



University of
Zurich^{UZH}

Flavor hierarchies and B-anomalies from 5D

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Based on work with G. Isidori, J. Fuentes Martín, N. Selimović and B. A. Stefanek

[\[2203.01952\]](#)

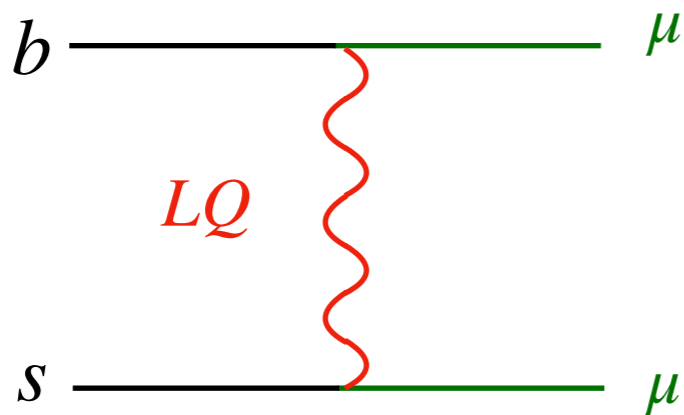
Planck - 2022

B-anomalies

Gino's talk

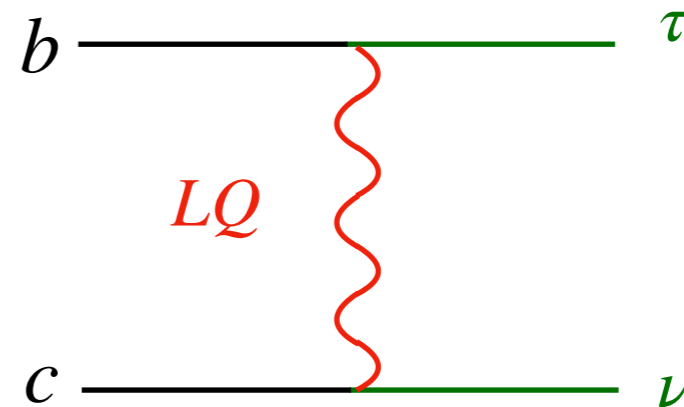
$$b \rightarrow sll$$

- Non-universality in e/μ , $> 4\sigma$



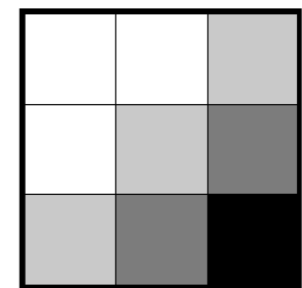
$$b \rightarrow c\tau\nu$$

- Non universality in $\tau/\mu, e$, $\sim 3\sigma$



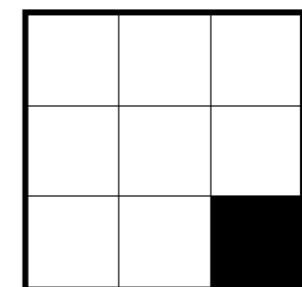
Single mediator: $U_{1\mu} \sim (3, 1, 2/3)$

$$\beta_L^{ql} \sim$$



$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \left[\beta_L^{i\alpha} (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) + \beta_R^{i\alpha} (\bar{d}_R^i \gamma_\mu e_R^\alpha) \right] + \text{h.c.}$$

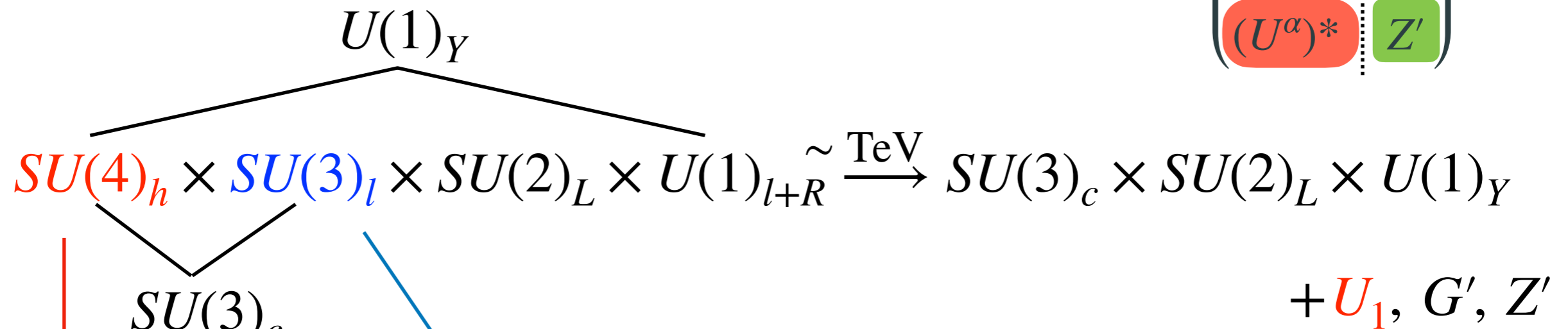
$$\beta_R^{ql} \sim$$



4321 model

4321 completion

$$SU(4) \sim \begin{pmatrix} G^a & U^\alpha \\ (U^\alpha)^* & Z' \end{pmatrix}$$



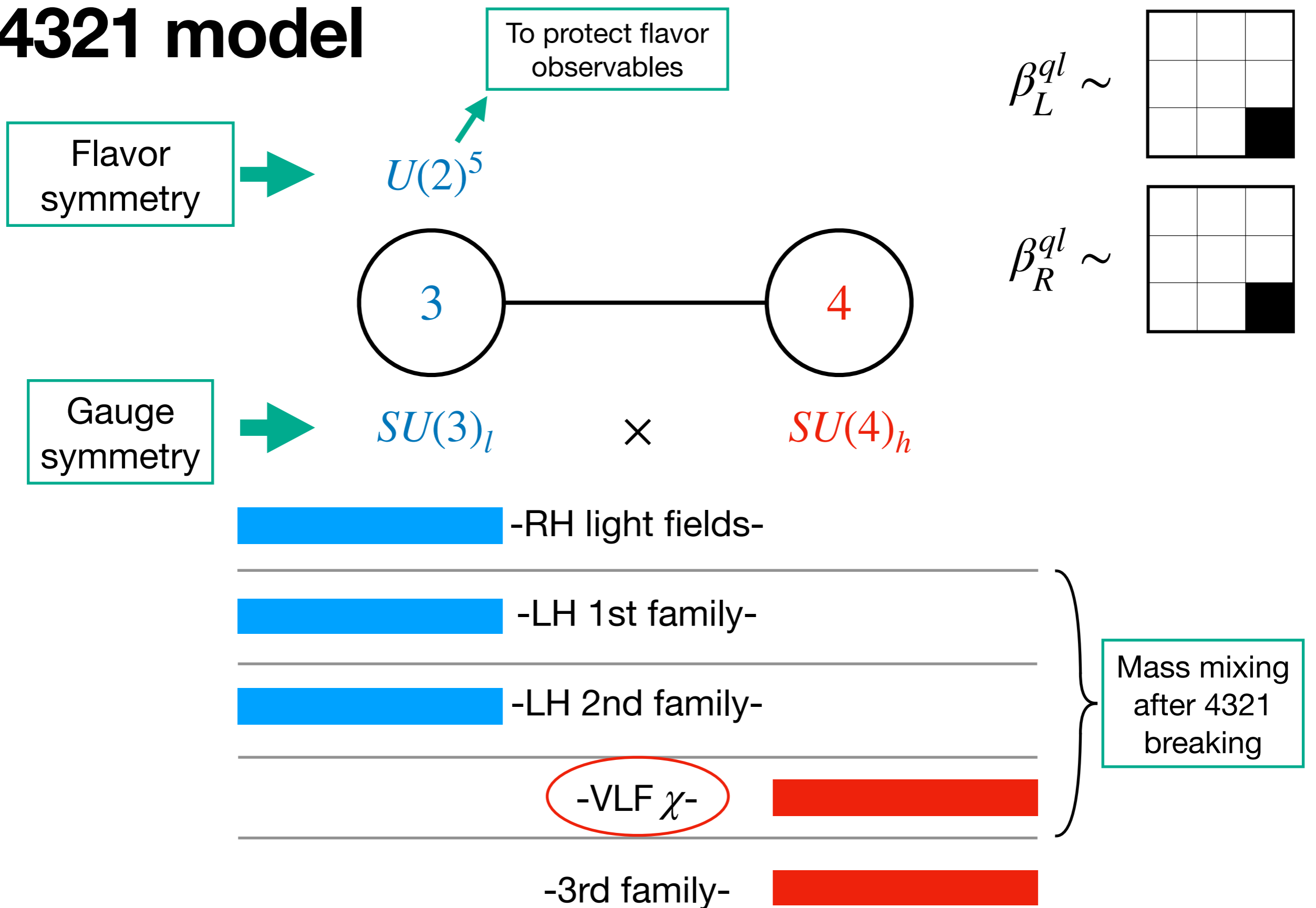
3rd family
Pati-Salam
unification

$$\Psi_{L/R}^{(3)} = \begin{pmatrix} Q_{L,R}^1 \\ Q_{L,R}^2 \\ Q_{L,R}^3 \\ L_{L,R} \end{pmatrix}$$

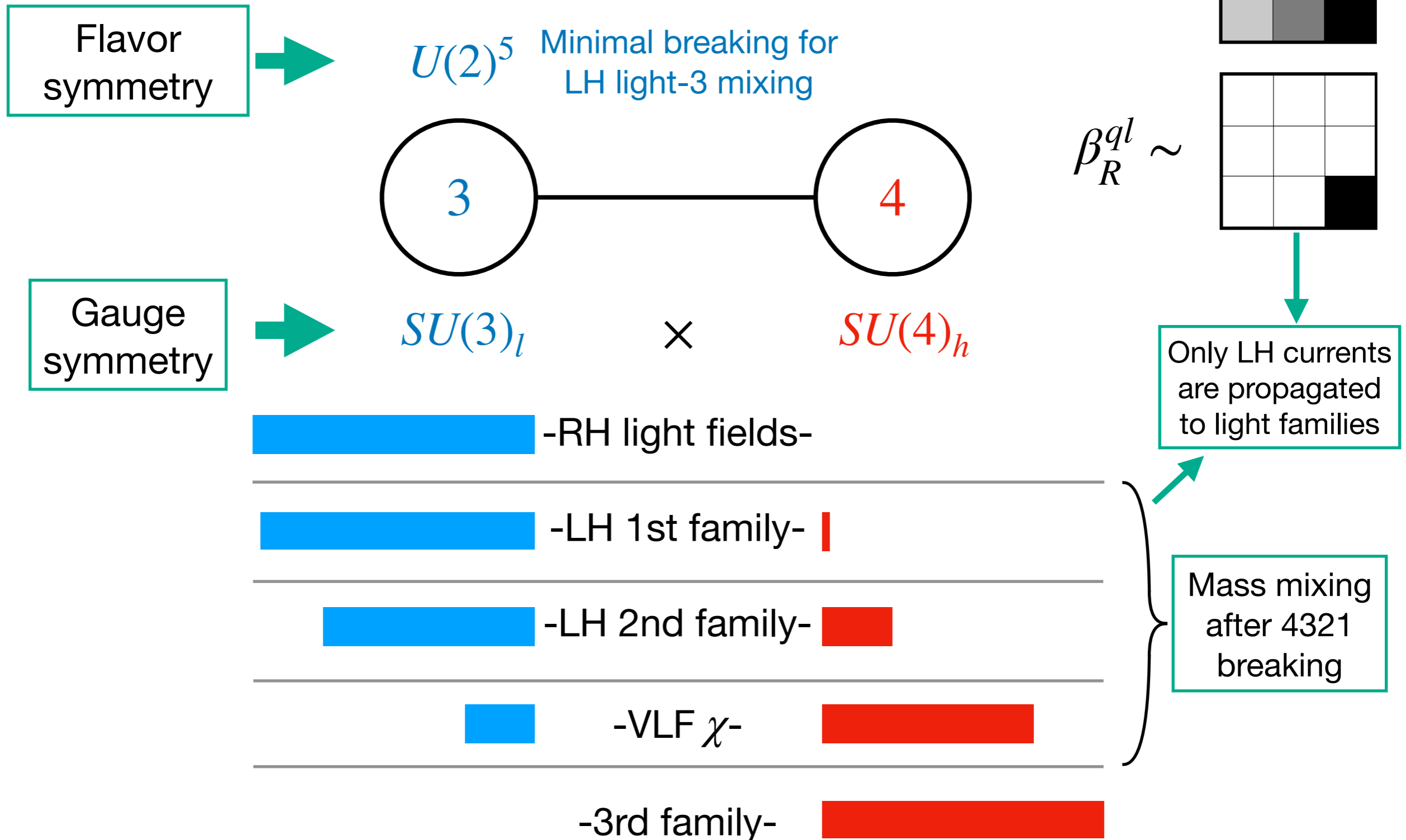
Light families

Flavor non-universal gauge theory

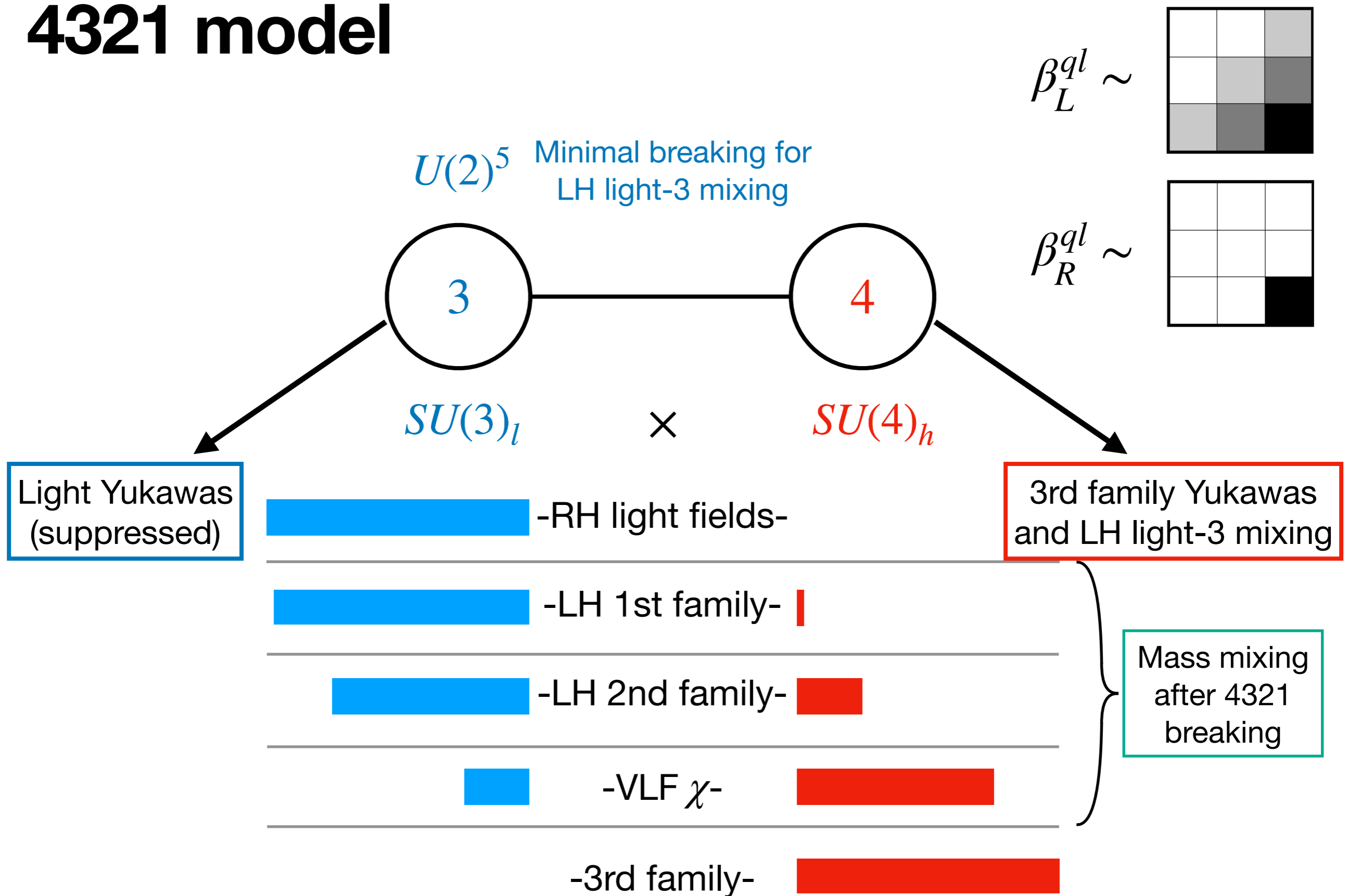
4321 model



4321 model



4321 model



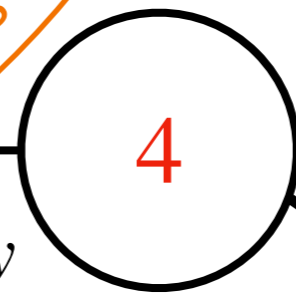
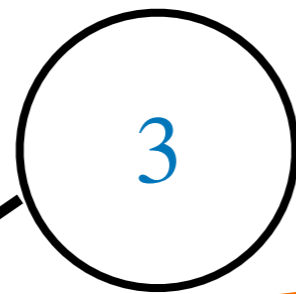
4321 model

Connection with Higgs stabilization?

Localization of fermion families in an extra dimension?

$U(2)^5$ Minimal breaking for LH light-3 mixing

Higgs



y

$SU(3)_l$

\times

$SU(4)_h$

Light Yukawas (suppressed)

-RH light fields-

3rd family Yukawas and LH light-3 mixing

-LH 1st family-

-LH 2nd family-

-VLF χ -

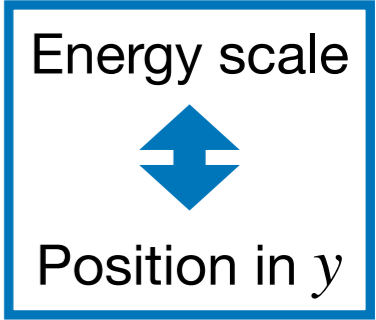
-3rd family-

Mass mixing after 4321 breaking

A 5D model: a first attempt

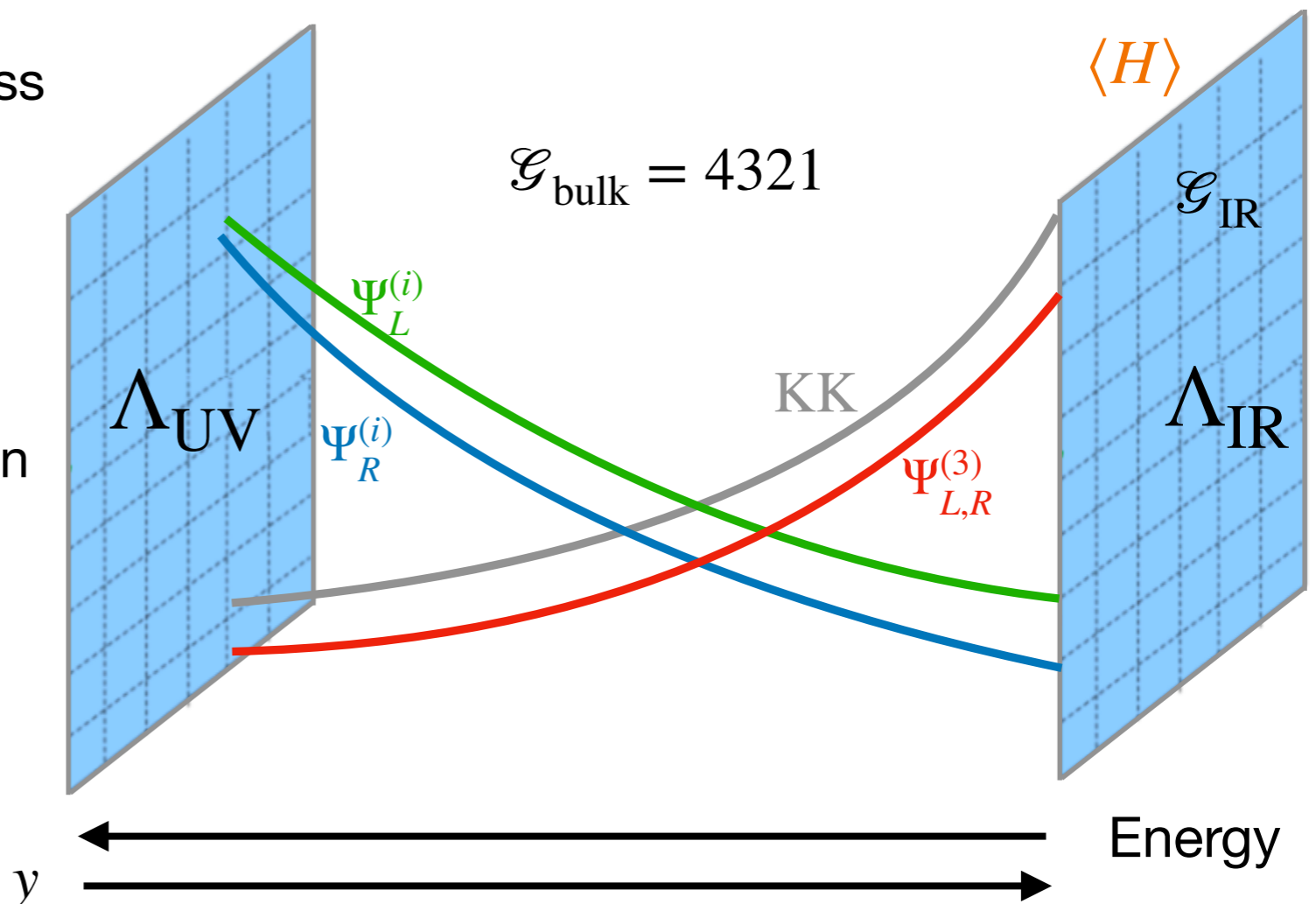
Curvature of the AdS slice

- Warped 5D geometry (RS): $ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2$
[Randall, Sundrum, [hep-ph/9905221](https://arxiv.org/abs/hep-ph/9905221)]
- Holography \Rightarrow Dual to a strongly coupled sector $\mathcal{G}_{\text{bulk}} \rightarrow \mathcal{G}_{\text{IR}}$
- The strong dynamics can be used to break 4321 [Fuentes-Martin, Stangl [2004.11376](https://arxiv.org/abs/2004.11376)]



- Anarchic partial compositeness paradigm in RS
- Emerging $U(2)$ symmetry at the TeV scale...
- But in principle, broken both in the LH and RH light sectors

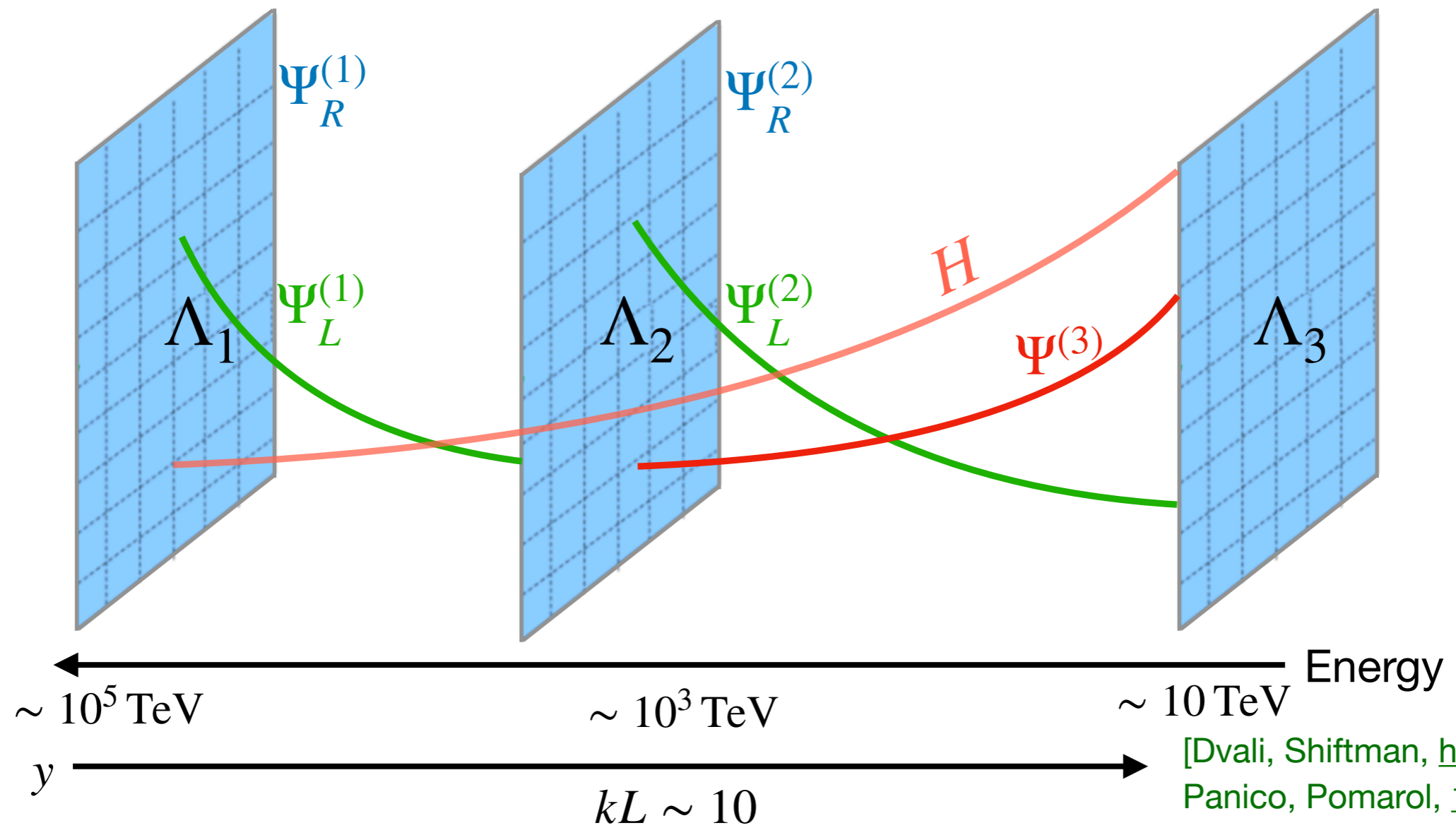
$$\mathcal{L} \supset \bar{\Psi}_L^{(3)} H \Psi_R^{(1,2)}$$



[Gherghetta, Pomarol, [arXiv:hep-ph/0003129](https://arxiv.org/abs/hep-ph/0003129)]

A multiscale 5D model

- Multi-brane construction: flavor hierarchies from different scales.
- \Rightarrow Emerging $U(2)$ symmetry minimally broken.



A 5D model that...

- Reduces to 4321 below the KK scale
- Explains flavour hierarchies from a multi-scale origin
- Realises the Higgs as a pNGB

Fuentes-Martin, Isidori, JML, Selimovic, Stefanek, [2203.01952](#)]

Gauge sector

A_5 of broken generators
dual to NGBs

Quark-lepton unification of light families

$$\mathcal{G}_{\text{bulk}}^{12} = SU(4)_h \times SU(4)_l \times SO(5)$$

↓ Λ_2 (6 broken)

$$\mathcal{G}_{\text{bulk}}^{23} = SU(4)_h \times SU(3)_l \times U(1)_l \times SO(5)$$

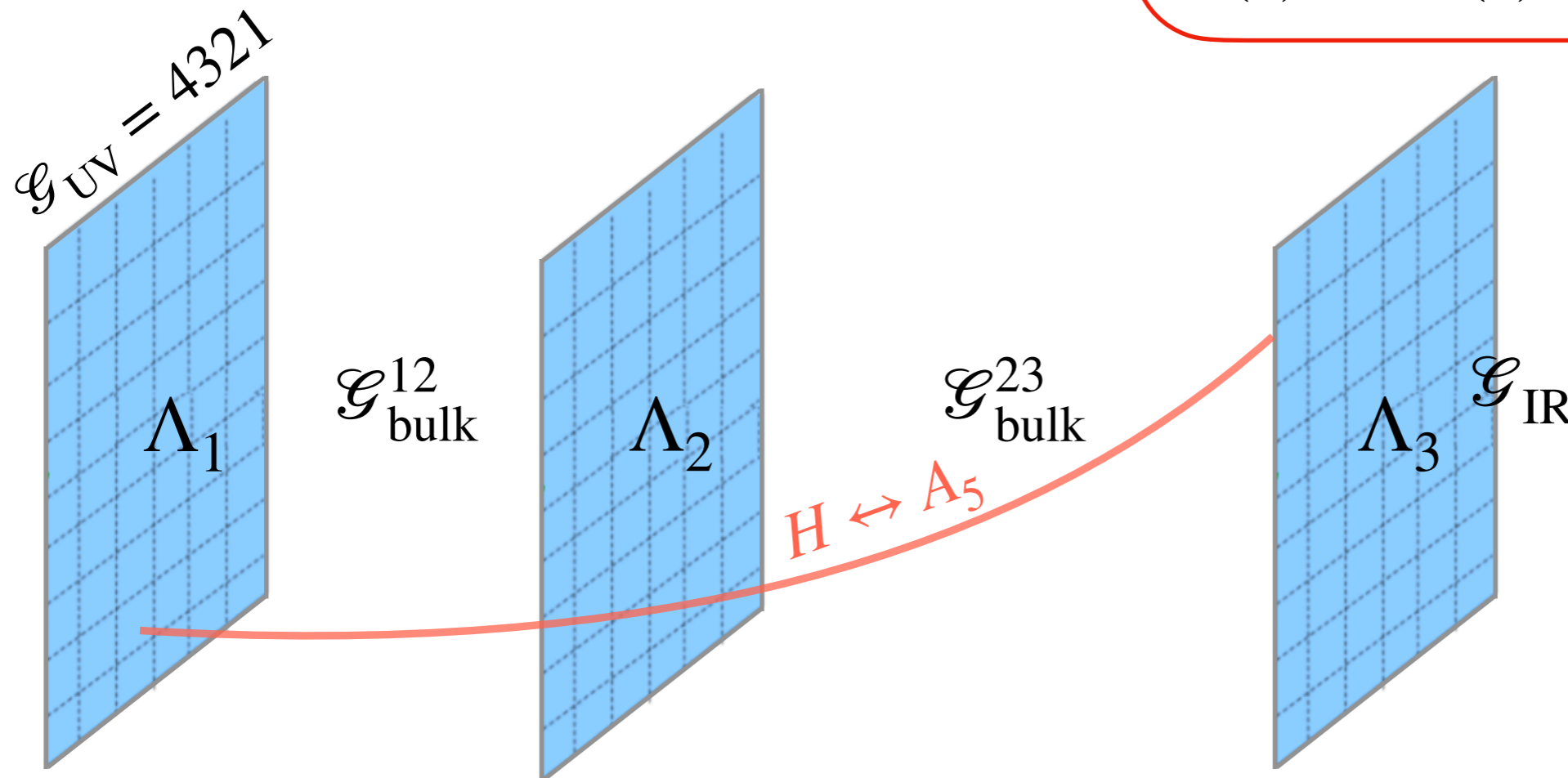
↓ $\Lambda_3 = \Lambda_{\text{IR}}$ (15 + 4 broken)

$$\mathcal{G}_{\text{IR}} = SU(3)_c \times U(1)_{B-L} \times SO(4)$$

15 eaten by $U_1, G', Z', M_{4321} \sim \frac{M_{\text{KK}}}{\sqrt{2kL}}$

4 as pNGB Higgs

$$SO(5) \rightarrow SO(4) = SU(2)_L \times SU(2)_R$$



Minimal composite
Higgs (MCHM)

[Agashe, Contino, Pomarol,
[hep-ph/0412089](https://arxiv.org/abs/hep-ph/0412089)]

Fermion and scalar sector

| | Field | $SU(4)_h$ | $SU(4)_l$ | $SO(5)$ | |
|---------------------|---------------------------------------|-----------|-----------|---------|----------|
| SM fermions and VLF | Ψ^3 Ψ_d^3 $\chi^{(\prime)}$ | 4 | 1 | 4 | Fermions |
| | Ψ^j $\Psi_{u,d}^j$ | 1 | 4 | 4 | |
| | \mathcal{S}^i | 1 | 1 | 1 | |
| For neutrinos | Φ | 1 | 1 | 1 | Scalars |
| | Ω | 1 | 4 | 4 | |
| For light Yukawas | Σ | 1 | 1 | 5 | |

[Fuentes-Martin, Isidori, Pages, Stefaneke, [2012.10492](#)]

Top Yukawa

| Field | $SU(4)_h$ | $SU(4)_l$ | $SO(5)$ |
|----------|-----------|-----------|---------|
| Ψ^3 | 4 | 1 | 4 |

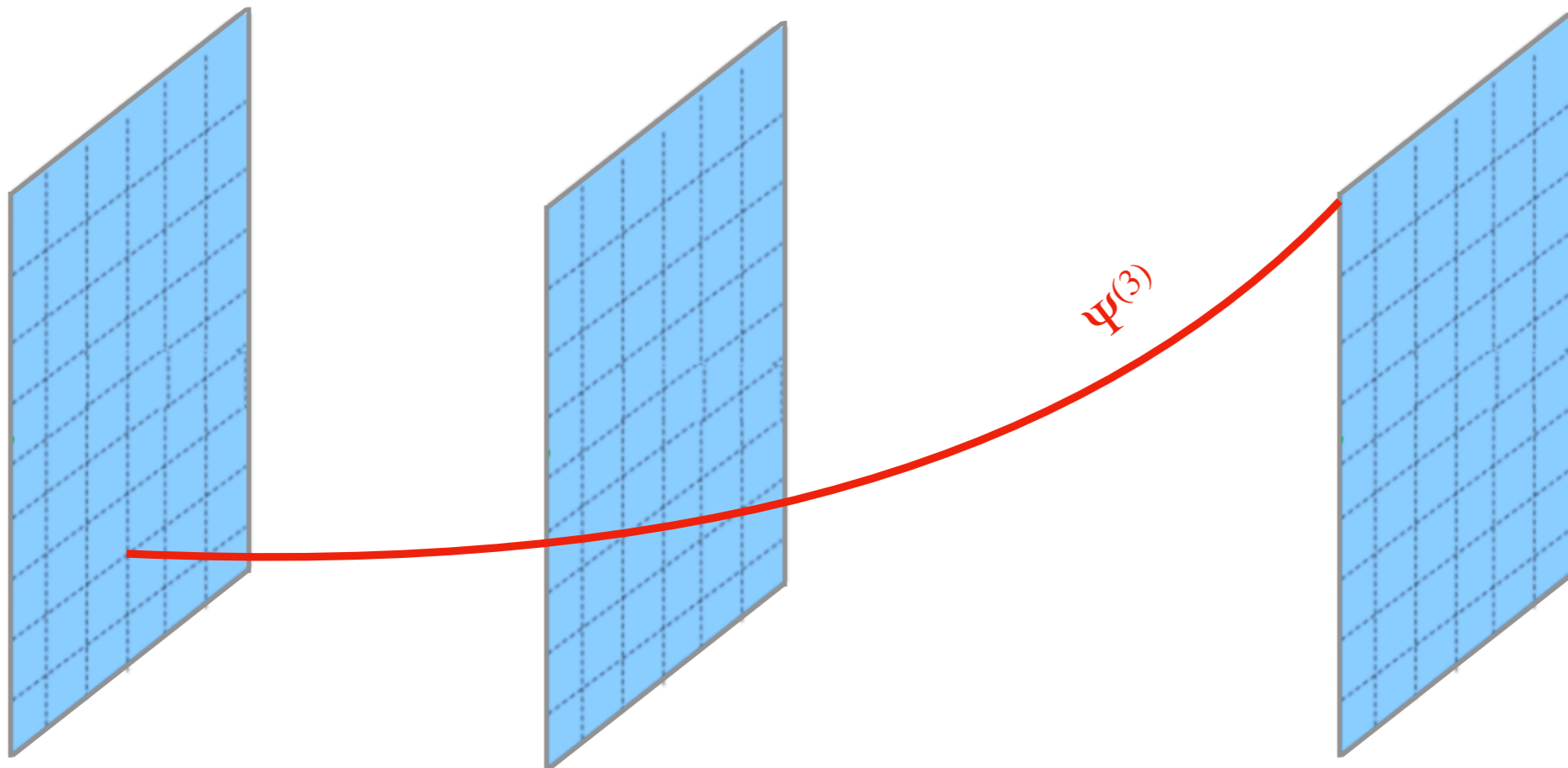
$$\Psi^3 = \left[\begin{array}{c} q_L \\ t_R \\ \times \end{array} \right] \left. \begin{array}{l} SU(2)_L \\ \\ SU(2)_R \end{array} \right\}$$

Top Yukawa from $\bar{\Psi}^3 A_5 \Psi^3$ coupling in the bulk

$$y_t = \frac{g_*}{2\sqrt{2}} P(M_{\Psi^3})$$

$$(g_*^2 = g_5^2 k)$$

For y_t : $g_* \geq 2.2$

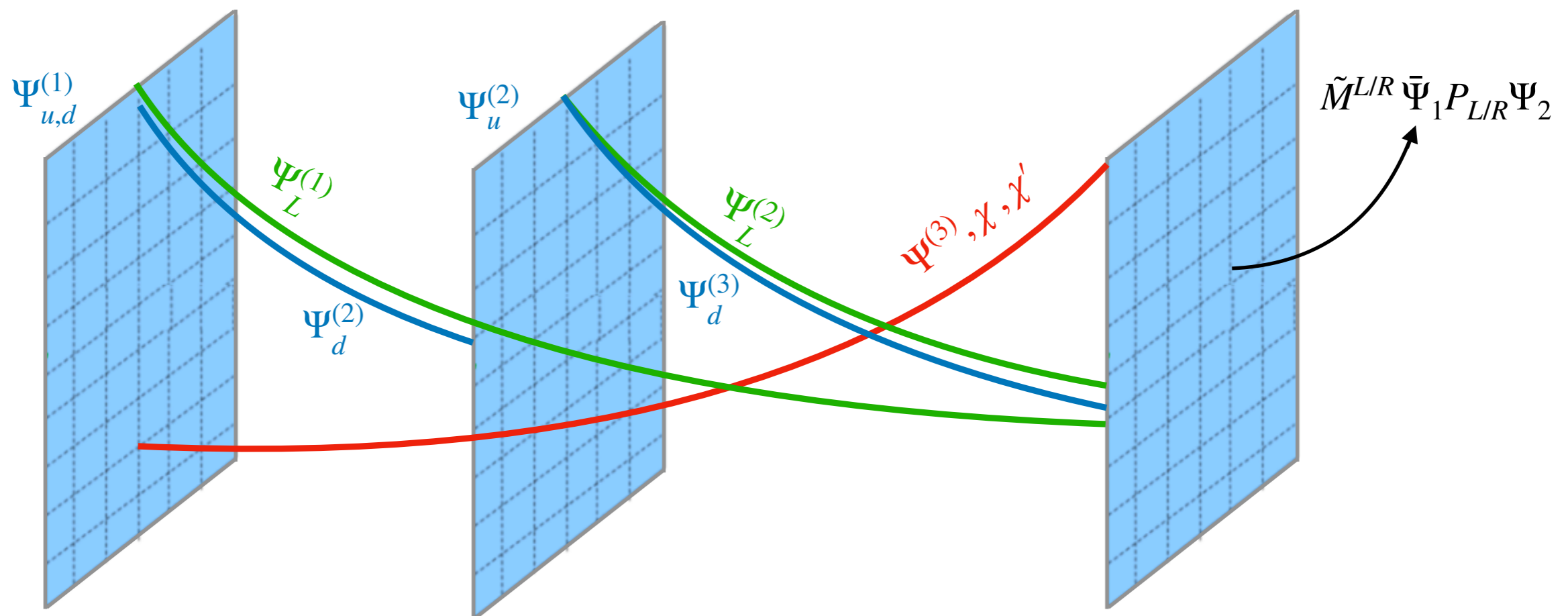


Other 3rd family Yuk. and light-heavy mixing

| Field | $SU(4)_h$ | $SU(4)_l$ | $SO(5)$ |
|---------------------------------------|-----------|-----------|---------|
| $\Psi^3, \Psi_d^3, \mathcal{X}^{(l)}$ | 4 | 1 | 4 |
| $\Psi^j, \Psi_{u,d}^j$ | 1 | 4 | 4 |

VLF mass, mass mixing of light families with VLF, and other 3rd family Yukawas from masses in the IR brane

$$y_{f_1 f_2} = \frac{g^*}{2\sqrt{2}} (\tilde{M}^L - \tilde{M}^R) \times (\text{profile suppression})$$

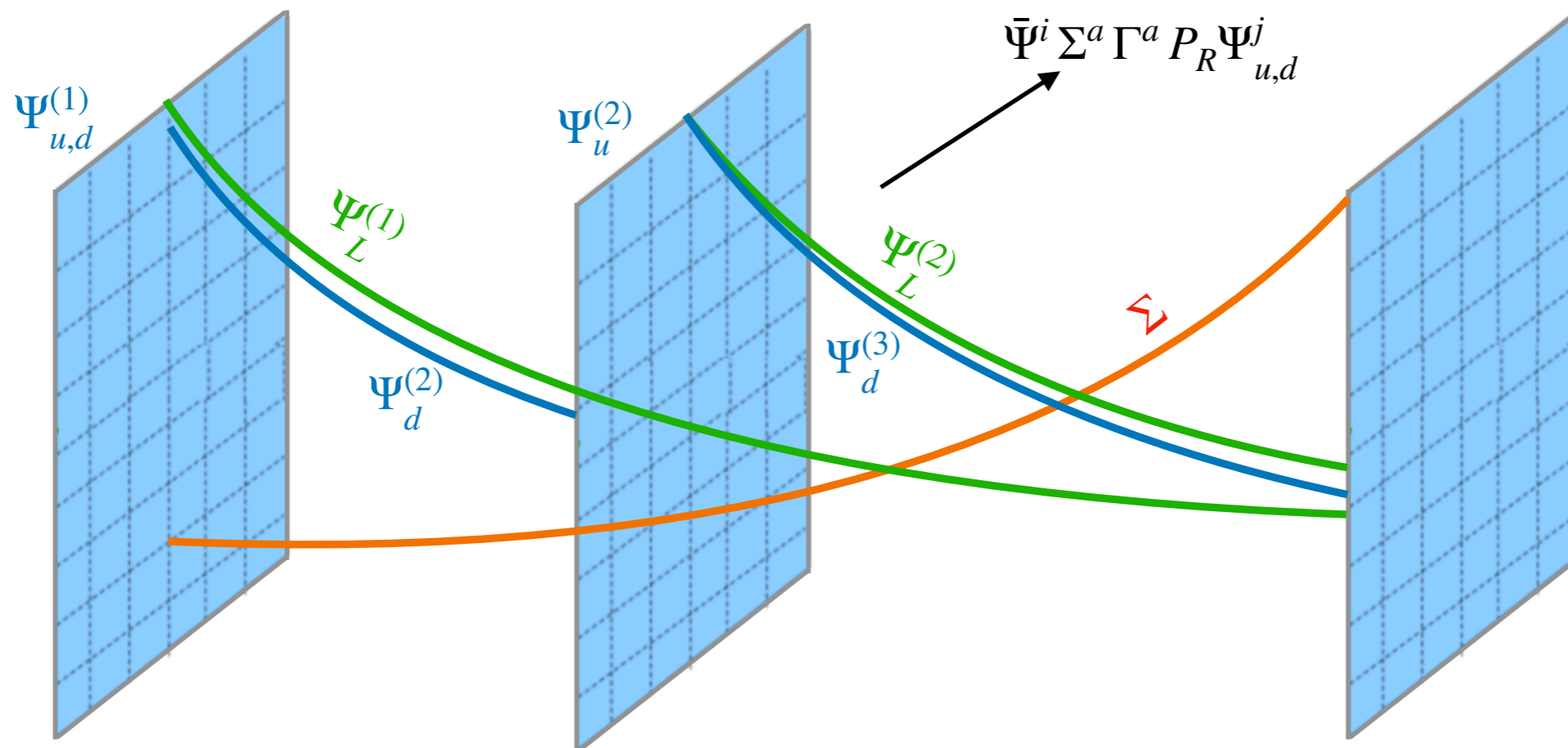


Light Yukawas

| Field | $SU(4)_h$ | $SU(4)_l$ | $SO(5)$ |
|------------------------|-----------|-----------|----------|
| $\Psi^j, \Psi_{u,d}^j$ | 1 | 4 | 4 |
| Σ | 1 | 1 | 5 |

$\Sigma^T \sim (H' \phi)$ takes a VEV along the singlet direction and propagates the breaking of $SO(5)$ into the bulk

$$y_{u,d}^{ij} = \frac{g^*}{2\sqrt{2}} \tilde{Y}_{u,d}^{ij} \frac{\langle \Sigma_{\text{IR}} \rangle}{\Lambda_{\text{IR}}} \times (\text{profile suppression})$$



Higgs potential

Higgs potential fully calculable

Contributions:

- **Tree level** from scalars with a VEV in the bulk breaking $SO(5)$: Σ , Ω
- **One loop** from top and gauge fields

Higgs decay constant:

$$V(h) \approx \alpha \cos\left(\frac{h}{f_h}\right) - \beta \sin^2\left(\frac{h}{f_h}\right)$$

\downarrow
 Ψ^3, Ω

\downarrow
 Ψ^3, Σ, W, Z

$\cos\left(\frac{\langle h \rangle}{f_h}\right) = -\frac{\alpha}{2\beta}$

$m_h^2 = \frac{2\beta \langle h \rangle^2}{f_h^4}$

$f_h = \frac{2\Lambda_{\text{IR}}}{g_*}$

All contributions of the correct order, up to some little-hierarchy tuning

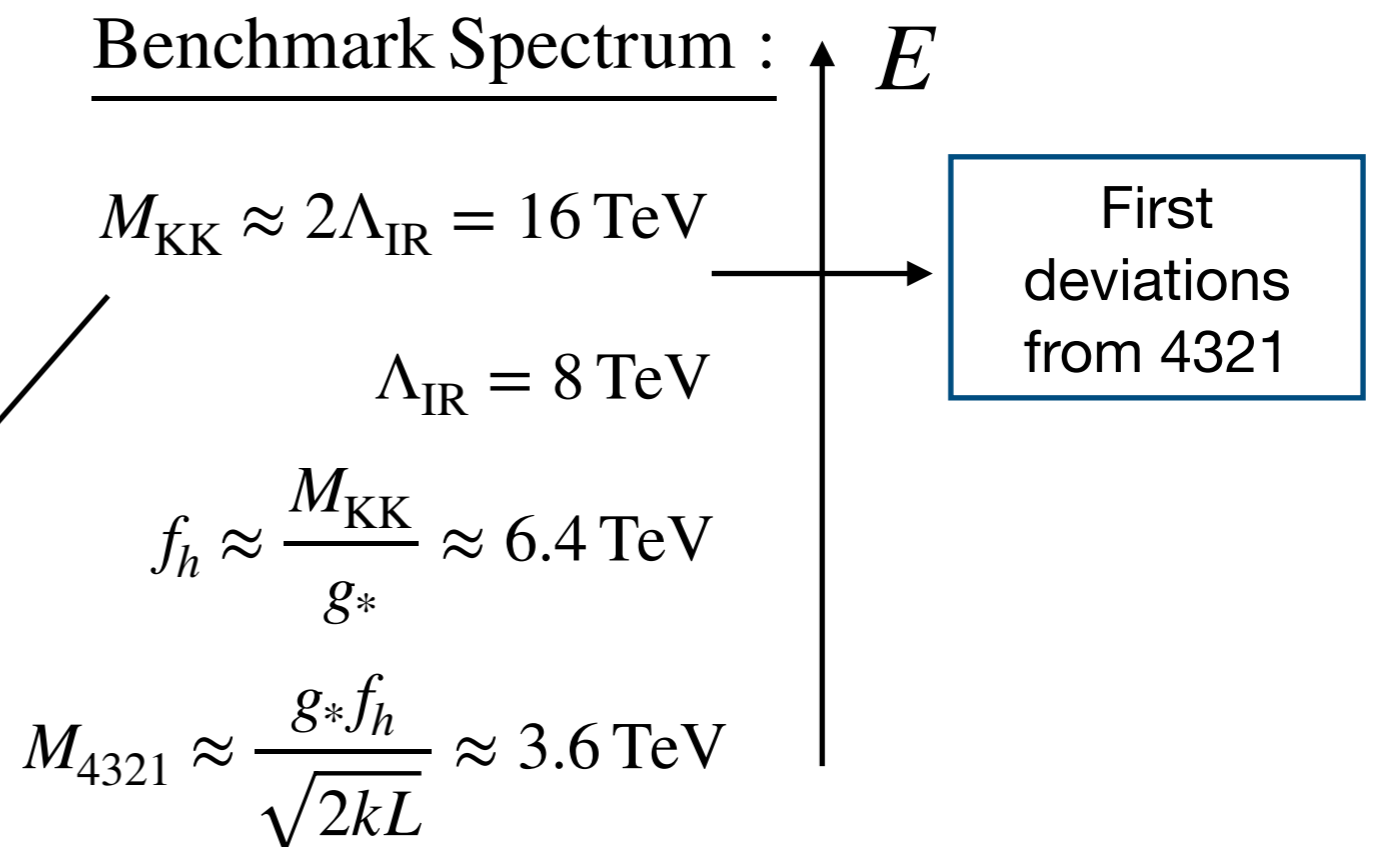
β of the right size for $g_* \approx 2.5$,
compatible with the top Yukawa

Low-energy phenomenology

- Below KK scale, same phenomenology as 4321 (B-anomalies)
- Main experimental limit coming from coloron direct searches:

$$M_{4321} \gtrsim 3.5 \text{ TeV}$$

[Cornella et al., [2103.16558](#)]



Conclusions

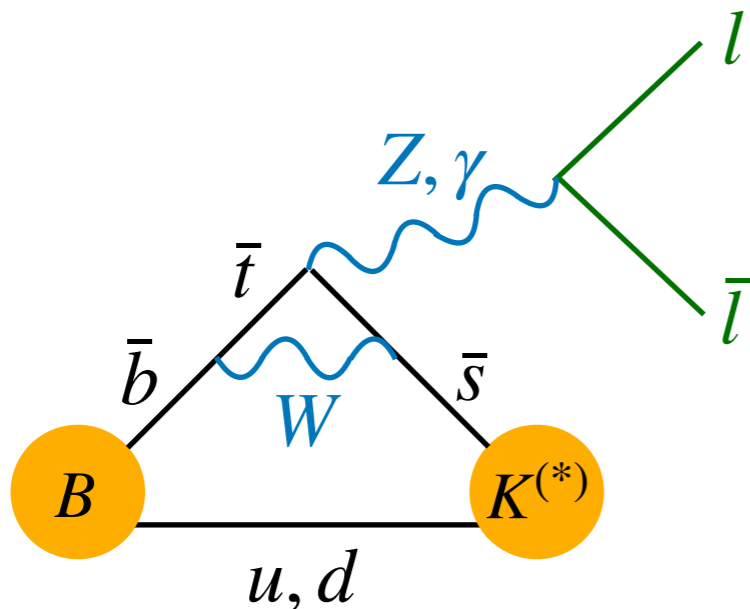
- B-anomalies hint new physics at the TeV scale.
- 4321 is a consistent UV model that can explain the B-anomalies.
- A completion of 4321 à la Randall-Sundrum points towards a multi-scale origin of the flavour hierarchies, where a $U(2)$ flavour symmetry at the TeV minimally broken can be implemented.
- We have presented a 5D model realising this idea where, in addition, the Higgs emerges as a pNGB from the same strong dynamics that breaks 4321.

Thank you!

Backup

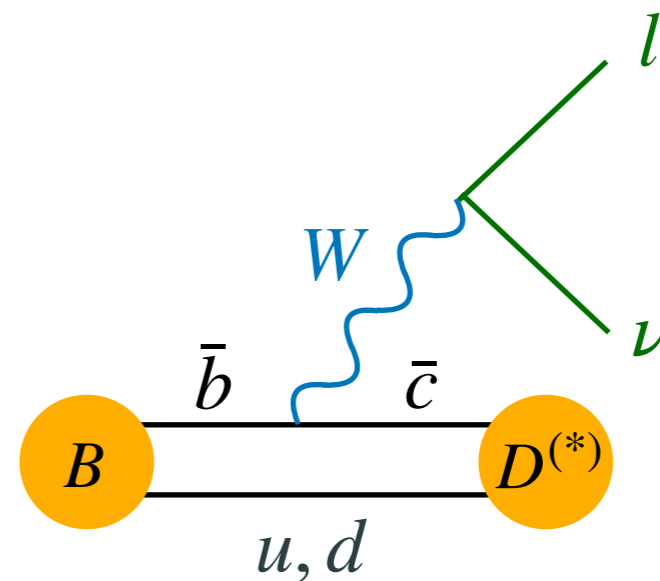
$$b \rightarrow sll$$

- $R_{K^{(*)}} = \frac{Br(B \rightarrow K^{(*)}\mu\mu)}{Br(B \rightarrow K^{(*)}ee)}$
- $B_s \rightarrow \mu\mu$
- $B \rightarrow Kll$, angular distributions, etc...
- Non-universality in e/μ , $> 4\sigma$



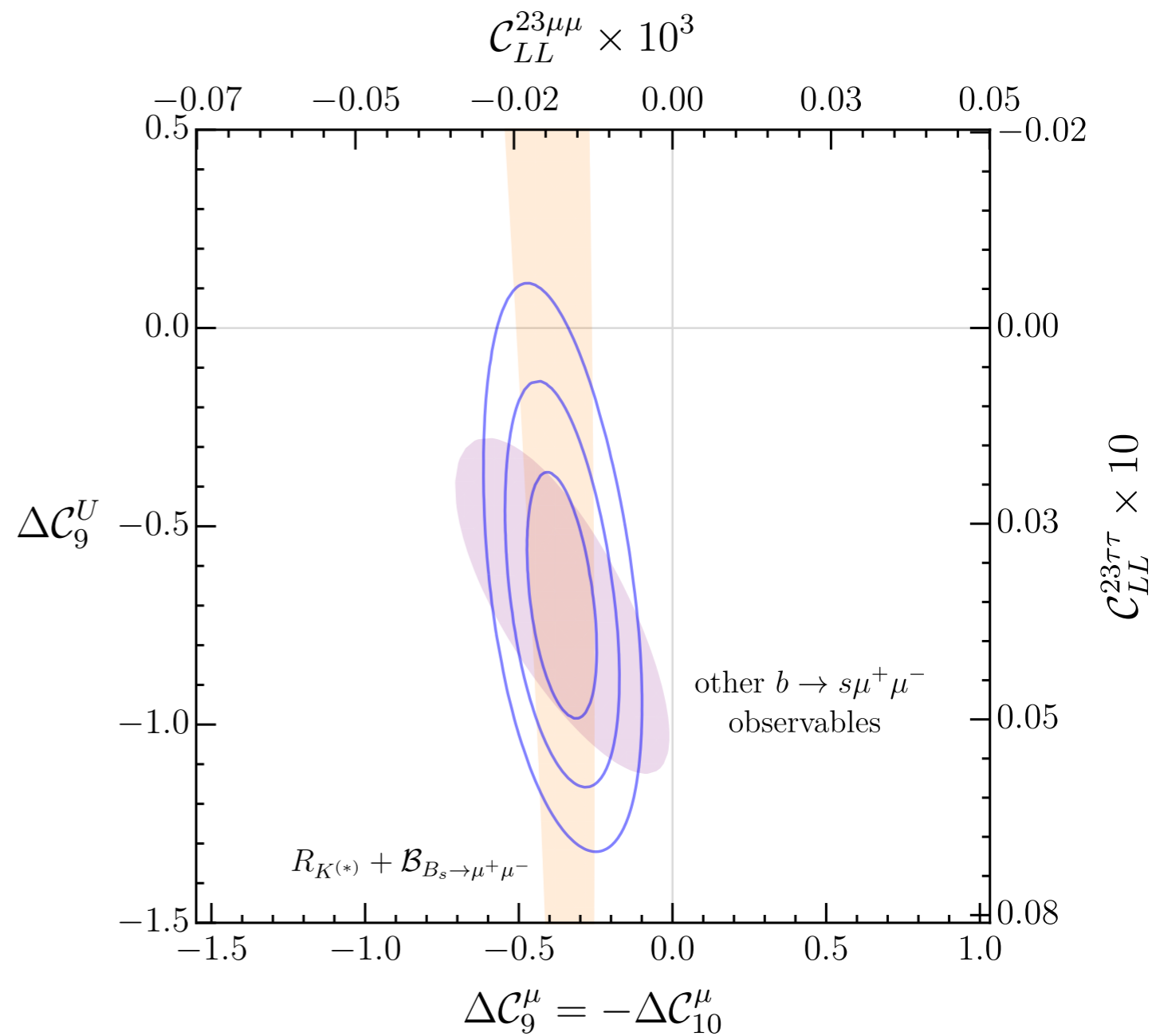
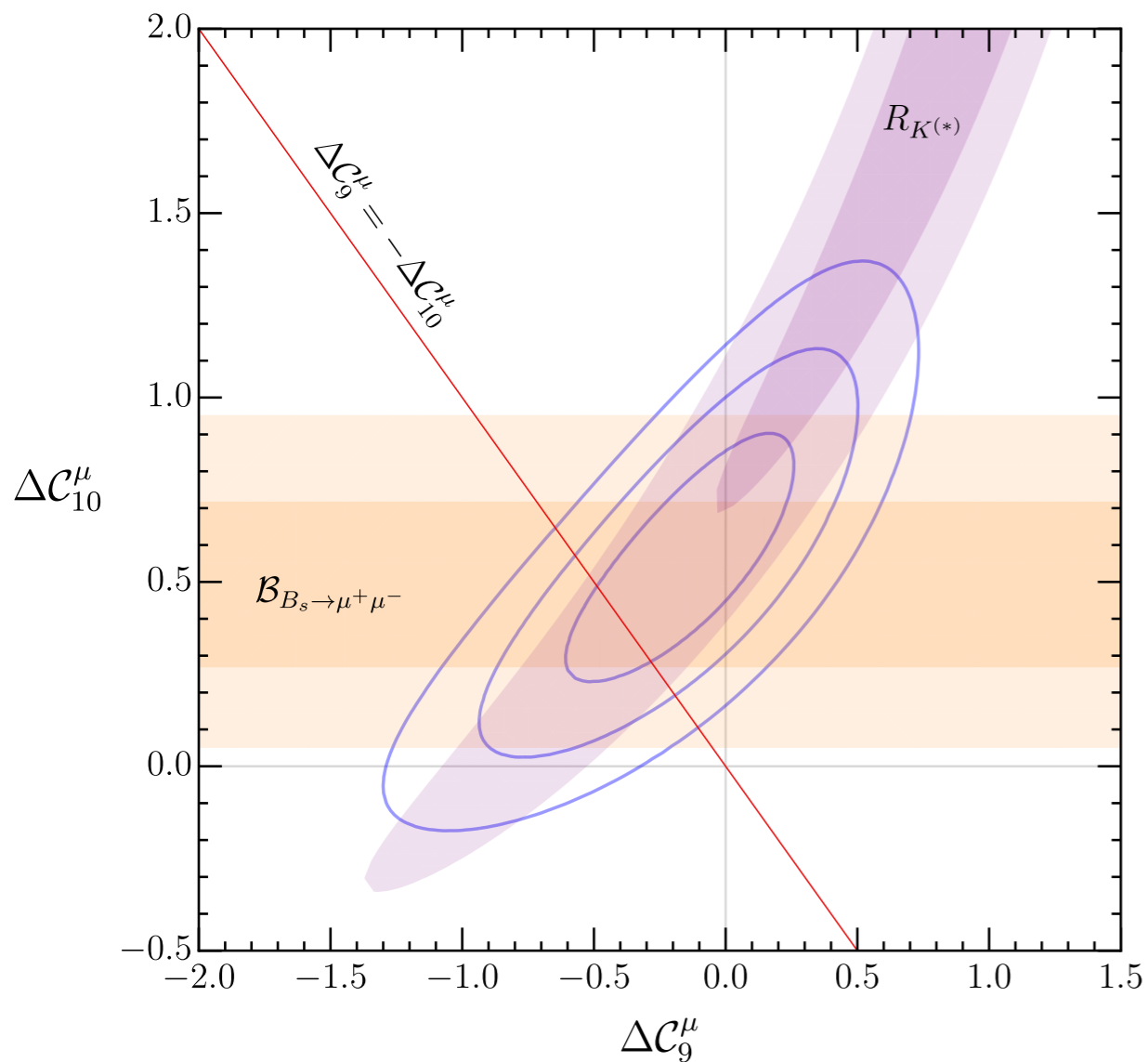
$$b \rightarrow c\tau\nu$$

- $R_{D^{(*)}} = \frac{Br(B \rightarrow D^{(*)}\tau\nu)}{Br(B \rightarrow D^{(*)}l\nu)}$
- Non universality in $\tau/\mu, e$, $\sim 3\sigma$



Backup

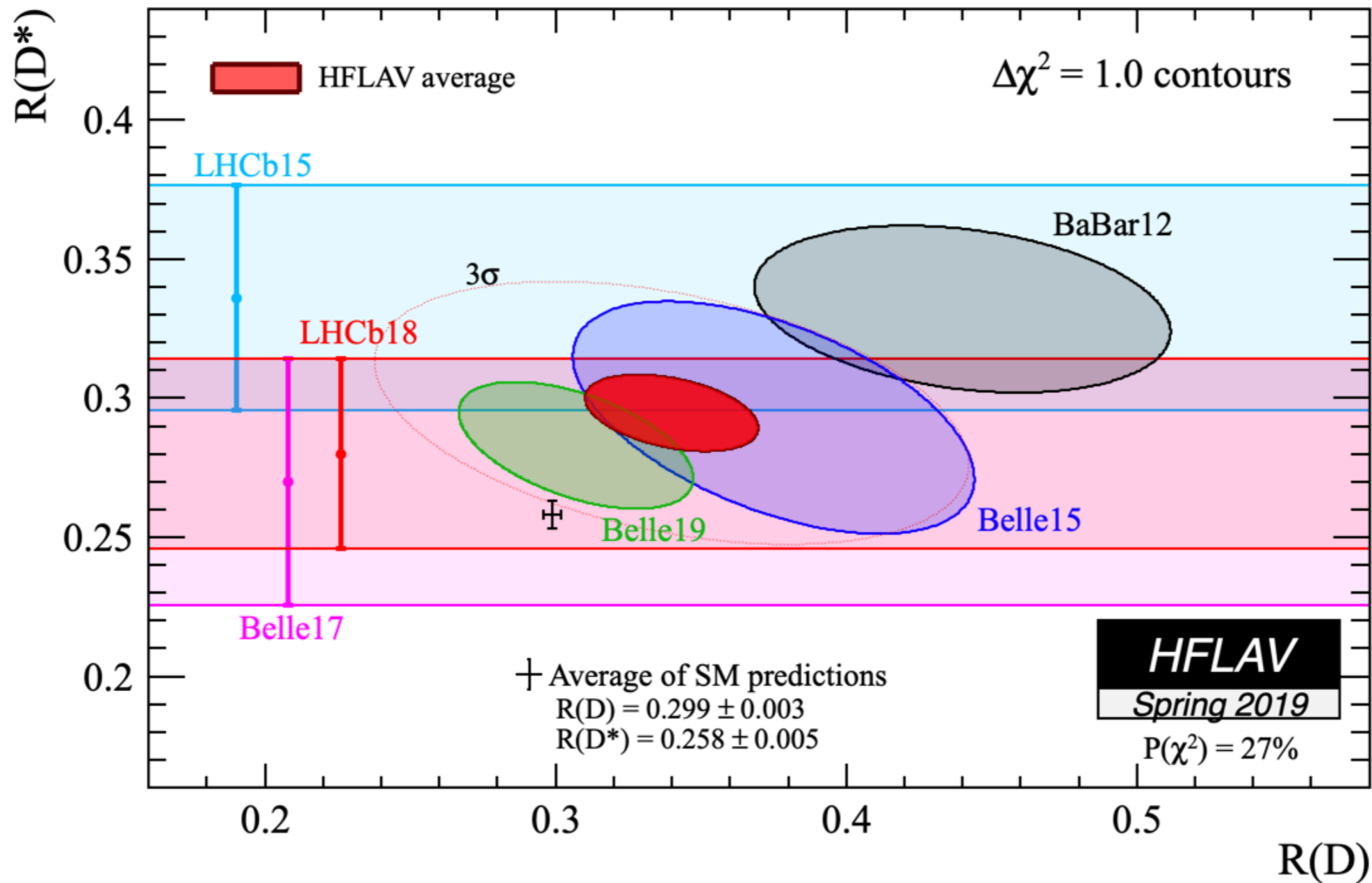
$b \rightarrow sll$



[Cornella et al., [2103.16558](#)]

Backup

$$b \rightarrow c\tau\nu$$



Backup

| Field | $SU(4)_h$ | $SU(4)_l$ | $SO(5)$ |
|---------------------------------------|-----------|-----------|---------|
| $\Psi^3, \Psi_d^3, \mathcal{X}^{(1)}$ | 4 | 1 | 4 |
| $\Psi^j, \Psi_{u,d}^j$ | 1 | 4 | 4 |
| \mathcal{S}^i | 1 | 1 | 1 |
| Σ | 1 | 1 | 5 |
| Ω | 1 | 4 | 4 |
| Φ | 1 | 1 | 1 |

$$\Psi^3 = \begin{bmatrix} \psi^3 (+, +) \\ \psi_u^3 (-, -) \\ \tilde{\psi}_d^3 (+, -) \end{bmatrix},$$

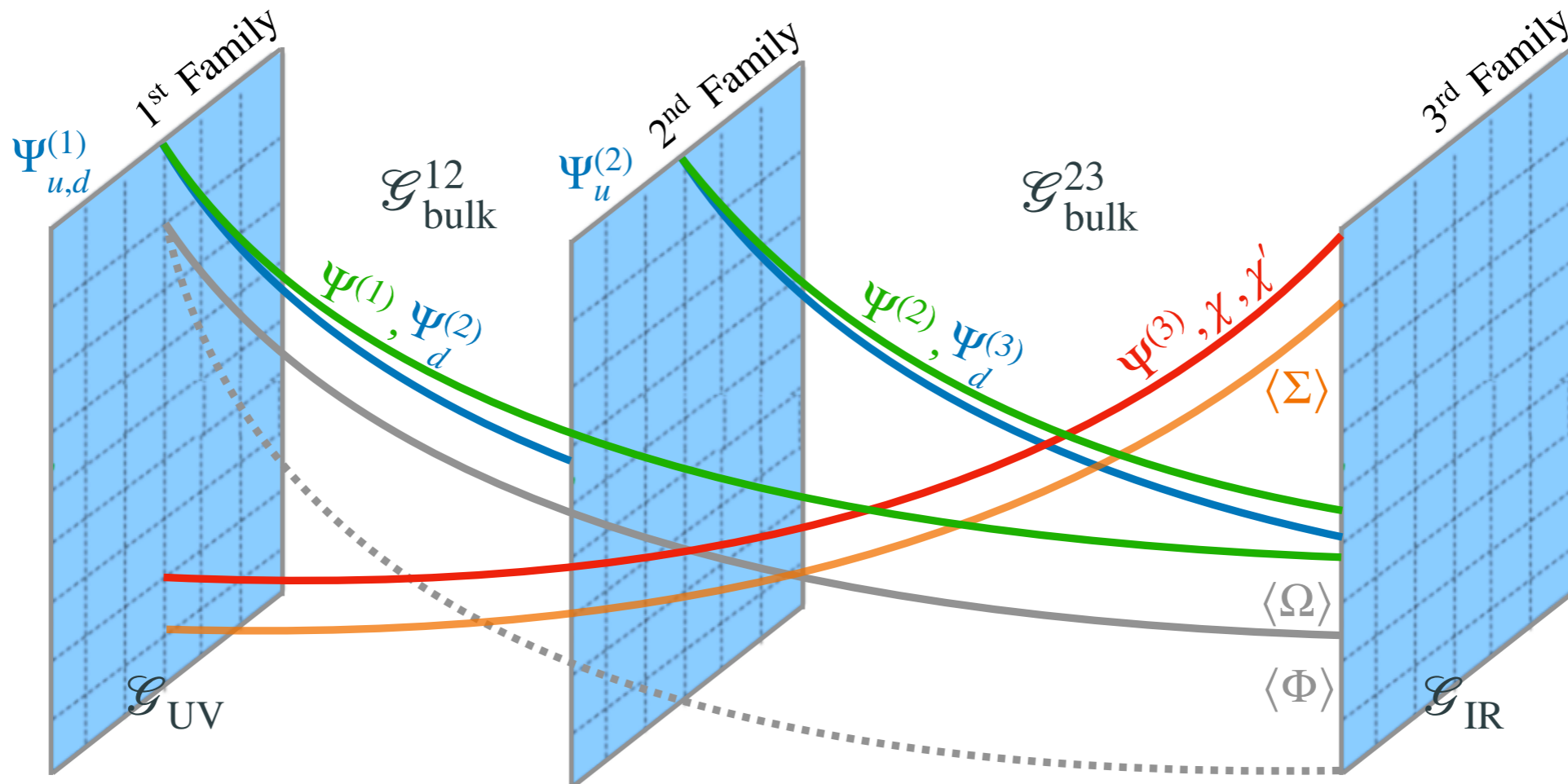
$$\Psi_d^3 = \begin{bmatrix} \tilde{\psi}^3 (+, -) \\ \tilde{\psi}_u^3 (+, -) \\ \psi_d^3 (-, -) \end{bmatrix},$$

$$\mathcal{X}^{(j)} = \begin{bmatrix} \chi^{(j)} (\pm, \pm) \\ \chi_u^{(j)} (\mp, \pm) \\ \chi_d^{(j)} (\mp, \pm) \end{bmatrix},$$

$$\Psi^j = \begin{bmatrix} \psi^j (+, +) \\ \tilde{\psi}_u^j (-, +) \\ \tilde{\psi}_d^j (-, +) \end{bmatrix},$$

$$\Psi_u^j = \begin{bmatrix} \tilde{\psi}^j (+, -) \\ \psi_u^j (-, -) \\ \hat{\psi}_d^j (+, -) \end{bmatrix},$$

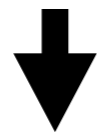
$$\Psi_d^j = \begin{bmatrix} \hat{\psi}^j (+, -) \\ \hat{\psi}_u^j (+, -) \\ \psi_d^j (-, -) \end{bmatrix},$$



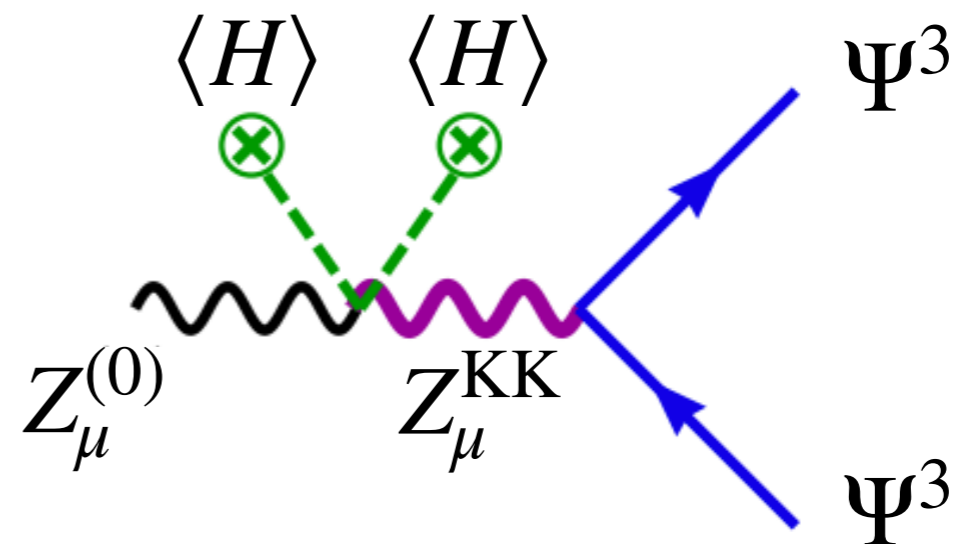
Backup

- The most constraining EW observable is $Z \rightarrow \bar{\tau}\tau$, affected by the mixing of Z and Z^{KK} :

$$\frac{\delta g_{Z\Psi^3\Psi^3}}{g_{Z\Psi^3\Psi^3}} \approx -0.3 \frac{m_Z^2}{M_{KK}^2} \frac{g_*^2}{g_L^2} \approx -\frac{0.3}{4c_W^2} \frac{\langle h \rangle^2}{f^2} \lesssim 10^{-3}$$



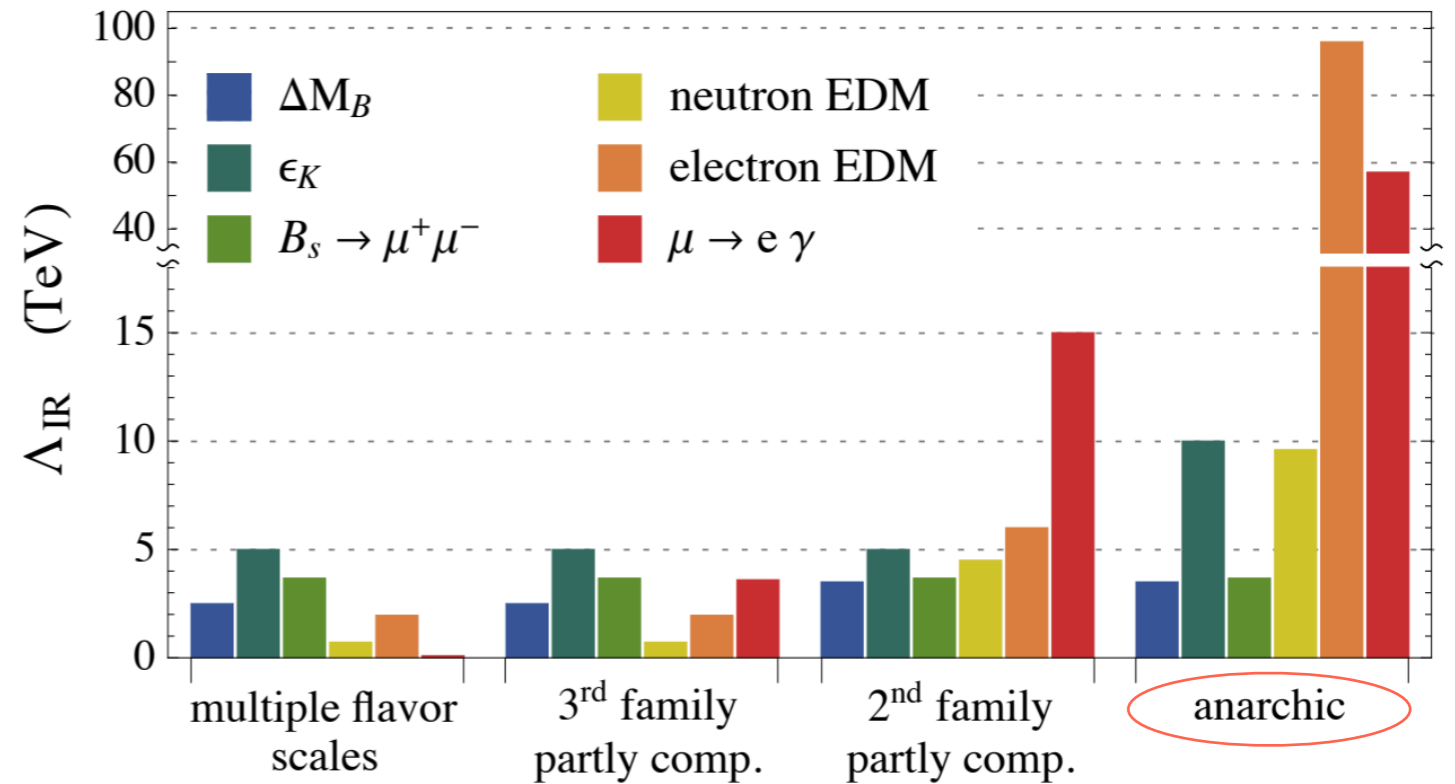
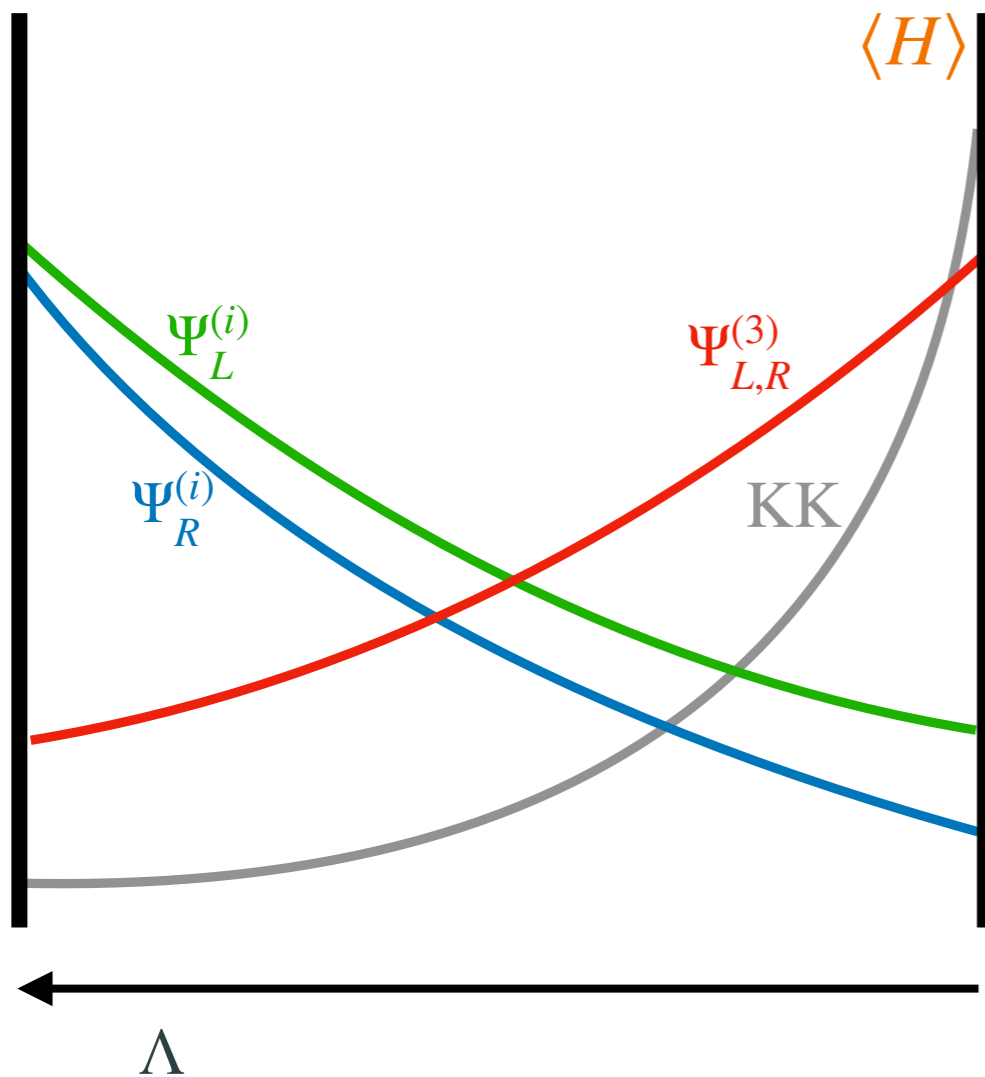
$$f > 2.5 \text{ TeV}, M_{KK} > 6 \text{ TeV}$$



Backup

- Anarchic partial compositeness paradigm in RS

Dangerous dipoles (among others) generated at the IR scale $\sim \frac{g_*^2}{16\pi^2} \frac{m_e}{\Lambda_{\text{IR}}^2} \bar{e}_L \sigma_{\mu\nu} e_R F^{\mu\nu}$



[Panico, Pomarol, 1603.06609]

Backup

Minimal composite Higgs (MCHM)

[Agashe, Contino, Pomarol, [hep-ph/0412089](https://arxiv.org/abs/hep-ph/0412089)]

- Breaking by a composite sector [Fuentes-Martin, Stangl [2004.11376](https://arxiv.org/abs/2004.11376)]

| | |
|-----------------|--|
| Global symmetry | $\mathcal{G}_{\text{global}} = SU(4)_h \times SU(4)_l \times SO(5)$ |
| Gauge symmetry | $\mathcal{G}_{\text{gauge}} = SU(4)_h \times SU(3)_l \times SU(2)_L \times U(1)_{l+R}$ |

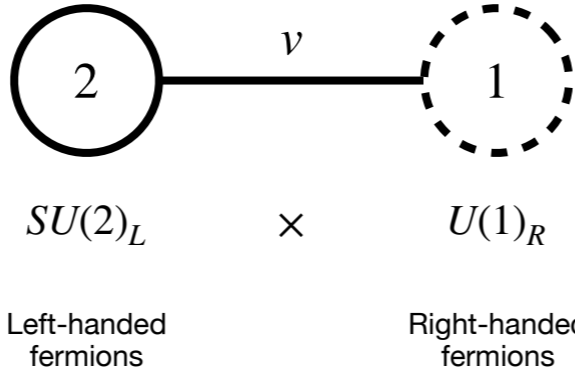
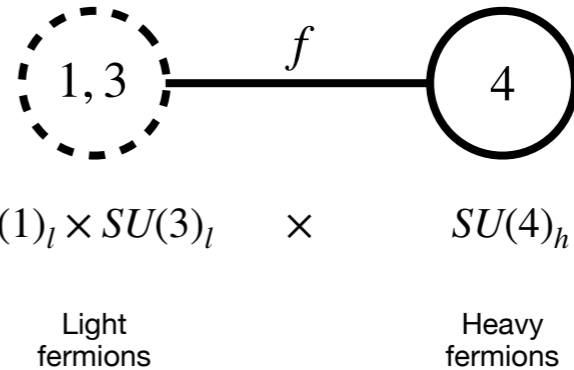
Spontaneously broken by a condensate at some IR scale

| | |
|------------|--|
| Global SSB | $\mathcal{G}_{\text{IR}} = SU(4)_D \times SU(2)_L \times SU(2)_R$ |
| Gauge SSB | $\mathcal{G}_0 = \mathcal{G}_{\text{IR}} \cap \mathcal{G}_{\text{gauge}} = SU(3)_c \times SU(2)_L \times U(1)_Y$ |
| Goldstones | 15 (eaten by U_1, G', Z') + 4 (NGB Higgs) |

SM Higgs emerges as a Nambu-Goldstone boson of the same (strong) dynamics breaking 4321 gauge symmetry

Backup

[Fuentes-Martin, Stangl 2004.11376]

| | SM Higgs Sector | 4321 Models |
|-----------------|--|---|
| Global symmetry | $SU(2)_L \times SU(2)_R$ | $SU(4)_l \times SU(4)_h$ |
| Gauge symmetry |  <p> $SU(2)_L \times U(1)_R$ Left-handed fermions Right-handed fermions </p> |  <p> $U(1)_l \times SU(3)_l \times SU(4)_h$ Light fermions Heavy fermions </p> |
| Global SSB | $SU(2)_V$ | $SU(4)_D$ |
| Gauge SSB | $U(1)_V$ | $U(1)_{B-L} \times SU(3)_c$ |
| Goldstones | 3 (3 eaten) | 15 (15 eaten) |

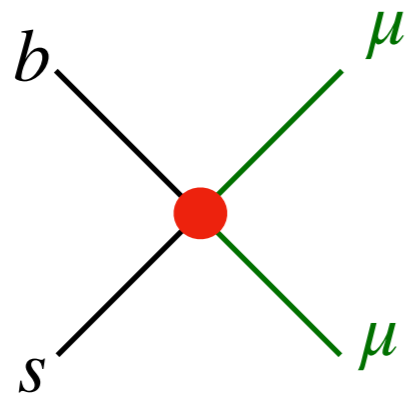
Backup

The two sites are connected by the gauging

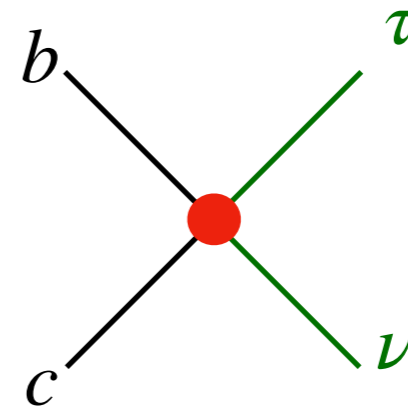
| | SM Higgs Sector | 4321 Models |
|-----------------|--|--------------------------|
| Global symmetry | $SU(2)_L \times SU(2)_R$ | $SU(4)_l \times SU(4)_h$ |
| Gauge symmetry | <p> $SU(2)_L$ × $U(1)_R$ × $U(1)_l \times SU(3)_l$ × $SU(4)_h$ </p> <p> Left-handed fermions Right-handed fermions Light fermions Heavy fermions </p> | |
| Global SSB | $SU(2)_V$ | $SU(4)_D$ |
| Gauge SSB | $U(1)_{em} \times SU(3)_c$ | |
| Goldstones | 3 (W, Z) | 15 (U_1, G', Z') |

B-anomalies

$$b \rightarrow sll$$



$$b \rightarrow c\tau\nu$$



- Non-universality in e/μ , $> 4\sigma$

$$c \sim (40 \text{ TeV})^{-2}$$

$$3_q \rightarrow 2_q 2_l 2_l$$

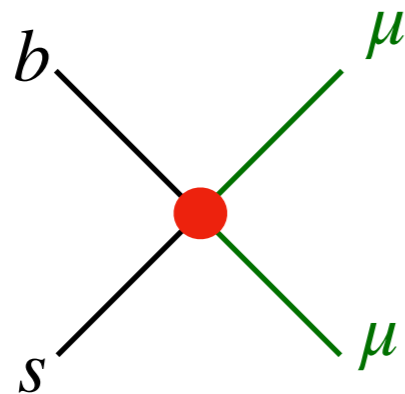
- Non universality in $\tau/\mu, e$, $\sim 3\sigma$

$$c \sim (3 \text{ TeV})^{-2}$$

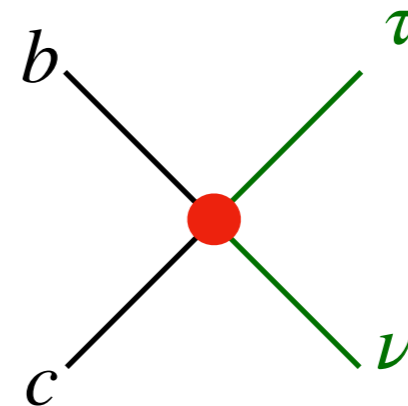
$$3_q \rightarrow 2_q 3_l 3_l$$

B-anomalies

$$b \rightarrow sll$$



$$b \rightarrow c\tau\nu$$



- Non-universality in e/μ , $> 4\sigma$

- Non universality in $\tau/\mu, e$, $\sim 3\sigma$

$$c \sim \epsilon_q \epsilon_l^2 \text{TeV}^{-2}$$

$$3_q \rightarrow 2_q 2_l 2_l$$

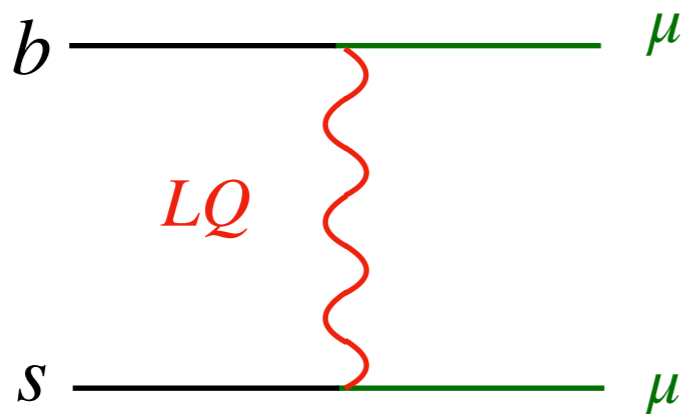
$$\epsilon_q, \epsilon_l \sim 0.1$$

$$c \sim \epsilon_q \text{TeV}^{-2}$$

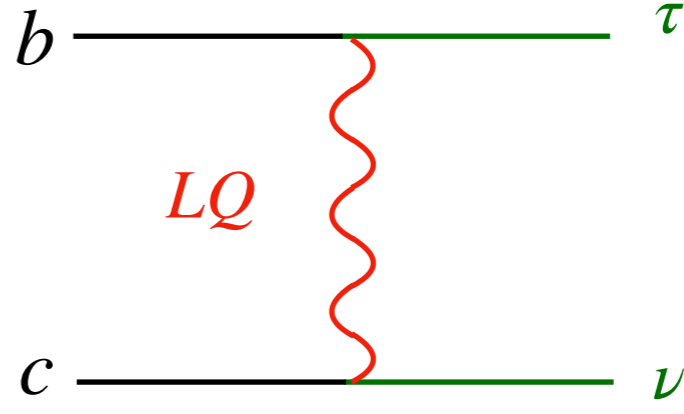
$$3_q \rightarrow 2_q 3_l 3_l$$

B-anomalies

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- Non-universality in e/μ , $> 4\sigma$

- Non universality in $\tau/\mu, e$, $\sim 3\sigma$

$$c \sim \epsilon_q \epsilon_l^2 \text{TeV}^{-2}$$

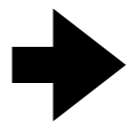
$$3_q \rightarrow 2_q 2_l 2_l$$

$$\epsilon_q, \epsilon_l \sim 0.1$$

$$c \sim \epsilon_q \text{TeV}^{-2}$$

$$3_q \rightarrow 2_q 3_l 3_l$$

LQ mostly coupled to the third family



$U(2)^5$ in light families
to protect flavour observables