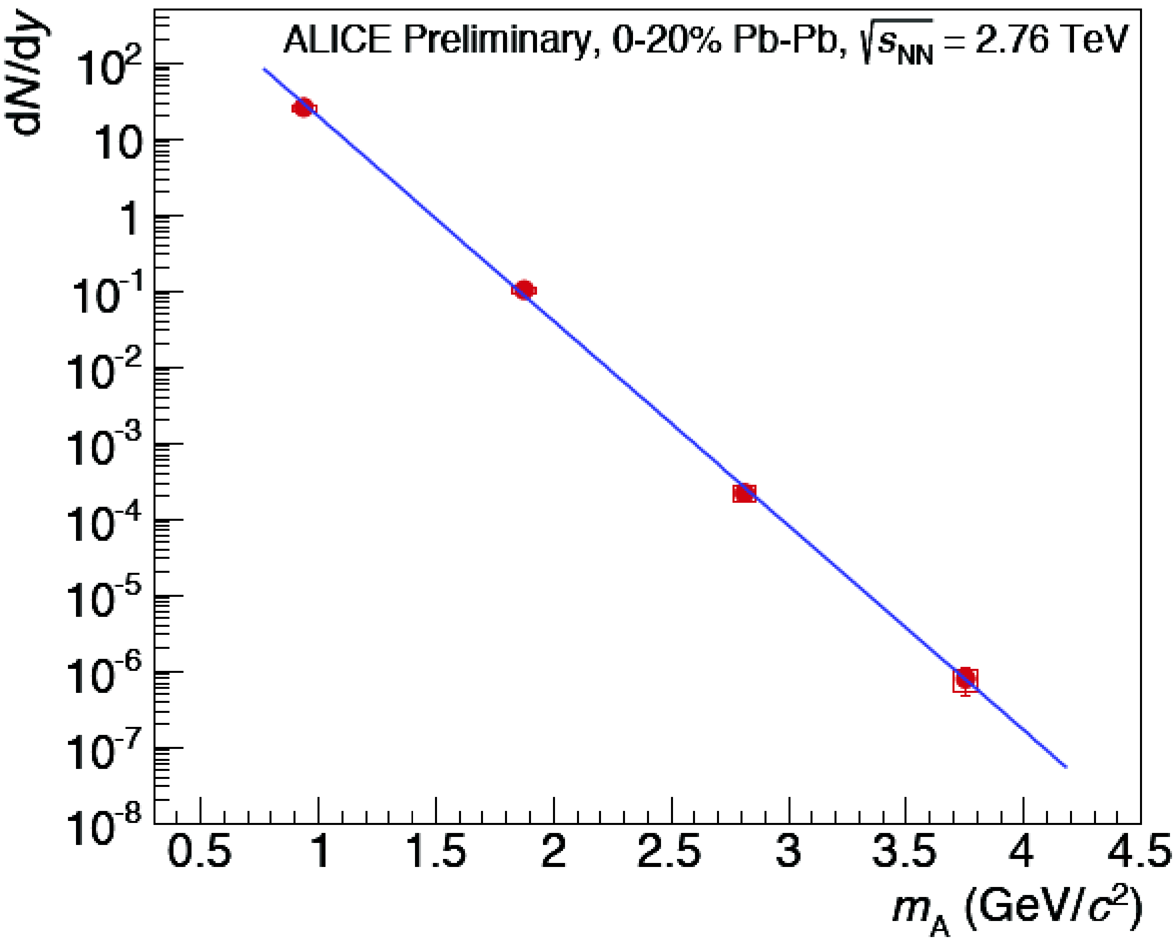
<u>GC statistical model:</u>

$R_p \cong \exp[(m_n \pm \mu_b)/T]$

for T=124 MeV and μ_b=537 MeV
$R_p = 24$ good agreement
also good for **antideuterons**:
data: $R_p = 2 \pm 1 \cdot 10^5$ SM: $1.3 \cdot 10^5$

P. Braun-Munzinger, J. Stachel,
J. Phys. G28 (2002) 1971

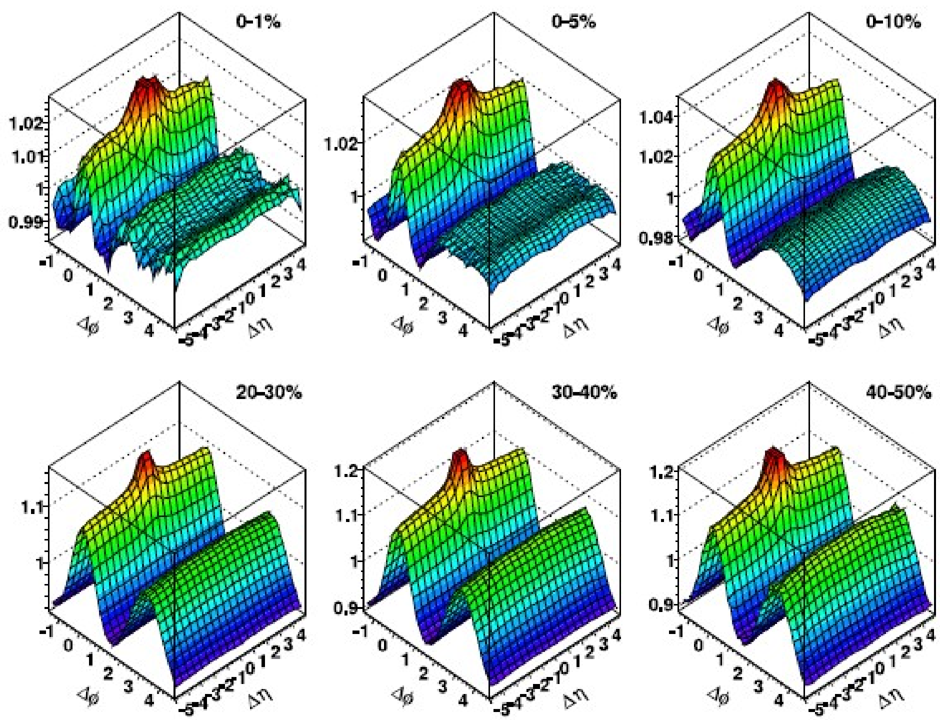
Production of light anti-nuclei at LHC energy

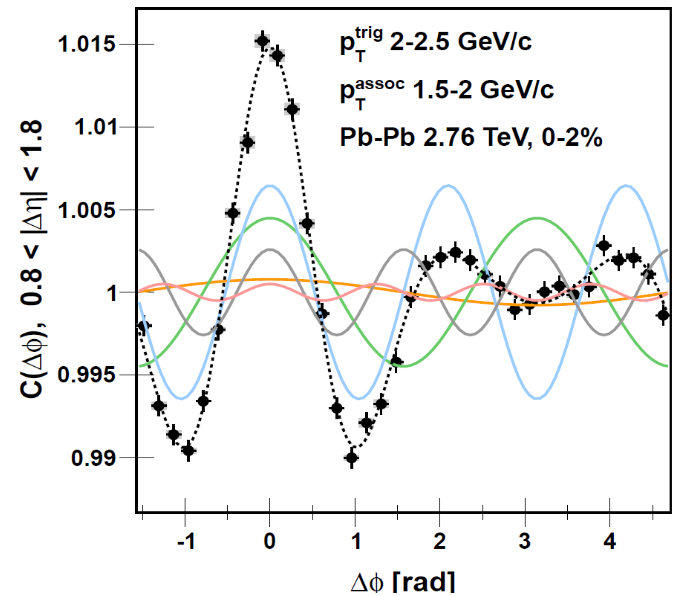


penalty factor exp(-m/T) ≈ 300 for nuclei and anti-nuclei as $\mu_b = 0$ at LHC
compared to 24 for nuclei at top AGS energy
and 140 000 for anti-nuclei with $\mu_b = 537$ and T=124 MeV

Propagation of sound in the quark-gluon plasma

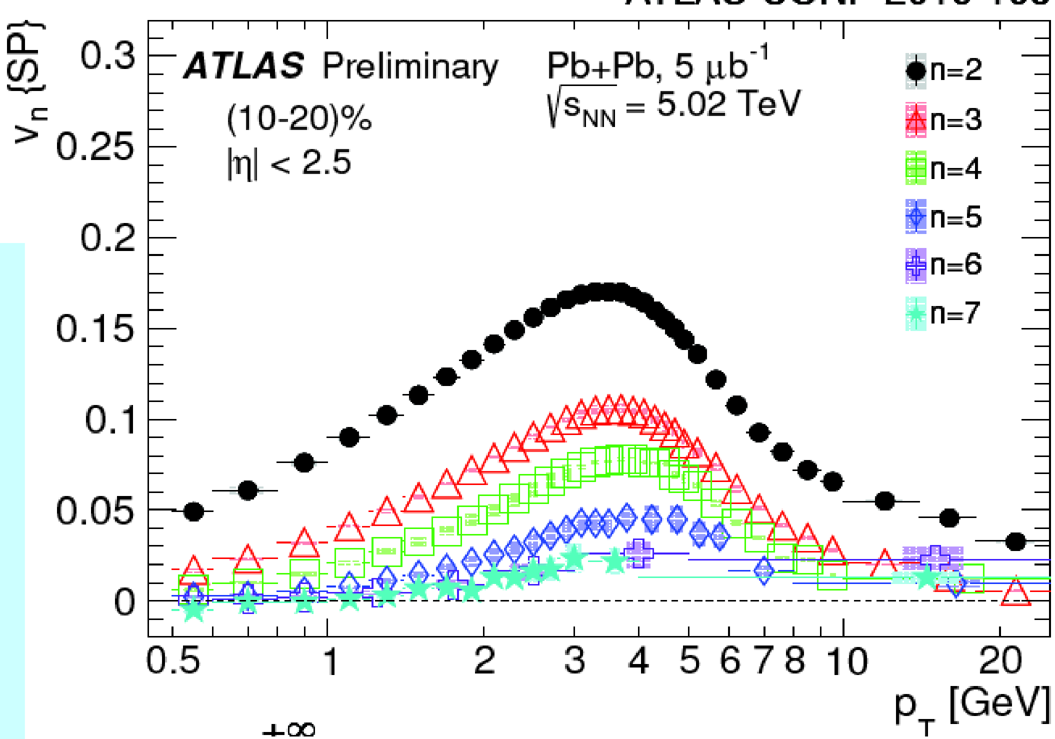
ATLAS-CONF-2011-074





ALICE, PLB 708 (2012)249

ATLAS-CONF-2016-105



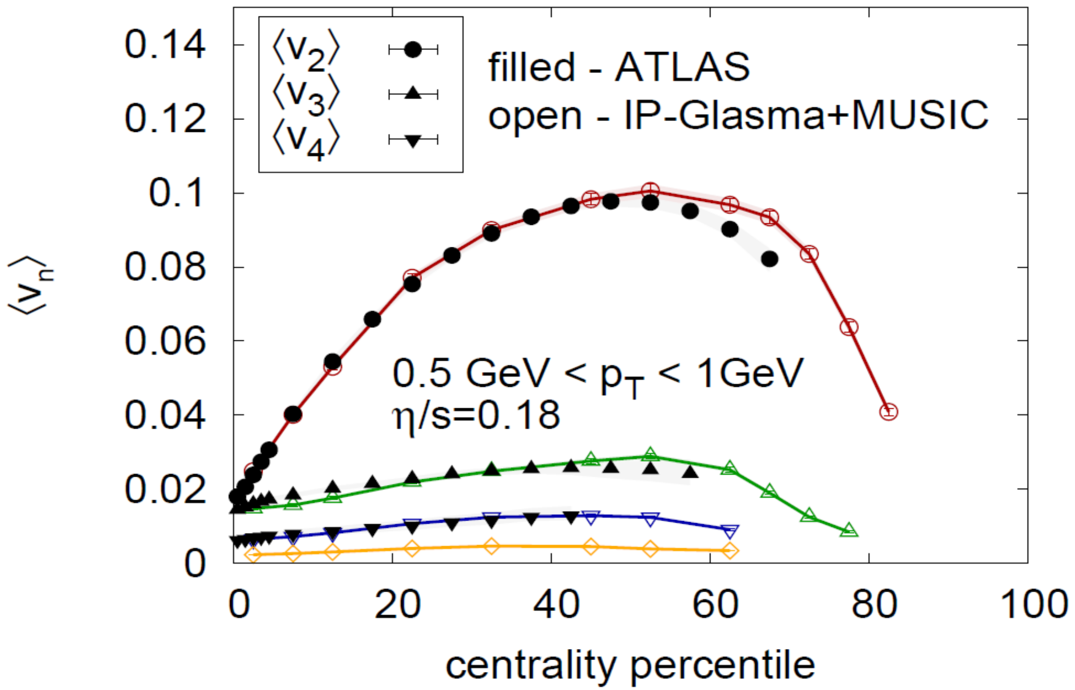
long-range rapidity correlations

understanding: higher harmonics (3,4,5,...) are due to initial inhomogeneities caused by granularity of binary parton-parton collisions survive the 10 fm/c hydrodynamic expansion phase

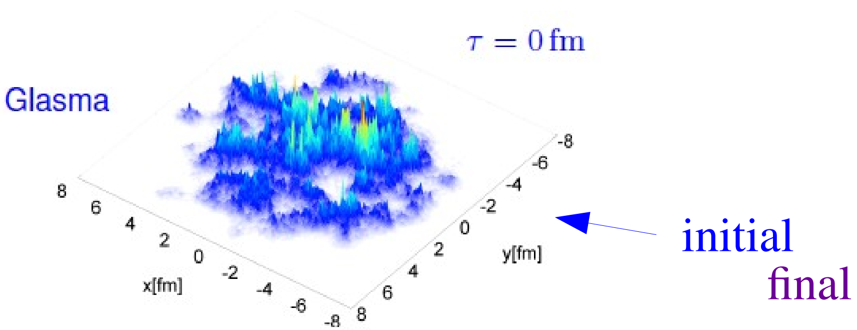
M. Luzum PLB 696 (2011) 499

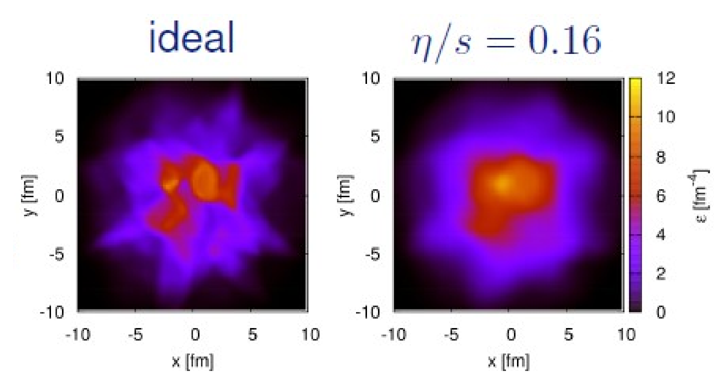
Higher flow harmonics and their fluctuations

data: ATLAS JHEP 1311 (2013) 183
calc: B. Schenke, R. Venugopalan, Phys. Rev. Lett. 113 (2014) 102301



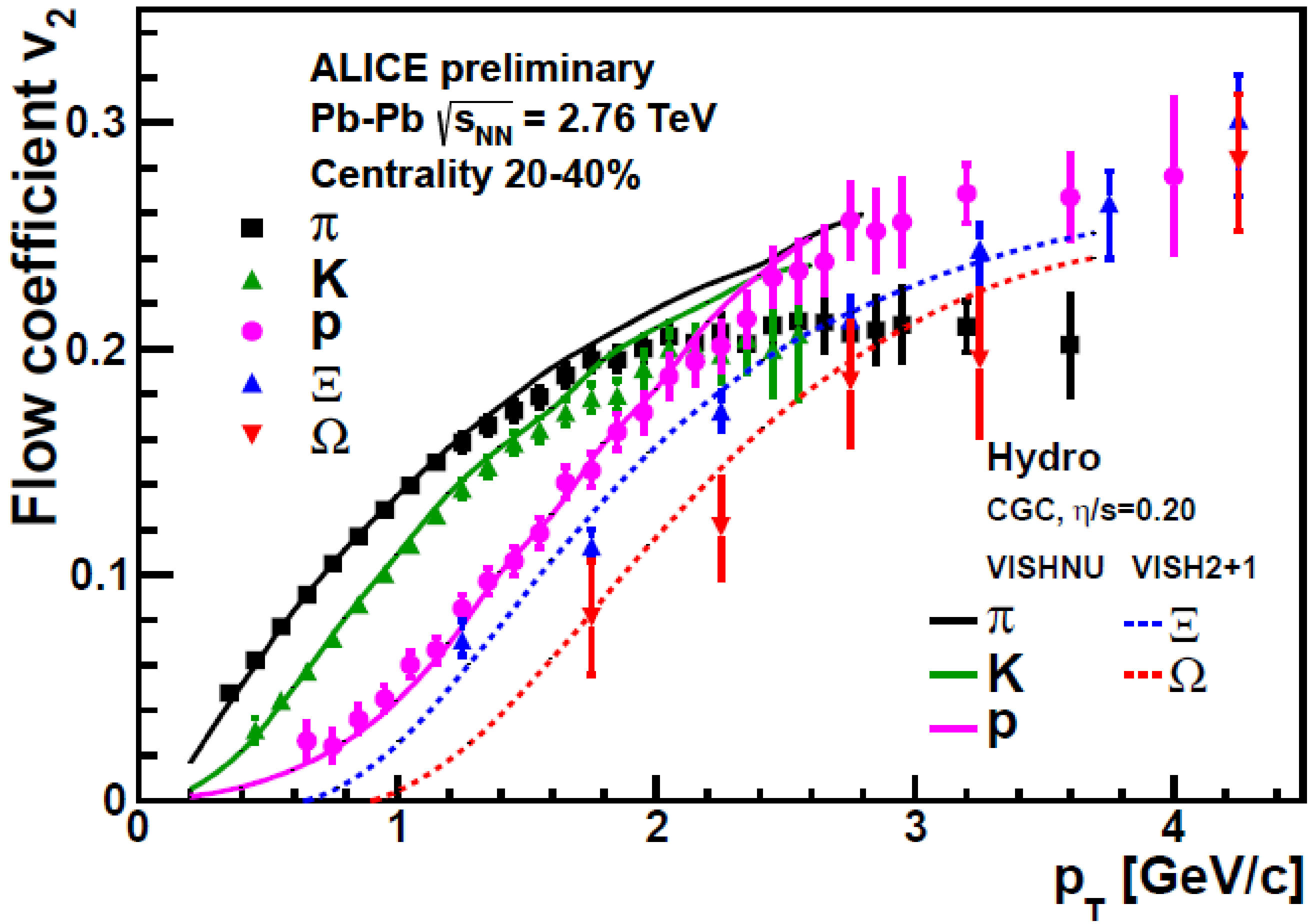
ratios of v_2/v_n and their fluctuations depend on initial condition





very well reproduced by viscous hydrodynamics (MUSIC) with fluctuating IP Glasma initial condition (including initial quantum fluctuations of gluon fields)
for LHC $\eta/s = 0.18$ for RHIC $\eta/s = 0.12$
indication of temperature dependence of η/s?

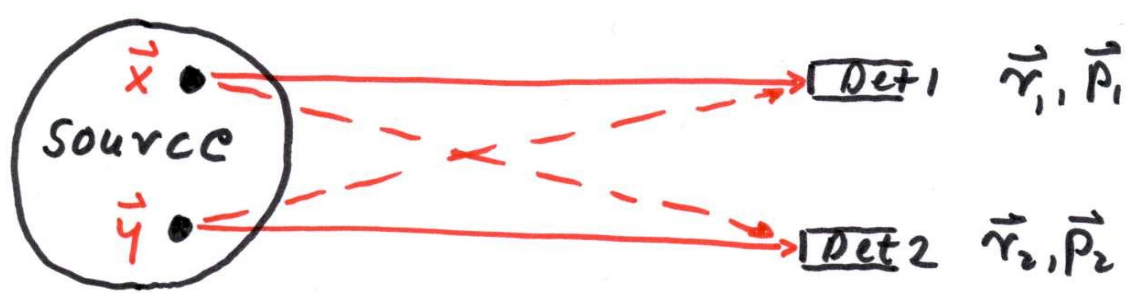
Elliptic Flow in PbPb Collisions at $\sqrt{s_{NN}}$ = 2.76 TeV



rapidly rising v_2 with p_t and mass ordering typical features of hydrodyn. expansion
same hydrodynamics calc. with small eta/s reproduces data

Bose-Einstein correlations and space-time extent of fireball

stochastic emission from extended source
consider 2 identical bosons (photons, pions, …)
2 detectors in locations r_1, r_2 observe identical bosons of momenta p_1 and p_2



cannot distinguish solid and dashed paths because of identical particles
for plane waves, the probability amplitude for detection of the pair is

$$A_{12} = \frac{1}{\sqrt{2}}[e^{ip_1(r_1-x)}e^{ip_2(r_2-y)} + e^{ip_1(r_1-y)}e^{ip_2(r_2-x)}]$$

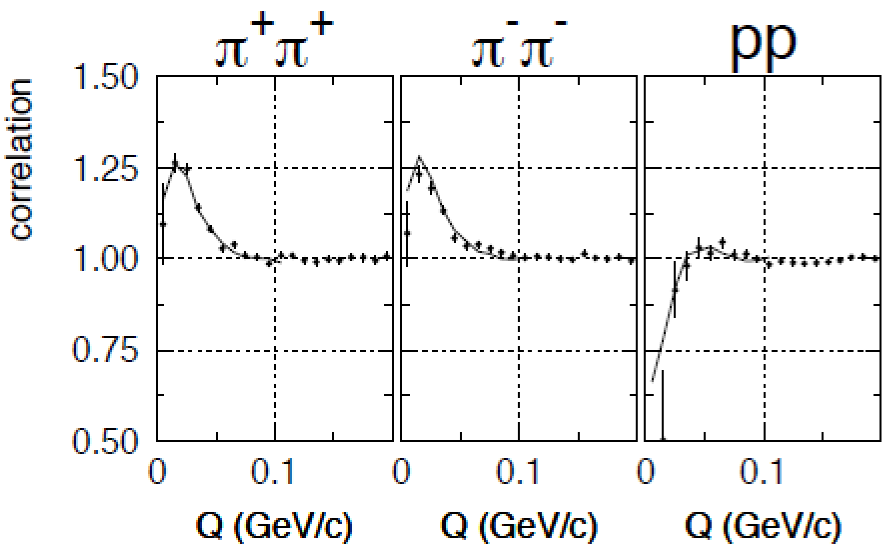
with 4-vectors p,r,x,y (to be general for nonstatic source)

square of amplitude: intensity $\longrightarrow$ "intensity interferometry"

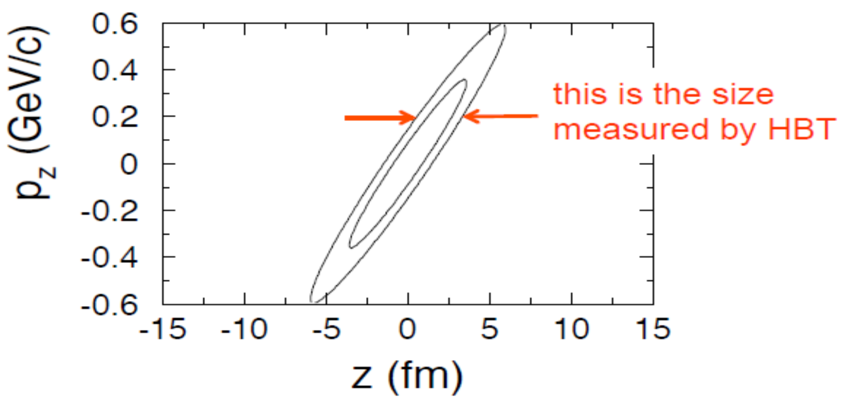
technique of intensity interferometry developed by Hanbury-Brown and Twiss in astrophysics as a means to determine size of distant objects

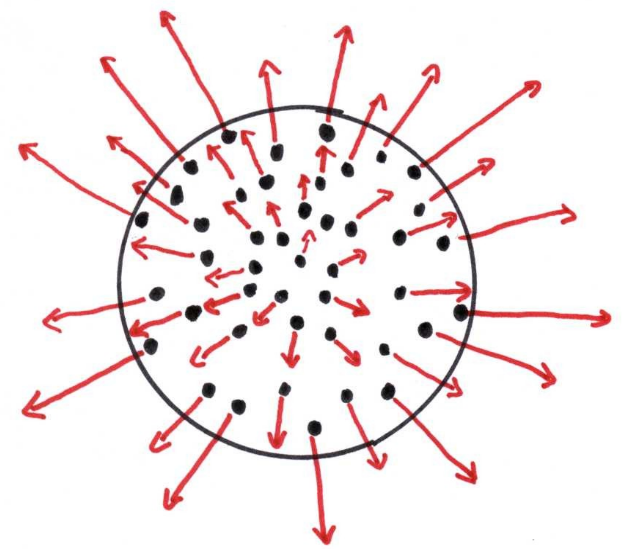
Hanbury-Brown/Twiss correlations to measure the space-time extent of the fireball

Au + Au at 10.8 A GeV
E877 data compared to RQMD



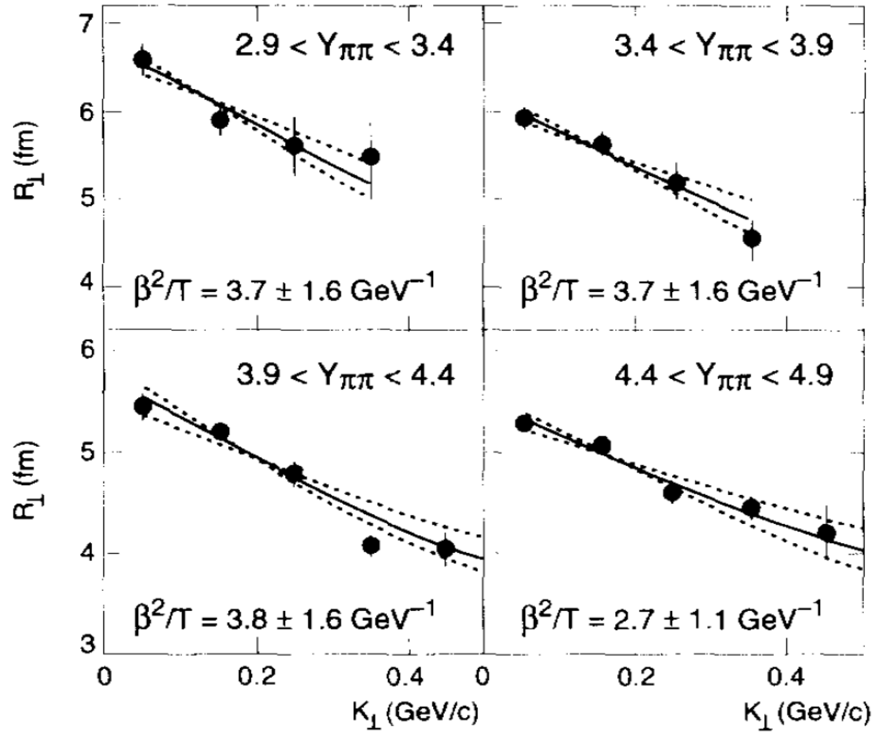
first a puzzle: small apparent radii (2 fm/c)
then a discovery: are due to collective expansion of fireball
space – momentum correlations → only part of the source is 'visible'
predicted by Mahklin/Sinyukov





2-Pion Hanbury-Brown/Twiss correlations → Radius Parameters as Function of Pair Transverse Momentum

158 A GeV PbPb – NA49 – Nucl. Phys. A638 (1998) 91c

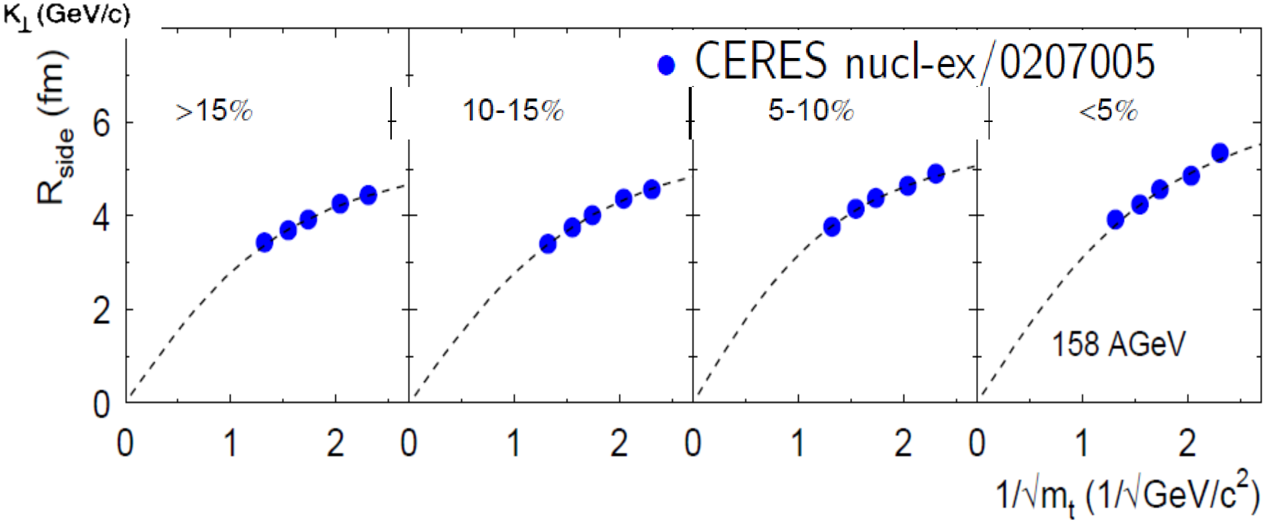


transverse mom. dependence shows typical shape for hydrodynamically expanding source

$$R_{\rm side} \approx R_{\rm geo}/(1 + m_t \cdot F(T_f, \beta_t))^{\frac{1}{2}}$$

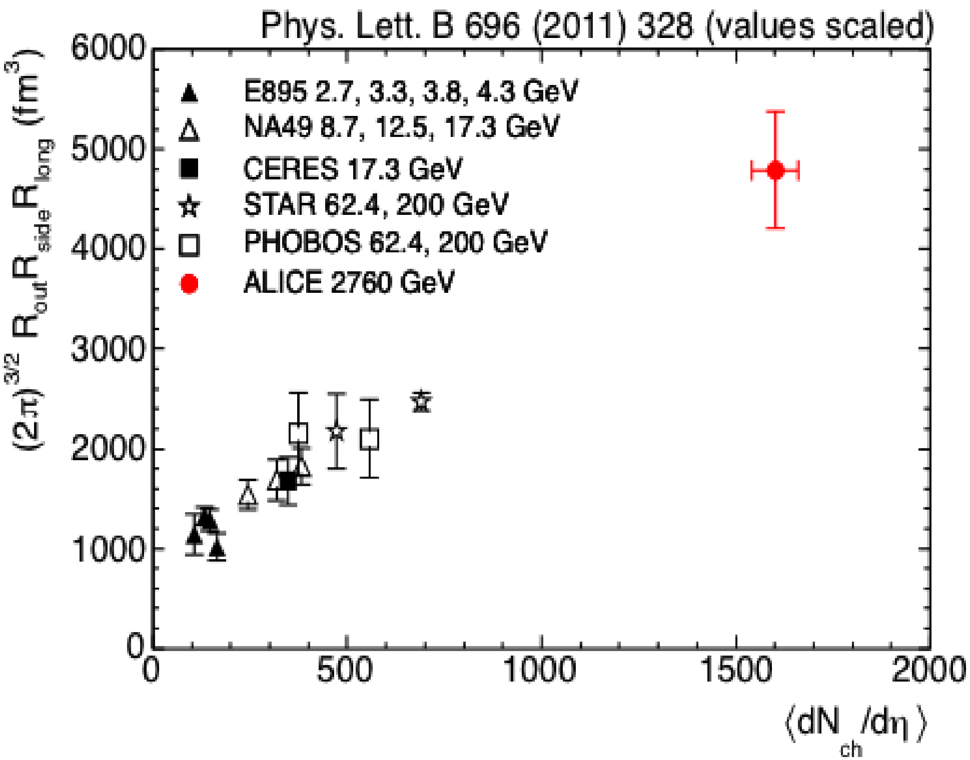
U. Heinz *et al.*

$\beta_t \approx 0.55$ for $T_f = 120$ MeV



Freeze-out volume and duration of expansion

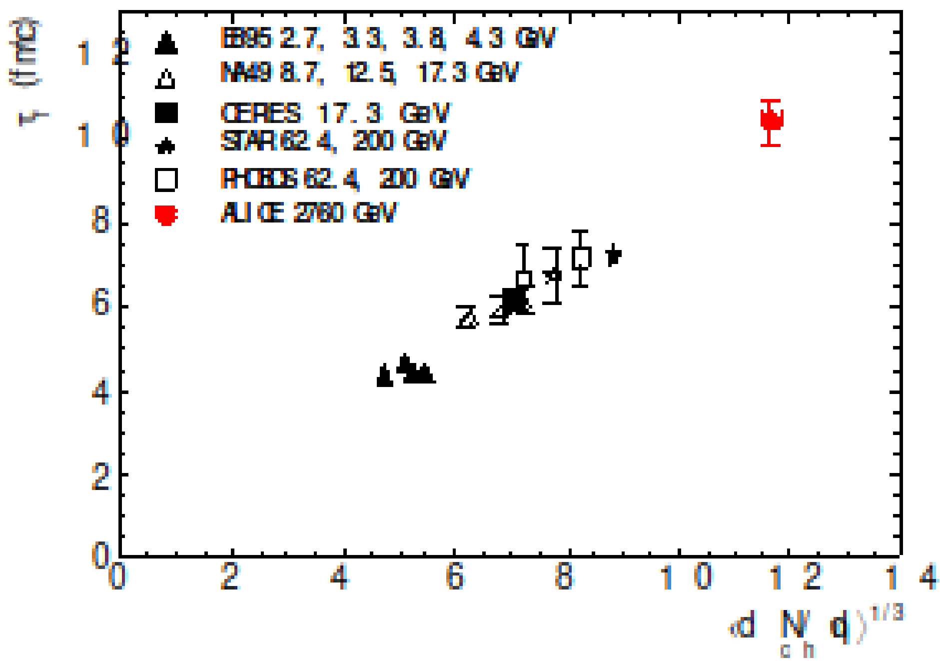
coherence volume $V = (2\pi)^{3/2} R_{side}^2 R_{long}$



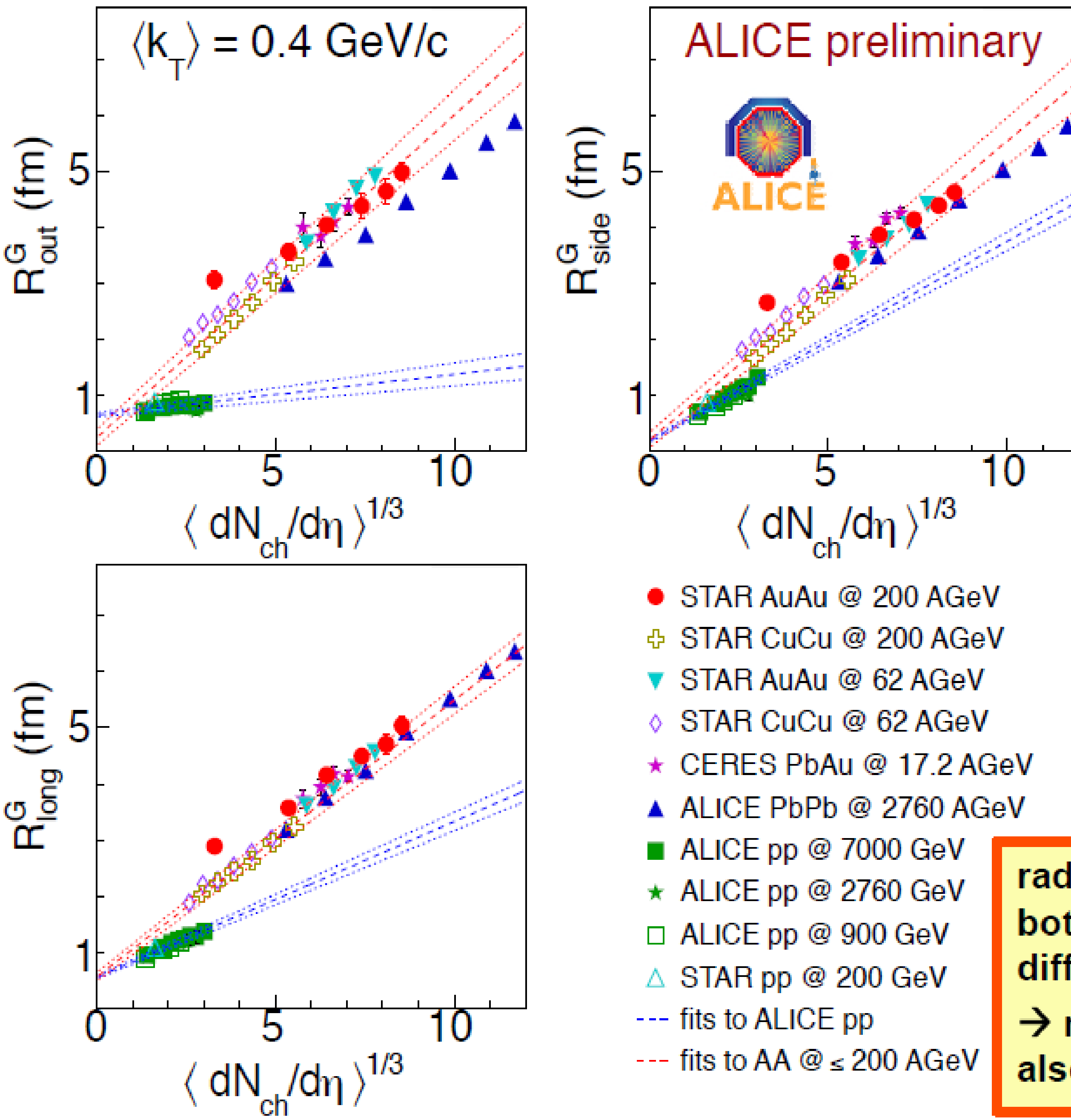
huge growth with √s at all energies
larger than overlap volume – reflects
strong expansion of fireball
at surface at LHC velocity ¾ c

from R_{long}: duration of expansion
4.5 fm/c at AGS to 10 fm/c at LHC

$$R_{long} = \tau_f \sqrt{T/m_t}$$



pion HBT

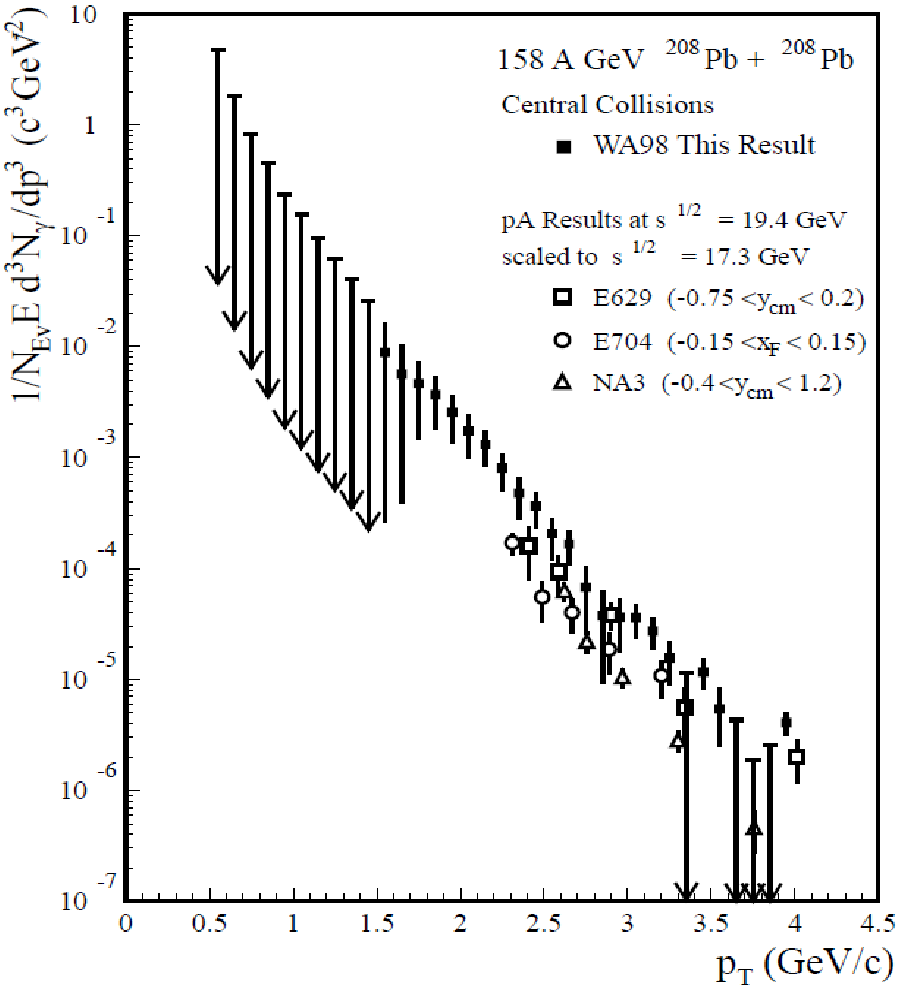


radii increase with multiplicity both in pp and Pb-Pb but with different slopes

→ not only final multiplicity but also initial geometry matters

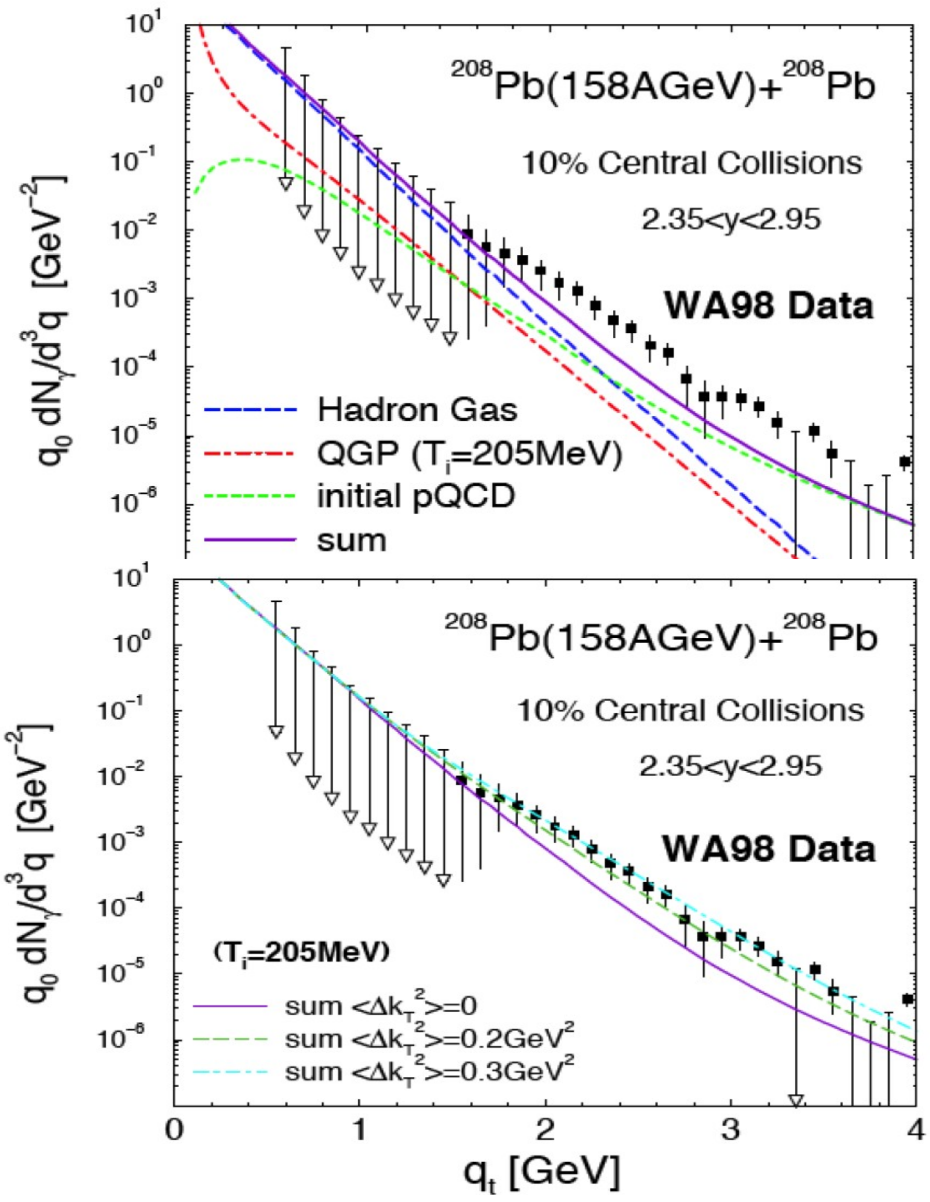
Direct photons: give access to entire time evolution

WA98



$\lambda_{mfp} \gg$ medium
$\rightarrow$ access to early QGP-phase

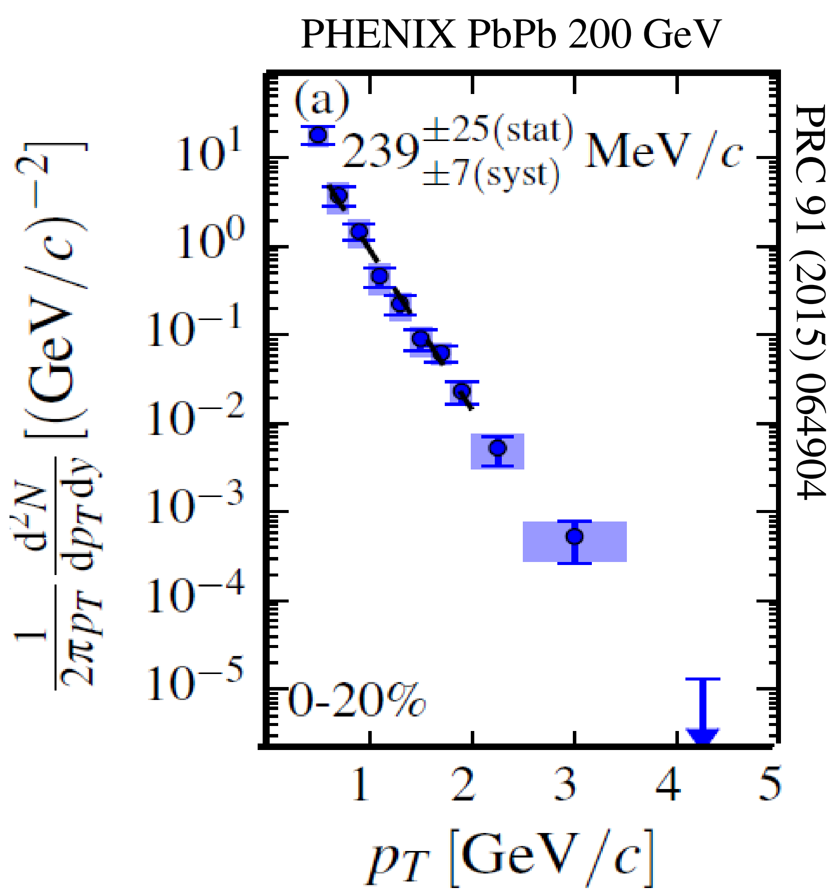
Direct photons: give access to entire time evolution

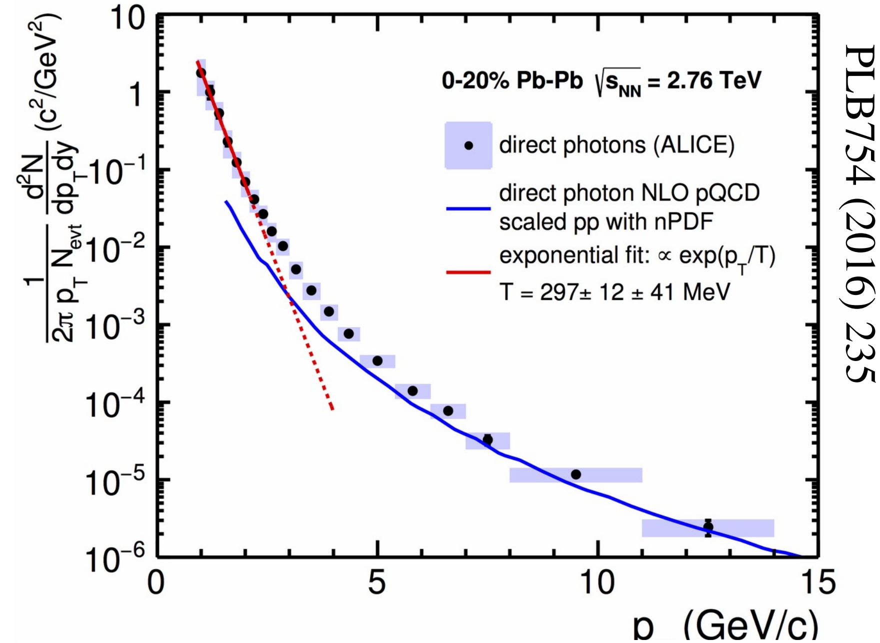


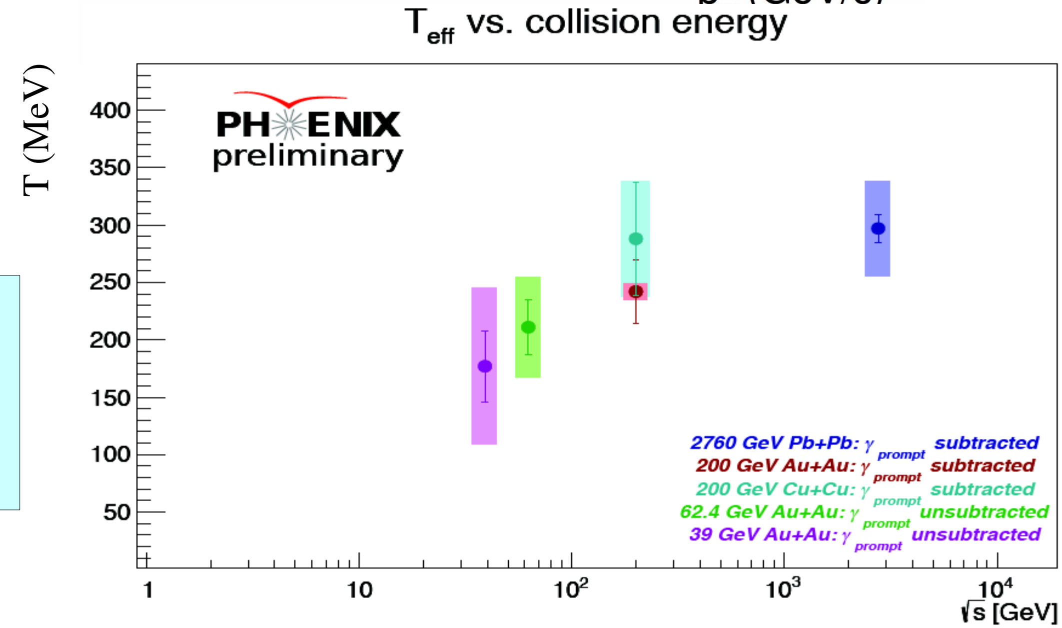
$\lambda_{mfp} \gg$ medium
→ access to early QGP-phase

first significant measurement in PbPb collisions: WA98 at SPS
- data consistent with QGP formation (T_i = 200-270 MeV)
- but also purely hadronic scenario w. Cronin enhancement accounts for data

Direct photons at RHIC and LHC





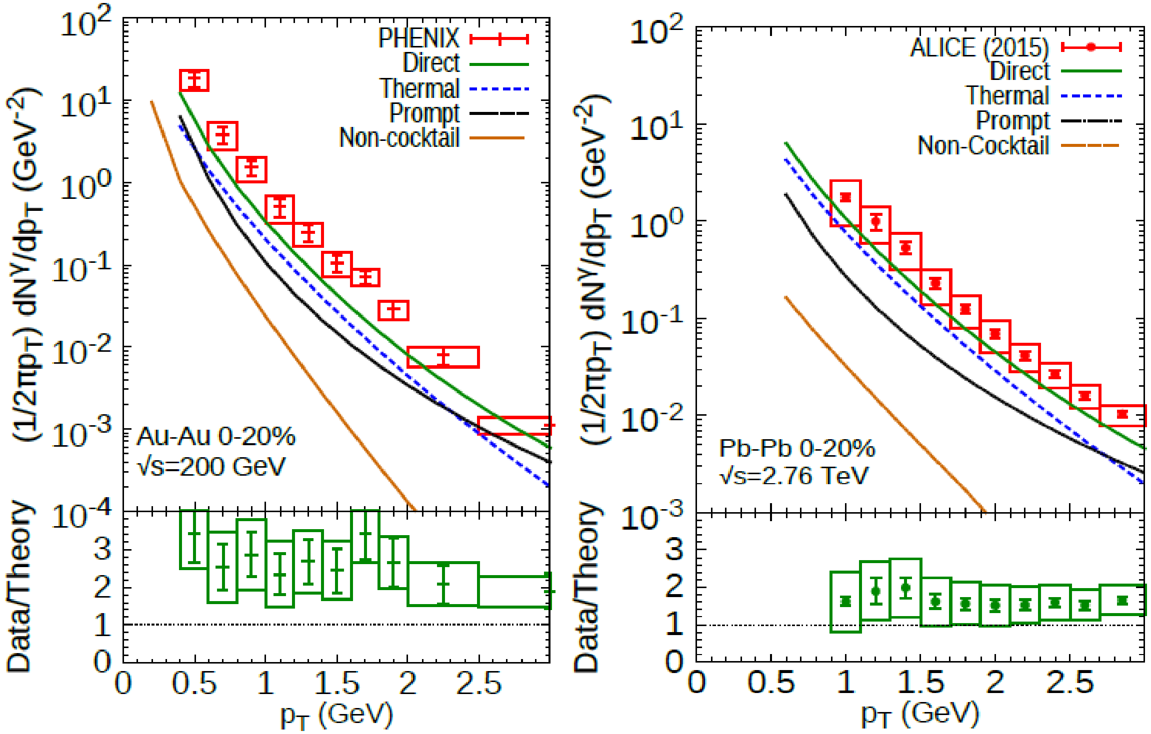


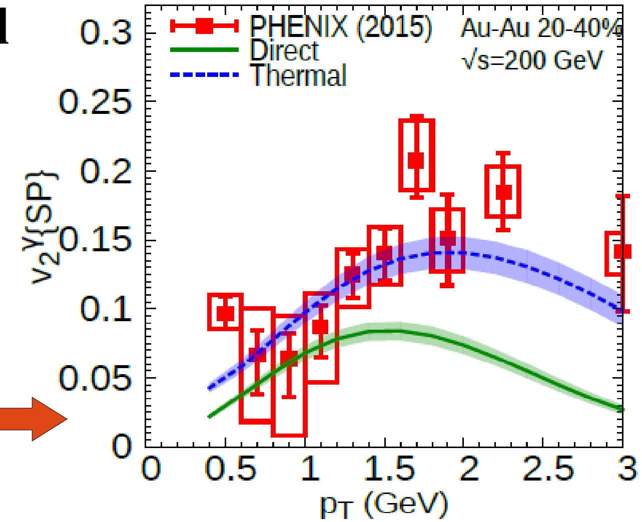
direct photons of fireball exhibit higher apparent temperatures with increasing √s

Direct photons at RHIC & LHC exhibit strong elliptic flow

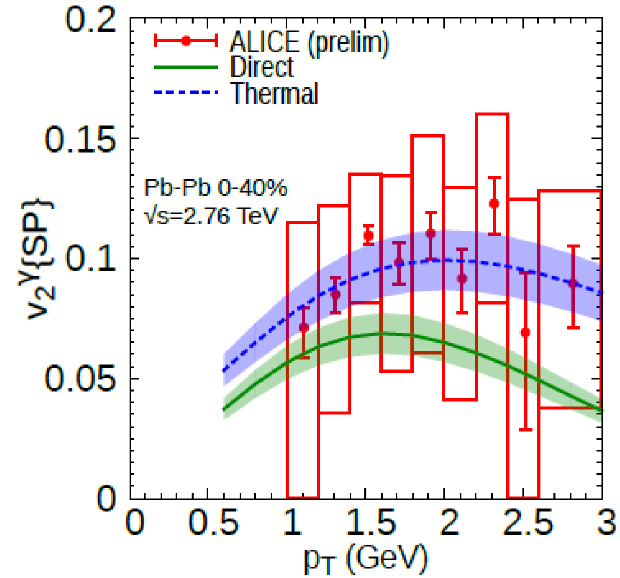
photon radiation of hydrodynamically expanding fireball

theory: J.F. Paquet et al., PRC93 (2016) 044906





PRC94 (2016) 064901

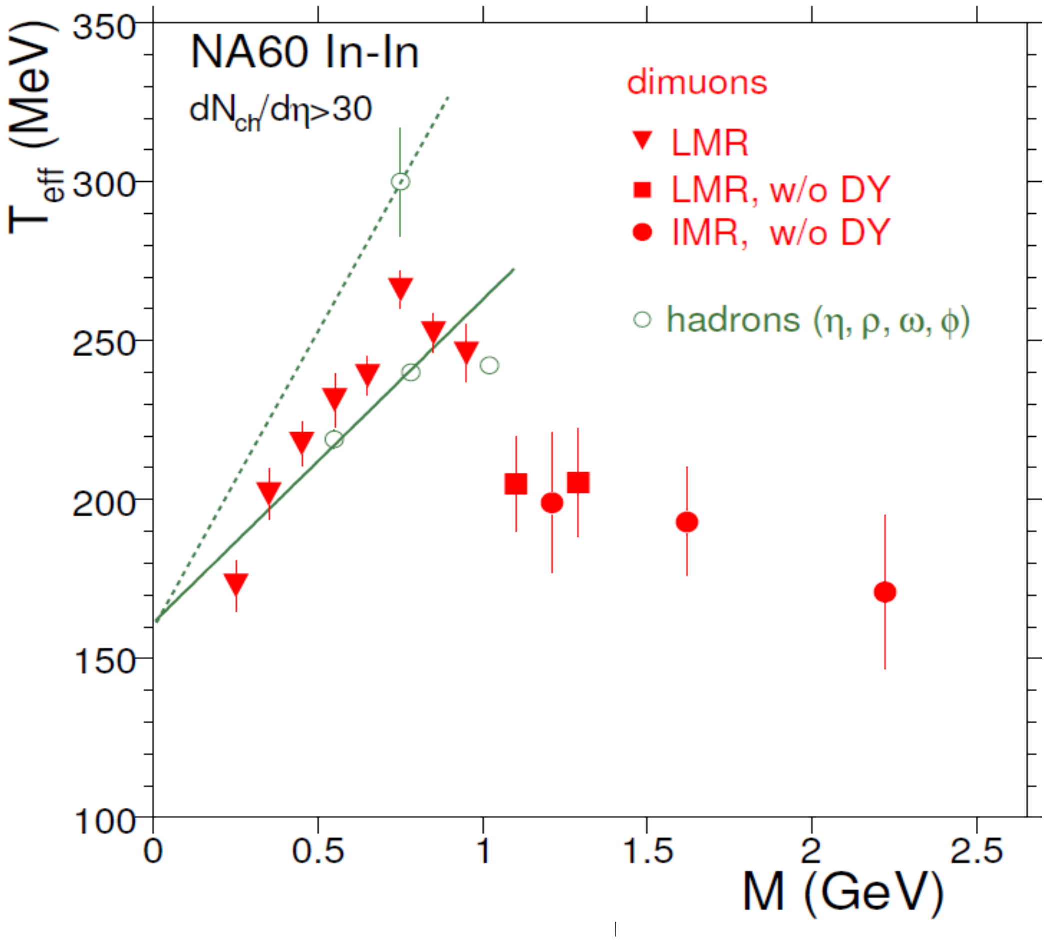


PLB754 (2016) 235

direct-photon puzzle: challenge of simultaneous description of spectra and v_2

rate $\propto$ T^2 - but flow takes time to evolve

Low and intermediate mass lepton pairs

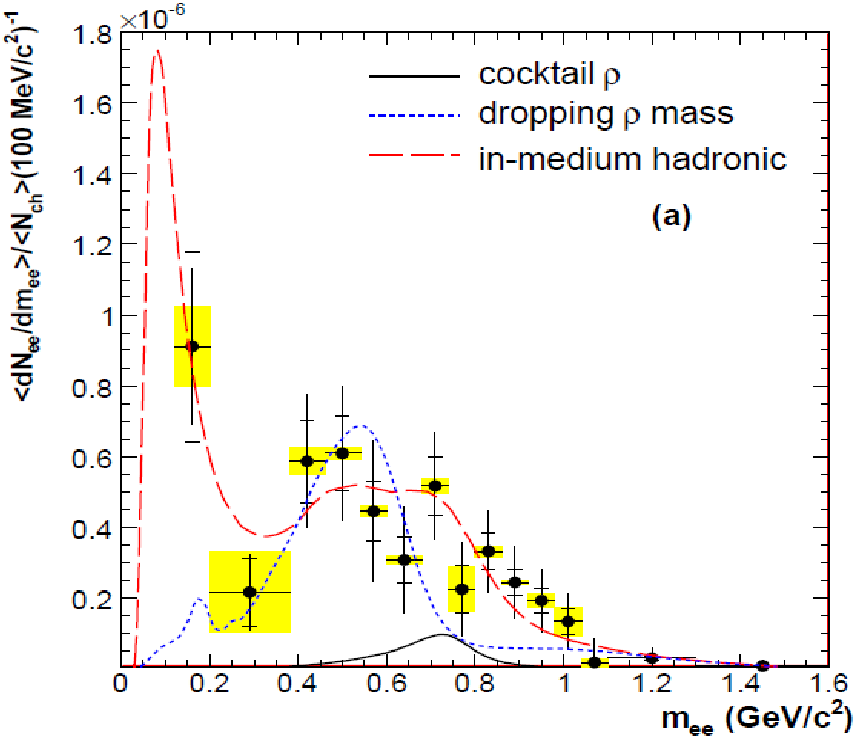


EPJ C61 (2009) 711

- up to mass ≈1.0 GeV: radial flow of a hadron-like di-lepton source
- above: thermal component with T = 205±12 MeV
- virtual photons vs real photons above

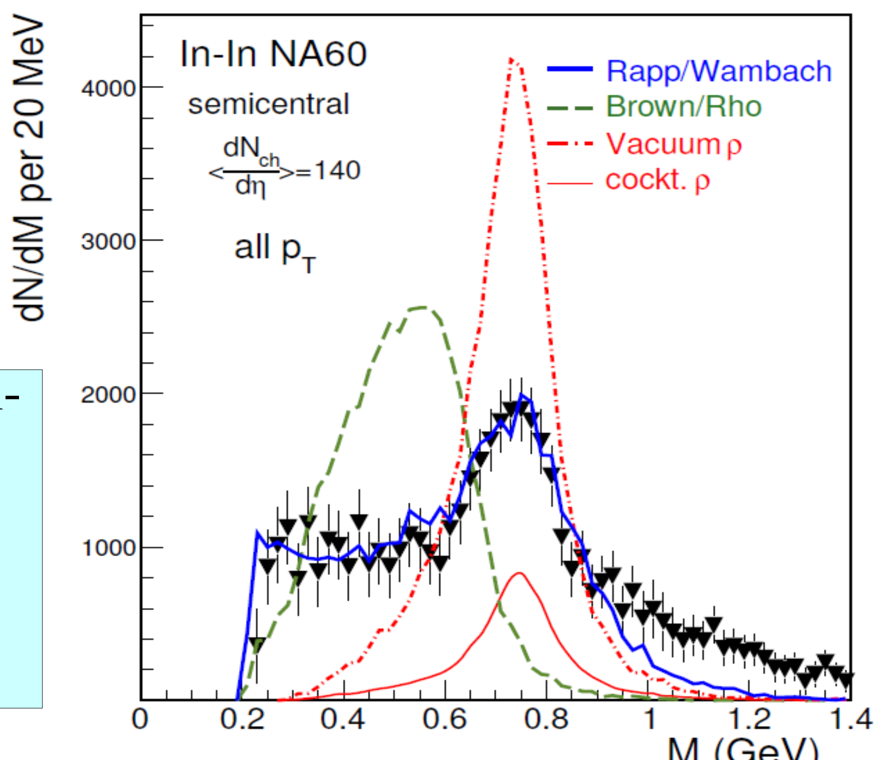
Low and intermediate mass lepton pairs

158 A GeV AuPb – CERES – Phys. Lett. B666 (2008) 425



ρ meson spectral function provides access to **restoration of chiral symmetry** at T_c degeneracy with chiral partner a_1

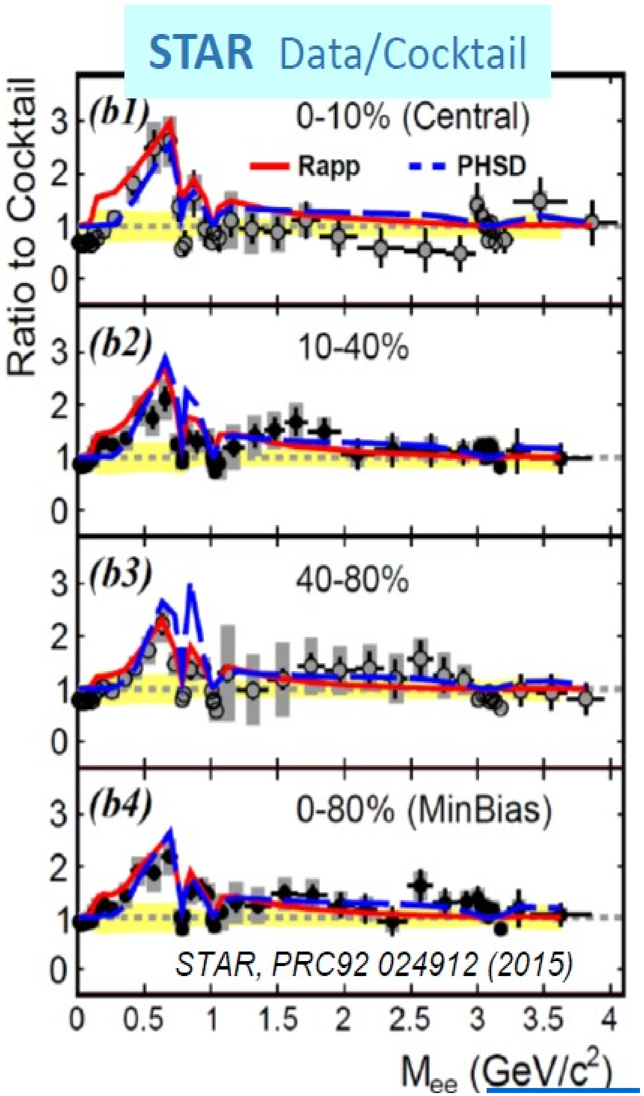
Eur. J. Phys. C61 (2009) 711



exp data of e^+e^- in central PbPb and of $\mu^+\mu^-$ in semi-central InIn after subtraction of all hadronic contributions except ρ: theories need significant broadening of spectral function of ρ at high T

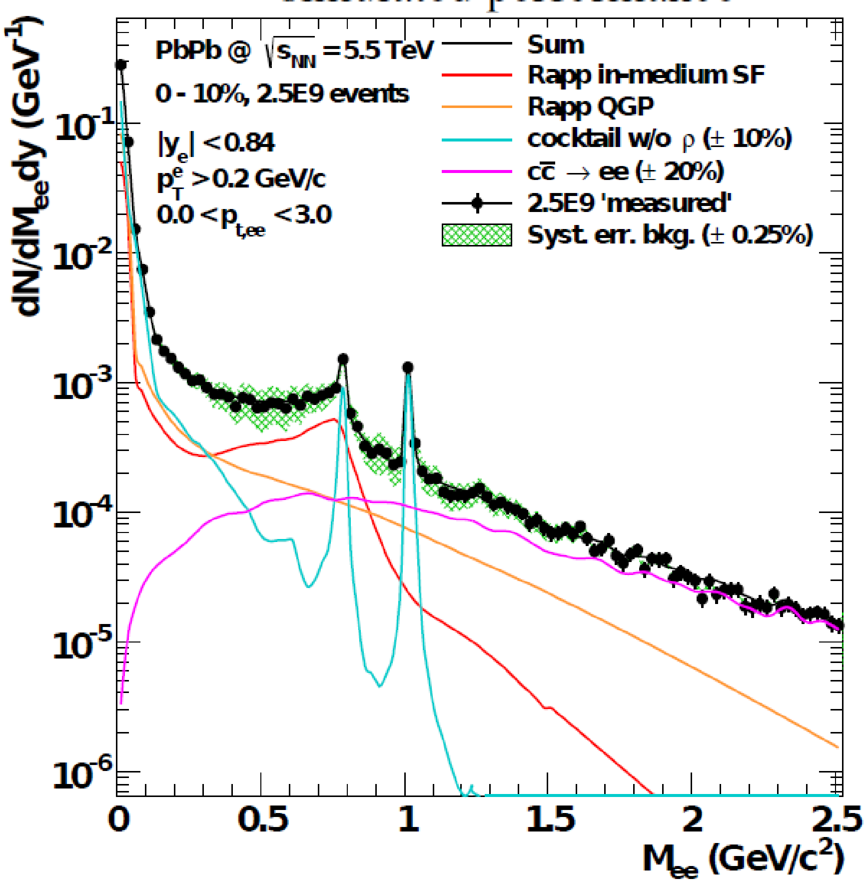
theory R. Rapp

Low and intermediate mass lepton pairs at colliders

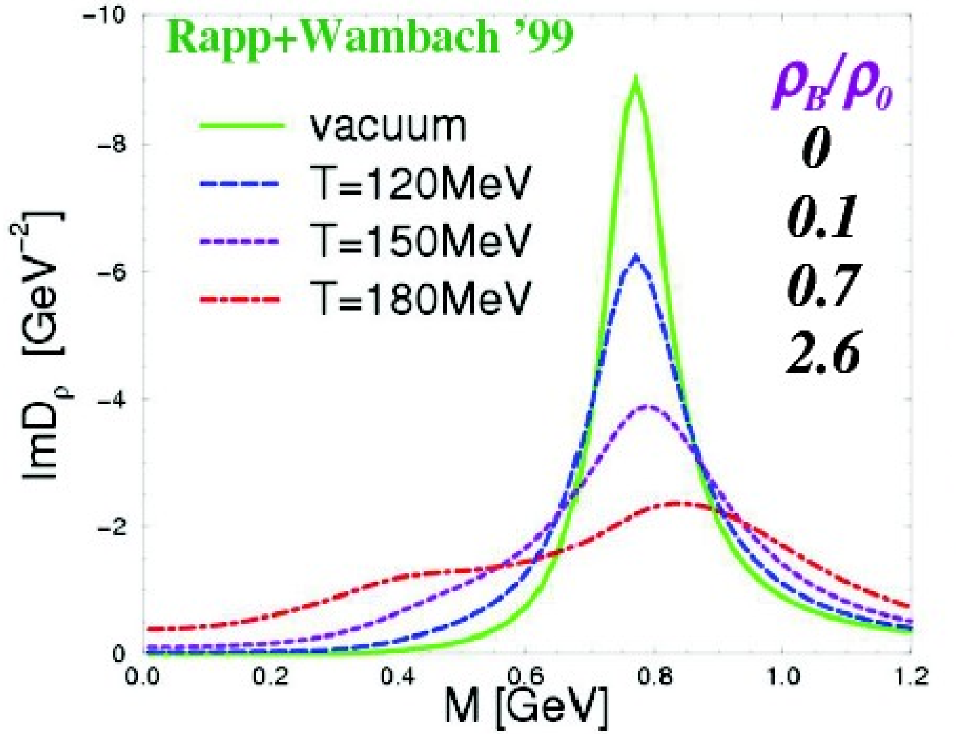


- at colliders much more difficult
- at RHIC after 15 years consolidated results between STAR and PHENIX described well by the same models as SPS data
- for ALICE very challenging project for Run3/Run4

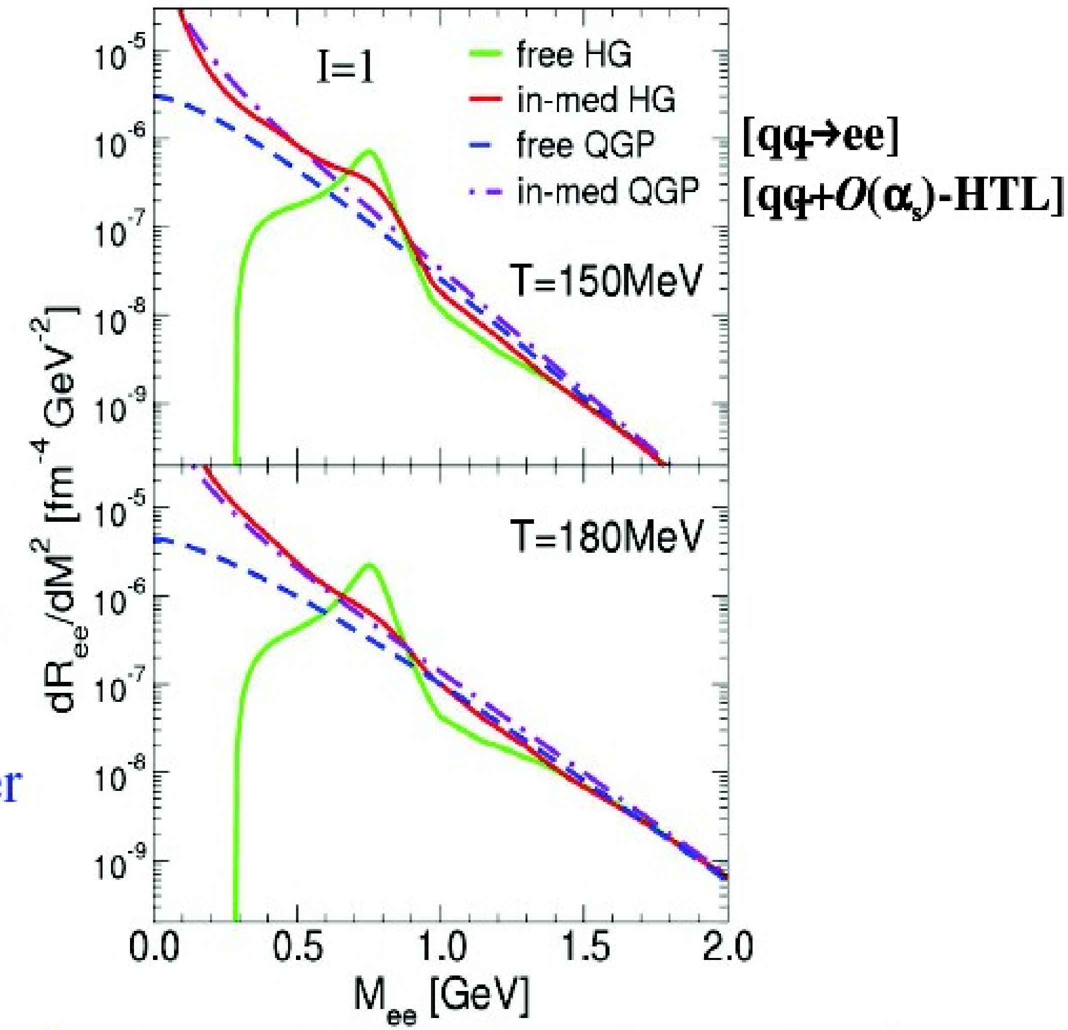
simulated performance



How does this modified ρ look like? integrate over space-time evolution of spectral function for ee mass spectrum



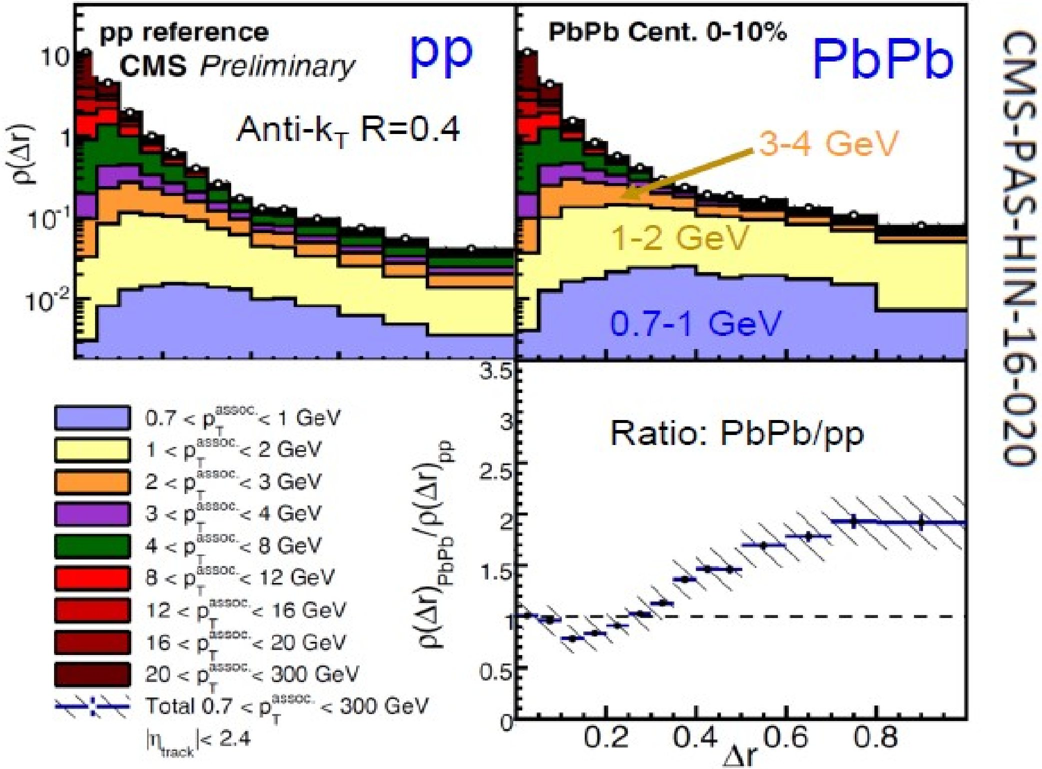
ρ-meson **"melts"** in hot and dense matter



in-med HG and QGP match
'quark – hadron duality?'

Where does lost energy go?

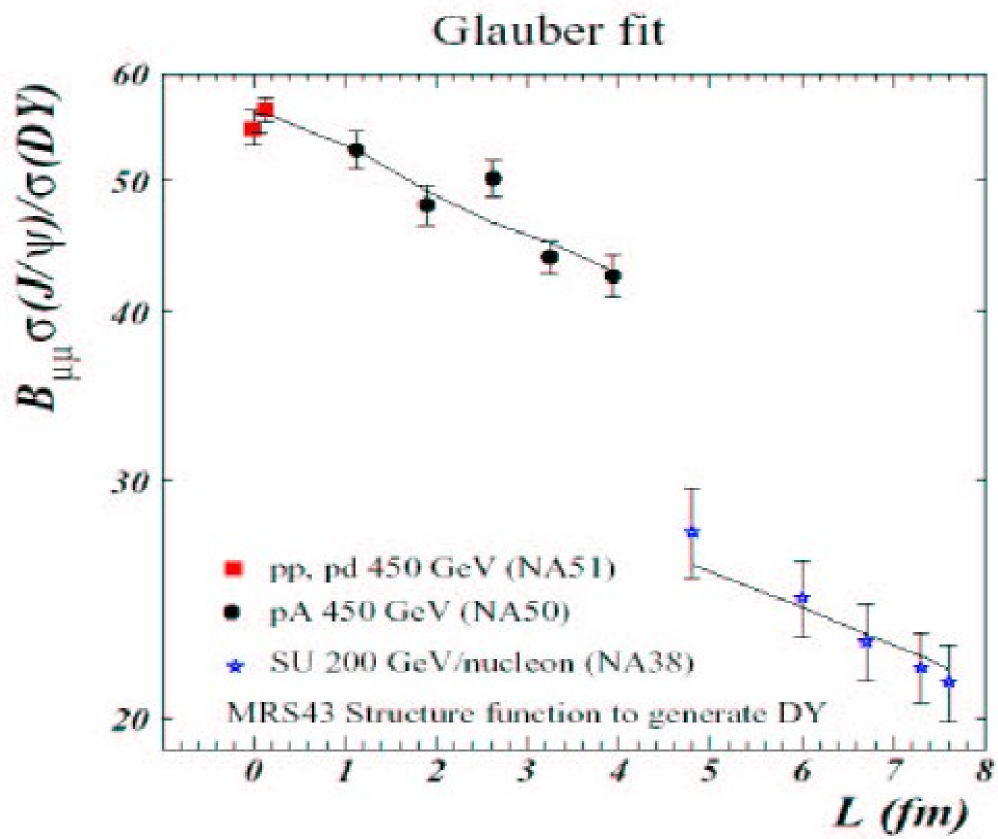
Jet-hadron correlations in pp and PbPb collisions at 5.02 TeV



low momentum particles
and at larger distance from jet core

J/ψ Suppression in pA Collisions

in pA and light nucl. coll. J/ψ production suppressed (NA38)



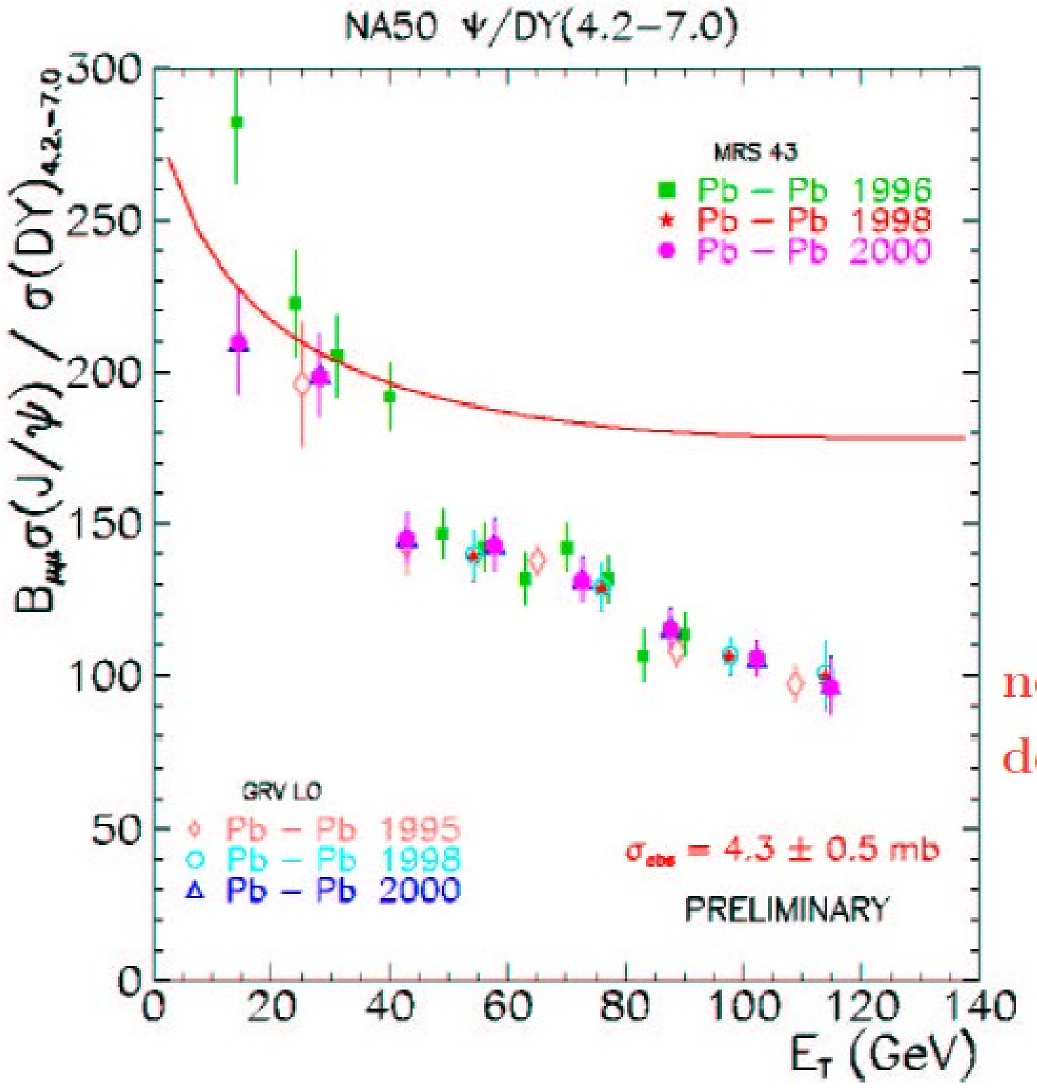
NA50,Phys.Lett.B553(2003)167

$$\sigma(J/\psi) \propto exp(-\rho\sigma_{abs}L)$$

with $\rho = 0.17/\text{fm}^3$ and $\sigma_{abs} = 4.3 \pm 0.6$ mb

Anomalous J/ψ Suppression in PbPb Collisions

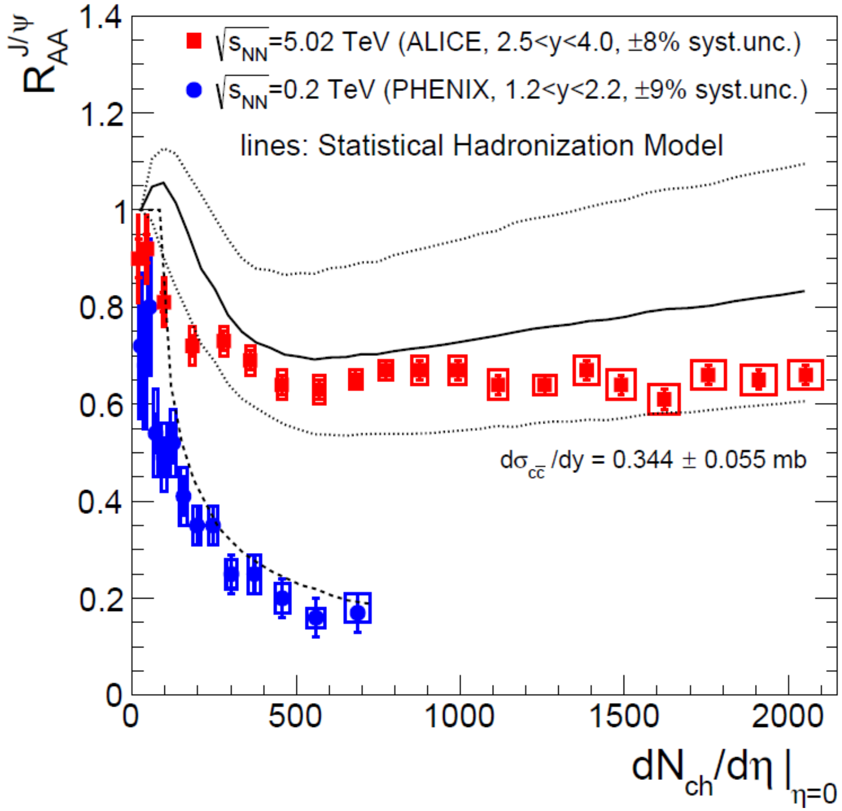
NA50, Phys. Lett. B447 (2000) 28 and Proc. Quarkmatter 2002, Nucl. Phys. A

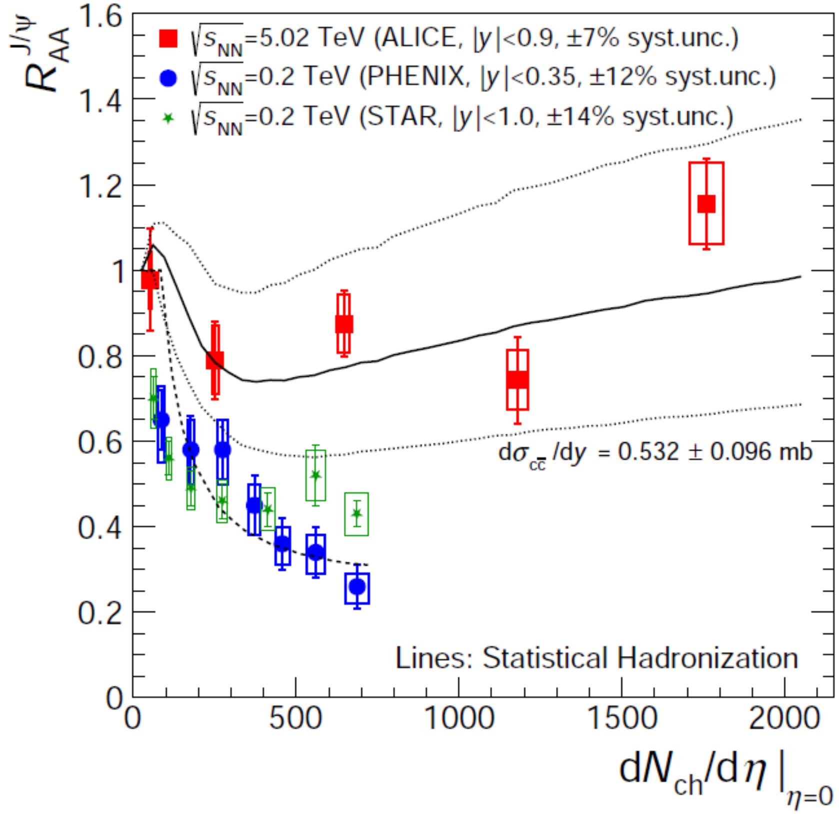


normal suppression as in pA does not descibe the data

J. Stachel

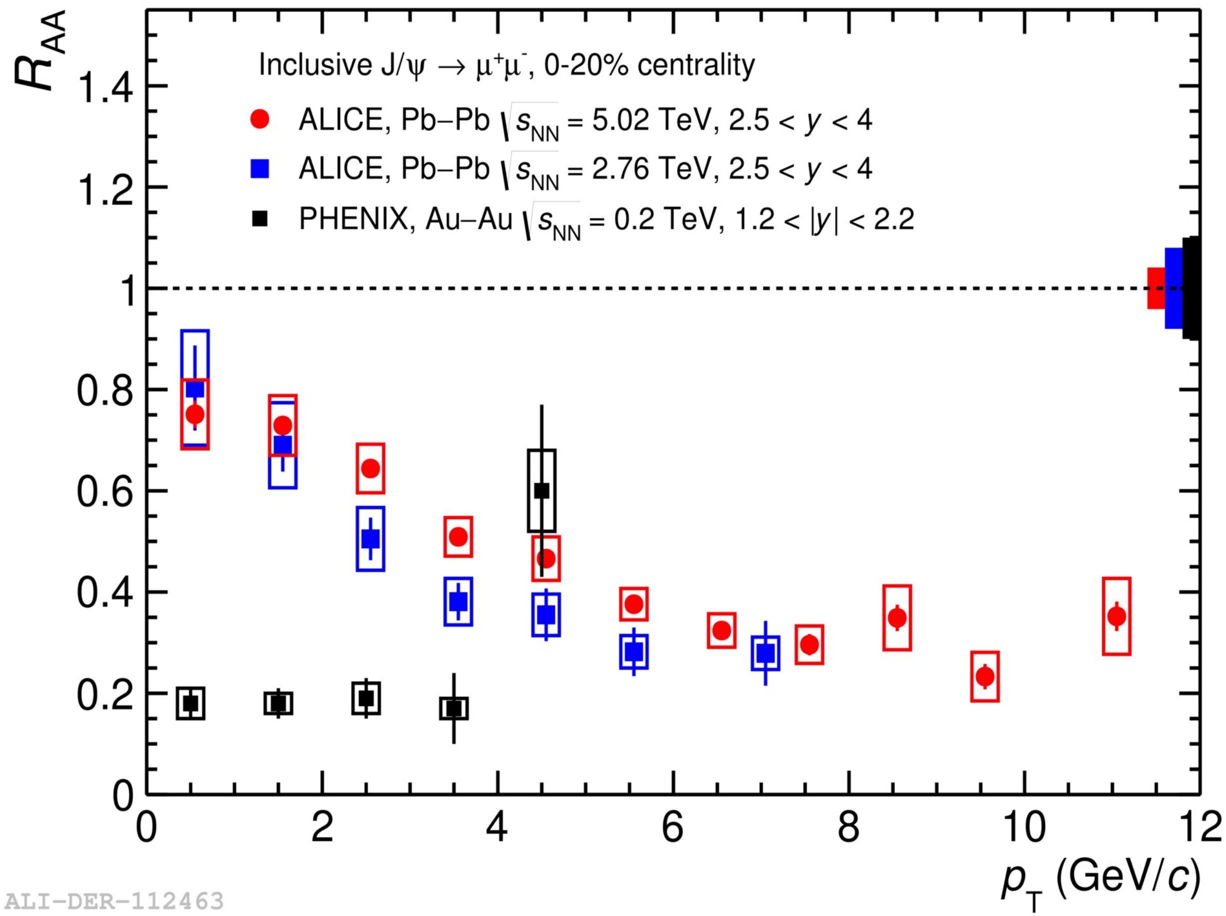
J/ψ and statistical hadronization





production in PbPb collisions at LHC consistent with deconfinement and subsequent statistical hadronization within present uncertainties
main uncertainties for models: open charm cross section, shadowing in Pb

transverse momentum spectrum



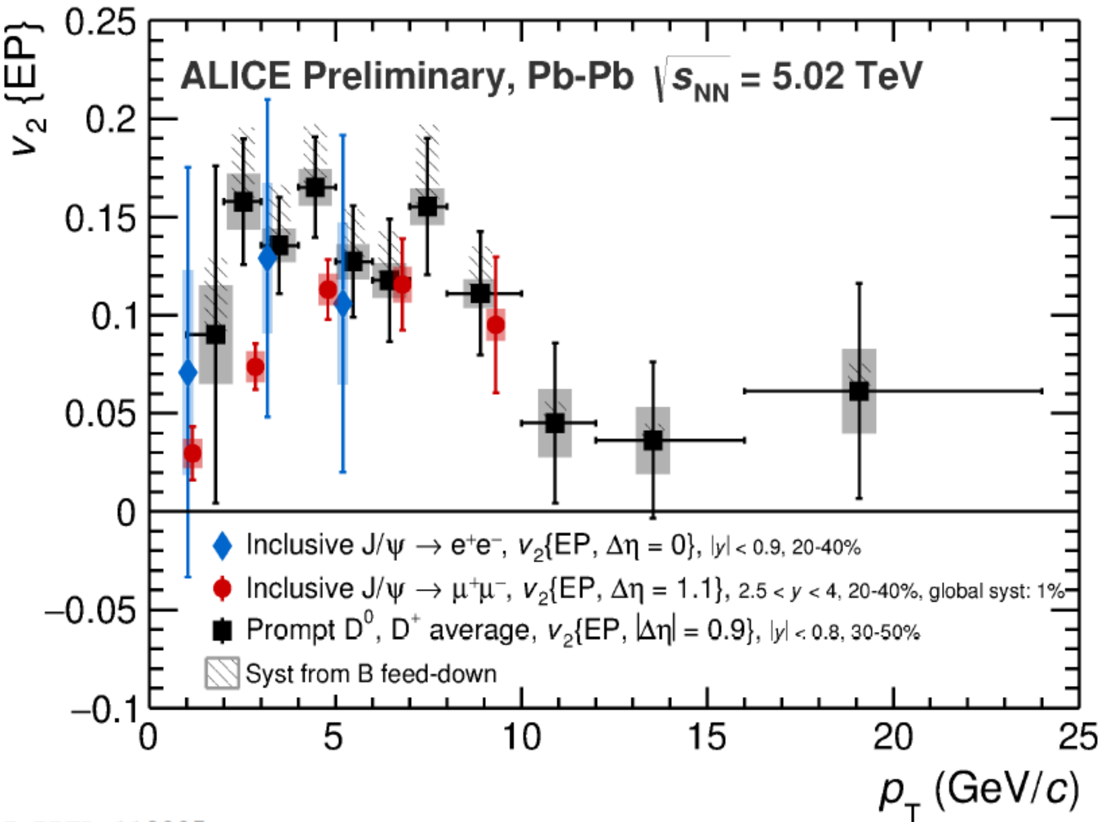
softer in PbPb as compared to pp

a qualitatively new feature as compared to RHIC where the trend is opposite

in line with thermalized charm in QGP at LHC, forming charmonia

elliptic flow of J/ψ vs p_t

arXiv:1705.05810



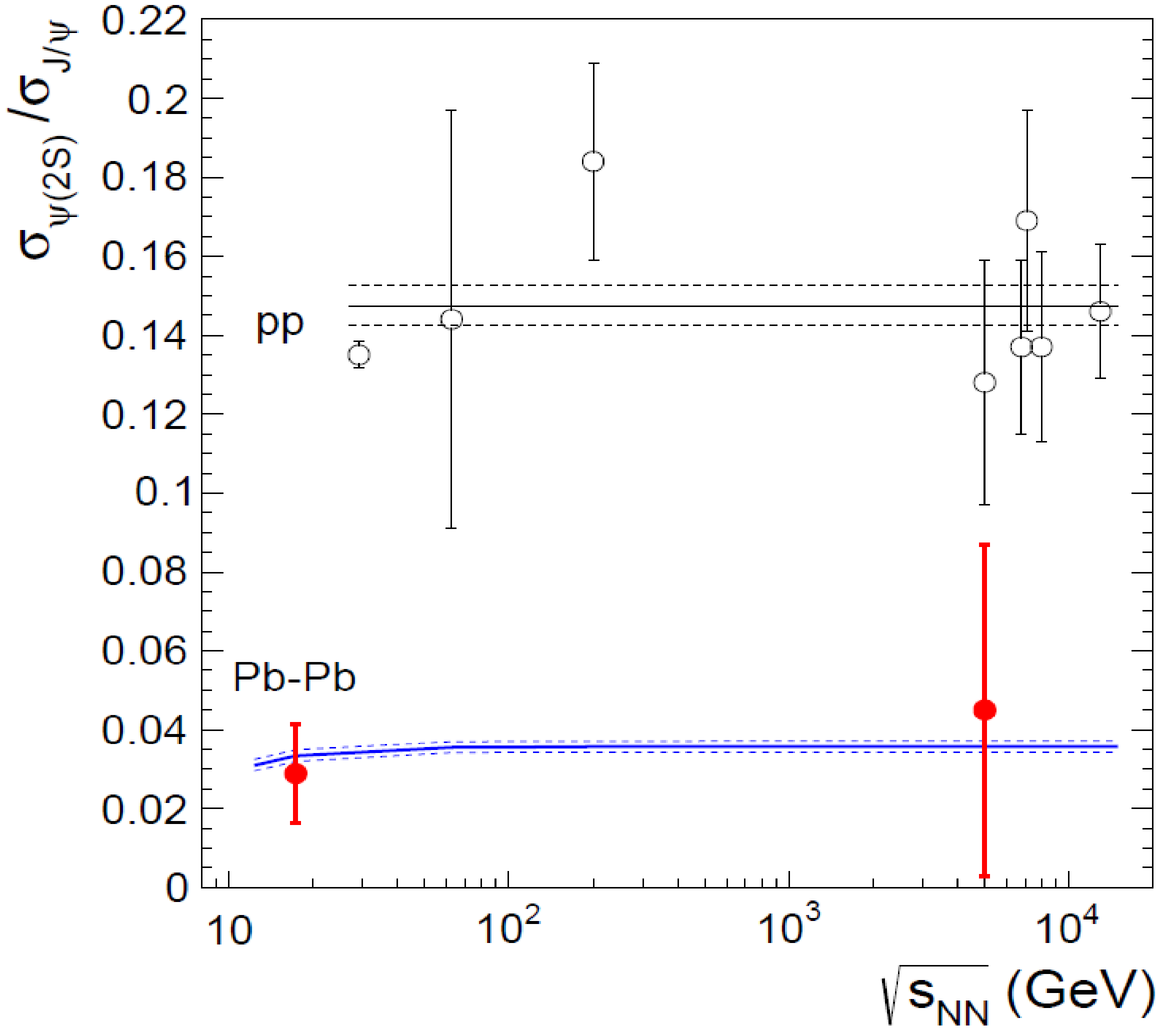
charm quarks thermalized in the QGP should exhibit the elliptic flow generated in this phase

- expect build-up with p_t as observed for π, p. K, Λ, … and vanishing signal for high p_t region where J/ψ not from hadronization of thermalized quarks

first observation of significant J/ψ v_2 in line with expectation from statistical hadronization

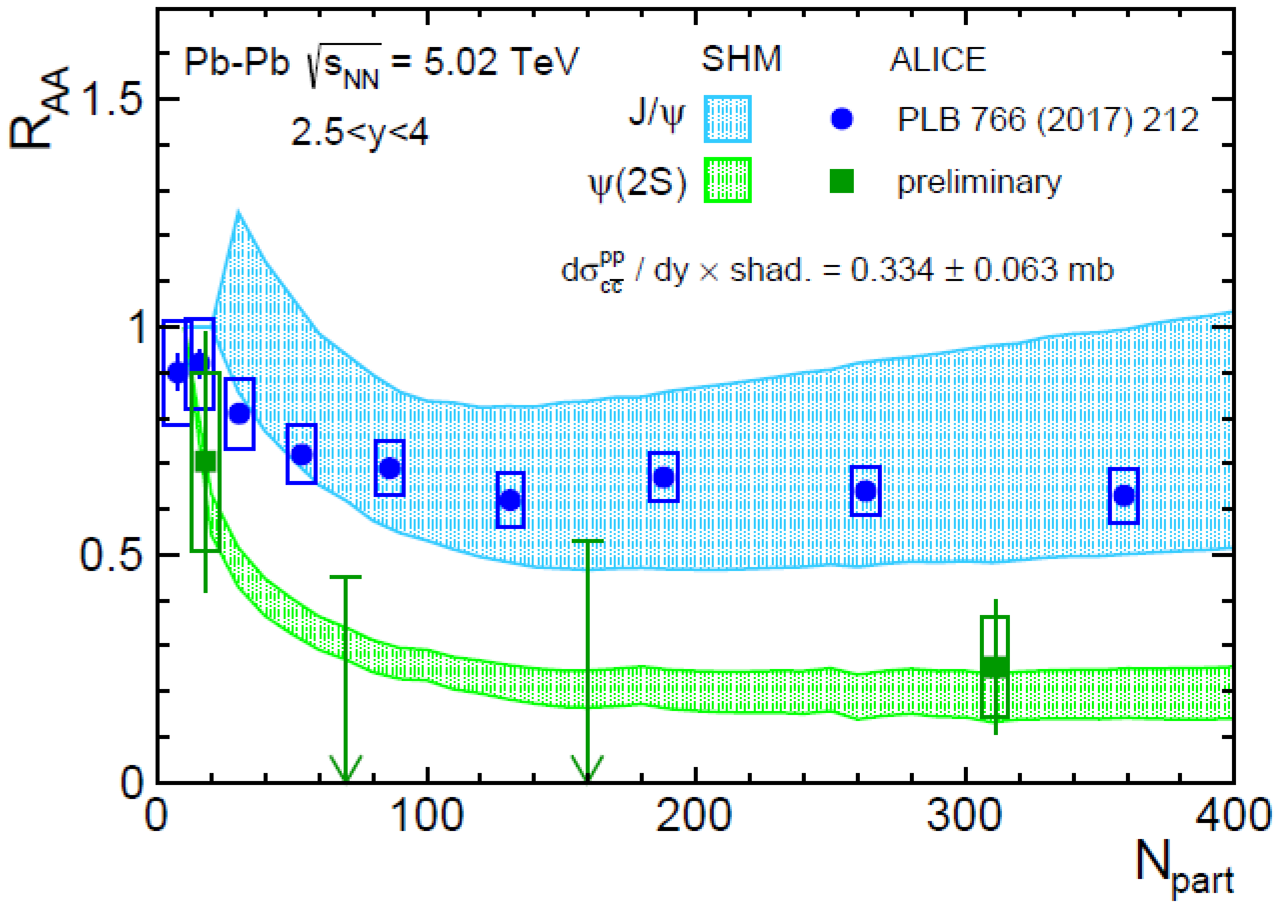
$\psi(2S)$

in picture where psi is created from deconfined quarks in QGP or at hadronization, psi(2S) is suppressed more than J/psi



What about ψ(2S)?

M. Köhler, A. Andronic, P. Braun-Munzinger, JS, arXiv:1807.01236

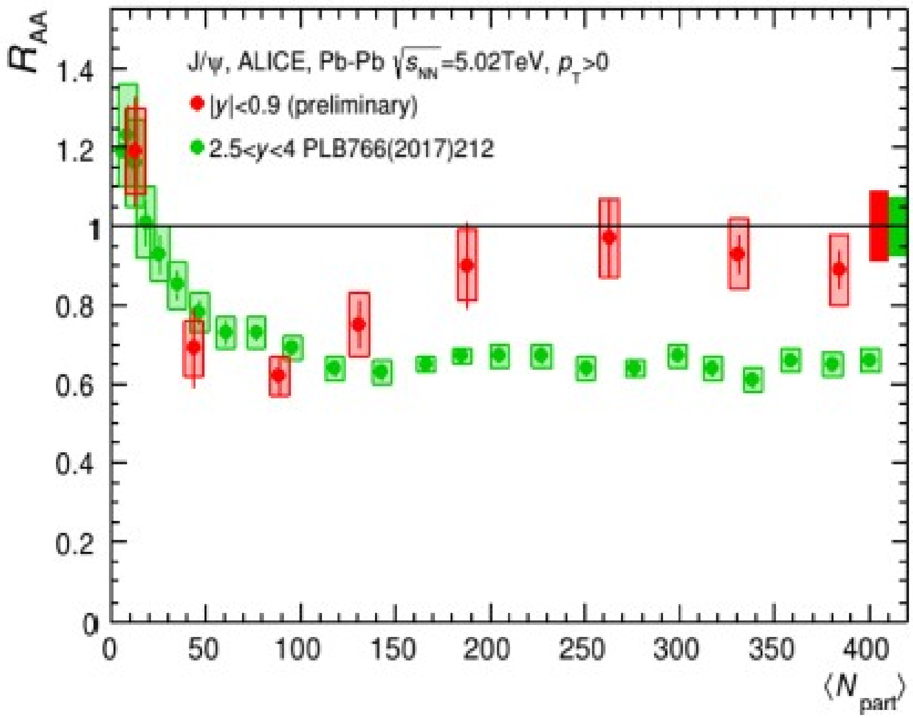


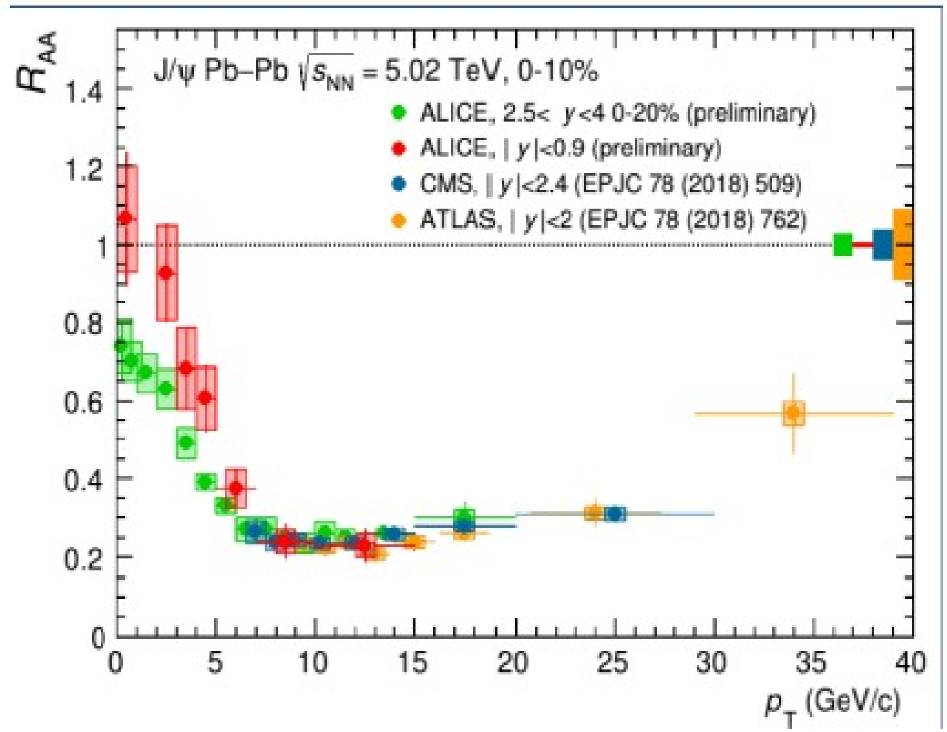
also excited state population completely in line, suppressed by Boltzmann factor

errors will decrease with more data in LHC Run3/4

Charmonium at LHC: peaks at mid-y and strong enhancement at low transverse momentum

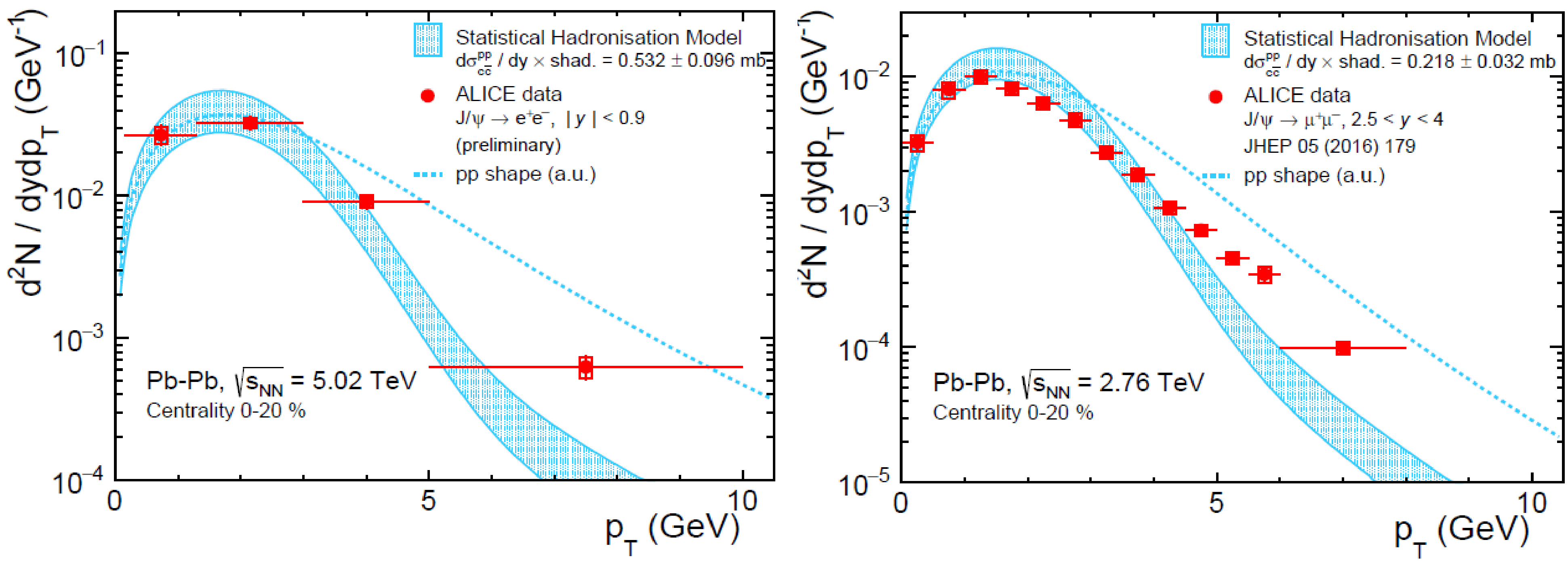
nuclear modification factor: $$R_{AA}(p_T) = \frac{dN^{AA}/dp_T}{\langle N_{coll}\rangle dN^{pp}/dp_T}$$





J/ψ transverse momentum spectra from stat. hadr.

M. Köhler, A. Andronic, P. Braun-Munzinger, JS, arXiv:1807.01236

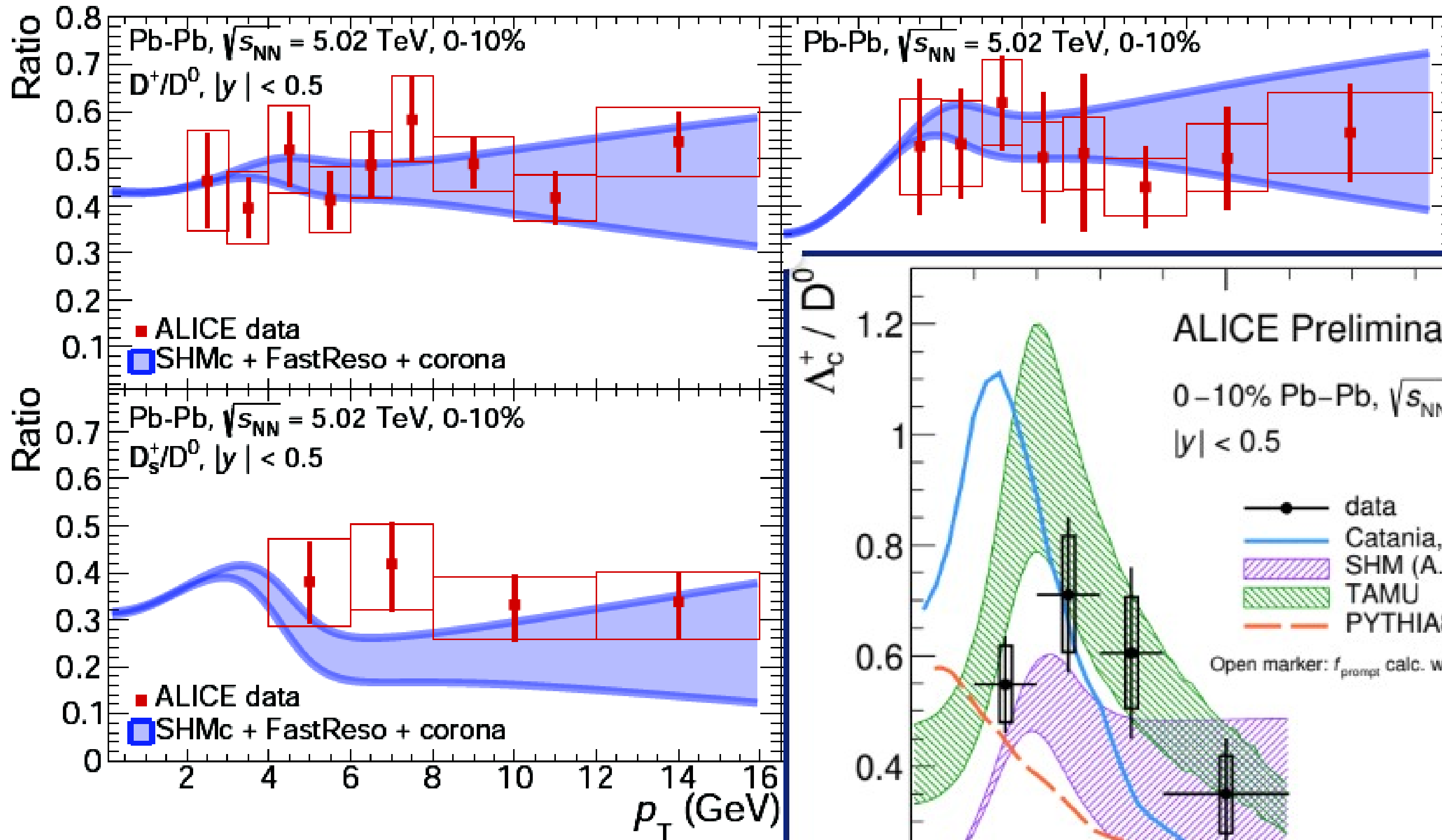


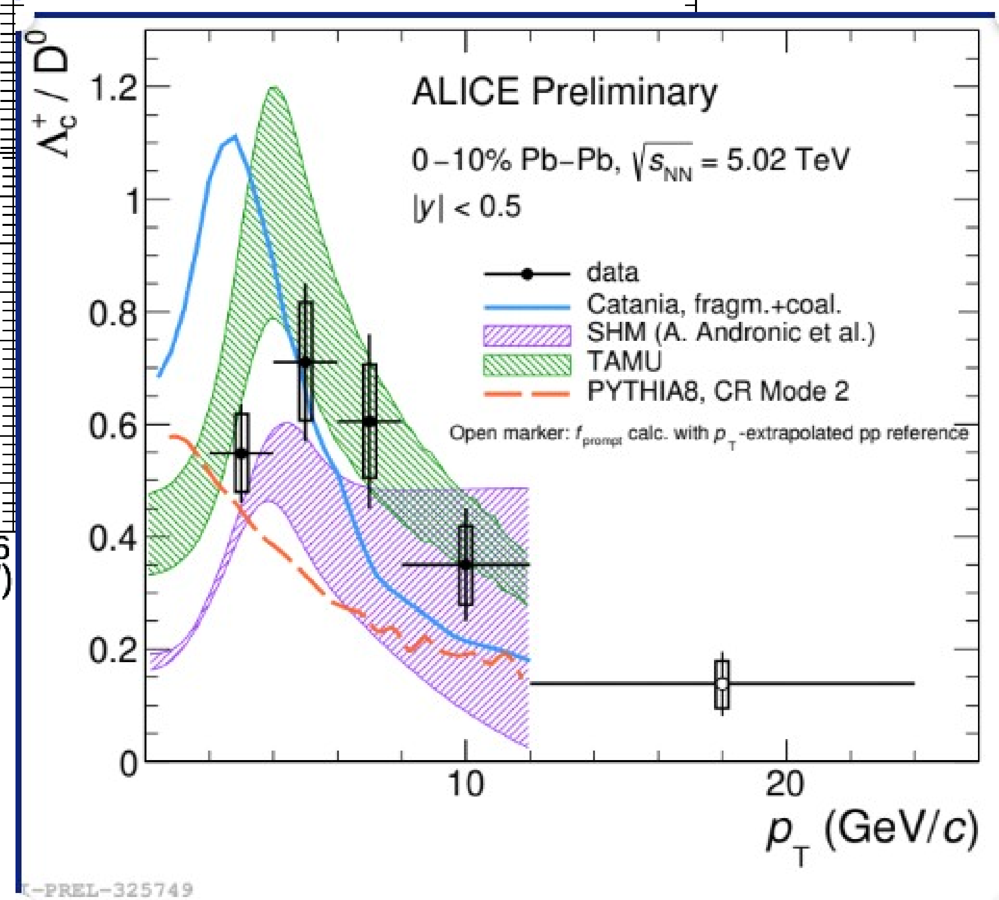
good agreement up to 5 GeV/c without any free parameters

J/ψ formed at hadronization at T_C from thermalized charm quarks flowing with the rest of the medium

Ratios of charm hadron to D^0 spectra

A. Andronic et al., JHEP07 (2021) 035





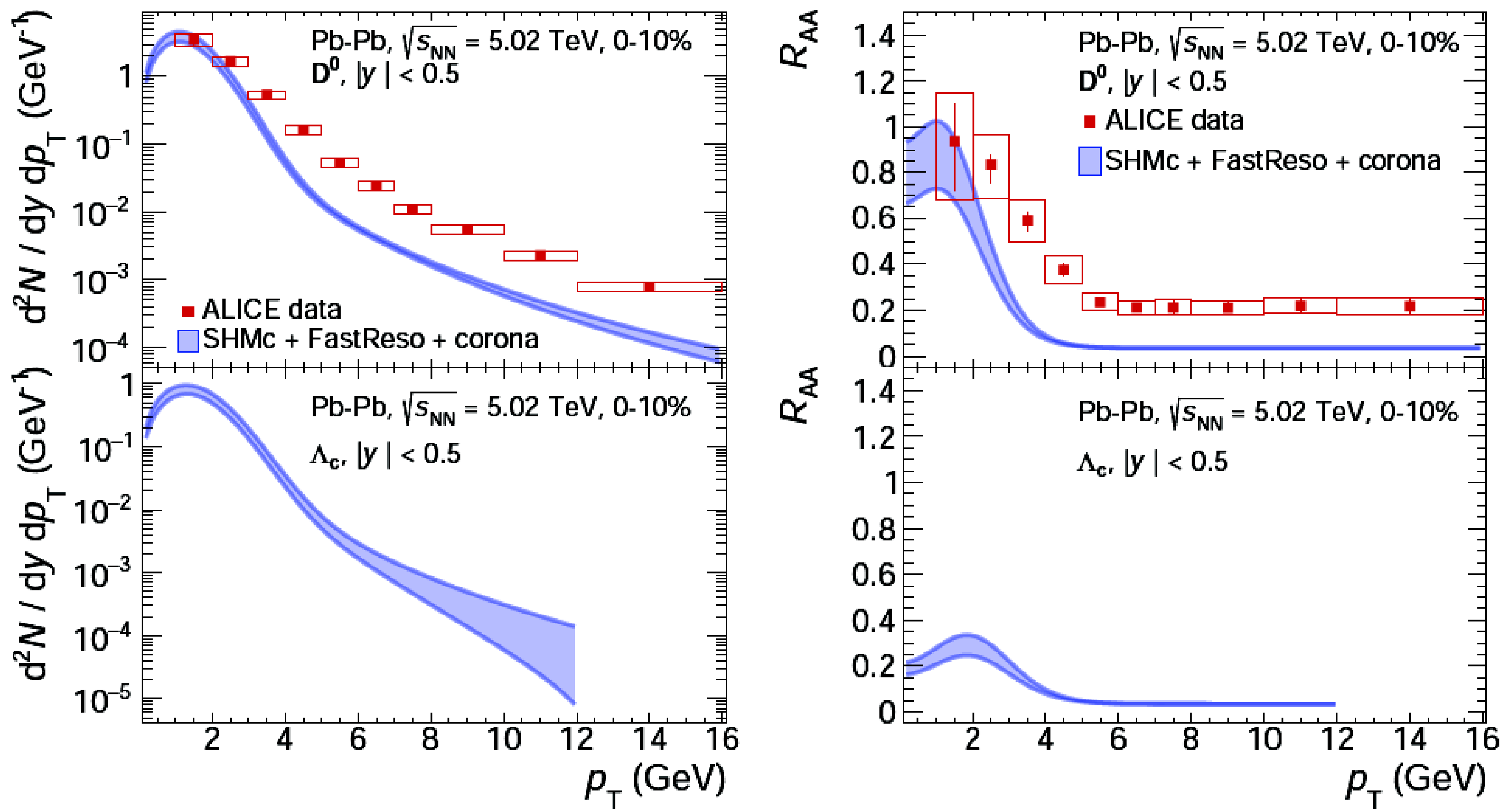
preliminary Lambda_c fits well into picture

Spectra and R_{AA} of D^0 mesons and Λ_c baryons

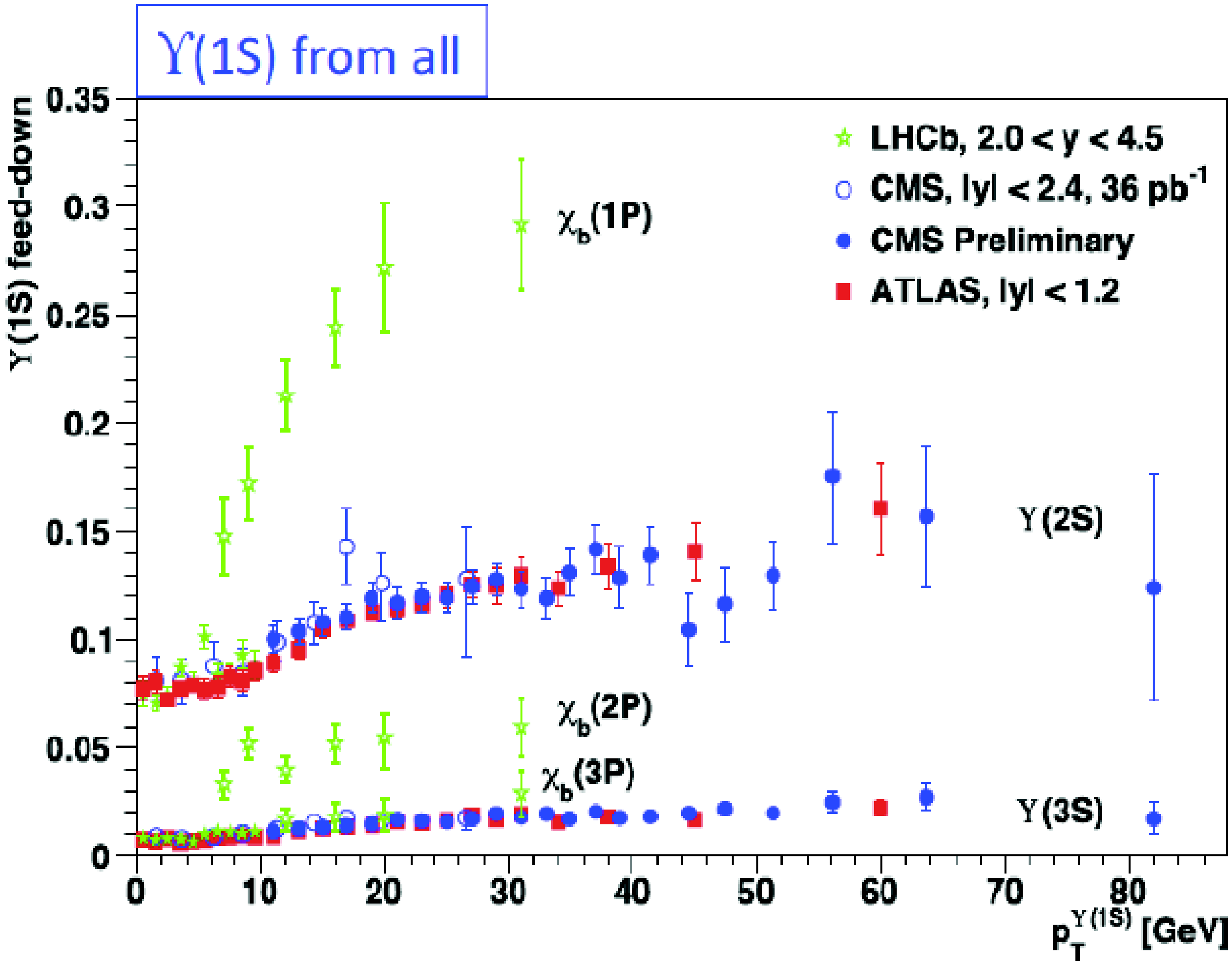
for open heavy flavor hadrons strong contribution from resonance decays
- include all known charm hadron states as of PDG2020 in SHMc
- compute decay spectra with FastReso: 76 2-body and 10 3-body decays

(A. Mazeliauskas, S. Floerchinger, E. Grossi, D. Teaney, EPJ C79 (2019) 284 arXiv: 1809.11049)

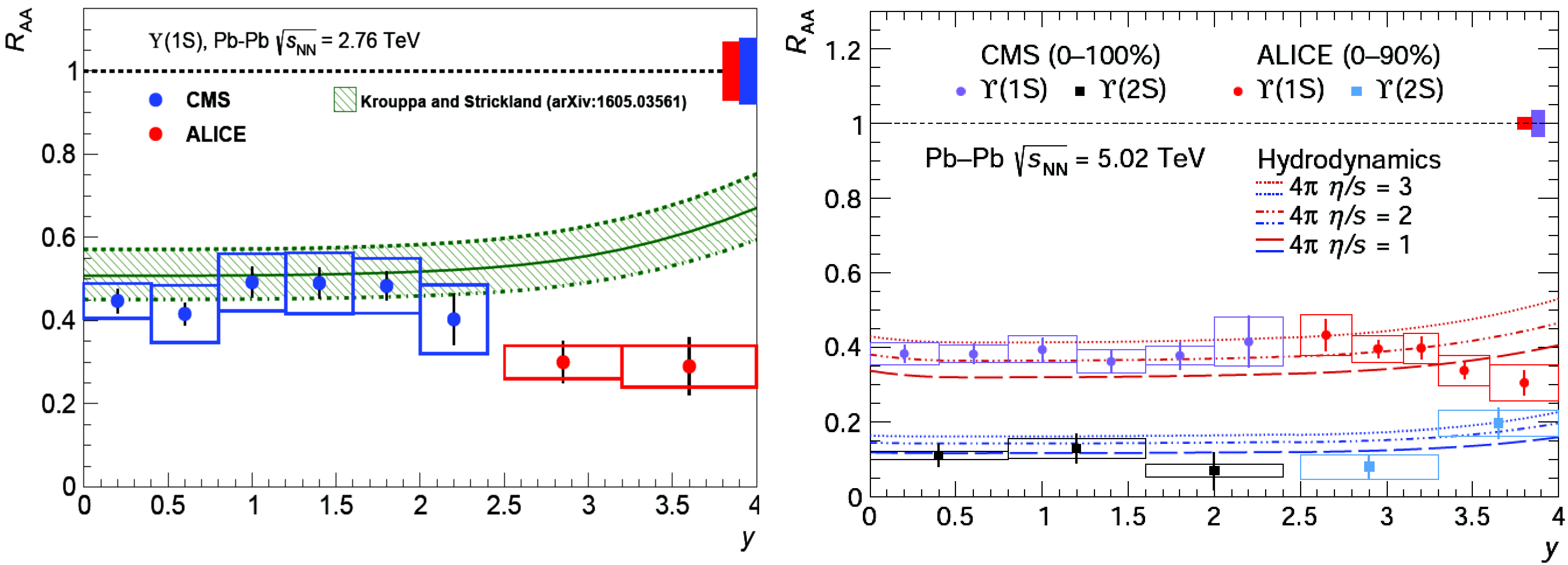
A. Andronic et al., JHEP07 (2021) 035



Feeding into Upsilon (1S)



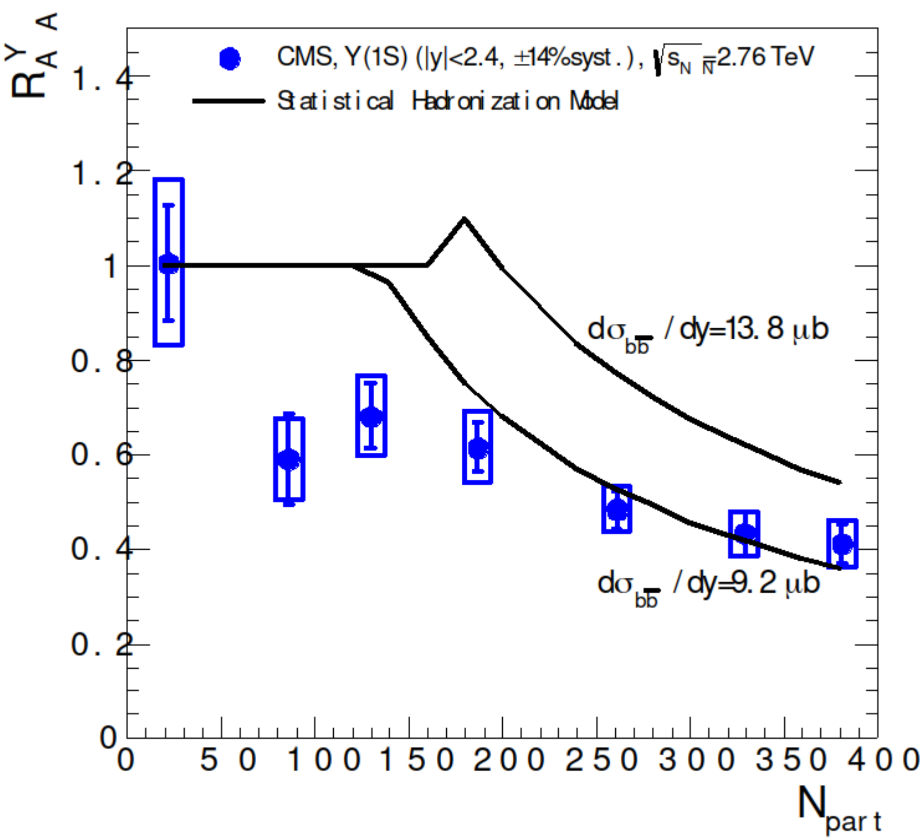
Upsilon R_{AA} rapidity dependence

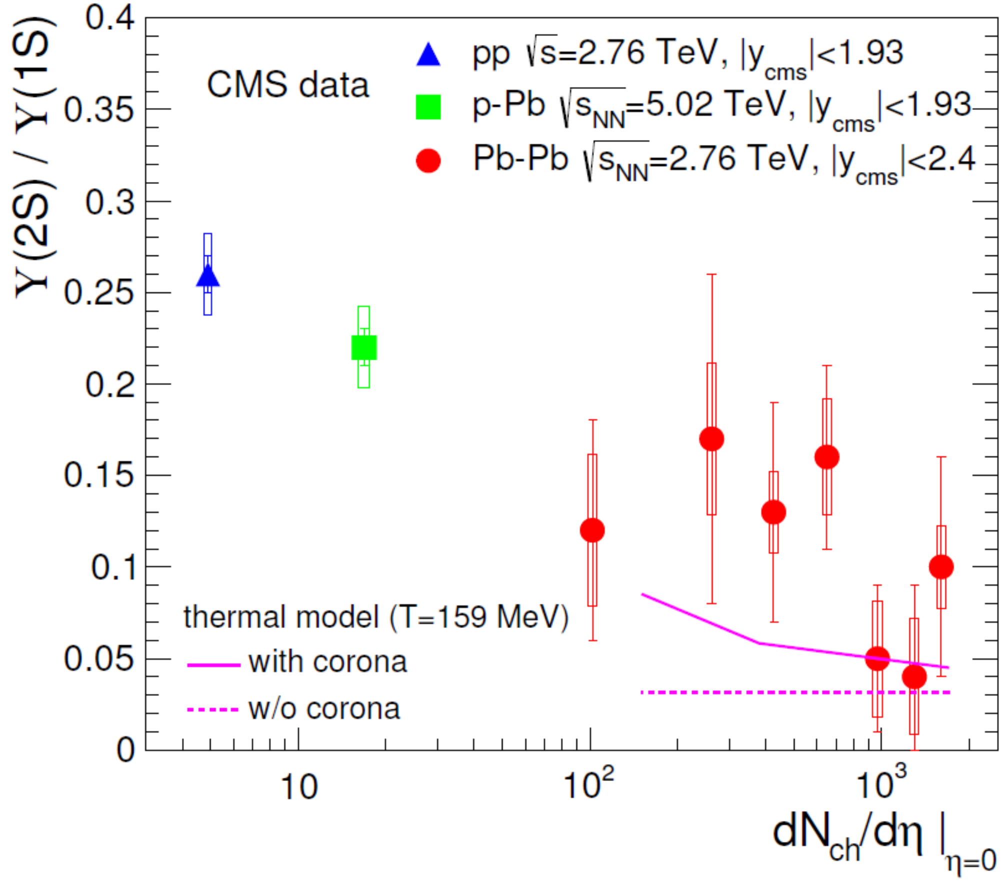


Indication: R_{AA} peaked at mid-y like for J/ψ
not in line with collisional damping in expanding medium

the Upsilon could also come from statistical hadronization

SHM/thermal model: Andronic et al.





in this picture, the entire Upsilon family is formed at hadronization
but: need to know first – do b-quark thermalize at all? spectra of B
- total b-cross section in PbPb