

Hawking radiation of primordial black holes

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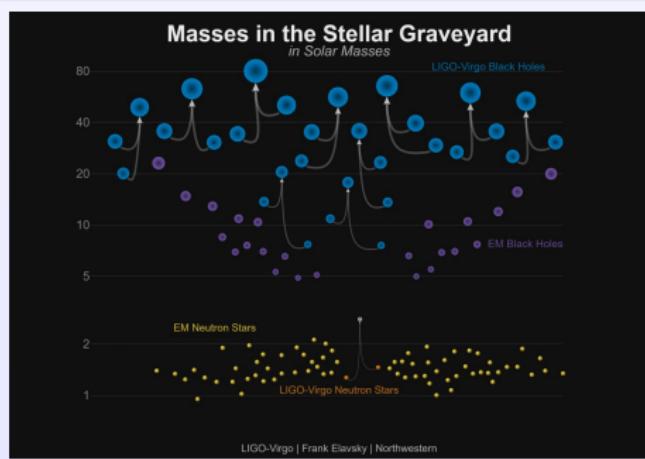
Institut Henri Poincaré – December 13th, 2021



Observed black holes

Three types of black holes have been discovered

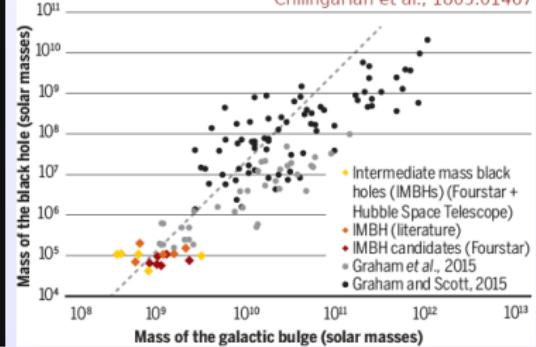
- Stellar black holes
BHs originated in the explosion of massive stars/supernovae, $\sim 3 - 100 M_{\odot}$
- Intermediate mass black holes (IMBH)
New class of recently discovered BHs, $\sim 10^3 - 10^6 M_{\odot}$
- supermassive black holes (SMBH)
BHs at the center of galaxies, $\sim 10^6 - 10^9 M_{\odot}$



Black hole growth chart

Black holes, including the newly discovered middleweights (color), have masses that correlate with the size of their host galaxy.

Chilingarian et al., 1805.01467



Origin of primordial black holes

Multiple inflationary origins

- collapse of large primordial overdensities
- phase transitions
- collapse of cosmic strings, domain walls

Mass predictions

Assuming that one PBH can be formed in a Hubble volume in the early Universe, one gets

$$M_{\text{PBH}} \sim M_{\text{Planck}} \times \frac{t_0}{t_{\text{Planck}}} \sim 10^{38} \text{ g} \times t_0(\text{s})$$

where t_0 is the creation time.

We get:

- $M \sim 10^{-5} \text{ g}$ for $t_0 \sim 10^{-43} \text{ s} \rightarrow$ Planck black holes
- $M \sim 10^{15} \text{ g}$ for $t_0 \sim 10^{-23} \text{ s} \rightarrow$ lightest black holes still (possibly) existing
- $M \sim 10^5 M_\odot$ for $t_0 \sim 1 \text{ s} \rightarrow$ IMHB? seeds for SMBH?

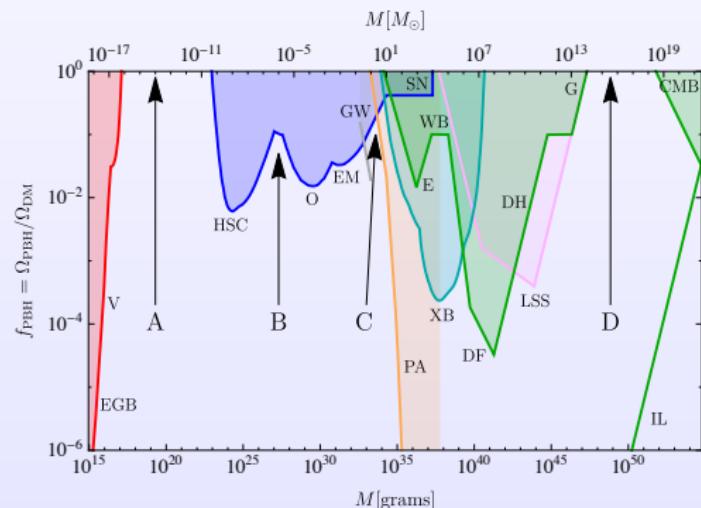


Constraints on Primordial Black Holes

Plausible dark matter candidates

- no need for Standard Model / General Relativity extension
- dynamically cold
- BH existence (somehow) proven
- mass ranges still available for BHs to represent all of dark matter

Constraints on PBHs – from Carr & Kuhnel, 2006.02838



red: evaporation
 blue: lensing
 gray: gravitational waves
 light blue: accretion
 orange: CMB distortions
 green: dynamical effects
 purple: large scale structure

A-D: possible open windows

Primordial black holes

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Hawking radiation

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Constraints

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Non-standard BHs

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BlackHawk

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Hawking radiation

Why are PBHs so special?

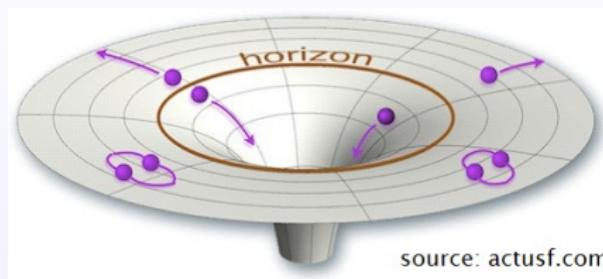
Light PBHs cannot be described only with General Relativity...



from B. Carr

... because they emit Hawking radiation and evaporate!

Black hole Hawking radiation



Fundamental equation for Kerr BHs

Rate of emission of Standard Model particles i at energy E by a BH of mass M and spin parameter a^* :

$$Q_i = \frac{d^2 N_i}{dt dE} = \frac{1}{2\pi} \sum_{\text{dof.}} \frac{\Gamma_i(M, E, a^*)}{e^{E'/T(M, a^*)} \pm 1}$$

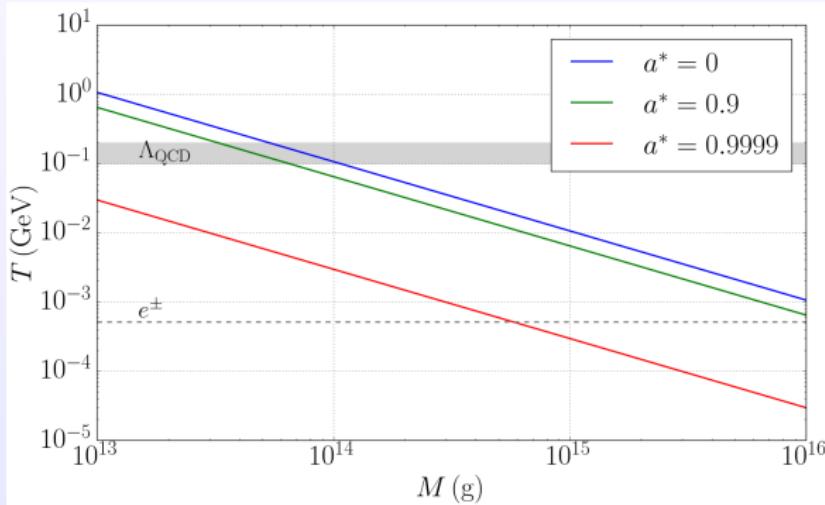
Γ_i is the greybody factor (\sim absorption coefficient in Planck's black-body law)

Hawking temperature

Hawking temperature for Kerr BHs

$$T(M, a^*) = \frac{1}{4\pi M} \left(\frac{\sqrt{1 - (a^*)^2}}{1 + \sqrt{1 - (a^*)^2}} \right) \xrightarrow[a^*=0]{\text{Schwarzschild}} \frac{1}{8\pi M}$$

Comparison with the e^\pm rest mass and QCD scale Λ_{QCD}



BHs get warmer when they lose energy via Hawking radiation.

Kerr Hawking radiation equations

Kerr metric

$$\begin{aligned} ds^2 = & \left(1 - \frac{2Mr}{\Sigma^2}\right) dt^2 + \frac{4a^* M^2 r \sin^2 \theta}{\Sigma^2} dt d\phi - \frac{\Sigma^2}{\Delta} dr^2 \\ & - \Sigma^2 d\theta^2 - \left(r^2 + (a^*)^2 M^2 + \frac{2(a^*)^2 M^3 r \sin^2 \theta}{\Sigma^2}\right) \sin^2 \theta d\phi^2 \end{aligned}$$

$$\Sigma \equiv r^2 + (a^*)^2 M^2 \cos^2 \theta \text{ and } \Delta \equiv r^2 - 2Mr + (a^*)^2 M^2$$

Equations of motion in free space

Dirac: $(i\not{\partial} - \mu)\psi = 0$ (fermions)

Proca: $(\square + \mu^2)\phi = 0$ (bosons)

μ = rest mass

Kerr Hawking radiation equations

Teukolsky radial equation

$$\frac{1}{\Delta^s} \frac{d}{dr} \left(\Delta^{s+1} \frac{dR}{dr} \right) + \left(\frac{K^2 + 2is(r-M)K}{\Delta} - 4isEr - \lambda_{slm} - \mu^2 r^2 \right) R = 0$$

R radial component of ψ/ϕ

$K \equiv (r^2 + a^2)E + am$, s = spin, I = angular momentum and m = projection

Transformation into a Schrödinger equation

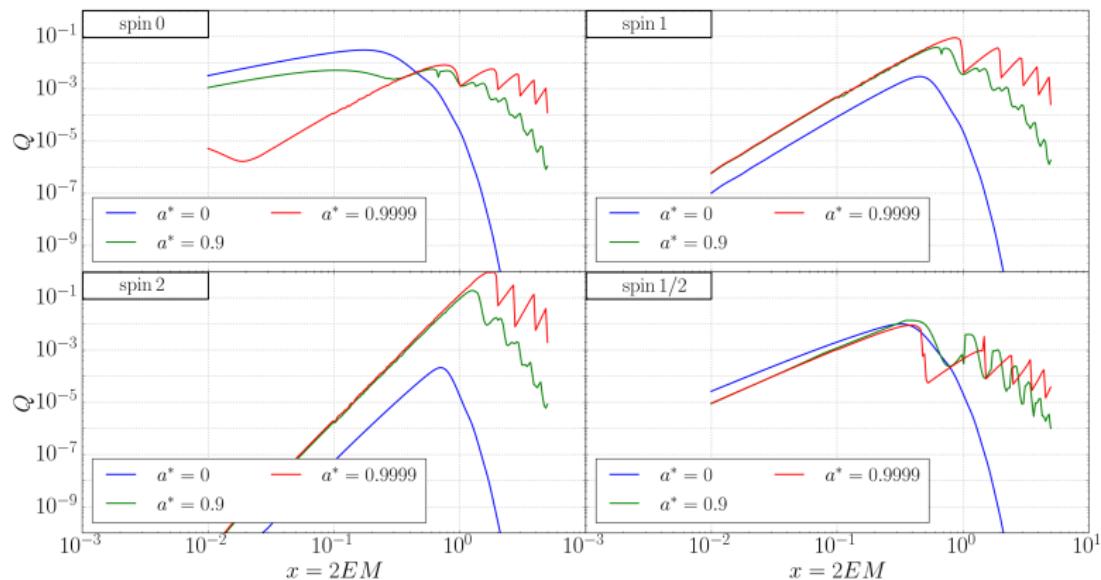
Change $\psi/\phi \rightarrow Z$ and $r \rightarrow r^*$ (generalized Eddington-Finkelstein coordinate system)
(Chandrasekhar & Detweiler 1970s)

$$\frac{d^2 Z}{dr^{*2}} + (E^2 - V(r^*))Z = 0$$

Solved with purely outgoing solution $Z \xrightarrow{r^* \rightarrow -\infty} e^{-iEr^*}$

Transmission coefficient $\Gamma \equiv |Z_{\text{out}}^{+\infty}/Z_{\text{out}}^{\text{horizon}}|^2$

Hawking radiation of particles



All particles can be emitted by a black hole!

Including gravitons / gravitational waves... and even new physics particles!

Hawking radiation is enhanced for particles of spin 1 or 2.

Black hole lifetime

Evolution equations

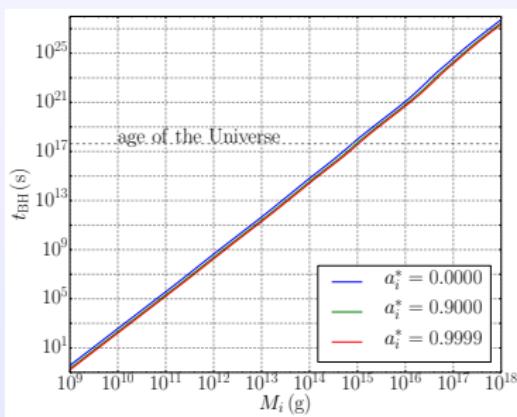
$$\frac{dM}{dt} = -\frac{f(M, a^*)}{M^2}$$

$$\frac{da^*}{dt} = \frac{a^*(2f(M, a^*) - g(M, a^*))}{M^3}$$

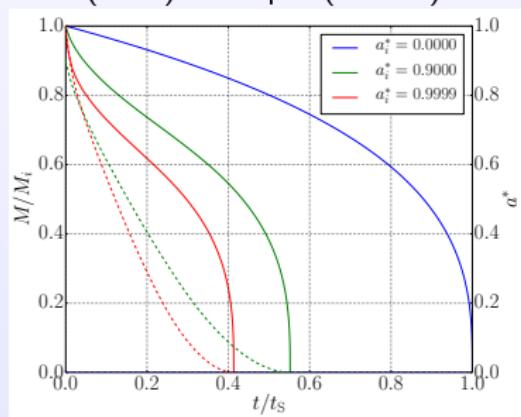
$$f \sim \int_E \text{ener.} \times \text{emiss.}$$

$$g \sim \int_E \text{ang. mom.} \times \text{emiss.}$$

BH lifetime



BH mass (solid) and spin (dotted) evolution



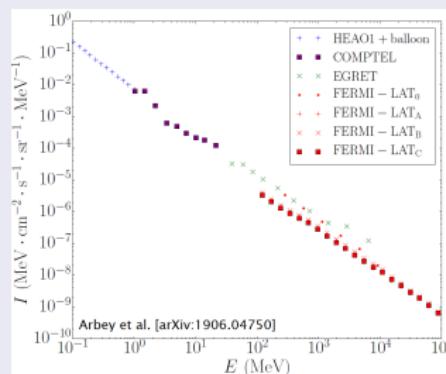
Constraints on primordial black holes

Isotropic gamma ray background (IGRB) constraints

Origin

Diffuse background +

- Active galactic nuclei
- Gamma ray bursts
- DM annihilation/decay?
- Hawking radiation?



Flux estimation for BHs

AA et al., arXiv:1906.04750

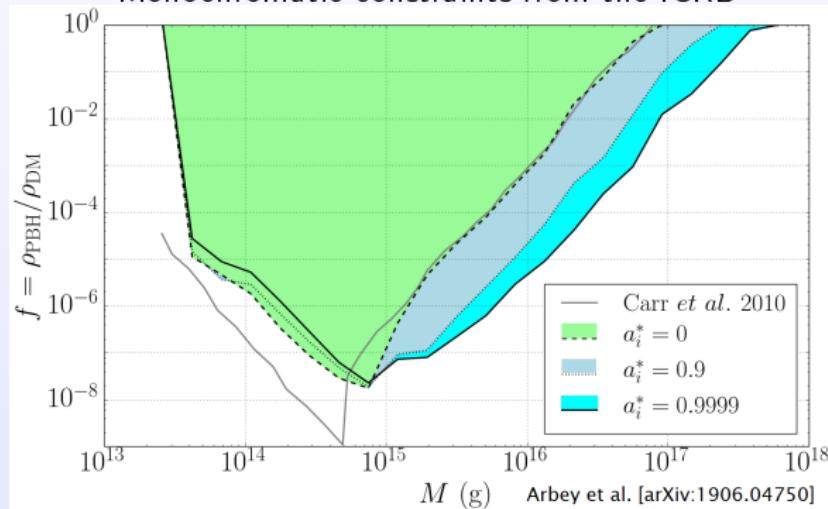
$$\begin{aligned} I &\approx \frac{1}{4\pi} E \int_{t_{\text{CMB}}}^{t_{\text{today}}} (1+z(t)) \\ &\quad \times \int_M \left[\frac{dn}{dM} \frac{d^2N}{dt dE} (M, (1+z(t))E) dM \right] dt \end{aligned}$$

IGRB and Kerr PBHs: monochromatic mass distributions

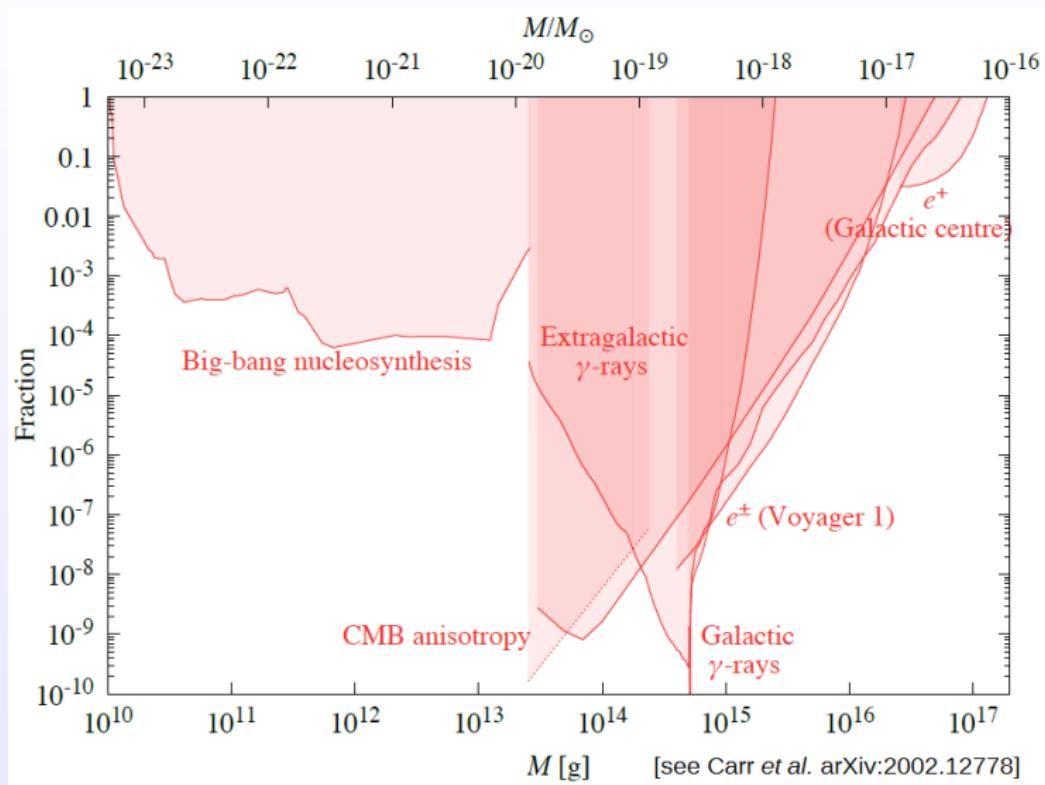
Main spin effects

- enhanced luminosity \Rightarrow stronger constraints
- reduced temperature \Rightarrow reduced emission energy \Rightarrow weaker constraints

Monochromatic constraints from the IGRB



Complementary constraints



Primordial black holes

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Hawking radiation

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Constraints

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Non-standard BHs

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BlackHawk

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Non-standard black holes

The case of static spherically-symmetric metrics

AA, J. Auffinger, M. Geiller, E. Livine, F. Sartini, arXiv:2101.02951 + 2107.03293

$$ds^2 = -G(r)dt^2 + \frac{1}{F(r)}dr^2 + H(r)d\Omega^2$$

Equations of motion:

$$\partial_{r^*}^2 Z + (\omega^2 - V(r(r^*)))Z = 0$$

where $dr^*/dr \equiv 1/\sqrt{FG}$ and the spin-dependent potentials V are related to F, G, H .

Transmission coefficient: $\Gamma_i \equiv |Z_{\text{out}}^{+\infty}/Z_{\text{out}}^{\text{horizon}}|^2$

Rate of emission of one degree of freedom i per unit time t and energy E :

$$\frac{d^2N_i}{dt dE} = \sum_{l,m} \frac{1}{2\pi} \frac{\Gamma_i(E, M, x_j)}{e^{E/T} - (-1)^{2s_i}},$$

where s_i is the spin of the particle i and T is its Hawking temperature given by

$$T = \frac{1}{4\pi} \left. \frac{F^{1/2}G'}{G^{1/2}} \right|_{\text{hor}}$$

where “hor” denotes the horizon $r = r_H$.

Extra-dimensional PBHs

Metric:

$$F(r) = G(r) \equiv 1 - \left(\frac{r_H}{r}\right)^{n+1}, \quad H(r) = r^2$$

where $n > 0$ is the number of extra dimensions.

Horizon radius is given by

$$r_H = \frac{1}{\sqrt{\pi} M_*} \left(\frac{M}{M_*} \right)^{1/(n+1)} \left(\frac{8\Gamma((n+3)/2)}{n+2} \right)^{1/(n+1)}$$

where Γ is the Euler gamma function and the rescaled Planck mass is

$$M_{\text{Pl}}^2 = M_*^{n+2} R^n$$

with R the radius of the extra-dimension(s).

Temperature:

$$T_n = \frac{\kappa_n}{2\pi} = \frac{n+1}{4\pi r_H}$$

Remark: Schwarzschild metric retrieved for $n = 0$.



Extra-dimensional PBHs

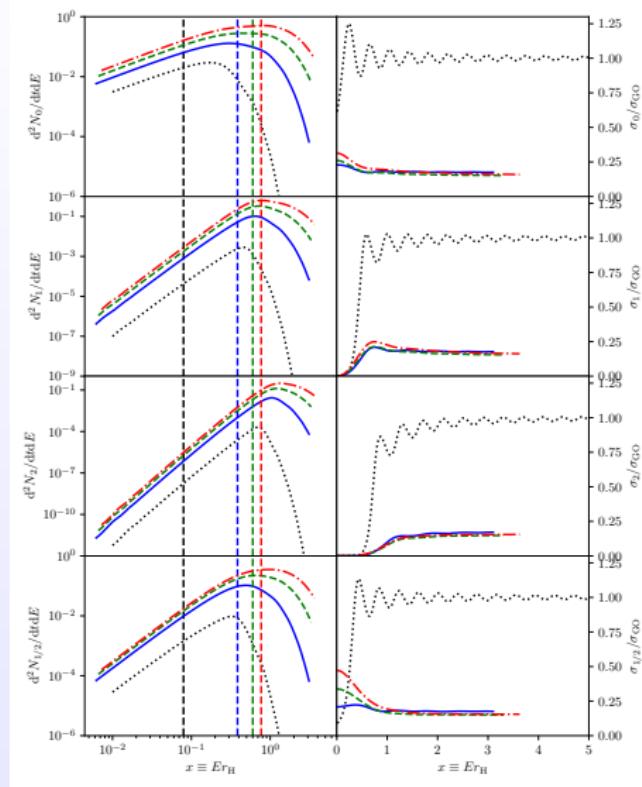
Hawking radiation spectra

dotted black: Schwarzschild

blue: $n = 2$

green: $n = 4$

red: $n = 6$



LQG-inspired PBHs

Schwarzschild BHs are singular at the center → non-physical → quantum gravity?

Loop Quantum gravity (LQG): structure of spacetime described by spin networks.

Metric inspired by LQG, regularising the naked singularity at the center of the Schwarzschild metric:

$$G = \frac{(r - r_+)(r - r_-)(r + \sqrt{r_+ r_-})^2}{r^4 + a_0^2}, \quad F = \frac{(r - r_+)(r - r_-)r^4}{(r + \sqrt{r_+ r_-})^2(r^4 + a_0^2)}, \quad H = r^2 + \frac{a_0^2}{r^2}.$$

Two roots: $r_+ = \frac{2M}{(1+P)^2}$ and $r_- = \frac{2MP(\varepsilon)^2}{(1+P)^2}$

M : (ADM) mass

$P(\varepsilon) = \frac{\sqrt{1+\varepsilon^2} - 1}{\sqrt{1+\varepsilon^2} + 1}$: polymerization factor

a_0 : minimal area in loop quantum gravity, typically of the Planck scale

$\varepsilon \geq 0$: typical scale of the geometry fluctuations (a priori independent parameter)

Temperature: $T_{\text{LQG}} = \frac{r_+^2(r_+ - r_-)}{4\pi(r_+^4 + a_0^2)}$

LQG-inspired PBHs

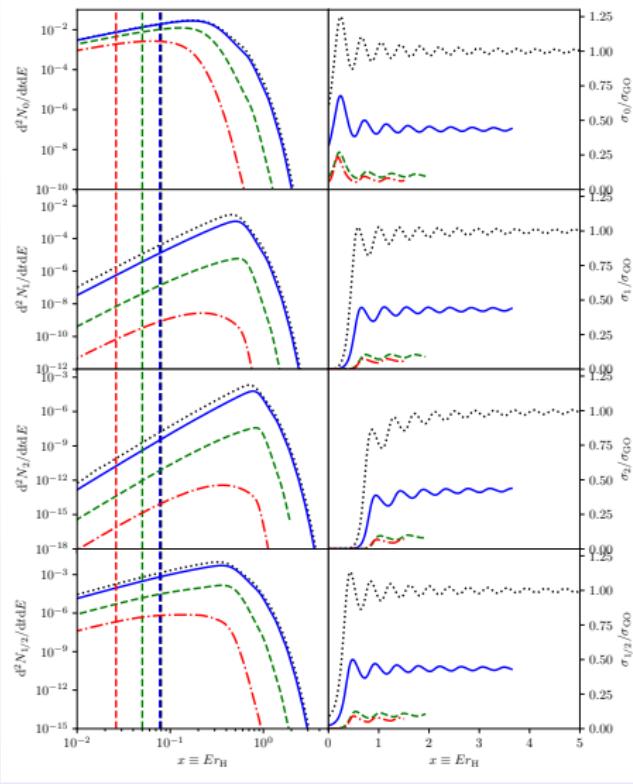
Hawking radiation spectra

dotted black: Schwarzschild

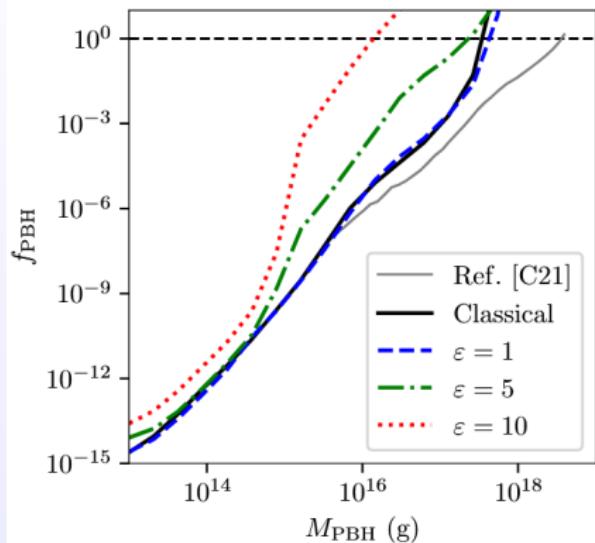
blue: $\epsilon = 1$

green: $\epsilon = 4$

red: $\epsilon = 10$



LQG-inspired PBHs



Predictions for the sensitivity of the AMEGO gamma-ray experiment to classical and polymerised LQG-inspired PBHs

ϵ : scale of deformation of spacetime due to polymerisation in Loop Quantum Gravity

AA, J. Auffinger, M. Geiller, E. Livine, F. Sartini, arXiv:2101.02951 + 2107.03293

→ Theories beyond General Relativity predict very different signals of Hawking radiation.

BlackHawk

Public C code computing Hawking radiation:

- Schwarzschild & Kerr PBHs + non-standard BHs
- primary spectra of all Standard Model fundamental particles
- secondary spectra of stable particles
- extended mass and spin functions
- time evolution of the PBHs

Download: <http://blackhawk.hepforge.org>

Manual: [arXiv:1905.04268](#), [Eur.Phys.J. C79, 693](#)
+ [arXiv:2108.02737](#), [Eur.Phys.J. C81, 10](#)

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BlackHawk

By Alexandre Arbey and Jérémie Auffinger

Calculation of the Hawking evaporation spectra of any black hole distribution

BlackHawk is a public C program for calculating the Hawking evaporation spectra of any black hole distribution. This program enables the users to compute the primary and secondary spectra of stable or long-lived particles generated by Hawking radiation of the distribution of black holes, and to study their evolution in time.

If you use BlackHawk to publish a paper, please cite:

A. Arbey and J. Auffinger, arXiv:1905.04268 [gr-qc]

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