



Berkeley
UNIVERSITY OF CALIFORNIA



Dark Radiation from Heavy QCD Axions

Planck 2022

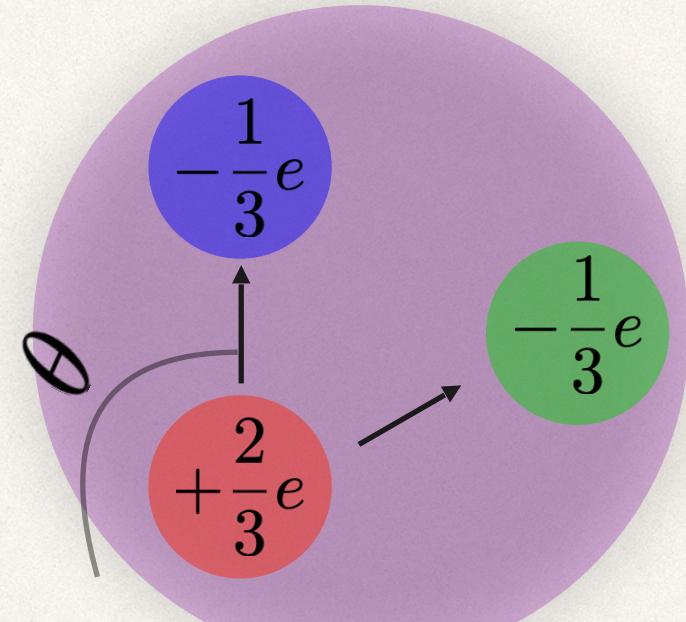
David I. Dunsky, Lawrence J. Hall, Keisuke Harigaya

arXiv: 2205.11540

Strong CP Problem

- # ❖ Why is neutron electric dipole moment so small?

Classically:



Hook 2018

Why θ so small is the Strong CP Problem

- # ❖ Quantum mechanically:

$$\mathcal{L}_{CP} = \frac{\theta}{32\pi^2} G\tilde{G}$$

Standard QCD Axion

- ❖ One of most popular solutions: promote θ to a field

$$\mathcal{L}_{CP} = \frac{\theta}{32\pi^2} G\tilde{G} \rightarrow \left(\theta + \frac{a}{f_a}\right) \frac{1}{32\pi^2} G\tilde{G} \quad (\text{Now } d_n \propto \theta + a/f_a)$$

- ❖ Below $\Lambda_{\text{QCD}} \sim 150 \text{ MeV}$, axion acquires a potential from strong dynamics

$$V \approx -\Lambda_{\text{QCD}}^4 \cos(\theta + a/f_a)$$

$$\frac{\partial V}{\partial a} = 0 \longrightarrow \langle d_n \rangle = 0$$

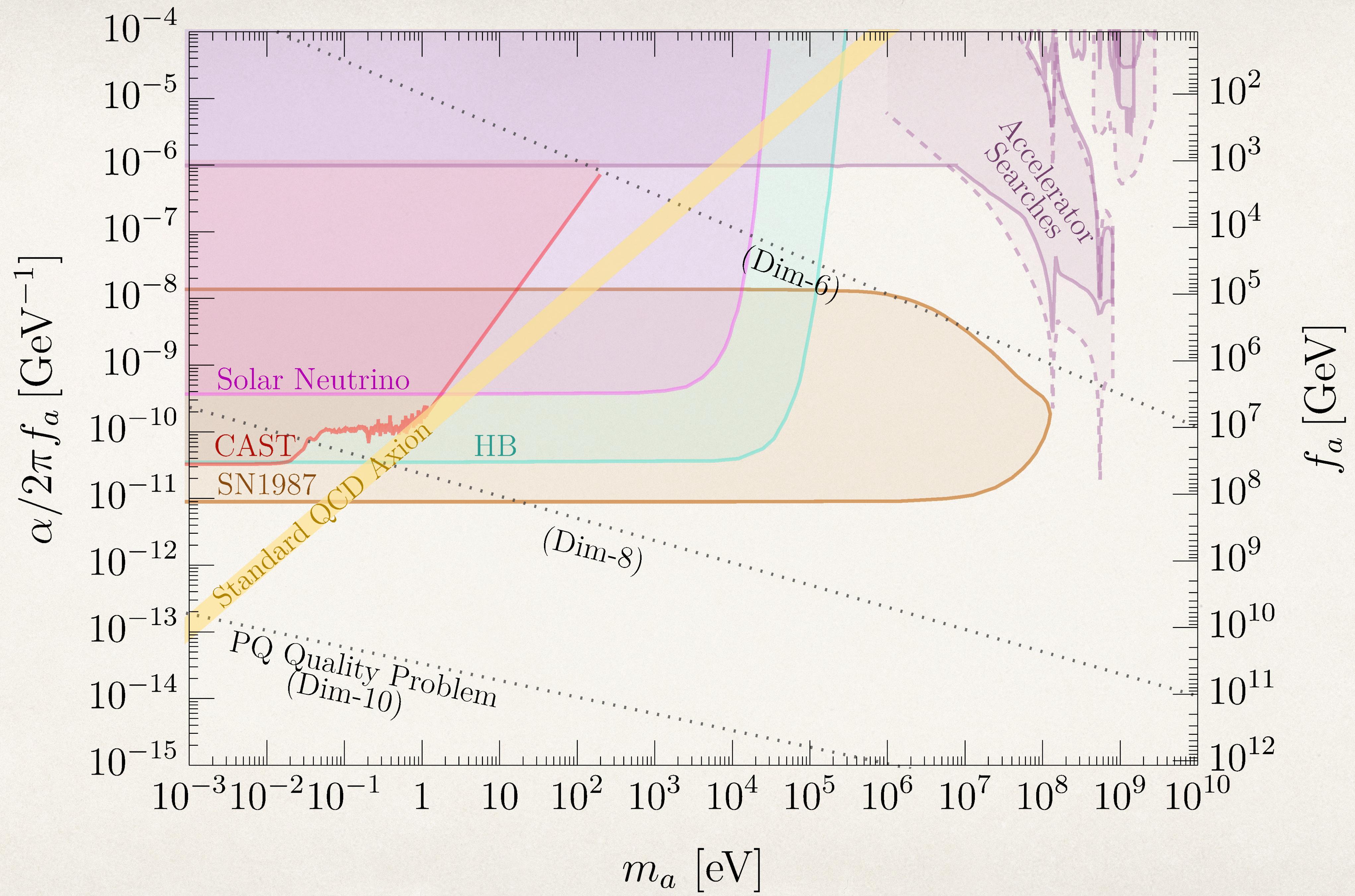
$$m_a^2 \equiv -\frac{\partial^2 V}{\partial a^2} \Big|_{\langle a \rangle} \simeq \frac{\Lambda_{\text{QCD}}^4}{f_a^2} \simeq \left(0.6 \text{ eV} \frac{10^7 \text{ GeV}}{f_a}\right)^2$$

Axion Quality Problem

- $U(1)_{\text{PQ}}$ symmetry not gauged \rightarrow expected to be broken by higher-dim operators
- For PQ breaking field $\phi = (f_a + s)e^{ia/f_a}$

$$V_{\cancel{\text{PQ}}} = \lambda_n \frac{\phi^n}{M_{\text{Pl}}^{n-4}} \quad \rightarrow \quad \text{Shifts minimum of axion potential}$$

$$m_a \gtrsim \text{MeV} \left(\frac{f_a}{10^5 \text{ GeV}} \right)^2 \left(\frac{f_a}{M_{\text{Pl}}} \right)^{\frac{n-6}{2}} \lambda_n \quad \begin{matrix} (\text{For axion to still solve}) \\ (\text{Strong CP Problem}) \end{matrix}$$



Origin of Heavy QCD Axions

- ✿ Heavy QCD axions can arise from:
 - ✿ SU(3) or embeddings of SU(3) becoming confined at high energies
Dimopolous 1979, Tye 1981, Holdom & Peskin 1982, Holdom 1985, Flynn and Randall 1987, Agrawal and Howe 2018
 - ✿ 5D instantons (see Tony Gherghetta's talk from Tuesday)
Gherghetta et al 2020
 - ✿ Mirror QCD theories
Rubakov 1997, Berezhiani 2001

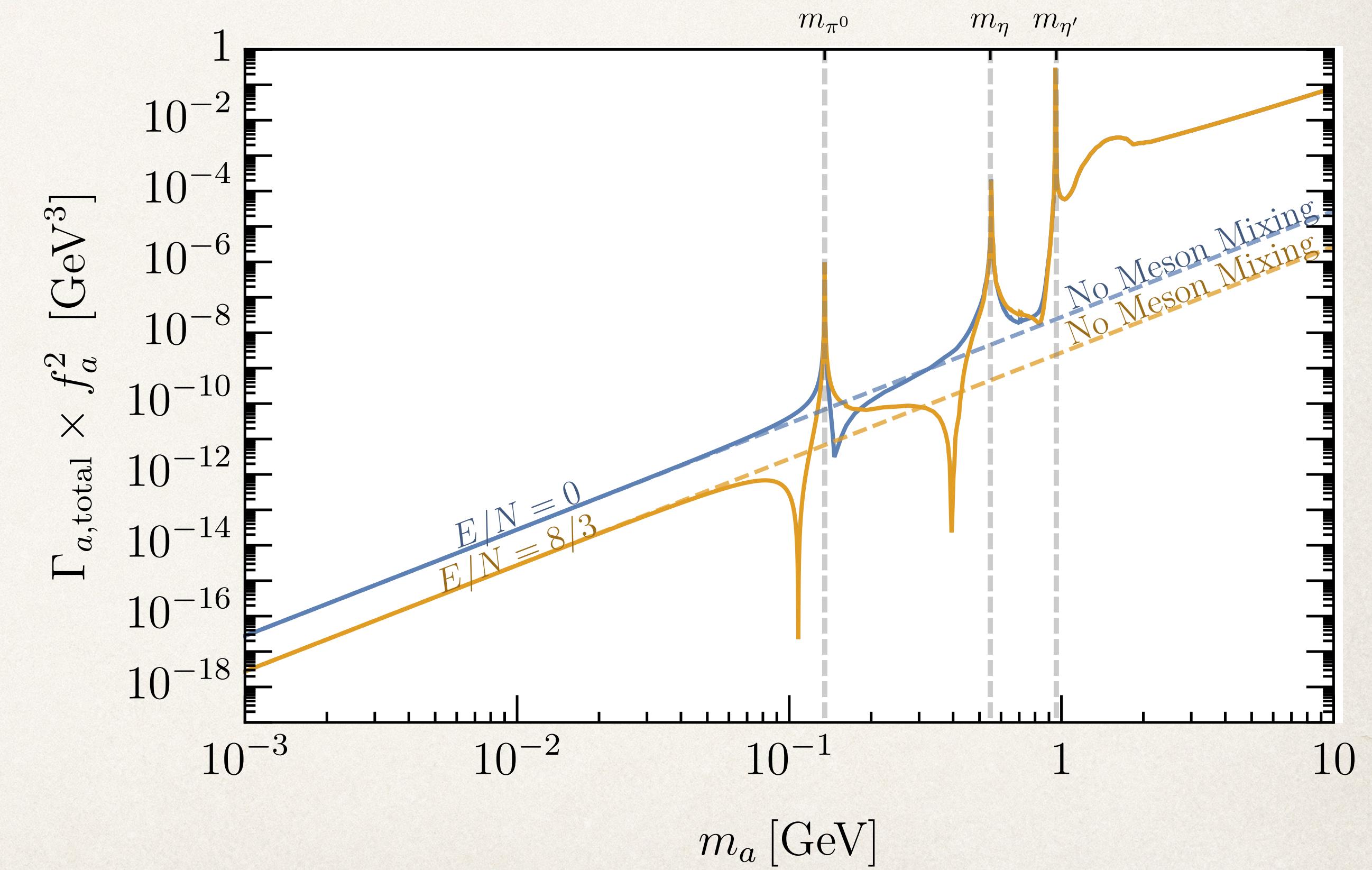
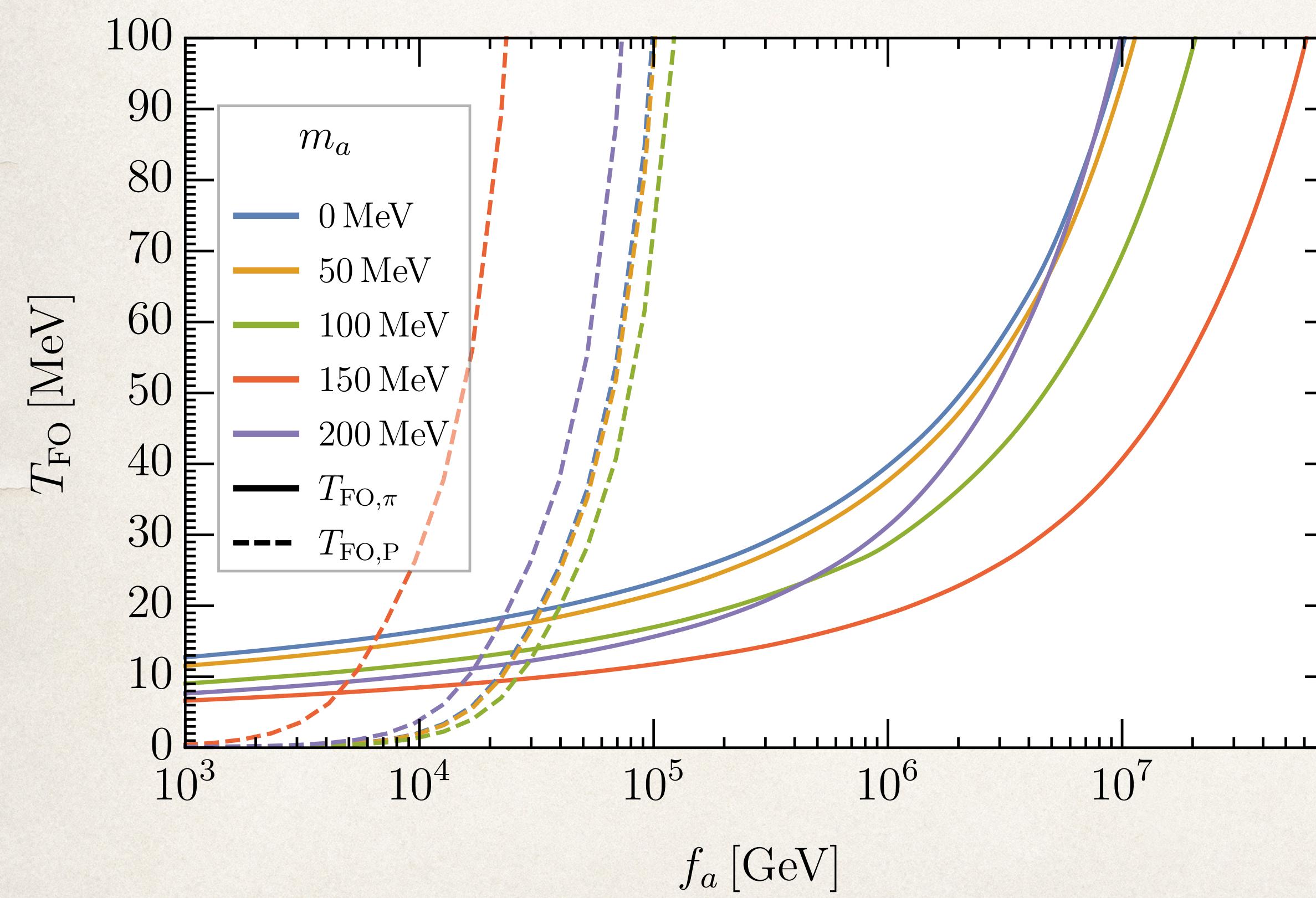
Cosmological Signals of Heavy QCD Axion

- ✿ For much of this parameter space, axion can decay around or after neutrino decoupling
- ✿ Axion decays heat up photons relative to neutrinos → *negative* ΔN_{eff} $(N_{\text{eff}} < 3.044)$ $(N_{\text{eff}} \propto \frac{\rho_\nu}{\rho_\gamma})$
- ✿ Accurate calculation requires momentum-space Boltzmann equation $f(\mathbf{p})$ because of:
 - ✿ Danger of some fraction of axions decaying late (high momentum, time dilation)
 - ✿ Precision CMB measurements

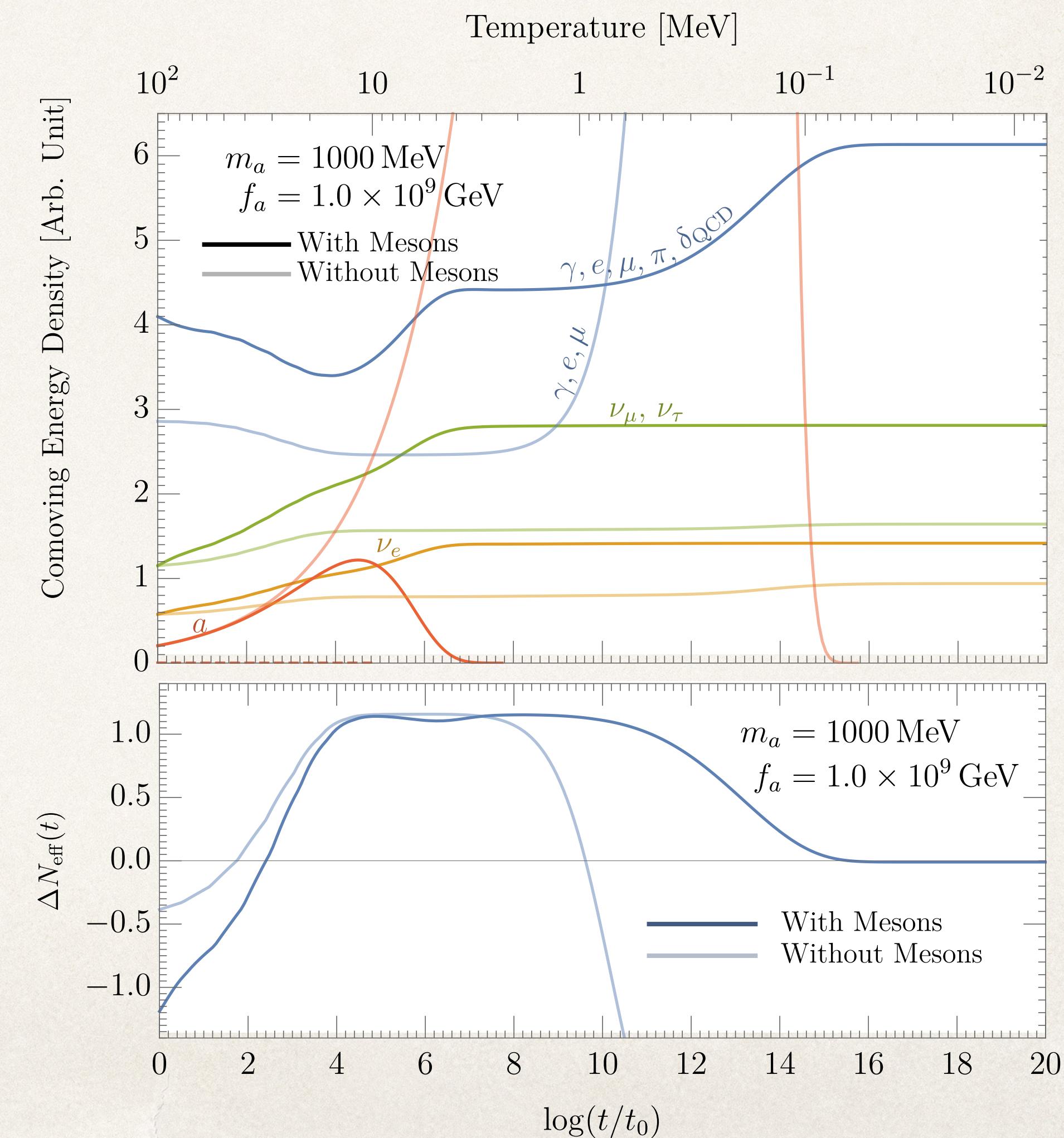
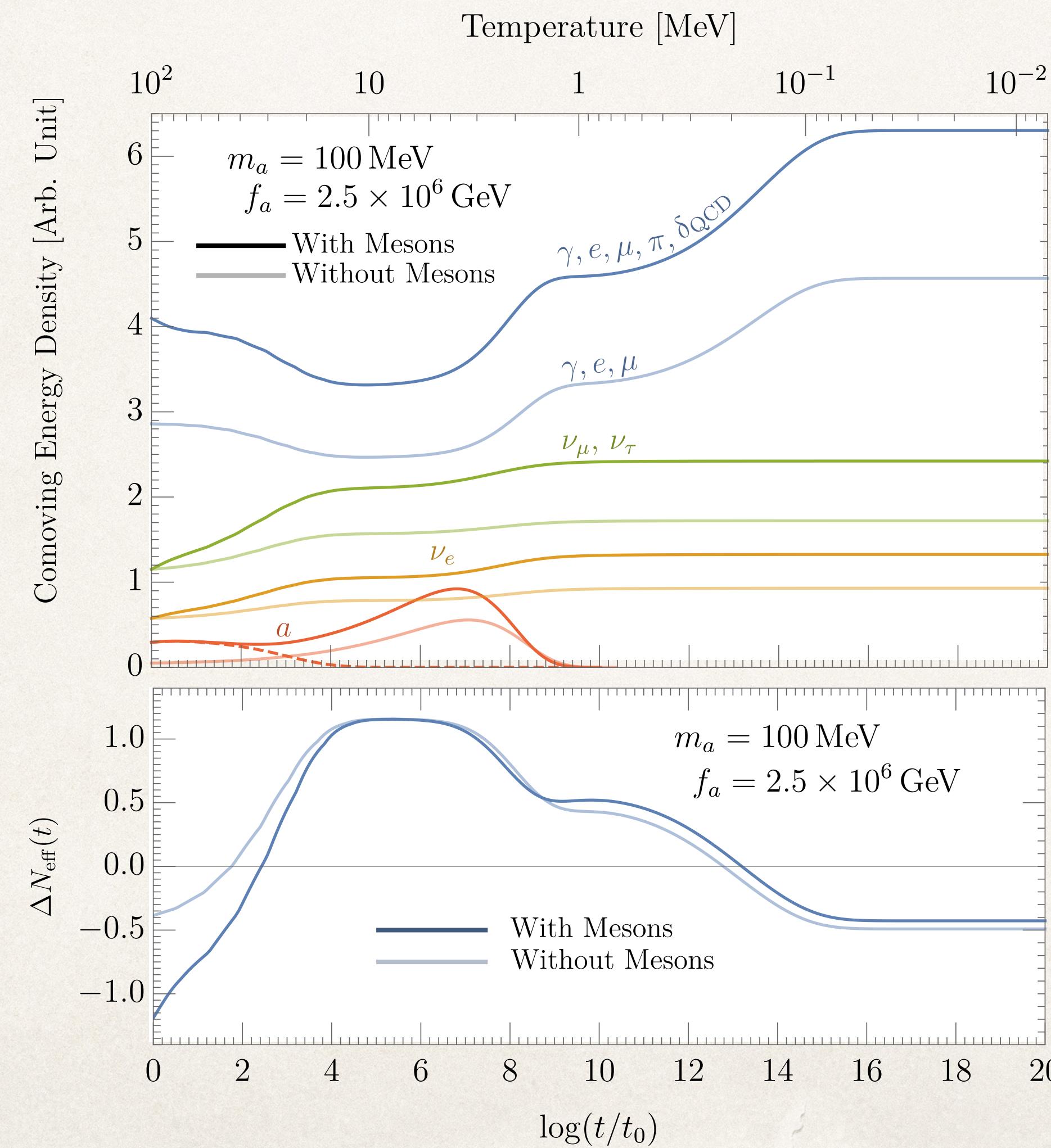
Important Effects for Heavy QCD Axions

- ❖ Previous work calculated N_{eff} for heavy Axion Like Particles (ALPs), *not* QCD axions
Cadamuro & Redondo 2012, Millea et al 2015, Depta et al 2020
- ❖ We must incorporate the following key effects:
 - ❖ Axion-pion scattering
 - ❖ Axion-meson mixing in photon coupling
Aloni et al 2019
 - ❖ Proper axion lifetime (QCD decay channels)
Aloni et al 2019
 - ❖ Enhanced frozen-out abundance from gluons
Salvio et al 2013, D'Eramo et al 2021
 - ❖ QCD contributions to the background energy density
Saikawa and Shirai 2018

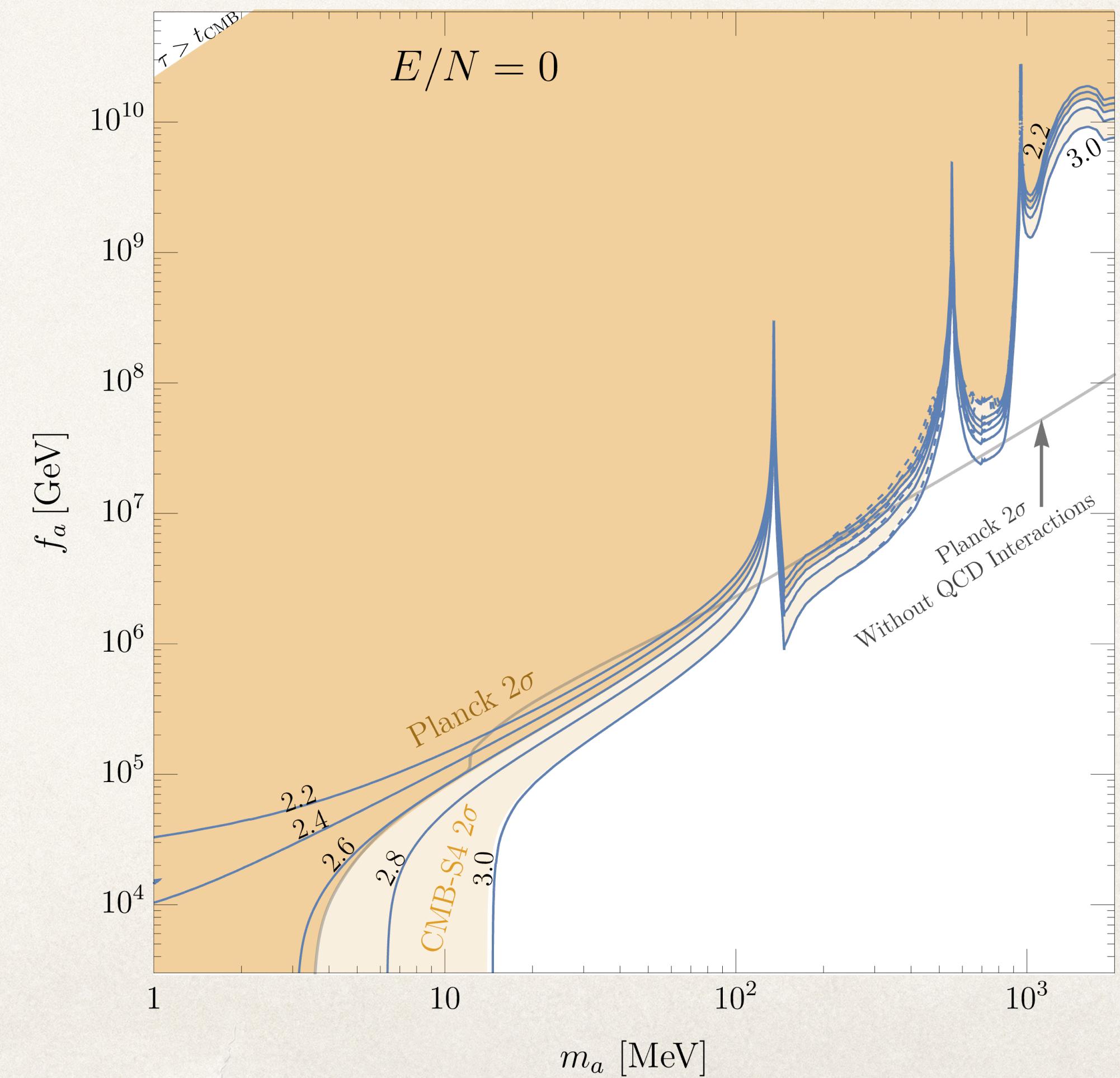
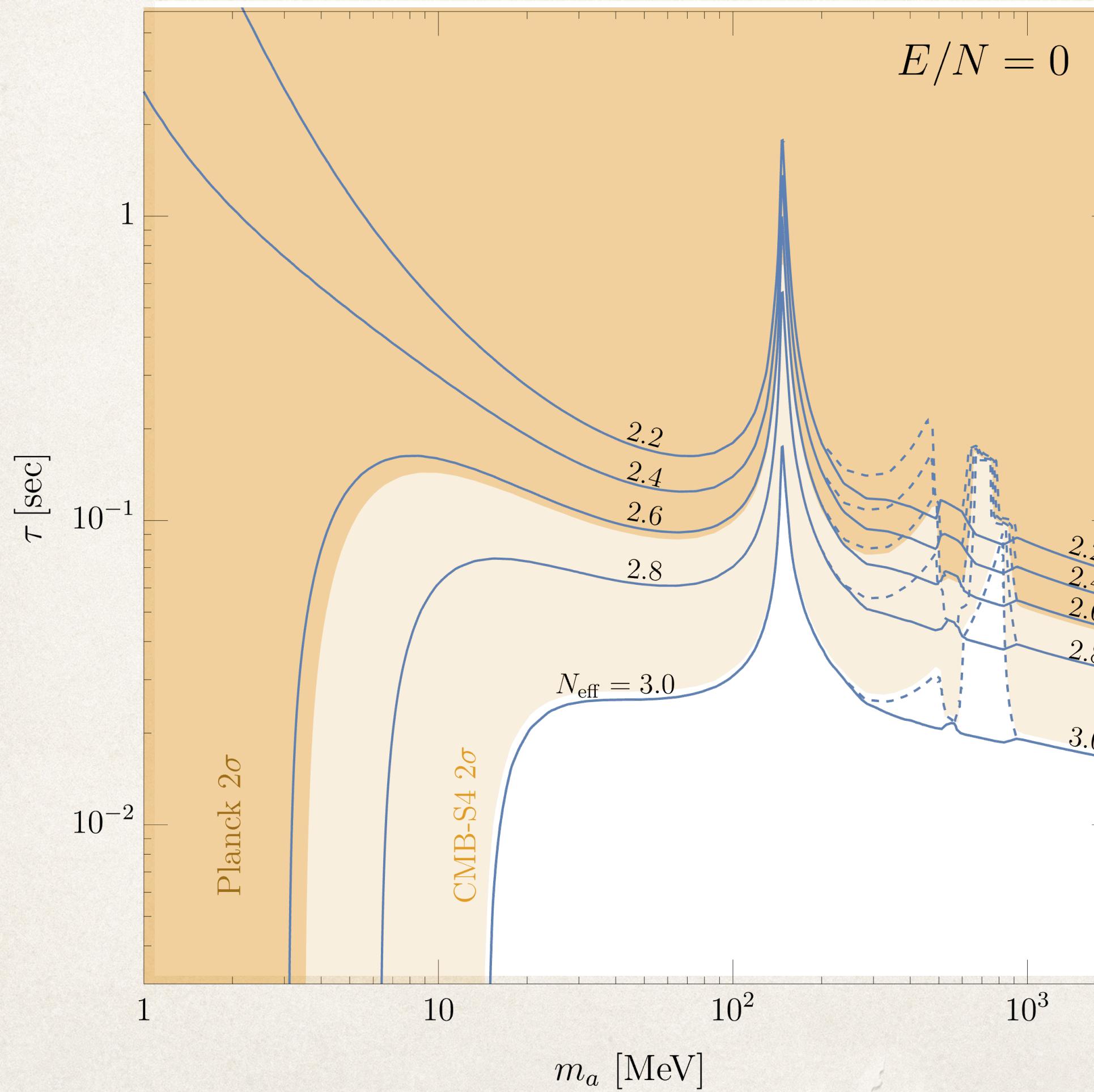
Differences Between Heavy QCD Axions and Heavy ALPs



Cosmological History of Heavy QCD Axion

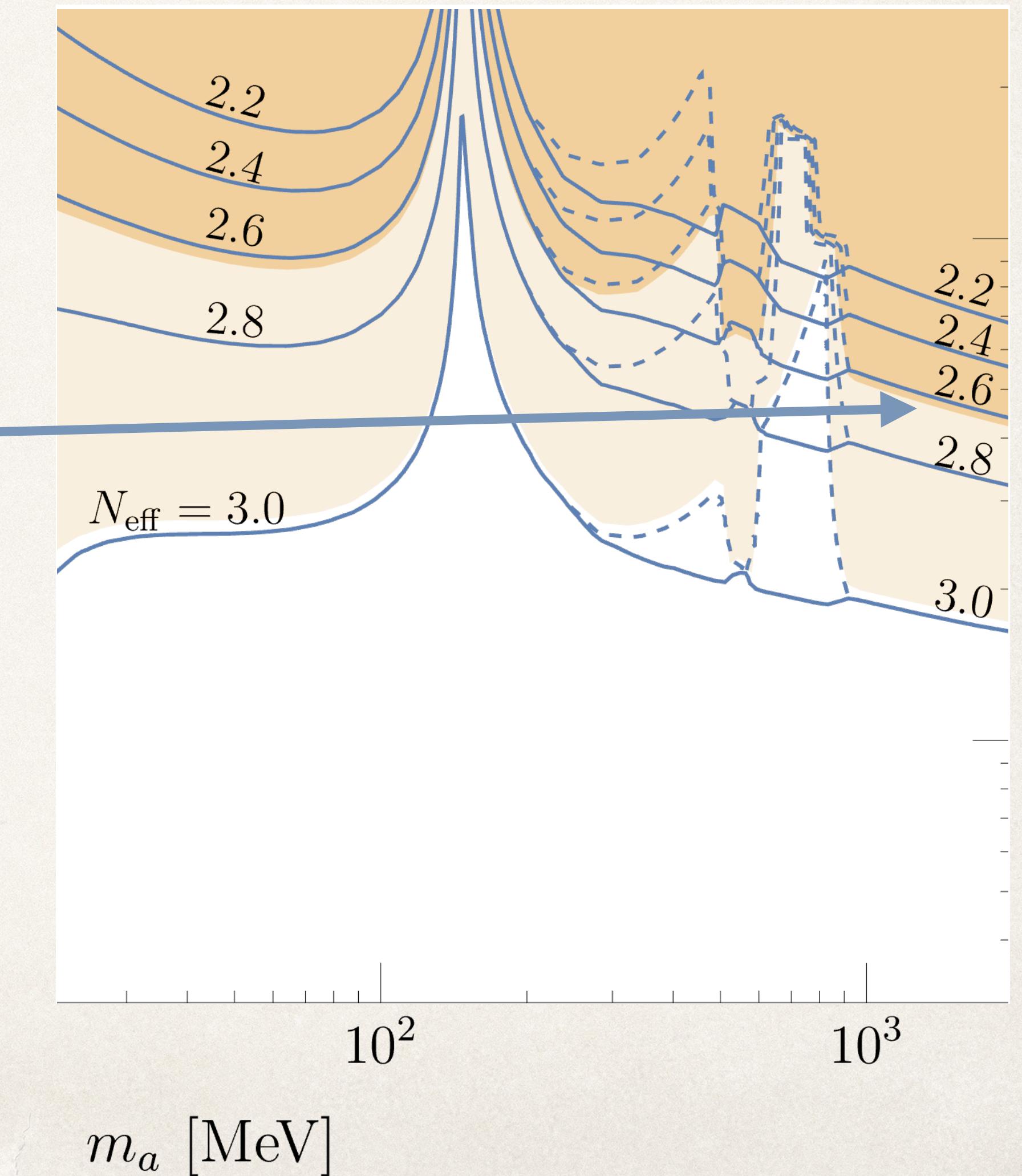


N_{eff} Results: Boltzmann Calculation

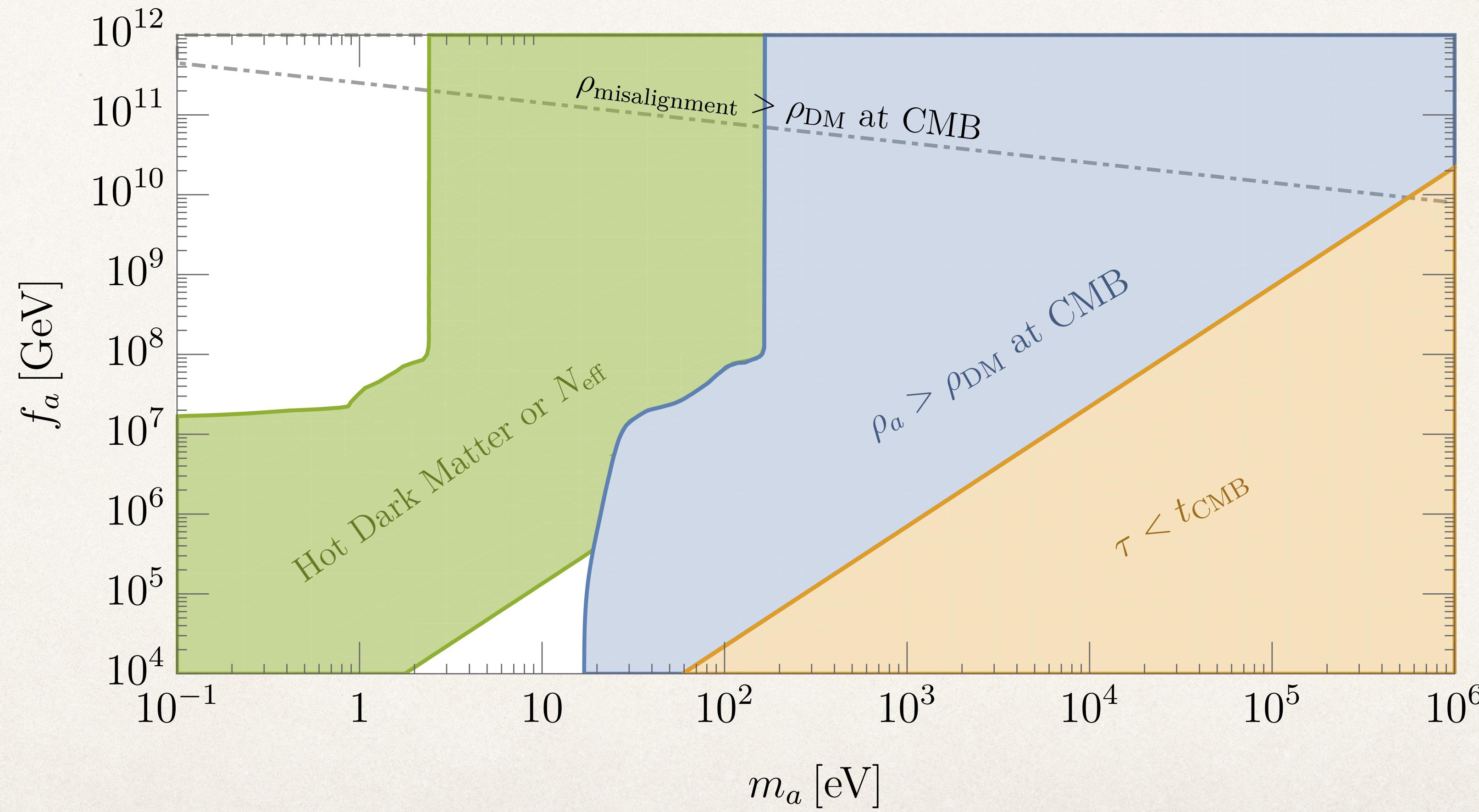


High m_a region

- For $m_a > 2 \text{ GeV}$, axion decays dominantly to gluons
- Note slight decrease in slope in $m_a - \tau$ plane
- For fixed, τ , axion energy density at decay increases with increasing mass
- Axion must lie on further tail of exponential suppression

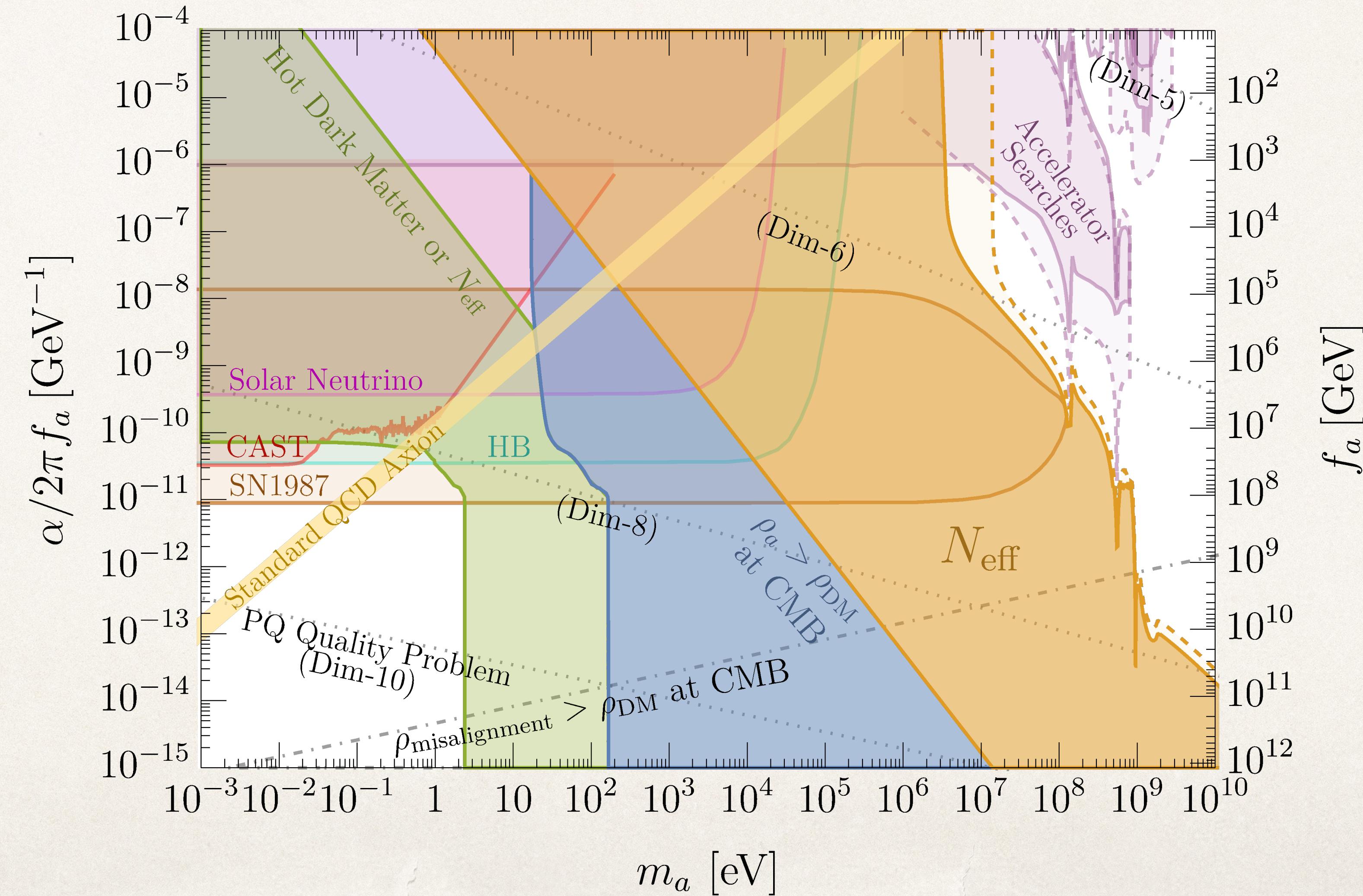


Low m_a region: Hot Dark Matter and N_{eff}



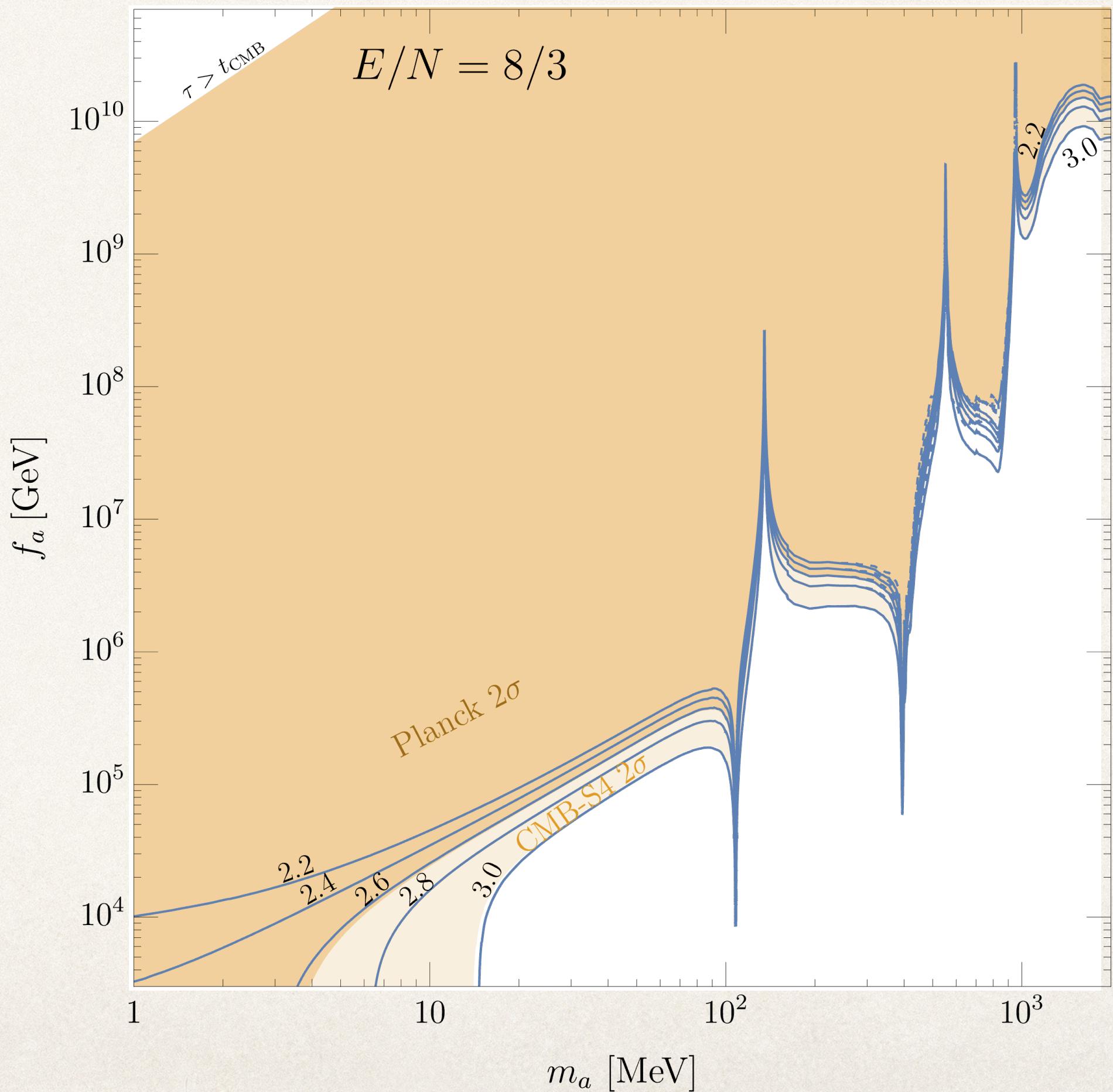
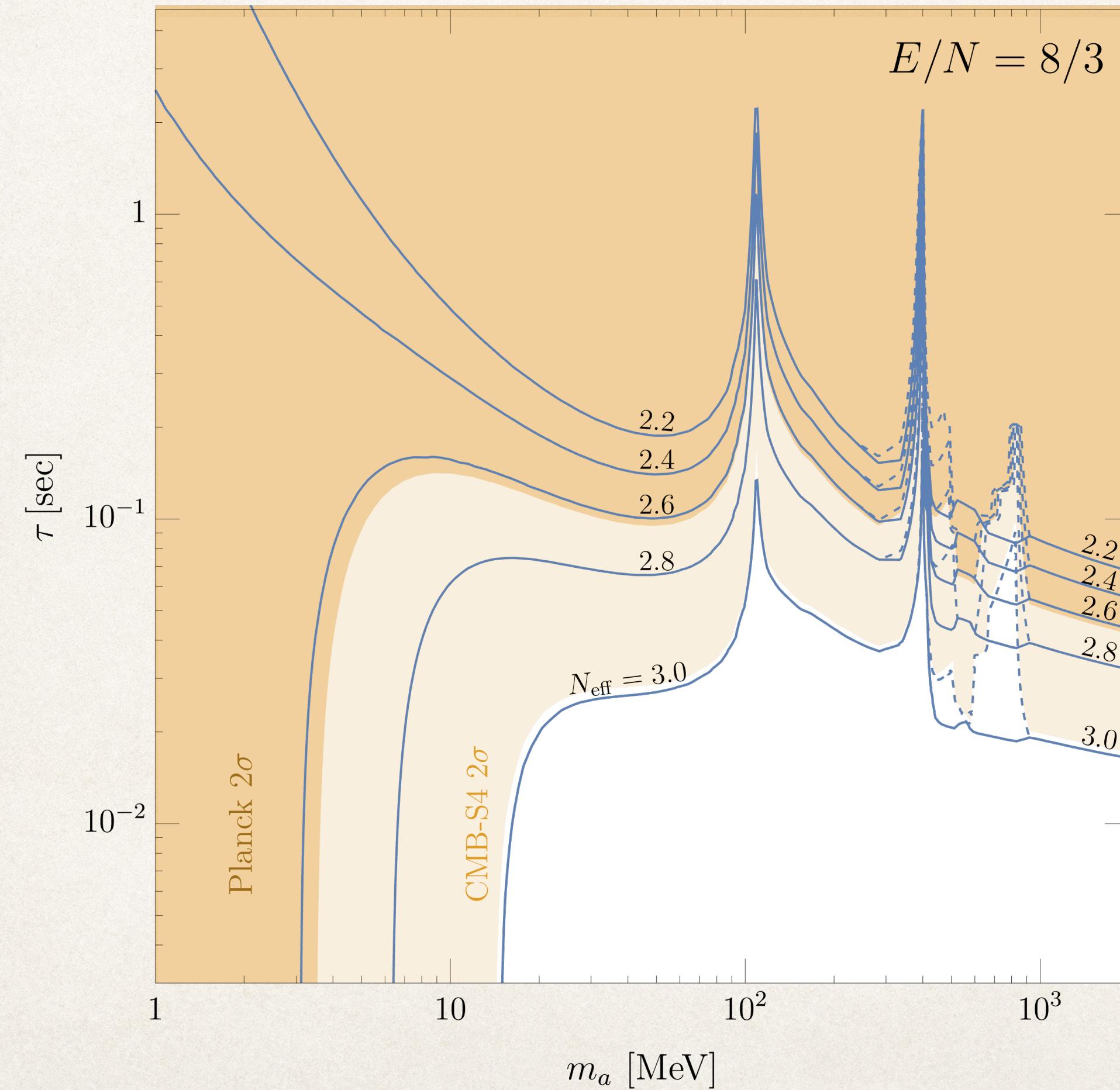
*Hot dark matter bounds
Reinterpreted from
Xu et al 2021*

Summary

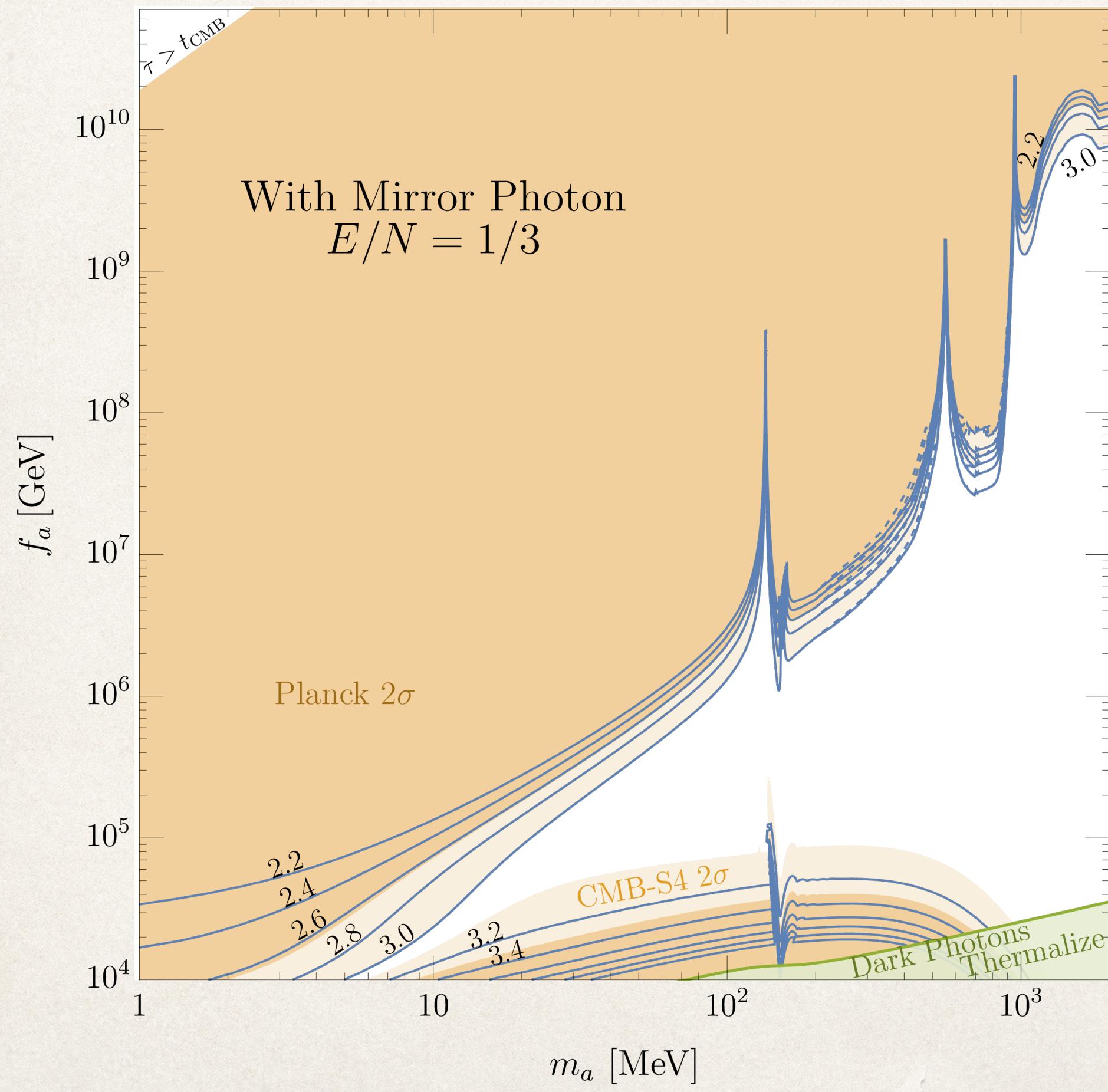


Thank you!

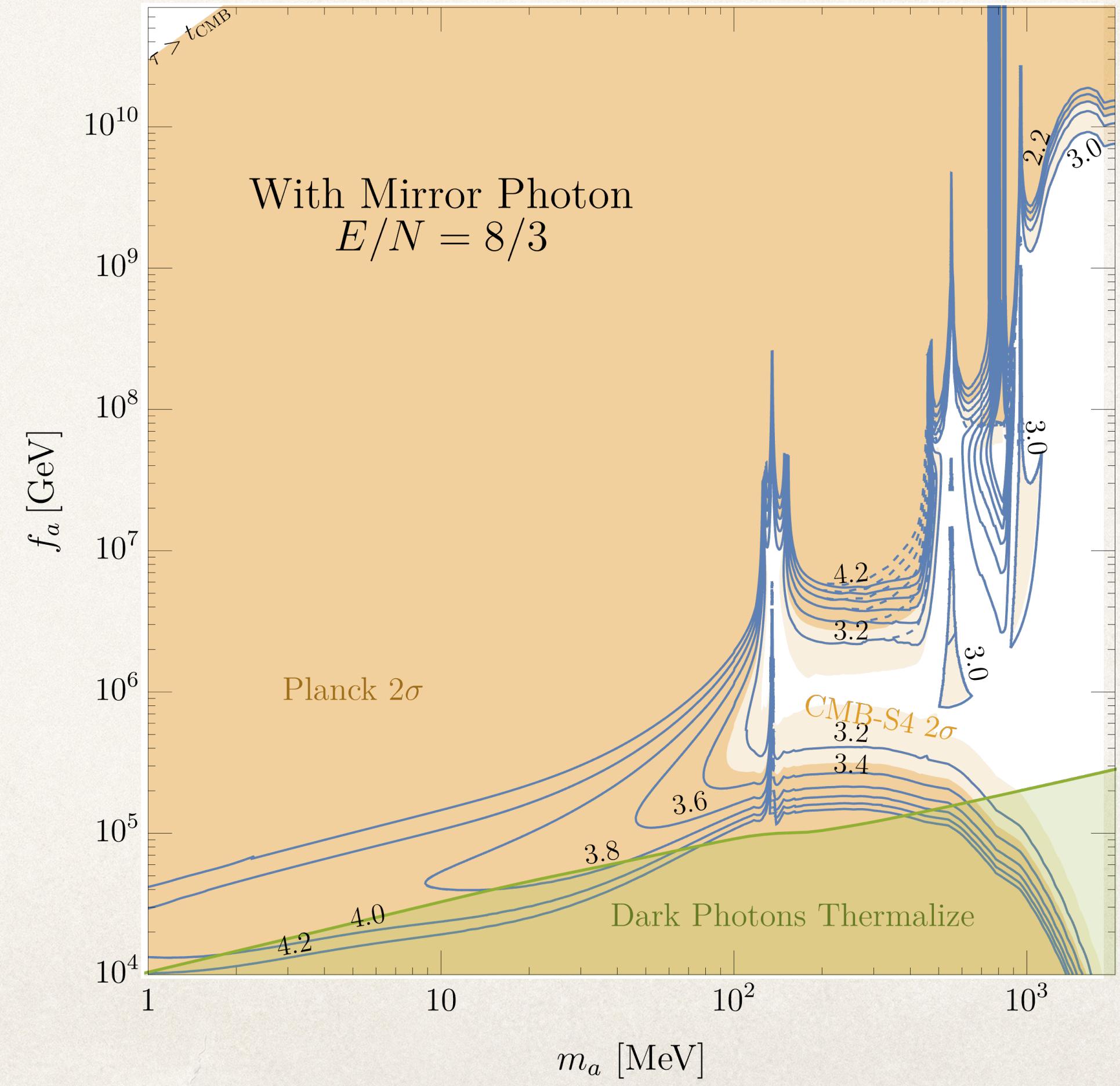
$$E/N = 8/3$$



Incorporation of Mirror Photon



With Mirror Photon
 $E/N = 1/3$



With Mirror Photon
 $E/N = 8/3$