



**Quantum simulation with arrays of superconducting qubits**  
**First applications in probing many-body localization**

Journée Théorie du CPTGA

Michele Filippone



# Historical

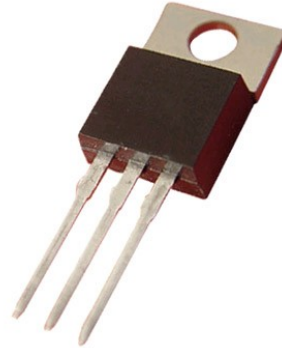
Silicon



band theory



Transistor



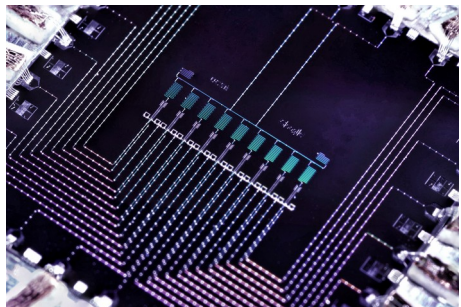
Classical bits

01010010 . . .



# Will history repeat ?

Synthetic  
Quantum Matter



Theory of  
Quantum Matter

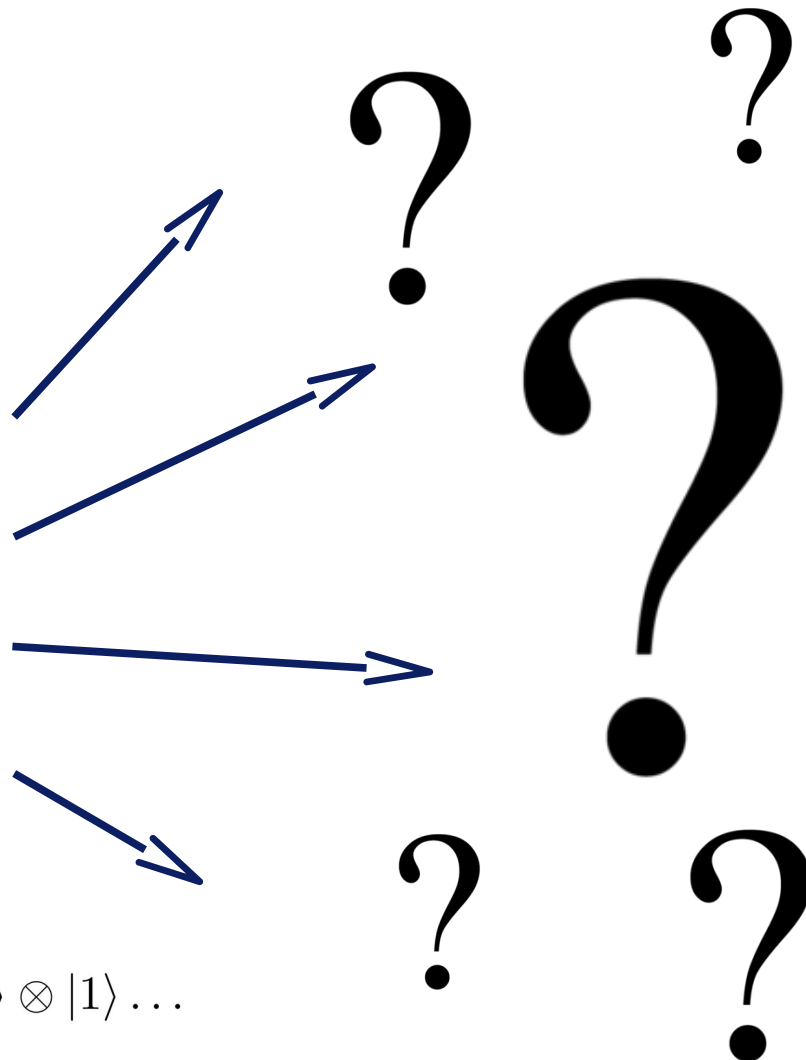


Miraculous  
device ?



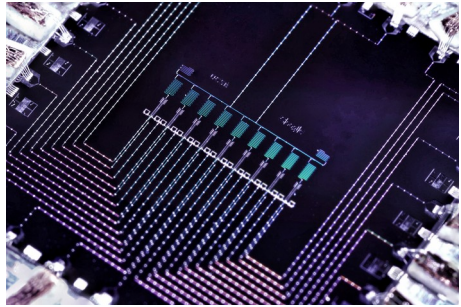
Quantum bits

$$\frac{|0\rangle + |1\rangle}{\sqrt{2}} \otimes \frac{|0\rangle + |1\rangle}{\sqrt{2}} \otimes |0\rangle \otimes |1\rangle \dots$$



# Will history repeat ?

Synthetic  
Quantum Matter



Theory of  
Quantum Matter

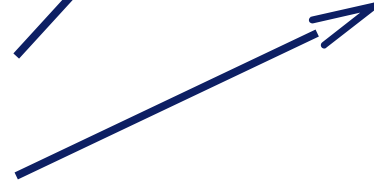
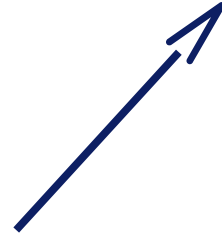


Miraculous  
device ?



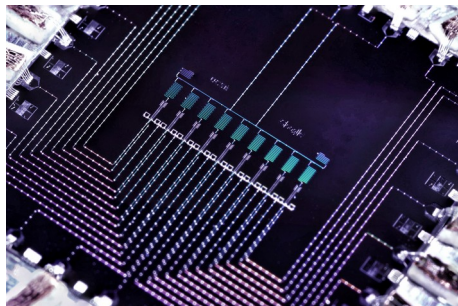
Quantum bits

$$\frac{|0\rangle + |1\rangle}{\sqrt{2}} \otimes \frac{|0\rangle + |1\rangle}{\sqrt{2}} \otimes |0\rangle \otimes |1\rangle \dots$$



# Will history repeat ?

Synthetic  
Quantum Matter



CPTGA - 2021



Theory of  
Quantum Matter

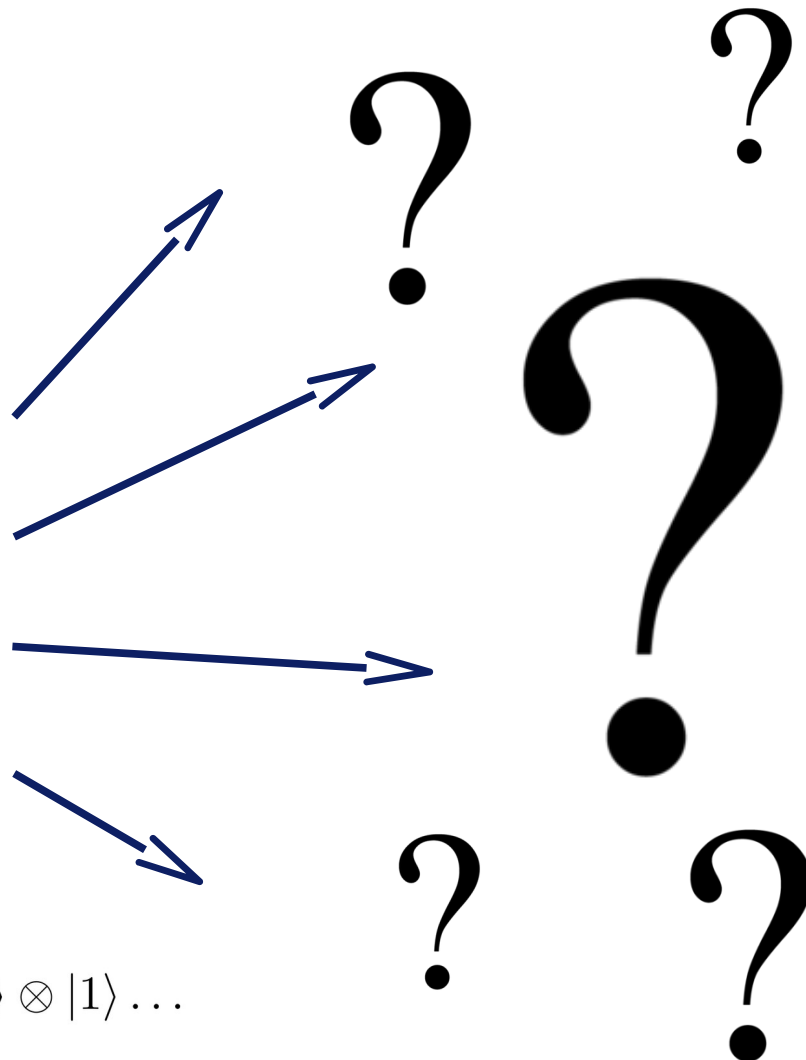


Miraculous  
device ?



Quantum bits

$$\frac{|0\rangle + |1\rangle}{\sqrt{2}} \otimes \frac{|0\rangle + |1\rangle}{\sqrt{2}} \otimes |0\rangle \otimes |1\rangle \dots$$



# Last two years ...

Article

## Quantum supremacy using a programmable superconducting processor

Nature | Vol 574 | 24 OCTOBER 2019 | 505

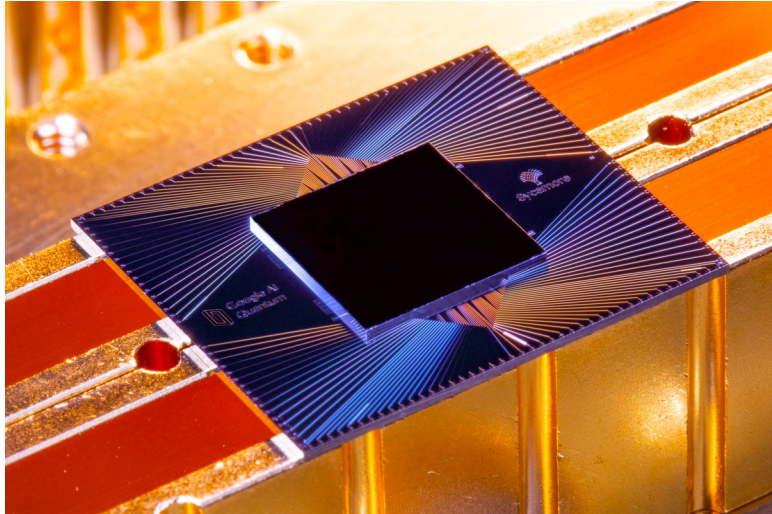


Photo credit: Eric Lucero



RESEARCH

QUANTUM COMPUTING

## Quantum computational advantage using photons

Zhong *et al.*, *Science* **370**, 1460–1463 (2020)



Photo credit: PhysicsWorld



# Last two years ...

Article

## Quantum supremacy using a programmable superconducting processor

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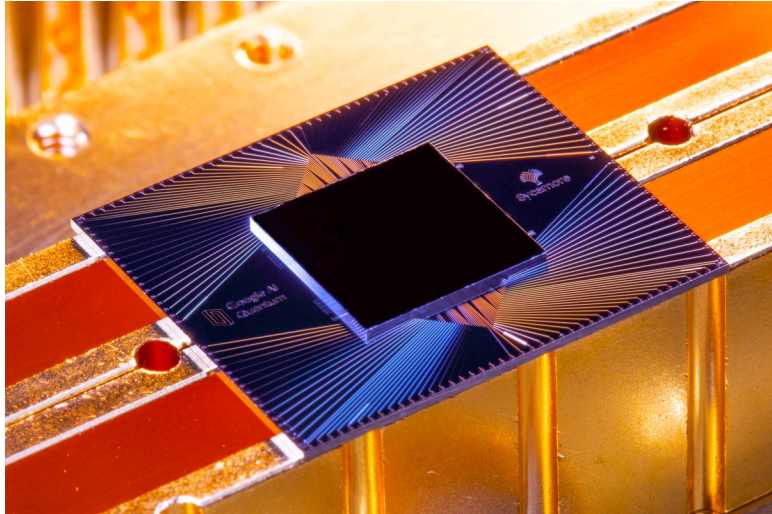


Photo credit: Eric Lucero

Google

RESEARCH

QUANTUM COMPUTING

## Quantum computational advantage using photons

Zhong *et al.*, *Science* **370**, 1460–1463 (2020)



Photo credit: PhysicsWorld



# Last two years ...

Article

## Quantum supremacy using a programmable superconducting processor

RESEARCH

PHYSICAL REVIEW X 10, 041038 (2020)

Featured in Physics

### What Limits the Simulation of Quantum Computers?

Yiqing Zhou<sup>1,2</sup>, E. Miles Stoudenmire<sup>1,2</sup>, and Xavier Waintal<sup>3</sup>  
<sup>1</sup>Department of Physics, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801, USA  
<sup>2</sup>Center for Computational Quantum Physics, Flatiron Institute, New York, New York 10010, USA  
<sup>3</sup>Univ. Grenoble Alpes, CEA, IRIG-Phelqs, 38054 Grenoble, France



Photons

370, 1460-1463 (2020)

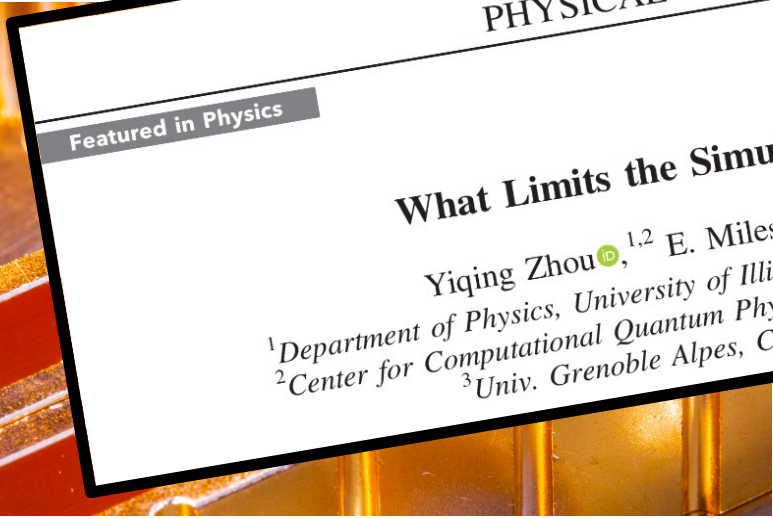


Photo credit: Eric Lucero

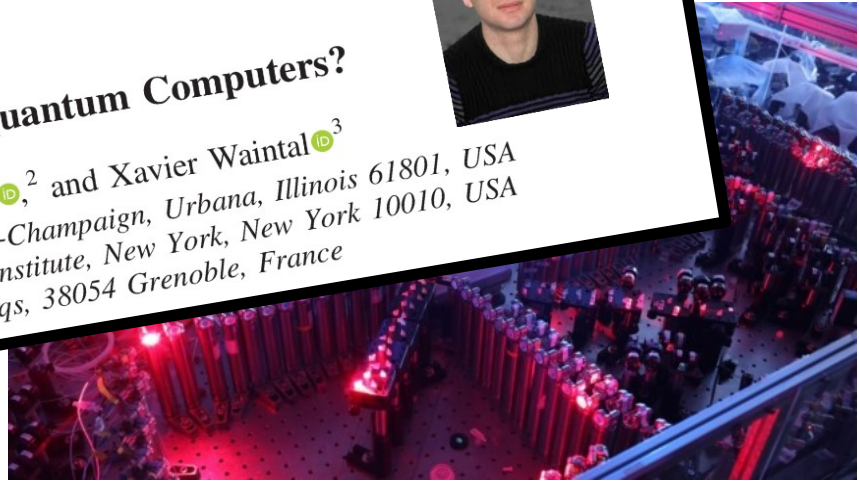


Photo credit: PhysicsWorld

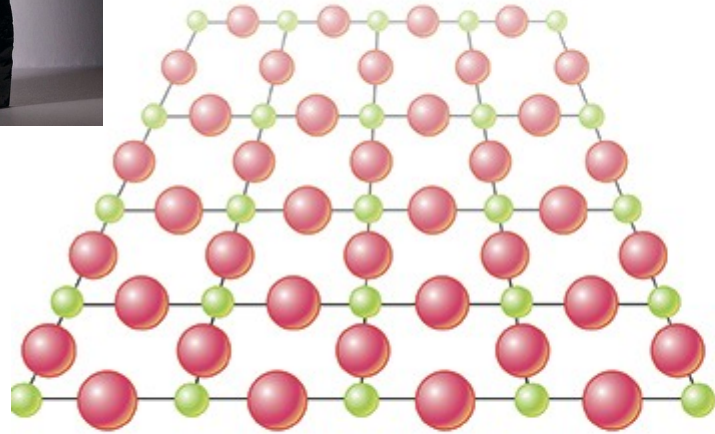






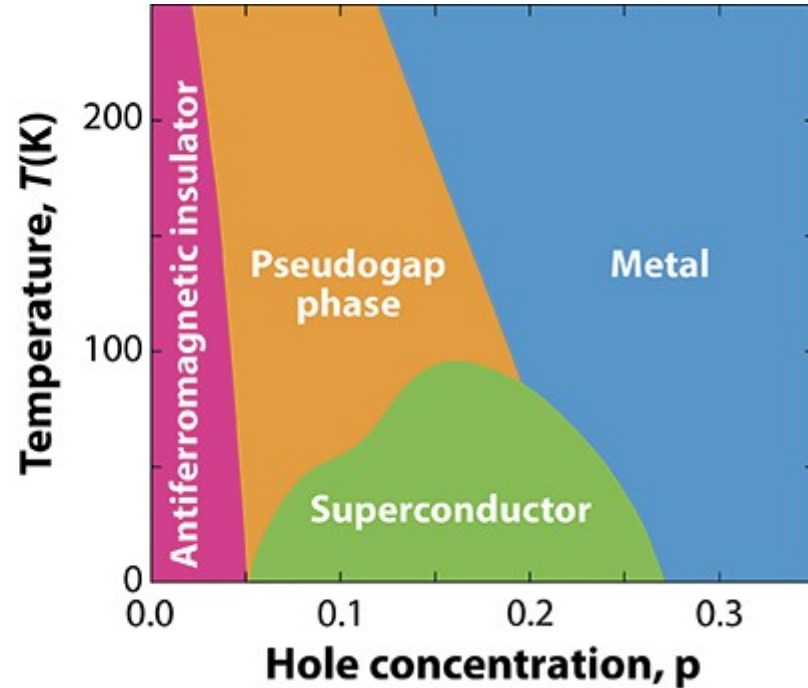
# Quantum simulation

# Does the Fermi-Hubbard model describe High-Tc Superconductors?



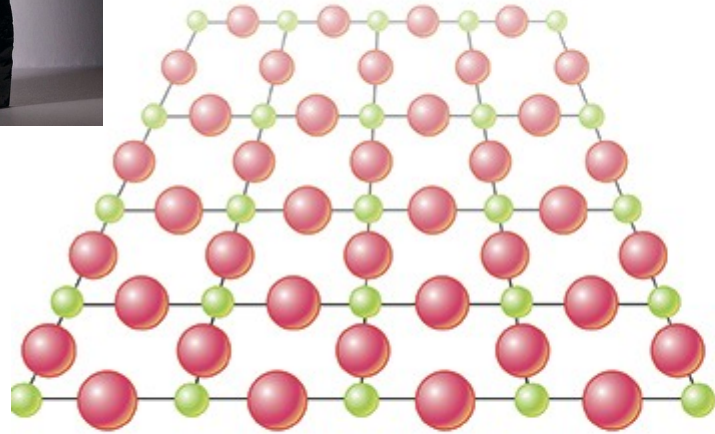
● Copper ● Oxygen

Electrons hop between the copper ions via the oxygen bonds



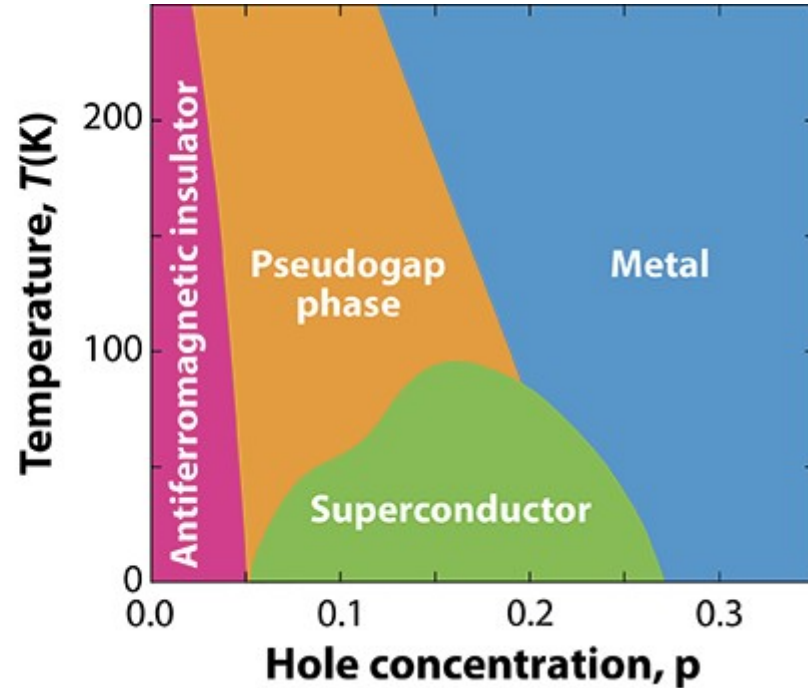
$$H_{\text{FH}} = t \sum_{\langle i,j \rangle, \sigma} \left[ c_{i,\sigma}^\dagger c_{j,\sigma} + \text{h.c.} \right] + U \sum_i n_{i,\uparrow} n_{i,\downarrow}$$

# Does the Fermi-Hubbard model describe High-Tc Superconductors?



● Copper ● Oxygen

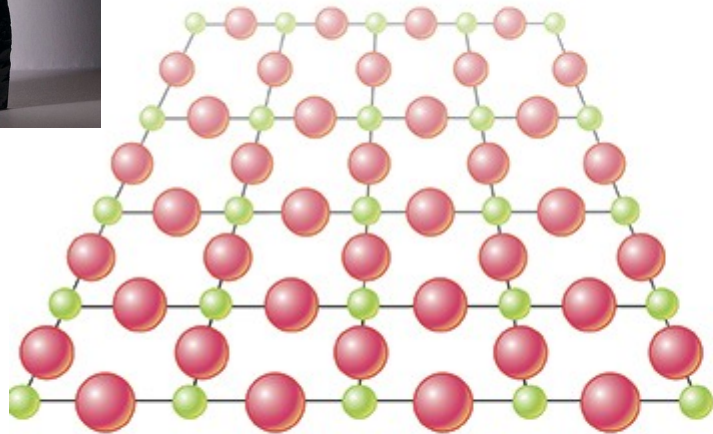
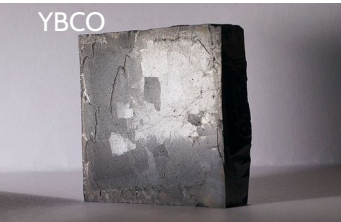
Electrons hop between the copper ions via the oxygen bonds



$$H_{\text{FH}} = t \sum_{\langle i,j \rangle, \sigma} \left[ c_{i,\sigma}^\dagger c_{j,\sigma} + \text{h.c.} \right] + U \sum_i n_{i,\uparrow} n_{i,\downarrow}$$

# Does the Fermi-Hubbard model explain High Temperature Superconductors?

YBCO



● Copper ● Oxygen

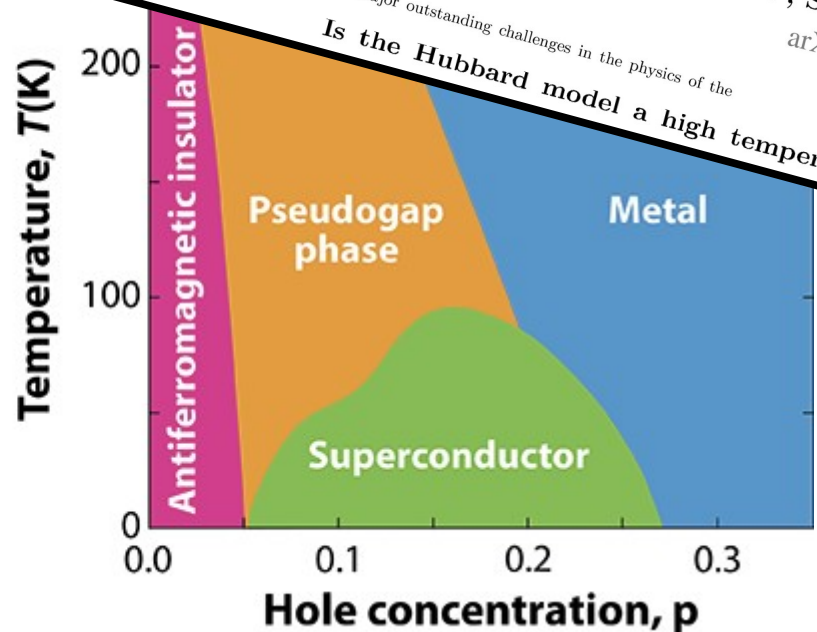
Electrons hop between the copper ions via the oxygen bonds

$$H_{\text{FH}} = t \sum_{\langle i,j \rangle, \sigma} [c_{i,\sigma}^\dagger c_{j,\sigma} + \text{h.c.}] + U \sum_i n_{i,\uparrow} n_{i,\downarrow}$$

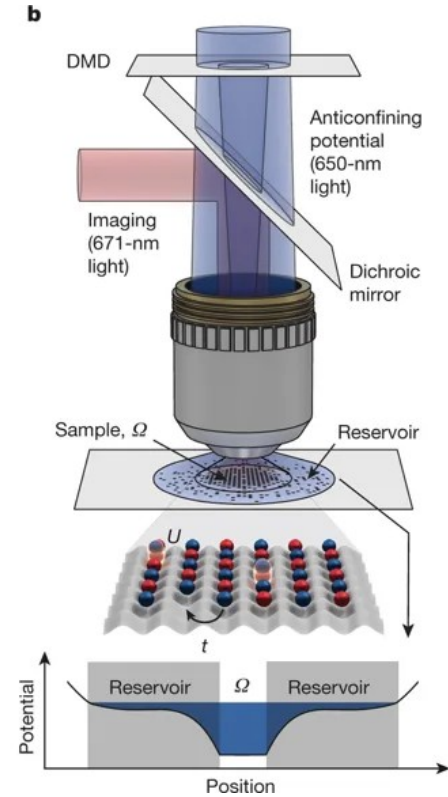
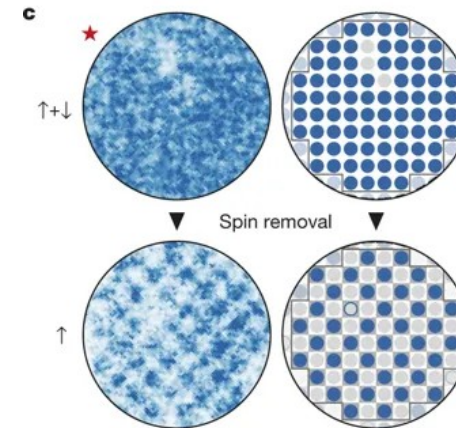
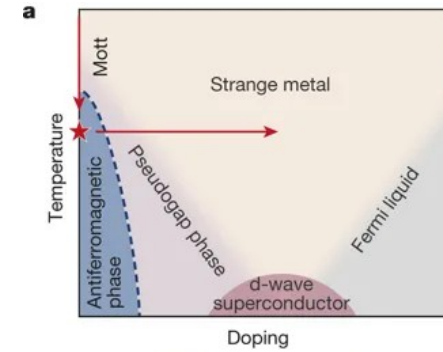
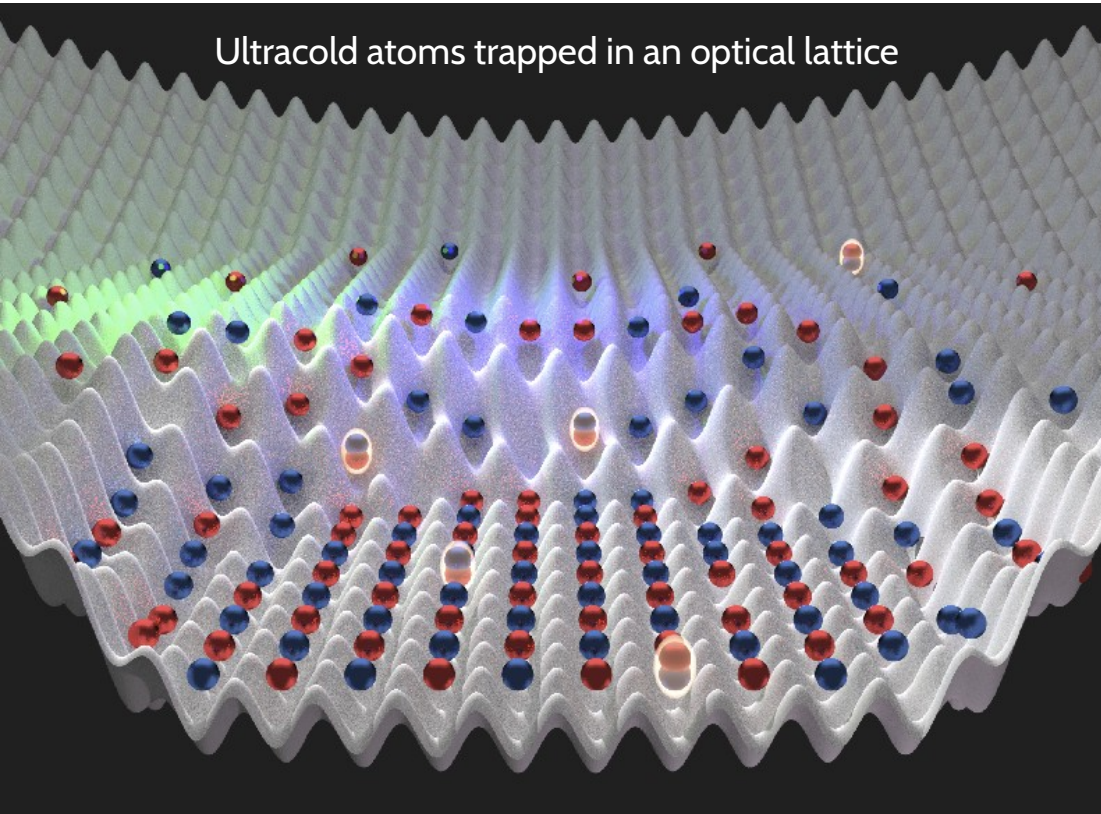
**The Hubbard Model**  
 Daniel P. Arovas<sup>1</sup>, Erez Berg<sup>2</sup>, Steven A. Kivelson<sup>3</sup>, Srinivas Raghu<sup>3,4</sup>  
 arXiv:2103.12097v2

**9. Important Open Questions**  
 We end by highlighting some of the major outstanding challenges in the physics of the Hubbard model.

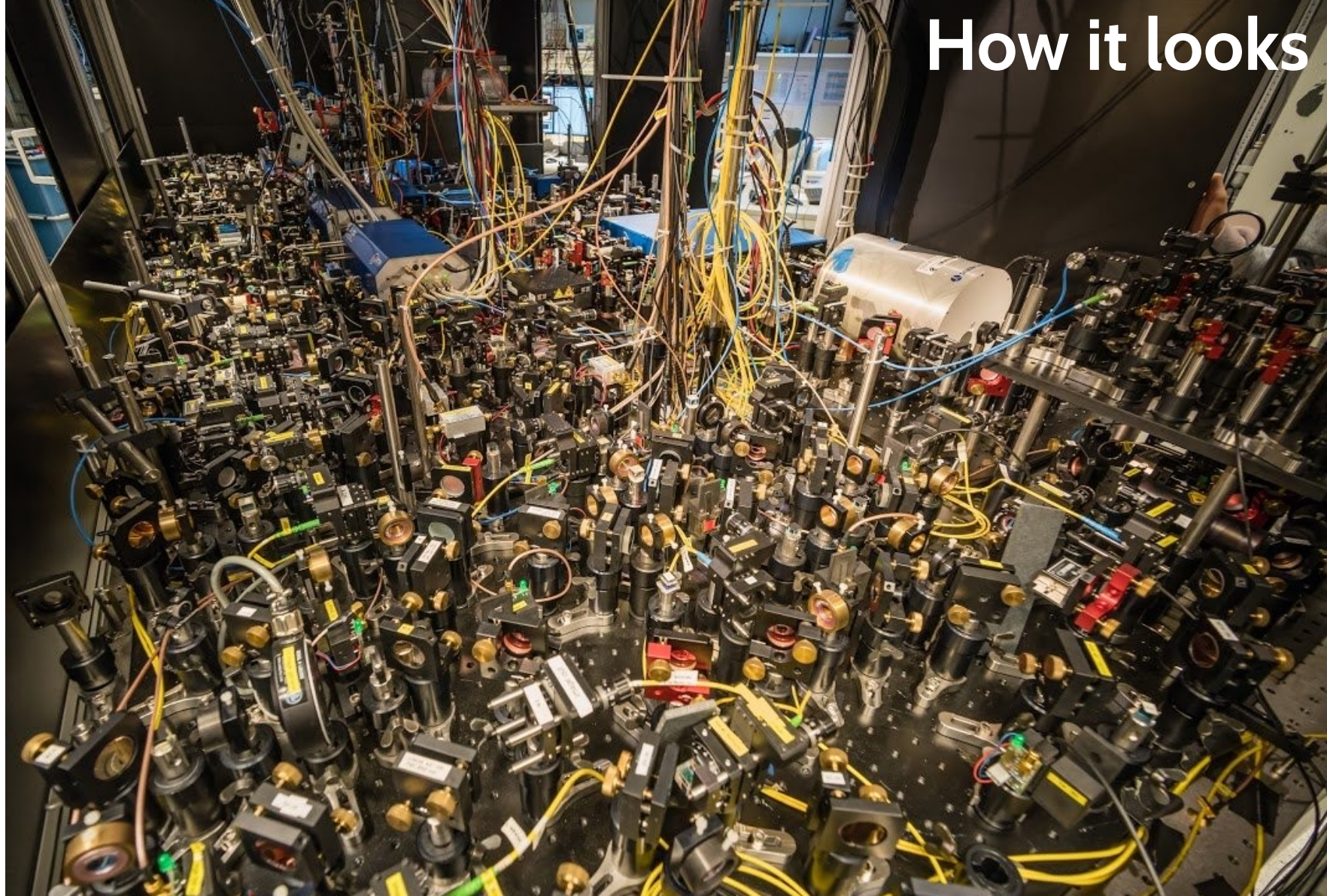
**Is the Hubbard model a high temperature SC?**



# Quantum simulation with cold atoms



How it looks



A blue-tinted close-up photograph of a printed circuit board (PCB). The image shows a dense network of fine, light-colored traces on a dark blue substrate. Several larger components, including integrated circuits and connectors, are visible. The text "And solid state ?" is overlaid in the center in a white, sans-serif font. The overall aesthetic is technical and futuristic.

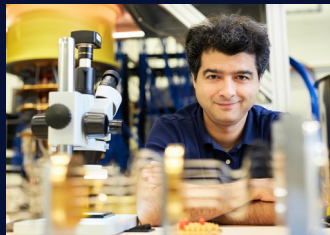
**And solid state ?**

# “Many”-body localization in a superconducting quantum simulator

**Ben Chiaro**



**Pedram Roushan**



**Annabelle Bohrdt**



**Michael Knap**

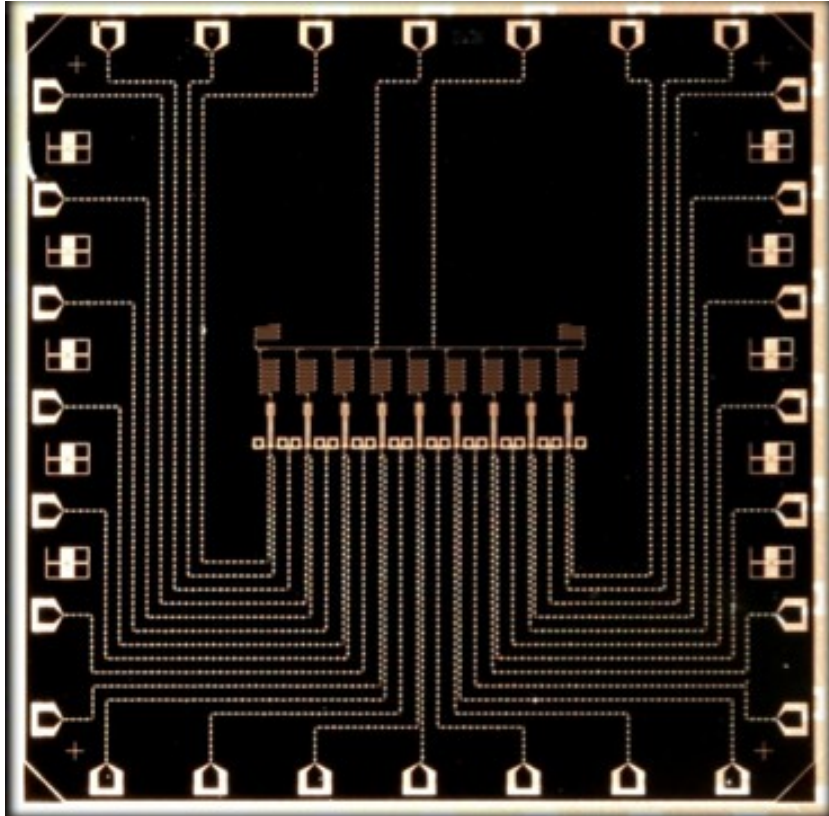


**Sarang Golapakrishnan Dmitry A. Abanin**



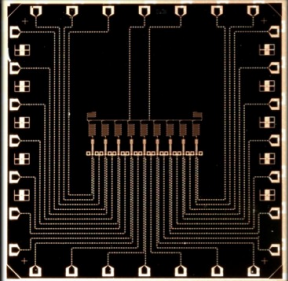


# Quantum simulation in a 9-Qubit system



**TODAY'S QUESTION:**

**Probing the exotic dynamics of  
disordered interacting systems  
(Many-Body Localization)**



UCSB

$\langle G|oogl|e \rangle$

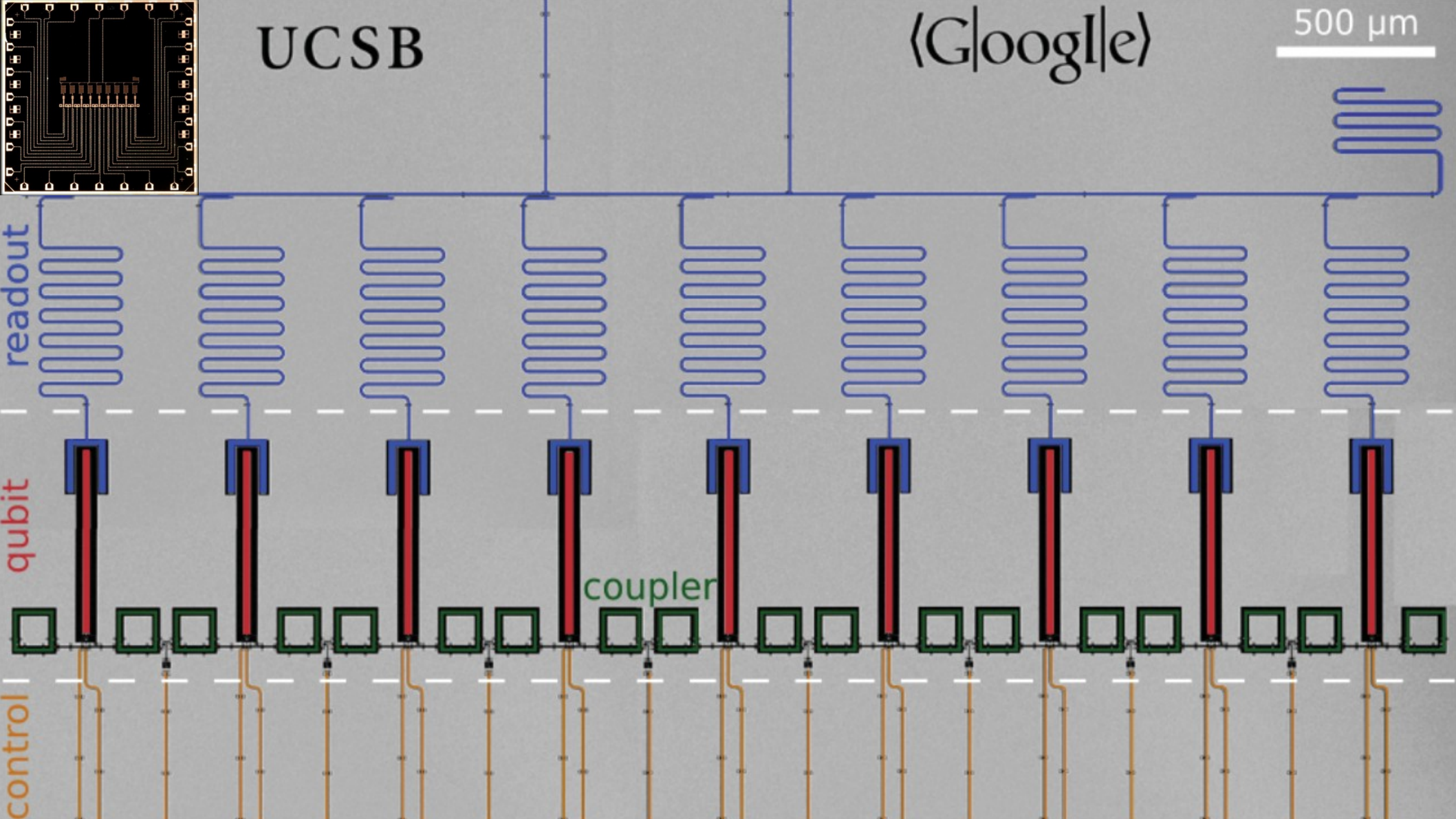
500  $\mu\text{m}$

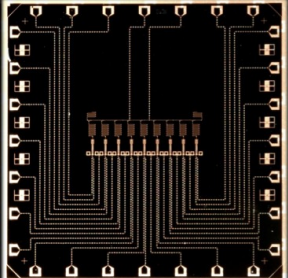
readout

qubit

coupler

control





UCSB

$\langle G|oogl|e \rangle$

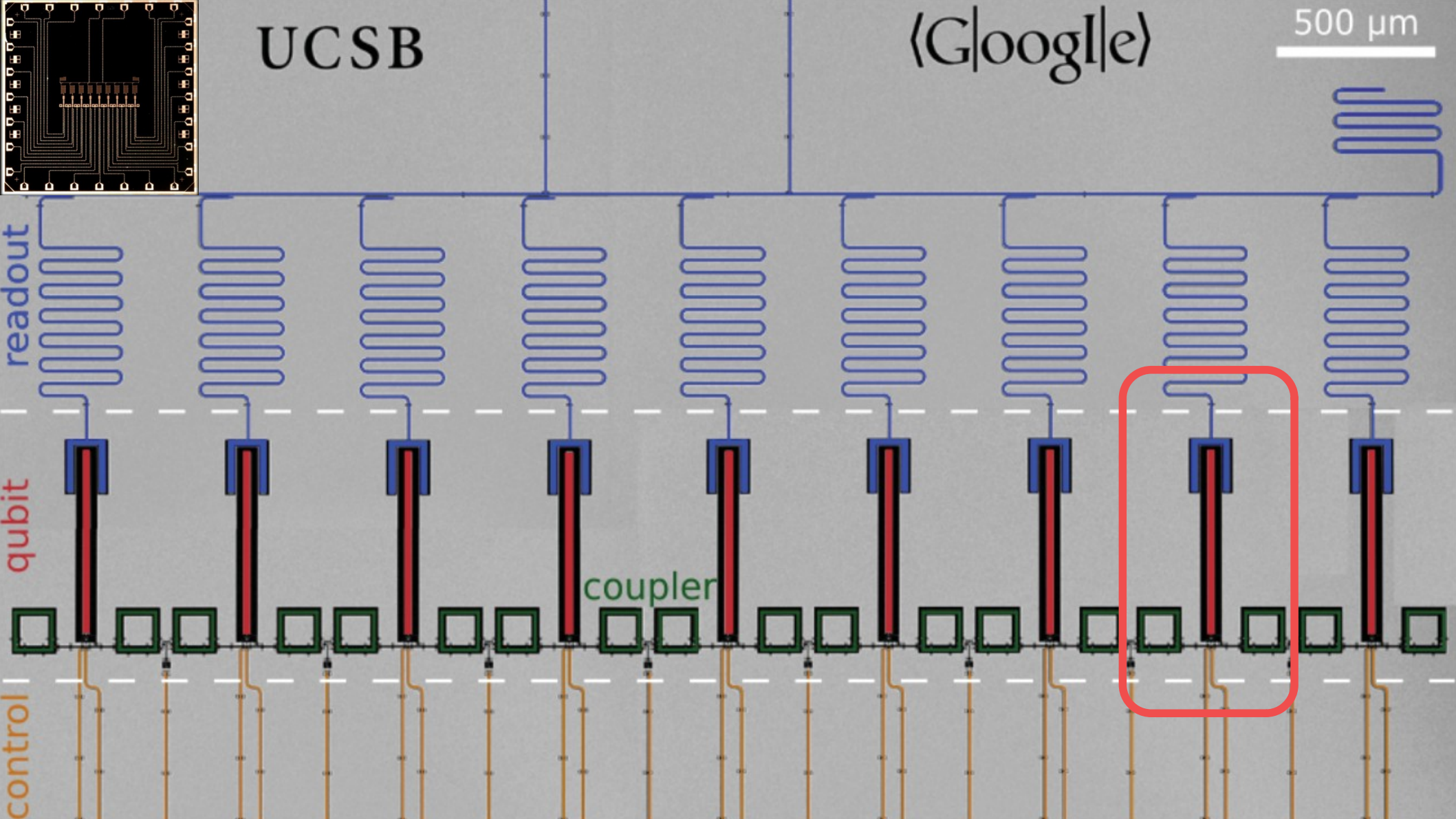
500  $\mu\text{m}$

readout

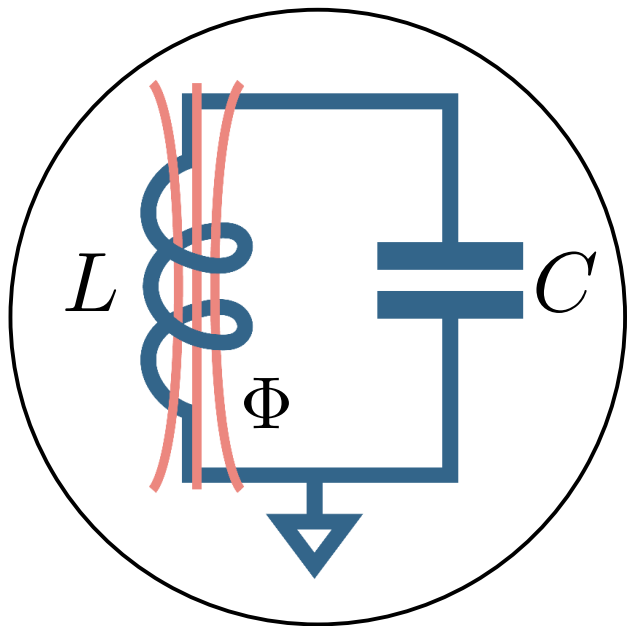
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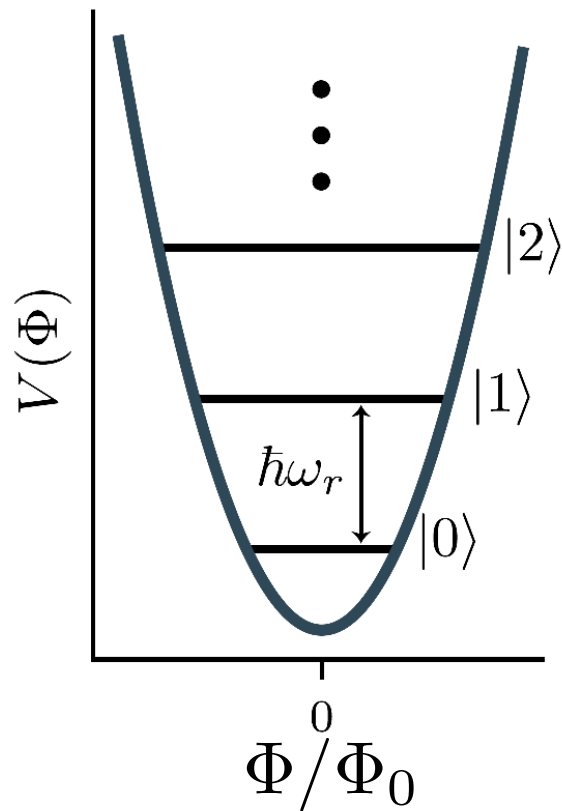


# Building artificial atoms with superconductors

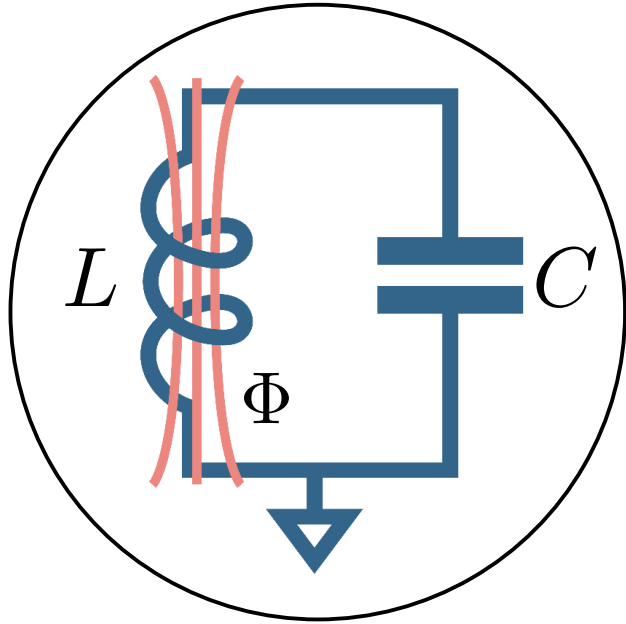


$$H_{LC} = \frac{Q^2}{2C} + \frac{1}{2} C \omega_r^2 \Phi^2$$

$$\omega_r = \frac{1}{\sqrt{LC}} \quad [\Phi, Q] = i\hbar$$

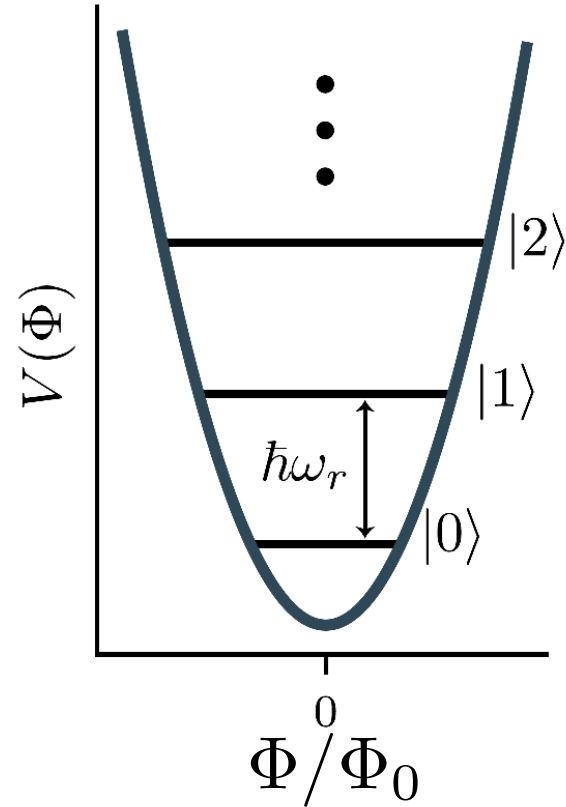


# Building artificial atoms with superconductors

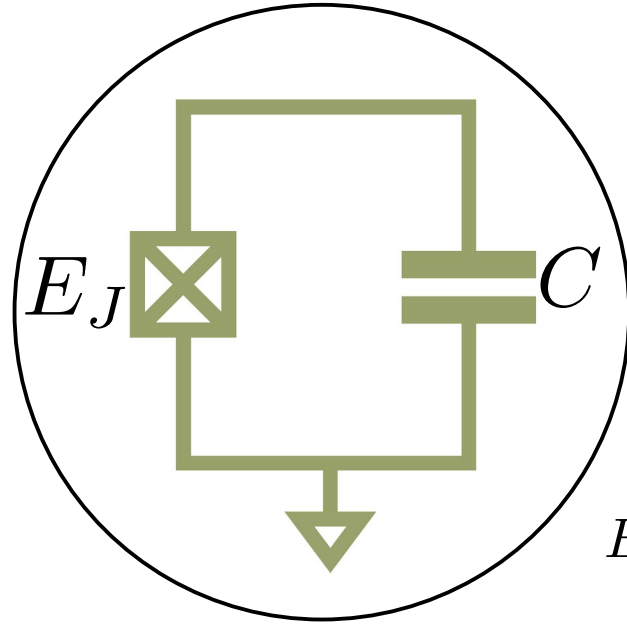


$$H_{LC} = \hbar\omega_r a^\dagger a = \hbar\omega_r n$$

$$\omega_r = \frac{1}{\sqrt{LC}} \quad [a, a^\dagger] = 1$$

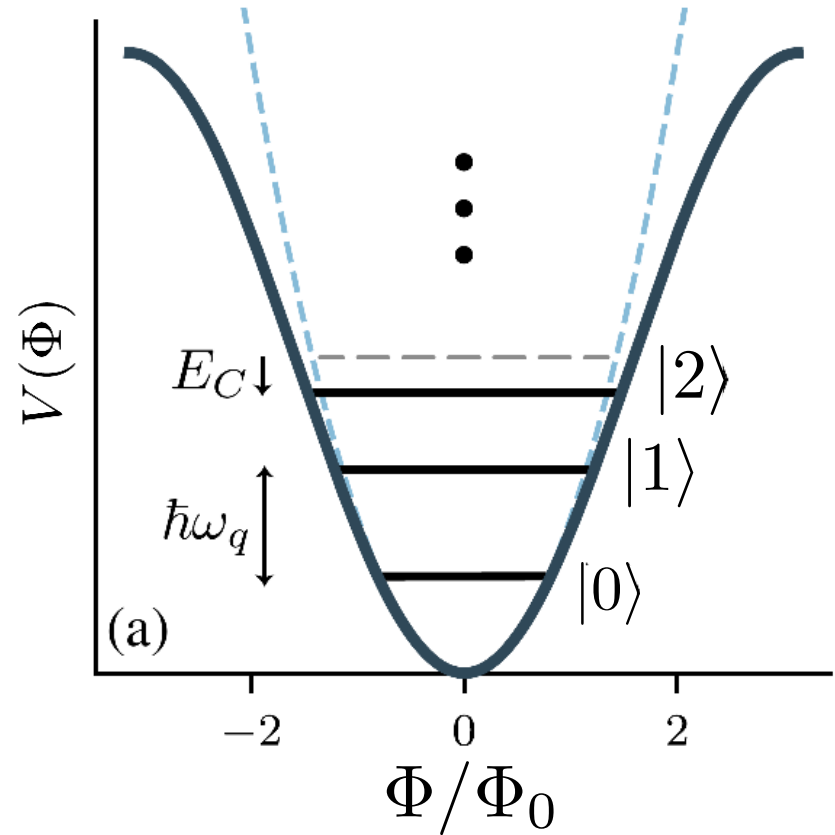


# Building artificial atoms with superconductors

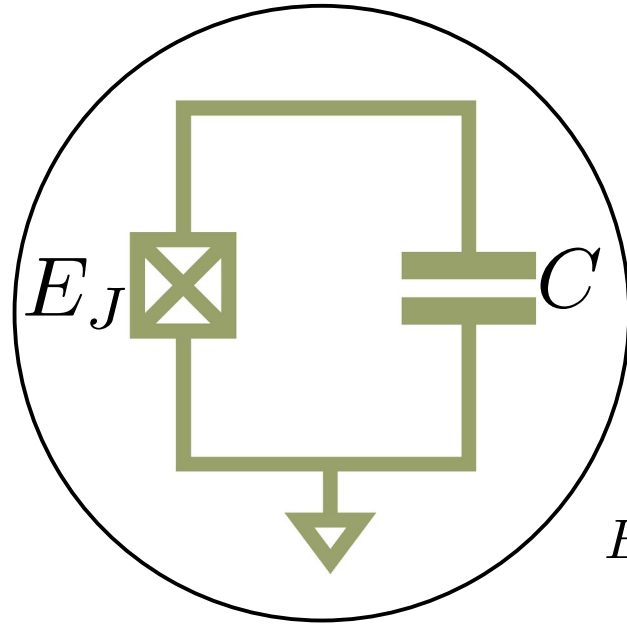


$$E_J = \frac{\Phi_0 I_c}{2\pi}$$

$$H = \frac{Q^2}{2C} - E_J \cos \left( 2\pi \frac{\Phi}{\Phi_0} \right)$$

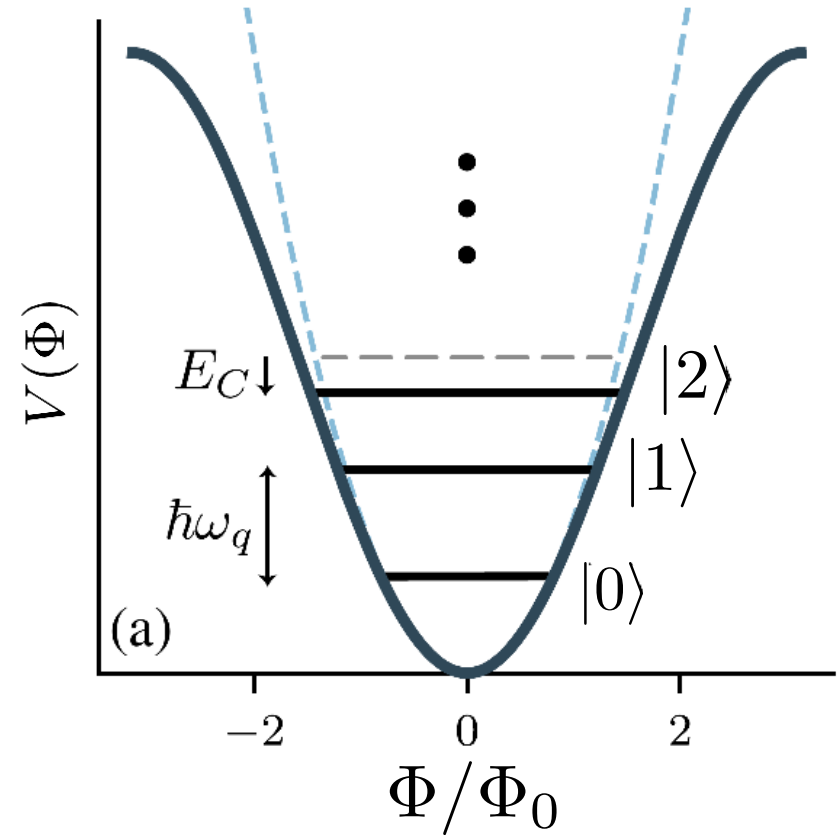


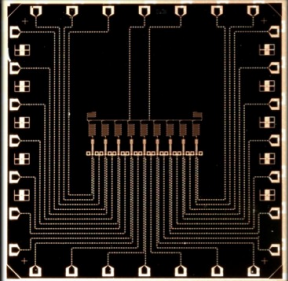
# Building artificial atoms with superconductors



$$E_J = \frac{\Phi_0 I_c}{2\pi}$$

$$H \simeq \hbar\omega_r n + \frac{U}{2} n(n-1)$$





UCSB

$\langle G|oogl|e \rangle$

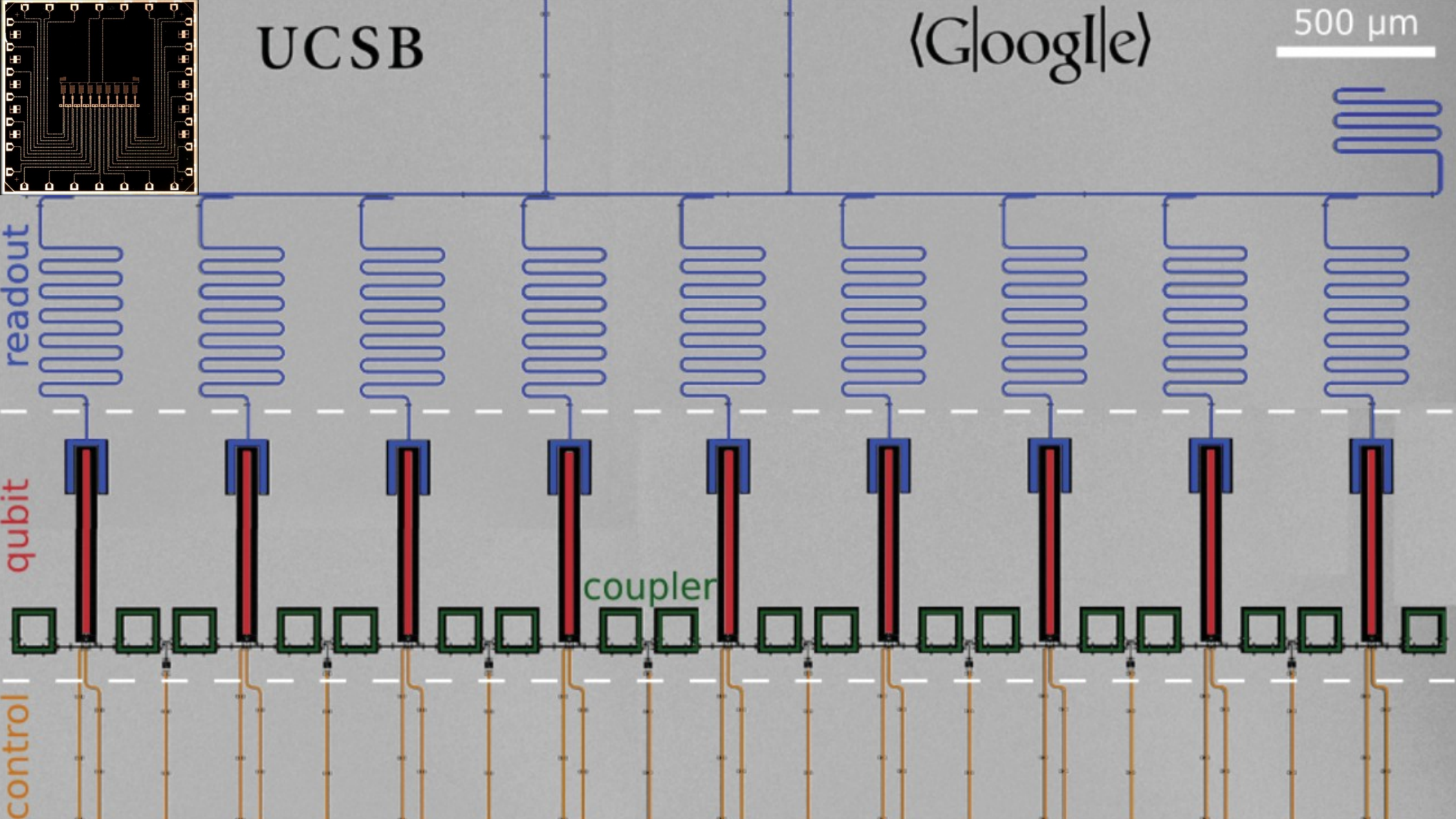
500  $\mu\text{m}$

readout

qubit

coupler

control







UCSB

$\langle \text{G}|\text{oogl}|e \rangle$

500  $\mu\text{m}$

Superconducting circuits simulate  
Hubbard-like interacting bosons

$$\mathcal{H} = \varepsilon n + \frac{U}{2} n(n-1)$$

Qubits are made out of the  
two bosonic Fock states

$$|n\rangle = |0\rangle, |1\rangle$$

readout

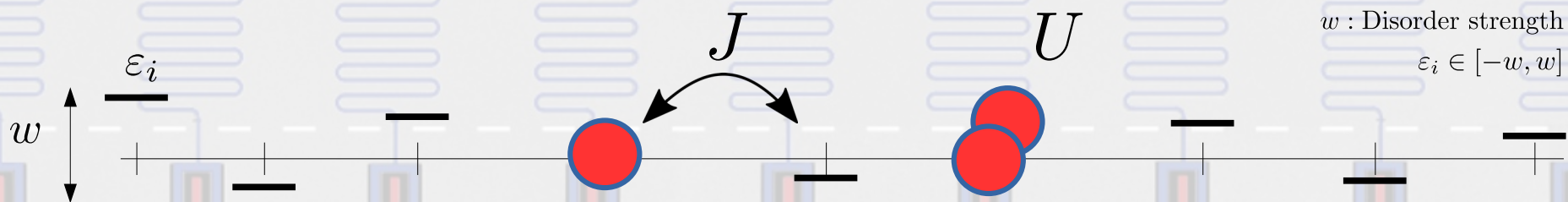
qubit

coupler

control

## Bose-Hubbard model

$$\mathcal{H}_{\text{BH}} = \sum_i^{N_Q} \varepsilon_i n_i + J \sum_{\langle i,j \rangle} (a_i^\dagger a_j + \text{h.c.}) + \frac{U}{2} \sum_i^{N_Q} n_i (n_i - 1)$$



## Quantum evolution of controlled initial states

$$|\Psi_0\rangle = |1\rangle \otimes |0\rangle \otimes |1\rangle \otimes |0\rangle \otimes \dots$$

$$|\Psi_0\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}} \otimes |0\rangle \otimes |1\rangle \otimes |0\rangle \otimes \dots$$

# Absence of Diffusion in Certain Random Lattices

P. W. ANDERSON

*Bell Telephone Laboratories, Murray Hill, New Jersey*

(Received October 10, 1957)



# Absence of Diffusion in Certain Random Lattices

P. W. ANDERSON

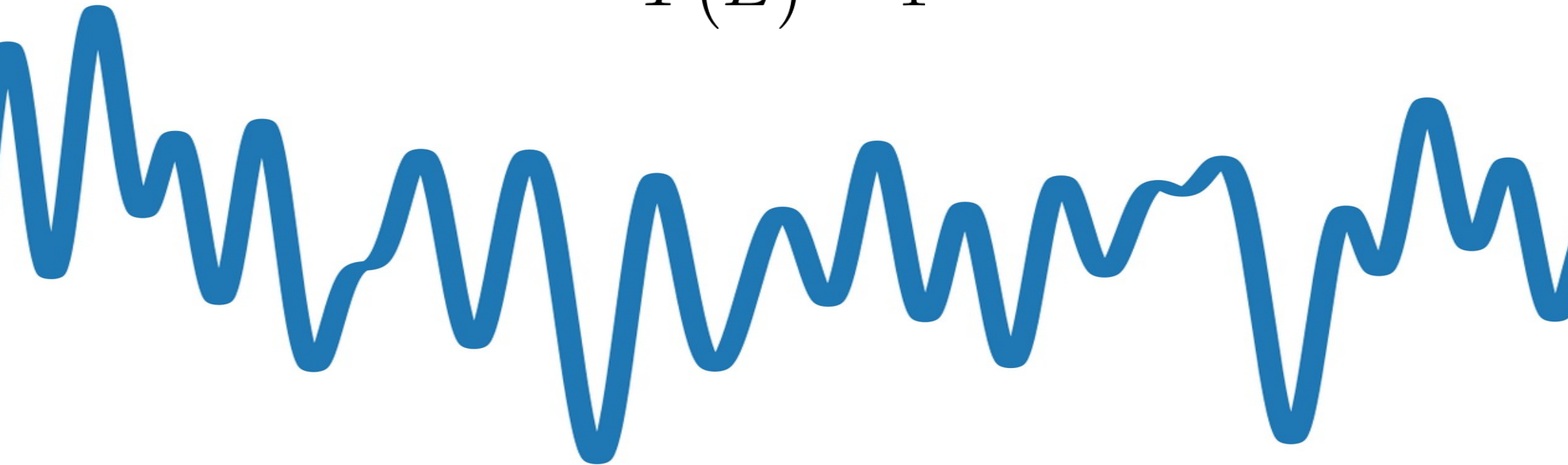
*Bell Telephone Laboratories, Murray Hill, New Jersey*

(Received October 10, 1957)



*Classical transmission probability*

$$T(E) = 1$$



# Absence of Diffusion in Certain Random Lattices

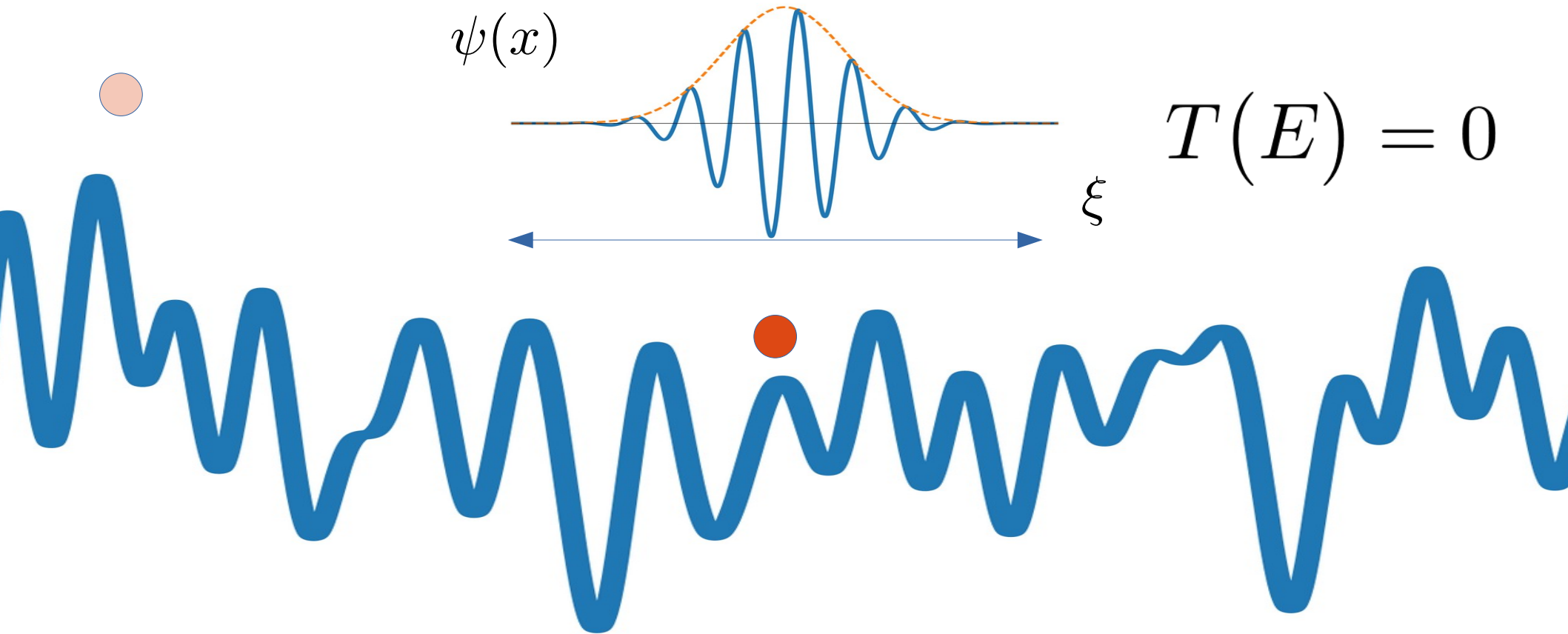
P. W. ANDERSON

*Bell Telephone Laboratories, Murray Hill, New Jersey*

(Received October 10, 1957)



**Quantum interference leads to particle localization**





# Is localization stable in the presence of interactions?

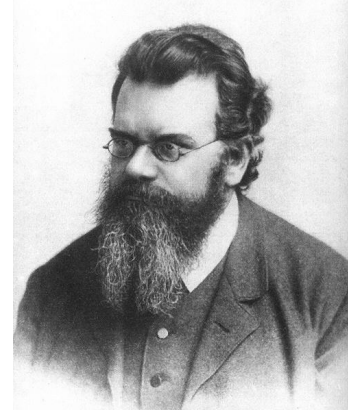
Giamarchi, Schulz (1988)

Basko, Aleiner, Altshuler, *Annals of Physics* (2006)

Pal, Huse, *PRB* (2010)

# Diffusion in interacting systems

*“Interacting systems explore all (micro) states allowed by energy conservation, they forget about their initial conditions and thermalize”*

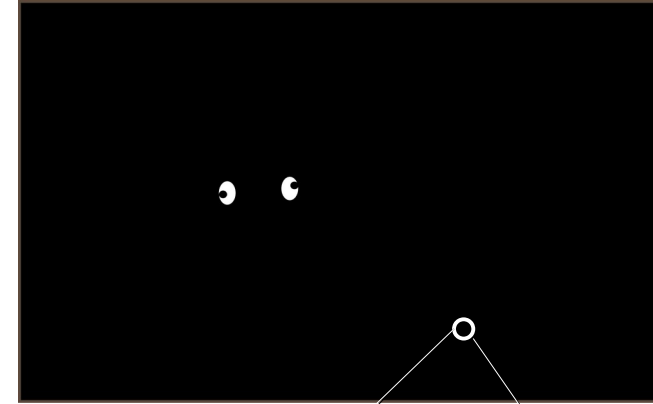
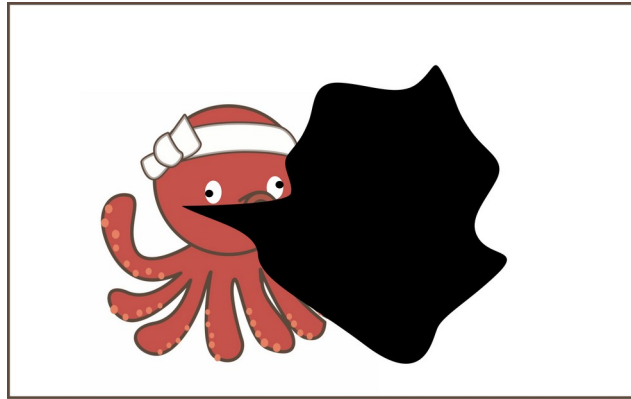
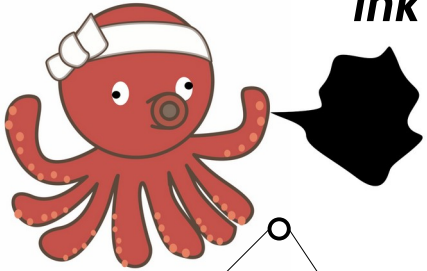


Ludwig E. Boltzmann

## Octopus in a box (Diffusion)

Octopus

ink



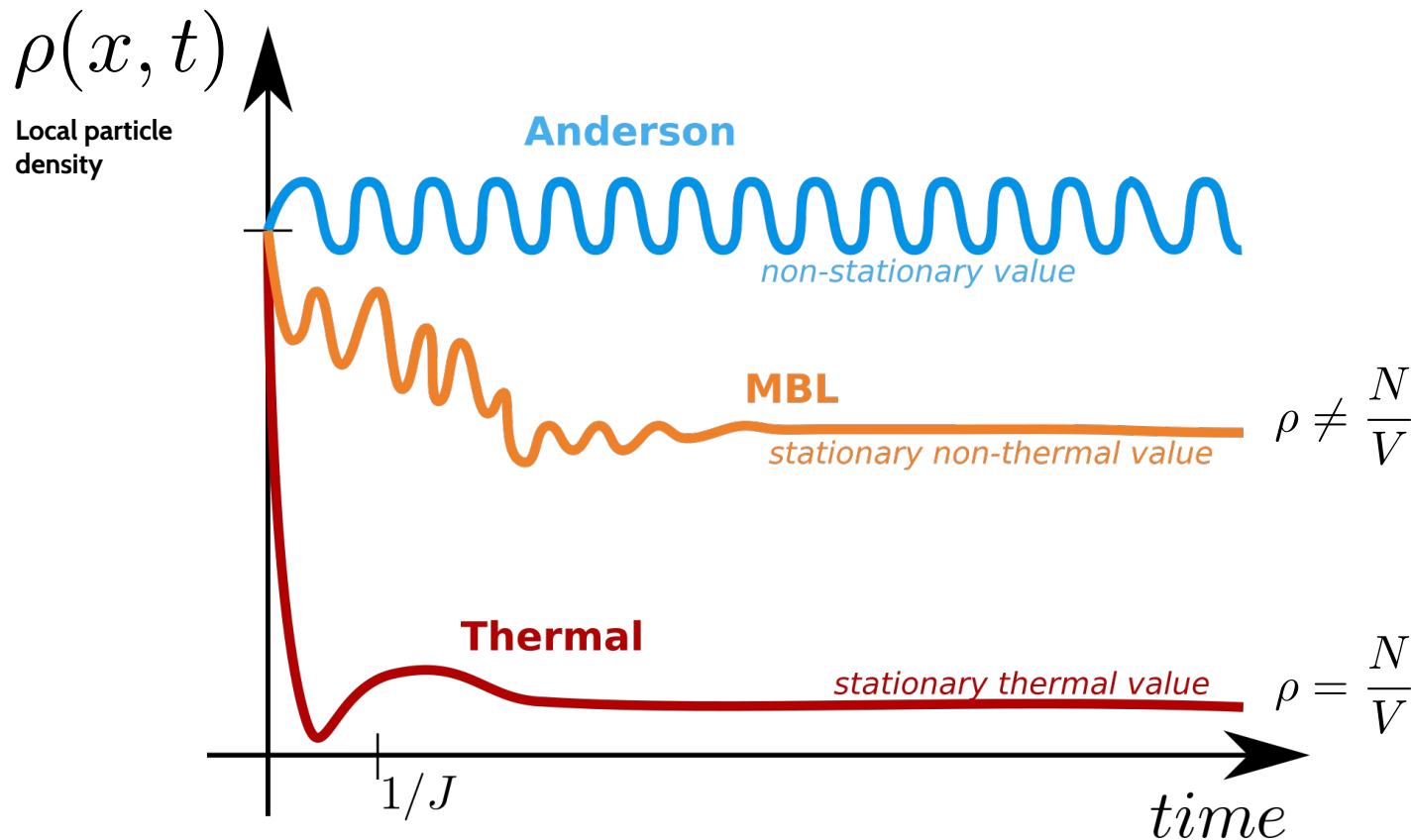
$$\rho \neq \frac{N}{V}$$

time

$$\rho = \frac{N}{V}$$

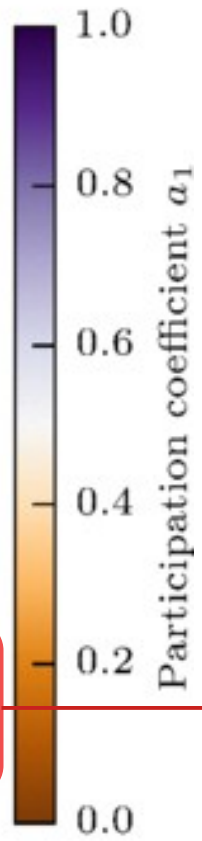
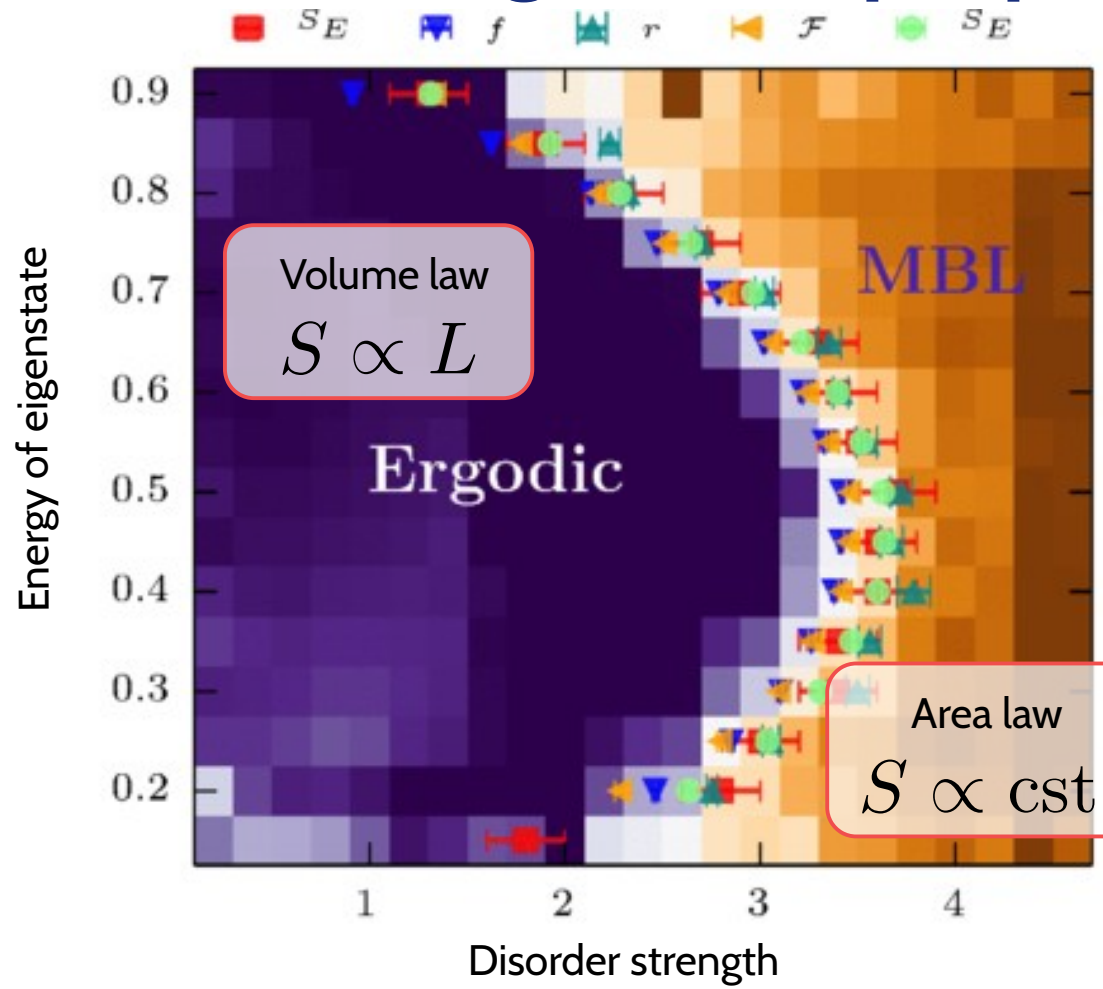
# Many-Body Localization (MBL)

MBL systems relax towards “non-thermal” states and keep memory about the initial state





# Exotic entanglement properties



**Quantifying entanglement**

**Benoît's talk !**

System  $A$  and Environment  $B$  are shown as a chain of sites.

- Reduced density matrix  $\rho_A = \text{Tr}_B(\rho)$
- Entanglement measures (pure states)
  - von-Neumann entropy  $S_A = -\text{Tr}_A[\rho_A \log \rho_A]$
  - Rényi entropy  $S_A^{(n)} = \frac{1}{1-n} \log \text{Tr}_A[\rho_A^n] \leq S_A$
- Entanglement entropies are entanglement monotones for pure states (cannot increase under local operation)

The many-body eigenstates are well described by Matrix Product States (MPS)

**Cécilia's talk !**

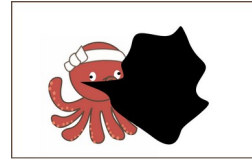
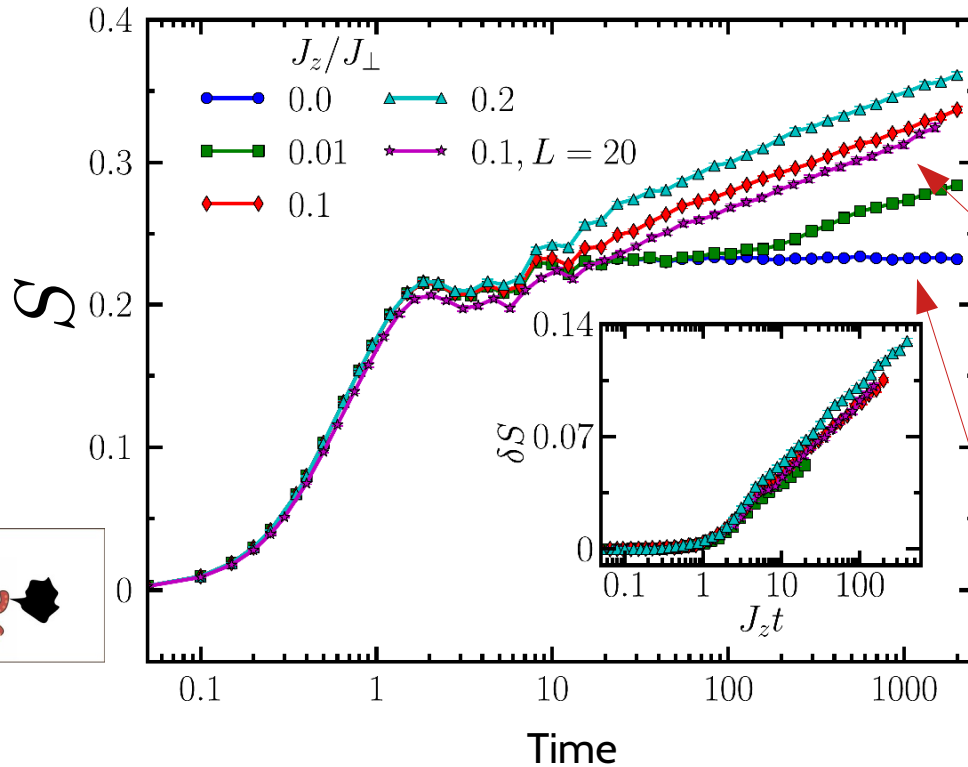
**Perspective: what about estimating the typical amount of entanglement?**

**Observation:** The amount of bipartite entanglement in a TNS can be upper bounded in terms of the bond dimension. Indeed, for any subregion having  $L$  boundary edges and  $V$  bulk vertices, its entanglement entropy is at most  $L \log q$ , which is usually much smaller than  $V \log d$ .

- entropy of reduced state  $\rightarrow$  area law (boundary dimension  $q^d$ )
- But what about the amount of genuinely multipartite entanglement?  $\rightarrow$  volume law (bulk dimension  $d^V$ )

Cécilia Lancien | Typical correlations and entanglement in random tensor network states | Journée théorie CPTGA - September 27 2021 | 19

# Exotic entanglement properties



Logarithmic entanglement growth

$$S \propto \ln(t)$$

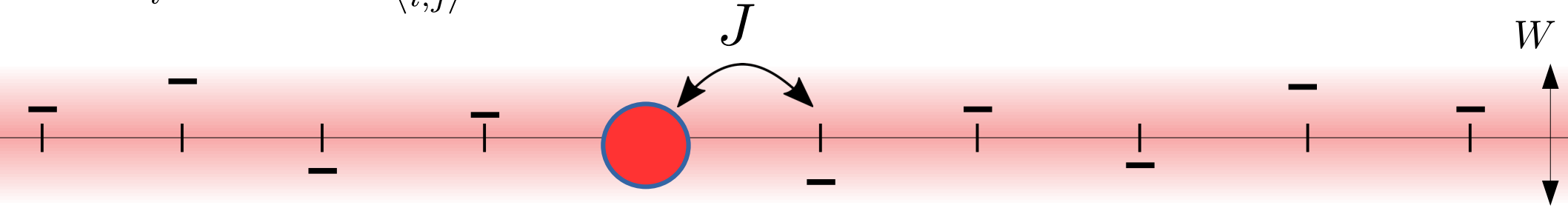
Anderson Insulator

# Emergent integrability : Local Integrals of Motion

$$H = \sum_i \varepsilon_i a_i^\dagger a_i + J \sum_{\langle i,j \rangle} (a_i^\dagger a_j + \text{h.c.})$$

$W$  : Disorder strength

$$\varepsilon_i \in [-W, W]$$

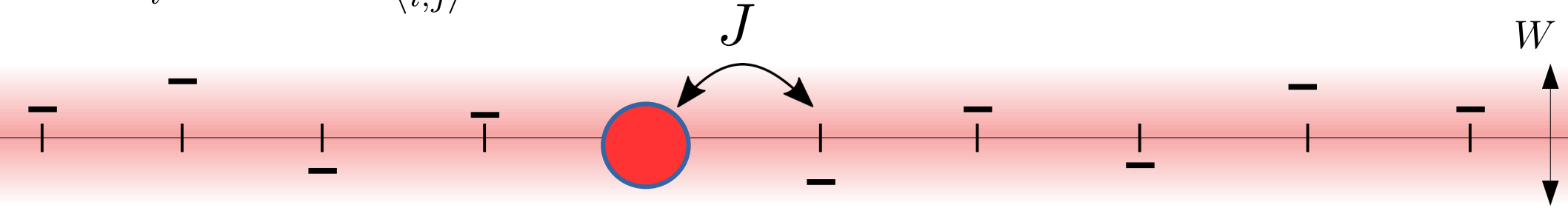


# Emergent integrability : Local Integrals of Motion

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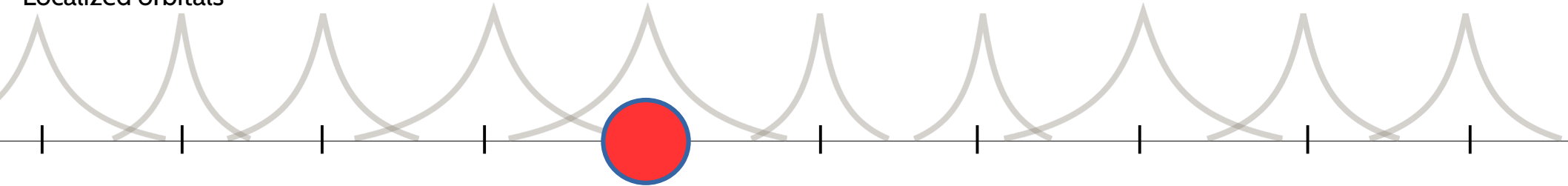
$$\varepsilon_i \in [-W, W]$$



Anderson localization

$$H = \sum_i \tilde{\varepsilon}_i \tilde{n}_i$$

Localized orbitals

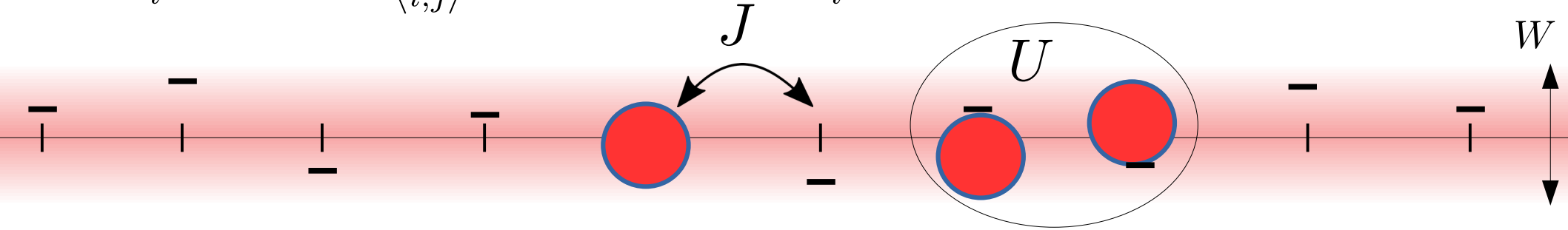


# Emergent integrability : Local Integrals of Motion

$$H = \sum_i \varepsilon_i a_i^\dagger a_i + J \sum_{\langle i,j \rangle} (a_i^\dagger a_j + \text{h.c}) + U \sum_i a_i^\dagger a_i a_{i+1}^\dagger a_{i+1}$$

$W$  : Disorder strength

$$\varepsilon_i \in [-W, W]$$



Anderson localization

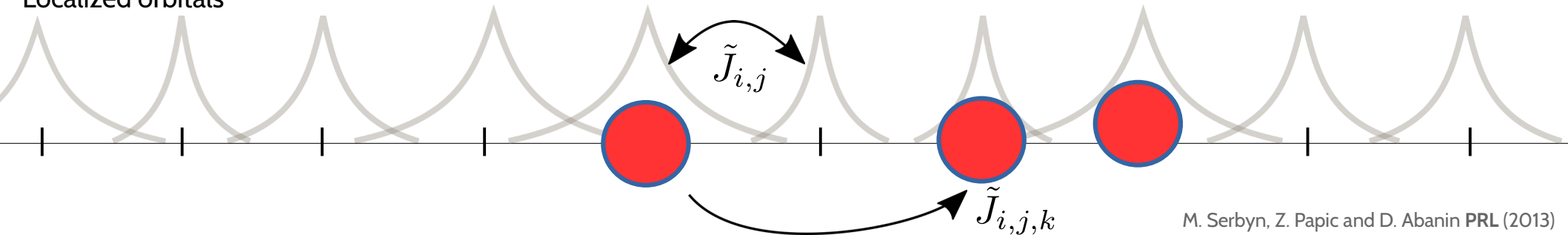
Non-local interactions

$$H = \sum_i \tilde{\varepsilon}_i \tilde{n}_i + \sum_{i,j} \tilde{J}_{i,j} \tilde{n}_i \tilde{n}_j + \sum_{i,j,k} \tilde{J}_{i,j,k} \tilde{n}_i \tilde{n}_j \tilde{n}_k + \dots$$

$\tilde{n}_i$  : Local Integral of Motion

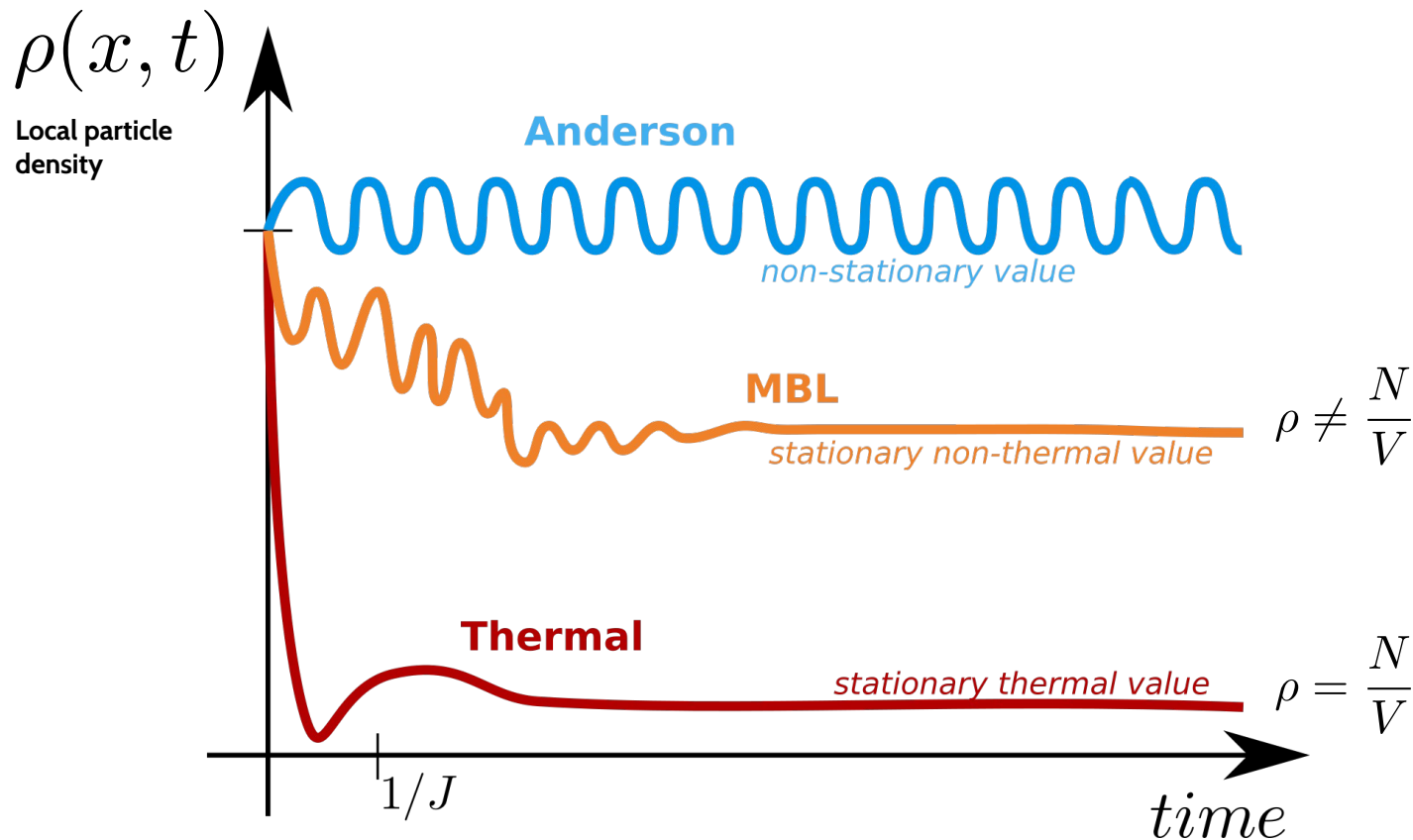
$$[\tilde{n}_i, H] = 0 \quad (\text{LIOM})$$

Localized orbitals



# Many-Body Localization (MBL)

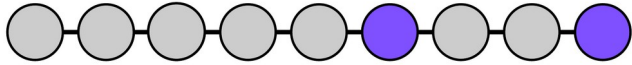
MBL systems relax towards “non-thermal” states and keep memory about the initial state



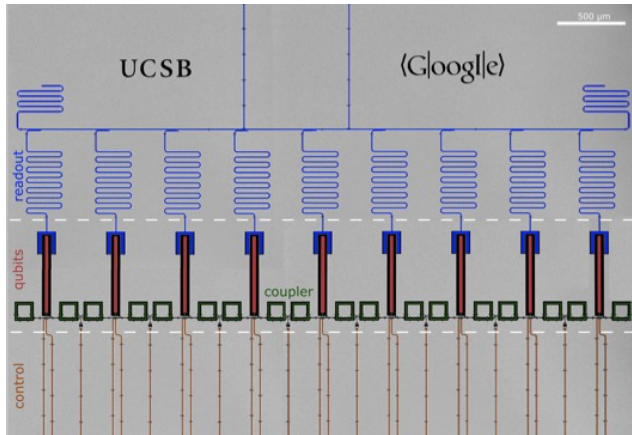
# Boson transport : Diffusion vs Localization

## Initial state

$$|\Psi_0\rangle = |0\rangle \otimes |0\rangle \otimes |0\rangle \otimes |0\rangle \otimes |0\rangle \otimes |1\rangle \otimes |0\rangle \otimes |0\rangle \otimes |1\rangle$$



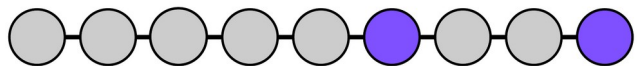
Measure density here



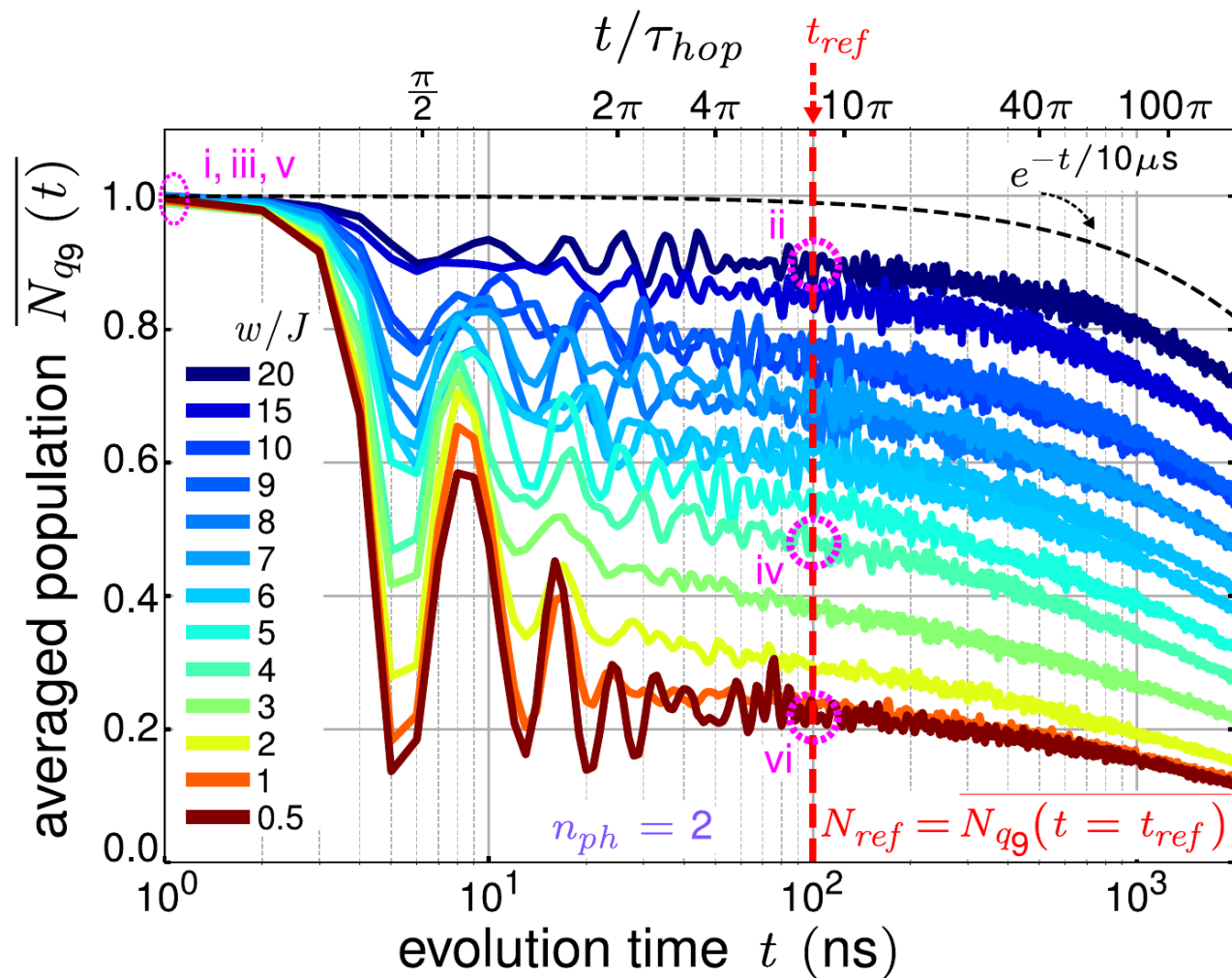
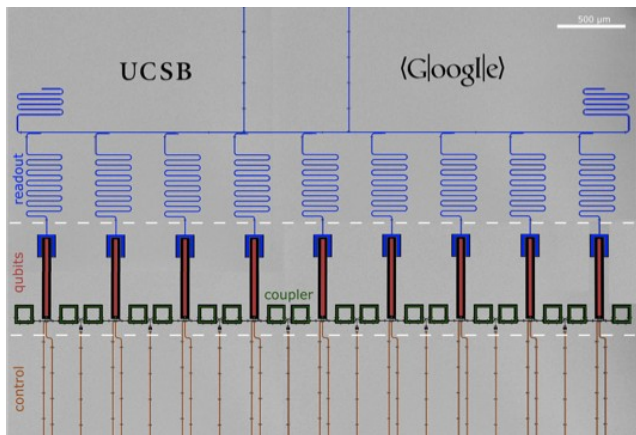
# Boson transport : Diffusion vs Localization

## Initial state

$$|\Psi_0\rangle = |0\rangle \otimes |0\rangle \otimes |0\rangle \otimes |0\rangle \otimes |0\rangle \otimes |1\rangle \otimes |0\rangle \otimes |0\rangle \otimes |1\rangle$$



Measure density here

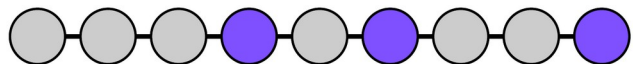




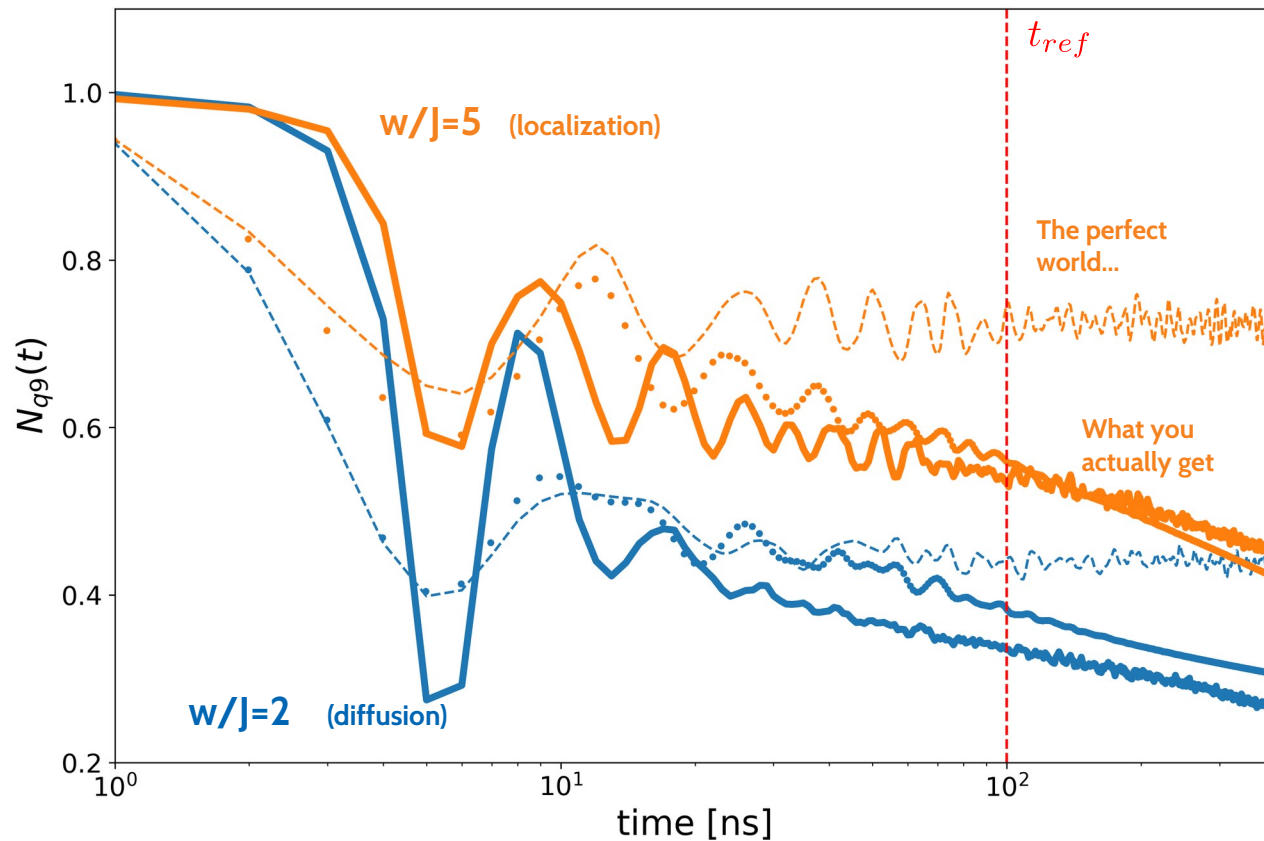
# The presence of dephasing and losses

## Initial state

$$|\Psi_0\rangle = |0\rangle \otimes |0\rangle \otimes |0\rangle \otimes |1\rangle \otimes |0\rangle \otimes |1\rangle \otimes |0\rangle \otimes |0\rangle \otimes |1\rangle$$

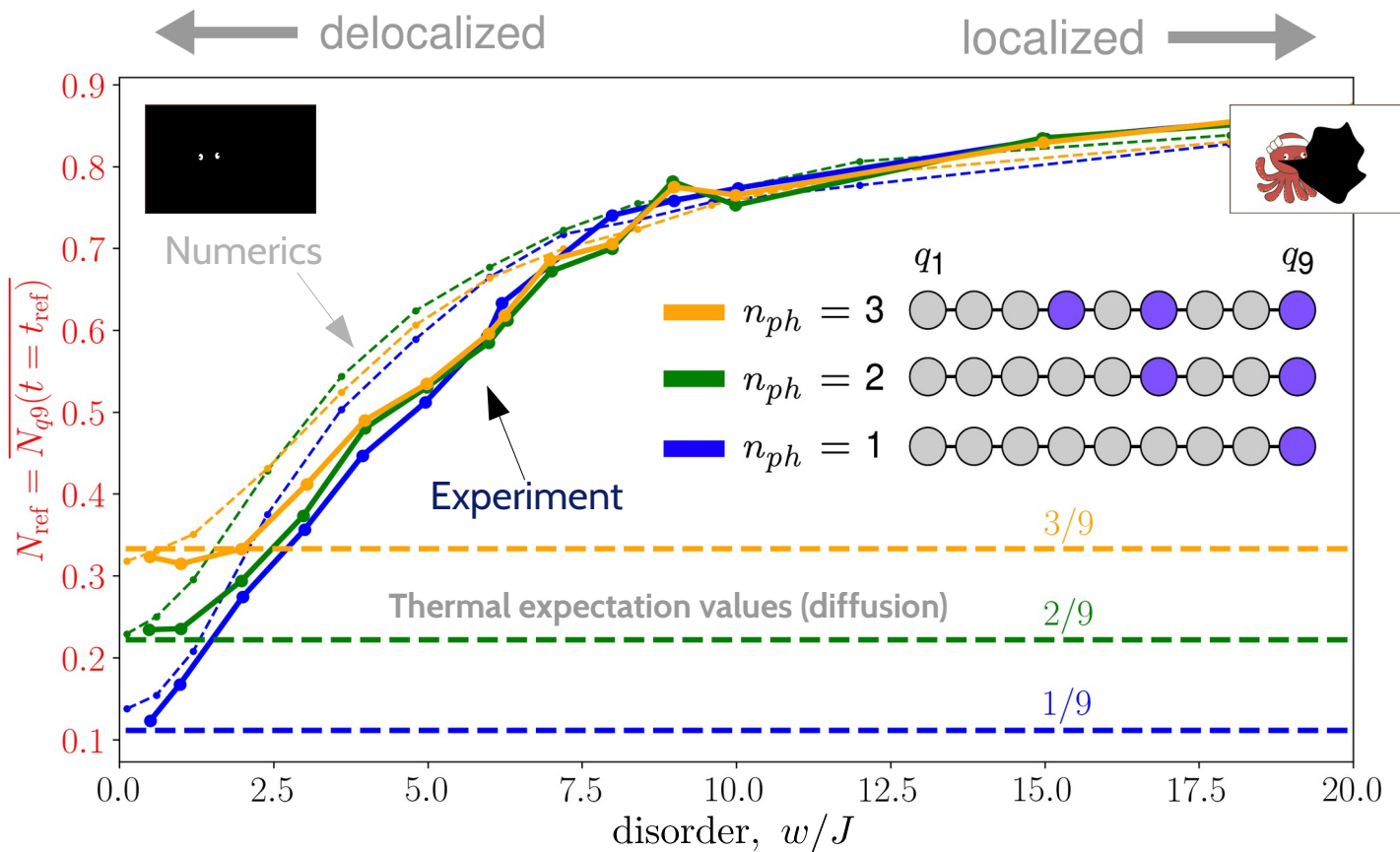
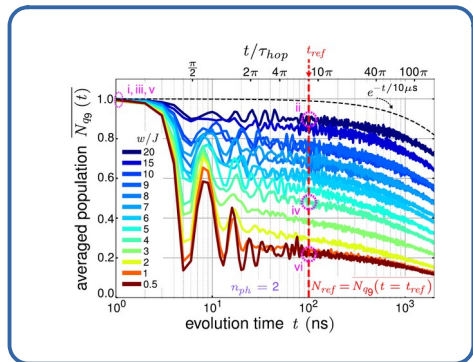


- Experimental Data  $w/J=2$
- Experimental Data  $w/J=5$
- - - Numerics isolated system
- - - Numerics isolated system
- Numerics with dissipation
- Numerics with dissipation



$$T_2 \sim 10\mu s$$

# Localization



# Slow dynamics in 2D

## Sycamore circuit

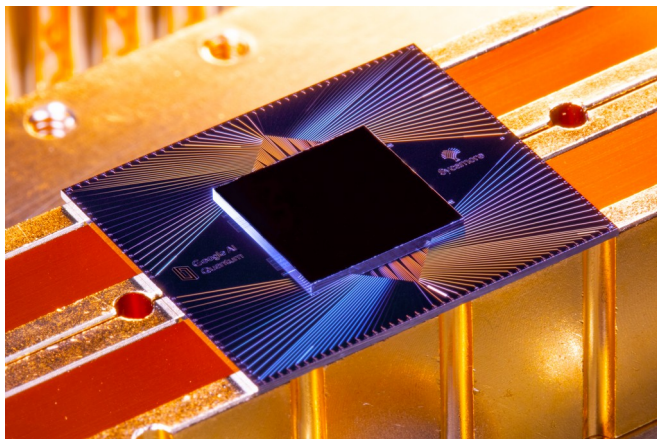
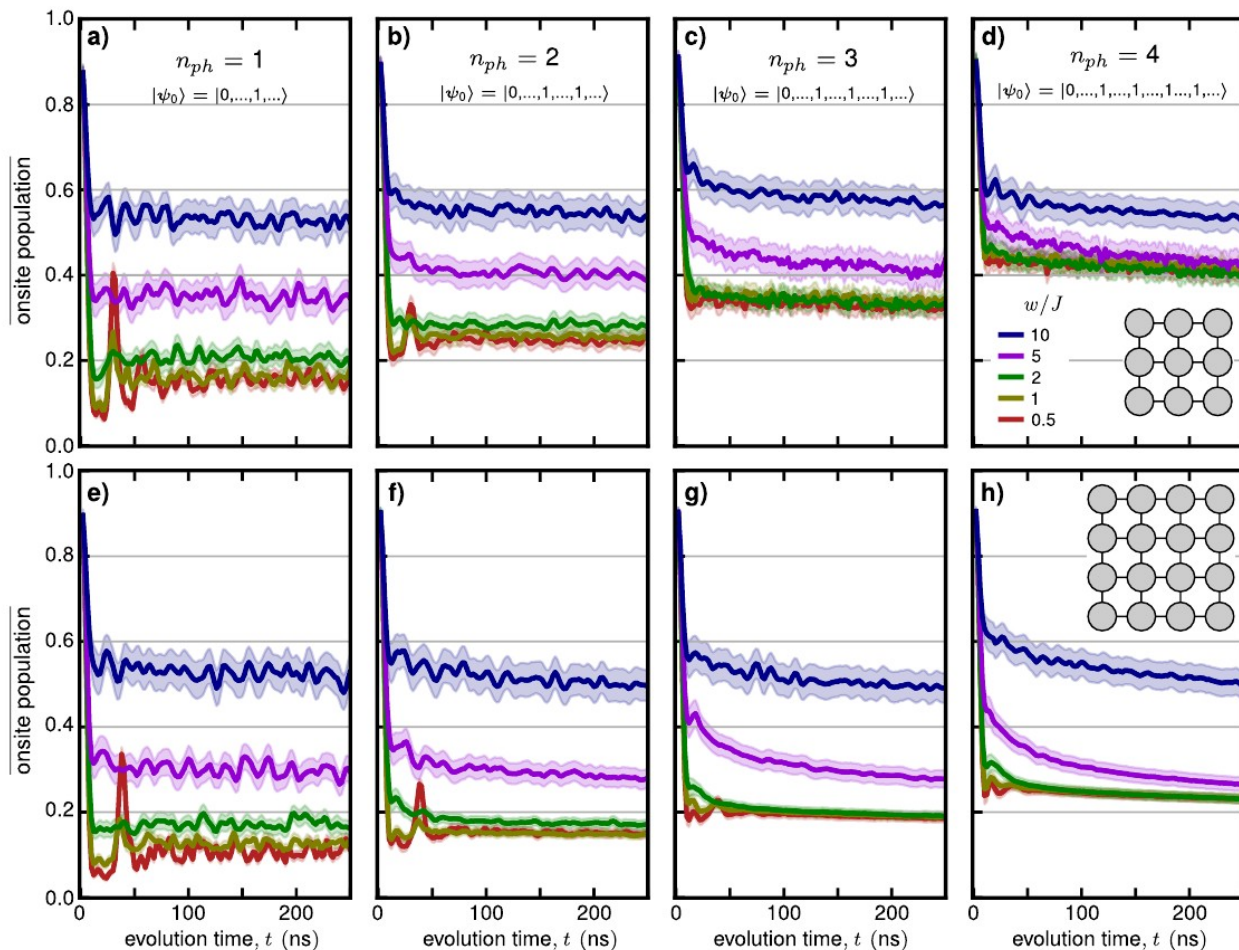


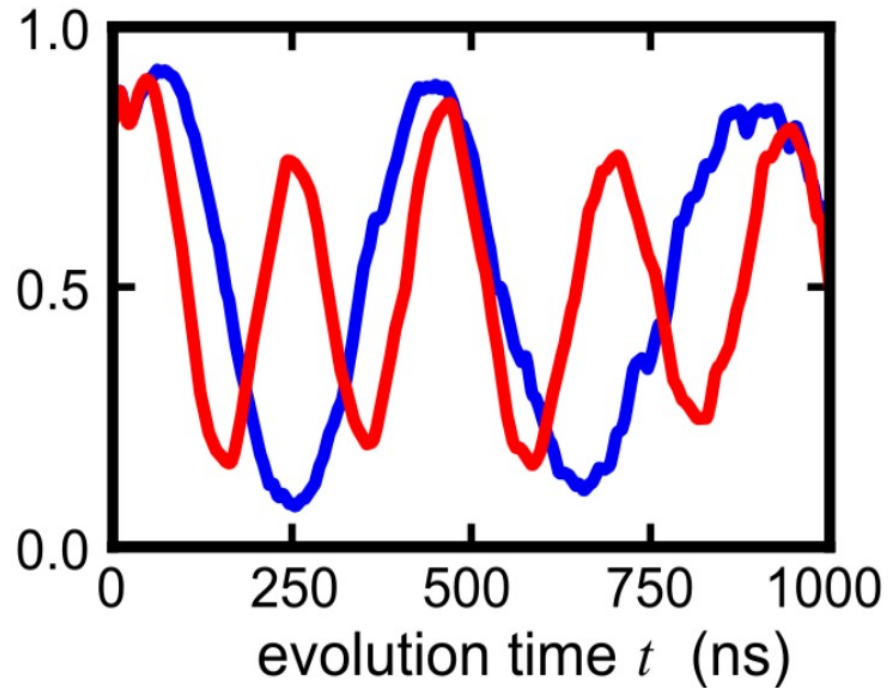
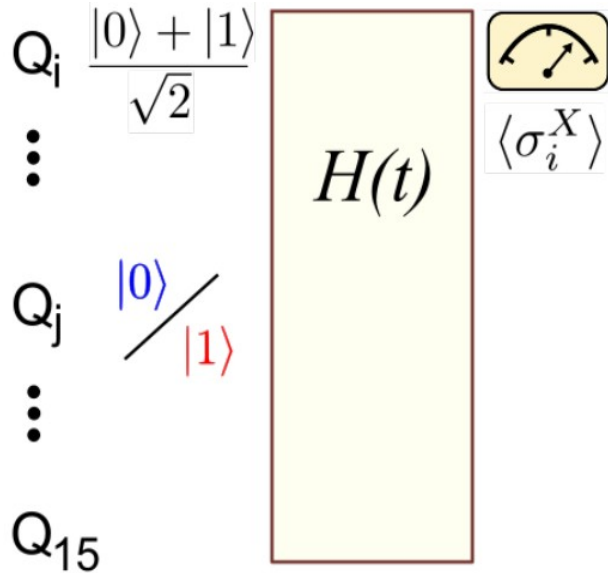
Photo credit: Eric Lucero

Arute, F. et al. "Quantum supremacy using a programmable superconducting processor", Nature (2019)



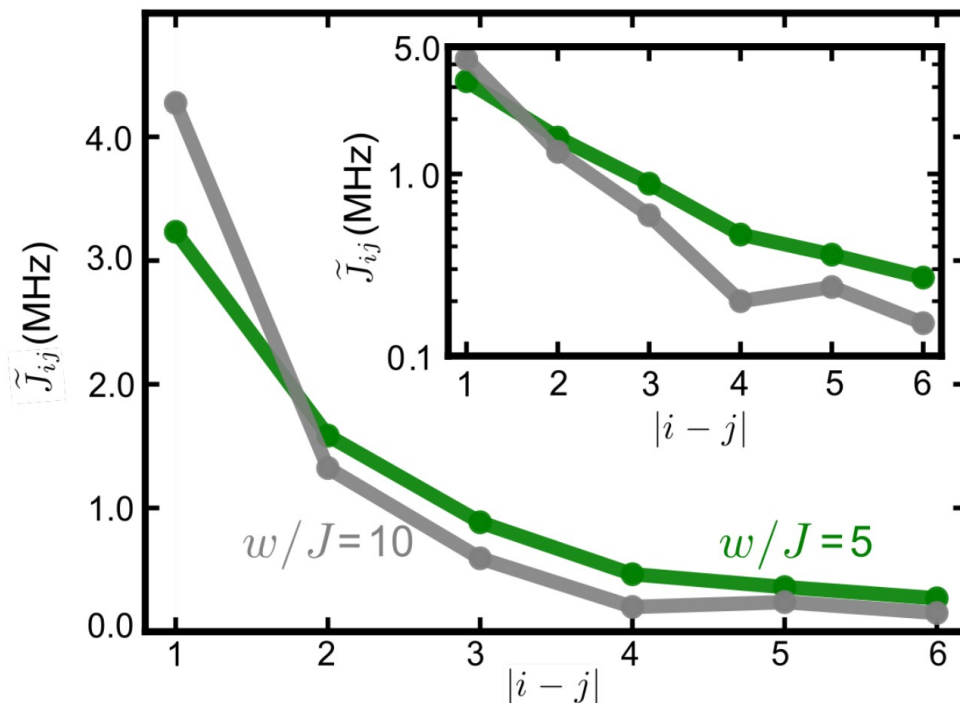
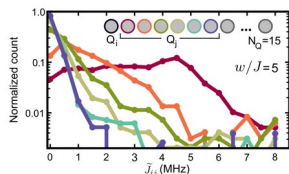
# Direct measurement of LIOMs

$$H = \sum_i \tilde{\epsilon}_i \tilde{n}_i + \sum_{i,j} \tilde{J}_{i,j} \tilde{n}_i \tilde{n}_j + \sum_{i,j,k} \tilde{J}_{i,j,k} \tilde{n}_i \tilde{n}_j \tilde{n}_k + \dots$$



# Exponential decay with distance

$$H = \sum_i \tilde{\epsilon}_i \tilde{n}_i + \sum_{i,j} \tilde{J}_{i,j} \tilde{n}_i \tilde{n}_j + \sum_{i,j,k} \tilde{J}_{i,j,k} \tilde{n}_i \tilde{n}_j \tilde{n}_k + \dots$$



# Growth and preservation of entanglement

## Sycamore circuit

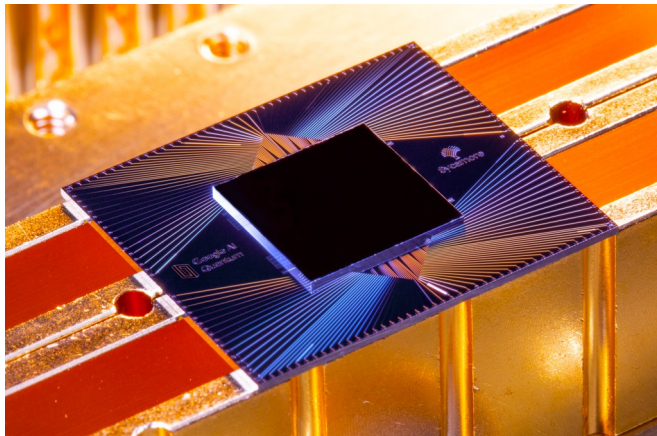
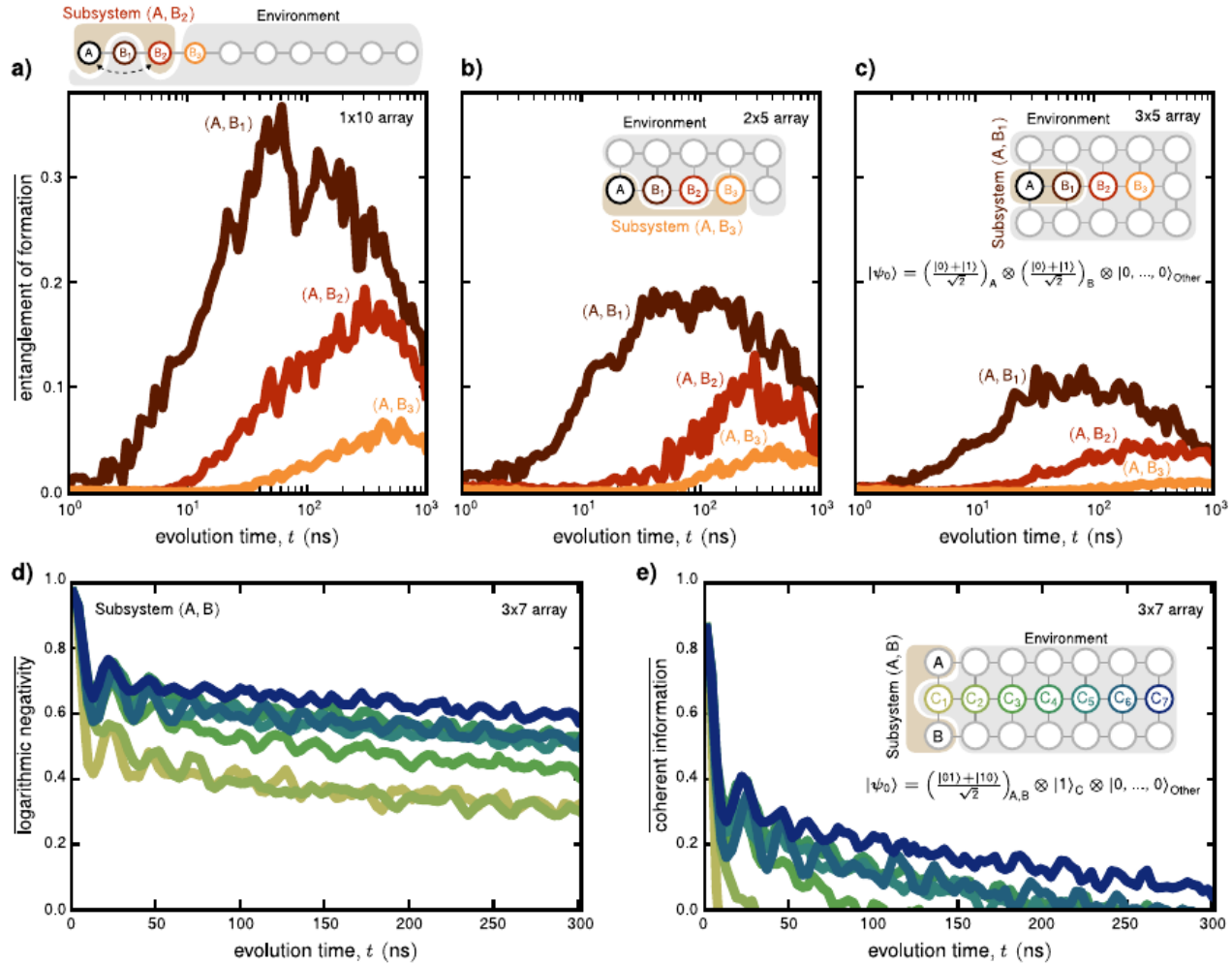


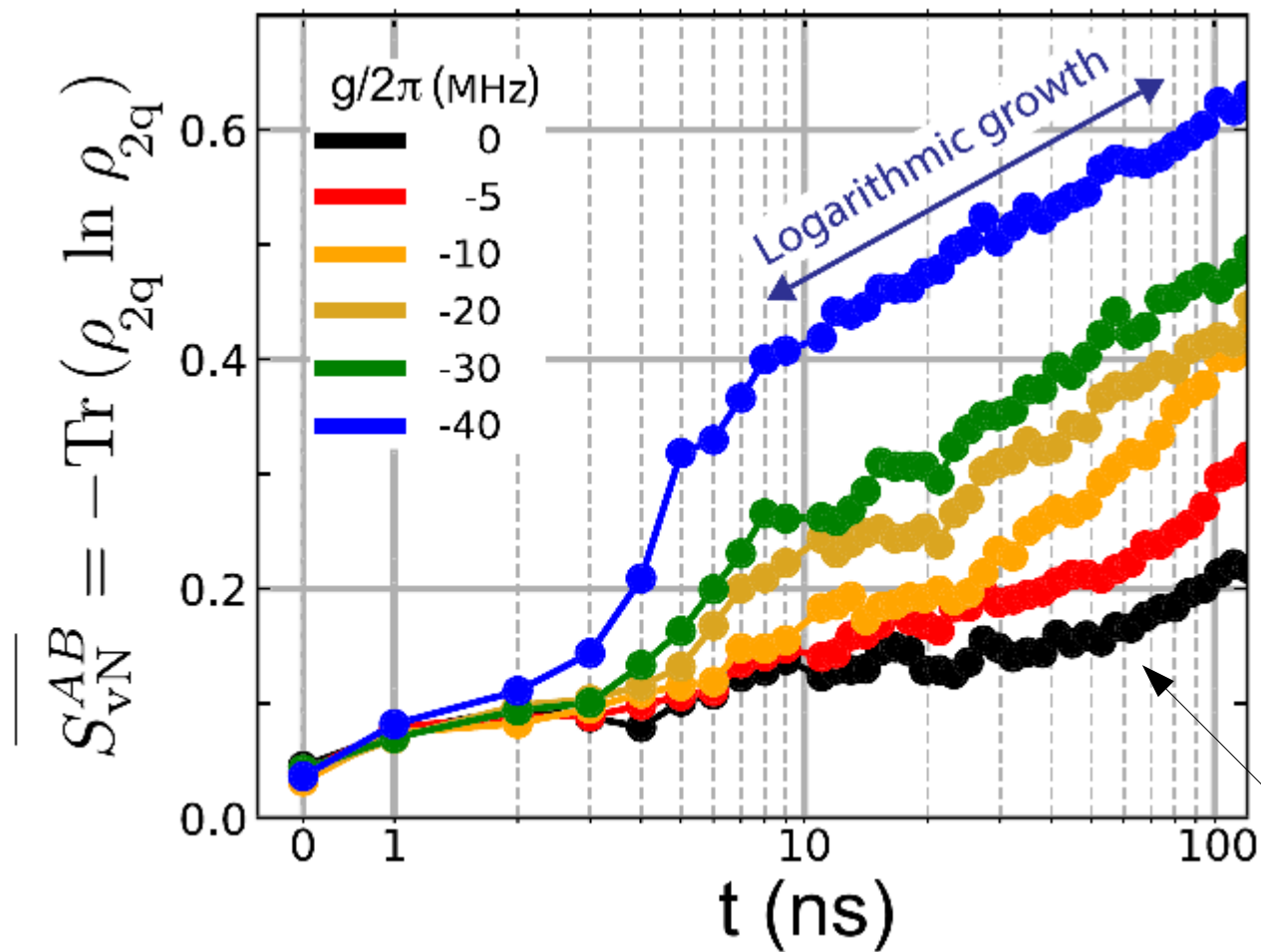
Photo credit: Eric Lucero

Arute, F. et al. "Quantum supremacy using a programmable superconducting processor", Nature (2019)

$$w/J = 10$$



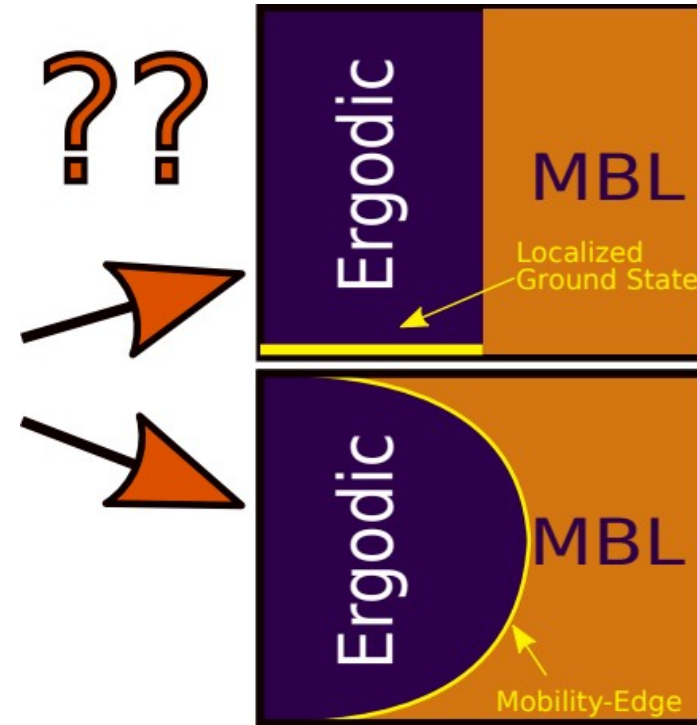
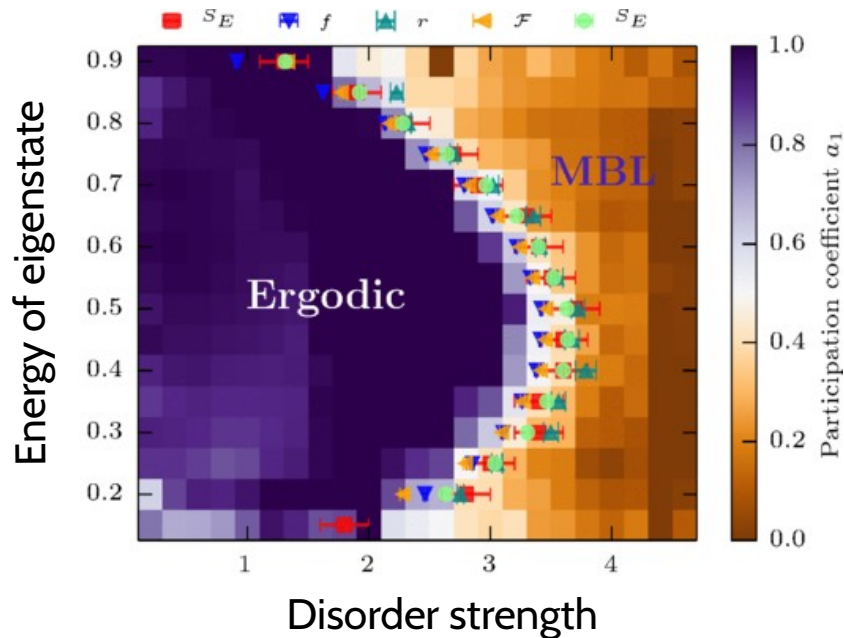
# Watch out the entanglement entropy !



!!!!!!

# Open questions about MBL

Ergodic-to-localized transition and mobility edge





# Nucleation of ergodicity by a single impurity in supercooled insulators

Ulrich Krause

Théo Pellegrin

Piet W. Brouwer

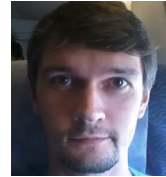
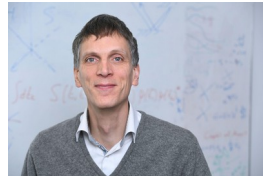
Dmitry Abanin

## supercooled insulators

PRL 2021



I asked him  
the photo  
too late :(

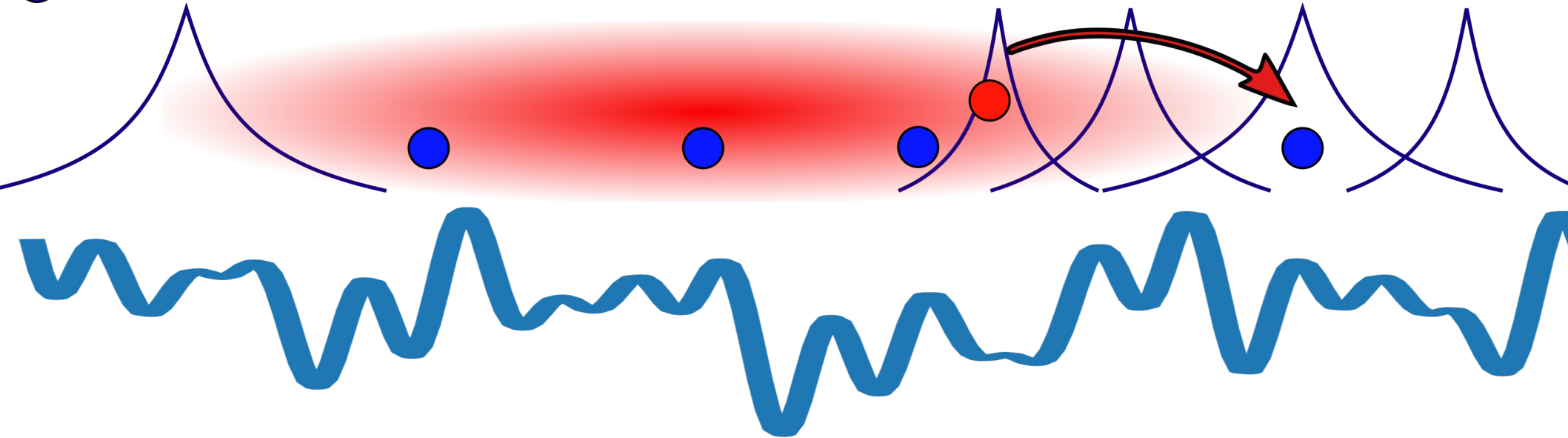


 Spin down

 Spin up

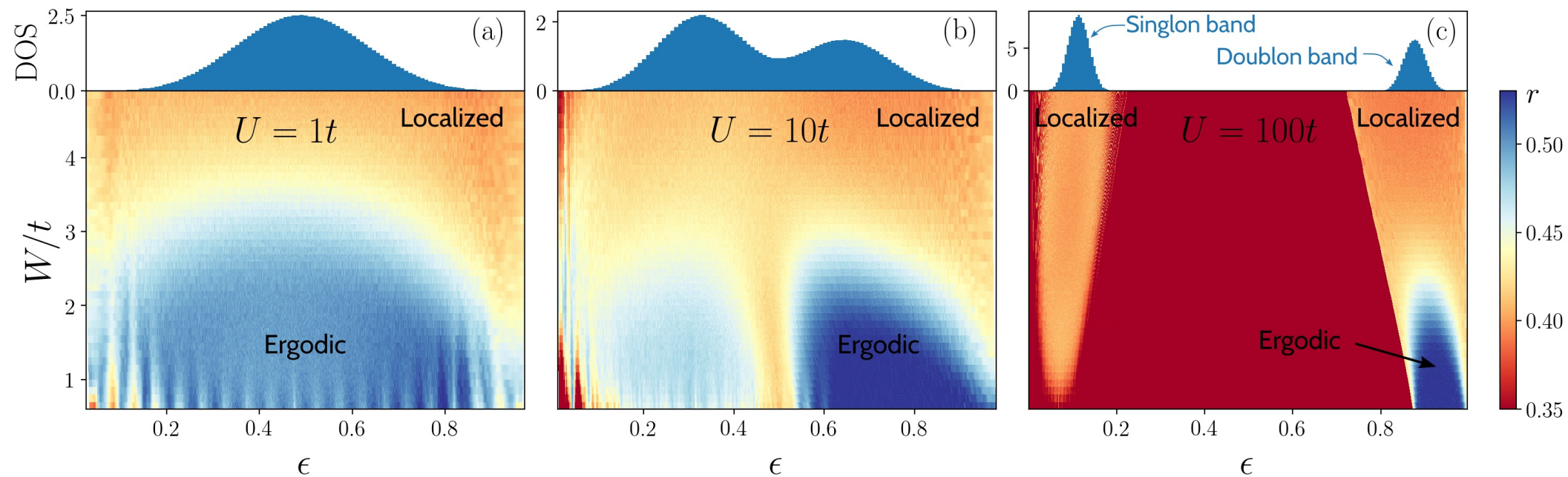
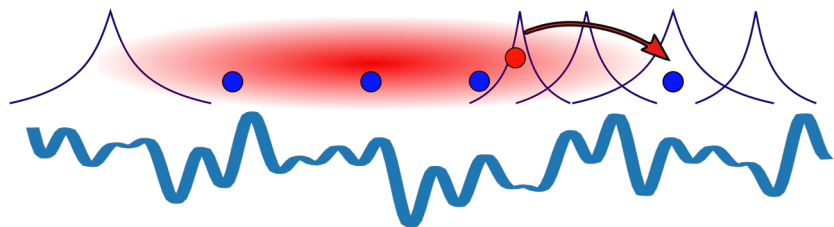
Ergodic bubble

Localized system



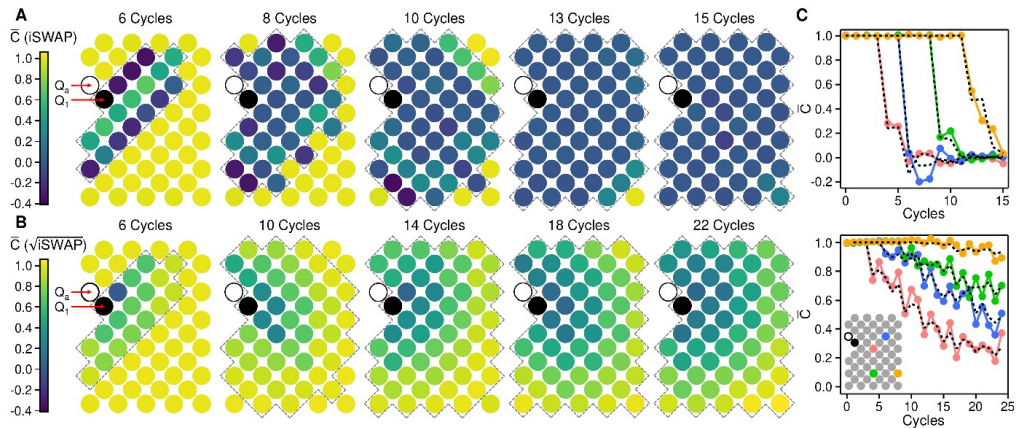
# Nucleation of ergodicity by a single impurity in supercooled insulators

PRL 2021

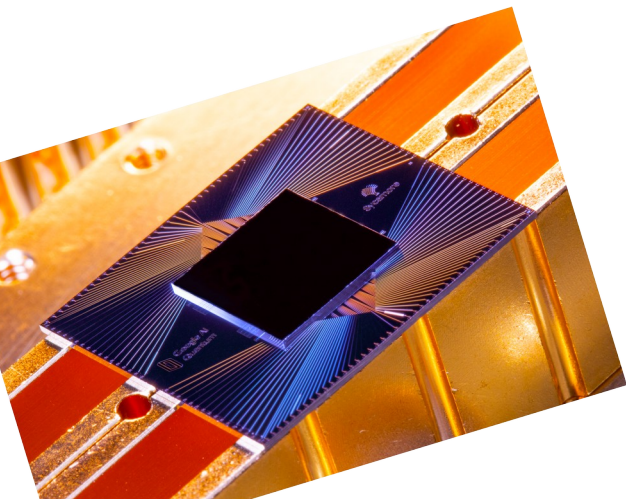
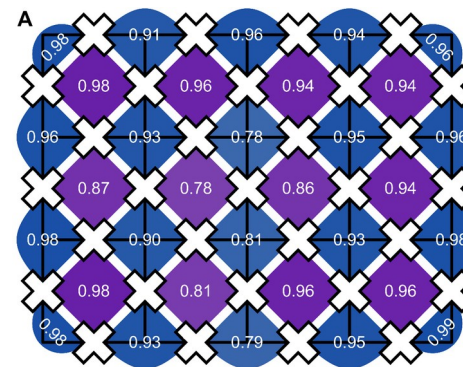


# Quantum dynamics in solid state simulators

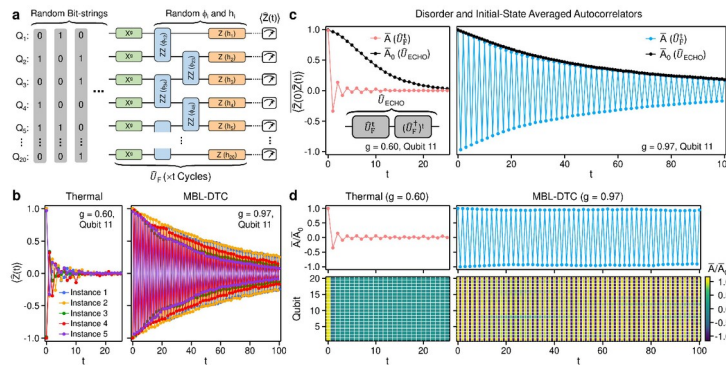
## Information Scrambling in Computationally Complex Quantum Circuits



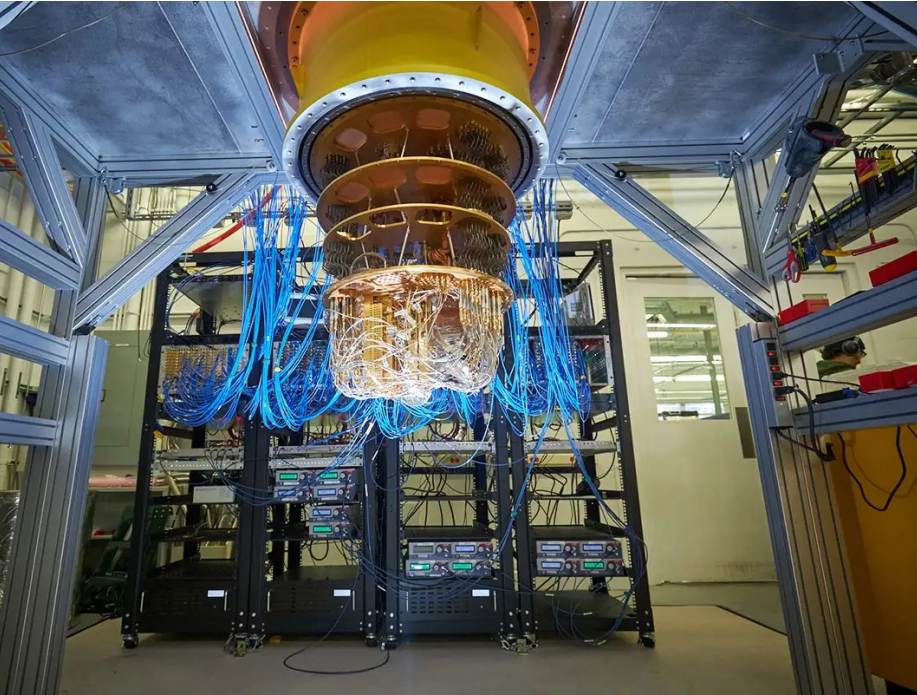
## Realizing topologically ordered states on a quantum processor



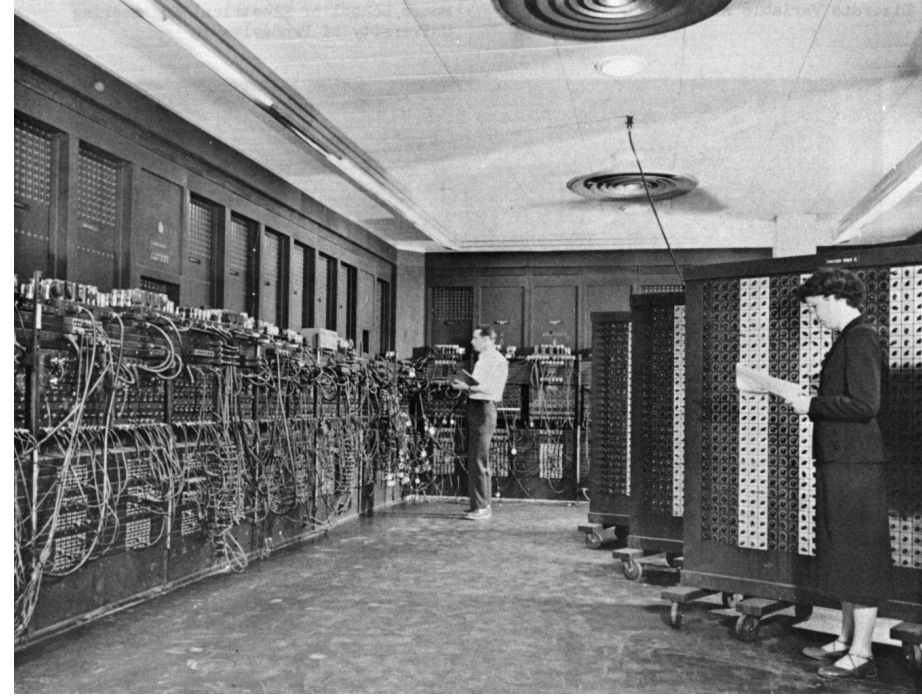
## Observation of Time-Crystalline Eigenstate Order on a Quantum Processor



Google Supercomputer (2019/ ?)



Eniac Supercomputer (1946/1955)



?

=

**THE SIMPSONS**



*"I predict that within 100 years, computers will be twice as powerful, 10,000 times larger, and so expensive that only the five richest kings in Europe will own them."*

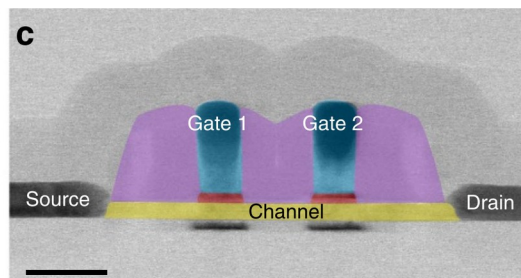
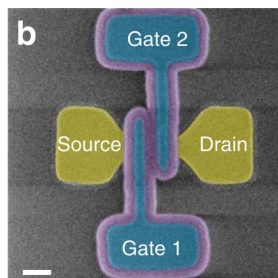
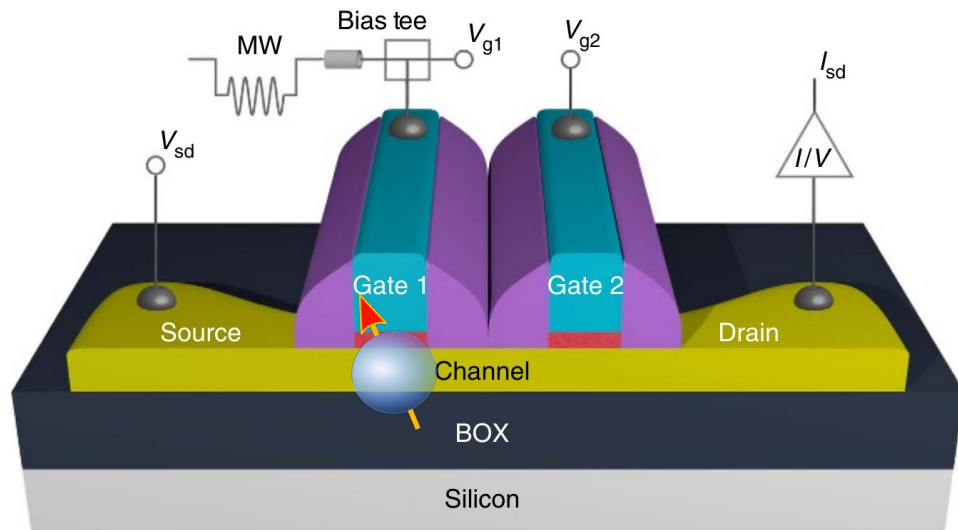
A blue-tinted close-up photograph of a printed circuit board (PCB). The image shows a dense network of white conductive traces on a dark substrate. Several components are visible, including a large integrated circuit (IC) in the center with multiple pins and a smaller component to its right. The text "Here in Grenoble ..." is overlaid in white, bold font in the center of the image.

**Here in Grenoble ...**

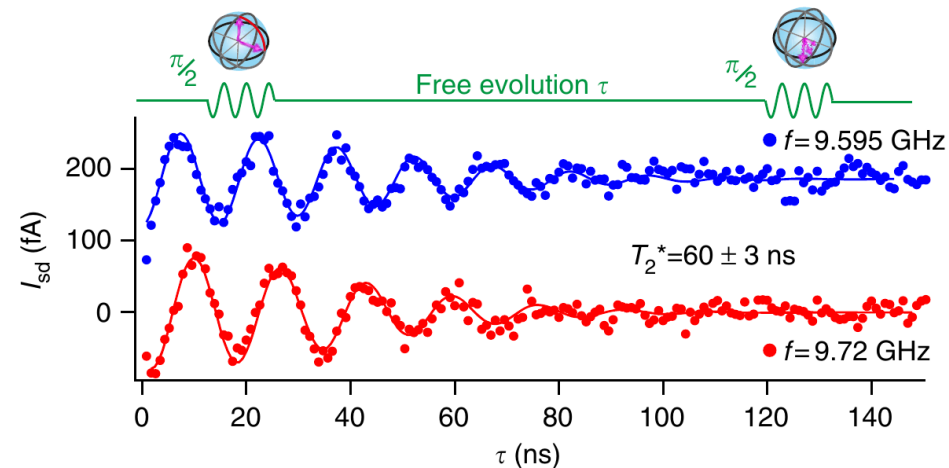
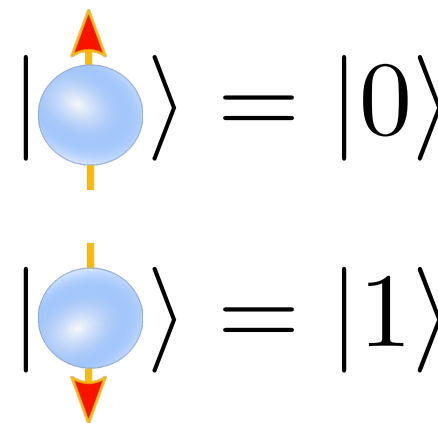


# A CMOS silicon spin qubit

R. Maurand<sup>1,2</sup>, X. Jehl<sup>1,2</sup>, D. Kotekar-Patil<sup>1,2</sup>, A. Corna<sup>1,2</sup>, H. Bohuslavskiy<sup>1,2</sup>, R. Laviéville<sup>1,3</sup>, L. Hutin<sup>1,3</sup>, S. Barraud<sup>1,3</sup>, M. Vinet<sup>1,3</sup>, M. Sanquer<sup>1,2</sup> & S. De Franceschi<sup>1,2</sup>

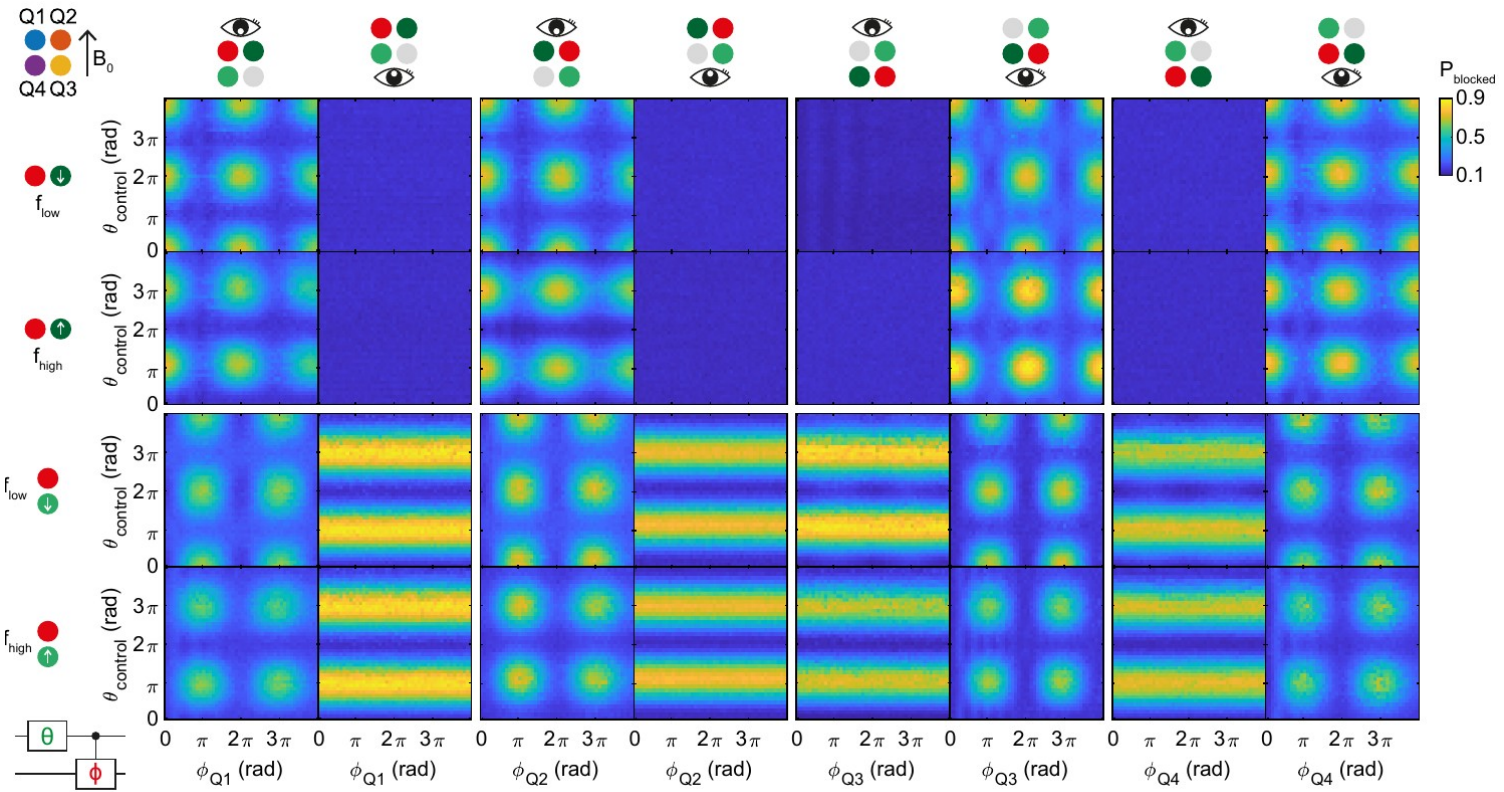
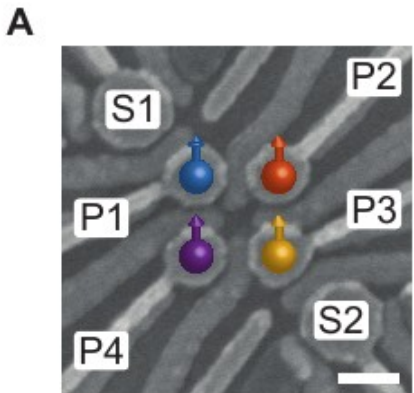


Qubits are now encoded in the spin state of single electrons confined in semiconductors



# A four-qubit germanium quantum processor

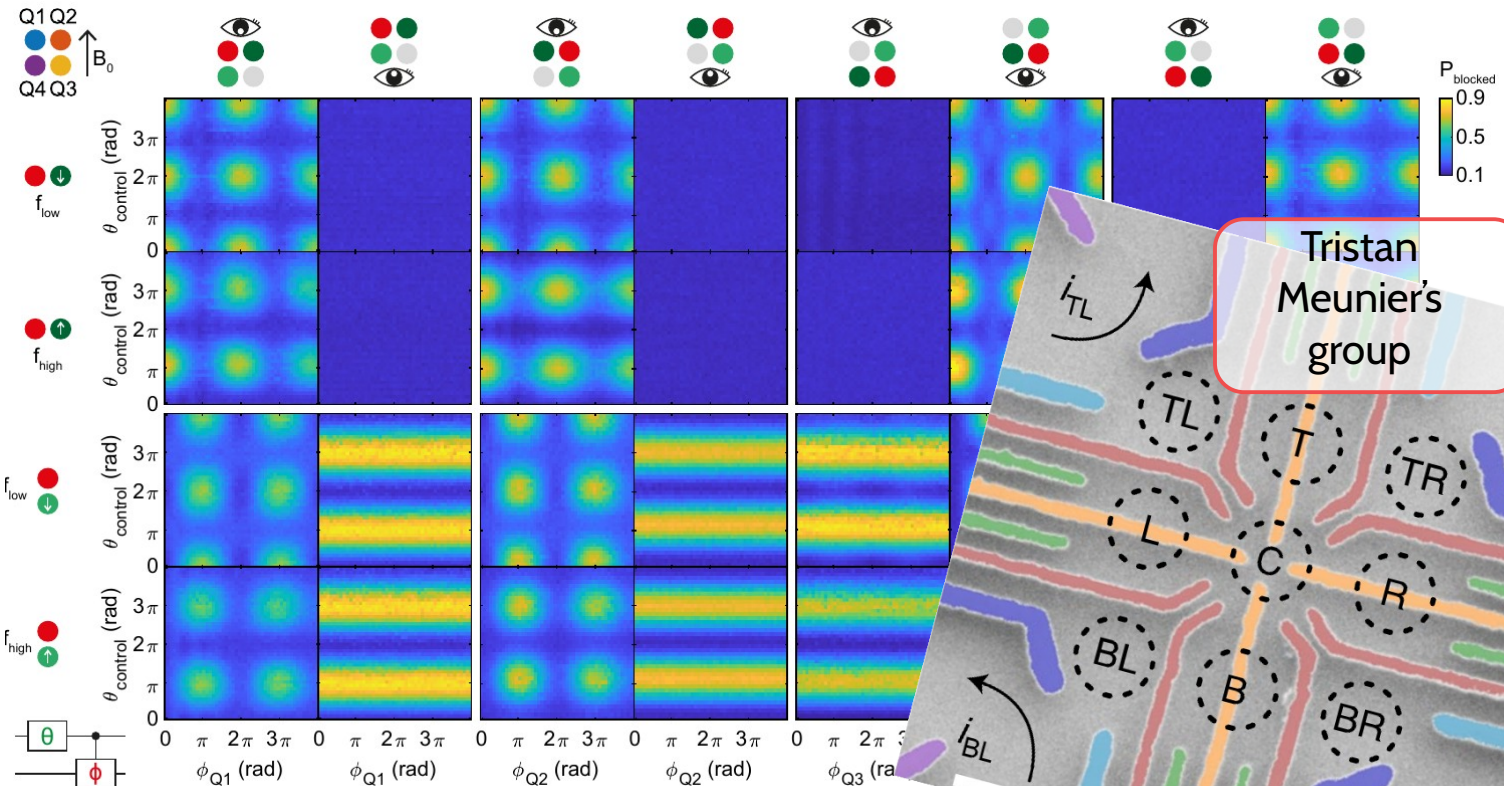
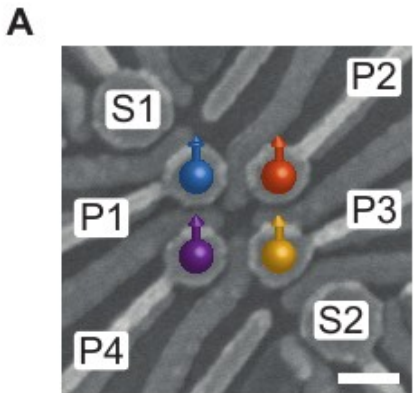
Nico W. Hendrickx<sup>1</sup>, William I. L. Lawrie<sup>1</sup>, Maximilian Russ<sup>1</sup>, Floor van Riggelen<sup>1</sup>, Sander L. de Snoo<sup>1</sup>, Raymond N. Schouten<sup>1</sup>, Amir Sammak<sup>2</sup>, Giordano Scappucci<sup>1</sup> & Menno Veldhorst<sup>1</sup>





# A four-qubit germanium quantum processor

Nico W. Hendrickx<sup>1</sup>, William I. L. Lawrie<sup>1</sup>, Maximilian Russ<sup>1</sup>, Floor van Riggelen<sup>1</sup>, Sander L. de Snoo<sup>1</sup>, Raymond N. Schouten<sup>1</sup>, Amir Sammak<sup>2</sup>, Giordano Scappucci<sup>1</sup> & Menno Veldhorst<sup>1</sup>



# Many thanks to many people

*Thierry Giamarchi*



*Piet W. Brouwer*



*Frank Hekking*



*Anna Minguzzi*



*Denis Basko*



*Cécile Repellin*



*Dima Abanin*

*Pedram Roushan*

*me with  
horrible haircut*





**Thank you for your attention!**