

Polynomial Inflation and Its Aftermath

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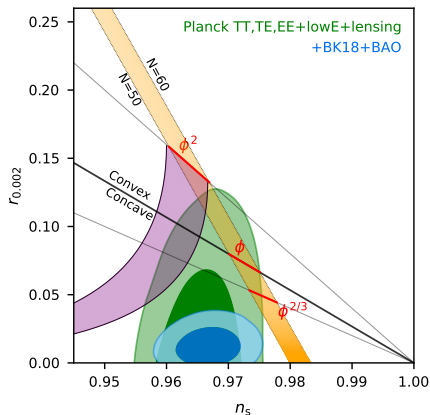
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Outline

1. Polynomial Inflation and Predictions
2. (P)reheating
3. Dark Matter Production
4. Baryogenesis
5. Summary

Current Status of Monomial Inflation [2110.00483]



- Monomial: $V(\phi) \sim \phi^p$, tensor-to-scalar ratio

$$r \sim \left(\frac{V'}{V} \right)^2 \sim \frac{4p}{N}, \quad N = \int_{t_*}^{t_e} H dt$$

- BICEP/Keck 2018: $r < 0.035 \Rightarrow p \lesssim 0.5$

Polynomial Inflation

- Alternative to monomial scenario

$$V(\phi) = \sum_{n=0}^4 \alpha_n \phi^n$$

- Avoid trans-Planckian $\Rightarrow \phi < M_p \Rightarrow$ Small Field
- Reasonable to insist on renormalizability
- $V(\phi)$: most general renormalizable inflaton potential
- Question:

Can $V(\phi)$ flat enough to match the CMB data?

Polynomial Inflation Analysis

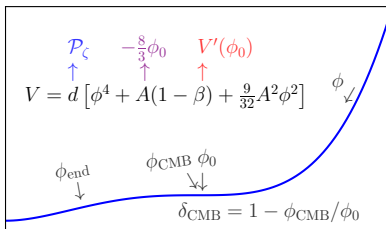
$$V(\phi) = d\phi^4 + c\phi^3 + b\phi^2$$

- Inflection-point at $\phi_0 = -\frac{3c}{8d}$ with $b = \frac{9c^2}{32d}$
- Reparametrization (d, A, β)

$$\begin{aligned} V(\phi) &= d \left[\phi^4 + \frac{c}{d} (1 - \beta) \phi^3 + \frac{9}{32} \left(\frac{c}{d} \right)^2 \phi^2 \right] \\ &\equiv d \left[\phi^4 + A(1 - \beta) \phi^3 + \frac{9}{32} A^2 \phi^2 \right] \end{aligned}$$

- $A \equiv c/d = -8/3\phi_0 \leftrightarrow$ location ϕ_0
- β : \leftrightarrow flatness $V(\phi_0)$; $V'(\phi_0) = 0$ if $\beta = 0$
- d : \leftrightarrow amplitude (power spectrum)

Slow-Roll Analysis



- $n_s \simeq 1 - 48\delta_{\text{CMB}}/\phi_0^2$
- $N_{\text{CMB}} \propto (\pi/2 - \arctan(\delta_{\text{CMB}}/\sqrt{2\beta}))$
- $r \propto \beta^2$
- $\alpha \simeq -576(2\beta + \delta^2)/\phi_0^4$
- $\mathcal{P}_\zeta \simeq d\phi_0^6/(\delta^2 + 2\beta)^2$
- Fix parameters: $n_s = 0.9649$, $N_{\text{CMB}} = 65$, $\mathcal{P}_\zeta = 2.1 \cdot 10^{-9} \Rightarrow$
 - $\delta_{\text{CMB}} = 7.31 \times 10^{-4} \phi_0^2$
 - $\beta = 9.73 \times 10^{-7} \phi_0^4$
 - $d = 6.61 \times 10^{-16} \phi_0^2$
- Predictions for r and α ?

Predictions for r and α

- r small

$$\frac{r}{7.09 \times 10^{-9} \phi_0^6} = 1 - 3.9 \cdot 10^{-2} (65 - N_{\text{CMB}}) + 15.0 (0.9649 - n_s) + 175 (0.9649 - n_s)^2$$

- α hopefully testable [via S4 CMB, see e.g. 1611.05883]

$$\alpha = -1.43 \cdot 10^{-3} - 5.56 \cdot 10^{-5} (65 - N_{\text{CMB}}) + 0.02 (0.9649 - n_s) - 0.25 (0.9649 - n_s)^2$$

- Central model parameters

$$\delta_{\text{CMB}} = 7.31 \times 10^{-4} \phi_0^2; \beta = 9.73 \times 10^{-7} \phi_0^4; d = 6.61 \times 10^{-16} \phi_0^2$$

- Inflaton mass and Inflationary scale:

$$m_\phi^2 = \left. \frac{\partial^2 V}{\partial \phi^2} \right|_{\phi=0} \simeq 4d \phi_0^2; H_{\text{inf}} = \sqrt{\frac{V(\phi_0)}{3}} \simeq 8.6 \cdot 10^{-9} \phi_0^3$$

- Question: What's the lower bound for ϕ_0 ?

Reheating \Rightarrow Lower Bound

- After inflation ends:
 - ϕ oscillates and transfers energy

- Decays:

$$\mathcal{L} \supset -g\phi|\phi'|^2 - y\phi\bar{\chi}\chi$$

- Rate (with $m_\phi \sim \phi_0^2$):

$$\Gamma_\phi \simeq \frac{g^2}{8\pi m_\phi}; \frac{y^2}{8\pi} m_\phi$$

- Reheating Tem:

$$T_{\text{rh}} \simeq 1.41 g_*^{-1/4} \Gamma_\phi^{1/2}$$

- BBN: $T_{\text{rh}} \gtrsim 4 \text{ MeV}$ [Hannestad '04] \Rightarrow

$$y\phi_0 \gtrsim 4.7 \times 10^{-17}; \frac{g}{\phi_0} \gtrsim 2.4 \times 10^{-24}$$

- **Remarks: Preheating not efficient here**

- EoM for ϕ' :

$$\ddot{\phi}'(\mathbf{k}, t) + (k^2/a^2 + g\phi)\phi'(\mathbf{k}, t) = 0$$

- Though $m_{\phi'}^2 \sim g\phi \Rightarrow$ tachyonic instability [Dufaux, Felder et al. '06], **however self-coupling $\lambda\phi'^4 \Rightarrow \lambda\langle\phi'^2\rangle\phi'^2 \Rightarrow$ positive mass term**

- Pauli blocking for χ

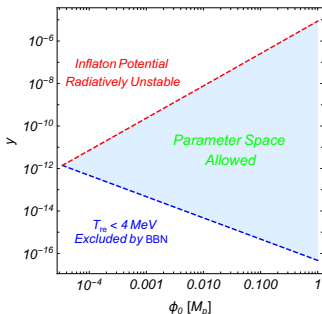
- **Question: upper bounds** for the couplings? \Rightarrow Radiative stability

Radiative Stability \Rightarrow Upper Bound

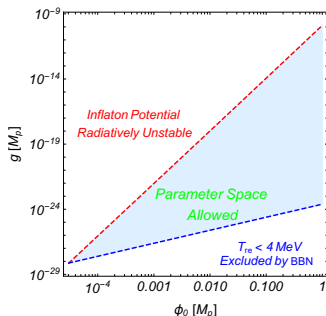
- Require:

$$\Delta V_{1\text{-loop}}(\phi_0) \ll V(\phi_0); \Delta V'_{1\text{-loop}}(\phi_0) \ll V'(\phi_0); \Delta V''_{1\text{-loop}}(\phi_0) \ll V''(\phi_0)$$

- Upper bound ($y\phi\bar{\chi}\chi$):

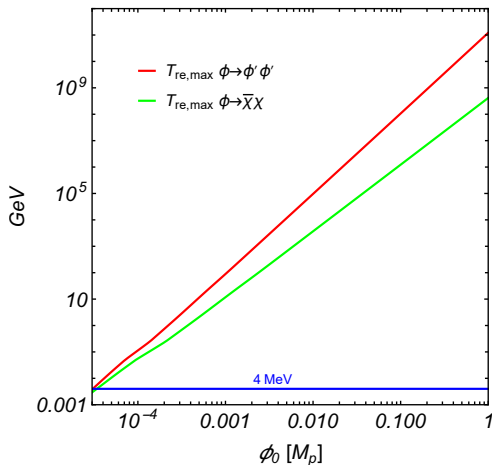


- Upper bound ($g\phi|\phi'|^2$):



- Radiative Stability + Reheating \Rightarrow Lower bound $\phi_0 > 3 \cdot 10^{-5} M_p$ [Drees, Xu '21]

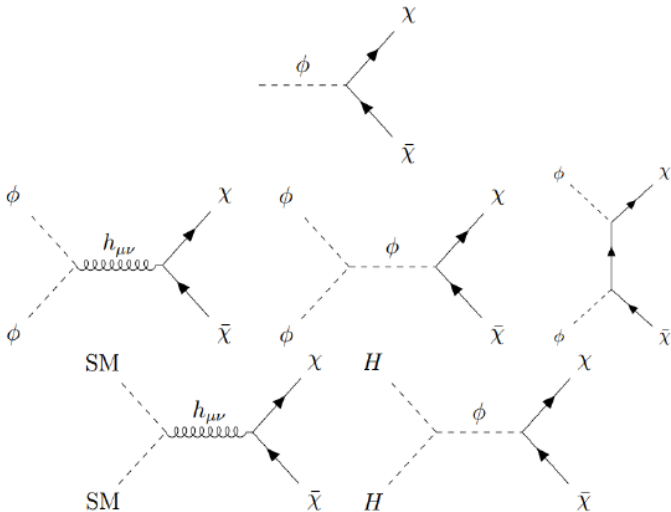
Reheating Temperature



- Bosonic: $4 \text{ MeV} \lesssim T_{\text{rh}} \lesssim 10^{11} \text{ GeV}$
- Fermionic: $4 \text{ MeV} \lesssim T_{\text{rh}} \lesssim 10^8 \text{ GeV}$

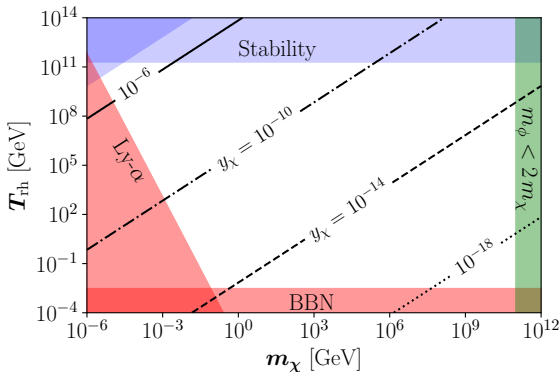
DM Production: After Polynomial Inflation

- Consider e.g. Fermionic DM $\mathcal{L}_\chi \supset y_\chi \phi \bar{\chi} \chi \Rightarrow 6$ possible channels:



Inflaton direct decay

- Parameter space: $y_\chi \simeq 1.2 \times 10^{-13} \sqrt{\frac{T_{\text{rh}}}{m_\chi}}$



- Bounds:
 - Stability: $T_{\text{rh}} < 1.2 \times 10^{11}$ GeV (Higgs loop), $y_\chi < 10^{-5}$ (DM loop)
 - BBN: $T_{\text{rh}} \gtrsim 4$ MeV
 - Ly α on cold DM: $v_\chi = \frac{p_0}{m_\chi} \lesssim 10^{-8} c \Leftrightarrow \frac{m_\chi}{\text{keV}} \gtrsim 2 \frac{m_\phi}{T_{\text{rh}}}$
- DM mass: $\mathcal{O}(10^{-5})$ GeV $\lesssim m_\chi \lesssim \mathcal{O}(10^{11})$ GeV [Bernal, Xu '21]

Baryogenesis

- A simple and attractive scenario: Leptogenesis [Fukugita, Yanagida 1986]

$$\mathcal{L}_N \supset - \left(\frac{1}{2} M_N \overline{N}_i^c N_i + h.c. \right) - \left(Y_{\alpha i} \bar{L}_\alpha \tilde{H} N_i + h.c. \right)$$

- Lower bound on M_1 [Davidson Ibarra '02, Buchmuller, Bari, Plumacher '04]

$$Y_B \sim \frac{28}{79} Y_{B-L} \sim \frac{28}{79} \epsilon_1 10^{-4} \gtrsim 10^{-10} \Rightarrow M_1 \gtrsim 10^{10} \text{ GeV}$$

with CP asymmetry parameter: $\epsilon_1 \lesssim 10^{-5} \left(\frac{M_1}{10^{11} \text{ GeV}} \right)$

- Recall $T_{\text{rh}} \lesssim 10^{11} \text{ GeV} \Rightarrow$ Thermal leptogenesis works (need high T_{rh})
- Question

Can one has leptogenesis with lower T_{rh} ?

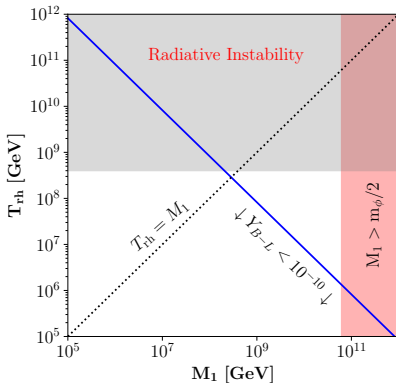
Non-thermal Leptogenesis [in preparation, Drees, Xu '22]

- Inflaton couples to RHN

$$\mathcal{L}_N \supset - (y_I \phi \overline{N}_I^c N_I + h.c.)$$

- Lepton yield

$$Y_{B-L} = \left(\frac{3}{2} \frac{T_{\text{rh}}}{m_\phi} \right) \cdot \epsilon_I$$



Summary

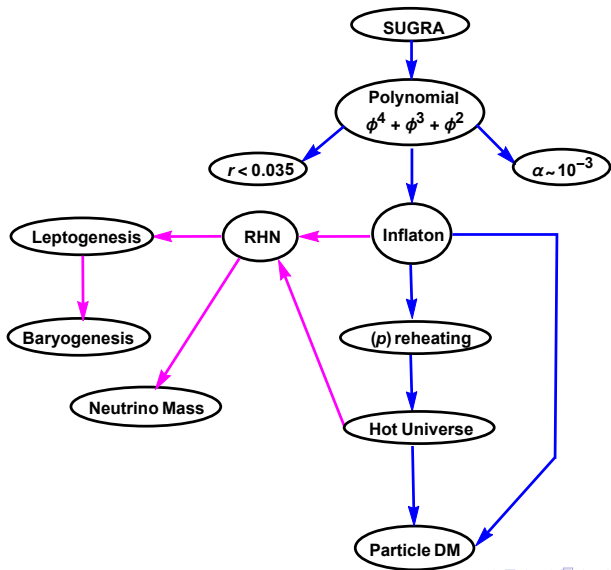
- A simple polynomial model fits data very well:

$$V \equiv d \left[\phi^4 + A(1 - \beta) \phi^3 + 9/32 A^2 \phi^2 \right]$$

with $A = -8/3\phi_0$; $\beta = 9.73 \times 10^{-7} \phi_0^4/M_p^4$; $d = 6.61 \times 10^{-16} \phi_0^2/M_p^2$.

- Parameter space: Reheating + Radiative Stability $\Rightarrow \phi_0 > 3 \cdot 10^{-5} M_p$
- Predictions:
 1. $r \simeq 7.1 \cdot 10^{-9} \phi_0^6/M_p^6$
 2. $\alpha \simeq -1.43 \cdot 10^{-3} \Rightarrow$ testable in future [S4 CMB]
- Implications:
 1. Inflationary scale: $H_{\text{inf}} \simeq 8.6 \cdot 10^{-9} \phi_0^3/M_p^2 \Rightarrow H_{\text{inf}}$ as low as 1 MeV!
 2. Reheating Tem: $T_{\text{re}} \in [4 \text{ MeV}, 10^{11} \text{ GeV}]$
- Dark Matter: $\mathcal{O}(10^{-5}) \text{ GeV} \lesssim m_\chi \lesssim \mathcal{O}(10^{11}) \text{ GeV}$
- Leptogenesis:
 1. thermal: $\mathcal{O}(10^{10}) \text{ GeV} \lesssim M_1 \lesssim \mathcal{O}(10^{11}) \text{ GeV}$
 2. non-thermal: $\mathcal{O}(10^8) \text{ GeV} \lesssim M_1 \lesssim \mathcal{O}(10^{10}) \text{ GeV}$

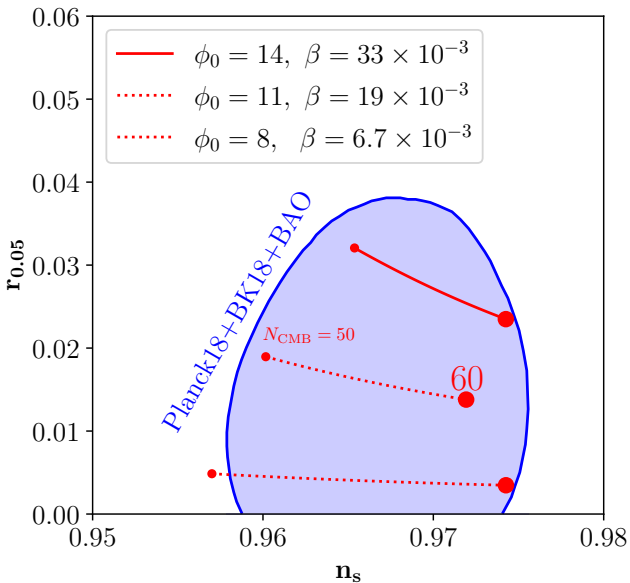
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Thank
you
for
your
attention
!

Backup: Large Field Scenario

[to appear soon, Drees, Xu '22]



Backup: Polynomial Inflation in Supergravity

- Scalar Potential

$$V = e^K \left[(D_i W) K_{i\bar{j}}^{-1} (D_{\bar{j}} \bar{W}) - 3|W|^2 \right],$$

with $D_i W = \frac{\partial W}{\partial \Phi_i} + \frac{\partial K}{\partial \Phi_i} W$ and $K_{i\bar{j}} = \frac{\partial^2 K}{\partial \Phi_i \partial \bar{\Phi}_j}$

- η problem see, e.g. [1101.2488] for review

$$V \sim (1 + |\Phi_i|^2) \left| \frac{\partial W}{\partial \Phi_i} \right|^2 = V^{\text{global}} + |\Phi_i|^2 V^{\text{global}} \Rightarrow \eta = V''/V \sim 1$$

- Consider e.g. a Superpotential and Kähler potential [Nakayama, Takahashi and Yanagida '13]:

$$W = X(\alpha_1 \Phi + \alpha_2 \Phi^2); K = \frac{1}{2}(\Phi + \Phi^\dagger)^2 + |X|^2$$

- Kähler potential admits a shift symmetry: $\Phi \rightarrow \Phi + iC$ [Kawasaki, Yamaguchi and Yanagida '00]
- $\phi \equiv \text{Im}(\Phi)$ not appear in $K \Rightarrow$ free from η problem
- Reproduce the (single field) polynomial inflaton potential:

$$v(\phi) = \left(1 + \frac{1}{2}(\phi^2 + \phi^{\dagger 2} + 2\phi\phi^\dagger) \right) (\alpha_1 \Phi + \alpha_2 \Phi^2) (\alpha_1^* \Phi^\dagger + \alpha_2^* \Phi^{\dagger 2}) \supset \left(\frac{|\alpha_1|^2}{2} \phi^2 - \frac{\sqrt{2}|\alpha_1||\alpha_2|\sin\theta}{2} \phi^3 + \frac{|\alpha_2|^2}{4} \phi^4 \right)$$