

DYNAMICAL MINIMAL FLAVOUR VIOLATING INVERSE SEESAW

FERNANDO ARIAS ARAGÓN



Based on: 2204.04672,

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Fundamentals and Motivation

Low Scale Seesaws and
Minimal Lepton Flavour Violation

Low Scale Seesaws

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- Neutrino masses require New Physics: Seesaw Mechanism

Mohapatra and Senjanovic, PRD 23 (1981) 165
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Yanagida, CPC 7902131 (1979) 95-99
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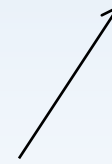
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Dinamical Inverse Seesaw in MFV

Setup, Phenomenology and
Impact of the CDF Measurement of M_W

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- Three scenarios:

CASE A

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$$\mu, \mu', \Lambda \propto \mathbb{1} \quad \mathcal{Y}_\nu \sim \mathcal{Y}'_\nu \sim (\mathbf{3}, 1, \bar{\mathbf{3}})$$

CASE B

$$\mathcal{G}_F^{\text{NA}} = SU(3)_V \times SU(3)_{e_R}$$

$$N_R, S_R \sim (\mathbf{3}, 1)$$

$$\mathcal{Y}_\nu, \mathcal{Y}'_\nu \propto \mathbb{1} \quad \mu \sim \mu' \sim \Lambda \sim (\bar{\mathbf{6}}, 1)$$

Dynamical ISS in MFV – Setup

- Froggatt-Nielsen + MFV: $U(1)_{FN} \equiv U(1)_{PQ} \supset \mathcal{G}_F$ FAA, L. Merlo, 1709.07039

- Flavon Φ introduced for PQ invariance, $x_\Phi = -1$

$$\mathcal{L}_Y \supset y_\psi \bar{\psi}_L H \psi_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{x_{\psi_R} - x_{\psi_L}}, \quad \varepsilon \equiv \frac{v_\Phi}{\sqrt{2}\Lambda_\Phi}, \quad m_\psi = y_\psi \frac{v}{\sqrt{2}} \varepsilon^{x_{\psi_R} - x_{\psi_L}}$$

$$x_u - x_Q = 0, \quad x_d - x_u \simeq \log_\varepsilon \frac{m_b}{m_t}, \quad x_e - x_\ell \simeq \log_\varepsilon \frac{m_\tau}{m_t}, \quad x_d - x_u = x_e - x_\ell = \begin{cases} 1 & \text{for } \varepsilon = 0.01, \\ 3 & \text{for } \varepsilon = 0.23, \end{cases}$$

- Three scenarios:

CASE A

$$\mathcal{G}_F^{\text{NA}} = SU(3)_{\ell_L} \times SU(3)_{e_R} \times SO(3)_{N_R+S_R}$$

$$\mu, \mu', \Lambda \propto \mathbb{1} \quad \mathcal{Y}_\nu \sim \mathcal{Y}'_\nu \sim (\mathbf{3}, 1, \bar{\mathbf{3}})$$

CASE B

$$\mathcal{G}_F^{\text{NA}} = SU(3)_V \times SU(3)_{e_R}$$

$$N_R, S_R \sim (\mathbf{3}, 1)$$

$$\mathcal{Y}_\nu, \mathcal{Y}'_\nu \propto \mathbb{1} \quad \mu \sim \mu' \sim \Lambda \sim (\bar{\mathbf{6}}, 1)$$

CASE C

$$\mathcal{G}_F^{\text{NA}} = SU(3)_V \times SU(3)_{e_R}$$

$$N_R \sim (\mathbf{3}, 1), \bar{S}_R \sim (\mathbf{3}, 1)$$

$$\mathcal{Y}_\nu, \Lambda \propto \mathbb{1} \quad \mu' \sim \mathcal{Y}'_\nu \sim \mu^\dagger \sim (\bar{\mathbf{6}}, 1)$$

Dynamical ISS in MFV – Case A

$$\mathcal{G}_F^{\text{NA}} = SU(3)_{\ell_L} \times SU(3)_{e_R} \times SO(3)_{N_R+S_R}$$

Dynamical ISS in MFV – Case A

$$\mathcal{G}_F^{\text{NA}} = SU(3)_{\ell_L} \times SU(3)_{e_R} \times SO(3)_{N_R+S_R}$$

$$\begin{aligned}
 -\mathcal{L}_Y^A = & \bar{\ell}_L H Y_e e_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{x_e - x_\ell} + \bar{\ell}_L \tilde{H} Y_\nu N_R + c_\nu \bar{\ell}_L \tilde{H} Y_\nu S_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell} + \\
 & + \frac{1}{2} c_N \bar{N}_R^c N_R \Phi \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \frac{1}{2} c_S \bar{S}_R^c S_R \Phi^\dagger \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \Lambda \bar{N}_R^c S_R + \text{h.c.}
 \end{aligned}
 \quad \mathcal{M}_X = \begin{pmatrix} 0 & \frac{v}{\sqrt{2}} Y_\nu & c_\nu \frac{v}{\sqrt{2}} \varepsilon^{2x_\ell} Y_\nu \\ \frac{v}{\sqrt{2}} Y_\nu^T & c_N \frac{v_\Phi}{\sqrt{2}} \varepsilon^{2x_\ell - 1} & \Lambda \\ c_\nu \frac{v}{\sqrt{2}} \varepsilon^{2x_\ell} Y_\nu^T & \Lambda & c_S \frac{v_\Phi}{\sqrt{2}} \varepsilon^{2x_\ell - 1} \end{pmatrix}$$

Dynamical ISS in MFV – Case A

$$\mathcal{G}_F^{\text{NA}} = SU(3)_{\ell_L} \times SU(3)_{e_R} \times SO(3)_{N_R+S_R}$$

$$\begin{aligned}
 -\mathcal{L}_Y^A = & \bar{\ell}_L H Y_e e_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{x_e - x_\ell} + \bar{\ell}_L \tilde{H} Y_\nu N_R + c_\nu \bar{\ell}_L \tilde{H} Y_\nu S_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell} + \\
 & + \frac{1}{2} c_N \bar{N}_R^c N_R \Phi \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \frac{1}{2} c_S \bar{S}_R^c S_R \Phi^\dagger \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \Lambda \bar{N}_R^c S_R + \text{h.c.}
 \end{aligned}$$

$$\mathcal{M}_X = \begin{pmatrix} 0 & \frac{v}{\sqrt{2}} Y_\nu & c_\nu \frac{v}{\sqrt{2}} \varepsilon^{2x_\ell} Y_\nu \\ \frac{v}{\sqrt{2}} Y_\nu^T & c_N \frac{v_\Phi}{\sqrt{2}} \varepsilon^{2x_\ell - 1} & \Lambda \\ c_\nu \frac{v}{\sqrt{2}} \varepsilon^{2x_\ell} Y_\nu^T & \Lambda & c_S \frac{v_\Phi}{\sqrt{2}} \varepsilon^{2x_\ell - 1} \end{pmatrix}$$

Dynamical ISS in MFV – Case A

$$\mathcal{G}_F^{\text{NA}} = SU(3)_{\ell_L} \times SU(3)_{e_R} \times SO(3)_{N_R+S_R}$$

$$\begin{aligned}
 -\mathcal{L}_Y^A = & \bar{\ell}_L H Y_e e_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{x_e - x_\ell} + \bar{\ell}_L \tilde{H} Y_\nu N_R + c_\nu \bar{\ell}_L \tilde{H} Y_\nu S_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell} + \\
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Dynamical ISS in MFV – Case A

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 -\mathcal{L}_Y^A = & \bar{\ell}_L H Y_e e_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{x_e - x_\ell} + \bar{\ell}_L \tilde{H} Y_\nu N_R + c_\nu \bar{\ell}_L \tilde{H} Y_\nu S_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell} + \\
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 \end{aligned}$$

$$\mathcal{M}_X = \begin{pmatrix} 0 & \frac{v}{\sqrt{2}} Y_\nu & \boxed{c_\nu} \frac{v}{\sqrt{2}} \varepsilon^{2x_\ell} \boxed{Y_\nu} \\ \frac{v}{\sqrt{2}} Y_\nu^T & \boxed{c_N} \frac{v_\Phi}{\sqrt{2}} \varepsilon^{2x_\ell - 1} & \boxed{\Lambda} \\ c_\nu \frac{v}{\sqrt{2}} \varepsilon^{2x_\ell} Y_\nu^T & \Lambda & \boxed{c_S} \frac{v_\Phi}{\sqrt{2}} \varepsilon^{2x_\ell - 1} \end{pmatrix}$$

Dynamical ISS in MFV – Case A

$$\mathcal{G}_F^{\text{NA}} = SU(3)_{\ell_L} \times SU(3)_{e_R} \times SO(3)_{N_R+S_R}$$

$$-\mathcal{L}_Y^A = \bar{\ell}_L H Y_e e_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{x_e - x_\ell} + \bar{\ell}_L \tilde{H} Y_\nu N_R + c_\nu \bar{\ell}_L \tilde{H} Y_\nu S_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell} +$$

$$+ \frac{1}{2} c_N \bar{N}_R^c N_R \Phi \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \frac{1}{2} c_S \bar{S}_R^c S_R \Phi^\dagger \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \Lambda \bar{N}_R^c S_R + \text{h.c.}$$

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$$Y_\nu = \frac{1}{f^{1/2}} U \hat{m}_\nu^{1/2} \mathcal{H}^T$$

$$U \equiv R_{23}(\theta_{23}) \cdot R_{13}(\theta_{13}, \delta_{\text{CP}}) \cdot R_{12}(\theta_{12}) \cdot \text{diag}(1, e^{i\alpha_2}, e^{i\alpha_3})$$

$$f \equiv \frac{v^2 \varepsilon^{2x_\ell - 1}}{2\sqrt{2}\Lambda^2} \left(c_S v_\Phi - 2\sqrt{2} c_\nu \varepsilon \Lambda \right)$$

$$\mathcal{H} \equiv e^{i\phi} = \mathbb{1} - \frac{\cosh r - 1}{r^2} \phi^2 + i \frac{\sinh r}{r} \phi$$

$$M_N \simeq \Lambda$$

$$\phi = \begin{pmatrix} 0 & \phi_1 & \phi_2 \\ -\phi_1 & 0 & \phi_3 \\ -\phi_2 & -\phi_3 & 0 \end{pmatrix}$$

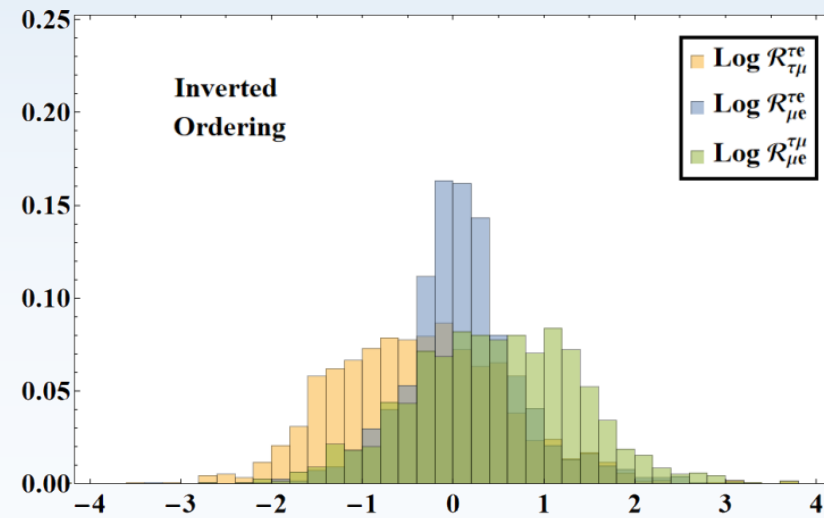
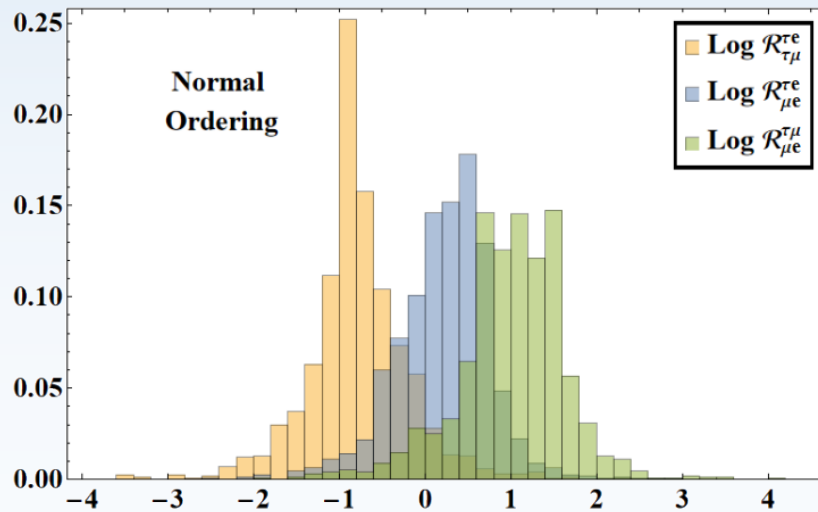
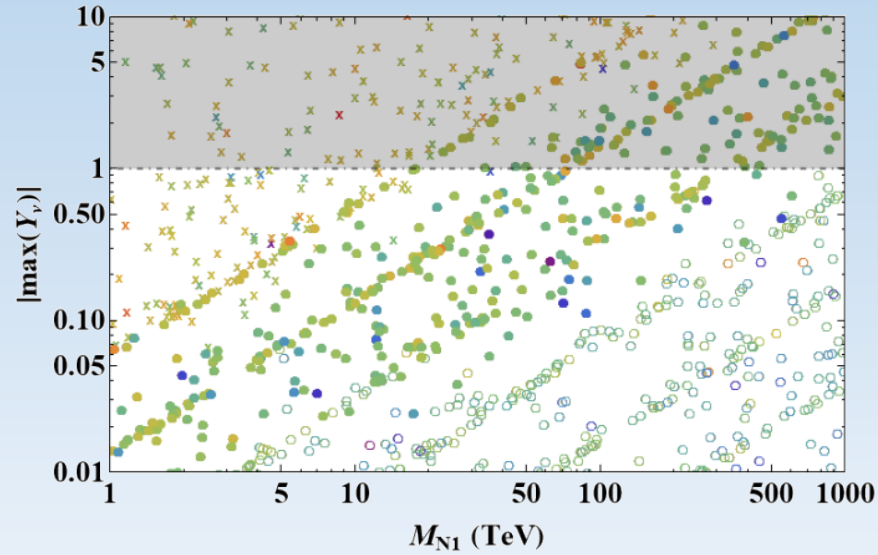
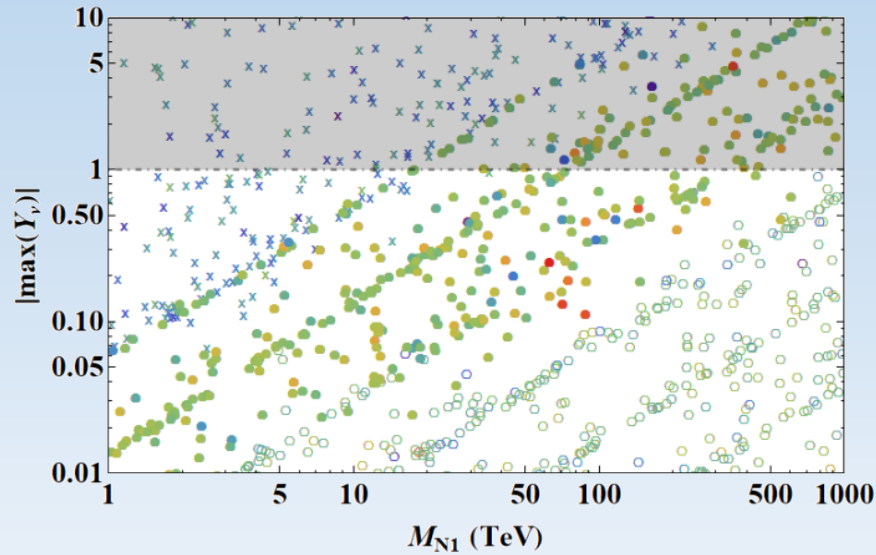
$$r \equiv \sqrt{\phi_1^2 + \phi_2^2 + \phi_3^2}$$

Dynamical ISS in MFV – Case A

$$\mathcal{N} = (\mathbb{1} - \eta) U \quad \eta = \frac{v^2}{4f\Lambda^2} U \widehat{m}_\nu^{1/2} \mathcal{H}^T \mathcal{H}^* \widehat{m}_\nu^{1/2} U^\dagger$$

Dynamical ISS in MFV – Case A

$$\mathcal{N} = (\mathbb{1} - \eta) U \quad \eta = \frac{v^2}{4f\Lambda^2} U \widehat{m}_\nu^{1/2} \mathcal{H}^T \mathcal{H}^* \widehat{m}_\nu^{1/2} U^\dagger$$



Dynamical ISS in MFV – Case B

$$\mathcal{G}_F^{\text{NA}} = SU(3)_V \times SU(3)_{e_R} \quad N_R, S_R \sim (\mathbf{3}, 1)$$

Dynamical ISS in MFV – Case B

$$\mathcal{G}_F^{\text{NA}} = SU(3)_V \times SU(3)_{e_R} \quad N_R, S_R \sim (\mathbf{3}, 1)$$

$$\begin{aligned}
 -\mathcal{L}_Y^{\text{B}} = & \bar{\ell}_L H Y_e e_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{x_e - x_\ell} + c_{\nu N} \bar{\ell}_L \tilde{H} N_R + c_{\nu S} \bar{\ell}_L \tilde{H} S_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell} + \\
 & + \frac{1}{2} c_N \bar{N}_R^c Y_N N_R \Phi \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \frac{1}{2} \bar{S}_R^c Y_N S_R \Phi^\dagger \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \\
 & + \Lambda \bar{N}_R^c Y_N S_R + \text{h.c.},
 \end{aligned}
 \quad \mathcal{M}_\chi = \begin{pmatrix} 0 & c_{\nu N} \frac{v}{\sqrt{2}} & c_{\nu S} \frac{v}{\sqrt{2}} \varepsilon^{2x_\ell} \\ c_{\nu N} \frac{v}{\sqrt{2}} & c_N \frac{v_\Phi}{\sqrt{2}} \varepsilon^{2x_\ell - 1} Y_N & \Lambda Y_N \\ c_{\nu S} \frac{v}{\sqrt{2}} \varepsilon^{2x_\ell} & \Lambda Y_N & \frac{v_\Phi}{\sqrt{2}} \varepsilon^{2x_\ell - 1} Y_N \end{pmatrix}$$

Dynamical ISS in MFV – Case B

$$\mathcal{G}_F^{\text{NA}} = SU(3)_V \times SU(3)_{e_R} \quad N_R, S_R \sim (\mathbf{3}, 1)$$

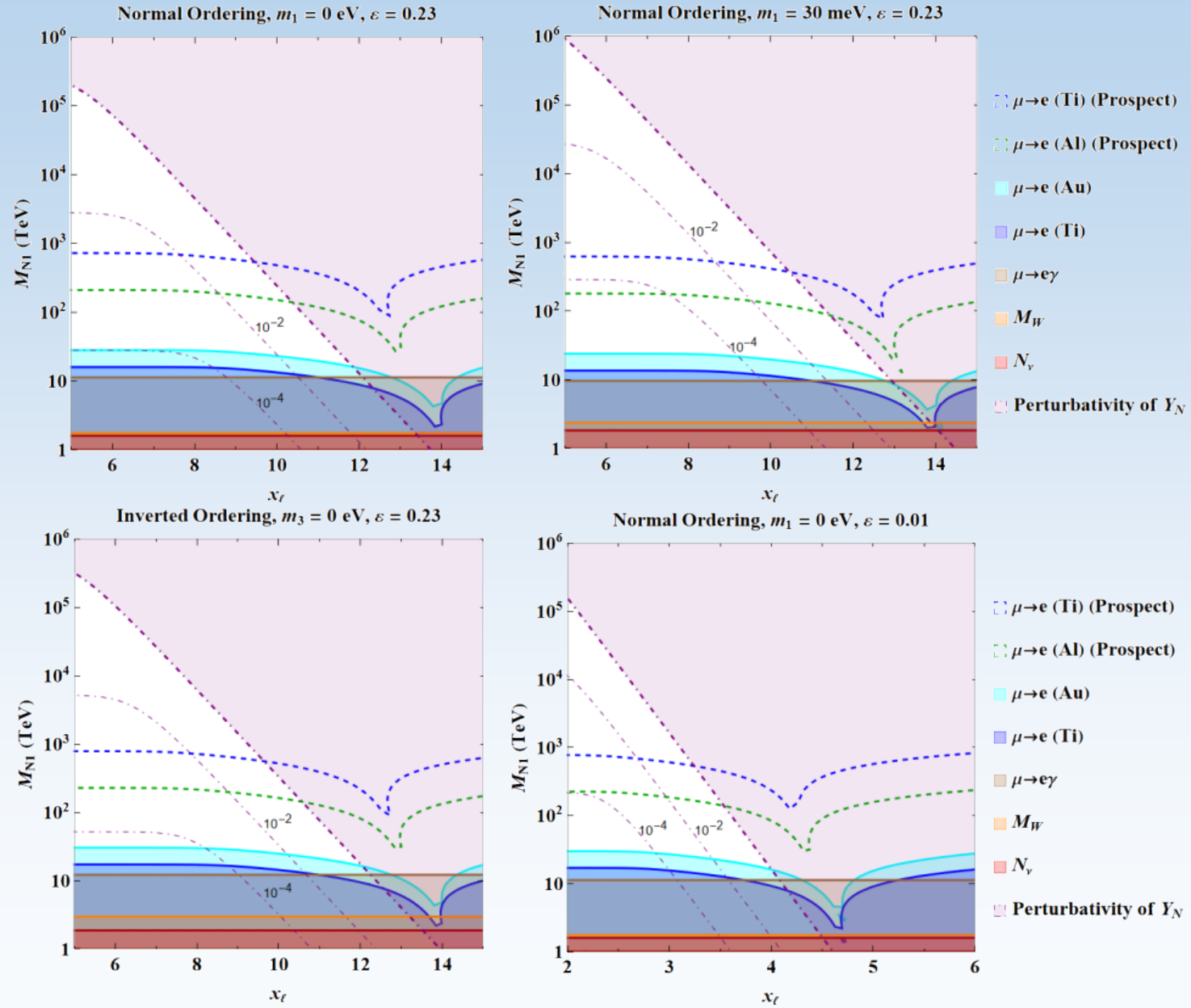
$$\begin{aligned}
 -\mathcal{L}_Y^{\text{B}} = & \bar{\ell}_L H Y_e e_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{x_e - x_\ell} + c_{\nu N} \bar{\ell}_L \tilde{H} N_R + c_{\nu S} \bar{\ell}_L \tilde{H} S_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell} + \\
 & + \frac{1}{2} c_N \bar{N}_R^c Y_N N_R \Phi \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \frac{1}{2} \bar{S}_R^c Y_N S_R \Phi^\dagger \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \\
 & + \Lambda \bar{N}_R^c Y_N S_R + \text{h.c.},
 \end{aligned}
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$$Y_N = f U^* \hat{m}_\nu^{-1} U^\dagger, \quad f \equiv \frac{v^2}{2} \varepsilon^{2x_\ell - 1} \left(\frac{c_{\nu N}^2 v_\Phi}{\sqrt{2} \Lambda^2} - \frac{2c_{\nu N} c_{\nu S} \varepsilon}{\Lambda} \right) \quad M_N \simeq \Lambda Y_N$$

Dynamical ISS in MFV – Case B

$$\mathcal{N} = (\mathbb{1} - \eta) U$$

$$\eta = \frac{c_{\nu N}^2 v^2}{4f^2 \Lambda^2} U \hat{m}_\nu^2 U^\dagger$$



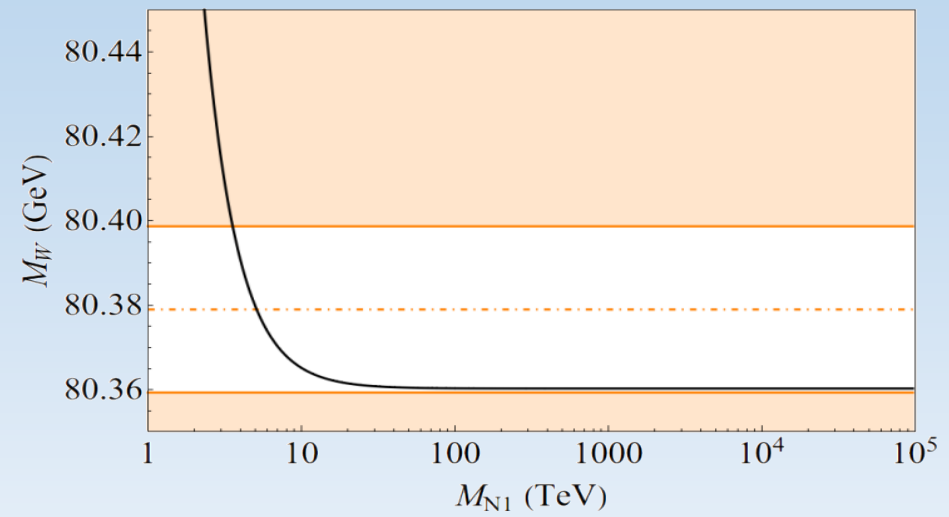
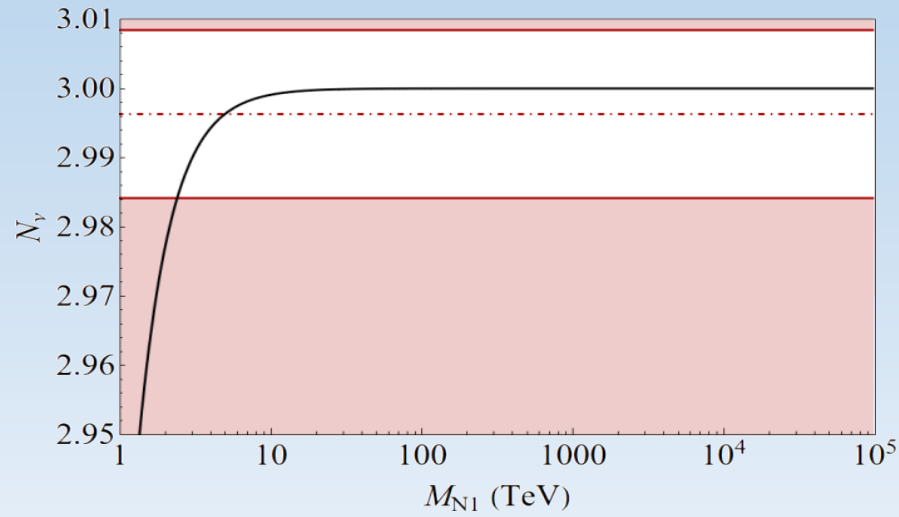
Dynamical ISS in MFV – Case C

$$\mathcal{N} = (1 - \eta) U \quad \eta = \frac{c_{\nu N}^2 v^2}{4\Lambda^2} \mathbb{1}$$

Dynamical ISS in MFV – Case C

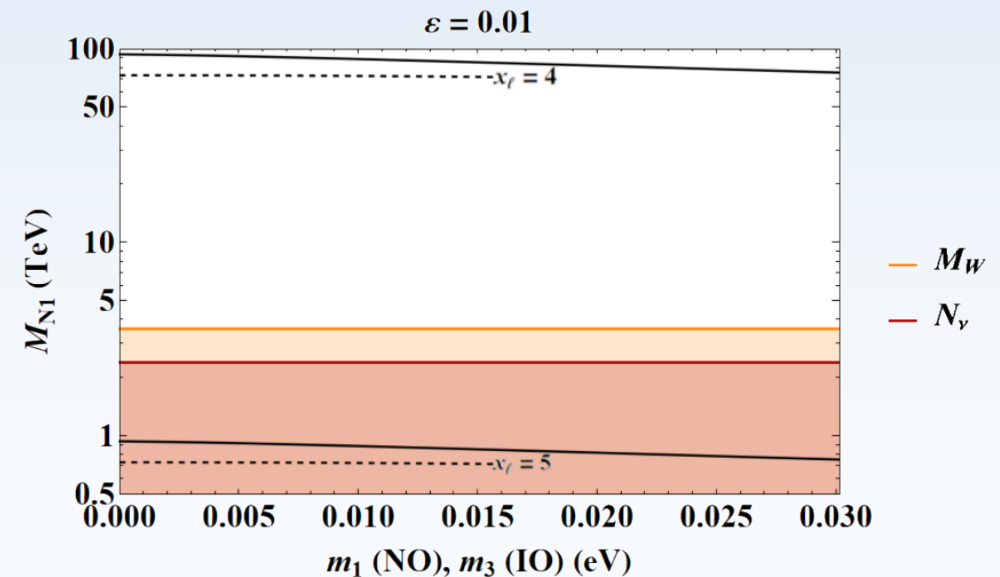
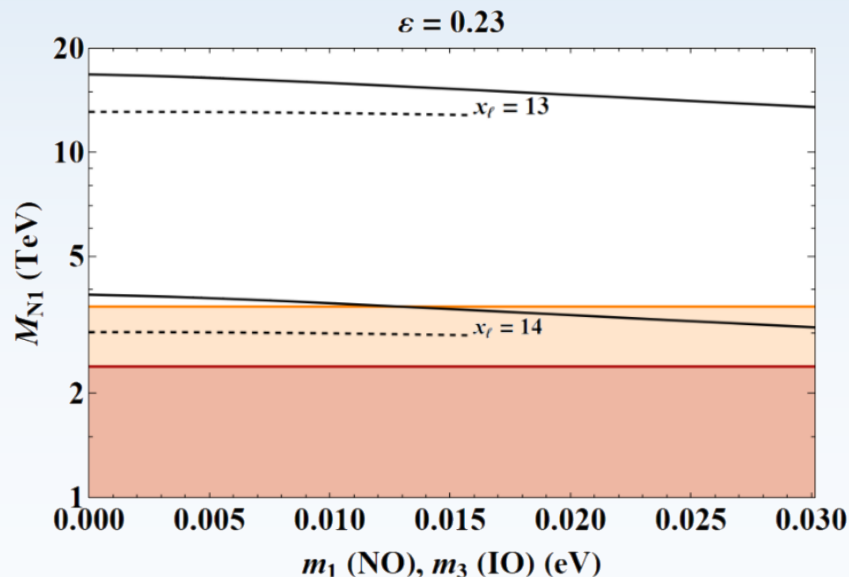
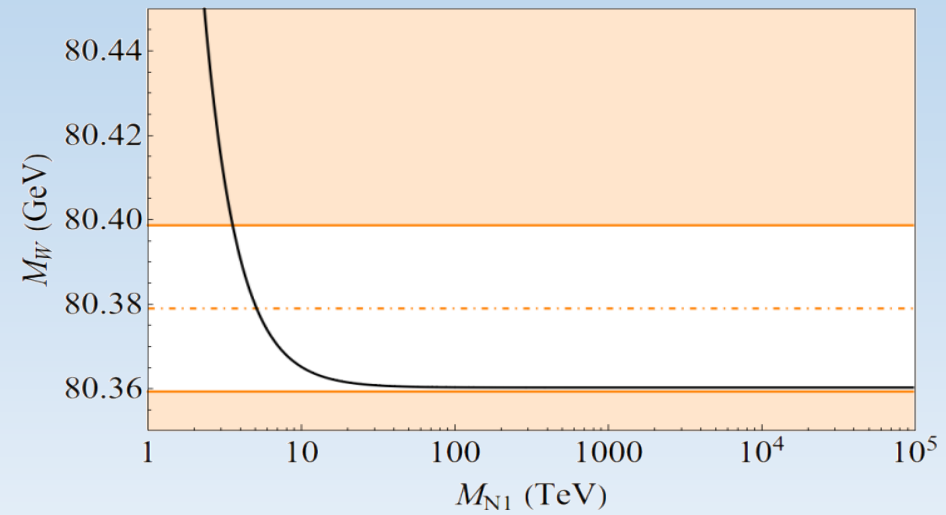
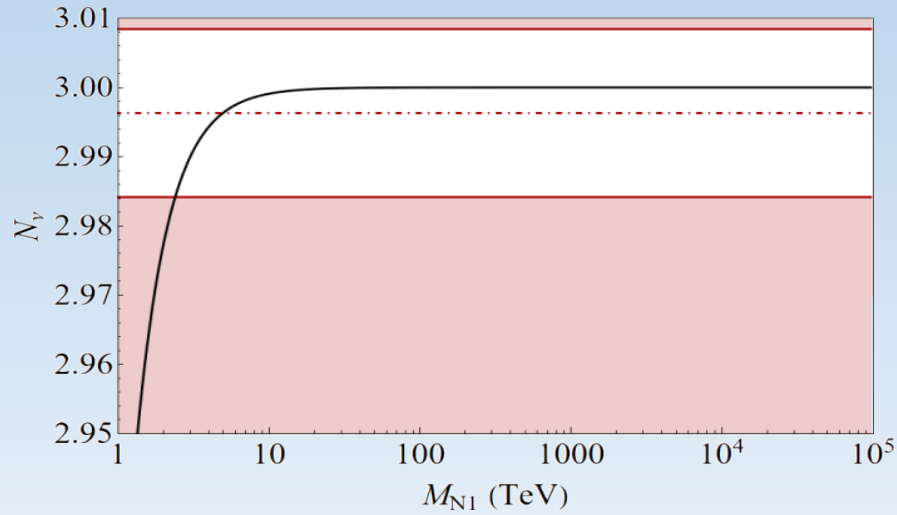
$$\mathcal{N} = (1 - \eta)U$$

$$\eta = \frac{c_{\nu N}^2 v^2}{4\Lambda^2} \mathbb{1}$$



Dynamical ISS in MFV – Case C

$$\mathcal{N} = (1 - \eta)U \quad \eta = \frac{c_{\nu N}^2 v^2}{4\Lambda^2} \mathbb{1}$$



Overview

Overview

- Inverse Seesaw embedded dynamically within MLFV

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- Strong CP Problem and Dark Matter addressed by the resulting axion

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Case A			
Case B			
Case C			

Overview

- Inverse Seesaw embedded dynamically within MLFV
- Strong CP Problem and Dark Matter addressed by the resulting axion

	LFV	M_{N_1} (TeV)	Best Probes
Case A			
Case B			
Case C			

Overview

- Inverse Seesaw embedded dynamically within MLFV
- Strong CP Problem and Dark Matter addressed by the resulting axion

	LFV	M_{N_1} (TeV)	Best Probes
Case A	✓	$\gtrsim \mathcal{O}(1)$	Colliders Indirect
Case B			
Case C			

Overview

- Inverse Seesaw embedded dynamically within MLFV
- Strong CP Problem and Dark Matter addressed by the resulting axion

	LFV	M_{N_1} (TeV)	Best Probes
Case A	✓	$\gtrsim \mathcal{O}(1)$	Colliders Indirect
Case B	✓	$\gtrsim \mathcal{O}(10)$	Indirect
Case C			

Overview

- Inverse Seesaw embedded dynamically within MLFV
- Strong CP Problem and Dark Matter addressed by the resulting axion

	LFV	M_{N_1} (TeV)	Best Probes
Case A	✓	$\gtrsim \mathcal{O}(1)$	Colliders Indirect
Case B	✓	$\gtrsim \mathcal{O}(10)$	Indirect
Case C	✗	$\gtrsim \mathcal{O}(1)$ [2.4,2.9](CDF)	Colliders

THANK YOU FOR YOUR
ATTENTION

BACKUP SLIDES

Dynamical ISS in MFV – Case C

$$\mathcal{G}_F^{\text{NA}} = SU(3)_V \times SU(3)_{e_R} \quad N_R \sim (\mathbf{3}, 1), \bar{S}_R \sim (\mathbf{3}, 1)$$

$$\begin{aligned}
 -\mathcal{L}_Y^{\text{C}} = & \bar{\ell}_L H Y_e e_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{x_e - x_\ell} + c_{\nu N} \bar{\ell}_L \tilde{H} N_R + c_{\nu S} \bar{\ell}_L \tilde{H} Y_N^\dagger S_R \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell} + \\
 & + \frac{1}{2} c_N \bar{N}_R^c Y_N N_R \Phi \left(\frac{\Phi}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \frac{1}{2} \bar{S}_R^c Y_N^\dagger S_R \Phi^\dagger \left(\frac{\Phi^\dagger}{\Lambda_\Phi} \right)^{2x_\ell - 1} + \\
 & + \Lambda \bar{N}_R^c S_R + \text{h.c.},
 \end{aligned}
 \quad \mathcal{M}_\chi = \begin{pmatrix} 0 & c_{\nu N} \frac{v}{\sqrt{2}} & c_{\nu S} \frac{v}{\sqrt{2}} \varepsilon^{2x_\ell} Y_N^\dagger \\ c_{\nu N} \frac{v}{\sqrt{2}} & c_N \frac{v_\Phi}{\sqrt{2}} \varepsilon^{2x_\ell - 1} Y_N & \Lambda \\ c_{\nu S} \frac{v}{\sqrt{2}} \varepsilon^{2x_\ell} Y_N^\dagger & \Lambda & \frac{v_\Phi}{\sqrt{2}} \varepsilon^{2x_\ell - 1} Y_N^\dagger \end{pmatrix}$$

$$Y_N = \frac{1}{f} U^* \hat{m}_\nu U^\dagger \quad f \equiv \frac{v^2}{2} \varepsilon^{2x_\ell - 1} \left(\frac{c_{\nu N}^2 v_\Phi}{\sqrt{2} \Lambda^2} - \frac{2c_{\nu N} c_{\nu S} \varepsilon}{\Lambda} \right) \quad M_N \simeq \Lambda$$

Dynamical ISS in MFV – Setup

$$M_W = M_Z \sqrt{\frac{1}{2} + \sqrt{\frac{1}{4} - \frac{\pi\alpha_{\text{em}}(1 - \eta_{\mu\mu} - \eta_{ee})}{\sqrt{2}G_\mu M_Z^2}}}$$

$$\Gamma_{Z\text{-inv}} = \frac{G_\mu M_Z^3}{12\sqrt{2}\pi} (3 - 4\eta_{\tau\tau} - \eta_{ee} - \eta_{\mu\mu})$$

$$\text{BR}(l_i \rightarrow l_j \gamma) \equiv \frac{\Gamma(l_i \rightarrow l_j \gamma)}{\Gamma(l_i \rightarrow l_j \nu \bar{\nu})} = \frac{3\alpha_{\text{em}}}{2\pi} |\eta_{l_j l_i}|^2$$

$$R_{\mu \rightarrow e} = \frac{\sigma(\mu^- X \rightarrow e^- X)}{\sigma(\mu^- X \rightarrow \text{Capture})} \simeq \frac{G_\mu^2 \alpha_{\text{em}}^5 m_\mu^5}{2s_w^4 \pi^4 \Gamma_{\text{capt}}} \frac{Z_{\text{eff}}^4}{Z} |\eta_{e\mu}|^2 F_p^2 \left[(A + Z) F_u + (2A - Z) F_d \right]^2$$

$$F_u = \frac{2}{3} s_W^2 \frac{16 \log\left(\frac{M_N^2}{M_W^2}\right) - 31}{12} - \frac{3 + 3 \log\left(\frac{M_N^2}{M_W^2}\right)}{8}$$

$$F_d = -\frac{1}{3} s_W^2 \frac{16 \log\left(\frac{M_N^2}{M_W^2}\right) - 31}{12} - \frac{3 - 3 \log\left(\frac{M_N^2}{M_W^2}\right)}{8}$$

Nucleus	$\frac{A}{Z}N$	Z_{eff}	F_p	$\Gamma_{\text{capt}} (10^6 \text{ s}^{-1})$
${}_{13}^{27}\text{Al}$		11.5	0.64	0.7054
${}_{22}^{48}\text{Ti}$		17.6	0.54	2.59
${}_{79}^{197}\text{Au}$		33.5	0.16	13.07

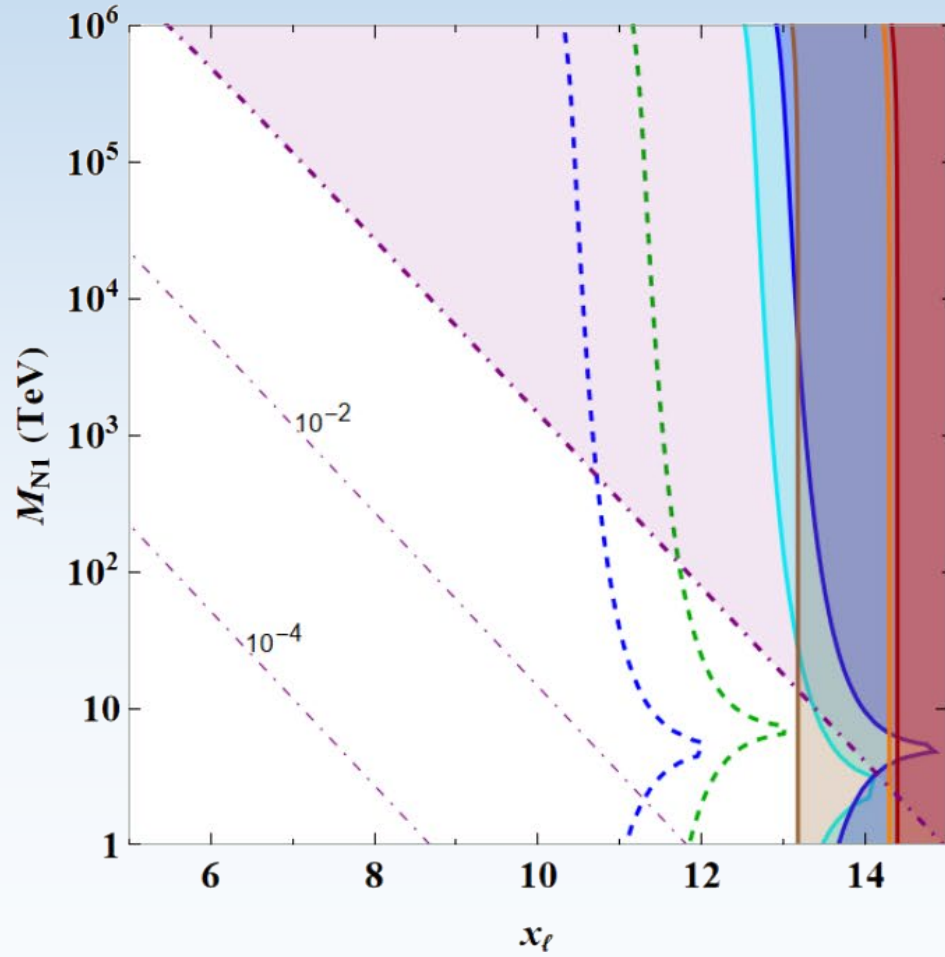
Dynamical ISS in MFV – Setup

Observable	Normal Ordering	Inverted Ordering	Observable	Experimental Bound	Future Sensitivity
$\sin^2 \theta_{12}$	$0.304^{+0.012}_{-0.012}$	$0.304^{+0.013}_{-0.012}$	$\text{BR}(\mu \rightarrow e\gamma)$	4.2×10^{-13}	6×10^{-14}
$\sin^2 \theta_{23}$	$0.450^{+0.019}_{-0.016}$	$0.570^{+0.016}_{-0.022}$	$\text{BR}(\tau \rightarrow e\gamma)$	1.9×10^{-7}	9×10^{-9}
$\sin^2 \theta_{13}$	$0.02246^{+0.00062}_{-0.00062}$	$0.02241^{+0.00074}_{-0.00062}$	$\text{BR}(\tau \rightarrow \mu\gamma)$	2.5×10^{-7}	6.9×10^{-9}
$\delta_{\text{CP}}/^\circ$	-130^{+36}_{-25}	-82^{+22}_{-30}	$R_{\mu \rightarrow e}(\text{Al})$	–	6×10^{-17}
$\frac{\Delta m_{\text{sol}}^2}{10^{-5} \text{ eV}^2}$	$7.42^{+0.21}_{-0.20}$	$7.42^{+0.21}_{-0.20}$	$R_{\mu \rightarrow e}(\text{Ti})$	4.3×10^{-12}	10^{-18}
$\frac{ \Delta m_{\text{atm}}^2 }{10^{-3} \text{ eV}^2}$	$2.510^{+0.027}_{-0.027}$	$2.490^{+0.026}_{-0.028}$	$R_{\mu \rightarrow e}(\text{Au})$	7×10^{-13}	–

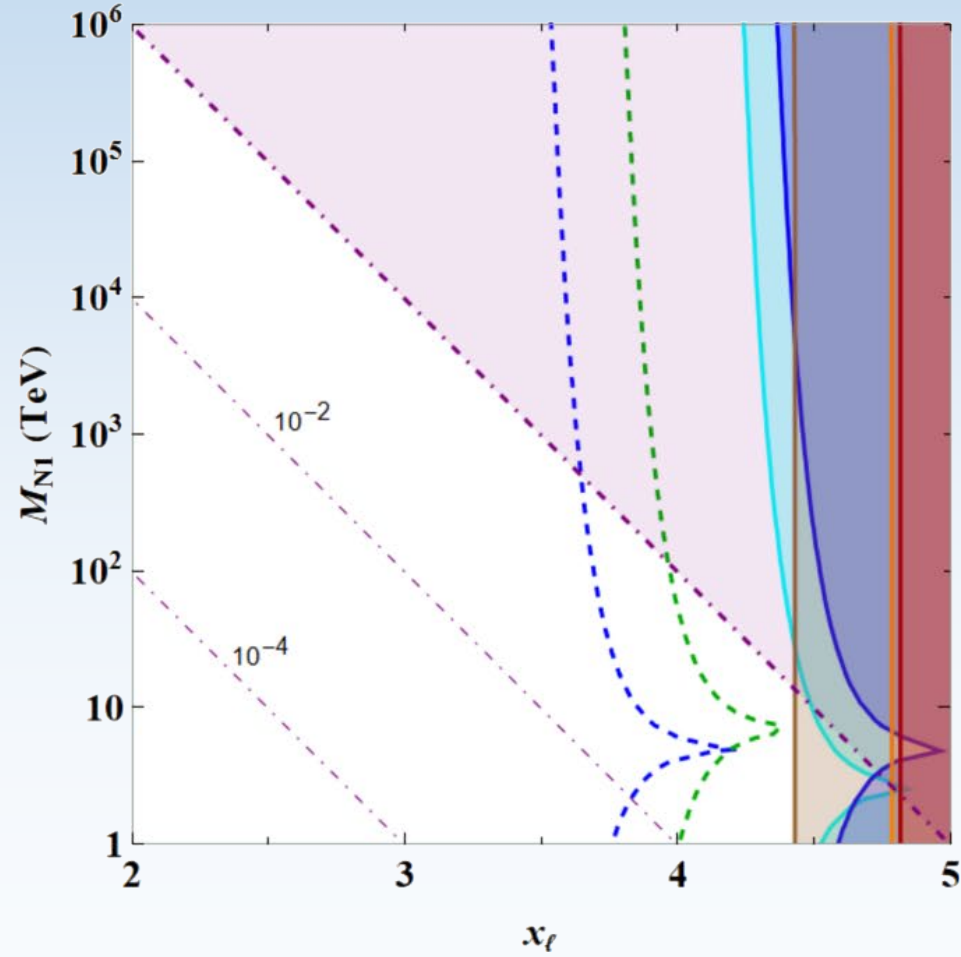
Dynamical ISS in MFV – Case A

$$\mathcal{N} = (\mathbb{1} - \eta) U \quad \eta = \frac{v^2}{4f\Lambda^2} U \widehat{m}_\nu^{1/2} \mathcal{H}^T \mathcal{H}^* \widehat{m}_\nu^{1/2} U^\dagger$$

Normal Ordering, $m_1 = 0$ eV,
 $\varepsilon = 0.23, r = 0$



Normal Ordering, $m_1 = 0$ eV,
 $\varepsilon = 0.01, r = 0$



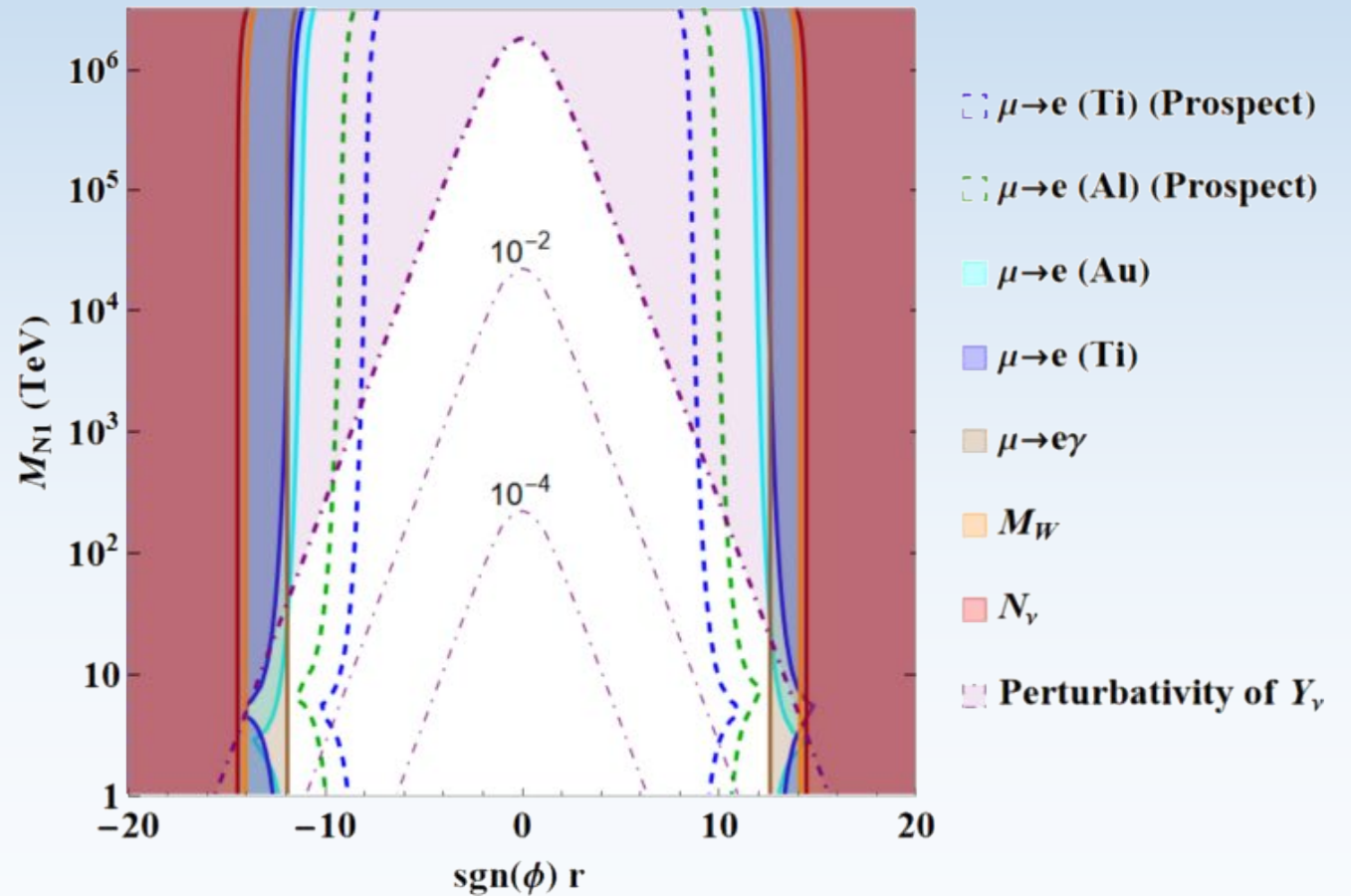
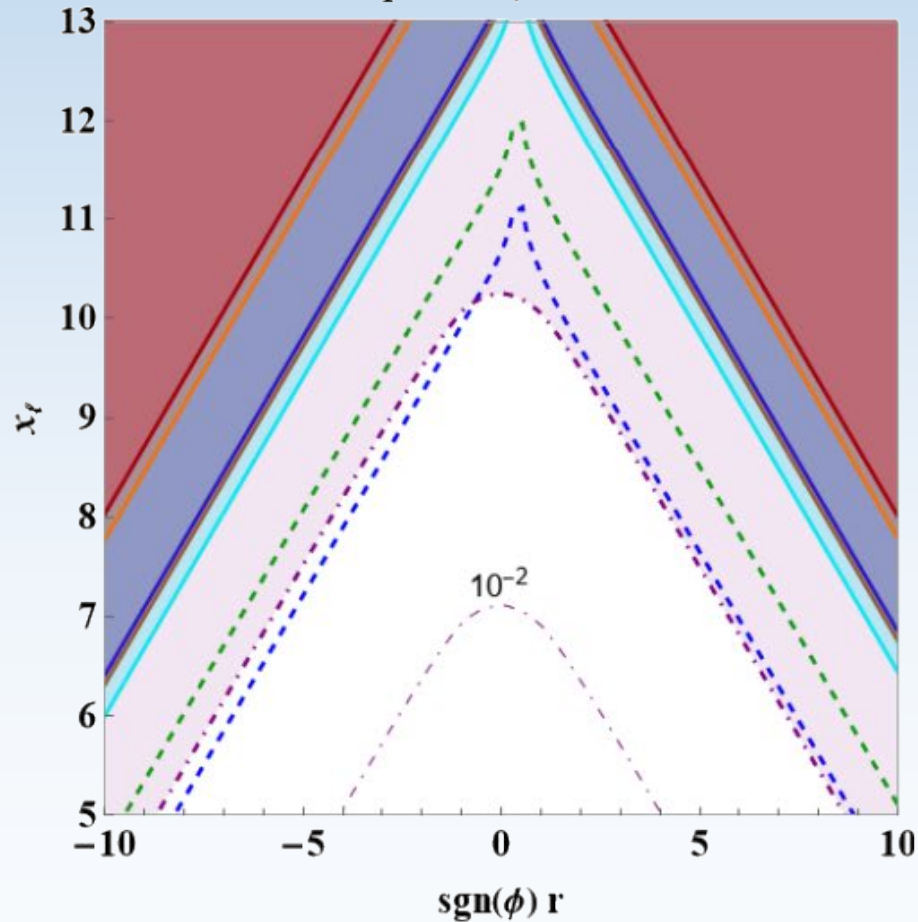
- - - $\mu \rightarrow e$ (Ti) (Prospect)
- - - $\mu \rightarrow e$ (Al) (Prospect)
- $\mu \rightarrow e$ (Au)
- $\mu \rightarrow e$ (Ti)
- $\mu \rightarrow e \gamma$
- M_W
- N_ν
- Perturbativity of Y_ν

Dynamical ISS in MFV – Case A

$$\mathcal{N} = (\mathbb{1} - \eta) U \quad \eta = \frac{v^2}{4f\Lambda^2} U \widehat{m}_\nu^{1/2} \mathcal{H}^T \mathcal{H}^* \widehat{m}_\nu^{1/2} U^\dagger$$

Normal Ordering, $M_{N1} = 10^3$ TeV,
 $m_1 = 0$ eV, $\varepsilon = 0.23$

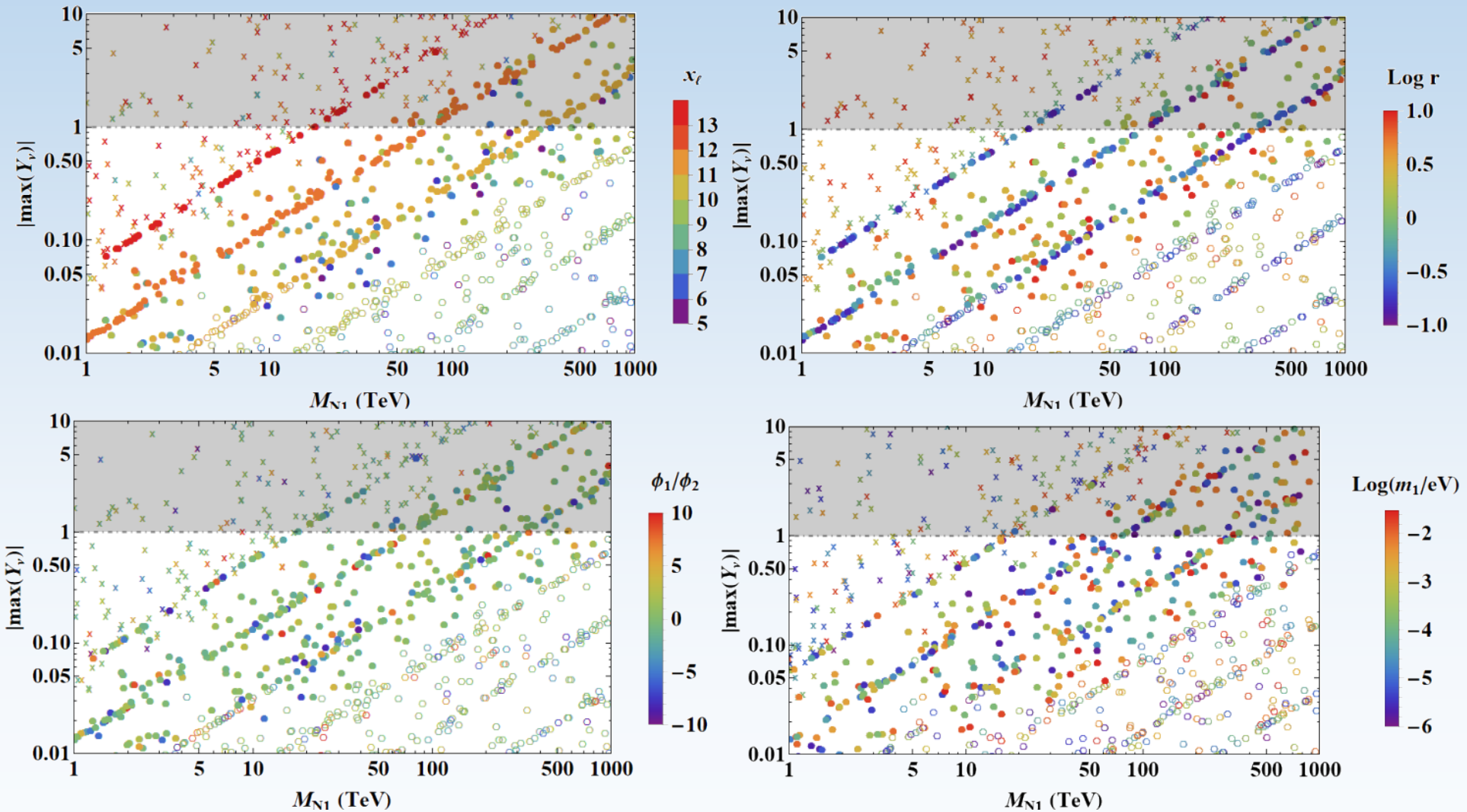
Normal Ordering, $x_\ell = 5$,
 $m_1 = 0$ eV, $\varepsilon = 0.23$



- $\mu \rightarrow e$ (Ti) (Prospect)
- $\mu \rightarrow e$ (Al) (Prospect)
- $\mu \rightarrow e$ (Au)
- $\mu \rightarrow e$ (Ti)
- $\mu \rightarrow e \gamma$
- M_W
- N_ν
- Perturbativity of Y_ν

Dynamical ISS in MFV – Case A

$$\mathcal{N} = (\mathbb{1} - \eta) U \quad \eta = \frac{v^2}{4f\Lambda^2} U \widehat{m}_\nu^{1/2} \mathcal{H}^T \mathcal{H}^* \widehat{m}_\nu^{1/2} U^\dagger$$

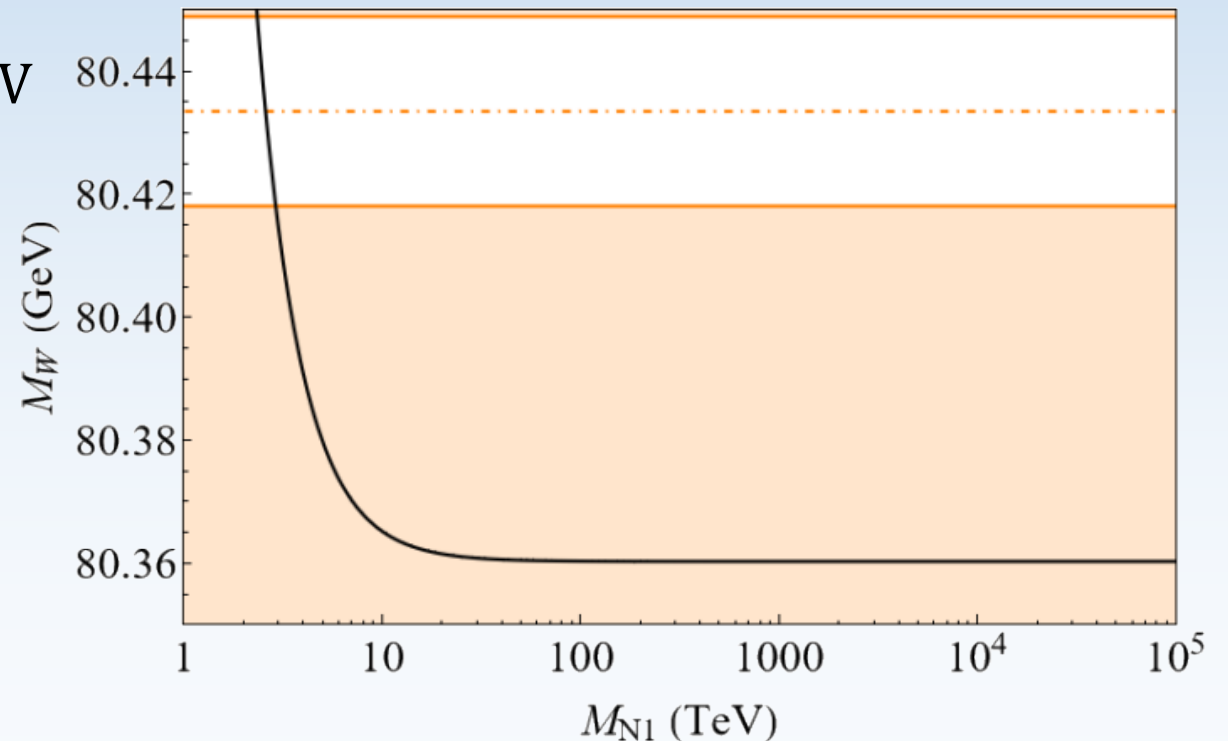


Impact of the CDF Measurement of M_W

- New very precise measurement of $M_W = 80,4335(94)$ GeV

CDF Collaboration, Science 376 (2022), no. 6589 170-176

- Cases A and B not compatible with the new value
- Case C is not constrained by LFV processes
- Sharp prediction for $M_{N1} \in [2.4, 2.9]$ TeV



The MFVA – Setup

- The axion arises as the angular part of Φ

$$\Phi = \frac{\rho + v_\Phi}{\sqrt{2}} e^{ia/v_\Phi}$$

- After integrating out ρ , the axion couplings read

$$- e^{i(x_u - x_q)a/v_\Phi} \bar{q}_L \tilde{H} Y_u u_R - e^{i(x_d - x_q)a/v_\Phi} \bar{q}_L H Y_d d_R - e^{i(x_e - x_l)a/v_\Phi} \bar{l}_L H Y_e e_R$$



$$c_{a\psi} \frac{\partial_\mu a}{2v_\Phi} \bar{\psi} \gamma^\mu \gamma_5 \psi$$

$$c_{au} = x_q - x_u$$

$$c_{ad} = x_q - x_d$$

$$c_{ae} = x_l - x_e$$

$$c_{agg} = 3(c_{au} + c_{ad}), \quad c_{aWW} = \frac{3}{2s_\theta^2} (3(c_{au} + c_{ad}) + c_{ae}),$$

$$c_{aZZ} = \frac{t_\theta^2}{4} (17c_{au} + 5c_{ad} + 15c_{ae}) + \frac{3}{4t_\theta^2} (3(c_{au} + c_{ad}) + c_{ae}),$$

$$c_{a\gamma Z} = \frac{t_\theta}{4} (17c_{au} + 5c_{ad} + 15c_{ae}) - \frac{3}{4t_\theta} (3(c_{au} + c_{ad}) + c_{ae}),$$

$$c_{a\gamma\gamma} = 2(4c_{au} + c_{ad} + 3c_{ae})$$

$$c_{agg} \neq 0$$

$$\frac{c_{agg}}{c_{a\gamma\gamma}} = 8/3$$

The MFVA – Phenomenology

$$f_a = \frac{v_\Phi}{c_{agg}}$$

	x_l	x_e	c_{au}	c_{ad}	c_{ae}	c_{agg}	$c_{a\gamma\gamma}$	c_{aZZ}	$c_{a\gamma Z}$	c_{aWW}
$S0$	0	3	0	-3	-3	-9	-24	-35.8	8.8	-81
$S1$	1	4	0	-3	-3	-9	-24	-35.8	8.8	-81

Jaeckel and Spannowsky, 1509.00476; Bauer et al., 1708.00443

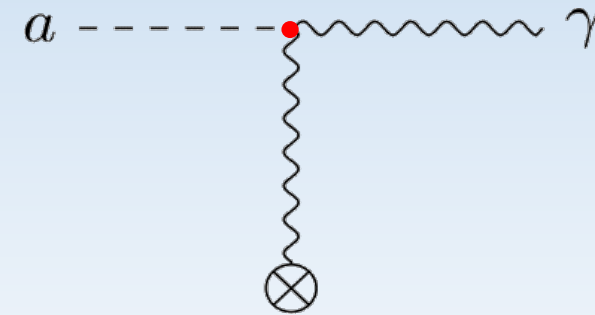
- Astrophysical and cosmological bounds on photon coupling

$$f_a \gtrsim 1.2 \times 10^7 \text{ GeV} \quad \text{for} \quad m_a \lesssim 10 \text{ meV},$$

$$f_a \gtrsim 8.7 \times 10^6 \text{ GeV} \quad \text{for} \quad 10 \text{ meV} \lesssim m_a \lesssim 10 \text{ eV},$$

$$f_a \gg 8.7 \times 10^8 \text{ GeV} \quad \text{for} \quad 10 \text{ eV} \lesssim m_a \lesssim 0.1 \text{ GeV},$$

$$f_a \gtrsim 3 \text{ GeV} \quad \text{for} \quad 0.1 \text{ GeV} \lesssim m_a \lesssim 1 \text{ TeV}$$

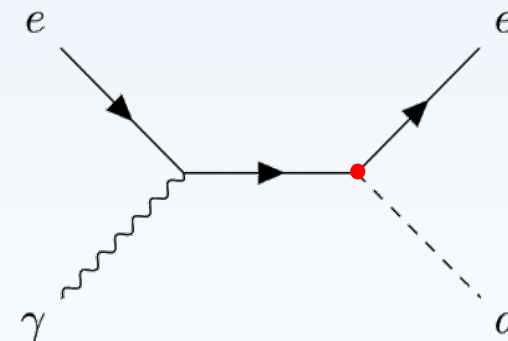


- Astrophysical bounds on electron coupling

Borexino Collaboration, Bellini et al., 1203.6258; Armengaud et al. 1307.1488;
Viaux et al., 1311.1669

$$f_a \gtrsim 3.9 \times 10^8 \text{ GeV} \quad \text{for} \quad m_a \lesssim 1 \text{ eV},$$

$$f_a \gtrsim 6.4 \times 10^6 \text{ GeV} \quad \text{for} \quad 1 \text{ eV} \lesssim m_a \lesssim 10 \text{ MeV}$$



The MFVA – Phenomenology

$$f_a = \frac{v_\Phi}{c_{agg}}$$

	x_l	x_e	c_{au}	c_{ad}	c_{ae}	c_{agg}	$c_{a\gamma\gamma}$	c_{aZZ}	$c_{a\gamma Z}$	c_{aWW}
$S0$	0	3	0	-3	-3	-9	-24	-35.8	8.8	-81
$S1$	1	4	0	-3	-3	-9	-24	-35.8	8.8	-81

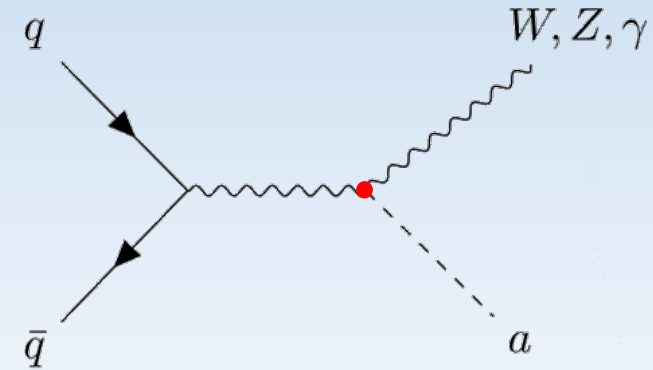
Brivio et al., 1701.05379

- Collider bounds on massive gauge bosons couplings ($0,1 \text{ GeV} \lesssim m_a \lesssim 1 \text{ GeV}$)

$$(aWW) \quad f_a \gtrsim 6.4 \text{ GeV}$$

$$(aZZ) \quad f_a \gtrsim 5.7 \text{ GeV}$$

$$(aZ\gamma) \quad f_a \gtrsim 17.8 \text{ GeV}$$

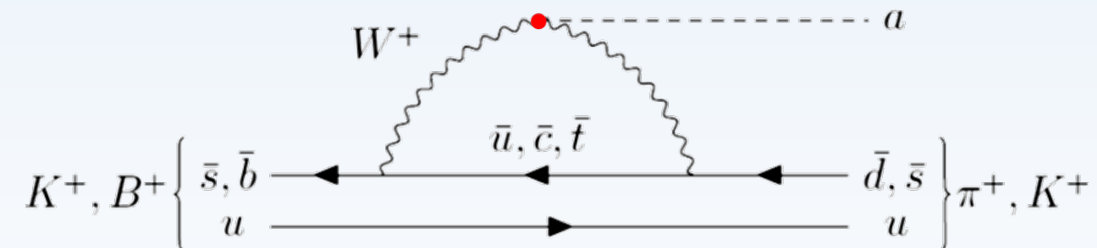


- Flavour bounds on aWW coupling

Izaguirre et al., 1611.09355

$$f_a \gtrsim 3.5 \times 10^3 \text{ GeV} \quad \text{for} \quad m_a \lesssim 0.2 \text{ GeV}$$

$$f_a \gtrsim 105 \text{ GeV} \quad \text{for} \quad 0.2 \text{ GeV} \lesssim m_a \lesssim 5 \text{ GeV}$$



The MFVA – Phenomenology

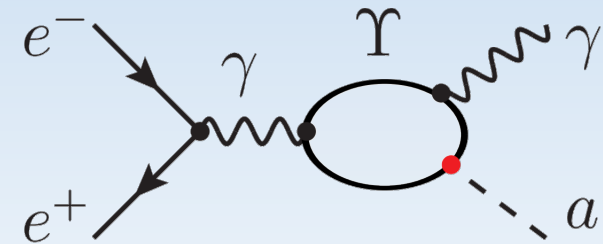
$$f_a = \frac{v_\Phi}{c_{agg}}$$

	x_l	x_e	c_{au}	c_{ad}	c_{ae}	c_{agg}	$c_{a\gamma\gamma}$	c_{aZZ}	$c_{a\gamma Z}$	c_{aWW}
$S0$	0	3	0	-3	-3	-9	-24	-35.8	8.8	-81
$S1$	1	4	0	-3	-3	-9	-24	-35.8	8.8	-81

- Flavour bound on bottom coupling through $\Upsilon \rightarrow a\gamma$ ($m_a \sim 1$ GeV)

Merlo et al., 1905.03259

$$f_a \gtrsim 830 \text{ GeV}$$



- Axion-bottom coupling bound from CLEO ($0,4 \lesssim m_a \lesssim 4,8$ GeV, decaying axion)

CLEO Collaboration, PRL 80 (1998) 1150-1155

$$f_a \gtrsim 667 \text{ GeV}$$