







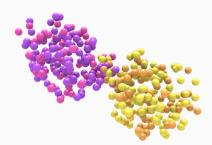


PHQMD: Effects of EOS: Flow, Clustering

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Subatech, CNRS, University of Nantes JRJC 22/10/2021



PHQMD

PHQMD(Parton Hadron Quantum Molecular Dynamics)¹ is a simulation code that is based on PHSD(Parton Hadron String Dynamics)² (quasi particle description of QGP of PHSD)but uses N-body dynamics for propagation of nucleons

Quantum Molecular Dynamics

Using n-body theory to track nucleon by nucleon the collisions, and calculate the two body interactions between all of the nucleons.

- IQMD³(Limited to lower energies 1.5 GeV/A)
- UrQMD⁴(Relativistic, used for coalescence, no potential)

¹J. Aichelin and, E. Bratkovskaya et al. *PR C101, 044905*

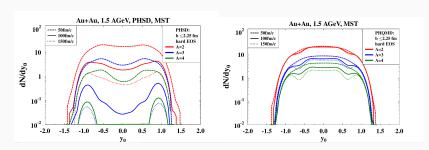
²W. Cassing and E. Bratkovskaya PR C78,034919

³C. Hartnack et al. Eur. Phys. J. A1:151

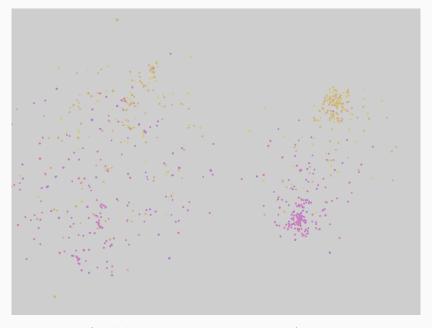
⁴S. A. Bass et al. NP A41:225

Why is it needed?

Current models are not well adapted to address cluster formation, and current QMD models are not usable at relativistic energies



Ice in Fire / Hyper-nuclei production / Cluster production



b = 2.00

b = 7.00

Future experimental data

New energy range from 2 GeV/A - 100 GeV/A, will be used to investigate the first order phase transition from hadronic to QGP matter, and degrees of freedom of hadronic matter (strangeness) .

- FAIR
- NICA

Experimental results so far

- 0.6 [GeV/A] ALADIN 5
- 1.23 [GeV/A] HADES 6
- 1.5 [GeV/A] FOPI 7

⁷A. Schüttauf et al *NP A607,457*

⁸J. Adamczewski-Musch et al NP A982

⁹W. Reisdorf et al. NP A876

Brief PHQMD model overview

Generalised Ritz variational principle

$$\delta \int_{t_1}^{t_2} dt \langle \Psi(t) | i \frac{d}{dt} - H | \Psi(t) \rangle = 0$$
 (1)

Gaussian test wave function which yields the Wigner density:

$$f(\vec{r_i}, \vec{p_i}, \vec{r_{i0}}, \vec{p_{i0}}, t) = \frac{1}{\pi^3 \hbar^3} \exp\left[-\frac{2}{L} (\vec{r_i} - \vec{r_{i0}})^2\right] \exp\left[-2L(\vec{p_i} - \vec{p_{i0}})^2\right]$$
(2)

Classical type equations of motion for expectation value of the Hamiltonian for Gaussian wavefunctions

$$\dot{r}_{i0} = \frac{\partial \langle H \rangle}{\partial p_{i0}} \qquad \dot{p}_{i0} = -\frac{\partial \langle H \rangle}{\partial r_{i0}} \tag{3}$$

Brief PHQMD model overview

Hamiltonian composed of Kinetic term and Two body potentials

$$\langle H \rangle = \sum_{i}^{N} \langle H_{i} \rangle = \sum_{i}^{N} \left(\langle T_{i} \rangle + \sum_{i \neq j}^{N} \langle V_{i,j} \rangle \right)$$
 (4)

Potentials (Coulomb , Asymmetry , Skyrme and Momentum dependent Skyrme)

$$\langle V \rangle = \langle V_c \rangle (\vec{r}) + \langle V_a \rangle (\vec{r}) \begin{cases} \langle V_s \rangle (\rho_{int}) \text{ Static} \\ \langle V_s \rangle (\rho_{int}) + \langle U_{opt} \rangle (\Delta \vec{p}^2) \text{ Momentum dependent} \end{cases}$$
(5)

Interaction density

$$\rho_{int}(i,t) = C \sum_{i,i \neq i}^{N} \exp\left[-\frac{1}{L}(\vec{r_i} - \vec{r_j})^2\right]$$
 (6)

E.o.S parametrisation

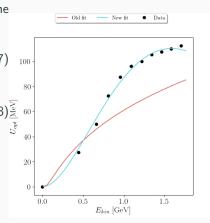
Momentum dependent Skyrme parametrisation

$$\langle V_{tot} \rangle = \langle V_s \rangle \left(\rho_{int} \right) \qquad (7) \qquad 100 - 4$$

$$+ \langle U_{opt} \rangle \left(\Delta p^2 \right) \qquad 80 - 4$$

$$U_{opt}(\Delta p^2) = \exp \left[-c\sqrt{\Delta p} \right] \qquad (8) \stackrel{\text{S}}{\underset{\text{opt}}{\overset{\text{S}}{\geq}}} \qquad 60 - 4$$

$$\left(a\Delta p + b\Delta p^2 \right) \rho(\vec{r}) \qquad 20 - 4$$



 $^{^1\}mbox{S}.$ Hama, et al., Phys. Rev. C 41, 2737

Optical potential fit to expermimental data⁸

E.o.S parametrisation

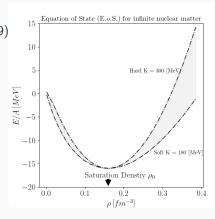
Static Skyrme parametrisation

$$\left\langle V_{s} \right
angle \left(
ho_{int}
ight) \propto lpha
ho + eta
ho^{\gamma}$$

Constraints for the potentials

$$\begin{cases} \rho_0 = 0.1695 [fm^{-3}] \\ E_0 = -16 [MeV] \\ K_0 = 200 \text{ or } 380 [MeV] \\ (\text{Soft/Hard}) \end{cases}$$

$$K_0 = 9 \frac{\partial^2 E/A}{\partial \rho^2} \Big|_{\rho = \rho_0}$$



Equations of state parametrised for soft and hard nuclear matter compressibility

E.o.S parametrisation

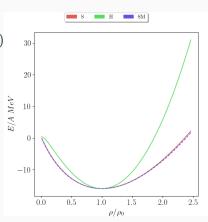
Static Skyrme parametrisation

$$\langle V_s \rangle (\rho_{int}) \propto \alpha \rho + \beta \rho^{\gamma}$$
 (10)

Constraints for the potentials

$$\begin{cases} \rho_0 = 0.1695 [fm^{-3}] \\ E_0 = -16 [MeV] \\ K_0 = 200 \text{ or } 380 [MeV] \\ (\text{Soft/Hard}) \end{cases}$$

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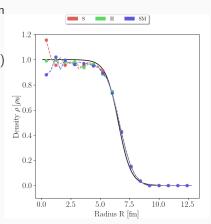
Equations of state parametrised for soft and hard nuclear matter compressibility

Nucleon distribution

Wood-Saxon distribution in position space of nucleons

$$\rho^{WS}(r) = \frac{\rho_0}{1 + \exp\left[\frac{r - R_A}{a}\right]}$$
(11)

$$\begin{cases} R_A = r_0 A^{1/3} \\ r_0 = 1.125[fm] \\ a = 0.535[fm] \end{cases}$$



Initial Wood-Saxon position distribution for the three E.o.S.

Binding energy

Total binding energy of our nucleus $(< B_F > = -10[MeV])$

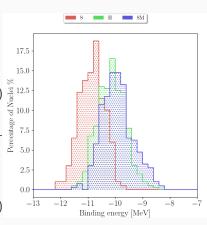
$$B_{E} >= -10[\textit{IMeV}]$$

$$B_{E} = E_{kin}(\vec{p}) + E_{c}(\vec{r}) + E_{a}(\vec{r})$$

$$\begin{cases} E_{s}(\vec{r}) \\ E_{s}(\vec{r}) + E_{m}(\vec{r}, \Delta p) \end{cases}$$
(12)
$$\begin{cases} E_{s}(\vec{r}) + E_{m}(\vec{r}, \Delta p) \\ 0 \end{cases}$$
instraint on momentum distribungan (10.0)

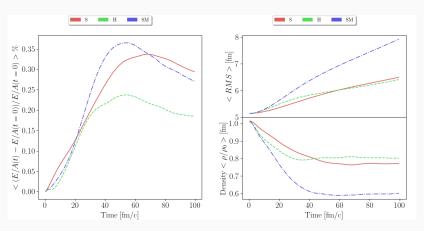
Constraint on momentum distribution

$$0 \le \sqrt{m^2 + \vec{p}_i^2} - m \le -V(\rho_{int})$$
 (13)



Initial binding energy of the nuclei

Propagation



Conservation of the total system energyEvolution of the average radial distance of the nucleons, and the average density of the nuclei

Cluster formation

SACA⁹ and MST¹⁰

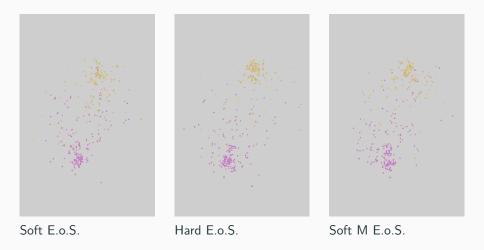
To form the cluster we use two methods

- MST: simple form of spanning tree to create clusters based on distance from other nucleons
- SACA: complex simulated annealing of all possible cluster patterns of the nucleons to find the lowest sum of binding energies of all clusters

All the following results are for the SACA method

¹⁰R. K. Puri, C. Hartnack and J. Aichelin PR C54, R28

¹¹ J. Aichelin, Phys. Rept 202, 233



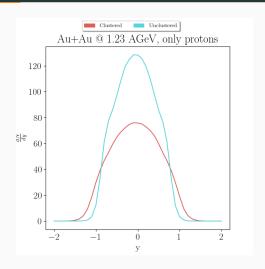
Cluster formation

SACA¹¹ and MST¹²

To form the cluster we use two methods

- MST :
- SACA:

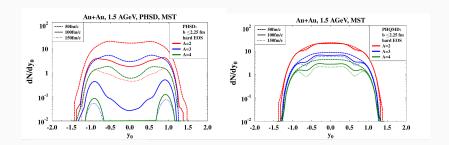
Clustering is important at these energies up to 40 % of protons are in clusters



¹⁰R. K. Puri, C. Hartnack and J. Aichelin *PR C54, R28*

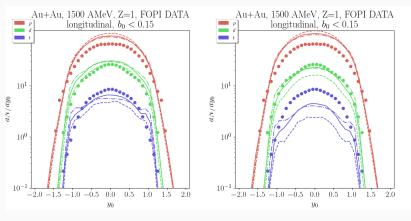
¹¹ J. Aichelin, Phys. Rept 202, 233

Extraction tine



Multiple extraction time possible Cluster stability and rapidity dependence

FOPI results for AuAu E = 1.5 [GeV/A] Multiplicity

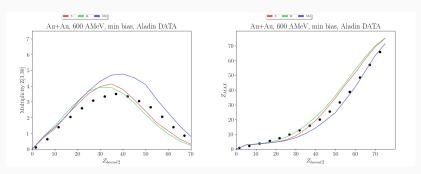


MST identified clusters

SACA identified clusters

Aladin results for AuAu E = 0.6 [GeV/A]

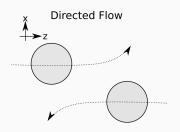
 Z_{bound2} the sum of the charges all fragments with $Z \geq 2$, and Z_{MAX} the charge of the largest cluster

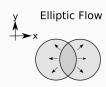


Rise and Fall curve for all three equations of state

Zmax curve for all three equations of state

Directed and Elliptic Flow





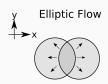
$$E\frac{\mathrm{d}^{3}N}{\mathrm{d}^{3}p} = \frac{1}{2\pi} \frac{\mathrm{d}^{2}N}{p_{T}\mathrm{d}p_{T}\mathrm{d}y}$$

$$\left(1 + 2\sum_{n=1}^{\infty} v_{n} \cos\left[n(\varphi - \Psi_{RP})\right]\right) \quad (14)$$

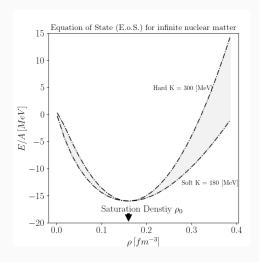
A Fourier expansions allows us to obtain the coefficients:

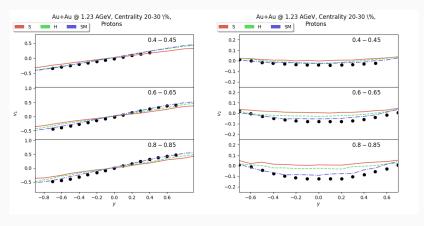
- Directed Flow $v_1 = \left\langle \frac{p_x}{p_T} \right\rangle$
- Elliptic Flow $v_2 = \left\langle \left(\frac{p_x}{p_T}\right)^2 \left(\frac{p_y}{p_T}\right)^2 \right\rangle$

Directed and Elliptic Flow

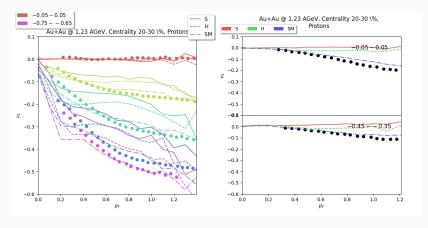


Thus More sensitive to the variation in density thus E.o.S. choice

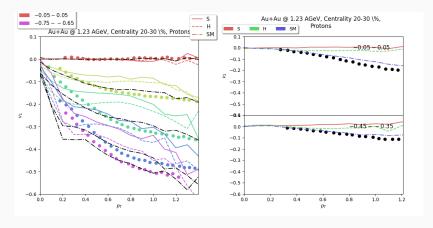




Directed proton flow (v_1) as a function Elliptical proton flow (v_2) as a function of rapidity for three p_T classes compared with the HADES data compared with the HADES data

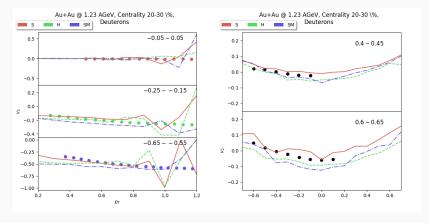


Directed proton flow (v_1) as a function Elliptical deuteron flow (v_2) as a of p_T for various rapidity classes function of rapidity for two p_T classes compared with the HADES data



Directed proton flow (v_1) as a function Elliptical deuteron flow (v_2) as a of p_T for various rapidity classes compared with the HADES data

function of rapidity for two p_T classes compared with the HADES data

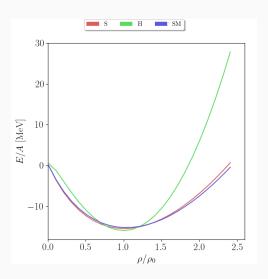


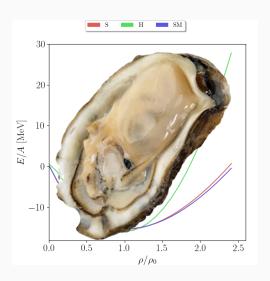
Directed deuteron flow (v_1) as a Elliptical deuteron flow (v_2) as a function of p_T for three various rapidity function of rapidity for two p_T classes classes compared with the HADES datacompared with the HADES data

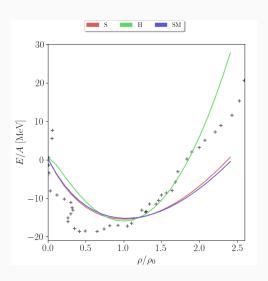
Conclusions

- PHQMD allows use to study cluster formation in a unique and novel way
- It permits us to study the effect of nuclear EOS on a host of observables
- In most cases for the FOPI data the SM EOS give better agreement with the data In some cases the HADES data is better replicated for the SM EOS

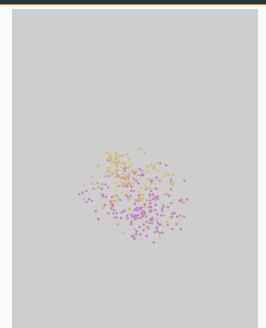
Thank you

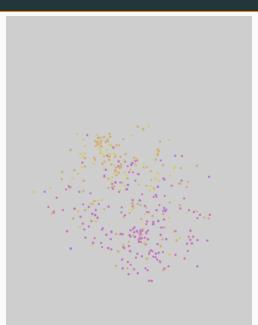


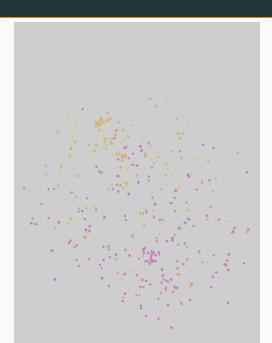






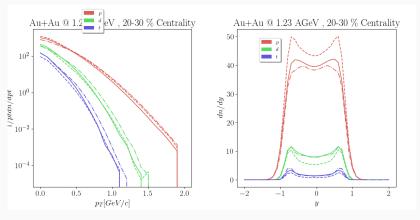






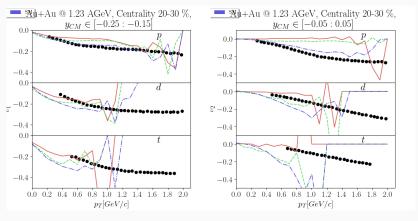


AuAu E = 1.23 [GeV/A] Multiplicity



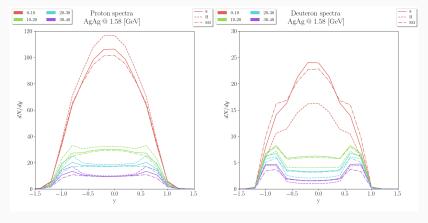
Multiplicity of Z=1 isotopes function of Multiplicity of Z=1 isotopes function of rapidity for Au+AU @ 1.23 AGeV p_T for Au+AU @ 1.23 AGeV

AuAu E = 1.23 [GeV/A] Flow



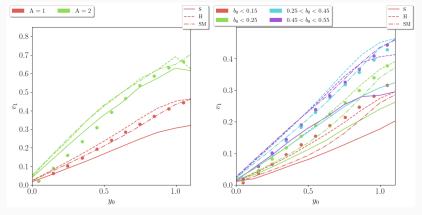
Directed isotope flow of Z=1 isotopes function of p_T for Au+AU @ 1.23 AGeV

Elliptical Z=1 isotopes flow function of p_T for Au+AU @ 1.23 AGeV



Proton multiplicity Ag Ag @ 1.58 AGeVDeuteron multiplicity Ag Ag @ 1.58 function rapidity AGeV function rapidity

FOPI results for AuAu E = 1.5 [GeV/A] Flow



Directed flow for proton (v_1) and Directed proton flow (v_1) for 4 different deuteron with a threshold on transverse impact parameter classes for all protons momentum with a threshold on transverse momentum