

# Capture of primordial black holes by neutron stars

Yoann Genolini

A work in collaboration with :  
Pasquale Serpico & Peter Tinyakov

Based on : Phys. Rev. D 102, 083004 (2020)

The logo for LAPTh (Laboratoire d'Annecy-le-Vieux de Physique Théorique) is displayed in white text on a dark teal square background. The letters 'L', 'A', and 'P' are in a standard sans-serif font, while 'T' and 'h' are in a stylized, cursive-like font. A small orange star is positioned above the 'h'.

TUG, December 2021

LAPTh/USMB

# Outline

Overview and motivations

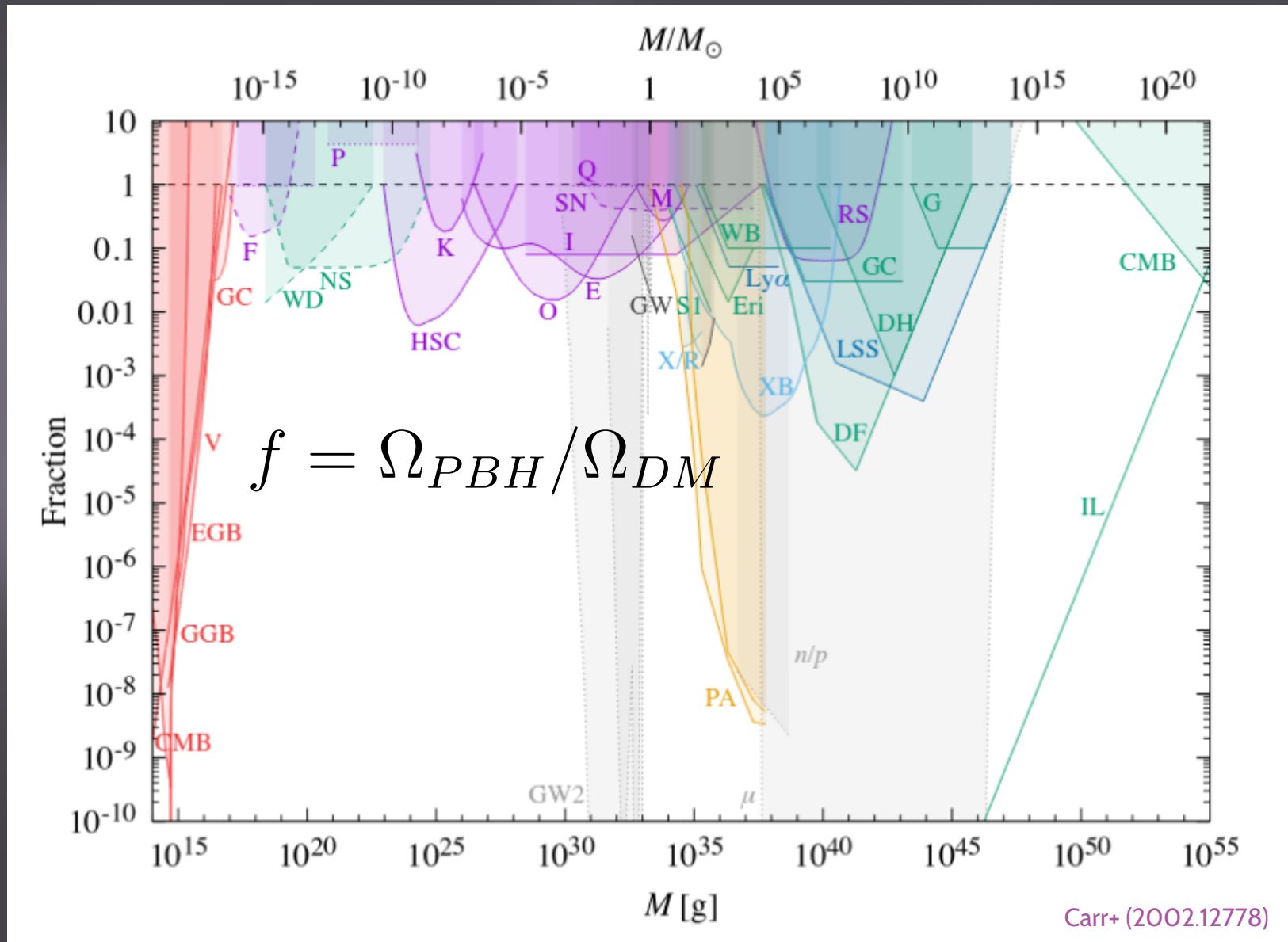
PBH interactions with a NS

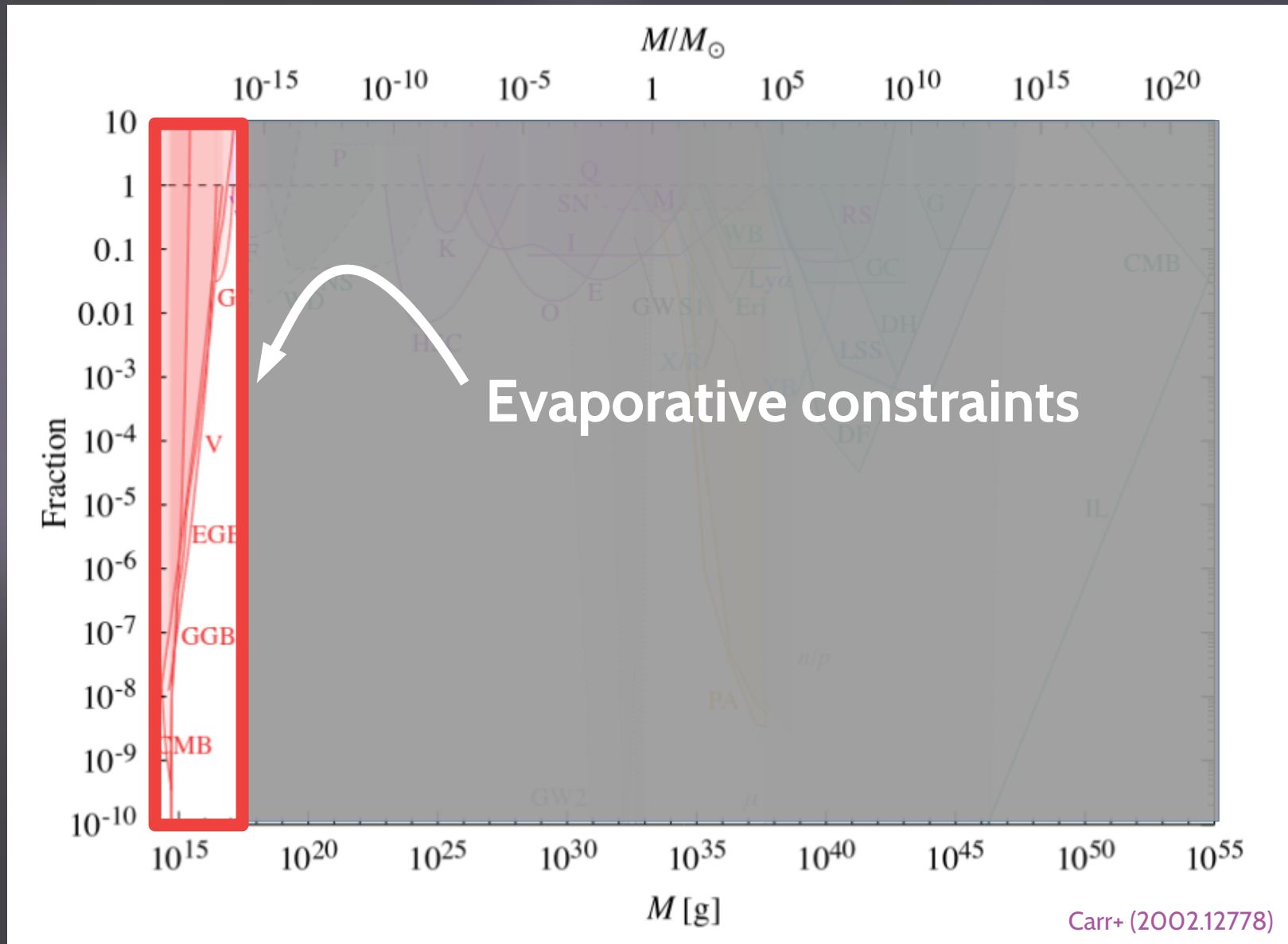
Capture of a PBH

Post capture dynamic

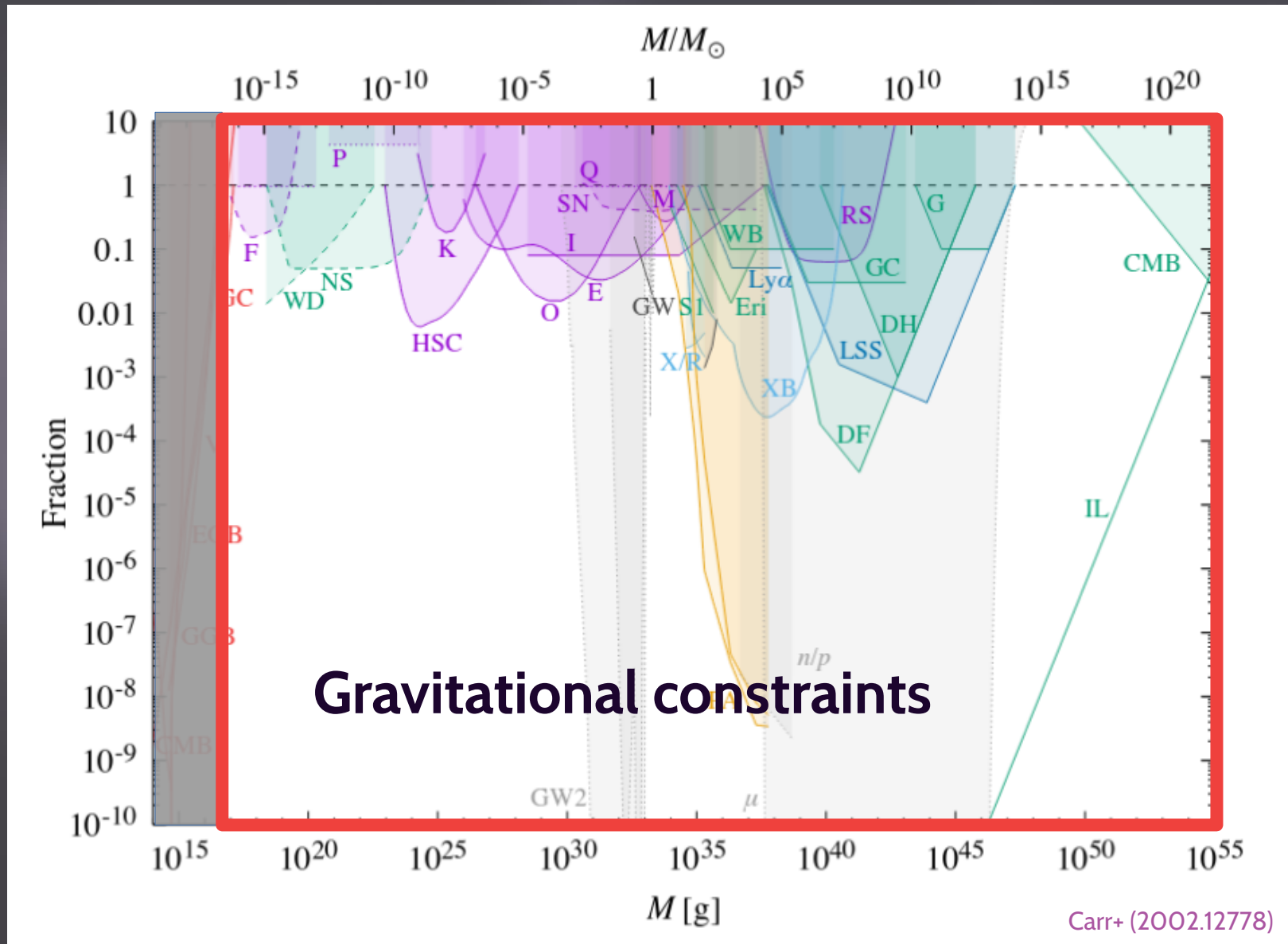
Signatures

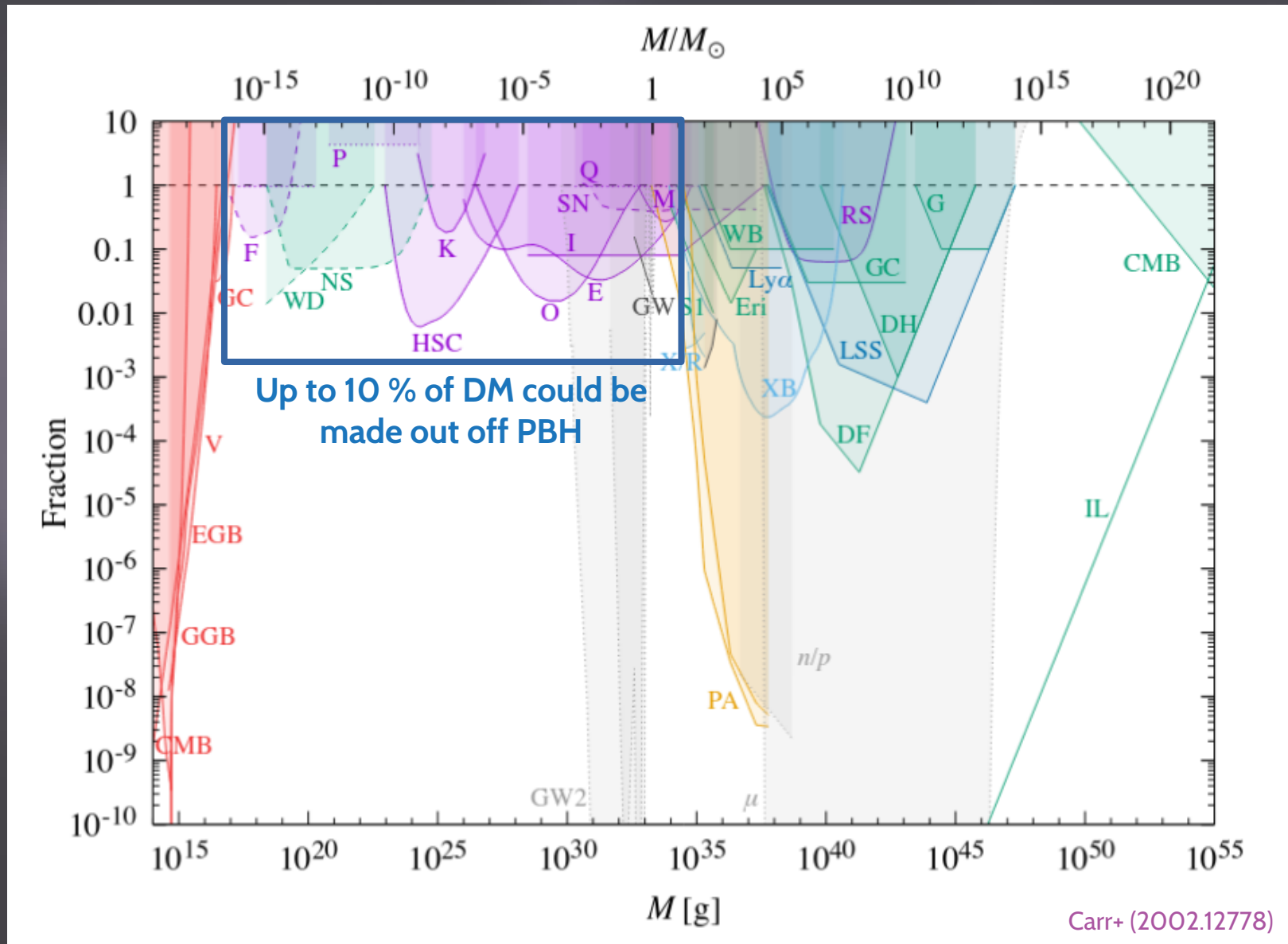
# PBH interactions with a NS - The scenario

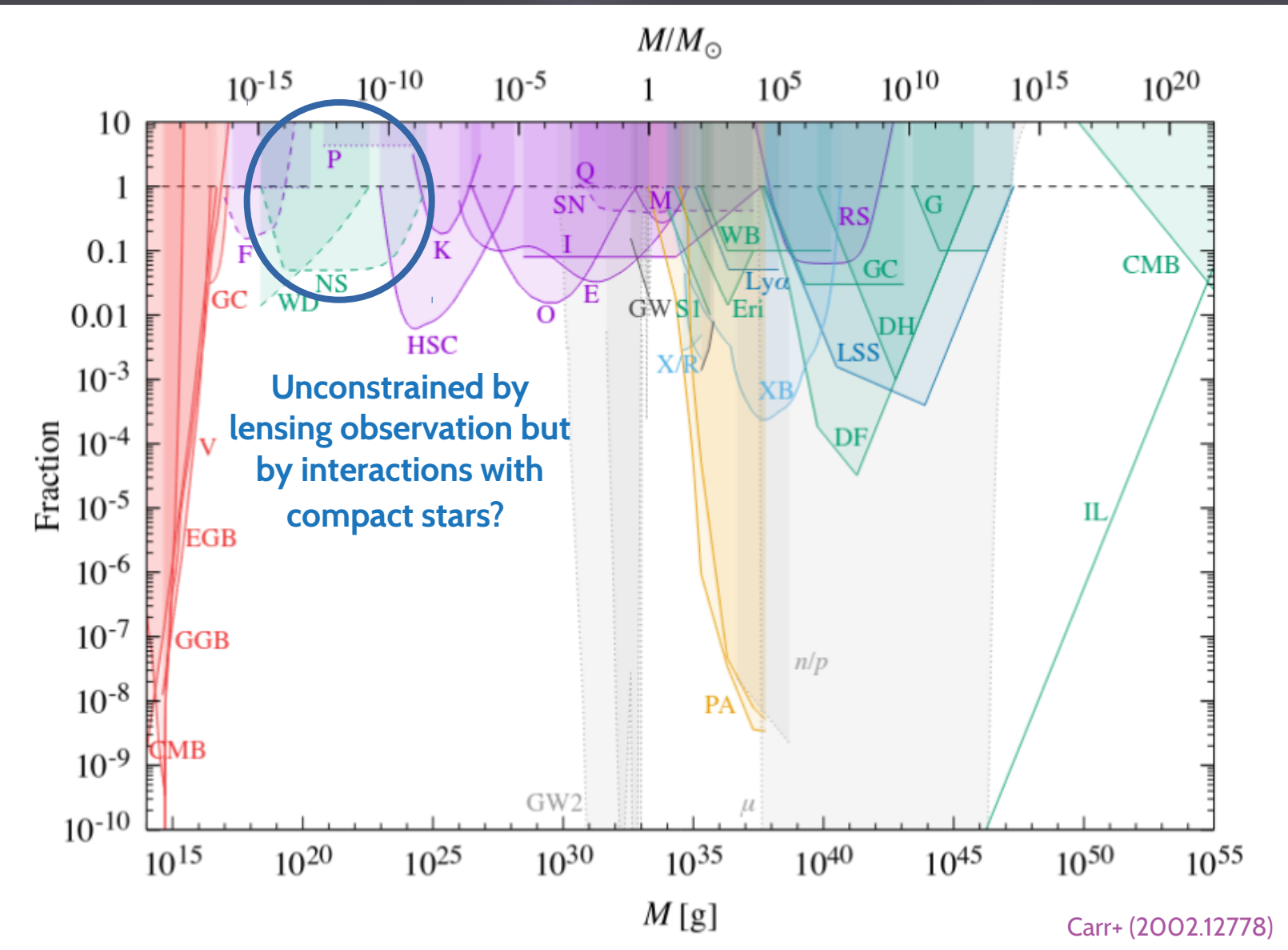




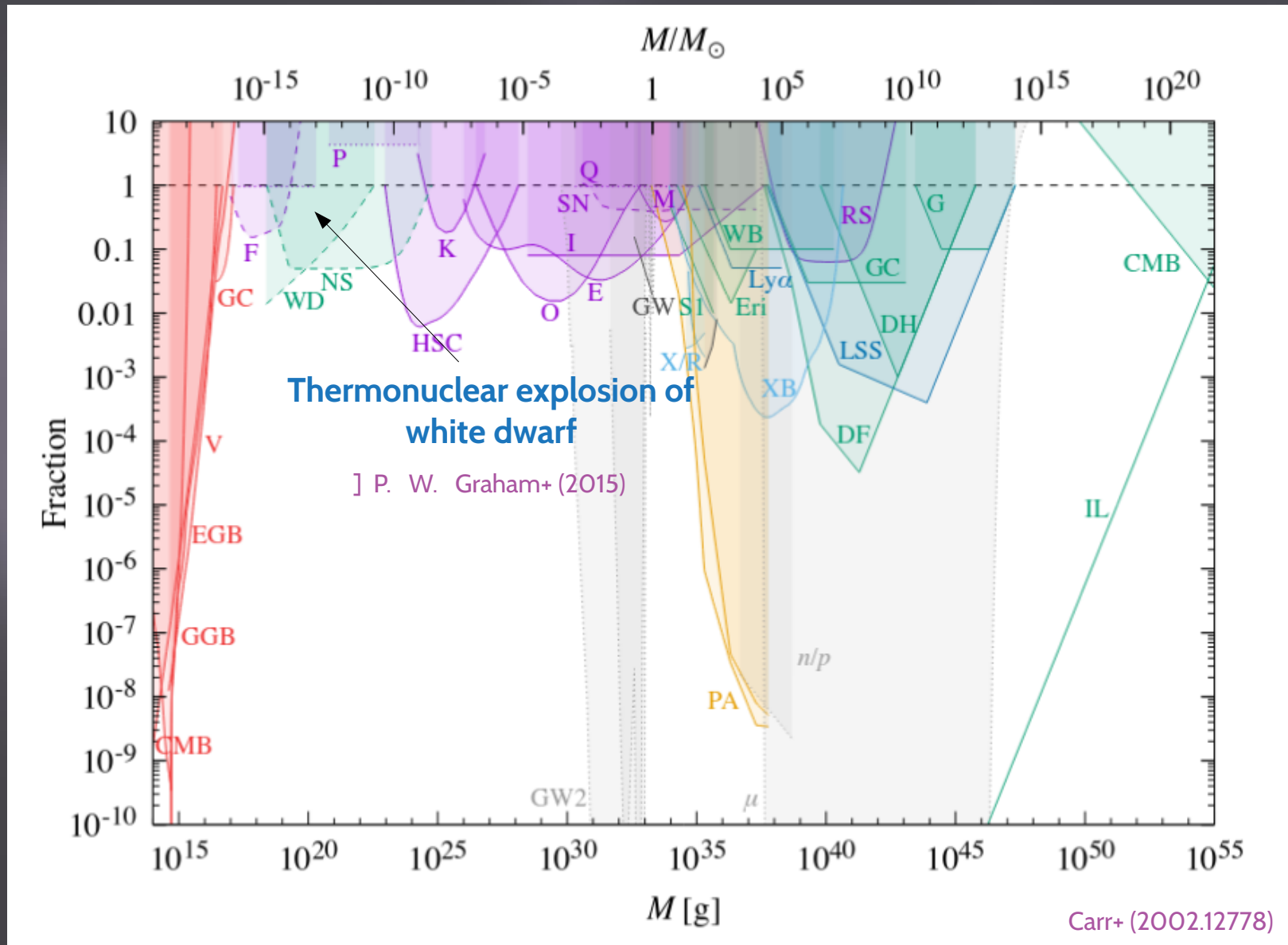
Carr+ (2002.12778)





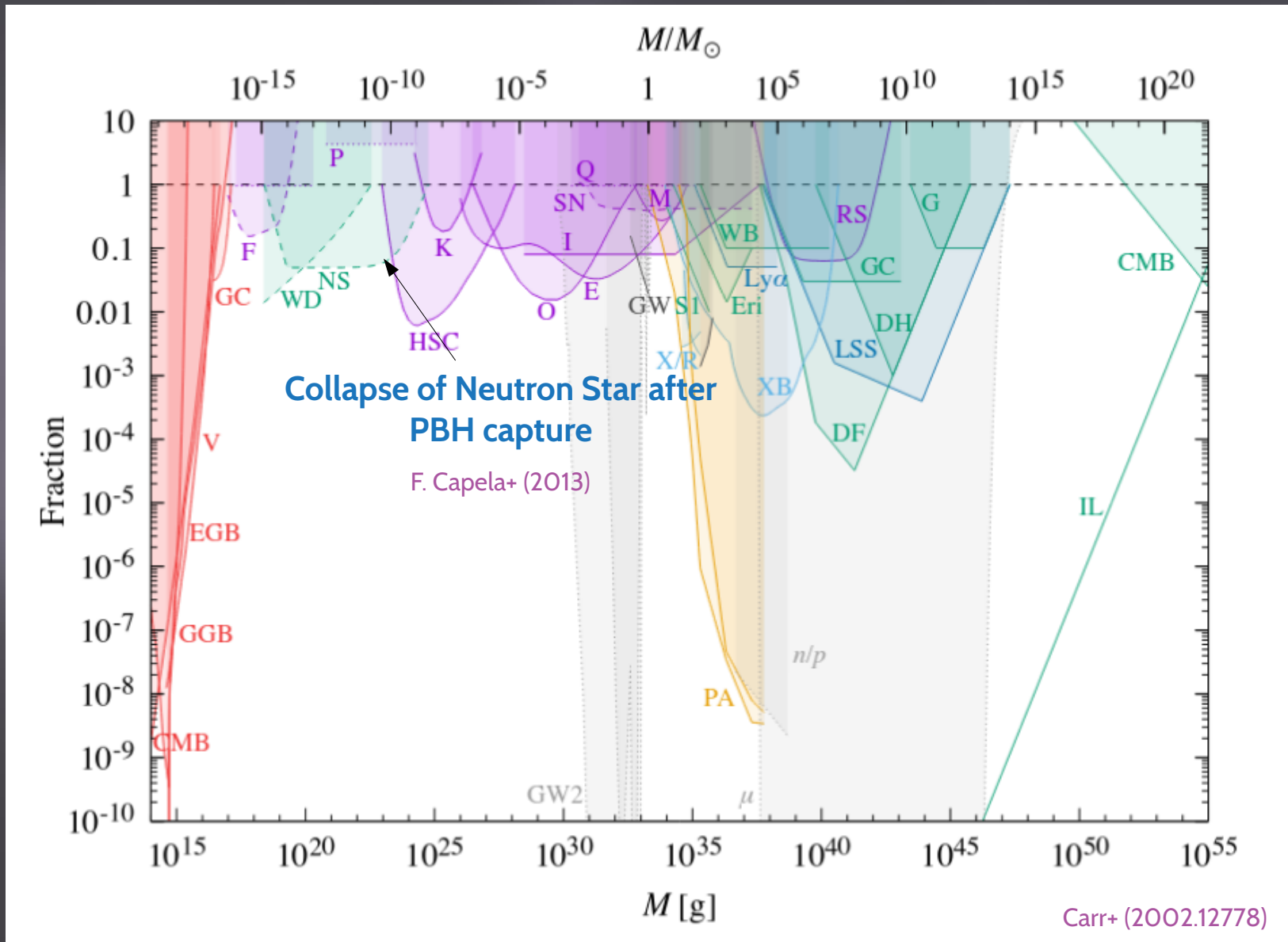




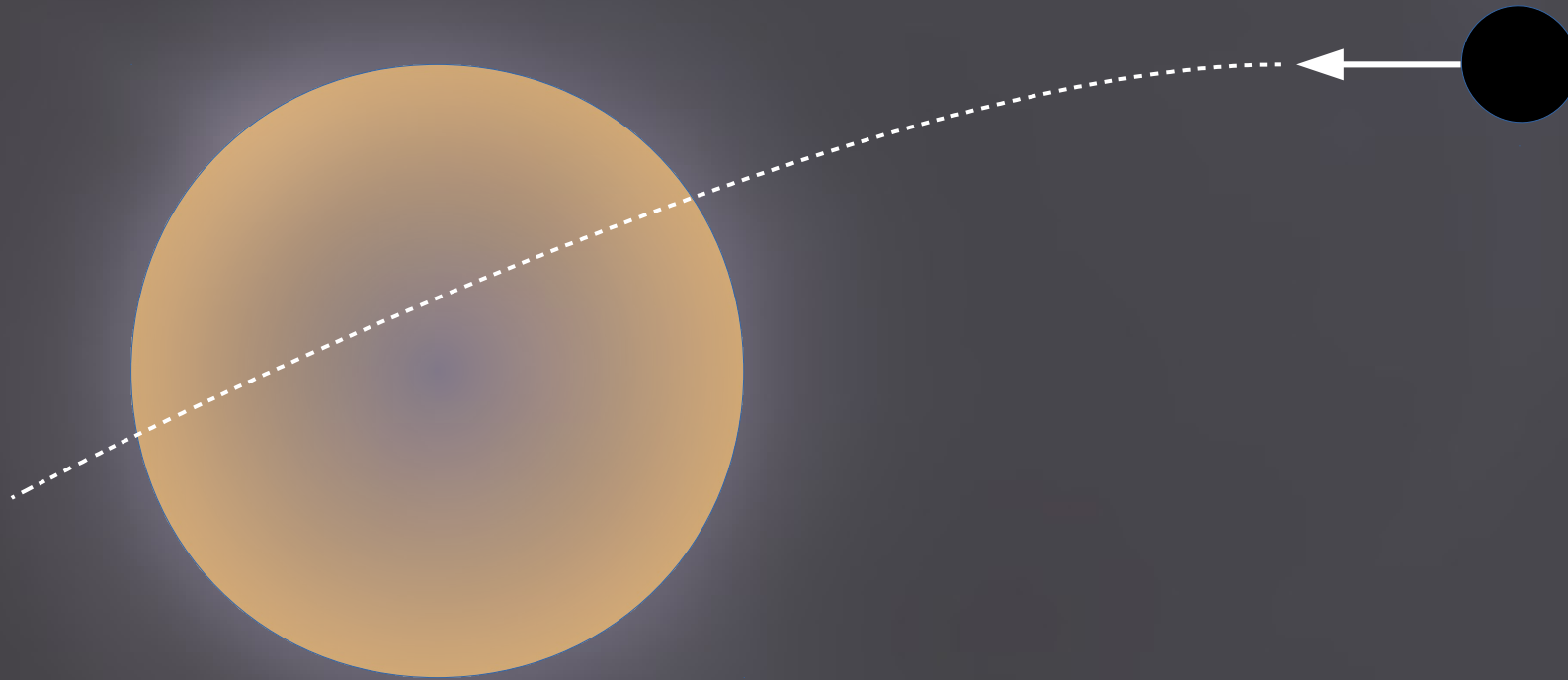




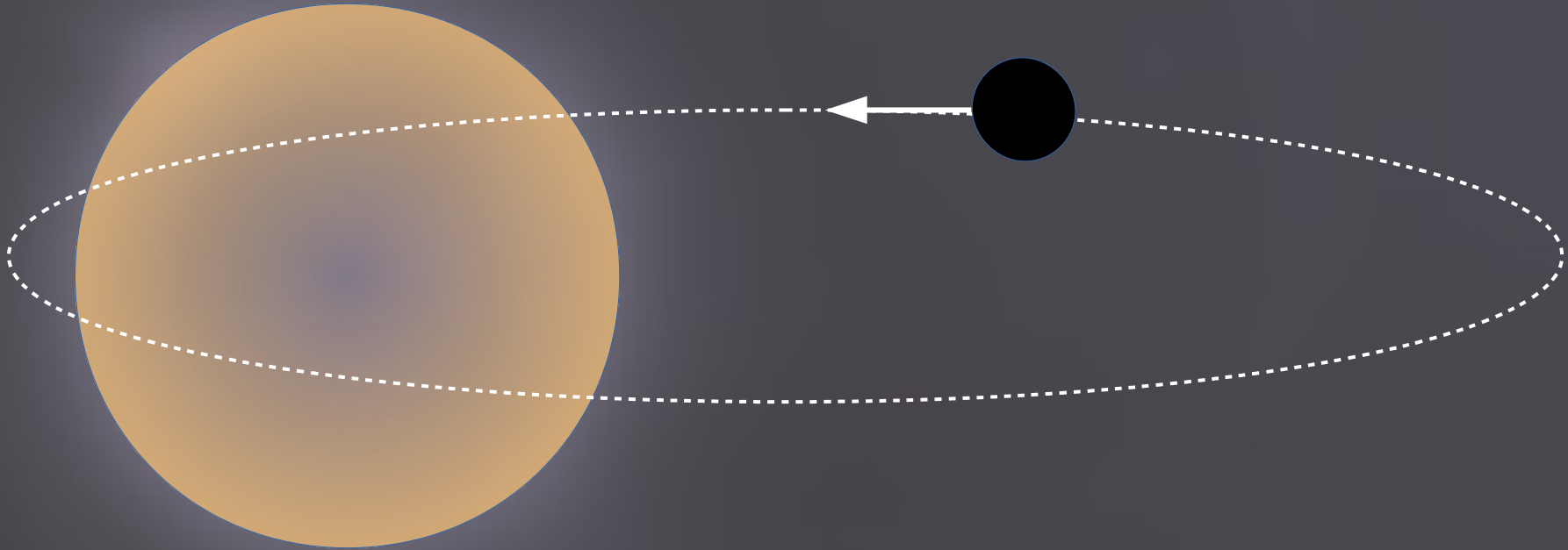
# PBH interactions with a NS - The scenario



1



2



3

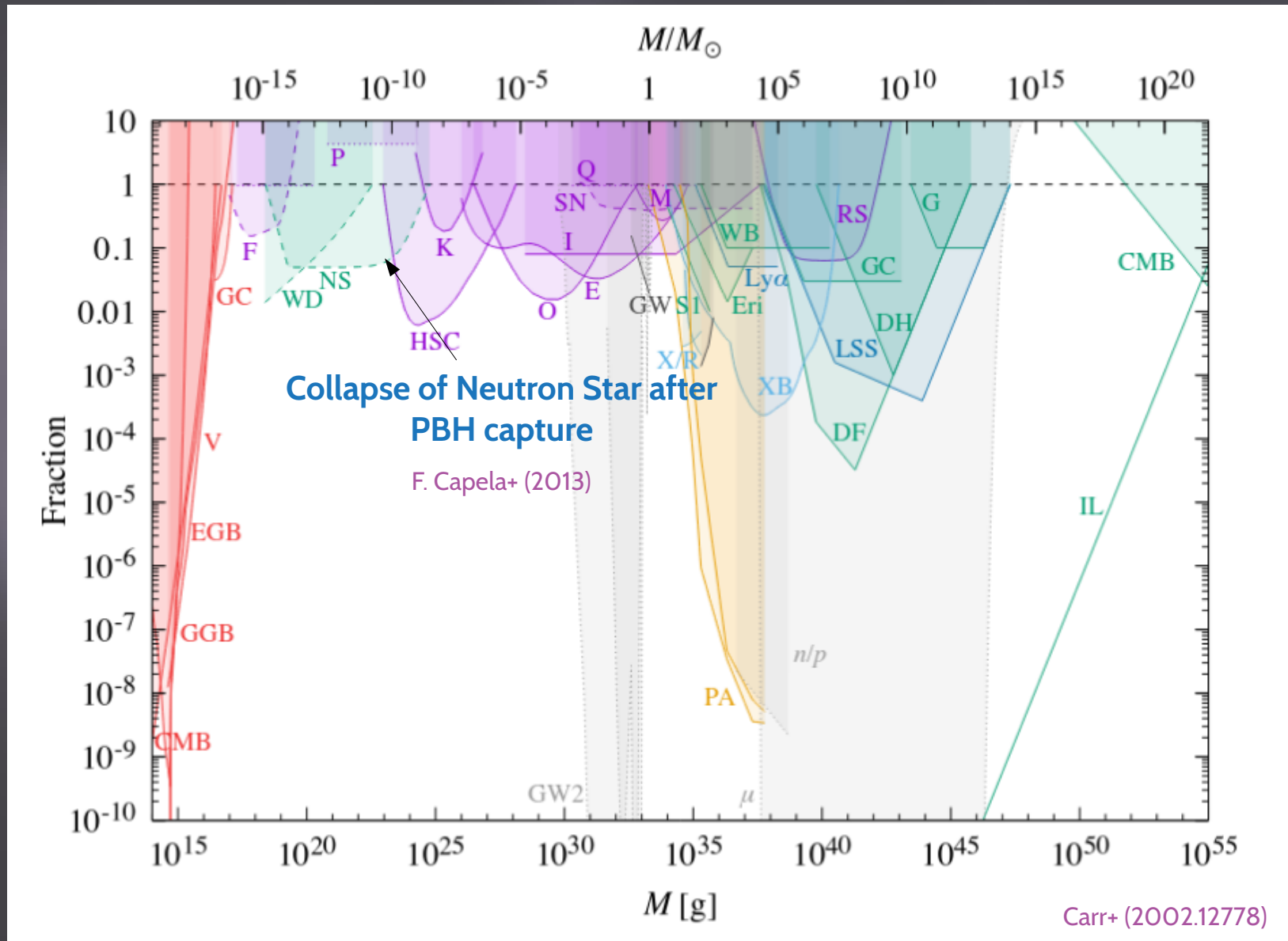


Observation of old NS in PBH-rich environment.

$$\tau_{old}^{NS} = 10 \text{ Gyr}$$

➔ Constraints on  $f_{PBH}$

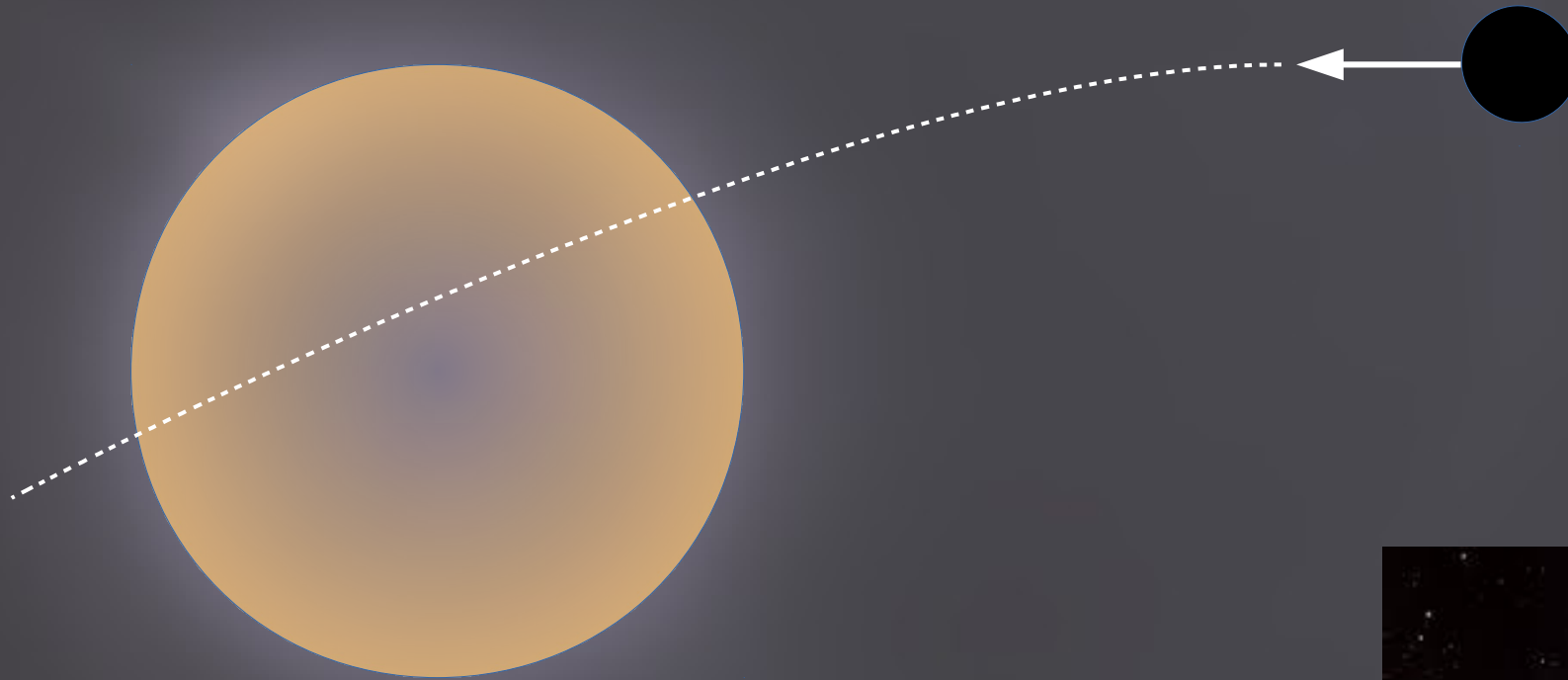
# PBH interactions with a NS - The scenario





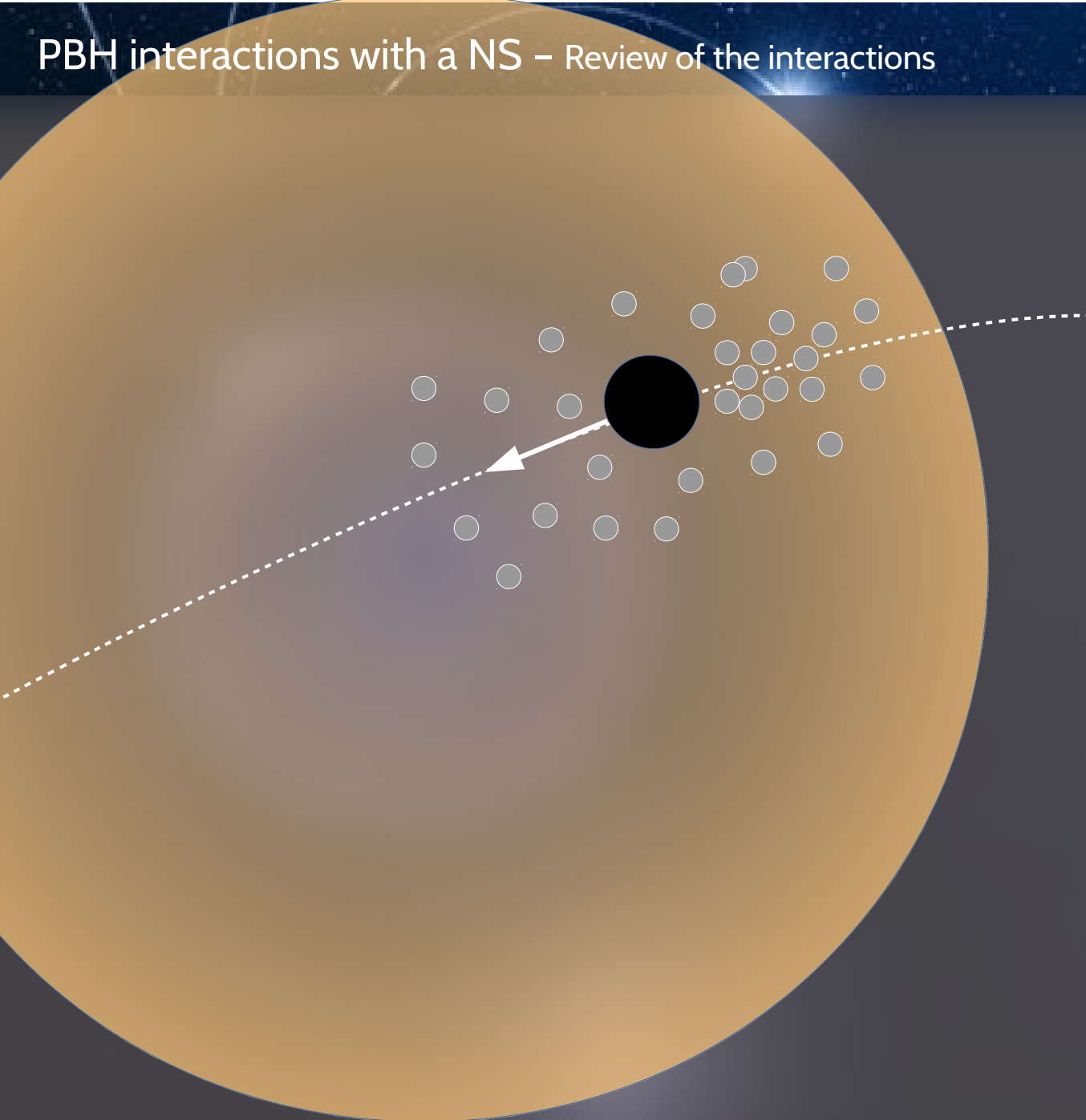
Yet, such a catastrophic event should be observable!

# 1 - Dynamical Friction

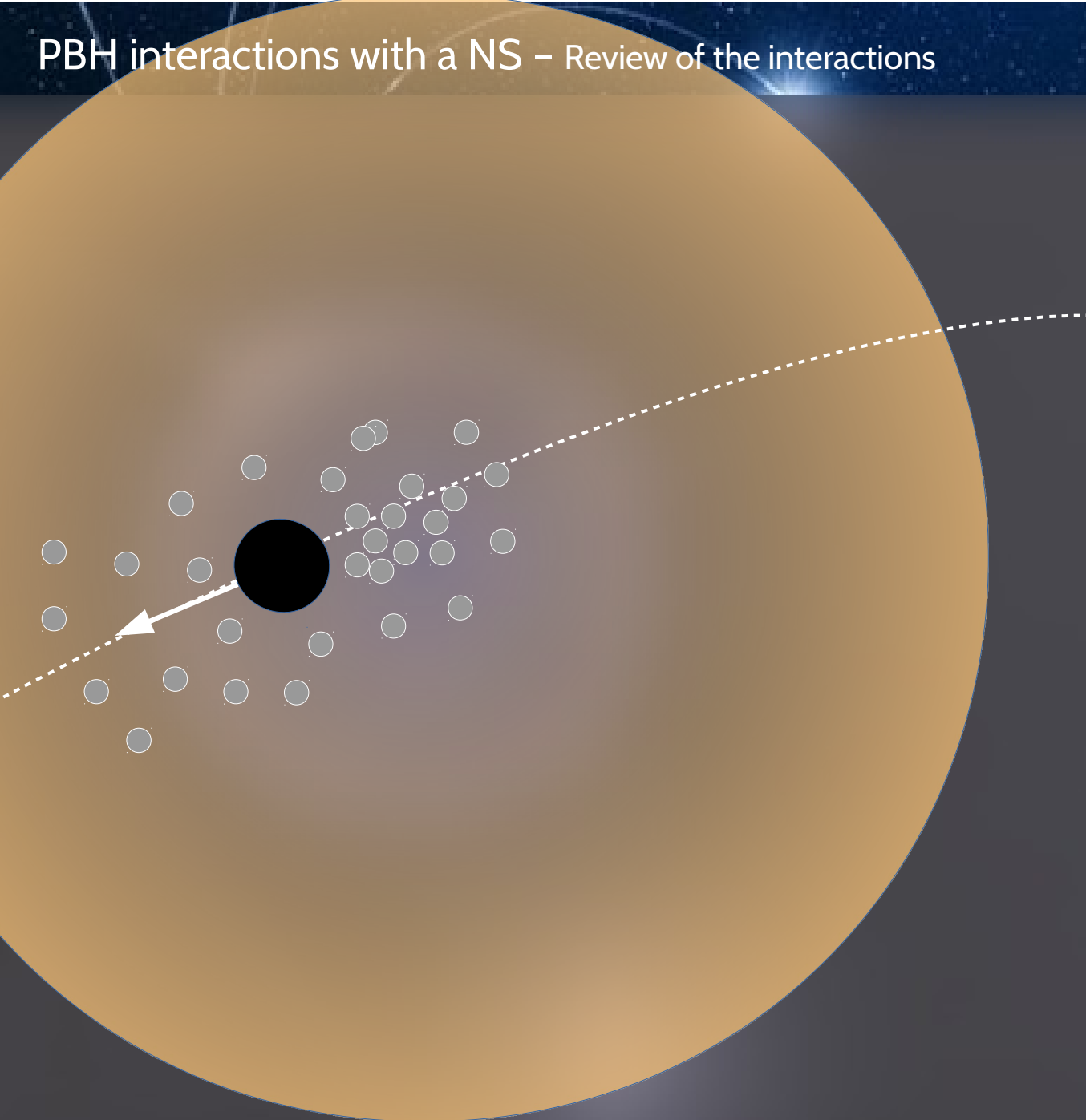


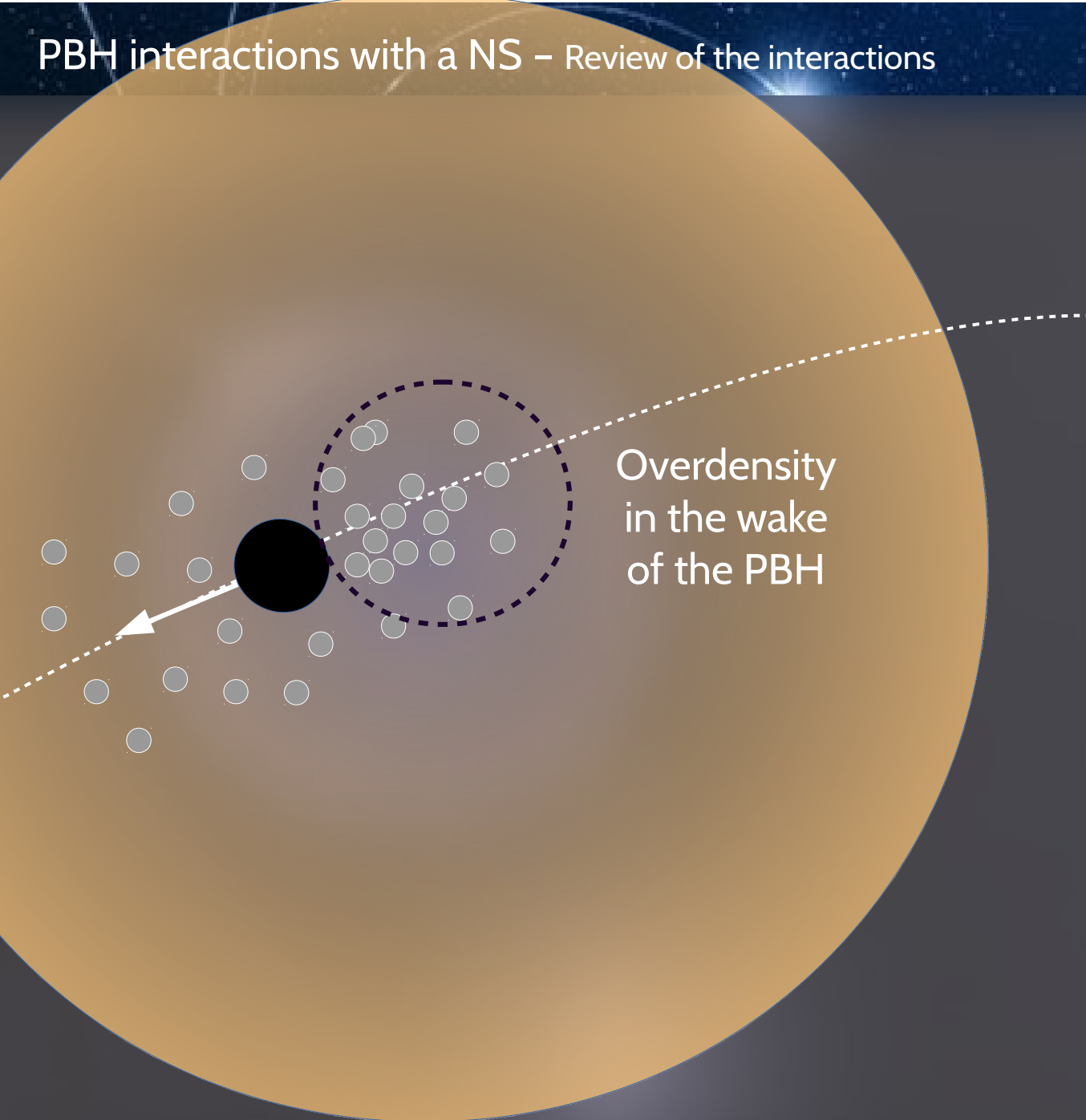


# PBH interactions with a NS – Review of the interactions

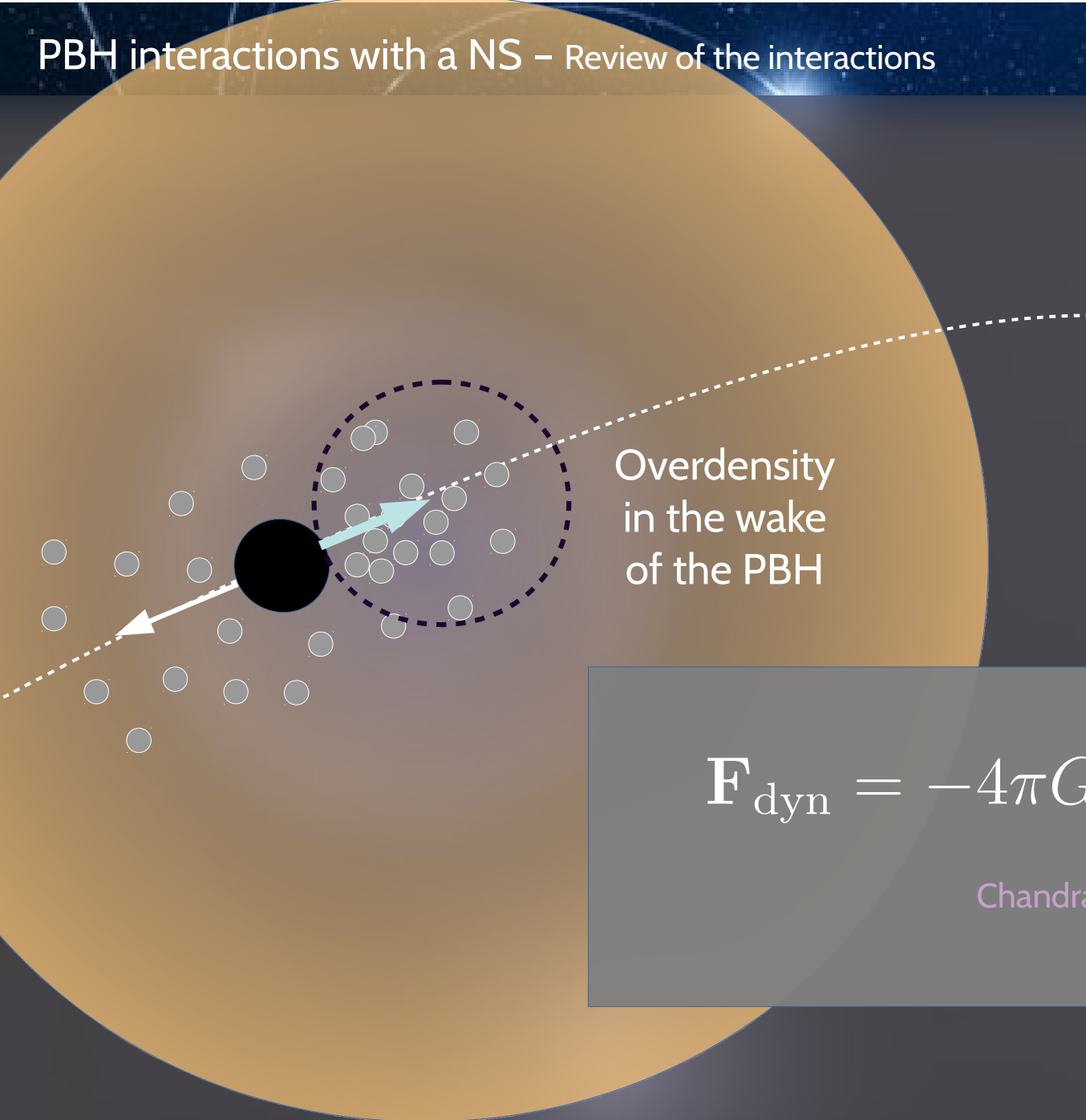


# PBH interactions with a NS – Review of the interactions





Overdensity  
in the wake  
of the PBH

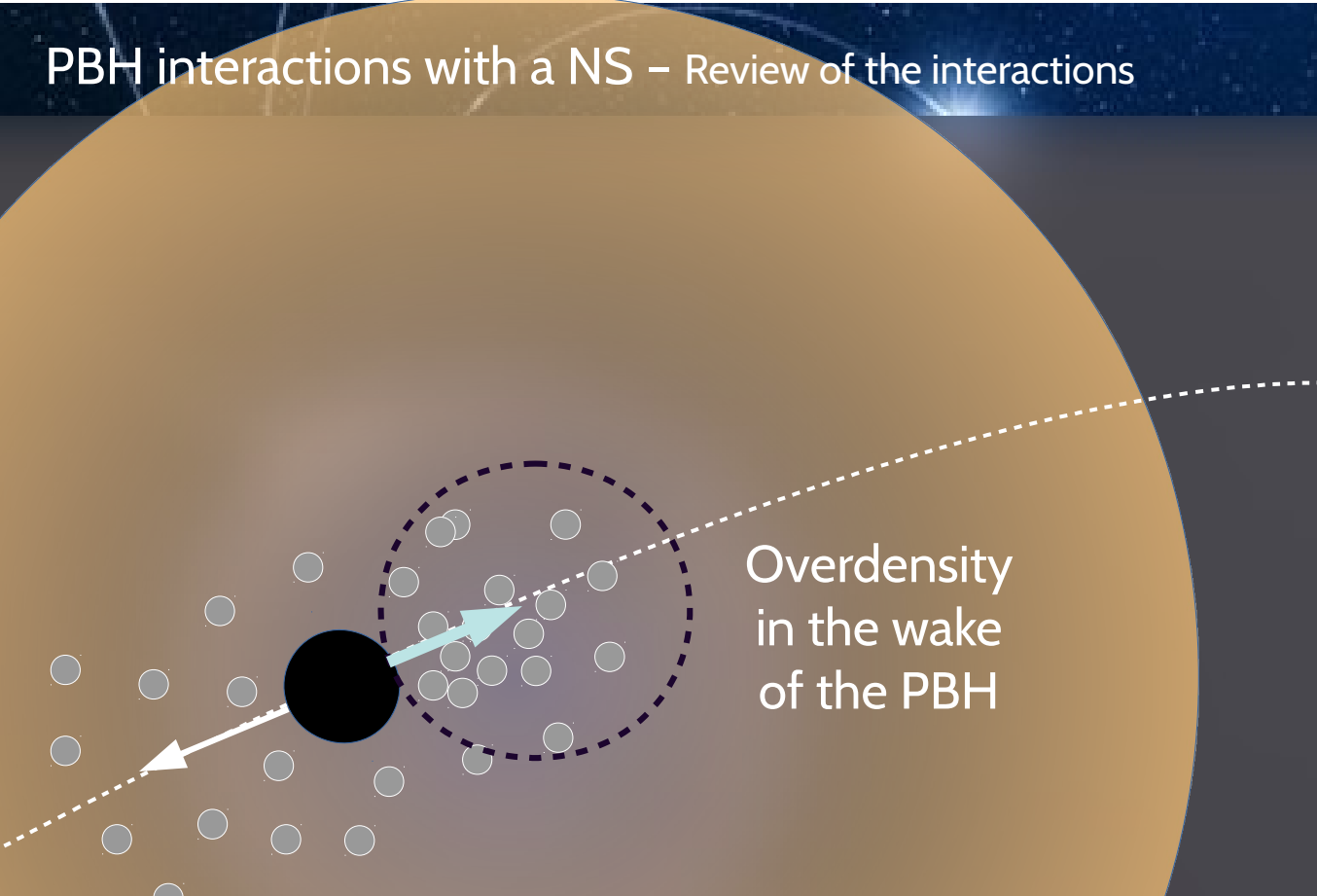


Overdensity  
in the wake  
of the PBH

$$\mathbf{F}_{\text{dyn}} = -4\pi G^2 m^2 \rho \ln \Lambda_{\text{dyn}}(v) \frac{\mathbf{v}}{v^3}$$

Chandrasekhar (1949)

# PBH interactions with a NS – Review of the interactions



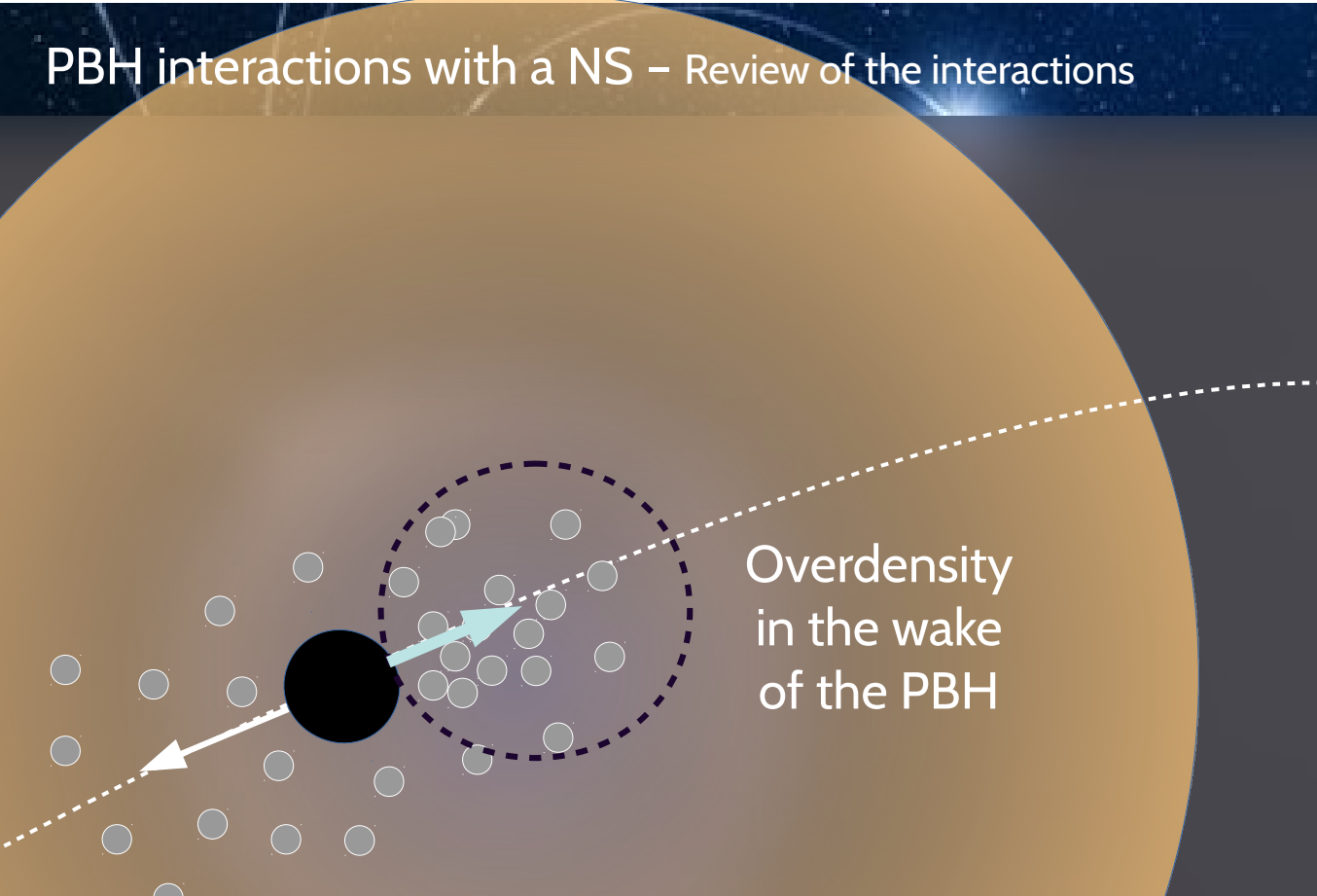
$\text{PBH } v$       $\text{PBH } v - \Delta v$   
 $E_p$  neutron      $E_p + q_0$

$$f_p(E_p, r) = \frac{1}{e^{(E_p - \mu_F(r))/T(r)} + 1}$$

Highly degenerate fermi gaz !

$T^{NS} \sim 10^{-8} \text{ GeV}$   
 $\mu_F^{NS} \sim 0.5 \text{ GeV}$

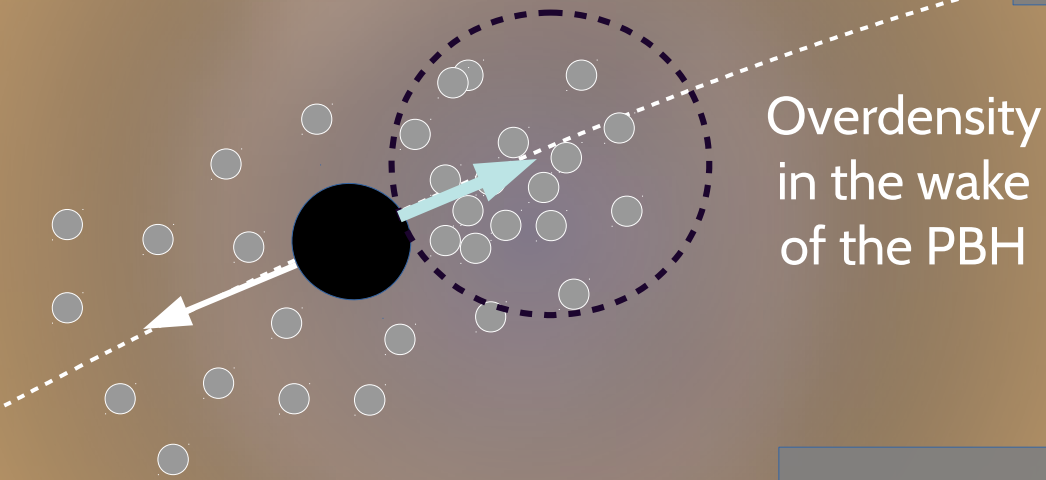
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$\text{PBH } v \rightarrow \text{PBH } v - \Delta v$   
 $E_p \text{ neutron} \rightarrow E_p + q_0$

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$E_p$



$$\mathbf{F}_{\text{dyn}} = -4\pi G^2 m^2 \rho \ln \Lambda_{\text{dyn}}(v) \frac{\mathbf{v}}{v^3}$$

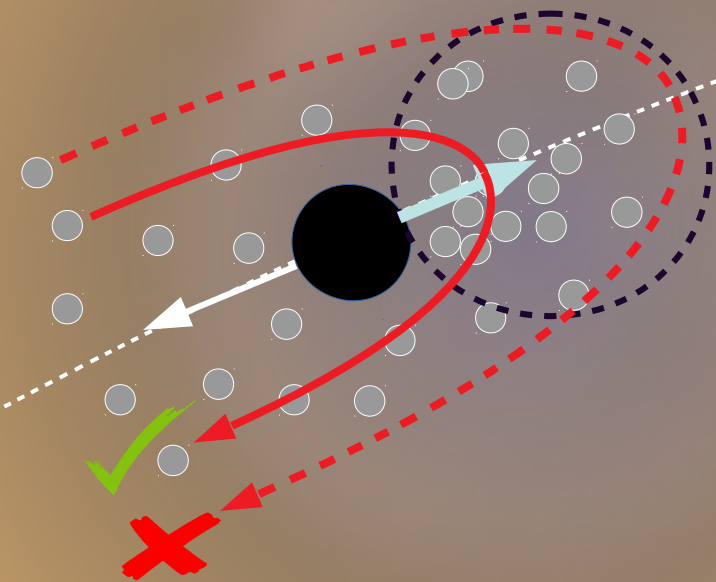
Chandrasekhar (1949)

$$\ln \Lambda_{\text{dyn}}(v) = v^4 \gamma^2 \frac{2}{R_g^2} \int_{d_{\text{crit}}}^{d_{\text{max}}} dx x (1 - \cos \varphi(x))$$

Capela+ (2013)



# PBH interactions with a NS – Review of the interactions



Overdensity  
in the wake  
of the PBH

Fermi- suppressed  
scatterings

$$\mathbf{F}_{\text{dyn}} = -4\pi G^2 m^2 \rho \ln \Lambda_{\text{dyn}}(v) \frac{\mathbf{v}}{v^3}$$

Chandrasekhar (1949)

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Capela+ (2013)

-> DF is suppressed by a factor of a few, up to 10

Derived for a collisionless medium

$$\mathbf{F}_{\text{dyn}} = -4\pi G^2 m^2 \rho \ln \Lambda_{\text{dyn}}(v) \frac{\mathbf{v}}{v^3}$$

Chandrasekhar (1949)

NS = strongly interacting neutron fluid

in the wake  
of the PBH

Fermi-suppressed  
process

Dynamical friction is suppressed by a factor of a few, up to 10.

Derived for a collisionless medium

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Chandrasekhar (1949)

NS = strongly interacting neutron fluid

in the wake  
of the PBH

Collisionless if  $\tau_{\text{gravitation}} \ll \tau_{\text{causal}}$

$$v \gg c_s$$

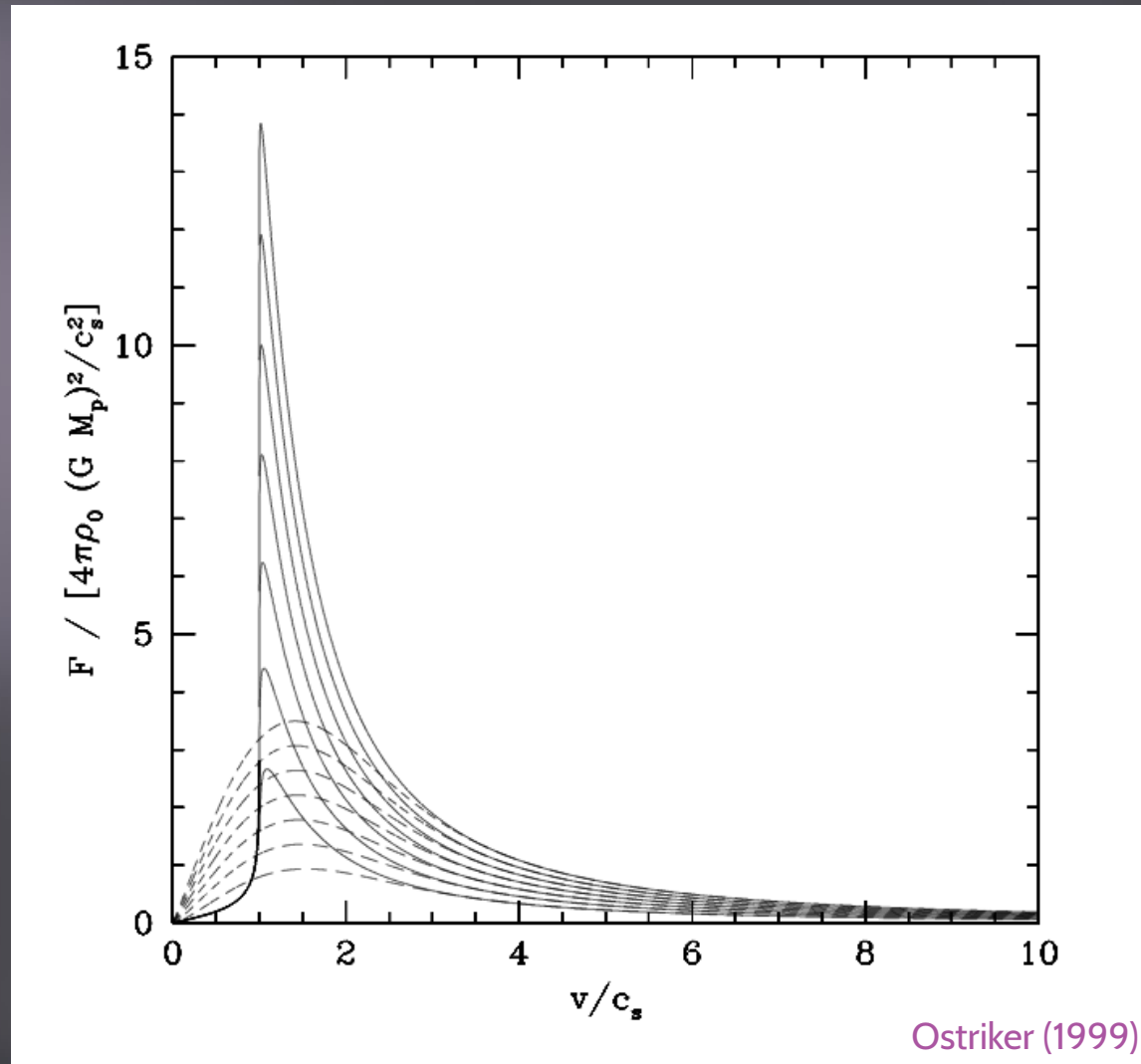
$$\mathcal{M} = v/c_s \gg 1$$

Fermi-suppressed  
process

Dynamical friction is suppressed by a factor of a few, up to 10.

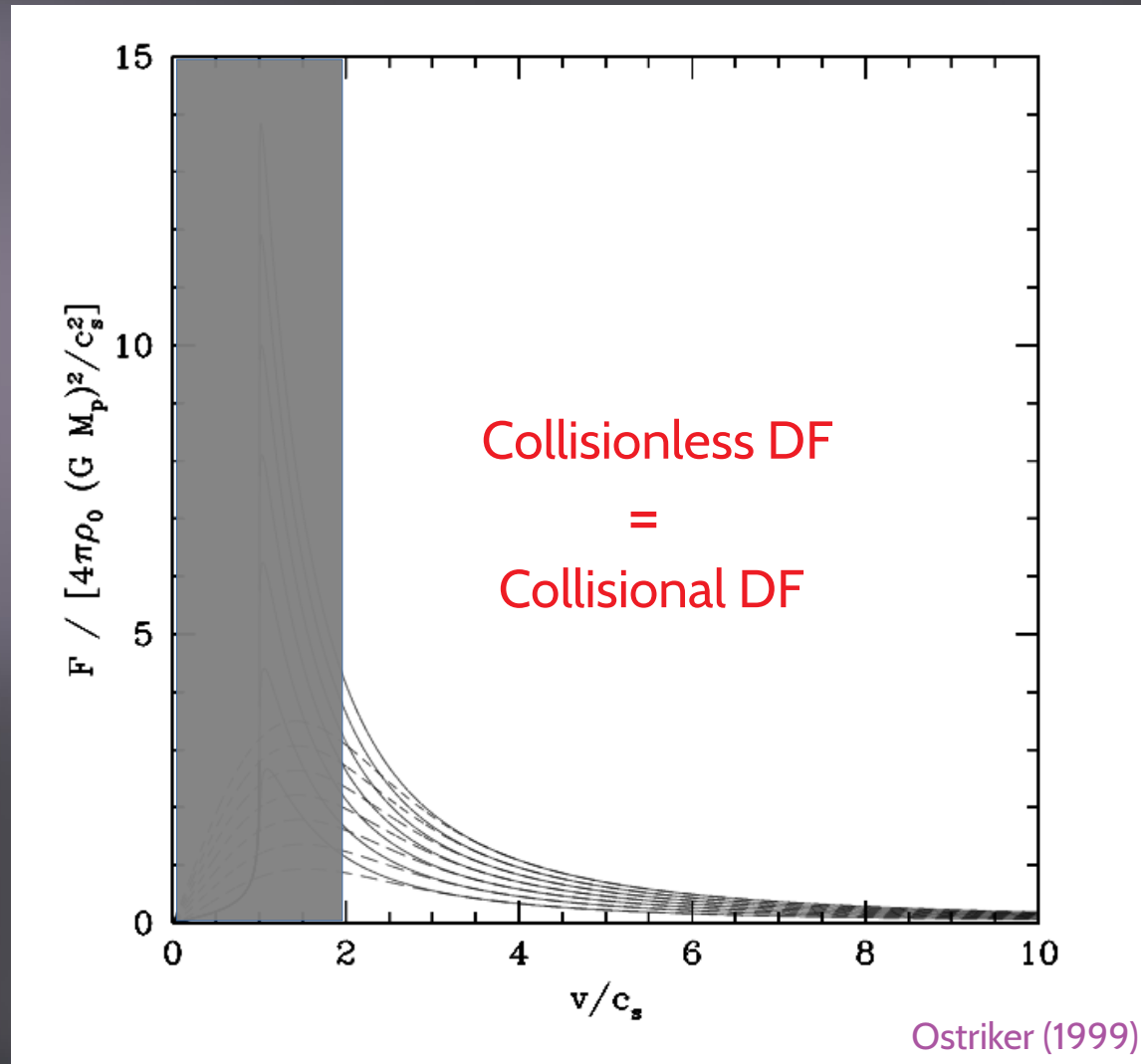
# 1 - Dynamical Friction:

In a collisionless or a collisional medium?



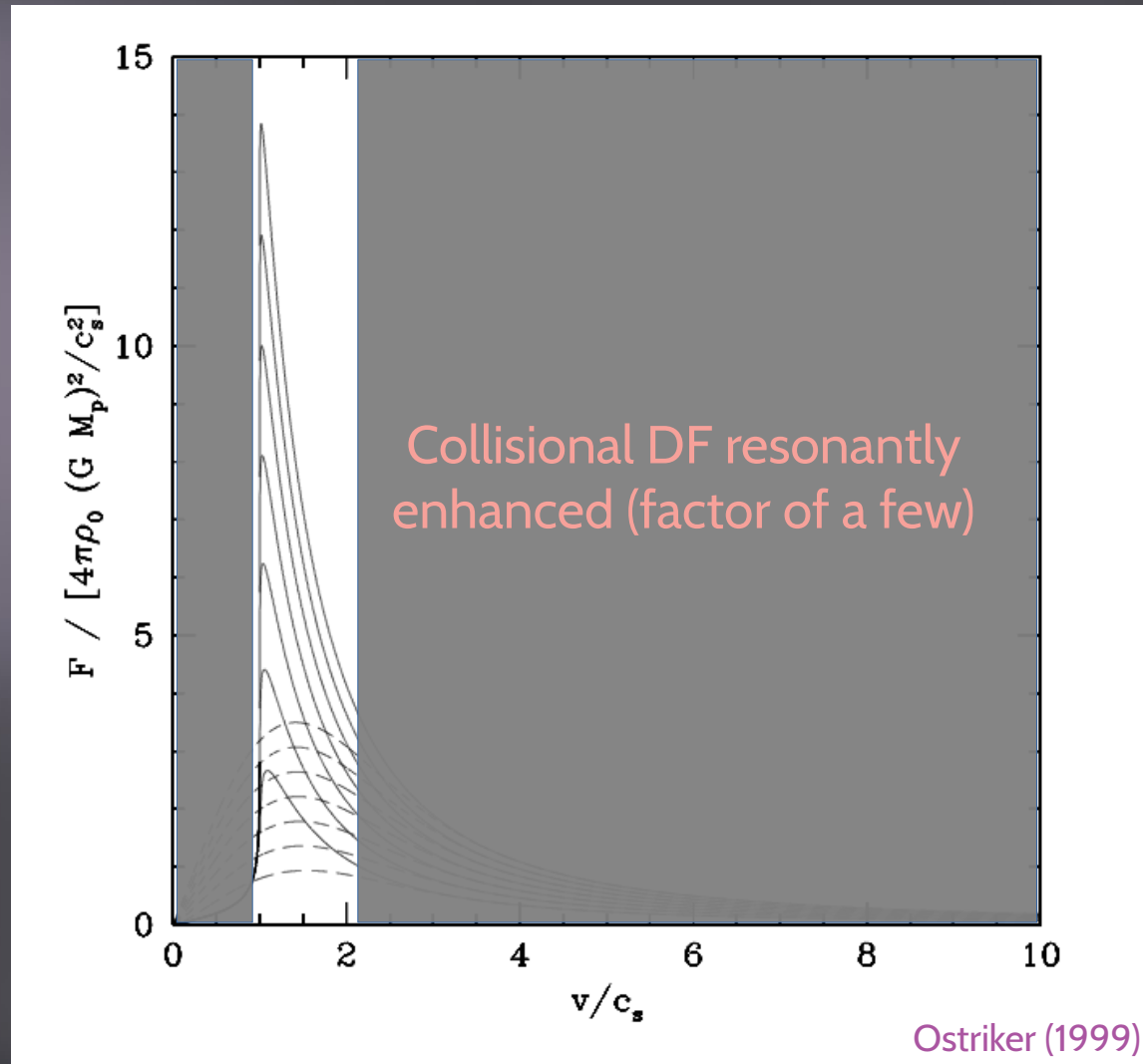
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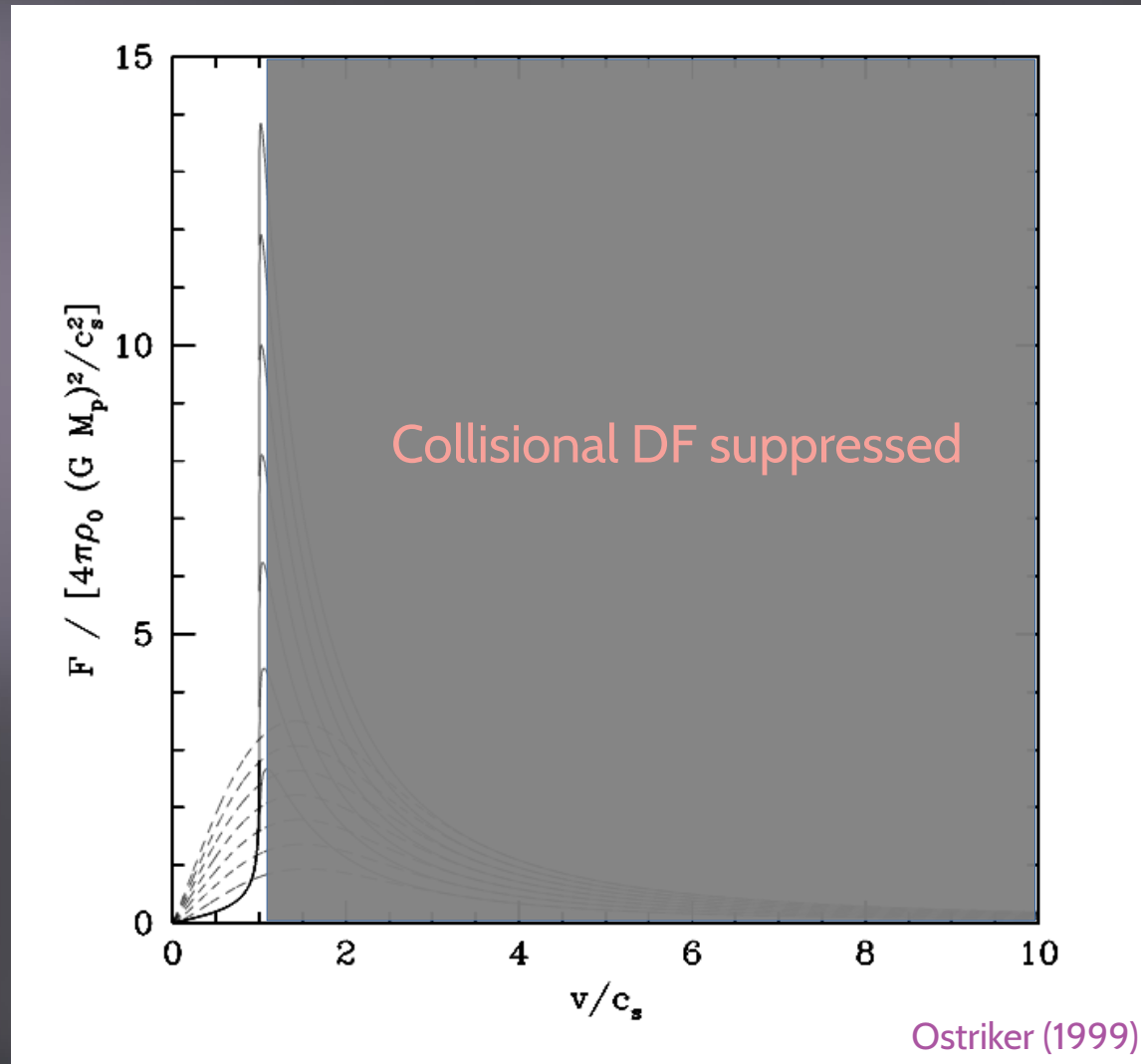
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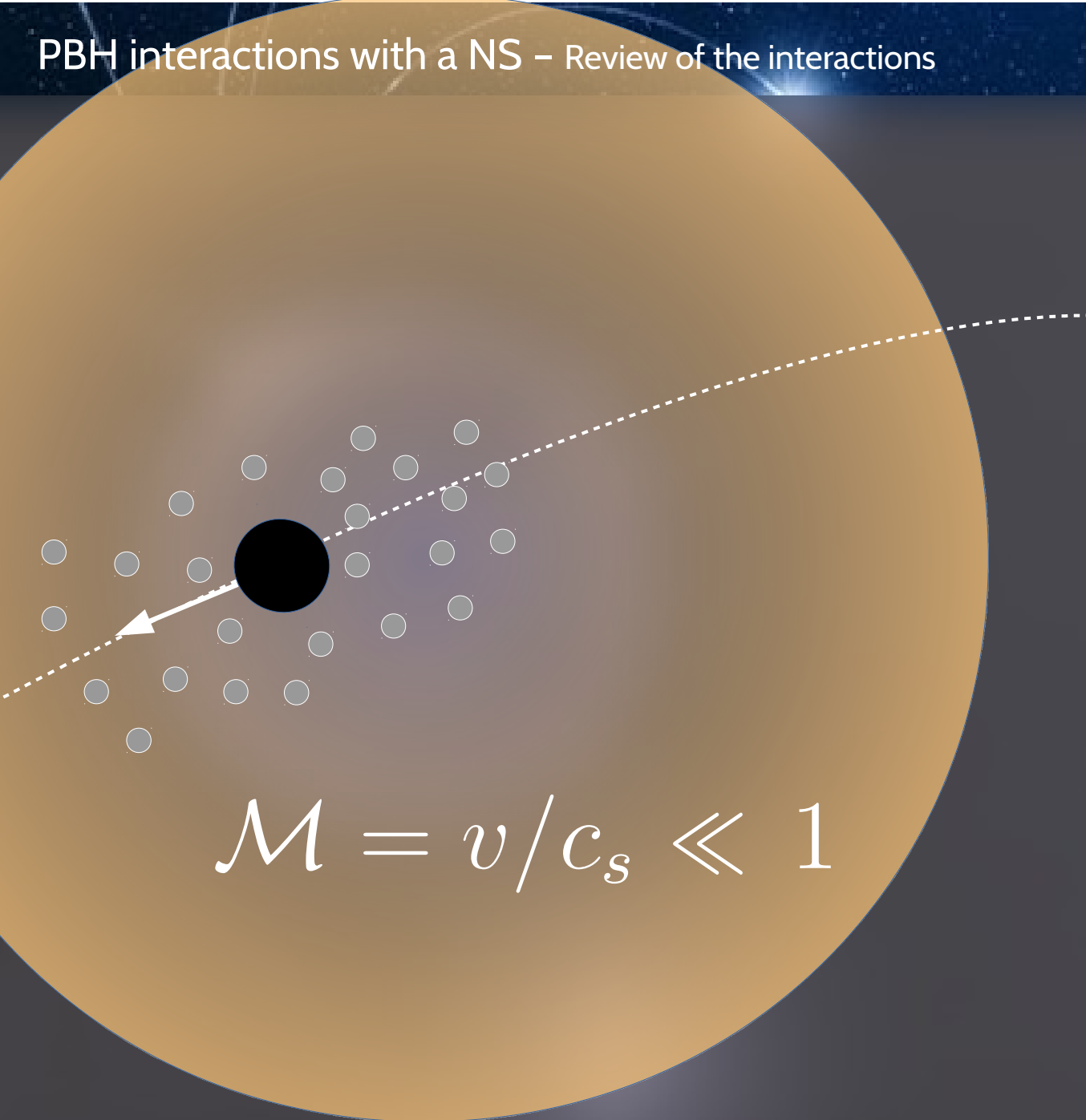


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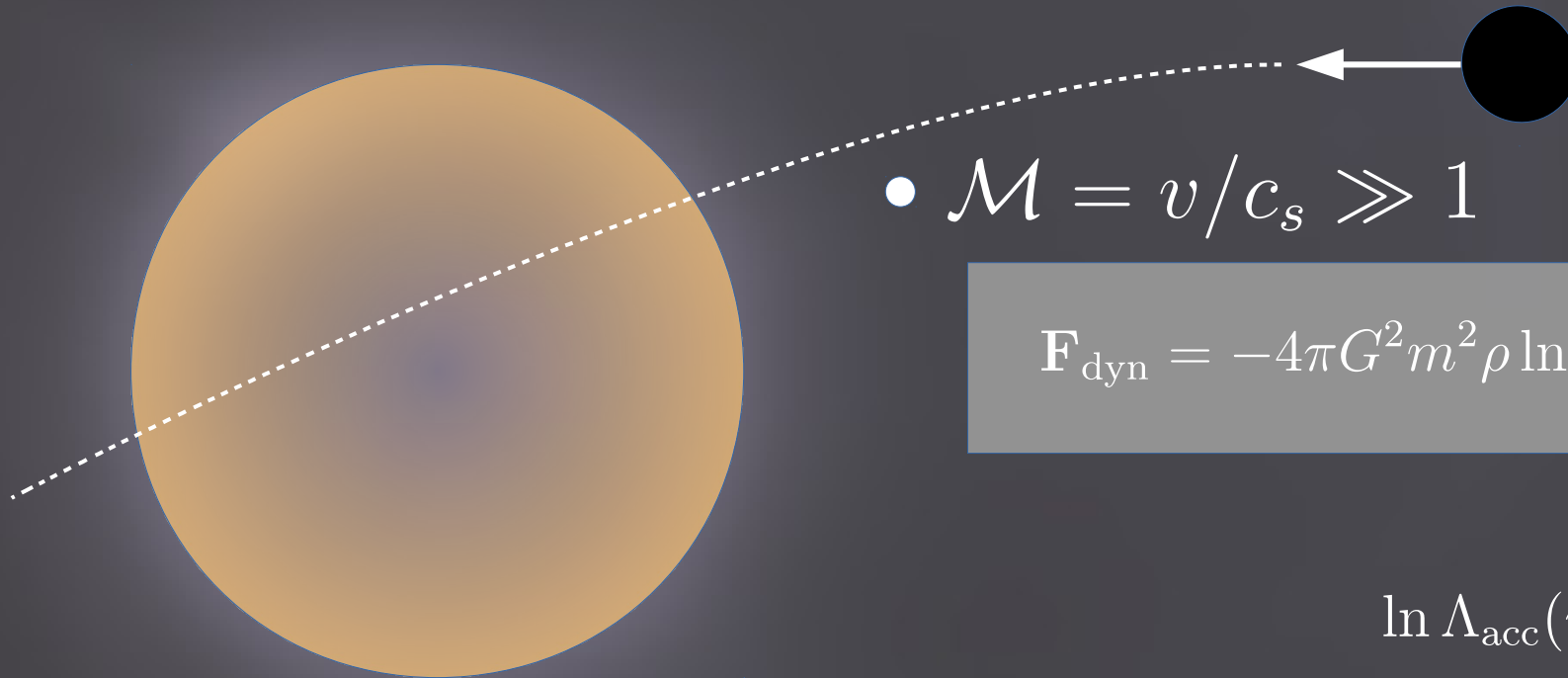






$$\mathcal{M} = v/c_s \ll 1$$

## 2 - Accretion



- $\mathcal{M} = v/c_s \gg 1$

$$\mathbf{F}_{\text{dyn}} = -4\pi G^2 m^2 \rho \ln \Lambda_{\text{dyn}}(v) \frac{\mathbf{v}}{v^3}$$

Capela+ (2013)

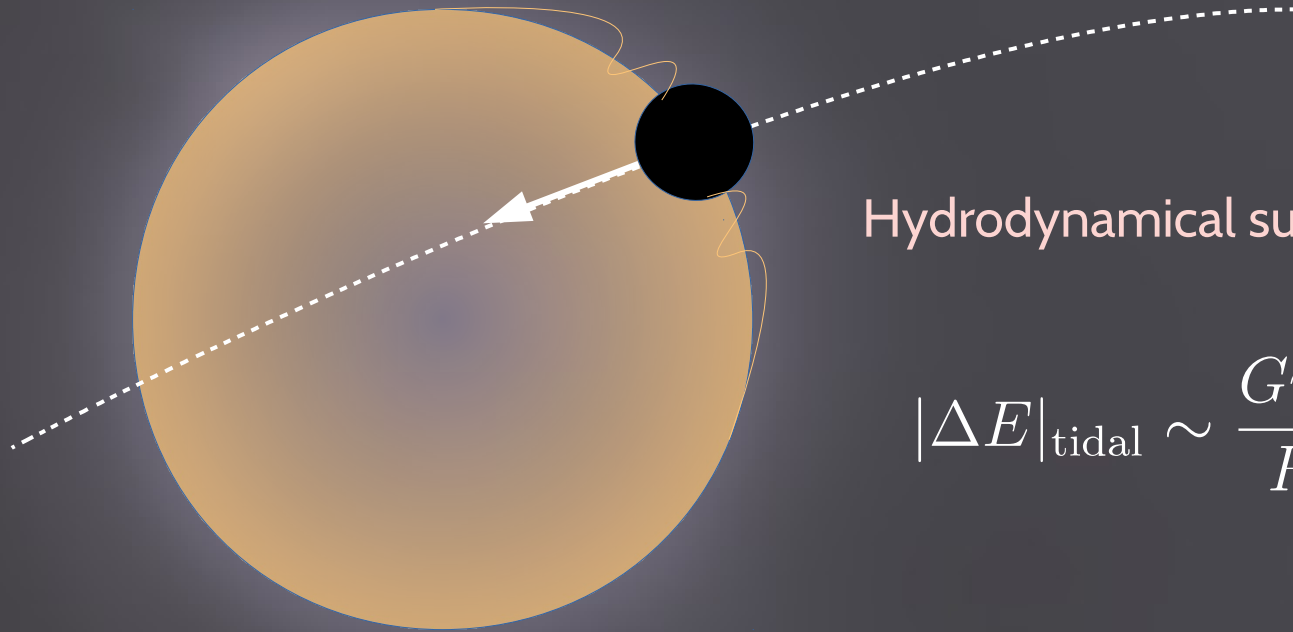
$$\ln \Lambda_{\text{acc}}(v) = v^4 \gamma^2 \frac{d_{\text{crit}}^2}{R_g^2}$$

- $\mathcal{M} = v/c_s \ll 1$

$$\mathbf{F}_{\text{drag}} = -\dot{m} \mathbf{v} = -4\pi G^2 m^2 \rho \frac{\mathbf{v}}{c_s^3}$$

Y.G. et al. PRD (2020)

### 3 - Surface waves



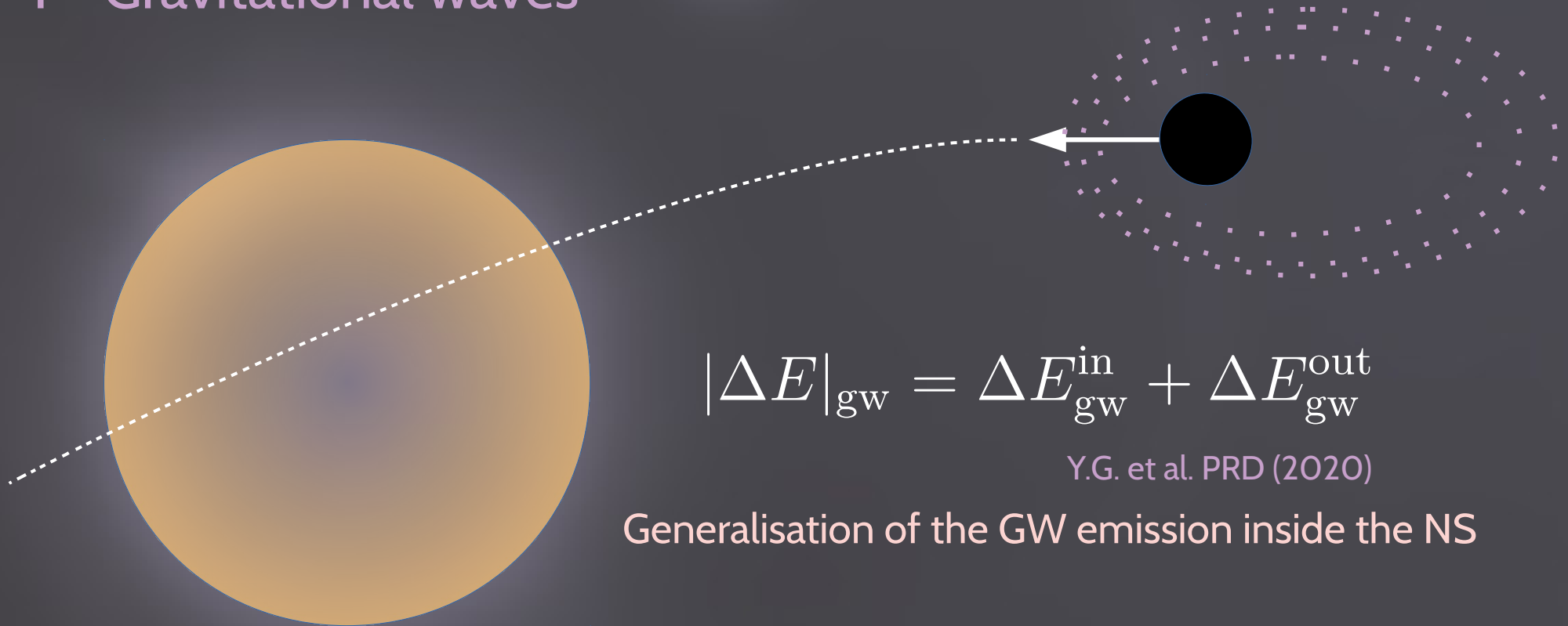
Hydrodynamical surface waves:

$$|\Delta E|_{\text{tidal}} \sim \frac{Gm^2}{R_\star} \sum_{l=2}^{\infty} \left( \frac{R_\star}{r_{\text{min}}} \right)^{2l+2} T_l,$$

Defillon+ (2014)

Press&Teukolsky (1977)

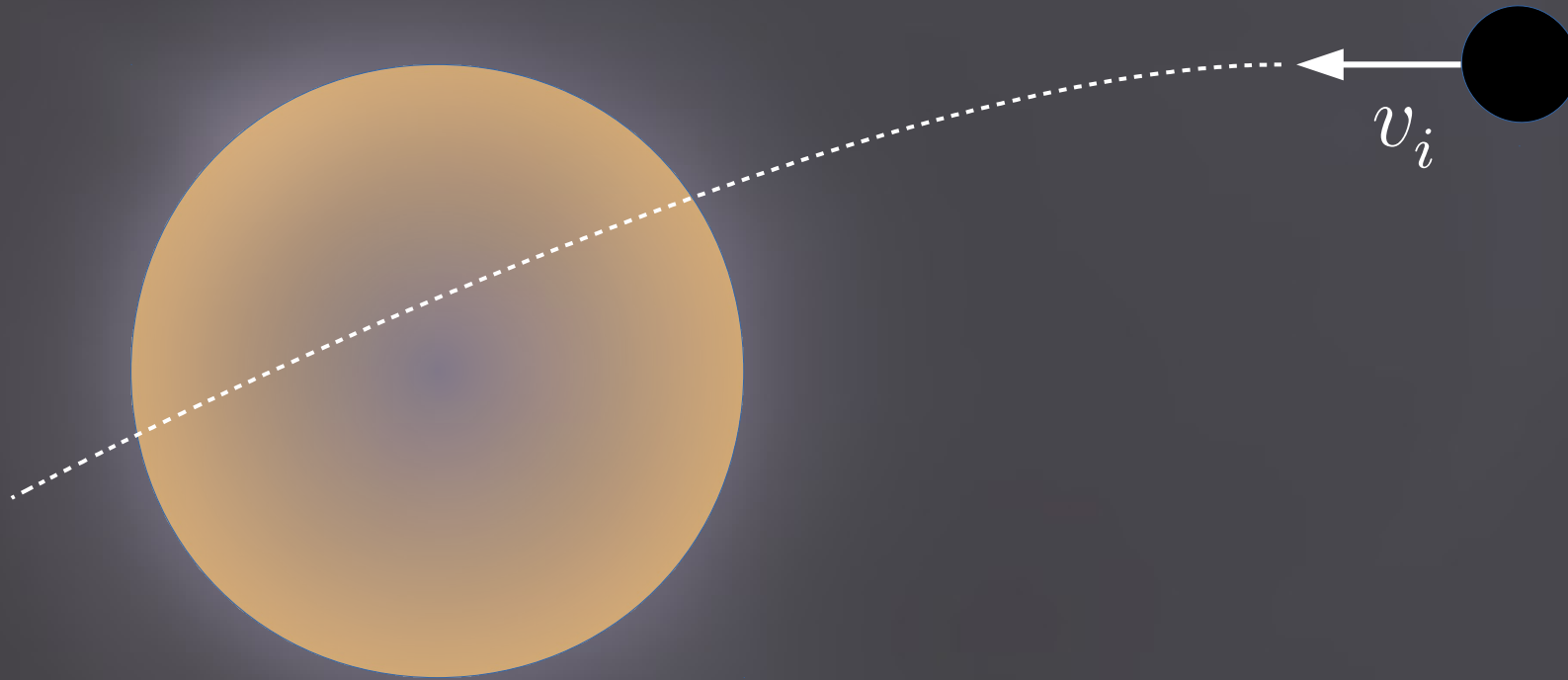
## 4 - Gravitational waves

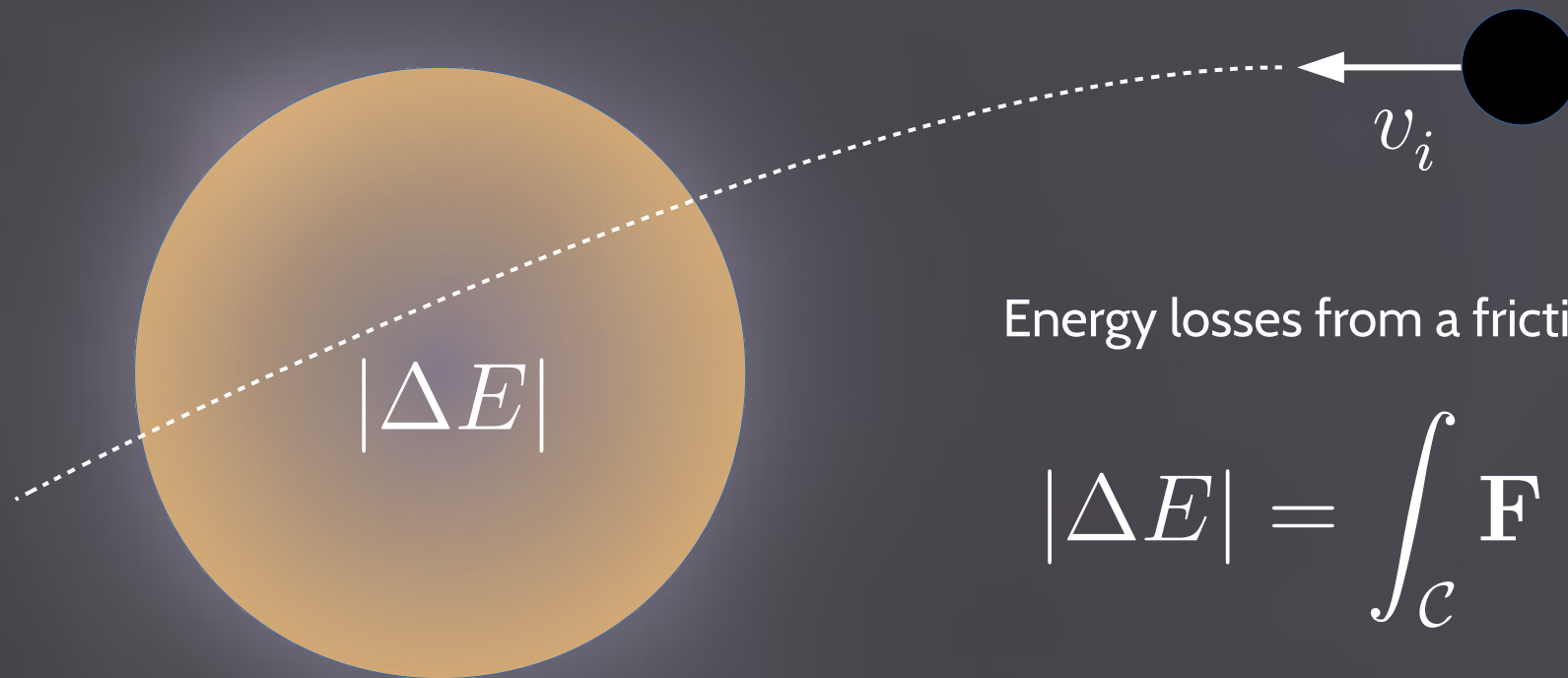


$$|\Delta E|_{\text{gw}} = \Delta E_{\text{gw}}^{\text{in}} + \Delta E_{\text{gw}}^{\text{out}}$$

Y.G. et al. PRD (2020)

Generalisation of the GW emission inside the NS



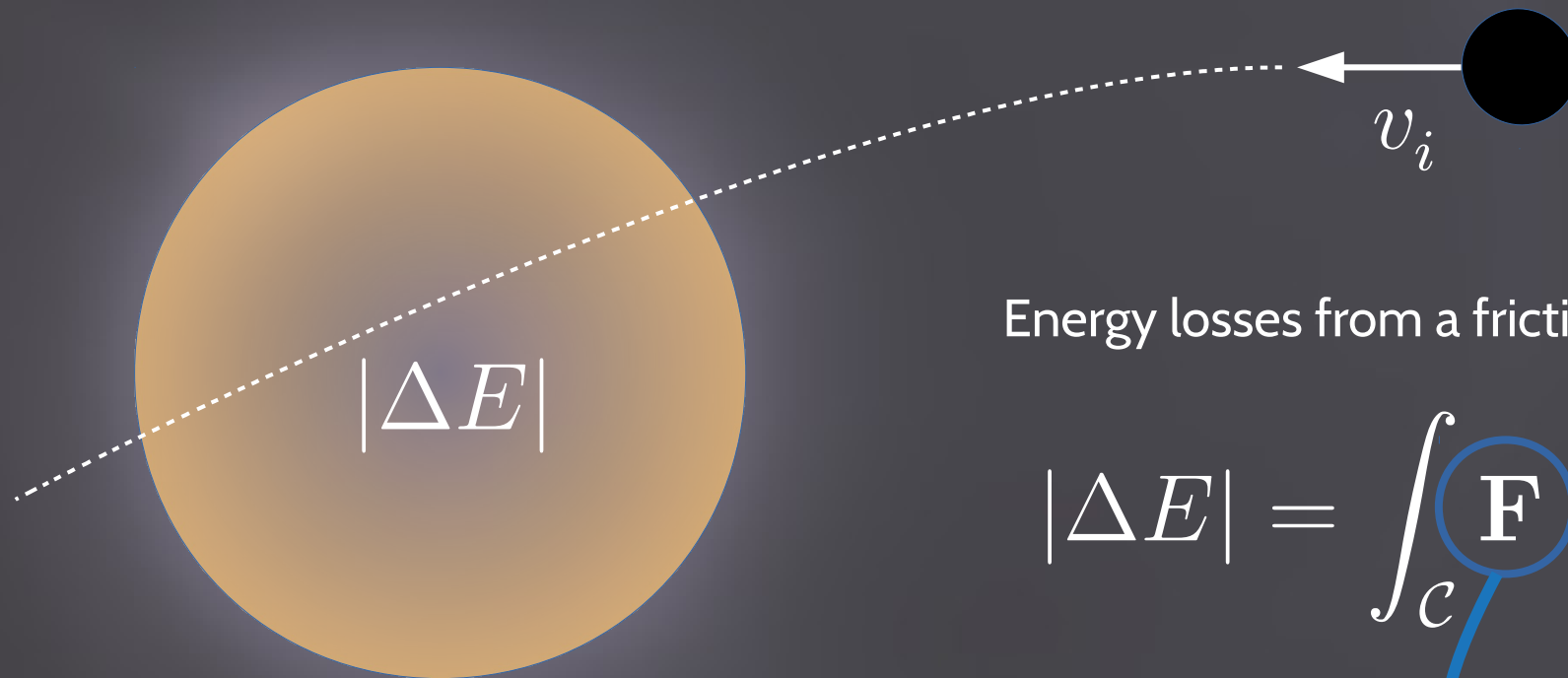


Energy losses from a friction force

$$|\Delta E| = \int_{\mathcal{C}} \mathbf{F} \cdot d\mathbf{l}$$

Capture condition

$$|\Delta E| > E_i = \frac{1}{2} m v_i^2$$



Energy losses from a friction force

$$|\Delta E| = \int_c \mathbf{F} \cdot d\mathbf{l}$$

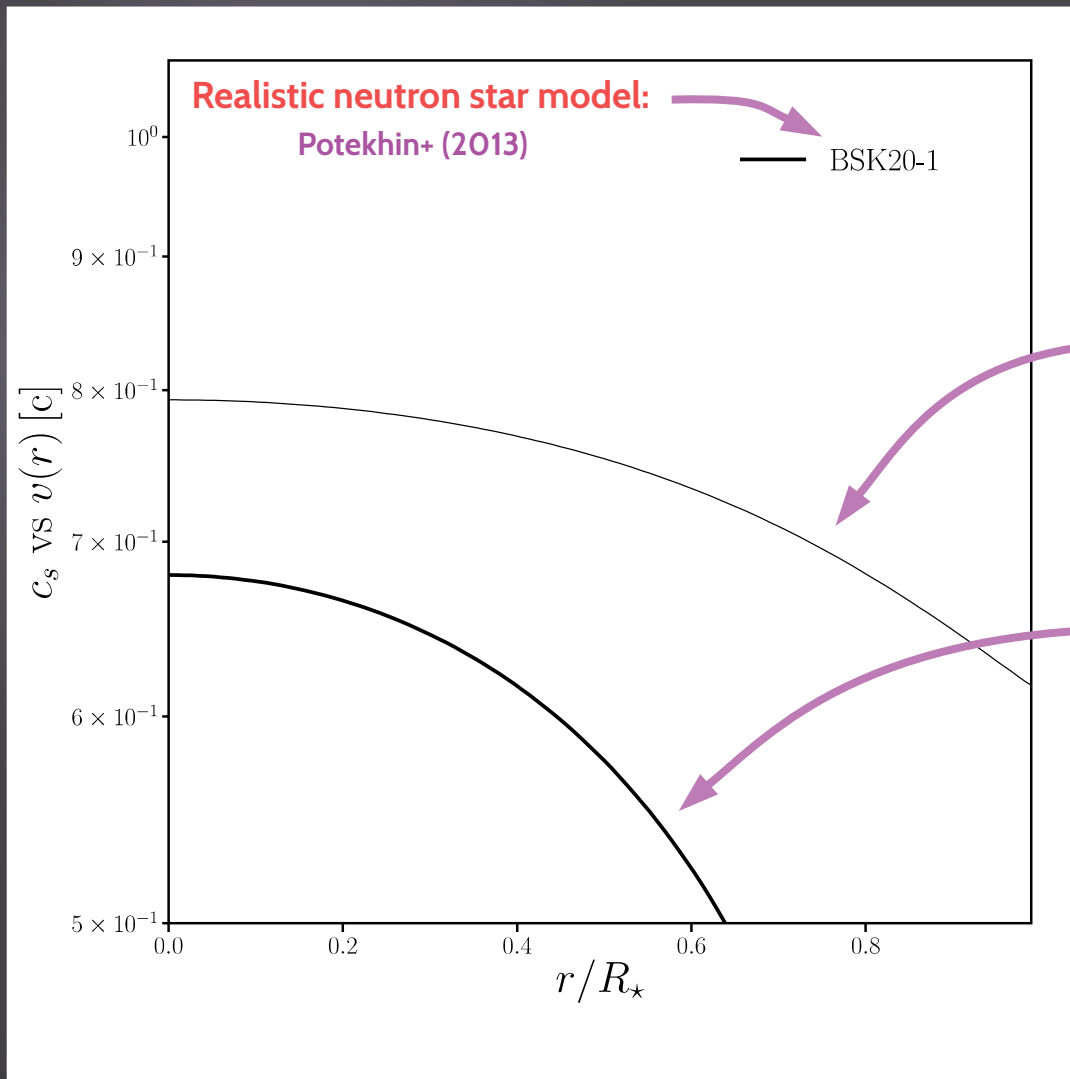
Forces we saw

Capture condition

$$|\Delta E| > E_i = \frac{1}{2} m v_i^2$$



# What is the speed regime for capture ?

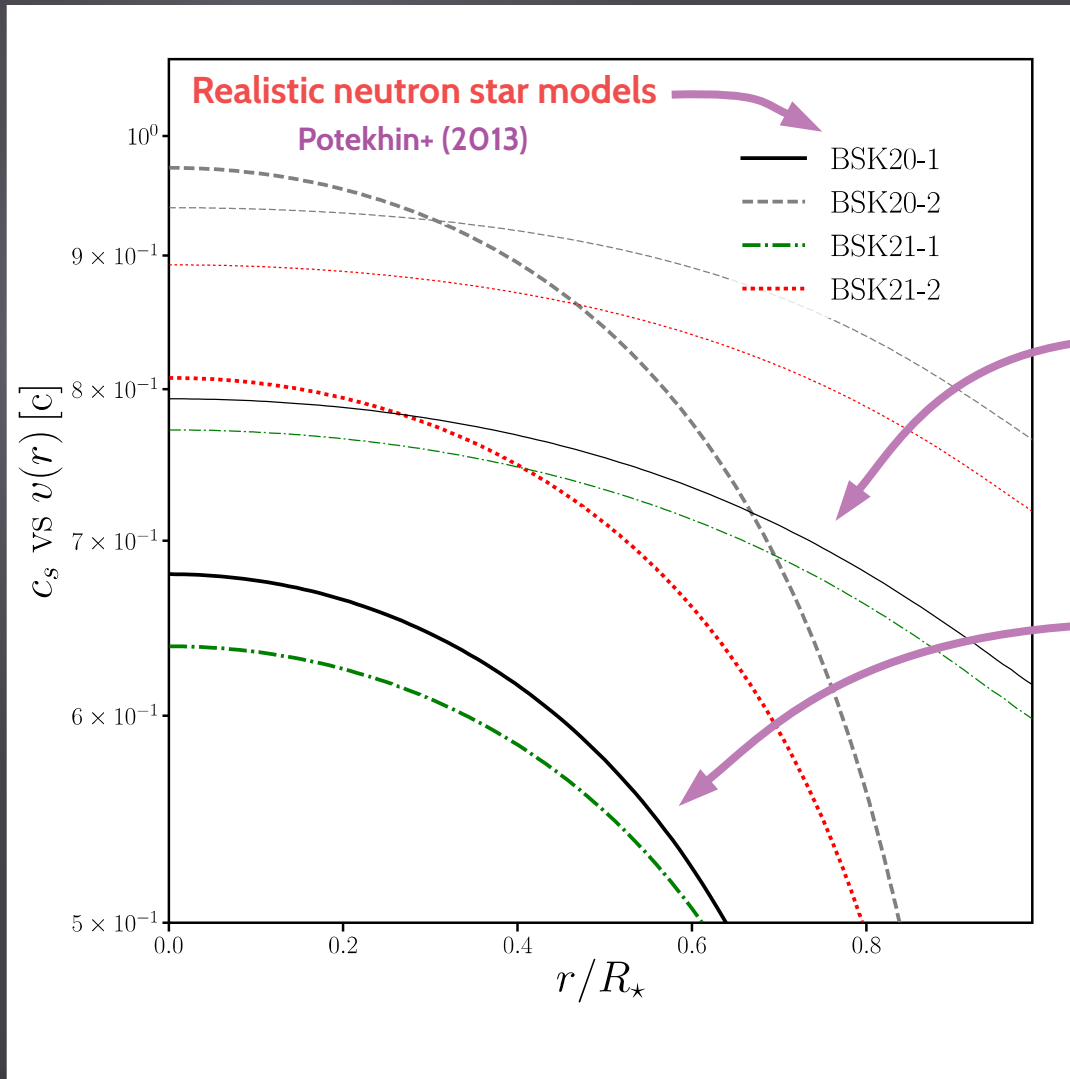


$$v(r) = \sqrt{1 - e^{2(\Phi(\infty) - \Phi(r))}}$$

$$c_s(r)$$

Capture happens in the supersonic regime!

# What is the speed regime for capture ?

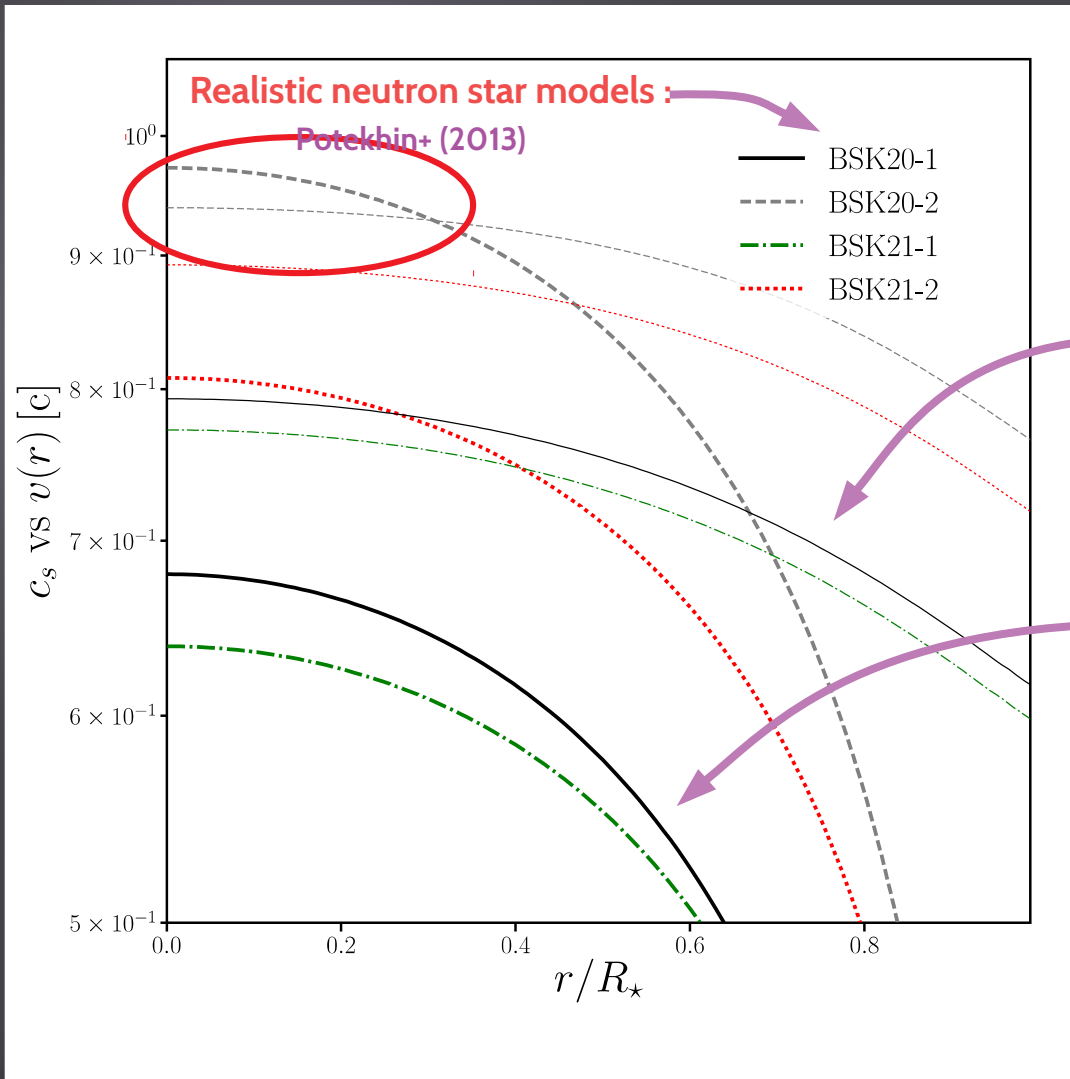


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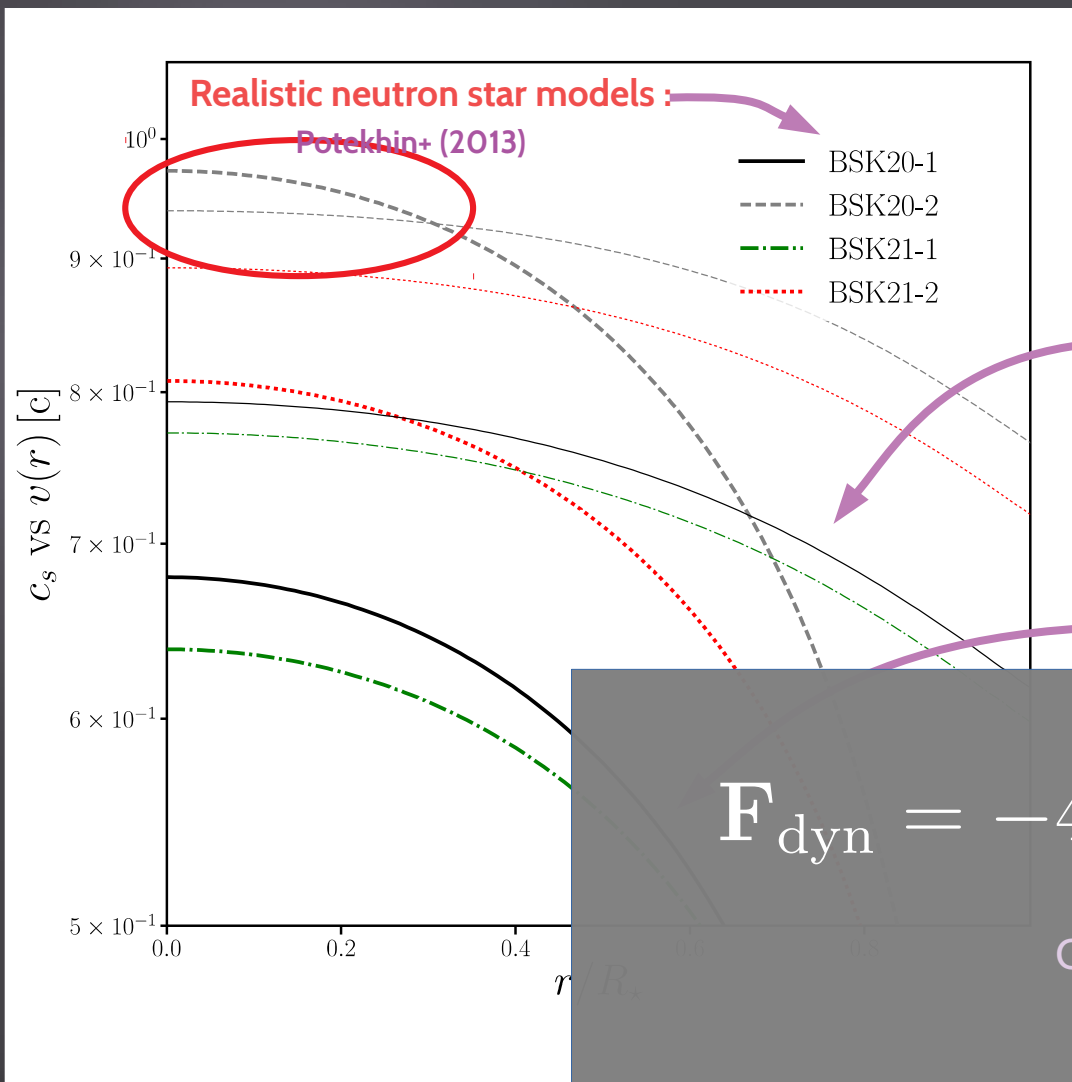
$$v(r) = \sqrt{1 - e^{2(\Phi(\infty) - \Phi(r))}}$$

$$c_s(r)$$

Capture can also happens in the subsonic regime!

But most of the time in the supersonic regime.

# What is the speed regime for capture ?



$$v(r) = \sqrt{1 - e^{2(\Phi(\infty) - \Phi(r))}}$$

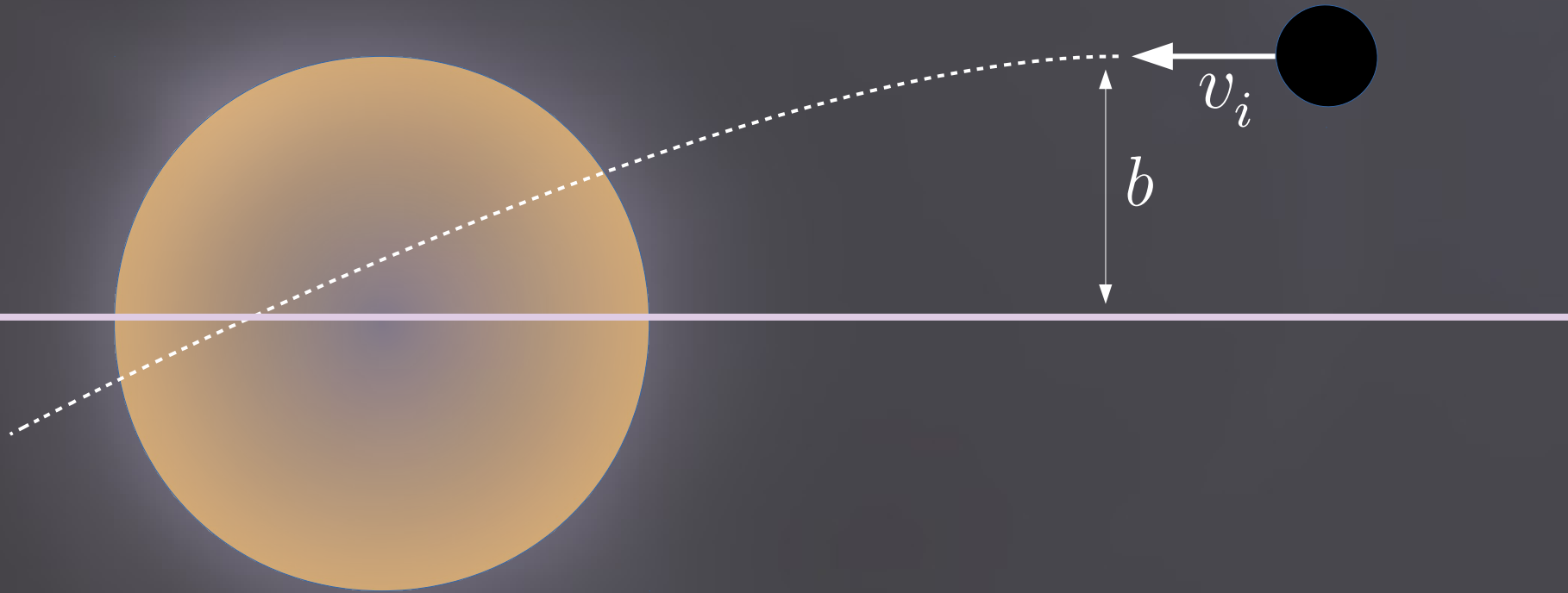
$$c_s(r)$$

$$\mathbf{F}_{\text{dyn}} = -4\pi G^2 m^2 \rho \ln \Lambda_{\text{dyn}}(v) \frac{\mathbf{v}}{v^3}$$

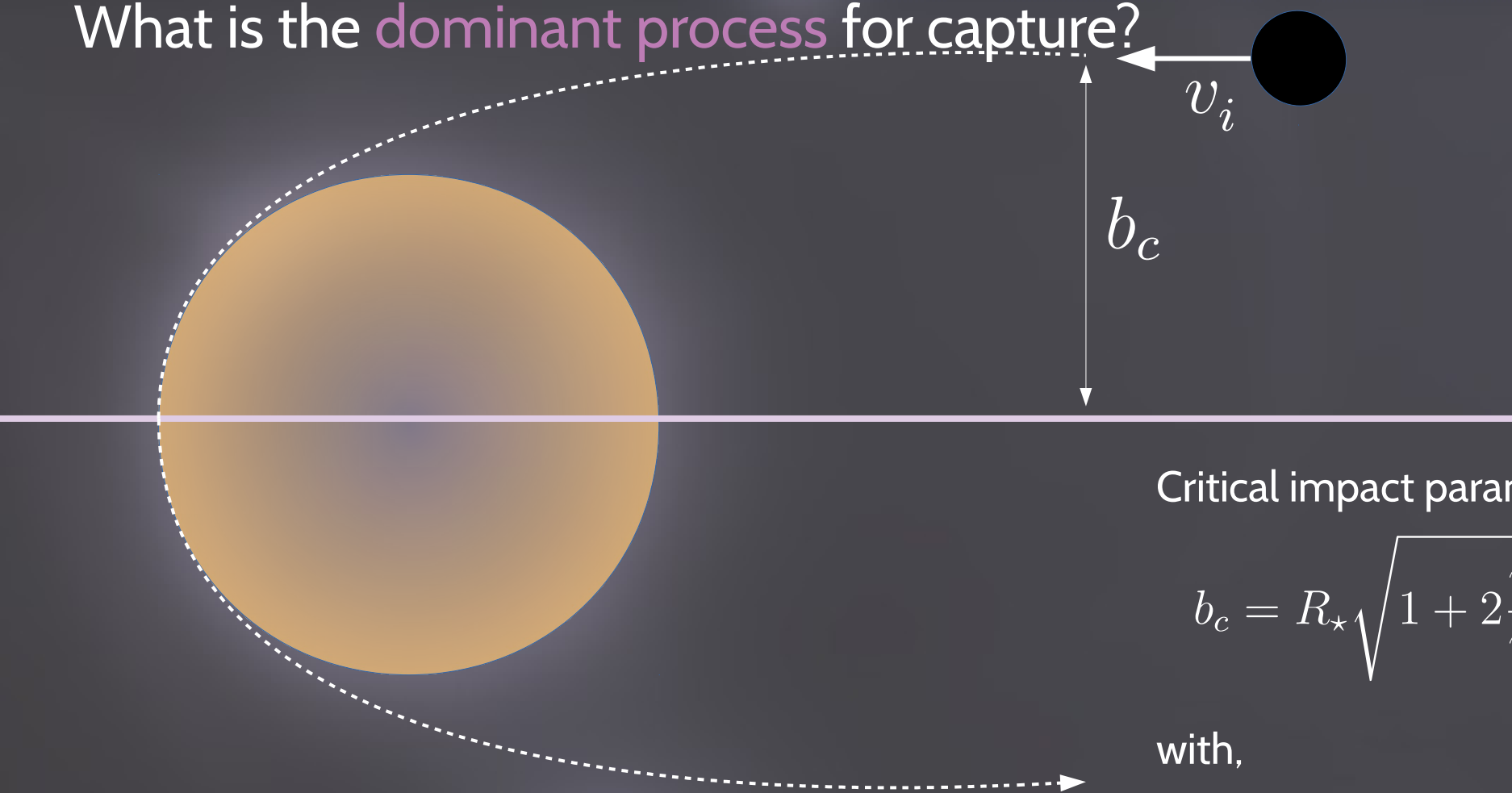
Chandrasekhar (1949)

most of the time in the supersonic regime

What is the **dominant process** for capture?



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Critical impact parameter

$$b_c = R_* \sqrt{1 + 2 \frac{v_*^2}{v_i^2}}$$

with,

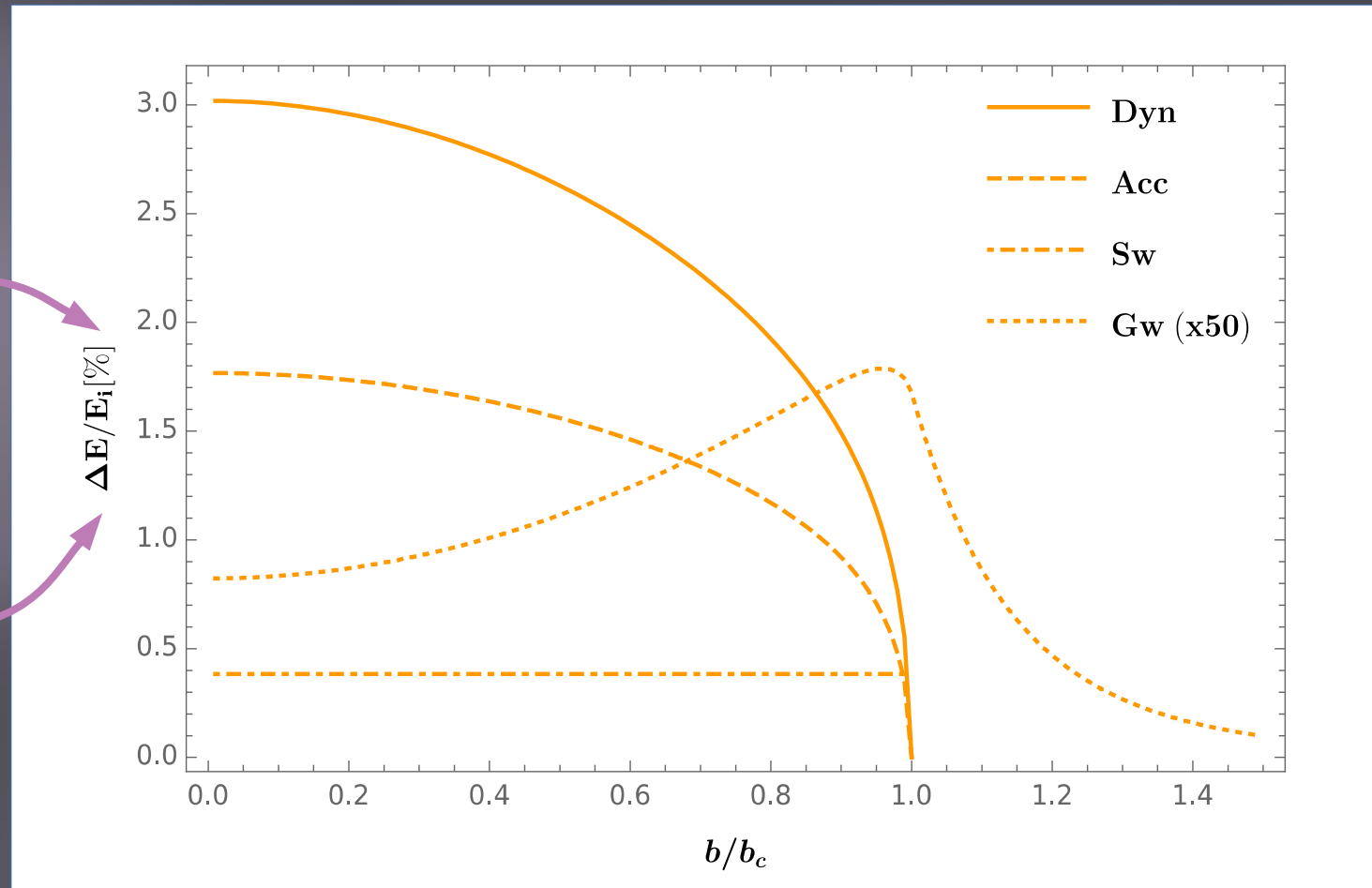
$$v_* = \sqrt{\frac{GM_*}{R_*}}$$

# What is the dominant process for capture?

$$E_i = \frac{1}{2} m v_i^2$$

(Here for  $v_i = 10^{-3}$ )

$$|\Delta E| = \int_c \mathbf{F} \cdot d\mathbf{l}$$



## Estimate of the number of event

The PBH distribution follows a Maxwellian in velocities

$$d^3n = n_{\text{PBH}} \left( \frac{3}{2\pi\bar{v}^2} \right)^{3/2} \exp \left\{ \frac{-3v^2}{2\bar{v}^2} \right\} d^3v,$$

Rate of NS-PBH encounter leading to capture

$$\mathcal{G}_\star = \int \frac{d^3n}{dv^3} \mathcal{S}(v) v d^3v \quad \text{with:} \quad \mathcal{S}(v) = \pi b_G^2$$



## Estimate of the number of event **in the Galaxy**

Rate of NS-PBH encounter leading to capture

$$N_{\star} \simeq 10^9$$

$$\mathcal{G}_{\star} N_{\star} \simeq 0.021 \left( \frac{\rho_{\text{PBH}}}{\text{GeV cm}^{-3}} \right) \left( \frac{10^{-3}}{\bar{v}} \right)^3 \mathcal{C} [X] \text{ Myr}^{-1}$$

with  $X = X(m, \bar{v}) \equiv \left( \frac{m}{10^{25} \text{g}} \right) \left( \frac{10^{-3}}{\bar{v}} \right)^2$

Within  $\tau_U = 10^{10} \text{yr}$ , few  $\sim 100$  of NS transmuted into BH.

Compare with the rate of NS-PBH encounter

$$\Gamma_{\star} \mathcal{N}_{\star} \simeq 0.38 \left( \frac{\rho_{\text{BH}}}{\text{GeV cm}^{-3}} \right) \left( \frac{10^{25} \text{g}}{m} \right) \left( \frac{10^{-3}}{\bar{v}} \right) \text{ Myr}^{-1}$$

Similar to the GRB rate in the Galaxy

## Estimate of the number of event **in the Galaxy**

Rate of NS-PBH encounter leading to capture

$$\mathcal{G}_* N_* \simeq 0.021 \left( \frac{\rho_{\text{PBH}}}{\text{GeV cm}^{-3}} \right) \left( \frac{10^{-3}}{\bar{v}} \right)^3 \mathcal{C} [X] \text{ Myr}^{-1}$$

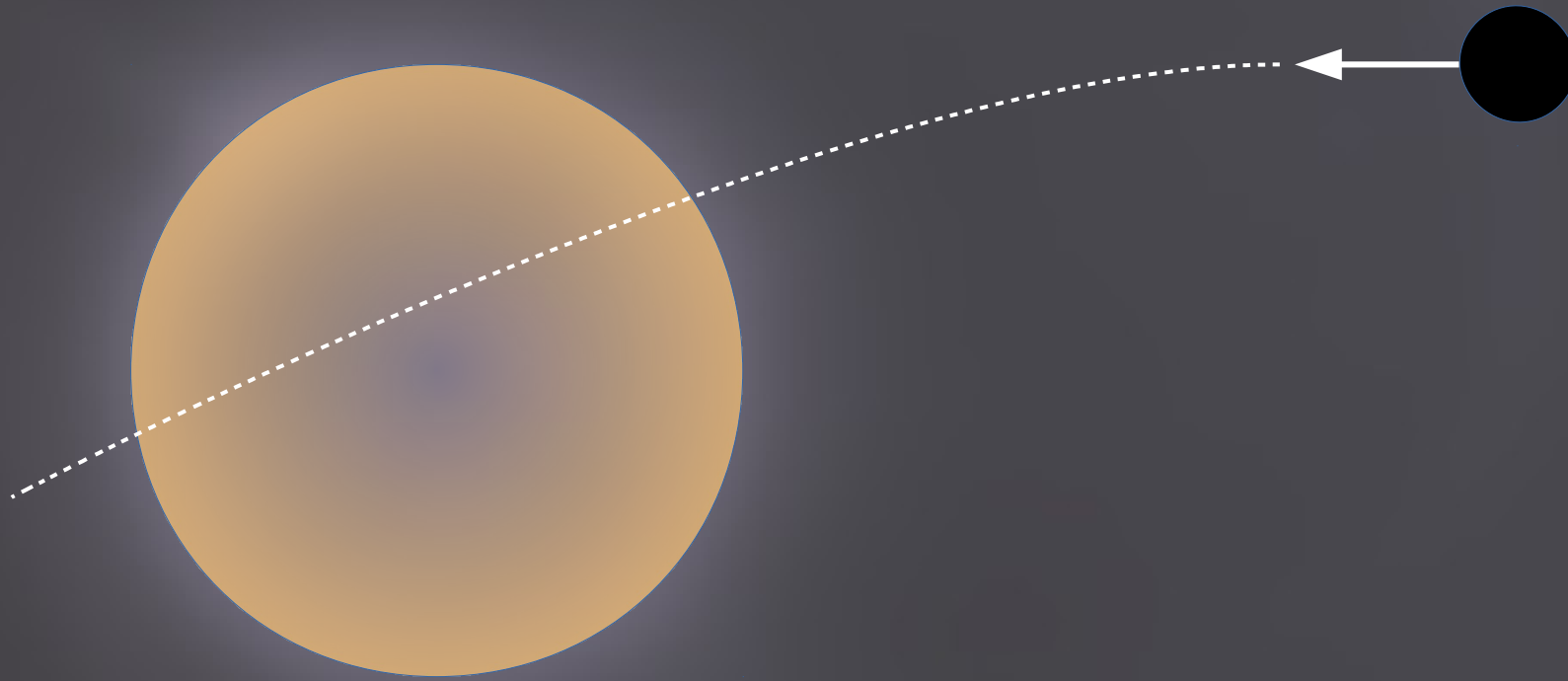
$$N_* \simeq 10^9$$

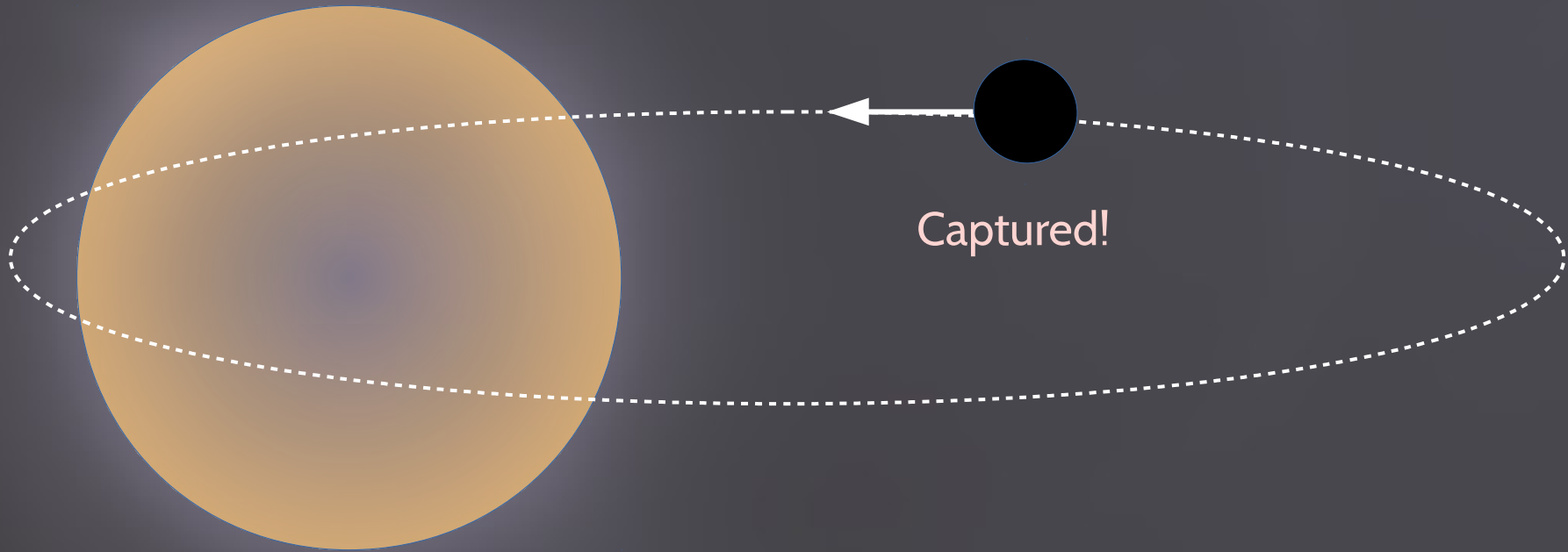
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 Rare events !

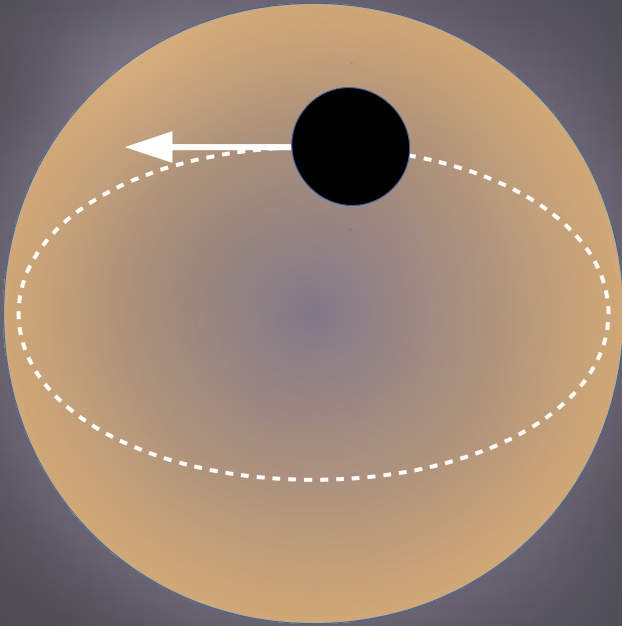




### Settling time within the NS

$$t_{\text{settle}} \lesssim 4 \times 10^4 \left( \frac{m}{10^{22} \text{ g}} \right)^{-3/2} \text{ yr}$$

Capela+ (2013)



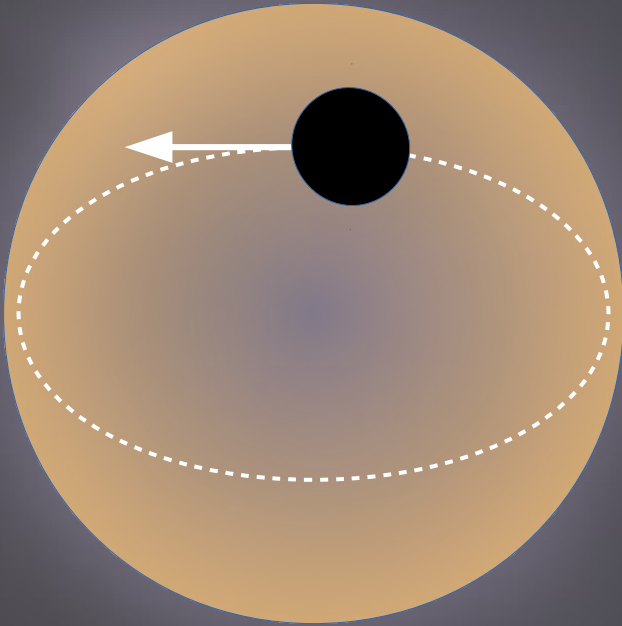
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The motion becomes subsonic for

$$r \lesssim R_{\star} \frac{c_s}{v_{\star}}$$



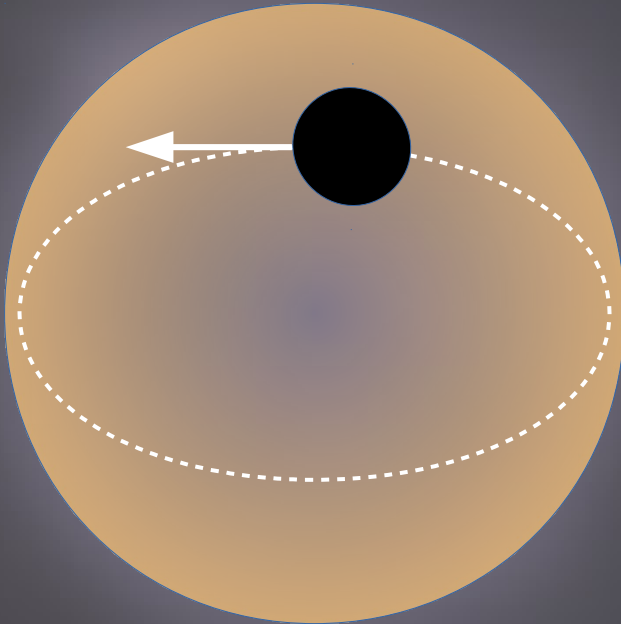
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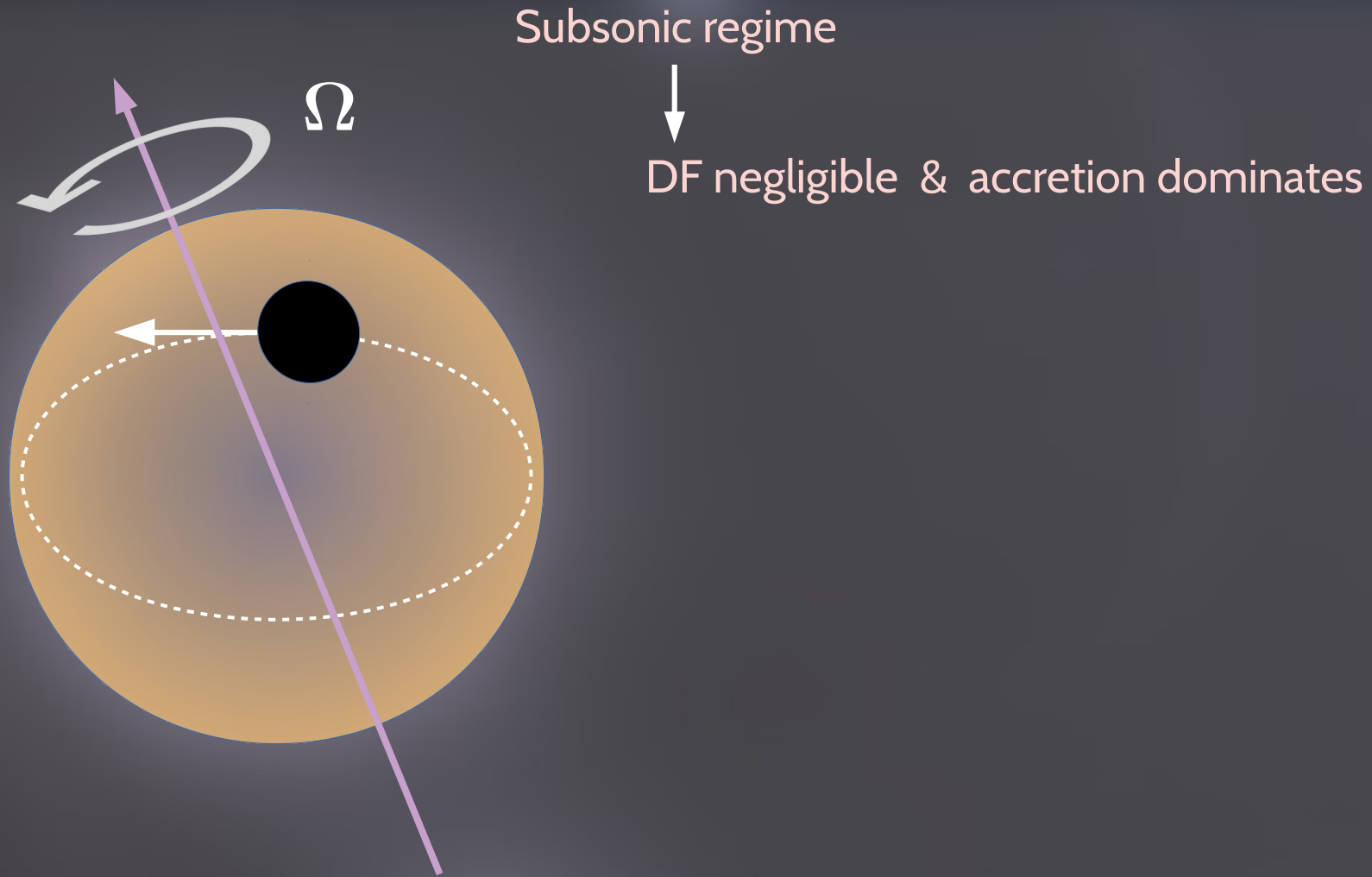
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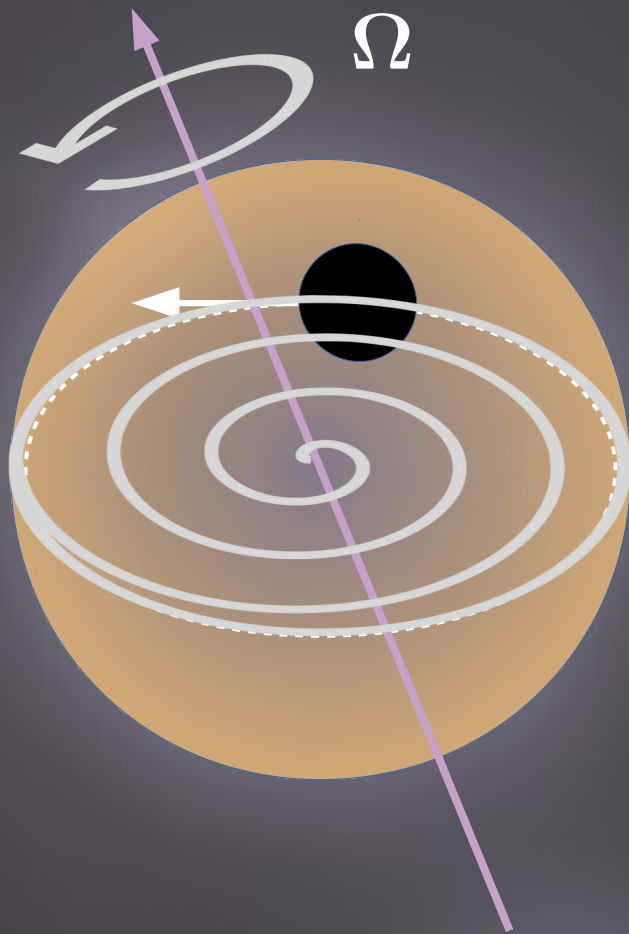


Model	BSK-20-1	BSK-20-2	BSK 21-1	BSK 21-2
Radius $R_{\star}$ [km]	11.6	10.7	12.5	12.0
Mass $M_{\star}$ [ $M_{\odot}$ ]	1.52	2.12	1.54	2.11
$v_{\star}$ [ $c$ ]	0.44	0.54	0.43	0.50
$f_{\star} = 1/T_{\star}$ [kHz]	1.8	2.4	1.6	2.0
$c_s$ (core) [ $c$ ]	0.68	0.97	0.64	0.81
$\mu_n$ (core) [GeV]	0.27	0.81	0.24	0.51

Realistic neutron star models Potekhin+ (2013)







Subsonic regime



DF negligible & accretion dominates

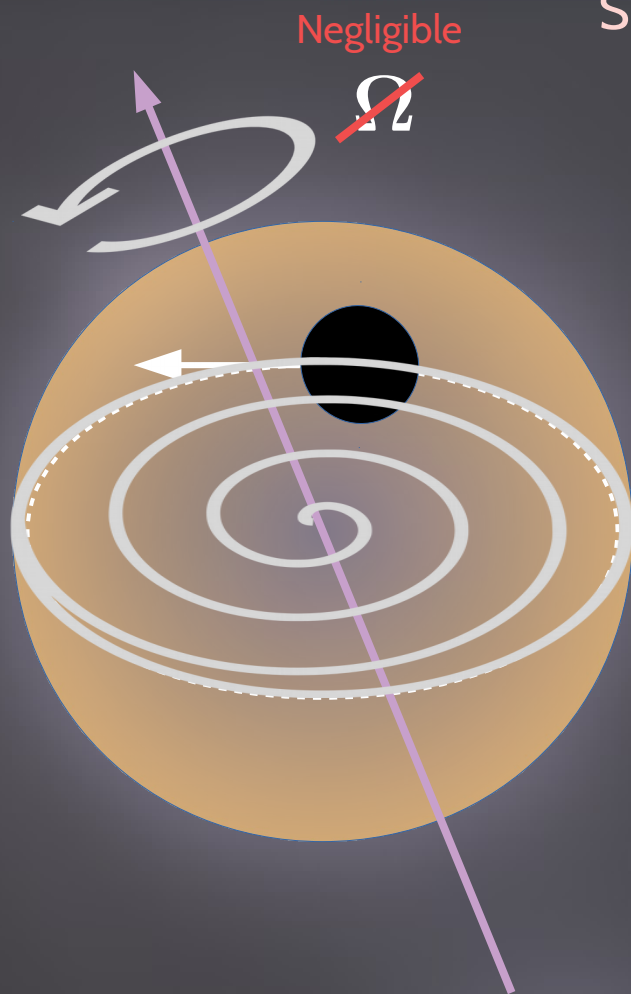


Equation of motion

$$\ddot{\mathbf{r}} + \mathcal{D}(t) [\dot{\mathbf{r}} - \boldsymbol{\Omega} \times \mathbf{r}] + \omega_{\star}^2 \mathbf{r} = 0$$

Y.G. et al. PRD (2020)

# PBH interactions with a NS – Post capture dynamic



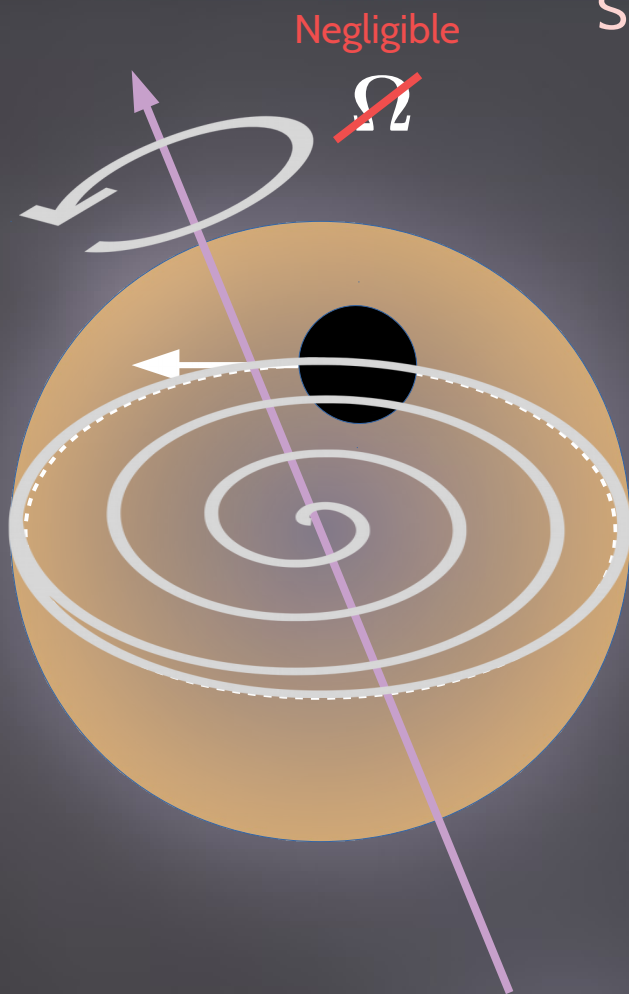
Subsonic regime

↓  
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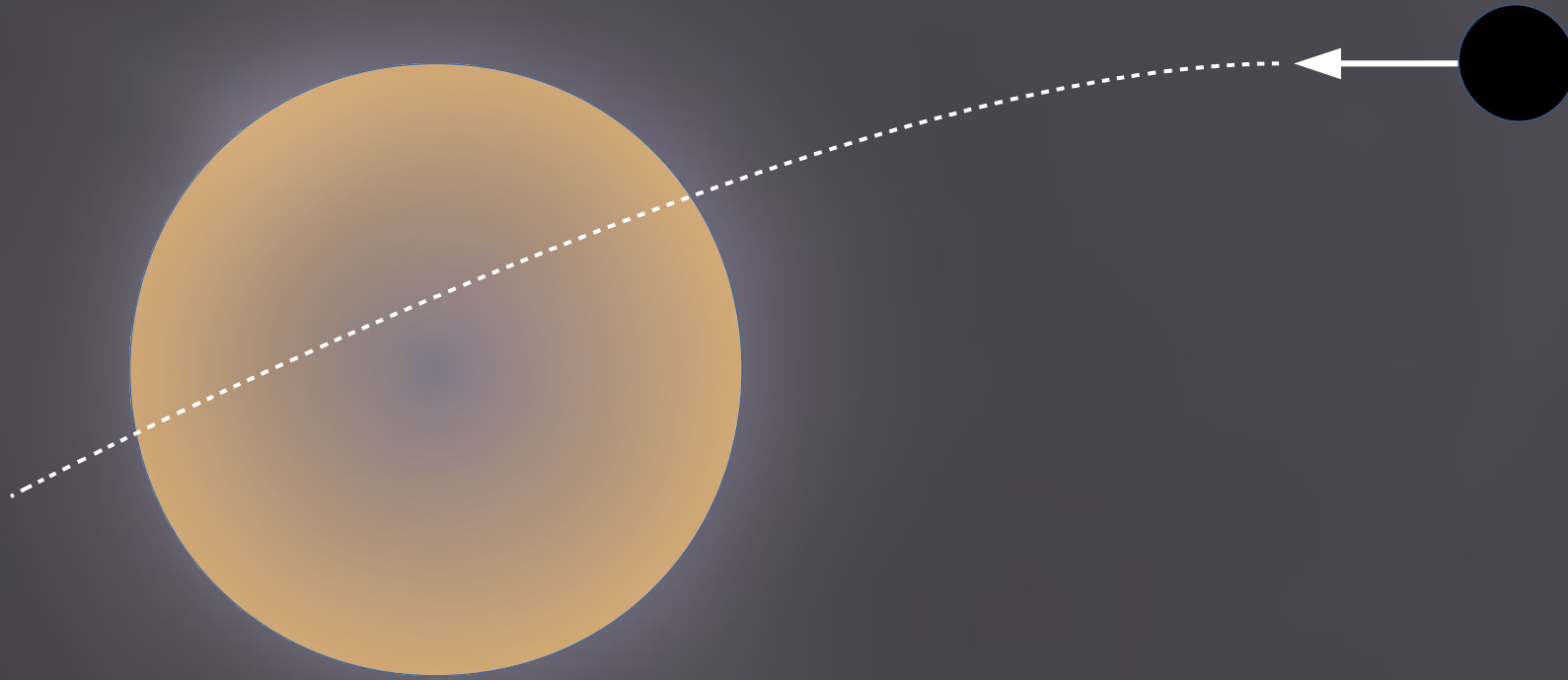
Y.G. et al. PRD (2020)

↓  
For  $\frac{\mathcal{D}}{\omega_{\star}} \sim 2.8 \times 10^{-12} \left( \frac{m}{10^{22} \text{g}} \right) \ll 1$

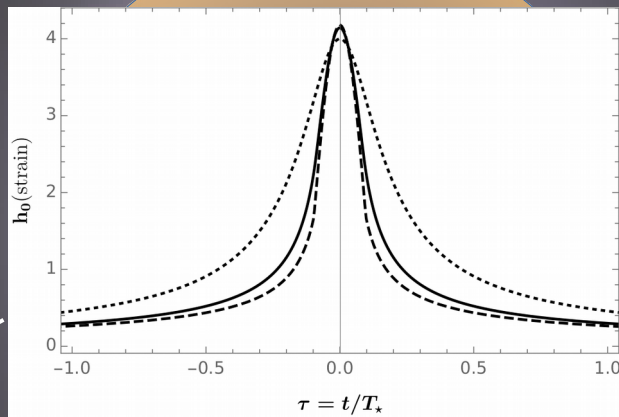
conserved quantity  $m r^2 = \text{const}$

whatever accretion regime

Signatures PBH – NS encounter



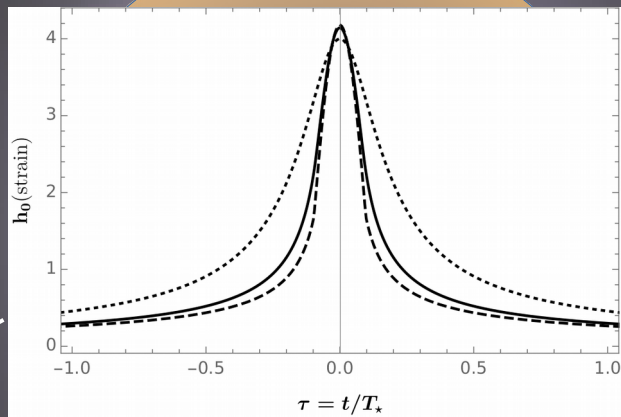
## Signatures PBH – NS encounter



→ Gravitational wave burst

$$h_0 \sim 10^{-25} \left( \frac{m}{10^{25} \text{ g}} \right) \left( \frac{1 \text{ kpc}}{d} \right)$$

## Signatures PBH – NS encounter



→ Gravitational wave burst

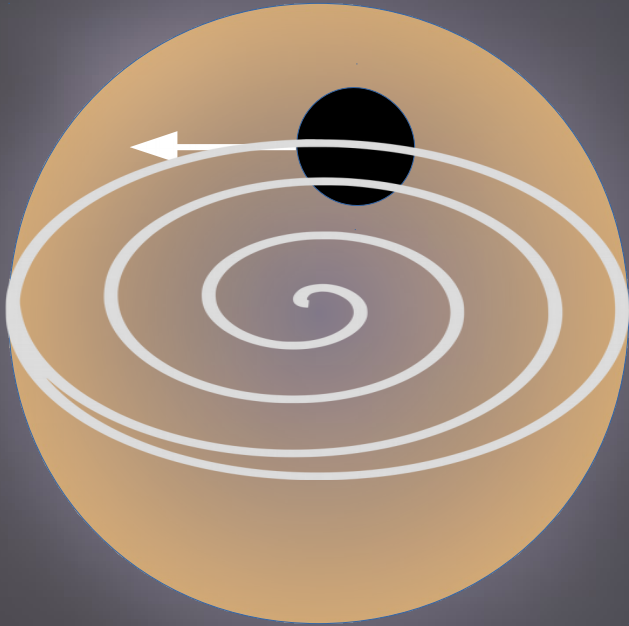
$$h_0 \sim 10^{-25} \left( \frac{m}{10^{25} \text{g}} \right) \left( \frac{1 \text{ kpc}}{d} \right)$$

→ Gravitational wave background

$$\sqrt{\langle h_c^2 \rangle} \simeq 3 \times 10^{-20} \left( \frac{10^{-10} \text{ Hz}}{f} \right)^2$$

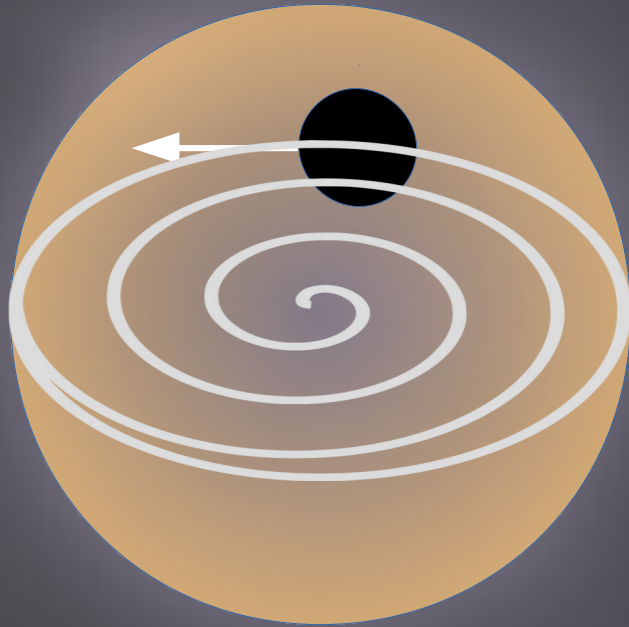
far below SKA sensitivity

Signatures captured PBH



## Signatures captured PBH

→ GW emission from the inspiral motion



$$h_0 = \frac{4\sqrt{2}G}{dc^4} m r^2 \omega_*^2 \approx 2.5 \times 10^{-25} \left( \frac{m}{10^{25}\text{g}} \right) \left( \frac{1 \text{ kpc}}{d} \right)$$

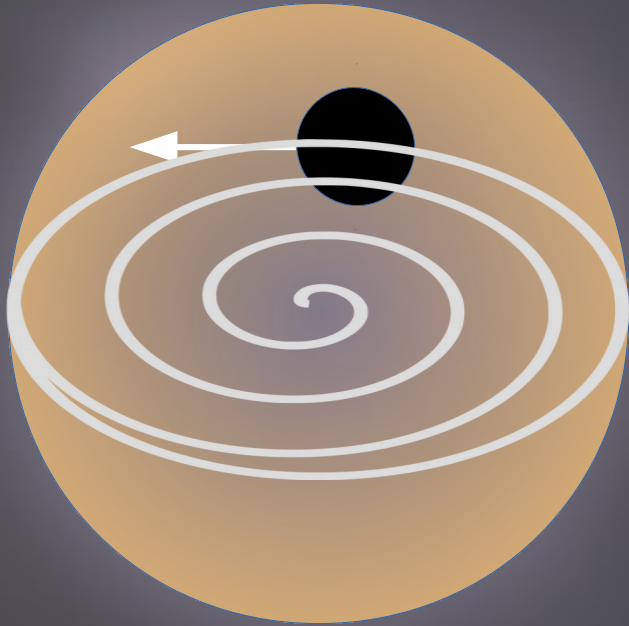
$$f_* \sim \text{kHz}$$

$$m r^2 = \text{const}$$



## Signatures captured PBH

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Emission sustained during the all accretion phase

$$t_B = \frac{c_s^3 R_*^3}{3G^2 M_* m} \approx 9 \left( \frac{10^{25}\text{g}}{m} \right) \text{ hours}$$

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→ GW emission from the inspiral motion

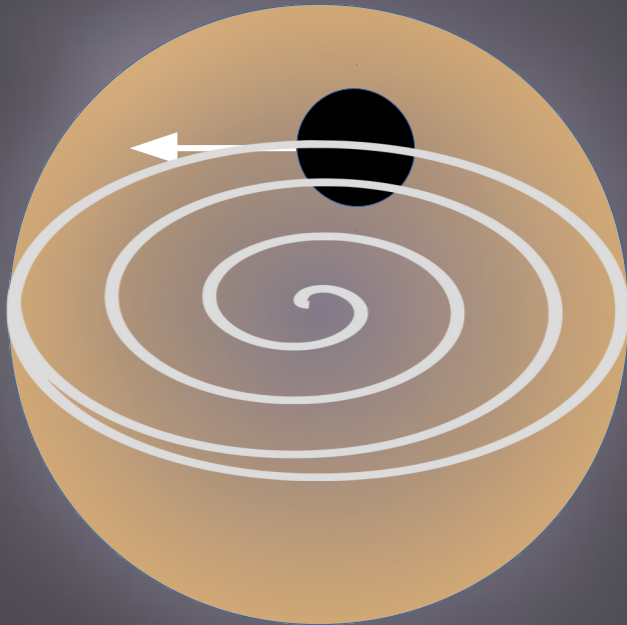
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→ Multiwavelength signature from the final collapse  
Might depend on the final asymmetry

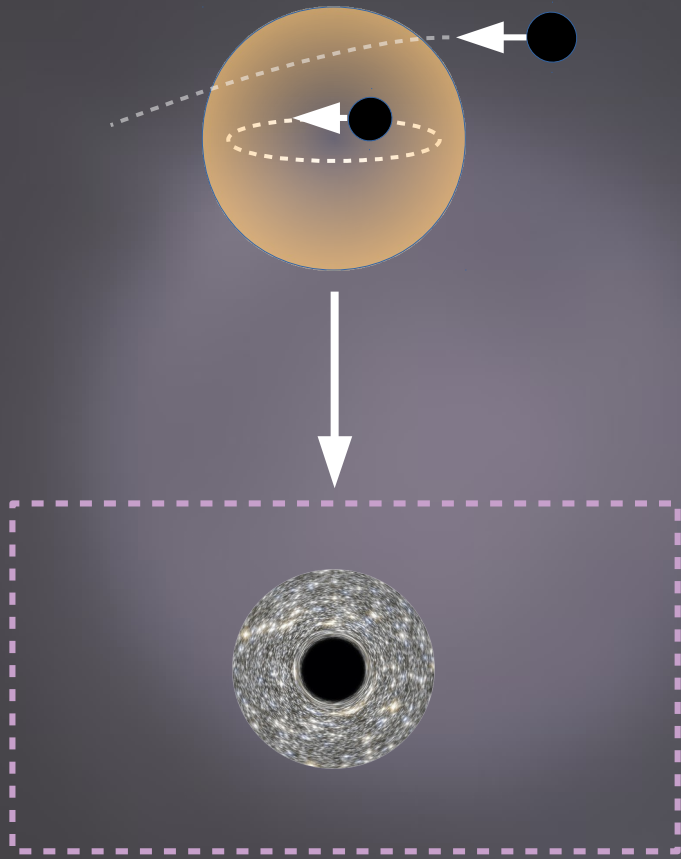


Final radius

$$R_f = R_* \sqrt{\frac{m_i}{f M_*}}$$

Initial PBH mass

Fraction of the star mass accreted



The collapse

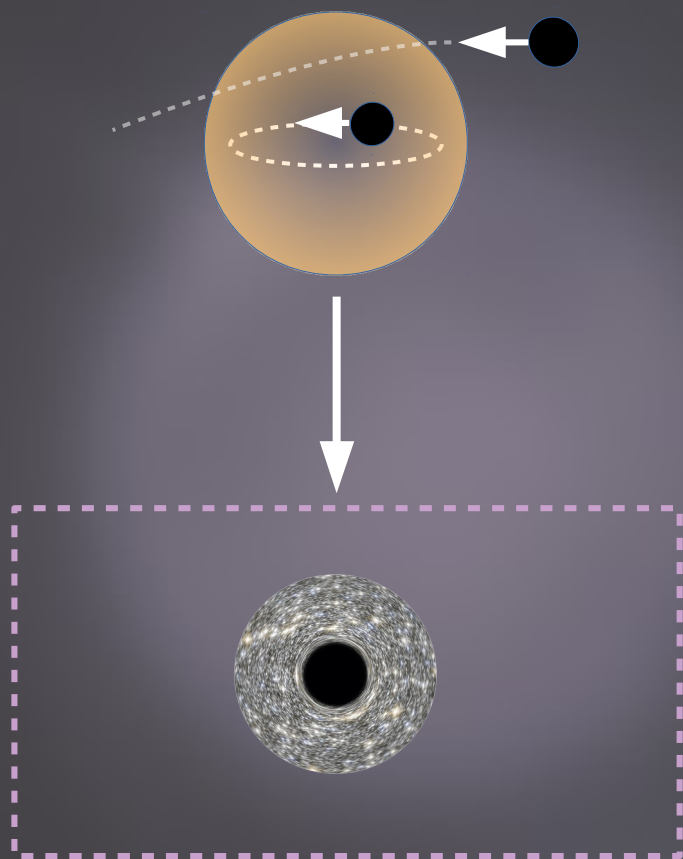
## Direct emissions

→ Electromagnetic waves: promising ! GRB? FRB?

No-hair theorem

$$E_B = \frac{B^2}{8\pi} \frac{4\pi}{3} R_{\star}^3 \simeq 2 \times 10^{41} \left( \frac{B}{10^{12} \text{G}} \right)^2 \left( \frac{R_{\star}}{10 \text{ km}} \right)^3 \text{ erg}$$

Fuller&Ott (2015), Abramowicz+ (2018), Chirenti+ (2019),...



The collapse

## Direct emissions

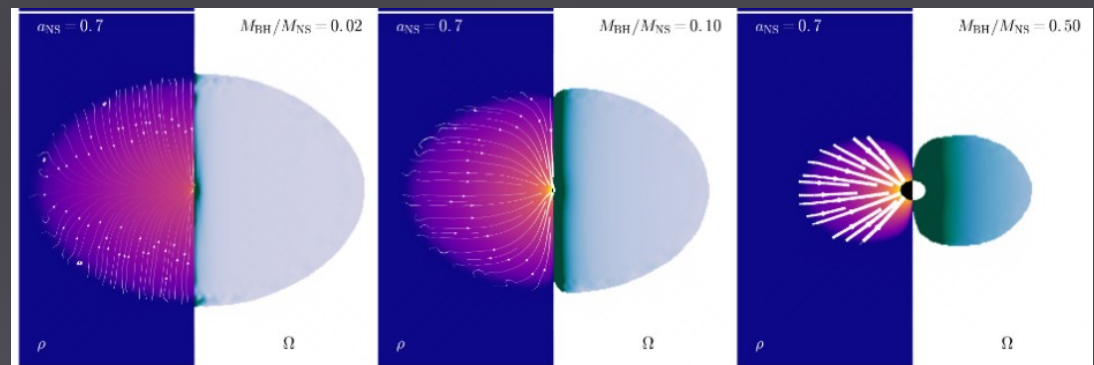
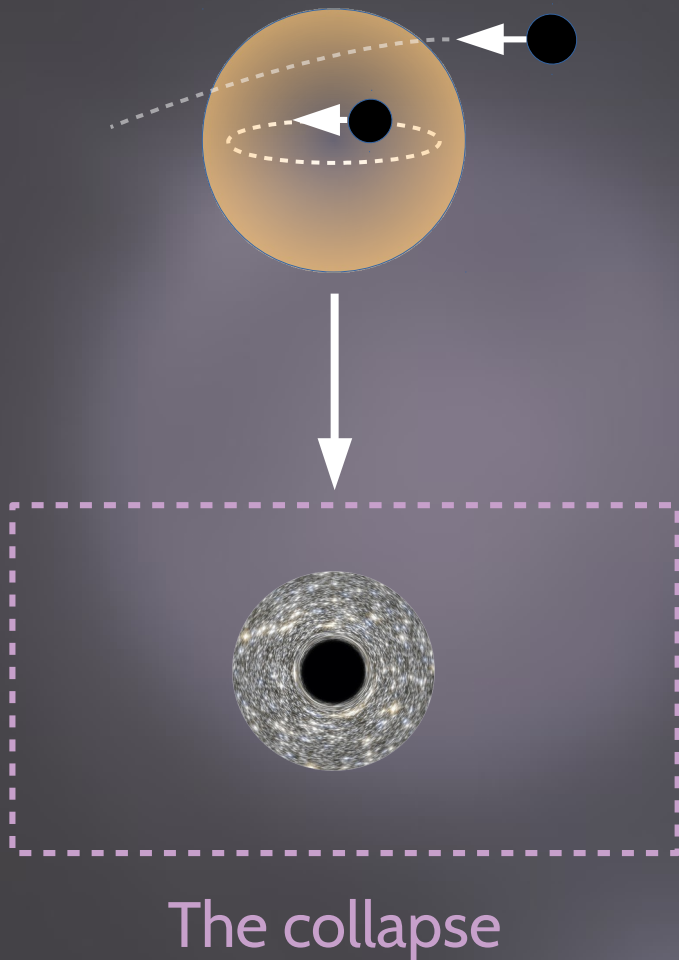
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Fuller&Ott (2015), Abramowicz+ (2018), Chirenti+ (2019),...

→ Gravitational waves: unpromising from simulations?



East+ (2019)

But PBH at the center and no magnetic field

## Direct emissions

→ Electromagnetic waves: promising ! GRB? FRB?

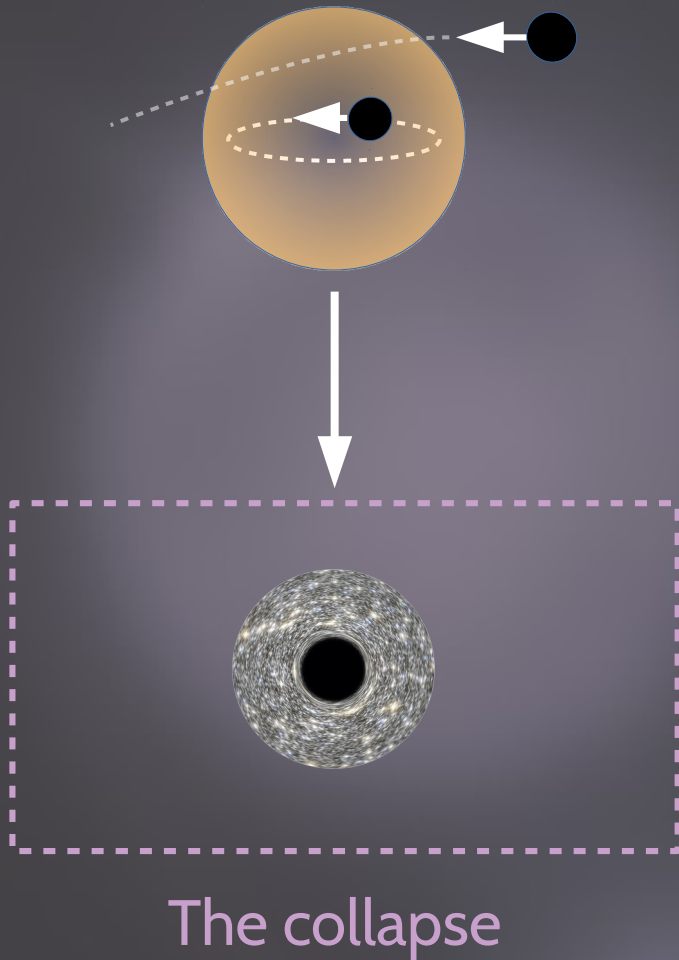
e.g. Fuller&Ott (2015), Abramowicz+ (2018), Chirenti+ (2019),...

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e.g. East+ (2019)

→ Observing quiet kilonovae?

e.g. Bramante+ (2016,2017)



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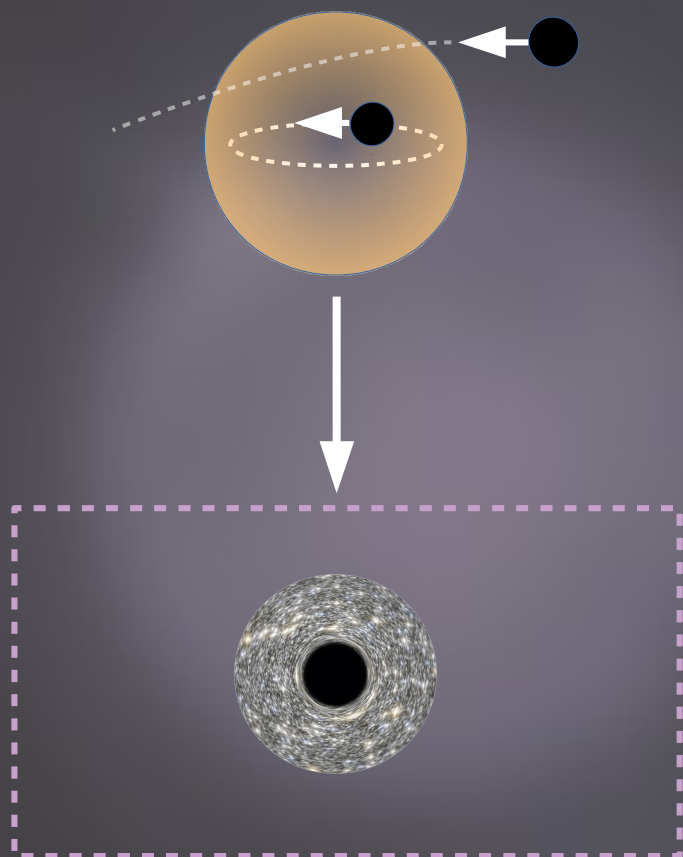
## Later detection

→ Leading mechanism for « light » BH formation?

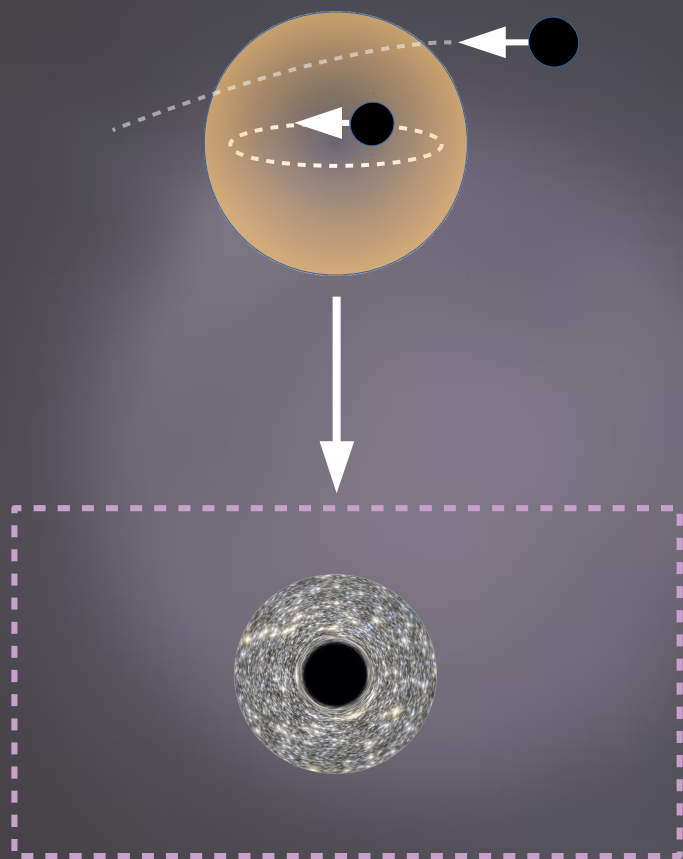
e.g. Takhistov+ (2021), Dexter+(2014)

→ Solving the missing pulsar problem?

e.g. Bramante+ (2016,2017)



The collapse



The collapse

## Direct emissions

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Thank you for listening!



Backups!

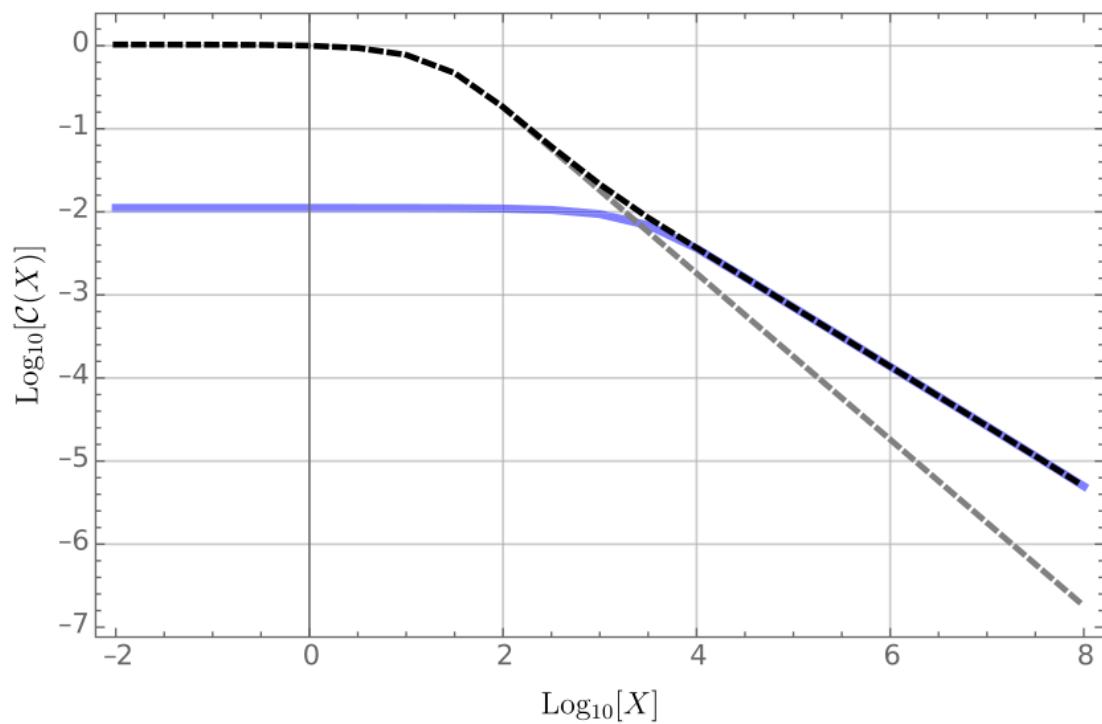


FIG. 3. Evolution of the function  $\mathcal{C}$  (black-dashed line) of Eq. (27) as a function of  $X$  defined in Eq. (28). The sole contribution of GW capture is displayed in blue and the difference with the total is shown with a dashed gray line.

