

# QCD axion coupling at finite density

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FÜR PHYSIK



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- **Strong CP problem of QCD:** By observations of the neutron EDM the theta parameter is known to be  $|\bar{\theta}| \lesssim 10^{-10}$

$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{q} \left( i\gamma^\mu D_\mu - m_q e^{i\theta_q} \right) q - \frac{1}{4} G^{a\mu\nu} G_{\mu\nu}^a + \theta \frac{g_s^2}{32\pi^2} G^{a\mu\nu} \tilde{G}_{\mu\nu}^a, \quad \bar{\theta} = \theta + \theta_q$$

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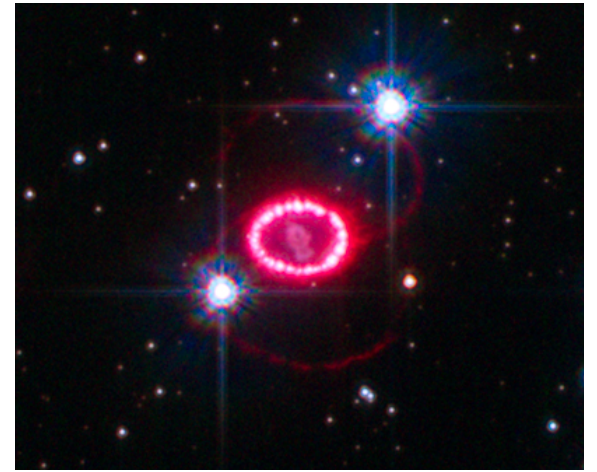
- $\rightarrow \kappa$  can be chosen such that the  $\bar{\theta}$  term is removed
- **Phenomenology determined by one parameter  $f_a$**
- Many ongoing experiments try to search for the (QCD) axion
- Strong bounds on  $f_a$  from SN and NS cooling

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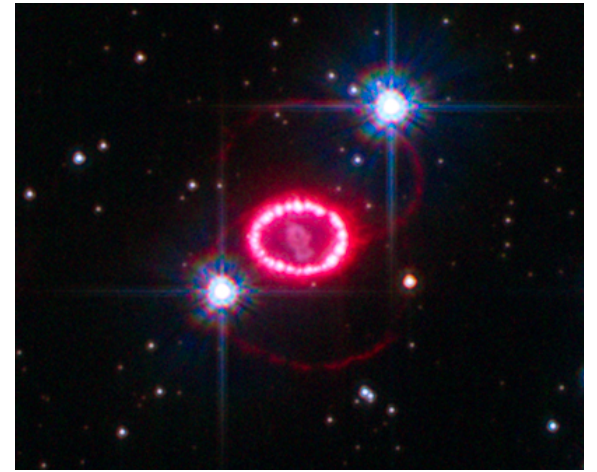
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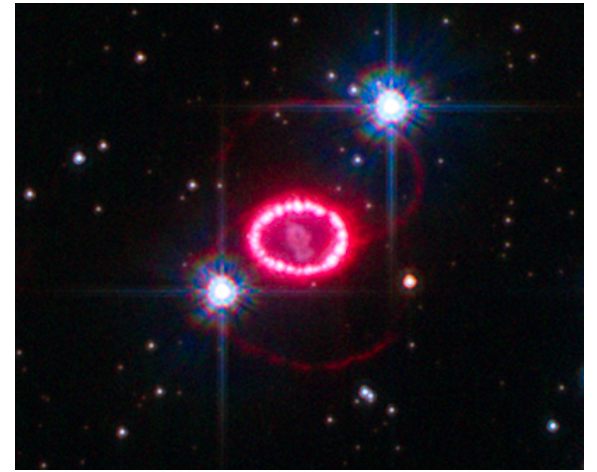
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- For the KSVZ axion this corresponds to  $m_a \lesssim 15 \text{ meV}$

# Supernova bound

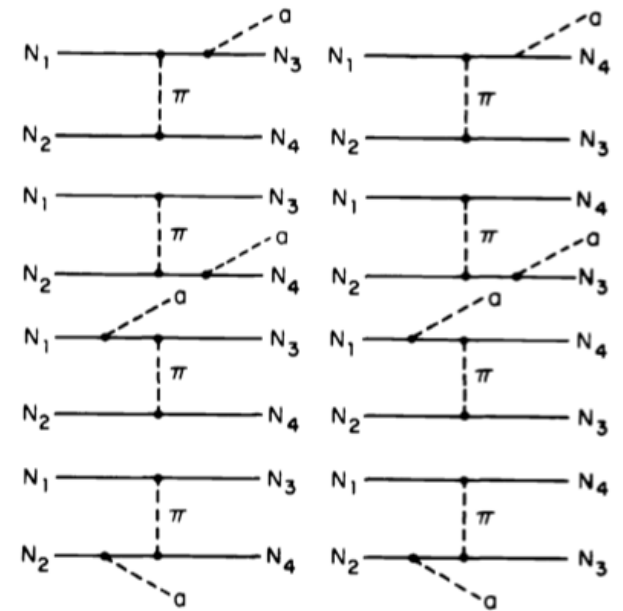
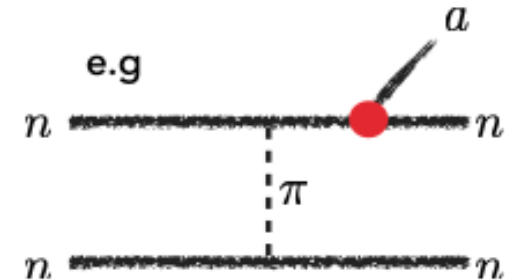
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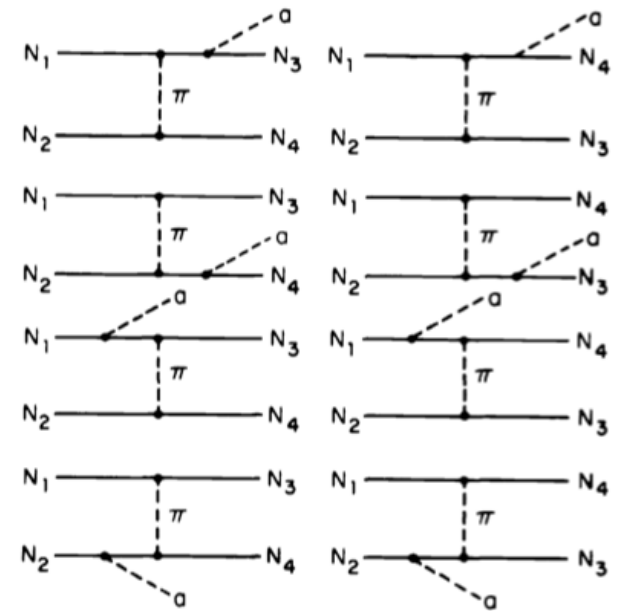
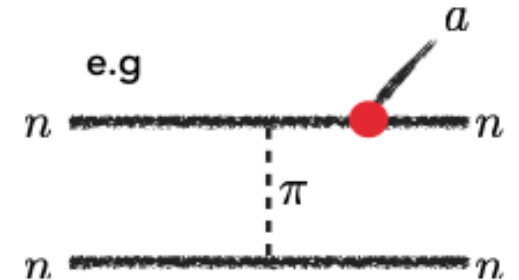
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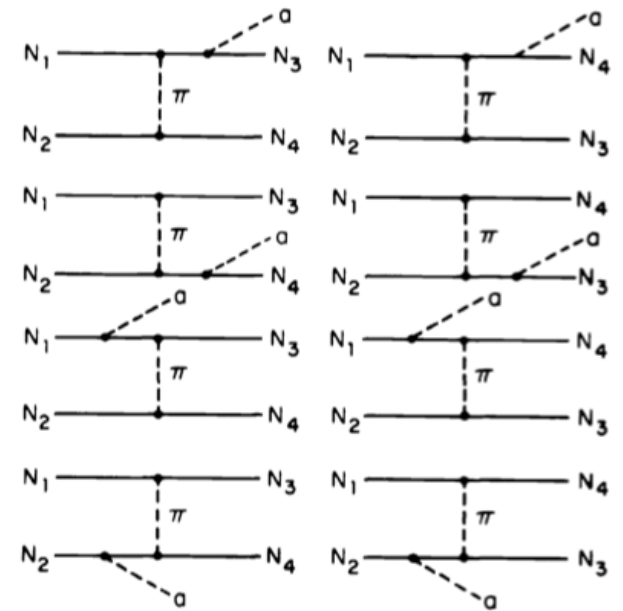
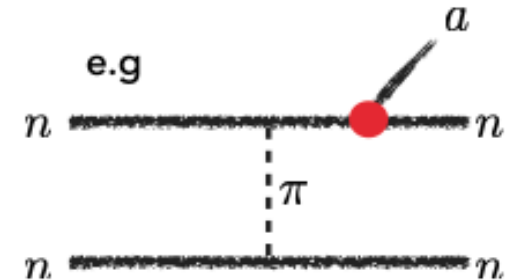
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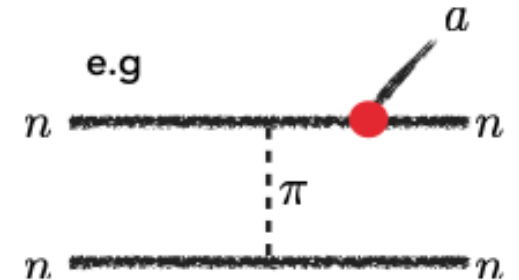
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- Density effects are also highly relevant for neutron star cooling



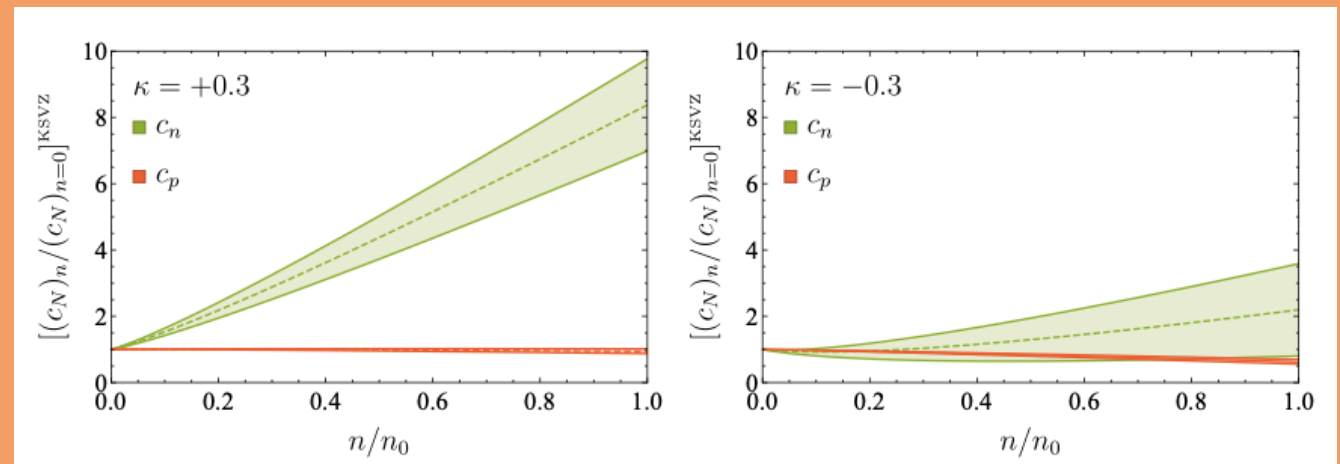
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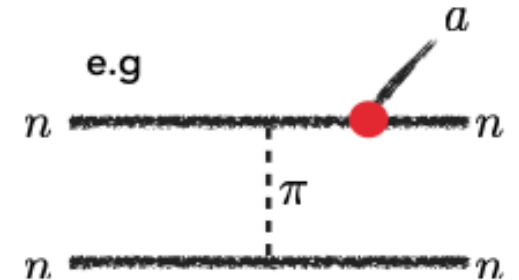
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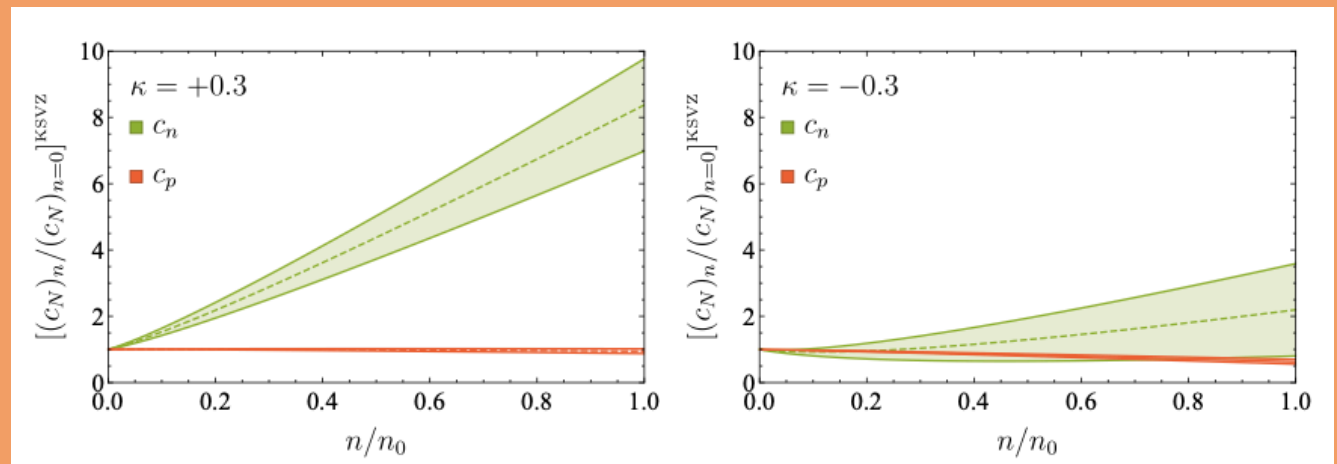
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**➔ We now calculate this systematically**



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- Degrees of freedom are mesons (and baryons)
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- Gauge bosons and other fields (e.g. axion, neutrino) can be consistently added to the theory

# Heavy baryon ChPT + finite density

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- Gives a systematic expansion in density  $\frac{k_f^3}{(4\pi f_\pi)^2 \Lambda_\chi} \sim \frac{n}{(4\pi f_\pi)^2 \Lambda_\chi}$



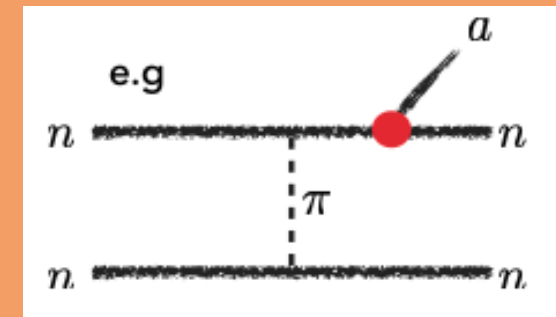
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**→ We use this to calculate the one-pion-exchange (OPE) process  $(N + N \rightarrow N + N + a)$  at leading order finite density**



(Tree level OPE process)

# Coupling axion-nucleon (KSVZ)

- Tree level Lagrangian

$$\mathcal{L}_{\pi N}^{(1)} \supset g_A \bar{N} S^\mu u_\mu N + g_0^i \bar{N} S^\mu \hat{u}_\mu^i N, \quad N = (p, n)^T$$

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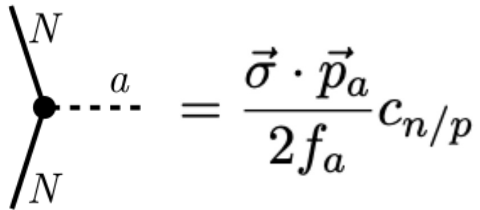
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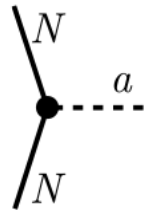
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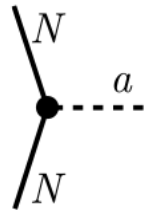
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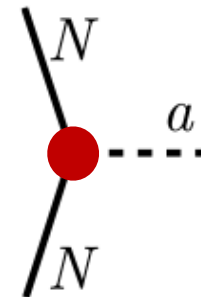
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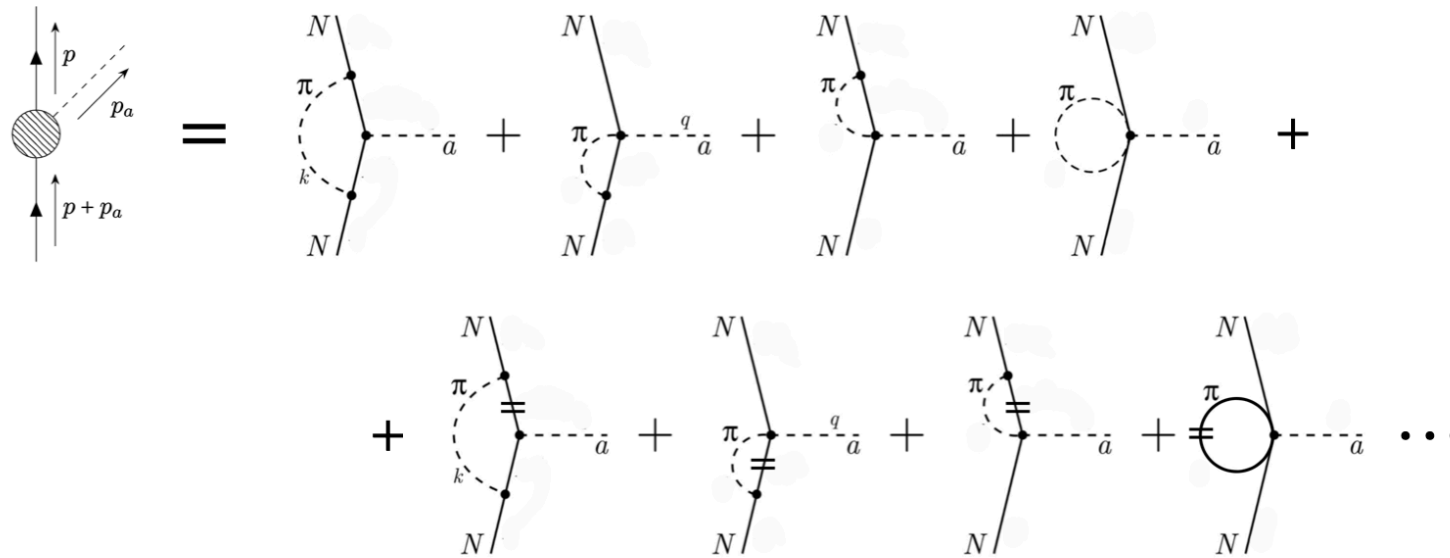
Q: How do they look at finite density?

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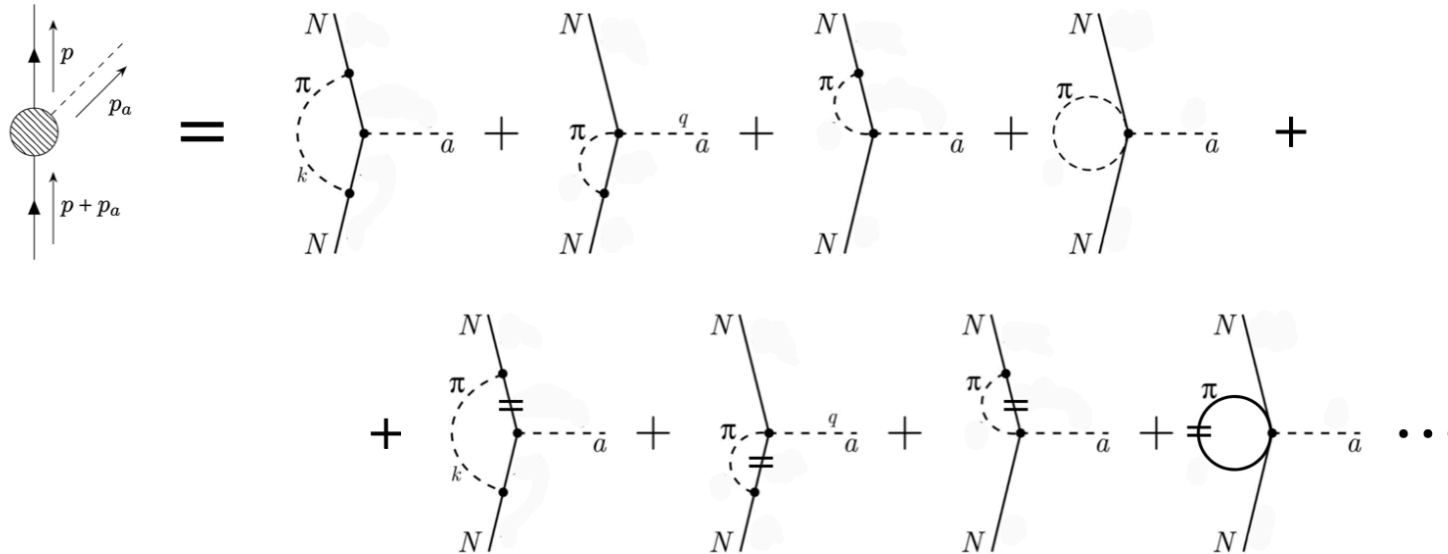


# Vertex corrections





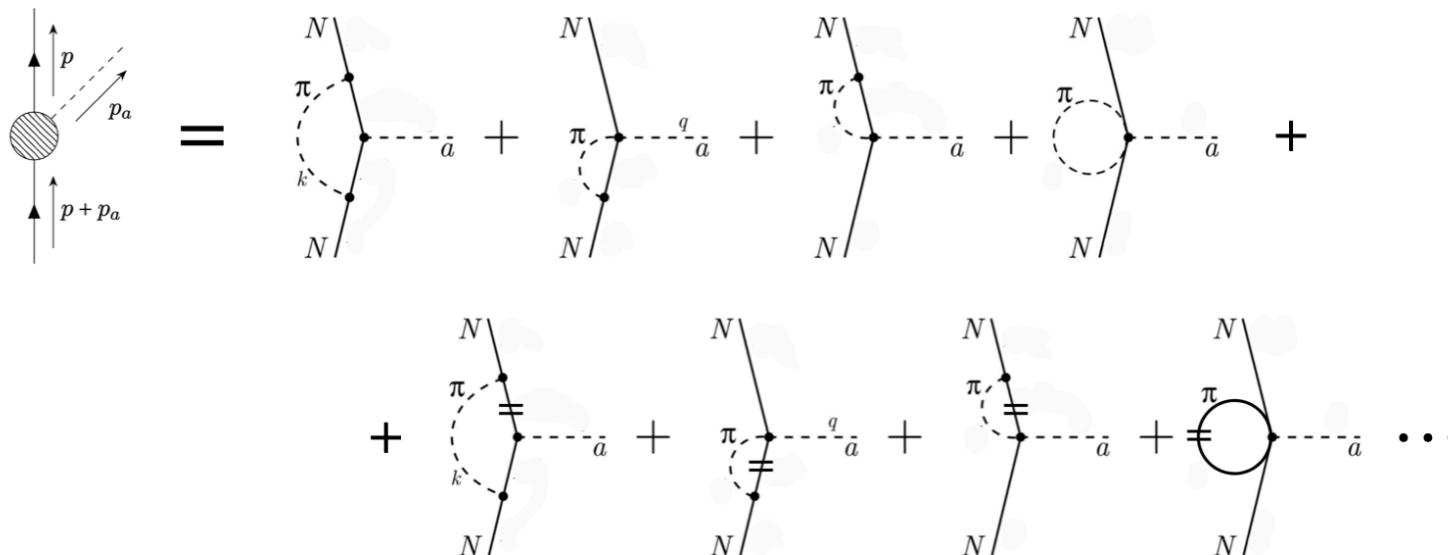
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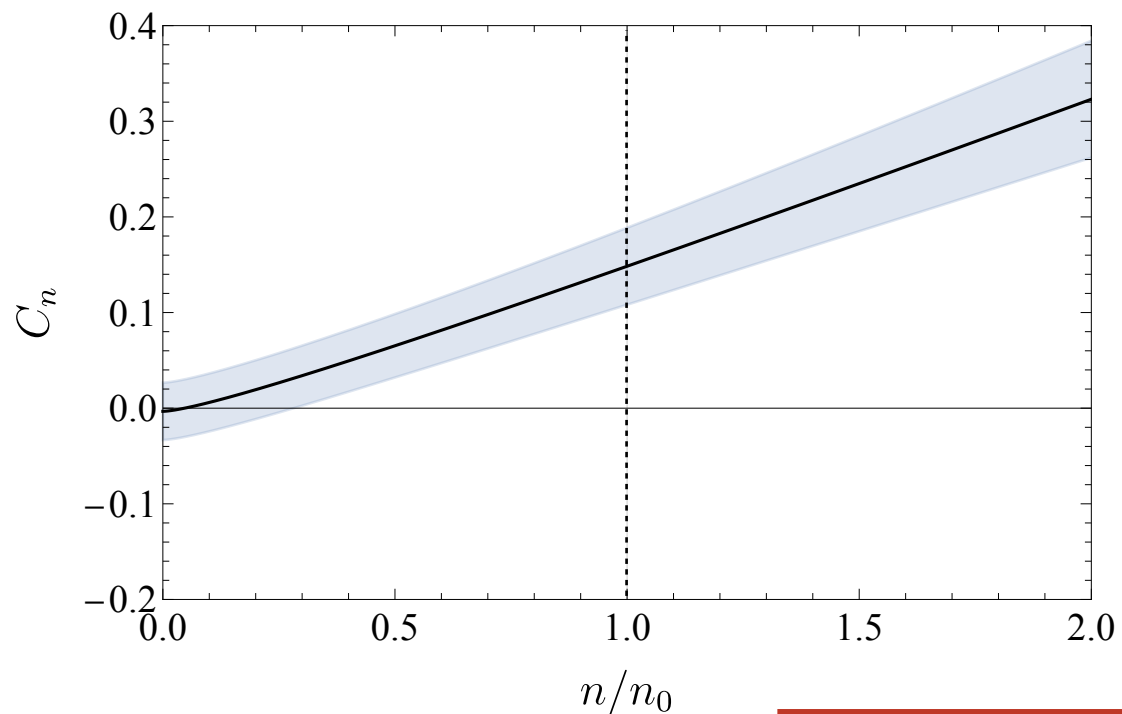
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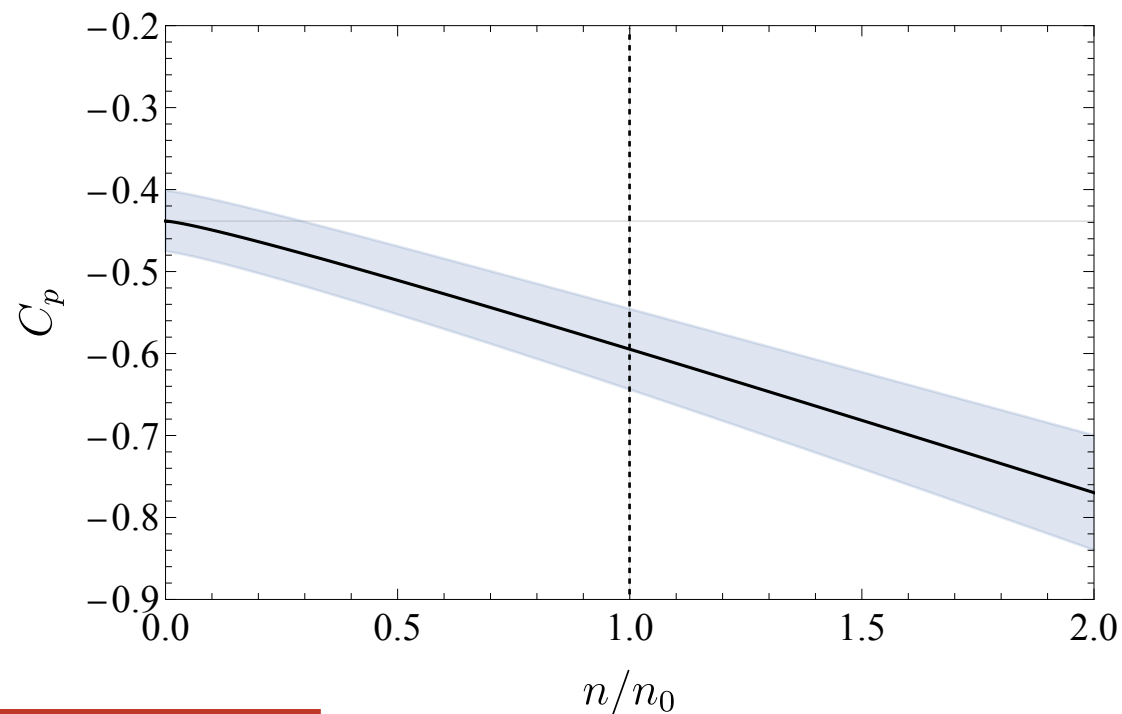
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**Neutron**



**Proton**



$n_0$ : nuclear saturation density

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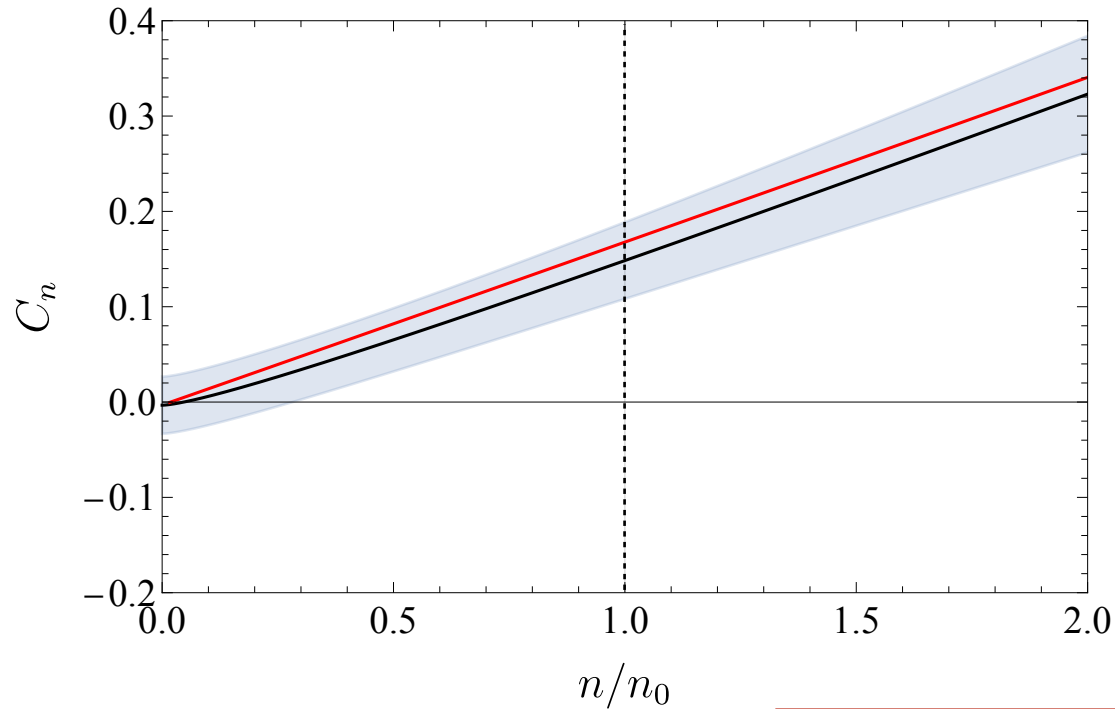


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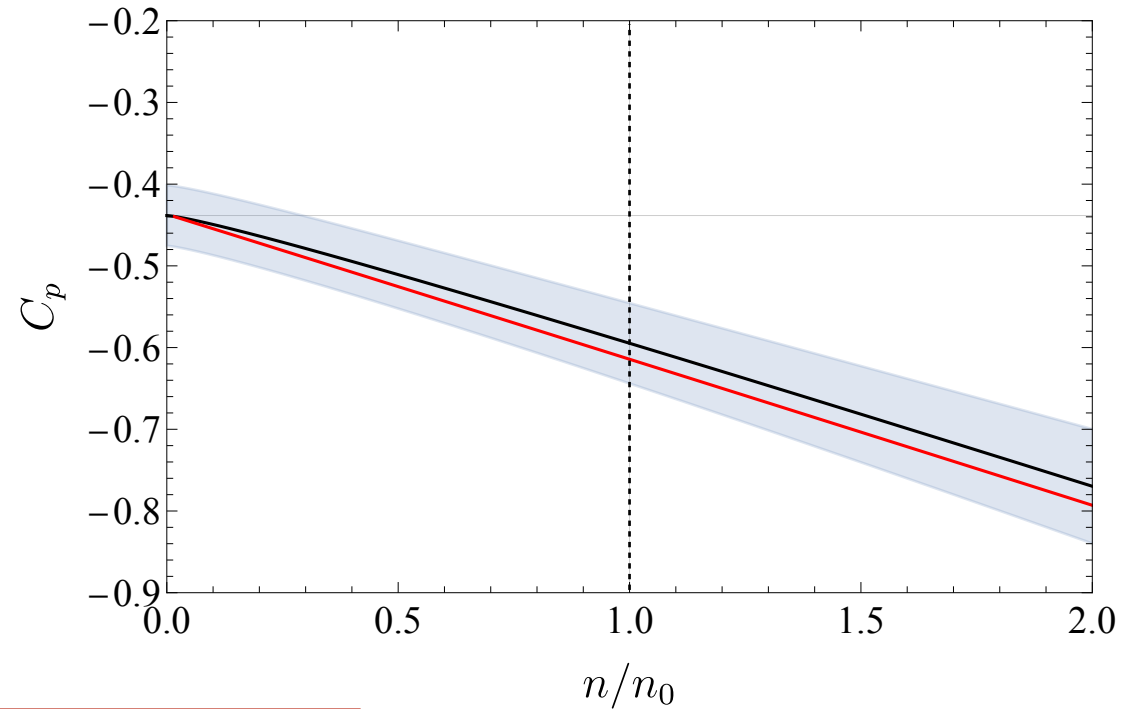
Black:  $T = 0$

Red:  $T = 50 \text{ MeV}$

### Neutron



### Proton



$n_0$  ... nuclear saturation density

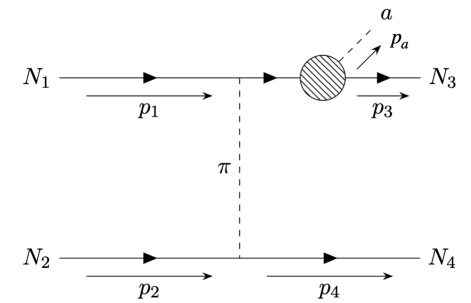
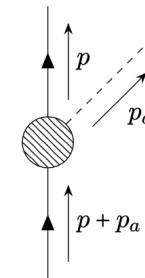
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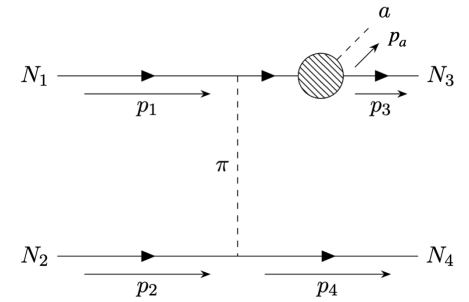
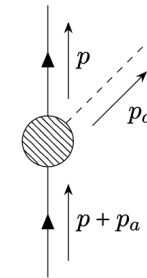
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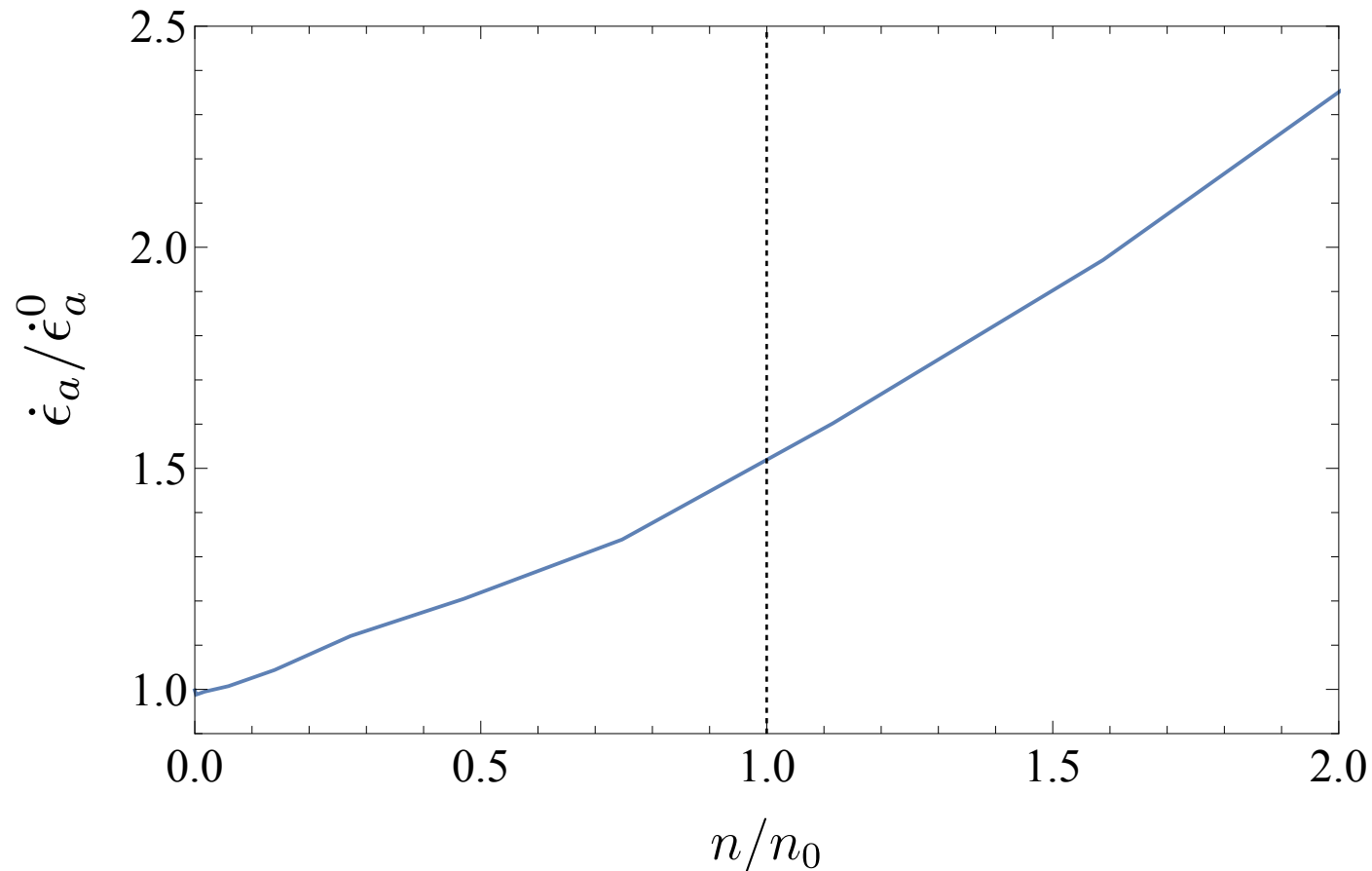
- This alters the axion emissivity  $\dot{\varepsilon}_a$ :

$$\dot{\varepsilon}_a = \int d\Pi_1 d\Pi_2 d\Pi_3 d\Pi_4 d\Pi_a (2\pi)^4 S |M|^2 \delta^{(4)}(p_1 + p_2 - p_3 - p_4 - p_a) E_a f_1 f_2 (1 - f_3) (1 - f_4)$$

  
 Density dependent



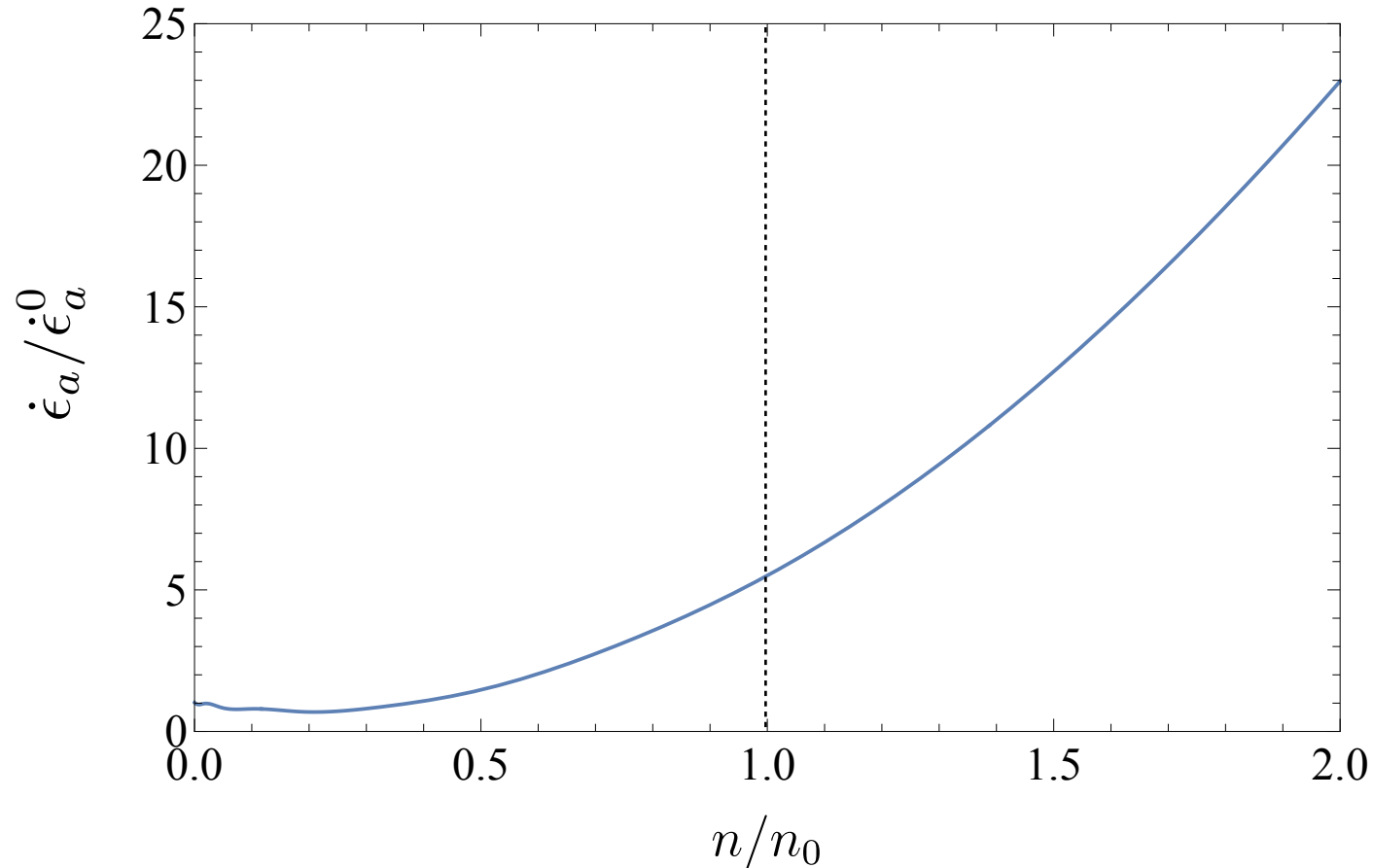
# Emissivity - proton-neutron matter (e.g. SN)



Mixed nuclear matter (e.g. SN)

$$T = 30 \text{ MeV and } \mu/T = -3$$

# Emissivity - pure neutron matter (e.g. NS)



Pure neutron matter (e.g. NS)

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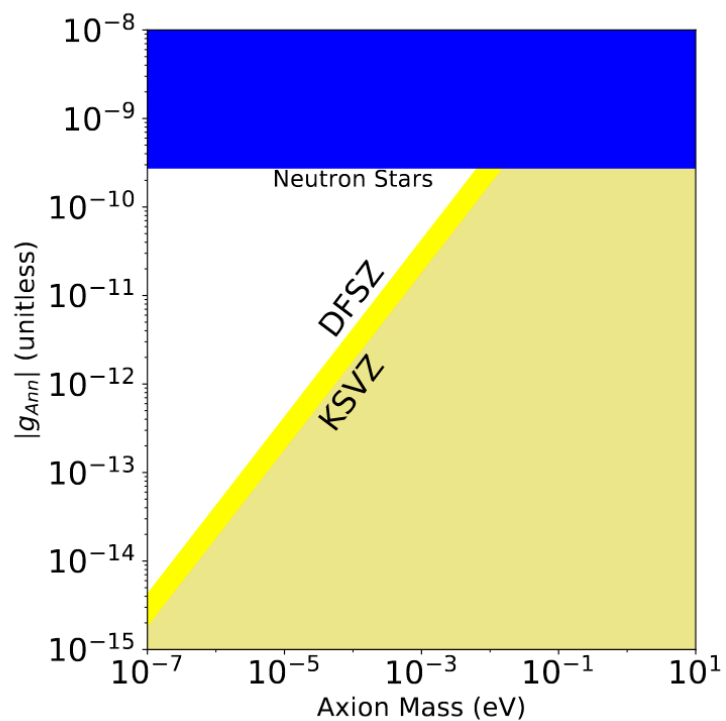
$$g_{Ann} = C_n \frac{m_N}{f_a}$$



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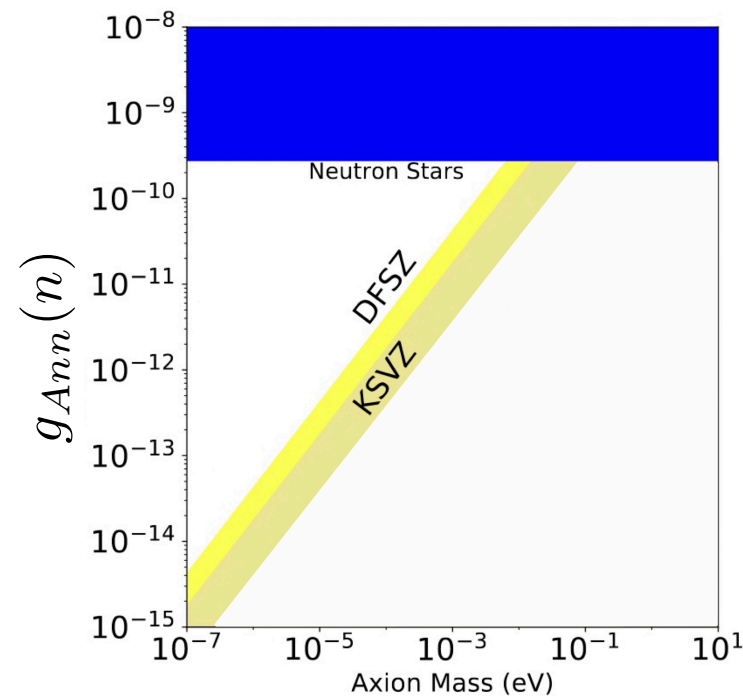
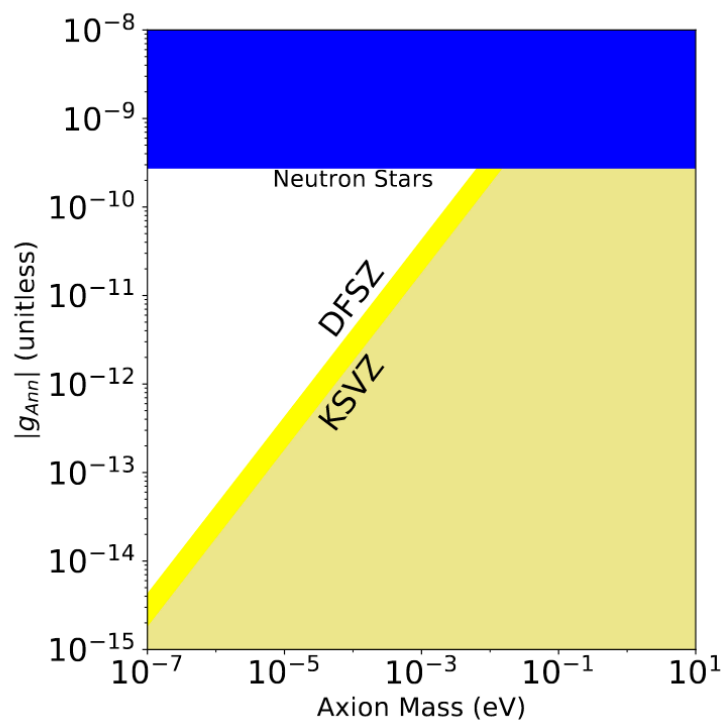
$$g_{Ann} = C_n \frac{m_N}{f_a}$$



# Consequences for NS cooling

- Axion neutron coupling at finite density is no longer compatible with zero

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- Long term goal: calculate the whole  $N + N \rightarrow N + N + a$  process @ finite density