



SUSY (g-2)_µ With & Without Neutralino Dark Matter

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$(g - 2)_{\mu}$ anomaly

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Motivation

• There are many BSM scenarios that can explain the $(g-2)_{\mu}$ anomaly:

Leptoquarks, Z', VLL, 2HDM, axion, ...

• Supersymmetry is particularly motivated since it offers:

Coupling Unification, Radiative EWSB, Baryogenesis, DM, ...

• There are many studies on SUSY g-2 already:

[Athrona, Balazsa, Jacoba, Kotlarskic, Stockingerc, Stockinger-Kim]; [Chakraborti, Heinemeyer, Saha]; [Endo,Hamaguchi,Iwamoto,Kitahara]; [Cox, Han, Yanagida]; [Baum, Carena, Shah, Wagner]; [Badziak, KS]; [Hagiwara,Ma,Mukhopadhyay'18], ...

- Most studies assume the neutralino is the Lightest SUSY Particle (LSP) and stable.
 - Q: What happens if neutralino is unstable? (e.g. RPV, Gravitino LSP)
 - A: DM constraints go away, but LHC constraints change. How?















- Bino has very small annihilation cross-section
 - → Tend to produce too much DM
- Large off-diagonal term in stau mass matrix:
 - charge breaking vacuum: $m_{stau1}^2 > 0$
 - LEP bound: $m_{stau1} > 90 \text{ GeV}$
 - stau LSP: m_{stau1} > m_{neutralino1}
 - Vacuum (meta-)stability



($\tilde{\tau}$ mass matrix) ~







SUSY g-2 has a tension with:

- DM Direct Detection
- (Bino-like) DM overproduction
- lepton + large E_Tmiss @ LHC
- Vacuum stability (for BLR)

consequence of stable neutralino

How the situation improves / deteriorates if **neutralino is unstable**?

Analysis



List of ATLAS & CMS searches included in our analysis

13 TeV

8 TeV

| Name | $E/{\rm TeV}$ | $\mathcal{L}/\mathrm{fb}^{-1}$ | Description |
|------------------------|---------------|--------------------------------|--|
| atlas_1604_01306 | 13 | 3.2 | Monophoton |
| $atlas_{1605_{09318}}$ | 13 | 3.3 | 3 b-jets + 0.1 lepton + MET |
| $atlas_{1609_{01599}}$ | 13 | 36 | Monophoton |
| $atlas_1704_03848$ | 13 | 36 | Monophoton |
| atlas_conf_2015_082 | 13 | 3.2 | $2 \text{ leptons } (\mathbf{Z}) + \text{jets} + \text{MET}$ |
| atlas_conf_2016_013 | 13 | 3.2 | 1 lepton + jets (4 tops, VVL quarks) |
| atlas_conf_2016_050 | 13 | 13.3 | 1 lepton + (b) jets + MET |
| $atlas_conf_2016_054$ | 13 | 13.3 | 1 lepton + (b) jets + MET |
| atlas_conf_2016_076 | 13 | 13.3 | 2 lepton + jets + MET |
| $atlas_conf_2016_096$ | 13 | 13.3 | Multi-lepton + MET |
| atlas_conf_2017_060 | 13 | 36 | Monojet |
| $atlas_conf_2016_066$ | 13 | 13.3 | Photons, jets and MET |
| atlas_1712_08119 | 13 | 36 | soft leptons (compressed EWKinos) |
| atlas_1712_02332 | 13 | 36 | squarks and gluinos, 0 lepton, 2-6 jets |
| $atlas_1709_04183$ | 13 | 36 | Jets + MET (stops) |
| $atlas_{1802_{03158}}$ | 13 | 36 | search for GMSB with photons |
| atlas_1708_07875 | 13 | 36 | EWKino search with taus and MET |
| atlas_1706_03731 | 13 | 36 | Multilepton + Jets + MET (RPC and RPV) |
| atlas_1908_08215 | 13 | 36 | 2 leptons + MET (EWKinos) |
| atlas_1909_08457 | 13 | 139 | SS lepton + MET (squark, gluino) |
| atlas_conf_2019_040 | 13 | 139 | Jets + MET (squark, gluino) |
| atlas_conf_2019_020 | 13 | 139 | 3 leptons (EWKino) |
| atlas_ 1803_02762 | 13 | 36 | 2 or 3 leptons (EWKino) |
| $atlas_conf_2018_041$ | 13 | 80 | Multi- <i>b</i> -jets (stops, sbottoms) |
| atlas_2101_01629 | 13 | 139 | 1 lepton + jets + MET |
| $atlas_conf_2020_048$ | 13 | 139 | Monojet |
| atlas_2004_14060 | 13 | 139 | $t\bar{t} + MET$ |
| atlas_1908_03122 | 13 | 139 | Higgs bosons $+ b$ -jets $+ MET$ |
| atlas_2103_11684 | 13 | 139 | 4 or more leptons (RPV, GMSB) |
| atlas_2106_09609 | 13 | 139 | Multijets + leptons (RPV) |
| atlas_1911_06660 | 13 | 139 | Search for Direct Stau Production |
| $cms_pas_sus_15_011$ | 13 | 2.2 | 2 leptons + jets + MET |
| $cms_sus_16_039$ | 13 | 35.9 | electrowekinos in multilepton final state |
| $cms_sus_16_025$ | 13 | 12.9 | electroweakino and stop compressed spectra |
| $cms_sus_16_048$ | 13 | 35.9 | two soft opposite sign leptons |

| Name | E/TeV | $\mathcal{L}/\mathrm{fb}^{-1}$ | Description |
|--------------------------|-------|--------------------------------|---|
| atlas_1308_1841 | 8 | 20.3 | $0 \text{ lepton} + \ge 7 \text{ jets} + \text{MET}$ |
| atlas_1308_2631 | 8 | 20.1 | 0 leptons + 2 b-jets + MET |
| $atlas_{-}1402_{-}7029$ | 8 | 20.3 | 3 leptons + MET (chargino+neutralino) |
| atlas_1403_4853 | 8 | 20.3 | 2 leptons + MET (direct stop) |
| atlas_1403_5222 | 8 | 20.3 | stop production with Z boson and b-jets |
| atlas_ 1404_2500 | 8 | 20.3 | Same sign dilepton or 3 lepton |
| atlas_ 1405_7875 | 8 | 20.3 | 0 lepton + 2-6 jets + MET |
| atlas_1407_0583 | 8 | 20.3 | ATLAS, 1 lepton $+$ (b-)jets $+$ MET (stop) |
| atlas_1407_0608 | 8 | 20.3 | Monojet or charm jet (stop) |
| $atlas_{1411_{1559}}$ | 8 | 20.3 | monophoton plus MET |
| atlas_1501_07110 | 8 | 20.3 | 1 lepton + 125 GeV Higgs + MET |
| $atlas_{1502}01518$ | 8 | 20.3 | Monojet + MET |
| atlas_ 1503_03290 | 8 | 20.3 | 2 leptons + jets + MET |
| atlas_ 1506_08616 | 8 | 20.3 | di-lepton and $2b$ -jets + lepton |
| atlas_ 1507_05493 | 8 | 20.3 | photonic signatures of gauge-mediated SUSY |
| $atlas_conf_2012_104$ | 8 | 20.3 | $1 \text{ lepton} + \ge 4 \text{ jets} + \text{MET}$ |
| $atlas_conf_2013_024$ | 8 | 20.3 | 0 leptons + 6 (2 b-) jets + MET |
| atlas_conf_2013_049 | 8 | 20.3 | 2 leptons + MET |
| atlas_conf_2013_061 | 8 | 20.3 | $0-1 \text{ leptons} + \ge 3 \text{ b-jets} + \text{MET}$ |
| $atlas_conf_2013_089$ | 8 | 20.3 | 2 leptons (razor) |
| $atlas_conf_2015_004$ | 8 | 20.3 | invisible Higgs decay in VBF |
| atlas_1403_5294 | 8 | 20.3 | 2 leptons + MET, (SUSY electroweak) |
| atlas_higg_2013_03 | 8 | 20.3 | 2 leptons + MET, (invisible Higgs) |
| $atlas_{1502}05686$ | 8 | 20.3 | search for massive sparticles decaying to many jets |
| cms_1303_2985 | 8 | 11.7 | α_T + b-jets |
| cms_1408_3583 | 8 | 19.7 | monojet + MET |
| cms_1502_06031 | 8 | 19.4 | 2 leptons, jets, MET (only on-Z) |
| cms_1504_03198 | 8 | 19.7 | 1 lepton, \geq 3 jets, \geq 1 b-jet, MET (DM + 2 top) |
| cms_sus_13_016 | 8 | 19.5 | OS lepton $3+$ b-tags |
| | | | |











All g-2 region will be probed by the next generation DM-DD experiments



Unstable Neutralino

We study 2 example-scenarios with unstable neutralino



UDD RPV

stable neutralino

 $\tilde{m}_{l_{L}} = \min(M_{2}, |\mu|) + 20 \text{GeV}, \ \tan \beta = 50, \ A = 0, \ \tilde{m}_{l_{R}} = M_{1} = 10 \text{TeV}$



 M_2 [GeV]

ATLAS multijet+l [2106.09609]

CMS multilepton [1709.05406]

ATLAS jets +E^{Tmiss} [ATLAS-CONF-2019-040]



UDD RPV

stable neutralino







Gravitino LSP



 γ, Z, h $\tilde{h}^{0,\pm}, \tilde{W}^{0,\pm} \qquad \tilde{G}$ $\tilde{h}^{0,\pm}, \tilde{W}^{0,\pm} \qquad \tilde{G}$ γ, Z, h



WHL plane:

$$(M_2 \text{ vs } \mu) \text{ with } \tilde{m}_{l_L} = \min(M_2, \mu) + 20 \text{ GeV} \implies m_{l_L} = \min(M_2, \mu) - 20 \text{ GeV}$$

BHL plane:
$$\widetilde{\nu}_L \text{ NLSP}$$

$$(M_1 \text{ vs } \mu) \text{ with } \tilde{m}_{l_L} = \min(M_1, \mu) + 20 \text{ GeV} \implies m_{l_L} = \min(M_2, \mu) - 20 \text{ GeV}$$

BHR plane:

$$(M_1 \text{ vs } \mu) \text{ with } \tilde{m}_{l_R} = \min(M_1, |\mu|) + 20 \text{ GeV} \implies m_{l_R} = \min(M_1, \mu) - 20 \text{ GeV} \quad \begin{cases} \tilde{e}_R, \tilde{\mu}_R, \tilde{\tau}_R \text{ NLSP} \end{cases}$$

BLR plane:

$$(\tilde{m}_{l_L} \text{ vs } \tilde{m}_{l_R}) \text{ with } M_1 = m_{\tilde{\tau}_1} - 20 \,\text{GeV} \implies M_1 = m_{\tilde{\tau}_1} + 20 \,\text{GeV}$$

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MSSM with stable neutralino:





UDD RPV :



Summary

- SUSY might be a solution to the $(g-2)_{\mu}$ anomaly

 - slepton-gaugino-Higgsino are light \implies stringent constraint from DM-DD detection
 - LR slepton and Bino are light \implies Bino overproduction
- If $\tilde{\chi}_1^0$ is not stable LSP, DM constraints go away, and LHC signature changes.
 - (1) RPV with UDD \implies LHC constraints from multijet + lepton
 - 2 Gravitino LSP with $\tilde{\chi}_1^0$ NLSP \implies (g-2)_µ region excluded by $\gamma + E_T$ channel
 - ③ Gravitino LSP with *non* $\tilde{\chi}_1^0$ NLSP \implies LHC constraints from soft lepton/tau

Explanation for $(g-2)_{\mu}$ anomaly is possible for the scenarios (1) and (3)





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Understanding the Early Universe: interplay of theory and collider experiments

Joint research project between the University of Warsaw & University of Bergen

Parameter planes definition

| name | axes | range [TeV] | other parameters | $\tan \beta$ |
|-------------------------------|---|--------------------------|---|--------------|
| \mathbf{WHL}_{μ} | (M_2,μ) | ([0.2,4],[0.2,4]) | $\tilde{m}_{l_{ m L}} = \min(M_2, \mu) + 20 { m GeV}, \ M_1 = \tilde{m}_{l_{ m R}} = 10 { m TeV}$ | 50 |
| \mathbf{WHL}_L | $(M_2,	ilde{m}_{l_{ m L}})$ | ([0.2,4],[0.2,2]) | $\mu = \min(M_2, \tilde{m}_{l_{\rm L}}) - 20 { m GeV}, \ M_1 = \tilde{m}_{l_{ m R}} = 10 { m TeV}$ | 50 |
| $\mathbb{B}\mathrm{HL}_{\mu}$ | (M_1,μ) | ([0.12,0.6],[0.12,0.35]) | $\tilde{m}_{l_{ m L}} = \min(M_1, \mu) + 20 { m GeV}, \ M_2 = \tilde{m}_{l_{ m R}} = 10 { m TeV}$ | 50 |
| \mathbf{BHL}_L | $(M_1, 	ilde{m}_{l_{ m L}})$ | ([0.12,0.8],[0.14,0.22]) | $\mu = \min(M_1, \tilde{m}_{l_{\rm L}}) - 20 { m GeV}, \ M_2 = \tilde{m}_{l_{ m R}} = 10 { m TeV}$ | 50 |
| \mathbf{BHR}_{μ} | (M_1, μ) | ([0.12,0.7],[0.12,0.7]) | $\tilde{m}_{l_{\mathrm{R}}} = \min(M_1, \mu) + 20 \mathrm{GeV}, \ M_2 = \tilde{m}_{l_{\mathrm{L}}} = 10 \mathrm{TeV}$ | 50 |
| \mathbf{BHR}_L | $(M_1,	ilde{m}_{l_{ m R}})$ | ([0.12,0.8],[0.14,0.25]) | $-\mu = \min(M_1, \tilde{m}_{l_{\mathrm{R}}}) - 20 \mathrm{GeV}, \ M_2 = \tilde{m}_{l_{\mathrm{L}}} = 10 \mathrm{TeV}$ | 50 |
| \mathbf{BLR}_{50} | $(ilde{m}_{l_{ m L}},	ilde{m}_{l_{ m R}})$ | ([0.15,0.6],[0.12,1.2]) | $M_1 = m_{	ilde{	au}_1} - 20{ m GeV}, \;\; \mu = \mu_{ m max}, \;\; M_2 = 10{ m TeV}$ | 50 |
| \mathbf{BLR}_{10} | $(ilde{m}_{l_{ m L}},	ilde{m}_{l_{ m R}})$ | ([0.15,0.6],[0.12,1.2]) | $M_1 = m_{	ilde{	au}_1} - 20 { m GeV}, \ \ \mu = \mu_{ m max}, \ \ M_2 = 10 { m TeV}$ | 10 |

Table 1: The parameter planes and choices of the other parameters. μ_{max} is defined as the maximum value allowed by the vacuum stability constraint.

For GMSB we modify the planes to ensure that slepton/stau/sneutrino is the NLSP.

QEDHVPEW
$$a_{\mu}^{\text{theo}} = 0.00$$
116591810(43) $a_{\mu}^{\exp} = 0.00$ 116592061(41)

• The deviation is size of the EW correction in SM:

$$a_{\mu}^{\exp} - a_{\mu}^{\text{theo}} \simeq (25 \pm 6) \times 10^{-10} \sim \mathcal{O}\left(\Delta a_{\mu}^{\text{SM,EW}}\right)$$

• We need very light BSM particles **OR** enhancement from couplings

$$\Delta a_{\mu}^{\text{BSM}} \sim \Delta a^{\text{SM,EW}} \cdot \left(\frac{m_W^2}{m_{\text{BSM}}^2}\right) \cdot \left(\frac{g_{\text{BSM}}}{g_{\text{SM}}}\right)$$

$$\underbrace{\mathcal{O}(1)}$$

Chiral (tanß) enhancement in SUSY

• (g-2) operator requires chirality flip:

$$\mathscr{L}_{\rm eff} \ni i\widetilde{a}_{\mu} \cdot \bar{\psi}_{L} \sigma^{\mu\nu} \psi_{R} F_{\mu\nu}$$

$$\vec{\mu} = g\left(\frac{e}{2m}\right)\vec{s}$$
$$a_{\mu} = \frac{(g-2)}{2} \equiv m_{\mu}\tilde{a}_{\mu}$$

SM: $\widetilde{a}_{\mu}^{\rm SM} \propto Y_{\mu} \langle H \rangle = m_{\mu}$

Chiral (tanβ) enhancement in SUSY

• (g-2) operator requires chirality flip:

$$\mathscr{L}_{\rm eff} \ni i \widetilde{a}_{\mu} \cdot \bar{\psi}_{\underline{L}} \sigma^{\mu\nu} \psi_{\underline{R}} F_{\mu\nu}$$

$$\vec{\mu} = g\left(\frac{e}{2m}\right)\vec{s}$$
$$a_{\mu} = \frac{(g-2)}{2} \equiv m_{\mu}\tilde{a}_{\mu}$$

SM:
$$\widetilde{a}_{\mu}^{\rm SM} \propto Y_{\mu} \langle H \rangle = m_{\mu}$$

SUSY:
$$\Delta \widetilde{a}_{\mu}^{\text{SUSY}} \propto Y_{\mu} \langle H_{u} \rangle = m_{\mu} \cdot \tan \beta$$

$$\begin{pmatrix} m_{\mu} = Y_{\mu} \langle H_{d} \rangle & \tan \beta \equiv \frac{\langle H_{u} \rangle}{\langle H_{d} \rangle} \end{pmatrix}$$

 $\begin{array}{c} \langle H_{u} \rangle & \gamma \\ \tilde{H}_{u}^{+} & \tilde{W}^{+} & \tilde{W}^{+} \\ \tilde{H}_{d}^{+} & \tilde{\nu}_{\mu} & \tilde{W}^{+} \\ \hline \mu_{R} & \tilde{\nu}_{\mu} & \mu_{L} \end{array}$

$$\frac{\langle H_u \rangle^2 + \langle H_d \rangle^2}{\uparrow} = \langle H \rangle^2$$

$$(246 \,\text{GeV})^2$$

Chiral (tanß) enhancement in SUSY



- Due to strong LHC constraints, we decouple coloured SUSY particles (they do not contribute to (g-2)_μ anyway).
- a_{μ}^{SUSY} depends on 5 mass parameters and $tan\beta$:

| M_1 : Bino mass | $ m_{\tilde{l}_R} \equiv \widetilde{m}_{\tilde{e}_R}^2 = \widetilde{m}_{\tilde{\mu}_R}^2 = \widetilde{m}_{\tilde{\tau}_R}^2 $ |
|-----------------------|---|
| M_2 : Wino mass | $ \sqrt{m_{\tilde{l}_L}} \equiv \widetilde{m}_{\tilde{\nu}_e} = \widetilde{m}_{\tilde{\nu}_{\mu}} = \widetilde{m}_{\tilde{\nu}_{\tau}} = \widetilde{m}_{\tilde{e}_L} = \widetilde{m}_{\tilde{\mu}_L} = \widetilde{m}_{\tilde{\tau}_L} $ |
| μ : Higgsino mass | $\tan\beta \equiv \langle H_u \rangle / \langle H_d \rangle$ |

no LFV due to universal soft masses: avoid strong constraint from $\mu \rightarrow e \gamma$

$$\Delta a_{\mu}^{\text{SUSY}} = \Delta a_{\mu}^{\text{WHL}} + \Delta a_{\mu}^{\text{BHL}} + \Delta a_{\mu}^{\text{BHR}} + \Delta a_{\mu}^{\text{BLR}}$$



$$\Delta a_{\mu}^{\text{WHL}}(M_{2},\mu,m_{\tilde{l}_{L}}) = \frac{\alpha_{W}}{8\pi} \frac{m_{\mu}^{2}}{M_{2}\mu} \tan\beta \cdot f_{W}(\{\mathbf{m}\})$$
$$\Delta a_{\mu}^{\text{BHL}}(M_{1},\mu,m_{\tilde{l}_{L}}) = \frac{\alpha_{Y}}{8\pi} \frac{m_{\mu}^{2}}{M_{1}\mu} \tan\beta \cdot f_{N}(\{\mathbf{m}\})$$
$$\Delta a_{\mu}^{\text{BHR}}(M_{1},\mu,m_{\tilde{l}_{R}}) = -\frac{\alpha_{Y}}{8\pi} \frac{m_{\mu}^{2}}{M_{1}\mu} \tan\beta \cdot f_{N}(\{\mathbf{m}\})$$



 $M_1 : \text{Bino } (\tilde{B}) \text{ mass}$ $M_2 : \text{Wino } (\tilde{W}) \text{ mass}$ $\mu : \text{Higgsino } (\tilde{H}_u, \tilde{H}_d) \text{ mass}$

Large gaugino-Higgsino mixing leads to a large cross-section for DM Direct Detection:



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Constraints:

Stau mass² becomes negative or too small!

 $(\tilde{\tau} \text{ mass matrix}) \sim \begin{pmatrix} m_{\tilde{\tau}_R}^2 & Y_{\tau} \mu \langle H_u \rangle \\ Y_{\tau} \mu \langle H_u \rangle & m_{\tilde{\tau}_T}^2 \end{pmatrix}$

- charge breaking vacuum: $m_{stau1}^2 > 0$

- LEP bound: m_{stau1} > 90 GeV
- stau LSP: m_{stau1} > m_{neutralino1}
- Vacuum (meta-)stability:

$$\left| m_{\tilde{\ell}_{LR}}^2 \right| \leq \left[1.01 \times 10^2 \,\text{GeV}\sqrt{m_{\tilde{\ell}_L} m_{\tilde{\ell}_R}} + 1.01 \times 10^2 \,\text{GeV}(m_{\tilde{\ell}_L} + 1.03m_{\tilde{\ell}_R}) - 2.27 \times 10^4 \,\text{GeV}^2 + \frac{2.97 \times 10^6 \,\text{GeV}^3}{m_{\tilde{\ell}_L} + m_{\tilde{\ell}_R}} - 1.14 \times 10^8 \,\text{GeV}^4 \left(\frac{1}{m_{\tilde{\ell}_L}^2} + \frac{0.983}{m_{\tilde{\ell}_R}^2} \right) \right]$$

[Kitahara, Yoshinaga 13]; [Endo, Hamaguchi, Kitahara, Yoshinaga 13]

* Overproduction of Bino-like neutralinos in the early universe: $\Omega_{\tilde{\chi}_1^0} < \Omega_{DM}$

slepton-coannihilation needed $\implies m_{slepton} \sim m_{Bino}$

Unstable Neutralino (Gravitino, RPV)

$$\Delta a_{\mu}^{\text{SUSY}} = \Delta a_{\mu}^{\text{WHL}} + \Delta a_{\mu}^{\text{BHL}} + \Delta a_{\mu}^{\text{BHR}} + \Delta a_{\mu}^{\text{BLR}}$$

 $\Delta a_{\mu}^{\text{WHL}}(M_{2}, \mu, m_{\tilde{l}_{L}})$ $\Delta a_{\mu}^{\text{BHL}}(M_{1}, \mu, m_{\tilde{l}_{L}})$ $\Delta a_{\mu}^{\text{BHR}}(M_{1}, \mu, m_{\tilde{l}_{R}})$

Higgsino, one gaugino, one slepton all must be light: \Rightarrow LHC constraint with large $\not{E_T} \leftarrow$ Modified gaugino-Higgsino mixing \Rightarrow DM direct detection

 $\Delta a_{\mu}^{\text{BLR}}(M_{1}, m_{\tilde{l}_{L}}, m_{\tilde{l}_{R}}; \mu)$ $\uparrow \quad -$ large

Bino and both L and R sleptons must be light: \Rightarrow LHC constraint with large $\not{E}_{1} \leftarrow Modified$ \Rightarrow Bino abundance $\Omega_{\chi_{1}^{0}} \leftarrow \Omega_{DM}$ \Rightarrow Charged LSP, Vacuum stability

- These terms give mass to quarks and reptons.
- These terms give mass to quarks and leptons.



R-Parity Violation; UDD



No missing energy, but multi-jet

LHC signature is the most challenging:
 no leptons, no b-jets in the neutralino decay



R-Parity Violation; UDD





| Bin | Final state | Definition |
|-----|---------------------------|---|
| 1 | 2 SS leptons | 0 jets, $M_{\rm T}$ > 100 GeV and $p_{\rm T}^{\rm miss}$ > 140 GeV |
| 2 | 2 SS leptons | 1 jet , $M_{ m T} < 100{ m GeV}$, $p_{ m T}^{\ell\ell} < 100{ m GeV}$ and $p_{ m T}^{ m miss} > 200{ m GeV}$ |
| 3 | 3 light leptons | $M_{\rm T} > 120 {\rm GeV}$ and $p_{\rm T}^{\rm miss} > 200 {\rm GeV}$ |
| 4 | 3 light leptons | $p_{\rm T}^{\rm miss} > 250 { m GeV}$ |
| 5 | 2 light leptons and 1 tau | $M_{\rm T2}(\ell_1, \tau) > 50 {\rm GeV} \text{ and } p_{\rm T}^{\rm miss} > 200 {\rm GeV}$ |
| 6 | 1 light lepton and 2 taus | $M_{\rm T2}(\ell, \tau_1) > 50 {\rm GeV}$ and $p_{\rm T}^{\rm miss} > 200 {\rm GeV}$ |
| 7 | 1 light lepton and 2 taus | $p_{\rm T}^{\rm miss} > 75 { m GeV}$ |
| 8 | more than 3 leptons | $p_{\rm T}^{\rm miss} > 200 { m GeV}$ |

not be applied. a gravitino, $c\tau_{\tilde{\chi}_1^0} < 1 \,\mathrm{mm}$. In the lower right region, the NSL

Graphied D for the lightest neutralino into the gravitino are given by [13,35]

• In the gauge-mediated SUSY breaking (GMSB) scenario, Night gravitino is motivated by naturalness: $\Gamma(\tilde{\chi}_{1}^{0} \to \tilde{G}Z) \tilde{\chi}_{1}^{0} \to \tilde{G}Z) \tilde{\chi}_{1}^{0} \to \tilde{G}W_{12}c_{W}^{-1} + \frac{N_{12}}{2} |\tilde{\chi}_{11}|_{SW}^{0}|^{2} + \frac{N_{14}s_{W}}{2} |\tilde{\chi}_{FS}^{0} \to N_{14}s_{\beta}|^{2})$

 $\sum_{\substack{\text{Left: } c_{\tau_{z^0}} \\ \text{Left: } c_{\tau_{z^0}}$ The contrained sufficiency promptly into the contrained sufficiency promptly into the contrained at the bill of the second promptly into the contrained at the bill of the second promptly into the contrained at the bill of the second promptly into the contrained and our analysis may here N_{ij} is the neutralino mixing matrix and $\tilde{\chi}_{ij}^{1} \approx 1 \text{ mm}$. In the lower right region, the NSLP is long-lived and our analysis may here N_{ij} is the neutralino can be calculated. (For light $\int_{100}^{5} \left(\frac{m_{3/2}}{10 \text{ eV}}\right)^{1/2}$. The decay rate of the NLSP neutralino into the gravitino can be calculated. (For light $\int_{100}^{5} \left(\frac{m_{3/2}}{10 \text{ eV}}\right)^{1/2}$. est neutraline into the gravitine are given by [13,35] the neutralino decays are prompt/2 $M_{\rm pl}^2$ In the left pane $\begin{array}{c} \widetilde{\chi}_{1}^{0} \xrightarrow{\rightarrow} \widetilde{G} \widetilde{\gamma} \\ \widetilde{\chi}_{1}^{0} \xrightarrow{\rightarrow} \widetilde{G} \widetilde{\gamma} \end{array} \begin{array}{c} = \\ = \end{array}$ $N_{11}c_W + N_{12}s_W |_{2}^{\bar{2}}A$ 800 $\tilde{G} \xrightarrow{\tilde{G}} \tilde{G} \xrightarrow{\tilde{G}} \tilde{Z} \xrightarrow{\tilde{V}_{1}} \tilde{G} \xrightarrow{\tilde{U}_{1}} \tilde{G} \xrightarrow{\tilde{U}$ Bino: $c\tau_{\tilde{B}} = 1 \,\mathrm{mm}$ 700 Wino: $c\tau_{\tilde{W}} = 1 \,\mathrm{mm}$ $\tilde{\tilde{G}}_{1}^{0} \xrightarrow{} \tilde{\tilde{G}}_{1}^{R} \xrightarrow{} \tilde{\tilde{G}}_{R}^{R} = \frac{\left(\left| N_{12}c_{W} - N_{14}s_{\mathcal{B}} \right|^{2} - N_{14}s_{\mathcal{B}} \right)^{2} \right)^{2} + \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \right|^{2} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \right|^{2} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \right)^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{\tilde{\chi}(4)}^{2}} \right)^{4} \mathcal{A}, \frac{1}{2} \left| N_{13}c_{\beta} - N_{14}s_{\beta} \right|^{2} \left(1 - \frac{m_{Z}^{2}}{m_{Z}^{2}} \right)^{2} \left(1 - \frac{m_{Z}^{2}}{m_{Z}^{2}$ Higgsino: $c\tau_{\tilde{h}} = 1 \text{ mm}$ 600 part of the plat $\Gamma(\tilde{\chi}_1^0 \to \tilde{G}h) = \frac{1}{2} |N_{13}c_\beta + N_{14}s_\beta|^2 \left(1 - \frac{m_h^2}{m_{\tilde{\chi}_1^0}^2}\right)^4 \mathcal{A},$ s the neutralino mixing matrix and lighter than to (U 500 Prompt massless particl 400 dSTN*W* $c\tau_{\rm NLSP} < 1\,{\rm mm}$ $\mathcal{A} \equiv \frac{\frac{m_{20}}{m_{20}}}{\frac{1}{10\pi m_{2}^{2}/2}M_{10}^{2}} \approx \frac{\frac{1}{10\pi m_{10}} \left(\frac{m_{20}}{100 \text{ GeV}}\right)^{\frac{5}{3}} \left(\frac{m_{3/2}}{100 \text{ GeV}}\right)^{\frac{-2}{3}}$ 1 eV throughton $\Lambda - m_{\rm NLSP}$ plasses of Figure 1 we plot contours of a fixed neutralino litetime $C_{F_{2}}$ = mm in Vht ef Figure 1 we plot contours et a fixed neutraline lifetime $C_{F_{2}}$ = mm in Vht Non-Prompt 200 $-m_{\mathrm{N}}$ $c\tau_{\rm NLSP} > 1\,{\rm mm}$ TŁ 100 + 10⁰ 101 10² 10³ $m_{3/2}$ [eV] ifficie in dealing with its kinematics at colliders and we conveniently fix



Analysis Framework

SUSY g-2: 1-loop + leading 2-loop GM2Calc [Eur.Phys.J. C76 (2016) no.2, 62]

Neutralino abundance, Direct Detection: MicrOMEGAs [2003.08621]

Decay of SUSY particles: SUSY-HIT [hep-ph/0609292]

LHC constraints:

- **MSSM:** (1) Mapping simplified model limits to the model point (σ BR)
- **RPV:** ② Pythia 8 + CheckMATE 2 [1907.09874], [1611.09856]
- Gravitino LSP: Both 1) and 2)

Parameter planes



| | 2D planes | | |
|------|--|---|-------|
| | | | |
| name | axes | other parameters | aneta |
| WHL | (M_2,μ) | $\tilde{m}_{l_{\rm L}} = \min(M_2, \mu) + 20 \text{GeV}, \ M_1 = \tilde{m}_{l_{\rm R}} = 10 \text{TeV}$ | 50 |
| BHL | (M_1,μ) | $\tilde{m}_{l_{\rm L}} = \min(M_1, \mu) + 20 \text{GeV}, \ M_2 = \tilde{m}_{l_{\rm R}} = 10 \text{TeV}$ | 50 |
| BHR | (M_1, μ) | $\tilde{m}_{l_{\mathrm{R}}} = \min(M_1, \mu) + 20 \mathrm{GeV}, \ M_2 = \tilde{m}_{l_{\mathrm{L}}} = 10 \mathrm{TeV}$ | 50 |
| BLR | $(ilde{m}_{l_{ m L}}, 	ilde{m}_{l_{ m R}})$ | $M_1 = m_{\tilde{\tau}_1} - 20 \text{GeV}, \ \mu = \mu_{\text{max}}, \ M_2 = 10 \text{TeV}$ | 50 |





Non $\tilde{\chi}_1^0$ NLSP (Short Summary)



• small $|\mu|$ region is compatible with $(g-2)_{\mu}$