

# D-branes on local $\mathbb{P}^2$ revisited

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- **“Old story”**: Douglas, M. R., Fiol, B., Römelsberger, C. The spectrum of BPS branes on a noncompact Calabi-Yau, 2000.
  - ▶ Not known at the time: the description of the Kähler moduli space as a modular curve, the precise mathematical notion of Bridgeland stability condition and the wall-crossing formula.
- **Present work**: revisit the story from a particular viewpoint given by scattering diagrams and flow trees.
  - ▶ Related: P. Bousseau, Scattering diagrams, stability conditions, and coherent sheaves on  $\mathbb{P}^2$ , 2019, arXiv:1909.02985.

# Introduction: the general picture

- $X$  Calabi-Yau 3-fold, IIA string on  $\mathbb{R}^{1,3} \times X$ 
  - ▶ BPS spectrum of the resulting  $\mathcal{N} = 2$  4d theory: D0-D2-D4-D6 branes wrapping holomorphic cycles in  $X$ .
  - ▶ Dependence on the Kähler moduli  $\sigma \in \mathcal{M}_K(X)$ , expect

$$\widetilde{\mathcal{M}}_K(X) \rightarrow \text{Stab}(D^b \text{Coh}(X))$$

- ▶  $\Omega(\gamma, \sigma)$ : counts of BPS states of charge  $\gamma =$  DT invariants counting  $\sigma$ -stable objects of class  $\gamma$  in  $D^b \text{Coh}(X)$ . Jump as a function of  $\sigma$  described by the wall-crossing formula.
- Mirror:  $Y$  Calabi-Yau 3-fold, IIB string on  $\mathbb{R}^{1,3} \times Y$ 
  - ▶ BPS spectrum of the resulting  $\mathcal{N} = 2$  4d theory: D3 branes wrapping special Lagrangians in  $Y$ .
  - ▶ Dependence on the complex moduli  $\sigma \in \mathcal{M}_C(Y)$ , expect

$$\widetilde{\mathcal{M}}_C(Y) \rightarrow \text{Stab}(\text{Fuk}(Y))$$

- ▶  $\Omega(\gamma, \sigma)$ : counts of BPS states of charge  $\gamma =$  DT invariants counting  $\sigma$ -stable objects of class  $\gamma$  in  $\text{Fuk}(Y)$ . Jump as a function of  $\sigma$  described by the wall-crossing formula.

- 1) For  $X = K_{\mathbb{P}^2}$ , the description of

$$\widetilde{\mathcal{M}}_K(X) \rightarrow \text{Stab}(D^b \text{Coh}(X))$$

- 2) Determination of the BPS spectrum  $\Omega(\gamma, \sigma)$  by a one-parameter family of "scattering diagrams" made of flow trees in  $\widetilde{\mathcal{M}}_K(X)$

- $X = K_{\mathbb{P}^2} = \mathcal{O}_{\mathbb{P}^2}(-3)$  non-compact Calabi-Yau 3-fold
  - ▶ Projection  $\pi: X \rightarrow \mathbb{P}^2$
  - ▶ Zero section  $\iota: \mathbb{P}^2 \hookrightarrow X$
- $D_{\mathbb{P}^2}(X)$ : bounded derived category of sheaves on  $X$  set-theoretically supported on  $\mathbb{P}^2$  (not necessarily scheme-theoretically!)
  - ▶  $\iota_* : D^b \text{Coh}(\mathbb{P}^2) \rightarrow D_{\mathbb{P}^2}(X)$  (not fully faithful nor essentially surjective!)
  - ▶  $\mathcal{O}(n) := \iota_* \mathcal{O}_{\mathbb{P}^2}(n)$  (D4-branes with  $n$  units of D2-charges)

# The Kähler moduli space $\mathcal{M}_K(X)$ of local $\mathbb{P}^2$

- Mirror symmetry prediction/definition: the universal cover of the Kähler moduli space is the upper half-plane

$$\widetilde{\mathcal{M}_K(X)} = \mathbb{H} = \{\tau \in \mathbb{C} \mid \text{Im}\tau > 0\}$$

- Kähler moduli space: modular curve

$$\mathcal{M}_K(X) = \mathbb{H}/\Gamma_1(3)$$

- Modular group:

$$\Gamma_1(3) := \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2, \mathbb{Z}) \mid \begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} 1 & * \\ 0 & 1 \end{pmatrix} \pmod{3} \right\}$$

- Genus 0 modular curve with two cusps:

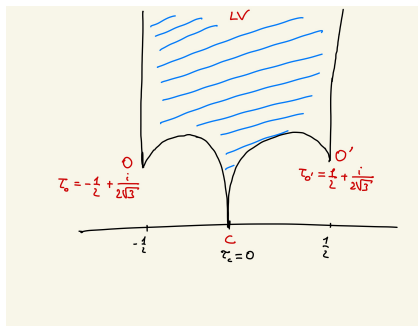
- ▶ "Large volume": class of  $\tau = i\infty$  and every  $p/q \in \mathbb{Q}$  with  $q \equiv 0 \pmod{3}$

- ▶ "Conifold": class of  $\tau = p/q \in \mathbb{Q}$  with  $q \not\equiv 0 \pmod{3}$

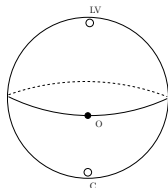
- $\mathbb{Z}/3\mathbb{Z}$ -orbifold point: class of  $\tau_O = -\frac{1}{2} + \frac{i}{2\sqrt{3}}$

# The Kähler moduli space $\mathcal{M}_K(X)$ of local $\mathbb{P}^2$

A fundamental domain  $F_C$  of  $\Gamma_1(3)$  acting on  $\mathbb{H}$ :



The modular curve  $\mathcal{M}_K(X) = \mathbb{H}/\Gamma_1(3)$ :



# Map to the space of stability conditions

- $Stab(D_{\mathbb{P}^2}(X))$ : space of Bridgeland stability conditions on  $D_{\mathbb{P}^2}(X)$ , complex manifold of dimension 3
- Bayer-Macri (2009):

$$\widetilde{M}_K(X) = \mathbb{H} \rightarrow Stab(D_{\mathbb{P}^2}(X))$$

$$\tau \mapsto (\mathcal{A}(\tau), Z(\tau))$$

- Data for a stability condition:
  - ▶  $\mathcal{A}(\tau) \subset D_{\mathbb{P}^2}(X)$ : abelian category, heart of a bounded t-structure
  - ▶ Central charge, additive map:

$$Z(\tau) : \Gamma = K_0(D_{\mathbb{P}^2}(X)) = \mathbb{Z}^3 \rightarrow \mathbb{C}$$

$$\gamma \mapsto Z_\gamma(\tau)$$



- **Mirror:** family of mirror curves, induced by the universal elliptic curve

$$\{E_\tau = \mathbb{C}/(\mathbb{Z} + \tau\mathbb{Z})\}_{\tau \in \mathbb{H}} \rightarrow \mathbb{H} = \{\tau\}$$

- **Mirror symmetry prediction for the central charge:**

$$Z_\gamma(\tau) = \int_\gamma \lambda_\tau$$

$\gamma \in H_1(E_\tau, \mathbb{Z})$ ,  $\lambda_\tau$ : particular meromorphic 1-form on  $E_\tau$ .

- **New formula for the central charge:**

$$Z_E(\tau) = -r(E)T_D(\tau) + d(E)T(\tau) - ch_2(E)$$

where

$$T(\tau) = -\frac{1}{2} + \int_{\tau_0}^{\tau} C(u)du, \quad T_D(\tau) = \frac{1}{3} + \int_{\tau_0}^{\tau} uC(u)du$$

where  $C$ : weight 3 modular form for  $\Gamma_1(3)$ , Eisenstein series:

$$C(\tau) = \frac{\eta(\tau)^9}{\eta(3\tau)^3} = \prod_{n \geq 1} \frac{(1 - q^n)^9}{(1 - q^{3n})^3} = 1 - 9 \sum_{n \geq 1} \chi(n) \frac{q^n}{1 - q^n}$$

$$\text{Dirichlet character } \chi(n) = \begin{cases} \pm 1 & \text{if } n = \pm 1 \pmod{3} \\ 0 & \text{otherwise} \end{cases}$$

# What is the abelian category $\mathcal{A}(\tau)$ ?

- Basic requirement for a Bridgeland stability condition  $(\mathcal{A}(\tau), Z(\tau))$ :

$$Z_E(\tau) \in \{z \in \mathbb{C} \mid \text{Im}z > 0 \text{ or } z \in \mathbb{R}_{<0}\}$$

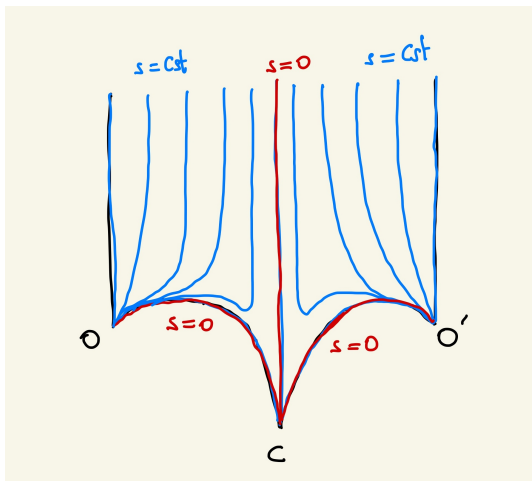
for every  $E \in \mathcal{A}(\tau)$  non-zero.

- ▶ By definition of  $Z(\tau)$ :  $\text{Im}(Z_E(\tau)) = -r\text{Im}(T_D) + d\text{Im}(T)$
- For  $E$  a sheaf on  $X$  (in particular  $r \geq 0$ ),  $\text{Im}(Z_E(\tau)) > 0$  is equivalent to

$$\mu = \frac{d}{r} > s := \frac{\text{Im}(T_D)}{\text{Im}(T)}$$

- ▶ If  $\mu(E) > s$ , want  $E \in \mathcal{A}(\tau)$ .
- ▶ If  $\mu(E) \leq s$ , want  $E[1] \in \mathcal{A}(\tau)$  (clear if  $<$ , why if  $=$ ?)
- ▶ For  $\tau$  in the fundamental domain  $F_C$ , away from the lower boundary,  $\mathcal{A}(\tau) := \mathcal{A}_s$  tilt of the heart of coherent sheaves.

# Curves $s = \text{constant}$



# Proof of $Z_E(\tau) < 0$ if $\mu(E) = s$

- For polynomial stability conditions on surfaces

$$Z_E(s, t) := - \int_{\mathbb{P}^2} e^{-(s+it)H} ch(E)$$

(much of the focus in the math literature, works for any surface), follows from the Bogomolov-Gieseker inequality

$$\Delta(E) := ch_2(E) - 2ch_0(E)ch_1(E) \geq 0$$

for every  $E$  slope-semistable torsion free sheaf.

- Here, "physics" stability condition on  $X = K_{\mathbb{P}^2}$ , need to know something stronger:

$$\Delta(E) \geq \frac{3}{8}$$

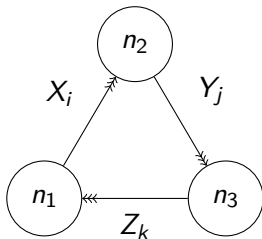
for every  $E$  slope-semistable torsion free sheaf on  $\mathbb{P}^2$  which is not a line bundle (follows from Drézet-Le Potier study of sheaves on  $\mathbb{P}^2$ ).

## What if $\tau$ is on the lower boundary of $F_C$ ?

- On the OC arc,  $s = 0$ ,  $\mathcal{A}(\tau) = \mathcal{A}_0^+$ .
- At the orbifold point  $O$ ,  $\text{Im} T = 0$ ,  $s$  is not well-defined.

$$\mathcal{A}(\tau_O) = \text{Coh}_0(\mathbb{C}^3/(\mathbb{Z}/3\mathbb{Z})) = \text{Rep}^{\text{nilp}}(Q, W)$$

induced by the exceptional collection  $\mathcal{O}, \mathcal{O}(1), \mathcal{O}(2)$  on  $\mathbb{P}^2$ .



Potential  $W = \sum_{i,j,k} \epsilon_{ijk} Z_k Y_j X_i$  with  $\epsilon_{ijk}$  the totally antisymmetric tensor with  $\epsilon_{123} = 1$ .

- There exists a group homomorphism

$$\Gamma_1(3) \longrightarrow \text{Aut}(D_{\mathbb{P}^2}(X))$$

- ▶ Monodromy around the large volume point:

$$\begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} \longmapsto (E \mapsto E \otimes \mathcal{O}(1))$$

- ▶ Monodromy around the conifold point:

$$\begin{pmatrix} 1 & 0 \\ -3 & 1 \end{pmatrix} \longmapsto (ST_{\mathcal{O}} : E \mapsto \text{Cone}(\text{RHom}(\mathcal{O}, E) \otimes \mathcal{O} \rightarrow E))$$

- To summarize:

$$\widetilde{M}_K(X) = \mathbb{H} \rightarrow \text{Stab}(D_{\mathbb{P}^2}(X))$$
$$\tau \mapsto (\mathcal{A}(\tau), Z(\tau))$$

- We can then do DT theory.

- ▶ Moduli spaces

$$M(\gamma, \tau) = \{\tau\text{-semistable objects in } \mathcal{A}(\tau) \text{ of class } \gamma\}$$

- ▶ DT/BPS invariants:

$$\Omega(\gamma, \tau) \in \mathbb{Z} \text{ or } \mathbb{Z}[y^{\pm}]$$

- ▶ Wall-crossing as a function of  $\tau \in \mathbb{H}$ .

- Goal: study of the DT/BPS invariants using flow trees organized in

“scattering diagrams” in  $\widetilde{M}_K(X) = \mathbb{H}$

- ▶ supergravity attractor picture
- ▶ Kontsevich-Soibelman wall-structure on base of complex integrable systems.



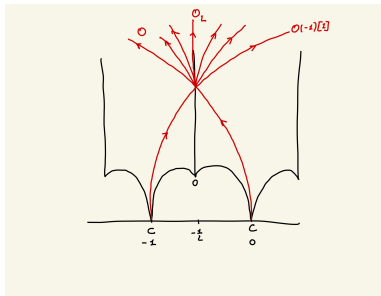
# Scattering diagrams

- Pick a phase  $\psi \in \mathbb{R}/2\pi\mathbb{Z}$

- ▶ For every  $\gamma \in \Gamma$ , consider the 1-dimensional locus, "rays":

$$\mathcal{R}_\gamma^+(\psi) := \{\tau \in \mathbb{H} \mid \text{Arg}(Z_\gamma(\tau)) = \psi + \frac{\pi}{2}, \Omega(\gamma, \tau) \neq 0\} \subset \mathbb{H}$$

- ▶ Orient rays such that  $\text{Im}(e^{-i\psi} Z_\gamma(\tau)) = |Z_\gamma(\tau)|$  increases.
- ▶ "K-wall" in Gaiotto-Moore-Neitzke terminology.



- “Scattering diagram of phase  $\psi$ ”:

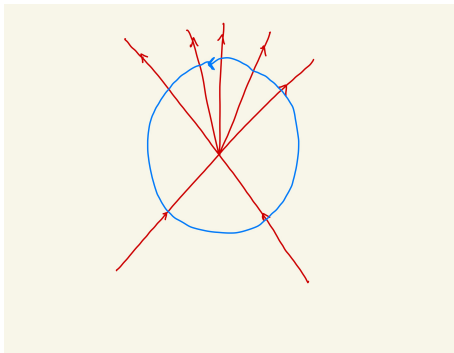
$$\mathcal{D}(\psi) = \bigcup_{\gamma \in \Gamma} (\mathcal{R}_{\gamma}^{+}(\psi), \exp(\bar{\Omega}(\gamma, \tau)z^{\gamma}))$$

- ▶ Rational DT invariants

$$\bar{\Omega}(\gamma, \tau)(y) = \sum_{\gamma = k\gamma'} \frac{1}{k} \frac{1}{y^k - y^{-k}} \Omega(\gamma', \tau)(y^k)$$

- ▶ Quantum torus variables:  $z^{\gamma}z^{\gamma'} = y^{\langle \gamma, \gamma' \rangle} z^{\gamma + \gamma'}$ , where  $\langle -, - \rangle$ : Euler form on  $\Gamma$

- Reformulation of the wall-crossing formula:  $\mathcal{D}(\psi)$  is consistent
  - ▶ for every point  $\tau \in \mathbb{H}$ , the ordered product of quantum torus elements of attached to the rays passing through  $\tau$  is the identity.



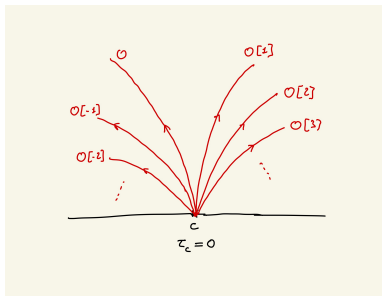
## Theorem

For every  $\psi \in \mathbb{R}/2\pi\mathbb{Z}$ , the scattering diagram  $\mathcal{D}(\psi)$  can be uniquely reconstructed from:

- *Explicit initial rays coming from the conifold points.*
  - *Scatterings imposed by the consistency condition.*
- 
- Application: algorithmic reconstruction of the full BPS spectrum (except pure D0) at any point of the physical space of stability conditions.

# Initial rays

At the conifold point  $\tau_O = 0$ ,  $Z_O(\tau_O) = 0$ . Infinitely many initial rays corresponding to the objects  $\mathcal{O}[k]$ ,  $k \in \mathbb{Z}$ .



General conifold point: apply  $\Gamma_1(3)$ , spherical object  $E$  becoming massless, infinitely many initial rays corresponding to the objects  $E[k]$ ,  $k \in \mathbb{Z}$ .

# Reconstruction from initial rays

- Rays of  $\mathcal{D}(\psi)$  are gradient flow lines of  $\text{Im}(e^{-i\psi} Z_\gamma(\tau))$ . Then,  $\text{Re}(e^{-i\psi} Z_\gamma(\tau)) = Cst$ , and we recover the field theory limit of attractor flow trees (Denef).
- Key point: for every  $\gamma \in \Gamma$ , the holomorphic function

$$\mathbb{H} \rightarrow \mathbb{C}$$

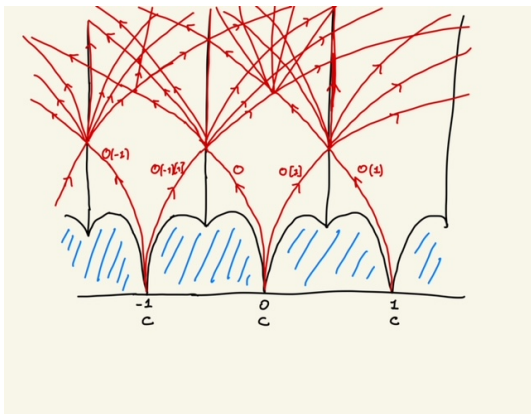
$$\tau \mapsto Z_\gamma(\tau)$$

has no critical point on  $\mathbb{H}$ :

$$\frac{d}{d\tau} Z_\gamma(\tau) = (-r\tau + d)C(\tau) \neq 0$$

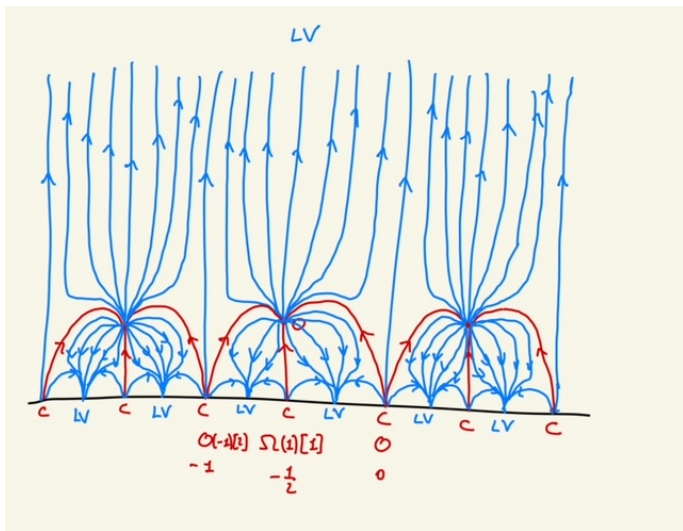
- Study of the boundary behavior:  $C(\tau) \rightarrow 0$  when  $\tau$  goes to a conifold point, not otherwise.
- For compact Calabi-Yau 3-folds, normalized non-holomorphic central charge  $Z_\gamma^{norm} = e^{K/Z} Z_\gamma = \int_\gamma \Omega / \sqrt{|\int_\gamma \Omega \wedge \overline{\Omega}|}$ , non-trivial "attractor points".

# The scattering diagram $\mathcal{D}(\psi)$ for $\psi = 0$



Recovers the scattering diagram of arXiv:1909.02985, computing the BPS spectrum at large volume, and so the scattering diagram of Carl-Pumperla-Siebert counting holomorphic curves in  $(\mathbb{P}^2, E)$  for mirror symmetry (hyperkähler rotated version of the special Lagrangians).

# The scattering diagram $\mathcal{D}(\psi)$ for $\psi = \frac{\pi}{2}$





# The scattering diagram $\mathcal{D}(\psi)$ for $\psi = \frac{\pi}{2}$

## Theorem

*The local picture of  $\mathcal{D}(\pi/2)$  near the orbifold point  $\tau_0$  can be identified with the slice  $\theta_1 + \theta_2 + \theta_3 = 0$  of the stability scattering diagram in  $\mathbb{R}^3$  for  $(Q, W)$ .*

## Corollary

*The stability scattering diagram of  $(Q, W)$  equals the corresponding cluster scattering diagram (up to the  $D_0$  wall), and so the attractor invariants of  $(Q, W)$  are trivial (up to  $D_0$  branes).*

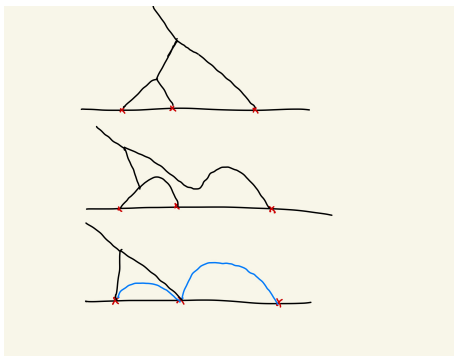
Alternative algebraic proof of the Corollary: see Descombes's talk.

# The scattering diagram $\mathcal{D}(\psi)$ for $\psi = \frac{\pi}{2}$

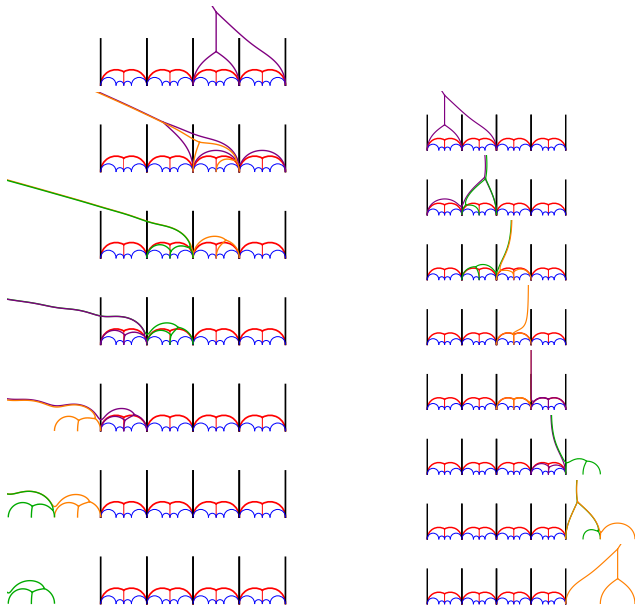
- The global picture of  $\mathcal{D}(\pi/2)$  also give a clear description of the correspondence between normalized ( $-1 < \mu \leq 0$ ) torsion free Gieseker semi-stable sheaves on  $\mathbb{P}^2$  and representations of the Beilinson quiver.
- For these objects  $\mathcal{D}(0)$  gives a path from the large volume point to the orbifold point avoiding the walls of marginal stability.

# Interpolation between $\psi = 0$ and $\psi = \frac{\pi}{2}$

- Interesting discrete jumps.
- Infinite sequence of critical phases.
- Physics interpretation?



# Jumps in $\mathcal{D}(\psi)$ as function of $\psi$



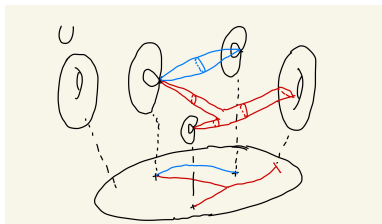
- Both (exponential) spectral networks and scattering diagrams compute the same DT/BPS numbers, “orthogonal approaches”:
  - ▶ Spectral networks drawn on a projection of the mirror curves.
  - ▶ Scattering diagrams drawn on the base of the family of mirror curves.
- Analogy: scattering diagrams as “4d spectral network”.

# Analogy with spectral networks

- “2d spectral networks”:
  - ▶ capture BPS spectrum of a  $\mathcal{N} = (2, 2)$  2d theory
  - ▶ rays interact according to the 2d Cecotti-Vafa wall-crossing formula, geometrically: special Lagrangians in  $T^*C$  with boundary on spectral curve  $\Sigma \rightarrow C$  ( $n:1$  cover).
  - ▶ Wall-crossing structure for  $\mathfrak{gl}_n$ .
  - ▶ Jumps in  $\psi$ : BPS spectrum of  $\mathcal{N} = 2$  4d theory containing the  $\mathcal{N} = (2, 2)$  2d theory as a surface defect.
- Scattering diagrams = “4d spectral networks”:
  - ▶ capture BPS spectrum of a  $\mathcal{N} = 2$  4d theory
  - ▶ rays interact according to the 4d Kontsevich-Soibelman wall-crossing formula, geometrically: special Lagrangians in a torus fibration with boundary on a torus fiber.
  - ▶ Wall-crossing structure for the quantum torus Lie algebra.
  - ▶ Jumps in  $\psi$ ?

## 4d-6d wall-crossing?

- $U$ : elliptic fibration with base a 3:1 cover of  $\mathcal{M}_K(X)$ . HyperKähler 4-manifold.
- Known string duality: IIA on  $\mathbb{R}^{1,3} \times X$  dual to IIA on  $\mathbb{R}^{1,5} \times U$  with a NS5-brane on  $\mathbb{R}^{1,3} \times T^2$ , where  $T^2$  is an elliptic fiber of  $U$ .
- 4d BPS spectrum: D2-branes/special Lagrangians in  $U$  ending on the NS5 brane. Scattering diagrams: projections on the base (string junctions).
- $\mathcal{N} = 2$  4d theory as a defect in the  $\mathcal{N} = (1, 1)$  6d theory obtained by IIA on  $\mathbb{R}^{1,5} \times U$ . Jumps of the scattering diagrams as a function of  $\psi$ : closed D2-branes/special Lagrangians in  $U$ , i.e. 6d BPS states.
- 4d-6d wall-crossing formula?



Thank you for your attention !