

# On the dispersive evaluation of the HVP contribution to $a_\mu$ and $\alpha_{\text{QED}}$ , and implications for the EW fit

Bogdan MALAESCU

LPNHE, CNRS



*Based on:* [1908.00921\(DHMZ\)](#), [2006.04822\(WP Theory Initiative\)](#), [2008.08107\(BM,MS\)](#)

*In collaboration & useful discussions with:*

Michel DAVIER, Andreas HOECKER, Zhiqing ZHANG (DHMZ)

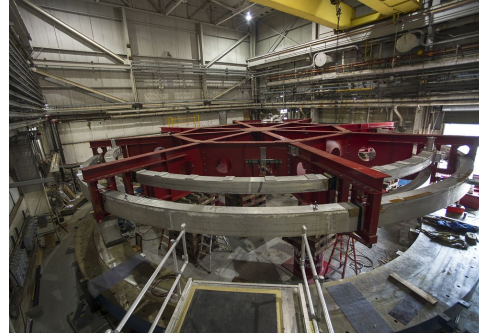
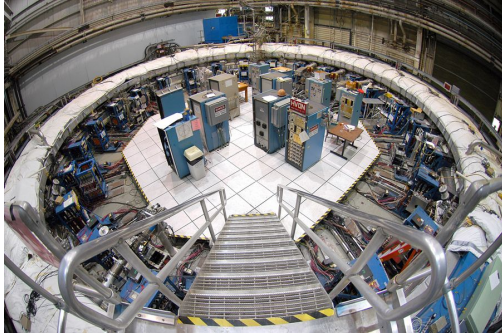
Laurent LELLOUCH, Matthias SCHOTT; the Gfitter group

*LIO international conference on Future colliders and the origin of mass*

25/06/2021

# How it all started (again)...

BNL → 1 month long trip for the g-2 storage ring



*This is NOT an UFO !!! ;-)*

→ Fermilab  
July 26, 2013



# Content of the talk

- Introduction: the  $(g-2)_\mu$  experiment & theoretical prediction
- Data on  $e^+e^- \rightarrow$  hadrons
- Combination of all  $e^+e^-$  data:  
focus on the combination procedure  
(HVPTools and fit based on analyticity & unitarity)
- Indications of uncertainties on uncertainties and on correlations & their implications for combinations
- Results on  $a_\mu$
- Impact of correlations between  $a_\mu$  and  $\alpha_{\text{QED}}$  on the EW fit
- Conclusions

# The $(g-2)_\mu$ : definition & experimental measurement

- Magnetic dipole moment of a charged lepton:  $\vec{\mu} = g \frac{e}{2m} \vec{s}$
- “anomaly” = deviation w.r.t. Dirac’s prediction:  $a = \frac{g-2}{2}$

- Experimental “ingredients” to measure  $a_\mu$ :

→ Polarised muons from pion decays (parity violation)

→ “Anomalous frequency”

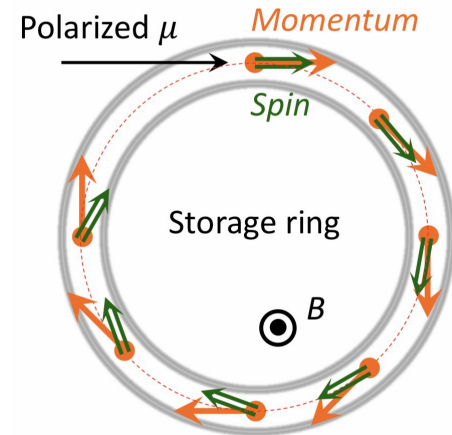
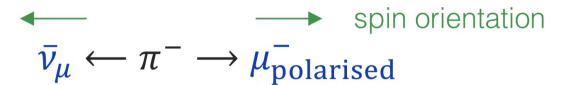
(difference between spin precession and cyclotron frequency)

proportional to  $a_\mu$  for the “magic  $\gamma$ ”

$$\vec{\omega}_a = \frac{e}{m_\mu c} \left[ a_\mu \vec{B} - \left( a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} \right] \approx \frac{e}{m_\mu c} a_\mu \vec{B}$$

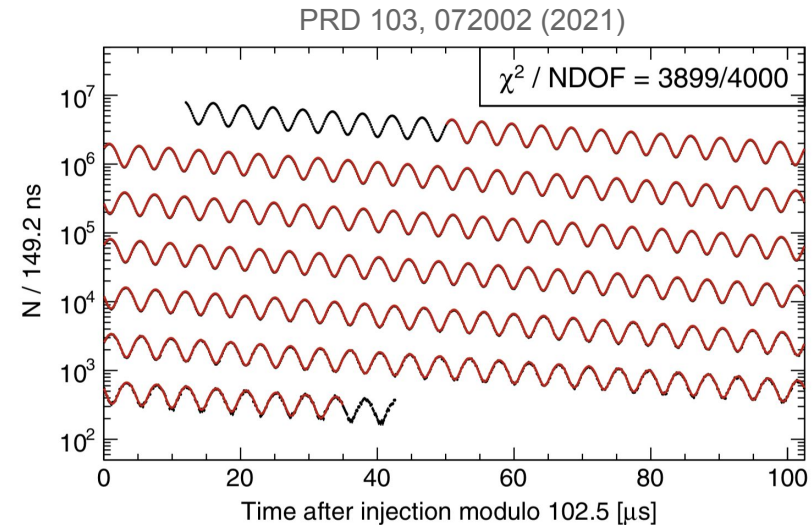
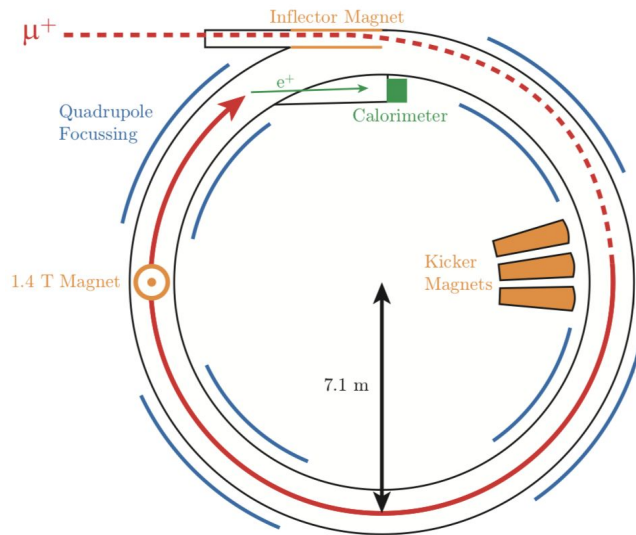
→ Parity violation in muon decays

(electron emitted in the direction opposite to the muon spin)



$$\mu^-_{\text{polarised}} \rightarrow e^- + \bar{\nu}_e + \nu_\mu$$

# The $(g-2)_\mu$ experiment



$$a_\mu^{\text{Exp}}(\text{BNL}): (11\,659\,208.9 \pm 6.3) \cdot 10^{-10}$$

→ Expected uncertainty reduction by a factor 4 with the experiment at Fermilab

- improved apparatus and enhanced statistics: more intense (x20) and pure muon beam; B-field mapped every 3 days with special trolley with probes pulled through beampipe (homogeneity ~ ppm); tracking system for electron detectors etc.

- first publication: similar precision & good agreement with BNL (7th of April 2021) PRL 126, 141801 (2021)

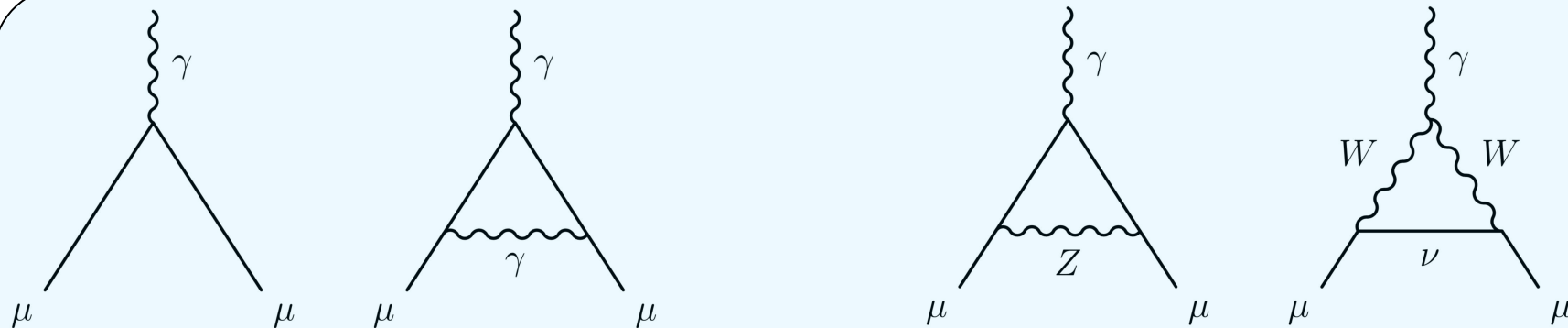
$$a_\mu^{\text{Exp}}(\text{Fermilab}): (11\,659\,204.0 \pm 5.1 \pm 1.8) \cdot 10^{-10} \rightarrow \text{so far only 6\% of the total data}$$

$$a_\mu^{\text{Exp}}(\text{Fermilab} + \text{BNL}): (11\,659\,206.1 \pm 4.1) \cdot 10^{-10} \text{ (0.35 ppm)} \rightarrow \text{One of the most precise quantities ever measured}$$

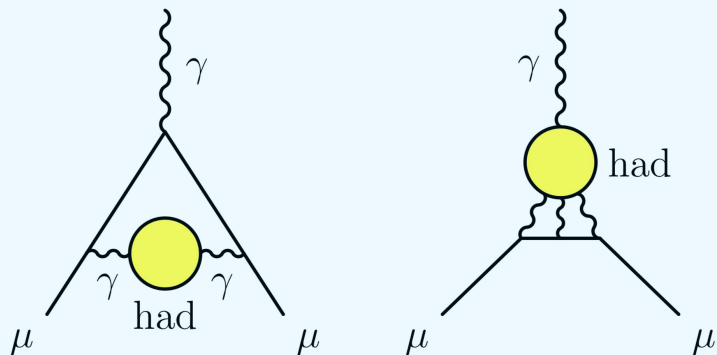
→ Initiative for a measurement using slow muons (KEK, Japan)

# Theoretical prediction

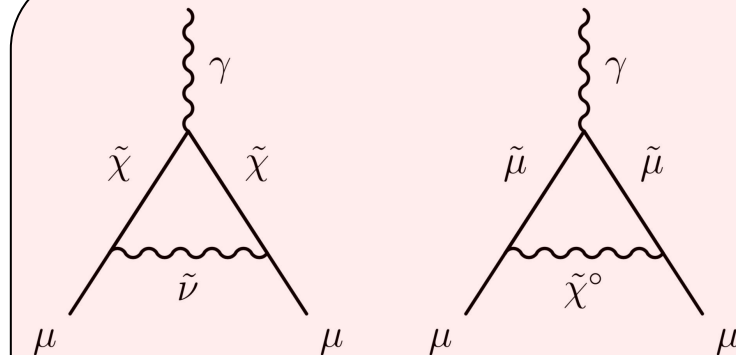
Why is it (so) complicated to compute one number ? (*very precisely*)



+ Many other diagrams at higher orders...



Dominant uncertainties: non-perturbative...



???

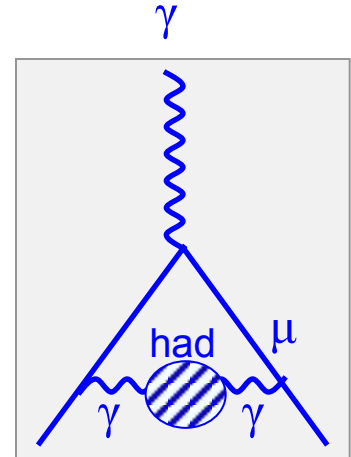
# Hadronic Vacuum Polarization and Muon $(g-2)_\mu$

Dominant uncertainty for the theoretical prediction: from lowest-order HVP piece  
 Cannot be calculated from QCD (low mass scale), but one can use experimental data on  $e^+e^- \rightarrow$  hadrons cross section

Born:  $\sigma^{(0)}(s) = \sigma(s)(\alpha/\alpha(s))^2$

$$12\pi \operatorname{Im}\Pi_\gamma(s) = \frac{\sigma^0 [e^+e^- \rightarrow \text{hadrons} (\gamma_{FSR})]}{\sigma_{pt}} \equiv R(s)$$

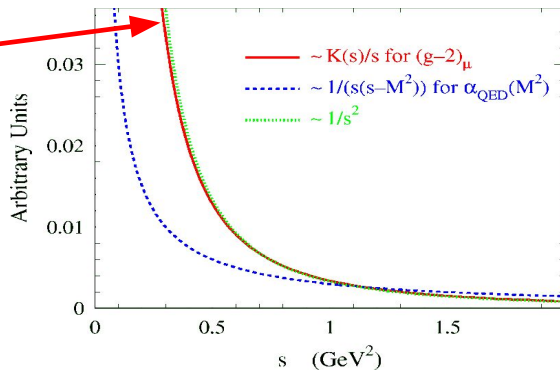
$\operatorname{Im}[\text{Diagram}] \propto |\text{Diagram} \rightarrow \text{hadrons}|^2$



$$a_\mu^{\text{had}} = \frac{\alpha^2}{3\pi^2} \int_{4m_\pi^2}^{\infty} ds \frac{K(s)}{s} R(s)$$

Dispersion relation

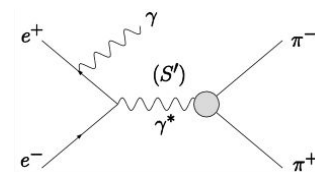
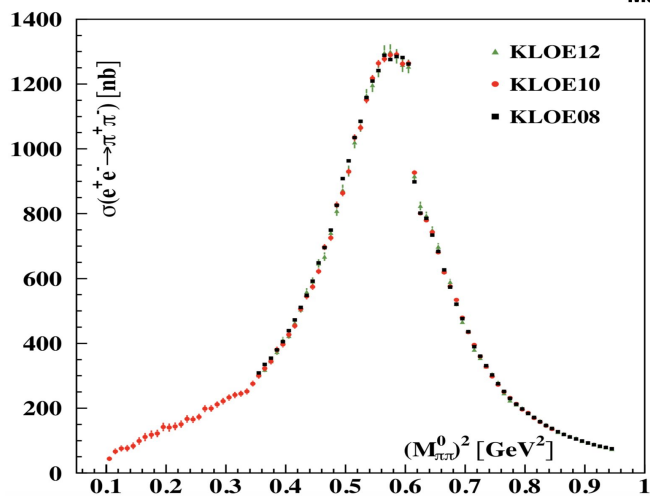
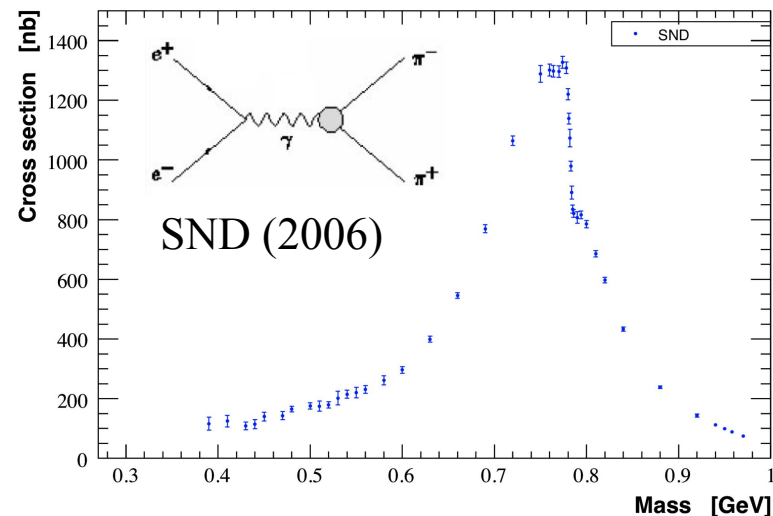
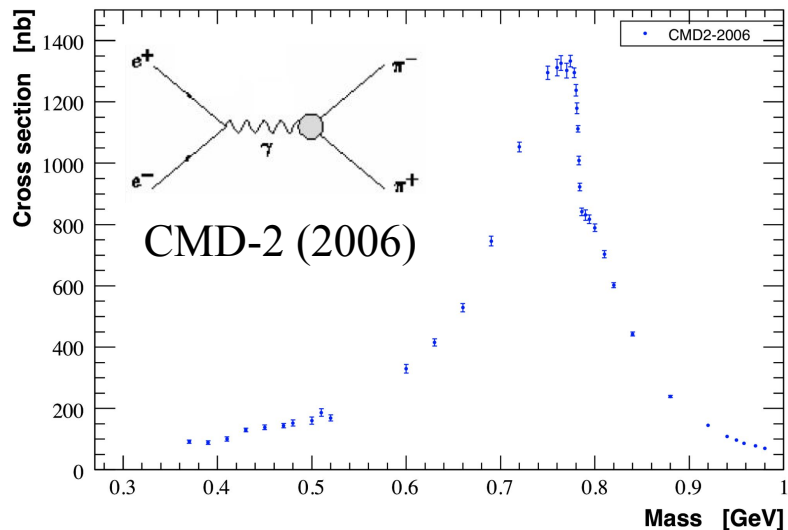
Bouchiat and Michel, 1961



→ Precise  $\sigma(e^+e^- \rightarrow \text{hadrons})$  measurements at low energy are very important

→ Do not use hadronic  $\tau$  decays data anymore (less precise + theory uncertainties)

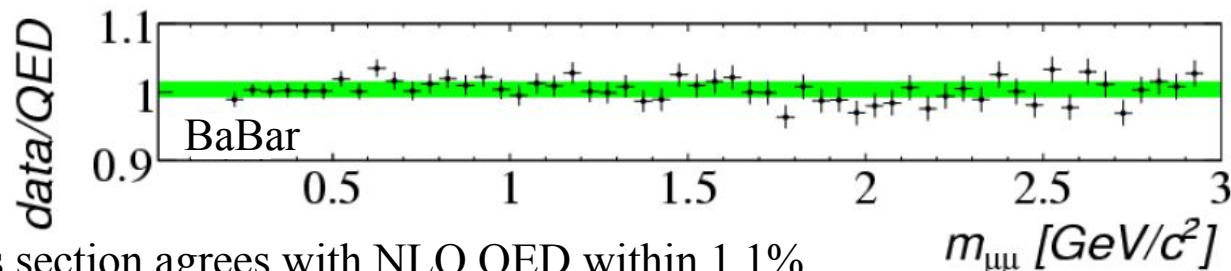
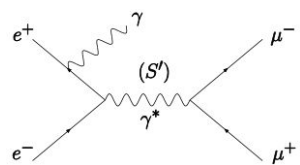
# HVP: Data on $e^+e^- \rightarrow \text{hadrons}$



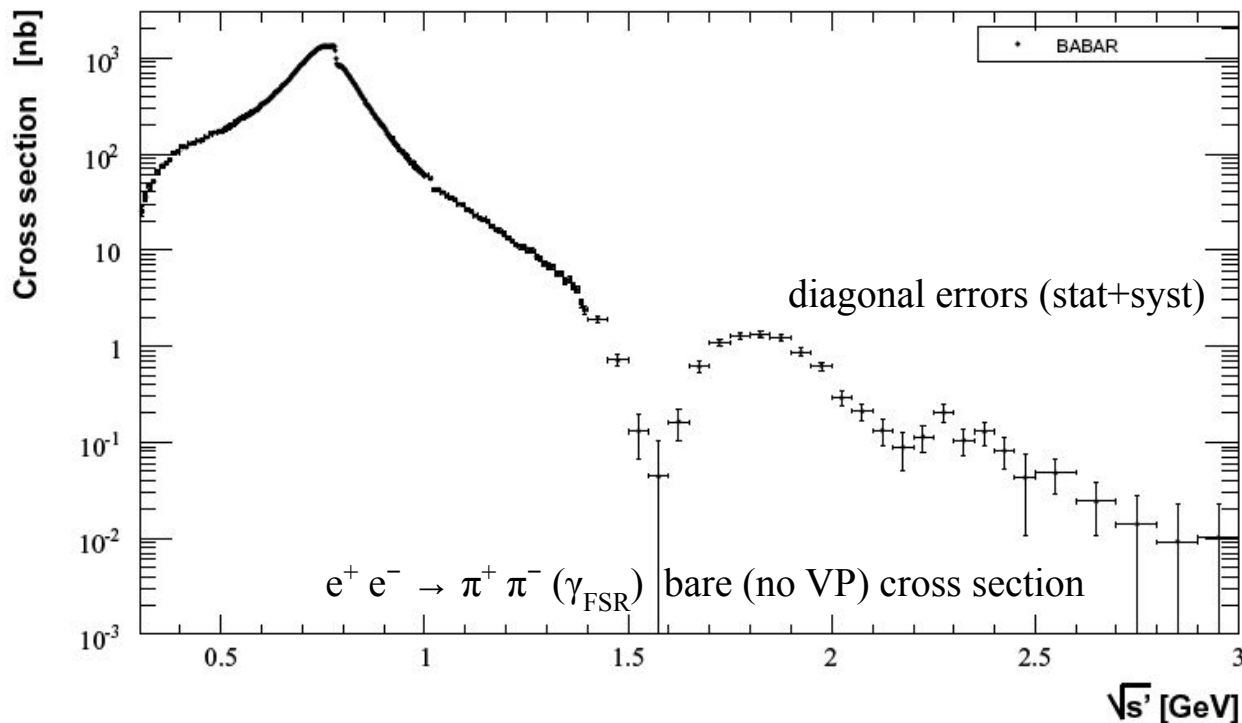
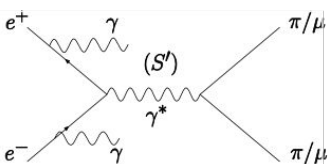
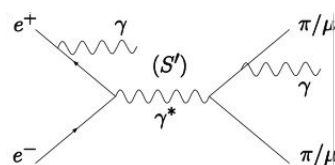
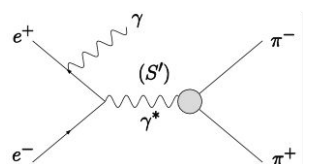
KLOE (08&10) +  $\mu\mu$  (12) (ISR)



# BaBar results (arXiv:0908.3589, PRL 103, 231801 (2009); arXiv:1205.2228(PRD))



Absolute  $\mu^+\mu^-$  cross section agrees with NLO QED within 1.1%



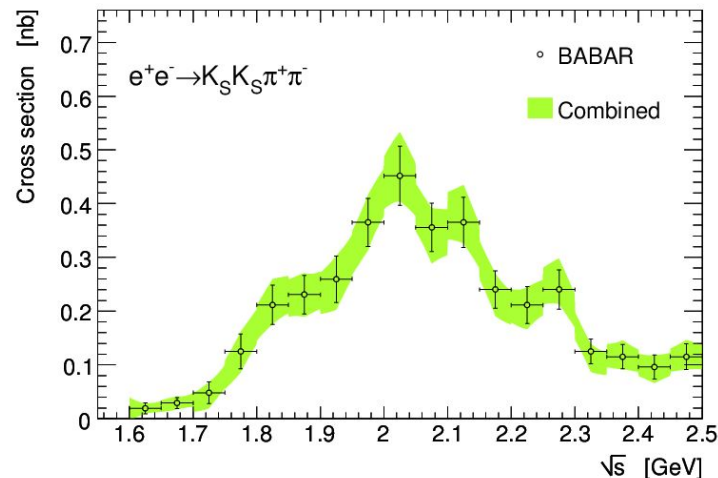
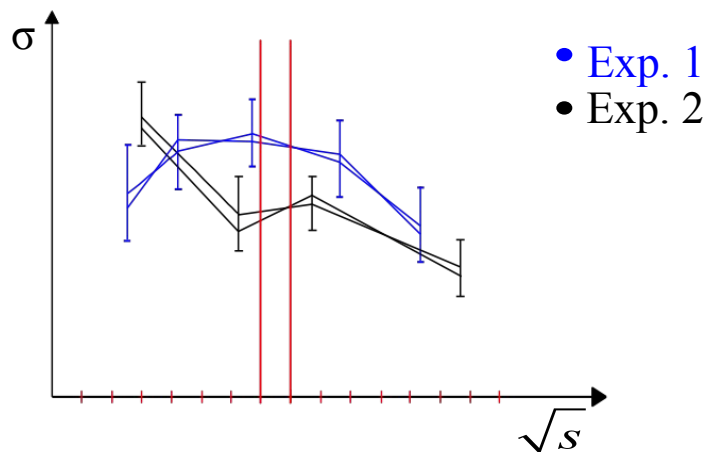
# Combine cross section data: goal and requirements

→ Goal: combine experimental spectra with arbitrary point spacing / binning

→ Requirements:

- Properly propagate uncertainties and correlations  
(*1<sup>st</sup> motivation for using DHMZ result as “baseline” in the TI White Paper*)
  - *Between measurements (data points/bins) of a given experiment*  
(covariance matrices and/or detailed split of uncertainties in sub-components)
  - *Between experiments (common systematic uncertainties, e.g. VP)*  
based on detailed information provided in publications
  - *Between different channels* – motivated by understanding of the meaning of systematic uncertainties and **identifying the common ones**  
BABAR luminosity (ISR or Bhabha), efficiencies (photon, Ks, Kl, modeling);  
BABAR radiative corrections;  $4\pi 2\pi^0 - \eta\omega$   
CMD2  $\eta\gamma - \pi^0\gamma$ ; CMD2/3 luminosity; SND luminosity;  
FSR; hadronic VP (old experiments)
- Minimize biases
- Optimize g-2 integral uncertainty  
(*without overestimating the precision with which the uncertainties of the measurements are known*)

# Combination procedure implemented in HVPTools software



- Define a (fine) final binning (to be filled and used for integrals etc.)
- Linear/quadratic splines to interpolate between the points/bins of each experiment
  - for binned measurements: preserve integral inside each bin
  - closure test: replace nominal values of data points by Gounaris-Sakurai model and re-do the combination
    - (non-)negligible bias for (linear)quadratic interpolation
- Fluctuate data points taking into account correlations & re-do the splines for each (pseudo-)experiment
  - each uncertainty fluctuated coherently for all the points/bins that it impacts
  - eigenvector decomposition for (statistical) covariance matrices
- In each fine bin: minimize  $\chi^2$  and get average coefficients

# Combination procedure implemented in HVPTools software

## For each final bin:

- Compute an average value for each measurement and its uncertainty
- Compute correlation matrix between experiments
- Minimize  $\chi^2$  and get average coefficients (weights)
- Compute average between experiments and its uncertainty

## Evaluation of integrals and propagation of uncertainties:

- Integral(s) evaluated for nominal result and for each set of toy pseudo-experiments; uncertainty of integrals from RMS of results for all toys
- The pseudo-experiments also used to derive (statistical & systematic) covariance matrices of combined cross sections → Integral evaluation
- Uncertainties also propagated through  $\pm 1\sigma$  shifts of each uncertainty:
  - allows to account for correlations between different channels (for integrals and spectra)
- *Checked consistency between the different approaches*

# Combination procedure: weights of various measurements

For each final bin:

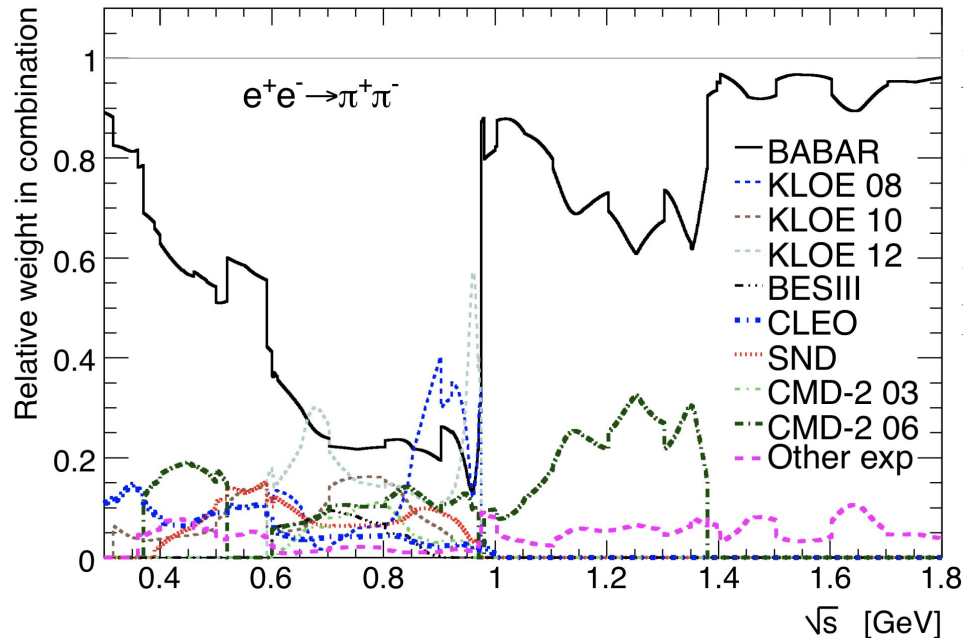
→ Minimize  $\chi^2$  and get average coefficients

Note: average weights must account for bin sizes / point spacing of measurements

(do not over-estimate the weight of experiments with large bins)

→ weights in fine bins evaluated using a common (large) binning for measurements + interpolation

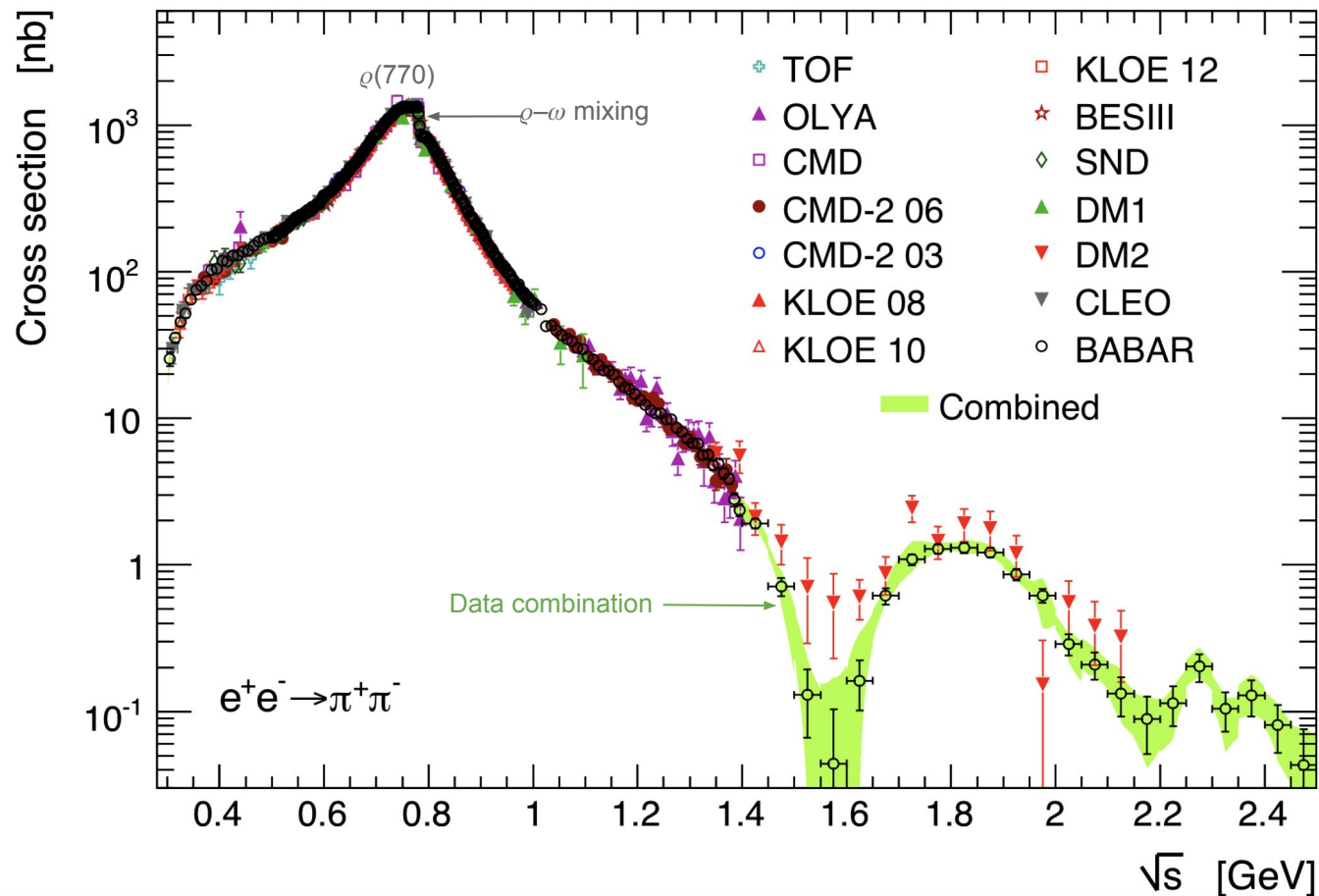
→ compare the precisions on the same footing



→ Bins used by KLOE larger than the ones by BABAR in  $\rho$ - $\omega$  interference region (factor  $\sim 3$ )

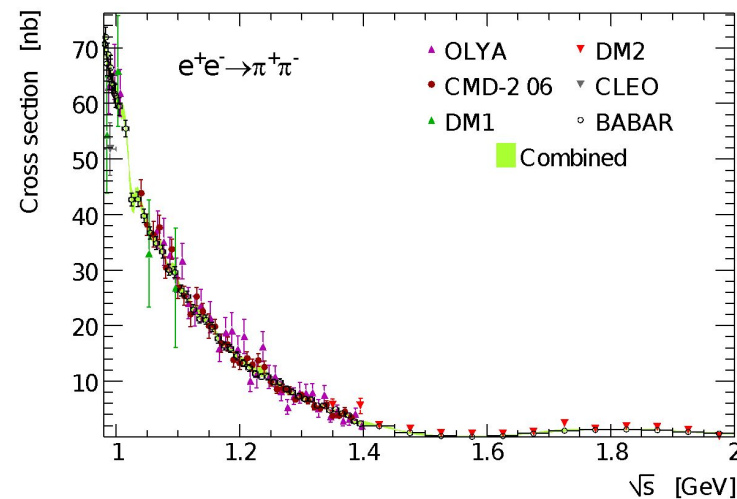
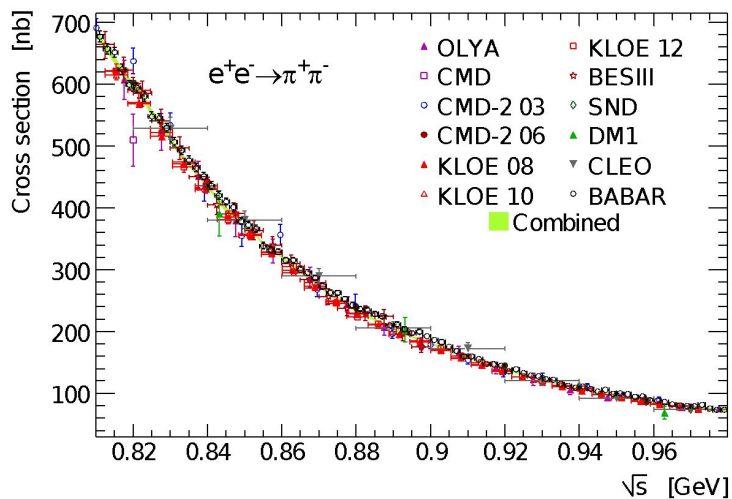
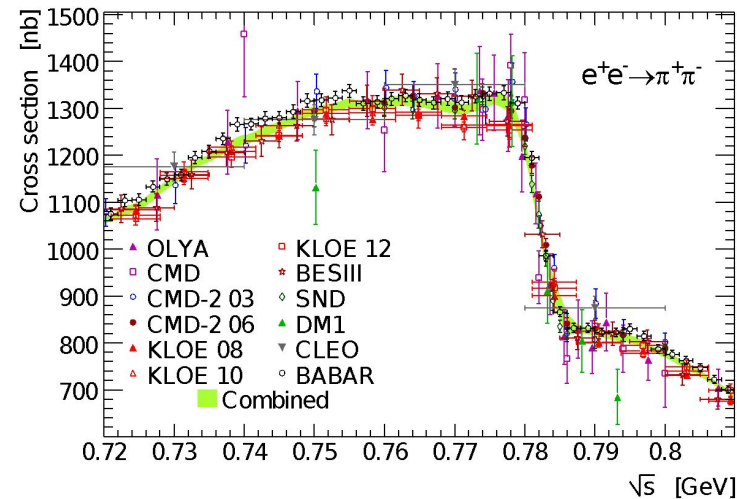
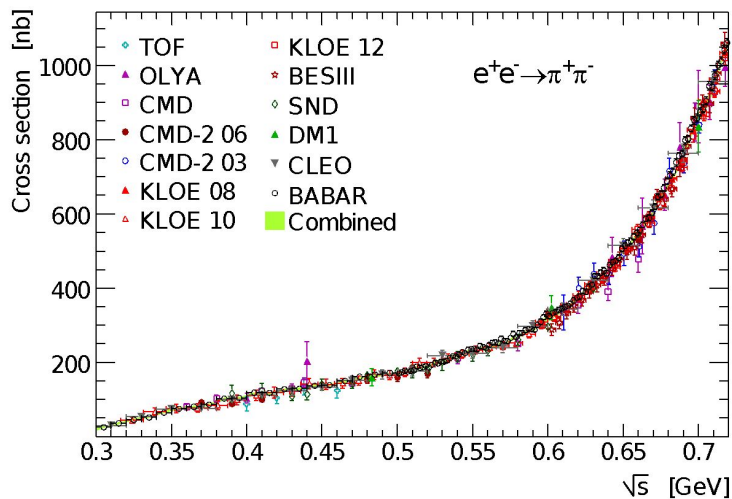
→ Average dominated by BaBar and KLOE, BaBar covering full range

# Combination for the $e^+e^- \rightarrow \pi^+\pi^-$ channel

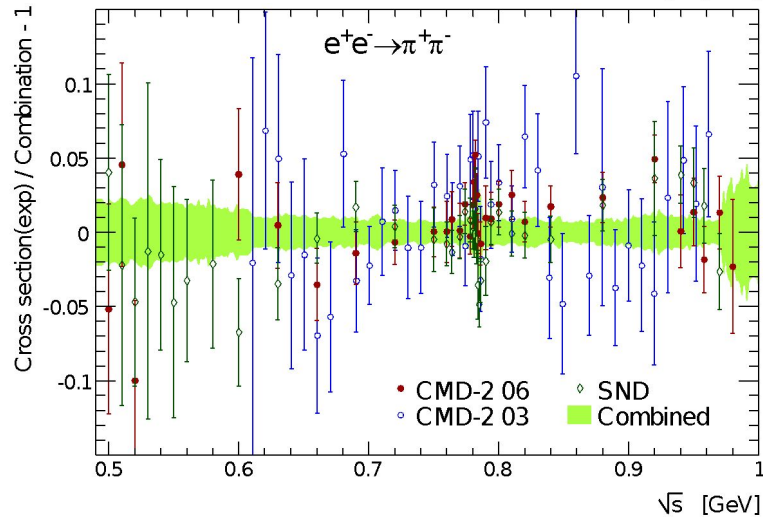
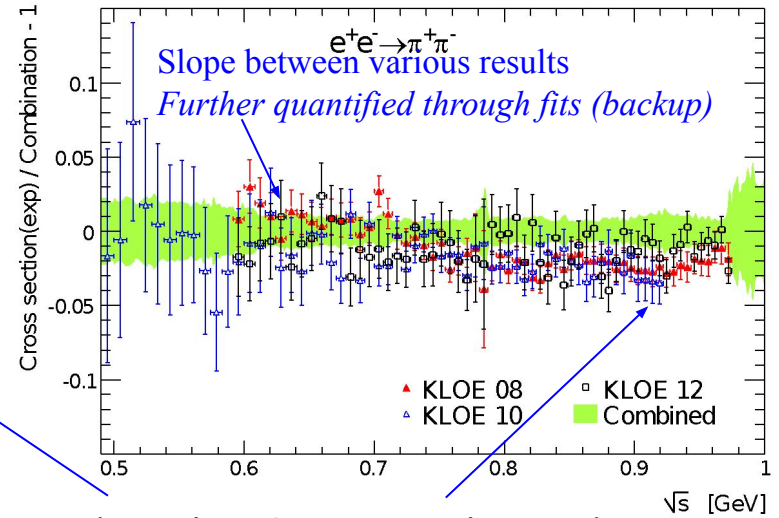
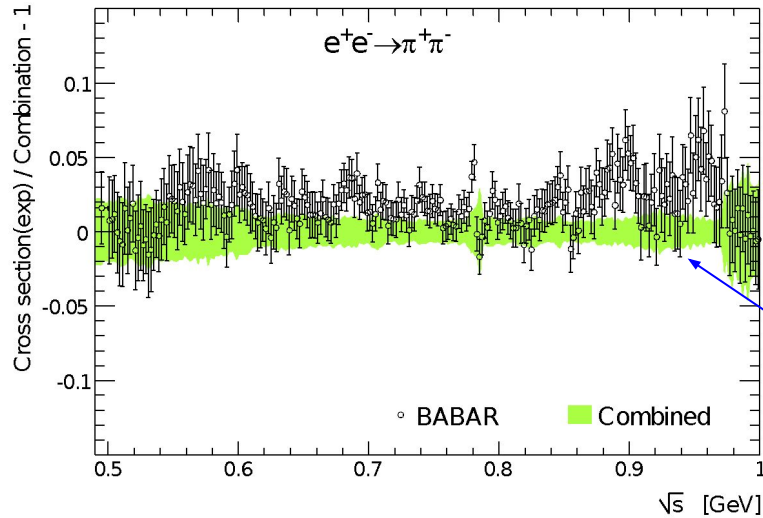


Procedure and software (HVPTools) for combining cross section data with arbitrary point spacing/binning

# Combination for the $e^+e^- \rightarrow \pi^+\pi^-$ channel



# Combination for the $e^+e^- \rightarrow \pi^+\pi^-$ channel



Local tension & systematic trends  
Indication of “uncertainties on uncertainties”  
(i.e. unaccounted biases)

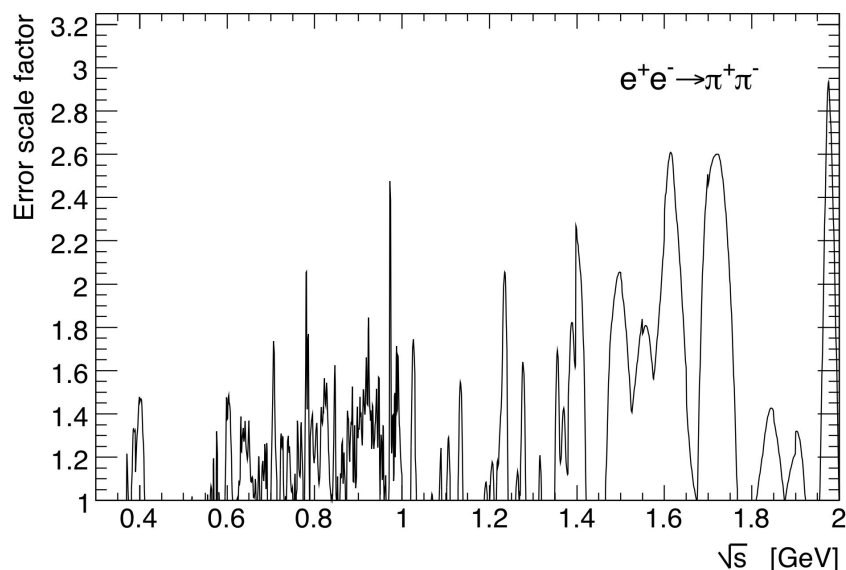
Other experiments not yet precise enough  
to discriminate  
(see however recent update from SND)



# Combination procedure: compatibility between measurements

For each final bin:

- $\chi^2 / \text{ndof}$ : test locally the level of agreement between input measurements, *taking into account the correlations*
- Scale uncertainties in bins where  $\chi^2 / \text{ndof} > 1$  (PDG): *locally conservative*
- Observed tension between BABAR and KLOE measurements



- Tension between measurements:  
*indication of underestimated uncertainties*  
Motivates conservative uncertainty treatment  
in combination fit (evaluation of weights)

- Included extra (dominant) uncertainty: difference between integrals without either BABAR or KLOE measurement to account for systematic deviations  
(*2<sup>nd</sup> motivation for using DHMZ result as “baseline” in White Paper*)

# Improving $a_\mu$ through fits for the $e^+e^- \rightarrow \pi^+\pi^-$ channel

→ Fit bare form-factor using 6 param. model based on *analyticity* and *unitarity*

$$|F_\pi^0|^2 = |R(s) \times J(s)|^2$$

$$R(s) = 1 + \alpha_V s + \frac{\kappa s}{m_\omega^2 - s - im_\omega \Gamma_\omega} \quad (1611.09359, \text{C. Hanhart et al.})$$

$$J(s) = e^{1 - \frac{\delta_1(s_0)}{\pi}} \left(1 - \frac{s}{s_0}\right)^{\left[1 - \frac{\delta_1(s_0)}{\pi}\right] \frac{s_0}{s}} \left(1 - \frac{s}{s_0}\right)^{-1} e^{\frac{s}{\pi} \int_{4m_\pi^2}^{s_0} dt \frac{\delta_1(t)}{t(t-s)}} \quad \text{Omnès integral}$$

(hep-ph/0402285, F.J. Yndurain et al.)

$$\cot \delta_1(s) = \frac{\sqrt{s}}{2k^3} (m_\rho^2 - s) \left[ \frac{2m_\pi^3}{m_\rho^2 \sqrt{s}} + B_0 + B_1 \omega(s) \right]$$

$$k = \frac{\sqrt{s - 4m_\pi^2}}{2}$$

(1102.2183, F.J. Yndurain et al.)

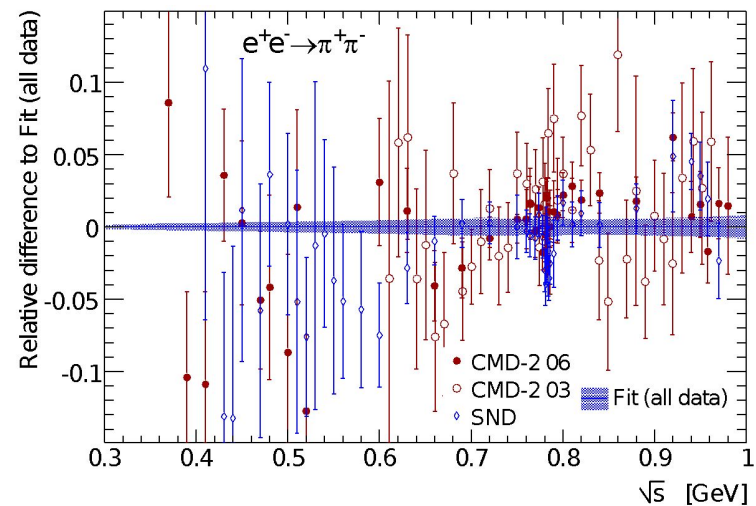
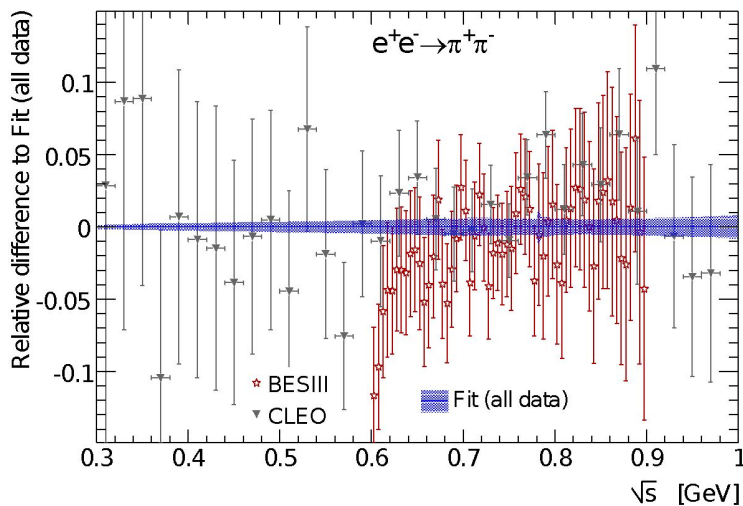
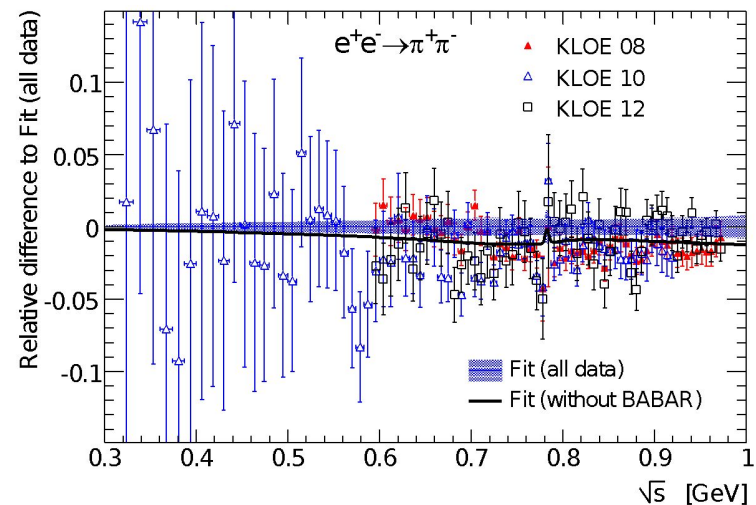
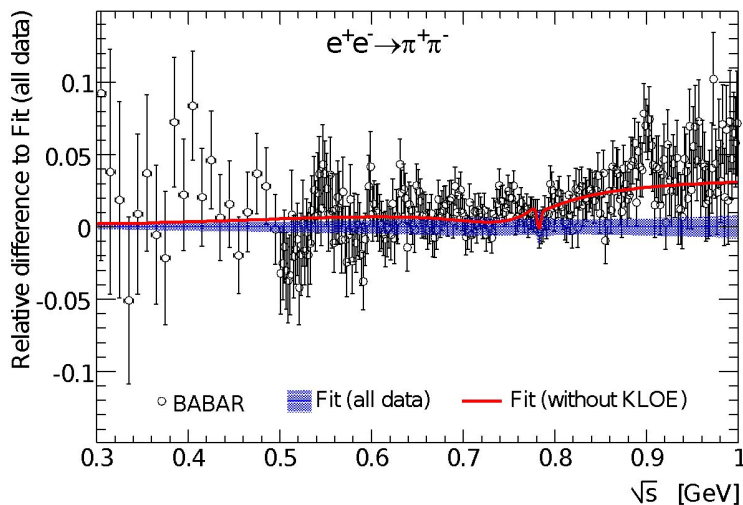
$$\omega(s) = \frac{\sqrt{s} - \sqrt{s_0 - s}}{\sqrt{s} + \sqrt{s_0 - s}} \quad \sqrt{s_0} = 1.05 \text{ GeV}$$

→ Conservative  $\chi^2$  (diagonal matrix) & local rescaling of input uncertainties

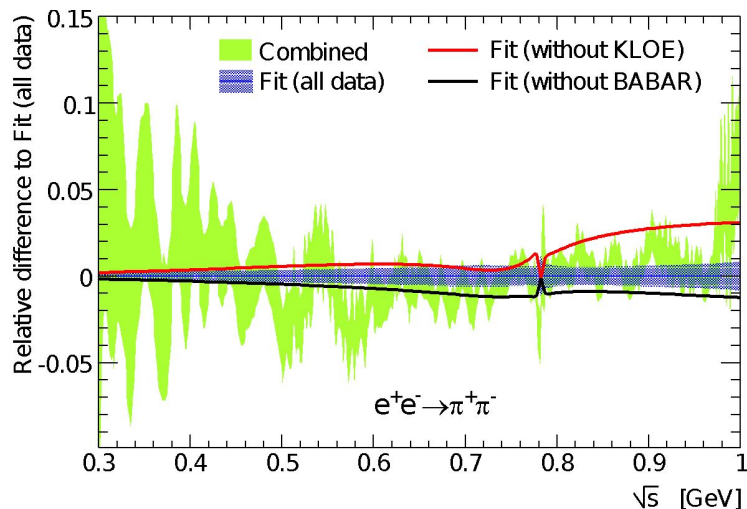
→ Full propagation of uncertainties & correlations using pseudo-experiments

DHMZ - 1908.00921

# Fit performed up to 1 GeV: comparison with data



# Fit performed up to 1 GeV, Result used up to 0.6 GeV



→ Use fit only below 0.6 GeV for  $a_\mu$  integral:

- where data is less precise and scarce
- less impacted by potential uncertainties of inelastic effects

$\sqrt{s}$ range [GeV]	$a_\mu^{\text{had}} [10^{-10}]$ Fit	$a_\mu^{\text{had}} [10^{-10}]$ Data Integration
0.3 - 0.6	$109.80 \pm 0.37_{\text{exp}} \pm 0.36_{\text{para}^*}$	$109.6 \pm 1.0_{\text{exp}}$

→ The difference  $0.2 \pm 0.9$   
(72% correlation accounted for)

→ The fit improves the precision by a factor  $\sim 2$

(\*) Parameter uncertainty corresponds to variations with/without the  $B_1$  term in the phase shift formula and  $\sqrt{s}_0$  varied from 1.05 GeV to 1.3 GeV (absolute values summed linearly), *checked to be statistically significant*

# Combined results: Fit [ $<0.6\text{GeV}$ ] + Data[ $0.6-1.8\text{GeV}$ ]

→ Full uncertainty propagation using the same pseudo-experiments as for the spline-based combination: 62% correlation among the two contributions

$\sqrt{s}$ range [GeV]	$a_{\mu}^{\text{had}} [10^{-10}]$ All data	$a_{\mu}^{\text{had}} [10^{-10}]$ All but BABAR	$a_{\mu}^{\text{had}} [10^{-10}]$ All but KLOE
threshold - 1.8	$506.9 \pm 1.9_{\text{total}}$	$505.0 \pm 2.1_{\text{total}}$	$510.6 \pm 2.2_{\text{total}}$

→ The difference “All but BABAR” and “All but KLOE” = 5.6  
to be compared with 1.9 uncertainty with “All data”

- The local error inflation is not sufficient to amplify the uncertainty
- Global tension (normalisation/shape) not previously accounted for
- Potential underestimated uncertainty in at least one of the measurements?
- Other measurements not precise enough to discriminate BABAR / KLOE

→ Given the fact we do not know which dataset is problematic, we decide to:

- Add half of the discrepancy (2.8) as an additional uncertainty  
(correcting the local PDG inflation to avoid double counting)
- Take (“All but BABAR” + “All but KLOE”) / 2 as central value

## Uncertainties on uncertainties and on correlations

*Topic of general interest, in other fields too (see backup)*

[1908.00921\(DHMZ\)](#), [2006.04822\(WP Theory Initiative\)](#)

# Two different approaches for combining ( $e^+e^-$ ) data

DHMZ:

- $\chi^2$  computed locally (in each fine bin), taking into account correlations between measurements (see previous slides)
- used to determine the weights on the measurements in the combination and their level of agreement
- uncertainties and correlations propagated using pseudo-experiments or  $\pm 1\sigma$  shifts of each uncertainty component

KNT:

- $\chi^2$  computed globally (for full mass range)

$$\chi_I^2 = \sum_{i=1}^{N_{\text{tot}}} \sum_{j=1}^{N_{\text{tot}}} (R_i^{(m)} - \mathcal{R}_m^{i,I}) \mathbf{C}_I^{-1}(i^{(m)}, j^{(n)}) (R_j^{(n)} - \mathcal{R}_n^{j,I}) \quad \text{KNT (1802.02995)}$$

$$\chi^2 = \sum_{i=1}^{195} \sum_{j=1}^{195} (\sigma_{\pi\pi(\gamma)}^0(i) - \bar{\sigma}_{\pi\pi(\gamma)}^0(m)) \mathbf{C}^{-1}(i^{(m)}, j^{(n)}) (\sigma_{\pi\pi(\gamma)}^0(j) - \bar{\sigma}_{\pi\pi(\gamma)}^0(n)) \quad \text{KLOE-KMT (1711.03085)}$$

- relies on description of correlations on long ranges

- *One of the main sources of differences for the uncertainty on  $a_\mu$*

# Evaluation of uncertainties and correlations ( $e^+e^-$ )

	$\sigma_{\pi\pi\gamma}$	$\sigma_{\pi\pi}^0$	$F_\pi$	$\Delta^{\pi\pi} a_\mu$
Reconstruction Filter	negligible			
Background subtraction	Tab. 1		0.3%	
Trackmass	0.2%			
Pion cluster ID	negligible			
Tracking efficiency	0.3%			
Trigger efficiency	0.1%			
Acceptance	Tab. 2		0.2%	
Unfolding	Tab. 3		negligible	
L3 filter	0.1%			
$\sqrt{s}$ dependence of $H$	-	Tab. 4		0.2%
Luminosity	0.3%			
Experimental systematics				0.6%
FSR resummation	-	0.3%		
Radiator function $H$	-	0.5%		
Vacuum Polarization	-	0.1%	-	0.1%
Theory systematics				0.6%

→ Systematics *evaluated* in  $\sim$ wide mass ranges with sharp transitions

$M_{\pi\pi}^2$ range (GeV <sup>2</sup> )	Systematic error (%)
$0.35 \leq M_{\pi\pi}^2 < 0.39$	0.6
$0.39 \leq M_{\pi\pi}^2 < 0.43$	0.5
$0.43 \leq M_{\pi\pi}^2 < 0.45$	0.4
$0.45 \leq M_{\pi\pi}^2 < 0.49$	0.3
$0.49 \leq M_{\pi\pi}^2 < 0.51$	0.2
$0.51 \leq M_{\pi\pi}^2 < 0.64$	0.1
$0.64 \leq M_{\pi\pi}^2 < 0.95$	-

KLOE 08 (0809.3950)

KLOE 10 (1006.5313)

	$\sigma_{\pi\pi\gamma}$	$\sigma_{\pi\pi}^{\text{bare}}$	$ F_\pi ^2$	$\Delta a_\mu^{\pi\pi}$
	threshold ; $\rho$ -peak			(0.1 - 0.85 GeV <sup>2</sup> )
Background Filter	0.5% ; 0.1%			negligible
Background subtraction	3.4% ; 0.1%			0.5%
$f_0 + \rho\pi$ bkg.	6.5% ; negl.			0.4%
$\Omega$ cut	1.4% ; negl.			0.2%
Trackmass cut	3.0% ; 0.2%			0.5%
$\pi$ -e PID	0.3% ; negl.			negligible
Trigger	0.3% ; 0.2%			0.2%
Acceptance	1.9% ; 0.3%			0.5%
Unfolding	negl. ; 2.0%			negligible
Tracking	0.3%			
Software Trigger (L3)	0.1%			
Luminosity	0.3%			
Experimental syst.				1.0%
FSR treatment	-	7% ; negl.		0.8%
Radiator function $H$	-	0.5%		
Vacuum Polarization	-	Ref. 34	-	0.1%
Theory syst.				0.9%



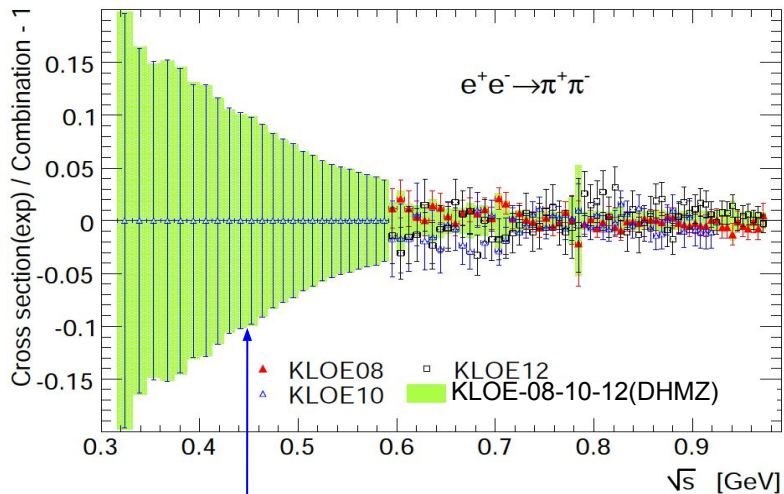
# Evaluation of uncertainties and correlations ( $e^+e^-$ )

Sources	0.3-0.4	0.4-0.5	0.5-0.6	0.6-0.9	0.9-1.2	1.2-1.4	1.4-2.0	2.0-3.0
trigger/ filter	5.3	2.7	1.9	1.0	0.7	0.6	0.4	0.4
tracking	3.8	2.1	2.1	1.1	1.7	3.1	3.1	3.1
$\pi$ -ID	10.1	2.5	6.2	2.4	4.2	10.1	10.1	10.1
background	3.5	4.3	5.2	1.0	3.0	7.0	12.0	50.0
acceptance	1.6	1.6	1.0	1.0	1.6	1.6	1.6	1.6
kinematic fit ( $\chi^2$ )	0.9	0.9	0.3	0.3	0.9	0.9	0.9	0.9
correl $\mu\mu$ ID loss	3.0	2.0	3.0	1.3	2.0	3.0	10.0	10.0
$\pi\pi/\mu\mu$ non-cancel.	2.7	1.4	1.6	1.1	1.3	2.7	5.1	5.1
unfolding	1.0	2.7	2.7	1.0	1.3	1.0	1.0	1.0
ISR luminosity	3.4	3.4	3.4	3.4	3.4	3.4	3.4	3.4
sum (cross section)	13.8	8.1	10.2	5.0	6.5	13.9	19.8	52.4

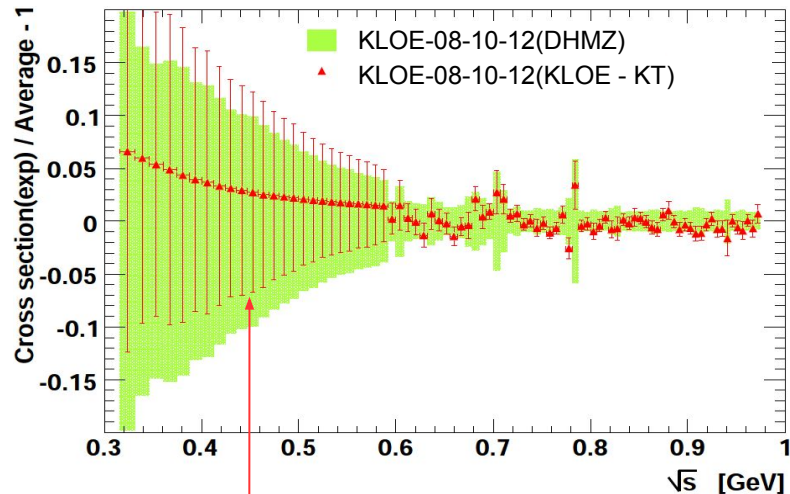
BABAR (1205.2228)

→ Systematics *evaluated* in  $\sim$ wide mass ranges with sharp transitions  
(statistics limitations when going to narrow ranges)

# Combining the 3 KLOE measurements



Local combination (DHMZ)



Information propagated between mass regions, through shifts of systematics - relying on correlations, amplitudes and shapes of systematics (KLOE-KT)

# Combining the 3 KLOE measurements - $a_\mu^{\pi\pi}$ contribution

KLOE08  $a_\mu[0.6 ; 0.9] : 368.3 \pm 3.2 [10^{-10}]$

KLOE10  $a_\mu[0.6 ; 0.9] : 365.6 \pm 3.3$

KLOE12  $a_\mu[0.6 ; 0.9] : 366.8 \pm 2.5$

→ Correlation matrix:

	08		10		12	
--	----	--	----	--	----	--

-----  
08 |      1    0.70    0.35

10 |    0.70      1    0.19

12 |    0.35    0.19      1

→ Amount of independent information provided by each measurement

→ KLOE-08-10-12(DHMZ) -  $a_\mu[0.6 ; 0.9] : 366.5 \pm 2.8$  (Without  $\chi^2$  rescaling:  $\pm 2.2$ )

→ Conservative treatment of uncertainties and correlations (*not perfectly known*) in weight determination

→ KLOE-08-10-12(KLOE-KT) -  $a_\mu[0.6 ; 0.9]\text{GeV} : 366.9 \pm 2.2$  (Includes  $\chi^2$  rescaling)

→ Assuming perfect knowledge of the correlations to minimize average uncertainty

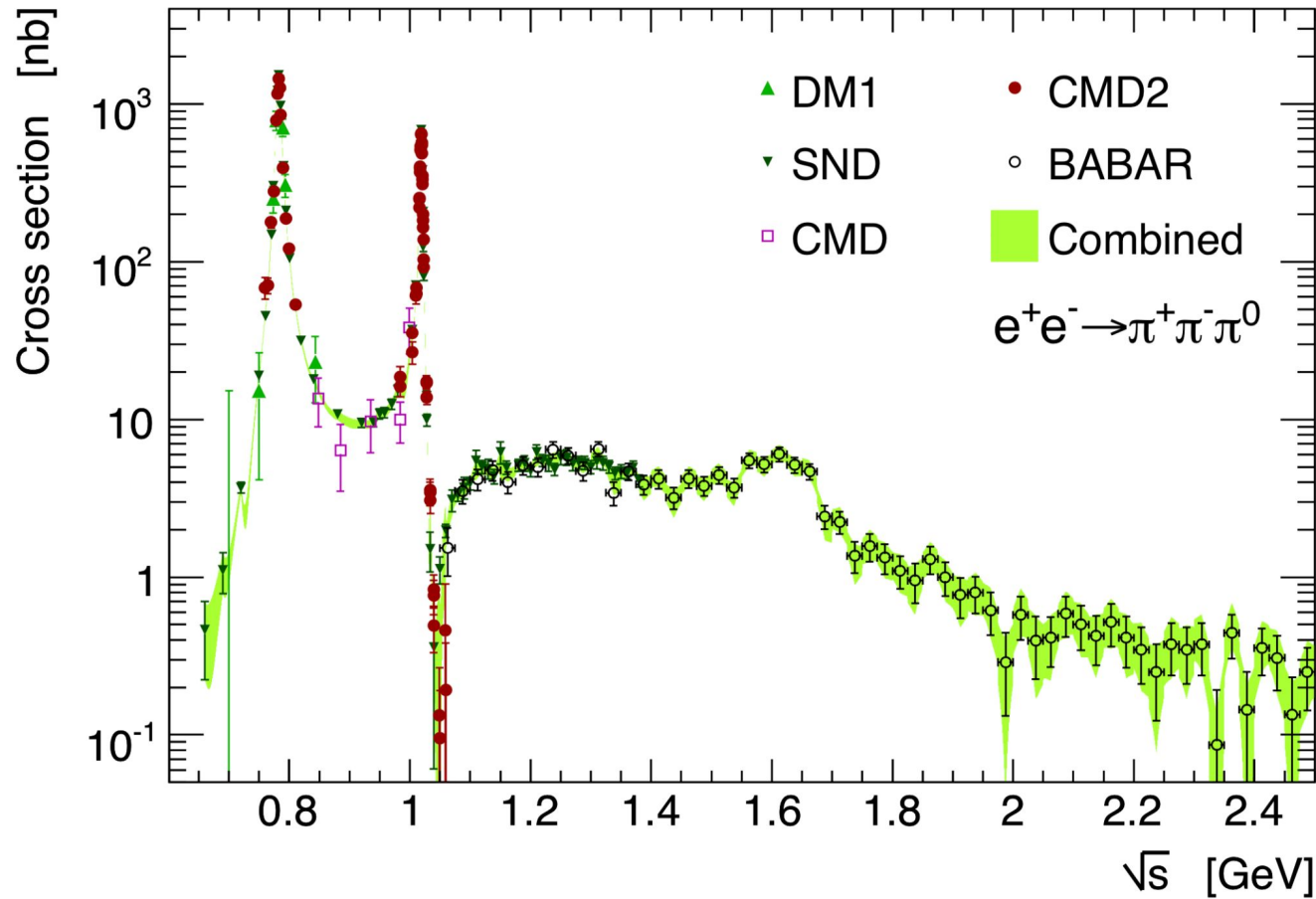
# Uncertainties on uncertainties and correlations

*Numerous indications of uncertainties on uncertainties and on correlations, with a direct impact on combination fits*

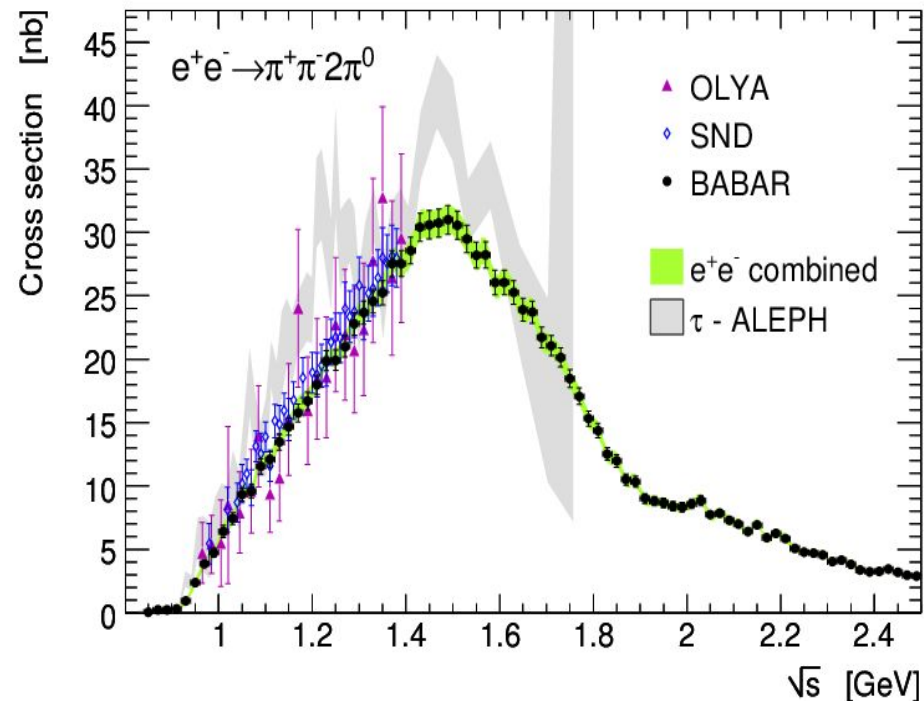
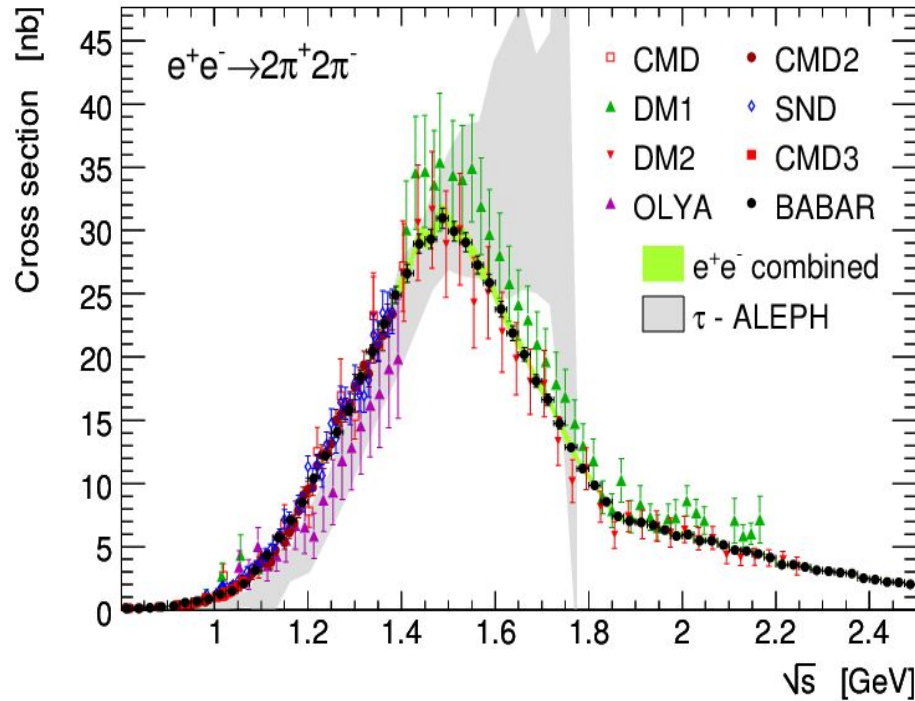
- Shapes of systematic uncertainties *evaluated* in  $\sim$ -wide mass ranges with sharp transitions
- One standard deviation is statistically not well defined for systematic uncertainties
- Systematic uncertainties like acceptance, tracking efficiency, background etc. not necessarily fully correlated between low and high mass
- Are all systematic uncertainty components fully independent between each-other? (e.g. tracking and trigger)
- *Yield uncertainties on uncertainties and on correlations*
- Tensions between measurements (BABAR/KLOE; 3 KLOE results etc.): *experimental indications of underestimated uncertainties*
- *Statistical methods* ( $\chi^2$  with correlations, likelihood fits, ratios of measured quantities etc.) *should not over-exploit the information on the amplitude and correlations of uncertainties*

# Combination of measurements for various channels and total HVP contribution

# Combination for the $e^+e^- \rightarrow \pi^+\pi^-\pi^0$ channel

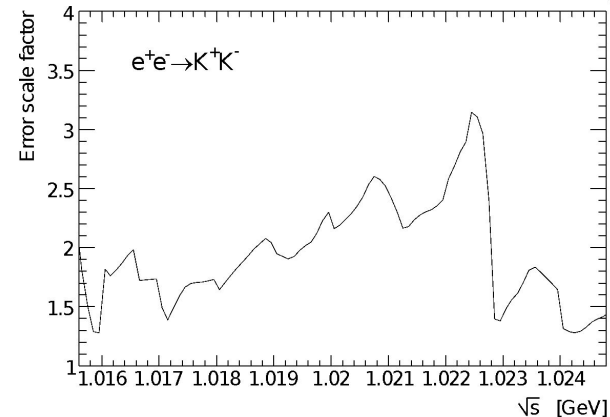
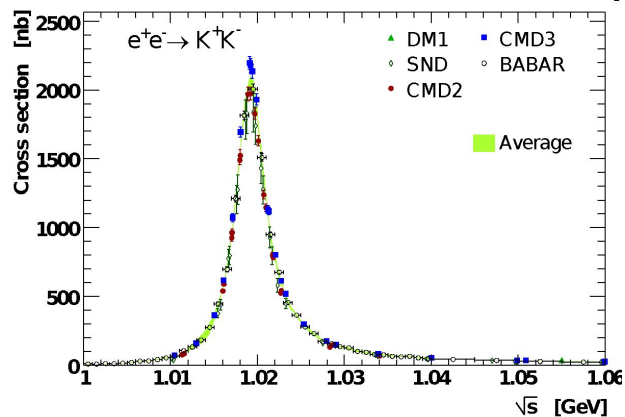
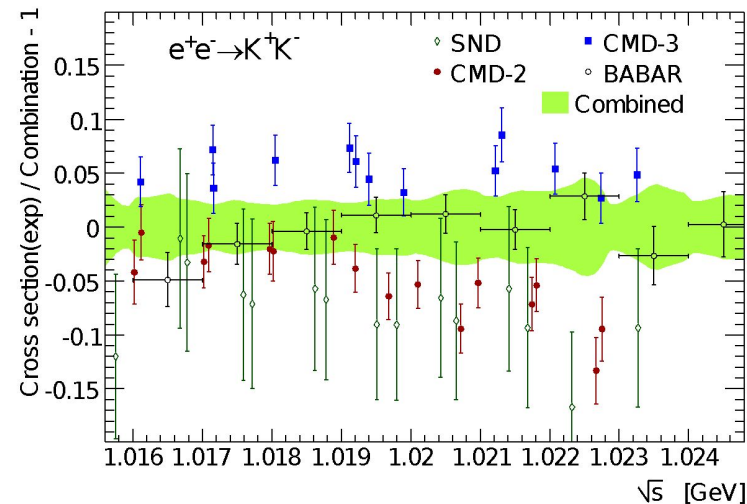
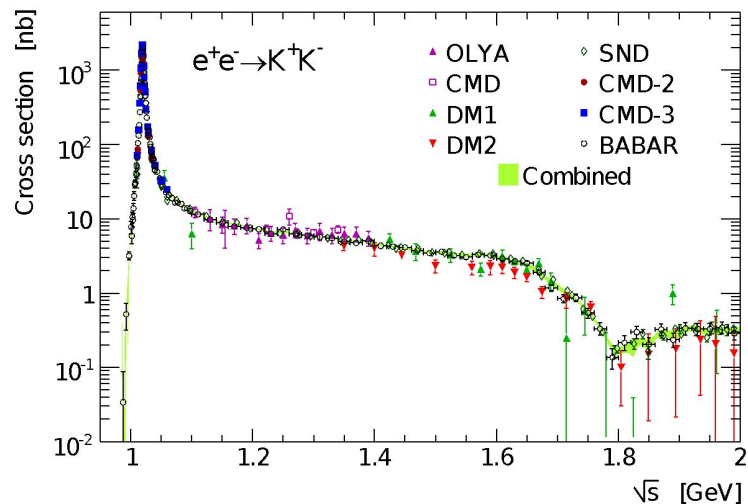


$$e^+e^- \rightarrow \pi^+\pi^-\pi^+\pi^-, e^+e^- \rightarrow \pi^+\pi^-\pi^0\pi^0$$



→ Essentially normalization differences w.r.t.  $\tau$  data: *cross-checks very desirable*

# Combination for the $e^+e^- \rightarrow K^+K^-$ channel

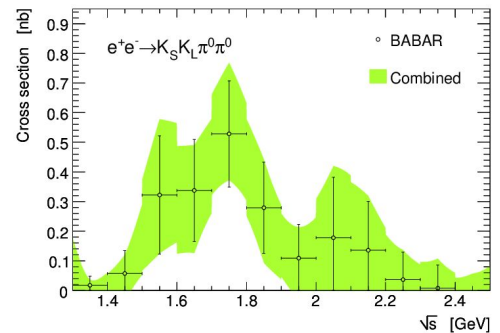
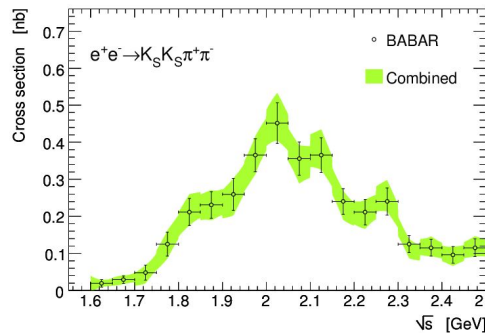
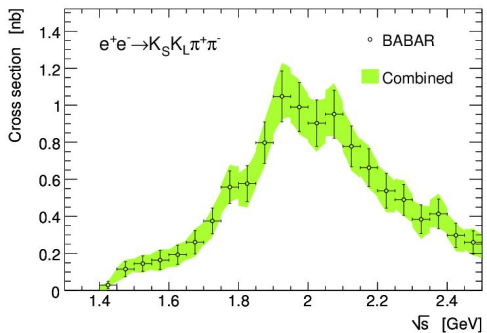
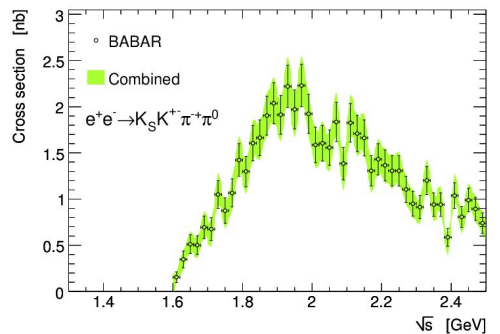
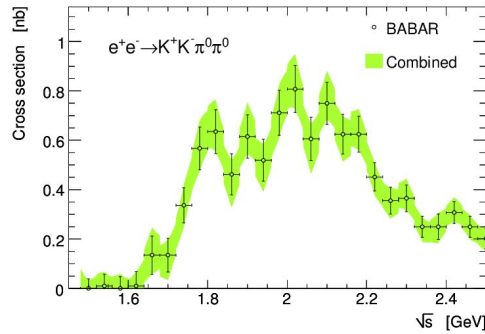
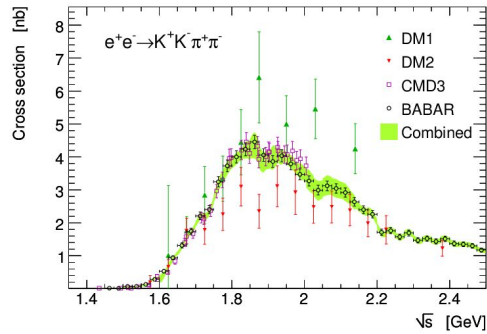
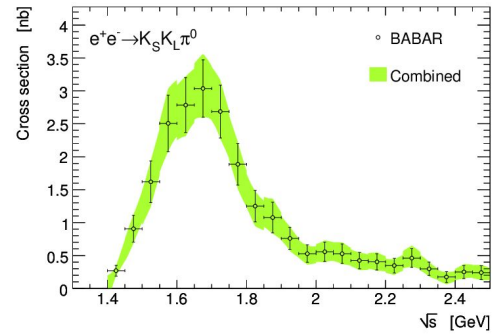
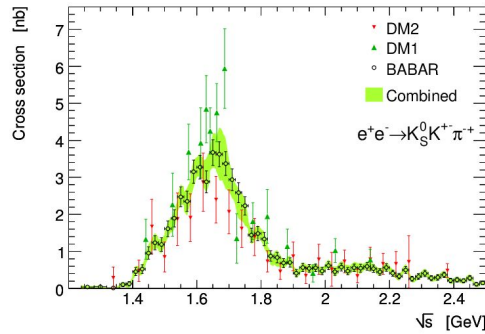
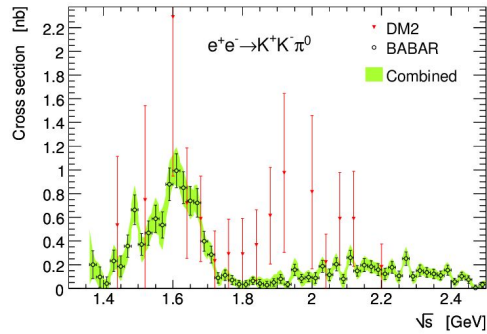


→ Tension between measurements

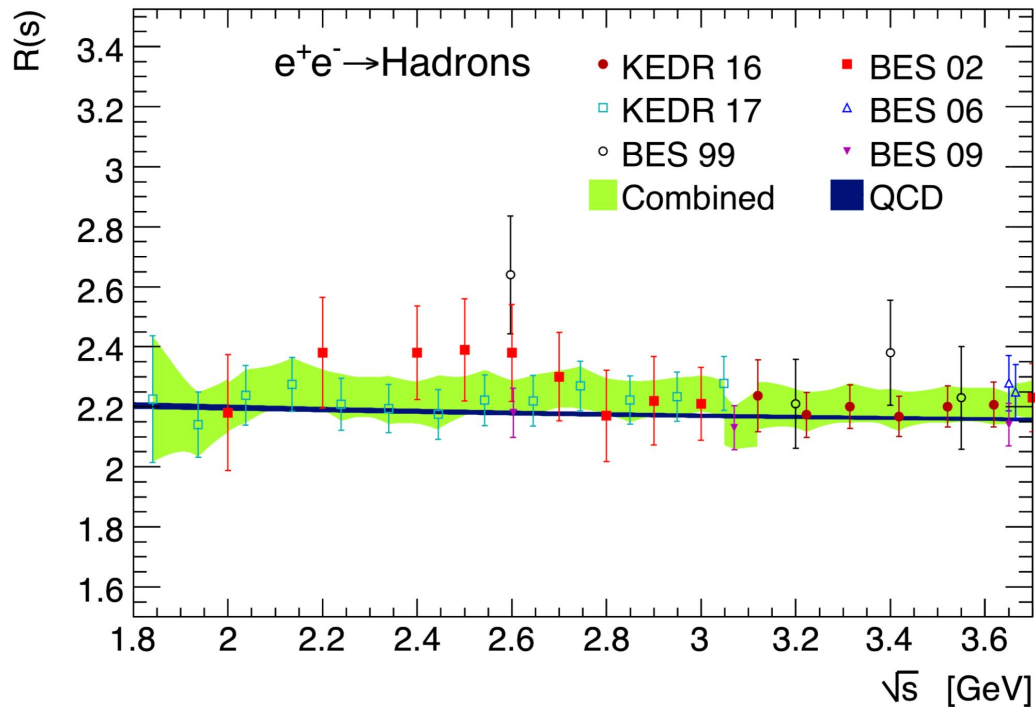
→  $a_\mu[\rightarrow 1.8\text{GeV}]$ :  $23.08 \pm 0.20$  (stat.)  $\pm 0.40$  (syst.) [ $10^{-10}$ ] (enhancement x 2.2)



# Combination for the $e^+e^- \rightarrow KK\pi$ and $KK2\pi$ channels

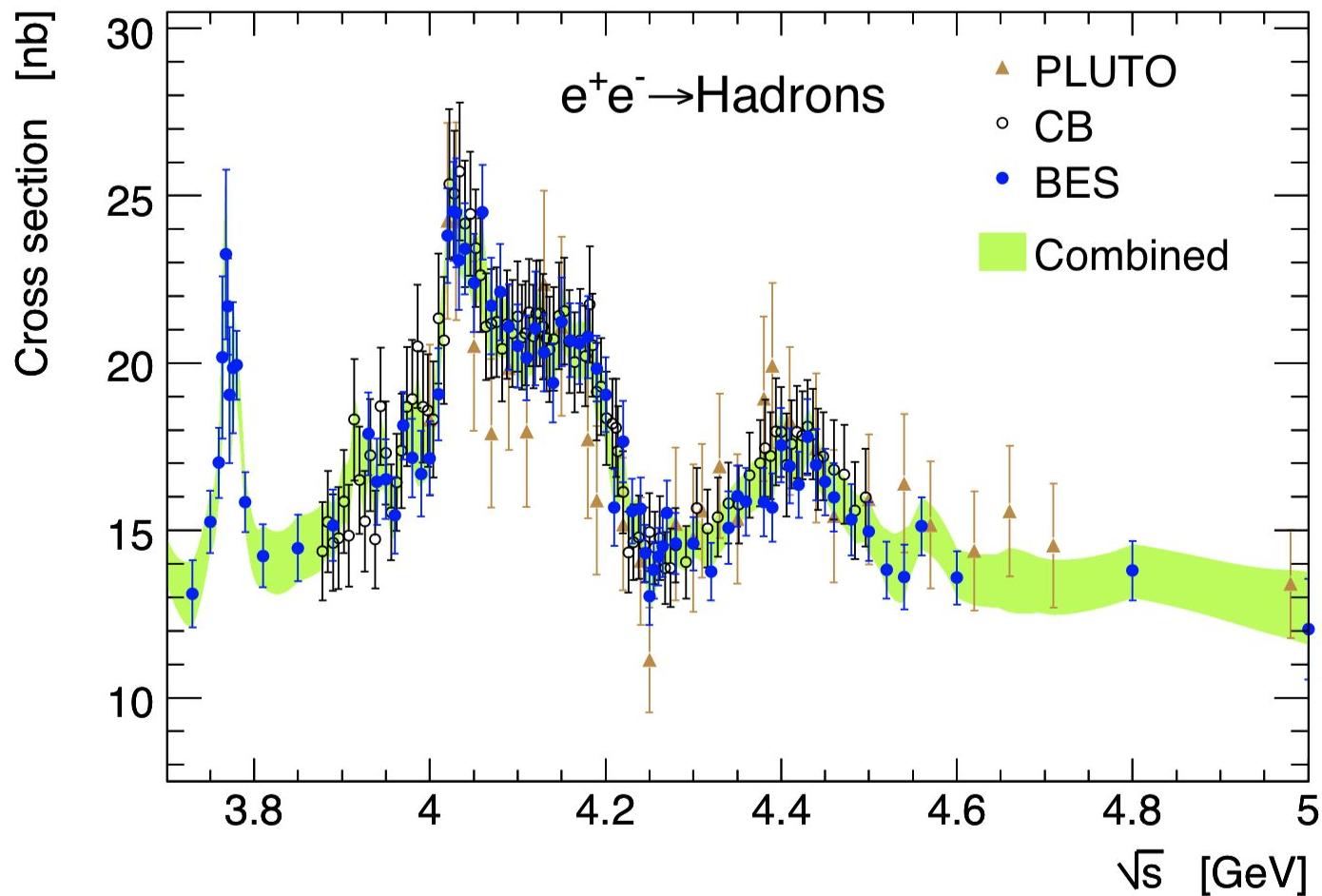


# Contributions from the 1.8 – 3.7 GeV region



- Contribution evaluated from pQCD (4 loops) +  $O(\alpha_s^2)$  quark mass corrections
- Uncertainties:  $\alpha_s$ , truncation of perturbative series, CIPT/FOPT,  $m_q$
- 1.8-2.0 GeV:  $7.65 \pm 0.31$  (data excl.);  $8.30 \pm 0.09$  (QCD); added syst.  $0.65 [10^{-10}]$
- 2.0-3.7 GeV:  $25.82 \pm 0.61$  (data);  $25.15 \pm 0.19$  (QCD); agreement within  $1\sigma$

# Contributions from the charm resonance region



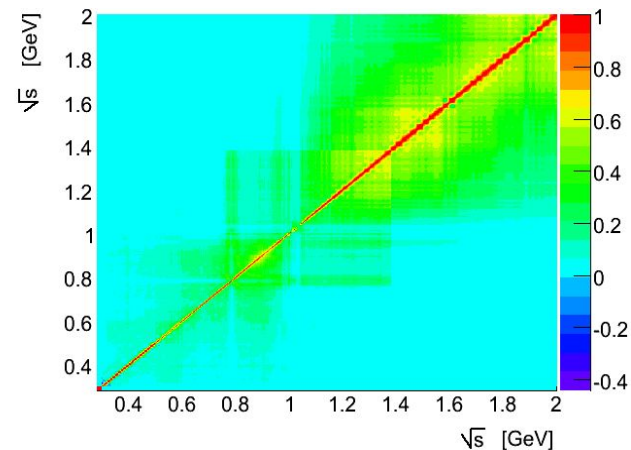
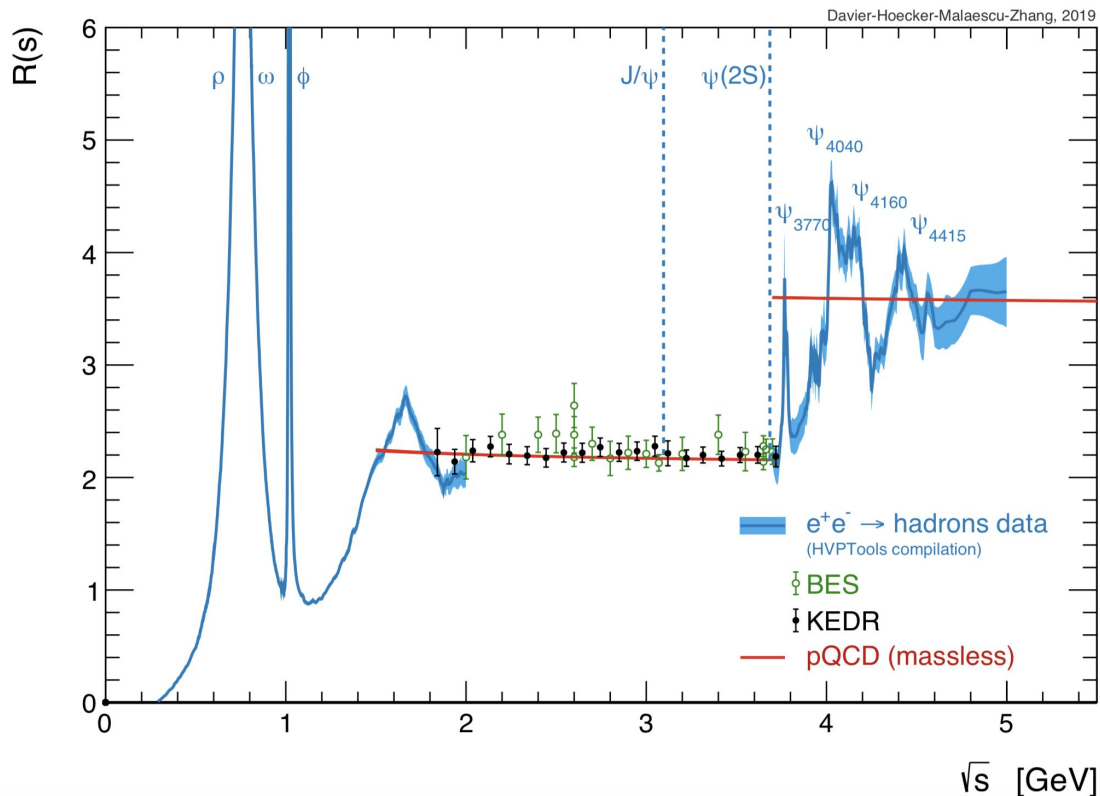
# Situation in arXiv:1908.00921 (EPJC)

Channel	$a_\mu^{\text{had,LO}} [10^{-10}]$	$\Delta\alpha_{\text{had}}(m_Z^2) [10^{-4}]$
$\pi^0\gamma$	$4.41 \pm 0.06 \pm 0.04 \pm 0.07$	$0.35 \pm 0.00 \pm 0.00 \pm 0.01$
$\eta\gamma$	$0.65 \pm 0.02 \pm 0.01 \pm 0.01$	$0.08 \pm 0.00 \pm 0.00 \pm 0.00$
$\pi^+\pi^-$	$507.85 \pm 0.83 \pm 3.23 \pm 0.55$	$34.50 \pm 0.06 \pm 0.20 \pm 0.04$
$\pi^+\pi^-\pi^0$	$46.21 \pm 0.40 \pm 1.10 \pm 0.86$	$4.60 \pm 0.04 \pm 0.11 \pm 0.08$
$2\pi^+2\pi^-$	$13.68 \pm 0.03 \pm 0.27 \pm 0.14$	$3.58 \pm 0.01 \pm 0.07 \pm 0.03$
$\pi^+\pi^-2\pi^0$	$18.03 \pm 0.06 \pm 0.48 \pm 0.26$	$4.45 \pm 0.02 \pm 0.12 \pm 0.07$
$2\pi^+2\pi^-\pi^0$ ( $\eta$ excl.)	$0.69 \pm 0.04 \pm 0.06 \pm 0.03$	$0.21 \pm 0.01 \pm 0.02 \pm 0.01$
$\pi^+\pi^-3\pi^0$ ( $\eta$ excl.)	$0.49 \pm 0.03 \pm 0.09 \pm 0.00$	$0.15 \pm 0.01 \pm 0.03 \pm 0.00$
$3\pi^+3\pi^-$	$0.11 \pm 0.00 \pm 0.01 \pm 0.00$	$0.04 \pm 0.00 \pm 0.00 \pm 0.00$
$2\pi^+2\pi^-2\pi^0$ ( $\eta$ excl.)	$0.71 \pm 0.06 \pm 0.07 \pm 0.14$	$0.25 \pm 0.02 \pm 0.02 \pm 0.05$
$\pi^+\pi^-4\pi^0$ ( $\eta$ excl., isospin)	$0.08 \pm 0.01 \pm 0.08 \pm 0.00$	$0.03 \pm 0.00 \pm 0.03 \pm 0.00$
$\eta\pi^+\pi^-$	$1.19 \pm 0.02 \pm 0.04 \pm 0.02$	$0.35 \pm 0.01 \pm 0.01 \pm 0.01$
$\eta\omega$	$0.35 \pm 0.01 \pm 0.02 \pm 0.01$	$0.11 \pm 0.00 \pm 0.01 \pm 0.00$
$\eta\pi^+\pi^-\pi^0$ (non- $\omega, \phi$ )	$0.34 \pm 0.03 \pm 0.03 \pm 0.04$	$0.12 \pm 0.01 \pm 0.01 \pm 0.01$
$\eta 2\pi^+2\pi^-$	$0.02 \pm 0.01 \pm 0.00 \pm 0.00$	$0.01 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega\eta\pi^0$	$0.06 \pm 0.01 \pm 0.01 \pm 0.00$	$0.02 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega\pi^0$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.94 \pm 0.01 \pm 0.03 \pm 0.00$	$0.20 \pm 0.00 \pm 0.01 \pm 0.00$
$\omega 2\pi$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.07 \pm 0.00 \pm 0.00 \pm 0.00$	$0.02 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega$ (non- $3\pi, \pi\gamma, \eta\gamma$ )	$0.04 \pm 0.00 \pm 0.00 \pm 0.00$	$0.00 \pm 0.00 \pm 0.00 \pm 0.00$
$K^+K^-$	$23.08 \pm 0.20 \pm 0.33 \pm 0.21$	$3.35 \pm 0.03 \pm 0.05 \pm 0.03$
$K_S K_L$	$12.82 \pm 0.06 \pm 0.18 \pm 0.15$	$1.74 \pm 0.01 \pm 0.03 \pm 0.02$
$\phi$ (non- $K\bar{K}, 3\pi, \pi\gamma, \eta\gamma$ )	$0.05 \pm 0.00 \pm 0.00 \pm 0.00$	$0.01 \pm 0.00 \pm 0.00 \pm 0.00$
$K\bar{K}\pi$	$2.45 \pm 0.05 \pm 0.10 \pm 0.06$	$0.78 \pm 0.02 \pm 0.03 \pm 0.02$
$K\bar{K}2\pi$	$0.85 \pm 0.02 \pm 0.05 \pm 0.01$	$0.30 \pm 0.01 \pm 0.02 \pm 0.00$
$K\bar{K}\omega$	$0.00 \pm 0.00 \pm 0.00 \pm 0.00$	$0.00 \pm 0.00 \pm 0.00 \pm 0.00$
$\eta\phi$	$0.33 \pm 0.01 \pm 0.01 \pm 0.00$	$0.11 \pm 0.00 \pm 0.00 \pm 0.00$
$\eta K\bar{K}$ (non- $\phi$ )	$0.01 \pm 0.01 \pm 0.01 \pm 0.00$	$0.00 \pm 0.00 \pm 0.01 \pm 0.00$
$\omega 3\pi$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.06 \pm 0.01 \pm 0.01 \pm 0.01$	$0.02 \pm 0.00 \pm 0.00 \pm 0.00$
$7\pi$ ( $3\pi^+3\pi^-\pi^0$ + estimate)	$0.02 \pm 0.00 \pm 0.01 \pm 0.00$	$0.01 \pm 0.00 \pm 0.00 \pm 0.00$
<hr/>		
$J/\psi$ (BW integral)	$6.20 \pm 0.11$	$7.00 \pm 0.13$
$\psi(2S)$ (BW integral)	$1.56 \pm 0.05$	$2.48 \pm 0.08$
<hr/>		
$R$ data [3.7 – 5.0] GeV	$7.29 \pm 0.05 \pm 0.30 \pm 0.00$	$15.79 \pm 0.12 \pm 0.66 \pm 0.00$
<hr/>		
$R_{\text{QCD}} [1.8 - 3.7 \text{ GeV}]_{uds}$	$33.45 \pm 0.28 \pm 0.65_{\text{dual}}$	$24.27 \pm 0.18 \pm 0.28_{\text{dual}}$
$R_{\text{QCD}} [5.0 - 9.3 \text{ GeV}]_{udsc}$	$6.86 \pm 0.04$	$34.89 \pm 0.18$
$R_{\text{QCD}} [9.3 - 12.0 \text{ GeV}]_{uds cb}$	$1.20 \pm 0.01$	$15.53 \pm 0.04$
$R_{\text{QCD}} [12.0 - 40.0 \text{ GeV}]_{uds cb}$	$1.64 \pm 0.00$	$77.94 \pm 0.13$
$R_{\text{QCD}} [ > 40.0 \text{ GeV}]_{uds cb}$	$0.16 \pm 0.00$	$42.70 \pm 0.05$
$R_{\text{QCD}} [ > 40.0 \text{ GeV}]_t$	$0.00 \pm 0.00$	$-0.72 \pm 0.01$
<hr/>		
<b>Sum</b>	$694.0 \pm 1.0 \pm 3.5 \pm 1.6 \pm 0.1_\psi \pm 0.7_{\text{QCD}}$	$275.29 \pm 0.15 \pm 0.72 \pm 0.23 \pm 0.15_\psi \pm 0.55_{\text{QCD}}$

→ 32 exclusive channels are integrated up to 1.8 GeV

→ Only  $0.016 \pm 0.016\%$  in missing (estimated) channels for  $a_\mu$

# $R_{e^+e^-} \rightarrow \text{Hadrons}$



Sum of *32 exclusive channels* with  
*full propagation of correlations*

→ Performed non-trivial check:

$a_\mu$  from sum of individual channels and from Ree integral < 1.8 GeV

# Theory initiative: prepare the Standard Model prediction for $(g-2)_\mu$

## UPCOMING WORKSHOPS

Fourth Plenary Workshop of the Muon  $g-2$  Theory Initiative

<https://www-conf.kek.jp/muong-2theory/>

A virtual workshop hosted by KEK (Tsukuba, Japan), to be held from 28 June - 02 July 2021.

## PAST WORKSHOPS

The hadronic vacuum polarization from lattice QCD at high precision

<https://indico.cern.ch/event/956699/>

A virtual topical workshop of the Muon  $g-2$  Theory Initiative, 16-20 Nov 2020.

Hadronic contributions to  $(g-2)_\mu$

<https://indico.fnal.gov/event/21626/>

held at the Institute for Nuclear Theory, University of Washington, Seattle, WA, 9-13 September 2019

Second workshop of the Muon  $g-2$  Theory Initiative

<https://wwwth.kph.uni-mainz.de/g-2/>

held at the Helmholtz Institute Mainz, University of Mainz, Mainz, Germany, 18-22 June 2018

Muon  $g-2$  Theory Initiative Hadronic Light-by-Light working group workshop

<https://indico.phys.uconn.edu/event/1/>

held at the University of Connecticut, Storrs, CT, 12-14 March 2018

Workshop on Hadronic Vacuum Polarization Contributions to Muon  $g-2$

<https://www-conf.kek.jp/muonHVPws/>

held at KEK, Tsukuba, Japan, 12-14 Feb 2018

First workshop of the Muon  $g-2$  Theory Initiative

<https://indico.fnal.gov/event/13795/>

held in St. Charles, IL, USA, 3-6 June 2017

Put together in a *coherent & conservative* way the results of various groups, *before the Fermilab result*

<https://muon-gm2-theory.illinois.edu>

White Paper: arXiv:2006.04822 (Phys. Rept.)

# Theory initiative white paper: executive summary

Contribution	Section	Equation	Value $\times 10^{11}$	References
Experiment (E821)		Eq. (8.13)	116 592 089(63)	Ref. [1]
HVP LO ( $e^+e^-$ )	Sec. 2.3.7	Eq. (2.33)	6931(40)	Refs. [2–7]
HVP NLO ( $e^+e^-$ )	Sec. 2.3.8	Eq. (2.34)	−98.3(7)	Ref. [7]
HVP NNLO ( $e^+e^-$ )	Sec. 2.3.8	Eq. (2.35)	12.4(1)	Ref. [8]
HVP LO (lattice, $udsc$ )	Sec. 3.5.1	Eq. (3.49)	7116(184)	Refs. [9–17]
HLbL (phenomenology)	Sec. 4.9.4	Eq. (4.92)	92(19)	Refs. [18–30]
HLbL NLO (phenomenology)	Sec. 4.8	Eq. (4.91)	2(1)	Ref. [31]
HLbL (lattice, $uds$ )	Sec. 5.7	Eq. (5.49)	79(35)	Ref. [32]
HLbL (phenomenology + lattice)	Sec. 8	Eq. (8.10)	90(17)	Refs. [18–30, 32]
QED	Sec. 6.5	Eq. (6.30)	116 584 718.931(104)	Refs. [33, 34]
Electroweak	Sec. 7.4	Eq. (7.16)	153.6(1.0)	Refs. [35, 36]
HVP ( $e^+e^-$ , LO + NLO + NNLO)	Sec. 8	Eq. (8.5)	6845(40)	Refs. [2–8]
HLbL (phenomenology + lattice + NLO)	Sec. 8	Eq. (8.11)	92(18)	Refs. [18–32]
Total SM Value	Sec. 8	Eq. (8.12)	116 591 810(43)	Refs. [2–8, 18–24, 31–36]
Difference: $\Delta a_\mu := a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$	Sec. 8	Eq. (8.14)	279(76)	

→ Dominant uncertainty: HVP LO → Based on *merging of model-independent methods*

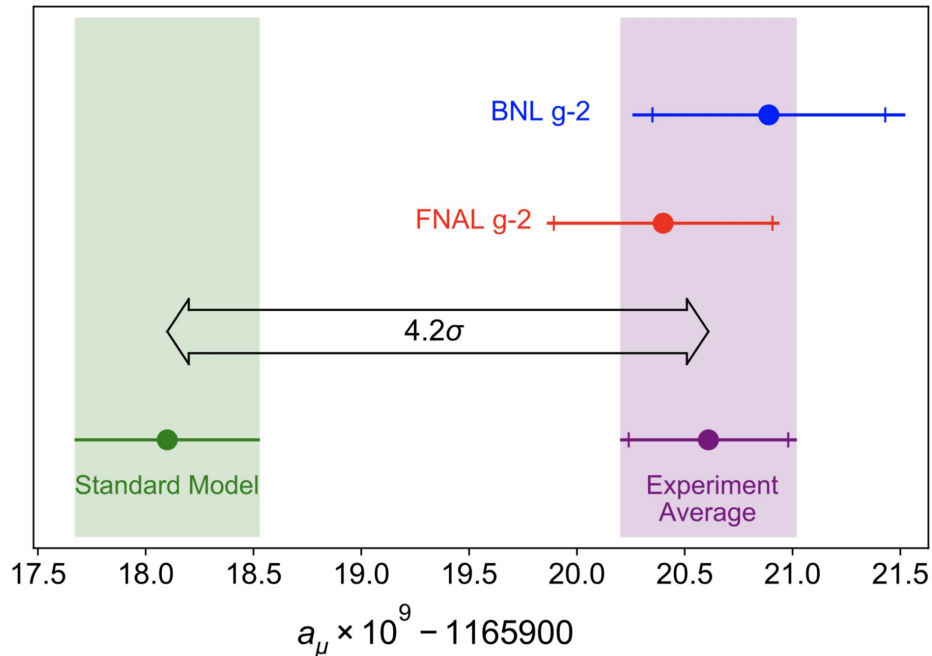
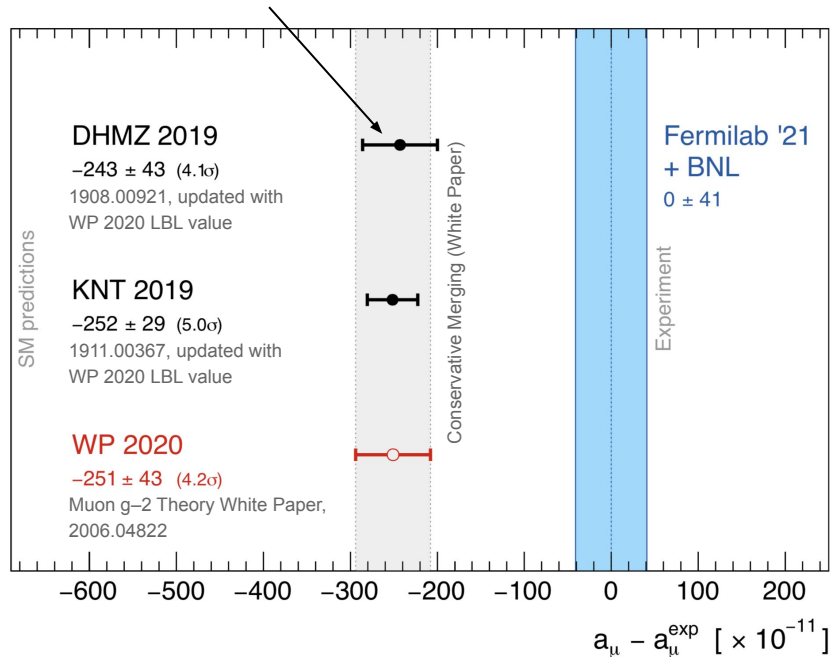
→ HLbL also has an important uncertainty

→ Lattice results become more and more interesting (*see next talk*)

→ A tension between the BNL measurement and SM prediction:  $\sim 3.7 \sigma$

# Status of $a_\mu$

*Important to account for BABAR-KLOE diff. & inter-channel correlations*



→ Caution about significance:

- statistics-dominated measurement
- prediction uncertainty limited by non-Gaussian systematic effects

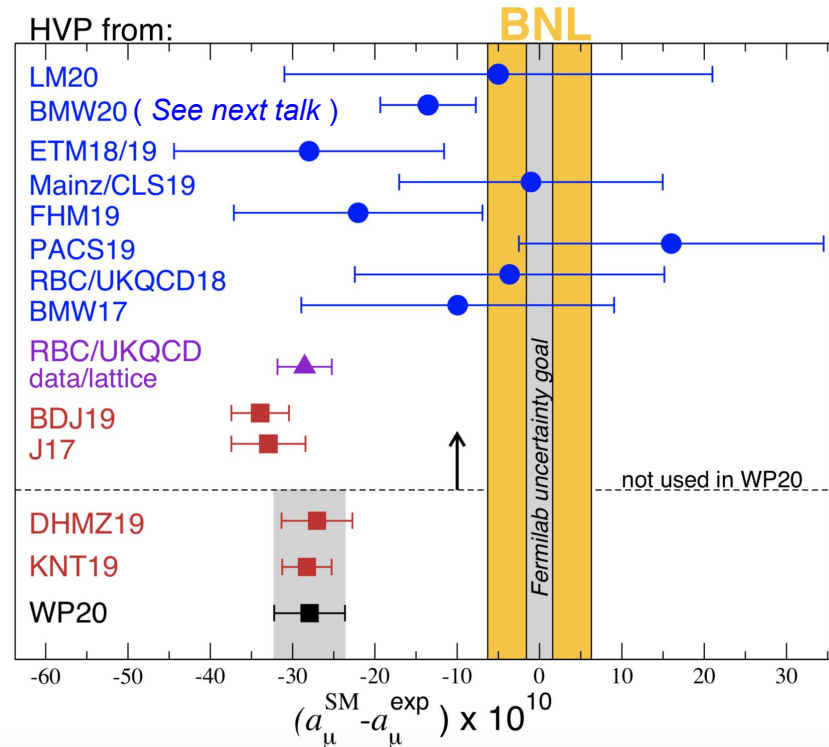
→ Nevertheless, large discrepancy between measurement and SM



# Impact of correlations between $a_\mu$ and $\alpha_{\text{QED}}$ on the EW fit

[2008.08107](#)(BM, Matthias Schott)

**See also:** Crivellin et al, 2003.04886;  
 Keshavarzi et al., 2006.12666 ;de Rafael,  
 2006.13880; Colangelo et al, 2010.07943



# Approaches considered for treating the $a_\mu - \alpha_{\text{QED}}$ correlations

Studied approaches probing different hypotheses concerning the possible source(s) of the  $a_\mu$  tension(s) :

(0) Scaling factor applied to the HVP contribution from some energy range of the hadronic spectrum

→ Approaches taking into account (*for the first time*) the full correlations between the uncertainties of the HVP contributions to  $a_\mu$  and  $\alpha_{\text{QED}}$ , based on input from DHMZ 19 (arXiv:1908.00921):

correlations between points/bins of a measurement in a given channel, between different measurements in the same channel, between different channels; full treatment of the BABAR-KLOE tension in the  $\pi^+\pi^-$  channel

Computation (Energy range)	$a_\mu^{\text{HVP, LO}} [10^{-10}]$	$\Delta\alpha_{\text{had}}(M_Z^2) [10^{-4}]$	$\rho$
Phenomenology (Full HVP)	$694.0 \pm 4.0$	$275.3 \pm 1.0$	44%
Phenomenology ([Th.; 1.8 GeV])	$635.5 \pm 3.9$	$55.4 \pm 0.4$	86%
Phenomenology ([Th.; 1 GeV])	$539.8 \pm 3.8$	$36.3 \pm 0.3$	99.5%
Lattice (Full HVP)BMW 20 (v1)	$712.4 \pm 4.5$	-	-

(1) Cov. matrix of  $a_\mu$  and  $\alpha_{\text{QED}}$  (Pheno) described by a nuisance parameter (NP<sub>1</sub>) impacting both quantities (used to shift  $a_\mu$  to some “target” value - coherent shift applied to  $\alpha_{\text{QED}}$ ) and another one (NP<sub>2</sub>) impacting only  $\alpha_{\text{QED}}$  (used in the EW fit)

Note: “target” values chosen in order to reach agreement with the BMW 20 prediction / Experimental  $a_\mu (\pm 1\sigma)$

Uncertainty components	$a_\mu^{\text{HVP, LO}}$	$\Delta\alpha_{\text{had}}(M_Z^2)$
NP <sub>1</sub>	$\sigma(a_\mu^{\text{HVP, LO}})$	$\sigma(\Delta\alpha_{\text{had}}(M_Z^2)) \cdot \rho$
NP <sub>2</sub>	0	$\sigma(\Delta\alpha_{\text{had}}(M_Z^2)) \cdot \sqrt{1 - \rho^2}$

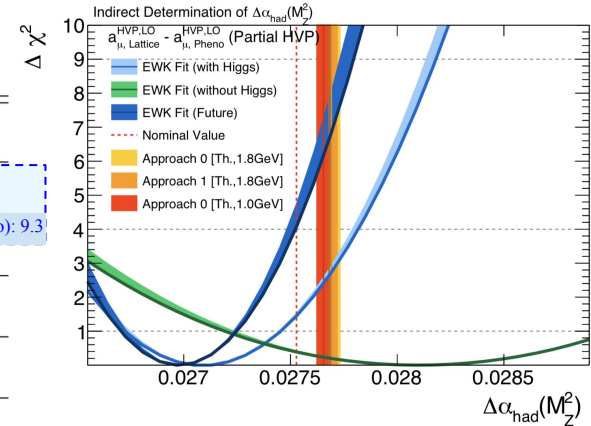
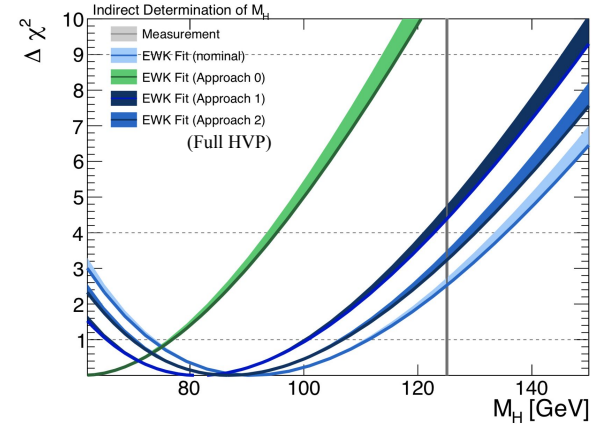
(2) Include the HVP contribution to  $a_\mu$  as extra parameter in the EW fit, constrained by the Pheno & BMW 20 values

Note: Also accounted for the coherent impact of  $\alpha_s$  on the HVP contribution and on the EW fit

# Results: comparing the Phenomenology & BMW 20 values

$a_\mu^{\text{HVP, LO}}$ shift (Energy range)	Approach 0		Approach 1		
	Scaling factor	$\Delta' \alpha_{\text{had}}(M_Z^2)$	Shift NP <sub>1</sub>	$\sigma'(\Delta \alpha_{\text{had}}(M_Z^2))$	$\Delta' \alpha_{\text{had}}(M_Z^2)$
$a_\mu^{\text{HVP, LO}}(\text{Lattice}) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ (Full HVP)	1.027	0.02826	4.6	$9.0 \cdot 10^{-5}$	0.02774
$(a_\mu^{\text{HVP, LO}}(\text{Lattice}) - 1\sigma) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ (Full HVP)	1.020	0.02808	3.5	$9.0 \cdot 10^{-5}$	0.02769
$a_\mu^{\text{HVP, LO}}(\text{Lattice}) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ ([Th.; 1.8 GeV])	1.029	0.02769	4.7	$9.5 \cdot 10^{-5}$	0.02768
$(a_\mu^{\text{HVP, LO}}(\text{Lattice}) - 1\sigma) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ ([Th.; 1.8 GeV])	1.022	0.02765	3.5	$9.5 \cdot 10^{-5}$	0.02764
$a_\mu^{\text{HVP, LO}}(\text{Lattice}) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ ([Th.; 1 GeV])	1.034	0.02765	-	-	-
$(a_\mu^{\text{HVP, LO}}(\text{Lattice}) - 1\sigma) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ ([Th.; 1 GeV])	1.026	0.02762	-	-	-

→ Large scaling factors (w.r.t. exp. uncertainties) & significant shifts of NP<sub>1</sub>



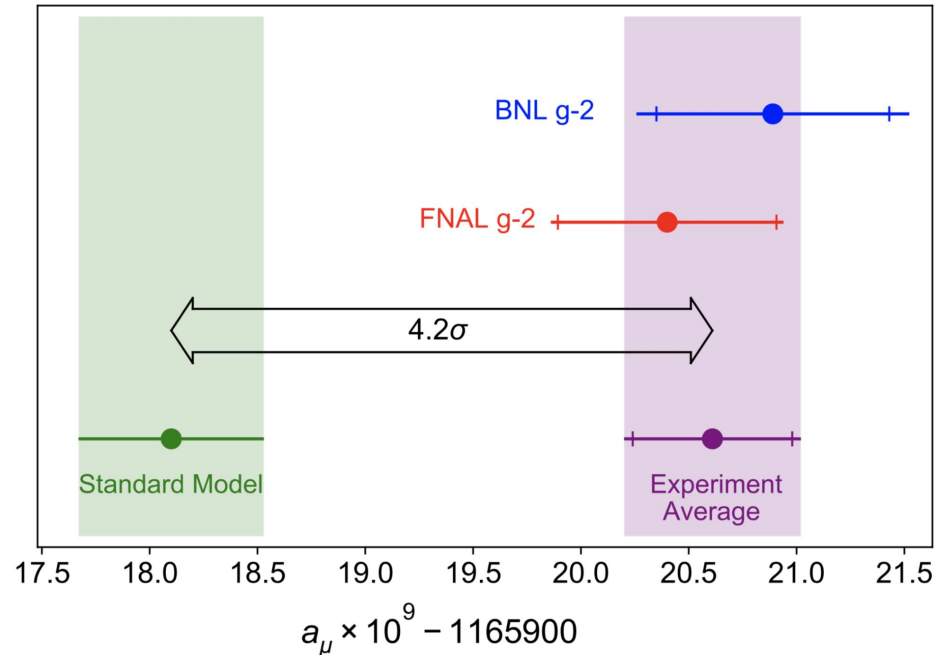
$a_\mu^{\text{HVP, LO}}$ shift (Energy range)	Nominal		Approach 0		Approach 1		Approach 2	
	$\Delta' \alpha_{\text{had}}(M_Z^2)$	$\chi^2/\text{ndf}$	$\Delta' \alpha_{\text{had}}(M_Z^2)$	$\chi^2/\text{ndf}$	$\Delta' \alpha_{\text{had}}(M_Z^2)$	$\chi^2/\text{ndf}$	$\Delta' \alpha_{\text{had}}(M_Z^2)$	$\chi^2/\text{ndf}$
	0.02753	18.6/16 (p=0.29)	-	-	-	-	0.02753	28.1/17 (p=0.04)
$a_\mu^{\text{HVP, LO}}(\text{Lattice}) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ (Full HVP)	-	-	0.02826	27.6/16 (p=0.04)	0.02774	20.3/16 (p=0.21)	-	$\chi^2(\text{BMW20-Pheno}): 9.3$
$a_\mu^{\text{HVP, LO}}(\text{Lattice}) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ ([Th.; 1.8 GeV])	-	-	0.02769	19.9/16 (p=0.22)	0.02768	19.8/16 (p=0.23)	-	-
$a_\mu^{\text{HVP, LO}}(\text{Lattice}) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ ([Th.; 1.0 GeV])	-	-	0.02765	19.6/16 (p=0.24)	-	-	-	-

→ Addressing the BMW 20 - Pheno difference for  $a_\mu$  has little impact on the EW fit, except for the unrealistic scenario rescaling the full HVP contribution

Note: Similar conclusions for the comparison with the Experimental  $a_\mu$  value (see backup)

# Conclusion

We have an interesting, long standing, multifaceted problem to solve...



???

# Backup

# Lepton Magnetic Anomaly: from Dirac to QED

$$\vec{\mu} = g \frac{e}{2m} \vec{s} \quad a = \frac{g-2}{2}$$

Dirac (1928)  $g_e=2$   $a_e=0$

anomaly discovered:

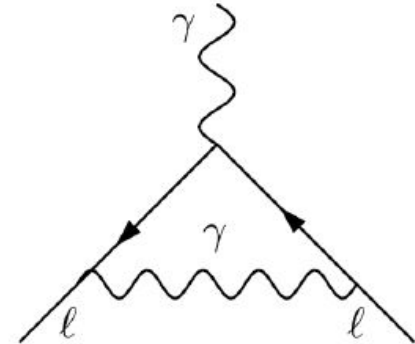
Kusch-Foley (1948)  $a_e = (1.19 \pm 0.05) 10^{-3}$

and explained by  $O(\alpha)$  QED contribution:

Schwinger (1948)  $a_e = \alpha/2\pi = 1.16 10^{-3}$

first triumph of QED

$\Rightarrow a_e$  sensitive to quantum fluctuations of fields



# More Quantum Fluctuations

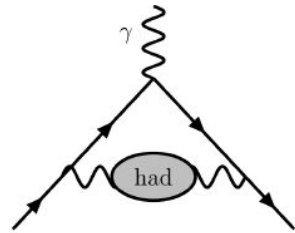
$$a = a^{\text{QED}} + a^{\text{had}} + a^{\text{weak}} + ? a^{\text{new physics}} ?$$

typical contributions:

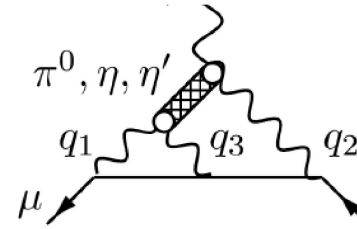
QED up to  $O(\alpha^5)$  (Kinoshita et al.)



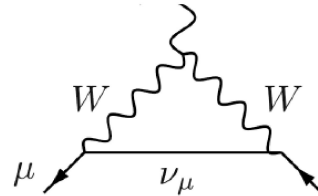
Hadrons vacuum polarization



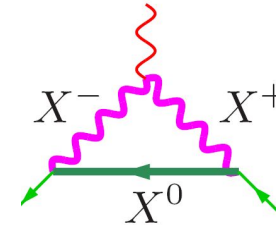
light-by-light (dispersive & lattice QCD)



Electroweak

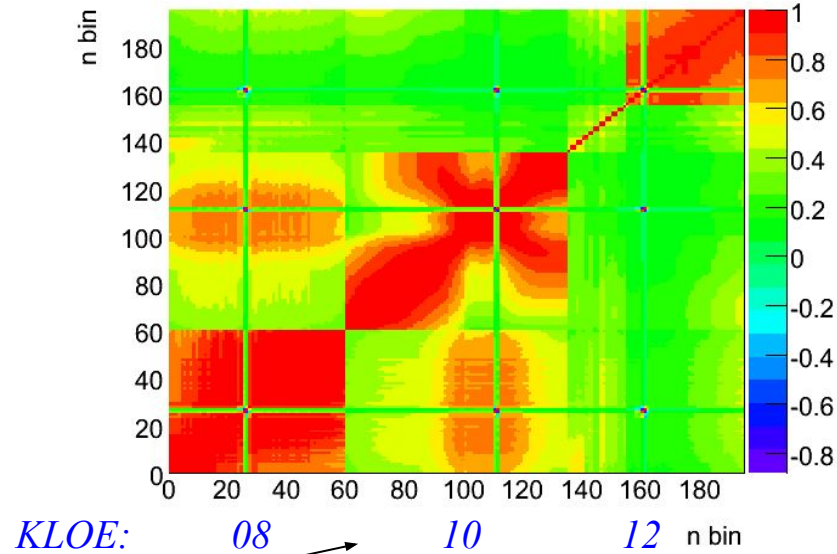
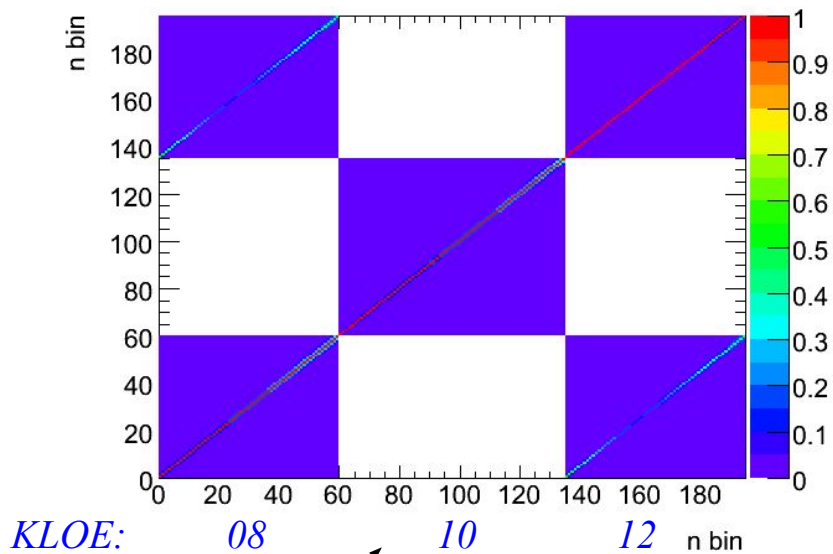


new physics at high mass scale



$$\delta a_l \propto \frac{m_l^2}{M^2} \Rightarrow a_\mu \text{ much more sensitive to high scales}$$

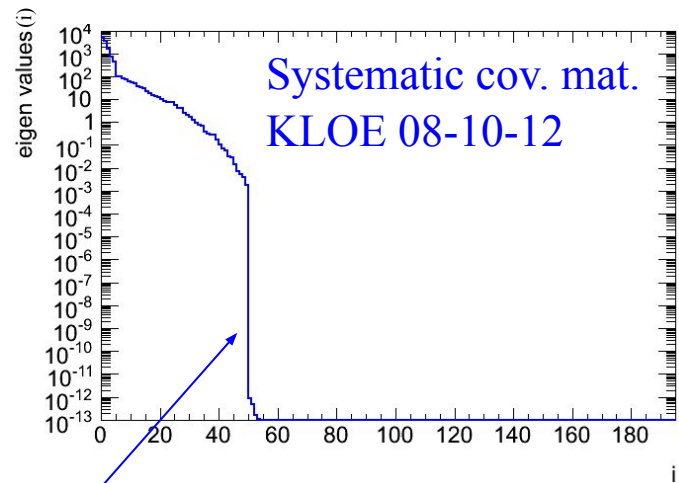
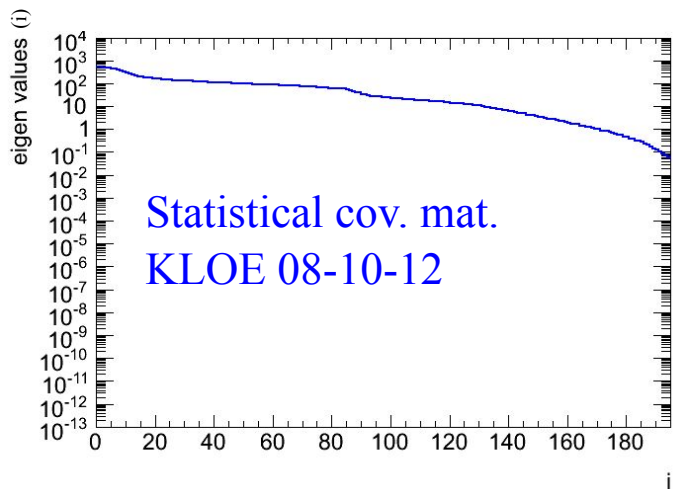
# Treatment of the KLOE correlation matrices



→ Statistical and systematic correlation matrices among the 3 measurements



# Treatment of the KLOE data – eigenvector decomposition



→ “counting” the number of independent components (50) used to build the covariance matrix

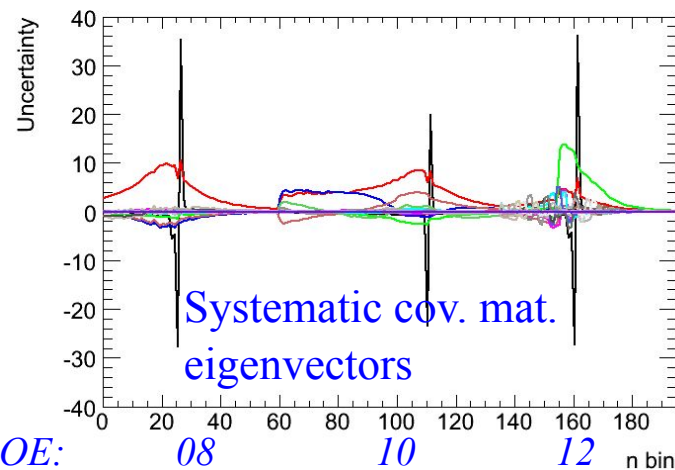
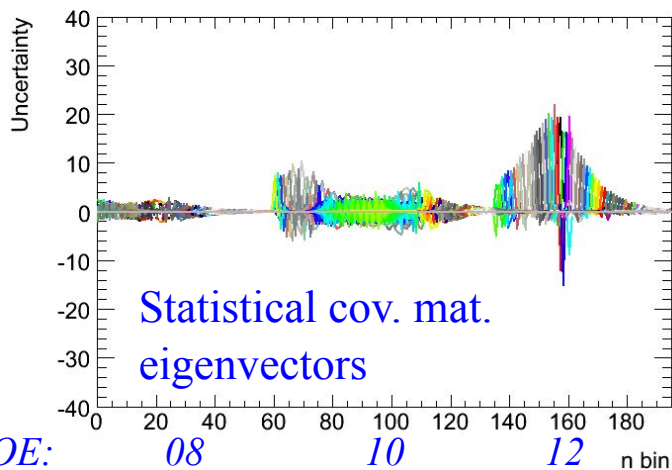
$$C = S \cdot D \cdot S^T$$

$$D = \begin{pmatrix} \diagdown & 0 & 0 \\ 0 & \sigma_i^2 & 0 \\ 0 & 0 & \diagdown \end{pmatrix}$$

$$S = \begin{pmatrix} V_1 & \dots & V_n \\ \vdots & & \vdots \end{pmatrix}$$

→ Problem of negative eigenvalues for previous systematic covariance matrix solved (informed KLOE collaboration about the problem in summer 2016)

# Treatment of the KLOE data – eigenvector decomposition

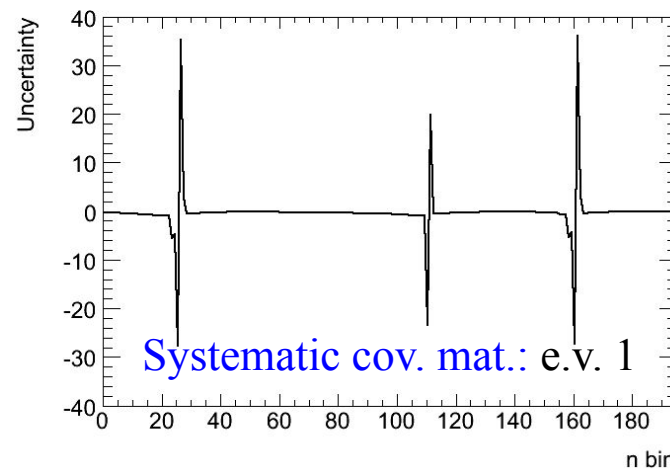
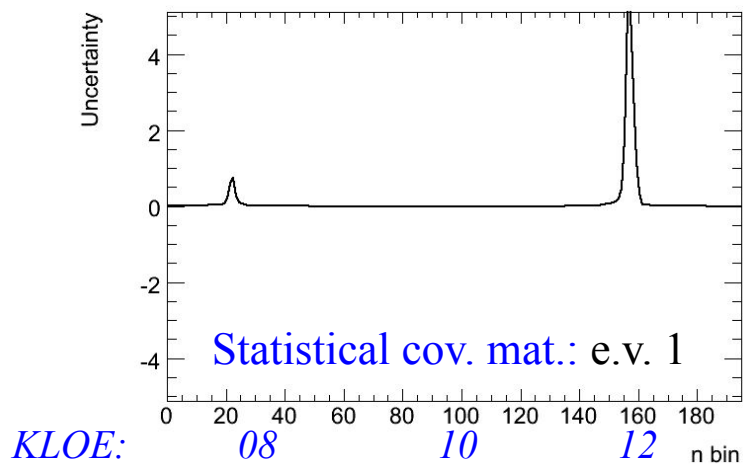


- Each normalized eigenvector ( $\sigma_i \cdot V_i$ ) treated as an uncertainty fully correlated between the bins
- All these uncertainties are independent between each-other

$$C = \sum_{i=1}^{N_{bins}} \sigma_i^2 \cdot C(V_i)$$

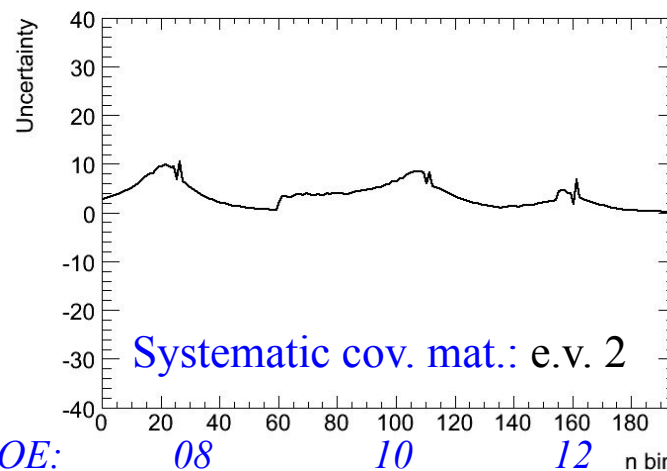
- Checked exact matching with the original matrices + with all  $a_\mu$  integrals and uncertainties published by KLOE

# Treatment of the KLOE data – eigenvector decomposition

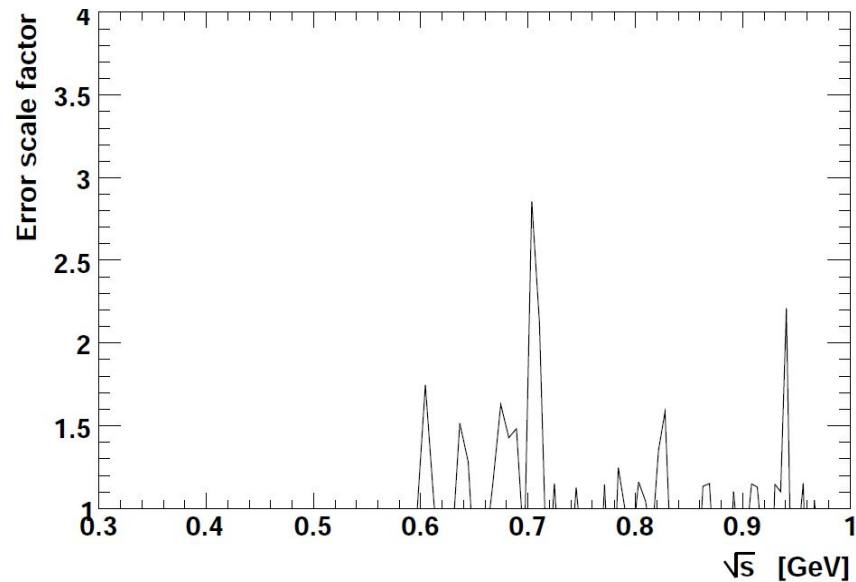


→ Eigenvectors carry the general features of the correlations:

- long-range for systematics
- ~short-range for statistical uncertainties + correlations between KLOE 08 & 12



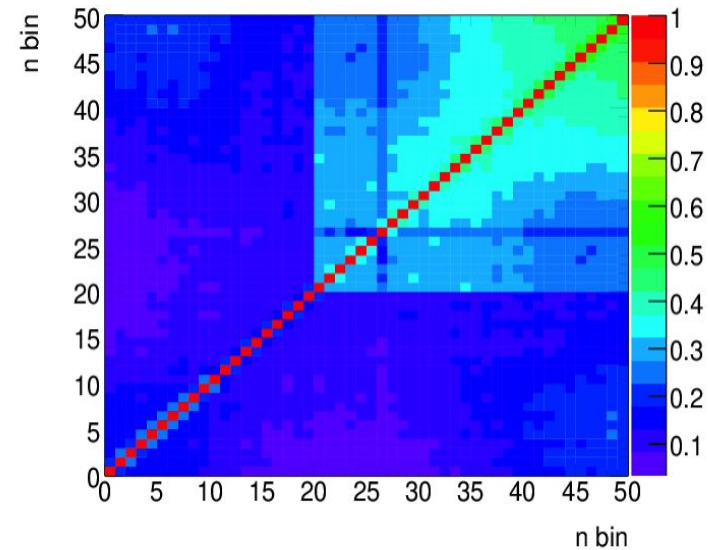
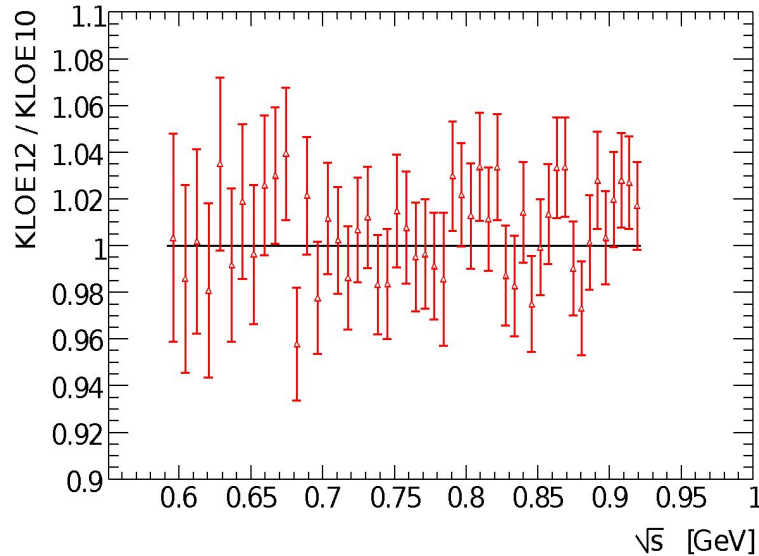
# Local comparison of the 3 KLOE measurements



- Local  $\chi^2$ /ndof test of the local compatibility between KLOE 08 & 10 & 12, taking into account the correlations: some tensions observed
- Does not probe general trends of the difference between the measurements (e.g. slopes in the ratio)

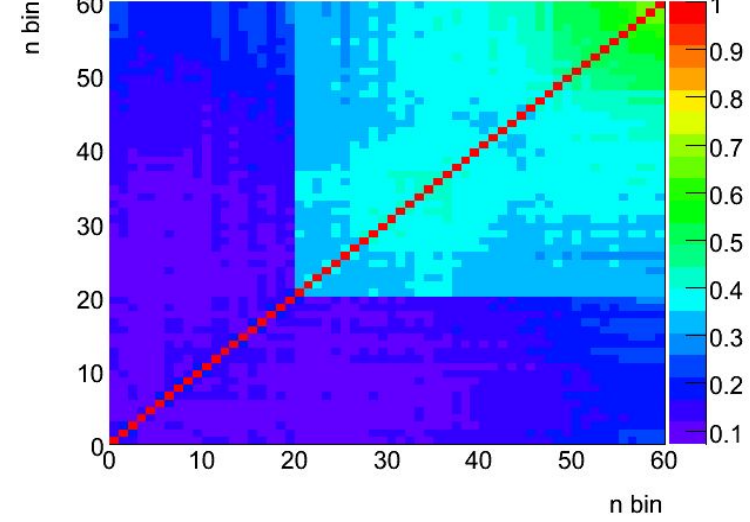
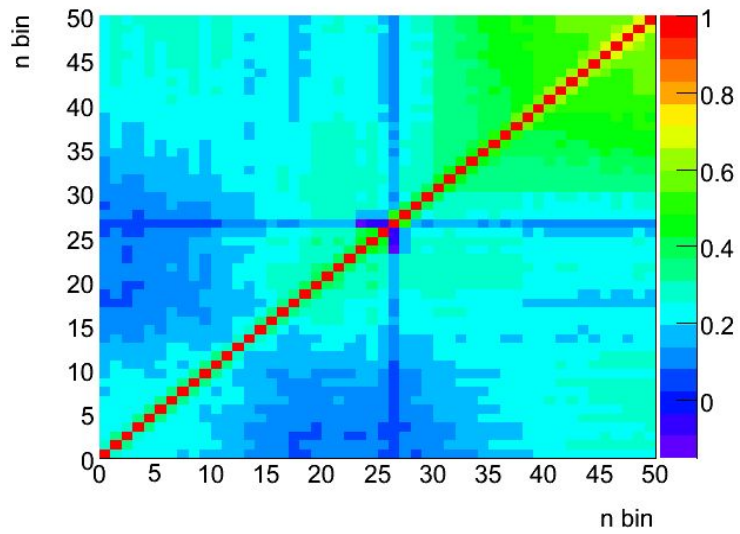
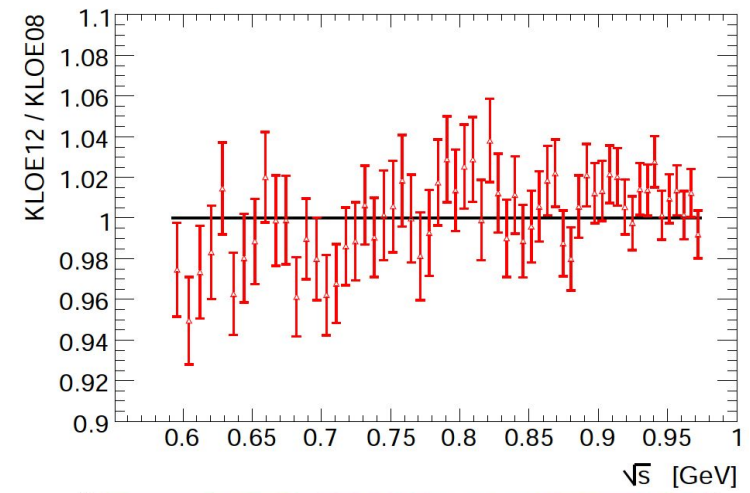
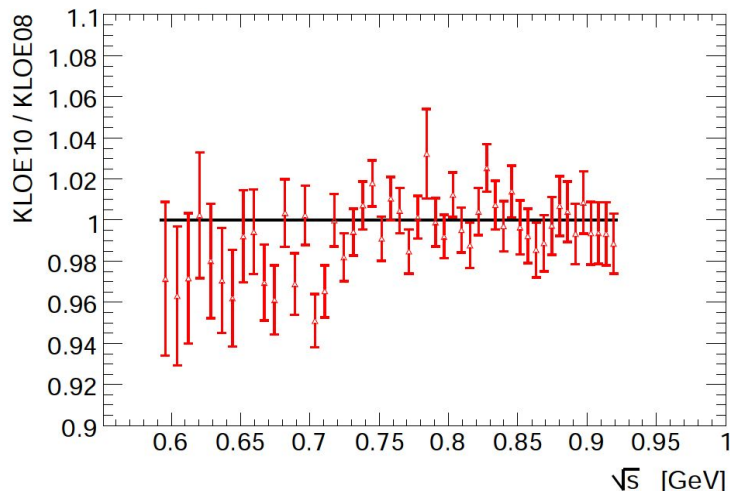
# Ratios between measurements

- Compute ratio between pairs of KLOE measurements
- Full propagation of uncertainties and correlations using pseudo-experiments (agreement with analytical linear uncertainty propagation)



→ Good agreement between KLOE 10 and KLOE 12

# Ratios between measurements



# Direct comparison of the 3 KLOE measurements

→ Quantitative comparison between the ratios and unity, taking into account correlations

## KLOE 10 / KLOE 08

$$\chi^2 [0.35;0.85] \text{ GeV}^2 : 79.0 / 50(\text{DOF})$$

p-value= 0.0056

$$\chi^2 [0.35;0.58] \text{ GeV}^2 : 46.2 / 23(\text{DOF})$$

p-value= 0.0028

$$\chi^2 [0.58;0.85] \text{ GeV}^2 : 29.7 / 27(\text{DOF})$$

p-value= 0.33

$$\chi^2 [0.64;0.85] \text{ GeV}^2 : 20.7 / 21(\text{DOF})$$

p-value= 0.47

## KLOE 12 / KLOE 08

$$\chi^2 [0.35;0.95] \text{ GeV}^2 : 73.7 / 60(\text{DOF})$$

p-value= 0.11

$$\chi^2 [0.35;0.58] \text{ GeV}^2 : 21.8 / 23(\text{DOF})$$

p-value= 0.53

$$\chi^2 [0.35;0.64] \text{ GeV}^2 : 27.5 / 29(\text{DOF})$$

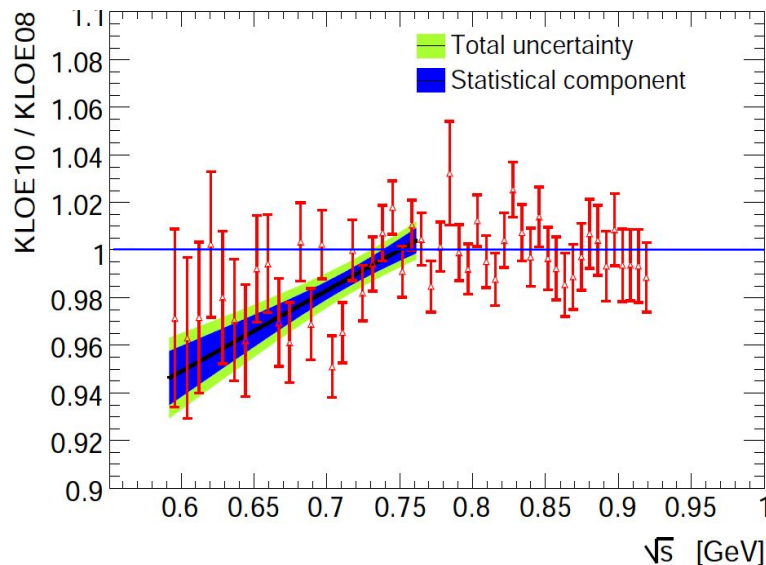
p-value= 0.55

$$\chi^2 [0.64;0.95] \text{ GeV}^2 : 39.4 / 31(\text{DOF})$$

p-value= 0.14

# Quantitative comparisons of the KLOE measurements

- Quantitative comparison between the ratios and unity, taking into account correlations
- Fitting the ratio taking into account correlations
- Full propagation of uncertainties and correlations – 3 methods yielding consistent results:  $\pm 1\sigma$  shifts of each uncertainty, pseudo-experiments and fit uncertainties from Minuit



Comparison with Unity:

$$\chi^2 [0.35; 0.85] \text{ GeV}^2 : 79.0 / 50(\text{DOF})$$

$$p\text{-value} = 0.0056$$

$$\chi^2 [0.35; 0.58] \text{ GeV}^2 : 46.2 / 23(\text{DOF})$$

$$p\text{-value} = 0.0028$$

$$\chi^2 [p_0 + p_1\sqrt{s}] : 36.1 / 21(\text{DOF})$$

$$p\text{-value} = 0.02$$

$$p_0 : 0.745 \pm 0.085$$

$$p_1 : 0.341 \pm 0.117$$

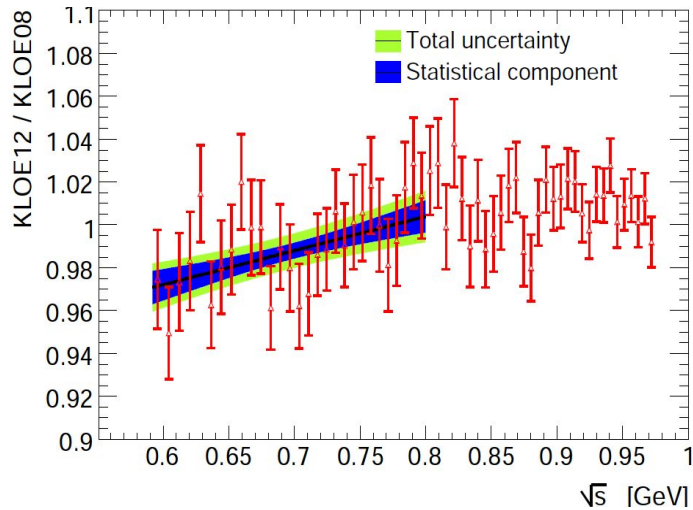
- Significant shift & slope ( $\sim 2.5\text{-}3\sigma$ ) at low  $\sqrt{s}$ , no significant shift at high  $\sqrt{s}$   
Similar shift & slope for KLOE 12 / KLOE 08 (*see below*)
- Should motivate conservative treatment of uncertainties and correlations in combination



# Direct comparison of the 3 KLOE measurements

→ Fitting the ratio taking into account correlations

→ Full propagation of uncertainties and correlations – 3 methods yielding consistent results:  
 $\pm 1\sigma$  shifts of each uncertainty, pseudo-experiments and fit uncertainties from Minuit

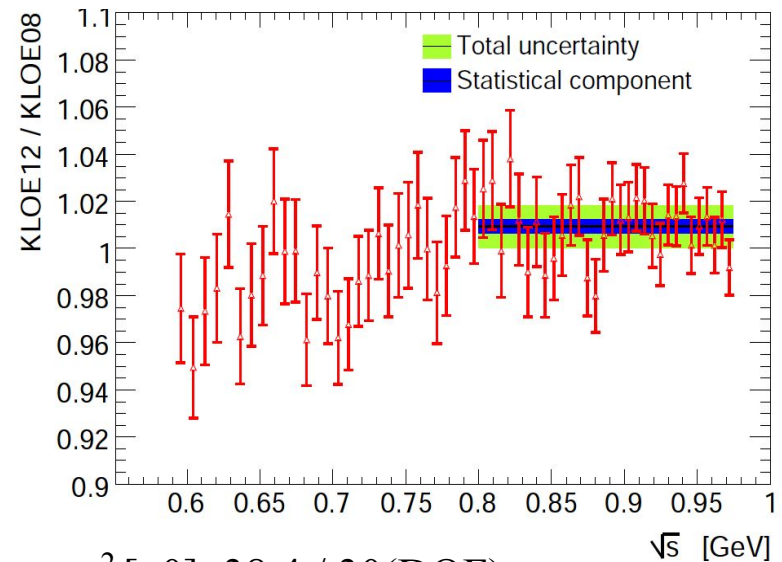


$$\chi^2 [p0 + p1\sqrt{s}]: 20.7 / 27(\text{DOF})$$

$$p\text{-value} = 0.80$$

$$p0 : 0.876 \pm 0.056$$

$$p1 : 0.159 \pm 0.081$$



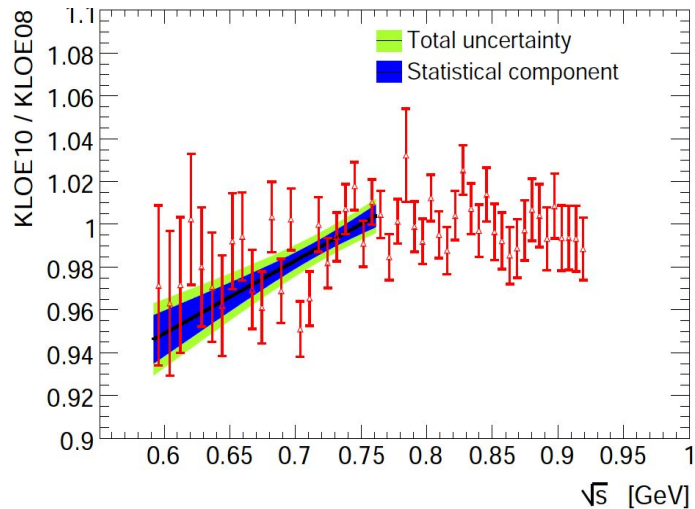
$$\chi^2 [p0]: 38.4 / 30(\text{DOF})$$

$$p\text{-value} = 0.14$$

$$p0 : 1.009 \pm 0.009$$

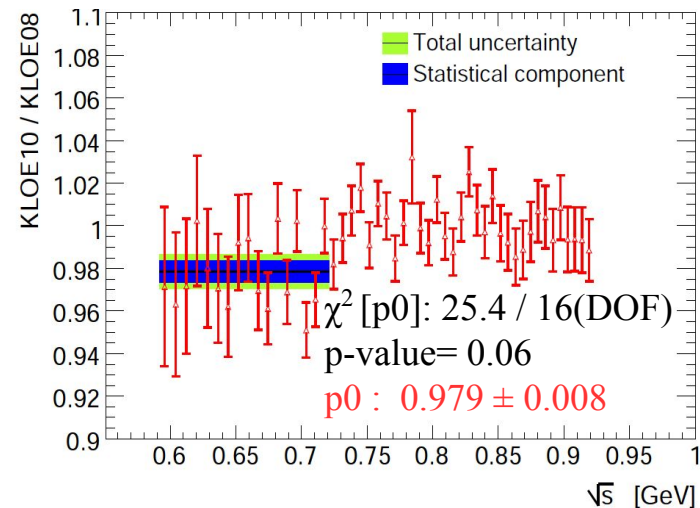
→ Significant shift and slope ( $\sim 2\sigma$ ) at low  $\sqrt{s}$ , no significant shift at high  $\sqrt{s}$

# Direct comparison of the 3 KLOE measurements

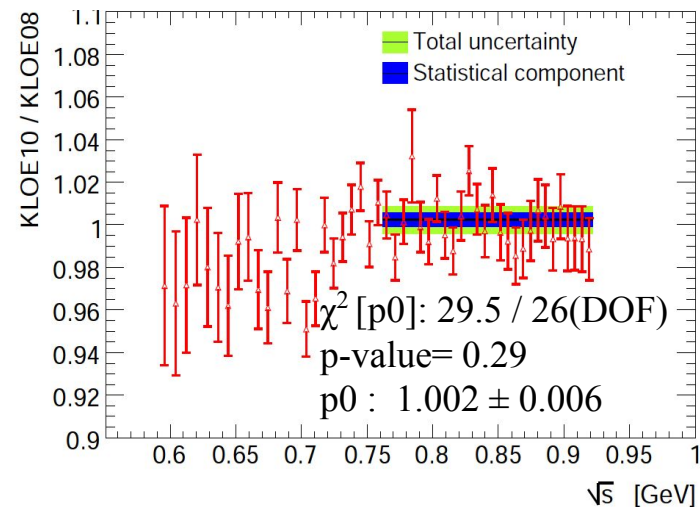


$\chi^2$  [p0 + p1 $\sqrt{s}$ ]: 36.1 / 21(DOF)  
 p-value= 0.02  
 p0 : 0.745  $\pm$  0.085  
 p1 : 0.341  $\pm$  0.117

→ Significant shift and slope ( $\sim 2.5$ - $3\sigma$ ) at low  $\sqrt{s}$ ,  
 no significant shift at high  $\sqrt{s}$

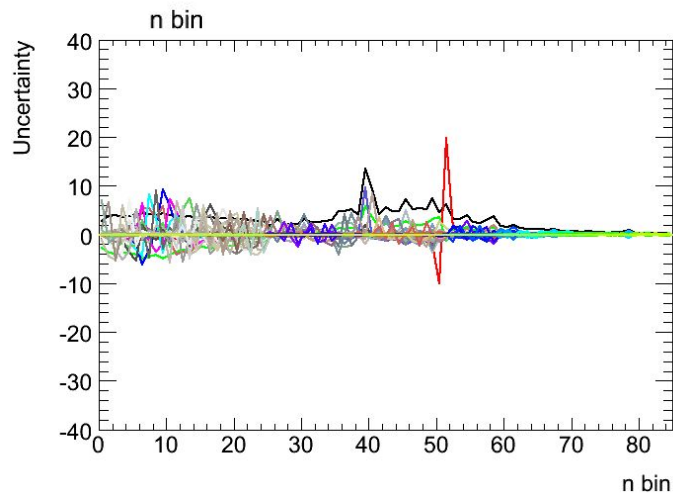
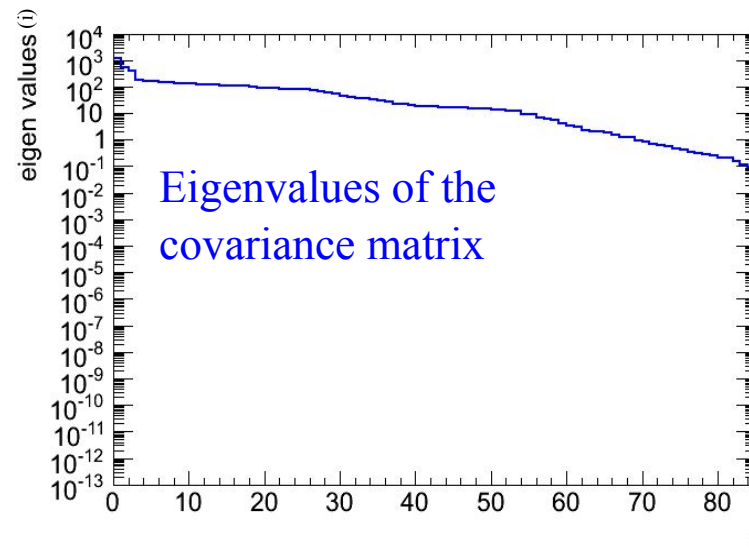
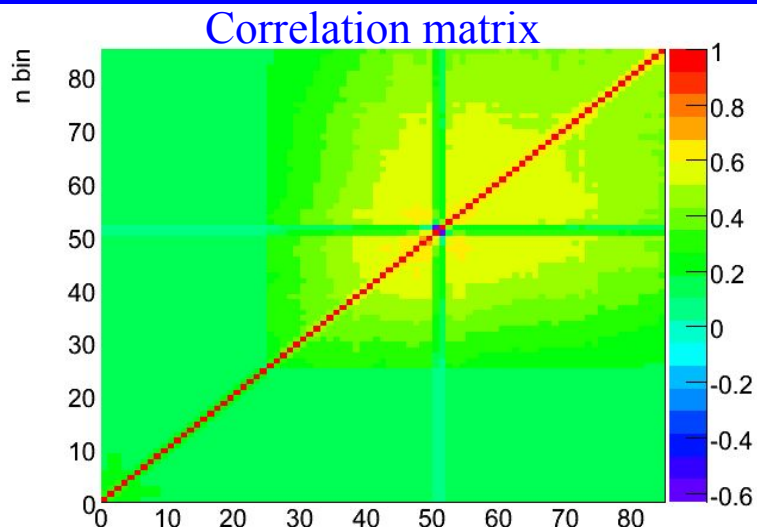


$\chi^2$  [p0]: 25.4 / 16(DOF)  
 p-value= 0.06  
 p0 : 0.979  $\pm$  0.008

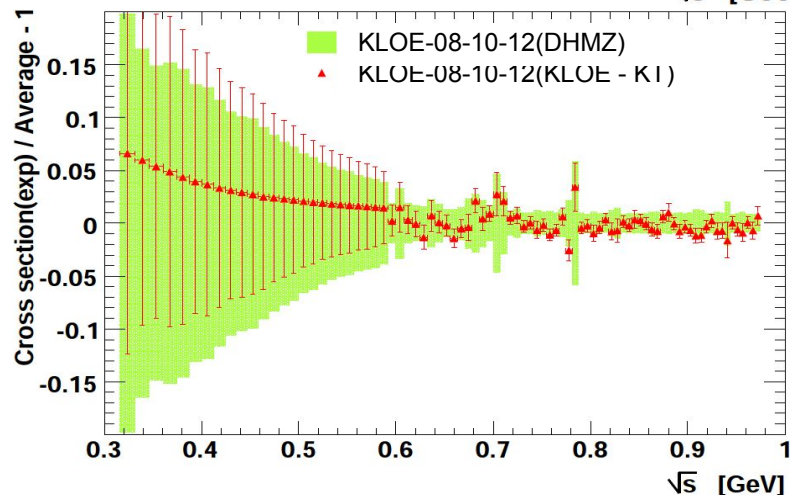
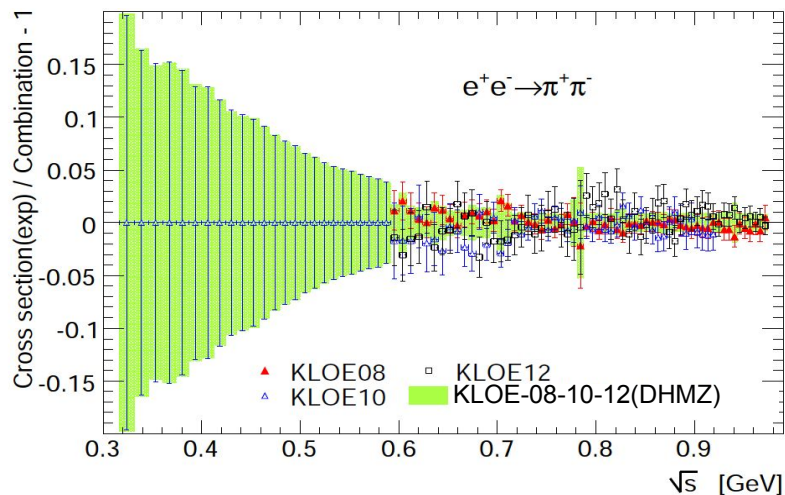
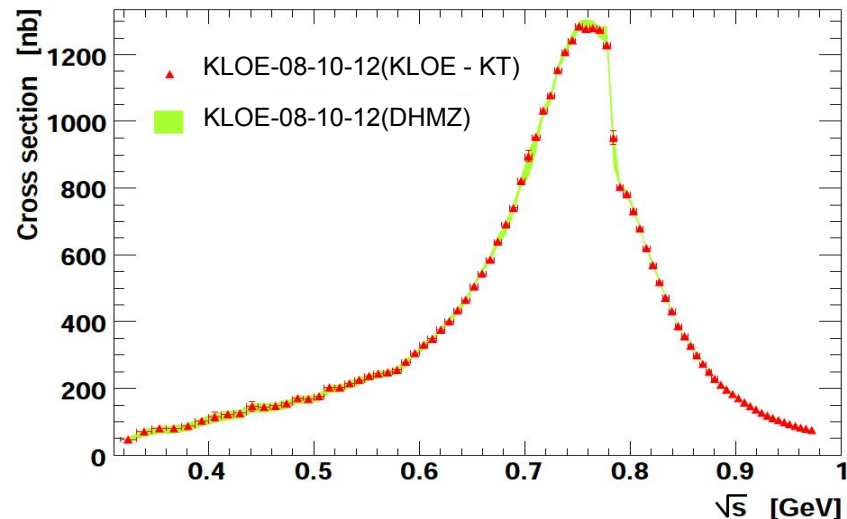
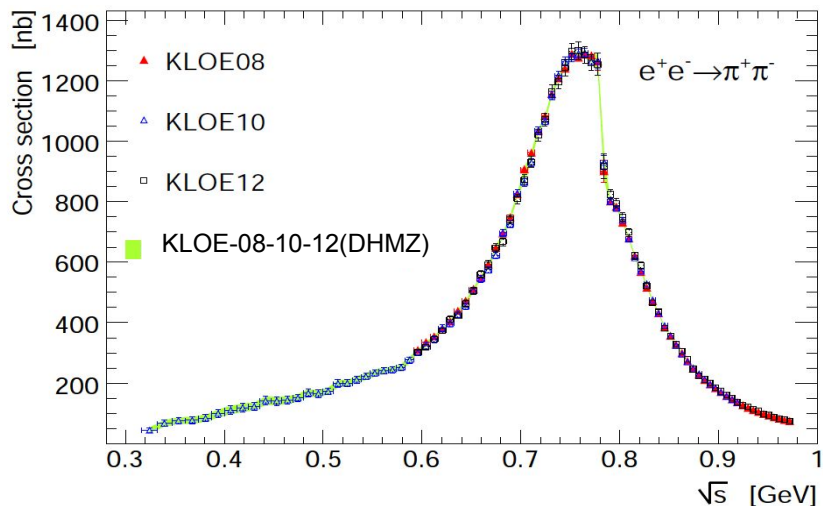


$\chi^2$  [p0]: 29.5 / 26(DOF)  
 p-value= 0.29  
 p0 : 1.002  $\pm$  0.006

# Treatment of the combined KLOE data



# Combining the 3 KLOE measurements



# $a_{\mu}^{\pi\pi}$ contribution [0.28; 1.8] GeV – spline-based (2018)

→ Updated result:

$$506.70 \pm 2.32 \left( \pm 1.01 \text{ (stat.)} \pm 2.08 \text{ (syst.)} \right) [10^{-10}]$$

(after uncertainty enhancement by  $\sim 14\%$  caused by the tension between inputs, taken into account through a local rescaling)

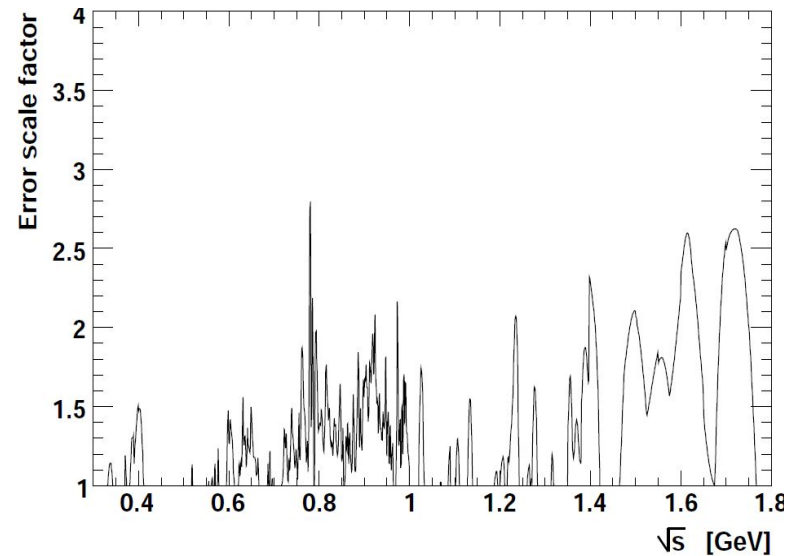
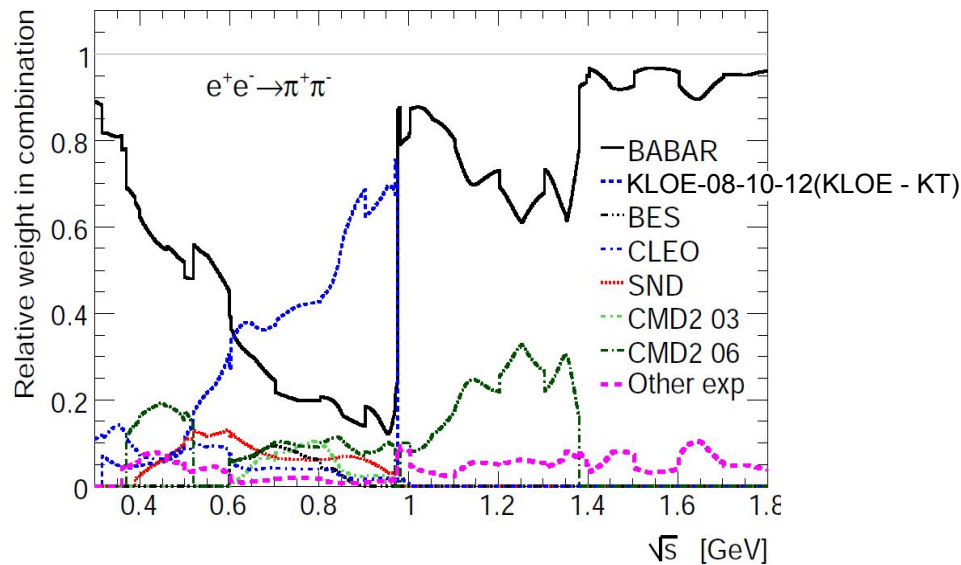
Total uncertainty: 5.9 (2003)  $\rightarrow$  2.8 (2011)  $\rightarrow$  2.6 (2017)  $\rightarrow$  2.3 (2018)

# $a_\mu^{\pi\pi}$ contribution [0.28; 1.8] GeV – spline-based (2018)

→ with KLOE-08-10-12 (KLOE-KT) used as input:  $506.55 \pm 2.38 [10^{-10}]$

(after uncertainty enhancement by 18% caused by the tension between inputs, taken into account through a local rescaling)

→ Compensation between uncertainty reduction for KLOE-08-10-12 (KLOE-KT), inducing a change of weights in DHMZ combination, and tension enhancement



# Fit parameters, uncertainties and correlations $e^+e^- \rightarrow \pi^+\pi^-$

	$\alpha_V$	$\kappa[10^{-4}]$	$B_0$	$B_1$	$m_\rho$ [MeV]	$m_\omega$ [MeV]
$\alpha_V$	$0.133 \pm 0.020$	0.52	-0.45	-0.97	0.90	-0.25
$\kappa[10^{-4}]$		$21.6 \pm 0.5$	-0.33	-0.57	0.64	-0.08
$B_0$			$1.040 \pm 0.003$	0.40	-0.40	0.29
$B_1$				$-0.13 \pm 0.11$	-0.96	0.20
$m_\rho$ [MeV]					$774.5 \pm 0.8$	-0.17
$m_\omega$ [MeV]						$782.0 \pm 0.1$

→  $\kappa$  corresponds to a Br ( $\omega \rightarrow \pi^+\pi^-$ ) of  $(2.09 \pm 0.09) \cdot 10^{-2}$ , in agreement with the result extracted from the fit of arXiv:1810.00007,  $(1.95 \pm 0.08) \cdot 10^{-2}$ . Both values disagree with the PDG average  $(1.51 \pm 0.12) \cdot 10^{-2}$ , dominated by the result of arXiv:1611.09359 which uses fits to essentially the same data.

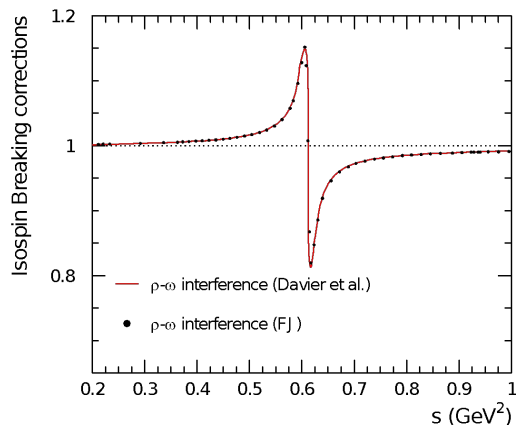
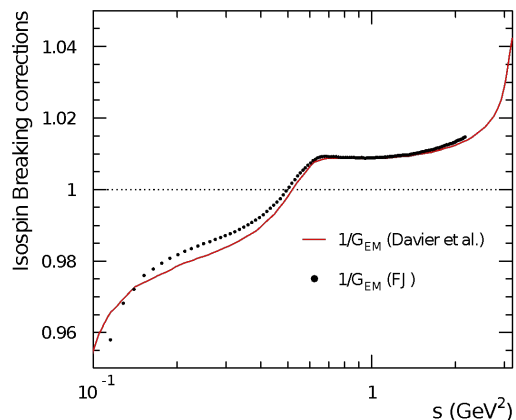
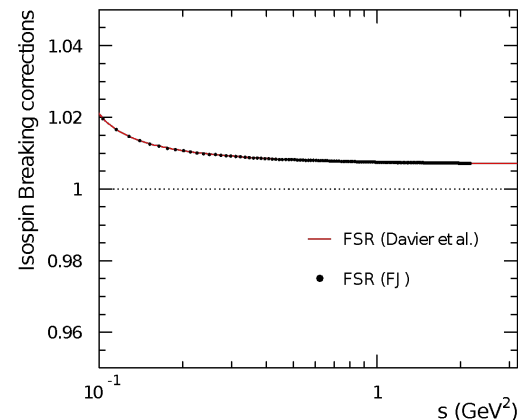
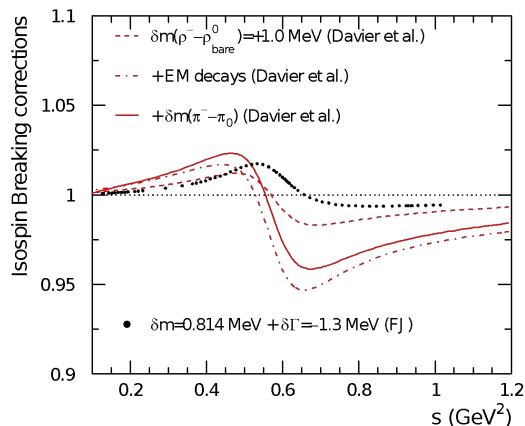
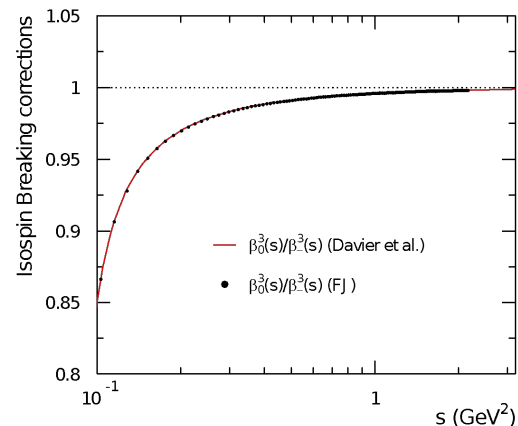
→ The fitted  $\omega$  mass is found to be lower than the PDG average obtained from  $3\pi$  decays by  $(0.65 \pm 0.12 \pm 0.12_{\text{PDG}})$  MeV, in agreement with previous fits of the  $\rho - \omega$  interference in the  $2\pi$  spectrum (see e.g. arXiv:1205.2228 and arXiv:1810.00007).

# Comparison with IB-corrected $\tau$ data

$$v_{1,X^-}(s) = \frac{m_\tau^2}{6|V_{ud}|^2} \frac{\mathcal{B}_{X^-}}{\mathcal{B}_e} \frac{1}{N_X} \frac{dN_X}{ds} \times \left(1 - \frac{s}{m_\tau^2}\right)^2 \left(1 + \frac{2s}{m_\tau^2}\right)^{-1} \frac{R_{\text{IB}}(s)}{S_{\text{EW}}}$$

$$R_{\text{IB}}(s) = \frac{\text{FSR}(s)}{G_{\text{EM}}(s)} \frac{\beta_0^3(s)}{\beta_-^3(s)} \left| \frac{F_0(s)}{F_-(s)} \right|^2$$

→ Comparing corrections used by Davier et al. with the ones by F. Jegerlehner



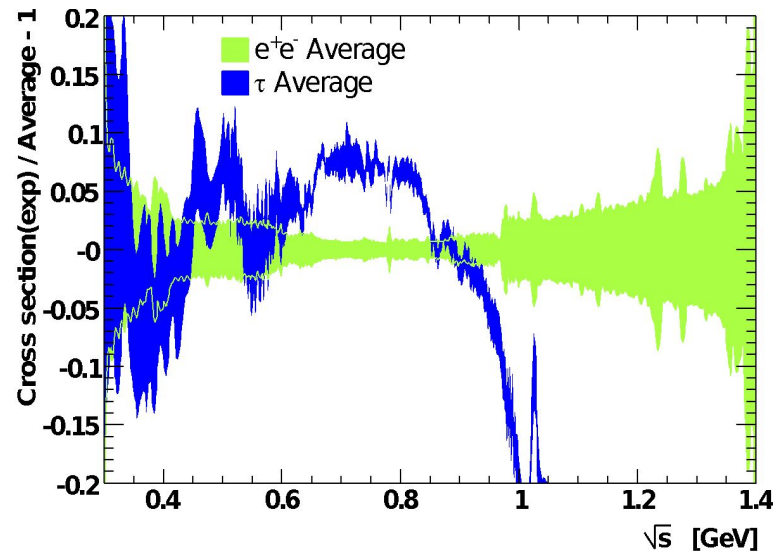
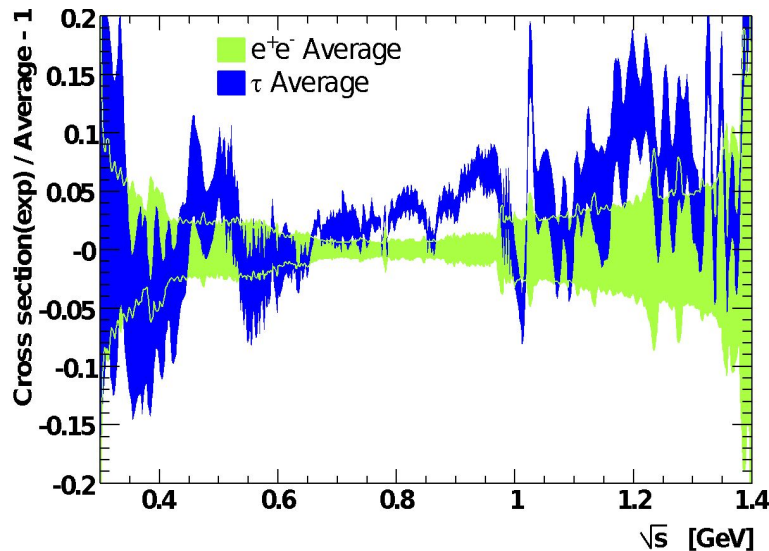
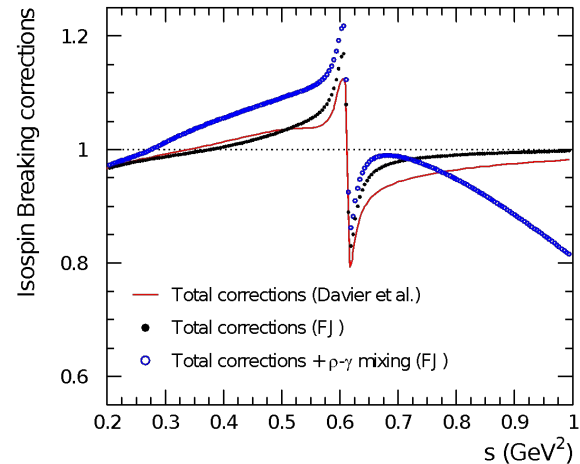
Plots by Z. Zhang, based on private communication with F. Jegerlehner



# Comparison with IB-corrected $\tau$ data

→ for  $a_{\mu}^e$ ,  $e^+e^- - \tau$  difference of  $2.2 \sigma$   
(Davier et al.)

→ the  $\rho$ - $\gamma$  mixing correction proposed in  
arXiv:1101.2872 (FJ) seems to over-estimate  
the  $e^+e^- - \tau$  difference



# $\chi^2$ definitions and properties

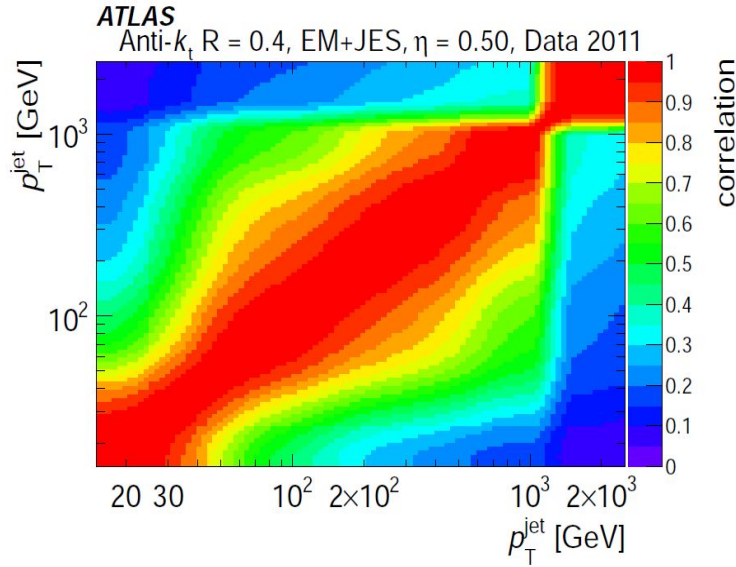
$$\chi^2(\mathbf{d}; \mathbf{t}) = \sum_{i,j} (d_i - t_i) \cdot [C^{-1}(\mathbf{t})]_{ij} \cdot (d_j - t_j) \quad C_{ij} = C_{ij}^{stat} + \sum_k s_i^k \cdot s_j^k$$

$$\chi^2(\mathbf{d}; \mathbf{t}) = \min_{\beta_a} \left\{ \sum_{i,j} \left[ d_i - \left( 1 + \sum_a \beta_a \cdot (\epsilon_a^\pm(\beta_a))_i \right) t_i \right] \cdot [C_{su}^{-1}(\mathbf{t})]_{ij} \cdot \left[ d_j - \left( 1 + \sum_a \beta_a \cdot (\epsilon_a^\pm(\beta_a))_j \right) t_j \right] + \sum_a \beta_a^2 \right\},$$

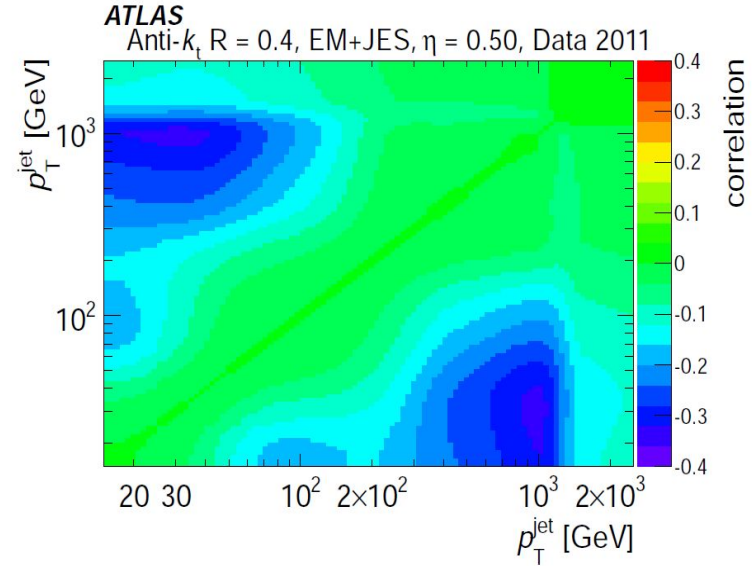
- Two  $\chi^2$  definitions, with systematic uncertainties included in covariance matrix or treated as fitted “nuisance parameters”
- Equivalent for symmetric Gaussian uncertainties  
(1312.3524 - ATLAS)
- *Both approaches assume the knowledge of the amplitude, shape (phase-space dependence) and correlations of systematic uncertainties*

# Example: published uncertainties on correlations

1406.0076 – ATLAS jet energy scale uncertainties



*Nominal correlation scenario*



*Weaker - stronger correlation scenarios*

# Scaling factors and NP shifts

$a_\mu^{\text{HVP, LO}}$ shift (Energy range)	Approach 0		Approach 1		
	Scaling factor	$\Delta' \alpha_{\text{had}}(M_Z^2)$	Shift NP <sub>1</sub>	$\sigma' (\Delta \alpha_{\text{had}}(M_Z^2))$	$\Delta' \alpha_{\text{had}}(M_Z^2)$
$a_\mu^{\text{HVP, LO}} - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ (Full HVP)	1.027	0.02826	4.6	$9.0 \cdot 10^{-5}$	0.02774
$(a_\mu^{\text{HVP, LO}} - 1\sigma) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ (Full HVP)	1.020	0.02808	3.5	$9.0 \cdot 10^{-5}$	0.02769
$a_\mu^{\text{HVP, LO}} - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ ([Th.; 1.8 GeV])	1.029	0.02769	4.7	$9.5 \cdot 10^{-5}$	0.02768
$(a_\mu^{\text{HVP, LO}} - 1\sigma) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ ([Th.; 1.8 GeV])	1.022	0.02765	3.5	$9.5 \cdot 10^{-5}$	0.02764
$a_\mu^{\text{HVP, LO}} - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ ([Th.; 1 GeV])	1.034	0.02765	-	-	-
$(a_\mu^{\text{HVP, LO}} - 1\sigma) - a_\mu^{\text{HVP, LO}}(\text{Pheno})$ ([Th.; 1 GeV])	1.026	0.02762	-	-	-
$a_\mu^{\text{Exp}} - a_\mu^{\text{SM}}(\text{Pheno})$ (Full HVP)	1.037	0.02856	6.6	$9.0 \cdot 10^{-5}$	0.02782
$(a_\mu^{\text{Exp}} - 1\sigma) - a_\mu^{\text{SM}}(\text{Pheno})$ (Full HVP)	1.028	0.02831	5.0	$9.0 \cdot 10^{-5}$	0.02775
$a_\mu^{\text{Exp}} - a_\mu^{\text{SM}}(\text{Pheno})$ ([Th.; 1.8 GeV])	1.041	0.02776	6.6	$9.5 \cdot 10^{-5}$	0.02774
$(a_\mu^{\text{Exp}} - 1\sigma) - a_\mu^{\text{SM}}(\text{Pheno})$ ([Th.; 1.8 GeV])	1.031	0.02770	5.0	$9.5 \cdot 10^{-5}$	0.02769
$a_\mu^{\text{Exp}} - a_\mu^{\text{SM}}(\text{Pheno})$ ([Th.; 1 GeV])	1.048	0.02771	-	-	-
$(a_\mu^{\text{Exp}} - 1\sigma) - a_\mu^{\text{SM}}(\text{Pheno})$ ([Th.; 1 GeV])	1.036	0.02766	-	-	-

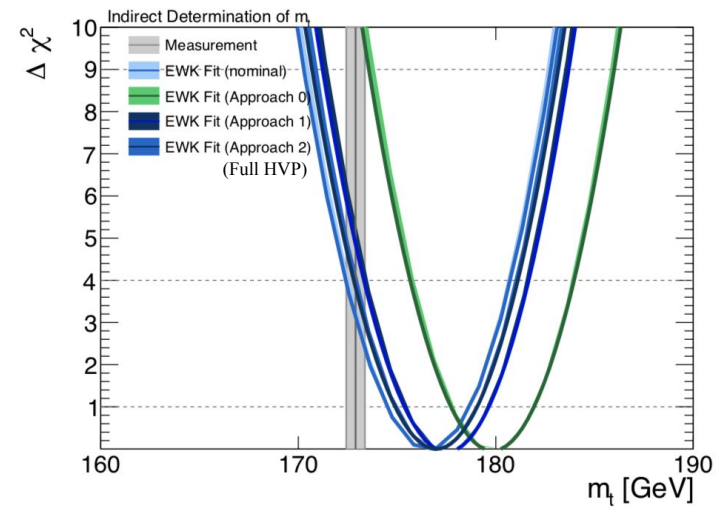
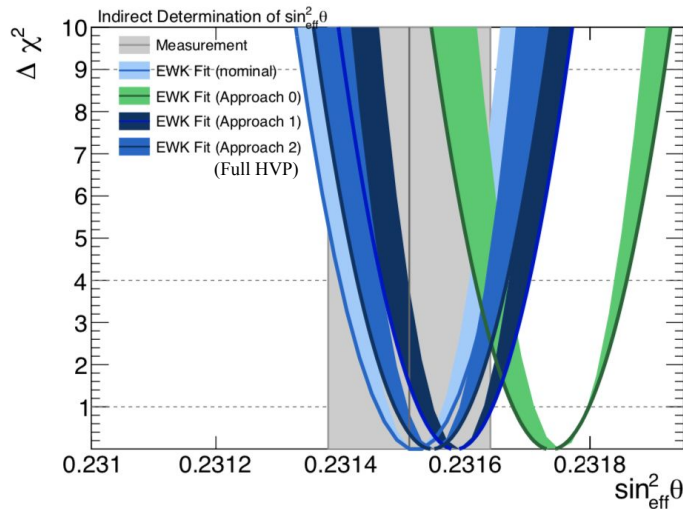
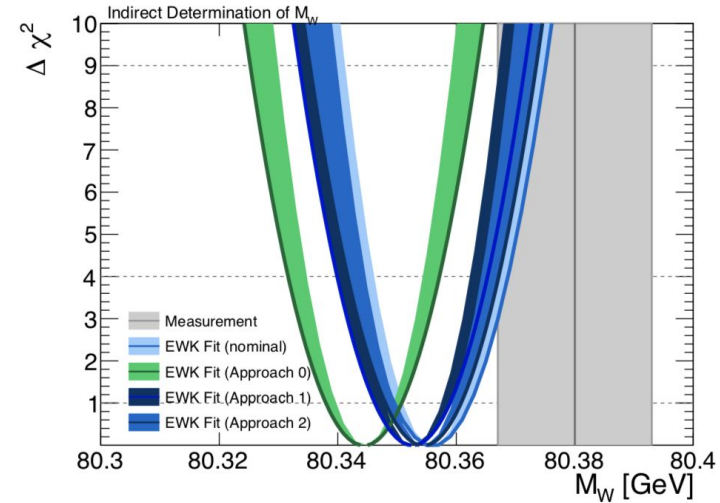
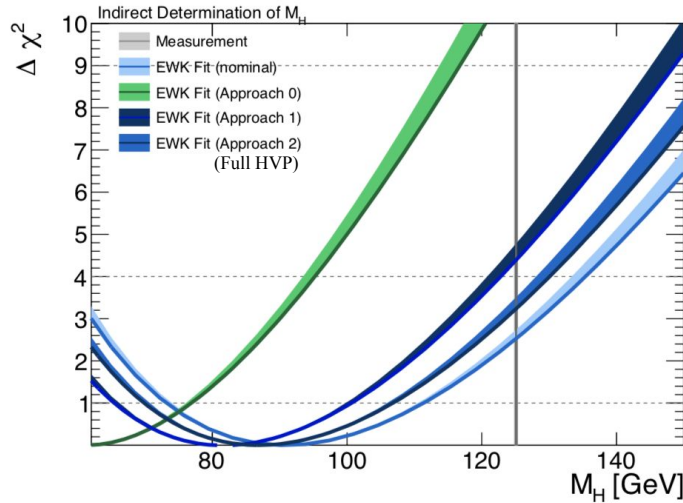
→ Large scaling factors (w.r.t. uncertainties) & significant shifts of NP<sub>1</sub>

# EW fit inputs and $\chi^2$ results

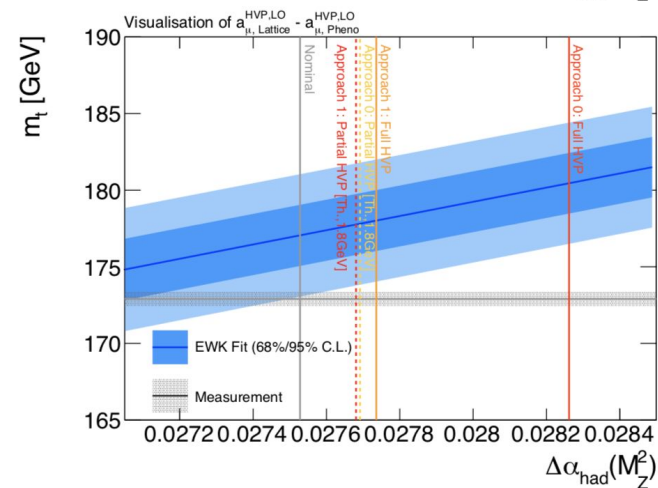
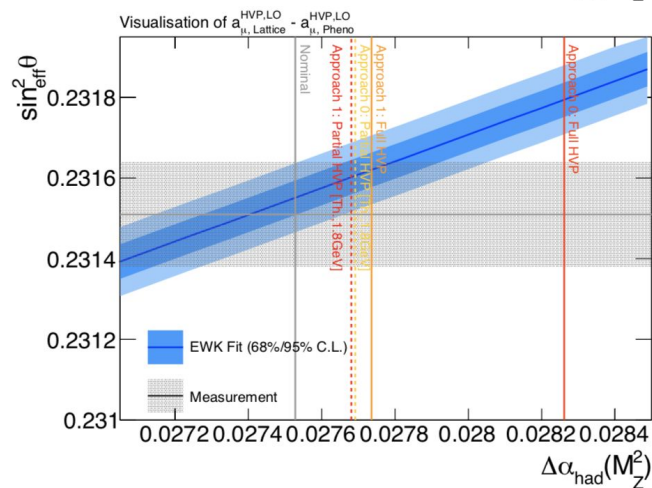
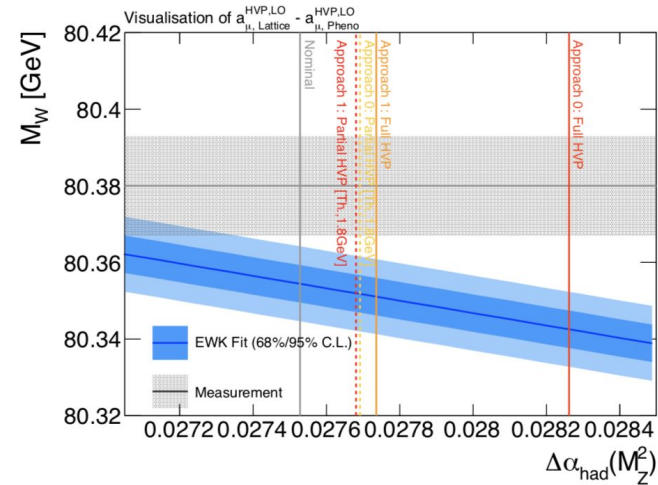
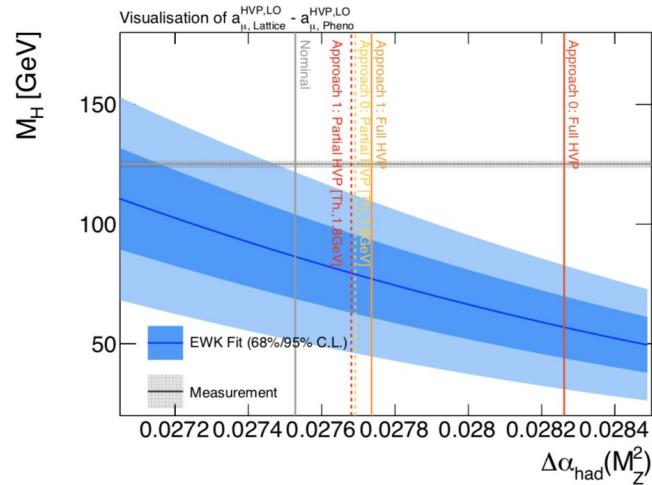
LEP/LHC/Tevatron					
$M_Z$ [GeV]	$91.188 \pm 0.002$	$R_c^0$	$0.1721 \pm 0.003$	$M_H$ [GeV]	$125.09 \pm 0.15$
$\sigma_{\text{had}}^0$ [nb]	$41.54 \pm 0.037$	$R_b^0$	$0.21629 \pm 0.00066$	$M_W$ [GeV]	$80.380 \pm 0.013$
$\Gamma_Z$ [GeV]	$2.495 \pm 0.002$	$A_c$	$0.67 \pm 0.027$	$m_t$ [GeV]	$172.9 \pm 0.5$
$A_l$ (SLD)	$0.1513 \pm 0.00207$	$A_l$ (LEP)	$0.1465 \pm 0.0033$	$\sin^2 \theta_{\text{eff}}^l$	$0.2314 \pm 0.00023$
$A_{\text{FB}}^l$	$0.0171 \pm 0.001$	$m_c$ [GeV]	$1.27_{-0.11}^{+0.07}$ GeV	After HL-LHC	
$A_{\text{FB}}^c$	$0.0707 \pm 0.0035$	$m_b$ [GeV]	$4.20_{-0.07}^{+0.17}$ GeV	$M_W$ [GeV]	$80.380 \pm 0.008$
$A_{\text{FB}}^b$	$0.0992 \pm 0.0016$	$\alpha_s(M_Z)$	$0.1198 \pm 0.003$	$\sin^2 \theta_{\text{eff}}^l$	$0.2314 \pm 0.00012$
$R_l^0$	$20.767 \pm 0.025$	$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) [10^{-5}]$	$2760 \pm 9$	$m_t$ [GeV]	$172.9 \pm 0.3$

$a_\mu^{\text{HVP, LO}}$ shift (Energy range)	Nominal		Approach 0		Approach 1		Approach 2	
	$\Delta' \alpha_{\text{had}}(M_Z^2)$	$\chi^2/\text{ndf}$	$\Delta' \alpha_{\text{had}}(M_Z^2)$	$\chi^2/\text{ndf}$	$\Delta' \alpha_{\text{had}}(M_Z^2)$	$\chi^2/\text{ndf}$	$\Delta' \alpha_{\text{had}}(M_Z^2)$	$\chi^2/\text{ndf}$
	0.02753	18.6/16 (p=0.29)	-	-	-	-	0.02753	28.1/17 (p=0.04)
$a_\mu^{\text{HVP, LO}} - a_\mu^{\text{HVP, LO}}$ (Full HVP)	-	-	0.02826	27.6/16 (p=0.04)	0.02774	20.3/16 (p=0.21)	-	$\chi^2(\text{BMW20-Pheno}): 9.3$
$a_\mu^{\text{HVP, LO}} - a_\mu^{\text{HVP, LO}}$ ([Th.; 1.8 GeV])	-	-	0.02769	19.9/16 (p=0.22)	0.02768	19.8/16 (p=0.23)	-	-
$a_\mu^{\text{HVP, LO}} - a_\mu^{\text{HVP, LO}}$ ([Th.; 1.0 GeV])	-	-	0.02765	19.6/16 (p=0.24)	-	-	-	-
$a_\mu^{\text{Exp}} - a_\mu^{\text{SM}}$ (Pheno) (Full HVP)	-	-	0.02856	33.6/16 (p=0.01)	0.02782	21.2/16 (p=0.17)	-	-
$a_\mu^{\text{Exp}} - a_\mu^{\text{SM}}$ (Pheno) ([Th.; 1.8 GeV])	-	-	0.02776	20.6/16 (p=0.19)	0.02774	20.4/16 (p=0.20)	-	-
$a_\mu^{\text{Exp}} - a_\mu^{\text{SM}}$ (Pheno) ([Th.; 1.0 GeV])	-	-	0.02771	20.1/16 (p=0.22)	-	-	-	-

# EW fit results: $\chi^2$ scans



# EW fit results: parameter scans for varying $\Delta\alpha_{\text{had}}(M_Z^2)$



# EW fit results: indirect determination of $\Delta\alpha_{\text{had}}(M_Z^2)$

