

# Flipping the Universe

bridging the  
gap between  
String Theory and  
Particle Physics?

*John Ellis*

**KING'S**  
College  
LONDON

*IgnatiosFest*  
*Planck 2022*

# Early Papers with Ignatios

35 years already!

Striving to link string  
theory with accessible  
particle  
physics:

Model-building,  
time-dependent  
cosmological solutions

<b>The Low-energy Effective Field Theory From Four-dimensional Superstrings</b> Ignatios Antoniadis (CERN and Crete U.), John R. Ellis (CERN), E. Floratos (Crete U.), Dimitri V. Nanopoulos (Wisconsin U., Madison), T. Tomaras (Crete U.) (Feb, 1987) Published in: <i>Phys.Lett.B</i> 191 (1987) 96-102 DOI cite 89 citatio
<b>On the Possibility of Avoiding Singularities by Dilaton Emission</b> Ignatios Antoniadis (CERN), G.F.R. Ellis (CERN), John R. Ellis (CERN), C. Kounnas (LBL, Berkeley), Dimitri V. Nanopoulos (Wisconsin U., Madison) (Feb, 1987) Published in: <i>Phys.Lett.B</i> 191 (1987) 393-398 DOI cite 9 citatio
<b>Supersymmetric Flipped SU(5) Revitalized</b> Ignatios Antoniadis (CERN), John R. Ellis (CERN), J.S. Hagelin (Maharishi U. of Management), Dimitri V. Nanopoulos (Wisconsin U., Madison) (May, 1987) Published in: <i>Phys.Lett.B</i> 194 (1987) 231-235 DOI cite 584 citatio
<b>Universality of the Mass Spectrum in Closed String Models</b> Ignatios Antoniadis (CERN), John R. Ellis (CERN), Dimitri V. Nanopoulos (Wisconsin U., Madison) (Aug, 1987) Published in: <i>Phys.Lett.B</i> 199 (1987) 402-406 DOI cite 50 citatio
<b>GUT Model Building with Fermionic Four-Dimensional Strings</b> Ignatios Antoniadis (CERN), John R. Ellis (CERN and SLAC), J.S. Hagelin (Maharishi U. of Management), Dimitri V. Nanopoulos (Wisconsin U., Madison) (Dec, 1987) Published in: <i>Phys.Lett.B</i> 205 (1988) 459-465 pdf links DOI cite 218 citatio
<b>An Improved SU(5) x U(1) Model from Four-Dimensional String</b> Ignatios Antoniadis (CERN), John R. Ellis (CERN), John S. Hagelin (Maharishi U. of Management), Dimitri V. Nanopoulos (Wisconsin U., Madison) (Mar, 1988) Published in: <i>Phys.Lett.B</i> 208 (1988) 209-215, <i>Phys.Lett.B</i> 213 (1988) 562 (addendum) DOI cite 232 citatio
<b>Cosmological String Theories and Discrete Inflation</b> Ignatios Antoniadis (CERN), C. Bachas (CERN), John R. Ellis (CERN), Dimitri V. Nanopoulos (Wisconsin U., Madison) (May, 1988) Published in: <i>Phys.Lett.B</i> 211 (1988) 393-399 DOI cite 311 citatio
<b>An Expanding Universe in String Theory</b> Ignatios Antoniadis (Ecole Polytechnique), C. Bachas (Ecole Polytechnique), John R. Ellis (CERN), Dimitri V. Nanopoulos (Texas A-M) (Nov, 1988) Published in: <i>Nucl.Phys.B</i> 328 (1989) 117-139 DOI cite 319 citatio
<b>The Flipped SU(5) x U(1) String Model Revamped</b> Ignatios Antoniadis (CERN), John R. Ellis (CERN), J.S. Hagelin (CERN), Dimitri V. Nanopoulos (CERN) (Jul, 1989) Published in: <i>Phys.Lett.B</i> 231 (1989) 65-74 DOI cite 565 citatio

# Generalised No-Scale Structure

Volume 191, number 1,2

PHYSICS LETTERS B

4 June 1987

## THE LOW-ENERGY EFFECTIVE FIELD THEORY FROM FOUR-DIMENSIONAL SUPERSTRINGS

I. ANTONIADIS<sup>1,2</sup>, John ELLIS, E. FLORATOS<sup>1</sup>, D.V. NANOPOULOS<sup>3</sup> and T. TOMARAS<sup>1</sup>  
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Received 25 February 1987

We derive the low-energy effective supergravity field theory obtained from a fermionic formulation of four-dimensional superstrings. It contains a single dilaton supermultiplet parametrizing an  $SU(1, 1)/U(1)$  manifold, three self-conjugate sets of matter supermultiplets parametrizing  $SO(m, 2)/SO(m) \times SO(2)$  manifolds, and chiral matter supermultiplets parametrizing a product of  $SU(M, 1)/SU(M) \times U(1)$  manifolds. We derive the Yukawa couplings of the theory and identify some of the self-conjugate matter fields as Higgses. We also discuss the possible pattern of supersymmetry breaking.

Crucial for cosmological model-building  
(no Planck-scale holes in the effective potential):  
Basis for supersymmetric models of inflation (later)

# Our Biggest Hit!

Volume 194, number 2

PHYSICS LETTERS B

6 August 1987

## **SUPERSYMMETRIC FLIPPED SU(5) REVITALIZED**

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and

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Received 16 May 1987

Resurrection of flipped SU(5):  
motivated by possibility of  
derivation in weakly-coupled  
heterotic string theory,  
because no need of adjoint (or  
larger) Higgs representation

We describe a simple  $N=1$  supersymmetric GUT based on the group  $SU(5) \times U(1)$  which has the following virtues: the gauge group is broken down to the  $SU(3)_C \times SU(2)_L \times U(1)_Y$  of the standard model using just **10**, **10** Higgs representations, and the doublet-triplet mass splitting problem is solved naturally by a very simple missing-partner mechanism. The successful supersymmetric GUT prediction for  $\sin^2 \theta_w$  can be maintained, whilst there are no fermion mass relations. The gauge group and representation structure of the model may be obtainable from the superstring.

# Flipped SU(5) GUT

Antoniadis, JE, Hagelin & Nanopoulos, 1987

- Extend GUT SU(5) with additional U(1) [motivated by string theory]
- “Flipped” fermion assignments to representations:

$$\bar{f}_i(\bar{\mathbf{5}}, -3) = \{U_i^c, L_i\} , \quad F_i(\mathbf{10}, 1) = \{Q_i, D_i^c, N_i^c\} , \quad l_i(\mathbf{1}, 5) = E_i^c , \quad i = 1, 2, 3$$

- Break GUT symmetry with 10-dimensional Higgses, electroweak symmetry with 5-dimensional Higgses:

$$H(\mathbf{10}, 1) = \{Q_H, D_H^c, N_H^c\} , \quad \bar{H}(\bar{\mathbf{10}}, -1) = \{\bar{Q}_H, \bar{D}_H^c, \bar{N}_H^c\}$$

$$h(\mathbf{5}, -2) = \{T_{H_c}, H_d\} , \quad \bar{h}(\bar{\mathbf{5}}, 2) = \{\bar{T}_{\bar{H}_c}, H_u\}$$

- Superpotential:

$$W = \lambda_1^{ij} F_i F_j h + \lambda_2^{ij} F_i \bar{f}_j \bar{h} + \lambda_3^{ij} \bar{f}_i \ell_j^c h + \lambda_4 H H h + \lambda_5 \bar{H} \bar{H} \bar{h} \\ + \lambda_6^{ia} F_i \bar{H} \phi_a + \lambda_7^a h \bar{h} \phi_a + \lambda_8^{abc} \phi_a \phi_b \phi_c + \mu_\phi^{ab} \phi_a \phi_b ,$$

# First Version of Stringy Flipped SU(5)

Volume 205, number 4

PHYSICS LETTERS B

5 May 1988

## GUT MODEL-BUILDING WITH FERMIONIC FOUR-DIMENSIONAL STRINGS <sup>★</sup>

I. ANTONIADIS <sup>a,1</sup>, John ELLIS <sup>a,b</sup>, J.S. HAGELIN <sup>c</sup> and D.V. NANOPOULOS <sup>d</sup>

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Received 7 January 1988

We report a first attempt at model-building using the fermionic formulation of string theories directly in four dimensions. An example is presented of a supersymmetric flipped  $SU(5) \times U(1)$  model with three generations and an adjustable hidden sector gauge group. The simplest version of the model contains most of the Yukawa couplings required by phenomenology, but not all those needed to give masses to quarks or conjugate neutrinos. These defects may be remedied in a more general version of the model.

First candidate unified “Theory of Everything”:  
derived using fermionic formulation of weakly-coupled heterotic string  
(Antoniadis, Bachas and Kounnas<sup>†</sup>)

# Refined Version of Stringy Flipped SU(5)

Volume 231, number 1,2

PHYSICS LETTERS B

2 November 1989

## THE FLIPPED $SU(5) \times U(1)$ STRING MODEL REVAMPED

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Received 21 July 1989

- We present a refined version of our three-generation flipped  $SU(5) \times U(1)$  string model with the following properties.
- The complete massless spectrum is derived and shown to be free of all gauge and mixed anomalies apart from a single anomalous  $U(1)$ .
  - The imaginary part of the dilaton supermultiplet is eaten by the anomalous  $U(1)$  gauge boson, and the corresponding D-term is cancelled by large VEVs for singlet fields that break surplus  $U(1)$  gauge factors, leaving a supersymmetric vacuum with an  $SU(5) \times U(1)$  visible gauge group and an  $SO(10) \times SO(6)$  hidden gauge group.
  - There are sufficient Higgs multiplets to break the visible gauge symmetry down to the standard model in an essentially unique way.
  - All trilinear superpotential couplings have been calculated and there are in particular some giving  $m_t, m_b, m_\tau \neq 0$ .
  - A renormalization group analysis shows that  $m_t < 190$  GeV and  $m_b \simeq 3m_\tau$ .
  - Light Higgs doublets are split automatically from heavy Higgs triplets, leaving no residual dimension-five operators for baryon decay, and the baryon lifetime  $\tau_B \sim 2 \times 10^{34 \pm 2}$  yr.
  - There are no tree-level flavour-changing neutral currents, but  $\mu \rightarrow e\gamma$  may occur at a detectable level:  $B(\mu \rightarrow e\gamma) \sim 10^{-11} - 10^{-14}$ .

Our second-biggest hit!

# Meanwhile Back in the Universe

## AN EXPANDING UNIVERSE IN STRING THEORY

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Received 14 April 1989

Time-dependent  
model of  
cosmology based  
on non-critical  
string theory

We present solutions of the bosonic, heterotic and type II string theories, whose space-time manifold is a linearly expanding homogeneous and isotropic universe. These solutions are obtained by giving a background charge to the time coordinate on the world-sheet. We find the spectrum, demonstrate positivity of the Hilbert space up to the second excited level, and construct modular-invariant partition functions. The central charge of transverse excitations is a free parameter that controls the asymptotic density of states; the critical dimension and gauge group can in particular be made arbitrarily large. We show how to construct dual, factorizable, energy-conserving amplitudes in this background, discuss their interpretation and comment on the initial singularity and flatness problems in the light of our results.



# Not Forgetting ...!

## LEP data and the light gluino window

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Received 28 March 1991

## String threshold corrections and flipped SU(5)

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Received 8 July 1991

## Unification bounds on the possible $N = 2$ supersymmetry-breaking scale <sup>★</sup>

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Received 17 January 1997

Editor: R. Gatto

Key issues in the phenomenology of strings and supersymmetry

# Daedalus or Icarus?



DEDALE PERD SON FILS ICARE

Flipped

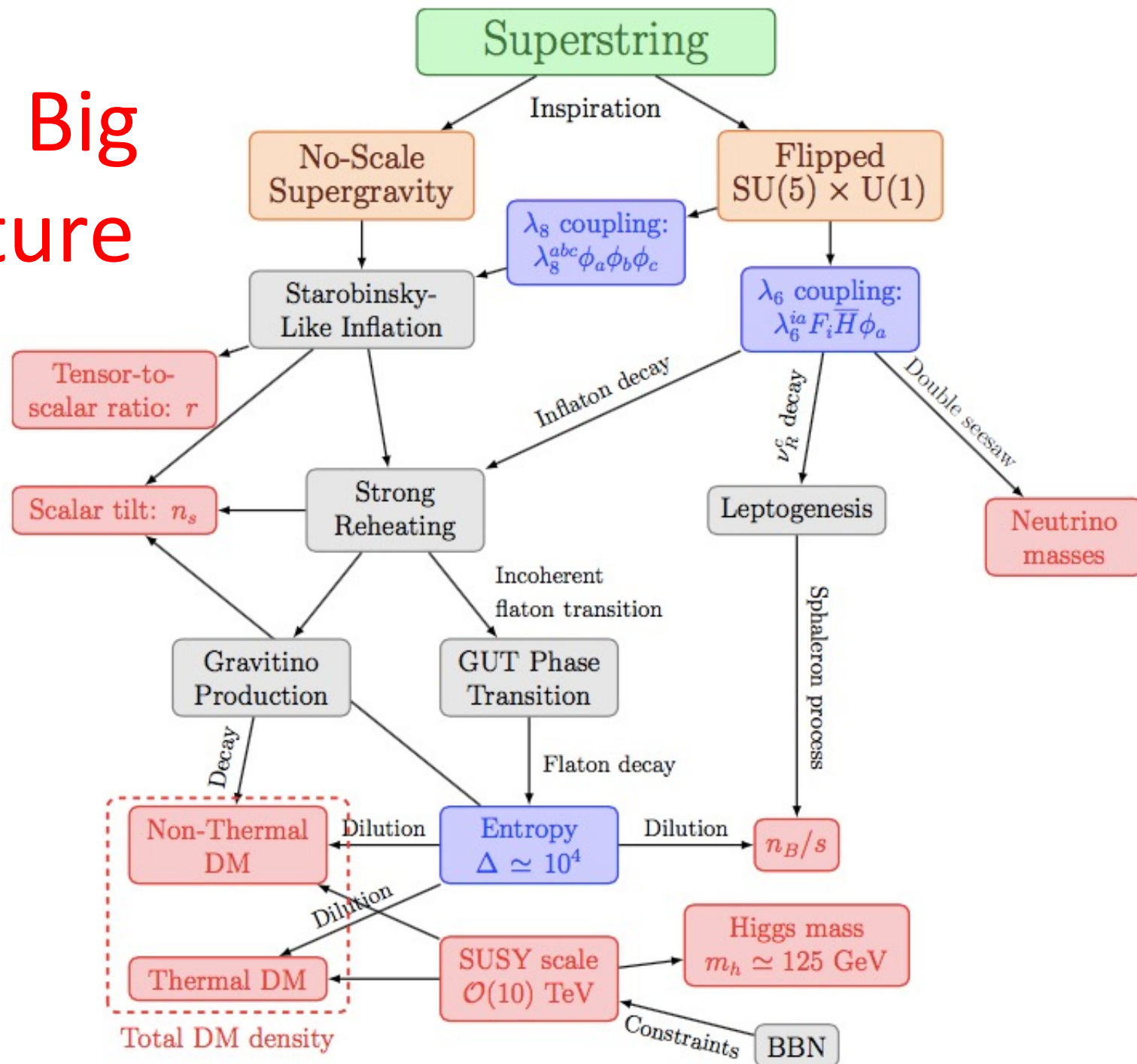
Almost

# A Model of Everything

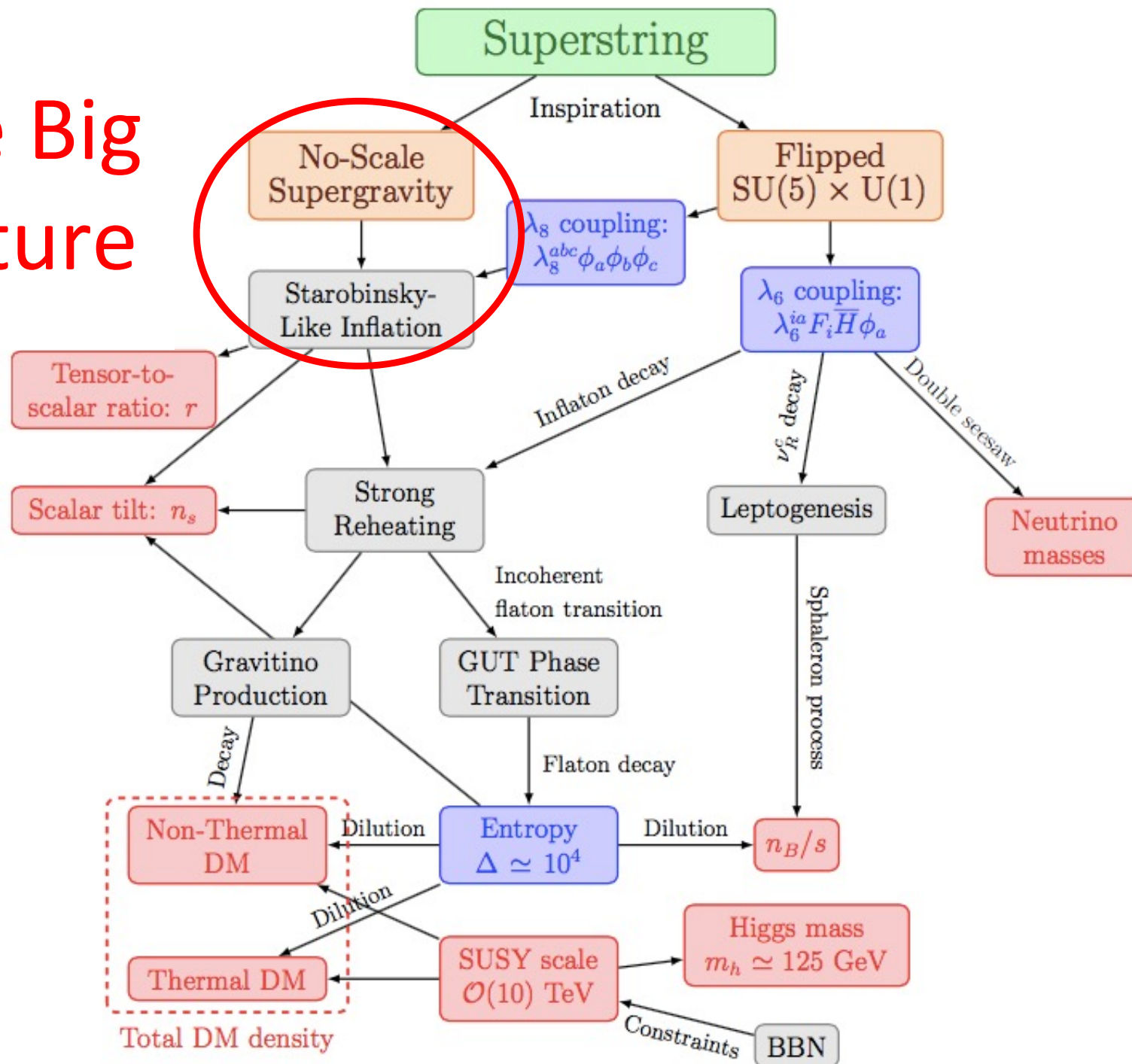
Below the Planck Scale

- Simple GUT models ( $SU(5)$ ,  $SO(10)$ ) not obtained from weakly-coupled string
  - They need adjoint Higgs, ...
- **Flipped  $SU(5) \times U(1)$  derived**, has advantages
  - Small (5-, 10-dimensional) Higgs representations
  - Long-lived proton, neutrino masses, leptogenesis, ...
- Construct model of Starobinsky-like inflation within flipped  $SU(5) \times U(1)$  framework

# The Big Picture



# The Big Picture



# Inflation Cries out for Supersymmetry

- Want “elementary” scalar field  
(at least looks elementary at energies  $\ll M_p$ )
- To get right magnitude of perturbations  
prefer mass  $\ll M_p$   
( $\sim 10^{13}$  GeV in simple  $\phi^2$  models)
- And/or prefer small self-coupling  $\lambda \ll 1$
- **Both technically natural with supersymmetry**

# Inflation cries out for Supergravity

- Stabilize 'elementary' scalar inflaton  
(needs mass  $\ll m_p$  and/or small coupling)
- **Supersymmetry**
- The only good symmetry is a local symmetry  
(cf, gauge symmetry in Standard Model)
- **Local supersymmetry = supergravity**
- Early Universe cosmology needs gravity
- **Supersymmetry + gravity = supergravity**

# No-Scale Supergravity Inflation

- **Supersymmetry + gravity = Supergravity**
- Include conventional matter?
- Potentials in generic supergravity models have ‘holes’ with depths  $\sim -M_p^4$
- Exception: **no-scale supergravity**
- **Appears in compactifications of string**
- Flat directions, scalar potential  $\sim$  global model + controlled corrections

Cremmer, Ferrara, Kounnas  
& Nanopoulos, 1983

Witten, 1985

JE, Enqvist, Nanopoulos, Olive & Srednicki, 1984

JE, Nanopoulos & Olive, arXiv:1305.1247, 1307.3537, ...



# Old No-Scale Supergravity Model of Inflation

Volume 152B, number 3,4

PHYSICS LETTERS

7 March 1985

## SU( $N$ , 1) INFLATION

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Received 7 December 1984

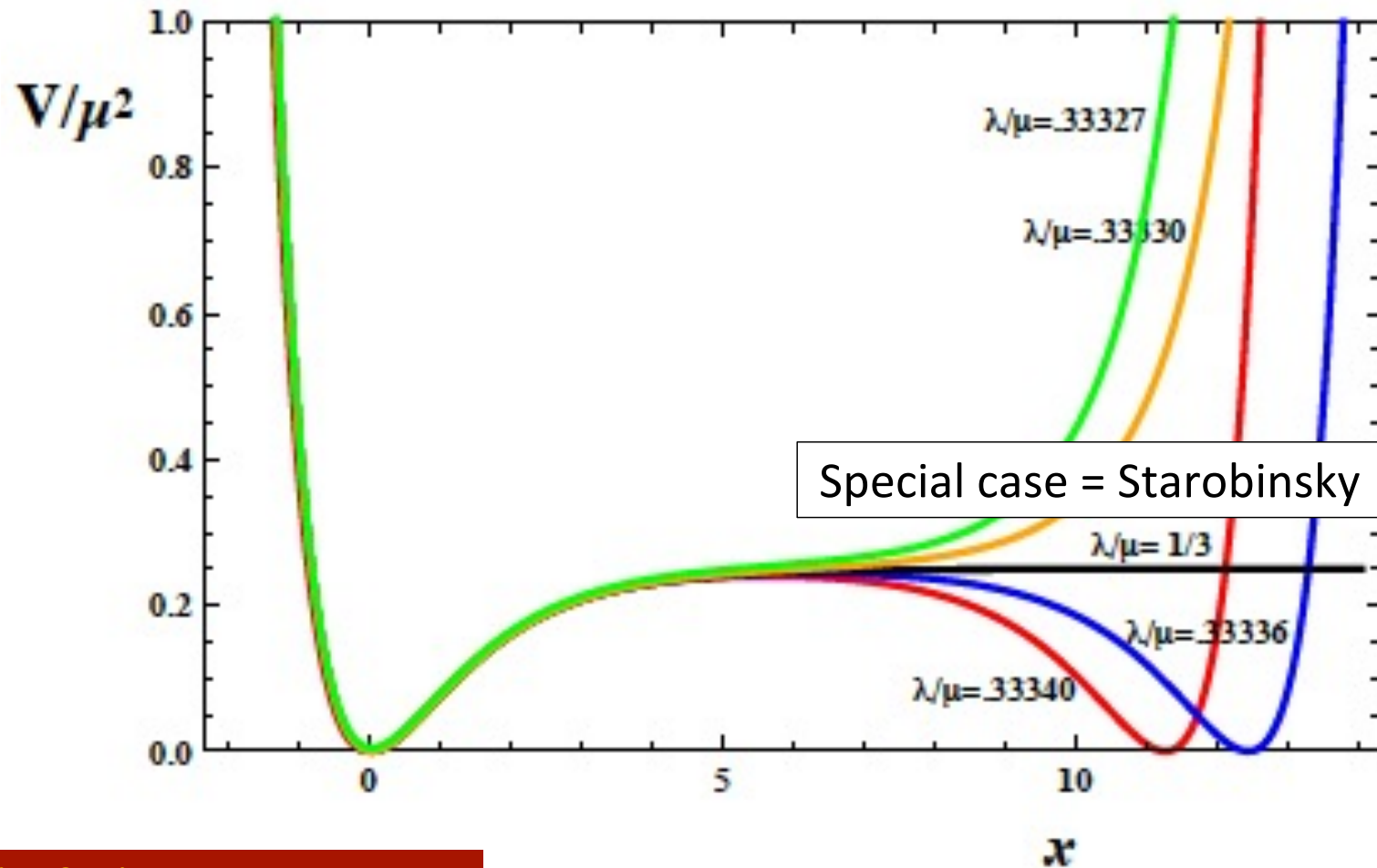
No 'holes' in  
effective potential  
with negative  
cosmological  
constant

JE, Enqvist, Nanopoulos, Olive & Srednicki, 1984

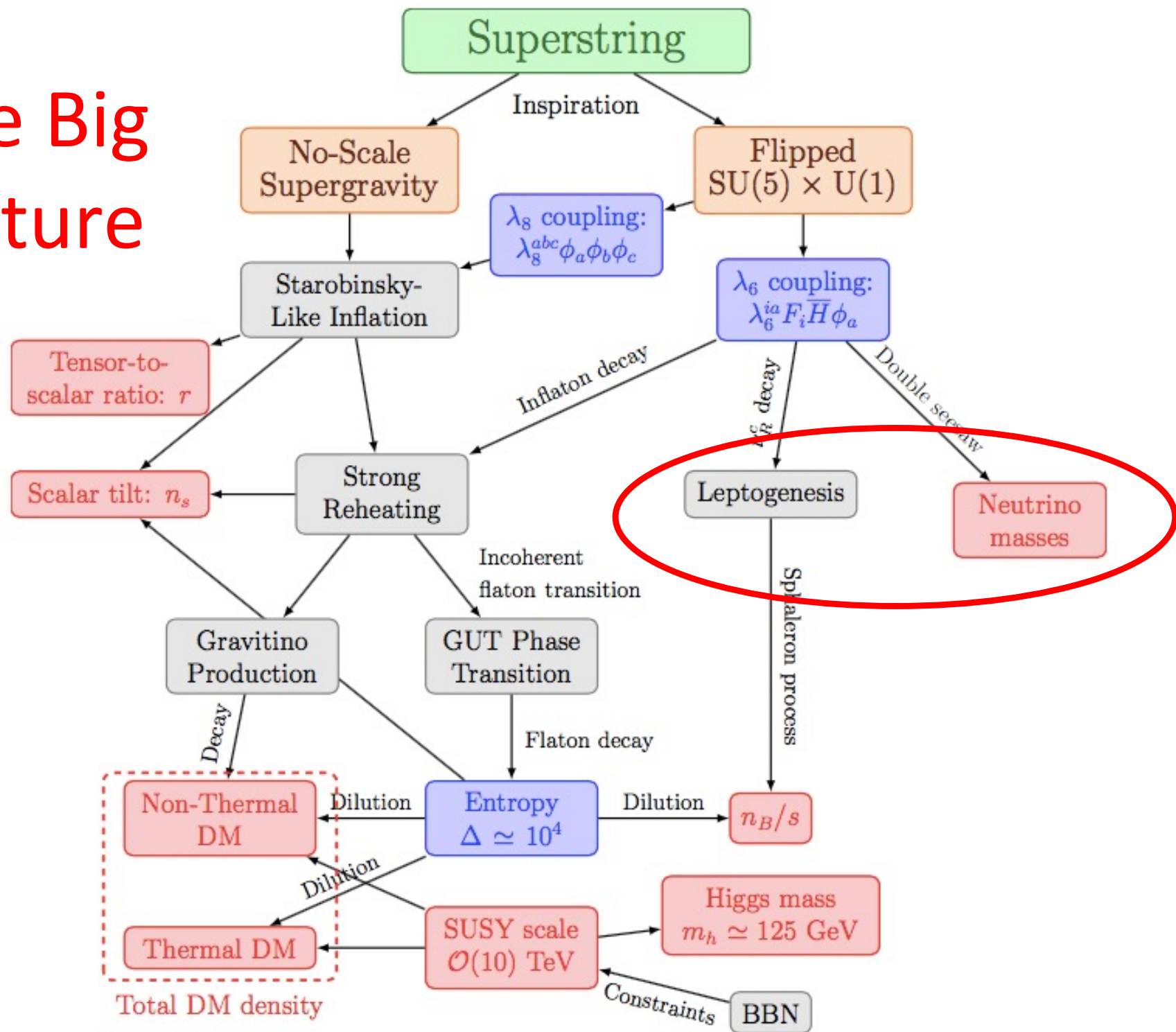
We present a simple model for primordial inflation in the context of SU( $N$ , 1) no-scale  $n = 1$  supergravity. Because the model at zero temperature very closely resembles global supersymmetry, minima with negative cosmological constants do not exist, and it is easy to have a long inflationary epoch while keeping density perturbations of the right magnitude and satisfying other cosmological constraints. We pay specific attention to satisfying the thermal constraint for inflation, i.e. the existence of a high temperature minimum at the origin.

# No-Scale Supergravity Inflation

- Inflationary potential for  $\lambda \simeq \mu/3$



# The Big Picture



# Neutrino Masses & Mixing

- Neutrinos couple to singlet scalars:

$$W = \sum_{i=1}^3 \lambda_2^i \nu_i^c L_i H_d + \sum_{a=0}^2 \mu^a \phi_a^2 + \sum_{i,a} \lambda_6^{ia} \nu_i^c \nu_{\bar{H}}^c \tilde{\phi}_a$$

- Neutrino mass matrix:

$\lambda_6$  key parameter in flipped SU(5)

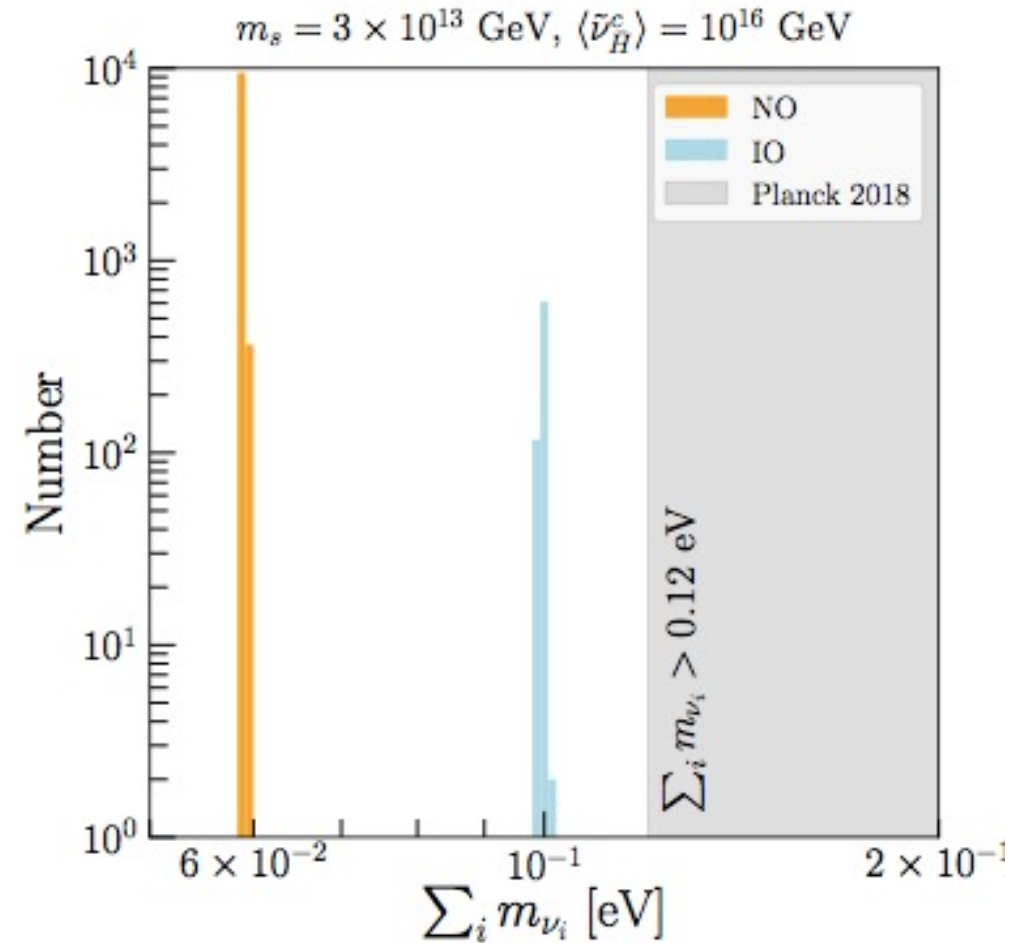
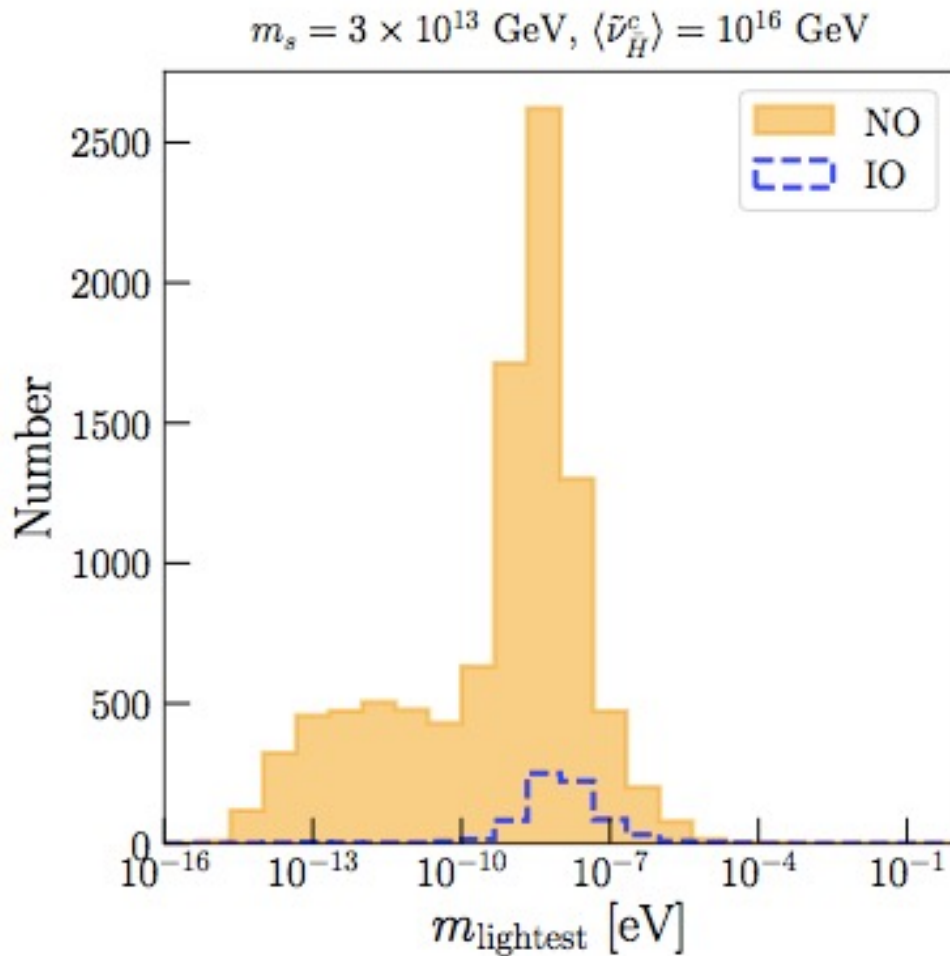
$$\mathcal{L}_{\text{mass}} = -\frac{1}{2} \begin{pmatrix} \nu_i & \nu_j^c & \tilde{\phi}_a \end{pmatrix} \begin{pmatrix} 0 & \lambda_2^{ij} \langle \bar{h}_0 \rangle & 0 \\ \lambda_2^{ij} \langle \bar{h}_0 \rangle & 0 & \lambda_6^{ja} \langle \tilde{\nu}_{\bar{H}}^c \rangle \\ 0 & \lambda_6^{ja} \langle \tilde{\nu}_{\bar{H}}^c \rangle & \mu^a \end{pmatrix} \begin{pmatrix} \nu_i \\ \nu_j^c \\ \tilde{\phi}_a \end{pmatrix} + \text{h.c.}$$

- Double seesaw for heavy & light neutrinos:

$$(m_{\nu^c})_{ij} = \sum_{a=0,1,2} \frac{\lambda_6^{ia} \lambda_6^{ja}}{\mu^a} \langle \tilde{\nu}_{\bar{H}}^c \rangle^2 \quad (m_{\nu})_{ij} = \sum_k \frac{\lambda_2^i \lambda_2^j (U_{\nu^c})_{ik} (U_{\nu^c})_{jk} \langle \bar{h}_0 \rangle^2}{(m_{\nu^c}^D)_k}$$

# Light Neutrino Masses

Normal ordering vs inverse ordering



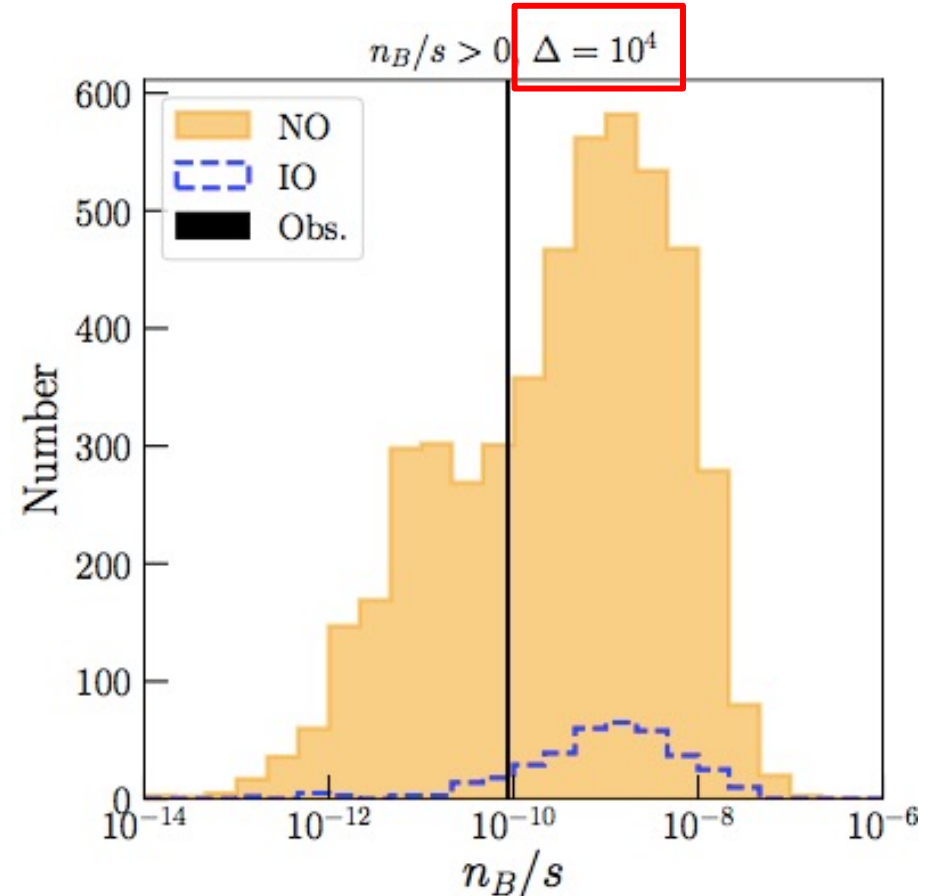
# Cosmological Baryon Density

- Generated by CP violation in heavy neutrino decay

$$\frac{n_B}{s} = -\frac{28}{79} \cdot \frac{135\zeta(3)}{4\pi^4 g_{\text{reh}} \Delta} \sum_{i=1,2,3} \epsilon_i,$$

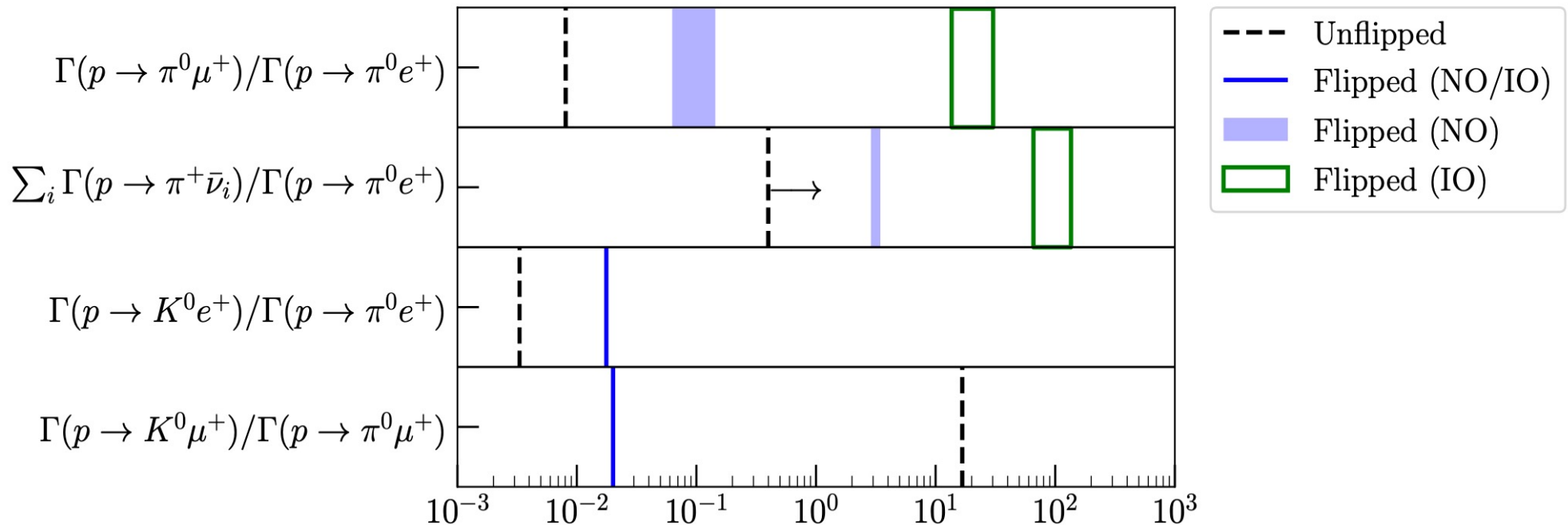
$$\epsilon_i = \frac{1}{2\pi} \frac{\sum_{j \neq i} \text{Im} \left[ \left( U_{\nu^c}^\dagger (\lambda_2^D)^2 U_{\nu^c} \right)_{ji}^2 \right]}{\left[ U_{\nu^c}^\dagger (\lambda_2^D)^2 U_{\nu^c} \right]_{ii}} g \left( \frac{m_{\nu_j^c}^2}{m_{\nu_i^c}^2} \right)$$

$$g(x) \equiv -\sqrt{x} \left[ \frac{2}{x-1} + \ln \left( \frac{1+x}{x} \right) \right].$$



Correct amount with appropriate entropy gain  $\Delta \geq 10^4$

# B Decay in Flipped SU(5) vs Unflipped



Probing different decay modes can distinguish between different models

# Hyper-Kamiokande Experiment

## Water Čerenkov detector

Being built to measure CP violation in neutrino oscillations

Access tunnel and cavern

Tank  
(Liner and Support structure for photo-detection system)

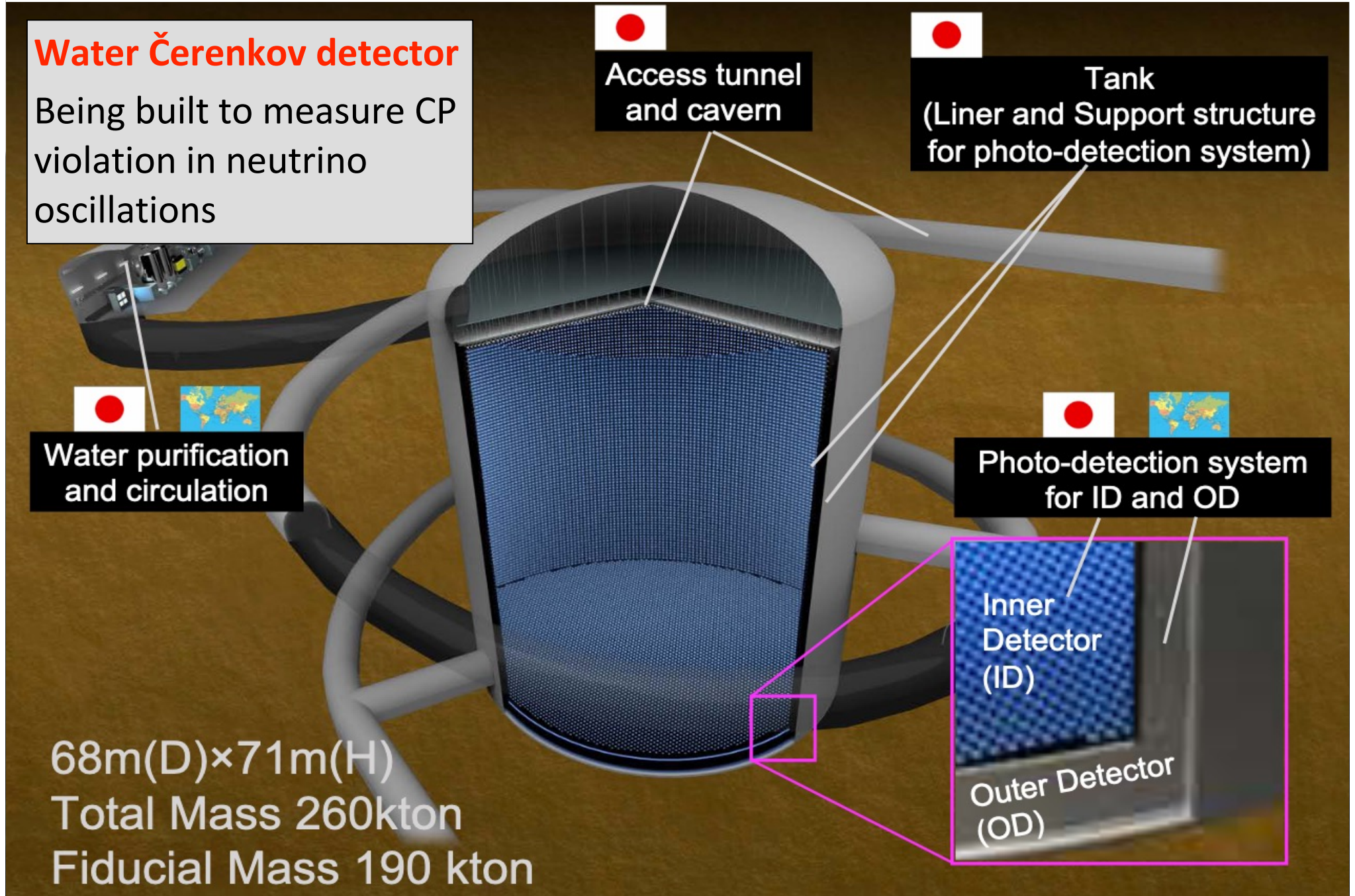
Water purification and circulation

Photo-detection system for ID and OD

Inner Detector (ID)

Outer Detector (OD)

68m(D)×71m(H)  
Total Mass 260kton  
Fiducial Mass 190 kton



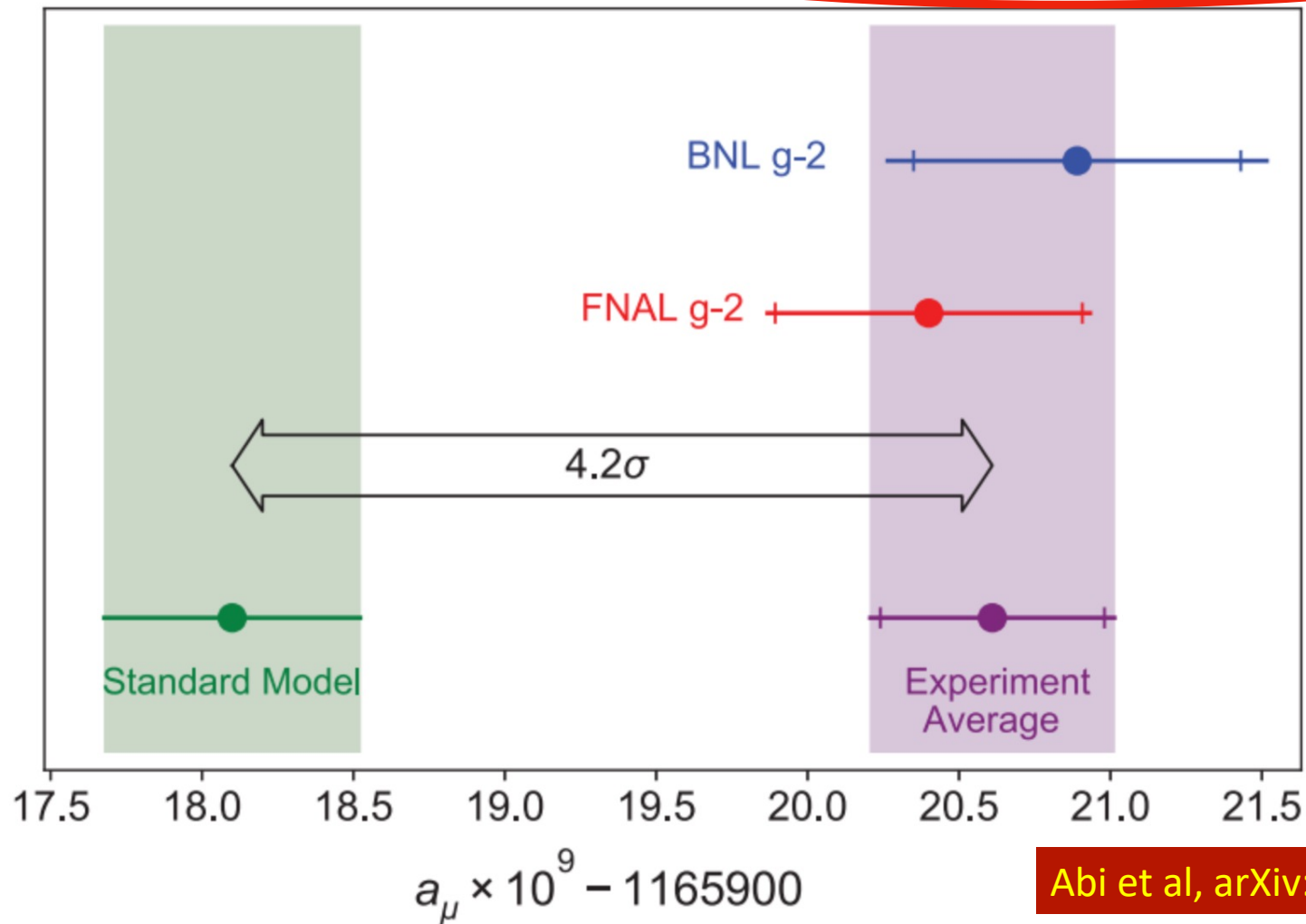


# Fermilab Measurement of $g_\mu - 2$

FNAL result:  $a_\mu(\text{FNAL}) = 116\,592\,040(54) \times 10^{-11}$  (0.46 ppm)

Combined result:  $a_\mu(\text{Exp}) = 116\,592\,061(41) \times 10^{-11}$  (0.35 ppm)

Difference from Standard Model:  $a_\mu(\text{Exp}) - a_\mu(\text{SM}) = (251 \pm 59) \times 10^{-11}$

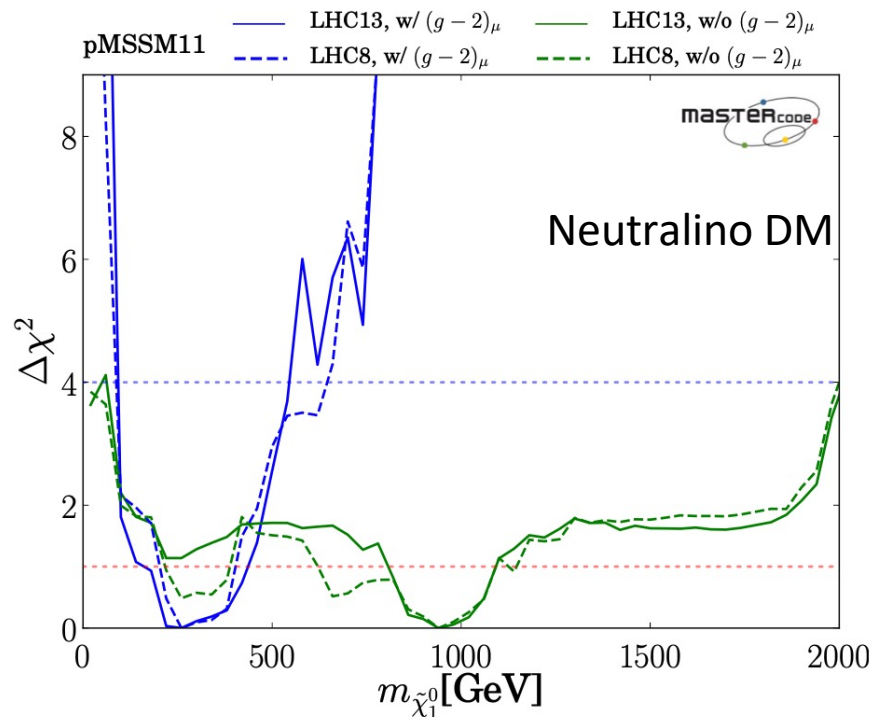


# Interpretation Papers

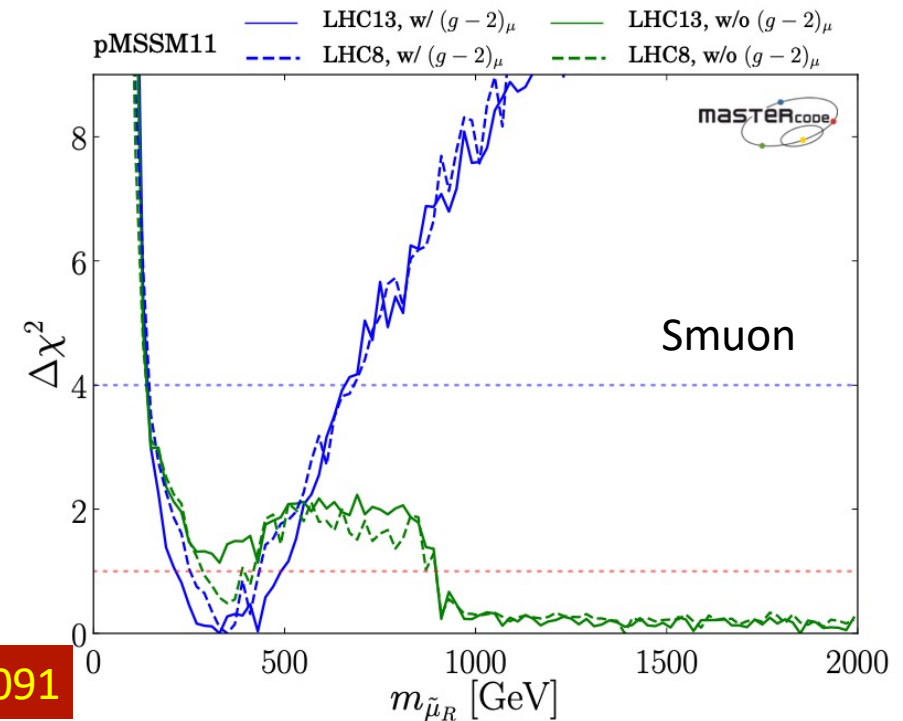
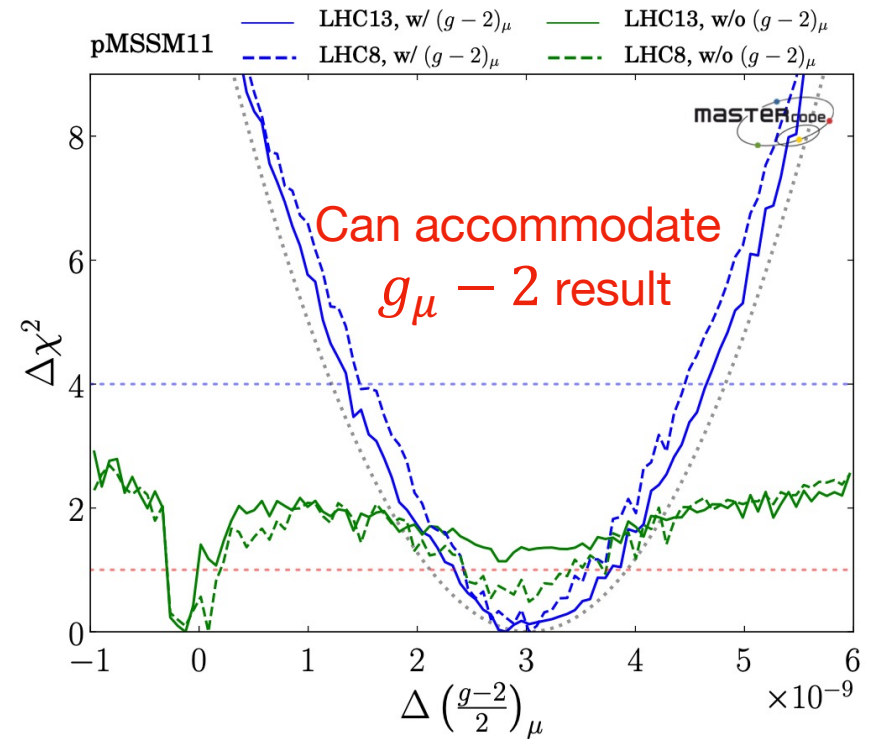
2104.05685	Vector LQ	B	Du				
5656	$L_{\mu} - L_{\tau}$	DM	Borah				
5006	$B_q - L_{\mu}$	B	Cen		Leptoquarks		
4494	LFV	LFV	Li				
4503	Pseudoscalar	DM, H decays	Lu		Extra U(1)		
4456	2HDM	DM	Arcadi				
3542	B-LSSM	H decays	Yang		Extra Higgs		
3701	Leptophilic spin 0	H factory	Chun				
3839	SUSY	HL-LHC	Aboubrahim		Supersymmetry		
3691	Survey	DM, LHC	Athron				
3705	Seesaw	$g_e$	Escribano		Axion		
3699	Gauged 2HDM	B	Chen				
3239	SUSY	Gravitino DM	Gu				
3284	NMSSM	DM	Cao				
3262	GUT-constrained SUSY	DM, LHC	Wang				
3292	MSSM	CPV	Han				
3296	lepton mass matrix	Flavour	Calibbi				
3280	$Z_d$	Cs weak charge	Cadeddu				
3334	$E_6$ 3-3-1	H stability	Li				
3242	$\mu$ - $\tau$ -philic H	$\tau$ decays, LHC	Wang				
3259	Anomaly mediation	DM	Yin				
3245	pMSSM	DM, fine-tuning	Van Beekveld				
3274	NMSSM	DM, AMS-02 $p\bar{p}$	Abdughani				
3290	MSSM	DM	Cox				
3367	2HDM	V-like leptons	Ferreira				
3267	Axion	Low-scale	Buen-Abad				
3340	$L_{\mu} - L_{\tau}$	AMS-02 positrons	Zu				
3282	ALP	V-like fermions	Brdar				
3301	Lepton portal	DM	Bai				
3276	Dark axion portal	Dark photon	Ge				
3491	GmSUGRA	LHC	Ahmed				
3227	2HDM	LHC	Han				
3302	SUSY	small $\mu$	Baum				
3238	Scalar	DM, $p$ radius	Zhu				
3489	$\mu$ $\nu$ SSM	B, H decays	Zhang				
3287	pMSSM	ILC	Chakraborti				
3228	DM	B, H decays	Arcadi				
890	Radiative seesaw						Chiang
2103.13991	Scalar LQ	B, H decays					Greljo
2012.11766	DM						D'Agnolo
2012.07894	Axions						Darmé
1812.06851	Charmphilic LQ						Kowalska
2104.04458	GUT-constrained SUSY	DM					Chakraborti
5730	LQ + charged singlet	B, Cabibbo					Marzocca
6320	L-R symmetry						Boyarkin
6858	$L_{\mu} - L_{\tau}$	$\nu$ masses					Zhou
6854	D-brane	U(1), Regge					Anchordoqui
6656	vector LQ	B					Ban
7597	SUSY	LHC, landscape					Baer
7047	3HDM	Fermion masses					Carcamo
7680	Leptophilic $Z'$	Global analysis					Buras
8289	Custodial symmetry	Light scalar + pseudoscalar					Balkin
9205	U(1)D	Neutrino mass					Dasgupta
8819	Lepton non-universality	Naturalness					Cacciapaglia
8640	$2 \times 2 \times 1$	Higgses, heavy $\nu$ s					Boyarkina
8293	Multi-TeV sleptons in FSSM	Extended H, $\tau$ decays					Altmannshofer
10114	SO(10)	Yukawa unification					Aboubrahim
7681	U(1)B-L	DUNE					Dev
10324	Gauged lepton number	Dark matter					Ma
10175	2HDM	Lighter Higgs?					Jueid
11229	LQ	Matter unification					Fileviez
15136	U(1)	HE neutrinos, H tension					Alonso
2105.00903	Anomalous 3-boson vertex	W mass					Arbuzov
7655	U(1)T3R	RK(*)					Dutta
8670	Leptoquark	$\nu$ mass, LFV					Zhang

# $g_\mu - 2$ in Phenomenological Supersymmetry (pMSSM11)

No relation between squark/gluino masses  
and slepton/neutralino masses

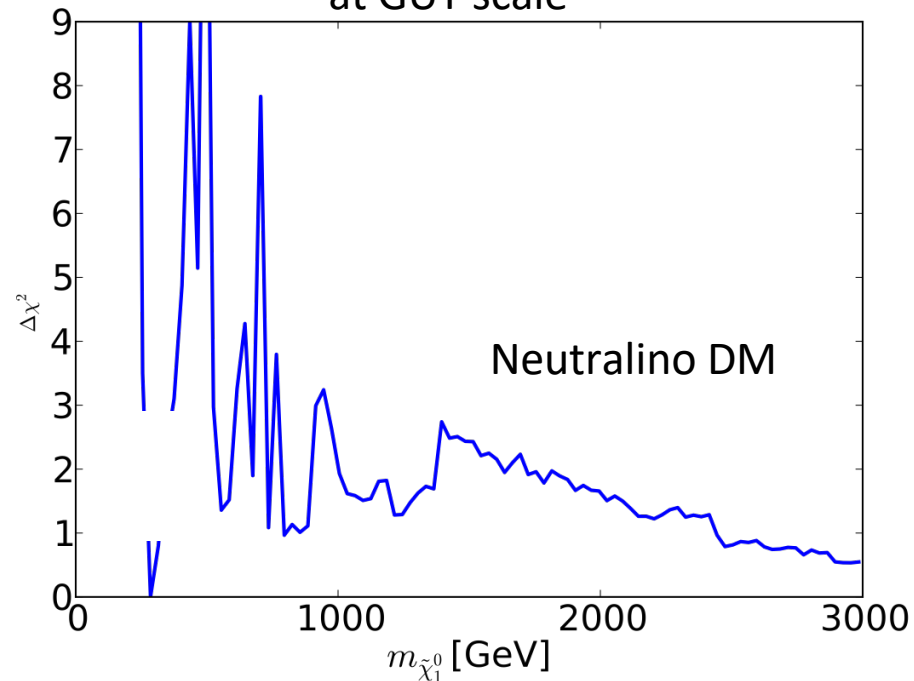


No problem accommodating BNL/FNAL result  
Neutralino DM, smuon masses  $\sim 300/400$  GeV

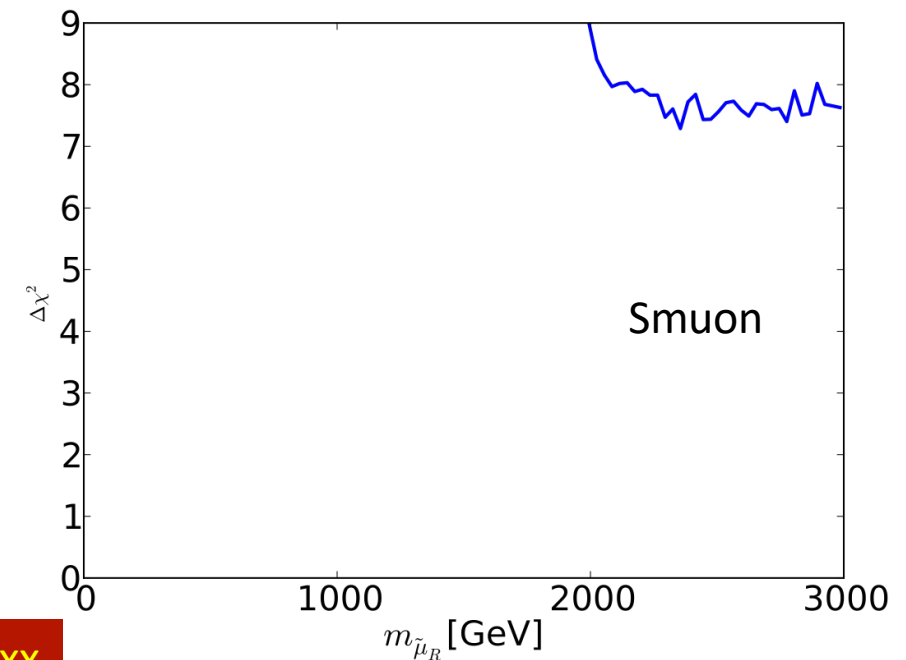
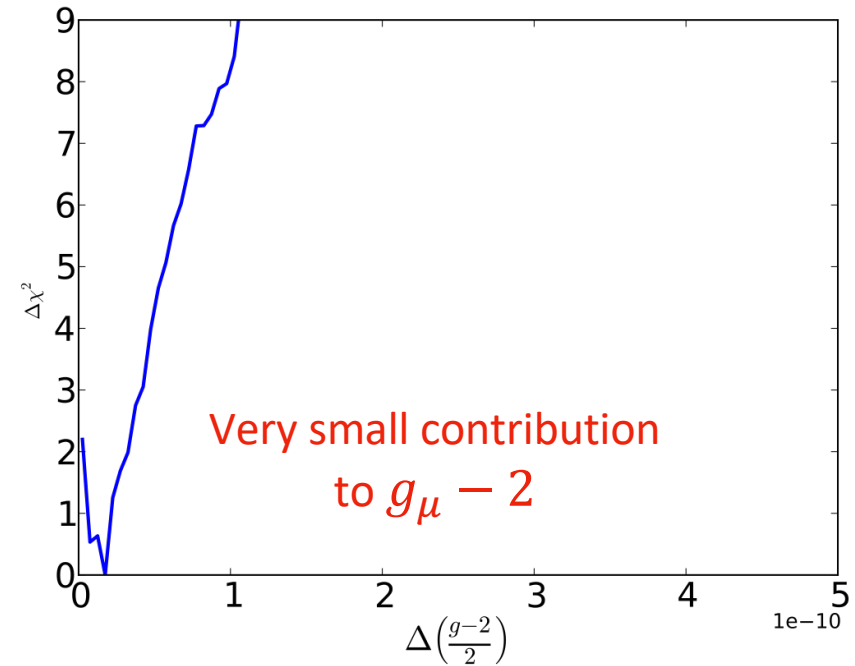


# $g_\mu - 2$ in Supersymmetric SU(5) GUT (CMSSM)

Assume universality between squark & slepton,  
and between gluino and electroweakino masses  
at GUT scale



Scenario relates squark/gluino masses  
to slepton/neutralino masses  
Cannot accommodate BNL/FNAL result  
Smuon masses  $\gtrsim 4$  TeV



# Flipped SU(5) GUT

- Extend GUT SU(5) with additional U(1) [motivated by string theory]

Antoniadis, JE, Hagelin & Nanopoulos, 1987

- “Flipped” fermion assignments to representations:

$$\bar{f}_i(\bar{\mathbf{5}}, -3) = \{U_i^c, L_i\}, \quad F_i(\mathbf{10}, 1) = \{Q_i, D_i^c, N_i^c\}, \quad l_i(\mathbf{1}, 5) = E_i^c, \quad i = 1, 2, 3$$

- Break GUT symmetry with 10-dimensional Higgses, electroweak symmetry with 5-dimensional Higgses:

$$H(\mathbf{10}, 1) = \{Q_H, D_H^c, N_H^c\}, \quad \bar{H}(\bar{\mathbf{10}}, -1) = \{\bar{Q}_H, \bar{D}_H^c, \bar{N}_H^c\}$$

$$h(\mathbf{5}, -2) = \{T_{H_c}, H_d\}, \quad \bar{h}(\bar{\mathbf{5}}, 2) = \{\bar{T}_{\bar{H}_c}, H_u\}$$

Lightest neutralino  
& lighter smuon  
can have small masses

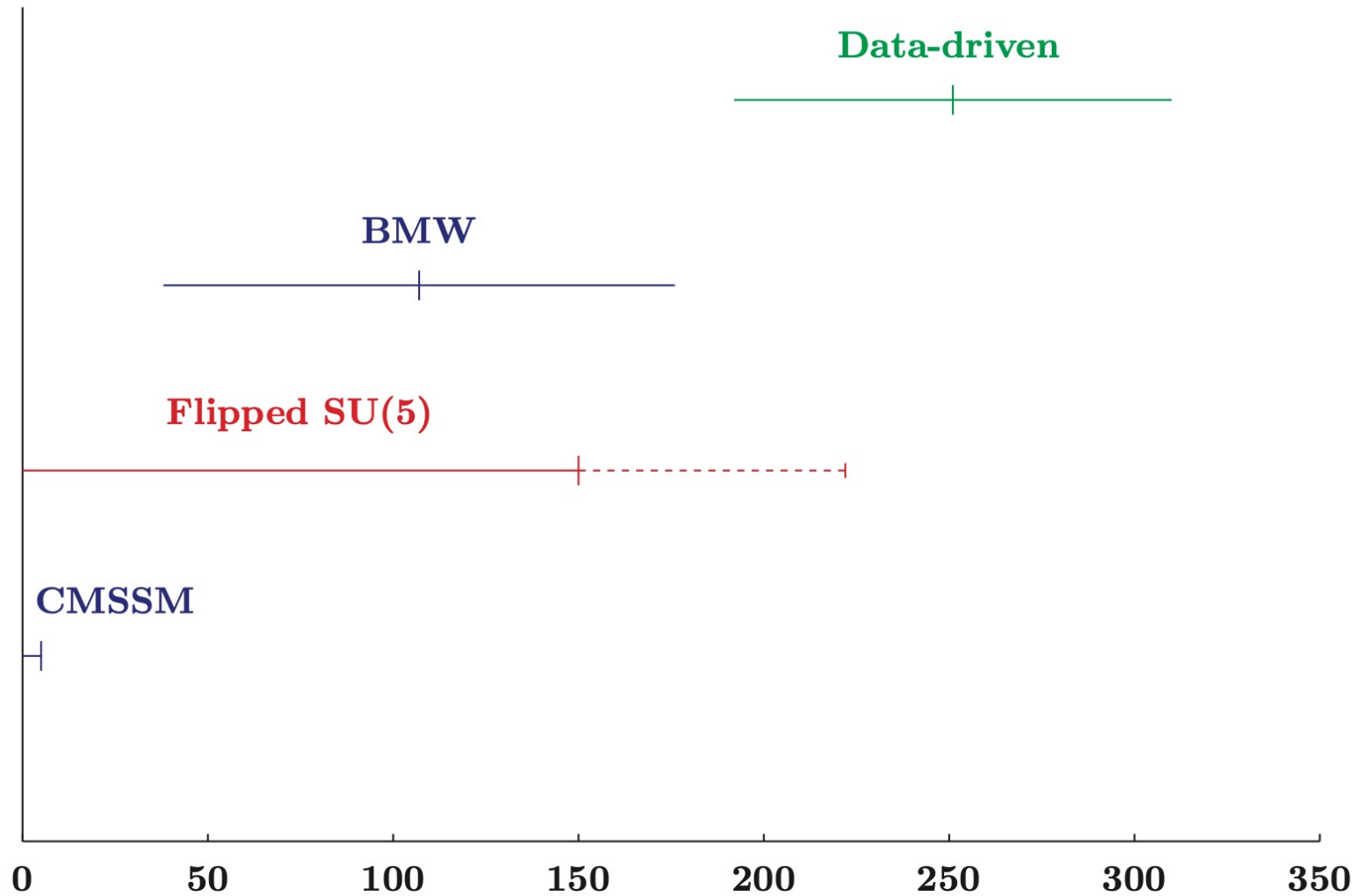
- Superpotential:

$$W = \lambda_1^{ij} F_i F_j h + \lambda_2^{ij} F_i \bar{f}_j \bar{h} + \lambda_3^{ij} \bar{f}_i \ell_j^c h + \lambda_4 H H h + \lambda_5 \bar{H} \bar{H} \bar{h} \\ + \lambda_6^{ia} F_i \bar{H} \phi_a + \lambda_7^a h \bar{h} \phi_a + \lambda_8^{abc} \phi_a \phi_b \phi_c + \mu_\phi^{ab} \phi_a \phi_b,$$

- Scan free parameters of model:

$$M_5, M_{X1}, m_{10}, m_5, m_1, \mu, M_A, A_0, \tan \beta$$

# $g_\mu - 2$ in CMSSM & Flipped SU(5) vs BMW Lattice, Data-Driven Calculations



$\Delta a_\mu (\times 10^{11})$ : GUT models vs Standard Model calculations

# $g_\mu = 2$ in Flipped SU(5)

## Parameters & predictions at best-fit point

Input GUT parameters (masses in units of $10^{16}$ GeV)		
$M_{GUT} = 1.00$	$M_X = 0.79$	$V = 1.13$
$\lambda_4 = 0.1$	$\lambda_5 = 0.3$	$\lambda_6 = 0.001$
$g_5 = 0.70$	$g_X = 0.70$	$m_{\nu_3} = 0.05$ eV
Input supersymmetry parameters (masses in GeV units)		
$M_5 = 2460$	$M_1 = 240$	$\mu = 4770$
$m_{10} = 930$	$m_{\bar{5}} = 450$	$m_1 = 0$
$M_A = 2100$	$A_0/M_5 = 0.67$	$\tan \beta = 35$
MSSM particle masses (in GeV units)		
$m_\chi = 84$	$m_{\tilde{t}_1} = 4030$	$m_{\tilde{g}} = 5090$
$m_{\chi_2} = 2160$	$m_{\chi_3} = 5080$	$m_{\chi_4} = 5080$
$m_{\tilde{\mu}_R} = 101$	$m_{\tilde{\mu}_L} = 1600$	$m_{\tilde{\tau}_1} = 1010$
$m_{\tilde{q}_L} = 4470$	$m_{\tilde{d}_R} = 4250$	$m_{\tilde{u}_R} = 4170$
$m_{\tilde{t}_2} = 4410$	$m_{\tilde{b}_1} = 4170$	$m_{\tilde{b}_2} = 4400$
$m_{\chi^\pm} = 2160$	$m_{H,A} = 2100$	$m_{H^\pm} = 2100$
Other observables		
$\Delta a_\mu = 150 \times 10^{-11}$	$\Omega_\chi h^2 = 0.13$	$m_h = 122$ GeV
Normal-ordered $\nu$ masses:	$\tau_{p \rightarrow e^+ \pi^0} _{\text{NO}} = 1.1 \times 10^{36}$ yrs	$\tau_{p \rightarrow \mu^+ \pi^0} _{\text{NO}} = 1.1 \times 10^{36}$ yrs
Inverse-ordered $\nu$ masses:	$\tau_{p \rightarrow e^+ \pi^0} _{\text{IO}} = 3.2 \times 10^{37}$ yrs	$\tau_{p \rightarrow \mu^+ \pi^0} _{\text{IO}} = 2.3 \times 10^{36}$ yrs

Opportunities to search for light smuon, neutralino at LHC  
Other sparticles too heavy?

Get good CDM density without even trying!

Refrain from “Supersymmetry” by Arcade Fire:

*I know you're living in my mind  
It's not the same as being alive*

**My Message to Ignatios:**

**Keep on Flipping!**