

Sharp turns in axion monodromy: primordial black holes and gravitational waves

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[arXiv: 2205.06065 \(SB, Ivonne Zavala\)](https://arxiv.org/abs/2205.06065)

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Motivation

- Inflation: single field models can simply and elegantly explain the scale invariant power spectra suggested by CMB $\rightarrow n_s, r.$ (1303.3787)

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 - Many theoretically inspired single field models ruled out by better constraints on scalar tilt and tensor fluctuations in Planck (Planck:Inflation). Future CMB surveys will constrain more.
 - Features via deviation from slow-roll \rightarrow require to implement fine-tuned inflection points/bumps in potential (1702.03901, 1911.00057). Interesting consequences: PBH, GW
- **Why multifield?**
 - Effective single field behaviour at large scales (near CMB pivot) can still satisfy Planck constraints.
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- **Goal and why?**
 - **Is it possible to construct a multifield model in supergravity with minimal fine-tuning?**
 - **Features with sharp turns in field space, but small field-space curvature. Consequences for the power spectra? PBH, induced GW?**

Multifield inflation: The model

- Can avoid η problem. (1908.09797)
- Can be consistent with swampland conjecture. (1901.08603)
- Construct a model: (i) CMB consistent at large scales; (ii) features lead to growth of power spectrum at small scales \rightarrow abundant PBH formation, large induced GW spectra.

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$$V = e^{K/M_{\text{Pl}}^2} \left(K^{i\bar{j}} D_i W \overline{D_{\bar{j}} W} - 3|W|^2 M_{\text{Pl}}^{-2} \right)$$

$$K/M_{\text{Pl}}^2 = -\alpha \log[(\Phi + \bar{\Phi})/M_{\text{Pl}} - \beta S \bar{S}/M_{\text{Pl}}]$$

$$W = S(M\Phi + i\lambda e^{-b\Phi})$$

with $\rho \equiv \text{Re}(\Phi)$ and $\theta \equiv \text{Im}(\Phi)$ leading to

$$V = \frac{M^2}{\beta} \left(\rho^2 + \theta^2 + \frac{2\lambda}{M} e^{-b\rho} \left[\theta \cos(b\theta) + \rho \sin(b\theta) + \frac{\lambda}{2M} e^{-b\rho} \right] \right)$$

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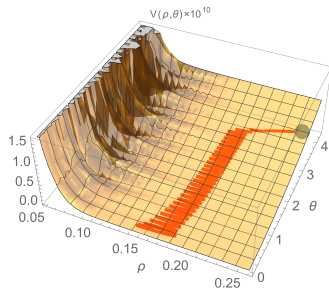
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- The potential is not just the sum of two potentials, but they are coupled: allows gentle cliffs and plateaus when the saxion is slightly away from its minimum.
- Turning rate: $\omega \equiv \frac{\Omega}{H} \equiv -N_a(D_t T^a)/H = \frac{V_{,N}}{H\dot{\phi}}$, where $T^a = \frac{\dot{\phi}^a}{\dot{\phi}}$.

Multifield inflation: The model

- Multiple sharp turns in (ρ, θ) space, without requiring large field space curvature.
- Repeated inflection points due to oscillations in the potential \rightarrow slow-roll violation $\eta > 1 \rightarrow$ large turns $\omega > 1$ and sharp turns $\nu \gg 1$.



$$\epsilon \equiv -\frac{\dot{H}}{H^2} = \frac{\dot{\varphi}^2}{2M_{Pl}^2 H^2}$$

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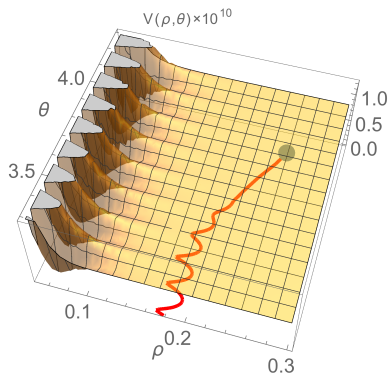
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$$\frac{V_{TT}}{3H^2} = \frac{\Omega^2}{3H^2} + 2\epsilon - \frac{\eta}{2} - \frac{\xi_\varphi}{3}$$

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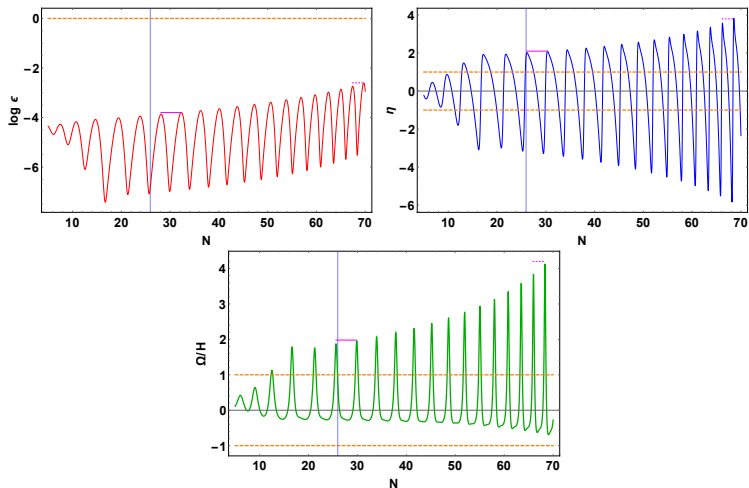
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Slow Roll Parameters

For $\lambda/M = 80$, $b = 50$:



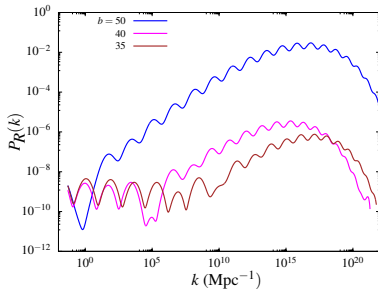
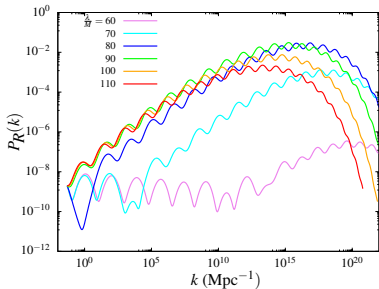
Oscillations with decreasing period.

Power spectra

$$\ddot{Q}_T + 3H\dot{Q}_T + \left(\frac{k^2}{a^2} + m_T^2\right) Q_T = (2\omega H Q_N) \dot{} - \left(\frac{\dot{H}}{H} + \frac{V_T}{\dot{\phi}}\right) 2\omega H Q_N,$$

$$\ddot{Q}_N + 3H\dot{Q}_N + \left(\frac{k^2}{a^2} + m_N^2\right) Q_N = -2\omega\dot{\phi}\dot{R}$$

$$\mathcal{R} = \frac{H}{\dot{\phi}} Q_T.$$

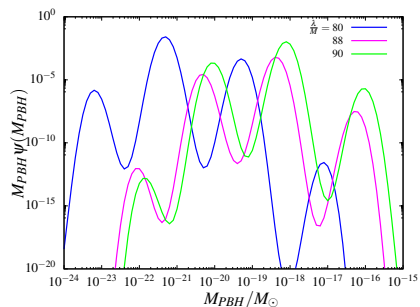


For these examples, $N_{\text{inf}} \sim 55 - 65$, $r \sim 0.010 - 0.024$, $V_{\text{inf}}^{1/4} \sim 0.003 - 0.03 M_{\text{Pl}}$. For CMB normalization, $\beta = 1$, $M \sim 10^{-8} - 10^{-6} M_{\text{Pl}}$.

Light PBHs in abundance

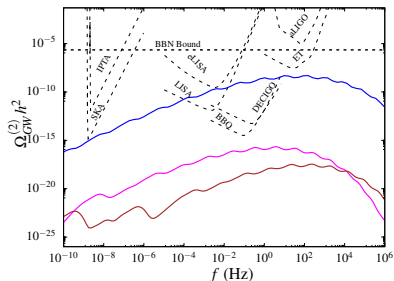
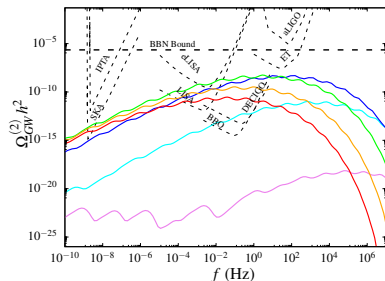
$$\psi(M_{\text{PBH}}) = \frac{\gamma}{T_{\text{eq}}} \left(\frac{g_s(T_1)}{g_s(T_{\text{eq}})} \right)^{\frac{1}{3}} \left(\frac{\Omega_m h^2}{\Omega_c h^2} \right) \left(\frac{90 M_{\text{Pl}}^2}{\pi^2 g_*(T_1)} \right)^{\frac{1}{4}} (4\pi\gamma M_{\text{Pl}}^2)^{\frac{1}{2}} \frac{\beta_{\text{PBH}}(M_{\text{PBH}})}{M_{\text{PBH}}^{\frac{3}{2}}}.$$

- $M_{\text{PBH}} \sim k^{-2}$ in RD.
- $f_{\text{PBH}} \sim 10^{-3} - 10^{-2}$.
- BBN constraint : $f_{\text{PBH}} \lesssim 10^{-4}$; CMB and γ -ray observations constrain $f_{\text{PBH}} < 10^{-10}$ for $10^{-20} M_{\odot} \lesssim M_{\text{PBH}} \lesssim 10^{-17} M_{\odot}$ (monochromatic $\psi(M_{\text{PBH}})$).
- Some target M_{PBH} are produced more than the others.



Induced GW

- $P_h(k, \tau) \propto \int du \int dv \mathcal{I}_R D P_{\mathcal{R}}(ku) P_{\mathcal{R}}(kv)$.
- $\Omega_{\text{GW}}^{(2)}(k, \tau) \sim k^2 \tau^2 P_h(k, \tau)$.



- Wide and large GW spectra, can be probed by multiple future surveys together.
- Spectral shape is characteristic, can be analysed with GW surveys.

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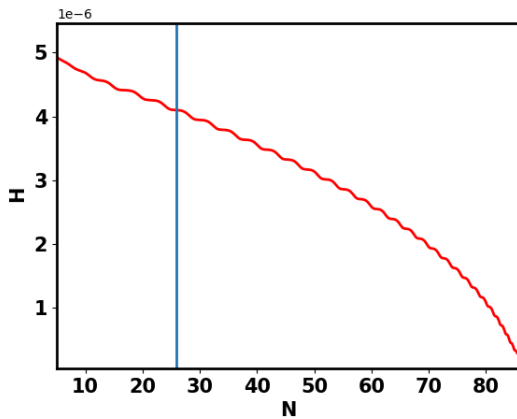
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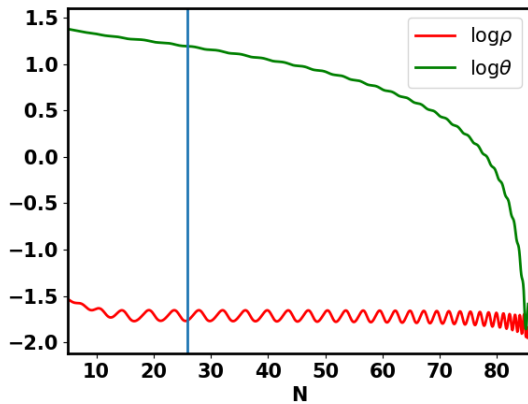
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- Induced GW spectra $\Omega_{\text{GW}}^{(2)}$ has a large and wide profile, sometimes appearing in multiple future sensitivity curves. Spectral shape may be probed with future GW surveys.

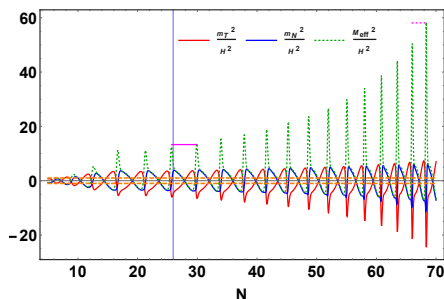
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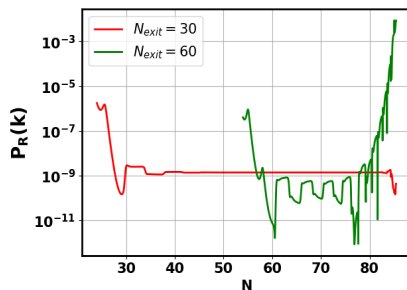




$$\frac{m_T^2}{H^2} \equiv -\frac{3}{2}\eta - \frac{1}{4}\eta^2 - \frac{1}{2}\epsilon\eta - \frac{1}{2}\frac{\dot{\eta}}{H}$$

$$\frac{m_N^2}{H^2} = \frac{V_{NN}}{H^2} + M_{\text{Pl}}^2 \epsilon \mathbb{R} - \omega^2$$

$$M_{\text{eff}}^2 \equiv m_N^2 + 4H^2\omega^2 = V_{NN} + M_{\text{Pl}}^2 \epsilon \mathbb{R} H^2 + 3H^2\omega^2.$$



$$\frac{V_{TT}}{3H^2} = \frac{\Omega^2}{3H^2} + 2\epsilon - \frac{\eta}{2} - \frac{\xi_\varphi}{3} \quad (1)$$

$$\frac{V_{TN}}{3H^2} = \omega \left(1 - \epsilon + \frac{\eta}{3} + \frac{\nu}{3} \right). \quad (2)$$

M	λ/M	b	ρ_{ini}	θ_{ini}	N_{inf}	r	$V_{\text{inf}}^{1/4}$
2.52×10^{-6}	60	50	0.250	4.20	64.77	0.010	0.0029
2.73×10^{-6}	70	50	0.250	4.20	62.32	0.016	0.0030
2.15×10^{-6}	80	50	0.245	4.20	59.48	0.018	0.0027
6.41×10^{-7}	90	50	0.250	4.20	57.49	0.020	0.0015
1.10×10^{-7}	100	50	0.250	4.20	56.07	0.022	0.0006
1.25×10^{-8}	110	50	0.250	4.20	55.06	0.024	0.0002
1.60×10^{-6}	80	40	0.250	4.50	63.63	0.011	0.0026
1.60×10^{-6}	80	35	0.400	5.50	56.99	0.012	0.0026