

Matching resummed endpoint and continuum γ -ray spectra from dark-matter annihilation

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based on

[arXiv:2203.01692](https://arxiv.org/abs/2203.01692)

with M. Beneke and M. Vollmann

also see earlier works

[1903.08702/1912.02034](https://arxiv.org/abs/1903.08702) – Sudakov logs

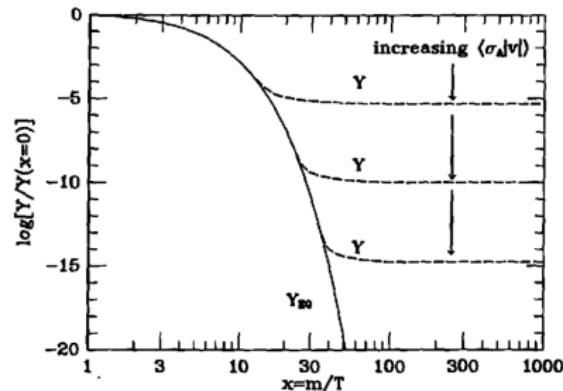
[1909.04584/2108.07285](https://arxiv.org/abs/1909.04584) – NLO potentials



Motivation

TeV-scale electroweak WIMPs

- ▶ “WIMP” miracle: TeV-scale particle with weak scale cross section
 $\xrightarrow{\text{freeze-out}}$ observed DM relic abundance
$$\Omega h^2 \sim 0.1 \frac{3 \times 10^{26} \text{cm}^3 \text{s}^{-1}}{\langle \sigma v \rangle} \sim 0.1 \frac{\pi \alpha_2^2 / m_\chi^2}{\langle \sigma v \rangle}$$
- ▶ two main minimal DM scenarios:
 - ▶ SU(2)_L-triplet (wino)
thermal mass $m_\chi \approx (2.7 - 3)$ TeV
 - ▶ SU(2)_L-doublet w. hypercharge (higgsino)
thermal mass $m_\chi \approx 1$ TeV
- ▶ Results and framework also applicable to broader class of models (e.g. MSSM, ...)
- ▶ Direct detection: around neutrino floor
- ▶ LHC: too heavy to be seen
- ▶ Indirect detection: higgsino/wino testable



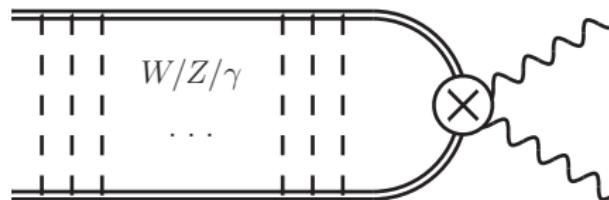
[Kolb, Turner '90]

Motivation

Sommerfeld effect

Non-relativistic DM particles exchange EW gauge bosons
(freeze-out $v \sim 0.1$ / indirect detection $v \sim 10^{-3}$)

[Hisano et al. '04/'06]



each ladder rung suppressed by (potential region)

$$\alpha_2 \frac{m_\chi}{m_W} \quad \text{or} \quad \frac{\alpha_2}{v} \quad \sim \mathcal{O}(1)$$

results in static force in form of Yukawa/Coulomb potentials

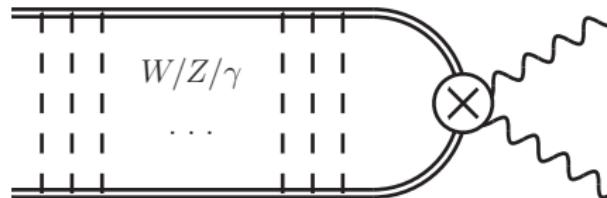
$$V(r) \sim -\alpha_2 \frac{e^{-m_{W/Z} r}}{r} \quad \text{or} \quad -\frac{\alpha}{r}$$

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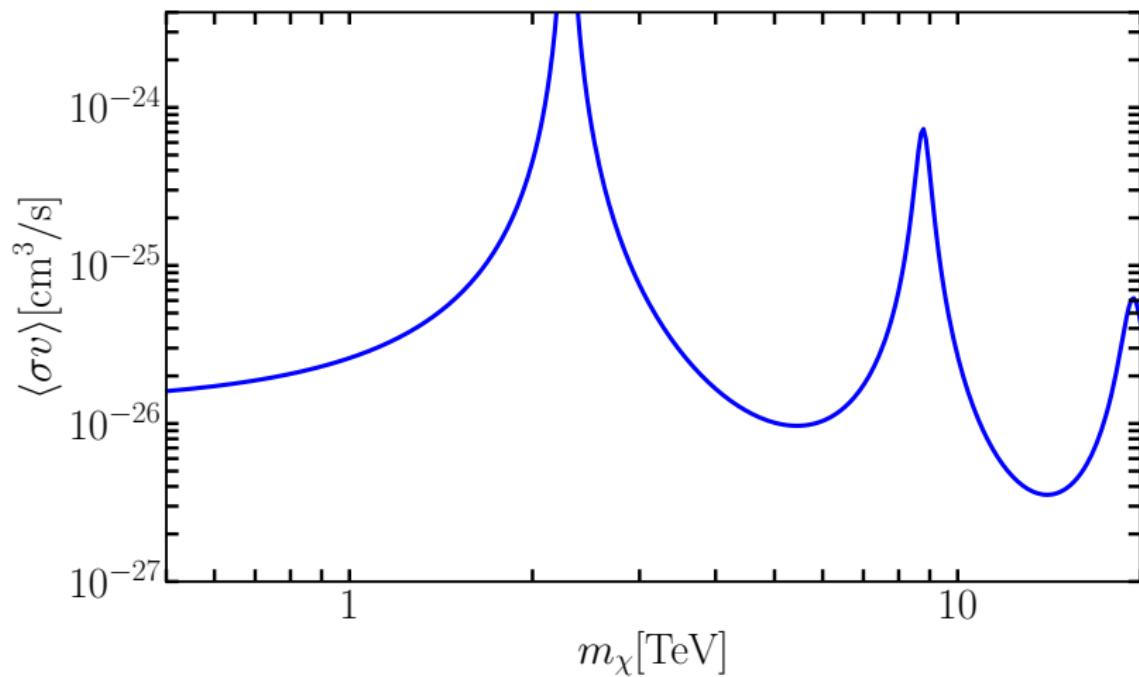
Example: Tree-level wino potential - neutral sector 1S_0

$$V_{\text{LO}}(r) = \begin{pmatrix} 0 & -\sqrt{2} \alpha_2 \frac{e^{-m_W r}}{r} \\ -\sqrt{2} \alpha_2 \frac{e^{-m_W r}}{r} & 2\delta m - \frac{\alpha}{r} - \alpha_2 c_W^2 \frac{e^{-m_Z r}}{r} \end{pmatrix} \quad \begin{matrix} (00) \\ (+-) \end{matrix}$$

Solve Schrödinger equation \Rightarrow cross-section enhancement

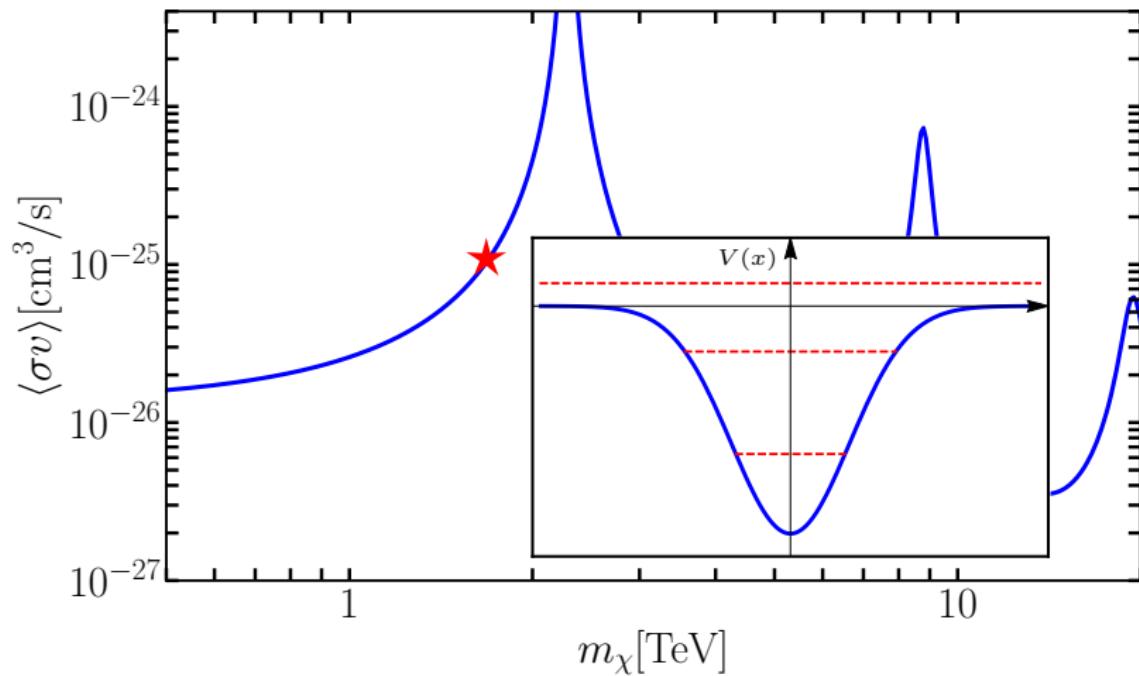
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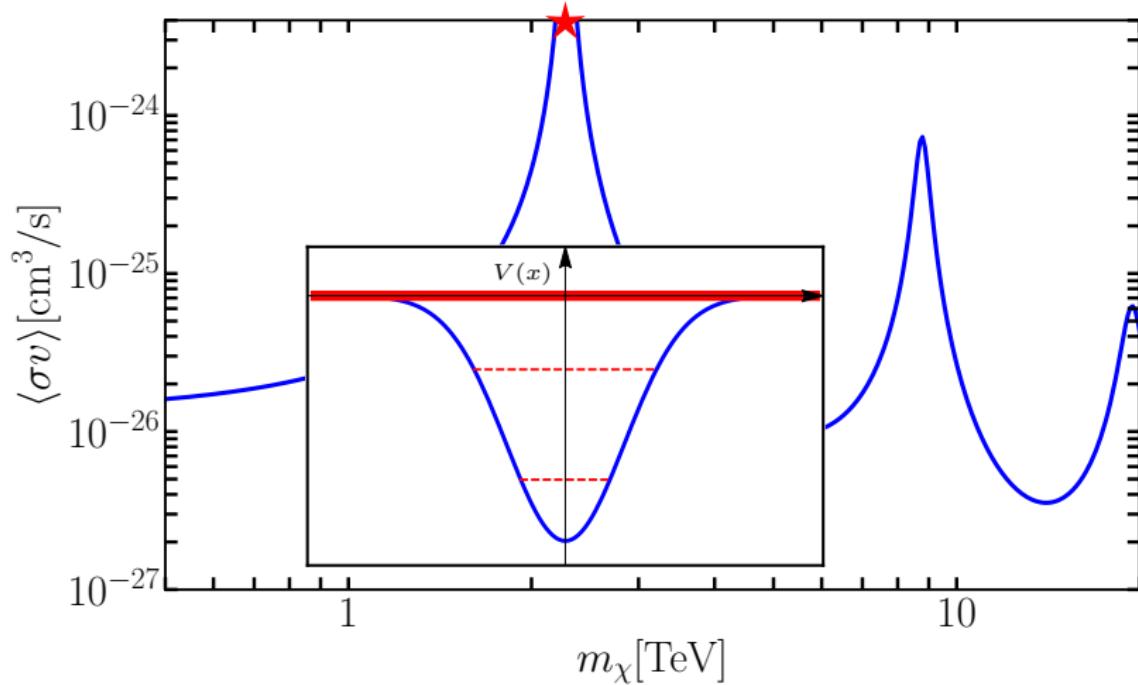
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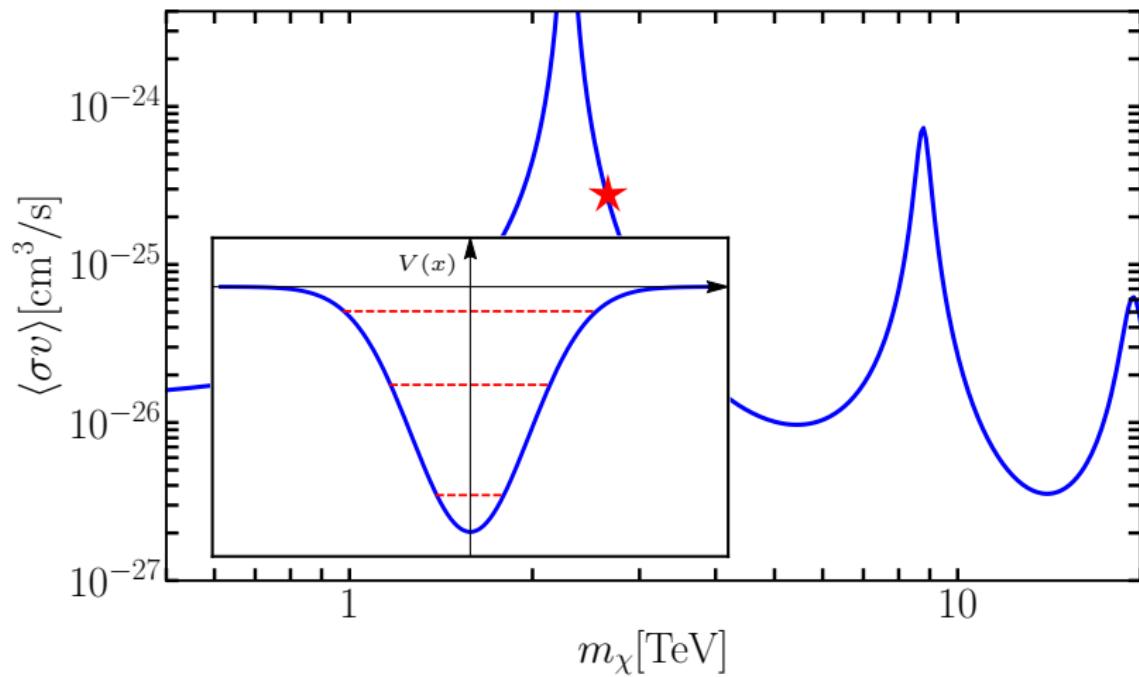
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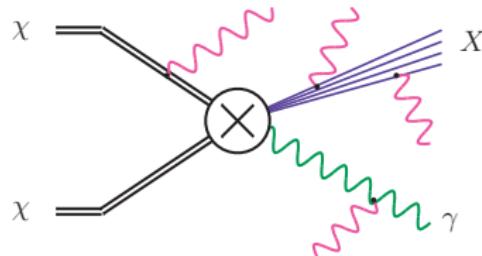
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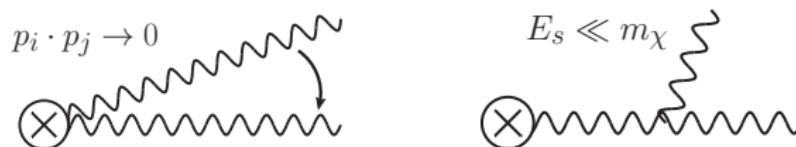


Motivation

Sudakov logarithms



Kinematics force soft/collinear limits for the soft radiation

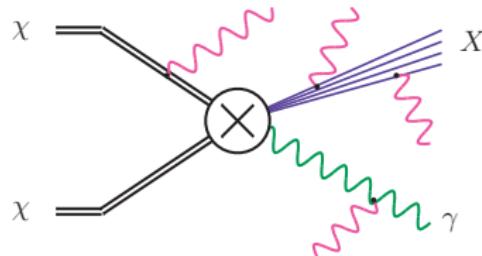


IR divergencies cause large Sudakov double logs at all orders in perturbation theory

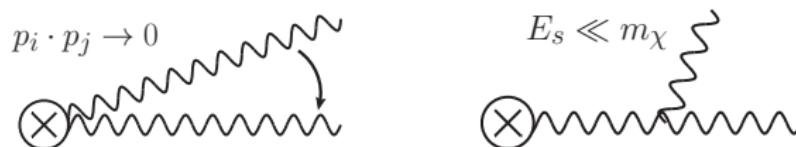
$$\frac{\alpha_2}{\pi} \log^2 \frac{4m_\chi^2}{m_W^2} \stackrel{m_\chi = 3 \text{ TeV}}{\approx} 0.83$$

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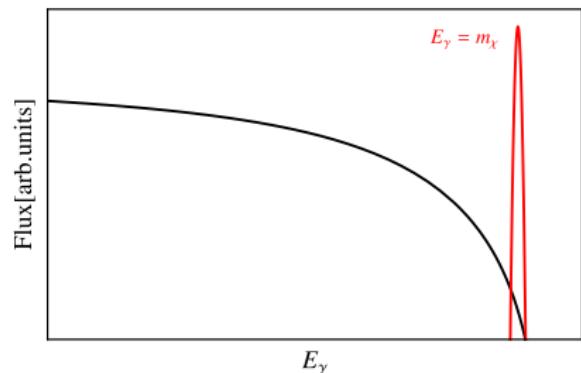
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⇒ Employ full QFT/EFT machinery (NRDM \otimes SCET)
(full NLO, factorization, resummation, ...)

Motivation

Line signal

- ▶ Heavy WIMP non-relativistic today ($v \sim 10^{-3}$)
⇒ annihilation at threshold / photon energies $E_\gamma \approx m_\chi$
- ▶ large corrections in exclusive case [Bauer et al. '15, Ovanesyan et al. '15/'17] and semi-inclusive case [Baumgart et al. '15/'16]



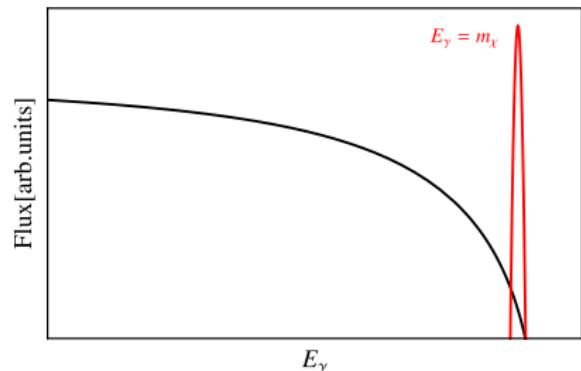
Line signal count rate in experiment

$$N_{jk}^S = T \frac{1}{8\pi m_\chi^2} J_k \int_{\Delta E_k} dE \int dE' A_{\text{eff}}(E') G(E', E) \frac{d\langle\sigma v\rangle}{dE}$$

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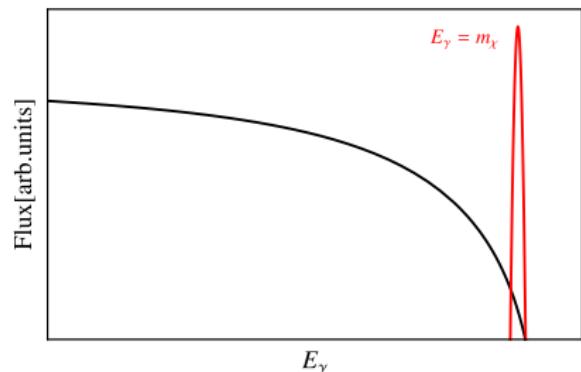
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differential cross section
(particle physics only)

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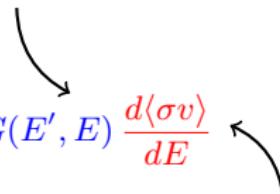
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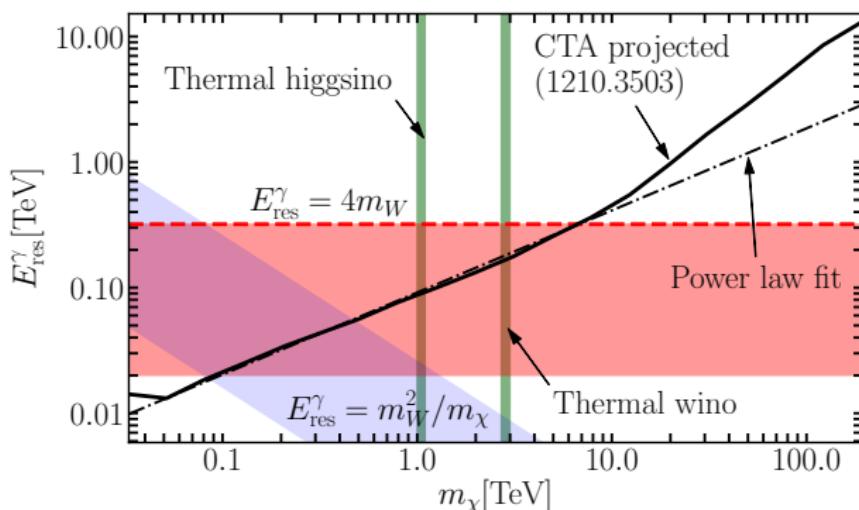
Detector resolution function



differential cross section
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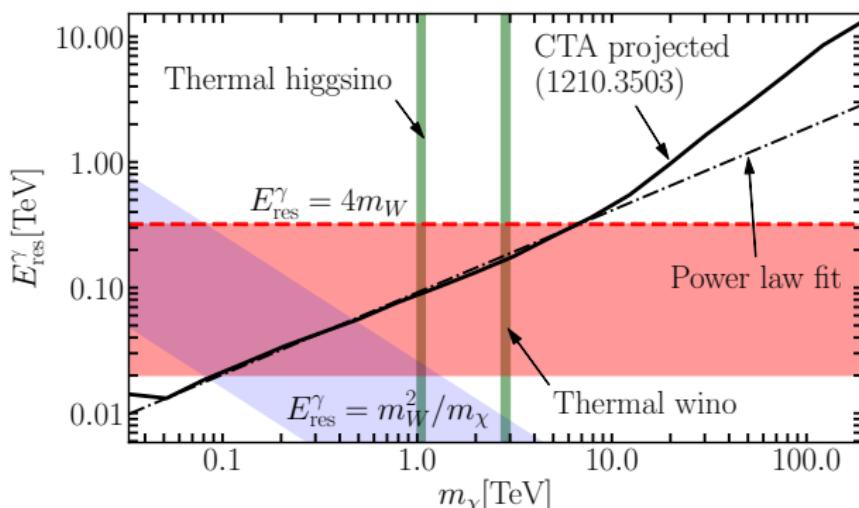
Finite energy resolution - $\chi^0 \chi^0 \rightarrow \gamma + X$

- Finite energy resolution of detector E_{res}^γ (leads to different EFT setups)
 - line signal - $E_{\text{res}}^\gamma \sim 0$
[Bauer et al. '14, Ovanesyan et al. '14/'16, Baumgart et al. '14]
 - narrow resolution - $E_{\text{res}}^\gamma \sim m_W^2/m_\chi$ [Beneke et al. '18]
 - intermediate resolution - $E_{\text{res}}^\gamma \sim m_W$
[Beneke, Hasner, KU, Vollmann, 1912.02034 + Broggio, 1903.08702]
 - wide resolution - $E_{\text{res}}^\gamma \gg m_W$ [Baumgart et al. '17/'18]



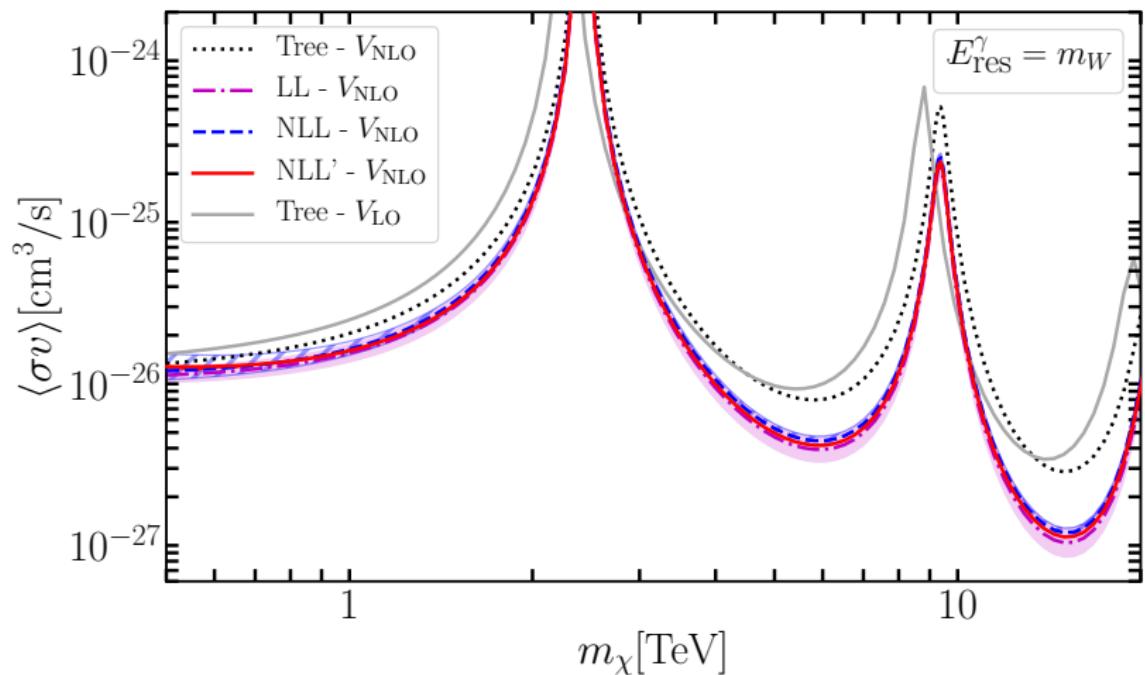
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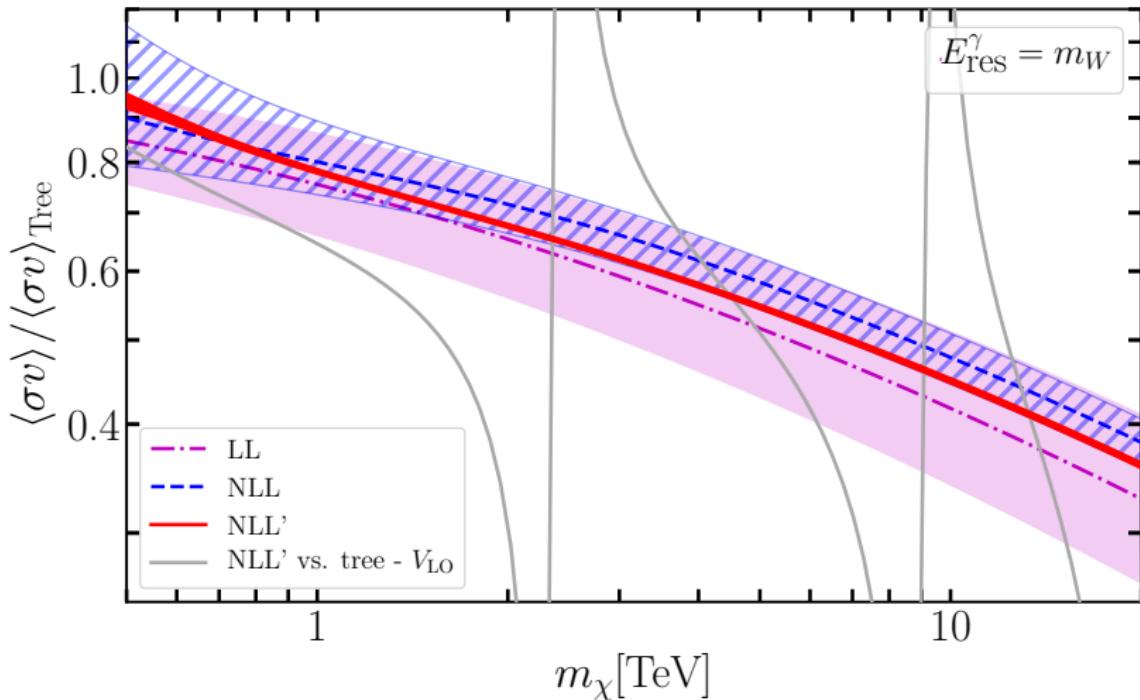
Wino indirect detection – $\chi^0\chi^0 \rightarrow \gamma + X$

[1903.08702/1912.02034]



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Extending energy spectrum beyond the endpoint

- ▶ For real experimental analysis, and thorough estimate of impact on theoretical predictions need differential spectrum also beyond kinematic endpoint.
- ▶ Two small parameters in the problem:

$$\epsilon = \frac{m_W}{2m_\chi}, \quad 1 - x = 1 - \frac{E_\gamma}{m_\chi}$$

- ▶ Away from endpoint, i.e., $1 - x \sim \mathcal{O}(1)$, regime of parton showers (in this talk PPPC4 DM ID [Cirelli et al. '10])
- ▶ Logarithmic discussion independent of specific parton shower used.

Compare splitting functions:

$$\frac{dN_{WW}^{IJ}}{dx} \equiv m_\chi \left. \frac{\Gamma_{IJ}^{\text{res}}}{\hat{\Gamma}_{IJ,\text{tree}}^{\text{WW}}} \right|_{\mathcal{O}(\hat{\alpha}_2)} + \mathcal{O}(\hat{\alpha}_2^2) \quad (x < 1)$$

Splittings wino $\chi^0 \chi^0 \rightarrow W^+ W^- \gamma$

intermediate resolution ($1 - x \sim \epsilon \ll \mathcal{O}(1)$)

$$\frac{dN_{WW}}{dx} \Big|_{\chi^0 \chi^0 \rightarrow WW\gamma}^{\text{int.}} = \frac{2\alpha_{\text{em}}}{\pi} \left[\frac{1}{1-x} \ln \left(1 + \frac{(1-x)^2}{\epsilon^2} \right) - \frac{1-x}{\epsilon^2 + (1-x)^2} \right]$$

collinear approximation – PPPC4DM ($\epsilon \ll 1 - x \sim \mathcal{O}(1)$)

$$\frac{dN_{WW}}{dx} \Big|_{\text{PPPC4DM}} = \frac{2\alpha_{\text{em}}}{\pi} \left[\frac{x}{1-x} \ln \frac{(1-x)^2}{\epsilon^2} - \left(\frac{1-x}{x} + x(1-x) \right) \ln \epsilon^2 \right]$$

one-loop fixed order expanded for small ϵ , and $\delta m_\chi = m_{\chi^+} - m_{\chi^0}$
 $(1 - x \gg \epsilon \text{ and } x \gg \epsilon)$

$$\begin{aligned} \frac{dN_{WW}}{dx} \Big|_{\chi^0 \chi^0 \rightarrow WW\gamma}^{\text{f.o.}} &= \frac{2\alpha_{\text{em}}}{\pi} \left[\frac{(1-x+x^2)^2}{(1-x)x} \ln \frac{1}{\epsilon^2} \right. \\ &\quad - \frac{(4 - 12x + 19x^2 - 22x^3 + 20x^4 - 10x^5 + 2x^6)}{(2-x)^2(1-x)x} \\ &\quad \left. + \frac{8 - 24x + 42x^2 - 37x^3 + 16x^4 - 3x^5}{(2-x)^3(1-x)x} \ln(1-x) \right] \end{aligned}$$

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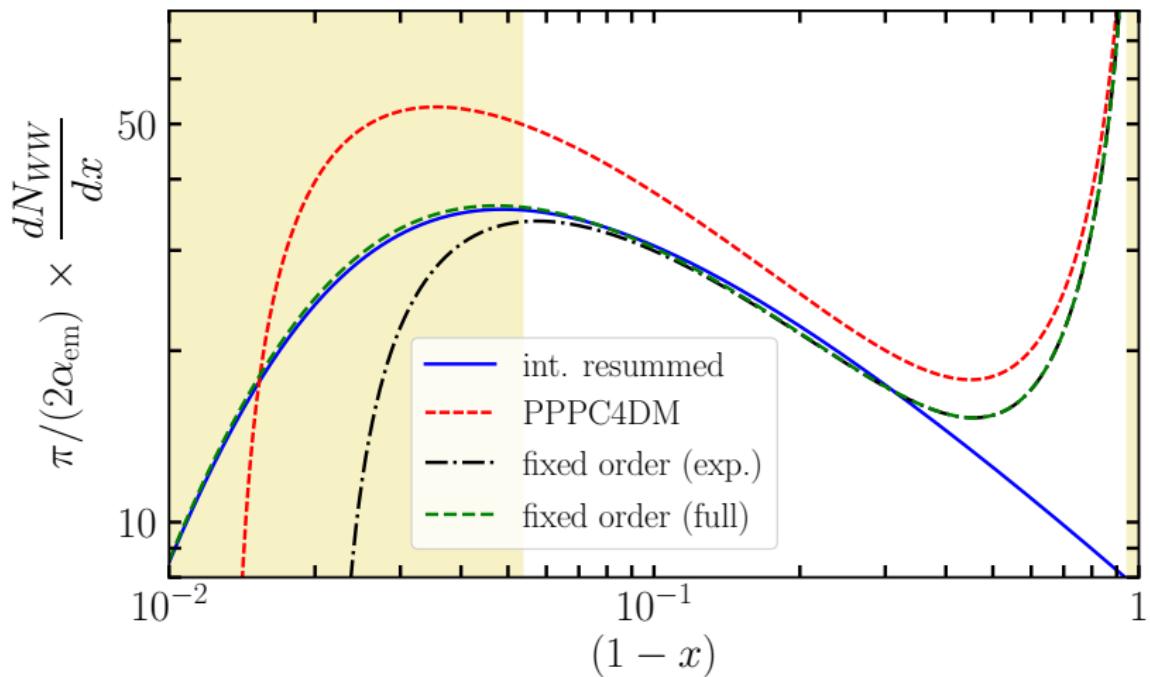
Leading $1/\epsilon$ term

$$\frac{dN_{WW}}{dx} = \frac{2\alpha_{\text{em}}}{\pi} \frac{1}{\epsilon} \left\{ \begin{array}{ll} \frac{\ln(1+\beta^2)}{\beta} - \frac{\beta}{1+\beta^2} & \text{int. res.} \\ \frac{\ln \beta^2}{\beta} & \text{PPPC4DM} \\ \frac{\ln \beta^2 - 1}{\beta} & \text{full fixed order} \end{array} \right\} + \mathcal{O}(\epsilon^0).$$

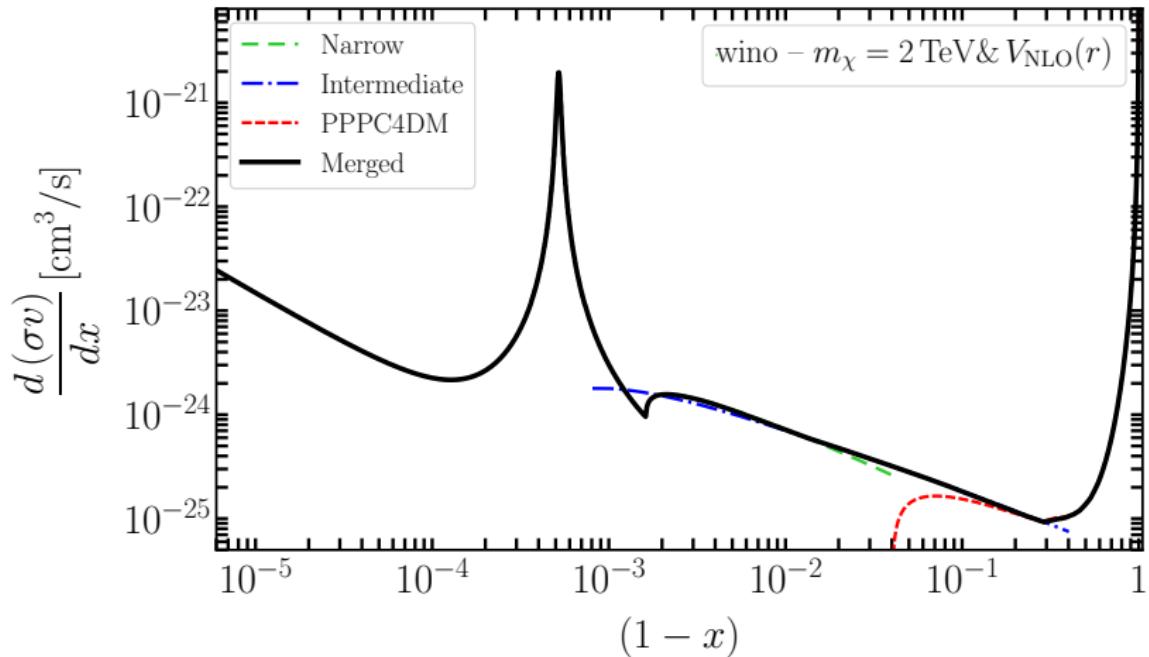
with $1-x = \beta\epsilon$.

Splittings wino $\chi^0\chi^0 \rightarrow W^+W^-\gamma$

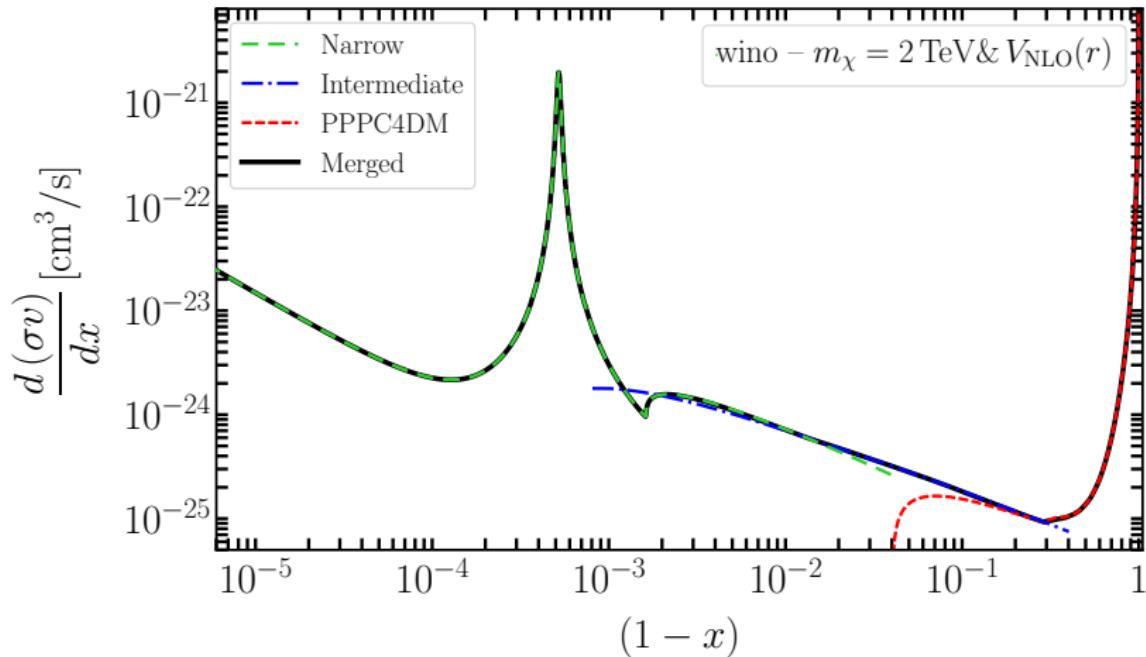
wino - $m_\chi = 3 \text{ TeV}$ - $\chi^0\chi^0 \rightarrow W^+W^-\gamma$



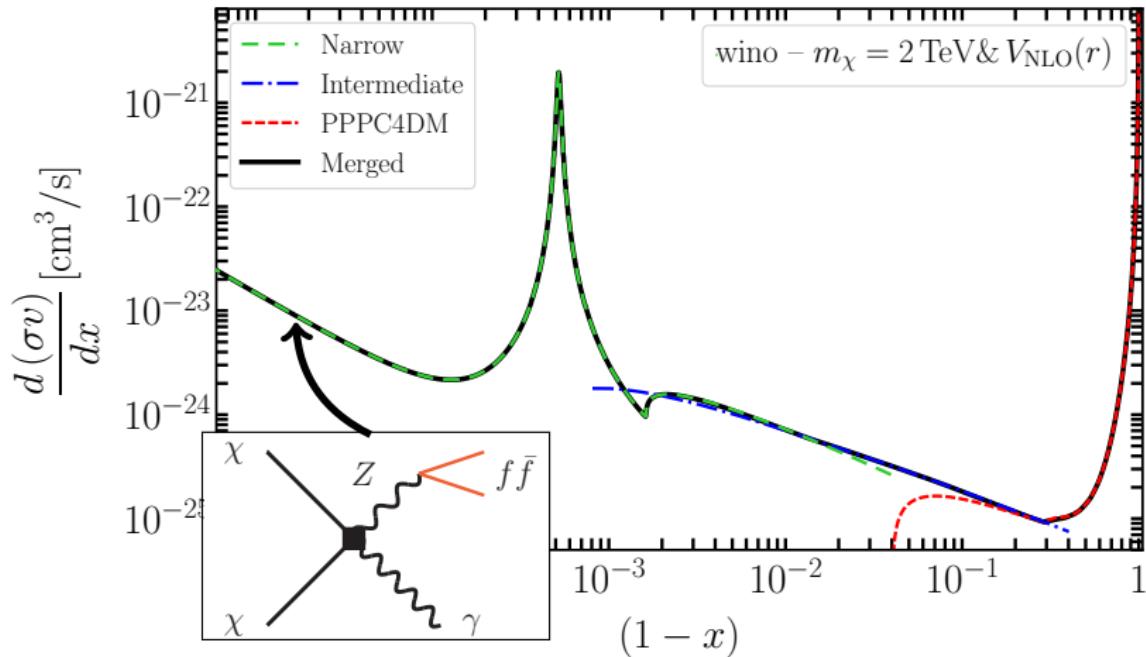
Wino differential spectrum – $\chi^0\chi^0 \rightarrow \gamma + X$



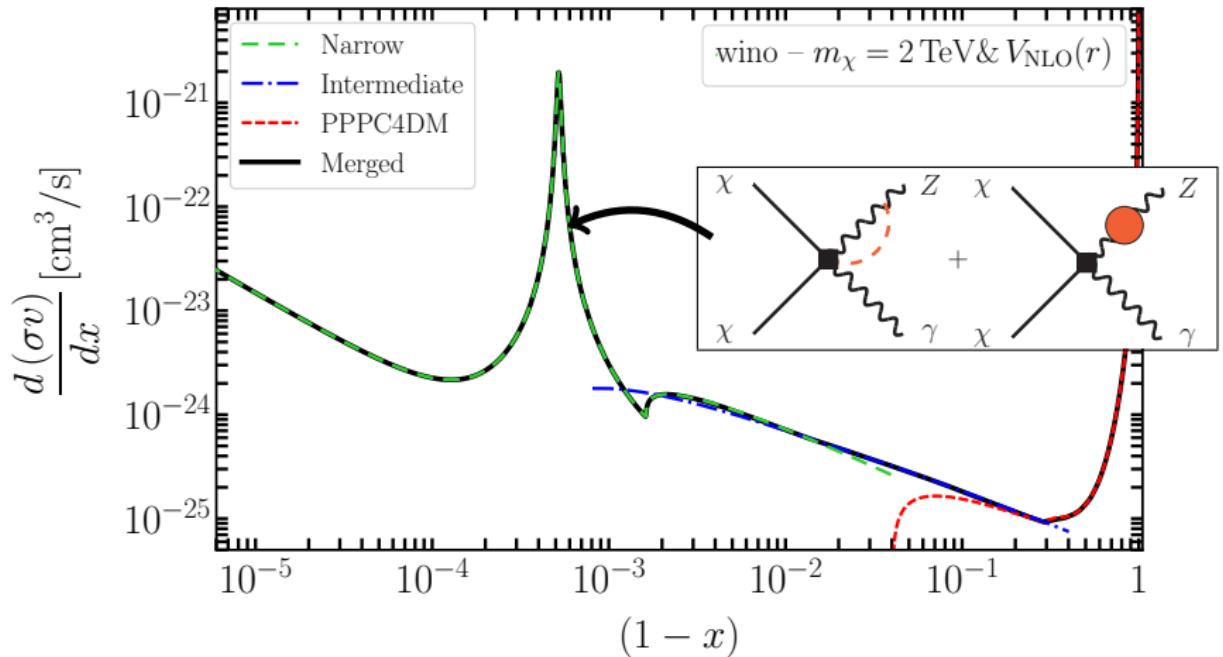
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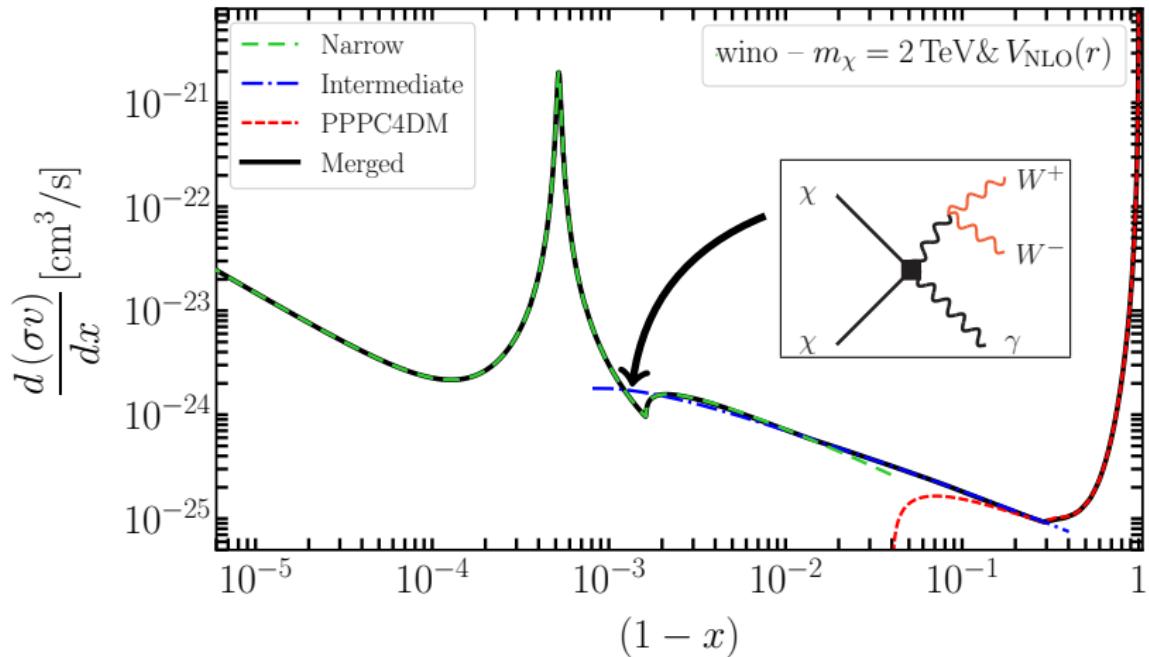
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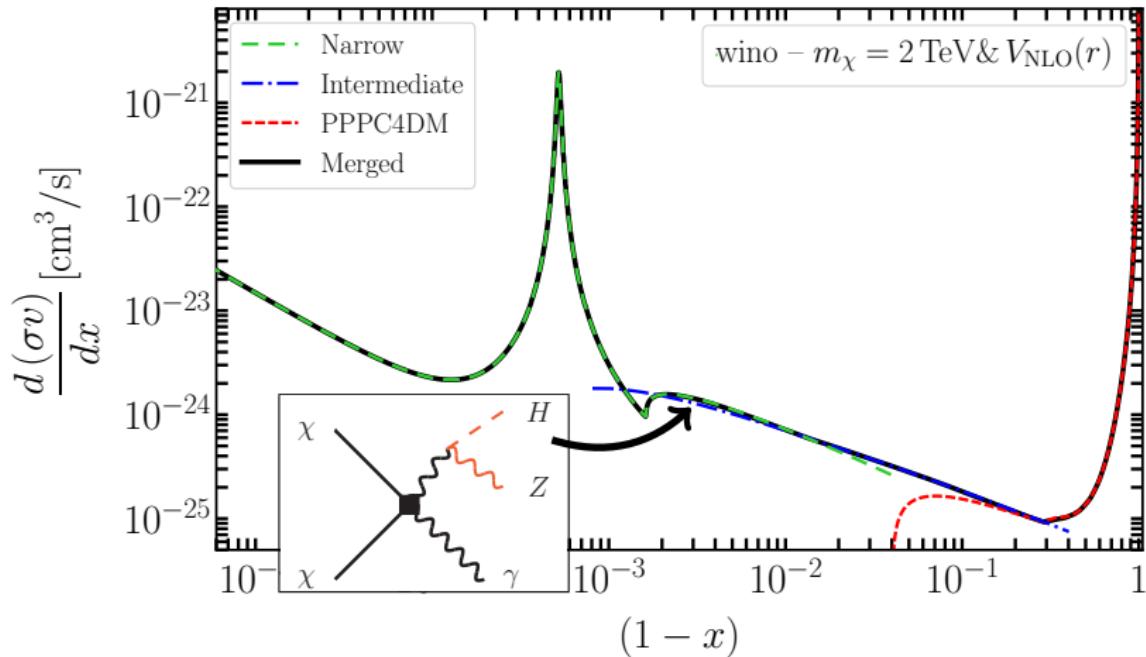
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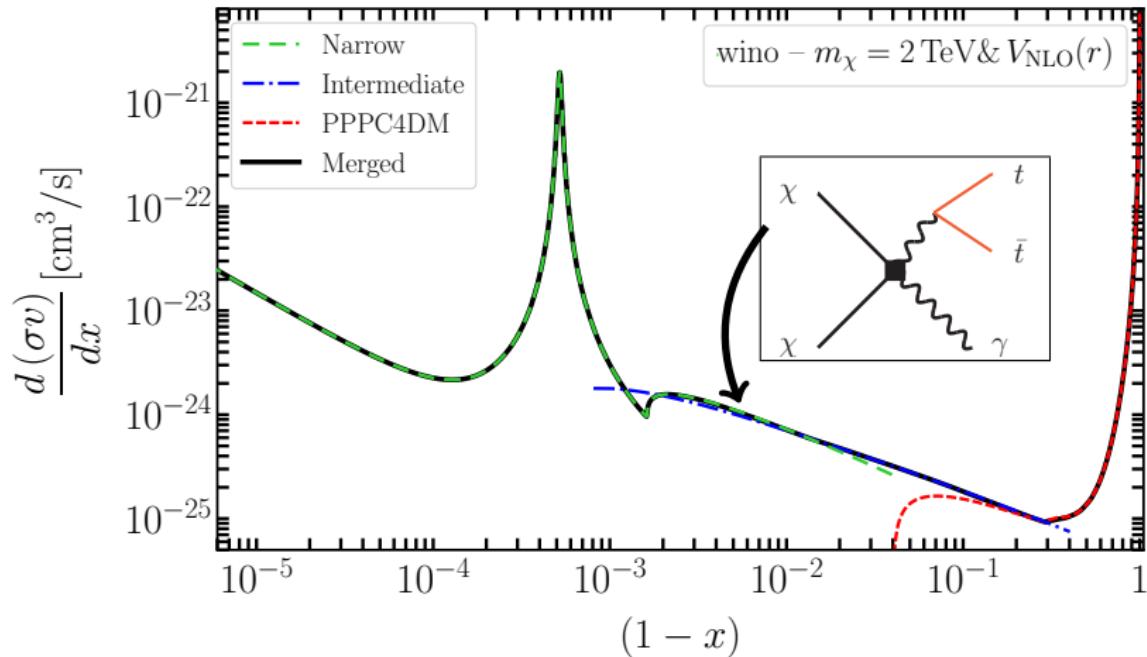
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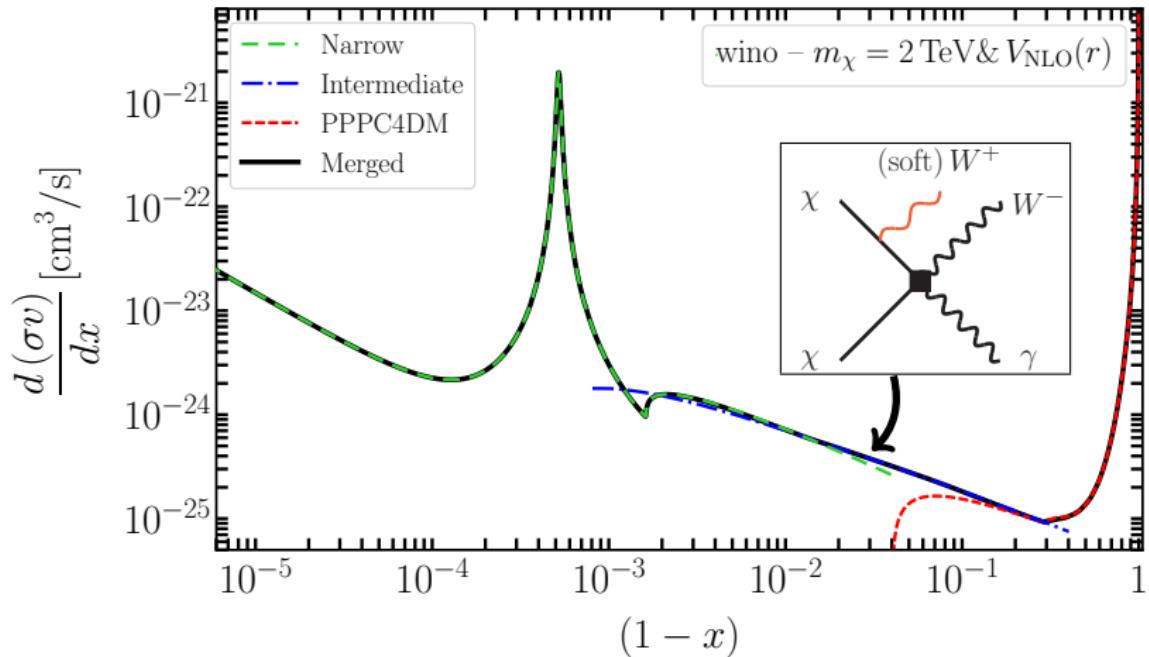
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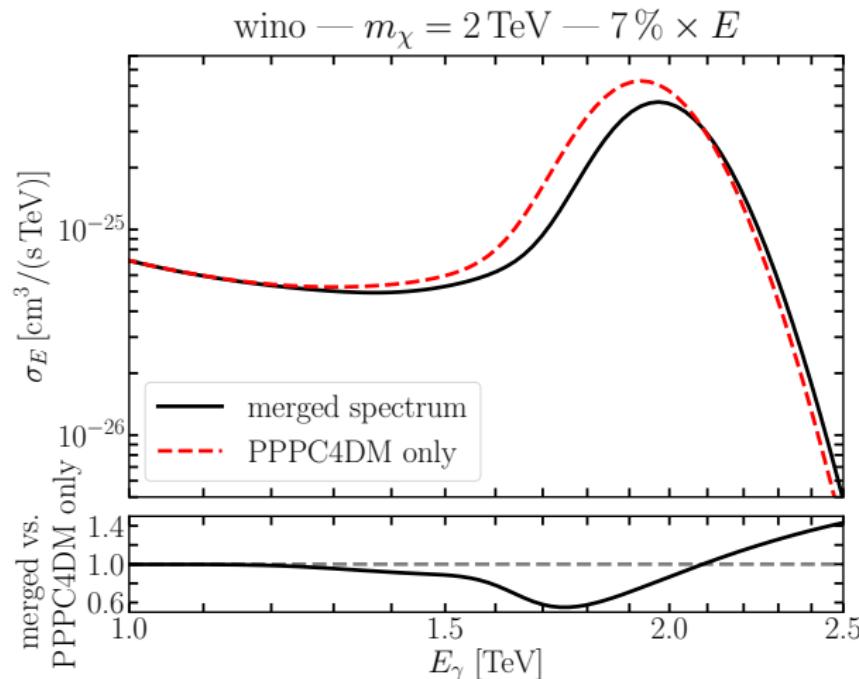
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Estimated effect in experiment

Folding with Gaussian resolution

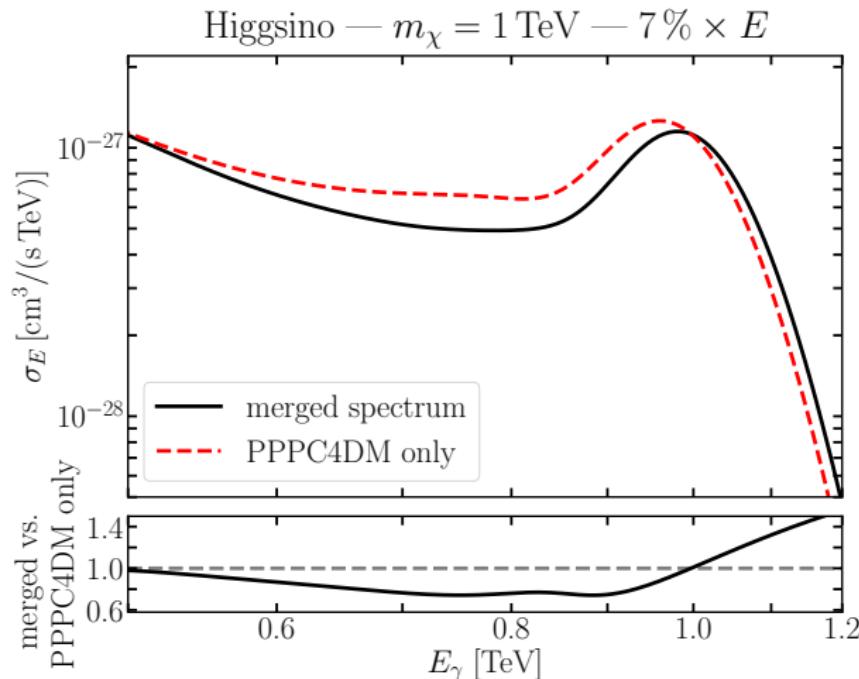
$$7\% \times E \cong \int_0^{m_\chi} dE'_\gamma \frac{d(\sigma v)}{dE'_\gamma} \frac{1}{\sqrt{2\pi} \cdot (7\%) \cdot E'_\gamma} e^{-\frac{(E_\gamma - E'_\gamma)^2}{2 \cdot (7\%)^2 \cdot E_\gamma^2}}$$



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Summary and Outlook

Main results:

- ▶ Achieved merging between endpoint resummed (NLL') and parton shower calculation
- ▶ Effect several 10's % suppression around nominal endpoint w.r.t. naive line prediction.
- ▶ Accurate predictions over entire possible range of photon energies including NLO Sommerfeld results for wino and Higgsino DM
- ▶ Code available `DMγspec`, check out dmyspec.hepforge.org

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- ▶ further models with richer structure (e.g., relevant bound states)
- ▶ CTA mock analysis
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