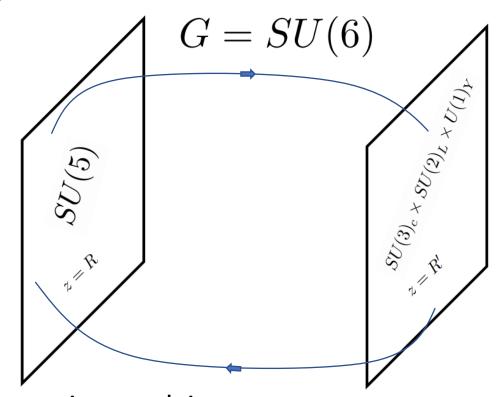
MAX-PLANCK-INSTITUT FÜR KERNPHYSIK HEIDELBERG



SU(6) Gauge-Higgs Grand Unification

Andreas Bally
Max-Planck-Institut für Kernphysik
30/05/2022: Planck 2022 Paris



Based on Phys. Rev. D 105, 035026 (2104.07366) and upcoming work in collaboration with **Andrei Angelescu** (MPIK), **Simone Blasi**(VUB) & **Florian Goertz**(MPIK)

Gauge-Higgs Unification:

Gauge field in 5 Dimensions

Manton 1979, Hosotani, Fairlie, Hatanaka, Inami, Lim...

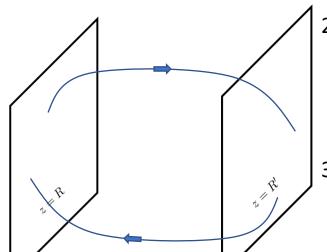
$$A_M = (A_\mu, A_5)$$

 No tree-level mass term by 5D gauge invariance: solution to the HP (in warped space)

Why Gauge-Higgs Uni.?

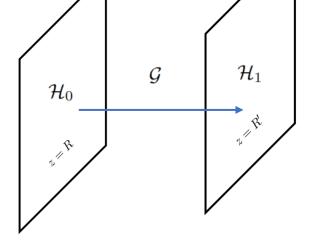
$$\frac{1}{2}m_{A_5}^2 A_5^2 \notin F_{MN} F^{MN}$$

2. The scalar sector unified within the gauge sector: break the gauge symmetry by boundary conditions

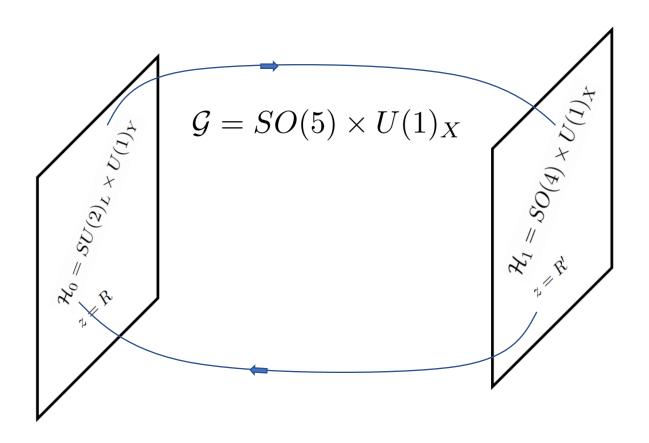


B. Dynamical origin of EWSB: Hosotani Mechanism

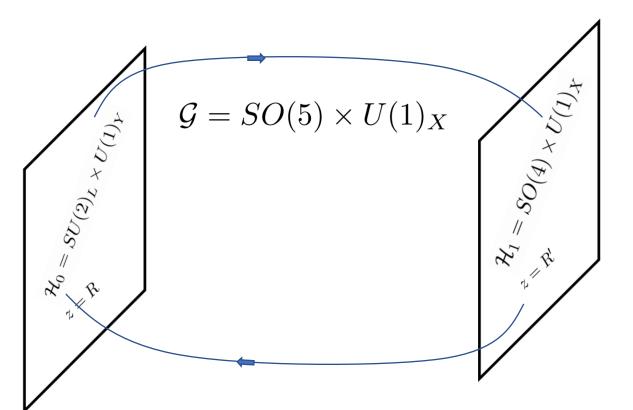
$$m_h \propto 1/R' \quad \left(= ke^{-kL} = m_{\rm KK} \right)$$

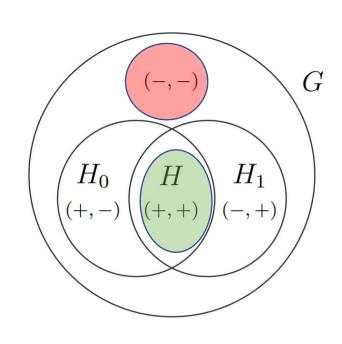


Electroweak Gauge-Higgs Unification



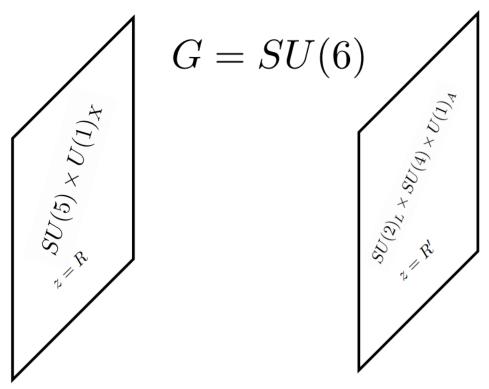
Electroweak Gauge-Higgs Unification





Massless Gauge Bosons: $T^a \in \mathcal{H} = \mathcal{H}_0 \cap \mathcal{H}_1 = SU(2)_L \times U(1)_Y \quad \left(= A^a_\mu(+,+)\right)$

Massless Scalars: $T^a\in \mathcal{G}/\mathcal{H}_0\cap \mathcal{G}/\mathcal{H}_1=(1,2)_{1/2} \quad \left(=A_\mu^{\dot{a}}(-,-)\right)$



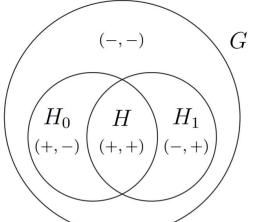
$$A_{\mu} = \begin{pmatrix} (++) & (++) & (++) & (+-) & (+-) & (+-) & (--) \\ (++) & (++) & (++) & (+-) & (+-) & (--) \\ \hline (+-) & (+-) & (++) & (++) & (++) & (-+) \\ (+-) & (+-) & (++) & (++) & (++) & (-+) \\ \hline (--) & (--) & (-+) & (-+) & (-+) & (++) \end{pmatrix}$$

Can be obtained from orbifold breaking

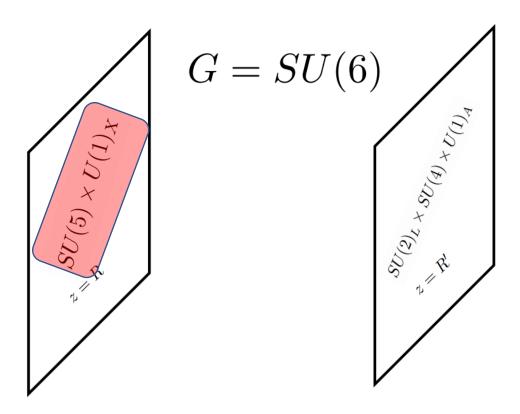
$$P = diag(1, 1, 1, 1, 1, -1)$$

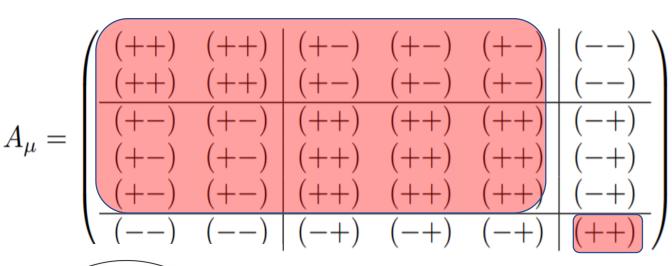
$$P' = diag(1, 1, -1, -1, -1, -1)$$

Exactly what we need!



$$SU(2)_L \times SU(3) \times U(1)_Y \times U(1)_X$$



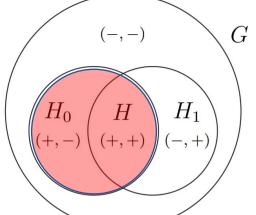


Can be obtained from orbifold breaking

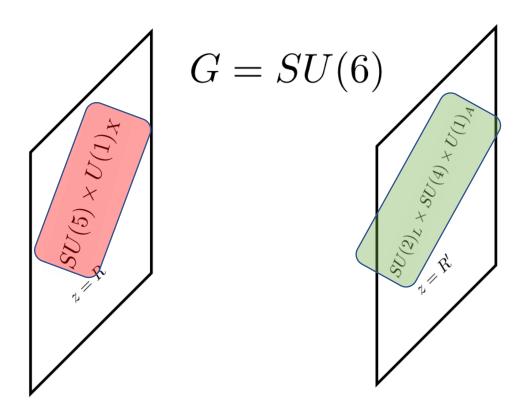
$$P = diag(1, 1, 1, 1, 1, -1)$$

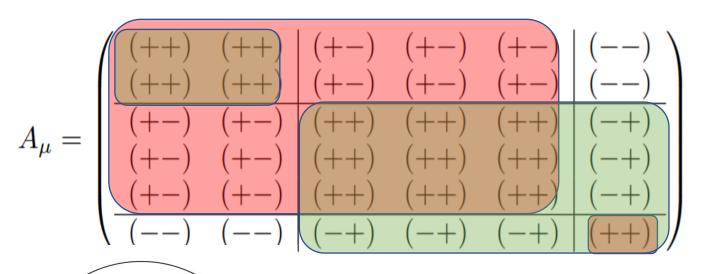
$$P' = diag(1, 1, -1, -1, -1, -1)$$

Exactly what we need!



$$SU(2)_L \times SU(3) \times U(1)_Y \times U(1)_X$$



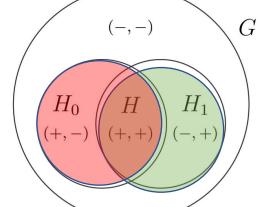


Can be obtained from orbifold breaking

$$P = diag(1, 1, 1, 1, 1, -1)$$

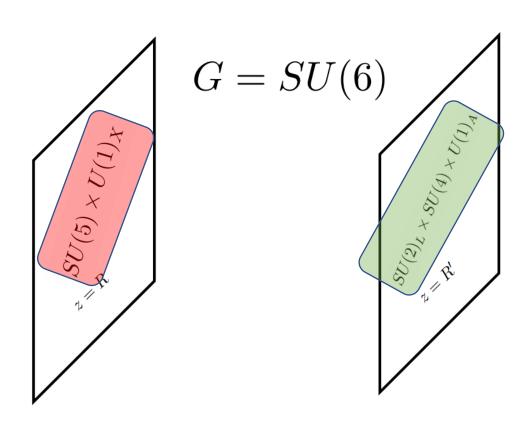
$$P' = diag(1, 1, -1, -1, -1, -1)$$

Exactly what we need!



$$SU(2)_L \times SU(3) \times U(1)_Y \times U(1)_X$$

$$({f 1},{f 2})_{1/2}$$

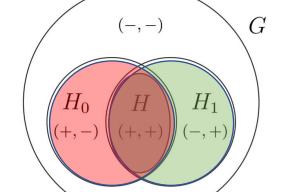


Can be obtained from orbifold breaking

$$P = diag(1, 1, 1, 1, 1, -1)$$

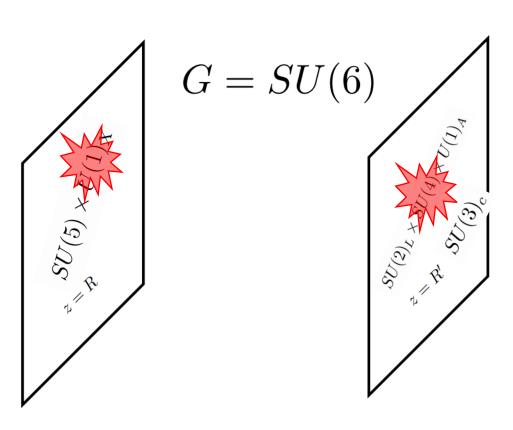
$$P' = diag(1, 1, -1, -1, -1, -1)$$

Exactly what we need!

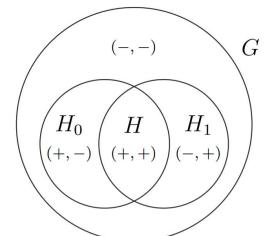


$$SU(2)_L \times SU(3) \times U(1)_Y \times U(1)_X$$

SU(6) GHGUT: Extra UV and IR breaking



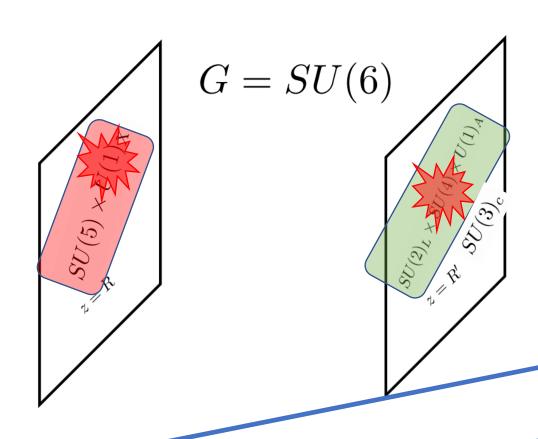
$$A_{\mu} = \begin{pmatrix} (++) & (++) & (+-) & (+-) & (+-) & (--) \\ (++) & (++) & (+-) & (+-) & (+-) & (--) \\ \hline (+-) & (+-) & (++) & (++) & (++) & (--) \\ (+-) & (+-) & (++) & (++) & (++) & (--) \\ \hline (--) & (--) & (--) & (--) & (--) & (--) \end{pmatrix}$$



$$SU(2)_L \times SU(3)_c \times U(1)_Y$$

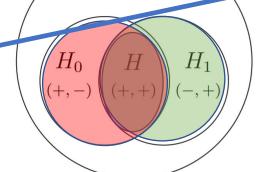
SU(6) GHGUT: Extra UV and IR breaking

 $(\mathbf{1},\mathbf{2})_{1/2}$



$$A_{\mu} = \begin{pmatrix} (++) & (++) & (++) & (+-) & (+-) & (--)$$

 $({f 3},{f 1})_{-1/3} \oplus ({f 1},{f 1})_0$



$$SU(2)_L \times SU(3)_c \times U(1)_Y$$

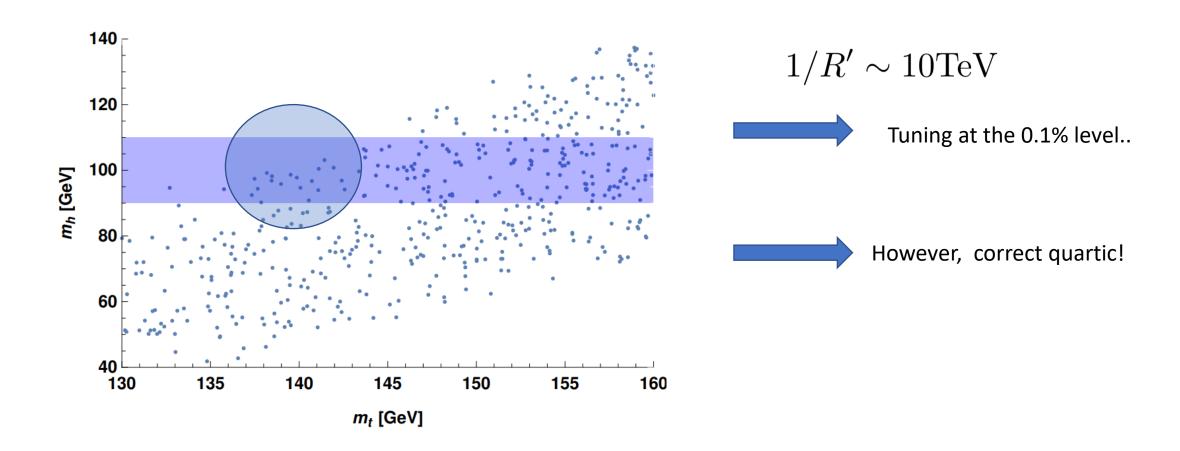
SU(6) GHGUT with Extra UV/IR Breaking: Fermionic content

$$\begin{aligned} \mathbf{20_L} &= \begin{cases} \mathbf{10_L} \rightarrow & (\mathbf{3}, \mathbf{2})_{1/6}^{-,+} \oplus (\mathbf{3}^*, \mathbf{1})_{-2/3}^{-,+} \oplus (\mathbf{1}, \mathbf{1})_{1}^{-,+} \\ \mathbf{10_L^*} \rightarrow & (\mathbf{3}^*, \mathbf{2})_{-1/6}^{-,+} \oplus u_R(\mathbf{3}, \mathbf{1})_{2/3}^{-,-} \oplus (\mathbf{1}, \mathbf{1})_{-1}^{-,+} \end{cases} \\ \mathbf{15_L} &= \begin{cases} \mathbf{10_L} \rightarrow & q_L(\mathbf{3}, \mathbf{2})_{1/6}^{+,+} \oplus (\mathbf{3}^*, \mathbf{1})_{-2/3}^{+,-} \oplus e_R^c(\mathbf{1}, \mathbf{1})_{1}^{+,+} \\ \mathbf{5_L} \rightarrow & (\mathbf{3}, \mathbf{1})_{-1/3}^{-,+} \oplus (\mathbf{1}, \mathbf{2})_{1/2}^{-,-} \end{cases} & \text{More breaking: allows for brane masses!} \\ \mathbf{6_L} &= \begin{cases} \mathbf{5_L} \rightarrow & d_R(\mathbf{3}, \mathbf{1})_{-1/3}^{-,-} \oplus l_L^c(\mathbf{1}, \mathbf{2})_{1/2}^{-,-} \\ \mathbf{1_L} \rightarrow & \nu_R^c(\mathbf{1}, \mathbf{1})_0^{+,+} \end{cases} & S_{UV} &= \int \mathrm{d}^4 x \left(M_u \psi_{\mathbf{20}, 10} \chi_{\mathbf{15}, 10} + \mathrm{h.c.} \right) \\ \mathbf{1_L} &= (\mathbf{1}, \mathbf{1})_0^{-,+} \end{cases} & S_{IR} &= \int \mathrm{d}^4 x \left(\frac{R}{R'} \right)^4 \left(M_{\tilde{u}} \psi_{\mathbf{15}, (3^*, 1)} \chi_{\mathbf{20}, (3^*, 1)} + M_d \chi_{\mathbf{15}, (3, 1)} \psi_{\mathbf{6}, (3, 1)} + M_d \chi_{\mathbf{15}, (1, 2)} \psi_{\mathbf{6}, (1, 2)} + M_{\nu} \chi_{\mathbf{6}, 1} \psi_{\mathbf{1}} + \mathrm{h.c.} \right) \end{aligned}$$

Reproduce all SM masses of the three generations + no light exotics!

Higgs potential

$$V_r(v, c, s) = \frac{N_r}{(4\pi)^2} \int_0^\infty dp \, p^3 \log(\rho_r(-p^2, v, c, s))$$



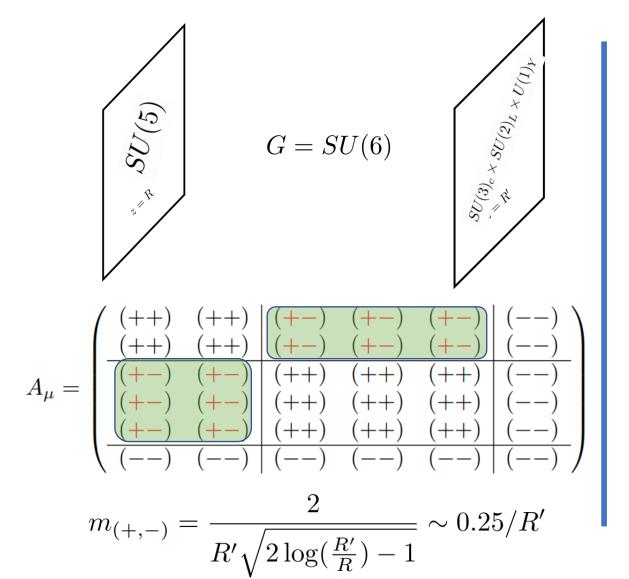
Exotic scalar spectrum: $(\mathbf{3},\mathbf{1})_{-1/3} \oplus (\mathbf{1},\mathbf{1})_0$

m_{LQ} [GeV]

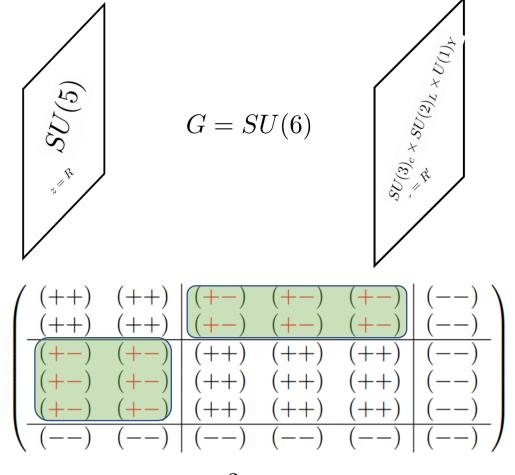
$$1/R' \sim 10 \text{TeV}$$

No vev for colored scalar!

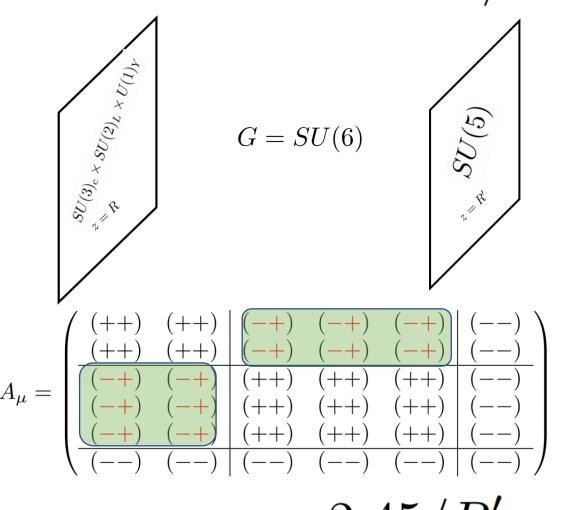
What about the X,Y gauge bosons? $(\mathbf{3},\mathbf{2})_{5/6}$



What about the X,Y gauge bosons? $(\mathbf{3},\mathbf{2})_{5/6}$



$$m_{(+,-)} = \frac{2}{R'\sqrt{2\log(\frac{R'}{R}) - 1}} \sim 0.25/R'$$



Little Hierarchy Problem : $1/R' \gtrsim 10 \text{TeV}$

Constraints on leptoquarks coupling to first generation fermions (Crivellin, Schnell arXiv: 210406417)

$$m_{X,Y} \gtrsim 2.5 \text{ TeV}$$

Unique to GHGUT

Constraints from the colored scalar: QCD double production

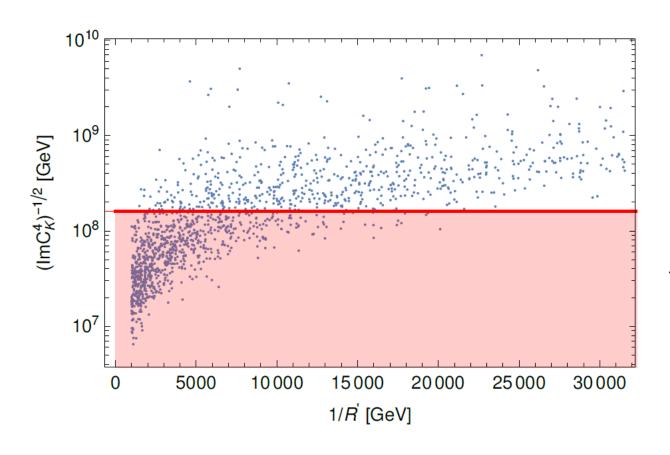
$$m_{S_1} \gtrsim 2 \text{ TeV}$$

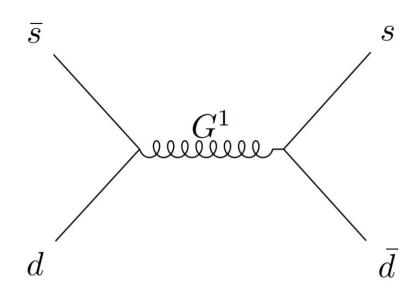
All extra dim models

Constraints from FCNCs (Csaki, Falkowski, Weiler arXiv:08041954)

$$m_{KK, \text{Gluon}} \gtrsim 25 \text{ TeV}$$

Tree-level flavor violation: Meson Mixing





$$\mathcal{H} \supset C_4^K Q_4^{ds} \sim \left(\frac{g^{ds}}{m_{G^1}}\right)^2 (\bar{d}_L^{\alpha} s_{R,\alpha}) (\bar{d}_R^{\beta} s_{L,\beta})$$

Loop-level flavor violation:

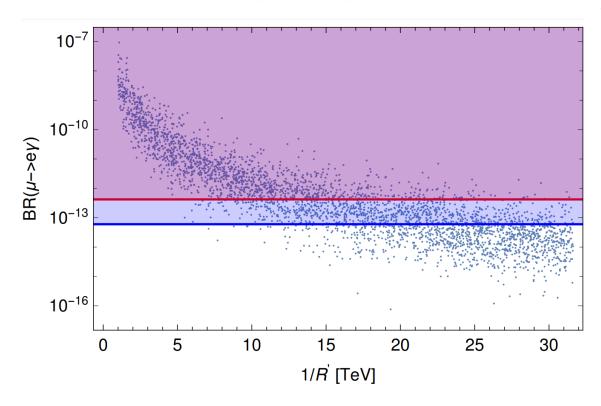
 $\mu \to e \gamma$

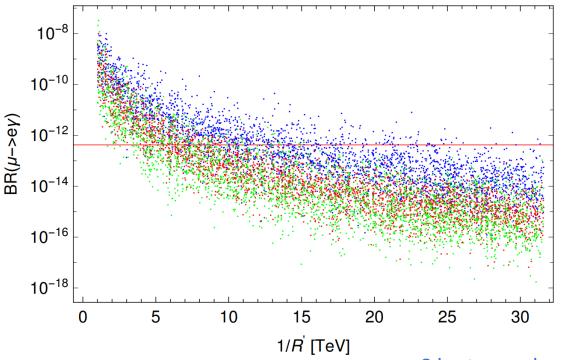
T T S T

$$Br(\mu \to e\gamma) < 4.2 \times 10^{-13}$$

MEG II (20??)

$$Br(\mu \to e\gamma) < 6 \times 10^{-14}$$





- S leptoquark
- Higgs
- Z boson

Proton Decay: no perturbative decay

There is no proton decay if one can consistently extend baryon number to each new field such that it is conserved at each vertex.

$$X^{\dagger}, Y^{\dagger} \sim (3^*, 2)_{5/6}$$

$$\mathbf{10} \to q_L(\mathbf{3}, \mathbf{2})_{1/6} \oplus u_R^c(\mathbf{3}^*, \mathbf{1})_{-2/3} \oplus e_R^c(\mathbf{1}, \mathbf{1})_1$$

$$X, Y \sim (3, 2)_{-5/6}$$

Proton Decay: no perturbative decay

There is no proton decay if one can consistently extend baryon number to each new field such that it is conserved at each vertex.

$$X^{\dagger}, Y^{\dagger} \sim (3^*, 2)_{5/6}$$

$$\mathbf{10} o q_L(\mathbf{3}, \mathbf{2})_{1/6} \oplus u_R^c(\mathbf{3}^*, \mathbf{1})_{-2/3} \oplus e_R^c(\mathbf{1}, \mathbf{1})_1$$

$$X, Y \sim (3, 2)_{-5/6}$$

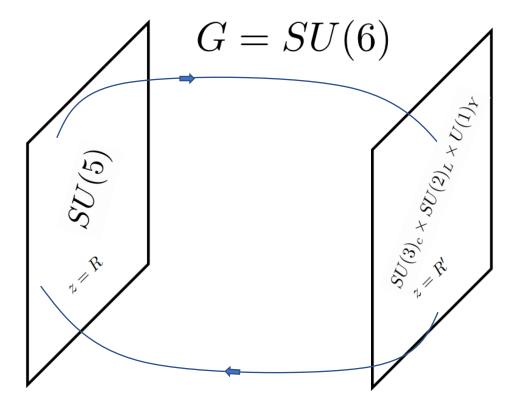
In Gauge-Higgs GUTS:
$$\mathbf{20} = \begin{cases} \mathbf{10} \rightarrow & q_L'(\mathbf{3}, \mathbf{2})_{1/6}^{-,+} \oplus (\mathbf{3}^*, \mathbf{1})_{-2/3}^{-,+} \oplus e_R^{c\prime}(\mathbf{1}, \mathbf{1})_1^{-,+} \\ \mathbf{10}^* \rightarrow & (\mathbf{3}^*, \mathbf{2})_{-1/6}^{-,+} \oplus u_R(\mathbf{3}, \mathbf{1})_{2/3}^{-,-} \oplus (\mathbf{1}, \mathbf{1})_{-1}^{-,+} \end{cases}$$

What about effective operators mediating proton decay? Turns out this baryon number can be gauged in the model 10

Conclusions

- Viable SU(6) GHGUT by introducing more breaking on the boundaries
- Minimal fermionic content leads to SM spectrum without light exotics
- Proton decay forbidden, exotics scalars, light X,Y bosons!
- Lots of directions to explore: flavor hierarchies, gauge coupling running, baryogenesis...

Thank you!



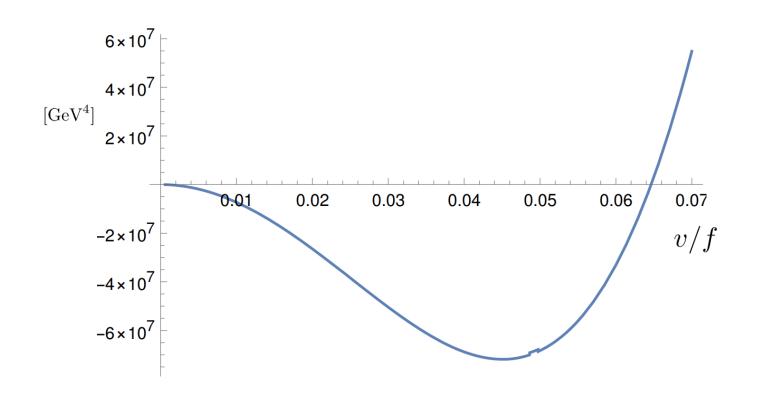
MAX-PLANCK-INSTITUT

FÜR KERNPHYSIK

HEIDELBERG

Andreas Bally

Higgs potential



$$1/R' \sim 10 \text{TeV}$$

$$f = \frac{2}{g\sqrt{\log R'/RR'}} \sim 5.3 \text{TeV}$$

$$v_{SM}/f \sim 0.05$$

SU(6) GHGUT: Extra UV Breaking $SU(5) \times U(1)_X \rightarrow SU(5)$

Via UV brane scalar

$$S_{UV} = \int d^4x \left((D_\mu \Phi_X)^\dagger (D^\mu \Phi_X) - V(\Phi^\dagger \Phi) + M_u \Phi^\dagger \psi_{\mathbf{20},10} \chi_{\mathbf{15},10} + \text{h.c.} \right)$$

$$\partial_5 A_\mu^X(z=R) = v_X A_\mu^X(z=R)$$

Via gauge boundary conditions

$$A_{\mu} = \begin{pmatrix} (++) & (++) & (+-) & (+-) & (+-) & (--) \\ (++) & (++) & (++) & (+-) & (+-) & (--) \\ \hline (+-) & (+-) & (++) & (++) & (++) & (-+) \\ (+-) & (+-) & (++) & (++) & (++) & (-+) \\ \hline (+-) & (+-) & (++) & (++) & (++) & (-+) \\ \hline (--) & (--) & (-+) & (-+) & (-+) & (-+) \end{pmatrix}$$

$$egin{aligned} \mathbf{20_L} = egin{cases} \mathbf{10_L}
ightarrow & (\mathbf{3}, \mathbf{2})_{1/6}^{-,+} \oplus (\mathbf{3}^*, \mathbf{1})_{-2/3}^{-,+} \oplus (\mathbf{1}, \mathbf{1})_{1}^{-,-} \ \mathbf{10_L^*}
ightarrow & (\mathbf{3}^*, \mathbf{2})_{1/6}^{+,+} \oplus \mathbf{\textit{u}_R} (\mathbf{3}, \mathbf{1})_{2/3}^{-,-} \oplus (\mathbf{1}, \mathbf{1})_{-1}^{-,+} \end{cases} \end{aligned}$$

$$\mathbf{15_L} = egin{cases} \mathbf{10_L}
ightarrow & q_L(\mathbf{3},\mathbf{2})_{1/6}^{+,+} \oplus (\mathbf{3}^*,\mathbf{1})_{-2/3}^{+,-} \oplus e_R^c(\mathbf{1},\mathbf{1})_1^{+,+} \ \mathbf{5_L}
ightarrow & d_R(\mathbf{3},\mathbf{1})_{-1/3}^{-,-} \oplus (\mathbf{1},\mathbf{2})_{1/2}^{-,+}, \end{cases}$$

Solves problems:

- Massless U(1)_X is eliminated!
- Up quark is massive!
- Still too constrained..

SU(6) GHGUT: Extra IR Breaking

$$SU(2)_L \times SU(4) \times U(1)_A \rightarrow SU(2)_L \times SU(3)_c \times U(1)_V$$

Via IR brane scalar

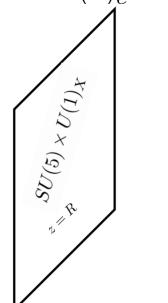
$$S_{IR} = \int d^4x \left((D_{\mu} \Phi_A)^{\dagger} (D^{\mu} \Phi_A) - V(\Phi_A^{\dagger} \Phi_A) \right)$$

7 (pseudo) Nambu Goldstone bosons

$$SU(4) \times U(1)_A / SU(3)_c \times U(1)_Y = (3,1)_{-1/3} \oplus (3^*,1)_{1/3} \oplus (1,1)_0$$

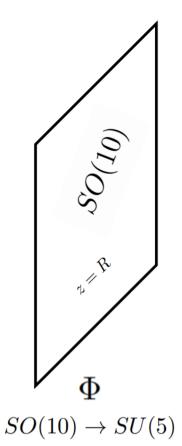
Via gauge boundary conditions

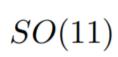
$$A_{\mu} = \begin{pmatrix} (++) & (++) & (++) & (+-) & (+-) & (--) \\ (++) & (++) & (++) & (+-) & (+-) & (--) \\ \hline (+-) & (+-) & (++) & (++) & (++) & (--) \\ \hline (+-) & (+-) & (++) & (++) & (++) & (--) \\ \hline (--) & (--) & (--) & (--) & (--) & (--) \end{pmatrix}$$

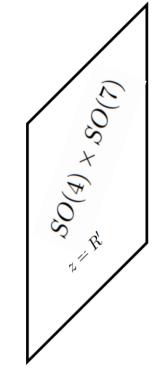


SO(11)

A custodial GHGUT





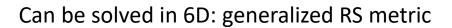


Higgs too light

Top quark	Bulk parameters			Brane parameters			Higgs
$m_t[{ m GeV}]$	c_0	c_1	c_2	μ_2	μ_3	μ_6	$m_H[{ m GeV}]$
165.0	0.3696	0.4286	0.2970	9.05×10^{10}	21.8	0.00249	50.96
170.0	0.3559	0.4293	0.3120	5.20×10^{10}	36.8	0.00420	51.77
175.0	0.3496	0.4286	0.3270	2.95×10^{10}	62.8	0.00719	53.52

• Light exotics

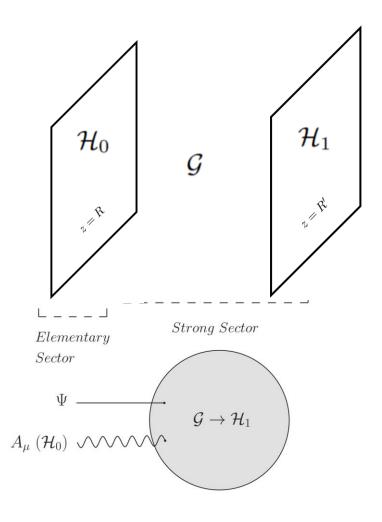
$$\frac{m_{\hat{u}}}{m_u} = \cot \frac{1}{2}\theta_H$$



Hosotani, Yamatsu arXiv: 1710.04811 Hosotani, Yamatsu arXiv: 1706.03503

Holographic dictionary

Extra dimensional models/Composite Higgs models



The 5D model gives a calculable model for the strong sector.

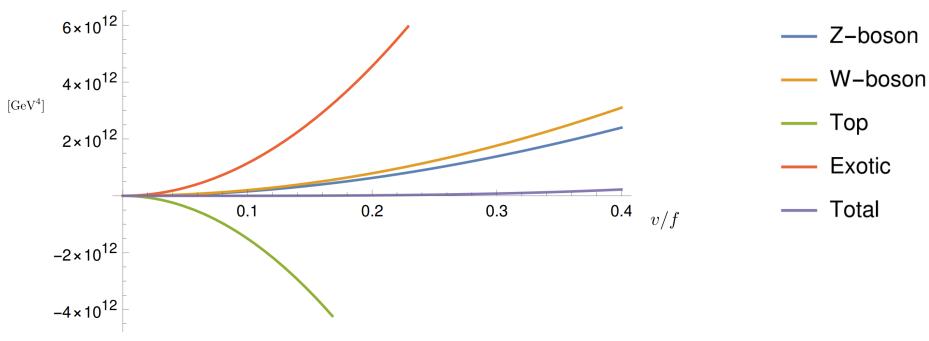
Otherwise, we are left with form factors that we have to approximate using the large SU(N) formalism

$$\frac{1}{N} \Longleftrightarrow \frac{g_5^2}{16\pi^2} \ll 1$$

Higgs potential

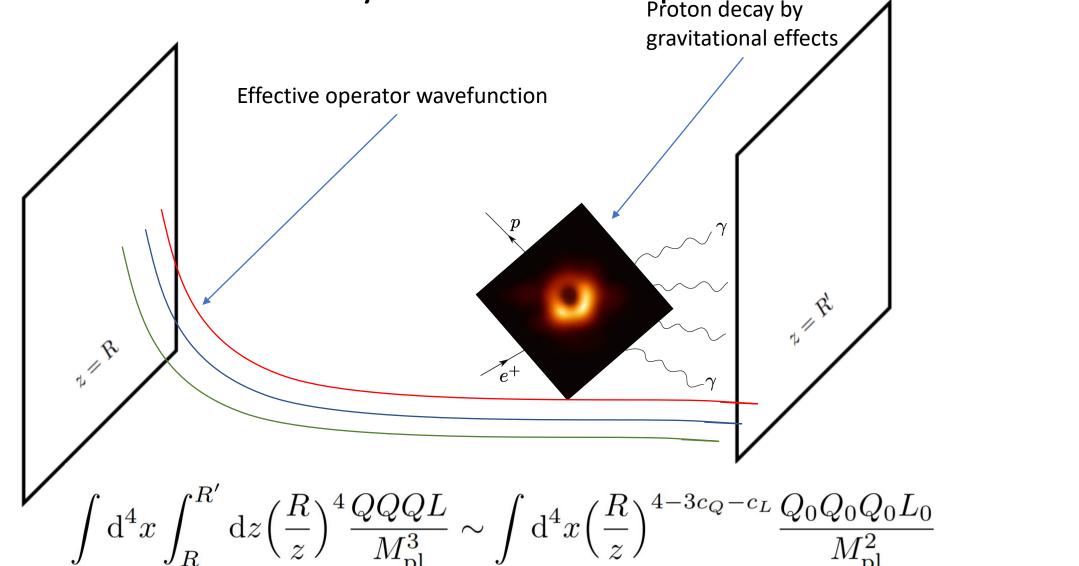
Little Hierarchy





$$f = \frac{2}{g\sqrt{\log R'/R}R'} \sim 5.3 \text{TeV}$$

Proton decay: effective operators in the bulk



U(6): Gauged Baryon number

$$SU(6) \rightarrow U(6) = SU(6) \times U(1)_C \rightarrow G_{SM} \times U(1)_B$$

$$U(1)_C: \quad ^{\Psi_6 o \, e^{ilpha}\Psi_6}$$

$$\Psi_6 \to e^{i\alpha} \Psi_6$$

$$\Psi_{15} \to e^{2i\alpha} \Psi_{15}$$

$$\Psi_{20} \rightarrow e^{3i\alpha} \Psi_{20}$$

C is (spontaneously) broken by the brane masses and B is the remaining symmetry

$$T_C = diag(1, 1, 1, 1, 1, 1)$$

$$T_B \to \frac{1}{3} \text{diag}(0, 0, 1, 1, 1, 0)$$

$$B\Psi_{6} = \begin{cases} \mathbf{5} \to e^{i\alpha/3} d_{R}(\mathbf{3}, \mathbf{1})_{-1/3}^{-,-} \oplus l_{L}^{c}(\mathbf{1}, \mathbf{2})_{1/2}^{-,-} \\ \mathbf{1} \to v_{R}^{c}(\mathbf{1}, \mathbf{1})_{0}^{+,+} \end{cases}$$

$$B\Psi_{20} = \begin{cases} \mathbf{10} \to & e^{i\alpha/3} q'_L(\mathbf{3}, \mathbf{2})^{-,+}_{1/6} \oplus e^{2i\alpha/3} (\mathbf{3}^*, \mathbf{1})^{-,+}_{-2/3} \oplus e^{c'}_R(\mathbf{1}, \mathbf{1})^{-,+}_1 \\ \mathbf{10}^* \to & e^{2i\alpha/3} (\mathbf{3}^*, \mathbf{2})^{-,+}_{-1/6} \oplus e^{i\alpha/3} u_R(\mathbf{3}, \mathbf{1})^{-,-}_{2/3} \oplus e^{i\alpha} (\mathbf{1}, \mathbf{1})^{-,+}_{-1} \end{cases}$$

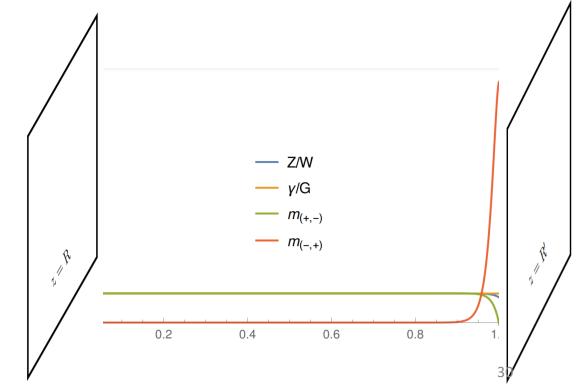
What about the X,Y gauge bosons? $(\mathbf{3},\mathbf{2})_{5/6}$

$$m_{(+,-)} = \frac{2}{R'\sqrt{2\log(\frac{R'}{R}) - 1}} \sim 0.25/R'$$

$$m_{(-,+)} \sim 2.45/R'$$

$$A_{\mu} = \begin{pmatrix} (++) & (++) & (++) & (+-) & (+-) & (--) \\ (++) & (++) & (++) & (+-) & (+-) & (--) \\ \hline (+-) & (++) & (++) & (++) & (++) & (--) \\ (+-) & (+-) & (++) & (++) & (++) & (--) \\ \hline (--) & (--) & (--) & (--) & (--) & (--) \end{pmatrix}$$

Flipping the gauge symmetries on the UV and IR brane switches the X,Y gauge bosons from (+,-) modes to (-,+) modes



Loop-level flavor violation: Suppressed by small neutrino masses and PMNS unitarity Higgs FCNCs due to small non-alignment of Higgs couplings and mass basis Color factor enhancement Even larger FCNCs than Higgs