



CLUSTER OF EXCELLENCE
QUANTUM UNIVERSE

The order parameters of CP violation in the SMEFT

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Work in progress
with E. Gendy, C. Grojean and J. Ruderman

Motivation

What is the (basis-independent) parameter space of CPV observables in the SM(EFT) ?

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In this talk : **when does the SMEFT conserve CP ?**

CP breaking in the SM

When does the SM break CP ?

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$$\mathcal{L} \supset \frac{i}{\sqrt{2}} \overline{u}_L \gamma^\mu W_\mu^+ V_{CKM} d_L - \overline{u}_L \text{diag}(m_{u_i}) u_R + (d) + h.c.$$

CP breaking in the SM

complex matrix !

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$$\begin{pmatrix} \frac{3}{13} & \frac{4}{13} & \frac{12}{13} \\ -\frac{692}{845} & -\frac{381}{845} & \frac{60}{169} \\ \frac{444}{845} & -\frac{708}{845} & \frac{25}{169} \end{pmatrix} \text{CP} \checkmark$$

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$$\mathcal{J} = \text{Im}(V_{CKM,us} V_{CKM,cb} V_{CKM,ub}^* V_{CKM,cs}^*) = 0$$

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If $m_u = m_c$ (or $m_u = m_t$ or ...) **CP** ✓

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$$(y_t^2 - y_c^2)(y_t^2 - y_u^2)(y_c^2 - y_u^2)(y_b^2 - y_s^2)(y_b^2 - y_d^2)(y_s^2 - y_d^2) \mathcal{J}$$

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$$J_4 = \text{Im Tr} \left[Y_u Y_u^\dagger, Y_d Y_d^\dagger \right]^3 = 6(y_t^2 - y_c^2)(y_t^2 - y_u^2)(y_c^2 - y_u^2)(y_b^2 - y_s^2)(y_b^2 - y_d^2)(y_s^2 - y_d^2) \mathcal{J}$$

$$\mathcal{L} \supset -\bar{Q}_L Y_u u_R \tilde{H} + (d) + h.c.$$

$$Y_u = \text{diag}(y_u, y_c, y_t) , \quad Y_d = V_{CKM} \text{diag}(y_d, y_s, y_b)$$

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Flavour invariant

	$SU(3)_Q$	$SU(3)_u$	$SU(3)_d$	$SU(3)_L$	$SU(3)_e$
Q_L	3	1	1	1	1
Y_u	3	$\bar{3}$	1	1	1
Y_d	3	1	$\bar{3}$	1	1
Y_e	1	1	1	3	$\bar{3}$

$$Y_u = \begin{pmatrix} Y_{u,11} & Y_{u,12} & Y_{u,13} \\ Y_{u,21} & Y_{u,22} & Y_{u,23} \\ Y_{u,31} & Y_{u,32} & Y_{u,33} \end{pmatrix}$$

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**CP is conserved in the SM
iff
 $J_4=0$**

[Jarlskog '85]

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$$J_4 = \text{Im Tr} \left[Y_u Y_u^\dagger, Y_d Y_d^\dagger \right]^3$$

\implies No CP breaking in the SM
with $N_f=2$ **[Cabibbo '63,
Kobayashi/Maskawa '73]**

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⇒ No CP breaking in the SM
with $N_f=2$ [Cabibbo '63,
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⇒ CP breaking is a **collective effect**
Explains suppressions

$$d_e^{\text{Fig.1a}} \sim e \mathcal{J} \frac{m_e m_c^2 m_s^2}{m_W^6} \frac{\alpha_W^3 \alpha_s}{(4\pi)^4} \quad [\text{Pospelov/Ritz '13}]$$

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CP breaking BSM ?

CP breaking in the SMEFT

No new particle at the LHC, no clear deviation from SM couplings : renewed interest in **effective field theories** (EFTs)

$$\mathcal{L} = \mathcal{L}_{d \leq 4} + \sum_i \frac{c_i}{\Lambda^{d_i-4}} \mathcal{O}_i$$

If no light BSM d.o.f. : generic framework to interpret experimental results

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The SMEFT

$$A_{\mu}^{a,i,Y} \quad \psi_{i,L/R} \quad H$$

$$SU(3) \times SU(2) \times U(1)$$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^{d_i-4}} \mathcal{O}_i$$

$$\mathcal{O}_i = C_{H\psi_L,ij}^{(1)} \psi_L^i \gamma^{\mu} \psi_L^j \left(H^{\dagger} \overleftrightarrow{D}_{\mu} H \right)$$

[Buchmüller/Wyler '85, Grzadkowski et al '10]

CP breaking in the SMEFT

When does the SMEFT break CP
at dimension-six ?

CP breaking in the SMEFT

When does the SMEFT break CP
at dimension-six ?

Type of op.	# of ops	# of real coeffs	# of imaginary coeffs
Yukawa	3	27	27
Dipoles	8	72	72
current-current	8	51	30
all bilinears	19	150	129
LLLL	5	171	54
RRRR	7	255	195
LLRR	8	360	288
LRRL	1	81	81
LRLR	4	324	324
all 4-Fermi	25	1191	1014
all		1341	1143

[Grzadkowski/Iskrzynski/Misiak/Rosiek '10,
Alonso/Manohar/Jenkins '13]

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CP-odd

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Collective aspects, basis-independence: use **flavour-invariants**

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When does the SMEFT break CP at dimension-six ?

**CP is conserved in the
SMEFT
iff
? = 0**

CP breaking in the SMEFT

When does the SMEFT break CP at dimension-six ?

CP is conserved in the SMEFT iff $\mathcal{C} = 0$

Observation 1: CP conservation follows the EFT power counting

CP breaking in the SMEFT

When does the SMEFT break CP at dimension-six?

CP is conserved in the SMEFT

iff

$$\mathbf{J_4 = 0 \ \& \ ? = 0}$$

Observation 1: CP conservation follows the EFT power counting

\implies CP should be conserved order by order

CP breaking in the SMEFT

When does the SMEFT break CP at dimension-six ?

**CP is conserved in the SMEFT
iff
 $J_4 = 0$ & $\theta = 0$**

Observation 1: CP conservation follows the EFT power counting

\implies CP should be conserved order by order

\implies Some imaginary parts are unphysical at first order

$$|\mathcal{A}|^2 = |\mathcal{A}_{\text{SM}}|^2 + \frac{2}{\Lambda^2} \text{Re} (\mathcal{A}_{\text{SM}} \mathcal{A}_{\text{BSM}}^*)$$

CP breaking in the SMEFT

When does the SMEFT break CP at dimension-six ?

CP is conserved in the SMEFT iff $J_4 = 0$ & $? = 0$

Observation 1: CP conservation follows the EFT power counting

⇒ CP should be conserved order by order

⇒ Some imaginary parts are unphysical at first order

Type of op.	# of ops	# real	# im.	# CP-odd inv.
Yukawa	3	27	27	21
Dipoles	8	72	72	60
current-current	8	51	30	21
all bilinears	19	150	129	102
LLLL	5	171	126	54
RRRR	7	255	195	126
LLRR	8	360	288	174
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all 4-Fermi	25	1191	1014	597
all		1341	1143	699

CP breaking in the SMEFT

When does the SMEFT break CP at dimension-six ?

CP is conserved in the SMEFT iff $J_4 = 0$ & $? = 0$

Observation 1: CP conservation follows the EFT power counting

⇒ CP should be conserved order by order

⇒ Some imaginary parts are unphysical at first order

⇒ the more CP is conserved, the less conditions should be enforced

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Ex : $C_{Hu,ij} \left(iH^\dagger \overleftrightarrow{D}_\mu H \right) \bar{u}_i \gamma^\mu u_j$

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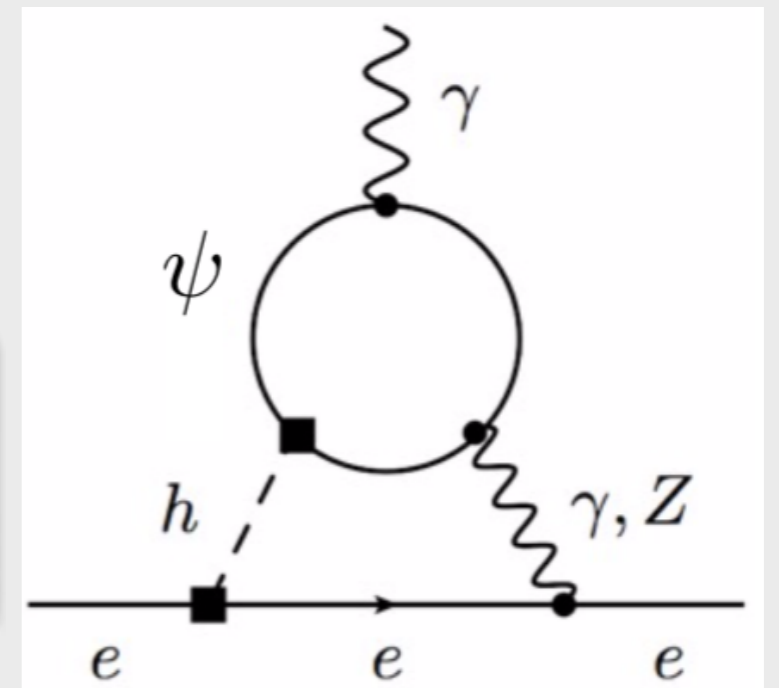
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Example : Barr-Zee two-loop contribution to the electron EDM

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{C_{uH}}{\Lambda^2} |H|^2 \overline{Q}_L u_R \tilde{H} + h.c.$$

$$\frac{d_e}{e} = -\frac{1}{48\pi^2} \frac{v m_e m_u}{m_h^2} \frac{\text{Im}(C_{uH})}{\Lambda^2} F_1 \left(\frac{m_u^2}{m_h^2}, 0 \right)$$

[Barr/Zee '90, Brod/Haisch/Zupan '13]



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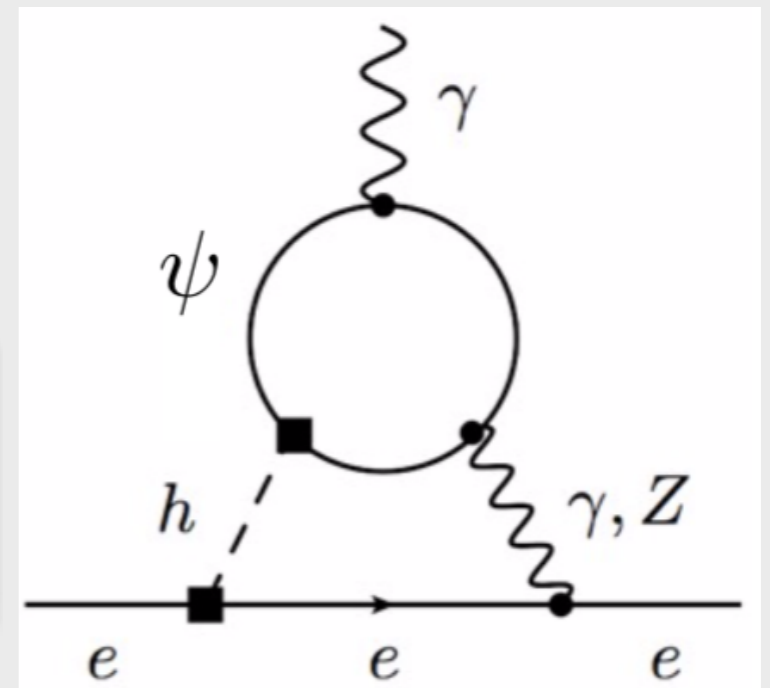
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one-to-one correspondance with invariants

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$$Y_u = \text{diag}(a_u \lambda^8, y_c \lambda^4, y_t)$$

$$Y_d = V_{\text{CKM}} \text{diag}(y_d \lambda^7, y_s \lambda^5, y_b \lambda^3)$$

$$V_{\text{CKM}} = \begin{pmatrix} 1 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

$$L_1 = \text{ImTr} \left(Y_u Y_u^\dagger Y_d Y_d^\dagger Y_u C_{Hu} Y_u^\dagger \right) \sim \lambda^{12}$$

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Ex : $C_{Hu,ij} \left(iH^\dagger \overleftrightarrow{D}_\mu H \right) \bar{u}_i \gamma^\mu u_j$

Froggatt-Nielsen models

$$Y_u \sim \begin{pmatrix} \lambda^8 & \lambda^5 & \lambda^3 \\ \lambda^7 & \lambda^4 & \lambda^2 \\ \lambda^5 & \lambda^2 & 1 \end{pmatrix} \quad C_{HQ} \sim \begin{pmatrix} 1 & \lambda & \lambda^3 \\ \lambda & 1 & \lambda^2 \\ \lambda^3 & \lambda^2 & 1 \end{pmatrix}$$

$$L_1 = \text{ImTr} \left(Y_u Y_u^\dagger Y_d Y_d^\dagger Y_u C_{Hu} Y_u^\dagger \right) \sim \lambda^{14}$$

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Use: - parametrize CPV observables in a basis-independent way
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Ex : 2HDM in the decoupling limit, most invariants prop. to $J4$

Outlook

We found **necessary and sufficient conditions for CP conservation at dimension-six in the SMEFT.** They define the new sources of CPV in a basis-independent way, and connect easily to UV models insensitive to preferred IR bases.

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Other topics :

- CPV with $J_4 \neq 0$ and parameter space of CPV observables
- preferred invariants given a UV model
- ...


THANK YOU

CP-odd flavour invariants for 4-Fermi ops

For 4-Fermi operators : « A-type » $\text{Im} \left(M_{ij}^{uH} M_{kl}^{dH} C_{ijkl}^{QuQd} \right)$

« B-type » $\text{Im} \left(M_{kj}^{uH} M_{il}^{dH} C_{ijkl}^{QuQd} \right)$

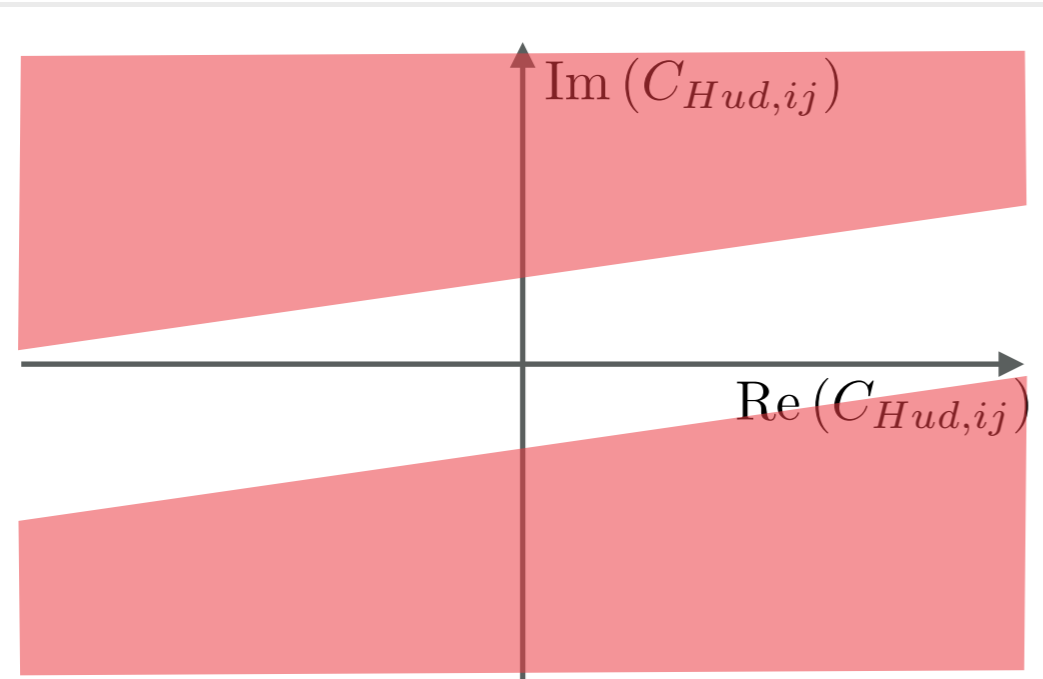
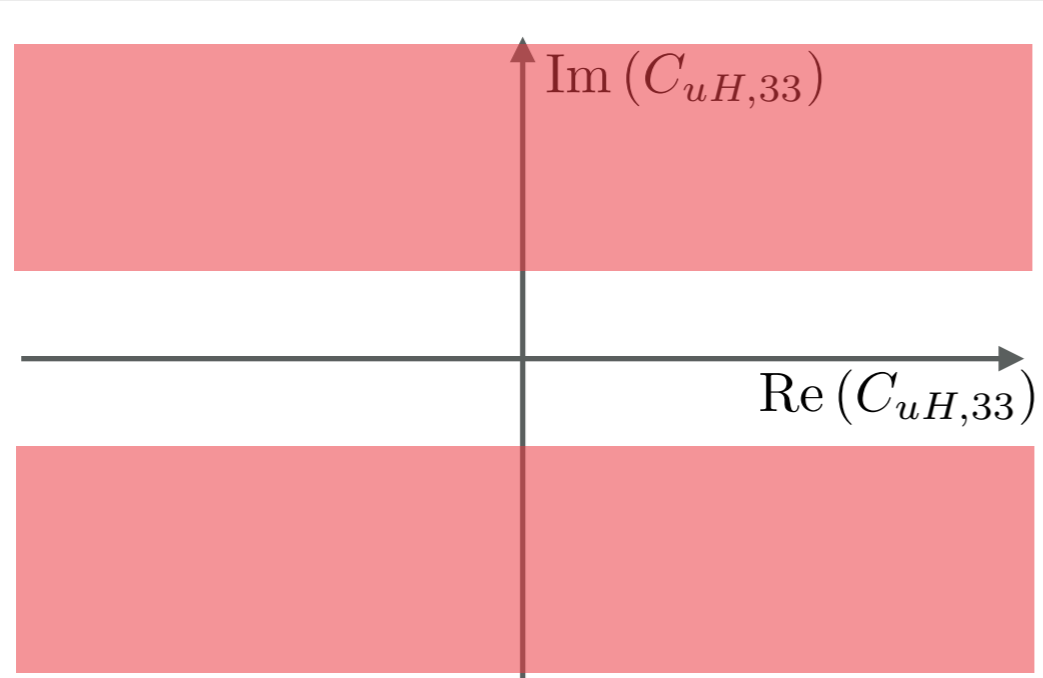
$Y_u^\dagger, Y_u^\dagger Y_d Y_d^\dagger, \dots$



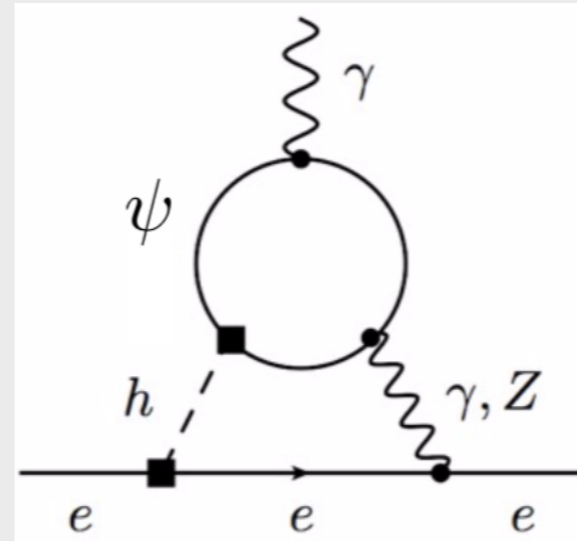
How many « CP-odd directions » are there ?

How many parameters enter CP-odd observables ?

With non-zero J_4 , real and imaginary parts violate CP

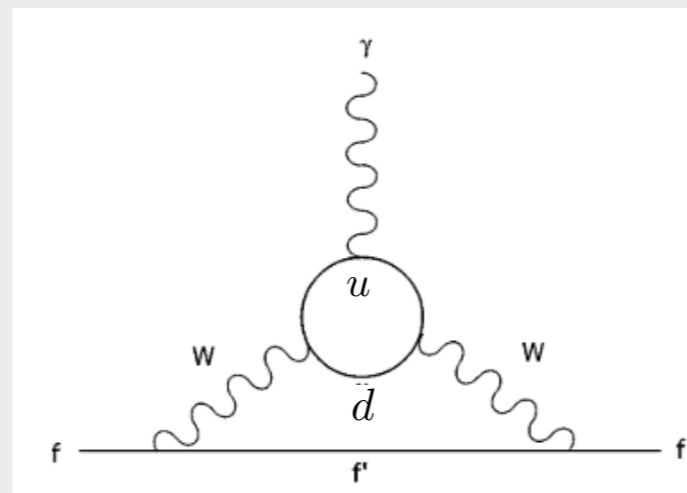


$$\mathcal{L} = \mathcal{L}_{SM} + \frac{C_{uH}}{\Lambda^2} |H|^2 \bar{Q}_L u_R \tilde{H} + h.c.$$



[Brod/Haisch
/Zupan '13]

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{C_{Hud,mn}}{\Lambda^2} i \tilde{H}^\dagger D_\mu H \bar{u}_{R,m} \gamma^\mu d_{R,n}$$



[Kadoyoshi/
Oshimo '97]

Theta QCD

$$\mathcal{L}_{\text{QCD}} \supset -\theta_{\text{QCD}} \frac{g_s^2}{16\pi^2} \text{Tr}(G\tilde{G})$$

	$SU(3)_{Q_L}$	$U(1)_{Q_L}$	$SU(3)_{u_R}$	$U(1)_{u_R}$	$SU(3)_{d_R}$	$U(1)_{d_R}$
Q_L	3	1	1	0	1	0
u_R	1	0	3	1	1	0
d_R	1	0	1	0	3	1
Y_u	3	1	$\bar{3}$	-1	1	0
Y_d	3	1	1	0	$\bar{3}$	-1
$e^{i\theta_{\text{QCD}}}$	1	6	1	-3	1	-3

$$\bar{\theta} \equiv \theta_{\text{QCD}} - \arg \det (Y_u Y_d)$$

$$\text{Im} \left(e^{-i\theta_{\text{QCD}}} \epsilon^{ABC} \epsilon^{abc} \epsilon^{DEF} \epsilon^{def} Y_{u,Aa} Y_{u,Bb} C_{QuQd,CcDd} Y_{d,Ee} Y_{d,Ff} \right)$$

In the UV, suppressed by $e^{-\frac{8\pi^2}{g_s^2}} \approx \lambda^{37-38}$. Relevant in the IR?