

Measuring the distortion of time at cosmological scales

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Current status

- ◆ Mystery to solve: what is causing the **acceleration** of our Universe
 - Cosmological constant
 - Dynamical dark energy
 - Theory of gravity
- ◆ Plethora of new **data**: DESI, Euclid, SKA, LSST
- ◆ **Goal**: test the validity of Λ CDM and of General Relativity

Testing General Relativity

What do we want to **test**?

◆ At late time our Universe is described as:

Homogeneous and **isotropic** background + **fluctuations**

Described by 4 fields

◆ Perturbations in the **geometry**

gravitational potentials

$$ds^2 = -a^2 \left[(1 + 2\Psi) d\eta^2 + (1 - 2\Phi) \delta_{ij} dx^i dx^j \right]$$

◆ Perturbations in the universe's **content**:

density fluctuations $\delta = \frac{\delta\rho}{\rho}$
peculiar velocity V

Testing General Relativity

- ◆ General Relativity provides **relations** between the **fields**

δ	Continuity	V
Poisson		Euler
Φ	Anisotropic stress	Ψ

- ◆ Ideally, we want to **measure** the 4 fields and **compare** them
- ◆ Currently not possible: we have **only 3** measurements
 - δ and V from the distribution of galaxies
 - $\Phi + \Psi$ from gravitational lensing

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$$\begin{array}{ccc} \delta & \text{Continuity} & V \\ \text{Poisson} & & \text{Euler} \\ \Phi & = & \Psi \\ & \Lambda\text{CDM} & \end{array}$$

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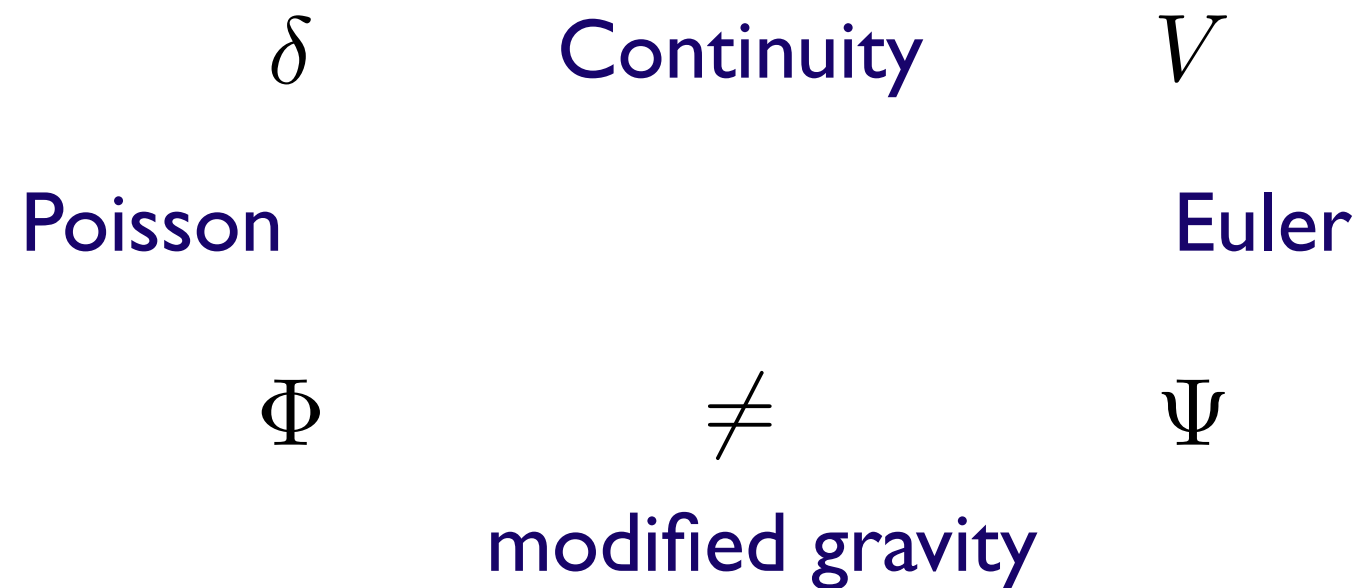
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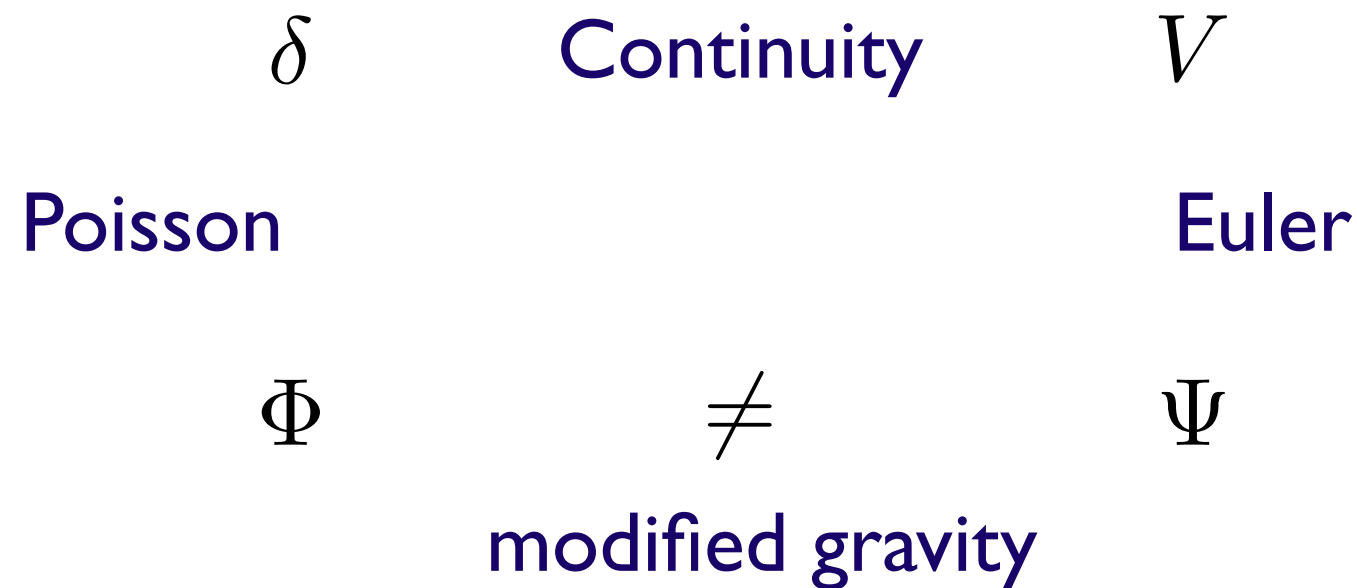
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Testing General Relativity

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We cannot test 4 relations with 3 observables

- δ and V from the distribution of galaxies
- $\Phi + \Psi$ from gravitational lensing

Current methods

- ◆ Assume that **Euler** and **continuity** equations are **valid**



- ◆ Measure $V \rightarrow \Psi$ **compare** with $\Phi + \Psi$
- ◆ Evolution equation for growth of structure \rightarrow constrain **Poisson** equation with galaxy clustering

Results consistent with General Relativity

- ◆ This is valid only if Euler and continuity eqs. are valid:
restrictive assumption for **dark matter**

Equivalence principle

- ◆ Euler equation encodes the **weak equivalence principle**
- ◆ It tells us that all objects fall in the same way in a gravitational potential
- ◆ It has been precisely **tested** for standard matter and photons, but not for dark matter
- ◆ If we want to test gravity in a **model-independent** way, we should **relax** the assumption that dark matter obeys the weak equivalence principle

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We cannot test **Poisson** equation

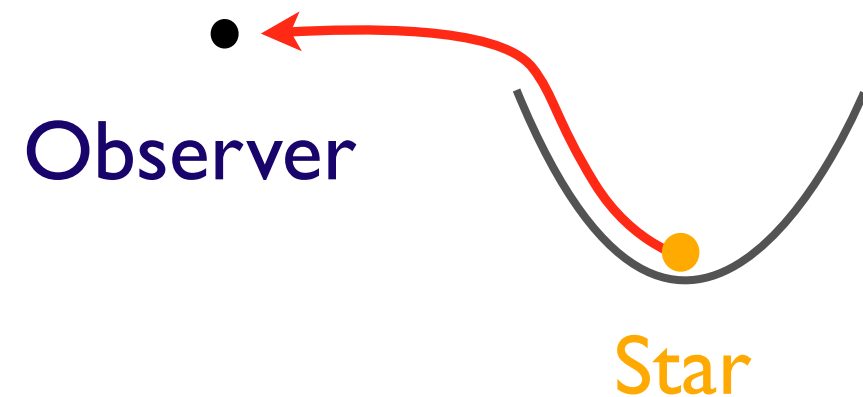
We cannot **compare** Φ and Ψ

Outline

- ◆ **Way out:** we can measure the **distortion** of **time** Ψ with future galaxy survey
 - Additional observable to test gravity
- ◆ **Method** to measure time distortion with galaxy clustering
- ◆ **Forecasts** with SKA
- ◆ Tests of gravity: weak **equivalence** principle

Distortion of time

- ◆ We measure the **redshift** of photons escaping a gravitational potential
- ◆ First test done with **white dwarfs** in 1925



Gravitational redshift

- ◆ Here: same test but at **cosmological** distances, ~ 50 Mpc
- ◆ **Method:** gravitational redshift modifies the observed **clustering** of galaxies

Distortion of time

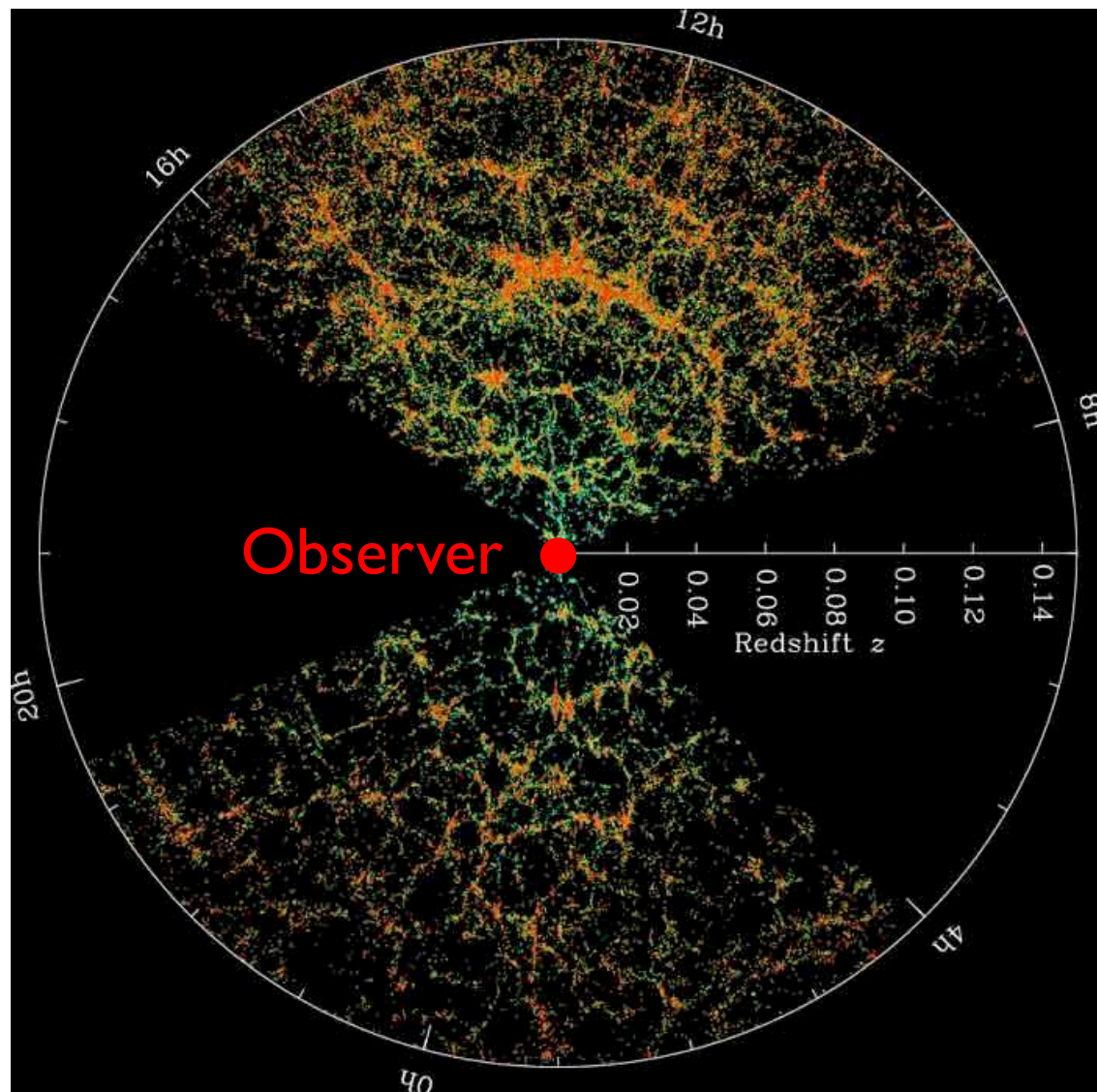
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Galaxy clustering

Credit: M. Blanton, SDSS



- ◆ Redshift as indicator of **distance**

Distant galaxies more affected by expansion

- ◆ **Gravitational redshift** is present

→ slight **distortion** in maps due to the gravitational potentials

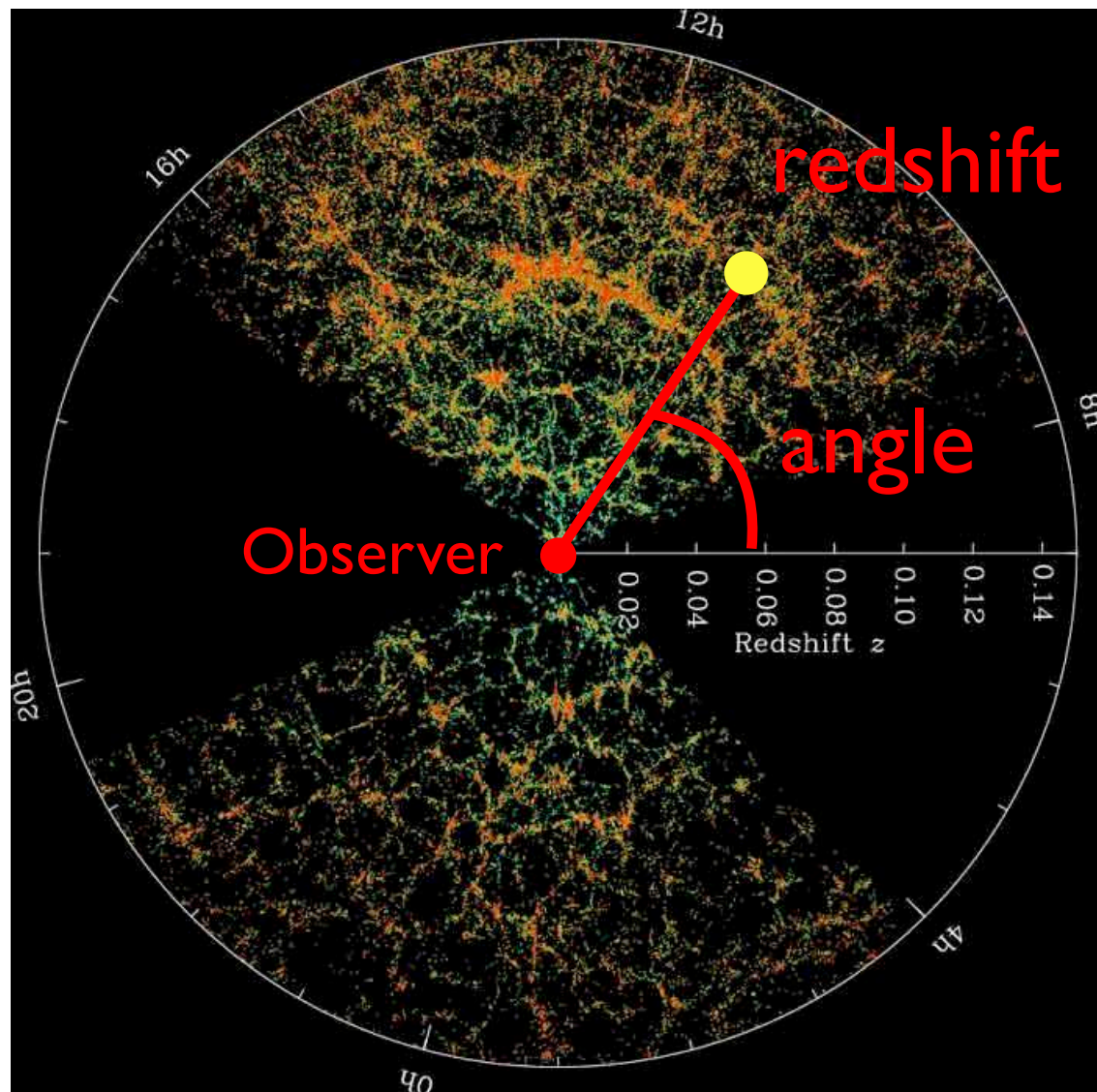
- ◆ Fluctuations in number of galaxies $\Delta = \frac{N - \bar{N}}{\bar{N}}$

$\Delta =$ intrinsic distortions + gravitational redshift

- ◆ How do we **isolate** gravitational redshift?

Galaxy clustering

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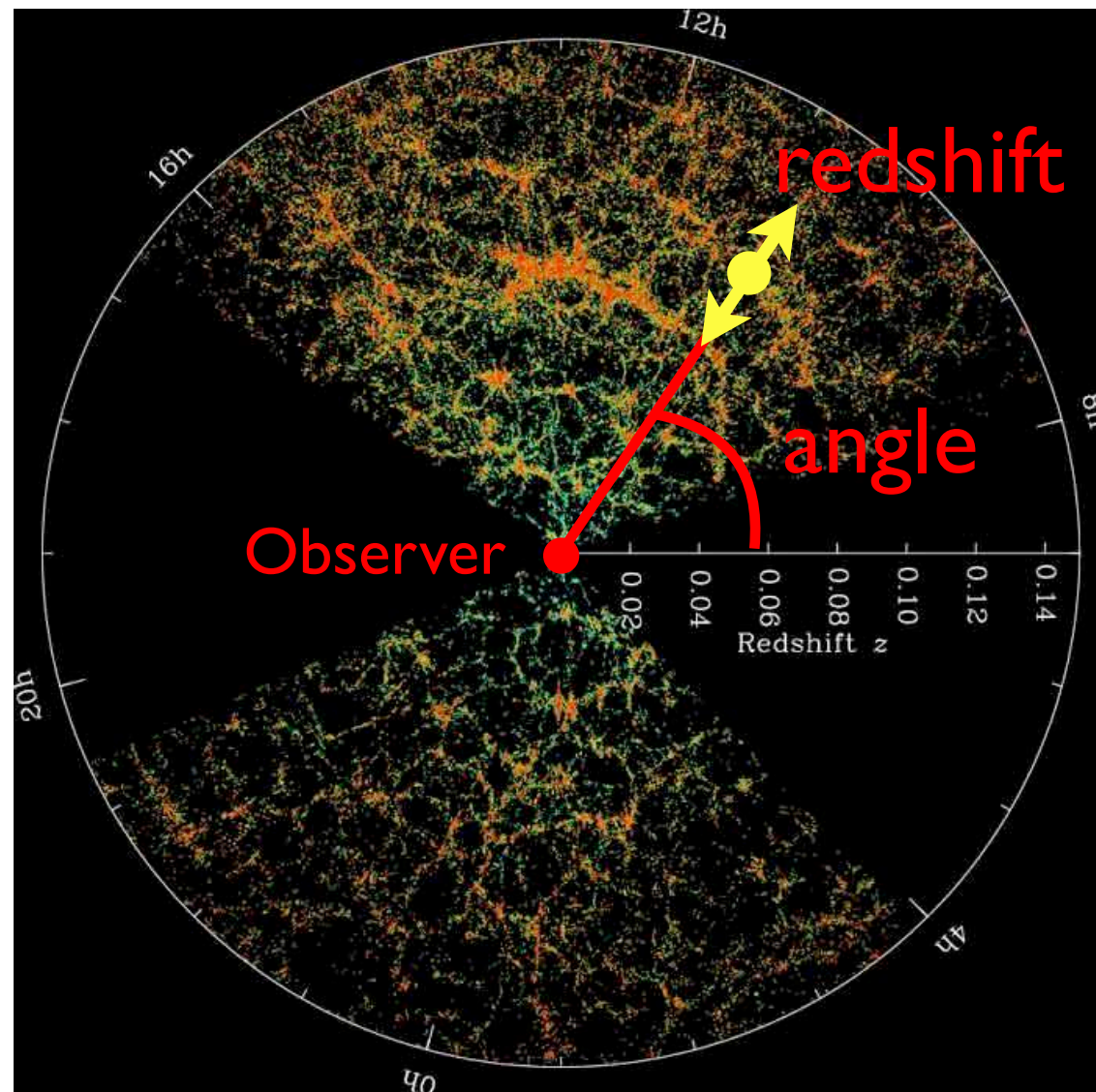
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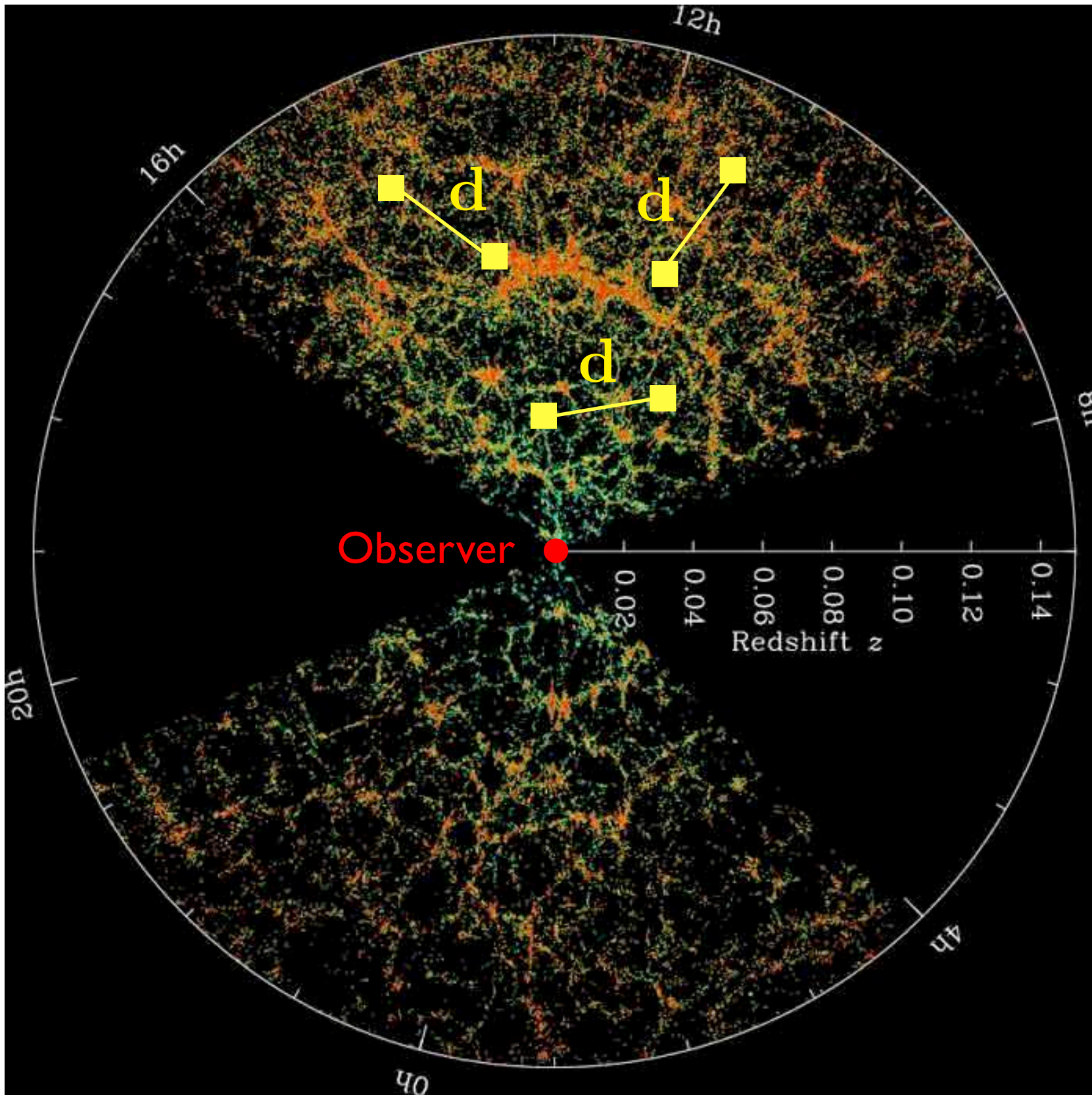
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Correlation function

Credit: M. Blanton, SDSS



$$\langle \Delta(\mathbf{x}) \Delta(\mathbf{x}') \rangle$$

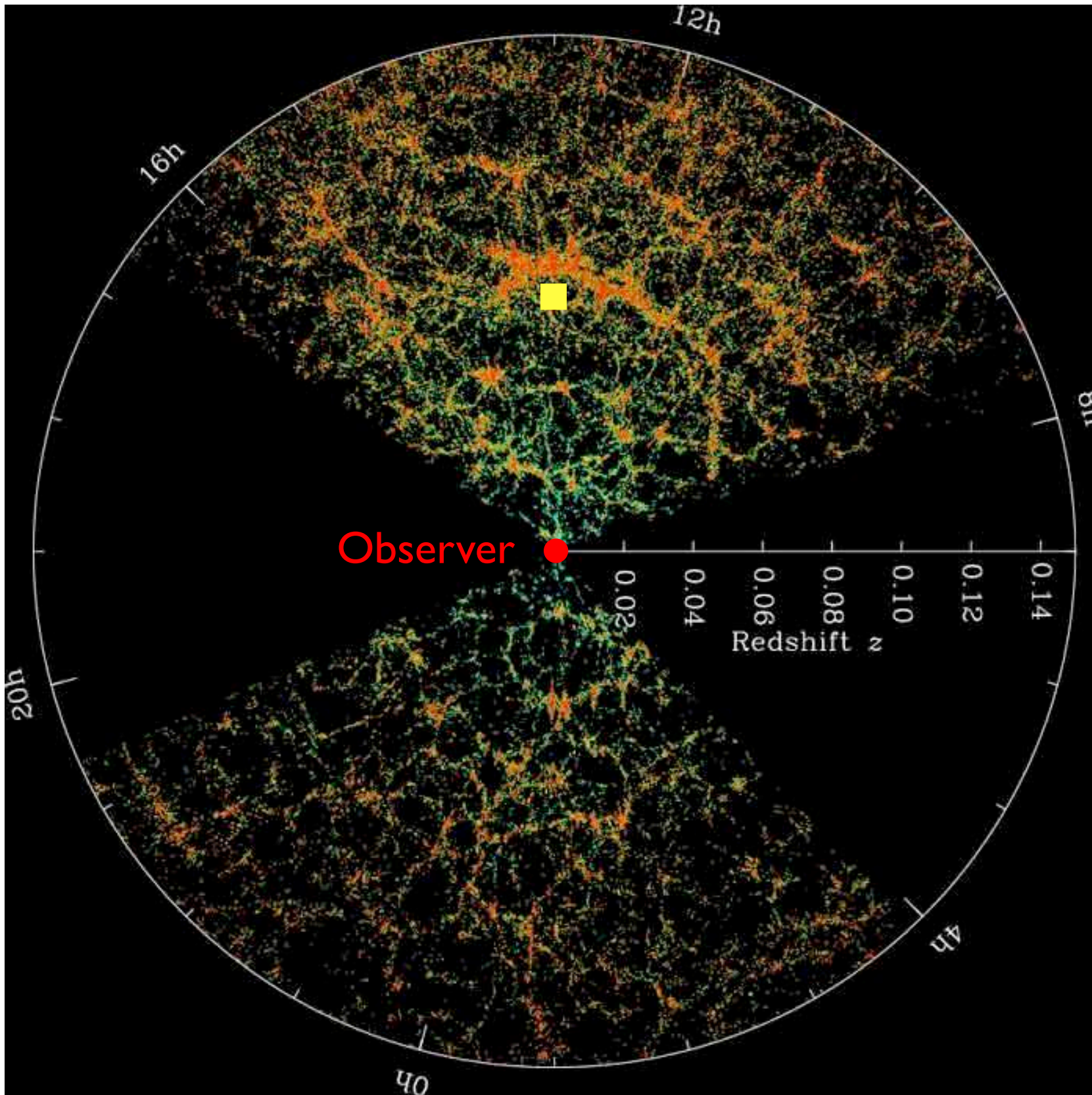


Different pixels

- ◆ Matter fluctuations generate **isotropic** correlations
- ◆ Gravitational redshift **breaks** the **symmetry**

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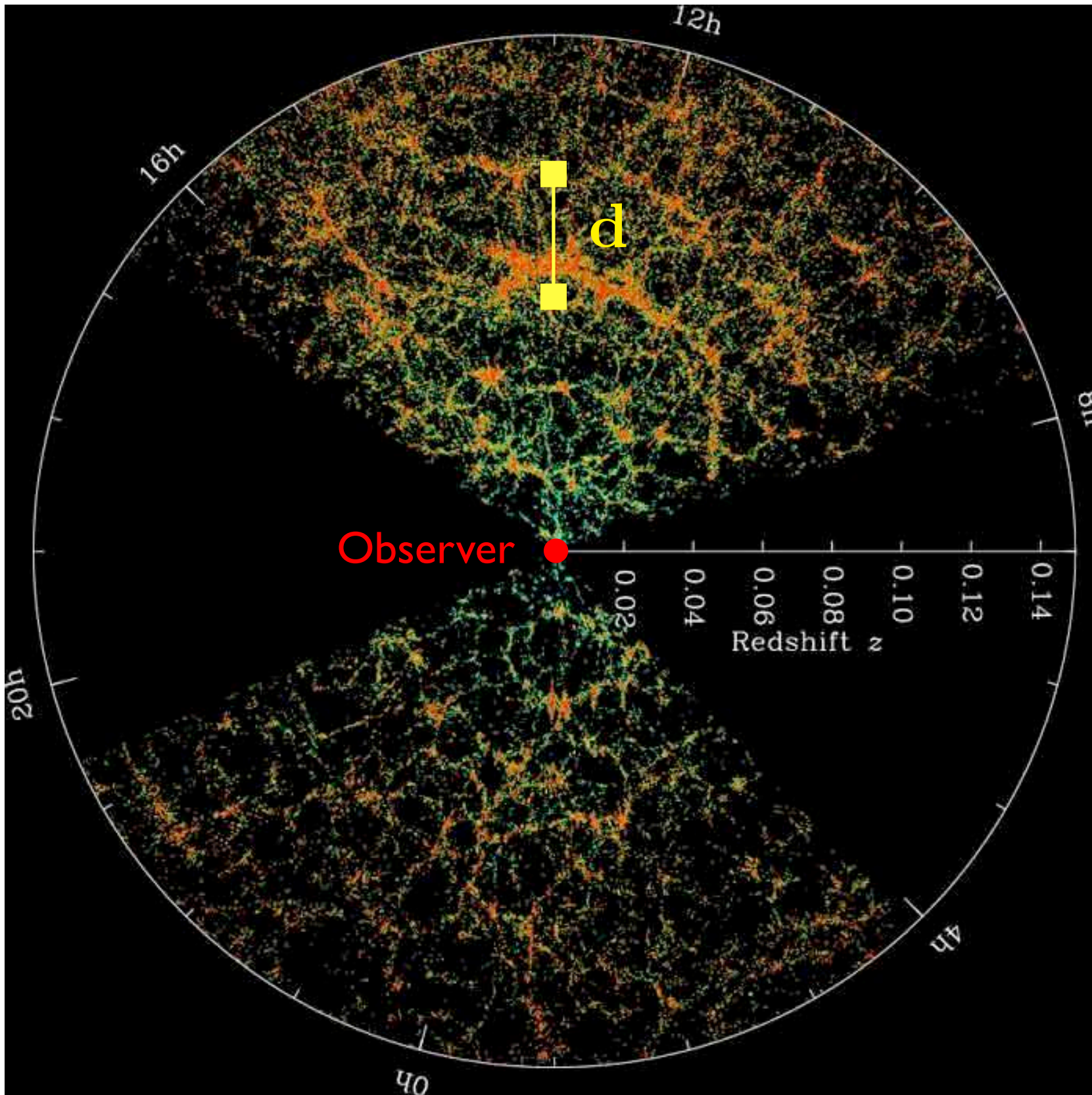


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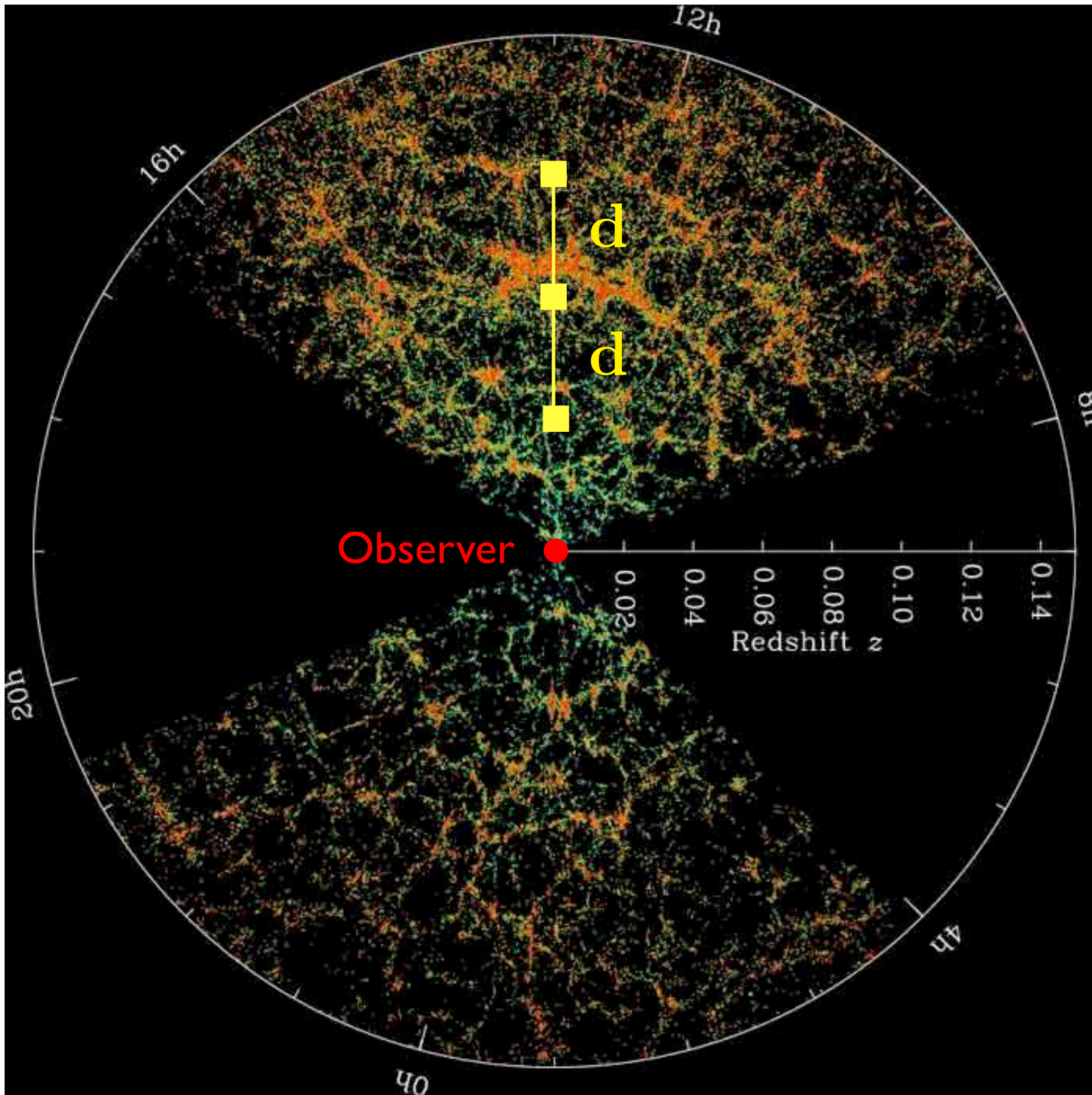


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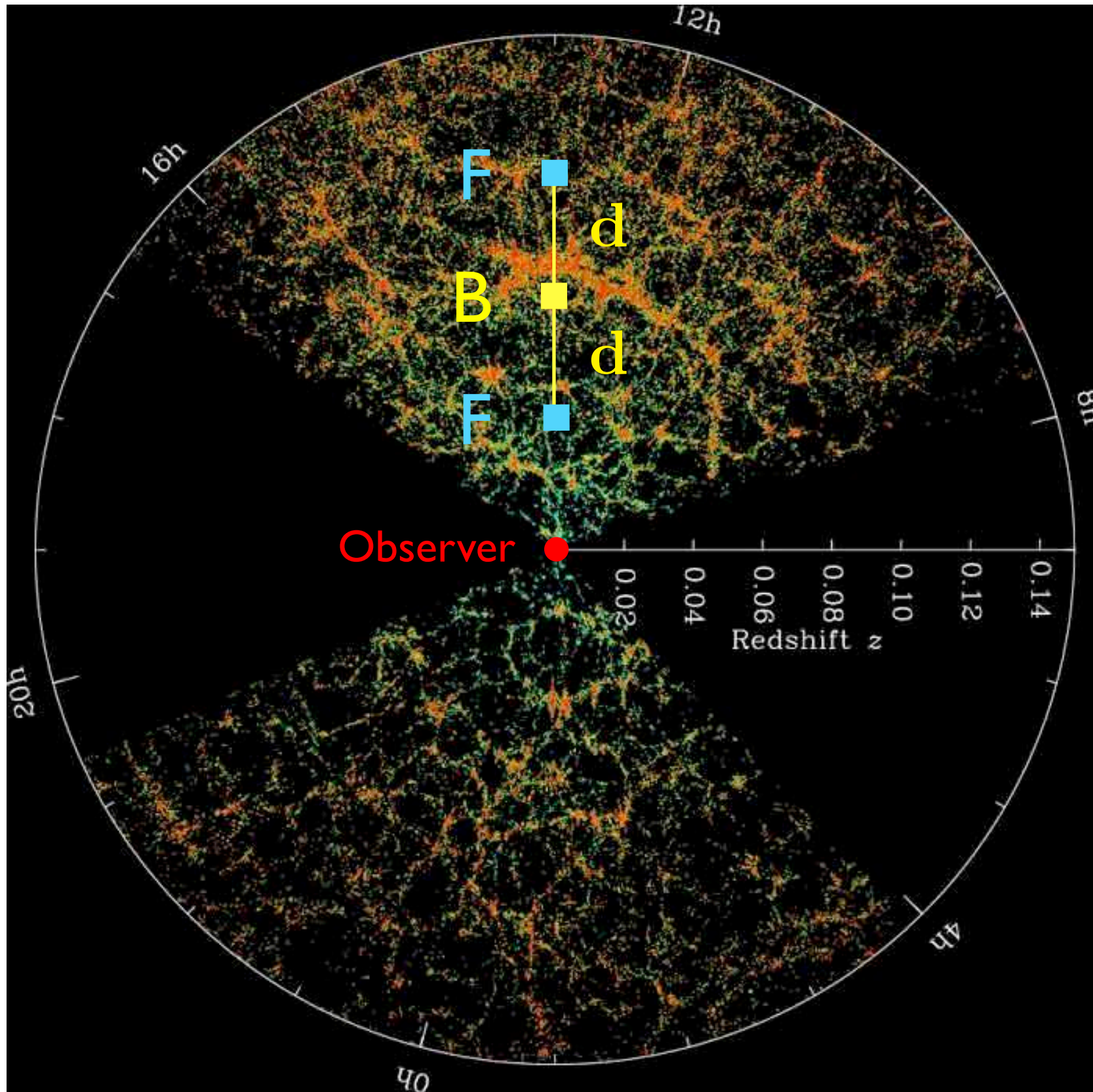


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Two populations

Credit: M. Blanton, SDSS



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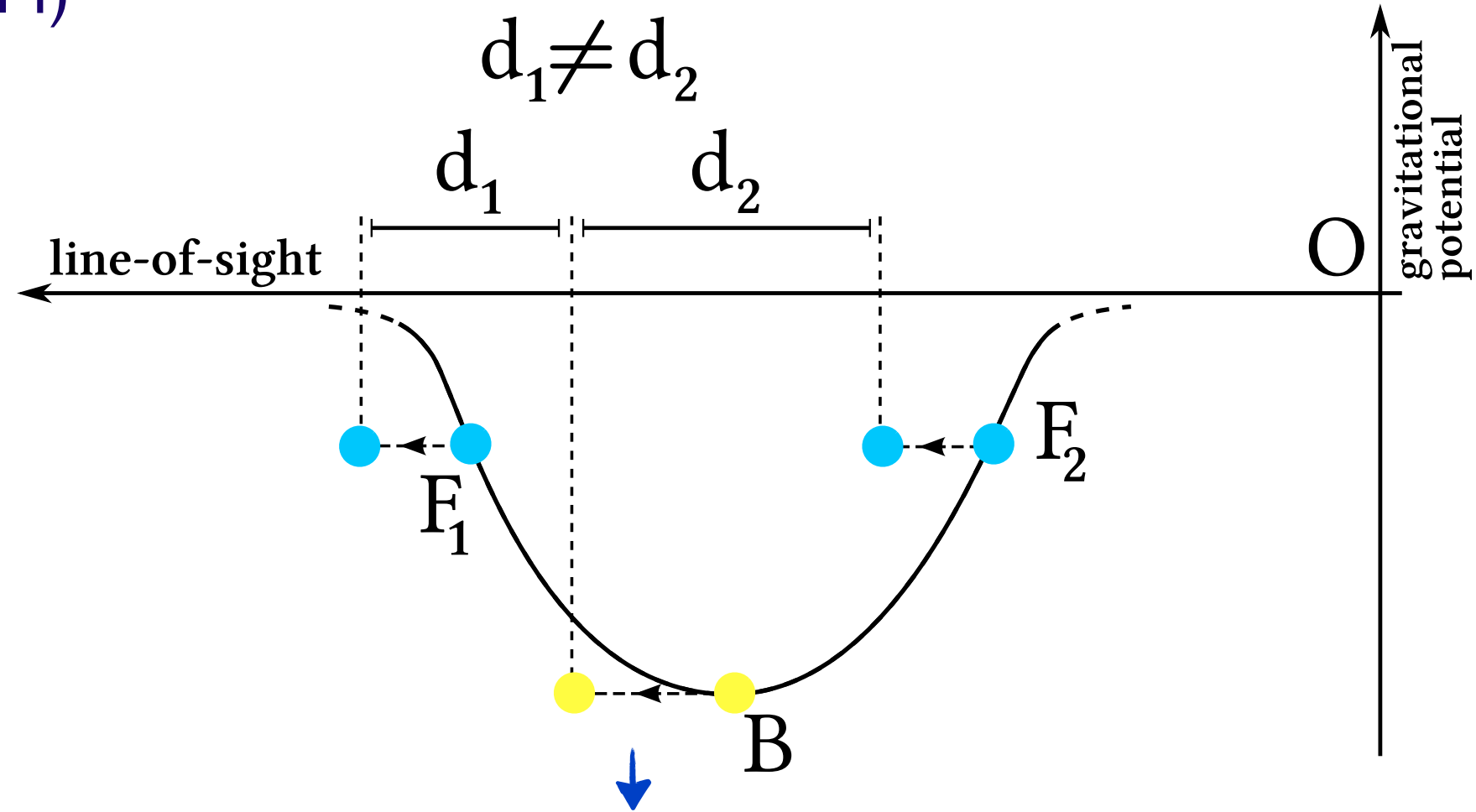


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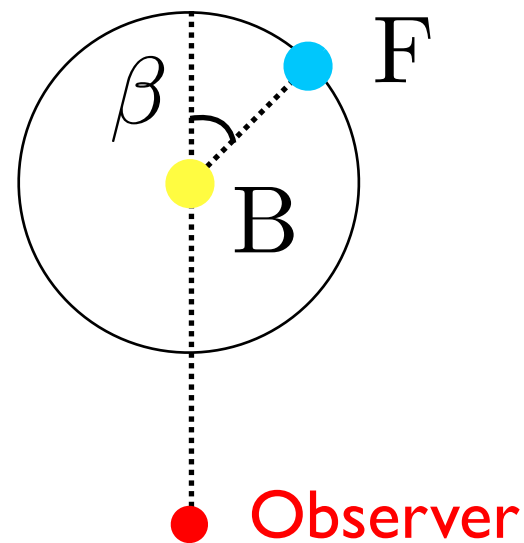
Breaking of symmetry from gravitational redshift

CB, Hui & Gaztanaga (2014)



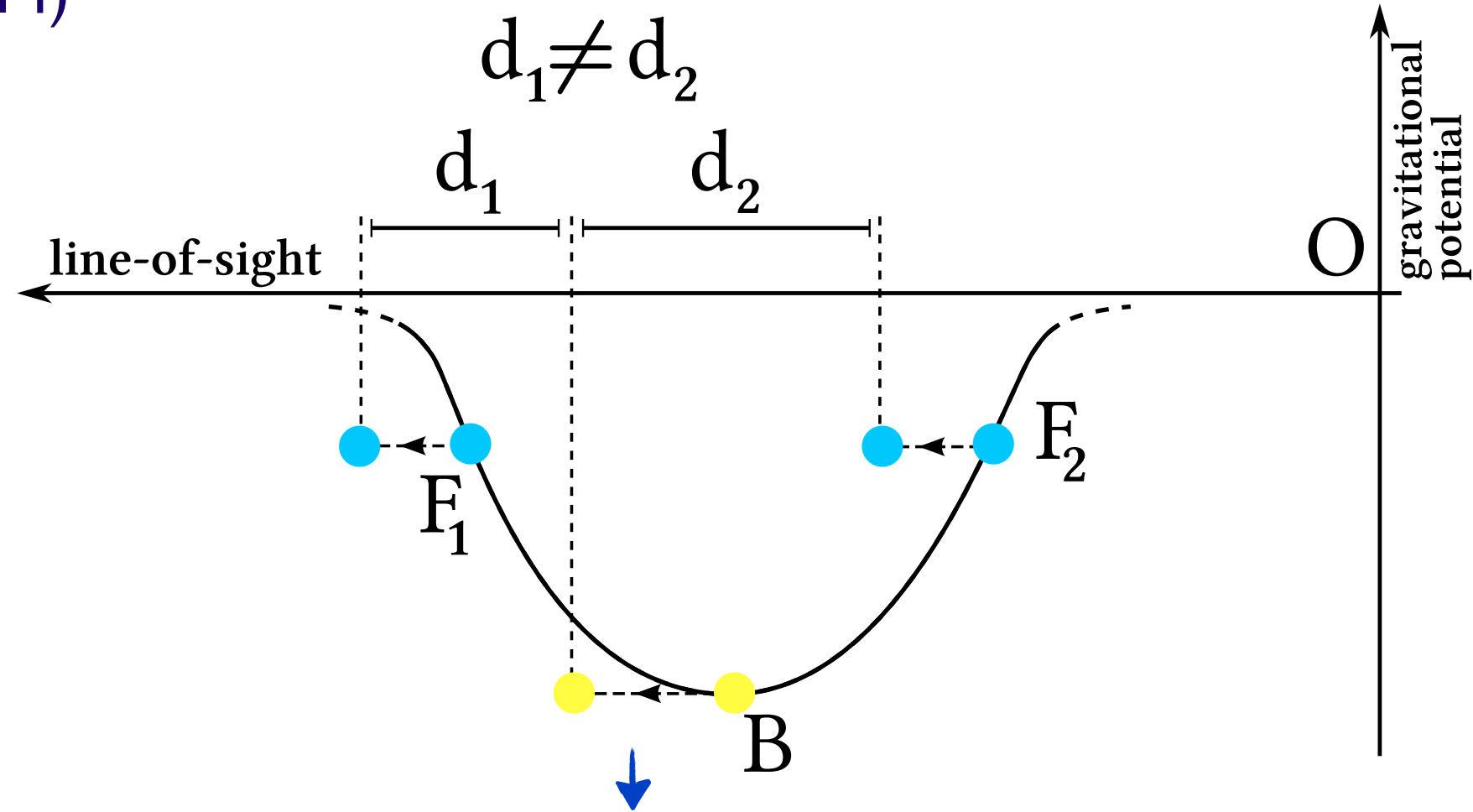
shift in position due to gravitational redshift

Taking all pairs of galaxies into account: **dipolar** modulation



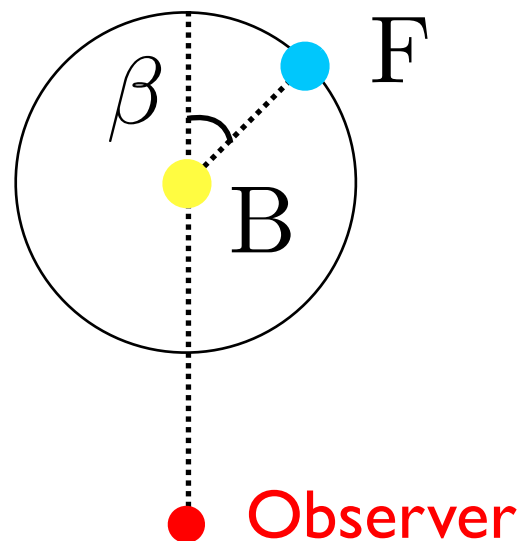
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We can **isolate** the effect by fitting for a dipole

What we really observe

Yoo et al (2010)
 CB and Durrer (2011)
 Challinor and Lewis (2011)

Matter fluctuations

$$\begin{aligned}
 \Delta(z, \mathbf{n}) = & \boxed{b \cdot \delta} - \frac{1}{\mathcal{H}} \partial_r (\mathbf{V} \cdot \mathbf{n}) \\
 & + (5s - 2) \int_0^r dr' \frac{r - r'}{2rr'} \Delta_\Omega(\Phi + \Psi) \\
 & + \left(1 - 5s - \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} + \frac{5s - 2}{r\mathcal{H}} \right) \mathbf{V} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \dot{\mathbf{V}} \cdot \mathbf{n} + \boxed{\frac{1}{\mathcal{H}} \partial_r \Psi} \\
 & + \frac{2 - 5s}{r} \int_0^r dr' (\Phi + \Psi) + 3\mathcal{H} \nabla^{-2} (\nabla \mathbf{V}) + \Psi + (5s - 2) \Phi \\
 & + \frac{1}{\mathcal{H}} \dot{\Phi} + \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^2} + \frac{2 - 5s}{r\mathcal{H}} + 5s \right) \left[\Psi + \int_0^r dr' (\dot{\Phi} + \dot{\Psi}) \right]
 \end{aligned}$$

Gravitational redshift

What we really observe

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$$\Delta(z, \mathbf{n}) = \frac{1}{\mathcal{H}} \partial_r (\Psi) + (5s - 2) \int_0^r \frac{1}{\omega} \frac{r - r'}{2rr'} \Delta_{\mathcal{L}}(\mathbf{I} + \mathbf{I})$$

Dipole

$$+ \left(1 - 5s - \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} + \frac{5s - 2}{r\mathcal{H}} \right) \mathbf{V} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \dot{\mathbf{V}} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \partial_r \Psi$$

~~$$+ \frac{2 - 5s}{r} \int_0^r \frac{1}{\omega} (\mathbf{I} + \mathbf{I}) + 3\mathcal{H} \nabla^2 (\nabla \mathbf{V}) + \mathbf{I} + (5s - 2) \mathbf{I}$$~~

~~$$+ \frac{1}{\mathcal{H}} \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^2} + \frac{2 - 5s}{r\mathcal{H}} + 5s \right) \left[\mathbf{I} + \int_0^r \frac{1}{\omega} (\dot{\mathbf{I}} + \dot{\mathbf{I}}) \right]$$~~

Isolating gravitational redshift

We combine the **dipole**, with measurements of redshift-space distortions (monopole, quadrupole and hexadecapole)

- ◆ Dipole $\rightarrow \Psi$ and V
- ◆ Redshift-space distortions $\rightarrow V$ and δ

Forecasts for SKA2

Redshift	0.35	0.45	0.55	0.65	0.75	0.85	0.95
Constraints	23%	24%	28%	33%	40%	48%	60%

Testing gravity

δ From RSD V

From lensing $\Phi + \Psi$

Test of the weak equivalence principle

$$V' + V - \frac{k}{\mathcal{H}}\Psi = 0$$

$$k^2\Psi = -\frac{3}{2}\mathcal{H}^2\mu\delta$$

Modified Poisson

$$\Phi = \eta\Psi$$

Non-zero anisotropic stress

→ standard

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Testing gravity

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$$V' + V - \frac{k}{\mathcal{H}}\Psi = -\Theta V + \frac{k}{\mathcal{H}}\Gamma\Psi$$

Friction
Fifth force

CB and Fleury (2018)

$$k^2\Psi = -\frac{3}{2}\mathcal{H}^2\mu\delta$$

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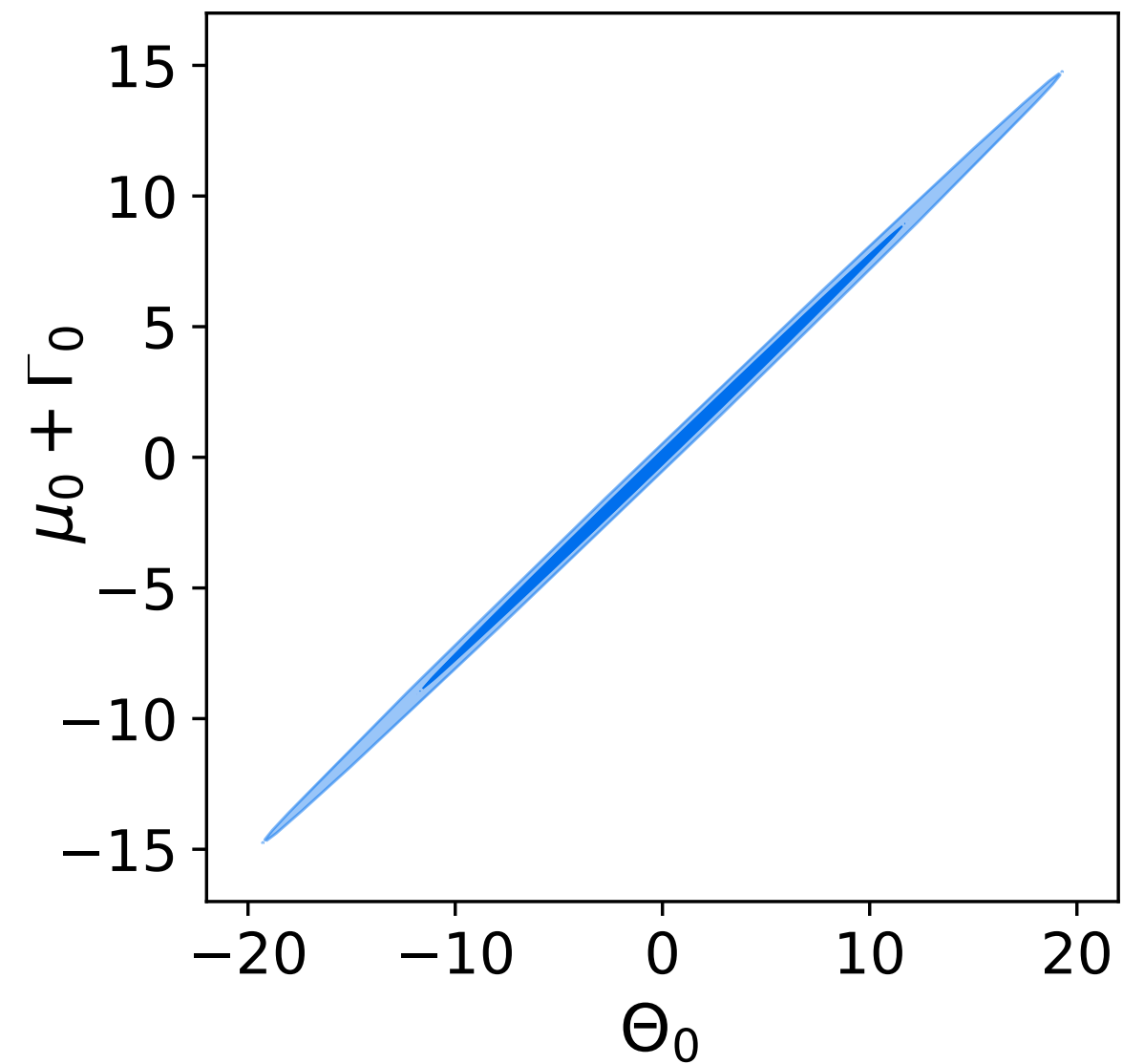
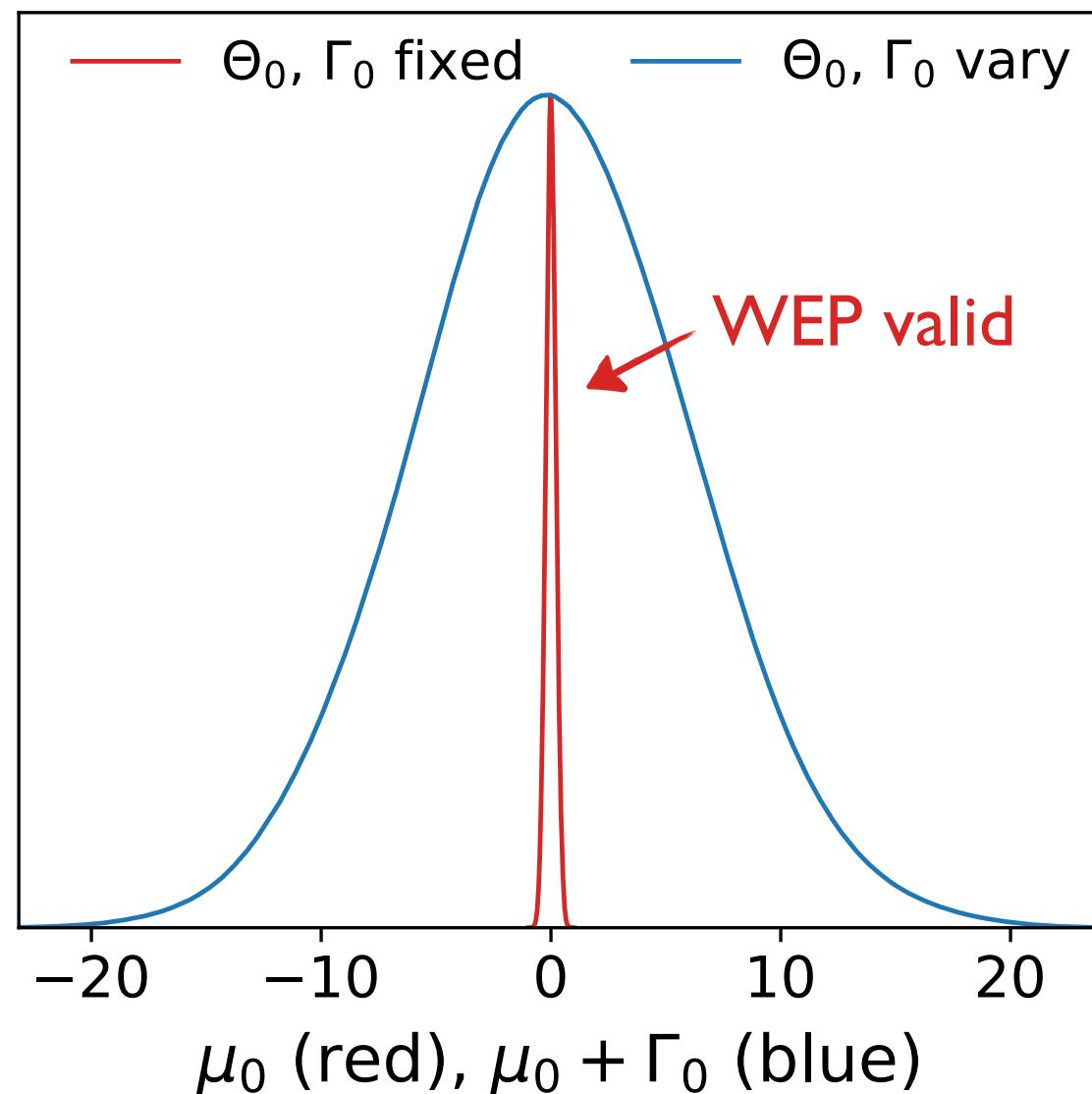
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Current constraints from SDSS

Castello, Grimm and CB (2022)

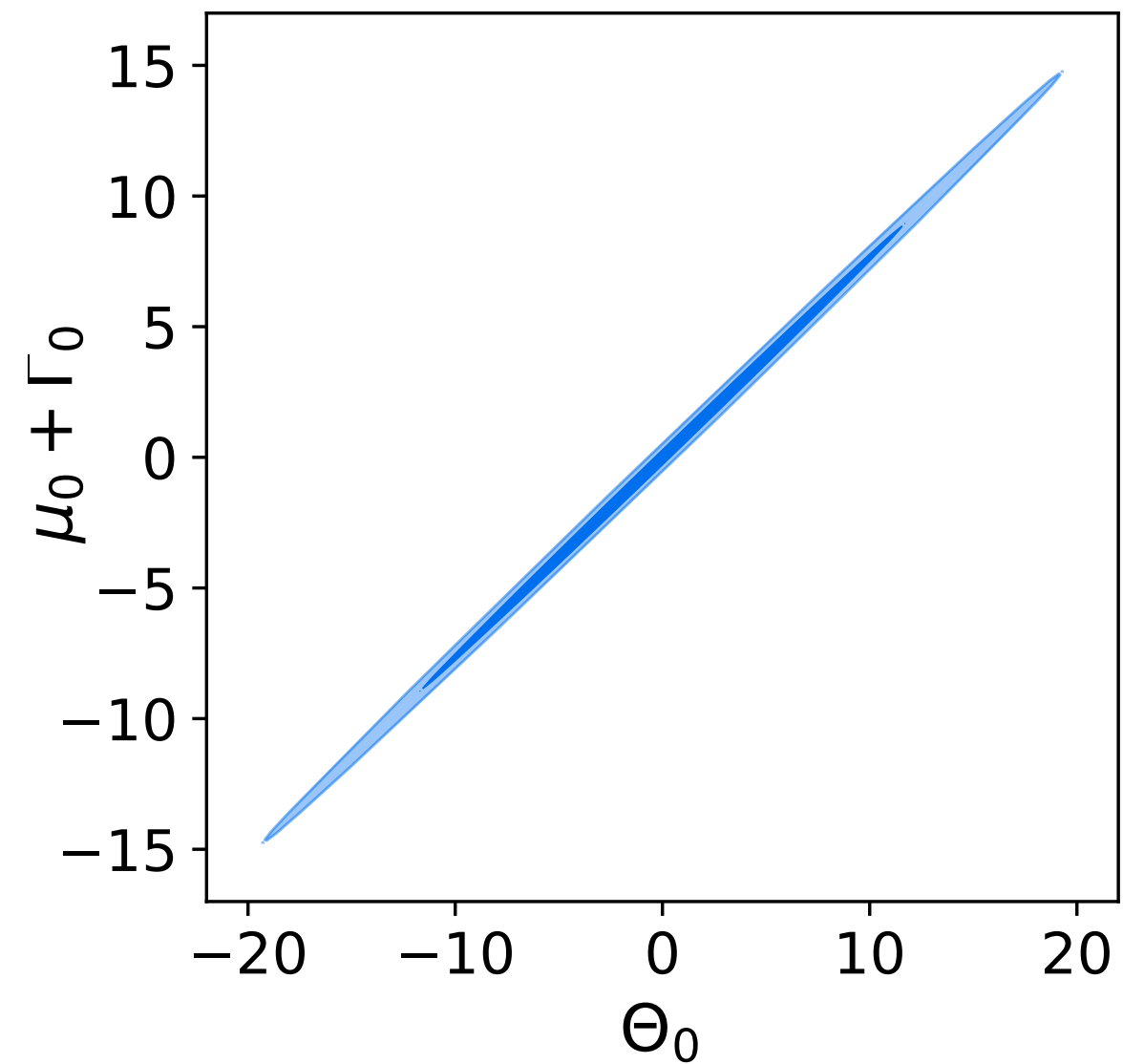
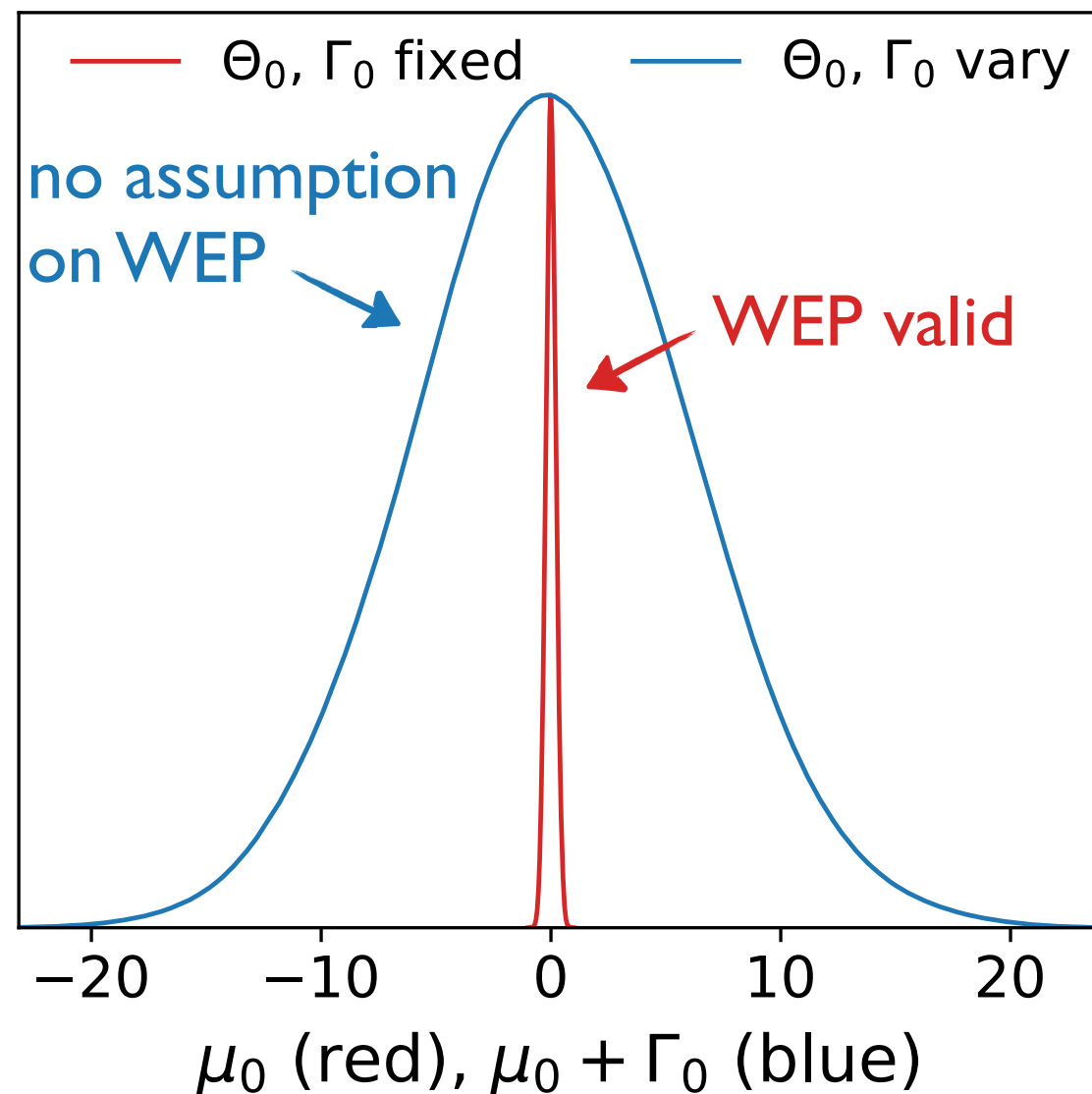
With redshift-space distortion only, we **cannot** test the weak equivalence principle \rightarrow **degeneracies**



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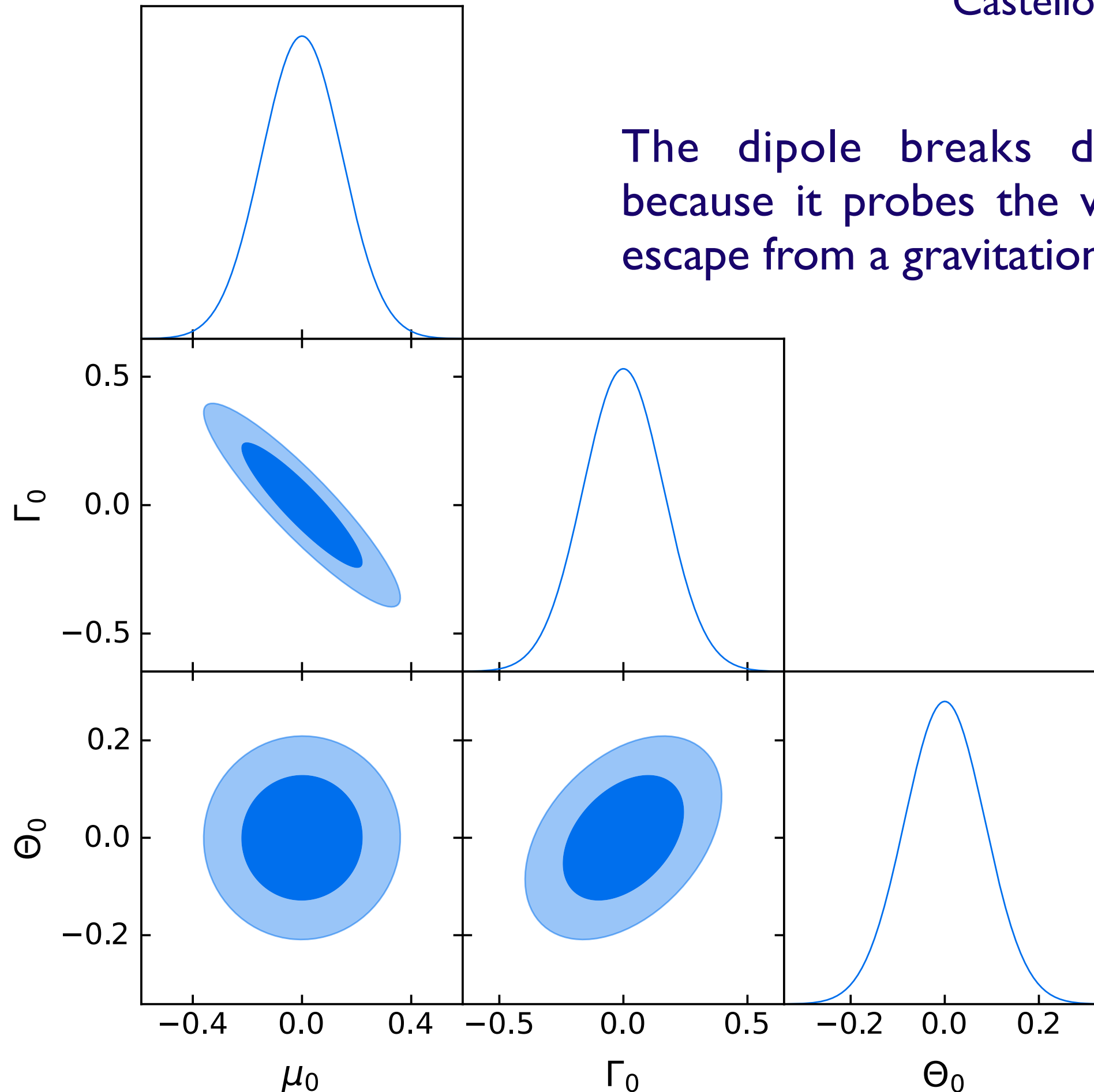
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Forecasts with dipole with SKA2

Castello, Grimm and CB (2022)



The dipole breaks degeneracies because it probes the way photons escape from a gravitational potential

Conclusion

- ◆ Mapping the distribution of galaxies allows us to measure **time distortion** at cosmological distance
- ◆ To **isolate** the effect, we look for a **dipole** in the cross-correlation of bright and faint galaxies
- ◆ Comparing this with redshift-space distortion provides a way of testing the validity of the **weak equivalence principle**

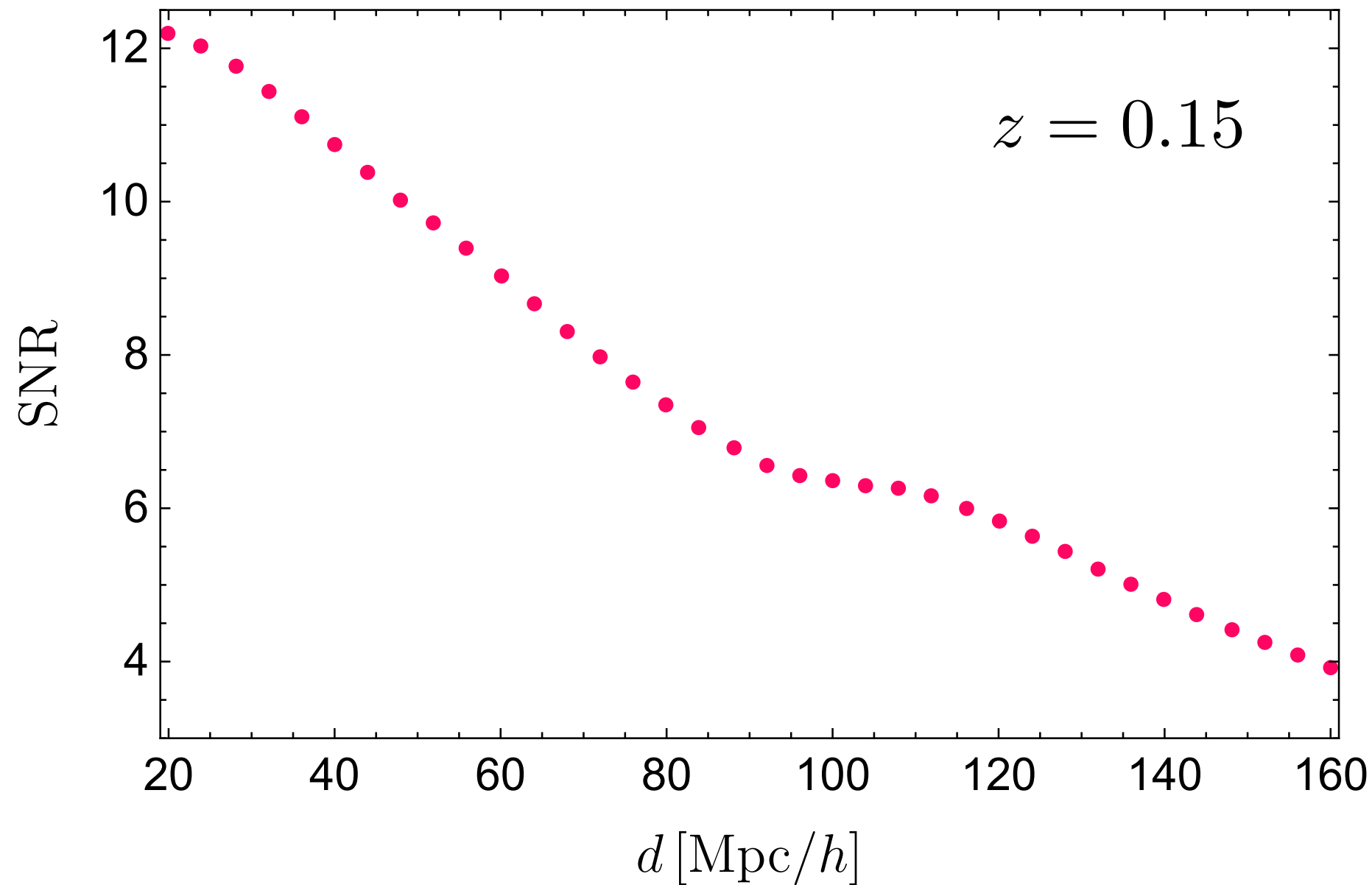
Backup slides

Evolution equation

$$\delta'' + \left(1 + \frac{\mathcal{H}'}{\mathcal{H}} + \Theta\right) \delta' - \frac{3}{2} \frac{\Omega_m}{a} \left(\frac{\mathcal{H}_0}{\mathcal{H}}\right)^2 \mu(\Gamma + 1) \delta = 0$$

Forecasts

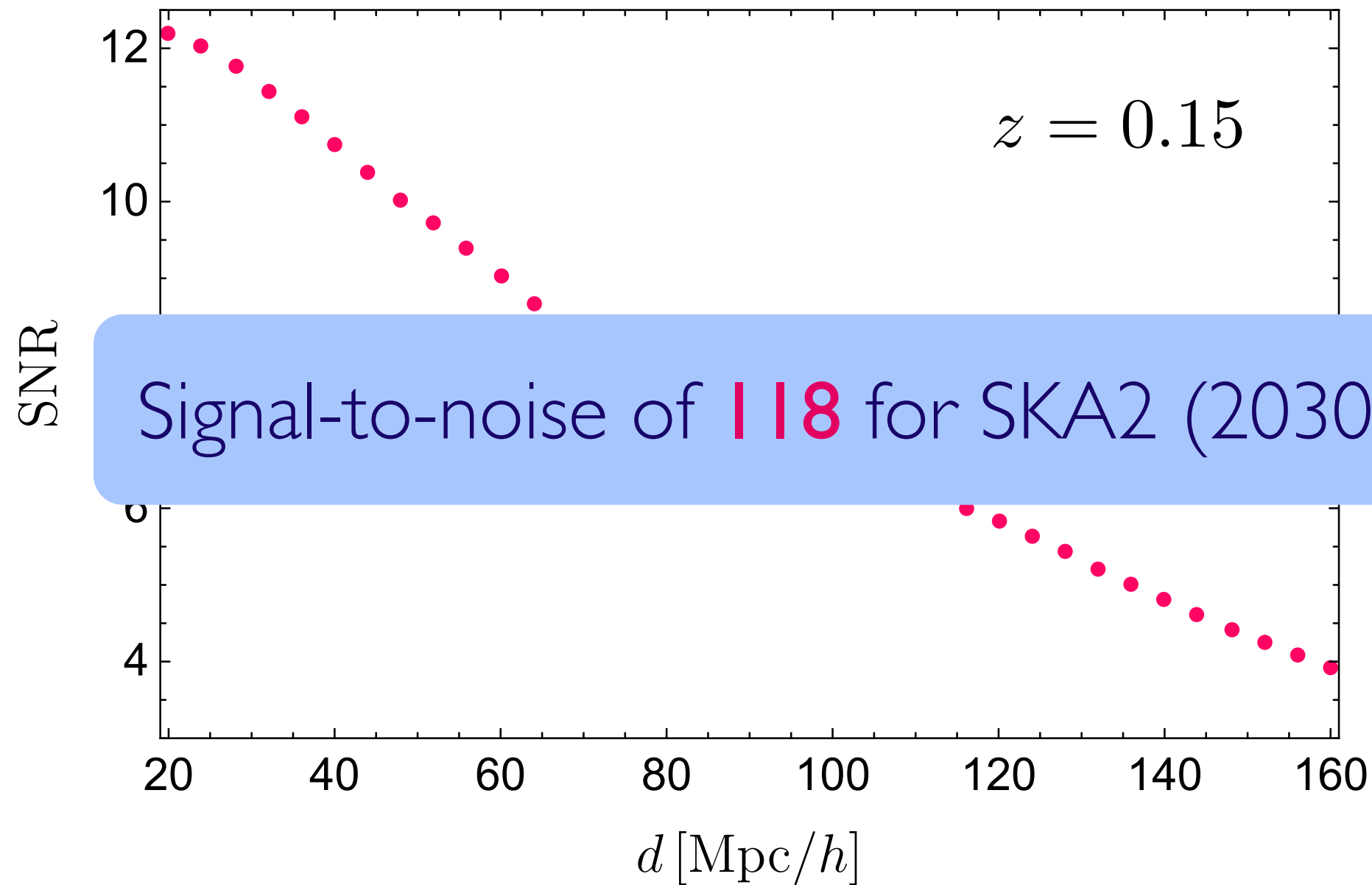
DESI Bright Galaxy Sample: 10 million galaxies at $z \leq 0.5$



Cumulative signal-to-noise: 33

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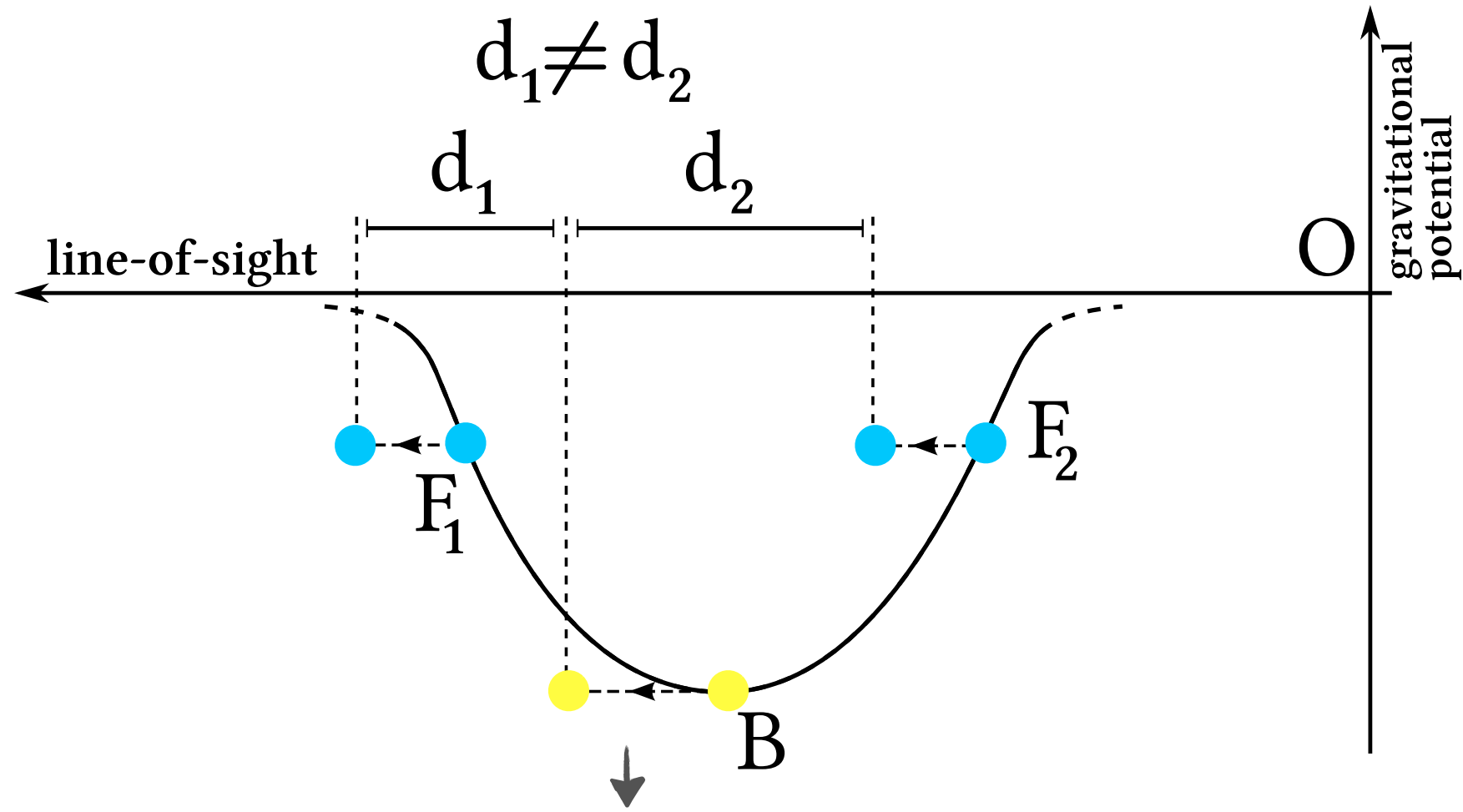
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shift in position due to gravitational redshift

$$\xi = \frac{\mathcal{H}}{\mathcal{H}_0} \left(\frac{D_1}{D_{10}} \right)^2 \left[(b_B - b_F) \left(\frac{2}{r\mathcal{H}} + \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} \right) + 3(s_F - s_B) f^2 \left(1 - \frac{1}{r\mathcal{H}} \right) + 5(b_B s_F - b_F s_B) f \left(1 - \frac{1}{r\mathcal{H}} \right) \right] \nu_1(d) \cos(\beta)$$