

Probing new physics solutions to the B-anomalies at high-energy colliders

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3rd Topical Discussion Session:
Interplay of quark and lepton flavour at Belle II and the LHC,
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Institut de Physique de l'Univers

Overview

■ Introduction

- B-anomalies

■ Flavor at high-energy pp-collider

- Di-taus • Di-muons • LFV tails

■ B-anomalies and Leptoquark solutions

- High-pT and low energy complementarity

■ UV model for the vector leptoquark

Flavor sector of the SM

The SM Lagrangian

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi + h.c. \\ & + \bar{\chi}_i Y_{ij} \chi_j \phi + h.c. \\ & + |\nabla_\mu \phi|^2 - V(\phi)\end{aligned}$$

Gauge sector

Flavor universal
Well tested experimentally

$$\mathcal{G}_{\text{flavor}} \equiv U(3)^5$$

$$U(3)_Q \otimes U(3)_L \otimes U(3)_u \otimes U(3)_d \otimes U(3)_e$$

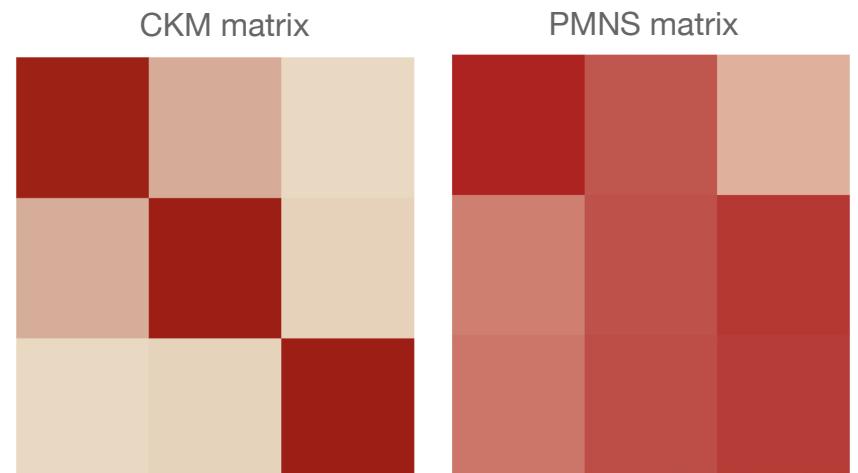
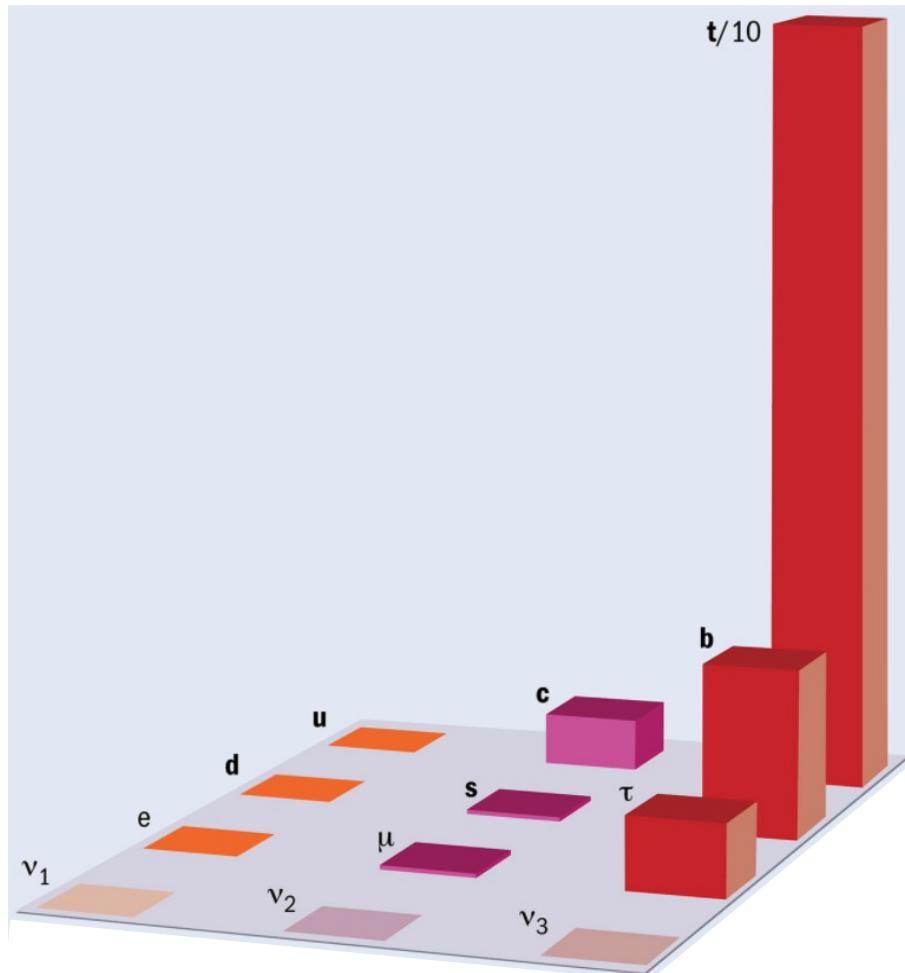
Yukawa sector:

$$\mathcal{L}_{\text{Yukawa}} = -Y_d^{ij} \bar{Q}^i \Phi d^j - Y_u^{ij} \bar{Q}^i \tilde{\Phi} u^j - Y_e^{ij} \bar{L}^i \Phi e^j \implies U(3)^5$$

- Contains most SM free parameters (fermion masses and mixings) extracted from data
- Lepton flavor universality (LFU) is broken in the SM by the fermion masses
- BSM could introduce new sources of LFU breaking
- Good test for BSM physics in flavor experiments!

The SM Flavor puzzle

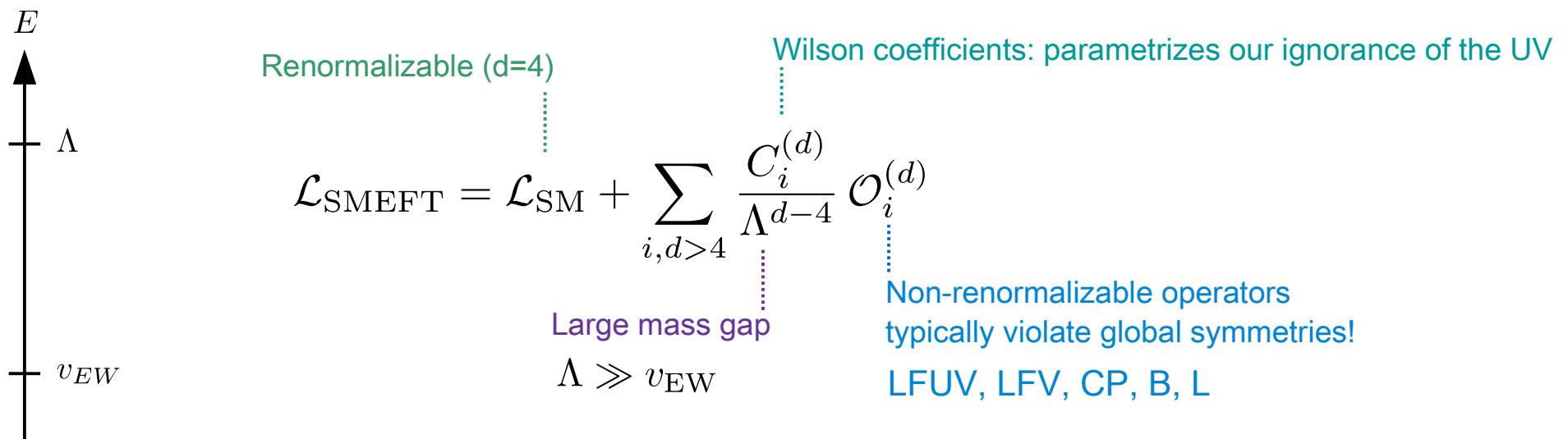
Non-generic pattern for the SM fermion masses and mixings



A clear hierarchy suggesting an underlying theory of Flavor!

Th SMEFT

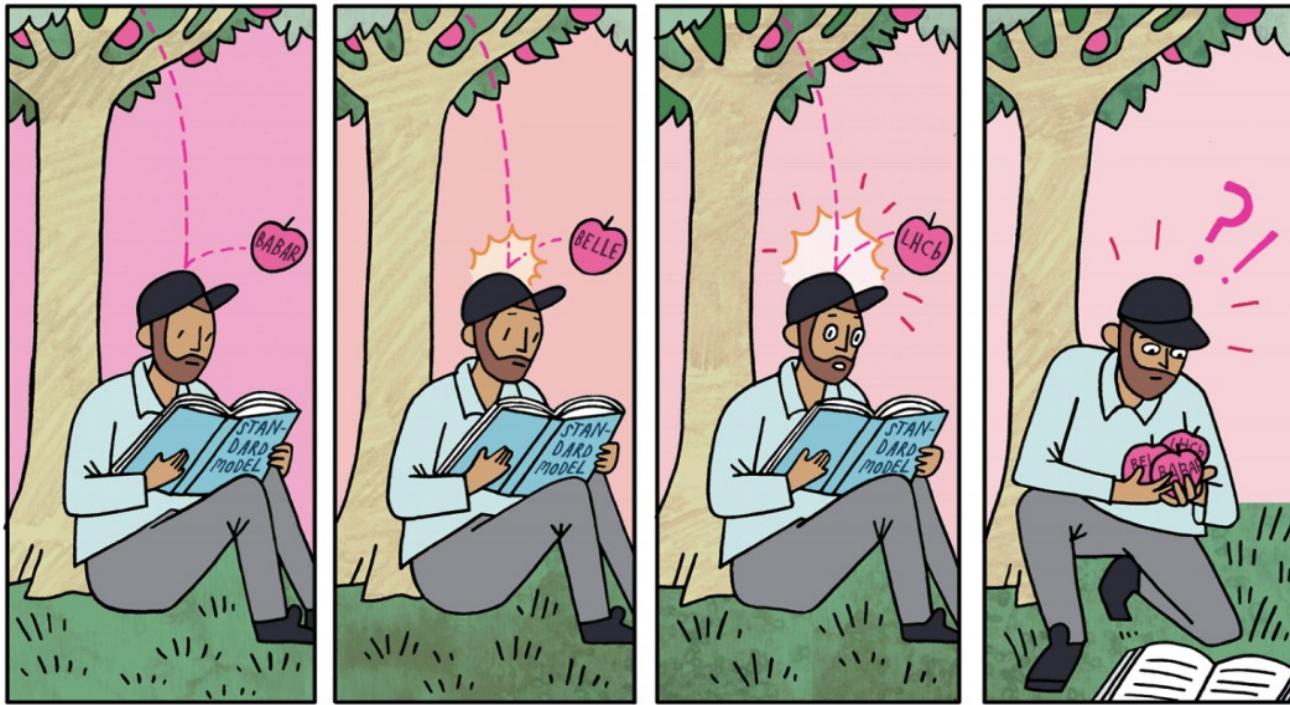
- SM as a low energy effective theory:



$$d = 6 \implies N_{\mathcal{O}} = 59, \quad N_C = 2599$$

- Testing accidental SM symmetries is of fundamental importance!

A Decade of B-meson Anomalies



Artwork by Sandbox Studio, Chicago with Corinne Mucha

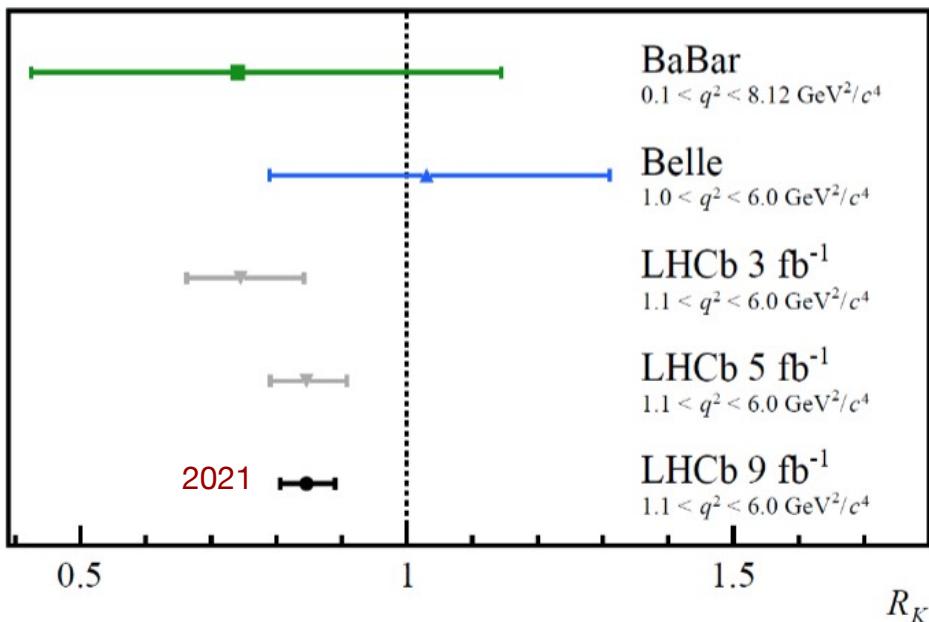
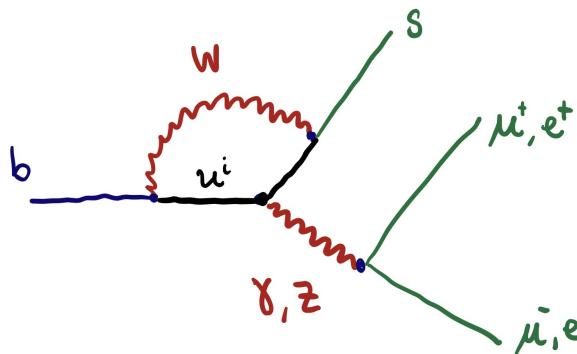
- Evidence of **LFU violation** in semi-leptonic B-decays: different experiments / different observables
- All B-anomalies seem compatible with each other!
- No single measurement is statistically relevant to claim discovery... yet...
- Anomalies come and go... Need more data
- Completely unexpected from theory side...

B-meson Anomalies – Neutral currents

$$b \rightarrow s\ell\bar{\ell}$$

- LFU ratio tau vs light leptons

$$R_{K^{(*)}} = \frac{\text{Br}(B \rightarrow K^{(*)} \mu\bar{\mu})}{\text{Br}(B \rightarrow K^{(*)} e\bar{e})}$$



Evidence for LFUV in R_K [3.1σ]
 [LHCb, [2103.11769](https://arxiv.org/abs/2103.11769)]



$$R_{K^{(*)}}^{\text{exp}} < R_{K^{(*)}}^{\text{SM}} \quad \text{Deficit!}$$

- Other anomalies (muon specific):

$$\begin{array}{ll} B_s \rightarrow \mu\mu & B \rightarrow K\mu\mu \\ B \rightarrow \phi\mu\mu & P'_5 \end{array}$$

All consistent with each other!

Low energy effective theory - RK(*)

Effective Lagrangian:

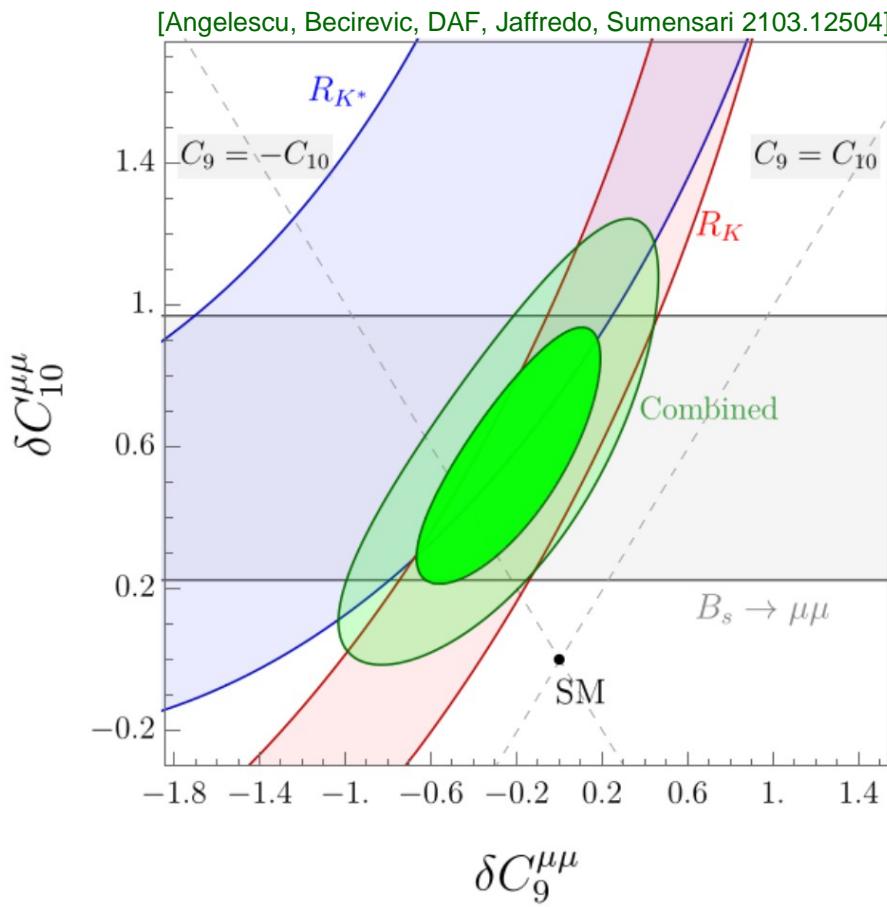
$$\mathcal{L}_{b \rightarrow s \ell \bar{\ell}} = \frac{4 G_F}{\sqrt{2}} V_{ts}^* V_{tb} \sum_i C_i \mathcal{O}_i^\ell + \text{h.c.}$$

- (pseudo) scalar operators very constrained
- Clean observables: R_K , R_{K^*} , $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-)$

$$\mathcal{O}_9^\ell = \frac{\alpha}{4\pi} (\bar{s}_L \gamma^\mu b_L) (\bar{\ell} \gamma^\mu \ell)$$

$$\mathcal{O}_{10}^\ell = \frac{\alpha}{4\pi} (\bar{s}_L \gamma^\mu b_L) (\bar{\ell} \gamma^\mu \gamma^5 \ell)$$

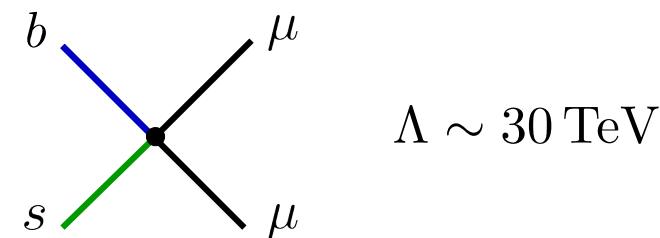
LFU in SM



- Completely consistent V-A solution:

$$\delta C_9^{\mu\mu} = -\delta C_{10}^{\mu\mu} = -0.41 \pm 0.09$$

NP preferred over SM 4.6σ



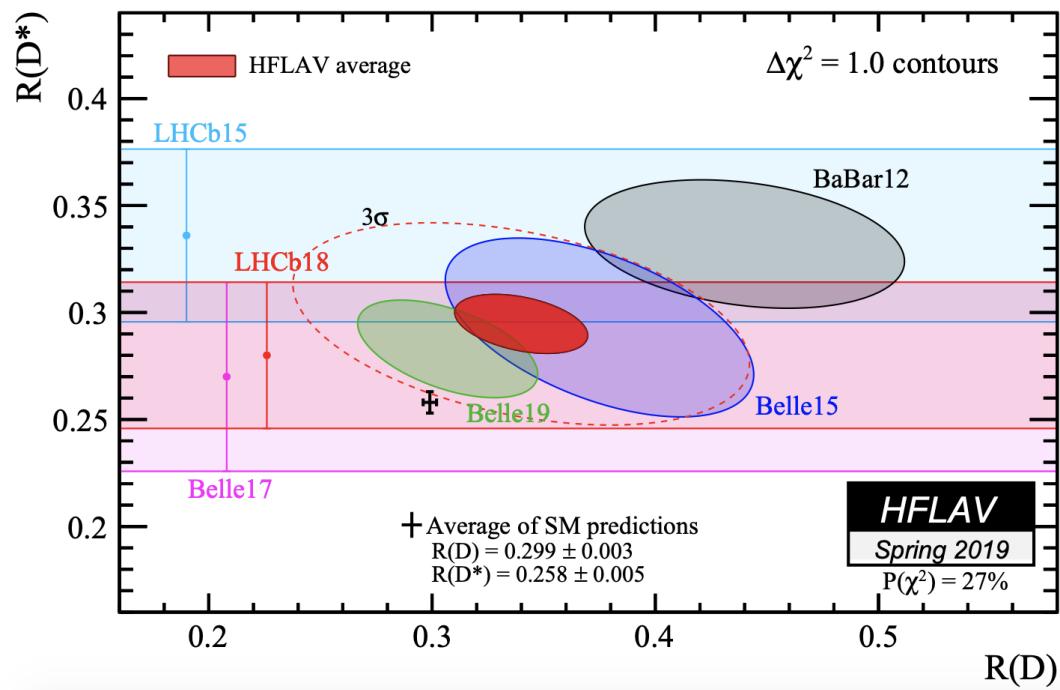
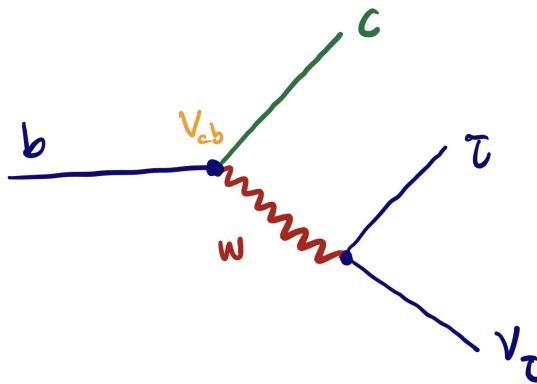
Large NP scale! (At tree level)

B-meson Anomalies – Charged currents

$$b \rightarrow c \ell \nu$$

- LFU ratio tau vs light leptons

$$R_{D^{(*)}} = \frac{\text{Br}(B \rightarrow D^{(*)} \tau \bar{\nu})}{\text{Br}(B \rightarrow D^{(*)} \ell \bar{\nu})} \Big|_{\ell=e,\mu}$$



$R_{D^{(*)}}^{\text{exp}} > R_{D^{(*)}}^{\text{SM}}$ Excess!



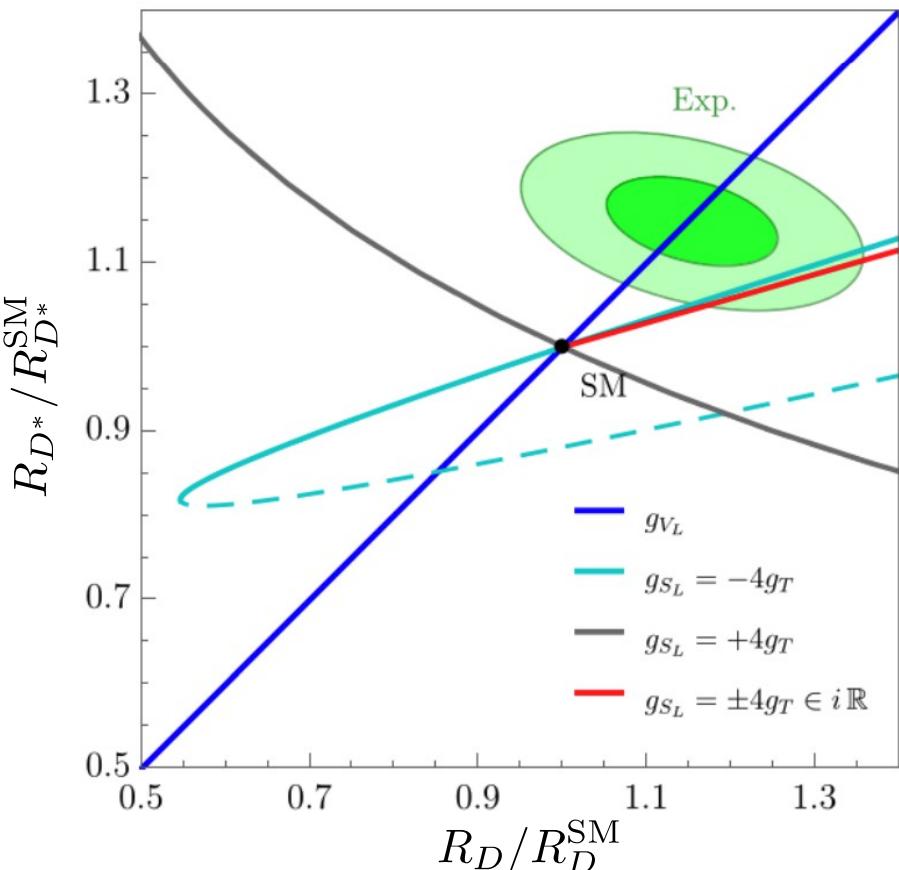
Soon joining the party!

Low energy effective theory - RD(*)

$$\mathcal{L}_{b \rightarrow c\tau\nu} = -2\sqrt{2}G_F V_{cb} \left[(1 + g_{V_L}) \mathcal{O}_{V_L} + g_{V_R} \mathcal{O}_{V_R} + g_{S_L} \mathcal{O}_{S_L} + g_{S_R} \mathcal{O}_{S_R} + g_T \mathcal{O}_T \right]$$

$\mathcal{O}_{V_L} = (\bar{c}_L \gamma^\mu b_L) (\bar{\tau}_L \gamma^\mu \nu_\tau)$	$\mathcal{O}_{S_L} = (\bar{c}_R b_L) (\bar{\tau}_R \nu_\tau)$	$\mathcal{O}_T = (\bar{c}_R \sigma^{\mu\nu} b_L) (\bar{\tau}_R \sigma^{\mu\nu} \nu_\tau)$
$\mathcal{O}_{V_R} = (\bar{c}_R \gamma^\mu b_R) (\bar{\tau}_L \gamma^\mu \nu_\tau)$	$\mathcal{O}_{S_R} = (\bar{c}_L b_R) (\bar{\tau}_R \nu_\tau)$	

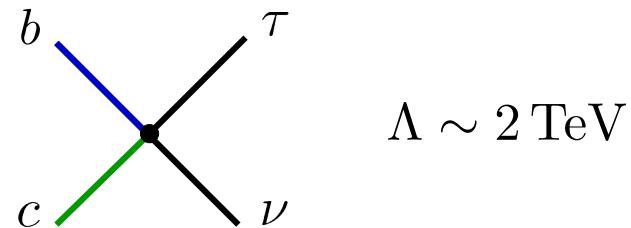
- Fit to R_D , R_{D^*} , $\mathcal{B}(B_c \rightarrow \tau\nu)$



- **V-A** only single operator fitting the anomaly

$$g_{V_L} \in [0.05, 0.09]$$

- 2 operators: **Scalar + Tensor** solutions are also fine

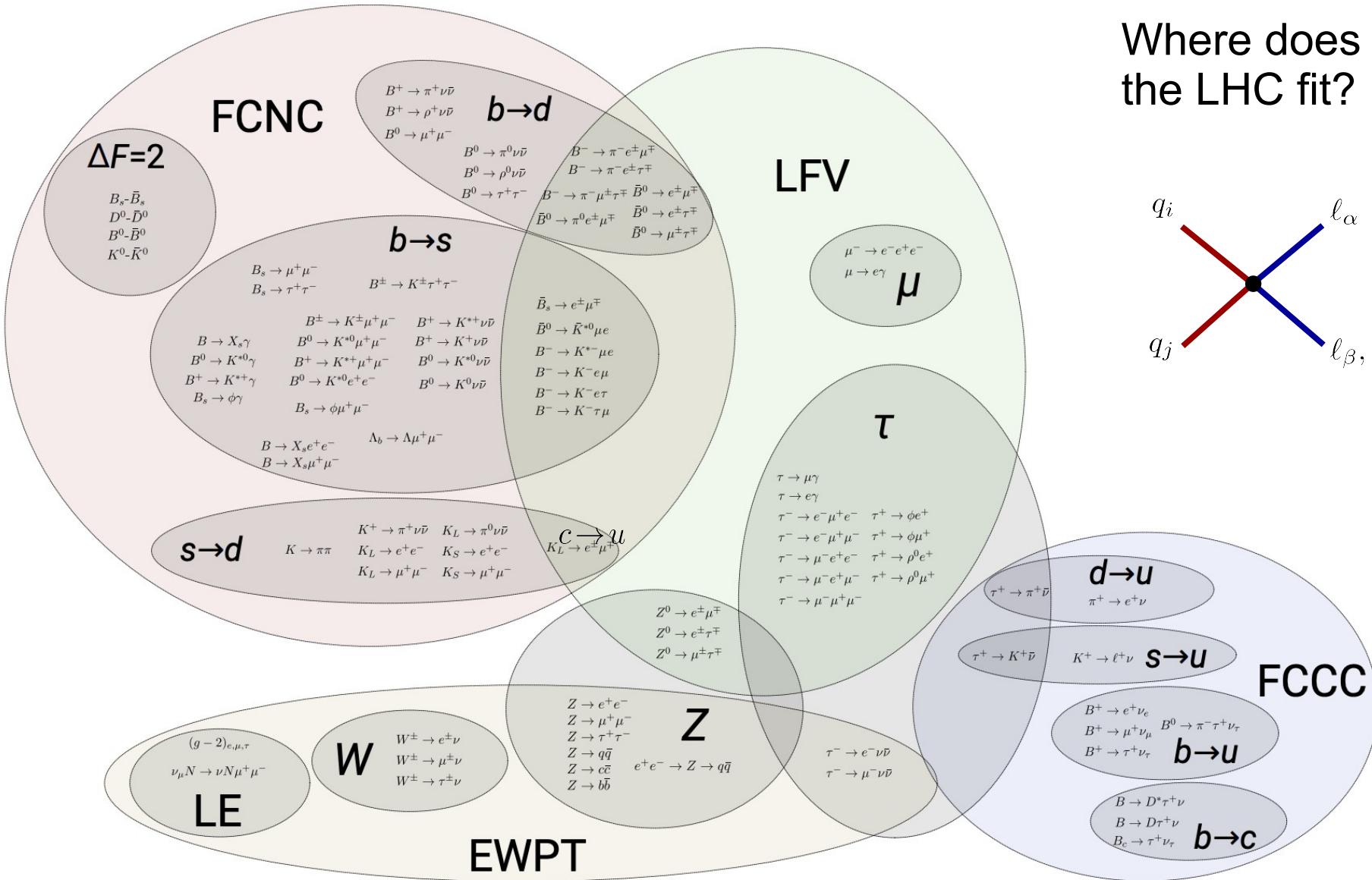


$$\Lambda \sim 2 \text{ TeV}$$

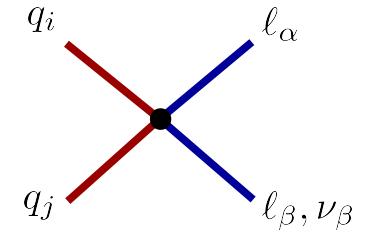
Strong physics case
for High-pT LHC!!

- Flavor @ LHC -

Semileptonics: observable Landscape



Where does
the LHC fit?



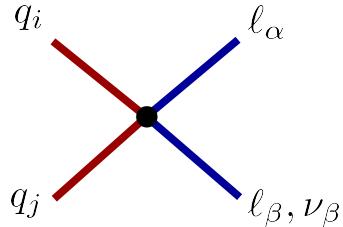
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Framework

- Semi-leptonics at the LHC: Drell-Yan process

$$q_i \bar{q}_j \rightarrow \ell_\alpha^\pm \ell_\beta^\mp$$

$$q_i \bar{q}_j \rightarrow \ell_\alpha^\pm \nu_\beta$$



$$\mathcal{L}_{\text{eff}} \supset \sum_I \sum_{ij,\alpha\beta} \left(\frac{C_I^{ij\alpha\beta}}{v^2} \right) \mathcal{O}_I^{ij\alpha\beta}$$

Eff. coeff.	Operator	SMEFT
$C_{V_{LL}}^{ij\alpha\beta}$	$(\bar{q}_{Li} \gamma_\mu q_{Lj}) (\bar{\ell}_{L\alpha} \gamma^\mu \ell_{L\beta})$	$\mathcal{O}_{lq}^{(1)}, \mathcal{O}_{lq}^{(3)}$
$C_{V_{RR}}^{ij\alpha\beta}$	$(\bar{q}_{Ri} \gamma_\mu q_{Rj}) (\bar{\ell}_{R\alpha} \gamma^\mu \ell_{R\beta})$	$\mathcal{O}_{ed}, \mathcal{O}_{eu}$
$C_{V_{LR}}^{ij\alpha\beta}$	$(\bar{q}_{Li} \gamma_\mu q_{Lj}) (\bar{\ell}_{R\alpha} \gamma^\mu \ell_{R\beta})$	\mathcal{O}_{qe}
$C_{V_{RL}}^{ij\alpha\beta}$	$(\bar{q}_{Ri} \gamma_\mu q_{Rj}) (\bar{\ell}_{L\alpha} \gamma^\mu \ell_{L\beta})$	$\mathcal{O}_{lu}, \mathcal{O}_{ld}$
$C_{S_R}^{ij\alpha\beta}$	$(\bar{q}_{Ri} q_{Lj}) (\bar{\ell}_{L\alpha} \ell_{R\beta}) + \text{h.c.}$	\mathcal{O}_{ledq}
$C_{S_L}^{ij\alpha\beta}$	$(\bar{q}_{Li} q_{Rj}) (\bar{\ell}_{L\alpha} \ell_{R\beta}) + \text{h.c.}$	$\mathcal{O}_{lequ}^{(1)}$
$C_T^{ij\alpha\beta}$	$(\bar{q}_{Li} \sigma_{\mu\nu} q_{Rj}) (\bar{\ell}_{L\alpha} \sigma^{\mu\nu} \ell_{R\beta}) + \text{h.c.}$	$\mathcal{O}_{lequ}^{(3)}$

- Two sources of flavor in high-energy proton-proton collisions:

$$\sigma(pp \rightarrow \ell^\alpha \ell^\beta) = \mathcal{L} \otimes \mathcal{H}^{\alpha\beta} = \sum_{ij} \int \frac{d\tau}{\tau} \mathcal{L}_{q_i \bar{q}_j}(\tau) \hat{\sigma}^{ij\alpha\beta}(\tau) \quad \tau \equiv \frac{\hat{s}}{s} = \frac{m_{\ell_\alpha \ell_\beta}^2}{s}$$

Parton-parton Luminosities
(non-perturbative)

$$\mathcal{L}_{q_i \bar{q}_j}(\tau) = \tau \int_\tau^1 \frac{dx}{x} [f_{q_i}(x, \mu_F) f_{\bar{q}_j}(\tau/x, \mu_F) + (q_i \leftrightarrow \bar{q}_j)]$$

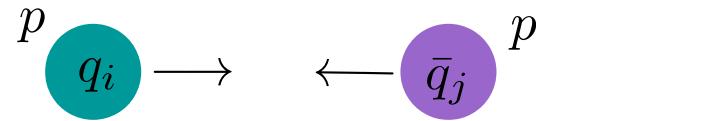
'PDFs'

Hard process

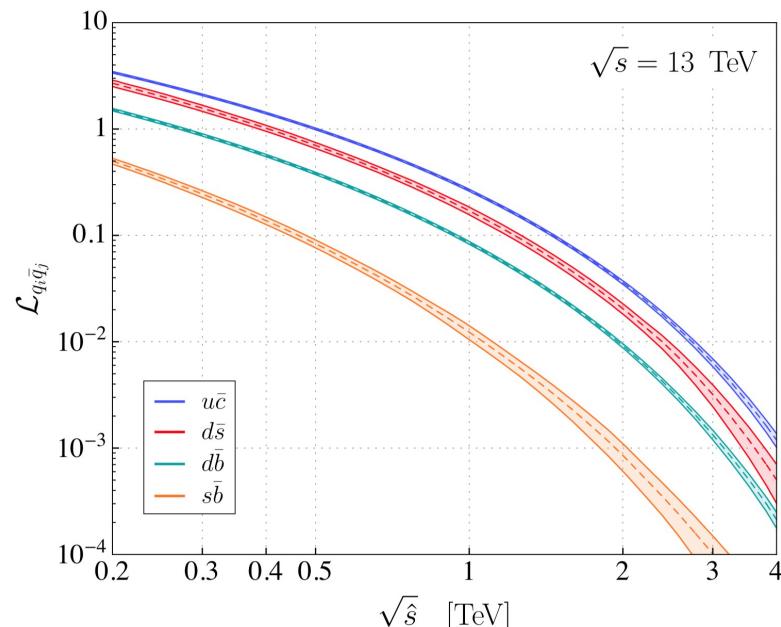
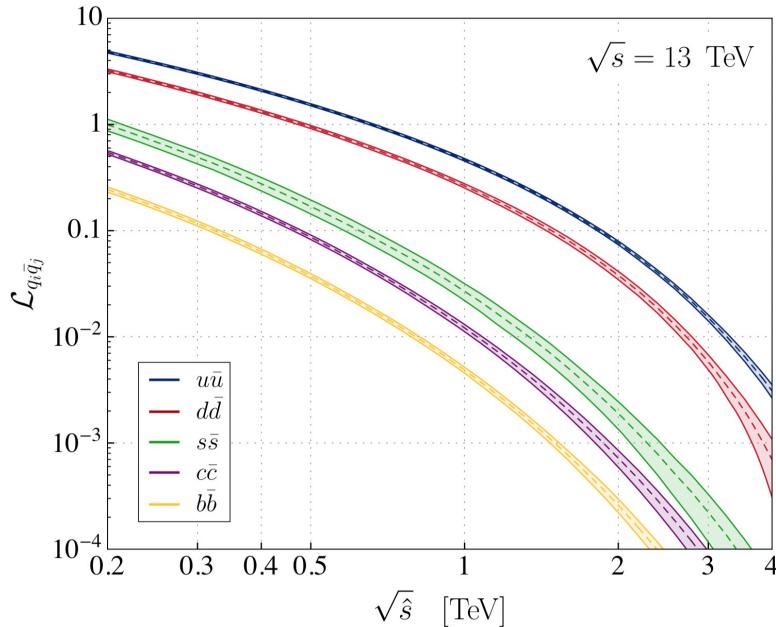
$$\hat{\sigma}^{ij\alpha\beta} \equiv \hat{\sigma}(q_i \bar{q}_j \rightarrow \ell_\alpha \ell_\beta) = \frac{\hat{s}}{v^4} \sum_{IJ} C_I^{ij\alpha\beta} C_J^{ij\alpha\beta*} M_{IJ}$$

$$\mathcal{G} = \mathcal{G}_q \otimes \mathcal{G}_\ell$$

- Luminosity functions: source of flavor hierarchies



$$\mathcal{L}_{q_i \bar{q}_j}(\tau) = \tau \int_{\tau}^1 \frac{dx}{x} [f_{q_i}(x, \mu_F) f_{\bar{q}_j}(\tau/x, \mu_F) + (q_i \leftrightarrow \bar{q}_j)] \quad \tau = \hat{s}/s = m_{\ell_\alpha \ell_\beta}^2/s$$



- example: LFV

$$\sigma(pp \rightarrow \mu\tau) = \frac{s}{144\pi v^4} \sum_{ij} |C_{V_L}^{ij\mu\tau}|^2 K_{ij} \quad (\text{V-A})$$

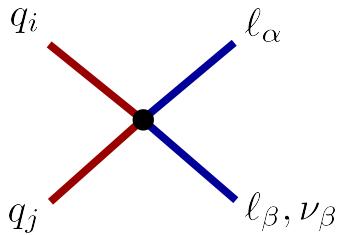
$$K_{ij} \equiv \int_{\tau_{\min}}^{\tau_{\max}} d\tau \mathcal{L}_{q_i \bar{q}_j}(\tau) \quad \xrightarrow{m_{\mu\tau} \in [1 \text{ TeV}, 5 \text{ TeV}]}$$

“PDF” tensor containing the proton flavor hierarchies

$$K_u = \kappa \begin{pmatrix} u & c \\ 1 & 0.5 & 0 \\ 0.03 & 0.01 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad K_d = \frac{\kappa}{2} \begin{pmatrix} d & s & b \\ 1 & 0.6 & 0.3 \\ 0.1 & 0.07 & 0.03 \\ 0.04 & 0.02 & 0.01 \end{pmatrix}$$

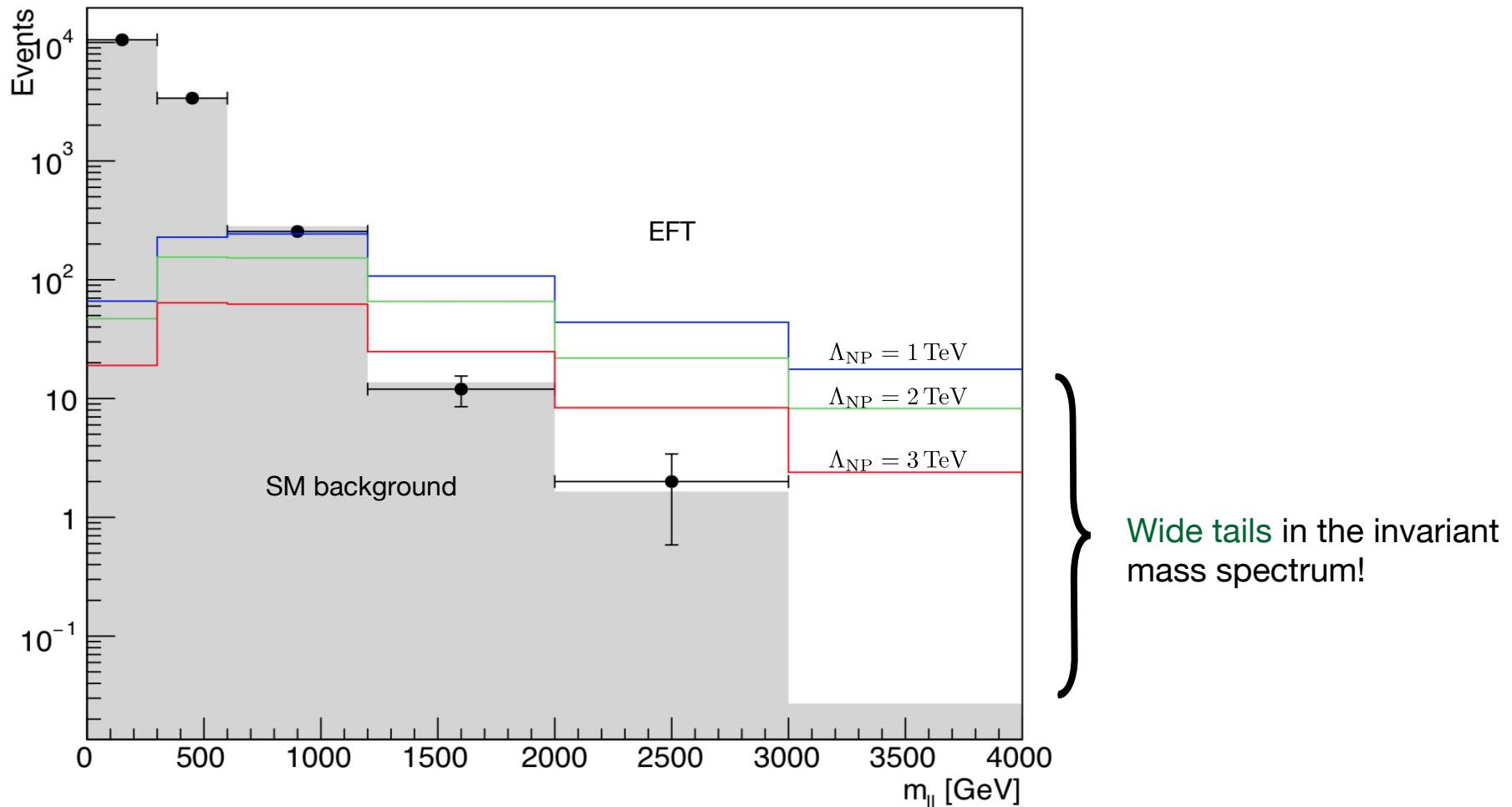
Very hierarchical!

'High-pT' Tails



$$q_i \bar{q}_j \rightarrow \ell_\alpha^\pm \ell_\beta^\mp$$

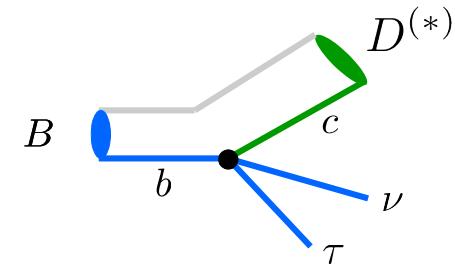
$$q_i \bar{q}_j \rightarrow \ell_\alpha^\pm \nu_\beta$$



Di-tau tails - RD(*)

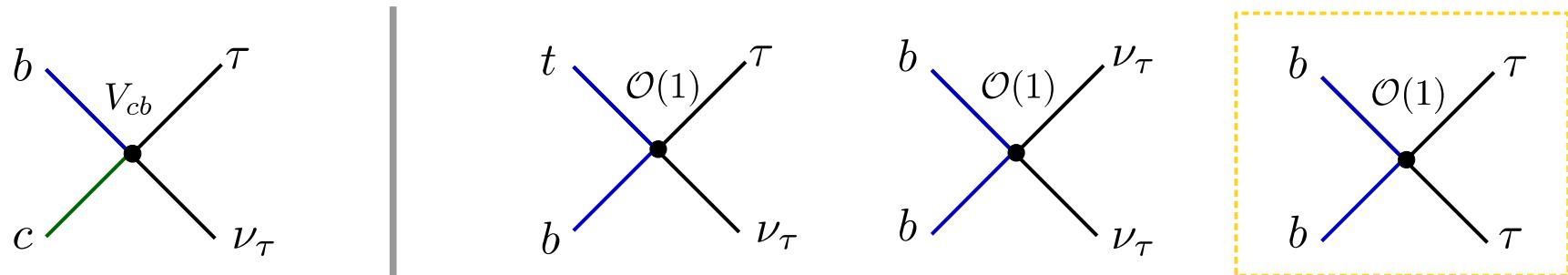
- SM EFT (d=6):

$$\begin{aligned}\mathcal{L}_{eff} \supset & C_{lq}^{(3) \ ij\alpha\beta} (\bar{Q}^i \gamma_\mu \sigma^a Q^j)(\bar{L}^\alpha \gamma^\mu \sigma^a L^\beta) \\ & + C_{ledq}^{ij\alpha\beta} (\bar{d}^i Q^j)(\bar{L}^\alpha e^\beta) \\ & + C_{lequ}^{(1) \ ij\alpha\beta} (\bar{Q}_i u^j) i\sigma^2 (\bar{L}^\alpha e^\beta) \\ & + C_{lequ}^{(3) \ ij\alpha\beta} (\bar{Q}^i \sigma_{\mu\nu} u^j)(\bar{L}^\alpha \sigma_{\mu\nu} e^\beta)\end{aligned}$$



- (Minimal) Flavor structure assumptions:

1. Alignment with **down quarks & charged leptons**: $Q_i = (V_{CKM}^{ij*} u_L^j, d_L^i)^T$ $L_i = (U_{PMNS}^{ij*} \nu_L^j, \ell_L^i)^T$
2. Dominant effects in **3rd generation fermions**: $C^{ij\alpha\beta} \approx C \delta_{i3} \delta_{j3} \delta_{3\alpha} \delta_{3\beta}$



V_{cb}^{-1} enhancement with respect to $b \rightarrow c$ transitions

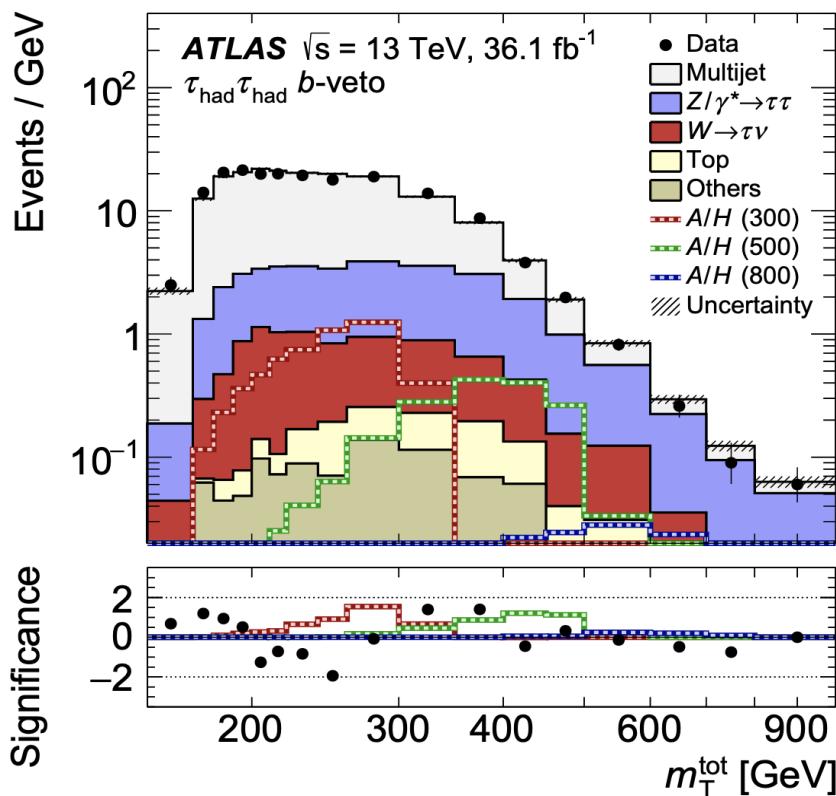
- Enhanced ditau production $b\bar{b} \rightarrow \tau\tau$ compared to mono-tau $b\bar{c} \rightarrow \tau\nu$

- We can probe generic NP solution to RD(*) in **Di-Tau** production at the LHC! $pp \rightarrow \tau^+ \tau^- + X$

Di-tau tails - RD(*)

Existing $\tau\tau$ resonance searches @ LHC

Z' Sequential SM	8 TeV	19.5 fb^{-1}
	13 TeV	3.2 fb^{-1}
MSSM Higgs	13 TeV	36.1 fb^{-1}
	13 TeV	139 fb^{-1}



Recast Z' di-Tau searches:

- Inclusive category $pp \rightarrow \tau^+ \tau^- + X$

- Hadronic mode $\tau_{\text{had}} \tau_{\text{had}}$

$$m_T^{\text{tot}} \equiv \sqrt{m_T^2(\tau_1, \tau_2) + m_T^2(E_T^{\text{miss}}, \tau_1) + m_T^2(E_T^{\text{miss}}, \tau_2)}$$

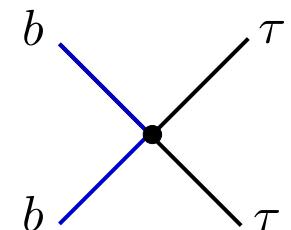
- High-mass tail cut $m_T^{\text{tot}} > 800 \text{ GeV}$

- Results for EFT limits @ 13 TeV, 3.2 fb^{-1} :

$$\mathcal{L}_{eff} \supset C_{lq}^{(3) \ ij\alpha\beta} (\bar{Q}^i \gamma_\mu \sigma^a Q^j)(\bar{L}^\alpha \gamma^\mu \sigma^a L^\beta)$$

$$|C_{lq}^{(3) bb\tau\tau}| < \frac{2.6}{(1 \text{ TeV})^2}$$

DAF, Greljo, Kamenik '16



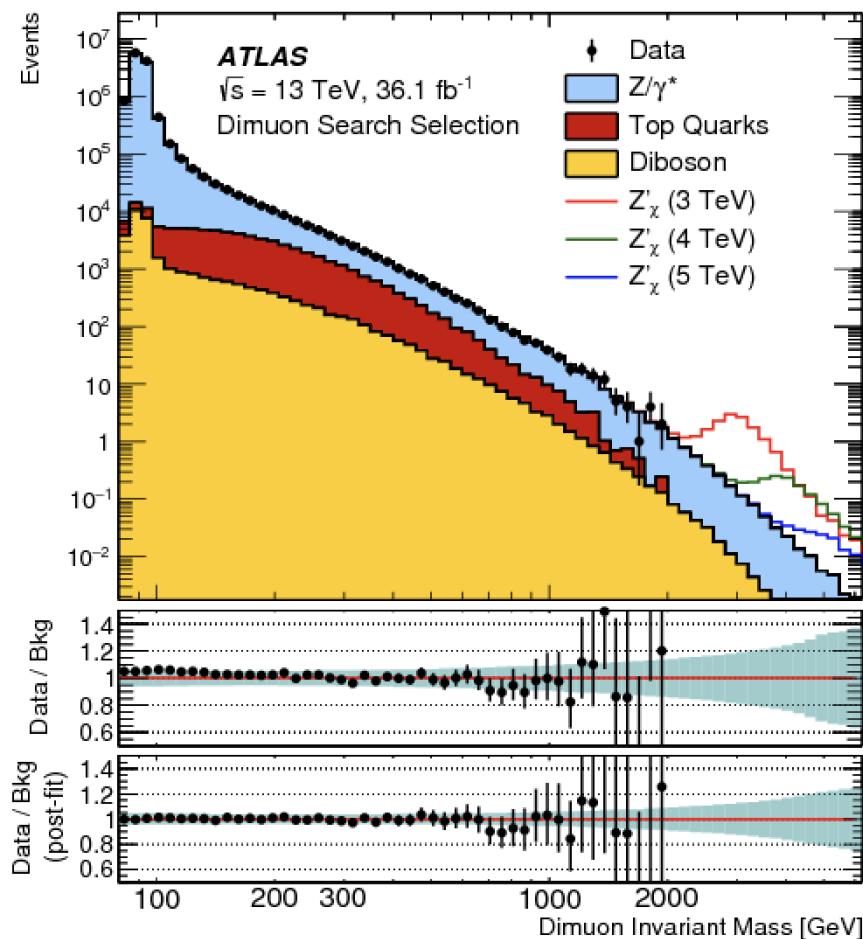
- LHC is becoming sensitive to the B-anomaly mass scale!

Di-muons tails - RK(*)

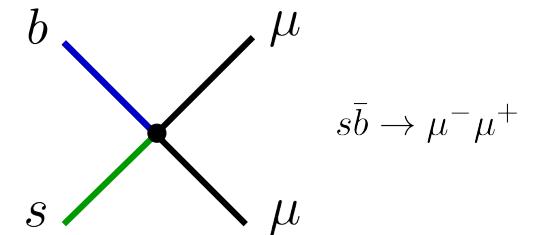
Greljo, Marzocca '17

$$\mathcal{L}^{\text{eff}} \supset \frac{C_{ij}^{U\mu}}{v^2} (\bar{u}_L^i \gamma_\mu u_L^j) (\bar{\mu}_L \gamma^\mu \mu_L) + \frac{C_{ij}^{D\mu}}{v^2} (\bar{d}_L^i \gamma_\mu d_L^j) (\bar{\mu}_L \gamma^\mu \mu_L)$$

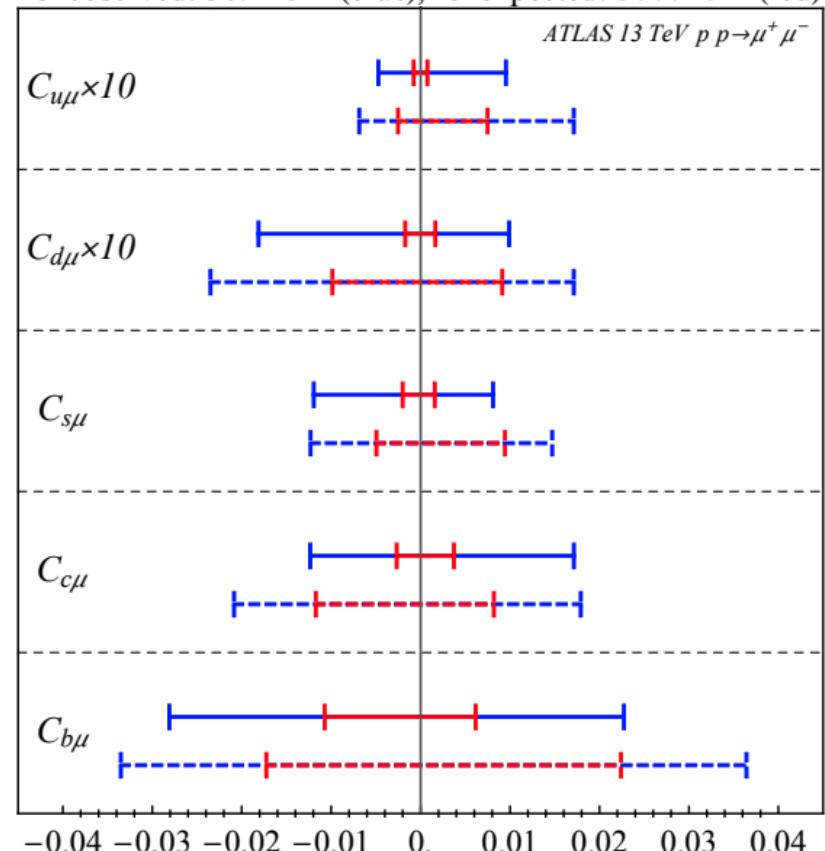
Light-lepton tails



$$C_{ij}^{D\mu} = \begin{pmatrix} C_{d\mu} & 0 & 0 \\ 0 & C_{s\mu} & C_{bs\mu}^* \\ 0 & C_{bs\mu} & C_{b\mu} \end{pmatrix}$$



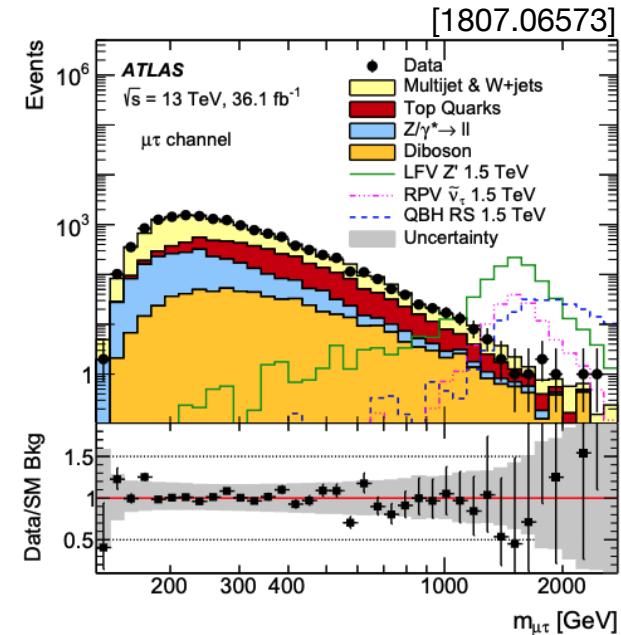
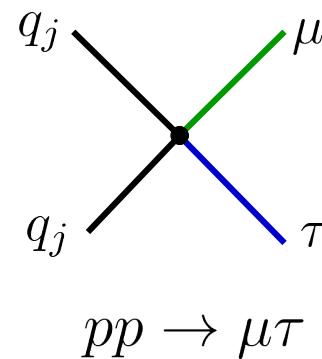
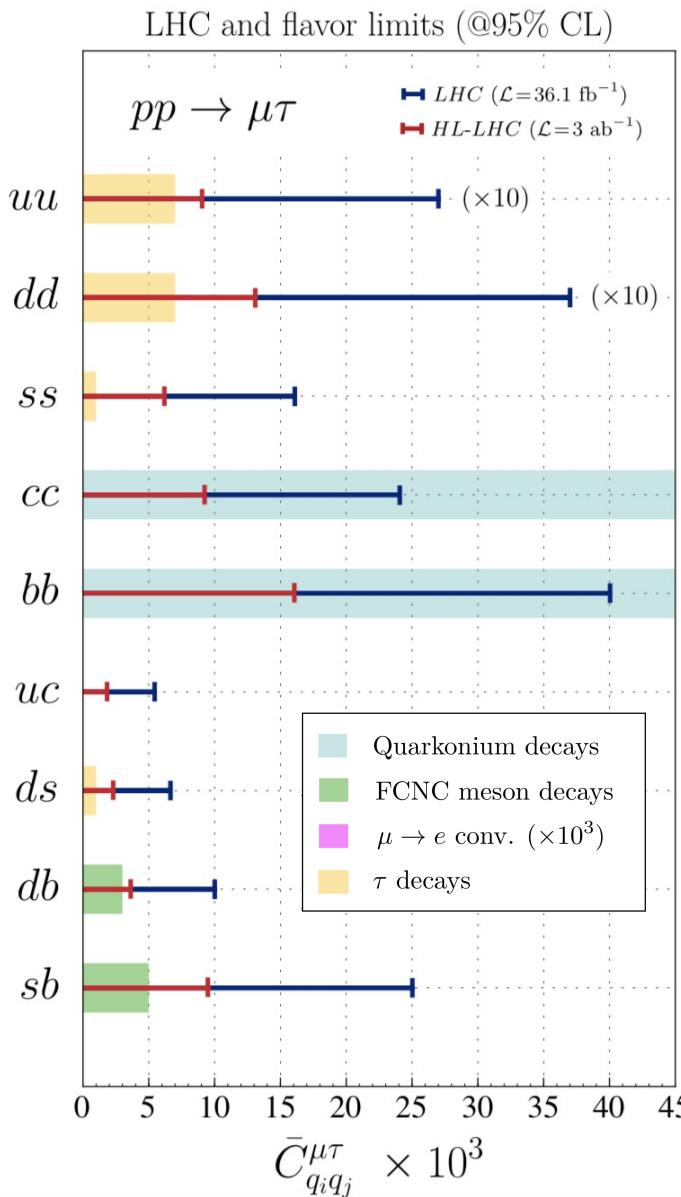
2 σ observed: 36.1 fb^{-1} (blue), 2 σ expected: 3000 fb^{-1} (red)



$\Lambda_{NP}(\text{valence}) > 8 \text{ TeV}$

$\Lambda_{NP}(\text{heavy sea}) > 1.5 \text{ TeV}$

LFV tails!



$\Lambda_{NP}(\text{valence}) > 8 \text{ TeV}$

- Recast results:

$\Lambda_{NP}(\text{sea}) > 1.5 \text{ TeV}$

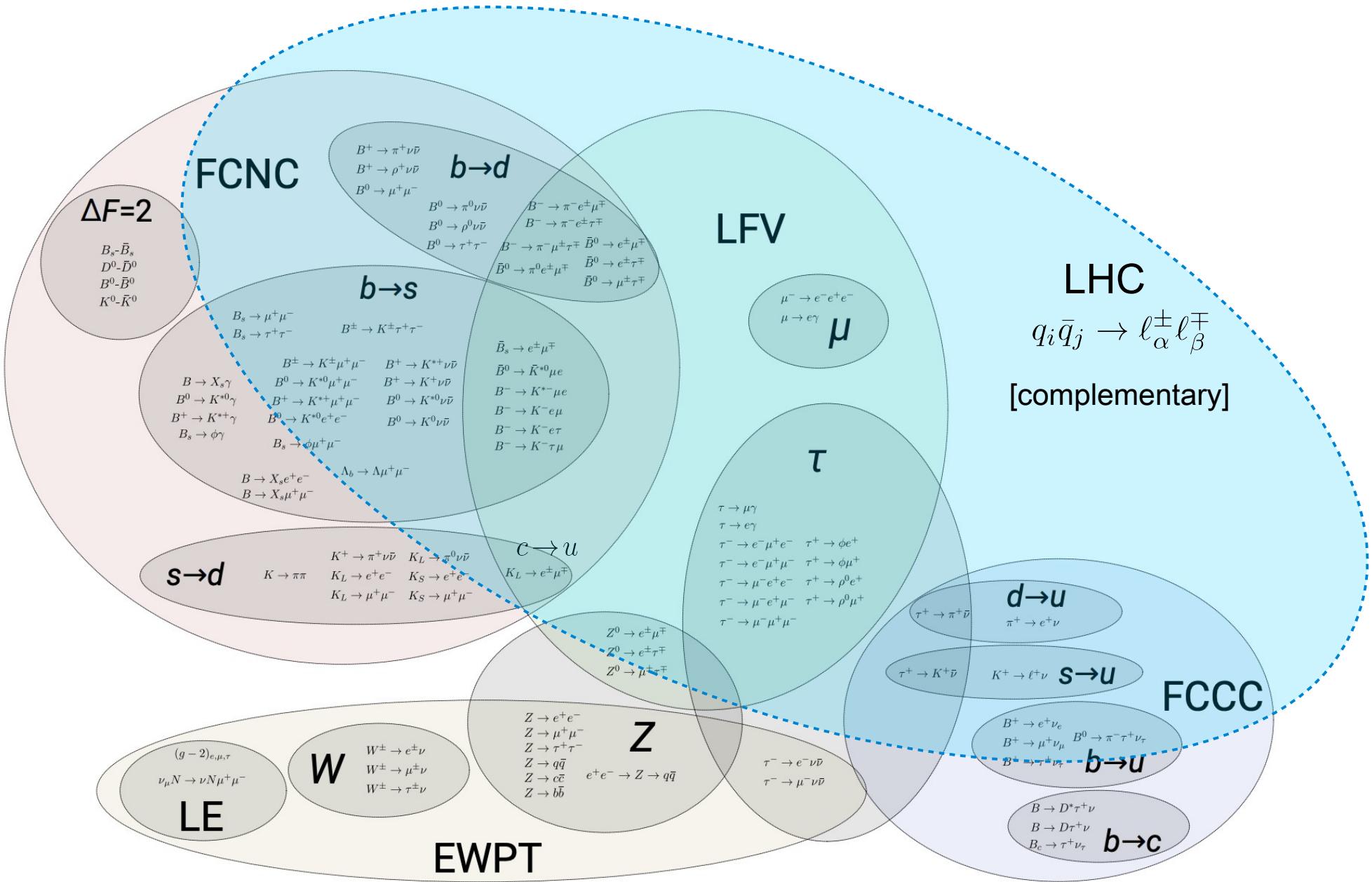
- Similar for results for $pp \rightarrow e\mu, e\tau$

$q\bar{q} \rightarrow \ell_\alpha^\pm \ell_\beta^\mp$ tails perform better than Flavor experiments.

$q\bar{q}' \rightarrow \ell_\alpha^\pm \ell_\beta^\mp$ tails worst than FCNC meson and tau decays.

Remarkable LHC / flavor complementarity!

Semileptonics: observables Landscape

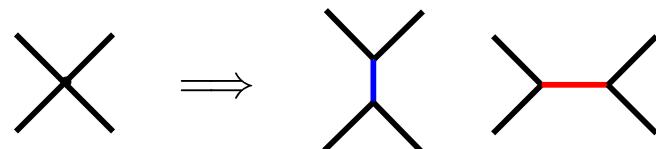


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B-anomalies:

- Leptoquark solutions -

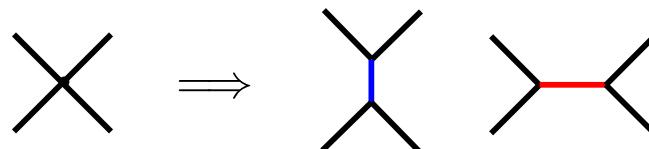
Simplified models - RD(*)



Necessary for reliable high-pT collider studies
Less parameters...
Paves the way towards UV complete models

mediators:	Colorless	Colorful
Spin 0	2HDM	Scalar Leptoquark
Spin 1	Vector Triplet W'	Vector Leptoquark
Flavor		
$b \rightarrow c$		
LHC phenomenology		
$b\bar{b} \rightarrow \tau^+\tau^-$		
Signature	di-Tau Resonance	Non-resonant excess in di-Tau tails

Simplified models - RD(*)



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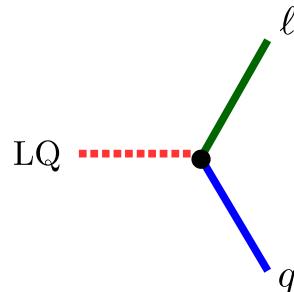
mediators:	Colorless		Colorful
Spin 0	2HDM		Scalar Leptoquark
Spin 1	Vector Triplet W'		Vector Leptoquark
Flavor	$b \rightarrow c$		
LHC phenomenology	$b\bar{b} \rightarrow \tau^+\tau^-$		
Signature	di-Tau Resonance		Non-resonant excess in di-Tau tails

Leptoquarks

- Quintessential 'semi-leptonic' state!

Hypothetical Scalar / Vector boson

Color triplet, $Q_{em} = \{\pm \frac{1}{3}, \pm \frac{2}{3}, \pm \frac{4}{3}, \pm \frac{5}{3}\}$



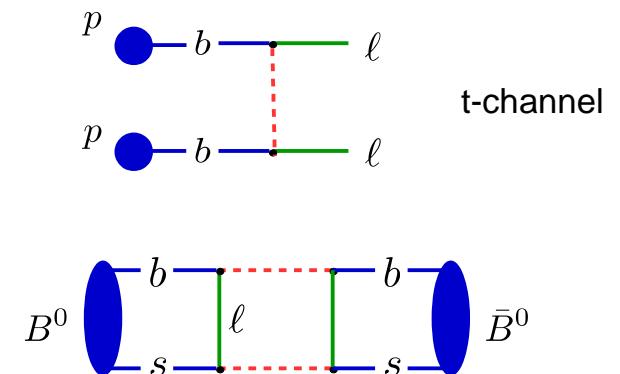
Leptoquark Bestiary

$(SU(3), SU(2), U(1))$	Spin	Symbol
$(\bar{\mathbf{3}}, \mathbf{3}, 1/3)$	0	S_3
$(\mathbf{3}, \mathbf{2}, 7/6)$	0	R_2
$(\mathbf{3}, \mathbf{2}, 1/6)$	0	\tilde{R}_2
$(\bar{\mathbf{3}}, \mathbf{1}, 4/3)$	0	\tilde{S}_1
$(\bar{\mathbf{3}}, \mathbf{1}, 1/3)$	0	S_1
$(\bar{\mathbf{3}}, \mathbf{1}, -2/3)$	0	\bar{S}_1
$(\mathbf{3}, \mathbf{3}, 2/3)$	1	U_3
$(\bar{\mathbf{3}}, \mathbf{2}, 5/6)$	1	V_2
$(\bar{\mathbf{3}}, \mathbf{2}, -1/6)$	1	\tilde{V}_2
$(\mathbf{3}, \mathbf{1}, 5/3)$	1	\tilde{U}_1
$(\mathbf{3}, \mathbf{1}, 2/3)$	1	U_1
$(\mathbf{3}, \mathbf{1}, -1/3)$	1	\bar{U}_1

- Some general features:

- Non-resonant Drell-Yan production
- No 4-quark / 4-lepton effective interactions at tree level
- Induces semi-leptonic tree level LFV $K_L \rightarrow e\mu$
- Some leptoquarks can mediate proton decay...

Relevant LQs for B-anomalies

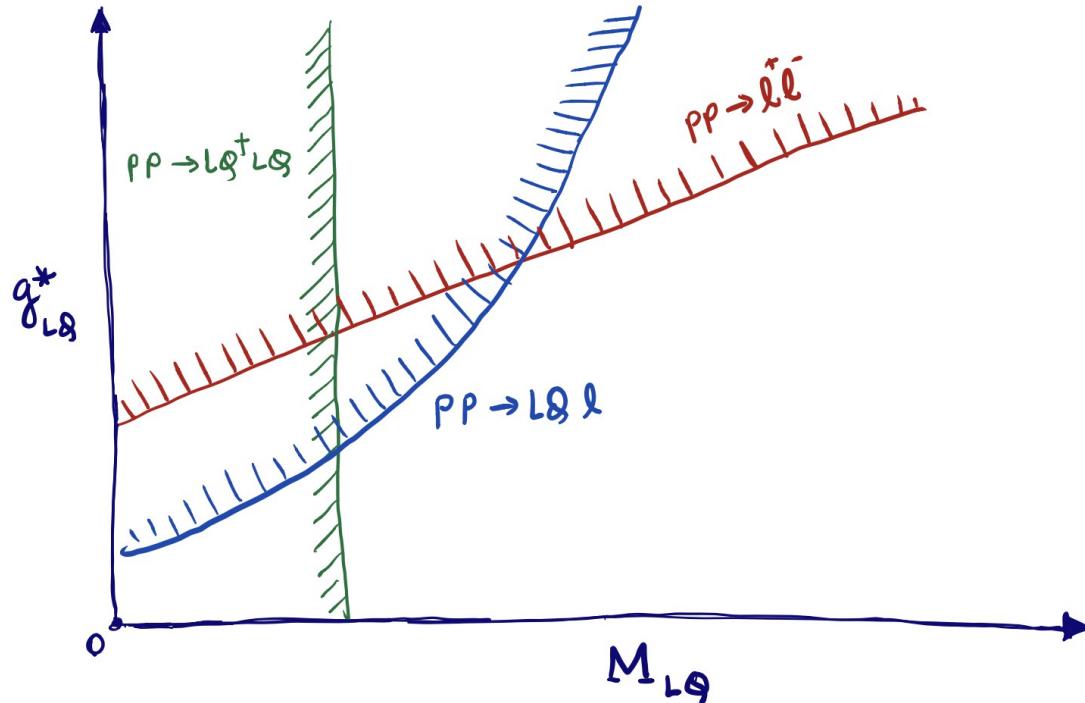
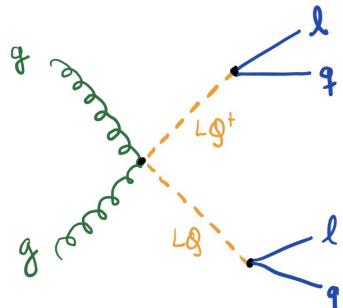


Leptoquarks @ LHC

LQ pair production

$$pp \rightarrow LQ^* LQ \rightarrow \ell\ell qq$$

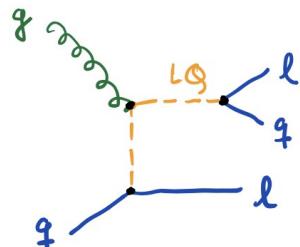
$$\sigma \sim g_*^0$$



Single LQ production

$$pp \rightarrow LQ \ell \rightarrow \ell\ell q$$

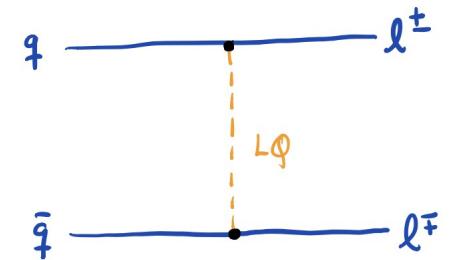
$$\sigma \sim g_*^2$$



Non-resonant Drell-Yan

$$pp \rightarrow \mu\mu, \tau\tau$$

$$\sigma \sim g_*^4$$



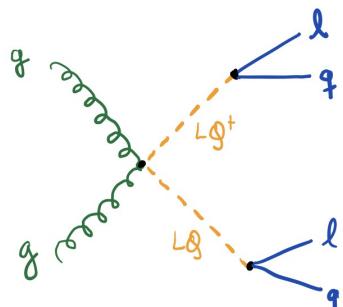
Leptoquarks @ LHC

Run 2 pair production limits

LQ pair production

$$pp \rightarrow \text{LQ}^*\text{LQ} \rightarrow \ell\ell qq$$

$$\sigma \sim g_*^0$$



\implies

Decays	Scalar LQ limits	Vector LQ limits	$\mathcal{L}_{\text{int}} / \text{Ref.}$
$jj\tau\bar{\tau}$	—	—	—
$b\bar{b}\tau\bar{\tau}$	1.0 (0.8) TeV	1.5 (1.3) TeV	36 fb^{-1} [39]
$t\bar{t}\tau\bar{\tau}$	1.4 (1.2) TeV	2.0 (1.8) TeV	140 fb^{-1} [40]
$jj\mu\bar{\mu}$	1.7 (1.4) TeV	2.3 (2.1) TeV	140 fb^{-1} [41]
$b\bar{b}\mu\bar{\mu}$	1.7 (1.5) TeV	2.3 (2.1) TeV	140 fb^{-1} [41]
$t\bar{t}\mu\bar{\mu}$	1.5 (1.3) TeV	2.0 (1.8) TeV	140 fb^{-1} [42]
$jj\nu\bar{\nu}$	1.0 (0.6) TeV	1.8 (1.5) TeV	36 fb^{-1} [43]
$b\bar{b}\nu\bar{\nu}$	1.1 (0.8) TeV	1.8 (1.5) TeV	36 fb^{-1} [43]
$t\bar{t}\nu\bar{\nu}$	1.2 (0.9) TeV	1.8 (1.6) TeV	140 fb^{-1} [44]

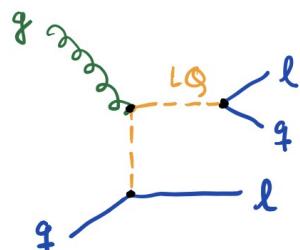


$$\text{BR} = 1 (0.5)$$

Single LQ production

$$pp \rightarrow \text{LQ} \ell \rightarrow \ell\ell q$$

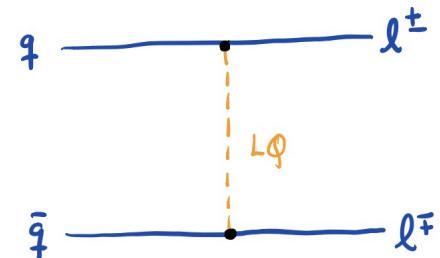
$$\sigma \sim g_*^2$$



Non-resonant Drell-Yan

$$pp \rightarrow \mu\mu, \tau\tau$$

$$\sigma \sim g_*^4$$

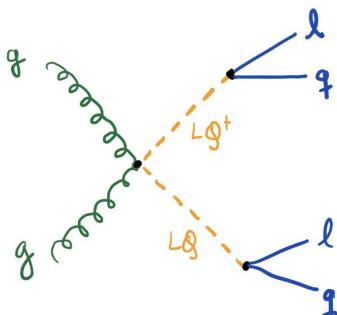


Leptoquarks @ LHC

LQ pair production

$$pp \rightarrow \text{LQ}^*\text{LQ} \rightarrow \ell\ell qq$$

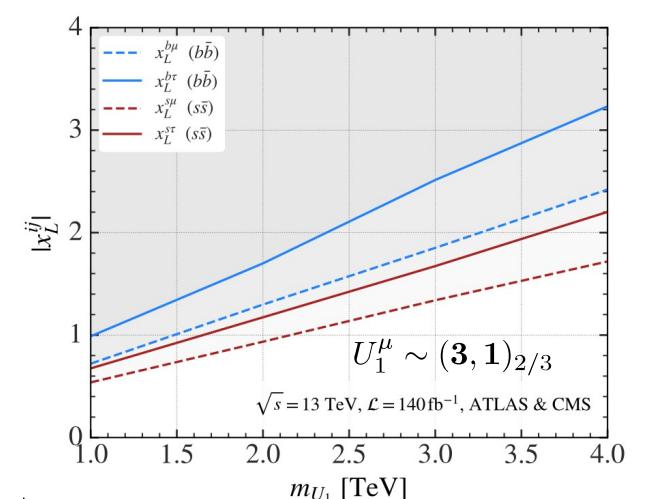
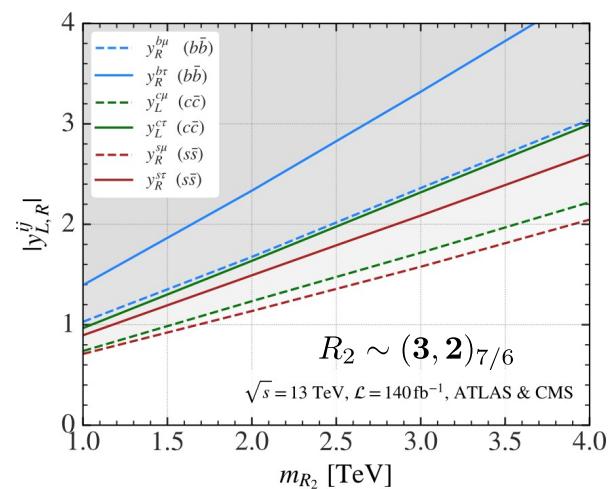
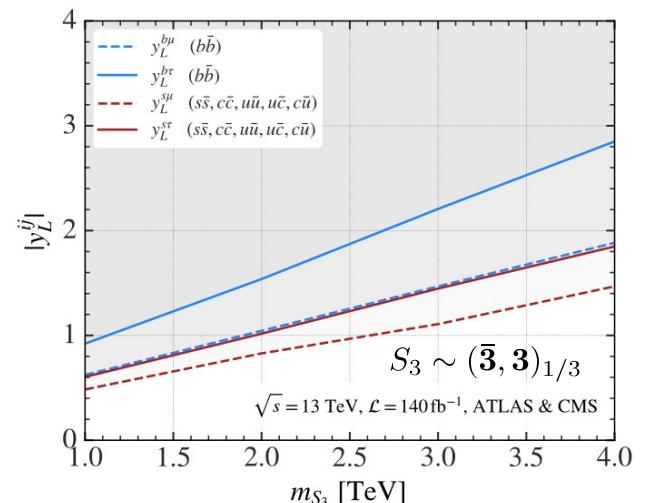
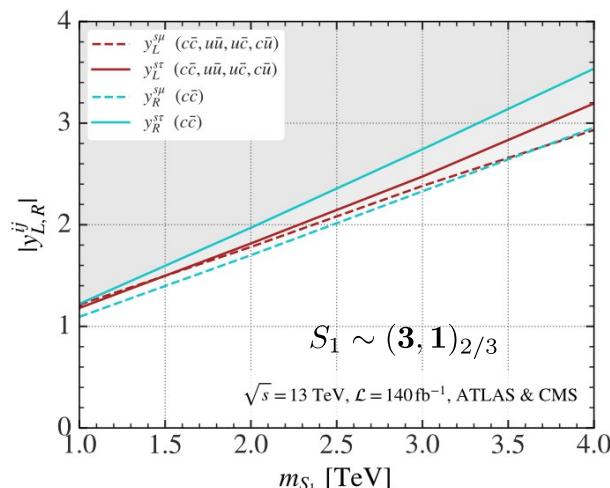
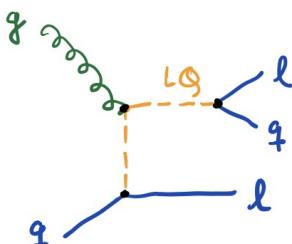
$$\sigma \sim g_*^0$$



Single LQ production

$$pp \rightarrow \text{LQ} \ell \rightarrow \ell\ell qq$$

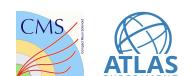
$$\sigma \sim g_*^2$$



Non-resonant Drell-Yan

$$pp \rightarrow \mu\mu, \tau\tau$$

$$\sigma \sim g_*^4$$

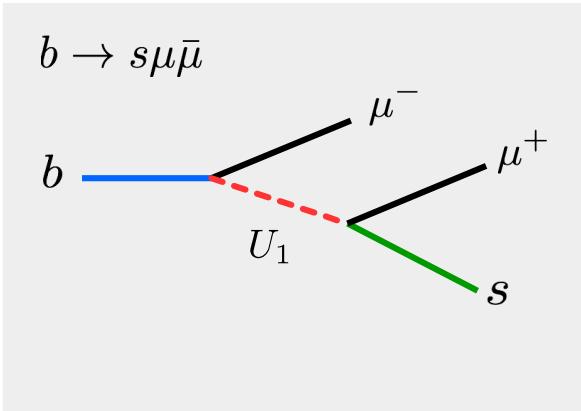


Run-2 Drell-Yan Recast limits

Include previously ignored
interference effects

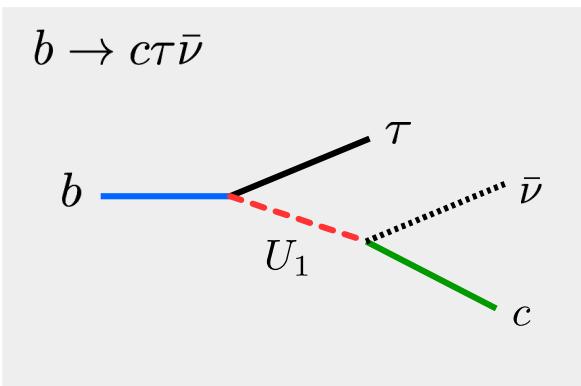
U₁ Vector LQ

$$U_1^\mu \sim (\mathbf{3}, \mathbf{1})_{2/3} \quad \mathcal{L}_{U1} = U_1^\mu \left[x_L^{ij} (\bar{d}_{Li} \gamma_\mu \ell_{Lj}) + (V x_L)^{ij} (\bar{u}_{Li} \gamma_\mu \nu_{Lj}) - x_R^{ij} (\bar{d}_{Ri} \gamma_\mu \ell_{Rj}) \right] + \text{h.c.}$$



$$x_L = \begin{pmatrix} 0 & 0 & 0 \\ 0 & x_L^{s\mu} & 0 \\ 0 & x_L^{b\mu} & 0 \end{pmatrix} \quad C_9 = -C_{10} = -\frac{\pi v^2}{V_{tb} V_{ts}^* \alpha} \frac{x_L^{s\mu} (x_L^{b\mu})^*}{m_U^2}$$

$$x_R = 0 \quad R_{K^{(*)}} < R_{K^{(*)}}^{\text{SM}}$$



$$x_L = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & x_L^{b\tau} \end{pmatrix} \quad g_{V_L} = -\frac{v^2 |x_L^{b\tau}|^2}{2m_U^2}$$

$$x_R = 0 \quad R_{D^{(*)}} > R_{D^{(*)}}^{\text{SM}}$$

- This solution can accomodate both anomalies! [Butazzo, Greljo, Isidori, Marzocca \[1706.07808\]](#)
- This solution has no (tree-level) contribution to $b \rightarrow s\nu\bar{\nu}$

Leptoquark ‘catalog’ (spring 2021 edition)

Model	$R_{K(*)}$	$R_{D(*)}$	$R_{K(*)} \& R_{D(*)}$
$S_1 = (3, 1)_{-1/3}$	✗	✓	✗
$R_2 = (3, 2)_{7/6}$	✗	✓	✗
$\tilde{R}_2 = (3, 2)_{1/6}$	✗	✗	✗
$S_3 = (3, 3)_{-1/3}$	✓	✗	✗
$U_1 = (3, 1)_{2/3}$	✓	✓	✓
$U_3 = (3, 3)_{2/3}$	✓	✗	✗

 Scalars
Vectors

Angelescu, Becirevic, DAF, Sumensari [1808.08179]

V-A structure

$$x_L = \begin{pmatrix} 0 & 0 & 0 \\ 0 & x_L^{s\mu} & x_L^{s\tau} \\ 0 & x_L^{b\mu} & x_L^{b\tau} \end{pmatrix}$$

Only state that can solve both B-anomalies!

Butazzo et al [1706.07808]

- UV completion necessary: e.g. ‘4321’ models

$$G_{4321} = SU(4) \times SU(3)_{c'} \times SU(2)_L \times U(1)_{Y'} \rightarrow G_{SM}$$

Di Luzio et al [1708.08450]

Bordone, Cornella, Fuentes-Martin, Isidori ‘18

No single scalar LQ solves both B-anomalies.

Two (or more) scalar LQ needed:

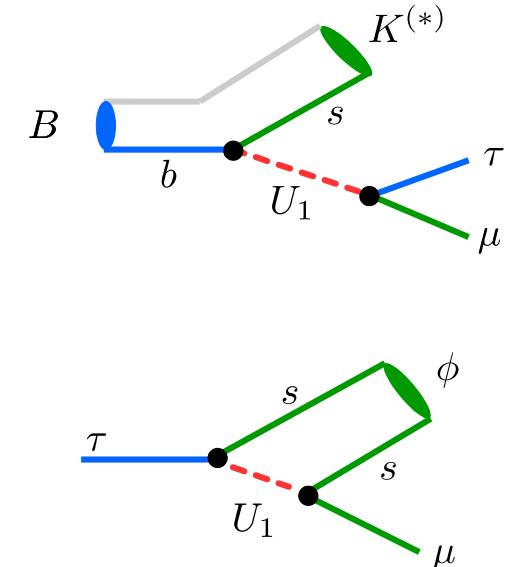
R_2 & S_3 (Scalar + Tensor & V-A) e.g. GUT inspired model Becirevic et al [1808.08179]

S_1 & S_3 (V-A) e.g. Strongly coupled model Marzocca [1803.10972]

Cornering U_1

Can we fully test the U_1 model using tree-level observables?

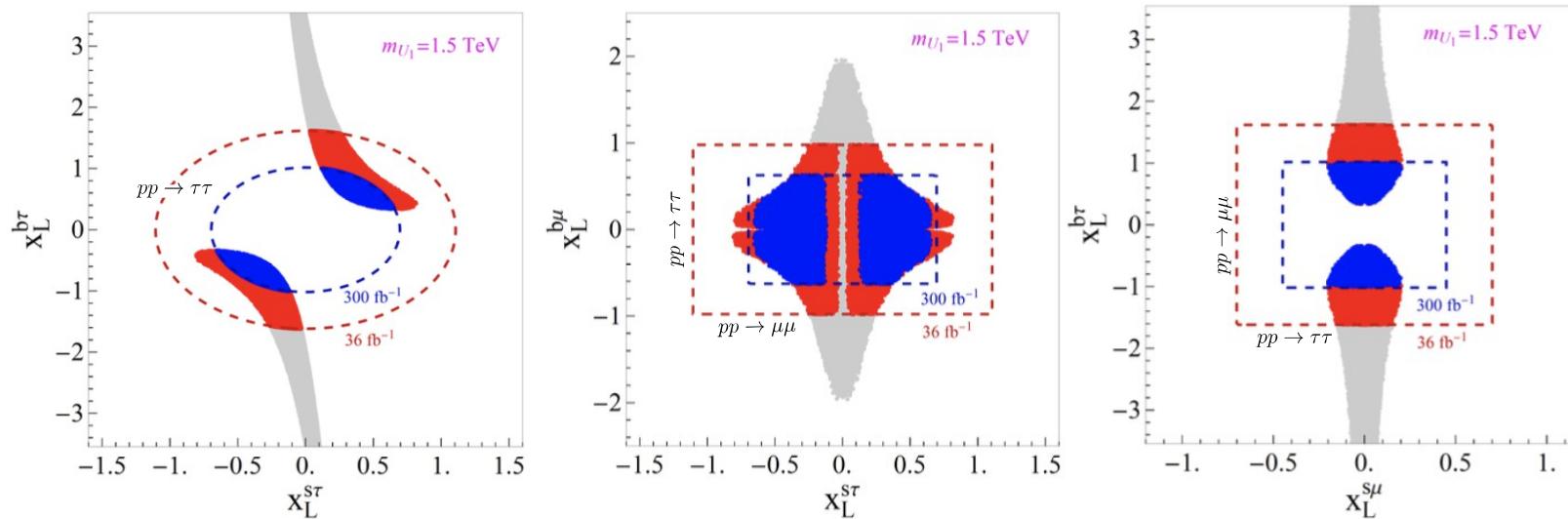
Complementarity between **LFV low-energy** and **LFC high-pT** observables



$pp \rightarrow \tau^+ \tau^- + X$	$B \rightarrow K \mu \tau$	$\tau \rightarrow \mu \phi$
$\begin{pmatrix} 0 & 0 & 0 \\ 0 & x_L^{s\mu} & x_L^{s\tau} \\ 0 & x_L^{b\mu} & x_L^{b\tau} \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 \\ 0 & x_L^{s\mu} & x_L^{s\tau} \\ 0 & x_L^{b\mu} & x_L^{b\tau} \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 \\ 0 & x_L^{s\mu} & x_L^{s\tau} \\ 0 & x_L^{b\mu} & x_L^{b\tau} \end{pmatrix}$

Fixed mass fit with tree-level observables

Angelescu, Becirevic, DAF, Sumensari [1808.08179]



Observable
$b \rightarrow s \mu \mu$
$b \rightarrow c \tau \nu$
$\mathcal{B}(\tau \rightarrow \mu \phi)$
$\mathcal{B}(B \rightarrow \tau \nu)$
$\mathcal{B}(D_s \rightarrow \mu \nu)$
$\mathcal{B}(D_s \rightarrow \tau \nu)$
$r_K^{e/\mu}$
$r_K^{\tau/\mu}$
$R_D^{\mu/e}$

- LHC di-lep allowed regions at **36/fb (blue)** & **300/fb (red)**

- U_1 - strange - tau coupling must be non-zero!

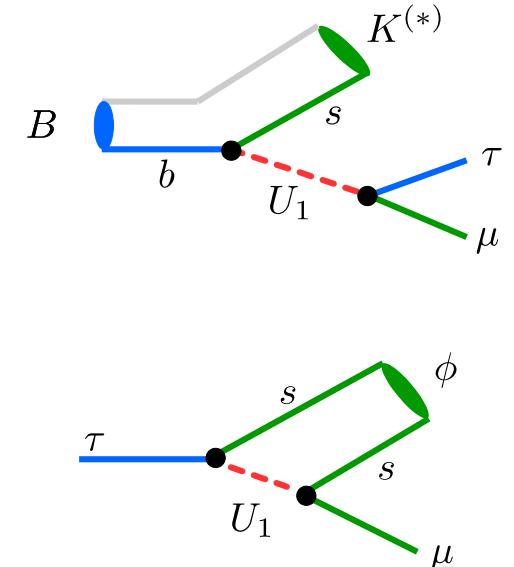
$$g_{VL} = -\frac{v^2}{2m_U^2} \left[|x_L^{b\tau}|^2 + \frac{V_{cs}}{V_{cb}} x_L^{s\tau} (x_L^{b\tau})^* \right]$$

$$x_{s\tau} = \mathcal{O}(1) \times V_{cb}$$

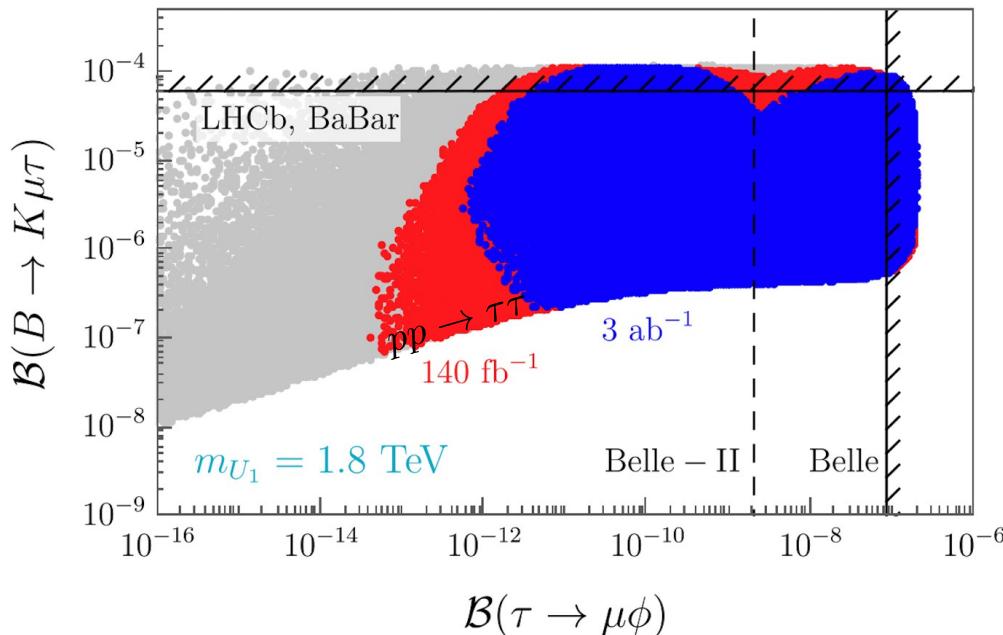
Cornering U_1

Can we fully test the U_1 model using tree-level observables?

Complementarity between **LFV low-energy** and **LFC high-pT** observables



- Predictions for LFV low-energy observables for U_1



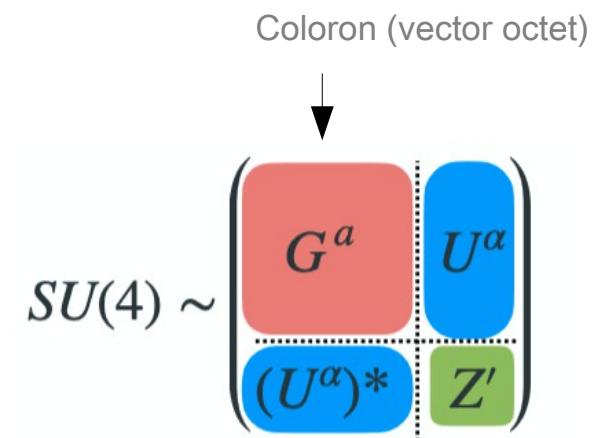
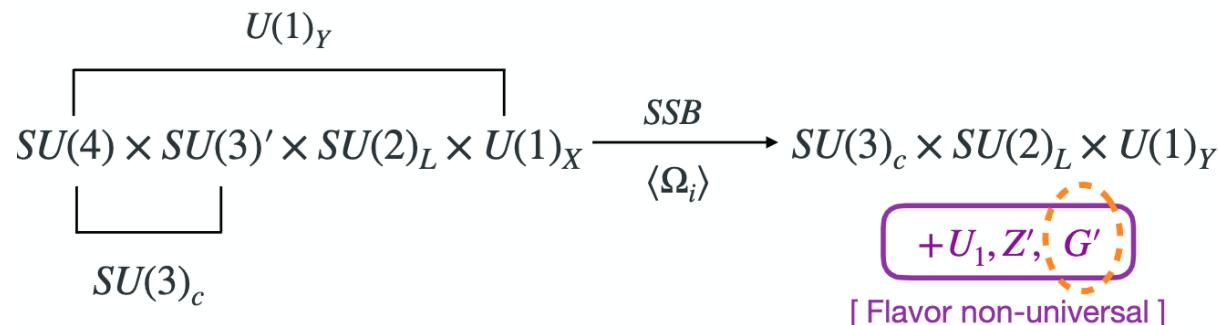
LHC Di-Taus tails push lower bounds on both LFV observables!

Model Potentially within reach at LHCb & Belle 2!

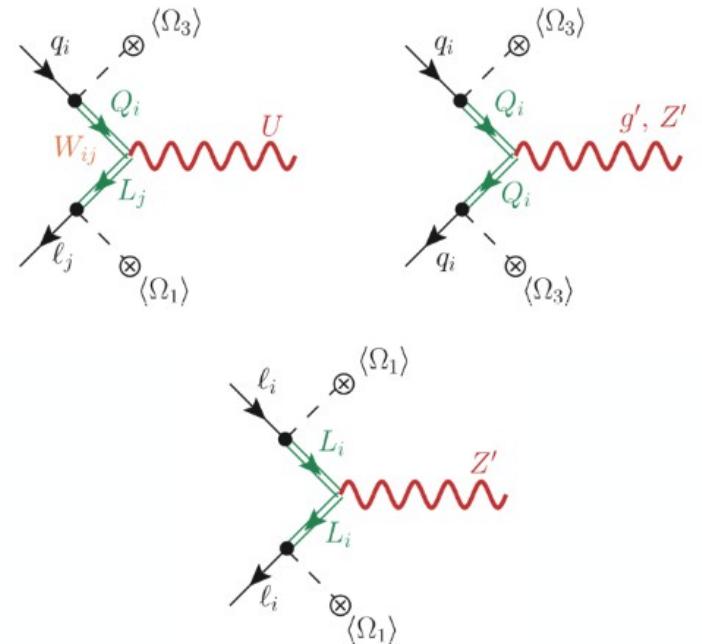
- a UV model -

'4321' model

Pati-Salam Leptoquark $U_1^\mu \sim (\mathbf{3}, \mathbf{1})_{2/3}$

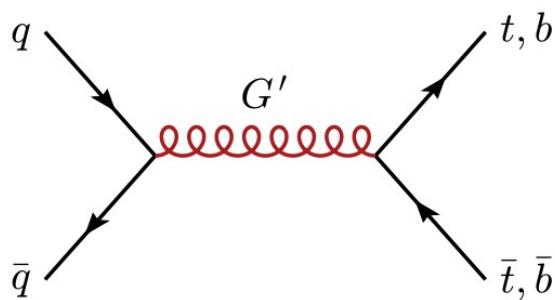


	Field	$SU(4)$	$SU(3)'$	$SU(2)_L$	$U(1)_X$
SM fields	q_L^i	1	3	2	$1/6$
	u_R^i	1	3	1	$2/3$
	d_R^i	1	3	1	$-1/3$
	ℓ_L^i	1	1	2	$-1/2$
	e_R^i	1	1	1	$-1/2$
	$\chi_{L,R}^i$	4	1	2	0
4321 SSB scalars	H	1	1	2	$1/2$
	Ω_1	4	1	1	$-1/2$
	Ω_3	4	3	1	$1/6$
	Ω_{15}	15	1	1	0



Coloron search

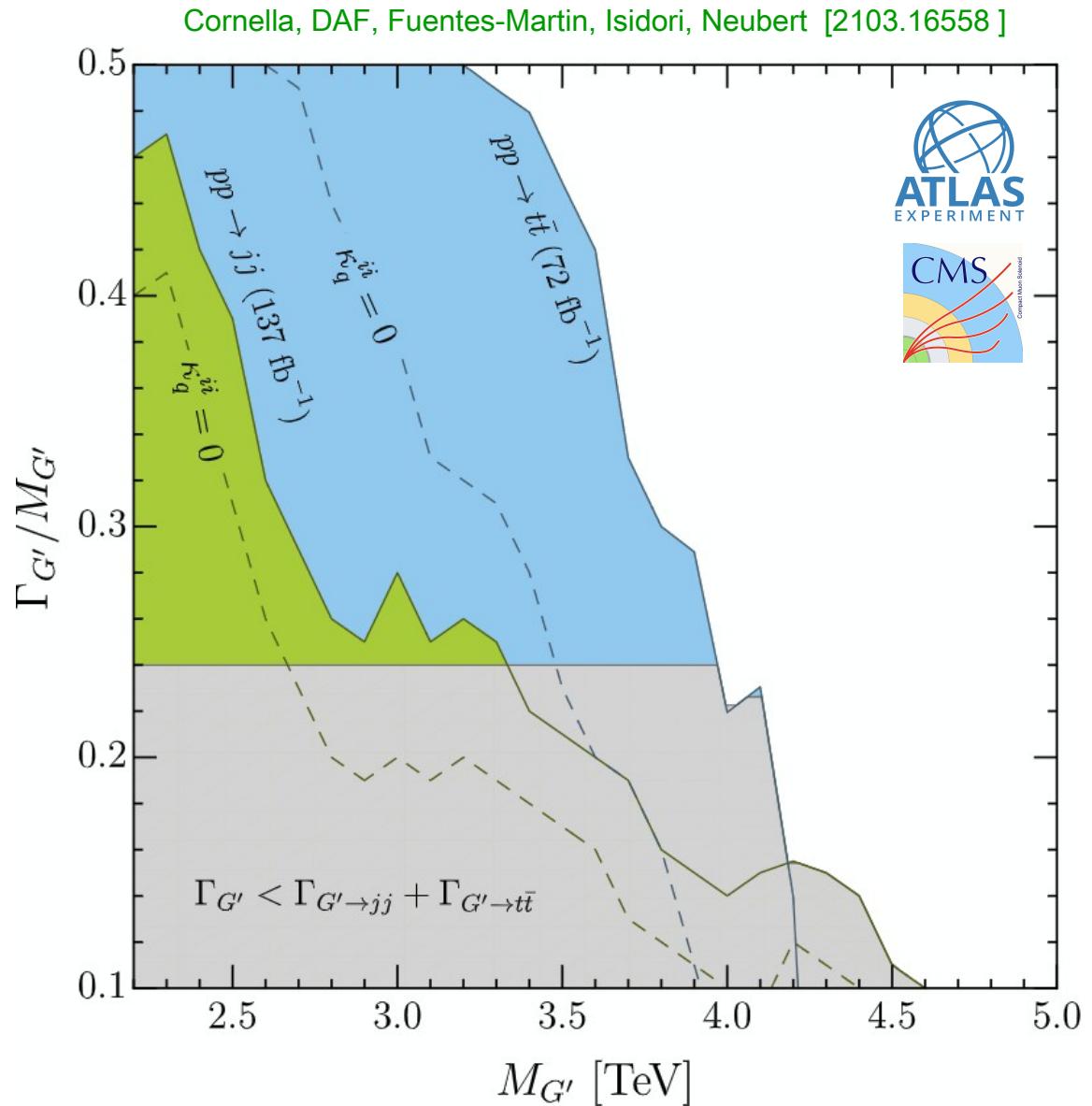
$$pp \rightarrow G' \rightarrow t\bar{t}, b\bar{b}$$



- 4321 models have broad resonances
- Recast of ditop/dijet searches

High mass normalized
invariant mass distribution

Best limits on the 4321 model!



Conclusions

- B-anomalies are alive and kicking: LHCb recently announced evidence of LFUV in R_K
- Combined anomalies in neutral currents are approaching the 5 sigma level.
- B-anomalies opened the doors to explore flavor physics at high-energy pp-colliders.

$$q_i \bar{q}_j \rightarrow \ell_\alpha^\pm \ell_\beta^\mp \quad q_i \bar{q}_j \rightarrow \ell_\alpha^\pm \nu_\beta$$

- We described a remarkable complementarity between high-pT and low energy LFV flavor observables for effective operators.
- High-pT observables are a crucial input for model building efforts \implies Leptoquarks
- Leptoquark solutions exhibit specific **complementary** between dijets and LFV searches:

$$pp \rightarrow \tau^+ \tau^- + X \quad \longleftrightarrow \quad B \rightarrow K \mu \tau \quad \longleftrightarrow \quad \tau \rightarrow \mu \phi$$

ATLAS/CMS

LHCb

Belle 2

- UV and simplified models can be efficiently tested at colliders!

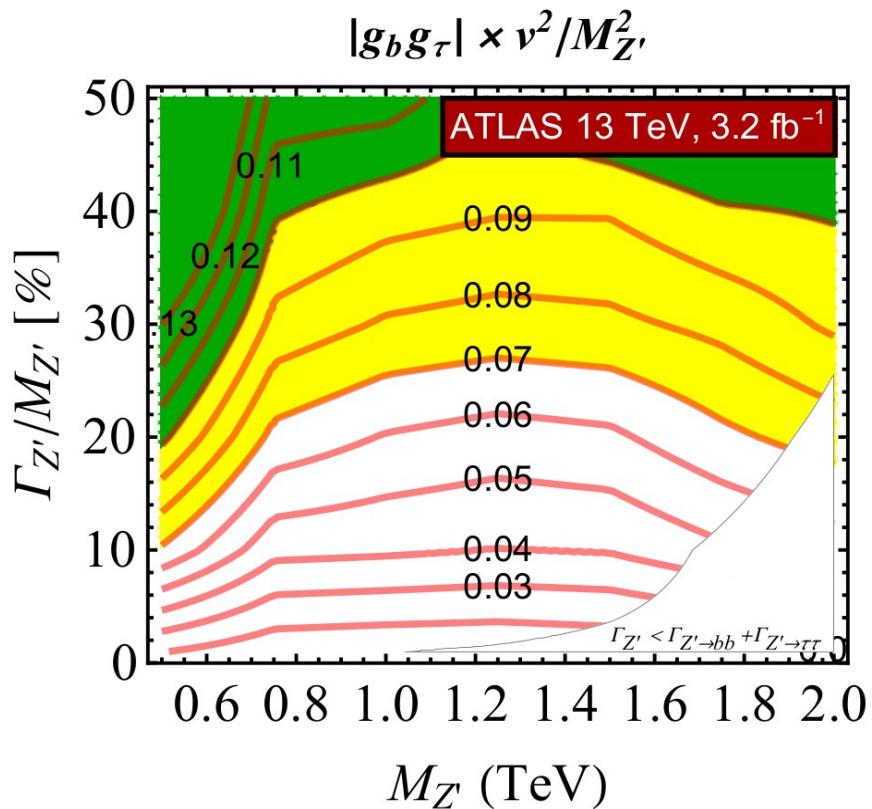
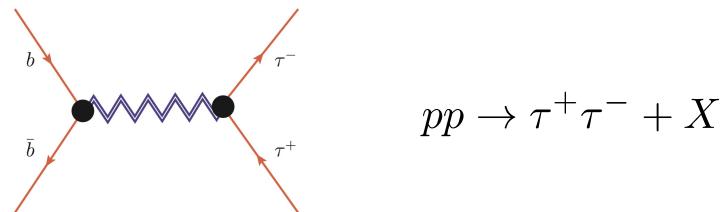
- extra slides -

Di-tau tails & colorless mediators

Vector triplet model:

$$\mathcal{L}_{W'} = -\frac{1}{4} W'_{\mu\nu}^a W'^{a\mu\nu} + \frac{M_{W'}^2}{2} W'^{a\mu} W'^a_\mu + W'^a_\mu J^{a\mu}_{W'}$$

$$J^{a\mu}_{W'} \equiv g_b \bar{Q}_3 \gamma_\mu \tau^a Q_3 + g_\tau \bar{L}_3 \gamma_\mu \tau^a L_3$$



Recast Di-tau search for Z' with arbitrary width
JHEP 1801 055 (2018)



High-pT di-tau searches excludes **colorless** solution to B-anomalies!

See DAF, Greljo, Kamenik Phys. Lett. B 764(2017)126-134

Di-lepton tails - RK(*)

EUROPEAN ORGANISATION FOR NUCLEAR RESEARCH (CERN)



Submitted to: Phys. Rev. Lett.



CERN-EP-2021-065
May 31, 2021

Search for new phenomena in final states with two leptons and one or no b -tagged jets at $\sqrt{s} = 13$ TeV using the ATLAS detector

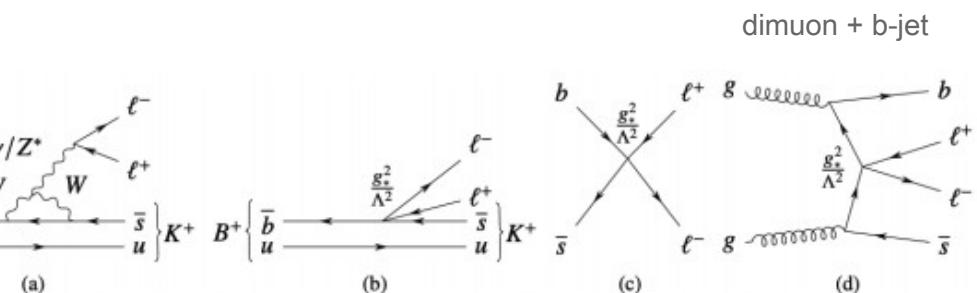
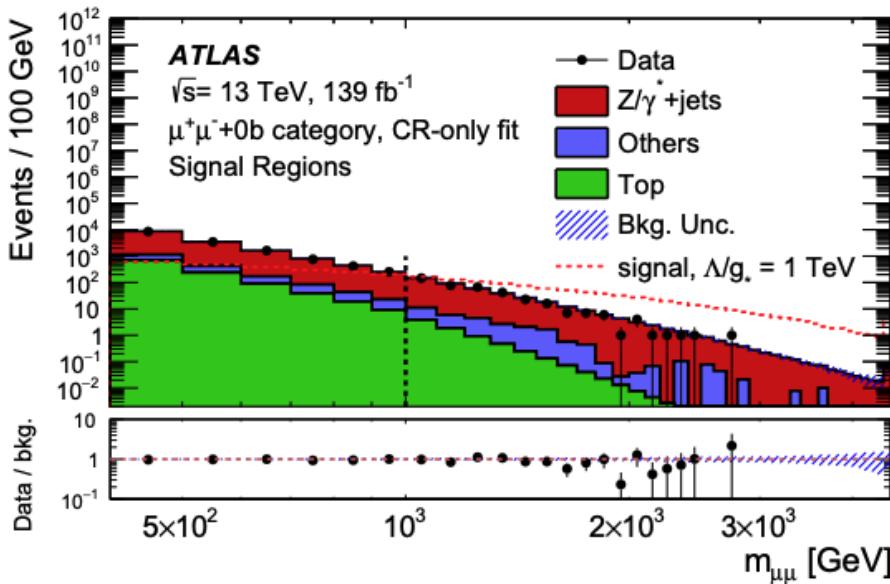
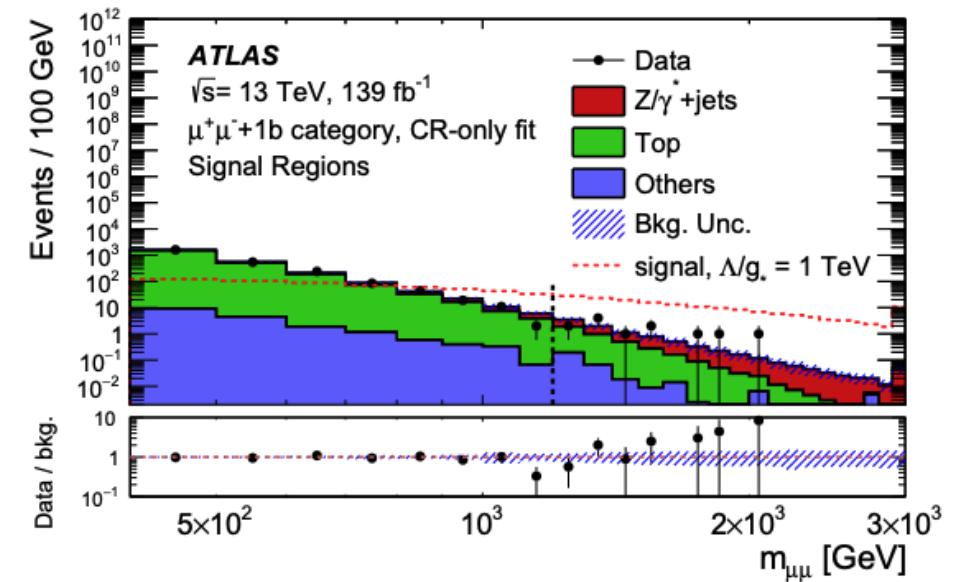


Figure 1: Representative Feynman diagrams for the decay of a B^+ meson to a K^+ meson in association with two leptons (a) in the SM and (b) in the EFT approach, and for production of two leptons via a $bs\ell\ell$ contact interaction in $p\bar{p}$ collisions (c) without and (d) with a b -jet in the final state.



- First limits on contact interaction $bs\mu\mu$

$$\Lambda/g^* > 2 \text{ TeV}$$

Di-lepton SMEFT

Greljo, Marzocca '17

$$\mathcal{L}^{\text{eff}} \supset \frac{\mathbf{C}_{ij}^{U\mu}}{v^2} (\bar{u}_L^i \gamma_\mu u_L^j) (\bar{\mu}_L \gamma^\mu \mu_L) + \frac{\mathbf{C}_{ij}^{D\mu}}{v^2} (\bar{d}_L^i \gamma_\mu d_L^j) (\bar{\mu}_L \gamma^\mu \mu_L)$$

$$\mathbf{C}_{ij}^{D\mu} = \begin{pmatrix} C_{d\mu} & 0 & 0 \\ 0 & C_{s\mu} & C_{bs\mu}^* \\ 0 & C_{bs\mu} & C_{b\mu} \end{pmatrix}$$

Minimal Flavor Violation (MFV) for $R_{K^{(*)}}$

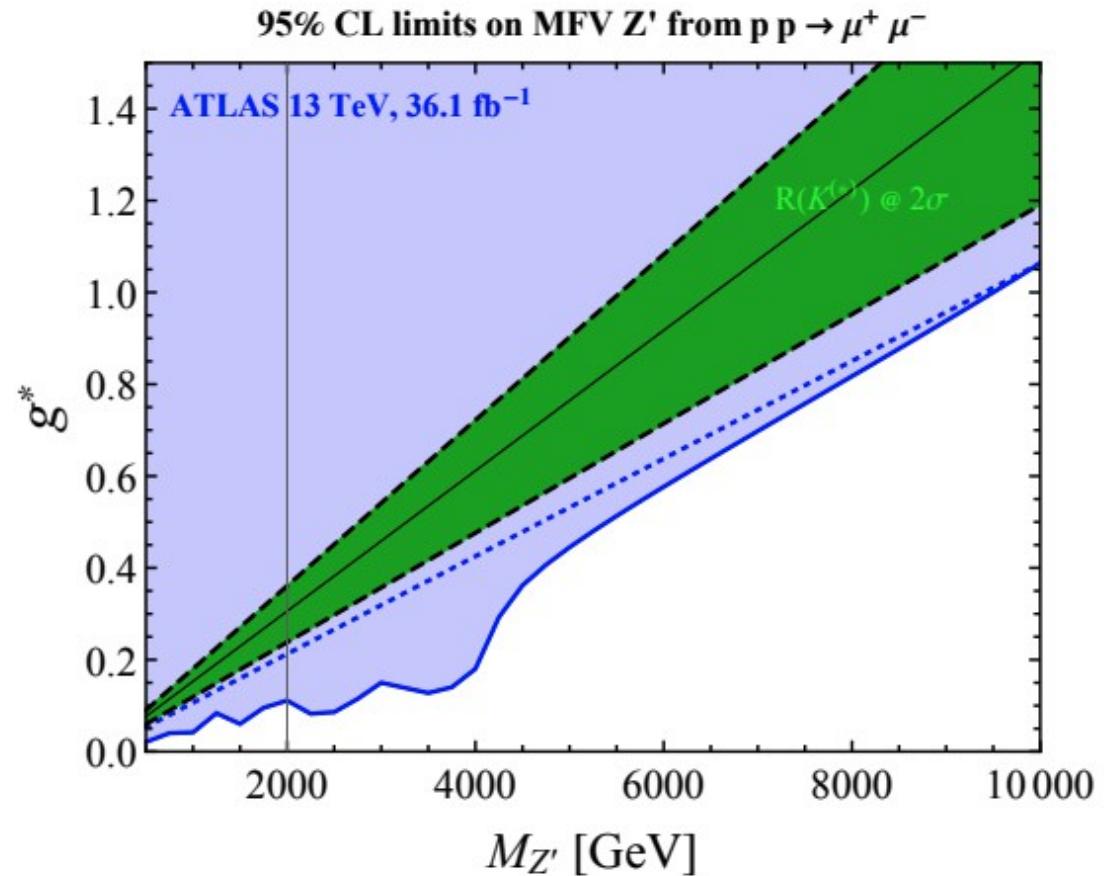
$$C_{d\mu} = C_{s\mu} = C_{b\mu} \equiv C_{D\mu}$$

$$|C_{bs\mu}| \sim |V_{tb} V_{ts}^* y_t^2 C_{D\mu}|$$



$$\frac{C}{v^2} = \frac{g_*^2}{M_{Z'}^2}$$

Scenario excluded by dileptons!



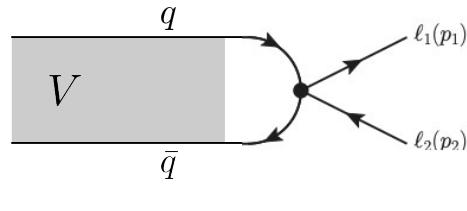
Di-lepton tails - LFV

Four relevant LFV low energy experiments:

- 1) $\mu \rightarrow e$ conversion rate in Nuclei [Only relevant for valence quarks...]

- 2) Quarkonia LFV decays:

$$V \rightarrow \ell_\alpha \ell_\beta$$

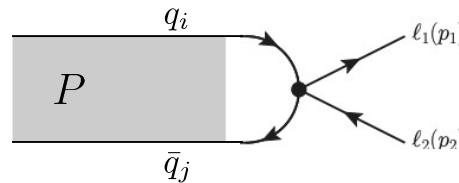


$$\begin{aligned} \mathcal{B}(V \rightarrow \ell_\alpha^- \ell_\beta^+) = & \frac{\tau_V m_V^3 f_V^2}{24\pi v^4} \left(1 - \frac{m_{\ell_\beta}^2}{m_V^2}\right)^2 \times \left\{ \left[|\mathcal{C}_{VL}|^2 + |\mathcal{C}_{VR}|^2 \right] \left(1 + \frac{m_{\ell_\beta}^2}{2m_V^2}\right) \right. \\ & + 6 \frac{f_V^T}{f_V} \frac{m_{\ell_\beta}}{m_V} \text{Re}(\mathcal{C}_T \mathcal{C}_{VR}^* + \hat{\mathcal{C}}_T \mathcal{C}_{VL}^*) \\ & \left. + 2 \left(\frac{f_V^T}{f_V} \right)^2 \left[|\mathcal{C}_T|^2 + |\hat{\mathcal{C}}_T|^2 \right] \left(1 + \frac{2m_{\ell_\beta}^2}{m_V^2}\right) \right\} \end{aligned}$$

- 3) FCNC Pseudo-scalar meson decays:

$$\begin{aligned} P \rightarrow \ell_\alpha \ell_\beta \\ P \rightarrow M \ell_\alpha \ell_\beta \end{aligned}$$

$$P = K_{L,S}, D^0, B_s$$

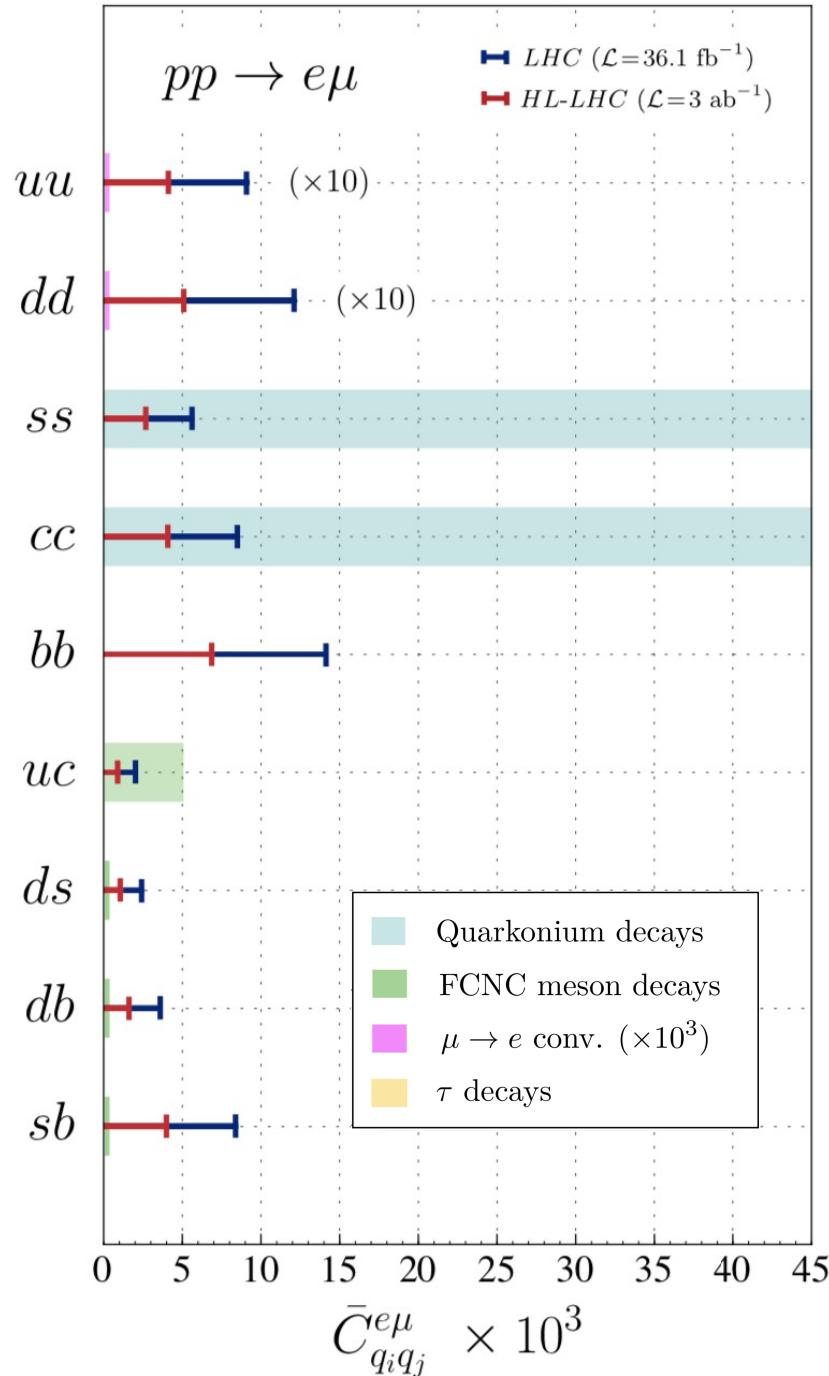


[Helicity suppression]

- 4) LFV tau decays:

$$\tau \rightarrow P \ell_\alpha, \quad \tau \rightarrow V \ell_\alpha$$

LHC and flavor limits (@95% CL)



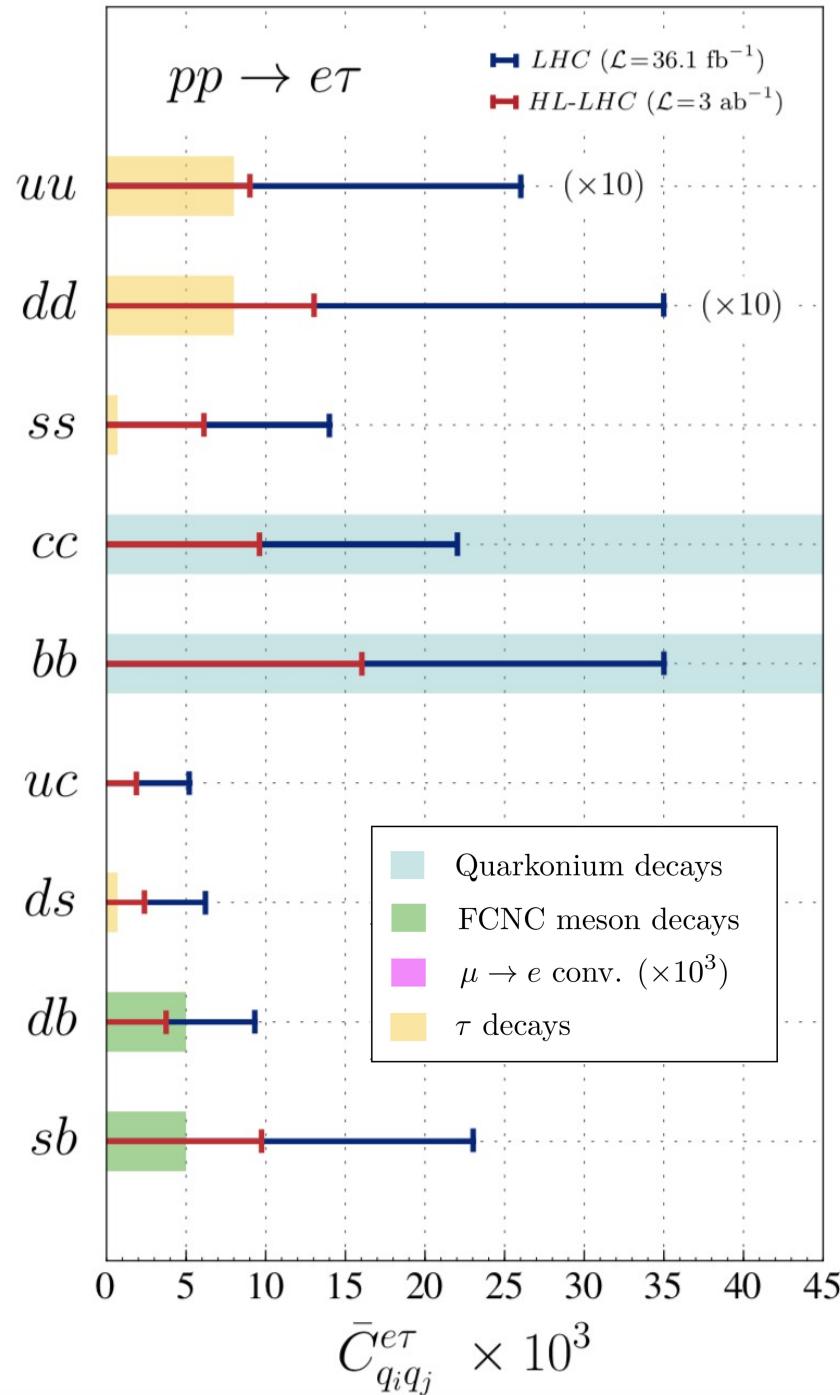
LHC vs Flavor $e\mu$

Decay mode	Exp. limit
$K_L \rightarrow \mu^\mp e^\pm$	6.1×10^{-12}
$K^+ \rightarrow \pi^+ \mu^+ e^-$	1.7×10^{-11}
$\phi \rightarrow \mu^\pm e^\mp$	2.6×10^{-6}
$D \rightarrow \mu^\pm e^\mp$	1.6×10^{-8}
$J/\psi \rightarrow \mu^\pm e^\mp$	2.1×10^{-7}
$B_d \rightarrow \mu^\mp e^\pm$	1.3×10^{-9}
$B^+ \rightarrow \pi^+ \mu^\mp e^\pm$	2.2×10^{-7}
$B_s \rightarrow \mu^\mp e^\pm$	6.3×10^{-9}
$B^+ \rightarrow K^+ \mu^+ e^-$	8.8×10^{-9}
$B^0 \rightarrow K^* \mu^\mp e^\pm$	2.3×10^{-7}

Selected LHC limits (left-handed scenario)		
Decay mode	Current (36 fb^{-1})	Future (3 ab^{-1})
$\phi \rightarrow \mu^\pm e^\mp$	6.2×10^{-18}	1.4×10^{-18}
$D^0 \rightarrow \mu^\pm e^\mp$	2.8×10^{-9}	5.0×10^{-10}
$J/\psi \rightarrow \mu^\pm e^\mp$	7.1×10^{-12}	1.6×10^{-12}
$B_d \rightarrow \mu^\mp e^\pm$	7.8×10^{-8}	1.4×10^{-9}
$B^+ \rightarrow \pi^+ \mu^\mp e^\pm$	3.5×10^{-5}	6.4×10^{-6}
$B_s \rightarrow \mu^\mp e^\pm$	6.4×10^{-7}	1.4×10^{-7}
$B^+ \rightarrow K^+ \mu^\mp e^\pm$	2.8×10^{-4}	6.2×10^{-5}
$B^0 \rightarrow K^* \mu^\mp e^\pm$	5.8×10^{-4}	1.3×10^{-4}
$\Upsilon(3S) \rightarrow \mu^\pm e^\mp$	5.6×10^{-9}	1.3×10^{-9}

LHC vs Flavor $e\tau$

LHC and flavor limits (@95% CL)



Decay mode	Exp. limit
$\tau \rightarrow e\rho$	2.3×10^{-8}
$\tau \rightarrow eK^*$	4.2×10^{-8}
$\tau \rightarrow e\phi$	4.0×10^{-8}
$J/\psi \rightarrow \tau^\pm e^\mp$	1.1×10^{-5}
$B_d \rightarrow \tau^\pm e^\mp$	3.6×10^{-5}
$B^+ \rightarrow \pi^+ \tau^\pm e^\mp$	9.7×10^{-5}
$B^+ \rightarrow K^+ \tau^\pm e^\mp$	3.9×10^{-5}
$\Upsilon(3S) \rightarrow \tau^\pm e^\mp$	5.4×10^{-6}

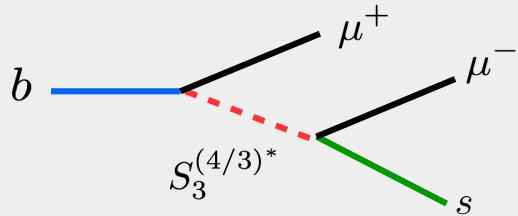
Selected LHC limits (left-handed scenario)		
Decay mode	Current (36 fb^{-1})	Future (3 ab^{-1})
$D^0 \rightarrow \tau^\pm e^\mp$	5.1×10^{-8}	6.8×10^{-9}
$J/\psi \rightarrow \tau^\pm e^\mp$	2.6×10^{-11}	4.9×10^{-12}
$B_d \rightarrow \tau^\pm e^\mp$	1.2×10^{-4}	1.9×10^{-5}
$B^+ \rightarrow \pi^+ \tau^\pm e^\mp$	1.6×10^{-4}	2.5×10^{-5}
$B_s \rightarrow \tau^\mp e^\pm$	1.1×10^{-3}	2.0×10^{-4}
$B^+ \rightarrow K^+ \tau^\pm e^\mp$	1.4×10^{-3}	2.5×10^{-4}
$B^0 \rightarrow K^* \tau^\pm e^\mp$	2.4×10^{-3}	4.2×10^{-4}
$\Upsilon(3S) \rightarrow \tau^\pm e^\mp$	3.4×10^{-8}	7.0×10^{-9}

S_3 Scalar Triplet

$$S_3 \sim (\bar{\mathbf{3}}, \mathbf{3})_{1/3}$$

$$\begin{aligned} \mathcal{L}_{S_3} = & -y_L^{ij} d_L^{\bar{c}}{}_i \nu_{Lj} S_3^{(1/3)} - \boxed{\sqrt{2} y_L^{ij} d_L^{\bar{c}}{}_i \ell_{Lj} S_3^{(4/3)}} \\ & + \sqrt{2} (V^* y_L)_{ij} u_L^{\bar{C}}{}_i \nu_{Lj} S_3^{(-2/3)} - (V^* y_L)_{ij} u_L^{\bar{C}}{}_i \ell_{Lj} S_3^{(1/3)} + \text{h.c.} \end{aligned}$$

$$b \rightarrow s\mu\bar{\mu}$$



Minimal texture

$$y_L = \begin{pmatrix} 0 & 0 & 0 \\ 0 & y_L^{s\mu} & 0 \\ 0 & y_L^{b\mu} & 0 \end{pmatrix}$$

$$\mathcal{C}_9 = -\mathcal{C}_{10} = \frac{\pi v^2}{V_{ts}^* V_{tb}} \frac{y_L^{b\mu} (y_L^{s\mu})^*}{m_S^2}$$

$$R_{K^{(*)}} < R_{K^{(*)}}^{\text{SM}}$$

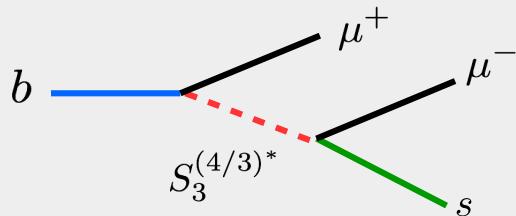
explains anomaly

S_3 Scalar Triplet

$$S_3 \sim (\bar{\mathbf{3}}, \mathbf{3})_{1/3}$$

$$\mathcal{L}_{S_3} = -y_L^{ij} d_L^{\bar{c}}{}_i \nu_{Lj} S_3^{(1/3)} - \sqrt{2} y_L^{ij} d_L^{\bar{c}}{}_i \ell_{Lj} S_3^{(4/3)} + \sqrt{2} (V^* y_L)_{ij} u_L^{\bar{C}}{}_i \nu_{Lj} S_3^{(-2/3)} - (V^* y_L)_{ij} u_L^{\bar{C}}{}_i \ell_{Lj} S_3^{(1/3)} + \text{h.c.}$$

$$b \rightarrow s\mu\bar{\mu}$$



Minimal texture

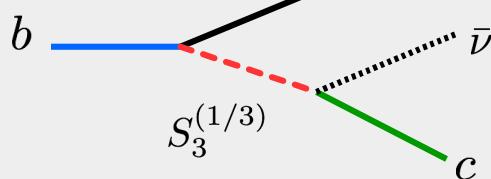
$$y_L = \begin{pmatrix} 0 & 0 & 0 \\ 0 & y_L^{s\mu} & 0 \\ 0 & y_L^{b\mu} & 0 \end{pmatrix}$$

$$\mathcal{C}_9 = -\mathcal{C}_{10} = \frac{\pi v^2}{V_{ts}^* V_{tb}} \frac{y_L^{b\mu} (y_L^{s\mu})^*}{m_S^2}$$

$$R_{K^{(*)}} < R_{K^{(*)}}^{\text{SM}}$$

explains anomaly

$$b \rightarrow c\tau\bar{\nu}$$



$$y_L = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & y_L^{s\tau} \\ 0 & 0 & y_L^{b\tau} \end{pmatrix}$$

$$g_{V_L} = -\frac{v^2 y_L^{b\tau} (V y_L^*)_{c\nu}}{4V_{cb} m_{S_3}^2} = -\frac{v^2}{4m_{S_3}^2} \left[|y_L^{b\tau}|^2 + \frac{V_{cs}}{V_{cb}} y_L^{b\tau} (y_L^{s\tau})^* \right].$$

V-A structure but

$$g_{V_L} < 0$$

$$R_{D^{(*)}} < R_{D^{(*)}}^{\text{SM}}$$

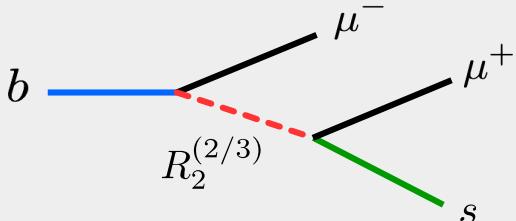
wrong direction

B-anomalies: R_2 scalar leptoquark

$$R_2 \sim (\mathbf{3}, \mathbf{2})_{7/6}$$

$$\begin{aligned} \mathcal{L}_{R_2} = & y_R^{ij} (\bar{u}_{Li} \ell_{Rj}) R_2^{(5/3)} + \boxed{y_R^{ij} (\bar{d}_{Li} \ell_{Rj}) R_2^{(2/3)}} \\ & + y_L^{ij} (\bar{u}_{Ri} \nu_{Lj}) R_2^{(2/3)} - y_L^{ij} (\bar{u}_{Ri} \ell_{Lj}) R_2^{(5/3)} \end{aligned}$$

$b \rightarrow s\mu\bar{\mu}$ tree-level solution



$$y_R = \begin{pmatrix} 0 & 0 & 0 \\ 0 & x_R^{s\mu} & 0 \\ 0 & x_R^{b\mu} & 0 \end{pmatrix} \quad \mathcal{C}_9 = \mathcal{C}_{10} = -\frac{\pi v^2}{2V_{tb}V_{ts}^*\alpha} \frac{y_R^{s\mu}(y_R^{b\mu})^*}{m_R^2}$$

$$R_{K^{(*)}} > R_{K^{(*)}}^{\text{SM}}$$

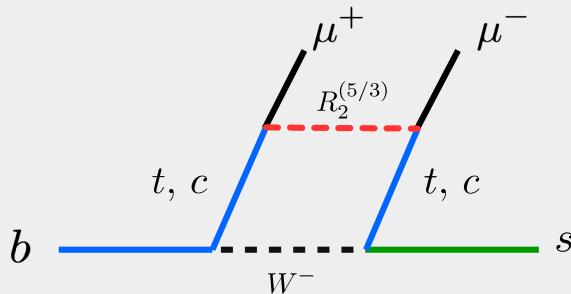
wrong
direction

R_2 Scalar LQ

$$R_2 \sim (\mathbf{3}, \mathbf{2})_{7/6}$$

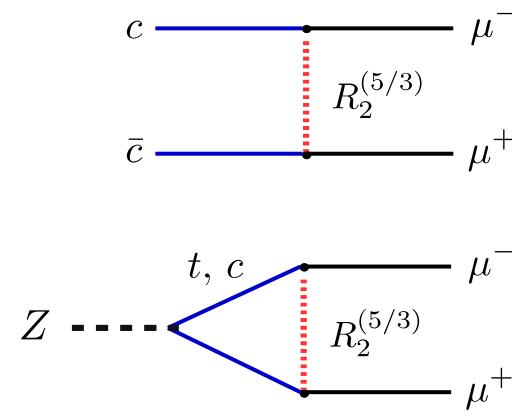
$$\begin{aligned} \mathcal{L}_{R_2} = & \boxed{y_R^{ij}(\bar{u}_{Li}\ell_{Rj})R_2^{(5/3)}} + y_R^{ij}(d_{Li}\ell_{Rj})R_2^{(2/3)} \\ & + y_L^{ij}(\bar{u}_{Ri}\nu_{Lj})R_2^{(2/3)} - y_L^{ij}(u_{Ri}\ell_{Lj})R_2^{(5/3)} \end{aligned}$$

$b \rightarrow s\mu\bar{\mu}$ loop-level solution



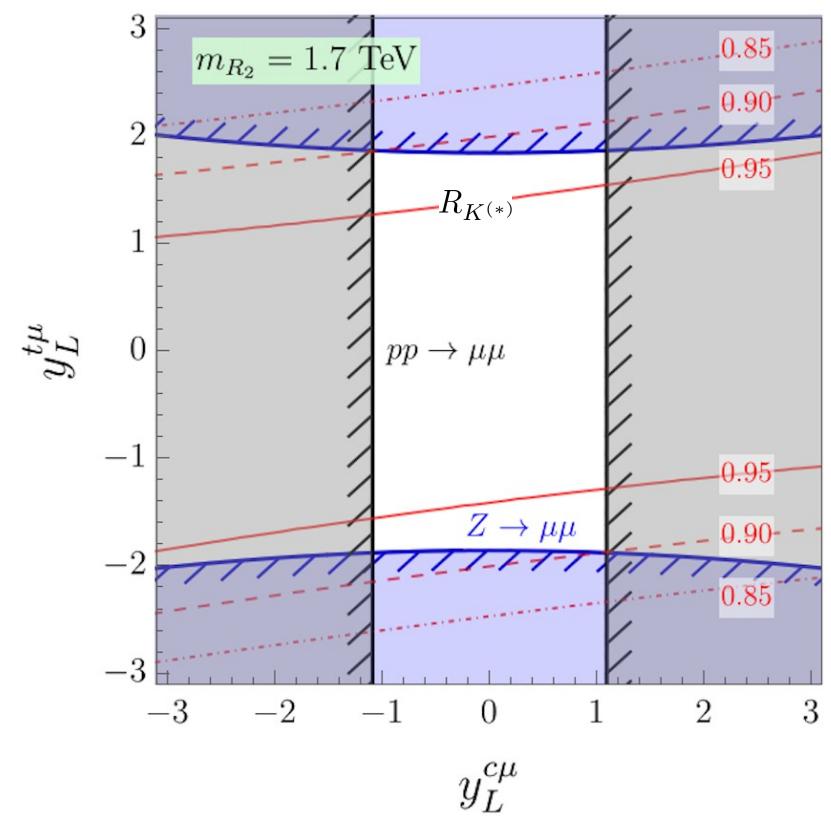
[Becirevic, Sumensari 2016]

Unfavoured by
Drell-Yan + Z-pole



$$y_L = \begin{pmatrix} 0 & 0 & 0 \\ 0 & x_R^{c\mu} & 0 \\ 0 & x_R^{t\mu} & 0 \end{pmatrix}$$

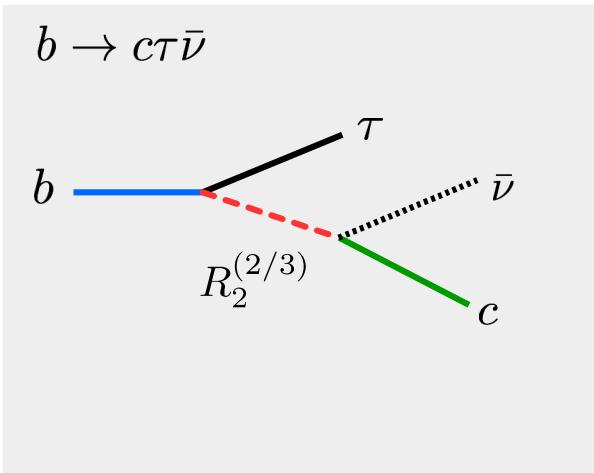
$$\boxed{R_{K^{(*)}} < R_{K^{(*)}}^{\text{SM}}}$$



R₂ Scalar LQ

$$R_2 \sim (\mathbf{3}, \mathbf{2})_{7/6}$$

$$\begin{aligned} \mathcal{L}_{R_2} = & y_R^{ij} (\bar{u}_{Li} \ell_{Rj}) R_2^{(5/3)} + \boxed{y_R^{ij} (d_{Li} \ell_{Rj}) R_2^{(2/3)}} \\ & + \boxed{y_L^{ij} (\bar{u}_{Ri} \nu_{Lj}) R_2^{(2/3)}} - y_L^{ij} (u_{Ri} \ell_{Lj}) R_2^{(5/3)} \end{aligned}$$



$$\begin{aligned} y_L &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & y_L^{c\tau} \\ 0 & 0 & 0 \end{pmatrix} & g_{S_L} = 4g_T = \frac{v^2}{4V_{cb}} \frac{y_L^{c\tau} (y_R^{b\tau})^*}{m_R^2} \\ y_R &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & y_R^{b\tau} \end{pmatrix} & \boxed{R_{D^{(*)}} > R_{D^{(*)}}^{\text{SM}}} \end{aligned}$$

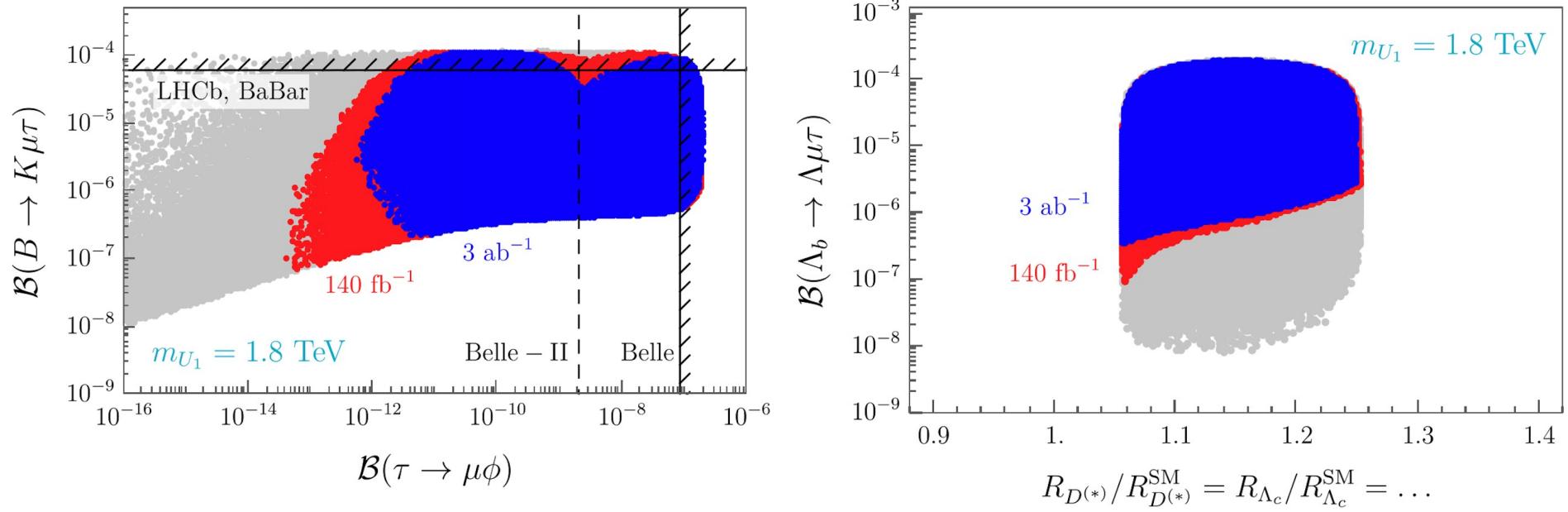
$$y_L^{c\tau} \in \mathbb{R} \quad y_R^{b\tau} \in \mathbb{C}$$

Scalar + Tensor solution!

Becirevic, Dorsner, Fajfer, DAF,
Kosnik, Sumensari [1808.08179]

Non-V-A solution to anomaly!

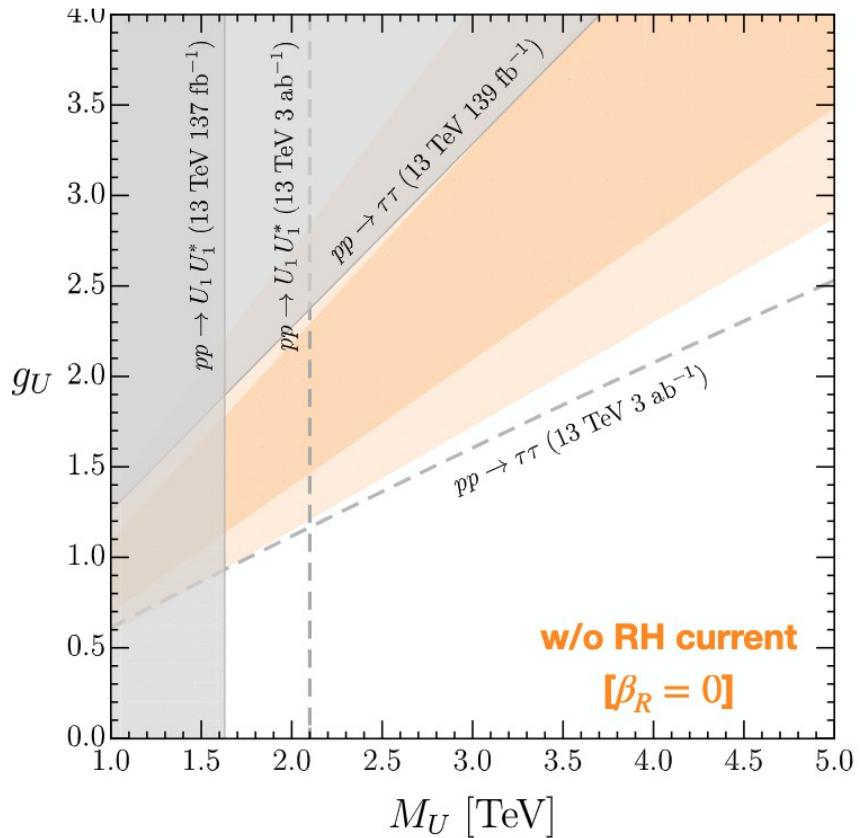
- Predictions for LFV low-energy observables for U_1



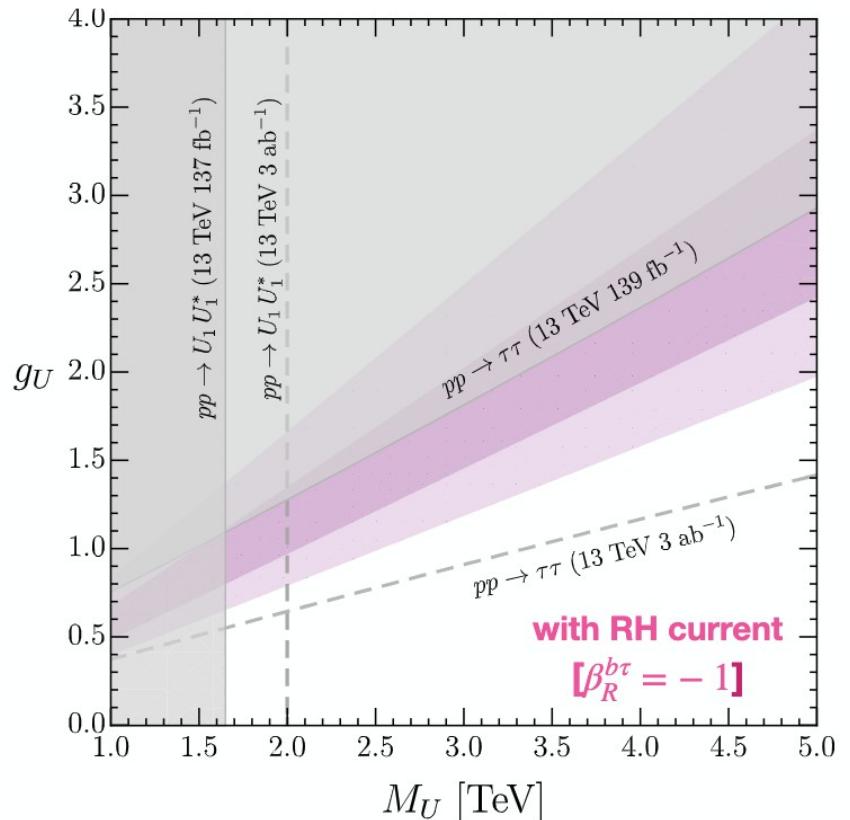
Angelescu, Becirevic, DAF, Jaffredo, Sumensari, [2021]

LHC Di-Taus tails push lower bounds on both LFV observables!

Model Potentially within reach at LHCb & Belle 2!



[Cornella et al., 2103.16558]



[$pp \rightarrow \tau\tau$ for U_1 originally proposed in
Faroughy, Greljo, Kamenik, 1609.07138]

GUT-inspired leptoquarks

GUT Leptoquarks: $R_2 \sim (3,2)_{7/6}$ and $S_3 \sim (3,3)_{-1/3}$

Becirevic, Dorsner, Fajfer,
DAF, Kosnik, Sumensari
[1808.08179]

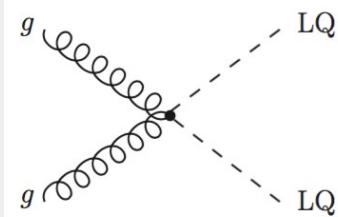
\downarrow \downarrow

for $R_{D^{(*)}}$ for $R_{K^{(*)}}$

Good Benchmark:

R_2 with $\sim O(1)$ Yukawa couplings
Couples to both **charm** and **bottom**
 $m_R \sim 800$ GeV

LQs decaying into tauonic channels

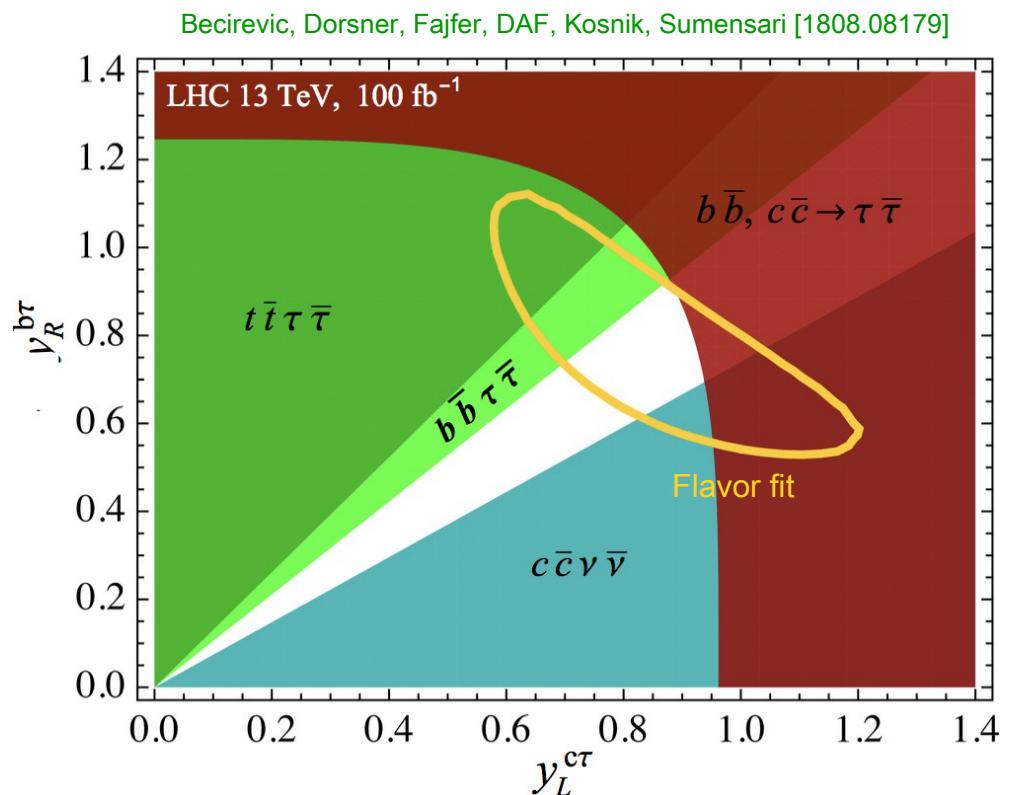


$$R_2^{\frac{2}{3}} \rightarrow \tau b, \nu c$$
$$R_2^{\frac{5}{3}} \rightarrow \tau t, \tau c$$

LQ pair production & DY di-Tau searches

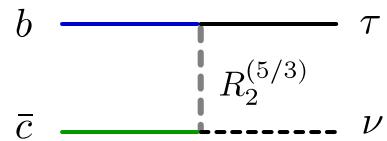
Model can be put to test at the LHC!

Complementarity:
Flavor obs. & **high- p_T di-Taus**

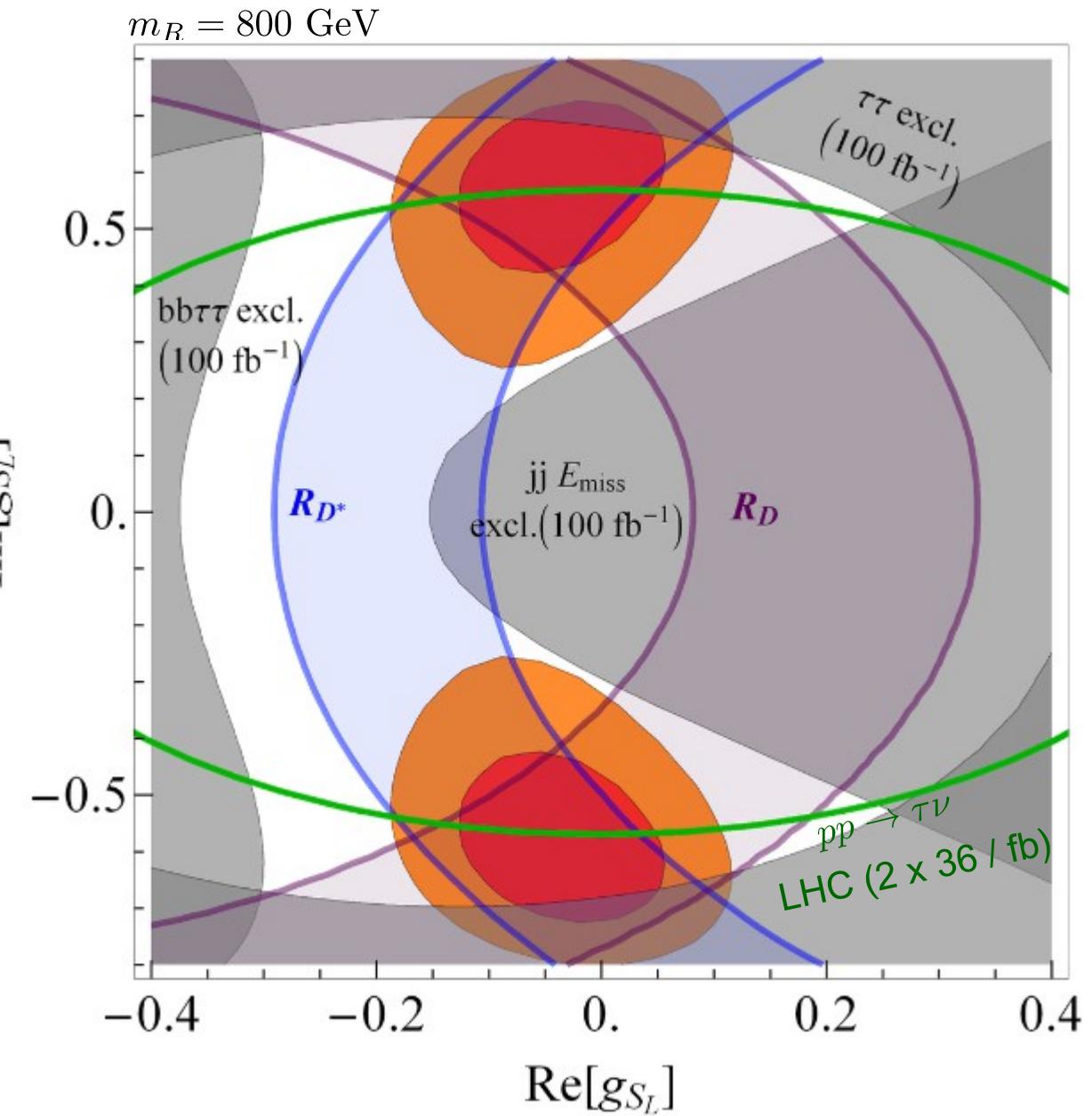


Mono-Tau searches

$$pp \rightarrow \tau\nu + X$$



Greljo,Camalich, Ruiz-Alvarez '18



LFV low-energy observables for other LQ models:

