

Fingerprints of freeze-in dark matter in an early matter-dominated era

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Based on 2204.03670

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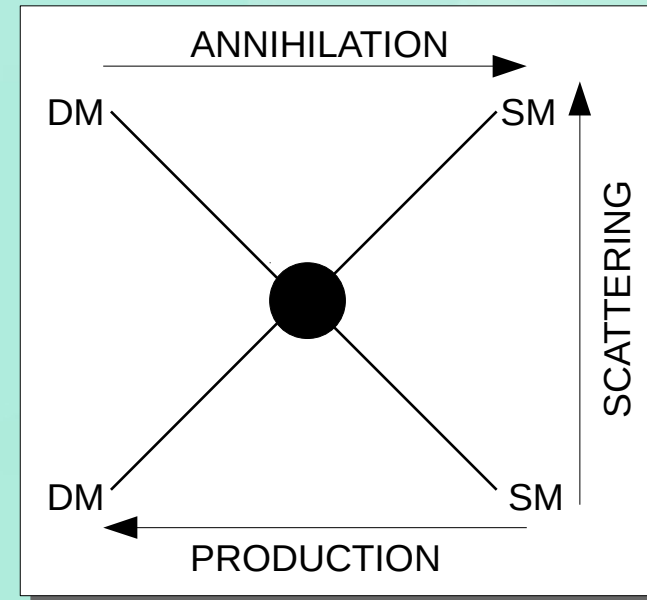
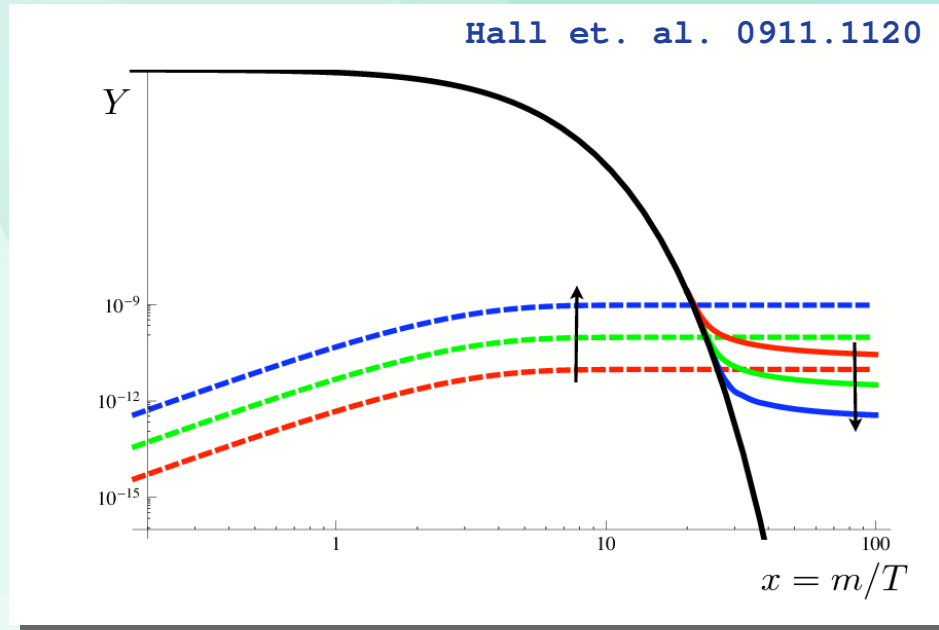
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Freeze-out vs. Freeze-in



| Freeze-out | Freeze-in |
|--|--|
| DM-SM in thermal equilibrium, Large coupling required | DM-SM never in thermal equilibrium, Extremely small coupling |
| At high temperature DM has thermal abundance | Initial abundance of DM at the end of inflation is negligible |
| Cosmology at high temperature is irrelevant | Cosmology from the end of inflation till today impacts relic |

Early matter domination

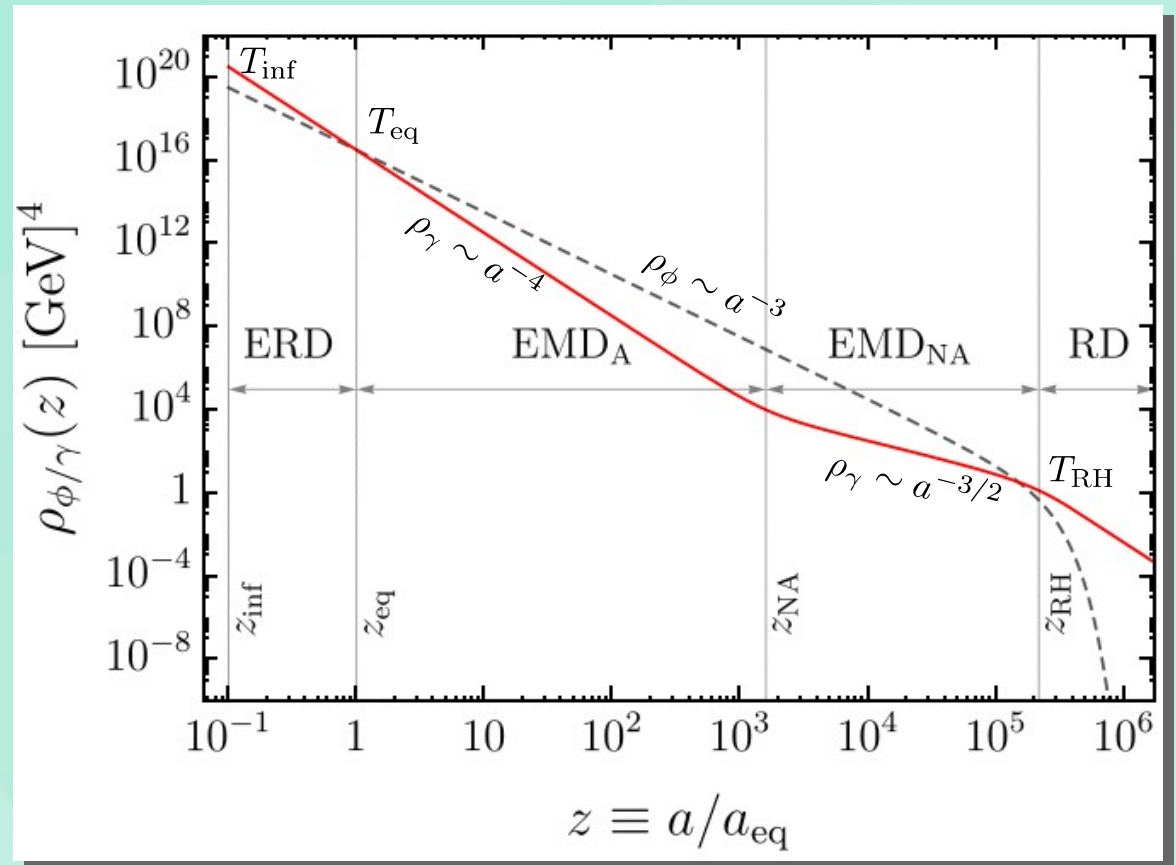
BSM motivations for EMD

- Meta-stable matter fields
- Oscillating scalar fields
- Moduli
- SUSY condensates
- Dilaton
- Q-balls
- Curvaton

$$\dot{\rho}_\phi + 3(1 + \omega)H\rho_\phi = -(1 + \omega)\Gamma_\phi\rho_\phi$$

$$\dot{\rho}_\gamma + 4H\rho_\gamma = (1 + \omega)\Gamma_\phi\rho_\phi$$

$$H = \frac{1}{\sqrt{3}M_p} \sqrt{\rho_\phi + \rho_\gamma}$$



Constraints from BBN: $T_{RH} \gtrsim \text{few MeV}$

Evolution is dependent on the dissipation rate

Matter dissipation rate

In general depends on the temperature and the expansion of the universe

$$\Gamma_\phi = \hat{\Gamma} \left(\frac{T}{T_{\text{eq}}} \right)^n \left(\frac{a}{a_{\text{eq}}} \right)^k$$

Examples:

**Oscillating scalar fields
with $V(\phi) \sim \phi^p$ potential**

$$\Gamma_{\phi \rightarrow f\bar{f}} \propto m_\phi(t) \propto a^{-3(p-2)/(p+2)}, \text{ (for fermionic decay),}$$

$$\Gamma_{\phi \rightarrow bb} \propto m_\phi^{-1}(t) \propto a^{3(p-2)/(p+2)}, \text{ (for bosonic decay),}$$

[Garcia et. al. 2012.10756](#)

[See talk by Keith Olive](#)

Moduli decay: $\Gamma_\phi \propto \frac{T^3}{M_p^2}$

[Bodeker, hep-ph/0605030](#)

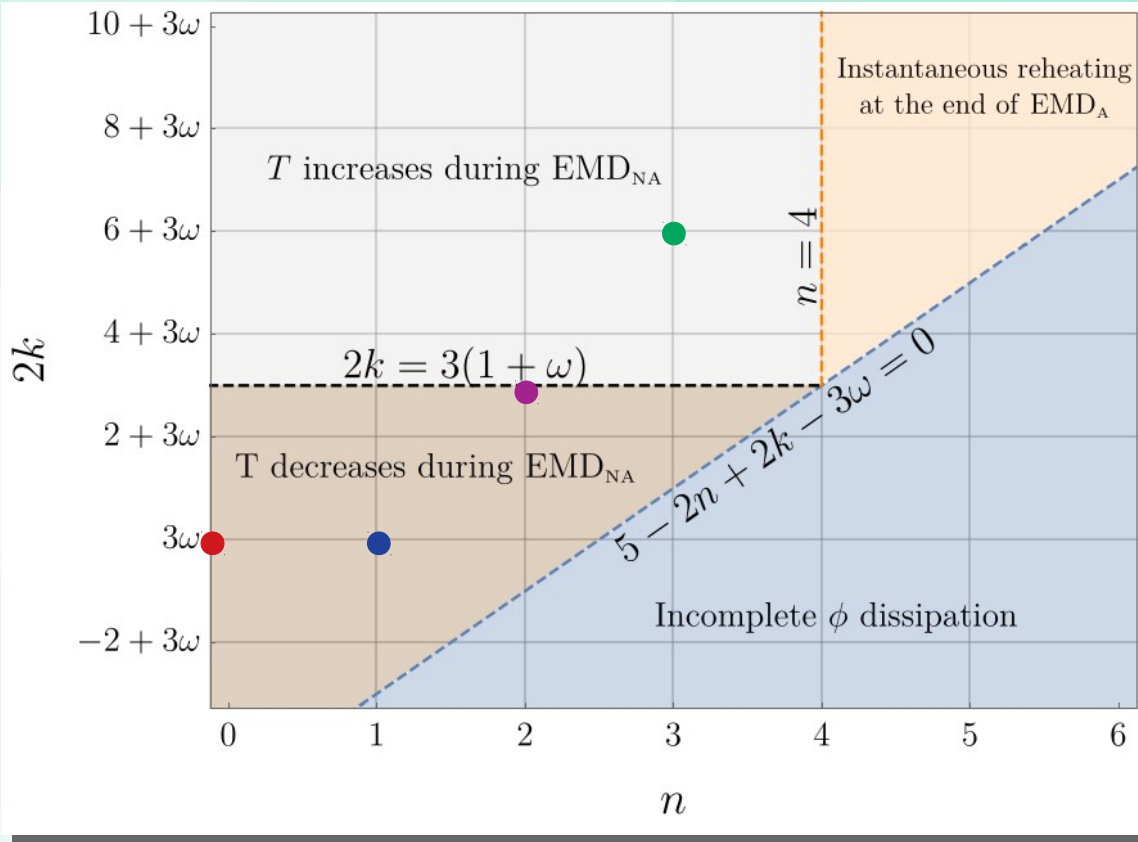
More Examples:

| Γ_ϕ | (n, k, ω) | $T(z)$ during EMD _{NA} |
|--------------------------------------|------------------|---------------------------------|
| const. | (0, 0, 0) | decreases with z |
| T | (1, 0, 0) | decreases with z |
| $\langle \phi \rangle^{-2}$ | (0, 3, 0) | increases with z |
| $\frac{T^3}{\langle \phi \rangle^2}$ | (3, 3, 0) | increases with z |
| $\frac{T^2}{\langle \phi \rangle}$ | (2, 3/2, 0) | remains constant |
| $\frac{T^2}{\langle \phi \rangle}$ | (2, 6/5, 1/5) | decreases with z |

[Mukaida et. al. 1208.3399, 1212.4985](#)

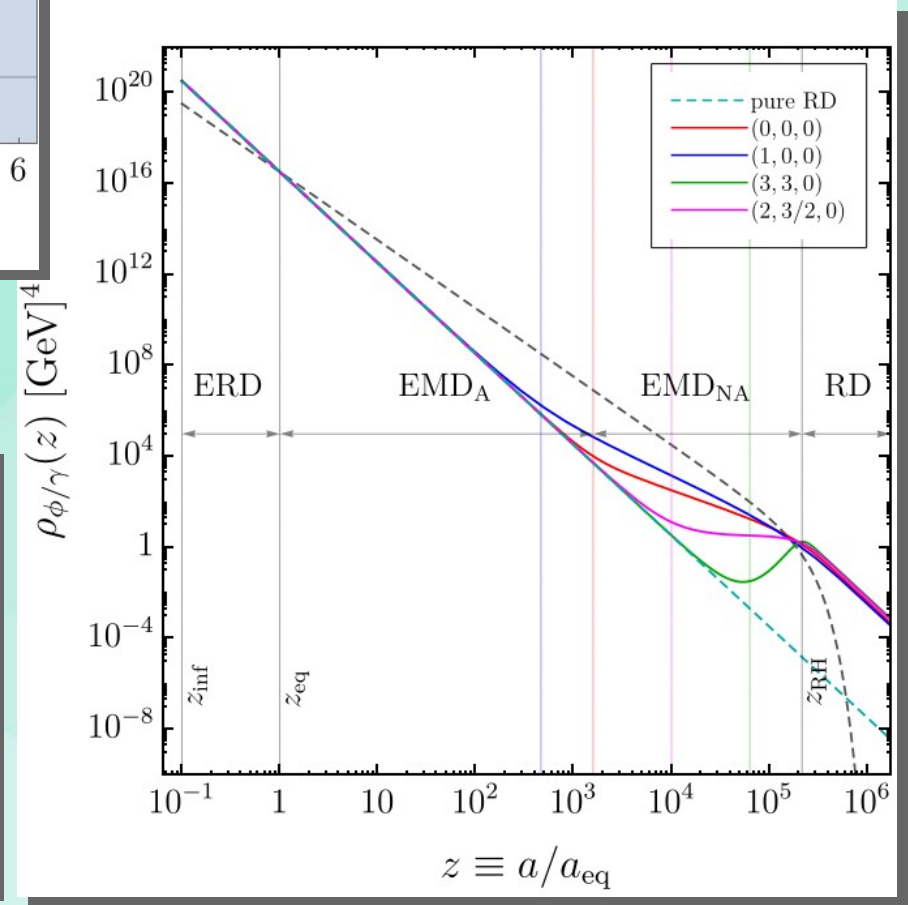
[Drewes, 1406.6243](#)

[Co et. al. 2007.04328](#)



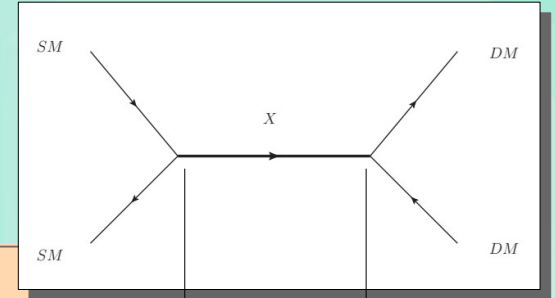
$$\Gamma_\phi = \hat{\Gamma} \left(\frac{T}{T_{\text{eq}}} \right)^n \left(\frac{a}{a_{\text{eq}}} \right)^k$$

| Epoch | z | $T(z)$ | $H(z)$ |
|-------------------|-------------------------------------|---|---|
| ERD | $z_{\text{inf}} < z < 1$ | $\frac{T_{\text{eq}}}{z}$ | $\sqrt{\frac{\rho_\gamma(T_{\text{eq}})}{3M_p^2}} z^{-2}$ |
| EMD _A | $1 < z < z_{\text{NA}}$ | $\frac{T_{\text{eq}}}{z}$ | $\sqrt{\frac{\rho_\gamma(T_{\text{eq}})}{3M_p^2}} z^{-\frac{3}{2}(1+\omega)}$ |
| EMD _{NA} | $z_{\text{NA}} < z < z_{\text{RH}}$ | $T_{\text{RH}} \left(\frac{z}{z_{\text{RH}}} \right)^{\frac{\delta-8+2n}{8-2n}}$ | $\sqrt{\frac{\rho_\gamma(T_{\text{eq}})}{3M_p^2}} z^{-\frac{3}{2}(1+\omega)}$ |
| RD | $z_{\text{RH}} < z$ | $T_{\text{eq}} z_{\text{RH}}^{\frac{1-3\omega}{4}} z^{-1}$ | $\frac{\sqrt{\rho_\gamma(T_{\text{RH}})}}{\sqrt{3}M_p} \left(\frac{z}{z_{\text{RH}}} \right)^{-2}$ |



Freeze-in production rate

$$\dot{n}_\chi + 3Hn_\chi = -\langle\sigma v\rangle (n_\chi^2 - n_\chi^{\text{eq}2}) = R(T)$$

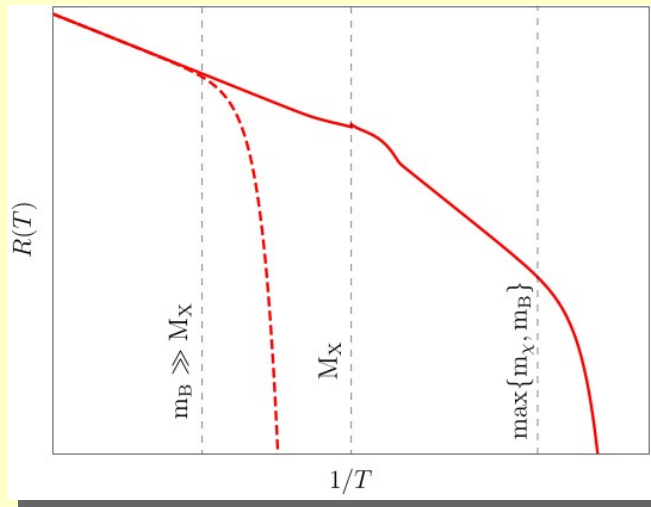


$$g_{X\text{-SM}} \ll 1 \quad g_{X\text{-DM}} \sim \mathcal{O}(1)$$

T_{inf}

$$R(T) \sim \frac{T^{p+4}}{M_X^4} e^{-\frac{2\max\{m_\chi, m_B\}}{T}}$$

M_X

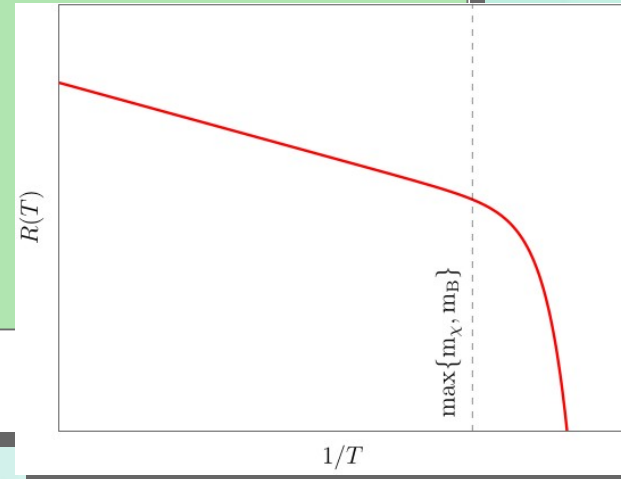


$$R(T) \sim \begin{cases} T^p e^{-\frac{2m_B}{T}}, & m_B \gg M_X \\ T^p, & m_B \ll M_X \ll T \\ \frac{\pi T M_X^p}{\Gamma_X} K_1\left(\frac{M_X}{T}\right), & m_B \ll M_X \sim T \\ \frac{T^{p+4}}{M_X^4} e^{-\frac{2m_B}{T}}, & m_B, T \ll M_X \end{cases}$$

$$M_X = 2m_\chi$$

$$R(T) \sim T^p e^{-\frac{2\max\{m_\chi, m_B\}}{T}}$$

m_χ



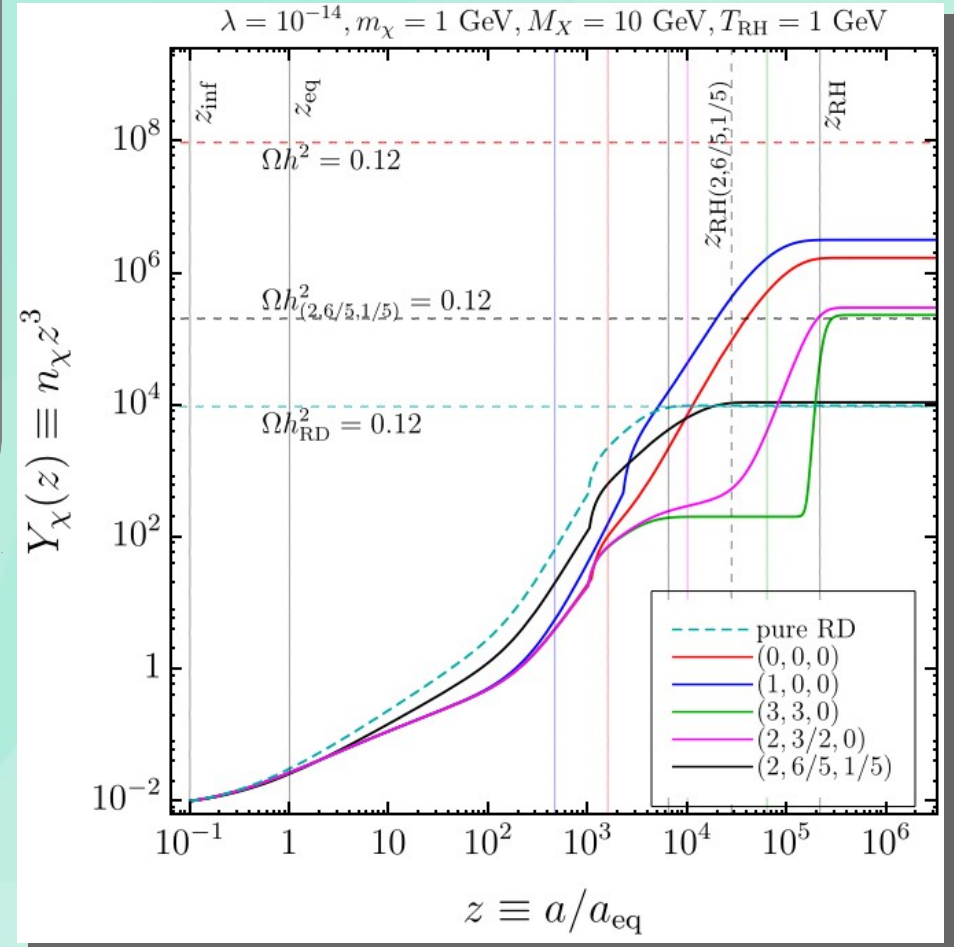
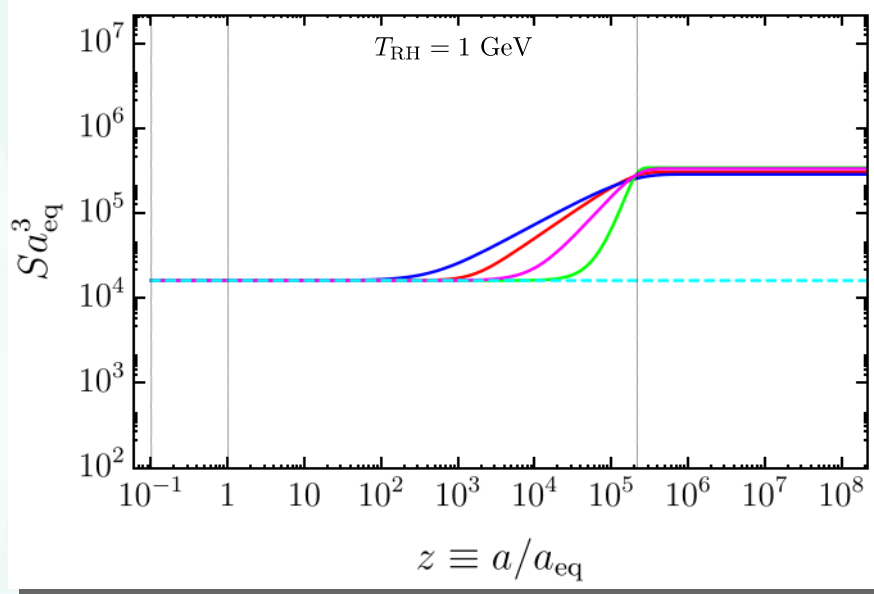
Freeze-in DM yield during EMD

- DM yield dilutes due to entropy production
- Non-trivial temperature evolution also changes the DM production rate during non-adiabatic EMD

$$\frac{\Omega h^2}{\Omega h_{RD}^2} = \frac{Y_\chi(z_0)}{Y_\chi^{RD}(z_0)} \left(\frac{z_0^{RD}}{z_0} \right)^3 = \frac{Y_\chi(z_0)}{Y_\chi^{RD}(z_0)} \left(\frac{T_{RH}}{T_{eq}} \right)^{\frac{1-3\epsilon}{1+\epsilon}}$$

$$\sim 10^{-2} - 10^{-3}$$

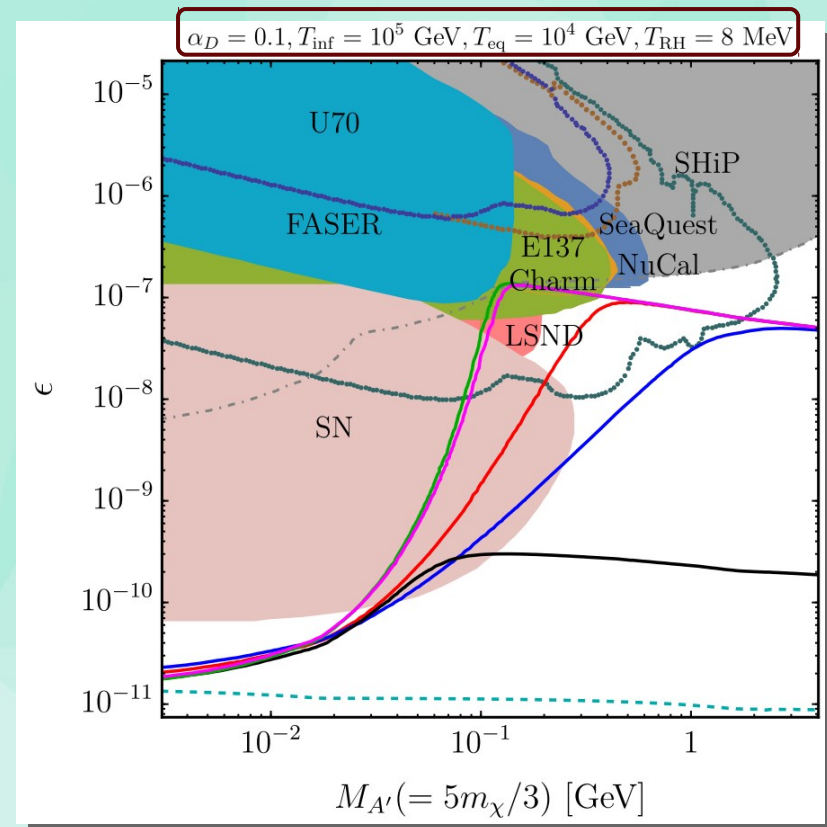
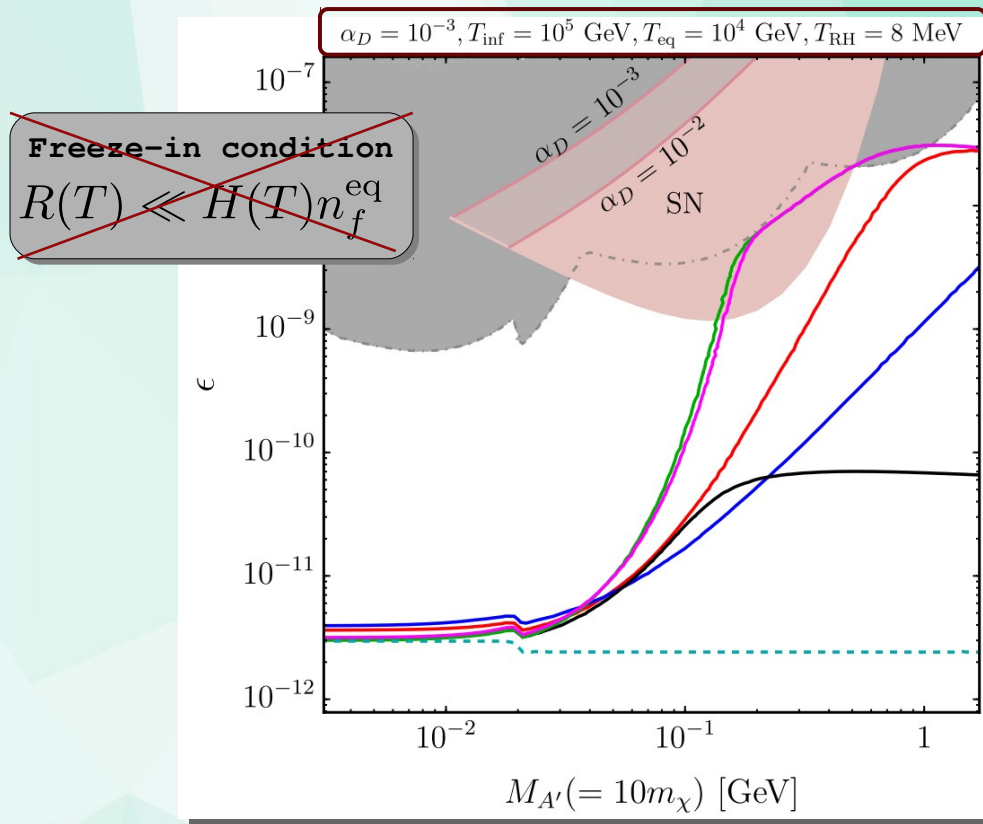
$$\frac{dY_\chi(z)}{dz} = \frac{z^2 R(T(z))}{H(z)}$$



Larger coupling is required to saturate freeze-in relic in presence of EMD epoch

Dark photon portal

$$\mathcal{L} = -\frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} - \frac{\epsilon}{2}F'_{\mu\nu}F^{\mu\nu} + \frac{1}{2}M_{A'}^2 A'_\mu A'^\mu + \bar{\chi} (i\not{\partial} - m_\chi) \chi + g_D \bar{\chi} \gamma^\mu A'_\mu \chi$$



Parameter space satisfying observed relic is accessible to experiments in the presence of an early matter dominated era

Summary

- Freeze-in DM relic depends on the **non-standard epochs of cosmology** at high temperatures
- An epoch of **pre-BBN early matter domination** leads to **freeze-in with larger couplings**
- Details depend crucially on the **temperature and expansion dependent dissipation rate** of the dominating matter field
- **Dark photon portal dark matter** model may come under the experimental radar in presence of early matter domination

THANK YOU

