

Three-body effects in waveforms

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In collaboration with : Enrico Trincherini, Francesco Serra

SCUOLA
NORMALE
SUPERIORE



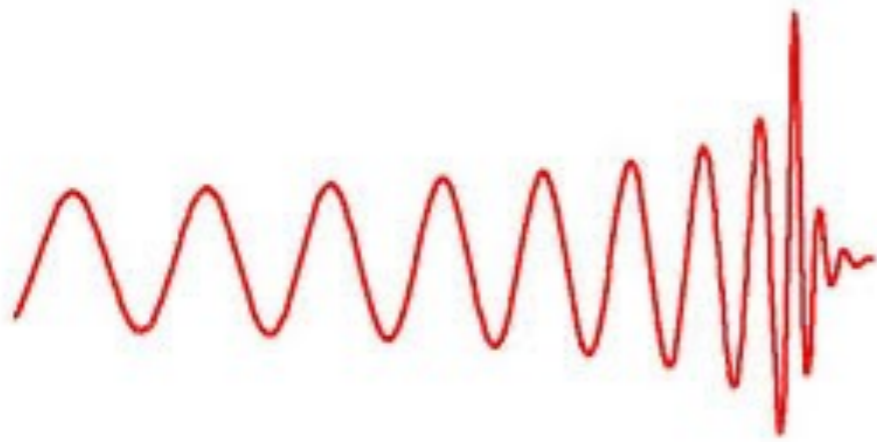
13/12/2021



Istituto Nazionale di Fisica Nucleare

INTRODUCTION

Detection of GW beautifully corresponds to two-body systems in isolation



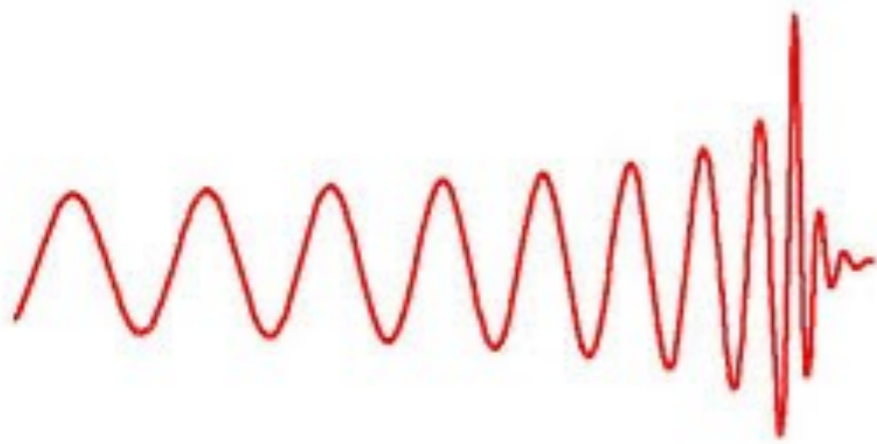
$$\Phi(f) = \phi_0 + 2\pi f t_0 + \sum_{k=0}^7 \alpha_k f^{(k-5)/3}$$

m_1, m_2, χ_1, χ_2

An arrow points from the parameters m_1, m_2, χ_1, χ_2 to the coefficient α_k in the summation term of the equation above.

INTRODUCTION

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$$\Phi(f) = \phi_0 + 2\pi f t_0 + \sum_{k=0}^7 \alpha_k f^{(k-5)/3}$$

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Three-body systems are also quite common !

- 90% of low-mass binaries are expected to belong to a 'hierarchical' triple system

Tokovinin et al. 2006

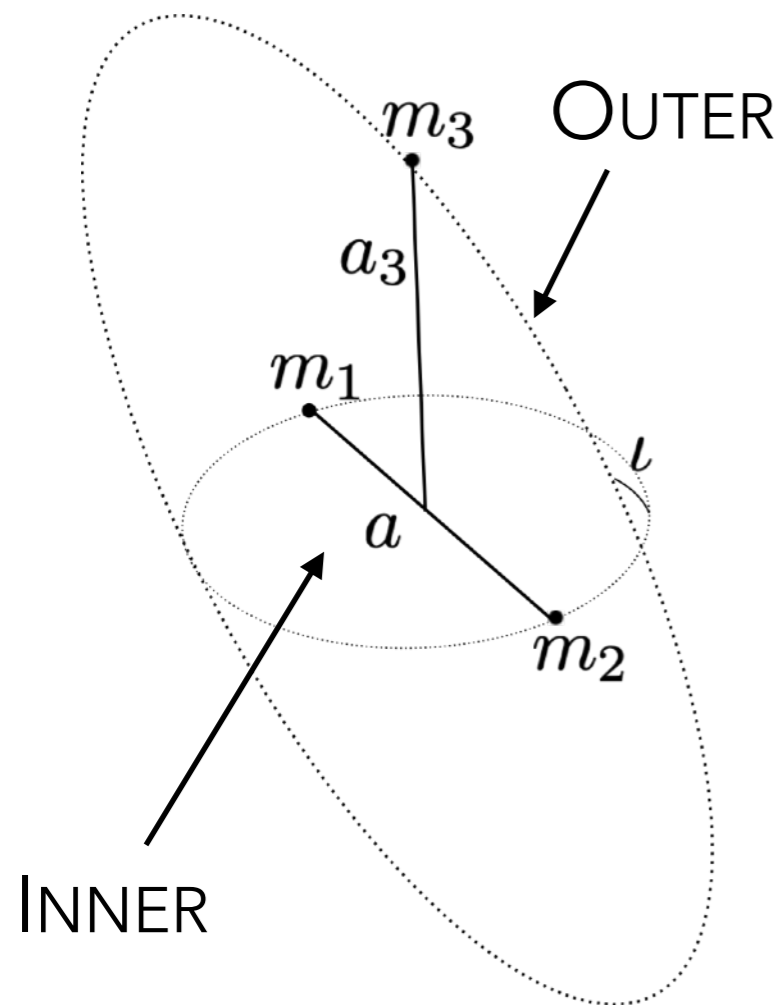
- 'Migration traps' around SMBH at $R \sim 20 - 600 R_{\text{sch}}$

Bellovary et al. 2015



Can we detect and measure parameters of the third body from waveform ?

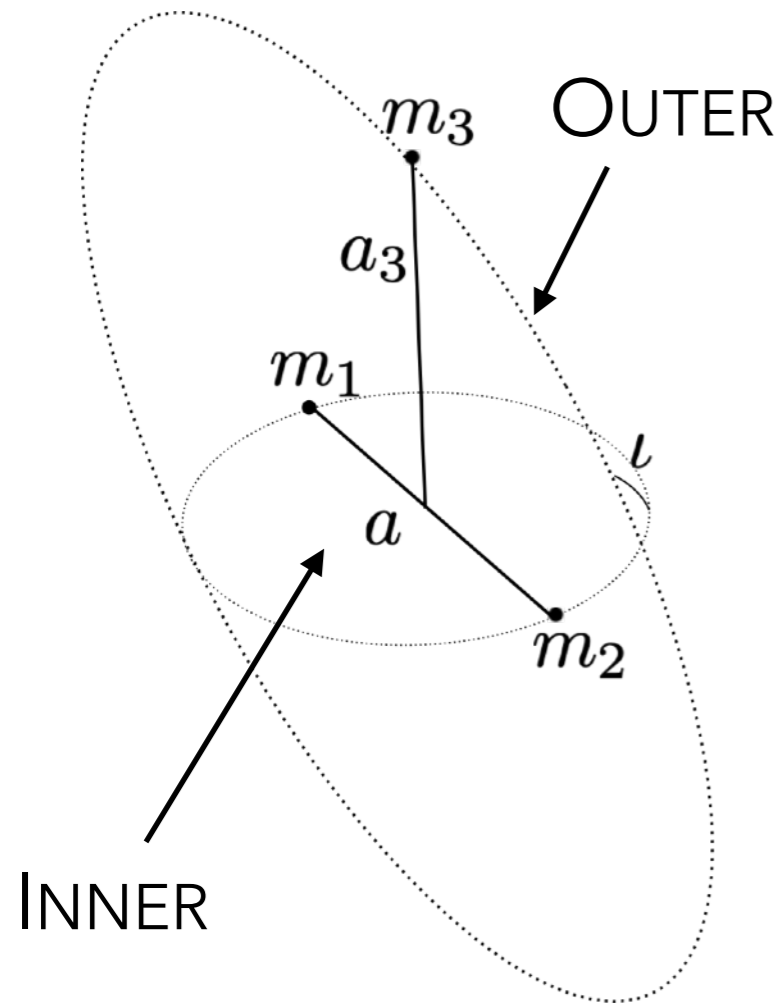
HIERARCHICAL THREE-BODY SYSTEMS



$$m = m_1 + m_2$$

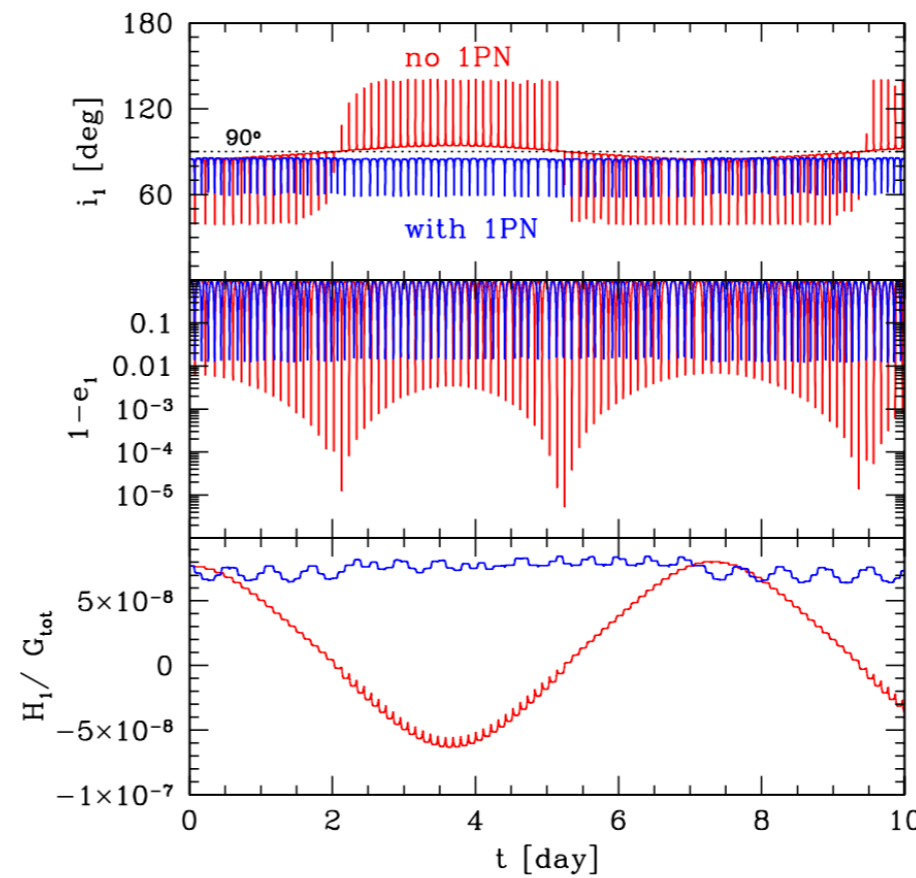
$$M = m_1 + m_2 + m_3$$

HIERARCHICAL THREE-BODY SYSTEMS



Timescales:

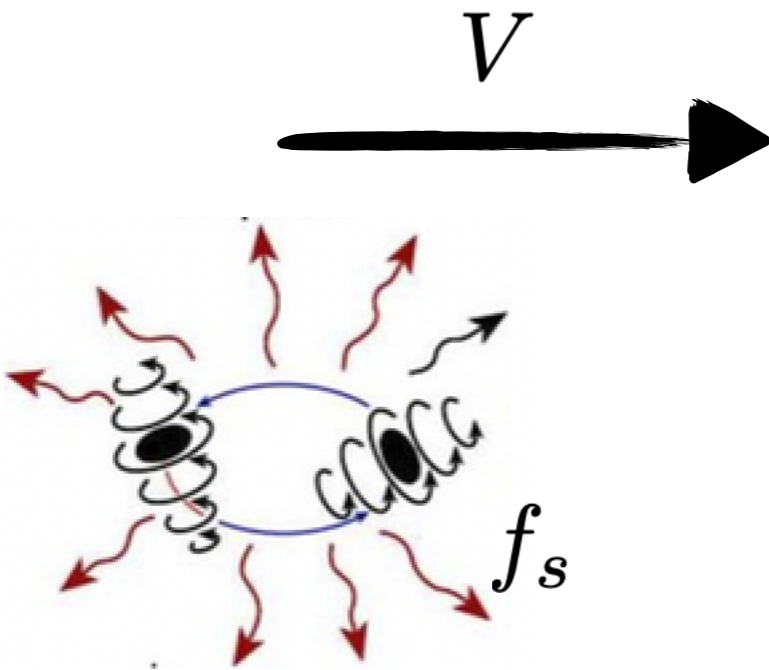
- t_{quad} Kozai-Lidov oscillations
- t_{PN} Perihelion precession



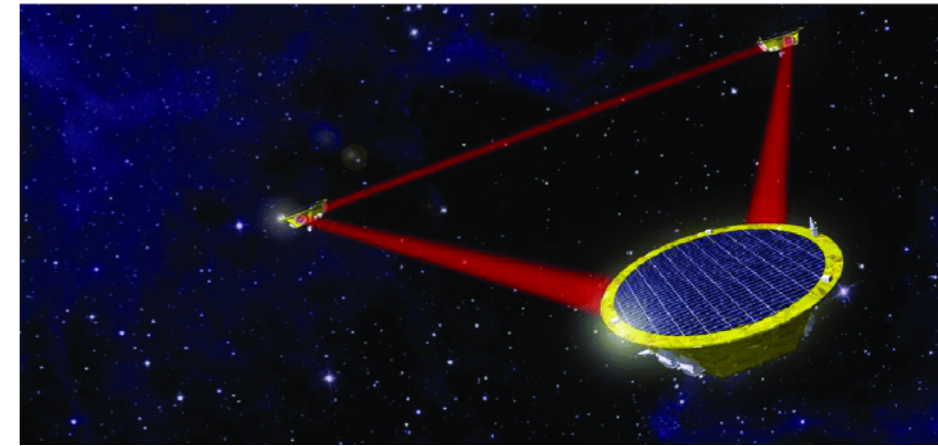
Naoz et al 2013

Outer object cannot influence the eccentricity of very relativistic inner binary

DOPPLER EFFECT



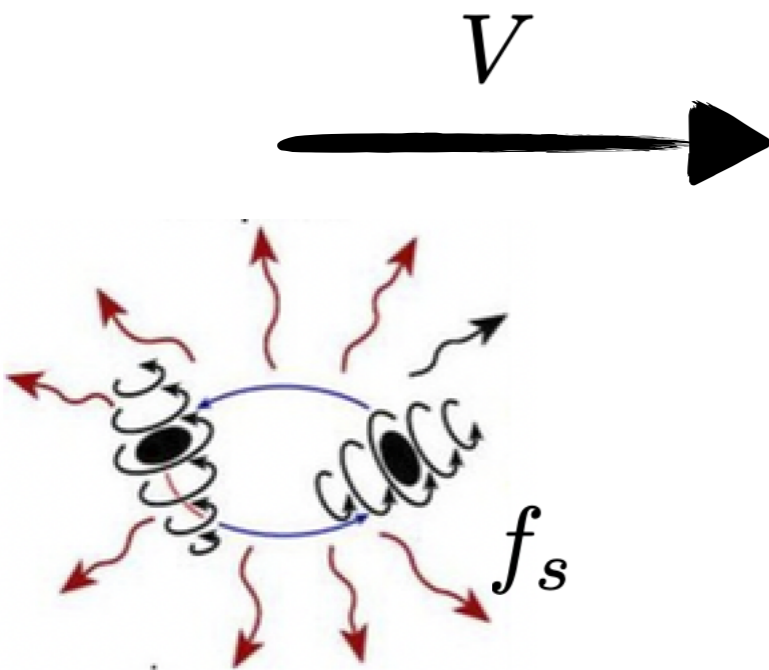
$$f_s (1 + V)$$



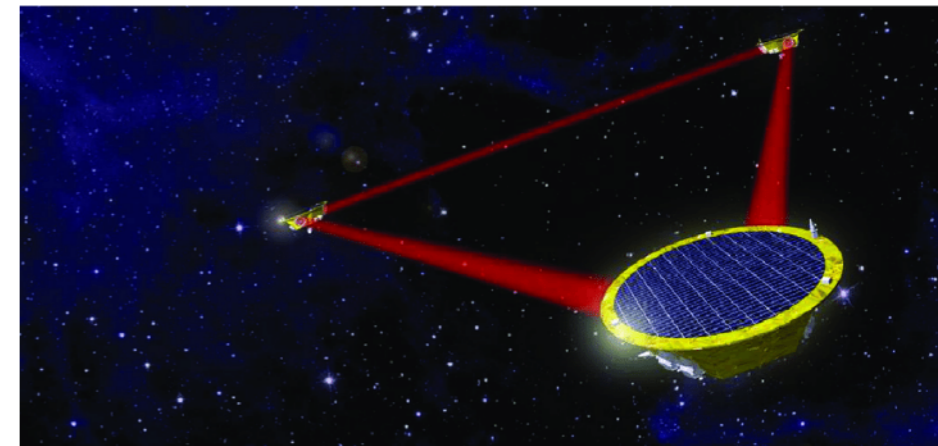
If $V = \text{Cst}$: similar to redshift $\begin{cases} m \rightarrow m(1 + z) \\ f \rightarrow f/(1 + z) \end{cases}$

Unobservable !

DOPPLER EFFECT



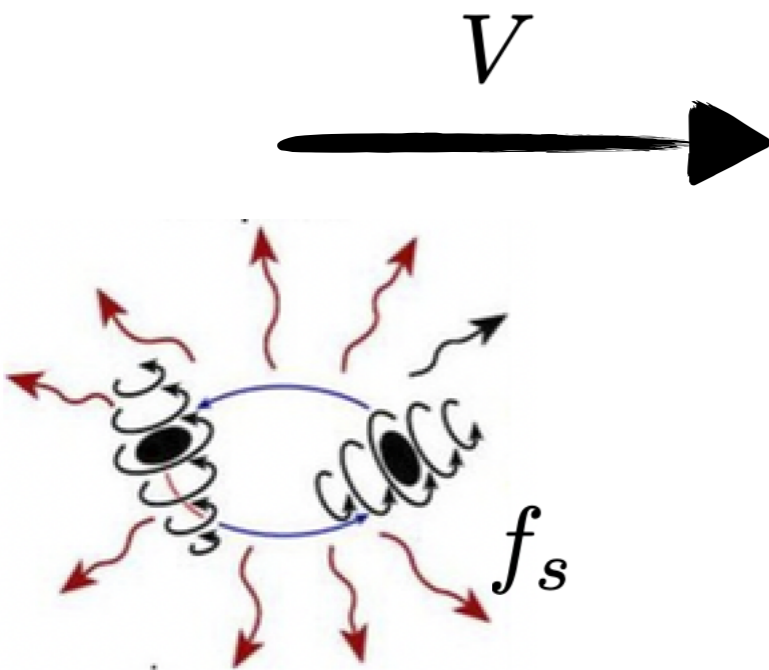
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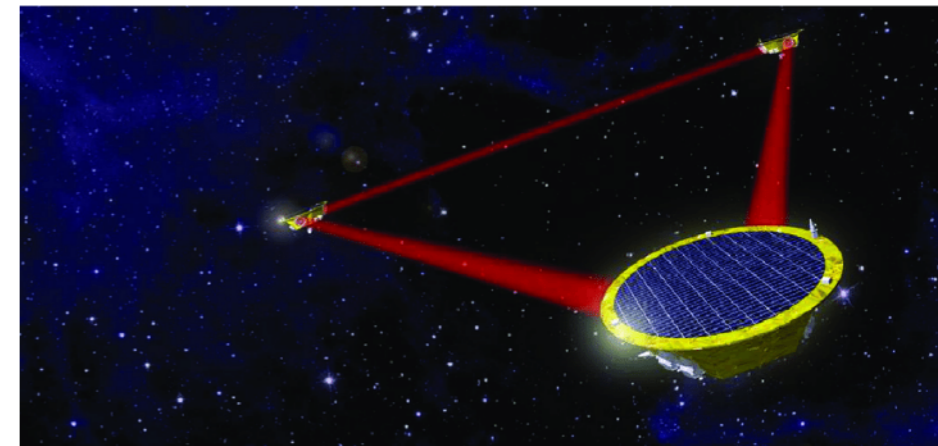
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Time-dependence: $V \sim |V| \cos(2\pi t/P_{\text{out}} + \psi) \sim |V|(1 + t/P_{\text{out}})$
 \swarrow $T_{\text{obs}} \ll P_{\text{out}}$

DOPPLER EFFECT



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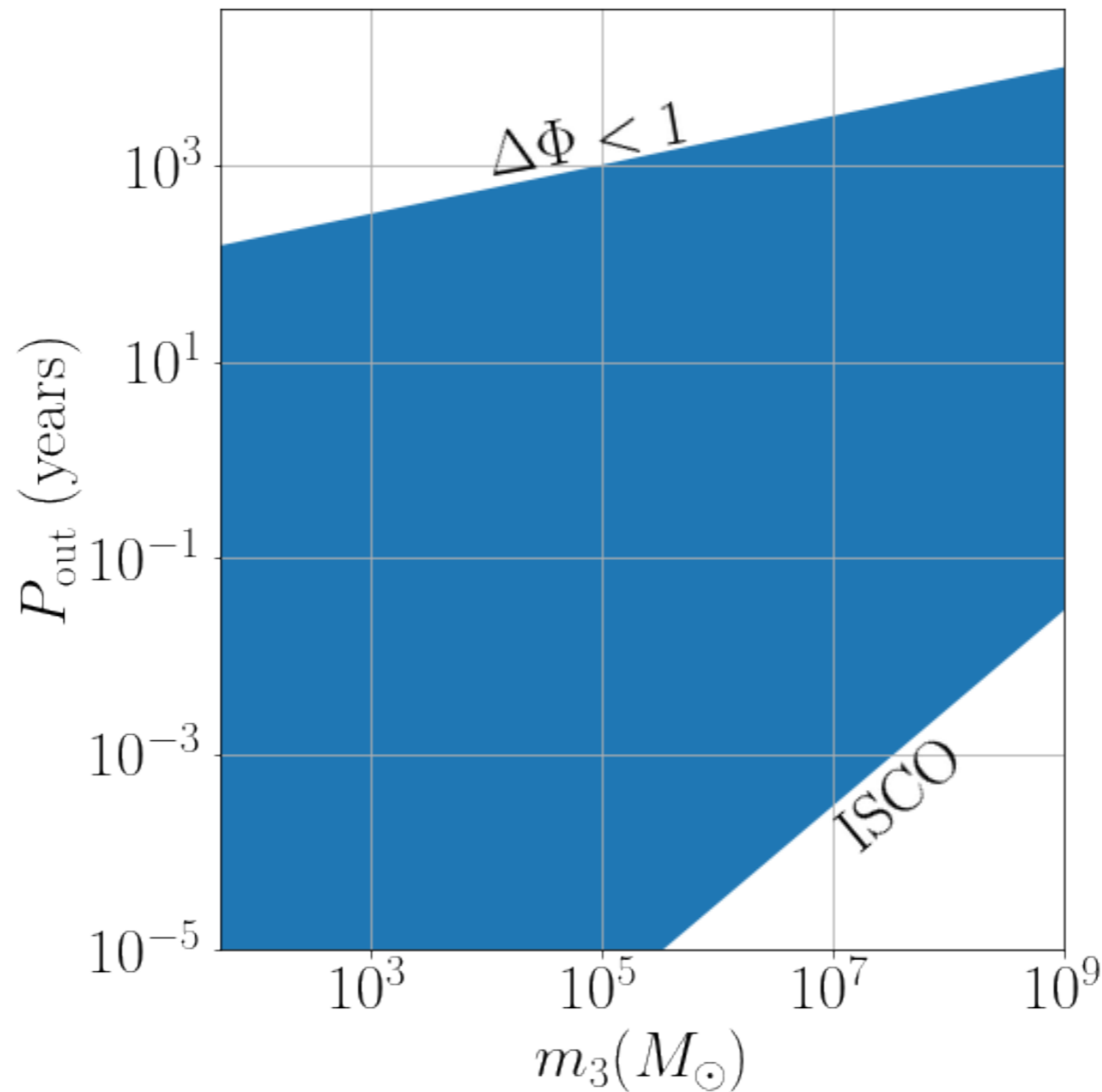
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 \swarrow
 $T_{\text{obs}} \ll P_{\text{out}}$

Phase shift: $\Phi = 2 \int dt f_s (1 + V) \sim 2 \int dt f_s (1 + |V|t/P_{\text{out}})$

A « -4PN » effect since $t_{\text{coal}} - t = \frac{5Gm}{256\nu} \left(\frac{Gm}{a} \right)^{-4}$

DOPPLER PHASE IN LISA

$$\Delta\Phi \sim |V| \frac{T_{\text{obs}}}{P_{\text{out}}} f_s T_{\text{obs}} \sim \left(\frac{GM}{P_{\text{out}}^4} \right)^{1/3} f_s T_{\text{obs}}^2$$



For $m = 50M_{\odot}$ that
will merge within
 $T_{\text{obs}} = 4\text{yr}$

LIMITS OF DOPPLER

- If $P_{\text{out}} \gg T_{\text{obs}}$, we measure only one -4PN parameter while outer orbit depends on

$$m_3, P_{\text{out}}, e_{\text{out}}, \iota$$

- If $P_{\text{out}} \lesssim T_{\text{obs}}$ one can break degeneracies [Yu&Chen 2020](#)
- In any case there remains the degeneracy $m_3 \sin \iota$

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- If the perturber is close enough (e.g. $P_{\text{out}} \lesssim T_{\text{obs}}$), one has to take into account other effects for accurate waveform modelling! (i.e. solving a relativistic three-body problem)

These two shortcomings may cure each other!

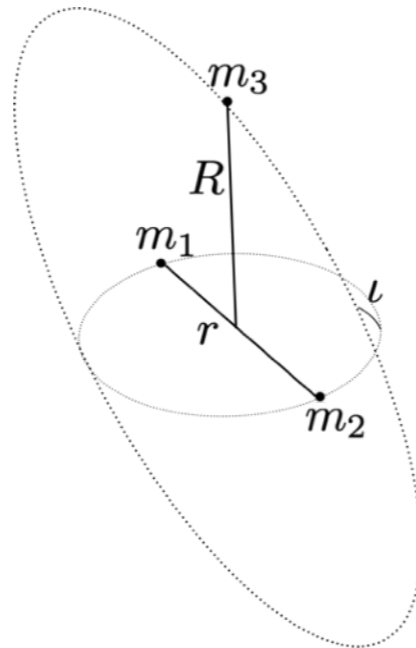
RELATIVISTIC THREE-BODY PROBLEM

GR CORRECTIONS AT 1PN

The 'hardcore' way: use EOM and expand in the CM frame to quadrupole order

Will (2014) Lim and Rodriguez (2020)

$$\begin{aligned}
 \mathbf{a}_a = & -\sum_{b \neq a} \frac{Gm_b \mathbf{x}_{ab}}{r_{ab}^3} + \frac{1}{c^2} \sum_{b \neq a} \frac{Gm_b \mathbf{x}_{ab}}{r_{ab}^3} \left[4 \frac{Gm_b}{r_{ab}} + 5 \frac{Gm_a}{r_{ab}} + \sum_{c \neq a,b} \frac{Gm_c}{r_{bc}} + 4 \sum_{c \neq a,b} \frac{Gm_c}{r_{ac}} - \frac{1}{2} \sum_{c \neq a,b} \frac{Gm_c}{r_{bc}^3} (\mathbf{x}_{ab} \cdot \mathbf{x}_{bc}) - v_a^2 + 4\mathbf{v}_a \cdot \mathbf{v}_b \right. \\
 & \left. - 2v_b^2 + \frac{3}{2} (\mathbf{v}_b \cdot \mathbf{n}_{ab})^2 \right] - \frac{7}{2c^2} \sum_{b \neq a} \frac{Gm_b}{r_{ab}} \sum_{c \neq a,b} \frac{Gm_c \mathbf{x}_{bc}}{r_{bc}^3} + \frac{1}{c^2} \sum_{b \neq a} \frac{Gm_b}{r_{ab}^3} \mathbf{x}_{ab} \cdot (4\mathbf{v}_a - 3\mathbf{v}_b)(\mathbf{v}_a - \mathbf{v}_b), \tag{3.1}
 \end{aligned}$$



GR CORRECTIONS AT 1PN

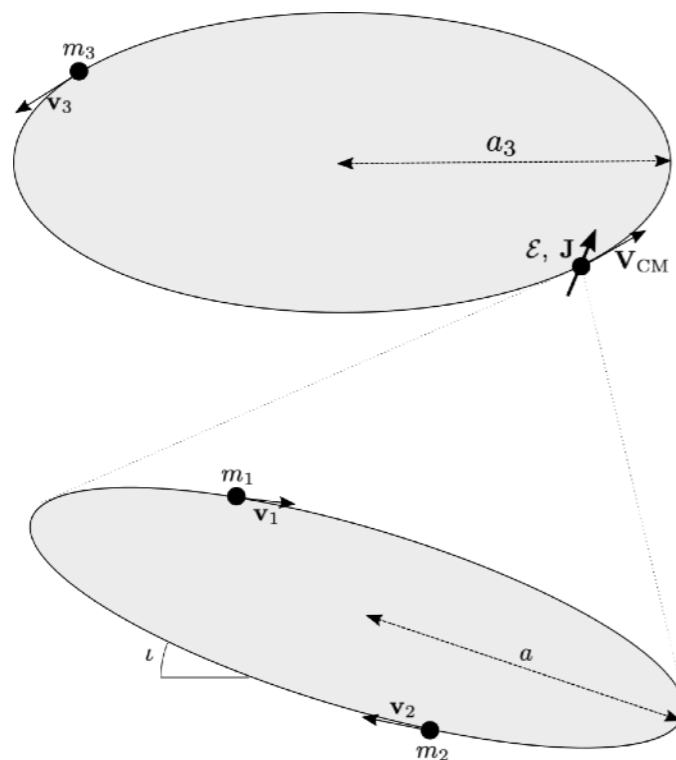
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The 'intuitive' way: 'EFFECTIVE TWO-BODY'

AK, F. Serra, E. Trincherini 2021



3-body motion = 2-body with spin !

THE BINARY EFT

AK, F. Serra, E. Trincherini 2021

PROPER TIME



$$\mathcal{L}_{\text{full}} = \sum_{A=1}^3 -m_A \sqrt{-g_{\mu\nu} v_A^\mu v_A^\nu} \quad \Leftrightarrow \quad \mathcal{L}_{\text{EFT}} = -\mathcal{E} \sqrt{-g_{\mu\nu} V_{\text{CM}}^\mu V_{\text{CM}}^\nu} + \frac{1}{2} J_{\mu\nu} \Omega^{\mu\nu} - m_3 \sqrt{-g_{\mu\nu} v_3^\mu v_3^\nu}$$

The equivalence principle fixes nearly everything!

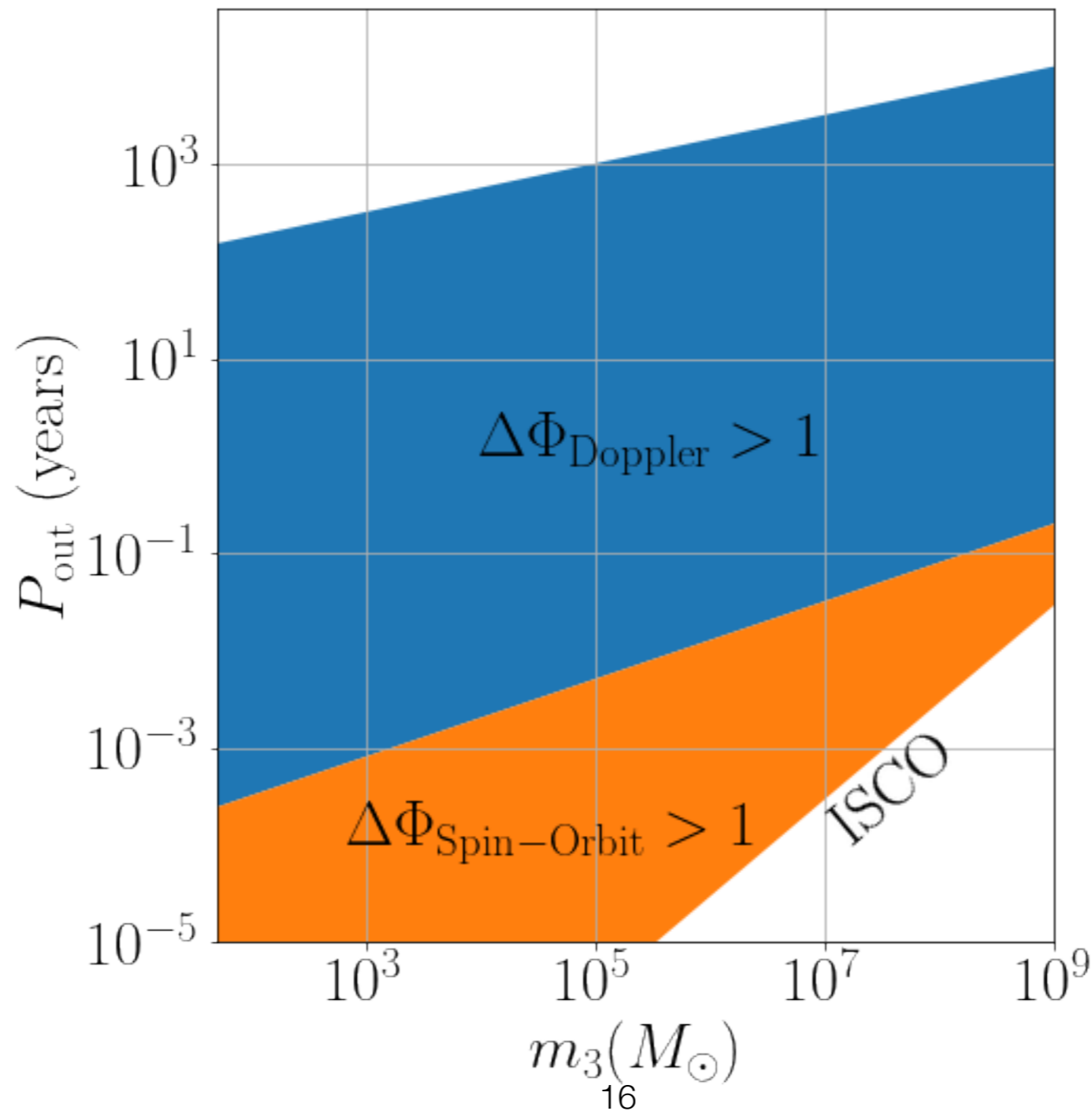
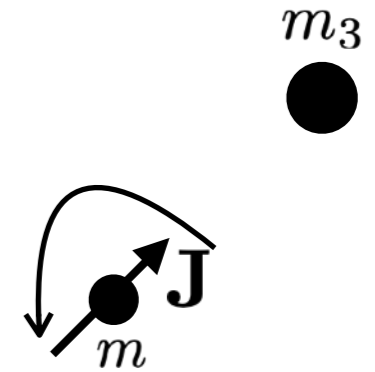
$$\mathcal{E} = m - \frac{G_N m \mu}{2a}, \quad J_{ij} = \epsilon_{ijk} J^k, \quad \mathbf{J} = \sqrt{G_N m a (1 - e^2)} \hat{\mathbf{j}}, \quad \mathbf{\Omega} = \hat{\mathbf{e}} \times \dot{\hat{\mathbf{e}}}$$

$$\mathbf{\Omega}_{ij} = \epsilon_{ijk} \Omega^k,$$

$\hat{\mathbf{e}} \equiv$ UNIT RUNGE-LENZ VECTOR

SPIN-ORBIT COUPLING

$$\mathcal{H}_{\text{SO}} = \frac{4m + 3m_3}{2m} \frac{G}{a_3^3 (1 - e_3^2)^{3/2}} \mathbf{J} \cdot \mathbf{J}_3$$



Yu&Chen 2020

CONCLUSIONS

- LISA will see triple effects in waveforms
- Doppler effect can be probed for perturber at large distance but mass is degenerate with inclination
- Further relativistic three-body effects may allow to break this degeneracy