Unraveling the Hubble constant tension

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The University Of Sheffield.

What is H0?

The Hubble constant H0 describes the expansion rate of the Universe today.

This can be obtained in mainly two ways: measuring the distance and the recessional velocity of standard candles, and computing the proportionality factor.

1.



Jha, S. (2002) Ph.D. thesis (Harvard Univ., Cambridge, MA).

What is H0?

The Hubble constant H0 describes the expansion rate of the Universe today.

This can be obtained in mainly two ways:

- 1. measuring the distance and the recessional velocity of standard candles, and computing the proportionality factor.
- 2. considering early universe measurements, and assuming a model for the expansion history of the universe.

For example, we have CMB measurements and we assume the standard model of cosmology, i.e. the LCDM scenario.

1st Friedmann equations describes the expansion history of the universe:

$$H^2(z)=H_0^2\left(\Omega_m(1+z)^3+\Omega_k(1+z)^2+\Omega_\Lambda
ight).$$



The H0 tension at 5o!!

The H0 tension is the most statistically significant, long-lasting and widely persisting disagreement between:















CMB constraints



From the map of the CMB anisotropies we can extract the temperature angular power spectrum.

Planck 2018, Astron.Astrophys. 641 (2020) A6



Cosmological parameters: (Ω_bh^2 , Ω_mh^2 , H0, n_s, τ , As)



We choose a set of cosmological parameters that describes our theoretical model and compute the angular power spectra.

Because of the correlations present between the parameters, variation of different quantities can produce similar effects on the CMB.



Cosmological parameters: (Ω_bh^2 , Ω_mh^2 , H0, n_s, τ , As)



Theoretical model

We compare the angular power spectra we computed with the data and, using a bayesian analysis, we get a combination of cosmological parameter values in agreement with these.



Planck 2018, Astron.Astrophys. 641 (2020) A6



Parameter constraints

CMB constraints

Parameter	TT+lowE 68% limits	TE+lowE 68% limits	EE+lowE 68% limits	TT,TE,EE+lowE 68% limits	TT,TE,EE+lowE+lensing 68% limits	TT,TE,EE+lowE+lensing+BAO 68% limits
$\Omega_{\rm b}h^2$	0.02212 ± 0.00022	0.02249 ± 0.00025	0.0240 ± 0.0012	0.02236 ± 0.00015	0.02237 ± 0.00015	0.02242 ± 0.00014
$\Omega_{\rm c}h^2$	0.1206 ± 0.0021	0.1177 ± 0.0020	0.1158 ± 0.0046	0.1202 ± 0.0014	0.1200 ± 0.0012	0.11933 ± 0.00091
100 <i>θ</i> _{MC}	1.04077 ± 0.00047	1.04139 ± 0.00049	1.03999 ± 0.00089	1.04090 ± 0.00031	1.04092 ± 0.00031	1.04101 ± 0.00029
τ	0.0522 ± 0.0080	0.0496 ± 0.0085	0.0527 ± 0.0090	$0.0544^{+0.0070}_{-0.0081}$	0.0544 ± 0.0073	0.0561 ± 0.0071
$\ln(10^{10}A_{\rm s})\ldots\ldots\ldots$	3.040 ± 0.016	$3.018^{+0.020}_{-0.018}$	3.052 ± 0.022	3.045 ± 0.016	3.044 ± 0.014	3.047 ± 0.014
<i>n</i> _s	0.9626 ± 0.0057	0.967 ± 0.011	0.980 ± 0.015	0.9649 ± 0.0044	0.9649 ± 0.0042	0.9665 ± 0.0038
$H_0 [\mathrm{kms^{-1}Mpc^{-1}}]$	66.88 ± 0.92	68.44 ± 0.91	69.9 ± 2.7	67.27 ± 0.60	67.36 ± 0.54	67.66 ± 0.42
Ω_{Λ}	0.679 ± 0.013	0.699 ± 0.012	$0.711^{+0.033}_{-0.026}$	0.6834 ± 0.0084	0.6847 ± 0.0073	0.6889 ± 0.0056
$\Omega_{\rm m}$	0.321 ± 0.013	0.301 ± 0.012	$0.289^{+0.026}_{-0.033}$	0.3166 ± 0.0084	0.3153 ± 0.0073	0.3111 ± 0.0056
$\Omega_{\rm m}h^2$	0.1434 ± 0.0020	0.1408 ± 0.0019	$0.1404^{+0.0034}_{-0.0039}$	0.1432 ± 0.0013	0.1430 ± 0.0011	0.14240 ± 0.00087
$\Omega_{\rm m}h^3$	0.09589 ± 0.00046	0.09635 ± 0.00051	$0.0981^{+0.0016}_{-0.0018}$	0.09633 ± 0.00029	0.09633 ± 0.00030	0.09635 ± 0.00030
<i>σ</i> ₈	0.8118 ± 0.0089	0.793 ± 0.011	0.796 ± 0.018	0.8120 ± 0.0073	0.8111 ± 0.0060	0.8102 ± 0.0060
$S_8 \equiv \sigma_8 (\Omega_{\rm m}/0.3)^{0.5} . \label{eq:ssecond}$	0.840 ± 0.024	0.794 ± 0.024	$0.781^{+0.052}_{-0.060}$	0.834 ± 0.016	0.832 ± 0.013	0.825 ± 0.011

Planck 2018, Astron.Astrophys. 641 (2020) A6

2018 Planck results are a wonderful confirmation of the flat standard ΛCDM cosmological model, but are **model dependent**!

- The cosmological constraints are obtained assuming a cosmological model.
- The results are affected by the degeneracy between the parameters that induce similar effects on the observables.

Are there other H0 estimates?

The H0 tension

CMB Polarization Measurements with SPTpol

On the same side of Planck, i.e. preferring smaller values of H0 we have:

Ground based CMB telescope

Nicholas Harrington UC Berkeley

 $\frac{\text{SPT-3G}}{\text{H0} = 68.8 \pm 1.5 \text{ km/s/Mpc} \text{ in } \Lambda\text{CDM}}$

LCDM - dependent

SPT-3G, Dutcher et al., Phys.Rev.D 104 (2021) 2, 022003



The H0 tension

On the same side of Planck, i.e. preferring smaller values of H₀ we have:



Ground based CMB telescope



ACT collaboration, Aiola et al., JCAP 12 (2020) 047

ACT-DR4: H0 = 67.9 ± 1.5 km/s/Mpc in Λ CDM

 $\frac{\text{ACT-DR4} + \text{WMAP}}{\text{H0} = 67.6 \pm 1.1 \text{ km/s/Mpc} \text{ in } \text{ACDM}}$

LCDM - dependent

The H0 tension

On the same side of Planck, i.e. preferring smaller values of H₀ we have:

BAO+Pantheon+BBN+ $\theta_{MC, Planck}$: H0 = 67.9 ± 0.8 km/s/Mpc

Planck 2018, Aghanim et al., Astron.Astrophys. 641 (2020) A6

BAO+BBN from BOSS and eBOSS: $H_0 = 67.35 \pm 0.97 \text{ km/s/Mpc}$

eBOSS, Alam et al., Phys.Rev.D 103 (2021) 8, 083533



eBOSS, Alam et al., Phys.Rev.D 103 (2021) 8, 083533

LCD/M - dependent



Balkenhol et al. (2021), Planck 2018+SPT+ACT : 67.49 ± 0.5 Pogosian et al. (2020), eBOSS+Planck mH2: 69.6 ± 1.8 Aghanim et al. (2020), Planck 2018: 67.27 ± 0.60 nim et al. (2020), Planck 2018+CMB lensing: 67.36 ± 0.54 Ade et al. (2016), Planck 2015, H0 = 67.27 ± 0.60

CMR w

Dutcher et al. (2021), SPT: 68.8 ± 1.5 Aiola et al. (2020). ACT: 67.9 + 1 Aiola et al. (2020), WMAP9+ACT: 67.6 + 1 Zhang Huang (2019) WMAP9+BAO: 68 36+9-5 Henning et al. (2018), SPT: 71.3 ± 2. Hinshaw et al. (2013), WMAP9: 70.0 ± 2.

No CMB, with BBN

Zhang et al. (2021), BOSS correlation function+BAO+BBN: 68.19±0.99 Chen et al. (2021). P+BAO+BBN: 69 23+0 7 Philcox et al. (2021), P+Bispectrum+BAO+BBN: 68.31+0.8 D' Amico et al. (2020). BOSS DR12+BBN: 68.5 + 2 Colas et al. (2020), BOSS DR12+BBN: 68.7 ± 1.5 Ivanov et al. (2020), BOSS+BBN: 67.9 ± 1. Alam et al. (2020), BOSS+eBOSS+BBN: 67.35 ± 0.97

CMB lensin Baxter et al. (2020): 73.5 ± 5

Philcox et al. (2020), P1(k)+CMB lensing: 70.6-3 LSS teg standard rule Farren et al. (2021): 69 5+3

SNIa-Cenheid

Riess et al. (2022), R22: 73.04 ± 1.04 Camarena Marra (2021): 74 30 + 1 45 Riess et al. (2020), R20: 73.2 ± 1. Breuval et al. (2020): 72.8 ± 2.7 Riess et al. (2019), R19: 74.03 ± 1.4 Camarena, Marra (2019): 754+1

SNIa-TRGB

Dhawan et al. (2022): 76.94 ± 6.4 Iones et al. (2022): 72.4 + 3.3 Anand, Tully, Rizzi, Riess, Yuan (2021): 71.5 ± 1. Freedman (2021): 69.8 ± 1. Kim, Kang, Lee, Jang (2021): 69.5 ± 4.2 Soltis Casertano Riess (2020): 72.1 + 2.0 Freedman et al. (2020): 69.6 ± 1. Reid Pesce Riess (2019) SH0ES: 71.1 + 1.99 Yuan et al. (2019): 72.4 + 2

> SNIa-Miras Huang et al. (2019): 73.3 ± 4.0

Blakeslee et al. (2021) IR-SBF w/ HST: 73.3 ± 2.5 Khetan et al. (2020) w/ LMC DEB: 71.1 + 4.1 Cantiello et al. (2018): 71.9 + 7.

> de Jacquer et al. (2022): 75.4+3de Jaeger et al. (2020): 75.8-3 c

Masers Pesce et al. (2020): 73.9 + 3.0

Tully Fisher Kourkchi et al. (2020): 76.0 ± 2.0

Schombert, McGaugh, Lelli (2020): 75.1 ± 2.8 HII galaxy Fernandez Arenas et al. (2018): 71.0 ± 3 Wang, Meng (2017): 76.12+3-42

Lensing related, mass model dener

Denzel et al. (2021): 71.8_ Birrer et al. (2020) TDCOSMO: 74.54 Birrer et al. (2020), TDCOSMO+SLACS: 67.4_ Yang, Birrer, Hu (2020): 73.65+1-9 Millon et al. (2020), TDCOSMO: 74.2 ± 1. Qi et al. (2020): 73.6+ Liao et al. (2020): 72.8+ Liao et al. (2019): 72.2 ± Shajib et al. (2019), STRIDES: 74.2 Wong et al. (2019), H0LiCOW 2019: 73.3+

GW re

GW related Mukherjee et al. (2022), GW170817+GWTC-3: 67⁺⁹, Abbott et al. (2021), GWTC-3: 68⁺¹ Palmese et al. (2021), GW170817: 72.77⁺¹ -2.55 Gavathri et al. (2020), GW190521+GW170817: 73.4+ Mukherjee et al. (2020), GW170817+ZTF: 67.6 Mukherice et al. (2019), GW170817+VLBI: 68 3+4 Hotokezaka et al. (2019): 70.3+2

Cosmic chronometer Moresco et al. (2022), flat Λ CDM with systematics: 66.5 ± 5. Moresco et al. (2022), open wCDM with systematics: 67.8+8-7



Hubble constant measurements made by different astronomical missions and groups over the years.

The orange vertical band corresponds to the H0 value from SHOES Team and the light pink vertical band corresponds to the H0 value as reported by Planck 2018 team within a ACDM scenario.

A sample code for producing similar figures with any choice of the data is made publicly available online at github.com/lucavisinelli/H0TensionRealm.



Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]

High precision measurements of HO

The high precision and consistency of the data at both ends present strong challenges to the possible solution space and demands a hypothesis with enough rigor to explain multiple observations – whether these invoke new physics, unexpected largescale structures or multiple, unrelated errors.







Freedman et al., arXiv:2002.01550



Freedman et al., arXiv:2002.01550

Soltis et al., arXiv:2012.09196

 $H0 = 69.6 \pm 1.9 \text{ km/s/Mpc}$





de Jaeger et al., arXiv:2006.03412





H0 = 76.00 ± 2.55 km/s/Mpc Kourkchi et al. arXiv:2004.14499

H0 = 75.10 ± 2.75 km/s/Mpc Schombert et al. arXiv:2006.08615

Tully-Fisher Relation (based on the correlation between the rotation rate of spiral galaxies and their absolute luminosity, and using as calibrators Cepheids and TRGB)

SNIa-Cepheid Late Universe Riess et al. (2022), R22: 73.04 ± 1.04 Camarena, Marra (2021): 74.30 ± 1.45 Riess et al. (2020) R20: 73.2 + 1.3 Breuval et al. (2020): 72.8 ± 2.7 SNIa-TRGB Dhawan et al. (2022): 76.94 + 6.4 Jones et al. (2022): 72.4 ± 3.3 Anand, Tully, Rizzi, Riess, Yuan (2021); 71.5 ± 1.8 Freedman (2021): 69.8 ± 1.7 Kim, Kang, Lee, Jang (2021): 69.5 ± 4.2 Soltis, Casertano, Riess (2020): 72.1 ± 2.0 Freedman et al. (2020): 69.6 ± 1.9 Blakeslee et al. (2021) IR-SBF w/ HST: 73.3 ± 2.5 Khetan et al. (2020) w/ LMC DEB: 71.1 ± 4.1 SNII de Jaeger et al. (2022): 75.4+3.5 de Jaeger et al. (2020): 75.8+5:2 Masers Pesce et al. (2020): 73.9 ± 3.0 Tully Fisher Kourkchi et al. (2020): 76.0 ± 2.6 Schombert, McGaugh, Lelli (2020): 75.1 ± 2.8 Lensing related, mass model dependen Denzel et al. (2021): 71.8+3.9 Birrer et al. (2020), TDCOSMO: 74.5+2 Birrer et al. (2020), TDCOSMO+SLACS: 67.4⁴²/₋₃ Yang, Birrer, Hu (2020): 73.65-1-9-2 Millon et al. (2020). TDCOSMO: 74.2 + 1.0 Qi et al. (2020): 73.6+ Liao et al. (2020): 72.8+

Abdalla et al., JHEAp 34 (2022) 49-211

Model Dependent

H0 = 72.8 + 1.6 - 1.7 km/s/MpcLiao et al. arXiv:2002.10605 H0 = 73.6 + 1.8 - 1.6 km/s/MpcQi et al. arXiv:2011.00713 H0 = 73.65 + 1.95 - 2.26 km/s/MpcYang et al. arXiv:2003.03277 **TDCOSMO** H0 = 74.5 + 5.6 - 6.1 km/s/Mpc**TDCOSMO+SLACS** H0 = 67.4 + 4.1 - 3.2 km/s/MpcBirrer et al. arXiv:2007.02941 H0 = 71.8 + 3.9 - 3.3 km/s/MpcDenzel et al. arXiv:2007.14398

 Strong Lensing measurements of the time delays of multiple images of quasar systems caused by the strong gravitational lensing from a foreground galaxy.
 Uncertainties coming from the lens mass profile.









Freedman, Astrophys. J. 919 (2021) 1, 16

In the past the tension was within the same types of measurements and at the same redshifts and thus pointing directly to systematics.

High precision measurements of Ho

Now there are no late universe measurements below the early ones and vice versa.



It is hard to conceive of a single type of systematic error that would apply to the measurements of the disparate phenomena we saw before as to effectively resolve the Hubble constant tension. Because the tension remains with the removal of the measurements of any single type of object, mode or calibration, it is challenging to devise a single error that would suffice. While multiple, unrelated systematic errors have a great deal more flexibility to resolve the tension but become less likely by their inherent independence.

Since the indirect constraints are model dependent, we can try to expand the cosmological scenario and see which extensions work in solving the tensions between the cosmological probes.

The **ACDM** model

Among a number of cosmological models introduced in the literature, the Lambda Cold Dark Matter (ACDM) cosmological model is the mathematically simplest model, and has now practically been selected as the "standard" cosmological scenario, because it provides a remarkable description of a wide range of astrophysical and cosmological probes.

However, despite its marvelous fit to the available observations, <u>ACDM harbours large areas of phenomenology and ignorance</u>. For example, it still cannot explain key pillars in our understanding of the structure and evolution of the Universe, namely, <u>Dark Energy, Dark Matter and Inflation</u>.

The ΛCDM model

In the ACDM paradigm these three pillars are our simplest guesses.

- DE assumes its simplest form, that is the cosmological constant, without any strong physical basis.
- The nature of DM is still a mystery except for its gravitational interaction, as suggested by the observational evidence. We know, however, that DM is essential for structure formation in the late Universe, so most of it must be pressure-less, cold, and stable on cosmological time scales. Moreover, despite the significant efforts in the last decades to investigate DM and the physics beyond the SM of particle physics, in laboratory experiments and from devised astrophysical observations, no evidence pointing to the dark matter particle has been found.
- Finally, even though the theory of inflation has solved a number of crucial puzzles related to the early evolution of the Universe, in the standard model this is given by a single, minimally coupled, slow-rolling scalar field.
The **ACDM** model

Therefore, the 6 parameter ACDM model lacks the deep underpinnings a model requires to approach fundamental physics laws. It can be rightly considered, at best, as an effective theory of an underlying physical theory, yet to be discovered. In this situation, we must be careful not to cling to the model too tightly or to risk missing the appearance of departures from the paradigm.

With the improvement of the number and the accuracy of the observations, deviations from ACDM may be expected. And, actually, discrepancies among key cosmological parameters of the models have emerged with different statistical significance.

While some proportion of these discrepancies may have a systematic origin, their persistence across probes should require multiple and unrelated errors, strongly hinting at cracks in the standard cosmological scenario and the necessity of new physics.

These tensions can indicate a failure of the canonical ACDM model.

Let's modify the ΛCDM model...

The Neutrino effective number

We can consider modifications in the dark matter sector.

A classical extension is the effective number of relativistic degrees of freedom, i.e. additional relativistic matter at recombination, corresponding to a modification of the expansion history of the universe at early times.

The Neutrino effective number

The expected value is Neff = 3.044, if we assume standard electroweak interactions and three active massless neutrinos. If we measure a Neff > 3.044, we are in presence of extra radiation.

If we compare the Planck 2015 constraint on Neff at 68% cl

$N_{\rm eff} =$	3.13 ± 0.32	Planck TT+lowP,
$N_{\rm eff} =$	3.15 ± 0.23	Planck TT+lowP+BAO,

with the new Planck 2018 bound,

 $N_{\rm eff} = 2.92^{+0.36}_{-0.37}$ (95%, *Planck* TT, TE, EE+lowE),

we see that the neutrino effective number is now very well constrained.

H0 passes from 68.0 \pm 2.8 km/s/Mpc (2015) to 66.4 \pm 1.4 km/s/Mpc (2018), and the tension with SH0ES increases from 1.7 σ to 3.9 σ also varying Neff.



Planck collaboration, 2015



The Dark energy equation of state

For example, we can consider modifications in the dark energy sector.

A classical extension is a varying dark energy equation of state, that is a modification of the expansion history of the universe at late times.

The Dark energy equation of state

Changing the dark energy equation of state w, we are changing the expansion rate of the Universe:

$$H^{2} = H_{0}^{2} \left[\Omega_{m} (1+z)^{3} + \Omega_{r} (1+z)^{4} + \Omega_{de} (1+z)^{3(1+w)} + \Omega_{k} (1+z)^{2} \right]$$

w introduces a geometrical degeneracy with the Hubble constant that will be unconstrained using the CMB data only, resulting in agreement with SH0ES.

We have in 2018 w = $-1.58^{+0.52}_{-0.41}$ with H0 > 69.9 km/s/Mpc at 95% c.l.

Planck data prefer a phantom dark energy, with an energy component with w < -1, for which the density increases with time in an expanding universe that will end in a Big Rip. A phantom dark energy violates the energy condition $\rho \ge |p|$, that means that the matter could move faster than light and a comoving observer measure a negative energy density, and the Hamiltonian could have vacuum instabilities due to a negative kinetic energy.

Formally successful models in solving H0

tension $\leq 2\sigma$ "Good models"	tension $\leq 3\sigma$ "Promising models"
Early Dark Energy [235]	Early Dark Energy [229]
Phantom Dark Energy [11]	Decaying Warm DM [474]
Dynamical Dark Energy [11,281,309]	Neutrino-DM Interaction [506]
GEDE [397]	Interacting dark radiation [517]
Vacuum Metamorphosis [402]	Self-Interacting Neutrinos [700, 701]
IDE [314, 653, 656, 661, 663, 670]	IDE [656]
Critically Emergent Dark Energy [997]	Unified Cosmologies [747]
$f(\mathcal{T})$ gravity [814]	Scalar-tensor gravity [856]
Über-gravity [59]	Modified recombination [986]
Reconstructed PPS [978]	Super ΛCDM [1007]
	Coupled Dark Energy [650]
	tension $\leq 2\sigma$ "Good models" Early Dark Energy [235] Phantom Dark Energy [11] Dynamical Dark Energy [11, 281, 309] GEDE [397] Vacuum Metamorphosis [402] IDE [314, 653, 656, 661, 663, 670] Critically Emergent Dark Energy [997] $f(\mathcal{T})$ gravity [814] Über-gravity [59] Reconstructed PPS [978]

Table B1. Models solving the H_0 tension with R20 within the 1σ , 2σ and 3σ Planck only confidence levels considering the *Planck* dataset only.

Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]

Let's see an example...

Parker Vacuum Metamorphosis

There is a model considered in the early days of dark energy investigations that possesses the phenomenological properties needed to solve the H0 tension, but is based on a sound theoretical foundation: the vacuum metamorphosis model of Parker and Raval, Phys. Rev. D 62, 083503 (2000), Parker and Vanzella, Phys. Rev. D 69, 104009 (2004),

Caldwell, Komp, Parker and Vanzella, Phys. Rev. D 73, 023513 (2006), which has a phase transition in the nature of the vacuum.

Vacuum metamorphosis arises from a nonperturbative summation of quantum gravity loop corrections due to a massive scalar field.

We found that the Parker vacuum metamorphosis model, physically motivated by quantum gravitational effects, with the same number of parameters as LCDM, but not nested with it, can remove the H₀ tension, because can mimic a phantom DE behaviour at low redshifts.

First principles theory

Parker Vacuum Metamorphosis

When the Ricci scalar evolves during cosmic history to reach the scalar field mass squared, then a phase transition occurs and R freezes with

$$R = 6(\dot{H} + 2H^2 + ka^{-2}) = m^2$$
 and defining $M = m^2/(12H_0^2)$

The expansion behaviour above and below the phase transition is

$$H^{2}/H_{0}^{2} = \Omega_{m}(1+z)^{3} + \Omega_{r}(1+z)^{4} + \Omega_{k}(1+z)^{2} + M \left\{ 1 - \left[3\left(\frac{4}{3\Omega_{m}}\right)^{4}M(1-M-\Omega_{k}-\Omega_{r})^{3} \right]^{-1} \right\}, \ z > z_{t}$$
$$H^{2}/H_{0}^{2} = (1-M-\Omega_{k})(1+z)^{4} + \Omega_{k}(1+z)^{2} + M, \quad z \le z_{t}$$

with $z_t = -1 + rac{3\Omega_m}{4(1-M-\Omega_k-\Omega_r)}$

We see that above the phase transition, the universe behaves as one with matter (plus radiation plus spatial curvature) plus a constant, and after the phase transition it effectively has a dark radiation component that rapidly redshifts away leaving a de Sitter phase.

Parker Vacuum Metamorphosis

When the Ricci scalar evolves during cosmic history to reach the scalar field mass squared, then a phase transition occurs and R freezes with

$$R = 6(\dot{H} + 2H^2 + ka^{-2}) = m^2$$
 and defining $M = m^2/(12H_0^2)$

The expansion behaviour above and below the phase transition is

$$\begin{split} H^2/H_0^2 &= \Omega_m (1+z)^3 + \Omega_r (1+z)^4 + \Omega_k (1+z)^2 + M \left\{ 1 - \left[3 \left(\frac{4}{3\Omega_m} \right)^4 M (1-M-\Omega_k - \Omega_r)^3 \right]^{-1} \right\}, \ z > z_t \\ H^2/H_0^2 &= (1-M-\Omega_k)(1+z)^4 + \Omega_k (1+z)^2 + M \,, \quad z \le z_t \end{split}$$

with

$$z_t = -1 + \frac{3\Omega_m}{4(1 - M - \Omega_k - \Omega_r)}$$

The original model did not include an explicit high redshift cosmological constant; we see that this implies that

$$\Omega_m = \frac{4}{3} \left[3M(1 - M - \Omega_k - \Omega_r)^3 \right]^{1/4}$$

i.e. the parameter M is fixed and depends on the matter density, and this model has the same number of degrees of freedom as ACDM.

A Vacuum Phase Transition Solves the H₀ Tension

Parameters	CMB	CMB+lensing	CMB+BAO	CMB+Pantheon	CMB+R19	CMB+BAO+Pantheon	CMB+BAO+R19
$\Omega_b h^2$	0.02238 ± 0.00014	0.02242 ± 0.00013	0.02218 ± 0.00012	0.02201 ± 0.00013	0.02221 ± 0.00012	0.02213 ± 0.00012	0.02217 ± 0.00012
$100 heta_{MC}$	1.04091 ± 0.00030	1.04097 ± 0.00029	1.04060 ± 0.00029	1.04033 ± 0.00031	1.04063 ± 0.00029	1.04053 ± 0.0029	1.04060 ± 0.00029
au	0.0524 ± 0.0078	0.0510 ± 0.0078	$0.0458\substack{+0.0083\\-0.0067}$	$0.039\substack{+0.010\\-0.007}$	0.0469 ± 0.0075	$0.0449\substack{+0.0079\\-0.0065}$	$0.0456\substack{+0.0083\\-0.0068}$
M	$0.9363\substack{+0.0055\\-0.0044}$	0.9406 ± 0.0034	0.9205 ± 0.0023	$0.8996^{+0.0081}_{-0.0073}$	$0.9230\substack{+0.0042\\-0.0036}$	0.9163 ± 0.0028	0.9198 ± 0.0020
$\ln(10^{10}A_s)$	3.041 ± 0.016	3.036 ± 0.015	$3.035\substack{+0.017\\-0.014}$	$3.027^{+0.020}_{-0.014}$	3.036 ± 0.016	$3.035\substack{+0.017\\-0.014}$	$3.035\substack{+0.017\\-0.015}$
n_s	0.9643 ± 0.0039	0.9663 ± 0.0036	0.9572 ± 0.0031	0.9511 ± 0.0036	0.9585 ± 0.0033	0.9560 ± 0.003	0.9571 ± 0.0031
$H_0[\rm km/s/Mpc]$	81.1 ± 2.1	82.9 ± 1.5	75.44 ± 0.69	70.1 ± 1.8	76.3 ± 1.2	74.21 ± 0.66	75.22 ± 0.60
σ_8	0.9440 ± 0.0077	0.9392 ± 0.0067	$0.9430_{-0.0070}$	$0.9419_{-0.0069}^{+0.0008}$	0.9457 ± 0.0075	$0.9401_{-0.0068}$	$0.9457^{+0.0082}_{-0.0073}$
S_8	0.805 ± 0.022	0.783 ± 0.014	0.865 ± 0.010	0.927 ± 0.023	0.856 ± 0.015	0.880 ± 0.010	0.8675 ± 0.0098
Ω_m	$0.218\substack{+0.010\\-0.012}$	0.2085 ± 0.0076	0.2510 ± 0.0046	0.291 ± 0.015	$0.2458\substack{+0.0074\\-0.0084}$	0.2593 ± 0.0046	0.2525 ± 0.0040
$\chi^2_{ m bf}$	2767.74	2776.23	2806-22	3874 13	2777~04	3910.01	2803.34
$\Delta \widetilde{\chi^2_{ m bf}}$	-4.91	-5.81	+26.51	+66.63	-14.80	+95.83	+11.29



We don't solve the tension, we do obtain H0~73-74 km/s/Mpc !!

H0 is exactly in agreement with SH0ES even if BAO and Pantheon are included. However, this worsen considerably the fit of the data because the model fails in recover the shape of H(z) at low redshifts.

Di Valentino et al., Phys.Dark Univ. 30 (2020) 100733

What about BAO+Pantheon?

BAO+Pantheon measurements constrain the product of H0 and the sound horizon r_s.

In order to have a higher H0 value in agreement with SH0ES, we need r_s near 137 Mpc. However, Planck by assuming Λ CDM, prefers r_s near 147 Mpc. Therefore, a cosmological solution that can increase H0 and at the same time can lower the sound horizon inferred from CMB data is the most promising way to put in agreement all the measurements.



Knox and Millea, Phys. Rev. D 101 (2020) 4, 043533

Early vs late time solutions

Here we can see the comparison of the 2 σ credibility regions of the CMB constraints and the measurements from late-time observations (SN + BAO + H0LiCOW + SH0ES).

We see that the late time solutions, as wCDM, increase H0 because they decrease the expansion history at intermediate redshift, but leave r_s unaltered.



Arendse et al., Astron.Astrophys. 639 (2020) A57

Early vs late time solutions

Here we can see the comparison of the 2 σ credibility regions of the CMB constraints and the measurements from late-time observations (SN + BAO + H0LiCOW + SH0ES).

However, the early time solutions, as Neff or Early Dark Energy, move in the right direction both the parameters, but can't solve completely the H0 tension between Planck and SH0ES.



Arendse et al., Astron.Astrophys. 639 (2020) A57

Early Dark Energy



ACT collaboration, Hill et al. arXiv:2109.04451

Considering ACT only data or combined with Planck TT up to multipoles 650, there is an evidence for EDE > 3σ , solving completely the Hubble tension.

Early Dark Energy



ACT collaboration, Hill et al. arXiv:2109.04451

The evidence for EDE > 3σ persists with the inclusion of Planck lensing + BAO data, but shifting H0 towards a lower value.

Early Dark Energy



ACT collaboration, Hill et al. arXiv:2109.04451

Once the full Planck data are considered, the evidence for EDE disappears and H0 is again in tension with SH0ES.

The Planck damping tail is in disagreement with EDE different from zero.

ACT DR4 vs Planck: LCDM

Planck

Handley and Lemos, arXiv:2007.08496 [astro-ph.CO]

ACT

SPT

0.024 $\Omega_b h^2$ 0.022 Dataset combination tension \boldsymbol{p} ACT vs *Planck* 0.86% 2.63σ ACT vs SPT 1.8% 2.37σ 0.120 $\Omega_c h^2$ *Planck* vs SPT 16.8% 1.38σ ACT vs Planck+SPT 0.52% 2.79σ 0.1051.044 $100 \theta_{MC}$ 1.0401.0360.120.080.04 $\ln(10^{10}A_s)$ 3.13.02.91.02 n_s 0.96 $0.105 \ 0.120 \ 1.036 \ 1.040 \ 1.044 \ 0.04 \ 0.08 \ 0.122.9$ 3.03.10.960.0220.0241.02 $\Omega_b h^2$ $\Omega_c h^2$ $100\theta_{MC}$ $\ln(10^{10}A_{s})$ au n_{s}

Global tensions between CMB datasets.

For each pairing of datasets this is the tension probability p that such datasets would be this discordant by (Bayesian) chance, as well as a conversion into a Gaussianequivalent tension.

Between Planck and ACT there is a 2.6σ tension.

Assuming LCDM

Formally successful models in solving H0

tension $\leq 1\sigma$ "Excellent models"	tension $\leq 2\sigma$ "Good models"	tension $\leq 3\sigma$ "Promising models"	
Early Dark Energy [228, 235, 240, 250]	Early Dark Energy [212, 229, 236, 263]	DE in extended parameter spaces [289]	
Exponential Acoustic Dark Energy [259]	Rock 'n' Roll [242]	Dynamical Dark Energy [281, 309]	
Phantom Crossing [315]	New Early Dark Energy [247]	Holographic Dark Energy [350]	
Late Dark Energy Transition [317]	Acoustic Dark Energy [257]	Swampland Conjectures [370]	
Metastable Dark Energy [314]	Dynamical Dark Energy [309]	MEDE [399]	
PEDE [394]	Running vacuum model [332]	Coupled DM - Dark radiation [534]	
Vacuum Metamorphosis [402]	Bulk viscous models [340, 341]	Decaying Ultralight Scalar [538]	
Elaborated Vacuum Metamorphosis [401, 402]	Holographic Dark Energy [350]	$BD-\Lambda CDM$ [852]	
Sterile Neutrinos [433]	Phantom Braneworld DE [378]	Metastable Dark Energy [314]	
Decaying Dark Matter [481]	PEDE [391, 392]	Self-Interacting Neutrinos [700]	
Neutrino-Majoron Interactions [509]	Elaborated Vacuum Metamorphosis [401]	Dark Neutrino Interactions [716]	
IDE [637, 639, 657, 661]	IDE [659,670]	$IDE \ [634-636, 653, 656, 663, 669]$	
DM - Photon Coupling [685]	Interacting Dark Radiation [517]	Scalar-tensor gravity [855,856]	
$f(\mathcal{T})$ gravity theory [812]	Decaying Dark Matter [471, 474]	Galileon gravity [877,881]	
BD- Λ CDM [851]	DM - Photon Coupling [686]	Nonlocal gravity [886]	
Über-Gravity [59]	Self-interacting sterile neutrinos [711]	Modified recombination [986]	
Galileon Gravity [875]	$f(\mathcal{T})$ gravity theory [817]	Effective Electron Rest Mass [989]	
Unimodular Gravity [890]	Über-Gravity [871]	Super ACDM [1007]	
Time Varying Electron Mass [990]	VCDM [893]	Axi-Higgs [991]	
M CDM [995]	Primordial magnetic fields [992]	Self-Interacting Dark Matter [479]	
Ginzburg-Landau theory [996]	Early modified gravity [859]	Primordial Black Holes [545]	
Lorentzian Quintessential Inflation [979]	Bianchi type I spacetime [999]		
Holographic Dark Energy [351]	f(T) [818]		

combination of datasets **Table B2.** Models solving the H_0 tension with R20 within 1σ , 2σ and 2σ *Planck* in combination with additional cosmological probes. datasets are discussed in the main text.

Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]

Let's see another example...

In the standard cosmological framework, DM and DE are described as separate fluids not sharing interactions beyond gravitational ones. At the background level, the conservation equations for the pressureless DM and DE components can be decoupled into two separate equations with an inclusion

of an arbitrary function, Q, known as the coupling or interacting function:

$$\dot{\rho}_c + 3\mathcal{H}\rho_c = Q,$$

$$\dot{\rho}_x + 3\mathcal{H}(1+w)\rho_x = -Q,$$

and we assume the phenomenological form for the interaction rate:

$$Q = \xi \mathcal{H} \rho_x$$

proportional to the dark energy density ρ_x and the conformal Hubble rate \mathcal{H} , via a negative dimensionless parameter ξ quantifying the strength of the coupling, to avoid early-time instabilities.

In this scenario of IDE the tension on H0 between the Planck satellite and R19 is completely solved. The coupling could affect the value of the present matter energy density $\Omega_{\rm m}$. Therefore, if within an interacting model Ω_m is smaller (because for negative ξ the dark matter density will decay into the dark energy one), a larger value of H0 would be required in order to satisfy the peaks structure of CMB observations, which accurately determine the value of $\Omega_m h^2$.

Parameter		Planck	Planck+R19	
	$\Omega_{ m b}h^2$	0.02239 ± 0.00015	0.02239 ± 0.00015	
	$\Omega_{ m c} h^2$	< 0.105	< 0.0615	
	n_s	0.9655 ± 0.0043	0.9656 ± 0.0044	
	$100\theta_{s}$	$1.0458\substack{+0.0033\\-0.0021}$	1.0470 ± 0.0015	
	au	0.0541 ± 0.0076	0.0534 ± 0.0080	
	ξ	$-0.54^{+0.12}_{-0.28}$	$-0.66\substack{+0.09\\-0.13}$	
H_0	$[{\rm kms^{-1}Mpc^{-1}}]$	$72.8^{+3.0}_{-1.5}$	$74.0^{+1.2}_{-1.0}$	

TABLE I. Mean values with their 68% C.L. errors on selected cosmological parameters within the $\xi\Lambda$ CDM model, considering either the *Planck* 2018 legacy dataset alone, or the same dataset in combination with the *R19* Gaussian prior on H_0 based on the latest local distance measurement from *HST*. The quantity quoted in the case of $\Omega_c h^2$ is the 95% C.L. upper limit.

Di Valentino et al., Phys.Dark Univ. 30 (2020) 100666

Therefore we can safely combine the two datasets together, and we obtain a nonzero dark matter-dark energy coupling ξ at more than FIVE standard deviations.



Di Valentino et al., Phys.Dark Univ. 30 (2020) 100666

Parameter	CMB+BAO	CMB+FS	CMB + BAO + FS
ω_c	$0.094\substack{+0.022\\-0.010}$	$0.101\substack{+0.015\\-0.009}$	$0.115\substack{+0.005\\-0.001}$
ξ	$-0.22^{+0.18}_{-0.09}$ [> -0	[0.48] > -0.35	> -0.12
$H_0[{ m km/s/Mpc}]$	$69.55\substack{+0.98\\-1.60}$	$69.04\substack{+0.84 \\ -1.10}$	$68.02\substack{+0.49\\-0.60}$
Ω_m	$0.243\substack{+0.054\\-0.030}$	$0.261\substack{+0.038\\-0.025}$	$0.299\substack{+0.015\\-0.007}$

The addition of low-redshift measurements, as BAO data, still hints to the presence of a coupling, albeit at a lower statistical significance. Also for this data sets the Hubble constant values is larger than that obtained in the case of a pure LCDM scenario, enough to bring the H0 tension at 2.1 σ with SH0ES.



Nunes, Vagnozzi, Kumar, Di Valentino, and Mena, arXiv:2203.08093 [astro-ph.CO]

The addition of low-redshift measurements, as BAO data, still hints to the presence of a coupling, albeit at a lower statistical significance. Also for this data sets the Hubble constant values is larger than that obtained in the case of a pure LCDM scenario, enough to bring the H0 tension at 2.1σ with SH0ES.

However, the IDE model does not survive to the additional information coming from the full shape (FS) power spectrum of the BOSS DR12 galaxies.

Baryon Acoustic Oscillations

BAO is formed in the early universe, when baryons are strongly coupled to photons, and the gravitational collapse due to the CDM is counterbalanced by the radiation pressure. Sound waves that propagate in the early universe imprint a characteristic scale on the CMB. Since the scale of these oscillations can be measured at recombination, BAO is considered a "standard ruler". These fluctuations have evolved and we can observe BAO at low redshifts in the distribution of galaxies. Since the data reduction process leading to these measurements requires assumptions about the fiducial cosmology, BAO is model dependent.

In other words, the tension between Planck+BAO or Planck+FS and SH0ES could be due to a statistical fluctuation in this case.

Actually, BAO and FS data are extracted under the assumption of LCDM, and the modified scenario of interacting dark energy could affect the result.

In fact, the full procedure which leads to the BAO and FS datasets carried out by the different collaborations might be not necessarily valid in extended DE models with important perturbations in the non-linear scales.

BAO and FS datasets (both the pre- and post- reconstruction measurements) might need to be revised in a non-trivial manner when applied to constrain more exotic dark energy cosmologies.

Additional complication: the models proposed to alleviate the H0 tension increase the S8 tension!



$$S_8 \equiv \sigma_8 \sqrt{\Omega_m/0.3}$$

A tension on S8 is present between the Planck data in the ΛCDM scenario and the cosmic shear data.



KiDS-1000, Heymans et al., arXiv:2007.15632 [astro-ph.CO]

The S8 tension is present at 3.4σ between Planck assuming ΛCDM and KiDS+VIKING-450 and BOSS combined together, or 3.1σ with KiDS-1000.

 $S_8 = 0.834 \pm 0.016$ Planck 2018, Aghanim et al., arXiv:1807.06209 [astro-ph.CO]

> $S_8 = 0.728 \pm 0.045$ Troster et al., arXiv:1909.11006 [astro-ph.CO]

> > $S_8 = 0.766^{+0.020}_{-0.014}$

KiDS-1000, Heymans et al., arXiv:2007.15632 [astro-ph.CO]



DES-Y3, Amon et al., arXiv:2105.13543 [astro-ph.CO]

The S8 tension is present at 2.5σ between Planck assuming Λ CDM and DES-Y3.

 $S_8 = 0.834 \pm 0.016$

Planck 2018, Aghanim et al., arXiv:1807.06209 [astro-ph.CO]

 $S_8 = 0.776^{+0.017}_{-0.017}$

DES-Y3, Abbott et al., arXiv:2105.13549 [astro-ph.CO]

 $S_8 = 0.759^{+0.025}_{-0.025}$

DES-Y3 fiducial, Amon et al., arXiv:2105.13543 [astro-ph.CO]

• CMB Planck TT.TE.EE+lowE+lensing · CMB ACT+WMAP Aiola et al. (2020) Late Universe • WL KiDS-1000 Asgari et al. (2021) • WL KiDS+VIKING+DES-Y1 Asgari et al. (2020) • WL KiDS+VIKING+DES-Y1 Joudaki et al. (2020) • WL KiDS+VIKING-450 Wright et al. (2020) • WL KiDS+VIKING-450 • WL KiDS-450 • WL KiDS-450 • WL DES-Y3 • WL DES-Y1 Troxel et al. (2018) • WL HSC-TPCF Hamana et al. (2020) • WL HSC-pseudo-Cl Hikage et al. (2019) • WL CFHTLenS Joudaki et al. (2017) • WL+GC HSC+BOSS • WL+GC+CMBL KiDS+DES+eBOSS+Planck • WL+GC KiDS-1000 3×2pt • WL+GC KiDS-450 3×2pt • WL+GC DES-Y3 3×2pt Abbott et al. (2021) • WL+GC DES-Y1 3×2pt • WL+GC KiDS+VIKING-450+BOSS Tröster et al. (2020) • WL+GC KiDS+GAMA 3x2pt • GC BOSS DR12 bispectrum Philcox et al. (2021) • GC BOSS+eBOSS Ivanov et al. (2021) • GC BOSS power spectra Chen et al. (2021) • GC BOSS DR12 Tröster et al. (2020) GC BOSS galaxy power spectrum Ivanov et al. (2020) • GC+CMBL DELS+Planck White et al. (2022) • GC+CMBL unWISE+Planck • CC AMICO KiDS-DR3 Lesci et al. (2021) · CC DES-Y1 Abbott et al. (2020d) CC SDSS-DR8 • CC XMM-XXL Pacaud et al. (2018) • CC ROSAT (WtG) Mantz et al. (2015) • CC SPT tSZ Bocquet et al. (2019) CC Planck tSZ Salvati et al. (2018) CC Planck tSZ Ade et al. (2016d) • RSD Benisty (2021) • RSD 0.2 0.4 0.6 0.8 1.0 $S_8 \equiv \sigma_8 \sqrt{\Omega_m / 0.3}$

CMB Planck TT.TE.EE+lowE

Aghanim et al. (2020d) Aghanim et al. (2020d)

Early Universe



García-García et al. (2021) Heymans et al. (2021) Joudaki et al. (2018) Abbott et al. (2018d) van Uitert et al. (2018)

Krolewski et al. (2021)

Costanzi et al. (2019)

Kazantzidis and Perivolaropoulos (2018)

1.2

See Di Valentino et al. Astropart. Phys. 131 (2021) 102604 and Abdalla et al., arXiv:2203.06142 [astro-ph.CO] for a summary of the possible candidates proposed to solve the S8 tension.

Abdalla et al., JHEAp 34 (2022) 49-211

Early solutions to the H0 tension

Actually, a dark energy model that merely changes the value of rd would not completely resolve the tension, since it will affect the inferred value of Ω m and transfer the tension to it.

This is a plot illustrating that achieving a full agreement between CMB, BAO and SH0ES through a reduction of rd requires a higher value of $\Omega_m h^2$.



Jedamzik et al., Commun.in Phys. 4 (2021) 123

Early solutions to the H0 tension

Model 2 is defined by the simultaneous fit to BAO and CMB acoustic peaks at $\Omega_m h^2 = 0.155$, while model 3 has $\Omega_m h^2 = 0.167$

The sound horizon problem should be considered not only in the plane H0–rd, but it should be extended to the parameters triplet H0–rd– Ω m.

The figure shows that when attempting to find a full resolution of the Hubble tension, with CMB, BAO and SH0ES in agreement with each other, one exacerbates the tension with DES and KiDS.



Jedamzik et al., Commun.in Phys. 4 (2021) 123

Successful models?

This is the density of the proposed cosmological models:

At the moment no specific proposal makes a strong case for being highly likely or far better than all others !!!



Di Valentino et al., Class.Quant.Grav. (2021), arXiv:2103.01183 [astro-ph.CO]
...but a systematic error in Planck could explain S8...

A_L internal anomaly

The lensing effect on the power spectrum is the smoothing of the acoustic peaks, increasing AL.

Interesting consistency checks is if the amplitude of the smoothing effect in the CMB power spectra matches the theoretical expectation AL = 1 and whether the amplitude of the smoothing is consistent with that measured by the lensing reconstruction.

If AL =1 then the theory is correct, otherwise we have a new physics or systematics.



Calabrese et al., Phys. Rev. D, 77, 123531

A_L : a failed consistency check

The Planck lensing-reconstruction power spectrum is consistent with the amplitude expected for LCDM models that fit the CMB spectra, so the Planck lensing measurement is compatible with AL = 1.

However, the distributions of AL inferred from the CMB power spectra alone indicate a preference for AL > 1.

The joint combined likelihood shifts the value preferred by the TT data downwards towards AL = 1, but the error also shrinks, increasing the significance of AL > 1 to 2.8σ .

The preference for high AL is not just a volume effect in the full parameter space, with the best fit improved by $\Delta\chi^2 \sim 9$ when adding AL for TT+lowE and 10 for TTTEEE+lowE.



 $A_{\rm L} = 1.243 \pm 0.096$ (68 %, *Planck* TT+lowE), $A_{\rm L} = 1.180 \pm 0.065$ (68 %, *Planck* TT,TE,EE+lowE),

A_L can explain the S8 tension



Di Valentino, Melchiorri and Silk, JCAP 2001 (2020) no.01, 013

A_L that is larger than the expected value at about 3 standard deviations even when combining the Planck data with BAO and supernovae type Ia external datasets.

A_L can be a systematic?

Di Valentino and Melchiorri, 2022 ApJL 931 L18

Constraints at 68% CL	$\Sigma m_{ u} [{ m eV}]$
${\rm Planck}(+A_{\rm lens})$	< 0.51
$\mathrm{Planck}{+}\mathrm{BAO}~(+A_{\mathrm{lens}})$	< 0.19
$\mathrm{Planck+Pantheon}~(+A_{\mathrm{lens}})$	< 0.25
$\mathrm{Planck+Lensing}\;(+A_{\mathrm{lens}})$	$0.41\substack{+0.17 \\ -0.25}$
ACT-DR4+WMAP	0.68 ± 0.31
ACT-DR4+WMAP+BAO	< 0.19
ACT-DR4+WMAP+Pantheon	< 0.25
ACT-DR4+WMAP+Lensing	0.60 ± 0.25
SPT-3G+WMAP	$0.46\substack{+0.14 \\ -0.36}$
SPT-3G+WMAP+BAO	$0.22\substack{+0.056\\-0.14}$
${ m SPT-3G+WMAP+Pantheon}$	$0.25\substack{+0.052\\-0.19}$
SPT-3G+WMAP+Lensing	< 0.37

A combination of Planck CMB+Lensing constrain $\Sigma m_V = 0.41^{+0.17}_{-0.25} \text{ eV}$ at 68% CL when variation in the Alens parameter are considered.

Di Valentino and Melchiorri, 2022 ApJL 931 L18

Constraints at 68% CL	$\Sigma m_{ u} [{ m eV}]$
${\rm Planck}(+A_{\rm lens})$	< 0.51
$\mathrm{Planck}{+}\mathrm{BAO}~(+A_{\mathrm{lens}})$	< 0.19
$\mathrm{Planck+Pantheon}~(+A_{\mathrm{lens}})$	< 0.25
${ m Planck+Lensing}\;(+A_{ m lens})$	$0.41\substack{+0.17 \\ -0.25}$
ACT-DR4+WMAP	0.68 ± 0.31
ACT-DR4+WMAP+BAO	< 0.19
ACT-DR4+WMAP+Pantheon	< 0.25
ACT-DR4+WMAP+Lensing	0.60 ± 0.25
SPT-3G+WMAP	$0.46\substack{+0.14 \\ -0.36}$
SPT-3G+WMAP+BAO	$0.22\substack{+0.056\\-0.14}$
${ m SPT-3G+WMAP+Pantheon}$	$0.25\substack{+0.052\\-0.19}$
SPT-3G+WMAP+Lensing	< 0.37

We found that both the ACT-DR4 and SPT-3G data, when combined with WMAP, mildly suggest a neutrino mass with $\Sigma m_V = 0.68 \pm 0.31$ eV and $\Sigma m_V = 0.46^{+0.14}$ -0.36 eV at 68% CL,

respectively.

Constraints at 68% CL	$\Sigma m_{ u}$ [eV]
$\mathrm{Planck}\;(+A_{\mathrm{lens}})$	< 0.50
$\mathrm{Planck}{+}\mathrm{BAO}~(+A_{\mathrm{lens}})$	< 0.22
$Planck+Pantheon (+A_{lens})$	< 0.47
$\mathrm{Planck+Lensing}~(+A_{\mathrm{lens}})$	$0.38\substack{+0.12 \\ -0.28}$
ACT-DR4+WMAP	0.81 ± 0.28
ACT-DR4+WMAP+BAO	< 0.27
ACT-DR4+WMAP+Pantheon	0.71 ± 0.28
ACT-DR4+WMAP+Lensing	0.56 ± 0.21
ACT-DR4+WMAP+R20	0.83 ± 0.230
ACT-DR4+WMAP+F21	$0.85\substack{+0.27\\-0.33}$
ACT-DR4+WMAP+BAO+R20	$0.39\substack{+0.13\\-0.25}$
ACT-DR4+WMAP+BAO+F21	< 0.34
SPT-3G+WMAP	< 0.56
SPT-3G+WMAP+BAO	< 0.28
SPT-3G+WMAP+Pantheon	$0.46\substack{+0.11 \\ -0.39}$
SPT-3G+WMAP+Lensing	< 0.39
SPT-3G+WMAP+R20	$0.49\substack{+0.12\\-0.42}$
SPT-3G+WMAP+F21	< 0.60
SPT-3G+WMAP+BAO+R20	$0.37\substack{+0.13 \\ -0.25}$
SPT-3G+WMAP+BAO+F21	< 0.32

Di Valentino and Melchiorri, 2022 ApJL **931** L18

What about the 10 parameters extended model?

ACT-DR4 suggests a neutrino mass with $\Sigma m_V = 0.81 \pm 0.28 \text{ eV}$ and SPT-3G $\Sigma m_V < 0.56 \text{ eV}$ at 68% CL.



Di Valentino and Melchiorri, 2022 *ApJL* **931** L18

What about the 10 parameters extended model?

We can notice that the inclusion of the R20 or F21 priors does not affect the total neutrino mass constraints, because H0 and Σm_V do not show any correlation.

Moreover, in this case it is alleviated also the S₈ tension with the weak lensing measurements. We find for ACT-DR4 + WMAP + R20 S₈ = 0.726 \pm 0.037 at 68% C.L., while for SPT-3G + WMAP + R20 S₈ = 0.732 \pm 0.037 at 68% C.L..

Constraints at 68% CL	$\Sigma m_{ u}$ [eV]
$\rm Planck~(+A_{lens})$	< 0.50
$\mathrm{Planck}{+}\mathrm{BAO}~(+A_{\mathrm{lens}})$	< 0.22
${ m Planck+Pantheon} \; (+A_{ m lens})$	< 0.47
$\mathrm{Planck+Lensing}~(+A_{\mathrm{lens}})$	$0.38\substack{+0.12 \\ -0.28}$
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ACT-DR4+WMAP+Lensing	0.56 ± 0.21
ACT-DR4+WMAP+R20	0.83 ± 0.230
ACT-DR4+WMAP+F21	$0.85\substack{+0.27 \\ -0.33}$
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ACT-DR4+WMAP+BAO+F21	< 0.34
SPT-3G+WMAP	< 0.56
SPT-3G+WMAP+BAO	< 0.28
SPT-3G+WMAP+Pantheon	$0.46\substack{+0.11 \\ -0.39}$
SPT-3G+WMAP+Lensing	< 0.39
SPT-3G+WMAP+R20	$0.49\substack{+0.12 \\ -0.42}$
SPT-3G+WMAP+F21	< 0.60
SPT-3G+WMAP+BAO+R20	$0.37\substack{+0.13 \\ -0.25}$
SPT-3G+WMAP+BAO+F21	< 0.32

Di Valentino and Melchiorri, 2022 ApJL **931** L18

What about the 10 parameters extended model?

ACT-DR4 suggests a neutrino mass with $\Sigma m_V = 0.81 \pm 0.28 \text{ eV}$ and SPT-3G $\Sigma m_V < 0.56 \text{ eV}$ at 68% CL.

Interestingly, this indication for a higher neutrino mass is present also the BAO and Pantheon data are considered.



Di Valentino and Melchiorri, 2022 ApJL **931** L18

What about the 10 parameters extended model?

When CMB and BAO constraints are considered in these extended cosmologies, they provide constraints on the Σm_V vs H0 plane that clearly show a correlation between these two parameters, that is exactly the opposite of what is obtained under standard LCDM.

10 parameters

standard LCDM



Di Valentino and Melchiorri, arXiv:2112.02993 [astro-ph.CO]

Di Valentino et al. Phys.Rev. D93 (2016) no.8, 083527



Di Valentino and Melchiorri, 2022 ApJL **931** L18

What about the 10 parameters extended model?

Therefore, in extended cosmologies that can solve the Hubble tension, a neutrino mass is preferred by the cosmological data: ACT-DR4+BAO+R20 gives $\Sigma m_V = 0.39^{+0.13}$ -0.25 eV, while SPT-3G+BAO+R20 $\Sigma m_V = 0.60^{+0.44}$ -0.50 eV at 68% CL.

Concluding

There are many anomalies and tensions involving the CMB data:

- H0 tension
- S8 tension
- $A_L > 1$ for Planck
- Σm_V for ACT and SPT
- EDE for ACT

presenting a serious limitation to the precision cosmology.

Are we sure that the CMB results are still a confirmation of the flat standard ACDM cosmological model?

At this point, given the quality of all the analyses, probably these discrepancies are indicating a problem with the underlying cosmology and our understanding of the Universe, rather than the presence of systematic effects.

These cosmic discordances

call for new observations and stimulate the investigation of alternative theoretical models and solutions.

Thank you!

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