

Searches for heavy neutral leptons: beyond simplified scenarios

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Local organizing committee:

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S. Loucatos (IRFU/APC)
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Photo by Karim Benakli

Gioacchino Piazza

IJCLab, Pôle Théorie, CNRS and Paris Saclay U.

In collab. with A. Abada, X. Marcano and P. Escribano



Outline of the talk

- Collider searches of Heavy Neutral Leptons (HNL) (*1 HNL + single mixing hypothesis*)
- How to reinterpret the bounds on $|V_{\ell N}|^2$ in two cases:
 - One HNL *mixing with 3 active flavors* (e, μ, τ)
 - *Two HNLs interfering* in the LNC/LNV processes ($W^+ \rightarrow \ell_1^+ \ell_2^\pm q \bar{q}'$)

Neutrino window to new physics

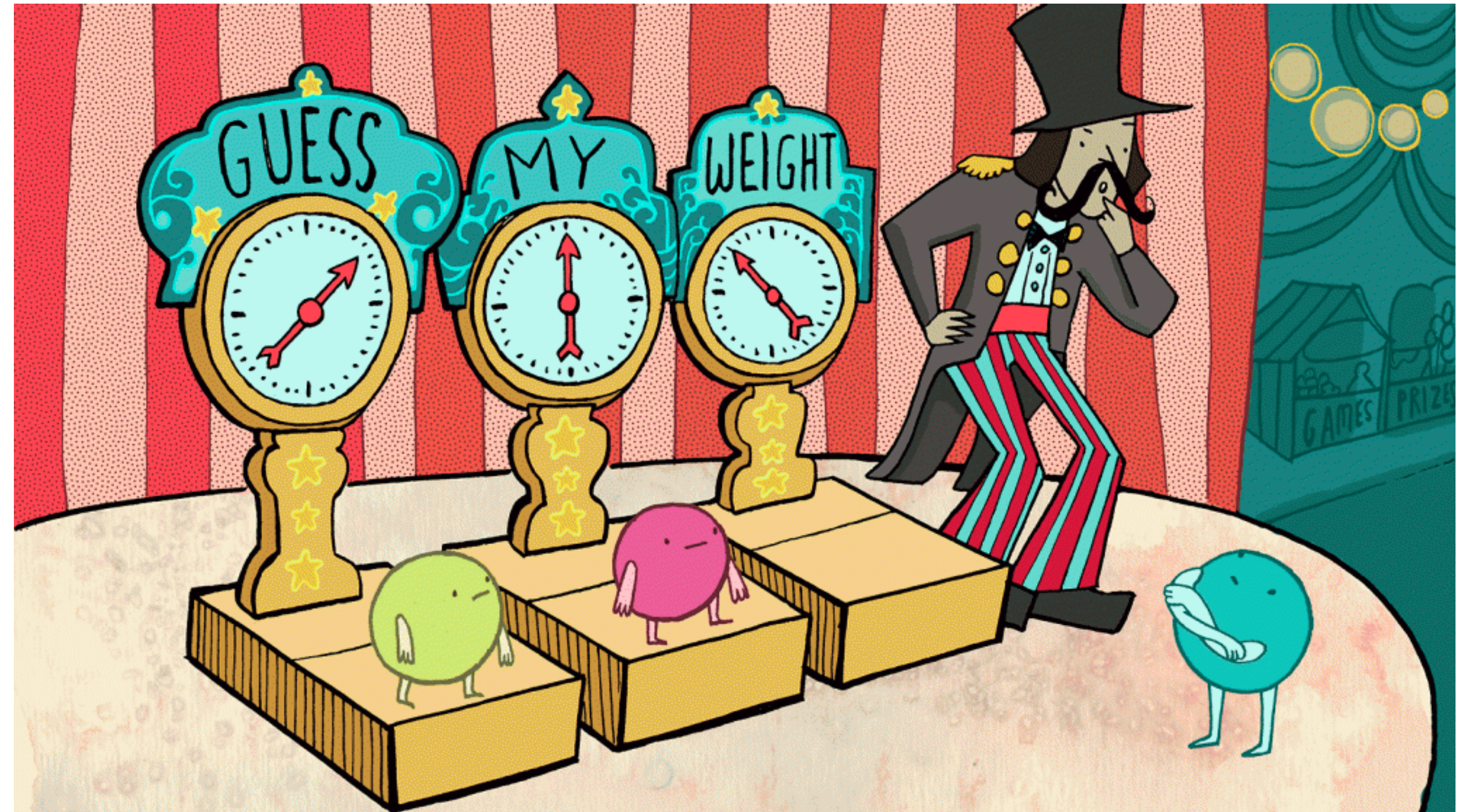
Neutrino oscillations establish that at least **two** of the SM neutrinos have masses $\neq 0$

$$\theta_{12} \simeq 33^\circ \quad \theta_{23} \simeq 49^\circ \quad \theta_{13} \simeq 8^\circ$$

$$\Delta m_{12}^2 \simeq 7 \times 10^{-5} \text{ eV}^2$$

$$\Delta m_{3\ell}^2 \simeq \pm 2.5 \times 10^{-3} \text{ eV}^2$$

see [2111.03086]



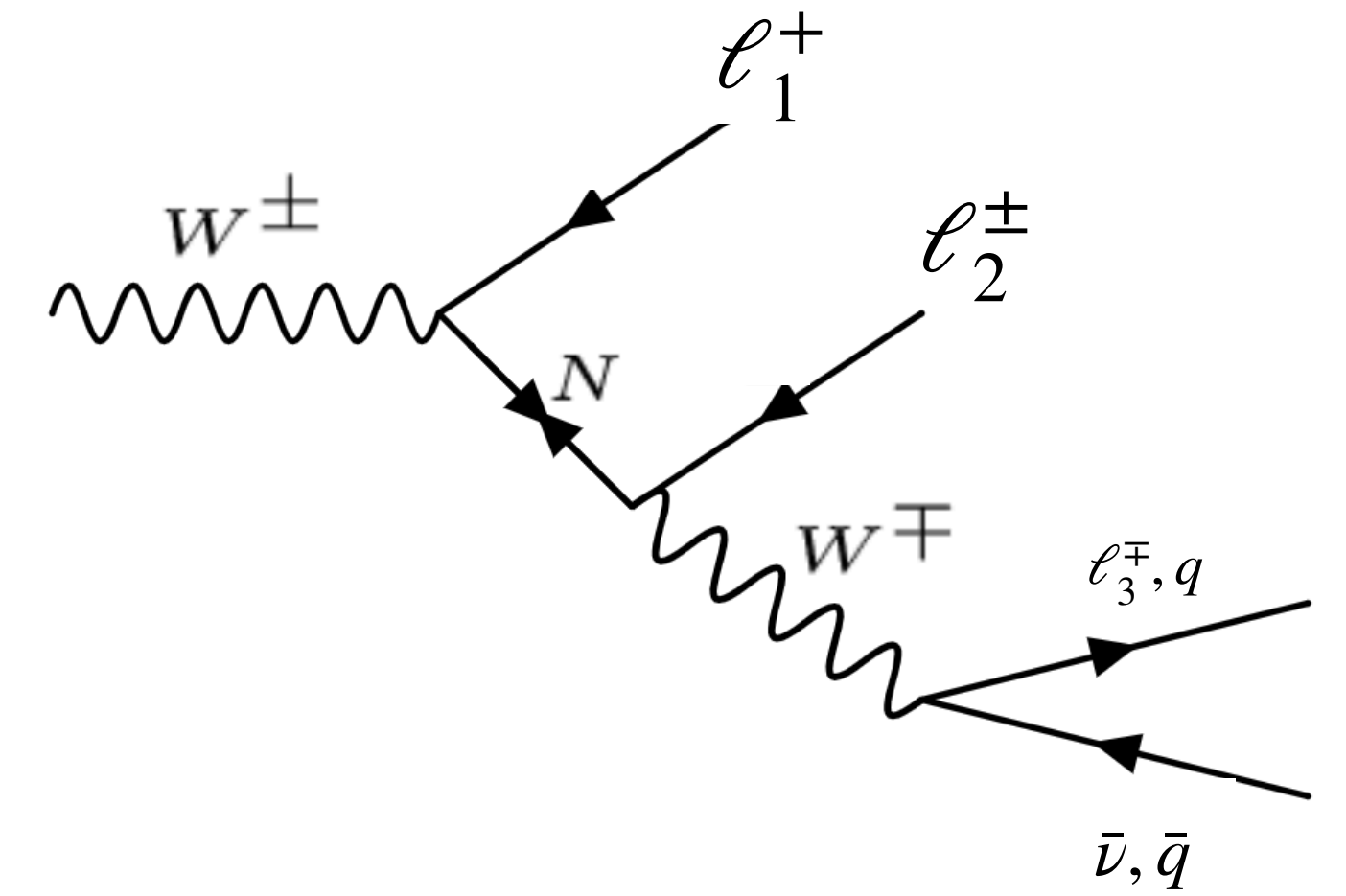
New neutral fermions are common byproducts of ν mass generation.

The case with **2 HNLs** is the minimal phenomenologically viable when embedding a seesaw.

HNL collider searches

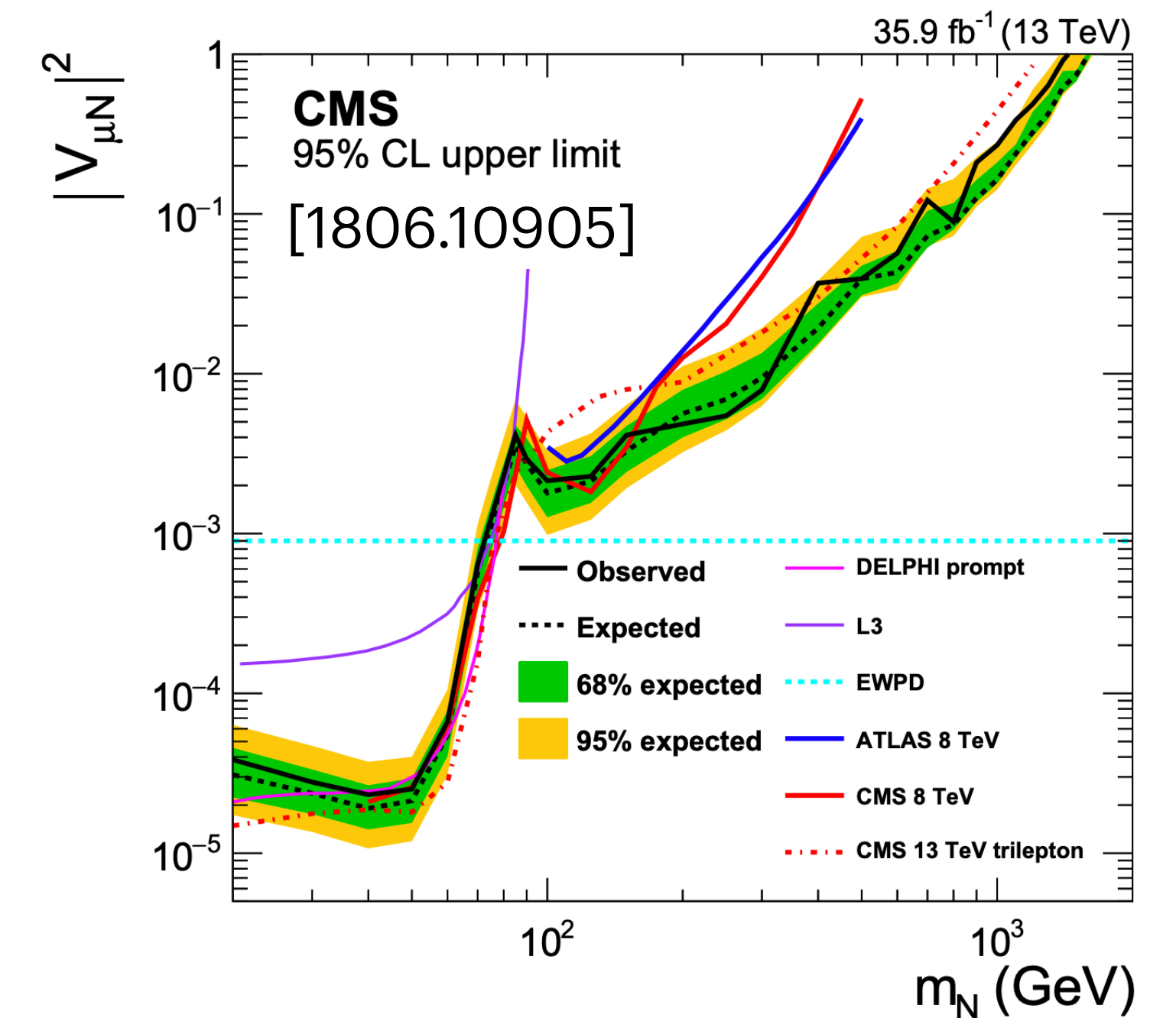
3 leptons + ME

CMS [1802.02965]	LNV, LFV or LFC	$pp \rightarrow \ell_1^+ \ell_2^+ \ell_3^- + \text{ME}$ (1,2,3 = μ, e)	[1,1200] GeV
ATLAS [1905.09787]	LNV, LFC	$pp \rightarrow \ell^+ \ell^+ \ell'^- + \text{ME}$ ($\ell^{(\prime)}$ = μ, e)	[5,50] GeV

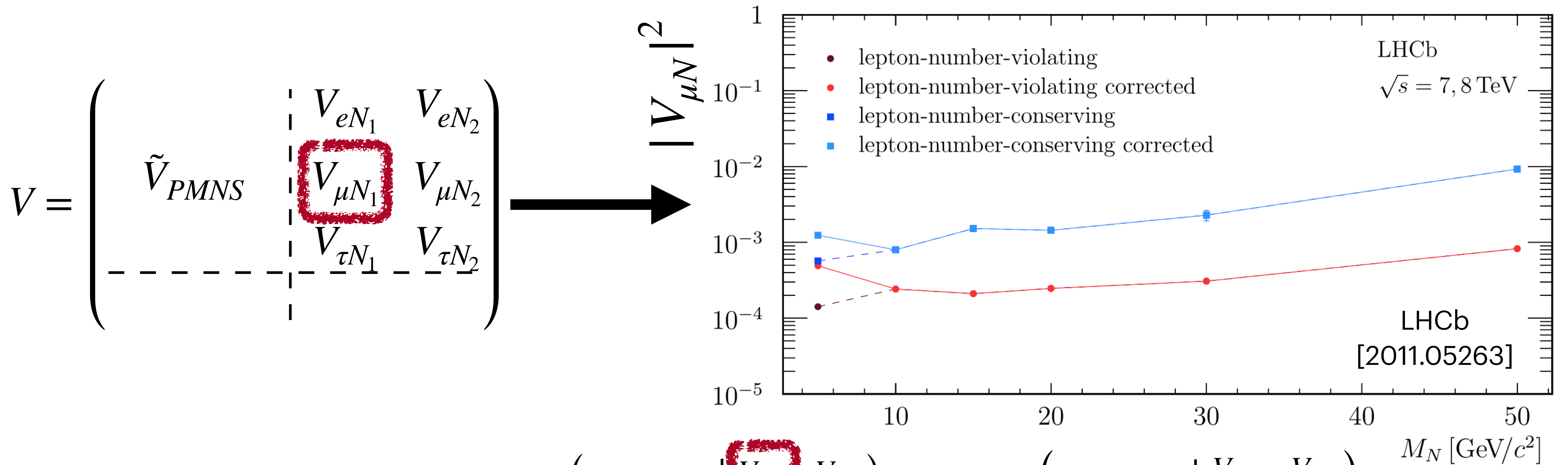


2 leptons + jets

ATLAS [1506.06020]	LNV, LFC	$pp \rightarrow \ell_1^\pm \ell_2^\pm j$ (1,2 = μ, e)	[100,500] GeV
CMS [1806.10905]	LNV, LFV or LFC	$pp \rightarrow \ell^\pm \ell^\pm j$ (ℓ = μ, e)	[20,1600] GeV
LHCb [2011.05263]	LNC or LNV, LFC	$pp \rightarrow \mu^+ \mu^\pm j$	[5,50] GeV



Experimental searches generally assume **one** generic HNL
and **single mixing hypothesis** ($V_{\alpha N}$)



Relax these assumptions:

$$V = \begin{pmatrix} \tilde{V}_{PMNS} & V_{eN_1} & V_{eN_2} \\ \tilde{V}_{PMNS} & V_{\mu N_1} & V_{\mu N_2} \\ \tilde{V}_{PMNS} & V_{\tau N_1} & V_{\tau N_2} \end{pmatrix} \quad \text{or} \quad V = \begin{pmatrix} \tilde{V}_{PMNS} & V_{eN_1} & V_{eN_2} \\ \tilde{V}_{PMNS} & V_{\mu N_1} & V_{\mu N_2} \\ \tilde{V}_{PMNS} & V_{\tau N_1} & V_{\tau N_2} \end{pmatrix}$$

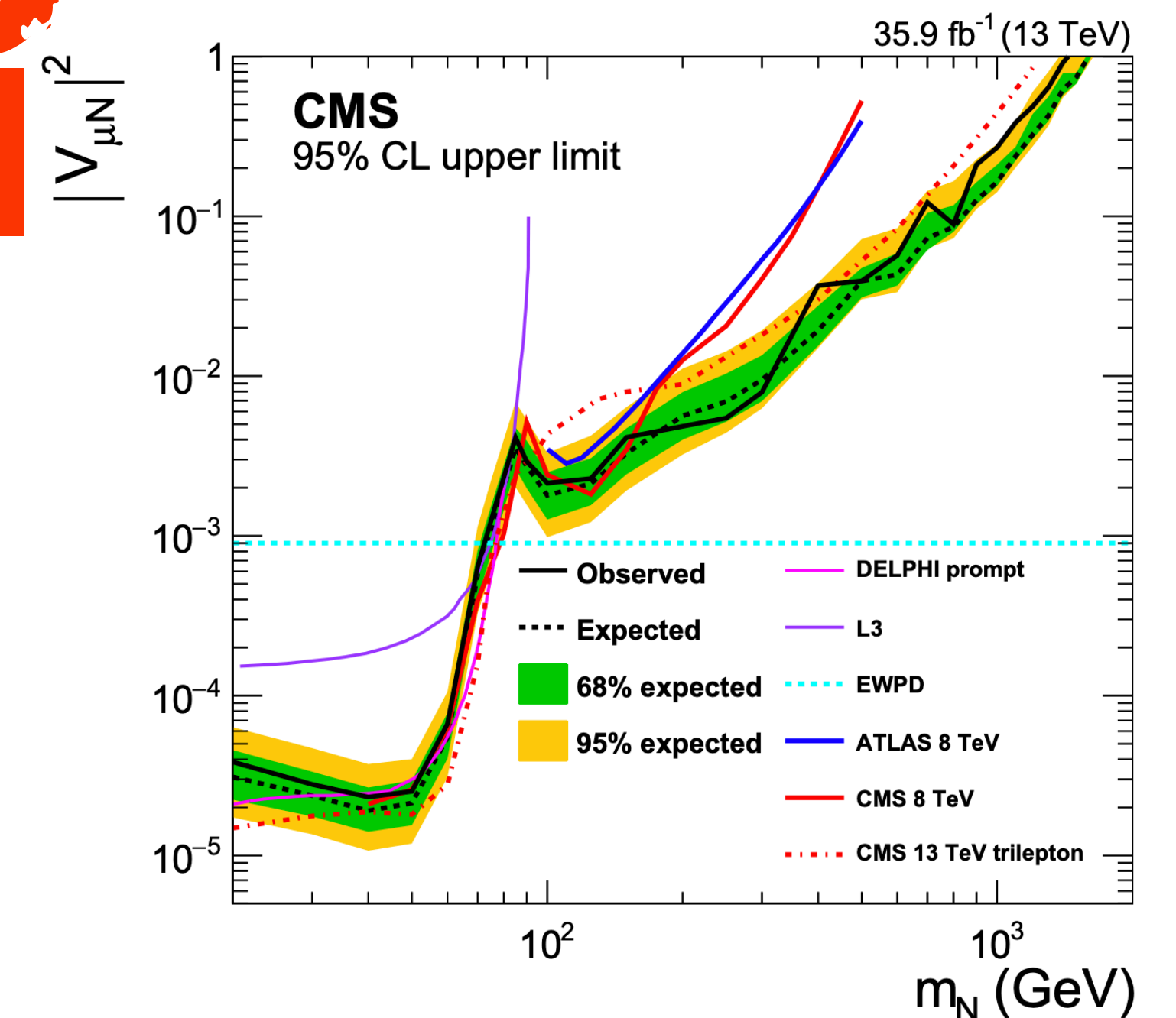
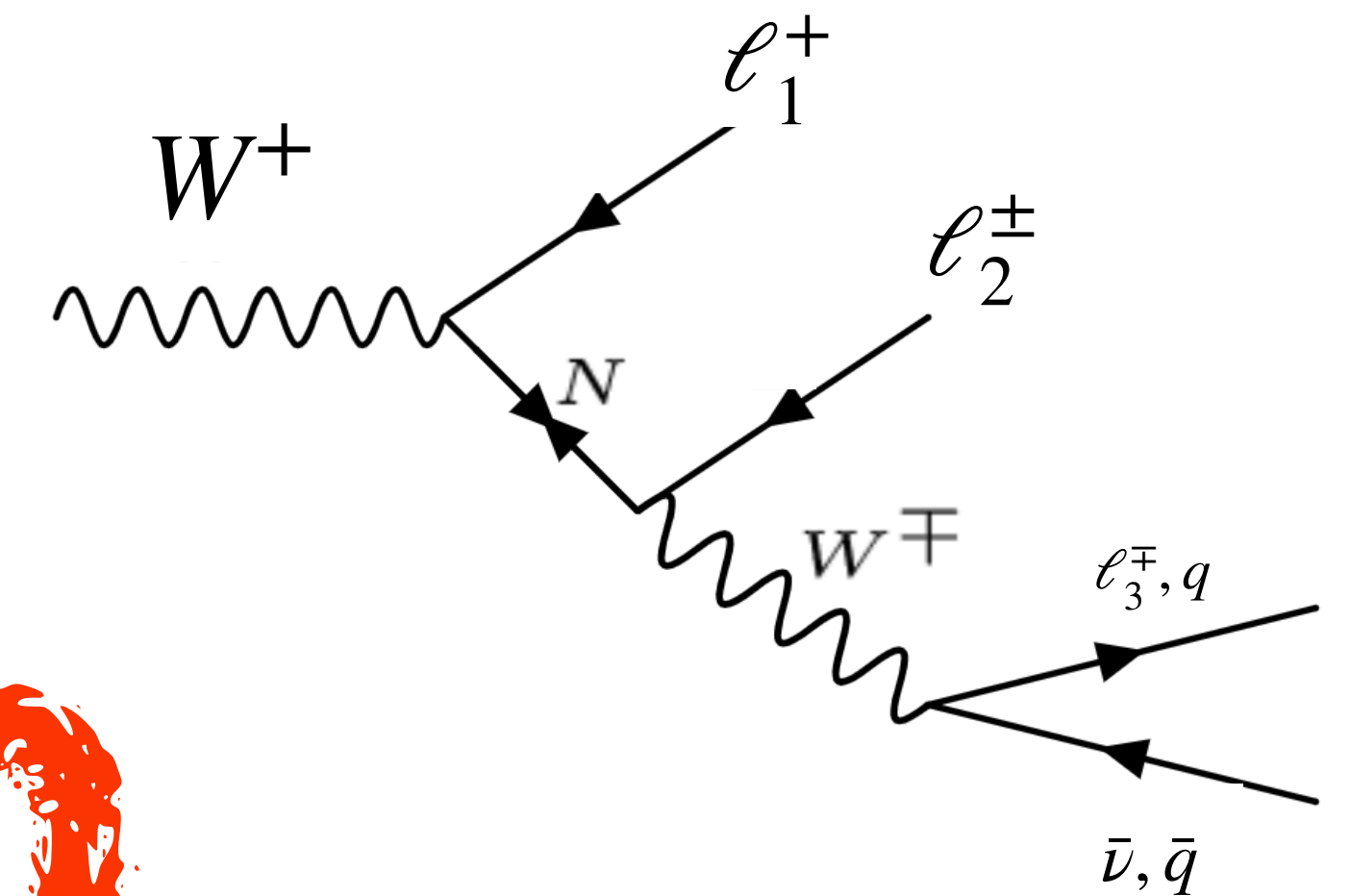
HNL collider searches

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See [Tastet et al. 2107.12980]

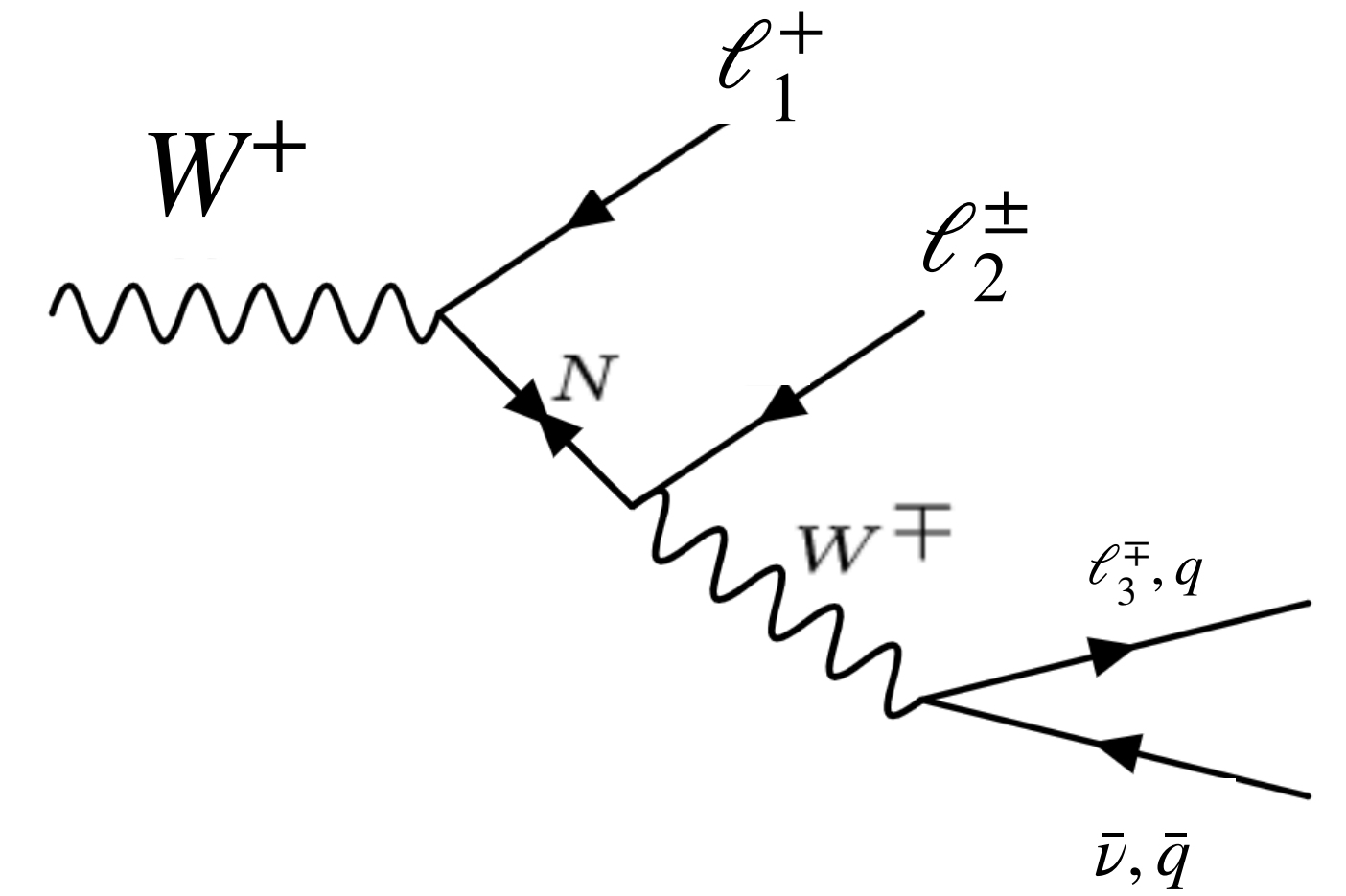
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HNL collider searches

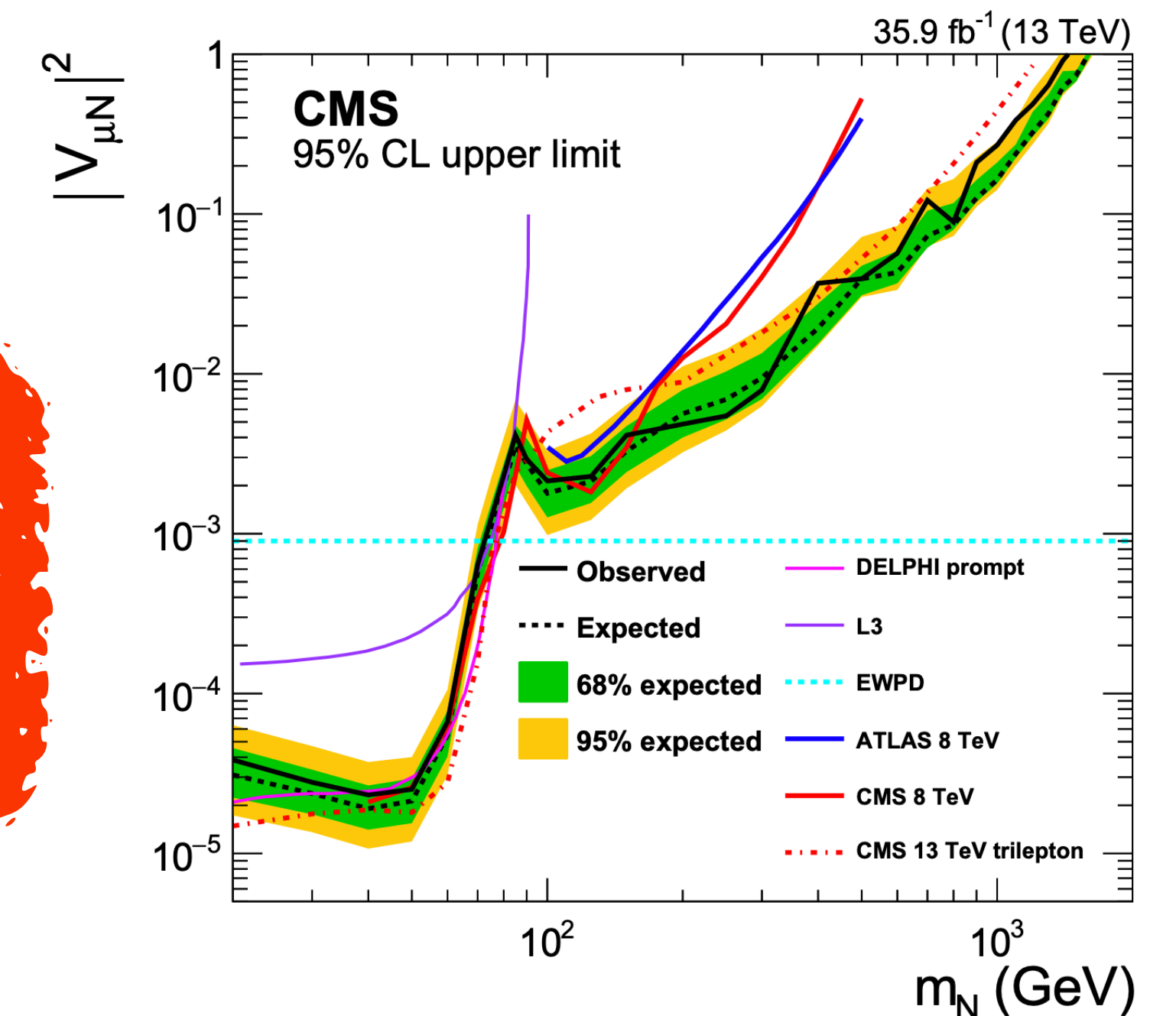
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2 leptons + jets

ATLAS [1510.06626]	LNV, LFC	$pp \rightarrow \ell^\pm \ell^\pm j$ (1,2 = μ, e)	[100,500] GeV
CMS [1806.10905]	LNV, LFV or LFC	$pp \rightarrow \ell^\pm \ell^\pm j$ ($\ell = \mu, e$)	[20,1600] GeV
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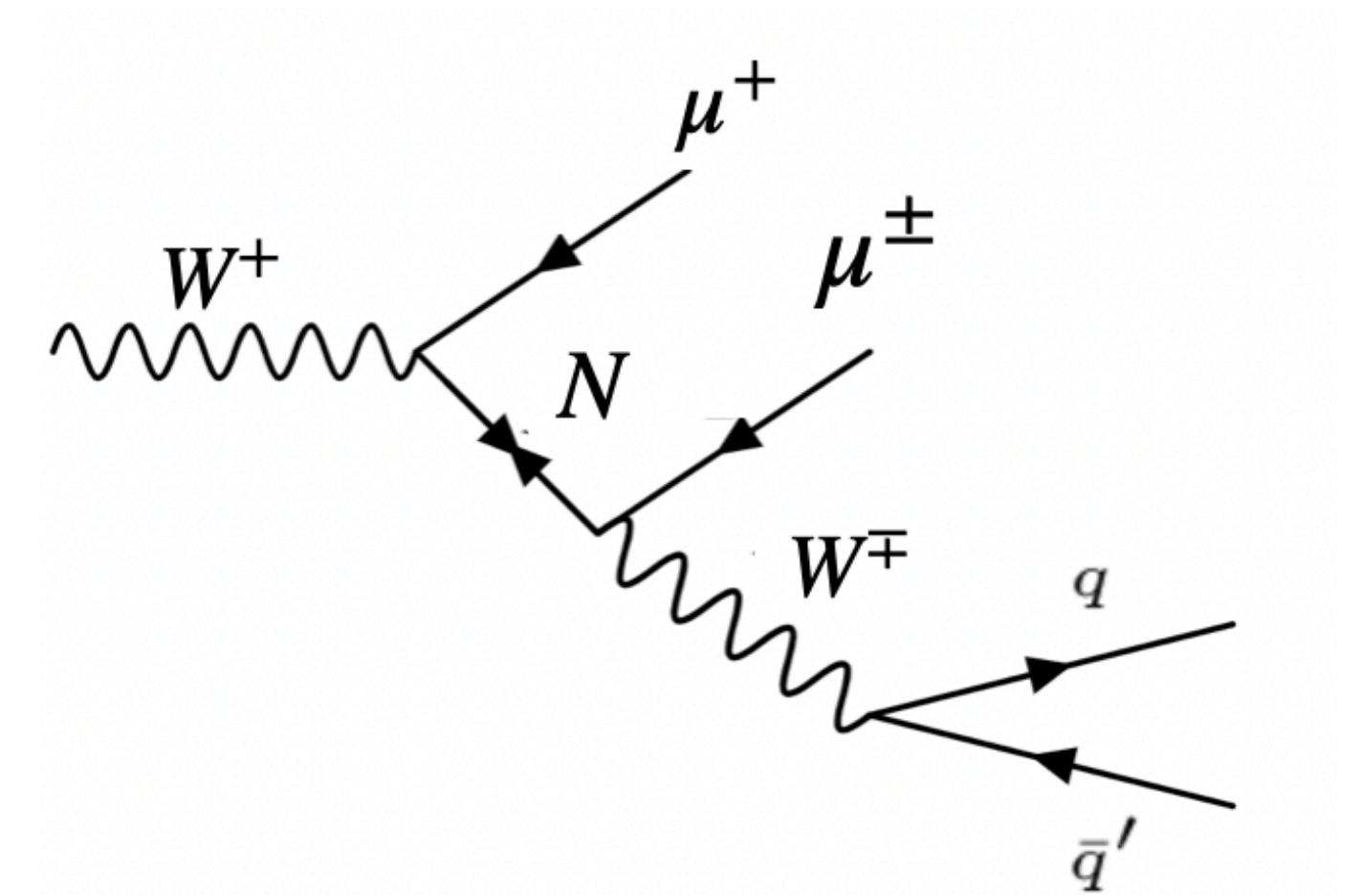
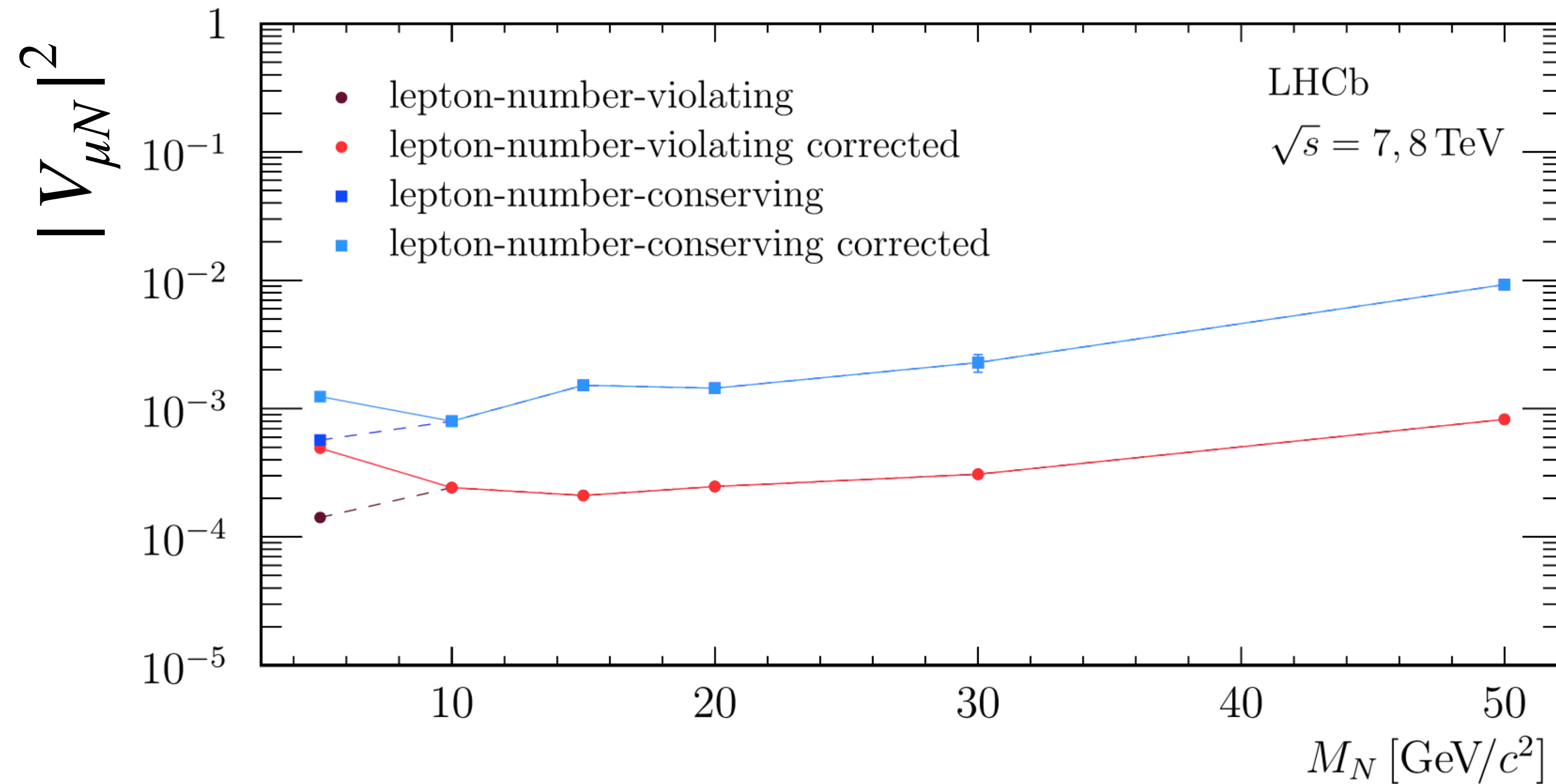


Short-lived HNL

2 leptons + jets: **generic mixing pattern (1HNL)**

$$V = \left(\begin{array}{c|cc} \tilde{V}_{PMNS} & V_{eN_1} & V_{eN_2} \\ & V_{\mu N_1} & V_{\mu N_2} \\ & V_{\tau N_1} & V_{\tau N_2} \\ \hline & & \end{array} \right)$$

2 leptons + jets: generic mixing pattern (1HNL)



2 leptons + jets: generic mixing pattern (1HNL)

$$\Gamma(W^+ \rightarrow \ell_\alpha^+ \ell_\beta^\pm q \bar{q}') = \Gamma(W^+ \rightarrow \ell_\alpha^+ N) \times \text{Br}(N \rightarrow \ell_\beta^\pm q \bar{q}') \propto |V_{\alpha N}|^2 \times \text{Br}(N \rightarrow \ell_\beta^\pm q \bar{q}')$$

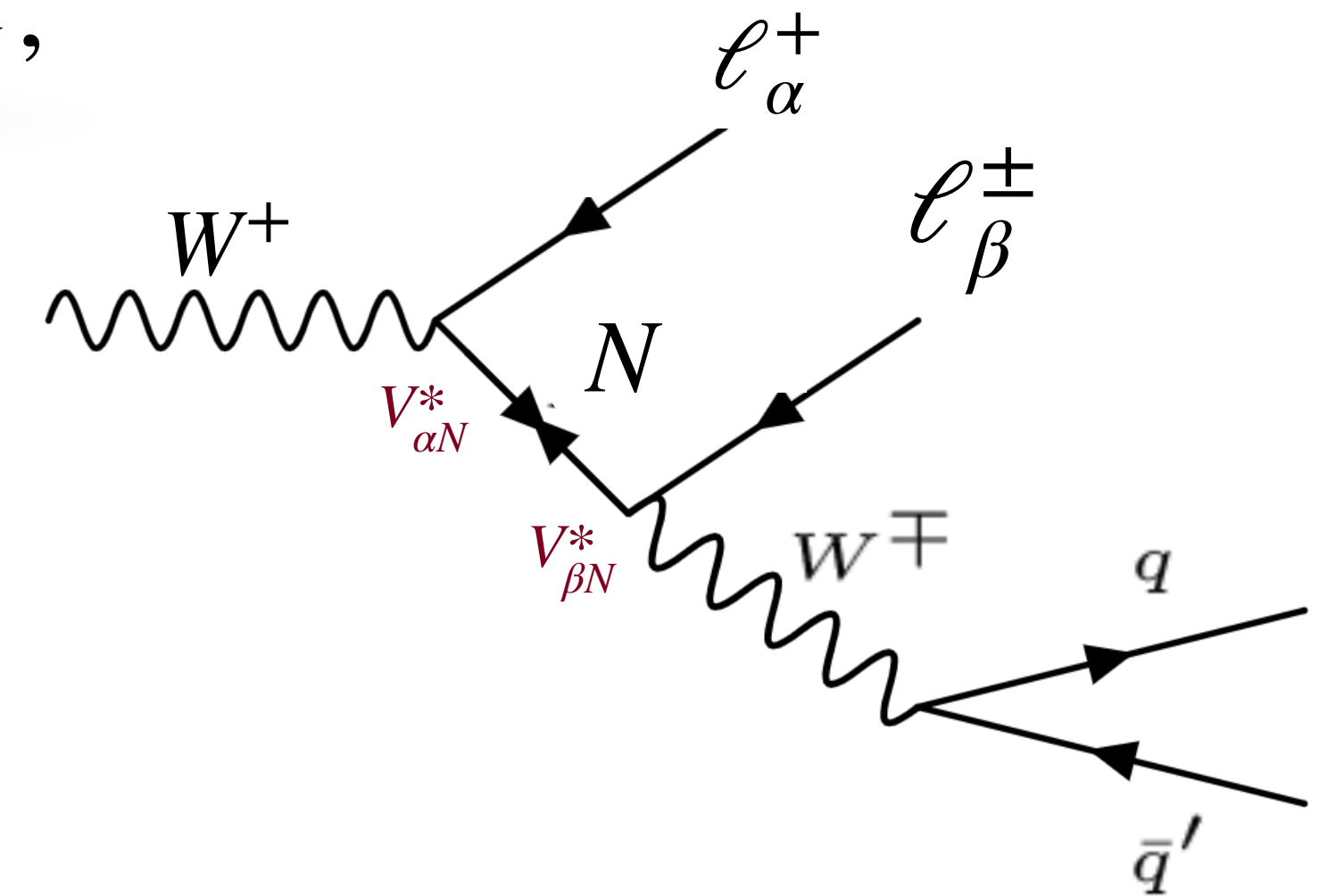
- Single mixing ($\alpha = \beta$): $\Gamma_N^{\text{Sing}} = |V_{\alpha N}|^2 \tilde{\Gamma}_N^\alpha$

$$\frac{\Gamma(N \rightarrow \ell_\beta^\pm q \bar{q}')}{\Gamma_N}$$

- Generic mixing: $\Gamma_N^{\text{Gen}} = |V_{eN}|^2 \tilde{\Gamma}_N^e + |V_{\mu N}|^2 \tilde{\Gamma}_N^\mu + |V_{\tau N}|^2 \tilde{\Gamma}_N^\tau$,

$$V = \begin{pmatrix} \tilde{V}_{PMNS} & \begin{matrix} | & V_{eN_1} & V_{eN_2} \\ | & V_{\mu N_1} & V_{\mu N_2} \\ | & V_{\tau N_1} & V_{\tau N_2} \\ | & & \end{matrix} \end{pmatrix}$$

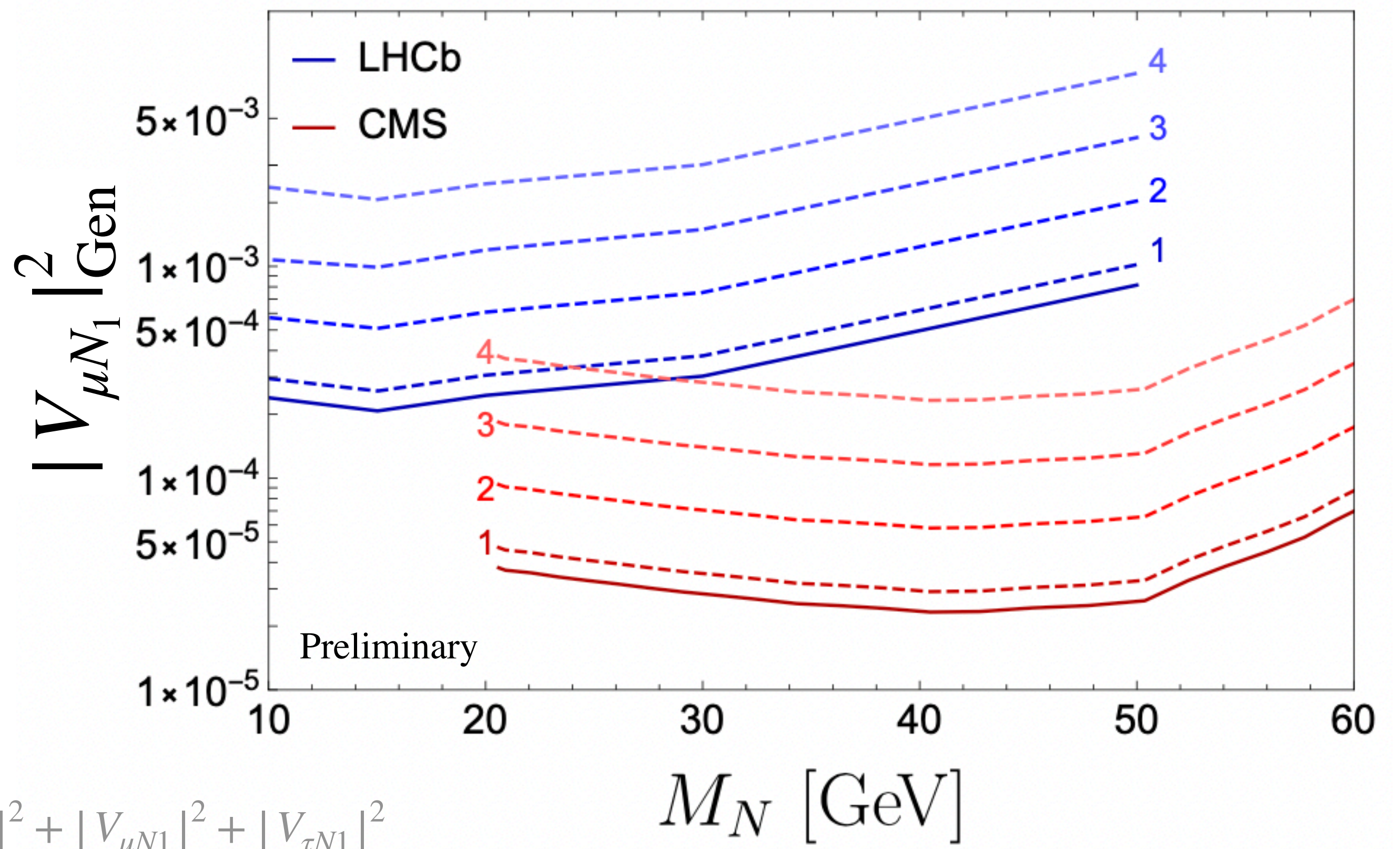
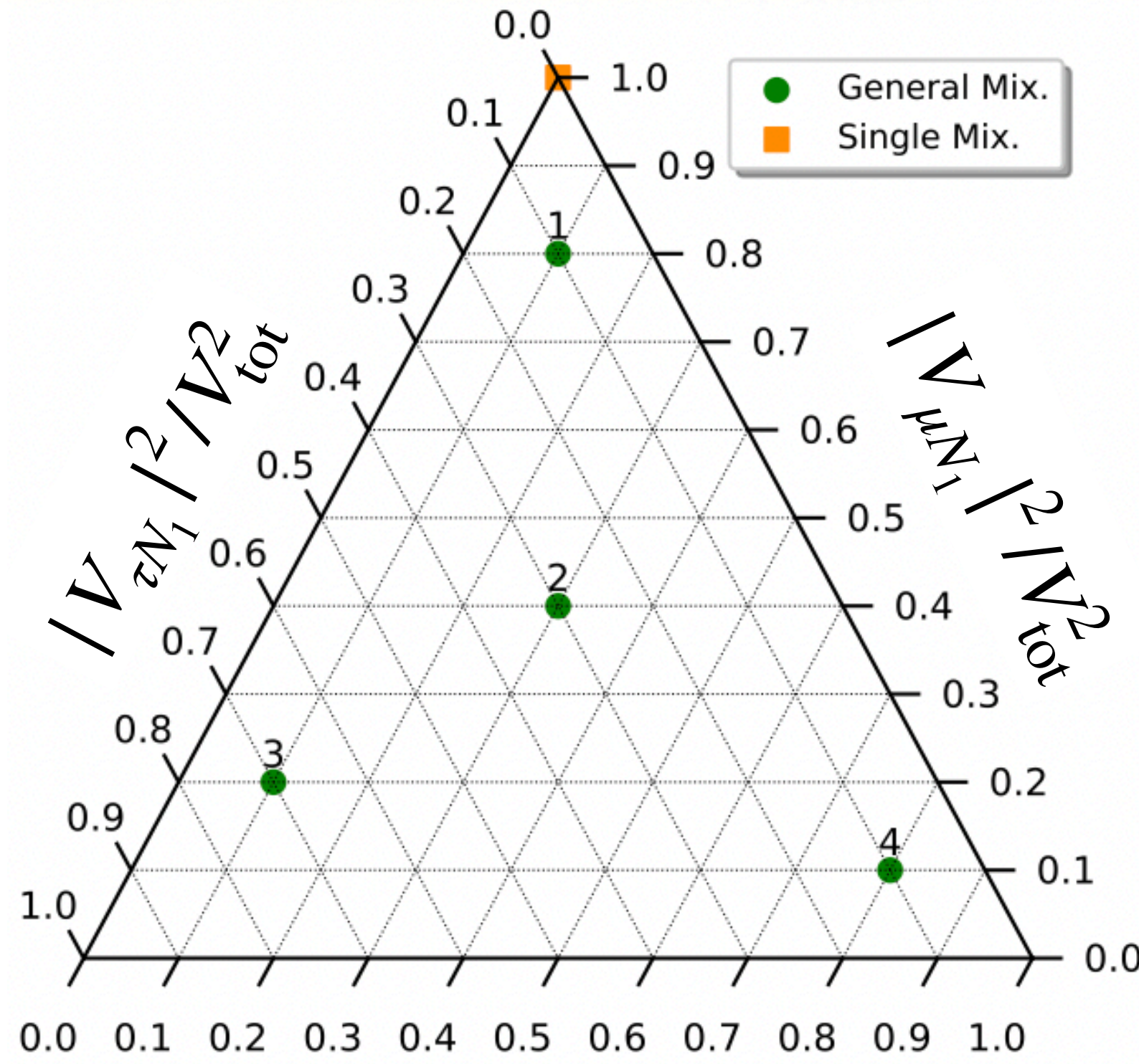
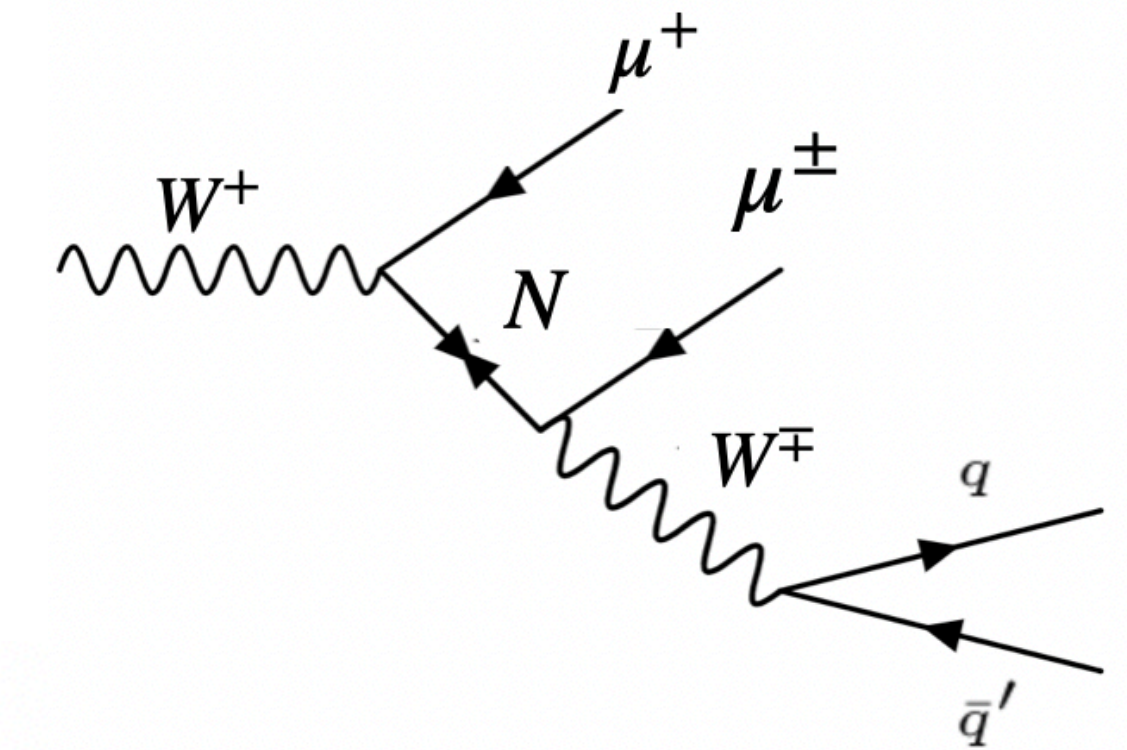
$$|V_{\mu N_1}|_{\text{Gen}}^2 = |V_{\mu N_1}|_{\text{Sing}}^2 \frac{\Gamma_N^{\text{Gen}}}{\Gamma_N^{\text{Sing}}}$$



2 leptons + jets: generic mixing pattern (1HNL)

$$|V_{\mu N_1}|_{\text{Gen}}^2 = |V_{\mu N_1}|_{\text{Sing}}^2 \frac{\Gamma_N^{\text{Gen}}}{\Gamma_N^{\text{Sing}}}$$

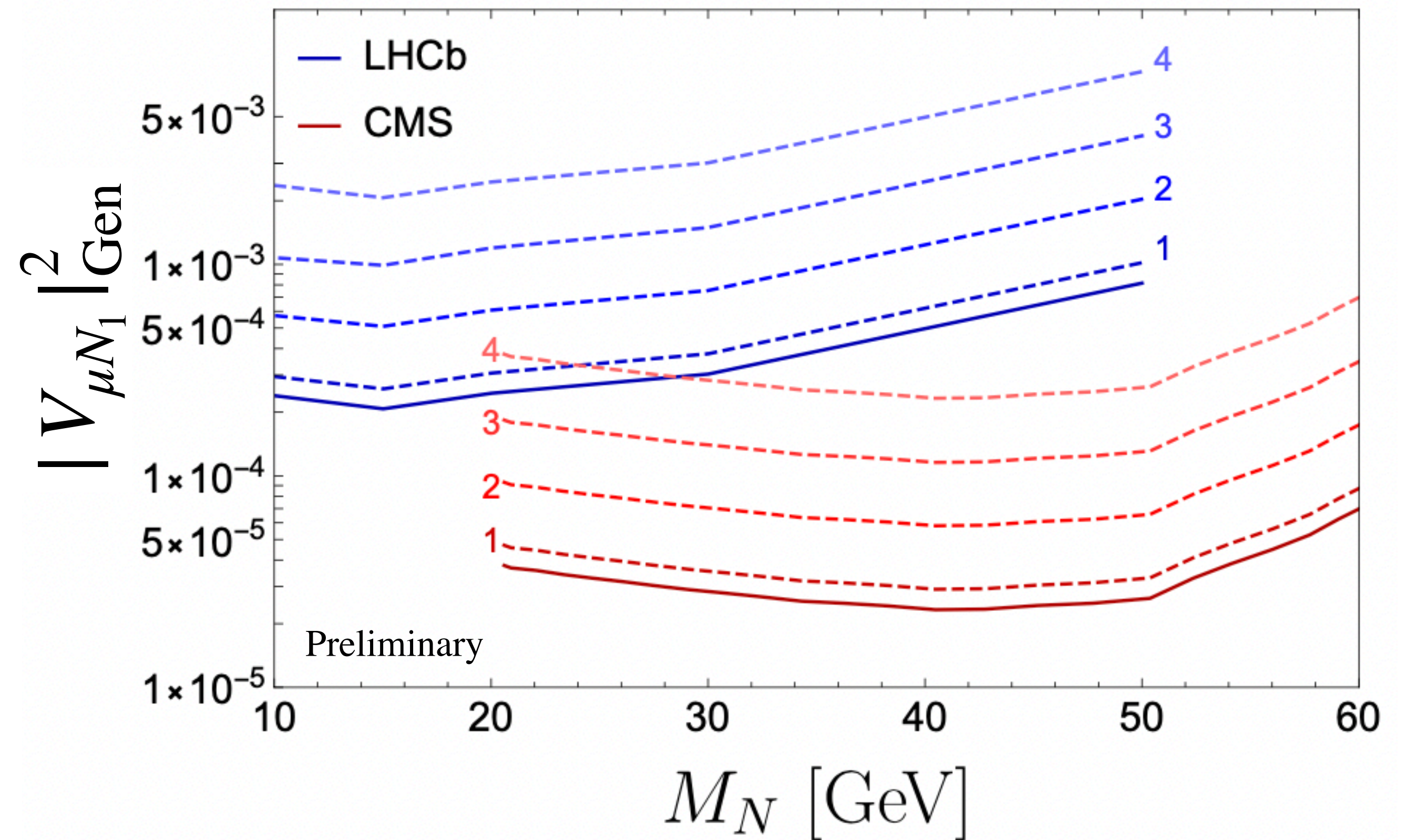
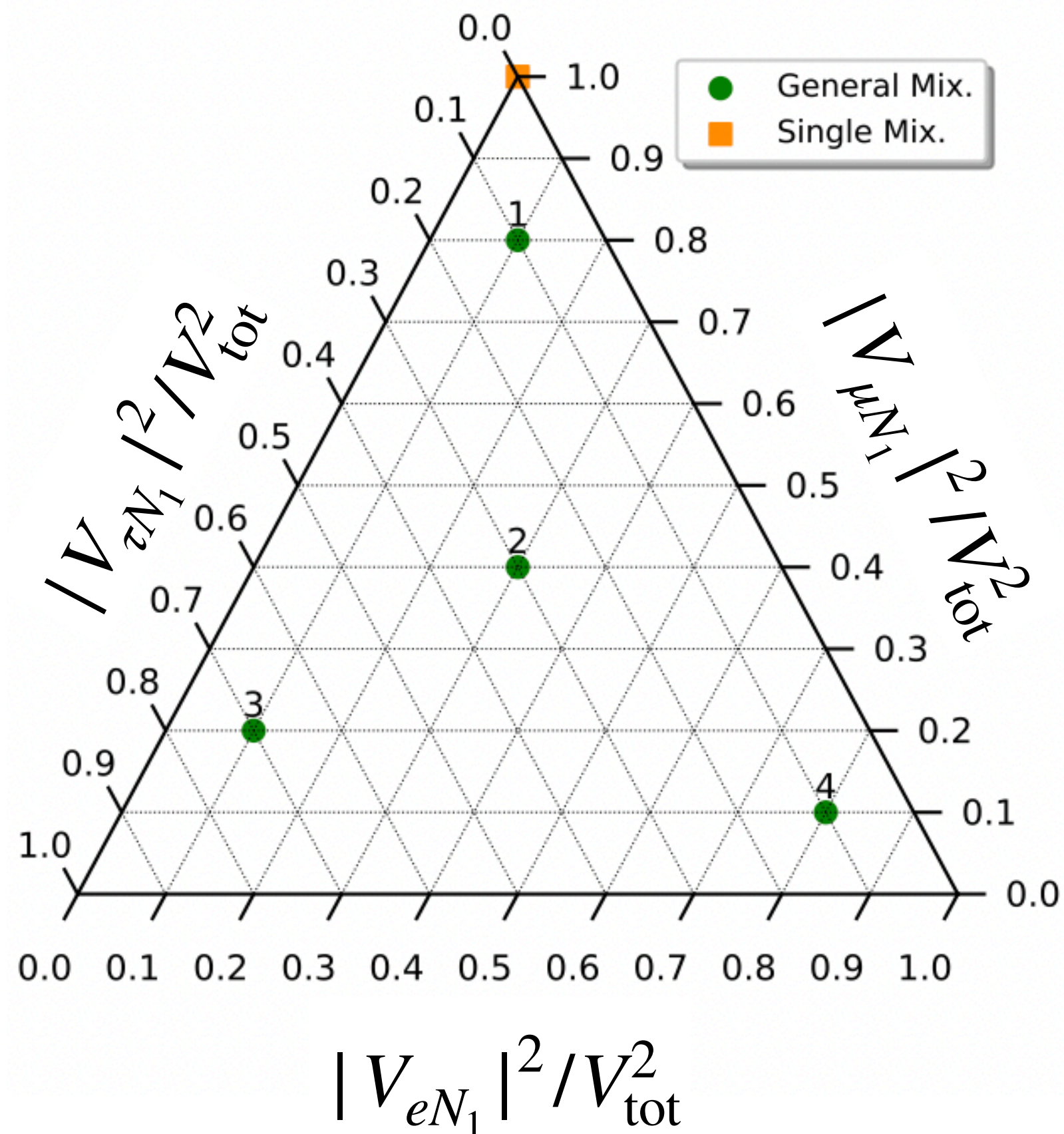
$$V = \begin{pmatrix} \tilde{V}_{PMNS} & \begin{matrix} V_{eN_1} \\ V_{\mu N_1} \\ V_{\tau N_1} \end{matrix} & \begin{matrix} V_{eN_2} \\ V_{\mu N_2} \\ V_{\tau N_2} \end{matrix} \end{pmatrix}$$



$$|V_{eN_1}|^2/V_{\text{tot}}^2 \quad |V_{\text{tot}}|^2 = |V_{eN_1}|^2 + |V_{\mu N_1}|^2 + |V_{\tau N_1}|^2$$

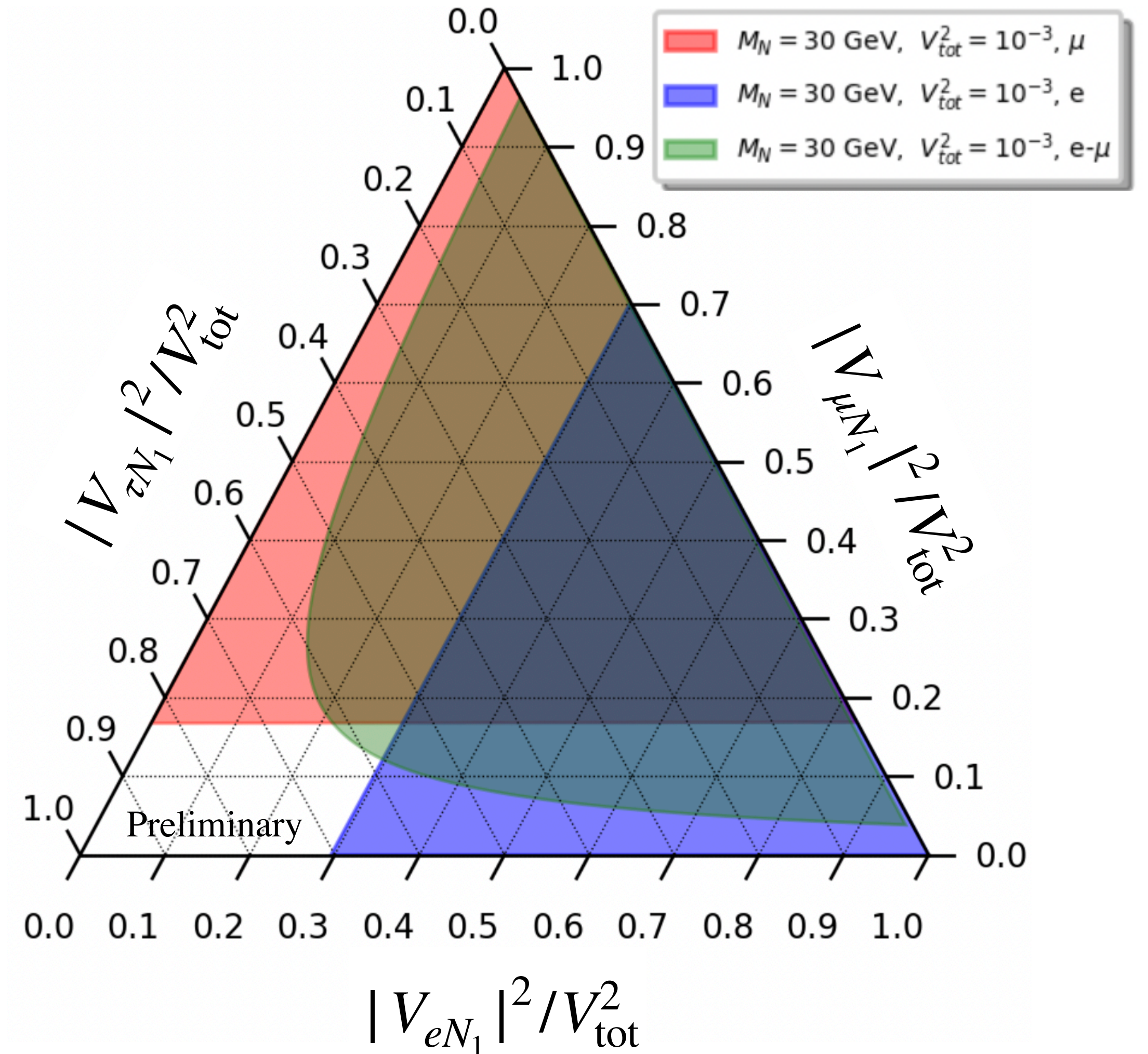
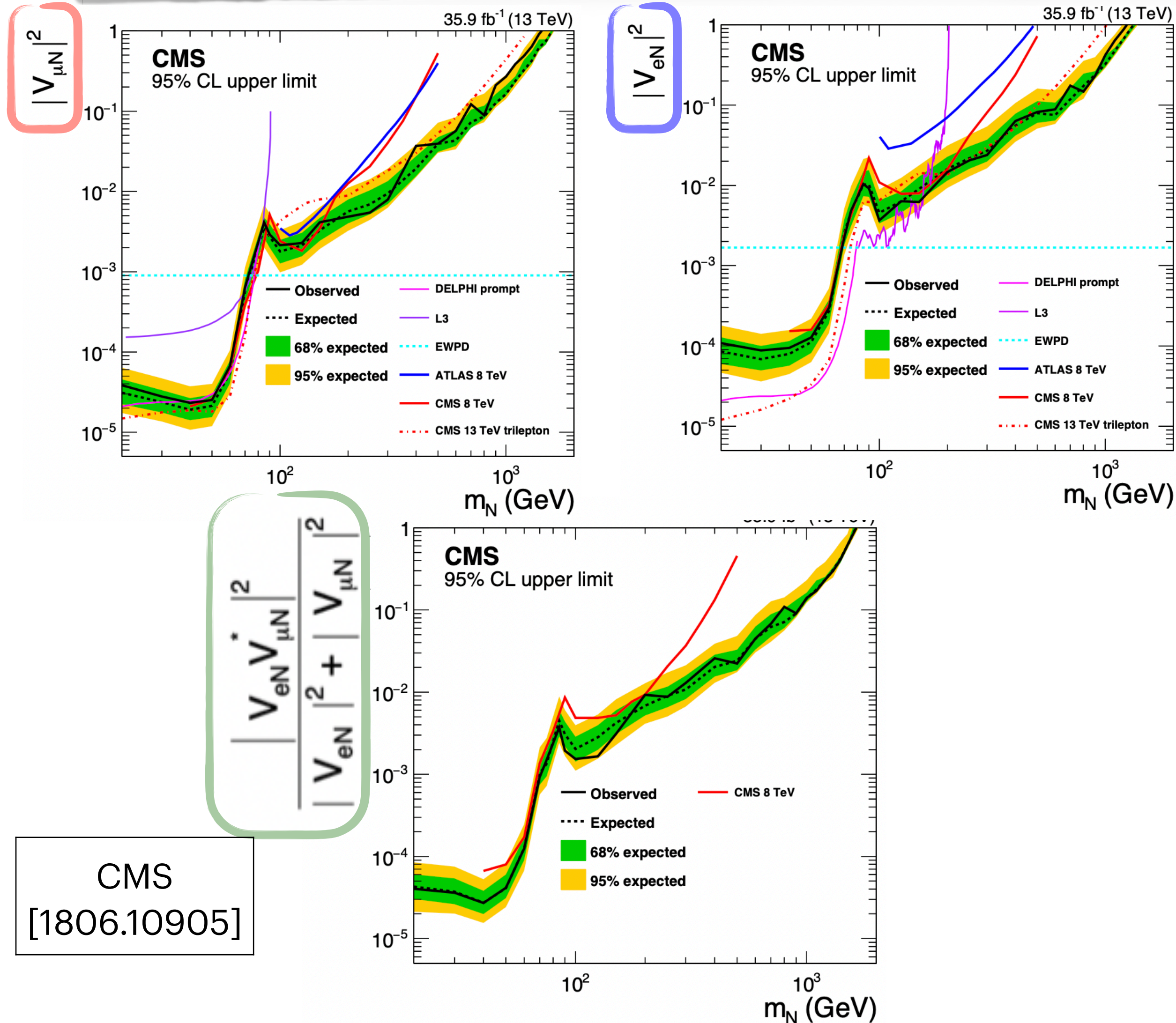
2 leptons + jets: generic mixing pattern (1HNL)

Fixing (M_N, V_{tot}^2) and scanning the parameter space in the ternary plot, check which points are allowed after the rescaling of the bound.



2 leptons + jets: generic mixing pattern (1HNL)

Fixing (M_N, V_{tot}^2) , and combining the searches on $(e, \mu, e - \mu)$:



CMS
[1806.10905]

2 leptons + jets: 2 HNLs (single mixing)

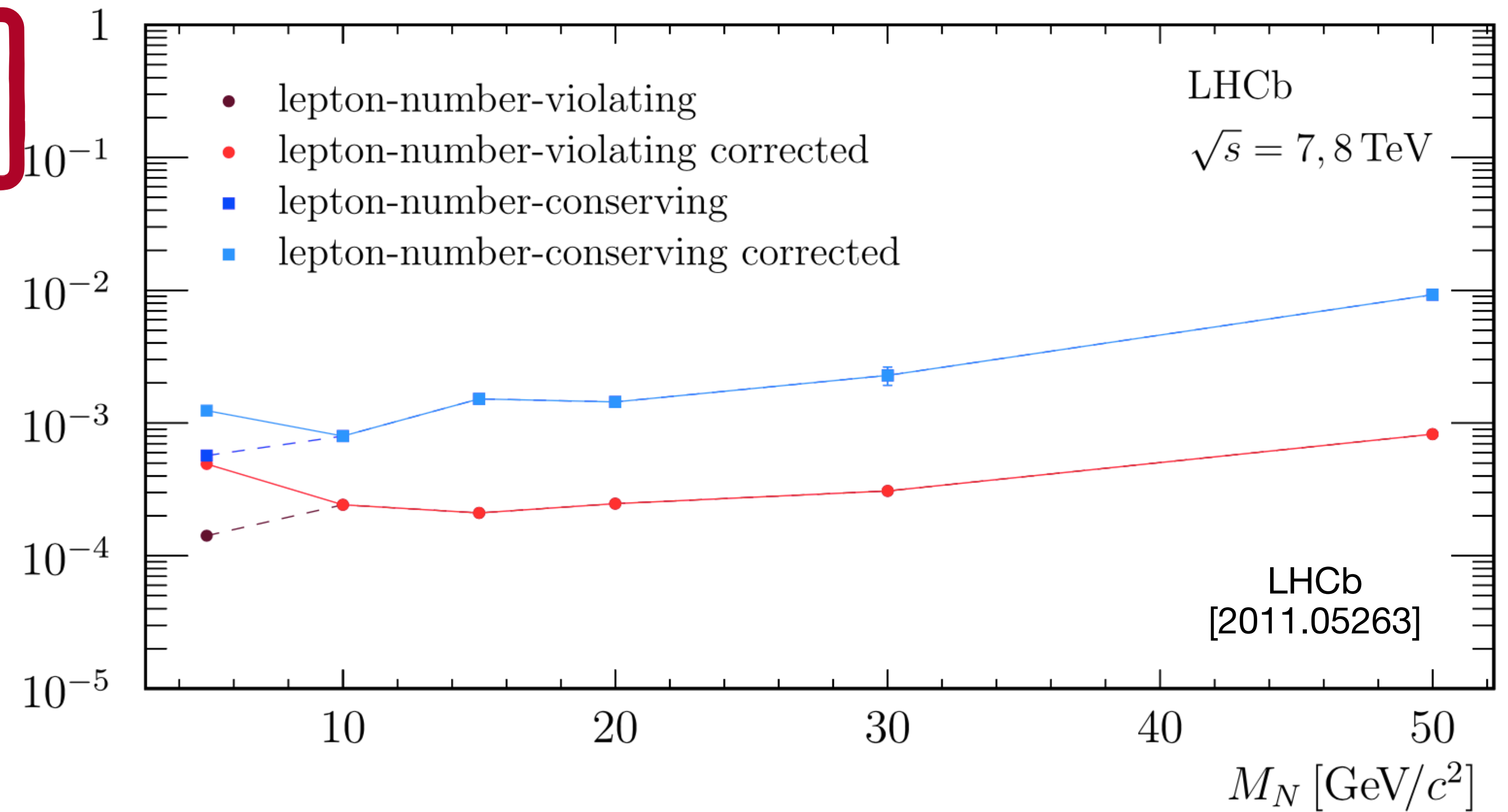
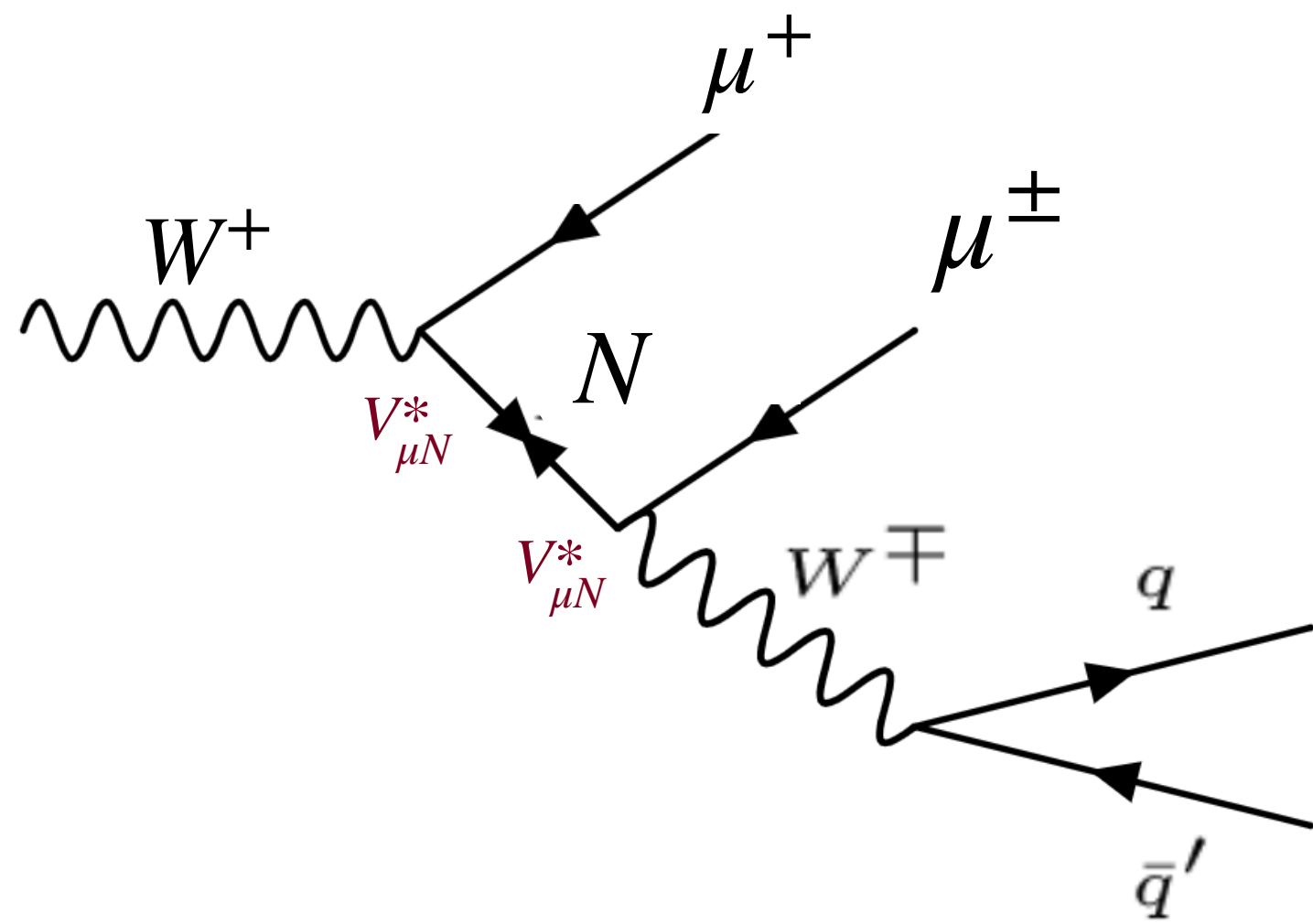
$$V = \left(\begin{array}{c|cc} & V_{eN_1} & V_{eN_2} \\ \tilde{V}_{PMNS} & V_{\mu N_1} & V_{\mu N_2} \\ \hline & V_{\tau N_1} & V_{\tau N_2} \end{array} \right)$$

2 leptons + jets: 2 HNLs (single mixing)

$$V = \begin{pmatrix} \tilde{V}_{PMNS} & \begin{matrix} V_{eN_1} & V_{eN_2} \\ V_{\mu N_1} & V_{\mu N_2} \\ V_{\tau N_1} & V_{\tau N_2} \end{matrix} \end{pmatrix}, \quad \text{OK if NO Interference!}$$

$|V_{\mu N_1}|^2 + |V_{\mu N_2}|^2$

$|V_{\mu N}|^2$



2 leptons + jets: 2 HNLs Interfering

Defining $V_{\ell_\alpha N_j} = |V_{\ell_\alpha N_j}| e^{i\phi_{\alpha j}}$, and assuming $|V_{\ell N_1}|^2 \simeq |V_{\ell N_2}|^2, M_1 \simeq M_2 \equiv M, \Delta M_{12} \neq 0, \Gamma_1 \simeq \Gamma_2 \equiv \Gamma$

$$\Gamma(W^+ \rightarrow \ell_\alpha^+ \ell_\beta^\pm q \bar{q}')|_{N_1 \& N_2} = \Gamma(W^+ \rightarrow \ell_\alpha^+ \ell_\beta^\pm q \bar{q}')|_{N_1} \times 2 \mathcal{K}(y, \delta\phi^\pm) !$$

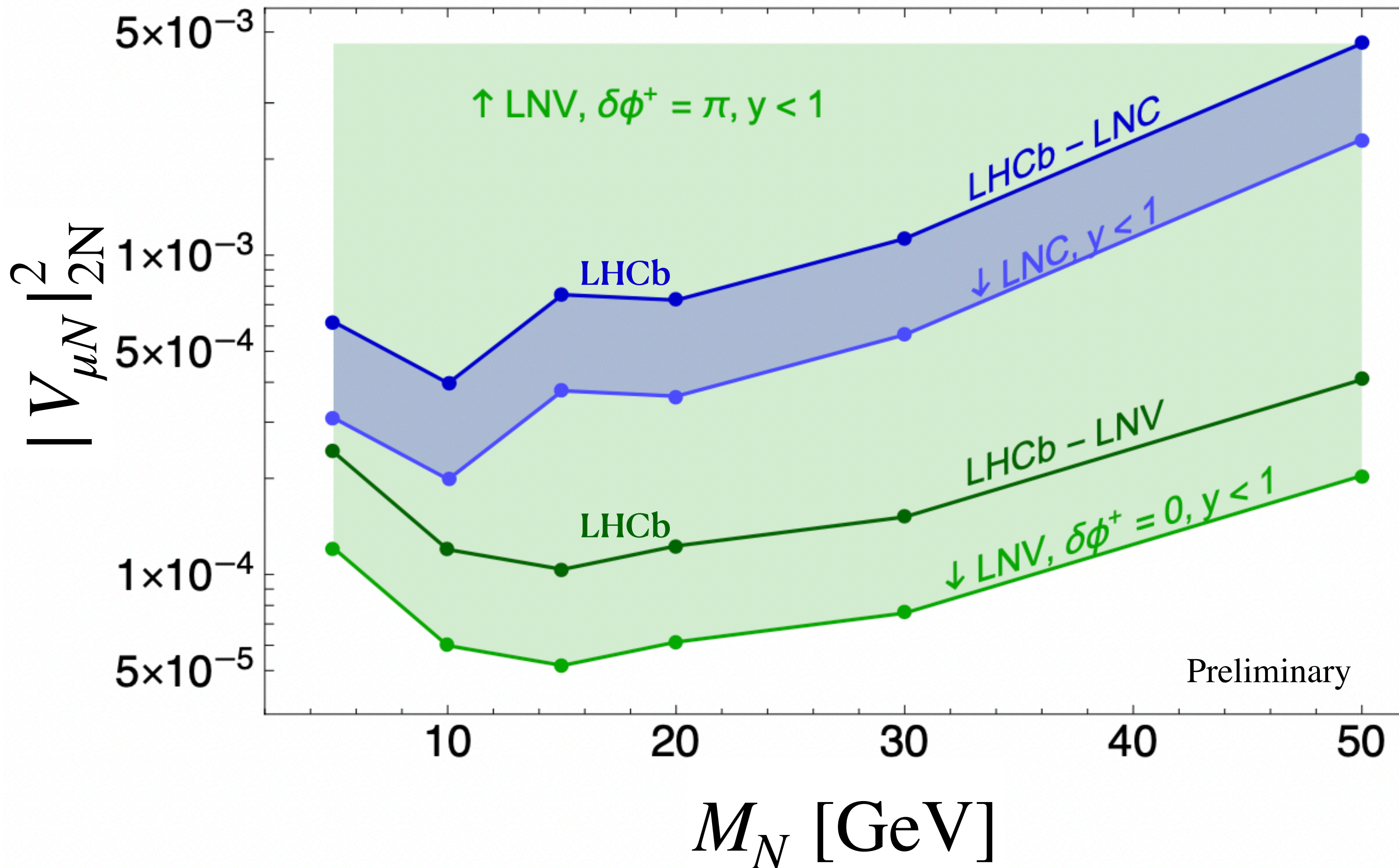
$$\mathcal{K}(y, \delta\phi^\pm) = \left(1 + \cos \delta\phi^\pm \frac{1}{1 + y^2} - \sin \delta\phi^\pm \frac{y}{1 + y^2} \right)$$

$$y \equiv \frac{\Delta M_{12}}{\Gamma}, \quad \delta\phi^\pm = (\phi_{\alpha 2} - \phi_{\alpha 1}) \pm (\phi_{\beta 2} - \phi_{\beta 1}),$$

$$\delta\phi^- = 0 \text{ for } \alpha = \beta$$

2 leptons + jets: 2 HNLs (single mixing)

$$|V_{\mu N_1}|_{2N}^2 = |V_{\mu N}|_{1N}^2 / [2\mathcal{K}(y, \delta\phi)] \quad \mathcal{K}(y, \delta\phi^\pm) = \left(1 + \cos \delta\phi^\pm \frac{1}{1+y^2} - \sin \delta\phi^\pm \frac{y}{1+y^2} \right)$$



$$y \equiv \frac{\Delta M_{12}}{\Gamma}, \quad \delta\phi^\pm = (\phi_{\alpha 2} - \phi_{\alpha 1}) \pm (\phi_{\beta 2} - \phi_{\beta 1})$$

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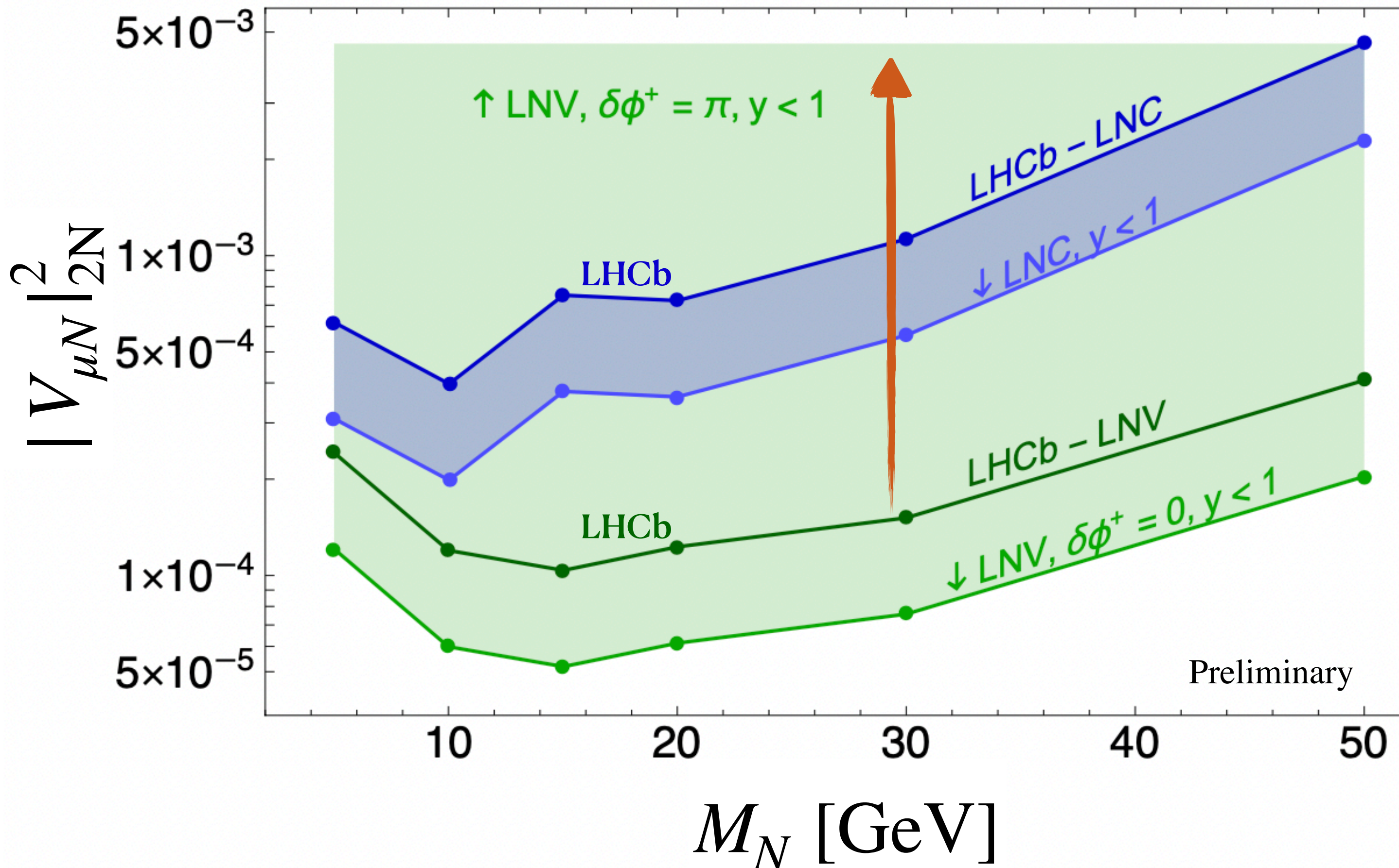
2 leptons + jets: 2 HNLs (single mixing)

$$V = \begin{pmatrix} \tilde{V}_{PMNS} & \begin{matrix} V_{eN_1} & V_{eN_2} \\ V_{\mu N_1} & V_{\mu N_2} \\ \dots & \dots \\ V_{\tau N_1} & V_{\tau N_2} \end{matrix} \end{pmatrix}$$

$$|V_{\mu N_1}|_{2N}^2 = |V_{\mu N}|_{1N}^2 / [2\mathcal{K}(y, \delta\phi)]$$

$$\mathcal{K}(y, \delta\phi^\pm) = \left(1 + \cos \delta\phi^\pm \frac{1}{1+y^2} - \sin \delta\phi^\pm \frac{y}{1+y^2} \right)$$

$$y \equiv \frac{\Delta M_{12}}{\Gamma}, \quad \delta\phi^\pm = (\phi_{\alpha 2} - \phi_{\alpha 1}) \pm (\phi_{\beta 2} - \phi_{\beta 1})$$



It is possible to **relax the LNV** bound.
 (N_1, N_2 PseudoDirac Pair \rightarrow LNV forbidden)

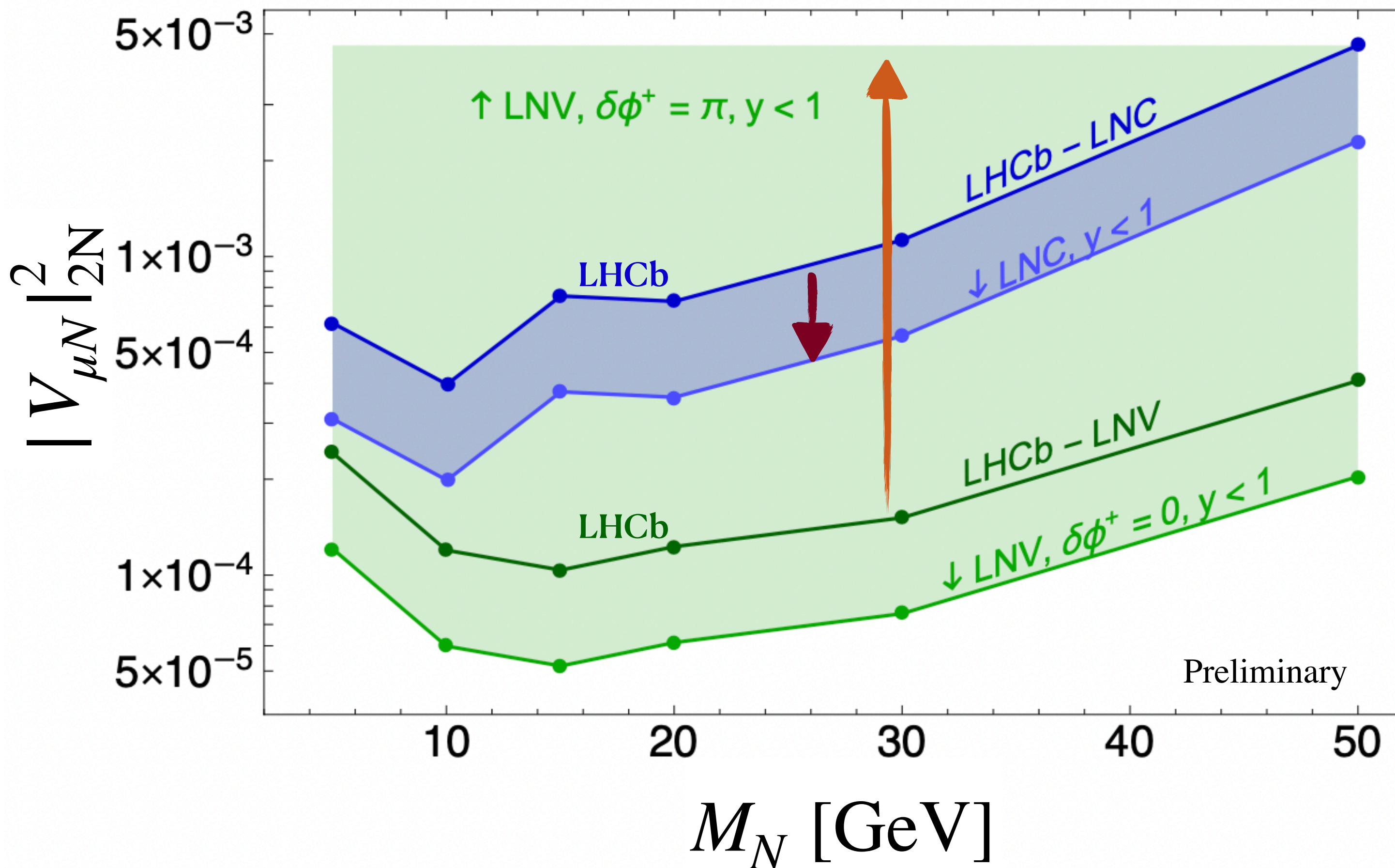
2 leptons + jets: 2 HNLs (single mixing)

$$V = \begin{pmatrix} \tilde{V}_{PMNS} & \begin{matrix} V_{eN_1} & V_{eN_2} \\ V_{\mu N_1} & V_{\mu N_2} \\ V_{\tau N_1} & V_{\tau N_2} \end{matrix} \end{pmatrix}$$

$$|V_{\mu N_1}|_{2N}^2 = |V_{\mu N}|_{1N}^2 / [2\mathcal{K}(y, \delta\phi)]$$

$$\mathcal{K}(y, \delta\phi^\pm) = \left(1 + \cos \delta\phi^\pm \frac{1}{1+y^2} - \sin \delta\phi^\pm \frac{y}{1+y^2} \right)$$

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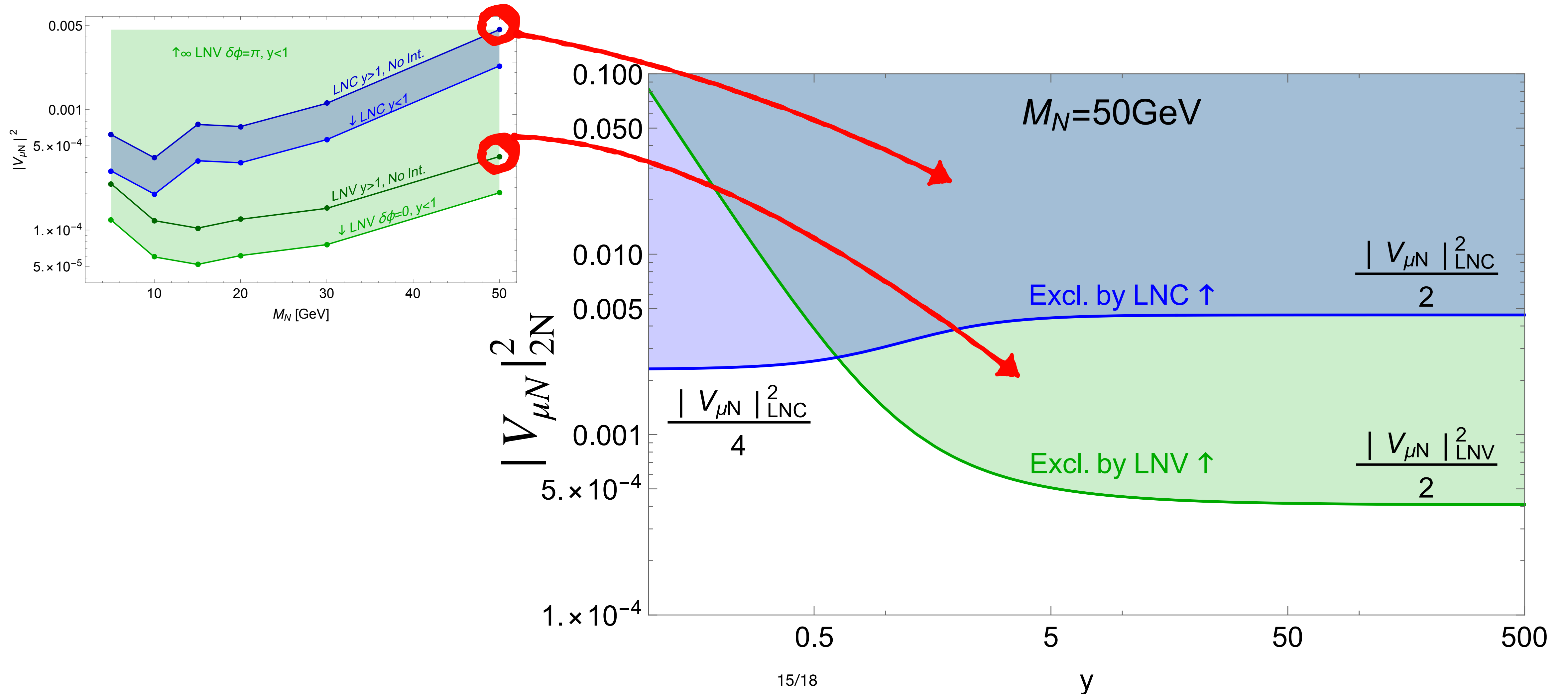


It is possible to **relax the LNV** bound.
(N_1, N_2 PseudoDirac Pair \rightarrow LNV forbidden)

However, for $y < 1$ **LNC bound stronger**

Combining LNV & LNC, What is the maximum allowed value for $|V_{\mu N}|^2$?

We can maximise LNC and LNV bounds over $[y, \delta\phi]$.

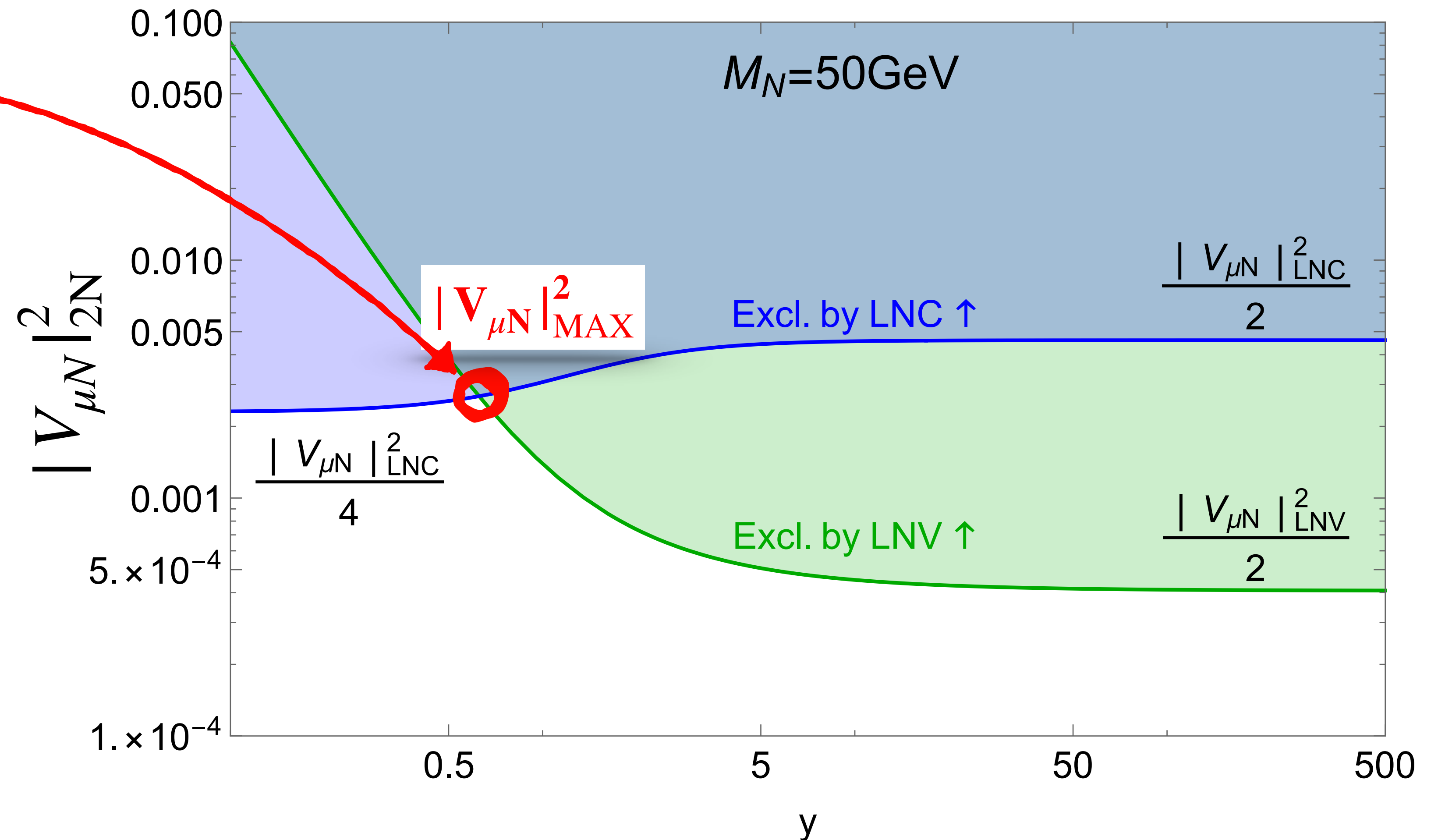


How to get a conservative bound?

We can maximise LNC and LNV bounds over $[y, \delta\phi]$.

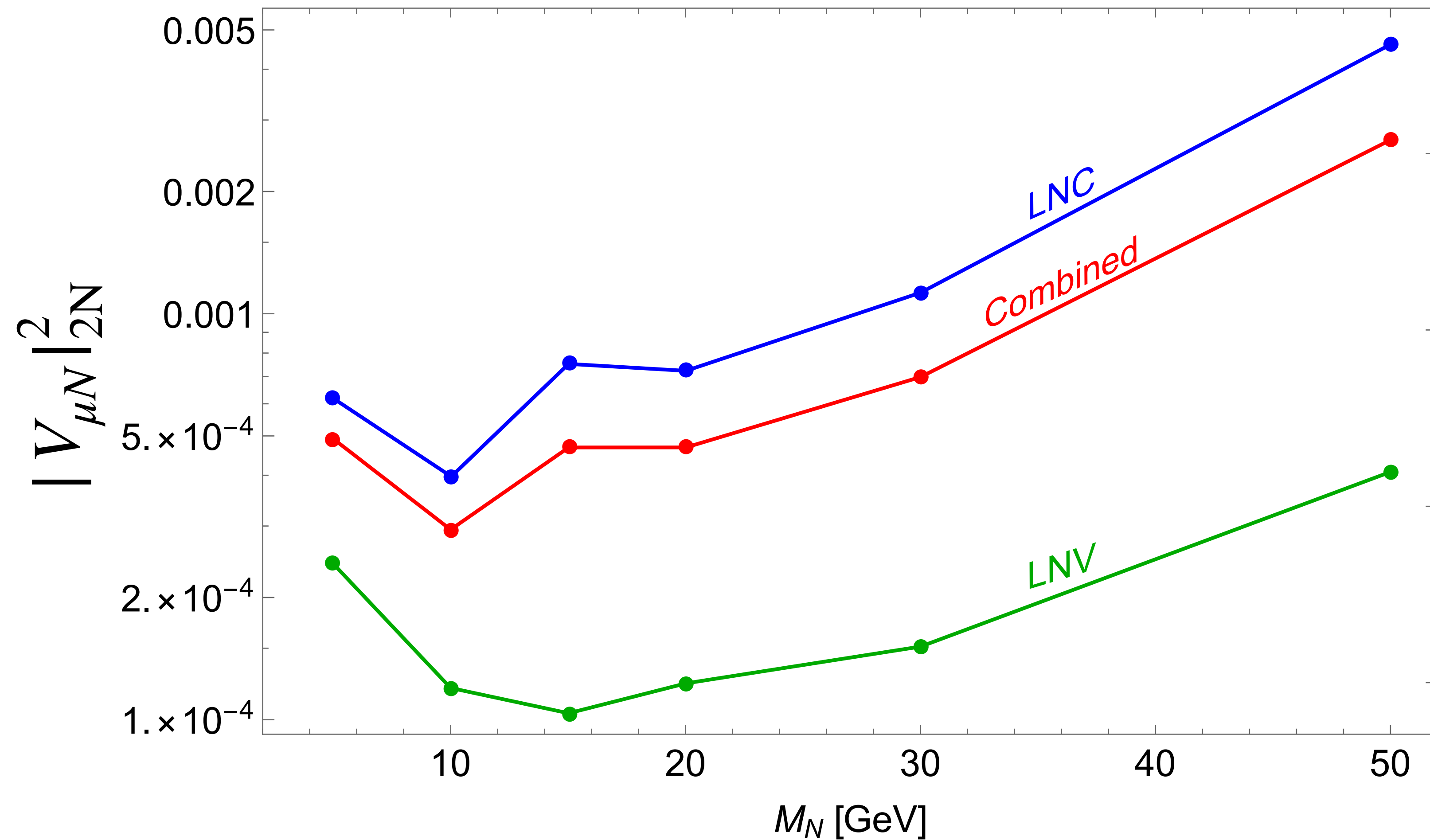
Combined
conservative
bound.

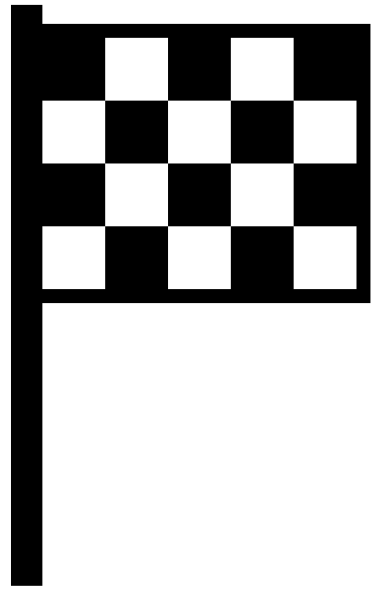
Impossible to relax
the bound above
this value, even
with interferences.



Combined bound with interferences

Identifying $|V_{\mu N}|_{\text{MAX}}^2$ for each M_N , we get:





Conclusions

- Bounds on HNLs in the simplified scenarios are often **over-constraining** if naively applied to realistic models;
- Actual bounds are **model and benchmark dependent**, and must be recast;
- It is crucial for experiments to perform **BOTH** LNV & LNC searches to get combined bounds.

Thank you!

This project has received funding from the European Union's Horizon 2020 research and innovation programme under the Marie Curie grant No 860881-HIDDeN

Backup

Outlook: Inferring details on Low-Scale SeeSaw

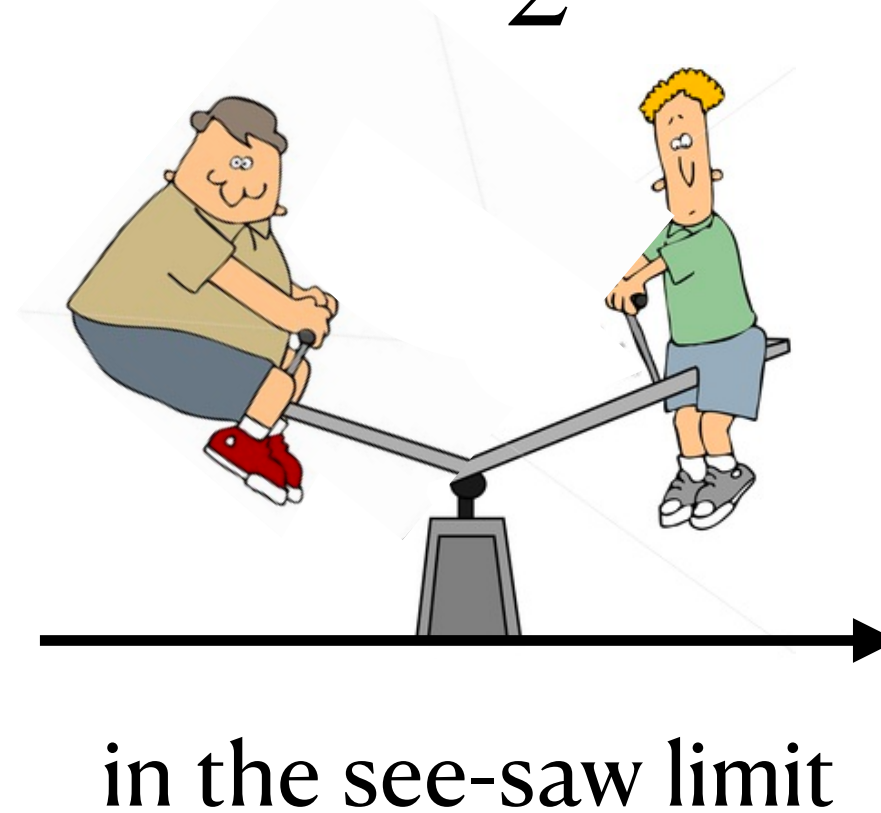
Our approximations are justified in the L-S ss

Degeneracy N_4 & N_5 : $M_4 \simeq M_5$, $\Gamma_4 \simeq \Gamma_5$

$$\mathcal{L}_{LS} = -m_D \bar{\nu}_R \nu_L - M \bar{\nu}_R^c \nu_s - \frac{1}{2} \mu_S \bar{\nu}_s \nu_s^c - \frac{1}{2} \mu_R \bar{\nu}_R \nu_R^c + \text{h.c.}$$

$$\nu = \begin{pmatrix} \nu_L \\ \nu_R^c \\ \nu_s^c \end{pmatrix}$$

$$M_{LS} = \begin{pmatrix} 0 & m_D & 0 \\ m_D^T & \mu_R & M \\ 0 & M & \mu_S \end{pmatrix}$$



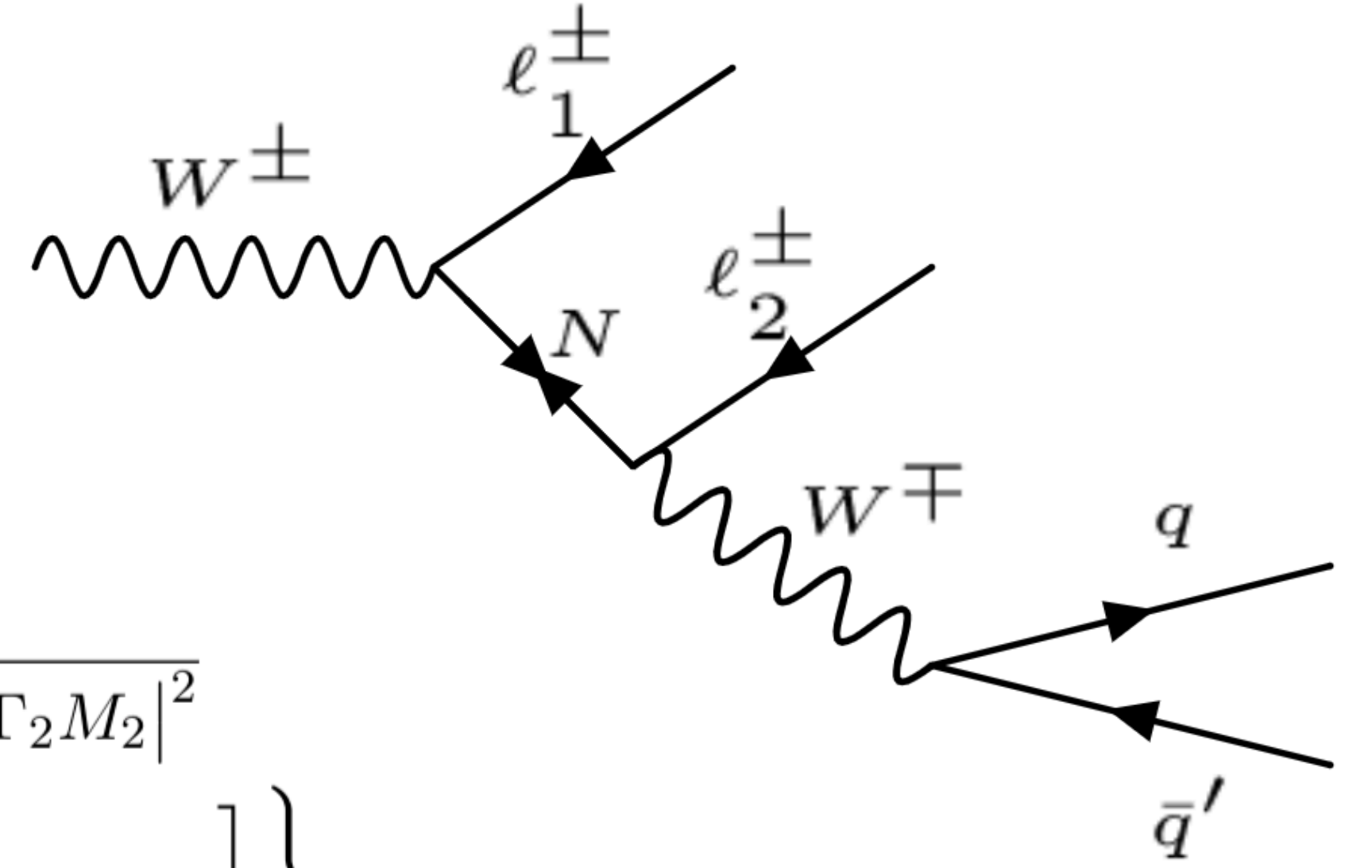
$$M_{\text{light}} \approx \mu \frac{m_D^2}{M^2}$$

$$M_{N_4, N_5} \approx \pm M + \mathcal{O}(\mu)$$

N_4, N_5 almost degenerate and with opposite CP phases $\rightarrow \Delta M \ll M, \delta\phi_V = \pi$ ✓

LNV & LNC searches can test LowScale SS models

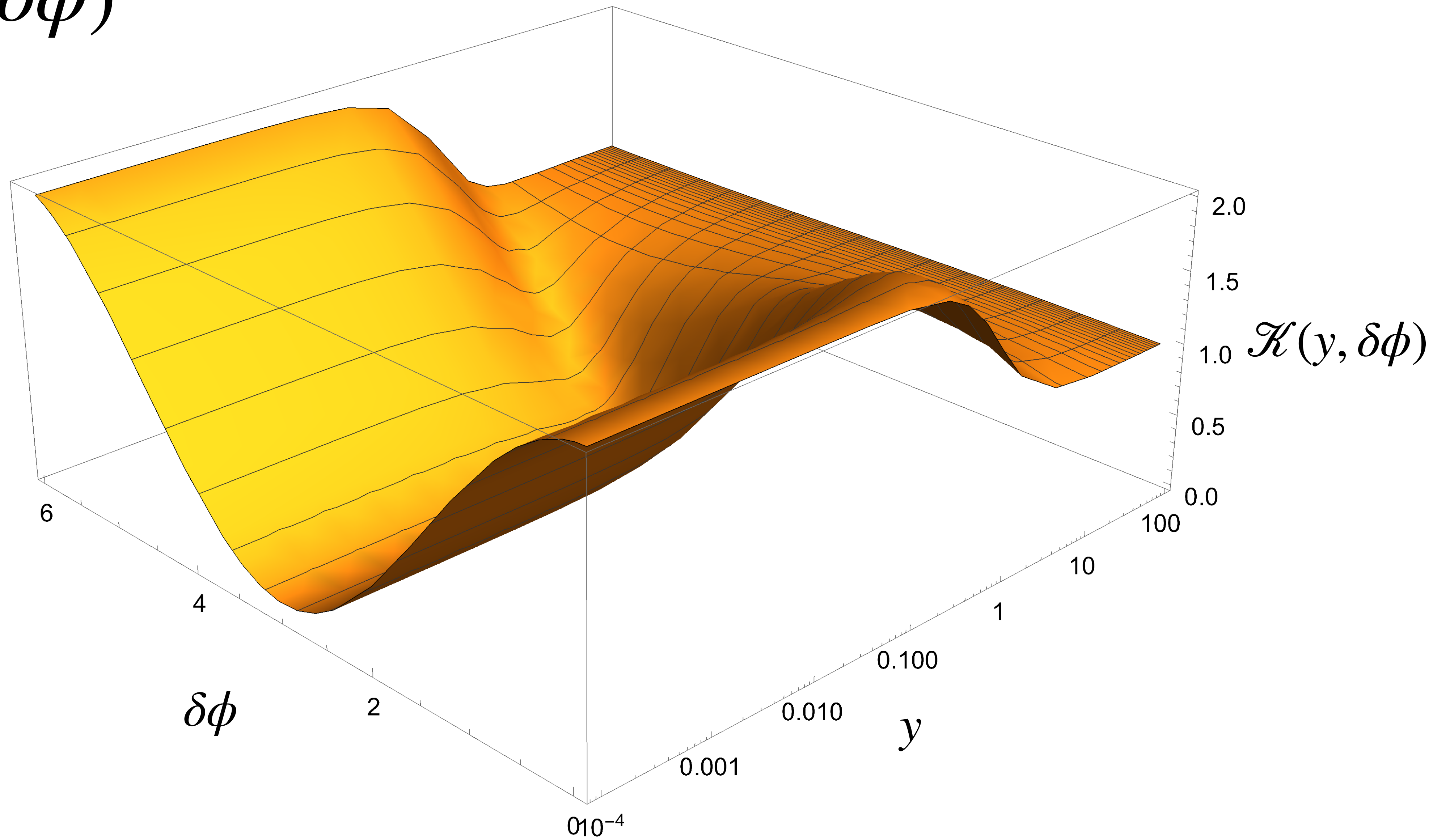
$$\begin{aligned}
|\overline{M}|^2 = & \frac{1}{16} \left[\frac{g^3}{2\sqrt{2}M_W^2} \right]^2 16 (p_{\ell_2} \cdot p_q) (2E_{\ell_1} E_{q'} + p_{\ell_1} \cdot p_{q'}) \left\{ \right. \\
& |U_{\ell_1 N_1}|^2 |U_{\ell_2 N_1}|^2 \frac{M_1^2}{|p_N^2 - M_1^2 + i\Gamma_1 M_1|^2} + |U_{\ell_1 N_2}|^2 |U_{\ell_2 N_2}|^2 \frac{M_2^2}{|p_N^2 - M_2^2 + i\Gamma_2 M_2|^2} \\
& \left. + 2 \operatorname{Re} \left[|U_{\ell_1 N_1}| |U_{\ell_1 N_2}| |U_{\ell_2 N_1}| |U_{\ell_2 N_2}| e^{i\delta\phi} \frac{M_1 M_2}{(p_N^2 - M_1^2 + i\Gamma_1 M_1)(p_N^2 - M_2^2 - i\Gamma_2 M_2)} \right] \right\}
\end{aligned}$$



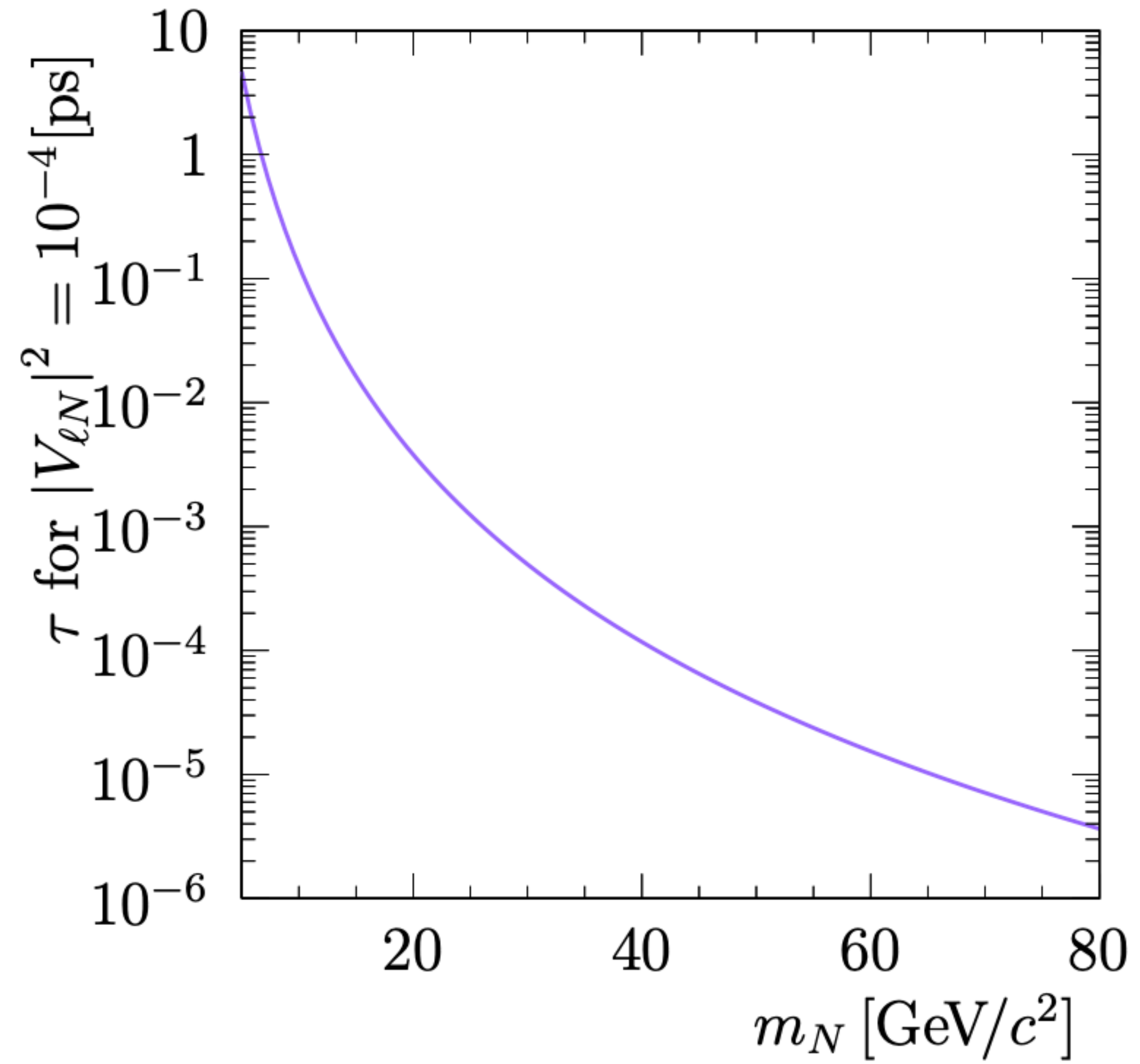
$$\begin{aligned}
|\overline{M}|^2 \simeq & \left[\frac{g^3}{2\sqrt{2}M_W^2} \right]^2 \pi (p_{\ell_2} \cdot p_q) (2E_{\ell_1} E_{q'} + p_{\ell_1} \cdot p_{q'}) \delta(p_N^2 - M^2) \frac{M}{\Gamma} |U_{\ell_1 N_1}|^2 |U_{\ell_2 N_1}|^2 \left\{ \right. \\
& \left. 1 + \xi^2 + 4\xi \left[2 \cos \delta\phi \frac{M^2 \Gamma^2}{(\Delta M^2)^2 + 4\Gamma^2 M^2} - \sin \delta\phi \frac{M\Gamma \Delta M^2}{(\Delta M^2)^2 + 4\Gamma^2 M^2} \right] \right\}
\end{aligned}$$

with $\xi = |U_{\ell_1 N_2}| |U_{\ell_2 N_2}| / |U_{\ell_1 N_1}| |U_{\ell_2 N_1}|$

$$\mathcal{K}(y, \delta\phi)$$



N lifetime



LHCb
[2011.05263]